



**HAL**  
open science

## Forecasting of turbulence impact on optical link from geostationary satellite

Emile Klotz, Sidonie Lefebvre, Nicolas Védrenne, Christian Musso, Thierry Fusco, Sylvain Poulenard, Laurent Coret, Alexis Louis

► **To cite this version:**

Emile Klotz, Sidonie Lefebvre, Nicolas Védrenne, Christian Musso, Thierry Fusco, et al.. Forecasting of turbulence impact on optical link from geostationary satellite. International Conference on Space Optics - ICSO 2022, Oct 2022, Dubrovnik, Croatia. pp.1277758, 10.1117/12.2690991 . hal-04187990

**HAL Id: hal-04187990**

**<https://hal.science/hal-04187990v1>**

Submitted on 25 Aug 2023

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Forecasting of turbulence impact on optical link from geostationary satellite

Emile Klotz<sup>a,1</sup>, Sidonie Lefebvre<sup>a</sup>, Nicolas Védrenne<sup>a</sup>, Christian Musso<sup>a</sup>, Thierry Fusco<sup>a</sup>,  
Sylvain Poulénard<sup>b</sup>, Laurent Coret<sup>b</sup>, Alexis Louis<sup>b</sup>

<sup>a</sup>DOTA, ONERA, Université Paris Saclay, F-91123 Palaiseau, France

<sup>b</sup>Airbus Defence and Space, Toulouse, France

## ABSTRACT

With the development of optical links for space-ground communications comes the need to mitigate the effects of the atmospheric turbulence to guarantee a lossless connection. By having a network of addressable ground stations we want to guarantee to always target a point where the link is available. For this to work, we need to be able to predict the forthcoming link availability for each station while keeping costs low. We have developed a method that allows, in the geostationary case, to obtain the power margin on the link without exhaustive knowledge of the turbulence state. In this work, we show that the sole knowledge of 4 integrated parameters of the turbulence profile ( $C_n^2$ ) and associated wind profile, which can be measured with low-complexity instruments, provides us with enough information to entirely describe the statistics of the received optical power after an adaptive optics correction. We further develop the method to take into account digital mitigation techniques (interleaving and numerical error correction) and obtain the link power margin with a maximum error compatible with current assumptions made in commonly used link budgets.

**Keywords:** Turbulence, Adaptive Optics, Forecasting, Artificial Intelligence (AI), Optical Communications, Optical Satellite Feeder Link

## 1. INTRODUCTION

With the always-increasing demand for data and bandwidth, and the potential saturation of the radio-frequency spectrum, digital optical links at  $1.55\mu\text{m}$  become a valuable alternative for ground-satellite data transmission. Many projects are ongoing to establish such a link between a GEO satellite and a fixed Earth ground station for missions such as Internet delivery and/or data repatriation from point to point on earth. However, the development of these atmospheric optical links remains conditioned by their availability, which is highly dependent on the atmospheric channel (absorption, scattering and turbulence).

High precision availability assessment of the impact of atmospheric turbulence on optical communications is necessary to manage handover between optical ground stations of a network and thus guarantee continuity of the service.

In this paper, we present a typical mission scenario (GEO satellite at  $30^\circ$  elevation) and the Tenerife site where the ESA optical ground station is located. Focusing on the downlink, we investigate the possibility to assess the availability of the link with high precision, taking into account every aspect of the reliability mechanisms (adaptive optics, coding and interleaving).

For our method to be suitable for a massive deployment of ground stations around the world, we wish to use exclusively data provided by low complexity instruments. Following this approach, we showed that the sole knowledge of integrated parameters of the turbulence profile ( $C_n^2$ ) and associated wind profile provides us with enough information to describe the statistics of the received optical power into a single mode fibered receiver after an Adaptive Optics correction (AO).

---

<sup>1</sup> [emile.klotz@onera.fr](mailto:emile.klotz@onera.fr)

Using machine learning on a database of atmospheric profiles (turbulence and wind) over Tenerife and a physical AO modelling tool, we propose here the methodology to build a metamodel to assess the probability density function of the injected power into the single mode fiber of the receiver of the optical ground station. We further develop this model to be able to estimate as well the temporal autocorrelation function of the injected power.

With the knowledge acquired on the distribution and temporal autocorrelation of injected power, we propose a method to determine the power margin of an optical downlink using only a limited set of atmospheric parameters.

## 2. MISSION SCENARIO

The considered mission scenario is the feeder link from/to a geostationary satellite.

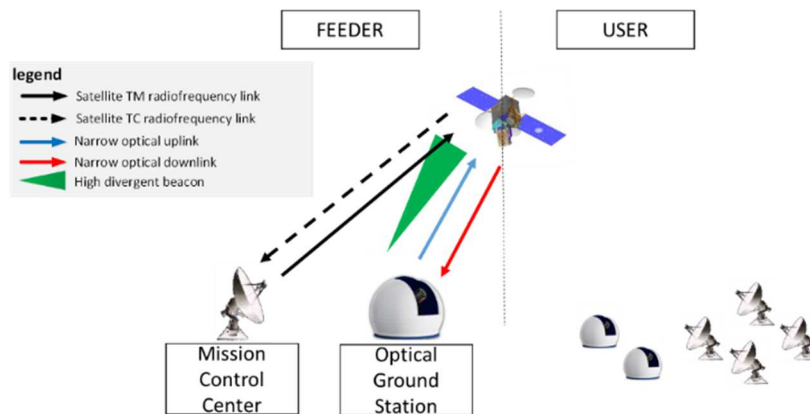


Figure 1: Mission overview.

## 3. OGS SITES

Airbus NL is developing optical ground station solution via the ScyLight activity CREOLA for such GEO feeder link<sup>1</sup>. In a first step, Airbus NL will operate:

- The already validated ESA OGS with support of its partner GA-Synopta as the receiving Optical Ground Terminal;
- Its own adaptive optics corrected TX-OGS in a transportable configuration, together with partners ASA and TNO.

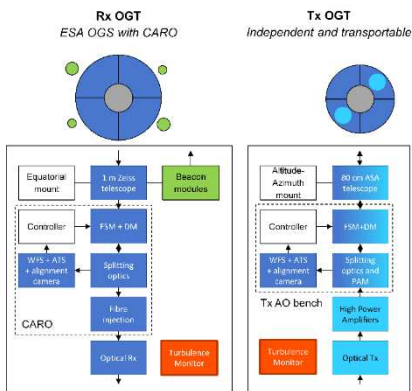


Figure 2: Bi-static OGS approach for Phases 1 & 2.



