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Repulsive force for Micro-&Nano- non-contact manipulation

Amélie Cot,^{1,2} Patrick Rougeot,¹ Sophie Lakard,² Michaël Gauthier^{1*} and Jérôme Dejeu^{1,3*}

¹ FEMTO-ST Institute, Univ. Franche-Comté, CNRS, ENSMM, 24 rue Savary, F-25000 Besançon, France

² UTINAM Institute, Univ. Franche-Comté, 16 Route de Gray, 25030 Besançon

³ Univ. Grenoble Alpes, CNRS, DCM UMR-5250, 38000 Grenoble, France

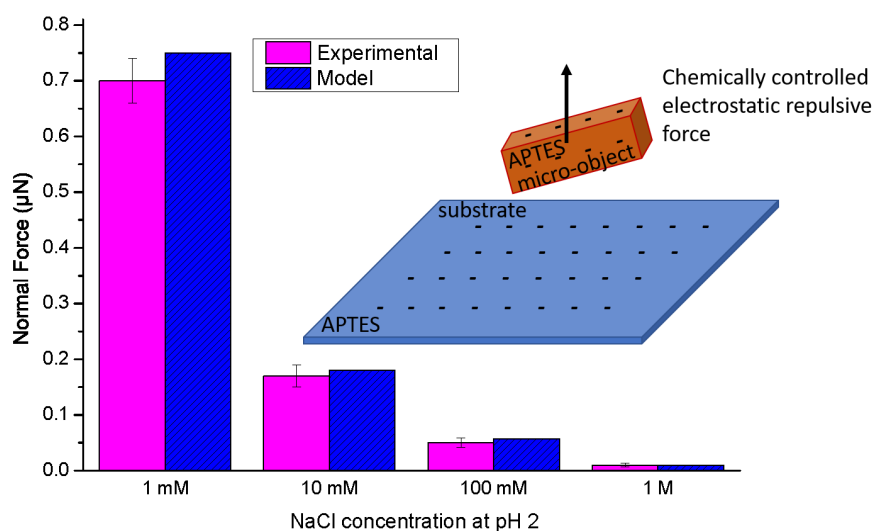
* Correspondence: JD: jerome.dejeu@univ-grenoble-alpes.fr and MG: michael.gauthier@femto-st.fr

Featured This functionalization is able to provide solutions in order to improve robotic nano or micromanipulation.

Abstract: Non-contact positioning of micro-objects using electric fields has been widely explored, based on several physical principles such as electrophoresis, dielectrophoresis (DEP) or optical dielectrophoresis (ODEP), in which the actuation force is induced by an electric charge or an electric dipole placed in an electric field. In this paper, we introduce a new way to control charges in non-contact positioning of micro-objects using chemical functionalization (3-aminopropyl)triethoxysilane - APTES) able to localize charges on a substrate and/or on a micro-object. We demonstrate that this functionalization in a liquid with a low ionic strength is able to concentrate a significant amount of electric charges on surfaces generating an electric field over a long distance (about 10 microns), also call a large exclusion zone (EZ). A model is proposed and validated with electrostatic force measurements between substrate and microparticles (diameter up to 40 μ m). We demonstrate that the magnitude of the force and the force range decrease rapidly when the ionic strength of the medium increases. Based on the proposed model, we show that this new way to localize charges on micro-objects may be used for non-contact positioning.

Keywords: APTES grafting, Exclusion Zone (EZ), Force modelling, Non-contact-manipulation, PANI electropolymerisation, Repulsive force, Surface functionalization.

Graphical Abstract:



1. Introduction

Non-contact micro-nanomanipulation or microsorting consists in manipulating or sorting microparticles using forces generated via a long range physical field (magnetic field, electrostatic field, acoustic field). The behavior and the design of the devices are significantly modified by the well-known scale effects [1]. When the scale reduces, the physical effects' magnitude is drastically modified: when the lengths are divided by 10, the volume effects (e.g. weight mass) are divided by 10^3 , and the surface forces (e.g., van der Waals force) are divided only by 10^2 . Therefore, the effect of gravity thus decreases more rapidly than surface forces' effect during miniaturization [2]. So, on the microscale, the surface forces are predominant compared to the weight and the objects tend to stick to the surfaces (e.g., adhesion [3]). Therefore, the manipulation (move and position) of a micro-object is usually performed without touching it but using "non-contact manipulation" proposed to avoid adhesion disturbance [4]. These manipulation methods are usually propelled by electrostatic [5], laser induced thermal gradient [6], optical trapping [7,8], magnetic [8,9] or thermocapillary [10] forces. This article presents a new way to perform micromanipulation using electrostatic forces. In this field, the electric field is usually generated by (micro)electrodes placed in a liquid in order to induce dielectrophoresis (DEP) [11], optically induced dielectrophoresis (ODEP) [12], electro-rotation [13] or electrophoresis [14]. Indeed, applying a voltage on electrodes induces charges generating an electric field up to the electrodes. The microparticles located in the electric field experience a force of several tens of micrometers away from the electrodes depending on the applied electric field generated by the charges of the electrodes, the particle's size and electrical properties. This general principle is usually used to sort particles [12, 15].

The objective of this paper is to propose a new way to control non-contact manipulations based on electric fields. Indeed, we propose to generate electrophoretic force using electric charges based on chemical principles in spite of generating electric charges generated with an external voltage. Concretely, a local chemical functionalization (typically amine groups) enables to locally concentrate electric charges generating an electric field up to the substrate in the liquid [16]. In order to be usable, the electric field generated by chemical functions has to be able to induce significant electrostatic forces on microparticles on a long range (typ. greater than 1 micron).

We are going to show that the interaction distance is highly impacted by the scale effect and that nanoparticles and microparticles have significant different behaviors.

