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# Er<sup>3+</sup> Doped Nanoparticles as Upconversion Thermometer Probe in confined fluids

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**Abstract:** Non-Contact temperature measurement at the nanoscale by photoluminescence using a nano-sensor in a confined fluid has been developed in the present work. Upconversion lanthanide-doped nanoparticles applied to ratiometric thermometry could be considered as a self-referenced nanosensor. Gadolinium orthovanadate (GdVO<sub>4</sub>) nanoparticles, doped with Yb<sup>3+</sup> and Er<sup>3+</sup> were synthesized and then dispersed in an ester-based fluid. Rheological measurements show that the viscosity of the dispersed NPs suspension remains unchanged up to the shear rate of 10<sup>-4</sup> s<sup>-1</sup> at 393 K. The NPs suspension allows luminescence intensity ratio (LIR) thermometry up to 473 K with a relative sensitivity of 1.17%. K<sup>-1</sup> with the NIR laser. Then, the temperature calibration by coupling the high pressure (1.08 GPa max) confirmed the applicability of NPs as a thermosensor in a variable pressure environment. According to these results, the fluid containing GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> nanoparticles can be used for temperature sensing in a pressurized environment for further application in tribology.

Keywords: optical thermometry; tribology, contactless thermo-sensor

## 1. Introduction

Temperature and pressure are fundamental parameters affecting the physicochemical properties of materials. Their measurements are crucial for both scientific and industrial applications. In many applications, they have to be measured locally in submicron-sized areas: nano-manometry and manothermometry in biomedical applications, microfluidics, catalysis ...as well as in tribology. In tribological measurements, temperature and pressure are crucial parameters for understanding and predicting the behavior of the interface between solids in contact. Local temperature and pressure changes between two contacting friction surfaces affect the rheological properties of the interface, which has an important influence on the life and performance of the tribological system.<sup>1,2</sup> To understand the friction between two materials in contact, local temperature and pressure measurement

with high accuracy and efficiency has become a major topic in tribology research. Well-known techniques for measuring temperature in friction include (1) thermocouple inserted in the material 2 mm below the sliding surface; (2) infrared thermography; (3) Raman micro-spectroscopy.<sup>3,4</sup> However, reliable measurement with good spatial and temporal resolution of temperature, or other physical parameters in frictional contact is not a simple matter so far.

The application of Raman spectroscopy to tribological measurements has recently been studied for local pressure mapping in the case of lubricated contact with a fluid by Jubault *et al.*<sup>5,6</sup> and Yagi *et al.*<sup>7</sup> In such experiments, this technique provides a precise and accurate pressure measurement with a higher spatial and temporal resolution than other temperature measurement techniques. However, it requires the selection of a fluid that exhibits strong Raman scattering, which is not the case for most lubricants. Infrared thermography collects the IR emission from the tribological contact. It is then possible to measure the temperature all over the contact area by using a full-field IR camera. A drawback of this technique is the calibration process. Indeed, different sources (solid surfaces and bodies, lubricant film or the interface and the environment) contribute to the overall IR radiation and must be separated by applying appropriate coatings to the transparent disc and/or using different filters. These issues can be solved using non-contact nano-manometers and nanothermometers. Recent studies have been carried out by Albahrani *et al.*<sup>8</sup> and Seoudi *et al.*<sup>9</sup> on the application of a nano-sensor for measuring temperature and pressure in a lubricating fluid. In their studies, CdSe/CdS/ZnS core/shell/shell quantum dots (QDs) were used as a calibrated nano-sensor for temperature and pressure (under 0.2 to 1.3 GPa between 0 and 100 °C) by their change in emission energy shift. The main drawback of this technique was the coupled dependence of the emission energy to pressure and temperature, while pressure and temperature generally vary independently in tribological contact.

It is therefore essential to find out the suitable optical thermometric sensor for temperature measurement under multivariable conditions. In tribological measurement, pressure independence should be required. Meanwhile, to reduce the measurement scale to achieve high spatial resolution, the miniaturized size of the sensor is one of the most important conditions. Over the past decades, many photoluminescent nanothermometers have been studied, such as fluorescent dyes,<sup>10-12</sup> CdSe-based<sup>9,13</sup> and perovskite based<sup>14,15</sup> quantum dots. Temperature sensitivity could be demonstrated by their spectroscopic properties, such as emission spectral shift, variation in emission intensity, change in emitter lifetime... Among the spectral variables of thermometry, the luminescent intensity ratio (LIR) is well known as a self-referenced measurement method, which will not be influenced by the absolute intensity of the emission.<sup>16-25</sup> This feature is particularly needed in dynamic measurement, where the sensor concentration cannot remain constant. Nanoparticles doped with lanthanide ions ( $\text{Ln}^{3+}$ )<sup>26,27</sup> are one of the main families of photoluminescent nanothermometers, and are commonly used as LIR thermometers. Thanks to the special property of  $\text{Er}^{3+}$ , it provides thermally coupled energy levels ( $^2\text{H}_{11/2}$  and  $^4\text{S}_{3/2}$ ) that are suitable for temperature measurement in the range 293-473 K, with high sensitivity.<sup>28-32</sup> The relatively small gap between the energy levels (about 600–800  $\text{cm}^{-1}$ ) allows excited electrons to be thermally promoted to the higher level. The proportion of thermally promoted electrons follows the Boltzmann distribution,<sup>33</sup> which allows the temperature to be calibrated by measuring the ratio of the emission intensity of its two thermally coupled energy levels. Besides the forementioned favorable luminescent properties, they exhibit also relatively low cytotoxicity, resistance to photobleaching, thermal decomposition and oxidation processes. Therefore, they can be used in various applications including biomedical ones.

