



On the Fisher infinitesimal model without variability

Amic Frouvelle, Cécile Taing

► To cite this version:

Amic Frouvelle, Cécile Taing. On the Fisher infinitesimal model without variability. 2023. hal-04168272v2

HAL Id: hal-04168272

<https://hal.science/hal-04168272v2>

Preprint submitted on 29 Sep 2023

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

On the Fisher infinitesimal model without variability

Amic Frouvelle^{*†¶}

Cécile Taing^{‡§¶}

October 2, 2023

Abstract

We study the long-time behavior of solutions to a model of sexual populations structured in phenotypes. The model features a nonlinear integral reproduction operator derived from the Fisher infinitesimal operator and a trait-dependent selection term. The reproduction operator describes here the inheritance of the mean parental traits to the offspring without variability. We show that, under assumptions on the growth of the selection rate, Dirac masses are stable around phenotypes for which the difference between the selection rate and its minimum value is less than $\frac{1}{2}$. Moreover, we prove the convergence in some Fourier-based distance of the centered and rescaled solution to a stationary profile under some conditions on the initial moments of the solution.

1 Introduction

1.1 The population model

We study a model of sexual populations structured by a phenotypic trait, represented by a continuous variable $x \in \mathbb{R}$. Denoting by $f(t, \cdot)$ the population density at time $t \geq 0$ in the trait space, the model we are interested in is the following :

$$\begin{cases} \partial_t f(t, x) = B_0[f(t, \cdot)](x) - m(x)f(t, x), & t \geq 0, \\ f(0, x) = f^0(x), \end{cases} \quad (1)$$

where B_0 is the trait inheritance operator, defined as, for $x \in \mathbb{R}$,

$$B_0[f](x) := \iint_{\mathbb{R}^2} \delta_0\left(x - \frac{z_1 + z_2}{2}\right) f(z_1) \frac{f(z_2)}{\int_{\mathbb{R}} f(z') dz'} dz_1 dz_2, \quad (2)$$

with δ_0 the Dirac measure. We interpret this last quantity as a the number of newborns with trait x per unit of time, and thus we refer to it as the reproduction term. With the operator B_0 , it is assumed that newborns inherit exactly the mean of the parental traits $\frac{1}{2}(z_1 + z_2)$. This mixing operator also features a normalization by the total mass of the population density $\int_{\mathbb{R}} f(z') dz'$, to illustrate that the choice of a partner is made at constant rate in time, uniformly among the whole population. We assume the selection rate m to be bounded below.

In the present work, we investigate the long time behavior of f . More precisely, we prove the local stability of some particular Dirac masses, and the convergence of a rescaled formulation of f to a stationary state with a measure-adapted Fourier distance.

^{*}CEREMADE, CNRS, Université Paris-Dauphine, Université PSL, 75016 Paris, France.

[†]The work of A.F. has been (partially) supported by the Project EFI ANR-17-CE40-0030 of the French National Research Agency.

[‡]Laboratoire de Mathématiques et Applications, Université de Poitiers, CNRS, F-86073 Poitiers, France.

[§]C.T. has received funding from the PEPS JCJC 2021 and 2022 programs of the INSMI (CNRS).

[¶]Emails: frouvelle@ceremade.dauphine.fr, cecile.taing@univ-poitiers.fr

The model (1) is motivated by works on the Fisher infinitesimal model [4, 22], with the following operator B_ε :

$$B_\varepsilon[f](x) := \frac{1}{\varepsilon\sqrt{\pi}} \iint_{\mathbb{R}^2} \exp \left[-\frac{1}{\varepsilon^2} \left(x - \frac{z_1 + z_2}{2} \right)^2 \right] f(z_1) \frac{f(z_2)}{\int_{\mathbb{R}} f(z') dz'} dz_1 dz_2.$$

which describes the offspring traits as normally distributed around the mean of the parental traits with small variance of order ε^2 . This model has been used in theoretical evolutionary biology [8]. From a mathematical point of view, the Fisher infinitesimal operator appears in the works [31, 37], where a different scaling than the small variance is used, and a spatial structure is added to study invasions. The authors show a derivation of the Kirkpatrick-Barton system [27] at a limit of a large reproduction rate. Also the infinitesimal operator is included in a selection-competition model in [18] with a population evolving between two habitats and a regime of small phenotypic variance compared to the environment heterogeneity. In [38], the Fisher infinitesimal operator is combined with a multiplicative selection one. The author showed local uniqueness of a steady state and exponential convergence of the solution for weak selection effects, thanks to a contraction property of the reproduction operator for the 2-Wasserstein distance, under the assumption of a compactly supported selection rate. A time-discrete version of the Fisher infinitesimal model, with a quadratic selection rate and non-overlapping generations, is analyzed in [10, 11], within a framework other than the Wasserstein one.

The study of the asymptotic behavior of the Fisher infinitesimal model with an appropriate time scaling has been tackled in [9, 33]. The authors study the concentration of the population distribution around some particular traits. More precisely, in the first work the authors studied the special stationary states at a regime of small variance, while in the second one the author considered the associated Cauchy problem in the same regime, showing that solutions can be approximated by Gaussian profiles with small variance. The asymptotic framework was built on the spirit of [20] with small variance limit for asexual population models.

Also, similar operators as B_ε have been studied in different contexts : alignment [17], protein exchanges [29], among other works. Mathematical models of sexual populations have recently received some attention, with variations of the reproduction term, to account for mating preferences [16, 28], asymmetrical inheritance [34] and allelic structure [15, 19].

In the present study, the main ingredients, coming from kinetic theory, are the estimation on the evolution of moments and the Fourier distance for measures. Indeed, the operator B_0 shares similarities with kinetic models of inelastic collisions. The idea of using the Fourier transform for Boltzmann-type equations goes back to the works of A. V. Bobylev in [6, 7] and the resulting equation has also been studied in [35]. Then, the Fourier distance has been employed in [25] to investigate the trend to equilibrium of the solutions to the Boltzmann equation for Maxwell molecules, and then has been adapted for other kinetic models of Boltzmann type, e.g. [5, 21, 30, 32, 36], where the formation of overpopulated tails is observed. Indeed, by the use of the Fourier transform of the collision operators, the equation on \widehat{f} is simpler than the one for f . This is also the case here for the operator B_0 defined in (2) since we have

$$\widehat{B_0(f)}(\xi) = \widehat{f}(\xi/2)^2,$$

which makes the use of Fourier metrics possible. Other distances have been proved to be non-expanding for measure solutions to structured populations PDE models, as Monge-Kantorovich type ones (see [23, 24]).

Apart from the trait mixing operator, our model also differs from the ones considered by these mentioned works in the selection component. This component is described by a trait-dependent mortality rate and is considered as being a constraint on the phenotypic variability

of the population. This affects the properties of the solution : there is no more conservation of mass, neither of the center of mass, and the evolution equation for the second moment is not closed. In the mentioned works on kinetic models, the proof of the convergence of the solution strongly relies on a contraction property of the collision operator in the Fourier distance and the estimation of the moments of the distribution. In the present work, we take the same approach with, in addition, the estimation of quantities derived from the selection rate, which helps to tackle the loss of conservation of the center of mass.

The work is organized as follows. In Section 2, we establish the existence and uniqueness of the solution of (1) in the sense of measures, we also define the moments of the centered and normalized solution and state some properties that they satisfy. Finally, we also derive the equation on the Fourier transform of the centered and rescaled formulation of the solution. In Section 3, we provide a local stability result of Dirac masses and derive estimates on the exponential decay in time of the moments. In Section 4, we prove the long time convergence of the centered and rescaled solution to a stationary self-similar profile.

1.2 Assumptions and main results

We outline in this section the assumptions we use and the main results that we prove throughout the present work.

First, we state the existence and uniqueness of the solution to (1) in the sense of measures. Denoting \mathcal{M}_+ the set of nonnegative Borel measures with finite mass, we use the following definition of a measure solution.

Definition 1.1. *If $f^0 \in \mathcal{M}_+$, we say that $f \in C([0, T], \mathcal{M}_+)$ is a weak solution to the model (1) if, for any Borel set $A \subset \mathbb{R}$, we have for all $t \in [0, T]$:*

$$\int_A f(t, x) dx = \int_A e^{-m(x)t} f^0(x) dx + \int_0^t \int_A e^{-m(x)(t-s)} B_0[f(s, \cdot)](x) dx ds. \quad (3)$$

The existence and uniqueness result is given by the following theorem that we prove in Section 2.1.

Theorem 1.1. *If m is a measurable function and is bounded below, then for any $f^0 \in \mathcal{M}_+$, and for any $T > 0$, there exists a unique weak solution $f \in C([0, T], \mathcal{M}_+)$ to the model (1), in the sense of Definition 1.1.*

The proof of Theorem 1.1 relies on standard arguments of fixed point construction using a distance based on the total variation of a measure.

Next, we focus on the study of the moments of f . As a first step, we derive a control on the time propagation of the initial moments of f , which enables to define the moments of f at any time. Namely, defining for any finite nonnegative measure ν , when $\int_{\mathbb{R}} |x|^k \nu(x) dx$ is finite, its moment of order k by

$$\mu_k(\nu) := \int_{\mathbb{R}} x^k \nu(x) dx,$$

the moments of $f(t, \cdot)$, denoted by $\mu_k(f(t, \cdot))$ are well defined for all time as soon as they are finite initially. Moreover, we prove that these moments are continuous and, for k even, that they satisfy some kind of differential inequalities, which will be used to establish estimates on the centered moments of the normalized density, denoted by

$$g(t, x) = \frac{f(t, x)}{\int_{\mathbb{R}} f(t, z) dz}.$$

We aim to get some regularity on the centered moments of g in order to determine the long time behavior of its centered and standardized formulation, which we will also call self-similar profile throughout this work. For this purpose, we add the following assumption of the selection rate m .

Assumption 1.1. *We suppose that m is measurable and that there exist constants $K \geq 0$ and $C \geq 0$ such that for all $x \in \mathbb{R}$, we have*

$$-K \leq m(x) \leq C(1 + x^2).$$

A first consequence of this assumption is that the time derivative of the moment of order k of f can be expressed with the lower order moments at any time, provided that the initial data has a finite moment of order $k + 2$. This implies that, in this case, the moments of f are continuously differentiable. Note that we do not impose the selection rate to be nonnegative. This may be interpreted as a supplementary asexual reproduction rate favored by the environment.

Then, we define the centered moments of g and the centered variations of m by the quantities

$$M_k(t) = \int_{\mathbb{R}} (x - \bar{x}(t))^k g(t, x) dx, \quad (4)$$

$$S_k(t) = \int_{\mathbb{R}} (x - \bar{x}(t))^k (m(x) - m(\bar{x}(t))) g(t, x) dx, \quad (5)$$

with $\bar{x}(t)$ the center of mass at time t ,

$$\bar{x}(t) := \int_{\mathbb{R}} x g(t, x) dx = \frac{\mu_1(f(t, \cdot))}{\rho(t)}, \quad (6)$$

and we denote their initial values by M_k^0, S_k^0, \bar{x}^0 . We obtain differential equations on the centered moments of g for any $k \geq 2$, assuming that f^0 has a finite moment of order $k + 2$. In the case where $k \geq 4$ is even and f^0 has only a moment of order k , we get a differential inequality on M_k .

A quantity that appears to be important in our study of these differential inequalities is the shifted deviation of m from its minimum, defined as

$$\eta(x) = \inf_{\mathbb{R}} m + \frac{1}{2} - m(x), \quad (7)$$

which is a quantity not exceeding $\frac{1}{2}$ (and which is positive if the excess of mortality is less than $\frac{1}{2}$).

For instance, when f has initially a fourth moment, we obtain

$$\frac{d}{dt} \bar{x} = -S_1, \quad (8)$$

$$\frac{d}{dt} M_2 = -\frac{1}{2} M_2 - S_2 + S_0 M_2, \quad (9)$$

$$\frac{d^+}{dt} M_4 \leq -\left(\frac{3}{8} + \eta(\bar{x}) - S_0\right) M_4 + \frac{3}{8} M_2^2 + 4S_1 M_3, \quad (10)$$

where we have written $\frac{d^+}{dt} M_4(t) := \limsup_{\tau \rightarrow 0, \tau > 0} \frac{M_4(t+\tau) - M_4(t)}{\tau}$.

One can see that we always have $S_2 \geq (\eta(\bar{x}) - \frac{1}{2}) M_2$, and therefore the differential equation (9) for M_2 gives $\frac{d}{dt} M_2 \leq -(\eta(\bar{x}) - S_0) M_2$. We expect the variation S_0 to be small, and we will see that a criterion for exponential decay of M_2 is indeed that $\eta(\bar{x})$ is initially positive.

With these differential equations and inequations, we aim to derive time estimates on the centered moments M_k . To do so, we need to have some bounds concerning the regularity and growth of the selection rate, which we state in the following assumption, which is for instance satisfied when m is C^2 with bounded second derivative, or when m is Lipschitz. Notice that this allows to cover the cases of selection rates with quadratic growth.

Assumption 1.2. *We suppose that there exists a continuous nonnegative function α and a constant $\beta \geq 0$ such that for all $x, y \in \mathbb{R}$, we have*

$$|m(x) - m(y)| \leq \alpha(y)|x - y| + \beta|x - y|^2.$$

With this last assumption on the growth of m , we prove the stability of Dirac masses around positions \bar{x}^0 for which the excess of mortality compared to its minimum value is less than $\frac{1}{2}$, that is to say $\eta(\bar{x}^0) > 0$. We show that if the center of mass at initial time is at such a position and if M_4^0 is small enough (which means that the initial normalized profile g^0 is close to a Dirac mass), then the center of mass converges to some limit, and the moments M_2 and M_4 decrease exponentially to 0 at long time, as stated in the following theorem.

Theorem 1.2. *Under Assumption 1.2, we suppose that initially we have $\eta(\bar{x}_0) > 0$. Then, for all δ such that $0 < \delta < \eta(\bar{x}_0)$, if M_4^0 is sufficiently small, we have for all $t > 0$:*

$$\eta(\bar{x}(t)) \geq \delta, \quad (11)$$

$$M_2(t) \leq M_2^0 e^{-\delta t}, \quad (12)$$

$$M_4(t) \leq M_4^0 e^{-\delta t}. \quad (13)$$

Furthermore, in that case, $\bar{x}(t)$ converges exponentially fast towards some $\bar{x}_\infty \in \mathbb{R}$. Consequently it means that if g is initially sufficiently close to a Dirac mass located at \bar{x}_0 (in 4-Wasserstein distance), then it converges exponentially fast to the Dirac mass at \bar{x}_∞ .

Thanks to this stability result and the regularity on the moments, we can define the Fourier transform of the centered and rescaled formulation of g and then characterize its convergence to a limit profile at long time, as announced in the beginning of this work.

Indeed, in order to investigate the long time behavior of the solution to (1), we make use of the following Fourier distance

$$d_s(\gamma_1, \gamma_2) := \sup_{\xi \neq 0} \frac{|\widehat{\gamma}_1(\xi) - \widehat{\gamma}_2(\xi)|}{|\xi|^s}, \quad (14)$$

which is defined when γ_1 and γ_2 are some measures that have same moments up to order $p \in \mathbb{N}$, with respective Fourier transforms $\widehat{\gamma}_1$ and $\widehat{\gamma}_2$, and $s \geq p$ to be determined.

For all $\xi \in \mathbb{R}$, we can define the Fourier transform of f solution to (1) on the trait variable

$$\widehat{f}(t, \xi) := \widehat{f(t, \cdot)}(\xi) = \int_{\mathbb{R}} e^{-i\xi x} f(t, x) dx.$$

Thus, to apply the Fourier distance (14) on measures that have same moments up to order 2, we define $\gamma(t, \cdot)$ as the rescaled and centered formulation of $g(t, \cdot)$, or associated self-similar profile,

$$\gamma(t, x) := \sqrt{M_2(t)} g\left(t, \sqrt{M_2(t)}x + \bar{x}(t)\right). \quad (15)$$

Then, we obtain that the Fourier transform satisfies the equation

$$\partial_t \widehat{\gamma}(t, \xi) = \widehat{\gamma}(t, \frac{\xi}{2})^2 - \widehat{\gamma}(t, \xi) + \frac{1}{4} \xi \partial_\xi \widehat{\gamma}(t, \xi) + R(t, \xi),$$

in which the term $R(t, \xi)$ comprises all the quantities generated by the selection rate m . Thus, R depends on some of the moments M_k and S_k , because of the formulation of γ and of the equations (8) and (9).

So, this last equation on $\widehat{\gamma}$ is divided into an operator applied only on $\widehat{\gamma}$, which comes from the Fourier formulation of the reproduction operator B_0 for probability distributions, and the remaining term R that we expect to become small at long time, because of the stability result on the center of mass of g and the estimates on the moments stated in Theorem 1.2.

In fact, it appears that the estimates on M_k obtained at this point are not sufficient to control the remaining term R . More precisely, we need to specifically quantify the exponential decay of the moments M_2 and M_4 , and also to derive a control on the ratio $\frac{M_{2k}}{M_2}$. With the assumptions of Theorem 1.2, we are able to improve the rates of convergence of the moments. First, we start by refining the estimates on M_2 and M_4 alone thanks to (9)-(10), then we obtain better rates of convergence for higher order moments, which finally are propagated back in the differential equations for lower moments to obtain precise estimates for all moments.

However, this is not sufficient to get a lower bound on the behaviour of M_2 , which happens to be crucial to control the remainder R . To this aim, we make a stronger assumption on the initial condition, by starting with an initial profile g^0 such that $\frac{M_{2k_0}^0}{M_2^0}$ is small (for some k_0 large enough). This is not necessarily the case for any profile close to a Dirac mass, but if we take any profile centered around \bar{x}^0 (with finite moment of order $2k_0$) and we shrink it by some large factor, we satisfy this assumption. In this framework we are able to obtain precise rates of convergence, as follows.

