# **Collisional excitation of radical molecules** in the interstellar medium

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Figure: The oxygen chemistry in molecular clouds  $_{2/16}$ 

# Application to the NH( ${}^{3}\Sigma$ )-H<sub>2</sub>( ${}^{1}\Sigma$ ) collisional system

# **Non-reactive collisions**

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Isotopologue ND is also detected

- $\rightarrow$  Probe chemical pathways  $\rightarrow$  Trace the  $\frac{[ND]}{[NH]}$  ratio
- $\rightarrow$  Trace the  $\overline{[NH]}$  ratio  $\rightarrow$  Trace history of molecular cloud

formation

# **Excitation mechanisms in the ISM**

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Without rate coefficients: application of local thermodynamic equilibrium (LTE) models

 $\Rightarrow$  generally not fullfilled

Lack of data for NH–H $_2$  and ND–H $_2$  collisional systems

<sup>1</sup>A. Bacmann et al., A&A 587, A26 (2016)



Figure: Observed spectrum of NH with a LTE model as red solid  ${\sf lines}^1$ 

### Obtention of the rate coefficients

Resolution of the Time-independant Schrödinger equation through Born-Oppenheimer approximation

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Most accurate approach  $\checkmark$  Most expensive one  $\times$ 

# Potential energy surface

- Linear rigid rotor approximation
- 4D PES<sup>1</sup>
- RCCSD(T)-F12/aVTZ method
- Calculations are done with the  ${\rm MOLPRO}\ code^2$
- The PES was expanded in bispherical harmonics

<sup>1</sup>Pirlot Jankowiak *et al.*, J. Chem. Phys. 155, 134303 (2021) <sup>2</sup>Werner, Molpro, a package of ab initio programs, version 2015.1, 2015



Figure: NH-H<sub>2</sub> representation in Jacobi coordinates

# Energetic structure of $NH(^{3}\Sigma)$

 $\begin{array}{l} \text{Open-shell molecule} \rightarrow \textbf{S} = 1 \\ \rightarrow \textbf{j}_1 = \textbf{N} + \textbf{S} \\ \text{Rotational levels } \textit{N} \rightarrow \text{Fine structure levels } \textit{N}_{j_1} \end{array}$ 

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NH and ND possess nonzero nuclear spins:

- $\mathbf{I}_{H/D} \rightarrow \mathbf{F}_1 = \mathbf{j}_1 + \mathbf{I}_{H/D}$
- $I_N \rightarrow F = F_1 + I_N$

Hyperfine structure levels  $N_{j_1,F_1,F}$ 

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$$\mathsf{NH}(\mathsf{N}_{j_1,\mathsf{F}_1,\mathsf{F}}) + \mathsf{H}_2(j_2) \to \mathsf{NH}(\mathsf{N}_{j_1',\mathsf{F}_1',\mathsf{F}'}) + \mathsf{H}_2(j_2)$$

<sup>1</sup> Klaus et al., Astron. Astrophys. 322, L1–L4 (1997)



Figure: NH first hyperfine energy levels<sup>1</sup>

# Fine structure excitation of NH by H<sub>2</sub>



- Resolution of the close-coupling equations
- Calculations have been done using HIBRIDON scattering code<sup>1</sup>

Figure: Temperature dependance of rate coefficients of NH with *para*- (solid) and *ortho*-H<sub>2</sub>(dashed) for  $N_j \rightarrow N'_{j'}$  transitions

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 $\rightarrow$  Most of rate coefficients deviate within a factor of  $\sim 2-3$  for NH-ortho-H2 transitions

 $\overset{0}{\rightarrow} \mbox{ Some of the most dominant} \\ {}^{0} \mbox{ transitions (> 4 \times 10^{-12} \mbox{ cm}^3 \mbox{s}^{-1}) \mbox{ are} \\ \mbox{ larger by a factor > 10 !} \\ {}^{1} \mbox{ Alexander 2023}$ 

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 $\delta r = 0.11 a_0 + \text{rotational structures} \rightarrow \text{deviation of the rate coefficients up to a factor of 3}$ 

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Type of interaction	Memory/task (Gb)	CPU time (hours)
Atom-Molecule ( $I_1, I_2 \neq 0$ )	$\leq$ 7	> 1000000
Molecule-Molecule $(I_1  eq 0)$	$\sim 10$	$\sim 1000000$
Molecule-Molecule $(I_1, I_2 \neq 0)$	$\sim 70$	×

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NH environments  $\rightarrow$  T > 50 K The internal structure of H\_2 MUST be taken into account for NH excitation

Figure: Temperature evolution of the  $o/p(H_2)$  ratio

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$$k_{i \to f}^{IOS} \propto \sum_{l} H_{l}(i, f) k_{1,0 \to l+1,l}^{CC}$$
$$WMEF = \frac{\sum_{i} k_{i}^{rec} r_{i}}{\sum_{i} k_{i}^{rec}}$$
$$r_{i} = max(k_{i}^{rec} / k_{i}^{NG}, k_{i}^{NG} / k_{i}^{rec})$$



<sup>1</sup>Dumouchel *et al.*, J. Chem. Phys. 137, 114306 (2012)

Figure: Systematic comparison of NH–He rate coefficients with recoupling  $ones^1$ 



Collisional excitation of open-shell molecules with molecular hydrogen is a true theoretical and numerical challenge

- First fine structure resolved rate coefficients for NH–H $_2$  and ND–H $_2$  collisional systems (5 300 K)
- Hyperfine structure treatment still needs improvement
- $\rightarrow$  The internal structure of  $H_2$  is not negligible
- $\rightarrow$  The isotopic subsitution has an impact on the order of magnitude of rate coefficents

# Perspectives: The $OH^+(^{3}\Sigma) - H_2$ collisional system

 $OH^+$  is a very reactive molecule  $\rightarrow$  rapidly destroyed after its formation **Competition between reactive and inelastic collisions** 

$$\mathsf{OH}^+(N,j_1) + \mathsf{H}_2(j_2) \rightarrow \begin{cases} \mathsf{OH}^+(N',j_1') + \mathsf{H}_2(j_2) & \text{inelastic} \\ \mathsf{H}_2\mathsf{O}^+(N',j_1') + \mathsf{H} & \text{reactive} \end{cases}$$

 $\rightarrow$  Both inelastic and reactive channels will be taken into account with a statistical approach



Figure: Reaction path of the  $OH^+ - H_2$ collisional system (a) and the representation of the complex in Jacobi coordinates (b)<sup>1</sup>

<sup>1</sup> Song *et al.*, J. Phys. Chem. A 2016, 120, 4742-4748

# Acknowledgments



# Comparison of $NH-H_2$ with NH-He



 $k_{H_2} = 1.3335 \times k_{He}$ 

# The Statistical Adiabatic Channels Model (SACM)

In the time-independant Close-Coupling equations:

$$\frac{-\hbar^2}{2\mu}\frac{d^2F(\alpha)}{dR^2} + \sum_{\alpha'} \langle \alpha' | H_{int} + V + \frac{L^2}{2\mu R} | \alpha \rangle F(\alpha) = EF(\alpha)$$

- The kinetic term is removed
- Calculations are independent of collisional energy  $\rightarrow$  Adiabatic curves<sup>1</sup>



$$|S_{ij}(E,j)|^2 = \begin{cases} rac{1}{N(E,j)} & \text{open channels} \\ 0 & \text{closed channels} \end{cases}$$

<sup>1</sup> Loreau, Lique, and Faure, ApJL, 853 L5 (2018)