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Zeyu Chi, Tamar Tchelidze, Corinne Sartel, Tsothe Gamsakhurdashvili, Ismail Madaci, et al.. Assessment of Large Critical Electric Field in Ultra-wide Bandgap p- type Spinel ZnGa<sub>2</sub>O<sub>4</sub>. Journal of Physics D: Applied Physics, 2023, 56 (10), pp.105102. 10.1088/1361-6463/acbb14 . hal-03997451

**HAL Id: hal-03997451**

**<https://hal.science/hal-03997451v1>**

Submitted on 20 Feb 2023

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# Assessment of Large Critical Electric Field in Ultra-wide Bandgap *p*-type Spinel $\text{ZnGa}_2\text{O}_4$

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## Abstract:

The spinel Zinc gallate  $\text{ZnGa}_2\text{O}_4$  stands out among the emerging ultra-wide bandgap ( $\sim 5$  eV) semiconductors as the ternary complex oxide with the widest gap where bipolar conductivity has been demonstrated. For power and energy electronics, a fundamental property of the material is its critical electric field ( $E_{CR}$ ) although, for  $\text{ZnGa}_2\text{O}_4$ , is yet unknown. In this work, highly resistive *p*-type  $\text{ZnGa}_2\text{O}_4$  thin films on sapphire and Si substrates were grown by metal organic chemical vapor deposition (MOCVD) to determine both, the remote acceptor concentration and vertical breakdown voltage. Hall Effect measurements confirmed a low carrier concentration at room temperature of  $\sim 10^{11} \text{ cm}^{-3}$ . From vertical metal-semiconductor-metal structures the average  $E_{CR}$  has been estimated to be of at least 5.3 MV/cm, which already is significantly larger than the one of SiC and GaN.

Keywords: ultra-wide band gap,  $\text{ZnGa}_2\text{O}_4$ , electrical properties, critical electric field

## 1. Introduction

Ultra-wide bandgap (UWBG) semiconductor oxides have recently attracted significant attention as a novel platform for ultra-high power electronics and deep ultraviolet opto-electronics [1]. The spinel zinc gallate ( $\text{ZnGa}_2\text{O}_4$ ) stands out within the family of oxide UWBG ( $E_g \sim 5$  eV) as an advanced semiconductor platform for applications such as photonics (i.e. photodetector [2], phototransistors [3–5]), gas sensors [6], thin-film-transistors [7–9], or optical memory application [10], thus, heralding novel spinel electronics. Spinel oxides have, indeed, a peculiar chemical formula  $\text{AB}_2\text{O}_4$ , where A and B cations occupy octahedral ( $O_h$ ) and tetrahedral ( $T_d$ ) sites formed by oxygen atoms. Thus, spinel oxides obey different doping rules from traditional binary semiconductor oxides because of their multivalent elements and different sites, such as the antisite defects due to the inverted distribution of cations [11]. The disordered distribution of ions and vacancies in the material, (caused by the so-called self-doping effect), can lead to intrinsic electron or hole conductivity [12].

Consequently, a wide range of intrinsic conductivities have been reported for  $\text{ZnGa}_2\text{O}_4$ . The resistivity of native  $n$ - $\text{ZnGa}_2\text{O}_4$  can be as low as  $2.7 \times 10^{-2} \Omega \cdot \text{cm}$  at room temperature (e.g., [7,13]), while intrinsic  $p$ - $\text{ZnGa}_2\text{O}_4$  is, in general, rather insulating at room temperature. Indeed, Chikoidze *et al.* [14] reported native  $p$ -type  $\text{ZnGa}_2\text{O}_4$  thin films grown by metal organic chemical vapor deposition (MOCVD), where the intrinsic hole conductivity at high temperature was attributed to the antisite defects  $\text{Zn}_{\text{Ga}}$ . Therefore, the intrinsic conductivity of  $\text{ZnGa}_2\text{O}_4$  thin films can be controlled well over 10 orders of magnitude by tuning the growth conditions (cation ratio, oxygen pressure, etc.) thanks to this self-doping effect [15]. In agreement with these empirical observations, Sabino *et al.* [16] suggested by ab-initio simulations that both,  $n$ - and  $p$ -type conductivities could be achieved in  $\text{ZnGa}_2\text{O}_4$ , and the hole conductivity is likely achieved by a polaronic band due to excess of Zn. It should be noted, however, that the electronic and optoelectronic devices based on  $\text{ZnGa}_2\text{O}_4$ , are mostly unipolar (i.e.,  $n$ -type) (e.g. [17]) and studies on  $p$ -type  $\text{ZnGa}_2\text{O}_4$  devices are still very limited. However, to exploit the full potential of  $\text{ZnGa}_2\text{O}_4$  spinel electronics, bipolar devices (such as  $p$ - $n$  junction (photo) diodes, PiN rectifiers, lateral and vertical power MOSFETs) and  $p$ -type conductivity should be achieved.

