

Exact controllability for a Rayleigh beam with piezoelectric actuator

Yubo Bai, Christophe Prieur, Zhiqiang Wang

▶ To cite this version:

Yubo Bai, Christophe Prieur, Zhiqiang Wang. Exact controllability for a Rayleigh beam with piezo-electric actuator. Fudan University; Gipsa-Lab. 2023. hal-03964846v1

HAL Id: hal-03964846 https://hal.science/hal-03964846v1

Submitted on 31 Jan 2023 (v1), last revised 24 Apr 2023 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

EXACT CONTROLLABILITY FOR A RAYLEIGH BEAM WITH PIEZOELECTRIC ACTUATOR*

YUBO BAI[†], CHRISTOPHE PRIEUR[‡], AND ZHIQIANG WANG[§]

Abstract. In this paper, exact controllability for a Rayleigh beam with piezoelectric actuator is considered. Controllability results show that the space of controllable initial data depends on the regularity of the control function and the location of the actuator. Two different spaces of control function, $L^2(0,T)$ and $(H^1(0,T))'$, correspond to two different controllability, L^2 -controllability and $(H^1)'$ -controllability. The approach to prove controllability results is based on Hilbert Uniqueness Method. Some non-controllability results are also obtained. In particular, non-controllability in short control time is studied by using Riesz basis property of exponential family in $L^2(0,T)$. Due to this non-controllability result, minimal time for the exact controllability is obtained.

Key words. exact controllability, Rayleigh beam, piezoelectric actuator

MSC codes. 93C20, 93B05, 35B65

1 2

1. Introduction and main results.

1.1. History and problem statement. In recent decades, there have been a large number of papers concerning the study of flexible structures. Three main directions of research can be considered, i.e., the modelling problem, the controllability problem and the stabilization problem. Modelling a flexible structure as a beam equation or a plate equation is an essential research field. In [11], the author summarized and analysed four types of model for the transversely vibrating uniform beam, i.e., the Euler-Bernoulli beam, Rayleigh beam, shear beam and Timoshenko beam. In the past few decades, the study of elastic structures with a piezoelectric actuator or sensor has gained a lot of attention. The paper on modelling elastic structures with a piezoelectric actuator or sensor as a PDE can refer to [7, 9].

Concerning controllability for PDEs, [17] proposed an important method, Hilbert Uniqueness Method (HUM), to study controllability for infinite-dimensional systems. Using HUM, boundary controllability for wave equation and plate equation was studied in [17]. There were plenty of works on controllability for beam and plate based on HUM. In [13, 17], boundary controllability for Kirchhoff plate equation was fully investigated. Exact controllability was obtained in sufficiently large control time with a single boundary control (active on a sufficiently large portion of the boundary) in the case of clamped boundary conditions. As for Euler-Bernoulli beam equation, in 1996, [21] firstly considered the exact controllability for Euler-Bernoulli beam hinged at both ends with piezoelectric actuator. Since the space dimension is one, Fourier series was used in [21]. Then [8] studied the exact controllability for the same beam equation with piezoelectric actuator in a different physical configuration: the clamped-free boundary conditions, i.e. a beam clamped at one end and free at the other end. In [19], Ingham inequality (see [4, 12]) was used to obtain the exact controllability for Rayleigh beam equation with a single boundary control among four different boundary

^{*}Submitted to the editors DATE.

[†]School of Mathematical Sciences, Fudan University, Shanghai 200433, China (ybbai21@m.fudan.edu.cn).

[‡]Univ. Grenoble Alpes, CNRS, Grenoble-INP, GIPSA-lab, F-38000, Grenoble, France (christophe.prieur@gipsa-lab.fr).

[§]School of Mathematical Sciences and Shanghai Key Laboratory for Contemporary Applied Mathematics, Fudan University, Shanghai 200433, China (wzq@fudan.edu.cn).

42

43

44

49

54

56

58

59

60

62

63

64

65

66

67 68

69 70

71

72

73

74 75

78

81

conditions. 40

> In this paper, we consider the control problem modelling the transverse deflection of a Rayleigh beam which is subject to the action of an attached piezoelectric actuator. If we suppose that the beam is hinged at both ends, the equation of Rayleigh beam can be written as (see, for instance, [7, 9])

(1.1a)

45
$$w_{tt}(x,t) - \alpha w_{xxtt}(x,t) + w_{xxxx}(x,t) = u(t) \frac{\mathrm{d}}{\mathrm{d}x} [\delta_{\eta}(x) - \delta_{\xi}(x)], \quad 0 < x < \pi, \ t > 0,$$

46 (1.1b)
$$w(0,t) = w(\pi,t) = w_{xx}(0,t) = w_{xx}(\pi,t) = 0, \quad t > 0,$$

$$45 (1.1c) w(x,0) = w^{0}(x), w_{t}(x,0) = w^{1}(x), 0 < x < \pi.$$

In the equations above w represents the transverse deflection of the beam, $\alpha > 0$ is a physical constant, ξ and η stands for the ends of the actuator $(0 < \xi < \eta < \pi)$, and δ_{η} is the Dirac mass at the point y. The control is given by the function $u:[0,T]\to\mathbb{R}$ standing for the time variation of the voltage applied to the actuator.

Our main purpose is to find the initial data that can be steered to rest by means of the control function u. Here we give some precise definitions.

DEFINITION 1.1. The initial data (w^0, w^1) is exactly L^2 -controllable in (ξ, η) at time T if there exists u in $L^2(0,T)$ such that the solution w of (1.1) satisfies the condition

$$w(x,T) = w_t(x,T) = 0, \quad 0 < x < \pi.$$

DEFINITION 1.2. The initial data (w^0, w^1) is exactly $(H^1)'$ -controllable in (ξ, η) at time T if there exists u in $(H^1(0,T))'$ such that the solution w of (1.1) satisfies the condition

$$w(x,T) = w_t(x,T) = 0, \quad 0 < x < \pi.$$

In Definition 1.2, $(H^1(0,T))'$ is the dual space of $H^1(0,T)$ with respect to the space $L^2(0,T)$. This definition is inspired by [17] which studied the controllability of changing the norm for wave equation and plate equation. Note that the system (1.1) is a time-reversible linear system, so the exact controllability is equivalent to null controllability (see [6]).

The paper is organized as follows. In the remaining part of this section, we present the main results of the paper, i.e. controllability, non-controllability and minimal time for the exact controllability. In Section 2 we give some preliminaries on the theory of diophantine approximation and Riesz basis property of exponential family. The well-posedness results for the control problem (1.1) are showed in Section 3. The main results are proved in Section 4, respectively. Appendix A provides the proof of a technical lemma which is used in the proof of non-controllability in short control time.

1.2. Controllability results. To state the exact controllability results, let us 76 introduce for any ω in \mathbb{R} the functional space Y_{ω} as follows. Let $Y_0 = L^2(0,\pi)$. For 77 $\omega > 0$, let Y_{ω} be the closure in $H^{\omega}(0,\pi)$ of the y in $C^{\infty}([0,\pi])$ satisfying the conditions

79 (1.2)
$$y^{(2n)}(0) = y^{(2n)}(\pi) = 0 \quad \forall n \ge 0.$$

For $\omega < 0$, let Y_{ω} be the dual space of $Y_{-\omega}$ with respect to the space Y_0 . 80

Our exact controllability results are the following.

- THEOREM 1.3. 1. There exists a set A contained in (0,1) (uncountable but with Lebesgue measure zero, defined in Section 2) such that for every ξ and η with $\frac{\eta+\xi}{2\pi}$ and $\frac{\eta-\xi}{2\pi}$ belonging to the set A and for any $T>2\pi\sqrt{\alpha}$, all initial data in $Y_4\times Y_3$ are exactly L^2 -controllable in (ξ,η) at time T.
- 2. Let $\varepsilon > 0$, there exists a set B_{ε} contained in (0,1) (with Lebesgue measure equal to 1, defined in Section 2) such that for every ξ and η with $\frac{\eta + \xi}{2\pi}$ and $\frac{\eta \xi}{2\pi}$ belonging to the set B_{ε} and for any $T > 2\pi\sqrt{\alpha}$, all initial data in $Y_{4+\varepsilon} \times Y_{3+\varepsilon}$ are exactly L^2 -controllable in (ξ, η) at time T.

The method for proving Theorem 1.3 is inspired by the ideas and methods used in [21] for Euler-Bernoulli beam with piezoelectric actuator. Theorem 1.3 gives us two exact L^2 -controllability results. The first result of Theorem 1.3 shows that, for the end of the piezoelectric actuator in an uncountable zero measure set, we have the exact L^2 -controllability in space $Y_4 \times Y_3$. The second result of Theorem 1.3 shows that, for almost all choices of the end of the piezoelectric actuator, we have the exact L^2 -controllability in Sobolev spaces more regular than $Y_4 \times Y_3$.

In order to obtain the exact controllability in less regular space, we propose the following Theorem.

- Theorem 1.4. 1. Let the set A is same as the set in Theorem 1.3. Then for every ξ and η with $\frac{\eta+\xi}{2\pi}$ and $\frac{\eta-\xi}{2\pi}$ belonging to the set A and for any $T>2\pi\sqrt{\alpha}$, all initial data in $Y_3\times Y_2$ are exactly $(H^1)'$ -controllable in (ξ,η) at time T.
- 2. Let $\varepsilon > 0$ and the set B_{ε} is same as the set in Theorem 1.3. Then for every ξ and η with $\frac{\eta + \xi}{2\pi}$ and $\frac{\eta \xi}{2\pi}$ belonging to the set B_{ε} and for any $T > 2\pi\sqrt{\alpha}$, all initial data in $Y_{3+\varepsilon} \times Y_{2+\varepsilon}$ are exactly $(H^1)'$ -controllable in (ξ, η) at time T.

As we have already mentioned, the definition of exact $(H^1)'$ -controllability is inspired by [17]. Similar to Theorem 1.3, the method for proving Theorem 1.4 is inspired by [21]. The differences between Theorem 1.4 and Theorem 1.3 are the space of the control function and the space of the controllable initial data. In Theorem 1.4, the control function belongs to $(H^1(0,T))'$ rather than $L^2(0,T)$ and the space of the controllable initial data is larger than the space in Theorem 1.3 with the same choice of ξ and η . Roughly speaking, the larger (less regular) the space of control function is, the larger (less regular) the space of controllable initial data is. To the best knowledge of the authors, such a result has not been developed yet for beam equation with piezoelectric actuator or interior control.

