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# Mass transfer enhancement for CO<sub>2</sub> chemical absorption in a spiral baffle

# embedded microchannel

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**Abstract**: This paper presents an experimental study on  $CO_2$  chemical absorption by Monoethanolamine (MEA) solution in microchannel. In particular, a novel spirally distributed baffle structure is proposed to be embedded into the microchannel so as to generate the swirling flow pattern and to enhance the mass transfer. The effect of spiral baffle structure on the two phase flow pattern transition, the  $CO_2$  gas void fraction, the pressure drop, and the mass transfer performance are characterized and discussed, under varied gas and liquid flow rates, and MEA concentrations. Results demonstrate that the spiral baffle structure could significantly enhance the gas-liquid mass transfer, with overall liquid side volumetric mass transfer rate ( $k_L a$ ) up to 2.35 s<sup>-1</sup> and enhancement number E up to 1.387 at small gas and liquid Reynolds numbers ( $Re_G < 5.5$ ;  $Re_L < 14$ ). New predictive correlations have also been proposed by considering both the structure enhancement factor and the chemical reaction factor.

**Keywords**: Microchannel, spiral baffle, gas-liquid two-phase flow, mass transfer enhancement, carbon capture, CO<sub>2</sub> chemical absorption

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Nomencl	ature			
A	Mass transfer interface area $(m^2)$	Greek syn	nbols	
а	Specific interface area $(m^{-1})$	$\delta_0$	Liquid film thickness (m)	
$C_{MEA}$	MEA volumetric concentration	μ	Dynamic viscosity (Pa.s)	
C*	$CO_2$ equilibrium concentration ( $mol.m^{-3}$ )	σ	Surface tension $(N. m^{-1})$	
D	Diffusion coefficient $(m^2s^{-1})$	ρ	Density $(kg.m^{-3})$	
$d_h$	Hydraulic diameter (m)	β	Dimensionless ratio of gas-liquid flow rate	
E	Enhancement number	Δ	Difference	
Еc	Energy consumption ratio	$\Psi$	Intensified structure number	
f	Frequency	ε	Gas void fraction	
H	Henry constant $(mol. m^{-3}. Pa^{-1})$	Ø	Relative uncertainties	
$k_{ov}$	Overall reaction rate $(s^{-1})$	η	Effective energy efficiency	
$k_L$	Liquid side mass transfer coefficient $(m. s^{-1})$			
b	Liquid side physical mass transfer coefficient			
$k_{L,P}$	$(m. s^{-1})$			
l	Length (m)			
n	Molar mass (mol)			
r	fillet radius	Dimensionless number		
R	Gas constant $(J. mol^{-1}K^{-1})$	Са	Capillary number, $Ca = \mu U/\sigma$	
P	Pressure (Pa)	Da	Damköhler number, $Da = k_{ov} \cdot d_h^2 / D$	
Q	Volumetric flow rate $(ml. h^{-1})$	На	Hatta number, $Ha = \sqrt{D k_{ov}/k_{L,p}^2}$	
T	Experimental temperature (K)	Re	Reynolds number, $Re = \rho U d_h / \mu$	
t	Absorption duration time (s)	Sh	Sherwood number, $Sh = k_L d_h / D_{AB}$	
U	Superficial velocity $(m. s^{-1})$	Sc	Schmidt number, $SC = \mu_L/D_{AB} \cdot \rho_L$	
V	Volume $(m^3)$	We	Weber number, $We = \rho U^2 d_h / \sigma$	
w	Microchannel width (m)			
Subscripts				
В	Bubble	in	Inlet	
C	Channel	ini	Initial	
G	Gas	out	Outlet	
L	Liquid			

#### 1. Introduction

The increase of CO<sub>2</sub> concentration in atmosphere is one of the main causes of global warming, whose disastrous consequences have been well recognized (Florides and Christodoulides, 2009) (Shakun et al., 2012). The development and implementation of efficient CO<sub>2</sub> capture technologies receive much more attention nowadays, aiming at the Carbon neutrality target by year 2050 (Schreyer et al., 2020). Conventional CO<sub>2</sub> capture methods include chemical absorption, physical absorption or adsorption, membrane separation, cryogenic distillation, etc. (Koytsoumpa et al., 2018; Lamb, 2017; Murai and Fujioka, 2008; Stewart and Hessami, 2005). Among them, the chemical CO<sub>2</sub> absorption method has shown to be an adapted and widely used one for large-scale industrial application, such as for power plant flue CO<sub>2</sub> treatment. But this method usually involves voluminous equipment, inevitably bringing some problems such as high investment costs, huge amount of absorbent needed, high energy consumption, and relatively low mass transfer performance. These drawbacks may be solved through the miniaturization, more precisely by the employment and implementation of micro-structured absorbers/reactors (Lin et al., 2019; Mc Carogher et al., 2021). The channel characteristic size of the microreactor is usually in the order of tens to hundreds of microns, greatly increasing the specific interfacial area for mass transfer. Owing to this benefit, numerous studies have been conducted on the microchannel based CO<sub>2</sub> chemical absorption over the past decades (Zanfir et al., 2005; TeGrotenhuis et al., 2000), the latest advances are summarized in recent review papers (Cheng et al., 2023; Pasha et al., 2022). However, some limitations still exist, mainly due to the laminar regime of two-phase flow at small Reynolds numbers (Re) of liquid and gas in the microchannel. The diffusion dominates the mass transfer between CO<sub>2</sub> bubble and the liquid solvent, restraining the absorption performance. In brief, researches are still needed on the microchannel-based CO<sub>2</sub> absorbers for performance improvement (Ma et al., 2020; Tan et al., 2012).

In this context, many intensification measures have been proposed, implemented and tested, working either on the global geometry (Pang et al., 2021) or inner structure of the microchannel (Chen et al., 2019), or on some special operational modes (Dong et al., 2015). Among these methods, installing flow disturbance structures in the microchannel could induce the local eddy

flow and deform the gas-liquid interface, thereby constituting an effective and economic way to improve the  $\mathrm{CO}_2$  absorption performance. A number of experimental studies have been performed with different channel and obstacle geometries (micro grooves, micro protrusions, etc.), as summarized in Table 1. For example, Zhang et al. (2021) proposed the sudden expansion microstructure accompanied microchannel to enhance the gas-liquid interface disturbance. Their results showed that the liquid side volume mass transfer rate ( $k_L a$ ) could be increased by a factor of 2.4. In particular, rectangular-shape baffle structure has been applied for the mass transfer enhancement. Yin et al. (2020) fabricated a rectangular microchannel with staggered rectangular baffles and investigated the effect of baffle parameters (length, arrangement, etc.) on the  $\mathrm{CO}_2$ /water mass transfer performance. Their results showed that the  $k_L a$  could be augmented by a factor of 2.8 owing to the baffle structure. Recently, Yin et al. (2022a) proposed a continuous sudden expansion microchannel structure accompanied with rectangular obstacles in the middle (cf. Table 1) to create the bubble flow splitting and recombining. Their results demonstrated that the  $k_L a$  could be increased by 2.74 times.

All these studies have verified the intensified CO<sub>2</sub> absorption performance by obstacle structures, but at the cost of an increased pressure drop. Further optimizing the obstacle structure embedded microchannel to intensify the mass transfer with a reasonable increase of pressure drop is a key issue in this field. This is especially true for small gas/liquid *Re* number operation for which the mass transfer is rather limited but interesting in terms of low presure loss. The enhancement of mass transfer by proper obstacle geometry and arrangement is of great practical significance, but systematic researches on this topic are more than needed. Besides, for CO<sub>2</sub> chemical absorption in microchannel, the conventional correlations based on common physical dimensionless numbers for the prediction of gas-liquid mass transfer performance become less applicable. Some new correlations considering the effect of chemical reaction have been proposed, showing better prediction accuracy. Nevertheless, the introduction of chemical reaction factor into the correlations is still complicated, letting alone the prediction of additional enhancement effect by flow disturbance structures.

