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# AN EXTERIOR OPTIMAL TRANSPORT PROBLEM 

JULES CANDAU-TILH, MICHAEL GOLDMAN, AND BENOIT MERLET


#### Abstract

This paper deals with a variant of the optimal transportation problem. Given $f \in L^{1}\left(\mathbb{R}^{d},[0,1]\right)$ and a cost function $c \in C\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)$ of the form $c(x, y)=k(y-x)$, we minimise $\int c d \gamma$ among transport plans $\gamma$ whose first marginal is $f$ and whose second marginal is not prescribed but constrained to be smaller than $1-f$. Denoting by $\Upsilon(f)$ the infimum of this problem, we then consider the maximisation problem $\sup \left\{\Upsilon(f): \int f=m\right\}$ where $m>0$ is given. We prove that maximisers exist under general assumptions on $k$, and that for $k$ radial, increasing and coercive these maximisers are the characteristic functions of the balls of volume $m$.


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## 1. Introduction

In this paper, we study the optimization problems associated with functionals which favour dispersion and are based on some Wasserstein energies. These functionals correspond to the non-local term of the energy studied in [4].

We denote by $\mathcal{M}_{+}\left(\mathbb{R}^{d}\right)$ the set of positive Radon measures on $\mathbb{R}^{d}$. Given a cost function $c$ and $\mu, \nu \in \mathcal{M}_{+}\left(\mathbb{R}^{d}\right)$, we let $\mathcal{T}_{c}(\mu, \nu)$ be the $c$-transport cost between $\mu$ and $\nu$ (see Section 2 for the exact definition of $\mathcal{T}_{c}$ ). Given a measurable set $E \subset \mathbb{R}^{d}$ with finite volume, we consider the optimisation problem

$$
\begin{equation*}
\Upsilon_{\text {set }}(E):=\inf \left\{\mathcal{T}_{c}(E, F): F \subset \mathbb{R}^{d} \text { Lebesgue measurable, }|F|=|E|,|F \cap E|=0\right\} \tag{1.1}
\end{equation*}
$$

where we identify $E$ with the restriction of the Lebesgue measure on $E$. Given $m>0$, we introduce the maximisation problem

$$
\begin{equation*}
\mathcal{E}_{\mathrm{set}}(m):=\sup _{|E|=m} \Upsilon_{\mathrm{set}}(E) \tag{1.2}
\end{equation*}
$$

The main goal of the article is to investigate the existence of maximisers for this problem and to characterise these latter.
If we apply the direct method of the Calculus of Variations, we obtain that, up to extraction, any maximising sequence $E_{n}$ converges weakly to some function $u_{\infty} \in L^{1}\left(\mathbb{R}^{d},[0,1]\right)$. However, there is no guarantee at this point that $u_{\infty}$ is a characteristic function or has mass $m$. Our strategy is to extend the functional $\Upsilon_{\text {set }}$ as a functional $\Upsilon$ defined on $L^{1}\left(\mathbb{R}^{d},[0,1]\right)$. Applying the bathtub principle (see Proposition 4.11) to a maximiser of the relaxed problem, we show that the supremum in (1.2) is actually reached (see Corollary 1.2). This relaxation approach is not new: it was successfully applied to several variational problems in the last few years (see for instance $[6,1,20,2]$ ).

Given $f \in L^{1}\left(\mathbb{R}^{d},[0,1]\right)$, the set of admissible exterior transport plans is defined as

$$
\Pi_{f}:=\left\{\gamma \in \mathcal{M}_{+}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right): \gamma_{x}=f, \gamma_{y} \leq 1-f\right\}
$$

Here, the measures $f d x$ and $(1-f) d y$ are identified with their respective densities and $\gamma_{x}$ and $\gamma_{y}$ denote respectively the first and second marginals of $\gamma$. We then define the primal problem

$$
\Upsilon(f):=\inf \left\{\int_{1} c d \gamma: \gamma \in \Pi_{f}\right\}
$$

We have $\Upsilon\left(\chi_{E}\right)=\Upsilon_{\text {set }}(E)$ under mild assumptions on $c$ (see Theorem 4.4). Given $m>0$, our maximisation problem is now

$$
\mathcal{E}(m):=\sup \left\{\Upsilon(f): f \in L^{1}\left(\mathbb{R}^{d},[0,1]\right), \int f d x=m\right\} .
$$

By abuse of notation and when no confusion is possible, we refer to the variational problems by the values they attain (e.g. we write $\Upsilon_{\text {set }}(E)$ for (1.1)).

### 1.1. Main results.

The first important result of this article is that maximisers of $\mathcal{E}(m)$ exist whenever $c$ is of the form $c(x, y)=k(y-x)$ for some $k: \mathbb{R}^{d} \rightarrow \mathbb{R}_{+}$and satisfies
(H1) $k \in C\left(\mathbb{R}^{d}, \mathbb{R}_{+}\right), k(0)=0$ and $k(x) \rightarrow \infty$ as $|x| \rightarrow \infty$,
(H2) $\forall x \neq 0$,

$$
\limsup _{r \rightarrow 0} \frac{1}{r^{d}}\left|B_{r}(x) \cap\left\{y \in \mathbb{R}^{d}, k(y)<k(x)\right\}\right|>0,
$$

(H3) $\forall \sigma \in \mathbb{S}^{d-1}, r \mapsto k(r \sigma)$ is increasing on $\mathbb{R}_{+}$.
Notice that $k$ is not assumed to be strictly convex, so that our results hold in cases where the existence of an optimal transport map is not guaranteed. Also observe that all the costs of the form $k(z)=|z|^{p}$ with $0<p<\infty$ satisfy the above hypotheses. However, radial symmetry is not required and the costs $k(z)=|z|^{p} h(z /|z|)$ with $h$ positive and Lipschitz continuous on $\mathbb{S}^{d-1}$ are also admissible.
Theorem 1.1. Assume that $c(x, y)=k(y-x)$ for $x, y \in \mathbb{R}^{d}$ with $k$ satisfying (H1),(H2) $\mathcal{Z}$ (H3). Then, for any $m>0$ the supremum in $\mathcal{E}(m)$ is attained. Moreover, there exists $R_{*}=R_{*}(m)$ such that (up to translation) any maximiser is supported in the ball $\bar{B}_{R_{*}}$.
Once the existence of maximisers for $\mathcal{E}(m)$ is established, the bathtub principle (see Proposition 4.11) and a saturation result (see Theorem 4.4) imply that (1.2) admits solutions.
Corollary 1.2. Assume that $c$ satisfies the hypotheses of Theorem 1.1. Then, (1.2) admits a maximiser and $\mathcal{E}_{\text {set }}(m)=\mathcal{E}(m)$ for any $m>0$.

As a second main result, we establish that if $k$ is furthermore radially symmetric then $\mathcal{E}(m)$ and $\mathcal{E}_{\text {set }}(m)$ are uniquely maximised by balls of volume $m$.
Theorem 1.3. Assume that $c(x, y)=k(|y-x|)$ for some $k \in C\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)$increasing and such that $k(0)=0$ and $k(x) \rightarrow \infty$ as $x \rightarrow \infty$. Then, for any $m>0$, the maximisers of $\mathcal{E}(m)$ (and consequently those of $\left.\mathcal{E}_{\text {set }}(m)\right)$ are the balls of volume $m$.

We point out that cost functions satisfying the hypotheses of Theorem 1.3 also satisfy hypotheses (H1),(H2)\&(H3). Let us briefly sketch the proofs of these three results. They all strongly rely on the properties of the dual problem

$$
\Upsilon^{*}(f):=\sup \left\{\int(f \varphi+(1-f) \psi) d x:(\varphi, \psi) \in \Phi\right\}
$$

where

$$
\Phi:=\left\{(\varphi, \psi) \in C_{b}\left(\mathbb{R}^{d}\right) \times C_{b}\left(\mathbb{R}^{d}\right), \psi \leq 0, \varphi(x)+\psi(y) \leq c(x, y) \forall(x, y) \in \mathbb{R}^{d} \times \mathbb{R}^{d}\right\}
$$

We establish Theorem 1.1 using the direct method of Calculus of Variations. The main difficulty is to establish compactness of maximising sequences. If we refer to the concentrationcompactness principle [17], we have to prove that given a maximising sequence $f_{n}$, no mass escapes at infinity. To do so we establish two crucial results. The first one is that $m \mapsto \mathcal{E}(m) / m$ is increasing (see Proposition 4.7). This implies that $m \mapsto \mathcal{E}(m)$ is strictly superadditive, i.e. that for $m>m^{\prime}>0$,

$$
\begin{equation*}
\mathcal{E}(m)+\mathcal{E}\left(m-m^{\prime}\right)<\mathcal{E}(m) \tag{1.3}
\end{equation*}
$$

Notice that this is the counterpart of the strict subbadditivity inequality (also called binding inequality) which is known to provide compactness in minimisation problems, see e.g. [17, 9, 10]. Using the dual formulation $\Upsilon^{*}$ of $\Upsilon$, we obtain the second crucial result for Theorem 1.3: a monotonicity principle on the sum of marginals of maximisers $\gamma$ of $\Upsilon(f)$ (see Corollary 4.6). Combining this and (1.3), we prove that if $f$ is almost maximising then most of its mass must remain in a bounded region (see Proposition 4.9). This gives tightness of maximising sequences for $\mathcal{E}(m)$.

To prove Corollary 1.2, we consider a maximiser $f$ of $\mathcal{E}(m)$ provided by Theorem 1.1 and a pair of potentials $(\varphi, \psi)$ optimal for the dual problem $\Upsilon^{*}(f)$. Using the definition of $\Upsilon^{*}$ we see that $f$ is a maximiser of

$$
\sup \left\{\int \tilde{f}(\varphi-\psi): 0 \leq \widetilde{f} \leq 1, \int \widetilde{f}=m\right\}
$$

By the bathtub principle, $f=\chi_{\{\varphi-\psi>t\}}+\theta$ for some $t \in \mathbb{R}$ and some $\theta \in L^{1}\left(\mathbb{R}^{d},[0,1]\right)$ supported in $\{\varphi-\psi=t\}$. Then for any measurable subset $G \subset\{\varphi-\psi=t\}$ with $|G|=\int \theta$, the characteristic function of $E:=\{\varphi-\psi>t\} \cup G$ is also a maximiser for $\Upsilon(f)$. By Theorem 4.4 and Corollary 4.5 applied to $E$, there exists $F \subset \mathbb{R}^{d}$ such that any minimiser $\gamma$ of $\Upsilon\left(\chi_{E}\right)$ satisfies $\gamma_{y}=\chi_{F}$. This finally implies that $E$ maximises (1.2).

Regarding Theorem 1.3, as explained in Section 5, we may assume without loss of generality that $m=\omega_{d}$, the volume of the unit ball. Combining Theorem 1.1 and Lemma 4.10 yields that

$$
\begin{equation*}
\sup _{\int=\omega_{d}}\left\{\sup _{\left(\psi^{c}, \psi\right) \in \Phi}\left\{\int f\left(\psi^{c}-\psi\right)+\int \psi\right\}\right\} \tag{1.4}
\end{equation*}
$$

coincides with $\mathcal{E}(m)$ and admits a solution $\left(f, \psi^{c}, \psi\right)$, where $\psi^{c}$ is the $c$-transform of $\psi$ (see Definition 2.2). To show that balls are maximisers of $\mathcal{E}(m)$, we establish that each term in (1.4) is improved by replacing $f$ by $\chi_{B_{1}}$ and $\psi$ by its symmetric increasing rearrangement $\psi_{*}$ (see Definition 5.1). As $\int \psi=\int \psi_{*}$, the third term in (1.4) does not change under rearrangement. Regarding the second term, combining the Hardy-Littlewood inequality (see [16, Theorem 3.4]) and the bathtub principle yields (recall that $\psi \leq 0$ )

$$
-\int f \psi \leq-\int f^{*} \psi_{*} \leq-\int \chi_{B_{1}} \psi_{*}
$$

where $f^{*}$ is the symmetric decreasing rearrangement of $f$ (see Definition 5.1). The study of the first term $\int f \psi^{c}$ is more involved. Using the Brunn-Minkowski inequality, we obtain the following crucial comparison:

$$
\left(\psi^{c}\right)^{*} \leq\left(\psi_{*}\right)^{c}
$$

Combining this inequality with the Hardy-Littlewood inequality yields

$$
\begin{equation*}
\int f \psi^{c} \leq \int f^{*}\left(\psi^{c}\right)^{*} \leq \int f^{*}\left(\psi_{*}\right)^{c} \tag{1.5}
\end{equation*}
$$

Additionally, as $\left(\psi_{*}\right)^{c}$ is non-increasing, $\chi_{B_{1}}$ is a maximiser of

$$
\begin{equation*}
\sup \left\{\int \widetilde{f}\left(\psi_{*}\right)^{c}: 0 \leq \widetilde{f} \leq 1, \int \widetilde{f}=\omega_{d}\right\} \tag{1.6}
\end{equation*}
$$

so that $\int f^{*}\left(\psi_{*}\right)^{c} \leq \int\left(\psi_{*}\right)^{c} \chi_{B_{1}}$. Lastly, by (1.5), $\int f \psi^{c} \leq \int \chi_{B_{1}}\left(\psi_{*}\right)^{c}$. This eventually proves that unit balls maximise $\mathcal{E}\left(\omega_{d}\right)$.

As for uniqueness, the key property to establish is that $\left(\psi_{*}\right)^{c}$ is decreasing on $B_{1}$ (see Lemma 5.3). Indeed, by [16, Theorem 3.4], this implies that $\chi_{B_{1}}$ is the unique maximiser of (1.6). Combining this with the fact that the inequalities in (1.5) are now equalities, we obtain that $f^{*}=\chi_{B_{1}}$, so that $f=\chi_{E}$ for some $E \subset \mathbb{R}^{d}$. Using the equality case of the Brunn-Minkowski inequality, we then show that (up to a translation) $f=\chi_{B_{1}}$, concluding the proof.

### 1.2. Motivation.

In [3], the following variational problem was introduced:

$$
\begin{equation*}
\inf _{|E|=\omega_{d}}\left\{P(E)+\alpha \Upsilon_{p}(E)\right\} \tag{1.7}
\end{equation*}
$$

where $\alpha>0$ and where $\Upsilon_{p}$ is the functional $\Upsilon$ defined in (1.1) with the cost $c(x, y)=|x-y|^{p}$. Such a variational problem may be used to model the formation of bi-layer biological membranes (see $[21,18]$ ). Existence of minimisers were obtained in the series of work $[3,23,19,4]$.

Notice that (1.7) is an isoperimetric problem with a non-local term $\Upsilon_{p}$. One of the best-known examples of this type of problem is Gamow's liquid drop model for the atomic nucleus. Since the beginning of the 2010s (see [5] for an historical perspective), this model has received a lot of attention from the mathematical community, and several versions of it have been studied, see for instance $[15,13,14,12]$. In this framework, the perimeter term represents the local attractive forces while the repulsive non-local term is given by the Riesz potential

$$
V_{\beta}(E):=\int_{E} \int_{E} \frac{d x d y}{|x-y|^{d-\beta}}
$$

where $\beta \in(0, d)$. A consequence of Riesz's rearrangement inequality is that balls are the volume constrained maximisers of $V_{\beta}$. This illustrates the competition between the perimeter and the Riesz potential. It is thus natural to investigate whether similar properties hold for (1.7). In our case, the proof is much more involved since the rearrangement argument does not seem to work well for the primal problem. We consider instead the dual problem $\Upsilon^{*}$ and study the (fortunately favourable) interplay between rearrangement and $c$-transforms.

As a closing remark, we point out that the functional $\Upsilon$ is a particular case of the optimal partial transport problem studied in [8, 7].

### 1.3. Organization of the article.

The paper is structured as follows. In Section 2, we introduce the notation and review standard facts related to optimal transport in complete separable metric spaces. In Section 3, we obtain preliminary results on the functional $\Upsilon$ defined in compact spaces. In Section 4, we establish Theorem 1.1. Eventually, in Section 5, we prove Theorem 1.3.

## 2. Notation and preliminary Results

### 2.1. Notation.

Let $\left(X, d_{X}\right)$ be a Polish space endowed with a positive Radon measure $\lambda$.
Given a function $f: X \rightarrow \mathbb{R}$, we decompose it as:

$$
f=f_{+}+f_{-} \quad \text { with } \quad f_{+}:=\max (0, f):=0 \vee f \quad \text { and } \quad f_{-}:=\min (0, f):=0 \wedge f
$$

Let us stress that $f_{-}$is non-positive, contrary to the classical decomposition of a function into its positive and negative parts.

We endow $\mathcal{M}_{+}(X)$ with the topology of weak-* convergence, that is the topology induced by duality with $C_{b}(X)$. The convergence of a sequence $\mu_{n} \in \mathcal{M}_{+}(X)$ to $\mu \in \mathcal{M}_{+}$is written: $\mu_{n} \stackrel{*}{\rightharpoonup} \mu$ as $n \rightarrow \infty$ 。

Given a measure $\mu \in \mathcal{M}_{+}(X)$ and a set $A \subset X$, the restriction of $\mu$ to $A$ is the measure $\mu\llcorner A$ defined as $\mu\llcorner A(B):=\mu(B \cap A)$ for every Borel set $B$ of $X$. The support of $\mu$, denoted by $\operatorname{supp} \mu$, is the closed set defined by

$$
\operatorname{supp} \mu:=\{x \in X: \mu(A)>0 \text { for all open set A containing } x\}
$$

Given $f \in L^{1}(X, \lambda)$ the support of $f$ is defined as the support of the measure $f d \lambda$ and denoted by $\operatorname{supp} f$. We identify the measure $f d \lambda$ with its density $f$ and write $f_{n} \stackrel{*}{\square} f$ as $n \rightarrow \infty$ to signify that $\int f_{n} \xi$ converges to $\int f \xi$ for every $\xi \in C_{b}(X)$.

Given a function $f \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{d}, \mathbb{R}\right)$, we denote by $\operatorname{Leb}(f)$ the set of its Lebesgue points.

Given $x \in \mathbb{R}^{d}$ and $r>0, B_{r}(x)$ denotes the open ball of radius $r$ centred at $x$, and $B_{r}$ denotes the open ball of radius $r$ centred at 0 . The closed ball of radius $r$ centred at $x$ is denoted by $\bar{B}_{r}(x)$. The volume of the unit ball in $\mathbb{R}^{d}$ is denoted by $\omega_{d}$.

Given two sets $A, B$ of $\mathbb{R}^{d}$, we define their sum $A+B:=\{a+b, a \in A, b \in B\}$. The gap between $A$ and $B$ is $d(A, B):=\inf \{|a-b|, a \in A, b \in B\}$.

### 2.2. Optimal transport theory.

In this subsection, we recall some results regarding standard optimal transport theory. Most of the material presented here comes from [22, Chapter 1].

Let $\left(X, d_{X}\right)$ be a complete separable metric space (i.e. a Polish space) and let $c: X \times X \rightarrow \mathbb{R}$ be measurable. Given $\mu, \nu \in \mathcal{M}_{+}(X)$ such that $\mu(X)=\nu(X)$, the Kantorovitch problem with marginals $\mu$ and $\nu$ and cost $c$ is

$$
\begin{equation*}
\mathcal{T}_{c}(\mu, \nu):=\inf \left\{\int c d \gamma: \gamma \in \Pi(\mu, \nu)\right\} \tag{2.1}
\end{equation*}
$$

where $\Pi(\mu, \nu)$ is the set of transport plans between $\mu$ and $\nu$, i.e.

$$
\Pi(\mu, \nu):=\left\{\gamma \in \mathcal{M}_{+}(X \times X): \gamma_{x}=\mu, \gamma_{y}=\nu\right\}
$$

Problem (2.1) admits a dual formulation given by

$$
\begin{equation*}
\mathcal{T}_{c}^{*}(\mu, \nu):=\sup \left\{\int \varphi d \mu+\int \psi d \nu: \varphi, \psi \in C_{b}(X), \varphi \oplus \psi \leq c\right\} \tag{2.2}
\end{equation*}
$$

where the function $\varphi \oplus \psi$ is defined on $X \times X$ by $(\varphi \oplus \psi)(x, y):=\varphi(x)+\psi(y)$.
Theorem 2.1 (Theorem 1.7 of [22]). Let $c: X \times X \rightarrow \mathbb{R}$ be lower semi-continuous and bounded from below and let $\mu, \nu \in \mathcal{M}_{+}(X)$ with $\mu(X)=\nu(X)$. Then (2.1) admits a solution and

$$
\mathcal{T}_{c}(\mu, \nu)=\mathcal{T}_{c}^{*}(\mu, \nu)
$$

Using the notion of $c$-transform of a function, the maxima of (2.2) can be further characterised.
Definition 2.2. Given a function $\xi: X \rightarrow \mathbb{R} \cup\{+\infty\}$, we define its $c$-transform (or $c$-conjugate) $\xi^{c}: X \rightarrow \mathbb{R} \cup\{-\infty\}$ by

$$
\xi^{c}(y):=\inf _{x \in X}\{c(x, y)-\xi(x)\}
$$

Denoting $\bar{c}(y, x):=c(x, y)$, the $\bar{c}$-transform of $\zeta: X \rightarrow \mathbb{R} \cup\{+\infty\}$ is given by

$$
\zeta^{\bar{c}}(x):=\inf _{y \in X}\{\bar{c}(y, x)-\zeta(y)\}
$$

A function $\psi: X \rightarrow \mathbb{R} \cup\{-\infty\}$ is said to be $\bar{c}$-concave if there exists $\xi: X \rightarrow \mathbb{R} \cup\{+\infty\}$ such that $\psi=\xi^{c}$ (the definition of $c$-concavity is analogous).

Definition 2.3. Let $\left(X, d_{X}\right)$ be a metric space and $\omega \in C\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)$be increasing and such that $\omega(0)=0$. A function $\varphi: X \rightarrow \mathbb{R}$ is $\omega$-continuous if for all $x, x^{\prime} \in X$,

$$
\left|\varphi(x)-\varphi\left(x^{\prime}\right)\right| \leq \omega\left(d_{X}\left(x, x^{\prime}\right)\right)
$$

Similarly, we say that $c: X \times X \rightarrow \mathbb{R}$ is $\omega$-continuous if for all $x, x^{\prime}, y, y^{\prime} \in X$,

$$
\left|c(x, y)-c\left(x^{\prime}, y^{\prime}\right)\right| \leq \omega\left(d_{X}\left(x, x^{\prime}\right)+d_{X}\left(y, y^{\prime}\right)\right)
$$

Proposition 2.4. Let $\varphi, \psi: X \rightarrow \mathbb{R}$ be fixed and assume that $\varphi^{c}$ and $\psi^{\bar{c}}$ take real values. The following statements hold:
(i) If $c$ is $\omega$-continuous, then $\varphi^{c}$ is also $\omega$-continuous,
(ii) $\varphi^{c \bar{c}} \geq \varphi$, and $\varphi^{c \bar{c}}=\varphi$ if and only if $\varphi$ is c-concave,
(iii) $\varphi^{c}$ is the largest function $\psi$ compatible with the constraint $\varphi \oplus \psi \leq c$ and $\psi^{\bar{c}}$ is the largest function $\varphi$ compatible with the constraint $\varphi \oplus \psi \leq c$.

