



HAL
open science

Large deviations for out of equilibrium correlations in the symmetric simple exclusion process

Thierry Bodineau, Benoit Dagallier

► **To cite this version:**

Thierry Bodineau, Benoit Dagallier. Large deviations for out of equilibrium correlations in the symmetric simple exclusion process. 2022. hal-03912965v1

HAL Id: hal-03912965

<https://hal.science/hal-03912965v1>

Preprint submitted on 26 Dec 2022 (v1), last revised 2 May 2024 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Large deviations for out of equilibrium correlations in the symmetric simple exclusion process

Thierry Bodineau*

Benoit Dagallier†

December 26, 2022

Abstract

For finite size Markov chains, the Donsker-Varadhan theory fully describes the large deviations of the time averaged empirical measure. We are interested in the extension of the Donsker-Varadhan theory for a large size non-equilibrium system: the one-dimensional symmetric simple exclusion process connected with reservoirs at different densities. The Donsker-Varadhan functional encodes a variety of scales depending on the observable of interest. In this paper, we focus on the time-averaged two point correlations and investigate the large deviations from the steady state behaviour. To control two point correlations out of equilibrium, the key input is the construction of a simple approximation to the invariant measure. This approximation is quantitative in time and space as estimated through relative entropy bounds building on the work of Jara and Menezes [JM18].

1 Introduction

For a fluid in thermal equilibrium, spatial correlations are expected to have fast decay, in such a way that, roughly speaking, each macroscopic portion of the fluid is basically independent from the rest. For fluids driven out of equilibrium, e.g. by contact with reservoirs at two different temperatures, the picture is quite different: the fluid settles in a steady state where heat and/or matter are transported at a macroscopic level. The transport induces long-range correlations, which can be modelled by a variety of approaches and that have been observed experimentally, see [Spo83][Gar+90] and references therein. These general predictions are part of the results of the Macroscopic Fluctuation Theory (see the review [Ber+15]), which proposes a framework to study out of equilibrium fluids at a macroscopic level.

The derivation, from a microscopic model, of the steady state correlations, which are of a genuinely dynamical nature, is usually a difficult problem. Rigorous results are mostly obtained for certain simple interacting particle systems on a lattice. The Symmetric Simple Exclusion Process connected with reservoirs (henceforth open SSEP) is a paradigmatic example for which this

*I.H.E.S., Université Paris-Saclay, CNRS, Laboratoire Alexandre Grothendieck. 35 Route de Chartres, 91440 Bures-sur-Yvette (France). E-mail: bodineau@ihes.fr.

†University of Cambridge, DPMMS, Statistical Laboratory. E-mail: bd444@cam.ac.uk.

correlation structure can be analysed [Der+05][LMO06]. In the open SSEP, defined in Section 2, particles follow symmetric random walks interacting by an exclusion rule on a finite subdomain of \mathbb{Z} . Reservoirs pump particles in and out of the system, fixing a certain density of particles in their vicinity. When reservoirs are at the same density, the open SSEP dynamics is reversible. However, when connected with reservoirs which enforce a different density of particles, this dynamics settles in long time in a non-equilibrium steady state, characterised by a macroscopic current of particles. The strength of this current is proportional to the density difference between the reservoirs [Spo83][ELS90][FLM11]. Two-point correlations in the steady state are known exactly [Spo83], as well as all higher cumulants [DLS07][Der07][DR13] in dimension one. The correlation structure of the steady state of the open SSEP is conjectured to be representative of a large class of out of equilibrium systems [Spo83][Ber+15]. However, much less is known rigorously about steady state correlations for general lattice gases.

Our goal is to estimate the asymptotic probability of observing a correlation structure that is different from the one of the steady state, thereby also gaining information on this invariant measure. When the value N of the scaling parameter is fixed, this question has already received a comprehensive answer by Donsker and Varadhan [DV75]. For a general, irreducible Markovian dynamics on a finite state space Ω_N , they study the time empirical measure $\tilde{\pi}_T$, defined for each $T > 0$ as a probability measure on the configuration space Ω_N by:

$$\forall \eta \in \Omega_N, \quad \tilde{\pi}^T = \frac{1}{T} \int_0^T \delta_{\eta_t} dt. \quad (1.1)$$

The quantity $\tilde{\pi}^T(\eta)$ then corresponds to the proportion of time spent at a configuration $\eta \in \Omega_N$. A full large deviation principle with speed T and rate function I_{DV}^N is then provided in [DV75] for the time empirical measure $\tilde{\pi}_T$, in the sense that, if μ^N is a probability measure on Ω_N and \mathbb{P} denotes the probability associated with the dynamics:

$$\lim_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P}(\tilde{\pi}^T \approx \mu^N) = -I_{DV}^N(\mu^N), \quad (1.2)$$

where \approx means proximity in the weak topology of probability measures on Ω_N . The rate function I_{DV}^N vanishes only at the invariant measure π_{inv}^N of the dynamics, and is defined through a complicated variational problem involving the generator L of the Markov chain:

$$I_{DV}^N(\mu^N) := \sup_{h: \Omega_N \rightarrow \mathbb{R}} \mu^N \left(e^{-h} (-L) e^h \right). \quad (1.3)$$

We are interested in the macroscopic behaviour of the system, i.e. the large N limit of the probability in (1.2). To obtain these asymptotics, one possibility is to study the limit of (1.3) when $N \rightarrow \infty$. When the underlying dynamics is reversible, this can be carried out: the variational problem (1.3) can be solved, and I_{DV}^N is expressed in terms of the Dirichlet form of the dynamics (see Section 2.2 below). Such computations are carried out in Section 2.2, in the case of the open SSEP with reversible dynamics (where the scaling parameter N is roughly the number of sites in the model). These computations highlight an important fact: one cannot naively take the large N limit in (1.2) without losing information, because not all the information contained in a measure μ^N is stored at the same scale in N . By this we mean e.g. that observing a macroscopic density

different from that of the steady state, or observing different two point-correlation but with the same density, are not events that have the same scaling in terms of N . Informally, it is shown in Section 2.2 that observing a macroscopic density profile different from the one of the invariant measure in the reversible open SSEP has a probability that scales like $e^{-TN^{-2} \cdot N}$ in the large T , then large N limit, up to sub-exponential corrections. In contrast, changing the two-point correlation structure only requires a cost of order $e^{-TN^{-2}}$. To study the scaling limit of (1.2), one therefore has to choose a scale. Out of equilibrium, the rate function I_{DV}^N is not known explicitly, and the equilibrium heuristics cannot be used, but we prove in Theorem 2.3 below that scales are still separated in the same way.

In this article, we focus on the scale corresponding to two-point correlations, and quantify the probability of observing anomalous two-point correlations in the one dimensional, out of equilibrium open SSEP in the large T, N limits. We establish a large deviation principle for the time-averaged two-point correlation field, in Theorem 2.3 below. We do not start from the Donsker-Varadhan asymptotics, but instead provide quantitative estimates on the dynamics, as a function of time and the system size. Note that density large deviations are well understood since the seminal paper [KOV89] (see also Chapter 10 of [KL99] for a review). The main difficulty of the article is to generalise these ideas to the estimate of two-point correlations, that are objects living on a much finer scale than the density. To illustrate this, recall that the two-point correlations in the steady state of the open SSEP are long range, and scale like $O(N^{-1})$, compared to $O_N(1)$ for the density of particle at a given site. For this reason, and while model-dependent estimates on correlations have been obtained e.g. in [Spo83][Der+05][LMO06][Gon+20], to our knowledge there is no general method to study the out of equilibrium behaviour of the two-point correlation field in the long-time, large N limits.

The proof of our result on two-point correlations, Theorem 2.3, builds upon a refinement of the relative entropy method obtained by Jara and Menezes [JM18]-[JM20]. This method, originally introduced by Yau [Yau91], consists in quantifying, at each time and in terms of the relative entropy, the proximity of the law of the Markovian dynamics in an interacting particle system with a known reference measure. The idea behind the method is that, locally, the dynamics in large microscopic boxes equilibrates much faster than the typical time-scale at which the system evolves macroscopically. If one has an ansatz for the evolution of macroscopic variables of interest, say, henceforth, the density in a lattice gas; one then expects that the corresponding microscopic variables, when averaged over a sufficiently large microscopic box, are close to their macroscopic counterpart. This property, known as local equilibrium, has recently been shown quantitatively even for mesoscopic boxes, see [Gon+22].

From the local equilibrium heuristics, one can build a reference measure in terms of the evolution of the macroscopic density only. If one is interested in the evolution of the macroscopic density or its fluctuations, it can be shown that local equilibrium holds and this reference measure is indeed a good enough approximation of the law of the dynamics, see Chapter 6 in [KL99], [JM18] and references therein. In particular, the reference measure does not need to contain any information on correlations.

To study two point correlations, however, the reference measure has to also contain information on the dynamical correlations. Adding such a correlation term in the reference measure is our key input. In the case of the open SSEP, since density fluctuations around the typical density profile at each time (and in the steady state [LMO06]) are known to be Gaussian [JM18], our

candidates for reference measures are discrete Gaussian measures, see (2.58). One expects that a good choice of discrete Gaussian measure will contain all leading order information about two point correlations. This statement is made precise in Theorem 2.6 where we obtain, for a family of exclusion dynamics that occur in the proof of the large deviation result of Theorem 2.3, a characterisation of long time, large N correlations as the solution of a certain partial differential equation.

The approach used in this paper is not restricted to the symmetric simple exclusion process. Much like the usual relative entropy method, it can be used for a large class of one-dimensional diffusive interacting particle systems satisfying the so-called gradient condition (see Section 8 in [JM18]). There are however some technical difficulties to be expected when generalising the present approach. All these points are discussed further in Section 2.6.

Let us however mention that, at equilibrium, the behaviour of various n -point correlation fields has come under much scrutiny in the past few years. In [Ass07], a two-point correlation field is studied in the SSEP on \mathbb{Z} . It is not the same object as in [GJ19], where two point correlations are studied on the one-dimensional torus as a means to defining squares of distributions arising in certain ill-posed stochastic partial differential equations. In [ACR21] and [CS21], interacting particle systems enjoying a self-duality property are investigated in all dimensions. In that context, equilibrium fluctuation fields involving n -point functions are investigated for any n .

In the same direction but using different techniques, long time large deviations for the density and the current have recently been considered in [BGL21]. Both the long diffusive time limit starting from the dynamical large deviation functional, and the long time, then large N limits using Donsker and Varadhan's formula are investigated. This last limit is the same as the one studied in the present article at the level of two point correlations. The approach is however different: here, we provide a quantitative (i.e. non asymptotic) control on the dynamics at the microscopic level. On the other hand, in the case of the density and the current, the microscopic model considered in [BGL21] may have dynamical phase transitions. To capture this very subtle phenomenon, the Donsker-Varadhan variational principle (1.3) in [BGL21] is solved indirectly in the large N limit, by looking at large deviations at process level (so-called level 3 large deviations), then using a contraction argument. A related paper by the same authors [BGL22b] uses a similar approach in the settings of diffusions with small noise. Note that different kind of asymptotics for the Donsker-Varadhan functional have been considered, e.g. to study metastability [BGL22a; Lan22].

The rest of the article is structured as follows. In Section 2, we present the model and results. Section 3 gives the main microscopic tool for the study of two-point correlations: the relative entropy estimate when the reference measure is a certain discrete Gaussian measure. Properties of these measures are established in Appendix A. The relative entropy bounds allow for the computation of the Radon-Nikodym derivative between the open SSEP and the tilted processes introduced to estimate rare events. This requires sharp estimates collected in the appendices. The large deviations are then established, in Section 4 for the upper bound, and 5 for the lower bound. For the lower bound, control of the open SSEP dynamics in long-time is obtained via the study of certain Poisson equations. Well-posedness of these equations is investigated in Appendix F, while Appendix E gathers useful topological facts.

Acknowledgements: This work has been motivated by many discussions with Bernard Derrida on the structure of correlations in non-equilibrium particle systems. We are extremely grateful to

him for sharing his insights. We would also like to thank Stefano Olla for very useful suggestions and discussions at various stages of the writing process.

2 Notations and results

2.1 Notations and definition of the microscopic model

2.1.1 The microscopic model.

For $N \in \mathbb{N}^*$, let $\Lambda_N := \{-N + 1, \dots, N - 1\}$ and $\Omega_N = \{0, 1\}^{2N-1}$. Elements of Ω_N , denoted by the letter η , will be called configurations. We say that there is a particle at site $i \in \Lambda_N$ if $\eta_i = 1$, and no particle if $\eta_i = 0$. The variable η_i is called the occupation number (of site i). On Ω_N , we consider the dynamics given by the Symmetric Simple Exclusion Process connected to reservoirs at position $\pm N$ (henceforth open SSEP), which we now describe. For a survey of particle systems in contact with reservoirs, we refer to [Der07; Ber+15].

Let $\rho_- < \rho_+ \in (0, 1)$ be the densities of the left (for ρ_-) and the right (ρ_+) reservoirs. The open SSEP is defined through its generator $N^2L := N^2(L_0 + L_- + L_+)$. It is made up of two parts, the bulk and boundary dynamics, corresponding to L_0 and L_\pm respectively. The operators L_0 and L_\pm act on $f : \Omega_N \rightarrow \mathbb{R}$ as follows:

$$\forall \eta \in \Omega_N, \quad N^2L_0f(\eta) = \frac{N^2}{2} \sum_{i < N-1} c(\eta, i, i+1) [f(\eta^{i,i+1}) - f(\eta)], \quad (2.1)$$

$$N^2L_\pm f(\eta) = \frac{N^2}{2} c(\eta, \pm(N-1)) [f(\eta^{\pm(N-1)}) - f(\eta)]. \quad (2.2)$$

Above, the jump rates c are defined, for each $\eta \in \Omega_N$, by:

$$\begin{aligned} c(\eta, i, i+1) &= \eta_{i+1}(1 - \eta_i) + \eta_i(1 - \eta_{i+1}), & i < N - 1, \\ c(\eta, \pm(N-1)) &= (1 - \rho_\pm)\eta_{\pm(N-1)} + \rho_\pm(1 - \eta_{\pm(N-1)}), \end{aligned} \quad (2.3)$$

and for $i, j \in \Lambda_N$ and $\eta \in \Omega_N$, the configurations η^i and $\eta^{i,j}$ read:

$$\eta_\ell^i = \begin{cases} \eta_\ell & \text{if } \ell \neq i, \\ 1 - \eta_i & \text{if } \ell = i, \end{cases} \quad \eta_\ell^{i,j} = \begin{cases} \eta_\ell & \text{if } \ell \notin \{i, j\}, \\ \eta_j & \text{if } \ell = i, \\ \eta_i & \text{if } \ell = j. \end{cases} \quad (2.4)$$

We write $\mathbb{P}^\eta, \mathbb{E}^\eta$ for the probability/expectation under this dynamics starting from $\eta \in \Omega_N$.

2.1.2 The invariant measure and the correlation field.

For each $N \in \mathbb{N}^*$, let π_{inv}^N denote the unique invariant measure of the open SSEP. If $\rho_+ = \rho_- = \rho \in (0, 1)$, π_{inv}^N is simply the Bernoulli product measure on Λ_N with parameter ρ . If $\rho_- < \rho_+$, however, the measure π_{inv}^N is not product. The average occupation number at each site was computed in [Spo83]: it is given in terms of an affine function $\bar{\rho}$, with:

$$\forall i \in \Lambda_N, \quad \bar{\rho}(i/N) := \mathbb{E}^{\pi_{inv}^N}[\eta_i] = \left(1 - \frac{i}{N}\right) \frac{\rho_-}{2} + \left(1 + \frac{i}{N}\right) \frac{\rho_+}{2}. \quad (2.5)$$

Note that, as $\bar{\rho}$ is affine, $\bar{\rho}' = (\rho_+ - \rho_-)/2$ is a constant. As π_{inv}^N is invariant, one has, for $F : \Omega_N \Rightarrow \mathbb{R}$:

$$\forall t \geq 0, \quad \pi_{inv}^N(F) := \mathbb{E}^{\pi_{inv}^N}[F(\eta_t)] = \sum_{\eta \in \Omega_N} \pi_{inv}^N(\eta) F(\eta). \quad (2.6)$$

In the following, the expectation $\mathbb{E}_\mu[F]$ with respect to a measure μ is abbreviated $\mu(F)$.

The measure π_{inv}^N exhibits long-range correlations. To make this statement precise, let us first define the main object of interest in this article, the *correlation field* Π^N . It is a distribution, acting on test functions $\phi : (-1, 1)^2 \rightarrow \mathbb{R}$ according to:

$$\forall \eta \in \Omega_N, \quad \Pi^N(\phi) = \Pi^N(\phi)(\eta) = \frac{1}{4N} \sum_{i \neq j \in \Lambda_N} \bar{\eta}_i \bar{\eta}_j \phi(i/N, j/N), \quad \bar{\eta} \cdot := \eta \cdot - \bar{\rho}(\cdot/N). \quad (2.7)$$

Note the scaling of (2.7) with N : the sum contains N^2 terms, and the normalisation is only proportional to N^{-1} as $\Pi^N(\phi)$ measures fluctuations of a central limit order. Indeed the correlation field Π^N is strongly related to the fluctuation field Y^N defined for any bounded test function $\psi : [-1, 1] \rightarrow \mathbb{R}$ by:

$$\forall \eta \in \Omega_N, \quad Y^N(\psi) = \frac{1}{\sqrt{N}} \sum_{i \in \Lambda_N} \bar{\eta}_i \psi(i/N). \quad (2.8)$$

Restricting to product test functions $\phi(x, y) = \psi_1(x)\psi_2(y)$, we note that:

$$\Pi^N(\phi)(\eta) = \frac{1}{4} Y^N(\psi_1) Y^N(\psi_2) - \frac{1}{4N} \sum_{i \in \Lambda_N} (\bar{\eta}_i)^2 \psi_1(i/N) \psi_2(i/N). \quad (2.9)$$

The fluctuation field has been extensively studied in the hydrodynamic scaling and at each time t , the fluctuation field $Y_t^N(\psi)$ can be proven to converge to a Gaussian random variable (see Chapter 11 in [KL99] in the equilibrium case, or [Der+05][LMO06][JM18] when $\rho_- \neq \rho_+$). The two-point correlations $Y_t(\psi_1)Y_t(\psi_2)$ are thus also of order 1 in N (see e.g. [GJ19] for the reversible SSEP on the torus).

As mentioned in the introduction, the correlation field Π^N is a natural quantity to study non equilibrium properties. For SSEP, the average value of the correlation field under the invariant measure was obtained in [Spo83][DLS02] in the large N limit :

$$\lim_{N \rightarrow \infty} \pi_{inv}^N(\Pi^N(\phi)) = \frac{1}{4} \int_{(-1, 1)^2} \phi(x, y) k_0(x, y) dx dy, \quad (2.10)$$

where the kernel k_0 is a symmetric continuous function on $[-1, 1]^2$ given by :

$$\forall (x, y) \in (-1, 1)^2, \quad k_0(x, y) = -\frac{(\bar{\rho}')^2}{2} \left[(1+x)(1-y) \mathbf{1}_{x \leq y} + (1+y)(1-x) \mathbf{1}_{x \geq y} \right]. \quad (2.11)$$

Thus when $\bar{\rho}' = (\rho_+ - \rho_-)/2 \neq 0$, the system is driven out of equilibrium by the reservoirs and long range correlations arise. These correlations come from the diffusive transport of particles and, as a consequence, k_0 is obtained as the solution of the Laplace equation restricted to the triangle $\triangleright = \{(x, y) \in [-1, 1]^2 : x < y\}$ with fixed normal derivative on the diagonal $D = \{(x, x) : x \in (-1, 1)\}$:

$$\forall (x, y) \in \triangleright, \quad \Delta k_0(x, y) = 0, \quad \text{with} \quad (\partial_1 - \partial_2) k_0(x_\pm, x) = \pm (\bar{\rho}')^2, \quad x \in (-1, 1), \quad (2.12)$$

with 0 Dirichlet boundary conditions on $\{x = -1\}$ and $\{y = 1\}$. Above, we used the notation $x_{\pm} = \lim_{h \downarrow 0} (x \pm h)$.

In the next section, we provide further insight on the correlation field and explain how Π^N is related to the Donsker-Varadhan functional.

Notation. *Throughout this paper, we write $\bar{\rho}_i := \bar{\rho}(i/N)$ for $i \in \Lambda_N$. More generally, for a function $\phi : [-1, 1]^p \rightarrow \mathbb{R}$, $p \in \mathbb{N}^*$, we write ϕ_{i_1, \dots, i_p} for $\phi(i_1/N, \dots, i_p/N)$, $(i_1, \dots, i_p) \in \Lambda_N^p$. The letters i, j, ℓ, \dots , when used as indices, index elements of Λ_N ; while x, y, z are used for continuous variables.*

More generally, when we speak of n -point correlations ($n \in \mathbb{N}^$), it will always mean products of centred variables, the $\bar{\eta}$'s, of the form $\bar{\eta}_{i_1} \dots \bar{\eta}_{i_n}$ for some $(i_1, \dots, i_n) \in \Lambda_N^n$. When considering a trajectory $\eta_t \in \Omega_N, t \geq 0$, we write Π_t^N for $\Pi^N(\cdot)(\eta_t)$.*

2.2 Heuristic scaling of the Donsker-Varadhan functional

In this section, we consider the SSEP dynamics at equilibrium as given in (2.1)–(2.2), i.e. with two reservoirs at equal densities $\rho_- = \rho_+ = \rho \in (0, 1)$. The dynamics is reversible with respect to the Bernoulli product measure ν_{ρ}^N with parameter ρ . In this setting, the Donsker-Varadhan rate function is explicit and is given by the Dirichlet form of the dynamics [DV75]: if μ^N is a probability measure on the state space Ω_N , then one has for each fixed N :

$$\lim_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P}_{\rho}^N \left(\frac{1}{T} \int_0^T \delta_{\eta_t} dt \approx \mu^N \right) = -I_{DV}^N(\mu^N) = -\nu_{\rho}^N(\sqrt{f}(-N^2 L)\sqrt{f}), \quad f := \frac{d\mu^N}{d\nu_{\rho}^N}, \quad (2.13)$$

where the \approx sign denotes proximity in the weak topology of measures on Ω_N .

In the next two sections, we use the explicit form (2.13) of the Donsker-Varadhan functional to derive its asymptotics as $N \rightarrow \infty$. We will show that the scalings in N are different for the cost of observing an anomalous macroscopic density (in Section 2.2.1) or the cost of observing anomalous macroscopic two-point correlations (in Section 2.2.2). Thus different types of information are intertwined in the Donsker-Varadhan functional and one has to zoom at the correct scaling to extract the relevant physical information on a given observable. The exact computation of Section 2.2.2 justifies our choice of focusing on the non-equilibrium large deviations of the two-point correlation field Π^N (see Section 2.4).

2.2.1 Changing the macroscopic density

Consider a smooth density profile $\hat{\rho} : [-1, 1] \rightarrow (0, 1)$. For simplicity, assume that the density close to the boundaries is unchanged, i.e. that $\hat{\rho}(x) = \rho$ for x in an open neighbourhood of ± 1 . Define then:

$$\mu^N = \bigotimes_{i \in \Lambda_N} \text{Ber}(\hat{\rho}(i/N)). \quad (2.14)$$

For a function $q : (-1, 1) \rightarrow \mathbb{R}$, write for short $q_i := q(i/N)$ for $i \in \Lambda_N$. Introduce the chemical potential λ and its discrete derivative $\partial^N \lambda_i$, defined by:

$$\forall x \in (-1, 1), \quad \lambda(x) := \log \left(\frac{\hat{\rho}(x)}{1 - \hat{\rho}(x)} \right), \quad \partial^N \lambda_i = N[\lambda_{i+1} - \lambda_i], \quad i < N - 1. \quad (2.15)$$

Recall that $\mu^N[\cdot]$ (or $\mu^N(\cdot)$) denotes expectation under the measure μ^N . Elementary computations using (2.13) then give, for each $N \in \mathbb{N}^*$:

$$I_{DV}^N(\mu^N) = -\frac{N^2}{4}\mu^N\left[\sum_{i < N-1} c(\eta, i, i+1)\left(\exp\left[-\frac{(\eta_{i+1}-\eta_i)}{2N}\partial^N\lambda_i\right]-1\right)^2\right]. \quad (2.16)$$

Note that there is no contribution from the boundary dynamics, since $\hat{\rho}$ is constant and equals to ρ close to ± 1 by assumption. Expanding the right-hand side of (2.16), one finds, with the notation $\sigma(r) = r(1-r)$ for $r \in [0, 1]$:

$$\begin{aligned} \lim_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P}_\rho^N\left(\frac{1}{T} \int_0^T \delta_{\eta_t} dt \approx \mu^N\right) &= -\frac{1}{16}\mu^N\left[\sum_{|i| < N-1} c(\eta, i, i+1)[\partial^N\lambda_i]^2\right] + O_N(1) \\ &= -\frac{N}{8} \int_{(-1,1)} \sigma(\hat{\rho}(x)) |\nabla\lambda(x)|^2 dx + O_N(1), \end{aligned} \quad (2.17)$$

where we used the smoothness of $\hat{\rho}$, and $\mu^N(c(\eta, i, i+1)) = 2\sigma(\hat{\rho}_i) + O(N^{-1})$ for each $i < N-1$, with the $O(N^{-1})$ uniform on i . It follows that a macroscopic change of density is observed with probability of order e^{-TN} in the large T , then large N limit.

As a remark, notice that, up to factors of N, T , the right-hand side of (2.17) is the same as the one given by the dynamical rate functional obtained in diffusive time in [Ber+03]. To see it, recall that the rate functional I_{SSEP} evaluated at the constant profile $\hat{\rho}$ on the time interval $[0, T]$ for $T > 0$ is given by:

$$I_{SSEP}((\hat{\rho}(x)dx)_{t \leq T}) = \frac{1}{2} \int_0^T \int_{(-1,1)} |\nabla h(t, x)|^2 \sigma(\hat{\rho}(x)) dx dt, \quad (2.18)$$

where h is the bias such that $h(t, \pm 1) = 0$ for each $t \in [0, T]$, and:

$$\partial_t \hat{\rho} = 0 = \frac{1}{2} \Delta \hat{\rho} - \nabla \cdot (\sigma(\hat{\rho}) \nabla h). \quad (2.19)$$

In particular, integrating the divergence operator, there is a divergence-free function j on $(-1, 1)$ (the current), i.e. a constant in our one-dimensional setting, such that:

$$\nabla h = \frac{(1/2)\nabla \hat{\rho} + j}{\sigma(\hat{\rho})} \Rightarrow I_{SSEP}((\hat{\rho}(x)dx)_{t \leq T}) = \frac{T}{2} \int_{(-1,1)} \frac{((1/2)\nabla \hat{\rho} + j)^2}{\sigma(\hat{\rho}(x))} dx. \quad (2.20)$$

In the present case, $h(\cdot, \pm 1) = 0$ and $\nabla \lambda = \nabla \hat{\rho} / \sigma(\hat{\rho})$ implies that $j = 0$, with λ defined in (2.15). As a result:

$$\forall t \geq 0, x \in (-1, 1), \quad \nabla h(t, x) = \nabla \lambda(x) / 2. \quad (2.21)$$

Thus replacing h by $\nabla \lambda / 2$ in (2.18), the functional (2.17) is recovered. In other words, the long time, large N limit and the long diffusive time limits coincide :

$$\lim_{N \rightarrow \infty} \lim_{T \rightarrow \infty} \frac{1}{NT} \log \mathbb{P}_\rho^N\left(\frac{1}{T} \int_0^T \delta_{\eta_t} dt \approx \mu^N\right) = \lim_{T \rightarrow \infty} \lim_{N \rightarrow \infty} \frac{1}{NT} \log \mathbb{P}_\rho^N\left((\pi_t^N)_{t \leq T} \approx (\hat{\rho}(x)dx)_{t \leq T}\right). \quad (2.22)$$

It is not difficult to adapt the proof of [Ber+03], derived in the hydrodynamic regime, to recover (2.22), i.e. to take the long time limit first before taking N large. Indeed the large deviation functional of the SSEP is convex, so that the optimal way to observe an averaged density profile in the long diffusive time limit is obtained by a time-independent tilt : there is no dynamical phase transitions. The exchange of limits (2.22) then remains valid in non equilibrium in the absence of a dynamical phase transition. The much more delicate proof that the two limits coincide even in the presence of a dynamical phase transition is carried out in [BGL21], where the joint deviations of the current and density are investigated.

2.2.2 Changing the macroscopic correlations

Consider again the open SSEP at equilibrium at density $\rho \in (0, 1)$. Our aim is to consider the large size asymptotics for a measure with the equilibrium density, but different correlations. Recall the definition (2.7) of the (off diagonal) correlation field Π^N , acting on a bounded test function $\phi : [-1, 1]^2 \rightarrow \mathbb{R}$ according to:

$$\forall \eta \in \Omega_N, \quad \Pi^N(\phi) = \frac{1}{4N} \sum_{i \neq j \in \Lambda_N} \bar{\eta}_i \bar{\eta}_j \phi\left(\frac{i}{N}, \frac{j}{N}\right). \quad (2.23)$$

In view of the above discussion, to find a measure that is close to ν_ρ^N but with a different correlation structure, it is reasonable to look at:

$$\mu^N = \nu_{\rho, \phi}^N := \frac{1}{\mathcal{Z}_{\rho, \phi}^N} e^{2\Pi^N(\phi)} \nu_\rho^N, \quad \mathcal{Z}_{\rho, \phi}^N \text{ a normalisation factor.} \quad (2.24)$$

Assume $\|\phi\|_\infty$ is sufficiently small and ϕ is smooth, symmetric, i.e.:

$$\forall (x, y) \in [-1, 1]^2, \quad \phi(x, y) = \phi(y, x). \quad (2.25)$$

Assume also that $\phi(x, \cdot) = 0$ for x in an open neighbourhood of ± 1 . Then the macroscopic density is still given by ρ , but the measure μ^N now features long-range correlations. Indeed, using the same kind of arguments as in Appendix A.1 (see also [Dag22]), one can show that there is a limiting kernel $k : \square \rightarrow \mathbb{R}$ such that:

$$\sup_{i \neq j \in \Lambda_N} |\mu^N(\bar{\eta}_i \bar{\eta}_j) - k(i/N, j/N)| = o(N^{-1}), \quad \sup_{i \in \Lambda_N} |\mu^N(\bar{\eta}_i \bar{\eta}_i) - \sigma(\rho)| = o_N(1). \quad (2.26)$$

The limiting covariance can therefore be described by an operator C acting on $\psi_1, \psi_2 \in \mathbb{L}^2((-1, 1))$ according to (recall the definition (3.26) of the fluctuation field Y^N):

$$\lim_{N \rightarrow \infty} \mu^N(Y^N(\psi_1)Y^N(\psi_2)) = \int_{(-1, 1)} \psi_1(x)(C\psi_2)(x) dx, \quad (2.27)$$

with:

$$\forall x \in (-1, 1), \quad C\psi_2(x) := \sigma(\rho)\psi_2(x) + \int_{(-1, 1)} k(x, y)\psi_2(y) dy. \quad (2.28)$$

The operator C is obtained from ϕ as the inverse of the operator U_ϕ , defined for $\psi \in \mathbb{L}^2((-1, 1))$ by:

$$\forall x \in (-1, 1), \quad U_\phi\psi(x) := \sigma(\rho)^{-1}\psi(x) - \int_{(-1, 1)} \phi(x, y)\psi(y) dy. \quad (2.29)$$

Intuitively, this relation means that for large N , the density under the measure μ^N in (2.24) behaves as a Gaussian field of covariance C . A similar structure will be derived out of equilibrium and the proofs will be given then, see for instance Theorem 2.6 where it is stated that the out of equilibrium SSEP dynamics stays very close to a measure of the form μ^N at each time, in the sense that it has the same two-point correlation structure. To summarise, the macroscopic density is still given by ρ , but the measure μ^N has now long-range correlations parametrised by ϕ . Let us again compute the Donsker-Varadhan rate functional for the measure μ^N with large N .

Lemma 2.1. *For ϕ and μ^N as in (2.24), one has with the notation (2.28) :*

$$\lim_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P} \left(\frac{1}{T} \int_0^T \delta_{\eta_t} dt \approx \mu^N \right) = -\frac{1}{8} \int_{(-1,1)} \sigma(\rho) \langle \partial_1 \phi(z, \cdot), C \partial_1 \phi(z, \cdot) \rangle dz + o_N(1), \quad (2.30)$$

where $\langle \cdot, \cdot \rangle$ is the scalar product in $\mathbb{L}^2((-1, 1))$.

Proof. Plugging the expression (2.24) of the measure μ^N in the Donsker-Varadhan functional (2.13), we obtain a formula similar to (2.16) where the variation of the chemical potential is now replaced by the non local expression

$$\Pi^N(\phi)(\eta^{i,i+1}) - \Pi^N(\phi)(\eta) = -\frac{(\eta_{i+1} - \eta_i)}{2N^2} \sum_{j \notin \{i, i+1\}} \bar{\eta}_j \partial_1^N \phi_{i,j}, \quad (2.31)$$

with $\partial_1^N \phi_{i,j} = N[\phi_{i+1,j} - \phi_{i,j}]$ bounded uniformly in i, j, N . Expanding the exponential, we get as in (2.17)

$$\begin{aligned} \lim_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P}_\rho^N \left(\frac{1}{T} \int_0^T \delta_{\eta_t} dt \approx \mu^N \right) \\ = -\frac{1}{16N^2} \mu^N \left[\sum_{\substack{|i| < N-1 \\ j, \ell \notin \{i, i+1\}}} c(\eta, i, i+1) \bar{\eta}_j \bar{\eta}_\ell \partial_1^N \phi_{i,j} \partial_1^N \phi_{i,\ell} \right] + o_N(1). \end{aligned} \quad (2.32)$$

The limit (2.30) can be easily guessed from the above formula, as $2\sigma(\rho)$ arises from $c(\eta, i, i+1)$ and the correlations $\bar{\eta}_j \bar{\eta}_\ell$ are approximated by the limiting covariance (2.26) of μ^N . Similar computations will be carried out numerous times in Section 3.2, so we give no details here. \square

Compared with the asymptotics of the density large deviations (2.17), the cost (2.30) of modifying only the correlations has a different scaling in N . Our aim is to derive similar results for systems driven out of equilibrium by reservoirs when the Donsker-Varadhan rate function is no longer given by the explicit formula (2.13) but by the variational principle (1.3). The computations in the reversible case highlight the fact that:

- the whole correlation structure of the invariant measure is not contained at the same scale in N , and therefore:
- one has to focus on observables at a specific scale to get a non trivial limit when N is large.

In the following, we focus on the large deviations of the correlation field Π^N introduced in (2.7), and generalise formula (2.30) to a non-equilibrium situation in Theorem 2.3.

2.3 The topology for correlations

Motivated by the heuristics of Section 2.2, we focus in this article on the next scale after the density, and consider the two point correlation field Π^N , defined in (2.7). Zooming at the level of correlations amounts to rewriting the asymptotics of Lemma 2.1 as follows: compute the asymptotics of observing a given correlation field Π

$$\mathbb{P}\left(\frac{1}{T} \int_0^T \Pi_t^N dt \stackrel{weak^*}{\approx} \Pi\right) \quad \text{when } T \text{ and then } N \text{ are large.} \quad (2.33)$$

In (2.33), $\stackrel{weak^*}{\approx}$ means proximity in the weak* topology. In this section, we start by defining the functional space to which Π^N belongs and the associated weak* topology.

Let us start with a few observations. By definition (2.7), Π^N can be seen as a linear form on several function spaces. Let $\square = (-1, 1)^2$ and notice that Π^N is symmetric in the following sense:

$$\forall \phi : \square \rightarrow \mathbb{R}, \quad \Pi^N(\phi) = \Pi^N(\phi_s), \quad \phi_s(x, y) = \phi(x, y)/2 + \phi(y, x)/2, \quad (x, y) \in \square. \quad (2.34)$$

In other words, Π^N could really be defined on the triangle $\{x, y \in (-1, 1) : x < y\}$, but we work on the square for symmetry reasons. Note also that any symmetric function ϕ , i.e. $\phi = \phi_s$, that is C^1 on the whole of \square , satisfies:

$$\forall (x, y) \in \square, \quad \partial_1 \phi(x, y) = \partial_2 \phi(y, x) \quad \Rightarrow \quad \forall x \in (-1, 1), \quad (\partial_1 - \partial_2) \phi(x_{\pm}, x) = 0. \quad (2.35)$$

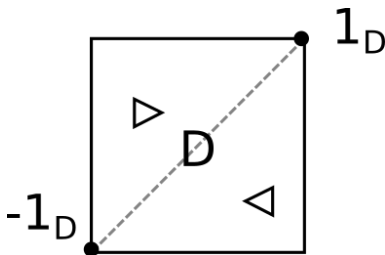


Figure 1: The domain $\square = (-1, 1)^2$ and the diagonal D (in dashed lines) are depicted, with $\square := \square \setminus D$. According to notation (2.37), the lower part is \triangleleft and the upper part \triangleright . The extremities of the diagonal 1_D and -1_D are defined in (2.43).

In view of (2.11), the two-point correlations are symmetric functions with singularities on the diagonal D of \square , defined by:

$$D := \{(x, x) : x \in (-1, 1)\}. \quad (2.36)$$

We therefore cannot only consider symmetric ϕ that are smooth on the whole of \square . Finally, to account for the reservoirs, we require ϕ to be continuous on $\partial\square$ as well, and set $\phi|_{\partial\square} = 0$.

Let us now define the test functions Π^N will act on. Split \square as follows (see Figure 1) :

$$\square = \triangleright \cup D \cup \triangleleft, \quad \triangleright := \{(x, y) \in \square : x < y\}, \quad \triangleleft := \{(x, y) \in \square : x > y\}, \quad \square := \triangleright \cup \triangleleft. \quad (2.37)$$

For $n \in \mathbb{N}$ and $p \geq 1$, let $\mathbb{W}^{n,p}(\square) := \mathbb{W}^{n,p}(\triangleright) \cap \mathbb{W}^{n,p}(\triangleleft)$ be the Sobolev space of functions with distributional derivatives up to order n in $\mathbb{L}^2(\square)$. Properties of these spaces are recalled in Appendix E. Note that $\mathbb{L}^p(\square) = \mathbb{L}^p(\square)$ since the diagonal has vanishing two-dimensional Lebesgue

measure. The difference between functions on \square and \boxminus arises in the integration by parts formula defining their weak derivatives. If $p = 2$, we simply write $\mathbb{H}^n(\triangleright) := \mathbb{W}^{n,2}(\triangleright)$. Define then the set \mathcal{T} of test functions:

$$\mathcal{T} := \mathbb{H}^2(\boxminus) = \mathbb{H}^2(\triangleright) \cap \mathbb{H}^2(\triangleleft). \quad (2.38)$$

The set \mathcal{T} is a separable Hilbert space, and $\mathcal{T} \subset C^0(\bar{\triangleright}) \cap C^0(\bar{\triangleleft})$ by Sobolev embedding, see Appendix E. Denote then by \mathcal{T}' the set of bounded linear forms on \mathcal{T} , and by $\mathcal{T}'_s \subset \mathcal{T}'$ the subset of those forms that are symmetric (recall (2.34)):

$$\mathcal{T}'_s = \{\Pi \in \mathcal{T}' : \forall \phi \in \mathcal{T}, \quad \Pi(\phi_s) = \Pi(\phi)\}. \quad (2.39)$$

To keep topology-related issues as simple as possible, we equip \mathcal{T}'_s with the weak* topology, i.e. $(\Pi_n)_n \subset \mathcal{T}'_s$ converges to $\Pi \in \mathcal{T}'_s$ if and only if $\lim_{n \rightarrow \infty} \Pi_n(\phi) = \Pi(\phi)$ for each $\phi \in \mathcal{T}$ (or, equivalently, each $\phi \in \mathcal{T}_\triangleright$). To avoid ambiguities, we write $(\mathcal{T}'_s, *)$ when we explicitly refer to the weak* topology.

As a bounded linear form on (a closed subset of) the Hilbert space $\mathbb{H}^2(\boxminus)$, the Riesz representation theorem allows each $\Pi \in \mathcal{T}'_s$ to be written as:

$$\Pi(\phi) := \frac{1}{4} \langle k_\Pi, \phi \rangle, \quad \phi \in \mathcal{T}. \quad (2.40)$$

Above, $\langle \cdot, \cdot \rangle$ denotes the standard scalar product on $\mathbb{L}^2(\boxminus)$ and duality pairing between elements of $\mathbb{H}^n(\boxminus)$ and $(\mathbb{H}^n(\boxminus))'$, $n \in \mathbb{N}$. The norm on $\mathbb{L}^2(\square) = \mathbb{L}^2(\boxminus)$ is denoted by $\|\cdot\|_2$. We use both Π and k_Π indifferently in the following.

2.4 Large deviations for time-averaged correlations

Our main result concerns the large deviation behaviour of the probability (2.33) in the large T , then large N limits. To state it, we need more notations.

Compared with the techniques used for the large deviations of the density (see e.g. Chapter 10 in [KL99]), producing atypical correlations requires modifying the jump rates of the dynamics by adding a long range interaction. Thus we consider the generator L_h parametrised by a function $h : [-1, 1]^2 \rightarrow \mathbb{R}$ which is a *non-local bias*, with the corresponding modified jump rates given by:

$$\begin{aligned} \forall \eta \in \Omega_N, \forall i \in \{\pm(N-1)\}, \quad c_h(\eta, i) &= c(\eta, i) \exp[\Pi^N(h)(\eta^i) - \Pi^N(h)(\eta)], \\ \forall j < N-1, \quad c_h(\eta, j, j+1) &= c(\eta, j, j+1) \exp[\Pi^N(h)(\eta^{j,j+1}) - \Pi^N(h)(\eta)]. \end{aligned} \quad (2.41)$$

We write $\mathbb{P}_h, \mathbb{E}_h$ for the probability/expectation under this dynamics, and $\mathbb{P}_h^{\mu^N}, \mathbb{E}_h^{\mu^N}$ when starting from the measure μ^N on Ω_N . The strategy is to find the correct bias h so that the rare event of observing the correlation Π in (2.33) becomes typical :

$$\mathbb{P}_h^{\mu^N} \left(\frac{1}{T} \int_0^T \Pi_t^N dt \stackrel{weak^*}{\approx} \Pi \right) = 1 \quad \text{when } T \text{ and then } N \text{ are large.} \quad (2.42)$$

Before stating our main Theorem 2.3, we have to introduce some restrictions, in particular on the size of the bias h .

Main assumption and characterisation of the biases. In theory, one could define the open SSEP dynamics with any value of the reservoir densities ρ_-, ρ_+ , and consider any sufficiently regular $h \in \mathcal{T}$ and any $\Pi \in \mathcal{T}'_s$ in (2.42). In practice, to focus on the key ideas and avoid many technical issues, we will restrict the range of ρ_-, ρ_+ as well as the size of the biases h , see Theorems 2.3–2.6 below. This restriction is discussed in Section 2.6. Let us now define the set of biases.

Let $\pm 1_{\triangleright}$ denote the two corners of the triangle \triangleright (see Figure 1) corresponding to extremities of the diagonal D , defined in (2.36):

$$1_D = (1, 1), \quad -1_D = (-1, -1). \quad (2.43)$$

For $\varepsilon > 0$, define:

$$\mathcal{S}(\varepsilon) = \left\{ h \in \mathcal{T} : h \in \mathbb{W}^{4,p}(\triangleright) \text{ for some } p > 2, h \text{ is symmetric,} \right. \\ \left. \|h\|_\infty, \|\partial_1 h\|_\infty \leq \varepsilon, \lim_{(x,y) \rightarrow \pm 1_D} (\partial_1 - \partial_2)h(x, y) = 0 \right\}. \quad (2.44)$$

The condition on $(\partial_1 - \partial_2)h$ at $\pm 1_D$ is purely technical. Introduce also the set $\mathcal{S}(\infty)$ of biases without size constraints:

$$\mathcal{S}(\infty) := \bigcup_{\varepsilon > 0} \mathcal{S}(\varepsilon). \quad (2.45)$$

The rate function. Introduce the bilinear mapping \mathcal{M} from $\mathbb{L}^2(\square)^2$ to $\mathbb{L}^2(\square)$, defined for $(u, v) \in \mathbb{L}^2(\square)^2$ by:

$$\forall (x, y) \in \square, \quad \mathcal{M}(u, v)(x, y) = \int_{(-1,1)} u(z, x) \bar{\sigma}(z) v(z, y) dz, \quad (2.46)$$

where $\bar{\sigma}$ is defined as

$$\forall x \in (-1, 1), \quad \bar{\sigma}(x) := \sigma(\bar{\rho}(x)) \quad \text{with} \quad \sigma(r) = r(1-r), \quad r \in [0, 1]. \quad (2.47)$$

For $\varepsilon > 0$, let h belong to the set $\mathcal{S}(\varepsilon)$ defined in (2.44). Introduce the functional J_h , defined for $\Pi \in \mathcal{T}'_s \cap \mathbb{H}^1(\square)$ by:

$$J_h(\Pi) = -\frac{1}{2} \Pi(\Delta h + \mathcal{M}(\partial_1 h, \partial_1 h)) + \frac{1}{4} \int_{(-1,1)} \text{tr}_D(k_\Pi)(x) (\partial_2 - \partial_1)h(x_+, x) dx \\ + \frac{(\bar{\rho}')^2}{4} \int_{(-1,1)} h(x, x) dx - \frac{1}{8} \int_{\square} \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 h(x, y)]^2 dx dy, \quad (2.48)$$

and:

$$J_h(\Pi) = +\infty \quad \text{if } \Pi \notin \mathcal{T}'_s \cap \mathbb{H}^1(\square). \quad (2.49)$$

In (2.48), $\text{tr}_D(k_\Pi)$ is the trace of k_Π on the diagonal D (defined in (2.36)), with k_Π related to Π via (2.40). It is well defined for Π in $\mathbb{H}^1(\square)$, see Theorem 1.5.1.3 in [Gri11]. Moreover, $\Pi \in \mathcal{T}'_s \cap \mathbb{H}^1(\square)$ implies that k_Π is symmetric, thus its trace on either side of the diagonal is the same and the notation $\text{tr}_D(k_\Pi)$ is not ambiguous.

Define then the functionals $\mathcal{I}_\varepsilon, \mathcal{I}_\infty : (\mathcal{T}'_s, *) \rightarrow \mathbb{R}_+$ ($\varepsilon > 0$) as follows:

$$\mathcal{I}_\varepsilon = \sup_{h \in \mathcal{S}(\varepsilon)} J_h, \quad \mathcal{I}_\infty = \sup_{h \in \mathcal{S}(\infty)} J_h, \quad (2.50)$$

where $\varepsilon > 0$ stands for the restriction on the size of the biases. To demystify the expression (2.48) of J_h , let $\Pi = \frac{1}{4}\langle k, \cdot \rangle \in \mathcal{T}'_s$ with $k \in C^3(\mathbb{T})$ a regular kernel. The associated correlation operator $C_k = \bar{\sigma} + k$ is defined for any test function $\phi \in \mathbb{L}^2((-1, 1)^2)$ by:

$$C_k \phi(x) = \bar{\sigma}(x)\phi(x) + \int_{(-1,1)} k(x, y)\phi(y)dy, \quad x \in [-1, 1], \quad (2.51)$$

where the diagonal part $\bar{\sigma}(x)$ corresponds to the variance at a single site, and k encodes the long range correlations. In comparison, in the steady state (corresponding to $h = 0$), the correlation operator in the large size limit is $C_{k_0} = \bar{\sigma} + k_0$, with the kernel k_0 introduced in (2.10).

The following proposition connects the bias h and the correlation kernel k . It is a classical result, proven in Appendix F.1.

Proposition 2.2 (Euler-Lagrange equation). *Assume that the supremum in the definition (2.50) is reached at some $h \in \mathcal{S}(\infty)$. Then k satisfies the Euler-Lagrange equation: for each $\phi \in \mathcal{T}$,*

$$\frac{1}{2} \int_{\mathbb{Z}} \nabla(k - k_0)(x, y) \cdot \nabla \phi(x, y) dx dy + \frac{1}{2} \int_{(-1,1)} \bar{\sigma}(z) \langle \partial_1 h(z, \cdot), C_k \partial_1 \phi(z, \cdot) \rangle dz = 0. \quad (2.52)$$

The expression (2.50) of the rate function \mathcal{I}_∞ then simplifies for $\Pi = \frac{1}{4}\langle k, \cdot \rangle$:

$$\mathcal{I}_\infty((1/4)\langle k, \cdot \rangle) = \frac{1}{8} \int_{(-1,1)} \bar{\sigma}(z) \langle \partial_1 h(z, \cdot), C_k \partial_1 h(z, \cdot) \rangle dz. \quad (2.53)$$

This generalises the asymptotics of the Donsker-Varadhan functional of Lemma 2.1 in the reversible case. In addition, the expression (2.53) has the familiar form of an \mathbb{L}^2 norm of the gradient of the bias in a suitable weighted space in terms of the target distribution (see [KL99, Chapter 10]).

The next theorem gives a large deviation result for the law of $\frac{1}{T} \int_0^T \Pi_t^N dt$. A more general claim is discussed in Section 2.6.