In nanoscale, the repulsive forces between colloids and a flat surface were already measured by different researchers in liquid media. The measures of repulsion forces enable to identify the hydration force [17] or electrostatic repulsions in inorganic solvents with different types of spheres glued on Atomic Force Microscope (AFM) tips: alumina or silica [18], silicon nitride [19], and gold [20] tips. More recently, some repulsive electrostatic forces have been measured in water [21]. In all these cases, the size of the sphere diameter is less than $5\ \mu\text{m}$ and the interaction distance is near 20 nm.

Our paper focuses on larger microspheres whose diameter is higher than several micrometers. We are going to show that it is possible to generate significant electrostatic forces on these microparticles over a long range (around ten of micron).

The section 2 introduces the materials and methods used in the technical parts. The section 3 reports some experimental force measurement illustrating the long-range interaction forces. The section 4 focuses on a model of the force whose comparison with experiments is described in section 5. The simulation of potential uses of these long range force in non-contact manipulation of microparticle is introduced in section 6 and the general results are discussed in section 7 before the conclusion.

2. Materials and Methods

2.1. Salt

Lithium perchlorate (LiClO₄), sodium tetrafluoroborate (NaBF₄), sodium nitrate (NaNO₃) and p-toluenesulfonic acid (PTS) come from Sigma-Aldrich

2.2. Media preparation:

AFM measurements were performed in water at pH 2 prepared before each day before the series of measurements. At the end of the experiments, the pH was controlled to validate the measurements performed. pH of the solution was measured with a pH-meter (Sartorius, PT-10) and an electrode (Sartorius, PY-P22) and adjusted at pH 2 by addition hydrochloric acid 1M just before the measurement to protonate all the amine functions. For all experiments, ionic strength was about 10⁻³ M except for experiments concerning ionic strength's influence (controlled by NaCl addition).

2.3. Surface Functionalization

Before being functionalized, silicon wafers (purchased from Tracit) were cleaned by immersion in a Piranha solution (2 parts H₂SO₄ and 1 part H₂O₂) for 25 min at 70 °C. Then, wafers were rinsed in Milli-Q water and ethanol before functionalization.

2.4. Silanization

Solutions were freshly prepared by direct dissolution of silanes (3-aminopropyl) triethoxysilane, APTES) in ethanol. The final silane concentration was 1%. The surfaces were functionalized by immersion in solutions for one night at room temperature. In the silane solution, the molecules were grafted on the substrate (through covalent bonds). The excess of ungrafted silanes was removed by ultrasonication for 2 min in ethanol. The mechanism of SAM formation during the silanization process has already been described by Wasserman et al. [22]. The mechanism of self-assembled monolayer formation during the silanization process and takes place in four steps [22, 23]. The first step is physisorption, in which the silane molecules become physisorbed at the hydrated silicon surface. In the second step, the silane head-groups arrive close to the substrate hydrolyse, in the presence of the adsorbed water layer on the surface, into highly polar trihydroxysilane Si(OH)₃ for triethoxysilane Si(OEt)₃ (APTES). These polar groups, (Si(OH)₃), form covalent bonds with the hydroxyl groups on the SiO₂ surface (third step); subsequently, condensation reaction (release of water molecules) goes on between silanol functions of neighbour molecules. Self-assembly is driven by lipophilic interactions between the linear alkane. During the initial period, only a few molecules will adsorb (by steps 1–3) on the surface and the monolayer will definitely be in a disordered (or liquid) state. However, at longer times, surface coverage eventually reaches the point where a well-ordered and compact (or crystalline) monolayer is obtained (step 4), by the condensation reaction between the APTES molecules.

The grafting was controlled by contact angle measurements. The contact angle before functionalisation was inferior to 10° and increased from 60° to 80° after APTES grafting. These values were concordant with previous experiments [24].

2.4. Electrochemical deposition of thin films

The pyrrole and 3-aminopropyltriethoxysilan were from Acros organics (99% pure) and were distilled under reduced pressure before use. Lithium perchlorate was from Sigma Aldrich and used as electrolytic salt. Electrolytes were composed of 0.1M pyrrole in an aqueous solution of 0.1 M LiClO₄ or other salts (PTS, NaBF₄, NaNO₃). Electrochemical experiments were performed with a PGZ 100 potentiostat (Tacussel-Radiometer Analytical SA-France) controlled by the VoltaMaster 4 software. A standard three-electrode system was relied to the potentiostat and composed by a Saturated Calomel Electrode (SCE) as the reference electrode, a platinum sheet as the counter-electrode, and a working electrode which was a silicon substrate previously covered by sprayed chrome and gold to enhance its conductivity. The working electrode was cleaned one hour using an UV-

Ozone treatment (Bioforce UV/Ozone Procleaner) before electrochemical deposition process. All electrochemical experiments were carried out at room temperature (293K). Firstly, cyclic voltammetry technique was used in order to define the potential of polymerization of pyrrole. Then, potentiostatic chronoamperometries were performed to obtain thin films of polypyrrole. Previous work of Patois et al. [25] shown that polypyrrole could be deposited from +0.7 V/SCE with an oxidation peak appearing at +1.0 V/SCE in the cyclic voltammogram, so we decide to work à +0.7V/SCE. The films obtained has previously characterized by SEM and AFM [26].

