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# Observation of a Plastic Crystal in Water-Ammonia Mixtures under High Pressure and Temperature

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## Abstract

Solid mixtures of ammonia and water, the so-called ammonia hydrates, are thought to be major components of solar and extra-solar icy planets. We present here a thorough characterization of the recently reported high pressure (P) - temperature (T) phase VII of ammonia monohydrate (AMH) using Raman spectroscopy, x-ray diffraction and quasi-elastic neutron scattering (QENS) experiments in the range 4-10 GPa,

450-600 K. Our results show that AMH-VII exhibits common structural features with the disordered ionic-molecular alloy (DIMA) phase, stable above 7.5 GPa at 300 K: both present a substitutional disorder of water and ammonia over the sites of a body-centered cubic lattice and are partially ionic. The two phases however markedly differ in their **hydrogen** dynamics, and QENS measurements show that AMH-VII is characterized by free molecular rotations around the lattice positions which are quenched in the DIMA phase. AMH-VII is thus a peculiar crystalline solid in that it combines three types of disorder: substitutional, compositional and rotational.

The physics of water, ammonia and other molecules under extreme P-T conditions has been a topic of intense research in the last decades, for their importance in many fields and in particular in the modeling of the interior of icy planets and their satellites.<sup>1,2</sup> One remarkable finding is that for pressures above a few GPa and T below 300 K, the configurations of the oxygen and nitrogen sublattices – body-centered cubic (*bcc*) for water ice, quasi hexagonal close-packed (*hcp*) for solid ammonia – are kept intact at least up to the Mbar (100 GPa).<sup>3-6</sup> By contrast, pressure has a strong effect on the average position and dynamics of H atoms, driving the symmetrization of H bonds in H<sub>2</sub>O<sup>7-10</sup> and ionization of NH<sub>3</sub><sup>11-13</sup> at low T. Coupling high temperatures to high pressures gives rise to other dynamic regimes for the **hydrogen atoms** where plastic states, characterized by free rotations of the molecules around their lattice positions,<sup>14-19</sup> and superionic states, characterized by fast diffusion of the protons through the oxygen or nitrogen lattice,<sup>19-25</sup> can be promoted.

In contrast with the pure ices, there is still little known on the properties and phase diagram of their mixtures under extreme P-T conditions despite the fact that they are more relevant systems for planetary modelling. At ambient pressure, the ammonia-water mixtures crystallize in three stoichiometric compounds, known as ammonia hydrates, with, respectively 1:2, 1:1 and 2:1 ammonia to water ratio. We focus here on the 1:1 compound, ammonia monohydrate (AMH). Early neutron studies have mapped the phase diagram of AMH up to 7 GPa and below 300 K, finding several molecular phases termed I to VI.<sup>26,27</sup>

AMH-VI, formed above 6.5 GPa and 270 K, is of particular interest as it was described as a disordered molecular alloy (DMA) where water and ammonia randomly occupy the sites of a simple *bcc* structure. More recently, Liu et al.<sup>28</sup> found that the DMA phase is actually a partially ionic phase, where ammonia and water coexist with  $\text{NH}_4^+$  and  $\text{OH}^-$  ions, in about 2:1 molecular to ionic ratio at 10 GPa and 300 K. DMA was thus renamed DIMA for disordered ionico-molecular alloy. Liu et al.<sup>28</sup> also found that DIMA phase may coexist with a purely ionic and ordered phase of *P4/nmm* space group, predicted as the ground state by calculations<sup>29</sup> but only seen in minor proportion in experiments. AMH-VI is thus a variable mixture of the DIMA and *P4/nmm* phases, and the question whether one is more stable than the other has not been settled yet.

At high T, Zhang et al.<sup>30</sup> recently reported that another AMH phase, named AMH-VII, becomes stable above 3.6 GPa–324 K and up to at least 8.3 GPa – 675 K. AMH-VII melts congruently, showing that it has the same 1:1 stoichiometry as the liquid but no other structural information has been made available to date. Interestingly, computer simulations based on the density functional theory (DFT)<sup>29,31–33</sup> have predicted the existence of plastic and superionic states in ammonia hydrates appearing at P-T conditions compatible with the range of stability of AMH-VII. Directly probing these exotic states by experiment is very challenging as this requires techniques sensitive to H atoms mobility, such as neutron scattering, proton NMR and conductivities measurements, which are difficult to combine with extreme P-T conditions.