However, considering the scattered medium for the  $\text{Er}^{3+}$  based fluorescent sensor in temperature measurement, the current excitation by laser of  $\text{Er}^{3+}$  under green light can cause interfering emission from the scattered medium. To solve this problem, the  $\text{Er}^{3+}$  based fluorescent sensor could be excited via an upconversion (UC) system, with low-energy absorption to induce high-energy emission.<sup>34,35</sup>  $\text{Er}^{3+}$  combined with a sensitizing ion such as  $\text{Yb}^{3+}$ , can be excited by a near-infrared (NIR) laser via the energy transferred by  $\text{Yb}^{3+}$ . The use of NIR excitation source avoids the possible emission of scattered media during the visible excitation.<sup>36,37</sup> In the UC system, several experimental parameters need to be carefully controlled to obtain accurate  $\text{Er}^{3+}$  thermometry, such as the effects induced by the NIR laser,<sup>29,38-41</sup> the optical quenching effect<sup>42-45</sup> and the chemical properties of the sensor.<sup>46-49</sup> These parameters can be affected by other physical influencing factors. As an example, in tribological

applications, local pressure can cause potential damage or material degradation. Under high pressure, the nano-sensor could be modified and thus its optical properties could change.<sup>50,51</sup> Studies in the literature have shown that high pressure (up to 25 GPa) can cause variations in Er<sup>3+</sup>-doped materials due to structural change such as shortening of interionic distances, shifting energy levels. These variations modify the probability of transitions and therefore have an impact on the position and intensity of the photoluminescence peak.<sup>52-58</sup> Since the inorganic structure is stable under compression, a relative high pressure is usually required to cause a structural variation. Under low-range pressure (below 6 GPa), the pressure sensitivity on SrF<sub>2</sub>: Yb<sup>3+</sup>, Er<sup>3+</sup> has been reported in Runowski's works in 2017, variations in lifetime, emission intensity and peak centroid were observed under compression up to 5.5 GPa.<sup>57</sup> Moreover, the pressure-sensing properties of upconversion Yb<sup>3+</sup>/Er<sup>3+</sup> doped fluoride nanoparticles have also been found to vary the red-to-green ratio, lifetime and brightness of Er<sup>3+</sup> with pressure up to 5.4 GPa by McLellan *et al.*<sup>58</sup>

Nonetheless, pressure sensitivity up to 1 GPa has not yet been reported by Ln-doped nanoparticles. As mentioned previously, for further application in tribology, the temperature measurement must be stable under specific conditions, such as high pressure (up to 3 GPa) and high shear rates (up to 10<sup>6</sup> s<sup>-1</sup>). Indeed, besides temperature, pressure and shear stress are the other two key parameters that vary in tribological contact. In this work, we studied the effect of a more moderated pressure (up to 1 GPa) on the LIR thermometry of Er<sup>3+</sup>-doped nanoparticles. As the nano-sensor will be used in a specific medium for further tribological measurements, the GdVO<sub>4</sub>:Yb<sup>3+</sup>/Er<sup>3+</sup> upconversion nano-sensor was synthesized and then dispersed in a lubricant (ester-based oil). Among the various possible matrices, it offers a number of advantages, including very simple synthesis in the form of very small particles (< 50 nm). The hydrophilic surface of the nanoparticles after co-precipitation synthesis means that dispersion of the nanoparticles in the lubricant can be optimised for future tribological applications. Moreover, the vanadate matrix provides very good sensitivity for Er<sup>3+</sup>/Yb<sup>3+</sup> thermometry.<sup>59,60</sup> It has the disadvantage of having a fairly high phonon energy,<sup>61</sup> but our previous study has already shown that this disadvantage is negligible in low concentration suspension.<sup>62</sup> The thermometric characteristics have also been studied as a function of dispersion and laser power conditions.<sup>62</sup> As GdVO<sub>4</sub> is known to exhibit a phase transition from a zircon-like to a scheelite-like structure at 6.8 GPa, this nanosensor should be used exclusively in a pressure range below 6 GPa.<sup>63</sup> The concentration of NPs GdVO<sub>4</sub>:Yb<sup>3+</sup>/Er<sup>3+</sup> in the suspension was selected according to hydrodynamic size and rheological properties. The photoluminescent thermometry of the dispersed GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> nano-sensor was characterized up to 423K. The effect of pressure on the dispersed suspension of the GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> nano-sensor was studied in a diamond anvil cell (DAC) at pressures up to 1 GPa and at different temperatures.

## 2. Materials and Methods

### 2.1. Material

**GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> (x(Yb) = 0.10; x(Er) = 0.02) Nanoparticles** were synthesized by a chemical co-precipitation technique.<sup>16,26,64-66</sup> The precise synthesis method can be found in our previous work.<sup>62</sup>

**Suspension of NPs GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup>** was prepared by adding the ground powder in an ester-based oil (Nycobase 5750: conventional mono- pentaerythritol lubricating oil, supplied by NYCO, Paris, France) at various weight concentrations, followed by mechanical stirring and ultrasonic treatment at room temperature for 30 minutes.

### 2.2. Characterizations

The crystalline phase of the synthesized nanoparticles was characterized by an X-ray diffractometer (Bruker D8 Advance diffractometer) with Cu K $\alpha$  radiation ( $\lambda = 1.54060 \text{ \AA}$ ) operating at 45 kV and 30 mA at room temperature. The scanning electron microscopy (Zeiss Merlin Compact SEM) is performed for the analysis of the morphology, with a secondary electron detector at a low accelerating voltage up to 5 kV. The morphology of the nanoparticles was analyzed by transmission electron microscopy (TEM), JEOL 2100F 200 kV) with a secondary electron detector. The

hydrodynamic size of the NPs in the suspension was calculated by the Dynamic Light Scattering (DLS, Malvern-Zetasizer) at different concentrations.

The viscosity of the colloidal suspension was determined by a rheometer (Physica MCR301, Anton Paar) with a parallel-plate geometry (PP40), at different temperatures (293—373K,  $\pm 0.03$ K) and imposing a shear rate (range between 1 to  $10^4$  s<sup>-1</sup>). During the measurement, a steel plate is rotating against a glass plate, the gap between the two parallel plates was fixed at 0.2 mm. Transparent glass plate will be used for further optical measurement. The uncertainty on viscosity measurement is about 5%.