Theorem 2.3. *Let $\rho_- \in (0, 1)$. There is then $\varepsilon_B = \varepsilon_B(\rho_-)$, defined in Theorem 2.6 below, such that, if $\bar{\rho}' \leq \varepsilon_B$, then the following holds. Let $\mathcal{O}, \mathcal{C} \subset (\mathcal{T}'_s, *)$ respectively be an open, closed set. Then:*

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\pi_{inv}^N}(\mathcal{C}) \leq - \inf_{\mathcal{C}} \mathcal{I}_{\varepsilon_B}, \quad (2.54)$$

The rate function correctly captures the behaviour of the correlation field for kernels close to the kernel k_0 of the steady state, in the following sense. Let $\mathcal{C}_B \subset \mathcal{T}'_s$ be the set of correlation kernels associated with a regular bias $h \in \mathcal{S}(\varepsilon_B)$:

$$\mathcal{C}_B = \left\{ k \in \mathcal{T}'_s : k \text{ solves the Euler-Lagrange equation (2.52) for some } h \in \mathcal{S}(\varepsilon_B) \right\}. \quad (2.55)$$

Then, for $k \in \mathcal{C}_B$, one has $\mathcal{I}_\infty(\frac{1}{4}\langle k, \cdot \rangle) = \mathcal{I}_{\varepsilon_B}(\frac{1}{4}\langle k, \cdot \rangle)$, and:

$$\liminf_{N \rightarrow \infty} \liminf_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\pi_{inv}^N}(\mathcal{O}) \geq - \inf_{\mathcal{O} \cap \mathcal{C}_B} \mathcal{I}_{\varepsilon_B} = - \inf_{\mathcal{O} \cap \mathcal{C}_B} \mathcal{I}_\infty. \quad (2.56)$$

The statement of Theorem 2.3 and extensions (in particular a lower bound for non-regular kernels) are discussed in Section 2.6. The fact that, for $k \in \mathcal{C}_B$, the associated h is a global optimiser and thus $\mathcal{I}_\infty, \mathcal{I}_{\varepsilon_B}$ agree at $\frac{1}{4}\langle k, \cdot \rangle$ is proven in Section 5.2.

Remark 2.4. Recall that the kernel k_0 of the steady state π_{inv}^N of the open SSEP in the large N limit, defined in (2.11), was observed to be smooth away from the diagonal. As Theorem 2.3 shows, by definition of the rate function (see (2.49)), it is in fact a general property that the time-average of Π^N is much more regular than an element of \mathcal{T}'_s when T, N are large: it belongs to $\mathbb{H}^1(\square)$. ■

2.5 The relative entropy method

In this section, we explain the method used to establish Theorem 2.3, i.e. to study, for some $\Pi \in \mathcal{T}_s$, the probability:

$$\mathbb{P}^{\pi_{inv}^N} \left(\frac{1}{T} \int_0^T \Pi_t^N dt \approx \Pi \right) \quad \text{when } T, \text{ then } N \text{ are large.} \quad (2.57)$$

Understanding correlations out of equilibrium ($\rho_- \neq \rho_+$) is notoriously difficult. Existing results in the literature deal either with the equilibrium case $\rho_- = \rho_+$ without bias (i.e $h = 0$), see e.g. [GJ19]; or use methods particular to the $h = 0$ case, which cannot easily be generalised to $h \neq 0$ [Gon+20]. The methods rely on explicit knowledge of the invariant measure of the dynamics. However, for $h \neq 0$, the invariant measure $\pi_{inv,h}^N$ of the tilted dynamics \mathbb{P}_h defined in (2.41) is not known explicitly. Even for $h = 0$, where the invariant measure π_{inv}^N is well understood [DLS02], its complexity makes the study of the probability in (2.57) difficult. To study (2.57) in full generality, we therefore need a different approach.

The key idea is to find an approximation of the invariant measure, which is sufficiently close to control the large time behaviour of correlations, yet simple enough to make explicit computations possible. The proximity of the law of the dynamics to this approximate invariant measure is quantified through the relative entropy method, relying on the beautiful generalisation by Jara and Menezes [JM18] of the ideas of Yau [Yau91]. The relative entropy method is presented in more mathematical terms in Section 3.1, and here we only describe it informally. In our context, the relative entropy method consists in finding a measure μ^N on the state space Ω_N , that is both sufficiently simple to perform explicit computations, and as close as possible to the invariant measure $\pi_{inv,h}^N$. This closeness to the invariant measure aims at ensuring that, if the dynamics \mathbb{P}_h starts from μ^N then, at time $t \geq 0$, the law $f_t \mu^N$ of the dynamics is still close to μ^N . The proximity to the invariant measure is quantified by the relative entropy $H(f_t \mu^N | \mu^N)$. The level of precision needed on this relative entropy depends both on the quantity to study - e.g. the density, the density fluctuations, or in our case the correlations; and on the time range one wishes to probe.

Here, we improve on the estimates of Jara and Menezes [JM18], obtaining sufficiently precise relative entropy estimates to study the probability (2.57). To do so, for a bias h in the set $\mathcal{S}(\varepsilon_B)$, defined in (2.44), we compare the law of the dynamics at each time with a discrete Gaussian measure $\nu_{g_h}^N$ of the following form:

$$\forall \eta \in \Omega_N, \quad \nu_{g_h}^N(\eta) = (\mathcal{Z}_{g_h}^N)^{-1} \exp [2\Pi^N(g_h)] \bar{\nu}^N(\eta), \quad \text{with } \bar{\nu}^N(\eta) := \bigotimes_{i \in \Lambda_N} \text{Ber}(\bar{\rho}_i). \quad (2.58)$$

Above, for $\rho \in [0, 1]$, $\text{Ber}(\rho)$ is the Bernoulli measure on $\{0, 1\}$ with parameter ρ . The partition function $\mathcal{Z}_{g_h}^N$ is a normalisation factor. The function $g_h : \square \rightarrow \mathbb{R}$, which solves a partial differential equation depending on the dynamical bias h , is the function that allows us to minimise the entropy production $\partial_t H(f_t \nu_{g_h}^N | \nu_{g_h}^N)$, as stated in Theorem 2.6 below.

Let us first consider $h = 0$ and give a heuristic reason why the invariant measure π_{inv}^N can be approximated by a measure of the form (2.58) for a well chosen function $g_0 := g_{h=0}$. Under the product Bernoulli measure $\bar{\nu}^N$ (which has the same density profile as the steady state, but no correlations), for each test function $\phi \in C^0([-1, 1])$, the fluctuation field $Y^N(\phi) = N^{-1/2} \sum_i \bar{\eta}_i \phi_i$ converges to a Gaussian field with covariance $\bar{\sigma}$. Assuming for a moment that each $\bar{\eta}_i$ is a continuous variable, the law of fluctuations under ν_g^N is thus close to a Gaussian field with covariance $\bar{\sigma}^{-1} - g$. This claim is made rigorous in the forthcoming paper [Dag22, Appendix A]. Thus, in order for ν_g^N to have the same correlations k_0 (defined in (2.10)) as the steady state π_{inv}^N when N is large, one can take $g = g_0$ to be the inverse correlation kernel of the steady state:

$$C_{k_0} = \bar{\sigma} + k_0 =: (\bar{\sigma}^{-1} - g_0)^{-1} \quad \Leftrightarrow \quad g_0 := \bar{\sigma}^{-1} - (\bar{\sigma} + k_0)^{-1}. \quad (2.59)$$

Note that C_{k_0} is indeed invertible as, by construction, k_0 is a negative definite operator on $\mathbb{L}^2(\square)$:

$$\forall \phi \in \mathbb{L}^2((-1, 1)) \setminus \{0\}, \quad \int_{(-1, 1)^2} \phi(x) k_0(x, y) \phi(y) dx dy < 0. \quad (2.60)$$

The fact that g_0 is indeed a kernel operator follows from $g_0 C_{k_0} = \bar{\sigma}^{-1} k_0$, see the first theorem in section 4.6.1 in [BKM96]. This identity also implies that g_0 is a negative definite kernel. Moreover, as a function on \square , g_0 inherits the regularity of k_0 (smooth all the way up to the diagonal D , but with normal derivative having a jump across D). The following proposition and theorem provide a systematic way of approximating the invariant measure for a non-local bias $h \in \mathcal{S}(\varepsilon)$ for small enough $\varepsilon > 0$.

Proposition 2.5 (Main equation). *Let $0 < \rho_-, \rho_+ < 1$, take a bias $h \in \mathcal{S}(\infty)$ and consider the following problem with unknown g , referred to as the main equation:*

$$\left\{ \begin{array}{l} \Delta(g - h)(x, y) + \frac{\bar{\sigma}'(x)}{\bar{\sigma}(x)} \partial_1(2g - h)(x, y) + \frac{\bar{\sigma}'(y)}{\bar{\sigma}(y)} \partial_2(2g - h)(x, y) \quad \text{for } (x, y) \in \square \\ + \int_{(-1, 1)} \bar{\sigma}(z) [\partial_1 g(z, x) \partial_1(g - h)(z, y) + \partial_1 g(z, y) \partial_1(g - h)(z, x)] dz = 0, \\ g = 0 \quad \text{on } \partial \square, \\ (\partial_2 - \partial_1)(h - g)(x_+, x) = (\partial_1 - \partial_2)(h - g)(x_-, x) = \frac{(\bar{\rho}')^2}{\bar{\sigma}(x)^2} \quad \text{for } x \in (-1, 1). \end{array} \right. \quad (2.61)$$

There is $\varepsilon_0(\rho_-) > 0$ and a function $\delta(\cdot)$ with $\lim_{x \downarrow 0} \delta(x) = 0$ such that, for any $\varepsilon \in (0, \varepsilon_0(\rho_-)]$, $\bar{\rho}' \leq \varepsilon$ and $h \in \mathcal{S}(\varepsilon)$ implies that (2.61) has a unique solution $g_h \in g_0 + \mathcal{S}(\delta(\varepsilon))$.

Proposition 2.5 is proven in Appendix F. The next theorem provides the key control on the dynamics and determines the parameter ε_B mentioned in Theorem 2.3.

Theorem 2.6. *Let $0 < \rho_- \leq \rho_+ < 1$. There is $\varepsilon_B = \varepsilon_B(\rho_-)$ such that, if $\bar{\rho}' \leq \varepsilon_B$ and $h \in \mathcal{S}(\varepsilon_B)$, then the measure $\nu_{g_h}^N$ defined in (2.58) is a good approximation of the invariant measure of the*

dynamics \mathbb{P}_h with bias h in the following sense. Let $f_t \nu_{g_h}^N$ denote the law of \mathbb{P}_h at time $t \geq 0$. There is then $C, K > 0$ depending on h, ρ_{\pm} such that:

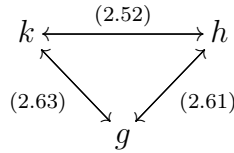
$$\forall t \geq 0, \quad H(f_t \nu_{g_h}^N | \nu_{g_h}^N) \leq e^{-Kt} H(f_0 \nu_{g_h}^N | \nu_{g_h}^N) + \frac{C}{N^{1/2}}. \quad (2.62)$$

Theorem 2.6 is proven in Section 3. The conditions that the parameter ε_B must satisfy are summarised in Definition F.10.

Remark 2.7. For each bias $h \in \mathcal{S}(\varepsilon_B)$, we prove in Appendix F.2 that obtaining a solution g_h to the main equation (2.61) with the desired regularity is equivalent to obtaining a classical solution k_h of the Euler-Lagrange (2.52) with the same regularity, with g_h and k_h related through the following identity of the associated operators:

$$\bar{\sigma} + k_h = (\bar{\sigma}^{-1} - g_h)^{-1}. \quad (2.63)$$

In particular this validates the heuristics behind the choice (2.59) of g_0 . The following diagram summarises the relationships between h , k and g .



■

2.6 Conclusion and perspectives

2.6.1 Extensions of the large deviation principle

The large deviation result of Theorem 2.3 is stated starting from the invariant measure π_{inv}^N . In fact any choice of initial condition is possible, with no change to the proof.

An advantage of our very precise, quantitative microscopic estimates is that Theorem 2.3 also holds if one takes a diverging sequence T_N of times.

Theorem 2.3 gives large deviation bounds in the weak-* topology. In fact, the theorem also holds in the strong dual topology. This can be seen to hold with no change to the proof for the upper bound. For the lower bound, one needs to be more careful, so we chose to work in a weaker topology to avoid technicalities. An extension of the lower bound to non-regular correlation fields is also possible, and sketched in Section 5.3.

2.6.2 Restriction on the biases and reservoir densities

In Theorem 2.3, restrictions are imposed on both the biases h and the difference $\rho_+ - \rho_-$ of the reservoir densities.

- Some restriction has to be imposed on the bias size h , otherwise one expects that not only the correlation structure, but also the typical density of the corresponding biased dynamics changes. Nevertheless, the conditions in (2.44) are far from sharp.

- The restriction on the slope $\bar{\rho}' = (\rho_+ - \rho_-)/2$ has two technical purposes. First, this is convenient to prove existence of g_h solving the main equation (2.61) (even though it should be possible to remove this assumption). Secondly, it is used in the derivation of the relative entropy bound of Theorem 2.6 to control the relative entropy uniformly in time: one has to make sure that error terms are bounded by cN uniformly on the state space for sufficiently small c in order to estimate their exponential moments, and c depends on $\bar{\rho}'$. Again, one should be able to relax this assumption, for instance through a priori estimate on the size of the error terms using large deviation results for the density. All other uses of the fact that $\bar{\rho}'$ is small can be relaxed without additional work, but at the cost of more technicalities.

Let us however stress again that the large deviation result of Theorem 2.3 is sharp for correlation kernels close to the kernel k_0 of the steady state (2.11), as in that case the rate functions $\mathcal{I}_{\varepsilon_B}, \mathcal{I}_\infty$ (recall (2.50)) coincide.

2.6.3 The relative entropy method

The relative entropy approach used to obtain large deviations can be applied to many other settings, still in dimension one. Let us list a few.

A first direction is the study of more general diffusive gradient systems with the following restrictions. For systems more complicated than the SSEP, the behaviour of correlations at the boundary may be problematic, and it is unclear whether one could still get a relative entropy estimate. However, results of the present paper should carry over to any diffusive gradient system on the torus (see the discussion in Section 8.1. in [JM18]). Additional work would be needed, in particular some large deviation estimates for the density to ensure that higher order correlations on the diagonal cannot be very large. The relative entropy bound would then only be of size $o_N(1)$, instead of $O(N^{-1/2})$ as in Theorem 2.6.

In a forthcoming paper [Dag22], we use the present refinement of the relative entropy method to study fluctuations in the SSEP on the one-dimensional discrete torus, constrained to produce a macroscopic current on the time interval $[0, T]$. The presence of this current is known to create a rich correlation structure. An expression of these correlations in the long time limit has been conjectured in [Bod+08]. In [Dag22], this conjecture is proved and the fluctuations of this process are described. In a related manner, Derrida and Sadhu [DS19] study the density large deviations for a non equilibrium SSEP conditioned to have an atypical macroscopic current.

Another direction of inquiry concerns the extension of the relative entropy results to higher dimensions. This seems to be a complicated problem. One major hurdle is the fact that, in higher dimensions, correlations are not smooth on the diagonal (see e.g. [Spo83]), while the relative entropy method requires smoothness.

3 Main ingredient: the entropic estimate

In this section, we provide the key microscopic estimates to study the long-time behaviour of the process $(\Pi_t^N)_{t \geq 0}$, i.e. we prove Theorem 2.6. We follow the strategy of Jara and Menezes [JM18]–[JM20] and improve the controls by tuning correlations of the reference measure. The same kind of

computations give the expression of the Radon-Nikodym derivative $D_h = d\mathbb{P}_h/d\mathbb{P}$ for $h \in \mathcal{S}(\varepsilon_B)$, stated in Proposition 3.18 at the end of the section.

3.1 The relative entropy method and Feynman-Kac inequality

Let us now recall the main features of the relative entropy method. Let $(\omega_t)_{t \geq 0}$ be a Markov chain on a state space Ω , assumed to be finite for simplicity. Let \mathbb{P}, \mathbb{E} denote the associated probability and expectation. Let $V : \Omega \rightarrow \mathbb{R}$. One would like to estimate quantities of the form:

$$\mathbb{E}[V(\omega_t)], \quad t > 0. \quad (3.1)$$

The entropy inequality provides a tool to estimate (3.1). Let μ be any probability measure on Ω and $f_t \mu$ be the law of ω_t , $t \geq 0$. Then, for any $\gamma > 0$:

$$\forall t \geq 0, \quad \mathbb{E}[V(\omega_t)] \leq \gamma^{-1} H(f_t \mu | \mu) + \gamma^{-1} \log \mu(\exp[\gamma V]). \quad (3.2)$$

Remark 3.1. *The symbol \mathbb{E} always denotes dynamical expectations. In contrast, static expectations with respect to a measure μ are denoted by $\mu[\cdot]$ (or $\mu(\cdot)$).* ■

Through (3.2), the dynamical estimate of V in (3.1) is reduced to a static problem: a relative entropy estimate, and a concentration-of-measure result under μ . The relative entropy method aims at finding a measure μ which satisfies the following two criteria: exponential moments must be under control, and the relative entropy $H(f_t \mu | \mu)$, $t \geq 0$ must be sufficiently small, the size depending on the kind of observables V one is interested in.

The analysis in [JM18]-[JM20] greatly improves the existing method to control $H(f_t \mu | \mu)$, $t \geq 0$. As a starting point, Jara and Menezes revisit Yau's entropy bounds in the following form.

Lemma 3.2 (Lemma A.1 in [JM18]). *Let $(\omega_t)_{t \geq 0}$ be a Markov chain on a finite state space Ω , with jump rates $(c(\omega, \omega'))_{(\omega, \omega') \in \Omega^2}$. Denote by L its generator and by Γ the corresponding carré du champ operator:*

$$\forall \omega \in \Omega, \forall f : \Omega \rightarrow \mathbb{R}, \quad \Gamma f(\omega) = \sum_{\omega' \in \Omega} c(\omega, \omega') [f(\omega') - f(\omega)]^2. \quad (3.3)$$

Let μ be a probability measure on Ω satisfying $\inf_{\omega \in \Omega} \mu(\omega) > 0$. Let $f_t \mu$ be the law of the process $(\omega_s)_{s \geq 0}$ at time $t \geq 0$. Then:

$$\forall t \geq 0, \quad \partial_t H(f_t \mu | \mu) \leq -\mu(\Gamma(\sqrt{f_t})) + \mu(f_t L^* \mathbf{1}), \quad (3.4)$$

where L^ is the adjoint of L in $\mathbb{L}^2(\mu) = \{f : \Omega \rightarrow \mathbb{R} : \mu(f^2) < \infty\}$. It acts on $f : \Omega \rightarrow \mathbb{R}$ according to:*

$$\forall \omega \in \Omega, \quad L^* f(\omega) = \sum_{\omega' \in \Omega} \left[c(\omega', \omega) f(\omega') \frac{\mu(\omega')}{\mu(\omega)} - c(\omega, \omega') f(\omega) \right]. \quad (3.5)$$

Since the adjoint L^* is known explicitly in terms of μ , (3.4) provides a way to estimate $\partial_t H(f_t \mu | \mu)$, $t \geq 0$. An estimate of $H(f_t \mu | \mu)$ follows by applying the entropy and Gronwall inequalities.

The same estimates used to bound the relative entropy will allow us, together with a log-Sobolev inequality (see Lemma 3.4), to get estimates on exponential moments. This is a consequence of a bound involving the Feynman-Kac formula, stated now for future reference.

Lemma 3.3 (Feynman-Kac inequality, Lemma A.2. in [JM18]). *For $V : \Omega \rightarrow \mathbb{R}$ and $T \geq 0$,*

$$\log \mathbb{E}^\mu \left[\exp \left(\int_0^T V(\omega_t) dt \right) \right] \leq T \sup_{f \geq 0: \mu(f)=1} \left\{ \mu(fV) - \frac{1}{2} \mu(\Gamma(\sqrt{f})) + \frac{1}{2} \mu(fL^* \mathbf{1}) \right\}. \quad (3.6)$$

The main difficulty to use (3.4)–(3.6) is to control the term $L^* \mathbf{1}$. Notice that $L^* \mathbf{1} = 0$ if and only if $\mu = \pi$ is the invariant measure. The quantity $L^* \mathbf{1}$ thus appears as a way to quantify the proximity of μ to the invariant measure π . This serves as an informal guiding principle for the choice of μ :

$$\text{for each } f \geq 0 \text{ with } \mu(f) = 1, \quad \mu(fL^* \mathbf{1}) \text{ must be small.} \quad (3.7)$$

We are going to apply Lemma 3.2 to the dynamics \mathbb{P}_h , defined in (2.41), for $h \in \mathcal{S}(\varepsilon_B)$. Let us emphasize again that "small" must always be understood in comparison with the size of the V 's one wishes to estimate as in (3.1). Typically, if one wants to study the hydrodynamic limit of the density of particles in a d -dimensional open SSEP on a lattice of side-length N , related observables are of the form $\sum_i \eta_i \phi_i \approx N^d$ for a test function ϕ . The rule of thumb is then that one needs $o(N^d)$ bounds on the relative entropy. These can be achieved if μ is a product measure with the same densities as those of the invariant measure at each site, i.e:

$$\mu = \bar{\nu}^N = \bigotimes_{i \in \Lambda_N} \text{Ber}(\bar{\rho}_i), \quad (3.8)$$

with $\text{Ber}(\rho)$ the Bernoulli measure on $\{0, 1\}$ with parameter $\rho \in (0, 1)$, and $\bar{\rho}$ the steady state density profile in the large N limit, see (2.5). In contrast, consider the fluctuation field:

$$\forall \phi : (-1, 1) \rightarrow \mathbb{R}, \quad Y_t^N(\phi) = \frac{1}{N^{1/2}} \sum_{i \in \Lambda_N} \bar{\eta}_i(t) \phi(i/N). \quad (3.9)$$

The typical observables to study Y^N should be of the form $N^{1/2} Y^N(\phi)$ which is typically of order $N^{1/2}$, so one needs $o(N^{1/2})$ bounds on the relative entropy (in fact $o(N^{d/2})$ bounds in dimension d). Remarkably, while one could expect that some information on the correlation structure of the invariant state should be necessary to study Y^N , Jara and Menezes [JM18] managed to obtain such bounds on the relative entropy by still taking μ product as in (3.8). To do so, they set up a general renormalisation scheme to bound $\mu(fL^* \mathbf{1})$, for a μ -density f , in terms of the carré du champ, and objects that can be estimated by the entropy inequality. Precisely, they manage to prove bounds of the form:

$$H(f_t \mu | \mu) \leq C(T) a_d(N) N^{d-2} \quad \text{with} \quad a_d(N) = \begin{cases} N & \text{if } d = 1 \\ \log N & \text{if } d = 2 \\ 1 & \text{if } d \geq 3. \end{cases} \quad (3.10)$$

This is enough to study fluctuations in dimension $d < 4$. They also argue that these bounds are the best possible when μ is product.

Let us come back to the study of the correlation process Π^N in the (one-dimensional) open SSEP. Observables, of the form $\Pi^N(\phi)$, should be bounded with N , so we need $o_N(1)$ bounds on

the relative entropy at each time. The measure μ therefore cannot be taken product: one needs to include information on the correlations under the invariant measure in μ . With (3.7) in mind, we look for μ that has both the same density at each site, and the same two-point correlations as the invariant measure - which are in general not known - when N is large. We tune these unknown correlations in an indirect way, taking a smooth function $g : \square \rightarrow \mathbb{R}$, and looking for the optimal choice of g such that the measure $\mu = \nu_g^N$ satisfies (3.7), with:

$$\nu_g^N := \frac{1}{\mathcal{Z}_g^N} \exp [2\Pi^N(g)] \bar{\nu}^N, \quad \mathcal{Z}_g^N \text{ a normalisation factor, and } \bar{\nu}^N \text{ as in (3.8),} \quad (3.11)$$

For each bias $h \in \mathcal{S}(\varepsilon)$ for sufficiently small $\varepsilon > 0$, the optimal $g = g_h$ arises as the solution of a certain partial differential equation, that we call the main equation (2.61). For this g_h , the method of [JM18], adapted to this context, and a logarithmic-Sobolev inequality yield the bound of Theorem 2.6:

$$\exists C, K > 0, \forall t \geq 0, \quad H(f_t \nu_{g_h}^N | \nu_{g_h}^N) \leq e^{-Kt} H(f_0 \nu_{g_h}^N | \nu_{g_h}^N) + CN^{-1/2}. \quad (3.12)$$

The exponent $-1/2$ improves on (3.10). It is related to the size of exponential moments of three-point and four-point correlation functions for product measures (see Section A.2). It thus cannot be improved without adding a correction to $\nu_{g_h}^N$. The proof of (3.12) is the main technical result of the article. The precise statement of the result is the content of the next two lemmas, a more comprehensive reformulation of Theorem 2.6. For $h \in \mathcal{S}(\infty)$ (defined in (2.45)) and $f : \Omega_N \rightarrow \mathbb{R}_+$, let $\Gamma_h(\sqrt{f})$ be the carré du champ operator associated with the generator L_h biased by h , with jump rates c_h defined in (2.41):

$$\begin{aligned} \forall \eta \in \Omega_N, \quad \Gamma_h(\sqrt{f})(\eta) &= \frac{1}{4} \sum_{i < N-1} c_h(\eta, i, i+1) [\sqrt{f}(\eta^{i,i+1}) - \sqrt{f}(\eta)]^2 \\ &+ \frac{1}{4} \sum_{i \in \{\pm(N-1)\}} c_h(\eta, i) [\sqrt{f}(\eta^i) - \sqrt{f}(\eta)]^2. \end{aligned} \quad (3.13)$$

Lemma 3.4 (Log-Sobolev inequality). *Let $0 < \rho_- \leq \rho_+ < 1$. There is a constant $C_{LS} = C_{LS}(\rho_\pm)$ such that, for each $\varepsilon, \varepsilon' \in (0, 1/4)$, each $h \in \mathcal{S}(\varepsilon)$ and each $g \in g_0 + \mathcal{S}(\varepsilon')$, the following inequality holds. For any density f for ν_g^N :*

$$\forall N \in \mathbb{N}^*, \quad H(f \nu_g^N | \nu_g^N) \leq C_{LS} N^2 \nu_g^N(\Gamma_h(\sqrt{f})). \quad (3.14)$$

The ε_B appearing in the next lemma is the same as the one of Theorems 2.3–2.6.

Lemma 3.5 (Approximation of the invariant measure). *Let $0 < \rho_- \leq \rho_+ < 1$. Let $h \in \mathcal{S}(\infty)$ and assume that the main equation (2.61) has a solution $g_h \in g_0 + \mathcal{S}(\infty)$ (the set $\mathcal{S}(\infty)$ is defined in (2.44)). For $N \in \mathbb{N}^*$, the reference measure $\nu_{g_h}^N$ is defined by (3.11).*

There is then $\varepsilon_B = \varepsilon_B(\rho_-) > 0$ such that, if $\bar{\rho}' := \frac{\rho_+ - \rho_-}{2} \leq \varepsilon_B$ and $h \in \mathcal{S}(\varepsilon_B)$, then g_h is a negative kernel (as defined in (2.60)) and the following is true.

• *For any $N \in \mathbb{N}^*$, there is a function $\mathcal{E} : \Omega_N \rightarrow \mathbb{R}$ such that, for any $\nu_{g_h}^N$ -density f , the adjoint L_h^* of L_h in $\mathbb{L}^2(\nu_{g_h}^N)$ satisfies:*

$$\nu_{g_h}^N(f N^2 L_h^* \mathbf{1}) \leq \nu_{g_h}^N(f \mathcal{E}) + \frac{N^2}{2} \nu_{g_h}^N(\Gamma_h(\sqrt{f})). \quad (3.15)$$

- There are constants $\gamma > 8C_{LS}, C > 0$ depending on h and the reservoir densities ρ_{\pm} , such that:

$$\forall N \in \mathbb{N}^*, \quad \gamma^{-1} \log \nu_{g_h}^N \left(\exp [\gamma |\mathcal{E}|] \right) \leq \frac{C}{N^{1/2}}. \quad (3.16)$$

- As a consequence, for any density f for $\nu_{g_h}^N$,

$$\begin{aligned} \nu_{g_h}^N(f L_h^* \mathbf{1}) - N^2 \nu_{g_h}^N(\Gamma_h(\sqrt{f})) &\leq \frac{H(f \nu_{g_h}^N | \nu_{g_h}^N)}{8C_{LS}} + \frac{C}{N^{1/2}} - \frac{N^2}{2} \nu_{g_h}^N(\Gamma_h(\sqrt{f})) \\ &\leq -\frac{H(f \nu_{g_h}^N | \nu_{g_h}^N)}{8C_{LS}} + \frac{C}{N^{1/2}} - \frac{N^2}{4} \nu_{g_h}^N(\Gamma_h(\sqrt{f})), \end{aligned} \quad (3.17)$$

and if $f_t \nu_{g_h}^N$ denotes the law of the dynamics at time t , then:

$$\forall t \geq 0, \forall N \in \mathbb{N}^*, \quad \partial_t H(f_t \nu_{g_h}^N | \nu_{g_h}^N) \leq -\frac{H(f_t \nu_{g_h}^N | \nu_{g_h}^N)}{8C_{LS}} + \frac{C}{N^{1/2}}. \quad (3.18)$$

Let us now prove Theorem 2.6 using Lemmas 3.4–3.5. Let g_h solve the main equation (2.61). Recalling (3.18) and applying Gronwall inequality to $t \mapsto H(f_t \nu_{g_h}^N | \nu_{g_h}^N)$ yields Theorem 2.6:

$$\forall t \geq 0, \quad H(f_t \nu_{g_h}^N | \nu_{g_h}^N) \leq H(f_0 \nu_{g_h}^N | \nu_{g_h}^N) e^{-(8C_{LS})^{-1}t} + \frac{8C_{LS}C}{N^{1/2}} \left(1 - e^{-(8C_{LS})^{-1}t} \right). \quad (3.19)$$

Lemma 3.4 is proven in Appendix A.3. The proof of the key ingredient, Lemma 3.5, takes up the next four subsections. The fact that g_h is a negative kernel if $h \in \mathcal{S}(\varepsilon)$ and $\bar{\rho}' \leq \varepsilon$ for small enough $\varepsilon > 0$ is a consequence of Proposition F.5, where it is shown that $\|g_h - g_0\|_2$ vanishes with ε .

A function $h \in \mathcal{S}(\infty)$ and a negative kernel $g \in g_0 + \mathcal{S}(\infty)$ are fixed throughout the rest of the section, and we highlight where we need to restrict the size of h and $\bar{\rho}'$. Recall that h, g are symmetric functions, and that their restrictions to \triangleright are in $C^3(\bar{\triangleright})$. We do not a priori assume that g solves the main equation (2.61), and explain along the proof where this comes into play.

3.2 Estimates on $L_h^* \mathbf{1}$

To prove Lemma 3.5, we need to compute $L_h^* \mathbf{1}$. The computation of $L_h^* \mathbf{1}$ will give rise to many different objects which, roughly speaking, will either contribute to leading order in N , or be sub-leading order error terms. In this section, we define precisely how to estimate the size of a function in terms of N , and formulate criteria to identify which terms are error terms.

3.2.1 Size of error terms

Consider a density f for ν_g^N , and a function $X_N : \Omega_N \rightarrow \mathbb{R}$. Our main tool to estimate the scaling of X_N with N is the entropy inequality:

$$\forall \gamma > 0, \quad \nu_g^N(f X_N) \leq \frac{H(f \nu_g^N | \nu_g^N)}{\gamma} + \frac{1}{\gamma} \log \nu_g^N \left(\exp [\gamma |X_N|] \right). \quad (3.20)$$

Informally, we will say that X_N is small if its moment generating function under ν_g^N vanishes with N for γ in a neighbourhood of 0. This is the kind of characterisation of smallness that is used to

estimate the size of the function \mathcal{E} in Lemma 3.5. In some cases, typically when dealing with the effect of the reservoirs, we will encounter an X_N that is not small, but can be transformed into some \tilde{X}_N that is indeed small, up to a cost estimated by the carré du champ operator. The next definition formalises these considerations, and examples are given in Lemma 3.9 below.

Definition 3.6. Let $a_N \in \mathbb{R}_+^*$, $N \in \mathbb{N}^*$. A family $X_N : \Omega_N \rightarrow \mathbb{R}$, $N \in \mathbb{N}^*$ of functions is said to be:

- Controllable with size a_N if there are γ, K independent of N such that:

$$\forall N \in \mathbb{N}^*, \quad \frac{1}{\gamma} \log \nu_g^N \left(\exp [\gamma |X_N|] \right) \leq K a_N. \quad (3.21)$$

By convention, if $a_N(p)$ depends on an additional parameter p , then the constant K in (3.21) will not depend on p . By the entropy inequality (3.20), (3.21) implies, for each density f for ν_g^N :

$$\forall N \in \mathbb{N}^*, \quad \nu_g^N(f|X_N) \leq \frac{H(f\nu_g^N|\nu_g^N)}{\gamma} + K a_N. \quad (3.22)$$

- Γ -controllable with size a_N if one can transform X_N , using the carré du champ Γ_h , into a controllable function with size a_N . More precisely: X_N is Γ -controllable with size a_N if there are controllable functions \tilde{X}_\pm^N with size a_N such that, for each $\delta > 0$, each $N \in \mathbb{N}^*$ and each density f for ν_g^N :

$$\nu_g^N(f(\pm X_N)) \leq \delta N^2 \nu_g^N(\Gamma_h(f^{1/2})) + \frac{1}{\delta} \nu_g^N(f \tilde{X}_\pm^N). \quad (3.23)$$

The entropy inequality (3.20) then again implies that there are $\gamma > 0, K > 0$ independent of δ, N such that:

$$\nu_g^N(f(\pm X_N)) \leq \delta N^2 \nu_g^N(\Gamma_h(f^{1/2})) + \frac{H(f\nu_g^N|\nu_g^N)}{\gamma\delta} + \frac{K a_N}{\gamma\delta}. \quad (3.24)$$

- An error term with size a_N , or error term for short, if it is either controllable or Γ -controllable with size a_N , and $a_N = o_N(1)$.

Remark 3.7. (Γ -)controllability behaves well with respect to multiplication by a small constant in the following sense. Assume that X_N is (Γ -)controllable with size a_N and let $b_N \in [0, 1]$ ($N \in \mathbb{N}^*$) satisfy $b_N = o_N(1)$. Then $b_N X_N$ is (Γ -)controllable with size $b_N a_N$ by Jensen inequality. ■

To illustrate the notion of controllability, the following proposition, proven in Appendix C.2, states its consequence on the dynamical behaviour of observables.

Proposition 3.8. Let $h \in \mathcal{S}(\varepsilon_B)$ with ε_B given by Lemma 3.5, and let g_h be the associated solution of the main equation (2.61) as in Lemma 3.5. Let $E^N : \Omega_N \rightarrow \mathbb{R}$ be an error term with size $a_N = o_N(1)$, and let F^N be (Γ -)controllable with size 1. There are then γ, C and $\gamma', C' > 0$ independent of N, T such that:

$$\begin{aligned} \forall T > 0, \quad \frac{1}{T} \log \mathbb{E}^{\nu_{g_h}^N} \left[\exp \left| \gamma \int_0^T E^N(\eta_t) dt \right| \right] &\leq C a_N, \\ \sup_N \frac{1}{T} \log \mathbb{E}^{\nu_{g_h}^N} \left[\exp \left| \gamma' \int_0^T F^N(\eta_t) dt \right| \right] &\leq C'. \end{aligned} \quad (3.25)$$

Let $\phi : \Lambda_N \rightarrow \mathbb{R}$. To determine whether a field is an error term or not, one must keep in mind the following heuristics: the measures ν_g^N are discrete Gaussian measures, in the sense that the fluctuation field $Y^N(\phi)$, which reads:

$$Y^N(\phi) := \frac{1}{N^{1/2}} \sum_{i \in \Lambda_N} \bar{\eta}_i \phi(i), \quad (3.26)$$

is close to a Gaussian random variable when N is large, provided $\|\phi\|_\infty < \infty$. In particular, $\lambda \mapsto \nu_g^N(\exp[\lambda Y^N(\phi)])$ is bounded uniformly in N in a neighbourhood $[0, \gamma(\phi))$ of 0 for some $\gamma(\phi) > 0$. By (3.21), this means that $Y^N(\phi)$ is controllable with size 1 (or size $C(\phi)$ for some constant $C(\phi) > 0$ if we want to keep track of the dependence on ϕ). In analogy with Gaussian random variables, one can prove that $Y^N(\phi)^2$ is controllable with size 1, but $Y^N(\phi)^n$ for $n \geq 3$ is not. Similarly, the quantity:

$$Z^N(\phi) := \frac{1}{N^{1/2}} \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} \phi(i) \quad (3.27)$$

should not have worse concentration properties than $Y^N(\phi)$. Contrary to genuine Gaussian random variable, however, $Y^N(\phi)$ and $Z^N(\phi)$ are bounded, by $C\|\phi\|_\infty N^{1/2}$ for some $C > 0$. As a result, it is always possible to find a_N small enough such that $a_N(Y^N(\phi))^n$ is controllable with size 1. This discussion is summarised in the next lemma.

Lemma 3.9. *For $n \in \mathbb{N}^*$, let $\phi_n : \Lambda_N^n \rightarrow \mathbb{R}$ satisfy $\sup_N \|\phi_n\|_\infty < \infty$. Define, for $N \in \mathbb{N}^*$ and either $J = \{0\}$ or $J = \{0, 1\}$, the functions (an empty product is by convention equal to 1):*

$$\begin{aligned} \forall \eta \in \Omega_N, \quad X_{n,J}^{\phi_n}(\eta) &= \sum_{i_0 < N-|J|} \sum_{i_1, \dots, i_{n-1} \in \Lambda_N} \phi_n(i_0, \dots, i_{n-1}) \left(\prod_{j \in J} \bar{\eta}_{i_0+j} \right) \left(\prod_{a=1}^{n-1} \bar{\eta}_{i_a} \right), \\ U_0^\varepsilon(\eta) &= \bar{\eta}_{\varepsilon(N-1)}, \quad U_1^\varepsilon(\eta) = \bar{\eta}_{\varepsilon(N-1)} \sum_{i \neq \varepsilon(N-1)} \bar{\eta}_i \phi_1(i), \quad \varepsilon \in \{-, +\}. \end{aligned} \quad (3.28)$$

Then:

- the function U_0^\pm is Γ -controllable with size N^{-1} , and U_1^\pm is Γ -controllable with size $\|\phi_1\|_\infty N^{-1}$. Moreover, $N^{-1/2} X_{1,J}^{\phi_1}$ is controllable with size $\|\phi_1\|_\infty^2$.
- For $n \geq 2$, there are constants $\gamma_n > 0$ that depend only on n , but not on ϕ_n , such that:

$$\log \nu_g^N \left(\exp \left[\frac{\gamma_n}{\|\phi_n\|_\infty} N^{-(n-1)} X_{n,J}^{\phi_n} \right] \right) \leq \frac{3}{N^{\frac{n-2}{2}}}. \quad (3.29)$$

This implies (but is stronger) that $N^{-(n-1)} X_{n,J}^{\phi_n}$ is controllable with size $\|\phi_n\|_\infty N^{-(n-2)/2}$. In addition, if $\phi_n(i_0, \dots, i_{n-1})$ vanishes whenever an index appears twice in the collection $i_0 + J, i_1, \dots, i_{n-1}$, then one can replace $\|\phi_n\|_\infty$ by the weaker norm $\|\phi_n\|_{2,N}$ in (3.29), where:

$$\|\phi_n\|_{2,N} := \left(\frac{1}{N^n} \sum_{i_0, \dots, i_{n-1}} \phi_n(i_0, \dots, i_{n-1})^2 \right)^{1/2}. \quad (3.30)$$

- For $n = 2$ and $|J| > 1$, say $J = \{0, 1\}$, the previous estimate can be improved: $N^{-1}X_{2,\{0,1\}}^{\phi_2}$ is Γ -controllable with size $\|\phi_n\|_\infty N^{-1/2}$ (or $\|\phi_n\|_{2,N} N^{-1/2}$ if $\phi_2(i, i) = 0 = \phi_2(i, i+1)$ for each i).

Remark 3.10. To help clarify the definition of $X_{n,J}^{\phi_n}$, take $J = \{0, 1\}$ and, for $n \in \mathbb{N}^*$, let $\phi_n(i_0, \dots, i_{n-1}) = \prod_{\ell=0}^{n-1} \phi_1(i_\ell)$. Then:

$$X_{n,J}^{\phi_n} = N^{n/2} (Y^N(\phi_1))^{n-1} Z^N(\phi_1), \quad (3.31)$$

with $Y^N(\phi)$, $Z^N(\phi)$ defined in (3.26)-(3.27). Moreover, for any $\phi_2 : \Lambda_N^2 \rightarrow \mathbb{R}$ with $\phi_2(i, i) = 0$ for $i \in \Lambda_N$, then the correlation field (2.7) can be recovered: $X_{2,\{0\}}^{\phi_2} = 4N\Pi^N(\phi_2)$. \blacksquare

Lemma 3.9 is proven in Appendix A.2 for the variables $X_{n,J}^{\phi_n}$. The statement for U_0^\pm, U_1^\pm is proven in Appendix B.2, and the last item corresponds to Proposition C.1.

3.2.2 Bounding the entropy

After classifying each term arising in the computation of $L_h^* \mathbf{1}$ according to the categories in Definition 3.6, we will obtain the estimate (3.15): for a good choice g_h depending on h and any density f for $\nu_{g_h}^N$,

$$\nu_{g_h}^N(f L_h^* \mathbf{1}) \leq \nu_{g_h}^N(f \mathcal{E}) + \frac{N^2}{2} \nu_{g_h}^N(\Gamma_h(\sqrt{f})). \quad (3.32)$$

Above, \mathcal{E} will be a controllable error term with size $N^{-1/2}$: there is $\gamma, C > 0$ depending on ρ_\pm, h such that

$$\forall N \in \mathbb{N}^*, \quad \log \nu_{g_h}^N \left(\exp[\gamma |\mathcal{E}|] \right) \leq \frac{C}{N^{1/2}}. \quad (3.33)$$

To turn this estimate on the adjoint into the entropy estimate of Theorem 2.6, there is one more step, which involves the log-Sobolev inequality of Lemma 3.4, see (3.19). To apply the log-Sobolev inequality, we need to ensure that the constant γ in (3.33) is sufficiently large compared to the log-Sobolev constant C_{LS} , in the sense that we want to have:

$$\begin{aligned} \nu_{g_h}^N(f L_h^* \mathbf{1}) - N^2 \nu_{g_h}^N(\Gamma_h(f^{1/2})) &\leq \frac{H(f \nu_{g_h}^N | \nu_{g_h}^N)}{\gamma} + \frac{1}{\gamma} \log \nu_{g_h}^N \left(\exp[\gamma |\mathcal{E}|] \right) - \frac{N^2}{2} \nu_{g_h}^N(\Gamma_h(f^{1/2})) \\ &\leq -\frac{H(f \nu_{g_h}^N | \nu_{g_h}^N)}{8C_{LS}} + \frac{C}{\gamma N^{1/2}} - \frac{N^2}{4} \nu_{g_h}^N(\Gamma_h(f^{1/2})), \end{aligned} \quad (3.34)$$

where $(8C_{LS})^{-1}$ in the last line could be replaced by any smaller positive number. The last bound requires $\gamma > 8C_{LS}$, and is the reason for the introduction of the following terminology.

Definition 3.11. • A (Γ) -controllable error term X_N of size a_N is said to be of vanishing type if there is a sequence $(\gamma_N)_N$ of positive numbers such that $\lim_N \gamma_N = \infty$, and $\gamma_N X_N$ is (Γ) -controllable with size 1.

- A function X_N (error term or not) is said to be of large type if, for some $\gamma > 1$ independent of N , γX_N has unbounded exponential moment as $N \rightarrow \infty$.

- A controllable function is said to be of LS type if one can take $\gamma > 2^{10}C_{LS}$ in the definition (3.21) of controllability (this is in particular always true for N large in the case of error terms of vanishing type). A Γ -controllable function is said to be of LS type if the controllable functions Y_{\pm}^N associated via (3.23) are of LS type.

Remark 3.12. The constant 2^{10} is not at all optimal (any large enough constant would also work), but is sometimes convenient later on. ■

Let us clarify this notion by classifying the functions appearing in Lemma 3.9.

Lemma 3.13. For $n \in \mathbb{N}$, $J = \{0\}$ or $\{0, 1\}$ and $\phi_n : \Lambda_N \rightarrow \mathbb{R}$ with $\sup_N \|\phi_n\|_{\infty}$, recall the definitions of $X_{n,J}^{\phi_n}, U_n^{\pm}$ in Lemma 3.9.

- $N^{1/2}U_0^{\pm}, N^{1/2}U_1^{\pm}$ are Γ -controllable with size 1, so both U_0^{\pm} and U_{\pm}^1 are error terms of vanishing type. Similarly, for any sequence $\varepsilon_N > 0$ ($N \in \mathbb{N}^*$) with $\varepsilon_N = o_N(1)$, $\varepsilon_N N^{-1/2}X_{1,J}^{\phi_1}$ and $\varepsilon_N N^{-(n-1)}X_{n,J}^{\phi_n}$ for $n \geq 2$ are error terms of vanishing type.
- For $n \geq 2$, $\gamma N^{-(n-1)}X_{n,J}^{\phi_n}$ has exponential moment bounded with N only if $\gamma > 0$ is small enough. Thus $N^{-(n-1)}X_{n,J}^{\phi_n}$ is of large type.
- For $n \geq 2$, there is a numerical constant $\zeta_n > 0$ such that $\frac{\zeta_n}{\|\phi_n\|_{\infty}}N^{-(n-1)}X_{n,J}^{\phi_n}$ is of LS type. If in addition $\phi_n(i_0, \dots, i_{n-1})$ vanishes whenever the same index appears twice in the collection $i_0 + J, i_1, \dots, i_{n-1}$, then the same statement is true replacing $\|\phi_n\|_{\infty}$ with $\|\phi_n\|_{2,N}$ (defined in (3.30)).
- If a function X_N is of LS type and $\varepsilon \in (0, 1)$, then εX_N is of LS type. Moreover, $(1 + \alpha)X_N$ is of LS type for sufficiently small $\alpha \in (0, 1)$ independent of N .

Lemma 3.13 is obtained as a consequence of the proof of Lemma 3.9, carried out in Corollary A.4.

To see why the claim of Lemma 3.13 is reasonable, consider again $Y^N(\phi_1) = N^{-1/2} \sum_i \bar{\eta}_i \phi_1(i)$ (ϕ_1 is bounded) as an approximately Gaussian random variable when N is large. Similarly, one should see $N^{-(n-1)}X_{n,J}^{\phi_n}$ as approximately $N^{-(n-2)/2}$ times a product of n Gaussian random variable ($n \geq 2$) for bounded ϕ_n . For $n \geq 3$, the prefactor $N^{-(n-2)/2}$ vanishes, thus $N^{-(n-1)}X_{n,J}^{\phi_n}$ should become an error term. Error terms, however, are defined in terms of smallness of exponential moments, and the moment generating function $M_n(\gamma)$ of a product of n Gaussian random variables is unbounded for γ large enough ($n = 2$) or for any $\gamma > 0$ if $n \geq 3$. For $n \geq 3$, the vanishing prefactor $N^{-(n-2)/2}$ when considering $N^{-(n-1)}X_{n,J}^{\phi_n}$ as a product of n Gaussians is precisely right to ensure that its exponential moments are bounded with N in a neighbourhood of 0.

For an error term of LS type, a straightforward adaptation of the proof of Proposition 3.8 (see Section C.2) yields the following result.

Corollary 3.14. Let E^N be a (Γ -)controllable function with size $a_N = O_N(1)$, and assume that E^N is of LS type. Then one can take $\gamma = 1$ in Proposition 3.8, i.e. there is $C = C(\rho_{\pm}) > 0$ such that:

$$\forall N \in \mathbb{N}^*, \forall T > 0 \quad \frac{1}{T} \log \mathbb{E}_{g_h}^{\nu_N} \left[\exp \left| \int_0^T E^N(\eta_t) dt \right| \right] \leq C a_N. \quad (3.35)$$

In particular, this applies to $E^N = \zeta_2 \|\phi_2\|_{2,N}^{-1} \Pi^N(\phi)$ with $a_N = 1$ for any bounded $\phi : \square \rightarrow \mathbb{R}$, with ζ_2 the constant in item 3 of Lemma 3.13.

Equipped with these notations, we now turn to computing $L_h^* \mathbf{1}$, in the next two subsections. We split L_h^* into $L_h^* = L_{h,0}^* + L_{h,\pm}^*$, the adjoint dynamics respectively in the bulk and at the boundaries, and study each contribution separately.