The critical electric field ( $E_{CR}$ ) of a material is a crucial parameter for assessing a semiconductor material's potential as a platform for power and energy electronics applications. In particular, the Baliga's figure of merit [18] is informative of the on-state/blocking device performance for power switching applications which primarily depends on carrier mobility and the cube of the critical electric field. Representative values of the critical electric field for the most common power electronics semiconductors, such as silicon (0.3 MV/cm), 4H-SiC (3.5 MV/cm) [19], GaN (3.9 MV/cm) [20], or AlGaN (5.9 MV/cm [21]) are well known. For other UWBG such as AlN, the  $E_{CR}$  values have also been reported (12 MV/cm [22]) while a generally agreed theoretical value for diamond is 10 MV/cm [23]. For  $\beta$ - $\text{Ga}_2\text{O}_3$ , a common estimation for the upper limit given in the literature is  $\sim 8$  MV/cm [24], while this value may be well underestimated as it has recently been experimentally observed to be at least 8.32 MV/cm and 13.2 MV/cm for  $p$ -NiO/ $\beta$ - $\text{Ga}_2\text{O}_3$  junctions [25] and Zn-doped  $\beta$ - $\text{Ga}_2\text{O}_3$  [26], respectively). However, to the best of our knowledge, critical electric field values for  $\text{ZnGa}_2\text{O}_4$  are not yet known.

Therefore, the objective of the present work was to define ZnGa<sub>2</sub>O<sub>4</sub> thin films onto insulating and conducting substrates for assessing the remote concentration and to establish a preliminary value of the ZnGa<sub>2</sub>O<sub>4</sub> critical electric field, respectively. According to the kinetic impact ionization theory and previous studies on Ga<sub>2</sub>O<sub>3</sub> [26], to achieve large critical electric fields, it is mandatory to obtain a sample with low remote carrier concentration to reduce the onset of the impact ionization process. In the present work, we studied ZnGa<sub>2</sub>O<sub>4</sub> thin films, which are highly resistive *p*-type in nature due to their deep acceptor native defects as we will show further on.

## 2. Experimental details

ZnGa<sub>2</sub>O<sub>4</sub> thin film were grown in a radio-frequency (RF) heated horizontal MOCVD reactor with separate inlets to avoid premature reactions in the manifold between oxygen and organometallics precursors. Trimethylgallium (TMGa), diethylzinc (DEZn) and, 5.5 N pure oxygen were used as gallium, zinc, and oxygen sources, respectively. The TMGa and DEZn bubbler temperatures were fixed at -10°C and 0°C, respectively. During the growth, the flow rates of TMGa, DEZn, and oxygen were kept at 11 μmol/min, 7 μmol/min, and 1200 sccm, respectively. The growth temperature and the reactor pressure were set at 775°C, and 38 torr, respectively. The growth time was 90 minutes. Argon was used as carrier gas. To analyse the ZnGa<sub>2</sub>O<sub>4</sub> crystal phases, X-Ray diffraction profiles were recorded in  $\theta/2\theta$  configuration between 15° and 115°, using a Rigaku SmartLab apparatus equipped with a Cu-K $\alpha$  source ( $\lambda = 0.1541$  nm). The thicknesses of the thin films were determined by means of scanning electron microscopy (SEM) cross-sectional images from a JEOL JSM 7001F scanning electron microscope.

For Hall effect measurements, Ohmic contacts were prepared with silver paint at the four corners of the sample. The resistivity and Hall effect measurements were performed, in a Van der Pauw configuration with perpendicularly to current applied magnetic fields varying from 0 to 1.6 T, in 550 to 850 K temperature range. The majority charge carrier sign, concentration, and mobility were then determined. A Keysight B1505A Power Device Analyzer has been used to carry out the vertical breakdown voltage measurements at room temperature. Negative bias has been applied on an isolated Ohmic contact with the grounded substrate to access the breakdown voltage.

