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LOCAL CONTROLLABILITY OF THE KORTEWEG-DE VRIES EQUATION WITH THE RIGHT DIRICHLET CONTROL

HOAI-MINH NGUYEN

ABSTRACT. The Korteweg-de Vries (KdV) equation with the right Dirichlet control is small time, locally, exactly controllable for all non-critical lengths and its linearized system is not controllable for *all* critical lengths. In this paper, we give a definitive picture of the local controllability properties of this control problem for all critical lengths. In particular, we show that the unreachable space of the linearized system is always of dimension 1 and the KdV system with the right Dirichlet control is not locally null controllable in small time for any critical length. We also give a criterion to determine whether the system is locally exactly controllable in finite time or *not* locally null controllable in any positive time for *all* critical lengths. Consequently, we show that there exist critical lengths such that the system is *not* locally null controllable in small time but is locally exactly controllable in finite time.

Key words: Controllability, KdV equations, critical lengths, unreachable space, power series expansion, Hilbert uniqueness method

AMS subject classification: 35Q53, 93B05, 93B07, 93C20, 35C20.

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1. INTRODUCTION AND THE STATEMENT OF THE MAIN RESULTS

This paper is devoted to the local controllability of the Korteweg-de Vries (KdV) equation using the right Dirichlet control. More precisely, we consider the following control problem, for $T > 0$,

$$(1.1) \quad \begin{cases} y_t + y_x + y_{xxx} + yy_x = 0 & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, T), \\ y(\cdot, L) = u & \text{in } (0, T), \\ y(0, \cdot) = y_0 & \text{in } (0, L). \end{cases}$$

Here y is the state, $y_0 \in L^2(0, L)$ is an initial datum, and u is a control, belonging to an appropriate functional space. The KdV equation has been introduced by Boussinesq [14] and Korteweg and de Vries [31] as a model for the propagation of surface water waves along a channel. This equation also furnishes a very useful nonlinear approximation model including a balance between weak nonlinearity and weak dispersive effects, see e.g. [50, 37, 29]. The KdV equation has been investigated from various aspects of mathematics, including the well-posedness, the existence, and stability of solitary waves, the integrability, the long-time behavior, etc., see e.g. [50, 37, 29, 48, 33].

1.1. State of the art. The local controllability for the KdV equation has been studied extensively in the literature, see, e.g., the surveys [44, 17] and the references therein. We briefly review here some results concerning boundary controls. When the controls are $y(\cdot, 0)$, $y(\cdot, L)$, $y_x(\cdot, L)$, Russell and Zhang [46] proved that the KdV equation is small time, locally exactly controllable. The case of left boundary control ($y(\cdot, L) = y_x(\cdot, L) = 0$) was investigated by Rosier [43] (see also [26]). The small-time local, null controllability holds in this case. The exact controllability does not hold for initial and final data in the $L^2(0, L)$ due to the regularisation effect but holds for a subclass of infinitely smooth initial and final data [35].

A very close setting to the one considered here is the setting in which one controls the right Neumann boundary, i.e., $y(\cdot, 0) = y(\cdot, L) = 0$ and $y_x(\cdot, L)$ is a control. For initial and final data in $L^2(0, L)$, and controls in $L^2(0, T)$, Rosier [42] proved that the KdV system with the right Neumann control is small time, locally, exactly controllable provided that the length L is not critical, i.e., $L \notin \mathcal{N}_N$ ¹, where²

$$(1.2) \quad \mathcal{N}_N := \left\{ 2\pi \sqrt{\frac{k^2 + kl + l^2}{3}}; k, l \in \mathbb{N}_* \right\}.$$

To this end, Rosier studied the controllability of the corresponding linearized system and showed that the linearized system is exactly controllable if $L \notin \mathcal{N}_N$. He as well established that when $L \in \mathcal{N}_N$, the linearized system is not controllable. More precisely, Rosier showed that there exists a non trivial, finite dimensional subspace \mathcal{M}_N of $L^2(0, L)$ such that its orthogonal space is reachable from 0 for small time whereas \mathcal{M}_N is not for any time. To tackle the control problem for a critical length $L \in \mathcal{N}_N$ with initial and final data in $L^2(0, L)$ and controls in $L^2(0, T)$, Coron and Crépeau introduced the power series expansion method [23]. The idea is to take into account the effect of the nonlinear term yy_x absent in the corresponding linearized system. Using this method, Coron

¹The letter N stands for the Neumann boundary control.

² $\mathbb{N}_* = \mathbb{N} \setminus \{0\}$.

and Crépeau showed [23] (see also [22, section 8.2]) that the KdV system is small time, locally, exactly controllable when $\dim \mathcal{M}_N = 1$. Cerpa [16] developed the analysis in [23] to prove that the KdV system is *finite time*, locally, exactly controllable in the case $\dim \mathcal{M}_N = 2$. Later, Crépeau and Cerpa [18] succeeded to extend the ideas in [16] to obtain the local, exact controllability in *finite time* for all other critical lengths. Recently, with Coron and Koenig [24], we prove that such a system is not small time, locally, null controllable for a class of critical lengths. This fact is surprising when compared with known results on internal controls for the KdV equation. It is known, see [15, 36, 41], that the KdV system (1.1) with $u = 0$ is small time, locally controllable using internal controls *whenever* the control region contains an *arbitrary*, open subset of $(0, L)$. The power expansion method is also a starting point of the analysis in [24]. Part of the analysis is to characterize all controls that bring 0 at 0 to 0 at time T for the corresponding linearized system. This idea is then used in the study of the water tank problem [25]. It is interesting to note that there are other types of boundary controls for the KdV equation for which there is no critical length, see [42, 43, 27, 17]. There are also results on the internal controllability for the KdV equation, see [46, 32, 15] and the references therein. A minimal time of the null controllability is also required for some linear partial differential equations. This is the case for equations with a finite speed of propagation, such as the transport equation, the wave equation, or the hyperbolic system, see, e.g., [22] and the references therein. But this can also happen for equations with the infinite speed of propagation, such as some parabolic systems [9], Grushin-type equations [2, 7, 30, 3], Kolmogorov-type equations [4], and the references therein.

We now turn back to the control problem (1.1). This control problem was first investigated by Glass and Guerrero [27]. To this end, in the spirit of Rosier's work mentioned above, they introduced the corresponding set of critical lengths ³

$$(1.3) \quad \mathcal{N}_D = \left\{ L \in \mathbb{R}_+; \exists z_1, z_2 \in \mathbb{C} : (1.4) \text{ holds} \right\},$$

where

$$(1.4) \quad z_1 e^{z_1} = z_2 e^{z_2} = -(z_1 + z_2) e^{-(z_1 + z_2)} \quad \text{and} \quad L^2 = -(z_1^2 + z_1 z_2 + z_2^2).$$

They proved that the set \mathcal{N}_D is infinite and has no accumulation point. Concerning (1.1), Glass and Guerrero proved that the corresponding linearized KdV system is small time, exactly controllable with initial and final data in $H^{-1}(0, L)$ using controls in $L^2(0, T)$ if $L \notin \mathcal{N}_D$. Developing this result, they also established that the KdV system (1.1) is small-time locally controllable for initial and final data in $L^2(0, L)$ and controls in $H^{1/6-4}$ for *non-critical* lengths, i.e., $L \notin \mathcal{N}_D$.

The KdV equation describes the waves moving to the right. Consequently, the control problem using boundary controls on the right is more interesting and more difficult than the one using boundary controls on the left in general. It is known that the KdV equation is locally controllable in small time using left Dirichlet controls for any length since the corresponding linearized system is always controllable; there is no critical length for the left Dirichlet controls [43, 26]. This is in contrast with the right Dirichlet controls considered in the paper.

Even the local controllability of the KdV system using boundary controls has been investigated extensively in the last two decades, to our knowledge, there is no result on the local controllability of system (1.1) for critical lengths. In comparison with the set of critical lengths \mathcal{N}_N , the set \mathcal{N}_D is less explicit. Moreover, the unreachable space for the linearized system related to (1.1) has not been determined due to the lack of an appropriate observability inequality for the linearized system in a suitable functional setting. As revealed later in this paper, the control properties related to the unreachable space for the right Dirichlet control are quite distinct from the one for the right Neumann control.

³The letter D stands for the Dirichlet boundary control.

⁴Controls in $H^{1/6-}$ means controls in $H^{1/6-\varepsilon}$ for all $\varepsilon > 0$.

1.2. Statement of the main results. The main goal of this paper is to give a rather complete picture of the local controllability of the control KdV system (1.1). The first main result of this paper is the characterization of all the critical lengths and the corresponding unreachable space \mathcal{M}_D for the corresponding linearized system. The first result in this direction is.

Theorem 1.1. *Let $L \in \mathcal{N}_D$. There exists a unique pair $(a, b) \in \mathbb{R}^2$ such that $a < 0$, $b > 0$,*

$$(1.5) \quad (a + ib)e^{a+ib} = (a - ib)e^{a-ib} = -2ae^{-2a},$$

and

$$(1.6) \quad L^2 = b^2 - 3a^2.$$

Set

$$(1.7) \quad q = -\frac{2a(a^2 + b^2)}{L^3} > 0,$$

$$(1.8) \quad \phi(x) = -\beta e^{\alpha x} \cos(\beta x) + \beta e^{-2\alpha x} + 3\alpha e^{\alpha x} \sin(\beta x) \text{ for } x \in [0, L],$$

with

$$(1.9) \quad \alpha = -a/L \quad \text{and} \quad \beta = -b/L,$$

and

$$(1.10) \quad \Phi(t, x) = e^{qt} \phi(x) \text{ for } (t, x) \in \mathbb{R} \times [0, L].$$

Then Φ is a solution of the system

$$(1.11) \quad \begin{cases} \Phi_t + \Phi_x + \Phi_{xxx} = 0 & \text{in } \mathbb{R} \times (0, L), \\ \Phi(\cdot, 0) = \Phi_x(\cdot, 0) = \Phi(\cdot, L) = \Phi_{xx}(\cdot, L) = 0 & \text{in } \mathbb{R}, \end{cases}$$

and the unreachable space \mathcal{M}_D of the linearized system of (1.1) is given by

$$(1.12) \quad \mathcal{M}_D = \text{span} \left\{ \phi \right\}.$$

Consequently,

$$(1.13) \quad \dim \mathcal{M}_D = 1.$$

Remark 1.1. The main part of Theorem 1.1 is to show that the complex numbers z_1, z_2 from the definition of \mathcal{N}_D in (1.3) can be chosen as

$$z_1 = a + ib \quad \text{and} \quad z_2 = a - ib,$$

for some $a < 0$ and $b > 0$. This is non-trivial, see Section 4.

Remark 1.2. Implicit information given in Theorem 1.1 is that the function ϕ , defined in (1.8), satisfies the following boundary condition:

$$\phi(0) = \phi_x(0) = \phi(L) = \phi_{xx}(L) = 0.$$

Using Theorem 1.1, we can precisely describe the set \mathcal{N}_D in the following result.

Proposition 1.1. *Let $L > 0$. Then $L \in \mathcal{N}_D$ if and only if $L = L_n$ for some $n \in \mathbb{N}$ where L_n is given by*

$$L_n^2 = 4a_n^2(e^{-6a_n} - 1),$$

with $a_n < 0$ being uniquely determined by

$$b_n \cos b_n + a_n \sin b_n = 0,$$

where b_n satisfies

$$\pi + 2n\pi < b_n < 3\pi/2 + 2n\pi \quad \text{and} \quad b_n^2 = 4a_n^2(e^{-6a_n} - 1/4).$$

After presenting the results on the unreachable space of the linearized system and the characterization of the critical lengths, we next discuss the local controllability of (1.1). Set, for $T > 0$,

$$(1.14) \quad X_T := C([0, T]; L^2(0, L)) \cap L^2((0, T); H^1(0, L))$$

equipped with the corresponding norm:

$$\|y\|_{X_T} = \|y\|_{C([0, T]; L^2(0, L))} + \|y\|_{L^2((0, T); H^1(0, L))}.$$

We later show that System (3.4) is well-posed in X_T if $\|y_0\|_{L^2(0, L)}$ and $\|u\|_{H^{1/3}(0, T)}$ are sufficiently small, see Proposition 3.4.

Before presenting our result on the local controllability properties of (1.1), we recall several definitions related to these properties.

Definition 1.1. Let $T > 0$, and let $H \subset L^2(0, L)$ and $H_c \subset H^{1/3}(0, T)$ be Hilbert spaces with the corresponding norm $\|\cdot\|_H$ and $\|\cdot\|_{H_c}$. System (1.1) is locally exactly controllable in time T for initial and final data in H and controls in H_c if for any $\eta > 0$ (small), there exists $\delta > 0$ (small) such that for all $y_0, y_T \in H$ with $\|y_0\|_H, \|y_T\|_H \leq \delta$, there exists $u \in H_c$ with $\|u\|_{H_c} \leq \eta$ such that $y(T, \cdot) = y_T$ where $y \in X_T$ is the unique solution of (1.1).

Definition 1.2. If for all $T > 0$, System (1.1) is locally exactly controllable in time T for initial and final data in H and controls in H_c , then System (1.1) is said to be small-time locally exactly controllable for initial and final data in H and controls in H_c .

Definition 1.3. If System (1.1) is locally exactly controllable in some time $T > 0$ for initial and final data in H and controls in H_c , then System (1.1) is said to be locally exactly controllable in finite time for initial and final data in H and controls in H_c .

Similar properties are associated with the local null-controllability of System (1.1) for which y_T is taken to be 0 in the previous definitions.

The second main result of this paper is that System (1.1) is not locally null controllable in small time for controls in $H^{1/2}(0, T)$ for all critical lengths.

Theorem 1.2. *Let $L \in \mathcal{N}_D$. There exist $T = T_0 > 0$ and $\varepsilon_0 > 0$ such that for all solutions $y \in X_T$ of the system*

$$(1.15) \quad \begin{cases} y_t + y_x + y_{xxx} + yy_x = 0 & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, T), \\ y(\cdot, L) = u & \text{in } (0, T), \\ y(0, \cdot) = \varepsilon\phi & \text{in } (0, L), \end{cases}$$

with $0 < \varepsilon < \varepsilon_0$ and $\|u\|_{H^{1/2}(0, T)} < \varepsilon_0$, we have

$$y(T, \cdot) \neq 0.$$

Recall that ϕ is defined in (1.8).

Remark 1.3. The function ϕ defined in (1.8) belongs to $C^\infty([0, L])$ and satisfies $\phi(0) = \phi_x(L) = 0$ (see Remark 1.2).

We next present a criterion on the local controllability property of (1.1) in finite time. Set

$$(1.16) \quad \Omega(z) = \int_0^L \left| \sum_{j=1}^3 (\lambda_j e^{\lambda_j L} - \lambda_{j+1} e^{\lambda_{j+1} L}) e^{\lambda_{j+2} x} \right|^2 \phi_x(x) dx,$$

where $\lambda_j = \lambda_j(z)$ with $j = 1, 2, 3$ are the three solutions of the equation

$$\lambda^3 + \lambda + iz = 0$$

and the convention $\lambda_{j+3} = \lambda_j$ for $j \geq 1$ is used. It is clear that Ω is continuous with respect to z . One can show that, see Remark 7.1,

$$\lim_{|z| \rightarrow \infty, z \in \mathbb{R}} \Omega(z + iq/2) = +\infty,$$

where q is defined by (1.7).

The third main result of this paper is the criterion on the local controllability property of (1.1) in finite time for initial and final data in $L^2(0, L)$ when $L \in \mathcal{N}_D$.

Theorem 1.3. *Let $L \in \mathcal{N}_D$ and let $q > 0$ be defined by (1.7). Set*

$$\omega = \min_{z \in \mathbb{R}} \Omega(z + iq/2),$$

where Ω is defined by (1.16). The following two facts hold.

i) If $\omega \geq 0$, then system (1.1) is not locally null controllable in any positive time with controls in $H^{1/2}$. More precisely, given $T > 0$ arbitrary, there exists $\varepsilon_T > 0$ such that for all solution of $y \in X_T$ of (1.15) with $0 < \varepsilon < \varepsilon_T$ and $\|u\|_{H^{1/2}(0, T)} < \varepsilon_T$, we have

$$y(T, \cdot) \neq 0.$$

ii) If $\omega < 0$, then system (1.1) is locally exactly controllable in finite time with controls in $H^{1/3}$, and initial and final data in $L^2(0, L)$. More precisely, there exist $T_0 > 0$ and $\varepsilon_0 > 0$ such that for all $y_0, y_1 \in L^2(0, L)$ with $\|y_0\|_{L^2(0, L)} \leq \varepsilon_0$ and $\|y_1\|_{L^2(0, L)} \leq \varepsilon_0$, there exists $u \in H^{1/3}(0, T_0)$ with $\|u\|_{H^{1/3}(0, T_0)} \leq C(\|y_0\|_{L^2(0, L)} + \|y_1\|_{L^2(0, L)})^{1/2}$ such that

$$y(T_0, \cdot) = y_1,$$

where $y \in X_{T_0}$ is the unique solution of (1.1). Here C denotes a positive constant independent of y_0 and y_1 .

Remark 1.4. Using Scilab, the program is given in Appendix D, one can show that

- if $n = 0$, then $L = 4.5183604$, $a = -0.5065520$, $b = 4.6027563$, $q = 0.2354919$.
- if $n = 1$, then $L = 10.866906$, $a = -0.6903700$, $b = 10.932497$, $q = 0.1291104$.
- if $n = 2$, then $L = 17.177525$, $a = -0.7947960$, $b = 17.232599$, $q = 0.0933315$, and $\Omega(iq/2 + 0.2668750) = -0.4687287$.
- if $n = 3$, then $L = 23.476776$, $a = -0.8687610$, $b = 23.524949$, $q = 0.0744156$, and $\Omega(iq/2 + 0.1419531) = -2.3917110$.

Parameters of the program (k, j, j_0) are given for $n = 2$. The parameters for $n = 0, 1, 3$ are also given there.

As a consequence of Theorem 1.2 and Theorem 1.3 (see also Remark 1.4), there exist critical lengths for which system (1.1) is not locally null controllable in small time but locally exactly controllable in finite time even the unreachable space of the corresponding linearized system is of dimension 1. This result is surprising and distinct when compared with known related results on the local controllability of the KdV equation. First, it is known that the corresponding KdV system with the right zero Dirichlet is locally exactly controllable in small time using internal controls. Second, the unreachable space of the linearized system of the corresponding KdV system with the right Neumann control might be of arbitrary dimension. Third, the KdV system with the right Neumann control is locally exactly controllable in small time if the corresponding unreachable space is of dimension 1.

Finally, in the case $L \notin \mathcal{N}_D$, we can prove the following result, which sharpens the results in [27] mentioned previously.

Theorem 1.4. *Let $L \notin \mathcal{N}_D$ and $T > 0$. There exist $C > 0$ and $\varepsilon_0 > 0$ depending only on L and T such that for all $y_0, y_1 \in L^2(0, L)$ with $\|y_0\|_{L^2(0, L)} \leq \varepsilon_0$ and $\|y_1\|_{L^2(0, L)} \leq \varepsilon_0$, there exists $u \in H^{1/3}(0, T)$ with $\|u\|_{H^{1/3}(0, T)} \leq C(\|y_0\|_{L^2(0, L)} + \|y_1\|_{L^2(0, L)})$ such that*

$$y(T, \cdot) = y_1,$$

where $y \in X_T$ is the unique solution of (1.1).

1.3. Ideas of the proofs. The characterization of the critical lengths and the unreachable space is based on the study of system (1.4). In comparison with the characterization of the set \mathcal{N}_N , the study of \mathcal{N}_D is more complex. The main goal and the difficulty in this direction, see Section 4, are to obtain/characterize *all* solutions of (1.4). The existence of a smaller class of solutions is easier as observed previously in [27] (see also Remark 4.2). As a consequence, we derive that $\dim \mathcal{M}_D = 1$ which is completely different from what is known for the right Neumann control.

Our approaches to Theorem 1.2 and Theorem 1.3 are inspired by the power series expansion method introduced by Coron and Crépeau [23]. This method has also been applied to other contexts such as local, bilinear control of Schrödinger's equations [5, 8, 13], decay of the solutions of the KdV equations [19, 47, 39], local controllability of Burgers' equations [34], local controllability of water tank [25], and the references therein.

The strategy is to use the nonlinear effect to move the solution in the direction prohibited by the linearized system. The idea of this method is to search for/understand a control u of the form

$$u = \varepsilon u_1 + \varepsilon^2 u_2 + \dots$$

The corresponding solution then formally has the form

$$y = \varepsilon y_1 + \varepsilon^2 y_2 + \dots,$$

and the non-linear term yy_x can be written as

$$yy_x = \varepsilon^2 y_1 y_{1,x} + \dots$$

One then obtains the following systems for y_1 and y_2 :

$$(1.17) \quad \begin{cases} y_{1,t} + y_{1,x} + y_{1,xxx} = 0 & \text{in } (0, T) \times (0, L), \\ y_1(\cdot, 0) = y_{1,x}(\cdot, L) = 0 & \text{in } (0, T), \\ y_1(\cdot, L) = u_1 & \text{in } (0, T), \end{cases}$$

and

$$(1.18) \quad \begin{cases} y_{2,t} + y_{2,x} + y_{2,xxx} + y_1 y_{1,x} = 0 & \text{in } (0, T) \times (0, L), \\ y_2(\cdot, 0) = y_{2,x}(\cdot, L) = 0 & \text{in } (0, T), \\ y_2(\cdot, L) = u_2 & \text{in } (0, T). \end{cases}$$

The aim is then to find the corresponding controls u_1 and u_2 such that if $y_1(0, \cdot) = y_2(0, \cdot) = 0$, then $y_1(T, \cdot) = 0$ and the L^2 -orthogonal projection of $y_2(T, \cdot)$ on \mathcal{M}_D is a given (non-zero) element in \mathcal{M}_D (the space of the directions prohibited by the linearized system, see Proposition 5.1). To this end, in [23, 16, 18], the authors used delicate contradiction arguments to capture the structure of their studied KdV system.

We here use the ideas in the spirit of the joint work with Coron and Koenig [24]. The starting point of the analysis is also the power series expansion method. The strategy is to characterize all possible u_1 which steers 0 at time 0 to 0 at time T (Proposition 6.1). This is done by taking the Fourier transform with respect to time of the solution y_1 and applying Paley-Wiener's theorem.

We then investigate the projection of y_2 into the unreachable space \mathcal{M}_D (in Section 6). This is where our analysis deviates from [24] to take advantage of the fact that q is real and to use the characterization via Paley-Wiener's theorem to construct u_1 for which the projection of y_2 into \mathcal{M}_D has a desired direction (see the proof of Assertion *ii*) of Theorem 1.3). Concerning the projection of y_2 into \mathcal{M}_D , we use the information of $|\hat{u}(z + iq/2)|^2$ instead of $\hat{u}(z)\overline{\hat{u}(z + iq)}$ as considered in [24]. This new proposal has two advantages. First, the way to obtain the projection of y_2 into the unreachable space \mathcal{M}_D is less complex. More importantly, it allows us to obtain the criterion on the local controllability of (1.1) in finite time given in Theorem 1.3 by taking into account the fact that $|\hat{u}(z + iq/2)|^2 \geq 0$. Another important part in the proof of the local controllability in Theorem 1.2 and Theorem 1.3 is to obtain the observability for the linearized system with initial/final data in \mathcal{M}_D^\perp . To this end, one needs to establish the optimal results on the stability and the well-posedness of the linearized system of (1.1) and related ones in which the boundary conditions are in fractional Sobolev spaces, see Section 3. The proof of the local controllability in finite time given in Assertion *ii*) of Theorem 1.3 is based on a new strategy to handle the lack of symmetry in our control system in comparison with the one using the right Neumann control. Concerning the KdV system using the right Neumann control, one can bring a state forward or backward easily thanks to the symmetry of the boundary conditions. This is not the case for (1.1) since the Dirichlet condition on the right is controlled and the Dirichlet condition on the left is imposed to be 0. To overcome this, we introduce a new approach that involves a KdV system with a new type boundary condition, investigated in Section 3, in a backward way. Our analysis uses the Banach fixed point arguments while the Brouwer fixed point theorem is usually applied to related contexts. It is worth noting that the analysis in this paper is based on the expansion up to the second order even if the unreachable space is of dimension 1. This explains the cost of the control given in Assertion *ii*) of Theorem 1.3. This way is different from what has been done for the KdV system with the right Neumann control where the expansion up to the third order is required.

Here are other comments on the analysis.

i) Various interpolation inequalities involving fractional Sobolev spaces and their dual spaces are established and used in the proof of Theorem 1.2 and Theorem 1.3.

ii) The analysis for the well-posedness of the KdV equation with various boundary conditions is in the spirit of the joint work with Coron and Koenig [24]. This partly involves a connection between the linear KdV equation and the linear KdV-Burgers equation as previously used by Bona et al. [11] and inspired by the work of Bourgain [12], and Molinet and Ribaud [38]. In this work, to deal with the boundary conditions containing the information of y_{xx} and derive *optimal* estimates on y , y_x , and y_{xx} in fractional Sobolev spaces and their dual spaces, new ideas and technique are implemented (see Remark 3.6). It is worth noting that similar estimates for y_{xx} are not known for the KdV equation in the real line setting and it is not clear whether such estimates hold for that setting. Our estimates particularly allow us to get the observability inequality for the critical lengths, which plays a role in determining the unreachable space for the linearized KdV system (see Section 3). The unreachable space is not known previously and this determination is also a contribution of the paper.

iii) The arguments used in this paper to disprove the small time, local controllability improve the ones in [24]. One of the key steps in the proof is to establish new (optimal) estimates for solutions of the KdV equations with various boundary conditions (see Section 3). The proof also relies on the positivity of a scalar product which comes naturally in the study of the controllability of the KdV system for small time. The positivity of a scalar product for small time was also a crucial point of several lacks of small time, local controllability results for systems with infinite propagation speed. Nevertheless, the previous ways used to derive the positivity are different, see, e.g., [21, 8, 34, 6, 13].

1.4. Organisation of the paper. The paper is organized as follows. Section 2 is devoted to the study of the linear KdV-Burgers equation with the periodic boundary condition. Section 3 is devoted to the study of the (linear and nonlinear) KdV equations with various boundary conditions. Section 4 is devoted to the study of critical lengths and some properties of the unreachable space \mathcal{M}_D . We there prove (1.5), (1.6), and (1.11), and establish Proposition 1.1. In Section 5, we prove (1.12) and establish the corresponding observability inequality (Lemma 5.2). The proof of Theorem 1.4 is given at the end of this section. In Section 6, we study properties of the controls which steer 0 at time 0 to 0 at time T. Using results in Section 6, we study attainable directions for small time in Section 7. The proofs of Theorem 1.2 and Theorem 1.3 are given in Section 8 and Section 9, respectively. Several technical results are stated and proved in the appendix. The Scilab program is also given there.

2. LINEAR KdV-BURGERS EQUATIONS

This section is devoted to the study of the linear KdV-Burgers equations with the periodic boundary condition. Here is the main result of this section, which plays an essential role in the study of the linearized KdV equation with various boundary conditions investigated in Section 3.

Proposition 2.1. *Let $g \in L^1(\mathbb{R}_+; L^2(0, L))$ with $\int_0^L g(t, x) dx = 0$ for a.e. $t > 0$, and let $y_0 \in L^2(0, L)$ be such that $\int_0^L y_0(x) dx = 0$. There exists a unique solution $y \in C([0, +\infty); L^2(0, L)) \cap L^2([0, +\infty); H^1([0, L]))$ which is periodic in space of the system*

$$(2.1) \quad \begin{cases} y_t + 4y_x + y_{xxx} - 3y_{xx} = g & \text{in } \mathbb{R}_+ \times (0, L), \\ y(0, \cdot) = y_0 & \text{in } (0, L). \end{cases}$$

Moreover, for $t \in \mathbb{R}_+$ and $x \in [0, L]$,

$$(2.2) \quad \|y(t, \cdot)\|_{L^2(0, L)} + \|y(\cdot, x)\|_{H^{1/3}(\mathbb{R}_+)} + \|y_x(\cdot, x)\|_{L^2(\mathbb{R}_+)} + \|y_{xx}(\cdot, x)\|_{[H^{1/3}(\mathbb{R}_+)]^*} \\ \leq C \left(\|y_0\|_{L^2(0, L)} + \|g\|_{L^1(\mathbb{R}_+; L^2(0, L))} \right),$$

and

$$(2.3) \quad \|y(\cdot, x)\|_{L^2(\mathbb{R}_+)} + \|y_x(\cdot, x)\|_{[H^{1/3}(\mathbb{R}_+)]^*} \leq C \left(\|y_0\|_{H^{-1}(0, L)} + \|g\|_{L^1(\mathbb{R}_+; H^{-1}(0, L))} \right)$$

for some positive constant C depending only on L .

The rest of this section consisting of two subsections is organized as follows. In the first subsection, we present and prove a useful lemma in the spirit of Bourgain's. The second subsection is devoted to the proof of Proposition 2.1.