Figure 3: Miratlas Integrated Sky Monitor.

The turbulence monitoring capability will include a Miratlas Integrated Sky Monitor<sup>2</sup>, an instrument providing characterization of optical atmospheric conditions. It delivers a number of integrated parameters of the turbulence and wind profile as well as local meteorological information.

## 4. DATABASES

The elaboration of the metamodel that relates integrated parameters to optical link performance relies on the exploitation of a large enough database of atmospheric conditions. Several possibilities can be considered to access such database: either from high resolution measurements<sup>3</sup> or from numerical models<sup>4</sup>. High resolution  $C_n^2$  and wind profile measurements are nowadays rarely available and if available, they are limited to night-time and/or astronomical sites. In addition to this, they present the risk to be affected by measurement noise (hence influencing the metamodel). On the other hand data provided by numerical models present the advantage to precisely control underlying hypothesis and input parameters at the expense of a more questionable relevance. As the first goal of this work being to demonstrate the pertinence of the possibility to rely on a metamodel for performance assessment in relevant conditions, the possibility to cover a large scope of atmospheric conditions justifies in itself to exploit data obtained from a numerical model.

### 4.1 Turbulence and wind database

The profiles database has been provided by Durham University. Wind and  $C_n^2$  profiles were obtained through a global turbulence model capable of converting global meteorological re-analysis models (in this case ERA5<sup>5</sup>) into 3-Dimensional optical turbulence maps. This model is similar to the one described in Osborn and Sarazin<sup>6</sup> but was improved to include a separate boundary layer and enable stronger turbulence strength near the ground to be modelled<sup>7</sup>.

The ECMWF model from which the turbulence is being calculated has a spatial resolution of  $0.3^\circ$  along latitude and longitude and a temporal resolution of 1 hour. We chose a grid of 11 by 11 points around Tenerife's island, which is a site of interest for a potential future ground station, and focused on the first 19 days of march 2018. It leaves us with 121 simultaneous measurements for each hour, with some missing values. On the overall 19 days considered it leaves us with 37059 profiles on 113 pressure levels each.

In the following we consider that these profiles are representative of field data and could be obtained with instruments directly measuring turbulence.



Figure 4: Tenerife with OGS-Rx accurately localized by red dot. The purple line indicate the real line of sight at  $33,2^\circ$ . Each color squared is a point of simulation at a given time.

### 4.2 Power injected time series database and limitation

Correlated time series of received optical power in a single mode fiber are obtained from the  $C_n^2$  and wind profile using SAOST (Simplified Adaptive Optics Simulation Tool)<sup>8,9,10</sup> a pseudo-analytic adaptive optics simulation tools.

This tool assumes decorrelation between phase and scintillation perturbations. Scintillation influence is simulated assuming the small perturbation approximation in the Rytov regime. In practice this constrains the validity of the approach to limited Rytov variance (for horizontal propagation, typically  $\sigma_\chi^2 < 0.3$ ). Correlated time series of  $\rho_\phi$ , the coupling efficiency neglecting scintillation, and  $\rho_I$ , the term of scintillation are generated independently. The received optical power (ROP) attenuation is obtained through the product of these two quantities.

The residual phase error is computed in SAOST using a Monte-Carlo approach. Following the algorithm described in Roddier<sup>11</sup> random occurrences of Zernike coefficients are sampled to describe the corrected phase. The temporal correlation of  $\rho_\phi$  is obtained by filtering the raw Zernike coefficients by their temporal power spectral model<sup>12</sup> in the Fourier

domain. Some hypothesis have to be made on the turbulence condition of which we can cite the use of von Kármán wind turbulence model and a fixed outer scale. More details and comparisons with an end-to-end simulation can be found in Canuet<sup>13</sup>.

The temporal impact of the scintillation contribution  $\rho_I$  to the overall attenuation is approximated by:

$$\rho_I = \exp(-\sigma_\chi^2) \exp(2\chi_{AP}(t))$$

**Eq. 1**

in the weak fluctuation regime of the scintillation index  $\chi(r, t)$  where  $\sigma_\chi^2$  is the variance of the punctual log amplitude and  $\chi_{AP}(t)$  is the log-amplitude averaged on the receiver aperture.

Under weak irradiance fluctuation hypothesis,  $\chi_{AP}(t) \sim \mathcal{N}(-\sigma_{\chi_{AP}}^2, \sigma_{\chi_{AP}}^2)$ <sup>14,15</sup>. Series are randomly generated following this normal law and the temporal correlation is obtained by filtering with the spectrum given in Robert<sup>16</sup>.

In the following all simulations have been made for a state of the art mid-range adaptive optics system. The number of corrected radial order is fixed to 10 and the AO frequency to 2kHz. The diameter of the ground station telescope is 60cm.

### 4.3 Choice of relevant integrated parameters

We have based our choice of integrated parameters on the integrated parameters regularly used to describe the error budget of an adaptive optics.

The first chosen parameter is Fried's parameter<sup>17</sup>, a key parameter when one is interested in the effects of turbulence. Fried's parameter  $r_0$  is defined as the typical diameter of a telescope that would be limited by atmospheric turbulence. In the case of a plane wave considering a Kolmogorov spectrum, we have along the line of sight:

$$r_0 = \left[ 0.42 \left( \frac{2\pi}{\lambda} \right)^2 \int_0^\infty C_n^2(z) dz \right]^{-\frac{3}{5}}$$

**Eq. 2**

Where  $\lambda$  is the wavelength (1.55 $\mu$ m). Experimentally, Fried parameter can be estimated either from the amplitude jitter of a star at the focal plane of an imager or more robustly thanks to differential imaging such as performed with a DIMM<sup>18</sup>.

