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Enhanced performance of planar perovskite solar cells by doping the SnO₂ electron transport layer with guanidinium chloride

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Tin (IV) oxide is a highly promising electron transport layer (ETL) for lead halide perovskite solar cells due to its high conductivity, transparency, wide band gap, and the possibility of low-temperature processing. Nonetheless, charge carrier recombination processes at the SnO₂/perovskite interface diminish the device performance. Here, we demonstrate that SnO₂ doping with guanidine hydrochloride (G-SnO₂) leads to efficient surface passivation and a larger band offset between the ETL and the perovskite layer, resulting in reduced voltage losses and faster electron transfer. Moreover, G-SnO₂ facilitates the growth of highly crystalline perovskite layers. Consequently, a power conversion efficiency of up to 23.48% and a high open-circuit voltage of 1.18 V are obtained in solar cells incorporating the G-SnO₂ ETL. These devices also exhibited negligible hysteresis and maintained more than 96% of their initial power conversion efficiency after 1,250 h exposure to the air without encapsulation.

KEYWORDS

perovskite solar cells, electron transport layer, SnO₂, guanidinium chloride, grain size

1 Introduction

The electron transport layer (ETL) is one essential part of perovskite solar cells (PSCs), assuring electron extraction from the perovskite layer and blocking holes to avoid undesired recombination processes. (Ke et al., 2015; Anaraki et al., 2016; Jiang et al., 2016). An ideal ETL should possess high electron mobility, favorable energy level alignment with the active layer, and good processability for the preparation of high-quality films. Amongst the N-type semiconductor used for the fabrication of ETLs, WO_x, NbO_x, CeO_x, and BaSnO₃ can be found along with the more widely used materials TiO₂, ZnO, and SnO₂. (Murugadoss et al., 2016; Lee et al., 2017; Zhu et al., 2017). In this group, SnO₂ is a highly promising candidate owing to (i) its high electron mobility (over 420 cm²/V.s), favoring efficient transport; (ii) its wide energy band gap, giving rise to fewer losses due to parasitic light absorption; (iii) its low annealing temperature, normally less than 150 °C, which enables the fabrication of solar cells at a lower cost and on various types of substrates including flexible ones. (Tu et al., 2019; Wang et al., 2020a; Xiong et al., 2022).

To overcome limitations originating from undesired recombination processes at the SnO₂/perovskite interface, which are related to sub-optimal energy level alignment and interfacial defect states, (Sutanto et al., 2021; Liu et al., 2022; Odunmbaku et al., 2022),

surface modification of the ETL is a commonly used approach. Various kinds of passivating agents have been reported, for example, inorganic salts, long-chain polymers, and small organic molecules. (Zhao et al., 2018; Jiang et al., 2019; Fu et al., 2020). As an example, Huang and coworkers demonstrated that passivation of the SnO₂ ETL with potassium ions (from KOH) eliminated hysteresis effects and influenced the growth of the perovskite grains, resulting in enhanced efficiency and stability. (Bu et al., 2018). Potassium chloride (KCl) has also been used for passivating SnO₂ thin films. (Jalebi et al., 2018; Wang et al., 2018; Qin et al., 2019). More recently, Yu et al. reported the use of guanidinium chloride for the treatment of the SnO₂ ETL to eliminate detrimental O₂⁻ species formed at its surface during thermal annealing, which led to an improvement of the power conversion efficiency (PCE) of the obtained perovskite solar cells from 15.33% to 18.46%. (Yu et al., 2020).

Compared with these post-deposition surface treatment procedures of the annealed ETL layer, *in situ* bulk doping of the ETL *via* the addition of appropriate molecular species during its preparation is an alternative approach to fine-tune its properties. (Jung et al., 2013; Ren et al., 2017; Huang et al., 2018). In particular, it can give access to increased conductivity of the ETL due to electrical doping, which is often accompanied by the positive side effect of lower annealing temperatures required for the ETL preparation. Moreover, the judicious choice of additives can lead to better control of the band alignment at the ETL/perovskite interface favoring efficient charge transfer. As an example, Li and coworkers adopted electron-donating triphenylphosphine oxide (TPPO) as a dopant of the SnO₂ layer, which enhanced its conductivity and decreased the energy barrier at perovskite/SnO₂ interface from 0.55 to 0.39 eV, resulting in a PCE of 20.69%. (Deng et al., 2020). In another example, Zheng and coworkers applied KF-doped SnO₂ to improve the performance of all-inorganic CsPbI₂Br solar cells showing an open-circuit voltage (V_{oc}) of 1.31 V (Zheng et al., 2021). In our previous work, we investigated ammonium chloride as a dopant for SnO₂, employing a double-layer ETL with different doping levels. This approach led to a significant improvement in the performance of planar PSCs due to a more favorable energy level alignment at the ETL/perovskite interface, resulting in a V_{oc} of 1.21 V and a PCE of 21.75% (pristine ETL: 1.14 V, 18.54%). (Ye et al., 2021).