Theorems 1.3 and 1.4 give some sufficient conditions for exact controllability. All the results show the dependence of the space of exactly controllable initial data on the location of the actuator. The proofs of these two Theorems are quite similar. First we use the HUM to claim the controllability is equivalent to the observability of its adjoint problem (see Propositions 4.1 and 4.4). Then we use the preliminaries introduced in Section 2 to prove the observability inequalities. The main difference between the proofs of Theorem 1.3 and Theorem 1.4 is that the measurements in Theorem 1.3 and Theorem 1.4 are different (see Propositions 4.1 and 4.4).

1.3. Non-controllability results. After the controllability results, we show some non-controllability results. In Section 4, from Propositions 4.1 and 4.4 and the solution (3.5) of the adjoint problem, we can see that condition

126 (1.3)
$$\frac{\eta - \xi}{2\pi}, \, \frac{\eta + \xi}{2\pi} \in \mathbb{R} \setminus \mathbb{Q}$$

is necessary to have any exact controllability result. Inspired by the study of [21] for Euler-Bernoulli beam with piezoelectric actuator, we prove the following non-

controllability results. They show that the condition (1.3) is not sufficient in the sense that there are ξ and η satisfying (1.3) that do not allow the control of arbitrary regular initial data in any time T.

- THEOREM 1.5. 1. For any $\beta \geq -1$, there exist ξ and η satisfying (1.3) such that for any T > 0, the space $Y_{\beta+3} \times Y_{\beta+2}$ contains initial data that are not exactly L^2 -controllable in (ξ, η) at time T.
- 2. For any $\beta \geq -2$, there exist ξ and η satisfying (1.3) such that for any T > 0, the space $Y_{\beta+3} \times Y_{\beta+2}$ contains initial data that are not exactly $(H^1)'$ -controllable in (ξ, η) at time T.

Remark 1.6. In Theorem 1.5, the initial data space is different between exact L^2 -controllability and exact $(H^1)'$ -controllability because the well-posedness result is in space $Y_2 \times Y_1$ as u belongs to $L^2(0,T)$ while the well-posedness result is in space $Y_1 \times Y_0$ as u belongs to $(H^1(0,T))'$ (see subsections 3.1 and 3.2).

Theorem 1.3 gives no information on the exact L^2 -controllability of initial data in $Y_{\beta+3} \times Y_{\beta+2}$ for $\beta < 1$ and Theorem 1.4 gives no information on the exact $(H^1)'$ -controllability of initial data in $Y_{\beta+3} \times Y_{\beta+2}$ for $\beta < 0$. A partial answer is given by the following results.

THEOREM 1.7. Suppose that $\varepsilon > 0$, T > 0 and ξ , η in $(0, \pi)$ are arbitrary.

- 1. The set $Y_{3-\varepsilon} \times Y_{2-\varepsilon}$ contains some initial data that are not exactly L^2 -controllable in (ξ, η) at time T.
- 2. The set $Y_{2-\varepsilon} \times Y_{1-\varepsilon}$ contains some initial data that are not exactly $(H^1)'$ controllable in (ξ, η) at time T.

Notice that Theorems 1.3 and 1.4 require $T>2\pi\sqrt{\alpha}$, however, in [21] the exact controllability results for Euler-Bernoulli beam have no requirement for control time. Consequently, a huge difference between Rayleigh beam and Euler-Bernoulli beam is revealed and the reason lies in various distributions of their eigenvalues. More precisely, under same boundary condition (1.1b), the eigenvalues of Rayleigh beam equation are $\frac{k^4}{1+\alpha k^2}$ for k in \mathbb{N}^* (see subsection 3.1) while the eigenvalues of Euler-Bernoulli beam equation are k^4 for k in \mathbb{N}^* (see [21]). Roughly speaking, this fact makes that Rayleigh beam equation possesses finite propagation speed and that Euler-Bernoulli beam equation possesses infinite propagation speed. For this reason, the exact controllability results of Rayleigh beam all require $T>2\pi\sqrt{\alpha}$ while the exact controllability results of Euler-Bernoulli beam hold for all T>0 (see [21]). Based on this fact, we give the non-controllability results for $0 < T < 2\pi\sqrt{\alpha}$. Inspired by [2], we propose and prove Theorem 1.8. In [2], the simultaneous controllability for two elastic strings was studied and the lack of approximate controllability was obtained when the control time is smaller than the critical time.

THEOREM 1.8. Assume $0 < T < 2\pi\sqrt{\alpha}$ and ξ , η in $(0,\pi)$ are arbitrary.

- 1. For any $\beta \geq -1$, the space $Y_{\beta+3} \times Y_{\beta+2}$ contains initial data that are not exactly L^2 -controllable in (ξ, η) at time T.
- 2. For any $\beta \geq -2$, the space $Y_{\beta+3} \times Y_{\beta+2}$ contains initial data that are not exactly $(H^1)'$ -controllable in (ξ, η) at time T.

171 Remark 1.9. For the question $T=2\pi\sqrt{\alpha}$, whether the exact controllability still holds remains open.

Notice that in Theorem 1.5, the lack of controllability holds for some special ξ and η which are related to the space of initial data. However, in Theorems 1.7 and 1.8, non-controllability holds for any ξ and η . From Theorem 1.8, we can see that

 $T \geq 2\pi\sqrt{\alpha}$ is necessary for exact controllability for Rayleigh beam equation and that minimal time for the exact controllability is obtained. As far as we know, this is the first result stating a lack of controllability for Rayleigh beam in short control time.

So far, the exact controllability for Rayleigh beam equation with piezoelectric actuator has been fully considered. As said in Remark 1.9, exact controllability in critical time is an open problem. Controllability for other type of beam equation with piezoelectric actuator, such as shear beam equation, also remains open.

2. Preliminaries. In this section, we provide some known results on the theory of diophantine approximation (see [5, 14]) and Riesz basis property of exponential family (see [3]).

For a real number ρ , we denote by $\|\rho\|_{\mathbb{Z}}$ the difference, taken positively, between ρ and the nearest integer, i.e.,

$$\|\rho\|_{\mathbb{Z}} = \min_{n \in \mathbb{Z}} |\rho - n|.$$

- Let us denote by A the set of all irrationals ρ in (0,1) such that if $[0,a_1,\ldots,a_n\ldots]$ is the expansion of ρ as a continued fraction, then (a_n) is bounded. Its Lebesgue measure is equal to zero (see [5]). The following property of this set is essentially
- useful in our work (see [14]).

179

180

181

182

183

184 185

186

187

208

PROPOSITION 2.1. A number ρ is in A if and only if there exists a constant C > 0 such that

$$||q\rho||_{\mathbb{Z}} \ge \frac{C}{q}$$

- 196 for all strictly positive integer q.
- The next proposition, which is proved in [5], shows that an inequality slightly weaker than (2.1) holds for almost all points in (0,1). This proposition is the definition of set B_{ε} .
- PROPOSITION 2.2. For any $\varepsilon > 0$ there exists a set $B_{\varepsilon} \subseteq (0,1)$ having Lebesgue measure equal to 1 and a constant C > 0, such that for any ρ in B_{ε} ,

$$||q\rho||_{\mathbb{Z}} \ge \frac{C}{q^{1+\varepsilon}}$$

- 203 for all strictly positive integer q.
- The following proposition on simultaneous approximation (see [5]) used in [21] is quite important to prove Theorem 1.7.
- PROPOSITION 2.3. Let ρ_1, \ldots, ρ_k be k irrationals in (0,1). Then there exists a strictly increasing sequence of natural numbers q_n such that

$$q_n^{\frac{1}{k}} \max_{i=1,\dots,k} (\|q_n \rho_1\|_{\mathbb{Z}}, \dots, \|q_n \rho_i\|_{\mathbb{Z}}, \dots, \|q_n \rho_k\|_{\mathbb{Z}}) \le \frac{k}{k+1} \quad \forall n \ge 1.$$

- The next proposition (see Theorem II.4.18 of [3]) on Riesz basis property of exponential family in $L^2(0,T)$ is essential for us to prove Theorem 1.8.
- PROPOSITION 2.4. Let $\{\lambda_n\}_{n\in\mathbb{Z}}$ be a sequence of complex numbers such that

$$\sup_{n\in\mathbb{Z}}|\mathrm{Im}\lambda_n|<\infty,\quad \inf_{n\neq m}|\lambda_m-\lambda_n|>0.$$

213 Set

223

224

225

$$N(x,r) := \sharp \{\lambda_n | x \le \operatorname{Re}\lambda_n < x + r\}, \quad x \in \mathbb{R}, r > 0,$$

where $\sharp A$ is the number of elements in the set A. Assume that for some T > 0,

$$\frac{N(x,r)}{r} \to \frac{T}{2\pi}, \quad as \ r \to \infty$$

- uniformly relative to $x \in \mathbb{R}$. Then for any T' in (0,T), family $\{e^{i\lambda_n t}\}_{n\in\mathbb{Z}}$ contains a subfamily $G_{T'}$ that forms a Riesz basis in $L^2(0,T')$.
- 3. Well-posedness of (1.1). In subsection 3.1, we show the well-posedness result of system (1.1) with $L^2(0,T)$ control function which has been proved in [23]. In subsection 3.2, we prove the well-posedness and regularity results of system (1.1) with $(H^1(0,T))'$ control function.
 - 3.1. Well-posedness of (1.1) with $L^2(0,T)$ control function. We state the well-posedness result and show the proof here, because the process of the proof is also used in other sections.
- THEOREM 3.1. Suppose that (w^0, w^1) belongs to $Y_2 \times Y_1$. For any u in $L^2(0,T)$ and for any ξ and η in $(0,\pi)$, the initial and boundary value problem (1.1) admits a unique solution having the regularity

229 (3.1)
$$w \in C([0,T]; Y_2) \cap C^1([0,T]; Y_1).$$

In order to prove Theorem 3.1, let us first consider the adjoint problem of (1.1)

231 (3.2a)
$$\phi_{tt}(x,t) - \alpha \phi_{xxtt}(x,t) + \phi_{xxxx}(x,t) = 0, \quad 0 < x < \pi, t > 0,$$

232 (3.2b)
$$\phi(0,t) = \phi(\pi,t) = \phi_{xx}(0,t) = \phi_{xx}(\pi,t) = 0, \quad t > 0,$$

$$\phi(x,0) = \phi^0(x), \ \phi_t(x,0) = \phi^1(x), \quad 0 < x < \pi.$$

- The following lemma proved in [23] shows the well-posedness of the adjoint problem (3.2) and some trace regularities needed in the proof of Theorem 3.1.
- LEMMA 3.2. For any initial data (ϕ^0, ϕ^1) in $Y_2 \times Y_1$, there exists a unique weak solution ϕ of (3.2) in the class $C([0,T];Y_2) \cap C^1([0,T];Y_1)$. Moreover, for all b in $(0,\pi)$ we have $\phi_x(b,\cdot)$ belongs to $H^1(0,T)$ and there exist C, C' > 0 such that