Theorem 1.3. *Under the assumptions of Theorem 1.2, we take $k_0 \geq 2$ such that $\frac{1}{2^{2k_0-1}} < \delta$, and we suppose that $\frac{M_{2k_0}^0}{M_2^0}$ is sufficiently small. Then,*

- *there exists $C_{k_0}(\frac{M_{2k_0}^0}{M_2^0}) > 0$, converging to 1 as $\frac{M_{2k_0}^0}{M_2^0} \rightarrow 0$ such that*

$$M_2(t) \geq C_{k_0}(\frac{M_{2k_0}^0}{M_2^0}) M_2^0 e^{-\frac{t}{2}},$$

- *for all $k \geq 2$ (with $k \leq k_0$) and for all $\lambda < \min(\frac{1}{2} - \frac{1}{2^{2k-1}}, \delta_0 - \frac{1}{2^{2k_0-1}})$, there exists $\widetilde{C}_{\lambda,k}(\frac{M_{2k_0}^0}{M_2^0}) > 0$, converging to 0 as $\frac{M_{2k_0}^0}{M_2^0} \rightarrow 0$ such that*

$$\frac{M_{2k}(t)}{M_2(t)} \leq \widetilde{C}_{\lambda,k}(\frac{M_{2k_0}^0}{M_2^0}) e^{-\lambda t}.$$

Interestingly, we obtain that the rate of convergence for M_2 is $\frac{1}{2}$, even when δ is very small.

Finally, the last theorem states the convergence in the Fourier distance of the profile γ to a stationary profile γ_∞ , which is identified in [3, 12], when t goes to $+\infty$, as

$$\gamma_\infty(x) = \frac{2}{\pi(1+x^2)^2}, \tag{16}$$

or given by

$$\widehat{\gamma}_\infty(\xi) = (1 + |\xi|)e^{-|\xi|}.$$

Theorem 1.4. *Under the assumptions of Theorem 1.3, if there exists some $k_0 \geq 3$ such that $\frac{1}{2^{2k_0-1}} < \delta$ and $\frac{M_{2k_0}^0}{M_2^0}$ is small enough, then we have that, for some values of $s \in (2, 3)$, there exists a constant $L > 0$ such that*

$$d_s(\gamma, \gamma_\infty)(t) \leq (d_s(\gamma_0, \gamma_\infty) + L)e^{-\lambda_s t},$$

with $\lambda_s := 1 - \frac{s}{4} - 2^{1-s}$, which is positive for $s \in (2, 3)$.

Remark 1.1. When $\delta - \frac{1}{2^{2k_0-1}} > \frac{1}{4}$, this estimation is valid for all $s \in (2, 3)$. Otherwise, there exist $2 < s_0 < \bar{s} < 3$ such that this estimation is valid on $(2, s_0)$. For $s = s_0$ we obtain $d_{s_0}(\gamma, \gamma_\infty)(t) \leq (d_{s_0}(\gamma_0, \gamma_\infty) + Lt)e^{-\lambda_{s_0}t}$, and for $s \in (s_0, \bar{s})$, we obtain that there exists a constant $c_s \in (0, \lambda_s)$ such that

$$d_s(\gamma, \gamma_\infty)(t) \leq d_s(\gamma_0, \gamma_\infty)e^{-\lambda_s t} + Le^{-c_s t}.$$

In all these situations, the constant L depends on δ, s and the initial moments M_{2k}^0 for $1 \leq k \leq k_0$.

To prove this last theorem, we prove a contraction property on the reproduction operator B_0 in the distance d_s , which is due to the quadratic formulation of its Fourier transform, and combine it with the results of Theorems 1.2 and 1.3 that enable to control the remainder R .

2 General results

2.1 Global existence and uniqueness of measure solutions

We want to give a meaning to our model (1) for an arbitrary nonnegative measure of finite mass on \mathbb{R} . If f is a smooth solution of (1), thinking of $B_0[f(t, \cdot)](x)$ as a prescribed source term, we can interpret (1) pointwise for every $x \in \mathbb{R}$ as an ordinary differential equation in time, therefore for all $t \geq 0$ and for all $x \in \mathbb{R}$ we have

$$f(t, x) = f^0(x)e^{-m(x)t} + \int_0^t e^{-m(x)(t-s)} B_0[f(s, \cdot)](x) ds,$$

and this will be the starting point of our definition of a measure solution. In the following, we will often use the abusive notation $f(x)dx$ even if f is only a measure (or $f(t, x)dx$ if the measure depends on time). Similarly, if h a measurable function in \mathbb{R} , we may write $\varphi(x) = h(x)f(x)$ to define a signed measure φ , even if f is only a measure.

We denote by \mathcal{M} the space of finite signed (Borel) measures on \mathbb{R} . Using the total variation norm defined by

$$\|\nu\|_{TV} = \sup \sum_{i=1}^n |\nu(A_i)|,$$

where the supremum is taken over all finite partitions of \mathbb{R} by Borel sets $(A_i)_{1 \leq i \leq n}$, this turns \mathcal{M} into a Banach space [14].

Furthermore, we have that for any measurable function h taking values in $[-H, H]$, we have

$$\left| \int_{\mathbb{R}} h(x) d\nu(x) \right| \leq H \|\nu\|_{TV},$$

and therefore we obtain that $\|\nu\|_{TV} = \sup \left| \int_{\mathbb{R}} h(x) d\nu(x) \right|$ where the supremum is taken over all measurable function h taking values in $[-1, 1]$.

The set $\mathcal{M}_+ \subset \mathcal{M}$ of nonnegative measures with finite mass is a closed subspace of \mathcal{M} . Indeed, if ν_n converges to ν in total variation norm, then we have $|\nu_n(A) - \nu(A)| \leq \|\nu_n - \nu\|_{TV} \rightarrow 0$ for all Borel set $A \subset \mathbb{R}$, thus $\nu_n(A) \rightarrow \nu(A)$. So, if ν_n is a nonnegative measure for all n , then ν is also nonnegative.

Notice that if $\nu \in \mathcal{M}_+$, then $B_0[\nu]$ also belongs to \mathcal{M}_+ , has same mass, and for all bounded (or nonnegative) and measurable h we have

$$\int_{\mathbb{R}} h(x) B_0[\nu](x) dx = \frac{\int_{\mathbb{R} \times \mathbb{R}} h\left(\frac{y+z}{2}\right) \nu(y) dy \nu(z) dz}{\int_{\mathbb{R}} \nu(x) dx}, \quad (17)$$

when ν has positive mass (thus the conservation of mass is obtained with $h = 1$), and $B_0[\nu] = 0$ if ν is the zero measure.

We are therefore ready to study the properties of measure solutions of our model (1), given by Definition 1.1. Notice that all terms in (3) are well defined as nonnegative integrals, and if m is measurable and bounded below, they are finite, since $x \mapsto e^{-m(x)t}$ is measurable and bounded for $t \geq 0$. Consequently, if f is a weak solution and h is a measurable bounded (resp. nonnegative) function, being a uniform (resp. monotone) limit of a sequence of measurable simple functions, we have :

$$\int_{\mathbb{R}} h(x) f(t, x) dx = \int_{\mathbb{R}} h(x) e^{-m(x)t} f^0(x) dx + \int_0^t \int_{\mathbb{R}} h(x) e^{-m(x)(t-s)} B_0[f(s, \cdot)](x) dx ds. \quad (18)$$

Thanks to this property, we can have a useful semigroup property for the solutions of our model.

Proposition 2.1. *Let $t_0 > 0$. If f is a weak solution to the model (1), then $t \mapsto f(t_0 + t, \cdot)$ is also a solution for the initial condition $f(t_0, \cdot)$.*

Proof. We fix $t > 0$. If A is a Borel set, we have

$$\begin{aligned} \int_A f(t_0 + t, x) dx &= \int_A e^{-m(x)(t_0+t)} f^0(x) dx + \int_0^{t_0+t} \int_A e^{-m(x)(t_0+t-s)} B_0[f(s, \cdot)](x) dx ds \\ &= \int_A e^{-m(x)(t_0+t)} f^0(x) dx + \int_0^{t_0} \int_A e^{-m(x)(t_0+t-s)} B_0[f(s, \cdot)](x) dx ds \\ &\quad + \int_0^t \int_A e^{-m(x)(t-s)} B_0[f(t_0 + s, \cdot)](x) dx ds. \end{aligned}$$

Since the function $\mathbf{1}_A(x) e^{-m(x)t}$ is measurable (and nonnegative), we have thanks to (18) that the first two terms of the right-hand side of this last equality combine and we get

$$\int_A f(t_0 + t, x) dx = \int_A e^{-m(x)t} f(t_0, x) dx + \int_0^t \int_A e^{-m(x)(t-s)} B_0[f(t_0 + s, \cdot)](x) dx ds,$$

and this ends the proof. \square

We proceed to the proof of Theorem 1.1, which provides global existence and uniqueness of solutions to (1) in \mathcal{M}_+ .

Proof. First of all, we need to study some contraction properties of B_0 . If h is a measurable function from \mathbb{R} to $[-H, H]$, and $\nu, \tilde{\nu} \in \mathcal{M}_+$, we write $\rho = \int_{\mathbb{R}} \nu(x) dx$ and $\tilde{\rho} = \int_{\mathbb{R}} \tilde{\nu}(x) dx$ and we have, when $\rho, \tilde{\rho} > 0$,

$$\begin{aligned} &\left| \int_{\mathbb{R}} h(x) B_0[\nu](x) dx - \int_{\mathbb{R}} h(x) B_0[\tilde{\nu}](x) dx \right| \\ &= \left| \frac{\int_{\mathbb{R} \times \mathbb{R}} h\left(\frac{y+z}{2}\right) \nu(y) dy \nu(z) dz}{\rho} - \frac{\int_{\mathbb{R} \times \mathbb{R}} h\left(\frac{y+z}{2}\right) \tilde{\nu}(y) dy \tilde{\nu}(z) dz}{\tilde{\rho}} \right| \\ &\leq \left| \frac{\int_{\mathbb{R} \times \mathbb{R}} h\left(\frac{y+z}{2}\right) \nu(y) dy \nu(z) dz}{\rho} - \frac{\int_{\mathbb{R} \times \mathbb{R}} h\left(\frac{y+z}{2}\right) \nu(y) dy \tilde{\nu}(z) dz}{\rho} \right| \\ &\quad + \left| \frac{\int_{\mathbb{R} \times \mathbb{R}} h\left(\frac{y+z}{2}\right) \nu(y) dy \tilde{\nu}(z) dz}{\rho} - \frac{\int_{\mathbb{R} \times \mathbb{R}} h\left(\frac{y+z}{2}\right) \tilde{\nu}(y) dy \tilde{\nu}(z) dz}{\tilde{\rho}} \right| \\ &\quad + \left| \frac{\int_{\mathbb{R} \times \mathbb{R}} h\left(\frac{y+z}{2}\right) \tilde{\nu}(y) dy \tilde{\nu}(z) dz}{\tilde{\rho}} - \frac{\int_{\mathbb{R} \times \mathbb{R}} h\left(\frac{y+z}{2}\right) \tilde{\nu}(y) dy \nu(z) dz}{\tilde{\rho}} \right| \end{aligned}$$

$$\begin{aligned}
&\leq \frac{\int_{\mathbb{R}} \left| \int_{\mathbb{R}} h\left(\frac{y+z}{2}\right) \nu(z) dz - \int_{\mathbb{R}} h\left(\frac{y+z}{2}\right) \tilde{\nu}(z) dz \right| \nu(y) dy}{\rho} \\
&\quad + \left| \frac{1}{\rho} - \frac{1}{\tilde{\rho}} \right| \int_{\mathbb{R} \times \mathbb{R}} \left| h\left(\frac{y+z}{2}\right) \right| \nu(y) dy \tilde{\nu}(z) dz \\
&\quad + \frac{\int_{\mathbb{R}} \left| \int_{\mathbb{R}} h\left(\frac{y+z}{2}\right) \nu(y) dy - \int_{\mathbb{R}} h\left(\frac{y+z}{2}\right) \tilde{\nu}(y) dy \right| \tilde{\nu}(z) dz}{\tilde{\rho}} \\
&\leq 2H \|\nu - \tilde{\nu}\|_{TV} + H|\tilde{\rho} - \rho|.
\end{aligned}$$

We actually have $\rho = \|\nu\|_{TV}$ (since ν is nonnegative) and $\tilde{\rho} = \|\tilde{\nu}\|_{TV}$, therefore we obtain

$$\left| \int_{\mathbb{R}} h(x) B_0[\nu](x) dx - \int_{\mathbb{R}} h(x) B_0[\tilde{\nu}](x) dx \right| \leq 3H \|\nu - \tilde{\nu}\|_{TV}, \quad (19)$$

and one easily checks that this inequality is still valid if $\rho = 0$ or $\tilde{\rho} = 0$.

We now fix $f^0 \in \mathcal{M}_+$ and $T > 0$. For any fixed $L > 0$ (to be determined later on), the space $\mathcal{X} = C([0, T], \mathcal{M})$ with the norm

$$\|f\|_{\mathcal{X}} = \sup_{t \in [0, T]} e^{-Lt} \|f(t, \cdot)\|_{TV},$$

is a Banach space. The subset $\mathcal{X}_+ = C([0, T], \mathcal{M}_+)$ of nonnegative finite measures varying continuously in time is a closed subset of \mathcal{X} , and we will construct a fixed point by Picard iteration on this complete metric space.

For $f \in C([0, T], \mathcal{M}_+)$, we denote $\Phi(f)$ the time-dependent measure defined for all Borel set $A \subset \mathbb{R}$ by

$$\int_A \Phi(f)(t, x) dx = \int_A e^{-m(x)t} f^0(x) dx + \int_0^t \int_A e^{-m(x)(t-s)} B_0[f(s, \cdot)](x) dx ds, \quad (20)$$

and we are therefore looking for a fixed point in \mathcal{X}_+ of the operator Φ . Let us first check that Φ is well-defined from \mathcal{X}_+ to itself.

If $h : \mathbb{R} \rightarrow [-1, 1]$ is a measurable function, being a uniform limit of a sequence of measurable simple functions, we have

$$\int_{\mathbb{R}} h(x) \Phi(f)(t, x) dx = \int_{\mathbb{R}} h(x) e^{-m(x)t} f^0(x) dx + \int_0^t \int_{\mathbb{R}} h(x) e^{-m(x)(t-s)} B_0[f(s, \cdot)](x) dx ds.$$

Then, taking $K \geq 0$ such that m is bounded below by $-K$, we obtain, for $t, t' \geq 0$,

$$\begin{aligned}
&\left| \int_{\mathbb{R}} h(x) \Phi(f)(t', x) dx - \int_{\mathbb{R}} h(x) \Phi(f)(t, x) dx \right| \\
&\leq \int_{\mathbb{R}} |e^{-m(x)t'} - e^{-m(x)t}| f^0(x) dx + \int_t^{t'} \int_{\mathbb{R}} |e^{-m(x)(t'-s)} - e^{-m(x)(t-s)}| B_0[f(s, \cdot)](x) dx ds \\
&\quad + \int_0^t \int_{\mathbb{R}} |e^{-m(x)(t'-s)} - e^{-m(x)(t-s)}| B_0[f(s, \cdot)](x) dx ds \\
&\leq K e^{KT} |t' - t| \|f^0\|_{TV} + \int_0^t K e^{KT} |t' - t| \|f(s, \cdot)\|_{TV} ds + \int_t^{t'} e^{KT} \|f(s, \cdot)\|_{TV} ds,
\end{aligned}$$

where we used the fact that $e^{-m(x)t}$ is $K e^{KT}$ -Lipschitz in time for all $x \in \mathbb{R}$, and the fact that the total mass of $B_0[f(s, \cdot)]$ (which is its total variation since f is nonnegative) is the same as

the one of $f(s, \cdot)$. Furthermore we have $\|f(s, \cdot)\|_{TV} \leq e^{Ls} \|f\|_{\mathcal{X}} \leq e^{LT} \|f\|_{\mathcal{X}}$ for all $s \in [0, T]$. Therefore we obtain

$$\|\Phi(f)(t', \cdot) - \Phi(f)(t, \cdot)\|_{TV} \leq (Ke^{KT} \|f^0\|_{TV} + (KT + 1)e^{(K+L)T} \|f\|_{\mathcal{X}}) |t' - t|,$$

showing that $t \mapsto \Phi(f)(t, \cdot)$ is continuous in time (and actually Lipschitz). Notice that since all the measures in the definition of Φ are nonnegative, we get that $\Phi(f)(t, \cdot)$ is a nonnegative measure for all time, and the operator Φ therefore sends \mathcal{X}_+ into itself.

For $f, \tilde{f} \in \mathcal{X}_+$, we fix a measurable function $h : \mathbb{R} \rightarrow [-1, 1]$, and we compute similarly, thanks to (19)

$$\begin{aligned} & \left| \int_{\mathbb{R}} h(x) \Phi(f)(t, x) dx - \int_{\mathbb{R}} h(x) \Phi(\tilde{f})(t, x) dx \right| \\ & \leq \int_0^t \left| \int_{\mathbb{R}} h(x) e^{-m(x)(t-s)} B_0[f(s, \cdot)](x) dx - \int_{\mathbb{R}} h(x) e^{-m(x)(t-s)} B_0[\tilde{f}(s, \cdot)](x) dx \right| ds \\ & \leq \int_0^t 3e^{KT} \|f(s, \cdot) - \tilde{f}(s, \cdot)\|_{TV} ds \\ & \leq \int_0^t 3e^{KT} \|f - \tilde{f}\|_{\mathcal{X}} e^{Ls} ds = \frac{3e^{KT}}{L} (e^{Lt} - 1) \|f - \tilde{f}\|_{\mathcal{X}}, \end{aligned}$$

which shows that

$$\|\Phi(f)(t, \cdot) - \Phi(\tilde{f})(t, \cdot)\|_{TV} \leq \frac{3e^{KT}}{L} e^{Lt} \|f - \tilde{f}\|_{\mathcal{X}},$$

and therefore

$$\|\Phi(f) - \Phi(\tilde{f})\|_{\mathcal{X}} \leq \frac{3e^{KT}}{L} \|f - \tilde{f}\|_{\mathcal{X}}.$$

So, for $L > 3e^{KT}$, Φ is a contraction and therefore has a unique fixed point in \mathcal{X}_+ . \square

2.2 Evolution of moments

In order to investigate the large time behavior of a solution f , we will need to control the moments. The key observation is that any bound on finite moments is preserved by B_0 .