Remark that if $X$ is compact and $\varphi, \psi$ and $c$ are bounded then $\varphi^{c}$ and $\psi^{\bar{c}}$ take real values. Moreover, if $c$ is continuous, say $\omega$-continuous, the proposition states that $\varphi^{c}$ and $\psi^{\bar{c}}$ are $\omega$ continuous. This yields the following existence result for (2.2).

Theorem 2.5 (Proposition 1.11 of [22]). Let $X$ be a compact metric space and $c: X \times X \rightarrow \mathbb{R}$ be continuous. Then there exists a solution $(\varphi, \psi)$ to $(2.2)$, where $\varphi$ is c-concave and $\psi=\varphi^{c}$. In particular,

$$
\mathcal{T}_{*}^{c}(\mu, \nu)=\max \left\{\int \varphi d \mu+\int \varphi^{c} d \nu: \varphi c \text {-concave }\right\}
$$

A pair of functions maximising (2.2) is called a pair of Kantorovitch potentials.

## 3. Study of $\Upsilon$ in compact metric spaces

Let $\left(X, d_{X}\right)$ be a compact metric space and let $c: X \times X \rightarrow \mathbb{R}$ be a continuous cost function. We endow $\left(X, d_{X}\right)$ with a measure $\lambda \in \mathcal{M}_{+}(X)$ such that $\lambda(X)>0$ and denote by $L^{1}(X)$ the set of $\mathbb{R}$-valued functions integrable with respect to $\lambda$. Given $f \in L^{1}(X)$, we define the set of admissible transport plans

$$
\Pi_{f}:=\left\{\gamma \in \mathcal{M}_{+}(X \times X): \gamma_{x}=f, \gamma_{y} \leq 1-f\right\}
$$

and the primal problem

$$
\begin{equation*}
\Upsilon(f):=\inf \left\{\int c d \gamma: \gamma \in \Pi_{f}\right\} \tag{3.1}
\end{equation*}
$$

Notice that $\Pi_{f}$ is empty whenever $f$ does not satisfy $0 \leq f \leq 1$ or when $\int f d \lambda>\lambda(X) / 2$. In the other cases, there exists $g \in L^{1}(X)$ such that $g \geq 0, f+g \leq 1$ and $\int g d \lambda=\int f d \lambda$. Thus,

$$
\gamma:=\frac{1}{\int f d \lambda}(f d \lambda) \otimes(g d \lambda) \in \Pi_{f}
$$

and $\Pi_{f}$ is not empty. We now fix $0<m \leq \lambda(X) / 2$ and define

$$
L_{m}^{1}:=\left\{f \in L^{1}(X,[0,1]): \int f \leq m\right\}
$$

Given $f \in L_{m}^{1}$ and $\varphi, \psi \in C(X)$, we set

$$
\begin{equation*}
K_{f}(\varphi, \psi):=\int(f \varphi+(1-f) \psi) d \lambda \tag{3.2}
\end{equation*}
$$

and define the dual problem

$$
\begin{equation*}
\Upsilon^{*}(f):=\sup \left\{K_{f}(\varphi, \psi):(\varphi, \psi) \in \Phi\right\} \tag{3.3}
\end{equation*}
$$

where

$$
\Phi:=\{(\varphi, \psi) \in C(X) \times C(X), \psi \leq 0, \varphi \oplus \psi \leq c\}
$$

For the remainder of the section we fix $f \in L_{m}^{1}$. As in the classical theory of optimal transport, a simple application of the direct method of Calculus of Variations shows that (3.1) admits a minimiser.

Proposition 3.1. Assume that $X$ is a compact metric space and that $c \in C(X \times X, \mathbb{R})$. Then, the infimum in (3.1) is a minimum.

Remark 3.2. If we let $f \in L_{m}^{1}$, by Proposition 3.1, there exists $\gamma \in \mathcal{M}_{+}(X \times X)$ optimal for $\Upsilon(f)$. Notice that $\gamma$ solves the classical optimal transport problem from $f$ towards $g:=\gamma_{y}$ defined by (2.1). Moreover, we have the identity $\Upsilon(f)=\mathcal{T}_{c}(f, g)$.

Let us now show that $\Upsilon^{*}(f)=\Upsilon(f)$ and that (3.3) admits a maximising pair $(\varphi, \psi)$. We first establish that we can reduce the set of competitors for (3.3). To simplify the notation we denote by $\varphi_{-}^{c}$ the function $\left(\varphi^{c}\right)_{-}:=\varphi^{c} \wedge 0$.

Lemma 3.3. Assume that $X$ is a compact metric space and that $c \in C(X \times X, \mathbb{R})$. Then, there holds

$$
\begin{equation*}
\Upsilon^{*}(f)=\sup \left\{K_{f}\left(\psi^{\bar{c}}, \psi\right): \psi=\varphi_{-}^{c} \text { for some } \varphi \in \Phi^{\prime}\right\} \tag{3.4}
\end{equation*}
$$

where

$$
\begin{equation*}
\Phi^{\prime}:=\left\{\varphi \in C(X), \varphi=\left(\varphi_{-}^{c}\right)^{\bar{c}}, \max \varphi^{c} \geq 0\right\} . \tag{3.5}
\end{equation*}
$$

Proof.
Step 1. We can replace $\psi$ by $\varphi_{-}^{c}$ - and assume that $\max \varphi^{c} \geq 0$.
Let $(\varphi, \psi) \in \Phi$. By Proposition 2.4 (iii), $\psi \leq \varphi^{c}$, so that $\psi \leq \varphi^{c} \wedge 0=\varphi_{-}^{c}$. As $1-f \geq 0$, $K_{f}\left(\varphi, \varphi_{-}^{c}\right) \geq K_{f}(\varphi, \psi)$. Therefore, we can restrict the maximisation to the pairs $\left(\varphi, \varphi_{-}^{c}\right)$ in the supremum (3.3). Now, if $\max \varphi^{c}=-t<0$ we set $\widetilde{\varphi}:=\varphi-t$ so that $\widetilde{\varphi}^{c}=\varphi^{c}+t$. Consequently, $\max \widetilde{\varphi}^{c}=0$ and in particular, $\widetilde{\varphi}_{-}^{c}=\widetilde{\varphi}^{c}$ so that $\left(\widetilde{\varphi}, \widetilde{\varphi}^{c}\right) \in \Phi$. We then compute

$$
\begin{aligned}
K_{f}\left(\widetilde{\varphi}, \widetilde{\varphi}^{c}\right) & =\int f\left(\widetilde{\varphi}-\widetilde{\varphi}^{c}\right) d \lambda+\int \widetilde{\varphi}^{c} d \lambda \\
& =\int f\left(\varphi-\varphi^{c}\right) d \lambda+\int \varphi^{c} d \lambda+t(\lambda(X)-2 m) \\
& =K_{f}\left(\varphi, \varphi_{-}^{c}\right)+t(\lambda(X)-2 m) .
\end{aligned}
$$

As $2 m \leq \lambda(X)$ we obtain $K_{f}\left(\widetilde{\varphi}, \widetilde{\varphi}_{-}^{c}\right) \geq K_{f}\left(\varphi, \varphi_{-}^{c}\right)$. Hence

$$
\Upsilon^{*}(f)=\sup \left\{K_{f}\left(\varphi, \varphi_{-}^{c}\right): \varphi \in C(X), \max \varphi^{c} \geq 0\right\} .
$$

Step 2. There holds $\varphi=\left(\varphi_{-}^{c}\right)^{\bar{c}}$.
Let us introduce the mapping $P: C(X) \rightarrow C(X)$ defined by $P(\varphi):=\left(\varphi_{-}^{c}\right)^{\bar{c}}$. For $\varphi \in C(X)$, $\varphi^{c} \geq \varphi_{-}^{c}$, so that $P(\varphi)=\left(\varphi_{-}^{c}\right)^{\bar{c}} \geq \varphi^{c \bar{c}}$. By Proposition 2.4 (ii), $\varphi^{c \bar{c}} \geq \varphi$, hence

$$
\begin{equation*}
P(\varphi) \geq \varphi \tag{3.6}
\end{equation*}
$$

By Proposition 2.4 (ii) again there holds $P(\varphi)^{c}=\left(\varphi_{-}^{c}\right)^{\bar{c} c} \geq \varphi_{-}^{c}$. Taking the negative part yields

$$
\begin{equation*}
P(\varphi)_{-}^{c} \geq \varphi_{-}^{c} . \tag{3.7}
\end{equation*}
$$

We deduce from (3.6) and (3.7) that

$$
K_{f}\left(P(\varphi), P(\varphi)_{-}^{c}\right) \geq K_{f}\left(\varphi, \varphi_{-}^{c}\right) .
$$

Now, we observe that if $\max \varphi^{c} \geq 0$ we also have $\max \varphi^{c}{ }_{-}=0$ and, by (3.7), $\max P(\varphi)^{c}{ }_{-}=0$ which implies that $\max P(\varphi)^{c} \geq 0$. Hence,

$$
\begin{equation*}
\Upsilon^{*}(f)=\sup \left\{K_{f}\left(\widetilde{\varphi}, \widetilde{\varphi}_{-}^{c}\right): \widetilde{\varphi} \in C(X), \max \widetilde{\varphi}^{c} \geq 0, \widetilde{\varphi}=P(\varphi) \text { for some } \varphi \in C(X)\right\} . \tag{3.8}
\end{equation*}
$$

To conclude, we show that $P(P(\varphi))=P(\varphi)$ for any $\varphi \in C(X)$. By (3.6), $P(P(\varphi)) \geq P(\varphi)$. Taking the $\bar{c}$-transform in (3.7) yields $P(P(\varphi)) \leq P(\varphi)$ and we have indeed $P(P(\varphi))=P(\varphi)$. Hence we have $\tilde{\varphi} \in \Phi^{\prime}$ in (3.8) and we get

$$
\begin{equation*}
\Upsilon^{*}(f)=\sup \left\{K_{f}\left(\tilde{\varphi}, \tilde{\varphi}_{-}^{c}\right): \tilde{\varphi} \in \Phi^{\prime}\right\} . \tag{3.9}
\end{equation*}
$$

Finally, by definition $\tilde{\varphi}=\left(\tilde{\varphi}_{-}^{c}\right)^{\bar{c}}$ for $\tilde{\varphi} \in \Phi^{\prime}$ and (3.4) follows from (3.9) by letting $\psi:=\tilde{\varphi}_{-}^{c}$.
We can now establish that the supremum in (3.4) is reached.
Proposition 3.4. Assume that $X$ is a compact metric space and that $c \in C(X \times X, \mathbb{R})$. Then, the set $\Phi^{\prime}$ is compact in $\left(C(X),\|\cdot\|_{\infty}\right)$ and the suprema in (3.4) and (3.3) are attained.

Proof.
Let us show that $\Phi^{\prime}$ is compact. Let $\varphi_{n}$ be a sequence in $\Phi^{\prime}$. The function $c$ is $\omega$-continuous for some modulus of continuity $\omega \in C\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)$, so that by Proposition 2.4 (ii) for every $n \geq 0$, $\varphi_{n}^{c}:=\left(\varphi_{n}\right)^{c}$ and $\varphi_{n-}^{c}:=\left(\left(\varphi_{n}\right)^{c}\right)_{-}$are $\omega$-continuous. By definition of $\Phi^{\prime}, \varphi_{n}=\left(\varphi_{n-}^{c}\right)^{\bar{c}}$, so that $\varphi_{n}$ is also $\omega$-continuous for every $n \geq 0$. Let us show that the sequences $\varphi_{n}$ and $\varphi_{n-}^{c}$ are uniformly
bounded in $\left(C(X),\|\cdot\|_{\infty}\right)$. We observe that for every $n \geq 0, \max \varphi_{n}^{c} \geq 0$. Denoting by $x_{n}$ a point of $X$ such that $\varphi_{n-}^{c}\left(x_{n}\right)=0$, by $\omega$-continuity we have for $x \in X$ and $n \geq 0$,

$$
-\omega(\operatorname{diam}(X)) \leq-\omega\left(\left|x-x_{n}\right|\right) \leq \varphi_{n-}^{c}(x)-\varphi_{n-}^{c}\left(x_{n}\right)=\varphi_{n-}^{c}(x) \leq 0
$$

Thus the sequence $\varphi_{n-}^{c}$ is uniformly bounded in $\left(C(X),\|\cdot\|_{\infty}\right)$. By definition of the $c$-transform

$$
\min _{X \times X} c-\max _{X} \varphi_{n-}^{c} \leq\left(\varphi_{n-}^{c}\right)^{\bar{c}} \leq \max _{X \times X} c-\min _{X} \varphi_{n-}^{c}
$$

Hence the sequence $\varphi_{n}$ is also uniformly bounded. By Arzelá-Ascoli's theorem, there exists a pair $(\varphi, \psi) \in C(X) \times C(X)$ such that, up to extraction, $\left(\varphi_{n}, \varphi_{n-}^{c}\right)$ converges uniformly to $(\varphi, \psi)$.

Let us show that $\varphi \in \Phi^{\prime}$. By Proposition 2.4 (iii) and by uniform convergence $\varphi_{n}^{c} \rightarrow \varphi^{c}$ as $n \rightarrow \infty$ so that

$$
\begin{equation*}
\varphi_{n-}^{c} \rightarrow \varphi_{-}^{c} \quad \text { uniformly as } n \rightarrow \infty \tag{3.10}
\end{equation*}
$$

which yields $\psi=\varphi_{-}^{c}$. From (3.10) and the uniform continuity of $c$, we deduce that

$$
\left(\varphi_{n-}^{c}\right)^{\bar{c}}=\varphi_{n} \rightarrow\left(\varphi_{-}^{c}\right)^{\bar{c}} \quad \text { uniformly as } n \rightarrow \infty
$$

Since $\varphi_{n} \rightarrow \varphi$ as $n \rightarrow \infty$, we obtain $\varphi=\left(\varphi_{-}^{c}\right)^{\bar{c}}$. Lastly, by uniform convergence, the fact that $\max \varphi_{n}^{c} \geq 0$ for all $n \geq 0$ implies that $\max \varphi^{c} \geq 0$, so that $\varphi \in \Phi^{\prime}$. This shows that $\Phi^{\prime}$ is a compact subset of $\left(C(X),\|\cdot\|_{\infty}\right)$.

Let now $\psi_{n}$ be a maximising sequence for (3.4). For all $n \geq 0$, there exists $\varphi_{n} \in \Phi^{\prime}$ such that $\psi_{n}=\varphi_{n-}^{c}$. By compactness of $\Phi^{\prime}, \varphi_{n} \rightarrow \varphi$ as $n \rightarrow \infty$ for some $\varphi \in \Phi^{\prime}$. Setting $\psi=\varphi_{-}^{c}$, we have $\psi_{n} \rightarrow \psi$ and $\psi_{n}^{\bar{c}} \rightarrow \psi^{\bar{c}}$ as $n \rightarrow \infty$. The functional $K_{f}$ being continuous with respect to uniform convergence, we obtain

$$
K_{f}\left(\psi^{\bar{c}}, \psi\right)=\lim K_{f}\left(\psi_{n}^{\bar{c}}, \psi_{n}\right)=\Upsilon^{*}(f)
$$

This proves that $\psi$ is a maximiser for (3.4) and by Lemma 3.3, $\psi$ also maximises (3.3).
We are now ready to prove that there is no duality gap between (3.1) and (3.3). The proof is an adaptation of [22, Section 1.6.3].

Proposition 3.5. Assume that $X$ is a compact metric space and that $c \in C(X \times X, \mathbb{R})$. Then,

$$
\Upsilon^{*}(f)=\Upsilon(f)
$$

Proof.
Step 1. Definition of $H$ and first properties.
For $p \in C(X \times X)$, we define

$$
H(p):=-\sup \left\{\int(f \varphi+(1-f) \psi) d \lambda:(\varphi, \psi) \in \Phi_{p}\right\}
$$

where

$$
\Phi_{p}:=\{(\varphi, \psi) \in C(X) \times C(X), \psi \leq 0, \varphi \oplus \psi \leq c-p\}
$$

We first observe that $c-p$ is continuous and bounded from below. Thus, by applying Proposition 3.4 with $c-p$ in place of $c$, we see that the above supremum is a maximum.

Let us now show that $H$ is convex. Let $p_{0}, p_{1} \in C(X \times X)$ and $\theta \in[0,1]$ and let us set $p:=(1-\theta) p_{0}+\theta p_{1}$. We denote by $\left(\varphi_{0}, \psi_{0}\right)$ and $\left(\varphi_{1}, \psi_{1}\right)$ two maximising pairs associated with $p_{0}$ and $p_{1}$ and set $\varphi:=(1-\theta) \varphi_{0}+\theta \varphi_{1}, \psi:=(1-\theta) \psi_{0}+\theta \psi_{1}$. We see that $(\varphi, \psi)$ is an admissible pair $(\psi \leq 0$ and $\varphi \oplus \psi \leq c-p)$, so that

$$
H(p) \leq-\int(f \varphi+(1-f) \psi) d \lambda=(1-\theta) H\left(p_{0}\right)+\theta H\left(p_{1}\right)
$$

This proves that $H$ is convex.
Next, we establish that $H$ is lower semi-continuous in $\left(C(X \times X),\|\cdot\|_{\infty}\right)$. Let $p_{n}$ and $p$ be elements of $C(X \times X)$ such that $p_{n} \rightarrow p$ uniformly as $n \rightarrow \infty$. The sequence $c-p_{n}$ is uniformly
equi-continuous. Therefore, proceeding as in the proof of Proposition 3.4, there exists a sequence of uniformly bounded and equi-continuous admissible pairs $\left(\varphi_{n}, \psi_{n}\right)$ such that

$$
H\left(p_{n}\right)=-\int\left(f \varphi_{n}+(1-f) \psi_{n}\right) d \lambda \quad \text { for every } n \geq 0
$$

We first extract a subsequence $p_{n^{\prime}}$ such that $\lim _{n^{\prime}} H\left(p_{n^{\prime}}\right)=\liminf _{n} H\left(p_{n}\right)$. By Arzelà-Ascoli's theorem, there exists $(\varphi, \psi) \in C(X) \times C(X)$ such that $\varphi_{n^{\prime}} \rightarrow \varphi$ and $\psi_{n^{\prime}} \rightarrow \psi$ uniformly as $n^{\prime} \rightarrow \infty$. By pointwise convergence, $\psi \leq 0$ and $\varphi \oplus \psi \leq c-p$. Passing to the limit yields

$$
H(p) \leq-\int(f \varphi+(1-f) \psi) d \lambda=-\lim _{n^{\prime}} \int\left(f \varphi_{n^{\prime}}+(1-f) \psi_{n^{\prime}}\right) d \lambda=\liminf _{n} H\left(p_{n}\right)
$$

Thus $H$ is lower semi-continuous.

## Step 2. Absence of duality gap.

Since $H$ is convex and lower semi-continuous on the Banach space $\left(C(X \times X),\|\cdot\|_{\infty}\right)$, there holds $H=H^{* *}$ on $C(X \times X)$. Here, for a Banach space $\mathcal{X}$ and a function $F: \mathcal{X} \rightarrow \mathbb{R} \cup\{+\infty\}$, $F^{*}$ denotes the Legendre transform of $F$ defined on the topological dual $\mathcal{X}^{*}$ of $\mathcal{X}$ by

$$
F^{*}\left(x^{*}\right):=\sup \left\{x^{*}(x)-F(x): x \in \mathcal{X}\right\}
$$

In particular,

$$
\begin{equation*}
\Upsilon^{*}(f)=-H(0)=-H^{* *}(0)=\inf \left\{H^{*}(\gamma): \gamma \in \mathcal{M}(X \times X)\right\} \tag{3.11}
\end{equation*}
$$

We now compute $H^{*}$. Let $\gamma \in \mathcal{M}(X \times X)$. By definition,

$$
H^{*}(\gamma)=\sup _{p \in C(X \times X)}\left\{\int p d \gamma+\sup _{(\varphi, \psi) \in \Phi_{p}}\left\{\int(f \varphi+(1-f) \psi) d \lambda\right\}\right\}
$$

Let us first assume that there exists $q \in C\left(X \times X, \mathbb{R}_{+}\right)$such that $t:=-\int q d \gamma>0$. We set $\varphi=\min c, \psi=0$ and $p_{n}:=-n q$ for $n \geq 1$. We obtain

$$
H^{*}(\gamma) \geq n t-m \min c \rightarrow \infty \quad \text { as } n \rightarrow \infty
$$

Thus, when computing $H^{*}(\gamma)$, we may assume that $\gamma \geq 0$. We rewrite $H^{*}(\gamma)$ as

$$
\begin{align*}
H^{*}(\gamma)=\int c d \gamma+\sup _{p \in C(X \times X)} \sup _{(\varphi, \psi) \in \Phi_{p}} & \left\{\int(p-c+\varphi \oplus \psi) d \gamma\right.  \tag{3.12}\\
& \left.\left.+\int \varphi d\left(f \lambda-\gamma_{x}\right)+\int \psi d\left((1-f) \lambda-\gamma_{y}\right)\right)\right\}
\end{align*}
$$

Let us set

$$
G(\gamma):=\sup \left\{\int \varphi d\left(f \lambda-\gamma_{x}\right)+\int \psi d\left((1-f) \lambda-\gamma_{y}\right):(\varphi, \psi) \in C(X) \times C(X), \psi \leq 0\right\}
$$

On the one hand, given $(\varphi, \psi) \in \Phi_{p}$ and $\gamma \geq 0$,

$$
\int(p-c+\varphi \oplus \psi) d \gamma \leq 0
$$

Therefore, $H^{*}(\gamma) \leq \int c d \gamma+G(\gamma)$. On the other hand, given $(\varphi, \psi)$ admissible for $G(\gamma)$, setting $p=c-\varphi \oplus \psi$ yields the converse inequality thanks to (3.12). Hence

$$
\begin{equation*}
H^{*}(\gamma)=\int c d \gamma+G(\gamma) \tag{3.13}
\end{equation*}
$$

Given $\gamma \in \mathcal{M}_{+}(X \times X)$, we have $G(\gamma)=0$ if $\gamma \in \Pi_{f}$ and $G(\gamma)=+\infty$ otherwise. Combining this with (3.13), we obtain that for $\gamma \in \mathcal{M}(X \times X)$,

$$
H^{*}(\gamma)= \begin{cases}\int c d \gamma & \text { if } \gamma \in \Pi_{f} \\ +\infty & \text { in the other cases. }\end{cases}
$$

Taking the infimum with respect to $\gamma \in \mathcal{M}(X \times X)$ and recalling (3.11), we get

$$
\Upsilon^{*}(f)=\inf \left\{H^{*}(\gamma): \gamma \in \mathcal{M}(X \times X)\right\}=\inf \left\{\int c d \gamma: \gamma \in \Pi_{f}\right\}=\Upsilon(f)
$$

which concludes the proof.
Remark 3.6. There is still no duality gap between (3.1) and (3.3) if we only assume $c$ to be lower semi-continuous. This result can be obtained by approximating $c$ pointwise from below by a non-decreasing sequence of continuous functions.