240 (3.3)
$$\|\phi_x(b,\cdot)\|_{H^1(0,T)}^2 \le C(\|\phi^0\|_{H^2(0,\pi)}^2 + \|\phi^1\|_{H^1(0,\pi)}^2),$$

$$\|\phi_x(b,\cdot)\|_{L^2(0,T)}^2 \le C'(\|\phi^0\|_{H^1(0,\pi)}^2 + \|\phi^1\|_{L^2(0,\pi)}^2).$$

- *Proof.* It is easy to see, by the semigroup method, that the problem (3.2) is well-posed in the space $Y_2 \times Y_1$ (see [20]).
- Next we prove (3.3) and (3.4). Since $\{x \mapsto \sin(kx)\}_{k \in \mathbb{N}^*}$ are eigenfunctions of (3.2a) and (3.2b), we put

$$\phi^{0}(x) = \sum_{k \ge 1} a_k \sin(kx), \quad \phi^{1}(x) = \sum_{k \ge 1} b_k \sin(kx),$$

248 with (k^2a_k) and (kb_k) in $l^2(\mathbb{R})$. Obviously, we have

249 (3.5)
$$\phi(x,t) = \sum_{k>1} \left[a_k \cos\left(\frac{k^2}{\sqrt{1+\alpha k^2}}t\right) + \frac{b_k \sqrt{1+\alpha k^2}}{k^2} \sin\left(\frac{k^2}{\sqrt{1+\alpha k^2}}t\right) \right] \sin(kx).$$

Then for all T > 0, $\phi_x(b,\cdot)$ belongs to $H^1(0,T)$ and

$$\int_0^T |\phi_{xt}(b,t)|^2 \mathrm{d}x \le C \sum_{k \ge 1} k^2 \left(a_k^2 \frac{k^4}{1 + \alpha k^2} + b_k^2 \right) \le C \sum_{k \ge 1} k^2 (a_k^2 k^2 + b_k^2),$$

252 which yields (3.3). And simultaneously we have

$$\int_0^T |\phi_x(b,t)|^2 \mathrm{d}x \le C' \sum_{k \ge 1} k^2 \left(a_k^2 + b_k^2 \frac{1 + \alpha k^2}{k^4} \right) \le C' \sum_{k \ge 1} (a_k^2 k^2 + b_k^2),$$

- 254 which clearly yields (3.4).
- 255 Proof of Theorem 3.1. Thanks to Lemma 3.2, the following backward adjoint 256 problem is well-posed in $Y_2 \times Y_1$ for every $\tau > 0$ and g in Y_1 .

257 (3.6a)
$$v_{tt}(x,t) - \alpha v_{xxtt}(x,t) + v_{xxxx}(x,t) = 0, \quad 0 < x < \pi, t \in (0,\tau),$$

258 (3.6b)
$$v(0,t) = v(\pi,t) = v_{xx}(0,t) = v_{xx}(\pi,t) = 0, \quad t \in (0,\tau),$$

369 (3.6c)
$$v(x,\tau) = 0, v_t(x,\tau) = g(x), \quad 0 < x < \pi.$$

Moreover, for any b in $(0, \pi)$ we have

$$||v_x(b,\cdot)||_{L^2(0,\tau)} \le C||g||_{Y_0}.$$

Since (1.1a) is linear, it is enough to consider the case $w^0 = w^1 = 0$. Suppose again g belongs to $C_0^{\infty}(0,\pi)$, and let v be the solution of (3.6). Define a linear operator $\mathcal{L} := I - \alpha \partial_{xx}$. It is well-known that operator \mathcal{L} is an isomorphism from Y_2 to Y_0 and an isomorphism from Y_1 to Y_{-1} by Lax-Milgram Theorem. If we multiply (1.1a) by v and integrate by parts we obtain

268 (3.8)
$$\int_{0}^{\pi} \mathcal{L}w(x,\tau)g(x)dx = \int_{0}^{\tau} u(t)(v_{x}(\eta,t) - v_{x}(\xi,t))dt.$$

269 Trace regularity (3.7) implies that

$$\left| \int_0^\tau u(t)(v_x(\eta, t) - v_x(\xi, t)) dt \right| \le C ||u||_{L^2(0, T)} ||g||_{Y_0},$$

so by (3.8), we obtain $\mathcal{L}w(\cdot,\tau)$ belongs to Y_0 , and hence $w(\cdot,\tau)$ belongs to Y_2 , for all τ in [0,T]. By replacing τ by $\tau+h$ in (3.8) we easily get that

273 (3.9)
$$w \in C([0,T]; Y_2).$$

Denote $\mathcal{R} := (I - \alpha \partial_{xx})^{-1}$. It follows from Lax-Milgram Theorem that operator \mathcal{R} is

275 an isomorphism from Y_{-2} to Y_0 and an isomorphism from Y_{-1} to Y_1 . Applying $\mathcal R$ to

both sides of (1.1a) yields

277 (3.10)
$$w_{tt}(x,t) + \mathcal{R}w_{xxxx}(x,t) = u(t)\mathcal{R}\frac{\mathrm{d}}{\mathrm{d}x}[\delta_{\eta}(x) - \delta_{\xi}(x)].$$

278 Regularity (3.9) implies that

279 (3.11)
$$\mathcal{R}w_{xxxx} \in C([0,T]; Y_0).$$

- 280 As w satisfies (3.10) and $\frac{\mathrm{d}\delta_b}{\mathrm{d}x}$ belongs to Y_{-2} for all b in $(0,\pi)$, we obtain from (3.11)
- 281 that
- 282 (3.12) $w_{tt} \in L^2(0,T;Y_0).$
- From (3.9) and (3.12), by applying the intermediate derivative theorem (see [18]) it
- 284 follows that
- $285 (3.13) w_t \in L^2(0,T;Y_1).$
- The conclusion (3.1) is now a consequence of (3.9) and (3.13) and of the general lifting result from [16].
- 3.2. Well-posedness of (1.1) with $(H^1(0,T))'$ control function. As control function u belongs to $(H^1(0,T))'$, we need to define the solution of (1.1) in the weak form. First notice that u is not a distribution. The definition of a function u in $(H^1(0,T))'$ is given by following.
- Definition 3.3. A function u is said to belong to $(H^1(0,T))'$ if there exist some functions u_0 and u_1 in $L^2(0,T)$ such that for all ϕ in $H^1(0,T)$, we have

294 (3.14)
$$\langle u, \phi \rangle = \int_0^T (u_0 \phi + u_1 \phi_t) dt.$$

295 We denote

301

302

303

304

305

306

307

308

309

$$296 (3.15) u = u_0 - \frac{\partial u_1}{\partial t}.$$

- Note that in (3.15), $\frac{\partial u_1}{\partial t}$ is not taken in the sense of distribution, but taken in the sense between $H^1(0,T)$ and its dual space.
- Inspired by [17], we define the weak solution of (1.1) by transposition and prove the well-posedness. We explain the results in three steps.
 - 1. We prove the well-posedness and trace regularity of a non-homogeneous problem (3.16).
 - 2. Using the solution of non-homogeneous problem (3.16), we define the weak solution of (1.1) by transposition.
 - 3. We prove the well-posedness of (1.1) with control function u in $(H^1(0,T))'$.
 - Step 1. In order to define the weak solution of (1.1) by transposition, we need the following well-posedness and trace regularity of a non-homogeneous problem. Assume f belongs to $L^1(0, T; Y_{-1})$ and (θ^0, θ^1) belongs to $Y_2 \times Y_1$, let us consider the following backward non-homogeneous problem
- 310 (3.16a) $\theta_{tt}(x,t) \alpha \theta_{xxtt}(x,t) + \theta_{xxxx}(x,t) = f(x,t), \quad 0 < x < \pi, t \in (0,T),$
- 311 (3.16b) $\theta(0,t) = \theta(\pi,t) = \theta_{xx}(0,t) = \theta_{xx}(\pi,t) = 0, \quad t \in (0,T),$
- $\theta(x,T) = \theta^0, \ \theta_t(x,T) = \theta^1, \quad 0 < x < \pi.$
- The following proposition proves the well-posedness and trace regularity of (3.16)
- which is needed in the definition of weak solution of (1.1). Our approach to prove
- this proposition is inspired by ideas and methods used in [1, 23].
- PROPOSITION 3.4. For any initial data (θ^0, θ^1) in $Y_2 \times Y_1$ and f in $L^1(0, T; Y_{-1})$,
- there exists a unique weak solution θ of (3.16) in the class $C([0,T];Y_2) \cap C^1([0,T];Y_1)$.
- Moreover, for all b in $(0,\pi)$ we have $\theta_x(b,\cdot)$ belongs to $H^1(0,T)$ and there exists C>0 such that
- 321 (3.17) $\|\theta_x(b,\cdot)\|_{H^1(0,T)} \le C(\|\theta^0\|_{H^2(0,\pi)} + \|\theta^1\|_{H^1(0,\pi)} + \|f\|_{L^1(0,T;Y_{-1})}).$

322 *Proof.* Applying \mathcal{R} to both sides of (3.16a) to obtain

$$\theta_{tt}(x,t) + \mathcal{R}\theta_{xxxx}(x,t) = \mathcal{R}f(x,t).$$

- Notice that $\mathcal{R}f$ belongs to $L^1(0,T;Y_1)$, then the problem (3.16) admits a unique
- solution θ in $C([0,T];Y_2) \cap C^1([0,T];Y_1)$ by the classical semigroup method (see [20]).
- 326 And there exists a constant $C_T > 0$ such that

327 (3.18)
$$\|\theta\|_{C([0,T];Y_2)} \le C_T(\|\theta^0\|_{H^2(0,\pi)} + \|\theta^1\|_{H^1(0,\pi)} + \|f\|_{L^1(0,T;Y_{-1})}).$$

- Then we need to prove trace regularity (3.17). The following lemma proved in [1] and [23] shows that operator $\mathcal{R}\partial_{xxxx}$ is "similar" to a elliptic operator $-\frac{1}{\alpha}\partial_{xx}$.
- Lemma 3.5. The linear operator $L = -\frac{1}{\alpha}\partial_{xx} \mathcal{R}\partial_{xxxx}$ is bounded from Y_2 to Y_2 .
- Using this lemma, we can reduce the proof of (3.17) to a regularity property for a string equation. We consider the initial value problem