Lemma 2.1. *If $\nu \in \mathcal{M}_+$, then for all $k \in \mathbb{N}$, $\int_{\mathbb{R}} |x|^k B_0[\nu](x) dx \leq \int_{\mathbb{R}} |x|^k \nu(x) dx$. Furthermore, if $\int_{\mathbb{R}} |x|^k \nu(x) dx$ is finite and ν has positive mass, all the moments up to order k are well defined and we have*

$$\mu_k(B_0[\nu]) = \sum_{j=0}^k \binom{k}{j} \frac{\mu_j(\nu) \mu_{k-j}(\nu)}{2^k \mu_0(\nu)}. \quad (21)$$

Proof. Let us remark that both sides of the first inequality are well defined as nonnegative integrals. We suppose that $\int_{\mathbb{R}} |x|^k \nu(x) dx$ is finite and that ν has positive mass, thanks to (17) we get :

$$\begin{aligned} \int_{\mathbb{R}} |x|^k B_0[\nu](x) dx &= \frac{\int_{\mathbb{R} \times \mathbb{R}} \frac{|y+z|^k}{2^k} \nu(y) dy \nu(z) dz}{\int_{\mathbb{R}} \nu(x) dx} \\ &\leq \frac{\sum_{j=0}^k \binom{k}{j} \int_{\mathbb{R}} |y|^j \nu(y) dy \int_{\mathbb{R}} |z|^{k-j} \nu(z) dz}{2^k \int_{\mathbb{R}} \nu(x) dx}. \end{aligned}$$

Then, by Hölder inequality, for $0 \leq j \leq k$ we obtain

$$\int_{\mathbb{R}} |y|^j \nu(y) dy \leq \left(\int_{\mathbb{R}} \nu(x) dx \right)^{1-\frac{j}{k}} \left(\int_{\mathbb{R}} |x|^k \nu(x) dx \right)^{\frac{j}{k}},$$

and we therefore get

$$\int_{\mathbb{R}} |x|^k B_0[\nu](x) dx \leq \frac{\sum_{j=0}^k \binom{k}{j} \left(\int_{\mathbb{R}} \nu(x) dx \right) \left(\int_{\mathbb{R}} |x|^j \nu(x) dx \right)}{2^k \int_{\mathbb{R}} \nu(x) dx} = \int_{\mathbb{R}} |x|^k \nu(x) dx.$$

Now, by using the dominated convergence theorem, we can approximate x^k by a bounded measurable function $h_N(x)$, for instance by setting

$$h_N(x) = \begin{cases} (-N)^k & \text{if } x \leq -N \\ x^k & \text{if } -N \leq x \leq N \\ N^k & \text{if } x \geq N, \end{cases} \quad (22)$$

and obtain, by applying (17) to h_N , that in the limit $N \rightarrow \infty$, we have

$$\begin{aligned} \int_{\mathbb{R}} x^k B_0[\nu](x) dx &= \frac{\int_{\mathbb{R} \times \mathbb{R}} \frac{(y+z)^k}{2^k} \nu(y) dy \nu(z) dz}{\int_{\mathbb{R}} \nu(x) dx}, \\ &= \frac{\sum_{j=0}^k \binom{k}{j} \int_{\mathbb{R}} y^j \nu(y) dy \int_{\mathbb{R}} z^{k-j} \nu(z) dz}{2^k \int_{\mathbb{R}} \nu(x) dx}. \end{aligned}$$

which is then well-defined and provides the formula (21). \square

Thanks to this property, we can now give the proof of the following proposition, which shows that finiteness of moments is propagated in time.

Proposition 2.2. *Let $f^0 \in \mathcal{M}_+$, m be measurable and bounded below by $-K \leq 0$, and f be the global solution to the model (1) given by Theorem 1.1. If for some $k \in \mathbb{N}$, $\int_{\mathbb{R}} |x|^k f^0(x) dx < +\infty$, then for all $t \geq 0$, we have*

$$\int_{\mathbb{R}} |x|^k f(t, x) dx \leq e^{(K+1)t} \int_{\mathbb{R}} |x|^k f^0(x) dx.$$

Proof. Back to the Picard iteration sequence $(f_n)_n$ with $f_{n+1} = \Phi(f_n)$ and $f_0 = f^0$, where Φ is defined in (20) in the proof of Theorem 1.1, we proceed by induction and assume as the induction hypothesis for $n \geq 0$ that

$$\int_{\mathbb{R}} |x|^k f_n(t, x) dx \leq e^{(K+1)t} \int_{\mathbb{R}} |x|^k f^0(x) dx.$$

Then we have, since $h_N(x)$ defined in (22) is bounded and measurable,

$$\begin{aligned} \int_{\mathbb{R}} |h_N(x)| f_{n+1}(t, x) dx &= \int_{\mathbb{R}} e^{-m(x)t} |h_N(x)| f^0(x) dx \\ &\quad + \int_0^t \int_{\mathbb{R}} e^{-m(x)(t-s)} |h_N(x)| B_0[f_n(s, \cdot)](x) dx ds \\ &\leq e^{Kt} \int_{\mathbb{R}} |x|^k f^0(x) dx + \int_0^t e^{K(t-s)} \int_{\mathbb{R}} |h_N(x)| B_0[f_n(s, \cdot)](x) dx ds. \end{aligned}$$

By Lemma 2.1 and the induction hypothesis, we obtain

$$\begin{aligned} \int_{\mathbb{R}} |h_N(x)| f_{n+1}(t, x) dx &\leq \left(e^{Kt} + \int_0^t e^{K(t-s)} e^{(K+1)s} ds \right) \int_{\mathbb{R}} |x|^k f^0(x) dx \\ &= \left(1 + \int_0^t e^s ds \right) e^{Kt} \int_{\mathbb{R}} |x|^k f^0(x) dx = e^{(K+1)t} \int_{\mathbb{R}} |x|^k f^0(x) dx. \end{aligned}$$

By construction of f , passing to the limit $n \rightarrow \infty$ we obtain

$$\int_{\mathbb{R}} |h_N(x)| f(t, x) dx \leq e^{(K+1)t} \int_{\mathbb{R}} |x|^k f^0(x) dx.$$

By monotone convergence theorem, we can pass to the limit $N \rightarrow \infty$ to get

$$\int_{\mathbb{R}} |x|^k f(t, x) dx \leq e^{(K+1)t} \int_{\mathbb{R}} |x|^k f^0(x) dx. \quad \square$$

Thanks to these bounds, we will obtain more precise properties on the evolution of the (even) moments of the solution. In some cases, we do not necessarily have that the moments are differentiable in time, but we obtain estimates similar to differential inequalities. A good framework to work with continuous functions of time is set in the following lemma.

Lemma 2.2. *Given $t \mapsto y(t)$ a function of time, we write*

$$\frac{d^+}{dt} y(t) = \limsup_{\tau \rightarrow 0, \tau > 0} \frac{y(t + \tau) - y(t)}{\tau}.$$

Then, for any $T > 0$, if y is continuous and satisfies $\frac{d^+}{dt} y(t) \leq 0$ on $[0, T]$, then it is nonincreasing on $[0, T]$.

Consequently, if, for some integrable functions a and b , we have a continuous function y satisfying the following Grönwall-like estimate for all time $t \in [0, T]$:

$$\frac{d^+}{dt} y(t) \leq a(t)y(t) + b(t),$$

then we have for all $t \in [0, T]$,

$$y(t) \leq y(0)e^{\int_0^t a(s)ds} + \int_0^t b(u)e^{\int_u^t a(s)ds} du.$$

Proof. The proof (by contraposition) is based on a result similar to the mean value theorem : if there exist some times $t_0 < t_1$ such that y is continuous on $[t_0, t_1]$ and $y(t_0) < y(t_1)$, then there exists $t_* \in [t_0, t_1)$ such that $\liminf_{\tau \rightarrow 0, \tau > 0} \frac{y(t_* + \tau) - y(t_*)}{\tau} \geq \frac{y(t_1) - y(t_0)}{t_1 - t_0}$. Thus, we can conclude that $\frac{d^+}{dt} y(t_*) \geq \frac{y(t_1) - y(t_0)}{t_1 - t_0} > 0$. Surprisingly, we see that we only need the liminf of the slope to be nonpositive to show that a continuous function is nonincreasing (the proof also works for lower-semicontinuous functions). This result may be classical, but we did not find any reference, this is why we include the proof here.

To show the claimed result, if there exist such times t_0 and t_1 , we introduce

$$z(s) = y((1-s)t_0 + st_1) - (1-s)y(t_0) - sy(t_1).$$

The function z is continuous on $[0, 1]$, and satisfies $z(0) = z(1) = 0$. Therefore it reaches its minimum on $s_* \in [0, 1]$, and without loss of generality we may assume $s_* \neq 1$ (since in that case the minimum is also reached at $s = 0$). We therefore get $\liminf_{s \rightarrow s_*, s > s_*} \frac{z(s) - z(s_*)}{s - s_*} \geq 0$, which corresponds to the claim, with $t_* = (1 - s_*)t_0 + s_*t_1$.

Finally, the resolution of the Grönwall-like estimate is a direct consequence of the first part of the lemma, by setting $z(t) = y(t)e^{-\int_0^t a(s)ds} - \int_0^t b(u)e^{-\int_0^u a(s)ds} du$. Then the continuous function z is such that $\frac{d^+}{dt} z \leq 0$, therefore nonincreasing. \square

We are now ready to provide this kind of estimates for the even moments of the solution, which will be useful when we derive estimates for the moments of the centered and normalized distribution. We proceed to the proof of the following proposition.

Proposition 2.3. *Let $f^0 \in \mathcal{M}_+$, m be measurable and bounded below by $-K \leq 0$, and f be the global solution to the model (1) given by Theorem 1.1.*

For $k \geq 0$, if $\int_{\mathbb{R}} |x|^k f^0 dx < +\infty$, then $t \mapsto \mu_k(f(t, \cdot))$ is continuous. Furthermore, if k is even, for all $t \geq 0$, we have

$$\frac{d^+}{dt} \mu_k(f(t, \cdot)) \leq \mu_k(B_0[f(t, \cdot)]) + K \mu_k(f(t, \cdot)). \quad (23)$$

Proof. We fix $T > 0$. For $0 \leq t < t + \tau < T$, and h_N as in (22), we write, thanks to (18) and to Proposition 2.1 :

$$\begin{aligned} \int_{\mathbb{R}} h_N(x) f(t + \tau, x) dx &= \int_{\mathbb{R}} e^{-m(x)\tau} h_N(x) f(t, x) dx \\ &\quad + \int_0^\tau \int_{\mathbb{R}} e^{-m(x)(\tau-s)} h_N(x) B_0[f(t + s, \cdot)](x) dx ds. \end{aligned}$$

By dominated convergence as $N \rightarrow +\infty$, thanks to Lemma 2.1, Proposition 2.2 and the fact that $m(x) \geq -K$, we obtain

$$\mu_k(f(t + \tau, \cdot)) = \int_{\mathbb{R}} e^{-m(x)\tau} x^k f(t, x) dx + \int_0^\tau \int_{\mathbb{R}} e^{-m(x)(\tau-s)} x^k B_0[f(t + s, \cdot)](x) dx ds. \quad (24)$$

In particular, by setting $t = 0$ and writing t instead of τ , we get that for all $t \geq 0$,

$$\mu_k(f(t, \cdot)) = \int_{\mathbb{R}} e^{-m(x)t} x^k f^0(x) dx + \int_0^t \int_{\mathbb{R}} e^{-m(x)(t-s)} x^k B_0[f(s, \cdot)](x) dx ds. \quad (25)$$

We write this expression under the form $\mu(f(t, \cdot)) = y(t) + \int_0^T z(t, s) ds$, by setting

$$\begin{aligned} y(t) &= \int_{\mathbb{R}} e^{-m(x)t} x^k f^0(x) dx, \\ z(t, s) &= \begin{cases} \int_{\mathbb{R}} e^{-m(x)(t-s)} x^k B_0[f(s, \cdot)](x) dx & \text{if } s < t \\ 0 & \text{if } s \geq t \end{cases}. \end{aligned}$$

By dominated convergence, since $e^{-m(x)t} |x|^k \leq e^{KT} |x|^k$ which is f^0 integrable, we obtain that y is continuous on $[0, T]$. Similarly $t \mapsto z(t, s)$ is continuous on $(s, T]$ (thanks to Proposition 2.2, Lemma 2.1 and the fact that $m(x) \geq -K$) and we also have

$$|z(t, s)| \leq e^{K(t-s)} \int_{\mathbb{R}} |x|^k f(s, x) dx \leq e^{K(t-s) + (K+1)s} \int_{\mathbb{R}} |x|^k f^0(x) dx.$$

Consequently, $z(t, s)$ is uniformly bounded on $[0, T]^2$ by some $C > 0$.

Obviously, $t \mapsto z(t, s)$ is also continuous on $[0, s)$, since it is equal to zero. We now want to prove that $t \mapsto \int_0^T z(t, s) ds$ is continuous. We fix $t_0 \in [0, T]$, and we have that for all $s \neq t_0$, $z(t, s)$ converges to $z(t_0, s)$ as $t \rightarrow t_0$, by the continuity properties of z . So by dominated convergence, we get that $\int_0^T z(t, s) ds$ converges to $\int_0^T z(t_0, s) ds$ as $t \rightarrow t_0$. This ends the proof of the time continuity of $\mu_k(f(t, \cdot))$.

We now suppose that k is even. From (24), we obtain

$$\begin{aligned} \mu_k(f(t + \tau, \cdot)) &\leq \mu_k(f(t, \cdot)) e^{K\tau} + \int_0^\tau e^{K(\tau-s)} \mu_k(B_0[f(t + s, \cdot)]) ds \\ &\leq e^{K\tau} \left(\mu_k(f(t, \cdot)) + \int_0^\tau \mu_k(B_0[f(t + s, \cdot)]) ds \right). \end{aligned}$$

We therefore obtain

$$\frac{\mu_k(f(t+\tau, \cdot)) - \mu_k(f(t, \cdot))}{\tau} \leq \frac{e^{K\tau} - 1}{\tau} \mu_k(f(t, \cdot)) + \frac{e^{K\tau}}{\tau} \int_0^\tau \mu_k(B_0[f(t+s, \cdot)]) ds. \quad (26)$$

By the formula (21) and the fact that all moments of f of order $j \leq k$ are continuous in time, we get that $t \mapsto \mu_k(B_0[f(t, \cdot)])$ is continuous, and therefore the right-hand side of (26) converges to $K\mu_k(f(t, \cdot)) + \mu_k(B_0[f(t, \cdot)])$ as $\tau \rightarrow 0$, and this proves (23). \square

2.3 Centered moments of the normalized population density

To obtain estimates on the centered moments of f , the inequalities provided by Proposition 2.3 will not be sufficient ; we will first need a more precise description of the evolution of the moments μ_k (in particular we will have to compute the time derivatives of μ_0 and μ_1). To this aim, we make the stronger assumption 1.1 on the growth of m .

Under this assumption, if f^0 has finite second moment, then the evolution of quantities of the form $\int h(x)f(t, x)dx$ (when h is measurable and bounded) provides a simpler weak formulation of (1), as stated in the next proposition.

Proposition 2.4. *Under assumption 1.1, we suppose that f^0 has finite second moment and denote by f the global solution to the model (1) given by Theorem 1.1. Then for all measurable and bounded function h , the quantity $\int_{\mathbb{R}} h(x)f(t, x)dx$ is differentiable in time and we have*

$$\frac{d}{dt} \left(\int_{\mathbb{R}} h(x)f(t, x)dx \right) = \int_{\mathbb{R}} h(x)(B_0[f(t, \cdot)](x) - m(x)f(t, x))dx. \quad (27)$$

Proof. From Proposition 2.2 and Assumption 1.1, for all measurable and bounded function h , we want to obtain dominations of the time derivatives of the integrands appearing in (18) on $[0, T]$. First of all, we have

$$|h(x)m(x)e^{-m(x)t}| \leq Ce^{KT}(1+x^2) \sup_{\mathbb{R}} |h|,$$

and the right-hand side is f^0 -integrable, therefore by dominated convergence theorem we get that the term $\int_{\mathbb{R}} h(x)e^{-m(x)t}f^0(x)dx$ is differentiable in time. We also have for all $s \in [0, T]$, and for all $t \geq s$,

$$|h(x)m(x)e^{-m(x)(t-s)}| \leq Ce^{KT}(1+x^2) \sup_{\mathbb{R}} |h|,$$

which is $B_0[f(s, \cdot)]$ -integrable for all $s \in [0, T]$, thanks to Lemma 2.1 and Proposition 2.2. Therefore by dominated convergence the derivative of $\int_{\mathbb{R}} h(x)e^{-m(x)(t-s)}B_0[f(s, \cdot)](x)dx$ with respect to variable t is well defined, and uniformly bounded, on $[s, T]$. By Leibniz rule on the formula (18), the derivative of $\int_{\mathbb{R}} h(x)f(t, x)dx$ is well defined for all $t > 0$ and we have

$$\begin{aligned} \frac{d}{dt} \left(\int_{\mathbb{R}} h(x)f(t, x)dx \right) &= - \int_{\mathbb{R}} h(x)m(x)e^{-m(x)t}f^0(x)dx + \int_{\mathbb{R}} h(x)B_0[f(t, \cdot)](x)dx \\ &\quad - \int_0^t \int_{\mathbb{R}} h(x)m(x)e^{-m(x)(t-s)}B_0[f(s, \cdot)](x)dx ds. \end{aligned}$$

Once more, by dominated convergence theorem, we can approximate $h(x)m(x)$ by a bounded measurable function $\tilde{h}_N(x)$ and obtain, applying (18) to \tilde{h}_N , that in the limit $N \rightarrow \infty$, we have

$$\begin{aligned} \int_{\mathbb{R}} h(x)m(x)f(s, x)dx &= \int_{\mathbb{R}} h(x)m(x)e^{-m(x)t}f^0(x)dx \\ &\quad + \int_0^t \int_{\mathbb{R}} h(x)m(x)e^{-m(x)(t-s)}B_0[f(s, \cdot)](x)dx ds. \end{aligned}$$

We therefore obtain

$$\frac{d}{dt} \left(\int_{\mathbb{R}} h(x) f(t, x) dx \right) = \int_{\mathbb{R}} h(x) B_0[f(s, \cdot)](x) dx - \int_{\mathbb{R}} h(x) m(x) f(s, x) dx.$$

Hence, the second weak formulation (27). \square

Remark 2.1. We chose in Assumption 1.1 to have at most quadratic growth of the mortality rate m . If instead we suppose that it does not grow faster than $|x|^k$, the same reasoning would apply, and Proposition 2.4 would still apply, replacing the requirement on f^0 by asking $\int_{\mathbb{R}} |x|^k f^0(x) dx$ to be finite.

