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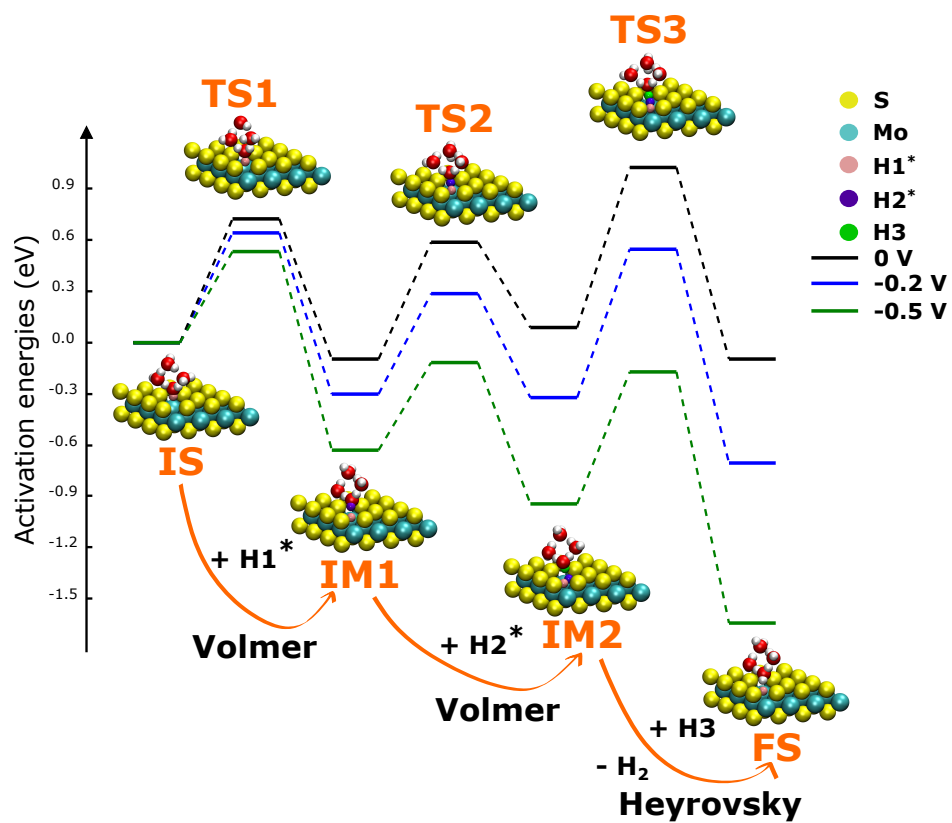
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# Graphical Abstract

## Potential and Support-Dependent Hydrogen Evolution Reaction Activation Energies on Sulfur Vacancies of MoS<sub>2</sub> from GC-DFT

Nawras Abidi, Audrey Bonduelle-Skrzypczak, Stephan N. Steinmann

### HER on sulfur vacancy of MoS<sub>2</sub>



## Highlights

### **Potential and Support-Dependent Hydrogen Evolution Reaction Activation Energies on Sulfur Vacancies of MoS<sub>2</sub> from GC-DFT**

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- Potential-dependent activation energies identify HER on S-vacancies of MoS<sub>2</sub> as a Volmer-Heyrovsky mechanism in presence of a spectator H\*
- The defect density has little influence on the S-vacancy based HER kinetics
- Support and nanosheet size effects are found to be small, suggesting large specific surface area carbon supports are adequate for MoS<sub>2</sub>

# Potential and Support-Dependent Hydrogen Evolution Reaction Activation Energies on Sulfur Vacancies of MoS<sub>2</sub> from GC-DFT

Nawras Abidi<sup>a</sup>, Audrey Bonduelle-Skrzypczak<sup>b</sup>, Stephan N. Steinmann<sup>a</sup>

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## Abstract

We present a detailed mechanistic study of HER at the sulfur vacancy  $V_S$  of 2H-MoS<sub>2</sub>. We evaluate the Volmer, Tafel, and Heyrovsky transition states for the different possible reaction steps, determining the activation energy as a function of the electrochemical potential via grand-canonical density functional theory. The results show that the Volmer and Heyrovsky steps depend on the electrochemical potential and the activation energies decrease for more negative potentials, while this is not the case for the Tafel step, for which the activation energy is constant. From the activation energies at -0.2 V vs SHE, it can be concluded that during HER on  $V_S$  a first hydrogen atom is adsorbed as a spectator via a Volmer step. Then, the catalytic cycle consists of a Volmer and a Heyrovsky step, with the latter being rate determining. In addition, we investigate for the first time the effect of a conductive support on the HER activity of these sulfur vacancies. Our results show that copper, gold and graphite supports have little effects on the activation energies of all steps. Hence, we conclude that cheap, acid-stable, high-surface area carbon supports are well suited for MoS<sub>2</sub>-based HER catalysts.