## 3. Results and discussion

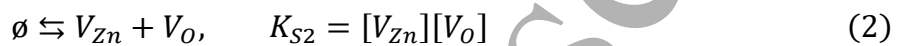
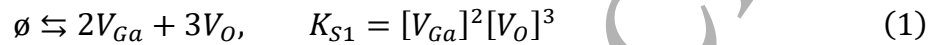
### 3.1. Thermodynamic Analysis of Point Defects and Free Carriers

We conducted a thermodynamic analysis of point defects and free carriers for equilibrium condition to establish the optimal growth parameters of the ZnGa<sub>2</sub>O<sub>4</sub> close to the stoichiometry composition. This thermodynamic analysis is carried out for the ZnGa<sub>2</sub>O<sub>4</sub> (crystal) - O<sub>2</sub> (gas) system. It is a three-component two-phase system. Thus, according to Gibbs' phase rule, its degree of freedom is 3, so three thermodynamic parameters can vary freely. All material

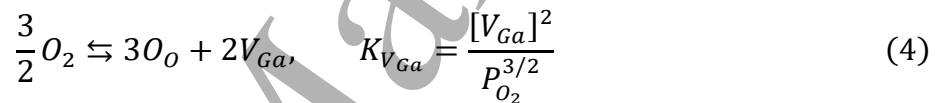
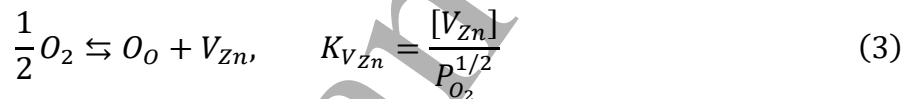
properties, including defects, free carriers' concentrations, and conductivity, depend on them. We chose the following independent variables: temperature  $T$ , oxygen partial pressure  $P_{O_2}$ , and the ratio of gallium and zinc partial pressures  $\chi = \frac{P_{Ga}}{P_{Zn}}$ .

To find the dependence of concentrations of defects and free carriers on these parameters we used the Kroger method of quasi-chemical reactions [27]. Special attention is paid to discern if the concentration of antisite defects  $Ga_{Zn}$  and  $Zn_{Ga}$  can be controlled in such a way to impact  $n$ - and  $p$ -type conductivities. The processes we considered with the corresponding mass action law are the following:

The creation of a Schottky defect pairs:



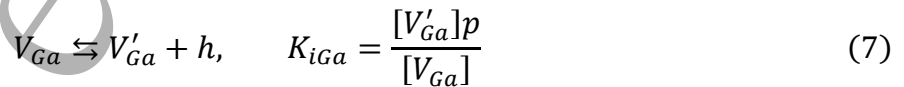
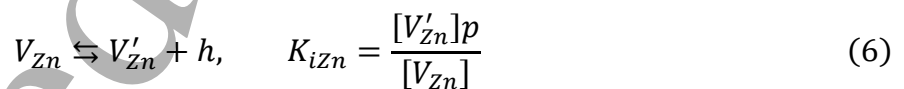
Where  $V_{Zn}$ ,  $V_{Ga}$ , and  $V_O$  are zinc, gallium, and oxygen vacancies, respectively;  $[V_{Zn}]$ ,  $[V_{Ga}]$ , and  $[V_O]$  are their concentrations, respectively. The incorporation of oxygen atoms in a crystal from the surrounding atmosphere, when zinc and gallium vacancies are created:



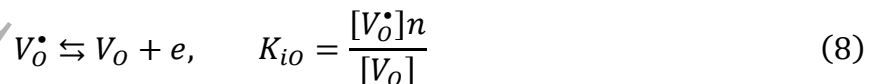
The lattice thermal ionization:



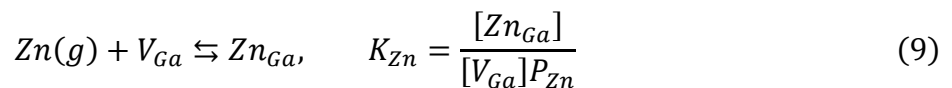
Where  $e$  and  $h$  are electrons and holes, respectively;  $n$  and  $p$  are their concentrations, respectively. The ionization of zinc, gallium and oxygen vacancies:

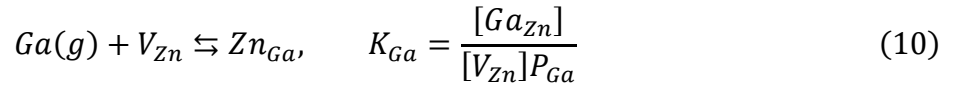


Where  $V'_{Zn}$  and  $V'_{Ga}$  denote positively single-charged zinc and gallium vacancies, respectively;  $[V'_{Zn}]$  and  $[V'_{Ga}]$  are their concentrations, respectively. The ionization of oxygen vacancy:

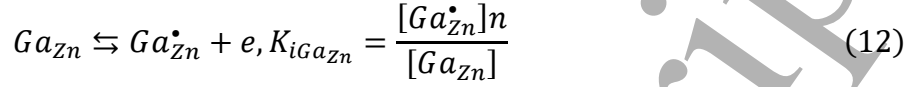
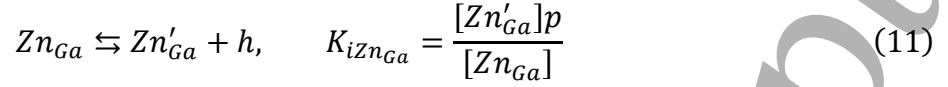