2.1. A useful lemma. The following result in the spirit of Bourgain's is used in the proof of Proposition 2.1.

Lemma 2.1. *Let $(a_n)_{n \in \mathbb{Z} \setminus \{0\}} \subset \mathbb{R}$ be such that $\sum_{n \in \mathbb{Z} \setminus \{0\}} a_n^2 < +\infty$. For $-2 < s \leq 3$, there exists a positive constant $C = C_s$ such that*

$$(2.4) \quad \int_{\mathbb{R}} (1 + |z|)^{2s/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n| |n|^{3-s}}{(z + 4n - n^3)^2 + n^4} \right)^2 dz \leq C \sum_{n \in \mathbb{Z} \setminus \{0\}} a_n^2.$$

Proof. For $s > -5/2$ and for $-1 \leq z \leq 1$, one has

$$\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n| |n|^{3-s}}{(z + 4n - n^3)^2 + n^4} \leq C \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} a_n^2 \right)^{1/2}.$$

It hence suffices to prove that, for $-2 < s \leq 3$,

$$(2.5) \quad \int_{z \in \mathbb{R}; |z| > 1} (1 + |z|)^{2s/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n| |n|^{3-s}}{(z + 4n - n^3)^2 + n^4} \right)^2 dz \leq C \sum_{n \in \mathbb{Z} \setminus \{0\}} a_n^2.$$

We have

$$(2.6) \quad \int_{z > 1} (1 + |z|)^{2s/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n| |n|^{3-s}}{(z + 4n - n^3)^2 + n^4} \right)^2 dz \\ \leq C \sum_{m \in \mathbb{N}_*} \int_{m^3}^{(m+1)^3} m^{2s} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n| |n|^{3-s}}{(z + 4n - n^3)^2 + n^4} \right)^2 dz.$$

For $m^3 \leq z \leq (m+1)^3$, one gets

$$(2.7) \quad \sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n| |n|^{3-s}}{(z + 4n - n^3)^2 + n^4} \\ = \sum_{k; m+k \in \mathbb{Z} \setminus \{0\}} \frac{|m+k|^{3-s} |a_{m+k}|}{(z + 4(m+k) - (m+k)^3)^2 + (m+k)^4} \\ \leq C \sum_{\substack{k; |k| \leq 2m \\ m+k \in \mathbb{Z} \setminus \{0\}}} \frac{|m+k|^{3-s} |a_{m+k}|}{m^4 (|k|+1)^2} + C \sum_{\substack{k; |k| \geq 2m+1 \\ m+k \in \mathbb{Z} \setminus \{0\}}} \frac{|k|^{3-s} |a_{m+k}|}{k^4 (|k|+1)^2}.$$

Combining (2.6) and (2.7) yields

$$(2.8) \quad \int_{z > 1} (1 + |z|)^{2s/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n| |n|^{3-s}}{(z + 4n - n^3)^2 + n^4} \right)^2 dz \\ \leq C \sum_{m \in \mathbb{N}_*} \left(\sum_{\substack{k; |k| \leq 2m \\ m+k \in \mathbb{Z} \setminus \{0\}}} \frac{|m+k|^{3-s} |a_{m+k}|}{m^{3-s} (|k|+1)^2} + C \sum_{\substack{k; |k| \geq 2m+1 \\ m+k \in \mathbb{Z} \setminus \{0\}}} \frac{|m|^{1+s} |a_{m+k}|}{|k|^{1+s} (|k|+1)^2} \right)^2 \\ \leq C \sum_{m \in \mathbb{N}_*} \left(\sum_{\substack{k \\ m+k \in \mathbb{Z} \setminus \{0\}}} \frac{|a_{m+k}|^2}{(|k|+1)^{1+\varepsilon}} \right) \left(\sum_{\substack{k; |k| \leq 2m \\ m+k \in \mathbb{Z} \setminus \{0\}}} \frac{|m+k|^{6-2s}}{m^{6-2s} (|k|+1)^{3-\varepsilon}} + \sum_{\substack{k; |k| \geq 2m+1 \\ m+k \in \mathbb{Z} \setminus \{0\}}} \frac{m^{2+2s}}{|k|^{2+2s} (|k|+1)^{3-\varepsilon}} \right).$$

Here $0 < \varepsilon \leq 1$ is fixed such that $s > (\varepsilon - 4)/2$ (the fact $s > -2$ is used). We have, for $m \in \mathbb{N}_*$, $|k| \leq 2m$, and $m + k \in \mathbb{Z} \setminus \{0\}$,

$$\frac{|m+k|^{6-2s}}{m^{6-2s}} \leq C \text{ for } s \leq 3,$$

and, for $m \in \mathbb{N}_*$,

$$\sum_{\substack{k; |k| \geq 2m+1 \\ m+k \in \mathbb{Z} \setminus \{0\}}} \frac{m^{2+2s}}{|k|^{2+2s}(|k|+1)^{3-\varepsilon}} = m^{2+2s} \sum_{\substack{k; |k| \geq 2m+1 \\ m+k \in \mathbb{Z} \setminus \{0\}}} \frac{1}{k^{2+2s}(|k|+1)^{3-\varepsilon}} \leq \frac{C}{m^{2-\varepsilon}} \leq C \text{ for } s > (\varepsilon - 4)/2.$$

Using the fact, for $m \in \mathbb{N}_*$,

$$\sum_{k; m+k \in \mathbb{Z} \setminus \{0\}} \frac{1}{(|k|+1)^{3-\varepsilon}} < +\infty \text{ for } 0 < \varepsilon < 1,$$

we derive from (2.8) that

$$(2.9) \quad \int_{z>1} (1+|z|)^{2s/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n||n|^{3-s}}{(z+4n-n^3)^2+n^4} \right)^2 dz \\ \leq C \sum_{m \in \mathbb{N}_*} \sum_{k; m+k \in \mathbb{Z} \setminus \{0\}} \frac{|a_{m+k}|^2}{(|k|+1)^{1+\varepsilon}} \leq C \sum_{k \in \mathbb{Z}} \frac{1}{(|k|+1)^{1+\varepsilon}} \sum_{m; m+k \in \mathbb{Z} \setminus \{0\}} |a_{m+k}|^2,$$

which yields

$$(2.10) \quad \int_{z>1} (1+|z|)^{2s/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n||n|^{3-s}}{(z+4n-n^3)^2+n^4} \right)^2 dz \leq C \sum_{n \in \mathbb{Z} \setminus \{0\}} |a_n|^2.$$

Similarly, we have

$$(2.11) \quad \int_{z<-1} (1+|z|)^{2s/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n||n|^{3-s}}{(z+4n-n^3)^2+n^4} \right)^2 dz \leq C \sum_{n \in \mathbb{Z} \setminus \{0\}} |a_n|^2.$$

Estimate (2.5) now follows from (2.10) and (2.11). The proof is complete. \square

2.2. Proof of Proposition 2.1. We only derive the estimates. The uniqueness follows from the estimates and the existence follows from the proof of these estimates as well.

Multiplying the equation of y by y and integrating by parts, we have

$$\frac{1}{2} \frac{d}{dt} \int_0^L |y(t, x)|^2 dx + 3 \int_0^L |y_x(t, x)|^2 dx = \int_0^L g(t, x) y(t, x) dx.$$

This yields

$$\frac{1}{2} \frac{d}{dt} \int_0^L |y(t, x)|^2 dx \leq \|g(t, \cdot)\|_{L^2(0, L)} \|y(t, \cdot)\|_{L^2(0, L)}.$$

Applying the Grönwall lemma, we obtain the desired estimate for $\|y(t, \cdot)\|_{L^2(0, L)}$.

We next establish the estimates for $\|y(\cdot, x)\|_{H^{1/3}(\mathbb{R}_+)}$, $\|y_x(\cdot, x)\|_{L^2(\mathbb{R}_+)}$, and $\|y_{xx}(\cdot, x)\|_{[H^{1/3}(\mathbb{R}_+)]^*}$. For notational ease, we assume that $L = 2\pi$. It suffices to consider the case $g \equiv 0$ and the case $y_0 \equiv 0$ separately.

We first consider the case $g \equiv 0$. Write the solution under the form

$$y(t, x) = \sum_{n \in \mathbb{Z}} a_n(t) e^{inx} \text{ in } \mathbb{R}_+ \times (0, L) = \mathbb{R}_+ \times (0, 2\pi).$$

We then derive that

$$y_0(x) = \sum_{n \in \mathbb{Z}} a_n(0) e^{inx} \text{ for } x \in [0, 2\pi]$$

and

$$a'_n(t) = (-3n^2 - i(4n - n^3))a_n(t) \text{ for } t \in \mathbb{R}_+.$$

We thus have

$$a_n(t) = e^{(-3n^2 - i(4n - n^3))t} a_n(0) \text{ for } t \in \mathbb{R}_+, n \in \mathbb{Z}.$$

Since $\int_0^L y_0 dx = 0$, it follows that $a_0 = 0$. This in turn implies that $a_0(t) = 0$ for $t \in \mathbb{R}_+$. We hence obtain

$$y(t, x) = \sum_{n \in \mathbb{Z} \setminus \{0\}} e^{inx} e^{(-3n^2 - i(4n - n^3))t} a_n(0) \text{ in } \mathbb{R}_+ \times (0, 2\pi).$$

Extend $y(t, x)$ for $t < 0$ by

$$(2.12) \quad y(t, x) = \sum_{n \in \mathbb{Z} \setminus \{0\}} e^{inx} e^{(3n^2 - i(4n - n^3))t} a_n(0) \text{ in } \mathbb{R}_- \times (0, 2\pi),$$

and still denote this extension by $y(t, x)$. We have, for $z \in \mathbb{R}$ and $n \in \mathbb{Z} \setminus \{0\}$,

$$(2.13) \quad \int_0^\infty e^{(-3n^2 - i(4n - n^3))t} e^{-itz} dt + \int_{-\infty}^0 e^{(3n^2 - i(4n - n^3))t} e^{-itz} dt \\ = \frac{1}{3n^2 + i(4n - n^3 + z)} + \frac{1}{3n^2 - i(4n - n^3 + z)} = \frac{6n^2}{9n^4 + (4n - n^3 + z)^2}.$$

This implies

$$(2.14) \quad \|y(\cdot, x)\|_{H^{1/3}(\mathbb{R}_+)}^2 \leq C \int_{z \in \mathbb{R}} (1 + |z|)^{2/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{n^2 |a_n(0)|}{n^4 + (z + 4n - n^3)^2} \right)^2 dz.$$

Applying Lemma 2.1 with $s = 1$, we obtain

$$(2.15) \quad \|y(\cdot, x)\|_{H^{1/3}(\mathbb{R}_+)}^2 \leq C \|y_0\|_{L^2(0, 2\pi)}^2.$$

Similarly, we have

$$(2.16) \quad \|y_x(\cdot, x)\|_{L^2(\mathbb{R}_+)}^2 \leq C \int_{z \in \mathbb{R}} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|n|^3 |a_n(0)|}{(z + 4n - n^3)^2 + n^4} \right)^2$$

and

$$(2.17) \quad \|y_{xx}(\cdot, x)\|_{[H^{1/3}(\mathbb{R}_+)]^*}^2 \leq C \int_{z \in \mathbb{R}} (1 + |z|)^{-2/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{n^4 |a_n(0)|}{(z + 4n - n^3)^2 + n^4} \right)^2.$$

Applying Lemma 2.1 with $s = 0$ and $s = -1$, we get

$$(2.18) \quad \|y_x(\cdot, x)\|_{L^2(\mathbb{R}_+)}^2 \leq C \|y_0\|_{L^2(0, 2\pi)}^2.$$

and

$$(2.19) \quad \|y_{xx}(\cdot, x)\|_{[H^{1/3}(\mathbb{R}_+)]^*}^2 \leq C \|y_0\|_{L^2(0, 2\pi)}^2.$$

The proof of (2.2) in the case $g \equiv 0$ is complete.

We next deal with (2.3). From (2.12) and (2.13), we have

$$\|y(\cdot, x)\|_{L^2(\mathbb{R}_+)}^2 \leq C \int_{z \in \mathbb{R}} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|n|^3(|a_n(0)|/|n|)}{n^4 + (z + 4n - n^3)^2} \right)^2 dz$$

and

$$(2.20) \quad \|y_x(\cdot, x)\|_{[H^{1/3}(\mathbb{R}_+)]^*}^2 \leq C \int_{z \in \mathbb{R}} (1 + |z|)^{-2/3} \left(\sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|n|^4(|a_n(0)|/|n|)}{(z + 4n - n^3)^2 + n^4} \right)^2.$$

Applying Lemma 2.1 with $s = 0$ and $s = -1$, we obtain

$$(2.21) \quad \|y(\cdot, x)\|_{L^2(\mathbb{R}_+)}^2 \leq C \sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n(0)|^2}{n^2} \leq C \|y_0\|_{H^{-1}(0, 2\pi)}^2$$

and

$$(2.22) \quad \|y_x(\cdot, x)\|_{[H^{1/3}(\mathbb{R}_+)]^*}^2 \leq C \sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{|a_n(0)|^2}{n^2} \leq C \|y_0\|_{H^{-1}(0, 2\pi)}^2.$$

This yields (2.3).

We next deal with the case $y_0 \equiv 0$. The proof, in this case, can be derived from the previous case as follows. For $t > 0$, let $W(t)$ be the mapping which maps y_0 to $y(t, \cdot)$ with $g \equiv 0$. We then have ⁵

$$(2.23) \quad y(t, x) = \int_0^t W(t-s)g(s, x) ds.$$

This implies

$$(2.24) \quad y_x(t, x) = \int_0^t \partial_x(W(t-s)g(s, x)) ds = \int_0^\infty \mathbf{1}_{(0,t)}(s) \partial_x(W(t-s)g(s, x)) ds.$$

Hence

$$\begin{aligned} \|y_x(\cdot, x)\|_{L_t^2(\mathbb{R}_+)} &\leq \int_0^{+\infty} \|\mathbf{1}_{(0,t)}(s) \partial_x(W(t-s)g(s, x))\|_{L_t^2(\mathbb{R}_+)} ds \\ &= \int_0^{+\infty} \left(\int_0^\infty |\mathbf{1}_{(0,t)}(s) \partial_x(W(t-s)g(s, x))|^2 dt \right)^{1/2} ds \\ &= \int_0^{+\infty} \left(\int_s^\infty |\partial_x(W(t-s)g(s, x))|^2 dt \right)^{1/2} ds. \end{aligned}$$

By applying the results in the previous case, we have

$$\left(\int_s^\infty |\partial_x(W(t-s)g(s, x))|^2 dt \right)^{1/2} \leq C \|g(s, x)\|_{L_x^2(0, L)}.$$

We thus obtain

$$(2.25) \quad \|y_x(\cdot, x)\|_{L_t^2(\mathbb{R}_+)} \leq C \int_0^{+\infty} \|g(s, x)\|_{L_x^2} ds.$$

⁵This identity can be derived from (2.28) and (2.29) in Remark 2.3.

By the same arguments, we have

$$(2.26) \quad \|y_{xx}(\cdot, x)\|_{[H_t^{1/3}(\mathbb{R}_+)]^*} \leq \int_0^{+\infty} \|\mathbf{1}_{(0,t)}(s)\partial_{xx}(W(t-s)g(s,x))\|_{[H_t^{1/3}(\mathbb{R}_+)]^*} ds \\ = \int_0^{+\infty} \|\partial_{xx}(W(t-s)g(s,x))\|_{[H_t^{1/3}(s,+\infty)]^*} ds \leq C \int_0^{+\infty} \|g(s,x)\|_{L_x^2(0,L)} ds.$$

Similarly, we obtain

$$(2.27) \quad \|y(\cdot, x)\|_{H_t^{1/3}(\mathbb{R}_+)} \leq \int_0^{+\infty} \|\mathbf{1}_{(0,t)}(s)W(t-s)g(s,x)\|_{H_t^{1/3}(\mathbb{R}_+)} ds \\ = \int_0^{+\infty} \|W(t-s)g(s,x)\|_{H_t^{1/3}(s,+\infty)} ds \leq C \int_0^{+\infty} \|g(s,x)\|_{L_x^2(0,L)} ds.$$

Assertion (2.2) in the case $y_0 \equiv 0$ now follows from (2.25), (2.26), and (2.27).

Assertion (2.3) in the case $y_0 \equiv 0$ follows similarly and the details are omitted.

The proof is complete. \square

Remark 2.1. The proof gives as well that

$$y \in C([0, L]; H^{1/3}(0, +\infty)), \quad y_x \in C([0, L]; L^2(0, +\infty)), \quad \text{and} \quad y_{xx} \in C([0, L]; [H^{1/3}(0, +\infty)]^*).$$

Remark 2.2. Assume that $g = 0$ in $\mathbb{R}_+ \times (0, L)$ and $y_0 \in C^\infty([0, L])$ is such that $\int_0^L y_0(x) dx = 0$, $y_0 - c \in C_c^\infty((0, L))$ for some constant c . Using (2.12), one can show that $y \in C^\infty([0, +\infty) \times [0, L])$. Moreover, using the equation of y , one can show that $\partial_t^k y(0, 0) = \partial_t^k y(0, L) = 0$ for all $k \geq 1$.

Remark 2.3. Assume that $y_0 = 0$ in $(0, L)$ and $g \in C_c^\infty((0, +\infty) \times [0, L])$ being such that $\int_0^L g(t, x) dx = 0$ for $t > 0$. One can prove that the solution is then smooth and is 0 around the time 0. Indeed, for notational ease, assume that $L = 2\pi$. One then can show that

$$(2.28) \quad y(t, x) = \sum_{n \in \mathbb{Z} \setminus \{0\}} e^{inx} \int_0^t e^{(-3n^2 - i(4n - n^3))(t-s)} g_n(s) ds,$$

where

$$(2.29) \quad g(t, x) = \sum_{n \in \mathbb{Z} \setminus \{0\}} g_n(t) e^{inx}.$$

The properties of y follow.

Remark 2.4. There is no term $\|y_{xx}\|_{[H^{2/3}(\mathbb{R}_+)]^*}$ in (2.3) since we applied Lemma 2.1, which requires $s > -2$.

Remark 2.5. In [24, Lemma 4.1], we proved that, when $y_0 = 0$,

$$(2.30) \quad \|y(\cdot, x)\|_{L^2(\mathbb{R}_+)} \leq C \|g\|_{L^1(\mathbb{R}_+ \times (0, L))}.$$

Estimate (2.3) improves this.

3. LINEARIZED KDV EQUATIONS

In this section, we establish several well-posedness results (Proposition 3.1, Proposition 3.2, and Proposition 3.3) for the linearized KdV equation equipped with various boundary conditions in the energy space X_T defined in (1.14) and present some of their consequences. Proposition 3.1 implies the local well-posedness of (1.1) in X_T (Proposition 3.4). Proposition 3.1 and Proposition 3.2 are the starting points of our analysis in deriving the unreachable space and establishing the corresponding observability inequality in Section 5 (Proposition 5.1). Proposition 3.3 and its consequence (Proposition 3.5) will be used in the proof of Theorem 1.3.

Here is the first main result of this section.

Proposition 3.1. *Let $L > 0$, $T > 0$, $(h_1, h_2, h_3) \in H^{1/3}(0, T) \times H^{1/3}(0, T) \times L^2(0, T)$, $f \in L^1((0, T); L^2(0, L))$, and $y_0 \in L^2(0, L)$. There exists a unique solution $y \in X_T$ of the system*

$$(3.1) \quad \begin{cases} y_t + y_x + y_{xxx} = f & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = h_1, y(\cdot, L) = h_2, y_x(\cdot, L) = h_3 & \text{in } (0, T), \\ y(0, \cdot) = y_0 & \text{in } (0, L). \end{cases}$$

Moreover, for $x \in [0, L]$,

$$(3.2) \quad \begin{aligned} & \|y\|_{X_T} + \|y(\cdot, x)\|_{H^{1/3}(0, T)} + \|y_x(\cdot, x)\|_{L^2(0, T)} + \|y_{xx}(\cdot, x)\|_{[H^{1/3}(0, T)]^*} \\ & \leq C_{T, L} \left(\|y_0\|_{L^2(0, L)} + \|f\|_{L^1((0, T); L^2(0, L))} + \|(h_1, h_2)\|_{H^{1/3}(0, T)} + \|h_3\|_{L^2(0, T)} \right), \end{aligned}$$

and

$$(3.3) \quad \begin{aligned} & \|y(\cdot, x)\|_{L^2(0, T)} + \|y_x(\cdot, x)\|_{[H^{1/3}(0, T)]^*} \\ & \leq C_{T, L} \left(\|y_0\|_{[H^1(0, L)]^*} + \|f\|_{L^1((0, T); [H^1(0, L)]^*)} + \|(h_1, h_2)\|_{L^2(0, T)} + \|h_3\|_{[H^{1/3}(0, T)]^*} \right), \end{aligned}$$

where $C_{T, L}$ denotes a positive constant independent of x , y_0 , f , and h_1, h_2, h_3 .

Here and in what follows, $[H^s]^*$ denotes the dual space of H^s for $s > 0$ and it is equipped with the standard corresponding norm.

Remark 3.1. In Proposition 3.1, we implicitly admit that

$$y \in C([0, L]; H^{1/3}(0, T)), \quad y_x \in C([0, L]; L^2(0, T)), \quad \text{and} \quad y_{xx} \in C([0, L]; [H^{1/3}(0, T)]^*).$$

These facts are derived from the proof. Results and analysis which are related to Proposition 3.1 will be discussed in Remark 3.6.

The next result is on the well-posedness of the linearized KdV system for which the Dirichlet condition and the second derivative in x on the right are described.

Proposition 3.2. *Let $L > 0$, $T > 0$, $(h_1, h_2, h_3) \in H^{1/3}(0, T) \times H^{1/3}(0, T) \times [H^{1/3}(0, T)]^*$, $f \in L^1((0, T); L^2(0, L))$, and $y_0 \in L^2(0, L)$. There exists a unique solution $y \in X_T$ of the system*

$$(3.4) \quad \begin{cases} y_t + y_x + y_{xxx} = f & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = h_1, y(\cdot, L) = h_2, y_{xx}(\cdot, L) = h_3 & \text{in } (0, T), \\ y(0, \cdot) = y_0 & \text{in } (0, L). \end{cases}$$

Moreover, for $0 \leq x \leq L$,

$$(3.5) \quad \|y\|_{X_T} + \|y(\cdot, x)\|_{H^{1/3}(0,T)} + \|y_x(\cdot, x)\|_{L^2(0,T)} + \|y_{xx}(\cdot, x)\|_{[H^{1/3}(0,T)]^*} \\ \leq C_{T,L} \left(\|y_0\|_{L^2(0,T)} + \|f\|_{L^1((0,T);L^2(0,L))} + \|(h_1, h_2)\|_{H^{1/3}(0,T)} + \|h_3\|_{[H^{1/3}(0,T)]^*} \right),$$

and

$$(3.6) \quad \|y(\cdot, x)\|_{L^2(0,T)} + \|y_x(\cdot, x)\|_{[H^{1/3}(0,T)]^*} \\ \leq C_{T,L} \left(\|y_0\|_{[H^1(0,L)]^*} + \|f\|_{L^1((0,T);[H^1(0,L)]^*)} + \|(h_1, h_2)\|_{L^2(0,T)} + \|h_3\|_{[H^{2/3}(0,T)]^*} \right),$$

for some positive constant $C_{T,L}$ independent of x , y_0 , f , and (h_1, h_2, h_3) .

Here is the third main result of this section.

Proposition 3.3. *Let $L > 0$, $T > 0$, $(h_1, h_2, h_3) \in L^2(0, T) \times H^{1/3}(0, T) \times L^2(0, T)$, $f \in L^1((0, T); L^2(0, L))$, and $y_0 \in L^2(0, L)$. There exists a unique solution $y \in X_T$ of the system*

$$(3.7) \quad \begin{cases} y_t + y_x + y_{xxx} = f & \text{in } (0, T) \times (0, L), \\ y_x(\cdot, 0) = h_1, y(\cdot, L) = h_2, y_x(\cdot, L) = h_3 & \text{in } (0, T), \\ y(0, \cdot) = y_0 & \text{in } (0, L). \end{cases}$$

Moreover, for $0 \leq x \leq L$,

$$(3.8) \quad \|y\|_{X_T} + \|y(\cdot, x)\|_{H^{1/3}(0,T)} + \|y_x(\cdot, x)\|_{L^2(0,T)} + \|y_{xx}(\cdot, x)\|_{[H^{1/3}(0,T)]^*} \\ \leq C_{T,L} \left(\|y_0\|_{L^2(0,T)} + \|f\|_{L^1((0,T);L^2(0,L))} + \|h_2\|_{H^{1/3}(0,T)} + \|(h_1, h_3)\|_{L^2(0,T)} \right)$$

and

$$(3.9) \quad \|y(\cdot, x)\|_{L^2(0,T)} + \|y_x(\cdot, x)\|_{[H^{1/3}(0,T)]^*} \\ \leq C_{T,L} \left(\|y_0\|_{[H^1(0,L)]^*} + \|f\|_{L^1((0,T);[H^1(0,L)]^*)} + \|(h_1, h_2)\|_{L^2(0,T)} + \|h_3\|_{[H^{1/3}(0,T)]^*} \right),$$

for some positive constant $C_{T,L}$ independent of x , y_0 , f , and (h_1, h_2, h_3) .

As a consequence of Proposition 3.1, we obtain the following well-posedness result, which particularly yields the well-posedness for system (1.1).

Proposition 3.4. *Let $L > 0$ and $T > 0$. There exists $\varepsilon_0 > 0$ such that for $(h_1, h_2, h_3) \in H^{1/3}(0, T) \times H^{1/3}(0, T) \times L^2(0, T)$, $f \in L^1((0, T); L^2(0, L))$, and $y_0 \in L^2(0, L)$ with*

$$\|y_0\|_{L^2(0,L)} + \|f\|_{L^1((0,T);L^2(0,L))} + \|(h_1, h_2)\|_{H^{1/3}(0,T)} + \|h_3\|_{L^2(0,T)} \leq \varepsilon_0,$$

there exists a unique solution $y \in X_T$ of the system

$$(3.10) \quad \begin{cases} y_t + y_x + y_{xxx} + yy_x = f & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = h_1, y(\cdot, L) = h_2, y_x(\cdot, L) = h_3 & \text{in } (0, T), \\ y(0, \cdot) = y_0 & \text{in } (0, L). \end{cases}$$

Moreover, we have

$$(3.11) \quad \begin{aligned} & \|y\|_{X_T} + \|y(\cdot, x)\|_{H^{1/3}(0,T)} + \|y_x(\cdot, x)\|_{L^2(0,T)} + \|y_{xx}(\cdot, x)\|_{[H^{1/3}(0,T)]^*} \\ & \leq C_{T,L} \left(\|y_0\|_{L^2(0,L)} + \|f\|_{L^1((0,T);L^2(0,L))} + \|(h_1, h_2)\|_{H^{1/3}(0,T)} + \|h_3\|_{L^2(0,T)} \right), \end{aligned}$$

for some positive constant $C_{T,L}$ independent of x , y_0 , f , and (h_1, h_2, h_3) .

As a consequence of Proposition 3.3, we obtain the following well-posedness result which will be used in the proof of Assertion *ii*) of Theorem 1.3.

Proposition 3.5. *Let $L > 0$ and $T > 0$. There exists $\varepsilon_0 > 0$ such that for $(h_1, h_2, h_3) \in L^2(0, T) \times H^{1/3}(0, T) \times L^2(0, T)$, $f \in L^1((0, T); L^2(0, L))$, and $y_0 \in L^2(0, L)$ with*

$$\|y_0\|_{L^2(0,L)} + \|f\|_{L^1((0,T);L^2(0,L))} + \|h_2\|_{H^{1/3}(0,T)} + \|(h_1, h_3)\|_{L^2(0,T)} \leq \varepsilon_0,$$

there exists a unique solution $y \in X_T$ of the system

$$(3.12) \quad \begin{cases} y_t + y_x + y_{xxx} + yy_x = f & \text{in } (0, T) \times (0, L), \\ y_x(\cdot, 0) = h_1, y(\cdot, L) = h_2, y_x(\cdot, L) = h_3, & \text{in } (0, T), \\ y(0, \cdot) = y_0 & \text{in } (0, L). \end{cases}$$

Moreover, for $0 \leq x \leq L$,

$$(3.13) \quad \begin{aligned} & \|y\|_{X_T} + \|y(\cdot, x)\|_{H^{1/3}(0,T)} + \|y_x(\cdot, x)\|_{L^2(0,T)} + \|y_{xx}(\cdot, x)\|_{[H^{1/3}(0,T)]^*} \\ & \leq C_{T,L} \left(\|y_0\|_{L^2(0,L)} + \|f\|_{L^1((0,T);L^2(0,L))} + \|h_2\|_{H^{1/3}(0,T)} + \|(h_1, h_3)\|_{L^2(0,T)} \right), \end{aligned}$$

for some positive constant $C_{T,L}$ independent of x , y_0 , f , and u .