*Keywords:* MoS<sub>2</sub> basal plane, Sulfur vacancies, Hydrogen evolution reaction, Grand-canonical DFT, Electrochemical potential, Activation energies, Support effect

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## 1. Introduction

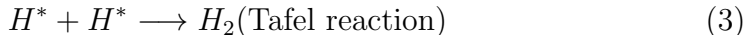
Hydrogen is expected to be one of the future clean energy vectors. Even though the electrochemical formation of  $H_2$ , the hydrogen evolution reaction (HER), is well studied for many decades, open questions still remain, as the detailed mechanism depends on the electrocatalyst and the reaction conditions.[1, 2] In acidic conditions, HER is, in general, a two-step process: The first step is the adsorption of protons from the acidic solution to form adsorbed H, denoted as  $H^*$ . This step is the Volmer reaction:



The second step depends on the electrocatalyst and the current densities. It can be a Heyrovsky step:



where a solvated proton reacts with  $H^*$  to form  $H_2(g)$ . Both the Volmer and the Heyrovsky reactions are "electrochemical" elementary steps, as they are formally coupled proton-electron transfers. Alternatively,  $H_2$  can also be formed via a formally "chemical" step: the recombination of two  $H^*$  to form  $H_2$ , known as the Tafel reaction:



The rate of the overall reaction depends critically on the free energy of adsorption of the intermediate ( $H^*$ ) on the catalyst surface ( $\Delta G_{H^*}$ ), but the activation energies for the three elementary steps are responsible for the actual mechanism and the experimentally measured Tafel slopes (i.e., the relation between the overpotential and the current density).[3] According to the Sabatier principle (volcano diagram),  $\Delta G_{H^*}$  should be close to 0 eV.[4]

Among the noble-metal free HER catalysts, transition metal dichalcogenides,[5, 6, 7, 8, 9] and in particular  $MoS_2$ ,[10, 11, 12, 13, 14] are most promising for large-scale, low-cost applications in acidic medium, while phosphides are promising in alkaline solutions [15]. The development of these catalyst is one of the archetypal examples of a successful collaboration between computations and experiments.[4, 14] Several studies have shown that the basal plane and edges of the most stable  $MoS_2$  polymorph (2H- $MoS_2$ )[16] have contrasting electrocatalytic properties: the perfect 2H- $MoS_2$  basal plane is catalytically inert but the edges are highly active.[17]

To make 2H-MoS<sub>2</sub> catalytically more reactive, one can resort to methods of generating active sites by phase modifications,[18] doping [19, 20, 21, 22], or by exposing more edges [23, 24]. Furthermore, several studies have shown that point defects (and in particular the sulfur vacancies that have the lowest formation energy) also significantly affect the electrocatalytic properties of MoS<sub>2</sub>. [25, 26, 27] The presence of these vacancy defects in MoS<sub>2</sub> effectively exposes the undercoordinated Mo atoms that act as active sites for the hydrogen evolution reaction (HER) as demonstrated by joint computational/experimental work. [28, 29, 30, 31, 32] These studies leave little doubt that such point defects are active sites for HER on MoS<sub>2</sub>, even though the precise contribution between edge and basal plane activity is certainly highly dependent on the synthesis conditions and catalyst formulation. Experimentally, various efforts have been made to generate and precisely control the formation of sulfur defects on MoS<sub>2</sub> to catalyze HER, exploiting various techniques such as electrochemical desulfurization, chemical etching, oxygen/hydrogen plasma treatment, and high-temperature hydrogen annealing [16, 28, 33, 29, 31, 32, 34].

Since MoS<sub>2</sub> is a semiconductor, the HER activity of the basal plane is also limited by the electron transfer from the electrode to the catalyst.[35] Therefore, depositing MoS<sub>2</sub> nanosheets on a conductive support is an attractive strategy to improve the electron transfer kinetics and thus of the overall HER, in analogy to ideas in photocatalytic alkaline water splitting [36]. Furthermore, the support imparts mechanical stability to the catalyst and augments its dispersion, while limiting the size of the nanosheets during the synthesis. This is particularly beneficial for catalysts with high edge-site activities. The support can, however, also affect the electronic structure and coordination environment of MoS<sub>2</sub>. [37, 38] Strong catalyst-support interactions have been found to reduce  $\Delta G_{H^*}$  and thus the support allows to tune the edge activity. Experimentally, a weakly interacting graphite support was demonstrated to lead to an increase in HER activity for the edges.[39] The high activity of Au nanostars encapsulated in a MoS<sub>2</sub> shell was rationalized by a slight improvement of the S-vacancy activity on a more strongly interacting Au support, but the the high electrical conductivity of the support and the high density of active sites might also contribute.[40]

While the use of the thermodynamics of HER, i.e.,  $\Delta G_{H^*}$ , as a first approximation to the HER activity of MoS<sub>2</sub> was successful [41, 42, 43, 44], it does not provide insight into the reaction kinetics and the precise mechanisms involved. Activation energies  $\Delta E_{act}$  are required to quantitatively

compute reaction rates and to elucidate the reaction mechanism. However, the computation of  $\Delta E_{act}$  is not straightforward for electrocatalytic systems. Several methods have been developed to determine the transition states and activation energies, but no generally accepted, easily applicable scheme exists.[45] In our experience, grand canonical density functional theory (GC-DFT)[46, 47, 48] is the most versatile approach to determine potential-dependent activation energies. Previously, we have applied GC-DFT to elucidate the atomistic origin of the electrochemical promotion of catalysis.[49, 50, 51]

Using simplified models, i.e., that do not explicitly take the electrochemical potential into account, Li et al. [52] investigated the HER mechanism on MoS<sub>2</sub> with explicit consideration of the solvent effect. The results show that the solvent H<sub>2</sub>O can significantly affect the barriers of HER. They proposed that HER at the MoS<sub>2</sub> (perfect) basal plane proceeds via the Volmer-Heyrovsky mechanism, and that the overall rate is determined by the Volmer step with an activation energy of 0.65 eV, which is at odds with the catalytic inertness of this surface. The HER mechanism at the defective basal plane of 2H-MoS<sub>2</sub>, has only rarely been studied. Still using a simplified model (H<sub>3</sub>O<sup>+</sup> is used as the reactant), Kong et al.[53] investigated the effects of sulfur vacancies distribution on the HER mechanism. Their computations suggest that HER proceeds via the Volmer-Heyrovsky mechanism with the Heyrovsky reaction as a rate-determined step (RDS) with an energy barrier of 0.8 eV on isolated defects. The study by Huang et al. [54] is the only one to date that has explicitly considered the effects of electrochemical potential on the activation energies of sulfur defects on MoS<sub>2</sub>. Using grand-canonical potential kinetics, they conclude that the rate-determining step in both acidic and basic environments is the Volmer reaction, in which the second hydrogen is adsorbed from solution. This second hydrogen at the sulfur vacancy site is most active for the hydrogen evolution reaction. In acidic media, H<sub>2</sub> is generated via the Heyrovsky step with a barrier of 0.38 eV at -0.5 V. However, they used H<sub>3</sub>O<sup>+</sup> as the reactant, which is likely[52] to be under-solvated by their implicit solvent.

