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# An exterior calculus framework for polytopal methods

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## Abstract

We develop in this work the first polytopal complexes of differential forms. These complexes, inspired by the Discrete De Rham and the Virtual Element approaches, are discrete versions of the de Rham complex of differential forms built on meshes made of general polytopal elements. Both constructions benefit from the high-level approach of polytopal methods, which leads, on certain meshes, to leaner constructions than the finite element method. We establish commutation properties between the interpolators and the discrete and continuous exterior derivatives, prove key polynomial consistency results for the complexes, and show that their cohomologies are isomorphic to the cohomology of the continuous de Rham complex.

**Key words.** Discrete de Rham complex, Virtual Element Method, differential forms, exterior calculus, polytopal methods

**MSC2020.** 65N30, 65N99, 14F40

## 1 Introduction

This work is a first step towards merging two extremely successful avenues of research in numerical analysis: finite element differential forms and arbitrary-order polytopal methods.

The well-posedness of important classes of partial differential equations (PDEs), and the development of stable approximations thereof, hinges on the properties of underlying Hilbert complexes [23]. The best-known example is provided by the de Rham complex which, for an open connected polyhedral domain  $\Omega \subset \mathbb{R}^3$ , reads

$$\{0\} \longrightarrow H^1(\Omega) \xrightarrow{\text{grad}} \mathbf{H}(\mathbf{curl}; \Omega) \xrightarrow{\text{curl}} \mathbf{H}(\text{div}; \Omega) \xrightarrow{\text{div}} L^2(\Omega) \longrightarrow \{0\}, \quad (1.1)$$

where  $H^1(\Omega)$  is the space of scalar-valued functions over  $\Omega$  that are square-integrable along with their gradient, while  $\mathbf{H}(\mathbf{curl}; \Omega)$  and  $\mathbf{H}(\text{div}; \Omega)$  are the spaces of vector-valued functions over  $\Omega$  that are square-integrable along with their curl and divergence, respectively. Using the framework of differential

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forms (see Appendix A), the de Rham complex (1.1) can be generalised to a domain  $\Omega$  of any dimension  $n$  as:

$$\{0\} \longrightarrow H\Lambda^0(\Omega) \xrightarrow{d^0} \dots \xrightarrow{d^{k-1}} H\Lambda^k(\Omega) \xrightarrow{d^k} \dots \xrightarrow{d^{n-1}} H\Lambda^n(\Omega) \longrightarrow \{0\}. \quad (1.2)$$

In what follows, we shall possibly omit the index  $k$  from exterior derivatives and spaces in (1.2) when no ambiguity can arise.

The de Rham complex enters the well-posedness analysis of PDEs through its cohomology spaces  $\text{Ker } d^k / \text{Im } d^{k-1}$ . A classical result links these spaces to the topological features of the domain and their dimensions to its Betti numbers. Preserving such homological structures at the discrete level leads to *compatible* methods and is key to the design of stable numerical schemes.

The compatible finite element approximation of the vector-valued spaces appearing in the de Rham complex (1.1) arose as a research subject in the late 70s [53, 54]. In the late 80s, links with Whitney forms were identified [17]. More recently, the development of Finite Element Exterior Calculus (FEEC) [2, 4, 5] has provided a unified perspective on the generation and analysis of finite element approximations of the de Rham complex (1.2). Finite Element Systems (FES) are a generalisation of FEEC covering spaces which are not necessarily piecewise polynomial inside mesh elements (but can be, for example, piecewise polynomial on subdivisions of these elements); see [29–31]. FEEC and FES led to the unification of several families of finite elements and heavily hinge on the notion of subcomplex, which makes them naturally geared towards conforming approximations.

While conforming methods are still widely used, their construction can only be carried out on conforming meshes, typically made of elements with simple shapes (e.g., tetrahedra or hexahedra); extensions to more general meshes, such as the barycentric dual of a simplicial mesh, have been considered, e.g., in [26]. In recent years, significant efforts have been made to develop and analyse numerical methods that support more general meshes including, e.g., general polytopal elements and non-matching interfaces; a representative but by no means exhaustive list of contributions includes [1, 6, 7, 13, 16, 18, 21, 22, 24, 33, 36, 38–42, 45, 46, 48]. Polytopal technologies typically introduce some degree of non-conformity, either because they are formulated in a fully discrete setting (like Hybrid High-Order [36, 42] or Discrete de Rham – DDR methods [33, 38]) or through the use of projections (as in Virtual Element Methods – VEM [7]).

Despite their non-conformity, polytopal technologies can be used to develop compatible frameworks. Polytopal discretisations of the de Rham complex (1.1) have been proposed, e.g., in [10, 33, 38], and applied to a variety of models, such as magnetostatics [8, 34], the Stokes equations [11], and the Yang–Mills equations [47]; they have also inspired further developments, based on the same principles, for other complexes of interest such as variants of the de Rham complex with increased regularity [32, 55], elasticity complexes [19, 44], and the Stokes complex [12, 14, 49]. Polytopal complexes have additionally been used to construct methods that are robust with respect to the variations of physical parameters, in particular for the Stokes problem [11], for the Reissner–Mindlin equation [43], or the Brinkman model [33]. Many of these models have also been tackled using finite element complexes and related methods (see, e.g., [2, 3, 27, 31]). However, due to their higher-level design, which does not require explicit expressions for the basis functions, polytopal methods offer distinctive advantages over finite elements. These include, in addition to the support of general meshes, the possibility to reduce the dimension of discrete spaces, sometimes below their finite element counterparts [35, Table 3], through systematic processes such as enhancement or serendipity [9, 35]

The purpose of the present work is to take one step further and show how exterior calculus can be used to generalise the construction and analysis of polytopal complexes. More specifically, we present two discrete de Rham complexes in arbitrary dimension and with arbitrary approximation degree that generalise those introduced in [33] (DDR) and [8] (VEM). Three key features set these constructions apart from Finite Element complexes:

- No explicit spaces of globally conforming differential forms (i.e., subspaces of  $H\Lambda(\Omega)$ ) are sought. Instead, we work with *fully discrete spaces* made of vectors of polynomial components on the mesh cells (of various dimensions). The meaning of these components is provided by the interpolators on the fully discrete spaces.
- Due to the absence of explicit underlying conforming spaces, the differential operator of the complex cannot be the exterior derivative. Instead, a *discrete exterior derivative* is constructed combining the polynomial components to mimic the Stokes formula.
- *Discrete potentials* are also designed, again mimicking the Stokes formula. They are piecewise (discontinuous) polynomial forms on the mesh used, in particular, to define an  $L^2$ -structure on the discrete spaces (an essential tool to discretise PDEs written in weak form).

The choice of the polynomial components in the spaces and the design of discrete exterior derivatives and potentials revolve around two key properties: *polynomial consistency*, which is related to the ability to reproduce polynomial differential forms up to a selected polynomial degree, and *compatibility*, linked to the existence of an isomorphism between the cohomology of the discrete and continuous de Rham complexes. While both the DDR- and VEM-inspired constructions heavily rely on discrete versions of the Stokes formula, they do so in a radically different spirit: in the DDR construction, the choice of components in the discrete spaces is inspired by the formula to reconstruct a discrete exterior derivative, which is then used to construct discrete potentials. In the VEM construction, on the other hand, the space components (and, in particular, those associated with differentials) are chosen based on the formula used to define a discrete potential. While the choice in the DDR construction leads to leaner spaces, the study of its properties is more elaborated. Notice that, at this early stage, we haven't tried to identify the virtual (conforming) spaces that underlie the VEM-inspired construction and we have made no effort whatsoever in trying to reduce the dimension of the discrete spaces through serendipity.

The rest of this work is organised as follows. In Section 2 we establish the setting. In Section 3 we present and analyse the discrete complex generalising the DDR construction of [33]. Section 4 contains the definition and analysis of the complex generalising the VEM construction of [8]. In Section 5, we discuss in greater detail similarities and differences with respect to the FEEC, FES and Distributional Differential Forms frameworks.

Differential forms of any degree in dimensions 2 and 3 have interpretations in terms of vector fields. To make the exposition self-contained and improve the legibility for the reader not used to differential forms, we recall some facts on these so-called vector proxies in Appendix A, and we include throughout the exposition a series of examples to illustrate the development in the differential forms framework through vector calculus operators.

## 2 Setting

We present here the main notions used in the construction of the polytopal complexes of differential forms. For the reader not used to the framework of differential forms, we recall in Appendix A some basic concepts and definitions.

### 2.1 Spaces of differential forms

Let  $M$  denote an  $n$ -dimensional manifold. In what follows,  $M$  will typically be a cell of a polytopal mesh (see Section 2.5 below), and thus a relatively open set in a subspace of  $\mathbb{R}^m$  for some  $m \geq n$ . For any natural number  $\ell$  such that  $0 \leq \ell \leq n$ , we will denote by  $\Lambda^\ell(M)$  the space of differential  $\ell$ -forms (often just called  $\ell$ -forms) on  $M$  without explicit regularity requirements. When relevant, regularity is made explicit by prepending the appropriate space (e.g.,  $L^2\Lambda^\ell(M)$  stands for square-integrable  $\ell$ -forms).

## 2.2 Integration by parts

We recall the following integration by parts (Stokes) formula:

$$\int_M d\omega \wedge \mu = (-1)^{\ell+1} \int_M \omega \wedge d\mu + \int_{\partial M} \text{tr}_{\partial M} \omega \wedge \text{tr}_{\partial M} \mu \quad \forall (\omega, \mu) \in C^1 \Lambda^\ell(\overline{M}) \times C^1 \Lambda^{n-\ell-1}(\overline{M}), \quad (2.1)$$

where, for any form degree  $m$ ,  $\text{tr}_{\partial M} : C^0 \Lambda^m(\overline{M}) \rightarrow C^0 \Lambda^m(\partial M)$  is the trace operator. Formula (2.1) will provide the starting point to define discrete counterparts of the exterior derivative and of the corresponding potentials on mesh cells. It will also drive the choice of the components in the discrete spaces, geared at ensuring that the reconstructions preserve certain polynomial differential forms.

## 2.3 Hodge star

Assume now that  $M$  is an open set in a subspace of  $\mathbb{R}^m$ . We denote by  $\star : \Lambda^\ell(M) \rightarrow \Lambda^{n-\ell}(M)$  the Hodge star operator, and we set

$$\star^{-1} := (-1)^{\ell(n-\ell)} \star, \quad (2.2)$$

a notation justified observing that, for any  $\omega \in \text{Alt}^\ell(V)$ ,  $\star^{-1} \star \omega = \omega$  (see (A.3) in the appendix).

## 2.4 $L^2$ -orthogonal projectors

Integrating the inner product of  $\text{Alt}^\ell(V)$  over  $M$  yields the inner product of  $L^2 \Lambda^\ell(M)$ . For any closed subspace  $\mathcal{X}$  of  $L^2 \Lambda^\ell(M)$ , we therefore have an  $L^2$ -orthogonal projector  $\pi_{\mathcal{X}} : L^2 \Lambda^\ell(M) \rightarrow \mathcal{X}$  on  $\mathcal{X}$ , defined by the following relation: For all  $\omega \in L^2 \Lambda^\ell(M)$ ,  $\pi_{\mathcal{X}} \omega \in \mathcal{X}$  satisfies

$$\int_M \pi_{\mathcal{X}} \omega \wedge \star \mu = \int_M \omega \wedge \star \mu \quad \forall \mu \in \mathcal{X}. \quad (2.3)$$

To improve legibility, we will introduce in the following sections specific notations to  $\pi_{\mathcal{X}}$  for some polynomial subspaces  $\mathcal{X}$  that are particularly relevant to our construction.

For future use, we note the following property.

**Lemma 1** (Projectors on subspaces of differential forms). *Let  $(k, d)$  be integers such that  $k \leq d \leq n$ ,  $f \in \Delta_d(\mathcal{M}_h)$  and  $\mathcal{X}$  be a closed subspace of  $L^2 \Lambda^{d-k}(f)$ . Then, it holds: For all  $\omega \in L^2 \Lambda^k(f)$  and all  $\mu \in \mathcal{X}$ ,*

$$\int_f \star^{-1} \pi_{\mathcal{X}}(\star \omega) \wedge \mu = \int_f \mu \wedge \star \pi_{\mathcal{X}}(\star \omega) = \int_f \omega \wedge \mu. \quad (2.4)$$

*Proof.* The first relation in (2.4) follows from (A.4). To prove the second relation, we write

$$\int_f \mu \wedge \star \pi_{\mathcal{X}}(\star \omega) = \int_f \pi_{\mathcal{X}}(\star \omega) \wedge \star \mu = \int_f \mu \wedge (\star \star \omega) = \int_f \omega \wedge \mu,$$

where the first equality follows from (A.4) (with  $(\omega, \mu) \leftarrow (\pi_{\mathcal{X}}(\star \omega), \mu)$ ), the cancellation of the projector is justified by its definition (2.3), the second equality is obtained using (A.4) again, and the conclusion follows from (A.3) and the anticommutativity (A.1) of  $\wedge$ .  $\square$

## 2.5 Polytopal mesh

From this point on,  $\Omega$  will denote a polytopal domain of  $\mathbb{R}^n$ . We let  $\mathcal{M}_h$  denote a *polytopal mesh* of  $\Omega$ , i.e., a collection of disjoint polytopal sets (mesh entities) of dimensions in  $[0, n]$ , relatively open in their spanned affine space, such that the boundary of each  $d$ -cell (polytopal set of dimension  $d$ ) is the union of mesh entities of dimension  $< d$ , and such that any  $d$ -cell for  $d < n$  is contained in the boundary of some  $(d+1)$ -cell. For any  $d \in [0, n]$ , the set collecting all  $d$ -cells of  $\mathcal{M}_h$  is denoted by  $\Delta_d(\mathcal{M}_h)$ . Notice that this notion of polytopal mesh essentially coincides with that of CW-complex in

algebraic topology. Thus, when  $\Omega$  is a domain in dimension  $n = 3$ ,  $\mathcal{M}_h$  gathers the vertices collected in the set  $\mathcal{V}_h := \Delta_0(\mathcal{M}_h)$ , the edges collected in the set  $\mathcal{E}_h := \Delta_1(\mathcal{M}_h)$ , the faces collected in the set  $\mathcal{F}_h := \Delta_2(\mathcal{M}_h)$ , and the elements collected in the set  $\mathcal{T}_h := \Delta_3(\mathcal{M}_h)$ . For all  $f \in \mathcal{M}_h$ , we select a point  $\mathbf{x}_f \in f$  which, when  $\mathcal{M}_h$  belongs to a refined mesh sequence, is assumed at a distance from the boundary of  $f$  comparable to the meshsize.

If  $f \in \Delta_d(\mathcal{M}_h)$  and  $d' \leq d$  is an integer, we denote by  $\Delta_{d'}(f)$  the set of subcells of  $f$  of dimension  $d'$ . Hence, if  $n = d = 3$ , so that  $f = T \in \mathcal{T}_h$  is a polyhedral element of the mesh,  $f \in \Delta_{d'}(T)$  is a vertex of  $T$  if  $d' = 0$ , an edge of  $T$  if  $d' = 1$ , a polygonal face of  $T$  if  $d' = 2$ , or  $T$  itself if  $d' = 3$ .

## 2.6 Local polynomial spaces of differential forms

Let  $f \in \Delta_d(\mathcal{M}_h)$ ,  $0 \leq d \leq n$ . For any integer  $r \geq 0$ , we denote by  $\mathcal{P}_r \Lambda^\ell(f)$  the space of polynomial  $\ell$ -forms of total degree  $\leq r$  on  $f$ . We also adopt the standard convention  $\mathcal{P}_{-1} \Lambda^\ell(f) := \{0\}$ . We denote by  $\pi_{r,f}^\ell : L^2 \Lambda^\ell(f) \rightarrow \mathcal{P}_r \Lambda^\ell(f)$  the  $L^2$ -orthogonal projector on  $\mathcal{P}_r \Lambda^\ell(f)$ , defined by (2.3) with  $\mathcal{X} = \mathcal{P}_r \Lambda^\ell(f)$ .

The Koszul differential on  $f$  (translated by  $\mathbf{x}_f$ ) is denoted by  $\kappa$  so that, for all  $\omega \in \Lambda^\ell(f)$ ,  $\kappa\omega \in \Lambda^{\ell-1}(f)$  satisfies  $(\kappa\omega)_x(\mathbf{v}_1, \dots, \mathbf{v}_{\ell-1}) = \omega_x(\mathbf{x} - \mathbf{x}_f, \mathbf{v}_1, \dots, \mathbf{v}_{\ell-1})$  for all vectors  $\mathbf{v}_1, \dots, \mathbf{v}_{\ell-1}$  tangent to  $f$ . For any  $f \in \Delta_d(\mathcal{M}_h)$ ,  $1 \leq d \leq n$ , any integer  $\ell \in [0, d]$ , and any polynomial degree  $r \geq 0$ , we define the Koszul complement space as

$$\mathcal{K}_r^\ell(f) := \kappa \mathcal{P}_{r-1} \Lambda^{\ell+1}(f). \quad (2.5)$$

The indices  $r$  and  $\ell$  in this notation serve as a reminder that elements in  $\mathcal{K}_r^\ell(f)$  are polynomial  $\ell$ -forms of polynomial degree  $r$ . Note also that, since  $\mathcal{P}_{-1} \Lambda^\ell(f) = \{0\}$  and  $\Lambda^{d+1}(f) = \{0\}$ , we have

$$\mathcal{K}_0^\ell(f) = \mathcal{K}_r^d(f) = \{0\} \text{ for all } \ell \text{ and all } r, \text{ respectively.} \quad (2.6)$$

Moreover, since  $\kappa \Lambda^0(f) = \{0\}$ , we adopt the convention  $\mathcal{K}_r^{-1}(f) := \{0\}$  for all  $r$ . We denote by  $\pi_{r,f}^{\mathcal{K},\ell}$  the  $L^2$ -orthogonal projector  $L^2 \Lambda^\ell(f) \rightarrow \mathcal{K}_r^\ell(f)$ , defined by (2.3) with  $\mathcal{X} = \mathcal{K}_r^\ell(f)$ .

For all integers  $r \geq 0$  and  $\ell \in [0, d]$ , the following direct decomposition holds (see [4, Eq. (3.11)] for  $\ell \geq 1$ , the case  $\ell = 0$  can be directly checked):

$$\mathcal{P}_r \Lambda^0(f) = \mathcal{P}_0 \Lambda^0(f) \oplus \mathcal{K}_r^0(f), \quad (2.7a)$$

$$\mathcal{P}_r \Lambda^\ell(f) = d \mathcal{P}_{r+1} \Lambda^{\ell-1}(f) \oplus \mathcal{K}_r^\ell(f) \quad \text{if } \ell \geq 1. \quad (2.7b)$$

Since  $d \circ d = 0$  and  $d \mathcal{P}_0 \Lambda^0(f) = \{0\}$ , this shows that

$$d \mathcal{P}_r \Lambda^\ell(f) = d \mathcal{K}_r^\ell(f). \quad (2.8)$$

Applying this relation to  $(r+1, \ell-1)$  instead of  $(r, \ell)$  and recalling that  $d$  is one-to-one on  $\mathcal{K}_{r+1}^{\ell-1}(f)$  (see [4, Theorem 3.2]), this shows that, for  $\ell \geq 1$ , the following mapping is an isomorphism:

$$\begin{aligned} \mathcal{K}_{r+1}^{\ell-1}(f) \times \mathcal{K}_r^\ell(f) &\xrightarrow{\cong} \mathcal{P}_r \Lambda^\ell(f), \\ (\mu, \nu) &\mapsto d\mu + \nu. \end{aligned} \quad (2.9)$$

**Example 2** (Interpretation in terms of vector proxies). *In the case  $n = 3$ , thanks to the links between differential forms and vector proxies (see Appendix A), we can associate to each space of polynomial differential forms a space of (vector- or scalar-valued) polynomial fields. Let us consider decomposition (2.7b). We denote by  $f_d$  a  $d$ -cell of  $\mathcal{M}_h$ , and we use a notation analogous to that of [33] for*

polynomial spaces and vector calculus differential operators (with the exception that polynomial degrees are in subscripts instead of superscripts). Then, by definition (2.5) of the Koszul space, when  $f_3 = T \in \mathcal{T}_h = \Delta_3(\mathcal{M}_h)$  is a mesh element, we have

$$\begin{aligned} d\mathcal{P}_{r+1}\Lambda^0(f_3) &\leftrightarrow \mathcal{G}_r(T) := \mathbf{grad} \mathcal{P}_{r+1}(T), & \mathcal{K}_r^1(f_3) &\leftrightarrow \mathcal{G}_r^c(T) := (\mathbf{x} - \mathbf{x}_T) \times \mathcal{P}_{r-1}(T), \\ d\mathcal{P}_{r+1}\Lambda^1(f_3) &\leftrightarrow \mathcal{R}_r(T) := \mathbf{curl} \mathcal{P}_{r+1}(T), & \mathcal{K}_r^2(f_3) &\leftrightarrow \mathcal{R}_r^c(T) := (\mathbf{x} - \mathbf{x}_T) \mathcal{P}_{r-1}(T), \\ d\mathcal{P}_{r+1}\Lambda^2(f_3) &\leftrightarrow \mathbf{div} \mathcal{P}_{r+1}(T) = \mathcal{P}_r(T), & \mathcal{K}_r^3(f_3) &= \{0\}, \end{aligned}$$

where the first identity in the last line results from the surjectivity of the divergence operator.

On the other hand, when  $f_2 = F \in \mathcal{F}_h = \Delta_2(\mathcal{M}_h)$  is a mesh face, we obtain the following pair of possible correspondences:

$$d\mathcal{P}_{r+1}\Lambda^0(f_2) \leftrightarrow \mathcal{G}_r(F) := \mathbf{grad}_F \mathcal{P}_{r+1}(F), \quad \mathcal{K}_r^1(f_2) \leftrightarrow \mathcal{G}_r^c(F) := (\mathbf{x} - \mathbf{x}_F)^\perp \mathcal{P}_{r-1}(F) \quad (2.10)$$

or

$$d\mathcal{P}_{r+1}\Lambda^0(f_2) \leftrightarrow \mathcal{R}_r(F) := \mathbf{rot}_F \mathcal{P}_{r+1}(F), \quad \mathcal{K}_r^1(f_2) \leftrightarrow \mathcal{R}_r^c(F) := (\mathbf{x} - \mathbf{x}_F) \mathcal{P}_{r+1}(F), \quad (2.11)$$

where, for any  $\mathbf{v} \in \mathbb{R}^2$ ,  $\mathbf{v}^\perp = \varrho_{-\pi/2} \mathbf{v}$  is the clockwise rotation of  $\mathbf{v}$  with respect to the orientation of  $F$ . The existence of two possible correspondences between polynomial 1-forms and polynomial vector fields is due to the fact that, when  $d = 2$ , one can identify a 1-form either with a vector field  $\mathbf{v} = (v_1, v_2)$  or with its rotation through a right angle (cf. [2, Chapter 6]); in particular, we choose to identify it with the clockwise rotation  $\mathbf{v}^\perp = (v_2, -v_1)$  (see Appendix A for further details). By (2.6), we have  $\mathcal{K}_r^2(f_2) = \{0\}$  and, according to whether we consider the vector proxy leading to (2.10) or (2.11),

$$d\mathcal{P}_{r+1}\Lambda^1(f_2) \leftrightarrow \mathbf{rot}_F \mathcal{P}_{r+1}(F) = \mathcal{P}_r(F) \quad \text{or} \quad d\mathcal{P}_{r+1}\Lambda^1(f_2) \leftrightarrow \mathbf{div}_F \mathcal{P}_{r+1}(F) = \mathcal{P}_r(F).$$

Hence, since both 1-forms and 2-forms in  $\mathbb{R}^3$  can be identified with vector fields, and accounting for the two-fold identification of 1-forms in  $\mathbb{R}^2$ , the decomposition (2.7b) reads, in terms of proxies,

$$\mathcal{P}_r(f_d) = \mathcal{G}_r(f_d) \oplus \mathcal{G}_r^c(f_d) = \mathcal{R}_r(f_d) \oplus \mathcal{R}_r^c(f_d), \quad d \in \{2, 3\},$$

i.e., the same expressions as [33, Eqs. (2.4) and (2.6)]. On the other hand, concerning 0-forms, the decomposition (2.7a) reads, in terms of proxies,

$$\mathcal{P}_r(f_d) = \mathcal{P}_0(f_d) \oplus \mathcal{P}_r^b(f_d), \quad d \in \{0, \dots, 3\},$$

where we have introduced the notation  $\mathcal{P}_r^b(f) := (\mathbf{x} - \mathbf{x}_f) \cdot \mathcal{P}_{r-1}(f)$  for any  $f \in \Delta_d(\mathcal{M}_h)$ .

## 2.7 Trimmed local polynomial spaces

We recall the following local trimmed polynomial spaces (see e.g. [4, Theorem 3.5]): For any  $f \in \Delta_d(\mathcal{M}_h)$ ,  $1 \leq d \leq n$ ,

$$\mathcal{P}_r^- \Lambda^0(f) = \mathcal{P}_r \Lambda^0(f), \quad (2.12a)$$

$$\mathcal{P}_r^- \Lambda^\ell(f) = d\mathcal{P}_r \Lambda^{\ell-1}(f) \oplus \mathcal{K}_r^\ell(f) \quad \text{for } \ell \geq 1. \quad (2.12b)$$

In (2.12b), comparing with the decompositions (2.7), we have decreased by one the polynomial degree of the first space in the direct sum. Note that this definition leads to the choice

$$\mathcal{P}_r^- \Lambda^0(f) := \mathcal{P}_r \Lambda^0(f) \cong \mathbb{R} \quad \forall f \in \Delta_0(\mathcal{M}_h). \quad (2.13)$$

The  $L^2$ -orthogonal projector  $L^2\Lambda^\ell(f) \rightarrow \mathcal{P}_r^-\Lambda^\ell(f)$  is denoted by  $\pi_{r,f}^{-,\ell}$ , and is defined by (2.3) with  $\mathcal{X} = \mathcal{P}_r^-\Lambda^\ell(f)$ .