Based on these results, we explore here guanidine hydrochloride (H₂N)₂C = NH₂⁺Cl⁻, also known as guanidinium chloride (Gua-Cl), as a bulk dopant for the SnO₂ ETL. Gua has been previously used mainly for improving the stability of PSCs by incorporating this large cation into the perovskite layer. (Ishibashi et al., 2017; Wu et al., 2019). In the case of methylammonium iodide (MAI)-based PSCs, the enhanced stability of the devices was attributed to the larger number of hydrogen bonds, increasing from 1-2 for MA to 6 for Gua. (Jodlowski et al., 2017). When using Gua-Cl for doping the ETL, a synergistic beneficial effect can be expected of the chloride anion, which can act as an n-type dopant in SnO₂, and the Gua cation, enabling efficient hydrogen bonding with the halide ions from the perovskite layer. Furthermore, the delocalization of the positive charge over the three nitrogen atoms and the resulting resonance stabilization are expected to result in an improved electronic coupling with the perovskite and enhanced interfacial charge transfer. In this work, we developed an effective strategy for

the preparation of guanidinium chloride-doped SnO₂ ETLs (G-SnO₂) at temperatures as low as 80°C, i.e., fully compatible with flexible substrates or with the integration in tandem devices. Our results reveal that the use of G-SnO₂ minimizes defects at the perovskite/ETL interface, which in turn enhances charge carrier transfer. The G-SnO₂ layer also facilitates perovskite crystal growth and offers better stability of the perovskite film by passivating undercoordinated ions. The PSCs prepared with G-SnO₂ showed a PCE up to 23.48% with negligible hysteresis and high V_{oc} (1.18 V), while PSCs using pristine SnO₂ (P-SnO₂) had a maximum PCE of 20.91%. Moreover, G-SnO₂-based devices maintained over 96% of their initial performance after being kept in the air for 1,250 h without encapsulation.

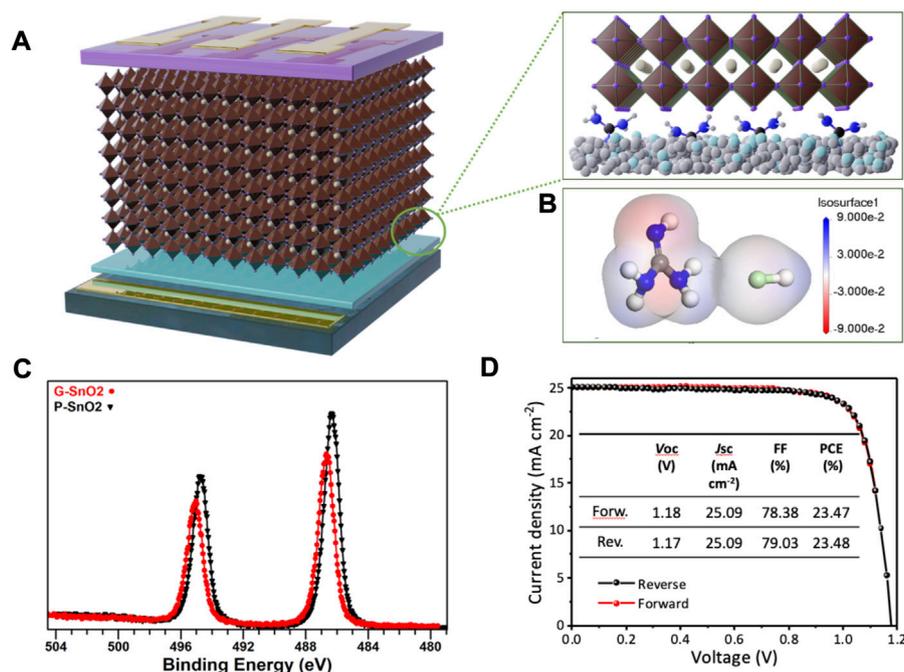
2 Materials and methods

2.1 Materials

Indium tin oxide (ITO) coated glass substrates were purchased from Xiangcheng Technology (6 Ωsq⁻¹). Lead iodide (PbI₂) and Spiro-OMeTAD were purchased from Xi'an Polymer Light Technology Corp. 4-tert-Butylpyridine (TBP, 96%), bis(trifluoromethylsulfonyl)-imide lithium salt (Li-TFSI), FK209 Co(III) PF₆ salt (FK209), chlorobenzene (anhydrous, 99.8%), dimethylsulfoxide (DMSO, anhydrous, ≥99.9%), N,N-dimethylformamide (DMF, anhydrous, 99.8%), isopropyl alcohol (IPA, anhydrous, 99.5%) and guanidine hydrochloride (99%) were purchased from Sigma Aldrich. Formamidinium iodide (FAI), methylammonium bromide (MABr), phenethylammonium iodide (PEAI) and methylammonium chloride (MACl) were procured from GreatCell Solar. The SnO₂ colloid solution (tin (IV) oxide, 15% in H₂O colloidal dispersion) was purchased from Alfa Aesar.