Here, we investigate the kinetics of the HER on sulfur vacancies on the basal plane of 2H-MoS<sub>2</sub>. Using GC-DFT we determine the activation energies governing the HER as a function of the electrochemical potential, using a hydronium ion (H<sub>3</sub>O<sup>+</sup>), microsolvated by three water molecules, giving H<sub>9</sub>O<sub>4</sub><sup>+</sup>. In addition, since in experimental setups a conductive support is necessary, we assess the effect of different supports (Au, Cu, and graphite)

compared to pure MoS<sub>2</sub> on these activation energies barriers.

## 2. Computational methods and models

### 2.1. Computational details

The computations were carried out relying on density functional theory with the Vienna ab initio simulation package (VASP) [55]. We adopted the PBE-dDsC [56, 57] density functional, which belongs to the family of the Generalized Gradient Approximation and is supplemented with a density-dependent dispersion correction.[58] The projector augmented wave method [59] is used to treat the ion–electron interaction. The electronic density was calculated with a 500 eV cutoff energy. Due to the size of the super cells,  $\Gamma$  point sampling of the Brillouin zone was found to be sufficient. The precision setting of VASP is set to "accurate". The convergence criterion for the self-consistency process is set to  $10^{-5}$  eV for the optimization of the wave function. The maximum forces are converged to 0.025 eV/Å during the geometry optimization. The transition states (TS) were identified by the climbing image nudged-elastic band[60] and the dimer method [61]. Upon geometric convergence, a frequency analysis was performed to confirm the located TS and to verify the number of imaginary frequencies: one/two for as/symmetric slabs.

### 2.2. Energy evaluations

We are interested in evaluating the HER pathway and reaction kinetics of 2H-MoS<sub>2</sub>. Therefore, we compute the hydrogen adsorption energy and the activation energy to determine the barrier for each step of the HER reaction.

For H<sub>2</sub> we use the following approximation for its free energy:

$$G_{\text{H}_2} = E_{\text{H}_2} - TS \tag{4}$$

where  $T$  is the room temperature (298.15 K) and  $S$  is the contribution of the rotational translational entropy at standard pressure.

To properly account for the electrochemical conditions at a given applied voltage, we use the grand canonical DFT approach. The number of electrons in the system is varied to probe the influence of the electrochemical potential. Counterions are included in the implicit solvation model via the linearized Poisson-Boltzmann equation to neutralize the system and provide a combined solvent-slab free energy.[62] This approach is implemented in

VASPsol[63], and we obtained a relationship between the grand canonical energy and the electrochemical potential fitted to an analytical formula that is a parabola for conductors.[64] An important parameter for the electrochemical computations is the iso-density value defining the solvent-solute interface. Following our previous studies [64], we have chosen  $0.00025 \text{ e}^-/\text{\AA}^{-3}$ , which is one tenth of the default value and ensures the absence of implicit solvent between weakly bonded layers such as graphite and  $\text{MoS}_2$ .

The activation energy is calculated as a function of electrochemical potential by:

$$\Delta E_{act}(U) = E_{IS}(U) - E_{TS}(U) \quad (5)$$

Where  $E_{IS}(U)$  and  $E_{TS}(U)$  are the potential-dependent energies of the initial state and the transition state of the various steps obtained from the grand canonical energy as a function of the electrochemical potential.

As shown in Fig. S5, the potential of zero charge (corresponding to the neutral "conventional" surfaces) changes dramatically between the reactant ( $\text{H}_9\text{O}_4^+ + \text{e}^-$ ) and the transition state, where the electron is no longer located on the surface, but is partially used to form the surface-H bond. Therefore, relying on the constant potential approach, i.e., GC-DFT, is key for these proton-electron transfer transition states. GC-DFT also allows to determine the charge transfer coefficient which determines the sensitivity of the activation energy on the applied potential, known from the empirical Butler-Volmer equation. Indeed, in the GC-DFT framework, the charge transfer coefficient is simply obtained as the slope of the activation energy reported against the electrochemical potential. This allows an atomistic understanding of this otherwise empirical parameter.

### 2.3. Structural Models

The basal plane of  $\text{MoS}_2$  consists of two hexagonal planes of sulfur (S) atoms and an intercalated hexagonal plane of molybdenum (Mo) atoms bonded to the S atoms. In the first part of our study, we work with the supercell of  $\text{MoS}_2$   $p(4 \times 4)$  with three layers, where the bottom layer is fixed. A vacuum of about  $30 \text{ \AA}$  is added in the  $z$ -direction (See Fig. 1a). A sulfur vacancy ( $V_s$ ) is created by removing one sulfur atom. To model the proton/electron transfer reactions - Volmer and Heyrovsky reactions - we use a micro-solvated structure of the hydronium ( $\text{H}_3\text{O}^+$ ) ion, where it is stabilized by three additional water molecules to form a  $\text{H}_9\text{O}_4^+$  complex. This approach

is inspired by the work of Skulasson,[65] but in our hands the original  $\text{H}_9\text{O}_4^+$  structure underwent a stabilizing rearrangement, forming the four-membered ring shown in Fig. 1a. For the Tafel step, since there is no electron transfer, only a surface slab and adsorbed hydrogens need to be included in the atomic model.