In this letter, we report on the structural and dynamic properties of AMH-VII determined by a combination of x-ray diffraction (XRD), Raman scattering and quasi elastic neutron scattering experiments (QENS). Recent advances in QENS techniques<sup>34,35</sup> enabled us to measure the self-dynamics structure factor of AMH to the record pressure of 8.1 GPa. Our results demonstrate that AMH-VII is a crystalline plastic form of the DIMA phase where molecular ( $\text{H}_2\text{O}$ ,  $\text{NH}_3$ ) and ionic species ( $\text{OH}^-$  and  $\text{NH}_4^+$ ) are randomly arranged on a body-centered cubic (*bcc*) lattice and present a dynamic rotational disorder of the **hydrogen atoms**.

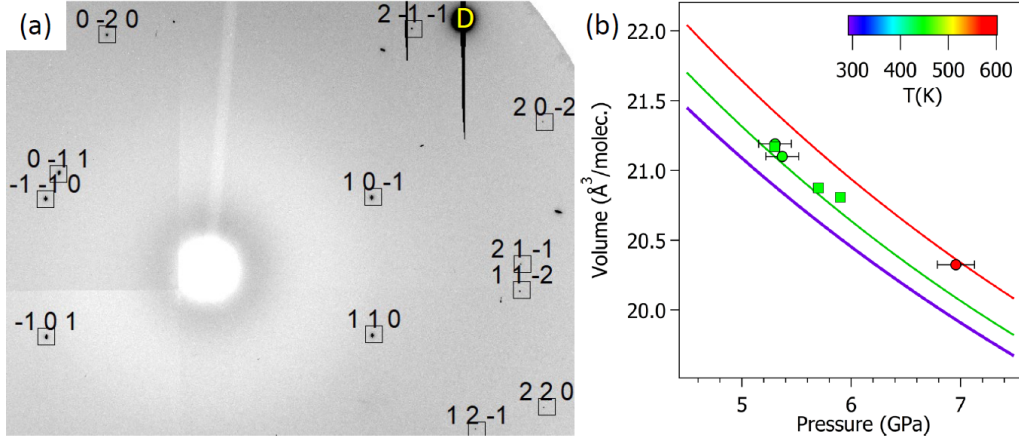


Figure 1: (a) Details of the diffraction image of a single crystal AMH-VII at 5.4 GPa and 453 K overlaid with the indexing of reflections. The saturated spots noted D are from the diamond anvils. (b) Molecular volume vs. pressure in AMH. The symbols represent the measured volumes of AMH-VII at high temperature. The purple curve is a fit to experimental data at 300 K (see SI section S2.2). The green and red lines are the calculated EOS at 450 K and 550 K, respectively, using the thermal dilatation coefficient of ice VII<sup>36</sup> as an estimate for that of AMH.

Experimental methods are presented in details in the Supplementary Materials (SM). *In-situ* HP measurements were made in diamond anvil cells (DAC) for XRD and Raman scattering, and in the Paris-Edinburg press (PE) for neutron scattering. Synchrotron XRD experiments were performed at the PSICHE beamline of SOLEIL (Saint Aubin, France) and the ID27 beamline of ESRF (Grenoble, France), and QENS experiments were performed at the IN6-Sharp beamline of ILL (Grenoble, France).

In order to determine the structure of AMH-VII, we grew single crystals by slowly compressing the liquid at temperatures above 324 K until a single crystal germ crystallized and filled the sample chamber. Part of a diffraction image collected at 5.4 GPa and 453 K is shown in Fig.1(a). All observed reflections could be indexed by a unique *bcc* unit cell with  $a = 3.4809(1)$   $\text{\AA}$ . A single structural solution was found by ShelXT<sup>37</sup> in space group *Im-3m*, which contains one atomic site for oxygen and nitrogen atoms at (0,0,0), replicated by symmetry at (0.5,0.5,0.5). This implies substitutional disorder for O and N. Using an equal O/N occupation as imposed by stoichiometry, the refinement with ShelXL<sup>38</sup> converged to a

R1-factor of 4.7% (see SI Table 1). We note however that the refinement is not sensitive to the O/N occupation, due to the small difference in the number of electrons ( $\Delta Z = 1$ ) and likely to the absence of H atoms in the structural model. As a matter of fact, we could not locate any H atom from the difference Fourier map, which suggests that they are all disordered. As far as the position of the O/N atoms is concerned, the structure of AMH-VII is thus identical to that reported for the DMA<sup>26</sup> or DIMA<sup>28</sup> phases at room temperature.