For the optical characterizations, the absorption spectrum of the NPs powder was acquired by LAMBDA 365 UV/Vis (Perkin Elmer) Spectrophotometer from 350 to 1100 nm, in an integrating sphere. The emission spectra were obtained by a modular spectrofluorometer, upon excitation of a continuous-wave NIR laser (970 nm, MDL-2W) focused on the sample to a spot size of 4 mm<sup>2</sup>. The spectrometer is equipped with a detector (CCD, Short Focal Length Triple Grating Imaging Spectrographs, TRIAX 320 HORIBA) and an optical fiber. At ambient pressure, the NPs powder were deposited on a silicon slice for better thermal transfer, the fluid was poured into a sapphire bowl closed with a glass lid for the emission acquisition. Data were collected after 10 s of temperature and laser stabilization, from 293 to 573K (each 5K) for the powder and from 298 to 423K (each 5K) for the fluid. The temperature was controlled by a thermo-cell (LINKAM Optical DSC600 cell).

Temperature calibrations under pressure were carried out using a high-pressure diamond anvil cell (DAC), as shown in Figure S1, with a working hydrostatic pressure range up to 1.3 GPa in Nycobase. According to the method described in a previous work,<sup>9</sup> the DAC (MDAC type BHP for Biology 100° symmetrical aperture, Betsa) consists of a 1.4 mm culet diamond anvil and a 400 mm thick diamond window, with a nickel gasket between them. The nickel gasket thickness was about 0.2 mm, and comprises a hole of 0.5 mm in diameter. The pressure inside the DAC was increased and maintained in the cell by inflating an internal membrane with helium. A heating ring (External Heating System, Betsa) was placed around the DAC for temperature regulation. A thermocouple inserted next to the metal seal in the DAC was used to measure the temperature of the sample with a precision of  $\pm 0.1$  K. During the experiments, a small amount of suspended sample ( $< 1\mu\text{L}$ ) was introduced into the hole in the seal. The hydrostatic pressure inside the DAC was determined by introducing ruby microparticles ( $> 3\mu\text{m}$ ) as a pressure sensor for calibration using their photoluminescent wavelength shift R1. These ruby microparticles were excited by a 460 nm laser. The wavelength position of the emission spectrum was referenced with an Argon lamp (Avalight-CAL-Mini,) with the peak at 696.54 nm. The calibration of the hydrostatic pressure was based on previous literature.<sup>67,68</sup> Based on previous work, the pressure uncertainty can be considered to be of the order of 0.05 GPa.<sup>9</sup>

### 2.3. Data analysis.

**Luminescence Intensity Ratio (LIR) calculation.** The Er<sup>3+</sup> thermometry is performed by the LIR variation in function of temperature. According to the Boltzmann distribution, the population ratio of the two thermally coupled levels (TCLs) ( $N_{L_2}$  and  $N_{L_1}$ ) varies with temperature following equation (1).

$$(1) \quad LIR = \frac{N_{L_2}}{N_{L_1}} = \frac{I_{L_2}}{I_{L_1}} = \frac{g_2 \sigma_2 \omega_2}{g_1 \sigma_1 \omega_1} \exp\left(-\frac{E_2 - E_1}{k_B T}\right) = B \exp\left(-\frac{\Delta E}{k_B T}\right)$$

where  $N_i$  denotes the population of excited ions, as the population of the energy level (N) is directly proportional to the emitted intensity (I),  $I_i$  is thus the emitted intensity for a given level  $i$ . The pre-exponential constant B includes the  $g_i$ ,  $\sigma_i$ ,  $\omega_i$ , which are the degeneracy, the effective emission cross section, and the angular frequency of the transitions for a given level  $i$ , respectively.  $\Delta E$  is the energy gap between the two TCLs,  $k$  is the Boltzmann constant, and  $T$  is the absolute temperature.<sup>33</sup> The emitted intensity was calculated by the integrated band area to average the total emitted intensities for each sublevels' transition.<sup>69</sup> The integrated area are estimated by MATLAB software according to Simpson's rule, from 540 to 567 nm for level  $L_1$  ( $^4S_{3/2}$  to  $^4I_{15/2}$ ) and from 512 to 540 nm for level  $L_2$  ( $^2H_{11/2}$  to  $^4I_{15/2}$ ) (Figure S2).  $\Delta E$  can be then determined by the calibration curve of  $\text{Ln}(LIR)$  versus  $1/T$ .

Moreover, the  $S_R$  (relative thermal sensitivity), defined by equation (2), is an important indicator of the reliability of the measurement.<sup>70</sup>

$$S_R = \frac{1}{LIR} \times \frac{d(LIR)}{dT} = \frac{\Delta E}{k_B T^2} \quad (2)$$

**LIR experimental uncertainty determination.** The intrinsic experimental uncertainty in the emission intensity ( $\delta I/I$ ) is detector dependent, and is mainly determined by the value of the signal-to-noise ratio (SNR)<sup>71</sup> (equation 1, Supporting information). The SNR of the CCD detector used for our measurement is stated by the manufacturer to be at best 0.05%. However, the detector operates with greater uncertainty and, therefore, the  $\delta I/I$  value is always higher than the indicated intrinsic SNR. The SNR is estimated experimentally for each spectrum by dividing the baseline spectral fluctuation by the peak intensity ( $N_{L2}$  and  $N_{L1}$ ). The experimental uncertainty is then determined by equation (3).