333
$$\theta_{1,tt}(x,t) - \frac{1}{\alpha}\theta_{1,xx}(x,t) = \mathcal{R}f(x,t), \quad 0 < x < \pi, \ t \in (0,T),$$
334
$$\theta_{1}(0,t) = \theta_{1}(\pi,t) = 0, \quad t \in (0,T),$$
335
$$\theta_{1}(x,T) = \theta^{0}, \ \theta_{1,t}(x,T) = \theta^{1}, \quad 0 < x < \pi.$$

337 The relations above imply that $\theta_2 = \theta - \theta_1$ satisfies

338
$$\theta_{2,tt}(x,t) - \frac{1}{\alpha}\theta_{2,xx}(x,t) = L\theta, \quad 0 < x < \pi, \ t \in (0,T),$$
339
$$\theta_{2}(0,t) = \theta_{2}(\pi,t) = 0, \quad t \in (0,T),$$
340
$$\theta_{2}(x,T) = 0, \ \theta_{2,t}(x,T) = 0, \quad 0 < x < \pi.$$

- 342 Since θ belongs to $C([0,T];Y_2)$ and L is bounded from Y_2 to Y_2 , $L\theta$ belongs to
- $C([0,T];Y_2)$. Then by the classical theory for evolution equations of hyperbolic type
- 344 (see [15]), we obtain that θ_2 belongs to $C([0,T]; H^3(0,\pi)) \cap C^1([0,T]; H^2(0,\pi))$ and
- there exists a constant $C_T > 0$ such that

$$\|(\theta_2, \theta_{2,t})\|_{C([0,T];H^3(0,\pi)\times H^2(0,\pi))} \le C_T \|\theta\|_{C([0,T];Y_2)}.$$

- This inequality, combined with (3.18) and the standard trace theorem, implies that
- 348 for any b in $(0, \pi)$,

349 (3.21)
$$\|\theta_{2,x}(b,\cdot)\|_{H^1(0,T)} \le C_T(\|\theta^0\|_{H^2(0,\pi)} + \|\theta^1\|_{H^1(0,\pi)} + \|f\|_{L^1(0,T;Y_{-1})}).$$

- As for θ_1 , it is already proved by using multiplier methods in [10] that there exists a
- 351 constant $C_T > 0$ such that

352 (3.22)
$$\|\theta_{1,x}(b,\cdot)\|_{H^{1}(0,T)} \leq C_{T}(\|\theta^{0}\|_{H^{2}(0,\pi)} + \|\theta^{1}\|_{H^{1}(0,\pi)} + \|\mathcal{R}f\|_{L^{1}(0,T;Y_{1})})$$
$$\leq C_{T}(\|\theta^{0}\|_{H^{2}(0,\pi)} + \|\theta^{1}\|_{H^{1}(0,\pi)} + \|f\|_{L^{1}(0,T;Y_{-1})}).$$

- Then trace regularity (3.17) follows from (3.21) and (3.22).
- Step 2. Now we give the definition of the weak solution of (1.1). Denote by \mathcal{X} a Hilbert space consisting of θ , the solution of (3.16). And give \mathcal{X} a natural Hilbert

structure such that $\{f, \theta^0, \theta^1\} \to \theta$ is an isomorphism from $L^1(0, T; Y_{-1}) \times Y_2 \times Y_1$ to \mathcal{X} . From Proposition 3.4 we have

358 (3.23)
$$\mathcal{X} \subseteq C([0,T]; Y_2) \cap C^1([0,T]; Y_1),$$

$$\theta \in \mathcal{X} \Rightarrow \theta_x(b, \cdot) \in H^1(0,T) \quad \forall b \in (0,\pi).$$

Assume that $\{f, \theta^0, \theta^1\}$ belongs to $L^1(0, T; Y_{-1}) \times Y_2 \times Y_1$ and that $\{u, w^0, w^1\}$ belongs to $L^2(0, T) \times Y_2 \times Y_1$. Denote by θ the solution of (3.16) and by w the solution of (1.1) given by Theorem 3.1. Multiplying (1.1a) by θ , we obtain

$$\int_0^T \int_0^{\pi} (w_{tt} - \alpha w_{xxtt} + w_{xxxx}) \theta dx dt = \int_0^T \int_0^{\pi} u(t) \frac{d}{dx} [\delta_{\eta}(x) - \delta_{\xi}(x)] \theta dx dt.$$

363 Integrating by parts we get

$$\int_{0}^{T} \int_{0}^{\pi} w f dx dt + \int_{0}^{\pi} [\mathcal{L}w_{t}(T)\theta^{0} - \mathcal{L}w(T)\theta^{1}] dx$$

$$= -\int_{0}^{T} u(t)(\theta_{x}(\eta, t) - \theta_{x}(\xi, t)) dt + \int_{0}^{\pi} [\mathcal{L}w^{1}\theta(0) - \mathcal{L}w^{0}\theta_{t}(0)] dx.$$

Now relaxing the assumption of u belonging to $L^2(0,T)$ to u belonging to $(H^1(0,T))'$, we obtain from (3.15) that

$$\int_{0}^{T} \int_{0}^{\pi} w f dx dt + \int_{0}^{\pi} [\mathcal{L}w_{t}(T)\theta^{0} - \mathcal{L}w(T)\theta^{1}] dx$$

$$= -\int_{0}^{T} [u_{0}(t)(\theta_{x}(\eta, t) - \theta_{x}(\xi, t)) + u_{1}(t)(\theta_{xt}(\eta, t) - \theta_{xt}(\xi, t))] dt$$

$$+ \int_{0}^{\pi} [\mathcal{L}w^{1}\theta(0) - \mathcal{L}w^{0}\theta_{t}(0)] dx.$$

Now we set (3.24) as the definition of weak solution.

Definition 3.6. Let T>0, u in $(H^1(0,T))'$ and (w^0,w^1) in $Y_1\times Y_0$ be given.

A solution of the (1.1) is a function w in $C([0,T];Y_1)$ such that, for every $\{f,\theta^0,\theta^1\}$ in $L^1(0,T;Y_{-1})\times Y_2\times Y_1$, (3.24) holds and $(w(T),w_t(T))$ belongs to $Y_1\times Y_0$.

372 Step 3. Then we are able to prove the well-posedness of (1.1) when u belongs to 373 $(H^1(0,T))'$.

THEOREM 3.7. Suppose (w^0, w^1) belongs to $Y_1 \times Y_0$. For any u in $(H^1(0,T))'$ and for any ξ and η in $(0,\pi)$, the initial and boundary value problem (1.1) admits a unique weak solution in sense of Definition 3.6. And the map $\{w^0, w^1, u\} \mapsto \{w, w(T), w_t(T)\}$ is linear and continuous with respect to the corresponding norm.

Proof. Since u belongs to $(H^1(0,T))'$, there exist u_0 , u_1 in $L^2(0,T)$ such that $u=u_0-\frac{\partial u_1}{\partial t}$, where $\frac{\partial}{\partial t}$ is in sense of $(H^1(0,T))'$ derivative. Moreover, $\{f,\theta^0,\theta^1\}\to\theta$ is an isomorphism from $L^1(0,T;Y_{-1})\times Y_2\times Y_1$ to \mathcal{X} . Therefore, we define a linear form Γ on \mathcal{X} such that

$$\Gamma(\theta) = -\int_{0}^{T} [u_{0}(t)(\theta_{x}(\eta, t) - \theta_{x}(\xi, t)) + u_{1}(t)(\theta_{xt}(\eta, t) - \theta_{xt}(\xi, t))]dt + \int_{0}^{\pi} [\mathcal{L}w^{1}\theta(0) - \mathcal{L}w^{0}\theta_{t}(0)]dx.$$

- Since (w^0, w^1) belongs to $Y_1 \times Y_0$, then $(\mathcal{L}w^0, \mathcal{L}w^1)$ belongs to $Y_{-1} \times Y_{-2}$. From
- Proposition 3.4 we have Γ is a continuous linear form on \mathcal{X} . Denote by (f_1, f_2) the
- linear form between Y_{γ} and $Y_{-\gamma}$ for any $\gamma \geq 0$ and denote by $\langle g_1, g_2 \rangle$ the linear form
- between $L^{\infty}(0,T;Y_1)$ and $L^1(0,T;Y_{-1})$. Therefore, for the linear form Γ in \mathcal{X}' , there
- exists unique $\{w, \zeta_*, \zeta\}$ in $L^{\infty}(0, T; Y_1) \times Y_{-2} \times Y_{-1}$ such that

388 (3.26)
$$\langle w, f \rangle + (\zeta_*, \theta^0) + (-\zeta, \theta) = \Gamma(\theta) \quad \forall \theta \in \mathcal{X}.$$

- Next we claim that w above is actually the weak solution of (1.1). It is sufficient to
- 390 prove that w satisfies (1.1) in weak sense, $\mathcal{L}w(T) = \zeta$ and $\mathcal{L}w_t(T) = \zeta_*$.
- Notice that $\{x \mapsto \sin(kx)\}_{k \in \mathbb{N}^*}$ are eigenfunctions of $\mathcal{R}\partial_{xxxx}$. Let $m(x) = \sin(kx)$
- 392 for some k in N*. Firstly we set $\mathcal{R}f(t)=h(t)m, \ \theta^0=0$ and $\theta^1=0$, then f(t)=0
- 393 $h(t)\mathcal{L}m$. Denote by $\lambda = k^4/(1+\alpha k^2)$ the corresponding eigenvalue. Then we obtain
- 394 from Proposition 3.4 for t in (0,T)

$$\theta(t) = q(t)m,$$

- 396 where q satisfies
- 397 $q_{tt} + \lambda q = h, \quad q(T) = q_t(T) = 0,$
- 398 namely

$$q(t) = -\frac{1}{\sqrt{\lambda}} \int_{t}^{T} \sin((t - \sigma)\sqrt{\lambda}) h(\sigma) d\sigma.$$

400 Then

$$\Gamma(\theta) = -(m_x(\eta) - m_x(\xi)) \int_0^T [u_0(t)q(t) + u_1(t)q_t(t)] dt + q(0)(\mathcal{L}w^1, m) - q_t(0)(\mathcal{L}w^0, m),$$

402 and we have (3.27)

$$\langle w, f \rangle = \int_0^T (w, \mathcal{L}m)(q_{tt} + \lambda q) dt = -(m_x(\eta) - m_x(\xi)) \int_0^T [u_0(t)q(t) + u_1(t)q_t(t)] dt + q(0)(\mathcal{L}w^1, m) - q_t(0)(\mathcal{L}w^0, m).$$