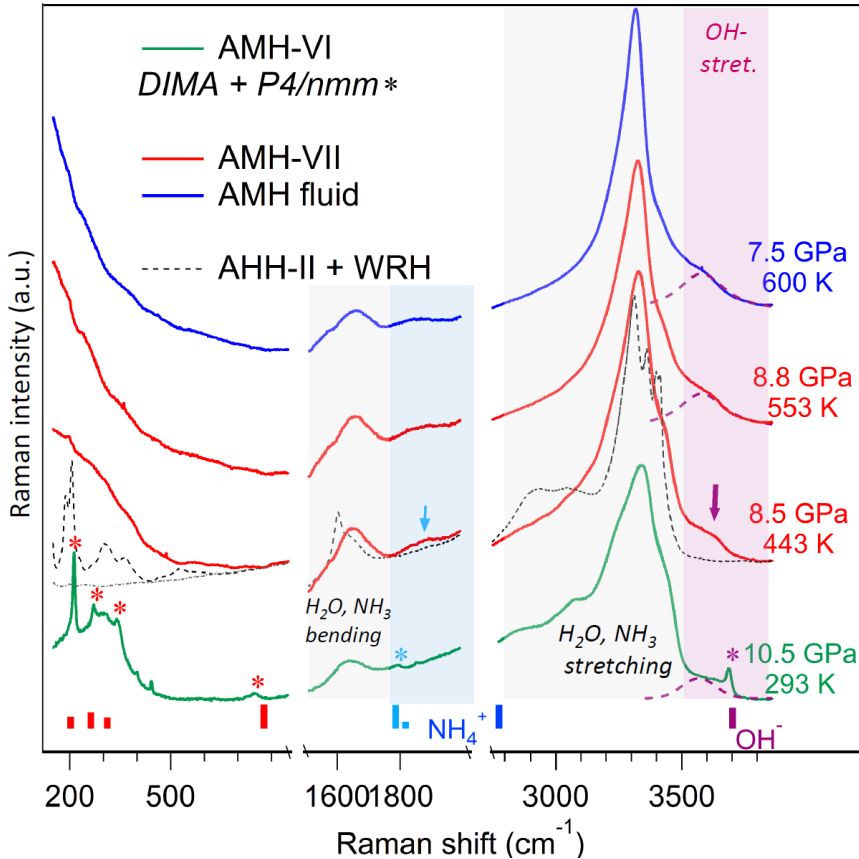


Figure 2: Comparison between the measured Raman spectra of AMH-VI (green), AMH-VII (red), AMH-fluid (blue) and AHH-II+WRH (gray dotted line) at P-T conditions indicated in the figure. The frequency windows 1200–1400 cm<sup>-1</sup> and 2300–2700 cm<sup>-1</sup> dominated by, respectively, the first and second-order Raman signal from the diamond anvils have been omitted. The bottom ticks and asterisks represent the peak positions for the *P4/nmm* phase predicted by DFT and experimentally observed,<sup>28</sup> respectively. The arrows emphasize bands assigned to ionic species in AMH-VII.

For all studied AMH-VII samples – either single crystals formed as above, or polycrystals obtained by heating powder samples (Fig. S1 in SI for a diffraction pattern of a polycryst-

talline sample), no evidence of the  $P4/nmm$  phase observed to coexist with DIMA at 300 K<sup>28</sup> was found. The measured molecular volumes in the intervals 5-7 GPa and 433-570 K are shown in Fig. 1(b). When taking thermal expansion into account, they agree very well with the equation of state of AMH-DIMA at 300 K (see SI section S2.2 and Fig. S2), confirming that AMH-VII is a single phase with the same 1:1 stoichiometry as the loaded mixture .

As recalled above, the unique difference between the structural models of DMA and DIMA is that the former is a purely molecular alloy while the latter also contains  $\text{OH}^-$  and  $\text{NH}_4^+$  ions. The questions thus arises whether AMH-VII is purely molecular (DMA-like) or a mixed molecular-ionic (DIMA-like) alloy. To answer this, we measured the Raman spectrum of AMH-VII at various P-T. The Raman spectrum of AMH-VII at 8.23 GPa-499 K is compared in Fig. 2 to those of AMH-VI (DIMA+ $P4/nmm$ ) at 10 GPa-293 K, the equimolar liquid at 7.5 GPa-600 K and the demixed sample AHH-II+WRH<sup>30</sup> at 8.7 GPa-415 K.