In the remainder of the section, we focus on the properties of the potentials $\left(\psi^{\bar{c}}, \psi\right)$ maximising (3.4). We first show that the sign of $\psi^{\bar{c} c}$ enforces constraints on the local values of the marginals of any plan $\gamma$ optimal for $\Upsilon(f)$.

Proposition 3.7. Assume that $X$ is a compact metric space and that $c \in C(X \times X, \mathbb{R})$. Let $\left(\psi^{\bar{c}}, \psi\right)$ be a maximiser of (3.3) and let $\gamma$ be a minimiser of (3.1). We set $g:=\gamma_{y}$. Then, $\gamma$ is a minimiser of (2.1) with $(\mu, \nu)=(f, g)$ and $\left(\psi^{\bar{c}}, \psi^{\bar{c} c}\right)$ is a pair of Kantorovitch potentials realising the maximum in (2.2). Moreover, up to $\lambda$-negligible sets,

$$
\begin{equation*}
f+g \equiv 1 \quad \text { on }\left\{\psi^{\bar{c} c}<0\right\} \quad \text { and } \quad g \equiv 0 \quad \text { on }\left\{\psi^{\bar{c} c}>0\right\} . \tag{3.14}
\end{equation*}
$$

Proof.
By Remark 3.2, $\gamma$ realises the minimum in (2.1) and $\Upsilon(f)=\mathcal{T}_{c}(f, g)$. As there is no duality gap in (3.1) nor in (2.1), $\Upsilon^{*}(f)=\mathcal{T}_{c}^{*}(f, g)$. Additionally, $\left(\psi^{\bar{c}}, \psi^{\bar{c} c}\right)$ is admissible for $\mathcal{T}_{c}^{*}(f, g)$ and $\psi=\left(\psi^{\bar{c} c}\right)_{-}$. Thus

$$
\begin{equation*}
K_{f}\left(\psi^{c}, \psi\right)=\int f \psi^{\bar{c}} d \lambda+\int(1-f)\left(\psi^{\bar{c} c}\right)_{-} d \lambda=\mathcal{T}_{c}^{*}(f, g) \geq \int f \psi^{\bar{c}} d \lambda+\int g \psi^{\bar{c} c} d \lambda \tag{3.15}
\end{equation*}
$$

Hence

$$
\int(1-f-g)\left(\psi^{\bar{c} c}\right)_{-} d \lambda \geq \int g\left(\psi^{\bar{c} c}\right)_{+} d \lambda
$$

Since $(1-f-g)\left(\psi^{\bar{c} c}\right)_{-} \leq 0$ and $g\left(\psi^{\bar{c} c}\right)_{+} \geq 0$, the integrands must vanish $\lambda$-almost everywhere: we deduce (3.14). Additionally, the inequality in (3.15) is an equality. Consequently, $\left(\psi^{\bar{c}}, \psi^{\bar{c} c}\right)$ is a pair of Kantorovitch potentials for (2.2).

To end this section, we establish a comparison principle on the potentials maximising (3.4). We say that a set $\Psi \subset C(X)$ admits a minimal (respectively maximal) element for the relation $\leq$ if there exists $\psi_{0} \in \Psi$ such that for any $\psi \in \Psi, \psi_{0} \leq \psi\left(\right.$ respectively $\left.\psi_{0} \geq \psi\right)$.

Proposition 3.8. Assume that $X$ is a compact metric space and that $c \in C(X \times X, \mathbb{R})$. Let $f \in L_{m}^{1}$ and let us define

$$
\Psi_{f}:=\left\{\psi, \psi=\varphi_{-}^{c} \text { for some } \varphi \in \Phi^{\prime} \text { and } K_{f}\left(\psi^{\bar{c}}, \psi\right)=\Upsilon^{*}(f)\right\}
$$

where $K_{f}$ is defined in (3.2) and $\Phi^{\prime}$ in Lemma 3.3.
Then:
(i) $\Psi_{f}$ admits a maximal element for the relation $\leq$, denoted by $\psi_{f}$ in the sequel,
(ii) For $f_{1}, f_{2} \in L_{m}^{1}$, there holds $f_{1} \leq f_{2} \Longrightarrow \psi_{f_{1}} \geq \psi_{f_{2}}$.

Proof.
Step 1. Sufficient condition and preliminary claim.
Notice that $\Psi_{f}$ is not empty by Proposition 3.4. To obtain $(i)$, we prove that the set

$$
\Phi_{f}:=\left\{\varphi \in \Phi^{\prime}, K_{f}\left(\varphi, \varphi_{-}^{c}\right)=\Upsilon^{*}(f)\right\}
$$

admits a minimal element $\varphi_{f}$ and then $\psi_{f}:=\left(\varphi_{f}^{c}\right)_{-}$is the desired maximal element of $\Psi_{f}$. Let us make a preliminary observation.

Claim. Let $f_{1}, f_{2} \in L_{m}^{1}$ with $f_{1} \leq f_{2}$ and set $\varphi_{i} \in \Phi_{f_{i}}$ for $i \in\{1,2\}$. Then $\varphi_{\wedge}:=\varphi_{1} \wedge \varphi_{2} \in \Phi_{f_{1}}$.

Let us first prove that $\varphi_{\wedge} \in \Phi^{\prime}$. In the sequel we write $\varphi_{i}^{c}:=\left(\varphi_{i}\right)^{c}$ and $\varphi_{i-}^{c}:=\left(\left(\varphi_{i}\right)^{c}\right)_{-}$for $i \in\{1,2, \wedge\}$. We observe that $\varphi_{\wedge} \in C(X)$. By definition of the $c$-transform, we obtain

$$
\begin{equation*}
\varphi_{\wedge}^{c}=\left(\varphi_{1} \wedge \varphi_{2}\right)^{c} \geq \varphi_{1}^{c} \vee \varphi_{2}^{c} \tag{3.16}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(\varphi_{1} \vee \varphi_{2}\right)^{c}=\varphi_{1}^{c} \wedge \varphi_{2}^{c} . \tag{3.17}
\end{equation*}
$$

Since for $i \in\{1,2\}, \max \varphi_{i}^{c} \geq 0$ we have by (3.16) that $\max \varphi_{\wedge}^{c} \geq 0$.
We now prove that $\left(\varphi_{\wedge_{-}}^{c}\right)^{\bar{c}}=\varphi_{\wedge}$. We observe that $\varphi_{\wedge_{-}}^{c} \leq \varphi_{\wedge}^{c}$ which implies $\left(\varphi_{\wedge-}^{c}\right)^{\bar{c}} \geq \varphi_{\wedge}^{c \bar{c}}$. By Proposition 2.4 (ii), $\varphi_{\wedge}^{c \bar{c}} \geq \varphi_{\wedge}$ so that $\left(\varphi_{\wedge-}^{c}\right)^{\bar{c}} \geq \varphi_{\wedge}$. Conversely, taking the negative part of (3.16), we have $\varphi_{\wedge_{-}}^{c} \geq\left(\varphi_{1}^{c} \vee \varphi_{2}^{c}\right)_{-}=\varphi_{1-}^{c} \vee \varphi_{2-}^{c}$. Taking the $\bar{c}$-transform and using (3.17) yields

$$
\left(\varphi_{\wedge-}^{c}\right)^{\bar{c}} \leq\left(\varphi_{1-}^{c} \vee \varphi_{2-}^{c}\right)^{\bar{c}}=\left(\varphi_{1-}^{c}\right)^{\bar{c}} \wedge\left(\varphi_{2-}^{c}\right)^{\bar{c}}=\varphi_{1} \wedge \varphi_{2}=\varphi_{\wedge} .
$$

Hence $\left(\varphi_{\wedge-}^{c}\right)^{\bar{c}}=\varphi_{\wedge}$ and $\varphi_{\wedge} \in \Phi^{\prime}$.
We now show that the pair $\left(\varphi_{\wedge}, \varphi_{\wedge-}^{c}\right)$ maximises $\Upsilon^{*}\left(f_{1}\right)$. We set

$$
\Delta_{K}:=K_{f_{1}}\left(\varphi_{\wedge}, \varphi_{\wedge-}^{c}\right)-K_{f_{1}}\left(\varphi_{1}, \varphi_{1-}^{c}\right)=\int f_{1}\left(\varphi_{\wedge}-\varphi_{1}\right)+\int\left(1-f_{1}\right)\left(\varphi_{\wedge-}^{c}-\varphi_{1-}^{c}\right)
$$

By optimality of $\varphi_{1}, \Delta_{K} \leq 0$. Let us prove the converse inequality. Substituting $f_{1}=f_{2}+f_{1}-f_{2}$ in the definition of $\Delta_{K}$, we obtain

$$
\Delta_{K}=\int f_{2}\left(\varphi_{\wedge}-\varphi_{1}\right)+\int\left(1-f_{2}\right)\left(\varphi_{\wedge-}^{c}-\varphi_{1-}^{c}\right)+\int\left(f_{2}-f_{1}\right)\left(\varphi_{1}-\varphi_{\wedge}+\varphi_{\wedge-}^{c}-\varphi_{1-}^{c}\right)
$$

We have $f_{2}-f_{1} \geq 0$ and $\varphi_{1}-\varphi_{\wedge} \geq 0$. Additionally, $\varphi_{\wedge}^{c} \geq \varphi_{1}^{c}$, so that $\varphi_{\wedge_{-}}^{c} \geq \varphi_{1-}^{c}$. Thus the last integral in $\Delta_{K}$ is non-negative. Adding and subtracting $f_{2} \varphi_{2}$ in the first integral yields

$$
\begin{equation*}
\Delta_{K} \geq \int f_{2} \varphi_{2}+\int f_{2}\left(\varphi_{\wedge}-\varphi_{1}-\varphi_{2}\right)+\int\left(1-f_{2}\right)\left(\varphi_{\wedge-}^{c}-\varphi_{1-}^{c}\right) \tag{3.18}
\end{equation*}
$$

Let us set $\varphi_{\vee}:=\varphi_{1} \vee \varphi_{2}$. By optimality of $\varphi_{2}$, we have $K_{f_{2}}\left(\varphi_{2}, \varphi_{2-}^{c}\right) \geq K_{f_{2}}\left(\varphi_{\vee}, \varphi_{\vee-}^{c}\right)$, which rewrites as

$$
\int f_{2} \varphi_{2} \geq \int f_{2} \varphi_{\vee}+\int\left(1-f_{2}\right)\left(\varphi_{\vee-}^{c}-\varphi_{2-}^{c}\right)
$$

Injecting this inequality in the first term of the right-hand side of (3.18) yields

$$
\begin{equation*}
\Delta_{K} \geq \int f_{2}\left(\varphi_{\wedge}+\varphi_{\vee}-\varphi_{1}-\varphi_{2}\right)+\int\left(1-f_{2}\right)\left(\varphi_{\vee-}^{c}+\varphi_{\wedge-}^{c}-\varphi_{1-}^{c}-\varphi_{2-}^{c}\right) \tag{3.19}
\end{equation*}
$$

The integrand in the first integral of (3.19) vanishes. Regarding the second term, using (3.17) and (3.16) we obtain

$$
\varphi_{\vee-}^{c}+\varphi_{\wedge-}^{c} \geq \varphi_{1-}^{c} \wedge \varphi_{2-}^{c}+\varphi_{1-}^{c} \vee \varphi_{2-}^{c}=\varphi_{1-}^{c}+\varphi_{2-}^{c}
$$

Hence the integrand in the second integral is non-negative. We conclude that $\Delta_{K} \geq 0$ and finally that $\Delta_{K}=0$ so that the claim is proved.

Step 2. Construction of the minimal element of $\Phi_{f}$.
By Lemma 3.4, $\Phi^{\prime}$ is compact. As $\varphi \mapsto K_{f}\left(\varphi, \varphi_{-}^{c}\right)$ is continuous for the norm of uniform convergence, $\Phi_{f}$ is compact as well. Let $\left(\varphi_{j}\right)_{j \geq 0}$ be a dense subset of $\Phi_{f}$. For $x \in X$ and $j \geq 0$, we define $\widetilde{\varphi}_{j}$ and $\varphi_{f}$ by

$$
\widetilde{\varphi}_{j}(x):=\min \left(\varphi_{0}(x), \ldots, \varphi_{j}(x)\right) \quad \text { and } \quad \varphi_{f}(x):=\inf \left\{\varphi(x), \varphi \in \Phi_{f}\right\}
$$

Using our preliminary claim with $f_{1}=f_{2}=f$ recursively, we obtain that for any $j \geq 0, \widetilde{\varphi}_{j} \in \Phi_{f}$. As $\Phi_{f}$ is compact and $\widetilde{\varphi}_{j} \rightarrow \varphi_{f}$ pointwise, we obtain that $\widetilde{\varphi}_{j} \rightarrow \varphi_{f}$ uniformly and $\varphi_{f} \in \Phi_{f}$, so that $\varphi_{f}$ is the desired minimal element of $\Phi_{f}$.

## Step 3. Conclusion.

Taking $\psi_{f}:=\varphi_{f_{-}}^{c}$ proves $(i)$. Let $f_{1} \leq f_{2}$ as given in the statement of (ii). By the previous step, there exist $\varphi_{1}, \varphi_{2}$ respective minimal elements for $\Phi_{f_{1}}$ and $\Phi_{f_{2}}$ such that $\psi_{1}:=\varphi_{1-}^{c}$
and $\psi_{2}:=\varphi_{2-}^{c}$ are respective maximal elements for $\Psi_{f_{1}}$ and $\Psi_{f_{2}}$. By the preliminary claim, $\varphi_{1} \wedge \varphi_{2} \in \Phi_{f_{1}}$ and by minimality of $\varphi_{1}$ we have $\varphi_{1} \leq \varphi_{1} \wedge \varphi_{2}$, so that $\varphi_{2} \geq \varphi_{1}$. Hence $\psi_{2} \leq \psi_{1}$.

## 4. Existence of maximisers of (4.2) for translation invariant costs in $\mathbb{R}^{d}$

We now assume that $X=\mathbb{R}^{d}$, that $\lambda$ is the Lebesgue measure and that $c(x, y)=k(y-x)$, with $k: \mathbb{R}^{d} \rightarrow \mathbb{R}_{+}$. We recall the following hypotheses on $k$.
(H1) $k \in C\left(\mathbb{R}^{d}, \mathbb{R}_{+}\right), k(0)=0$ and $k(x) \rightarrow \infty$ as $|x| \rightarrow \infty$,
(H2) $\forall x \neq 0$,

$$
\limsup _{r \rightarrow 0} \frac{1}{r^{d}}\left|B_{r}(x) \cap\left\{y \in \mathbb{R}^{d}, k(y)<k(x)\right\}\right|>0
$$

(H3) $\forall \sigma \in \mathbb{S}^{d-1}, r \mapsto k(r \sigma)$ is increasing on $\mathbb{R}_{+}$.
Notice that under hypotheses (H1) \& (H2), there holds $k(x)>0$ for $x \neq 0$.
The primal problem is now defined as

$$
\begin{equation*}
\Upsilon(f):=\inf \left\{\int c d \gamma: \gamma \in \Pi_{f}\right\} \tag{4.1}
\end{equation*}
$$

where

$$
\Pi_{f}:=\left\{\gamma \in \mathcal{M}_{+}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right): \gamma_{x}=f, \gamma_{y} \leq 1-f\right\}
$$

The goal of this section is to prove that for every $m>0$ the energy

$$
\begin{equation*}
\mathcal{E}(m):=\sup \left\{\Upsilon(f): f \in L^{1}\left(\mathbb{R}^{d},[0,1]\right), \int f=m\right\} \tag{4.2}
\end{equation*}
$$

admits a maximiser.

### 4.1. First properties of $\Upsilon$ and saturation theorem.

In this subsection, we collect some properties of the functional $\Upsilon$ defined in $\mathbb{R}^{d}$ and establish a saturation property (Theorem 4.4), namely that if $\gamma$ is a minimiser for $\Upsilon(f)$ then $\gamma_{y}(x) \in$ $\{f(x), 1-f(x)\}$ for almost every $x \in \mathbb{R}^{d}$.

We start by proving that minimisers of (4.1) exist. The proof of this result is similar to the proof of [4, Proposition 2.1], but with weaker assumptions on the cost $c$ and in the context of functions taking values in $[0,1]$ rather than in $\{0,1\}$.

Proposition 4.1. Assume that $k$ satisfies (H1). Then, for any $m>0$ and $f \in L_{m}^{1}$, the infimum in (4.1) is attained. Additionally, given any minimiser $\gamma$ of (4.1) we have $\Upsilon(f)=\mathcal{T}_{c}(f, g)$, where $g:=\gamma_{y}$.

Lastly, there exists $R=R(m)$ non-decreasing in $m$ such that for any $f \in L_{m}^{1}$,

$$
\begin{equation*}
\Upsilon(f)=\min \left\{\int c d \gamma: \gamma \in \Pi_{f}, \forall(x, y) \in \operatorname{supp} \gamma,|x-y| \leq R\right\} \tag{4.3}
\end{equation*}
$$

and for any minimiser $\gamma$ of (4.1), there holds $|x-y| \leq R$ on supp $\gamma$.
Proof.
The strategy of the proof is to first establish (4.3) with an infimum in place of the minimum. Then we use this property to derive compactness for (4.1).
Step 1. Restricting the set of competitors for (4.1).
We let $\gamma \in \Pi_{f}$ and set $g:=\gamma_{y}$. We want to build a competitor $\widetilde{\gamma}$ for $\Upsilon(f)$ such that for some $R>0,|x-y| \leq R$ for every $(x, y) \in \operatorname{supp} \tilde{\gamma}$. For $R>0$, we define

$$
\Gamma_{R}:=\left\{(x, y) \in \mathbb{R}^{d} \times \mathbb{R}^{d},|x-y| \geq R\right\} .
$$

We consider a standard partition of $\mathbb{R}^{d}$ into cubes $\left(Q_{i}\right)_{i \geq 0}$ with side-length $\rho_{1}(m):=(3 m)^{1 / d}$. We define

$$
I:=\left\{i \geq 0, m_{i}:=\gamma\left(\Gamma_{R} \cap\left(Q_{i} \times \mathbb{R}^{d}\right)\right)>0\right\}
$$

and for $i \in I$, we set

$$
\gamma_{\mathrm{bad}, i}:=\chi_{\Gamma_{R} \cap\left(Q_{i} \times \mathbb{R}^{d}\right)} \gamma .
$$

As $\left|Q_{i}\right|-\int f-\int g \geq m \geq m_{i}$, there exists a positive measure $\mu_{i} \ll \chi_{Q_{i}} \lambda$ such that

$$
\mu_{i} \leq \chi_{Q_{i}}(1-f-g) \quad \text { and } \quad \mu_{i}\left(\mathbb{R}^{d}\right)=\gamma_{\mathrm{bad}, i}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)
$$

Denoting by $\theta_{i}$ the first marginal of $\gamma_{\text {bad }, i}$ we set

$$
\widetilde{\gamma}_{i}:=\chi_{Q_{i} \times \mathbb{R}^{d}} \gamma-\gamma_{\mathrm{bad}_{, i}}+\frac{1}{m_{i}} \theta_{i} \otimes \mu_{i} \geq 0 .
$$

For $i \in I^{c}$, we simply define $\widetilde{\gamma}_{i}:=\chi_{\left(Q_{i} \times \mathbb{R}^{d}\right)} \gamma$. As a consequence, $\widetilde{\gamma}:=\sum_{i \geq 0} \widetilde{\gamma}_{i}$ is a transport plan whose first marginal is $f$ and second marginal $\widetilde{g}$ verifies $\widetilde{g} \leq g \leq 1-\widetilde{f}$. By construction, for $R>\sqrt{d} \rho_{1}(m)$, we have $\widetilde{\gamma}\left(\Gamma_{R}\right)=0$.

Let us now compare the transportation cost of $\gamma$ and $\widetilde{\gamma}$. We compute:

$$
\begin{aligned}
\int c d \widetilde{\gamma}-\int c d \gamma & =\sum_{i \in I} \int_{Q_{i} \times \mathbb{R}^{d}} c d\left(\frac{\theta_{i} \otimes \mu_{i}}{m_{i}}-\gamma_{\mathrm{bad}, i}\right) \\
& \leq\left(\sum_{i \in I} m_{i}\right)\left(\max _{z \in \bar{Q}_{\rho_{1}(m)}} k(z)-\inf _{|z| \geq R} k(z)\right) .
\end{aligned}
$$

Let us set $M:=\max \left\{k(z), z \in \bar{Q}_{\rho_{1}(m)}\right\}$. By (H1), there exists $R>\sqrt{d} \rho_{1}(m)$ such that if $|z|>R$, then $k(z)>M$. With this choice of $R$ we have $\int c d \widetilde{\gamma} \leq \int c d \gamma$. Lastly, whenever $\gamma\left(\Gamma_{R}\right)>0$,

$$
\begin{equation*}
\int c d \widetilde{\gamma}<\int c d \gamma \tag{4.4}
\end{equation*}
$$

Step 2: Lower semi-continuity of the transportation cost.
This step is classical. To prove that $\gamma \mapsto \int c d \gamma$ is lower semi-continuous with respect to weak convergence, we proceed by approximation. Let us assume that $\gamma_{n} \stackrel{*}{\rightharpoonup} \gamma$ as $n \rightarrow \infty$. For $j \geq 0$, we define $c_{j}:=c \wedge j$. The sequence $c_{j}$ is non-decreasing and converges pointwise to $c$. For every $j \geq 0, c_{j} \in C_{b}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)$, so that

$$
\int c_{j} d \gamma=\lim _{n} \int c_{j} d \gamma_{n} \leq \liminf _{n} \int c d \gamma_{n}
$$

By the monotone convergence theorem,

$$
\int c d \gamma=\lim _{j} \int c_{j} d \gamma \leq \liminf _{n} \int c d \gamma_{n}
$$

which concludes the second step of the proof.
Step 3. $\Upsilon(f)$ admits a minimiser.
Let $\gamma_{n}$ be a minimising sequence for (4.1). Let us show that the sequence $\gamma_{n}$ is tight. By the first step, we can assume that there exists $R=R(m)$ such that for any $n \geq 0$ there holds $|x-y| \leq R$ on supp $\gamma_{n}$. Now, because $\int f \leq m<\infty$, there exists $R^{\prime}=R^{\prime}(m)>0$ such that $\int_{\mathbb{R}^{d} \backslash B_{R^{\prime}}} f \leq \varepsilon$. Hence,

$$
\gamma_{n}\left(\mathbb{R}^{d} \times \mathbb{R}^{d} \backslash\left(B_{R^{\prime}} \times B_{R+R^{\prime}}\right)\right)=\gamma_{n}\left(B_{R^{\prime}} \times\left(\mathbb{R}^{d} \backslash B_{R+R^{\prime}}\right)\right)+\gamma_{n}\left(\left(\mathbb{R}^{d} \backslash B_{R^{\prime}}\right) \times \mathbb{R}^{d}\right) \leq 0+\varepsilon
$$

which proves that the sequence $\gamma_{n}$ is tight. Together with the second step, this shows that (4.1) admits a minimiser. Moreover, by (4.4) for any minimiser $\gamma$ of (4.1) there holds $|x-y| \leq R$ on $\operatorname{supp} \gamma$. Lastly, setting $g:=\gamma_{y}$ the identity $\Upsilon(f)=\mathcal{T}_{c}(f, g)$ is immediate.