The second parameter, denoted  $\bar{h}$ <sup>19</sup>, is a measure of the height dispersion of atmospheric layers, homogeneous to an altitude:

$$\bar{h} = \left[ \frac{\int_0^\infty z^{\frac{5}{3}} C_n^2(z) dz}{\int_0^\infty C_n^2(z) dz} \right]^{\frac{3}{5}}$$

**Eq. 3**

It provides an assessment of the physical origin for the angular decorrelation of the phase perturbations (the influence of distant turbulence layers) while being independent from the turbulence strength. It is related to the Fried parameter and the isoplanatic patch  $\theta_0$ <sup>20</sup> by:

$$\theta_0 = 0.314 \frac{r_0}{\bar{h}}$$

**Eq. 4**

The estimation of  $\bar{h}$  can therefore be performed thanks to a measurement of  $\theta_0$  which is derived for instance from limited aperture averaged scintillation by night-time<sup>21</sup> or thanks to a Shabar measurement by day-time<sup>22</sup>. The last parameter, denoted  $\bar{v}$  describes an average wind speed over the turbulent layers<sup>23</sup>:

$$\bar{v} = \left[ \frac{\int_0^\infty v(z)^{\frac{5}{3}} C_n^2(z) dz}{\int_0^\infty C_n^2(z) dz} \right]^{\frac{3}{5}}$$

Eq. 5

It is related to the turbulence coherence time  $\tau$  according to:

$$\tau = 0.314 \frac{r_0}{\bar{v}}$$

Eq. 6

The turbulence coherence time can be extracted for instance from the temporal analysis of the jitter of a bright enough point source image. This evaluation can then be exploited to provide an estimation of  $\bar{v}$ .

In the following the point source log-amplitude variance  $\sigma_\chi^2$  is supposed to be known (measured for instance from the scintillation of a bright point source with a small diameter instrument).

We report on figure 5 the distribution and cumulative distribution of the Fried parameter for all the profiles of the database. Same quantities for  $\bar{h}$  and  $\bar{v}$  are plotted figure 6 and figure 7 respectively. Point source log-amplitude variance distribution and cumulative distribution are plotted figure 8. Evaluations are calculated at  $1.55 \mu\text{m}$  for a  $33.2^\circ$  elevation. Point source log-amplitude variance is very rarely above 0.3 which tends to consolidate the Rytov regime hypothesis.

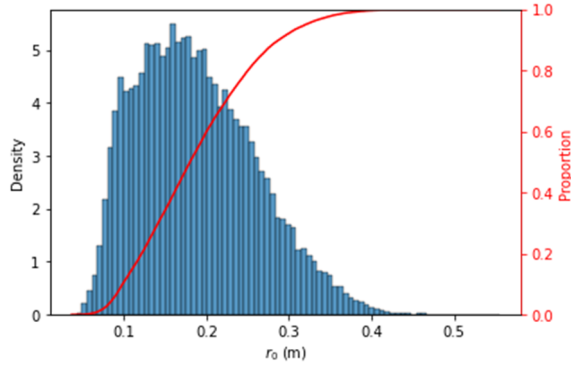


Figure 5: Distribution and cumulative distribution of  $r_0$  (m).

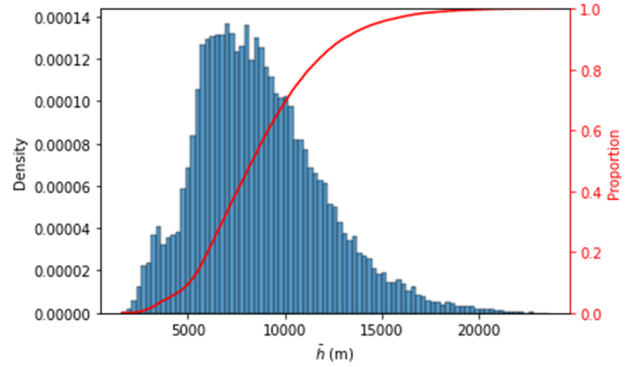


Figure 6: Distribution and cumulative distribution of  $\bar{h}$  (m) for all profiles.

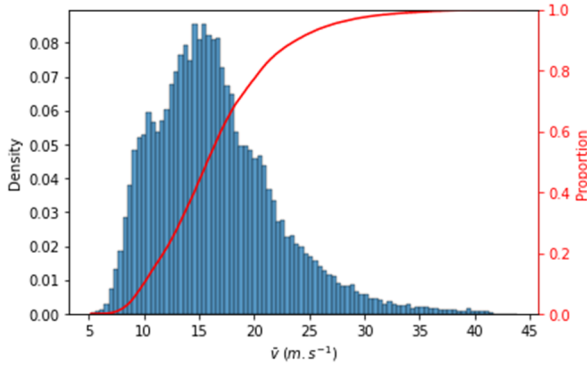


Figure 7: Distribution and cumulative distribution of  $\bar{v}$  ( $\text{m.s}^{-1}$ ) for all profiles.

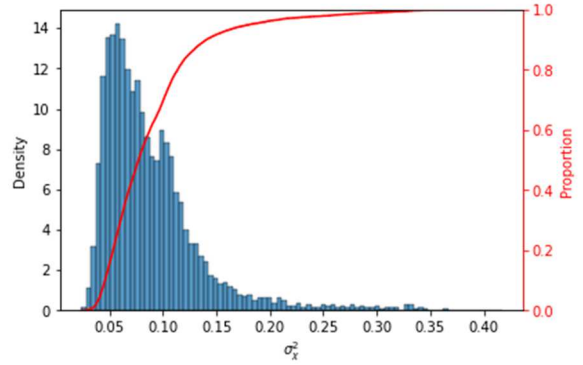


Figure 8: Distribution and cumulative distribution of  $\sigma_\chi^2$ .

## 5. PREDICTION OF STATISTICS OF DOWNLINK INJECTED POWER

Now having a database of over 37,000 time series of received optical power and the associated integrated parameters of the turbulence, we investigate the possibility to estimate the statistics of coupled flux using only these moments. We adopt the same approach as in SAOST, i.e. to consider independently the effects due to the phase error and those due to the scintillation; the benefit of this dissociation being to ease the understanding of the results.

### 5.1 Method

Using the same notation as in 4.2, we want a machine-learning algorithm that would give us the probability density function (PDF) of  $\rho_\phi$  and of  $\rho_I$ . In a second step, we will also want to estimate the half correlation time of  $\rho_\phi$  and of  $\rho_I$ . A usual and effective way to describe a distribution using machine learning is to parameterize the PDF and then use a model to estimate the parameters. In 4.2, we gave the parametric form of  $\rho_I$  that depends only on one unknown parameter  $\sigma_{\chi_{AP}}^2$  assuming the point source log-amplitude variance is measured. For  $\rho_\phi$ , it is a bit more complex as there is no known parametric form so we conducted a statistical study to determine which density function best fits the data.