We are now ready to give the evolution equations for the moments of the solution f , which we state in the following proposition.

Proposition 2.5. Let $k \in \mathbb{N}$. Under Assumption 1.1, if f^0 has a moment of order $k + 2$ and f is the unique weak solution to the model (1), then $\mu_k(f)$ is continuously differentiable in time and we have for all $t \geq 0$

$$\frac{d}{dt} \mu_k(f(t, \cdot)) = \sum_{j=0}^k \binom{k}{j} \frac{\mu_j(f(t, \cdot)) \mu_{k-j}(f(t, \cdot))}{2^k \mu_0(f(t, \cdot))} - \int_{\mathbb{R}} m(x) x^k f(t, x) dx. \quad (28)$$

Proof. Let us remark that by Definition 1.1 we obtain that $\mu_0(f(t, \cdot)) > 0$ for any time. If not, we would have $\int_{\mathbb{R}} e^{-m(x)t} f^0(x) dx = 0$ and therefore $e^{-m(x)t}$ would be zero almost everywhere with respect to the measure f^0 , giving a contradiction. By Proposition 2.4, using again the function h_N defined in (22), we obtain

$$\begin{aligned} \int_{\mathbb{R}} h_N(x) f(t, x) dx &= \int_{\mathbb{R}} h_N(x) f^0(x) dx \\ &+ \int_0^t \left(\int_{\mathbb{R}} h_N(x) B_0[f(s, \cdot)](x) dx - \int_{\mathbb{R}} h_N(x) m(x) f(s, x) dx \right) ds. \end{aligned}$$

Thanks to Proposition 2.2 and to Assumption 1.1, $m(\cdot) f(s, \cdot)$ has a moment of order k , uniformly on $[0, t]$, therefore we obtain by dominated convergence theorem

$$\mu_k(f(t, \cdot)) = \mu_k(f^0) + \int_0^t (\mu_k(B_0[f(s, \cdot)]) - \mu_k(m(\cdot) f(s, \cdot))) ds.$$

Then, as in the proof of Proposition (2.3) for the time continuity of $\mu_k(f(t, \cdot))$, we get similarly as in (25) that

$$\mu_k(m(\cdot) f(t, \cdot)) = \int_{\mathbb{R}} e^{-m(x)t} m(x) x^k f^0(x) dx + \int_0^t \int_{\mathbb{R}} e^{-m(x)(t-s)} m(x) x^k B_0[f(s, \cdot)](x) dx ds.$$

And we proceed in the same way, by dominated convergence thanks to Assumption 1.1, to prove that $\mu_k(m(\cdot) f(t, \cdot))$ is continuous in time. Thanks to Propositions 2.2 and (2.3), we have the time continuity of $\mu_k(B_0[f(t, \cdot)])$ and therefore $\mu_k(f(t, \cdot))$ is continuously differentiable in time and satisfies (28). \square

From now on, we suppose that Assumption 1.1 is satisfied and that f^0 (of positive mass ρ^0) has a finite second moment. We are interested in the normalized population distribution. Let $\rho(t)$ be the total population mass at time t

$$\rho(t) = \mu_0(f) = \int_{\mathbb{R}} f(t, x) dx.$$

From Proposition 2.4, we obtain the evolution equation for $\rho(t)$

$$\frac{d}{dt}\rho(t) = \rho(t) - \int_{\mathbb{R}} m(x)f(t, x) dx. \quad (29)$$

As noted in the proof of Proposition 2.5, we have that ρ is positive for all time, and we can therefore consider the normalized population distribution g that is defined in (1.2).

Thus, from Proposition 2.4, we obtain a weak formulation for the following equation on g

$$\partial_t g(t, x) = B_0[g](t, x) - m(x)g(t, x) - \left(1 - \int_{\mathbb{R}} m(z)g(t, z) dz\right) g(t, x).$$

Remark 2.2. *This evolution equation for the normalized population distribution has the same form as the replicator-mutator equation, with the linear mutator term replaced by a sexual reproduction one. Replicator-mutator type equations have been quite studied in the form of nonlocal integro-differential equations (see e.g. [1, 2, 26]). Recently, the authors of [13] have proved concentration results and new convergence estimates in the non-singular case for measure solutions of the “house of cards” replicator-mutator model.*

Then, we can also define the mean trait, or center of mass, which we denote by \bar{x} and which is formulated by (6). Thus, we can focus on the study of the centered moments M_k and the centered variations of the selection S_k , for $k \in \mathbb{N}$, which formulations are given in (4)-(5).

Indeed, the moment M_k is well-defined for all time whenever f^0 has a moment of order k and S_k is well-defined if f^0 has a moment of order $k+2$. From now on, we will write $M_k^0 = M_k(0)$.

In particular we have $M_0(t) = 1$ and $M_1(t) = 0$ for all time. Using (28) (for $k = 1$) and (29) we obtain the equation of evolution of \bar{x} given by (8) whenever f^0 has a moment of order 3 :

$$\begin{aligned} \frac{d\bar{x}}{dt} &= \frac{d}{dt} \left(\frac{\mu_1(f(t, \cdot))}{\rho(t)} \right) = - \int_{\mathbb{R}} m(x)xg(t, x)dx + \frac{\mu_1(f(t, \cdot))}{\rho(t)} \int_{\mathbb{R}} m(x)g(t, x)dx \\ &= \int_{\mathbb{R}} m(x)(\bar{x}(t) - x)g(t, x)dx = -S_1(t), \end{aligned}$$

which was claimed in (8).

From now on, we actually assume that f^0 has a moment of order 4 (we will only compute estimates for even order moments), so that this mean trait \bar{x} is differentiable, and we can then compute the evolution of the centered moments M_k , thanks to Proposition 2.5 (or Proposition 2.3) and to the evolutions (29) and (8) of ρ and \bar{x} .

The next proposition states the differential equations and inequalities that M_k satisfies. They will be used in the following sections to prove stability results of Dirac masses. Notice that when f^0 has a moment of order 4, this proposition provides the results claimed in (9)-(10) for the evolutions of the moments M_2 and M_4 .

Proposition 2.6. *Let $k \geq 2$. Under Assumption 1.1, if f^0 has a moment of order $k+2$, then M_k is continuously differentiable in time and we have*

$$\frac{d}{dt}M_k = -\left(1 - \frac{1}{2^{k-1}}\right)M_k + \frac{1}{2^k} \sum_{i=2}^{k-2} \binom{k}{i} M_{k-i}M_i - S_k + S_0M_k + kS_1M_{k-1}. \quad (30)$$

If $k \geq 4$ is even and f^0 has only a moment of order k , then M_k is continuous in time and we have for all $t \geq 0$,

$$\frac{d^+}{dt}M_k \leq -\left(\frac{1}{2} - \frac{1}{2^{k-1}} - S_0 + \eta(\bar{x})\right)M_k + \frac{1}{2^k} \sum_{i=2}^{k-2} \binom{k}{i} M_{k-i}M_i + kS_1M_{k-1}. \quad (31)$$

Proof. Let us first remark that if we fix $x_0 \in \mathbb{R}$, then $\tilde{f} = f(\cdot, \cdot + x_0)$ satisfies the same equation with m replaced by $\tilde{m} = m(\cdot + x_0)$.

$$\begin{aligned}\rho(t)M_k(t) &= \int_{\mathbb{R}} (x - x_0 + x_0 - \bar{x}(t))^k f(t, x) dx = \sum_{i=0}^k \binom{k}{i} (x_0 - \bar{x}(t))^i \mu_{k-i}(\tilde{f}(t, \cdot)) \\ &= \mu_k(\tilde{f}(t, \cdot)) - k(\bar{x}(t) - x_0)\mu_{k-1}(\tilde{f}(t, \cdot)) + \mathcal{O}(|\bar{x}(t) - x_0|^2),\end{aligned}$$

where the bounds for the term $\mathcal{O}(|\bar{x}(t) - x_0|^2)$ are uniform on $[0, T]$, thanks to Proposition 2.2. Therefore if $x_0 = \bar{x}(t_0)$ for some $t_0 \in [0, T]$, since $t \mapsto \mu_{k-1}(\tilde{f}(t, \cdot))$ is continuous in time (and in particular at t_0), and using (8), we have that

$$\frac{\rho(t)M_k(t) - \mu_k(\tilde{f}(t, \cdot))}{t - t_0} \xrightarrow{t \rightarrow t_0} -k \frac{d\bar{x}}{dt}(t_0) \mu_{k-1}(\tilde{f}(t_0, \cdot)) = kS_1(t_0) \rho(t_0) M_{k-1}(t_0).$$

Therefore, setting $y(t) = \rho(t)M_k(t) - \mu_k(\tilde{f}(t, \cdot))$, we get that y is differentiable at t_0 with

$$\frac{dy}{dt} \Big|_{t=t_0} = kS_1(t_0) \rho(t_0) M_{k-1}(t_0).$$

When f has a moment of order $k+2$ (initially), this is the same for \tilde{f} and therefore we can use Proposition 2.5 to get the derivative of $\mu_k(\tilde{f}(t, \cdot))$ at t_0 . We obtain

$$\begin{aligned}\frac{d}{dt} \mu_k(\tilde{f}(t, \cdot)) \Big|_{t=t_0} &= \sum_{j=0}^k \binom{k}{j} \frac{\mu_j(\tilde{f}(t_0, \cdot)) \mu_{k-j}(\tilde{f}(t_0, \cdot))}{2^k \mu_0(\tilde{f}(t_0, \cdot))} - \int_{\mathbb{R}} \tilde{m}(x) x^k \tilde{f}(t_0, x) dx \\ &= \frac{1}{2^k} \sum_{j=0}^k \binom{k}{j} \rho(t_0) M_j(t_0) M_{k-j}(t_0) - \rho(t_0) (S_k(t_0) + m(\bar{x}(t_0)) M_k(t_0)).\end{aligned}$$

Since $M_k(t) = \frac{1}{\rho(t)}(y(t) + \mu_k(\tilde{f}(t, \cdot)))$ and since ρ is differentiable at t_0 , we get that

$$\frac{d}{dt} M_k \Big|_{t=t_0} = \frac{1}{\rho(t_0)} \frac{d}{dt} (y(t) + \mu_k(\tilde{f}(t, \cdot))) \Big|_{t=t_0} - \frac{M_k(t_0)}{\rho(t_0)} \frac{d\rho}{dt} \Big|_{t=t_0}.$$

The derivative of ρ at t_0 is given in (29) by $\rho(t_0)(1 - S_0(t_0) - m(\bar{x}(t_0)))$, therefore we obtain (30) at time t_0 , using the fact that $M_0(t_0) = 1$ and $M_1(t_0) = 0$. Since the right-hand side of (30) is continuous in time, as proved in Proposition 2.5, we obtain that M_k is continuously differentiable in time.

When f has only (initially) a moment of order k with k even and $k \geq 4$, we use similarly Proposition 2.3 instead of 2.5 to get, by setting $K = -\inf_{\mathbb{R}} m$,

$$\begin{aligned}\frac{d^+}{dt} \mu_k(\tilde{f}(t, \cdot)) \Big|_{t=t_0} &\leq \sum_{j=0}^k \binom{k}{j} \frac{\mu_j(\tilde{f}(t_0, \cdot)) \mu_{k-j}(\tilde{f}(t_0, \cdot))}{2^k \mu_0(\tilde{f}(t_0, \cdot))} + K \mu_k(\tilde{f}(t_0, \cdot)) \\ &\leq \frac{1}{2^k} \sum_{j=0}^k \binom{k}{j} \rho(t_0) M_j(t_0) M_{k-j}(t_0) + \rho(t_0) K M_k(t_0).\end{aligned}$$

Since $k \geq 4$, ρ and y are still differentiable at t_0 , and since ρ is positive, we get that

$$\frac{d^+}{dt} M_k \Big|_{t=t_0} = \frac{1}{\rho(t_0)} \left(\frac{dy}{dt} \Big|_{t=t_0} + \frac{d^+ \mu_k(\tilde{f}(t, \cdot))}{dt} \Big|_{t=t_0} \right) - \frac{M_k(t_0)}{\rho(t_0)} \frac{d\rho}{dt} \Big|_{t=t_0}.$$

The computations are therefore similar to the previous ones, and we obtain

$$\frac{d^+}{dt}M_k \leq -\left(1 - \frac{1}{2^{k-1}} - S_0 - K - m(\bar{x})\right)M_k + \frac{1}{2^k} \sum_{i=2}^{k-2} \binom{k}{i} M_{k-i}M_i + kS_1M_{k-1},$$

which corresponds to (31). \square

Remark 2.3. *If we suppose that m grows at most linearly, we only need to have a second moment to control the derivative of ρ , and therefore (31) is also valid for $k = 2$.*

2.4 Fourier formulation, and rescaled profile

Before tackling the long time behavior of \bar{x} and of the centered moments M_k , we finish this section by stating the equation satisfied by the Fourier transform of the centered and rescaled formulation of $g(t, \cdot)$. This equation will be further investigated in Section 4.

From Definition 15, we have

$$\int_{\mathbb{R}} \gamma(t, x) dx = 1, \int_{\mathbb{R}} x\gamma(t, x) dx = 0, \text{ and } \int_{\mathbb{R}} x^2\gamma(t, x) dx = 1. \quad (32)$$

We also define the signed measure φ

$$\varphi(t, x) := [m(\sqrt{M_2(t)}x + \bar{x}(t)) - m(\bar{x}(t))]\gamma(t, x), \quad (33)$$

Now we prove the following result.

Proposition 2.7. *Under Assumption 1.1, if f^0 has a moment of order 4, then $\hat{\gamma}$ is differentiable and we have the equation on $\hat{\gamma}$*

$$\partial_t \hat{\gamma}(t, \xi) = \hat{\gamma}(t, \frac{\xi}{2})^2 - \hat{\gamma}(t, \xi) + \frac{1}{4} \xi \partial_\xi \hat{\gamma}(t, \xi) + R(t, \xi), \quad (34)$$

with $R(t, \xi)$ given by

$$R(t, \xi) := -iS_1(t) \frac{\xi}{\sqrt{M_2(t)}} \hat{\gamma}(t, \xi) - \hat{\varphi}(t, \xi) + S_0(t) \hat{\gamma}(t, \xi) + \frac{1}{2} \left(\frac{S_2(t)}{M_2(t)} - S_0(t) \right) \xi \partial_\xi \hat{\gamma}(t, \xi). \quad (35)$$

Proof. From Assumption 1.1 and Proposition 2.4, and assuming that f^0 has finite moments up to order 4, \hat{f} is differentiable in time, and using (17), we obtain

$$\partial_t \hat{f}(t, \xi) = \frac{\left(\hat{f}(t, \xi/2)\right)^2}{\hat{f}(t, 0)} - \widehat{mf}(t, \xi).$$

Thus, we obtain the equation on \hat{g} , using (29):

$$\partial_t \hat{g}(t, \xi) = (\hat{g}(t, \xi/2))^2 - \hat{g}(t, \xi) - (1 - \widehat{mg}(t, 0)) \hat{g}(t, \xi). \quad (36)$$

Then, $\hat{\gamma}$ is well defined and we have, for all $\xi \in \mathbb{R}$,

$$\hat{\gamma}(t, \xi) = e^{i \frac{\xi}{\sqrt{M_2(t)}} \bar{x}(t)} \hat{g}\left(t, \frac{\xi}{\sqrt{M_2(t)}}\right).$$

Thus, from Proposition 2.6 and (36), $\hat{\gamma}$ is differentiable in time and we can compute

$$\begin{aligned}\partial_t \hat{\gamma}(t, \xi) = & i \frac{\xi}{\sqrt{M_2}} \frac{d\bar{x}}{dt} \hat{\gamma}(t, \xi) - i \frac{\xi \bar{x}}{2M_2^{3/2}} \frac{dM_2}{dt} \hat{\gamma}(t, \xi) + e^{i \frac{\xi}{\sqrt{M_2(t)}} \bar{x}(t)} \partial_t \hat{g}(t, \xi / \sqrt{M_2}) \\ & - e^{i \frac{\xi}{\sqrt{M_2(t)}} \bar{x}} \frac{1}{2M_2^{3/2}} \frac{dM_2}{dt} \xi \partial_\xi \hat{g}(t, \xi / \sqrt{M_2}),\end{aligned}$$

with $\partial_\xi \hat{g}$ well defined since g has moment of order one. Using (8) and (36), we get

$$\begin{aligned}\partial_t \hat{\gamma}(t, \xi) = & (\hat{\gamma}(t, \xi/2))^2 - \hat{\gamma}(t, \xi) - e^{i \frac{\xi}{\sqrt{M_2(t)}} \bar{x}(t)} \widehat{mg}(t, \frac{\xi}{\sqrt{M_2}}) + \widehat{mg}(t, 0) \hat{\gamma}(t, \xi) \\ & - i S_1(t) \frac{\xi}{\sqrt{M_2}} \hat{\gamma}(t, \xi) - \frac{1}{2} \frac{dM_2}{dt} \frac{\xi}{M_2^{3/2}} \left(i \bar{x} \hat{\gamma}(t, \xi) + e^{i \frac{\xi}{\sqrt{M_2(t)}} \bar{x}(t)} \partial_\xi \hat{g}(t, \frac{\xi}{\sqrt{M_2}}) \right).\end{aligned}$$