3.3 Adjoint at the boundary

In this section, we compute $L_{h,\pm}^* \mathbf{1}$, the part of the adjoint $L_h^* \mathbf{1}$ of L_h with respect to ν_g^N corresponding to the dynamics at the boundary. By (3.5), it reads:

$$\begin{aligned} N^2 L_{h,\pm}^* \mathbf{1}(\eta) &= \frac{N^2}{2} \sum_{i \in \{\pm(N-1)\}} \left[c_h(\eta^i, i) \frac{\nu_g^N(\eta^i)}{\nu_g^N(\eta)} - c_h(\eta, i) \right] \\ &= \frac{N^2}{2} \sum_{i \in \{\pm(N-1)\}} \left[c(\eta^i, i) \left(\frac{\bar{\rho}_i}{1 - \bar{\rho}_i} \right)^{1-2\eta_i} \exp \left[\frac{(1-2\eta_i)}{2N} \sum_{j \neq i} \bar{\eta}_j (2g - h)_{i,j} \right] \right. \\ &\quad \left. - c(\eta, i) \exp \left[\frac{(1-2\eta_i)}{2N} \sum_{j \neq i} \bar{\eta}_j h_{i,j} \right] \right]. \end{aligned} \quad (3.36)$$

The jump rates $c(\eta, i)$, $i \in \{\pm(N-1)\}$ are defined in (2.3). To compute (3.36), recall that both h, g satisfy $h(\pm 1, \cdot) = 0 = g(\pm 1, \cdot)$ by hypothesis. It follows that the arguments of the exponentials in (3.36) are bounded by $O(N^{-1})$. Moreover, introduce:

$$\forall i \in \Lambda_N = \{-N+1, \dots, N-1\}, \quad \lambda_i := \log \left(\frac{\bar{\rho}_i}{1 - \bar{\rho}_i} \right) \quad \text{as well as} \quad \lambda_{\pm N} = \log \left(\frac{\rho_{\pm}}{1 - \rho_{\pm}} \right). \quad (3.37)$$

With this notation and by reversibility, one has for $i \in \{\pm(N-1)\}$:

$$c(\eta^i, i) = c(\eta, i) \left(\frac{\rho_{\text{sign}(i)}}{1 - \rho_{\text{sign}(i)}} \right)^{2\eta_i - 1} = \exp \left[\frac{(2\eta_i - 1)}{N} \lambda_{\text{sign}(i)N} \right]. \quad (3.38)$$

As a result, with the notation $\partial^N \lambda_\ell = N(\lambda_{\ell+1} - \lambda_\ell)$, we get :

$$c(\eta^i, i) \left(\frac{\bar{\rho}_i}{1 - \bar{\rho}_i} \right)^{1-2\eta_i} = c(\eta, i) \exp \left[\frac{(1-2\eta_i)}{N} (\mathbf{1}_{i=-(N-1)} \partial^N \lambda_{-N} - \mathbf{1}_{i=N-1} \partial^N \lambda_{N-1}) \right]. \quad (3.39)$$

Write for short the boundary term as:

$$\mathcal{B}_i := \mathbf{1}_{i=-(N-1)} \partial^N \lambda_{-N} - \mathbf{1}_{i=N-1} \partial^N \lambda_{N-1}. \quad (3.40)$$

Using (3.39) and the existence of $C_{h,g,\bar{\rho}} = C(\|h\|_\infty, \|g\|_\infty, \|\bar{\rho}\|_\infty) > 0$ such that

$$|e^x - 1 - x^2/2| \leq C_{h,g,\bar{\rho}} |x|^3 \quad \text{for} \quad |x| \leq 2(\|h\|_\infty + 2\|g\|_\infty + \|\lambda\|_\infty), \quad (3.41)$$

Equation (3.36) can be expanded with an error term:

$$\begin{aligned} \delta_{\pm}^{N,1}(\eta) &:= N^2 L_{h,\pm}^* \mathbf{1}(\eta) - \frac{N}{2} \sum_{i \in \{\pm(N-1)\}} c(\eta, i) (1-2\eta_i) \left[\frac{1}{N} \sum_{j \neq i} \bar{\eta}_j N (g - h)_{i,j} + \mathcal{B}_i \right] \\ &\quad - \frac{1}{4} \sum_{i \in \{\pm(N-1)\}} c(\eta, i) \left[\left(\frac{1}{2N} \sum_{j \neq i} \bar{\eta}_j N (2g - h)_{i,j} + \mathcal{B}_i \right)^2 - \left(\frac{1}{2N} \sum_{j \neq i} \bar{\eta}_j N h_{i,j} \right)^2 \right], \end{aligned} \quad (3.42)$$

and $\delta_{\pm}^{N,1}$ satisfies $|\delta_{\pm}^{N,1}| \leq C_{h,g,\bar{\rho}}/N$. To compute the two terms in the left-hand side of (3.42), let $i \in \{\pm(N-1)\}$ and let us first rewrite the jump rate in terms of $\bar{\eta}_i$ in two different ways. One has:

$$c(\eta, i) := (1 - \eta_i)\rho_{\text{sign}(i)} + (1 - \rho_{\text{sign}(i)})\eta_i = a_{\text{sign}(i)} + \bar{\eta}_i(1 - 2\rho_{\text{sign}(i)}), \quad (3.43)$$

with:

$$a_{\text{sign}(i)} := (1 - \bar{\rho}_i)\rho_{\text{sign}(i)} + (1 - \rho_{\text{sign}(i)})\bar{\rho}_i. \quad (3.44)$$

Moreover, it also holds that:

$$\begin{aligned} (1 - 2\eta_i)c(\eta, i) &= -(\eta_i - \rho_{\text{sign}(i)}) = -\bar{\eta}_i + (\rho_{\text{sign}(i)} - \bar{\rho}_i) \\ &= -\bar{\eta}_i - \frac{1}{N} [\mathbf{1}_{i=-(N-1)}\partial^N \bar{\rho}_{-N} - \mathbf{1}_{i=N-1}\partial^N \bar{\rho}_{N-1}]. \end{aligned} \quad (3.45)$$

Using (3.45) in the first sum in (3.42), one finds:

$$\begin{aligned} \frac{N}{2} \sum_{i \in \{\pm(N-1)\}} c(\eta, i)(1 - 2\eta_i) &\left[\frac{1}{N} \sum_{j \neq i} \bar{\eta}_j N(g - h)_{i,j} + (\mathbf{1}_{i=-(N-1)}\partial^N \lambda_{-N} - \mathbf{1}_{i=N-1}\partial^N \lambda_{N-1}) \right] \\ &= -\frac{N\partial^N \lambda_{-N}}{2} \left[\bar{\eta}_{-(N-1)} + \frac{\partial^N \bar{\rho}_{-N}}{N} \right] + \frac{N\partial^N \lambda_{N-1}}{2} \left[\bar{\eta}_{N-1} - \frac{\partial^N \bar{\rho}_{N-1}}{N} \right] + \delta_{\pm}^{N,2}(\eta), \end{aligned} \quad (3.46)$$

where $\delta_{\pm}^{N,2}$ is an error term that reads:

$$\delta_{\pm}^{N,2}(\eta) = -\frac{1}{2} \sum_{i \in \{\pm(N-1)\}} \left(\bar{\eta}_i + \frac{1}{N} [\mathbf{1}_{i=-(N-1)}\partial^N \bar{\rho}_{-N} - \mathbf{1}_{i=N-1}\partial^N \bar{\rho}_{N-1}] \right) \sum_{j \neq i} \bar{\eta}_j N(g - h)_{i,j}. \quad (3.47)$$

Since $N(g - h)_{i,j}$ is of order 1 for i close to the boundary, the term involving $\bar{\eta}_i$ above is of the same form as U_1^{\pm} in Lemma 3.9, and recall that U_1^{\pm} is Γ -controllable with size N^{-1} and of vanishing type (as defined in Lemma 3.13). The other term is of the form $N^{-1/2}Y^N(\phi)$ for a bounded ϕ (recall (3.26)), and $Y^N(\phi)$ is controllable with size 1, thus $N^{-1/2}Y^N(\phi)$ is also of vanishing type. It follows that $\delta_{\pm}^{N,2}$ is Γ -controllable with size N^{-1} and of vanishing type.

Consider now the second sum in (3.42). Using (3.43) and recalling the definition (3.44) of a_{\pm} , it reads:

$$\begin{aligned} \frac{1}{4} \sum_{i \in \{\pm(N-1)\}} c(\eta, i) &\left[\left(\frac{1}{2N} \sum_{j \neq i} \bar{\eta}_j N(2g - h)_{i,j} + (\mathbf{1}_{i=-(N-1)}\partial^N \lambda_{-N} - \mathbf{1}_{i=N-1}\partial^N \lambda_{N-1}) \right)^2 \right. \\ &\left. - \left(\frac{1}{2N} \sum_{j \neq i} \bar{\eta}_j N h_{i,j} \right)^2 \right] = \frac{a_-(\partial^N \lambda_{-N})^2}{4} + \frac{a_+(\partial^N \lambda_{N-1})^2}{4} + \delta_{\pm}^{N,3}(\eta), \end{aligned} \quad (3.48)$$

where $\delta_{\pm}^{N,3}$ is an error term that contains all other contributions:

$$\begin{aligned} \delta_{\pm}^{N,3}(\eta) &= \frac{1}{4} \sum_{i \in \{\pm(N-1)\}} c(\eta, i) \left[\left(\frac{1}{2N} \sum_{j \neq i} \bar{\eta}_j N(2g - h)_{i,j} \right)^2 - \left(\frac{1}{2N} \sum_{j \neq i} \bar{\eta}_j N h_{i,j} \right)^2 \right] \\ &+ \frac{1}{4N} \sum_{i \in \{\pm(N-1)\}} c(\eta, i) \left[(\mathbf{1}_{i=-(N-1)}\partial^N \lambda_{-N} - \mathbf{1}_{i=N-1}\partial^N \lambda_{N-1}) \sum_{j \neq i} \bar{\eta}_j N(2g - h)_{i,j} \right] \\ &+ \bar{\eta}_{-(N-1)}(1 - 2\rho_-) \frac{(\partial^N \lambda_{-N})^2}{4} + \bar{\eta}_{N-1}(1 - 2\rho_+) \frac{(\partial^N \lambda_{N-1})^2}{4}. \end{aligned} \quad (3.49)$$

With the notations of Lemma 3.9 and bounding $c(\eta, \cdot)$ by 1, the first line of $\delta_{\pm}^{N,3}$ is bounded by a term of the form $N^{-2}X_{2,\{0\}}^{\phi_2}$, thus N^{-1} times a quantity controllable with size 1. Similarly, the second line is bounded by a term of the form $N^{-1/2}|Y^N(\phi)|$ (recall that Y^N is defined in (3.26)) for bounded ϕ , and $|Y^N(\phi)|$ is controllable with size 1. Finally, the third line is of the form U_0^{\pm} , and U_0^{\pm} is Γ -controllable with size N^{-1} and of vanishing type. The quantity $\delta_{\pm}^{N,3}$ is therefore Γ -controllable with size $N^{-1/2}$ and of vanishing type.

Putting together (3.46) and (3.48), we have obtained the following expression of the adjoint at the boundary:

$$\begin{aligned} N^2 L_{h,\pm}^* \mathbf{1}(\eta) &= -\frac{N\partial^N \lambda_{-N}}{2} \left[\bar{\eta}_{-(N-1)} + \frac{\partial^N \bar{\rho}_{-N}}{N} \right] + \frac{N\partial^N \lambda_{N-1}}{2} \left[\bar{\eta}_{N-1} - \frac{\partial^N \bar{\rho}_{N-1}}{N} \right] \\ &\quad + \frac{a_-(\partial^N \lambda_{-N})^2}{4} + \frac{a_+(\partial^N \lambda_{N-1})^2}{4} + \sum_{q=1}^3 \delta_{\pm}^{N,q}(\eta). \end{aligned} \quad (3.50)$$

It remains to notice that the constant terms in the last equation compensate each other to obtain the final expression for $N^2 L_{h,\pm}^* \mathbf{1}(\eta)$. Indeed, for each $i \in \{\pm(N-1)\}$, a Taylor expansion yields:

$$a_{\text{sign}(i)} = 2\bar{\sigma}_i + O(N^{-1}), \quad \partial^N \lambda_i = \frac{\partial^N \bar{\rho}_i}{\bar{\sigma}_i} + O(N^{-1}). \quad (3.51)$$

It follows that there is a configuration-independent error term $\delta^{N,4}$, with $\delta^{N,4} = O(N^{-1})$, such that:

$$N^2 L_{h,\pm}^* \mathbf{1}(\eta) = -\frac{N\partial^N \lambda_{-N}}{2} \bar{\eta}_{-(N-1)} + \frac{N\partial^N \lambda_{N-1}}{2} \bar{\eta}_{N-1} + \delta_{\pm}^N(\eta), \quad \delta_{\pm}^N := \sum_{q=1}^4 \delta_{\pm}^{N,q}(\eta). \quad (3.52)$$

The quantities δ_{\pm}^N are, by definition, Γ -controllable with size N^{-1} and of vanishing type.

3.4 Adjoint in the bulk

We now compute $L_{h,0}^* \mathbf{1}$. For each $i < N-1$, define B_i^h, C_i^h, D_i^h as follows:

$$B_i^h(\eta) = \frac{1}{2N} \sum_{j \notin \{i, i+1\}} \bar{\eta}_j \partial_1^N h_{i,j}, \quad D_i^h(\eta) = \frac{\partial^N \bar{\rho}_i h_{i,i+1}}{2N}, \quad C_i^h = B_i^h + D_i^h, \quad (3.53)$$

where, for $u : \Lambda_N^2 \rightarrow \mathbb{R}$:

$$\partial_1^N u(i, j) = N[u(i+1, j) - u(i, j)], \quad i < N-1, j \in \Lambda_N. \quad (3.54)$$

With these definitions,

$$\forall i < N-1, \quad \Pi^N(h)(\eta^{i,i+1}) - \Pi^N(h)(\eta) = -\frac{(\eta_{i+1} - \eta_i)}{N} C_i^h(\eta). \quad (3.55)$$

Define similarly C_i^g , and notice that, since h, g are regular:

$$\sup_{N \in \mathbb{N}^*} \sup_{\eta \in \Omega_N} \sup_{i < N-1} (|C_i^h(\eta)| + |C_i^g(\eta)|) < \infty. \quad (3.56)$$

By (3.5), the adjoint $L_{h,0}^*$ in the bulk reads, by definition:

$$N^2 L_{h,0}^* \mathbf{1}(\eta) = \frac{N^2}{2} \sum_{i < N-1} \left[c_h(\eta^{i,i+1}, i, i+1) \frac{\nu_g^N(\eta^{i,i+1})}{\nu_g^N(\eta)} - c_h(\eta, i, i+1) \right]. \quad (3.57)$$

With the above notations, this becomes:

$$N^2 L_{h,0}^* \mathbf{1}(\eta) = \frac{N^2}{2} \sum_{i < N-1} c(\eta, i, i+1) \left[\exp \left[\frac{(\eta_{i+1} - \eta_i)}{N} (C_i^h - 2C_i^g) \right] \left[\frac{\bar{\rho}_i(1 - \bar{\rho}_{i+1})}{\bar{\rho}_{i+1}(1 - \bar{\rho}_i)} \right]^{\eta_{i+1} - \eta_i} - \exp \left[- \frac{(\eta_{i+1} - \eta_i)}{N} C_i^h \right] \right]. \quad (3.58)$$

To compute (3.58), recall the definition of λ :

$$\forall i \in \Lambda_N, \quad \lambda_i := \log \left(\frac{\bar{\rho}_i}{1 - \bar{\rho}_i} \right). \quad (3.59)$$

Notice that $C_i^h - 2C_i^g = C_i^{h-2g}$. Moreover, $c(\eta, i, i+1) = (\eta_{i+1} - \eta_i)^2$ for each $i < N-1$. With these notations, (3.58) reads:

$$N^2 L_{h,0}^* \mathbf{1}(\eta) = \frac{N^2}{2} \sum_{i < N-1} (\eta_{i+1} - \eta_i)^2 \left[\exp \left[\frac{(\eta_{i+1} - \eta_i)}{N} (C_i^{h-2g} - \partial^N \lambda_i) \right] - \exp \left[- \frac{(\eta_{i+1} - \eta_i)}{N} C_i^h \right] \right]. \quad (3.60)$$

To compute (3.60), we expand the above exponentials. Write $(N^2 L_{h,0}^* \mathbf{1})_{\text{order } p}$ for the term of order $p \in \mathbb{N}$. From the existence of $C_{h,g} = C(\|h\|_\infty, \|g\|_\infty) > 0$ such that $|e^x - 1 - x - x^2/2 - x^3/6| \leq C_{h,g} x^4$ when $|x| \leq 2(\|h\|_\infty + 2\|g\|_\infty)$, one has $|\delta_{0, \text{order} \geq 4}^N(\eta)| \leq 2C_{h,g}/N$, with:

$$N^2 L_{h,0}^* \mathbf{1}(\eta) = \frac{N}{2} \sum_{i < N-1} (\eta_{i+1} - \eta_i) \left[2C_i^{h-g} - \partial^N \lambda_i \right] \quad (3.61)$$

$$+ \frac{1}{4} \sum_{i < N-1} (\eta_{i+1} - \eta_i)^2 \left[(C_i^{h-2g} - \partial^N \lambda_i)^2 - (C_i^h)^2 \right] \quad (3.62)$$

$$+ \frac{1}{12N} \sum_{i < N-1} (\eta_{i+1} - \eta_i) \left[[C_i^{h-2g} - \partial^N \lambda_i]^3 + (C_i^h)^3 \right] + \delta_{0, \text{order} \geq 4}^N(\eta). \quad (3.63)$$

The sum in the last line (3.63) will later be found to be an error term, in Section 3.4.3. The important terms are therefore the sums in (3.61)-(3.62), which we will see impose conditions on the choice of g .

To highlight the structure of $L_{h,0}^* \mathbf{1}$, let us rewrite the sums in (3.61)-(3.62) by grouping together terms involving n -point correlations, $n \in \mathbb{N}^*$. By (3.53), C is the sum of B , which involves one-point correlations (i.e. one $\bar{\eta}$); and of D , which is configuration-independent, like λ . Moreover, the sum in (3.61) will have to be integrated by parts to remove the N factor. To do so, write:

$$\forall i < N-1, \quad \eta_{i+1} - \eta_i = \bar{\eta}_{i+1} - \bar{\eta}_i + \bar{\rho}_{i+1} - \bar{\rho}_i = \bar{\eta}_{i+1} - \bar{\eta}_i + N^{-1} \partial^N \bar{\rho}_i. \quad (3.64)$$

The sum in (3.61) therefore contains constant terms, fluctuations and two-point correlations. Let us similarly analyse the second line (3.62). The jump rate $(\eta_{i+1} - \eta_i)^2$, $i < N - 1$ can be expressed in terms of $\bar{\eta}_i$ and $\bar{\eta}_{i+1}$ as follows:

$$\forall i < N - 1, \quad (\eta_{i+1} - \eta_i)^2 = a_i + \sigma'(\bar{\rho}_i)\bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1})\bar{\eta}_i - 2\bar{\eta}_i\bar{\eta}_{i+1}, \quad (3.65)$$

where $a_i = \bar{\rho}_{i+1}(1 - \bar{\rho}_i) + \bar{\rho}_i(1 - \bar{\rho}_{i+1})$, $i < N - 1$. The sum in (3.62) therefore involves constant terms and n -point correlations for each $1 \leq n \leq 4$. In Section 3.4.3, we prove that three-point and four-point correlations lead to an error term $\delta_{0,3-4}^N$, while the sum in the third line (3.63) is an error term $\delta_{0,\text{order } 3}^N$. The adjoint in the bulk thus reads:

$$N^2 L_{h,0}^* \mathbf{1} = \text{Const} + \text{Fluct} + \text{Corr} + \delta_{0,3-4}^N + \delta_{0,\text{order } 3}^N + \delta_{0,\text{order} \geq 4}^N, \quad (3.66)$$

where Const, Fluct, Corr respectively denote the constant terms, the fluctuations and the correlations. The expression of these terms is given in the next three sections. Informally, these are small only when ν_g^N satisfies specific conditions. Namely, the fluctuations term Fluct is small because, by definition, ν_g^N has the same average occupation number as the invariant measure in the large N limit, see Section 3.4.1. On the other hand, the correlations Corr are small provided g solves the partial differential equation (2.61), as shown in Section 3.4.2. Finally, the constant Const is small provided all other terms are, as established in Section 3.4.4. We will repeatedly use the following estimates (recall the definition (3.53) of D):

$$\sup_{i < N-1} |D_i| = O(N^{-1}), \quad \sup_{i < N-1} \left| \frac{a_i \partial^N \lambda_i}{2} - \partial^N \bar{\rho}_i \right| = O(N^{-1}) = \sup_{i < N-1} \left| \partial^N \lambda_i - \frac{\partial^N \bar{\rho}_i}{\bar{\sigma}_i} \right|. \quad (3.67)$$

3.4.1 The fluctuations

Here, we estimate the fluctuations term Fluct in (3.66), which we recall accounts for all terms with a single $\bar{\eta}$ in the two sums (3.61)-(3.62). Recalling (3.65), it reads:

$$\begin{aligned} \text{Fluct} &= \frac{1}{2} \sum_{i < N-1} \left[N(\bar{\eta}_{i+1} - \bar{\eta}_i)(2D_i^{h-g} - \partial^N \lambda_i) + 2\partial^N \bar{\rho}_i B_i^{h-g} \right. \\ &\quad + \frac{a_i}{2} [2B_i^{h-2g}(D_i^{h-2g} - \partial^N \lambda_i) - 2B_i^g D_i^g] \\ &\quad \left. + \frac{1}{2} [\sigma'(\bar{\rho}_i)\bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1})\bar{\eta}_i] [(D_i^{h-2g} - \partial^N \lambda_i)^2 - (D_i^g)^2] \right]. \end{aligned} \quad (3.68)$$

To estimate the size of each term above, recall from Lemma 3.9 that a term of the form $Y^N(\phi)$ (defined in (3.26)), with $\phi : (-1, 1) \rightarrow \mathbb{R}$ bounded, is controllable with size 1. Using (3.67), (3.68) thus turns into:

$$\begin{aligned} \text{Fluct} &= \frac{1}{2} \sum_{i < N-1} \left[N(\bar{\eta}_{i+1} - \bar{\eta}_i)(2D_i^{h-g} - \partial^N \lambda_i) + 2\partial^N \bar{\rho}_i B_i^{h-g} - a_i \partial^N \lambda_i B_i^{h-2g} \right. \\ &\quad \left. + \frac{(\partial^N \lambda_i)^2}{2} [\sigma'(\bar{\rho}_i)\bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1})\bar{\eta}_i] \right] + \delta_{0,1}^{N,1}(\eta), \end{aligned} \quad (3.69)$$

where $\delta_{0,1}^{N,1}(\eta)$ reads:

$$\delta_{0,1}^{N,1}(\eta) := \frac{1}{2} \sum_{i < N-1} \left[a_i [B_i^{h-2g} D_i^{h-2g} - B_i^g D_i^g] + [\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] \right. \\ \left. \times [(D_i^{h-2g})^2 - 2D_i^{h-2g} \partial^N \lambda_i - (D_i^g)^2] \right]. \quad (3.70)$$

Recall from (3.53) that each B_i for $i < N-1$ is of the form $N^{-1/2} Y^N(\phi)$ for a bounded function ϕ , and that $\sup_i |D_i| = O(N^{-1})$. As a result, all terms composing $\delta_{0,1}^{N,1}$ are of the form $N^{-1/2} Y^N(\psi)$ or $N^{-3/2} Y^N(\psi)$ for a bounded ψ , and $Y^N(\psi)$ is controllable with size 1 by Lemma 3.9. It follows that $\delta_{0,1}^{N,1}$ is controllable with size N^{-1} and of vanishing type (recall Definitions 3.6–3.11).

Let us compute (3.69). We start by integrating its first term by parts. From (3.53) and the regularity of h, g , one draws, for each $i < N-1$:

$$D_{i-1}^{h-g} - D_i^{h-g} = -\frac{1}{2N^2} \left[(h-g)_{i,i+1} \Delta^N \bar{\rho}_i + N [(h-g)_{i,i+1} - (h-g)_{i-1,i}] \partial^N \bar{\rho}_{i-1} \right], \quad (3.71)$$

where:

$$\Delta^N \bar{\rho}_i = \partial^N (\partial^N \bar{\rho}_{i-1}) = N^2 [\bar{\rho}_{i+1} + \bar{\rho}_{i-1} - 2\bar{\rho}_i]. \quad (3.72)$$

Remark 3.15. *In the present case, $\Delta^N \bar{\rho}_i = 0$ and $\partial^N \bar{\rho}_i = \bar{\rho}'$ for each i , which could be used to simplify (3.71) and several other expressions below. We chose not to use the properties of $\bar{\rho}$ until the end of this section to highlight the structure of the Fluct term: it will contain a discrete PDE involving $\bar{\rho}$. Had we defined ν_g^N in terms of a density function $\rho : [-1, 1] \rightarrow [0, 1]$, this PDE would determine the choice of density ρ in the measure ν_g^N in order to obtain an optimal bound on the adjoint. \blacksquare*

As a result, $\sup_i |D_{i-1}^{h-g} - D_i^{h-g}| = O(N^{-2})$. Moreover, $g(\pm 1, \cdot) = 0 = h(\pm 1, \cdot)$, which implies that $D_{-(N-1)}^{h-g} = O(N^{-2}) = D_{N-2}^{h-g}$. An integration by parts therefore turns (3.69) into:

$$\text{Fluct} = \frac{1}{2} \sum_{|i| < N-1} \bar{\eta}_i \Delta^N \lambda_i - \frac{N}{2} \bar{\eta}_{N-1} \partial^N \lambda_{N-2} + \frac{N}{2} \bar{\eta}_{-(N-1)} \partial^N \lambda_{-(N-1)} \\ + \sum_{i < N-1} \left[\partial^N \bar{\rho}_i B_i^{h-g} - \frac{a_i \partial^N \lambda_i}{2} B_i^{h-2g} + \frac{(\partial^N \lambda_i)^2}{4} [\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] \right] + \sum_{q=1}^2 \delta_{0,1}^{N,q}(\eta), \quad (3.73)$$

where $\delta_{0,1}^{N,2}$ is a controllable error term with size N^{-1} of vanishing type, as it reads:

$$\delta_{0,1}^{N,2}(\eta) = \frac{1}{N} \sum_{|i| < N-1} \bar{\eta}_i N^2 [D_{i-1}^{h-g} - D_i^{h-g}] + \bar{\eta}_{N-1} D_{N-2}^{h-g} - \bar{\eta}_{-(N-1)} D_{-(N-1)}^{h-g}. \quad (3.74)$$

Using (3.67) to express $a \cdot \partial^N \lambda$ in terms of $\partial^N \bar{\rho}$ in the terms involving B in the second line, (3.73) becomes:

$$\text{Fluct} = \frac{1}{2} \sum_{|i| < N-1} \bar{\eta}_i \Delta^N \lambda_i - \frac{N}{2} \bar{\eta}_{N-1} \partial^N \lambda_{N-2} + \frac{N}{2} \bar{\eta}_{-(N-1)} \partial^N \lambda_{-(N-1)} \\ + \sum_{i < N-1} \left[\partial^N \bar{\rho}_i B_i^g + \frac{(\partial^N \lambda_i)^2}{4} [\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] \right] + \sum_{q=1}^3 \delta_{0,1}^{N,q}(\eta), \quad (3.75)$$

with $\delta_{0,1}^{N,3}(\eta)$ of the form $N^{-1/2}Y^N(\phi)$ (recall (3.26)) with ϕ bounded, thus controllable with size N^{-1} and of vanishing type. Indeed, recalling the estimates (3.67) and the definition (3.53) of B , $\delta_{0,1}^{N,3}$ reads:

$$\delta_{0,1}^{N,3}(\eta) = -\frac{1}{2N} \sum_{j \in \Lambda_N} \bar{\eta}_j \left(\frac{1}{N} \sum_{i \notin \{j-1, j, N-1\}} N \left[\frac{a_i \partial^N \lambda_i}{2} - \partial^N \bar{\rho}_i \right] \partial_1^N (h - 2g)_{i,j} \right), \quad (3.76)$$

and the term between parenthesis is bounded uniformly in $j \in \Lambda_N$ and $N \in \mathbb{N}^*$. Let us now compute the term involving B in (3.75). Recall that $g \in C^3(\mathbb{S})$ and is symmetric, thus $\sup_{|j| < N-1} N |g_{j-1, j} - g_{j+1, j}|$ is bounded uniformly in N . Integrating by parts and using $g(\pm 1, \cdot) = 0$, and the symmetry of g in the last line below; we find:

$$\begin{aligned} \sum_{i < N-1} \partial^N \bar{\rho}_i B_i^g &= \frac{1}{2N} \sum_{i < N-1} \partial^N \bar{\rho}_i \sum_{j \notin \{i, i+1\}} \bar{\eta}_j \partial_1^N g_{i,j} = \frac{1}{2N} \sum_{j \in \Lambda_N} \bar{\eta}_j \sum_{i \notin \{j-1, j, N-1\}} \partial^N \bar{\rho}_i \partial_1^N g_{i,j} \\ &= -\frac{1}{2N} \sum_{j \in \Lambda_N} \bar{\eta}_j \left[\sum_{\substack{|i| < N-1 \\ |j-i| > 1}} \Delta^N \bar{\rho}_i g_{i,j} + \partial^N \bar{\rho}_{j-2} (N g_{j-1, j}) - \partial^N \bar{\rho}_{j+1} (N g_{j+1, j}) \right. \\ &\quad \left. + \partial^N \bar{\rho}_{N-2} (N g_{N-1, j}) - \partial^N \bar{\rho}_{-N+1} (N g_{-N+1, j}) \right] \\ &=: -\frac{1}{2N} \sum_{j \in \Lambda_N} \bar{\eta}_j \sum_{\substack{|i| < N-1 \\ |j-i| > 1}} \Delta^N \bar{\rho}_i g_{i,j} + \delta_{0,1}^{N,4}(\eta) \\ &=: -\frac{1}{2} \sum_{j \in \Lambda_N} \bar{\eta}_j (N^{-1} M_g)(\Delta \bar{\rho})(j) + \delta_{0,1}^{N,5}(\eta), \end{aligned} \quad (3.77)$$

where $\delta_{0,1}^{N,4}, \delta_{0,1}^{N,5}$ are of the form $N^{-1/2}Y^N(\phi)$ ($Y^N(\phi)$ is defined in (3.26)) for bounded ϕ , and therefore $\delta_{0,1}^{N,4}, \delta_{0,1}^{N,5}$ are controllable with size N^{-1} and of vanishing type. M_g is the matrix $(g_{i,j})_{(i,j) \in \Lambda_N^2}$ and $(\Delta \bar{\rho})$ the vector $(\Delta \bar{\rho}_i)_{i \in \Lambda_N}$, so that $\delta_{0,1}^{N,5}$ accounts for the replacement of $\Delta^N \bar{\rho}$ by $\Delta \bar{\rho}$ (this cost vanishes in our case since $\bar{\rho}$ is linear, but we do not use this fact at this point), as well as the addition of missing terms in the sum on i :

$$\delta_{0,5}^N(\eta) = -\frac{1}{2N} \sum_{j \in \Lambda_N} \bar{\eta}_j \sum_{\substack{-N+1 < i < N-1 \\ |j-i| > 1}} [\Delta^N \bar{\rho}_i - \Delta \bar{\rho}_i] g_{i,j} - \frac{1}{2N} \sum_{j \in \Lambda_N} \bar{\eta}_j \sum_{i \in \{\pm(N-1), j, j \pm 1\}} \Delta \bar{\rho}_i g_{i,j}. \quad (3.78)$$

Consider now the sums involving λ in (3.75). Elementary computations give, for each $i < N-1$:

$$\Delta^N \lambda_i = N[\partial^N \lambda_i - \partial^N \lambda_{i-1}] = \frac{\Delta^N \bar{\rho}_i}{\bar{\sigma}_i} - \frac{(\partial^N \bar{\rho}_i)^2 \sigma'(\bar{\rho}_i)}{(\bar{\sigma}_i)^2} + \varepsilon_i^N, \quad \sup_{|i| < N-1} |\varepsilon_i^N| = O(N^{-1}). \quad (3.79)$$

By (3.67), we also know $\sup_i |\partial^N \lambda_i - \partial^N \bar{\rho}_i / \bar{\sigma}_i| = O(N^{-1})$. As a result:

$$\begin{aligned}
& \frac{1}{2} \sum_{|i| < N-1} \bar{\eta}_i \Delta^N \lambda_i + \sum_{i < N-1} \frac{(\partial^N \lambda_i)^2}{4} [\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] \\
& \quad - \frac{N}{2} \bar{\eta}_{N-1} \partial^N \lambda_{N-2} + \frac{N}{2} \bar{\eta}_{-(N-1)} \partial^N \lambda_{-(N-1)} \\
& = \frac{1}{2} \sum_{i \in \Lambda_N} \bar{\eta}_i \frac{\Delta \bar{\rho}_i}{\bar{\sigma}_i} - \frac{N}{2} \bar{\eta}_{N-1} \partial^N \lambda_{N-2} + \frac{N}{2} \bar{\eta}_{-(N-1)} \partial^N \lambda_{-(N-1)} + \delta_{0,1}^{N,6}(\eta), \tag{3.80}
\end{aligned}$$

where $\delta_{0,1}^{N,6}$ reads:

$$\begin{aligned}
\delta_{0,1}^{N,6}(\eta) & := \frac{1}{2} \sum_{|i| < N-1} \bar{\eta}_i \left[\Delta^N \lambda_i + \frac{1}{4} [(\partial^N \lambda_{i-1})^2 \sigma'(\bar{\rho}_{i-1}) + (\partial^N \lambda_i)^2 \sigma'(\bar{\rho}_{i+1})] - \frac{\Delta \bar{\rho}_i}{\bar{\sigma}_i} \right] \\
& \quad + \frac{\bar{\eta}_{N-1} (\partial^N \lambda_{N-2})^2 \sigma'(\bar{\rho}_{N-2})}{4} + \frac{\bar{\eta}_{-(N-1)} (\partial^N \lambda_{-(N-1)})^2 \sigma'(\bar{\rho}_{-(N-2)})}{4} \\
& \quad - \bar{\eta}_{N-1} \frac{\Delta \bar{\rho}_{N-1}}{\bar{\sigma}_{N-1}} - \bar{\eta}_{-(N-1)} \frac{\Delta \bar{\rho}_{-(N-1)}}{\bar{\sigma}_{-(N-1)}}. \tag{3.81}
\end{aligned}$$

The function $\delta_{0,1}^{N,6}$ involves terms of the form $\bar{\eta}_{\pm(N-1)} u(\pm(N-1))$ with u bounded (the last two lines), and $N^{-1/2} Y^N(\phi)$ for bounded $\phi : (-1, 1) \rightarrow \mathbb{R}$ (the first line). It follows that $\delta_{0,1}^{N,6}$ is Γ -controllable with size N^{-1} and of vanishing type by Lemma 3.13. Note that the last line actually vanishes since $\Delta \bar{\rho} = 0$. This last line is an error term anyway, so we do not need this fact. Putting (3.75), (3.77) and (3.80) together, we have computed the fluctuations term (3.68) in $N^2 L_{h,0}^* \mathbf{1}$:

$$\begin{aligned}
\text{Fluct} & = \frac{1}{2} \sum_{i \in \Lambda_N} \bar{\eta}_i (\bar{\sigma}^{-1} - N^{-1} M_g) (\Delta \bar{\rho}) (i) - \frac{N}{2} \bar{\eta}_{N-1} \partial^N \lambda_{N-2} \\
& \quad + \frac{N}{2} \bar{\eta}_{-(N-1)} \partial^N \lambda_{-(N-1)} + \sum_{q=1}^6 \delta_{0,1}^{N,q}(\eta). \tag{3.82}
\end{aligned}$$

Since $\bar{\rho}$ is the steady state profile satisfying $\Delta \bar{\rho} = 0$, the first sum vanishes, and:

$$\text{Fluct} = -\frac{N}{2} \bar{\eta}_{N-1} \partial^N \lambda_{N-2} + \frac{N}{2} \bar{\eta}_{-(N-1)} \partial^N \lambda_{-(N-1)} + \delta_{0,1}^N(\eta), \quad \delta_{0,1}^N := \sum_{q=1}^6 \delta_{0,1}^{N,q}. \tag{3.83}$$

By definition of error terms, see Definition 3.6, we have proven the following: for any $\theta, \gamma > 0$, there is $C(\theta, \gamma) > 0$ such that, for any density f for ν_g^N and N larger than some $N(\gamma)$:

$$\begin{aligned}
\nu_g^N(f \cdot \text{Fluct}) & \leq \frac{N}{2} \nu_g^N(f [-\bar{\eta}_{N-1} \partial^N \lambda_{N-2} + \bar{\eta}_{-(N-1)} \partial^N \lambda_{-(N-1)}]) \\
& \quad + \theta N^2 \nu_g^N(\Gamma_h(\sqrt{f})) + \frac{H(f \nu_g^N | \nu_g^N)}{\gamma} + \frac{C(\theta, \gamma)}{N}. \tag{3.84}
\end{aligned}$$

The expectation in the first line of (3.84) is not an error term, but it will cancel out with the boundary term obtained in (3.52). This cancellation will be proven in Section 3.5, where all contributions to $L_h^* \mathbf{1}$ will be summed.

3.4.2 The correlations

In this section, we compute the Corr term in (3.66) and obtain the partial differential equation that an optimal g must solve. Recall that Corr corresponds to all terms in (3.61)-(3.62) that involve products of two $\bar{\eta}$'s. It reads:

$$\begin{aligned} \text{Corr} = & \sum_{i < N-1} \left[N(\bar{\eta}_{i+1} - \bar{\eta}_i) B_i^{h-g} - a_i B_i^g B_i^{h-g} - \frac{1}{2} \bar{\eta}_i \bar{\eta}_{i+1} [(D_i^{h-2g} - \partial^N \lambda_i)^2 - (D_i^h)^2] \right. \\ & \left. + \frac{1}{2} [\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] [B_i^{h-2g} (D_i^{h-2g} - \partial^N \lambda_i) - B_i^h D_i^h] \right] \end{aligned} \quad (3.85)$$

Recall from Lemma 3.9 that terms of the form $\Pi^N(u)$, or $N^{-1/2} X_{1,\{0,1\}}^v = N^{-1/2} \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} v_i$, are controllable with size 1 as soon as the test functions u, v are bounded. Multiplying them by ε_N with $\varepsilon_N = o_N(1)$ therefore turns them into controllable error terms with size $\varepsilon_N, \varepsilon_N^2$ respectively by Lemma 3.9. As in Section 3.4.1, we first use the estimate (3.67) on the size of D to remove some terms from (3.85):

$$\begin{aligned} \text{Corr} = & \sum_{i < N-1} \left[N(\bar{\eta}_{i+1} - \bar{\eta}_i) B_i^{h-g} - a_i B_i^g B_i^{h-g} - \frac{1}{2} \bar{\eta}_i \bar{\eta}_{i+1} (\partial^N \lambda_i)^2 \right. \\ & \left. - \frac{1}{2} [\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] \partial^N \lambda_i B_i^{h-2g} \right] + \delta_{0,2}^{N,1}(\eta), \end{aligned} \quad (3.86)$$

with $\delta_{0,2}^{N,1}$ controllable with size N^{-1} of vanishing type, defined by:

$$\begin{aligned} \delta_{0,2}^{N,1}(\eta) := & -\frac{1}{2N} \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} N [(D_i^{h-2g})^2 - 2D_i^{h-2g} \partial^N \lambda_i - (D_i^h)^2] \\ & + \frac{1}{2N} \sum_{i < N-1} [\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] N [B_i^{h-2g} D_i^{h-2g} - B_i^h D_i^h]. \end{aligned} \quad (3.87)$$

Let us integrate by parts the term involving $\bar{\eta}_{i+1} - \bar{\eta}_i$ in (3.86), $i < N - 1$. To do so, notice first that, for each i with $|i| < N - 1$:

$$\begin{aligned} N[B_{i-1}^{h-g} - B_i^{h-g}] = N[B_i^{g-h} - B_{i-1}^{g-h}] = & \frac{1}{2N} \sum_{j:|j-i|>1} \bar{\eta}_j \Delta_1^N (g-h)_{i,j} \\ & + \frac{1}{2} [\bar{\eta}_{i-1} \partial_1^N (g-h)_{i,i-1} - \bar{\eta}_{i+1} \partial_1^N (g-h)_{i-1,i+1}], \end{aligned} \quad (3.88)$$

with $\Delta_1^N u(i, j) = \partial_1^N (\partial_1^N u(i-1, j))$ for $u : \mathbb{Z}^2 \rightarrow \mathbb{R}$ and $(i, j) \in \mathbb{Z}^2$. As a result:

$$\begin{aligned} \sum_{i < N-1} N(\bar{\eta}_{i+1} - \bar{\eta}_i) B_i^{h-g} = & N \bar{\eta}_{N-1} B_{N-2}^{h-g} - N \bar{\eta}_{-(N-1)} B_{-(N-1)}^{h-g} + \sum_{|i| < N-1} \bar{\eta}_i N [B_{i-1}^{h-g} - B_i^{h-g}] \\ = & \delta_{0,2}^{N,2}(\eta) + \frac{1}{2N} \sum_{|i| < N-1} \sum_{j:|j-i|>1} \bar{\eta}_i \bar{\eta}_j \Delta_1^N (g-h)_{i,j} \\ & + \frac{1}{2} \sum_{|i| < N-1} \bar{\eta}_i [\bar{\eta}_{i-1} \partial_1^N (g-h)_{i,i-1} - \bar{\eta}_{i+1} \partial_1^N (g-h)_{i-1,i+1}], \end{aligned} \quad (3.89)$$

where $\delta_{0,2}^{N,2}(\eta) := N\bar{\eta}_{N-1}B_{N-2}^{h-g} - N\bar{\eta}_{-(N-1)}B_{-(N-1)}^{h-g}$. $\delta_{0,2}^{N,2}$ involves correlations between the reservoir and the bulk, of the same form as the function U_1^\pm defined in Lemma 3.9. $\delta_{0,2}^{N,2}$ is thus Γ -controllable with size N^{-1} and of vanishing type. On the other hand, the last line of (3.89) can be integrated by parts once again:

$$\begin{aligned} & \sum_{|i|<N-1} \bar{\eta}_i [\bar{\eta}_{i-1} \partial_1^N (g-h)_{i,i-1} - \bar{\eta}_{i+1} \partial_1^N (g-h)_{i-1,i+1}] \\ &= \sum_{-(N-1)<i<N-2} \bar{\eta}_i \bar{\eta}_{i+1} [\partial_1^N (g-h)_{i+1,i} - \partial_1^N (g-h)_{i-1,i+1}] + \delta_{0,2}^{N,3}(\eta), \end{aligned} \quad (3.90)$$

with:

$$\delta_{0,2}^{N,3}(\eta) := \bar{\eta}_{-(N-1)} \bar{\eta}_{-(N-2)} \partial_1^N (g-h)_{-(N-2),-(N-1)} - \bar{\eta}_{N-2} \bar{\eta}_{N-1} \partial_1^N (g-h)_{N-3,N-1}. \quad (3.91)$$

The quantity $\delta_{0,2}^{N,3}$ is again a Γ -controllable error term with size N^{-1} of vanishing type, as it is of the same form as the function U_1^\pm of Lemma 3.9. Equation (3.89) thus reads:

$$\begin{aligned} \sum_{i<N-1} N(\bar{\eta}_{i+1} - \bar{\eta}_i) B_i^{h-g} &= \frac{1}{2N} \sum_{|i|<N-1} \sum_{j:|j-i|>1} \bar{\eta}_i \bar{\eta}_j \Delta_1^N (g-h)_{i,j} \\ &+ \frac{1}{2} \sum_{-(N-1)<i<N-2} \bar{\eta}_i \bar{\eta}_{i+1} [\partial_1^N (g-h)_{i+1,i} - \partial_1^N (g-h)_{i-1,i+1}] + \sum_{q=2}^3 \delta_{0,2}^{N,q}(\eta). \end{aligned} \quad (3.92)$$

The other terms in (3.86) are simpler. Indeed, recall that:

$$\sup_{i<N-1} |a_i - 2\bar{\sigma}_i| = O(N^{-1}), \quad \sup_{i<N-1} \left| \partial^N \lambda_i - \frac{\partial^N \bar{\rho}_i}{\bar{\sigma}_i} \right| = O(N^{-1}). \quad (3.93)$$

Using these estimates in (3.86), Corr becomes:

$$\begin{aligned} \text{Corr} &= \frac{1}{2N} \sum_{|i|<N-1} \sum_{j:|j-i|>1} \bar{\eta}_i \bar{\eta}_j \Delta_1^N (g-h)_{i,j} \\ &+ \frac{1}{2} \sum_{-(N-1)<i<N-2} \bar{\eta}_i \bar{\eta}_{i+1} \left[\partial_1^N (g-h)_{i+1,i} - \partial_1^N (g-h)_{i-1,i+1} - \frac{(\partial^N \bar{\rho}_i)^2}{(\bar{\sigma}_i)^2} \right] \\ &- \sum_{i<N-1} \left[[\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] \frac{\partial^N \bar{\rho}_i}{2\bar{\sigma}_i} B_i^{h-2g} + 2\bar{\sigma}_i B_i^g B_i^{h-g} \right] + \sum_{q=1}^4 \delta_{0,2}^{N,q}(\eta), \end{aligned} \quad (3.94)$$

where $\delta_{0,2}^{N,4}$ reads:

$$\begin{aligned} \delta_{0,2}^{N,4}(\eta) &= -\frac{1}{2N} \sum_{i<N-1} [\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] \left[\partial^N \lambda_i - \frac{\partial^N \bar{\rho}_i}{\bar{\sigma}_i} \right] B_i^{h-2g} - \frac{1}{N} \sum_{i<N-1} N[a_i - 2\bar{\sigma}_i] B_i^g B_i^{h-g} \\ &- \frac{1}{2N} \sum_{-(N-1)<i<N-2} \bar{\eta}_i \bar{\eta}_{i+1} N \left[(\partial^N \lambda_i)^2 - \frac{(\partial^N \bar{\rho}_i)^2}{\bar{\sigma}_i^2} \right] \\ &- \frac{1}{2} \left[\bar{\eta}_{-(N-1)} \bar{\eta}_{-(N-2)} (\partial^N \lambda_{-(N-1)})^2 + \bar{\eta}_{N-2} \bar{\eta}_{N-1} (\partial^N \lambda_{N-2})^2 \right]. \end{aligned} \quad (3.95)$$

The last line comes from the fact that the sum involving $\bar{\eta}_i \bar{\eta}_{i+1}$ in (3.86) and in (3.92) do not have the same range. It is of the same form as U_1^+ in Lemma 3.9, thus its multiplication by $N^{1/2}$ gives a Γ -controllable term with size 1. The first line above is of the form $N^{-2} X_{2,\{0\}}^{\phi_2}$, while the second line reads $N^{-1/2} X_{1,\{0,1\}}^{\phi_1}$; for bounded tensors ϕ_1, ϕ_2 . As a result, $\delta_{0,2}^{N,4}$ is Γ -controllable with size N^{-1} and of vanishing type by Lemmas 3.9–3.13.

To conclude on the expression of the correlations, it remains to take care of the two terms involving B in (3.94). Recalling the definition (3.53) of B , using the regularity of h, g and $\bar{\rho}$ and changing indices, one can write:

$$\begin{aligned} \frac{1}{2} \sum_{i < N-1} [\sigma'(\bar{\rho}_i) \bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1}) \bar{\eta}_i] \frac{\partial^N \bar{\rho}_i}{\bar{\sigma}_i} B_i^{2g-h} \\ = \frac{1}{2N} \sum_{|i| < N-1} \frac{\sigma'(\bar{\rho}_i) \partial^N \bar{\rho}_i}{\bar{\sigma}_i} \bar{\eta}_i \sum_{j \notin \{i, i+1\}} \bar{\eta}_j \partial_1^N (2g-h)_{i,j} + \delta_{0,2}^{N,5}(\eta), \end{aligned} \quad (3.96)$$

where $\delta_{0,2}^{N,5}$ is an error term that reads:

$$\begin{aligned} \delta_{0,2}^{N,5}(\eta) &= \frac{1}{2N} \sum_{|i| < N-1} \bar{\eta}_i N [\sigma'(\bar{\rho}_{i+1}) - \sigma'(\bar{\rho}_i)] \frac{\partial^N \bar{\rho}_i}{\bar{\sigma}_i} B_i^{2g-h} \\ &+ \frac{1}{2N} \sum_{|i| < N-1} \bar{\eta}_i N \left[\sigma'(\bar{\rho}_{i-1}) \frac{\partial^N \bar{\rho}_{i-1}}{\bar{\sigma}_{i-1}} B_{i-1}^{2g-h} - \sigma'(\bar{\rho}_i) \frac{\partial^N \bar{\rho}_i}{\bar{\sigma}_i} B_i^{2g-h} \right] \\ &+ \bar{\eta}_{N-1} \frac{\sigma'(\bar{\rho}_{N-2}) \partial^N \bar{\rho}_{N-2}}{2\bar{\sigma}_{N-2}} B_{N-2}^{2g-h} + \bar{\eta}_{-(N-1)} \frac{\sigma'(\bar{\rho}_{-(N-2)}) \partial^N \bar{\rho}_{-(N-1)}}{2\bar{\sigma}_{-(N-1)}} B_{-(N-1)}^{2g-h}. \end{aligned} \quad (3.97)$$

$\delta_{0,2}^{N,5}$ is of the form $N^{-1} X_{2,\{0\}}^{\phi_2}$ for a bounded ϕ_2 for the first two lines, and $N^{-1} U_1^\pm$ for the third line. By Lemmas 3.9–3.13, $\delta_{0,2}^{N,5}$ is therefore Γ -controllable with size N^{-1} and of vanishing type.