Where  $V_O^\bullet$  denotes negatively single-charged oxygen vacancy;  $[V_O^\bullet]$  is its concentration. Creation of antisite defects via incorporation of metal atoms from the gas phase into vacancies:





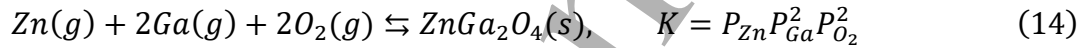
Where  $Zn_{Ga}$  and  $Ga_{Zn}$  are zinc in gallium site and gallium in zinc site anti-site defects, respectively;  $[Zn_{Ga}]$  and  $[Ga_{Zn}]$  are their concentrations, respectively. And their ionizations:



Where  $Zn'_{Ga}$  and  $Ga^*_{Zn}$  are positively and negatively single-charged anti-site defects, respectively;  $[Zn'_{Ga}]$  and  $[Ga^*_{Zn}]$  their concentrations, respectively. To find the concentration of defects and free carriers by the system of equations (1) - (12), the charge neutrality condition must be added:

$$n + [V'_{Zn}] + [V'_{Ga}] + [Zn'_{Ga}] \rightleftharpoons p + [V^*_O] + [Ga^*_{Zn}] \quad (13)$$

The connection between zinc, gallium, and oxygen partial pressures are defined by means of the constant of the zinc gallate creation reaction:



Having solved the system of equations (1) - (14), one can find the concentrations of defects and free carriers versus temperature, partial pressure of one component and ratio of gallium and zinc pressures. The calculated concentrations of charged defects and free carriers in dependence of oxygen partial pressure are shown in Figure 1. The temperature was chosen to be 775 °C, the ratio  $x = \frac{P_{Ga}}{P_{Zn}}$  equals to 3. These parameters correspond to our actual growth conditions. The  $x = P_{Ga}/P_{Zn}$  ratio is roughly estimated on the base of TMGa and DEZn flow rates and their formation enthalpies:

$$x = \frac{P_{Ga}}{P_{Zn}} = \frac{J_{Ga}\sqrt{2m_{Ga}kT} \exp(-\Delta H_{Ga}/kT)}{J_{Zn}\sqrt{2m_{Zn}kT} \exp(-\Delta H_{Zn}/kT)}$$

Here  $J_{Ga}$  and  $J_{Zn}$  are flow rates of TMGa and DEZn, respectively,  $m_{Ga}$  and  $m_{Zn}$  masses of Ga and Zn atoms,  $\Delta H_{Ga}$  and  $\Delta H_{Zn}$  formation enthalpies of TMGa and DEZn.

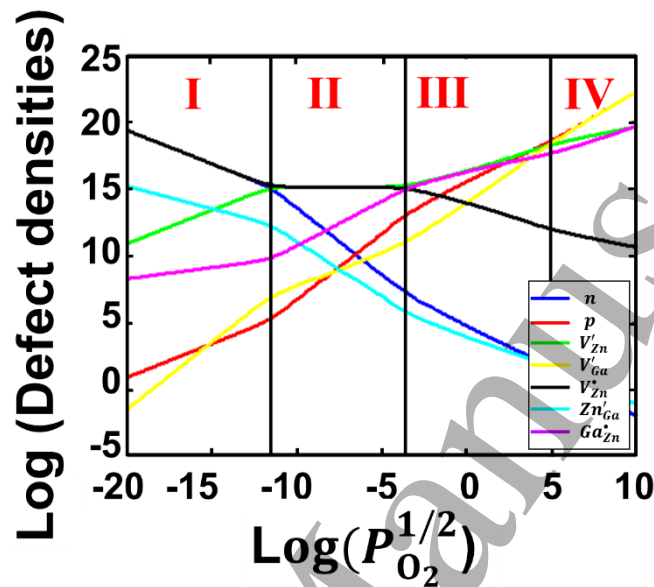
One can distinguish the following regions in the figure that:

- I. The region of intrinsic  $n$ -conductivity, origin of which are oxygen vacancies  $V_O$ . In this region  $[V^*_O] = n$ , and they are dominant;
- II. The first region of self-compensation, here oxygen vacancies  $[V^*_O]$  and zinc vacancies  $[V'_{Zn}]$  are dominant species and  $[V^*_O] = [V'_{Zn}]$ ;
- III. The second region of self-compensation where zinc vacancies  $[V'_{Zn}]$  and gallium in zinc sites  $[Ga^*_{Zn}]$  are dominant, and  $[V'_{Zn}] = [Ga^*_{Zn}]$ . In this region, the concentration of holes exceeds the concentration of electrons significantly. Around  $P_{O_2} = 1 \text{ atm}$ , the concentration of holes achieves  $10^{15} \text{ cm}^{-3}$ . It is remarkable that the concentration of holes at a whole range of pressure tails negatively charged gallium vacancies  $[V'_{Ga}]$  - the concentrations of  $V'_{Ga}$  and holes change

the same way when oxygen partial pressure changes and difference between these concentrations is quite small along the whole pressure range.

