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Treating inequality constraint in a continous medium with a coupled X-FEM/Level-Set strategy

N. Bonfils¹, N. Chevaugeon¹, N. Moës¹

¹ GeM, Ecole Centrale de Nantes, 1 rue de la Noë 44321 Nantes, {nicolas.bonfils, nicolas.chevaugeon, nicolas.moes}@ec-nantes.fr

Introduction

The aim of this study is to propose a new approach to deal with continuum media subject to a kinematic inequality constraint $\gamma(\underline{\underline{K}}) \leq 0$, where $\underline{\underline{K}}$ may be any kinematic second-order tensor. This kind of problem is similar to a constrained minimization problem with inequality constraints. It is equivalent to :

$$\inf_{\underline{u}\in S}\Pi(\underline{u})$$

with $S = \{\underline{u} \mid \gamma(\underline{\underline{K}}) \leq 0\}$ and $\Pi(\underline{u})$ the total potential energy of the sytem. The common approaches to imposing the inequality constraints classically the lagrange multiplier method [1], the penalty method [2], and the augmented lagrangian method [3]. Excluding the penalty method, all this approaches lead to variational inequalities because we do not know a priori the shape of the active constraint and unactive constraint zone in the medium. To take into account such a constraint the elastic strain energy function Ψ can be split into two distinct parts :

$$\Psi(\underline{C}) = \Psi^{r}(\underline{C}) + I_{-}(\gamma(\underline{K}))$$

with ψ^r , the parent or regular potential, $\underline{\underline{C}}$, the right Cauchy-Green tensor, and $I_{-}(x)$, the characteristic function of \mathbb{R}^- such that :

$$I_{-}(x) = \begin{cases} 0 & \text{if } x \le 0\\ +\infty & \text{else} \end{cases}$$

The level-set approach

The strategy proposed here is to split the whole domain Ω into two distinct subdomains separated by Γ . The first domain denoted by Ω_c is the constrained zone or active zone where an equality constraint is imposed. The second domain denoted by Ω_l is the unconstrained or unactive zone, where no constraints are imposed.

The two subdomains are separated by the interface Γ , located by the iso-zero of a level-set function [4]. On each subdomains, we have:

$$\begin{split} \Psi(\underline{\underline{C}}) &= \Psi^{r}(\underline{\underline{C}}) + I_{0}(\gamma(\underline{\underline{K}}) \ \forall \underline{\underline{x}} \in \Omega_{c} \\ \Psi(\underline{\underline{C}}) &= \Psi^{r}(\underline{\underline{C}}) \ \forall \underline{\underline{x}} \in \Omega_{l} \end{split}$$

with,

$$I_0(x) = \begin{cases} 0 & \text{if } x = 0\\ +\infty & \text{else} \end{cases}$$

The location of the interface Γ becomes an unknown of the problem and, thus, we are looking for the exact position of the interface such that the total potential energy is stationary :

$$\inf_{u \in S'} \Pi(\underline{u}, \Gamma) \text{ avec } S' = \{ \underline{u} \mid \gamma(\underline{\underline{K}}) = 0 \text{ sur } \Omega_c \}$$
(1a)

$$D_{\delta q_n} \Pi(\underline{u}(\Gamma), \Gamma) = 0 \,\forall \, \delta q_n \tag{1b}$$

where $D_{\delta q_n} \Pi$ indicates the directional derivative of Π in the direction of δq_n , where δq_n denotes a velocity field modifying the location of the interface Γ .

Starting from an initial position Γ_0 of the interface, we search the exact position iteratively. At each step of the algorithm a velocity is computed form the stationary condition (1) and a level-set equation is solved to find the new location of the interface.

First results and future prospects

We have applied the method described above to solve a simple one dimensional problem in the small strain elastostatic setting. We wish to limit the uniaxial strain of a truss by an arbitrary value α under a prescribed load. Since we know the exact solution, we can compare it to the numerical one. The first results are promising since the presented strategy provides a convergence rate equals to the optimal rate of convergence in spite of the lack of regularity of the exact solution. The level-set localization of the interface and a correct enrichment of the finite element approximation [5] permits the optimal convergence. Compared to an augmented Lagrangian method coupled with an active-set algorithm [6], we obtain a significant gain in accuracy. Preliminary results have been obtained for two dimensional cases and will be presented.

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