2.5. Force measurements

The force measurement experiments were performed with an AFM tip on which a functionalized borosilicate sphere was glued according to the procedure previously described [27]. In order to characterize surface functionalization, a Smena S7 Atomic Force Microscope (AFM) from NTMDT has been used. The silicon rectangular AFM cantilever (from Novascan Technologies) has a stiffness of 0.3 N/m. The cantilever is fixed while the substrate moves vertically. Most of the AFM force measurements have been made with the tip whose diameter is several tens of nanometers. In order to evaluate the interaction between a micrometer scaled robot and a substrate, the interaction between a microsphere and a substrate has been considered. Consequently, a borosilicate sphere (from 1 to 40 μm of diameter) has been glued on the cantilever. Force-distance curves were obtained by the measurement's exploitation of the measurement of the AFM cantilever's deformation of the AFM cantilever with a laser beam and a sensitive four-quadrants photodiode. The measurements were performed at the driving speed of 200 nm/s, to stave off the influence of the hydrodynamic drag forces, in 10 different points minimum with different surfaces and borosilicate sphere (at least five) for all the conditions tested (ionic strength and borosilicate diameter). All measurements were done in a liquid medium using the pH 2 to protonate the chemical groups of interest.

3. Results

Substrate and spheres of several sizes from 1 to 40 μm diameter have been functionalized by grafting (3-aminopropyl)triethoxysilane (APTES). The forces between the two surfaces were measured in different points on a sample and on different APTES-modified substrate and AFM tips at pH 2. The surface was kept in the solution for 2 min before starting measurements in order to stabilize the system. Then, the adhesion force was measured with an AFM in which a sphere was glued on the tipless cantilever extremity.

3.1. Influence of the sphere size on the repulsive force measured

The force distance measurements obtained for APTES-modified surface, with borosilicate spheres of different diameters (between 1 and 40 μm), structured surfaces are presented in Figure 1-A

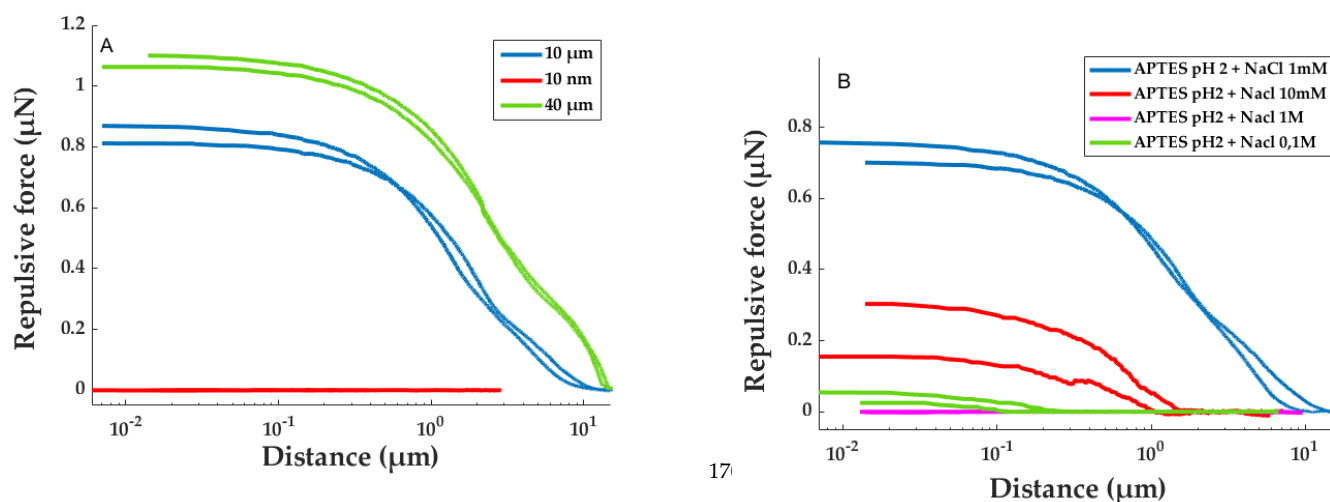


Figure 1. AFM force-distance curves at pH 2 between two APTES surface (one flat surface and one sphere glued on the tipless cantilever): A) Influence of the functionalized borosilicate sphere: 10nm (red), 10 μm (blue) and 40 μm (green) ; B) Influence of the ionic strength for a 10μm borosilicate sphere diameter glued on the tipless: no salt added(blue), 10mM (red), 0.1 M (green) and 1M (pink). All the surfaces were functionalized by APTES.

In Figure 1-A, a cantilever deformation is observed on a long distance (typically several micrometers) when the sphere was approaching the surface. This distance increases with the size of the sphere from few nanometers, for a 10 nm sphere, to 12 μm, for a 40 μm sphere. The variation of the repulsive force is similar to the one observed during previous experiments [28]. The size of the sphere glued on the tipless extremity influenced the repulsive force. Indeed, for a lower sphere diameter (10 nm) no significant repulsion was measured but when the sphere diameter was increased, the repulsive force appeared and it increased with the sphere diameter (Figure 1-A). It reached around 800 nN and 1.1 μN for 10 μm and 40 μm spheres, respectively. The repulsion can be explained by the electrostatic repulsion of the positive charges, at pH 2, of the amine grafted on the tip and on the surface.

3.2. Influence of the ionic strength

To confirm the origin of the repulsion (electrostatic charge), the ionic strength was modified by adding NaCl salt ranging from 10 mM to 1 M in the measuring medium. Then, the force-distance curve was recorded with a 10 μm sphere (figure 1-B). When the medium did not contain any salt, a repulsive force of 700 nN is obtained. The introduction of a small amount of NaCl, to achieve a concentration of 10 mM, results in a 3.5 times reduction of the repulsive force (200 nN). By increasing this concentration to 0.1 M the force does not exceed 50 nN, and finally for a concentration of 1 M, it becomes null. Each force curve presents an hysteresis behavior inducing a measure incertitude. However, the incertitude is still below the observed force reduction inducing by the increase of the ionic strength. So, we can conclude that the origin of the repulsion force is electrostatic.