Let us note a few properties of trimmed polynomial spaces. For  $r = 0$ , only the space (2.12a) is non-trivial, that is,  $\mathcal{P}_0^-\Lambda^\ell(f) = \{0\}$  if  $\ell \in [1, d]$ . Applying, if  $r \geq 1$  and  $\ell \geq 1$ , (2.7b) with  $r - 1$  instead of  $r$  and noting that  $\mathcal{K}_{r-1}^\ell(f) \subset \mathcal{K}_r^\ell(f)$ , we obtain the equality

$$\mathcal{P}_r^-\Lambda^\ell(f) = \mathcal{P}_{r-1}\Lambda^\ell(f) + \mathcal{K}_r^\ell(f). \quad (2.14)$$

This equality, which obviously also holds for  $\ell = 0$  (see (2.7a)), shows that trimmed polynomial spaces sit between full polynomial spaces:

$$\mathcal{P}_{r-1}\Lambda^\ell(f) \subset \mathcal{P}_r^-\Lambda^\ell(f) \subset \mathcal{P}_r\Lambda^\ell(f).$$

Recalling that  $\mathcal{K}_r^d(f) = \{0\}$  and that  $d\mathcal{P}_r\Lambda^{d-1}(f) = \mathcal{P}_{r-1}\Lambda^d(f)$  (by exactness of the tail of the polynomial de Rham sequence [2, Corollary 7.3]), it holds

$$\mathcal{P}_r^-\Lambda^d(f) = \mathcal{P}_{r-1}\Lambda^d(f). \quad (2.15)$$

Applying (2.8) with  $\ell - 1$  instead of  $\ell$ , we moreover have

$$\mathcal{P}_r^-\Lambda^\ell(f) = d\mathcal{K}_r^{\ell-1}(f) + \mathcal{K}_r^\ell(f) \quad \text{for } \ell \geq 1. \quad (2.16)$$

Since  $d$  is one-to-one on  $\mathcal{K}_r^{\ell-1}(f)$ , this gives the following isomorphism, whenever  $\ell \geq 1$ :

$$\begin{aligned} \mathcal{K}_r^{\ell-1}(f) \times \mathcal{K}_r^\ell(f) &\xrightarrow{\cong} \mathcal{P}_r^-\Lambda^\ell(f), \\ (\mu, \nu) &\mapsto d\mu + \nu. \end{aligned} \quad (2.17)$$

**Example 3** (Interpretation of (2.12b) in terms of vector proxies). *Let  $n = 3$ . For  $d \in \{2, 3\}$ , denoting again by  $f_d$  a  $d$ -cell, we define the (local) Nédélec and Raviart–Thomas spaces*

$$\mathcal{N}_r(f_d) := \mathcal{G}_{r-1}(f_d) + \mathcal{G}_r^c(f_d), \quad \mathcal{RT}_r(f_d) := \mathcal{R}_{r-1}(f_d) + \mathcal{R}_r^c(f_d).$$

*Notice that, when  $d = 3$ , the Nédélec and Raviart–Thomas spaces can be obtained as polynomial spaces of vector proxies of (2.12b) for  $\ell = 1$  and  $\ell = 2$ , respectively. On the other hand, when considering  $d = 2$ , both spaces can be obtained by taking the same value  $\ell = 1$  in (2.12b). Again, this is linked to the two-fold interpretation of 1-forms in terms of vector proxies in  $\mathbb{R}^2$ , as already mentioned in Example 2, and corresponds to the well-known fact that two-dimensional Nédélec elements coincide with two-dimensional Raviart–Thomas elements rotated by a right angle.*

The following result generalises [33, Proposition 8].

**Lemma 4** (Traces of trimmed polynomial spaces). *The trace preserves trimmed spaces: For all integers  $d \in [0, n]$ ,  $d' \in [0, d]$  and  $\ell \in [0, d']$ , all  $f \in \Delta_d(\mathcal{M}_h)$  and all  $f' \in \Delta_{d'}(f)$ , we have*

$$\text{tr}_{f'} \mathcal{P}_r^-\Lambda^\ell(f) \subset \mathcal{P}_r^-\Lambda^\ell(f').$$

*Proof.* We first notice that the case  $\ell = 0$  is obvious since, in this case, trimmed spaces are full polynomial spaces (see (2.12a)), and the trace preserves full polynomial spaces. We therefore assume in the rest of the proof that  $\ell \geq 1$ . As the Koszul operators on differential forms on  $f$  and  $f'$  are not the same (due to the translation by  $\mathbf{x}_f$  and  $\mathbf{x}_{f'}$ , respectively), we temporarily denote them in this proof by  $\kappa_f$  and  $\kappa_{f'}$ .



The trace is a pullback, so it commutes with  $d$ , and we thus have

$$\mathrm{tr}_{f'}(d\mathcal{P}_r\Lambda^{\ell-1}(f)) = d(\mathrm{tr}_{f'}\mathcal{P}_r\Lambda^{\ell-1}(f)) \subset d\mathcal{P}_r\Lambda^{\ell-1}(f'),$$

where the inclusion holds since the trace preserves full polynomial spaces. Given the definition (2.12b) of the trimmed spaces, the lemma follows if we show that

$$\mathrm{tr}_{f'}\mathcal{K}_r^\ell(f) \subset \mathcal{P}_r^-\Lambda^\ell(f') = \mathcal{P}_{r-1}\Lambda^\ell(f') + \mathcal{K}_r^\ell(f') \quad (2.18)$$

(where the equality follows from (2.14) applied to  $f'$  instead of  $f$ ). Let  $\omega \in \mathcal{P}_{r-1}\Lambda^{\ell+1}(f)$ . The definitions of  $\mathrm{tr}_{f'}$  and  $\kappa_f$  give, for any  $\mathbf{x} \in f'$  and  $\mathbf{v}_1, \dots, \mathbf{v}_\ell$  tangent to  $f'$ ,

$$\begin{aligned} \mathrm{tr}_{f'}(\kappa_f\omega)_\mathbf{x}(\mathbf{v}_1, \dots, \mathbf{v}_\ell) &= \omega_\mathbf{x}(\mathbf{x} - \mathbf{x}_f, \mathbf{v}_1, \dots, \mathbf{v}_\ell) \\ &= \omega_\mathbf{x}(\mathbf{x}_{f'} - \mathbf{x}_f, \mathbf{v}_1, \dots, \mathbf{v}_\ell) + \omega_\mathbf{x}(\mathbf{x} - \mathbf{x}_{f'}, \mathbf{v}_1, \dots, \mathbf{v}_\ell) \\ &= \alpha_\mathbf{x}(\mathbf{v}_1, \dots, \mathbf{v}_\ell) + (\kappa_{f'}\mathrm{tr}_{f'}\omega)_\mathbf{x}(\mathbf{v}_1, \dots, \mathbf{v}_\ell), \end{aligned}$$

where we have used the linearity of  $\omega_\mathbf{x}$  with respect to its first argument to obtain the second equality, and introduced the differential form  $\alpha := \omega(\mathbf{x}_{f'} - \mathbf{x}_f, \cdot)$  in the third equality. Hence,  $\mathrm{tr}_{f'}(\kappa_f\omega) = \alpha + \kappa_{f'}\mathrm{tr}_{f'}\omega$ , which proves (2.18) since  $\alpha \in \mathcal{P}_{r-1}\Lambda^\ell(f')$  (as  $\mathbf{x}_{f'} - \mathbf{x}_f$  is constant) and  $\mathrm{tr}_{f'}\omega \in \mathcal{P}_{r-1}\Lambda^{\ell+1}(f')$ .  $\square$

### 3 Discrete de Rham complex

We define in this section a discrete counterpart of the de Rham complex of differential forms (1.2) in the spirit of [33, 38]. Let, from this point on, an integer  $r \geq 0$  be fixed corresponding to the polynomial degree of the discrete sequence. The general idea is, for each form degree  $k \in [0, n]$ , to select the polynomial components of the discrete spaces in order to reconstruct, on each  $d$ -cell  $f$  and iteratively on the dimension  $d$ :

- A *discrete exterior derivative* in the full polynomial space  $\mathcal{P}_r\Lambda^{k+1}(f)$  that can reproduce exactly the exterior derivative of differential forms in  $\mathcal{P}_{r+1}^-\Lambda^k(f)$ ;
- Based on this discrete exterior derivative and on traces on  $(d-1)$ -cells (either directly available or reconstructed), a *discrete potential* in  $\mathcal{P}_r\Lambda^k(f)$  that can reproduce exactly differential forms belonging to this same space.

#### 3.1 Definition

##### 3.1.1 Discrete spaces

The discrete counterpart  $\underline{X}_{r,h}^k$  of the space  $H\Lambda^k(\Omega)$ ,  $0 \leq k \leq n$ , is defined as

$$\underline{X}_{r,h}^k := \bigtimes_{d=k}^n \bigtimes_{f \in \Delta_d(\mathcal{M}_h)} \mathcal{P}_r^-\Lambda^{d-k}(f). \quad (3.1)$$

We define the restrictions of the global space (3.1) to a mesh entity or its boundary as follows: For all integers  $k$  and  $d$  such that  $0 \leq k \leq d \leq n$  and all  $f \in \Delta_d(\mathcal{M}_h)$ ,

$$\underline{X}_{r,f}^k := \bigtimes_{d'=k}^d \bigtimes_{f' \in \Delta_{d'}(f)} \mathcal{P}_r^-\Lambda^{d'-k}(f') \quad \text{and} \quad \underline{X}_{r,\partial f}^k := \bigtimes_{d'=k}^{d-1} \bigtimes_{f' \in \Delta_{d'}(f)} \mathcal{P}_r^-\Lambda^{d'-k}(f') \quad \text{if } d \geq 1.$$

We shall use the notation  $\underline{\omega}_h = (\omega_f)_{f \in \Delta_d(\mathcal{M}_h), d \in [k, n]} \in \underline{X}_{r,h}^k$  for a generic element of the global discrete space of  $k$ -forms and  $\underline{\omega}_f = (\omega_{f'})_{f' \in \Delta_{d'}(f), d' \in [k, d]} \in \underline{X}_{r,f}^k$  (resp.,  $\underline{\omega}_{\partial f} = (\omega_{f'})_{f' \in \Delta_{d'}(f), d' \in [k, d-1]} \in$

| $k \backslash d$ | 0   | 1                                  | 2                                  | 3                                  |
|------------------|---|------------------------------------|------------------------------------|------------------------------------|
| 0                | $\mathbb{R} = \mathcal{P}_r \Lambda^0(f_0)$ | $\mathcal{P}_{r-1} \Lambda^1(f_1)$ | $\mathcal{P}_{r-1} \Lambda^2(f_2)$ | $\mathcal{P}_{r-1} \Lambda^3(f_3)$ |
| 1                |   | $\mathcal{P}_r \Lambda^0(f_1)$     | $\mathcal{P}_r^- \Lambda^1(f_2)$   | $\mathcal{P}_r^- \Lambda^2(f_3)$   |
| 2                |   |                                    | $\mathcal{P}_r \Lambda^0(f_2)$     | $\mathcal{P}_r^- \Lambda^1(f_3)$   |
| 3                |   |                                    |                                    | $\mathcal{P}_r \Lambda^0(f_3)$     |

  

| $k \backslash d$ | 0                                 | 1                        | 2                        | 3                        |
|------------------|-----------------------------------|--------------------------|--------------------------|--------------------------|
| 0                | $\mathbb{R} = \mathcal{P}_r(f_0)$ | $\mathcal{P}_{r-1}(f_1)$ | $\mathcal{P}_{r-1}(f_2)$ | $\mathcal{P}_{r-1}(f_3)$ |
| 1                |                                   | $\mathcal{P}_r(f_1)$     | $\mathcal{RT}_r(f_2)$    | $\mathcal{RT}_r(f_3)$    |
| 2                |                                   |                          | $\mathcal{P}_r(f_2)$     | $\mathcal{N}_r(f_3)$     |
| 3                |                                   |                          |                          | $\mathcal{P}_r(f_3)$     |

Table 1: Polynomial components attached to each mesh entity  $f_d$  of dimension  $d \in \{0, \dots, 3\}$  for the space  $\underline{X}_{r,h}^k$  for  $k \in \{0, \dots, 3\}$  (top) and counterpart through vector proxies (bottom).

$\underline{X}_{r,\partial f}^k$ ) for its restriction to  $f$  (resp.,  $\partial f$ ), obtained collecting the components on the mesh entities  $f' \in \Delta_{d'}(f)$ ,  $d' \in [k, d]$  (resp.,  $d' \in [k, d-1]$ ). As a generic convention in this article, underlined letters denote spaces or vectors made of polynomial components on mesh entities. Table 1 gives an overview of the polynomial unknowns in  $\underline{X}_{r,f}^k$ , along with their vector proxies, in dimensions 0 to 3.

### 3.1.2 Interpolators and interpretation of the polynomial components

The precise meaning of the components in each DDR space is provided by the corresponding interpolator. For  $f \in \Delta_d(\mathcal{M}_h)$  and  $k \leq d$ , the interpolator  $\underline{I}_{r,f}^k : C^0 \Lambda^k(\bar{f}) \rightarrow \underline{X}_{r,f}^k$  is defined by: For all  $\omega \in C^0 \Lambda^k(\bar{f})$ ,

$$\underline{I}_{r,f}^k \omega := (\pi_{r,f'}^-, d'^{-k} (\star \text{tr}_{f'} \omega))_{f' \in \Delta_{d'}(f), d' \in [k, d]}. \quad (3.2)$$

In other words, a discrete  $k$ -form on the mesh is made of polynomial forms attached to each mesh entity of dimension  $d \geq k$ ; on each such entity, the form is of degree  $d - k$  as it corresponds to the Hodge star of an underlying  $k$ -form. The Hodge star operator is used in the definition of the polynomial components to ensure that the full space  $\mathcal{P}_r \Lambda^0(f)$  (see (2.12a)) is attached to the lowest-dimensional cells  $f \in \Delta_k(\mathcal{M}_h)$ .

### 3.1.3 Local discrete potentials and discrete exterior derivative

Let  $0 \leq k \leq n$  be a fixed integer. For all  $f \in \Delta_d(\mathcal{M}_h)$  with  $d \geq k$ , we define the *discrete potential*  $P_{r,f}^k : \underline{X}_{r,f}^k \rightarrow \mathcal{P}_r \Lambda^k(f)$  and, if  $d \geq k+1$ , the *discrete exterior derivative*  $d_{r,f}^k : \underline{X}_{r,f}^k \rightarrow \mathcal{P}_r \Lambda^{k+1}(f)$  recursively on the dimension  $d$  as follows:

- If  $d = k$ , then the discrete potential on  $f$  is directly given by the component of  $\underline{\omega}_f$  on  $f$ :

$$P_{r,f}^k \underline{\omega}_f := \star^{-1} \omega_f \in \mathcal{P}_r \Lambda^d(f). \quad (3.3)$$

- If  $k+1 \leq d \leq n$ :

1. First, the discrete exterior derivative is defined by: For all  $\underline{\omega}_f \in \underline{X}_{r,f}^k$ ,

$$\int_f d_{r,f}^k \underline{\omega}_f \wedge \mu = (-1)^{k+1} \int_f \star^{-1} \omega_f \wedge d\mu + \int_{\partial f} P_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \mu \quad \forall \mu \in \mathcal{P}_r \Lambda^{d-k-1}(f), \quad (3.4)$$

where we have introduced the piecewise polynomial boundary potential  $P_{r,\partial f}^k : \underline{X}_{r,\partial f}^k \rightarrow \Lambda^k(\partial f)$  such that  $(P_{r,\partial f}^k)|_{f'} := P_{r,f'}^k$  for all  $f' \in \Delta_{d-1}(f)$  ( $P_{r,f'}^k$  being the discrete potential on the  $(d-1)$ -cell  $f'$  defined at the previous step).

2. Then, the discrete potential on the  $d$ -cell  $f$  is given by: For all  $\underline{\omega}_f \in \underline{X}_{r,f}^k$ ,

$$\begin{aligned} (-1)^{k+1} \int_f P_{r,f}^k \underline{\omega}_f \wedge (d\mu + \nu) \\ = \int_f d_{r,f}^k \underline{\omega}_f \wedge \mu - \int_{\partial f} P_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \mu + (-1)^{k+1} \int_f \star^{-1} \omega_f \wedge \nu \\ \forall (\mu, \nu) \in \mathcal{K}_{r+1}^{d-k-1}(f) \times \mathcal{K}_r^{d-k}(f). \end{aligned} \quad (3.5)$$

Some remarks are in order.

*Remark 5* (Definitions (3.4) and (3.5)). The fact that condition (3.4) defines  $d_{r,f}^k \underline{\omega}_f$  uniquely is an immediate consequence of the Riesz representation theorem for  $\mathcal{P}_r \Lambda^{k+1}(f)$  equipped with the  $L^2$ -product  $(\rho, \beta) \ni L^2 \Lambda^{k+1}(f) \times L^2 \Lambda^{k+1}(f) \mapsto \int_f \rho \wedge \star \beta \in \mathbb{R}$ , after observing that (3.4) can be equivalently reformulated as follows (notice the change in the degree of the test differential form, with  $\beta$  below corresponding to  $\star^{-1} \mu$  in (3.4)):

$$\int_f d_{r,f}^k \underline{\omega}_f \wedge \star \beta = (-1)^{k+1} \int_f \omega_f \wedge \star d \star \beta + \int_{\partial f} P_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \star \beta \quad \forall \beta \in \mathcal{P}_r \Lambda^{k+1}(f),$$

where we have additionally used (A.4) for the first term in the right-hand side. Similar considerations apply to the definition (3.5) of  $P_{r,f}^k$ , applying the isomorphism (2.9) with  $\ell = d - k \geq 1$ .

*Remark 6* (Validity of (3.5)). For  $k+1 \leq d \leq n$ , equation (3.5) actually holds for all  $\mu \in \mathcal{P}_{r+1}^- \Lambda^{d-k-1}(f)$ . To prove this assertion, since (3.5) holds for  $\mu \in \mathcal{K}_{r+1}^{d-k-1}(f)$ , it suffices to show that it also holds for  $\nu = 0$  and  $\mu$  belonging to  $\mathcal{P}_0 \Lambda^0(f)$  if  $d = k+1$  (see (2.12a) and (2.7a)) or  $d\mathcal{P}_{r+1} \Lambda^{d-k-2}(f)$  if  $d \geq k+2$  (see (2.12b)). In both cases, we have  $d\mu = 0$ , so that the left-hand side of (3.5) vanishes; since  $\mu \in \mathcal{P}_r \Lambda^{d-k-1}(f)$ , the right-hand side of (3.5) also vanishes due to the definition (3.4) of the discrete exterior derivative, which concludes the argument.

*Remark 7* (Potential for  $k = 0$ ). In the case  $k = 0$ , we can define an improved potential  $P_{r+1,f}^0 : \underline{X}_{r+1,f}^0 \rightarrow \mathcal{P}_{r+1} \Lambda^0(f)$  of polynomial degree  $r+1$  (instead of  $r$ ) as follows: For all  $\underline{\omega}_f \in \underline{X}_{r,f}^0$ ,

- If  $d = 0$ , then  $P_{r+1,f}^0 \underline{\omega}_f = \star^{-1} \omega_f \in \mathcal{P}_r \Lambda^0(f) \cong \mathbb{R} \cong \mathcal{P}_{r+1} \Lambda^0(f)$  (since  $f$  has dimension 0);
- If  $1 \leq d \leq n$ ,

$$- \int_f P_{r+1,f}^0 \underline{\omega}_f \wedge d\mu = \int_f d_{r,f}^0 \underline{\omega}_f \wedge \mu - \int_{\partial f} P_{r+1,\partial f}^0 \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \mu \quad \forall \mu \in \mathcal{K}_{r+2}^{d-1}(f). \quad (3.6)$$

This definition is justified by the isomorphism (2.9) with  $\ell = d$  and  $r+1$  instead of  $r$  (recalling that  $\mathcal{K}_{r+1}^d(f) = \{0\}$ ), and it can easily be checked, testing (3.5) and (3.6) with  $\mu \in \mathcal{K}_{r+1}^{d-1}(f)$ , that  $\pi_{r,f}^0 P_{r+1,f}^0 \underline{\omega}_f = P_{r,f}^0 \underline{\omega}_f$ . We will moreover see in Remark 18 that  $P_{r+1,f}^0$  enjoys optimal consistency properties.

*Remark 8* (Space of DDR potentials). The space of DDR reconstructed potentials, that is,

$$\{(P_{r,f}^k \underline{\omega}_h)_{f \in \Delta_d(\mathcal{M}_h), d \in [k,n]} : \underline{\omega}_h \in \underline{X}_{r,h}^k\}$$

cannot be considered as a space of differential forms with global regularity, as the reconstructed polynomials do not have any compatibility condition of the traces; they are inherently piecewise discontinuous polynomials.

**Example 9** (Interpretation in terms of vector proxies). *We start by considering  $k = 0$ . In this case, formula (3.3) means that (constant) real values are attached to the vertices  $f_0 = V \in \mathcal{V}_h = \Delta_0(\mathcal{M}_h)$  of the mesh, so that an iterative procedure can be initialised to reconstruct discrete gradients and related traces/potentials over higher-dimensional cells. Indeed, formula (3.4) reconstructs a (scalar) gradient over edges  $f_1 = E \in \mathcal{E}_h = \Delta_1(\mathcal{M}_h)$  (i.e., the derivative along the direction given by the orientation of  $E$ ) based on the values at the vertices and the value on the edge itself. This edge gradient, in turn, enters (3.5) to define a scalar edge trace over  $E$ . When  $d$  takes the values 2 and 3, the successive application of formulas (3.4)–(3.5) defines, respectively, the pairs (face gradient, scalar face trace) on mesh faces  $f_2 = F \in \mathcal{F}_h = \Delta_2(\mathcal{M}_h)$ , and (element gradient, scalar element potential) on mesh elements  $f_3 = T \in \mathcal{T}_h = \Delta_3(\mathcal{M}_h)$ .*

Let us now turn to the case  $k = 1$ , for which we provide more details. The vector proxy for the space  $\underline{X}_{r,h}^1$  is the space

$$\underline{X}_{\text{curl},h}^r = \prod_{E \in \mathcal{E}_h} \mathcal{P}_r(E) \times \prod_{F \in \mathcal{F}_h} \mathcal{RT}_r(F) \times \prod_{T \in \mathcal{T}_h} \mathcal{RT}_r(T)$$

and, with standard DDR notation, we denote by  $\underline{X}_{\text{curl},Y}^r$  its restriction to a mesh element or face  $Y \in \mathcal{T}_h \cup \mathcal{F}_h$ . By (3.3) with  $d = k = 1$ , the reconstruction process is initialised by 1-forms, whose vector proxies are scalar-valued polynomials of degree  $r$  over edges  $f_1 = E \in \mathcal{E}_h$  that play the role of edge tangential traces.

Then, for each mesh face  $f_2 = F \in \mathcal{F}_h$ , we sequentially reconstruct a scalar face curl  $C_F^r : \underline{X}_{\text{curl},F}^r \rightarrow \mathcal{P}_r(F)$  by (3.4) with  $d = k + 1 = 2$  and a vector face tangential trace  $\gamma_{\text{t},F}^r : \underline{X}_{\text{curl},F}^r \rightarrow \mathcal{P}_r(F)$  by (3.5).  $C_F^r$  is such that, for all  $\underline{v}_F = ((v_E)_{E \in \mathcal{E}_F}, \mathbf{v}_F) \in \underline{X}_{\text{curl},F}^r$ ,

$$\int_F C_F^r \underline{v}_F q = \int_F \mathbf{v}_F \cdot \mathbf{rot}_F q + \sum_{E \in \mathcal{E}_F} \varepsilon_{FE} \int_E v_E q \quad \forall q \in \mathcal{P}_r(F),$$

where, for all  $E \in \mathcal{E}_F$  (the set of edges of  $F$ ),  $\varepsilon_{FE} \in \{-1, +1\}$  denotes the orientation of  $E$  relative to  $F$ , while  $\gamma_{\text{t},F}^r$  satisfies, for all  $\underline{v}_F \in \underline{X}_{\text{curl},F}^r$ ,

$$\int_F \gamma_{\text{t},F}^r \underline{v}_F \cdot (\mathbf{rot}_F q + \mathbf{w}) = \int_F C_F^r \underline{v}_F q - \sum_{E \in \mathcal{E}_F} \varepsilon_{FE} \int_E v_E q + \int_F \mathbf{v}_F \cdot \mathbf{w},$$

$$\forall (q, \mathbf{w}) \in \mathcal{P}_{r+1}^b(F) \times \mathcal{R}_r^c(F).$$

The alternative interpretation of 1-forms in dimension  $d = 2$  results in a rotation of  $\underline{X}_{\text{curl},F}^r$  by a right angle. Correspondingly, (3.4) yields a face divergence (see (A.9) and (A.10) in Appendix A.2).

Next, for each mesh element  $f_3 = T \in \mathcal{T}_h$ , (3.4) defines the element curl  $C_T^r : \underline{X}_{\text{curl},T}^r \rightarrow \mathcal{P}_r(T)$  such that, for all  $\underline{v}_T = ((v_E)_{E \in \mathcal{E}_T}, (\mathbf{v}_F)_{F \in \mathcal{F}_T}, \mathbf{v}_T) \in \underline{X}_{\text{curl},T}^r$ ,

$$\int_T C_T^r \underline{v}_T \cdot \mathbf{w} = \int_T \mathbf{v}_T \cdot \mathbf{curl} \mathbf{w} + \sum_{F \in \mathcal{F}_T} \varepsilon_{TF} \int_F \gamma_{\text{t},F}^r \underline{v}_F \cdot (\mathbf{w} \times \mathbf{n}_F) \quad \forall \mathbf{w} \in \mathcal{P}_r(T),$$

where, for all  $F \in \mathcal{F}_T$  (the set of faces of  $T$ ),  $\varepsilon_{TF} \in \{-1, +1\}$  denotes the orientation of  $F$  relative to  $T$ , while (3.5) defines the vector potential  $\mathbf{P}_{\text{curl},T}^r : \underline{\mathbf{X}}_{\text{curl},T}^r \rightarrow \mathcal{P}_r(T)$  such that, for all  $\underline{\mathbf{v}}_T \in \underline{\mathbf{X}}_{\text{curl},T}^r$ ,

$$\int_T \mathbf{P}_{\text{curl},T}^r \underline{\mathbf{v}}_T \cdot (\text{curl } \mathbf{w} + \mathbf{z}) = \int_T \mathbf{C}_T^r \underline{\mathbf{v}}_T \cdot \mathbf{w} - \sum_{F \in \mathcal{F}_T} \varepsilon_{TF} \int_F \gamma_{\mathbf{v},F}^r \underline{\mathbf{v}}_F \cdot (\mathbf{w} \times \mathbf{n}_F) + \int_T \mathbf{v}_T \cdot \mathbf{z}$$

$$\forall (\mathbf{w}, \mathbf{z}) \in \mathcal{G}_{r+1}^c(T) \times \mathcal{R}_r^c(T).$$

When  $k = 2$ , (3.4) reconstructs on mesh elements  $f_3 = T \in \mathcal{T}_h$  a discrete divergence of order  $r$  based on the polynomial scalar trace defined by (3.3), which plays the role of a normal trace on the face  $f_2 = F \in \mathcal{F}_T$ . Then, (3.5) defines a vector potential of degree  $r$  over  $T$ .