2.2 Device fabrication

ITO-coated glass substrates were cleaned in detergent solution, deionized water, acetone, ethanol and isopropyl alcohol in order before exposing them to UV-ozone for 30 min to remove residual organic materials. The SnO₂ precursor solution was prepared by diluting 1 ml of the colloidal dispersion of SnO₂ nanoparticles 6 times with deionized water. In the case of the G-SnO₂ layer, 50 mmol of the dopant guanidine hydrochloride was added just before use. The obtained SnO₂ precursor solutions were spin-coated on ITO glass substrates twice with no heat treatment between the two steps using 100 μl solution, 2000 rpm acceleration, 4,000 rpm speed, and 30 s duration, and then the films were annealed at 80°C for 1 h. For the deposition of the CsFAMA perovskite film, the precursor solution was prepared by mixing PbI₂ (1.45 M), FAI (1.15 M), PbBr₂ (0.21 M), MABr (0.21 M) and MACl (0.45 M) in a mixed solvent of DMF/DMSO (8:1 v/v, 1 ml) and finally adding 17.5 μl of CsI solution (2 M in DMSO). The obtained solution was spin-coated onto the ETL substrate by two consecutive spin-coating steps at 1,000 and 4,000 rpm for 10 and 30 s, respectively. During the second spin-coating step, 150 μl of chlorobenzene were dropped onto the substrate after 20 s before transferring it to a hotplate for thermal treatment at 105°C for 60 min within a glove-box. After



cooling to room temperature, 60 μ l of a 20 mM PEAI solution in IPA were spin-coated onto the perovskite surface at a speed of 4,000 rpm. For the preparation of the HTL, spiro-OMeTAD was dissolved in 1 mL of chlorobenzene (80 mg/mL), with the addition of 29.8 μ l of 4-tertbutylpyridine, 18 μ l of Li-TFSI (520 mg/mL in acetonitrile) and 12 μ l FK209 (250 mg/mL in acetonitrile) to increase the conductivity. Around 40 μ l of the resulting precursor solution were spin-coated on top of the perovskite/PEAI layer at a speed of 4,000 rpm for 30 s. Finally, a 100-nm Au electrode was thermally evaporated on top of the stack.

2.3 Characterization

Morphological characterization (top-view and cross-sectional images) of the perovskite layer and full devices was performed using a high-resolution scanning electron microscope (Zeiss Ultra 55). AFM images were obtained with a Bruker Dimension Icon Atomic Force Microscope. GIXRD and XRD data were acquired on a Panalytical Empyrean diffractometer equipped with a Cobalt anode beam tube (Co $\lambda_{K\alpha 1}$ = 1.7890 Å, Co $\lambda_{K\alpha 2}$ = 1.7929 Å), a Göbel mirror, and a 1D Pixcel detector; to simplify the comparison with the literature, the data is presented for Cu K α radiation. UV-vis absorbance spectra were recorded on a Shimadzu UV-1800 spectrometer. Conductivity values were obtained using an Ecopia HMS-3000 Hall Measurement System, a magnetic field of 0.5T, and a current from -200 mA to +200 mA. The current-voltage characteristics of the devices were assessed with a Keithley 2,400 Source meter under AM 1.5G illumination using a

calibrated solar simulator (Newport AAA). The area of the solar cells was delimited with a mask (0.09 cm²) and a scan rate of 20 mVs⁻¹ was used for both forward and reverse scans. Steady-state and time-resolved PL analyses were performed using a FLUOROLOG spectrometer (HORIBA FL3) with a 450 nm continuous or 425 nm pulsed laser excitation (NanoLED-425L). X-ray photoelectron spectroscopy (XPS) analyses were carried out with a Versa Probe II spectrometer (ULVAC-PHI) equipped with a monochromated Al K α source ($h\nu$ = 1,486.6 eV). The core level peaks were recorded with a constant pass energy of 23.3 eV. The XPS spectra were fitted with CasaXPS 2.3 software using Shirley background. Binding energies (BEs) are referenced with respect to adventitious carbon (C 1s BE = 284.8 eV). The UPS measurements were performed on an M-XPS spectrometer from Omicron Nanotechnology equipped with an ultraviolet He lamp (He I = 21.2 eV). The samples were electrically biased at -5.0 V for accurate determination of the low-kinetic energy cutoff. Photoelectrons were collected under normal geometry by a hemispherical analyzer, with an angular acceptance of 8°.