Finally, recall that $(\bar{\eta})^2 = \bar{\sigma} + \sigma'(\bar{\rho}) \bar{\eta}$. Separating diagonal and off-diagonal contributions, the term involving $B^g B^{h-g}$ in (3.94) reads:

$$\begin{aligned} -2 \sum_{i < N-1} \bar{\sigma}_i B_i^g B_i^{h-g} &= \frac{1}{2N^2} \sum_{i < N-1} \bar{\sigma}_i \sum_{j, \ell \notin \{i, i+1\}} \bar{\eta}_j \bar{\eta}_\ell \partial_1^N g_{i,j} \partial_1^N (g-h)_{i,\ell} \\ &= \frac{1}{2N} \sum_{\substack{|j| < N-1 \\ \ell \neq j}} \bar{\eta}_j \bar{\eta}_\ell \left(\frac{2}{2N} \sum_{i \notin \{j-1, j, \ell-1, \ell, N-1\}} \bar{\sigma}_i \partial_1^N g_{i,j} \partial_1^N (g-h)_{i,\ell} \right) \\ &+ \frac{1}{2N^2} \sum_{i < N-1} \sum_{j \notin \{i, i+1\}} \bar{\sigma}_i \bar{\sigma}_j \partial_1^N g_{i,j} \partial_1^N (g-h)_{i,j} + \delta_{0,2}^{N,6}(\eta), \end{aligned} \quad (3.98)$$

where $\delta_{0,2}^{N,6}$ is the sum of error terms of the form $N^{-1/2} Y^N(\phi)$ (recall (3.26)) and U_1^\pm , thus $\delta_{0,2}^{N,6}$ is

Γ -controllable with size N^{-1} and of vanishing type. It is given by :

$$\begin{aligned} \delta_{0,2}^{N,6}(\eta) &= \frac{1}{2N^2} \sum_{i < N-1} \sum_{j \notin \{i, i+1\}} \bar{\sigma}_i \sigma'(\bar{\rho}_j) \bar{\eta}_j \partial_1^N g_{i,j} \partial_1^N (g-h)_{i,j} \\ &+ \sum_{j \in \{\pm(N-1)\}} \sum_{\ell \neq j} \bar{\eta}_j \bar{\eta}_\ell \left(\frac{2}{2N} \sum_{i \notin \{j-1, j, \ell-1, \ell, N-1\}} \bar{\sigma}_i \partial_1^N g_{i,j} \partial_1^N (g-h)_{i,\ell} \right). \end{aligned} \quad (3.99)$$

The correlations (3.85) have so far been rewritten as follows:

$$\text{Corr} = \sum_{\substack{|i| < N-1 \\ j \neq i}} \bar{\eta}_i \bar{\eta}_j \left\{ \frac{1}{2N} \left[\mathbf{1}_{|i-j| > 1} \Delta_1^N (g-h)_{i,j} + \mathbf{1}_{j \neq i+1} \partial^N \bar{\rho}_i \frac{\sigma'(\bar{\rho}_i)}{\bar{\sigma}_i} \partial_1^N (2g-h)_{i,j} \right. \right. \quad (3.100)$$

$$\left. + \frac{1}{N} \sum_{\ell \notin \{i-1, i, j-1, j, N-1\}} \bar{\sigma}_\ell \partial_1^N g_{\ell,i} \partial_1^N (g-h)_{\ell,j} \right] \\ + \mathbf{1}_{j=i+1 < N-1} \frac{1}{2} \left[\partial_1^N (g-h)_{i+1,i} - \partial_1^N (g-h)_{i-1,i+1} - \frac{\partial^N \bar{\rho}_i}{(\bar{\sigma}_i)^2} \right] \left. \right\} \quad (3.101)$$

$$+ \frac{1}{2N^2} \sum_{i < N-1} \sum_{j \notin \{i, i+1\}} \bar{\sigma}_i \bar{\sigma}_j \partial_1^N g_{i,j} \partial_1^N (g-h)_{i,j} + \sum_{q=1}^6 \delta_{0,2}^{N,q}(\eta). \quad (3.102)$$

We claim that the curly bracket is a discrete version of the partial differential equation (2.61). To see it, first use the symmetry of g, h and exchange i, j . Recall then that, by assumption, $h, g \in W^{4,s}(\square)$ for some $s > 2$. By Sobolev embedding, $W^{4,s}(\square) \subset C^3(\bar{\square}) \cap C^3(\bar{\square})$, see Appendix E. As a result, approximating discrete derivatives by continuous ones and the Riemann sum in the second line of (3.101) by an integral, there is an error term $\delta_{0,2}^{N,7}$, controllable with size N^{-1} and of vanishing type, of the form $N^{-1} \Pi^N(u) + N^{-1} X_{1,\{0,1\}}^v$ for bounded u, v , such that:

$$\begin{aligned} (3.100) + (3.101) &= \Pi^N \left(\Delta(g-h) + \frac{(\bar{\sigma})'}{\bar{\sigma}} \partial_1(2g-h) + \partial_2(2g-h) \frac{(\bar{\sigma})'}{\bar{\sigma}} \right) \\ &+ \Pi^N \left(\mathcal{M}(\partial_1 g, \partial_1(g-h)) + \mathcal{M}(\partial_1(g-h), \partial_1 g) \right) \\ &+ \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} \left[\partial_1(g-h)_{i+,i} - \partial_1(g-h)_{i-,i} - \frac{(\bar{\rho}')^2}{\bar{\sigma}_i^2} \right] + \delta_{0,2}^{N,7}(\eta), \end{aligned} \quad (3.103)$$

with \mathcal{M} the bilinear operator defined in (2.46) and the convention, for $(x, y) \in \square$:

$$w\phi(x, y) = w(x)\phi(x, y), \quad \phi w(x, y) = \phi(x, y)w(y), \quad \phi : \square \rightarrow \mathbb{R}, w : (-1, 1) \rightarrow \mathbb{R}. \quad (3.104)$$

If g is chosen as the solution g_h of the main equation (2.61), then the right-hand side of (3.103) reduces to $\delta_{0,2}^{N,7}(\eta)$, whence:

$$\text{Corr} = \frac{1}{2N^2} \sum_{i < N-1} \sum_{j \notin \{i, i+1\}} \bar{\sigma}_i \bar{\sigma}_j \partial_1^N (g_h)_{i,j} \partial_1^N (g_h - h)_{i,j} + \delta_{0,2}^N(\eta), \quad \delta_{0,2}^N = \sum_{q=1}^7 \delta_{0,2}^{N,q}. \quad (3.105)$$

The first term is independent of the configuration, and will be estimated in Section 3.4.4. By definition of error terms (see Definition 3.6), for any $\theta, \gamma > 0$, there is thus $C(\theta, \gamma) > 0$ such that, for any density f for $\nu_{g_h}^N$ and any N large enough depending on γ :

$$\begin{aligned} \nu_{g_h}^N(f \cdot \text{Corr}) &\leq \frac{1}{2N^2} \sum_{i < N-1} \sum_{j \notin \{i, i+1\}} \bar{\sigma}_i \bar{\sigma}_j \partial_1^N(g_h)_{i,j} \partial_1^N(g_h - h)_{i,j} \\ &\quad + \theta N^2 \nu_{g_h}^N(\Gamma_h(\sqrt{f})) + \frac{H(f \nu_{g_h}^N | \nu_{g_h}^N)}{\gamma} + \frac{C(\theta, \gamma)}{N}. \end{aligned} \quad (3.106)$$

Remark 3.16. *The choice of $g = g_h$ cancelling the curly bracket (3.100)–(3.101) is optimal in the following sense: for another g , the first two terms in the right-hand side of (3.103) are not even error terms, but only controllable with size 1. As a result, for $g \neq g_h$, (3.106) can at best only be true with $\gamma_\theta N^{-1/2}, C_\theta$ respectively replacing $\gamma_\theta, C_\theta N^{-1/2}$ in the right-hand side, which breaks the $o_N(1)$ bound on the entropy obtained in Section 3.5. \blacksquare*

3.4.3 Higher-order correlations

In this section, we again assume that $g = g_h$ solves the main equation (2.61), and we estimate the third order term $(N^2 L_h^* \mathbf{1})_{\text{order } 3}$ in the development (3.63) of the exponentials making up the adjoint in the bulk, as well as three-point and four-point correlations arising in (3.62).

Consider first $(N^2 L_h^* \mathbf{1})_{\text{order } 3}$, which by (3.63) reads:

$$(N^2 L_h^* \mathbf{1}(\eta))_{\text{order } 3} = \frac{1}{12N} \sum_{i < N-1} (\eta_{i+1} - \eta_i) \left[[C_i^{h-2g_h} - \partial^N \lambda_i]^3 + [C_i^h]^3 \right]. \quad (3.107)$$

For $i < N-1$, write $\eta_{i+1} - \eta_i = \bar{\eta}_{i+1} - \bar{\eta}_i + N^{-1} \partial^N \bar{\rho}_i$ as before, and recall from (3.56) that C^{h-2g_h}, C^h are bounded with N . As a result:

$$\left| (N^2 L_h^* \mathbf{1}(\eta))_{\text{order } 3} - \frac{1}{12N} \sum_{i < N-1} (\bar{\eta}_{i+1} - \bar{\eta}_i) \left[[C_i^{h-2g_h} - \partial^N \lambda_i]^3 + [C_i^h]^3 \right] \right| \leq \frac{C(h)}{N}. \quad (3.108)$$

One need not even integrate by parts to find that the sum in (3.108) is an error term. Indeed, developing the cubes and recalling that $C = B + D$ (see (3.53)) with ND bounded, one finds:

$$\frac{1}{12N} \sum_{i < N-1} (\bar{\eta}_{i+1} - \bar{\eta}_i) \left[[C_i^{h-2g_h} - \partial^N \lambda_i]^3 + [C_i^h]^3 \right] = \sum_{n=1}^4 N^{-n} X_{n, \{0\}}^{\phi_n}(\eta), \quad (3.109)$$

for bounded ϕ_n , $1 \leq n \leq 4$, depending on ρ_\pm, h . By Lemmas 3.9–3.13, $N^{-n} X_{n, \{0\}}^{\phi_n}$ is controllable with size $N^{-3/2}$ at most for $1 \leq n \leq 4$, and of vanishing type. This observation and (3.108) yield:

$$\left| (N^2 L_h^* \mathbf{1}(\eta))_{\text{order } 3} \right| \leq \delta_{0, \text{order } 3}^N(\eta) := \frac{C(h)}{N} + \left| \sum_{n=1}^4 N^{-n} X_{n, \{0\}}^{\phi_n}(\eta) \right|, \quad (3.110)$$

and $\delta_{0, \text{order } 3}^N$ is controllable with size N^{-1} and of vanishing type.

Consider now three-and four-point correlations arising in the terms (3.63)–(3.64)–(3.65). They read:

$$\delta_{0,3-4}^N(\eta) := \sum_{i < N-1} \left[- [\sigma'(\bar{\rho}_i)\bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1})\bar{\eta}_i] B_i^{g_h} B_i^{h-g_h} + \bar{\eta}_i \bar{\eta}_{i+1} B_i^{g_h} (D_i^{2h-2g_h} - \partial^N \lambda_i) \right] \quad (3.111)$$

$$+ \bar{\eta}_i \bar{\eta}_{i+1} B_i^{h-g_h} (D_i^{2g_h} + \partial^N \lambda_i) \Big] + 2 \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} B_i^{g_h} B_i^{h-g_h}. \quad (3.112)$$

Recall that $\sup_i |D_i^\phi| = O(N^{-1})$ for any bounded $\phi : [-1, 1]^2 \rightarrow \mathbb{R}$, and that $\sup_i |\partial^N \lambda_i - \bar{\rho}'/\bar{\sigma}_i| = O(N^{-1})$. There are then bounded functions $\phi_n, \psi_n : (-1, 1)^n \rightarrow \mathbb{R}$ ($n \in \{2, 3\}$) such that, in the notations of Lemma 3.9:

$$\sum_{i < N-1} \left[- [\sigma'(\bar{\rho}_i)\bar{\eta}_{i+1} + \sigma'(\bar{\rho}_{i+1})\bar{\eta}_i] B_i^{g_h} B_i^{h-g_h} \right] =: N^{-2} X_{3,\{0\}}^{\phi_3}, \quad (3.113)$$

$$\sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} \left[B_i^{g_h} D_i^{2h-2g_h} + B_i^{h-g_h} D_i^{2g_h} \right] =: N^{-2} X_{2,\{0,1\}}^{\phi_2}, \quad (3.114)$$

$$\sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} B_i^{h-2g_h} \partial^N \lambda_i =: N^{-1} X_{2,\{0,1\}}^{\psi_2}, \quad (3.115)$$

$$2 \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} B_i^{g_h} B_i^{h-g_h} =: N^{-2} X_{3,\{0,1\}}^{\psi_3}, \quad (3.116)$$

so that $\delta_{0,3-4}^N$ reads:

$$\delta_{0,3-4}^N = N^{-2} X_{2,\{0,1\}}^{\phi_2} + N^{-2} X_{3,\{0\}}^{\phi_3} + N^{-2} X_{3,\{0,1\}}^{\psi_3} + N^{-1} X_{2,\{0,1\}}^{\psi_2}. \quad (3.117)$$

The first three terms are controllable error terms with size $N^{-1/2}$ at most, with the first term of vanishing type. The second and third term, involving ϕ_3, ψ_3 are of large type, and ϕ_3, ψ_3 are bounded as follows:

$$\max \left\{ \|\phi_3\|_\infty, \|\psi_3\|_\infty \right\} \leq \frac{1}{2} \|\partial_1 g_h\|_\infty \|\partial_1 (g_h - h)\|_\infty. \quad (3.118)$$

By taking $\bar{\rho}' \leq \varepsilon$ and $h \in \mathcal{S}(\varepsilon)$ for sufficiently small ε , the right-hand side can be made as small as needed, see Appendix F.5. Lemma 3.13 therefore ensures that $N^{-2} X_{3,\{0\}}^{\phi_3}$ and $N^{-2} X_{3,\{0,1\}}^{\psi_3}$ are of LS type.

The last term in (3.117) is $N^{-1} X_{2,\{0,1\}}^{\psi_2}$ given in (3.115). ψ_2 satisfies:

$$\|\psi_2\|_\infty \leq \frac{C \bar{\rho}' \|\partial_1 (h - 2g_h)\|_\infty}{\min\{\rho_-(1 - \rho_-), \rho_+(1 - \rho_+)\}} \xrightarrow{\bar{\rho}' \rightarrow 0} 0. \quad (3.119)$$

At first glance, the second item of Lemma 3.9 only yields that $N^{-1} X_{2,\{0,1\}}^{\psi_2}$ is controllable with size $\|\psi_2\|_\infty N^{-1/2}$. The last item of Lemma 3.9 in fact yields that $N^{-1} X_{2,\{0,1\}}^{\psi_2}$ is Γ -controllable with size $N^{-1/2}$. This improvement is obtained in Appendix C.1, through the renormalisation scheme of Jara and Menezes [JM18]. In view of the bound (3.119), taking $\bar{\rho}' \leq \varepsilon$ and $h \in \mathcal{S}(\varepsilon)$ for sufficiently small ε ensures that $N^{-1} X_{2,\{0,1\}}^{\psi_2}$ is of LS type. All terms in the decomposition (3.117) have now

been treated, and we conclude that, for small enough $\varepsilon > 0$ such that $\bar{\rho}' \leq \varepsilon$, $h \in \mathcal{S}(\varepsilon)$, $\delta_{0,3-4}^N$ is a Γ -controllable error term of LS type. Thus, for all large enough N and each $\delta > 0$:

$$\nu_g^N(f\delta_{0,3-4}^N) \leq \delta N^2 \nu_g^N(\Gamma_h(\sqrt{f})) + \frac{H(f\nu_g^N|\nu_g^N)}{2^{10}C_{LS}\delta} + \frac{C'(g,h)}{N^{1/2}\delta}. \quad (3.120)$$

Remark 3.17. *In addition to the boundary terms, the term $N^{-1}X_{2,\{0,1\}}^{\psi_2}$ is the only one for which a renormalisation scheme using the carre du champ is needed.* ■

3.4.4 The constant terms

Here, we prove that the configuration-independent terms appearing in the full adjoint $N^2L_h^*\mathbf{1}$ are small when $g = g_h$, with g_h the solution of the main equation (2.61). The terms in question correspond to various constant terms bounded by $O(N^{-1})$ encountered in the previous subsections and the computation of $N^2L_{h,\pm}^*\mathbf{1}$, which already are error terms; and the sum of the constant term in (3.105), as well as the Const term of (3.66), which reads:

$$\text{Const} := \sum_{i < N-1} \left[\partial^N \bar{\rho}_i \left(D_i^{h-g} - \frac{\partial^N \lambda_i}{2} \right) + \frac{a_i}{4} \left[(D_i^{h-2g} - \partial^N \lambda_i)^2 - (D_i^h)^2 \right] \right]. \quad (3.121)$$

By definition of $N^2L_h^*\mathbf{1}$, one has:

$$\nu_{g_h}^N(N^2L_h^*\mathbf{1}) = \nu_{g_h}^N(N^2L_h\mathbf{1}) = 0. \quad (3.122)$$

We can also estimate $\nu_{g_h}^N(N^2L_h^*\mathbf{1})$ through the expression (3.52) of the adjoint at the boundary and the expansion (3.63) of the adjoint in the bulk. Indeed, Lemma A.2 yields the estimates:

$$\forall n \geq 1, \quad \sup_{\substack{J \subset \Lambda_N \\ |J|=n}} \nu_{g_h}^N \left(\prod_{j \in J} \bar{\eta}_j \right) = O(N^{-n/2}), \quad (3.123)$$

$$\forall j \notin \{\pm(N-1)\}, \quad \nu_{g_h}^N(\bar{\eta}_{\pm(N-1)}\bar{\eta}_j) = O(N^{-2}), \quad \nu_{g_h}^N(\bar{\eta}_{\pm(N-1)}) = O(N^{-3/2}). \quad (3.124)$$

These bounds can be used on the error term δ_{\pm}^N involving the adjoint at the boundary (3.52), the error term $\delta_{\text{order} \geq 3,0}^N$ defined in (3.110), the error term $\delta_{3-4,0}^N$ accounting for three and four point correlations defined in (3.111)–(3.112), the estimates (3.83) of the fluctuations and (3.105) of the correlations. They yield:

$$\begin{aligned} \nu_{g_h}^N(N^2L_h^*\mathbf{1}) = 0 &= \text{Const} + \frac{1}{2N^2} \sum_{i < N-1} \sum_{j \notin \{i, i+1\}} \bar{\sigma}_i \bar{\sigma}_j \partial_1^N g_{i,j} \partial_1^N (g-h)_{i,j} + O(N^{-1}) \\ &=: \delta_{0,0}^N + O(N^{-1}). \end{aligned} \quad (3.125)$$

The configuration-independent terms $\delta_{0,0}^N$ arising in $N^2L_{h,0}^*\mathbf{1}$ are thus bounded by $O(N^{-1})$.

3.5 Conclusion

Let us put together the estimates obtained so far to conclude the proof of Lemma 3.5. The expression of the adjoint at the boundary was obtained in (3.52) for general g , while the adjoint

in the bulk has been estimated in the last three sections, provided one takes $g = g_h$, with g_h the solution of the main equation (2.61). One has therefore:

$$\begin{aligned}
N^2 L_h^* \mathbf{1}(\eta) &= N^2 L_{h,\pm}^* \mathbf{1}(\eta) + N^2 L_{h,0}^* \mathbf{1}(\eta) \\
&= \frac{N}{2} \bar{\eta}_{-(N-1)} \left(\partial^N \lambda_{-N} - \partial^N \lambda_{-(N-1)} \right) - \frac{N}{2} \bar{\eta}_{N-1} \left(\partial^N \lambda_{N-1} - \partial^N \lambda_{N-2} \right) + \delta^N(\eta) \\
&= \frac{\bar{\eta}_{-(N-1)}}{2} \Delta^N \lambda_{-(N-1)} - \frac{\bar{\eta}_{N-1}}{2} \Delta^N \lambda_{N-1} + \delta^N(\eta),
\end{aligned} \tag{3.126}$$

with δ^N a Γ -controllable error term with size $N^{-1/2}$, given by:

$$\delta^N(\eta) = \delta_{\pm}^N(\eta) + \delta_{0,0}^N + \delta_{0,1}^N(\eta) + \delta_{0,2}^N(\eta) + \delta_{0,3-4}^N(\eta) + \delta_{0,\text{order } 3}^N(\eta) + \delta_{0,\text{order } \geq 4}^N(\eta). \tag{3.127}$$

Since $\sup_N |\Delta^N \lambda_{\pm(N-1)}| < \infty$, the quantity:

$$\frac{\bar{\eta}_{-(N-1)}}{2} \Delta^N \lambda_{-(N-1)} - \frac{\bar{\eta}_{N-1}}{2} \Delta^N \lambda_{N-1} \tag{3.128}$$

is Γ -controllable with size N^{-1} by Lemma B.3, and of vanishing type. It follows that $N^2 L_h^* \mathbf{1}$ is Γ -controllable with size $N^{-1/2}$, and equal to a sum of terms of vanishing type, plus $\delta_{0,3-4}^N$ which is of LS type. Thus, by Definition 3.6 of Γ -controllability and taking $\delta = 1/2$ there, there is a controllable function \mathcal{E} with size $N^{-1/2}$ and $C(\rho_{\pm}, h)$ satisfying, for some $\gamma > 8C_{LS}$:

$$\frac{1}{\gamma} \log \nu_{g_h}^N \left(\exp [\gamma \mathcal{E}] \right) \leq \frac{C(\rho_{\pm}, h)}{N^{1/2}}, \tag{3.129}$$

and for any density f for $\nu_{g_h}^N$:

$$\begin{aligned}
\nu_{g_h}^N (f N^2 L_h^* \mathbf{1}) &\leq \nu_{g_h}^N (f \mathcal{E}) + \frac{N^2}{2} \nu_{g_h}^N (\Gamma_h(\sqrt{f})) \\
&\leq \frac{H(f \nu_{g_h}^N | \nu_{g_h}^N)}{8C_{LS}} + \frac{N^2}{2} \nu_{g_h}^N (\Gamma_h(\sqrt{f})) + \frac{C(\rho_{\pm}, h)}{N^{1/2}}.
\end{aligned} \tag{3.130}$$

This concludes the proof of Lemma 3.5. □

3.6 The Radon-Nikodym derivative

The computations in the previous subsections can be used to obtain an expression of the Radon-Nikodym derivative $D_h = d\mathbb{P}_h/d\mathbb{P}$ ($h \in \mathcal{S}(\infty)$) on each fixed time interval. By definition, for trajectories up to time $T > 0$, D_h reads (see Appendix A.7 in [KL99]):

$$\log D_h((\eta_t)_{t \leq T}) = \Pi_T^N(h) - \Pi_0^N(h) - N^2 \int_0^T e^{-\Pi_t^N(h)} L e^{\Pi_t^N(h)} dt. \tag{3.131}$$

The correlation field Π^N is defined in (2.7).

Proposition 3.18. *Let $h \in \mathcal{S}(\varepsilon_B)$, and recall that $\mathcal{M}(u, v)(x, y) = \int_{(-1,1)} u(z, x)\bar{\sigma}(z)u(z, y)dz$ for $u, v \in \mathbb{L}^2(\square)$ and $(x, y) \in \square$. Then, for each $\eta \in \Omega_N$:*

$$\begin{aligned} N^2 e^{-\Pi^N(h)} L e^{\Pi^N(h)}(\eta) &= \frac{1}{2} \Pi^N \left(\Delta h + \mathcal{M}(\partial_1 h, \partial_1 h) \right) - \frac{(\bar{\rho}')^2}{4} \int_{(-1,1)} h(x, x) dx \\ &+ \frac{1}{4} \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} (\partial_1 h_{i+,i} - \partial_1 h_{i-,i}) + \frac{1}{8} \int_{\square} \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 h(x, y)]^2 dx dy + \varepsilon^N(h), \end{aligned} \quad (3.132)$$

where $\varepsilon^N(h)$ is a Γ -controllable error term with size $N^{-1/2}$ and of LS type. It thus satisfies by Corollary 3.14, for some $C(h, \rho_{\pm}) > 0$, each large enough N and each $T > 0$:

$$\mathbb{E}_h^{\nu_h^N} \left[\exp \left[\int_0^T \varepsilon_t^N(h) dt \right] \right] \leq \exp \left[\frac{C(h, \rho_{\pm}) T}{N^{1/2}} \right]. \quad (3.133)$$

Remark 3.19. *A bias $h \in \mathcal{S}(\varepsilon_B)$ is a symmetric function by definition. As a result, for each $(x, y) \in \square$:*

$$\partial_1 h(x, y) = \partial_2 h(y, x) \quad \Rightarrow \quad \partial_1 h(x_+, x) - \partial_1 h(x_-, x) = (\partial_1 - \partial_2) h(x_+, x). \quad (3.134)$$

The first term in the second line of (3.132) thus corresponds to a contribution of the derivative of h normal to the diagonal. \blacksquare

The following corollary will be useful in the proof of lower bound large deviations.

Corollary 3.20. *Consider the \mathbb{P}_h -martingale $M^{N,\phi}$, defined for $T \geq 0$ and $\phi \in \mathcal{T}$ (\mathcal{T} is defined in (2.38)) by:*

$$M_T^{N,\phi} = \Pi_T^N(\phi) - \Pi_0^N(\phi) - \int_0^T N^2 L_h \Pi_t^N(\phi) dt. \quad (3.135)$$

If additionally ϕ is a symmetric function in $C^3(\bar{\square})$, there is a Γ -controllable error term $\tilde{\varepsilon}^N(h, \phi)$ with size $N^{-1/2}$ such that, for any $T \geq 0$:

$$\begin{aligned} M_T^{N,\phi} &= \Pi_T^N(\phi) - \Pi_0^N(\phi) - \frac{1}{2} \int_0^T \Pi_t^N \left(\Delta \phi + \mathcal{M}(\partial_1 \phi, \partial_1 h) + \mathcal{M}(\partial_1 h, \partial_1 \phi) \right) dt \\ &+ \frac{1}{4} \int_0^T \sum_{i < N-1} \bar{\eta}_i(t) \bar{\eta}_{i+1}(t) (\partial_1 \phi_{i+,i} - \partial_1 \phi_{i-,i}) dt + \frac{(\bar{\rho}')^2 T}{4} \int_{(-1,1)} \phi(x, x) dx \\ &- \frac{T}{4} \int_{\square} \bar{\sigma}(x) \bar{\sigma}(y) \partial_1 \phi(x, y) \partial_1 h(x, y) dx dy + \int_0^T \tilde{\varepsilon}_t^N(h, \phi) dt. \end{aligned} \quad (3.136)$$

When $\phi = h$, one has in particular:

$$\begin{aligned} M_T^{N,h} - \log D_h((\eta_t)_{t \leq T}) &= -\frac{1}{2} \int_0^T \Pi_t^N (\mathcal{M}(\partial_1 h, \partial_1 h)) dt - \frac{T}{8} \int_{\square} \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 h(x, y)]^2 dx dy \\ &+ \int_0^T \hat{\varepsilon}_t^N(h) dt, \end{aligned} \quad (3.137)$$

for a Γ -controllable error term $\hat{\varepsilon}^N(h)$ with size $N^{-1/2}$ of LS type.

4 Long-time behaviour: upper bound

In this section, we establish the upper bound in Theorem 2.3: for the ε_B of Theorem 2.6 and any closed set \mathcal{F} in $(\mathcal{T}'_s, *)$,

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P}^{\pi_{inv}^N} \left(\frac{1}{T} \int_0^T \Pi_t^N dt \in \mathcal{F} \right) \leq - \inf_{\Pi \in \mathcal{F}} \mathcal{I}_{\varepsilon_B}(\Pi), \quad (4.1)$$

with $\mathcal{I}_{\varepsilon_B}$ the functional defined in (2.50). A bound for compact sets is established in Section 4.1, relying on a regularity estimate in the space $\mathbb{H}^1(\square)$, proven in Section 4.2. Exponential tightness, which yields the bound for closed sets, is obtained in Section 4.3.

Before we start, let us make some remarks and fix notations. For $N \in \mathbb{N}^*$, $T > 0$, a probability measure μ^N on the state space Ω_N and each measurable set $\mathcal{B} \subset \mathcal{T}'_s$, introduce the notation:

$$\mathbb{Q}_T^{\mu^N}(\mathcal{B}) := \mathbb{P}^{\mu^N} \left(\frac{1}{T} \int_0^T \Pi_t^N dt \in \mathcal{B} \right). \quad (4.2)$$

For short, we will also write:

$$\mathcal{B}^T = \mathcal{B}_N^T := \left\{ \frac{1}{T} \int_0^T \Pi_t^N dt \in \mathcal{B} \right\}. \quad (4.3)$$

Recall that, for $N \in \mathbb{N}^*$, $\nu_{g_0}^N$ is the discrete Gaussian measure (2.58) built from the inverse correlation kernel g_0 of the steady state of the open SSEP. Changing initial condition from π_{inv}^N to the measure $\nu_{g_0}^N$ (defined in (3.11), with g_0 given by (2.59)) has a cost independent of T :

$$\forall T > 0, \forall N \in \mathbb{N}^*, \quad \mathbb{Q}_T^{\pi_{inv}^N}(\mathcal{B}) \leq \max_{\eta \in \Omega_N} \frac{\pi_{inv}^N(\eta)}{\nu_{g_0}^N(\eta)} \times \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B}). \quad (4.4)$$

Upon taking $T^{-1} \log$ and the large T , then large N limit, the initial condition of the dynamics does not change the value of the left-hand side in (4.1). The dynamics will therefore be started from the measure $\nu_{g_0}^N$.

4.1 Upper bound for open and compact sets

To estimate (4.2), we consider dynamics with tilted two-point correlations as in (2.41) and use the martingale method presented in Chapter 10 of [KL99]. It relies on the computation of the Radon-Nikodym derivatives $D_h = d\mathbb{P}_h/d\mathbb{P}$, $h \in \mathcal{S}(\varepsilon_B)$. In Section 4.1.1, a first upper bound on compact sets with a rate function $\tilde{\mathcal{I}}_{\varepsilon_B} \leq \mathcal{I}_{\varepsilon_B}$ is established. The bound is then improved to $\mathcal{I}_{\varepsilon_B}$ in Section 4.1.2 with the help of the regularity estimate of Section 4.2.

4.1.1 A first upper bound

Here, we build a functional $\tilde{\mathcal{I}}_{\varepsilon_B} : (\mathcal{T}'_s, *) \rightarrow \mathbb{R}_+$ such that, if $\mathcal{K} \subset (\mathcal{T}'_s, *)$ is a compact set,

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq - \inf_{\mathcal{K}} \tilde{\mathcal{I}}_{\varepsilon_B}. \quad (4.5)$$

We first prove an upper bound for general Borel sets, then specify to compactn sets. Let $h \in \mathcal{S}(\varepsilon_B)$. In Proposition 3.18, we proved that, for any $T > 0$ and any trajectory $(\eta(t))_{t \leq T}$:

$$\begin{aligned} \log D_h((\eta(t))_{t \leq T}) &= \Pi_T^N(h) - \Pi_0^N(h) - \frac{1}{2} \int_0^T \Pi_t^N \left(\Delta h + \mathcal{M}(\partial_1 h, \partial_1 h) \right) dt \\ &\quad - \frac{1}{4} \int_0^T \sum_{i < N-1} \bar{\eta}_i(t) \bar{\eta}_{i+1}(t) (\partial_1 h_{i+,i} - \partial_1 h_{i-,i}) + \frac{(\bar{\rho}')^2 T}{4} \int_{(-1,1)} h(x, x) dx \quad (4.6) \\ &\quad - \frac{T}{8} \int_{\square} \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 h(x, y)]^2 dx dy - \int_0^T \varepsilon_t^N(h) dt, \end{aligned}$$

with $\varepsilon^N(h)$ a Γ -controllable error term with size $N^{-1/2}$ of LS type (see Definitions 3.6–3.11). For any Borel set \mathcal{B} in \mathcal{T}'_s , one can thus write:

$$\mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B}) = \mathbb{E}_h^{\nu_{g_0}^N} \left[\mathbf{1}_{\mathcal{B}^T}(D_h)^{-1} \right], \quad (4.7)$$

and the point is to build the functional $\tilde{\mathcal{I}}_{\varepsilon_B}$ from D_h^{-1} ($h \in \mathcal{S}(\varepsilon_B)$).

Closed expression. Up to the error term $\varepsilon^N(h)$, the expression (4.6) is nearly closed in terms of the distribution $\int_0^T \Pi_t^N dt$ applied to regular test functions. The only problematic term is the diagonal term involving the $\bar{\eta}_i \bar{\eta}_{i+1}$, $i < N - 1$. Call it W_h^{Neu} :

$$W_h^{\text{Neu}}(\eta) = \frac{1}{4} \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} (\partial_1^N h_{i+,i} - \partial_1^N h_{i-,i}), \quad \eta \in \Omega_N. \quad (4.8)$$

In Appendix D, we estimate the cost of rewriting W_h^{Neu} in terms of the correlation field Π^N applied to a smooth test function supported in a small neighbourhood of the diagonal D . To state this estimate, consider a function $\chi^\varepsilon \in C^\infty(\bar{\square})$ with $\chi^\varepsilon = 0$ on $\partial \square$, $0 \leq \chi^\varepsilon \leq 2/\varepsilon$, such that $\chi^\varepsilon(x, \cdot)$ is supported on $(x, x + \varepsilon) \cap (-1, 1)$ for each $x \in (-1, 1)$, and:

$$\forall x < 1 - \varepsilon, \quad \int_{(x, x+\varepsilon)} \chi^\varepsilon(x, y) dy = 1, \quad \forall x \geq 1 - \varepsilon, \quad \int_{(x, x+\varepsilon)} \chi^\varepsilon(x, y) dy \leq 1. \quad (4.9)$$

Define then $\mathcal{N}_h^\varepsilon(x, y)$ for $(x, y) \in \square$ as follows:

$$\mathcal{N}_h^\varepsilon(x, y) = \frac{\bar{\sigma}(x)}{\bar{\sigma}(y)} \chi^\varepsilon(x, y) (\partial_1 - \partial_2) h(x_+, x). \quad (4.10)$$

Then $\mathcal{N}_h^\varepsilon \in \mathcal{T}$ (defined in (2.38)), and we prove the following in Proposition D.1. Let $\theta > 0$ and $\varepsilon \in (0, 1)$ be smaller than some $\varepsilon_0(\rho_\pm, h, \theta) > 0$. There are constants $C_1(\rho_\pm, h, \theta), C_2(\rho_\pm, h, \varepsilon, \theta) > 0$ such that, for each N large enough depending on ε, A and each $T > 0$:

$$\mathbb{E}_h^{\nu_{g_0}^N} \left[\exp \left[\theta \int_0^T [W_h^{\text{Neu}}(\eta_t) - \Pi_t^N(\mathcal{N}_h^\varepsilon)] dt \right] \right] \leq \exp \left[C_1(h, \rho_\pm, \theta) \varepsilon^{1/2} T + \frac{C_2(h, \rho_\pm, \varepsilon, \theta) T}{N^{1/2}} \right]. \quad (4.11)$$

The same bound is valid starting from $\nu_{g_0}^N$ up to an additional $e^{C(\rho_\pm, h)N}$ factor in the right-hand side (in fact only a factor $e^{C(\rho_\pm, h)}$, i.e. bounded with N , but this does not matter for our purposes).

Since we take T large first, this $e^{C(\rho_{\pm}, h)N}$ multiplicative constant is not a problem. Introduce then the following continuous functional J_h^ε on $(\mathcal{T}'_s, *)$:

$$\begin{aligned} \forall \Pi \in \mathcal{T}'_s, \quad J_h^\varepsilon(\Pi) &= -\frac{1}{2}\Pi\left(\Delta h + \mathcal{M}(\partial_1 h, \partial_1 h)\right) - \Pi(\mathcal{N}_h^\varepsilon) \\ &\quad + \frac{(\bar{\rho}')^2}{4} \int_{(-1,1)} h(x, x) dx - \frac{1}{8} \int_{\square} \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 h(x, y)]^2 dx dy. \end{aligned} \quad (4.12)$$

With this definition, the Radon-Nikodym derivative (4.6) becomes:

$$\begin{aligned} \log D_h((\eta(t))_{t \leq T}) &= \Pi_T^N(h) - \Pi_0^N(h) + T J_h^\varepsilon\left(\frac{1}{T} \int_0^T \Pi_t^N dt\right) \\ &\quad - \int_0^T [W_h^{\text{Neu}}(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^h)] dt - \int_0^T \varepsilon_t^N(h) dt. \end{aligned} \quad (4.13)$$

Estimates on error terms. Let us estimate the contribution of $\varepsilon^N(h) + W_h^{\text{Neu}}(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^h)$. By Proposition 3.18, $\varepsilon^N(h)$ will not contribute to the large deviations as it has small exponential moment. Indeed, Proposition 3.18 together with the last item of Lemma 3.13 give that, for some small enough $\alpha > 0$:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P}^{\nu_{g_h}^N} \left[\exp \left[(1 + \alpha) \int_0^T \varepsilon_t^N(h) dt \right] \right] = 0. \quad (4.14)$$

Since a change of initial condition has no influence in the large T limit, the above estimate is also valid starting from $\nu_{g_0}^N$. Thus, by Hölder inequality:

$$\begin{aligned} \frac{1}{T} \log \mathbb{P}_h^{\nu_{g_0}^N} \left[\exp \left[\int_0^T [W_h^{\text{Neu}}(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^h)] dt + \int_0^T \varepsilon_t^N(h) dt \right] \right] \\ \leq \frac{\alpha}{(1 + \alpha)T} \log \mathbb{P}_h^{\nu_{g_0}^N} \left[\exp \left[\int_0^T \frac{(1 + \alpha)}{\alpha} [W_h^{\text{Neu}}(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^h)] dt \right] \right] \\ + \frac{1}{(1 + \alpha)T} \log \mathbb{P}_h^{\nu_{g_0}^N} \left[\exp \left[\int_0^T (1 + \alpha) \varepsilon_t^N(h) dt \right] \right]. \end{aligned} \quad (4.15)$$

Combining (4.11)–(4.14), we find that there is a constant $C = C(h, \rho_{\pm}, \alpha) > 0$ such that:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P}_h^{\nu_{g_0}^N} \left[\exp \left[\int_0^T [W_h^{\text{Neu}}(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^h)] dt + \int_0^T \varepsilon_t^N(h) dt \right] \right] \leq C\varepsilon^{1/2}. \quad (4.16)$$

We can now try to obtain the upper bound (4.5). Take a Borel set \mathcal{B} . Starting from (4.7), one has:

$$\begin{aligned} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B}) &= \frac{1}{T} \log \mathbb{E}_h^{\nu_{g_0}^N} \left[\mathbf{1}_{\mathcal{B}^T}(D_h)^{-1} \right] \\ &\leq \frac{1}{T} \log \mathbb{E}_h^{\nu_{g_0}^N} \left[\exp \left[-\Pi_T^N(h) + \Pi_0^N(h) + \int_0^T [W_h^{\text{Neu}}(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^h)] dt + \int_0^T \varepsilon_t^N(h) dt \right] \right] \\ &\quad + \sup_{\Pi \in \mathcal{B}} (-J_h^\varepsilon(\Pi)) \\ &\leq \frac{C(h)N}{T} + \frac{1}{T} \log \mathbb{E}_h^{\nu_{g_0}^N} \left[\exp \left[\int_0^T [W_h^{\text{Neu}}(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^h)] dt + \int_0^T \varepsilon_t^N(h) dt \right] \right] + \sup_{\Pi \in \mathcal{B}} (-J_h^\varepsilon(\Pi)). \end{aligned} \quad (4.17)$$

Taking the large T , then large N limits and estimating the error terms via (4.16), one finds that, for all $\varepsilon > 0$ smaller than some $\varepsilon_0 = \varepsilon_0(\rho_{\pm}, h, \alpha)$:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B}) \leq C(h, \rho_{\pm})\varepsilon^{1/2} + \sup_{\Pi \in \mathcal{B}} (-J_h^\varepsilon(\Pi)). \quad (4.18)$$

Taking the infimum on $\varepsilon \in (0, \varepsilon_0)$ and $h \in \mathcal{S}(\varepsilon_B)$ yields a first bound:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B}) \leq \inf_{h \in \mathcal{S}(\varepsilon_B)} \inf_{\varepsilon \in (0, \varepsilon_0)} \sup_{\Pi \in \mathcal{B}} (-J_h^\varepsilon(\Pi)). \quad (4.19)$$

We now work out a way to exchange supremum and infima in (4.19) when \mathcal{B} is a compact set. The argument is standard and relies on Lemmas 3.2 and 3.3 in Appendix 2 of [KL99]. Let $\mathcal{K} \subset (\mathcal{T}', *)$ be compact. We wish to prove:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq \sup_{\mathcal{K}} \inf_{h \in \mathcal{S}(\varepsilon_B)} \liminf_{\varepsilon \rightarrow 0} (-J_h^\varepsilon). \quad (4.20)$$

Since $(J_h^\varepsilon)_{h, \varepsilon}$ is a family of continuous functionals on $(\mathcal{T}', *)$, Lemmas 3.2 and 3.3 in Appendix 2 of [KL99] allow for the exchange of the infima on h, ε and the supremum on (open covers of) \mathcal{K} :

$$\limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq \sup_{\mathcal{K}} \inf_{h \in \mathcal{S}(\varepsilon_B)} \inf_{\varepsilon \in (0, \varepsilon_0)} J_h^\varepsilon. \quad (4.21)$$

Since $\inf_{\varepsilon < \varepsilon_0}$ increases when ε_0 shrinks, (4.21) yields:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq \sup_{\Pi \in \mathcal{K}} \inf_{h \in \mathcal{S}(\varepsilon_B)} \liminf_{\varepsilon \rightarrow 0} (-J_h^\varepsilon(\Pi)). \quad (4.22)$$

This yields a first bound on compact sets and proves (4.5):

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq -\inf_{\mathcal{K}} \tilde{\mathcal{I}}_{\varepsilon_B}, \quad \tilde{\mathcal{I}}_{\varepsilon_B} = \sup_{h \in \mathcal{S}(\varepsilon_B)} \limsup_{\varepsilon \rightarrow 0} J_h^\varepsilon. \quad (4.23)$$

4.1.2 Refinement of the upper bound to restrict to more regular correlations

To improve the expression for the functional $\tilde{\mathcal{I}}_{\varepsilon_B}$, we would like to take the limit in ε in (4.23). This is however not possible in general. Indeed, recall that $\Pi \in \mathcal{T}'_s$ is a distribution, and taking ε to 0 amounts to asking for Π to have a well-defined trace on the diagonal D of the square \square . This is possible only if Π has some regularity.

In this section, we explain how to improve the upper bound (4.23) so that it is finite only on correlation fields Π that have a well defined trace. More precisely, write $\Pi = \frac{1}{4} \langle k_{\Pi}, \cdot \rangle \in \mathcal{T}'_s$. we prove that one can restrict to Π 's with $k_{\Pi} \in \mathbb{H}^1(\square)$. Note that, as $\Pi \in \mathcal{T}'_s \cap \mathbb{H}^1(\square)$, k_{Π} is a symmetric function, and the traces on both sides of the diagonal coincide. We thus write $\text{tr}_D(k_{\Pi})$ for the trace of k_{Π} on the diagonal without ambiguity. By definition of the trace, one has then:

$$\lim_{\varepsilon \rightarrow 0} \Pi(\mathcal{N}_h^\varepsilon) = \lim_{\varepsilon \rightarrow 0} \frac{1}{4} \langle k_{\Pi}, \mathcal{N}_h^\varepsilon \rangle = \frac{1}{4} \int_{(-1,1)} \text{tr}(k_{\Pi})(x, x) (\partial_2 - \partial_1) h(x_+, x) dx. \quad (4.24)$$

Thus, for $\Pi \in \mathcal{T}'_s \cap \mathbb{H}^1(\square)$, $\limsup_{\varepsilon \rightarrow 0} |J_h^\varepsilon(\Pi) - J_h(\Pi)| = 0$, with J_h the functional equal to $+\infty$ outside of $\mathcal{T}'_s \cap \mathbb{H}^1(\square)$, and:

$$\begin{aligned} \forall \Pi \in \mathcal{T}'_s \cap \mathbb{H}^1(\square), \quad J_h(\Pi) = & -\frac{1}{2}\Pi\left(\Delta h + \mathcal{M}(\partial_1 h, \partial_1 h)\right) + \frac{1}{4} \int_{(-1,1)} \text{tr}_D(k_\Pi)(\partial_2 - \partial_1)h(x_+, x)dx \\ & + \frac{(\bar{\rho}')^2}{4} \int_{(-1,1)} h(x, x)dx - \frac{1}{8} \int_{\square} \bar{\sigma}(x)\bar{\sigma}(y)[\partial_1 h(x, y)]^2 dx dy. \end{aligned} \quad (4.25)$$

As a result, for $\Pi \in \mathcal{T}'_s \cap \mathbb{H}^1(\square)$, $\tilde{\mathcal{I}}_{\varepsilon_B}(\Pi) = \mathcal{I}(\Pi)$, with $\tilde{\mathcal{I}}_{\varepsilon_B}$ defined in (4.23), and $\mathcal{I}_{\varepsilon_B} \geq \tilde{\mathcal{I}}_{\varepsilon_B}$ the better rate function (recall its definition (2.50)). The goal of the section is summarised in the following lemma.

Lemma 4.1. *Let $\mathcal{K} \subset (\mathcal{T}'_s, *)$ be a compact set. Then:*

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq -\inf_{\mathcal{K}} \mathcal{I}_{\varepsilon_B}. \quad (4.26)$$

To prove Lemma 4.1, consider the functionals \mathcal{Q} and \mathcal{Q}_ϕ , defined for each $\phi \neq 0$ in $C_c^\infty(\square)$, the set of compactly supported, C^∞ functions on \square , by:

$$\forall \Pi \in \mathcal{T}'_s, \quad \mathcal{Q}_\phi(\Pi) = \frac{\Pi(\partial_1 \phi)}{\|\phi\|_2}, \quad \mathcal{Q} = \sup_{\phi \in C_c^\infty(\square) \setminus \{0\}} \mathcal{Q}_\phi, \quad (4.27)$$

Note that \mathcal{Q} is weak* lower semi-continuous on $(\mathcal{T}'_s, *)$. Indeed, it is a supremum over the \mathcal{Q}_ϕ for $\phi \in C_c^\infty(\square) \setminus \{0\}$, and \mathcal{Q}_ϕ is weak* continuous since $\|\phi\|_2 \mathcal{Q}_\phi(\Pi)$ is the evaluation of Π at $\partial_1 \phi \in \mathcal{T}$ for $\Pi \in \mathcal{T}'_s$.

It is classical that \mathcal{Q} controls the regularity of elements of \mathcal{T}'_s , as stated in the next lemma.

Lemma 4.2. *For $\Pi \in \mathcal{T}'_s$, $\Pi \in \mathbb{H}^1(\square)$ if and only if $\mathcal{Q}(\Pi) < \infty$.*

Proof of Lemma 4.1. We now begin the proof of the large deviation bound (4.26). Consider a sequence $\phi_j \in C_c^\infty(\square) \setminus \{0\}$, $j \in \mathbb{N}^*$, dense for the norm of $\mathbb{H}^3(\square)$. Introduce then, for each $\ell \in \mathbb{N}^*$ and each $A > 0$:

$$U(\ell, A) = \left\{ \max_{1 \leq j \leq \ell} \mathcal{Q}_{\phi_j} \leq A \right\}. \quad (4.28)$$

In Proposition 4.3, we prove the existence of $C = C(\rho_\pm) > 0$ such that, for A larger than some $A_0 > 0$ and each $\ell \in \mathbb{N}^*$:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}((U(\ell, A))^c) \leq -CA. \quad (4.29)$$

Notice also that, for $\Pi \in \mathcal{T}'_s$, if (ϕ_{j_n}) converges to $\phi \in C_c^\infty(\square)$ in the norm of $\mathbb{H}^3(\square)$, then $\partial_1 \phi_{j_n}$ converges to $\partial_1 \phi$ in \mathcal{T} , so that $\lim_n \mathcal{Q}_{\phi_{j_n}}(\Pi) = \mathcal{Q}_\phi(\Pi)$, and:

$$\mathcal{Q}(\Pi) = \sup_{j \in \mathbb{N}^*} \mathcal{Q}_{\phi_j}(\Pi). \quad (4.30)$$

The regularity of the correlation fields can therefore be controlled through the \mathcal{Q}_{ϕ_j} ($j \in \mathbb{N}^*$). Let $\ell \in \mathbb{N}^*$ and $A > A_0$. Recall the notation:

$$U(\ell, A)^T := \left\{ \frac{1}{T} \int_0^T \Pi_t dt \in U(\ell, A) \right\}, \quad (4.31)$$

so that $\mathbb{P}(U(\ell, A)^T) = \mathbb{Q}_T(U(\ell, A))$. Let $\mathcal{B} \subset (\mathcal{T}'_s, *)$ be a Borel set. For ease of writing, let us abbreviate limits in N, T as follows:

$$\limsup_{N, T \rightarrow \infty} := \limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty}. \quad (4.32)$$

We again estimate $\mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B})$, starting from the bound:

$$\begin{aligned} \limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B}) &\leq \max \left\{ \limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B} \cap U(\ell, A)), \limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(U(\ell, A)^c) \right\} \\ &\leq \max \left\{ \limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B} \cap U(\ell, A)), -CA \right\}, \end{aligned} \quad (4.33)$$

where (4.29) was used to obtain the second inequality. Proceeding as in (4.17)–(4.19) to estimate the first probability, we find:

$$\limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B}) \leq \max \left\{ C(h, \rho_{\pm})\varepsilon + \sup_{\Pi \in \mathcal{B} \cap U(\ell, A)} (-J_h^\varepsilon(\Pi)), -CA \right\}. \quad (4.34)$$

For each admissible h, ε, ℓ, A , let $J_{h, \ell, A}^\varepsilon$ be equal to $+\infty$ in $U(\ell, A)^c$, and on $U(\ell, A)$:

$$J_{h, \ell, A}^\varepsilon := \max \left\{ C(h, \rho_{\pm})\varepsilon - J_h^\varepsilon(\Pi), -CA \right\}. \quad (4.35)$$

Minimising (4.34) on $\varepsilon \in (0, \varepsilon_0)$, $h \in \mathcal{S}(\varepsilon_B)$, $\ell \in \mathbb{N}^*$ and $A > A_0$, we have therefore obtained the upper bound:

$$\limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{B}) \leq \inf_{A, \ell, h, \varepsilon} \sup_{\mathcal{B}} \bar{J}_{h, \ell, A}^\varepsilon. \quad (4.36)$$

Let us obtain a bound on compact sets from (4.36). Let $\mathcal{K} \subset (\mathcal{T}'_s, *)$ be compact. Since $U(\ell, A)$ is weak* closed by continuity of each \mathcal{Q}_{f_j} ($1 \leq j \leq \ell$), $(J_{h, \ell, A}^\varepsilon)_{h, \varepsilon, \ell, A}$ is a family of weak* upper semi-continuous functionals. Lemmas A.2.3.2 and A.2.3.3 in [KL99] thus give as before:

$$\limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq \sup_{\mathcal{K}} \inf_{h \in \mathcal{S}(\varepsilon_B)} \inf_{\ell, A, \varepsilon} J_{h, \ell, A}^\varepsilon, \quad (4.37)$$

with the infimum still taken on $\ell \in \mathbb{N}^*$, $A > A_0$ and $\varepsilon \in (0, \varepsilon_0)$.