This conclusion was concordant with the literature where it has been established that the formation of repulsion on a long distance could be explained by a consequent generation of electric field [16, 28, 29-31]. However, the Williams' review concluded that a more complete understanding of the mechanisms behind EZ phenomena will assist in understanding their possible roles in biology as well as their possible engineering applications such as microfluidics and filtration [32]. In this context, we decide to model the

experimental force based on electrical layer theory developed by Gouy, Chapman, Stern and Grahame [33, 34]

4. Model Development

Usually, force measurements are conducted between a sphere and a planar substrate. A coarse model of the interaction between a charged micro-object and a charged surface has been proposed previously [28] based on major assumptions: (i) the object is a sphere; (ii) charges on the object are localized in the centre of the sphere and (iii) the surface is an infinite plane. In this paper, we propose a more precise numerical model able to predict the interaction forces on objects whatever their shape is.

4.1. General case

The presence of a charged surface in an ionic solution induces a specific modification of the medium. If the surface is positively charged, a digressive layer of anions appears around the contact with the surface until the return to the electric equilibrium in the bulk solution. The modeling of object-surface interaction is based on electrical layer theory developed by Gouy, Chapman, Stern and Grahame [33, 34]. An electrical layer's formation, namely the formation of a compact layer of charged ions opposite to that of the surface, was modelled by representing the surface as a set of electric dipoles (see Figure 2).

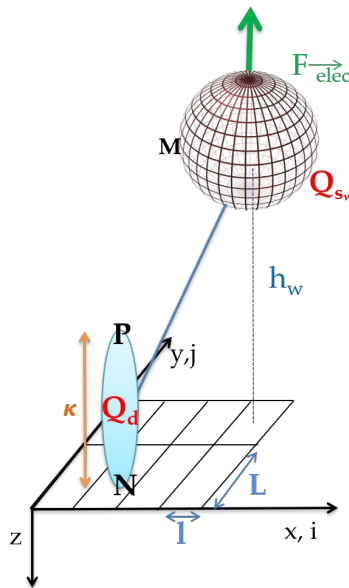


Figure 2. Geometric modeling between charged micro-object and surface. The relative scale was not respected: distance κ is weaker than the diameter of the sphere.

Each elementary dipole on the surface is represented by two electrical charges $+d_{Q_d}$ and $-d_{Q_d}$ separated by a distance κ . These parameters can be determined for each experimental condition depending on the ionic strength (see equation SI-1 distance). Each elementary dipole induces an electric field up to the substrate whose component dE_z along the vertical axis z is defined by (projection of Coulomb law on z axis):

$$dE_z = \frac{1}{4\pi\epsilon_0\epsilon_r} \left(\frac{h_w - \kappa}{MP^3} - \frac{h_w}{MN^3} \right) dQ_d, \tag{1}$$

Each charged particle located in this electric field experiences an electrostatic force in a very similar principle as electrophoresis. Considering an elementary charge dQ_w placed

at a point M on an object (figure 2), the vertical force applied by the elementary dipole on the elementary charge is directly obtained from (1):

$$dF_{elec} = d_{Q_w} dE_z = \frac{1}{4\pi\epsilon_0\epsilon_r} \left(\frac{h_w - \kappa}{MP^3} - \frac{h_w}{MN^3} \right) d_{Q_d} d_{Q_w}, \quad (2)$$

Considering that both the substrate and the object have a uniform charge density (expressed in Coulomb.m⁻²) respectively noted Γ_d and Γ_w , the total vertical force F_{elec} can be written as the integral of the elementary force on the substrate surface S_d of the plane and the surface of the object S_w :

$$F_{elec} = \frac{\Gamma_D \Gamma_w}{4\pi\epsilon_0\epsilon_r} \iint_{S_d} \iint_{S_w} \left(\frac{h_w - \kappa}{MP^3} - \frac{h_w}{MN^3} \right) dS_w dS_d \quad (3)$$

4.2. Sphere-plan case modeling

In order to compare our model with experimental results, a sphere-plane numerical modeling has been developed. Both the surfaces of the microsphere and of the substrate have been sampled in order to calculate numerically the integral in (3) (see figure SI-1). The total force F_{elec} along z -axis applied by the substrate on the total sphere is given by:

$$F_{elec} = \frac{\Gamma_D \Gamma_w}{4\pi\epsilon_0\epsilon_r} \sum_w \sum_{ij} \left(\frac{h_w - \kappa}{MP(i,j)^3} - \frac{h_w}{MN(i,j)^3} \right) \delta S_d \cdot \delta S_{w_{elec}} \quad (4)$$

where MP (respectively MN) represents the distance between the sphere and the top (respectively the bottom) of the dipole (Figure 2) which are defined by:

$$MP(i,j) = \sqrt{(h_w - \kappa)^2 + (l.i)^2 + (L.j)^2}, \quad (5)$$

$$MN(i,j) = \sqrt{h_w^2 + (l.i)^2 + (L.j)^2}, \quad (6)$$

where l and L are the distances of the sampling of the substrate surface along x -axis (vector i) and y -axis (vector j), respectively; w is the number of samples considered on the object surface ($w=5000$ in the next sections).

4.3. Influence of the parameters

The dipole thickness, κ , depends on the ionic strength and also of the electrical theory layer. Indeed, the triple electrical layer (Figure SI-2) is the most complex theory since it takes account of all ions present in the solution in the vicinity of the charged surface. It was modeled as a set of dipoles with different lengths contrary to a simplified model, call the double electrical layer (Figure SI-3) where all the dipole are similar. The major difference between these two models is that the triple layer takes into account the specific adsorption of ions on the surface which creates a division in the compact layer. We sought to determine whether it was possible to avoid the calculation of the exponential decay's calculation of the electric potential in the compact layer. The difference between these two theories leads to figure 3-A. The variation of this force F_{elec} versus the distance z is presented, for the two theories. For numerical application, different sizes of dipoles with a distribution based on the Debye-Hückel approximation were taken into account: a charge density is fixed at 1 charge.nm⁻² for the silicon substrate Γ_d and the borosilicate sphere Γ_w .