Comparison of the Raman spectra of phases VII and VI reveals both strong similarities and singularities. First, the Raman peaks assigned to the molecular species  $\text{NH}_3$  and  $\text{H}_2\text{O}$  (the O-H and N-H stretching modes at 3000-3500  $\text{cm}^{-1}$  and the bending modes around 1600  $\text{cm}^{-1}$ ), are located in the same frequency ranges in phases VII and VI and have similar shapes. Their broad bandwidths are typical of **hydrogen** disordered phases (see SI Fig. S3 for a tentative deconvolution of the stretching band). Second, the sharp peak centered at 3700  $\text{cm}^{-1}$  in AMH-VI is absent in AMH-VII. The former was assigned to the  $\text{OH}^-$  stretching mode in the ionic  $P4/nmm$  phase at 300 K<sup>28</sup> (see Fig. 2), thus its absence in AMH-VII samples is consistent with our XRD observations above. On contrary, the broad band centered at 3600  $\text{cm}^{-1}$  (pink peak in Fig 2) is present in both AMH-VI and AMH-VII but not in the molecular phases. Calculations in Ref.<sup>28</sup> showed that this band most likely corresponds to  $\text{OH}^-$  stretching modes in the DIMA phase, which is also consistent with the fact that it hardens with pressure as in  $P4/nmm$  (see SI Fig. S4). Its broader bandwidth in the DIMA (RT) and AMH-VII (HT) phases compared to  $P4/nmm$  can be explained by the

substitutional disorder in the former two, leading to various possible environments for the ions, and/or by a shorter lifetime of these ions.

Finding the spectroscopic signature of  $\text{NH}_4^+$  ions is more difficult as their N-H stretching frequencies are expected around  $2800\text{-}2900\text{ cm}^{-1}$ , that is close to the strong second-order Raman band of the diamond anvils. The  $\text{NH}_4^+$  twisting modes, by contrast, are well isolated (around  $1850\text{ cm}^{-1}$  in  $P4/nmm$ ) and we observe a weak and broad band in this region for AMH-VII but not for the molecular phases. Here again the broad bandwidth is assigned to disorder.

Moving to the low-frequency range, a marked difference can be observed : AMH-VII presents a strong, featureless Raman activity extending from  $500\text{ cm}^{-1}$  down to the lowest measurable frequency ( $90\text{ cm}^{-1}$ ), while the Raman spectra of AMH-VI contains three sharp peaks assigned to the lattice modes of  $P4/nmm$ ,<sup>28</sup> riding over a broad band [ $100\text{-}400\text{ cm}^{-1}$ ] from the DIMA phase. The intensity difference (AMH-VII - AMH-VI) in the lattice region, shown in SI Fig. S5, evidences an additional Raman activity in AMH-VII at frequencies below  $200\text{ cm}^{-1}$  whose intensity increases with temperature. Such a low-frequency signal, together with the broadening or disappearance of the lattice modes, have been identified as signatures of a plastic crystal in  $\text{NH}_3\text{-II}$ ,<sup>39</sup>  $\beta\text{-N}_2$ <sup>40</sup> and in other molecular systems.<sup>41,42</sup> It can also be seen in Fig. 2 that the Raman spectrum of AMH-VII is very similar to that of the liquid phase at  $7.5\text{ GPa}$ , a feature which has been previously noted in pure ammonia between the liquid and the plastic phases II and III.<sup>43</sup>

To obtain direct evidence of plasticity, we performed QENS measurements on an equimolar mixture of  $\text{D}_2\text{O}$  and  $\text{NH}_3$  in the range  $2\text{-}8.6\text{ GPa}$ ,  $300\text{-}523\text{ K}$ . Owing to the very large incoherent scattering cross section of H atoms compared to that of oxygen, nitrogen or deuterium, the incoherent quasi-elastic scattered signal from a hydrogenated sample is dominated by the scattering processes from the H atoms in the molecules. QENS experiments are thus well suited for studying the single-particle dynamics of hydrogenated molecules, which includes both the molecular scale translational diffusion and rotational relaxation. The use



of deuterated water was chosen here both to reduce the amount of multiple scattering and auto-absorption of the sample, and to focus on the ammonia molecules reorientational dynamics.

A selection of QENS spectra as a function of exchanged energy  $E$  and wavevector  $Q$  are shown in Fig. 3 for the liquid and AMH-VII (see SI Fig. S6 for additional data), and the investigated P-T points are indicated in the phase diagram of Fig. 4(a). At high T and low  $Q$ , the vibrational, translational and reorientational motions of the molecules are decoupled<sup>44</sup> and the QENS spectra can be modelled by a sum of Lorentzian functions whose width is proportional to the time scales of molecular diffusion and rotation. The details of the data reduction, fit function and procedure are given in the SI section S4. The best fit curves and their components are reported in Fig. 3 (see also SI Fig. S6).