- IV. The region of hole conductivity, here  $[V'_{Ga}]$  and holes are dominant species. However, the oxygen pressure is unrealistically high here.

Notably, there is a very weak dependence of concentrations of free carriers on  $x$ . For  $x = 10^{-3}$ , that is, when the partial pressure of Zn exceeds the partial pressure of Ga three orders of magnitude, the concentration of holes increases only ten times, compared to  $x = 3$ .



**Figure 1.** The dependence of concentration of point defects and free carriers on oxygen partial pressure, for  $T = 775$  °C, and  $x = \frac{P_{Ga}}{P_{Zn}} = 3$ .

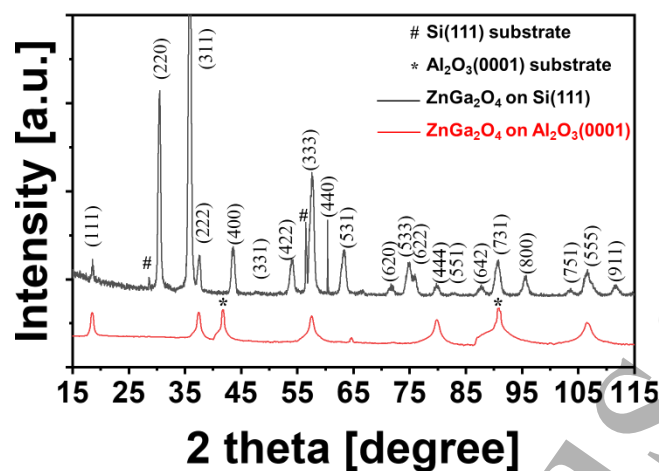
Therefore, the growth of both  $n$ - and  $p$ -type  $ZnGa_2O_4$  thin films are feasible by tuning the growth parameters such as, mainly oxygen partial pressure and temperature. The cations concentration ratio weakly affects the conductivity (as experimentally observed [14]). In this work, we concentrate on the regions predicted to correspond to  $p$ -type conductivity. Then, the highly-resistive  $p$ -type  $ZnGa_2O_4$  thin layer was grown in the conditions which are input from thermodynamic analysis.

### 3.2. $ZnGa_2O_4$ Thin Film Growth and Structural Properties

$ZnGa_2O_4$  thin films were grown on two kinds of substrates simultaneously: on conducting  $p$ -type Si(111) and  $c$ -oriented sapphire  $Al_2O_3(0001)$  (insulating) substrates.  $ZnGa_2O_4/Si$  sample was used to fabricate vertical parallel plane capacitors, and characterized vertically for breakdown voltage capability. While the  $ZnGa_2O_4/sapphire$  thin film was used for the investigation of the electrical transport properties (resistivity, carrier, mobility) at elevated temperatures in a Van Der Pauw configuration.

Figure 2 shows the Cu-K $\alpha$  X-ray diffraction (XRD) spectra of  $ZnGa_2O_4$  thin films grown on Si(111) and  $Al_2O_3(0001)$ . According to the PDF 28-1240 file, it is pure  $ZnGa_2O_4$  spinel without a secondary phase. Film on Si(111) is polycrystalline without any preferential texturization, similar to

XRD pattern reported by Galazka *et al.* [28] for their melt-grown bulk  $\text{ZnGa}_2\text{O}_4$  spinel, milled in a randomly oriented powder. For films on  $\text{Al}_2\text{O}_3(0001)$  substrates, Bragg peaks appearing at  $18.6^\circ$ ,  $37.6^\circ$ ,  $57.6^\circ$ ,  $79.7^\circ$ , and  $106.6^\circ$  correspond to the (111), (222), (333), (444), and (555) planes of the cubic spinel  $\text{ZnGa}_2\text{O}_4$ , respectively, as reported by previous literature [14,29].



**Figure 2.** *Cu-K $\alpha$*  XRD spectra of  $\text{ZnGa}_2\text{O}_4$  thin film grown on  $\text{Si}(111)$  (black), and  $\text{Al}_2\text{O}_3(0001)$  (red). \* and # correspond respectively to  $\text{Al}_2\text{O}_3(0001)$  and  $\text{Si}(111)$  Bragg reflections.