We now establish some basic properties of the functional $\Upsilon$. The results here are similar to [4, Proposition 2.2 \& Lemma 2.4].

Proposition 4.2. Assume that $k$ satisfies (H1). Given $m>0$ and $f_{1}, f_{2} \in L_{m}^{1}$ we have:
(i) If $f_{1}+f_{2} \leq 1$, then

$$
\Upsilon\left(f_{1}+f_{2}\right) \geq \Upsilon\left(f_{1}\right)+\Upsilon\left(f_{2}\right)
$$

As a consequence, if $f_{1} \leq f_{2}$, then $\Upsilon\left(f_{1}\right) \leq \Upsilon\left(f_{2}\right)$.
(ii) There exists $R=R(m)$ such that if $d\left(\operatorname{supp} f_{1}, \operatorname{supp} f_{2}\right) \geq R$, then

$$
\Upsilon\left(f_{1}+f_{2}\right)=\Upsilon\left(f_{1}\right)+\Upsilon\left(f_{2}\right) .
$$

(iii) There exists $C=C(m)>0$ such that

$$
\left|\Upsilon\left(f_{1}\right)-\Upsilon\left(f_{2}\right)\right| \leq C\left\|f_{1}-f_{2}\right\|_{L^{1}}
$$

(iv) Let $f, f_{n} \in L^{1}\left(\mathbb{R}^{d},[0,1]\right)$ be such that the sequence $f_{n} \lambda$ is tight and $f_{n} \stackrel{*}{\rightharpoonup} f$. Then $\Upsilon\left(f_{n}\right) \rightarrow$ $\Upsilon(f)$.

Proof.
Step 1. Proof of (i) $\mathcal{G}(i i)$.
To prove $(i)$, we consider a transport plan $\gamma$ optimal for $\Upsilon\left(f_{1}+f_{2}\right)$ whose existence is guaranteed by Proposition 4.1. We would like to extract from $\gamma$ two plans $\gamma^{1}$ and $\gamma^{2}$ admissible for $\Upsilon\left(f_{1}\right)$ and $\Upsilon\left(f_{2}\right)$ respectively. Using the convention $0 / 0=0$, we define $\gamma^{1}$ and $\gamma^{2}$ through

$$
d \gamma^{1}(x, y):=\frac{f_{1}(x)}{\left(f_{1}+f_{2}\right)(x)} d \gamma(x, y) \quad \text { and } \quad d \gamma^{2}(x, y):=\frac{f_{2}(x)}{\left(f_{1}+f_{2}\right)(x)} d \gamma(x, y)
$$

By construction, $\gamma_{x}^{1}=f_{1}$ and $\gamma_{x}^{2}=f_{2}$. We also have $\gamma^{1} \leq \gamma$, so that

$$
\gamma_{y}^{1} \leq \gamma_{y} \leq 1-\left(f_{1}+f_{2}\right) \leq 1-f_{1}
$$

Likewise, $\gamma_{y}^{2} \leq 1-f_{2}$. Therefore, $\gamma^{1}$ and $\gamma^{2}$ are admissible for $\Upsilon\left(f_{1}\right)$ and $\Upsilon\left(f_{2}\right)$ respectively. Moreover

$$
\Upsilon\left(f_{1}\right)+\Upsilon\left(f_{2}\right) \leq \int c d \gamma^{1}+\int c d \gamma^{2}=\int c d \gamma=\Upsilon\left(f_{1}+f_{2}\right)
$$

which is the desired conclusion.
To prove (ii), we consider transport plans $\gamma^{1}$ and $\gamma^{2}$ which are optimal for $\Upsilon\left(f_{1}\right)$ and $\Upsilon\left(f_{2}\right)$ respectively. We define $g_{1}:=\gamma_{y}^{1}$ and $g_{2}:=\gamma_{y}^{2}$. If we set $\gamma:=\gamma^{1}+\gamma^{2}$, we have $\gamma_{x}=f_{1}+f_{2}$ and $\gamma_{y}=g_{1}+g_{2}$. Moreover, by Proposition 4.1, if $d\left(\operatorname{supp} f_{1}, \operatorname{supp} f_{2}\right) \geq R$ for $R=R(m)$ large enough, then the supports of $g_{1}$ and $g_{2}$ are also disjoint. Consequently, $g_{1}+g_{2} \leq 1-\left(f_{1}+f_{2}\right)$, so that $\gamma$ is admissible for $\Upsilon\left(f_{1}+f_{2}\right)$ and we have the desired converse inequality

$$
\Upsilon\left(f_{1}+f_{2}\right) \leq \int c d\left(\gamma^{1}+\gamma^{2}\right) \leq \Upsilon\left(f_{1}\right)+\Upsilon\left(f_{2}\right)
$$

Step 2. Proof of (iii).
Exchanging the roles of $f_{1}$ and $f_{2}$, it is enough to prove the estimate

$$
\begin{equation*}
\Upsilon\left(f_{2}\right)-\Upsilon\left(f_{1}\right) \leq C\left\|f_{2}-f_{1}\right\|_{L^{1}} \tag{4.5}
\end{equation*}
$$

Let $\gamma^{1}$ be a minimiser of $\Upsilon\left(f_{1}\right)$ and let us set $g_{1}:=\gamma_{x}^{1}$. In the next substeps, we build from $\gamma^{1}$ an exterior transport plan $\gamma^{2}$ for $f_{2}$ with controlled cost.
Step 2.a. Transporting most of $f_{1} \wedge f_{2}$.
Using the convention $0 / 0=0$, we define a plan $\gamma^{\prime}$ by

$$
d \gamma^{\prime}(x, y):=\frac{\left(f_{1} \wedge f_{2}\right)(x)}{f_{1}(x)} d \gamma_{1}(x, y)
$$

We set $g^{\prime}:=\gamma_{y}^{\prime}$. Notice that $\gamma^{\prime} \leq \gamma^{1}$, which implies $g^{\prime} \leq g_{1}$. Additionally, $\gamma_{x}^{\prime}=f_{1} \wedge f_{2}$, so that

$$
\begin{equation*}
\gamma^{1}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)-\gamma^{\prime}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)=\int\left(f_{1}-f_{1} \wedge f_{2}\right)=\int\left(f_{1}-f_{2}\right)_{+} \tag{4.6}
\end{equation*}
$$

Heuristically, $\gamma^{\prime}$ corresponds to sending through $\gamma^{1}$ as much mass from $f_{2}$ as possible. However, we have to remove some of this mass because the constraint $g^{\prime} \leq 1-f_{2}$ might not hold true everywhere. Let

$$
u:=\left(f_{2}+g^{\prime}-1\right)_{+}
$$

and define $\gamma^{\prime \prime}$ as

$$
d \gamma^{\prime \prime}(x, y):=\frac{g^{\prime}(y)-u(y)}{g^{\prime}(y)} d \gamma^{\prime}(x, y)
$$

We set $f^{\prime \prime}:=\gamma_{x}^{\prime \prime}$ and $g^{\prime \prime}:=\gamma_{y}^{\prime \prime}$. By construction, $g^{\prime \prime}=g^{\prime}-u$ so $g^{\prime \prime} \leq 1-f_{2}$ as desired. Since $\gamma^{\prime \prime} \leq \gamma^{\prime}$, we also have $f^{\prime \prime} \leq f_{1} \wedge f_{2} \leq f_{2}$. Now since $g^{\prime} \leq g_{1} \leq 1-f_{1}$ we have $u \leq\left(f_{2}-f_{1}\right)_{+}$ from which we infer

$$
\gamma^{\prime}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)-\gamma^{\prime \prime}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)=\int\left(g^{\prime}-g^{\prime \prime}\right)=\int u \leq \int\left(f_{2}-f_{1}\right)_{+}
$$

Summing this and (4.6) yields

$$
\begin{equation*}
\gamma^{1}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)-\gamma^{\prime \prime}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right) \leq\left\|f_{2}-f_{1}\right\|_{L^{1}} \tag{4.7}
\end{equation*}
$$

Eventually since $\gamma^{\prime \prime} \leq \gamma^{\prime} \leq \gamma^{1}$ and $c \geq 0$ we have

$$
\begin{equation*}
\int c d \gamma^{\prime \prime}-\int c d \gamma^{1} \leq 0 \tag{4.8}
\end{equation*}
$$

Step 2.b. Final construction.
We are now ready to build an admissible transport plan $\gamma^{2}$ for $\Upsilon\left(f_{2}\right)$. Noticing that $f_{2}-f^{\prime \prime} \geq 0$ we write $f_{2}=f^{\prime \prime}+\left(f_{2}-f^{\prime \prime}\right)$. By (4.7) we have

$$
\begin{equation*}
\int\left(f_{2}-f^{\prime \prime}\right)=\int\left(f_{2}-f_{1}\right)+\int\left(f_{1}-f^{\prime \prime}\right) \leq 2\left\|f_{2}-f_{1}\right\|_{L^{1}} \tag{4.9}
\end{equation*}
$$

Arguing as in the proof of Proposition 4.1, we can find a function $0 \leq g^{\prime \prime \prime} \leq 1-f_{2}-g^{\prime \prime}$ with $\int g^{\prime \prime \prime} \leq \int\left(f_{2}-f^{\prime \prime}\right)$ and a transport plan $\gamma^{\prime \prime \prime}$ between $f_{2}-f^{\prime \prime}$ and $\overline{g^{\prime \prime \prime}}$ such that for some $C=C(m)>0$,

$$
\begin{equation*}
\int c d \gamma^{\prime \prime \prime} \leq C \int\left(f_{2}-f^{\prime \prime}\right) \stackrel{(4.9)}{\leq} C\left\|f_{2}-f_{1}\right\|_{L^{1}} \tag{4.10}
\end{equation*}
$$

Finally we define $\gamma^{2}:=\gamma^{\prime \prime}+\gamma^{\prime \prime \prime}$ which is admissible for $\Upsilon\left(f_{2}\right)$ by construction. Summing (4.8) and (4.10), we get

$$
\Upsilon\left(f_{2}\right) \leq \int c d \gamma^{2} \leq \int c d \gamma_{1}+C\left\|f_{2}-f_{1}\right\|_{L^{1}}=\Upsilon\left(f_{1}\right)+C\left\|f_{2}-f_{1}\right\|_{L^{1}}
$$

This proves (4.5) and thus point (iii).
Step 3. Proof of (iv).
Let $f_{n}$ and $f_{n}$ be as in the statement of the proposition. By weak convergence, we have $f_{n}, f \in L_{m}^{1}$ for some $m>0$. Using the Lipschitz continuity of $\Upsilon$ with respect to $L^{1}$ convergence, we may assume without loss of generality that $f_{n}$ (and thus also $f$ ) are supported in $\bar{B}_{R_{0}}$ for some $R_{0}>0$. Applying Proposition 4.1 we get that minimisers of $\Upsilon\left(f_{n}\right)$ and $\Upsilon(f)$ are supported in $\bar{B}_{R} \times \bar{B}_{R}$ for some $R>R_{0}>0$. We may thus restrict these problems to the compact set $\bar{B}_{R}$. Using Proposition 3.5 we have $\Upsilon\left(f_{n}\right)=\Upsilon^{*}\left(f_{n}\right)$ and it is thus enough to prove the continuity of $\Upsilon^{*}$ with respect to the weak-* topology.
By Proposition 3.4, for every $n \geq 0$ there exists a pair of potentials $\left(\varphi_{n}, \psi_{n}\right)$ maximising $\Upsilon^{*}\left(f_{n}\right)$. Since for every $n, \varphi_{n}$ belongs to $\Phi^{\prime}$ (where $\Phi^{\prime}$ is defined by (3.5)) and since this set is compact by Proposition 3.4 we have that a subsequence $\varphi_{n^{\prime}}$ of $\varphi_{n}$ converges in $C\left(\bar{B}_{R}\right)$ to some $\varphi \in \Phi^{\prime}$. Arguing as in the proof of Proposition 3.4 we see that $\psi_{n^{\prime}}$ also converges to $\psi$ with $(\varphi, \psi)$ admissible for $\Upsilon^{*}(f)$. By weak-strong convergence we then have

$$
\lim \sup \Upsilon^{*}\left(f_{n^{\prime}}\right)=\lim \sup \int f_{n^{\prime}} \varphi_{n^{\prime}}+\left(1-f_{n^{\prime}}\right) \psi_{n^{\prime}}=\int f \varphi+(1-f) \psi \leq \Upsilon^{*}(f)
$$

Similarly, if $(\varphi, \psi)$ are optimal potentials for $\Upsilon^{*}(f)$, they are admissible for $\Upsilon^{*}\left(f_{n}\right)$ and thus

$$
\liminf \Upsilon^{*}\left(f_{n}\right) \geq \liminf \int f_{n} \varphi+\left(1-f_{n}\right) \psi=\int f \varphi+(1-f) \psi=\Upsilon^{*}(f)
$$

We then have $\lim \Upsilon^{*}\left(f_{n^{\prime}}\right)=\Upsilon^{*}(f)$ and by uniqueness of the limit we see that the extraction was not necessary. This establishes (iv) and ends the proof of the proposition.

The next lemma and theorem state very important saturation properties satisfied by the optimal exterior transport plan. These results extend [7, Lemma 5.1 \& Proposition 5.2] to more general costs $c$.

Lemma 4.3. Assume that $k$ satisfies (H1) $\mathcal{G}(H 2)$. For $f \in L_{m}^{1}$ let $\gamma$ be optimal for $\Upsilon(f)$. Then for every $\left(x_{0}, y_{0}\right) \in \operatorname{supp} \gamma$ there holds $f+\gamma_{y} \equiv 1$ almost everywhere on the saturation set

$$
S\left(x_{0}, y_{0}\right):=\left\{y \in \mathbb{R}^{d}, k\left(y-x_{0}\right)<k\left(y_{0}-x_{0}\right)\right\}
$$

Proof.
In the proof we set $g:=\gamma_{y}$ and $h:=f+g$. Let $\left(x_{0}, y_{0}\right) \in \operatorname{supp} \gamma$ and assume without loss of generality that $x_{0}=0$. We suppose by contradiction that there exists $\varepsilon>0$ such that the set

$$
S_{\varepsilon}:=\{h<1\} \cap\left\{y \in \mathbb{R}^{d}, k(y)<k\left(y_{0}\right)-\varepsilon\right\}
$$

has positive Lebesgue measure. Notice that by (H1), $k(x) \rightarrow \infty$ as $|x| \rightarrow \infty$ so that $S_{\varepsilon}$ is bounded. Therefore,

$$
m_{\varepsilon}:=\int_{S_{\varepsilon}}(1-h) \in(0, \infty)
$$

We now exhibit an exterior transport plan $\bar{\gamma}$ whose transportation cost is strictly smaller than the one of $\gamma$. Given $r>0$, we define the measure $\gamma^{0}:=\gamma\left\llcorner\left(B_{r} \times B_{r}\left(y_{0}\right)\right)\right.$. As $\left(0, y_{0}\right) \in \operatorname{supp} \gamma$, for every $r>0$,

$$
\begin{equation*}
0<\gamma^{0}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right) \leq \int_{B_{r}} f \leq\left|B_{r}\right| \tag{4.11}
\end{equation*}
$$

Thus, by the last inequality in (4.11) there exists $r_{\varepsilon}>0$ such that for every $r \in\left(0, r_{\varepsilon}\right]$,

$$
\gamma^{0}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)=\alpha m_{\varepsilon}
$$

for some $0<\alpha \leq 1$. Let us fix $r \in\left(0, r_{\varepsilon}\right]$. We define a competitor $\widetilde{\gamma}$ for $\Upsilon(f)$ by setting $\widetilde{\gamma}:=\gamma-\gamma^{0}+\eta$, where

$$
\eta:=\gamma_{x}^{0} \otimes \frac{1-h}{m_{\varepsilon}} \chi_{S_{\varepsilon}}
$$

By construction, $\widetilde{\gamma}_{x}=\gamma_{x}=f$. We also have

$$
f+\widetilde{\gamma}_{y} \leq f+g+\alpha(1-h) \chi_{S_{\varepsilon}} \leq h+\alpha(1-h)=1-(1-h)(1-\alpha) \leq 1
$$

so that $\widetilde{\gamma}$ is admissible for $\Upsilon(f)$. We compute

$$
\int c d \widetilde{\gamma}-\int c d \gamma=\int c d \eta-\int c d \gamma^{0} \leq \alpha m_{\varepsilon}\left(\max _{\bar{B}_{r} \times \bar{S}_{\varepsilon}} c(x, y)-\min _{\bar{B}_{r} \times \bar{B}_{r}\left(y_{0}\right)} c(x, y)\right)
$$

By continuity of $c$ there exists $r_{\varepsilon}>0$ such that for $0<r \leq r_{\varepsilon}$,

$$
\max _{\bar{B}_{r} \times \bar{S}_{\varepsilon}} c(x, y) \leq k\left(y_{0}\right)-\varepsilon / 2 \quad \text { and } \quad \min _{\bar{B}_{r} \times \bar{B}_{r}\left(y_{0}\right)} c(x, y) \geq k\left(y_{0}\right)-\varepsilon / 4
$$

Thus for $0<r \leq r_{\varepsilon}$,

$$
\int c d \widetilde{\gamma}-\int c d \gamma \leq-\alpha \varepsilon m_{\varepsilon} / 4<0
$$

which contradicts the fact that $\gamma$ is a minimiser for $\Upsilon(f)$.

Theorem 4.4. Assume that $k$ satisfies (H1) $\mathcal{Z}(H 2)$. For $f \in L_{m}^{1}$, let $\gamma \in \Pi_{f}$ be a minimiser of (4.1) and set $g:=\gamma_{y}$. Then, defining

$$
E:=\{x: \exists y \neq x \text { such that }(x, y) \in \operatorname{supp} \gamma \text { or }(y, x) \in \operatorname{supp} \gamma\},
$$

the set $E$ is Lebesgue measurable and we have the identity $g=(1-f) \chi_{E}+f \chi_{E^{c}}$.
Proof.
Step 1. A preliminary claim.
We first prove the following. Let $\mu, \nu \in \mathcal{M}_{+}\left(\mathbb{R}^{d}\right)$ be such that $\mu\left(\mathbb{R}^{d}\right)=\nu\left(\mathbb{R}^{d}\right)$, and let $\gamma \in \Pi(\mu, \nu)$. Assume that $\mu$ and $\nu$ are absolutely continuous with respect to the Lebesgue measure. If we define the set

$$
\mathcal{A}(\gamma):=\{x: \exists y \neq x \text { such that }(x, y) \in \operatorname{supp} \gamma\},
$$

then $\mu \leq \nu$ on $\mathcal{A}(\gamma)^{c}$.
To prove the claim, let us first show that $\mathcal{A}(\gamma)$ is Lebesgue measurable. We define

$$
\mathcal{D}(\gamma):=\operatorname{supp} \gamma \backslash\left\{(x, x): x \in \mathbb{R}^{d}\right\},
$$

which is a Borel set of $\mathbb{R}^{d} \times \mathbb{R}^{d}$. If we denote by $p_{x}: X \times X \rightarrow X$ the canonical projection on the first variable, we have

$$
p_{x}(\mathcal{D}(\gamma))=\left\{x \in \mathbb{R}^{d}: \exists y \neq x,(x, y) \in \operatorname{supp} \gamma\right\}=\mathcal{A}(\gamma) .
$$

We deduce that $\mathcal{A}(\gamma)$ is Lebesgue measurable as the continuous image of a Borel set.
We now show that $\mu \leq \nu$ on $\mathcal{A}(\gamma)^{c}$. Let $\phi \in C_{c}\left(\mathbb{R}^{d}, \mathbb{R}_{+}\right)$. By definition of $\mathcal{A}(\gamma)$, if $(x, y) \in$ $\operatorname{supp} \gamma$ and $x \in \mathcal{A}(\gamma)^{c}$ then $x=y$. Therefore

$$
\begin{aligned}
\int_{\mathcal{A}(\gamma)^{c}} \phi d \mu=\int \phi(x) \chi_{\mathcal{A}(\gamma)^{c}}(x) d \gamma(x, y) & =\int \phi(y) \chi_{\mathcal{A}(\gamma)^{c}}(y) \chi_{\mathcal{A}(\gamma)^{c}}(x) d \gamma(x, y) \\
& \leq \int \phi(y) \chi_{\mathcal{A}(\gamma)^{c}}(y) d \gamma(x, y)=\int_{\mathcal{A}(\gamma){ }^{c}} \phi d \nu,
\end{aligned}
$$

and the claim is proved.
Step 2. Construction of $E$.
We now consider an optimal exterior transport plan $\gamma$ for $\Upsilon(f)$ and set $g:=\gamma_{y}, h:=f+g$. By Proposition 4.1, $\gamma$ is an optimal transport plan from $f$ to $g$. Let $\bar{\gamma}$ be the image of $\gamma$ through the map $(x, y) \mapsto(y, x)$ and define

$$
E:=\mathcal{A}(\gamma) \cup \mathcal{A}(\bar{\gamma}) .
$$

We have $E^{c}=\mathcal{A}(\gamma)^{c} \cap \mathcal{A}(\bar{\gamma})^{c}$ and by the first step there holds $f \leq g$ and $g \leq f$ almost everywhere on $E^{c}$. Hence,

$$
g \chi_{E^{c}}=f \chi_{E^{c}} .
$$

To conclude the proof, we have to show that $g \equiv 1-f$ on $E$ or equivalently that up to Lebesgue negligible sets $\mathcal{A}(\gamma)$ and $\mathcal{A}(\bar{\gamma})$ are included in $\{g=1-f\}$.

On the one hand, if $x_{0} \in \mathcal{A}(\gamma)$ there exists $y_{0} \neq x_{0}$ such that $\left(x_{0}, y_{0}\right) \in \operatorname{supp} \gamma$. By Lemma 4.3, denoting

$$
S\left(x_{0}, y_{0}\right):=\left\{y \in \mathbb{R}^{d}, k\left(y-x_{0}\right)<k\left(y_{0}-x_{0}\right)\right\},
$$

we have $g=1-f$ on $S\left(x_{0}, y_{0}\right)$. Notice that $S\left(x_{0}, y_{0}\right)$ is an open set and that $x_{0} \in S\left(x_{0}, y_{0}\right)$ (since for $x \neq 0, k(x)>0=k(0))$. Hence $g\left(x_{0}\right)=1-f\left(x_{0}\right)$ and $\mathcal{A}(\gamma) \subset\{g=1-f\}$.