Finally, in the case  $k = 3$ , (3.3) simply yields a polynomial over mesh elements  $f_3 = T \in \mathcal{T}_h$ .

### 3.1.4 Global discrete exterior derivative and DDR complex

To arrange the spaces  $\underline{\mathbf{X}}_{r,h}^k$  into a sequence that mimics the continuous de Rham complex, for any form degree  $k$  such that  $0 \leq k \leq n-1$ , we introduce the *global discrete exterior derivative*  $\underline{\mathbf{d}}_{r,h}^k : \underline{\mathbf{X}}_{r,h}^k \rightarrow \underline{\mathbf{X}}_{r,h}^{k+1}$  defined as follows:

$$\underline{\mathbf{d}}_{r,h}^k \omega_h := (\pi_{r,f}^{-,d-k-1} (\star \underline{\mathbf{d}}_{r,f}^k \omega_f))_{f \in \Delta_d(\mathcal{M}_h), d \in [k+1, n]}. \quad (3.7)$$

In what follows, given a  $d$ -cell  $f \in \Delta_d(\mathcal{M}_h)$  with  $d \in [k+1, n]$ , we denote by  $\underline{\mathbf{d}}_{r,f}^k$  the *local discrete exterior derivative* collecting the components of  $\underline{\mathbf{d}}_{r,h}^k$  on  $f$  and its boundary. The DDR sequence reads

$$\text{DDR}(r) := \{0\} \longrightarrow \underline{\mathbf{X}}_{r,h}^0 \xrightarrow{\underline{\mathbf{d}}_{r,h}^0} \underline{\mathbf{X}}_{r,h}^1 \longrightarrow \dots \longrightarrow \underline{\mathbf{X}}_{r,h}^{n-1} \xrightarrow{\underline{\mathbf{d}}_{r,h}^{n-1}} \underline{\mathbf{X}}_{r,h}^n \longrightarrow \{0\}. \quad (3.8)$$

The main results concerning this sequence are stated hereafter.

**Theorem 10** (Cohomology of the Discrete de Rham complex). *The DDR sequence (3.8) is a complex and its cohomology is isomorphic to the cohomology of the continuous de Rham complex (1.2).*

*Proof.* See Section 3.5. □

**Theorem 11** (Polynomial consistency of the discrete potential and exterior derivative). *For all integers  $0 \leq k \leq d \leq n$  and all  $f \in \Delta_d(\mathcal{M}_h)$ , it holds*

$$\mathbf{P}_{r,f}^k \underline{\mathbf{I}}_{r,f}^k \omega = \omega \quad \forall \omega \in \mathcal{P}_r \Lambda^k(f), \quad (3.9)$$

and, if  $d \geq k+1$ ,

$$\underline{\mathbf{d}}_{r,f}^k \underline{\mathbf{I}}_{r,f}^k \omega = \text{d}\omega \quad \forall \omega \in \mathcal{P}_{r+1}^- \Lambda^k(f). \quad (3.10)$$

*Proof.* See Section 3.4. □

**Remark 12** (Consistency of traces). The above theorem actually implies that, for any  $d$ -face  $f \in \Delta_d(\mathcal{M}_h)$ ,  $d \in [k, n]$ , any  $\omega \in \mathcal{P}_r \Lambda^k(f)$ , and any integer  $d' \in [k, d]$ ,

$$\mathbf{P}_{r,f'}^k \underline{\mathbf{I}}_{r,f'}^k \omega = \text{tr}_{f'} \omega \quad \forall f' \in \Delta_{d'}(f).$$

This can be easily seen noticing that  $\underline{\mathbf{I}}_{r,f'}^k \omega = \underline{\mathbf{I}}_{r,f}^k \omega$ ,  $\text{tr}_{f'} \omega$  and  $\text{tr}_{f'} \omega \in \mathcal{P}_r \Lambda^k(f')$ , and invoking (3.9) with  $(f, \omega) \leftarrow (f', \text{tr}_{f'} \omega)$ .

### 3.1.5 Discrete $L^2$ -products

Using the potentials built in Section 3.1.3, we can define, for all  $k \in [0, n]$ , an inner product  $(\cdot, \cdot)_{k,h} : \underline{X}_{r,h}^k \times \underline{X}_{r,h}^k \rightarrow \mathbb{R}$  that induces an  $L^2$ -structure on  $\underline{X}_{r,h}^k$ . Specifically, we set: For all  $(\underline{\omega}_h, \underline{\mu}_h) \in \underline{X}_{r,h}^k \times \underline{X}_{r,h}^k$ ,

$$(\underline{\omega}_h, \underline{\mu}_h)_{k,h} := \sum_{f \in \Delta_n(\mathcal{M}_h)} (\underline{\omega}_f, \underline{\mu}_f)_{k,f} \quad (3.11)$$

$$\text{with } (\underline{\omega}_f, \underline{\mu}_f)_{k,f} := \int_f P_{r,f}^k \underline{\omega}_f \wedge \star P_{r,f}^k \underline{\mu}_f + s_{k,f}(\underline{\omega}_f, \underline{\mu}_f) \text{ for all } f \in \Delta_n(\mathcal{M}_h),$$

where  $s_{k,f} : \underline{X}_{r,f}^k \times \underline{X}_{r,f}^k \rightarrow \mathbb{R}$  is the stabilisation bilinear form such that

$$s_{k,f}(\underline{\omega}_f, \underline{\mu}_f) = \sum_{d'=k}^{n-1} h_f^{n-d'} \sum_{f' \in \Delta_{d'}(f)} \int_{f'} (\text{tr}_{f'} P_{r,f}^k \underline{\omega}_f - P_{r,f'}^k \underline{\omega}_{f'}) \wedge \star (\text{tr}_{f'} P_{r,f}^k \underline{\mu}_f - P_{r,f'}^k \underline{\mu}_{f'}),$$

with  $h_f$  denoting the diameter of  $f$ . The first term in the right-hand side of  $(\cdot, \cdot)_{k,f}$  is responsible for consistency, while the second one ensures the positivity of this bilinear form. More specifically, by Theorem 11 and Remark 12 it holds, for all  $f \in \Delta_n(\mathcal{M}_h)$ ,

$$(\underline{I}_{r,f}^k \underline{\omega}, \underline{\mu})_{k,f} = \int_f \underline{\omega} \wedge \star P_{r,f}^k \underline{\mu} \quad \forall \underline{\omega} \in \mathcal{P}_r \Lambda^k(f), \quad \forall \underline{\mu} \in \underline{X}_{r,f}^k.$$

Additionally, by (3.20) below, the mapping  $\underline{X}_{r,f}^k \ni \underline{\omega}_f \mapsto \|\underline{\omega}_f\|_{k,f} := (\underline{\omega}_f, \underline{\omega}_f)_{k,f}^{1/2} \in \mathbb{R}$  defines a norm on  $\underline{X}_{r,f}^k$ . Numerical schemes for linear PDEs related to the de Rham complex are typically obtained replacing continuous spaces and  $L^2$ -products with their discrete counterparts, according to the principles illustrated, e.g., in [33, Section 7].

*Remark 13* (Stabilisation). A more general expression for the local  $L^2$ -product in (3.11) is obtained replacing  $s_{k,f}$  with

$$s_{\mathcal{B},k,f}(\underline{\omega}_f, \underline{\mu}_f) = \mathcal{B}_f(\underline{I}_{r,f}^k P_{r,f}^k \underline{\omega}_f - \underline{\omega}_f, \underline{I}_{r,f}^k P_{r,f}^k \underline{\mu}_f - \underline{\mu}_f),$$

with  $\mathcal{B}_f : \underline{X}_{r,f}^k \times \underline{X}_{r,f}^k \rightarrow \mathbb{R}$  denoting a symmetric positive definite bilinear form inducing a norm that scales in  $h_f$  as  $\|\cdot\|_{k,f}$  defined above. Crucially,  $s_{\mathcal{B},k,f}$  depends on its arguments only through the difference operator  $\underline{X}_{r,f}^k \ni \underline{\omega}_f \mapsto \underline{I}_{r,f}^k P_{r,f}^k \underline{\omega}_f - \underline{\omega}_f \in \underline{X}_{r,f}^k$ , which guarantees that it vanishes whenever one of its arguments is the interpolate of a differential form in  $\mathcal{P}_r \Lambda^k(f)$  (as can be checked in the same spirit as [36, Lemma 2.11]).

## 3.2 Complex property

We introduce, for all integers  $d \in [1, n]$ , the piecewise polynomial boundary exterior derivative  $d_{r,\partial f}^k : \underline{X}_{r,\partial f}^k \rightarrow \Lambda^{k+1}(\partial f)$  such that  $(d_{r,\partial f}^k)_{|f'} := d_{r,f'}^k$  for all  $f' \in \Delta_{d-1}(f)$  ( $d_{r,f'}^k$  being the discrete exterior derivative on the  $(d-1)$ -cell  $f'$  defined by (3.4)). The following lemma is a generalisation of the links, in the DDR framework based on vector proxies, between element gradients (resp., curls) and face gradients (resp., curls), see [33, Propositions 1 and 4].

**Lemma 14** (Link between discrete exterior derivatives on subcells). *It holds, for all  $d \geq k+2$ , all  $f \in \Delta_d(\mathcal{M}_h)$ , and all  $\underline{\omega}_f \in \underline{X}_{r,f}^k$ ,*

$$\int_f d_{r,f}^k \underline{\omega}_f \wedge d\alpha = (-1)^{k+1} \int_{\partial f} d_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \alpha \quad \forall \alpha \in \mathcal{P}_{r+1}^- \Lambda^{d-k-2}(f). \quad (3.12)$$

*Proof.* Take  $\mu = d\alpha \in \mathcal{P}_r \Lambda^{d-k-1}(f)$  in (3.4) and use  $d \circ d = 0$  and  $\text{tr}_{\partial f} d = d \text{tr}_{\partial f}$  (since the trace is a pullback, it commutes with the exterior derivative) to get

$$\int_f d_{r,f}^k \underline{\omega}_f \wedge d\alpha = \int_{\partial f} P_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge d \text{tr}_{\partial f} \alpha. \quad (3.13)$$

For each  $f' \in \Delta_{d-1}(f)$  forming  $\partial f$ , by Lemma 4 we have  $\text{tr}_{f'} \alpha \in \mathcal{P}_{r+1}^- \Lambda^{d-k-2}(f')$  so, by (3.5) applied to  $f'$  instead of  $f$  with test function  $(\mu, \nu) = (\text{tr}_{f'} \alpha, 0)$  (see Remark 6), we have

$$(-1)^{k+1} \int_{f'} P_{r,f'}^k \underline{\omega}_{f'} \wedge d \text{tr}_{f'} \alpha = \int_{f'} d_{r,f'}^k \underline{\omega}_{f'} \wedge \text{tr}_{f'} \alpha - \int_{\partial f'} P_{r,\partial f'}^k \underline{\omega}_{\partial f'} \wedge \text{tr}_{\partial f'} (\text{tr}_{f'} \alpha).$$

Use  $\text{tr}_{\partial f'} \circ \text{tr}_{f'} = \text{tr}_{\partial f}$ , sum these relations over  $f' \in \Delta_{d-1}(f)$ , invoke  $\partial(\partial f) = 0$  (which implies that the integrals over  $(\partial f')_{f' \in \Delta_{d-1}(f)}$  cancel out due to compatible orientations), and plug the result into (3.13) to conclude.  $\square$

**Theorem 15** (Link between discrete potentials and exterior derivatives, complex property). *It holds, for all integers  $k \in [1, n]$  and  $d \geq k$ , all  $f \in \Delta_d(\mathcal{M}_h)$ , and all  $\underline{\omega}_f \in \underline{X}_{r,f}^{k-1}$ ,*

$$P_{r,f}^k (d_{r,f}^{k-1} \underline{\omega}_f) = d_{r,f}^{k-1} \underline{\omega}_f, \quad (3.14)$$

and, if  $d \geq k + 1$ ,

$$d_{r,f}^k (d_{r,f}^{k-1} \underline{\omega}_f) = 0. \quad (3.15)$$

As a consequence, the sequence (3.8) defines a complex.

*Proof.* The proof is done by induction on  $\rho := d - k$ .

If  $\rho = 0$  (i.e.,  $d = k$ ), by the definitions (3.3) of the discrete potential and (3.7) of the global discrete exterior derivative with  $k - 1$  instead of  $k$ , we have  $P_{r,f}^k (d_{r,f}^{k-1} \underline{\omega}_f) = \star^{-1} (\star d_{r,f}^{k-1} \underline{\omega}_f) = d_{r,f}^{k-1} \underline{\omega}_f$  (notice that, in the first passage, we can omit the projector found in the definition of  $d_{r,f}^{k-1} \underline{\omega}_f$  in front of  $\star d_{r,f}^{k-1} \underline{\omega}_f$  since this quantity sits in  $\mathcal{P}_r \Lambda^0(f) = \mathcal{P}_r^- \Lambda^0(f)$ , and is therefore left unchanged by  $\pi_{r,f}^{-,0}$ ). This proves (3.14), and the relation (3.15) is irrelevant here since  $d = k$ .

Let us now assume that (3.14) and (3.15) hold for a given  $\rho \geq 0$ , and let us consider  $d$  and  $k$  such that  $d - k = \rho + 1$ . We start by considering (3.15) (which we need to prove since  $d \geq k + 1$  in the present case). Let us take  $f \in \Delta_d(\mathcal{M}_h)$ . Applying (3.4) with  $d_{r,f}^{k-1} \underline{\omega}_f$  instead of  $\underline{\omega}_f$  and a generic  $\mu \in \mathcal{P}_r \Lambda^{d-k-1}(f)$ , we have, expanding the local discrete exterior derivative  $d_{r,f}^{k-1} \underline{\omega}_f$  according to its definition (i.e., the restriction to  $f$  of (3.7) with  $k - 1$  instead of  $k$ ),

$$\begin{aligned} \int_f d_{r,f}^k (d_{r,f}^{k-1} \underline{\omega}_f) \wedge \mu &= (-1)^{k+1} \int_f \star^{-1} (\pi_{r,f}^{-,d-k} (\star d_{r,f}^{k-1} \underline{\omega}_f)) \wedge d\mu \\ &\quad + \int_{\partial f} P_{r,\partial f}^k (d_{r,\partial f}^{k-1} \underline{\omega}_{\partial f}) \wedge \text{tr}_{\partial f} \mu. \end{aligned} \quad (3.16)$$

By the induction hypothesis, (3.14) holds on each  $f' \in \Delta_{d-1}(f)$  (since  $(d - 1) - k = \rho$ ), and thus

$$P_{r,\partial f}^k (d_{r,\partial f}^{k-1} \underline{\omega}_{\partial f}) = d_{r,\partial f}^{k-1} \underline{\omega}_{\partial f}. \quad (3.17)$$

Invoking then (2.4) with  $(\mathcal{X}, \omega, \mu) \leftarrow (\mathcal{P}_r^- \Lambda^{d-k}(f), d_{r,f}^{k-1} \underline{\omega}_f, d\mu)$ , noticing that  $d\mu \in \mathcal{P}_{r-1} \Lambda^{d-k}(f) \subset \mathcal{P}_r^- \Lambda^{d-k}(f)$  (by (2.14) with  $\ell = d - k$ ) to handle the first term in the right-hand side of (3.16), we infer

$$\int_f d_{r,f}^k (d_{r,f}^{k-1} \underline{\omega}_f) \wedge \mu = (-1)^{k+1} \int_f d_{r,f}^{k-1} \underline{\omega}_f \wedge d\mu + \int_{\partial f} d_{r,\partial f}^{k-1} \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \mu = 0, \quad (3.18)$$

where the conclusion follows from the link (3.12) between discrete exterior derivatives on subcells applied with  $k - 1$  instead of  $k$  and  $\alpha = \mu \in \mathcal{P}_r \Lambda^{d-k-1}(f) \subset \mathcal{P}_{r+1}^- \Lambda^{d-(k-1)-2}(f)$ . Since  $\mu$  is arbitrary in  $\mathcal{P}_r \Lambda^{d-k-1}(f)$ , (3.18) proves (3.15).

We next prove (3.14). For any  $(\mu, \nu) \in \mathcal{K}_{r+1}^{d-k-1}(f) \times \mathcal{K}_r^{d-k}(f)$ , the definition (3.5) of the potential applied to  $\underline{d}_{r,f}^{k-1} \underline{\omega}_f$  gives

$$\begin{aligned} (-1)^{k+1} \int_f P_{r,f}^k(\underline{d}_{r,f}^{k-1} \underline{\omega}_f) \wedge (d\mu + \nu) &= \int_f d_{r,f}^k(\underline{d}_{r,f}^{k-1} \underline{\omega}_f) \wedge \mu \\ &\quad - \int_{\partial f} P_{r,\partial f}^k(\underline{d}_{r,\partial f}^{k-1} \underline{\omega}_{\partial f}) \wedge \text{tr}_{\partial f} \mu + (-1)^{k+1} \int_f \star^{-1} \pi_{r,f}^{-,d-k}(\star \underline{d}_{r,f}^{k-1} \underline{\omega}_f) \wedge \nu, \end{aligned}$$

where we have additionally used, in the last term, the definition of the local discrete exterior derivative  $\underline{d}_{r,f}^{k-1} \underline{\omega}_f$ , corresponding to the restriction to  $f$  of (3.7) with  $k - 1$  instead of  $k$ . Using the complex property (3.15) that we have just proved, we have  $d_{r,f}^k(\underline{d}_{r,f}^{k-1} \underline{\omega}_f) = 0$ . Moreover, the induction hypothesis (3.17) yields  $P_{r,\partial f}^k(\underline{d}_{r,\partial f}^{k-1} \underline{\omega}_{\partial f}) = d_{r,\partial f}^{k-1} \underline{\omega}_{\partial f}$ . Hence, invoking (3.12) with  $k - 1$  instead of  $k$  and  $\alpha = \mu$  (notice that  $\mu \in \mathcal{K}_{r+1}^{d-k-1}(f) \subset \mathcal{P}_{r+1}^- \Lambda^{d-(k-1)-2}(f)$  by (2.12)) and applying (2.4) with  $(X, \omega, \mu) \leftarrow (\mathcal{P}_r^- \Lambda^{d-k}(f), d_{r,f}^{k-1} \underline{\omega}_f, \nu)$ , which is valid since  $\nu \in \mathcal{K}_r^{d-k}(f) \subset \mathcal{P}_r^- \Lambda^{d-k}(f)$  by (2.12b) with  $\ell = d - k \geq 1$ , we obtain

$$(-1)^{k+1} \int_f P_{r,f}^k(\underline{d}_{r,f}^{k-1} \underline{\omega}_f) \wedge (d\mu + \nu) = -(-1)^k \int_f d_{r,f}^{k-1} \underline{\omega}_f \wedge d\mu + (-1)^{k+1} \int_f d_{r,f}^{k-1} \underline{\omega}_f \wedge \nu.$$

Simplifying by  $(-1)^{k+1}$  and recalling the isomorphism (2.9) concludes the proof of (3.14).  $\square$

### 3.3 Commutation

The following lemma shows that the reconstructed potential  $P_{r,f}^k \underline{\omega}_f$  on a  $d$ -cell  $f$  is built by adding a high-order correction to  $\star^{-1} \omega_f$ ; this correction is designed to obtain a polynomial consistency unachievable by the component alone (see (3.9)).

**Lemma 16** (Links between component and potential reconstruction). *For all integers  $d \in [0, n]$  and  $k \leq d$ , if  $f \in \Delta_d(\mathcal{M}_h)$  and  $\underline{\omega}_f \in \underline{X}_{r,f}^k$ , then it holds*

$$\begin{aligned} (-1)^{k+1} \int_f P_{r,f}^k \underline{\omega}_f \wedge (d\mu + \nu) &= (-1)^{k+1} \int_f \star^{-1} \omega_f \wedge (d(\pi_{r,f}^{d-k-1} \mu) + \nu) \\ &\quad + \int_f d_{r,f}^k \underline{\omega}_f \wedge (\mu - \pi_{r,f}^{d-k-1} \mu) - \int_{\partial f} P_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} (\mu - \pi_{r,f}^{d-k-1} \mu) \\ &\quad \forall (\mu, \nu) \in \mathcal{K}_{r+1}^{d-k-1}(f) \times \mathcal{K}_r^{d-k}(f). \end{aligned} \quad (3.19)$$

As a consequence,

$$\pi_{r,f}^{-,d-k}(\star P_{r,f}^k \underline{\omega}_f) = \omega_f. \quad (3.20)$$

*Proof.* If  $d = k$ , the relation (3.19) follows from  $\mathcal{K}_{r+1}^{d-k-1}(f) = \mathcal{K}_{r+1}^{-1}(f) = \{0\}$  and  $P_{r,f}^k \underline{\omega}_f = \star^{-1} \omega_f$  (see (3.3)), which also establishes (3.20) since  $\pi_{r,f}^{-,0} = \text{Id}$  on  $\mathcal{P}_r \Lambda^0(f) = \mathcal{P}_r^- \Lambda^0(f)$ .

Consider now  $d \geq k + 1$  and take  $(\mu, \nu) \in \mathcal{K}_{r+1}^{d-k-1}(f) \times \mathcal{K}_r^{d-k}(f)$ . Inserting  $\pm \pi_{r,f}^{d-k-1} \mu$  into the



definition (3.5) of  $P_{r,f}^k \underline{\omega}_f$  we have

$$\begin{aligned}
& (-1)^{k+1} \int_f P_{r,f}^k \underline{\omega}_f \wedge (d\mu + \nu) \\
&= \int_f d_{r,f}^k \underline{\omega}_f \wedge \pi_{r,f}^{d-k-1} \mu + \int_f d_{r,f}^k \underline{\omega}_f \wedge (\mu - \pi_{r,f}^{d-k-1} \mu) - \int_{\partial f} P_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \mu \\
&\quad + (-1)^{k+1} \int_f \star^{-1} \omega_f \wedge \nu.
\end{aligned} \tag{3.21}$$

On the other hand, the definition (3.4) of  $d_{r,f}^k \underline{\omega}_f$  applied to  $\pi_{r,f}^{d-k-1} \mu$  yields

$$\int_f d_{r,f}^k \underline{\omega}_f \wedge \pi_{r,f}^{d-k-1} \mu = (-1)^{k+1} \int_f \star^{-1} \omega_f \wedge d(\pi_{r,f}^{d-k-1} \mu) + \int_{\partial f} P_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} (\pi_{r,f}^{d-k-1} \mu).$$

Substituting this relation into (3.21) yields (3.19).

To prove (3.20) we apply (3.19) with  $(\mu, \nu) \in \mathcal{K}_r^{d-k-1}(f) \times \mathcal{K}_r^{d-k}(f)$  and notice that  $\mu = \pi_{r,f}^{d-k-1} \mu$  since  $\mathcal{K}_r^{d-k-1}(f) \subset \mathcal{P}_r \Lambda^{d-k-1}(f)$ , to get

$$\int_f P_{r,f}^k \underline{\omega}_f \wedge (d\mu + \nu) = \int_f \star^{-1} \omega_f \wedge (d\mu + \nu). \tag{3.22}$$

The isomorphism (2.17) with  $\ell = d - k \geq 1$  shows that  $d\mu + \nu$  spans  $\mathcal{P}_r^- \Lambda^{d-k}(f)$  when  $(\mu, \nu)$  span  $\mathcal{K}_r^{d-k-1}(f) \times \mathcal{K}_r^{d-k}(f)$ . Hence, (3.22) gives

$$\int_f \star^{-1} \omega_f \wedge \alpha = \int_f P_{r,f}^k \underline{\omega}_f \wedge \alpha \stackrel{(2.4)}{=} \int_f \star^{-1} \pi_{r,f}^{-,d-k} (\star P_{r,f}^k \underline{\omega}_f) \wedge \alpha \quad \forall \alpha \in \mathcal{P}_r^- \Lambda^{d-k}(f),$$

proving (3.20) since  $\star^{-1} \omega_f \in \mathcal{P}_r^- \Lambda^{d-k}(f)$  and  $\star^{-1}$  is an isomorphism.  $\square$

**Theorem 17** (Commutation property for the local discrete exterior derivative). *For all integers  $d \in [1, n]$  and  $k \leq d - 1$ , and for all  $f \in \Delta_d(\mathcal{M}_h)$ , recalling the definition (3.2) of the interpolators, it holds*

$$\underline{d}_{r,f}^k (\underline{I}_{r,f}^k \omega) = \underline{I}_{r,f}^{k+1} (d\omega) \quad \forall \omega \in C^1 \Lambda^k(\bar{f}), \tag{3.23}$$

expressing the commutativity of the following diagram:

$$\begin{array}{ccc}
C^1 \Lambda^k(\bar{f}) & \xrightarrow{d} & C^0 \Lambda^{k+1}(\bar{f}) \\
\downarrow \underline{I}_{r,f}^k & & \downarrow \underline{I}_{r,f}^{k+1} \\
\underline{X}_{r,f}^k & \xrightarrow{\underline{d}_{r,f}^k} & \underline{X}_{r,f}^{k+1}.
\end{array}$$

*Proof.* Given the definitions of the interpolator and of the local discrete exterior derivative (see (3.2) and (3.7)), we have to prove that, for all  $f' \in \Delta_{d'}(f)$  with  $d' \in [k + 1, d]$ ,  $\pi_{r,f'}^{-,d'-k-1} (\star d_{r,f'}^k \underline{I}_{r,f'}^k \omega) = \pi_{r,f'}^{-,d'-k-1} (\star \text{tr}_{f'}(d\omega))$ . Recalling the definition of the projector  $\pi_{r,f'}^{-,d'-k-1}$  (i.e., (2.3) with  $\mathcal{X} = \mathcal{P}_r^- \Lambda^{d'-k-1}(f')$ ), we need to prove that, for any  $\mu \in \mathcal{P}_r^- \Lambda^{d'-k-1}(f')$

$$\int_{f'} \star d_{r,f'}^k \underline{I}_{r,f'}^k \omega \wedge \star \mu = \int_{f'} \star \text{tr}_{f'}(d\omega) \wedge \star \mu.$$

Applying (A.4), this amounts to proving that

$$\int_{f'} d_{r,f'}^k \underline{I}_{r,f'}^k \omega \wedge \mu = \int_{f'} \text{tr}_{f'}(d\omega) \wedge \mu. \tag{3.24}$$

Using the definitions (3.4) of the discrete exterior derivative on  $f'$  and (3.2) of  $\underline{I}_{r,f'}^k$ , we have

$$\int_{f'} d_{r,f'}^k \underline{I}_{r,f'}^k \omega \wedge \mu = (-1)^{k+1} \int_{f'} \star^{-1} \pi_{r,f'}^{-,d'-k} (\star \underline{\text{tr}}_{f'} \omega) \wedge d\mu + \int_{\partial f'} P_{r,\partial f'}^k \underline{I}_{r,\partial f'}^k \omega \wedge \text{tr}_{\partial f'} \mu, \quad (3.25)$$

where the substitution is justified by (2.4) with  $(\mathcal{X}, \omega, \mu) \leftarrow (\mathcal{P}_r^- \Lambda^{d'-k}(f'), \text{tr}_{f'} \omega, d\mu)$ , since  $d\mu \in d\mathcal{P}_r^- \Lambda^{d'-k-1}(f') \subset d\mathcal{P}_r \Lambda^{d'-k-1}(f') \subset \mathcal{P}_r^- \Lambda^{d'-k}(f')$  (see (2.12b)). For all  $f'' \in \Delta_{d'-1}(f')$  we have  $\text{tr}_{f''} \mu \in \mathcal{P}_r^- \Lambda^{d'-1-k}(f'')$  (see Lemma 4), so

$$\begin{aligned} \int_{f''} P_{r,f''}^k \underline{I}_{r,f''}^k \omega \wedge \text{tr}_{f''} \mu &= \int_{f''} \star^{-1} \pi_{r,f''}^{-,d'-1-k} (\star P_{r,f''}^k \underline{I}_{r,f''}^k \omega) \wedge \text{tr}_{f''} \mu && \text{Eq. (2.4)} \\ &= \int_{f''} \star^{-1} \pi_{r,f''}^{-,d'-1-k} (\star \text{tr}_{f''} \omega) \wedge \text{tr}_{f''} \mu && \text{Eqs. (3.20), (3.2)} \\ &= \int_{f''} \text{tr}_{f''} \omega \wedge \text{tr}_{f''} \mu. && \text{Eq. (2.4)} \end{aligned}$$

Summing this relation over  $f'' \in \Delta_{d'-1}(f')$  and substituting the result into (3.25) we obtain

$$\int_{f'} d_{r,f'}^k \underline{I}_{r,f'}^k \omega \wedge \mu = (-1)^{k+1} \int_{f'} \text{tr}_{f'} \omega \wedge d\mu + \int_{\partial f'} \text{tr}_{\partial f'} \omega \wedge \text{tr}_{\partial f'} \mu.$$

The proof of (3.24) is concluded invoking the integration by part formula (2.1) and writing  $d \text{tr}_{f'} = \text{tr}_{f'} d$  (since the trace is a pullback).  $\square$

### 3.4 Polynomial consistency

*Proof of Theorem 11.* The proof is made, as in Theorem 15, by induction on  $\rho := d - k$ .