And using (9), we then compute

$$\begin{aligned}\partial_t \hat{\gamma}(t, \xi) = & (\hat{\gamma}(t, \xi/2))^2 - \hat{\gamma}(t, \xi) - e^{i \frac{\xi}{\sqrt{M_2(t)}} \bar{x}(t)} \widehat{mg}(t, \frac{\xi}{\sqrt{M_2}}) + \widehat{mg}(t, 0) \hat{\gamma}(t, \xi) - i S_1(t) \frac{\xi}{\sqrt{M_2}} \hat{\gamma}(t, \xi) \\ & + \frac{\xi}{2} \left(\frac{1}{2} \frac{1}{\sqrt{M_2}} + \frac{S_2}{M_2^{3/2}} - \frac{S_0}{\sqrt{M_2}} \right) \left(i \bar{x} \hat{\gamma}(t, \xi) + e^{i \frac{\xi}{\sqrt{M_2(t)}} \bar{x}(t)} \partial_\xi \hat{g}(t, \frac{\xi}{\sqrt{M_2}}) \right).\end{aligned}$$

Then, noticing the expressions

$$\begin{aligned}\partial_\xi \hat{\gamma}(t, \xi) = & \frac{1}{\sqrt{M_2}} \left(i \bar{x} \hat{\gamma}(t, \xi) + e^{i \frac{\xi}{\sqrt{M_2(t)}} \bar{x}(t)} \partial_\xi \hat{g}(t, \frac{\xi}{\sqrt{M_2}}) \right), \\ \widehat{mg}(t, 0) = & S_0(t) - m(\bar{x}(t)),\end{aligned}$$

the definition (33) enables to write

$$e^{i \frac{\xi}{\sqrt{M_2(t)}} \bar{x}(t)} \widehat{mg}(t, \frac{\xi}{\sqrt{M_2}}) - \widehat{mg}(t, 0) \hat{\gamma}(t, \xi) = \widehat{\varphi}(t, \xi) - S_0(t) \hat{\gamma}(t, \xi),$$

and we obtain the equation on $\hat{\gamma}$

$$\begin{aligned}\partial_t \hat{\gamma}(t, \xi) = & (\hat{\gamma}(t, \xi/2))^2 - \hat{\gamma}(t, \xi) + \frac{1}{4} \xi \partial_\xi \hat{\gamma}(t, \xi) - i S_1(t) \frac{\xi}{\sqrt{M_2}} \hat{\gamma}(t, \xi) - \widehat{\varphi}(t, \xi) \\ & + S_0(t) \hat{\gamma}(t, \xi) + \frac{1}{2} \left(\frac{S_2}{M_2} - S_0 \right) \xi \partial_\xi \hat{\gamma}(t, \xi).\end{aligned}$$

Finally, with the definition of R (35), we get (34). \square

3 Long-time behaviour of normalized moments

3.1 Local stability of some Dirac masses

We begin this section by proving the local stability result of Theorem 1.2, which states that a sufficient condition for the convergence of g towards a Dirac mass is to be initially close enough to the Dirac mass located at $\bar{x}_0 = \bar{x}(0)$, when $\eta(\bar{x}_0) > 0$, i.e. when the initial selection rate at the center of mass satisfies $m(\bar{x}_0) < \inf_{\mathbb{R}} m + \frac{1}{2}$.

Proof. Thanks to the definitions (4), (5) and (7), we obtain

$$S_2 \geq (\eta(\bar{x}) - \frac{1}{2}) M_2, \quad S_4 \geq (\eta(\bar{x}) - \frac{1}{2}) M_4. \quad (37)$$

By Assumption 1.2, using Cauchy-Schwarz inequalities in the definitions (4)-(5), we obtain

$$M_2 \leq \sqrt{M_4}, \quad (38)$$

$$|M_3| \leq \sqrt{M_2 M_4}, \quad (39)$$

$$|S_0| \leq \alpha(\bar{x})\sqrt{M_2} + \beta M_2, \quad (40)$$

$$|S_1| \leq \alpha(\bar{x})M_2 + \beta\sqrt{M_2 M_4}. \quad (41)$$

Therefore we obtain, thanks to (8), (9), and (10):

$$\left| \frac{d\bar{x}}{dt} \right| \leq \alpha(\bar{x})M_2 + \beta\sqrt{M_2 M_4}, \quad (42)$$

$$\frac{d}{dt}M_2 \leq (-\eta(\bar{x}) + \alpha(\bar{x})\sqrt{M_2} + \beta M_2)M_2, \quad (43)$$

$$\frac{d^+}{dt}M_4 \leq (-\eta(\bar{x}) + 5(\alpha(\bar{x})\sqrt{M_2} + \beta M_2))M_4, \quad (44)$$

where we have used (38) and (39) to get that $|M_2 M_3| \leq \sqrt{M_4}\sqrt{M_2 M_4} = \sqrt{M_2}M_4$.

We now fix δ_0 such that $0 < \delta < \delta_0 < \eta(\bar{x}_0)$, and by continuity of m we find $r > 0$ such that $\eta(x) \geq \delta_0$ for all $x \in [\bar{x}_0 - r, \bar{x}_0 + r]$. We set $\tilde{\alpha} = \sup_{[\bar{x}_0 - r, \bar{x}_0 + r]} \alpha$, which is finite by continuity of α . Using (38), we then fix M_4^0 sufficiently small in order to have

$$\begin{aligned} 5(\tilde{\alpha}\sqrt{M_2^0} + \beta M_2^0) &< \delta_0 - \delta, \\ \tilde{\alpha}M_2^0 + \beta\sqrt{M_2^0 M_4^0} &< r\delta. \end{aligned}$$

Finally we define T as the supremum of all $t > 0$ such that for all $s \in [0, t]$, we have

$$\begin{cases} |\bar{x}(s) - \bar{x}_0| \leq r, \\ 5\tilde{\alpha}\sqrt{M_2(s)} + \beta M_2(s) \leq \delta_0 - \delta, \\ \tilde{\alpha}M_2(s) + \beta\sqrt{M_2(s)M_4(s)} \leq r\delta, \end{cases} \quad (45)$$

and the set of such $t > 0$ is non empty, therefore $T > 0$ (we may have $T = +\infty$).

For all $t \in [0, T]$, we get from (43)-(44) that $\frac{d}{dt}M_2 \leq -\delta M_2$ and $\frac{d^+}{dt}M_4 \leq -\delta M_4$, giving that M_2 and M_4 are nonincreasing in time and satisfying estimates (12)-(13), thanks to Lemma 2.2. Thanks to (42), we obtain

$$\begin{aligned} |\bar{x}(t) - \bar{x}_0| &\leq \int_0^t (\tilde{\alpha}M_2^0 + \beta\sqrt{M_2^0 M_4^0})e^{-\delta s} ds \\ &\leq (\tilde{\alpha}M_2^0 + \beta\sqrt{M_2^0 M_4^0})\frac{1 - e^{-\delta t}}{\delta} < (1 - e^{-\delta t})r. \end{aligned}$$

Let us prove that $T = +\infty$ by contradiction. If T is finite, then $|\bar{x}(T) - \bar{x}_0| \leq (1 - e^{-\delta T})r < r$, and by monotonicity of M_2 and M_4 we also get

$$\begin{aligned} 5\tilde{\alpha}\sqrt{M_2(T)} + \beta M_2(T) &\leq 5\tilde{\alpha}\sqrt{M_2^0} + \beta M_2^0 < \delta_0 - \delta, \\ \tilde{\alpha}M_2(T) + \beta\sqrt{M_2(T)M_4(T)} &\leq \tilde{\alpha}M_2^0 + \beta\sqrt{M_2^0 M_4^0} < r\delta. \end{aligned}$$

Therefore, by continuity in time of M_2 , M_4 and \bar{x} , there exists $\tau > 0$ such that for all s in $[T - \tau, T + \tau]$, we have the estimations (45). Thanks to the definition of T , these estimations

are true for all $s \in [0, T)$, and therefore also for all $s \in [0, T + \tau]$, in contradiction with the definition of T as a supremum.

Since the derivative of \bar{x} is decaying exponentially fast, we get that $\bar{x}(t)$ converges to some \bar{x}_∞ in \mathbb{R} , and more precisely we have

$$|\bar{x}(t) - \bar{x}_\infty| \leq \int_t^\infty (\tilde{\alpha} M_2^0 + \beta \sqrt{M_2^0 M_4^0}) e^{-\delta s} ds = \frac{\tilde{\alpha} M_2^0 + \beta \sqrt{M_2^0 M_4^0}}{\delta} e^{-\delta t}.$$

This also shows that the smaller the value of M_4^0 , the closer the location of \bar{x}_∞ relatively to \bar{x}_0 .

Finally, the convergence in Wasserstein distance comes from the fact that $W_4(g, \delta_{\bar{x}}) = M_4(t)^{\frac{1}{4}}$, and therefore

$$W_4(g, \delta_{\bar{x}_\infty}) \leq (M_4^0)^{\frac{1}{4}} e^{-\frac{\delta}{4}t} + W_4(\delta_{\bar{x}(t)}, \delta_{\bar{x}_\infty}) = (M_4^0)^{\frac{1}{4}} e^{-\frac{\delta}{4}t} + |\bar{x}(t) - \bar{x}_\infty|. \quad \square$$

Remark 3.1. If m is Lipschitz ($\beta = 0$ and α constant in Assumption 1.2), then we only need to suppose M_2^0 sufficiently small to get the estimates (11), (12) and the exponential convergence of \bar{x} . The proof is exactly the same, only using the estimate on the evolution of M_2 given by (31) (see Remark 2.3), and the convergence is in 2-Wasserstein distance.

3.2 Improved estimates on even order moments

We can now use this result of stability of Theorem 1.2 together with the equation for higher order moments to improve the bound on the rate of convergence. In this subsection, we start from the assumptions made in Theorem 1.2:

Assumption 3.1. We suppose that the selection function m satisfies Assumption 1.2, and we take \bar{x}_0 and δ such that $0 < \delta < \eta(\bar{x}_0)$. Finally, we suppose that the initial profile g_0 is such that $\bar{x}(0) = \bar{x}_0$ and that M_4^0 is sufficiently small, ensuring that the conclusion of Theorem 1.2 is true. As in the proof of this theorem, we denote by $\tilde{\alpha}$ a uniform bound in time for $\alpha(\bar{x}(t))$ (independent of M_4^0 , only depending on η, δ and \bar{x}_0).

To obtain the estimates on M_{2k} , the idea is to use the equations (30) and (31), along with the stability result stated in Theorem 1.2. Apart from the exponential decay rates of the moments, we will also need a control on the constants in front of these estimates. Furthermore we will often have to deal with quantities controlled by $e^{-\tilde{\lambda}t}$ for all $\tilde{\lambda}$ strictly less than some $\lambda > 0$ to be identified. Therefore we introduce the following definition.

Definition 3.1. For $\lambda > 0$, we say that a nonnegative quantity $y(t)$ (which may depend on several parameters) has a decay controlled by the parameter $\sigma \geq 0$ and with rate $[\lambda]^-$ if for all $\tilde{\lambda} < \lambda$, there exists a constant $C(\sigma) > 0$ converging to 0 when $\sigma \rightarrow 0$, and such that for all $t > 0$, we have $y(t) \leq C(\sigma) e^{-\tilde{\lambda}t}$.

We will use several times the following lemma.

Lemma 3.1. If y is a nonnegative continuous function in time satisfying for all $t > 0$

$$\frac{d^+}{dt} y(t) \leq (-\lambda + v_1(t)) y(t) + v_0(t),$$

where $v_0(t)$ and $v_1(t)$ have a decay controlled by the parameter σ and with respective rates $[\omega_0]^-$ and $[\omega_1]^-$, then $y(t)$ has a decay controlled by $\max(y(0), \sigma)$ and with rate $[\min(\lambda, \omega_0)]^-$.

Proof. If $\tilde{\lambda} < \min(\lambda, \omega_0)$, we fix $\tilde{\omega}_0$ and $\tilde{\omega}_1$ such that $\tilde{\lambda} < \tilde{\omega}_0 < \omega_0$ and $0 < \tilde{\omega}_1 < \omega_1$, and we have for all $t > 0$

$$\frac{d^+}{dt}y(t) \leq (-\tilde{\lambda} + C_1(\sigma)e^{-\tilde{\omega}_1 t})y(t) + C_2(\sigma)e^{-\tilde{\omega}_0 t},$$

with $C_1(\sigma)$ and $C_2(\sigma)$ converging to 0 as $\sigma \rightarrow 0$. Therefore, thanks to Lemma 2.2, we obtain

$$\begin{aligned} y(t)e^{\tilde{\lambda}t + \frac{C_1(\sigma)}{\tilde{\omega}_1}(e^{-\tilde{\omega}_1 t} - 1)} &\leq y(0) + C_2(\sigma) \int_0^t e^{\tilde{\lambda}s + \frac{C_1(\sigma)}{\tilde{\omega}_1}(e^{-\tilde{\omega}_1 s} - 1)} e^{-\tilde{\omega}_0 s} ds \\ &\leq y(0) + C_2(\sigma) \int_0^t e^{\tilde{\lambda}s - \tilde{\omega}_0 s} ds = y(0) + \frac{C_2(\sigma)}{\tilde{\omega}_0 - \tilde{\lambda}} (1 - e^{-(\tilde{\omega}_0 - \tilde{\lambda})t}) \\ &\leq y(0) + \frac{C_2(\sigma)}{\tilde{\omega}_0 - \tilde{\lambda}}. \end{aligned}$$

Finally we get

$$y(t) \leq \left(y(0) + \frac{C_2(\sigma)}{\tilde{\omega}_0 - \tilde{\lambda}} \right) e^{\frac{C_1(\sigma)}{\tilde{\omega}_1}(1 - e^{-\tilde{\omega}_1 t}) - \tilde{\lambda}t} \leq \left(y(0) + \frac{C_2(\sigma)}{\tilde{\omega}_0 - \tilde{\lambda}} \right) e^{\frac{C_1(\sigma)}{\tilde{\omega}_1}} \cdot e^{-\tilde{\lambda}t},$$

and this ends the proof, since the constant converges to 0 as $y(0)$ and σ converge to 0. \square

First, we obtain the next proposition that presents links between rates of convergence of M_2 and M_4 , allowing to improve the rates provided by Theorem 1.2.

Proposition 3.1. *Under Assumption 3.1 we have that M_2 and M_4 have decay controlled by M_4^0 and with respective rates $[\min(\frac{1}{2}, \frac{3}{8} + \delta)]^-$ and $[\frac{3}{8} + \delta]^-$.*

Proof. We first suppose that M_2 and M_4 have decay controlled by M_4^0 and with respective rates $[\lambda_2]^-$ and $[\lambda_4]^-$, for $\lambda_2, \lambda_4 > 0$, and show that these rates can be improved. Thanks to Assumption 1.2, using once again Cauchy-Schwarz inequality in the definition (5), we obtain

$$|S_2| \leq \tilde{\alpha} \sqrt{M_2 M_4} + \beta M_4. \quad (46)$$

Therefore we obtain, thanks to (9) and the estimate (40),

$$\begin{aligned} \frac{d}{dt}M_2 &\leq \left(-\frac{1}{2} + \tilde{\alpha} \sqrt{M_2} + \beta M_2 \right) M_2 + \tilde{\alpha} \sqrt{M_2 M_4} + \beta M_4, \\ &\leq \left(-\frac{1}{2} + \varepsilon + \tilde{\alpha} \sqrt{M_2} + \beta M_2 \right) M_2 + \left(\frac{\tilde{\alpha}^2}{4\varepsilon} + \beta \right) M_4, \end{aligned}$$

where we used Young's inequality for some $\varepsilon > 0$. Similarly, using (10), (37) and (11) we obtain

$$\frac{d^+}{dt}M_4 \leq \left(-\delta - \frac{3}{8} + 5(\tilde{\alpha} \sqrt{M_2} + \beta M_2) \right) M_4 + \frac{3}{8} M_2^2.$$

Therefore, since $\tilde{\alpha} \sqrt{M_2} + \beta M_2$, M_4 and M_2^2 have decay controlled by the initial moment M_4^0 and respective rates $[\frac{1}{2}\lambda_2]^-$, $[\lambda_4]^-$ and $[2\lambda_2]^-$, by applying Lemma 3.1, M_4 have decay controlled by M_4^0 with rate $[\min(\frac{3}{8} + \delta, 2\lambda_2)]^-$, while M_2 has decay controlled by $\max(M_2^0, M_4^0)$ (and therefore, thanks to (38), by M_4^0 only) and with rate $[\min(\frac{1}{2} - \varepsilon, \lambda_4)]^-$. Since this is true for any $\varepsilon > 0$, we obtain that M_2 has rate $[\min(\frac{1}{2}, \lambda_4)]^-$. Therefore, by applying once more this property, we obtain that M_2 and M_4 have decay controlled by M_4^0 and with respective rates $[\min(\frac{1}{2}, \frac{3}{8} + \delta, 2\lambda_2)]^-$ and $[\min(\frac{3}{8} + \delta, 2\lambda_4)]^-$. Indeed $\min(\frac{3}{8} + \delta, 1, 2\lambda_4) = \min(\frac{3}{8} + \delta, 2\lambda_4)$ since $\delta < \eta(\bar{x}_0) \leq \frac{1}{2}$.