Remark 3.2. The meaning of the solutions considered in Proposition 3.1, Proposition 3.2, and Proposition 3.3 are given in Definition 3.1, Definition 3.2, Definition 3.3, respectively. These are motivated by the integration by parts arguments. The meaning of the solutions in Proposition 3.4 and Proposition 3.5 are understood in the same manner where the nonlinear term yy_x plays as a part of the source f .

Proposition 3.4 and Proposition 3.5 follows from Proposition 3.1 and Proposition 3.3 by standard arguments. We just mention here that

$$\begin{aligned} \|yy_x\|_{L^1((0,T);L^2(0,L))} & \leq \int_0^T \sup_{x \in [0,L]} |y(t, x)| \left(\int_0^L |y_x(t, x)|^2 \right)^{\frac{1}{2}} \\ & \leq C \int_0^T \left(\int_0^L |y(t, x)|^2 + |y_x(t, x)|^2 \right)^{\frac{1}{2}} \left(\int_0^L |y_x(t, x)|^2 \right)^{\frac{1}{2}} \leq C \|y\|_{X_T}^2. \end{aligned}$$

The details of the proof are left to the reader. A sharper estimate can be derived from Lemma 8.1

The rest of this section is organized as follows. The proof of Proposition 3.1, Proposition 3.2, and Proposition 3.3 are given in Section 3.1, Section 3.2, and Section 3.3, respectively.

3.1. Proof of Proposition 3.1. We first give the meaning of the solutions considered in Proposition 3.1.

Definition 3.1. Let $L > 0$, $T > 0$, $(h_1, h_2, h_3) \in H^{1/3}(0, T) \times H^{1/3}(0, T) \times L^2(0, T)$, $f \in L^1((0, T); L^2(0, L))$, and $y_0 \in L^2(0, L)$. A solution $y \in X_T$ of the system (3.1) is a function $y \in X_T$ such that

$$(3.14) \quad \int_0^T \int_0^L f(t, x) \varphi(t, x) dx dt + \int_0^L y_0(x) \varphi(0, x) dx + \int_0^T h_3(t) \varphi_x(t, L) dt \\ - \int_0^T h_2(t) \varphi_{xx}(t, L) dt + \int_0^T h_1(t) \varphi_{xx}(t, 0) dt = - \int_0^T \int_0^L y(\varphi_t + \varphi_x + \varphi_{xxx}) dx dt$$

for all $\varphi \in C^3([0, T] \times [0, L])$ with $\varphi(T, \cdot) = 0$ and $\varphi(\cdot, 0) = \varphi(\cdot, L) = \varphi_x(\cdot, 0) = 0$.

Remark 3.3. One can check that if y is a smooth solution of (3.1) then y is a solution of (3.1) in the sense of Definition 3.1 by standard integration by parts arguments. In fact, Definition 3.1 is motivated by such arguments.

The proof of Proposition 3.1 is divided into two parts given in the following two subsections. The first part is on the existence and the estimates for a constructed solution. The second part is on the uniqueness.

3.1.1. Existence and estimates. The proof of (3.2) for a solution constructed below is given in two steps.

Step 1: We first consider the case where $y_0 = 0$ and $f = 0$. By the linearity, it suffices to consider the three cases $(h_1, h_2, h_3) = (0, 0, h_3)$, $(h_1, h_2, h_3) = (h_1, 0, 0)$, and $(h_1, h_2, h_3) = (0, h_2, 0)$ separately. In what follows, we extend h_1, h_2, h_3 by 0 for $t > T$ and still denote y the corresponding solution and these extensions by h_1, h_2, h_3 .

In what follows in this proof, for an appropriate function v defined on $\mathbb{R} \times (0, L)$, we denote by \hat{v} its Fourier transform with respect to t , i.e., for $z \in \mathbb{C}$,

$$\hat{v}(z, x) = \frac{1}{\sqrt{2\pi}} \int_0^{+\infty} v(t, x) e^{-izt} dt.$$

Extend y and h_1, h_2, h_3 by 0 for $t < 0$ and still denote these extensions by y , and h_1, h_2, h_3 . Then

$$(3.15) \quad y_t + y_x + y_{xxx} = 0 \text{ in } \mathbb{R} \times (0, L).$$

Taking the Fourier transform with respect to t , we obtain, for $z \in \mathbb{R}$,

$$(3.16) \quad iz\hat{y} + \hat{y}_x + \hat{y}_{xxx} = 0 \text{ in } (0, L).$$

For $z \in \mathbb{C}$, let $\lambda_j = \lambda_j(z)$ with $j = 1, 2, 3$ be the three solutions of the equation $\lambda^3 + \lambda + iz = 0$. Set

$$Q(z) := \begin{pmatrix} 1 & 1 & 1 \\ e^{\lambda_1 L} & e^{\lambda_2 L} & e^{\lambda_3 L} \\ \lambda_1 e^{\lambda_1 L} & \lambda_2 e^{\lambda_2 L} & \lambda_3 e^{\lambda_3 L} \end{pmatrix},$$

and

$$\Xi = \Xi(z) := \det \begin{pmatrix} 1 & 1 & 1 \\ \lambda_1 & \lambda_2 & \lambda_3 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 \end{pmatrix},$$

with the convention $\lambda_{j+3} = \lambda_j$ for $j \geq 1$. It is useful to note that

$$(3.17) \quad H(z) := \det Q(z) \Xi(z) \text{ in } \mathbb{C}.$$

is an analytic function, see, e.g., [24, Lemma A1]. Moreover, H only has a finite number of zeros on the real line since $H(z) \neq 0$ for $z \in \mathbb{R}$ with large $|z|$.

We first consider the case $(h_1, h_2, h_3) = (0, 0, h_3)$. Taking into account the equation of \hat{y} , we search for the solution of the form

$$(3.18) \quad \hat{y}(z, \cdot) = \sum_{j=1}^3 a_j e^{\lambda_j x},$$

where $a_j = a_j(z)$ for $j = 1, 2, 3$. Taking the boundary condition, we then have

$$\begin{aligned} \sum_{j=1}^3 a_j &= 0, \\ \sum_{j=1}^3 a_j e^{\lambda_j L} &= 0, \\ \sum_{j=1}^3 a_j \lambda_j e^{\lambda_j L} &= \hat{h}_3, \end{aligned}$$

This implies, with the convention $\lambda_{j+3} = \lambda_j$,

$$(3.19) \quad a_j = \frac{e^{\lambda_{j+2}L} - e^{\lambda_{j+1}L}}{\det Q} \hat{h}_3 \text{ for } j = 1, 2, 3.$$

To estimate the solution, we proceed as follows. By Lemma C.1 in the appendix, there exists $g_3 \in C^\infty(\mathbb{R})$ with $\text{supp } g_3 \subset [T, 3T]$ such that if z is a real solution of the equation $H(z) = 0$ of order m then z is also a real solution of order m of $\hat{h}_3(z) - \hat{g}_3(z)$, and, for $k \geq 1$,

$$(3.20) \quad \|g_3\|_{H^k(\mathbb{R})} \leq C_k \|h_3\|_{H^{-1/3}(\mathbb{R})}.$$

We now establish (3.2). Let y_3 be the solution of (3.1) where (h_1, h_2, h_3) are $(0, 0, h_3 - g_3)$, and $f = 0$, and $y_0 = 0$. We have, by applying (3.19) to y_3 ,

$$(3.21) \quad \hat{y}_3(z, x) = \frac{\hat{h}_3(z) - \hat{g}_3(z)}{\det Q(z)} \sum_{j=1}^3 (e^{\lambda_{j+2}L} - e^{\lambda_{j+1}L}) e^{\lambda_j x} \text{ for a.e. } x \in (0, L).$$

We derive that, by the choice of g_3 , for $z \in \mathbb{R}$ and $|z| \leq \gamma$,

$$(3.22) \quad \left| \frac{\hat{h}_3(z) - \hat{g}_3(z)}{H(z)} \right| \left| \Xi(z) \sum_{j=1}^3 (e^{\lambda_{j+2}L} - e^{\lambda_{j+1}L}) e^{\lambda_j x} \right| \leq C_{T,\gamma} \|h_3 - g_3\|_{H^{-1/3}(\mathbb{R})},$$

and, by Lemma 3.1 below on the behaviors of $\lambda_j(z)$ for large z (the case large negative z can be obtained by considering the conjugate), we obtain, for $z \in \mathbb{R}$, $|z| \geq \gamma$ with sufficiently large γ ,

$$(3.23) \quad \left| \frac{1}{\det Q} \sum_{j=1}^3 (e^{\lambda_{j+2}L} - e^{\lambda_{j+1}L}) e^{\lambda_j x} \right| \leq \frac{C}{(1 + |z|)^{1/3}}.$$

Combining (3.22) and (3.23) yields

$$\|y_3(\cdot, x)\|_{H^{1/3}(\mathbb{R})} \leq C_{T,L} \|h_3 - g_3\|_{L^2(\mathbb{R})}$$

and

$$\|y_3(\cdot, x)\|_{L^2(\mathbb{R})} \leq C_{T,L} \|h_3 - g_3\|_{H^{-1/3}(\mathbb{R})}.$$

Similarly, we have

$$\|y_{3,x}(\cdot, x)\|_{L^2(\mathbb{R})} + \|y_{3,xx}(\cdot, x)\|_{H^{-1/3}(\mathbb{R})} \leq C_{T,L} \|h_3 - g_3\|_{L^2(\mathbb{R})}$$

and

$$\|y_{3,x}(\cdot, x)\|_{H^{-2/3}(\mathbb{R})} \leq C_{T,L} \|h_3 - g_3\|_{H^{-1/3}(\mathbb{R})}.$$

The estimates for $\|y(\cdot, x)\|_{H^{1/3}(0,T)}$, $\|y_x(\cdot, x)\|_{L^2(0,T)}$, and $\|y_{xx}(\cdot, x)\|_{[H^{1/3}(0,T)]^*}$ and the estimate for $\|y(\cdot, x)\|_{L^2(0,T)}$, $\|y_x(\cdot, x)\|_{[H^{1/3}(0,T)]^*}$ follow by noting that $y = y_3$ in $(0, T) \times (0, L)$ ⁶.

We also have, by integration by parts, for $0 \leq \tau_1 < \tau_2 \leq T$,

$$(3.24) \quad \frac{1}{2} \int_0^L |y(\tau_2, x)|^2 dx - \frac{1}{2} \int_0^L |y(\tau_1, x)|^2 dx + \frac{1}{2} \int_{\tau_1}^{\tau_2} \left(|y(t, L)|^2 - |y(t, 0)|^2 \right) dt \\ + \int_{\tau_1}^{\tau_2} \left(y_{xx}(t, L)y(t, L) - y_{xx}(t, 0)y(t, 0) \right) dt - \frac{1}{2} \int_{\tau_1}^{\tau_2} \left(|y_x(t, L)|^2 - |y_x(t, 0)|^2 \right) dt = 0.$$

Using the estimates for $\|y(\cdot, x)\|_{H^{1/3}(0,T)}$, $\|y_x(\cdot, x)\|_{L^2(0,T)}$, and $\|y_{xx}(\cdot, x)\|_{[H^{1/3}(0,T)]^*}$, we obtain the one for $\|y\|_{X_T}$.

The proof in the case $(h_1, h_2, h_3) = (h_1, 0, 0)$ or in the case $(h_1, h_2, h_3) = (0, h_2, 0)$ is similar after noting Lemma A.1 in the appendix. We mention here that the solution corresponding to the triple $(h_1, 0, 0)$ is given by

$$\hat{y}(z, x) = \frac{\hat{h}_1(z)}{\det Q} \sum_{j=1}^3 (\lambda_{j+2} e^{-\lambda_j L} - \lambda_{j+1} e^{-\lambda_j L}) e^{\lambda_j x} \text{ for a.e. } x \in (0, L),$$

and the solution corresponding to the triple $(0, h_2, 0)$ is given by

$$\hat{y}(z, x) = \frac{\hat{h}_2(z)}{\det Q} \sum_{j=1}^3 (\lambda_{j+1} e^{\lambda_{j+1} L} - \lambda_{j+2} e^{\lambda_{j+2} L}) e^{\lambda_j x} \text{ for a.e. } x \in (0, L).$$

The details are left to the reader.

Remark 3.4. If $h_1, h_2, h_3 \in C_c^\infty(\mathbb{R})$, the constructed solution is also smooth.

Step 2: We now deal with the general case. The starting point of the proof is a connection between the linearized KdV equation and the linear KdV-Burgers equation. Set $v(t, x) = e^{-2t+x} y(t, x)$, which is equivalent to $y(t, x) = e^{2t-x} v(t, x)$. One can check that if y satisfies the equation

$$y_t + y_x + y_{xxx} = f \text{ in } \mathbb{R}_+ \times (0, L),$$

if and only if

$$v_t + 4v_x + v_{xxx} - 3v_{xx} = f e^{-2t+x} \text{ in } \mathbb{R}_+ \times (0, L).$$

Set, in $\mathbb{R}_+ \times (0, L)$,

$$(3.25) \quad \psi(t, x) = \psi(t) := \frac{1}{L} \int_0^L f(t, \xi) e^{-2t+\xi} d\xi \quad \text{and} \quad g(t, x) := f(t, x) e^{-2t+x} - \psi(t, x).$$

Then

$$\int_0^L g(t, x) dx = 0 \text{ for } t > 0.$$

Let $y_1 \in C([0, +\infty); L^2(0, L)) \cap L_{\text{loc}}^2([0, +\infty); H^1(0, L))$ be the unique solution periodic in space of the system

$$(3.26) \quad y_{1,t} + 4y_{1,x} + y_{1,xxx} - 3y_{1,xx} = g \text{ in } (0, +\infty) \times (0, L),$$

and

$$(3.27) \quad y_1(0, \cdot) = y_0 e^x \text{ in } (0, L).$$

⁶More precisely, one can take $y = y_3$ in $(0, T) \times (0, L)$.

Set

$$\alpha = \frac{1}{L} \int_0^L y_1(0, x) dx.$$

By Proposition 2.1, we have, for $x \in [0, L]$,

$$(3.28) \quad \|y_1(\cdot, x) - \alpha\|_{H^{1/3}(\mathbb{R}_+)} + \|y_{1,x}(\cdot, x)\|_{L^2(\mathbb{R}_+)} + \|y_{1,xx}(\cdot, x)\|_{[H^{1/3}(\mathbb{R}_+)]^*} \\ \leq C \left(\|f\|_{L^1(\mathbb{R}_+; L^2(0, L))} + \|y_0\|_{L^2(0, L)} \right)$$

and

$$(3.29) \quad \|y_1(\cdot, x) - \alpha\|_{L^2(\mathbb{R}_+)} + \|y_{1,x}(\cdot, x)\|_{[H^{1/3}(\mathbb{R}_+)]^*} \leq C \left(\|y_0\|_{[H^1(0, L)]^*} + \|f\|_{L^1(\mathbb{R}_+; [H^1(0, L)]^*)} \right).$$

Let $y_2 \in X_T$ be the unique solution of

$$(3.30) \quad \begin{cases} y_{2,t} + y_{2,x} + y_{2,xxx} = 0 & \text{in } (0, T) \times (0, L), \\ y_2(t, 0) = h_1(t) - e^{2t} \left(y_1(t, 0) + \int_0^t \psi(s) ds \right) & \text{in } (0, T), \\ y_2(t, L) = h_2(t) - e^{2t-L} \left(y_1(t, L) + \int_0^t \psi(s) ds \right) & \text{in } (0, T), \\ y_{2,x}(t, L) = h_3(t) - \left(e^{2t-x} (y_1(t, \cdot) + \int_0^t \psi(s) ds) \right)_x(t, L) & \text{in } (0, T), \\ y_2(t = 0, \cdot) = 0 & \text{in } (0, L). \end{cases}$$

Applying the results of Step 1 to y_2 , and using (3.28) and (3.28), we derive that, for $x \in [0, L]$,

$$\|y_2\|_{X_T} + \|y_2(\cdot, x)\|_{H^{1/3}(0, T)} + \|y_{2,x}(\cdot, x)\|_{L^2(0, T)} + \|y_{2,xx}(\cdot, x)\|_{[H^{1/3}(0, T)]^*} \\ \leq C_T \left(\|(h_1, h_2)\|_{H^{1/3}(\mathbb{R}_+)} + \|h_3\|_{L^2(\mathbb{R}_+)} + \|f\|_{L^1([0, T]; L^2(0, L))} \right)$$

and

$$\|y_2(\cdot, x)\|_{L^2(0, T)} + \|y_{2,x}(\cdot, x)\|_{[H^{1/3}(0, T)]^*} \\ \leq C_T \left(\|(h_1, h_2)\|_{L^2(\mathbb{R}_+)} + \|h_3\|_{[H^{1/3}(\mathbb{R}_+)]^*} + \|f\|_{L^1([0, T]; [H^1(0, L)]^*)} \right).$$

The conclusion follows by noting that $y = e^{2t-x} \left(y_1 + \int_0^t \psi(s) ds \right) + y_2$ in $(0, T) \times (0, L)$.

The proof of the existence of one solution and its estimates is complete if one can show that using this process, one can construct a solution which verifies Definition 3.1. To this end, one first notes that if h_1, h_2, h_3, y_0, g are smooth and $\text{supp } h_1, \text{supp } h_2, \text{supp } h_3 \Subset (0, T)$, $\text{supp } y_0 \Subset (0, L)$ and $\text{supp } g \Subset (0, T) \times (0, L)$ then one can construct a solution with the desired estimates. Indeed, the solution y_1 given by (3.26) is smooth and has the property that $y_1(t, 0)$ and $y_1(t, L)$ are 0 for t close to 0 (see Remark 2.2 and Remark 2.3). One can then construct a smooth solution y_2 of (3.30) in the interval $(0, 2T)$ where the the expression of $y_2(t, 0)$, $y_2(t, L)$, and $y_{2,x}(t, L)$ are replaced by the expressions given in (3.30) multiplied by a cutoff function χ , which is 1 in $(0, T)$ and is 0 for $t > 3T/2$. The solution y_2 is also smooth (see Remark 3.4). Then $e^{2t-x} \left(y_1 + \int_0^t \psi(s) ds \right) + y_2$ is still a solution in the time interval $(0, T)$ with the desired estimates. Using this observation, by standard approximation arguments, one can construct a solution in the sense of Definition 3.1 with the required estimates. \square

In the proof of Proposition 3.1, we used the following elementary result.

Lemma 3.1. For $p \in \mathbb{C}$ and z in a sufficiently small conic neighborhood of \mathbb{R}_+ , let λ_j with $j = 1, 2, 3$ be the three solutions of the equation $\lambda^3 + \lambda + iz = 0$. Consider the convention $\Re(\lambda_1) < \Re(\lambda_2) < \Re(\lambda_3)$. We have, in the limit $|z| \rightarrow \infty$,

$$(3.31) \quad \lambda_j = \mu_j z^{1/3} - \frac{1}{3\mu_j} z^{-1/3} + O(z^{-2/3}) \quad \text{with } \mu_j = e^{-i\pi/6 - 2j\pi/3},$$

Here $z^{1/3}$ denotes the cube root of z with the real part positive.

Here and in what follows, for $s \in \mathbb{R}$, $O(z^s)$ denotes a quantity bounded by Cz^s for large positive z . Similar convention is used for $O(|z|^s)$ for $z \in \mathbb{C}$.

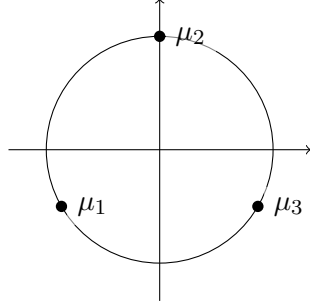


FIGURE 1. The roots λ_j of $\lambda^3 + \lambda + iz = 0$ satisfy, when $z > 0$ is large, $\lambda_j \sim \mu_j z^{1/3}$ where $\mu_j^3 = -i$.

The same proof as in Step 1 of the proof of Proposition 3.1 gives the following result.

Lemma 3.2. Let $T > 0$ and $(h_1, h_2, h_3) \in H^{1/3}(0, T) \times H^{1/3}(0, T) \times L^2(0, T)$. Then the unique solution $y \in X_T$ of the system (3.1) with $y_0 = 0$ and $f = 0$ satisfy, for $x \in [0, L]$,

$$(3.32) \quad \|y(\cdot, x)\|_{[H^{2/3}(0, T)]^*} \leq C_{T, L} \left(\|(h_1, h_2)\|_{[H^{2/3}(0, T)]^*} + \|h_3\|_{[H^1(0, T)]^*} \right),$$

where $C_{T, L}$ denotes a positive constant independent of x, y_0 , and h_1, h_2, h_3 .

3.1.2. Uniqueness. In this section, we establish the uniqueness of the solutions given in Proposition 3.1. The uniqueness is somehow known for a close definition of the solution which requires φ less regular. The definition considered here is somehow more advantageous, see Remark 3.5. We shall highlight the analysis in such a way that one can extend it to other settings considered in this paper; thus we do not need to repeat the arguments as much as possible. Let $y \in X_T$ be a solution with the zero data, i.e., $f = 0, y_0 = 0$, and $(h_1, h_2, h_3) = (0, 0, 0)$. Fix $\psi \in C_c^\infty((0, T) \times (0, L))$ (arbitrarily). Let $\tilde{y} \in X_T$ be a solution of the backward system

$$\begin{cases} \tilde{y}_t + \tilde{y}_x + y_{xxx} = \psi & \text{in } (0, T) \times (0, L), \\ \tilde{y}(\cdot, 0) = 0, \tilde{y}_x(\cdot, 0) = 0, \tilde{y}(\cdot, L) = 0 & \text{in } (0, T), \\ \tilde{y}(T, \cdot) = 0 & \text{in } (0, L). \end{cases}$$

Using the construction given in Section 3.1.1, one can assume that \tilde{y} is smooth (see the last part of Step 2 of Section 3.1.1). Taking $\varphi = \tilde{y}$ in the definition, we derive that

$$\int_0^T \int_0^L \psi(t, x) y(t, x) dt dx = 0.$$

Since $\psi \in C_c^\infty((0, T) \times (0, L))$ is arbitrary, we deduce that

$$y = 0 \text{ in } (0, T) \times (0, L).$$

The uniqueness is proved. \square

Remark 3.5. The proof of the uniqueness also gives the uniqueness of the solutions in $L^1((0, T) \times (0, L))$, i.e., one requires $y \in L^1((0, T) \times (0, L))$ instead of $y \in X_T$ in Definition 3.1.

Remark 3.6. We end this section with some comments on Proposition 3.1 and its proof.

i) The well-posedness of (3.1) in X_T is proved in [10, Theorem 2.10 and Proposition 2.16] for $L = 1$ when $(0, T) = \mathbb{R}_+$ and the estimate for $\|y_x(\cdot, x)\|_{L^2(0, T)}$ in this case is a consequence of their results. For $L = 1$, their results imply that, when $(h_1, h_2, h_3) \equiv (0, 0, 0)$,

$$(3.33) \quad \|y_x(\cdot, x)\|_{L^2(\mathbb{R}_+)} \leq C\|y_0\|_{L^2(0, 1)} \text{ for } x \in [0, 1].$$

Note that $L = 1$ is smaller than the smallest critical length in \mathcal{N}_N which is 2π . The estimate as (3.33) cannot hold for arbitrary L . Such an estimate is not valid for any critical length by considering a non-zero initial datum in \mathcal{M}_N .

ii) Similar estimates for $\|y(\cdot, x)\|_{H^{1/3}(0, T)}$ and $\|y_x(\cdot, x)\|_{L^2(0, T)}$ as in (3.2) in the real line space setting can be found in [20, 28]. To our knowledge, variants of the estimate for $\|y_{xx}(\cdot, x)\|_{[H^{1/3}(0, T)]^*}$ in (3.2) are not known even in the real line space setting. Our proof of Proposition 3.1 is in the spirit of [24], which involves the Fourier transform with respect to time of the solution, as in [10], and a connection between the linearized KdV and the linear KdV-Burger equations. However, in the study of the linear KdV-Burger equation with periodic boundary conditions, the singularity of the kernel is appropriately *compensated* (see the proof of Proposition 2.1), which is one of the novelties of the analysis. The proof given here is self-contained and is different from the ones in [20, 28] for the whole line setting, which are based on the Riemann-Liouville fractional integrals and the theory of Airy functions.

3.2. Proof of Proposition 3.2. We first give the definition of the solutions considered in Proposition 3.2.

Definition 3.2. Let $L > 0$, $T > 0$, $(h_1, h_2, h_3) \in H^{1/3}(0, T) \times H^{1/3}(0, T) \times [H^{1/3}(0, T)]^*$, $f \in L^1((0, T); L^2(0, L))$, and $y_0 \in L^2(0, L)$. A solution $y \in X_T$ of the system (3.4) is a function $y \in X_T$ such that

$$(3.34) \quad \int_0^T \int_0^L f(t, x)\varphi(t, x) + \int_0^L y_0(x)\varphi(0, x) dx - \int_0^T h_3(t)\varphi(t, L) dt \\ - \int_0^T h_2(t)\varphi_{xx}(t, L) dt + \int_0^T h_1(t)\varphi_{xx}(t, 0) dt = - \int_0^T \int_0^L y(\varphi_t + \varphi_x + \varphi_{xxx}) dx dt$$

for all $\varphi \in C^3([0, T] \times [0, L])$ with $\varphi(T, \cdot) = 0$, $\varphi(\cdot, 0) = \varphi_x(\cdot, 0) = \varphi_x(\cdot, L) = 0$.

The proof of Proposition 3.1 is similar to the one of Proposition 3.1. The details are left to the reader. We just present here the formula of the solution of the system when $y_0 = 0$, $h_1 = 0$, $h_2 = 0$, and $f = 0$ and discuss the uniqueness.

Extend h_3 by 0 for $t \notin [0, T]$ and still denote this extension by h_3 . Then $h_3 \in H^{-1/3}(\mathbb{R})$. Denote y be the corresponding solution for $t \geq 0$ and extend y by 0 for $t < 0$. Still denote this extension by y . One then has

$$(3.35) \quad \begin{cases} y_t + y_x + y_{xxx} = 0 & \text{in } \mathbb{R} \times (0, L), \\ y(\cdot, 0) = 0, y(\cdot, L) = 0, y_{xx}(\cdot, L) = h_3 & \text{in } \mathbb{R}. \end{cases}$$

Taking the Fourier transform with respect to t , from the system of y , we have

$$(3.36) \quad \begin{cases} iz\hat{y}(z, x) + \hat{y}_x(z, x) + \hat{y}_{xxx}(z, x) = 0 & \text{in } \mathbb{R} \times (0, L), \\ \hat{y}(z, 0) = \hat{y}(z, L) = 0 & \text{in } \mathbb{R}, \\ \hat{y}_{xx}(z, L) = \hat{h}_3(z) & \text{in } \mathbb{R}. \end{cases}$$

Taking into account the equation of \hat{y} , we search for the solution of the form

$$(3.37) \quad \hat{y}(z, \cdot) = \sum_{j=1}^3 a_j e^{\lambda_j x},$$

where $\lambda_j = \lambda_j(z)$ with $j = 1, 2, 3$ are the three solutions of the equation $\lambda^3 + \lambda + iz = 0$, and $a_j = a_j(z)$ for $j = 1, 2, 3$. We then have

$$\begin{aligned} \sum_{j=1}^3 a_j &= 0, \\ \sum_{j=1}^3 a_j e^{\lambda_j L} &= 0, \\ \sum_{j=1}^3 a_j \lambda_j^2 e^{\lambda_j L} &= \hat{h}_3, \end{aligned}$$

This implies, with the convention $\lambda_{j+3} = \lambda_j$,

$$(3.38) \quad a_j = \frac{e^{\lambda_{j+2}L} - e^{\lambda_{j+1}L}}{\sum_{k=1}^3 e^{-\lambda_k L} (\lambda_{k+2}^2 - \lambda_{k+1}^2)} \hat{h}_3 \text{ for } j = 1, 2, 3.$$

We now deal with the uniqueness. Let $y \in X_T$ be a solution with the zero data, i.e., $f = 0$, $y_0 = 0$, and $(h_1, h_2, h_3) = (0, 0, 0)$. Let $\psi \in C_c^\infty((0, T) \times (0, L))$. Let $\tilde{y} \in X_T$ be a solution of the backward system

$$\begin{cases} \tilde{y}_t + \tilde{y}_x + \tilde{y}_{xxx} = \psi & \text{in } (0, T) \times (0, L), \\ \tilde{y}(\cdot, 0) = 0, \tilde{y}_x(\cdot, 0) = 0, \tilde{y}_x(\cdot, L) = 0 & \text{in } (0, T), \\ \tilde{y}(T, \cdot) = 0 & \text{in } (0, L). \end{cases}$$

One can show that \tilde{y} is smooth (see the construction of the solutions given in Proposition 3.3 and the last part of Step 2 of Section 3.1.1; in fact, the arguments are simpler since the data on the boundary are 0 and the initial datum is 0). Taking $\varphi = \tilde{y}$ in Definition 3.1, we derive that

$$\int_0^T \int_0^L \psi(t, x) y(t, x) dt dx = 0.$$

Since $\psi \in C_c^\infty((0, T) \times (0, L))$ is arbitrary, we deduce that

$$y = 0 \text{ in } (0, T) \times (0, L).$$

The uniqueness is proved. \square

Remark 3.7. The proof of the uniqueness also gives the uniqueness of the solutions in $L^1((0, T) \times (0, L))$, i.e., one requires $y \in L^1((0, T) \times (0, L))$ instead of $y \in X_T$ in Definition 3.2.

