International Conference on Numerical Methods in Multiphase Flows

Algebric momentum preserving method with a fully coupled 3D parallel solver for the simulation of two-phase incompressible flows

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# Introduction

# Multi-scale flows





Impact of a liquid on a surface [http://physique.unice.fr/sem6/ 2012-2013/PagesWeb/PT/Splash/Html/Index.html]

Phase inversion between two immiscible liquids (to be submitted to JCP 2022)

# Introduction

# Multi-scale flows



Breaking wave (LEGI Laboratory) [images.cnrs.fr/en/photo/20150001\_0066]



Dam-Break flow in the presence of obstacle (to be submitted to JCP 2022)



# Challenges



- Presence of different scales in space and time.
- Complex interactions between capillary, inertial and viscous forces.
- Simulations may be unstable and velocities may be wrongly evaluated when high density/viscosity ratios are involved.
- Treatment of velocity-pressure coupling in incompressible flow
- Numerical difficulties with **ill-conditioned linear systems**.

### Objectives

In the framework of the homemade code Fugu, develop a parallel and robust solver dealing with two-phase flows, with high density and viscosity ratios.





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# Numerical Platform





- Implicit finite volumes discretized on staggered meshes.
- Three pressure-velocity coupling methods are considered, a **fully-coupled solver** published in El Ouafa et al. CICP2021, an augmented Lagrangian approach of Vincent el al. JCP2011 and a standard Goda JCP1978 projection method.
- The solver uses a **BiCGStabll solver** of Vorst et al. IAM1998, preconditionned with **MG** (HYPRElibrary) and home made preconditionner for fully-coupled approach (which also uses HYPRElibrary).
- Centered schemes are utilized for all spacial derivatives exept inertial term.
- The non-linear terms in the momentum equation are discretised with a **consistent mass-momentum transport** and linearized with an Adams-Bashforth scheme.
- The VOF method used is the one of Yue and Weymooth JCP2011.
- CSF models for surface tension with VOF smoothing Pianet et al. IJMF2010.
- MPI based parallelization.



# 2: Numerical methods

- Incompressible two phase flows with interface tracking
- Algebraic momentum preserving method
- Pressure-velocity coupling algorithms
- Fully-coupled solver and block preconditionning



Incompressible two phase flows with interface tracking (Kataoka IJMF 1986)

$$\nabla \cdot \mathbf{u} = 0 \tag{1}$$

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u}) + \overline{\overline{B}}_{penu} \left( f(\mathbf{u}) - \mathbf{u}_{\infty} \right) = -\nabla p + \nabla \cdot \overline{\overline{T}} + \underline{\sigma} \kappa \delta \mathbf{n} + \rho \mathbf{g}$$
Continuum Surface Force Brackbill et al. JCP1992

 $B_{penu}$  for specifying BC Angot et al. NM 1999

Temporal evolution of volume fraction given by

$$\frac{\partial C}{\partial t} + \mathbf{u} \cdot \nabla C = 0 \tag{3}$$

Recasting in the "conservative form Weymouth et al. JCP2011 " as

$$\frac{\partial C}{\partial t} + \nabla \cdot (\mathbf{u}C) = C\nabla \cdot \mathbf{u} \tag{4}$$

Material properties defined on basis of interface position



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(2)

## Algebraic momentum preserving method

The Navier-Stokes equations (conservative form) can be recast in operator notation

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u}) = -\nabla \rho + L \left(\overline{\overline{B}}_{\rho enu}, \overline{\overline{T}}, \sigma, \mathbf{g}\right)$$

$$L = L \left(\overline{\overline{B}}_{\rho enu}\right) + L \left(\overline{\overline{T}}\right) + L_{\sigma} + L_{g}$$
(7)