In this work, we propose a novel spiral baffle structure embedded into the microchannel to

intensify the CO<sub>2</sub> chemical absorption at small gas and liquid *Re* numbers. The spirally distributed obstacles at the inner walls of microchannel are expected to generate a swirling flow pattern and more spatially asymmetric vortices in liquid phase to enhance the mass transfer. The comparative experimental study of CO<sub>2</sub> absorption in Monoethanolamine (MEA) solvent in a spiral baffle embedded microchannel and in a smooth microchannel was performed. The influence of spiral baffle structure on the two-phase flow pattern, the gas void fraction, the pressure drop, and the CO<sub>2</sub> absorption rate were experimentally investigated under a varied gas and liquid flow rates and MEA concentrations. Based on the obtained experimental data, new correlations were also proposed for predicting the liquid side volumetric mass transfer rate by considering both the structure enhancement factor and the chemical reaction factor, more appropriate for the cases like CO<sub>2</sub> chemical absorption in structured microchannels. Furthermore, the energy consumption of the CO<sub>2</sub> absorber with enhancement structure were evaluated. The obtained results could contribute to the potential industrial application of minimicrochannel reactors for carbon capture.

Table 1. A summary of recent microchannel based CO<sub>2</sub> absorption enhancement studies by embedded obstacle structures

Study	Two-phase flow system	Operating range	Flow pattern/working conditions	Microchannel geometry & dimension (mm)	Intensification measure	Structure diagram	$k_L a$ improvement
Chen et al. (2021)	CO <sub>2</sub> [Bmim][BF <sub>4</sub> ] aqueous	$Q_L$ : 50 - 80 $ml/h$ $Q_G$ : 80 - 360 $ml/h$	Slug flow $298.15 \pm 1 \text{ K}$ Ambient pressure	T-type rectangular straight microchannel 0.4×0.8×35.2	0.1 mm thickness heart-shaped groove		2.2 times
Yin et al. (2020)	CO <sub>2</sub> Water	$U_L$ : 0.043 - 0.07 $m/s$ $U_G$ : 0.017 - 0.26 $m/s$	Slug/broken slug flow 298.15 K Ambient pressure	T-type rectangular straight microchannel $0.8\times0.8\times30.4$ $0.6\times0.6\times30.4$ $0.4\times0.4\times30.4$	Staggered rectangular baffles	S mm J <sub>b</sub> 0.2 mm J <sub>d</sub> 1.5 mm J <sub>d</sub> 30.4 mm	2.8 times
Yin et al. (2022a)	CO <sub>2</sub> AMP aqueous	$Q_L$ : $40 - 100 \ ml/h$ $Q_G$ : $20 - 500 \ ml/h$	Complete/partial breakup 293.15 K Ambient pressure	Optimized T-type rectangular straight microchannel 0.4×0.6×25	Split–and– recombine structure	Gas Inlet 8 12 5	2.74 times
Zhang et al. (2021)	CO <sub>2</sub> [Bmim][BF <sub>4</sub> ] aqueous	$Re_L$ : 1.67 - 27.22 $Re_G$ : 11.81 - 32.30	Slug/slug-bubbly flow 298.15 K Ambient pressure	T-type rectangular straight microchannel 0.4×0.4×34	Sudden expansion structure		2.4 times
Yin et al. (2019b)	CO <sub>2</sub> [Bmim][BF <sub>4</sub> ] MEA aqueous	$Q_L$ : 30 - 60 $ml/h$ $Q_G$ : 20 - 500 $ml/h$	Slug/broken slug/slug- bubbly flow 298.15 K Ambient pressure	T-type rectangular straight microchannel 0.4×0.8×30.4	Up-down embedded baffles	co. The state of t	1.5 times
Nie et al. (2022)	CO <sub>2</sub> [Bmim][BF <sub>4</sub> ] aqueous	$Q_L$ : 50 - 80 $ml/h$ $Q_G$ : 40 - 220 $ml/h$	Slug/broken slug flow 298.15 K Ambient pressure	T-type rectangular straight microchannel 0.8×0.8×30.4	Up-down embedded triangular obstacles		2.1 times

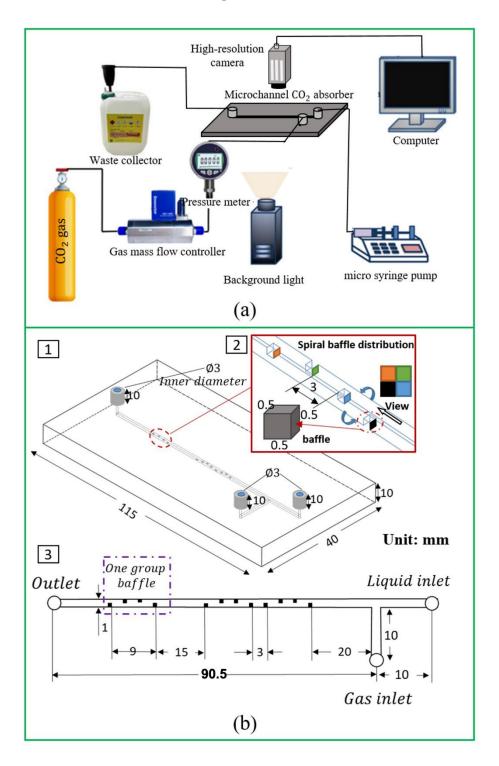
#### 2. Methodology

# 2.1 Experimental setup and microchannel reactor

Fig. 1a shows the experimental setup used for this study. CO<sub>2</sub> (Air Liquide, France, 99.9% purity) was used as the gas phase and its flow rate was controlled by a mass flow controller (Brooks SLA 5850). MEA (Arcane-industries, France, 99.9% purity) solution was used as the liquid phase and was delivered at constant flow rate by a micro syringe pump (kdScientific-267, 1% uncertainty). A camera with 4112 × 3008 -pixel resolution (Basler acA4112-30μm) equipped with lens (EDMUND optical 50mm/F2.0 59873, USA) was used to monitor and record the two-phase flow in microchannel after the two-phase flow pattern and pressure reached steady. The camera was operated at a frame rate of 30 fps. Basler video recording software was used to control and adjust the image parameters captured by high resolution camera. A 36w flat uniform cold light source (LED panel from LEDVANCE) was positioned close behind the microchannel to provide background illumination during images recording process. A pressure gauge (Kobold Pad FEE 3S2NS00 series) was installed at the gas inlet port to measure the pressure drop, while the pressure at the outlet was estimated as atmospheric pressure. The measuring range and precision of these instruments are given in Table 3.

Two microchannel reactors (absorbers) were tested and compared in this study for CO<sub>2</sub> chemical absorption: one smooth straight microchannel and another with intensification structure. Both microchannels have a square cross-section of 1 mm × 1 mm, the length of main channel being 90.5 mm. A T-type junction was designed for gas-liquid contact. The inlet branches have the same cross-section dimension as the main channel, and the length for each inlet branch is 10 mm. For the microchannel with intensification structure, three spiral baffle groups were embedded in the main channel. Each baffle group consists of 4 small cubes (0.5 mm side), arranged in rotation and quarterly distributed in the inner walls of the microchannel with 3 mm spacing (Fig. 1b). Such novel baffle structure is intended to create the swirling flow and the spatial asymmetric two-phase interface disturbance, so as to enhance the mass transfer. A space of 15 mm was left between the second and the third group, for better recognizing the bubble volume changes after passing through the first two groups which were successively

arranged. More details about the geometry and dimensions of the microchannel with intensification structure can be found in Fig. 1b.



**Figure 1** Schematic diagram of the experimental setup. (a) the test-rig; (b) 3-D diagram of spiral baffle embedded microchannel.

Both microchannel absorber prototypes (Fig. 1b) were realized with transparent resin

(WaterClear) using laser stereolithography. In this process, photopolymer was cured selectively at the local focus of ultraviolet laser. The external walls of the prototypes were further polished for better transparency.

The tested flow rate range was 12-168 ml/h for the gas phase and  $10-50 \, ml/h$  for the liquid phase, respectively. The corresponding Re number ( $Re = \rho U d_h/\mu$ ) range was 0.42-5.51 (for gas) and 2.78-13.9 (for liquid), respectively. Note that  $\rho$  is the density  $(kg/m^3)$ , U is the superficial velocity (m/s),  $d_h$  is the hydraulic diameter of the main microchannel (m) and  $\mu$  is the dynamic viscosity  $(Pa \cdot s)$ . Experiments were performed using MEA solution with three different concentrations, i.e.,  $C_{MEA} = 1\%$ ; 3% and 5% vol.