3 Results and discussion

We fabricated solar cell devices with n-i-p architecture comprising a stack of ITO/SnO₂/Perovskite/Spiro/Au (Figure 1A) using the triple-cation (cesium/formamidinium/methylammonium, CsFAMA) perovskite. The ETL was prepared by adding an optimized amount of Gua-Cl to commercial SnO₂ colloids in an aqueous solution followed by spin-coating and thermal annealing at

low temperature (*cf.* Experimental Section). Protonated Gua and chloride ions on SnO₂ promote photogenerated electron transfer from the perovskite layer because of electrostatic interactions. Due to the higher formation energy, Pb-Cl bonds are less likely to form, avoiding antisite defects and thus having little effect on interface recombination. (Grancini et al., 2015; Jiang et al., 2015; Chandrasekhar et al., 2016). To visualize the electron density distribution, electrostatic potential (ESP) analysis was performed. As shown in Figure 1B, the area of the imino moiety presents a high electron density making it prone to interaction with under-coordinated Pb²⁺ cations at the perovskite grain surface, which are the main sources of charge traps. These factors are favorable for reducing hysteresis effects and improving electron extraction. Moreover, guanidinium cations can be directly incorporated into the lattice of triple cation perovskites and form a thin layer of a 2D perovskite phase between SnO₂ and the 3D perovskite, (Zhang et al., 2019), contributing to thermal and environmental stability enhancement. XPS measurements (Figure 1C) were used to investigate the oxidation state of tin in SnO₂. The peaks visible at 486.3 and 494.7 eV in P-SnO₂, corresponding to the Sn 3d_{5/2} and Sn 3d_{3/2} states, exhibit a shift to higher binding energy by 0.4 eV in G-SnO₂. This shift signifies that Gua-Cl doping of SnO₂ leads to a lowering of electron density and thus a stronger oxidation of the Sn atoms, inductive of the filling of oxygen vacancies. The passivation of oxygen vacancies by the combined action of chloride ions and tetramethylammonium ions has recently been reported by Parida et al. for the case of fully inorganic CsPbI₂Br PSCs. (Parida et al., 2021). Hall Effect measurements show that, due to the better passivation and doping, the G-SnO₂ layer exhibits an 11% decreased resistivity and a 7.5% higher electron mobility than the G-SnO₂ ETL (*cf.* Supplementary Table S2). Gua-Cl doping thus improves the charge extraction and transport properties of the ETL, which should be directly visible in the solar cell performances and hysteresis behavior.

Figure 1D shows the *J/V* curves of the champion device employing the G-SnO₂ ETL, which yielded the best PCE of 23.48% with a *V*_{OC} of 1.18 V, a short-circuit current density (*J*_{SC}) of 25.1 mA/cm² and a fill factor (*FF*) of 78.4%. The integrated *J*_{SC} obtained from the EQE measurement was 24.69 mA/cm² for the device, which is in accordance with the value determined from the *J/V* scan (Fig. S2). Compared with the reference device PCE of 20.91% (Supplementary Table S1), the performance showed a 12.3% enhancement with suppressed hysteresis, which can be attributed to improved electron transport in the ETL and fewer electron accumulation at the interfaces of the ETL.

The structural properties of the SnO₂ ETL were first studied employing grazing incidence X-ray diffraction (GIXRD). At a low incident angle ω of 0.2° the signals related to the SnO₂ crystallites were observed with distinct diffraction peaks at 25.95°, 33.31° and 50.82° (Supplementary Figure S3A) whose positions are coherent with the tetragonal rutile structure of SnO₂ (ref. ICDD PDF 01-071-5,323). When increasing the incident angle ω from 0.2° to 0.4°, the ITO peaks became visible due to the larger depth probed, making it easy to distinguish them from the SnO₂-related peaks. Compared with P-SnO₂, films of G-SnO₂ showed no difference in the peak positions (Supplementary Figure S3B), confirming that Gua-Cl doping does not change the structural properties of the ETL. Next, ultraviolet photoelectron spectroscopy (UPS) was

carried out to analyze the energy level positions of P-SnO₂ and G-SnO₂. The work functions (Φ) of P-SnO₂ and G-SnO₂ were estimated from the secondary electron cutoff as -4.23 and -4.41 eV, respectively (Figure 2A). Plotting the UPS near the Fermi levels (*E*_F) allowed calculating the valence band maxima (VBM) with respect to *E*_F (-3.73 and -3.69 eV for G/P-SnO₂) (Figure 2B), resulting in of *E*_{VBM} = -8.14 and -7.92 eV vs. vacuum level, respectively (*cf.* Supplementary Figure S6). Notably, the Fermi level of the G-SnO₂ is situated deeper than that of the intrinsic P-SnO₂ layer, suggesting n-type doping and leading to a slight band bending at the interface with the perovskite layer. Figure 2C shows *T*_{auc} plots of each SnO₂ film on glass substrates. The optical bandgap values of G-SnO₂ and P-SnO₂ were deduced as 3.90 and 3.88 eV, respectively, *i.e.*, essentially unchanged after doping. From these results, and combined with the conduction band minimum (CBM) of CsFAMA perovskites of -4.05 eV calculated by Abate and coworkers, (Wang et al., 2020b), an energy level diagram was constructed as shown in Figure 2D. As one important feature, an energy offset at the ETL/perovskite interface of 0.19 eV is obtained with G-SnO₂, while P-SnO₂ results in a flat alignment. Therefore, electron extraction is enhanced with G-SnO₂ because of the higher energetic driving force, (Song et al., 2017), and a gradient-type band alignment between the perovskite layer and ITO is achieved.