If  $\rho = 0$ , then  $d = k$  and the definitions (3.3) of the discrete potential and (3.2) of the interpolator give  $P_{r,d}^k \underline{I}_{r,f}^k \omega = \star^{-1} \pi_{r,f}^{-,0} (\star \omega) = \star^{-1} \star \omega = \omega$ , where, to remove the projector, we have used the fact that  $\star \omega \in \mathcal{P}_r \Lambda^0(f) = \mathcal{P}_r^- \Lambda^0(f)$  (cf. (2.12a)).

Let us now assume that the lemma holds for a given  $\rho \geq 0$ , and let us consider  $d$  and  $k$  such that  $d - k = \rho + 1$ . We first look at (3.10). By the link (3.14) applied to  $k+1$  instead of  $k$  and the commutation property (3.23), we have, for  $\omega \in \mathcal{P}_{r+1}^- \Lambda^k(f)$ ,

$$d_{r,f}^k \underline{I}_{r,f}^k \omega = P_{r,f}^{k+1} (d_{r,f}^k \underline{I}_{r,f}^k \omega) = P_{r,f}^{k+1} \underline{I}_{r,f}^{k+1} (d\omega).$$

We have  $d\omega \in d\mathcal{P}_{r+1}^- \Lambda^k(f) \subset \mathcal{P}_r \Lambda^{k+1}(f)$ , and the pair  $(d, k+1)$  satisfies  $d - (k+1) = \rho$ . We can therefore apply the induction hypothesis to see that (3.9) holds for this pair and  $d\omega$  instead of  $\omega$ ; this gives  $d_{r,f}^k \underline{I}_{r,f}^k \omega = d\omega$  and proves (3.10).

We now turn to (3.9). For  $\omega \in \mathcal{P}_r \Lambda^k(f)$ , applying the property (3.10) that we have just proved to  $\omega$  and recalling the definitions (3.5) and (3.2) of the potential and of the interpolator, we find, for all  $(\mu, \nu) \in \mathcal{K}_{r+1}^{d-k-1}(f) \times \mathcal{K}_r^{d-k}(f)$ ,

$$\begin{aligned} &(-1)^{k+1} \int_f P_{r,f}^k \underline{I}_{r,f}^k \omega \wedge (d\mu + \nu) \\ &= \int_f d\omega \wedge \mu - \int_{\partial f} P_{r,\partial f}^k \underline{I}_{r,\partial f}^k \text{tr}_{\partial f} \omega \wedge \text{tr}_{\partial f} \mu + (-1)^{k+1} \int_f \star^{-1} \pi_{r,f}^{-,d-k} (\star \omega) \wedge \nu, \end{aligned}$$

the replacement being justified by (2.4) and  $\nu \in \mathcal{P}_r^- \Lambda^{d-k}(f)$  (see (2.12)). We can then apply the polynomial consistency (3.9) on each  $f' \in \Delta_{d-1}(f)$  (as  $(d-1) - k = \rho$ ) to write  $P_{r,\partial f}^k \underline{I}_{r,\partial f}^k \text{tr}_{\partial f} \omega = \text{tr}_{\partial f} \omega$ , and then integrate by parts to conclude, since  $\mu, \nu$  are generic elements, that  $P_{r,f}^k \underline{I}_{r,f}^k \omega = \omega$ .  $\square$

*Remark 18* (Consistency property of the improved potential for  $k = 0$ ). In the case  $k = 0$ , the improved potential defined in Remark 7 satisfies the following consistency property:

$$P_{r+1,f}^0 I_{r,f}^0 \omega = \omega \quad \forall \omega \in \mathcal{P}_{r+1} \Lambda^0(f).$$

To see this, first notice that when  $d = k = 0$  we have  $P_{r+1,f}^0 = P_{r,f}^0$  since  $\mathcal{P}_{r+1} \Lambda^0(f) = \mathcal{P}_r \Lambda^0(f) \cong \mathbb{R}$ , and then, for  $d \geq k + 1$ , invoke the definition (3.6) of  $P_{r+1,f}^0 I_{r,f}^0 \omega$ , apply (3.10) (since  $\mathcal{P}_{r+1}^- \Lambda^0(f) = \mathcal{P}_{r+1} \Lambda^0(f)$ ) and a recursion argument on  $d$ .

### 3.5 Cohomology

A strategy to establish the exactness of the de Rham complex (for a domain with trivial topology) is to design a Poincaré operator  $p : C^1 \Lambda^k(\bar{\Omega}) \rightarrow C^1 \Lambda^{k-1}(\bar{\Omega})$ , that satisfies  $dp + pd = \text{Id}$ . The Poincaré operator is built integrating a certain flow of contracted differential forms; see [31, 51] for details and applications to the design of finite element complexes. Extending such a construction to the context of fully discrete spaces is not trivial, as it is not clear how the discrete polynomial components on cells should evolve with such a flow. We therefore select an alternative approach, more suited to hierarchical discrete spaces.

The starting point is the following idea: if  $\eta \in C^1 \Lambda^k(\bar{\Omega})$  satisfies  $d\eta = 0$  and we have  $\omega \in C^2 \Lambda^{k-1}(\bar{\Omega})$  such that  $d\omega = \eta$ , then (2.1) shows that, for any  $d$ -cell  $f$ ,

$$(-1)^k \int_f \omega \wedge d\mu = \int_f \eta \wedge \mu - \int_{\partial f} \text{tr}_{\partial f} \omega \wedge \text{tr}_{\partial f} \mu \quad \forall \mu \in C^1 \Lambda^{d-k}(\bar{\Omega}). \quad (3.26)$$

In the discrete setting,  $\omega$  is built starting from the lowest-dimensional cells, and (3.26) thus gives a condition on  $\omega$  over  $f$  based on the already constructed  $\text{tr}_{\partial f} \omega$ . To start this process, we must fix the values of  $\omega$  on the lowest-dimensional cells, which is not an easy task in general. Actually, from the point of view of differential forms, the lowest-dimensional cells encode the topology of the domain, and thus the cohomology of the complex; for a generic  $\eta$ , the recursive construction of  $\omega$  can therefore only be fully complete if the complex is exact, and thus the topology trivial.

This limitation is circumvented by using the following idea: if  $\eta$  has zero average on  $k$ -cells, then  $\omega$  can be set to zero on  $(k - 1)$ -cells, which completes the construction above (see Lemma 19 below). This result is then exploited, through the extension/reduction strategy developed in [35, 37], to compare the cohomology of the arbitrary-order DDR( $r$ ) complex to that of the lowest-order DDR(0) complex, which is trivially isomorphic to the CW complex based on the mesh.

We therefore start by considering the subspace  $\underline{X}_{r,h,b}^k$  of  $\underline{X}_{r,h}^k$  made of vectors of differential forms whose integrals over cells of dimension  $d = k$  vanish:

$$\underline{X}_{r,h,b}^k := \left\{ \underline{\omega}_h = (\omega_f)_{f \in \Delta_d(\mathcal{M}_h), d \in [k,n]} : \int_f \star^{-1} \omega_f = 0 \quad \forall f \in \Delta_k(\mathcal{M}_h) \right\}.$$

**Lemma 19** (Exactness property for  $\underline{X}_{r,h,b}^k$ ). *For any integer  $k \in [0, n]$ , if  $\underline{\eta}_h \in \underline{X}_{r,h,b}^k$  satisfies  $\underline{d}_{r,h}^k \underline{\eta}_h = \underline{0}$ , then there exists  $\underline{\omega}_h \in \underline{X}_{r,h,b}^{k-1}$  such that  $\underline{\eta}_h = \underline{d}_{r,h}^{k-1} \underline{\omega}_h$ , where, in accordance with (3.8), we have set  $\underline{d}_{r,h}^{-1} = \underline{d}_{r,h}^n := \underline{0}$ .*

*Remark 20* (Exact sub-complex). It can easily be checked that  $\underline{d}_{r,h}^k : \underline{X}_{r,h,b}^k \rightarrow \underline{X}_{r,h,b}^{k+1}$ . As a consequence, the previous lemma shows that  $(\underline{X}_{r,h,b}^k, \underline{d}_{r,h}^k)_k$  is an exact sub-complex of DDR( $r$ ) (even if the latter complex is not exact).

*Proof.* We first notice that the case  $r = 0$  is trivial since, for all  $k$ ,  $\underline{X}_{0,h,b}^k = \{(0)_{f \in \Delta_k(\mathcal{M}_h)}\}$ . This comes from the fact that the space  $\underline{X}_{0,h}^k$  only has non-zero components (which are moreover constant) on cells

of dimension  $d = k$ ; to check this, notice that the spaces (2.12b) are all trivial since the first component vanishes for  $k$ -forms with constant coefficients, while the second is zero by (2.6). We can therefore assume that  $r \geq 1$ . The cases  $k = 0$  and  $k \geq 1$  have to be handled separately.

Case  $k = 0$ . We prove that, if  $\underline{\eta}_h \in \underline{X}_{r,h,b}^0$  and  $\underline{d}_{r,h}^0 \underline{\eta}_h = \underline{0}$ , then  $\eta_f = 0$  for all  $f \in \Delta_d(\mathcal{M}_h)$ ,  $d \in [0, n]$ . This is done by induction on  $d$ . The case  $d = 0$  follows immediately from the definition of  $\underline{X}_{r,h,b}^0$  which shows that the value of  $\star^{-1}\eta_f$  on any vertex  $f \in \Delta_0(\mathcal{M}_h)$  is zero. Assuming that all components of  $\underline{\eta}_h$  on cells of dimension  $d - 1 \geq 0$  vanish, we now prove that  $\eta_f = 0$  for all  $f \in \Delta_d(\mathcal{M}_h)$ . Note first that, by (3.14), the property  $\underline{d}_{r,f}^0 \underline{\eta}_f = \underline{0}$  implies  $\underline{d}_{r,f}^0 \eta_f = 0$ . Enforcing then  $\underline{\eta}_{\partial f} = \underline{0}$  (by induction hypothesis) in the definition (3.4) of  $\underline{d}_{r,f}^0 \eta_f$  gives

$$\int_f \star^{-1}\eta_f \wedge d\mu = 0 \quad \forall \mu \in \mathcal{P}_r \Lambda^{d-1}(f).$$

By definition (2.12b) of the trimmed space with  $\ell = d$ , and accounting for (2.6), we have  $d\mathcal{P}_r \Lambda^{d-1}(f) = \mathcal{P}_r^- \Lambda^d(f)$ , so the relation above and (A.4) with  $(\omega, \mu) \leftarrow (\eta_f, d\mu)$  and  $\rho = d\mu$  show that  $\int_f \eta_f \wedge \star \rho = 0$  for all  $\rho \in \mathcal{P}_r^- \Lambda^d(f)$ . Since  $\eta_f$  belongs to that space, we conclude that  $\eta_f = 0$ .

Case  $k \geq 1$ . Let  $\underline{\eta}_h \in \underline{X}_{r,h,b}^k$  be such that  $\underline{d}_{r,h}^k \underline{\eta}_h = \underline{0}$ , and let us construct  $\underline{\omega}_h \in \underline{X}_{r,h,b}^{k-1}$  such that  $\underline{d}_{r,h}^{k-1} \underline{\omega}_h = \underline{\eta}_h$ . This construction of  $\underline{\omega}_h$  is done by increasing dimension  $d \in [k-1, n]$  of the cells. For all  $f \in \Delta_{k-1}(\mathcal{M}_h)$ , we set  $\omega_f = 0$  (which ensures, in particular, that the zero-average condition embedded in the space  $\underline{X}_{r,h,b}^{k-1}$  is fulfilled). Assume now that the components of  $\underline{\omega}_h$  have been constructed up to cells of dimension  $d - 1 \geq k - 1$ , and consider  $f \in \Delta_d(\mathcal{M}_h)$ . We choose  $\omega_f \in \mathcal{P}_r^- \Lambda^{d-k+1}(f)$  such that the following relation holds:

$$(-1)^k \int_f \star^{-1}\omega_f \wedge d\mu = \int_f P_{r,f}^k \eta_f \wedge \mu - \int_{\partial f} P_{r,\partial f}^{k-1} \omega_{\partial f} \wedge \text{tr}_{\partial f} \mu \quad \forall \mu \in \mathcal{K}_r^{d-k}(f). \quad (3.27)$$

Notice that, since the construction is recursive on the dimension of the cells,  $\omega_{\partial f}$  has already been constructed at this stage. Owing to the isomorphism (2.17) with  $\ell = d - k + 1 \geq 1$ , this relation completely defines the projection of  $\omega_f$  on  $d\mathcal{K}_r^{d-k}(f) \subset \mathcal{P}_r^- \Lambda^{d-k+1}(f)$ . The projection of  $\omega_f$  on the remaining component  $\mathcal{K}_r^{d-k+1}(f)$  of  $\mathcal{P}_r^- \Lambda^{d-k+1}(f)$  is not relevant to the rest of the proof and can be set to 0.

Let us now prove that  $\underline{d}_{r,h}^{k-1} \underline{\omega}_h = \underline{\eta}_h$ . It suffices to show that

$$\underline{d}_{r,f}^{k-1} \omega_f = P_{r,f}^k \eta_f \quad \forall f \in \Delta_d(\mathcal{M}_h), \quad d \in [k, n]. \quad (3.28)$$

Indeed, applying  $\pi_{r,f}^{-,d-k} \star$  to this relation and using (3.20) yields  $\pi_{r,f}^{-,d-k} (\star \underline{d}_{r,f}^{k-1} \omega_f) = \eta_f$ ; using this relation for all cells  $f$ , and recalling the definition (3.7) of the global discrete exterior derivative (with  $k - 1$  instead of  $k$ ), then gives  $\underline{d}_{r,h}^{k-1} \underline{\omega}_h = \underline{\eta}_h$  as claimed.

The relation (3.28) is a direct consequence of the following property, for the same  $f$  and  $d$ :

$$\int_f \underline{d}_{r,f}^{k-1} \omega_f \wedge \mu = \int_f P_{r,f}^k \eta_f \wedge \mu \quad \forall \mu \in \mathcal{P}_r \Lambda^{d-k}(f). \quad (3.29)$$

Owing to (2.7), we only need to prove this relation for  $\mu \in \mathcal{K}_r^{d-k}(f)$ , and  $\mu \in \mathcal{P}_0 \Lambda^0(f)$  if  $d = k$  or  $\mu \in d\mathcal{P}_{r+1} \Lambda^{d-k-1}(f)$  if  $d \geq k + 1$ .

If  $\mu \in \mathcal{K}_r^{d-k}(f)$ , the definition (3.4) of  $\underline{d}_{r,f}^{k-1} \omega_f$  together with the property (3.27) immediately give (3.29).

Let us consider the case  $d = k$  and  $\mu \in \mathcal{P}_0 \Lambda^0(f)$ . Then  $d\mu = 0$ , so the definition (3.4) of  $\underline{d}_{r,f}^{k-1} \underline{\omega}_f$  and  $\underline{\omega}_{\partial f} = 0$  (by construction,  $\underline{\omega}_h$  vanishes on cells of dimension  $d - 1 = k - 1$ ) show that the left-hand side of (3.29) vanishes. Since  $P_{r,f}^k \underline{\eta}_f = \star^{-1} \eta_f$  (see (3.3)) and  $\int_f \star^{-1} \eta_f = 0$  as  $\underline{\eta}_h \in \underline{X}_{r,h,b}^k$ , the right-hand side of (3.29) vanishes as well, and this relation holds.

Finally, we turn to the case  $d \geq k + 1$  and  $\mu \in d\mathcal{P}_{r+1} \Lambda^{d-k-1}(f)$ , which is proved by induction on  $d$  (the base case  $d = k$  having already been covered). By (2.8) with  $(\ell, r) \leftarrow (d - k - 1, r + 1)$ , we have  $\mu \in d\mathcal{K}_{r+1}^{d-k-1}(f)$ , and we can therefore write  $\mu = d\alpha$  with  $\alpha \in \mathcal{K}_{r+1}^{d-k-1}(f) \subset \mathcal{P}_{r+1}^- \Lambda^{d-k-1}(f)$  (see (2.12)). Invoking the link (3.12) between discrete exterior derivatives on subcells (notice that  $d \geq (k - 1) + 2$ ), we obtain

$$\int_f \underline{d}_{r,f}^{k-1} \underline{\omega}_f \wedge \mu = (-1)^k \int_{\partial f} \underline{d}_{r,\partial f}^{k-1} \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \alpha = (-1)^k \int_{\partial f} P_{r,\partial f}^k \underline{\eta}_{\partial f} \wedge \text{tr}_{\partial f} \alpha,$$

where the second equality follows from the induction hypothesis that (3.29) holds on subcells of  $f$ . We have  $\underline{d}_{r,f}^k \underline{\eta}_f = 0$  and  $d \geq k + 1$ , so we can apply (3.14) with  $k + 1$  instead of  $k$  to get  $\underline{d}_{r,f}^k \underline{\eta}_f = 0$ ; the definition (3.5) of  $P_{r,f}^k \underline{\eta}_f$  (with  $(\mu, \nu) \leftarrow (\alpha, 0)$ , see Remark 6 for the validity of this choice of  $\mu$ ) allows us to continue with

$$\int_f \underline{d}_{r,f}^{k-1} \underline{\omega}_f \wedge \mu = -(-1)^k \times (-1)^{k+1} \int_f P_{r,f}^k \underline{\eta}_f \wedge d\alpha.$$

Recalling that  $d\alpha = \mu$  concludes the proof of (3.29).  $\square$

*Proof of Theorem 10.* As in [37, Lemma 4], it is straightforward to see that the (discrete) de Rham map establishes a chain isomorphism between the lowest-degree complex  $\text{DDR}(0)$  and the CW complex defined by  $\mathcal{M}_h$ . Since this CW complex has the same cohomology as the de Rham complex (1.2), the proof is complete if we show that the cohomology of  $\text{DDR}(r)$  is isomorphic to the cohomology of  $\text{DDR}(0)$ . This obviously means that we can assume  $r \geq 1$  in the following.

*Step 1: Reductions and extensions.* With the goal of applying [35, Proposition 2], we define reduction and extension maps between  $\text{DDR}(r)$  and  $\text{DDR}(0)$  as in (3.30).

$$\begin{array}{ccccccc} \text{DDR}(r) : & \cdots & \longrightarrow & \underline{X}_{r,h}^k & \xrightarrow{\underline{d}_{r,h}^k} & \underline{X}_{r,h}^{k+1} & \longrightarrow \cdots \\ & & & \begin{array}{c} \uparrow \\ E_h^k \end{array} & & \begin{array}{c} \uparrow \\ E_h^{k+1} \end{array} & \\ & & & \left( \begin{array}{c} \uparrow \\ \downarrow \end{array} \right) & & \left( \begin{array}{c} \uparrow \\ \downarrow \end{array} \right) & \\ & & & R_h^k & & R_h^{k+1} & \\ \text{DDR}(0) : & \cdots & \longrightarrow & \underline{X}_{0,h}^k & \xrightarrow{\underline{d}_{0,h}^k} & \underline{X}_{0,h}^{k+1} & \longrightarrow \cdots \end{array} \quad (3.30)$$

The reduction  $\underline{R}_h^k : \underline{X}_{r,h}^k \rightarrow \underline{X}_{0,h}^k$  is defined taking the average of components on the cells of dimension  $k$  (recall that vectors in  $\underline{X}_{0,h}^k$  only have components on these cells): For all  $\underline{\omega}_h \in \underline{X}_{r,h}^k$ ,

$$\underline{R}_h^k \underline{\omega}_h = (\pi_{0,f}^0 \omega_f)_{f \in \Delta_k(\mathcal{M}_h)}. \quad (3.31)$$

The extension  $\underline{E}_h^k : \underline{X}_{0,h}^k \rightarrow \underline{X}_{r,h}^k$  is defined by induction on the cell dimension: For all  $\underline{\eta}_h \in \underline{X}_{0,h}^k$ ,  $\underline{E}_h^k \underline{\eta}_h := (E_f^k \eta_f)_{f \in \Delta_d(\mathcal{M}_h), d \in [k, n]}$ , where

- If  $d = k$ ,

$$E_f^k \eta_f = \eta_f \in \mathcal{P}_0^- \Lambda^0(f) \subset \mathcal{P}_r^- \Lambda^0(f); \quad (3.32a)$$

- If  $d \geq k + 1$ ,  $E_f^k \underline{\eta}_f \in \mathcal{P}_r^- \Lambda^{d-k}(f)$  satisfies

$$\begin{aligned} & (-1)^{k+1} \int_f \star^{-1} E_f^k \underline{\eta}_f \wedge (d\mu + \nu) \\ &= \int_f \underline{d}_{0,f}^k \underline{\eta}_f \wedge \mu - \int_{\partial f} P_{r,\partial f}^k \underline{E}_{\partial f}^k \underline{\eta}_{\partial f} \wedge \text{tr}_{\partial f} \mu + (-1)^{k+1} \int_f P_{0,f}^k \underline{\eta}_f \wedge \nu \\ & \quad \forall (\mu, \nu) \in \mathcal{K}_r^{d-k-1}(f) \times \mathcal{K}_r^{d-k}(f), \end{aligned} \quad (3.32b)$$

where  $\underline{E}_{\partial f}^k \underline{\eta}_{\partial f} = (\underline{E}_{f'}^k \underline{\eta}_{f'})_{f' \in \Delta_{d-1}(f)}$  gathers the extensions already built at previous steps on the subcells of dimension  $d - 1$  of  $f$ . The isomorphism (2.17) with  $\ell = d - k \geq 1$  ensures that the relation above fully and properly defines  $E_f^k \underline{\eta}_f$ .

In passing, we notice that, for all  $f \in \Delta_d(\mathcal{M}_h)$  with  $d \geq k + 1$ , combining the definitions (3.4) of  $\underline{d}_{r,f}^k \underline{E}_f^k \underline{\eta}_f$  and (3.32b) of  $E_f^k \underline{\eta}_f$  (with  $\nu = 0$ ), we have

$$\int_f \underline{d}_{r,f}^k \underline{E}_f^k \underline{\eta}_f \wedge \mu = \int_f \underline{d}_{0,f}^k \underline{\eta}_f \wedge \mu \quad \forall \mu \in \mathcal{K}_r^{d-k-1}(f).$$

Since  $\underline{d}_{0,f}^k \underline{\eta}_f = P_{0,f}^{k+1} \underline{d}_{0,f}^k \underline{\eta}_f$  by (3.14), using the definition of  $E_f^{k+1} \underline{d}_{0,f}^k \underline{\eta}_f$  (namely, (3.32a) if  $d = k + 1$ , or (3.32b) with  $(\mu, \nu) \leftarrow (0, \mu)$  if  $d \geq k + 2$ ) we deduce that

$$\int_f \underline{d}_{r,f}^k \underline{E}_f^k \underline{\eta}_f \wedge \mu = \int_f \star^{-1} E_f^{k+1} \underline{d}_{0,f}^k \underline{\eta}_f \wedge \mu \quad \forall \mu \in \mathcal{K}_r^{d-k-1}(f). \quad (3.33)$$

Step 2: Proof of the theorem. To apply [35, Proposition 2], we need to prove that

$$\underline{d}_{0,h}^k = \underline{R}_h^{k+1} \underline{d}_{r,h}^k \underline{E}_h^k \quad (3.34)$$

and that [35, Assumption 1] holds, that is:

- (C1)  $\underline{R}_h^k \underline{E}_h^k = \text{Id}$  on  $\text{Ker } \underline{d}_{0,h}^k$ ;
- (C2)  $(\underline{E}_h^k \underline{R}_h^k - \text{Id})(\text{Ker } \underline{d}_{r,h}^k) \subset \text{Im } \underline{d}_{r,h}^{k-1}$ ;
- (C3) The graded maps  $\underline{E}_h^\bullet$  and  $\underline{R}_h^\bullet$  are cochain maps.