Notice that $(w, \mathcal{L}m) = (\mathcal{L}w, m)$, then (3.27) implies that

$$(\mathcal{L}w, m)_{tt} + \lambda(\mathcal{L}w, m) = -(m_x(\eta) - m_x(\xi)) \left(u_0 - \frac{\partial u_1}{\partial t}\right),$$

$$(\mathcal{L}w, m)(0) = (\mathcal{L}w^0, m),$$

$$(\mathcal{L}w, m)_t(0) = (\mathcal{L}w^1, m),$$

- where $\frac{\partial u_1}{\partial t}$ is taken in the sense of $(H^1(0,T))'$. Since $m(x) = \sin(kx)$ and $k \ge 1$ is an arbitrary natural number, then w satisfies (1.1) in weak sense.
- Now if we set f = 0, $\theta^0 = 0$ and $\theta^1 = -m$, then for t in (0, T)

$$\theta(t) = \frac{1}{\sqrt{\lambda}} \sin((T - t)\sqrt{\lambda})m.$$

410 Then (3.26) implies that

$$(\zeta, m) = \Gamma(\theta) = -\left(m_x(\eta) - m_x(\xi)\right) \int_0^T \left[u_0(t) \frac{1}{\sqrt{\lambda}} \sin((T - t)\sqrt{\lambda})\right] dt$$

$$-u_1(t) \cos((T - t)\sqrt{\lambda}) dt$$

$$+ (\mathcal{L}w^1, m) \frac{1}{\sqrt{\lambda}} \sin(T\sqrt{\lambda}) + (\mathcal{L}w^0, m) \cos(T\sqrt{\lambda}).$$

Moreover, (3.28) implies that

$$(\mathcal{L}w, m)(T) = (\mathcal{L}w^{0}, m)\cos(T\sqrt{\lambda}) + (\mathcal{L}w^{1}, m)\frac{\sin(T\sqrt{\lambda})}{\sqrt{\lambda}}$$

$$- (m_{x}(\eta) - m_{x}(\xi))\int_{0}^{T} \frac{1}{\sqrt{\lambda}}\sin((T-\sigma)\sqrt{\lambda})\left(u_{0} - \frac{\partial u_{1}}{\partial \sigma}\right)d\sigma.$$

- Note the definition of $\frac{\partial u_1}{\partial t}$ and compare to (3.29), we obtain $(\mathcal{L}w, m)(T) = (\zeta, m)$, which proves $\mathcal{L}w(T) = \zeta$.
- Next we set f = 0, $\theta^0 = m$ and $\theta^1 = 0$, then for t in (0, T)

$$\theta(t) = \cos((T - t)\sqrt{\lambda})m.$$

418 Then (3.26) implies that

$$(\zeta_*, m) = \Gamma(\theta) = -(m_x(\eta) - m_x(\xi)) \int_0^T [u_0(t)\cos((T - t)\sqrt{\lambda}) + u_1(t)\sqrt{\lambda}\sin((T - t)\sqrt{\lambda})]dt + (\mathcal{L}w^1, m)\cos(T\sqrt{\lambda}) - (\mathcal{L}w^0, m)\sqrt{\lambda}\sin(T\sqrt{\lambda}).$$

420 Moreover, (3.28) implies that

$$(\mathcal{L}w, m)_{t}(T) = -(\mathcal{L}w^{0}, m)\sqrt{\lambda}\sin(T\sqrt{\lambda}) + (\mathcal{L}w^{1}, m)\cos(T\sqrt{\lambda})$$

$$-(m_{x}(\eta) - m_{x}(\xi))\int_{0}^{T}\cos((T-\sigma)\sqrt{\lambda})\left(u_{0} - \frac{\partial u_{1}}{\partial \sigma}\right)d\sigma.$$

- Note the definition of $\frac{\partial u_1}{\partial t}$ and compare to (3.31), we obtain $(\mathcal{L}w, m)_t(T) = (\zeta_*, m)$, which implies $\mathcal{L}w_t(T) = \zeta_*$.
 - Now we have proved there exists unique $\{w, w(T), w_t(T)\}$ in $L^{\infty}(0, T; Y_1) \times Y_1 \times Y_0$ such that (3.24) holds and the map $\{w^0, w^1, u\} \to \{w, w(T), w_t(T)\}$ is linear and continuous with respect to the corresponding norm. In fact we have property w belongs to $C([0, T]; Y_1)$. Since when the known data $\{w^0, w^1, u\}$ belongs to $Y_2 \times Y_1 \times L^2(0, T)$, we have (3.1), then according to the results we obtained above, this property is preserved in the density argument. It concludes the proof of Theorem 3.7.
- 4. Proofs of the main results. In this section, we prove the main results. For exact controllability results, Theorems 1.3 and 1.4, we use the HUM, introduced in [17], to illustrate the controllability problem is equivalent to the observability problem of the adjoint system. Then we prove the observability inequality in each case. For non-controllability results, Theorems 1.5, 1.7, and 1.8, we show the observability inequalities are false. Essentially, we obtained some sufficient conditions and some necessary conditions for the observability inequality.
- 4.1. Exact L^2 -controllability (Proof of Theorem 1.3). Choose (ϕ^0, ϕ^1) in $(C^{\infty}[0, \pi])^2$ satisfying the compatible conditions (1.2) and denote by $\phi(x, t)$ the solution of (3.2) with initial value (ϕ^0, ϕ^1) .
- 440 Consider a backward adjoint system

(4.1a)

424

425

426

427

429

441
$$\psi_{tt}(x,t) - \alpha \psi_{xxtt}(x,t) + \psi_{xxxx}(x,t) = u(t) \frac{\mathrm{d}}{\mathrm{d}x} [\delta_{\eta}(x) - \delta_{\xi}(x)], \quad 0 < x < \pi, \ t > 0,$$

442 (4.1b)
$$\psi(0,t) = \psi(\pi,t) = \psi_{xx}(0,t) = \psi_{xx}(\pi,t) = 0, \quad t > 0,$$

$$\psi(x,T) = \psi_t(x,T) = 0, \quad 0 < x < \pi,$$

where u in $L^2(0,T)$ will be chosen later. Problem (4.1) is well-posed according to Theorem 3.1. Then, multiplying (4.1a) by ϕ and integrating by parts, we get

447 (4.2)
$$\int_0^{\pi} \phi^0(x) \mathcal{L} \psi_t(x,0) - \phi^1(x) \mathcal{L} \psi(x,0) dx = \int_0^T u(t) (\phi_x(\eta,t) - \phi_x(\xi,t)) dt.$$

Let $u(t) = \phi_x(\eta, t) - \phi_x(\xi, t)$. Since (3.4), u belongs to $L^2(0, T)$. Define a linear operator Λ satisfying

450 (4.3)
$$\Lambda(\phi^0, \phi^1) = (\mathcal{L}\psi_t(\cdot, 0), -\mathcal{L}\psi(\cdot, 0)).$$

Since $(\mathcal{L}\psi_t(\cdot,0), -\mathcal{L}\psi(\cdot,0))$ belongs to $Y_{-1} \times Y_0$ by Theorem 3.1, operator Λ is well defined. In particularly

453 (4.4)
$$\langle \Lambda(\phi^0, \phi^1), (\phi^0, \phi^1) \rangle = \int_0^T |\phi_x(\eta, t) - \phi_x(\xi, t)|^2 dt.$$

454 Therefore, we can define a seminorm

455 (4.5)
$$\|(\phi^0, \phi^1)\|_F := \left(\int_0^T |\phi_x(\eta, t) - \phi_x(\xi, t)|^2 dt \right)^{\frac{1}{2}},$$

- 456 for all (ϕ^0, ϕ^1) in $(C^{\infty}[0, \pi])^2$ satisfying the compatible conditions (1.2).
- 457 A classical argument in HUM implies the following proposition.
- PROPOSITION 4.1. All initial data in $Y_{\beta+3} \times Y_{\beta+2}$ are exactly L^2 -controllable in (ξ, η) at time T if and only if there exists a constant c > 0 such that

460 (4.6)
$$\int_0^T |\phi_x(\eta, t) - \phi_x(\xi, t)|^2 dt \ge c(\|\phi^0\|_{H^{-\beta}}^2 + \|\phi^1\|_{H^{-\beta-1}}^2)$$

- 461 for all (ϕ^0, ϕ^1) in $(C^{\infty}[0, \pi])^2$ satisfying the compatible conditions (1.2).
- Equation (4.6) is called observability inequality. As in the proof of Lemma 3.2, the solution ϕ of the adjoint problem (3.2) has the form of (3.5), which implies that

$$\int_{0}^{T} |\phi_{x}(\eta, t) - \phi_{x}(\xi, t)|^{2} dt$$

$$=4 \int_{0}^{T} \left| \sum_{k \geq 1} k \sin\left(\frac{k(\eta + \xi)}{2}\right) \sin\left(\frac{k(\eta - \xi)}{2}\right) \left\{ a_{k} \cos\left(\frac{k^{2}}{\sqrt{1 + \alpha k^{2}}}t\right) + \frac{b_{k}\sqrt{1 + \alpha k^{2}}}{k^{2}} \sin\left(\frac{k^{2}}{\sqrt{1 + \alpha k^{2}}}t\right) \right\} \right|^{2} dt.$$

- Remark 4.2. From (4.7), we can see that (1.3) is necessary for controllability.
- To prove observability inequality (4.6) for some β , we use the following Ingham inequality (see [12, 4]) to our problem.
- LEMMA 4.3. Let $(\nu_k)_{k\in\mathbb{Z}}$ be a strictly increasing sequence of real numbers and let $\gamma_{\infty} > 0$ be defined by

$$\gamma_{\infty} = \underline{\lim}_{|k| \to \infty} |\nu_{k+1} - \nu_k|.$$

For any real $T > 2\pi/\gamma_{\infty}$, there exist two constants $C_1, C_2 > 0$ such that, for any sequence $(x_k)_{k\in\mathbb{Z}}$ in $l^2(\mathbb{C})$

473
$$C_1 \sum_{k \in \mathbb{Z}} |x_k|^2 \le \int_0^T \left| \sum_{k \in \mathbb{Z}} x_k e^{i\nu_k t} \right|^2 dt \le C_2 \sum_{k \in \mathbb{Z}} |x_k|^2.$$

We apply Lemma 4.3 with 474

$$\nu_k = -\nu_{-k} = \frac{k^2}{\sqrt{1 + \alpha k^2}},$$

$$2x_k = 2\overline{x_{-k}} = \left(a_k - i\frac{b_k\sqrt{1 + \alpha k^2}}{k^2}\right)k\sin\left(\frac{k(\eta + \xi)}{2}\right)\sin\left(\frac{k(\eta - \xi)}{2}\right).$$