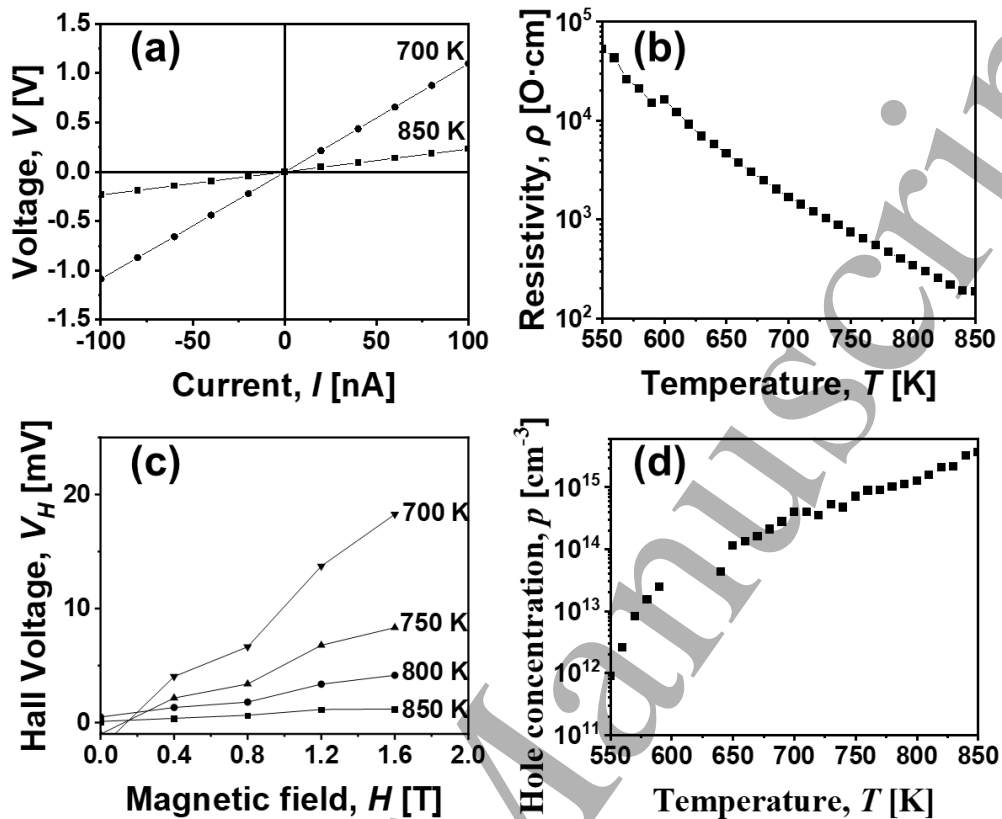
$\text{ZnGa}_2\text{O}_4$  thin films have already been grown polycrystalline on sapphire substrates by MOCVD [30], radio-frequency (RF) magnetron sputtering [31,32], pulsed laser deposition (PLD) [13], with a preferential orientation along [111] of cubic spinel and equivalent direction for  $\text{Al}_2\text{O}_3(0001)$ . Analogously,  $\text{ZnGa}_2\text{O}_4$  thin films have also been grown on Si substrates by atomic layer deposition (ALD) [29], PLD [33], and RF magnetron sputtering [32]. A larger lattice mismatch between  $\text{ZnGa}_2\text{O}_4$  and  $\text{Si}(100)$  (which may also contain a native ultra-thin  $\text{SiO}_2$  interlayer) results in a relatively weaker orientation of crystallite in the  $\text{ZnGa}_2\text{O}_4$  grown on Si substrate than that on sapphire (0001) [32]. The full width at half maximum (FWHM) values were measured as  $0.29^\circ$  and  $0.49^\circ$  for  $\text{ZnGa}_2\text{O}_4$  (111) on sapphire and  $\text{Si}(100)$  substrates, respectively, indicating also a smaller size of crystallites in the case of  $\text{Si}(111)$  substrates with respect to  $\text{Al}_2\text{O}_3(0001)$  ones.