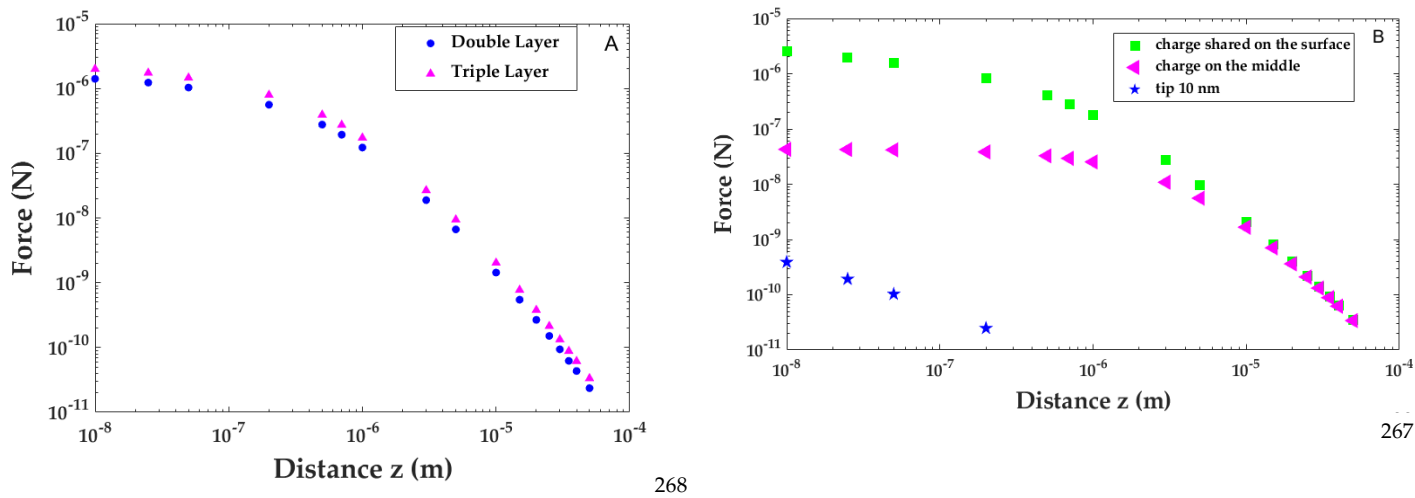


Figure 3. A) Modeling of the interaction between a surface covered by dipoles and a charged sphere of 10 μm diameter by considering the simplified double layer (blue sphere) or triple layer model (pink triangles); B) Modeling of the impact of the charge's discretization on sphere surface. Repulsive force obtained between a dipole surface and a 10 μm charged sphere with the charged fairly shared on the surface (green squares), center at the middle of the sphere as a point charge: the same charge of previously but located in the sphere center (pink triangles) or 10 nm charged tip (blue stars).

Differences between the two models have a low impact on the considered force and interaction distances. Both models vary only by a factor less than 2 and the difference is constant on the considered scale. Simplest model (double layer model) will be used for the following simulations. The results obtained with both models highlight a long interaction distance as the interaction force at 1 μm distance is only 10 times lower than the interaction force at the contact.

The double sum expressed in equation (4) was complex to compute due to long time calculation. In order to reduce this computational time, we also proposed a simplified model considering that the charges of the objects are located in its center of gravity (figure 3-B, pink triangles). This elementary charge has the same value as the sum of initial distributed charges on the entire sphere surface. For the 10 μm sphere, full (green squares) and simplified models (pink triangles) are similar until 10 μm from the substrate, then the simplified model underestimates the force. Indeed, in the full model, charges that are at the bottom of the sphere induce a bigger force than those located at the equator. Therefore, it is thus possible to simplify the simulations when the distances of interaction are high enough, at least the diameter of the considered sphere. Furthermore, it is necessary to use the complete model for the weak interaction distances.

As most of the AFM force measurements are done with AFM tips having an apparent radius of several nanometer (10 to 100 nm typically) [20, 22, 35], we propose to compare our simulated results on microspheres with a sphere of 10 nm, modelling the interaction with an AFM tip (figure 3-B, blue stars). The comparison between microsphere and nanosphere shows that the force is lower on nanosphere even for the same charge density. The most important difference is the interaction distance which is significantly lower on nanospheres (few tens of nm) than the distance on microspheres (few μm). This result is coherent with the current literature assuming that the electrostatic interaction between a substrate and a nanosphere cannot exceed few tens of nm.

5. Model and Experimental Measurements Comparison

The experimental adhesion forces measured were compared with the ones obtained with our model. In this model, the charge of the surface and of the sphere was necessary. So, the zeta potential of the borosilicate sphere with and without APTES functionalization was measured with a Nanosizer (Malvern). At pH 2, the zeta potential was 34.6 mV and 4.3 mV, corresponding to $0.29 \mu\text{C}/\text{cm}^2$ and $0.03 \mu\text{C}/\text{cm}^2$, with and without APTES functionalization, respectively. So, the APTES-modified sphere charge density was $0.02 \text{ charge}/\text{nm}^2$. In the model, we decide to fixe also the charge density on the surface at the same value, $0.02 \text{ charge}/\text{nm}^2$, for all the modelling. Using these parameters, the experimental and modeled data simulated using Matlab Simulink software was compared.

5.1. Influence of the sphere size on the repulsive force measured

The influence of the sphere size on the interaction distance and the maximum repulsive force predicted by the model (Figure SI-5) and measured are presented in table 1. The factor κ (calculated from the equation SI-1) for 10^{-3}M of ionic strength and r_{sw} are fixed (and not fitting) depending on the experimental conditions.