We were led to study  $L_{\phi(t)} = 10 \log_{10}(\rho_\phi(t))$ , the loss in power induced by the phase fluctuation in dB as it gives a bigger weight to the low values of  $\rho_\phi(t)$  that are the critical values for our application.

Studying the distribution of this loss on our 37k profiles, it appears that the tail of the distribution shows an exponential decay, which is consistent with the closed form of the distribution proposed by Canuet<sup>10</sup>. Among all light-tailed distributions, the one that statistically best fits our data is a Gumbel distribution of the following form:

$$\frac{1}{\beta} e^{-(z+e^z)}$$

**Eq. 7**

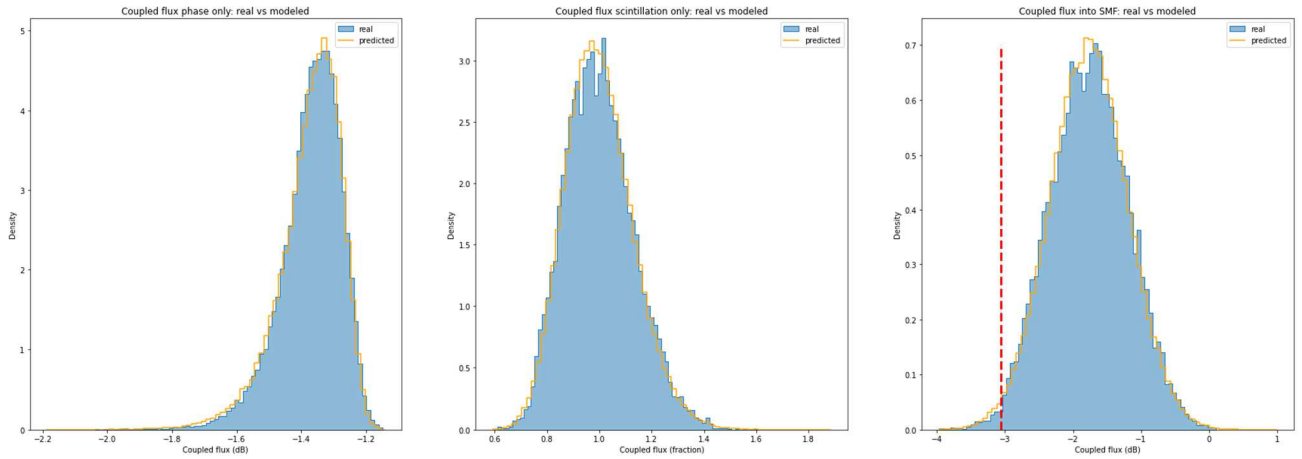
where  $z = \frac{x-\mu}{\beta}$ . Comparing real and fitted distribution quantiles, we show that for the 1% quantile the median relative error is 0.48% and never exceeds 2.01%, and the bigger the quantile the lower the error.

This result is particularly interesting as, with the Gumbel distribution being a good approximation of the density probability of phase related power attenuation we can describe the latest with only two parameters,  $\mu$  and  $\beta$  and then describe power attenuation distribution with only 3 parameters :  $\mu$ ,  $\beta$  and  $\sigma_{\chi_{AP}}$  (4 parameters if we include the point source log-amplitude variance which is supposed to be directly measured).

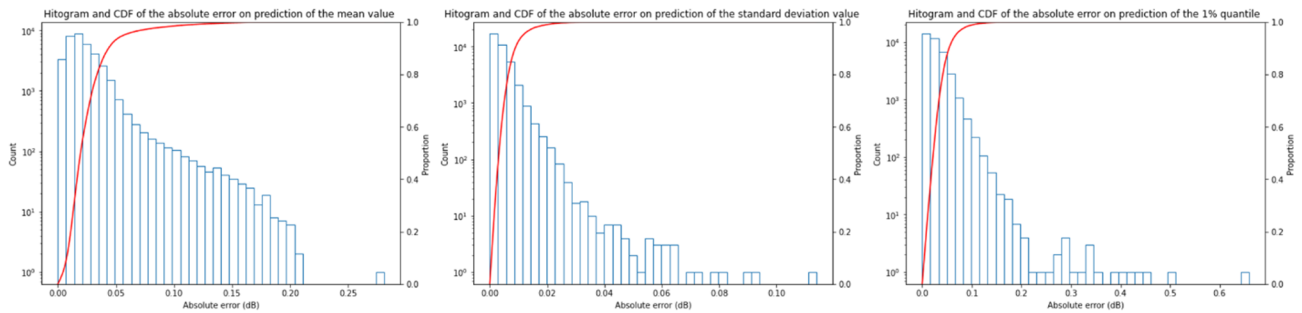
### 5.2 Result for probability density

Different machine learning metamodels have been tested of which we can cite gradient boosting, multilayer perceptron, support vector machine regression and Gaussian process<sup>24</sup>. Best results were obtained for Gaussian process<sup>25</sup>.

15% of our database was used for a training set and the 85% remaining for a test set. Using  $r_0$ ,  $\bar{h}$  and  $\bar{v}$  as inputs of the model we were able to assess  $\mu$  with a coefficient of determination of 0.9994,  $\beta$  with a coefficient of determination of 0.9992 and  $\sigma_{\chi_{AP}}$  with 0.9969.



**Figure 9 :** Example of PDF computed from SAOST (blue) time series and PDF computed from our model (yellow) with, from left to right:  $L_\varphi$ ,  $\rho_I$  and the received optical power.