3.3. Proof of Proposition 3.3. We first introduce the notion of the solutions considered in Proposition 3.3.

Definition 3.3. Let $L > 0$, $T > 0$, $(h_1, h_2, h_3) \in L^2(0, T) \times H^{1/3}(0, T) \times L^2(0, T)$, $f \in L^1((0, T); L^2(0, L))$, and $y_0 \in L^2(0, L)$. A solution $y \in X_T$ of the system (3.7) is a function $y \in X_T$ such that

$$(3.39) \quad \begin{aligned} & \int_0^T \int_0^L f(t, x) \varphi(t, x) + \int_0^L y_0(x) \varphi(0, x) dx + \int_0^T h_3(t) \varphi_x(t, L) dt \\ & - \int_0^T h_1(t) \varphi_x(t, 0) dt - \int_0^T h_2(t) \varphi_{xx}(t, L) dt = - \int_0^T \int_0^L y(\varphi_t + \varphi_x + \varphi_{xxx}) dx dt \end{aligned}$$

for all $\varphi \in C^3([0, T] \times [0, L])$ with $\varphi(T, \cdot) = 0$, $\varphi(\cdot, 0) = \varphi(\cdot, L) = \varphi_{xx}(\cdot, 0) = 0$.

The proof of Proposition 3.3 is similar to the one of Proposition 3.1. We only give here the formula of the solutions in the case $f = 0$, $y_0 = 0$, and $(h_1, h_2, h_3) = (0, 0, h_3)$. Taking the Fourier transform with respect to t , from the system of y , we have

$$(3.40) \quad \begin{cases} iz\hat{y}(z, x) + \hat{y}_x(z, x) + \hat{y}_{xxx}(z, x) = 0 & \text{in } \mathbb{R} \times (0, L), \\ \hat{y}_x(z, 0) = \hat{y}(z, L) = 0 & \text{in } \mathbb{R}, \\ \hat{y}_x(z, L) = \hat{h}_3(z) & \text{in } \mathbb{R}. \end{cases}$$

Taking into account the equation of \hat{y} , we search for the solution of the form

$$(3.41) \quad \hat{y}(z, \cdot) = \sum_{j=1}^3 a_j e^{\lambda_j x},$$

where $\lambda_j = \lambda_j(z)$ with $j = 1, 2, 3$ are the three solutions of the equation $\lambda^3 + \lambda + iz = 0$, and $a_j = a_j(z)$ for $j = 1, 2, 3$. We then have

$$\begin{aligned} \sum_{j=1}^3 \lambda_j a_j &= 0, \\ \sum_{j=1}^3 a_j e^{\lambda_j L} &= 0, \\ \sum_{j=1}^3 a_j \lambda_j e^{\lambda_j L} &= \hat{h}_3, \end{aligned}$$

This implies, with the convention $\lambda_{j+3} = \lambda_j$,

$$(3.42) \quad a_j = \frac{\lambda_{j+1} e^{\lambda_{j+2} L} - \lambda_{j+2} e^{\lambda_{j+1} L}}{\sum_{k=1}^3 \lambda_k e^{\lambda_k L} (\lambda_{k+1} e^{\lambda_{k+2} L} - \lambda_{k+2} e^{\lambda_{k+1} L})} \hat{h}_3 \text{ for } j = 1, 2, 3.$$

The rest of the proof is almost the same as the one of Proposition 3.1 and is left to the reader. \square

4. THE SET OF THE CRITICAL LENGTHS

The main goal of this section is to establish Theorem 1.1, except for (1.12) and its consequence (1.13), and Proposition 1.1. The results in this section, in particular in (4.1) will be complemented with the analysis in the next section (Section 5) to derive (1.12) and (1.13). We begin with

Proposition 4.1. *Assume that $L \in \mathcal{N}_D$. There exists $a < 0$ and $b > 0$ such that*

$$(4.1) \quad (a + ib)e^{(a+ib)} = (a - ib)e^{(a-ib)} = -2ae^{-2a}$$

and

$$(4.2) \quad (a + ib)^2 + (a - ib)^2 + (a + ib)(a - ib) = -L^2.$$

These conditions are equivalent to

$$(4.3) \quad b^2 = L^2 + 3a^2, \quad a^2 + b^2 = 4a^2 e^{-6a}, \quad b \cos b + a \sin b = 0,$$

$$(4.4) \quad a \cos b - b \sin b > 0.$$

Moreover,

$$(4.5) \quad a \text{ and } b \text{ are uniquely determined by } L.$$

Proof. The proof is divided into two steps.

Step 1: We prove (4.1) - (4.4).

Let $z_1 = a + ib$ and $z_2 = c + id$ where $a, b, c, d \in \mathbb{R}$. Set

$$z_3 = -(z_1 + z_2) = -(a + c) - i(b + d).$$

Note that, if $z_1 + z_2 + z_3 = 0$, then

$$(4.6) \quad z_1^2 + z_2^2 + z_3^2 + z_1z_2 = -(z_1z_2 + z_2z_3 + z_3z_1).$$

Condition (1.4) can be then written as

$$(4.7) \quad z_1 + z_2 + z_3 = 0, \quad z_1e^{z_1} = z_2e^{z_2} = z_3e^{z_3}, \quad \text{and} \quad z_1z_2 + z_2z_3 + z_3z_1 = L^2.$$

In the three numbers a , c , and $-(a+c)$, there are two numbers of the same sign. Note also that if (z_1, z_2, z_3) is a solution of (4.7), then $(\bar{z}_1, \bar{z}_2, \bar{z}_3)$ is also a solution. Without loss of generality, one might assume from now on that

$$(4.8) \quad ac \geq 0, \quad |c| \geq |a|,$$

and

$$(4.9) \quad b \geq 0.$$

From (1.4), we have

$$(a+ib)^2 + (c+id)^2 + (a+ib)(c+id) = -L^2.$$

This yields

$$a^2 - b^2 + c^2 - d^2 + ac - bd + i(2ab + 2cd + ad + bc) = -L^2.$$

We thus obtain

$$(4.10) \quad a^2 + ac + c^2 + L^2 = b^2 + bd + d^2$$

and

$$(4.11) \quad 2ab + 2cd + ad + bc = 0.$$

From (1.4), we have

$$(4.12) \quad (a+ib)e^{a+ib} = (c+id)e^{c+id} = -(a+c+i(b+d))e^{-(a+c)-i(b+d)}.$$

Considering the modulus in (4.12), we get

$$(4.13) \quad (a^2 + b^2)e^{2a} = (c^2 + d^2)e^{2c} = ((a+c)^2 + (b+d)^2)e^{-2(a+c)}.$$

Considering the imaginary parts in (4.12), we obtain

$$(4.14) \quad (b \cos b + a \sin b)e^a = (d \cos d + c \sin d)e^c = (-(b+d) \cos(b+d) + (a+c) \sin(b+d))e^{-(a+c)}.$$

This implies that $c \neq 0$ since otherwise a is also 0 (since $|a| \leq |c|$) and one reaches, by (4.13),

$$|b| = |d| = |b+d|,$$

which yields $b = d = 0$ and $L = 0$. This is impossible since $L > 0$.

Using (4.11) and the facts $ac \geq 0$ and $c \neq 0$, we have

$$(4.15) \quad d = -\frac{2a+c}{2c+a}b.$$

This implies

$$(4.16) \quad b+d = \frac{c-a}{2c+a}b.$$

Combining (4.9) and (4.15) yields that

$$(4.17) \quad b > 0$$

since $b^2 + bd + d^2 \geq L^2$.

The proof now is divided into four cases.

- Case 1: $a \geq 0$. Thus $c \geq a \geq 0$. Since $(a+c)^2 e^{-2(a+c)} \leq e^{-2}$ for $a+c \geq 0$, it follows that

$$\begin{aligned} (b+d)^2 + e^{-2} &\geq (b+d)^2 + (a+c)^2 e^{-2(a+c)} \\ &\geq ((a+c)^2 + (b+d)^2) e^{-2(a+c)} \stackrel{(4.13)}{=} (a^2 + b^2) e^{2a} \geq b^2 e^{2a} \geq b^2. \end{aligned}$$

In summary, one reaches

$$(4.18) \quad (b+d)^2 + e^{-2} \geq b^2.$$

One has, for $c \geq a \geq 0$,

$$(4.19) \quad (b+d)^2 \stackrel{(4.16)}{=} \left(\frac{c-a}{2c+a} \right)^2 b^2 \leq b^2/4.$$

Combining (4.18) and (4.19) yields

$$(4.20) \quad e^{-2} \geq \frac{3}{4} b^2; \text{ thus, by (4.17), } 0 < b \leq 0.425.$$

Since $a \geq 0$, it follows from (4.20) that

$$(4.21) \quad b \cos b + a \sin b > 0.$$

We have

$$\begin{aligned} a+c &\leq (2a^2 + 2c^2 + 2ac)^{1/2} \stackrel{(4.10)}{\leq} (2(b^2 + d^2 + bd))^{1/2} = (b^2 + d^2 + (b+d)^2)^{1/2} \\ &\stackrel{(4.15),(4.19)}{\leq} (2b^2 + b^2/4)^{1/2} = (9b^2/4)^{1/2} \stackrel{(4.20)}{\leq} (3e^{-2})^{1/2} \leq 0.638. \end{aligned}$$

Since $0 \leq b+d \leq b/2 \leq 0.225$ by (4.16), (4.19), and (4.20), it follows that

$$(4.22) \quad \begin{aligned} -(b+d) \cos(b+d) + (a+c) \sin(b+d) &\leq -(b+d) \cos(0.225) + (a+c)(b+d) \\ &\leq -0.97(b+d) + 0.638(b+d) \leq 0. \end{aligned}$$

Combining (4.14), (4.21), and (4.22), we derive that the case $a \geq 0$ does not happen.

- Case 2: $a < 0$ and $a \neq c$. Thus $c < a < 0$ since $|c| \geq |a|$.

* Case 2.1: $a < 0$ and $c \leq 2a$. We have

$$d^2 \stackrel{(4.15)}{=} \left(\frac{2a+c}{2c+a} \right)^2 b^2 \leq (4/5)^2 b^2.$$

It follows that

$$b^2 \leq (a^2 + b^2) \stackrel{(4.13)}{=} e^{2(c-a)} (c^2 + d^2) \leq c^2 e^c + (4/5)^2 b^2.$$

This yields, since $c \leq 0$,

$$(9/25)b^2 \leq c^2 e^c \leq 2^2 e^{-2}.$$

We obtain

$$(4.23) \quad b^2 \leq \frac{100}{9} e^{-2} \leq 1.51, \text{ which yields, by (4.17), } 0 < b \leq 1.23.$$

We have

$$(4.24) \quad a^2 + b^2 \stackrel{(4.13)}{=} ((a+c)^2 + (b+d)^2) e^{-4a-2c} \geq (a+c)^2 e^{-2a-2c}$$

since $a < 0$. It follows from (4.23) that

$$1.51 \geq b^2 \geq (a+c)^2 e^{-2a-2c} - a^2 \geq (a+c)^2 \left(e^{-2a-2c} - \frac{1}{9} \right).$$

Since $f_1(t) = t^2(e^{2t} - 1/9)$ is an strictly increasing function for $t > 0$ and

$$f_1(0.66) \geq 1.55,$$

it follows that

$$(4.25) \quad 0.66 \geq |a + c| = |c| + |a| \geq 3|a|.$$

Using (4.23) and (4.25), we obtain

$$(4.26) \quad b \cos b + a \sin b \geq b \cos 1.23 + ab > 0.33b + ab > 0.$$

and, since $0 < b + d \leq b/2 \leq 0.64$ by (4.16) and (4.23),

$$(4.27) \quad - (b + d) \cos(b + d) + (a + c) \sin(b + d) \\ \leq - (b + d) \cos 0.64 + (a + c)(b + d) \cos 0.64 \stackrel{(4.25)}{<} 0.$$

Here we used the fact $\sin x \geq x \cos \theta_0$ for $0 \leq x \leq \theta_0 \leq 1$.

Combining (4.14), (4.26), and (4.27), we derive that Case 2.1 does not happen.

* Case 2.2: $a < 0$, $2a \leq c < a$ with $|a| \leq 1$. We have

$$a^2 + b^2 \stackrel{(4.13)}{=} (c^2 + d^2)e^{2(c-a)} \stackrel{(4.15)}{\leq} (c^2 + b^2)e^{2(c-a)}.$$

This implies, with $s = c - a < 0$,

$$b^2(1 - e^{2s}) = b^2(1 - e^{2(c-a)}) \leq c^2 e^{2(c-a)} - a^2 \\ = (a + s)^2 e^{2s} - a^2 = a^2 e^{2s} + 2ase^{2s} + s^2 e^{2s} - a^2 \leq (2as + s^2)e^{2s}.$$

We derive that

$$(4.28) \quad b^2 \leq \frac{|2as + s^2|}{|1 - e^{2s}|} e^{2s}.$$

Set

$$g(t) = 2te^{-2t} - 1 + e^{-2t} \text{ for } t \geq 0.$$

Then

$$g'(t) = (2 - 4t - 2)e^{-2t} = -4te^{-2t} < 0 \text{ for } t > 0,$$

which yields $g(t) < g(0) = 0$ for $t > 0$. This implies

$$(4.29) \quad |s|e^{2s} \leq \frac{1}{2}(1 - e^{2s}).$$

Combining (4.28) and (4.29) yields

$$(4.30) \quad b^2 \leq \frac{1}{2}(2|a| + |s|) = \frac{1}{2}(|a| + |c|) \leq 3/2,$$

where in the last inequality we used the fact $|c| \leq 2|a|$ and $|a| \leq 1$.

We have

$$(4.31) \quad a^2 + b^2 \stackrel{(4.13)}{=} ((a + c)^2 + (b + d)^2)e^{-4a-2c} \geq (a + c)^2 e^{-2a-2c}.$$

It follows from (4.30) that

$$(4.32) \quad \frac{1}{2}(|a| + |c|) \geq b^2 \geq (a + c)^2 e^{-2a-2c} - a^2 \geq (a + c)^2 (e^{-2a-2c} - \frac{1}{4}), \\ \frac{1}{2} \geq |a + c|(e^{-2a-2c} - \frac{1}{4}),$$

Since $f_2(t) = t(e^{2t} - 1/4)$ is an increasing function for $t \geq 0$ and

$$f_2(0.32) \geq 0.52,$$

it follows from (4.32) that

$$(4.33) \quad |a| + |c| \leq 0.32, \text{ which yields } |a| \leq 0.16.$$

Using (4.30) and (4.32), we obtain

$$(4.34) \quad b \cos b + a \sin b \geq b \cos \sqrt{3/2} + ab > 0.33b + ab > 0,$$

and, since $0 < b + d \leq b/2 \leq \sqrt{3/2}/2 \leq 0.62$ by (4.16) and (4.30),

$$(4.35) \quad - (b + d) \cos(b + d) + (a + c) \sin(b + d) \\ \leq - (b + d) \cos 0.62 + (a + c)(b + d) \cos 0.62 \stackrel{(4.33)}{<} 0.$$

Here we used the fact $\sin x \geq x \cos \theta_0$ for $0 \leq x \leq \theta_0 \leq 1$.

Combining (4.14), (4.34), and (4.35), we derive that Case 2.2 does not happen.

* Case 2.3: $a < 0$ and $2a \leq c < a$ with $|a| > 1$. We have, by (4.13),

$$(a^2 + b^2)e^{2a} = (c^2 + d^2)e^{2c}$$

Since $|c| \geq |a| \geq 1$ and $d^2 \leq b^2$ by (4.15), and the function $f_3(t) = t^2 e^{-2t}$ is strictly decreasing in $[1, +\infty)$, we derive that

$$(a^2 + b^2)e^{2a} > (c^2 + d^2)e^{2c}.$$

Thus Case 2.3 does not happen.

• Case 3: $a < 0$ and $a = c$. We derive from (4.16) that

$$b + d = 0.$$

Assertion (4.1)-(4.4) then follow from (4.13), (4.14) and the fact that

$$\Re(a + ib)e^{(a+ib)} = e^a(a \cos b - b \sin b) > 0.$$

• Case 4: $a = 0$ and $c < 0$. It follows that

$$d = -b/2.$$

This case does not happen by (4.13).

Step 2: We prove (4.5). Since

$$L^2 = b^2 - 3a^2 = 4a^2(e^{-6a} - 1)$$

and the function $h(t) = 4t^2(e^{6t} - 1)$ is strictly increasing in $(0, +\infty)$, it follows that

$$a \text{ is uniquely determined by } L.$$

This in turn implies that

$$b \text{ is uniquely determined by } L$$

since $L^2 = b^2 - 3a^2$ and $b > 0$.

The proof is complete. □

We are ready to give

Proof of Proposition 1.1. By (4.3) and (4.4), we have

$$L^2 = 4a^2(e^{-6a} - 1),$$

where $a < 0$ is determined by, with $b^2 = 4a^2(e^{-6a} - 1/4)$, $b > 0$,

$$b \cos b + a \sin b = 0 \quad \text{and} \quad a \cos b - b \sin b > 0.$$

Since $b > 0$ and $a < 0$, it follows that

$$\cos b < 0 \quad \text{and} \quad \sin b < 0.$$

Hence $b = b_n$ for some $\pi + 2n\pi < b_n < 3\pi/2 + n2\pi$ with $n \in \mathbb{N}$. Since $b_n \geq \pi$ and $4a_n^2(e^{-6a_n} - 1/4) = b_n^2$, it follows that

$$|a_n| \geq 0.43.$$

This yields

$$(4.36) \quad |\cot b_n| = |a_n/b_n| = \frac{1}{4|a_n|(e^{6|a_n|} - 1/4)} \leq 0.045.$$

We thus derive that

$$|b_n - 3\pi/2| \leq 0.045.$$

Set

$$t_n = b_n - \pi - n2\pi.$$

Since

$$|b_n \cos b_n| = |a_n \sin b_n|,$$

it follows that

$$(4.37) \quad (t_n + \pi + n2\pi) \cos t_n = |a_n| \sin t_n.$$

For $t_n \in (\pi/4, \pi/2)$, the LHS of (4.37) is a strictly decreasing function of t_n and the RHS of (4.37) is a strictly increasing function of t_n since $|a_n|$ is a increasing function of t_n . Hence there exists a unique $t_n \in (\pi/4, \pi/2)$ such that the identity holds. One can check that the corresponding a_n fulfills all the requirements. \square

Remark 4.1. It is clear, as $n \rightarrow +\infty$, that $a_n/b_n \rightarrow 0$ so $b_n - 3\pi/2 - 2\pi n \rightarrow 0$ as $n \rightarrow +\infty$.

We end this section with the following result, which yields (1.8) and (1.11). Recall that α and β are defined by (1.9) and q is defined by (1.7).

Lemma 4.1. *Let $L \in \mathcal{N}_D$. Set*

$$(4.38) \quad \eta_1 = \alpha + i\beta, \quad \eta_2 = \alpha - i\beta, \quad \text{and} \quad \eta_3 = -2\alpha,$$

$$(4.39) \quad \varphi(x) = (\eta_3 - \eta_2)e^{\eta_1 x} + (\eta_1 - \eta_3)e^{\eta_2 x} + (\eta_2 - \eta_1)e^{\eta_3 x} \text{ in } [0, L],$$

and

$$\Psi(t, x) = e^{qt}\psi(x) \text{ in } \mathbb{R} \times [0, L].$$

Then $\varphi(x) = -2i\phi$ and Ψ satisfies (1.11). Consequently, (1.8) and (1.11) hold.

Proof. Since

$$\eta_1 e^{-\eta_1 L} = \eta_2 e^{-\eta_2 L} = \eta_3 e^{-\eta_3 L}$$

$$\eta_1 + \eta_2 + \eta_3 = 0, \quad \eta_1 \eta_2 + \eta_2 \eta_3 + \eta_3 \eta_1 = 1, \quad \text{and} \quad \eta_1 \eta_2 \eta_3 = -q,$$

it follows that

$$\begin{cases} \varphi_{xxx} + \varphi_x + q\varphi = 0 & \text{in } (0, L), \\ \varphi(0) = \varphi_x(0) = \varphi(L) = \varphi_{xx}(L) = 0 & \text{in } [0, L]. \end{cases}$$

We derive that Ψ is a solution of (1.11). A computation gives

$$\begin{aligned}\varphi(x) &= (-3\alpha + i\beta)e^{(\alpha+i\beta)x} + (3\alpha + i\beta)e^{(\alpha-i\beta)x} - 2i\beta e^{-2\alpha x} \\ &= -6i\alpha e^{\alpha x} \sin(\beta x) + 2i\beta e^{\alpha x} \cos(\beta x) - 2i\beta e^{-2\alpha x} \\ &= 2i\left(\beta e^{\alpha x} \cos(\beta x) - \beta e^{-2\alpha x} - 3\alpha e^{\alpha x} \sin(\beta x)\right) = -2i\phi.\end{aligned}$$

The proof is complete. \square

Remark 4.2. In this section, we characterize *all* solutions of (1.4) and obtain the corresponding unreachable space \mathcal{M}_D for the corresponding linearized system. Previously, a subclass of solutions of (1.4), which corresponding to even and large n in our notations, are given in [27, The proof of proposition 4]. Our analysis is different from and much more involved than the one given in [27].

5. THE UNREACHABLE SPACE OF THE LINEARIZED SYSTEM FOR A CRITICAL LENGTH

In this section, we prove that the unreachable space is given by (1.12) and study its controllability properties of the linearized KdV system:

$$(5.1) \quad \begin{cases} y_t + y_x + y_{xxx} = 0 & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, T), \\ y(\cdot, L) = u & \text{in } (0, T), \end{cases}$$

for a critical length. The main result of this section is Proposition 5.1 which is based on an observability inequality with initial data and final data in \mathcal{M}_D^\perp , where \mathcal{M}_D is defined by (1.12). In comparison with the right Neumann boundary control system, this part in the Dirichlet setting is more complex and technical. The proof of Proposition 5.1 involves the well-posedness and the estimates for solutions in X_T given in Proposition 3.1 and Proposition 3.2.

We begin with the following simple but useful result.

Lemma 5.1. *Let $L \in \mathcal{N}_D$ and $T > 0$. For $y_0 \in L^2(0, L)$ and $u \in H^{1/3}(0, T)$, let $y \in X_T$ be the unique solution of the system*

$$\begin{cases} y_t + y_x + y_{xxx} = 0 & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, T), \\ y(\cdot, L) = u & \text{in } (0, T), \\ y(0, \cdot) = y_0 & \text{in } (0, L). \end{cases}$$

Then

$$\int_0^L y(t, x) \Phi(t, x) dx \text{ is constant on } [0, T].$$

Recall that Φ is defined by (1.10).

Proof. Using the fact, by integration by parts,

$$\frac{d}{dt} \int_0^L y(t, x) \Phi(t, x) dx = 0,$$

the conclusion follows. \square

We are ready to state the main result of this section on the unreachable space for the linearized KdV system (5.1).

Proposition 5.1. *Let $L \in \mathcal{N}_D$ and $T > 0$. We have*

i) for $\varphi \in \mathcal{M}_D \setminus \{0\}$, there does not exist $u \in H^{1/3}(0, T)$ such that $y(T, \cdot) = \varphi$ where $y \in X_T$ is the unique solution of (5.1) with $y(0, \cdot) = 0$.

ii) there exists a linear continuous operator $\mathcal{L} : \mathcal{M}_D^\perp \rightarrow H^{1/3}(0, T)$ such that $y(T, \cdot) = \varphi$ where $y \in X_T$ is the unique solution of (5.1) with $y(0, \cdot) = 0$ and $u = \mathcal{L}(\varphi)$.

iii) There exists a linear continuous operator $\hat{\mathcal{L}} : \mathcal{M}_D^\perp \rightarrow H^{1/3}(0, T)$ such that $y(T, \cdot) = 0$ where $y \in X_T$ is the unique solution of (5.1) with $y(0, \cdot) = \varphi$ and $u = \hat{\mathcal{L}}(\varphi)$.

Proof. Assertion *i)* is just a consequence of Lemma 5.1.

We next deal with assertions *ii)* and *iii)*. Let $\psi \in \mathcal{M}_D^\perp$. Let $y^* \in X_T$ be the unique solution of the backward linear KdV system

$$(5.2) \quad \begin{cases} y_t^* + y_x^* + y_{xxx}^* = 0 & \text{in } (0, T) \times (0, L), \\ y^*(\cdot, 0) = y_x^*(\cdot, 0) = y^*(\cdot, L) = 0 & \text{in } (0, T), \\ y^*(T, \cdot) = \psi & \text{in } (0, L). \end{cases}$$

Applying the observability inequality in Lemma 5.2 below to $y^*(T - \cdot, L - \cdot)$, we obtain

$$(5.3) \quad \lambda^{-1} \|\psi\|_{L^2(0, L)} \leq \|y_{xx}^*(\cdot, L)\|_{[H^{1/3}(0, T)]^*} \leq \lambda \|\psi\|_{L^2(0, L)}$$

for some constant $\lambda \geq 1$. Fix a continuous linear mapping

$$(5.4) \quad \mathcal{L}_1 : [H^{1/3}(0, T)]^* \rightarrow H^{1/3}(0, T)$$

such that

$$(5.5) \quad \langle \mathcal{L}_1(h), h \rangle_{H^{1/3}(0, T); [H^{1/3}(0, T)]^*} \geq C \|h\|_{[H^{1/3}(0, T)]^*}^2 \quad \text{for all } h \in [H^{1/3}(0, T)]^*,$$

for some positive constant C . This can be done using the Fourier series or the Fourier transform appropriately.

We first prove assertion *ii)*. Equip \mathcal{M}_D^\perp with the $L^2(0, L)$ -scalar product. Define

$$\mathcal{A} : \mathcal{M}_D^\perp \rightarrow \mathcal{M}_D^\perp$$

by

$$\mathcal{A}(\psi) = y(T, \cdot),$$

where $y \in X_T$ is the unique solution of the following system

$$(5.6) \quad \begin{cases} y_t + y_x + y_{xxx} = 0 & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, T), \\ y(\cdot, L) = h & \text{in } (0, T), \\ y(0, \cdot) = 0 & \text{in } (0, L), \end{cases}$$

with $h = \mathcal{L}_1(y_{xx}^*(\cdot, L))$ where y^* is determined by (5.2) ($y(T, \cdot) \in \mathcal{M}_D^\perp$ since $y(0, \cdot) = 0$ by Lemma 5.1). An integration by part yields

$$(5.7) \quad \int_0^L y(T, x) z^*(T, x) dx = - \int_0^T y(t, L) z_{xx}^*(t, L) dt,$$

for all solutions $z^* \in X_T$ of (5.2) with $z^*(T, \cdot) \in \mathcal{M}_D^\perp$.

Using (5.3) and applying the Lax-Milgram theory, we derive that \mathcal{A} is linear continuous and its inverse is also linear continuous. The conclusion of *ii)* follows by taking $\mathcal{L}(\psi) = \mathcal{L}_1(y_{xx}^*(\cdot, L))$ and y^* is the solution of (5.2) with ψ being replaced by $\mathcal{A}^{-1}(\psi)$.

We now deal with assertion *iii*). Set

$$H = \left\{ y_{xx}^*(\cdot, L) \in [H^{1/3}(0, T)]^*; \text{ where } y^* \text{ is determined by (5.2) with } \psi \in \mathcal{M}_D^\perp \right\}.$$

It follows from (5.3) that H is a closed subspace of $[H^{1/3}(0, T)]^*$ so is a Hilbert space. We next consider the following bilinear form on H :

$$\hat{a}(u, v) = \langle \mathcal{L}_1(u), v \rangle_{H^{1/3}(0, T); [H^{1/3}(0, T)]^*}.$$

Using (5.3), we derive from the Lax-Milgram theorem that there exists a continuous linear application $\hat{\mathcal{A}}: \mathcal{M}_D^\perp \rightarrow H$ such that

$$(5.8) \quad \int_0^L \varphi(x) y^*(0, x) dx = \hat{a}(\hat{\mathcal{A}}\varphi, y_{xx}^*(\cdot, L))$$

for all solution $y^* \in X_T$ of (5.2) with $\psi \in \mathcal{M}_D^\perp$.