### 3.3. Electrical Transport Properties

Hall and electrical transport properties were studied for  $\text{ZnGa}_2\text{O}_4$  grown on  $\text{Al}_2\text{O}_3$  crystals, as the sapphire substrate is insulating up to high temperatures thus preventing parasitic leakage through the substrate. Figure 3-a shows a typical Ohmic current–voltage ( $I$ – $V$ ) behavior at different temperatures. The ohmic contact was made with silver paste painted at four corners of the sample in a square shape. The resistivity versus temperature for the spinel  $\text{ZnGa}_2\text{O}_4$  sample is shown in Figure 3-b. At 850 K, the resistivity was measured as  $1.9 \times 10^2 \Omega \cdot \text{cm}$ , and increased to  $5.3 \times 10^4 \Omega \cdot \text{cm}$  at 550 K. Hall voltage ( $V_H$ ) versus magnetic field (0 – 1.6 T) at 700 – 850 K for a  $\text{ZnGa}_2\text{O}_4$  sample is shown in Figure 3-c, the positive  $V_H$  increases with a perpendicularly applied magnetic field in a Van der Pauw configuration, which evidenced that our spinel  $\text{ZnGa}_2\text{O}_4$  thin film has  $p$ -type conductivity at elevated temperature. The temperature dependence for Hall hole concentration is shown in Figure 3-d for  $p$ -type  $\text{ZnGa}_2\text{O}_4$  film, the hole concentration was measured between  $8.8 \times 10^{11}$  and  $3.7 \times 10^{15} \text{ cm}^{-3}$ .



<sup>3</sup> at 550 – 850 K. Although the hopping behavior starts to pronounce below 600 K, extrapolating the band-activated regime of the hole concentration versus temperature at the higher temperatures, the hole density is below  $10^{11} \text{ cm}^{-3}$  at 300 K, confirming the low level of  $p$ -type doping with the applied growth conditions. The mobility of the  $p$ -type  $\text{ZnGa}_2\text{O}_4$  thin film varies between 9 and  $15 \text{ cm}^2/(\text{V}\cdot\text{s})$ .



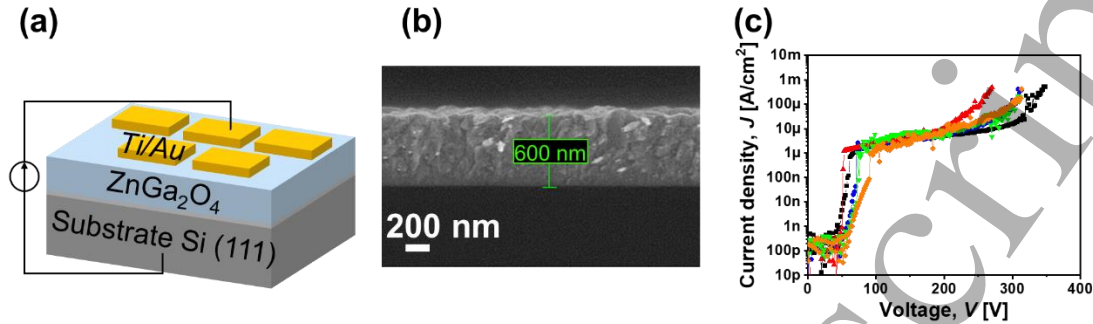
**Figure 3.** (a) Ohmic  $I$ – $V$  characteristics at 700K and 850 K (b) Resistivity versus temperature. (c) Hall voltage ( $V_H$ ) versus magnetic field at different temperatures (700 – 850 K). The Hall voltage is positive and increase with the increasing applied magnetic field, indicating the  $p$ -type conductivity in  $\text{ZnGa}_2\text{O}_4$  layer. (d) Temperature dependence for Hall free hole concentration for  $\text{ZnGa}_2\text{O}_4$  thin film.

### 3.4. $\text{ZnGa}_2\text{O}_4$ Critical Electric Field

To determine the vertical critical electric field experimentally, vertical parallel plane capacitors has been fabricated by depositing Ti/Au contacts (with a size of  $170 \mu\text{m} \times 300 \mu\text{m}$ ) onto  $p$ -type  $\text{ZnGa}_2\text{O}_4$  grown onto the conducting Si thin films (Figure 4-a). Figure 4-b shows the scanning electron microscopy (SEM) cross-section image revealing that the thickness of the  $\text{ZnGa}_2\text{O}_4/\text{Si}$  is 600 nm.

The vertical hard breakdown voltage  $V_{BR}$  was in the range of 270 to 350 V (Figure 4-c). It can be pointed out that a low leakage current below  $1 \text{ mA}/\text{mm}^2$  is observed all the way to the breakdown voltage. A thin  $\text{SiO}_2$  layer (2 nm) at the interface  $\text{ZnGa}_2\text{O}_4/\text{Si}$  can often be seen [32] after the growth. To exclude the influence of the  $\text{SiO}_2$  layer, a bare Si substrate with a native  $\text{SiO}_2$  (2 – 3 nm thick layer) has been characterized as well. The vertical breakdown voltage of around 10 V [26] is considered to produce a negligible contribution to the experimentally measured  $V_{BR}$  of  $p$ -type  $\text{ZnGa}_2\text{O}_4/\text{Si}$  structure.