As $\lim_{|k|\to\infty} |\nu_{k+1}-\nu_k|=1/\sqrt{\alpha}$, then for any real $T>2\pi\sqrt{\alpha}$, there exist two constants $C_1,\,C_2>0$ such that,

$$C_{1} \sum_{k \geq 1} k^{2} \left(a_{k}^{2} + \frac{b_{k}^{2}(1 + \alpha k^{2})}{k^{4}} \right) \left[\sin \left(\frac{k(\eta + \xi)}{2} \right) \sin \left(\frac{k(\eta - \xi)}{2} \right) \right]^{2}$$

$$\leq \int_{0}^{T} |\phi_{x}(\eta, t) - \phi_{x}(\xi, t)|^{2} dt$$

$$\leq C_{2} \sum_{k \geq 1} k^{2} \left(a_{k}^{2} + \frac{b_{k}^{2}(1 + \alpha k^{2})}{k^{4}} \right) \left[\sin \left(\frac{k(\eta + \xi)}{2} \right) \sin \left(\frac{k(\eta - \xi)}{2} \right) \right]^{2}.$$

When $\frac{\eta+\xi}{2\pi}$ and $\frac{\eta-\xi}{2\pi}$ belong to A, from (2.1) we see that there exists a constant C>0 such that for all $k\geq 1$ we have

$$\left| \sin \left(\frac{k(\eta \pm \xi)}{2} \right) \right| = \left| \sin \left\{ \pi \left[\frac{k(\eta \pm \xi)}{2\pi} - p \right] \right\} \right| \ge \left| \sin \left(\frac{\pi C}{k} \right) \right| \ge \frac{C}{k}.$$

Inequalities (4.8) and (4.9) imply that 482

483
$$\int_0^T |\phi_x(\eta, t) - \phi_x(\xi, t)|^2 dt \ge c \sum_{k \ge 1} \left(\frac{a_k^2}{k^2} + \frac{b_k^2 (1 + \alpha k^2)}{k^6} \right),$$

- which is exactly (4.6) when $\beta = 1$. This fact completes the proof of the first part of 484
- When $\frac{\eta+\xi}{2\pi}$ and $\frac{\eta-\xi}{2\pi}$ belong to B_{ε} , from (2.2) we see that there exists a constant C>0 such that for all $k\geq 1$ we have 486

$$\left| \sin \left(\frac{k(\eta \pm \xi)}{2} \right) \right| \ge \frac{C}{k^{1+\varepsilon}}.$$

Inequalities (4.8) and (4.10) imply that

490
$$\int_0^T |\phi_x(\eta, t) - \phi_x(\xi, t)|^2 dt \ge c \sum_{k > 1} \left(\frac{a_k^2}{k^{2+\varepsilon}} + \frac{b_k^2 (1 + \alpha k^2)}{k^{6+\varepsilon}} \right),$$

which is exactly (4.6) when $\beta = 1 + \varepsilon$. This completes the proof of the second part of 491

492 Theorem 1.3. 493 **4.2.** Exact $(H^1)'$ -controllability (Proof of Theorem 1.4). Similarly to sub-494 section 4.1, we use the HUM to study the controllability problem. Let (ϕ^0, ϕ^1) in 495 $(C^{\infty}[0,\pi])^2$ satisfy the compatible conditions (1.2). Denote by ϕ the solution of (3.2) 496 with initial value (ϕ^0, ϕ^1) .

Consider a backward system (4.1) where u in $(H^1(0,T))'$ will be chosen later. Problem (4.1) is well-posed according to Theorem 3.7. Then multiplying (4.1a) by ϕ and integrating by parts, we obtain

500 (4.11)
$$\int_0^{\pi} \phi^0(x) \mathcal{L} \psi_t(x,0) - \phi^1(x) \mathcal{L} \psi(x,0) dx = \int_0^T u(t) (\phi_x(\eta,t) - \phi_x(\xi,t)) dt.$$

Recall the definition (3.15) of u, this equation is equivalent to

$$\int_{0}^{\pi} \phi^{0}(x) \mathcal{L}\psi_{t}(x,0) - \phi^{1}(x) \mathcal{L}\psi(x,0) dx$$

$$= \int_{0}^{T} [u_{0}(t)(\phi_{x}(\eta,t) - \phi_{x}(\xi,t)) + u_{1}(t)(\phi_{xt}(\eta,t) - \phi_{xt}(\xi,t))] dt.$$

Let $u_0(t)=0$ and $u_1(t)=\phi_{xt}(\eta,t)-\phi_{xt}(\xi,t)$. Since (3.3), u_1 belongs to $L^2(0,T)$, and therefore, $u=-\frac{\partial u_1}{\partial t}$ belongs to $(H^1(0,T))'$, where the derivative $\frac{\partial}{\partial t}$ is taken in sense of $(H^1(0,T))'$. Define a linear operator Λ_* satisfying

506 (4.13)
$$\Lambda_*(\phi^0, \phi^1) = (\mathcal{L}\psi_t(\cdot, 0), -\mathcal{L}\psi(\cdot, 0)).$$

Since $(\mathcal{L}\psi_t(\cdot,0), -\mathcal{L}\psi(\cdot,0))$ belongs to $Y_{-2} \times Y_{-1}$ because of Theorem 3.7, operator Λ_* is well defined. In particularly,

509 (4.14)
$$\langle \Lambda_*(\phi^0, \phi^1), (\phi^0, \phi^1) \rangle = \int_0^T |\phi_{xt}(\eta, t) - \phi_{xt}(\xi, t)|^2 dt.$$

510 Therefore, we can define a seminorm

511 (4.15)
$$\|(\phi^0, \phi^1)\|_{F_*} := \left(\int_0^T |\phi_{xt}(\eta, t) - \phi_{xt}(\xi, t)|^2 dt \right)^{\frac{1}{2}},$$

- for all (ϕ^0, ϕ^1) in $(C^{\infty}[0, \pi])^2$ satisfying the compatible conditions (1.2).
- A classical argument in HUM shows the following proposition.
- PROPOSITION 4.4. All initial data in $Y_{\beta+3} \times Y_{\beta+2}$ are exactly $(H^1)'$ -controllable in (ξ, η) at time T if and only if there exists a constant c > 0 such that

516 (4.16)
$$\int_0^T |\phi_{xt}(\eta, t) - \phi_{xt}(\xi, t)|^2 dt \ge c \left(\|\phi^0\|_{H^{-\beta}}^2 + \|\phi^1\|_{H^{-\beta-1}}^2 \right)$$

- for all (ϕ^0, ϕ^1) in $(C^{\infty}[0, \pi])^2$ satisfying the compatible conditions (1.2).
- After a similar argument in subsection 4.1. We have the following result. Assume $T > 2\pi\sqrt{\alpha}$. When $\frac{\eta+\xi}{2\pi}$ and $\frac{\eta-\xi}{2\pi}$ belong to A, there exists a constant c > 0 such that

520 (4.17)
$$\int_0^T |\phi_{xt}(\eta, t) - \phi_{xt}(\xi, t)|^2 dt \ge c \sum_{k>1} \left(\frac{a_k^2 k^2}{1 + \alpha k^2} + \frac{b_k^2}{k^2} \right),$$

- which is exactly (4.16) when $\beta = 0$. This completes the proof of the first part of Theorem 1.4.
- When $\frac{\eta + \xi}{2\pi}$ and $\frac{\eta \xi}{2\pi}$ belong to B_{ε} , there exists a constant c > 0 such that

$$\int_0^T |\phi_{xt}(\eta, t) - \phi_{xt}(\xi, t)|^2 dt \ge c \sum_{k>1} \left(\frac{a_k^2 k^2}{(1 + \alpha k^2) k^{\varepsilon}} + \frac{b_k^2}{k^{2+\varepsilon}} \right),$$

- which is exactly (4.16) when $\beta = \varepsilon$. This completes the proof of the second part of Theorem 1.4.
- 4.3. Condition (1.3) is not sufficient (Proof of Theorem 1.5). As mentioned in Remark 4.2, condition (1.3) is necessary for controllability. Thanks to Propositions 4.1 and 4.4, to show condition (1.3) is not sufficient for L^2 -controllability, it is sufficient to show there exist ξ and η satisfying (1.3) such that (4.6) is false for any $\beta \geq -1$ and c > 0. Similarly, to show condition (1.3) is not sufficient for $(H^1)'$ -controllability, it is sufficient to show there exist ξ and η satisfying (1.3) such that (4.16) is false for any $\beta \geq -2$ and c > 0.
- For any fixed $\beta \ge -1$ and

535 (4.19)
$$\nu > \max\left(\frac{3}{2}\beta + 1, 2\right),$$

536 we choose

537 (4.20)
$$\frac{\eta + \xi}{2\pi} = \sum_{n=1}^{\infty} \frac{a_n}{10^{n!}},$$

- where a_n belongs to $\{0, 1, \dots, 9\}$ for all $n \geq 1$, and a_n is not identically zero for great
- 539 n. According to [22] the right-hand side of (4.20) is a Liouville number, i.e., it is
- transcendental and there exists a strictly increasing sequence of integers q_n such that

$$\left| \sin \left(q_n \frac{\eta + \xi}{2\pi} \right) \right| \le \frac{\pi}{q_n^{\nu}} \quad \forall n \ge 1.$$

Now consider the sequence of initial data

543 (4.22)
$$\phi_n^0(x) = q_n^\mu \sin(q_n x), \quad \phi_n^1(x) = 0 \quad \forall x \in (0, \pi),$$

- 544 where $\mu = \frac{3}{2}\beta$ if $\beta > 0$ and $\mu = 1$ if $-1 \le \beta \le 0$. Obviously (ϕ_n^0, ϕ_n^1) belongs to
- 545 $(C^{\infty}[0,\pi])^2$ and satisfies compatible conditions (1.2) and

546 (4.23)
$$\|\phi_n^0\|_{H^{-\beta}}^2 + \|\phi_n^1\|_{H^{-\beta-1}}^2 \to \infty \quad \forall \beta \ge -1.$$

547 So by (4.7), (4.19), and (4.21), we have

$$\int_{0}^{T} |\phi_{n,x}(\eta,t) - \phi_{n,x}(\xi,t)|^{2} dt$$

$$=4 \int_{0}^{T} \left| q_{n} \sin\left(q_{n} \frac{\eta + \xi}{2\pi}\right) \sin\left(q_{n} \frac{\eta - \xi}{2\pi}\right) q_{n}^{\mu} \cos\left(\frac{q_{n}^{2}}{\sqrt{1 + \alpha q_{n}^{2}}}t\right) \right|^{2} dt$$

$$\leq 4T q_{n}^{2(\mu+1)} \left| \sin\left(q_{n} \frac{\eta + \xi}{2\pi}\right) \right|^{2}$$

$$\leq 4\pi T q_{n}^{2(\mu+1-\nu)} \to 0, \quad \text{as } n \to \infty.$$

- Relations (4.23) and (4.24) show that (4.6) is false for any $\beta \geq -1$ and c > 0.
- Then we show that there exist ξ and η in $(0, \pi)$ satisfying (1.3) such that (4.16)
- is false for any $\beta \geq -2$ and c > 0. The proof is quite similar to the proof above in
- 552 this subsection. For any fixed $\beta \geq -2$, we only need to change ν as