Set

$$\hat{\mathcal{L}} = \mathcal{L}_1 \circ \hat{\mathcal{A}}.$$

The conclusion of *iii*) then follows after noting that if $y(0, \cdot) \in \mathcal{M}_D^\perp$ then $y(T, \cdot) \in \mathcal{M}_D^\perp$ by Lemma 5.1, and the fact

$$\int_0^L y(T, x) y^*(T, x) dx - \int_0^L y(0, x) y^*(0, x) dx = - \int_0^T y(t, L) y_{xx}^*(t, L) dt.$$

The proof is complete. \square

Here is the observability inequality used in the proof of Proposition 5.1.

Lemma 5.2. *Let $L \in \mathcal{N}_D$, $T > 0$, and let $y \in X_T$ be a solution of the linearized KdV equation in $(0, T) \times (0, L)$ with $y(0, L - \cdot) \in \mathcal{M}_D^\perp$ and with $y(\cdot, 0) = y(\cdot, L) = y_x(\cdot, L) = 0$. Then there exists $\Lambda \geq 1$ depending only on L and T such that*

$$(5.9) \quad \Lambda^{-1} \|y(0, \cdot)\|_{L^2(0, L)} \leq \|y_{xx}(\cdot, 0)\|_{[H^{1/3}(0, T)]^*} \leq \Lambda \|y(0, \cdot)\|_{L^2(0, L)}.$$

Proof. By Proposition 3.1, we have

$$(5.10) \quad \|y_{xx}(\cdot, 0)\|_{[H^{1/3}(0, T)]^*} \leq C \|y(0, \cdot)\|_{L^2(0, L)}.$$

We next prove

$$(5.11) \quad \|y(0, \cdot)\|_{L^2(0, L)} \leq C \|y_{xx}(\cdot, 0)\|_{[H^{1/3}(0, T)]^*}$$

by contradiction. Assume that there exists a sequence (φ_n) such that $\varphi_n(L - \cdot) \in \mathcal{M}_D^\perp$ and

$$(5.12) \quad \|y_{n,xx}(\cdot, 0)\|_{[H^{1/3}(0, T)]^*} \leq \frac{1}{n} \|y_n(0, \cdot)\|_{L^2(0, L)} = \frac{1}{n},$$

where $y_n \in X_T$ is the unique solution of the linearized KdV equation in $(0, T) \times (0, L)$ with $y_n(0, \cdot) = \varphi_n$ and $y_n(\cdot, 0) = y_n(\cdot, L) = y_{n,x}(\cdot, L) = 0$. Set

$$(5.13) \quad \mathbf{y}_n(t, x) = y_n(T - t, L - x).$$

Then $\mathbf{y}_n \in X_T$ is a solution of the equation $\mathbf{y}_{n,t} + \mathbf{y}_{n,x} + \mathbf{y}_{n,xxx} = 0$ in $(0, T) \times (0, L)$. By the regularizing effect of the linearized KdV equation, without loss of generality, one might assume that $\mathbf{y}_n(0, \cdot) = y_n(T, L - \cdot)$ converges in $L^2(0, L)$. Applying Proposition 3.2 to \mathbf{y}_n , one derives that $y_n(0, L - \cdot) = \mathbf{y}_n(T, \cdot)$ is a Cauchy sequence in $L^2(0, L)$. In other words, $\varphi_n = y_n(0, \cdot)$ is a Cauchy sequence in $L^2(0, L)$. Let φ be the limit of φ_n in $L^2(0, L)$ and denote $y \in X_T$ be the corresponding solution of the linearized KdV system. Then $\|\varphi\|_{L^2(0, L)} = 1$ and $\varphi(L - \cdot) \in \mathcal{M}_D^\perp$.

Set $\mathbf{y}(t, x) = y(T - t, L - x)$. Then $\mathbf{y} \in X_T$ is a solution of the system

$$(5.14) \quad \mathbf{y}_t + \mathbf{y}_x + \mathbf{y}_{xxx} = 0 \text{ in } (0, T) \times (0, L),$$

$$(5.15) \quad \mathbf{y}(\cdot, 0) = \mathbf{y}_x(\cdot, 0) = \mathbf{y}(\cdot, L) = \mathbf{y}_{xx}(\cdot, L) = 0 \text{ in } (0, T),$$

and

$$\mathbf{y}(0, \cdot) \in \mathcal{M}_D^\perp \text{ (by Lemma 5.1)} \quad \text{and} \quad \|\mathbf{y}(0, \cdot)\|_{L^2(0, L)} \stackrel{\text{Proposition 3.2}}{\geq} C\|\varphi\|_{L^2(0, L)} = C.$$

By Lemma 5.3 below, $\mathbf{y}(0, \cdot) = c\phi$ in $(0, L)$ for some $c \in \mathbb{R}$ where ϕ is defined by (1.8). Since $\mathbf{y}(0, \cdot) \in \mathcal{M}_D^\perp$, it follows that $c = 0$ and hence $\mathbf{y}(0, \cdot) = 0$ in $(0, L)$. We obtain a contradiction. Thus (5.11) is proved.

The conclusion now follows from (5.10) and (5.11). \square

In Lemma 5.2, we used the following result, which is also helpful in the proof of Assertion ii) of Theorem 1.3.

Lemma 5.3. *Let $L \in \mathcal{N}_D$ and $T > 0$. Assume that $y \in X_T$ is a solution of the system*

$$(5.16) \quad \begin{cases} y_t + y_x + y_{xxx} = 0 & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, 0) = y(\cdot, L) = y_{xx}(\cdot, L) = 0 & \text{in } (0, T). \end{cases}$$

There exists some $c \in \mathbb{R}$ such that

$$(5.17) \quad y = c\Phi \text{ in } (0, T) \times (0, L).$$

Proof. Set

$$(5.18) \quad \mathcal{V} = \left\{ \varphi \in L^2(0, T) \cap \mathcal{M}_D^\perp; \exists y \in X_T \text{ satisfying (5.16) and } y(0, \cdot) = \varphi \right\} \subset L^2(0, L).$$

We claim that

$$(5.19) \quad \mathcal{V} = \{0\}.$$

Admitting the claim, the conclusion now follows from the claim as follows. Set

$$\varphi = \text{proj}_{\mathcal{M}_D} y(0, \cdot).$$

Then

$$\varphi = c\phi \text{ for some } c \in \mathbb{R}.$$

Set

$$\tilde{y} = y - c\Phi \text{ in } (0, T) \times (0, L),$$

where Φ is defined by (1.10). Then $\tilde{y} \in X_T$ is a solution of (5.16) and $\tilde{y}(0, \cdot) \in \mathcal{M}_D^\perp$. By (5.19), one has

$$\tilde{y}(0, \cdot) = 0.$$

It follows that

$$y = c\Phi \text{ in } (0, T) \times (0, L),$$

which is the conclusion.

It remains to prove (5.19). We will prove that $\mathcal{V} = \{0\}$ by contradiction. Assume that $\mathcal{V} \neq \{0\}$. Using the regularity theory for the linear KdV equation and Proposition 3.2, one can show that

$$\mathcal{V} \subset C^\infty([0, L]).$$

For the same reason, one can show that any bounded sequence in \mathcal{V} (equipped $L^2(0, L)$ -norm) has a subsequence converging in \mathcal{V} . Thus \mathcal{V} is of finite dimension and is not $\{0\}$.

We can now involve the arguments via spectral theory in the spirit of [27] or even simpler (see also [1, 42]) to show that there exists $\varphi \in \mathcal{V} \setminus \{0\}$ and $\lambda \in \mathbb{C}$ such that

$$(5.20) \quad \lambda\varphi + \varphi_x + \varphi_{xxx} = 0 \text{ in } (0, L),$$

$$(5.21) \quad \varphi(\cdot, 0) = \varphi_x(\cdot, 0) = \varphi_x(\cdot, L) = \varphi_{xx}(\cdot, L) = 0.$$

Indeed, this can be done by considering

$$(5.22) \quad \begin{aligned} \mathcal{A} : \mathcal{V} &\rightarrow \mathcal{V} \\ \psi &\mapsto -(\psi_x + \psi_{xxx}) \end{aligned}$$

and taking $\lambda \in \mathbb{C}$ and $\varphi \in \mathcal{V} \setminus \{0\}$ such that $\mathcal{A}\varphi = \lambda\varphi$. The only point required to be checked is the fact that $\psi_x + \psi_{xxx} \in \mathcal{V}$ for $\psi \in \mathcal{V}$. To this end, one just notes that $-(\psi_x + \psi_{xxx}) = y_t(0, \cdot)$ where $y \in X_T$ is the corresponding solution (thus $y(0, \cdot) = \psi$).

Let η_j with $j = 1, 2, 3$ be the three solutions of the equations

$$\eta^3 + \eta + \lambda = 0.$$

Then φ has the form

$$\varphi = \gamma_1 e^{\eta_1 x} + \gamma_2 e^{\eta_2 x} + \gamma_3 e^{\eta_3 x} \text{ in } [0, L].$$

Using the boundary conditions of φ , we obtain

$$\begin{cases} \gamma_1 + \gamma_2 + \gamma_3 = 0, \\ \gamma_1 \eta_1 + \gamma_2 \eta_2 + \gamma_3 \eta_3 = 0, \\ \gamma_1 e^{\eta_1 L} + \gamma_2 e^{\eta_2 L} + \gamma_3 e^{\eta_3 L} = 0, \\ \gamma_1 \eta_1^2 e^{\eta_1 L} + \gamma_2 \eta_2^2 e^{\eta_2 L} + \gamma_3 \eta_3^2 e^{\eta_3 L} = 0. \end{cases}$$

As in [27, (30)], we then derive that

$$\eta_1 e^{-\eta_1 L} = \eta_2 e^{-\eta_2 L} = \eta_3 e^{-\eta_3 L}.$$

Without loss of generality, one can assume that

$$-\eta_1 L = z_1 \quad \text{and} \quad -\eta_2 L = z_2,$$

(here z_1 and z_2 are the complex numbers in the definition of \mathcal{N}_D) which yields

$$\varphi = c\phi$$

for some constant $c \in \mathbb{C}$. In other words,

$$\varphi \in \mathcal{M}_D.$$

We have a contradiction since $\varphi \neq 0$ and $\varphi \in \mathcal{M}_D^\perp$.

The proof is complete. \square

By the same arguments used in the proof of Lemma 5.2, we also have the following observability inequality in the case $L \notin \mathcal{N}_D$, which is the key point of the proof of Theorem 1.4.

Lemma 5.4. *Let $L \notin \mathcal{N}_D$, $T > 0$, and let $y \in X_T$ be a solution of the linearized KdV equation with $y(0, \cdot) \in L^2(0, L)$ and $y(\cdot, 0) = y(\cdot, L) = y_x(\cdot, 0) = 0$. Then, for some $\Lambda \geq 1$,*

$$(5.23) \quad \Lambda^{-1} \|y(0, \cdot)\|_{L^2(0, L)} \leq \|y_{xx}(\cdot, 0)\|_{[H^{1/3}(0, T)]^*} \leq \Lambda \|y(0, \cdot)\|_{L^2(0, L)}.$$

Here is a variant of Proposition 5.1 in the case $L \notin \mathcal{N}_D$.

Proposition 5.2. *Let $L \notin \mathcal{N}_D$ and $T > 0$. Then*

a) *There exists a linear continuous operator $\mathcal{L} : L^2(0, L) \rightarrow H^{1/3}(0, T)$ such that $y(T, \cdot) = \varphi$ where $y \in X_T$ is the unique solution of (5.1) with $y(0, \cdot) = 0$ and $u = \mathcal{L}(\varphi)$.*

b) *There exists a linear continuous operator $\hat{\mathcal{L}} : L^2(0, L) \rightarrow H^{1/3}(0, T)$ such that $y(T, \cdot) = 0$ where $y \in X_T$ is the unique solution of (5.1) with $y(0, \cdot) = \varphi$ and $u = \hat{\mathcal{L}}(\varphi)$.*

We are ready to give

Proof of Theorem 1.4. Theorem 1.4 follows from Proposition 3.4 and Lemma 5.4 as usual. The details are omitted. \square

6. PROPERTIES OF CONTROLS WHICH STEER 0 AT TIME 0 TO 0 AT TIME T

In this section, we study controls that steer 0 at time 0 to 0 at time T for the linearized KdV system of (1.1). To this end, it is convenient to introduce the following quantities.

Definition 6.1. For $z \in \mathbb{C}$, let $(\lambda_j)_{1 \leq j \leq 3} = (\lambda_j(z))_{1 \leq j \leq 3}$ be the three solutions of

$$(6.1) \quad \lambda^3 + \lambda + iz = 0.$$

Set

$$(6.2) \quad Q = Q(z) := \begin{pmatrix} 1 & 1 & 1 \\ e^{\lambda_1 L} & e^{\lambda_2 L} & e^{\lambda_3 L} \\ \lambda_1 e^{\lambda_1 L} & \lambda_2 e^{\lambda_2 L} & \lambda_3 e^{\lambda_3 L} \end{pmatrix},$$

$$(6.3) \quad P_D = P_D(z) := \sum_{j=1}^3 \lambda_j^2 (\lambda_{j+1} e^{\lambda_{j+1} L} - \lambda_{j+2} e^{\lambda_{j+2} L}),$$

and

$$(6.4) \quad \Xi = \Xi(z) := \det \begin{pmatrix} 1 & 1 & 1 \\ \lambda_1 & \lambda_2 & \lambda_3 \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 \end{pmatrix},$$

with the convention $\lambda_{j+3} = \lambda_j$ for $j \geq 1$.

Remark 6.1. The matrix Q and the quantities P_D and Ξ are antisymmetric with respect to λ_j ($j = 1, 2, 3$), and their definitions depend on a choice of the order of $(\lambda_1, \lambda_2, \lambda_3)$. Nevertheless, we later consider a product of either P_D , Ξ , or $\det Q$ with another antisymmetric function of (λ_j) , or deal with $|\det Q|$, and these quantities therefore make sense. The definitions of P , Ξ , and Q are only understood in these contexts.

Given $u \in H^{1/3}(0, +\infty)$, let $y \in C([0, +\infty); L^2(0, L)) \cap L_{\text{loc}}^2([0, +\infty); H^1(0, L))$ be the unique solution of the system

$$(6.5) \quad \begin{cases} y_t + y_x + y_{xxx} = 0 & \text{in } (0, +\infty) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, +\infty), \\ y(\cdot, L) = u & \text{in } (0, +\infty), \\ y(0, \cdot) = 0 & \text{in } (0, L). \end{cases}$$

In what follows in this section, we extend y and u by 0 for $t < 0$ and still denote these extensions by y and u , respectively. For an appropriate function v defined on $\mathbb{R} \times (0, L)$, let \hat{v} denote its Fourier transform with respect to t , i.e.,

$$\hat{v}(z, x) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} v(t, x) e^{-izt} dt.$$

From the system of y , we have

$$(6.6) \quad \begin{cases} iz\hat{y}(z, x) + \hat{y}_x(z, x) + \hat{y}_{xxx}(z, x) = 0 & \text{in } \mathbb{R} \times (0, L), \\ \hat{y}(z, 0) = \hat{y}_x(z, L) = 0 & \text{in } \mathbb{R}, \\ \hat{y}(z, L) = \hat{u}(z) & \text{in } \mathbb{R}. \end{cases}$$

This system has a unique solution outside a discrete set of z in \mathbb{C} , see [24, Lemma 2.1].

Taking into account the equation of \hat{y} , we search for the solution of the form

$$\hat{y}(z, \cdot) = \sum_{j=1}^3 a_j e^{\lambda_j x},$$

where $\lambda_j = \lambda_j(z)$ with $j = 1, 2, 3$ determined by (6.1).

Using the boundary conditions for \hat{y} , we require that

$$\begin{cases} \sum_{j=1}^3 a_j = 0, \\ \sum_{j=1}^3 e^{\lambda_j L} a_j = \hat{u}, \\ \sum_{j=1}^3 \lambda_j e^{\lambda_j L} a_j = 0. \end{cases}$$

This implies

$$Q(a_1, a_2, a_3)^\top = (0, \hat{u}, 0)^\top,$$

where $Q = Q(z)$ is given in Definition 6.1. We thus obtain

$$a_j = \frac{\hat{u}}{\det Q} (\lambda_{j+1} e^{\lambda_{j+1} L} - \lambda_{j+2} e^{\lambda_{j+2} L}).$$

This yields

$$(6.7) \quad \hat{y}(z, x) = \frac{\hat{u}}{\det Q} \sum_{j=1}^3 (\lambda_{j+1} e^{\lambda_{j+1} L} - \lambda_{j+2} e^{\lambda_{j+2} L}) e^{\lambda_j x}.$$

From (6.7), we derive that

$$(6.8) \quad \partial_{xx} \hat{y}(z, 0) = \frac{\hat{u}(z) P_D(z)}{\det Q(z)},$$

where $P_D(z)$ is given in Definition 6.1.

Set

$$(6.9) \quad G(z) = P_D(z)/\Xi(z) \quad \text{and} \quad H(z) = \det Q(z)/\Xi(z),$$

where $\Xi(z)$ is given in Definition 6.1. It is convenient to consider $\partial_{xx} \hat{y}(z, 0)$ under the form

$$(6.10) \quad \partial_{xx} \hat{y}(z, 0) = \frac{\hat{u}(z) G(z)}{H(z)},$$

By [24, Lemmas A1 and B1], $\partial_{xx} \hat{y}(z, 0)$ is a meromorphic function, and

$$(6.11) \quad G(z) \text{ and } H(z) \text{ are entire functions.}$$

We thus have just established the following result.

Lemma 6.1. *Let $u \in H^{1/3}(\mathbb{R}_+)$ and let $y \in C([0, +\infty); L^2(0, L)) \cap L_{\text{loc}}^2([0, +\infty); H^1(0, L))$ be the unique solution of*

$$(6.12) \quad \begin{cases} y_t + y_x + y_{xxx} = 0 & \text{in } (0, +\infty) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, +\infty), \\ y(\cdot, L) = u & \text{in } (0, +\infty), \end{cases}$$

with

$$(6.13) \quad y(0, \cdot) = 0 \text{ in } (0, L).$$

Outside a discrete set $z \in \mathbb{R}$, we have

$$(6.14) \quad \hat{y}(z, x) = \frac{\hat{u}}{\det Q} \sum_{j=1}^3 (\lambda_{j+1} e^{\lambda_{j+1}L} - \lambda_{j+2} e^{\lambda_{j+2}L}) e^{\lambda_j x} \text{ for a.e. } x \in (0, L).$$

Remark 6.2. Assume that $\hat{u}(z, \cdot)$ is well-defined for $z \in \mathbb{C}$ (e.g. when u has a compact support). Then the conclusions of Lemma 6.1 hold outside of a discrete set $z \in \mathbb{C}$.

We end this section with the following result, which is the starting point of our approach, and follows from Lemma 6.1 and Paley-Wiener's theorem, see e.g., [45].

Proposition 6.1. *Let $L > 0$, $T > 0$, and $u \in H^{1/3}(\mathbb{R}_+)$. Assume that u has a compact support in $[0, T]$, and u steers 0 at the time 0 to 0 at the time T , i.e., the unique solution y of (6.12) and (6.13) satisfies $y(T, \cdot) = 0$ in $(0, L)$. Then \hat{u} and $\hat{u}G/H$ satisfy the assumption of Paley-Wiener's theorem concerning the support in $[-T, T]$, i.e.,*

$$(6.15) \quad \hat{u} \text{ and } \hat{u}G/H \text{ are entire functions,}$$

and

$$(6.16) \quad |\hat{u}(z)| + \left| \frac{\hat{u}G(z)}{H(z)} \right| \leq C e^{T|z|},$$

for some positive constant C .

Remark 6.3. The computations in this section are in the spirit of the ones [24]. Nevertheless, in the conclusions of Proposition 6.1, we have/require that

$$\hat{y}(z, L) \text{ and } \partial_{xx}\hat{y}(z, 0) \text{ are entire functions satisfying (6.16).}$$

This is different with the one used in [24, Proposition 3.1] where one obtains that

$$\partial_x \hat{y}(z, L) \text{ and } \partial_x \hat{y}(z, 0) \text{ are entire functions satisfying a variant of (6.16).}$$

These differences are important to take into account different boundary conditions, see the proof of Assertion *ii*) of Theorem 1.3 in Section 9.2.

7. ATTAINABLE DIRECTIONS IN THE UNREACHABLE SPACE IN SMALL TIME

In this section, we investigate whether or not directions in \mathcal{M}_D , defined in (1.12), can be reached in small time. The starting point comes from the power series expansion approach. Let y_1 and y_2 be the solutions of

$$(7.1) \quad \begin{cases} y_{1,t} + y_{1,x} + y_{1,xxx} = 0 & \text{in } (0, T) \times (0, L), \\ y_1(\cdot, 0) = y_{1,x}(\cdot, L) = 0 & \text{in } (0, T), \\ y_1(\cdot, L) = u_1 & \text{in } (0, T), \\ y_1(0, \cdot) = 0 & \text{in } (0, L), \end{cases}$$

$$(7.2) \quad \begin{cases} y_{2,t} + y_{2,x} + y_{2,xxx} + y_1 y_{1,x} = 0 & \text{in } (0, T) \times (0, L), \\ y_2(\cdot, 0) = y_2(\cdot, L) = y_{2,x}(\cdot, L) = 0 & \text{in } (0, T), \\ y_2(0, \cdot) = 0 & \text{in } (0, L), \end{cases}$$

for some control u_1 . The key point of this approach is to first understand how one can choose u_1 so that

$$y_1(T, \cdot) = 0$$

and then analyse what the behavior of $\text{proj}_{\mathcal{M}_D} y_2(T, \cdot)$ is. To this end, we compute the quantity

$$(7.3) \quad \int_0^L y_2(T, x) \Phi(t, x) dx,$$

where Φ is defined in (1.10). Multiplying the equation of y_2 by Φ , integrating by parts in $(0, T) \times (0, L)$, we obtain, after using the boundary conditions and the initial conditions,

$$(7.4) \quad \int_0^L y_2(T, x) \Phi(t, x) dx = \frac{1}{2} \int_0^T \int_0^L y_1^2(t, x) \Phi_x(t, x) dx.$$

The goal is then to understand the value of the RHS of (7.4).

We will study the value of the RHS of (7.4) in a more general setting. Motivated by the definition of y_1 , we consider the unique solution $y \in X_T$ of the system, for $u \in H^{1/3}(0, T)$,

$$(7.5) \quad \begin{cases} y_t + y_x + y_{xxx} = 0 & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, T), \\ y(\cdot, L) = u & \text{in } (0, T), \\ y(0, \cdot) = 0 & \text{in } (0, L). \end{cases}$$

Guided by the definition of ϕ , as suggested in [24, 39], for $\eta_1, \eta_2, \eta_3 \in \mathbb{C} \setminus \{0\}$, we set

$$(7.6) \quad \varphi(x) = \sum_{j=1}^3 (\eta_{j+1} - \eta_j) e^{\eta_{j+2} x} \text{ for } x \in [0, L],$$

with the convention $\eta_{j+3} = \eta_j$ for $j \geq 1$. The following assumptions on η_j are used repeatedly throughout this section:

$$(7.7) \quad \eta_1 + \eta_2 + \eta_3 = 0, \quad \eta_1 \eta_2 + \eta_1 \eta_3 + \eta_2 \eta_3 = 1,$$

and

$$(7.8) \quad \eta_1 e^{-\eta_1 L} = \eta_2 e^{-\eta_2 L} = \eta_3 e^{-\eta_3 L}.$$

Extend y and u by 0 for $t > T$ and still denote these extensions by y and u , respectively. Then, by Lemma A.1 in the appendix,

$$\|u\|_{H^{1/3}(\mathbb{R}_+)} \leq C \|u\|_{H^{1/3}(0, T)}.$$

Assume that

$$y(T, \cdot) = 0.$$

Then the extension $y \in C([0, +\infty); L^2(0, L)) \cap L^2((0, +\infty); H^1(0, L))$ is also a solution of the linearized KdV system in $[0, +\infty) \times (0, L)$ using the control which is the extension of u (by 0 outside $(0, T)$), i.e.,

$$(7.9) \quad \begin{cases} y_t + y_x + y_{xxx} = 0 & \text{in } (0, +\infty) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, +\infty), \\ y(\cdot, L) = u & \text{in } (0, +\infty), \\ y(0, \cdot) = 0 & \text{in } (0, L). \end{cases}$$

In what follows in this section, we study this quantity, for $p \in \mathbb{R}$:

$$(7.10) \quad \int_0^T \int_0^L y^2(t, x) \varphi_x(x) e^{pt} dt dx \quad \left(= \int_0^{+\infty} \int_0^L y^2(t, x) \varphi_x(x) e^{pt} dt dx \right).$$

We have, by Lemma 6.1 (see also Remark 6.2), for $z \in \mathbb{C}$ outside a discrete set,

$$(7.11) \quad \hat{y}(z, x) = \hat{u}(z) \frac{\sum_{j=1}^3 (\lambda_j e^{\lambda_j L} - \lambda_{j+1} e^{\lambda_{j+1} L}) e^{\lambda_{j+2} x}}{\sum_{j=1}^3 (\lambda_{j+1} - \lambda_j) e^{-\lambda_{j+2} L}},$$

where $\lambda_j = \lambda_j(z)$ for $j = 1, 2, 3$ are determined by (6.1).

We begin with

Lemma 7.1. *Let $p \in \mathbb{R}$, $\eta_1, \eta_2, \eta_3 \in \mathbb{C} \setminus \{0\}$, and let φ be defined by (7.6). Set, for $(z, x) \in \mathbb{R} \times (0, L)$,*

$$B_D(z, x) = \left| \frac{\sum_{j=1}^3 (\lambda_j e^{\lambda_j L} - \lambda_{j+1} e^{\lambda_{j+1} L}) e^{\lambda_{j+2} x}}{\sum_{j=1}^3 (\lambda_{j+1} - \lambda_j) e^{-\lambda_{j+2} L}} \right|^2 \varphi_x(x),$$

where $\lambda_j = \lambda_j(z)$ with $j = 1, 2, 3$ are the three solutions of $\lambda^3 + \lambda + iz = 0$. Let $u \in H^{1/3}(0, +\infty)$ with compact support in $[0, +\infty)$ and let $y \in C([0, +\infty); L^2(0, L)) \cap L_{\text{loc}}^2([0, +\infty); H^1(0, L))$ be the unique solution of (7.9). Assume that $y(t, \cdot) = 0$ for large t . We have

$$(7.12) \quad \int_0^L \int_0^{+\infty} |y(t, x)|^2 \varphi_x(x) e^{pt} dt dx = \int_{\mathbb{R}} |\hat{u}(z + ip/2)|^2 \int_0^L B_D(z + ip/2, x) dx dz.$$

Proof. The conclusion is a direct consequence of Parseval's theorem and (7.11). \square

We next investigate the behavior of

$$\int_0^L B_D(z + ip/2, x) dx$$

for $z \in \mathbb{R}$ with large $|z|$. We have

Lemma 7.2. *Let $p \in \mathbb{R}$ and $\eta_1, \eta_2, \eta_3 \in \mathbb{C} \setminus \{0\}$. Assume (7.8). We have*

$$(7.13) \quad \int_0^L B_D(z + ip/2, x) dx = E_D |z|^{-1/3} + O(|z|^{-2/3}) \text{ for } z \in \mathbb{R} \text{ with large } |z|,$$

where $E_D = E_D(\eta_1, \eta_2, \eta_3)$ is defined by

$$(7.14) \quad E_D = \frac{1}{\sqrt{3}A} \sum_{j=1}^3 \eta_{j+2}^2 (\eta_{j+1} - \eta_j),$$

with ⁷

$$A = A(\eta_1, \eta_2, \eta_3) := \eta_j e^{-\eta_j L}.$$

Here and in what follows, for $s \in \mathbb{R}$, $O(z^s)$ denotes a quantity bounded by Cz^s for large positive z . Similar convention is used for $O(|z|^s)$ for $z \in \mathbb{C}$.

Proof. We first consider the case where z is positive and large. We use the following convention $\Re(\lambda_1) < \Re(\lambda_2) < \Re(\lambda_3)$.

We first look at the denominator of $B_D(z + ip/2, x)$. We have, by Lemma 3.1, at $(z + ip/2, x)$,

$$(7.15) \quad \frac{1}{\sum_{j=1}^3 (\lambda_{j+1} - \lambda_j) e^{-\lambda_{j+2} L}} \cdot \frac{1}{\sum_{j=1}^3 (\bar{\lambda}_{j+1} - \bar{\lambda}_j) e^{-\bar{\lambda}_{j+2} L}} \\ = \frac{e^{\lambda_1 L} e^{\bar{\lambda}_1 L}}{(\lambda_3 - \lambda_2)(\bar{\lambda}_3 - \bar{\lambda}_2)} \left(1 + O(e^{-C|z|^{1/3}}) \right).$$

⁷ A does not depend on j by (7.8).