All experiments were conducted under ambient temperature (293 K) and pressure. For each measurement, it took about 10 min for the stabilization of two-phase flow. The gas-liquid two-phase flow patterns observed and the absorption performance show good stability and reproducibility, which will be presented in section 3.1.

# 2.2 Chemical reaction theory

The CO<sub>2</sub> chemical absorption by MEA can be characterized by zwitterion mechanism (Lv et al., 2015). Several principal reaction steps could occur in the CO<sub>2</sub>-MEA-H<sub>2</sub>O absorption system. Firstly, the MEA molecules quickly react with CO<sub>2</sub> to form intermediate zwitterion (Eq. 1). Then the zwitterion is quickly neutralized to form carbamate (Eq. 2). The whole process can be combined as carbamate formation and expressed as Eq. 3.

$$CO_2 + R'NH_2 = R'NH_2^+COO^-$$
 (1)

$$R'NH_2^+COO^- + R'NH_2 = R'NHCOO^- + R'NH_3^+$$
 (2)

$$CO_2 + 2R'NH_2 = R'NHCOO^- + R'NH_3^+$$
 (3)

Where R' is  $CH_2CH_2OH$ ,  $R'NH_2^+COO^-$  and  $R'NHCOO^-$  are the zwitterion and the carbamate, respectively. Meanwhile,  $HCO_3^-$  is formed by the  $CO_2$  hydration under high  $CO_2$  loading (Eq. 4) and accompanied by the hydrolysis of carbamate (Eq. 5).

$$CO_2 + H_2O + R'NH_2 = HCO_3^- + R'NH_3^+$$
 (4)

$$CO_3^{2-} + H^+ + R'NH_2 = R'NHCOO^- + H_2O$$
 (5)

For CO<sub>2</sub> desorption, the carbamate is decomposed into MEA and CO<sub>2</sub> under heating condition and the solvent is regenerated (Eq. 6).

$$R'NHCOO^- + H^+ \xrightarrow{heating} R'NH_2 + CO_2(g)$$
 (6)

The chemical reaction rate of MEA-CO<sub>2</sub> is fast enough, so that the CO<sub>2</sub> concentration in bulk MEA solution is considered to be '0'.

Different MEA concentrations have been used in this study. The physicochemical properties are listed in Table 2. Note that the diffusion coefficient (D) and overall chemical reaction rate  $(k_{ov})$  of CO<sub>2</sub>- MEA solution reaction system are refer to papers (Yin et al., 2019a).

**Table 2.** Physicochemical properties of  $CO_2$ - MEA solution reaction system

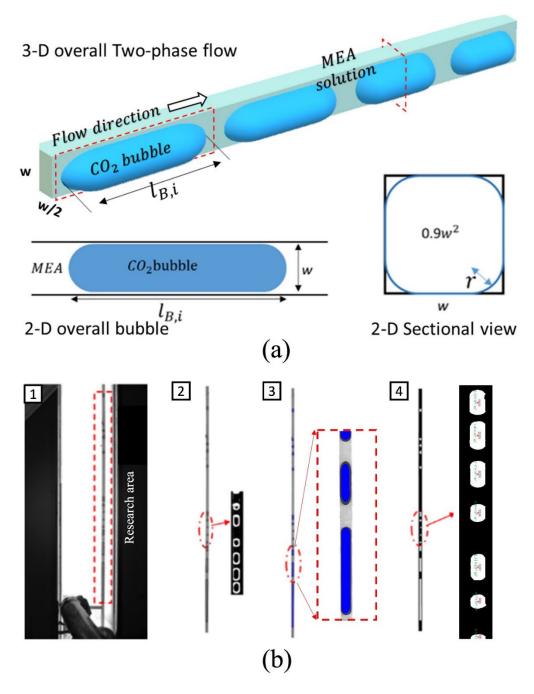
$C_{MEA}$	Density	Viscosity	Diffusivity	Chemical reaction rate
	$\rho_L({\rm kg/m^3})$	$\mu_L \times 10^3 (\text{Pa.s})$	$D \times 10^9 (\text{m}^2/\text{s})$	$k_{ov}(s^{-1})$
1%	998.6	0.978	1.99	451.06
3%	998.9	0.986	2.03	1403.18
5%	999.1	1.041	2.08	2355.24

#### 2.3 Calculation of bubble volume and surface area

The variation of CO<sub>2</sub> bubble volume and surface area over time in microchannel are key parameters for the evaluation of absorption performance. In this study, the optical capturing and image processing method has been used to acquire the transient bubble geometry and dimensions. For slug flow, a single CO<sub>2</sub> bubble was treated as 'two hemisphere caps + symmetrical main body', as shown in Fig. 2(a). The thickness of liquid film between the bubble and the channel wall can be estimated as (Youn et al., 2015):

$$\frac{\delta_0}{d_h} = \frac{0.67Ca^{2/3}}{3.13Ca^{2/3} + 1 + 0.504Ca^{0.672}Re_L^{0.589} - 0.352We^{0.629}} \quad (Ca < 0.04, Re_L < 2000)$$
 (7)

Where Ca is the capillary number  $(\mu U/\sigma)$  and We is the Weber number  $(\rho U^2 d_h/\sigma)$ , respectively.  $\sigma$  is the two-phase surface tension (N/m).



**Figure 2** Image processing for the calculation of bubble volume and surface area. (a) Schematic diagram of gasliquid two-phase flow in microchannel. (b) MATLAB based image processing method

For this study ( $Ca < 1 \times 10^{-3}$ ,  $Re_L < 15$ ), the dimensionless liquid film thickness ( $\delta_0/d_h$ ) is estimated to be smaller than  $6.5 \times 10^{-3}$ , thus could be neglected in the calculation of bubble volume and interfacial area. The bubble cross sectional area is estimated to be equal to  $0.9w^2$  (Zhu et al., 2017). The single slug bubble volume ( $V_{B,i}$ ) and surface area ( $A_{B,i}$ ) are then decided by the bubble length, calculated as follows:

$$V_{B,i} = \frac{\pi w^3}{6} + 0.9w^2 (l_{B,i} - w) \quad l_{B,i} > w$$
 (8)

$$V_{B,i} = \frac{\pi l_{B,i}^{3}}{6} \qquad l_{B,i} \le w \tag{9}$$

$$A_{B,i} = \pi w^2 + (2\pi r + 4w - 8r)(l_{B,i} - w) \quad l_{B,i} > w$$
 (10)

$$A_{B,i} = \pi l_{B,i}^{2} \qquad l_{B,i} \le w \tag{11}$$

The bubble length ( $l_{B,i}$ ) and the fillet radius (r) were obtained by analyzing the two-phase flow pictures captured by high-resolution camera using MATLAB (v2020a) tools, as explained in Fig. 2b. The original RGB images were converted to gray images and the two-phase flow areas were extracted by MATLAB. The gas-liquid interface was identified by using edge recognition algorithm 'Sobel'. Then the interface gray image was converted to black-white binary image to reduce the image background noise. In order to verify the processed result, the recognized bubble area was compared with the original images. Since the size and the scale of the pixels are known, the 2-D bubble size parameters could be acquired by counting the number of pixels in every single bubble. The real bubble instantaneous volume and surface area could then be calculated based on Eqs. 8-11.

#### 2.3 Mass transfer calculation

For single CO<sub>2</sub> bubble, the reduction of CO<sub>2</sub> molecules ( $\Delta n_{co_2}$ , mol) can be expressed by:

$$\Delta n_{co_2} = \frac{P_{in}V_{in} - P_{out}V_{out}}{RT} \tag{12}$$

Where R  $(J/mol \cdot K)$  is the ideal gas constant, T (K) is the working temperature,  $P_{in}$  (Pa)

is the gas pressure at microchannel inlet measured by pressure gauge,  $P_{out}$  (Pa) is the pressure at microchannel outlet (estimated as atmosphere pressure),  $V_{in}$  ( $m^3$ ) is the initial bubble volume and  $V_{out}$  ( $m^3$ ) is the bubble volume at the end of microchannel.