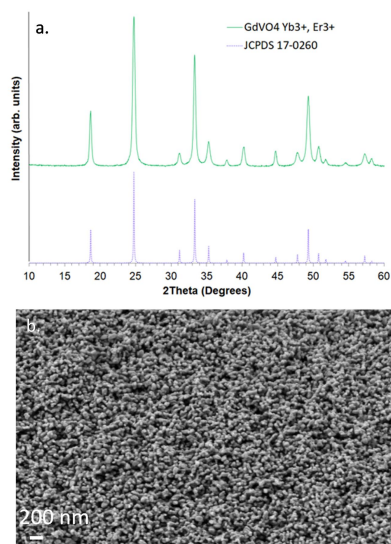
$$\frac{\delta(LIR)}{LIR} = \sqrt{\left(\frac{\delta I_1}{I_1}\right)^2 + \left(\frac{\delta I_2}{I_2}\right)^2} \quad (3)$$

### 3. Results

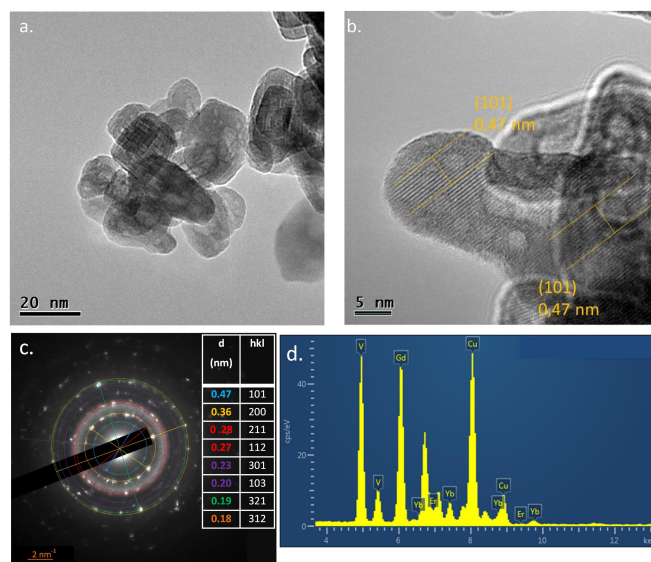
The results section will be divided into two parts. The first part deals with the characterizations of  $GdVO_4: Yb^{3+}/Er^{3+}$  powder and the second one with the characterizations of  $GdVO_4: Yb^{3+}/Er^{3+}$  dispersed at 0.5 wt.% in a lubricant oil (Nycobase 5750).

#### 3.1. Powder: $GdVO_4: Yb^{3+}/Er^{3+}$ NPs

**Structural characterizations.** The crystalline phase of the synthesized nanoparticles  $GdVO_4: Yb^{3+}/Er^{3+}$  has been determined to be a pure tetragonal phase (JCPDC-17-0260) in our previous work (Figure 1a).<sup>62</sup> The SEM image of NPs  $GdVO_4: Yb^{3+}/Er^{3+}$  show that the size of individual NP is about 30 nm (Figure 1b). The TEM image (Figure 2a) confirms that the particle size is of the order of 30 nm and shows the porosity of the particles induced by atomic diffusion during heat treatment. The HR TEM image (Figure 2b) shows reticular planes 0.47 nm apart, which is consistent with the plane distance (101) of  $GdVO_4$  (JCPDS-17-0260). The electron diffraction pattern (Figure 2c) confirms the presence of crystalline  $GdVO_4$  (JCPDS-17-0260). Finally, the EDS spectrum (Figure 2d) confirms the chemical composition of the probes, which contain Gd (28.9 at%), V (28.5 at%), Yb (4 at%) and Er (low content). This analysis leads to an Yb/Gd ratio of 0.14, consistent with the theoretical one (Yb/Gd=0.10). The agglomerates formed by NPs size about 400 nm, which has been confirmed by TEM image<sup>62</sup>, this result has also been observed in other's work.<sup>72</sup>



**Fig. 1** XRD pattern (a) and SEM image (b) of GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> NPs powder.

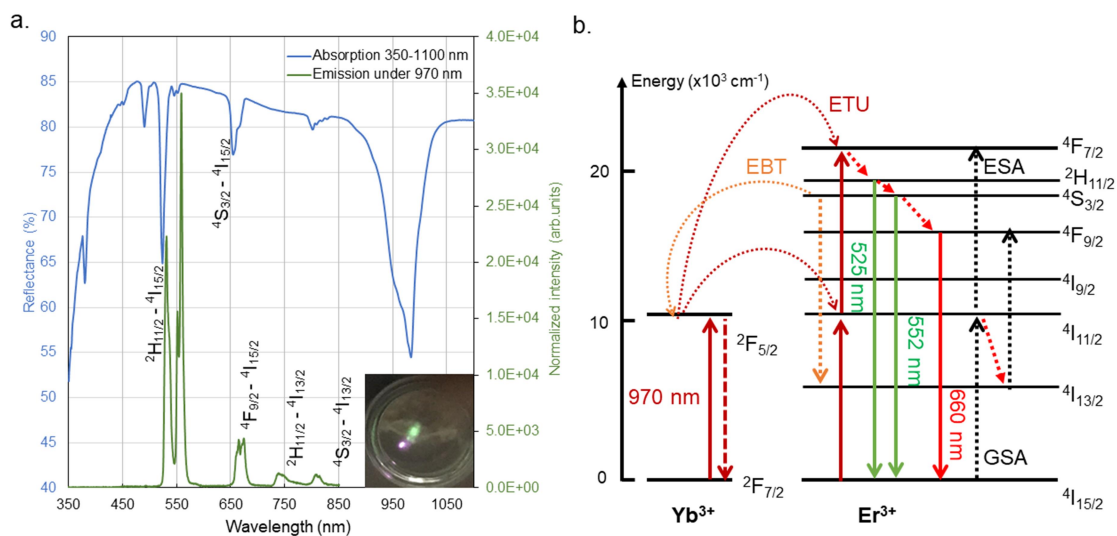


**Fig. 2** TEM picture (a), HRTEM picture (b) electron diffraction pattern (c), and EDS spectrum (d) of GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> NPs powder.