Hence, since M_2 and M_4 both have decay controlled by M_4^0 and with rate $[\delta]^-$, thanks to Theorem 1.2 and the estimate (38), we directly obtain by induction that for all $k \in \mathbb{N}$, M_2 and M_4 have decay controlled by M_4^0 and with respective rates $[\min(\frac{1}{2}, \frac{3}{8} + \delta, 2^k \delta)]^-$ and $[\min(\frac{3}{8} + \delta, 2^k \delta)]^-$, so for k sufficiently large, this ends the proof. \square

The next proposition uses these estimates on M_2 and M_4 to obtain estimates for the rates of convergence of all moments of even order, provided they are finite initially.

Proposition 3.2. *Under Assumption 3.1, for all $k \geq 2$, if M_{2k}^0 is finite, then $M_{2k}(t)$ has decay controlled by M_{2k}^0 and with rate $[\frac{1}{2} - \frac{1}{2^{2k-1}} + \delta]^-$.*

Proof. We proceed by induction, since, the case $k = 2$ is provided by Proposition 3.1. We fix $k > 2$, and suppose the result is true for $k - 1$.

First of all, if $2i \leq \ell \leq 2j$, using Hölder inequality in the definition (4), we get

$$\begin{aligned} |M_\ell| &\leq \int_{\mathbb{R}} |x - \bar{x}(t)|^\ell g(x) dx = \int_{\mathbb{R}} |x - \bar{x}(t)|^{2i \frac{2j-\ell}{2j-2i} + 2j \frac{\ell-2i}{2j-2i}} g(x) dx \\ &\leq M_{2i}^{\frac{2j-\ell}{2j-2i}} M_{2j}^{\frac{\ell-2i}{2j-2i}}. \end{aligned} \quad (47)$$

In particular, for $i = 0$ and $j = k$, we obtain

$$M_2 \leq M_{2k}^{\frac{1}{k}}, \quad M_4 \leq M_{2k}^{\frac{2}{k}}, \quad M_{2k-2} \leq M_{2k}^{1-\frac{1}{k}}, \quad \text{and} \quad |M_{2k-1}| \leq M_{2k}^{1-\frac{1}{2k}}. \quad (48)$$

Using (41) to get

$$|S_1 M_{2k-1}| \leq (\tilde{\alpha} M_2 + \beta \sqrt{M_2 M_4}) M_{2k}^{1-\frac{1}{2k}} \leq (\tilde{\alpha} \sqrt{M_2} + \beta \sqrt{M_4}) M_{2k}. \quad (49)$$

As previously, thanks to the definitions (4), (5), (7) and the property (11), we obtain

$$S_{2k} \geq \left(\delta - \frac{1}{2}\right) M_{2k}. \quad (50)$$

Now, using (47) with $i = 1$ and $j = k - 1$, we obtain, when $2 \leq \ell \leq 2k - 2$:

$$|M_\ell| |M_{2k-\ell}| \leq M_2 M_{2k-2}, \quad (51)$$

and therefore, we use (50), (40), and (49) in the differential inequality (31) for M_{2k} to get

$$\frac{d^+}{dt} M_{2k} \leq \left(-\frac{1}{2} + \frac{1}{2^{2k-1}} - \delta + v_1(t)\right) M_{2k} + \left(\frac{1}{2^{2k}} \sum_{\ell=2}^{2k-2} \binom{2k}{\ell}\right) M_2 M_{2k-2}, \quad (52)$$

where $v_1(t) = (2k+1)\tilde{\alpha}\sqrt{M_2} + \beta(2k\sqrt{M_4} + M_2)$. Now, using (48), we obtain that when M_{2k}^0 converges to 0, then M_4^0 and M_2^0 also converge to 0. Therefore by induction, and thanks to Proposition 3.1, $v_1(t)$ and $M_2 M_{2k-2}$ have decay controlled by M_{2k}^0 and with respective rates $[\min(\frac{1}{4}, \frac{3}{16} + \frac{1}{2}\delta)]^-$ and $[\omega_0]^-$, where

$$\omega_0 = \left(\min\left(\frac{1}{2}, \frac{3}{8} + \delta\right) + \left(\frac{1}{2} - \frac{1}{2^{2k-3}} + \delta\right)\right) > \left(\frac{3}{8} + \frac{1}{2} - \frac{1}{2^{2k-3}} + \delta\right) \geq \frac{3}{4} + \delta, \quad (53)$$

since $k \geq 3$. Therefore, since $\frac{3}{4} + \delta > \frac{1}{2} - \frac{1}{2^{2k-1}} + \delta$, we can apply Lemma 3.1 and this ends the proof. \square

We can now use the estimate on a moment of high order to obtain better rates for all lower order moments, as stated in the following proposition.

Proposition 3.3. *We suppose that Assumption 3.1 is satisfied. If $M_{2k_0}^0$ is finite (for $k_0 \geq 2$), then for all $k \leq k_0$, M_{2k} has decay controlled by $M_{2k_0}^0$ with rate $[\min(1 - \frac{1}{2^{2k-1}}, \frac{1}{2} - \frac{1}{2^{2k_0-1}} + \delta)]^-$.*

Proof. We proceed once more by induction, from $k = k_0$ down to $k = 1$. Indeed, since we have $\delta < \eta(\bar{x}_0) \leq \frac{1}{2}$, we get that $\min(1 - \frac{1}{2^{2k_0-1}}, \frac{1}{2} - \frac{1}{2^{2k_0-1}} + \delta) = \frac{1}{2} - \frac{1}{2^{2k_0-1}} + \delta$, which corresponds to the rate given by Proposition 3.2. We fix $k < k_0$ and suppose that the result is true for $k + 1$. We have, thanks to Assumption 1.2 and using Young's inequality for some $\varepsilon > 0$:

$$\begin{aligned} |S_{2k}| &\leq \tilde{\alpha} \int_{\mathbb{R}} |x - \bar{x}(t)|^{2k+1} g(x) dx + \beta M_{2k+2} \\ &\leq \int_{\mathbb{R}} |x - \bar{x}(t)|^{2k} \left(\varepsilon + \frac{\tilde{\alpha}^2}{4\varepsilon} |x - \bar{x}(t)|^2 \right) g(x) dx + \beta M_{2k+2} \leq \varepsilon M_{2k} + \left(\frac{\tilde{\alpha}^2}{4\varepsilon} + \beta \right) M_{2k+2}. \end{aligned}$$

Therefore, as in the proof of Proposition 3.2, using (51), (40), and (49) in the differential equation (30) for M_{2k} , we get

$$\frac{d}{dt} M_{2k} \leq \left(-1 + \frac{1}{2^{2k-1}} + \varepsilon + v_1(t) \right) M_{2k} + \left(\frac{1}{2^{2k}} \sum_{\ell=2}^{2k-2} \binom{2k}{\ell} \right) M_2 M_{2k-2} + \left(\frac{\tilde{\alpha}^2}{4\varepsilon} + \beta \right) M_{2k+2},$$

where $v_1(t) = (2k+1)\tilde{\alpha}\sqrt{M_2} + \beta(2k\sqrt{M_4} + M_2)$ has decay controlled by M_4^0 (and hence by $M_{2k_0}^0$) and rate $[\min(\frac{1}{4}, \frac{3}{16} + \frac{1}{2}\delta)]^-$. By induction M_{2k+2} has decay controlled by $M_{2k_0}^0$ and with rate $[\tilde{\omega}_0]^-$ where $\tilde{\omega}_0 = \min(1 - \frac{1}{2^{2k+1}}, \frac{1}{2} - \frac{1}{2^{2k_0-1}} + \delta)$. And we know by Propositions 3.1 and 3.2 that $M_2 M_{2k-2}$ has decay controlled by M_{2k}^0 (and hence by $M_{2k_0}^0$) and with rate $[\omega_0]^-$ with ω_0 given in (53), so when $k \geq 3$ we have $\omega_0 > \frac{3}{4} + \delta > \frac{1}{2} - \frac{1}{2^{2k_0-1}} + \delta \geq \tilde{\omega}_0$.

When $k = 2$ we also obtain $\omega_0 = \min(\frac{3}{4} + 2\delta, 1) > \frac{1}{2} - \frac{1}{2^{2k_0-1}} + \delta \geq \tilde{\omega}_0$ since $\delta < \frac{1}{2}$. Finally when $k = 1$ the sum on ℓ is empty.

So, in all cases we get thanks to Lemma 3.1 that M_{2k} has decay controlled by $\max(M_{2k}^0, M_{2k_0}^0)$ (and therefore by $M_{2k_0}^0$ only) and with rate $[\min(1 - \frac{1}{2^{2k-1}} - \varepsilon, \tilde{\omega}_0)]^-$, for any $\varepsilon > 0$, which ends the proof. \square

When the initial condition has enough moments, we then see thanks to Proposition 3.3 that all M_{2k} with $k \geq 2$ decay with rates strictly greater than $\frac{1}{2}$, and that M_2 decays with rate $[\frac{1}{2}]^-$. Therefore we expect all the moments M_{2k} to decay “faster” than M_2 , and we will see in the next section that this property is important to study the self-similar behaviour of the solution as time goes to infinity. However, to have such a property, we have to get a lower bound for $M_2(t)$. By using a refinement of Lemma 3.1, we could prove that there exists $C_2 \geq 0$ such that $|M_2(t) - C_2 e^{-\frac{1}{2}t}| \leq C e^{-\omega_0 t}$ with $\omega_0 > \frac{1}{2}$, but this is not sufficient to get a lower bound on M_2 , as we could have $C_2 = 0$.

To overcome this difficulty, instead of asking the fourth moment M_4 to be small initially, we will ask $\frac{M_{2k_0}}{M_2}$ to be small initially, for k_0 sufficiently large. This is for instance the case if we shrink the initial profile around \bar{x}_0 by a given parameter $\frac{1}{\varepsilon}$, which has the effect of scaling any initial moment M_{2k} by a factor ε^{2k} . Before proving that this condition is enough to get lower bounds on M_2 , we will need an adaptation of Lemma 3.1 in a nonlinear setting (the main difference being that we now require smallness of the initial condition $y(0)$ and of the parameter σ), which is stated in the following result.

Lemma 3.2. *We suppose that y is a nonnegative continuous function satisfying for all $t > 0$:*

$$\frac{d^+}{dt} y(t) \leq (-\lambda + v_1(y(t)))y(t) + v_0(t),$$

where $v_1(r) \leq Cr^\theta$ for all $r \geq 0$ sufficiently small (with $\theta > 0$ and $C > 0$), and $v_0(t)$ has a decay controlled by the parameter σ and with rate $[\omega_0]^-$.

Then, if $y(0)$ and σ are small enough, $y(t)$ has a decay controlled by $\max(y(0), \sigma)$ and with rate $[\min(\lambda, \omega_0)]^-$.

Proof. We fix $\tilde{\omega}_0 < \omega_0$, then $\tilde{\lambda} < \min(\lambda, \tilde{\omega}_0)$, and we pick $r > 0$ such that $\lambda - Cr^\theta > \tilde{\lambda}$ (and small enough so that $v_1(\tilde{r}) \leq C\tilde{r}^\theta$ as soon as $\tilde{r} \leq r$). Therefore, as soon as $y(t) \leq r$, we have

$$\frac{d^+}{dt}y(t) \leq -\tilde{\lambda}y(t) + \tilde{C}(\sigma)e^{-\tilde{\omega}_0 t},$$

where $\tilde{C}(\sigma) \rightarrow 0$ as $\sigma \rightarrow 0$. We suppose that $y(0) + \frac{\tilde{C}(\sigma)}{\tilde{\omega}_0 - \tilde{\lambda}} \leq r$.

We denote $T = \sup\{t > 0, \forall s \in [0, t], y(s) \leq r\}$. By continuity of y , since $y(0) < r$, we have $T > 0$ and therefore for all $t \in [0, T]$,

$$y(t) \leq \left(y(0) + \tilde{C}(\sigma) \frac{1 - e^{-(\tilde{\omega}_0 - \tilde{\lambda})t}}{\tilde{\omega}_0 - \tilde{\lambda}}\right) e^{-\tilde{\lambda}t} \leq re^{-\tilde{\lambda}t},$$

showing that $T = +\infty$. Indeed, if it was not the case, we would have $y(T) < r$, in contradiction with the definition of T , by continuity of y . Therefore we also obtain

$$v_1(y(t)) \leq C \left(y(0) + \tilde{C}(\sigma) \frac{1}{\tilde{\omega}_0 - \tilde{\lambda}}\right)^\theta e^{-\theta \tilde{\lambda}t},$$

showing that $v_1(y(t))$ has a rate controlled by $\max(y(0), \sigma)$ with rate $[\theta \tilde{\lambda}]^-$. We can therefore apply Lemma 3.2, as we now have a linear estimate (provided $y(0) + \frac{\tilde{C}(\sigma)}{\tilde{\omega}_0 - \tilde{\lambda}} \leq r$), and this ends the proof. \square

We are now ready to prove the following result, which leads to the statements of Theorem 1.3.

Proposition 3.4. *Under Assumption 3.1, if $\frac{M_{2k_0}^0}{M_2^0}$ is sufficiently small and if $\frac{1}{2^{2k_0-1}} < \delta$, then $\frac{M_{2k_0}}{M_2}$ has decay controlled by its initial value $\frac{M_{2k_0}^0}{M_2^0}$ and with rate $[\delta - \frac{1}{2^{2k_0-1}}]^-$.*

Furthermore, in that case, there exists $C_{k_0}(\frac{M_{2k_0}^0}{M_2^0})$, converging to 1 as $\frac{M_{2k_0}^0}{M_2^0} \rightarrow 0$ such that

$$M_2(t) \geq C_{k_0} \left(\frac{M_{2k_0}^0}{M_2^0}\right) M_2^0 e^{-\frac{t}{2}}.$$

Consequently, for all $k \geq 2$ (with $k \leq k_0$), $\frac{M_{2k}}{M_2}$ has decay controlled by $\frac{M_{2k_0}^0}{M_2^0}$ and with rate $[\min(\frac{1}{2} - \frac{1}{2^{2k-1}}, \delta - \frac{1}{2^{2k_0-1}})]^-$.

Proof. We compute

$$\frac{d^+}{dt} \left(\frac{M_{2k_0}}{M_2}\right) = \frac{1}{M_2} \frac{d^+ M_{2k_0}}{dt} - \frac{1}{M_2} \frac{dM_2}{dt} \left(\frac{M_{2k_0}}{M_2}\right).$$

Thanks to (9) and the estimates (40)–(46), we have

$$\frac{d}{dt} M_2 \geq \left(-\frac{1}{2} - \tilde{\alpha}\sqrt{M_2} - \beta M_2\right) M_2 - \tilde{\alpha}\sqrt{M_2 M_4} - \beta M_4, \quad (54)$$

and therefore, using (52), since $\frac{1}{2^{2k}} \sum_{\ell=2}^{2k-2} \binom{2k}{\ell} \leq 1$, we obtain

$$\frac{d^+}{dt} \left(\frac{M_{2k_0}}{M_2}\right) \leq \left(-\delta + \frac{1}{2^{2k_0-1}} + \tilde{w}(t)\right) \frac{M_{2k_0}}{M_2} + M_{2k_0-2}, \quad (55)$$

where

$$\tilde{w}(t) = \tilde{\alpha}((2k+2)\sqrt{M_2} + \sqrt{\frac{M_4}{M_2}}) + \beta(2k\sqrt{M_4} + 2M_2 + \frac{M_4}{M_2}).$$

Using (38), we get

$$M_2 \leq \sqrt{M_4} = \frac{M_4}{\sqrt{M_4}} \leq \frac{M_4}{M_2}, \quad (56)$$

and therefore we obtain

$$\tilde{w}(t) \leq (2k+3)\tilde{\alpha} \sqrt{\frac{M_4}{M_2}} + (2k+3)\beta \frac{M_4}{M_2}.$$

Now, using (47) with $i = 1$, $\ell = 4$ and $j = k_0$, we obtain

$$\frac{M_4}{M_2} \leq \frac{M_2^{\frac{2k_0-4}{2k_0-2}} M_2^{\frac{2}{2k_0-2}}}{M_2} = \left(\frac{M_{2k_0}}{M_2}\right)^{\frac{2}{2k_0-2}}, \quad (57)$$

and therefore we get

$$\tilde{w}(t) \leq (2k+3)\tilde{\alpha} \left(\frac{M_{2k_0}}{M_2}\right)^{\frac{1}{2k_0-2}} + (2k+3)\beta \left(\frac{M_{2k_0}}{M_2}\right)^{\frac{2}{2k_0-2}}. \quad (58)$$

Finally, thanks to Proposition 3.3, M_{2k_0-2} has decay controlled by the initial moment $M_{2k_0}^0$ and with rate $[\min(1 - \frac{1}{2^{2k_0-3}}, \frac{1}{2} - \frac{1}{2^{2k_0-1}} + \delta)]^-$, and since we have

$$M_{2k_0} \leq \left(\frac{M_{2k_0}}{M_2}\right) \cdot M_2 \leq \left(\frac{M_{2k_0}}{M_2}\right)^{1+\frac{2}{2k_0-2}},$$

we get that $M_{2k_0}^0 \rightarrow 0$ if $\frac{M_{2k_0}^0}{M_2^0} \rightarrow 0$, so M_{2k_0-2} has decay controlled by $\frac{M_{2k_0}^0}{M_2^0}$. Thanks to (55) and (58), we can apply Lemma 3.2 with $y(t) = \frac{M_{2k_0}}{M_2}$ and $\sigma = \frac{M_{2k_0}^0}{M_2^0} = y(0)$, the function w being given by $v_1(r) = (2k+3)(\tilde{\alpha} r^\theta + \beta r^{2\theta})$, where $\theta = \frac{1}{2k_0-2}$. We therefore obtain that if $\frac{M_{2k_0}^0}{M_2^0}$ is sufficiently small, then $\frac{M_{2k_0}}{M_2}$ has decay controlled by its initial value $\frac{M_{2k_0}^0}{M_2^0}$ and with rate $[\delta - \frac{1}{2^{2k_0-1}}]^-$.