We start by noticing that, since  $\text{DDR}(0)$  is already known to be a complex, (C1) and (C3) imply (3.34). Indeed, (C3) gives  $\underline{R}_h^{k+1} \underline{d}_{r,h}^k \underline{E}_h^k = \underline{R}_h^{k+1} \underline{E}_h^{k+1} \underline{d}_{0,h}^k$  and, by the complex property,  $\text{Im } \underline{d}_{0,h}^k \subset \text{Ker } \underline{d}_{0,h}^{k+1}$ , so (C1) applied to  $k + 1$  instead of  $k$  yields (3.34).

1. Proof of (C1). The definitions (3.31) and (3.32a) of the reduction and the extension components on the lowest dimensional cells directly shows that  $\underline{R}_h^k \underline{E}_h^k = \text{Id}$  on  $\underline{X}_{r,h}^k$ , which establishes a stronger result than (C1).

2. Proof of (C3) for the extension. We now turn to (C3), considering first the case of the extension. We have to show that, for all  $\underline{\eta}_h \in \underline{X}_{0,h}^k$  it holds  $\underline{d}_{r,h}^k \underline{E}_h^k \underline{\eta}_h = \underline{E}_h^{k+1} \underline{d}_{0,h}^k \underline{\eta}_h$ . Given the definitions (3.7) of the global discrete exterior derivative and of the extension, this boils down to showing that

$$\star^{-1} \pi_{r,f}^{-,d-k-1} (\star \underline{d}_{r,f}^k \underline{E}_f^k \underline{\eta}_f) = \star^{-1} E_f^{k+1} \underline{d}_{0,f}^k \underline{\eta}_f \quad \forall f \in \Delta_d(\mathcal{M}_h) \text{ with } d \geq k + 1$$

which, testing against  $\rho \in \mathcal{P}_r^- \Lambda^{d-k-1}(f)$  and recalling the relation (2.4), can be recast as

$$\int_f \mathbf{d}_{r,f}^k \underline{E}_f^k \underline{\eta}_f \wedge \rho = \int_f \star^{-1} E_f^{k+1} \underline{\mathbf{d}}_{0,f}^k \underline{\eta}_f \wedge \rho \quad \forall f \in \Delta_d(\mathcal{M}_h) \text{ with } d \geq k+1, \quad (3.35)$$

$$\forall \rho \in \mathcal{P}_r^- \Lambda^{d-k-1}(f).$$

We start by noticing that, by (3.33), the relation (3.35) holds for  $\rho \in \mathcal{K}_r^{d-k-1}(f)$ . The decompositions (2.7a) of  $\mathcal{P}_r \Lambda^0(f) = \mathcal{P}_r^- \Lambda^0(f)$  (if  $d = k+1$ ) and (2.16) of  $\mathcal{P}_r^- \Lambda^{d-k-1}(f)$  (if  $d \geq k+2$ ) then show that we only have to prove (3.35) for  $\rho \in \mathcal{P}_0 \Lambda^0(f)$  (if  $d = k+1$ ) or  $\rho \in \mathcal{d}\mathcal{K}_r^{d-k-2}(f)$  (if  $d \geq k+2$ ). This fact is proved by induction on  $d$ :

- Let us first consider  $d = k+1$  and take  $\rho \in \mathcal{P}_0 \Lambda^0(f)$ . We can use this polynomial form as a test function in the definition (3.4) of  $\mathbf{d}_{0,f}^k \underline{\eta}_f$  to get

$$\int_{\partial f} P_{0,\partial f}^k \underline{\eta}_{\partial f} \wedge \text{tr}_{\partial f} \rho = \int_f \mathbf{d}_{0,f}^k \underline{\eta}_f \wedge \rho = \int_f \star^{-1} E_f^{k+1} \underline{\mathbf{d}}_{0,f}^k \underline{\eta}_f \wedge \rho \quad \forall \rho \in \mathcal{P}_0 \Lambda^0(f), \quad (3.36)$$

where the second equality follows from (3.32a) with  $(k, \underline{\eta}_f) \leftarrow (k+1, \underline{\mathbf{d}}_{0,f}^k \underline{\eta}_f)$ . For all  $f' \in \Delta_k(f)$ , by definition (3.3) of  $P_{0,f'}^k$  and (3.32a) of  $E_{f'}^k$ , we have  $P_{0,f'}^k \underline{\eta}_{f'} = \star^{-1} \eta_{f'} = \star^{-1} E_{f'}^k \underline{\eta}_{f'} = P_{r,f'}^k \underline{E}_{f'}^k \underline{\eta}_{f'}$ , where the last relation follows applying the definition (3.3) of  $P_{r,f'}^k$ . We infer from this equality and (3.36) that

$$\int_{\partial f} P_{r,\partial f}^k \underline{E}_{\partial f}^k \underline{\eta}_{\partial f} \wedge \text{tr}_{\partial f} \rho = \int_f \star^{-1} E_f^{k+1} \underline{\mathbf{d}}_{0,f}^k \underline{\eta}_f \wedge \rho \quad \forall \rho \in \mathcal{P}_0 \Lambda^0(f).$$

Applying the definition (3.4) of  $\mathbf{d}_{r,f}^k \underline{E}_f^k \underline{\eta}_{\partial f}$  with  $\mu = \rho$  (which satisfies  $\mathbf{d}\rho = 0$ ) to the left-hand side then concludes the proof of (3.35).

- We now take  $d \geq k+2$  and  $\rho \in \mathcal{d}\mathcal{K}_r^{d-k-2}(f)$ , which we write  $\rho = \mathbf{d}\alpha$  with  $\alpha \in \mathcal{K}_r^{d-k-2}(f) \subset \mathcal{P}_r^- \Lambda^{d-k-2}(f)$ . Applying the link (3.12) between discrete exterior derivatives on  $f$  and  $\partial f$ , we have

$$\int_f \mathbf{d}_{r,f}^k \underline{E}_f^k \underline{\eta}_f \wedge \rho = (-1)^{k+1} \int_{\partial f} \mathbf{d}_{r,\partial f}^k \underline{E}_{\partial f}^k \underline{\eta}_{\partial f} \wedge \text{tr}_{\partial f} \alpha. \quad (3.37)$$

By Lemma 4, for all  $f' \in \Delta_{d-1}(f)$ ,  $\text{tr}_{f'} \alpha \in \mathcal{P}_r^- \Lambda^{d-k-2}(f')$ , so we can apply (3.35) on  $f'$  (by the induction hypothesis) to get

$$\int_{f'} \mathbf{d}_{r,f'}^k \underline{E}_{f'}^k \underline{\eta}_{f'} \wedge \text{tr}_{f'} \alpha = \int_{f'} \star^{-1} E_{f'}^{k+1} \underline{\mathbf{d}}_{0,f'}^k \underline{\eta}_{f'} \wedge \text{tr}_{f'} \alpha = \int_{f'} P_{r,f'}^{k+1} \underline{E}_{f'}^{k+1} \underline{\mathbf{d}}_{0,f'}^k \underline{\eta}_{f'} \wedge \text{tr}_{f'} \alpha,$$

the second equality being justified by (3.20) and (2.4) (with  $(\mathcal{X}, f, d, k) \leftarrow (\mathcal{P}_r^- \Lambda^{(d-1)-(k+1)}(f'), f', d-1, k+1)$ ) and the fact that  $\text{tr}_{f'} \alpha \in \mathcal{P}_r^- \Lambda^{(d-1)-(k+1)}(f')$ . Plugging this relation into (3.37) yields

$$\int_f \mathbf{d}_{r,f}^k \underline{E}_f^k \underline{\eta}_f \wedge \rho = (-1)^{k+1} \int_{\partial f} P_{r,\partial f}^{k+1} \underline{E}_{\partial f}^{k+1} \underline{\mathbf{d}}_{0,\partial f}^k \underline{\eta}_{\partial f} \wedge \text{tr}_{\partial f} \alpha.$$

Invoking then the definition (3.32b) of  $E_f^{k+1} \underline{\mathbf{d}}_{0,f}^k \underline{\eta}_f$  with  $(k, \mu, \nu, \underline{\eta}_f) \leftarrow (k+1, \alpha, 0, \underline{\mathbf{d}}_{0,f}^k \underline{\eta}_f)$ , and using the property  $\mathbf{d}_{0,f}^{k+1} \circ \underline{\mathbf{d}}_{0,f}^k = 0$  (consequence of (3.15) with  $k+1$  instead of  $k$ ) we infer

$$\int_f \mathbf{d}_{r,f}^k \underline{E}_f^k \underline{\eta}_f \wedge \rho = \int_f \star^{-1} E_f^{k+1} \underline{\mathbf{d}}_{0,f}^k \underline{\eta}_f \wedge \mathbf{d}\alpha$$

and (3.35) follows by recalling that  $\rho = \mathbf{d}\alpha$ .

3. Proof of (C3) for the reduction. To conclude the proof of (C3), it remains to show that  $\underline{R}_h^{k+1} \underline{d}_{r,h}^k \underline{\omega}_h = \underline{d}_{0,h}^k \underline{R}_h^k \underline{\omega}_h$  for all  $\underline{\omega}_h \in \underline{X}_{r,h}^k$ . Since vectors in  $\underline{X}_{0,h}^{k+1}$  only have constant components on cells of dimension  $k+1$ , and since  $\underline{R}_h^{k+1}$  is defined by (3.31), we only have to show that

$$\int_f \underline{d}_{r,f}^k \underline{\omega}_f \wedge \rho = \int_f \underline{d}_{0,f}^k \underline{R}_f^k \underline{\omega}_f \wedge \rho \quad \forall f \in \Delta_{k+1}(\mathcal{M}_h), \forall \rho \in \mathcal{P}_0 \Lambda^0(f). \quad (3.38)$$

Let  $\rho$  as above and apply the definition (3.4) of  $\underline{d}_{r,f}^k \underline{\omega}_f$  to  $\mu = \rho$ ; accounting for  $d\rho = 0$ , we obtain

$$\int_f \underline{d}_{r,f}^k \underline{\omega}_f \wedge \rho = \int_{\partial f} P_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \rho. \quad (3.39)$$

For each  $f' \in \Delta_k(f)$ , by definition (3.3) of  $P_{r,f'}^k$ , we can write

$$\int_{f'} P_{r,f'}^k \underline{\omega}_{f'} \wedge \text{tr}_{f'} \rho = \int_{f'} \star^{-1} \omega_f \wedge \text{tr}_{f'} \rho = \int_{f'} \star^{-1} \pi_{0,f}^0 \omega_f \wedge \text{tr}_{f'} \rho = \int_{f'} P_{0,f'}^k \underline{R}_{f'}^k \underline{\omega}_{f'} \wedge \text{tr}_{f'} \rho, \quad (3.40)$$

where we have used the fact that  $\text{tr}_{f'} \rho \in \mathcal{P}_0 \Lambda^0(f')$  to insert the projector in the second equality and the definitions (3.31) of  $\underline{R}_{f'}^k$  and (3.3) of  $P_{0,f'}^k$ , to conclude. Combining (3.39) and (3.40), we find

$$\int_f \underline{d}_{r,f}^k \underline{\omega}_f \wedge \rho = \int_{\partial f} P_{0,\partial f}^k \underline{R}_{\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \rho.$$

Applying the definition (3.4) of  $\underline{d}_{0,f}^k \underline{R}_f^k \underline{\omega}_f$  then concludes the proof of (3.38).

4. Proof of (C2). Finally, to prove (C2), we notice that if  $\underline{\omega}_h \in \underline{X}_{r,h}^k$ , then by (3.31) and (3.32a) the components of  $\underline{E}_h^k \underline{R}_h^k \underline{\omega}_h$  on the lowest dimensional cells  $f \in \Delta_k(\mathcal{M}_h)$  are just the averages of the components of  $\underline{\omega}_h$  on these cells; hence,  $\underline{E}_h^k \underline{R}_h^k \underline{\omega}_h - \underline{\omega}_h \in \underline{X}_{r,h,b}^k$ . Moreover, by the cochain map property (C3),  $\underline{d}_{r,h}^k (\underline{E}_h^k \underline{R}_h^k \underline{\omega}_h - \underline{\omega}_h) = \underline{E}_h^k \underline{R}_h^k \underline{d}_{r,h}^k \underline{\omega}_h - \underline{d}_{r,h}^k \underline{\omega}_h = \underline{0}$  whenever  $\underline{\omega}_h \in \text{Ker } \underline{d}_{r,h}^k$ . We can thus, for such an  $\underline{\omega}_h$ , apply Lemma 19 with  $\underline{\omega}_h \leftarrow \underline{E}_h^k \underline{R}_h^k \underline{\omega}_h - \underline{\omega}_h$  to see that this element belongs to  $\text{Im } \underline{d}_{r,h}^{k-1}$ , establishing (C2).  $\square$

## 4 A VEM-inspired complex

In this section we consider an alternative construction inspired by the Virtual Element complex of [8]. Notice that we make here no effort to reduce the polynomial degree of certain components of the discrete spaces, which is known to be possible; see, e.g., [9] and also [35] for a general framework with application to DDR methods. Notice also that we work in a fully discrete spirit, without attempting to identify the underlying virtual spaces (which are not needed for the purposes of the present work).

Let again a polynomial degree  $r \geq 0$  be fixed. The general principle to design the VEM-inspired sequence is to select polynomial components that make it possible to reconstruct, for each  $d$ -cell and inductively on the dimension  $d$ , a discrete potential capable of reproducing polynomial forms in  $\mathcal{P}_{r+1}^- \Lambda^k(f)$ . The main difference with respect to the DDR approach illustrated in Section 3 is that, with the exception of  $(k+1)$ -cells, the required information on the discrete exterior derivative is directly encoded in the discrete spaces.

Adopting this approach has several, far-reaching, consequences. The first one is that the discrete spaces contain a mix of both traces and exterior derivatives (which, in passing, requires higher regularity in the definition of the interpolators). The components on  $k$ - and  $(k+1)$ -cells in the discrete space of  $k$ -forms play a slightly different role than the others (and are, as a result, treated separately in the definition of the space). The second consequence is that the proofs of key properties (polynomial



| $k \backslash d$ | 0   | 1                                   | 2  | 3  |
|------------------|---|-------------------------------------|--|--|
| 0                | $\mathbb{R} = \mathcal{P}_r \Lambda^0(f_0)$ | $\{0\} \times \mathcal{K}_r^0(f_1)$ | $\{0\} \times \mathcal{K}_{r+1}^1(f_2)$                | $\{0\} \times \mathcal{K}_{r+1}^2(f_3)$                    |
| 1                |   | $\mathcal{P}_r \Lambda^0(f_1)$      | $\mathcal{K}_{r+1}^1(f_2) \times \mathcal{K}_r^0(f_2)$ | $\mathcal{K}_{r+1}^2(f_3) \times \mathcal{K}_{r+1}^1(f_3)$ |
| 2                |   |                                     | $\mathcal{P}_r \Lambda^0(f_2)$                         | $\mathcal{K}_{r+1}^1(f_3) \times \mathcal{K}_r^0(f_3)$     |
| 3                |   |                                     |  | $\mathcal{P}_r \Lambda^0(f_3)$                             |

  

| $k \backslash d$ | 0                                 | 1                                   | 2  | 3  |
|------------------|-----------------------------------|-------------------------------------|--|--|
| 0                | $\mathbb{R} = \mathcal{P}_r(f_0)$ | $\{0\} \times \mathcal{P}_r^b(f_2)$ | $\{0\} \times \mathcal{R}_{r+1}^c(f_2)$                | $\{0\} \times \mathcal{R}_{r+1}^c(f_3)$                    |
| 1                |                                   | $\mathcal{P}_r(f_1)$                | $\mathcal{R}_{r+1}^c(f_2) \times \mathcal{P}_r^b(f_2)$ | $\mathcal{R}_{r+1}^c(f_3) \times \mathcal{G}_{r+1}^c(f_3)$ |
| 2                |                                   |                                     | $\mathcal{P}_r(f_2)$                                   | $\mathcal{G}_{r+1}^c(f_3) \times \mathcal{P}_r^b(f_3)$     |
| 3                |                                   |                                     |  | $\mathcal{P}_r(f_3)$                                       |

Table 2: Polynomial components attached to each mesh entity  $f_d$  of dimension  $d \in \{0, \dots, 3\}$  for the space  $\underline{V}_{r,h}^k$  for  $k \in \{0, \dots, 3\}$  (top) and counterparts through vector proxies (bottom).

consistency, cohomology, etc.) are carried out by induction on the dimension (and not on the difference between the dimension and the form degree, as in Theorems 11 and 15). This leads to somewhat simpler arguments, at the cost of larger discrete spaces. Also, the commutation property is essentially obtained by definition of the local discrete exterior derivative (with the exception of lowest-dimensional cells).

## 4.1 Definition

### 4.1.1 Discrete spaces

We define the following discrete counterpart of  $H\Lambda^k(\Omega)$ ,  $0 \leq k \leq n$ :

$$\underline{V}_{r,h}^k := \bigtimes_{f \in \Delta_k(\mathcal{M}_h)} \mathcal{P}_r \Lambda^0(f) \times \bigtimes_{f \in \Delta_{k+1}(\mathcal{M}_h)} \mathcal{K}_{r+1}^1(f) \times \mathcal{K}_r^0(f) \times \bigtimes_{d=k+2}^n \bigtimes_{f \in \Delta_d(\mathcal{M}_h)} \mathcal{K}_{r+1}^{d-k}(f) \times \mathcal{K}_{r+1}^{d-k-1}(f). \quad (4.1)$$

Notice that, on  $(k+1)$ -cells, the second component has polynomial degree reduced by one compared to  $d$ -cells with  $d \geq k+2$ , i.e., we have  $\mathcal{K}_r^0(f)$  instead of  $\mathcal{K}_{r+1}^0(f)$ . A generic element of  $\underline{V}_{r,h}^k$  will be denoted by

$$\underline{\omega}_h = ((\omega_f)_{f \in \Delta_k(\mathcal{M}_h)}, (\omega_f, D_{\omega,f})_{f \in \Delta_d(\mathcal{M}_h), d \in [k+1, n]}). \quad (4.2)$$

The notation  $D_{\omega,f}$  is reminiscent of the fact that these polynomial components are interpreted as Hodge stars of exterior derivatives. We refer to Table 2 for an overview of the polynomial unknowns in  $\underline{V}_{r,f}^k$  in dimensions 0 to 3, as well as their vector proxies.

### 4.1.2 Interpolators

For all integers  $0 \leq k \leq d \leq n$  and any  $f \in \Delta_d(\mathcal{M}_h)$ , the local interpolator is such that, for all  $\omega \in C^1\Lambda^k(\bar{f})$ ,

$$\begin{aligned} \underline{I}_{r,f}^k \omega := & \left( (\pi_{r,f'}^0(\star \text{tr}_{f'} \omega))_{f' \in \Delta_k(f)}, \right. \\ & (\pi_{r+1,f'}^{\mathcal{K},1}(\star \text{tr}_{f'} \omega), \pi_{r,f'}^{\mathcal{K},0}(\star \text{tr}_{f'} d\omega))_{f' \in \Delta_{k+1}(f)}, \\ & \left. (\pi_{r+1,f'}^{\mathcal{K},d'-k}(\star \text{tr}_{f'} \omega), \pi_{r+1,f'}^{\mathcal{K},d'-k-1}(\star \text{tr}_{f'} d\omega))_{f' \in \Delta_{d'}(f), d' \in [k+2, d]} \right). \end{aligned} \quad (4.3)$$

*Remark 21* (Domain of the interpolator). Owing to the presence of polynomial components that are interpreted as exterior derivatives (compare (3.2) with (4.3)), the interpolator in the VEM-inspired construction requires higher regularity of the interpolated functions compared to the DDR complex presented in Section 3, namely  $C^1\Lambda^k(\bar{f})$  instead of  $C^0\Lambda^k(\bar{f})$ .

### 4.1.3 Global discrete exterior derivative and VEM complex

For all  $f \in \Delta_{k+1}(\mathcal{M}_h)$ , we define the *discrete exterior derivative*  $\underline{d}_{r,f}^k : \underline{V}_{r,f}^k \rightarrow \mathcal{P}_r\Lambda^{k+1}(f)$  such that, for all  $\underline{\omega}_f \in \underline{V}_{r,f}^k$ ,

$$\int_f \underline{d}_{r,f}^k \underline{\omega}_f \wedge (\mu + \nu) = \int_{\partial f} \star^{-1} \omega_{\partial f} \wedge \text{tr}_{\partial f} \mu + \int_f \star^{-1} D_{\omega,f} \wedge \nu \quad \forall (\mu, \nu) \in \mathcal{P}_0\Lambda^0(f) \times \mathcal{K}_r^0(f), \quad (4.4)$$

where, as before,  $\omega_{\partial f}$  is defined by  $(\omega_{\partial f})|_{f'} = \omega_{f'} \in \mathcal{P}_r\Lambda^0(f')$  for all  $f' \in \Delta_k(f)$ . Notice that the above equation defines  $\underline{d}_{r,f}^k \underline{\omega}_f$  uniquely as, by (2.7a),  $\mu + \nu$  spans  $\mathcal{P}_r\Lambda^0(f)$  as  $(\mu, \nu)$  spans  $\mathcal{P}_0\Lambda^0(f) \times \mathcal{K}_r^0(f)$ . Moreover, taking  $\mu = 0$  and letting  $\nu$  span  $\mathcal{K}_r^0(f)$ , we infer, using (2.4) with  $(X, \omega, \mu) \leftarrow (\mathcal{K}_r^0(f), \underline{d}_{r,f}^k \underline{\omega}_f, \nu)$ ,

$$D_{\omega,f} = \pi_{r,f}^{\mathcal{K},0}(\star \underline{d}_{r,f}^k \underline{\omega}_f) \quad \forall f \in \Delta_{k+1}(\mathcal{M}_h). \quad (4.5)$$

Unlike the DDR complex, the construction of a *global discrete exterior derivative* for the VEM complex does not require to first reconstruct traces on lower-dimensional cells, as all the necessary information is encoded in the polynomial components  $(D_{\omega,f})_{f \in \Delta_d(\mathcal{M}_h), d \in [k+2, n]}$  supplemented by  $(\underline{d}_{r,f}^k \underline{\omega}_f)_{f \in \Delta_{k+1}(\mathcal{M}_h)}$ . More specifically, for all integers  $k \in [0, n-1]$ , we let  $\underline{d}_{r,h}^k : \underline{V}_{r,h}^k \rightarrow \underline{V}_{r,h}^{k+1}$  be such that, for all  $\underline{\omega}_h \in \underline{V}_{r,h}^k$ ,

$$\underline{d}_{r,h}^k \underline{\omega}_h := ((\star \underline{d}_{r,f}^k \underline{\omega}_f)_{f \in \Delta_{k+1}(\mathcal{M}_h)}, (D_{\omega,f}, 0)_{f \in \Delta_d(\mathcal{M}_h), d \in [k+2, n]}) \quad (4.6)$$

(compare with (4.2) and notice the different positioning, compared to  $\underline{\omega}_h$ , of the polynomial components  $D_{\omega,f}$ ). As for the DDR complex, we will denote by  $\underline{d}_{r,f}^k$  the restriction of  $\underline{d}_{r,h}^k$  to  $f \in \Delta_d(\mathcal{M}_h)$  with  $d \in [0, n]$  such that  $k \leq d-1$ .

The VEM sequence of spaces and operators then reads

$$\text{VEM}(r) := \{0\} \longrightarrow \underline{V}_{r,h}^0 \xrightarrow{\underline{d}_{r,h}^0} \underline{V}_{r,h}^1 \longrightarrow \dots \longrightarrow \underline{V}_{r,h}^{n-1} \xrightarrow{\underline{d}_{r,h}^{n-1}} \underline{V}_{r,h}^n \longrightarrow \{0\}. \quad (4.7)$$

### 4.1.4 Local discrete potentials and discrete exterior derivatives

Given a form degree  $k \in [0, n]$ , for all  $f \in \Delta_d(\mathcal{M}_h)$ ,  $k \leq d \leq n$ , we define the *local discrete potential*  $P_{r,f}^k : \underline{V}_r^k(f) \rightarrow \mathcal{P}_{r+1}^-\Lambda^k(f)$  by induction on  $d$  as follows: For all  $\underline{\omega}_f \in \underline{V}_{r,f}^k$ ,

- If  $d = k$ , we simply set

$$P_{r,f}^k \underline{\omega}_f := \star^{-1} \omega_f \in \mathcal{P}_r \Lambda^d(f) = \mathcal{P}_{r+1}^- \Lambda^d(f), \quad (4.8)$$

where the last equality follows from (2.13) if  $d = 0$  (after noticing that  $\mathcal{P}_r \Lambda^d(f) \cong \mathbb{R} \cong \mathcal{P}_{r+1}^- \Lambda^d(f)$ ) and from (2.15) if  $d \geq 1$ ;

- If  $k + 1 \leq d \leq n$ , using the isomorphism (2.17) with  $\ell = d - k \geq 1$  and  $r$  replaced by  $r + 1$ , we define  $P_{r,f}^k \underline{\omega}_f \in \mathcal{P}_{r+1}^- \Lambda^k(f)$  as the unique solution of the following equation:

$$(-1)^{k+1} \int_f P_{r,f}^k \underline{\omega}_f \wedge (d\mu + \nu) = \int_f \star^{-1} \tilde{D}_{\omega,f} \wedge \mu - \int_{\partial f} P_{r,\partial f}^k \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \mu + (-1)^{k+1} \int_f \star^{-1} \omega_f \wedge \nu$$

$$\forall (\mu, \nu) \in \mathcal{K}_{r+1}^{d-k-1}(f) \times \mathcal{K}_{r+1}^{d-k}(f), \quad (4.9)$$

where

$$\tilde{D}_{\omega,f} := \begin{cases} \star d_{r,f}^k \underline{\omega}_f & \text{if } d = k + 1, \\ D_{\omega,f} & \text{if } d \geq k + 2, \end{cases} \quad (4.10)$$

and we have introduced the piecewise polynomial boundary potential  $P_{r,\partial f}^k : \underline{V}_{r,\partial f}^k \rightarrow \Lambda^k(\partial f)$  such that  $(P_{r,\partial f}^k)|_{f'} := P_{r,f'}^k$  for all  $f' \in \Delta_{d-1}(f)$ .

Leveraging the above-defined discrete potentials, we can define the *discrete exterior derivative*  $d_{r,f}^k : \underline{V}_{r,f}^k \rightarrow \mathcal{P}_r \Lambda^{k+1}(f)$  for all  $f \in \Delta_d(\mathcal{M}_h)$ ,  $k + 2 \leq d \leq n - 1$  (this object was previously only defined for  $d = k + 1$ , see (4.4)), setting:

$$d_{r,f}^k \underline{\omega}_f := P_{r,f}^{k+1} d_{r,f}^k \underline{\omega}_f \quad \forall \underline{\omega}_f \in \underline{V}_{r,f}^k. \quad (4.11)$$

Notice that these discrete exterior derivatives are not relevant in the definition of the VEM complex, but may be useful in practical applications.