We next deal with the numerator of $B_D(z + ip/2, x)$. Set, for $(z, x) \in \mathbb{R} \times (0, L)$,

$$(7.16) \quad f(z, x) = \sum_{j=1}^3 (\lambda_j e^{\lambda_j L} - \lambda_{j+1} e^{\lambda_{j+1} L}) e^{\lambda_{j+2} x}, \quad g(z, x) = \sum_{j=1}^3 (\bar{\lambda}_j e^{\bar{\lambda}_j L} - \bar{\lambda}_{j+1} e^{\bar{\lambda}_{j+1} L}) e^{\bar{\lambda}_{j+2} x},$$

and ⁸

$$f_m(z, x) = \lambda_3 e^{\lambda_3 L} e^{\lambda_2 x} - \lambda_2 e^{\lambda_2 L} e^{\lambda_3 x} - \lambda_3 e^{\lambda_3 L} e^{\lambda_1 x}, \quad g_m(z, x) = \bar{\lambda}_3 e^{\bar{\lambda}_3 L} e^{\bar{\lambda}_2 x} - \bar{\lambda}_2 e^{\bar{\lambda}_2 L} e^{\bar{\lambda}_3 x} - \bar{\lambda}_3 e^{\bar{\lambda}_3 L} e^{\bar{\lambda}_1 x}.$$

We have

$$\begin{aligned} & \int_0^L f(z + ip/2, x) g(z + ip/2, x) \varphi_x(x) dx \\ &= \int_0^L f_m(z + ip/2, x) g_m(z + ip/2, x) \varphi_x(x) dx + \int_0^L (f - f_m)(z + ip/2, x) g_m(z + ip/2, x) \varphi_x(x) dx \\ &+ \int_0^L f_m(z + ip/2, x) (g - g_m)(z + ip/2, x) \varphi_x(x) dx + \int_0^L (f - f_m)(z + ip/2, x) (g - g_m)(z + ip/2, x) \varphi_x(x) dx. \end{aligned}$$

By Lemma 3.1, we have

$$(7.17) \quad \begin{aligned} & \int_0^L |(f - f_m)(z + ip/2, x) g_m(z + ip/2, x) \varphi_x(x)| dx \\ &+ \int_0^L |(f - f_m)(z + ip/2, x) (g - g_m)(z + ip/2, x) \varphi_x(x)| dx \\ &+ \int_0^L |f_m(z + ip/2, x) (g - g_m)(z + ip/2, x) \varphi_x(x)| dx \leq C |e^{(\lambda_3 + \bar{\lambda}_3)L}| e^{-C|z|^{1/3}}. \end{aligned}$$

We next estimate

$$(7.18) \quad \begin{aligned} & \int_0^L f_m(z + ip/2, x) g_m(z + ip/2, x) \varphi_x(x) \\ &= \int_0^L f_m(z + ip/2, x) g_m(z + ip/2, x) \left(\sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) e^{\eta_{j+2} x} \right) dx. \end{aligned}$$

We first have, by Lemma 3.1, at $z + ip/2$,

$$(7.19) \quad \begin{aligned} & \int_0^L \left(-\lambda_3 e^{\lambda_3 L} e^{\lambda_2 x} \bar{\lambda}_2 e^{\bar{\lambda}_2 L} e^{\bar{\lambda}_3 x} - \lambda_2 e^{\lambda_2 L} e^{\lambda_3 x} \bar{\lambda}_3 e^{\bar{\lambda}_3 L} e^{\bar{\lambda}_2 x} + \lambda_2 e^{\lambda_2 L} e^{\lambda_3 x} \bar{\lambda}_2 e^{\bar{\lambda}_2 L} e^{\bar{\lambda}_3 x} \right) \\ & \times \left(\sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) e^{\eta_{j+2} x} \right) dx \stackrel{(7.8)}{=} A^{-1} e^{(\lambda_3 + \bar{\lambda}_3 + \lambda_2 + \bar{\lambda}_2)L} \left(S_1(z + ip/2) + O(e^{-C|z|^{1/3}}) \right), \end{aligned}$$

where

$$(7.20) \quad S_1(z) := \sum_{j=1}^3 \eta_{j+2}^2 (\eta_{j+1} - \eta_j) \left(\frac{\lambda_2 \bar{\lambda}_2}{\lambda_3 + \bar{\lambda}_3 + \eta_{j+2}} - \frac{\lambda_2 \bar{\lambda}_3}{\lambda_3 + \bar{\lambda}_2 + \eta_{j+2}} - \frac{\lambda_3 \bar{\lambda}_2}{\lambda_2 + \bar{\lambda}_3 + \eta_{j+2}} \right) (z).$$

⁸The index m stands the main part.

We next obtain, by Lemma 3.1, at $z + ip/2$,

$$(7.21) \quad \int_0^L \left(\lambda_3 e^{\lambda_3 L} e^{\lambda_1 x} \bar{\lambda}_3 e^{\bar{\lambda}_3 L} e^{\bar{\lambda}_1 x} - \lambda_3 e^{\lambda_3 L} e^{\lambda_1 x} \bar{\lambda}_3 e^{\bar{\lambda}_3 L} e^{\bar{\lambda}_2 x} - \lambda_3 e^{\lambda_3 L} e^{\lambda_2 x} \bar{\lambda}_3 e^{\bar{\lambda}_3 L} e^{\bar{\lambda}_1 x} \right) \\ \times \left(\sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) e^{\eta_{j+2} x} \right) dx = e^{(\lambda_3 + \bar{\lambda}_3)L} \left(S_2(z + ip/2) + O(e^{-C|z|^{1/3}}) \right),$$

where

$$(7.22) \quad S_2(z) := \sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) \left(-\frac{\lambda_3 \bar{\lambda}_3}{\lambda_1 + \bar{\lambda}_1 + \eta_{j+2}} + \frac{\lambda_3 \bar{\lambda}_3}{\lambda_1 + \bar{\lambda}_2 + \eta_{j+2}} + \frac{\lambda_3 \bar{\lambda}_3}{\lambda_2 + \bar{\lambda}_1 + \eta_{j+2}} \right) (z).$$

We have, at $z + ip/2$,

$$(7.23) \quad \int_0^L \lambda_3 e^{\lambda_3 L} e^{\lambda_2 x} \bar{\lambda}_3 e^{\bar{\lambda}_3 L} e^{\bar{\lambda}_2 x} \left(\sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) e^{\eta_{j+2} x} \right) dx = e^{(\lambda_3 + \bar{\lambda}_3)L} S_3(z + ip/2),$$

where

$$(7.24) \quad S_3(z) := \sum_{j=1}^3 \frac{\eta_{j+2} (\eta_{j+1} - \eta_j) \lambda_3 \bar{\lambda}_3 \left(e^{\lambda_2 L + \bar{\lambda}_2 L + \eta_{j+2} L} - 1 \right)}{\lambda_2 + \bar{\lambda}_2 + \eta_{j+2}} (z).$$

We finally get, by Lemma 3.1, at $z + ip/2$,

$$(7.25) \quad \left| \int_0^L \left(\lambda_3 e^{\lambda_3 L} e^{\lambda_1 x} \bar{\lambda}_2 e^{\bar{\lambda}_2 L} e^{\bar{\lambda}_3 x} + \lambda_2 e^{\lambda_2 L} e^{\lambda_3 x} \bar{\lambda}_3 e^{\bar{\lambda}_3 L} e^{\bar{\lambda}_1 x} \right) \left(\sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) e^{\eta_{j+2} x} \right) dx \right| \\ = |e^{(\lambda_3 + \bar{\lambda}_3)L}| O(e^{-Cz^{1/3}}).$$

By Lemma 3.1, we have, at $z + ip/2$,

$$(7.26) \quad \begin{cases} \lambda_1 + \bar{\lambda}_1 + \lambda_2 + \bar{\lambda}_2 + \lambda_3 + \bar{\lambda}_3 = O(z^{-1/3}), \\ \lambda_1 + \bar{\lambda}_1 + \lambda_3 + \bar{\lambda}_3 = O(z^{-1/3}), \\ (\lambda_3 - \lambda_2)(\bar{\lambda}_3 - \bar{\lambda}_2) = 3z^{2/3}(1 + O(z^{-1/3})). \end{cases}$$

We claim that

$$(7.27) \quad |z|^{-1/3} |S_1(z + ip/2)| + |S_2(z + ip/2)| + |S_3(z + ip/2)| = O(1) \text{ for large positive } z.$$

Admitting this, by combining (7.15), (7.19), (7.21), (7.23), (7.25), and (7.26), and using (7.27), we obtain

$$(7.28) \quad \int_0^L B_D(z + ip/2, x) dx = \frac{A^{-1} S_1(z + ip/2)}{3z^{2/3}} + O(|z|^{-2/3}).$$

We first derive the the asymptotic behavior of $S_1(z + ip/2)$. We have, by Lemma 3.1, at $z + ip/2$,

$$(7.29) \quad \lambda_2 \bar{\lambda}_2 = z^{2/3} + O(1), \quad \lambda_2 \bar{\lambda}_3 = z^{2/3} e^{i\pi/3} + O(1), \quad \lambda_3 \bar{\lambda}_2 = z^{2/3} e^{-i\pi/3} + O(1),$$

and

$$(7.30) \quad \frac{1}{\lambda_3 + \bar{\lambda}_3 + \eta_{j+2}} = \frac{1 + O(z^{-1/3})}{\sqrt{3}z^{1/3}}, \quad \frac{1}{\lambda_2 + \bar{\lambda}_3 + \eta_{j+2}} = \frac{1 + O(z^{-1/3})}{(e^{i\pi/6} + i)z^{1/3}},$$

$$\frac{1}{\lambda_3 + \bar{\lambda}_2 + \eta_{j+2}} = \frac{1 + O(z^{-1/3})}{(e^{-i\pi/6} - i)z^{1/3}}.$$

It follows that

$$(7.31) \quad S_1(z + ip/2) = \left(\frac{1}{\sqrt{3}} + 2\Re \frac{e^{i\pi/3}}{e^{i\pi/6} + i} \right) \sum_{j=1}^3 \eta_{j+2}^2 (\eta_{j+1} - \eta_j) |z|^{1/3} + O(1)$$

$$= \sqrt{3} \sum_{j=1}^3 \eta_{j+2}^2 (\eta_{j+1} - \eta_j) |z|^{1/3} + O(1).$$

We next deal with $S_2(z + ip/2)$. Since

$$\sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) = 0,$$

it follows that

$$S_2(z + ip/2) = \sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) \left(-\frac{\lambda_3 \bar{\lambda}_3}{\lambda_1 + \bar{\lambda}_1 + \eta_{j+2}} + \frac{\lambda_3 \bar{\lambda}_3}{\lambda_1 + \bar{\lambda}_1} \right) (z + ip/2)$$

$$+ \sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) \left(\frac{\lambda_3 \bar{\lambda}_3}{\lambda_1 + \bar{\lambda}_2 + \eta_{j+2}} - \frac{\lambda_3 \bar{\lambda}_3}{\lambda_1 + \bar{\lambda}_2} \right) (z + ip/2)$$

$$+ \sum_{j=1}^3 \eta_{j+2} (\eta_{j+1} - \eta_j) \left(\frac{\lambda_3 \bar{\lambda}_3}{\lambda_2 + \bar{\lambda}_1 + \eta_{j+2}} - \frac{\lambda_3 \bar{\lambda}_3}{\lambda_2 + \bar{\lambda}_1} \right) (z + ip/2).$$

Using Lemma 3.1, we derive that

$$(7.32) \quad S_2(z + ip/2) = O(1).$$

We next derive the asymptotic behavior of $S_3(z + ip/2)$. From Lemma 3.1, we have, at $z + ip/2$,

$$(7.33) \quad \lambda_2 + \bar{\lambda}_2 = O(z^{-2/3}).$$

Using (7.8), we derive from (7.24) that

$$S_3(z + ip/2) = z^{2/3} \sum_{j=1}^3 \frac{(\eta_{j+1} - \eta_j) \eta_{j+2}}{\eta_{j+2}} \left(A^{-1} \eta_{j+2} - 1 \right) + O(1)$$

$$= z^{2/3} \sum_{j=1}^3 (\eta_{j+1} - \eta_j) \left(A^{-1} \eta_{j+2} - 1 \right) + O(1).$$

Since $\sum_{j=1}^3 (\eta_{j+1} - \eta_j) = 0$ and $\sum_{j=1}^3 (\eta_{j+1} - \eta_j) \eta_{j+2} = 0$, it follows that

$$(7.34) \quad S_3(z + ip/2) = O(1).$$

Combing (7.28) and (7.31) yields

$$\int_0^L B_D(z + ip/2, x) dz = E_D |z|^{-1/3} + O(z^{-2/3}),$$

where

$$(7.35) \quad E_D = \frac{1}{\sqrt{3}A} \sum_{j=1}^3 \eta_{j+2}^2 (\eta_{j+1} - \eta_j).$$

The conclusion in the case where z is large and negative can be derived from the case where z is positive and large as follows. Define, for $(z, x) \in \mathbb{R} \times (0, L)$, with large $|z|$,

$$M_D(z, x) = \frac{\sum_{j=1}^3 (\lambda_j e^{\lambda_j L} - \lambda_{j+1} e^{\lambda_{j+1} L}) e^{\lambda_{j+2} x}}{\sum_{j=1}^3 (\lambda_{j+1} - \lambda_j) e^{-\lambda_{j+2} L}}.$$

Then

$$B_D(z + ip/2, x) = |M_D(z + ip/2, x)|^2 \varphi_x(x).$$

It is clear from the definition of M_D that

$$M_D(-z, x) = \overline{M_D(\bar{z}, x)}.$$

We then have

$$B_D(-z + ip/2, x) = |M_D(-z + ip/2, x)|^2 \varphi_x(x) = |M_D(z - ip/2, x)|^2 \varphi_x(x).$$

We thus obtain the result in the case where z is negative and large by the corresponding expression for large positive z in which p is replaced by $-p$. The conclusion follows. \square

As a consequence of Lemmas 7.1 and 7.2, we obtain

Proposition 7.1. *Let $L \in \mathcal{N}_D$. Let $u \in H^{1/3}(0, +\infty)$ with compact support in $[0, +\infty)$, and let $y \in C([0, +\infty); L^2(0, L)) \cap L^2_{\text{loc}}([0, +\infty); H^1(0, L))$ be the unique solution of (7.9). Assume that $y(t, \cdot) = 0$ for large t . We have*

$$(7.36) \quad \int_0^{+\infty} \int_0^L |y(t, x)|^2 \Phi_x(t, x) dx dt = \int_{\mathbb{R}} |\hat{u}(z + iq/2)|^2 \int_0^L B(z + iq/2, x) dx dz,$$

where

$$B(z, x) = \left| \frac{\sum_{j=1}^3 (\lambda_j e^{\lambda_j L} - \lambda_{j+1} e^{\lambda_{j+1} L}) e^{\lambda_{j+2} x}}{\sum_{j=1}^3 (\lambda_{j+1} - \lambda_j) e^{-\lambda_{j+2} L}} \right|^2 \phi_x(x),$$

with $\lambda_j = \lambda_j(z)$ ($j = 1, 2, 3$) being the three solutions of $\lambda^3 + \lambda + iz = 0$. Moreover,

$$(7.37) \quad \int_0^L B(z + iq/2, x) dx = E |z|^{-1/3} + O(|z|^{-2/3}),$$

where

$$E = -\frac{be^{2a}}{2\sqrt{3}aL^2} (b^2 + 9a^2) > 0.$$

Recall that q, a, b, ϕ , and Φ are defined in Theorem 1.1.

Remark 7.1. Note that

$$(7.38) \quad \Omega(z) = \left| \sum_{j=1}^3 (\lambda_{j+1} - \lambda_j) e^{-\lambda_{j+2} L} \right|^2 \int_0^L B(z, x) dx.$$

We then derive from (7.37) and Lemma 3.1 that

$$\lim_{|z| \rightarrow +\infty, z \in \mathbb{R}} \Omega(z + iq/2) = +\infty.$$

Proof of Proposition 7.1. We will apply Lemmas 7.1 and 7.2 with $p = q$,

$$\eta_1 = \alpha + i\beta, \quad \eta_2 = \alpha - i\beta, \quad \text{and} \quad \eta_3 = -2\alpha,$$

where α and β are defined in (1.9). Thus

$$(7.39) \quad A = \eta_3 e^{-\eta_3 L} = -2\alpha e^{2\alpha L} = \frac{2a}{L} e^{-2a}.$$

We have

$$\begin{aligned} \sqrt{3}AE_D &= \eta_1^2(\eta_3 - \eta_2) + \eta_2^2(\eta_1 - \eta_3) + \eta_3^2(\eta_2 - \eta_1) \\ &= (\alpha + i\beta)^2(-3\alpha + i\beta) + (\alpha - i\beta)^2(3\alpha + i\beta) - 8i\alpha^2\beta \\ &= -12i\alpha^2\beta + 2i\beta(\alpha^2 - \beta^2) - 8i\alpha^2\beta. \end{aligned}$$

This implies

$$(7.40) \quad \sqrt{3}AE_D = -2i(\beta^2 + 9\alpha^2)\beta = \frac{2ib}{L^3}(b^2 + 9a^2).$$

We obtain

$$E_D = \frac{ibe^{2a}}{\sqrt{3}aL^2}(b^2 + 9a^2).$$

Since, by Lemma 4.1,

$$\varphi = -2i\phi,$$

the conclusion now follows from Lemmas 7.1 and 7.2. \square

Using Proposition 7.1, we derive the following result which is the key ingredient for the analysis of the local controllability of the KdV system (1.1) in small time.

Proposition 7.2. *Let $L \in \mathcal{N}_D$ and $T < 1$. Let $u \in H^{1/3}(0, +\infty)$ and let $y \in C([0, +\infty); L^2(0, L)) \cap L^2_{\text{loc}}([0, +\infty); H^1(0, L))$ be the unique solution of (7.9). Assume that $u(t) = 0$ for $t > T$, and $y(t, \cdot) = 0$ for large t . Then*

$$(7.41) \quad \int_0^\infty \int_0^L |y(t, x)|^2 \Phi_x(t, x) dx dt = E \|ue^{q/2}\|_{[H^{1/6}(0, +\infty)]^*}^2 \left(1 + O(T^{2/9})\right).$$

Proof. We have, by Proposition 7.1 and Lemma B.2,

$$|\hat{u}(z + iq/2)|^2 \left| \int_0^L B(z + iq/2, x) dx \right| \leq \frac{C|\hat{u}(z + iq/2)|^2}{1 + |z|^{1/3}}.$$

Applying Proposition 7.1 again and using the fact $\int_{\mathbb{R}} dz = \int_{z \in \mathbb{R}; |z| < m} dz + \int_{z \in \mathbb{R}; |z| \geq m} dz$, we derive that

$$(7.42) \quad \left| \int_{\mathbb{R}} |\hat{u}(z + iq/2)|^2 \int_0^L B(z + iq/2, x) dx dz - E \int_{\mathbb{R}} \frac{|\hat{u}(z + iq/2)|^2}{(1 + |z|^2)^{1/6}} dz \right| \\ \leq C \int_{|z| < m} \frac{|\hat{u}(z + iq/2)|^2}{(1 + |z|^2)^{1/6}} + Cm^{-1/3} \int_{|z| > m} \frac{|\hat{u}(z + iq/2)|^2}{(1 + |z|^2)^{1/6}} dz.$$

Applying Lemma 7.3 below with $s = 1/6$, we have

$$(7.43) \quad |\hat{u}(z + iq/2)| \leq C(|z|T^{4/3} + T^{1/3}) \|ue^{q/2}\|_{[H^{1/6}(0, +\infty)]^*}.$$

Combining (7.42) and (7.43) yields

$$(7.44) \quad \left| \int_{\mathbb{R}} |\hat{u}(z + iq/2)|^2 \int_0^L B(z + iq/2, x) dx dz - E \int_{\mathbb{R}} \frac{|\hat{u}(z + iq/2)|^2}{(1 + |z|^2)^{1/6}} dz \right| \\ \leq C \|ue^{q/2}\|_{[H^{1/6}(0, +\infty)]^*}^2 (T^{8/3} m^{8/3} + T^{2/3} m^{2/3} + m^{-\frac{1}{3}}).$$

By choosing $m = T^{-2/3}$, we obtain

$$\left| \int_{\mathbb{R}} |\hat{u}(z + iq/2)|^2 \int_0^L B(z + iq/2, x) dx dz - E \int_{\mathbb{R}} \frac{|\hat{u}(z + iq/2)|^2}{(1 + |z|^2)^{1/6}} dz \right| \leq C \|ue^{q/2}\|_{H^{-1/6}(0, +\infty)}^2 T^{2/9}.$$

The conclusion follows. \square

In the proof of Proposition 7.2, we have used the following lemma.

Lemma 7.3. *Let $0 < T < 1$, $0 < s < 1/2$ and let $v \in [H^s(0, +\infty)]^*$ with $\text{supp } v \subset [0, T]$. Extend v by 0 for $t < 0$ and still denote this extension by v . There exists a positive constant C independent of T and v such that, for $z \in \mathbb{R}$,*

$$(7.45) \quad |\hat{v}(z)| \leq C(|z|T^{\frac{3}{2}-s} + T^{\frac{1}{2}-s}) \|v\|_{[H^s(0, +\infty)]^*}.$$

Proof. Fix $\chi \in C^1(\mathbb{R})$ such that $\chi = 1$ in $[0, T]$, $\chi = 0$ for $t \geq 2T$, and $0 \leq \chi \leq 1$ and $\chi' \leq C/T$ in \mathbb{R} . Then

$$(7.46) \quad |\hat{v}(z)| \leq \|v\|_{[H^s(0, +\infty)]^*} \|e^{-iz} \cdot \chi\|_{H^s(0, +\infty)}.$$

We have

$$(7.47) \quad C \|e^{-iz} \cdot \chi\|_{H^s(0, +\infty)}^2 \leq \int_0^{4T} \int_0^{4T} \frac{|\chi(x)|^2 |e^{-izx} - e^{-izy}|^2}{|x - y|^{1+2s}} dx dy \\ + \int_0^{4T} \int_0^{4T} \frac{|\chi(x) - \chi(y)|^2}{|x - y|^{1+2s}} dx dy + \int_0^{2T} \int_{4T}^{\infty} \frac{1}{|x - y|^{1+2s}} dx dy.$$

Using the fact, for $x \in \mathbb{R}$,

$$|e^{ix} - 1|^2 = 4 \sin^2(x/2) \leq x^2,$$

we derive that the RHS of (7.47) is bounded by

$$(7.48) \quad Cz^2 T^{3-2s} + CT^{1-2s} + CT^{1-2s}.$$

The conclusion now follows from (7.46) and (7.48). \square

8. SMALL TIME LOCAL CONTROLLABILITY PROPERTIES OF THE KDV SYSTEM - PROOF OF THEOREM 1.2

The main result of this section is the following, which implies, in particular, Theorem 1.2 after applying Proposition 7.2 and Lemma 8.2 at the end of this section. This result is also the key ingredient in the proof of Assertion i) of Theorem 1.3.

Theorem 8.1. *Let $L \in \mathcal{N}_D$, $T > 0$, and $\alpha > 0$. Assume that*

$$(8.1) \quad \int_0^T \int_0^L |\xi(t, x)|^2 \Phi(t, x) dx dt \geq \alpha \|v\|_{[H^{1/6}(0, T)]^*}^2$$

for all $v \in H^{1/3}(0, T)$ such that $\xi(T, \cdot) = 0$ in $(0, L)$ where $\xi \in X_T$ is the unique solution of the system

$$(8.2) \quad \begin{cases} \xi_t + \xi_x + \xi_{xxx} = 0 & \text{in } (0, T) \times (0, L), \\ \xi(\cdot, 0) = \xi_x(\cdot, L) = 0 & \text{in } (0, T), \\ \xi(\cdot, L) = v & \text{in } (0, T), \\ \xi(0, \cdot) = 0 & \text{in } (0, L). \end{cases}$$

There exists $\varepsilon_0 > 0$ depending only on α, T , and L such that for all $0 < \varepsilon < \varepsilon_0$, and for all solutions $y \in X_{T/2}$ of the system

$$(8.3) \quad \begin{cases} y_t + y_x + y_{xxx} + yy_x = 0 & \text{in } (0, T/2) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, T/2), \\ y(\cdot, L) = u & \text{in } (0, T/2), \\ y(0, \cdot) = \varepsilon\phi & \text{in } (0, L), \end{cases}$$

with $\|u\|_{H^{1/2}(0, T/2)} < \varepsilon_0$, we have

$$y(T/2, \cdot) \neq 0.$$

Recall that ϕ is defined in (1.8).

Proof. Let ε_0 be a small positive constant, which depends only on α, T , and L , and is determined later. We prove Theorem 8.1 by contradiction. Assume that there exists a solution $y \in C([0, +\infty); L^2(0, L)) \cap L^2_{\text{loc}}([0, +\infty); H^1(0, L))$ of (8.3) with $y(t, \cdot) = 0$ for $t \geq T/2$, for some $0 < \varepsilon < \varepsilon_0$, for some $u \in H^{1/3}(0, +\infty)$ with $\text{supp } u \subset [0, T/2]$ and $\|u\|_{H^{1/2}(0, T)} < \varepsilon_0$.

Using the fact $y(t, \cdot) = 0$ for $t \geq T/2$, from Proposition 3.4, we have, for ε_0 small,

$$(8.4) \quad \|y\|_{L^2(\mathbb{R}_+; H^1(0, L))} \leq C \left(\|y_0\|_{L^2(0, L)} + \|u\|_{H^{1/3}(\mathbb{R}_+)} \right),$$

which in turn implies, by Proposition 3.1 and Lemma 8.1 below,

$$(8.5) \quad \|y\|_{L^2(\mathbb{R}_+ \times (0, L))} \leq C \left(\|y_0\|_{L^2(0, L)} + \|u\|_{L^2(\mathbb{R}_+)} \right).$$

Here and in what follows, C denotes a positive constant depending only on T and L (C thus does not depend on α).