On the other hand, if $y_{0} \in \mathcal{A}(\bar{\gamma})$ there exists $x_{0} \neq y_{0}$ such that $\left(x_{0}, y_{0}\right) \in \operatorname{supp} \gamma$. Let us assume by contradiction that $g\left(y_{0}\right)<1-f\left(y_{0}\right)$. Without loss of generality, we can assume that $y_{0}$ is a point of Lebesgue density one of $\{g<1-f\}$. Then

$$
\lim _{r \rightarrow 0} \frac{1}{r^{d}}\left|\{g=1-f\} \cap B\left(y_{0}, r\right)\right|=0 .
$$

Thus by Lemma 4.3,

$$
\lim _{r \rightarrow 0} \frac{1}{r^{d}}\left|S\left(x_{0}, y_{0}\right) \cap B\left(y_{0}, r\right)\right|=0
$$

which contradicts $(\mathrm{H} 2)$ as $y_{0} \neq x_{0}$. Hence $g\left(y_{0}\right)=1-f\left(y_{0}\right)$ and $\mathcal{A}(\bar{\gamma}) \subset\{g=1-f\}$. This concludes the proof of the theorem.

An important corollary is the uniqueness of the second marginal of minimisers of (4.1).
Corollary 4.5. Assume that $k$ satisfies (H1) $\mathcal{E}(H 2)$. Let $f \in L^{1}\left(\mathbb{R}^{d}\right)$. Then all minimisers $\gamma$ of $\Upsilon(f)$ have the same second marginal $\gamma_{y}$.

## Proof.

We let $\gamma, \gamma^{\prime}$ be two minimisers of (4.1) and define $\widetilde{\gamma}:=\left(\gamma+\gamma^{\prime}\right) / 2$ which also minimises (4.1). We denote $g:=\gamma_{y}, g^{\prime}:=\gamma_{y}^{\prime}$ and $\widetilde{g}:=\widetilde{\gamma}_{y}$ and introduce the set

$$
F:=\left\{x \in \mathbb{R}^{d}: g(x), g^{\prime}(x), \widetilde{g}(x) \in\{f(x), 1-f(x)\}\right\} .
$$

Assuming by contradiction that $g(x) \neq g^{\prime}(x)$ for some $x \in F$, we have $1 / 2=\widetilde{g}(x) \in\{f(x), 1-$ $f(x)\}$ so that $f(x)=1 / 2$ and $g(x)=g^{\prime}(x)=1 / 2$, which is absurd. Hence $g=g^{\prime}$ on $F$ and since $F$ is of full measure by Theorem 4.4, the proof is complete.
4.2. Preliminary results for the existence of a maximiser of (4.12).

We now gather results which, combined with Theorem 4.4, allow us to prove existence of a maximiser for both $\mathcal{E}_{\text {set }}(m)$ and

$$
\begin{equation*}
\mathcal{E}(m):=\sup \left\{\Upsilon(f): f \in L^{1}\left(\mathbb{R}^{d},[0,1]\right), \int f=m\right\} \tag{4.12}
\end{equation*}
$$

We first establish a corollary of Theorem 4.4 regarding the monotonicity of the sum of the marginals of solutions to (4.1).

Corollary 4.6. Assume that $k$ satisfies (H1) $\mathcal{G}(H 2)$. Let $m>0$, let $f_{1}, f_{2} \in L_{m}^{1}$ be such that $f_{1} \leq f_{2}$ and let $\gamma^{1}, \gamma^{2}$ be respective minimisers of $\Upsilon\left(f_{1}\right)$ and $\Upsilon\left(f_{2}\right)$. Then setting $g_{1}:=\gamma_{y}^{1}$ and $g_{2}:=\gamma_{y}^{2}$, we have $f_{1}+g_{1} \leq f_{2}+g_{2}$.

Proof.
Let $f_{1}, f_{2} \in L_{m}^{1}$ be such that $f_{1} \leq f_{2}$. In the first three steps of the proof, we additionally assume that they are compactly supported. This condition is relaxed in the fourth and final step.
By Proposition 4.1, we can assume that the ambient space is a compact ball $\bar{B}_{R}$. Let $\gamma^{1}, \gamma^{2}$ be minimisers for $\Upsilon\left(f_{1}\right)$ and $\Upsilon\left(f_{2}\right)$ respectively. For $i \in\{1,2\}$ we define $g_{i}:=\gamma_{y}^{i}, h_{i}:=f_{i}+g_{i}$ and set

$$
F:=\left\{h_{1}>h_{2}\right\} .
$$

We shall prove that $|F|=0$. By Theorem 4.4, there exists $E_{1}, E_{2} \subset \bar{B}_{R}$ such that

$$
h_{1}=\chi_{E_{1}}+2 f_{1} \chi_{E_{1}^{c}} \quad \text { and } \quad h_{2}=\chi_{E_{2}}+2 f_{2} \chi_{E_{2}^{c}} .
$$

Since $h_{2} \geq 0, h_{1} \leq 1$ and $h_{2} \geq f_{2} \geq f_{1}$ we have

$$
\begin{equation*}
h_{1}>0, \quad h_{2}=2 f_{2}<1 \quad \text { and } \quad f_{1}<1 \quad \text { on } F . \tag{4.13}
\end{equation*}
$$

Step 1. $\left|E_{1}^{c} \cap F\right|=0$.
By definition of $E_{1}$ we have $h_{1}=2 f_{1}$ on $E_{1}^{c}$ and by (4.13) we have $h_{2}=2 f_{2}$ on $F$ and since $f_{1} \leq f_{2}$ we get $h_{1} \leq h_{2}$ on $E_{1}^{c} \cap F$. This contradicts the definition of $F$, hence $E_{1}^{c} \cap F=\emptyset$ and in particular $\left|E_{1}^{c} \cap F\right|=0$.

## Step 2. Intermediate claim.

Let $\psi_{1}$ be the maximal potential for $\Upsilon^{*}\left(f_{1}\right)$ given by Proposition 3.8. We define

$$
G:=\left\{\psi_{1}^{\bar{c} c}=0\right\} \cap E_{1} \cap F
$$

and claim that $|G|=0$. Let us assume by contradiction that $|G|>0$. First notice that on $F$,

$$
f_{1}+g_{1}=h_{1}>h_{2}=2 f_{2},
$$

so that

$$
g_{1}>2 f_{2}-f_{1} \geq f_{2} \geq f_{1} .
$$

Thus

$$
\begin{equation*}
G \subset E_{1} \cap F \subset\left\{g_{1}>f_{1}\right\} . \tag{4.14}
\end{equation*}
$$

Now recall that by Theorem 4.4,

$$
E_{1}=\left\{x: \exists y \neq x \text { such that }(x, y) \in \operatorname{supp} \gamma^{1} \text { or }(y, x) \in \operatorname{supp} \gamma^{1}\right\} .
$$

Together with (4.14) we obtain that for almost every $y_{0} \in G$ there exists $x_{0} \neq y_{0}$ with $\left(x_{0}, y_{0}\right) \in$ $\operatorname{supp} \gamma^{1}$. Without loss of generality, we assume that $y_{0}$ is a point of density of $G$ and we set

$$
S\left(x_{0}, y_{0}\right):=\left\{y \in \mathbb{R}^{d}: k\left(y-x_{0}\right)<k\left(y_{0}-x_{0}\right)\right\} .
$$

By (H2), we have $\left|G \cap S\left(x_{0}, y_{0}\right)\right|>0$. Let now $\widetilde{y} \in G \cap S\left(x_{0}, y_{0}\right)$. By Proposition 3.7, $\left(\psi_{1}^{\bar{c}}, \psi_{1}^{\bar{c} c}\right)$ forms a pair of Kantorovitch potentials for the optimal transport from $f_{1}$ to $g_{1}$. Thus

$$
\psi_{1}^{\bar{c}}\left(x_{0}\right)+\psi_{1}^{\bar{c} c}\left(y_{0}\right)=k\left(y_{0}-x_{0}\right) \quad \text { and } \quad \psi_{1}^{\bar{c}}\left(x_{0}\right)+\psi_{1}^{\bar{c} c}(\widetilde{y}) \leq k\left(\widetilde{y}-x_{0}\right) .
$$

However, $y_{0}, \widetilde{y} \in G$, so that $\psi_{1}^{\bar{c} c}\left(y_{0}\right)=\psi_{1}^{\bar{c} c}(\widetilde{y})=0$, hence

$$
\psi_{1}^{\bar{c}}\left(x_{0}\right)=k\left(y_{0}-x_{0}\right) \quad \text { and } \quad \psi_{1}^{\bar{c}}\left(x_{0}\right) \leq k\left(\widetilde{y}-x_{0}\right) .
$$

Eventually, as $\widetilde{y} \in S\left(x_{0}, y_{0}\right)$, we conclude that

$$
\psi_{1}^{\bar{c}}\left(x_{0}\right) \leq k\left(\widetilde{y}-x_{0}\right)<k\left(y_{0}-x_{0}\right)=\psi_{1}^{\bar{c}}\left(x_{0}\right),
$$

obtaining a contradiction. Thus $|G|=0$, which is the claim.
Step 3. $\left|E_{1} \cap F\right|=0$.
By Proposition 3.7,

$$
\left\{\psi_{1}^{\bar{c} c}>0\right\} \subset\left\{g_{1}=0\right\} \quad \text { so that } \quad\left\{g_{1}>0\right\} \subset\left\{\psi_{1}^{\bar{c} c} \leq 0\right\} .
$$

We observe that $g_{1}=1-f_{1}$ on $E_{1} \cap F$. By (4.13), $E_{1} \cap F \subset\left\{g_{1}>0\right\}$ and by the previous step, $E_{1} \cap F \subset\left\{\psi_{1}^{\bar{c} c} \neq 0\right\}$, hence $\psi_{1}^{\bar{c} c}<0$ almost everywhere on $E_{1} \cap F$.

Let $\psi_{2}$ be the maximal potential for $\Upsilon^{*}\left(f_{2}\right)$ given by Proposition 3.8. As $f_{1} \leq f_{2}$, we have $\psi_{1} \geq \psi_{2}$ so that $\psi_{1}^{\bar{c} c} \geq \psi_{2}^{\bar{c} c}$. Thus

$$
\psi_{2}^{\bar{c} c}<0 \text { on } E_{1} \cap F \text {. }
$$

By Proposition 3.7 we deduce that

$$
h_{2}=g_{2}+f_{2}=1 \text { on } E_{1} \cap F .
$$

But since $h_{2}<1$ on $F$ we get that $\left|E_{1} \cap F\right|=0$ and with the first step we conclude that $|F|=0$.
Step 4. Extension to the non-compact case.
Let $f_{1}, f_{2} \in L_{m}^{1}$ be such that $f_{1} \leq f_{2}$. For $i \in\{1,2\}$, we set $f_{i, R}=f_{i} \chi_{B_{R}}$, consider $\gamma_{R}^{i}$ an optimal exterior transport plan for $\Upsilon\left(f_{i, R}\right)$ and set $g_{i, R}:=\left(\gamma_{R}^{i}\right)_{y}$. Applying the previous steps to $f_{1, R}$ and $f_{2, R}$, we obtain

$$
\begin{equation*}
f_{1, R}+g_{1, R} \leq f_{2, R}+g_{2, R} . \tag{4.15}
\end{equation*}
$$

For $i \in\{1,2\}, f_{i, R} L^{1}$-converges to $f_{i}$ as $R \rightarrow \infty$. By Proposition $4.2(i i i), \Upsilon\left(f_{i, R}\right) \rightarrow \Upsilon\left(f_{i}\right)$ as $R \rightarrow \infty$. Additionally, $\gamma_{R}^{i}$ admits a subsequence converging weakly-* to some $\widetilde{\gamma}^{i}$ admissible for $\Upsilon\left(f_{i}\right)$. By lower semi-continuity of $\gamma \mapsto \int c d \gamma$ with respect to weak convergence, we get

$$
\int c d \widetilde{\gamma^{i}} \leq \liminf _{R} \int c d \gamma_{R}^{i}=\liminf _{R} \Upsilon\left(f_{i, R}\right)=\Upsilon\left(f_{i}\right) .
$$

Hence $\widetilde{\gamma}^{i}$ is optimal for $\Upsilon\left(f_{i}\right)$, so that by Corollary 4.5, $\widetilde{\gamma}_{y}^{i}=g_{i}$. Finally, as $\gamma_{R}^{i} \stackrel{*}{\gamma^{i}}$ as $R \rightarrow \infty$, $g_{j, R}$ converges in duality with $C_{c}\left(\mathbb{R}^{d}\right)$ to $g_{i}$ as $R \rightarrow \infty$. Multiplying (4.15) by $\phi \in C_{c}\left(\mathbb{R}^{d}, \mathbb{R}_{+}\right)$, integrating and passing to the limit we obtain that for any $\phi \in C_{c}\left(\mathbb{R}^{d}, \mathbb{R}_{+}\right)$,

$$
\int\left(f_{1}+g_{1}\right) \phi \leq \int\left(f_{2}+g_{2}\right) \phi .
$$

Hence $f_{1}+g_{1} \leq f_{2}+g_{2}$ which completes the proof.
We now prove that $\mathcal{E}$ is strictly superadditive.
Proposition 4.7. Assume that $k$ satisfies (H1) $\mathcal{G}(H 3)$. Let $m \in(0, \infty)$ and define $e(m):=$ $\mathcal{E}(m) / m$. Then, $e$ is increasing on $(0, \infty)$. As a consequence, given $0<m^{\prime}<m$,

$$
\mathcal{E}\left(m^{\prime}\right)+\mathcal{E}\left(m-m^{\prime}\right)<\mathcal{E}(m)
$$

Proof.
Let $M>m>0$. We have to establish that $\mathcal{E}(m)<(m / M) \mathcal{E}(M)$.
Step 1. $\mathcal{E}(m) \leq(m / M) \mathcal{E}(M)$.
For $R>0$ we set

$$
\Gamma_{R}:=\left\{(x, y) \in \mathbb{R}^{d} \times \mathbb{R}^{d}:|x-y|>R\right\}
$$

Let $0 \leq \varepsilon<\mathcal{E}(m) / 2$ and $f \in L_{m}^{1}$ of mass exactly $m$ and such that $\Upsilon(f) \geq \mathcal{E}(m)-\varepsilon$. We denote $\lambda:=(M / m)^{1 / d}>1$ and we set

$$
f_{\lambda}(x):=f(x / \lambda) \quad \text { for } x \in \mathbb{R}^{d},
$$

so that $\int f_{\lambda}=M$. Let $\gamma^{\lambda}$ be an optimal transport plan for $\Upsilon\left(f_{\lambda}\right)$. We define a Radon measure $\gamma$ by

$$
\int \xi(x, y) d \gamma(x, y):=\frac{m}{M} \int \xi(x / \lambda, y / \lambda) d \gamma^{\lambda}(x, y) \quad \text { for } \xi \in C_{c}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)
$$

Observe that $\gamma$ is admissible for $\Upsilon(f)$. By Proposition 4.1, there exists $R_{\lambda}=R_{\lambda}(M)$ such that $\gamma^{\lambda}\left(\Gamma_{R_{\lambda}}\right)=0$. Setting $R:=R_{\lambda} / \lambda$, we then have

$$
\begin{equation*}
\gamma\left(\Gamma_{R}\right)=\frac{m}{M} \gamma^{\lambda}\left(\Gamma_{R_{\lambda}}\right)=0 . \tag{4.16}
\end{equation*}
$$

Let us define

$$
\kappa(r):=\min \left\{k(z)-k(z / \lambda): r \leq|z| \leq R_{\lambda}\right\} .
$$

As $\lambda>1$ we have by (H3) that $\kappa(r)>0$ for $0<r<R_{\lambda}$. Additionally, $k(z / \lambda) \leq k(z)$ for any $z \in \mathbb{R}^{d}$. Consequently, for any $0<r<R_{\lambda}$,

$$
\begin{aligned}
\Upsilon(f) & \leq \int k(y-x) d \gamma(x, y) \\
& =\frac{m}{M} \int k\left(\frac{y-x}{\lambda}\right) d \gamma^{\lambda}(x, y) \\
& =\frac{m}{M} \int\left[k\left(\frac{y-x}{\lambda}\right)-k(y-x)\right] d \gamma^{\lambda}(x, y)+\frac{m}{M} \int k(y-x) d \gamma^{\lambda}(x, y) \\
& \leq \frac{m}{M} \int_{\Gamma_{r}}\left[k\left(\frac{y-x}{\lambda}\right)-k(y-x)\right] d \gamma^{\lambda}(x, y)+\frac{m}{M} \mathcal{E}(M) .
\end{aligned}
$$

In the integral over $\Gamma_{r}$, the term in brackets is smaller than $-\kappa(r)$. Hence, for every $0<r<R_{\lambda}$

$$
\begin{equation*}
\mathcal{E}(m)-\varepsilon \leq \Upsilon(f) \leq \frac{m}{M} \mathcal{E}(M)-\frac{m}{M} \kappa(r) \gamma^{\lambda}\left(\Gamma_{r}\right) . \tag{4.17}
\end{equation*}
$$

At this point we can send $\varepsilon$ to 0 and deduce that $\mathcal{E}(m) \leq(m / M) \mathcal{E}(M)$. However we need to establish a strict inequality. For this we prove in the next step that there exist $r^{*}, \delta>0$ not depending on $\varepsilon$ or $f$ such that $\gamma^{\lambda}\left(\Gamma_{r^{*}}\right) \geq \delta$.
Step 2. Conclusion.

For $r \geq 0$, we set

$$
\bar{k}(r):=\max \{k(z):|z| \leq r\}
$$

This function is increasing, continuous and there holds $\bar{k}(0)=0$. Notice that using a ball of mass $m$ as a candidate for the energy $\mathcal{E}(m)$, we see that $\mathcal{E}(m)>0$ for any $m>0$. Let us fix $0<r^{*}<R$ such that

$$
\begin{equation*}
m \bar{k}\left(r^{*}\right) \leq \mathcal{E}(m) / 4 \tag{4.18}
\end{equation*}
$$

By (4.16) we have $|x-y| \leq R$ for $(x, y) \in \operatorname{supp} \gamma$ and by definition $|x-y| \leq r^{*}$ for $(x, y) \notin \Gamma_{r^{*}}$. We deduce

$$
\begin{aligned}
\frac{\mathcal{E}(m)}{2}<\Upsilon(f)=\int_{\Gamma_{r^{*}}} c d \gamma+\int_{\Gamma_{r^{*}}^{c}} c d \gamma & \leq \gamma\left(\Gamma_{r^{*}}\right) \bar{k}(R)+\left(m-\gamma\left(\Gamma_{r^{*}}\right)\right) \bar{k}\left(r^{*}\right) \\
& \stackrel{(4.18)}{\leq} \gamma\left(\Gamma_{r^{*}}\right) \bar{k}(R)+\left(m-\gamma\left(\Gamma_{r^{*}}\right)\right) \frac{\mathcal{E}(m)}{4 m} .
\end{aligned}
$$

This implies

$$
\gamma\left(\Gamma_{r^{*}}\right)\left(\bar{k}(R)-\frac{\mathcal{E}(m)}{4 m}\right)>\frac{\mathcal{E}(m)}{4}
$$

Thus $4 m \bar{k}(R)>\mathcal{E}(m)$ and

$$
\frac{m}{M} \gamma^{\lambda}\left(\Gamma_{\lambda r^{*}}\right)=\gamma\left(\Gamma_{r^{*}}\right) \geq \frac{m \mathcal{E}(m)}{4 m \bar{k}(R)-\mathcal{E}(m)}=: m^{*}>0
$$

Plugging this in (4.17) with $r=\lambda r^{*}<R_{\lambda}$ we obtain

$$
\mathcal{E}(m)-\varepsilon \leq \frac{m}{M} \mathcal{E}(M)-m^{*} \kappa\left(r^{*}\right)
$$

Since $\varepsilon \in[0, \mathcal{E}(m) / 2]$ is arbitrary and $m^{*} \kappa\left(r^{*}\right)>0$, this proves the proposition.
We close this subsection with a lemma establishing that if a function $f$ nearly maximises $\mathcal{E}(m)$ for some $m>0$ then there exists a cube which is at least half filled by $f$.

Lemma 4.8. Let $m>0$. There exists a non-decreasing function $r_{0}: m \mapsto r_{0}(m)$ such that for $m>0$ and $f \in L_{m}^{1}$ with $\Upsilon(f) \geq \mathcal{E}(m) / 2$, there exists a cube $Q_{0}$ of side-length $r_{0}(m)$ such that:

$$
\int_{Q_{0}} f \geq \frac{\left|Q_{0}\right|}{2}
$$

## Proof.

Let $r_{0}>0$ to be fixed later and assume by contradiction that there exists a partition $\mathcal{Q}$ of $\mathbb{R}^{d}$ in cubes with side-length $r_{0}$ such that for every $Q \in \mathcal{Q}$,

$$
\int_{Q} f<\frac{|Q|}{2} .
$$

The strategy to get a contradiction from this hypothesis is to build an exterior transport plan for $f$ with too small transport cost. Let $Q \in \mathcal{Q}$. Since $\int_{Q}(1-f) \geq \int_{Q} f$ there exists a function $g_{Q} \geq 0$ supported in $Q$ such that $\int g_{Q}=\int_{Q} f$ and $f \chi_{Q}+g_{Q} \leq 1$. We then set

$$
\gamma_{Q}:=f \chi_{Q} \otimes \frac{g_{Q}}{\int_{Q} f} \quad \text { and } \quad \gamma:=\sum_{Q \in \mathcal{Q}} \gamma_{Q}
$$

Notice that $\gamma$ is a valid competitor for $\Upsilon(f)$. Next for $R>0$, we define

$$
\bar{k}(R):=\max \{k(x),|x| \leq R\} .
$$

We compute:

$$
\begin{array}{rl}
0<\frac{\mathcal{E}(m)}{2} \leq \Upsilon(f) \leq \sum_{Q \in \mathcal{Q}} \int_{Q \times Q} & k(y-x) d \gamma_{Q}  \tag{4.19}\\
& \leq \bar{k}\left(\sqrt{d} r_{0}\right) \sum_{Q \in \mathcal{Q}} \int_{Q} f=\bar{k}\left(\sqrt{d} r_{0}\right) \int f \leq \bar{k}\left(\sqrt{d} r_{0}\right) m
\end{array}
$$

Remarking that $\bar{k}: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$is continuous at 0 , increasing and with $\bar{k}(0)=0$, we set

$$
r_{0}:=\max \left\{r>0: \bar{k}(\sqrt{d} r) \leq \frac{\mathcal{E}(m)}{4 m}\right\}>0
$$

and obtain a contradiction with (4.19). This concludes the proof.

### 4.3. Existence of a maximiser for (4.12).

In the following subsection, we assume that (H1),(H2)\&(H3) hold and prove the existence of maximisers for (4.12).

We only have to prove that maximising sequences for $\mathcal{E}(m)$ are tight. However our result is more precise. We obtain that if $f$ nearly maximises $\mathcal{E}(m)$ then almost all its mass concentrates in a closed ball with radius $R_{*}=R_{*}(m)$. In the limit, maximisers are supported in such balls.