## 4.2 Main properties of the VEM complex

The main results for the VEM complex are stated below.

**Theorem 22** (Cohomology of the VEM complex). *The VEM sequence (4.7) is a complex and its cohomology is isomorphic to the cohomology of the continuous de Rham complex (1.2).*

*Proof.* See Section 4.6. □

**Theorem 23** (Polynomial consistency of the discrete potential and exterior derivative). *For all integers  $0 \leq k \leq d \leq n$  and all  $f \in \Delta_d(\mathcal{M}_h)$ , it holds*

$$P_{r,f}^k I_{r,f}^k \omega = \omega \quad \forall \omega \in \mathcal{P}_{r+1}^- \Lambda^k(f), \quad (4.12)$$

and, if  $d \geq k + 1$ ,

$$d_{r,f}^k I_{r,f}^k \omega = d\omega \quad \forall \omega \in \mathcal{P}_{r+1}^- \Lambda^k(f). \quad (4.13)$$

*Proof.* See Section 4.5. □

### 4.3 Complex property

**Lemma 24** (Complex property). *The sequence (4.7) defines a complex, i.e., for all integers  $k \in [1, n-1]$  and all  $\underline{\omega}_h \in \underline{V}_{r,h}^{k-1}$ ,*

$$\underline{d}_{r,h}^k(\underline{d}_{r,h}^{k-1}\underline{\omega}_h) = \underline{0}.$$

*Proof.* Applying the definition (4.6) of the global discrete exterior derivative first for  $k-1$  then for  $k$ , we obtain

$$\underline{d}_{r,h}^{k-1}\underline{\omega}_h = \left( (\star \underline{d}_{r,f}^{k-1}\underline{\omega}_f)_{f \in \Delta_k(\mathcal{M}_h)}, (D_{\omega,f}, 0)_{f \in \Delta_d(\mathcal{M}_h), d \in [k+1, n]} \right) \in \underline{V}_{r,h}^k, \quad (4.14)$$

which shows that, for all  $d \in [k+1, n]$  and all  $f \in \Delta_d(\mathcal{M}_h)$ , the exterior derivative components of  $\underline{d}_{r,h}^{k-1}\underline{\omega}_h$  are zero, and thus that

$$\underline{d}_{r,h}^k(\underline{d}_{r,h}^{k-1}\underline{\omega}_h) = \left( (\star \underline{d}_{r,f}^k(\underline{d}_{r,f}^{k-1}\underline{\omega}_f))_{f \in \Delta_{k+1}(\mathcal{M}_h)}, (0, 0)_{f \in \Delta_d(\mathcal{M}_h), d \in [k+2, n]} \right) \in \underline{V}_{r,h}^{k+1}.$$

The assertion is therefore proved if we show that  $\underline{d}_{r,f}^k(\underline{d}_{r,f}^{k-1}\underline{\omega}_f) = 0$  for all  $f \in \Delta_{k+1}(\mathcal{M}_h)$ . Applying the definition of the local discrete exterior derivative (see (4.4)) with  $\underline{\omega}_f$  replaced by  $\underline{d}_{r,f}^{k-1}\underline{\omega}_f$  obtained by restricting (4.14) to  $f$ , we get: For all  $(\mu, \nu) \in \mathcal{P}_0\Lambda^0(f) \times \mathcal{K}_r^0(f)$ ,

$$\int_f \underline{d}_{r,f}^k \underline{d}_{r,f}^{k-1} \underline{\omega}_f \wedge (\mu + \nu) = \int_{\partial f} \underline{d}_{r,\partial f}^{k-1} \underline{\omega}_{\partial f} \wedge \text{tr}_{\partial f} \mu = 0,$$

where the conclusion follows using the definition (4.4) of  $\underline{d}_{r,f}^{k-1}\underline{\omega}_f$  with  $(\mu, \nu) \leftarrow (\text{tr}_{f'} \mu, 0)$  for all  $f' \in \Delta_k(f)$  and noticing, as at the end of the proof of Lemma 14, that the sum over  $f'$  of the integrals over  $\partial f'$  is zero.  $\square$

### 4.4 Commutation

**Proposition 25** (Commutation property for the discrete exterior derivative in dimension  $d = k+1$ ). *For all  $f \in \Delta_{k+1}(\mathcal{M}_h)$ , it holds*

$$\underline{d}_{r,f}^k \underline{I}_{r,f}^k \omega = \star^{-1} \pi_{r,f}^0 (\star \underline{d}\omega) \quad \forall \omega \in C^1 \Lambda^k(\bar{f}), \quad (4.15)$$

expressing the commutativity of the following diagram:

$$\begin{array}{ccc} C^1 \Lambda^k(\bar{f}) & \xrightarrow{\underline{d}} & C^0 \Lambda^{k+1}(\bar{f}) \\ \downarrow \underline{I}_{r,f}^k & & \downarrow \star^{-1} \pi_{r,f}^0 \star \\ \underline{V}_{r,f}^k & \xrightarrow{\underline{d}_{r,f}^k} & \mathcal{P}_r \Lambda^{k+1}(f). \end{array}$$

*Proof.* Plugging the definition (4.3) of the interpolator into (4.4) we get, for all  $(\mu, \nu) \in \mathcal{P}_0\Lambda^0(f) \times \mathcal{K}_r^0(f)$ ,

$$\int_f \underline{d}_{r,f}^k \underline{I}_{r,f}^k \omega \wedge (\mu + \nu) = \int_{\partial f} \star^{-1} \pi_{r,\partial f}^0 (\star \text{tr}_{\partial f} \omega) \wedge \text{tr}_{\partial f} \mu + \int_f \star^{-1} \pi_{r,f}^0 (\star \underline{d}\omega) \wedge \nu,$$

where  $\pi_{r,\partial f}^0$  denotes the piecewise  $L^2$ -orthogonal projector obtained patching together the  $\pi_{r,f'}^0$ ,  $f' \in \Delta_k(f)$ . Using (2.4) with  $(\mathcal{X}, \omega, \mu) \leftarrow (\mathcal{K}_r^0(f'), \underline{d}\omega, \nu)$  for the second term and, for each  $f' \in \Delta_k(f)$ ,  $(\mathcal{X}, d, f) \leftarrow (\mathcal{P}_r \Lambda^0(f'), k, f')$  for the first term, the projectors can be removed. The Stokes formula (2.1) along with  $d\mu = 0$  (since  $\mu$  is constant) then yields

$$\int_f \underline{d}_{r,f}^k \underline{I}_{r,f}^k \omega \wedge (\mu + \nu) = \int_f \underline{d}\omega \wedge \mu + \int_f \underline{d}\omega \wedge \nu = \int_f \star^{-1} \pi_{r,f}^0 (\star \underline{d}\omega) \wedge (\mu + \nu),$$

where the conclusion follows from (2.4) with  $(\mathcal{X}, \omega, \mu) \leftarrow (\mathcal{P}_r \Lambda^0(f), \underline{d}\omega, \mu + \nu)$ . Since, by (2.7a),  $\mu + \nu$  spans  $\mathcal{P}_r \Lambda^0(f)$  as  $(\mu, \nu)$  spans  $\mathcal{P}_0\Lambda^0(f) \times \mathcal{K}_r^0(f)$ , this concludes the proof.  $\square$

**Proposition 26** (Commutation property for the local discrete exterior derivative). *For all integers  $d \in [1, n]$  and  $k \leq d - 1$ , and all  $f \in \Delta_d(\mathcal{M}_h)$ , it holds*

$$\underline{d}_{r,f}^k(I_{r,f}^k \omega) = I_{r,f}^{k+1}(d\omega) \quad \forall \omega \in C^2 \Lambda^k(\bar{f}), \quad (4.16)$$

expressing the commutativity of the following diagram:

$$\begin{array}{ccc} C^2 \Lambda^k(\bar{f}) & \xrightarrow{d} & C^1 \Lambda^{k+1}(\bar{f}) \\ \downarrow I_{r,f}^k & & \downarrow I_{r,f}^{k+1} \\ \underline{V}_r^k(f) & \xrightarrow{\underline{d}_{r,f}^k} & \underline{V}_r^{k+1}(f). \end{array}$$

*Proof.* Immediate consequence of (4.15) along with the definition (4.3) of the interpolator, and the property  $d \circ d = 0$ .  $\square$

#### 4.5 Polynomial consistency

*Proof of Theorem 23.* The proof proceeds by induction on the dimension  $d$ . When  $d = k$ , (4.12) is a direct consequence of the definitions (4.8) of the potential and (4.3) of the interpolator, which give  $P_{r,f}^k I_{r,f}^k \omega = \star^{-1} \pi_{r,f}^0(\star \omega) = \omega$ , where, to remove the projector, we have used the fact that  $\star \omega \in \mathcal{P}_r \Lambda^0(f)$ , since  $\omega \in \mathcal{P}_{r+1}^- \Lambda^d(f) = \mathcal{P}_r \Lambda^d(f)$  (see (2.15) with  $r + 1$  instead of  $r$ ).

We next prove (4.12) for  $d \geq k + 1$  assuming that it holds for  $d - 1$ . Writing the definition (4.9) of the potential for  $\underline{\omega}_f = I_{r,f}^k \omega$ , we get, for all  $(\mu, \nu) \in \mathcal{K}_{r+1}^{d-k-1}(f) \times \mathcal{K}_{r+1}^{d-k}(f)$ ,

$$\begin{aligned} & (-1)^{k+1} \int_f P_{r,f}^k(I_{r,f}^k \omega) \wedge (d\mu + \nu) \\ &= \int_f \star^{-1} \widetilde{D}_{\omega,f} \wedge \mu - \int_{\partial f} P_{r,\partial f}^k(I_{r,\partial f}^k \text{tr}_{\partial f} \omega) \wedge \text{tr}_{\partial f} \mu + (-1)^{k+1} \int_f \star^{-1} (\pi_{r+1,f}^{\mathcal{K},d-k} \star \omega) \wedge \nu \\ &= \int_f \star^{-1} \widetilde{D}_{\omega,f} \wedge \mu - \int_{\partial f} \text{tr}_{\partial f} \omega \wedge \text{tr}_{\partial f} \mu + (-1)^{k+1} \int_f \omega \wedge \nu, \end{aligned} \quad (4.17)$$

where we have used the induction hypothesis for the second term in the right-hand side after noticing that, by Lemma 4 with  $\ell = k$ ,  $\text{tr}_{f'} \omega \in \mathcal{P}_{r+1}^- \Lambda^k(f')$  for all  $f' \in \Delta_{d-1}(f)$ , together with (2.4) for the third one. Recalling the definition (4.10) of  $\widetilde{D}_{\omega,f}^k$ , we distinguish two cases for the first term in the right-hand side. If  $d = k + 1$ , (4.13) (immediate consequence of (4.15) after observing that  $d\mathcal{P}_{r+1}^- \Lambda^k(f) \subset \mathcal{P}_r \Lambda^{k+1}(f)$ ) gives  $\star^{-1} \widetilde{D}_{\omega,f} = \star^{-1} \star \underline{d}_{r,f}^k(I_{r,f}^k \omega) = d\omega$ . If, on the other hand,  $d \geq k + 2$ , recalling the definition (4.3) of the interpolator, we have  $\int_f \star^{-1} \widetilde{D}_{\omega,f} \wedge \mu = \int_f \star^{-1} (\pi_{r+1,f}^{\mathcal{K},d-k-1} \star d\omega) \wedge \mu \stackrel{(2.4)}{=} \int_f d\omega \wedge \mu$ . Plugging these relations into (4.17), using the Stokes formula (2.1), and simplifying, we get

$$\int_f P_{r,f}^k(I_{r,f}^k \omega) \wedge (d\mu + \nu) = \int_f \omega \wedge (d\mu + \nu),$$

which yields (4.12) for  $d \geq k + 1$  since, by (2.17) with  $\ell = d - k \geq 1$ ,  $d\mu + \nu$  spans  $\mathcal{P}_{r+1}^- \Lambda^{d-k}(f)$  as  $(\mu, \nu)$  spans  $\mathcal{K}_{r+1}^{d-k-1}(f) \times \mathcal{K}_{r+1}^{d-k}(f)$ .

We have already seen above that (4.13) holds for  $d = k + 1$ . To prove this relation for  $d \geq k + 2$ , it suffices to recall (4.11) and (4.16) to write

$$\underline{d}_{r,f}^k I_{r,f}^k \omega = P_{r,f}^{k+1}(\underline{d}_{r,f}^k I_{r,f}^k \omega) = P_{r,f}^{k+1}(I_{r,f}^{k+1} d\omega) = d\omega,$$

where the conclusion follows from (4.12) after observing that  $d\omega \in \mathcal{P}_r \Lambda^{k+1}(f) \subset \mathcal{P}_{r+1}^- \Lambda^{k+1}(f)$ .  $\square$

## 4.6 Cohomology

As in Section 3.5, given a form degree  $k \in [0, n]$ , we first consider the following subspace of  $V_{r,h}^k$ :

$$\underline{V}_{r,h,b}^k := \left\{ \underline{\omega}_h = ((\omega_f)_{f \in \Delta_k(\mathcal{M}_h)}, (\omega_f, D\omega_f)_{f \in \Delta_d(\mathcal{M}_h), d \in [k+1, n]}) : \int_f \star^{-1} \omega_f = 0 \quad \forall f \in \Delta_k(\mathcal{M}_h) \right\}.$$

**Lemma 27** (Exactness property for  $V_{r,h,b}^k$ ). *For all  $k \in [0, n]$ , if  $\underline{\eta}_h \in V_{r,h,b}^k$  satisfies  $\underline{d}_{r,h}^k \underline{\eta}_h = \underline{0}$ , then there exists  $\underline{\omega}_h \in V_{r,h,b}^{k-1}$  such that  $\underline{\eta}_h = \underline{d}_{r,h}^{k-1} \underline{\omega}_h$ , where, in accordance with the sequence (4.7), we have set  $\underline{d}_{r,h}^{-1} = \underline{d}_{r,h}^n := 0$ .*

*Proof.* Recalling the definition (4.6) of  $\underline{d}_{r,h}^k \underline{\eta}_h$ , we have

$$\underline{d}_{r,h}^k \underline{\eta}_h = ((\star \underline{d}_{r,f}^k \underline{\eta}_f)_{f \in \Delta_{k+1}(\mathcal{M}_h)}, (D\eta_{f,f}, 0)_{f \in \Delta_d(\mathcal{M}_h), d \in [k+2, n]}).$$

If  $k = 0$ , then  $\int_f \star^{-1} \eta_f = 0$  implies  $\eta_f = 0$  for all  $f \in \Delta_0(\mathcal{M}_h)$ ; moreover,  $\eta_f = 0$  for all  $f \in \Delta_d(\mathcal{M}_h)$ ,  $d \in [1, n]$ , by definition (4.1) of  $V_{r,h}^0$  (recall that  $\mathcal{K}_r^d(f) = \{0\}$  for all  $r$ , cf. (2.6)). The condition  $\underline{d}_{r,h}^k \underline{\eta}_h = \underline{0}$  together with (4.5) yields  $D\omega_{f,f} = 0$  for all  $f \in \Delta_d(\mathcal{M}_h)$ ,  $d \geq k+1$ , and thus

$$\underline{\eta}_h = ((0)_{f \in \Delta_0(\mathcal{M}_h)}, (0, 0)_{f \in \Delta_d(\mathcal{M}_h), d \in [k+1, n]}) = \underline{d}_{r,h}^{-1} 0.$$

If  $1 \leq k \leq n-1$ , on the other hand, from  $\underline{d}_{r,h}^k \underline{\eta}_h = \underline{0}$  and (4.5) we infer

$$\underline{\eta}_h = ((\eta_f)_{f \in \Delta_k(\mathcal{M}_h)}, (\eta_f, 0)_{f \in \Delta_d(\mathcal{M}_h), d \in [k+1, n]}), \quad (4.18a)$$

while, if  $k = n$ , we simply have

$$\underline{\eta}_h = (\eta_f)_{f \in \Delta_n(\mathcal{M}_h)}. \quad (4.18b)$$

Let now

$$\underline{\omega}_h = ((0)_{f \in \Delta_{k-1}(\mathcal{M}_h)}, (0, \pi_{r,f}^{\mathcal{K},0} \eta_f)_{f \in \Delta_k(\mathcal{M}_h)}, (0, \eta_f)_{f \in \Delta_d(\mathcal{M}_h), d \in [k+1, n]}) \in V_{r,h,b}^{k-1}.$$

To check that this  $\underline{\omega}_h$  is well defined, it suffices to notice that, if  $f \in \Delta_d(\mathcal{M}_h)$  with  $d \geq k+1 = (k-1)+2$ , then  $\eta_f \in \mathcal{K}_{r+1}^{d-k}(f) = \mathcal{K}_{r+1}^{d-(k-1)-1}(f)$  is a suitable choice for the corresponding component of  $\underline{\omega}_h$ . By definition (4.4) of  $\underline{d}_{r,f}^{k-1}$ , we have: For all  $f \in \Delta_k(\mathcal{M}_h)$  and all  $(\mu, \nu) \in \mathcal{P}_0 \Lambda^0(f) \times \mathcal{K}_r^0(f)$ , since  $\omega_{f'} = 0$  for all  $f' \in \Delta_k(f)$ ,

$$\int_f \underline{d}_{r,f}^{k-1} \underline{\omega}_f \wedge (\mu + \nu) = \int_f \star^{-1} \cancel{\pi_{r,f}^{\mathcal{K},0}} \eta_f \wedge \nu = \int_f \star^{-1} \eta_f \wedge (\mu + \nu),$$

where the cancellation of  $\pi_{r,f}^{\mathcal{K},0}$  is made possible by (2.4) with  $(\mathcal{X}, \omega, \mu) \leftarrow (\mathcal{K}_r^0(f), \star^{-1} \eta_f, \nu)$ , while the introduction of  $\mu$  in the last passage is justified observing that  $\underline{\eta}_h \in V_{r,h,b}^k$  implies  $\int_f \star^{-1} \eta_f = 0$  for all  $f \in \Delta_k(\mathcal{M}_h)$ . This relation gives  $\underline{d}_{r,f}^{k-1} \underline{\omega}_f = \star^{-1} \eta_f$  for all  $f \in \Delta_k(\mathcal{M}_h)$  which, combined with the definition (4.6) of the global discrete exterior derivative and the expression (4.18) of  $\underline{\eta}_h$ , readily yields  $\underline{\eta}_h = \underline{d}_{r,h}^{k-1} \underline{\omega}_h$  and concludes the proof.  $\square$

*Proof of Theorem 22.* Contrary to the DDR(0) complex, the VEM(0) complex is not isomorphic to the CW complex (the VEM spaces for  $r = 0$  do not have only constant polynomial components on the lowest-dimensional cells). As a consequence, designing extensions and reductions between the VEM( $r$ ) and VEM(0) complexes would not directly allow us, as it was the case for the DDR complex in the

proof of Theorem 10, to analyse the cohomology of the VEM complex. To circumvent this difficulty, we will instead design extensions  $\underline{E}_h^k : \underline{X}_{0,h}^k \rightarrow \underline{V}_{r,h}^k$  and reductions  $\underline{R}_h^k : \underline{V}_{r,h}^k \rightarrow \underline{X}_{0,h}^k$  between the VEM( $r$ ),  $r \geq 0$ , and the DDR(0) complexes, in order to show that their cohomologies are isomorphic. By Theorem 10, this will prove that the cohomology of VEM( $r$ ) is isomorphic to the continuous de Rham cohomology.

Throughout the rest of this proof,  $(P_{0,f}^k, d_{0,f}^k)$  and  $(P_{r,f}^k, d_{r,f}^k)$  denote, respectively, the couple (potential reconstruction, discrete exterior derivative) of the DDR(0) and VEM( $r$ ) complexes. We do not need to differentiate these notations, as the argument removes all ambiguity. For all form degrees  $k \in [0, d]$ , the reduction is obtained setting

$$\underline{R}_h^k \underline{\omega}_h := ((\pi_{0,f}^0 \omega_f)_{f \in \Delta_k(\mathcal{M}_h)}) \quad \forall \underline{\omega}_h \in \underline{V}_{r,h}^k, \quad (4.19)$$

while the extension is given by

$$\begin{aligned} \underline{E}_h^k \underline{\eta}_h &:= ((\eta_f)_{f \in \Delta_k(\mathcal{M}_h)}), \\ &(\pi_{r+1,f}^{\mathcal{K},1}(\star P_{0,f}^k \underline{\eta}_f), \pi_{r,f}^{\mathcal{K},0}(\star d_{0,f}^k \underline{\eta}_f))_{f \in \Delta_{k+1}(\mathcal{M}_h)}, \\ &(\pi_{r+1,f}^{\mathcal{K},d-k}(\star P_{0,f}^k \underline{\eta}_f), \pi_{r+1,f}^{\mathcal{K},d-k-1}(\star d_{0,f}^k \underline{\eta}_f))_{f \in \Delta_d(\mathcal{M}_h), d \geq k+2} \end{aligned} \quad \forall \underline{\eta}_h \in \underline{X}_{0,h}^k. \quad (4.20)$$

As in the proof of Theorem 10, we need to establish the properties (C1)–(C3) of [35, Assumption 1] to obtain the desired isomorphism in cohomology (also in this case, the relation (3.34) is an immediate consequence of (C1) and (C3)).

Proof of (C1). An inspection of the definitions (4.19) of the reduction and (4.20) of the extension shows that  $\underline{R}_h^k \underline{E}_h^k \underline{\eta}_h = \underline{\eta}_h$  for all  $\underline{\eta}_h \in \underline{X}_{0,h}^k$ , and thus (C1) holds a fortiori.

Proof of (C3). We need to prove that both the reduction and extension are cochain maps.

Let us start with the extension. We have to prove that, for any integer  $k \in [0, n-1]$  and all  $\underline{\eta}_h \in \underline{X}_{r,h}^k$ ,  $\underline{E}_h^{k+1}(\underline{d}_{0,h}^k \underline{\eta}_h) = \underline{d}_{r,h}^k(\underline{E}_h^k \underline{\eta}_h)$ . Owing to the definitions (4.20) of the extension, (3.7) of  $\underline{d}_{0,f}^k$ , and (4.6) of  $\underline{d}_{r,f}^k$ , and since  $\underline{d}_{0,f}^{k+1} \circ \underline{d}_{0,f}^k = 0$  (by (3.15) with  $r = 0$  and  $k+1$  instead of  $k$ ) this amounts to proving that

$$\star d_{0,f}^k \underline{\eta}_f = \star d_{r,f}^k(\underline{E}_f^k \underline{\eta}_f) \quad \forall f \in \Delta_{k+1}(\mathcal{M}_h), \quad (4.21)$$

$$\pi_{r+1,f}^{\mathcal{K},d-k-1}(\star P_{0,f}^{k+1} \underline{d}_{0,f}^k \underline{\eta}_f) = \pi_{r+1,f}^{\mathcal{K},d-k-1}(\star d_{0,f}^k \underline{\eta}_f) \quad \forall f \in \Delta_d(\mathcal{M}_h), d \geq k+2. \quad (4.22)$$

The relation (4.22) trivially follows from  $P_{0,f}^{k+1} \underline{d}_{0,f}^k = \underline{d}_{0,f}^k$ , see (3.14) with  $(k, r) \leftarrow (k+1, 0)$ . To prove (4.21), we take  $(\mu, \nu) \in \mathcal{P}_0 \Lambda^0(f) \times \mathcal{K}_r^0(f)$  and apply the definitions (4.4) of  $\underline{d}_{r,f}^k(\underline{E}_f^k \underline{\eta}_f)$  and (4.20) of  $\underline{E}_h^k \underline{\eta}_h$  to get

$$\begin{aligned} \int_f \underline{d}_{r,f}^k(\underline{E}_f^k \underline{\eta}_f) \wedge (\mu + \nu) &= \int_{\partial f} \star^{-1} \eta_{\partial f} \wedge \mu + \int_f \star^{-1} \pi_{r,f}^{\mathcal{K},0}(\star d_{0,f}^k \underline{\eta}_f) \wedge \nu \\ &= \int_f \underline{d}_{0,f}^k \underline{\eta}_f \wedge \mu + \int_f \underline{d}_{0,f}^k \underline{\eta}_f \wedge \nu = \int_f \underline{d}_{0,f}^k \underline{\eta}_f \wedge (\mu + \nu), \end{aligned}$$

where the second equality is obtained using the definition (3.4) of  $\underline{d}_{0,f}^k$  for the first term and (2.4) with  $(\mathcal{X}, \omega, \mu) \leftarrow (\mathcal{K}_r^0(f), \underline{d}_{0,f}^k \underline{\eta}_f, \nu)$  for the second one. By (2.7a), and since both  $\underline{d}_{r,f}^k(\underline{E}_f^k \underline{\eta}_f)$  and  $\underline{d}_{0,f}^k \underline{\eta}_f$  belong to  $\mathcal{P}_r \Lambda^{k+1}(f)$ , this relation gives  $\underline{d}_{r,f}^k(\underline{E}_f^k \underline{\eta}_f) = \underline{d}_{0,f}^k \underline{\eta}_f$ , thus proving (4.21).

Let us now turn to the reduction. We need to show that, for any integer  $k \in [0, n-1]$  and all  $\underline{\omega}_h \in \underline{V}_{r,h}^k$ ,  $\underline{R}_h^{k+1}(\underline{d}_{r,h}^k \underline{\omega}_h) = \underline{d}_{0,h}^k(\underline{R}_h^k \underline{\omega}_h)$ , i.e., accounting for the definitions (4.19) of the reduction, (3.7) of  $\underline{d}_{0,h}^k$  (additionally noticing that  $\pi_{0,f}^{-,0}$  coincides with  $\pi_{0,f}^0$  owing to (2.12a)), and (4.6) of  $\underline{d}_{r,h}^k$ ,

$$\pi_{0,f}^0(\star \underline{d}_{r,f}^k \underline{\omega}_f) = \star \underline{d}_{0,f}^k \underline{R}_f^k \underline{\omega}_f \quad \forall f \in \Delta_{k+1}(\mathcal{M}_h). \quad (4.23)$$

To check this relation, let  $f \in \Delta_{k+1}(\mathcal{M}_h)$  and write, for all  $\mu \in \mathcal{P}_0 \Lambda^0(f)$ ,

$$\begin{aligned} \int_f \star^{-1} \pi_{0,f}^0(\star \underline{d}_{r,f}^k \underline{\omega}_f) \wedge \mu &= \int_f \underline{d}_{r,f}^k \underline{\omega}_f \wedge \mu && \text{Eq. (2.4)} \\ &= \int_{\partial f} \star^{-1} \omega_{\partial f} \wedge \text{tr}_{\partial f} \mu && \text{Eq. (4.4)} \\ &= \int_{\partial f} \star^{-1} \pi_{0,\partial f}^0 \omega_{\partial f} \wedge \text{tr}_{\partial f} \mu \\ &= \int_f \underline{d}_{0,f}^k \underline{R}_f^k \underline{\omega}_f \wedge \mu, && \text{Eqs. (4.19), (3.4), (3.3)} \end{aligned}$$

where the third equality follows from (2.4) with  $(X, \omega, \mu) \leftarrow (\mathcal{P}_0 \Lambda^0(\partial f), \star^{-1} \omega_{\partial f}, \text{tr}_{\partial f} \mu)$ . This proves (4.23), and thus that the reductions form a cochain map.