Table 1. Comparison of the repulsive force and distance measured or modeled between silica surface and different diameter size of borosilicate sphere functionalized both by APTES. The data are collected during 3 experimental campaigns. For every campaign new bead and substrate are considered and force are measured in 10 different locations on a substrate.

Sphere diameter (μm)	Experimental		Predicted	
	Interaction distance (μm)	Interaction force (μN)	Interaction distance (μm)	Interaction force (μN)
0.01	0.004 ± 0.007	0.02 ± 0.01	N.d.	$1.10 \cdot 10^{-4}$
1	0.8 ± 0.2	0.07 ± 0.01	0.5 ± 0.05	0.03 ± 0.005
5	3.3 ± 0.5	0.6 ± 0.2	3 ± 0.3	0.34 ± 0.05
10	5.7 ± 0.7	0.7 ± 0.15	5 ± 0.5	0.64 ± 0.1
20	8.2 ± 0.6	0.9 ± 0.25	7.5 ± 0.75	1.2 ± 0.2
40	12.7 ± 0.29	1.1 ± 0.12	12.5 ± 1.2	2 ± 0.4

N.d. the distance was too weak in order to be estimated.

The experimental data (force values and distance of repulsion) are consistent with the results of the modeling. Indeed, in all the case, the same order of magnitude was obtained. The variation of the difference between the expected and the experimental repulsive force could be explained by a small variation of the charge density. Indeed, each sphere was functionalized with an individual APTES solution and so, the grafting percentage could be slightly changed between each sphere diameter. The charge density of each sphere was not measured individually. For the predicted value, we take an average value of the density charge measured on $10 \mu\text{m}$ radius borosilicate sphere functionalized by APTES. It should be noted that the forces obtained are substantially greater than the weight of the objects considered (11 pN for the $10 \mu\text{m}$ -diameter borosilicate sphere).

5.2. Influence of the ionic strength

As the behavior of these objects is essentially governed by these electrostatic forces, the impact of the ionic force on both the experimental data and the model has been tested and is reported (Figure SI-6) in table 2. For that, the distance κ was calculated from the equation SI-1 for each experimental condition and r_{sw} is fixed at $5 \mu\text{m}$.

Table 2. Comparison of the repulsive force and distance measured or modelled between silica surface and 10 μ m borosilicate sphere functionalized both by APTES for different ionic strength controlled by NaCl addition. The data are collected during 3 experimental campaigns. For every campaign new bead and substrate are considered and force are measured in 10 different locations on a substrate.

Ionic Strength : NaCl (M)	Experimental		Modeling	
	Interaction force (μ N)	Interaction distance (μ m)	Interaction force (μ N)	Interaction distance (μ m)
0	0.7 ± 0.15	5.7 ± 0.7	0.6 ± 0.1	5 ± 0.5
0.01	0.2 ± 0.1	0.5 ± 0.43	0.3 ± 0.045	2 ± 0.2
0.1	0.05 ± 0.01	0.2 ± 0.1	0.09 ± 0.02	0.7 ± 0.07
1	0.03 ± 0.01	0.1 ± 0.06	0.03 ± 0.005	0.2 ± 0.02

A good concordance was noted between the experimental and the predicted repulsive force. When NaCl salt is added in the medium, these two ions come into contact with the dipoles, on the opposite side of their charge and create an electric balance. This behavior has the effect of skimming the loads of the double layer and extending the load of dipoles, reducing electrostatic interactions with the sphere and decreasing the length of Debye. As the number of interactions is thus reduced and the length of Debye decreases, the electric field generated by the surface at the height of the sphere also decreases.

5.3. Repulsive force on polymer film

The equation (4) predicts also that the repulsive force increases with the density charge of the surface or of the sphere. In order to validate the model, the charge density of the surface was modified by changing the molecule deposited on the surface.

Previous studies have demonstrated that polymer film could be build-up a large repulsion distance to solutes when immersed in an aqueous solution greater than 200 μ m [16]. To increase the charge density of the deposited and to better localize it, an electropolymerization of a film was performed. The polypyrrole was chosen due to the presence of NH groups and its ability to be deposited locally on electrodes (figure SI-7). The electrodeposition was performed by potentiostatic chronoamperometry on a silicon substrate previously covered by sprayed chrome and gold to enhance the conductivity. Different counter ions (0,1M) were used to determine the impact of each other on the repulsive force. The evolution of the current intensity with time is similar whatever the supporting salt used (figure SI-8). However, the charge density values differ from one salt to the other: the highest charge density is obtained with tetrafluoroborate anions when the lowest one is obtained for toluenesulfonate anions which is consistent with literature [26]. The PPy/LiClO₄ film has granular structure and covers the whole surface of the substrate (figure SI-9) even if it can be seen that the thickness of the film is not totally uniform due to its surface roughness. The same tendency is observed with the other supporting salts used in this work.

Indeed, the counter-ion used during the synthesis of the polymer film plays an important role in the structure and thus on the film morphology (Figure SI-10) [25,26]. So, its influence on the repulsive force has been studied at pH 2, between a polypyrrole-modified surface and a 10 μ m sphere functionalized by APTES (figure SI-11) and summarize in Table 3. In Table 3, a repulsive force was measured with an amplitude between 2.5 and 4.9 μ N and a repulsive distance upper to 33 μ m which are higher than the previous experiment (Figure 1-A (blue curve), and table 1), with the aminosilane APTES. This increase can be explained by the higher density of the amine groups in polymer films compared to aminosilane-film.

Table 3. Repulsive force and distance measured at pH 2 for different counter-ions used during the pyrrole electrodeposition.