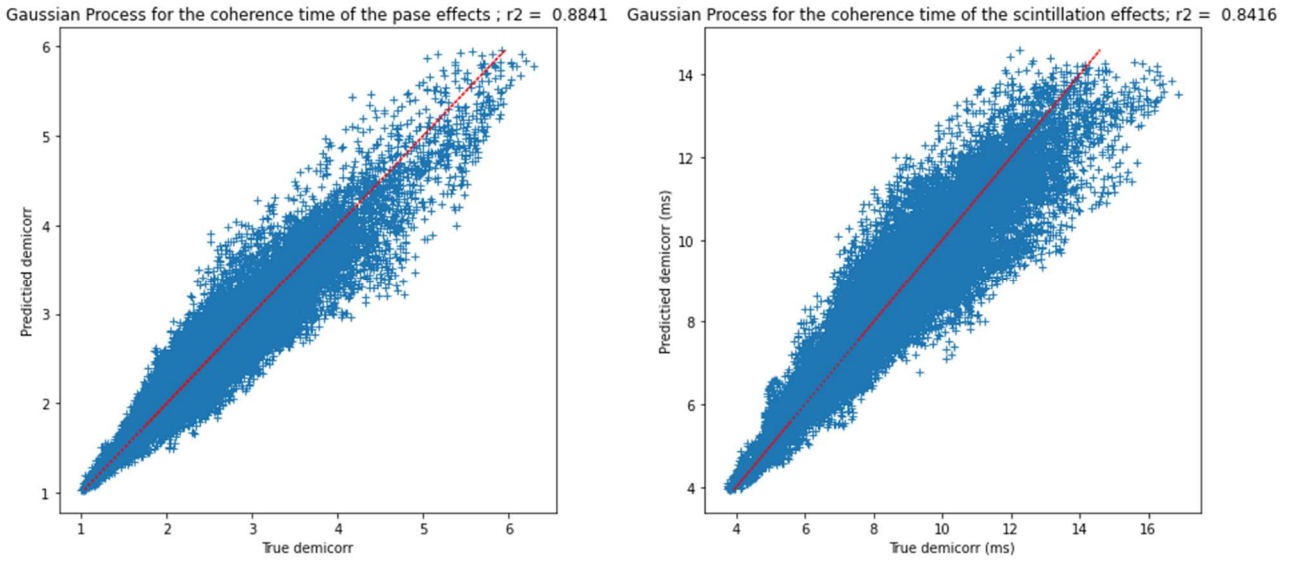
**Figure 10 :** Histogram and CDF of the absolute error on prediction of the mean, the standard deviation and the 1% quantile.

In order to characterize the goodness of our estimation on the received optical power's statistic, the absolute error made on the mean and standard deviation of the reconstructed PDF of the received optical power are computed. We also look at the absolute error on the 1% quantile of the received optical power (see red dashed line on figure 9) as we want a faithful reproduction of the tail of the distribution. Statistics on the absolute error made on these values on our 37k profiles can be seen in figure 10. On all profiles, the prediction error on the value of the 1% quantile is less than 0.7 dB.

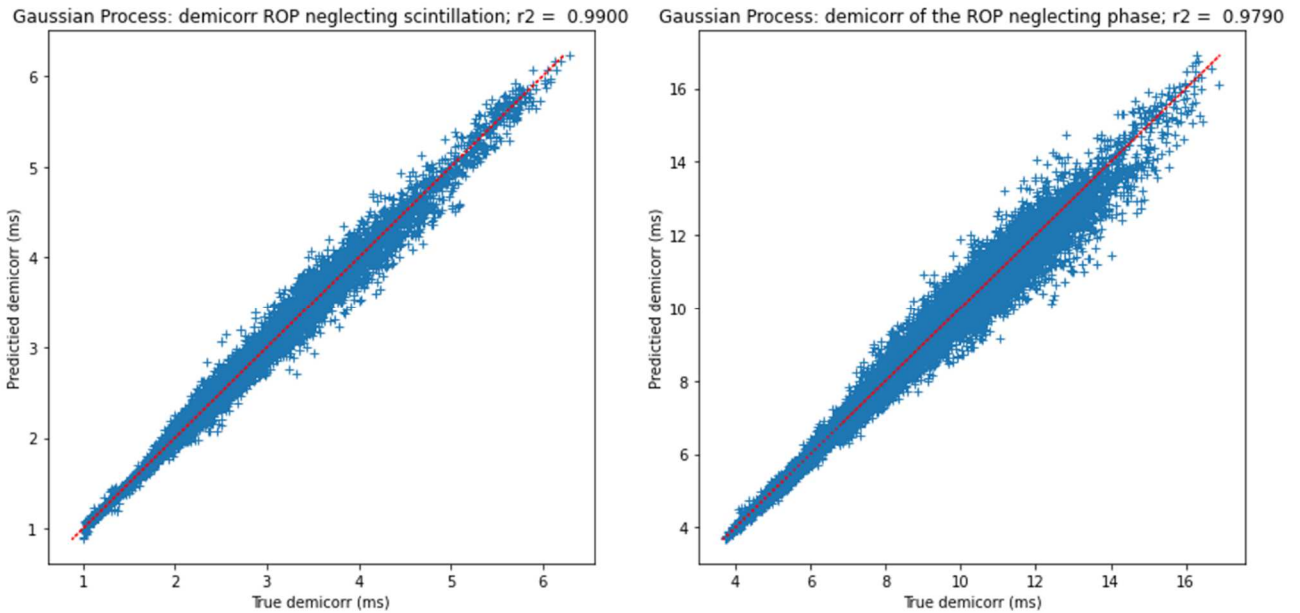
### 5.3 Result for temporal content

Temporal correlation of the power attenuation will affect the duration on which interleaving shall be performed to mitigate fading losses. We investigate here the possibility to estimate the half correlation of  $L_\varphi$  and  $\rho_I$  using only  $r_0$ ,  $\bar{h}$  and  $\bar{v}$ . Once again, best results were obtained with the use of a Gaussian process metamodel. On figure 11 where the estimated value is compared to the expected one, we can see that the coefficient of determination is 0.88 for the phase demi-correlation and 0.84 for the scintillation<sup>25</sup>.





**Figure 11:** Prediction of the half-correlation time using Gaussian process on inputs  $r_0$ ,  $\bar{h}$  and  $\bar{v}$ , x axis shows the true value and y the predicted one. On the left the half-correlation neglecting scintillation effects and on the right neglecting phase effects.



**Figure 12:** Prediction of the half-correlation time using Gaussian process on inputs  $r_0$ ,  $\bar{h}$ ,  $\bar{v}$  and information on the PSD of scintillation for a 5cm pupil. Same other parameters as for figure 11.

This low score is due to the lack of temporal information found in  $r_0$ ,  $\bar{h}$  and  $\bar{v}$ . One way to recover the missing temporal information is to add a measurement of the temporal power spectral density (PSD) of the scintillation to our inputs. As inputs of our model we added 5 moments of the PSD of scintillation measured by a 5 cm pupil instrument, these results were simulated using the formulas of Shen<sup>26</sup>. As seen on figure 12, adding information on the PSD to our model enabled much better results with coefficient of determination of 0.99 for the phase and 0.98 for the scintillation<sup>25</sup>.