Still denoting $\sigma = \frac{M_{2k_0}^0}{M_2^0}$, if we pick $\omega_0 < \delta - \frac{1}{2^{2k_0-1}}$, we then get that there exists $C(\sigma)$, converging to 0 as $\sigma \rightarrow 0$, such that we have

$$\frac{M_{2k_0}}{M_2} \leq C(\sigma) e^{-\omega_0 t}.$$

Therefore, using (54), (56) and (57), we obtain

$$\begin{aligned} \frac{d}{dt} M_2 &\geq \left(-\frac{1}{2} - 2\tilde{\alpha} \sqrt{\frac{M_4}{M_2}} - 2\beta \frac{M_4}{M_2}\right) M_2 \\ &\geq \left(-\frac{1}{2} - 2\tilde{\alpha} \left(\frac{M_{2k_0}}{M_2}\right)^\theta - 2\beta \left(\frac{M_{2k_0}}{M_2}\right)^{2\theta}\right) M_2 \\ &\geq \left(-\frac{1}{2} - \tilde{C}(\sigma) e^{-\theta \omega_0 t}\right) M_2, \end{aligned}$$

where

$$\tilde{C}(\sigma) = 2\tilde{\alpha} C(\sigma)^\theta + 2\beta C(\sigma)^{2\theta},$$

which still converges to 0 as $\sigma \rightarrow 0$. Solving this differential inequality for M_2 , we get

$$M_2(t) \geq e^{-\frac{t}{2} + \frac{\tilde{C}(\sigma)}{\omega_0 \theta} (e^{-\omega_0 \theta t} - 1)} M_2^0 \geq e^{-\frac{\tilde{C}(\sigma)}{\omega_0 \theta}} M_2^0 e^{-\frac{t}{2}},$$

which gives the expected lower bound on M_2 .

The last part of the proposition is a direct application of Proposition 3.3, using the fact that $M_{2k_0} \rightarrow 0$ as soon as $\frac{M_{2k_0}^0}{M_2^0} \rightarrow 0$. \square

4 Convergence to a self-similar profile

The goal of this section is to prove Theorem 1.4, which states the large time convergence of the centered and rescaled measure γ to a self-similar profile in the Fourier distance defined in (14), for some value of s .

We first discuss some prior results obtained on the following equation

$$\partial_t \widehat{\gamma}(t, \xi) = \widehat{\gamma}^2 \left(t, \frac{\xi}{2} \right) - \widehat{\gamma}(t, \xi) + \frac{1}{4} \xi \partial_\xi \widehat{\gamma}(t, \xi), \quad (59)$$

which is precisely equation (34) in the case $R \equiv 0$, that is when m is constant and then there is no selection. This equation on the Fourier transform of the rescaled distribution can also be derived from the kinetic model of Boltzmann-type studied in [32]. In this work, the authors proved a contraction property of the operator of equation (59) in the space of probability measures that satisfy (32) and have finite moments up to order $s = 2 + \delta$ with $\delta > 0$ well chosen, endowed with the Fourier distance d_s . Also, they proved the convergence in the Fourier distance of the rescaled measure towards the unique stationary solution of (59) with mass 1, centered, and second moment equal to 1, which is denoted by γ_∞ and defined in (16).

Here is the outline of the proof. First, we derive an estimate on the distance between γ and γ_∞ , under the assumption that the term $R(t, \xi)$ exponentially decreases in time and is also controlled by a $|\xi|^s$. Next, we prove that, for all $(t, \xi) \in \mathbb{R}_+ \times \mathbb{R}$, $R(t, \xi)$ satisfy theses conditions in the framework we set. We eventually combine these results to conclude to the statement of Theorem 1.4.

The first step of the proof is the following result.

Proposition 4.1. *We fix $s \in (2, 3)$ and we set $\lambda_s = 1 - \frac{s}{4} - 2^{1-s}$ (we have $\lambda_s > 0$ since $s \mapsto \lambda_s$ is strictly concave and $\lambda_2 = \lambda_3 = 0$). We suppose there exist $L > 0$ and $c > 0$ such that*

$$|R(t, \xi)| \leq |\xi|^s L e^{-ct}, \quad \forall \xi \in \mathbb{R}. \quad (60)$$

Then, for all time t ,

- *if $c = \lambda_s$, then we have*

$$d_s(\gamma, \gamma_\infty)(t) \leq d_s(\gamma_0, \gamma_\infty) e^{-\lambda_s t} + L t e^{-\lambda_s t},$$

- *otherwise, we have*

$$d_s(\gamma, \gamma_\infty)(t) \leq d_s(\gamma_0, \gamma_\infty) e^{-\lambda_s t} + L \frac{e^{-ct} - e^{-\lambda_s t}}{\lambda_s - c}.$$

Proof. From (34), we obtain

$$\frac{d}{dt} \left(e^t \widehat{\gamma}(t, e^{-\frac{t}{4}} \xi) \right) = e^t \widehat{\gamma}^2 \left(t, e^{-\frac{t}{4}} \frac{\xi}{2} \right) + e^t R \left(t, e^{-\frac{t}{4}} \xi \right),$$

and so we get the following Duhamel formulation

$$e^t \widehat{\gamma} \left(t, e^{-\frac{t}{4}} \xi \right) = \widehat{\gamma}_0 + \int_0^t e^\tau \widehat{\gamma}^2 \left(\tau, e^{-\frac{\tau}{4}} \frac{\xi}{2} \right) d\tau + \int_0^t e^\tau R \left(\tau, e^{-\frac{\tau}{4}} \xi \right) d\tau. \quad (61)$$

Since γ_∞ is a stationary solution to (34) with $R = 0$, applying then in this case (61) with $\widehat{\gamma}(t, \xi) = \widehat{\gamma}_\infty(\xi)$, we obtain

$$e^t \widehat{\gamma}_\infty \left(e^{-\frac{t}{4}} \xi \right) = \widehat{\gamma}_\infty + \int_0^t e^\tau \widehat{\gamma}_\infty^2 \left(e^{-\frac{\tau}{4}} \frac{\xi}{2} \right) d\tau. \quad (62)$$

First, notice that

$$\begin{aligned} |\widehat{\gamma}^2\left(\tau, e^{-\frac{\tau}{4}}\xi\right) - \widehat{\gamma_\infty}^2\left(e^{-\frac{\tau}{4}}\xi\right)| &\leq |\widehat{\gamma}\left(\tau, e^{-\frac{\tau}{4}}\xi\right) + \widehat{\gamma_\infty}\left(e^{-\frac{\tau}{4}}\xi\right)| |\widehat{\gamma}\left(\tau, e^{-\frac{\tau}{4}}\xi\right) - \widehat{\gamma_\infty}\left(e^{-\frac{\tau}{4}}\xi\right)| \\ &\leq 2|e^{-\frac{\tau}{4}}\xi|^s d_s(\gamma, \gamma_\infty)(\tau) = 2^{1-s} e^{-s\frac{\tau}{4}} |\xi|^s d_s(\gamma, \gamma_\infty)(\tau) \end{aligned}$$

since $\|\widehat{\gamma}\|_\infty \leq 1$ and $\|\widehat{\gamma_\infty}\|_\infty \leq 1$ (γ and γ_∞ are probability measures).

Besides, we have from (60)

$$|R(\tau, e^{-\frac{\tau}{4}}\xi)| \leq e^{-\frac{s\tau}{4}} |\xi|^s L e^{-c\tau}.$$

Therefore, we obtain the estimate on the last term in formula (61)

$$\int_0^t e^\tau R\left(\tau, e^{-\frac{\tau}{4}}\xi\right) d\tau \leq L|\xi|^s \int_0^t e^{(1-\frac{s}{4}-c)\tau} d\tau.$$

Now, by subtracting (62) from (61) and dividing by $|\xi|^s$, we obtain

$$e^t \frac{|\widehat{\gamma}\left(t, e^{-\frac{t}{4}}\xi\right) - \widehat{\gamma_\infty}\left(e^{-\frac{t}{4}}\xi\right)|}{|\xi|^s} \leq d_s(\gamma_0, \gamma_\infty) + 2^{1-s} \int_0^t e^{(1-\frac{s}{4})\tau} d_s(\gamma, \gamma_\infty)(\tau) d\tau + L \int_0^t e^{(1-\frac{s}{4}-c)\tau} d\tau.$$

Since we assumed that f^0 has finite moment up to order 4 and also γ_∞ up to order s , the distance $d_s(\gamma(t, \cdot), \gamma_\infty)$ is finite, then we can take the supremum on $\xi \neq 0$ in the estimate above to finally obtain

$$e^{(1-\frac{s}{4})t} d_s(\gamma, \gamma_\infty)(t) \leq d_s(\gamma_0, \gamma_\infty) + 2^{1-s} \int_0^t e^{(1-\frac{s}{4})\tau} d_s(\gamma, \gamma_\infty)(\tau) d\tau + L \int_0^t e^{(1-\frac{s}{4}-c)\tau} d\tau. \quad (63)$$

This Grönwall estimate can be solved classically by setting

$$z(t) = e^{-2^{1-s}t} \left(d_s(\gamma_0, \gamma_\infty) + 2^{1-s} \int_0^t e^{(1-\frac{s}{4})\tau} d_s(\gamma, \gamma_\infty)(\tau) d\tau \right),$$

thus we obtain

$$z'(t) = 2^{1-s} e^{-2^{1-s}t} e^{(1-\frac{s}{4})t} d_s(\gamma, \gamma_\infty)(t) - 2^{1-s} z(t) \leq 2^{1-s} L e^{-2^{1-s}t} \int_0^t e^{(1-\frac{s}{4}-c)\tau} d\tau,$$

which gives, since $z(0) = d_s(\gamma_0, \gamma_\infty)$,

$$z(t) \leq d_s(\gamma_0, \gamma_\infty) + L \left(-e^{-2^{1-s}t} \int_0^t e^{(1-\frac{s}{4}-c)\tau} d\tau + \int_0^t e^{(\lambda_s-c)\tau} d\tau \right). \quad (64)$$

Getting back to (63), we obtain

$$d_s(\gamma, \gamma_\infty)(t) \leq e^{-(1-\frac{s}{4})t} \left(e^{2^{1-s}t} z(t) + L \int_0^t e^{(1-\frac{s}{4}-c)\tau} d\tau \right)$$

and using (64), we get

$$\begin{aligned} d_s(\gamma, \gamma_\infty)(t) &\leq d_s(\gamma_0, \gamma_\infty) e^{-\lambda_s t} + L \left(e^{-\lambda_s t} \int_0^t e^{(\lambda_s-c)\tau} d\tau - e^{-(1-\frac{s}{4})t} \int_0^t e^{(1-\frac{s}{4}-c)\tau} d\tau \right) \\ &\quad + L e^{-(1-\frac{s}{4})t} \int_0^t e^{(1-\frac{s}{4}-c)\tau} d\tau. \end{aligned}$$

Hence we have

$$d_s(\gamma, \gamma_\infty)(t) \leq d_s(\gamma_0, \gamma_\infty) e^{-\lambda_s t} + L e^{-\lambda_s t} \int_0^t e^{(\lambda_s-c)\tau} d\tau,$$

which gives the result stated in Proposition 4.1. \square

Next, we establish some estimates on $R(t, \xi)$ and the moments of g that will help to derive an estimate of the type (60). From now on, we use the notation

$$\widehat{\phi}'(t, \xi) := \frac{\partial \widehat{\phi}}{\partial \xi}(t, \xi), \quad \widehat{\phi}''(t, \xi) := \frac{\partial^2 \widehat{\phi}}{\partial \xi^2}(t, \xi), \quad \widehat{\phi}'''(t, \xi) := \frac{\partial^3 \widehat{\phi}}{\partial \xi^3}(t, \xi),$$

for the sake of simplicity. We will also omit the dependency on time in the proofs of the two next stated results, since all the estimates are derived at a fixed time t . We will resume the dependency on time when we prove Theorem 1.4.

The lemma below provides some useful estimates on R .

Lemma 4.1. *We fix $t \geq 0$. For all $\xi \in \mathbb{R}$, we have the two following estimates:*

$$|R(t, \xi)| \leq \left(\frac{1}{2} \|\widehat{\varphi}(t, \cdot)\|_\infty + \|\widehat{\varphi}'(t, \cdot)\|_\infty + \|\widehat{\varphi}''(t, \cdot)\|_\infty \right) |\xi|^2, \quad (65)$$

$$|R(t, \xi)| \leq \left(\frac{1}{12} (\|\widehat{\varphi}(t, \cdot)\|_\infty + 3 \|\widehat{\varphi}''(t, \cdot)\|_\infty) \|\widehat{\gamma}'''(t, \cdot)\|_\infty + \frac{1}{2} \|\widehat{\varphi}'(t, \cdot)\|_\infty + \frac{1}{6} \|\widehat{\varphi}'''(t, \cdot)\|_\infty \right) |\xi|^3. \quad (66)$$

Proof. From the definition (35) of R , we obtain

$$R(t, \xi) = S_0(t) \widehat{\gamma}(t, \xi) - \widehat{\varphi}(t, \xi) - \frac{1}{2} \left(S_0(t) - \frac{S_2(t)}{M_2(t)} \right) \xi \widehat{\gamma}'(t, \xi) - i \frac{S_1(t)}{\sqrt{M_2(t)}} \xi \widehat{\gamma}(t, \xi). \quad (67)$$

Since $\gamma(t, \cdot)$ is a centered probability density with second moment 1, we get

$$\widehat{\gamma}(t, 0) = 1, \quad \widehat{\gamma}'(t, 0) = 0, \quad \widehat{\gamma}''(t, 0) = -1.$$

Furthermore, for all ξ , we have

$$|\widehat{\gamma}(t, \xi)| \leq \int_{\mathbb{R}} \gamma(x) dx = 1, \quad |\widehat{\gamma}''(t, \xi)| = |-\widehat{x^2 \gamma}(t, \xi)| \leq \int_{\mathbb{R}} |x|^2 \gamma(t, x) dx = 1. \quad (68)$$

We also have, by Cauchy-Schwarz inequality,

$$|\widehat{\gamma}'(t, \xi)| = | -i \widehat{x \gamma}(t, \xi) | \leq \int_{\mathbb{R}} |x| \gamma(t, x) dx \leq 1. \quad (69)$$

Now, a simple change of variable reads, for $k \in \mathbb{N}$,

$$\frac{S_k}{\sqrt{M_2}^k} = \int_{\mathbb{R}} x^k \varphi(x) dx = i^k \widehat{\varphi}^{(k)}(0),$$

(where the dependency on time t is omitted) and therefore, the expression (67) of $R(t, \xi)$ can be written under the form

$$R(t, \xi) = \widehat{\varphi}(0) \widehat{\gamma}(\xi) - \widehat{\varphi}(\xi) + \left(\widehat{\varphi}'(0) \widehat{\gamma}(\xi) - \frac{1}{2} (\widehat{\varphi}(0) + \widehat{\varphi}''(0)) \widehat{\gamma}'(\xi) \right) \xi.$$

From this, we easily remark that $R(0) = 0$ and we obtain

$$\begin{aligned} R'(t, \xi) &= \widehat{\varphi}'(0) \widehat{\gamma}(\xi) - \widehat{\varphi}'(\xi) + \frac{1}{2} (\widehat{\varphi}(0) - \widehat{\varphi}''(0)) \widehat{\gamma}'(\xi) \\ &\quad + \left(\widehat{\varphi}'(0) \widehat{\gamma}'(\xi) - \frac{1}{2} (\widehat{\varphi}(0) + \widehat{\varphi}''(0)) \widehat{\gamma}''(\xi) \right) \xi. \end{aligned} \quad (70)$$

Also, we remark once again that $R'(t, 0) = 0$ (since $\widehat{\gamma}'(0) = 0$), and we also have from (70)

$$\begin{aligned} R'(t, \xi) &= \int_0^\xi \left(\widehat{\varphi}'(0) \widehat{\gamma}'(\eta) - \widehat{\varphi}''(\eta) + \frac{1}{2} (\widehat{\varphi}(0) - \widehat{\varphi}''(0)) \widehat{\gamma}''(\eta) \right) d\eta \\ &\quad + \left(\widehat{\varphi}'(0) \widehat{\gamma}'(\xi) - \frac{1}{2} (\widehat{\varphi}(0) + \widehat{\varphi}''(0)) \widehat{\gamma}''(\xi) \right) \xi. \end{aligned}$$

Using (68)-(69), we obtain a first estimate on R'

$$|R'(t, \xi)| \leq (\|\widehat{\varphi}\|_\infty + 2\|\widehat{\varphi}'\|_\infty + 2\|\widehat{\varphi}''\|_\infty)|\xi|,$$

and since $R(t, 0) = 0$, by integration we get the estimation (65).