553 (4.25)
$$\nu > \max\left(\frac{3}{2}\beta + 2, 3\right),$$

- and to set $\mu = \frac{3}{2}\beta$ if $\beta > 0$ and $\mu = 1$ if $-2 \le \beta \le 0$. By similar calculation, we obtain that (4.16) is false for any $\beta \ge -2$ and c > 0.
- 556 **4.4.** Non-controllability for less regular initial data set (Proof of The-557 **orem 1.7).** Similar to last subsection, we focus on (4.6) and (4.16). By applying 558 Proposition 2.3, for any ξ and η in $(0, \pi)$ we obtain the existence of a strictly increas-
- ing sequence of positive integers $\{q_n\}_{n\geq 1}$ such that

$$\left| \sin \left(q_n \frac{\eta + \xi}{2} \right) \right| \le \frac{\pi}{\sqrt{q_n}}, \quad \left| \sin \left(q_n \frac{\eta - \xi}{2} \right) \right| \le \frac{\pi}{\sqrt{q_n}} \quad \forall n \ge 1.$$

- 561 First consider the sequence of initial data
- 562 (4.27) $\phi_n^0(x) = \sin(q_n x), \quad \phi_n^1(x) = 0 \quad \forall x \in (0, \pi).$
- 563 We note that

564 (4.28)
$$\|\phi_n^0\|_{H^{\varepsilon}}^2 + \|\phi_n^1\|_{H^{\varepsilon+1}}^2 = Cq_n^{2\varepsilon} \to \infty, \quad n \to \infty,$$

where C is a positive constant. By (4.7) and (4.26) we have

$$\int_{0}^{T} |\phi_{n,x}(\eta,t) - \phi_{n,x}(\xi,t)|^{2} dt$$

$$=4 \int_{0}^{T} \left| q_{n} \sin\left(q_{n} \frac{\eta + \xi}{2\pi}\right) \sin\left(q_{n} \frac{\eta - \xi}{2\pi}\right) \cos\left(\frac{q_{n}^{2}}{\sqrt{1 + \alpha q_{n}^{2}}} t\right) \right|^{2} dt \leq K \quad \forall n \geq 1,$$

- where K is a positive constant. So (4.28) and (4.29) show that (4.6) is false for $\beta = -\varepsilon$ and arbitrary c > 0.
- If we choose the sequence of initial data

570 (4.30)
$$\phi_n^0(x) = q_n^{-1}\sin(q_n x), \quad \phi_n^1(x) = 0 \quad \forall x \in (0, \pi).$$

571 We note that

572 (4.31)
$$\|\phi_n^0\|_{H^{\varepsilon+1}}^2 + \|\phi_n^1\|_{H^{\varepsilon+2}}^2 = Cq_n^{2\varepsilon} \to \infty, \quad n \to \infty,$$

where C is a positive constant. By (4.26) we have

$$\int_{0}^{T} |\phi_{n,xt}(\eta,t) - \phi_{n,xt}(\xi,t)|^{2} dt$$

$$=4 \int_{0}^{T} \left| q_{n} \sin\left(q_{n} \frac{\eta + \xi}{2\pi}\right) \sin\left(q_{n} \frac{\eta - \xi}{2\pi}\right) \frac{q_{n}^{2}}{\sqrt{1 + \alpha q_{n}^{2}}} \cos\left(\frac{q_{n}}{\sqrt{1 + \alpha q_{n}^{2}}} t\right) \right|^{2} dt$$

$$\leq K \quad \forall n \geq 1,$$

where K is a positive constant. So (4.31) and (4.32) show that (4.16) is false for

 $\beta = -\varepsilon - 1$ and arbitrary c > 0.

4.5. Non-controllability in short control time (Proof of Theorem 1.8).

We study the lack of controllability when $0 < T < 2\pi\sqrt{\alpha}$ in this subsection. For any

579 $\beta \geq -1$, we need to find $\{(\phi_m^0, \phi_m^1)\}_{m \in \mathbb{N}^*}$ such that

$$\int_{0}^{T} |\phi_{m,x}(\eta,t) - \phi_{m,x}(\xi,t)|^{2} dt \to 0, \text{ as } m \to \infty$$

581 and

$$\|\phi_m^0\|_{H^{-\beta}}^2 + \|\phi_m^1\|_{H^{-\beta-1}}^2 \ge c > 0$$

583 for any $m \ge 1$ and $\beta \ge -1$.

As in subsection 4.1, denote

$$585 \quad (4.33) \qquad \qquad \nu_k = -\nu_{-k} = \frac{k^2}{\sqrt{1 + \alpha k^2}}.$$

- Obviously, $\{\nu_k\}_{k\in\mathbb{Z}^*}$ is a strictly increasing sequence and $\lim_{|k|\to\infty} |\nu_{k+1}-\nu_k| =$
- 587 $1/\sqrt{\alpha} > 0$. Define N(x,r) as in Proposition 2.4 corresponding to $\{\nu_k\}_{k \in \mathbb{Z}^*}$, we have
- 588 the following lemma.

LEMMA 4.5. Let $\{\nu_k\}_{k\in\mathbb{Z}^*}$ and N(x,r) be defined above. We have

590 (4.34)
$$\frac{N(x,r)}{r} \to \sqrt{\alpha}, \quad as \ r \to \infty$$

- 591 uniformly relative to x in \mathbb{R} .
- We will prove this lemma in Appendix A. For any $0 < T < 2\pi\sqrt{\alpha}$, we can choose a
- 593 T' such that $0 < T < T' < 2\pi\sqrt{\alpha}$. Let f in $L^2(0, 2\pi\sqrt{\alpha})$ be a real valued function
- such that f(t) = 0 if $0 \le t \le T$ and $||f||_{L^2(0,T')} \ne 0$. According to Lemma 4.5
- and Proposition 2.4, the family $\{e^{i\nu_k t}\}_{k\in\mathbb{Z}^*}$ contains a subfamily $G_{T'}$ which forms a
- 596 Riesz basis in $L^2(0,T')$.
- 597 Remark 4.6. The algorithm proposed in the proof of Proposition 2.4 (see Theorem
- II.4.18 of [3]) allows us to construct a subfamily $G_{T'}$ with symmetric spectrum, i.e.,

$$e^{i\lambda t} \in G_{T'} \Rightarrow e^{-i\lambda t} \in G_{T'}.$$

- 600 This fact is used in our proof.
- 601 According to the remark above, there exists a strictly increasing sequence of integer
- 602 $\{q_n\}_{n\in\mathbb{Z}^*}$ satisfying $q_n=-q_{-n}$, such that $G_{T'}=\{e^{i\lambda_{q_n}t}\}_{n\in\mathbb{Z}^*}$. Then for f in $L^2(0,T')$
- defined above, there exists a sequence $\{l_n\}_{n\in\mathbb{Z}^*}$ in $l^2(\mathbb{C})$ such that

$$f(t) = \sum_{n \in \mathbb{Z}^*} l_n e^{i\lambda_{q_n} t} \quad \text{in } L^2(0, T')$$

605 and

606

$$0 < \sum_{n \in \mathbb{Z}^*} |l_n|^2 < \infty.$$

607 Since f(t) is a real valued function, $l_n = \overline{l_{-n}}$. Now we can defined a sequence

608 $\{(\phi_m^0, \phi_m^1)\}_{m \in \mathbb{N}^*}$ of initial data such that (4.35)

$$\phi_m^0(x) = 2\sum_{n=1}^m \operatorname{Re}(l_n) \left[q_n \sin\left(q_n \frac{\eta + \xi}{2}\right) \sin\left(q_n \frac{\eta - \xi}{2}\right) \right]^{-1} \sin(q_n x),$$

$$\phi_m^1(x) = -2\sum_{n=1}^m \text{Im}(l_n) \left[q_n \sin\left(q_n \frac{\eta + \xi}{2}\right) \sin\left(q_n \frac{\eta - \xi}{2}\right) \right]^{-1} \frac{q_n^2}{\sqrt{1 + \alpha q_n^2}} \sin(q_n x).$$

- As mentioned in Remark 4.2, we assume (1.3) holds. This fact implies the sequence $\{(\phi_m^0,\phi_m^1)\}_{m\in\mathbb{N}^*}$ of initial data is well-defined.
- Since $0 < \sum_{n \in \mathbb{Z}^*} |l_n|^2 < \infty$ and $l_n = \overline{l_{-n}}$, then there exists a $m_0 \ge 1$ such that $l_{m_0} \ne 0$. So for any $\beta \ge -1$ and $m \ge m_0$, we have 612

614 (4.36)
$$\|\phi_m^0\|_{H^{-\beta}}^2 + \|\phi_m^1\|_{H^{-\beta-1}}^2 \ge \|\phi_{m_0}^0\|_{H^{-\beta}}^2 + \|\phi_{m_0}^1\|_{H^{-\beta-1}}^2 = c > 0.$$

Moreover, thanks to (4.7), we have 615

616 (4.37)
$$\int_0^T |\phi_{m,x}(\eta,t) - \phi_{m,x}(\xi,t)|^2 dt = 4 \int_0^T \left| \sum_{n=-m,n\neq 0}^m l_n e^{\lambda_{q_n} t} \right|^2 dt.$$

Since $0 = f(t) = \sum_{n \in \mathbb{Z}^*} l_n e^{\lambda_{q_n} t}$ in $L^2(0,T)$, we obtain that

618 (4.38)
$$\int_0^T |\phi_{m,x}(\eta,t) - \phi_{m,x}(\xi,t)|^2 dt \to 0, \text{ as } m \to \infty.$$

- Relations (4.36) and (4.38) finish the proof of the lack of exact L^2 -controllability for 619
- any $\beta \geq -1$, $0 < T < 2\pi\sqrt{\alpha}$ and ξ , η in $(0, \pi)$. 620
- As for the lack of exact $(H^1)'$ -controllability for any $\beta \geq -2$, $0 < T < 2\pi\sqrt{\alpha}$ and 621
- ξ , η in $(0,\pi)$, we only need to change the initial data in (4.35) to 622

$$\phi_m^0(x) = 2\sum_{n=1}^m \operatorname{Re}(l_n) \left[q_n \sin\left(q_n \frac{\eta + \xi}{2}\right) \sin\left(q_n \frac{\eta - \xi}{2}\right) \right]^{-1} \frac{\sqrt{1 + \alpha q_n^2}}{q_n^2} \sin(q_n x),$$

$$\phi_m^1(x) = -2\sum_{n=1}^m \operatorname{Im}(l_n) \left[q_n \sin\left(q_n \frac{\eta + \xi}{2}\right) \sin\left(q_n \frac{\eta - \xi}{2}\right) \right]^{-1} \sin(q_n x).$$