The morphology of the perovskite absorber layer was studied by SEM. It can be seen from Figures 3A, B that the perovskite films on P-SnO₂ layers have crystallite sizes of around 200–400 nm, while much larger sizes of 500–800 nm were formed on G-SnO₂ layers, leading to the grain size distributions depicted in Supplementary Figure S7. As visible in the cross-section view, the perovskite films grown on both types of SnO₂ showed densely packed, highly crystalline grains (Figure 3D). The differences observed between the two cases clearly demonstrate the direct influence of the underlying ETL layer, generating nucleation centers and promoting perovskite growth, on the crystalline quality of the perovskite film as has been shown in many literature examples (Bouchard et al., 2017; Dunlap-Shohl et al., 2019; Medjahed et al., 2020). The different morphology of the perovskite films having the same thickness (Figure 3C, D) leads to a slight increase in the absorption of the layer grown on G-SnO₂ (Supplementary Figure S8), which is beneficial for enhancing *J*_{sc}. This behavior is in accordance with reported results comparing the light absorption properties of perovskite thin films constituted of different grain sizes. (Xiao et al., 2014).

To assess the surface morphology and roughness of the underlying ETL layer, atomic force microscopy (AFM) measurements were conducted. As shown in Figure 3F, the roughness of the G-SnO₂ film (RMS = 2.52 nm) is slightly larger than that (2.39 nm) of the P-SnO₂ film (*cf.* Supplementary Table S3), but the G-SnO₂ film shows a denser surface morphology with improved homogeneity. The upper perovskite layer exhibited a reduced RMS roughness (*cf.* Supplementary Table S4), from 23.6 nm grown on a P-SnO₂ film to 16.3 nm on G-SnO₂. It can therefore be concluded that guanidine hydrochloride doped SnO₂ favors the formation of a dense and smooth perovskite layer which is beneficial for the subsequent deposition of the HTL. Moreover, the large crystallite size contributes to the suppression of trap states generated at grain boundaries.

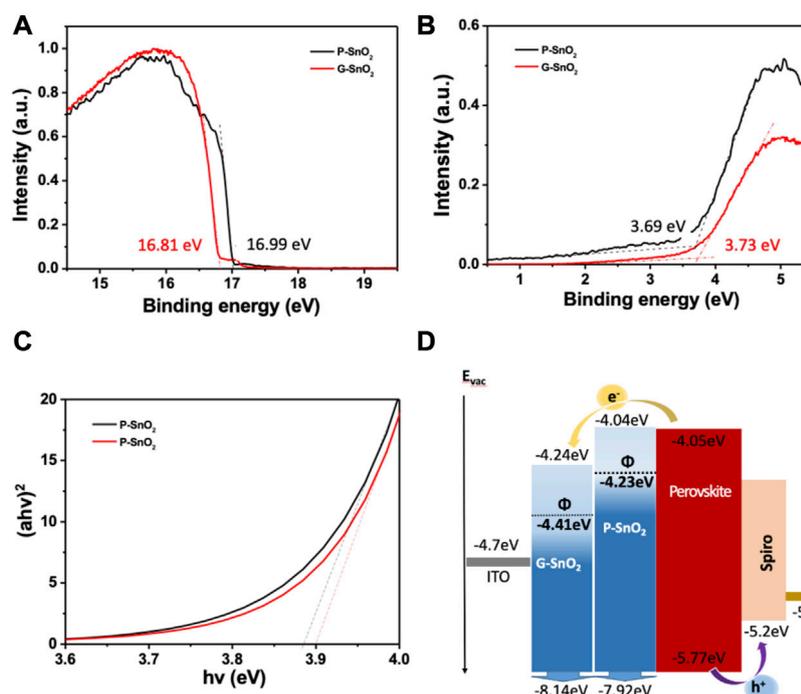


FIGURE 2

(A, B) UPS measurements on the pristine (P-SnO₂) and guanidine doped (G-SnO₂) films deposited on ITO. (A) secondary edge region; (B) Fermi level determination. (C) Tauc plot of P-SnO₂ and G-SnO₂ films deposited on ITO glass. (D) Energy level diagram derived from photoelectron and UV-vis absorption spectroscopy.