**Optical characterizations:** Figure 3a shows the powder absorption and emission (upon 970 nm excitation) spectra in blue and green line, respectively. The GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> NPs powder absorbs a broad band of light in the NIR range, around 985 nm, due to the presence of the Yb<sup>3+</sup> sensitizer of the Er<sup>3+</sup> emitter. In the emission spectrum upon NIR laser excitation (970 nm), the NPs are able to emit green light (525 & 552 nm), from the Er<sup>3+</sup> excited levels (<sup>2</sup>H<sub>11/2</sub> & <sup>4</sup>S<sub>3/2</sub>, respectively) to the ground state.<sup>28,72,73</sup> The inserted photograph shows as well, that the powder emits strong green light when using a NIR laser pen. The proposed excitation mechanism is thus presented in Figure 3b: upon excitation at 970 nm, Er<sup>3+</sup> can be excited to the <sup>4</sup>F<sub>7/2</sub> excited state by two-sequentially absorbed photons (Ground-state absorption (GSA) and excited-state absorption (ESA), respectively indicated by the black dotted arrow).<sup>60,74-77</sup> When Yb<sup>3+</sup> is present in the system, the energy transfer between Yb<sup>3+</sup> and Er<sup>3+</sup> would occur mainly in the Er<sup>3+</sup> excitation pathway. Since Yb<sup>3+</sup> has a much larger absorption



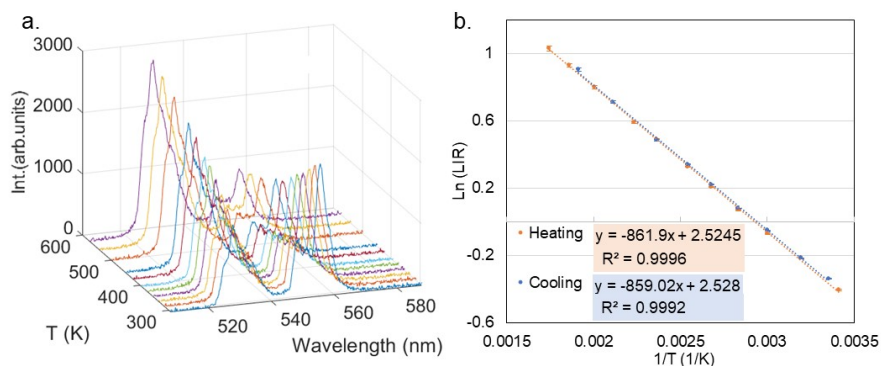
effective cross section in the NIR region than  $\text{Er}^{3+}$  ( ${}^2\text{I}_{15/2}$  to  ${}^4\text{I}_{11/2}$ ), the energy transfer upconversion process (ETU) becomes the main excitation of  $\text{Er}^{3+}$  to its upper levels ( ${}^4\text{F}_{7/2}$ ) to generate the same emissions. Other cross-relaxations and of some backward energy transfer (EBT: Er to Yb) can also occur in this system,<sup>76,78–80</sup> but they were not observed in our experiments. It is important to note that non-radiative relaxations are more noticed in the upconversion system, which induces a non-negligible laser-heating during high power density excitation. To prevent the laser-induced thermal effect from distorting the thermometric measurement, previous work<sup>62</sup> has shown that the laser intensity should be less than  $2.5 \text{ W/cm}^2$  for NPs powder.



**Fig. 3** (a) Absorption & emission spectra of the  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  powder sample. Inset: Photograph of the  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  powder sample with a 980 nm handheld laser pen. (b) The Upconversion process between  $\text{Er}^{3+}$  &  $\text{Yb}^{3+}$  under NIR excitation.

**Temperature calibration by photoluminescent thermometry:** The temperature was calibrated up to 573 K with the NPs powder under NIR laser at  $2.5 \text{ W/cm}^2$ . As shown in Figure 4a, the intensity of the peak  $\lambda_2$  increases while the peak  $\lambda_1$  decreases with temperature. The ratiometric thermometry following the Boltzmann distribution was proven by the temperature calibration, the LIR values varying with temperature (Figure 4b). According to the Boltzmann equation (1),  $\text{Ln}(\text{LIR})$  is linearly proportional to  $1/T$  up to 573 K, and more importantly, this variation is reversible. The slopes calculated by the fitted trendline of the calibration curves are used to estimate the energy gap ( $\Delta E$ ) in equation (1). The average of the slopes is used to estimate the  $\Delta E$  (Table 1). The thermal detection efficiency was evaluated by calculating the relative thermal sensitivity  $S_R$  (equation (2)), which was calculated with the average of the heating and cooling curves and is equal to  $0.81\% \cdot \text{K}^{-1}$  for the sample at 313 K. This value at low temperature is close to the literature data ( $0.47$  at  $298 \text{ K}$ <sup>73</sup> and  $1.11\% \cdot \text{K}^{-1}$  at  $307 \text{ K}$ <sup>81</sup>).



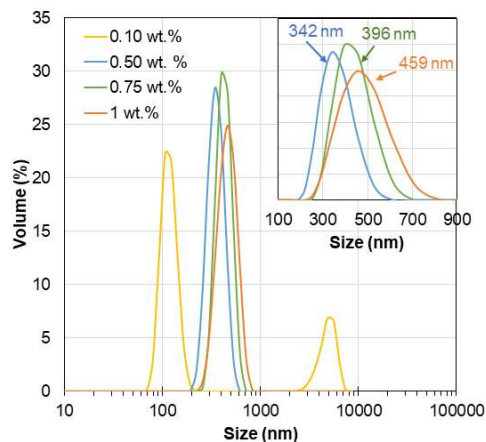


**Fig. 4** GdVO<sub>4</sub>:Yb<sup>3+</sup>/Er<sup>3+</sup> powder upon NIR laser at 2.5 W/cm<sup>2</sup>: (a) Emission spectra at 293 – 573 K. (b) Temperature calibration curve and fitted trendline equations, error bars determined by uncertainty (equation (3)).