For assessing the support effect, we prepare a  $\text{Mo}_{12}\text{S}_{24}$  triangular nanosheet from  $\text{MoS}_2$  with reconstructed Mo-edges represented in Fig. 1b in order to avoid the strain of the support or  $\text{MoS}_2$ . Gronborg et al.[66] determined with scanning tunneling microscopy (STM) that  $\text{MoS}_2$  nanosheets larger than 2 nm adopt a perfect triangular shape terminated by the Mo edges which is more stable than the S-edge under the sulfiding conditions of the synthesis. Then, we add the  $\text{H}_9\text{O}_4^+$  structure based on the optimized geometries from the periodic  $\text{MoS}_2$  surfaces. The nanosheet is positioned on different supports: Molybdenum disulfide ( $\text{MoS}_2$ ) as a reference, gold (Au), copper (Cu) and graphite (C). (See Fig. 2) The gold support is chosen because of its use in previous experimental studies,[11, 67] copper as a cheaper metallic analogue and graphite is our model for conducting high-area carbon supports that are widely studied in the fuel-cell literature and used on an industrial scale. Different positions of the nanosheet on the support were tested using ChemCat[68]: Top, Bridge, Hollow, fcc and hcp. The relative energy values are shown in Table S1. In all cases, the "top" position is the most stable.

To properly determine the workfunction of the various systems, we use symmetric slabs. Even though the use of asymmetric slabs has been proposed in the context of GC-DFT,[69] for the highly polar initial state and TS involving a proton, we have observed in Fig. S1 artifacts for the asymmetric potential dependent models. We suspect that these observations are due to the implicit solvent that is present on both surfaces, in contrast to ref. [69]. For the systems with support and only in the case of the first Volmer step ( $V^I$ ), we realize plane symmetrization for the  $\text{MoS}_2$  support and point symmetrization for Cu and Au. Graphite was not considered for these GC-DFT computations as the system is computationally excessively heavy. Our final symmetric slabs are made of five layers of the support, with the central layer frozen in its bulk position. (see Fig. S2).

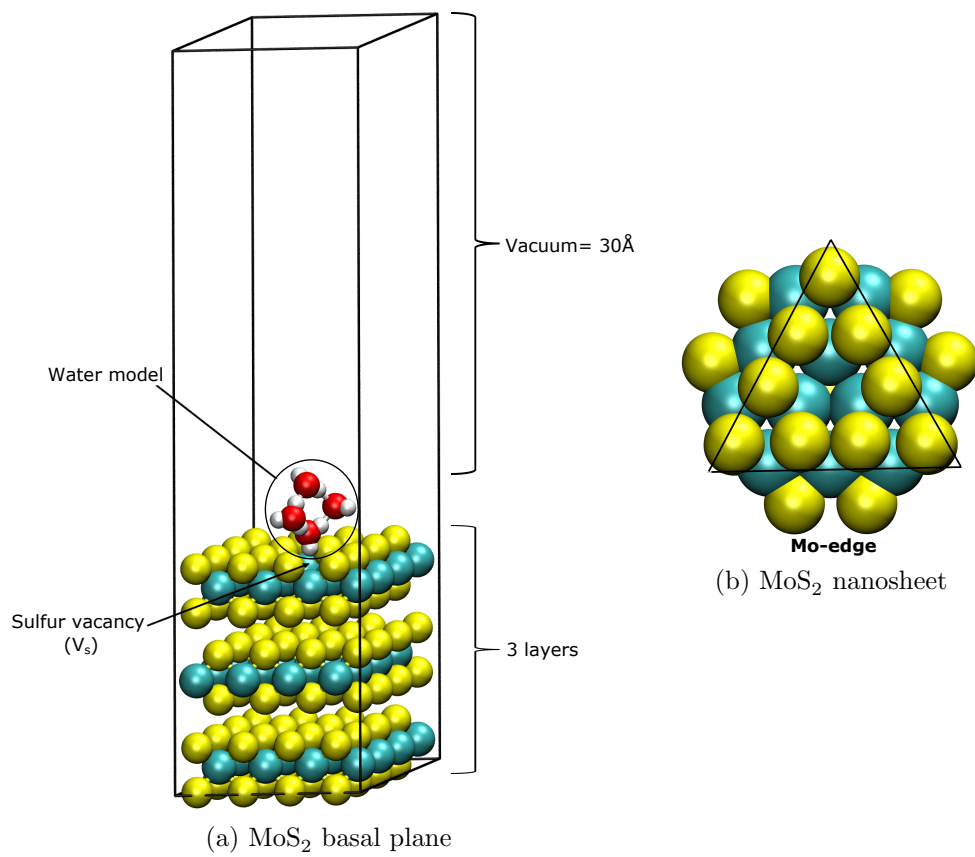
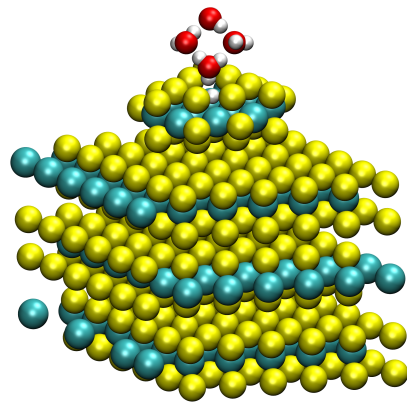
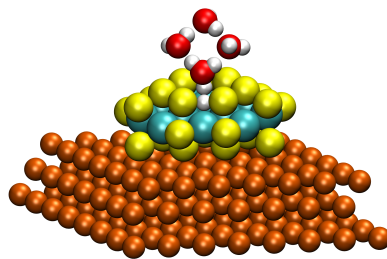