Similarly, (4.38) still holds for any $\beta \geq -2$ and $m \geq m_0$. Moreover, we have

625
$$\int_0^T |\phi_{m,xt}(\eta,t) - \phi_{m,xt}(\xi,t)|^2 dt = 4 \int_0^T \left| \sum_{n=-m,n\neq 0}^m l_n e^{\lambda_{q_n} t} \right|^2 dt \to 0, \quad \text{as } m \to \infty.$$

- Relations (4.36) and (4.40) finish the proof of the lack of exact $(H^1)'$ -controllability 626
- for any $\beta \geq -2$, $0 < T < 2\pi\sqrt{\alpha}$ and ξ , η in $(0, \pi)$.
- Appendix A. Proof of Lemma 4.5. Solve $\frac{k^2}{\sqrt{1+\alpha k^2}} < r$, we have k < 1628
- $[(\alpha r^2 + r\sqrt{\alpha^2 r^2 + 4})/2]^{\frac{1}{2}}$. Thus we define a real function $g:(0,+\infty)\to(0,+\infty)$ as 629
- $g(r) = [(\alpha r^2 + r\sqrt{\alpha^2 r^2 + 4})/2]^{\frac{1}{2}}$ and define $\{x\} = q(x) 1$, where $q(x) = \min_{q \in \mathbb{Z}} \{q \ge 1\}$ 630
- x. Obviously, $x-1 \leq \{x\} < x$. Then we have $N(0,r) = \{g(r)\}$. Notice that for 631
- $x \ge 0$, we have $N(x,r) = N(0,x+r) N(0,x) = \{g(x+r)\} \{g(x)\}$. Therefore, 632

633
$$g(x+r) - g(x) - 1 \le N(x,r) \le g(x+r) - g(x) + 1.$$

- Now we need to estimate g(x+r)-g(x) for x,r>0. Let $f(x)=\frac{x^2}{\sqrt{1+\alpha x^2}}$ for x>0. 634
- 635 Then

636
$$f'(x) = 2x(1+\alpha x^2)^{-\frac{1}{2}} - \alpha x^3(1+\alpha x^2)^{-\frac{3}{2}},$$

637
638
$$f''(x) = 2(1 + \alpha x^2)^{-\frac{1}{2}} - 5\alpha x^2 (1 + \alpha x^2)^{-\frac{3}{2}} + 3\alpha^2 x^4 (1 + \alpha x^2)^{-\frac{5}{2}}.$$

- Solve f''(x) > 0, we obtain $0 < x < \sqrt{2/\alpha}$. Therefore, there exists a N_0 in \mathbb{N} satisfying
- 640 $N_0 > \sqrt{2/\alpha} + 1$ such that for all $n \ge N_0$, $\lambda_{n+1} \lambda_n$ is decreasing to $1/\sqrt{\alpha}$. Then for
- all r > 0, N(x,r) is increasing in $x \ge N_0$. Then for $x \ge N_0$, $N(N_0,r) \le N(x,r) \le N(x,r)$
- 642 $\lim_{x\to+\infty} N(x,r)$. Simple calculation shows that $\lim_{x\to+\infty} [g(x+r)-g(x)] = \sqrt{\alpha}r$.
- 643 Therefore, we obtain that for $x \geq N_0$,

644
$$g(N_0 + r) - g(N_0) - 1 \le N(x, r) \le \sqrt{\alpha}r + 1.$$

645 For $0 \le x \le N_0$, we have

$$N(x,r) \le N(x,N_0-x) + N(N_0,r) = N(0,N_0) - N(0,x) + N(N_0,r)$$

$$\le N(0,N_0) + N(N_0,r) \le \sqrt{\alpha}r + 1 + g(N_0).$$

647 Assume that $r > N_0$, then we have

$$N(x,r) \ge N(N_0, r - N_0 + x) = N(0, x + r) - N(0, N_0)$$

$$\ge g(x+r) - g(N_0) - 1 \ge \min_{x \in [0, N_0]} g(x+r) - g(N_0) - 1.$$

Then for all $x \ge 0$ and $r > N_0$, we have

$$\min_{x \in [0, N_0]} g(x+r) - g(N_0) - 1 \le N(x, r) \le \sqrt{\alpha}r + 1 + g(N_0),$$

and hence,

652 (A.1)
$$\lim_{r \to \infty} \frac{N(x,r)}{r} = \sqrt{\alpha}$$

- 653 is uniformly relative to $x \geq 0$.
- For x < 0, if $|x| \le r$, then

$$N(0,|x|) + N(0,r-|x|) \le N(x,r) \le N(0,|x|) + N(0,r-|x|) + 1.$$

Let θ belong to [0, 1], and we consider $g(\theta r) + g((1-\theta)r)$. Notice that

657
$$g(\theta r) = \left(\frac{\alpha \theta^2 r^2 + \theta r \sqrt{\alpha^2 \theta^2 r^2 + 4}}{2}\right)^{\frac{1}{2}} = \frac{r}{\sqrt{2}} \left(\alpha \theta^2 + \theta \sqrt{\alpha^2 \theta^2 + 4/r^2}\right)^{\frac{1}{2}}.$$

- Therefore, we obtain $\lim_{r\to\infty} g(\theta r)/r = \sqrt{\alpha}\theta$. Consequently, (A.1) holds uniformly
- 659 relative to $-r \le x < 0$.
- If |x| > r, then

$$N(|x| - r, r) \le N(x, r) \le N(|x| - r, r) + 1.$$

- 662 Set $t = |x| r \ge 0$, then as same as in the situation $x \ge 0$, we have (A.1) holds
- uniformly relative to $t \geq 0$, which means that (A.1) holds uniformly relative to |x| > r.
- Combining all the situations, we obtain that (A.1) holds uniformly relative to x in \mathbb{R} . Lemma 4.5 is thus proved.

[1] K. Ammari, Z. Liu, and M. Tucsnak, Decay rates for a beam with pointwise force and moment feedback, Mathematics of Control Signals and Systems, 15 (2002), pp. 229–255. [2] S. AVDONIN AND M. TUCSNAK, Simultaneous controllability in sharp time for two elastic strings,
 ESAIM-Control Optimisation and Calculus of Variations, 6 (2001), pp. 259–273.

671

672

673

674

675

 $676 \\ 677$

678

679

680

681

682

683

684

685

686

687

688

689 690

691

692

700

 $701 \\ 702$

 $705 \\ 706$

707

 $708 \\ 709$

- [3] S. A. AVDONIN AND S. A. IVANOV, Families of exponentials. The method of moments in controllability problems for distributed parameter systems, Cambridge University Press, Cambridge; New York, 1995.
- [4] C. BAIOCCHI, V. KOMORNIK, AND P. LORETI, Ingham type theorem and applications to control theory, Bollettino Della Unione Matematica Italiana B, (1999).
- [5] J. W. S. Cassels, An introduction to diophantine approximation, Cambridge tracts in mathematics and mathematical physics, University Press, Cambridge Eng., 1957.
- [6] J.-M. CORON, Control and Nonlinearity, Mathematical surveys and monographs v. 136., American Mathematical Society, Providence, R.I, 2007.
- [7] E. CRAWLEY AND E. ANDERSON, Detailed models of piezoceramic actuation of beams, in 30th Structures, Structural Dynamics and Materials Conference, 1989.
- [8] E. CRÉPEAU AND C. PRIEUR, Control of a clamped-free beam by a piezoelectric actuator, ESAIM-Control Optimisation and Calculus of Variations, 12 (2006), pp. 545–563.
- P. Destuynder, I. Legrain, L. Castel, and N. Richard, Theoretical, numerical and experimental discussion on the use of piezoelectric devices for control-structure interaction, European Journal of Mechanics - A/Solids, 11 (1992), pp. 181–213.
- [10] C. Fabre and J.-P. Puel, Pointwise controllability as limit of internal controllability for the wave equation in one space dimension, Portugaliae Mathematica, 51 (1994).
- [11] S. M. HAN, H. BENAROYA, AND T. WEI, Dynamics of transversely vibrating beams using four engineering theories, Journal of Sound and Vibration, 225 (1999), pp. 935–988.
- [12] A. E. Ingham, Some trigonometrical inequalities with applications to the theory of series, Mathematische Zeitschrift, 41 (1936), pp. 367–379.
- [13] J. LAGNESE AND J. L. LIONS, Modelling analysis and control of thin plates, Recherches en mathématiques appliquées, Masson, Paris, 1988.
- [4] S. LANG, Introduction to diophantine approximations, Springer-Verlag, New York, new expanded ed., 1995.
- [697 [15] I. LASIECKA, J. L. LIONS, AND R. TRIGGIANI, Non homogeneous boundary value problems
 for second order hyperbolic operators, Journal de Mathématiques Pures et Appliqués, 65
 (1986), pp. 149–192.
 - [16] I. LASIECKA AND R. TRIGGIANI, A lifting theorem for the time regularity of solutions to abstract equations with unbounded operators and applications to hyperbolic equations, Proceedings of the American Mathematical Society, 104 (1988), pp. 745–755.
- 703 [17] J. L. LIONS, Exact controllability, stabilization and perturbations for distributed systems, SIAM 704 Review, 30 (1988), pp. 1–68.
 - [18] J. L. LIONS AND E. MAGENES, Non-Homogeneous Boundary Value Problems and Applications, Die Grundlehren der mathematischen Wissenschaften in Einzeldarstellungen, Springer-Verlag, Berlin, New York, 1972.
 - [19] A. OZKAN OZER AND S. W. HANSEN, Exact controllability of a Rayleigh beam with a single boundary control, Mathematics of Control, Signals and Systems, 23 (2011), pp. 199–222.
- [20] A. Pazy, Semigroups of Linear Operators and Applications to Partial Differential Equations,
 Applied mathematical sciences, Springer-Verlag, New York, 1983.
- [21] M. Tucsnak, Regularity and exact controllability for a beam with piezoelectric actuator, SIAM
 Journal on Control and Optimization, 34 (1996), pp. 922–930.
- 714 [22] G. Valiron, Théorie des fonctions, Masson, Paris, 1990.
- 715 [23] G. Weiss, O. J. Staffans, and M. Tucsnak, Well-posed linear systems—a survey with empha-716 sis on conservative systems, International Journal of Applied Mathematics and Computer 717 Science, 11 (2001), pp. 7–33.