Operator-split framework

$$\frac{\rho^{**}\mathbf{u}^{**} - \rho^n \mathbf{u}^n}{\Delta t} + \nabla \cdot \left(\rho^{*,n+1}\mathbf{u}^{*,n+1} \otimes \mathbf{u}^{*,n+1}\right) = 0$$
(8)

$$\frac{\rho^{n+1}\mathbf{u}^{n+1} - \rho^{**}\mathbf{u}^{**}}{\Delta t} = \nabla \rho + L\left(\overline{\overline{B}}_{penu}, \overline{\overline{T}}, \sigma, \mathbf{g}\right)$$
(9)

•  $\rho^{n+1}$  is obtained from the resolution of

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0 \tag{10}$$

strong stability preserving Runge-Kutta 3 (SPP-RK3) time integrator + WENO or CUI Cubista TVD schemes

Rudmannn (IJNM 1998), Desjardins Moureau (CTR SummerProgram 2010), Rhaessi, Pitsch, (CaF 2012), V. LeChenadec, H. Pitsch, (JCP 2013), Ghods, Herrmann, (PS 2013), Vaudor et al. (CaF2016) Owkes, Desjardins (JCP 2017), Nangia et al JCP (2019), Estivalezes et al JCP (2021), Arrufat et al C&F (2021)

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Pressure-velocity coupling algorithms

Find (u, p) such as the conservation equations are satisfied

#### Segregated approach

- Chorin type projection (MC1968) and its variants
- Kinematics Scalar Projection Method (KSP) (Caltagirone et al. MFD2015)
- Fractional time-step technique (Guermond et al. JCP2009
- SIMPLE, SIMPLEC, PISO..... (Patankar et al. 1983)



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Fully-coupled solver and block preconditionning (El Ouafa et al. CICP2021)

Suitable temporal and spatial discretization leads to:

$$\begin{pmatrix} M_{u}^{(\rho^{n+1})} + A_{u}^{(\mu^{n+1})} & B_{\rho}^{T} \\ B_{u} & 0 \end{pmatrix} \begin{pmatrix} \mathbf{u}_{\rho^{n+1}}^{n+1} \end{pmatrix} = \begin{pmatrix} \mathbf{f} \\ \mathbf{0} \end{pmatrix} \text{ or }$$

$$M_{u}^{(\rho^{n+1})} \mathbf{u}^{n+1} = \int_{\Omega} \left[ \frac{\rho^{n+1}}{\Delta t} \mathbf{u}^{n+1} + \overline{B}_{penu} f(\mathbf{u}^{n+1}) \right] dV$$

$$A_{u}^{(\mu)} \mathbf{u}^{n+1} = -\int_{\Sigma} \left[ \mu^{n+1} \left( \nabla \mathbf{u}^{n+1} + \nabla^{T} \mathbf{u}^{n+1} \right) \right] \cdot \mathbf{n} dS$$

$$B_{\rho}^{T} \rho^{n+1} = \int_{\Omega} \nabla \rho^{n+1} dV$$

$$B_{u} \mathbf{u}^{n+1} = \int_{\Omega} \nabla \cdot \mathbf{u}^{n+1} dV$$

$$\mathbf{f} = \int_{\Omega} \left( \frac{\rho^{**} \mathbf{u}^{**}}{\Delta t} + \overline{B}_{penu} \mathbf{u}_{\infty} + \mathbf{F}_{s} \right) dV$$



Ax = b

Iterative solver: BICGStabII (Vorst et al. IAM1987)  
Preconditionner  
Instead of solving, 
$$Ax = b$$
 we prefer to solve  $P^{-1}Ax = P^{-1}b$  for which  $cond(P^{-1}A) \ll cond(A)$ .  
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Block preconditionning techniques

build the preconditioner P of the problem by taking a block LU decomposition of the original matrix A

$$\begin{pmatrix} F_u & B_p^T \\ B_u & 0 \end{pmatrix} = \begin{pmatrix} I_u & 0 \\ B_u F_u^T & I_p \end{pmatrix} \begin{pmatrix} F_u & B_p^T \\ 0 & S_p \end{pmatrix} \text{ with } F_u = M_u + A_u$$