Two-film model (Kashid et al., 2011) is used for mass transfer calculation, assuming that there are a gas stagnant layer (gas film) and a liquid stagnant layer (liquid film) on each side of the gas-liquid interface. For quick chemical reaction and pure  $CO_2$  gas phase, the liquid-side mass transfer resistance is dominant while the gas-side mass transfer resistance can be ignored (Whitman, 1923). Therefore, the  $\Delta n_{CO_2}$  for a single  $CO_2$  bubble could also be calculated as:

$$\Delta n_{CO_2} = A_{B,i} k_{L,i} (C^* - C_{CO_2}) t \tag{13}$$

Where  $k_L$  (m/s) is the liquid side mass transfer coefficient, t (s) is the absorption duration,  $C^*$  ( $mol/m^3$ ) is the equilibrium CO<sub>2</sub> concentration in MEA solution which is decided by Henry's law.  $C_{CO_2}$  is the CO<sub>2</sub> concentration in liquid slug at t time (estimate as 0). Thus, for a single bubble, the volumetric mass transfer rate  $k_{L,i}a_i$  ( $s^{-1}$ ) could be calculated as follows:

$$k_{L,i}a_i = \frac{P_{in}V_{in} - P_{out}V_{out}}{V_CC^*tRT}$$
(14)

$$C^* = P_{CO_2} \times H \tag{15}$$

Where  $a_i$  (1/m) is the specific surface area of single bubble which is defined as  $A_{B,i}/V_C$ ,  $V_C$  is the microchannel volume,  $P_{CO_2}$  is the average CO<sub>2</sub> pressure in microchannel, calculated as:

$$P_{CO_2} = \frac{\ln(\frac{P_{in}}{P_{out}})}{P_{in} - P_{out}} \tag{16}$$

H is Henry's constant estimated to be  $3.78 \times 10^{-4}$  (mol/m<sup>3</sup>. Pa) (Versteeg and Van Swaaij, 1988).

For continuous and stable slug flow in microchannel, the initial CO<sub>2</sub> bubble volume  $(V_{in})$  and bubble generation frequency  $(f_{CO_2})$  could be estimated as constant. Thus, the overall liquid

side volumetric mass transfer rate  $k_L a$  (s<sup>-1</sup>) could be calculated as:

$$k_L a = \frac{(P_{in}V_{in} - P_{out}V_{out}) \cdot f_{CO_2}}{V_C C^* tRT}$$
(17)

# 2.5 Uncertainty Analysis

The estimated uncertainties of experimentally measured parameters and calculated results are listed in Table 3. The CO<sub>2</sub> mass flow controller (Brooks SLA 5850) and the syringe pump (kdScientific-267) were calibrated by standard volumetric method and the relative uncertainties are  $\pm 0.03$  ml/min and  $\pm 0.1$  ml/h, respectively. The experimental temperature was measured by a calibrated K-type thermocouple with uncertainty of  $\pm 0.5$  K. The uncertainties of calculated parameters were estimated based on error propagation method (Moffat, 1988). In this study, the bubble parameters were captured by high-resolution camera and used to calculate the  $V_{B,i}$ . The largest uncertainty is 2 pixels (66 µm) and its contribution to the relative uncertainties ( $\emptyset$ ) of  $V_{B,i}$  is showed in Fig. 3. With the decrease of bubble volume, the relative uncertainty becomes much larger. The initial bubble length is larger than 5 mm under the tested conditions and the maximum relative uncertainty of bubble volume is estimated to be smaller than  $\pm 4\%$  (Fig. 3a). The uncertainty of CO<sub>2</sub> gas phase void fraction ( $\varepsilon$ ) could be caused by microchannel fabrication inaccuracy and  $V_{B,i}$ . The dimension error due to 3D printing is estimated to be in the order of 0.05 mm which can be ignored regarding the whole microchannel volume. The relative uncertainty of  $\varepsilon$  is thereby estimated to be equal to that of bubble volume. Under the testing range of this study, the bubble volume change is smaller than  $4.9 \times 10^{-9}~m^3$ and the relative uncertainty of  $k_L a$  (Eq. (17)) is estimated to be  $\pm$  3%. With the increase of  $Q_G$ , the uncertainty of  $k_L a$  could gradually decrease and become smaller than  $\pm 1\%$  (Fig. 3b). The uncertainty for the enhancement number E  $((k_L a)_{structured}/(k_L a)_{smooth})$  is decided by  $\emptyset$  ( $k_L a$ ) and its maximum value is estimated to be  $\pm$  6.21%.

Table 3 Estimated uncertainties of measured and calculated parameters

Measured parameter		Uncertainty
Liquid volumetric flow rate	$Q_L$	±0.1 ml/h
Gas volumetric flow rate	$Q_G$	±0.03 ml/min
Pressure drop	$\Delta P$	±5.6 Pa
High speed camera pixels	-	2 pixels (±66 μm)
Experimental temperature	T	±0.5 K
Calculated parameters		Uncertainty
Initial single bubble volume	$V_{B,i}$	±3.91%
Void fraction	${\cal E}$	±3.91%
Liquid side volumetric mass transfer rate	$k_L a$	±3.01%

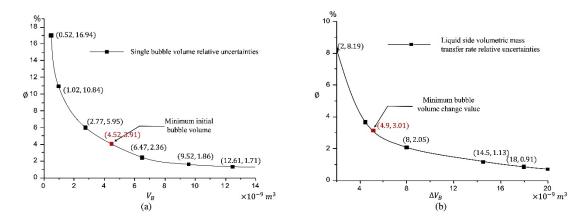


Figure 3 Range of relative uncertainties for calculated parameter. (a) Single bubble volume  $V_{B,i}$ ; (b) Liquid side volumetric mass transfer rate  $k_L a$ .

## 3. Results and discussion

# 3.1 Two-phase flow pattern

The gas-liquid two-phase flow pattern plays a determinant role on the mass transfer characteristics of CO<sub>2</sub> chemical absorption in microchannel (Niu et al. 2009). The chemical reaction accompanied two-phase flow patterns become much more complex because of the enhanced mass transfer dynamics. In this study, CO<sub>2</sub> chemical absorption has been performed in two above-presented microchannels. According to the changes of bubble shape along the

microchannel, various two-phase flow patterns were captured under the tested flow rate conditions, including bubbly flow, slug-bubbly flow, slug flow, and slug-annular flow. Some typical flow patterns are shown in Fig. 4 ( $C_{MEA}$ : 5% vol.) and briefly described below.

Bubbly flow: This flow pattern appears at low  $Q_G$  and high  $Q_L$  with short initial bubble length  $(l_{B,ini} < 5d_h)$ . Due to the fast chemical reaction accompanied mass transfer, the gas bubble quickly shrinks into a small bubble within a very short distance from the junction.

Slug-bubbly flow: With the increase of  $Q_G$ , the initial bubble has a longer length, the initial two-phase flow pattern shows slug flow pattern and gradually transits to bubbly flow due to the chemical absorption.