Similarly, we compute the second derivative of R and we obtain

$$\begin{aligned} R''(t, \xi) &= 2\widehat{\varphi}'(0)\widehat{\gamma}'(\xi) - \widehat{\varphi}''(\xi) - \widehat{\varphi}''(0)\widehat{\gamma}''(\xi) \\ &\quad + \left(\widehat{\varphi}'(0)\widehat{\gamma}''(\xi) - \frac{1}{2}(\widehat{\varphi}(0) + \widehat{\varphi}''(0))\widehat{\gamma}'''(\xi) \right) \xi. \end{aligned}$$

From this, we remark once more that $R''(t, 0) = 0$ (since $\widehat{\gamma}''(0) = -1$), and we also have

$$\begin{aligned} R''(t, \xi) &= \int_0^\xi (2\widehat{\varphi}'(0)\widehat{\gamma}''(\eta) - \widehat{\varphi}'''(\eta) - \widehat{\varphi}''(0)\widehat{\gamma}'''(\eta)) d\eta \\ &\quad + \left(\widehat{\varphi}'(0)\widehat{\gamma}''(\xi) - \frac{1}{2}(\widehat{\varphi}(0) + \widehat{\varphi}''(0))\widehat{\gamma}'''(\xi) \right) \xi. \end{aligned}$$

Using again (68)-(69), we obtain

$$|R''(t, \xi)| \leq \left(\frac{1}{2}(\|\widehat{\varphi}\|_\infty + 3\|\widehat{\varphi}''\|_\infty) \|\widehat{\gamma}'''\|_\infty + 3\|\widehat{\varphi}'\|_\infty + \|\widehat{\varphi}'''\|_\infty \right) |\xi|,$$

and since $R'(t, 0) = R(t, 0) = 0$, by two successive integrations we get the estimation (66) and this ends the proof. \square

Next, we prove estimates on the derivatives of φ with the moments M_k , using the results of Section 3. Once more, we fix $t \geq 0$ and omit the dependency on t in the notations, the norm $\|\cdot\|_\infty$ being the supremum with respect to the Fourier variable ξ only.

Proposition 4.2. *We have*

$$\|\widehat{\gamma}'''\|_\infty \leq \frac{\sqrt{M_4}}{M_2}, \quad (71)$$

and if m satisfies Assumption 1.2, we have

$$\begin{aligned} \|\widehat{\varphi}^{(k)}\|_\infty &\leq \alpha(\overline{x}) \sqrt{\frac{M_4}{M_2}} + \beta \frac{M_4}{M_2}, & \text{for } 0 \leq k \leq 2, \\ \|\widehat{\varphi}^{(k)}\|_\infty &\leq \alpha(\overline{x}) \frac{M_4}{M_2^{\frac{3}{2}}} + \beta \frac{\sqrt{M_4 M_6}}{M_2^{\frac{3}{2}}}, & \text{for } 0 \leq k \leq 3. \end{aligned}$$

Proof. By a simple change of variable, we have

$$\frac{M_k}{\sqrt{M_2}^k} = \int_{\mathbb{R}} x^k \gamma(x) dx,$$

and in particular, we get (71) by

$$|\widehat{\gamma}'''(\xi)| \leq \int_{\mathbb{R}} |x|^3 \gamma(x) dx \leq \sqrt{\int_{\mathbb{R}} x^2 \gamma(x) dx} \sqrt{\int_{\mathbb{R}} x^4 \gamma(x) dx} = \sqrt{1} \sqrt{\frac{M_4}{\sqrt{M_2}^4}} = \frac{\sqrt{M_4}}{M_2}.$$

Furthermore, the quantity $\int_{\mathbb{R}} |x|^k \gamma(x) dx$ is nondecreasing with k , from $k \geq 1$. Indeed, by induction, thanks to (69) we first have $\int_{\mathbb{R}} |x| \gamma(x) dx \leq 1 = \int_{\mathbb{R}} |x|^2 \gamma(x) dx$. And if for some $k \geq 1$, we have $\int_{\mathbb{R}} |x|^k \gamma(x) dx \leq \int_{\mathbb{R}} |x|^{k+1} \gamma(x) dx$ then by Cauchy-Schwarz inequality, we get:

$$\begin{aligned} \int_{\mathbb{R}} |x|^{k+1} \gamma(x) dx &\leq \sqrt{\int_{\mathbb{R}} |x|^k \gamma(x) dx} \sqrt{\int_{\mathbb{R}} |x|^{k+2} \gamma(x) dx} \\ &\leq \sqrt{\int_{\mathbb{R}} |x|^{k+1} \gamma(x) dx} \sqrt{\int_{\mathbb{R}} |x|^{k+2} \gamma(x) dx}, \end{aligned}$$

which proves that $\int_{\mathbb{R}} |x|^{k+1} \gamma(x) dx \leq \int_{\mathbb{R}} |x|^{k+2} \gamma(x) dx$.

If now m satisfies Assumption 1.2, we obtain by the definition (33) that for all $x \in \mathbb{R}$, we have

$$|\varphi| \leq [\alpha(\bar{x})\sqrt{M_2}|x| + \beta M_2|x|^2]\gamma(x).$$

Therefore we obtain

$$\begin{aligned} |\widehat{\varphi}^{(k)}(\xi)| &\leq \int_{\mathbb{R}} |x|^k |\varphi(x)| dx \\ &\leq \alpha(\bar{x})\sqrt{M_2} \int_{\mathbb{R}} |x|^{k+1} \gamma(x) dx + \beta M_2 \int_{\mathbb{R}} |x|^{k+2} \gamma(x) dx, \end{aligned}$$

which the gives

$$\begin{aligned} \|\widehat{\varphi}^{(k)}\|_{\infty} &\leq \alpha(\bar{x})\sqrt{M_2} \int_{\mathbb{R}} |x|^3 \gamma(x) dx + \beta M_2 \int_{\mathbb{R}} |x|^4 \gamma(x) dx && \text{for } 0 \leq k \leq 2, \\ \|\widehat{\varphi}^{(k)}\|_{\infty} &\leq \alpha(\bar{x})\sqrt{M_2} \int_{\mathbb{R}} |x|^4 \gamma(x) dx + \beta M_2 \int_{\mathbb{R}} |x|^5 \gamma(x) dx && \text{for } 0 \leq k \leq 3. \end{aligned}$$

We have $\int_{\mathbb{R}} |x|^3 \gamma(x) dx \leq \frac{\sqrt{M_4}}{M_2}$ thanks to (71), $\int_{\mathbb{R}} |x|^4 \gamma(x) dx = \frac{M_4}{M_2^2}$, and finally, by Cauchy-Schwarz inequality,

$$\int_{\mathbb{R}} |x|^5 \gamma(x) dx \leq \sqrt{\int_{\mathbb{R}} |x|^4 \gamma(x) dx \int_{\mathbb{R}} |x|^6 \gamma(x) dx} \leq \sqrt{\frac{M_4 M_6}{M_2^2 M_2^3}},$$

and we therefore obtain the desired estimates. \square

We can conclude with the proof of Theorem 1.4.

Proof. If m satisfies Assumption 1.2, then combining the results of Lemma 4.1 and Proposition 4.2 gives the following estimates for $|R(t, \xi)|$:

$$|R(t, \xi)| \leq \frac{5}{2} \left(\alpha(\bar{x})\sqrt{\frac{M_4}{M_2}} + \beta \frac{M_4}{M_2} \right) |\xi|^2,$$

and

$$\begin{aligned} |R(t, \xi)| &\leq \left(\alpha(\bar{x})\frac{M_4}{M_2^{\frac{3}{2}}} + \beta \left(\frac{1}{3} \frac{M_4^{\frac{3}{2}}}{M_2^2} + \frac{2}{3} \frac{\sqrt{M_4 M_6}}{M_2^{\frac{3}{2}}} \right) \right) |\xi|^3 \\ &\leq \left(\alpha(\bar{x})\frac{M_4}{M_2^{\frac{3}{2}}} + \beta \frac{\sqrt{M_4 M_6}}{M_2^{\frac{3}{2}}} \right) |\xi|^3, \end{aligned}$$

with the last inequality coming from the fact that $M_4 \leq \sqrt{M_2 M_6}$ by Cauchy-Schwarz inequality.

Then, we fix δ from Assumption 3.1. By Proposition 3.4, assuming that $\frac{M_{2k_0}^0}{M_2^0}$ is small enough for some $k_0 \geq 3$ (and such that $\frac{1}{2^{2k_0-1}} < \delta$), we have that, for each $\omega < \min(\frac{3}{8}, \delta - \frac{1}{2^{2k_0-1}})$, there exist constants C_1 and C_2 such that

$$|R(t, \xi)| \leq C_1 e^{-\frac{\omega}{2}t} |\xi|^2,$$

and

$$|R(t, \xi)| \leq C_2 e^{(\frac{1}{4}-\omega)t} |\xi|^3.$$

Thus, taking $s \in (2, 3)$ and defining $c_s := (3 - s)\frac{\omega}{2} - (s - 2)(\frac{1}{4} - \omega)$, we obtain

$$|R(t, \xi)| = |R(t, \xi)|^{(3-s)+(s-2)} \leq C_1^{3-s} C_2^{s-2} e^{-c_s t} |\xi|^s,$$

with $c_s > 0$ when $s < \bar{s} := 2 + \frac{\omega}{\frac{1}{2} - \omega}$.

We can distinguish two situations depending on the value of ω .

- If $\omega \geq \frac{1}{4}$, then we have $c_s > 0$ for all $s \in (2, 3)$. Moreover, we have that $c_s > \lambda_s$ for each value of s , which implies from Proposition 4.1

$$d_s(\gamma, \gamma_\infty)(t) \leq d_s(\gamma_0, \gamma_\infty) e^{-\lambda_s t} + L \frac{e^{-\lambda_s t}}{c_s - \lambda_s}.$$

Indeed, in this case, for $s = 3$ and $\omega = \frac{1}{4}$, we have $\lambda_3 - c_3 = 0$, and for $s \in (2, 3)$ the function $s \mapsto \lambda_s - c_s$ is increasing (it is strictly concave and its derivative at $s = 3$ is $\frac{1}{4}(\ln 2 - \frac{1}{2}) > 0$). Therefore for $\omega = \frac{1}{4}$ we have $\lambda_s < c_s$ for all $s \in (2, 3)$. Since c_s is increasing with respect to ω (and λ_s does not depend on ω), this provides the same result for $\omega \geq \frac{1}{4}$.

- If $\omega < \frac{1}{4}$, s must be in $(2, \bar{s})$ to get c_s positive. Furthermore the function $s \mapsto \lambda_s - c_s$ is still increasing on $(2, 3)$, since it is strictly concave and with positive derivative at $s = 3$. Since $\lambda_3 - c_3 > 0$ and $\lambda_2 - c_2 < 0$, we get that there is a unique value s_0 for which $\lambda_{s_0} = c_{s_0}$. Then, Proposition 4.1 can be applied. We obtain

$$d_{s_0}(\gamma, \gamma_\infty)(t) \leq d_{s_0}(\gamma_0, \gamma_\infty) e^{-\lambda_{s_0} t} + L t e^{-\lambda_{s_0} t}.$$

And for $s \neq s_0$, we have

$$d_s(\gamma, \gamma_\infty)(t) \leq d_s(\gamma_0, \gamma_\infty) e^{-\lambda_s t} + L' e^{-\min(\lambda_s, c) t},$$

with $L' = \frac{L}{|\lambda_s - c|}$.

Hence the statement of Theorem 1.4, together with the precisions given in Remark 1.1. \square

Acknowledgements

The authors are thankful to Florian Patout for the fruitful discussions that helped to achieve this work. A.F. wants to thank the hospitality of the LMA (Université de Poitiers) where most of this research was conducted.

References

- [1] M. Alfaro and R. Carles. Explicit solutions for replicator-mutator equations: extinction versus acceleration. *SIAM J. Appl. Math.*, 74(6):1919–1934, 2014.
- [2] M. Alfaro and R. Carles. Replicator-mutator equations with quadratic fitness. *Proc. Am. Math. Soc.*, 145(12):5315–5327, 2017.
- [3] A. Baldassarri, U. M. B. Marconi, and A. Puglisi. Kinetics Models of Inelastic Gases. *Math. Models Methods Appl. Sci.*, 12(07):965–983, 2002.
- [4] N. Barton, A. Etheridge, and A. Véber. The infinitesimal model: Definition, derivation, and implications. *Theor. Popul. Biol.*, 118:50–73, 2017.

- [5] M. Bisi, J. A. Carrillo, and G. Toscani. Contractive metrics for a Boltzmann equation for granular gases: Diffusive equilibria. *J. Stat. Phys.*, 118(1):301–331, 2005.
- [6] A. V. Bobylev. Fourier transform method in the theory of the Boltzmann equation for Maxwellian molecules. *Sov. Phys., Dokl.*, 20(1975)(6):820–822, 1976. translation from *Dokl. Akad. Nauk SSSR* 225:1041–1044, 1975.
- [7] A. V. Bobylev. The theory of nonlinear spatially uniform Boltzmann equation for Maxwell molecules. *Soviet Sci. Rev., Sect. C, Math. Phys. Rev.*, 7:111–233, 1988.
- [8] M. G. Bulmer. *The mathematical theory of quantitative genetics*. The Clarendon Press, Oxford University Press, New York, 1980.
- [9] V. Calvez, J. Garnier, and F. Patout. Asymptotic analysis of a quantitative genetics model with nonlinear integral operator. *J. Éc. Polytech., Math.*, 6:537–579, 2019.
- [10] V. Calvez, T. Lepoutre, and D. Poyato. Ergodicity of the Fisher infinitesimal model with quadratic selection. ArXiv e-prints, 2021.
- [11] V. Calvez, D. Poyato, and F. Santambrogio. Uniform contractivity of the Fisher infinitesimal model with strongly convex selection. ArXiv e-prints, 2023.
- [12] J. A. Carrillo and G. Toscani. Contractive probability metrics and asymptotic behavior of dissipative kinetic equations. *Riv. Mat. Univ. Parma (7)*, 6:75–198, 2007.
- [13] B. Cloez and P. Gabriel. Fast, slow convergence, and concentration in the house of cards replicator-mutator model. ArXiv e-prints, 2022.
- [14] D. L. Cohn. *Measure theory*. Birkhäuser Adv. Texts, Basler Lehrbüch. Birkhäuser/Springer, New York, NY, 2nd revised edition, 2013.
- [15] P. Collet, S. Méléard, and J. A. J. Metz. A rigorous model study of the adaptive dynamics of Mendelian diploids. *J. Math. Biol.*, 67(3):569–607, 2013.
- [16] C. Coron, M. Costa, H. Leman, and C. Smadi. A stochastic model for speciation by mating preferences. *J. Math. Biol.*, 76(6):1421–1463, May 2018.
- [17] P. Degond, A. Frouvelle, and G. Raoul. Local stability of perfect alignment for a spatially homogeneous kinetic model. *J. Stat. Phys.*, 157(1):84–112, 2014.
- [18] L. Dekens. Evolutionary dynamics of complex traits in sexual populations in a heterogeneous environment: how normal? *J. Math. Biol.*, 84(3):15, 2022.
- [19] L. Dekens and S. Mirrahimi. Dynamics of Dirac concentrations in the evolution of quantitative alleles with sexual reproduction. ArXiv e-prints, 2021.
- [20] O. Diekmann, P.-E. Jabin, S. Mischler, and B. Perthame. The dynamics of adaptation: an illuminating example and a Hamilton-Jacobi approach. *Theor. Popul. Biol.*, 67(4):257–271, 2005.
- [21] B. Düring, D. Matthes, and G. Toscani. A Boltzmann-type approach to the formation of wealth distribution curves. *Riv. Mat. Univ. Parma (8)*, 1:199–261, 2009.
- [22] R. A. Fisher. The correlation between relatives on the supposition of Mendelian inheritance. *Trans. R. Soc. Edinburgh*, 52(2):399–433, 1919.

- [23] N. Fournier and B. Perthame. Transport distances for PDEs: the coupling method. *EMS Surv. Math. Sci.*, 7(1):1–31, 2020.
- [24] N. Fournier and B. Perthame. A nonexpanding transport distance for some structured equations. *SIAM J. Math. Anal.*, 53(6):6847–6872, 2021.
- [25] E. Gabetta, G. Toscani, and B. Wennberg. Metrics for probability distributions and the trend to equilibrium for solutions of the Boltzmann equation. *J. Stat. Phys.*, 81(5):901–934, 1995.
- [26] M.-E. Gil, F. Hamel, G. Martin, and L. Roques. Mathematical properties of a class of integro-differential models from population genetics. *SIAM J. Appl. Math.*, 77(4):1536–1561, 2017.
- [27] M. Kirkpatrick and N. H. Barton. Evolution of a Species’ Range. *The American Naturalist*, 150(1):1–23, 1997.
- [28] H. Leman. A stochastic model for reproductive isolation under asymmetrical mating preferences. *Bull. of Math. Biol.*, 80(9):2502–2525, 2018.
- [29] P. Magal and G. Raoul. Dynamics of a kinetic model describing protein exchanges in a cell population. ArXiv e-prints, 2015.
- [30] D. Matthes and G. Toscani. On steady distributions of kinetic models of conservative economies. *J. Stat. Phys.*, 130(6):1087–1117, 2008.
- [31] S. Mirrahimi and G. Raoul. Dynamics of sexual populations structured by a space variable and a phenotypical trait. *Theor. Popul. Biol.*, 84:87–103, 2013.
- [32] L. Pareschi and G. Toscani. Self-similarity and power-like tails in nonconservative kinetic models. *J. Stat. Phys.*, 124(2-4):747–779, 2006.
- [33] F. Patout. The Cauchy problem for the infinitesimal model in the regime of small variance. ArXiv e-prints, 2020.
- [34] B. Perthame, M. Strugarek, and C. Taing. Selection–mutation dynamics with asymmetrical reproduction kernels. *Nonlinear Anal.*, 222:112947, 2022.
- [35] A. Pulvirenti. *Alcuni risultati sull’equazione di Boltzmann*. PhD thesis, Pavia, 1994.
- [36] A. Pulvirenti and G. Toscani. Asymptotic properties of the inelastic Kac model. *J. Stat. Phys.*, 114:1453–1480, 01 2004.
- [37] G. Raoul. Macroscopic limit from a structured population model to the Kirkpatrick-Barton model. ArXiv e-prints, 2017.
- [38] G. Raoul. Exponential convergence to a steady-state for a population genetics model with sexual reproduction and selection. ArXiv e-prints, 2021.