One can again bound the infimum on ε by a liminf:

$$\limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq \sup_{\mathcal{K}} \inf_{h \in \mathcal{S}(\varepsilon_B)} \inf_{\ell, A} \max \left\{ \liminf_{\varepsilon \rightarrow 0} (-J_{h, \ell, A}^\varepsilon), -CA \right\}. \quad (4.38)$$

Let $A > A_0$. As $U(\ell, A) \subset U(\ell', A)$ if $\ell \leq \ell'$, the argument of the supremum on \mathcal{K} in (4.38) is equal to $-\infty$ when evaluated at any $\Pi \notin \bigcap_{\ell \in \mathbb{N}^*} U(\ell, A)$. By definition of $U(\ell, A)$ in (4.28) and using (4.30), one has:

$$\bigcap_{\ell \in \mathbb{N}^*} U(\ell, A) = \{\mathcal{Q} \leq A\}. \quad (4.39)$$

Equation (4.38) thus becomes:

$$\limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq \sup_{\mathcal{K}} \inf_{h \in \mathcal{S}(\varepsilon_B)} \inf_{A > A_0} \max \left\{ \liminf_{\varepsilon \rightarrow 0} (-J_{h, A}^\varepsilon), -CA \right\}, \quad (4.40)$$

with, for each h, ε, A , $J_{h,A}^\varepsilon = +\infty$ on $\{\mathcal{Q} > A\}$, and $J_{h,A}^\varepsilon = J_h^\varepsilon$ on $\{\mathcal{Q} \leq A\}$. Consider again $A > A_0$. For each $\Pi \in \{\mathcal{Q} \leq A\}$, the k_Π associated with Π via $\Pi(\cdot) = \frac{1}{4}\langle k_\Pi, \cdot \rangle$ belongs to $\mathbb{H}^1(\square)$ by Lemma 4.2. In particular, by (4.24), if $h \in \mathcal{S}(\varepsilon_B)$ and $\Pi \in \{\mathcal{Q} \leq A\}$,

$$\liminf_{\varepsilon \rightarrow 0} (-J_{h,A}^\varepsilon(\Pi)) = -J_{h,A}(\Pi). \quad (4.41)$$

Above, $J_{h,A} = J_h$ on $\{\mathcal{Q} \leq A\}$, $J_{h,A} = +\infty$ outside, and J_h is defined in (4.25). Equation (4.40) thus becomes:

$$\limsup_{N, T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq \sup_{\mathcal{K}} \inf_{h \in \mathcal{S}(\varepsilon_B)} \inf_{A > A_0} \max \{ -J_{h,A}, -CA \}. \quad (4.42)$$

Finally, note that $J_{h,A} \geq J_h$ for $A > A_0$, since J_h may be finite on $\{\mathcal{Q} > A\}$ while $J_{h,A}$ may not. Lemma 4.1 is proven:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{Q}_T^{\nu_{g_0}^N}(\mathcal{K}) \leq -\inf_{\mathcal{K}} \mathcal{I}_{\varepsilon_B}, \quad \mathcal{I}_{\varepsilon_B} := \sup_{h \in \mathcal{S}(\varepsilon_B)} J_h. \quad (4.43)$$

□

4.2 Regularity estimate

In this section, we prove the energy estimate (4.29). The key argument is the following proposition.

Proposition 4.3. *Let $0 < \rho_- \leq \rho_+ < 1$ and assume $\bar{\rho}' \leq \varepsilon_B$, with ε_B given by Theorem 2.6. Let $A > 0$ be larger than some fixed $A_0 > 0$ and let $\phi \in C_c^2(\square)$, where the subscript c stands for compactly supported. There are constants $C_1(\rho_\pm), C_2(\rho_\pm) > 0$ and $C_3(\rho_\pm, \phi)$ such that, for each $T > 0$, each $A > 0$ and each N larger than some $N_0(\phi)$:*

$$\mathbb{P}^{\nu_{g_0}^N} \left(\left| \frac{1}{T} \int_0^T \Pi_t^N(\partial_1 \phi) dt \right| > A \|\phi\|_2 \right) \leq 2 \exp \left[-C_1(\rho_\pm)(A - C_2(\rho_\pm))T + \frac{C_3(\rho_\pm, \phi)T}{N^{1/2}} \right]. \quad (4.44)$$

In particular,

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P}^{\nu_{g_0}^N} \left(\left| \frac{1}{T} \int_0^T \Pi_t^N(\partial_1 \phi) dt \right| > A \|\phi\|_2 \right) \leq -C_1(\rho_\pm)(A - C_2(\rho_\pm)). \quad (4.45)$$

Assuming the proposition, (4.29) immediately follows by a union bound and the appropriate limits.

Proof of Proposition 4.3. Up to considering $-\phi$, it is enough to prove the result without the absolute value and the factor 2 in front of the probability. Up to taking N large enough depending on ϕ , we can assume that the support of ϕ is contained in $\{z \in \square : d(z, \partial \square) > 2/N\}$, so that $\phi_{i,i+1} = 0 = \partial_1 \phi_{i,i+1}$ for each $i < N - 1$. We may also assume without loss of generality that ϕ is symmetric owing to the identity $\Pi^N(\phi) = \Pi^N(\check{\phi})$, with $\check{\phi}(x, y) = \phi(y, x)$ for $x, y \in \square$. The idea is to use Feynman-Kac inequality (Lemma 3.3) and a microscopic integration by parts to rewrite, for each density f for $\nu_{g_0}^N$, the average $\nu_{g_0}^N(f \|\phi\|_2^{-1} \Pi^N(\partial_1 \phi) dt)$ as $\nu_{g_0}^N(f \Pi^N(F(\|\phi\|_2^{-1} \phi)))$ plus a term involving the carré du champ. The term $F(\|\phi\|_2^{-1} \phi)$ now involves only $\|\phi\|_2^{-1} \phi$, and not its

first partial derivative. The average $\nu_{g_0}^N(f\Pi^N(F(\|\phi_2\|_2^{-1}\phi)))$ is then estimated through the entropy inequality.

By Feynman-Kac inequality (3.6) together with the bound (3.17) on the adjoint, one has, for some $C = C(\rho_{\pm}) > 0$ and some $\kappa > 0$ to be chosen later:

$$\begin{aligned} & \frac{1}{T} \log \mathbb{P}^{\nu_{g_0}^N} \left(\frac{1}{T} \int_0^T \|\phi\|_2^{-1} \Pi_t^N(\partial_1 \phi) dt > A \right) \\ & \leq -A\kappa + \sup_{f \geq 0: \nu_{g_0}^N(f)=1} \left\{ \kappa \|\phi\|_2^{-1} \nu_{g_0}^N(f\Pi^N(\partial_1 \phi)) - \frac{N^2}{8} \nu_{g_0}^N(\Gamma(f^{1/2})) + \frac{C(\rho_{\pm})}{N^{1/2}} \right\}. \end{aligned} \quad (4.46)$$

To estimate the right-hand side, let us write out $\|\phi\|_2^{-1} \Pi^N(\partial_1 \phi)$. Fix $\eta \in \Omega_N$.

$$\|\phi\|_2^{-1} \Pi^N(\partial_1 \phi) = \frac{1}{4\|\phi\|_2 N} \sum_{i < N-1} \left[\sum_{j \notin \{i, i+1\}} \bar{\eta}_i \bar{\eta}_j \partial_1^N \phi_{i,j} + \bar{\eta}_i \bar{\eta}_{i+1} \partial_1 \phi_{i, i+1} \right] + \frac{1}{N\|\phi\|_2} \Pi^N(b), \quad (4.47)$$

where $b = N[\partial_1 \phi - \partial_1^N \phi]$ is the discretisation error, which is bounded with N . By assumption on the compact support of ϕ in \square , $\partial_1 \phi_{i, i+1} = 0$ for each $i < N - 1$, and (4.47) becomes:

$$\|\phi\|_2^{-1} \Pi^N(\partial_1 \phi) = \frac{1}{4\|\phi\|_2 N} \sum_{i < N-1} \sum_{j \notin \{i, i+1\}} \bar{\eta}_i \bar{\eta}_j \partial_1^N \phi_{i,j} + \frac{1}{N\|\phi\|_2} \Pi^N(b). \quad (4.48)$$

By Lemma 3.9, as b is bounded, the second term in the right-hand side above is controllable with size N^{-1} , and of vanishing type. Let us rewrite the first term through an integration by parts:

$$\begin{aligned} \frac{1}{4\|\phi\|_2} \sum_{i < N-1} \sum_{j \notin \{i, i+1\}} \bar{\eta}_i \bar{\eta}_j (\phi_{i+1, j} - \phi_{i, j}) &= \frac{1}{4\|\phi\|_2} \sum_{|i| < N-1} \sum_{j: |j-i| > 1} (\bar{\eta}_{i-1} - \bar{\eta}_i) \bar{\eta}_j \phi_{i, j} \\ &+ \frac{1}{4\|\phi\|_2} \sum_{|i| < N-1} \bar{\eta}_{i-1} (\bar{\eta}_{i+1} \phi_{i, i+1} - \bar{\eta}_i \phi_{i, i-1}) \\ &= \frac{1}{4\|\phi\|_2} \sum_{|i| < N-1} \sum_{j: |j-i| > 1} (\bar{\eta}_{i-1} - \bar{\eta}_i) \bar{\eta}_j \phi_{i, j} =: S, \end{aligned} \quad (4.49)$$

where the first equality makes use of $\phi_{\pm(N-1), \cdot} = 0$, while the second equality follows from $\phi_{i, i\pm 1} = 0$ for each $|i| < N - 1$. To estimate the supremum in (4.46), we see from (4.49) that we have to estimate $\nu_{g_0}^N(fS)$. This is done through the integration by parts Lemma B.2. This lemma is formulated with the variables $\omega_i = \bar{\eta}_i / \bar{\sigma}_i$ for $i \in \Lambda_N$, so we first rewrite (4.49) in terms of these variables. For $|i| < N - 1$, using the identity:

$$\begin{aligned} \bar{\eta}_{i-1} - \bar{\eta}_i &= \bar{\sigma}_{i-1}(\omega_{i-1} - \omega_i) - (\bar{\sigma}_i - \bar{\sigma}_{i-1})\omega_i \\ &= \bar{\sigma}_{i-1}(\omega_{i-1} - \omega_i) - \frac{\bar{\rho}'}{N} \left(\sigma'(\bar{\rho}_i) + \frac{\bar{\rho}'}{N} \right) \omega_i \end{aligned} \quad (4.50)$$

and the convention:

$$q\phi(x, y) = q(x)\phi(x, y) \quad \text{for } q: (-1, 1) \rightarrow \mathbb{R} \text{ and } (x, y) \in \square, \quad (4.51)$$

the right-hand side S of (4.49) can be rewritten as follows:

$$\begin{aligned} S &= \frac{1}{4\|\phi\|_2} \sum_{|i|<N-1} \sum_{j:|j-i|>1} \bar{\sigma}_{i-1}(\omega_{i-1} - \omega_i) \bar{\eta}_j \phi_{i,j} - \frac{1}{\|\phi\|_2} \Pi^N \left(\frac{\bar{\rho}'}{\bar{\sigma}} \left(\sigma'(\bar{\rho}) + \frac{\bar{\rho}'}{N} \right) \phi \right) \\ &=: S' + \Pi^N(Y^{(0)}), \quad Y^{(0)} = -\frac{1}{\|\phi\|_2} \frac{\bar{\rho}'}{\bar{\sigma}} \left(\sigma'(\bar{\rho}) + \frac{\bar{\rho}'}{N} \right) \phi. \end{aligned} \quad (4.52)$$

$\Pi^N(Y^{(0)})$ is of the form $\|\phi\|_2^{-1} \Pi^N(q\phi)$ with q bounded independently of ϕ , i.e. the kind of expression we were after. It remains to estimate S' . Define, for $|i| < N - 1$, a function v_i on Ω_N as follows:

$$\forall \eta \in \Omega_N, \quad v_i(\eta) = \frac{1}{4} \sum_{j:|j-i|>1} \bar{\eta}_j \bar{\sigma}_{i-1} \frac{\phi_{i,j}}{\|\phi\|_2}. \quad (4.53)$$

Recall also the notation C^{g_0} defined in (3.53):

$$\forall \eta \in \Omega_N, \forall i < N - 1, \quad \Pi^N(g_0)(\eta^{i,i+1}) - \Pi^N(g_0)(\eta) =: -\frac{(\eta_{i+1} - \eta_i)}{N} C_i^{g_0}. \quad (4.54)$$

With these notations, we can apply the integration by parts Lemma B.2 to each $|i| < N - 1$ with $u = v_i$, and obtain the existence of $C > 0$ such that, for each $\delta > 0$ and each density f for $\nu_{g_0}^N$:

$$\begin{aligned} \nu_{g_0}^N(fS') &\leq \delta N^2 \nu_{g_0}^N(\Gamma(f^{1/2})) + \frac{C}{\delta N^2} \sum_{|i|<N-1} \nu_{g_0}^N(f(v_i)^2) \\ &\quad - \sum_{|i|<N-1} (\bar{\rho}_i - \bar{\rho}_{i-1}) \nu_{g_0}^N(\omega_{i-1} \omega_i f e^{-(\eta_i - \eta_{i-1}) C_{i-1}^{g_0}/N} v_i) \\ &\quad + \sum_{|i|<N-1} \nu_{g_0}^N \left((\omega_i - \omega_{i-1}) (1 - e^{-(\eta_i - \eta_{i-1}) C_{i-1}^{g_0}/N}) f v_i \right). \end{aligned} \quad (4.55)$$

We express each term appearing in (4.55) as $\Pi^N(G(\phi))$ plus error terms, for explicit G 's. Consider first the second term on the first line. By definition (4.53) of v and using $(\bar{\eta})^2 = \bar{\sigma} + \sigma'(\bar{\rho})\bar{\eta}$, it reads:

$$\begin{aligned} &\frac{C}{16\delta N} \nu_{g_0}^N \left(f \sum_{j,\ell \in \Lambda_N} \frac{\bar{\eta}_j \bar{\eta}_\ell}{N} \sum_{\substack{|i|<N-1: \\ |i-j|>1, |i-\ell|>1}} \bar{\sigma}_i^2 \frac{\phi_{i,j} \phi_{i,\ell}}{\|\phi\|_2^2} \right) \\ &= \frac{1}{\delta} \nu_{g_0}^N(f \Pi^N(Y^{(1)})) + \frac{C}{16\delta} \int_{\square} \bar{\sigma}(x) \bar{\sigma}(y)^2 \frac{\phi(x,y)^2}{\|\phi\|_2^2} dx dy + \nu_{g_0}^N(f \theta_\delta^{N,1}(\phi)), \end{aligned} \quad (4.56)$$

with $Y^{(1)}$ the function recording the off-diagonal, $\ell \neq j$ contribution:

$$\forall (x, y) \in \square, \quad Y^{(1)}(x, y) = \frac{C}{4} \int_{(-1,1)} \bar{\sigma}(z)^2 \frac{\phi(z, x) \phi(z, y)}{\|\phi\|_2^2} dz. \quad (4.57)$$

The error term $\theta_\delta^{N,1}(\|\phi\|_2^{-1}\phi)$ in (4.56) involves discretisation errors and the diagonal, $\ell = j$ contributions. It is given for $\eta \in \Omega_N$ by:

$$\begin{aligned} \theta^{N,1}(\|\phi\|_2^{-1}\phi)(\eta) &= \frac{C}{16\delta} \int_{\mathbb{Z}} \bar{\sigma}(x)\bar{\sigma}(y)^2 \frac{\phi(x,y)^2}{\|\phi\|_2^2} dx dy - \frac{C}{16\delta N^2} \sum_{\substack{|i| < N-1 \\ |j-i| > 1}} \bar{\sigma}_i \bar{\sigma}_j^2 \frac{\phi_{i,j}^2}{\|\phi\|_2^2} \\ &+ \frac{C}{16\delta N^2} \nu_{g_0}^N \left(f \sum_{\substack{|i| < N-1 \\ |j-i| > 1}} \bar{\eta}_i \sigma'(\bar{\rho}_i) \bar{\sigma}_j^2 \frac{\phi_{i,j}^2}{\|\phi\|_2^2} \right) + \frac{1}{N} \Pi^N(c), \end{aligned} \quad (4.58)$$

where c is a discretisation error arising in the replacement of (4.56) by $Y^{(1)}$. We write $\theta^{N,1}(\|\phi\|_2^{-1}\phi)$ to emphasise the fact that θ only depends on $\|\phi\|_2^{-1}\phi$ rather than ϕ itself. The first line of $\theta^{N,1}(\|\phi\|_2^{-1}\phi)$ is configuration-independent, and bounded by $C(\phi)/N$. The first sum on the second line is of the form $N^{-1/2}Y^N(u)$ for $u : (-1, 1) \rightarrow \mathbb{R}$ bounded, with $Y^N(u)$ the fluctuations defined in (3.26). $\theta^{N,1}(\|\phi\|_2^{-1}\phi)$ is therefore controllable with size N^{-1} by Lemma 3.9, and of vanishing type by Lemma 3.13. For later use, note that the middle term in (4.56) is bounded by $C/(2^{10}\delta)$ for all large enough N , as $\bar{\sigma} \leq 1/4$.

Consider now line 2 of (4.55). Using the identity $e^x = 1 + \int_0^1 x e^{tx} dt$ for $x \in \mathbb{R}$, there is $C(\rho_\pm) > 0$ such that:

$$\begin{aligned} & \left| \sum_{|i| < N-1} (\bar{\rho}_i - \bar{\rho}_{i-1}) \nu_{g_0}^N (\omega_{i-1} \omega_i f e^{-(\eta_i - \eta_{i-1}) C_{i-1}^{g_0}/N} v_i) - \frac{\bar{\rho}'}{N} \sum_{|i| < N-1} \nu_{g_0}^N (\omega_{i-1} \omega_i f v_i) \right| \\ & \leq \frac{C(\rho_\pm)}{N^2} \sum_{|i| < N-1} \nu_{g_0}^N (f |v_i C_{i-1}^{g_0}|) \end{aligned} \quad (4.59)$$

The second sum in the left-hand side already involves three point correlations. It is shown to be Γ -controllable with size $N^{-1/2}$ in Proposition C.1, and is of large type. The second term is of the form $N^{-2}X_{2,\{0\}}^{v_2}$ in the notations of Lemma 3.9, and therefore controllable with size N^{-1} and of vanishing type. It follows that the first sum is a Γ -controllable error term with size $N^{-1/2}$ and of large type, depending only on $\|\phi\|_2^{-1}\phi$ rather than ϕ itself. We rewrite it as follows:

$$\theta^{N,2}(\|\phi\|_2^{-1}\phi) := - \sum_{|i| < N-1} (\bar{\rho}_i - \bar{\rho}_{i-1}) \nu_{g_0}^N (\omega_{i-1} \omega_i f e^{-(\eta_i - \eta_{i-1}) C_{i-1}^{g_0}/N} v_i). \quad (4.60)$$

By Lemma 3.13, there is a numerical constant $\gamma_2 > 0$ such that $\gamma_2 \theta^{N,2}(\|\phi\|_2^{-1}\phi)$ is of LS type. Consider finally line 3 of (4.55). Let $C(\bar{\rho}) > 0$ be such that:

$$\forall |i| < N-1, \quad |\omega_i - \omega_{i-1}| \leq C(\bar{\rho}). \quad (4.61)$$

Using this time the existence of $c(g_0) > 0$ such that $|e^x - 1 - x| \leq c(g_0)x^2$ for all $|x| \leq 2\|g_0\|_\infty$, one can write, for each $\eta \in \Omega_N$:

$$\begin{aligned} \sum_{|i| < N-1} (\omega_i - \omega_{i-1}) (1 - e^{-(\eta_i - \eta_{i-1}) C_{i-1}^{g_0}/N}) v_i &\leq \frac{1}{N} \sum_{|i| < N-1} (\omega_i - \omega_{i-1}) (\eta_i - \eta_{i-1}) C_{i-1}^{g_0} v_i \\ &+ \frac{c(g_0)C(\bar{\rho})}{N} \sum_{|i| < N-1} \frac{|v_i|}{N} (C_{i-1}^{g_0})^2. \end{aligned} \quad (4.62)$$

The last term is an average over i of terms of the form $N^{-3}|X_{3,\{0\}}^{w_3^i}|$ with the notations of Lemma 3.9, where the w_3^i satisfy $\sup_{N,i} \sup_{\Lambda_N^3} |w_3^i| < \infty$. It is therefore controllable with size $N^{-3/2}$ and of vanishing type. To estimate the first sum in the right-hand side of (4.62), we use the following elementary identity, valid for each $|i| < N - 1$:

$$(\omega_i - \omega_{i-1})(\eta_i - \eta_{i-1}) = \left[2 + (1 - \bar{\rho}_i - \bar{\rho}_{i-1})[\omega_i + \omega_{i-1}] - (\bar{\sigma}_{i-1} + \bar{\sigma}_i)\omega_{i-1}\omega_i \right]. \quad (4.63)$$

This identity can be obtained by making the following observation:

$$\forall i \in \Lambda_N, \quad \eta_i \omega_i := \eta_i \frac{(\eta_i - \bar{\rho}_i)}{\bar{\sigma}_i} = \frac{\eta_i}{\bar{\sigma}_i} (1 - \bar{\rho}_i) = \frac{\bar{\eta}_i}{\bar{\rho}_i} + 1. \quad (4.64)$$

Looking at (4.62), we see that the term $C_{i-1}^{g_0} v_i$ already contains two-point correlations for each $|i| < N - 1$. We therefore claim that only the constant term in the identity (4.63) will give something that is not an error term in (4.62). More precisely, we claim that one can obtain the following bound for line 3 of (4.62):

$$\begin{aligned} & \sum_{|i| < N-1} \nu_{g_0}^N \left((\omega_i - \omega_{i-1}) (1 - e^{-(\eta_i - \eta_{i-1}) C_{i-1}^{g_0}/N}) f v_i \right) \\ & \leq \nu_{g_0}^N (f \Pi^N (Y^{(2)})) + \frac{1}{4} \int_{\mathbb{Z}} \bar{\sigma}(x) \bar{\sigma}(y) \frac{\phi(x, y)}{\|\phi\|_2} \partial_1 g_0(x, y) dx dy + \nu_{g_0}^N (f \theta^{N,3} (\|\phi\|_2^{-1} \phi)), \end{aligned} \quad (4.65)$$

where $\theta^{N,3}(\|\phi\|_2^{-1} \phi)$ is controllable with size $N^{-1/2}$, of large type, and again depends on $\|\phi\|_2^{-1} \phi$ only rather than on ϕ . By Lemma 3.13, there is thus $\gamma_3 > 0$ such that $\theta^{N,3}(\|\phi\|_2^{-1} \phi)$ is of LS type. There is $C(g_0) = C(\rho_{\pm}) > 0$ independent of ϕ bounding the middle term in (4.65), which comes from the constant term in the identity $(\bar{\eta})^2 = \bar{\sigma} + \sigma'(\bar{\rho})\bar{\eta}$, and $Y^{(2)}$ is defined as:

$$\forall (x, y) \in \mathbb{Z}, \quad Y^{(2)}(x, y) = \int_{(-1,1)} \bar{\sigma}(z) \partial_1 g_0(z, x) \frac{\phi(z, y)}{\|\phi\|_2} dz. \quad (4.66)$$

To summarise, we have established the following. If, for $\delta > 0$, V_δ is the quantity:

$$V_\delta = \Pi^N (\delta^{-1} Y^{(0)} + Y^{(1)} + Y^{(2)}), \quad (4.67)$$

then, for each N large enough depending on ϕ :

$$\nu_{g_0}^N (f \|\phi\|_2^{-1} \Pi^N (\partial_1 \phi)) \leq \delta N^2 \nu_{g_0}^N (\Gamma(\sqrt{f})) + \nu_{g_0}^N (f V_\delta) + C(\rho_{\pm}) + \zeta_\delta^N (\|\phi\|_2^{-1} \phi). \quad (4.68)$$

The quantity $\zeta_\delta^N (\|\phi\|_2^{-1} \phi)$ is a Γ -controllable error term of size $N^{-1/2}$ of large type. One can then choose the κ appearing in (4.46) equal to some $\kappa_0(\rho_{\pm}) > 0$ independent of ϕ, δ such that, for N large enough depending on ϕ :

$$\kappa_0(\rho_{\pm}) \zeta_\delta^N (\|\phi\|_2^{-1} \phi) \text{ is of LS type.} \quad (4.69)$$

To conclude the estimate of the right-hand side of (4.46) using (4.68), it remains to bound the average of V_δ . Fix $\delta = 1/16$. Recall that $Y^{(0)}$ has 2-norm bounded by 4 and, by Cauchy-Schwarz inequality, that $Y^{(1)}$ has 2-norm bounded by $2^{-6} \delta^{-1} C = C/4$ and $Y^{(2)}$ by $\sqrt{2} \|\nabla g_0\|_2 / 8$. The $Y^{(i)}$

($0 \leq i \leq 2$) thus have 2-norm bounded independently of ϕ . Lemma 3.9 then implies that there is $\kappa'_0(\rho_\pm) > 0$ independent of ϕ and a numerical constant C such that $\kappa'_0(\rho_\pm)V_{1/16}$ is of LS type. Taking $\kappa(\rho_\pm) := \min\{\kappa_0(\rho_\pm), \kappa'_0(\rho_\pm)\}$, there are then constants $C(\rho_\pm), C(\rho_\pm, \phi) > 0$ such that, for each N large enough depending on ϕ :

$$\kappa(\rho_\pm)\nu_{g_0}^N \left(f[V_{1/16} + \zeta_{1/16}^N(\|\phi\|_2^{-1}\phi)] \right) \leq \frac{H(f\nu_{g_0}^N|\nu_{g_0}^N)}{2^9 C_{LS}} + C(\rho_\pm) + \frac{C(\rho_\pm, \phi)}{N^{1/2}} + \frac{N^2}{16}\nu_{g_0}^N(\Gamma(f^{1/2})). \quad (4.70)$$

Injecting this bound in (4.46) and using the log-Sobolev inequality of Lemma 3.4 concludes the proof:

$$\frac{1}{T} \log \mathbb{P}^{\nu_{g_0}^N} \left(\frac{1}{T} \int_0^T \|\phi\|_2^{-1} \Pi_t^N(\partial_1 \phi) dt > A \right) \leq -A\kappa(\rho_\pm) + C(\rho_\pm) + \frac{C(\rho_\pm, \phi)}{N^{1/2}}. \quad (4.71)$$

□

4.3 Exponential tightness

In this section, we prove that the upper bound (4.43) is also valid for closed sets. We refer to Appendix E for a characterisation of compact sets in $(\mathcal{T}'_s, *)$, and establish exponential tightness in the next proposition.

Proposition 4.4. *For each large enough $A > 0$, there is a compact set $\mathcal{K}_A \subset (\mathcal{T}'_s, *)$ such that:*

$$\sup_N \limsup_{T \rightarrow \infty} \frac{1}{T} \log \mathbb{P}^{\nu_{g_0}^N} \left(\frac{1}{T} \int_0^T \Pi_t^N dt \notin \mathcal{K}_A \right) \leq -A. \quad (4.72)$$

Proof. In Appendix E, we prove that there is a norm $\|\cdot\|_{\mathbb{T}, -2}$ on \mathcal{T}'_s such that the set $\{\|\cdot\|_{\mathbb{T}, -2} \leq B\}$ is relatively weak-* compact for each $B > 0$. The norm $\|\cdot\|_{\mathbb{T}, -2}$ involves a certain family of functions $\psi_m \in \mathcal{T}$ ($m \in \mathbb{N}^2$), and reads:

$$\|\Pi\|_{\mathbb{T}, -2}^2 := \sum_{m \in \mathbb{N}^2} \frac{1}{(1 + |m|^2)^2} |\Pi(\psi_m)|^2, \quad \Pi \in (\mathcal{T}'_s). \quad (4.73)$$

Above, $|m|^2 = m_1^2 + m_2^2$ for $m = (m_1, m_2) \in \mathbb{N}^2$. It is therefore enough to prove the existence of $c(\cdot) > 0$ with $\lim_{B \rightarrow \infty} c(B) = \infty$ such that:

$$\sup_{N \in \mathbb{N}^*} \mathbb{P}^{\nu_{g_0}^N} \left(\left\| \frac{1}{T} \int_0^T \Pi_t^N dt \right\|_{\mathbb{T}, -2} \geq A \right) \leq e^{-c(A)T}. \quad (4.74)$$

To prove (4.74), let $\varepsilon \in (0, 1)$ and define $c_\varepsilon := \sum_m (1 + |m|^2)^{1+\varepsilon} < \infty$. A union bound gives:

$$\mathbb{P}^{\nu_{g_0}^N} \left(\left\| \frac{1}{T} \int_0^T \Pi_t^N dt \right\|_{\mathbb{T}, -2} \geq A \right) \leq \sum_{m \in \mathbb{N}^2} \mathbb{P}^{\nu_{g_0}^N} \left(\left| \frac{1}{T} \int_0^T \Pi_t^N(\psi_m) dt \right| \geq \frac{A}{c_\varepsilon} (1 + |m|^2)^{\frac{1-\varepsilon}{2}} \right). \quad (4.75)$$

The ψ_m are just restrictions to \triangleright of the eigenvectors of the torus Laplacian on $[-2, 2]^2$ (see Appendix E). In particular, $\|\psi_m\|_\infty \leq C$ where C does not depend on N, m . Each probability in

the above sum is thus estimated by Corollary 3.14 according to:

$$\begin{aligned} \mathbb{P}^{\nu_{g_0}^N} \left(\left\| \frac{1}{T} \int_0^T \Pi_t^N dt \right\|_{\mathbb{T}, -2} \geq A \right) &\leq \sum_{m \in \mathbb{N}^2} \exp \left[- \frac{\zeta_2 A (1 + |m|^2)^{\frac{1-\varepsilon}{2}}}{c_\varepsilon^{1/2} \|\psi_m\|_\infty} T + C(\rho_\pm) T \right] \\ &\leq \sum_{m \in \mathbb{N}^2} \exp \left[- \frac{\zeta_2 A (1 + |m|^2)^{\frac{1-\varepsilon}{2}}}{c_\varepsilon^{1/2} C \|\psi_m\|_\infty} T + C(\rho_\pm) T \right], \end{aligned} \quad (4.76)$$

where the parameter ζ_2 is a universal constant defined in item 3 of Lemma 3.13. The expression in the right-hand side is summable for any $\varepsilon < 1$. It is moreover bounded by $e^{-c(A)T}$, with $c(A) > 0$ for A large enough and $\lim_{B \rightarrow \infty} c(B) = +\infty$. This completes the proof. \square

5 Lower bound for smooth trajectories

In this section, we give a lower bound on $\mathbb{Q}_T^{\pi_{inv}^N}(\mathcal{O})$ (defined in (4.2)) when \mathcal{O} is an open subset of $(\mathcal{T}'_s, *)$, in terms of the kernels $k_h, h \in \mathcal{S}(\varepsilon_B)$, with ε_B the quantity appearing in Theorems 2.3–2.6. As for standard large deviations (see Chapter 10 in [KL99]), we consider the tilted dynamics $\mathbb{P}_h, h \in \mathcal{S}(\varepsilon_B)$ such that $\frac{1}{4} \langle k_h, \cdot \rangle \in \mathcal{O}$, and obtain a lower bound by proving that the measure \mathbb{P}_h concentrates on \mathcal{O} . In the following, let $\mathbb{Q}_{T,h}$ denote the law of $\frac{1}{T} \int_0^T \Pi_t^N dt$ under \mathbb{P}_h .

We first change the initial measure to $\nu_{g_h}^N$, with g_h chosen according to (2.61), and tilt the dynamics by h . Using Jensen inequality to obtain the last line below, one finds, with $D_h = d\mathbb{P}_h/d\mathbb{P}$ (see (3.131)):

$$\begin{aligned} \log \mathbb{Q}_T^{\pi_{inv}^N}(\mathcal{O}) &= \log \mathbb{E}_h^{\nu_{g_h}^N} \left[\mathbf{1}_{\mathcal{O}} \frac{\pi_{inv}^N(\eta_0)}{\nu_{g_h}^N(\eta_0)} (D_h)^{-1} \right] \\ &= \log \mathbb{E}_{h, \mathcal{O}}^{\nu_{g_h}^N} \left[\frac{\pi_{inv}^N(\eta_0)}{\nu_{g_h}^N(\eta_0)} (D_h)^{-1} \right] + \log \mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O}) \\ &\geq \mathbb{E}_{h, \mathcal{O}}^{\nu_{g_h}^N} [- \log D_h] + \mathbb{E}_{h, \mathcal{O}}^{\nu_{g_h}^N} \left[\log \left(\frac{\pi_{inv}^N(\eta_0)}{\nu_{g_h}^N(\eta_0)} \right) \right] + \log \mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O}) \\ &\geq \mathbb{E}_{h, \mathcal{O}}^{\nu_{g_h}^N} [- \log D_h] - C(\rho_\pm, h)N + \log \mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O}). \end{aligned} \quad (5.1)$$

Above, $C(g, g_h) > 0$, while $\mathbb{P}_{h, \mathcal{O}}^{\nu_{g_h}^N}$ is the probability $\mathbb{P}_h^{\nu_{g_h}^N}$ conditional to $\left\{ \frac{1}{T} \int_0^T \Pi_t^N dt \in \mathcal{O} \right\}$:

$$\mathbb{P}_{h, \mathcal{O}}^{\nu_{g_h}^N}(\cdot) = \left(\mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O}) \right)^{-1} \mathbb{P}_h^{\nu_{g_h}^N} \left(\cdot \cap \left\{ \frac{1}{T} \int_0^T \Pi_t^N dt \in \mathcal{O} \right\} \right), \quad \mathbb{E}_{h, \mathcal{O}}^{\nu_{g_h}^N}[\cdot] = \int \cdot d\mathbb{P}_{h, \mathcal{O}}^{\nu_{g_h}^N}. \quad (5.2)$$

The terms appearing in (5.1) are of three types: the change of initial condition corresponding to the constant $-C(\rho_\pm, g_h)N$, which will vanish upon division by T when T is large; the dynamical part with $\log D_h$, and the term $\mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O})$. The latter is well controlled only if h is such that, under $\mathbb{Q}_{T,h}^{\nu_{g_h}^N}$, correlations are typically in \mathcal{O} when N, T are large. For such an h , upon dividing by T and taking the large T limit, only the dynamical part will contribute. The limit of $\mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O})$ is worked out in the next section, and the dynamical part is studied in Section 5.2.

5.1 Law of large numbers and Poisson equation

Proposition 5.1. *Let $h \in \mathcal{S}(\varepsilon_B)$, and let k_h be the large N limit of the correlations under $\nu_{g_h}^N$, where g_h solves the main equation (2.61). If $\mathcal{O} \subset (\mathcal{T}'_s, *)$ is an open set containing $\frac{1}{4}\langle k_h, \cdot \rangle$, then:*

$$\liminf_{N \rightarrow \infty} \liminf_{T \rightarrow \infty} \mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O}) = 1. \quad (5.3)$$

An open set in $(\mathcal{T}'_s, *)$ is a (possibly uncountable) union of finite intersections of sets of the form $\left\{ \left| \frac{1}{T} \int_0^T \Pi_t^N(\phi) dt - \frac{1}{4} \langle k_h, \phi \rangle \right| \in U \right\}$, for an open set $U \subset \mathbb{R}$ and $\phi \in \mathcal{T}$. It is therefore enough to prove (5.3) for those sets, with $U = (-\varepsilon, \varepsilon)$ for $\varepsilon > 0$.

Let us first check that it suffices to prove (5.3) for symmetric $\phi \in C^2(\bar{\square})$ with $\phi|_{\partial \square} = 0$. Recall the notation $\|\phi\|_{2,N}^2 = N^{-2} \sum_{i,j} \phi_{i,j}^2$. For each $a > 0$ and $\phi, \psi \in \mathcal{T}$, the entropy- and Markov inequalities together with the second item of Lemma 3.9 give, for some $C(h, \rho_{\pm}) > 0$:

$$\mathbb{P}_h^{\nu_{g_h}^N} \left(\left| \frac{1}{T} \int_0^T \Pi_t^N(\phi - \psi) dt \right| \geq a \right) \leq \frac{C(h, \rho_{\pm}) \|\phi - \psi\|_{2,N}}{a} = \frac{C(h, \rho_{\pm}) \|\phi - \psi\|_2}{a} + o_N(1). \quad (5.4)$$

In addition, $|\langle k_h, \phi - \psi \rangle| \leq \|k_h\|_2 \|\phi - \psi\|_2$. Since one can approximate $\phi \in \mathcal{T}$ in $\mathbb{L}^2(\square)$ with arbitrary precision by some $\psi \in \mathcal{T} \cap C^2(\bar{\square})$ with $\psi|_{\partial \square} = 0$, it is enough to focus on such ψ as claimed. This is done in the next proposition, by means of a Poisson problem associated with the large N limit of the generator $N^2 L_h$.

Proposition 5.2. *Let $\phi \in \mathcal{T} \cap C^2(\bar{\square})$ be a symmetric function with $\phi|_{\partial \square} = 0$, where we recall $\mathcal{T} = \mathbb{H}^2(\square)$. Let $h \in \mathcal{S}(\varepsilon_B)$. Then, for any $\varepsilon \in (0, 1)$ and any $T > 0$, there are positive constants $C(h, \phi), C'(h, \phi) > 0$ (independent of T) such that:*

$$\mathbb{P}_h^{\nu_{g_h}^N} \left(\left| \frac{1}{T} \int_0^T \Pi_t^N(\phi) dt - \frac{1}{4} \langle k_h, \phi \rangle \right| \geq \varepsilon \|\phi\|_2 \right) \leq \frac{C(h, \phi)}{T} (\varepsilon^{-2} + N\varepsilon^{-1}) + \frac{C'(h, \phi)}{\varepsilon^2 N^{1/2}}. \quad (5.5)$$

Proof. Fix a symmetric $\phi \in \mathcal{T} \cap C^2(\bar{\square})$. To prove Proposition 5.2, we express the difference appearing in the probability in (5.5) as a time integral involving the generator $N^2 L_h$, plus a martingale term. The martingale term is then proven to fluctuate like \sqrt{T} when N, T are large. It thus vanishes in the large T limit upon dividing by T . Recall that $\mathcal{M}(u, v)(x, y) = \int_{(-1,1)} u(z, x) \bar{\sigma}(z) v(z, y) dz$ for any $u, v \in \mathbb{L}^2(\square)$. The key ingredient is the following Poisson equation:

$$\begin{cases} \frac{1}{2} \Delta f(x, y) + \frac{1}{2} \mathcal{M}(\partial_1 f, \partial_1 h) + \frac{1}{2} \mathcal{M}(\partial_1 h, \partial_1 f) = \frac{\phi(x, y)}{\|\phi\|_2} & \text{for } (x, y) \in \square, \\ f = 0 & \text{on } \partial \square, \\ (\partial_1 - \partial_2) f(x_{\pm}, x) = 0 & \text{for } x \in (-1, 1). \end{cases} \quad (5.6)$$

For $\phi \in \mathcal{T} \cap C^2(\bar{\square})$ satisfying $\phi|_{\partial \square} = 0$, we prove in Appendix F that (5.6) has a unique solution $f_{\phi} \in \mathcal{T} \cap C^3(\bar{\square})$, a symmetric function on \square . It satisfies $\|f_{\phi}\|_{\infty} \leq C$ for a constant independent of ϕ , see Proposition F.5. The semi-martingale decomposition for $\Pi^N(f_{\phi})$ reads:

$$\forall T \geq 0, \quad \Pi_T^N(f_{\phi}) = \Pi_0^N(f_{\phi}) + \int_0^T N^2 L_h \Pi_t^N(f_{\phi}) dt + M_t^{N, f_{\phi}}. \quad (5.7)$$

Let us first use the Poisson equation (5.6) to express $N^2 L_h \Pi^N(f_{\phi})$ in terms of $\Pi^N(\phi)$.

Lemma 5.3.

$$N^2 L_h \Pi^N(f_\phi) = \frac{1}{\|\phi\|_2} \left(\Pi^N(\phi) - \frac{1}{4} \langle k_h, \phi \rangle \right) + \theta^N(f_\phi), \quad (5.8)$$

where $\theta^N(f_\phi)$ is an error term of size $N^{-1/2}$ of large type (recall Definitions 3.6–3.11 of error terms and of large type).

Assuming Lemma 5.3 for the moment, let us prove Proposition 5.2. For each $T > 0$ and $\varepsilon > 0$, integrate (5.8) between 0 and T and use the martingale decomposition (5.7) to find:

$$\begin{aligned} & \mathbb{P}_h^{\nu_{gh}^N} \left(\left| \frac{1}{T} \int_0^T \Pi_t^N(\phi) dt - \frac{1}{4} \langle k_h, \phi \rangle \right| \geq \varepsilon \|\phi\|_2 \right) \\ & \leq \mathbb{P}_h^{\nu_{gh}^N} \left(\frac{1}{T} \left| \Pi_T^N(f_\phi) - \Pi_0^N(f_\phi) - M_T^{N,f_\phi} - \int_0^T \theta_t^N(f_\phi) dt \right| > \varepsilon \right). \end{aligned} \quad (5.9)$$

Let us estimate each of the terms appearing in the last probability. Proposition 3.8 takes care of $\theta^N(f_\phi)$: There are $\gamma, C > 0$ depending on ρ_\pm, h , but not ϕ , such that for all $N \in \mathbb{N}^*$ and $T > 0$:

$$\mathbb{E}_h^{\nu_{gh}^N} \left[\left| \frac{1}{T} \int_0^T \theta_t^N(f_\phi) dt \right| \right] \leq \frac{1}{\gamma} \log \mathbb{E}_h^{\nu_{gh}^N} \left[\exp \left| \frac{\gamma}{T} \int_0^T \theta_t^N(f_\phi) dt \right| \right] \leq CN^{-1/2}. \quad (5.10)$$

By Markov- and Chebychev inequalities:

$$\begin{aligned} & \mathbb{P}_h^{\nu_{gh}^N} \left(\left| \frac{1}{T} \int_0^T \Pi_t^N(\phi) dt - \frac{1}{4} \langle k_h, \phi \rangle \right| \geq \varepsilon \|\phi\|_2 \right) \\ & \leq \frac{3C}{\varepsilon N^{1/2}} + \frac{3}{T\varepsilon} \mathbb{E}_h^{\nu_{gh}^N} \left[\left| \Pi_T^N(f_\phi) - \Pi_0^N(f_\phi) \right| \right] + \frac{9}{\varepsilon^2 T^2} \mathbb{E}_h^{\nu_{gh}^N} \left[\langle M^{N,f_\phi} \rangle_T \right]. \end{aligned} \quad (5.11)$$

The middle expectation involving the terms $\Pi_T^N(f_\phi), \Pi_0^N(f_\phi)$ is bounded by $\|f_\phi\|_\infty N$ (which does not depend on ϕ). Its contribution to (5.11) thus vanishes when T is large. Let us prove that the quadratic variation of M^{N,f_ϕ} has average bounded linearly in time for N large, which will be enough to conclude the proof of Proposition 5.2. The quadratic variation is given for each $t \geq 0$ by:

$$\langle M^{N,f_\phi} \rangle_t = \frac{1}{2} \int_0^t \sum_{i < N-1} c_h(\eta(s), i, i+1) \left[\frac{\bar{\rho}'}{4N} (f_\phi)_{i,i+1} + \frac{1}{2N} \sum_{j \notin \{i, i+1\}} \bar{\eta}_j(s) \partial_1^N (f_\phi)_{i,j} \right]^2 ds. \quad (5.12)$$

Recall from the definition (2.41) of c_h that $\sup_{N, \Omega_N} c_h \leq c(h)$ for some $c(h) > 0$. Using the inequality $(a+b)^2 \leq 2a^2 + 2b^2$ for $a, b \in \mathbb{R}$ and the fact that f_ϕ is bounded, one has:

$$\mathbb{E}_h^{\nu_{gh}^N} \left[\langle M^{N,f_\phi} \rangle_T \right] \leq c(h) \mathbb{E}_h^{\nu_{gh}^N} \left[\int_0^T \sum_{i < N-1} \left(\frac{1}{2N} \sum_{j \notin \{i, i+1\}} \bar{\eta}_j(t) \partial_1^N (f_\phi)_{i,j} \right)^2 dt \right] + \frac{Tc(h)\|f_\phi\|_\infty}{N}. \quad (5.13)$$

The integrand at each time $t \leq T$ is of the form $N^{-1} X_{2, \{0\}}^u$ in the notations of Lemma 3.9, with $u = (u_{i,j})_{i,j \in \Lambda_N}$ given by:

$$u_{i,j} = \frac{1}{N} \sum_{k \notin \{i-1, i, j-1, j, N-1\}} \partial_1^N (f_\phi)_{k,i} \partial_1^N (f_\phi)_{k,j} = \int_{(-1,1)} \partial_1 f_\phi(z, i/N) \partial_1 f_\phi(z, j/N) dz + \frac{v_{i,j}^N}{N}, \quad (5.14)$$

where v^N is a discretisation error bounded by $C(\phi) > 0$. $N^{-1}X_{2,\{0\}}^u$ is thus controllable with size 1, and Proposition 3.8 yields the existence of $C(h, \phi) > 0$, $C'(h, \phi) > 0$ such that:

$$\frac{9}{\varepsilon^2 T^2} \mathbb{E}_h^{\nu_{g_h}^N} \left[\langle M^{N, f_\phi} \rangle_T \right] \leq \frac{1}{\varepsilon^2} \left(\frac{C(h, \phi)}{T} + \frac{C(h, \phi)}{N} \right). \quad (5.15)$$

This estimate and (5.11) conclude the proof of Proposition 5.2 assuming Lemma 5.3, proven below. \square

Proof of Lemma 5.3. To prove Lemma 5.3, the starting point is the expression of $N^2 L_h \Pi^N(f_\phi)$ worked out in Corollary 3.20: there is a Γ -controllable error term $\tilde{\varepsilon}^N(h, f_\phi)$ with size $N^{-1/2}$, of large type, such that:

$$\begin{aligned} N^2 L_h \Pi^N(f_\phi) &= \frac{1}{2} \Pi^N \left(\Delta f_\phi + \mathcal{M}(\partial_1 f_\phi, \partial_1 h) + \mathcal{M}(\partial_1 h, \partial_1 f_\phi) \right) - \frac{(\bar{\rho}')^2}{4} \int_{(-1,1)} f_\phi(x, x) dx \\ &\quad + \frac{1}{4} \int_{\mathbb{Z}} \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 f_\phi(x, y) \partial_1 h(x, y)] dx dy + \tilde{\varepsilon}^N(h, f_\phi). \end{aligned} \quad (5.16)$$

Note the absence of the diagonal term $\sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} (\partial_1(f_\phi)_{i+,i} - \partial_1(f_\phi)_{i-,i})$: recall from Remark 3.19 that this term corresponds to the derivative of f_ϕ in the normal direction to the diagonal, which vanishes according to (5.6). By Proposition 3.8, there are $\gamma, C > 0$ depending on h, ϕ with:

$$\mathbb{E}_h^{\nu_{g_h}^N} \left[\left| \int_0^T \tilde{\varepsilon}_t^N(h, f_\phi) dt \right| \right] \leq \frac{1}{\gamma} \log \mathbb{E}_h^{\nu_{g_h}^N} \left[\exp \left| \gamma \int_0^T \tilde{\varepsilon}_t^N(h, f_\phi) dt \right| \right] \leq C(h, \phi) N^{-1/2}. \quad (5.17)$$

As f_ϕ solves (5.6), (5.16) can be written as:

$$\begin{aligned} N^2 L_h \Pi^N(f_\phi) &= \frac{1}{\|\phi\|_2} \Pi^N(\phi) - \frac{(\bar{\rho}')^2}{4} \int_{(-1,1)} f_\phi(x, x) dx \\ &\quad + \frac{1}{4} \int_{\mathbb{Z}} \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 f_\phi(x, y) \partial_1 h(x, y)] dx dy + \tilde{\varepsilon}^N(h, f_\phi). \end{aligned} \quad (5.18)$$

Equation (5.18) will correspond to (5.8) with $\theta^N(f_\phi) = \tilde{\varepsilon}^N(h, f_\phi)$, if we can prove:

$$-\frac{1}{4\|\phi\|_2} \langle k_h, \phi \rangle = -\frac{(\bar{\rho}')^2}{4} \int_{(-1,1)} f_\phi(x, x) dx + \frac{1}{4} \int_{\mathbb{Z}} \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 f_\phi(x, y) \partial_1 h(x, y)] dx dy. \quad (5.19)$$

By definition, $k_h + \bar{\sigma} = (\bar{\sigma}^{-1} - g_h)^{-1}$ with g_h solving the main equation (2.61). We prove in Appendix F that this is equivalent to saying that k_h solves the Euler-Lagrange equation (F.5). The identity (5.19) is then straightforwardly obtained by taking f_ϕ as a test function in the weak formulation (F.4) of the Euler-Lagrange equation, then integrating by parts. \square

5.2 Estimation of the dynamical part and conclusion of the lower bound

Let $h \in \mathcal{S}(\varepsilon_B)$. In this section, we estimate the term $\mathbb{E}_{h, \mathcal{O}}^{\nu_{g_h}^N} [-\log D_h]$ arising in (5.1), and prove:

$$\liminf_{N \rightarrow \infty} \liminf_{T \rightarrow \infty} \frac{1}{T} \mathbb{E}_{h, \mathcal{O}}^{\nu_{g_h}^N} [-\log D_h] = -J_h(k_h) = -\mathcal{I}_{\varepsilon_B}(k_h) = \mathcal{I}_\infty(k_h), \quad (5.20)$$

where \mathcal{I}_∞ is defined in (2.50) and $\mathcal{I}_{\varepsilon_B}(k_h)$ is short for $\mathcal{I}_{\varepsilon_B}(\frac{1}{4}\langle k_h, \cdot \rangle)$, idem for $\mathcal{I}_\infty(k_h)$ and $J_h(k_h)$ below. This concludes the proof of the lower bound for regular kernels close to that of the steady state.