## 6. POWER LINK MARGIN ESTIMATION

### 6.1 Inputs and limitation

The main challenge in the estimation of the power margin is to describe the impact of using an interleaver. Indeed, the interleaver acts as a convolutional filter on the mutual information time series seen by the receiver and this non-linear operation cannot be translated in simple terms by knowing only the half correlation and the probability density. In addition, we are looking for a valid method regardless of the size of the interleaver.

The best way to achieve our goal is to generate a stochastic process equivalent to any time series of ROP obtained with the same exact atmospheric and instrumental conditions. By being stochastically equivalent to the ground truth, the generated series is sharing the same finite-dimensional law and thus the same power margin.

We know that a Gaussian process is defined entirely by its auto covariance function<sup>27</sup>. However, this is not necessarily true for all types of stochastic processes and there is no reason it should be the case for our process, product of two non-Gaussian stochastic processes (phase and scintillation).

### 6.2 Downlink budget excluding turbulence

The optical downlink signal between the on-board terminal described in Berceau et al.<sup>28</sup> and the ESA OGS described section 3 is considered. The communication chain equipment (Laser Communication Electronic, Modems and Laser Power Electronics) considered are the one of Poulenard et al.<sup>1</sup>.

### 6.3 Method for link margin estimation

Even though it is very easy to generate a Gaussian process, there are very few methods for generation of stochastic process with any given marginal distribution. An intuitive method was developed by Tsoukalas<sup>29</sup> and comes with a R library anySim<sup>30</sup>. It is based on Nataf's joint distribution model where the distribution of a correlated random variable is mapped to a Gaussian distribution via its inverse cumulative distribution function.

Using this library and the previously estimated distribution and half-correlation, we generated time series of coupled flux with phase effects only and with scintillation effects only. Combining the two series with the downlink budget given in section 6.2, we obtained time series of ROP using no prior knowledge on the profiles but  $r_0$ ,  $\bar{h}$ ,  $\bar{v}$  and  $\sigma_\chi^2$ .

AnySim requires an autocorrelation function; the simplest interpolation was made using an exponential decay matching the predicted half correlation time.

### 6.4 Result for link margin estimation

The generation of random time series of 5 seconds with a 4 kHz sampling, i.e. 20.000 occurrences to ensure good statistical convergence, is very memory intensive. In the following, we limited our simulation to 5.000 pairs of profiles randomly selected in our database.

Generation of random time series was run two times, once using the half-correlation time computed using only  $r_0$ ,  $\bar{h}$ ,  $\bar{v}$  and  $\sigma_\chi^2$  and once with the moments of the temporal PSD added into the prediction of the half correlation time. Following results of power margin are given for an interleaver of 100ms and the MODCOD 9/10 with a required mutual information of 0.94 and MODCOD 3/10 with a required mutual information of 0.12.

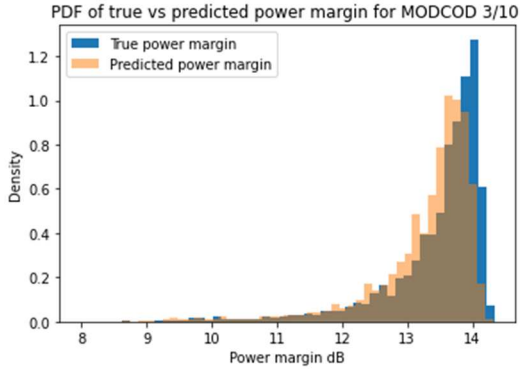


Figure 13: PDF true power margin vs predicted using only  $r_0$ ,  $\bar{h}$ ,  $\bar{v}$  and  $\sigma_\chi^2$  for MODCOD 3/10- interleaver 100ms.

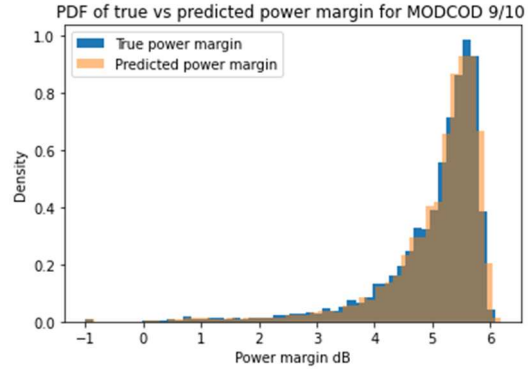


Figure 14 : PDF true power margin vs predicted using only  $r_0$ ,  $\bar{h}$ ,  $\bar{v}$  and  $\sigma_\chi^2$  for MODCOD 9/10- interleaver 100ms.

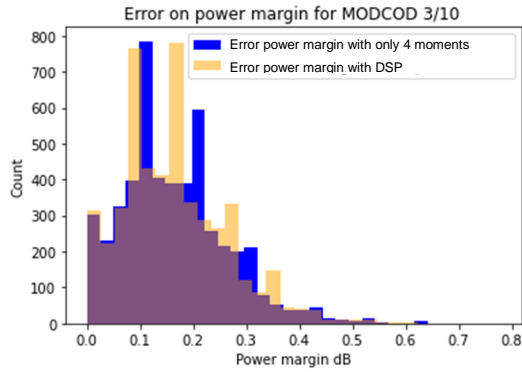


Figure 15: Distribution of power margin error [dB] for MODCOD 3/10 - interleaver 100ms. In blue using the model with  $r_0$ ,  $\bar{h}$ ,  $\bar{v}$  and  $\sigma_\chi^2$  in yellow when the PSD of scintillation is added.

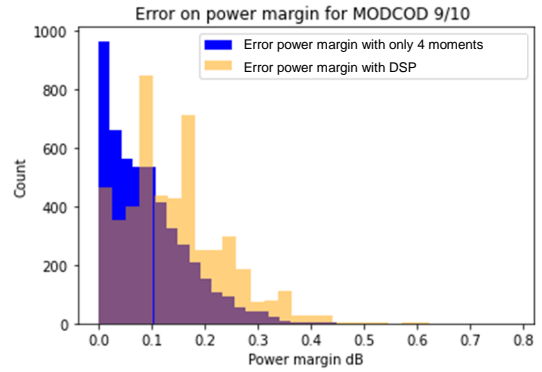


Figure16: Distribution and cumulative distribution of power margin error [dB] for MODCOD 9/10 - interleaver 100ms. In blue using the model with  $r_0$ ,  $\bar{h}$ ,  $\bar{v}$  and  $\sigma_\chi^2$  in yellow when the PSD of scintillation is added.