*Proof of (C2).* For all  $\underline{\omega}_h \in \underline{V}_{r,h}^k$ , by the definitions (4.19) and (4.20) of the reduction and extension, it holds  $\underline{E}_h^k \underline{R}_h^k \underline{\omega}_h - \underline{\omega}_h \in \underline{V}_{r,h,b}^k$ . The proof then continues as in point 4. of the proof of Theorem 10 (see Section 3.5) with  $\underline{X}_{r,h,b}^k$  replaced by  $\underline{V}_{r,h,b}^k$  and Lemma 19 replaced by Lemma 27.  $\square$

## 5 Links with the literature

We provide here some elements of comparison between the DDR and VEM constructions of Sections 3 and 4, and two other families of discrete complexes.

### 5.1 Finite Element Exterior Calculus and Finite Element Systems

Finite Element Exterior Calculus (FEEC) is the (conforming) finite element approach for the unified analysis of discrete complexes [2, 4]. It is based on the selection of piecewise polynomial subspaces of  $H\Lambda(\Omega)$  that form a subcomplex of the continuous complex (1.2). Finite Element Systems (FES) is a framework for designing such subcomplexes that generalises FEEC to cover finite dimensional spaces spanned by differential forms that may not be piecewise polynomial on the selected mesh [29, 31]. In FEEC/FES complexes, only the spaces of differential forms in the continuous complex are replaced with discrete counterparts: the graded map that links these spaces is the usual exterior derivative  $d$ .

A FES space is a space of  $k$ -forms on all  $d$ -cells with  $d \geq k$ , with a compatibility condition on the traces:

$$A^k(\mathcal{M}_h) = \left\{ \underline{v}_h \in \bigtimes_{\substack{f \in \Delta_d(\mathcal{M}_h) \\ d \in [k, n]}} A^k(f) : \text{tr}_{f'} v_f = v_{f'} \text{ for all } (f, f') \in \Delta_d(\mathcal{M}_h) \times \Delta_{d'}(f) \text{ with } k \leq d' \leq d \right\}, \quad (5.1)$$

where each  $A^k(f)$  is a finite-dimensional space of  $k$ -forms and  $v_f$  denotes the component of  $\underline{v}_h$  on  $f$ . Because of this compatibility, any element in  $A^k(\mathcal{M}_h)$  can be identified with an element of  $H\Lambda^k(\Omega)$ . FES contains the usual FEEC complexes (in which case  $A^k(f)$  are certain polynomial subspaces – typically full polynomial spaces or trimmed polynomial spaces depending on the considered finite element), but has also been used to develop other discrete complexes, e.g. based on macro-elements



(in which case  $A^k(f)$  is a space of piecewise polynomial forms on a subdivision of  $f$ ) or with higher inter-element regularity ( $C^1$  spaces, for example).

The concept of (faithful) *mirror system* plays the role of degrees of freedom in the FES framework. Mirror systems are constructed on a case-by-case basis for each FES, and are auxiliary tools in the framework: they are not required to design the FES spaces, but they identify (by duality) a basis of such spaces. A mirror system for  $A^k(\mathcal{M}_h)$  is a family of subspaces of linear forms:

$$Z^k(\mathcal{M}_h) = \bigtimes_{\substack{f \in \Delta_d(\mathcal{M}_h) \\ d \in [k, n]}} Z^k(f) \quad \text{with } Z^k(f) \subset A^k(f)^* \text{ for all } f \in \Delta(\mathcal{M}_h), \quad (5.2)$$

where  $A^k(f)^*$  is the dual space of  $A^k(f)$  (actually, to link mirror systems and interpolators, each  $Z^k(f)$  is chosen as a subspace of  $\hat{X}^k(f)^*$  with  $\hat{X}^k(f) \supset A^k(f)$ , but we won't need this in the discussion here). As can be seen in (5.2), a mirror system is built hierarchically on the mesh, and each  $Z^k(f)$  identifies the modes of the FES forms that are “interior” to  $f$ ; to obtain all the modes (interior and boundary) associated with  $f$ , one must consider  $\bigtimes_{f' \in \Delta_{d'}(f), d' \in [k, d]} Z^k(f')$ .

A particular case of interest in the present context is when  $Z^k(f) \subset L^2\Lambda^k(f)^*$  (see Remark 5). Using the Riesz representation theorem and applying the Hodge star transformation,  $Z^k(f)$  can then be identified with a family of subspaces of  $L^2$ -integrable  $(d - k)$ -forms:

$$Z^k(\mathcal{M}_h) \cong \bigtimes_{\substack{f \in \Delta_d(\mathcal{M}_h) \\ d \in [k, n]}} \tilde{Z}^{d-k}(f) \quad \text{with } \tilde{Z}^{d-k}(f) \subset L^2\Lambda^{d-k}(f). \quad (5.3)$$

Here, and contrary to (5.1), no compatibility condition of the traces is imposed: the spaces  $\tilde{Z}^{d-k}(f)$  are completely disconnected from each other.

FEEC and FES provide computable spaces – that is, in which functions are entirely described in an algebraic manner – only on certain types of meshes. This is due to the requirement of compatible traces. In the DDR and VEM constructions of Sections 3 and 4, on the other hand, this requirement of computability for conforming subspaces is relaxed. Actually, no such space even needs to be identified: polytopal methods can be entirely built using spaces of polynomial functions on the mesh, without any compatibility condition on the traces. These spaces are explicit, and their basis is directly given by the polynomial components.

Comparing (3.1) and (5.3) for example, we see that the DDR space plays the role of a mirror system, and puts discrete polynomial components at the center of the construction. A similar approach is also true for the VEM-inspired spaces (4.1), with, contrary to DDR, some polynomial components representing exterior derivatives; see the definition (4.3) of the interpolator.

A closer link between DDR and FES can be drawn by noticing that the FES [30, Section 2.1] has the DDR spaces as mirror system (in the sense of (5.3)). This FES space, based on liftings of harmonic functions on each cell, therefore identifies a space of conforming functions whose degrees of freedom correspond to the DDR polynomial components; note that, in the context of vector proxies, another such identification was done in [11, Section 6.2].

The analogies, however, seem to stop here. While the FES theory would then use this conforming space and the continuous exterior derivative  $d$  to construct a discrete de Rham complex, DDR reconstructs *directly from the unknowns (mirror system)* a discrete exterior derivative  $\underline{d}_{r,h}$ , whose link with the continuous derivative of the corresponding FES function is not immediate. The appeal of this fully discrete approach is that, even when the FES space may not be computable (e.g., on polytopal meshes) and thus not directly usable in a scheme, the DDR space, its discrete exterior derivative, and its potential reconstruction are always computable, and are polynomially consistent, thus ensuring their practical applicability and optimal approximation properties.

## 5.2 Distributional Differential Forms

The theory of Distributional Differential Forms (DDF) has been introduced in [52] as a generalisation of the construction in [20] for the a posteriori error analysis of Nédélec edge elements. DDF are built on triangulations of the domain and, using their relation with the underlying simplicial complexes (as well as the concept of double complexes), their cohomology was analysed in [52] for rather general boundary conditions. Poincaré–Friedrichs inequalities were later established in [28].

As is the case for the spaces appearing in the DDR and VEM complexes, DDF spaces are collections of differential forms on cells of various dimensions, with form degree depending on the dimension of the cell: if the domain  $\Omega$  has dimension  $n$ , the DDF space of degree  $k$  is made of  $(k - n + d)$ -forms on  $d$ -cells. No compatibility of the traces is enforced on these forms, which can be completely discontinuous between two  $d$ -simplices. The discrete distributional exterior derivative on the DDF space is then composed of two contributions: the exterior derivative inside the simplices, and a trace term. For example, focusing on the highest dimension  $d = n$ , if the DDF space of  $k$ -forms is

$$\hat{\Lambda}_{-2}^k(\Delta_n(\mathcal{M}_h)) = \hat{\Lambda}_{-1}^k(\Delta_n(\mathcal{M}_h)) \oplus \hat{\Lambda}_{-1}^{k-1}(\Delta_{n-1}(\mathcal{M}_h)), \quad (5.4)$$

(with  $\hat{\Lambda}_{-1}^\ell$  subspace of piecewise  $C^\infty \Lambda^\ell$  forms, the index  $-1$  expressing the absence of continuity properties at the interfaces), for a family  $\omega_{n,h} = (\omega_f)_{f \in \Delta_n(\mathcal{M}_h)} \in \hat{\Lambda}_{-1}^k(\Delta_n(\mathcal{M}_h))$ , we define the distributional derivative  $\hat{d}_h^k : \hat{\Lambda}_{-1}^k(\Delta_n(\mathcal{M}_h)) \rightarrow \hat{\Lambda}_{-2}^{k+1}(\Delta_n(\mathcal{M}_h))$  by

$$\hat{d}_h^k \omega_{n,h} = \left( (d^k \omega_f)_{f \in \Delta_n(\mathcal{M}_h)}, \left( - \sum_{f \in \Sigma_n(f')} \varepsilon_{ff'} \operatorname{tr}_{f'} \omega_f \right)_{f' \in \Delta_{n-1}(\mathcal{M}_h)} \right), \quad (5.5)$$

where  $\Sigma_n(f')$  is the set of  $n$ -simplices  $f$  that share  $f'$  (that is,  $f' \in \Delta_{n-1}(f)$ ), and  $\varepsilon_{ff'}$  is the relative orientation of the simplex  $f'$  with respect to the simplex  $f$ . Note that, in (5.5), we have adopted a presentation of the distributional derivative that distributes its two contributions (D and T in [52]) on the corresponding components  $(\hat{\Lambda}_{-1}^{k+1-i}(\Delta_{n-i}(\mathcal{M}_h)))_{i=0,1}$  of  $\hat{\Lambda}_{-2}^{k+1}(\Delta_n(\mathcal{M}_h))$  (see (5.4) with  $k+1$  instead of  $k$ ), instead of writing  $\hat{d}_h^k$  as a sum of elements in the global space  $\hat{\Lambda}_{-2}^{k+1}(\Delta_n(\mathcal{M}_h))$ ; this is to better compare with the definition (3.7). This definition of distributional derivative is a global one, obtained by testing the piecewise smooth form  $\omega_{n,h}$  against globally smooth forms, which classically results in a term inside each  $f \in \Delta_n(\mathcal{M}_h)$  corresponding to the standard exterior derivative (first component in (5.5)), and a jump across the  $(n-1)$ -sub-simplices based on the difference of traces on the two adjacent  $n$ -simplices (second component in (5.5)).

A crucial remark is that, in (5.5), the component  $(d^k \omega_f)_{f \in \Delta_n(\mathcal{M}_h)}$  of  $\hat{d}_h^k \omega_{n,h}$  on  $n$ -cells only depends on the values  $\omega_{n,h}$  of the discrete distributional differential form on  $n$ -cells, not on the values of these forms on lower-dimensional cells (e.g.,  $\hat{\Lambda}_{-1}^{k-1}(\Delta_{n-1}(\mathcal{M}_h))$  in (5.4)). This is in contrast with the discrete exterior derivatives in DDR and VEM complexes, whose definition on higher-dimensional cells depends on polynomial components on their sub-cells; see (3.4) and (4.6). Another difference between DDR and DDF can be seen when recasting the discrete exterior derivative: integrating by parts (3.4) yields the following characterisation:

$$\int_f d_{r,f}^k \omega_f \wedge \mu = - \int_f d(\star^{-1} \omega_f) \wedge \mu + \int_{\partial f} (P_{r,\partial f}^k \omega_{\partial f} - \operatorname{tr}_{\partial f}(\star^{-1} \omega_f)) \wedge \operatorname{tr}_{\partial f} \mu$$

$$\forall \mu \in \mathcal{P}_r \Lambda^{d-k-1}(f).$$

This relation reveals that  $d_{r,f}^k \omega_f$  is, as in DDF, composed of an exterior derivative term in the  $d$ -cell and a boundary term involving jumps. However, contrary to DDF, the jumps here are between the trace of the  $d$ -cell unknown and the potential  $P_{r,\partial f}^k \omega_{\partial f}$  reconstructed on  $(d-1)$ -cells (which depends on

the unknowns on all  $d'$ -subcells of  $f$ ,  $k \leq d' \leq d$ ), not between traces of two  $d$ -cells unknowns (as in (5.5) with  $d = n$ ). In this respect, the “jump” term in DDR relates more to the kind of face differences encountered in polytopal methods (e.g., the HHO method [36]) while the jump term in DDF is more akin to those arising in discontinuous Galerkin (DG) methods [41].

This comparison can be extended to the potential reconstructions themselves. Equation (3.19) shows that  $P_{r,f}^k \omega_f$  is obtained applying a higher-order correction to the cell component  $\star^{-1} \omega_f$ , designed from the discrete exterior derivative on  $f$  and the potentials on  $\partial f$ . This *enhancement* process ensures the high-order consistency of the method starting from lower-order polynomial unknowns. In the context of elliptic equations, it is commonly used in methods with unknowns in the elements and on the faces of the mesh, but it is not directly available in DG methods. In DDF, as in DG, the cell unknown itself must be used (e.g., in a scheme to discretise the source term), and the consistency is therefore limited by the degree of this unknown.

## A Differential forms and vector proxies

In this section, we briefly recall basic concepts on alternating (resp. differential) forms, and their representation in terms of vectors (resp. vector fields); these representations are often referred to as “vector proxies”. We refer the reader to [2, Chapter 6] for a presentation in the framework of Finite Element Exterior Calculus, and to [15], [25, Chapter 1], [50, Chapter 1] for an introduction in more general scientific and engineering contexts.

### A.1 Exterior algebra in $\mathbb{R}^n$

#### A.1.1 Alternating forms

Let  $\{e_i\}_{i \in [1,n]}$  be the canonical basis of  $\mathbb{R}^n$ , equipped with the standard inner product. A basis for the space of linear forms over  $\mathbb{R}^n$ , i.e., the dual space  $(\mathbb{R}^n)'$  of  $\mathbb{R}^n$ , is given by  $\{dx^i\}_{i \in [1,n]}$ , with  $dx^i(e_j) := \delta_j^i$  (Kronecker symbol), for all  $(i, j) \in [1, n]^2$ . The starting point of exterior calculus is to consider *alternating* multilinear forms, vanishing whenever they are applied to a set of linearly dependent vectors in  $\mathbb{R}^n$ . For any integer  $k \geq 1$ , the set of alternating  $k$ -linear forms on  $\mathbb{R}^n$  is denoted by  $\text{Alt}^k(\mathbb{R}^n)$ ; by convention, we set  $\text{Alt}^0(\mathbb{R}^n) := \mathbb{R}$ . We also note that  $\text{Alt}^1(\mathbb{R}^n) = (\mathbb{R}^n)'$  and that  $\text{Alt}^k(\mathbb{R}^n) = \{0\}$  if  $k > n$  (since families of  $k > n$  vectors are always linearly dependent). It can be checked that  $\dim \text{Alt}^k(\mathbb{R}^n) = \binom{n}{k}$ . In particular,  $\text{Alt}^n(\mathbb{R}^n)$  is the 1-dimensional space spanned by the determinant in the canonical basis  $\text{vol}$  (called the volume form).

#### A.1.2 Exterior product

Given two alternating multilinear forms  $\omega \in \text{Alt}^i(\mathbb{R}^n)$  and  $\mu \in \text{Alt}^j(\mathbb{R}^n)$ , their *exterior product*  $\omega \wedge \mu \in \text{Alt}^{i+j}(\mathbb{R}^n)$  is defined, for any vectors  $v_1, \dots, v_{i+j} \in \mathbb{R}^n$ , by

$$(\omega \wedge \mu)(v_1, \dots, v_{i+j}) := \sum_{\sigma \in \Sigma_{i,j}} \text{sign}(\sigma) \omega(v_{\sigma_1}, \dots, v_{\sigma_i}) \mu(v_{\sigma_{i+1}}, \dots, v_{\sigma_{i+j}}),$$

where  $\Sigma_{i,j}$  is the set of all permutations  $\sigma$  of the  $(i+j)$ -tuple  $(1, \dots, i+j)$  such that  $\sigma_1 < \dots < \sigma_i$  and  $\sigma_{i+1} < \dots < \sigma_{i+j}$ . The exterior product satisfies the anticommutativity law

$$\omega \wedge \mu = (-1)^{ij} \mu \wedge \omega, \quad (\text{A.1})$$

so that, in particular, we have  $dx^i \wedge dx^i = 0$  and  $dx^i \wedge dx^j = -dx^j \wedge dx^i$ . With these definitions, for  $k \in [1, n]$  a basis of the space  $\text{Alt}^k(\mathbb{R}^n)$  is  $\{dx^{\sigma_1} \wedge \dots \wedge dx^{\sigma_k}\}_{\sigma}$  where  $\sigma$  spans all strictly increasing functions  $[1, k] \rightarrow [1, n]$ . Hence, any  $\omega \in \text{Alt}^k(\mathbb{R}^n)$  can be written

$$\omega = \sum_{1 \leq \sigma_1 < \dots < \sigma_k \leq n} a_{\sigma} dx^{\sigma_1} \wedge \dots \wedge dx^{\sigma_k}, \quad a_{\sigma} \in \mathbb{R}. \quad (\text{A.2})$$

### A.1.3 Hodge star operator

The scalar product in  $\mathbb{R}^n$  induces a scalar product, denoted by  $\langle \cdot, \cdot \rangle$ , on  $\text{Alt}^{n-k}(\mathbb{R}^n)$  – namely, the scalar product for which the aforementioned basis  $\{dx^{\sigma_1} \wedge \cdots \wedge dx^{\sigma_{n-k}}\}_\sigma$  of  $\text{Alt}^{n-k}(\mathbb{R}^n)$  is orthonormal. The *Hodge star operator* is the unique linear mapping  $\star : \text{Alt}^k(\mathbb{R}^n) \rightarrow \text{Alt}^{n-k}(\mathbb{R}^n)$  such that, for all  $\omega \in \text{Alt}^k(\mathbb{R}^n)$ ,  $\langle \star\omega, \mu \rangle \text{vol} = \omega \wedge \mu$  for all  $\mu \in \text{Alt}^{n-k}(\mathbb{R}^n)$ . It can be checked that

$$\star(dx^{\sigma_1} \wedge \cdots \wedge dx^{\sigma_k}) = \text{sign}(\sigma, \tau)(dx^{\tau_1} \wedge \cdots \wedge dx^{\tau_{n-k}}),$$

where  $(\sigma, \tau) = (\sigma_1, \dots, \sigma_k, \tau_1, \dots, \tau_{n-k})$  is a permutation of  $(1, \dots, n)$  such that  $\sigma_1 < \cdots < \sigma_k$  and  $\tau_1 < \cdots < \tau_{n-k}$ . From the above identity, one can infer that

$$\star(\star\omega) = (-1)^{k(n-k)}\omega \quad \forall \omega \in \text{Alt}^k(\mathbb{R}^n) \quad (\text{A.3})$$

and, hence, that  $\langle \star\omega, \star\mu \rangle = \langle \omega, \mu \rangle$ , i.e.,  $\star$  is an isometry. Formula (A.3) justifies the definition (2.2) of  $\star^{-1}$ . The anticommutativity (A.1) of  $\wedge$ , the definition of  $\star$ , and the symmetry of  $\langle \cdot, \cdot \rangle$  then give

$$\star^{-1}\omega \wedge \mu = \mu \wedge \star\omega = \omega \wedge \star\mu \quad \forall \omega, \mu \in \text{Alt}^k(\mathbb{R}^n). \quad (\text{A.4})$$

**Example 28** (Hodge star operator in two and three dimensions). *If  $\omega \in \text{Alt}^2(\mathbb{R}^3)$ , i.e.,  $\omega = a_{12} dx^1 \wedge dx^2 + a_{13} dx^1 \wedge dx^3 + a_{23} dx^2 \wedge dx^3$  (see (A.2)), one obtains  $\star\omega \in \text{Alt}^1(\mathbb{R}^3)$  with*

$$\star\omega = a_{12} dx^3 - a_{13} dx^2 + a_{23} dx^1.$$

*If  $\omega \in \text{Alt}^1(\mathbb{R}^2)$ , i.e.,  $\omega = a_1 dx^1 + a_2 dx^2$ , then  $\star\omega \in \text{Alt}^1(\mathbb{R}^2)$  with*

$$\star\omega = a_1 dx^2 - a_2 dx^1.$$

### A.1.4 Vector proxies for alternating forms

As already mentioned in Section A.1.1,  $\text{Alt}^0(\mathbb{R}^n) = \mathbb{R}$  and  $\text{Alt}^n(\mathbb{R}^n) \cong \mathbb{R}$ . Using the Riesz representation theorem to identify  $(\mathbb{R}^n)'$  and  $\mathbb{R}^n$ , we can identify two further spaces of alternating forms. Specifically,  $\text{Alt}^1(\mathbb{R}^n) = (\mathbb{R}^n)' \cong \mathbb{R}^n$  and, writing  $\star\text{Alt}^{n-1}(\mathbb{R}^n) = \text{Alt}^1(\mathbb{R}^n) \cong \mathbb{R}^n$ , since  $\star$  is bijective, we obtain the identification  $\text{Alt}^{n-1}(\mathbb{R}^n) \cong \mathbb{R}^n$ .

Applied with  $n = 3$ , and recalling the formula for Hodge star transformations of 2-forms in Remark 28, these identifications lead to considering a vector  $\mathbf{v} = (a, b, c) \in \mathbb{R}^3$  as a *proxy* for both the alternating linear and bilinear forms

$$\text{Alt}^1(\mathbb{R}^3) \ni \omega = a dx^1 + b dx^2 + c dx^3 \text{ and } \text{Alt}^2(\mathbb{R}^3) \ni \mu = a dx^2 \wedge dx^3 - b dx^1 \wedge dx^3 + c dx^1 \wedge dx^2.$$

On the other hand, when  $n = 2$ , the discussion above gives two possible ways to identify  $\text{Alt}^1(\mathbb{R}^2) = \text{Alt}^{2-1}(\mathbb{R}^2)$  with  $\mathbb{R}^2$ . This leads to associating  $a dx^1 + b dx^2 = \omega \in \text{Alt}^1(\mathbb{R}^2)$  either to the vector  $\mathbf{v} = (a, b) \in \mathbb{R}^2$ , or to its rotation by a right angle  $\varrho_{-\pi/2}\mathbf{v} = (b, -a) \in \mathbb{R}^2$ .

Based on the the above identifications, when  $n = 3$ , one can interpret the exterior product of two alternating multilinear forms  $\omega \wedge \mu$  in terms of vector proxies  $(\mathbf{w}, \mathbf{v})$  as follows:

- the vector product  $\mathbb{R}^3 \times \mathbb{R}^3 \ni (\mathbf{w}, \mathbf{v}) \mapsto \mathbf{w} \times \mathbf{v} \in \mathbb{R}^3$  when  $(\omega, \mu) \in \text{Alt}^1(\mathbb{R}^3) \times \text{Alt}^1(\mathbb{R}^3)$ ;
- the dot product  $\mathbb{R}^3 \times \mathbb{R}^3 \ni (\mathbf{w}, \mathbf{v}) \mapsto \mathbf{w} \cdot \mathbf{v} \in \mathbb{R}$  when  $(\omega, \mu) \in \text{Alt}^1(\mathbb{R}^3) \times \text{Alt}^2(\mathbb{R}^3)$ .

On the other hand, if  $n = 2$  and  $\omega, \mu \in \text{Alt}^1(\mathbb{R}^2)$ , we can write  $\omega \wedge \mu = (a dx^1 + b dx^2) \wedge (f dx^1 + g dx^2) = (ag - bf) dx^1 \wedge dx^2$ . Considering the correspondences  $\omega \leftrightarrow \mathbf{w} = (a, b)$  and  $\mu \leftrightarrow \mathbf{v} = (f, g)$ , we obtain

$$\omega \wedge \mu = (\mathbf{w} \cdot \varrho_{-\pi/2}\mathbf{v}) dx^1 \wedge dx^2. \quad (\text{A.5})$$

### A.1.5 Contraction and trace

For a given vector  $\mathbf{v} \in \mathbb{R}^n$ , the *contraction*  $\omega \lrcorner \mathbf{v} \in \text{Alt}^{k-1}(\mathbb{R}^n)$  of  $\omega \in \text{Alt}^k(\mathbb{R}^n)$  with  $\mathbf{v}$  is defined, for any  $\mathbf{v}_1, \dots, \mathbf{v}_{k-1} \in \mathbb{R}^n$ , by

$$(\omega \lrcorner \mathbf{v})(\mathbf{v}_1, \dots, \mathbf{v}_{k-1}) := \omega(\mathbf{v}, \mathbf{v}_1, \dots, \mathbf{v}_{k-1}). \quad (\text{A.6})$$

In terms of vector proxies, in the case where  $n = 3$ , this contraction with  $\mathbf{v}$  corresponds to

- the scalar product  $\mathbb{R}^3 \ni \mathbf{w} \mapsto \mathbf{v} \cdot \mathbf{w} \in \mathbb{R}$  when  $\mathbf{w} \leftrightarrow \omega \in \text{Alt}^1(\mathbb{R}^3)$ ;
- the vector product  $\mathbb{R}^3 \ni \mathbf{w} \mapsto \mathbf{v} \times \mathbf{w} \in \mathbb{R}^3$  when  $\mathbf{w} \leftrightarrow \omega \in \text{Alt}^2(\mathbb{R}^3)$ ;
- the multiplication of a real number  $\mathbb{R} \ni w \mapsto w\mathbf{v} \in \mathbb{R}^3$  when  $w \leftrightarrow \omega \in \text{Alt}^3(\mathbb{R}^3)$ .