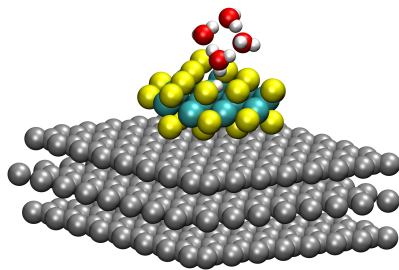
Figure 1: (a) The structure of non-symmetric MoS<sub>2</sub> basal plane with the water model and the sulfur vacancy. (b) MoS<sub>2</sub> nanosheet with sulfur vacancy and triangular shape terminated by Mo edges. The color code for atoms is white for H; red for O; yellow for S and greenish for Mo.



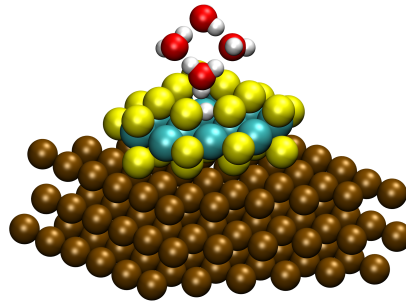
(a)  $\text{MoS}_2@ \text{MoS}_2$



(b)  $\text{MoS}_2@ \text{Cu}$



(c)  $\text{MoS}_2@ \text{C}$



(d)  $\text{MoS}_2@ \text{Au}$

Figure 2:  $3 \times 3$   $\text{MoS}_2$  nanosheet supported on molybdenum disulphide (a), copper (b), graphite (c) and gold (d). The structures presented correspond to the transition state of  $\text{V}^{II}$ . The color code for atoms is white for H; red for O; yellow for S, greenish for Mo, orange for Cu, grey for C and brown for Au.



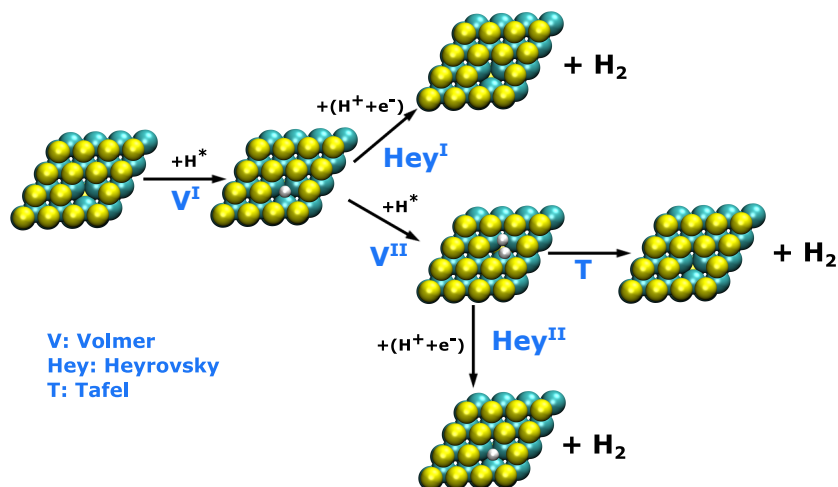


Figure 3: The different steps and the labels for the corresponding transition states of Volmer-Heyrovsky and Volmer-Tafel mechanisms to produce  $\text{H}_2$  on the sulfur vacancy  $V_S$ . The color code for atoms is white for H; yellow for S and greenish for Mo.

### 3. Results and Discussion

#### 3.1. HER mechanism on $V_S$

We start by the investigation of the mechanism of hydrogen production at the "isolated" sulfur vacancy ( $V_S$ ). As shown in our previous study [64],  $V_S$  is an active site with hydrogen adsorption energy close to 0 eV (two hydrogen atoms can be adsorbed). Fig. 3 presents the different intermediates for the possible HER mechanisms. The results of hydrogen adsorption energies and activation energies at zero charge (i.e., without GC-DFT) are shown in Table S2 and Table S3, respectively.

By using GC-DFT, Fig. S3 shows the evolution of the adsorption energies of the two hydrogen atoms on  $V_S$  as a function of the electrochemical potential. We find that the production of  $\text{H}_2$  is thermodynamically starting from -0.2 V vs SHE. Fig. 4 displays the optimized geometries of the different transition states for each step with the corresponding activation energies as a function of electrochemical potential. As shown,  $V^I$ ,  $V^{II}$ ,  $\text{Hey}^I$  and  $\text{Hey}^{II}$  depend strongly (0.34 eV/V) on the electrochemical potential, which is not the case for the formally chemical Tafel (T) step. This is in complete agreement with previous studies suggesting that Volmer and Heyrovsky are electrochemical steps while T is a chemical reaction.[54]