### 3.2. In suspension: GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> NPs

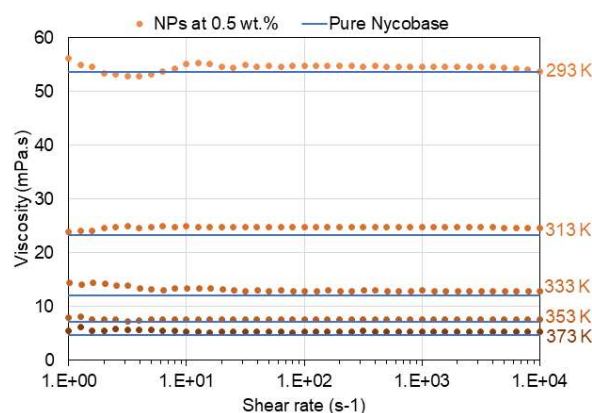
For further application in temperature measurement in tribology<sup>9</sup>, an ester-based oil (Nycobase 5750) was chosen as the fluid medium to disperse GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> nanoprobles. Nycobase 5750 is used as a based oil for lubricant in tribological studies due to its suitable properties: rheological, high thermal resistance, chemical stability, non-toxicity. The NPs were dispersed in the fluid and the suspension can be stable for more than 3 months. Furthermore, it was confirmed that Nycobase 5750 (Nyco) does not absorb any NIR light.<sup>62</sup> Thus, Nyco can be considered as a good fluid medium for the upconversion thermo-sensor, as it does not emit light upon NIR excitation and cannot be heated by the NIR laser.

**Hydrodynamic size measurements.** The NPs suspensions are prepared at different and relatively low concentrations: 0.1, 0.5, 0.75 and 1 wt.% for DLS analysis. As shown in Figure 5, the measured hydrodynamic size of the NPs suspension after one week settling. The hydrostatic size of the NPs aggregation is about 400 nm, which corresponds to the previous result observed in the TEM image.<sup>62</sup> However, although in the 0.1 wt.% suspension, the hydrodynamic size is the smallest among all, around 100 nm, the second peak at 5000 nm shows that the suspension is not homogeneous. The position of the peak centroid, as well as the width of the peak increases with concentration, which means that the hydrodynamic size increases while the monodispersity becomes more important with increasing concentration, respectively. Among the concentrations of 0.5, 0.75 and 1 wt.%, the suspension at 0.5 wt.% is the one that has both the smallest hydrodynamic size of NPs (less than 350nm) and the best monodispersity. It is important to note that the monodispersity and small size of nano probes are required for tribological applications to improve the spatial resolution.



**Fig. 5** The DLS diagram of GdVO<sub>4</sub>:Yb<sup>3+</sup>/Er<sup>3+</sup> suspension at 0.1, 0.5, 0.75 and 1 wt.%.

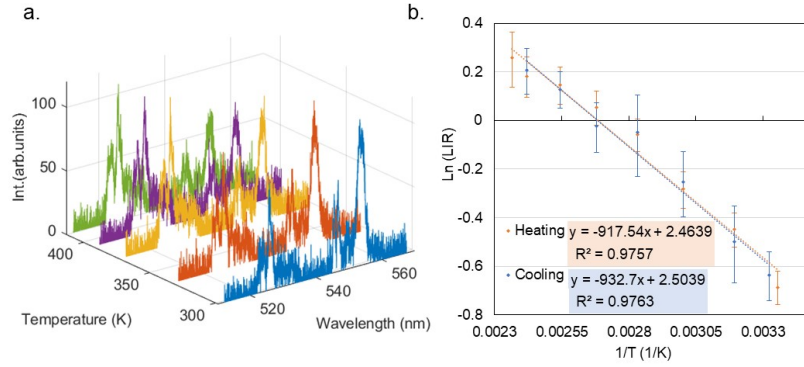
**Rheological properties characterizations.** Since it is essential to preserve the rheological properties of Nyco as much as possible for further friction tests, the viscosity of 0.5 wt.% suspensions is measured at different temperatures with pure Nycobase as a reference. Figures 6 show the measured viscosities of 0.5 wt.% and the reference of pure Nycobase at 293 to 373 K, respectively. As the pure Nycobase is a Newtonian fluid, its viscosity will not change with the shear rate, which is thus presented in line in Figure 6. The difference in viscosity between the 0.5 wt.% suspension and the pure Nycobase ( $< 2.97\%$ ) can hardly be observed in the experimental uncertainty range ( $< 5\%$ ). Moreover, at 293K, under high shear rate ( $> 8000 \text{ s}^{-1}$ ), the viscosity decreases very slightly with increasing shear rate for both suspensions and pure Nycobase, which is explained by shear-induced heating effect that decreases viscosity.<sup>8</sup>



**Fig. 6** Viscosity of  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  NPs suspension at 293 – 373 K under a shear rate of  $1 - 10^4 \text{ s}^{-1}$  for 0.5 wt.% NPs suspension (in dot) and pure Nycobase (in line).

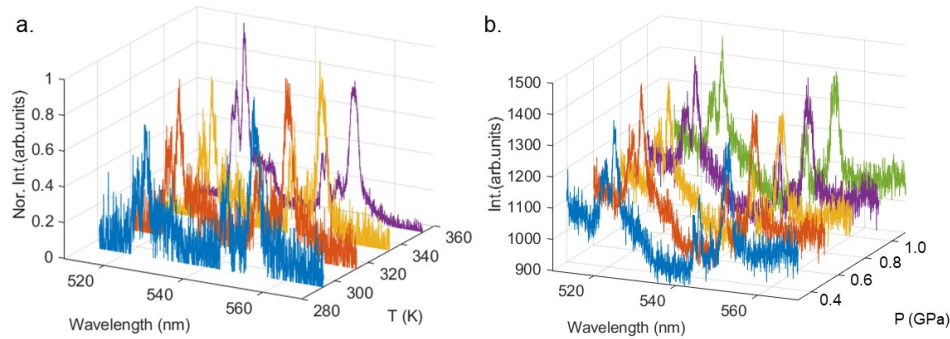
The viscosity of the 0.5 wt.% NPs suspension varies almost in the same way as that of the pure Nyco fluid. Thus, the dispersion of the NPs does not affect the rheological properties of the Nyco in the temperature and shear rate range studied, which is a crucial condition for using the NPs suspension as a lubricant. This means that the presence of  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  nanoprobes at 0.5 wt.% in the fluid will not change the tribological contact behavior. More importantly, for the 0.5 wt.% NPs suspension, it has been verified that there is no impact of NIR laser heating even at high laser intensity in our previous work.<sup>62</sup> Therefore, the 0.5 wt.%  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  nanoprobes suspension has been selected for further temperature calibration under atmospheric and under high pressure (up to 1.1 GPa) in the further study.