Let now  $V \subset W$  be finite dimensional subspaces of  $\mathbb{R}^n$ , and  $\iota_V : V \hookrightarrow W$  be the inclusion of  $V$  in  $W$ . The *trace*  $\text{tr}_V : \text{Alt}^k(W) \rightarrow \text{Alt}^k(V)$  is the pullback under  $\iota_V$ , that is: for any  $\mathbf{v}_1, \dots, \mathbf{v}_k \in V$ ,

$$\text{tr}_V \omega(\mathbf{v}_1, \dots, \mathbf{v}_k) := \omega(\iota_V \mathbf{v}_1, \dots, \iota_V \mathbf{v}_k). \quad (\text{A.7})$$

The trace respects the exterior product, i.e.,  $\text{tr}_V(\omega \wedge \mu) = \text{tr}_V \omega \wedge \text{tr}_V \mu$ .

It is easy to see that, through the vector proxy of  $\text{Alt}^1$  spaces,  $\text{tr}_V : \text{Alt}^1(W) \rightarrow \text{Alt}^1(V)$  is the orthogonal projection  $\pi_V : W \rightarrow V$  of a vector  $\mathbf{w} \in W$  onto  $V$ .

Let us fix an integer  $m \in [1, n]$  and suppose that  $\dim(W) = m$  and  $\dim(V) = m - 1$ , and that both spaces are oriented; let  $\mathbf{n}_V$  be the unit normal to  $V$  such that, given a positively oriented basis  $(\mathbf{e}_1, \dots, \mathbf{e}_{m-1})$  of  $V$ , the family  $(\mathbf{n}_V, \mathbf{e}_1, \dots, \mathbf{e}_{m-1})$  forms a positively oriented basis of  $W$ . Then, an identification of the trace  $\text{tr}_V : \text{Alt}^{m-1}(W) \rightarrow \text{Alt}^{m-1}(V)$  through vector proxies is the scalar product with the vector  $\mathbf{n}_V$ , that is,  $W \ni \mathbf{w} \mapsto \mathbf{w} \cdot \mathbf{n}_V \in \mathbb{R}$ .

## A.2 Exterior calculus in $\mathbb{R}^n$

### A.2.1 Differential forms

Let  $M$  be an  $n$ -dimensional flat manifold. When the coefficients in (A.2) are functions  $a_\sigma : M \rightarrow \mathbb{R}$ , the map  $\omega : M \rightarrow \text{Alt}^k(\mathbb{R}^n)$  is referred to as a *differential form*, or simply a  $k$ -form. Consistently with the notation adopted in Section 2.1, the space of  $k$ -forms over  $M$  without any specific smoothness requirement on the coefficients  $a_\sigma$  is denoted by  $\Lambda^k(M)$ . If  $\omega \in \Lambda^k(M)$ , the value of  $\omega$  at  $\mathbf{x} \in M$  is denoted by  $\omega_{\mathbf{x}} \in \text{Alt}^k(\mathbb{R}^n)$ .

If the coefficients  $a_\sigma$  in (A.2) are polynomial functions,  $\omega$  is said to be a *polynomial differential form*. Specifically, for an integer  $r \geq 0$ , the space of polynomial  $k$ -forms of degree  $\leq r$  is defined as

$$\mathcal{P}_r \Lambda^k(M) := \left\{ \sum_{1 \leq \sigma_1 < \dots < \sigma_k \leq n} p_\sigma dx^{\sigma_1} \wedge \dots \wedge dx^{\sigma_k} : p_\sigma \in \mathcal{P}_r(M) \right\},$$

where  $\mathcal{P}_r(M)$  is the space of scalar polynomials of degree  $\leq r$  over  $M$ . All the arguments concerning vector proxies presented in Section A.1 for alternating  $k$ -linear forms can be immediately extended to the case of  $k$ -forms. Hence, when  $n \in \{2, 3\}$ , their corresponding vector proxies are scalar fields over  $M$  when  $k \in \{0, n\}$ , and vector fields over  $M$  when  $k \in \{1, n - 1\}$ .

### A.2.2 Exterior derivative and de Rham complexes

Provided that the coefficients  $a_\sigma$  in (A.2) are smooth enough, the *exterior derivative* of a  $k$ -form  $\omega \in \Lambda^k(M)$  is the linear unbounded operator  $d : \Lambda^k(M) \rightarrow \Lambda^{k+1}(M)$  such that, in terms of standard coordinates on  $\mathbb{R}^n$ ,

$$d\omega = \sum_{1 \leq \sigma_1 < \dots < \sigma_k \leq n} \sum_{i=1}^n \frac{\partial a_\sigma}{\partial x_i} dx^i \wedge dx^{\sigma_1} \wedge \dots \wedge dx^{\sigma_k}.$$

The interpretation of the exterior derivative in terms of vector calculus operators, through vector proxies of alternating forms and when  $M$  is a domain  $\Omega$  of  $\mathbb{R}^3$ , is given in (A.8). We have used in this diagram the spaces defined in the introduction of the paper.

$$\begin{array}{ccccccc}
\text{Differential forms:} & H\Lambda^0(\Omega) & \xrightarrow{d} & H\Lambda^1(\Omega) & \xrightarrow{d} & H\Lambda^2(\Omega) & \xrightarrow{d} & H\Lambda^3(\Omega) \\
& \updownarrow & & \updownarrow & & \updownarrow & & \updownarrow \\
\text{Vector proxies:} & H^1(\Omega) & \xrightarrow{\text{grad}} & \mathbf{H}(\text{curl}; \Omega) & \xrightarrow{\text{curl}} & \mathbf{H}(\text{div}; \Omega) & \xrightarrow{\text{div}} & L^2(\Omega).
\end{array} \tag{A.8}$$

In the case  $n = 2$ , as we have two possible vector proxies for  $\text{Alt}^1(\mathbb{R}^2)$ . These interpretations are illustrated in (A.9) when  $\omega = a dx^1 + b dx^2 \in \text{Alt}^1(\mathbb{R}^2)$  is identified with  $\mathbf{v} = (a, b)$ , and in (A.10) when  $\omega \in \text{Alt}^{2-1}(\mathbb{R}^2)$  is identified with  $\varrho_{-\pi/2}\mathbf{v}$  (with  $\text{rot} = \text{div} \varrho_{-\pi/2}$  and  $\mathbf{rot} = \varrho_{-\pi/2} \mathbf{grad}$ , respectively, denoting the scalar and vector curls, and  $\mathbf{H}(\text{rot}; \Omega)$  the space of square-integrable vector-valued functions whose rot is also square-integrable).

$$\begin{array}{ccccccc}
\text{Differential forms:} & H\Lambda^0(\Omega) & \xrightarrow{d} & H\Lambda^1(\Omega) & \xrightarrow{d} & H\Lambda^2(\Omega) & & \\
& \updownarrow & & \updownarrow & & \updownarrow & & \\
\text{Vector proxies:} & H^1(\Omega) & \xrightarrow{\text{grad}} & \mathbf{H}(\text{rot}; \Omega) & \xrightarrow{\text{rot}} & L^2(\Omega). & & 
\end{array} \tag{A.9}$$

$$\begin{array}{ccccccc}
\text{Differential forms:} & H\Lambda^0(\Omega) & \xrightarrow{d} & H\Lambda^1(\Omega) & \xrightarrow{d} & H\Lambda^2(\Omega) & & \\
& \updownarrow & & \updownarrow & & \updownarrow & & \\
\text{Vector proxies:} & H^1(\Omega) & \xrightarrow{\text{rot}} & \mathbf{H}(\text{div}; \Omega) & \xrightarrow{\text{div}} & L^2(\Omega). & & 
\end{array} \tag{A.10}$$

Notice, finally, that the exterior derivative satisfies the *complex property*  $d \circ d = 0$ . This property translates, through vector proxies, into the well-known identities  $\mathbf{curl} \mathbf{grad} = \mathbf{0}$  and  $\text{div} \mathbf{curl} = 0$  for  $n = 3$ , and  $\text{rot} \mathbf{grad} = 0$ ,  $\text{div} \mathbf{rot} = 0$  when  $n = 2$ .

### A.2.3 Koszul differential

Given  $\mathbf{x}_M \in \mathbb{R}^n$ , the *Koszul differential*  $\kappa_M : \Lambda^k(M) \rightarrow \Lambda^{k-1}(M)$  is defined pointwise over  $M$  as follows: For all  $\mathbf{x} \in M$ , recalling the definition (A.6) of the contraction  $\lrcorner$ ,

$$(\kappa_M \omega)_{\mathbf{x}} := \omega_{\mathbf{x} \lrcorner} (\mathbf{x} - \mathbf{x}_M).$$

Its interpretation in terms of vector fields proxy is then analogous to that of a contraction of an alternating multilinear form with a vector, except that the contraction is made pointwise with the vector field  $\mathbb{R}^n \ni \mathbf{x} \mapsto \mathbf{x} - \mathbf{x}_M \in \mathbb{R}^n$ . The terminology ‘‘differential’’ is legitimate, as  $\kappa_M$  satisfies the complex property  $\kappa_M \circ \kappa_M = 0$  (since any alternating form applied to the same vector twice vanishes).

### A.2.4 Trace

If  $P \subset Q$  are (relatively) open sets in affine subspaces  $V \subset W$  of  $\mathbb{R}^n$ , the trace operator  $\text{tr}_P : C^0 \Lambda^k(\overline{Q}) \rightarrow C^0 \Lambda^k(\overline{P})$  on differential forms is defined pointwise, using the trace operator (A.7) on alternating forms: For all  $\omega \in C^0 \Lambda^k(\overline{Q})$ ,

$$(\text{tr}_P \omega)_{\mathbf{x}} := \text{tr}_V \omega_{\mathbf{x}} \quad \forall \mathbf{x} \in P.$$

Note that, in the case  $P = Q$ , the trace is simply the identity operator (and can be defined without any continuity assumption):  $\text{tr}_P \omega = \omega$  for all  $\omega \in \Lambda^k(P)$ .

Applying the same arguments as in Section A.1 pointwise over  $P$ , the trace operator in terms of vector fields proxy gives

- the restriction of functions, when  $k = 0$ ;
- the orthogonal projection onto  $V$  (that is,  $\text{tr}_P \omega \leftrightarrow \pi_V \mathbf{w}$  if  $\omega \leftrightarrow \mathbf{w}$ ), when  $k = 1$ ;
- the normal component on  $P$  along the direction  $\mathbf{n}$  (that is,  $\text{tr}_P \omega \leftrightarrow \mathbf{w} \cdot \mathbf{n}$  if  $\omega \leftrightarrow \mathbf{w}$ ), with  $\mathbf{n}$  unit normal vector field preserving the orientations of  $V$  and  $W$ , when  $k = \dim(P) = \dim(Q) - 1$ .

**Example 29** (Interpretation of the Stokes formula for  $\ell = 1$  and  $n = 3$ ). We rewrite here, for the reader's convenience, the integration by parts formula (2.1) for  $\ell = 1$  and  $n = 3$ :

$$\int_M d\omega \wedge \mu = \int_M \omega \wedge d\mu + \int_{\partial M} \text{tr}_{\partial M} \omega \wedge \text{tr}_{\partial M} \mu \quad \forall (\omega, \mu) \in \Lambda^1(M) \times \Lambda^1(M). \quad (\text{A.11})$$

Given the previous interpretations of the exterior derivative and product in terms of vector proxies, if  $\omega \leftrightarrow \mathbf{w}$  and  $\mu \leftrightarrow \mathbf{v}$ , then  $d\omega \wedge \mu \leftrightarrow \mathbf{curl} \mathbf{w} \cdot \mathbf{v}$  and  $\omega \wedge d\mu \leftrightarrow \mathbf{w} \cdot \mathbf{curl} \mathbf{v}$ . This leads to considering the integration by parts formula for the curl:

$$\int_M \mathbf{curl} \mathbf{w} \cdot \mathbf{v} = \int_M \mathbf{w} \cdot \mathbf{curl} \mathbf{v} + \int_{\partial M} (\mathbf{n} \times (\mathbf{w} \times \mathbf{n})) \cdot (\mathbf{v} \times \mathbf{n}), \quad (\text{A.12})$$

where  $\mathbf{n}$  is the outer unit normal vector field over  $\partial M$ . For any fixed  $\mathbf{x} \in \partial M$ , we have  $\mathbf{n}(\mathbf{x}) \times (\mathbf{w}(\mathbf{x}) \times \mathbf{n}(\mathbf{x})) = \pi_{T_{\mathbf{x}}\partial M} \mathbf{w}(\mathbf{x})$  (here,  $T_{\mathbf{x}}\partial M$  is the tangent space of  $\partial M$  at  $\mathbf{x}$ ), whereas  $\mathbf{v}(\mathbf{x}) \times \mathbf{n}(\mathbf{x}) = \varrho_{-\pi/2}(\pi_{T_{\mathbf{x}}\partial M} \mathbf{v}(\mathbf{x}))$ , where the rotation is considered with respect to the orientation of the tangent plane given by  $\mathbf{n}(\mathbf{x})$ . The boundary terms of (A.11) and (A.12) therefore coincide, through the vector proxy for the exterior product of 1-forms in dimension 2 (see (A.5)).

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## References

- [1] P. F. Antonietti, S. Giani, and P. Houston. “hp-version composite discontinuous Galerkin methods for elliptic problems on complicated domains”. In: *SIAM J. Sci. Comput.* 35.3 (2013), A1417–A1439. DOI: 10.1137/120877246.
- [2] D. Arnold. *Finite Element Exterior Calculus*. SIAM, 2018. DOI: 10.1137/1.9781611975543.
- [3] D. Arnold and K. Hu. “Complexes from Complexes”. In: *Foundations of Computational Mathematics* 6 (2021), pp. 1739–1774. DOI: 10.1007/s10208-021-09498-9.
- [4] D. N. Arnold, R. S. Falk, and R. Winther. “Finite element exterior calculus, homological techniques, and applications”. In: *Acta Numer.* 15 (2006), pp. 1–155. DOI: 10.1017/S0962492906210018.
- [5] D. N. Arnold, R. S. Falk, and R. Winther. “Finite element exterior calculus: from Hodge theory to numerical stability”. In: *Bull. Amer. Math. Soc. (N.S.)* 47.2 (2010), pp. 281–354. DOI: 10.1090/S0273-0979-10-01278-4.
- [6] F. Bassi, L. Botti, A. Colombo, D. A. Di Pietro, and P. Tesini. “On the flexibility of agglomeration based physical space discontinuous Galerkin discretizations”. In: *J. Comput. Phys.* 231.1 (2012), pp. 45–65. DOI: 10.1016/j.jcp.2011.08.018.

- [7] L. Beirão da Veiga, F. Brezzi, A. Cangiani, G. Manzini, L. D. Marini, and A. Russo. “Basic principles of virtual element methods”. In: *Math. Models Methods Appl. Sci. (M3AS)* 199.23 (2013), pp. 199–214. doi: 10.1142/S0218202512500492.
- [8] L. Beirão da Veiga, F. Brezzi, F. Dassi, L. D. Marini, and A. Russo. “A family of three-dimensional virtual elements with applications to magnetostatics”. In: *SIAM J. Numer. Anal.* 56.5 (2018), pp. 2940–2962. doi: 10.1137/18M1169886.
- [9] L. Beirão da Veiga, F. Brezzi, F. Dassi, L. D. Marini, and A. Russo. “Serendipity virtual elements for general elliptic equations in three dimensions”. In: *Chin. Ann. Math. Ser. B* 39.2 (2018), pp. 315–334. doi: 10.1007/s11401-018-1066-4.
- [10] L. Beirão da Veiga, F. Brezzi, L. D. Marini, and A. Russo. “ $H(\text{div})$  and  $H(\text{curl})$ -conforming VEM”. In: *Numer. Math.* 133 (2016), pp. 303–332. doi: 10.1007/s00211-015-0746-1.
- [11] L. Beirão da Veiga, F. Dassi, D. A. Di Pietro, and J. Droniou. “Arbitrary-order pressure-robust DDR and VEM methods for the Stokes problem on polyhedral meshes”. In: *Comput. Meth. Appl. Mech. Engrg.* 397.115061 (2022). doi: 10.1016/j.cma.2022.115061.
- [12] L. Beirão da Veiga, F. Dassi, and G. Vacca. “The Stokes complex for Virtual Elements in three dimensions”. In: *Math. Models Methods Appl. Sci.* 30 (2020), pp. 477–512. doi: 10.1142/S0218202520500128.
- [13] L. Beirão da Veiga, K. Lipnikov, and G. Manzini. *The mimetic finite difference method for elliptic problems*. Vol. 11. MS&A. Modeling, Simulation and Applications. Springer, Cham, 2014, pp. xvi+392. doi: 10.1007/978-3-319-02663-3.
- [14] L. Beirão da Veiga, D. Mora, and G. Vacca. “The Stokes complex for Virtual Elements with application to Navier–Stokes flows”. In: *J. Sci. Comput.* 81 (2019), pp. 990–1018. doi: 10.1007/s10915-019-01049-3.
- [15] J. Blair Perot and C. J. Zusi. “Differential forms for scientists and engineers”. In: *Journal of Computational Physics* 257 (2014). Physics-compatible numerical methods, pp. 1373–1393. doi: 10.1016/j.jcp.2013.08.007.
- [16] J. Bonelle and A. Ern. “Analysis of compatible discrete operator schemes for elliptic problems on polyhedral meshes”. In: *ESAIM: Math. Model. Numer. Anal.* 48 (2014), pp. 553–581. doi: 10.1051/m2an/2013104.
- [17] A. Bossavit. “Whitney forms: a class of Finite Elements for three-dimensional computation in electromagnetism”. In: *IEEE Proceedings A* 135 (1988), pp. 493–500.
- [18] L. Botti. “Influence of reference-to-physical frame mappings on approximation properties of discontinuous piecewise polynomial spaces”. In: *J. Sci. Comput.* 52.3 (2012), pp. 675–703. doi: 10.1007/s10915-011-9566-3.
- [19] M. Botti, D. A. Di Pietro, and M. Salah. “A serendipity fully discrete div-div complex on polygonal meshes”. In: *Comptes Rendus Mécanique* (2023). Published online. doi: 10.5802/crmeca.150.
- [20] D. Braess and J. Schöberl. “Equilibrated residual error estimator for edge elements”. In: *Math. Comp.* 77.262 (2008), pp. 651–672. doi: 10.1090/S0025-5718-07-02080-7.
- [21] F. Brezzi, A. Buffa, and K. Lipnikov. “Mimetic finite differences for elliptic problems”. In: *M2AN Math. Model. Numer. Anal.* 43.2 (2009), pp. 277–295. doi: 10.1051/m2an:2008046.
- [22] F. Brezzi, R. S. Falk, and L. D. Marini. “Basic principles of mixed virtual element methods”. In: *ESAIM Math. Model. Numer. Anal.* 48.4 (2014), pp. 1227–1240. doi: 10.1051/m2an/2013138.
- [23] J. Brüning and M. Lesch. “Hilbert complexes”. In: *J. Funct. Anal.* 108.1 (1992), pp. 88–132. doi: 10.1016/0022-1236(92)90147-B.



- [24] A. Cangiani, Z. Dong, E. H. Georgoulis, and P. Houston. *hp-version discontinuous Galerkin methods on polygonal and polyhedral meshes*. SpringerBriefs in Mathematics. Springer, Cham, 2017, pp. viii+131.
- [25] M. P. do Carmo. *Differential forms and applications*. Universitext. Translated from the 1971 Portuguese original. Springer-Verlag, Berlin, 1994, pp. x+118. DOI: 10.1007/978-3-642-57951-6.
- [26] S. H. Christiansen. “A construction of spaces of compatible differential forms on cellular complexes”. In: *Math. Models Methods Appl. Sci.* 18.5 (2008), pp. 739–757. DOI: 10.1142/S021820250800284X.
- [27] S. H. Christiansen, J. Gopalakrishnan, J. Guzmán, and K. Hu. *A discrete elasticity complex on three-dimensional Afeld splits*. arXiv preprint. Sept. 2020. URL: <https://arxiv.org/abs/2009.07744>.
- [28] S. H. Christiansen and M. W. Licht. “Poincaré–Friedrichs inequalities of complexes of discrete distributional differential forms”. In: *BIT Numerical Mathematics* 60.2 (2020), pp. 345–371.
- [29] S. H. Christiansen, H. Z. Munthe-Kaas, and B. Owren. “Topics in structure-preserving discretization”. In: *Acta Numerica* 20 (2011), pp. 1–119.
- [30] S. H. Christiansen and A. Gillette. “Constructions of some minimal finite element systems”. In: *ESAIM Math. Model. Numer. Anal.* 50.3 (2016), pp. 833–850. DOI: 10.1051/m2an/2015089.
- [31] S. H. Christiansen and K. Hu. “Generalized finite element systems for smooth differential forms and Stokes’ problem”. In: *Numer. Math.* 140.2 (2018), pp. 327–371. DOI: 10.1007/s00211-018-0970-6.
- [32] D. A. Di Pietro. *An arbitrary-order discrete rot-rot complex on polygonal meshes with application to a quad-rot problem*. Oct. 2022. arXiv: 2210.15581 [math.NA].
- [33] D. A. Di Pietro and J. Droniou. “An arbitrary-order discrete de Rham complex on polyhedral meshes: Exactness, Poincaré inequalities, and consistency”. In: *Found. Comput. Math.* 23 (2023), pp. 85–164. DOI: 10.1007/s10208-021-09542-8.
- [34] D. A. Di Pietro and J. Droniou. “An arbitrary-order method for magnetostatics on polyhedral meshes based on a discrete de Rham sequence”. In: *J. Comput. Phys.* 429.109991 (2021). DOI: 10.1016/j.jcp.2020.109991.
- [35] D. A. Di Pietro and J. Droniou. “Homological- and analytical-preserving serendipity framework for polytopal complexes, with application to the DDR method”. In: *ESAIM: Math. Model Numer. Anal.* 57.1 (2023), pp. 191–225. DOI: 10.1051/m2an/2022067.
- [36] D. A. Di Pietro and J. Droniou. *The Hybrid High-Order method for polytopal meshes. Design, analysis, and applications*. Vol. 19. Modeling, Simulation and Application. Springer International Publishing, 2020. DOI: 10.1007/978-3-030-37203-3.
- [37] D. A. Di Pietro, J. Droniou, and S. Pitassi. *Cohomology of the discrete de Rham complex on domains of general topology*. Aug. 2022. arXiv: 2209.00957 [math.NA].
- [38] D. A. Di Pietro, J. Droniou, and F. Rapetti. “Fully discrete polynomial de Rham sequences of arbitrary degree on polygons and polyhedra”. In: *Math. Models Methods Appl. Sci.* 30.9 (2020), pp. 1809–1855. DOI: 10.1142/S0218202520500372.
- [39] D. A. Di Pietro and A. Ern. “Arbitrary-order mixed methods for heterogeneous anisotropic diffusion on general meshes”. In: *IMA J. Numer. Anal.* 37.1 (2017), pp. 40–63. DOI: 10.1093/imanum/drw003.

- [40] D. A. Di Pietro and A. Ern. “Discrete functional analysis tools for discontinuous Galerkin methods with application to the incompressible Navier-Stokes equations”. In: *Math. Comp.* 79.271 (2010), pp. 1303–1330. doi: 10.1090/S0025-5718-10-02333-1.
- [41] D. A. Di Pietro and A. Ern. *Mathematical aspects of discontinuous Galerkin methods*. Vol. 69. Mathématiques & Applications (Berlin) [Mathematics & Applications]. Springer, Heidelberg, 2012, pp. xviii+384. doi: 10.1007/978-3-642-22980-0.
- [42] D. A. Di Pietro, A. Ern, and S. Lemaire. “An arbitrary-order and compact-stencil discretization of diffusion on general meshes based on local reconstruction operators”. In: *Comput. Meth. Appl. Math.* 14.4 (2014), pp. 461–472. doi: 10.1515/cmam-2014-0018.
- [43] D. A. Di Pietro and J. Droniou. “A discrete de Rham method for the Reissner-Mindlin plate bending problem on polygonal meshes”. In: *Comput. Math. Appl.* 125 (2022), pp. 136–149. doi: 10.1016/j.camwa.2022.08.041.
- [44] D. A. Di Pietro and J. Droniou. “A fully discrete plates complex on polygonal meshes with application to the Kirchhoff–Love problem”. In: *Math. Comp.* 92 (339 2023), pp. 51–77. doi: 10.1090/mcom/3765.
- [45] J. Droniou, R. Eymard, T. Gallouët, C. Guichard, and R. Herbin. *The gradient discretisation method*. Vol. 82. Mathematics & Applications. Springer, 2018, 511p. doi: 10.1007/978-3-319-79042-8.
- [46] J. Droniou, R. Eymard, T. Gallouët, and R. Herbin. “A unified approach to mimetic finite difference, hybrid finite volume and mixed finite volume methods”. In: *Math. Models Methods Appl. Sci. (M3AS)* 20.2 (2010), pp. 1–31. doi: 10.1142/S0218202510004222.
- [47] J. Droniou, T. A. Oliynyk, and J. J. Qian. “A polyhedral discrete de Rham numerical scheme for the Yang–Mills equations”. In: *J. Comput. Phys.* (2023), 26p. doi: 10.1016/j.jcp.2023.111955.
- [48] R. Eymard, T. Gallouët, and R. Herbin. “Discretization of heterogeneous and anisotropic diffusion problems on general nonconforming meshes. SUSHI: a scheme using stabilization and hybrid interfaces”. In: *IMA J. Numer. Anal.* 30.4 (2010), pp. 1009–1043. doi: 10.1093/imanum/drn084.
- [49] M.-L. Hanot. *An arbitrary-order fully discrete Stokes complex on general polyhedral meshes*. May 2022. arXiv: 2112.03125 [math.NA].
- [50] M.-L. Hanot. “Complexes discrets pour les fluides incompressibles”. In French. Université de Montpellier, Dec. 2022.
- [51] S. Lang. *Fundamentals of differential geometry*. Vol. 191. Graduate Texts in Mathematics. Springer-Verlag, New York, 1999, pp. xviii+535. doi: 10.1007/978-1-4612-0541-8.
- [52] M. W. Licht. “Complexes of discrete distributional differential forms and their homology theory”. In: *Found. Comput. Math.* 17.4 (2017), pp. 1085–1122. doi: 10.1007/s10208-016-9315-y.
- [53] J.-C. Nédélec. “Mixed finite elements in  $\mathbf{R}^3$ ”. In: *Numer. Math.* 35.3 (1980), pp. 315–341. doi: 10.1007/BF01396415.
- [54] P. A. Raviart and J. M. Thomas. “A mixed finite element method for 2nd order elliptic problems”. In: *Mathematical Aspects of the Finite Element Method*. Ed. by I. Galligani and E. Magenes. New York: Springer, 1977.
- [55] J. Zhao and B. Zhang. “The curl-curl conforming virtual element method for the quad-curl problem”. In: *Math. Models Methods Appl. Sci.* 31.8 (2021), pp. 1659–1690. doi: 10.1142/S0218202521500354.