In the liquid phase, both a translational and a rotational contribution are observed. The measured molecular translational  $D_T = 1.1(2) \times 10^5 \text{ cm}^2/\text{s}$  and reorientational  $D_R = 0.30(5) \text{ ps}^{-1}$  diffusion coefficients are, respectively, about 1/6 and 1/2 the ones measured in liquid water at similar P-T conditions.<sup>45</sup> To our knowledge, there is no corresponding data in the literature for ammonia at these P-T conditions.

For all P-T points corresponding to the AMH-VII solid, we observed that the translational diffusion is frozen but the rotational dynamics is still active, although the rotations slow down by about a factor 3 [see the evolution of  $D_R$  with pressure in Fig. 4(b)]. By contrast, no broadening of the elastic peak other than the instrumental resolution is observed at 8.1 GPa-300 K, showing that the molecules do not rotate in the ambient temperature phase VI (see SI Fig. S7). We may thus conclude that AMH-VII is a plastic solid to at least 8.6 GPa and 523 K, although further experiments are needed in order to probe the H motions in water molecules. **We note in this regard that pure water ice has been predicted<sup>14,15,46</sup> to become a plastic *bcc* ice in a similar P-T range above 352 K, 6 GPa,<sup>15</sup> so it is likely that water molecules also rotate in *bcc* AMH-VII.** As seen in Fig. 4(b),  $D_R$  is, within uncertainties, independent on pressure and temperature conditions in the plastic solid VII. This behavior

can be understood taking into account that in the limited  $P$  range investigated here, the hydrogen bond strength has a very weak variation due to the low compressibility of the crystal. Moreover, the number of first neighbours of molecules remains constant, both with pressure and temperature, and the molecular rotation is not phonon assisted. By contrast, in the liquid, the molecular rotation is coupled and assisted by the molecular translational diffusion, therefore the reorientational motion is faster, and has a stronger dependence on the thermodynamic conditions.<sup>45,47,48</sup>

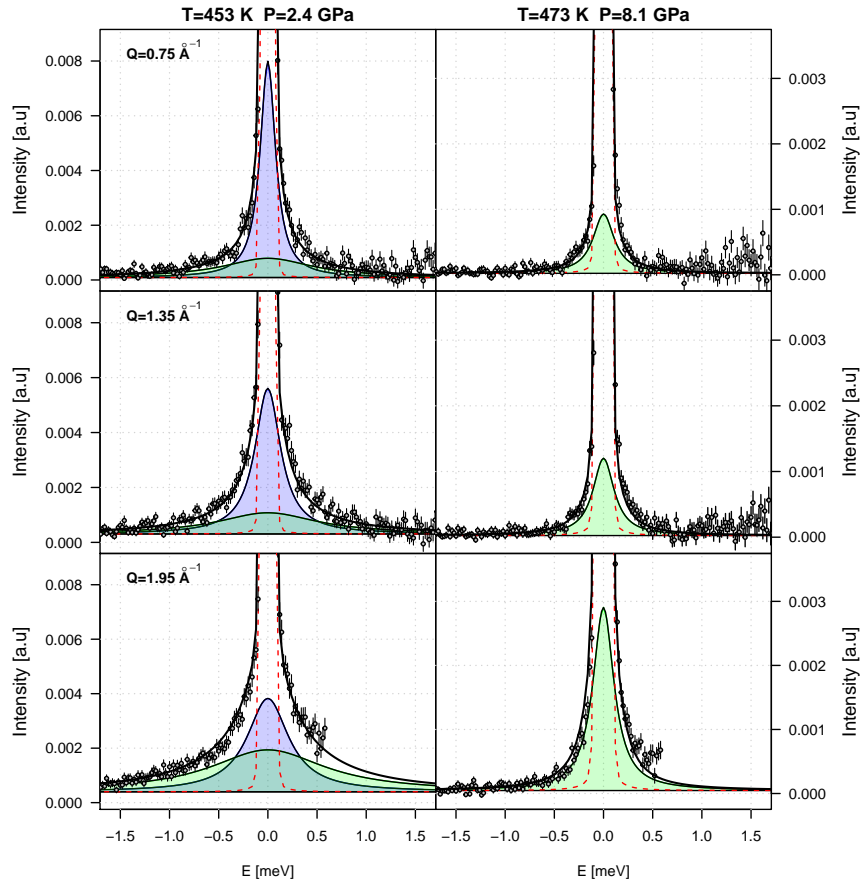


Figure 3: Examples of QENS spectra obtained in the liquid (left) at 2.4(6) GPa–450 K, and AMH-VII(right) phases, at 8.1(3) GPa–473 K for selected values of  $Q$ . Black symbols with error bars are experimental data, the red dotted line is the instrumental resolution, and the blue and green lines represent the translational and rotational components obtained from fits, respectively.