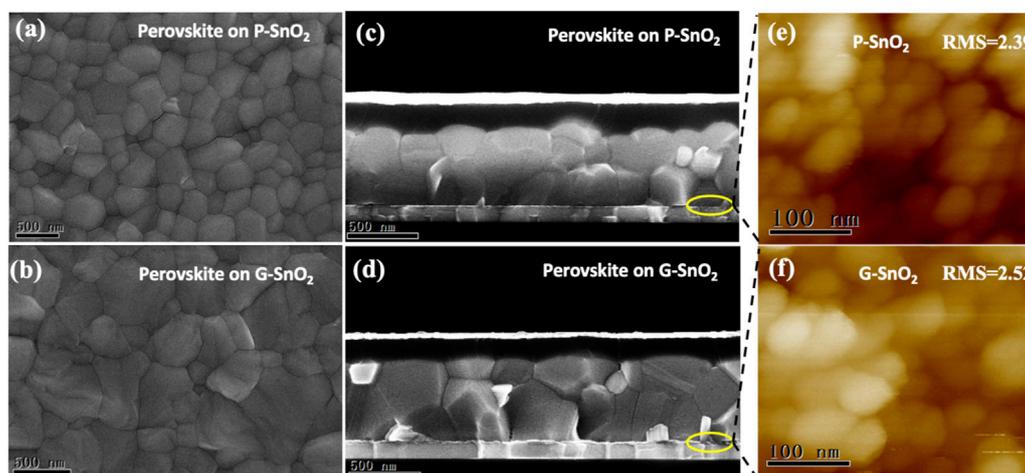


FIGURE 3

(A–D) SEM top and cross-section view of the perovskite layer grown on P- and G-SnO₂, respectively (E, F) AFM images of the tin oxide surface w/o Gua-Cl doping (scale bar: 100 nm).

X-ray diffraction patterns of the perovskite layers on the different ETLs (Figure 4A) exhibit the main peaks located at 14.06°, 28.43°, and 31.84°, which are assigned to the (001), (002), and (012) reflections of the cubic crystalline phase of CsFAMA (space group *m-3m*). A strong diffraction peak at 12.68° originates from excess PbI₂ used in the perovskite preparation, which is known

to passivate the grain boundaries. (Parida et al., 2021). Perovskite growth on P-SnO₂ resulted in a lattice constant of 6.278 (2) Å, while growth on G-SnO₂ led to a slightly decreased value of 6.274 (2) Å, as visible by the shift of the diffraction peaks to larger angles (Figure 4B). Figure 4C depicts the evolution of the (001) peak family for perovskites on P-SnO₂ and G-SnO₂ ETLs (*cf.*