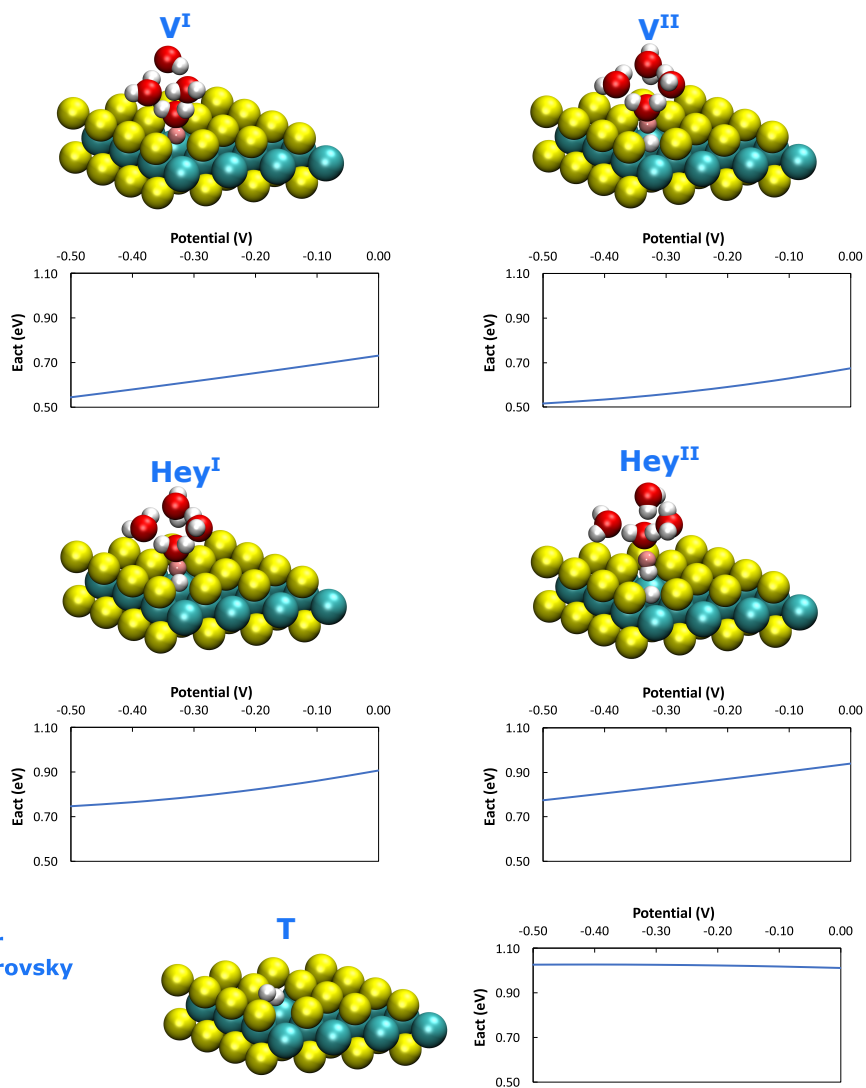


Figure 4: The structure of the different transition states and the variation of their activation energies as function of the electrochemical potential. The color code for atoms is white for H; pink for H coming from the water model; yellow for S and greenish for Mo.

As can be seen in Fig. 4, the hydrogen atom is adsorbed from the solution with a barrier of 0.54 eV at -0.5 V. This value can be compared to the corresponding activation energy of Huang et al.[54] which is much smaller (0.07 eV). Compared to our model, they used a smaller p(3 × 3) MoS<sub>2</sub> surface (instead of p(4 × 4) herein) and H<sub>3</sub>O<sup>+</sup> as the proton source (instead of H<sub>9</sub>O<sub>4</sub><sup>+</sup>). Fig. S4 shows the variation of the activation energy of V<sub>I</sub> as a function of electrochemical potential for p(4 × 4), our surface, and p(3 × 3) used by Huang et al.. We find that at a given electrochemical potential, the activation energy is lower by 0.05 eV at -0.2 V with p(4 × 4) than with p(3 × 3). Hence, this is not the origin of the difference between our barrier and the previous report. However, due to insufficiencies of the implicit solvent model, the proton is much less stabilised in H<sub>3</sub>O<sup>+</sup> compared with H<sub>9</sub>O<sub>4</sub><sup>+</sup>. This greater stability of the initial species explains the increase in the barrier and is in line with a previous study comparing H<sub>3</sub>O<sup>+</sup> and H<sub>5</sub>O<sub>2</sub><sup>+</sup> as proton sources.[52]

To identify the HER pathway, we compare the activation energies of the different steps at -0.2 V vs SHE, which is close to the onset potential. The first step is V<sup>I</sup>, the detachment of the hydrogen atom from H<sub>3</sub>O<sup>+</sup> to be adsorbed onto the sulfur vacancy with  $\Delta E_{act}$  of 0.65 eV. Then, H\* on V<sub>s</sub> can interact with the solution proton to produce H<sub>2</sub> via the Hey<sup>II</sup> reaction with a barrier of 0.82 eV. These results indicate that the first adsorbed H is not easily desorbed to generate H<sub>2</sub> because the sulfur vacancy has a strong adsorption energy for H\*. Therefore, before desorption of H<sub>2</sub>, another H<sup>+</sup> plus e<sup>-</sup> pair moves to the sulfur vacancy corresponding to V<sup>II</sup> with a  $\Delta E_{act}$  of 0.59 eV. Subsequently, H<sub>2</sub> is produced by Hey<sup>II</sup>, the reaction of H\* with one electron and proton, or the combination with the second H\* (T).  $\Delta E_{act}$  for Hey<sup>II</sup> and T are 0.87 eV and 1.02 eV, respectively. Thus, Tafel step has a higher activation energy compared to Hey<sup>II</sup>. Therefore, the HER occurs on the sulfur vacancy mainly via catalytic cycles between V<sup>II</sup> and Hey<sup>II</sup>. The Hey<sup>II</sup> reaction is found to be the rate-determining step (RDS), as it has the highest barrier energy. Overall, these barriers are in good agreement with a reaction occurring at room temperature: According to the Eyring equation, it corresponds to a turn-over frequency of 0.2 per second.