Recall from Corollary 3.20 the definition of the martingale $M_T^{N,h}$. Then $\mathbb{E}_{h,\mathcal{O}}^{\nu_{g_h}^N}[-\log D_h]$ reads:

$$\begin{aligned} \mathbb{E}_{h,\mathcal{O}}^{\nu_{g_h}^N}[-\log D_h] &= \left(\mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O}) \right)^{-1} \mathbb{E}_h^{\nu_{g_h}^N} \left[\mathbf{1}_{\mathcal{O}} \left(-M_T^{N,h} - \frac{1}{2} \int_0^T \Pi_t^N(\mathcal{M}(\partial_1 h, \partial_1 h)) dt \right) \right] \\ &\quad - \frac{T}{8} \int_{\square} \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 h(x, y)]^2 dx dy + \frac{\mathbb{E}_h^{\nu_{g_h}^N} \left[\mathbf{1}_{\mathcal{O}} \int_0^T \hat{\varepsilon}^N(h)(\eta_t) dt \right]}{\mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O})}, \end{aligned} \quad (5.21)$$

with $\hat{\varepsilon}^N(h)$ the error term defined in Corollary 3.20. Proposition 5.1 establishes the convergence of $\mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O})$ to 1 as T , then N become large. The error term $\hat{\varepsilon}^N(h)$ is controlled through Proposition 3.8:

$$\limsup_{N \rightarrow \infty} \sup_{T > 0} \mathbb{E}_h^{\nu_{g_h}^N} \left[\left| \frac{1}{T} \int_0^T \hat{\varepsilon}^N(h)(\eta_t) dt \right| \right] = 0. \quad (5.22)$$

In particular,

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \left| \frac{\mathbb{E}_h^{\nu_{g_h}^N} \left[\mathbf{1}_{\mathcal{O}} \int_0^T \hat{\varepsilon}^N(h)(\eta_t) dt \right]}{\mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O})} \right| = 0. \quad (5.23)$$

Consider now the expectation in the first line of (5.21). By Cauchy-Schwarz inequality and the fact that $\langle M^{N,h} \rangle_T$ grows at most linearly in T up to a small error vanishing with N (see (5.15)), the contribution of the martingale term vanishes:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \frac{\mathbb{E}_h^{\nu_{g_h}^N} [\mathbf{1}_{\mathcal{O}} |M_T^{N,h}|]}{T \mathbb{Q}_{T,h}^{\nu_{g_h}^N}(\mathcal{O})} = 0. \quad (5.24)$$

Consider now the last remaining term in (5.21), the term involving $\Pi^N(\mathcal{M}(\partial_1 h, \partial_1 h))$. To estimate it, we state the following moment bound, proven afterwards.

Lemma 5.4. *For any continuous and bounded $F : \square \rightarrow \mathbb{R}$, there is $C(h, \|F\|_\infty) > 0$ such that:*

$$\sup_{N \in \mathbb{N}^*} \sup_{t \geq 0} \mathbb{E}_h^{\nu_{g_h}^N} [|\Pi_t^N(F)|^{3/2}] \leq C(h, \|F\|_\infty), \quad (5.25)$$

Proposition 5.2 and the moment bound of Lemma 5.4 directly yield:

$$\limsup_{N \rightarrow \infty} \limsup_{T \rightarrow \infty} \left| \mathbb{E}_h^{\nu_{g_h}^N} \left[\mathbf{1}_{\mathcal{O}} \frac{1}{T} \int_0^T \Pi_t^N(\mathcal{M}(\partial_1 h, \partial_1 h)) dt \right] - \frac{1}{4} \langle k_h, \mathcal{M}(\partial_1 h, \partial_1 h) \rangle \right| = 0. \quad (5.26)$$

From equations (5.23) to (5.26), we obtain:

$$\begin{aligned} \liminf_{N \rightarrow \infty} \liminf_{T \rightarrow \infty} \frac{1}{T} \mathbb{E}_{h,\mathcal{O}}^{\nu_{g_h}^N}[-\log D_h] &= -\frac{1}{8} \int_{\square} dx dy \bar{\sigma}(x) \bar{\sigma}(y) [\partial_1 h(x, y)]^2 - \frac{1}{8} \langle k_h, \mathcal{M}(\partial_1 h, \partial_1 h) \rangle \\ &= -\frac{1}{8} \int_{(-1,1)} dz \bar{\sigma}(z) \langle \partial_1 h(z, \cdot), (\bar{\sigma} + k_h) \partial_1 h(z, \cdot) \rangle. \end{aligned} \quad (5.27)$$

Let us check that the right-hand side in (5.27) is indeed equal to $-\mathcal{I}_{\varepsilon_B}(k_h)$, and that $\mathcal{I}_\infty(k_h) = \mathcal{I}_{\varepsilon_B}(k_h)$. The weak formulation (2.52) of the Euler-Lagrange equation on k_h with test function h gives:

$$J_h(k_h) = \frac{1}{8} \int_{(-1,1)} dz \bar{\sigma}(z) \langle \partial_1 h(z, \cdot), (\bar{\sigma} + k_h) \partial_1 h(z, \cdot) \rangle. \quad (5.28)$$

Again using (2.52) for test functions $f \in \mathcal{S}(\varepsilon_B)$, the rate function $\mathcal{I}_{\varepsilon'}(k_h) := \sup_{f \in \mathcal{S}(\varepsilon')} J_f(k_h)$ reads, for any $\varepsilon' \geq \varepsilon_B$:

$$\begin{aligned} \mathcal{I}_{\varepsilon'}(k_h) &= \sup_{f \in \mathcal{S}(\varepsilon')} \left\{ \frac{1}{4} \int_{(-1,1)} dz \bar{\sigma}(z) \langle \partial_1 f(z, \cdot), (\bar{\sigma} + k_h) \partial_1 h(z, \cdot) \rangle \right. \\ &\quad \left. - \frac{1}{8} \int_{(-1,1)} dz \bar{\sigma}(z) \langle \partial_1 f(z, \cdot), (\bar{\sigma} + k_h) \partial_1 f(z, \cdot) \rangle \right\} \\ &= J_h(k_h) - \sup_{f \in \mathcal{S}(\varepsilon')} \left\{ \frac{1}{8} \int_{(-1,1)} dz \bar{\sigma}(z) \langle \partial_1 (f - h)(z, \cdot), (\bar{\sigma} + k_h) \partial_1 (f - h)(z, \cdot) \rangle \right\} = J_h(k_h). \end{aligned} \quad (5.29)$$

This in particular yields $J_h(k_h) = \mathcal{I}_{\varepsilon_B}(k_h) = \mathcal{I}_\infty(k_h)$ and, together with (5.27)–(5.28), proves (5.20).

Proof of Lemma 5.4. Let $F : \square \rightarrow \mathbb{R}$ be bounded. Let $t \geq 0$, and let $\varepsilon > 0$ that will eventually be chosen as $\varepsilon = \frac{1}{2}$. The moment bound is obtained by a careful application of the entropy inequality, putting to good use the $O(N^{-1/2})$ estimate on the size of the relative entropy of Theorem 2.6. Note that $|\Pi_t^N(F)| \leq \|F\|_\infty N$. As a result, fixing $c > 0$ to be chosen later and applying the entropy inequality to $c\lambda \mathbf{1}_{|\Pi^N(F)| > \lambda}$ for each $\lambda > 1$ in the second line below:

$$\begin{aligned} \mathbb{E}_h^{\nu_{g_h}^N} [|\Pi_t^N(F)|^{1+\varepsilon}] &\leq 1 + (1 + \varepsilon) \int_1^{\|F\|_\infty N} \lambda^\varepsilon \mathbb{P}_{g_h}^{\nu_{g_h}^N} (|\Pi_t^N(F)| > \lambda) d\lambda \\ &\leq 1 + (1 + \varepsilon) \int_1^{\|F\|_\infty N} c^{-1} \lambda^{-1+\varepsilon} \left[H(f_t \nu_{g_h}^N | \nu_{g_h}^N) + \log \left(1 + (e^{c\lambda} - 1) \nu_{g_h}^N (|\Pi^N(F)| > \lambda) \right) \right] d\lambda. \end{aligned} \quad (5.30)$$

By Theorem 2.6, $H(f_t \nu_{g_h}^N | \nu_{g_h}^N) \leq CN^{-1/2}$ for some $C = C(h, \rho_\pm) > 0$. Moreover, by Corollary A.4, the probability involving $\nu_{g_h}^N$ above is bounded by $C(h, \|F\|_\infty) e^{-c(h, \|F\|_\infty)\lambda}$ for some $C(h, \|F\|_\infty), c(h, \|F\|_\infty) > 0$. Choosing $c = c(h, \|F\|_\infty)/2$, one obtains the existence of $C'(h, \|F\|_\infty) > 0$ such that:

$$\begin{aligned} \mathbb{E}_h^{\nu_{g_h}^N} [|\Pi_t^N(F)|^{1+\varepsilon}] &\leq 1 + \frac{(1 + \varepsilon)C'(h, \|F\|_\infty)}{\varepsilon N^{1/2-\varepsilon}} \\ &\quad + \frac{2(1 + \varepsilon)}{c(h, \|F\|_\infty)} \int_1^{\|F\|_\infty N} \lambda^{-1+\varepsilon} e^{-c(h, \|F\|_\infty)\lambda/2} d\lambda. \end{aligned} \quad (5.31)$$

The integral is bounded with N whatever the choice of $\varepsilon > 0$. Overall, the right-hand side above is therefore bounded with N as soon as $\varepsilon \in (0, 1/2]$, which yields (5.25). This concludes the proof of the lemma. \square

5.3 Towards a lower bound for non regular correlations

In this section, we discuss how to extend the lower bound to non regular correlations, i.e. to kernels $k \in \mathbb{H}^1(\square)$ which solve the Euler-Lagrange equation (2.52) associated with a possibly non regular

bias h . As we shall see below, this extension is much simpler here than in the case of large deviations of the density for the open SSEP [Ber+03]. This is due to the fact that the Euler-Lagrange equation (5.32) is linear as an equation with fixed h and unknown k , and the smallness assumption on admissible biases h .

Let us sketch an informal proof of this extension. Take a symmetric $k \in \mathbb{H}^1(\square)$ with $k|_{\partial\square} = 0$ and recall that the Euler-Lagrange equation (2.52), viewed as an equation with unknown h , reads for all test functions $\phi \in \mathbb{H}^1(\square)$ with $\phi|_{\partial\square} = 0$:

$$\frac{1}{2} \int_{\square} \nabla(k - k_0)(x, y) \cdot \nabla\phi(x, y) dx dy + \frac{1}{2} \int_{(-1,1)} \bar{\sigma}(z) \langle \partial_1 h(z, \cdot), C_k \partial_1 \phi(z, \cdot) \rangle dz = 0. \quad (5.32)$$

For k close to k_0 , one has $\bar{\sigma} + k = (\bar{\sigma} + k_0) + (k - k_0) \geq \alpha \text{id}$ for some $\alpha > 0$, and the second term is a coercive bilinear form for the norm $\|\nabla \cdot\|_2$. There is thus a unique symmetric $h \in \mathbb{H}^1(\square)$ with $h|_{\partial\square} = 0$ associated with k . Using (5.32) with test function h , this h has small norm (recall $\bar{\sigma} \leq 1/4$ and $\|\nabla h\|_2^2 = 2\|\partial_1 h\|_2^2$ as h is symmetric):

$$\frac{\alpha}{2} \|\nabla h\|_2^2 \leq \alpha \int_{\square} \bar{\sigma}(z) \partial_1 h(x, z)^2 dx dz \leq \|\nabla h\|_2 \|\nabla(k - k_0)\|_2. \quad (5.33)$$

One would expect the estimate (5.33) to be enough to approximate h in $\|\nabla \cdot\|_2$ norm by more regular biases. However, to avoid technical difficulties, our set $\mathcal{S}(\varepsilon_B)$ of regular biases was defined with an assumption on the sup norm of derivatives rather than on $\|\nabla \cdot\|_2$. This assumption is technical and could be lifted at the cost of lengthening some arguments. Assuming therefore that one can take $(h_n) \in \mathcal{S}(\varepsilon_B)^{\mathbb{N}}$ with $\lim_n \|\nabla(h_n - h)\|_2 = 0$, let k_n denote the kernel associated with h_n ($n \in \mathbb{N}$) through (5.32). Equation (5.32), written for k and k_n with test function $k - k_n$, and the bounds $C_k = \bar{\sigma} + k \leq c \text{id}$, $\bar{\sigma} \leq 1/4$ and $\|\nabla h_n\|_2^2 = 2\|\partial_1 h_n\|_2^2$ give:

$$\|\nabla(k - k_n)\|_2^2 \leq c \|\nabla(h - h_n)\|_2 \|\nabla(k - k_n)\|_2 + \frac{1}{4\sqrt{2}} \|k - k_n\|_2 \|\nabla(k - k_n)\|_2 \|\nabla h_n\|_2. \quad (5.34)$$

By assumption k is close to k_0 . This means that $\|\nabla h\|_2$ is small by (5.33). Since $(\|\nabla h_n\|_2)_n$ converges to $\|\nabla h\|_2$, $\|\nabla h_n\|_2$ is also small for large enough n . Using Poincaré inequality on \square with 0 Dirichlet boundary conditions on $\partial\square$ to bound $\|k - k_n\|_2$ by $(2\pi)^{-1} \|\nabla(k - k_n)\|_2$, we thus find:

$$\|\nabla(k - k_n)\|_2 \left(1 - \frac{\|\nabla h_n\|_2}{8\pi\sqrt{2}}\right) \leq c \|\nabla(h - h_n)\|_2. \quad (5.35)$$

Assuming that k is close enough to k_0 for the parenthesis on the left-hand side to be, say, larger than $1/2$ when n is large enough; we conclude that $\lim_n \|\nabla(k_n - k)\|_2 = 0$. This yields the alleged general lower bound:

$$\mathcal{I}_{\varepsilon_B}(k_n) = \frac{1}{8} \int_{(-1,1)} dz \bar{\sigma}(z) \langle \partial_1 h_n(z, \cdot), (\bar{\sigma} + k_n) \partial_1 h_n(z, \cdot) \rangle \xrightarrow{n \rightarrow \infty} \mathcal{I}_{\varepsilon_B}(k). \quad (5.36)$$

A Correlations and concentration under discrete Gaussian measures

In this section, we investigate the measures ν_g^N , defined for $g : \square \rightarrow \mathbb{R}$ as follows:

$$\bar{\nu}^N = \bigotimes_{i \in \Lambda_N} \text{Ber}(\bar{\rho}_i), \quad \forall \eta \in \Omega_N, \quad \nu_g^N(\eta) = \frac{1}{\mathcal{Z}_g^N} \exp \left[\frac{1}{2N} \sum_{i \neq j \in \Lambda_N} g_{i,j} \bar{\eta}_i \bar{\eta}_j \right] \bar{\nu}^N(\eta), \quad (\text{A.1})$$

where the partition function \mathcal{Z}_g^N is a normalising constant. For simplicity, we assume throughout that $g \in \mathbb{L}^2(\square)$ is a negative kernel, i.e.:

$$\forall u \in \mathbb{L}^2((-1, 1)), \quad \int_{\square} u(x)g(x, y)u(y)dx dy \leq 0. \quad (\text{A.2})$$

None of the result in Appendix A would be modified if we instead assumed that $g = g_- + g_+$ where g_+ is a positive kernel, and $\|g_+\|_2 \leq c$ for a small enough constant c .

A.1 Bound on the partition function and correlations

Lemma A.1 (Bound on the partition function). *Let $g : \square \rightarrow \mathbb{R}$ be a continuous, bounded and symmetric function, and assume that g is a negative kernel (see (A.2)). Then:*

$$\sup_{N \geq 1} \mathcal{Z}_g^N < \infty, \quad \mathcal{Z}_g^N := \bar{\nu}^N \left[\exp \left[\frac{1}{2N} \sum_{i \neq j \in \Lambda_N} \bar{\eta}_i \bar{\eta}_j g_{i,j} \right] \right]. \quad (\text{A.3})$$

Proof. As g is continuous and a negative kernel, the matrix $(g_{i,j})_{(i,j) \in \Lambda_N^2}$ has negative eigenvalues, and one has:

$$\forall \eta \in \Omega_N, \quad 2\Pi^N(g) + \frac{1}{2N} \sum_{i \in \Lambda_N} (\bar{\eta}_i)^2 g_{i,i} = \frac{1}{2N} \sum_{(i,j) \in \Lambda_N^2} \bar{\eta}_i \bar{\eta}_j g_{i,j} \leq 0. \quad (\text{A.4})$$

The continuity of g on the diagonal and the bound $|\bar{\eta}_i| \leq 1$ imply:

$$\begin{aligned} \mathcal{Z}_g^N &= \bar{\nu}^N \left[\exp \left[\frac{1}{2N} \sum_{(i,j) \in \Lambda_N^2} \bar{\eta}_i \bar{\eta}_j g_{i,j} - \frac{1}{2N} \sum_{i \in \Lambda_N} (\bar{\eta}_i)^2 g_{i,i} \right] \right] \\ &\leq e^{\|g\|_\infty}. \end{aligned} \quad (\text{A.5})$$

□

Lemma A.2. *Let $g : \square \rightarrow \mathbb{R}$ satisfy the hypotheses of Lemma A.1. Then:*

$$\forall n \in \mathbb{N}^*, \quad \sup_{I \subset \Lambda_N : |I|=n} \left| \nu_g^N \left(\prod_{a \in I} \bar{\eta}_a \right) \right| = O(N^{-n/2}). \quad (\text{A.6})$$

Moreover, if $g \in C^1(\bar{\triangleright}) \cap C^1(\bar{\triangleleft})$ with $g(\pm 1, \cdot) = 0$, then, for each $n \in \mathbb{N}^*$ and each $\varepsilon \in \{+, -\}$:

$$\sup_{\substack{I \subset \Lambda_N \\ \varepsilon(N-1) \in I \text{ and } |I|=n}} \left| \nu_g^N \left(\bar{\eta}_{\varepsilon(N-1)} \prod_{a \in I \setminus \{\varepsilon(N-1)\}} \bar{\eta}_a \right) \right| = O(N^{-n/2-1}). \quad (\text{A.7})$$

Proof. Let $I_n \subset \Lambda_N$ with $|I| = n \in \mathbb{N}^*$. The proof relies on a development of the exponential defining ν_g^N in (A.1), and the observation that the measure $\bar{\nu}$ is product, so that, for each $p \in \mathbb{N}^*$,

$$\forall (i_1, \dots, i_p) \in \Lambda_N^p, \quad |\{i_1, \dots, i_p\}| > p/2 \quad \Rightarrow \quad \bar{\nu}^N \left(\prod_{j=1}^p \bar{\eta}_{i_j} \right) = 0. \quad (\text{A.8})$$

In other words, in a product $\bar{\eta}_{i_1} \dots \bar{\eta}_{i_n}$, each $\bar{\eta}$ must be paired with at least another $\bar{\eta}$ with the same index i_1, \dots, i_n . Since $\bar{\eta}$ is bounded, a direct consequence is that, if $1 \leq p \leq n$ and $w^p = (w_i^p)_{i \in \Lambda_N}$ satisfies $\sup_{\Lambda_N} |w^p| = O_N(1)$, then:

$$\bar{\nu}^N \left[\prod_{p=1}^n \left(\frac{1}{N} \sum_{j \in \Lambda_N} w_j^p \bar{\eta}_j \right) \right] = O(N^{-\lceil n/2 \rceil}). \quad (\text{A.9})$$

Write $I_p = \{i_1, \dots, i_p\}$ for each $1 \leq p \leq n$ (and by extension write $I_{i_p} := I_p$). Define, for $J \subset \Lambda_N$:

$$G_{J^c}(\eta) := \frac{1}{4N} \sum_{i \neq j \in \Lambda_N \setminus \{J\}} \bar{\eta}_i \bar{\eta}_j g_{i,j}, \quad \eta \in \Omega_N. \quad (\text{A.10})$$

Then, for each $\eta \in \Omega_N$:

$$\begin{aligned} 2\Pi^N(g) &= 2G_{I_1^c}(\eta) + \bar{\eta}_{i_1} \sum_{j \neq i_1} \bar{\eta}_j g_{i_1,j} \\ &= 2G_{I_p^c}(\eta) + \frac{1}{N} \sum_{q=1}^p \bar{\eta}_{i_q} \sum_{j \in \Lambda_N \setminus I_q} \bar{\eta}_j g_{i_q,j}, \quad 1 \leq p \leq n. \end{aligned} \quad (\text{A.11})$$

As a result, one can write:

$$\mathcal{Z}_g^N \nu_g^N \left(\prod_{p=1}^n \bar{\eta}_{i_p} \right) = \bar{\nu}^N \left[\left(\prod_{p=1}^n \bar{\eta}_{i_p} \right) e^{2G_{I_n^c}(\eta)} \exp \left[\frac{1}{N} \sum_{p=1}^n \bar{\eta}_{i_p} \sum_{j \in \Lambda_N \setminus I_p} \bar{\eta}_j g_{i_p,j} \right] \right], \quad (\text{A.12})$$

and $G_{I_n^c}$ does not involve any $\bar{\eta}_{i_p}$ ($1 \leq p \leq n$). Note also that $\sup_N \mathcal{Z}_g^N < \infty$ by Lemma A.1, thus establishing Lemma A.2 only boils down to estimating the right-hand side above. To do so, recall first the following identity:

$$\forall x \in \mathbb{R}, \quad e^x = \sum_{p=0}^{n-1} \frac{x^p}{p!} + \int_0^1 \frac{x^m}{(m-1)!} (1-t)^m e^{tx} dt. \quad (\text{A.13})$$

Expanding the second exponential in (A.12) to order n thus yields, writing $n_0 := \lfloor n/2 \rfloor$:

$$\mathcal{Z}_g^N \nu_g^N \left(\prod_{p=1}^n \bar{\eta}_{i_p} \right) = \bar{\nu}^N \left[\left(\prod_{p=1}^n \bar{\eta}_{i_p} \right) e^{2G_{I_n^c}(\eta)} \left\{ T_1(\eta) + T_2(\eta) + T_3(\eta) \right\} \right], \quad (\text{A.14})$$

where:

$$T_1(\eta) := \sum_{m=0}^{n_0-1} \frac{1}{m!} \left(\frac{1}{N} \sum_{q=1}^n \bar{\eta}_{i_q} \sum_{j \in \Lambda_N \setminus I_q} \bar{\eta}_j g_{i_q, j} \right)^m, \quad (\text{A.15})$$

$$T_2(\eta) := \sum_{m=n_0}^{n-1} \frac{1}{m!} \left(\frac{1}{N} \sum_{q=1}^n \bar{\eta}_{i_q} \sum_{j \in \Lambda_N \setminus I_q} \bar{\eta}_j g_{i_q, j} \right)^m, \quad (\text{A.16})$$

$$T_3(\eta) := \int_0^1 \frac{(1-t)^{n-1}}{(n-1)!} \left(\frac{1}{N} \sum_{q=1}^n \bar{\eta}_{i_q} \sum_{j \in \Lambda_N \setminus I_q} \bar{\eta}_j g_{i_q, j} \right)^n \exp \left[\frac{t}{N} \sum_{q=1}^n \bar{\eta}_{i_q} \sum_{j \in \Lambda_N \setminus I_q} \bar{\eta}_j g_{i_q, j} \right] dt. \quad (\text{A.17})$$

Let us estimate the average of T_1, T_2, T_3 separately. Consider first T_1 in (A.15). For each $0 \leq m \leq n_0 - 1$, developing the term elevated to the power m yields products of at most $2m = 2n_0 - 2 < n$ $\bar{\eta}$'s. In view of Equation (A.8), this is not enough to achieve a pairing of all indices in $\prod_{p=1}^n \bar{\eta}_{i_p}$, thus the average of T_1 vanishes identically.

Let us count the factors of $\bar{\eta}$'s appearing in T_2 in (A.16). For each $n_0 \leq m < n$, the term to the power m in (A.16) is a sum of terms of the form:

$$\prod_{i \in K} \left(\frac{\bar{\eta}_i}{N} \sum_{j \in \Lambda_N \setminus I_i} \bar{\eta}_j g_{i, j} \right)^{n_i}, \quad K \subset I_n, (n_i)_{i \in K} \in \mathbb{N}^{|K|} \text{ with } 1 \leq |K| \leq m, \sum_{i \in K} n_i = m. \quad (\text{A.18})$$

Again by (A.8), each term appearing in (A.18) gives a non-vanishing contribution to the average (A.14) provided each element i_1, \dots, i_n of I_n appears at least once in the expansion. Since the term (A.18) contains the product $\prod_{i \in K} \bar{\eta}_i$ with $|K| < n$, $n - |K|$ indices j of the sums in (A.18) must be fixed to an element of $I_n \setminus K$. A contribution $1/N$ arises each time an index j is singled out in a sum. It follows that the contribution of T_2 to the average (A.14) is bounded by a sum of terms of the form:

$$\sum_{m=n_0}^{n-1} \frac{C(m)}{N^{n-|K|}} \left(\prod_{i \in I_n \setminus K} \|g_{i, \cdot}\|_{\infty} \right) \bar{\nu}^N \left[e^{2G_{I_n^c}(\eta)} \prod_{i \in K} \left| \frac{1}{N} \sum_{j \in \Lambda_N \setminus I_i} \bar{\eta}_j g_{i, j} \right|^{n'_i} \right], \quad C(m) > 0, \quad (\text{A.19})$$

where $K \subset I_n$, and the n'_i satisfy:

$$0 \leq n'_i \leq n_i \quad \text{for } i \in K, \quad \text{and} \quad \sum_{i \in K} n'_i = (m + |K| - n) \mathbf{1}_{m+|K|-n \geq 0}. \quad (\text{A.20})$$

Let us now estimate the average in (A.19). Recall that g is continuous, and a negative kernel (as defined in (A.2)). As a result, $(g_{i, j})_{i, j \in \Lambda_N \setminus I_n}$ has negative eigenvalues, and for each $\eta \in \Omega_N$:

$$\begin{aligned} 2G_{I_n^c}(\eta) &= \frac{1}{2N} \sum_{i, j \in \Lambda_N \setminus I_n} \bar{\eta}_i \bar{\eta}_j g_{i, j} - \frac{1}{2N} \sum_{i \in \Lambda_N \setminus I_n} (\bar{\eta}_i)^2 g_{i, i} \\ &\leq -\frac{1}{2N} \sum_{i \in \Lambda_N \setminus I_n} (\bar{\eta}_i)^2 g_{i, i} \leq \|g\|_{\infty}. \end{aligned} \quad (\text{A.21})$$

Using then Cauchy-Schwarz inequality followed by the estimate (A.9) for moments of correlations under the product measure $\bar{\nu}^N$, the average in (A.19) is bounded by:

$$\begin{aligned} \bar{\nu}^N \left[e^{2G_{I_n^c}(\eta)} \prod_{i \in K} \left| \frac{1}{N} \sum_{j \in \Lambda_N \setminus K} \bar{\eta}_j g_{i,j} \right|^{n'_i} \right] &\leq e^{\|g\|_\infty} \bar{\nu}^N \left[\prod_{i \in K} \left(\frac{1}{N} \sum_{j \in \Lambda_N \setminus K} \bar{\eta}_j g_{i,j} \right)^{2n'_i} \right]^{1/2} \\ &= \mathbf{1}_{m+|K|-n \geq 0} O(N^{-(m+|K|-n)/2}). \end{aligned} \quad (\text{A.22})$$

Putting together (A.19)-(A.22) and summing on all possible choices for K , $(n_q)_{q \in K}$, $(n'_q)_{q \in K}$ yields:

$$\begin{aligned} \left| \bar{\nu}^N \left[e^{2G_{I_n^c}(\eta)} \left(\prod_{q \in I} \bar{\eta}_{i_q} \right) \cdot T_2(\eta) \right] \right| &= \sum_{m=n_0}^{n-1} \sum_{K: 1 \leq |K| \leq m} \sum_{(n_i), (n'_i)} O(N^{-(n-|K|)}) O(N^{-(m+|K|-n)/2}) \\ &= O(N^{-n/2}). \end{aligned} \quad (\text{A.23})$$

It remains to estimate the contribution of the average of T_3 , defined in (A.17). Since there are already n sums involving $\bar{\eta}$, it is enough to use the following bound, valid for each $\eta \in \Omega_N$:

$$\left| \left(\prod_{p=1}^n \bar{\eta}_{i_p} \right) T_3(\eta) \right| \leq \sum_{q=1}^n \left| \frac{1}{N} \sum_{j \in \Lambda_n \setminus I_q} \bar{\eta}_j g_{i_q, j} \right|^n \int_0^1 \frac{(1-t)^{n-1}}{(n-1)!} \exp \left[\frac{t}{N} \sum_{q=1}^n \bar{\eta}_{i_q} \sum_{j \in \Lambda_N \setminus I_q} \bar{\eta}_j g_{i_q, j} \right] dt. \quad (\text{A.24})$$

The integral on t is taken care of with the following identity: for each $t \geq 0$ and $\eta \in \Omega_N$,

$$2G_{I_n^c}(\eta) + \frac{t}{N} \sum_{q=1}^n \bar{\eta}_{i_q} \sum_{j \in \Lambda_N \setminus I_q} \bar{\eta}_j g_{i_q, j} = 2\Pi^N(g) + \frac{(t-1)}{N} \sum_{q=1}^n \bar{\eta}_{i_q} \sum_{j \in \Lambda_N \setminus I_q} \bar{\eta}_j g_{i_q, j}. \quad (\text{A.25})$$

Indeed, since g is a negative kernel, it follows that:

$$\begin{aligned} \left| \bar{\nu}^N \left[e^{2G_{I_n^c}(\eta)} \left(\prod_{q=1}^n \bar{\eta}_{i_q} \right) T_3 \right] \right| &\leq \frac{e^{(n+1)\|g\|_\infty}}{n!} \bar{\nu}^N \left[\left| \frac{1}{N} \sum_{q=1}^n \bar{\eta}_{i_q} \sum_{j \in \Lambda_N \setminus I_q} \bar{\eta}_j g_{i_q, j} \right|^n \right] \\ &= O(N^{-n/2}). \end{aligned} \quad (\text{A.26})$$

Putting (A.23)-(A.26) together yields the first part of Lemma A.2, i.e. (A.6).

Let us now prove the improved estimate (A.7) if correlations include points at the boundaries. If g is assumed to be C^1 with $g(\pm 1, \cdot) = 0$, then $\|g_{\pm(N-1), \cdot}\|_\infty = O(N^{-1})$. By (A.8), the index $\pm(N-1)$ must arise in the terms (A.15)-(A.17) to compensate the $\bar{\eta}_{\pm(N-1)}$ present in the product $\prod_{q \in I} \bar{\eta}_{i_q} \pm(N-1)$. It follows that the $O(N^{-n/2})$ bound in (A.23)-(A.26) is improved to $O(N^{-n/2-1})$, which concludes the proof of Lemma A.2. \square

A.2 Exponential moments of higher order correlations

Let $g \in g_0 + \mathcal{S}(\infty)$ be a negative kernel, where this set is defined in (2.45). In this section, we give bounds on the size of exponential moments, under ν_g^N , of random variables involving n -point correlations, $n \geq 1$. These are useful when applying the entropy inequality.

Such concentration results are established in the literature by means of a logarithmic Sobolev inequality, see [GSS+19]-[SS20]. That g be a negative kernel implies that, for any $F : \Omega_N \rightarrow \mathbb{R}$:

$$\nu_g^N(\exp[F]) \leq e^{\|g\|_\infty} (\mathcal{Z}_g^N)^{-1} \bar{\nu}^N(\exp[F]), \quad (\text{A.27})$$

and it is enough to estimate exponential moments under the product measure $\bar{\nu}^N$. To do so, let us fix some notations. For $d \in \mathbb{N}$, let $A : \Lambda_N^d \rightarrow \mathbb{R}$ be a tensor. Define its Hilbert-Schmidt norm by:

$$\|A\|_{HS} = \left(\sum_{(i_0, \dots, i_{d-1}) \in \Lambda_N^d} (A(i_0, \dots, i_{d-1}))^2 \right)^{1/2}. \quad (\text{A.28})$$

For $J \subset \mathbb{Z}$ containing 0, let $X_{d,J}^A$ be defined as:

$$\forall \eta \in \Omega_N, \quad X_{d,J}^A(\eta) = \sum_{\substack{(i_0, \dots, i_{d-1}) \in \Lambda_N^d \\ i_0 + J \subset \Lambda_N}} A(i_0, \dots, i_{d-1}) \bar{\eta}_{i_0+J} \prod_{p=1}^{d-1} \bar{\eta}_{i_p}, \quad \bar{\eta}_{i_0+J} := \prod_{j \in J} \bar{\eta}_{i_0+j}. \quad (\text{A.29})$$

The next theorem gives concentration estimates of $X_{d,J}^A$ under $\bar{\nu}^N$ for $J \subset \mathbb{Z}$ and $d \in \mathbb{N}^*$. The case $J = \{0\}$, $d \leq 4$ corresponds to Theorem 1.4 in [GSS+19], but their proof extends to any $d \in \mathbb{N}^*$. However, we also use in the article the case $J = \{0, 1\}$, for which a proof is needed.

Theorem A.3. *Let $J \subset \mathbb{Z}$ contain 0 and $d \in \mathbb{N}^*$. Assume that A is such that $A(i_0, \dots, i_{d-1})$ vanishes whenever the same site appears twice in $\bar{\eta}_{i_0+J} \prod_{p=1}^{d-1} \bar{\eta}_{i_p}$ for $\eta \in \Omega_N$, i.e. assume:*

$$\forall (i_0, \dots, i_{d-1}) \in \Lambda_N^d, \quad \left(\exists j \in J, |\{i_0 + j, i_1, \dots, i_{d-1}\}| < d \right) \Rightarrow A(i_0, \dots, i_{d-1}) = 0. \quad (\text{A.30})$$

There are then constants $c_d > 0$ depending only on d such that, for any $c \in (0, c_d)$ and any N with $J \subset \Lambda_N$:

$$\bar{\nu}^N \left(\exp \left[\frac{c |X_{d,J}^A|^{2/d}}{\|A\|_{HS}^{2/d}} \right] \right) \leq 2. \quad (\text{A.31})$$

Proof. The proof for general J and the $J = \{0\}$ case in [GSS+19] are very similar, so we only give a sketch. Without loss of generality, A can be assumed to be invariant under permutation of its last $d-1$ indices. The idea is to proceed by recursion on d , noticing that, for each $\ell \in \Lambda_N$ and $\eta \in \Omega_N$:

$$\begin{aligned} \nabla_\ell X_{d,J}^A(\eta) &:= X_{d,J}^A(\eta^\ell) - X_{d,J}^A(\eta) \\ &= (1 - 2\eta_\ell) \left[\sum_{i_0: i_0+J \ni \ell} A(i_0, \dots, i_{d-1}) \bar{\eta}_{(i_0+J) \setminus \{\ell\}} \prod_{a=1}^{d-1} \bar{\eta}_{i_a} \right. \\ &\quad \left. + (d-1) A(i_0, \dots, i_{d-2}, \ell) \bar{\eta}_{i_0+J} \prod_{a=1}^{d-2} \bar{\eta}_{i_a} \right]. \end{aligned} \quad (\text{A.32})$$

Fix $d \in \mathbb{N}$, $J \subset \mathbb{Z}$ with $0 \in J$, and N with $J \subset \Lambda_N$. For brevity, simply write X_d for $X_{d,J}^A$.

Step 1: reduction to moment bound. To prove (A.31), it is enough to prove the existence of $C_d > 0$ such that:

$$\forall p \in \mathbb{N}^*, \quad M_p(X_d) := \bar{\nu}^N [|X_d|^p]^{1/p} \leq C_d \|A\|_{HS} p^{\frac{d}{2}}. \quad (\text{A.33})$$

Indeed, assuming such a bound, one has, for each $c > 0$, using Jensen inequality when $d \geq 2$ for the convex function $f(x) = x^{d/2}$, $x \in \mathbb{R}$:

$$\bar{\nu}^N [\exp(c|X_d|^{2/d})] = 1 + \sum_{p=1}^{\infty} \frac{c^p}{p!} \bar{\nu}^N (|X_d|^{2p/d}) \leq 1 + \sum_{p=1}^{\infty} \frac{(cC_d p \|A\|_{HS}^{2/d})^p}{p!}. \quad (\text{A.34})$$

As $p^p \leq p!e^p$ for each $p \geq 1$, taking $c \leq c_d := (2C_d e)^{-1}$ yields (A.31).

Step 2: moment estimate. It is enough to prove (A.33) for $p \geq 2$, since the first moment can be estimated by Cauchy-Schwarz inequality. We will restrict to $p \geq 2$ at some point in the computation. For now, we treat p as a continuous variable in \mathbb{R}_+^* and differentiate $M_p(X_d)$. For each $p > 0$, one has:

$$\begin{aligned} \frac{dM_p(X_d)}{dp} &= \frac{d}{dp} \left(\exp \left[\frac{1}{p} \log \bar{\nu}^N (|X_d|^p) \right] \right) \\ &= - \frac{\log \bar{\nu}^N (|X_d|^p)}{p^2} M_p(X_d) + \frac{1}{p} \frac{\bar{\nu}^N (|X_d|^p \log |X_d|)}{\bar{\nu}^N (|X_d|^p)} M_p(X_d), \end{aligned} \quad (\text{A.35})$$

which can be written as:

$$\forall p > 0, \quad \frac{dM_p(X_d)}{dp} = \frac{M_p(X_d)^{1-p}}{p^2} \text{Ent}(|X_d|^p), \quad (\text{A.36})$$

with $\text{Ent}(F^2)$ the entropy of F^2 against $\bar{\nu}^N$, given by:

$$\forall F : \Omega_N \rightarrow \mathbb{R}, \quad \text{Ent}(F^2) = \bar{\nu}^N (F^2 \log F^2) - \bar{\nu}^N (F^2) \log \bar{\nu}^N (F^2). \quad (\text{A.37})$$

The entropy on the right-hand side of (A.36) is estimated by means of a logarithmic Sobolev inequality, satisfied by $\bar{\nu}^N$ for the Glauber dynamics on Ω_N (see e.g. Theorem A.1. in [DS96]): there is $C_{LS} > 0$, independent of N , such that:

$$\forall F : \Omega_N \rightarrow \mathbb{R}_+, \quad \text{Ent}(F^2) \leq C_{LS} \bar{\nu}^N \left(\sum_{i \in \Lambda_N} [\nabla_i F]^2 \right), \quad (\text{A.38})$$

where, for $i \in \Lambda_N$ and $\eta \in \Omega_N$, $\nabla_i F(\eta) = F(\eta^i) - F(\eta)$ for each $F : \Omega_N \rightarrow \mathbb{R}$. Similarly, a Poincare inequality holds with constant $C_{LS}/2$:

$$\forall F : \Omega_N \rightarrow \mathbb{R}_+, \quad \bar{\nu}^N (F^2) - \bar{\nu}^N (F)^2 \leq \frac{C_{LS}}{2} \bar{\nu}^N \left(\sum_{i \in \Lambda_N} [\nabla_i F]^2 \right). \quad (\text{A.39})$$

Injecting (A.38) in (A.36) and proceeding as in [GSS+19], one successively obtains, restricting to $p > 2$:

$$\begin{aligned} \forall p > 2, \quad \frac{dM_p(X_d)^2}{dp} &\leq \frac{2C_{LS} M_p(X_d)^{2-p}}{p^2} \bar{\nu}^N \left(\sum_{i \in \Lambda_N} [\nabla_i (|X_d|^{p/2})]^2 \right) \\ &\leq C_{LS} M_p(X_d)^{2-p} \bar{\nu}^N \left(|X_d|^{p-2} \sum_{i \in \Lambda_N} [\nabla_i |X_d|]^2 \right). \end{aligned} \quad (\text{A.40})$$

Applying Hölder inequality with exponents $(p/(p-2), p/2)$ then yields, for each $p > 2$:

$$\frac{dM_p(X_d)^2}{dp} \leq C_{LS} M_{p/2} \left(\sum_{i \in \Lambda_N} [\nabla_i |X_d|]^2 \right) \leq C_{LS} M_{p/2} \left(\sum_{i \in \Lambda_N} [\nabla_i X_d]^2 \right). \quad (\text{A.41})$$

The function $M_p(X_d)$ is increasing for $p > 0$ by (A.36). As a result, integrating between 2 and p and using the Poincaré inequality (A.39) to estimate $M_2(X_d)$ yields:

$$\begin{aligned} \forall p \geq 2, \quad M_p(X_d)^2 &\leq M_2(X_d)^2 + C_{LS}(p-2)M_{p/2} \left(\sum_{i \in \Lambda_N} [\nabla_i X_d]^2 \right) \\ &\leq C_{LS} p M_{p/2} \left(\sum_{i \in \Lambda_N} [\nabla_i X_d]^2 \right). \end{aligned} \quad (\text{A.42})$$

Step 3: recursion on d . Let $p \geq 2$. For $j \in \Lambda_N$ and $0 \leq a \leq d-1$, define $A^{(i_a=j)}$ as the $d-1$ -tensor $(A(i_0, \dots, i_{a-1}, j, i_{a+1}, \dots, i_{d-1}))_{(i_q)_{q \neq a}}$, and note that $A^{(i_a=j)}$ also satisfies the assumption (A.30). Recall that X_d was short for $X_{d,J}^A$. Let us prove by recursion on d :

$$\forall p \geq 2, \quad M_p(X_{d,J}^A)^2 \leq 2^{d-1} (C_{LS} p)^d (d!)^2 |J|^2 \|A\|_{HS}^2. \quad (\text{A.43})$$

Proving such a result would conclude the proof. In the $d=1$ case, (A.42) yields, for each $p \geq 2$:

$$M_p(X_{1,J}^A)^2 \leq C_{LS} p M_{p/2} \left(\sum_{i \in \Lambda_N} [\nabla_i X_{1,J}^A]^2 \right) \leq C_{LS} p \sum_{i \in \Lambda_N} M_p(\nabla_i X_{1,J}^A)^2. \quad (\text{A.44})$$

By (A.32), bounding $\bar{\eta}$ by 1, the result for $d=1$ is proven:

$$\forall p \geq 2, \quad M_p(X_{1,J}^A)^2 \leq C_{LS} p \sum_{\ell \in \Lambda_N} \left[\sum_{i_0+J \ni \ell} A(i_0) \right]^2 \leq C_{LS} p |J|^2 \|A\|_{HS}^2. \quad (\text{A.45})$$

For $d \geq 2$, (A.42) similarly gives, bounding $\bar{\eta}$ by 1:

$$\begin{aligned} M_{p/2} \left(\sum_{\ell \in \Lambda_N} [\nabla_\ell X_{d,J}^A]^2 \right) &\leq \sum_{\ell \in \Lambda_N} M_p(\nabla_\ell X_{d,J}^A)^2 \\ &\leq \sum_{\ell \in \Lambda_N} \left[(d-1) M_p(X_{d-1,J}^{A^{(i_{d-1}=\ell)}}) + \sum_{i:i+J \ni \ell} M_p(X_{d-1,\{0\}}^{A^{(i_0=i)}}) \right]^2. \end{aligned} \quad (\text{A.46})$$

For each $\ell \in \Lambda_N^d$, the recursion hypothesis at rank $d-1$ applied to $X_{d-1,J}^{A^{(i_{d-1}=\ell)}}$, and to $X_{d-1,\{0\}}^{A^{(i_0=i)}}$ for each $i+J \ni \ell$, yields:

$$\begin{aligned} M_{p/2} \left(\sum_{\ell \in \Lambda_N} [\nabla_\ell X_{d,J}^A]^2 \right) &\leq 2^{d-2} (C_{LS} p)^{d-1} ((d-1)!)^2 \sum_{\ell \in \Lambda_N} \left[|J|(d-1) \|A^{(i_{d-1}=\ell)}\|_{HS} + \sum_{i:i+J \ni \ell} \|A^{(i_0=i)}\|_{HS} \right]^2 \\ &\leq 2^{d-1} (C_{LS} p)^{d-1} ((d-1)!)^2 d^2 |J|^2 \|A\|_{HS}^2 = 2^{d-1} (C_{LS} p)^{d-1} (d!)^2 |J|^2 \|A\|_{HS}^2, \end{aligned} \quad (\text{A.47})$$

where we used $(a+b)^2 \leq 2a^2 + 2b^2$ for $a, b \in \mathbb{R}$ and bounded $(d-1)^2 + 1$ by d^2 to obtain the second inequality. Injecting this bound in (A.42) yields (A.43) at rank d , concluding the proof. \square

In the next corollary, we use Theorem A.3 to establish the controllability results of Lemma 3.9 on the variables $X_{d,J}^A$.

Corollary A.4. *Let $J \subset \mathbb{Z}$ contain 0 and, for $N \in \mathbb{N}^*$ and $d \in \mathbb{N}^*$, let $A = A(d, N)$ be a d -tensor. We do not assume that A satisfy (A.30), but instead that $\sup_{N \in \mathbb{N}^*} \|A\|_\infty < \infty$. There are constants $\gamma_d, C_d(g) > 0$ that are independent of A , with $\gamma_1 = +\infty$, such that, for each $\gamma \leq \gamma_d$:*

$$\begin{aligned} & \text{for each } N \in \mathbb{N}^* \text{ with } J \subset \Lambda_N, \\ & \begin{cases} \log \nu_g^N \left(\exp \left[\frac{\gamma}{\|A\|_\infty} N^{-1/2} X_{0,J}^A \right] \right) \leq C_0(g) \gamma^2 & \text{if } d = 1, \\ \log \nu_g^N \left(\exp \left[\frac{\gamma}{\|A\|_\infty} N^{-(d-1)} X_{d,J}^A \right] \right) \leq \frac{C_d(g) \gamma}{N^{\frac{d-2}{2}}} & \text{if } d \geq 2. \end{cases} \end{aligned} \quad (\text{A.48})$$

If instead A is assumed to satisfy (A.30), then (A.48) is valid with $\|A\|_\infty$ replaced by $\|A\|_{2,N} := \|N^{-d/2} A\|_{HS} \leq \|A\|_\infty$ everywhere.

Proof. Fix $J \subset \mathbb{Z}$ with $0 \in J$, and N such that $J \subset \Lambda_N$. We will use Theorem A.3 to obtain (A.48). Let us first explain why it is not necessary to assume the condition (A.30) on the tensor A if it is bounded. Assume that Corollary A.4 is proven for tensors satisfying the condition (A.30), and assume that A is bounded but does not satisfy the condition. Then we claim that $N^{-(d-1)} X_{d,J}^A$ can be written as $N^{-(d-1)} X_{d,J}^{\tilde{A}}$ with \tilde{A} satisfying condition (A.30), plus a sum of terms of the form $N^{-(d-1)} X_{d',J'}^B$ for bounded d' -tensors B , $1 \leq d' \leq d-1$, and $J' \subset J$. A recursion then gives Corollary A.4 also in the case of bounded tensors that do not satisfy (A.30) provided the corollary holds under (A.30). The existence of the decomposition of $X_{d,J}^A$ is not difficult to see: it is enough to sum separately on the indices i_0, \dots, i_{d-1} in $X_{d,J}^A$ for which the same site appears twice in the collection $I_J = (i_0 + J, i_1, \dots, i_{d-1})$, and use the formula $\bar{\eta}_i^2 = \bar{\sigma}_i + \bar{\eta}_i \sigma'(\bar{\rho}_i)$ ($i \in \Lambda_N$).

We thus prove Corollary A.4 for tensors A that satisfy (A.30).