Let y_1 and y_2 be the solution of the following systems

$$(8.6) \quad \begin{cases} y_{1,t} + y_{1,x} + y_{1,xxx} = 0 & \text{in } \mathbb{R}_+ \times (0, L), \\ y_1(\cdot, 0) = y_{1,x}(\cdot, L) = 0 & \text{in } \mathbb{R}_+, \\ y_1(\cdot, L) = u & \text{in } \mathbb{R}_+, \\ y_1(0, \cdot) = 0 & \text{in } (0, L), \end{cases}$$

$$(8.7) \quad \begin{cases} y_{2,t} + y_{2,x} + y_{2,xxx} + y_1 y_{1,x} = 0 & \text{in } \mathbb{R}_+ \times (0, L), \\ y_2(\cdot, 0) = y_{2,x}(\cdot, L) = y_2(\cdot, L) = 0 & \text{in } \mathbb{R}_+, \\ y_2(0, \cdot) = 0 & \text{in } (0, L). \end{cases}$$

Then, by Proposition 3.1,

$$(8.8) \quad \|y_1\|_{X_T} \leq C \|u\|_{H^{1/3}(\mathbb{R}_+)}$$

and

$$(8.9) \quad \|y_1\|_{L^2((0,T) \times (0,L))} \leq C \|u\|_{L^2(\mathbb{R}_+)},$$

which in turn imply, by Proposition 3.1 and Lemma 8.1 below,

$$(8.10) \quad \|y_2\|_{X_T} \leq C \|y_1 y_{1,x}\|_{L^1((0,T); L^2(0,L))} \leq C \|u\|_{L^2(\mathbb{R}_+)}^{1/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{3/2},$$

and

$$(8.11) \quad \|y_2\|_{L^2((0,T) \times (0,L))} \leq C \|y_1^2\|_{L^1((0,T); L^2(0,L))} \leq C \|u\|_{L^2(\mathbb{R}_+)}^{3/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2}.$$

Set

$$\delta y = y - y_1 - y_2 \text{ in } \mathbb{R}_+ \times (0, L).$$

We have

$$\begin{cases} \delta y_t + \delta y_x + \delta y_{xxx} + y \delta y_x + \delta y (y_1 + y_2)_x + (y_1 y_2)_x + y_2 y_{2,x} = 0 & \text{in } \mathbb{R}_+ \times (0, L), \\ \delta y(\cdot, 0) = \delta y_x(\cdot, L) = \delta y(\cdot, L) = 0 & \text{in } \mathbb{R}_+, \\ \delta y(0, \cdot) = \varepsilon \phi & \text{in } (0, L). \end{cases}$$

Applying Proposition 3.1 and using Lemma 8.1 below, we derive that

$$(8.12) \quad \|\delta y\|_{X_T} \leq C \left(\|y_1 y_{2,x}\|_{L^1((0,T); L^2(0,L))} + \|y_2 y_{1,x}\|_{L^1((0,T); L^2(0,L))} + \|y_2 y_{2,x}\|_{L^1((0,T); L^2(0,L))} \right) \\ \stackrel{(8.8)-(8.11)}{\leq} C \|u\|_{L^2(\mathbb{R}_+)} \|u\|_{H^{1/3}(\mathbb{R}_+)}^2 + C\varepsilon,$$

and

$$(8.13) \quad \|\delta y\|_{L^2((0,T) \times (0,L))} \leq C \left(\|y \delta y_x\|_{L^1((0,T); L^2(0,L))} + \|(y_1 + y_2) \delta y_x\|_{L^1((0,T); L^2(0,L))} \right) \\ + C \left(\|(y_1 + y_2) \delta y\|_{L^1((0,T); L^2(0,L))} + \|y_1 y_2\|_{L^1((0,T); L^2(0,L))} + \|y_2^2\|_{L^1((0,T); L^2(0,L))} \right) \\ \stackrel{(8.4)-(8.5), (8.8)-(8.12)}{\leq} C \|u\|_{L^2(\mathbb{R}_+)}^{3/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{5/2} + C \|u\|_{L^2(\mathbb{R}_+)}^2 \|u\|_{H^{1/3}(\mathbb{R}_+)} + C\varepsilon \\ \leq C \|u\|_{L^2(\mathbb{R}_+)}^{3/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{3/2} + C\varepsilon.$$

In (8.12), we absorbed the contribution in the RHS of

$$\|y \delta y_x\|_{L^1((0,T); L^2(0,L))} + \|\delta y (y_1 + y_2)_x\|_{L^1((0,T); L^2(0,L))},$$

which is bounded by $C\varepsilon_0 \|\delta y\|_{X_T}$. Combining (8.11) and (8.13) yields

$$(8.14) \quad \|y - y_1\|_{L^2((0,T) \times (0,L))} \leq C \|u\|_{L^2(\mathbb{R}_+)}^{3/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2} + C\varepsilon.$$

Since $y = 0$ for $t \geq T/2$ and $u = 0$ for $t \geq T/2$, after using the regularizing effect of the linear KdV equation, and considering the projection into \mathcal{M}_D^\perp , we derive that

$$(8.15) \quad \|y_1(T, \cdot)\|_{L^2(0,L)} \leq C \|u\|_{L^2(\mathbb{R}_+)}^{3/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2} + C\varepsilon.$$

Since $y_1(T, \cdot) \in \mathcal{M}_D^\perp$ by Lemma 5.1, it follows from Proposition 5.1 that there exists $u_1 \in H^{1/3}(0, T)$ such that

$$(8.16) \quad \|u_1\|_{H^{1/3}(0,T)} \leq C \|y_1(T, \cdot)\|_{L^2(0,L)}$$

and the solution $\tilde{y}_1 \in X_T$ of the system

$$\begin{cases} \tilde{y}_{1,t} + \tilde{y}_{1,x} + \tilde{y}_{1,xxx} = 0 & \text{in } (0, T) \times (0, L), \\ \tilde{y}_1(\cdot, 0) = \tilde{y}_{1,x}(\cdot, L) = 0 & \text{in } (0, T), \\ \tilde{y}_1(\cdot, L) = u_1 & \text{in } (0, T), \\ \tilde{y}_1(0, \cdot) = 0 & \text{in } (0, L), \end{cases}$$

satisfies

$$\tilde{y}_1(T, \cdot) = -y_1(T, \cdot).$$

Using (8.15), we derive from (8.16) that

$$(8.17) \quad \|u_1\|_{H^{1/3}(0,T)} \leq C \|u\|_{L^2(\mathbb{R}_+)}^{3/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2} + C\varepsilon,$$

which in turn implies, by Proposition 3.1,

$$(8.18) \quad \|\tilde{y}_1\|_{X_T} \leq C \|u\|_{L^2(\mathbb{R}_+)}^{3/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2} + C\varepsilon.$$

Set

$$\tilde{y} = y_1 + \tilde{y}_1 \text{ in } (0, T) \times (0, L).$$

Then

$$(8.19) \quad \begin{cases} \tilde{y}_t + \tilde{y}_x + \tilde{y}_{xxx} = 0 & \text{in } (0, T) \times (0, L), \\ \tilde{y}(\cdot, 0) = \tilde{y}_x(\cdot, L) = 0 & \text{in } (0, T), \\ \tilde{y}(\cdot, L) = u + u_1 & \text{in } (0, T), \\ \tilde{y}(0, \cdot) = \tilde{y}(T, \cdot) = 0 & \text{in } (0, L). \end{cases}$$

We have

$$(8.20) \quad \|y - \tilde{y}\|_{L^2((0,T) \times (0,L))} \leq \|y - y_1\|_{L^2((0,T) \times (0,L))} + \|\tilde{y}_1\|_{L^2((0,T) \times (0,L))} \\ \stackrel{(8.14), (8.18)}{\leq} C \|u\|_{L^2(\mathbb{R}_+)}^{3/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2} + C\varepsilon.$$

Multiplying the equation of y with $\Phi(t, x)$, integrating by parts on $[0, L]$, and using (1.11), we have

$$(8.21) \quad \frac{d}{dt} \int_0^L y(t, x) \Phi(t, x) dx - \frac{1}{2} \int_0^L y^2(t, x) \Phi_x(t, x) dx = 0.$$

Integrating (8.21) from 0 to T and using the fact $y(T, \cdot) = 0$ yield

$$(8.22) \quad \int_0^L y_0(x) \Phi(0, x) dx + \frac{1}{2} \int_0^T \int_0^L y^2(t, x) \Phi_x(t, x) dx dt = 0.$$

We have

$$(8.23) \quad \|y^2 - \tilde{y}^2\|_{L^1((0,T) \times (0,L))} \leq \|y - \tilde{y}\|_{L^2((0,T) \times (0,L))} \|y + \tilde{y}\|_{L^2((0,T) \times (0,L))} \\ \stackrel{(8.5), (8.20)}{\leq} C \|u\|_{L^2(\mathbb{R}_+)}^{5/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2} + C\varepsilon_0\varepsilon.$$

Combining (8.22) and (8.23) yields

$$(8.24) \quad \int_0^L y_0(x) \Phi(0, x) dx + \frac{1}{2} \int_0^T \int_0^L \tilde{y}^2(t, x) \Phi_x(t, x) dx dt \leq C \|u\|_{L^2(\mathbb{R}_+)}^{5/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2} + C\varepsilon_0\varepsilon.$$

Applying (8.1) to \tilde{y} after noting (8.19), we derive from (8.24) that, for ε_0 sufficiently small,

$$\varepsilon \|\phi\|_{L^2(0,L)}^2 + \frac{1}{2} \alpha \|u + u_1\|_{[H^{1/6}(\mathbb{R}_+)]^*}^2 \leq C \|u\|_{L^2(\mathbb{R}_+)}^{5/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2} + C \varepsilon_0 \varepsilon.$$

Using (8.17), it follows that, for sufficiently small ε_0 ,

$$\alpha \|u\|_{[H^{1/6}(\mathbb{R}_+)]^*}^2 \leq C \|u\|_{L^2(\mathbb{R}_+)}^{5/2} \|u\|_{H^{1/3}(\mathbb{R}_+)}^{1/2}.$$

We derive from Lemma A.1 that

$$\alpha \|u\|_{[H^{1/6}(0,T/2)]^*}^2 \leq C \|u\|_{L^2(0,T/2)}^{5/2} \|u\|_{H^{1/3}(0,T/2)}^{1/2}.$$

Note that

$$\|u\|_{L^2(0,T/2)} \leq C \|u\|_{[H^{1/6}(0,T/2)]^*}^{3/4} \|u\|_{H^{1/2}(0,T/2)}^{1/4}$$

and

$$\|u\|_{H^{1/3}(0,T/2)} \leq C \|u\|_{[H^{1/6}(0,T/2)]^*}^{1/4} \|u\|_{H^{1/2}(0,T/2)}^{3/4}.$$

This yields

$$\|u\|_{L^2(0,T/2)}^{5/2} \|u\|_{H^{1/3}(0,T/2)}^{1/2} \leq C \|u\|_{[H^{1/6}(0,T/2)]^*}^2 \|u\|_{H^{1/2}(0,T/2)}.$$

So, for fixed sufficiently small ε_0 ,

$$u = 0.$$

Hence $y(t, \cdot) = \varepsilon \Psi(t, \cdot) \not\equiv 0$ for all $t > 0$. We have a contradiction.

The proof is complete. □

In the proof of Theorem 8.1, we repeatedly used the following useful result.

Lemma 8.1. *Let $L > 0$ and $T > 0$ and let $f \in L^2((0, T); H^1(0, L))$ and $g \in L^2((0, T) \times (0, L))$. Then*

$$(8.25) \quad \|fg\|_{L^1((0,T);L^2(0,L))} \leq C \|f\|_{L^2((0,T)\times(0,L))}^{1/2} \|f\|_{L^2((0,T);H^1(0,L))}^{1/2} \|g\|_{L^2((0,T)\times(0,L))}$$

and

$$(8.26) \quad \|(fg)_x\|_{L^1((0,T);[H^1(0,L)]^*)} \leq C \left(\|fg\|_{L^1((0,T);L^2(0,T))} + \|(fg)(\cdot, 0)\|_{L^1(0,T)} + \|(fg)(\cdot, L)\|_{L^1(0,T)} \right),$$

for some positive constant C depending only on L . Consequently, for $f, g \in L^2((0, T); H^1(0, L))$,

$$(8.27) \quad \|fg_x\|_{L^1((0,T);L^2(0,L))} \leq C \|f\|_{L^2((0,T)\times(0,L))}^{1/2} \|f\|_{L^2((0,T);H^1(0,L))}^{1/2} \|g_x\|_{L^2((0,T)\times(0,L))}.$$

Proof. We first establish (8.25). We have

$$(8.28) \quad \|fg\|_{L^1((0,T);L^2(0,L))}^2 = \int_0^T \left(\int_0^L |fg|^2(s, x) dx \right)^{\frac{1}{2}} ds \\ \leq \int_0^T \sup_{x \in [0, L]} |f(s, x)| \left(\int_0^L |g|^2(s, x) dx \right)^{\frac{1}{2}} ds.$$

Using the fact

$$(8.29) \quad \sup_{x \in [0, L]} |f(s, x)| \leq C \left(\int_0^L |f(s, x)|^2 dx + \int_0^L |f(s, x)| |f_x(s, x)|^2 dx \right)^{\frac{1}{2}} \\ \leq C \|f(s, \cdot)\|_{L^2(0, L)}^{\frac{1}{2}} \|f(s, \cdot)\|_{H^1(0, L)}^{\frac{1}{2}},$$

we derive from (8.28) that

$$\|fg_x\|_{L^1((0, T); L^2(0, L))}^2 \leq C \int_0^T \|f(s, \cdot)\|_{L^2(0, L)}^{\frac{1}{2}} \|f(s, \cdot)\|_{H^1(0, L)}^{\frac{1}{2}} \|g(s, \cdot)\|_{L^2(0, L)} ds \\ \leq C \|f\|_{L^2((0, T) \times (0, L))}^{\frac{1}{2}} \|f\|_{L^2((0, T); H^1(0, L))}^{\frac{1}{2}} \|g\|_{L^2((0, T) \times (0, L))} \quad (\text{by Hölder's inequality}),$$

which is (8.25).

We next derive (8.26). We have, for $0 \leq t \leq T$,

$$\int_0^L (fg)_x(t, x) \varphi(x) dx = - \int_0^L fg(t, x) \varphi_x(x) dx + (fg)(t, L) \varphi(L) - (fg)(t, 0) \varphi(0) \quad \text{for } \varphi \in H^1(0, L).$$

This implies, by the trace theory applied to $\varphi \in H^1(0, L)$, for $0 \leq t \leq T$,

$$\|(fg)_x(t, \cdot)\|_{[H^1(0, L)]^*} \leq C \left(\|(fg)(t, \cdot)\|_{L^2(0, L)} + \|(fg)(t, 0)\|_{L^1(0, T)} + \|(fg)(t, L)\|_{L^1(0, T)} \right).$$

Estimate (8.26) follows.

The proof is complete. □

To derive Theorem 1.2 from Theorem 8.1, we also need to use the following simple lemma.

Lemma 8.2. *Let $T > 0$, $0 < s < 1$, $\varphi \in C^1([0, +\infty])$. Then*

$$\|\varphi v\|_{[H^s(\mathbb{R}_+)]^*} \leq C \|v\|_{[H^s(\mathbb{R}_+)]^*} \quad \text{for } v \in [H^s(\mathbb{R}_+)]^* \text{ with } \text{supp } v \subset [0, T],$$

where C is a positive constant independent of v .

Proof. Without loss of generality, one might assume that $\text{supp } \varphi \subset [0, 2T]$. Using the duality between $[H^s(\mathbb{R}_+)]^*$ and $H^s(\mathbb{R}_+)$, it suffices to prove that

$$\|\varphi V\|_{H^s(\mathbb{R}_+)} \leq C \|V\|_{H^s(\mathbb{R}_+)} \quad \text{for } V \in H^s(\mathbb{R}_+),$$

for some positive constant C independent of V . This follows immediately from the Gagliardo-Nirenberg characterization of the norm of $H^s(\mathbb{R}_+)$. □

9. LOCAL, EXACT CONTROLLABILITY OF THE KdV EQUATION IN A POSITIVE TIME - PROOF OF THEOREM 1.3

This section consists of two subsections. In the first one, we give the proof of Assertion *i*). The proof of Assertion *ii*) is given in the second subsection.

9.1. Proof of Assertion *i*) of Theorem 1.3. We first consider the case $\omega = 0$. We have

$$\begin{aligned} & \left| \sum_{j=1}^3 (\lambda_j e^{\lambda_j L} - \lambda_{j+1} e^{\lambda_{j+1} L}) e^{\lambda_{j+2} x} \right|^2 (z) \\ &= \sum_{j=1}^3 (\lambda_j e^{\lambda_j L} - \lambda_{j+1} e^{\lambda_{j+1} L} e^{\lambda_{j+2} x})(z) \sum_{j=1}^3 (\lambda_j e^{\lambda_j L} - \lambda_{j+1} e^{\lambda_{j+1} L} e^{\lambda_{j+2} x})(-\bar{z}) \end{aligned}$$

and similarly,

$$|\Xi(z)|^2 = \Xi(z)\Xi(-\bar{z}),$$

where Ξ is defined in (6.4). It follows that

$$|\Xi(\cdot + iq/2)|^2 |\Omega(\cdot + iq/2)|^2 \text{ is an analytic function in } z \in \mathbb{R}.$$

There thus exists a finite set $\{z_j; 1 \leq j \leq m\} \subset \mathbb{R}$ such that

$$\Omega(z_j + iq/2) = 0 \text{ for } 1 \leq j \leq m \quad \text{and} \quad \Omega(z + iq/2) \neq 0 \text{ for } z \in \mathbb{R} \setminus \{z_j; 1 \leq j \leq m\}.$$

For $\rho > 0$, set

$$J_\rho = \{z \in \mathbb{R}; \Omega(z + iq/2) \geq \rho\}.$$

Since $\Omega(\cdot + iq/2)$ is continuous in \mathbb{R} and $\Omega(z + iq/2) \rightarrow +\infty$ as $|z| \rightarrow +\infty$ by (7.37) of Proposition 7.1 and Lemma 3.1, we derive that

$$\delta(\rho) := |\mathbb{R} \setminus J_\rho| \rightarrow 0 \text{ as } \rho \rightarrow 0.$$

Let $T > 0$ be arbitrary and let y be a solution of (7.9) such that $y(t, \cdot) = 0$ for $t \geq T$. Applying Proposition 7.1 and taking into account the fact $\omega \geq 0$, we get

$$(9.1) \quad \int_0^{+\infty} \int_0^L |y(t, x)|^2 \Phi_x(t, x) dx dt = \int_{\mathbb{R}} \int_0^L |\hat{u}(z + iq/2)|^2 B(z + iq/2, x) dx dz,$$

where we extend u by 0 for $t < 0$ and $t > T$. This implies

$$\begin{aligned} (9.2) \quad & \int_0^{+\infty} \int_0^L |y(t, x)|^2 \Phi_x(t, x) dx dt \\ &= \int_{J_\rho} \int_0^L |\hat{u}(z + iq/2)|^2 B(z + iq/2, x) dx dz + \int_{\mathbb{R} \setminus J_\rho} \int_0^L |\hat{u}(z + iq/2)|^2 B(z + iq/2, x) dx dz \\ &\geq C_1 \rho \|ue^{q \cdot / 2}\|_{[H^{1/6}(0, T)]^*}^2 - C_2 \rho \int_{\mathbb{R} \setminus J_\rho} |\hat{u}(z + iq/2)|^2 dz \\ &\stackrel{\text{Lemma 8.2}}{\geq} C_1 \rho \|u\|_{[H^{1/6}(0, T)]^*}^2 - C_T \rho \delta(\rho) \|u\|_{[H^{1/6}(0, T)]^*}^2. \end{aligned}$$

By choosing ρ sufficiently small, we arrive that

$$(9.3) \quad \int_0^{+\infty} \int_0^L |y(t, x)|^2 \Phi_x(t, x) dx dt \geq C \rho \|u\|_{[H^{1/6}(0, T)]^*}^2.$$

Applying Theorem 8.1, we derive that system (1.1) is not locally null controllable at time $T/2$. Since T is arbitrary, the conclusion follows.

We next deal with the case $\omega > 0$. It is clear that (9.3) holds with $\rho = \omega$. The conclusion follows by Theorem 8.1 as before since T is arbitrary. \square

9.2. **Proof of Assertion *ii* of Theorem 1.3.** Set

$$(9.4) \quad \sigma(t) = e^{-\frac{1}{1-t^2}} \mathbf{1}_{(-1,1)},$$

where $\mathbf{1}_{(-1,1)}$ denotes the characteristic function of the interval $(-1, 1)$. Then $\sigma \in C_c^\infty(\mathbb{R})$ and

$$(9.5) \quad |\hat{\sigma}(z)| \leq c_1 e^{-c_2 \sqrt{|z|}} \text{ for } z \in \mathbb{R},$$

see e.g., [49]. Let $m \in \mathbb{N}$ be large and denote

$$\sigma_m(t) = \frac{1}{m} \sigma(t/m) \text{ for } t \in \mathbb{R}.$$

Let $z_0 \in \mathbb{R}$ be such that

$$\Omega(z_0 + iq/2) = \omega < 0.$$

Consider u_m defined by

$$(9.6) \quad \hat{u}_m(z) = e^{imz} \hat{\sigma}_m(z - iq/2 - z_0) \det Q(z) \Xi(z).$$

Applying Paley-Wiener's theorem, see e.g., [45], and using (9.5), one can check that the function v_m which is defined by $\hat{v}_m(z) = \hat{\sigma}_m(z - iq/2 - z_0) \det Q(z) \Xi(z)$ belongs to $L^2(\mathbb{R})$ with support in $[-m, m]$. We then derive that $u_m \in L^2(\mathbb{R})$ with support in $[0, 2m]$ since $u_m(t) = v_m(t + m)$ for $t \in \mathbb{R}$. Using (9.5), one can also derive that $u_m \in C^\infty(\mathbb{R})$.

Let $y = y_m$ be the unique solution of (6.12) with $u = u_m$. From the definition of u_m , we deduce that

$$\text{supp } y \subset [0, m] \times [0, L],$$

since, by Lemma 6.1,

$$\hat{y}_{xx}(z, 0) = \frac{\hat{u}_m(z) G(z)}{H(z)}$$

is an entire function with modulus is bounded by $Ce^{(m+\varepsilon)|z|}$ for all $\varepsilon > 0$, and $\hat{y}_{xx}(\cdot, 0) \in L^2(\mathbb{R})$. It follows that

$$\text{supp } y_{xx}(\cdot, 0) \subset [0, m].$$

Since $y(m, x) \in \mathcal{M}_D^\perp$, applying Lemma 5.3, we derive that

$$y(t, x) = 0 \text{ for } (t, x) \in [m, +\infty) \times [0, L].$$

We thus can apply Proposition 7.1 to y .

Since, by Proposition 7.1,

$$\int_0^\infty \int_0^{+\infty} |y(t, x)|^2 \Phi_x(t, x) dx dt = \int_{\mathbb{R}} \frac{|\hat{u}_m(z + iq/2)|^2}{|H(z + iq/2)|^2} \Omega(z + iq/2) dz,$$

it follows from (9.6) that

$$(9.7) \quad \int_0^\infty \int_0^{+\infty} |y(t, x)|^2 \Phi_x(t, x) dx dt = e^{-mq} \int_{\mathbb{R}} |\hat{\sigma}_m(z - z_0)|^2 |\Xi(z + iq/2)|^2 \Omega(z + iq/2) dz.$$

We have, by (9.5)

$$(9.8) \quad |\hat{\sigma}_m(z)| \leq c_1 e^{-c_2 \sqrt{m|z|}} \text{ for } z \in \mathbb{R}.$$

$$(9.9) \quad |\hat{\sigma}_m(z)| \geq c_3 \text{ for } z \in \mathbb{R} \text{ with } |z| < c_4/m,$$

$$(9.10) \quad \inf_{z \in \mathbb{R}} |\Xi(z + iq/2)| > 0, \quad |\Xi(z + iq/2)| \leq C(|z| + 1) \text{ for } z \in \mathbb{R},$$

$$(9.11) \quad |\Omega(z + iq/2)| \leq c_5 e^{c_6 |z|^{1/3}},$$

for some positive constants c_1, \dots, c_6 , independent of m .

Using (9.8), (9.9), (9.10), (9.11), and the fact $\Omega(z_0 + iq/2) = \omega < 0$, it follows from (9.7) that

$$\int_0^\infty \int_0^{+\infty} |y(t, x)|^2 \Phi_x(t, x) dx dt < 0,$$

for large m . Fix such an m and set

$$T = 2m.$$

We thus have just proved, after considering the real and/or imaginary parts of u_m and scaling, that there exists $U_1 \in H^{1/3}(0, T)$ real such that

$$(9.12) \quad Y_1(T, \cdot) = 0 \quad \text{and} \quad \text{proj}_{\mathcal{M}_D} Y_2(T, \cdot) = -\phi,$$

where $Y_1, Y_2 \in X_T$ are the unique solution of the systems

$$(9.13) \quad \begin{cases} Y_{1,t} + Y_{1,x} + Y_{1,xxx} = 0 & \text{in } (0, T) \times (0, L), \\ Y_1(\cdot, 0) = Y_{1,x}(\cdot, L) = 0 & \text{in } (0, T), \\ Y_1(\cdot, L) = U_1 & \text{in } (0, T), \\ Y_1(0, \cdot) = 0 & \text{in } (0, L), \end{cases}$$

$$(9.14) \quad \begin{cases} Y_{2,t} + Y_{2,x} + Y_{2,xxx} + Y_1 Y_{1,x} = 0 & \text{in } (0, T) \times (0, L), \\ Y_2(\cdot, 0) = Y_{2,x}(\cdot, L) = 0 & \text{in } (0, T), \\ Y_2(\cdot, L) = 0 & \text{in } (0, T). \end{cases}$$

We now establish the local controllability for the time T which is $2m$. In what follows, we consider $T = 2m$. Fix $y_0, y_T \in L^2(0, L)$ with small $L^2(0, L)$ -norms. We first consider the case

$$\rho = \|y_0\|_{L^2(0,L)} > 0,$$

$$(9.15) \quad \int_0^L y_0 \phi_1 dx \geq \mathbf{c}_1 \rho \int_0^L \phi^2 dx,$$

and

$$(9.16) \quad \|y_T\|_{L^2(0,L)} \leq \frac{\mathbf{c}_1 e^{qT} \rho}{10}.$$

for some fixed constant \mathbf{c}_1 independent of ρ .

For $r > 0$ and $y_0 \in L^2(0, L)$, denote $B_r(y_0)$ the open ball centered at 0 and of radius r in $L^2(0, L)$. We also denote $\overline{B_r(y_0)}$ its closure in $L^2(0, L)$.

Let c be a small positive constant determined later. For $\varphi \in \overline{B_{c\rho}(y_0)}$, let u_1 and u_2 be controls in $H^{1/3}(0, T)$ for which the solutions y_1 and y_2 in X_T of the systems

$$(9.17) \quad \begin{cases} y_{1,t} + y_{1,x} + y_{1,xxx} = 0 & \text{in } (0, T) \times (0, L), \\ y_1(\cdot, 0) = y_1(\cdot, L) = 0 & \text{in } (0, T), \\ y_{1,x}(\cdot, L) = u_1 & \text{in } (0, T), \\ y_1(0, \cdot) = 0 & \text{in } (0, L), \end{cases}$$

$$(9.18) \quad \begin{cases} y_{2,t} + y_{2,x} + y_{2,xxx} + y_1 y_{1,x} = 0 & \text{in } (0, T) \times (0, L), \\ y_2(\cdot, 0) = y_2(\cdot, L) = 0 & \text{in } (0, T), \\ y_{2,x}(\cdot, L) = u_2 & \text{in } (0, T), \\ y_2(0, \cdot) = \text{proj}_{\mathcal{M}_D} \varphi / \rho & \text{in } (0, L), \end{cases}$$

satisfy

$$y_1(T, \cdot) = 0 \quad \text{and} \quad y_2(T, \cdot) = y_T/\rho.$$

Moreover, one can choose u_1 and u_2 as Lipschitz functions of $\text{proj}_{\mathcal{M}_D} \varphi/\rho$ with the Lipschitz constants bounded by positive constants independent of ρ . For example, using Lemma 5.1 and (9.12), one can take ⁹

$$u_1 = (e^{-qT} \alpha_0 - \alpha_T) U_1,$$

where $\alpha_0 \phi = \text{proj}_{\mathcal{M}_D} \varphi/\rho$ and $\alpha_T \phi = \text{proj}_{\mathcal{M}_D} y_T/\rho$, and

$$u_2 = \mathcal{L}(-\text{proj}_{\mathcal{M}_D^\perp} y_3(T, \cdot) + \text{proj}_{\mathcal{M}_D^\perp} y_T/\rho) + \hat{\mathcal{L}}(\text{proj}_{\mathcal{M}_D} \varphi/\rho),$$

where $y_3 \in X_T$ is the unique solution of the system

$$(9.19) \quad \begin{cases} y_{3,t} + y_{3,x} + y_{3,xxx} + y_1 y_{1,x} = 0 & \text{in } (0, T) \times (0, L), \\ y_3(\cdot, 0) = y_3(\cdot, L) = y_{3,x}(\cdot, L) = 0 & \text{in } (0, T), \\ y_3(0, \cdot) = \text{proj}_{\mathcal{M}_D^\perp} \varphi/\rho & \text{in } (0, T). \end{cases}$$

Here \mathcal{L} and $\hat{\mathcal{L}}$ are the operators given in Proposition 5.1.

For $\varphi \in \overline{B_{c\rho}(y_0)}$, let $y \in X_T$ be the unique solution of the *backward* linear KdV system

$$(9.20) \quad \begin{cases} y_t + y_x + y_{xxx} + y y_x = 0 & \text{in } (0, T) \times (0, L), \\ y(\cdot, 0) = y_x(\cdot, L) = 0 & \text{in } (0, T), \\ y_x(\cdot, 0) = v & \text{in } (0, T), \\ y(T, \cdot) = y_T & \text{in } (0, L). \end{cases}$$

where $v(t) = \rho^{1/2} y_{1,x}(t, 0) + \rho y_{2,x}(t, 0)$ with y_1 and y_2 being defined by (9.17) and (9.18).

Note that $y \in X_T$ is well-defined if ρ is sufficiently small by Proposition 3.5. We will denote

$$(9.21) \quad H_0(\varphi) = y(0, \cdot) \text{ in } (0, L) \text{ and } H(\varphi) = y \text{ in } (0, T) \times (0, L).$$

Consider the map

$$\begin{aligned} \Lambda: \overline{B_{c\rho}(y_0)} &\rightarrow L^2(0, L) \\ \varphi &\mapsto \varphi - H_0(\varphi) + y_0. \end{aligned}$$

We will prove that

$$(9.22) \quad \Lambda(\varphi) \in \overline{B_{c\rho}(y_0)},$$

and

$$(9.23) \quad \|\Lambda(\varphi) - \Lambda(\tilde{\varphi})\|_{L^2(0, L)} \leq \lambda \|\varphi - \tilde{\varphi}\|_{L^2(0, L)},$$

for some $\lambda \in (0, 1)$. Assuming this, one derives from the contraction mapping theorem that there exists a unique $\varphi_0 \in \overline{B_{c\rho}(y_T)}$ such that $\Lambda(\varphi_0) = \varphi_0$. As a consequence, with $y = H(\varphi_0)$,

$$y(0, \cdot) = y_0,$$

and $y(\cdot, L)$ is hence a required control.

We next establish (9.22) and (9.23). Indeed, assertion (9.22) follows from the fact

$$\|\varphi - H_0(\varphi)\|_{L^2(0, L)} \leq C \|\varphi\|_{L^2(0, L)}^{3/2} \text{ for } \varphi \in \overline{B_{c\rho}(y_0)}.$$

This can be proved using the approximation via the power series method as follows. Applying Proposition 3.3 to y , we derive from (9.20) that

$$(9.24) \quad \|y\|_{X_T} \leq C \rho^{1/2}.$$

⁹It is useful to note that $e^{-qT} \alpha_0 - \alpha_T > 0$ by (9.15) and (9.16) if c is sufficiently small.