### 3.2. The support effect

One of the strategies to improve the electrocatalytic water splitting activity of non-metallic catalysts such as MoS<sub>2</sub> is to enhance the charge transfer from the conductive support to the catalyst surfaces.[70] Here, we investigate

for the first time the effect of different supports on the activation energies of the elementary reactions for the release of  $\text{H}_2$  on sulfur vacancies, including explicitly the electrochemical potential for one of the representative reactions. The corresponding structures are depicted in Fig. 2.

Identifying the transition states on all the supports has been a very tedious task, and despite significant efforts we have not been able to identify the transition state of  $\text{Hey}^{II}$  for any of them, so we exclude it from the following discussion. Fig. 5 shows the activation energies of  $\text{V}^I$ ,  $\text{V}^{II}$ ,  $\text{Hey}^I$  and T without considering the electrochemical potential. We find that the activation energies for  $\text{V}^I$ ,  $\text{V}^{II}$  and  $\text{Hey}^I$  barely depend on the supports, implying that the type of support has no effect on the activation energy. The only notable difference is the Tafel step, where the copper support increases the activation energy by 0.26 eV. As detailed in section S4 of the Supporting Information, various tests were performed related to the nanosheet size, the deposition of a periodic surface of  $\text{MoS}_2$  on the support instead of a nanosheet, and the defect density. The final result was always the same: the activation energies obtained are very close to those in Fig. 5. As a proxy for the electronic structure, we have analyzed the Hirshfeld atomic charges and their redistribution from the initial state to the TS, see Fig. S8 and S9. According to our analysis (see also the discussion in section S5), the atomic charges change only little as a function of the support and nanosheet vs. periodic  $\text{MoS}_2$  layer. We conclude that there is a very small interaction between the  $\text{MoS}_2$  nanosheet and the different supports and that neither the size of the nanosheet nor the number of (isolated) sulfur vacancies can contribute to the reduction of the barrier energies for the  $\text{V}^I$  step. Therefore, we conclude that for the release of  $\text{H}_2$  with  $\text{MoS}_2$  as electrocatalyst,  $\text{MoS}_2$  support, gold or graphite give the same energy barrier. However, copper is to be avoided for this type of reaction.

Given that the potentials of zero charge vary by up to 1 V between the initial and transition states (see Fig. S5), we have also assessed the effect of the electrochemical potential on the activation energy for a typical electrochemical process. We have chosen  $\text{V}^I$  as a representative transition state, which also was computationally well behaved. Hence, we have symmetrized the corresponding systems (see section 2.3), except the graphite supported one which turned out to be computationally intractable. The GC-DFT results are depicted in Fig. 6, where the strong dependence of the activation energy on the electrochemical potential (and thus non-negligible, but support in-dependent) charge-transfer coefficient) is retrieved. We first note that the

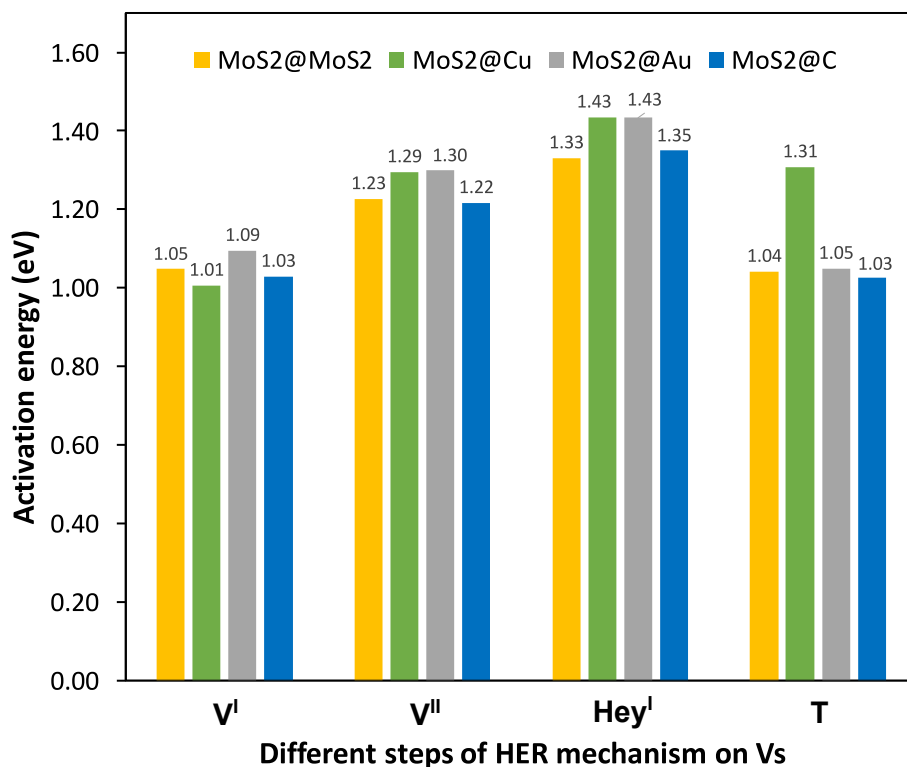


Figure 5: Comparison between the activation energies of HER mechanism steps on MoS<sub>2</sub> nanosheet with the different type of supports at zero charge, i.e., using canonical DFT, not GC-DFT. The colors used: yellow, green, grey and blue are for the supports MoS<sub>2</sub>, copper, gold and graphite, respectively.