**Temperature calibration under atmospheric pressure.** Temperature calibration by upconversion emission was performed on the NPs suspension between 298 to 423 K (Figure 7a). The maximum temperature was limited to avoid possible evaporation of Nyco during the measurement. At 0.5 wt.% of NPs, the emission of the suspension could be measured at a high pump power density ( $13.5 \text{ W/cm}^2$ ). However, the calculated uncertainty (equation (3)) can be not negligible and has been represented on the calibration curve by an error bar. Although the spectral uncertainty is notable, the temperature calibration curve of the 0.5 wt.% NPs  $\{\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}\}$  is linear and reversible up to  $150^\circ\text{C}$ , with little deviation between the heating and cooling curves (Figure 7b). The energy gap ( $\Delta E$ ) was then fitted by the curves in Table 1, and the relative sensibility ( $S_R$ ) is  $1.17\%K^{-1}$  at 313 K, which is consistent with the values obtained with the powder calibration curve.

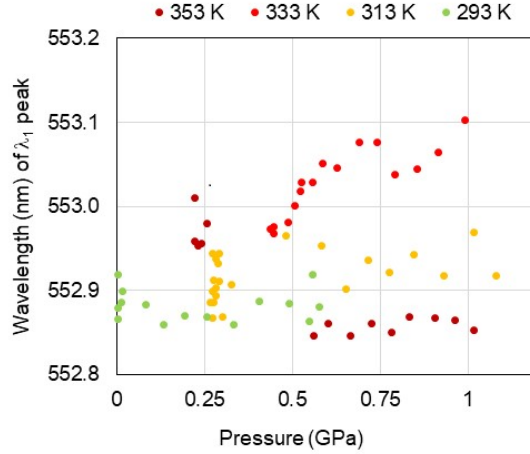


**Fig. 7**  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  suspended at 0.5 wt.% upon NIR laser at  $13.5 \text{ W/cm}^2$ : (a) Emission spectra at 298 - 423K. (b) Temperature calibration curve and fitted trendline equations, error bars determined by uncertainty (equation (3)).

**Temperature calibration under high pressure.** In order to apply the NPs suspension as a thermo-sensor lubricant for temperature measurement in tribological friction experiments, the photoluminescent thermometry of the suspension was studied under pressure (up to 1.1 GPa) in a diamond anvil cell (DAC) at every 20 K from 293 K to 353 K, under a  $13.5 \text{ W/cm}^2$  NIR laser as in the previous experiments. To study the effect of pressure on the LIR of NPs, two sets of data were collected in the DAC under isothermal conditions. The first group, called  $P_{\text{constant}}$ , is a repeated measurement at the same temperature and pressure (the lowest pressure for the corresponding temperature). The second group, called  $P_{\text{varying}}$  is a sequence of measurements at a given temperature for a variable pressure. The emission spectra are presented in Figure 8a at varying temperature and constant pressure ( $P = 0.6 \text{ GPa}$ ), showing that increasing temperature increases the emission intensity of peak  $\lambda_2$  and decreases that of the peak  $\lambda_1$ . Conversely, Figure 8b shows the negligible change in emission intensity under varying pressure at constant temperature. Besides the LIR measurement, the peak shift has been verified for peak 1 (at 553 nm). As shown in Figure 9, the corrected peak shift is not significant ( $< 0.2 \text{ nm/GPa}$ ) at a given temperature under varying pressure.

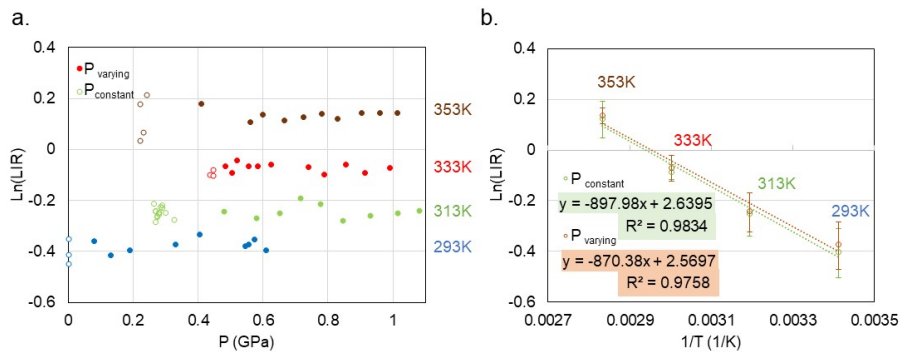


**Fig. 8**  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  suspended at 0.5 wt.% in DAC upon NIR laser at  $13.5 \text{ W/cm}^2$ : (a) Emission spectra at 293 - 353 K under 0.6 GPa; (b) Emission spectra at 313 K under 0.4 - 1.1 GPa.



**Fig. 9** GdVO<sub>4</sub>:Yb<sup>3+</sup>/Er<sup>3+</sup> suspended at 0.5 wt.% in DAC upon NIR laser at 13.5 W/cm<sup>2</sup>:  $\lambda_1$  peak position (corrected) under pressure at different temperatures.

As shown in Figure 10a, the Ln(LIR) under  $P_{\text{constant}}$  is plotted by a hollow point for each temperature, and the Ln(LIR) under  $P_{\text{varying}}$  is plotted by a solid symbol for each temperature as well. The LIR value under  $P_{\text{constant}}$  and  $P_{\text{varying}}$  do not change with increasing pressure but only with temperature. An analysis of variance is conducted to conclude the influence of pressure on the LIR variation. In Table 2, the variance (VAR) at  $P_{\text{constant}}$  (here the residual variance) and that at  $P_{\text{varying}}$  (here the explained variance) are calculated for each temperature. The uncertainties calculated by SNR of emission spectra are listed. The VAR( $P_{\text{varying}}$ ) is always lower than the VAR( $P_{\text{constant}}$ ). This means that the distribution of the LIR value is less dispersed when the pressure varies than when the pressure is constant. This means that the pressure has no effect on the LIR variation. To confirm this result, the ratio VAR( $P_{\text{varying}}$ )/VAR( $P_{\text{constant}}$ ) was calculated and found to be less than 1. This result shows that the slight change in the LIR value under isothermal conditions cannot be explained by the variation in pressure. Then, the LIR value is independent of pressure (up to 1.1 GPa) in the temperature range studied.