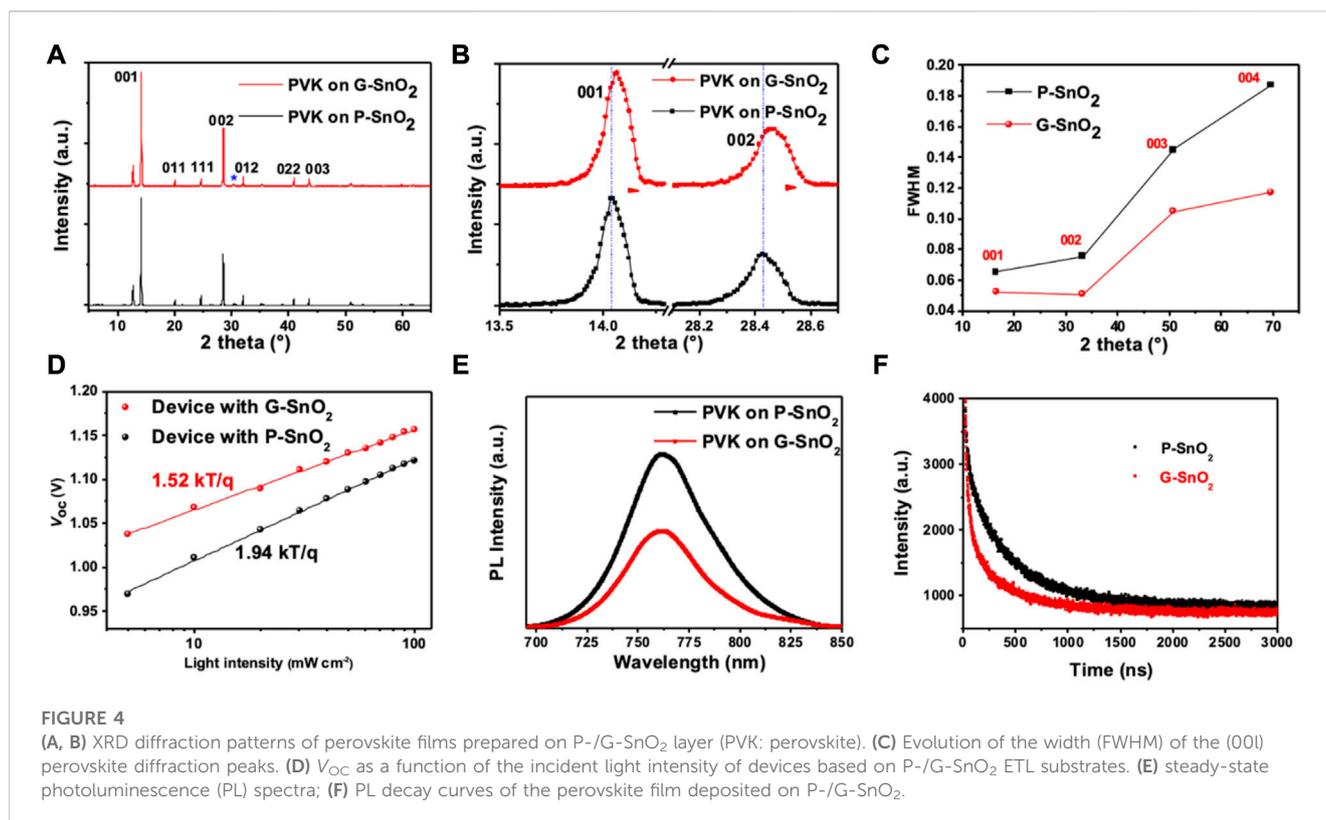


FIGURE 4

(A, B) XRD diffraction patterns of perovskite films prepared on P-/G-SnO₂ layer (PVK: perovskite). (C) Evolution of the width (FWHM) of the (001) perovskite diffraction peaks. (D) V_{OC} as a function of the incident light intensity of devices based on P-/G-SnO₂ ETL substrates. (E) steady-state photoluminescence (PL) spectra; (F) PL decay curves of the perovskite film deposited on P-/G-SnO₂.

Supplementary Table S5). Due to the non-negligible contribution of the instrumental resolution to the broadening of the Bragg peaks, a quantitative analysis was not possible. However, considering that the crystallite size is >100 nm as determined from SEM analyses, the observed increased FWHM of the diffraction peaks denotes the presence of distortion (strain) in the lattice for both types of ETL. Importantly, the smaller FWHM increase observed for the perovskite grown on G-SnO₂ is a sign of reduced strain and improved crystalline quality compared to the perovskite layer on P-SnO₂.

To further study the influence of the observed perovskite morphology and strain on the carrier recombination processes, we investigated the incident light intensity dependence of V_{OC} (Figure 4D). If the slope of V_{OC} versus light intensity is greater than kT/q (q: electron charge, k: Boltzmann constant, T: Kelvin temperature), V_{OC} is strongly dependent on the incident light intensity and additional recombination caused by strain is involved. Here, the P-SnO₂ devices exhibit a slope of 1.94 kT/q, while G-SnO₂ devices show a significantly smaller value (1.52 kT/q), indicating that trap-assisted recombination was effectively inhibited by reduced strain, which is consistent with the fitted recombination resistance from electrochemical impedance spectroscopy (Supplementary Figure S9). This assumption is further supported by steady-state and time-resolved photoluminescence (TRPL) spectroscopy, as shown in Figure 4E. In both cases, the perovskite film gives rise to a broad PL signal peaking at 762 nm, but the PL intensity of the film grown on the G-SnO₂ ETLs is significantly lower than that involving P-SnO₂, which is a sign of enhanced electron injection. Analysis of the TRPL data (Figure 4F)

results in a shorter excited state lifetime for the ITO/G-SnO₂/perovskite stack ($\tau = 318$ ns) than for the sample comprising the P-SnO₂ ETL ($\tau = 415$ ns), demonstrating the faster electron extraction from the perovskite film in the former case. Taken together, these data demonstrate that doping of the ETL with Gua-Cl leads to reduced charge accumulation and recombination losses at the perovskite/SnO₂ interface, which is highly beneficial for solar cell performance.

To assess the reproducibility of the PSCs, we tested 25 devices using each type of ETL (Figure 5A). All solar cell parameters of the devices incorporating G-SnO₂ are improved, while the main contribution stems from enhanced J_{SC} (23.4 vs. 24.9 mA/cm²), originating from the concomitant reduction of the trap state density and the modulated band structure as discussed above. The statistical analysis yields an average and champion PCE of 22.86% and 23.48%, respectively, compared to 19.85% and 20.91% for the P-SnO₂ cells. These results demonstrate the high reproducibility and effectiveness of the employed strategy. Figure 5B exhibits the stabilized power output at the maximum power point (MPP) in a nitrogen-filled glovebox, and no change was visible after MPP tracking for 120 s for both types of PSC. The stabilized PCE of the G-SnO₂-based cell is 23.2%, in agreement with the PCE obtained from the J/V curve. Finally, the stability of the non-encapsulated PSCs was studied under a dry air environment (Figure 5C). The G-SnO₂-based devices retained 96.1% of their initial PCE after 1,250 h, while the PCE of the P-SnO₂ based PSCs shrank to 89.9% under the same conditions. The observed stability enhancement is ascribed to the enhanced carrier extraction capacity of G-SnO₂ and the improved perovskite crystalline quality.