Taking the experimentally determined 600 nm thickness for the  $\text{ZnGa}_2\text{O}_4$  layer and extracting bare Si substrate related  $V_{BR}$  into account, the average vertical critical electric field across the film can be as high as  $E_{CR} = 5.3$  MV/cm, which is already significantly larger than the one from SiC and (single crystal) GaN. Therefore, the blank of breakdown field for  $\text{ZnGa}_2\text{O}_4$  in the literature for assessing WBG semiconductors [34] can be filled as shown in **Table I**.



**Figure 4.** (a) Sketch of the vertical heterojunction structure with Ti/Au metal contacts. (b) SEM cross-section image of  $\text{ZnGa}_2\text{O}_4$  polycrystalline film deposited on Si(111). The estimated thickness of the film is 600 nm. (c) Experimental  $I$ - $V$  curves of vertical breakdown voltage measurements through the  $\text{ZnGa}_2\text{O}_4/\text{Si}$  structure, indicating a hard breakdown voltage up to 350 V at room temperature.

**Table I.** Comparison of various wide-bandgap semiconductor materials properties [26,34,35]. \*: experimental  $E_{CR}$  values measured by authors. <sup>a</sup>Common value from literature. <sup>b</sup> $p$ -NiO/ $n$ -Ga<sub>2</sub>O<sub>3</sub> [25] <sup>c</sup> $p$ -Ga<sub>2</sub>O<sub>3</sub> [22]

	4H-SiC	GaN	$\beta$ -Ga <sub>2</sub> O <sub>3</sub>	ZnGa <sub>2</sub> O <sub>4</sub>
<b>Crystal Structure</b>	Hexagonal (Wurtzite)	Hexagonal (Wurtzite)	Monoclinic	Cubic (Spinel)
<b>Space Group</b>	P63mc	P63mc	C2/m	Fd-3m
<b>Direct bandgap, <math>E_g</math> (eV)</b>	3.2	3.4	4.8	4.8
<b>Dielectric constant, <math>\epsilon</math></b>	9.7	9	10 [36]	10.4 [37]
<b>Critical electric field, <math>E_{CR}</math> (MV/cm)</b>	3.5	3.9	8 <sup>a</sup> [24], 8.32 <sup>b</sup> , 5.7 <sup>c</sup> , 13.2 <sup>c</sup> [26]	<b>&gt; 5.3*</b>
<b>Electron mobility, <math>\mu</math> (cm<sup>2</sup>/(V·s))</b>	1000	1250	300	100
<b>Hole mobility, <math>\mu</math> (cm<sup>2</sup>/(V·s))</b>	106 (300 K) [38]	34 (300 K) [39]	~10 (high T) [26]	~10 (high T)
<b>Saturation velocity, <math>v_s</math> (10<sup>7</sup> cm/s)</b>	2	2.5	1.8 – 2	–
<b>Thermal conductivity, <math>\lambda</math> (W/cm·K)</b>	4.9	2.3	0.1 – 0.3	0.22

#### 4. Conclusion

In this work, following a thermodynamic analysis, highly resistive undoped polycrystalline  $\text{ZnGa}_2\text{O}_4$  thin film of 600 nm were grown on  $p$ -type silicon substrate by MOCVD. The conductivity and hole concentration of the film at 550 K were determined to be as low as  $5 \times 10^4 \Omega \cdot \text{cm}$  and  $9 \times 10^{11} \text{cm}^{-3}$ , respectively. The average breakdown voltage of  $\text{ZnGa}_2\text{O}_4/\text{Si}$  vertical structures was  $V_{BR} = 350$

V, resulting in a critical electric field  $E_{CR}$  of 5.3 MV/cm. This value is already near twice the breakdown electrical fields of SiC and GaN and is in line with previous demonstration of its binary counterpart, Ga<sub>2</sub>O<sub>3</sub>. We believe that this first experimental estimation of  $E_{CR}$  for polycrystalline ZnGa<sub>2</sub>O<sub>4</sub> will trigger further studies related to nanocrystal and high quality and homo-epitaxial films growth for future possible integration of spinel ZnGa<sub>2</sub>O<sub>4</sub> and doped relatives for electronic and optoelectronic devices.

## Acknowledgement

The present work is a part of “GALLIA” International Research Project, CNRS, France. GEMaC colleagues acknowledge financial support of French National Agency of Research (ANR), project “GOPOWER”, CE-50 N0015-01. The ICN2 is funded by the CERCA programme / Generalitat de Catalunya and by the Severo Ochoa programme of the Spanish Ministry of Economy, Industry and Competitiveness (MINECO, grant no. SEV-2017-0706).

## Conflicts of interest

There are no conflicts to declare.

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