Set ¹⁰

$$(9.25) \quad y_a = \rho^{1/2}y_1 + \rho y_2 \text{ in } (0, T) \times (0, L).$$

We have

$$(9.26) \quad (y - y_a)_t + (y - y_a)_x + (y - y_a)_{xxx} + (y - y_a)y_x + y_a(y - y_a)_x = h(t, x),$$

where

$$h(t, x) = -y_a y_{a,x} + \rho y_1 y_{1,x} = -\rho^{3/2}(y_1 y_2)_x - \rho^2 y_2 y_{2,x}.$$

Applying Proposition 3.3 to $y - y_a$, for small ρ , one can ignore (absorb) the contribution from the last two terms in the LHS of (9.26), to obtain

$$(9.27) \quad \|y - y_a\|_{X_T} \leq C \|h\|_{L^1((0,T);L^2(0,L))} \leq C \rho^{3/2}.$$

Assertion (9.22) follows since $H_0(\varphi) = y(0, \cdot)$ and $\varphi = y_a(0, \cdot)$.

We next establish (9.23). To this end, we estimate

$$\left(\varphi - H_0(\varphi) \right) - \left(\tilde{\varphi} - H_0(\tilde{\varphi}) \right).$$

For $\tilde{\varphi} \in \overline{B_{c\rho}(y_0)}$, denote \tilde{u}_1, \tilde{u}_2 and $\tilde{y}_1, \tilde{y}_2, \tilde{y}_a, \tilde{y}$ the corresponding functions which are defined in the same way as the functions u_1, u_2 , and y_1, y_2, y_a, y considered for φ .

We have

$$(9.28) \quad (y - \tilde{y})_t + (y - \tilde{y})_x + (y - \tilde{y})_{xxx} + y y_x - \tilde{y} \tilde{y}_x = 0,$$

$$(9.29) \quad (y_a - \tilde{y}_a)_t + (y_a - \tilde{y}_a)_x + (y_a - \tilde{y}_a)_{xxx} + y_a y_{a,x} - \tilde{y}_a \tilde{y}_{a,x} = g(t, x),$$

where

$$(9.30) \quad g(t, x) = \rho^{3/2} \left((y_1 y_2)_x - (\tilde{y}_1 \tilde{y}_2)_x \right) + \rho^2 \left(y_2 y_{2,x} - \tilde{y}_2 \tilde{y}_{2,x} \right).$$

Applying Proposition 3.5 to $y - \tilde{y}$, for ρ small, one can ignore (absorb) the contribution of the last two terms in the LHS of (9.28), to obtain

$$(9.31) \quad \|y - \tilde{y}\|_{X_T} \leq C \|y_{a,x}(\cdot, 0) - \tilde{y}_{a,x}(\cdot, 0)\|_{H^{1/3}(0,T)} \leq C \rho^{-1/2} \|\varphi - \tilde{\varphi}\|_{L^2(0,L)}.$$

Using (9.27) for $y - y_a$ and similar fact for $\tilde{y} - \tilde{y}_a$, we obtain

$$(9.32) \quad \|(y - y_a, \tilde{y} - \tilde{y}_a)\|_{X_T} \leq C \rho^{3/2}.$$

From the definition of g in (9.30), we deduce that

$$(9.33) \quad \|g\|_{L^1((0,T);L^2(0,L))} \leq C \rho^{1/2} \|\varphi - \tilde{\varphi}\|_{L^2(0,L)}.$$

Applying Proposition 3.1 to $y_a - \tilde{y}_a$ after ignoring (absorbing) the contribution of $y_a y_{a,x} - \tilde{y}_a \tilde{y}_{a,x}$, we derive from (9.29) that

$$(9.34) \quad \|y_a - \tilde{y}_a\|_{X_T} \leq C \rho^{-1/2} \|\varphi - \tilde{\varphi}\|_{L^2(0,L)} \leq C \rho^{1/2}.$$

Set

$$Y = y - y_a - (\tilde{y} - \tilde{y}_a) \text{ in } (0, T) \times (0, L).$$

Using (9.28) and (9.29), we have

$$(9.35) \quad \partial_t Y + \partial_x Y + \partial_{xxx} Y = f(t, x) \text{ in } (0, T) \times (0, L),$$

where

$$f(t, x) = -g(t, x) - \left(y y_x - \tilde{y} \tilde{y}_x - (y_a y_{a,x} - \tilde{y}_a \tilde{y}_{a,x}) \right).$$

¹⁰The index a stands the approximation.

From (9.31), (9.32), (9.34), and (9.33), we obtain

$$(9.36) \quad \|yy_x - \tilde{y}\tilde{y}_x - (y_a y_{a,x} - \tilde{y}_a \tilde{y}_{a,x})\|_{L^1((0,T);L^2(0,L))} \leq C\rho^{1/2}\|\varphi - \tilde{\varphi}\|_{L^2(0,L)}.$$

Using (9.33) and (9.36), and applying Proposition 3.5 to Y , we derive from (9.35) that

$$\|(y - y_a - \tilde{y} + \tilde{y}_a)(T, \cdot)\|_{L^2(0,L)} \leq C\rho^{1/2}\|\varphi - \phi\|_{L^2(0,L)}.$$

Assertion (9.23) follows.

We next consider the general case. One can bring this case to the previous case as follows. Set

$$\rho = \|y_0\|_{L^2(0,L)} + \|y_T\|_{L^2(0,L)}.$$

Without loss of generality, one might assume that $\rho > 0$ since otherwise, one just takes zero as a control.

Fix $\varepsilon > 0$ small. By Proposition 7.1, there exists $v_1 \in H^{1/3}(0, \varepsilon)$ such that if $y_1 \in X_\varepsilon$ is the solution of (9.17) with $y_1(\cdot, L) = v_1$ and $y_2 \in X_\varepsilon$ is the solution of (9.18) with $y_2(\cdot, L) = 0$ then

$$y_1(\varepsilon, \cdot) = 0, \quad \text{and} \quad \int_0^L y_2(\varepsilon, \cdot)\phi \geq c.$$

Set

$$(9.37) \quad u = \rho^{1/2}\gamma^{1/2}v_1 \text{ in } (0, \varepsilon).$$

Then the unique solution $y \in X_\varepsilon$ of the KdV system (1.1) verifies the properties of the previous case, for large positive γ . We are now in the position to apply the previous case with the initial datum $y(\varepsilon, \cdot)$. The proof is complete. \square

Remark 9.1. This paper gives a complete picture of the local controllability properties of the KdV equation with the right Dirichlet controls. The question of whether or not a large control might suffice for small-time controllability, even in the small data case, is a very interesting open question.

APPENDIX A. HARDY TYPE INEQUALITIES

This section is devoted to the following result related to the Hardy inequality.

Lemma A.1. *Let $-\infty < a < b < c \leq +\infty$, and $0 < s < 1/2$ and let $u \in H^s(a, b)$. Set*

$$v = \begin{cases} u & \text{in } (a, b), \\ 0 & \text{in } (b, c). \end{cases}$$

Then $v \in H^s(a, c)$ and

$$\|v\|_{H^s(a,c)} \leq C\|u\|_{H^s(a,b)},$$

for some positive constant C depending only on a, b, c , and s .

Proof. For notational ease, we assume that $a = -1$ and $b = 0$. Without loss of generality, we then can assume that $c = +\infty$. Let $V \in H^s(\mathbb{R})$ be an extension of u such that

$$\|V\|_{H^s(\mathbb{R})} \leq C\|u\|_{H^s(-1,0)}.$$

Applying [40, Theorem 1.1] with $\gamma = -s$, $\tau = 2$, $p = 2$ to V , one obtains

$$\||x|^{-s}V\|_{L^2(\mathbb{R})} \leq C\|V\|_{H^s(\mathbb{R})}.$$

The condition $s < 1/2$ is required here. This yields

$$(A.1) \quad \||x|^{-s}u\|_{L^2(-1,0)} \leq C\|u\|_{H^s(-1,0)}.$$

Using the equivalent Gagliardo-Nirenberg definition of the semi-norm H^s , we have

$$\|v\|_{H^s(-1,+\infty)}^2 \sim \int_{-1}^{\infty} \int_{-1}^{\infty} \frac{|v(x) - v(y)|^2}{|x - y|^{1+2s}} dx dy + \int_{-1}^{\infty} |v|^2 dx.$$

Since

$$\int_{-1}^{\infty} \int_{-1}^{\infty} \frac{|v(x) - v(y)|^2}{|x - y|^{1+2s}} dx dy \leq \int_{-1}^0 \int_{-1}^0 \frac{|u(x) - u(y)|^2}{|x - y|^{1+2s}} dx dy + C_s \int_{-1}^0 \frac{|u(x)|^2}{|x|^{2s}} dx,$$

and

$$\int_{-1}^{\infty} |v|^2 dx = \int_{-1}^0 |u|^2 dx,$$

we derive from (A.1) that

$$\|v\|_{H^s(-1,+\infty)} \leq C \|u\|_{H^s(-1,0)}.$$

The proof is complete. □

APPENDIX B. ON THE ZEROS OF $\det Q$

We begin this section with

Lemma B.1. *Let $L > 0$, $z \in \mathbb{C}$, and $\varphi \in C^\infty([0, L])$ be such that*

$$\varphi_{xxx} + \varphi_x + iz\varphi = 0 \text{ in } [0, L],$$

and

$$\varphi(0) = \varphi(L) = \varphi_x(L) = 0.$$

Then

$$\Im(z) \geq 0.$$

Proof. Set

$$\Psi(t, x) = \Re \left\{ e^{izt} \varphi(x) \right\} \text{ in } \mathbb{R}_+ \times [0, L].$$

Then Ψ is a solution of the linearized KdV system

$$(B.1) \quad \begin{cases} \Psi_t + \Psi_x + \Psi_{xxx} = 0 & \text{in } \mathbb{R}_+ \times [0, L], \\ \Psi(\cdot, 0) = \Psi(\cdot, L) = \Psi_x(\cdot, L) = 0 & \text{in } \mathbb{R}_+. \end{cases}$$

We then derive that

$$\frac{d}{dt} \int_0^L |\Psi(t, x)|^2 dx \leq 0.$$

This implies

$$\Im(z) \geq 0.$$

The proof is complete. □

The following result is useful.

Lemma B.2. *Let $L \in \mathcal{N}_D$. Then $\det Q(z + iq/2) \neq 0$ for $z \in \mathbb{R}$.*

Proof. We prove the assertion by contradiction. Assume that $\det Q(z + iq/2) = 0$ for some $z \in \mathbb{R}$. Then there exists $\varphi \in C^\infty[0, L]$ such that

$$\varphi_{xxx} + \varphi_x + i(z + iq/2)\varphi = 0 \text{ in } [0, L],$$

and

$$\varphi(0) = \varphi(L) = \varphi_x(L) = 0.$$

Applying Lemma B.1, we have

$$\Im(z + iq/2) = q/2 \leq 0.$$

We have a contradiction. □

APPENDIX C. A LEMMA RELATED TO THE MOMENT METHOD

The following lemma from [24] is used in the proof of Proposition 3.1.

Lemma C.1. *Let φ be an analytic function in \mathbb{C} such that φ has a finite number of zeros on the real line, and*

$$(C.1) \quad |\varphi(z)| \leq c_1 e^{c_2 |z|^\alpha} \text{ in } \mathbb{C},$$

for some $0 < \alpha < 1$, and $c_1, c_2 > 0$. Let $T_1, T_2 > 0$, and $h \in H^s(\mathbb{R})$ for some $s \leq 0$ with support in $(0, T_1)$. There exists $g \in C^\infty(\mathbb{R})$ with support in $[T_1, T_1 + T_2]$ such that if z is a real solution of order m of the equation $\varphi(z) = 0$, then z is also a real solution of order m of the equation $\hat{h} - \hat{g} = 0$, and

$$(C.2) \quad \|g\|_{H^k(\mathbb{R})} \leq C_k \|h\|_{H^s(\mathbb{R})} \text{ for } k \in \mathbb{N},$$

for some positive constant C_k depending only on k, T_1, T_2, s , and real zeros and their multiplicity of φ .

Proof. The proof of Lemma C.1 is as in [24], where a special case is considered. The construction of g , inspired by the moment method, see e.g. [49], can be done as follows. Set $\eta(t) = e^{-1/(t^2 - (T_2/2)^2)} \mathbf{1}_{|t| < T_2}$ for $t \in \mathbb{R}$. Assume that z_1, \dots, z_k are real, distinct solutions of the equation $\varphi(z) = 0$, and m_1, \dots, m_k are the corresponding orders. Set, for $z \in \mathbb{C}$,

$$\zeta(z) = \sum_{i=1}^k \left(\hat{\eta}(z - z_i) \prod_{\substack{j=1 \\ j \neq i}}^k (z - z_j)^{m_j} \left(\sum_{l=0}^{m_i} c_{i,l} (z - z_i)^l \right) \right),$$

where $c_{i,l} \in \mathbb{C}$ is chosen such that

$$\frac{d^l}{dz^l} \left(e^{i(T_1 + T_2/2)z} \zeta(z) \right)_{z=z_i} = \frac{d^l}{dz^l} \hat{h}_3(z_i) \text{ for } 0 \leq l \leq m_i, 1 \leq i \leq k.$$

This can be done since $\hat{\eta}(0) \neq 0$. Since

$$|\hat{\eta}(z)| \leq C e^{T_2 |\Im(z)|/2},$$

and, by [49, Lemma 4.3],

$$|\hat{\eta}(z)| \leq C_1 e^{-C_2 |z|^{1/2}} \text{ for } z \in \mathbb{R},$$

using (C.1), and applying Paley-Wiener's theorem, one can prove that ζ is the Fourier transform of a function ψ of class C^1 ; moreover, ψ has the support in $[-T_2/2, T_2/2]$. Set, for $z \in \mathbb{C}$,

$$g(t) = \psi(t + T_1 + T_2/2).$$

Using the fact $\hat{g}(z) = e^{i(T_1 + T_2/2)z} \zeta(z)$, one can check that $\hat{g} - \hat{h}$ has zeros z_1, \dots, z_k with the corresponding orders m_1, \dots, m_k . One can check that

$$\|\psi\|_{H^k(\mathbb{R})} \leq C_{T,L,k} \sum_{i=1}^k \sum_{l=0}^{m_i} \left| \frac{d^l}{dz^l} \hat{h}(z_i) \right|,$$

which yields

$$\|\psi\|_{H^k(\mathbb{R})} \leq C_{T,L,k} \|h\|_{H^s(\mathbb{R})}.$$

The required properties of g follow. □

APPENDIX D. SCILAB PROGRAM FOR COMPUTING L_n AND CHECKING THE LOCAL
CONTROLLABILITY PROPERTY FOR $n = 0, 1, 2, 3$.

Here is the Scilab program which gives the results in Remark 1.4.

```

clc
a=0;
b=0;
L=0;
q=0;
alpha=0;
beta=0;
value=0;

function [B] = B(x)
    B = (4*x^2*%e^(-6*x) - x^2)^(1/2);
endfunction

function [F] = F(x)
    F = B(x) * cos (B(x)) + x* sin (B(x));
endfunction

function [l] = l(x)
    l = (B(x)*B(x) - 3*x^2)^(1/2);
endfunction

//k=0; j=-0.44; j0=j;
//Here are the outcomes for n=0;
//value=0.9770641; L=4.5183604; a=-0.5065520; b=4.6027563; q=0.2354919;

//k=1; j=-0.67; j0=j;
// // Here are the outcomes for n=1:
// // value=3.8760737; L=10.866906; a=-0.6903700; b=10.932497; q=0.1291104;

k=2; j=-0.78; j0=j;
//// Here are the outcomes for n=2:
//// value=-0.4687287; L=17.177525; a=-0.7947960; b=17.232599; q=0.0933315;

//k=3; j=-0.86; j0=j;
//// Here are the outcomes for n=3:
//// value=-0.4687287; L=23.476776; a=-0.8687610; b=23.524949; q=0.0744156;

while j > -2
    j=j - 10^(-6);
    if (B(j)> %pi + 2*k*%pi) & (B(j)<3*%pi/2+ 2*k*%pi) & (F(j)<0)

```

```

    then j0=j0-10^(-6);
    else
        j=-3;
    end
end
end

a=j0
b=B(a)
L=l(a)
q= - 2*a*(a^2 + b^2)/ L^3
alpha=-a/L
beta=-b/L

function [p] = phix(x)
    p = -alpha*beta*%e^(alpha*x)*cos(beta*x) ..
    + beta^2*%e^(alpha*x)*sin(beta*x)..
    - 2*alpha*beta*%e^(-2*alpha*x) ..
    + 3*alpha^2*%e^(alpha*x)*sin(beta*x)..
    + 3*alpha*beta*%e^(alpha*x)*cos(beta*x);
endfunction
function [p] = P(x)
    p = inv_coeff([-q/2 + %i*x 1 0 1]);
endfunction

function [g] = G(z,x)
    lambda= roots(P(z));
    g = (lambda(1)*%e^(lambda(1)*L) ..
    - lambda(2)*%e^(lambda(2)*L))*%e^(lambda(3)*x)..
    + (lambda(2)*%e^(lambda(2)*L) ..
    - lambda(3)*%e^(lambda(3)*L))*%e^(lambda(1)*x)..
    +(lambda(3)*%e^(lambda(3)*L) ..
    - lambda(1)*%e^(lambda(1)*L))*%e^(lambda(2)*x);
endfunction

function [h] = H(z,x)
    h = abs(G(z,x))*abs(G(z,x))*phix(x);
endfunction

function [omega]=Omega(z)
    function [xi] = Xi(x)
        xi = abs(G(z,x))*abs(G(z,x))*phix(x);
    endfunction
    omega = intg(0,L,Xi);
end

value=Omega(0.1897656) //k=2

```

//value=Omega(0.1419531) //k=3

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REFERENCES

- [1] Claude Bardos, Gilles Lebeau, and Jeffrey Rauch, *Sharp sufficient conditions for the observation, control, and stabilization of waves from the boundary*, SIAM J. Control Optim. **30** (1992), no. 5, 1024–1065. MR 1178650
- [2] K. Beauchard, P. Cannarsa, and R. Guglielmi, *Null controllability of Grushin-type operators in dimension two*, J. Eur. Math. Soc. (JEMS) **16** (2014), no. 1, 67–101. MR 3141729
- [3] Karine Beauchard, Jérémie Dardé, and Sylvain Ervedoza, *Minimal time issues for the observability of Grushin-type equations*, Ann. Inst. Fourier (Grenoble) **70** (2020), no. 1, 247–312. MR 4105940
- [4] Karine Beauchard, Bernard Helffer, Raphael Henry, and Luc Robbiano, *Degenerate parabolic operators of Kolmogorov type with a geometric control condition*, ESAIM Control Optim. Calc. Var. **21** (2015), no. 2, 487–512. MR 3348409
- [5] Karine Beauchard and Camille Laurent, *Local controllability of 1D linear and nonlinear Schrödinger equations with bilinear control*, J. Math. Pures Appl. (9) **94** (2010), no. 5, 520–554. MR 2732927
- [6] Karine Beauchard and Frédéric Marbach, *Unexpected quadratic behaviors for the small-time local null controllability of scalar-input parabolic equations*, J. Math. Pures Appl. (9) **136** (2020), 22–91. MR 4076969
- [7] Karine Beauchard, Luc Miller, and Morgan Morancey, *2D Grushin-type equations: minimal time and null controllable data*, J. Differential Equations **259** (2015), no. 11, 5813–5845. MR 3397310
- [8] Karine Beauchard and Morgan Morancey, *Local controllability of 1D Schrödinger equations with bilinear control and minimal time*, Math. Control Relat. Fields **4** (2014), no. 2, 125–160. MR 3167929
- [9] Assia Benabdallah, Franck Boyer, and Morgan Morancey, *A block moment method to handle spectral condensation phenomenon in parabolic control problems*, Ann. H. Lebesgue **3** (2020), 717–793. MR 4149825
- [10] Jerry L. Bona, Shu Ming Sun, and Bing-Yu Zhang, *A nonhomogeneous boundary-value problem for the Korteweg-de Vries equation posed on a finite domain*, Comm. Partial Differential Equations **28** (2003), no. 7-8, 1391–1436. MR 1998942
- [11] Jerry L. Bona, Shu Ming Sun, and Bing-Yu Zhang, *A non-homogeneous boundary-value problem for the Korteweg-de Vries equation posed on a finite domain. II*, J. Differential Equations **247** (2009), no. 9, 2558–2596. MR 2568064
- [12] Jean Bourgain, *Fourier transform restriction phenomena for certain lattice subsets and applications to nonlinear evolution equations. II. The KdV-equation*, Geom. Funct. Anal. **3** (1993), no. 3, 209–262. MR 1215780
- [13] Mégane Bournois, *Quadratic behaviors of the 1D linear Schrödinger equation with bilinear control*, J. Differential Equations **351** (2023), 324–360. MR 4542546
- [14] Joseph Boussinesq, *Essai sur la théorie des eaux courantes*, Mémoires présentés par divers savants à l’Acad. des Sci. Inst. Nat. France, XXIII, pp. 1–680 (1877), 1–680.
- [15] Roberto A. Capistrano-Filho, Ademir F. Pazoto, and Lionel Rosier, *Internal controllability of the Korteweg-de Vries equation on a bounded domain*, ESAIM Control Optim. Calc. Var. **21** (2015), no. 4, 1076–1107. MR 3395756
- [16] Eduardo Cerpa, *Exact controllability of a nonlinear Korteweg-de Vries equation on a critical spatial domain*, SIAM J. Control Optim. **46** (2007), no. 3, 877–899. MR 2338431
- [17] Eduardo Cerpa, *Control of a Korteweg-de Vries equation: a tutorial*, Math. Control Relat. Fields **4** (2014), no. 1, 45–99. MR 3191303
- [18] Eduardo Cerpa and Emmanuelle Crépeau, *Boundary controllability for the nonlinear Korteweg-de Vries equation on any critical domain*, Ann. Inst. H. Poincaré Anal. Non Linéaire **26** (2009), no. 2, 457–475. MR 2504039
- [19] Jixun Chu, Jean-Michel Coron, and Peipei Shang, *Asymptotic stability of a nonlinear Korteweg-de Vries equation with critical lengths*, J. Differential Equations **259** (2015), no. 8, 4045–4085. MR 3369271
- [20] J. E. Colliander and C. E. Kenig, *The generalized Korteweg-de Vries equation on the half line*, Comm. Partial Differential Equations **27** (2002), no. 11-12, 2187–2266. MR 1944029
- [21] Jean-Michel Coron, *On the small-time local controllability of a quantum particle in a moving one-dimensional infinite square potential well*, C. R. Math. Acad. Sci. Paris **342** (2006), no. 2, 103–108. MR 2193655
- [22] Jean-Michel Coron, *Control and nonlinearity*, Mathematical Surveys and Monographs, vol. 136, American Mathematical Society, Providence, RI, 2007. MR 2302744
- [23] Jean-Michel Coron and Emmanuelle Crépeau, *Exact boundary controllability of a nonlinear KdV equation with critical lengths*, J. Eur. Math. Soc. (JEMS) **6** (2004), no. 3, 367–398. MR 2060480

- [24] Jean-Michel Coron, Armand Koenig, and Hoai-Minh Nguyen, *On the small-time local controllability of a KdV system for critical lengths*, J. Eur. Math. Soc. (2022), doi:10.4171/JEMS/1307.
- [25] Jean-Michel Coron, Armand Koenig, and Hoai-Minh Nguyen, *Lack of local controllability for a water-tank system when the time is not large enough*, (2022).
- [26] Olivier Glass and Sergio Guerrero, *Some exact controllability results for the linear KdV equation and uniform controllability in the zero-dispersion limit*, Asymptot. Anal. **60** (2008), no. 1-2, 61–100. MR 2463799
- [27] Olivier Glass and Sergio Guerrero, *Controllability of the Korteweg-de Vries equation from the right Dirichlet boundary condition*, Systems Control Lett. **59** (2010), no. 7, 390–395. MR 2724598
- [28] Justin Holmer, *The initial-boundary value problem for the Korteweg-de Vries equation*, Comm. Partial Differential Equations **31** (2006), no. 7-9, 1151–1190. MR 2254610
- [29] Tosio Kato, *On the Cauchy problem for the (generalized) Korteweg-de Vries equation*, Studies in applied mathematics, Adv. Math. Suppl. Stud., vol. 8, Academic Press, New York, 1983, pp. 93–128. MR 759907
- [30] Armand Koenig, *Non-null-controllability of the Grushin operator in 2D*, C. R. Math. Acad. Sci. Paris **355** (2017), no. 12, 1215–1235. MR 3730500
- [31] Diederik J. Korteweg and Gustave de Vries, *On the change of form of long waves advancing in a rectangular canal, and on a new type of long stationary waves*, Philos. Mag. (5) **39** (1895), no. 240, 422–443. MR 3363408
- [32] Camille Laurent, Lionel Rosier, and Bing-Yu Zhang, *Control and stabilization of the Korteweg-de Vries equation on a periodic domain*, Comm. Partial Differential Equations **35** (2010), no. 4, 707–744. MR 2753618
- [33] Felipe Linares and Gustavo Ponce, *Introduction to nonlinear dispersive equations*, second ed., Universitext, Springer, New York, 2015. MR 3308874
- [34] Frédéric Marbach, *An obstruction to small-time local null controllability for a viscous Burgers' equation*, Ann. Sci. Éc. Norm. Supér. (4) **51** (2018), no. 5, 1129–1177. MR 3942039
- [35] Philippe Martin, Ivonne Rivas, Lionel Rosier, and Pierre Rouchon, *Exact controllability of a linear Korteweg-de Vries equation by the flatness approach*, SIAM J. Control Optim. **57** (2019), no. 4, 2467–2486. MR 3981376
- [36] Gustavo Alberto Perla Menzala, Carlos Frederico Vasconcellos, and Enrique Zuazua, *Stabilization of the Korteweg-de Vries equation with localized damping*, Quart. Appl. Math. **60** (2002), no. 1, 111–129. MR 1878262
- [37] Robert M. Miura, *The Korteweg-de Vries equation: a survey of results*, SIAM Rev. **18** (1976), no. 3, 412–459. MR 404890
- [38] Luc Molinet and Francis Ribaud, *On the low regularity of the Korteweg-de Vries-Burgers equation*, Int. Math. Res. Not. (2002), no. 37, 1979–2005. MR 1918236
- [39] Hoai-Minh Nguyen, *Decay for the nonlinear KdV equations at critical lengths*, J. Differential Equations **295** (2021), 249–291. MR 4268729
- [40] Hoai-Minh Nguyen and Marco Squassina, *Fractional Caffarelli-Kohn-Nirenberg inequalities*, J. Funct. Anal. **274** (2018), no. 9, 2661–2672. MR 3771839
- [41] Ademir Fernando Pazoto, *Unique continuation and decay for the Korteweg-de Vries equation with localized damping*, ESAIM Control Optim. Calc. Var. **11** (2005), no. 3, 473–486. MR 2148854
- [42] Lionel Rosier, *Exact boundary controllability for the Korteweg-de Vries equation on a bounded domain*, ESAIM Control Optim. Calc. Var. **2** (1997), 33–55. MR 1440078
- [43] Lionel Rosier, *Control of the surface of a fluid by a wavemaker*, ESAIM Control Optim. Calc. Var. **10** (2004), no. 3, 346–380. MR 2084328
- [44] Lionel Rosier and Bing-Yu Zhang, *Control and stabilization of the Korteweg-de Vries equation: recent progresses*, J. Syst. Sci. Complex. **22** (2009), no. 4, 647–682. MR 2565262
- [45] Walter Rudin, *Real and complex analysis*, third ed., McGraw-Hill Book Co., New York, 1987. MR 924157
- [46] David L. Russell and Bing Yu Zhang, *Exact controllability and stabilizability of the Korteweg-de Vries equation*, Trans. Amer. Math. Soc. **348** (1996), no. 9, 3643–3672. MR 1360229
- [47] Shuxia Tang, Jixun Chu, Peipei Shang, and Jean-Michel Coron, *Asymptotic stability of a Korteweg-de Vries equation with a two-dimensional center manifold*, Adv. Nonlinear Anal. **7** (2018), no. 4, 497–515. MR 3871418
- [48] Terence Tao, *Nonlinear dispersive equations*, CBMS Regional Conference Series in Mathematics, vol. 106, Published for the Conference Board of the Mathematical Sciences, Washington, DC; by the American Mathematical Society, Providence, RI, 2006, Local and global analysis. MR 2233925
- [49] Gerald Tenenbaum and Marius Tucsnak, *New blow-up rates for fast controls of Schrödinger and heat equations*, J. Differential Equations **243** (2007), no. 1, 70–100. MR 2363470
- [50] Gerald Beresford Whitham, *Linear and nonlinear waves*, Wiley-Interscience [John Wiley & Sons], New York-London-Sydney, 1974, Pure and Applied Mathematics. MR 0483954

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