Proposition 4.9. Let $m>0$. There exist $R_{*}=R_{*}(m)>0$, $\varepsilon_{0}=\varepsilon_{0}(m)>0$ non-decreasing in $m$ with the following property. Let $0<\varepsilon \leq \varepsilon_{0}$ and let $f \in L_{m}^{1}$ such that $\int f=m$ and $\Upsilon(f) \geq \mathcal{E}(m)-\varepsilon$, then up to a translation there holds

$$
\int_{\mathbb{R}^{d} \backslash B_{R_{*}}} f \leq \frac{2 m}{\mathcal{E}(m)} \varepsilon .
$$

Proof.
Outline of the proof.
(Step 1) We start by using Lemma 4.8 to get a collection $\mathcal{Q}_{0}$ of cubes $Q$ of side-length $r_{0}=r_{0}(m)$ such that $\int_{Q} f+g \geq|Q| / 2$. We denote $\Omega_{0}:=\cup \mathcal{Q}_{0}$. We also consider the set $\Omega$ obtained by thickening $\Omega_{0}$ by adding the cubes closer than some distance $R=R(m)$. The real $R$ is chosen so that no mass of $f \chi_{\Omega_{0}}$ is sent outside $\Omega$ by any optimal exterior transport plan of $f$.
(Step 2) We build an exterior transport plan for $f$ whose cost is very close to $\Upsilon\left(f \chi_{\Omega}\right)$.
(Step 3) Next, we show that $\Omega$ concentrates almost all the mass of $f$. Using the strict superadditivity of $m \mapsto \mathcal{E}(m)$ and the previous step, we deduce that $m_{\Omega}:=\int f_{\chi_{\Omega}}$ is close to $m$.
(Step 4) Eventually, we show that the distance between cubes in $\mathcal{Q}_{0}$ is uniformly bounded. As the cardinal of $\mathcal{Q}_{0}$ is also bounded, we conclude that the diameter of $\Omega$ is bounded by a distance only depending on $m$.
Step 1. Construction of a collection of cubes on which $\int_{Q}(f+g) \geq|Q| / 2$.
Let $m>0$ and $f$ as in the statement of the proposition and assume that $\Upsilon(f) \geq \mathcal{E}(m) / 2$ so that

$$
\begin{equation*}
\varepsilon:=\mathcal{E}(m)-\Upsilon(f) \leq \mathcal{E}(m) / 2 . \tag{4.20}
\end{equation*}
$$

Let $\gamma$ be a minimiser for $\Upsilon(f)$ and let us set $g:=\gamma_{y}$. Let $r_{0}$ and $Q_{0}$ be given by Lemma 4.8. We denote by $\hat{\mathcal{Q}}$ the regular partition of $\mathbb{R}^{d}$ into cubes of side-length $r_{0}$ such that $Q_{0} \in \hat{\mathcal{Q}}$. For $j \geq 0$ to be fixed later, we set $r_{j}:=2^{-j} r_{0}$. Considering the partition $\mathcal{Q}$ of $\mathbb{R}^{d}$ into cubes of side-length $r_{j}$ obtained by refining $\hat{\mathcal{Q}}$, we define $\mathcal{Q}_{0}$ as the subset formed by the elements $Q \in \mathcal{Q}$ such that

$$
\int_{Q}(f+g) \geq \frac{|Q|}{4} .
$$

We remark that $\mathcal{Q}_{0}$ is not empty since

$$
\int_{Q}(f+g) \geq \int_{Q} f \geq \frac{|Q|}{2}
$$

for at least one of the $2^{j}$ sub-cubes of $Q_{0}$ in the partition $\mathcal{Q}$.
Let us define $\Omega_{0}:=\cup \mathcal{Q}_{0}$. By Proposition 4.1, there exists $R=R(m)$ such that $|x-y| \leq R$ on supp $\gamma$. We denote by $\mathcal{Q}_{R}$ the collection of cubes $Q \in \mathcal{Q}$ such that $d\left(Q, \Omega_{0}\right) \leq R$, and by $\Omega$ their union. By construction, there holds $\gamma\left(\Omega_{0} \times \Omega^{c}\right)=0$. We now define

$$
f_{\Omega}:=f \chi_{\Omega}, \quad \text { and } \quad m_{\Omega}:=\int f_{\Omega}
$$

and we let $\gamma_{\Omega}$ be an optimal exterior transport plan for $f_{\Omega}$, that is $\gamma \in \Pi_{f_{\Omega}}$ with $\int c d \gamma_{\Omega}=\Upsilon\left(f_{\Omega}\right)$.
By Proposition 4.1 again, we have (since $m_{\Omega} \leq m$ ) that

$$
\begin{equation*}
\gamma_{\Omega}\left(\Omega_{0} \times \Omega^{c}\right)=0 \tag{4.21}
\end{equation*}
$$

Step 2. Building a transport plan for $f$ whose cost is close to $\Upsilon\left(f_{\Omega}\right)$.
In this step we modify $\gamma_{\Omega}$ to build an exterior transport plan $\gamma$ for $f$ with a cost close to $\Upsilon\left(f_{\Omega}\right)$. More precisely, we require that for some constant $C=C\left(r_{j}\right)>0$ with $C\left(r_{j}\right) \rightarrow 0$ as $r_{j} \rightarrow 0$,

$$
\int c d \gamma-\int c d \gamma_{\Omega} \leq C\left(m-m_{\Omega}\right)
$$

The proof is a refinement of the proof of the Lipschitz continuity of $\Upsilon$, see Proposition 4.2 (iii). In the following we define successively the plans $\gamma^{0}, \gamma^{1}, \gamma^{2}, \gamma^{3}$ which satisfy in particular $\operatorname{supp} \gamma^{0} \subset \bar{\Omega} \times \bar{\Omega}, \quad \operatorname{supp} \gamma^{1} \subset \bar{\Omega} \times \overline{\Omega^{c}}, \quad \operatorname{supp} \gamma^{2} \subset \overline{\Omega \backslash \Omega_{0}} \times \overline{\Omega \backslash \Omega_{0}}, \quad \operatorname{supp} \gamma^{3} \subset \overline{\Omega^{c}} \times \overline{\Omega^{c}}$. First we set $\gamma^{0}:=\gamma_{\Omega}\left\llcorner\Omega \times \Omega\right.$ and denote $f^{0}:=\gamma_{x}^{0}, g^{0}:=\gamma_{y}^{0}$. We build the three remaining plans in the following substeps. These constructions will satisfy

$$
\left(\gamma^{1}+\gamma^{2}\right)_{x}=f_{\Omega}-\gamma_{x}^{0}=f \chi_{\Omega}-f^{0} \quad \text { and } \quad \gamma_{x}^{3}=f-f_{\Omega}=f \chi_{\Omega^{c}}
$$

We will set eventually $\widehat{\gamma}:=\gamma^{0}+\gamma^{1}+\gamma^{2}+\gamma^{3}$ which will be an optimal transport plan for $f$. The difficulty is to preserve the constraint $f+\widehat{\gamma}_{y} \leq 1$ while controlling the cost.
Step 2.a. Construction of $\gamma^{1}$.
Let us denote $\gamma_{\Omega}^{1}:=\gamma_{\Omega}\left\llcorner\Omega \times \Omega^{c}, f_{\Omega}^{1}:=\left(\gamma_{\Omega}^{1}\right)_{x}\right.$ and $g_{\Omega}^{1}:=\left(\gamma_{\Omega}^{1}\right)_{y}=\chi_{\Omega^{c}} g_{\Omega}$. We can not rule out the possibility that $f+g_{\Omega}>1$ in some part of $\Omega^{c}$ so that we cannot set $\gamma^{1}=\gamma_{\Omega}^{1}$. However, we will transport as much as possible mass through $\gamma_{\Omega}^{1}$. Let us define

$$
u:=\left(f+g_{\Omega}-1\right)_{+},
$$

which corresponds to the excess mass transported through $\gamma_{\Omega}^{1}$. Using the convention $0 / 0=0$, we define $\gamma^{1}$ by

$$
d \gamma^{1}(x, y):=\frac{g_{\Omega}(y)-u(y)}{g_{\Omega}(y)} d \gamma_{\Omega}^{1}(x, y)
$$

At this point, we have

$$
\begin{equation*}
\int c d\left(\gamma^{0}+\gamma^{1}\right) \leq \int c d \gamma_{\Omega}=\Upsilon\left(f_{\Omega}\right) \tag{4.22}
\end{equation*}
$$

Moreover setting $f^{1}:=\gamma_{x}^{1}$ and $g^{1}:=\gamma_{y}^{1}$, there holds supp $g^{1} \subset \overline{\Omega^{c}}$. Notice that since $f_{\Omega} \leq f$, by Corollary 4.6 we have $f_{\Omega}+g_{\Omega} \leq f+g$, so that $g_{\Omega} \leq f+g$ in $\Omega^{c}$ which implies $g^{1} \leq f+g$. Thus

$$
\begin{equation*}
\int_{Q}\left(f+g^{1}\right) \leq \int_{Q}(2 f+g)<\frac{|Q|}{2}, \quad \text { for every } Q \in \mathcal{Q} \backslash \mathcal{Q}_{R} \tag{4.23}
\end{equation*}
$$

where we used the definition of $\mathcal{Q}_{0}$ and the fact that $\left[\mathcal{Q} \backslash \mathcal{Q}_{R}\right] \cap \mathcal{Q}_{0}=\emptyset$.
Let us compute for later use the mass from $\Omega$ that still requires to be transported. By construction

$$
\begin{equation*}
\int_{\Omega}\left(f-f^{0}-f^{1}\right)=\int d\left(\gamma_{\Omega}^{1}-\gamma^{1}\right)=\int \frac{u(y)}{g_{\Omega}(y)} d \gamma_{\Omega}^{1}(x, y)=\int_{\Omega^{c}}\left(f+g_{\Omega}-1\right)_{+} \leq \int_{\Omega^{c}} f \tag{4.24}
\end{equation*}
$$

Step 2.b. Construction of $\gamma^{2}$.

We now define

$$
\gamma_{\Omega}^{2}:=\gamma_{\Omega}^{1}-\gamma^{1}=\gamma_{\Omega}-\gamma^{0}-\gamma^{1}
$$

Notice that by $(4.21), \gamma_{\Omega}^{1}\left(\Omega_{0} \times \Omega^{c}\right)=0$, so that

$$
\operatorname{supp} \gamma_{\Omega}^{2} \subset \overline{\Omega \backslash \Omega_{0}} \times \overline{\Omega^{c}}
$$

In particular, $f^{2}:=f_{\Omega}-f^{0}-f^{1}$ is supported in $\overline{\Omega \backslash \Omega_{0}}$. Let $Q \in \mathcal{Q}_{R} \backslash \mathcal{Q}_{0}$. Since $g^{0} \leq g_{\Omega}$ and $f=f_{\Omega}$ on $Q$, using Corollary 4.6 again we see that

$$
\int_{Q}\left(f+g^{0}\right) \leq \int_{Q}\left(f_{\Omega}+g_{\Omega}\right) \leq \int_{Q}(f+g) \leq \frac{|Q|}{4}
$$

Therefore for such $Q$ there exists a function $g_{Q}^{2}: Q \rightarrow \mathbb{R}_{+}$such that $f+g^{0}+g_{Q}^{2} \leq 1$ and $\int g_{Q}^{2}=\int_{Q} f^{2}$. Defining

$$
\gamma_{Q}^{2}:=\frac{1}{\int_{Q} f^{2}}\left[\chi_{Q} f^{2}\right] \otimes g_{Q}^{2} \quad \text { for } Q \in \mathcal{Q}_{R} \backslash \mathcal{Q}_{0} \quad \text { and then } \quad \gamma^{2}:=\sum_{Q \in \mathcal{Q}_{R} \backslash \mathcal{Q}_{0}} \gamma_{Q}^{2}
$$

we have $\gamma_{x}^{2}=f^{2}$ and $g^{2}:=\gamma_{y}^{2}=\sum_{Q} g_{Q}^{2}$. Hence $f+g^{0}+g^{2} \leq 1$ and

$$
\int c d \gamma^{2} \leq\left(\int f^{2}\right) \bar{k}\left(\sqrt{d} r_{j}\right)
$$

where as in the proof of Proposition 4.7 we denote $\bar{k}(r):=\max \{k(x):|x| \leq r\}$.
By construction $f^{2}=f_{\Omega}-f^{0}-f^{1}$, so by (4.24) there holds $\int f^{2} \leq m-m_{\Omega}$ which leads to the cost estimate

$$
\begin{equation*}
\int c d \gamma^{2} \leq\left(m-m_{\Omega}\right) \bar{k}\left(\sqrt{d} r_{j}\right) \tag{4.25}
\end{equation*}
$$

Step 2.c. Construction of $\gamma^{3}$.
We still have to transport the mass corresponding to $\chi_{\Omega^{c}} f$. For every $Q \in \mathcal{Q} \backslash \mathcal{Q}_{R}$ we have $\int_{Q} f \leq \int_{Q}(f+g) \leq|Q| / 4$, therefore, in view of (4.23), there exists a function $g_{Q}^{3}: Q \rightarrow \mathbb{R}_{+}$ such that $\int g_{Q}^{3}=\int_{Q} f$ and $f+g^{1}+g_{Q}^{3} \leq 1$. As in the previous step, we define

$$
\gamma_{Q}^{3}:=\frac{1}{\int_{Q} f}\left[\chi_{Q} f\right] \otimes g_{Q}^{3} \quad \text { and } \quad \gamma^{3}:=\sum_{Q} \gamma_{Q}^{3}
$$

By construction, $\left(\gamma^{3}\right)_{x}=\chi_{\Omega^{c}} f$ and denoting $g^{3}:=\left(\gamma^{3}\right)_{y}$, we have $\operatorname{supp} g^{3} \subset \overline{\Omega^{c}}$ as well as $f+g^{1}+g^{3} \leq 1$. Moreover

$$
\begin{equation*}
\int c d \gamma^{3} \leq\left(\int_{\Omega^{c}} f\right) \bar{k}\left(\sqrt{d} r_{j}\right)=\left(m-m_{\Omega}\right) \bar{k}\left(\sqrt{d} r_{j}\right) \tag{4.26}
\end{equation*}
$$

Step 2.d. Conclusion : definition and properties of $\widehat{\gamma}$.
Eventually, we set $\widehat{\gamma}:=\gamma^{0}+\gamma^{1}+\gamma^{2}+\gamma^{3}$ and $\widehat{g}:=\widehat{\gamma}_{y}$. There holds $\widehat{\gamma}_{x}=f^{0}+f^{1}+f^{2}+f^{3}=f$ and $f+\widehat{g} \leq 1$ so that $\widehat{\gamma}$ is an admissible exterior transport plan for $f$. Besides, collecting the estimates $(4.22),(4.25) \&(4.26)$ we get

$$
\begin{equation*}
\Upsilon(f) \leq \int c d \widehat{\gamma} \leq \Upsilon\left(f_{\Omega}\right)+2\left(m-m_{\Omega}\right) \bar{k}\left(\sqrt{d} r_{j}\right) \tag{4.27}
\end{equation*}
$$

Step 3. We show that $m-m_{\Omega} \leq C(m) \varepsilon$ (recall the definition (4.20) of $\varepsilon$ ).
As $\Upsilon\left(f_{\Omega}\right) \leq \mathcal{E}\left(m_{\Omega}\right)$ and $\mathcal{E}(m)-\varepsilon=\Upsilon(f),(4.27)$ yields

$$
\mathcal{E}(m)-\varepsilon \leq \mathcal{E}\left(m_{\Omega}\right)+2 \bar{k}\left(\sqrt{d} r_{j}\right)\left(m-m_{\Omega}\right)
$$

Additionally by Proposition $4.7, \mathcal{E}\left(m_{\Omega}\right) \leq \frac{m_{\Omega}}{m} \mathcal{E}(m)$. Hence

$$
\left(\frac{\mathcal{E}(m)}{m}-2 \bar{k}\left(\sqrt{d} r_{j}\right)\right)\left(m-m_{\Omega}\right) \leq \varepsilon
$$

By continuity of $k, \bar{k}\left(\sqrt{d} r_{j}\right) \rightarrow 0$ as $r_{j} \rightarrow 0$. Recalling that $r_{j}=2^{-j} r_{0}$, we fix $j \geq 0$ as the first integer such that $\bar{k}\left(\sqrt{d} r_{j}\right) \leq \mathcal{E}(m) / 4 m$ (notice that $j$ does not depend on $\varepsilon$ ). Therefore

$$
\begin{equation*}
m-m_{\Omega} \leq \frac{2 m \varepsilon}{\mathcal{E}(m)} \tag{4.28}
\end{equation*}
$$

This yields

$$
\begin{equation*}
\int_{\mathbb{R}^{d} \backslash \Omega} f=\int f-\int f_{\Omega}=m-m_{\Omega} \leq \frac{2 m \varepsilon}{\mathcal{E}(m)} \tag{4.29}
\end{equation*}
$$

For future use, let us also notice that injecting (4.28) into (4.27) we obtain

$$
\begin{equation*}
\mathcal{E}(m)-\varepsilon \leq \Upsilon(f) \leq \Upsilon\left(f_{\Omega}\right)+\varepsilon \tag{4.30}
\end{equation*}
$$

Step 4 : Bounding the diameter of $\Omega$.
We finally prove that $\Omega$ is uniformly bounded which would conclude the proof. For $Q_{-}$, $Q_{+} \in \mathcal{Q}_{0}$, we write $Q_{-} \sim Q_{+}$if there exists a finite chain

$$
\begin{equation*}
Q_{-}=Q_{0}, Q_{1}, \ldots, Q_{n}=Q_{+} \tag{4.31}
\end{equation*}
$$

such that $Q_{i} \in \mathcal{Q}_{0}$ and $d\left(Q_{i-1}, Q_{i}\right) \leq 4 R+\sqrt{d} r_{j}$ for $1 \leq i \leq n$. This defines an equivalence relation. Let us show that there exists only one equivalence class. We assume by contradiction that there exist at least two equivalence classes, and we let $\mathcal{C}^{1}$ be one of these classes and $\mathcal{C}^{2}$ be the union of the remaining classes. For $i \in\{1,2\}$, we then define $\Omega^{i}$ to be the reunion of the cubes $Q$ such that $d\left(Q, \mathcal{C}^{i}\right) \leq R$. By construction, $d\left(\Omega_{1}, \Omega_{2}\right)>2 R$. Recalling that $\Omega$ is the reunion of the cubes $Q$ such that $d\left(Q, \Omega_{0}\right) \leq R$, we have $\Omega^{1} \cup \Omega^{2}=\Omega$.

For $i \in\{1,2\}$, we set $f_{\Omega}^{i}:=f_{\Omega} \chi_{\Omega^{i}}$ and $m_{\Omega}^{i}=\int f_{\Omega}^{i}$. We have $m_{\Omega}^{1}+m_{\Omega}^{2}=m_{\Omega} \leq m$ and $m_{\Omega}^{1}, m_{\Omega}^{2} \geq 2^{-j d}\left|Q_{0}\right| / 4=2^{-j d-2}\left|Q_{0}\right|$. Additionally, by Proposition 4.2 (ii),

$$
\Upsilon\left(f_{\Omega}\right)=\Upsilon\left(f_{\Omega}^{1}\right)+\Upsilon\left(f_{\Omega}^{2}\right) \leq \mathcal{E}\left(m_{\Omega}^{1}\right)+\mathcal{E}\left(m_{\Omega}^{2}\right)
$$

Injecting this inequality into (4.30) yields

$$
\mathcal{E}(m)-\varepsilon \leq \Upsilon(f) \leq \Upsilon\left(f_{\Omega}\right)+\varepsilon \leq \mathcal{E}\left(m_{\Omega}^{1}\right)+\mathcal{E}\left(m_{\Omega}^{2}\right)+\varepsilon
$$

By definition of $e$ this rewrites as

$$
\begin{equation*}
m e(m) \leq m_{\Omega}^{1} e\left(m_{\Omega}^{1}\right)+m_{\Omega}^{2} e\left(m_{\Omega}^{2}\right)+2 \varepsilon \tag{4.32}
\end{equation*}
$$

As $m_{\Omega}^{1}+m_{\Omega}^{2} \leq m$ and for $i \in\{1,2\}, m_{\Omega}^{i} \geq 2^{-j d-2}\left|Q_{0}\right|$, we have $m_{\Omega}^{i} \leq m-2^{-j d-2}\left|Q_{0}\right|$. Recall that by Proposition 4.7, $e$ is increasing, so that $e\left(m_{\Omega}^{i}\right) \leq e\left(m-2^{-j d} m_{0}\right)$. Hence

$$
m_{\Omega}^{1} e\left(m_{\Omega}^{1}\right)+m_{\Omega}^{2} e\left(m_{\Omega}^{2}\right) \leq m e\left(m-2^{-j d} m_{0}\right)
$$

With (4.32), we obtain

$$
m e(m) \leq m e\left(m-2^{-j d} m_{0}\right)+2 \varepsilon
$$

which is absurd for $\varepsilon$ small enough because $e$ is increasing. It follows that for $\varepsilon>0$ small enough the relation $\sim$ has a single class. Recall that for all $Q \in \mathcal{Q}_{0}, \int_{Q}(f+g) \geq 2^{-j d-2}\left|Q_{0}\right|$. Thus the maximal length of a chain in (4.31) without any repetition is bounded by $N:=\left\lfloor 2^{j d+3} \mathrm{~m} /\left|Q_{0}\right|\right\rfloor$. Therefore, the diameter of $\Omega$ is bounded by $\left(4 R+2 \sqrt{d} r_{j}\right)(N+1)$ with $r_{j}$ and $N$ only depending on $m$, the dimension $d$ and the cost $c$. Together with (4.29) this proves the proposition.

We can now apply the direct method of Calculus of Variations to establish the existence of a maximiser for (4.12).

Proof of Theorem 1.1.
Let $f_{n}$ be a maximising sequence for (4.12) and let $R_{*}=R_{*}(m)$ be given by Proposition 4.9 so that up to translation,

$$
\int_{\mathbb{R}^{d} \backslash B_{R_{*}}} f_{n} \rightarrow 0 \quad \text { as } \quad n \rightarrow \infty
$$

Therefore, $f_{n} \lambda$ is a tight sequence of $\mathcal{M}_{+}\left(\mathbb{R}^{d}\right)$ and up to extraction of a subsequence it converges weakly-* to $f \lambda$ where $f$ is admissible for (4.12). By Proposition 4.2 (iv),

$$
\Upsilon(f)=\lim \Upsilon\left(f_{n}\right)=\mathcal{E}(m)
$$

so that $f$ is a maximiser for $\mathcal{E}(m)$.
Let now $f$ be any maximiser of $\mathcal{E}(m)$. Applying Proposition 4.9 to $f$ we have that up to a translation supp $f \subset \bar{B}_{R_{*}}$. This concludes the proof.

Let us show that when $f$ is compactly supported there exist Kantorovitch potentials for the problem (4.1) (this is the situation of interest as we have just established that the maximisers of $\mathcal{E}(m)$ are compactly supported in $\mathbb{R}^{d}$ ).