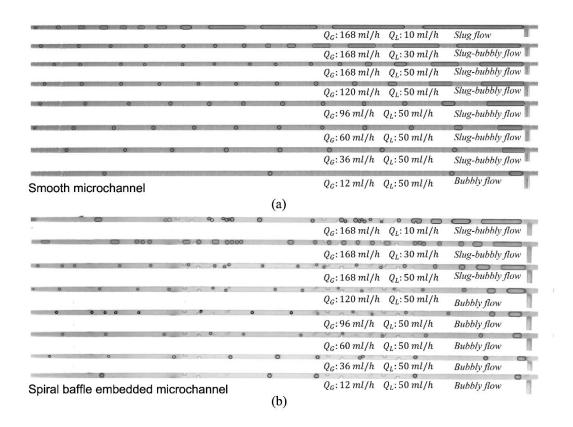


Figure 4 Typical gas-liquid two phase flow patterns for  $CO_2$  chemical absorption ( $C_{MEA}$ : 5% vol.). (a) Smooth microchannel; (b) Spiral baffle embedded microchannel

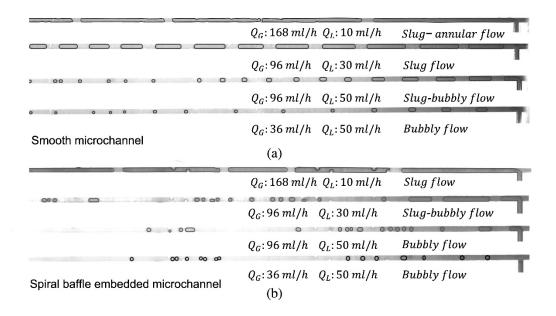


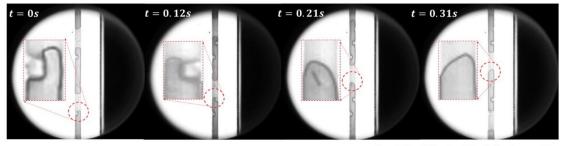
Figure 5 Typical gas-liquid two phase flow patterns for  $CO_2$  chemical absorption ( $C_{MEA}$ : 1% vol.). (a) Smooth microchannel; (b) Spiral baffle embedded microchannel

Slug flow: This flow pattern occurs at the higher  $Q_G$  and the lower  $Q_L$  in the microchannel without baffle (e.g.,  $Q_G > 100 \, ml/h$  and  $Q_L = 10 \, ml/h$ ). In this situation, the  $l_{B,ini}$  largely increases while the liquid-side mass transfer coefficient is limited. As a result, slug flow pattern persists along the microchannel. But the situation is different in spiral baffle embedded microchannel, which will be discussed later.

Slug-annular flow: At still higher  $Q_G$  and low  $Q_L$ , the liquid slug length obviously decreases, and gas slugs could merge together. In our study, this slug-annular flow pattern has only been observed in smooth microchannel with  $C_{MEA} = 1\%$  (Fig. 5) when  $Q_L$  is 10 ml/h and  $Q_G$  larger than 140 ml/h. Under the same condition, the slug flow has always been observed in the spiral baffle embedded microchannel.

Besides, the bubble deformation due to the baffle structure was observed, as shown in the images captured with 16x lens in Fig. 6. When bubble passes through the spiral baffle structure, the bubble cap is squeezed and asymmetrically deformed. The quarterly rotational distribution of baffle cubes generates the swirling pattern of the gas bubble. In the meantime, the secondary

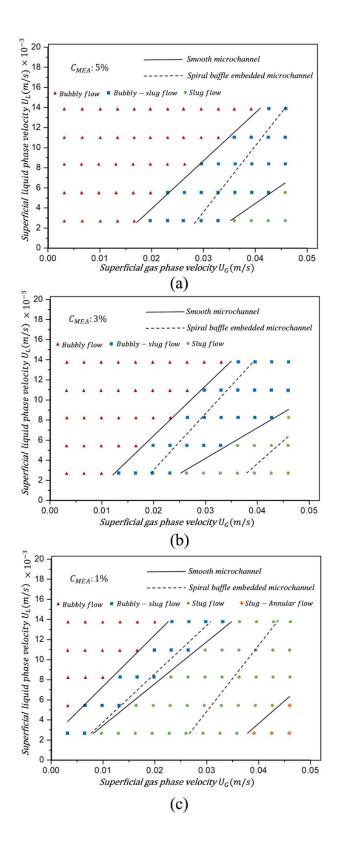
flow of the liquid phase could also be generated, especially close to the bubble caps where the mass transfer is the most intensive. In this way, the liquid side mass transfer resistance is reduced. After going through the spiral baffle structure, the length of bubble is obviously decreased owing to the enhanced mass transfer. Nevertheless, the broken slug flow pattern, as reported by Yin et al. (2019b), has not been observed under our tested conditions.



 $Q_G: 168ml/h; Q_L: 30ml/h; C_{MEA}: 1\%$ 

Figure 6 Typical bubble deformation in spiral baffle embedded microchannel captured by high-speed camera (Concitions:  $Q_G = 168 \ ml/h$ ;  $Q_L = 30 \ ml/h$ ;  $C_{MEA} = 1\%$ ).

Fig. 7 shows the two-phase flow pattern map with the  $U_G$  and  $U_L$  as coordinate for CO<sub>2</sub> chemical absorption in smooth and spiral baffle structure embedded microchannels. The effects of baffle structure and absorbent concentration on the transition of flow patterns can be clearly seen. Firstly, higher  $C_{MEA}$  accelerates the reaction rate between the CO<sub>2</sub> and the MEA solution. The bubble volume shrinks very quickly, the bubbly flow/bubbly-slug flow becoming dominant on the flow map. In contrast, at low  $C_{MEA}$  (Fig. 7c), slug-annular flow can still be present in smooth microchannel as discussed above. Secondly, the presence of baffle structure results in smaller  $l_{B,ini}$  (cf. Fig. 4 and Fig. 5) as well as the bubble deformation. As a result, the transition lines for the spiral baffle embedded microchannel are shifted to right compared to those for smooth microchannel, indicating the later transition from bubbly flow to bubbly-slug flow and from bubbly-slug flow to slug flow.



**Figure 7** Two-phase flow map for  $CO_2$  chemical absorption in smooth microchannel and in spiral baffle embedded microchannel. (a)  $C_{MEA}$ =5%; (b)  $C_{MEA}$ =3%; (c)  $C_{MEA}$ =1%.

#### 3.2 CO<sub>2</sub> gas phase void fraction

 $CO_2$  gas phase void fraction ( $\mathcal{E} = V_B/V_C$ ) is defined as the ratio of overall bubble volume ( $V_B$ ) to the whole microchannel volume ( $V_C$ ). It is an important parameter to predict the two-phase pressure drop and  $CO_2$  mass transfer performance. In this study, the  $\varepsilon$  value has been estimated by image analysis method with some assumptions on the geometry of the slug bubble as explained in section 2.2. Due to the generation and the rapid contraction of  $CO_2$  bubbles, the void fraction exhibits a repeated periodical fluctuation. As a result, the  $\varepsilon$  values reported here have been averaged within 20 s (including several periods from bubble generation to exhaust which is sufficiently long and representative). The results are shown in Fig. 8.

At a certain  $Q_L$ ,  $\varepsilon$  value increases with the increasing  $Q_G$ , mainly due to the larger  $l_{B,ini}$  as already shown in Fig. 4. Contrarily at a certain  $Q_G$ ,  $\varepsilon$  value shows an upward trend with the decreasing  $Q_L$ . This is because of the lower liquid-side mass transfer coefficient at low  $Q_L$ . By comparing Figs. 8a-c, the positive effect of  $C_{MEA}$  on the reduction of  $\varepsilon$  can be clearly seen, mainly due to the stronger chemical reaction. Moreover, under the same testing condition, the  $\varepsilon$  value of spiral baffle embedded microchannel is always smaller than that of smooth microchannel, implying the higher mass transfer rate and the better absorption performance.

Further investigating the void fraction is meaningful for comprehensively understanding the two-phase hydrodynamic in microchannel (Yin et al. 2022b). Many prediction models have then been proposed at early stage, such as by Ali et al. (1993) (Eq. 18), Kawahara et al. (2002) (Eq. 19) and Saisorn & Wongwises (2009) (Eq. 20).

$$\varepsilon = 0.8\beta \tag{18}$$

$$\varepsilon = \frac{0.03\beta^{0.5}}{1 - 0.97\beta^{0.5}} \tag{19}$$

$$\varepsilon = \frac{0.036\beta^{0.5}}{1 - 0.945\beta^{0.5}} \tag{20}$$

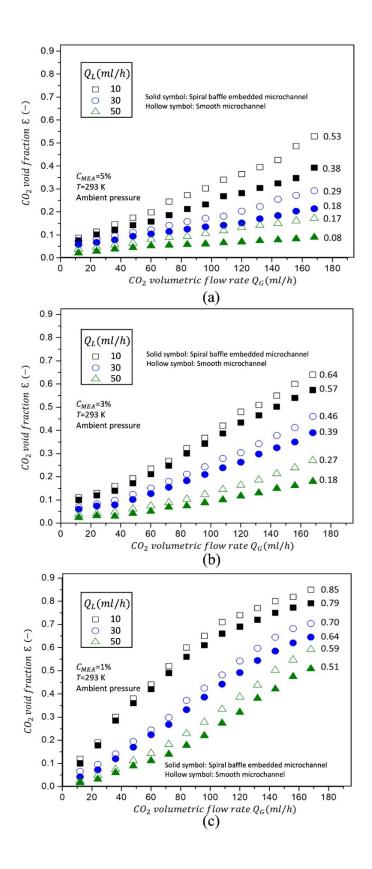


Figure 8 CO<sub>2</sub> gas phase void fraction ( $\varepsilon$ ) in smooth and spiral baffle embedded microchannels. (a)  $C_{MEA} = 5\%$ ; (b)  $C_{MEA} = 3\%$ ; (c)  $C_{MEA} = 1\%$ .