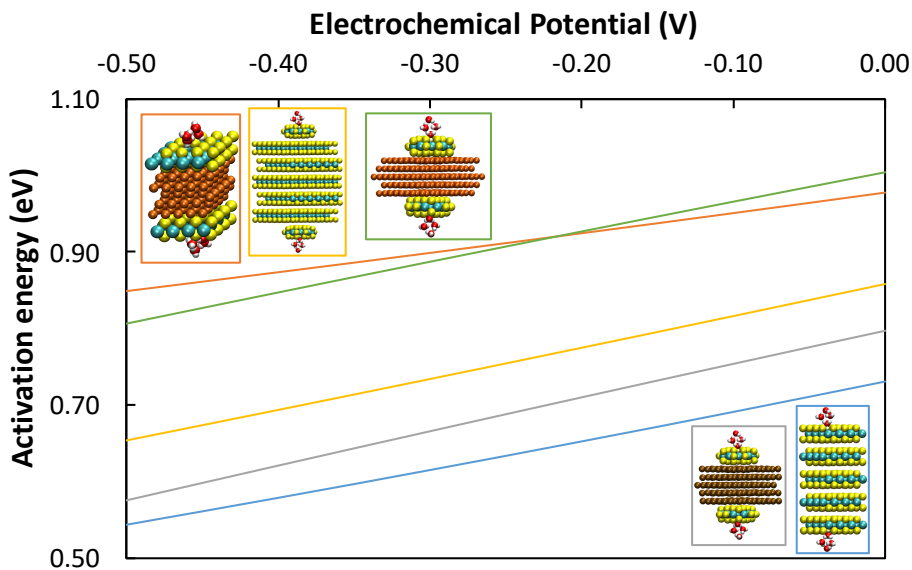


Figure 6: Variation of the activation energy as function of the electrochemical potential. The yellow, green and grey curves are for MoS<sub>2</sub> nanosheet supported on MoS<sub>2</sub>, copper and gold respectively. The blue curve is the periodic MoS<sub>2</sub> surface and the orange color is for periodic MoS<sub>2</sub> monolayer supported on Cu.

nanosheet model supported on MoS<sub>2</sub> and the periodic MoS<sub>2</sub> surface give similar results (within  $\sim 0.1$  eV) and the nanosheet supported on Au is right in the middle and hence close to the pure MoS<sub>2</sub> systems. For the Cu support, we find that the barrier energy increases for nanosheets or periodic surfaces of MoS<sub>2</sub>, confirming that this type of support should be avoided. Based on these results, we suggest that a conductive (graphitic) carbon support (from the non-symmetric results) is more advantageous compared to expensive gold electrodes.

From these various computations and models, we can draw several conclusions: First, the increasing or decreasing of the defect density has little effect on the HER activity: Starting from a rather high (1 per  $p(3\times 3)$  unit cell (or  $1.3\times 10^{13}$  per  $\text{cm}^2$ ) defect density down to 1 per  $p(6\times 6)$  unit cell (or  $3.1\times 10^{12}$  per  $\text{cm}^2$ ) leads to equivalent activation energies (see Fig. S4 and S7). Second, only symmetric slabs ensure reliable conclusions on the impact of the support. This originates in the difference in workfunction between the "occupied" and "unoccupied" support surfaces. The most telling example is MoS<sub>2</sub>@Cu: For an asymmetric system, the surface bound electrons, that

depend on the electrochemical potential, are shared among the "bare" and the "covered" surface. Hence, the overall workfunction is roughly the average, which impacts the energetics for reactions where the polarity changes dramatically, such as the Volmer step, where one moves from  $\text{H}_3\text{O}^+ + \text{e}^-$  to  $\text{H}^* + \text{H}_2\text{O}$ . Third, in the systems studied herein, the charge-transfer is "weak", in the sense that it is not due to a chemical interaction, but just due to the equalization of the Fermi-levels of the different materials. However, in electrocatalysis the Fermi-level is determined by the electrode potential. Hence, the GC-DFT framework reveals that the weak charge-transfer has no significant impact on electrocatalysts. This is best seen for the "size" effect (small nanosheet vs. monolayer) where the change in electronic structure is small enough to be largely counteracted by the electrode potential. However, "stronger" modifications, such as doping [71, 72] and defect engineering [73, 74] has, of course, a role to play even in electrocatalysis.

#### 4. Conclusion

We have investigated the kinetics of HER over sulfur vacancies on the  $\text{MoS}_2$  basal plane. We have identified the activation energies of the five plausible elementary steps (two consecutive Volmer steps, one Tafel step and two Heyrovsky steps) as a function of the electrochemical potential. In agreement with expectations based on the change of polarity between the reactant and the transition state, the Tafel step is a purely chemical step at the potentials of interest and does, therefore, not depend significantly on the electrochemical potential. The Volmer and Heyrovsky barriers, are, in contrast, dependent on the electrochemical potential, with a slope of about 0.34 eV/V. We then investigated means to tune the HER kinetics over these defects via support effects, size effects (nanosheet vs. periodic surface) and defect density. While in the absence of the electrochemical potential such a tuning seems possible, taking the electrochemical potential explicitly into account "washes out" these effects. In other words, the "intrinsic" activity of the sulfur vacancy is retrieved essentially independently on the size of the sheets, the defect density and the support. Hence, small nanosheets on a high surface area carbon support are likely to be an optimal choice, as it simply increases the number of active sites on a geometrical area basis and should, therefore, lead to a high geometrical current density.

## Author Contributions

Nawras Abidi: Investigation, Writing - Original Draft  
Audrey Bonduelle-Skrzypczak: Writing - Review  
Stephan N. Steinmann: Supervision, Conceptualization, Resources, Writing - Review & Editing

## Conflicts of interest

There are no conflicts to declare

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## Supporting Information

The supplementary material contains computational details and additional Figures and Tables. Coordinates and DFT raw data is freely available on NOMAD-lab: <https://dx.doi.org/10.17172/NOMAD/2022.09.07-1>.

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