On figure 13 and 14 we compared the probability density function of the “true” power margin (after SAOST) and the estimated power margin using only  $r_0$ ,  $\bar{h}$ ,  $\bar{v}$  and  $\sigma_\chi^2$  for MODCOD 3/10 and 9/10. When computing the error made on the power margin we see that it never exceeds 0.7 dB for MODCOD 3/10 and 0.6 dB for MODCOD 9/10 (in blue figure 13 and 14).

When the temporal PSD of scintillation is added to the input (yellow on figure 15 and 16), no significant improvement is observed, we even have a maximal error a tiny bit bigger with 0.8dB in both MODCOD case that might be attributed to the convergence error of the estimation due to the limited length of the considered time series.

In this specific case of 100ms interleaver, significantly bigger than the half correlation time (ranging in the 10<sup>th</sup> of milliseconds), better prediction of the half correlation does not translate into better prediction of the power margin which confirms the relevance of the interleaving process.

## 6.5 Model summary and sensitivity analysis

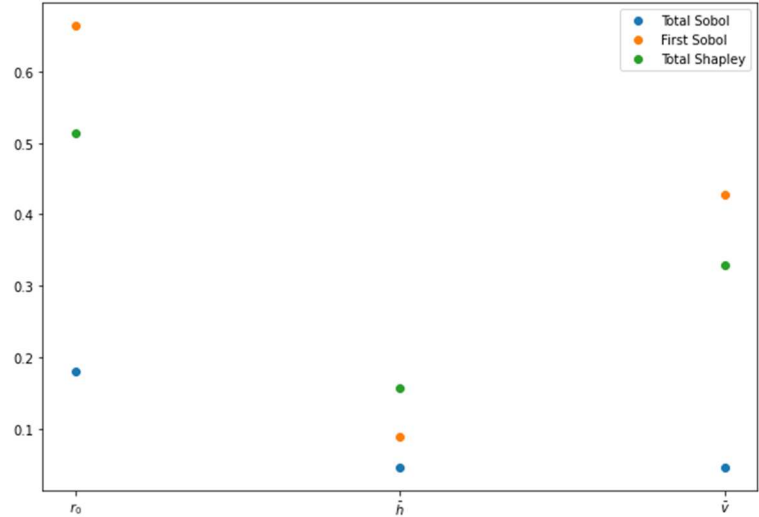
We proposed a model that takes as input 4 moments ( $r_0$ ,  $\bar{h}$ ,  $\bar{v}$  and  $\sigma_\chi^2$ ) and gives as output a power margin estimation with an error below 1dB. This model includes intermediate steps of estimating PDF and autocorrelation time as well as random time series generation.

An interesting thing to look at is which one of the inputs has the biggest impact on the power margin estimation. Knowledge on the impact of each of the moments would enable us to improve the precision on the measurement of a given moment or on the contrary to reduce the required resolution. To do so we use classical tools of sensitivity analysis that are first and total order Sobol indices as well as Shapley values.

The Sobol indices<sup>31</sup> are also known as variance-based sensitivity analysis. The variance of the output of our model is decomposed into fractions, which can be attributed to each of the moments used as inputs. The first order Sobol indices characterize the contribution of a given input to the output variance and the total Sobol indices measure the contribution to the output variance of the studied input, including all variance caused by its interactions, with other input variables.

Sobol indices are good to determine the sensitivity of the model caused by a specific input but works under the assumption of independence between inputs, their interpretation becomes hazardous in the case of correlated inputs. In our case,  $r_0$ ,  $\bar{h}$  and  $\bar{v}$  are strongly correlated being all moments of the same profiles. To deal with correlated inputs, methods have been developed around Shapley indices<sup>32,33</sup>.

On figure 17 we can see that  $r_0$  has the biggest impact on the predicted value of the power margin, accounting for more than half of the variance according to Shapley values. This result was expected, as in the downlink case,  $r_0$  intervenes in every contributors to the adaptive optics error budget considered here. We expect that  $\bar{h}$ , that has little impact here, will have a bigger impact on the uplink case as the anisoplanatic error (related to  $\bar{h}$ ) will become an important term in the error budget.



**Figure 17: Sensitivity analysis to describe the impact of  $r_0$ ,  $\bar{h}$  and  $\bar{v}$  on the prediction of the power margin.**

## 7. CONCLUSION AND NEXT STEP

We developed and studied a metamodel that enables us to obtain the optical power margin from a restricted set of integrated parameters. We demonstrated that the availability of an optical link can be described by a small number of integrated parameters, that can be measured by rather simple instrumentation. These results constitute a very encouraging step forward toward the development of future ground-stations networks and of their intelligence for automatic routing. Ultimately, the model will be extended to the uplink case. The impact of AO correction parameters and of typical measurements error on the integrated parameters on the robustness of the approach shall also be investigated. Further investigations will also be conducted to understand the impact of a deviation of the instrument aiming angle with respect to the link line of sight on which the optical link will effectively take place. Finally, tests on real world data will be conducted to demonstrate experimentally the relevance of the approach and investigate its practical limits.

## ACKNOWLEDGEMENTS

The authors would like to thank James Osborn from Durham University for the database without which this work would not have been possible. They would also like to thank Miratlas for their support in this project.