**Fig. 10** GdVO<sub>4</sub>:Yb<sup>3+</sup>/Er<sup>3+</sup> suspended at 0.5 wt.% upon NIR laser at 13.5 W/cm<sup>2</sup>: (a) Distribution of Ln(LIR) values at four temperatures (293 to 353 K) under constant pressure (hollow point) and pressure varying up to 1.1 GPa (solid point). (b) Temperature calibration curve and the fitted trendline equations under constant pressure (green) and pressure varying (brown), error bars determined by the standard deviation of overall data.

Therefore, the thermometric validity of NPs suspended under pressure was investigated by averaging Ln(LIR) under  $P_{\text{constant}}$  as well as  $P_{\text{varying}}$  for each temperature. As shown in Figure 10b, the temperature

calibration curve under  $P_{\text{constant}}$  and  $P_{\text{varying}}$  are both linear as a function of  $1/T$  and overlap with each other. The  $\Delta E_{\text{estimated}}$  by their slopes are therefore very close, 624 and 605  $\text{cm}^{-1}$  respectively (Table 1). This observation confirmed that the LIR of suspended NPs  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  is independent of pressure (up to 1GPa), and therefore the NPs are able to measure temperature by their emission without any effect caused by pressure variation.

**Table 1:** The estimated energy gap ( $\Delta E$ ) and the uncertainty of temperature calibration curves.

	Y offset	B	Slope	$\Delta E_{\text{estimated}}$	$\delta\text{LIR}/\text{LIR}$ (average)
NPs powder	2.5263	13	-860.46	598	1.92%
NPs in suspension	2.4839	12	-925.12	643	19.97%
NPs suspension ( $P_{\text{constant}}$ )	2.6395	14	-897.98	624	14.18%
NPs suspension ( $P_{\text{varying}}$ )	2.5697	13	-870.38	605	12.84%

**Table 2:** The average of LIR and variance of LIR for each temperature in DAC.

T (K)	$P_{\text{Constant}}$			$P_{\text{Varying}}$			VAR ratio*
	P (GPa)	SNR(%)	VAR	P max (GPa)	SNR(%)	VAR	
293	0.01	6.47	2.54E-03	0.61	6.48	6.30E-04	<b>0.25</b>
313	0.29	6.55	5.27E-04	1.08	5.97	5.00E-04	<b>0.95</b>
333	0.45	2.74	2.92E-04	0.99	4.22	2.81E-04	<b>0.96</b>
353	0.23	8.28	7.49E-03	1.02	3.47	4.06E-04	<b>0.05</b>

\*VAR ratio: the ratio of VAR  $P_{\text{Varying}}/ \text{VAR } P_{\text{Constant}}$

To summarize, if we compare the temperature calibration curve at the same laser pumping power (13.5  $\text{W}/\text{cm}^2$ ) of the 0.5 wt.% NPs suspension in our previous results (Table 1), the  $\Delta E_{\text{estimated}}$  obtained by the slope of the calibration curve could be considered as the same, which confirms the reliability of these NPs  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  for temperature measurement under a pressure up to 1 GPa. The uncertainty of the measurement depends on the intensity of the emission, the quantity of NPs in the suspension being much less important than in the powder, the uncertainty has become important in the suspension. Nevertheless, in the NPs suspension under atmospheric pressure, the 19% uncertainty caused only 7.5% error in the  $\Delta E_{\text{estimated}}$  value compared to the NPs powder. It is important to note that the configuration of the optical experiments is different between the temperature calibration under pressure in the DAC and at atmospheric pressure with a thermal controller, which may imply differences in the illuminations and therefore uncertainty in the emission intensity, and consequently in the LIR value. This may be the reason for a slight shift in the constant B in the Boltzmann equation. However, once the  $\Delta E$  is consistent with the previous result, the  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  NPs can be validated as thermometric and pressure independent probe.

#### 4. Conclusion

Non-contact, pressure-independent nanothermometer was investigated for tribological temperature measurement using  $\text{GdVO}_4:\text{Yb}^{3+}/\text{Er}^{3+}$  upconversion nanoparticles. The NPs were prepared by precipitation and thermal treatment. The NPs dispersed in the fluid lubricant has the same photoluminescent thermometric properties as the NPs powder. At low NPs dispersion concentration, the suspension retains the original rheological properties of the lubricant even at high temperature (373 K) and under a high shear rate ( $10^4 \text{ s}^{-1}$ ). In order to apply the NPs in frictional contact for temperature sensing, temperature calibration was performed under hydrostatic pressure (up to 1.1 GPa) in a diamond anvil cell. The result confirmed the independence of the LIR from pressure, and the

calibration curve under variable pressure is superimposed on that under constant pressure. Therefore, the incorporation of these GdVO<sub>4</sub>: Yb<sup>3+</sup>/Er<sup>3+</sup> nanosensors into the lubricant leads to the formation of a suspension that allows temperature measurement by the LIR method without being disturbed by the influence of pressure. The next step is the application in tribological contact, in particular in fluid lubrication.

## ASSOCIATED CONTENT

### Supporting information

#### Figures S1-S2

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C.J., L.B., D.P. and S.D. are the supervisors of projects. Y.Z., L.B. C.J and D.P. performed the experiments. G.L., D.P., N.B., M.M., Y.M., C.D, F.C. and D.F. provided input into the design of the experiments. Y.Z., G.L., L.B., N.B., D.P. and S.D. co-wrote the manuscript with input from other authors.

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