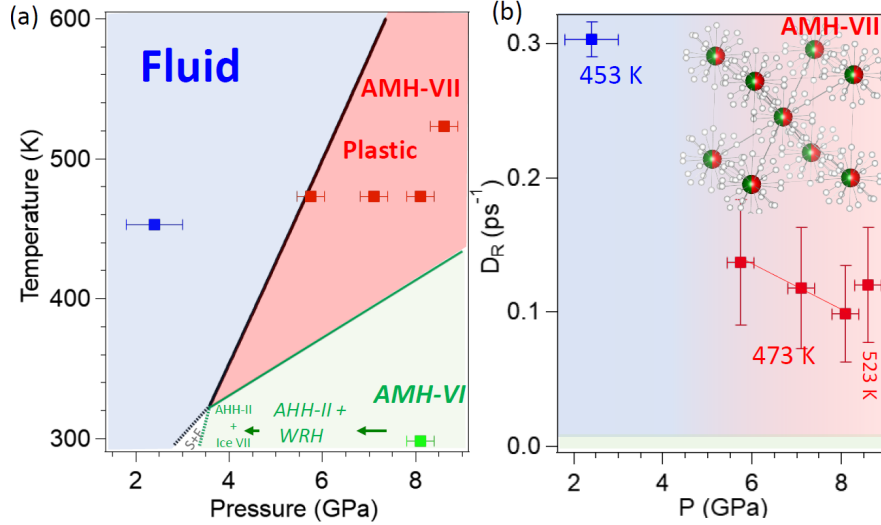


Figure 4: (a) Phase diagram of AMH. The symbols show P-T points where QENS data were collected. (b) Evolution of the rotational diffusion coefficient as a function of pressure for different temperatures. The insert shows a schematic description of the unit cell of the plastic AMH-VII.

In summary, the present work shows that AMH-VII is built on the same *bcc* structure with substitutional O/N disorder as the room-temperature DIMA, and is also composed of ions and molecules with molecular species in major proportions. Conversely, the rotational dynamics in AMH-VII markedly differs from that of DIMA, and we show through QENS experiments that the present species freely rotate as in a plastic solid. AMH-VII is thus a crystalline plastic form of the room-temperature DIMA phase, compiling three types of disorder: substitutional, compositional (molecules and ions) and rotational. Plasticity has been recently predicted in AMH by ab initio molecular dynamics simulations,<sup>32</sup> starting from the H-ordered and fully ionic  $P4/nmm$  phase. It would be interesting to determine how the initial substitutional and compositional disorder of DIMA influences the onset of plasticity, however such calculations are challenging as they require large simulation boxes to account for disorder. We found no evidence in the present QENS experiments for proton diffusive motion in AMH-VII up to 8 GPa and 523 K, as one would expect in the superionic state. This indicates that superionicity requires higher P-T conditions where both rotational and translational protonic motions are active, as observed in pure water ice.<sup>22</sup> Ref.<sup>32</sup> reported the

onset of superionic conduction in AHH around 10 GPa and 600 K, that is, slightly higher P-T conditions than achieved in present experiments, and probably within reach of the QENS technique in the near future. Finally, the P-T domain of stability of plastic AMH-VII is close to the present conditions in some Jovian satellites such as Titan, Ariel and Enceladus<sup>49,50</sup> and sub-Neptune size exoplanets.<sup>51</sup> Its presence during the formation of these bodies should thus be taken into consideration.

### **Supporting information :**

**The supporting Information is available free of charge at [link to the pdf SI]  
It contains additional experimental methods, details of X-ray diffraction experiments (including structure of AMH-VII, equation of state of AMH), Raman experiments (including Raman spectra of AMH-VII and VI, pressure evolution) and Quasi-Elastic Neutron Scattering experiments (including methods, data analysis, QENS spectra of AMH-VII and liquid and their comparison with AMH-VII,  $Q$  dependence of the HWHM of the rotational term and QISF of AMH-VII).**

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