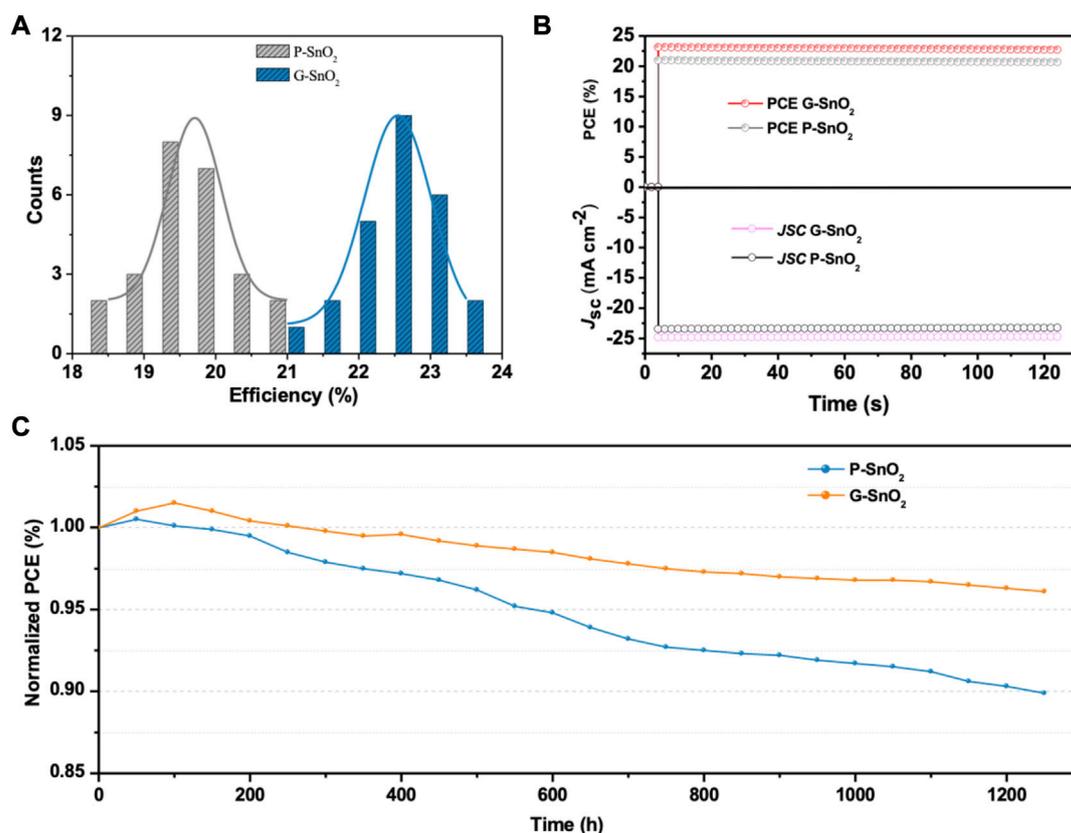


FIGURE 5

(A) Statistical distribution of the PCE of PSCs prepared with P-SnO₂ and G-SnO₂ (25 devices in each group). (B) Steady-state photocurrent density and device efficiency at maximum power point under continuous AM 1.5G illumination. (C) PCE decrease under dry air atmosphere of a reference device comprising the P-SnO₂ ETL and of a device with G-SnO₂ measured over a period of 1,250 h.

4 Conclusion

In summary, we presented a simple strategy for enhancing the performance of planar PSCs by doping the SnO₂ ETL with guanidinium chloride. The observed improvements are related to both the modification of the ETL intrinsic properties and the enhancement of the crystalline quality of the perovskite layer. More specifically, the modified G-SnO₂ ETL exhibited higher electron mobility and reduced defect state density than the pristine P-SnO₂ ETL, resulting in reduced interfacial charge carrier recombination. In addition, G-SnO₂ has a better-matched band alignment with the perovskite absorber further fostering efficient charge extraction. On the other hand, favored by the interaction between the guanidinium cation and halide anions from the perovskite precursors, the use of G-SnO₂ also induces the growth of a perovskite layer exhibiting significantly increased grain size (around 600 nm vs. 300 nm) and lower strain, while maintaining a dense packing and smooth surface. As a result of this synergistic behavior, the PSCs using G-SnO₂ exhibited an average PCE enhancement of 3% compared to devices with P-SnO₂, reaching a maximum value of 23.48%. Combined with appealing stability and the possibility of low-temperature (80°C) processing, the proposed strategy also facilitates the realization of

PSCs on flexible substrates and their integration into tandem devices.

Data availability statement

The original contributions presented in the study are included in the article/[Supplementary Material](#), further inquiries can be directed to the corresponding authors.

Author contributions

PR and JY conceived the study and JY, YLi, AAM, SP, and DA performed the experiments. All authors contributed to the data interpretation. The manuscript was drafted by PR, JY, and YLi, and all authors read and approved it.

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Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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Supplementary material

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fmats.2023.1118641/full#supplementary-material>

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