Where  $\beta = Q_G/(Q_G + Q_L)$ . However, these models are only based on the gas and liquid flow rate ratio, thereby suited for two-phase flow systems without the chemical reaction (or when its effect could be ignored). Yin et al. (2018) proposed a modified model by adding the dimensionless Hatta number (Ha):

$$\varepsilon = \frac{0.8\beta}{1 + 5.3 \times 10^{-3} Ha} \tag{21}$$

$$Ha = \sqrt{Dk_{ov}/k_{L,P}^2} \tag{22}$$

Where D  $(m^2/s)$  is the two-phase diffusion coefficient,  $k_{ov}$  (1/s) is the overall chemical reaction rate,  $k_{L,P}$  (m/s) is the liquid side physical mass transfer rate estimated by renewable mass transfer model.

Fig. 8 compares the experimental data of gas void fraction obtained in this study with the above-mentioned prediction models. A clear deviation can be seen. In fact, Ali's model is suitable for mini-scale flat plates channel reactor with stable two-phase split flow. The bubble volume decrease rate is small because of the limited mass transfer rate; thus the linear prediction correlation is proposed for this situation. Kawahara's model and Saisorn's model are proposed under the large gas and liquid superficial velocities  $(U_G: 0.1-60\ m/s,\ U_L: 0.01-4\ m/s)$ which lead to totally different flow patterns. For example, nearly single-phase flow was observed at high liquid flow rate for all gas flow rate, and for  $\beta$  < 0.5, the  $\varepsilon$  value remained smaller than 0.1. Thus, they cannot be used to predict the stable split flow (such as slug flow and bubbly flow) under small gas and liquid Re numbers tested in the present study. For correlation proposed by Yin et al. (2018), though the chemical reaction greatly promotes the two-phase mass transfer performance, the main channel length of their study is rather short (40 mm), resulting in the overestimation of  $\varepsilon$ . In this experiment with a longer microchannel accompanied with quick chemical reaction, the bubble volume change is realized within a shorter distance. For this reason, the current experimental  $\varepsilon$  data are generally smaller than the prediction values using Yin's model and shows a non-linear trend.

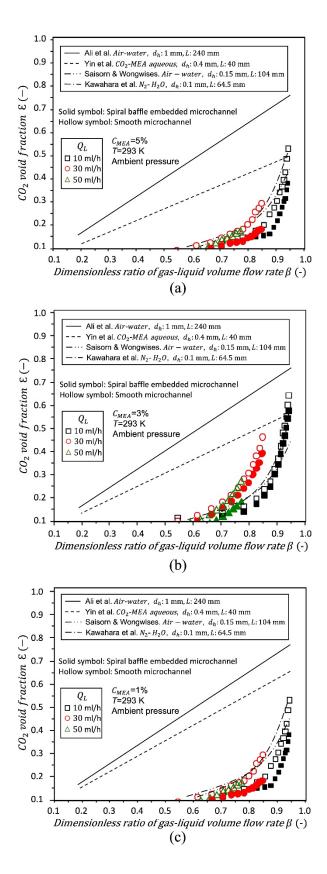


Figure 9 Comparison of  $CO_2$  void fraction between the experimental data and calculated values by prediction models. (a)  $C_{MEA}=5\%$ ; (b)  $C_{MEA}=3\%$ ; (c)  $C_{MEA}=1\%$ 

Given the need of a new model for more accurate prediction of  $\varepsilon$  for CO<sub>2</sub> chemical absorption in microchannel at small gas and liquid Re numbers, the dimensionless Damköhler number (Da) is proposed to be included in the correlation to consider the effect of chemical reaction. The Da number is defined as the ratio of chemical reaction rate  $(k_{ov})$  to the diffusion rate (D), expressed as follows (Ma et al., 2020):

$$Da = k_{ov} \cdot d_h^2 / D \tag{23}$$

For smooth microchannel, the new correlation could be written as follows:

$$\varepsilon_{smooth} = \frac{a\beta^b}{1 + c \cdot Da} \tag{24}$$

And based on the experimental data obtained in this study, the fitting factor a is 1.42, b is 3.22 and c is  $1.93 \times 10^{-6}$ .

For structured microchannel, the structure intensify factor  $(\Psi)$ ,  $Re_G$  and  $Re_L$  should also be considered to indicate the effects of flow disturbance structure on the CO<sub>2</sub> void fraction, written as:

$$\varepsilon_{structured} = (dRe_G^e Re_L^f)^{\Psi - 1} \varepsilon_{smooth}$$
 (25)

Where the fitting constant is 0.006 for d, 2.856 for e and 0.17 for f, respectively, based on the results of this study.  $\Psi$  value for the spiral baffle structure is decided by the micro baffle dimension, the number and the distribution mode. It is estimated to be equal to 1.1 in this research while  $\Psi = 1$  for the straight smooth microchannel case. Fig. 10 shows that the prediction values of  $\varepsilon$  using the new model are in good agreement with the experimental data, the maximum absolute error being smaller than  $\pm 25\%$ .

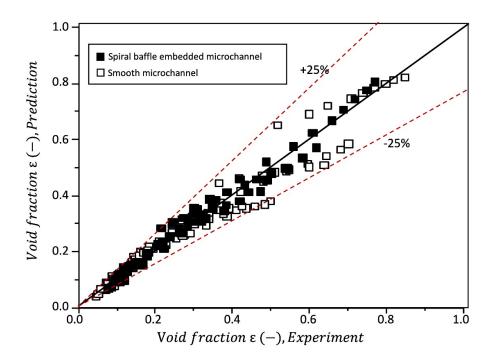
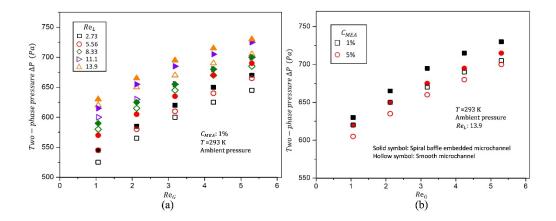


Figure 10 Predicted values of CO<sub>2</sub> void fraction compared with experimental data.

## 3.3 Pressure drop

Figure 11 shows the measured two-phase pressure drop ( $\Delta P$ ) for CO<sub>2</sub> chemical absorption by MEA solution in microchannel. From Fig. 11a, it can be seen that  $\Delta P$  increases with the increasing  $Re_G$  and  $Re_L$ . The presence of the flow disturbance structure (spiral baffle) would cause an additional increase of  $\Delta P$ , and this increase is more significant at high  $Re_G$ . The highest  $\Delta P$  recorded under the tested conditions is about 750 Pa and the pressure drop increase due to the embedded spiral baffle is no more than 50 Pa, which is less than 10% and may be considered as acceptable in various industrial applications.

Figure 11b shows the effect of MEA solution concentration on the  $\Delta P$ , indicating that the  $\Delta P$  decreases with the increasing  $C_{MEA}$  under the same  $Re_L$  and  $Re_G$  condition. Given the almost constant physical properties of MEA solution within the tested concentration range (cf. Table 2), this difference in  $\Delta P$  is mainly due to the different CO<sub>2</sub> absorption rate at different  $C_{MEA}$ . In fact, the lower  $\varepsilon$  value at high  $C_{MEA}$  would reduce the pressure drop caused by the bubble moving.