## REFERENCES

- [1] Poulénard et al., "10 Gbauds digital optical link and analog link from/to geostationary satellite," ICSO 2022.
- [2] [https://www.miratlas.com/products\\_services.html#ism](https://www.miratlas.com/products_services.html#ism)
- [3] Osborn J, Wilson RW, Sarazin M, et al (2018), "Optical turbulence profiling with Stereo-SCIDAR for VLT and ELT". *Monthly Notices of the Royal Astronomical Society* 478:825–834. <https://doi.org/10.1093/mnras/sty1070>
- [4] Masciadri E, Vermin J, Bougeault P (1999) 3D mapping of optical turbulence using an atmospheric numerical model: I. A useful tool for the ground-based astronomy. *Astron Astrophys Suppl Ser* 137:185–202. <https://doi.org/10.1051/aas:1999474>.
- [5] Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J., Nicolas, J., Peubey, C., Radu, R., Schepers, D., Simmons, A., Soci, C., Abdalla, S., Abellan, X., Balsamo, G., Bechtold, P., Biavati, G., Bidlot, J., Bonavita, M., et al., "The ERA5 global reanalysis," *Q. J. R. Meteorol. Soc.* 146(730), 1999–2049 (2020).
- [6] Osborn, J. and Sarazin, M., "Atmospheric turbulence forecasting with a general circulation model for Cerro Paranal," *Mon. Not. R. Astron. Soc.* 480(1), 1278–1299 (2018).
- [7] University, C. for A. I.-D., "Turbulence Forecasts," <<https://www.dur.ac.uk/cfai/sitecharacterisation/forecasts/>> (29 June 2022 ).
- [8] Vedrenne, N., Conan, J.-M., Petit, C. and Michau, V., "Adaptive optics for high data rate satellite to ground laser link," *Proc SPIE 9739 Free-Space Laser Commun. Atmospheric Propag. XXVIII 97390E 9739* (2016).
- [9] Conan, J.-M., Montmerle-Bonnefois, A., Védrenne, N., Lim, C. B., Petit, C., Michau, V., Velluet, M.-T., Sauvage, J.-F., Meimon, S., Robert, C., Montri, J., Mendez, F., Perrault, P., Artaud, G. and Benammar, B., "Adaptive Optics for GEO-Feeder Links: from Performance Analysis via Reciprocity Based Models to Experimental Demonstration," *COAT-2019 - Workshop Commun. Obs. Atmospheric Turbul. Charact. Mitig.*, ONERA, Châtillon, France (2019).
- [10] Canuet, L., Védrenne, N., Conan, J.-M., Petit, C., Artaud, G., Rissons, A. and Lacan, J., "Statistical properties of single-mode fiber coupling of satellite-to-ground laser links partially corrected by adaptive optics," *J. Opt. Soc. Am. A* 35(1), 148 (2018).
- [11] Roddier, N. A., "Atmospheric wavefront simulation using Zernike polynomials," *Opt. Eng.* 10(29) (1990).
- [12] Conan, J.-M., Rousset, G. and Madec, P.-Y., "Wave-front temporal spectra in high-resolution imaging through turbulence," *JOSA A* 12(7), 1559–1570 (1995).
- [13] Canuet, L., "Fiabilisation des transmissions optiques satellite-sol," thesis (2018).
- [14] Tatarski, V. I., [Wave Propagation In a Turbulent Medium], Dover Publications, Inc. New York (1961).
- [15] Vetelino, F. S., Young, C., Andrews, L. and Reclons, J., "Aperture averaging effects on the probability density of irradiance fluctuations in moderate-to-strong turbulence," *Appl Opt* 46(11), 2099–2108 (2007).
- [16] Robert, C., Conan, J.-M., Michau, V., Renard, J.-B., Robert, C. and Dalaudier, F., "Retrieving parameters of the anisotropic refractive index fluctuations spectrum in the stratosphere from balloon-borne observations of stellar scintillation," *J Opt Soc Am A* 25(2), 379–393 (2008).
- [17] D. L. Fried, "Optical Resolution Through a Randomly Inhomogeneous Medium for Very Long and Very Short Exposures", *J. Opt. Soc. Am.*, vol. 56, no 10, Art. no 10, oct. 1966, doi: 10.1364/JOSA.56.001372
- [18] A. Tokovinin, "From differential image motion to seeing", *The Publications of the Astronomical Society of the Pacific*, vol. 114, p. 1156-1166, 2002.
- [19] Roddier, F., "V The Effects of Atmospheric Turbulence in Optical Astronomy," E. Wolf, Ed., Elsevier, 281–376 (1981).
- [20] D. L. Fried, "Anisoplanatism in adaptive optics", *J. Opt. Soc. Am.*, 72(1) :52–61, Jan. 1982.
- [21] G. C. Loos et C. B. Hogge, "Turbulence of the upper atmosphere and isoplanatism", *Appl. Opt.*, AO, vol. 18, no 15, p. 2654-2661, août 1979, doi: 10.1364/AO.18.002654.
- [22] J. M. Beckers, "A Seeing Monitor for Solar and Other Extended Object Observations", *Experimental Astronomy*, vol. 12, p. 1-20, 2001.
- [23] F. Roddier, J. M. Gilli, et G. Lund, "On the origin of speckle boiling and its effects in stellar speckle interferometry", *jop*, vol. 13, no 5, Art. no 5, 1982.
- [24] Rasmussen, C. E. and Williams, C. K. I., [Gaussian Processes for Machine Learning], MIT press Cambridge, MA (2005).
- [25] Klotz, E., "Assessment of adaptive optics corrected optical links statistics from integrated turbulence parameters," *Prepr. Soon Available* (2022).

- [26] Shen, H., Yu, L. and Fan, C., “Temporal spectrum of atmospheric scintillation and the effects of aperture averaging and time averaging,” *Opt. Commun.* 330, 160–164 (2014).
- [27] Bishop, C. M. and Nasrabadi, N. M., [Pattern recognition and machine learning], Springer (2006).
- [28] Berceau et al., “Space Optical Instrument for GEO-Ground Laser Communications,” ICSOS 2022.
- [29] Tsoukalas, I., “Modelling and simulation of non-Gaussian stochastic processes for optimization of water-systems under uncertainty” (2018).
- [30] Tsoukalas, I., Kossieris, P. and Makropoulos, C., “Simulation of Non-Gaussian Correlated Random Variables, Stochastic Processes and Random Fields: Introducing the anySim R-Package for Environmental Applications and Beyond,” *Water* 12(6), 1645 (2020).
- [31] Sobol’, I. M., “Global sensitivity indices for nonlinear mathematical models and their Monte Carlo estimates,” *Math. Comput. Simul.* 55(1), 271–280 (2001).
- [32] Song, E., Nelson, B. L. and Staum, J., “Shapley Effects for Global Sensitivity Analysis: Theory and Computation,” *SIAMASA J. Uncertain. Quantif.* 4(1), 1060–1083 (2016).
- [33] Iooss, B. and Prieur, C., “Shapley effects for sensitivity analysis with correlated inputs: comparisons with Sobol’ indices, numerical estimation and applications,” *Int. J. Uncertain. Quantif.* 9(5), 493–514 (2019).