We first obtain estimates of tail probabilities of $X_d := X_{d,J}^A$ under ν_g^N . These are then used to obtain (A.48). By Theorem A.3, there is $\zeta_d > 0$ such that, for any $\zeta \in (0, \zeta_d)$:

$$\nu_g^N \left(\exp \left[\frac{\zeta |X_d|^{2/d}}{\|A\|_{HS}^{2/d}} \right] \right) \leq \frac{e^{\|g\|_\infty}}{\mathcal{Z}_g^N} \bar{\nu}^N \left(\exp \left[\frac{\zeta |X_d|^{2/d}}{\|A\|_{HS}^{2/d}} \right] \right) \leq 2e^{\|g\|_\infty} (\mathcal{Z}_g^N)^{-1}, \quad (\text{A.49})$$

Define:

$$\|A\|_{2,N} := \left(\frac{1}{N^d} \sum_{i_0, \dots, i_{d-1}} A(i_0, \dots, i_{d-1})^2 \right)^{1/2} = \|N^{-d/2} A\|_{HS}. \quad (\text{A.50})$$

By assumption on A ,

$$\sup_{N \in \mathbb{N}^*} \|A\|_{2,N} < \infty. \quad (\text{A.51})$$

For ζ as in (A.49), one has then, for each $t \geq 0$:

$$\nu_g^N \left(|X_d| > t \right) \leq 2(\mathcal{Z}_g^N)^{-1} e^{\|g\|_\infty} \exp \left[- \frac{\zeta t^{2/d}}{\|A\|_{HS}^{2/d}} \right] \leq 2(\mathcal{Z}_g^N)^{-1} e^{\|g\|_\infty} \exp \left[- \frac{\zeta t^{2/d}}{N \|A\|_{2,N}^{2/d}} \right]. \quad (\text{A.52})$$

At this point the proof is the same for each $d \geq 1$. We focus on the $d \geq 2$ case. Let $\gamma > 0$ and write:

$$\nu_g^N \left[\exp \left(\frac{\gamma}{\|A\|_{2,N}} N^{-(d-1)} |X_d| \right) \right] = 1 + \int_0^\infty e^t \nu_g^N \left(\frac{\gamma}{\|A\|_{2,N}} N^{-(d-1)} |X_d| > t \right) dt. \quad (\text{A.53})$$

Note that $N^{-(d-1)} X_d$ is bounded by $C \|A\|_{2,N} N$ for some numerical constant $C > 0$ and each $N \in \mathbb{N}^*$. Thus, with ζ given by (A.52):

$$\begin{aligned} \nu_g^N \left(\exp \left[\frac{\gamma}{\|A\|_{2,N}} N^{-(d-1)} |X_d| \right] \right) &\leq 1 + 2 (\mathcal{Z}_g^N)^{-1} e^{\|g\|_\infty} \int_0^{\gamma CN} \exp \left[t - \frac{\zeta t^{2/d} N^{(d-2)/d}}{\gamma^{2/d}} \right] dt \\ &=: 1 + 2 (\mathcal{Z}_g^N)^{-1} e^{\|g\|_\infty} \int_0^{\gamma CN} q_\gamma(t) dt. \end{aligned} \quad (\text{A.54})$$

If γ is small enough, we claim that the negative part of the exponential is dominant. Indeed, one has:

$$\begin{aligned} \forall t \leq \gamma CN, \quad q_\gamma(t) &\leq \exp \left[- \frac{\zeta t^{2/d} N^{(d-2)/d}}{2\gamma^{2/d}} \right] \\ \Leftrightarrow \quad \gamma &\leq \gamma_d, \quad \gamma_d = \gamma_d(\zeta) := \frac{C^{(2-d)/d} \zeta}{2}. \end{aligned} \quad (\text{A.55})$$

For any $\gamma < \gamma_d$, one has then:

$$\nu_g^N \left(\exp \left[\frac{\gamma}{\|A\|_{2,N}} N^{-(d-1)} |X_d| \right] \right) \leq 1 + 2 (\mathcal{Z}_g^N)^{-1} e^{\|g\|_\infty} \int_0^\infty \exp \left[- \frac{\zeta t^{2/d} N^{(d-2)/d}}{2\gamma^{2/d}} \right] dt. \quad (\text{A.56})$$

The change of variable $u = t \zeta^{d/2} \gamma^{-1} N^{(d-2)/2}$ and the boundedness of $(\mathcal{Z}_g^N)_N$ (see Lemma A.1) conclude the proof of Corollary A.4 for $d \geq 2$: for each $\gamma < \gamma_d$,

$$\nu_g^N \left(\exp \left[\frac{\gamma}{\|A\|_{2,N}} N^{-(d-1)} |X_d| \right] \right) \leq 1 + \frac{2\gamma\zeta^{-d/2}}{N^{(d-2)/2}} (\mathcal{Z}_g^N)^{-1} e^{\|g\|_\infty} \int_0^\infty \exp \left[- \frac{u^{2/d}}{2} \right] du. \quad (\text{A.57})$$

□

A.3 Log-Sobolev inequality under ν_g^N

Proposition A.5. *Let $0 \leq \rho_- \leq \rho_+ < 1$ and $h, g : \square \rightarrow \mathbb{R}$ be bounded with $\|g\|_\infty < 1/4$. There is $C_{LS}(\rho_\pm, \|g\|_\infty, \|h\|_\infty) > 0$ such that, for each density f for ν_g^N ,*

$$H(f \nu_g^N | \nu_g^N) \leq C_{LS} N^2 \nu_g^N(\Gamma_h(\sqrt{f})). \quad (\text{A.58})$$

If in particular $h \in \mathcal{S}(\varepsilon)$ and $g \in g_0 + \mathcal{S}(\varepsilon')$ with $\varepsilon, \varepsilon' \in (0, 1/4)$, then C_{LS} can be taken to depend only on ρ_\pm .

Proof. We shall see that $C_{LS}(\rho_\pm, \|g\|_\infty, \|h\|_\infty)$ can be taken to be an increasing function of $\|g\|_\infty, \|h\|_\infty$, which proves the second claim.

As $\|h\|_\infty < \infty$, the jump rates c_h satisfy:

$$\begin{aligned} c(\eta, i) &\leq e^{\|h\|_\infty} c(\eta, i), & i \in \{\pm(N-1)\}, \\ c(\eta, j, j+1) &\leq e^{\|h\|_\infty} c_h(\eta, j, j+1), & j < N-1. \end{aligned} \quad (\text{A.59})$$

It is therefore enough to prove the proposition for $h = 0$. By assumption $\|g\|_\infty < 1/4$, thus $\|g\|_2 \leq 2\|g\|_\infty$, thus the difference λ_g^N between largest and smallest eigenvalues of the matrix $(N^{-1}g_{i,j})_{i \neq j}$ is strictly below 1 for large enough N . By Theorem 1 in [BB19], the following log-Sobolev inequality for the Glauber dynamics associated with ν_g^N holds: for each density f with respect to ν_g^N ,

$$H(f\nu_g^N | \nu_g^N) \leq \frac{1}{2} \left(1 + \frac{2\lambda_g^N}{1 - \lambda_g^N}\right) \sum_{i \in \Lambda_N} \nu_g^N \left([\nabla_i \sqrt{f(\eta)}]^2 \right). \quad (\text{A.60})$$

The claim of Proposition A.5 is then an immediate consequence of (A.60) and the following bound: for some $c(\rho_\pm, \|g\|_\infty) > 0$,

$$\forall i \in \Lambda_N, \quad \nu_g^N \left([\nabla_i \sqrt{f(\eta)}]^2 \right) \leq c(\rho_\pm, \|g\|_\infty) N \nu_g^N (\Gamma(\sqrt{f})). \quad (\text{A.61})$$

Let us now prove (A.61). It is the claim of Lemma 4.2 in [Gon+21] when $g \equiv 0$. When $g \neq 0$, the proof is very similar and we only explain what changes. To prove (A.61), the idea is that changing the occupation number at a site i requires one to bring the particle or hole at i all the way to a reservoir (say the one on the right), perform an exchange, then bring back to i the new hole/particle. This is expressed rigorously by the following formula, which gives a recursive description of the above procedure: for any $F : \Omega_N \rightarrow \mathbb{R}$,

$$\nabla_i F(\eta) = \nabla_{i,i+1} F(\eta) + \nabla_{i+1} F(\eta^{i,i+1}) + \nabla_{i,i+1} F((\eta^{i,i+1})^{i+1}). \quad (\text{A.62})$$

Using the identity $(x+y+z)^2 \leq 2(1+\beta)(x^2+y^2) + (1+\beta^{-1})z^2$, valid for any $\beta > 0$ and $x, y, z \in \mathbb{R}$, one finds, after changes of variables:

$$\begin{aligned} \nu_g^N \left([\nabla_i F(\eta)]^2 \right) &\leq 2(1+\beta) \sum_{\eta \in \Omega_N} \nu_g^N(\eta) [\nabla_{i,i+1} F(\eta)]^2 \left[1 + \frac{\nu_g^N((\eta^i)^{i,i+1})}{\nu_g^N(\eta)} \right] \\ &\quad + (1+\beta^{-1}) \sum_{\eta \in \Omega_N} \nu_g^N(\eta) (\nabla_{i+1} F(\eta))^2 \frac{\nu_g^N(\eta^{i,i+1})}{\nu_g^N(\eta)}. \end{aligned} \quad (\text{A.63})$$

The computation of the above ratios is the only place where g plays a role. Since $\|g\|_\infty < \infty$, one can check that, for each $\eta \in \Omega_N$ (recall that $\bar{\nu}^N$ is the product Bernoulli measure defined in (2.58)):

$$\frac{\nu_g^N((\eta^i)^{i,i+1})}{\nu_g^N(\eta)} \leq \frac{\bar{\nu}^N((\eta^i)^{i,i+1})}{\bar{\nu}^N(\eta)} e^{2\|g\|_\infty}, \quad \frac{\nu_g^N(\eta^{i,i+1})}{\nu_g^N(\eta)} \leq \frac{\bar{\nu}^N(\eta^{i,i+1})}{\bar{\nu}^N(\eta)} e^{\|g\|_\infty}. \quad (\text{A.64})$$

The rest of the proof is then identical to that of Lemma 4.2 in [Gon+21]. In particular,

$$\frac{\bar{\nu}^N((\eta^i)^{i,i+1})}{\bar{\nu}^N(\eta)} \leq \frac{1}{\min\{1 - \rho_+, \rho_-\}} - 1, \quad \frac{\bar{\nu}^N(\eta^{i,i+1})}{\bar{\nu}^N(\eta)} \leq 1 + \frac{\bar{\rho}'}{\min\{1 - \rho_+, \rho_-\}^2 N}. \quad (\text{A.65})$$

This yields the following bound:

$$\begin{aligned} \nu_g^N \left([\nabla_i F(\eta)]^2 \right) &\leq 2(1 + \beta) \left[1 + \left(\frac{1}{\min\{1 - \rho_+, \rho_-\}} - 1 \right) e^{2\|g\|_\infty} \right] \nu_g^N \left([\nabla_{i,i+1} F(\eta)]^2 \right) \\ &\quad + (1 + \beta^{-1}) \left(1 + \frac{\bar{\rho}'}{\min\{1 - \rho_+, \rho_-\}^2 N} e^{\|g\|_\infty} \right) \nu_g^N \left([\nabla_{i+1} F(\eta)]^2 \right). \end{aligned} \quad (\text{A.66})$$

Iterating the bound with the choice $\beta = N$ concludes the proof of (A.61), thus of Proposition A.5. \square

B Integration by parts formulae

Fix $h \in \mathcal{S}(\infty)$ and $g \in g_0 + \mathcal{S}(\infty)$ (the set $\mathcal{S}(\infty)$ is defined in (2.45)) throughout. In this section, we provide integration by parts formulas under the measure ν_g^N , both in the bulk and close to the reservoirs. These formulas are key to the renormalisation procedure used to estimate error terms, in Section C. In particular, they are essential to proving the Γ -controllability of the variables $U_0^\pm, U_1^\pm, X_{2,\{0,1\}}^{\phi_2}$ encountered in Lemma 3.9.

B.1 Integration by parts in the bulk

Before stating the result, let us give some notations and explain what we mean by an integration by parts formula. Fix a density $f : \Omega_N \rightarrow \mathbb{R}$ for ν_g^N . For $i < N - 1$, let $\Gamma_h^{i,i+1}$ be defined as:

$$\forall \eta \in \Omega_N, \quad \Gamma_h^{i,i+1}(\sqrt{f})(\eta) = \frac{1}{2} c_h(\eta, i, i+1) [\nabla_{i,i+1} \sqrt{f}(\eta)]^2, \quad (\text{B.1})$$

with, for any $u : \Omega_N \rightarrow \mathbb{R}$ and any $i < N - 1$:

$$\forall \eta \in \Omega_N, \quad \nabla_{i,i+1} u(\eta) = u(\eta^{i,i+1}) - u(\eta). \quad (\text{B.2})$$

The jump rates c_h are defined in (2.41).

Consider a family $(\omega_i)_{i \in \Lambda_N}$ of functions on Ω_N . To estimate certain error terms in the adjoint $L_h^* \mathbf{1}$ in Section 3.2, a renormalisation scheme is used in Section C.1 below. For each density f for ν_g^N , some $\varepsilon > 0$ and $i \leq N - \varepsilon N$, this scheme consists in estimating the replacement:

$$\begin{aligned} \nu_g^N \left(f \left[\omega_i - \frac{1}{\varepsilon N} \sum_{j=i}^{i+\varepsilon N-1} \omega_j \right] \right) \\ = \nu_g^N \left(f \sum_{j=i}^{i+\varepsilon N-2} \frac{\varepsilon N - 1 - (j-i)}{\varepsilon N} [\omega_j - \omega_{j+1}] \right) \end{aligned} \quad (\text{B.3})$$

in terms of the entropy $H(f \nu_g^N | \nu_g^N)$ and the averaged carré du champ $\nu_g^N(\Gamma_h(f^{1/2}))$. The right-hand side in the last equation is obtained via a simple resummation. The key issue, then, is to understand how the space gradient $\omega_j - \omega_{j+1}$ ($j < N - 1$) can be turned into difference $f(\eta') - f(\eta)$, for a transition $\eta \rightarrow \eta'$ allowed by the dynamics (i.e. one of the differences contained in $\Gamma_h(f^{1/2})$).

Solving this issue is what we mean by finding an integration by parts formula, typically a formula of the form:

$$\sum_{\eta \in \Omega_N} (\omega_{i+1} - \omega_i) f(\eta) \nu_g^N(\eta) = \sum_{\eta \in \Omega_N} q(\eta) [f(\eta^{i,i+1}) - f(\eta)] \nu_g^N(\eta) + \nu_g^N(fX), \quad (\text{B.4})$$

where q, X are explicit functions, and the average of X can be estimated via the entropy inequality or another integration by parts formula. The natural choice for ω in our case is $\omega_i = \bar{\eta}_i$, however a simpler formula (B.4) is obtained, following [JM18], through the choice:

$$\forall i \in \Lambda_N, \quad \omega_i = \frac{\bar{\eta}_i}{\bar{\sigma}_i}, \quad \bar{\eta}_i = \eta_i - \bar{\rho}_i, \quad \bar{\sigma}_i = \bar{\rho}_i(1 - \bar{\rho}_i). \quad (\text{B.5})$$

Lemma B.1. *Let f be a ν_g^N -density. Fix $i < N - 1$ and let $u : \Omega_N \rightarrow \mathbb{R}$ be such that $\nabla_{i,i+1}u = 0$. Then:*

$$\begin{aligned} \nu_g^N [u(\omega_{i+1} - \omega_i)f] &= \nu_g^N (uq \nabla_{i,i+1}f) \\ &\quad - (\bar{\rho}_{i+1} - \bar{\rho}_i) \nu_g^N \left[\omega_i \omega_{i+1} e^{-(\eta_{i+1} - \eta_i)C_i^g/N} u f \right] \\ &\quad + \nu_g^N \left[(\omega_{i+1} - \omega_i) \left(1 - e^{-(\eta_{i+1} - \eta_i)C_i^g/N} \right) u f \right], \end{aligned} \quad (\text{B.6})$$

where the function $q = q_{i,i+1}(\eta)$ is given by:

$$\forall \eta \in \Omega_N, \quad q(\eta) = \frac{\eta_i(1 - \eta_{i+1})}{\bar{\rho}_i(1 - \bar{\rho}_{i+1})} e^{-(\eta_{i+1} - \eta_i)C_i^g/N}. \quad (\text{B.7})$$

Recall that, for each $i < N - 1$, $N^{-1}(\eta_i - \eta_{i+1})C_i^g = \nabla_{i,i+1}\Pi^N(g)$ is defined in (3.53), and satisfies $\max_i |C_i^g| \leq 2\|g\|_\infty$.

Proof. Let $i < N - 1$ and $q : \Omega_N \rightarrow \mathbb{R}$. Notice that, by definition of $\nabla_{i,i+1}$:

$$\nu_g^N [uq \nabla_{i,i+1}f] = \sum_{\Omega_N} uq \nabla_{i,i+1}f \nu_g^N = \sum_{\Omega_N} f \frac{\nabla_{i,i+1}(uq \nu_g^N)}{\nu_g^N} \nu_g^N = \nu_g^N \left[u f \frac{\nabla_{i,i+1}(q \nu_g^N)}{\nu_g^N} \right], \quad (\text{B.8})$$

where we used $\nabla_{i,i+1}u = 0$ to obtain the second equality. The gradient in the right-hand side reads:

$$\forall \eta \in \Omega_N, \quad \frac{\nabla_{i,i+1}(q \nu_g^N)(\eta)}{\nu_g^N(\eta)} = q(\eta^{i,i+1}) \exp \left[-\frac{(\eta_{i+1} - \eta_i)}{N} (2C_i^g + \partial^N \lambda_i) \right] - q(\eta), \quad (\text{B.9})$$

where $\lambda_i = \log(\bar{\rho}_i/(1 - \bar{\rho}_i))$. We need to choose a suitable q in order to have a difference $\omega_{i+1} - \omega_i$ arise above. In the $g = 0$ case, corresponding to [JM18], one can take:

$$\tilde{q}(\eta) := \frac{\eta_i(1 - \eta_{i+1})}{\bar{\rho}_i(1 - \bar{\rho}_{i+1})}. \quad (\text{B.10})$$

When $g \neq 0$, the exponential of C_i^g does not change things much, and if q is taken as in (B.7), then:

$$\frac{\nabla_{i,i+1}(q \nu_g^N)(\eta)}{\nu_g^N(\eta)} = \left[\frac{\eta_{i+1}(1 - \eta_i)}{\bar{\rho}_{i+1}(1 - \bar{\rho}_i)} - \frac{\eta_i(1 - \eta_{i+1})}{\bar{\rho}_i(1 - \bar{\rho}_{i+1})} \right] e^{-(\eta_{i+1} - \eta_i)C_i^g/N}. \quad (\text{B.11})$$

The variables ω . (see (B.5)) are tailored to give the above bracket a nice expression (see (A.3) in [JM18]):

$$\left[\frac{\eta_{i+1}(1-\eta_i)}{\bar{\rho}_{i+1}(1-\bar{\rho}_i)} - \frac{\eta_i(1-\eta_{i+1})}{\bar{\rho}_i(1-\bar{\rho}_{i+1})} \right] = \omega_{i+1} - \omega_i + (\bar{\rho}_{i+1} - \bar{\rho}_i)\omega_i\omega_{i+1}. \quad (\text{B.12})$$

This formula can be checked by looking for the left-hand side as a polynomial in ω_i, ω_{i+1} , of the form $a + b\omega_i + c\omega_{i+1} + d\omega_i\omega_{i+1}$ for real numbers a, b, c, d . Equation (B.11) then becomes:

$$\begin{aligned} \frac{\nabla_{i,i+1}(q\nu_g^N)}{\nu_g^N}(\eta) &= \omega_{i+1} - \omega_i + (\omega_{i+1} - \omega_i) \left(e^{-(\eta_{i+1}-\eta_i)C_i^g/N} - 1 \right) \\ &\quad + (\bar{\rho}_{i+1} - \bar{\rho}_i)\omega_i\omega_{i+1}e^{-(\eta_{i+1}-\eta_i)C_i^g/N}, \end{aligned} \quad (\text{B.13})$$

which proves the lemma when plugged into (B.8). \square

The next lemma is a rewriting of Lemma B.1 in terms of the carré du champ operator.

Lemma B.2. *Let $i < N-1$ and let $u : \Omega_N \rightarrow \mathbb{R}$ be such that $\nabla_{i,i+1}u = 0$. There is then a constant $C = C(h, g, \bar{\rho}) > 0$ such that, for any $\delta > 0$:*

$$\begin{aligned} \nu_g^N [u(\omega_{i+1} - \omega_i)f] &\leq \delta N^2 \nu_g^N [\Gamma_h^{i,i+1}(f^{1/2})] + \frac{C}{\delta N^2} \nu_g^N [fu^2] \\ &\quad - (\bar{\rho}_{i+1} - \bar{\rho}_i) \nu_g^N [\omega_i\omega_{i+1}e^{-(\eta_{i+1}-\eta_i)C_i^g/N} uf] \\ &\quad + \nu_g^N [(\omega_{i+1} - \omega_i) \left(1 - e^{-(\eta_{i+1}-\eta_i)C_i^g/N} \right) uf]. \end{aligned} \quad (\text{B.14})$$

Proof. Let $i < N-1$ and $\beta > 0$. In (B.6), write, for each $\eta \in \Omega_N$:

$$\nabla_{i,i+1}f(\eta)u(\eta)q(\eta) = \beta^{1/2}[f^{1/2}(\eta^{i,i+1}) - f^{1/2}(\eta)] \cdot \beta^{-1/2}u(\eta)q(\eta)[f^{1/2}(\eta^{i,i+1}) + f^{1/2}(\eta)]. \quad (\text{B.15})$$

Apply then Cauchy-Schwarz inequality to obtain:

$$\begin{aligned} \nu_g^N [u(\omega_{i+1} - \omega_i)f] &\leq \frac{\beta}{2} \nu_g^N [(\nabla_{i,i+1}f^{1/2})^2] + \frac{1}{\beta} \nu_g^N [(f(\eta) + f(\eta^{i,i+1}))u^2q^2] \\ &\quad - (\bar{\rho}_{i+1} - \bar{\rho}_i) \nu_g^N [\omega_i\omega_{i+1}e^{-(\eta_{i+1}-\eta_i)C_i^g/N} uf] + \nu_g^N [(\omega_{i+1} - \omega_i) \left(1 - e^{-(\eta_{i+1}-\eta_i)C_i^g/N} \right) uf]. \end{aligned} \quad (\text{B.16})$$

Changing variables $\eta \leftarrow \eta^{i,i+1}$, since $\bar{\rho}_i \in [\rho_-, \rho_+] \subset (0, 1)$, the second expectation in (B.16) reads:

$$\begin{aligned} \frac{1}{\beta} \nu_g^N [(f(\eta) + f(\eta^{i,i+1}))u^2q^2] &= \frac{1}{\beta} \sum_{\eta \in \Omega_N} f(\eta)u^2(\eta) \left[q^2(\eta) + q^2(\eta^{i,i+1}) \frac{\nu_g^N(\eta^{i,i+1})}{\nu_g^N(\eta)} \right] \nu_g^N(\eta) \\ &\leq \frac{e^{4\|g\|_\infty} M}{\beta} \nu_g^N [fu^2], \end{aligned} \quad (\text{B.17})$$

where $M = \max\{(\bar{\rho}_j(1-\bar{\rho}_{j+1}))^{-2}, (\bar{\sigma}_j\bar{\sigma}_{j+1})^{-1} : j < N-1\}$ depends only on ρ_\pm .

Consider now the first term in the right-hand side of (B.16). Since $c_h(\eta, i, i+1) \geq c(\eta, i, i+1)e^{-2\|h\|_\infty}$ for each $\eta \in \Omega_N$, it reads:

$$\frac{\beta}{2} \int [\nabla_{i,i+1}f^{1/2}]^2 d\nu_g^N \leq \beta e^{2\|h\|_\infty} \nu_g^N (\Gamma_h^{i,i+1}(\sqrt{f})). \quad (\text{B.18})$$

Taking $\beta = \delta N^2 e^{2\|h\|_\infty}$ concludes the proof. \square

B.2 Integration by parts at the boundary and boundary correlations

Here, we estimate dynamical correlations involving sites close the reservoirs, i.e. correlations of the form $\bar{\eta}_{\pm(N-1)}X_N$ for a function $X_N : \Omega_N \rightarrow \mathbb{R}$. Recall that $h \in \mathcal{S}(\infty)$, $g \in g_0 + \mathcal{S}(\infty)$ are fixed, and the definition (2.2) of the jump rates at the boundary. Define, for $f : \Omega_N \rightarrow \mathbb{R}$:

$$\forall \eta \in \Omega_N, \quad \Gamma_h^\pm(f) = \frac{c_h(\eta, \pm(N-1))}{2} [f(\eta^{\pm(N-1)}) - f(\eta)]^2, \quad (\text{B.19})$$

and observe:

$$c_h(\eta, \pm(N-1)) \geq e^{-2\|h\|_\infty} \min\{\rho_{\varepsilon_1}, (1 - \rho_{\varepsilon_2}) : \varepsilon_1, \varepsilon_2 \in \{-, +\}\}. \quad (\text{B.20})$$

Lemma B.3. *Fix $n \in \mathbb{N}^*$ and, for each $N \in \mathbb{N}^*$, let $\phi_n : \Lambda_N^n \rightarrow \mathbb{R}$ satisfy $\sup_{N \in \mathbb{N}^*} \|\phi_n\|_\infty < \infty$. For $\varepsilon \in \{+, -\}$, define $U_0^\varepsilon(\eta) = \bar{\eta}_{\varepsilon(N-1)}$ ($\eta \in \Omega_N$) and, for $n \geq 1$:*

$$\forall \eta \in \Omega_N, \quad U_n^\varepsilon(\eta) = \bar{\eta}_{\varepsilon(N-1)} V_n^\varepsilon(\eta), \quad V_n^\varepsilon(\eta) = \frac{1}{N^{n-1}} \sum_{i_1, \dots, i_n \neq \varepsilon(N-1)} \bar{\eta}_{i_1} \dots \bar{\eta}_{i_n} \phi_n(i_1, \dots, i_n). \quad (\text{B.21})$$

Then, for each $n \in \mathbb{N}$, U_n^ε is Γ -controllable with size N^{-1} and of vanishing type in the sense of Lemmas 3.9–3.13. More precisely, there is $C > 0$ such that, for any ν_g^N -density f and any $\delta > 0$:

$$\nu_g^N(f N^{1/2} U_n^\varepsilon) \leq \delta N^2 \nu_g^N(\Gamma_h^\varepsilon(\sqrt{f})) + \nu_g^N \left(f \frac{C}{N} \left[\frac{(V_n^\varepsilon)^2}{\delta} + \left| \frac{V_n^\varepsilon}{N^{1/2}} \sum_{j \neq \varepsilon(N-1)} \bar{\eta}_j (N g_{\varepsilon(N-1), j}) \right| \right] \right) + \alpha_n^N, \quad (\text{B.22})$$

setting $V_0^\varepsilon := 1$ by convention and with $\alpha_n^N = O(N^{-1/2})$.

Remark B.4. *The estimate on the size of U_n^ε is optimal only if $n \leq 1$. U_n^ε with $n \leq 1$ are used in the computation of the adjoint in Section 3.2, while U_2^ε is useful in Section C. \blacksquare*

Proof. Let $n \in \mathbb{N}^*$. Using the notations of Corollary A.4, the term in the expectation in the right-hand side of (B.22) is of the form $N^{-(2n-1)} X_{2n, \{0\}}^{\psi_{2n}} + N^{-n-\frac{1}{2}} X_{n+1, \{0\}}^{\psi_{n+1}}$ with $\sup_{N \in \mathbb{N}^*} \|\psi_\ell\|_\infty < \infty$, $\ell \in \{n, 2n\}$. It is thus controllable with size (at most) δ^{-1} . The estimate (B.22) then implies that U_n^ε is Γ -controllable with size N^{-1} by taking $\delta = N^{1/2}$, but also that $N^{1/2} U_n^\varepsilon$ is Γ -controllable with size 1 and thus of vanishing type.

Let us therefore prove (B.22). We do so for the left boundary $\varepsilon = -$, the proof for the right boundary being identical. The goal is to create a gradient of f of the form (B.19). We use the shorthand notation $b = -(N-1)$. Notice that $V_n^-(\eta^b) = V_n^-(\eta)$ for any η . The mapping $\eta \mapsto \eta^b$ is bijective on Ω_N , thus:

$$\nu_g^N(f \bar{\eta}_b V_n^-) = \frac{1}{2} \sum_{\eta \in \Omega_N} \nu_g^N(\eta) V_n^-(\eta) \left[f(\eta) (\eta_b - \bar{\rho}_b) + f(\eta^b) (1 - \eta_b - \bar{\rho}_b) \frac{\nu_g^N(\eta^b)}{\nu_g^N(\eta)} \right]. \quad (\text{B.23})$$

For $\eta \in \Omega_N$, let us compute the ratio $\nu_g^N(\eta^b)/\nu_g^N(\eta)$:

$$\begin{aligned} \frac{\nu_g^N(\eta^b)}{\nu_g^N(\eta)} &= \frac{(1 - \eta_b) \bar{\rho}_b + \eta_b (1 - \bar{\rho}_b)}{\eta_b \bar{\rho}_b + (1 - \eta_b) (1 - \bar{\rho}_b)} e^{2\Pi^N(g)(\eta^b) - 2\Pi^N(g)(\eta)} \\ &= \left(\frac{\bar{\rho}_b}{1 - \bar{\rho}_b} \right)^{1-2\eta_b} \exp \left[\frac{(1 - 2\eta_b)}{N} \sum_{j \neq b} \bar{\eta}_j g_{b,j} \right]. \end{aligned} \quad (\text{B.24})$$

For future reference, notice that (B.24) is bounded by $C(\rho_{\pm})e^{2\|g\|_{\infty}}$ for some $C(\rho_{\pm}) > 0$. Forgetting g for a second in (B.23), notice also that, for each $\eta \in \Omega_N$:

$$\begin{aligned} & f(\eta)(\eta_b - \bar{\rho}_b) + f(\eta^b)(1 - \eta_b - \bar{\rho}_b) \left(\frac{\bar{\rho}_b}{1 - \bar{\rho}_b} \right)^{1-2\eta_b} \\ &= \eta_b(1 - \bar{\rho}_b)[f(\eta) - f(\eta^b)] - (1 - \eta_b)\bar{\rho}_b[f(\eta) - f(\eta^b)] \\ &= \bar{\eta}_b(f(\eta) - f(\eta^b)), \end{aligned} \quad (\text{B.25})$$

which involves a gradient of f as desired. Coming back to (B.24), note that, since $g(\pm 1, \cdot) = 0$, the argument of the exponential in (B.24) is bounded by $O(N^{-1})$. Equation (B.25) and the existence of $C(g) > 0$ such that $|e^x - 1 - x| \leq C(g)x^2$ holds for $x \leq 2\|g\|_{\infty}$ therefore yield the bound:

$$\begin{aligned} & \left| N^{1/2} \nu_g^N(f \bar{\eta}_b V_n^-) - \frac{N^{1/2}}{2} \sum_{\eta \in \Omega_N} \nu_g^N(\eta) V_n^-(\eta) \bar{\eta}_b [f(\eta) - f(\eta^b)] - \frac{N^{1/2}}{2N^2} \sum_{\eta \in \Omega_N} \nu_g^N(\eta) V_n^-(\eta) f(\eta^b) \right. \\ & \quad \left. \times (1 - \eta_b - \bar{\rho}_b) \left(\frac{\bar{\rho}_b}{1 - \bar{\rho}_b} \right)^{1-2\eta_b} (1 - 2\eta_b) \sum_{j \neq b} \bar{\eta}_j(Ng_{b,j}) \right| \leq \frac{C(g) \|V_n^-\|_{\infty}}{N^{3/2}} =: \alpha_N^n. \end{aligned} \quad (\text{B.26})$$

Since $\|V_n^-\|_{\infty} = O(N)$, $\alpha_N^n = O(N^{-1/2})$ as claimed. It remains to estimate the second and third terms in the left-hand side. Consider first the third term. Using the bijection $\eta \mapsto \eta^b$ to turn $f(\eta^b)$ into $f(\eta)$, recalling that $V_n^-(\eta^b) = V_n^-(\eta)$ for each $\eta \in \Omega_N$, and bounding the ratio (B.24) by $C(\rho_{\pm})e^{2\|g\|_{\infty}}$, one finds:

$$\begin{aligned} & \left| \frac{N^{1/2}}{2N^2} \sum_{\eta \in \Omega_N} \nu_g^N(\eta) V_n^-(\eta) f(\eta^b) (1 - \eta_b - \bar{\rho}_b) \left(\frac{\bar{\rho}_b}{1 - \bar{\rho}_b} \right)^{1-2\eta_b} \sum_{j \neq b} \bar{\eta}_j(Ng_{b,j}) \right| \\ & \leq \frac{C(\rho_{\pm})e^{2\|g\|_{\infty}}}{N^{3/2}} \nu_g^N \left(f \left| V_n^- \sum_{j \neq b} \bar{\eta}_j(Ng_{b,j}) \right| \right), \end{aligned} \quad (\text{B.27})$$

which is one of the terms appearing in (B.22). Consider now the second term in the left-hand side of (B.26). For $\beta > 0$ and $\eta \in \Omega_N$, split $N^{1/2}[f(\eta) - f(\eta^b)]$ into $\beta N^{1/2}[f^{1/2}(\eta) - f^{1/2}(\eta^b)]\beta^{-1}[f^{1/2}(\eta) + f^{1/2}(\eta^b)]$ and apply Cauchy-Schwarz inequality twice to find, bounding $(\bar{\eta}_b)^2$ by 1:

$$\begin{aligned} & \frac{N^{1/2}}{2} \sum_{\eta \in \Omega_N} V_n^-(\eta) \bar{\eta}_b [f(\eta) - f(\eta^b)] \\ & \leq \frac{\beta N}{4} \sum_{\eta \in \Omega_N} [f^{1/2}(\eta) - f^{1/2}(\eta^b)]^2 + \frac{1}{2\beta} \sum_{\eta \in \Omega_N} \nu_g^N(\eta) [f(\eta) + f(\eta^b)] (V_n^-)^2(\eta). \end{aligned} \quad (\text{B.28})$$

As for (B.27), the last expectation is bounded by $\beta^{-1}C(g)\nu_g^N(f(V_n^-)^2)$. To conclude the proof, recall from (B.19) the expression of Γ_h^- and from (B.20) the lower bound $C(\rho_{\pm})e^{-\|h\|_{\infty}}$ on the jump rates. Choose $\beta = 2C(\rho_{\pm})e^{-2\|h\|_{\infty}}\delta N$ for $\delta > 0$. Then:

$$\frac{1}{2} \sum_{\eta \in \Omega_N} V_n^-(\eta) \bar{\eta}_b [f(\eta) - f(\eta^b)] \leq \delta N^2 \nu_g^N(\Gamma_h^-(\sqrt{f})) + \frac{C(g)C(\rho_{\pm})^{-1}e^{2\|h\|_{\infty}}}{2\delta N} \nu_g^N(f(V_n^-)^2). \quad (\text{B.29})$$

The last equation together with (B.26) and (B.27) is precisely the right-hand side of (B.22). \square

C Control of the error terms

Fix $h \in \mathcal{S}(\infty)$ (this set is defined in (2.45)) and an associated $g_h \in g_0 + \mathcal{S}(\infty)$, solution of the main equation (2.61). In this section we estimate, for each density f for $\nu_{g_h}^N$, the average against $f\nu_{g_h}^N$ of the function $X_{2,\{0,1\}}^\phi$, $\phi : \Lambda_N^2 \rightarrow \mathbb{R}$, defined below in (C.1). This proves the last item of Lemma 3.9. We also estimate the expectation of the time average of any error term encountered in the text, proving Proposition 3.8 and Corollary 3.14.

Proposition C.1. *Let $\phi \in \Lambda_N^2 \rightarrow \mathbb{R}$ satisfy $\sup_{N \in \mathbb{N}^*} \|\phi_2\|_\infty < \infty$. Recall that $X_{2,\{0,1\}}^\phi$, abbreviated as X_2 , was defined in (A.29) by:*

$$\forall \eta \in \Omega_N, \quad X_2(\eta) = \sum_{i < N-1} \sum_{j \notin \{i, i+1\}} \bar{\eta}_i \bar{\eta}_{i+1} \bar{\eta}_j \phi(i, j). \quad (\text{C.1})$$

Then $N^{-1}X_2$ is Γ -controllable with size $N^{-1/2}$ (recall Definition 3.6 of controllability).

The next proposition was stated as Proposition 3.8, and is proven in Section C.2.

Proposition C.2. *Let $0 < \rho_- \leq \rho_+ < 1$ and $\varepsilon_B(\rho_-)$ be chosen as in Lemma 3.5, and assume $\bar{\rho}' \leq \varepsilon_B$, $h \in \mathcal{S}(\varepsilon_B)$. Let $E^N : \Omega_N \rightarrow \mathbb{R}$ be an error term with size $a_N = o_N(1)$, and let F^N be controllable with size 1. There are then γ, C and $\gamma', C' > 0$ depending on h, ρ_\pm such that, for each $T > 0$ and $N \in \mathbb{N}^*$:*

$$\begin{aligned} \frac{1}{T} \log \mathbb{E}^{\nu_{g_h}^N} \left[\exp \left| \gamma \int_0^T E^N(\eta_t) dt \right| \right] &\leq C a_N, \\ \frac{1}{T} \log \mathbb{E}^{\nu_{g_h}^N} \left[\exp \left| \gamma' \int_0^T F^N(\eta_t) dt \right| \right] &\leq C'. \end{aligned} \quad (\text{C.2})$$

Propositions C.1-C.2 are proven in the next two sections.

C.1 Estimate of X_2

In this section, we prove Proposition C.1. Fix $\phi : \Lambda_N^2 \rightarrow \mathbb{R}$, $N \in \mathbb{N}^*$ as in the proposition. Fix also a density f for $\nu_{g_h}^N$ once and for all. The proof of Proposition C.1 being very technical, we first present its general structure.

The idea is to smoothen the product $\bar{\eta}_i \bar{\eta}_{i+1}$ into a quantity that depends on all $\bar{\eta}$'s in a box of size ℓ with ℓ sufficiently large, then use the entropic inequality to estimate the resulting term. The cost of this replacement will be estimated by an integration by parts formula, see Section B.

We need room between the indices $i, i+1$ and j in the definition (C.1) of X_2 to take averages in a box. Let I_ℓ be the segment $\{0, \dots, \ell-1\}$ and split the sum on j in (C.1) as follows:

$$\forall \eta \in \Omega_N, \quad \frac{1}{N} X_2(\eta) = \overrightarrow{X}_2^\ell + \overleftarrow{X}_2^\ell, \quad \overrightarrow{X}_2^\ell = \frac{1}{N} \sum_{i < N-1} \sum_{\substack{j \in \Lambda_N \\ j \notin i+1+I_\ell}} \bar{\eta}_i \bar{\eta}_{i+1} \bar{\eta}_j \phi(i, j), \quad (\text{C.3})$$

$$\overleftarrow{X}_2^\ell = \frac{1}{N} \sum_{i < N-1} \sum_{\substack{j \in \Lambda_N \setminus \{i, i+1\} \\ j \in i+1+I_\ell}} \bar{\eta}_i \bar{\eta}_{i+1} \bar{\eta}_j \phi(i, j). \quad (\text{C.4})$$

The direction of the arrow indicates the direction in which the replacement of $\bar{\eta}_i$ (\leftarrow) or $\bar{\eta}_{i+1}$ (\rightarrow) by averages on sites to the left of i (\leftarrow) or to the right of $i+1$ (\rightarrow) is going to be performed. Estimates for \overleftarrow{X}_2^ℓ and $\overrightarrow{X}_2^\ell$ are identical, so we only estimate the latter. In practice, the replacement is made thanks to the integration by parts Lemma B.2, which uses $\omega = \bar{\eta}/\bar{\sigma}$ as main variable. Write:

$$A(i, j) := \bar{\sigma}_{i+1} \phi(i, j), \quad i < N-1, j \in \Lambda_N. \quad (\text{C.5})$$

Then:

$$\forall \eta \in \Omega_N, \quad \overrightarrow{X}_2^\ell(\eta) = \frac{1}{N} \sum_{i < N-1} \sum_{\substack{j \in \Lambda_N \setminus \{i, i+1\} \\ j \notin i+1+I_\ell}} \bar{\eta}_i \omega_{i+1} \bar{\eta}_j A(i, j), \quad (\text{C.6})$$

and we replace ω_{i+1} by $\frac{1}{\ell} \sum_{a \in i+1+I_\ell} \omega_a$. If $i+1$ is too close to the reservoirs, i.e. if $i+\ell > N-1$, then this replacement does not make sense. In this case, we spread the unit mass at $i+1$ to $1/\ell$ at each site in $i+1+I_\ell \cap \Lambda_N$, and leave the remaining $\frac{N-1-(i+1)}{\ell}$ mass at the boundary. This is summarised in the following definition of the replacement $\overrightarrow{\omega}_{i+1}^\ell$ of ω_{i+1} :

$$\forall i < N-1, \quad \overrightarrow{\omega}_{i+1}^\ell = \frac{1}{\ell} \sum_{a=i+1}^{\min\{i+\ell, N-1\}} \omega_a + \mathbf{1}_{i+\ell > N-1} \left(1 - \frac{N-1-i}{\ell}\right) \omega_{N-1}. \quad (\text{C.7})$$

Choice of ℓ . Let $\overrightarrow{Y}_2^\ell$ denote the averaged version of $\overrightarrow{X}_2^\ell$:

$$\begin{aligned} \forall \eta \in \Omega_N, \quad \overrightarrow{Y}_2^\ell(\eta) &= \overrightarrow{Z}_2^\ell(\eta) \\ &+ \frac{1}{N} \sum_{i < N-1} \sum_{\substack{j \in \Lambda_N \setminus \{i, i+1\} \\ j \notin i+1+I_\ell}} \bar{\eta}_i \bar{\eta}_j \mathbf{1}_{i+\ell > N-1} \left(1 - \frac{N-1-i}{\ell}\right) \omega_{N-1} A(i, j), \end{aligned} \quad (\text{C.8})$$

with:

$$\forall \eta \in \Omega_N, \quad \overrightarrow{Z}_2^\ell(\eta) = \frac{1}{N} \sum_{i < N-1} \sum_{\substack{j \in \Lambda_N \setminus \{i, i+1\} \\ j \notin i+1+I_\ell}} \bar{\eta}_i \bar{\eta}_j \left(\frac{1}{\ell} \sum_{a=i+1}^{\min\{i+1+\ell, N-1\}} \omega_a \right) A(i, j). \quad (\text{C.9})$$

The last term in (C.8) is Γ -controllable with size N^{-1} by Lemma B.3. For the replacement of $\overrightarrow{X}_2^\ell$ by $\overrightarrow{Y}_2^\ell$ to be useful, $\overrightarrow{Z}_2^\ell$ should be controllable with size $o_N(1)$. This requirement will fix the choice of ℓ . Looking at Corollary A.4, we see that any ℓ such that $\ell = o(N)$ fails, so we take:

$$\ell := N. \quad (\text{C.10})$$

The entropy inequality is then effective on \overrightarrow{Z}_2^N . Indeed, it is of the form:

$$\forall \eta \in \Omega_N, \quad \overrightarrow{Z}_2^N(\eta) = \frac{1}{N^2} \sum_{\substack{(i, j, a) \in \Lambda_N^3 \\ |\{(i, j, a)\}|=3}} \bar{\eta}_i \bar{\eta}_j \bar{\eta}_a \tilde{A}(i, j, a), \quad (\text{C.11})$$

for some function \tilde{A} satisfying $|\tilde{A}(i, j, a)| \leq |A(i, j)|$ for each $(i, j, a) \in \Lambda_N^3$. By Corollary A.4, \overrightarrow{Z}_2^N is therefore controllable with size $N^{-1/2}$: there are $\gamma, C > 0$ such that:

$$\nu_{g_h}^N(f \overrightarrow{Z}_2^N) \leq \frac{H(f \nu_g^N | \nu_g^N)}{\gamma} + \frac{1}{\gamma} \log \nu_g^N \left[\exp(\gamma \overrightarrow{Z}_2^N) \right] \leq \frac{H(f \nu_g^N | \nu_g^N)}{\gamma} + \frac{C}{N^{1/2}}. \quad (\text{C.12})$$

Cost of the replacement. Let us estimate the cost of replacing \vec{X}_2^N by \vec{Y}_2^N , defined in (C.8). To do so, we use of the following integration by parts identity, which explicitly describes how to spread the unit mass at $i + 1$, $i < N - 1$, to $1/N$ on every site up to the boundary, where the remaining mass is then left. One has:

$$\omega_{i+1} - \vec{\omega}_{i+1}^N = \sum_{a=i+1}^{\min\{i+N, N-1\}-1} \phi_N(a - (i + 1))(\omega_a - \omega_{a+1}), \quad \phi_N(b) = \frac{N - 1 - b}{N} \mathbf{1}_{0 \leq b < N}, \quad b \in \mathbb{Z}. \quad (\text{C.13})$$

For brevity, for $a \in \Lambda_N$, let u_a denote the quantity:

$$\forall \eta \in \Omega_N, \quad u_a(\eta) = \frac{1}{N} \sum_{\substack{i < N-1 \\ a-N < i < a}} \phi_N(a - (i + 1)) \bar{\eta}_i \sum_{\substack{j \in \Lambda_N \setminus \{i, i+1\} \\ j \notin i+1+I_N}} \bar{\eta}_j A(i, j). \quad (\text{C.14})$$

Then, for each $\eta \in \Omega_N$:

$$\vec{X}_2^N - \vec{Y}_2^N = \sum_{a < N-1} (\omega_a - \omega_{a+1}) u_a(\eta). \quad (\text{C.15})$$

To estimate the expectation of the right-hand side above under $f\nu_g^N$, apply, for each $a < N - 1$, the integration by parts formula of Lemma B.2, with $u = -u_a$. There is thus a constant $C > 0$ such that, for each $\delta > 0$:

$$\sum_{a < N-1} \nu_{g_h}^N(f(\omega_a - \omega_{a+1})u_a) \leq \delta N^2 \nu_{g_h}^N(\Gamma_h(\sqrt{f})) + \frac{C}{\delta N^2} \sum_{a < N-1} \nu_{g_h}^N(u_a^2 f) \quad (:= \nu_{g_h}^N(fR_1)) \quad (\text{C.16})$$

$$+ \sum_{a < N-1} (\rho_{a+1} - \rho_a) \nu_{g_h}^N(\omega_a \omega_{a+1} e^{-(\eta_{a+1} - \eta_a) C_a^{g_h}/N} u_a f) \quad (:= \nu_{g_h}^N(fR_2)) \quad (\text{C.17})$$

$$- \sum_{a < N-1} \nu_{g_h}^N\left((\omega_{a+1} - \omega_a) \left(1 - e^{-(\eta_{a+1} - \eta_a) C_a^{g_h}/N}\right) u_a f\right) \quad (:= \nu_{g_h}^N(fR_3)). \quad (\text{C.18})$$

Let us estimate one by one each of (C.16)-(C.17)-(C.18). Consider first $\nu_{g_h}^N(fR_3)$, and note that:

$$\sup_N \sup_{a \in \Lambda_N} |(\omega_{a+1} - \omega_a)(\eta_{a+1} - \eta_a)| \leq C(\rho_{\pm}). \quad (\text{C.19})$$

As a result, using the identity $e^x = 1 + \int_0^1 x e^{tx} dt$ for $x \in \mathbb{R}$ and the fact that $|C^{g_h}| \leq 2\|g_h\|_{\infty}$, $\nu_{g_h}^N(fR_3)$ can be bounded as follows:

$$|\nu_{g_h}^N(fR_3)| \leq C(\rho_{\pm}) e^{2\|g\|_{\infty}} \nu_{g_h}^N\left(\frac{1}{N} \sum_{a < N-1} f |C_a^{g_h} u_a|\right). \quad (\text{C.20})$$

By definition of C^{g_h} (see e.g. Lemma B.1) and of u in (C.14), the product $C^{g_h} u$ is of the form:

$$\forall a \in \Lambda_N, \quad C_a^{g_h} u_a = \frac{1}{N^2} \sum_{(i,j,b) \in \Lambda_N^3} \bar{\eta}_i \bar{\eta}_j \bar{\eta}_b \psi_{i,j,b}^a + D_a^{g_h} u_a, \quad \sup D^{g_h} = O(N^{-1}), \quad (\text{C.21})$$

where the functions $\psi^a : (-1, 1)^3 \rightarrow \mathbb{R}$ are bounded uniformly in a . It follows by Corollary A.4 that R_3 is controllable with size $N^{-1/2}$.

Consider now R_2 , defined in (C.17). Again using $e^x = 1 + \int_0^1 xe^{tx} dt$ for $x \in \mathbb{R}$, we can bound $\nu_{g_h}^N(fR_2)$ as follows:

$$\nu_{g_h}^N(fR_2) \leq \frac{\bar{\rho}'}{N} \sum_{a < N-1} \nu_{g_h}^N(f\omega_a\omega_{a+1}u_a) + \frac{\bar{\rho}'e^{2\|g_h\|_\infty}}{N^2} \sum_{a < N-1} \frac{1}{\bar{\sigma}_a\bar{\sigma}_{a+1}} \nu_{g_h}^N\left(f|C_a^{g_h}u_a|\right). \quad (\text{C.22})$$

Recalling the definition of u from (C.14), the first term in (C.22) is of the form $N^{-2}X_{3,\{0,1\}}^B$ in the notations of Theorem A.3, i.e. of the form $N^{-2}\sum_{i,j,b} \bar{\eta}_i\bar{\eta}_{i+1}\bar{\eta}_j\bar{\eta}_b B(i,j,b)$, with B bounded. Corollary A.4 tells us that this function does not behave worse than a sum of three-point correlations, and is therefore controllable with size $N^{-1/2}$. In addition, the second term in (C.22) has the same structure as $N^{-1}R_3$, and is therefore controllable with size $N^{-3/2}$.