Lemma 4.10. Let $m>0$ and assume that $f \in L_{m}^{1}$ is compactly supported. Let $R=R(m)$ be given by Proposition 4.1 such that all maximisers $\gamma$ of $\Upsilon(f)$ are supported in $X:=\operatorname{supp} f+\bar{B}_{R}$. Then, there exists a pair $(\varphi, \psi) \in C_{c}\left(\mathbb{R}^{d}\right) \times C_{c}\left(\mathbb{R}^{d}\right)$ optimal for $\Upsilon^{*}(f)$. Additionally, $\varphi=\psi^{c}$, $\psi=\varphi_{-}^{c}$ and both $\varphi$ and $\psi$ are compactly supported in $X$.

Proof. Let us introduce $\tilde{c}:=c_{\mid X \times X}$ which is a continuous cost function on the compact set $X$. By Proposition 3.4, there exists $\widetilde{\psi} \in C(X)$ with $\widetilde{\psi}=(\widetilde{\psi} \tilde{c} \tilde{c})_{-}$such that $\Upsilon^{*}(f)=K_{f}\left(\widetilde{\psi}^{\tilde{c}}, \widetilde{\psi}\right)$. By Proposition 3.7,

$$
\left\{\widetilde{\psi}^{\tilde{c} \tilde{c}}<0\right\} \subset\{f+g=1\} \subset X
$$

Combining this with $\widetilde{\psi}=(\widetilde{\psi} \tilde{c} \tilde{c})_{-}$and $f+g=0$ on $\partial X$, we get $\widetilde{\psi}=0$ on $\partial X$. We extend the potentials on $\mathbb{R}^{d}$ by setting

$$
\psi:=\left\{\begin{array}{lll}
\widetilde{\psi} & \text { in } \quad X, \\
0 & \text { in } \quad X^{c},
\end{array} \quad \text { and for } x \in \mathbb{R}^{d}, \quad \varphi(x):=\psi^{c}(x)=\inf \left\{c(x, y)-\psi(y): y \in \mathbb{R}^{d}\right\}\right.
$$

We now show that the pair $(\varphi, \psi)$ satisfies the conclusion of the lemma.
Observe that $\psi$ is continuous and supported in $X$ and that $\psi \leq 0$. Hence $\varphi \geq 0$. Moreover, for $x \in \mathbb{R}^{d} \backslash \operatorname{supp} \psi$,

$$
\varphi(x) \leq c(x, x)-\psi(x)=0
$$

so that $\varphi$ is also supported in $X$.
Next, for $x \in X$,

$$
\varphi(x)=\min \left(\min _{y \in \mathbb{R}^{d} \backslash X} c(x, y), \min _{y \in X}\{c(x, y)-\widetilde{\psi}(y)\}\right)
$$

Let $x \in X$. For $y \in \mathbb{R}^{d} \backslash X$, there exists $\widetilde{y}$ in the intersection of the segment $[x, y]$ with $\partial X$. By continuity, $\widetilde{\psi}(\widetilde{y})=0=\widetilde{\psi}(y)$ and moreover by $(\mathrm{H} 3), c(x, \widetilde{y}) \leq c(x, y)$ so that $c(x, \widetilde{y}) \leq$ $c(x, y)-\widetilde{\psi}(y)$. We deduce that for $x \in X$ the above formula simplifies as

$$
\varphi(x)=\min _{y \in X}\{c(x, y)-\widetilde{\psi}(y)\}=\widetilde{\psi^{c}}(x)
$$

This proves that $\varphi$ is continuous and that $K_{f}(\varphi, \psi)=K_{f}(\widetilde{\psi} \tilde{c}, \widetilde{\psi})=\Upsilon^{*}(f)$. Moreover, using the same argument as above, we have $\varphi^{c}(y)=0$ for $y \notin X$. For $y \in X$,

$$
\varphi^{c}(y)=\min \left(\min _{x \in \mathbb{R}^{d} \backslash X} c(x, y), \min _{x \in X}\left\{c(x, y)-\widetilde{\psi}^{\tilde{c}}(y)\right\}\right)=\min _{x \in X}\left\{c(x, y)-\widetilde{\psi}^{\tilde{c}}(y)\right\}=\widetilde{\psi^{c} \tilde{c}}(y)
$$

We deduce $\varphi_{-}^{c}=0=\psi$ in $X^{c}$, and $\varphi_{-}^{c}=\left(\psi^{\tilde{c} \tilde{c}}\right)_{-}=\widetilde{\psi}$ in $X$. Thus $\varphi_{-}^{c}=\psi$ everywhere. This ends the proof of the lemma.

Let us now recall a variant of the bathtub principle, see [16, Theorem 1.14].
Proposition 4.11. Let $\xi: \mathbb{R}^{d} \rightarrow \mathbb{R}_{+}$be measurable and such that for all $t \geq 0,|\{\xi>t\}|<\infty$. Given $m>0$, let

$$
t:=\inf \{s \geq 0,|\{\xi>s\}| \leq m\}
$$

Then, the maximisers of

$$
\sup _{\widetilde{f}}\left\{\int \widetilde{f} \xi: \widetilde{f} \in L^{1}\left(\mathbb{R}^{d}\right), 0 \leq \widetilde{f} \leq 1, \int \widetilde{f}=m\right\}
$$

are the functions $f:=\chi_{\{\xi>t\}}+\theta$, where $\theta \in L^{1}\left(\mathbb{R}^{d},[0,1]\right)$ is supported in $\{\xi=t\}$ and satisfies

$$
\int \theta=m-|\{\xi>t\}|
$$

We are now ready to establish Corollary 1.2.
Proof of Corollary 1.2.
By Theorem 1.1, the optimisation problem (4.12) admits a compactly supported solution $f$. Let $(\varphi, \psi) \in C_{c}\left(\mathbb{R}^{d}\right) \times C_{c}\left(\mathbb{R}^{d}\right)$ be an optimal pair for $\Upsilon^{*}(f)$ provided by Lemma 4.10 , so that

$$
\Upsilon^{*}(f)=\int f(\varphi-\psi)+\int \psi .
$$

We see that $f$ is a maximiser of:

$$
\sup \left\{\int \tilde{f}(\varphi-\psi): \widetilde{f} \in L^{1}\left(\mathbb{R}^{d}\right), 0 \leq \widetilde{f} \leq 1, \int \tilde{f}=m\right\}
$$

Let us set $\xi:=\varphi-\psi \geq 0$. By Proposition 4.11 there exists $t \geq 0$ and $\theta \in L^{1}\left(\mathbb{R}^{d},[0,1]\right)$ supported in $\{\xi=t\}$ such that $f=\chi_{\{\xi>t\}}+\theta$. Notice in particular that since $\theta \in[0,1]$, we have

$$
|\{\xi=t\}| \geq \int \theta=m-|\{\xi>t\}|
$$

and there exist measurable subsets $G \subset\{\xi=t\}$ with $|G|=m-|\{\xi>t\}|$. For any such set, setting

$$
\bar{f}:=\chi_{\{\xi>t\}}+\chi_{G},
$$

we have $\Upsilon^{*}(\bar{f})=\Upsilon^{*}(f)$ and $\bar{f}$ is also a maximiser of (4.12). Since $\bar{f}$ is a characteristic function, by Theorem 4.4 and Corollary 4.5, there exists $F \subset \mathbb{R}^{d}$ such that any minimiser $\gamma$ of $\Upsilon(\bar{f})$ satisfies $\gamma_{y}=\chi_{F}$. Setting $E:=\{\xi>t\} \cup G$, we deduce that

$$
\Upsilon_{\mathrm{set}}(E)=\Upsilon^{*}(\bar{f})=\Upsilon(\bar{f})=\Upsilon(f)
$$

so that $\mathcal{E}(m)=\mathcal{E}_{\text {set }}(m)$, which concludes the proof.

## 5. Maximisers of (4.12) are characteristic functions of balls

In this section we prove Theorem 1.3. We assume that $c(x, y)=k(|y-x|)$ with $k \in C\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)$ increasing and coercive and with $k(0)=0$. In particular, we have now $c=\bar{c}$, so that the operations of $c$-transform and $\bar{c}$-transform coincide. Also notice that by Theorem 4.4, if $f=\chi_{E}$ for some Lebesgue measurable set $E$ then $\Upsilon_{\text {set }}(E)=\Upsilon\left(\chi_{E}\right)$. By abuse of notation, we write $\Upsilon(E)$ for $\Upsilon\left(\chi_{E}\right)$. Since the class of costs that we consider is invariant by scaling we assume without loss of generality that $m=\omega_{d}$.

We now recall the definition of symmetric rearrangement of functions with constant sign (see [16, Chapter 3] for more details on symmetric rearrangements).

## Definition 5.1.

(i) Given a measurable set $A \subset \mathbb{R}^{d}$, we define the symmetric rearrangement of $A$ as the open ball $A^{*}$ centred at the origin and of volume $|A|$.
(ii) Let $\varphi: \mathbb{R}^{d} \rightarrow \mathbb{R}_{+}$be measurable and such that for every $t \geq 0,|\{\varphi>t\}|<\infty$. Its symmetric decreasing rearrangement is defined by

$$
\varphi^{*}(x):=\int_{\mathbb{R}_{+}} \chi_{\{\varphi>t\}^{*}}(x) d t
$$

(iii) Let $\psi: \mathbb{R}^{d} \rightarrow \mathbb{R}_{-}$be measurable and such that for every $t \leq 0,|\{\psi<t\}|<\infty$. Its symmetric increasing rearrangement is defined by

$$
\psi_{*}(x):=-(-\psi)^{*}(x)=-\int_{\mathbb{R}_{-}} \chi_{\{\psi<t\}^{*}}(x) d t .
$$

The following lemma recalls some basic properties of the symmetric increasing rearrangement $\psi_{*}$ of a non-positive function $\psi$. All these properties but the continuity of $\psi_{*}$ follow immediately from the definition. The fact that continuity is preserved by symmetric rearrangement is wellknown but we have no reference for this at hand. We provide a short proof for the reader's convenience.

Lemma 5.2. Let $\psi: \mathbb{R}^{d} \rightarrow \mathbb{R}_{-}$be as in Definition 5.1. Then, $\psi_{*}$ is non-positive, radial, nondecreasing, and for any $t \leq 0,\left\{\psi_{*}<t\right\}=\{\psi<t\}^{*}$. Besides, if $\psi$ is supported in a compact set of diameter bounded by $2 R>0$ then $\psi_{*}$ is supported in $\bar{B}_{R}$. If moreover $\psi$ is continuous then $\psi_{*}$ is also continuous.
Proof of the last point. Let $\psi \in C_{c}\left(\mathbb{R}^{d}, \mathbb{R}_{-}\right)$. First, as the strict sublevels sets $\left\{\psi_{*}<t\right\}$ are the open balls $\{\psi<t\}^{*}, \psi_{*}$ is upper semi-continuous (note that this is true even when $\psi$ is not continuous).

Let us now establish that $\psi_{*}$ is lower semi-continuous, i.e. that for any $t \leq 0,\left\{\psi_{*} \leq t\right\}$ is closed. We first notice that $\left\{\psi_{*} \leq 0\right\}=\mathbb{R}^{d}$ is closed. Given $t<0$, let $t_{n}<0$ be a decreasing sequence converging to $t$. Observe that if for some $n \geq 0,\left\{\psi<t_{n}\right\}=\emptyset$, then $\left\{\psi_{*} \leq t\right\}=\emptyset$ is closed. Next, we assume that for every $n \geq 0$,

$$
\begin{equation*}
\left\{\psi<t_{n}\right\} \neq \emptyset \tag{5.1}
\end{equation*}
$$

We denote by $R_{n}$ the radius of the ball $\left\{\psi_{*}<t_{n}\right\}$. Notice that the sequence $R_{n}$ is non-increasing and bounded by 0 , so that $R_{n}$ converges to some $R \geq 0$.
Let us show that the sequence $R_{n}$ is decreasing. By contradiction, we assume that $R_{n}=R_{n+1}$ for some $n \geq 0$. Then $\left\{\psi<t_{n}\right\}^{*}=\left\{\psi<t_{n+1}\right\}^{*}$ and $\left|\left\{t_{n+1} \leq \psi<t_{n}\right\}\right|=0$. Using (5.1) and the fact that $\psi$ is compactly supported, there exists $x$ such that $\psi(x)<t_{n+1}$ and $y$ such that $\psi(y)>t_{n}$. Thus by continuity of $\psi$ there exists $z$ such that $\psi(z)=\left(t_{n+1}+t_{n}\right) / 2$. By continuity of $\psi$ again, there exists $\eta>0$ such that $B_{\eta}(z) \subset\left\{t_{n+1}<\psi<t_{n}\right\}$, contradicting the fact that $\left|\left\{t_{n+1} \leq \psi<t_{n}\right\}\right|=0$. As a conclusion, the sequence $R_{n}$ is decreasing and

$$
\left\{\psi_{*} \leq t\right\}=\bigcap_{n \geq 0}\left\{\psi_{*}<t_{n}\right\}=\bigcap_{n \geq 0} B_{R_{n}}=\bar{B}_{R}
$$

Hence $\psi_{*}$ is lower semi-continuous and therefore continuous.
To prove Theorem 1.3, we need a last lemma about final lemma characterising optimal potentials of $\Upsilon^{*}\left(\chi_{B_{1}}\right)$.
Lemma 5.3. Let $\left(\psi^{c}, \psi\right)$ be a pair of optimal potentials for $\Upsilon^{*}\left(\chi_{B_{1}}\right)$ such that $\psi$ is radially symmetric and non-decreasing. Then $\psi^{c}$ is radially symmetric and non-increasing. Besides, $\psi^{c}$ is radially decreasing on $B_{1}$.
Proof. Combining the facts that $k$ is continuous, that $k(r) \rightarrow \infty$ as $r \rightarrow \infty$ and that $\psi$ is bounded, we see that for any $x \in \mathbb{R}^{d}$,

$$
\psi^{c}(x)=\min \left\{k(|y-x|)-\psi(y): y \in \mathbb{R}^{d}\right\} .
$$

As $\psi$ is radially symmetric non-decreasing and $k$ is increasing, we easily see that

$$
\begin{equation*}
\psi^{c}(x)=\min \{k(|y-x|)-\psi(y): y, \exists \lambda \geq 1, y=\lambda x\} \tag{5.2}
\end{equation*}
$$

which in turn implies that $\psi^{c}$ is radially symmetric.
From now on, for radial functions $\zeta: \mathbb{R}^{d} \rightarrow \mathbb{R}$, we make the abuse of notation $\zeta(r)=\zeta(r \sigma)$ for $r \geq 0$ where $\sigma$ is some fixed element of $\mathbb{S}^{d-1}$. With this convention (5.2) reads

$$
\begin{equation*}
\psi^{c}(r)=\min _{s \geq r} k(s-r)-\psi(s) \tag{5.3}
\end{equation*}
$$

Let us prove that $\psi^{c}$ is non-decreasing. Let $0 \leq r_{1} \leq r_{2}$. By (5.3), there exists $r \geq r_{1}$ such that

$$
\begin{equation*}
\psi^{c}\left(r_{1}\right)=k\left(r-r_{1}\right)-\psi(r) \tag{5.4}
\end{equation*}
$$

If $r \leq r_{2}$, we use $\psi^{c}\left(r_{2}\right) \leq k(0)-\psi\left(r_{2}\right)=-\psi\left(r_{2}\right)$ and deduce from (5.4) and the fact that $\psi$ is non-decreasing that

$$
\psi^{c}\left(r_{2}\right)-\psi^{c}\left(r_{1}\right) \leq \psi(r)-\psi\left(r_{2}\right)-k\left(r-r_{1}\right) \leq 0
$$

If $r>r_{2}$, we use $\psi^{c}\left(r_{2}\right) \leq k\left(r-r_{2}\right)-\psi(r)$ to get

$$
\psi^{c}\left(r_{2}\right)-\psi^{c}\left(r_{1}\right) \leq k\left(r-r_{2}\right)-k\left(r-r_{1}\right) \leq 0
$$

because $r_{1} \leq r_{2}<r$ and $k$ is increasing. In both cases $\psi^{c}\left(r_{2}\right)-\psi^{c}\left(r_{1}\right) \leq 0$. Hence $\psi^{c}$ is non-increasing on $\mathbb{R}^{d}$.

We finally prove that $\psi^{c}$ is decreasing on $B_{1}$. Let $0<r_{1}<r_{2}<1$. Given $\gamma$ a minimiser for $\Upsilon\left(B_{1}\right)$, there exists $y \in \mathbb{R}^{d} \backslash B_{1}$ such that $\left(r_{1} \sigma, y\right) \in \operatorname{supp} \gamma$. By Proposition 3.7 , $\gamma$ is an optimal transport plan between $f$ and $g:=\gamma_{y}$ and $\left(\psi^{c}, \psi^{c c}\right)$ is a pair of Kantorovitch potentials for the transport between $f$ and $g$. Therefore,

$$
\begin{equation*}
\psi^{c}\left(r_{1}\right)+\psi^{c c}(|y|)=k\left(\left|y-r_{1} \sigma\right|\right) . \tag{5.5}
\end{equation*}
$$

Let us prove by contradiction that $y \in[1,+\infty) \sigma$. Assume it is not and let $y^{\prime}:=|y| \sigma$. Recalling that $|y| \geq 1>r_{1}$, we have $\left|y^{\prime}-r_{1} \sigma\right|=|y|-r_{1}<\left|y-r_{1} \sigma\right|$ and since $k$ is increasing we deduce

$$
k\left(\left|y^{\prime}-r_{1} \sigma\right|\right)<k\left(\left|y-r_{1} \sigma\right|\right) .
$$

Then, by definition of $\psi^{c c}$ and taking into account that it is radially symmetric we get

$$
\psi^{c}\left(r_{1}\right)+\psi^{c c}(|y|) \leq k\left(\left|y^{\prime}-r_{1} \sigma\right|\right)<k\left(\left|y-r_{1} \sigma\right|\right)
$$

which contradicts (5.5). Therefore, $y=r \sigma$ for some $r \geq 1$. By definition of the $c$-transform,

$$
\begin{equation*}
\psi^{c}\left(r_{2}\right)+\psi^{c c}(r) \leq k\left(r-r_{2}\right) \tag{5.6}
\end{equation*}
$$

Subtracting (5.5) to (5.6), we obtain

$$
\psi^{c}\left(r_{2}\right)-\psi^{c}\left(r_{1}\right) \leq k\left(r-r_{2}\right)-k\left(r-r_{1}\right)<0
$$

where we used $r_{1}<r_{2}<1 \leq r$. This shows that $\psi^{c}$ is decreasing on $B_{1}$.
We are now ready to prove Theorem 1.3.
Proof of Theorem 1.3.
Part I : Unit balls are maximisers of $\mathcal{E}\left(\omega_{d}\right)$.
By Theorem 1.1, there exists a compactly supported maximiser $f$ for (4.12) with $m=\omega_{d}$. By Lemma 4.10, there exists an optimal pair $\left(\psi^{c}, \psi\right) \in C_{c}\left(\mathbb{R}^{d}\right) \times C_{c}\left(\mathbb{R}^{d}\right)$ for problem $\Upsilon^{*}(f)$ such that $\psi=\left(\psi^{c c}\right)_{-}$.
Step 1. We build a radially symmetric maximiser for (4.12).
Let $\psi_{*}$ be the symmetric increasing rearrangement of $\psi$. By Lemma 5.2 , as $\psi \in C_{c}\left(\mathbb{R}^{d}\right)$, we also have $\psi_{*} \in C_{c}\left(\mathbb{R}^{d}\right)$. We denote by $\psi_{*}^{c}$ the function $\left(\psi_{*}\right)^{c}$. By definition, $\psi_{*}^{c} \oplus \psi_{*} \leq c$. Proceeding as in the proof of Lemma 4.10, we obtain $\psi_{*}^{c} \in C_{c}\left(\mathbb{R}^{d}\right)$. Thus $\left(\psi_{*}^{c}, \psi_{*}\right)$ is admissible for $\Upsilon^{*}\left(B_{1}\right)$.