**Figure 11** The measured two-phase pressure drop for CO<sub>2</sub> chemical absorption by MEA solution in smooth and spiral baffle embedded microchannels. (a) Influence of gas and liquid *Re* numbers; (b) Influence of MEA concentration.

# 3.4 Mass transfer performance

The effects of operating conditions (gas and liquid flow rates and MEA concentration) and spiral baffle structure on the mass transfer performance of  $CO_2$  chemical absorption in microchannel are discussed here. Fig. 12 show the variation of  $k_L a$  values under the tested conditions for smooth and spiral baffle embedded microchannels. The experimental results clearly show the dependence of  $k_L a$  values on  $Re_G$  and  $Re_L$ , i.e., higher gas and liquid flow rates (or superficial velocities) enhance the mass transfer. Firstly, increasing the  $Re_G$  augments the mass of  $CO_2$  on one hand, and promotes the transition from bubbly flow to slug flow on the other hand, which greatly increases the two-phase mass transfer surface area. Both effects would augment the  $k_L a$ . Moreover, increasing gas flow rate will shorten the liquid slug, which directly leads to an increase in shear force between liquid and solid, thereby promoting the liquid flow mixing at the meniscus position. Secondly, according to the surface renewal theory, the increased liquid superficial velocity could directly enhance the renewal rate of the absorbent around the  $CO_2$  bubbles, thereby reducing the liquid side mass transfer resistance.

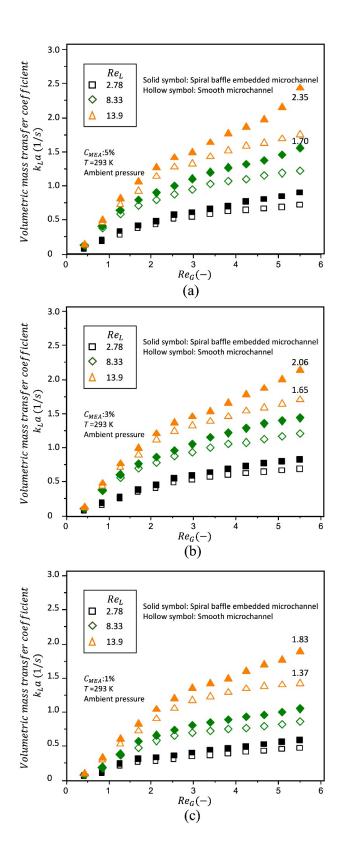
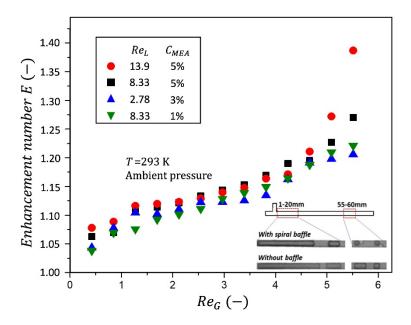


Figure 12 Effects of  $Re_G$ ,  $Re_L$ , absorbent concentration and structure factor on liquid side volumetric mass transfer rate  $k_L a$ . (a)  $C_{MEA} = 5\%$ ; (b)  $C_{MEA} = 3\%$ ; (c)  $C_{MEA} = 1\%$ 

In addition, the increasing  $C_{MEA}$  could directly increase the chemical reaction rate and decreasing the CO<sub>2</sub> molar concentration in the liquid phase correspondingly. The mass transfer process driven by the concentration difference is significantly enhanced and  $k_L a$  obviously increased, especially at high  $Re_G$ . For smaller  $Re_G$ , the effect of  $C_{MEA}$  on the increase of  $k_L a$  is less significant because of the decreased gas—liquid interfacial area for mass transfer.



**Figure 13** Enhancement number *E* achieved by spiral baffle structure.

Finally, the effect of spiral baffle structure on the enhancement of mass transfer is rather significant, indicated by the high values of E shown in Fig. 13. In smooth microchannel, the mass transfer is dominated by diffusion and the main convective mass transfer is happening at liquid phase meniscus position. The existence of spiral baffle structure could disturb the gasliquid mass transfer interface and promote its renewal, thereby strengthening the  $CO_2$  gas dissolve process. Besides, the spiral baffle structure could induce more spatially asymmetric vortices in liquid phase and this enhanced liquid mixing is more intensive at higher  $Re_L$  number, hence improving the chemical reaction rate and convective mass transfer. Both these two factors enhance the two-phase mass transfer performance. However, the enhancement achieved by the

spiral baffle structure varies with gas and liquid Re numbers. At low  $Re_L$  and  $Re_G$  numbers, the enhancement is relatively small because of the relatively weak disturbance when  $CO_2$  bubbles pass through the spiral baffle structure. With the continuous increase of  $Re_L$  and  $Re_G$  numbers,  $k_La$  and E values are significantly augmented and the maximum value of  $k_La=2.35\ s^{-1}$  and E=1.387 can be achieved. Such  $k_La$  values are 2-3 orders of magnitude higher than that of traditional devices (e.g., trickle bed reactor or packed column) at such small gas and liquid Re numbers ( $Re_G < 5, Re_L < 15$ ) (Table 4), highlighting the interests of microchannel-based  $CO_2$  chemical absorption devices. Besides, the novel spiral baffle structure also shows its effectiveness on the enhancement of mass transfer, especially for operations at small gas and liquid Re numbers.

**Table 4** Comparison of  $k_L a$  and E for different CO<sub>2</sub> chemical absorption devices within the given gas and liquid Re number range  $(Re_G < 5, Re_L < 15)$ .

Туре	$k_L a (s^{-1})$	E	Ref	
Trickle bed reactor	0.001-0.01		Larachi et al. (1998)	
Packed column	0.001-0.05	-	Zeng et al. (2011)	
Heart shape structured microreactor	0.53	1.12	Chen et al. (2021)	
Rectangular baffles structured microreactor	2.1	1.27	Yin et al. (2020)	
Spiral baffle embedded microchannel (This work)	Up to 2.35	Up to 1.387		

Plenty of semi-empirical correlations have been proposed for the prediction of mass transfer performance in mini/micro channel-based CO<sub>2</sub> absorber. For example, the correlation of Yue et al. (2007b) for slug flow (Eq. 26) and for slug-annular flow (Eq. 27), and the correlation of Ji et al. (2010) for slug flow (Eq. 28) are shown as below

$$Sh_L \cdot a \cdot d_h = 0.084 Re_G^{0.213} \cdot Re_L^{0.937} \cdot Sc_L^{0.5}$$
 (26)

$$Sh_L \cdot a \cdot d_h = 0.058 Re_G^{0.344} \cdot Re_L^{0.912} \cdot Sc_L^{0.5}$$
 (27)

$$Sh_L \cdot a \cdot d_h = 0.22 Re_G^{0.78} \cdot Re_L^{0.0535} \cdot Sc_L \cdot Ca_L^{0.7568}$$
 (28)

Where Sh is the Sherwood number  $(Sh_L = k_L \cdot d_h/D)$  and Sc is the Schmidt number  $(Sc_L = \mu_L/(D \cdot \rho_L)$ 

Nevertheless, these correlations were developped for  $CO_2$ – $H_2O$  system, thus can better predict the  $CO_2$  physical absorption. For  $CO_2$  chemical absorption (MEA solution) in microchannel, empirical correlations have also been proposed by Li et al. (2014) (Eqs. 29, 30) for the operation condition (26.73 <  $Re_G$  < 467.79 and 5.35 <  $Re_L$  < 37.42) and by Yin et al. (2019a) (Eq. 31) for the operation condition (1.29 <  $Re_G$  < 9.44 and 11.68 <  $Re_L$  < 31.68):

$$Sh_L \cdot a \cdot d_h = 5.0107 Re_G^{2.9423} \cdot Re_L^{-1.5583} \cdot Sc_L^{0.5} \cdot Ca_L^{0.0743} \qquad \frac{Re_G}{Re_L} < 1.7319$$
 (29)

$$Sh_L \cdot a \cdot d_h = 19.389 Re_G^{0.3872} \cdot Re_L^{0.8865} \cdot Sc_L^{0.5} \cdot Ca_L^{0.4176} \qquad \frac{Re_G}{Re_L} \ge 1.7319$$
 (30)

$$Sh_L \cdot a \cdot d_h = 0.81 \cdot Re_G^{0.78} \cdot Re_L^{0.41} \cdot Da^{0.35}$$
 (31)