Salt	interaction force (μN)	interaction distance (μm)
LiClO_4	4.9 ± 0.2	47 ± 4
PTS	2.5 ± 0.1	33 ± 3
NaNO_3	3.6 ± 0.1	38 ± 4
NaBF_4	4.1 ± 0.2	45 ± 5

Whatever the counter-ions used, the repulsive forces with the polypyrrole were at least 4 times higher than with APTES-modified surfaces which confirms that an increase of the charge density increases the force as predicted by the model. The repulsive force differences can partially be explained by the morphology of the polypyrrole film depending on the counter-ion used. Indeed, previous experiments have demonstrated that the electrochemical film was only influenced by the anion of the counter-ion whatever the cation Li^+ or Na^+ [25,26]. ClO_4^- , BF_4^- and NO_3^- anions have a similar size (approximately 6 Å) while PTS anions have a higher diameter (10 Å) [36]. This size difference impacts directly the film morphology: the use of a small anion leads to the formation of a film with a high roughness whereas the use of a bigger anion' use as the PTS leads to a more homogeneous and planar structure (figure SI-10). We can directly bind the roughness of the substrate at the interaction strengths. Indeed, the higher the roughness of the substrate is, the higher the specific surface is. Consequently, the number of interactions increased and so the induced force. So, the surface's roughness enables to enlarge electrostatic forces in liquid medium. It should be noticed, that for adhesion forces in the air, the impact is the opposite. The roughness reduces the size of the contact surface's size induced by local mechanical deformation and thus reduces the adhesion force [37-39].

6. Applications

The repulsive force generated between the surface and the object could be used for micromanipulation task. Indeed, if we compare the repulsive force generated by the interaction between a microsphere and a surface in pH 2 ($4.9 \mu\text{N}$, Table 3) with the weight of the microsphere of borosilicate (11 pN), it appears that it will be easy to place it easily in levitation up to a substrate. This may be interesting for the non-contact micromanipulation to guarantee that a manipulated object will never come in to contact with the substrate.

We may also imagine other ways to exploit this high density of charges on micro-objects in non-contact micromanipulation. Indeed, the chemical functionalization generates an important charge density on the sphere which may induce high electrophoresis force when located in the electric field. These properties can be used to control the trajectory of the object by electrophoresis. As an example, we consider a sphere functionalized with APTES located in an electric field controlled with electric voltage located on micro-electrodes. We consider, a working space made up of four square-electrodes of $100 \mu\text{m}$ each in aqueous middle, and a voltage applied to each of the electrodes one after the other every 2 seconds clockwise. The point of electrolysis of the water being 2 V, it is impossible to apply larger tensions (electrolysis bubbles would perturb object manipulation), we thus consider a voltage of 1.8 V. The results of the simulation, presented in figure 4-A, show the trajectory of the micro-sphere. The travel speed increases exponentially between each position (electrode) from $80 \mu\text{m}\cdot\text{s}^{-1}$ at the beginning of the movement to $170 \mu\text{m}\cdot\text{s}^{-1}$ at the end of the trajectory (Figure SI-12). As the electrophoresis and the drag forces are surface forces, the behavior of the micro-sphere and thus its speed will be the same whatever the radius of the micro-sphere. If we compare these results with the dielectrophoresis which generates a volume force [5] the manipulation speed varies from $25 \mu\text{m}\cdot\text{s}^{-1}$ to $1000 \mu\text{m}\cdot\text{s}^{-1}$ for an object diameter from $2 \mu\text{m}$ to $80 \mu\text{m}$, respectively. So the electrophoresis combined

with surface functionalization appears an interesting alternative to dielectrophoresis when the object size is lower than 10 μm .

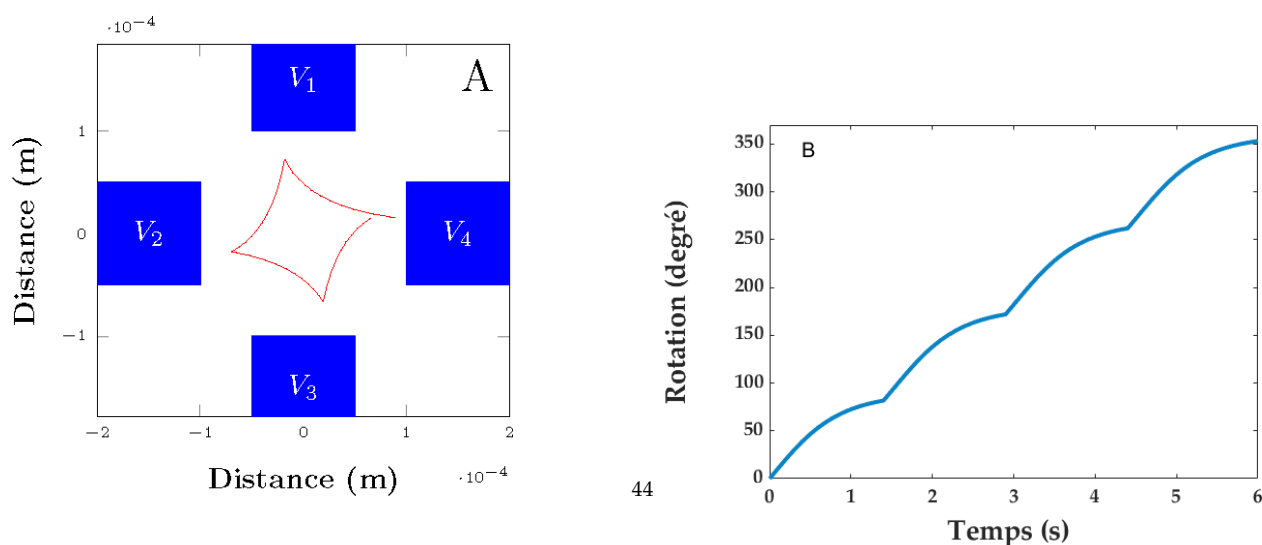


Figure 4. A) Trajectory modelling for a microsphere functionalized by APTES by dielectrophoresis; b) Rotation modelled for rectangular object (20 μm) functionalized by amine and carboxylic groups each on a side of the object.