Notice that $(f, \psi)$ solves the double supremum problem (recall the definition (3.2) of $K_{f}$ )

$$
\sup _{f} \sup _{\psi \in C_{c}\left(\mathbb{R}^{d}\right)}\left\{K_{f}\left(\psi^{c}, \psi\right): 0 \leq f \leq 1, \int f=\omega_{d}, \psi \leq 0\right\} .
$$

Hence
$\mathcal{E}\left(\omega_{d}\right)=K_{f}\left(\psi^{c}, \psi\right)=\int f\left(\psi^{c}-\psi\right)+\int \psi \geq \Upsilon^{*}\left(B_{1}\right) \geq K_{\chi_{B_{1}}}\left(\psi_{*}^{c}, \psi_{*}\right)=\int_{B_{1}}\left(\psi_{*}^{c}-\psi_{*}\right)+\int \psi_{*}$.
In the remainder of this step, we establish the converse inequality

$$
\begin{equation*}
K_{f}\left(\psi^{c}, \psi\right) \leq K_{\chi_{B_{1}}}\left(\psi_{*}^{c}, \psi_{*}\right), \tag{5.7}
\end{equation*}
$$

so that $B_{1}$ is a maximiser of $\mathcal{E}\left(\omega_{d}\right)$ and the first part of Theorem 1.3 is proved. Notice that (5.7) also implies that $\left(\psi_{*}^{c}, \psi_{*}\right)$ is a pair of optimal potentials for $\Upsilon^{*}\left(B_{1}\right)$. To establish (5.7), we first notice that by construction $\int \psi=\int \psi^{*}$ so that we only need to prove

$$
\begin{equation*}
\int f\left(\psi^{c}-\psi\right) \leq \int_{B_{1}}\left(\psi_{*}^{c}-\psi_{*}\right) . \tag{5.8}
\end{equation*}
$$

In Step 2 below we establish the inequality

$$
\begin{equation*}
\left(\psi^{c}\right)^{*} \leq\left(\psi_{*}\right)^{c}=\psi_{*}^{c}, \tag{5.9}
\end{equation*}
$$

where $\left(\psi^{c}\right)^{*}$ denotes the symmetric decreasing rearrangement of $\psi^{c}$. Admitting that (5.9) holds we deduce (5.8) as follows. Since $f$ is non-negative and compactly supported we have by the Hardy-Littlewood inequality (see [16, Theorem 3.4])

$$
\begin{equation*}
-\int f \psi \leq-\int f^{*} \psi_{*} \quad \text { and } \quad \int f \psi^{c} \leq \int f^{*}\left(\psi^{c}\right)^{*} \stackrel{(5.9)}{\leq} \int f^{*} \psi_{*}^{c} . \tag{5.10}
\end{equation*}
$$

Using that $-\psi_{*}$ and $\psi_{*}^{c}$ are radially symmetric and non-increasing, we may appeal to Proposition 4.11 and conclude that separately,

$$
\begin{equation*}
-\int f \psi \leq-\int \chi_{B_{1}} \psi_{*} \quad \text { and } \quad \int f \psi^{c} \leq \int \chi_{B_{1}} \psi_{*}^{c} . \tag{5.11}
\end{equation*}
$$

Summing these inequalities gives (5.8) and thus (5.7). This proves that $\chi_{B_{1}}$ is a maximiser for $\mathcal{E}\left(\omega_{d}\right)$ and then that $B_{1}$ is a maximiser for $\mathcal{E}_{\text {set }}\left(\omega_{d}\right)$.
Step 2. Proof of (5.9).
As $\psi_{*}^{c}$ and $\left(\psi^{c}\right)^{*}$ are both continuous radially symmetric functions, to prove (5.9) it is sufficient to establish that for any $t>0,\left\{\left(\psi^{c}\right)^{*}>t\right\} \subset\left\{\psi_{*}^{c}>t\right\}$, i.e. that

$$
\begin{equation*}
\left|\left\{\left(\psi^{c}\right)^{*}>t\right\}\right|=\left|\left\{\psi^{c}>t\right\}\right| \leq\left|\left\{\psi_{*}^{c}>t\right\}\right| . \tag{5.12}
\end{equation*}
$$

Recall that as $\psi \in C_{c}\left(\mathbb{R}^{d}\right)$ and $k \in C\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)$with $k(x) \rightarrow \infty$ as $x \rightarrow \infty$, for any $x \in \mathbb{R}^{d}$ the function $k(|y-x|)-\psi(y)$ admits a minimum on $\mathbb{R}^{d}$. Thus for any $x \in \mathbb{R}^{d}$ the infimum defining $\psi^{c}(x)$ (see Definition 2.2) is reached. Recalling that $k$ is also radially symmetric and increasing, we obtain

$$
\begin{aligned}
\left\{\psi^{c}>t\right\} & =\left\{x \in \mathbb{R}^{d}: \min \left\{k(|y-x|)-\psi(y): y \in \mathbb{R}^{d}\right\}>t\right\} \\
& =\left\{x \in \mathbb{R}^{d}:-\psi>t-k(r) \text { on } \bar{B}_{r}(x) \forall r \geq 0\right\} \\
& =\bigcap_{r \geq 0}\left\{x \in \mathbb{R}^{d}:-\psi>t-k(r) \text { on } \bar{B}_{r}(x)\right\} \\
& =\bigcap_{r \geq 0}\{-\psi>t-k(r)\}_{r},
\end{aligned}
$$

where for $\Omega \subset \mathbb{R}^{d}$ and $r \geq 0, \Omega_{r}$ is defined as $\Omega_{r}:=\left\{x \in \Omega: d\left(x, \mathbb{R}^{d} \backslash \Omega\right)>r\right\}$. In particular,

$$
\begin{equation*}
\left|\left\{\psi^{c}>t\right\}\right| \leq \inf _{r \geq 0}\left|\{-\psi>t-k(r)\}_{r}\right| . \tag{5.13}
\end{equation*}
$$

We observe that $\{-\psi>t-k(r)\}$ is an open set for any $t>0$ and $r \geq 0$. We also notice that (5.13) holds for all $\psi \in C_{c}\left(\mathbb{R}^{d}\right)$. In particular, it holds for $\psi_{*}$. Moreover, as $\psi_{*}$ is radially non-decreasing by construction, the sets $\left\{-\psi_{*}>t-k(r)\right\}_{r}$ are open balls centred at the origin and we have in fact

$$
\left|\left\{\psi_{*}^{c}>t\right\}\right|=\inf _{r \geq 0}\left|\left\{-\psi_{*}>t-k(r)\right\}_{r}\right|
$$

Let us now prove the following claim.
Claim. Let $s>0$ and $V>0$.
(i) If $V>\omega_{d} s^{d}$ then, among open sets $\Omega \subset \mathbb{R}^{d}$ of volume $V,\left|\Omega_{s}\right|$ is maximal if and only if $\Omega$ is a ball.
(ii) If $V \leq \omega_{d} s^{d}$ then $\left|\Omega_{s}\right|=0$ for any set of volume $V$.

Let $V>0$ and $s>0$ and let $\Omega \subset \mathbb{R}^{d}$ be an open set. We assume without loss of generality that $|\Omega|=V$ and $\left|\Omega_{s}\right|>0$. Notice that we always have $\Omega_{s}+B_{s} \subset \Omega$ (but the converse inclusion may fail). By the Brunn-Minkowski inequality (see for instance [11]) applied to $\Omega_{s}$ and $B_{s}$, there holds

$$
\begin{equation*}
V^{1 / d}=|\Omega|^{1 / d} \geq\left|\Omega_{s}+B_{s}\right|^{1 / d} \geq\left|\Omega_{s}\right|^{1 / d}+\left|B_{s}\right|^{1 / d} \tag{5.14}
\end{equation*}
$$

If $\Omega_{s}$ is a ball, then $\Omega$ is a ball of volume $V, \Omega=\Omega_{s}+B_{s}$ and we have equality in (5.14). Conversely if we have equality in (5.14), by the equality case of the Brunn-Minkowski inequality and the fact that $s>0, \Omega_{s}$ is a ball and $|\Omega|=\left|\Omega_{s}+B_{s}\right|$, so that $\Omega$ is a ball. This proves the first part of the claim.

Regarding the second part, we assume that $|\Omega| \leq \omega_{d} s^{d}$ and (by contradiction) that $\left|\Omega_{s}\right|>0$. The above reasoning applies and we have $\Omega=\Omega_{s}+B_{s}$ so that $\left|\Omega_{s}\right|>0$ implies $|\Omega|>\left|B_{s}\right|=\omega_{d} s^{d}$ and we get a contradiction. This proves the claim.

By definition, $\{-\psi>t-k(r)\}$ and $\left\{-\psi_{*}>t-k(r)\right\}$ have the same volume. As a consequence of the claim, for any $t>0$ and $r>0$,

$$
\begin{equation*}
\left|\{-\psi>t-k(r)\}_{r}\right| \leq\left|\left\{-\psi_{*}>t-k(r)\right\}_{r}\right| \tag{5.15}
\end{equation*}
$$

Notice that the previous inequality is an equality if $r=0$, as $\Omega_{0}=\Omega$ for any open set $\Omega$. Taking the infimum on $r \geq 0$ yields

$$
\begin{equation*}
\left|\left\{\psi^{c}>t\right\}\right| \leq \inf _{r \geq 0}\left|\{-\psi>t-k(r)\}_{r}\right| \leq \inf _{r \geq 0}\left|\left\{-\psi_{*}>t-k(r)\right\}_{r}\right|=\left|\left\{\psi_{*}^{c}>t\right\}\right| . \tag{5.16}
\end{equation*}
$$

This proves (5.12) which in turn implies (5.9).

Part II : Unit balls are the unique maximisers of $\mathcal{E}\left(\omega_{d}\right)$.
Step 1. Proof of $f=\chi_{\left\{\psi^{c}>\psi_{*}^{c}(1)\right\}}$ (exploiting the equality case in the bathtub principle).
We now show that any maximiser $f$ is of the form $\chi_{\left\{\psi^{c}>\psi_{*}^{c}(1)\right\}}$. By Lemma 5.3, $\psi_{*}^{c}$ is radially decreasing on $B_{1}$ and non-increasing on $\mathbb{R}^{d}$. Thus $\chi_{B_{1}}$ is the only function maximising

$$
\sup _{\widetilde{f}}\left\{\int \tilde{f} \psi_{*}^{c}: 0 \leq \tilde{f} \leq 1, \int \tilde{f}=\omega_{d}\right\}
$$

As $\Upsilon^{*}\left(\chi_{B_{1}}\right)=\mathcal{E}\left(\omega_{d}\right)$, the inequalities in (5.10) and (5.11) are in fact equalities (and (5.9) is also an equality in $\left.B_{1}\right)$. Namely, there hold

$$
\left(\psi^{c}\right)^{*}=\psi_{*}^{c}, \quad-\int f \psi=-\int f^{*} \psi_{*} \quad \text { and } \quad \int f \psi^{c}=\int f^{*} \psi_{*}^{c}
$$

This leads to

$$
\int f \psi^{c}=\int f^{*}\left(\psi^{c}\right)^{*}=\int f^{*} \psi_{*}^{c}=\int_{B_{1}} \psi_{*}^{c},
$$

and $f$ is a maximiser of

$$
\sup _{\tilde{f}}\left\{\int \tilde{f} \psi^{c}: 0 \leq \tilde{f} \leq 1, \int \tilde{f}=\omega_{d}\right\} .
$$

Let us now prove that $\left|\left\{\psi^{c}>\psi_{*}^{c}(1)\right\}\right|=\omega_{d}$ (which with Proposition 4.11 yields $\left.f=\chi_{\left\{\psi^{c}>\psi_{*}^{c}(1)\right\}}\right)$. Since $\left(\psi^{c}\right)^{*}=\psi_{*}^{c}$ in $B_{1}$, and $\psi_{*}^{c}$ is decreasing in $B_{1}$ by Lemma 5.3, there holds for $t \geq \psi_{*}^{c}(1)$,

$$
\begin{equation*}
\left|\left\{\psi^{c}>t\right\}\right|=\left|\left\{\left(\psi^{c}\right)^{*}>t\right\}\right|=\left|\left\{\psi_{*}^{c}>t\right\}\right| . \tag{5.17}
\end{equation*}
$$

Using this for $t=\psi_{*}^{c}(1)$ we get $\left|\left\{\psi^{c}>\psi_{*}^{c}(1)\right\}\right|=\omega_{d}$ and we conclude with Proposition 4.11 that $f=\chi_{\left\{\psi^{c}>\psi_{*}^{c}(1)\right\}}$.
Step 2. We prove that $\left\{\psi^{c}>t\right\}$ is a ball for $t>\psi_{*}^{c}(1)$ (exploiting the equality case in the Brunn-Minkowski inequality).
Step 2.a.
We fix $t>\psi_{*}^{c}(1)$. Combining (5.17) and (5.16), we get that

$$
\begin{equation*}
\left|\left\{\psi^{c}>t\right\}\right|=\inf _{r \geq 0}\left|\{-\psi>t-k(r)\}_{r}\right|=\inf _{r \geq 0}\left|\left\{-\psi_{*}>t-k(r)\right\}_{r}\right|=\left|\left\{\psi_{*}^{c}>t\right\}\right| . \tag{5.18}
\end{equation*}
$$

The following claim is established in Step 2.b below.
Claim. There exists $r_{*}=r_{*}(t)>0$ such that

$$
\left|\left\{\psi_{*}^{c}>t\right\}\right|=\inf _{r \geq 0}\left|\left\{-\psi_{*}>t-k(r)\right\}_{r}\right|=\mid\left\{-\psi_{*}>t-k\left(r_{*}\right)\right\}_{r_{*}} .
$$

Provisionally assuming the claim let us prove that $\left\{\psi^{c}>t\right\}$ is a ball.
We assume without loss of generality that $\left|\left\{\psi_{*}^{c}>t\right\}\right|>0$ (otherwise $\left|\left\{\psi^{c}>t\right\}\right| \leq\left|\left\{\psi_{*}^{c}>t\right\}\right|=0$ by (5.15) and the open set $\left\{\psi^{c}>t\right\}$ is empty). Next, the claim, (5.18) and (5.15) yield that $r_{*}(t)$ also minimises $\inf _{r \geq 0}\left|\{-\psi>t-k(r)\}_{r}\right|$. Thus by (5.18), $\left\{-\psi_{*}>t-k\left(r_{*}(t)\right)\right\}_{r_{*}(t)}$ is a ball of positive volume. As $r_{*}(t)>0$, by the equality case of the claim of Part I, Step 2, the set $\left\{-\psi>t-k\left(r_{*}(t)\right)\right\}_{r_{*}(t)}$ is also a ball. As $\left\{\psi^{c}>t\right\} \subset\left\{-\psi>t-k\left(r_{*}(t)\right)\right\}_{r_{*}(t)}$, by (5.18) the inclusion is actually an equality. Hence $\left\{\psi^{c}>t\right\}$ is a ball.

Step 2.b. Proof of the claim.
We first show that there exist $0<r_{t}<R_{t}<\infty$ such that

$$
\begin{equation*}
\inf _{r \geq 0}\left|\left\{-\psi_{*}>t-k(r)\right\}_{r}\right|=\inf _{r_{t} \leq r \leq R_{t}}\left|\left\{-\psi_{*}>t-k(r)\right\}_{r}\right| . \tag{5.19}
\end{equation*}
$$

We start with the upper bound on $r$. By (H1)\&(H3), there exists $R_{t}$ such that $k\left(R_{t}\right)=t+1$.
Hence, if $r>R_{t},\left\{-\psi_{*}>t-k(r)\right\}=\mathbb{R}^{d}$. We can thus only consider the radii $r \leq R_{t}$.
We now prove the lower bound on $r$. Recall that $t>\psi_{*}^{c}(1)$ and that $\psi_{*}^{c}$ is decreasing in $B_{1}$. Therefore there exists $R_{*}(t)<1$ such that

$$
\left\{\psi_{*}^{c}>t\right\}=B_{R_{*}(t)} .
$$

We set $r_{t}:=\frac{1-R_{*}(t)}{2}$ and claim that (5.19) holds for this value. To ease notation, let us set for $r>0$

$$
S_{r}:=\left\{-\psi_{*}>t-k(r)\right\}_{r}=\left\{x \in \mathbb{R}^{d}:-\psi_{*}>t-k(r) \text { on } \bar{B}_{r}(x)\right\} .
$$

We also define $\bar{R}:=\frac{1+R_{*}(t)}{2}$. In order to prove (5.19) it is enough to show that

$$
\begin{equation*}
\left\{\psi_{*}^{c}>t\right\}=\cap_{r \geq r_{t}} S_{r} . \tag{5.20}
\end{equation*}
$$

Recalling that the sets $S_{r}$ are centred balls and that $B_{R_{*}(t)} \subset B_{\bar{R}}$, we have

$$
\left\{\psi_{*}^{c}>t\right\}=\cap_{r \geq 0}\left(S_{r} \cap B_{\bar{R}}\right) .
$$

We now claim that

$$
\cap_{r \geq r_{t}}\left(S_{r} \cap B_{\bar{R}}\right) \subset \cap_{r<r_{t}}\left(S_{r} \cap B_{\bar{R}}\right),
$$

which is equivalent to

$$
\begin{equation*}
\cup_{r<r_{t}}\left(S_{r}^{c} \cap B_{\bar{R}}\right) \subset \cup_{r \geq r_{t}}\left(S_{r}^{c} \cap B_{\bar{R}}\right) \tag{5.21}
\end{equation*}
$$

To prove this let $x \in S_{r}^{c} \cap B_{\bar{R}}$ for some $r<r_{t}$. By definition of $S_{r}^{c}$,

$$
\min _{y \in \bar{B}_{r}(x)} k(r)-\psi_{*}(y) \leq t
$$

In particular since $k$ is increasing, there exists $y \in \bar{B}_{r}(x)$ such that

$$
k(|x-y|)-\psi^{*}(y) \leq t
$$

As $x \in B_{\bar{R}} \subset B_{1}$, and $\left(\psi_{*}^{c}, \psi_{*}^{c c}\right)$ are Kantorovitch potentials for the external transport minimising $\Upsilon\left(B_{1}\right)$ (see Proposition 3.7) there exists $z \in B_{1}^{c}$ such that $\psi_{*}^{c c}(z)=\psi_{*}(z)$ (by (3.14)) and

$$
\psi_{*}^{c}(x)=k(x-z)-\psi_{*}(z)=\min _{y} k(|x-y|)-\psi_{*}(y) \leq t
$$

Since $z \in B_{1}^{c}$ and $x \in B_{\bar{R}}$ we have

$$
r^{\prime}=|z-x| \geq 1-\bar{R}=\frac{1-R_{*}(t)}{2}=r_{t}
$$

and thus

$$
\min _{z \in \bar{B}_{r^{\prime}}(x)} k\left(r^{\prime}\right)-\psi_{*}(z) \leq t
$$

so that $x \in S_{r^{\prime}}$. This shows (5.21) which implies

$$
\left\{\psi_{*}^{c}>t\right\}=\cap_{r \geq r_{t}}\left(S_{r} \cap B_{\bar{R}}\right)
$$

Eventually, we must have $S_{r} \subset B_{\bar{R}}$ for some $r \geq r_{t}$ (otherwise $\left\{\psi_{*}^{c}>t\right\}=B_{\bar{R}}$ which is absurd). This concludes the proof of (5.20) and thus of (5.19).
Next, setting

$$
L(r):=\left|\left\{-\psi_{*}>t-k(r)\right\}_{r}\right|
$$

we still have to establish that the infimum of $L$ over $\left[r_{t}, R_{t}\right]$ is reached. For this we establish that $L$ is lower semi-continuous (together with (5.19) this will conclude the proof of the existence of $r_{*}(t)>0$ minimising $L$ over $\left.\mathbb{R}_{+}\right)$. We start by noticing that, $r \mapsto\left|\left\{-\psi_{*}>r\right\}\right|$ is lower semi-continuous on $\mathbb{R}_{+}$. Let us denote by $\rho_{t}(r)$ the radius of the ball $\left\{-\psi_{*}>t-k(r)\right\}$. As $k$ is continuous, the function $r \mapsto \rho_{t}(r)$ is also lower semi-continuous. Finally, as $L(r)=$ $\omega_{d}\left[\left(\rho_{t}(r)-r\right)_{+}\right]^{d}, L$ is lower semi-continuous as well. This ends the proof of the claim.
Step 3. Conclusion.
Let now $t_{n}$ be a decreasing sequence converging to $\psi_{*}^{c}(1)$. We have

$$
\begin{equation*}
\left\{\psi^{c}>\psi_{*}^{c}(1)\right\}=\bigcup_{n \geq 0}\left\{\psi^{c}>t_{n}\right\} \quad \text { and } \quad\left\{\psi_{*}^{c}>\psi_{*}^{c}(1)\right\}=\bigcup_{n \geq 0}\left\{\psi_{*}^{c}>t_{n}\right\}=B_{1} \tag{5.22}
\end{equation*}
$$

By (5.18), for every $n \geq 0,\left\{\psi^{c}>t_{n}\right\}=B_{r_{n}}\left(z_{n}\right)$, where $r_{n}$ is the radius of $\left\{\psi_{*}^{c}>t_{n}\right\}$ and $z_{n} \in \mathbb{R}^{d}$. Since $t_{n}$ is decreasing the sequence $B_{r_{n}}\left(z_{n}\right)$ is non-decreasing. Moreover, by (5.22) $r_{n} \rightarrow 1$ as $n \rightarrow \infty$. Hence there exists $z \in \mathbb{R}^{d}$ such that $\chi_{B_{r_{n}}} \rightarrow \chi_{B_{1}}(z)$ monotonically in $L^{1}\left(\mathbb{R}^{d}\right)$ as $n \rightarrow \infty$. Eventually (5.22) implies that $\left\{\psi^{c}>\psi_{*}^{c}(1)\right\}=B_{1}(z)$. Consequently, $f=\chi_{\left\{\psi^{c}>\psi_{*}^{c}(1)\right\}}=\chi_{B_{1}(z)}$. This concludes the proof of the fact that balls are the unique maximisers to (4.12).

## References

[1] Bonacini, M., Knüpfer, H., and Röger, M. Optimal distribution of oppositely charged phases: Perfect screening and other properties. SIAM Journal on Mathematical Analysis 48, 2 (2016), 1128-1154.
[2] Burchard, A., Choksi, R., and Topaloglu, I. Nonlocal shape optimization via interactions of attractive and repulsive potentials. Indiana University Mathematics Journal 67, 1 (2018), 375-395.
[3] Buttazzo, G., Carlier, G., and Laborde, M. On the Wasserstein distance between mutually singular measures. Adv. Calc. Var. 13, 2 (2020), 141-154.
[4] Candau-Tilh, J., and Goldman, M. Existence and stability results for an isoperimetric problem with a non-local interaction of Wasserstein type. ESAIM: COCV 28 (2022), 37.
[5] Choksi, R., Muratov, C. B., and Topaloglu, I. An old problem resurfaces nonlocally: Gamow's liquid drops inspire today's research and applications. Not. Am. Math. Soc. 64, 11 (2017), 1275-1283.
[6] Cicalese, M., Luca, L. D., Novaga, M., and Ponsiglione, M. Ground states of a two phase model with cross and self attractive interactions. SIAM Journal on Mathematical Analysis 48, 5 (2016), 3412-3443.
[7] De Philippis, G., Mészàros, A. R., Santambrogio, F., and Velichkov, B. BV estimates in optimal transportation and applications. Archive for Rational Mechanics and Analysis 219, 2 (2015), 829-860.
[8] Figalli, A. The optimal partial transport problem. Archive for Rational Mechanics and Analysis 195 (02 2010), 533-560.
[9] Frank, R. L., and Lieb, E. H. A compactness lemma and its application to the existence of minimizers for the liquid drop model. SIAM Journal on Mathematical Analysis 47, 6 (2015), 4436-4450.
[10] Frank, R. L., and Nam, P. T. Existence and nonexistence in the liquid drop model. Calculus of Variations and Partial Differential Equations 60, 6 (2021), 223.
[11] Gardner, R. J. The Brunn-Minkoswki inequality. Bull. Amer. Math. Soc. 39, 3 (2002), 355-405.
[12] Goldman, M., Merlet, B., and Pegon, M. Uniform $C^{1, \alpha}$-regularity for almost-minimizers of some nonlocal perturbations of the perimeter. preprint, arXiv:2209.11006 (2022).
[13] Goldman, M., Novaga, M., and Ruffini, B. Existence and stability for a non-local isoperimetric model of charged liquid drops. Arch. Ration. Mech. Anal. 217, 1 (2015), 1-36.
[14] Knupfer, H., Muratov, C. B., and Novaga, M. Low density phases in a uniformly charged liquid. Communications in Mathematical Physics 345, 1 (2016), 141-183.
[15] Knüpfer, H., and Muratov, C. B. On an isoperimetric problem with a competing nonlocal term i: The planar case. Communications on Pure and Applied Mathematics 66, 7 (2013), 1129-1162.
[16] Lieb, E., and Loss, M. Analysis. Crm Proceedings \& Lecture Notes. American Mathematical Society, 2001.
[17] Lions, P. L. The concentration-compactness principle in the calculus of variations. the locally compact case, part 1. Annales de l'I.H.P. Analyse non linéaire 1, 2 (1984), 109-145.
[18] Lussardi, L., Peletier, M. A., and Röger, M. Variational analysis of a mesoscale model for bilayer membranes. Journal of Fixed Point Theory and Applications 15, 1 (2014), 217-240.
[19] Novack, M., Topaloglu, I., and Venkatraman, R. Least Wasserstein distance between disjoint shapes with perimeter regularization. Journal of Functional Analysis 284, 1 (2023).
[20] Pegon, P., Santambrogio, F., and Xia, Q. A fractal shape optimization problem in branched transport. Journal de Mathématiques Pures et Appliquées 123 (2019), 244-269.
[21] Peletier, M. A., and Röger, M. Partial localization, lipid bilayers, and the elastica functional. Arch. Ration. Mech. Anal. 193, 3 (2009), 475-537.
[22] Santambrogio, F. Optimal transport for applied mathematicians, vol. 87 of Progress in Nonlinear Differential Equations and their Applications. Birkhäuser/Springer, Cham, 2015. Calculus of variations, PDEs, and modeling.
[23] Xia, Q., and Zhou, B. The existence of minimizers for an isoperimetric problem with Wasserstein penalty term in unbounded domains. Advances in Calculus of Variations (2021).
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