Specifically for intensified microchannels, the enhancement number E has been proposed to be added into the correlation Chu et al. (2019):

$$Sh_L \cdot a \cdot d_h = 1.3425 \cdot Re_G^{1.0968} \cdot Re_L^{-0.4705} \cdot Sc_L^{0.5} \cdot Ca_L^{-0.0302} \cdot E^{0.7198}$$
 (32)

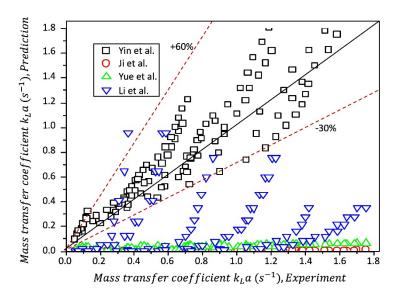


Figure 14 Comparison of  $k_L a$  value between the experimental data in smooth microchannel and prediction values by related correlations.

Fig. 14 shows a comparison of experimental data for smooth microchannel obtained in this study with the prediction results using the above-mentioned correlations. Large deviations can be seen for correlations based only on physical parameters. For Yin et al. (2019a)'s correlation

that includes the Da number, the prediction error could be significantly reduced, but could still reach -30% to 60%. This is because Yin et al. (2019a)'s correlation has been proposed for  $Re_L$  number range (11.68 – 31.68) and  $Re_G$  number range (1.29-9.44), two times larger than that of this study. Besides, the above mass transfer prediction model did not consider the effect of intensification micro-structure on the mass transfer performance. As a result, proper mass transfer prediction correlations that consider chemical reaction and intensification structure for small  $Re_G$  and  $Re_L$  numbers range are still lacking.

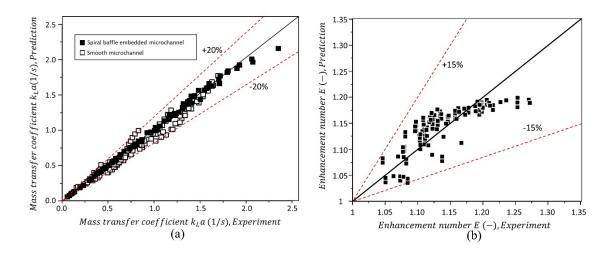
Therefore, new correlations are proposed based on the experimental data obtained in this study. In particular, the Da number and the E number have been included to consider the effect of chemical reaction and intensification structure on the enhancement of mass transfer, shown in Eq. 33 and Eq. 34.

$$Sh_L \cdot a \cdot d_h = 6.45 \cdot Re_G^{0.6234} \cdot Re_L^{0.628} \cdot Da^{0.16} \cdot E$$
 (33)

$$E = (Re_G^{0.512}Re_L^{-0.026}Da^{0.07})^{\Psi-1}$$
(34)

From these correlations, it can be seen that for the smooth channel, both  $Re_G$  and  $Re_L$  have almost the equal effect on the improvement of mass transfer. The positive impact of solvent concentration (Da) can also be well recognized. For the structured microchannel, the enhancement is much more sensitive to  $Re_G$  than others under our tested conditions.

Fig. 15 shows that the experimental data and the prediction results of the new model are in good agreement: the relative error is below  $\pm 15\%$  for E number and below  $\pm 20\%$  for  $k_L a$ , respectively.



**Figure 15** Comparison of experimental data with new prediction model. (a) Liquid side volumetric mass transfer rate  $k_L a$ . (b) Mass transfer enhancement number E.

## 3.5 Effective energy efficiency

Undoubtedly, adding micro intensification structures will increase the two-phase pressure drop, but it also enhances the two-phase mass transfer performance. In order to further estimate the mass transfer enhancement efficacy, the parameter effective energy efficiency ( $\eta$ ) is introduced here for analysis (Li et al., 2019).  $\eta$  is defined as the ratio of mass transfer enhancement number (E) to the energy consumption ratio, expressed in Eq.35. Ec is micro CO<sub>2</sub> absorber energy consumption and calculated as Eq.36 (Chen et al., 2021):

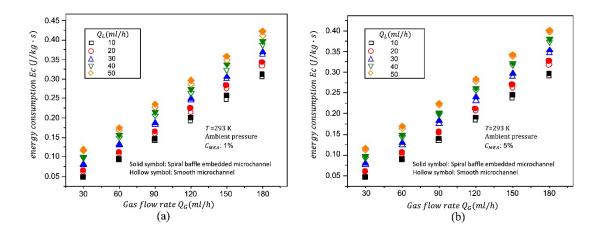
$$\eta = \frac{E}{Ec_{structured}/Ec_{smooth}}$$
 (35)

$$Ec = \frac{\Delta P \cdot (Q_G + Q_L)}{V_C \rho_L} \tag{36}$$

Where  $\Delta P$  is two-phase pressure drop,  $\rho_L$  is the liquid phase density and estimated here as constant. Fig.16a shows the effect of gas or liquid flow rate and absorbent concentration on energy consumption for two tested microchannels. It could be found that the Ec values increase with the increasing gas-liquid flow rates and the largest Ec is no more than 0.5 W / kg within our tested conditions. Compared with smooth microchannel, the Ec slightly increased. It is

worth noting that the impact of absorbent concentration on the energy consumption is far less significant than that of gas or liquid flow rate.

The influence of gas or liquid flow rate and absorbent concentration on the effective energy efficiency  $\eta$  is shown in Fig.17. It can be seen that the value of  $\eta$  clearly increases with the gas flow rate, owing to the quickly increased E.



**Figure 16** Effect of gas or liquid flow rate and absorbent concentration on the energy consumption Ec. (a)  $C_{MEA} = 1\%$ ; (b)  $C_{MEA} = 5\%$ .

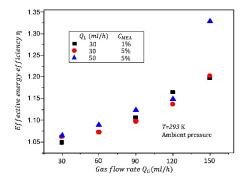


Figure 17 Effect of gas or liquid flow rate and absorbent concentration on the effective energy efficiency  $\eta$ .

# 4. Conclusion

In this study, CO<sub>2</sub>-MEA mass transfer performance in microchannel-based absorber has been experimentally evaluated. In particular, a novel spiral baffle structure has been embedded in the

straight microchannel to enhance the gas-liquid mass transfer. Results show that the existence of spiral baffle structure could induce the CO<sub>2</sub> bubble deformation, thereby disturbing the gasliquid mass transfer interface and promoting its renewal. Besides, more spatially asymmetric vortices in liquid phase could be formed to enhance the liquid mixing, hence improving the chemical reaction rate and convective mass transfer. With the increase of gas and liquid flow rates, the induced vortex intensity could also be augmented. Compared with the straight smooth microchannel, better CO<sub>2</sub> absorption performance can be achieved owing to the spiral baffle structure, with  $k_L a$  value up to 2.35  $s^{-1}$  and enhancement number E up to 1.387 under the tested conditions. This mass transfer enhancement achieved in small gas and liquid Re number range ( $Re_G < 5.5$  and  $Re_L < 14$ ) is particularly interesting because of the small pressure drop increase (<10%) compared to the smooth microchannel. The results show that largest energy consumption (Ec) is no more than 0.5 W / kg and effective energy efficiency ( $\eta$ ) gradually increased within the gas-liquid flow operating range. Based on the obtained experimental data, new models and correlations have been proposed for predicting the gas void fraction ( $\varepsilon$ ), the enhancement number (E) and the liquid-side mass transfer coefficient  $(k_L a)$  at small gas and liquid Re numbers. The effects of chemical reaction (solvent concentration) and intensification structure have been taken into account by including the Da number and  $\Psi$ factor into the models/correlations. Good agreement has been achieved between the experimental data and prediction results, with the maximum absolute error smaller than 25% under the tested conditions.

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