Here  $s_{\rho}$  is the Schur complement of the pressure block:  $s_{\rho} = -B_{u}F_{u}^{-1}B_{\rho}^{T}$ Block triangular preconditioner:  $P \approx \begin{pmatrix} F_{u} & B_{\rho}^{T} \\ 0 & S_{\rho} \end{pmatrix}$ not feasible to use the exact Schur complement  $S_{\rho}$  and the velocity block  $F_{u}$  as a part of the preconditionning operator [[[[[]]]]]

#### Block preconditionning techniques for the velocity block (El Ouafa et al. CICP2021)

using the Neumann series expansion at the order zero, the form of our velocity preconditioner reads as follows:

$$\hat{F}_{u} = \begin{pmatrix} F_{uu} & F_{uv} & F_{uw} \\ 0 & F_{vv} & F_{vw} \\ 0 & 0 & F_{ww} \end{pmatrix}$$
(11)

Block preconditionning techniques for the inverse of Schur complement

 $\hat{S}_p^{-1}$  by the pressure convection diffusion (PCD) preconditioning (Bootland et al. JCP 2017)

$$\hat{S}_{PCD}^{-1} = (M_{\rho}^{1/\mu})^{-1} + (A_{\rho}^{1/\rho})^{-1} (N_{\rho}^{(1)} + \frac{1}{\Delta t} M_{\rho}^{(1)}) (M_{\rho}^{(1)})^{-1}$$
(12)

 $\hat{P}^{-1}Ax = \hat{P}^{-1}b$  —> each matrix-vector product Ax = z will be followed by the calculation of  $\hat{P}^{-1}z = y$ 

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# 3: Solver validations

• Three dimensional test cases







Hight density droplet advection without viscous, gravity and capillary effect. Comparison between Algbraic momentum preserving and centred scheme with VOF on the mesh 64x64x64. The interface is ploted at time 0s (first column) 0.25s (seconde column) 0.5s (third column) and 1s (fourth column).



Simulation of the inversion of the phase for Re=221 and We=121(to be submitted to JCP 2022)

• 
$$Re = \frac{\rho_1 H U_g}{\mu_1}$$
 and  $We = \frac{\rho_1 H U_g^2}{\sigma}$  with  $U_g = \frac{\rho_2 - \rho_1}{\rho_1} \sqrt{\frac{gH}{2}}$ 

• HPC DNS calculations  $\approx 10^9$  cells – 12800 processors of the HPC Occigen cluster of CINES







0s



7s









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1.5s

**5s** CNMMF-4 <u>202</u>

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#### Comparaison of macroscopic quantities of interest with the reference data of Saeedipour et al. AM 2021

۰ HPC DNS calculations  $\approx 10^9$  cells – 12800 processors of the HPC Occigen cluster of CINES



Dam-Break Flow in the Presence of Obstacle  $\rho_2/\rho_1 = 10^3$  and  $\mu_2/\mu_1 = 10^2$  (to be submitted to JCP 2022)

- Domain D=[ $3.22 \times 1 \times 1$ ]
- mesh 200x80x80





Water elevation



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0.8747s

3.4585s

#### Strong scalability

#### Goal

-Find the point at which a reasonable computation time is achieved but still limit the overhead induced by a parallel process (Skylake and Jean Zay for a number of processors increasing from 160 to 12800 over 30 time iterations on 134 million cells ).

Speed-up and Efficiency

$$S(p) = \frac{T_{ref}}{T_{p}} N proc_{ref}$$
(13)

$$E(p) = \frac{S(p)}{p} \tag{14}$$



Conclusions and perspectives:

- Fully-coupled solver 2D/3D
- High density and viscosity ratios
- Simulations massively parellel
- Momentum preserving
- Extention to LES of turbulent two-phase flows and compressible flows



# Thank you for your attention