Moreover, several different chemical functions may be placed on different locations on a micro-object to generate more advanced behaviors. Concretely, the orientation of the object may also be controlled by the combination of localized surface functionalization and electric field. As an example, considering an object having both oxide surfaces and conductive surfaces, a polymer electrodeposition (APTES) can be localized on conductive surface (figure SI-7) whereas the silane molecules could be grafted on oxide surfaces. If the object is grafted with localized amine and carboxylic groups on each side, it generates a high electric dipole (Figure SI-13). The object will behave as a dipole in an electric field and experiences a high torque enabling to control its rotation. We demonstrated through simulations that the total rotation of a micro-sphere may be realized in 6 s with the electrodes defined previously and by applying to the opposite electrodes +0,9V and - 0,9V simultaneously every 1.5 seconds to maximize the speed (figure 4-B).

7. Discussion

Controlling the localization of a high number of charges on surface may have a lot of applications. This paper illustrates an original way to induce high electrostatic forces on micro-objects along several micrometers and opens the way to application in non-contact manipulation of micro-objects. The electrostatic force depicted in this article can be used in two different ways in non-contact manipulation.

The first approach consists in using the electrostatic force to place the object in levitation several micrometers up to a substrate and the use an other actuator to move the object parallel to the substrate. In such a case both the substrate and the micro-object have to be functionalized.

The second approach is closer to electrophoresis, where a functionalized object can be moved in an electric field induced by microelectrodes. The possible performances (typ. manipulation velocity) have been compared with more usual principles such as the dielectrophoresis and our approach shows a potential interest compared to the state of the art.

Moreover, both approaches (chemical and physical) could be combined in the same device in the future. As an example, the bottom face of the sorting microchannel can have a chemical pattern in order to generate passive levitation of microparticles in the channel avoiding sedimentation. The top face could be structured with electrodes enabling to actively sort microparticles.

The long distance of repulsion was already mentioned on several studies with a distance close to 200 μm [32]. Indeed, a large exclusion zone (EZ) was observed in the vicinity of metals [40], hydrogels [41], ion exchange polymer [42], biological tissues [43], white blood cell [29], Nafion polymer [16] and self-assembled monolayer (SAM) [28, 44]. Important research was performed to understand such peculiar phenomenon during the last years. Several hypotheses were developed as: i) water structuring [45], ii) a pH gradient and therefore charge separation [46], iii) chemotaxis driven by solute gradients (OH^- , H^+ and salts in solution) [30], iv) a combination of ion-exchange at the surface, diffusion of ions, and diffusiophoresis of particles in the resulting ionic gradients [31]. Recently, a review analyzed the different theory to explained the EZ and concluded several major problems with the theory that water in the EZ undergoes a phase change or significant reordering. They added that Schurr's theory [30] of macroscopic chemotaxis presents a compelling alternative theory which can explain experimental findings but there are still many open questions about exclusion zones.

8. Conclusion

In this paper, we have studied the interaction behavior, and most precisely the repulsive force, between a functionalized surface and borosilicate spheres. The experiments were performed as a function of the borosilicate sphere diameter, the medium ionic strength and the surface charge density. The experimental measurements were compared to a precise numeric model able to predict the interaction forces between a charged surface and a charged micro-object whatever their shape and a good agreement was observed. The surface functionalization by polymer electrodeposition way allowed to generate an electrostatic interaction through electrical charges of chemical origin connected to the ionic strength of the measure medium. This interaction can be characterized by two elements, its strength and its distance. In both cases, the results obtained are innovative and more raised than the repulsive strengths usually met in the chemical systems of typical colloidal suspensions. Because adhesion is the highest current disturbance in micromanipulation (positioning and releasing), the surface functionalization is a promising way to improve micro-robotics efficiency and accuracy, and to control electrostatic forces in non-contact microrobotic applications. A wide range of applications, in the fields of telecommunications, bioengineering, and more generally speaking MEMS can be also envisaged for these functionalized micro-grippers.

Supplementary Materials: The following supporting information can be downloaded at: www.mdpi.com/xxx/s1, Equation SI-1 to SI-, Figure S1: Band by band modeling of the sphere; Figure S2: triple electrical layer, Figure S3: Double Electrical Layer, Figure S4: Modeling of the impact of the sphere band number on the force for a borosilicate sphere with a 10 μm of diameter, Figure S5: Modelling of the repulsive forces and distance between APTES film on surface and on different diameter size of borosilicate sphere glued on the tip extremity, Figure S6: Modelling of the repulsive forces and distance between APTES film on surface and on 10 μm borosilicate sphere glued on the tip extremity for different ionic strength, Figure S7: SEM images of the electrodeposition localization of polypyrrole with LiClO_4 on a gold electrode, Figure S8: Evolution of charge density with time for different supporting salts, Figure S9: SEM image of the PPy/ LiClO_4 film, Figure S10: Polypyrrole film morphology versus the counter-ion used: A) ClO_4^- , B) PTS , C) NO_3^- and D) BF_4^- , Figure S11:

Repulsive force measured at pH 2 at three different points on the polypyrrole film electrodeposited with: A A) ClO₄, B) PTS, C) NO₃⁻ and D) BF₄⁻, Figure S12: Electrophoresis simulation of functionalized microsphere, Figure S13: Rotation simulation of a micro-object functionalized by two molecules of opposite charge.

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