Consider finally R_1 in (C.16). It reads:

$$\nu_{g_h}^N(fR_1) = \frac{C}{\delta N^4} \sum_{a < N-1} \sum_{\substack{i,j < N-1 \\ a-N < i,j < a}} \left[\phi_N(a-(i+1))\phi_N(a-(j+1))\bar{\eta}_i\bar{\eta}_j \right. \\ \left. \times \sum_{\substack{b,c \in \Lambda_N \setminus \{i,i+1\} \\ b \notin i+1+I_N \\ c \notin j+1+I_N}} \bar{\eta}_b\bar{\eta}_c A(i,b)A(j,c) \right]. \quad (\text{C.23})$$

In particular, it is of the form $N^{-3}X_{4,\{0\}}^{v_4}$ for a $v_4 : \Lambda_N^4 \rightarrow \mathbb{R}$ satisfying $\sup_N \sup_{\Lambda_N^4} |v_4| < \infty$, and therefore controllable with size N^{-1} by Corollary A.4. For each $\delta > 0$, we have proven the existence of a controllable function $R_\delta : \Omega_N \rightarrow \mathbb{R}$ such that:

$$\left| \nu_{g_h}^N(\overrightarrow{X}_2^N) \right| \leq 2\delta N^2 \nu_{g_h}^N(\Gamma_h(\sqrt{f})) + \nu_{g_h}^N(f(\overrightarrow{Y}_2^N + R_\delta)), \quad R_\delta := R_1^\delta + R_2 + R_3. \quad (\text{C.24})$$

The arguments above do not depend on the sign of A in the definition (C.4) of \overrightarrow{X}_2^N . This implies that \overrightarrow{X}_2^N is Γ controllable with size $N^{-1/2}$ in the sense of Definition 3.6. Since the same arguments also apply to \overleftarrow{X}_2^N , Proposition C.1 is proven.

C.2 Proof of Propositions 3.8 and Corollary 3.14

Proof. Let $T > 0$. Corollary 3.14 is obtained as a side product of the proof of Proposition C.2, which we focus on. Write G^N for either the error term E^N with size a_N , or the (Γ) -controllable function F^N with size 1. For each $\gamma > 0$, Feynman-Kac inequality (3.6) and the bound (3.15) on the adjoint $L_h^* \mathbf{1}$ in $\mathbb{L}^2(\nu_{g_h}^N)$ imply:

$$\frac{1}{T} \log \mathbb{E}_h^{\nu_{g_h}^N} \left[\exp \left[\gamma \int_0^T G^N(\eta_t) dt \right] \right] \leq \sup_{f \geq 0: \nu_{g_h}^N(f)=1} \left\{ \nu_{g_h}^N \left(f(\gamma G^N + \mathcal{E}/2) \right) - \frac{N^2}{4} \nu_{g_h}^N(\Gamma_h(\sqrt{f})) \right\}. \quad (\text{C.25})$$

In addition, the function \mathcal{E} satisfies, for some constant $C = C(\rho_\pm, h) > 0$:

$$\nu_{g_h}^N(f\mathcal{E}) \leq \frac{H(f\nu_{g_h}^N|\nu_{g_h}^N)}{4C_{LS}} + \frac{2C}{N^{1/2}} \leq \frac{N^2}{4} \nu_{g_h}^N(\Gamma_h(\sqrt{f})) + \frac{2C}{N^{1/2}}, \quad (\text{C.26})$$

where we used the entropy inequality first, then the log-Sobolev inequality of Proposition A.5 to get the right-hand side. Thus:

$$\frac{1}{T} \log \mathbb{E}_h^{\nu_{g_h}^N} \left[\exp \left[\gamma \int_0^T G^N(\eta_t) dt \right] \right] \leq \sup_{f \geq 0: \nu_{g_h}^N(f) = 1} \left\{ \nu_{g_h}^N(f \gamma G^N) - \frac{N^2}{8} \nu_{g_h}^N(\Gamma_h(\sqrt{f})) + \frac{C}{N^{1/2}} \right\}. \quad (\text{C.27})$$

Suppose first that G^N is controllable with size s_N (including both $s_N = O_N(1)$ and $s_N = o_N(1)$ cases): for some $\gamma_0 > 0$,

$$\forall N \in \mathbb{N}^*, \quad \log \nu_{g_h}^N [e^{\gamma_0 G^N}] \leq s_N. \quad (\text{C.28})$$

By the entropy- and log-Sobolev inequalities, the quantity inside the supremum in Equation (C.27) is bounded above, for each density f for $\nu_{g_h}^N$, by:

$$\left(\frac{\gamma C_{LS}}{\gamma_0} - \frac{1}{8} \right) N^2 \nu_{g_h}^N(\Gamma_h(\sqrt{f})) + \frac{\gamma s_N}{\gamma_0} + \frac{C}{N^{1/2}}. \quad (\text{C.29})$$

Taking any $\gamma < \frac{\gamma_0}{8C_{LS}}$ ensures that the first term is negative, and concludes the proof in the controllable case:

$$\forall N \in \mathbb{N}^*, \quad \frac{1}{T} \log \mathbb{E}_h^{\nu_{g_h}^N} \left[\exp \left[\gamma \int_0^T G^N(\eta_t) dt \right] \right] \leq \frac{\gamma s_N}{\gamma_0} + CN^{1/2}. \quad (\text{C.30})$$

If G^N is not only controllable but also of LS type (recall Definition 3.11), corresponding to the controllable case in Corollary 3.14, then $\gamma_0 > 2^{10}C_{LS} > 8C_{LS}$ by assumption, thus one can take $\gamma > 1$.

If now G^N is Γ -controllable with size s_N , the idea is the same, except that one first bounds $\nu_{g_h}^N(fG^N)$ from above using Definition 3.6 of Γ -controllability:

$$\forall \delta > 0, \quad \nu_{g_h}^N(fG^N) \leq \delta N^2 \nu_{g_h}^N(\Gamma_h(\sqrt{f})) + \frac{1}{\delta} \nu_g^N(Y_{G^N}), \quad (\text{C.31})$$

with Y_{G^N} controllable with size s_N . Choosing $\delta = \frac{1}{16}$, the proof becomes identical to the controllable case. \square

Remark C.3. *If one is interested only in estimating the expectation of the time integral of G^N rather than its exponential moment, then the log-Sobolev inequality is not necessary. One can instead directly rely on Theorem 2.6 that bounds the relative entropy along the dynamics.* \blacksquare

D The Neumann condition on the diagonal

Let $h \in \mathcal{S}(\infty)$ (this set is defined in (2.45)) Assume that $\bar{\rho}' \leq \varepsilon_B$, $h \in \mathcal{S}(\varepsilon_B)$ so that the solution $g_h \in g_0 + \mathcal{S}(\infty)$ of the main equation (2.61) exists and the conclusions of Lemma 3.5 hold. In this section, we rewrite the term:

$$\frac{1}{4} \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} (\partial_1 h_{i+,i} - \partial_1 h_{i-,i}) = \frac{1}{4} \sum_{i < N-1} \omega_i \omega_{i+1} \bar{\sigma}_i \bar{\sigma}_{i+1} (\partial_1 h_{i+,i} - \partial_1 h_{i-,i}), \quad \omega_i = \frac{\bar{\eta}_i}{\bar{\sigma}_i}, \quad (\text{D.1})$$

as a function of the two-point correlations field Π^N , defined in (2.7). This is necessary in the proof of upper-bound large deviations, in order to obtain a closed expression of the Radon-Nikodym derivative in terms of the field Π^N . It is done through the integration by parts Lemma B.1, replacing $\omega_i \omega_{i+1}$ by local averages of ω 's. As $h \in \mathcal{S}(\infty)$, the function $\delta_h(x) = \partial_1 h(x_+, x) - \partial_1 h(x_-, x)$, $x \in (-1, 1)$ can be extended into an element of $C^2([-1, 1])$, still denoted by δ_h , which satisfies $\delta_h(\pm 1) = 0$. Let $\varepsilon \in (0, 1) \in \mathbb{N}^*$ and $I_{\varepsilon N} := \{0, \dots, N\varepsilon - 1\}$, writing εN for $\lfloor \varepsilon N \rfloor$. In the large N limit, the correlation fields Π we consider act on \mathcal{T} , which only contains functions with a certain regularity. We cannot simply replace ω_{i+1} by a uniform average of ω on $i + 1 + I_{\varepsilon N}$ and obtain an element of \mathcal{T} (the indicator function $\mathbf{1}_{[0, \varepsilon]}$ is not regular enough). Consider instead a smooth function $\chi^\varepsilon \in C^\infty(\bar{\square})$ with $\chi^\varepsilon = 0$ on $\partial \square$, $0 \leq \chi^\varepsilon \leq 2/\varepsilon$, and such that $\chi^\varepsilon(x, \cdot)$ approximates $\varepsilon^{-1} \mathbf{1}_{(x, x+\varepsilon) \cap (-1, 1)}$ in the following sense: $\chi^\varepsilon(x, \cdot)$ is supported on $(x, x + \varepsilon) \cap (-1, 1)$ for each $x \in (-1, 1)$, and:

$$\int_{\square} |\chi^\varepsilon(x, y) - \varepsilon^{-1} \mathbf{1}_{(x, x+\varepsilon) \cap (-1, 1)}(y)|^2 dx dy \leq \varepsilon. \quad (\text{D.2})$$

Define then, recalling that $\bar{\sigma}(x) = \bar{\rho}(x)(1 - \bar{\rho}(x))$ for $x \in [-1, 1]$:

$$\forall (x, y) \in \square, \quad \mathcal{N}_\varepsilon^{\delta_h}(x, y) = \frac{\bar{\sigma}(x)}{\bar{\sigma}(y)} \delta_h(x) \chi^\varepsilon(x, y). \quad (\text{D.3})$$

Note that $\mathcal{N}_\varepsilon^{\delta_h}$ belongs to \mathcal{T} , defined in (2.38), thus $\Pi(\mathcal{N}_\varepsilon^{\delta_h})$ is now a well defined object for each $\Pi \in \mathcal{T}'_s$. At the microscopic level, Π^N also acts on functions with less regularity and we have the following more general result.

Proposition D.1. *Let $q \in C^0([-1, 1])$. Define, for $\eta \in \Omega_N$:*

$$W_q(\eta) = \frac{1}{4} \sum_{i < N-1} \bar{\eta}_i \bar{\eta}_{i+1} q_i. \quad (\text{D.4})$$

For each $\varepsilon \in (0, 1)$ smaller than some $\varepsilon_0(\rho_\pm, h, q) > 0$, there are constants $C_1(\rho_\pm, h, q)$, $C_2(\rho_\pm, h, q, \varepsilon) > 0$ such that, for each N larger than some constant depending on ε and each $T > 0$:

$$\mathbb{E}_h^{\nu_h^N} \left[\exp \left[\int_0^T [W_q(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^q)] dt \right] \right] \leq \exp \left[C_1(\rho_\pm, h, q) \varepsilon^{1/2} T + \frac{C_2(\rho_\pm, h, q, \varepsilon) T}{N^{1/2}} \right]. \quad (\text{D.5})$$

Proof. Let $T > 0$, $\varepsilon \in (0, 1)$, and write:

$$W_q - \Pi^N(\mathcal{N}_\varepsilon^q) = W_q - \vec{W}_q^{\varepsilon N} + \vec{W}_q^{\varepsilon N} - \Pi^N(\mathcal{N}_\varepsilon^q), \quad (\text{D.6})$$

where $\vec{W}_q^{\varepsilon N}$ corresponds to W_q in which the unit mass at each $i + 1 < N$ has been replaced by a mass $(\varepsilon N)^{-1}$ at each site in $\{i + 1, \dots, i + \varepsilon N \wedge N - 1\}$ (recall that $\varepsilon N = \lfloor \varepsilon N \rfloor$):

$$\vec{W}_q^{\varepsilon N} := \Pi^N(A_\varepsilon), \quad A_\varepsilon(x, y) := \frac{\bar{\sigma}(x)}{\bar{\sigma}(y)} q(x) \varepsilon^{-1} \mathbf{1}_{(x, x+\varepsilon) \cap (-1, 1)}(y), \quad (x, y) \in \square. \quad (\text{D.7})$$

Up to applying Cauchy-Schwarz inequality in the exponential moment in (D.5), it is enough to separately estimate the contribution of each difference in (D.6).

Consider first the contribution of $\vec{W}_q^{\varepsilon N} - \Pi^N(\mathcal{N}_\varepsilon^q)$. Corollary A.4 tells us that $\|A_\varepsilon - \mathcal{N}_\varepsilon^q\|_{2,N}^{-1} (\vec{W}_q^{\varepsilon N} - \Pi^N(\mathcal{N}_\varepsilon^q))$ satisfies, for some $\gamma, C > 0$ independent of ε, N, q :

$$\frac{1}{T} \log \mathbb{E}_h^{\nu_{gh}^N} \left[\exp \left[\gamma \int_0^T \|A_\varepsilon - \mathcal{N}_\varepsilon^q\|_{2,N}^{-1} |\vec{W}_q^{\varepsilon N}(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^q)| dt \right] \right] \leq C. \quad (\text{D.8})$$

Hölder inequality applied to (D.8) then yields:

$$\frac{1}{2T} \log \mathbb{E}_h^{\nu_{gh}^N} \left[\exp \left[2 \int_0^T |\vec{W}_q^{\varepsilon N}(\eta_t) - \Pi_t^N(\mathcal{N}_\varepsilon^q)| dt \right] \right] \leq \frac{1}{\gamma} \|A_\varepsilon - \mathcal{N}_\varepsilon^q\|_{2,N}. \quad (\text{D.9})$$

Let us estimate $\|A_\varepsilon - \mathcal{N}_\varepsilon^q\|_{2,N}$. Since χ^ε is smooth,

$$\max_{(i,j) \in \Lambda_N^2} \sup_{(r,s) \in [0,1]^2} \left| \chi_{i,j}^\varepsilon - \chi^\varepsilon \left(\frac{i+r}{N}, \frac{j+s}{N} \right) \right| \leq \frac{C(\varepsilon)}{N}. \quad (\text{D.10})$$

It is also not difficult to check that, for some different $C(\varepsilon) > 0$:

$$\frac{1}{N^2} \sum_{(i,j) \in \Lambda_N^2} \int_{[0,1]^2} \varepsilon^{-2} \left[\mathbf{1} \left\{ \frac{j}{N} \in \left(\frac{i}{N}, \frac{i}{N} + \varepsilon \right) \right\} - \mathbf{1} \left\{ \frac{j}{N} \in \left(\frac{i}{N}, \frac{i}{N} + \varepsilon \right) \right\} \right]^2 du dv \leq \frac{3}{\varepsilon^2 N}. \quad (\text{D.11})$$

From (D.10)–(D.11) and property (D.2) of χ^ε , using also the elementary bound $(a+b)^2 \leq 2a^2 + 2b^2$ for real a, b ; the difference $A_\varepsilon - \mathcal{N}_\varepsilon^q$ therefore satisfies:

$$\begin{aligned} \|A_\varepsilon - \mathcal{N}_\varepsilon^q\|_{2,N}^2 &:= \frac{1}{N^2} \sum_{(i,j) \in \Lambda_N^2} (A_\varepsilon - \mathcal{N}_\varepsilon^q)_{i,j}^2 \leq 2 \|A_\varepsilon - \mathcal{N}_\varepsilon^q\|_2^2 + \frac{C(\rho_\pm, q, \varepsilon)}{N} \\ &\leq 2\varepsilon + \frac{C(\rho_\pm, q, \varepsilon)}{N}. \end{aligned} \quad (\text{D.12})$$

This yields a bound on (D.9) of the same form as the right-hand side in Proposition D.1 for any $\varepsilon > 0$ and any N large enough depending on ε, ρ_\pm, q .

Consider now the contribution of $W_q - \vec{W}_q^{\varepsilon N}$ to (D.5). The idea is the same as in the proof of Proposition C.2: we express $W_q - \vec{W}_q^{\varepsilon N}$ in terms of the carré du champ operator and explicit controllable functions with size vanishing when ε is small. We start from the bound (C.27) on exponential moments:

$$\begin{aligned} &\frac{1}{T} \log \mathbb{E}_h^{\nu_{gh}^N} \left[\exp \left[\int_0^T [W_q(\eta_t) - \vec{W}_q^{\varepsilon N}(\eta_t)] dt \right] \right] \\ &\leq \sup_{f \geq 0: \nu_{gh}^N(f)=1} \left\{ \nu_{gh}^N \left(f [W_q - \vec{W}_q^{\varepsilon N}] \right) - \frac{N^2}{8} \nu_{gh}^N (\Gamma_h(\sqrt{f})) + \frac{C(\rho_\pm, h)}{N^{1/2}} \right\}. \end{aligned} \quad (\text{D.13})$$

Recall Definition 3.6 of controllability. To obtain the claim (D.8), it is enough to prove:

$$\varepsilon^{-1/2} [W_q - \vec{W}_q^{\varepsilon N}] \text{ is } (\Gamma\text{-})\text{controllable with size } 1. \quad (\text{D.14})$$

Indeed, if so, $W_q - \overrightarrow{W}_q^{\varepsilon N}$ will be (Γ) -controllable with size $\varepsilon^{1/2}$, and of LS type (see Definition 3.11) for ε small enough depending on q, ρ_{\pm}, h .

To prove (D.14), we use the integration by parts Lemma B.1. It is formulated with the variables $\omega_i = \bar{\eta}_i/\bar{\sigma}_i$, $i \in \Lambda_N$, for which W_q becomes:

$$\forall \eta \in \Omega_N, \quad W_q(\eta) = \frac{1}{4} \sum_{i < N-1} \omega_i \omega_{i+1} \bar{\sigma}_i \bar{\sigma}_{i+1} q_i = \frac{1}{4} \sum_{i < N-1} \omega_i \omega_{i+1} (\bar{\sigma}_i)^2 q_i + \theta^{N,0}(\eta), \quad (\text{D.15})$$

where $\theta^{N,0}$ is the error term:

$$\forall \eta \in \Omega_N, \quad \theta^{N,0}(\eta) = \frac{1}{4N} \sum_{i < N-1} \omega_i \omega_{i+1} N \bar{\sigma}_i [\bar{\sigma}_{i+1} - \bar{\sigma}_i] q_i. \quad (\text{D.16})$$

It is of the form $N^{-1} X_{1,\{0,1\}}^v$ with the notations of Theorem A.3, thus controllable with size N^{-1} and of vanishing type (recall the terminology of Definitions 3.6–3.11).

Recall from (C.7) the definition of the quantity $\overrightarrow{\omega}_{i+1}^{\varepsilon N}$:

$$\overrightarrow{\omega}_{i+1}^{\varepsilon N} = \frac{1}{\varepsilon N} \sum_{j=i+1}^{\min\{i+\varepsilon N, N-1\}} \omega_j + \mathbf{1}_{i+\varepsilon N > N-1} \left(1 - \frac{N-1-i}{\varepsilon N}\right) \omega_{N-1}, \quad i < N-1. \quad (\text{D.17})$$

For each $i < N-1$, one can write as before:

$$\omega_{i+1} - \overrightarrow{\omega}_{i+1}^{\varepsilon N} = \sum_{a=i+1}^{\min\{i+N, N-1\}-1} \phi_{\varepsilon N}(a-(i+1))(\omega_a - \omega_{a+1}), \quad \phi_{\varepsilon N}(b) = \frac{\varepsilon N - 1 - b}{\varepsilon N} \mathbf{1}_{0 \leq b < \varepsilon N}, \quad b \in \mathbb{Z}. \quad (\text{D.18})$$

Define then $u_j : \Omega_N \rightarrow \mathbb{R}$ for $j > -(N-1)$ similarly to (C.14):

$$\forall \eta \in \Omega_N, \quad u_j(\eta) = \sum_{i < j} \bar{\sigma}_i^2 q_i \omega_i \phi_{\varepsilon N}(j - (i+1)). \quad (\text{D.19})$$

With this definition, the quantity $W_q - \overrightarrow{W}_q^{\varepsilon N}$ reads, for each $\eta \in \Omega_N$:

$$\begin{aligned} W_q(\eta) - \overrightarrow{W}_q^{\varepsilon N}(\eta) &= \sum_{j > -(N-1)} (\omega_j - \omega_{j+1}) u_j(\eta) \\ &+ \omega_{N-1} \sum_{i < N-1} \mathbf{1}_{i+\varepsilon N > N-1} \left(1 - \frac{N-1-i}{\varepsilon N}\right) \omega_i (\bar{\sigma}_i)^2 + \theta^{N,0}(\eta). \end{aligned} \quad (\text{D.20})$$

Fix a density f for ν_{gh}^N . The first term in the second line involves boundary correlations. According to Lemma B.3, it is Γ -controllable with size N^{-1} and of vanishing type: for each $\delta > 0$, there is a function \mathcal{D}_δ , controllable with size N^{-1} and of vanishing type, such that:

$$\nu_{gh}^N \left(f \omega_{N-1} \sum_{i < N-1} \mathbf{1}_{i+\varepsilon N > N-1} \left(1 - \frac{N-1-i}{\varepsilon N}\right) (\bar{\sigma}_i)^2 q_i \omega_i \right) \leq \delta N^2 \nu_{gh}^N(\Gamma_h(\sqrt{f})) + \nu_{gh}^N(f \mathcal{D}_\delta). \quad (\text{D.21})$$

It therefore remains to estimate the other term in the right-hand side of (D.20). By the integration by parts Lemma B.2 applied to $-u_j$ for each $j < N - 1$, there is a constant $C > 0$ such that, for any $\delta > 0$:

$$\begin{aligned} \nu_{g_h}^N \left(f \sum_{j < N-1} (\omega_j - \omega_{j+1}) u_j \right) &\leq \delta N^2 \nu_{g_h}^N (\Gamma_h(\sqrt{f})) + \frac{C}{\delta N^2} \sum_{j < N-1} \int f |u_j|^2 d\nu_{g_h}^N \quad (:= \frac{1}{\delta} \nu_{g_h}^N (f \mathcal{N}_1)) \\ &+ \sum_{j < N-1} (\bar{\rho}_{j+1} - \bar{\rho}_j) \int \omega_j \omega_{j+1} e^{-(\eta_{j+1} - \eta_j) C_j^{g_h} / N} f u_j d\nu_{g_h}^N \quad (:= \nu_{g_h}^N (f \mathcal{N}_2)) \\ &- \sum_{j < N-1} \int (\omega_{j+1} - \omega_j) \left(1 - e^{-(\eta_{j+1} - \eta_j) C_j^{g_h} / N} \right) f u_j d\nu_{g_h}^N \quad (:= \nu_{g_h}^N (f \mathcal{N}_3)). \end{aligned} \quad (\text{D.22})$$

The functions \mathcal{N}_i ($1 \leq i \leq 3$) are then estimated as in (C.23)–(C.22)–(C.20) respectively, but not all of them are error terms and one has to be careful to get a bound that vanishes with ε . Let us check that each of them indeed satisfies (D.14). For \mathcal{N}_1 , (D.19) and the bound $\|u_j\|_\infty \leq \|q\|_\infty \varepsilon N$ imply that it is of the form $N^{-1} X_{2,\{0\}}^{v_2}$ with v_2 given by:

$$v_2(a, b) = \frac{C}{\delta N} \sum_{j < N-1} \mathbf{1}_{j > \max\{a, b\}} q_a q_b \bar{\sigma}_a^2 \bar{\sigma}_b^2 \phi_{\varepsilon N}(j - (a + 1)) \phi_{\varepsilon N}(j - (b + 1)) \leq \frac{C \|q\|_\infty^2 \varepsilon}{\delta}. \quad (\text{D.23})$$

It follows that $\varepsilon^{-1} \mathcal{N}_1$ satisfies (D.14), thus $\varepsilon^{-1/2} \mathcal{N}_1$ as well.

Consider now \mathcal{N}_2 . From (C.22), we get:

$$\nu_{g_h}^N (f \mathcal{N}_2) \leq \frac{\bar{\rho}'}{N} \sum_{j < N-1} \omega_j \omega_{j+1} u_j + \frac{\bar{\rho}' e^{2\|g_h\|_\infty}}{N} \sum_{j < N-1} \frac{1}{\bar{\sigma}_j \bar{\sigma}_{j+1}} \nu_{g_h}^N (f |C_j^{g_h} u_j|). \quad (\text{D.24})$$

Here we do not even need the Γ -controllability of the first term with size $N^{-1/2}$ established in Proposition C.2. Instead, recall that $C^{g_h} = B^{g_h} + D^{g_h}$ with $\sup D^{g_h} = O(N^{-1})$ and B^{g_h} given in (B.1). It is then enough to notice that the right-hand side of (D.24) is an average under $f \nu_{g_h}^N$ of $N^{-1} X_{2,\{0,1\}}^{w_2} + |N^{-1} X_{2,\{0\}}^{w'_2}| + Y_2$, where:

$$\begin{aligned} w_2(i, j) &= \mathbf{1}_{i < j} \frac{\bar{\rho}' \bar{\sigma}_i}{\bar{\sigma}_j} q_i \phi_{\varepsilon N}(j - (i + 1)), \\ w'_2(a, b) &= \frac{\bar{\rho}' e^{2\|g_h\|_\infty}}{2N} \sum_{\substack{j > b \\ j \notin \{a-1, a\}}} \bar{\sigma}_b q_b \phi_{\varepsilon N}(j - (b + 1)) (\partial_1^N g_h)_{j, a}, \\ Y_2 &= \frac{\bar{\rho}' e^{2\|g_h\|_\infty}}{N} \sum_{j < N-1} |u_j D_j^{g_h}| \leq C(\rho_\pm, q, g_h) \varepsilon. \end{aligned} \quad (\text{D.25})$$

Due to the fact that $\phi_{\varepsilon N}$ is non-zero only for εN different integers, both w_2 and w'_2 have $\|\cdot\|_{2, N}$ -norm bounded by $C \varepsilon^{1/2}$ for some $C = C(\rho_\pm, q, g_h) > 0$. It follows from Corollary A.4 that $\varepsilon^{-1/2} \mathcal{N}_2$ is controllable with size 1 and of large type, i.e. satisfies (D.14).

Consider finally \mathcal{N}_3 , defined in (D.22). The bound (C.20) shows that $\nu_{g_h}^N (f \mathcal{N}_3)$ is bounded by a constant times the second term in (D.24). It follows that $\varepsilon^{-1/2} \mathcal{N}_3$ also satisfies (D.14), which concludes the proof. \square

E Sobolev spaces

Definition E.1. Let $U \subset \mathbb{R}^2$ be a bounded open set with Lipschitz boundary. For $n \in \mathbb{N}$ and $p \geq 1$, let $\mathbb{W}^{n,p}(U)$ be the following space. If $n = 0$, it is simply $\mathbb{L}^p(U)$. If $n \geq 1$, $\mathbb{W}^{n,p}(U)$ is the set of functions $f \in \mathbb{L}^p(U)$ such that, for any $(n_1, n_2) \in \mathbb{N}^2$ with $n_1 + n_2 \leq n$, there is $f^{n_1, n_2} \in \mathbb{L}^p(U)$ satisfying:

$$\forall u \in C_c^\infty(U), \quad \int_U f(x, y) \partial_1^{n_1} \partial_2^{n_2} u(x, y) dx dy = (-1)^{n_1+n_2} \int_U f^{n_1, n_2}(x, y) u(x, y) dx dy. \quad (\text{E.1})$$

$\mathbb{W}^{n,p}(U)$ is a separable Banach space for the norm:

$$\forall f \in \mathbb{W}^{n,p}(U), \quad \|f\|_{\mathbb{W}^{n,p}(U)} = \left[\sum_{\substack{(n_1, n_2) \in \mathbb{N}^2 \\ n_1 + n_2 \leq n}} \|f^{n_1, n_2}\|_{\mathbb{L}^p(U)}^2 \right]^{1/2}. \quad (\text{E.2})$$

Moreover, the set $C^\infty(\bar{U})$ of restrictions of elements of $C^\infty(\mathbb{R}^2)$ to \bar{U} is dense in $\mathbb{W}^{n,p}(U)$ for $\|\cdot\|_{\mathbb{W}^{n,p}(U)}$. In the special case $p = 2$, define $\mathbb{H}^n(U) := \mathbb{W}^{n,p}(U)$. This is a Hilbert space.

Along the text, we make use of the following Sobolev embedding results (see Theorem 4.12 in [AF03] and Theorem 1.4.4.1 in [Gri11]).

Proposition E.2. Let $U \subset \mathbb{R}^2$ be a bounded set with Lipschitz boundary. The following embeddings hold.

- Let $p > 2$ and $n \in \mathbb{N}^*$, then $\mathbb{W}^{n,p}(U) \subset C^n(\bar{U})$.
- Let $p \geq 2$ and $n \in \mathbb{N}^*$, then $\mathbb{W}^{n,p}(U) \subset \mathbb{W}^{\ell,q}(U)$ for any $\ell \leq n - 1$ and any $q \geq 1$.

In our case, $U = \square = \triangleleft \cup \triangleright$, where we recall that $\square = (-1, 1)^2$, $\square = \square \setminus D$ and $\triangleright = \{(x, y) \in \square : x < y\} = \square \setminus \{\triangleleft\}$. We are interested in the subset \mathcal{T}'_s , defined in (2.39), of the topological dual \mathcal{T}' of \mathcal{T} .

Definition E.3. If $(X, \|\cdot\|_X)$ is a Banach space, let X' be its topological dual, equipped with the norm:

$$\forall L \in X', \quad N_X(L) = \sup_{\phi \in X \setminus \{0\}} \frac{|L(\phi)|}{\|\phi\|_X}. \quad (\text{E.3})$$

If $\phi : \square \rightarrow \mathbb{R}$, let $\phi_s(x, y) = [\phi(x, y) + \phi(y, x)]/2$ denote its symmetric part, and let $\mathcal{T}'_s \subset \mathcal{T}'$ be the subset of elements Π satisfying $\Pi(\phi) = \Pi(\phi_s)$ for any $\phi \in \mathcal{T}$. Then:

$$\forall \Pi \in \mathcal{T}'_s, \quad \mathcal{N}_{\mathcal{T}}(\Pi) := \sup_{\phi \in \mathcal{T} \setminus \{0\}} \frac{|\Pi(\phi)|}{\|\phi\|_{\mathbb{H}^2(\square)}} = \sup_{\phi \in \mathcal{T} \setminus \{0\}} \frac{|\Pi(\phi_s)|}{\|\phi_s\|_{\mathbb{H}^2(\square)}} = \sup_{\phi \in \mathcal{T} \setminus \{0\}} \frac{|\Pi(\phi|_{\triangleright})|}{\|\phi|_{\triangleright}\|_{\mathbb{H}^2(\triangleright)}}, \quad (\text{E.4})$$

where $\phi|_{\triangleright}$ is the restriction of f to \triangleright . \mathcal{T}'_s is closed for the norm (E.4).

The weak* topology on \mathcal{T}' is the topology of simple convergence: a sequence $(\Pi_n) \in (\mathcal{T}')^{\mathbb{N}}$ weak* converges to $\Pi \in \mathcal{T}'$ if and only if:

$$\forall \phi \in \mathcal{T}, \quad \lim_{n \rightarrow \infty} \Pi_n(\phi) = \Pi(\phi) \quad \Leftrightarrow \quad \forall \phi \in \mathcal{T}_{\triangleright}, \quad \lim_{n \rightarrow \infty} \Pi_n(\phi|_{\triangleright}) = \Pi(\phi|_{\triangleright}). \quad (\text{E.5})$$

The set \mathcal{T}'_s is also closed for the weak* topology. We write $(\mathcal{T}'_s, *)$ when explicitly referring to this topology.

E.1 Compact sets

In this section, we give a sufficient condition for compactness in $(\mathcal{T}'_s, *)$. Recall that $\mathcal{T} = \mathbb{H}^2(\square)$. Banach-Alaoglu's theorem characterises compactness in $(\mathcal{T}'_s, *)$:

Proposition E.4 (Banach-Alaoglu). *Let $\mathcal{K} \subset (\mathcal{T}'_s, *)$ be such that $\sup_{\mathcal{K}} \mathcal{N}_{\mathcal{T}} < \infty$. Then \mathcal{K} is relatively weak* compact.*

The norm $\mathcal{N}_{\mathcal{T}}$ is defined through a supremum, which is difficult to work with. Instead, we formulate a sufficient condition for compactness which involves a sum. Such a characterisation is known to hold when the underlying space is periodic, e.g. on the torus $\mathbb{T}^2_{-2,2} = [-2, 2]^2$: a linear form $\pi \in (\mathbb{H}^2(\mathbb{T}^2_{-2,2}))'$ is bounded if and only if

$$\|\pi\|_{\mathbb{T}, -2} := \sum_{m \in \mathbb{N}^2} (1 + |m|^2)^{-2} |\pi(\phi_m)|^2 < \infty, \quad (\text{E.6})$$

with $|m|^2 = m_1^2 + m_2^2$ for $m = (m_1, m_2) \in \mathbb{N}^2$, and where $(\phi_m)_{m \in \mathbb{N}^2}$ is an orthonormal basis of $\mathbb{L}^2(\mathbb{T}^2_{-2,2})$ made of real eigenvalues of the Laplacian: for $(x, y) \in \mathbb{T}^2_{-2,2}$,

$$\phi_m(x, y) = \varphi_{m_1}(x)\varphi_{m_2}(y), \quad \varphi_{m_1}(x) = \begin{cases} 1/2 & \text{if } m_1 = 0, \\ 2^{-1/2} \cos\left(\frac{m'\pi x}{2}\right) & \text{if } m_1 = 2m' - 1 \in \mathbb{N}^*, \\ 2^{-1/2} \sin\left(\frac{m'\pi x}{2}\right) & \text{if } m_1 = 2m' \in \mathbb{N}^*. \end{cases} \quad (\text{E.7})$$

The equivalence between (E.6) and $\mathcal{N}_{\mathcal{T}}$ in the periodic setting comes from the fact that a function has the same regularity as its Fourier transform. In our case, however, this property does not hold because of the boundaries in \square , and $\mathcal{N}_{\mathcal{T}}$ is not equivalent to the norm $\|\pi\|_{\mathbb{T}, -2}$.

We look for a sufficient condition for compactness that can nonetheless be stated in terms of the norm $\|\cdot\|_{\mathbb{T}, -2}$ defined in (E.6). To do so, note first that, by (E.4), it is sufficient to work on the triangle $\triangleright = \{(x, y) \in \square : x < y\}$. The idea is then to extend elements $\Pi \in \mathcal{T}'_s$ to linear forms acting on the larger space $\mathbb{H}^2(\mathbb{T}^2_{-2,2})$ of test functions, then check that the norms of Π and its extension are comparable. Define then:

$$\Pi^{\text{ext}}(u) = \Pi(u|_{\triangleright}), \quad u \in \mathbb{H}^2(\mathbb{T}^2_{-2,2}). \quad (\text{E.8})$$

Clearly, Π^{ext} is a linear form on $\mathbb{H}^2(\mathbb{T}^2_{-2,2})$, although it may not be bounded any more. The sufficient condition for compactness can now be stated.

Proposition E.5. *Let $A > 0$, and let $\mathcal{K}_A := \{\Pi \in \mathcal{T}' : \|\Pi^{\text{ext}}\|_{\mathbb{T}, -2} \leq A\}$. Then \mathcal{K}_A is weak* relatively compact in \mathcal{T}'_s .*

Proof. The goal is to bound the original norm $\mathcal{N}_{\mathcal{T}}$ by the norm $\|\cdot\|_{\mathbb{T}, -2}$ from (E.6), then use the Banach-Alaoglu theorem (Proposition E.4) to conclude. In view of (E.4), it is sufficient to work with test functions defined on the triangle. Consider:

$$\mathcal{T}_{\triangleright} = \{\phi|_{\triangleright} : \phi \in \mathcal{T}\}. \quad (\text{E.9})$$

We first explain how to embed $\mathcal{T}_{\triangleright}$ into $\mathbb{H}^2(\mathbb{T}^2_{-2,2})$ and bound $\mathcal{N}_{\mathcal{T}}$ by the strong dual norm $\|\cdot\|_{\mathbb{H}^{-2}(\mathbb{T}^2_{-2,2})}$ (recall (E.3)). Since this norm is equivalent to the norm (E.6), this will be enough to

conclude.

By Theorem 1.4.3.1 in [Gri11], there is a continuous linear extension P from $\mathcal{T}_\triangleright$ to $\mathbb{H}^2(\mathbb{R}^2)$, i.e. there is $C(\triangleright) > 0$ such that:

$$\forall u \in \mathcal{T}_\triangleright, \quad Pu \in \mathbb{H}^2(\mathbb{R}^2), \quad (Pu)|_\triangleright = u, \quad \|Pu\|_{\mathbb{H}^2(\mathbb{R}^2)} \leq C(\triangleright)\|u\|_{\mathbb{H}^2(\triangleright)}. \quad (\text{E.10})$$

Let $\chi \in C^\infty(\mathbb{R}^2)$ be equal to 1 on \triangleright , and be compactly supported in $(-2, 2)^2$. By Theorem 1.4.4.2 in [Gri11], $\mathbb{H}^n(\mathbb{R}^2)$ is a Banach algebra as soon as $n \geq 2$. It follows that there is $C(\chi, \triangleright) > 0$ such that:

$$\begin{aligned} \forall u \in \mathcal{T}_\triangleright, \quad \|\chi Pu\|_{\mathbb{H}^2(\mathbb{R}^2)} &\leq C(\triangleright)\|\chi\|_{\mathbb{H}^2(\mathbb{R}^2)}\|Pu\|_{\mathbb{H}^2(\mathbb{R}^2)} = C(\chi, \triangleright)\|u\|_{\mathbb{H}^2(\triangleright)} \\ &\leq C(\chi, \triangleright)\|\chi Pu\|_{\mathbb{H}^2(\mathbb{R}^2)}, \end{aligned} \quad (\text{E.11})$$

where the last inequality comes from the inclusion $\triangleright \subset (-2, 2)^2$ and the fact that $(\chi Pu)|_\triangleright = u$. The mapping χP is an embedding from $\mathcal{T}_\triangleright$ to $\mathcal{T}_{\triangleright}^{\text{ext}} := \mathbb{H}_0^2((-2, 2)^2)$, the closure of C^∞ , compactly supported functions on $(-2, 2)^2$ for the norm of $\mathbb{H}^2((-2, 2)^2)$. An element u of $\mathcal{T}_{\triangleright}^{\text{ext}}$ can be turned into a periodic function in $\mathbb{H}^2(\mathbb{T}_{-2,2}^2)$ with the same norm, by setting $u(\cdot + (a, b)[-2, 2]^2) = u(\cdot)$ for each $(a, b) \in \mathbb{Z}^2$.

Let us now compare the elements of \mathcal{T}'_s and their extensions to $\mathbb{H}^2(\mathbb{T}_{-2,2}^2)$. Take $\Pi \in \mathcal{T}'_s$, and extend it to a linear form Π^{ext} on $\mathbb{H}^2(\mathbb{T}_{-2,2}^2)$ (possibly unbounded) through (E.8). Then:

$$\begin{aligned} \mathcal{N}_{\mathcal{T}}(\Pi) &:= \sup_{u \in \mathcal{T}_\triangleright \setminus \{0\}} \frac{|\Pi(u)|}{\|u\|_{\mathbb{H}^2(\triangleright)}} = \sup_{u \in \mathcal{T}_{\triangleright}^{\text{ext}} \setminus \{0\}} \frac{|\Pi^{\text{ext}}(\chi Pu)|}{\|u\|_{\mathbb{H}^2(\triangleright)}} \\ &\leq C(\chi, \triangleright) \sup_{v \in \mathcal{T}_{\triangleright}^{\text{ext}} \setminus \{0\}} \frac{|\Pi^{\text{ext}}(v)|}{\|v\|_{\mathbb{H}^2(\mathbb{R}^2)}} \\ &\leq C(\chi, \triangleright) \sup_{w \in \mathbb{H}^2(\mathbb{T}_{-2,2}^2) \setminus \{0\}} \frac{|\Pi^{\text{ext}}(w)|}{\|w\|_{\mathbb{H}^2(\mathbb{T}_{-2,2}^2)}} =: C(\chi, \triangleright)\|\Pi^{\text{ext}}\|_{\mathbb{H}^{-2}(\mathbb{T}_{-2,2}^2)}. \end{aligned} \quad (\text{E.12})$$

The fact that $\|\cdot\|_{\mathbb{H}^{-2}(\mathbb{T}_{-2,2}^2)} \leq c\|\cdot\|_{\mathbb{T}, -2}$ for some $c > 0$ concludes the proof. \square

F Poisson equations

In this section, we give conditions for the existence and uniqueness of solutions of the various Poisson equations – among which the Euler-Lagrange equation (2.52) and the main equation (2.61) – encountered along the text. To do so, we prove that finding a kernel k that solves the (linear) Euler-Lagrange equation or a solution g of the (non-linear) main equation is the same, and rewrite all linear equations in a common framework.

F.1 Euler-Lagrange equation

Throughout Appendix F, a function $u \in \mathbb{L}^2(\square)$ is identified with the kernel operator $u\phi(\cdot) = \int u(\cdot, y)\phi(y) dy$ for $\phi \in \mathbb{L}^2((-1, 1))$. View also $\bar{\sigma}$ as a multiplication operator: if $\phi \in \mathbb{L}^2((-1, 1))$,

$$(\bar{\sigma}u)(x, y) = \bar{\sigma}(x)u(x, y), \quad (u\bar{\sigma})(x, y) = u(x, y)\bar{\sigma}(y), \quad (u\bar{\sigma}\phi)(x) = \int_{(-1, 1)} u(x, y)\bar{\sigma}(y)\phi(y) dy. \quad (\text{F.1})$$

For a symmetric function $f \in \mathcal{T}$, write δ_f for the operator:

$$\delta_f(x) = (\partial_2 - \partial_1)f(x, x_+) = \partial_2 f(x, x_+) - \partial_2 f(x, x_-) = -(\partial_2 - \partial_1)f(x_+, x), \quad x \in (-1, 1). \quad (\text{F.2})$$

The next proposition is classical, and proves Proposition 2.2.

Proposition F.1. *Let $\Pi \in \mathcal{T}'_s$ be associated with a kernel k via $\Pi = \frac{1}{4}\langle k, \cdot \rangle$, and write $C_k := \bar{\sigma} + k$. Assume $\mathcal{I}_\infty(\Pi) < \infty$. There is then a generalised bias $h \in \mathbb{L}^2(\square)$, with h a symmetric function admitting a weak derivative that satisfies:*

$$\int_{(-1,1)} \bar{\sigma}(z) \langle \partial_1 h(z, \cdot), C_k \partial_1 h(z, \cdot) \rangle dz < \infty, \quad (\text{F.3})$$

such that $\mathcal{I}_\infty(k) = J_h(k)$. Moreover, k and h are related through the following Euler-Lagrange equation: for any test function $\phi \in \mathbb{H}^1(\square)$ with $\phi|_{\partial\square} = 0$,

$$\frac{1}{2} \int_{\square} \nabla(k - k_0) \cdot \nabla \phi - \int_{(-1,1)} \bar{\sigma}(z) \langle \partial_1 h(z, \cdot), C_k \partial_1 \phi(z, \cdot) \rangle dz = 0. \quad (\text{F.4})$$

This is a weak formulation of:

$$\left\{ \begin{array}{l} \Delta k(x, y) - [\bar{\sigma}(x)\delta_h(x) + \bar{\sigma}(y)\delta_h(y)]k(x, y) \\ \quad - \int_{(-1,1)} \left[\partial_x(\bar{\sigma}(x)\partial_1 h(x, z))k(z, y) + k(x, z)\partial_y(\partial_2 h(z, y)\bar{\sigma}(y)) \right] dz \\ \quad = \partial_x(\bar{\sigma}(x)\bar{\sigma}(y)\partial_1 h(x, y)) + \partial_y(\bar{\sigma}(x)\bar{\sigma}(y)\partial_2 h(x, y)) \quad \text{for } (x, y) \in \square, \\ h|_{\partial\square} = k|_{\partial\square} = 0, \\ \bar{\sigma}(x)^2(\partial_2 - \partial_1)h(x_+, x) - (\partial_2 - \partial_1)k(x_+, x) = (\bar{\rho}')^2 \quad \text{for } x \in (-1, 1). \end{array} \right. \quad (\text{F.5})$$

Proof. Write for short $\mathcal{I}_\infty(k)$ for $\mathcal{I}_\infty(\Pi)$, and idem for $J_f(k)$, $f \in \mathcal{S}(\varepsilon)$. The existence (and uniqueness) of the generalised bias h such that $J_h(k) = \mathcal{I}_\infty(k)$ is classical, see Lemma 5.3 in Chapter 10 of [KL99]. Let us show that k satisfies the Euler-Lagrange equation with bias h . Since h is such that $J_h(k) = \inf_{\tilde{h} \in \mathcal{S}(\infty)} J_{\tilde{h}}(k)$, one has $\varepsilon^{-1}(J_{h \pm \varepsilon \phi} - J_h) \leq 0$ for any $\phi \in \mathbb{H}^1(\square)$ with $\phi|_{\partial\square} = 0$ and small enough $\varepsilon > 0$. Thus:

$$\begin{aligned} & \frac{1}{8} \int_{\square} \nabla k \cdot \nabla \phi + \frac{1}{4} \int_{(-1,1)} (\bar{\rho}')^2 \phi(x, x) dx \\ & - \frac{1}{4} \int_{(-1,1)} \bar{\sigma}(z) \langle \partial_1 h(z, \cdot), (\bar{\sigma} + k)(\partial_1 \phi(z, \cdot)) \rangle dz = 0. \end{aligned} \quad (\text{F.6})$$

Applying this equation in the $h = 0$ case corresponding to the steady-state kernel k_0 (or just recalling the expression (2.11)), one gets:

$$\frac{1}{8} \int_{\square} \nabla k_0 \cdot \nabla \phi = -\frac{1}{4} \int_{(-1,1)} (\bar{\rho}')^2 \phi(x, x) dx. \quad (\text{F.7})$$

This yields (F.5). Equation (2.52) follows by careful integration by parts. For instance, the term $\bar{\sigma}(x)\delta_h(x)k(x, y)$ comes from the contribution of k in the integral involving $C_k = \bar{\sigma} + k$:

$$\begin{aligned} & \int_{(-1,1)^2} \bar{\sigma}(z)\partial_1 h(z, x)(k\partial_1\phi)(z, x) dx dz = \int_{(-1,1)^2} \bar{\sigma}(z)\partial_1 h(z, x)\partial_1(k\phi)(z, x) dx dz \\ & = \int_{(-1,1)} \bar{\sigma}(x) \left[\partial_1 h(x_-, x) - \partial_1 h(x_+, x) \right] (k\phi)(x, x) dx \\ & \quad - \int_{(-1,1)^2} \partial_1(\bar{\sigma}(z)\partial_1 h(z, x))(k\phi)(z, x) dx dz. \end{aligned} \quad (\text{F.8})$$

Note that $k\phi(x, x)$ is simply $\int_{(-1,1)^2} k(x, y)\phi(x, y) dx dy$. Using the symmetry of h and the fact that $\partial_1 h(x_-, x) = \partial_1 h(x, x_+)$ (it is the same point in the same triangle), one has:

$$\partial_1 h(x_-, x) - \partial_1 h(x_+, x) = (\partial_2 - \partial_1)h(x_+, x) = -\delta_h(x), \quad x \in (-1, 1). \quad (\text{F.9})$$

□

F.2 Equivalence of the Euler-Lagrange equation and the main equation

Proposition F.2. *Let $h \in \mathcal{S}(\infty)$, where this set is defined in (2.45). Recall that $\square := (-1, 1)^2$, $\boxminus = \square \setminus D$ and $\triangleright = \{(x, y) \in \square : x < y\}$, $\triangleleft = \square \setminus \triangleright$.*

- *Let $k \in C^3(\triangleright) \cap C^3(\triangleleft)$ be symmetric and suppose k solves the Euler-Lagrange equation (F.5). Assume that k satisfies:*

$$\int_{\boxminus} \bar{\sigma}^{-1}(x)k(x, y)^2\bar{\sigma}^{-1}(y) dx dy < 1. \quad (\text{F.10})$$

Then the correlation operator $C_k = \bar{\sigma} + k$ is invertible in $\mathbb{L}^2(\boxminus)$. Define $g \in \mathbb{L}^2(\boxminus)$ through $C_k = \bar{\sigma} + k = (\bar{\sigma}^{-1} - g)^{-1}$. Then $g \in C^3(\triangleright) \cap C^3(\triangleleft)$ is symmetric and solves the main equation (2.61).

- *Conversely, let $g \in C^3(\triangleright) \cap C^3(\triangleleft)$ be symmetric and solve the main equation (2.61), and assume:*

$$\int \bar{\sigma}(x)g(x, y)^2\bar{\sigma}(y) dx dy < 1. \quad (\text{F.11})$$

Then $\bar{\sigma}^{-1} - g$ is invertible, and $k \in C^3(\triangleright) \cap C^3(\triangleleft)$ defined through $\bar{\sigma} + k = (\bar{\sigma}^{-1} - g)^{-1}$ is symmetric and solves the Euler-Lagrange equation (F.5).

Remark F.3. *Conditions (F.10)–(F.11) hold if $\bar{\rho}' \leq \varepsilon, h \in \mathcal{S}(\varepsilon)$ for small enough ε . ■*

Proof. We only show that $k \in C^3(\triangleright) \cap C^3(\triangleleft)$ symmetric solving the Euler-Lagrange equation implies that g solves the main equation, is symmetric and is in $C^3(\triangleright) \cap C^3(\triangleleft)$, the other implication being similar.

By assumption (F.10), g admits a series expansion:

$$\begin{aligned} g & = -\bar{\sigma}^{-1/2}(1 + \bar{\sigma}^{-1/2}k\bar{\sigma}^{-1/2})^{-1}\bar{\sigma}^{-1/2} + \bar{\sigma}^{-1} = -\bar{\sigma}^{-1/2} \circ \sum_{n \geq 1} (-1)^n (\bar{\sigma}^{-1/2}k\bar{\sigma}^{-1/2})^{\circ n} \circ \bar{\sigma}^{-1/2} \\ & = -\bar{\sigma}^{-1} \circ \sum_{n \geq 1} (k \circ \bar{\sigma}^{-1})^{\circ n}, \end{aligned} \quad (\text{F.12})$$

where \circ denotes composition and $\circ n$ n -times composition, $n \geq 1$. In particular, g is symmetric, satisfies (F.11), and it inherits the regularity of k : $g \in C^3(\overline{\square}) \cap C^3(\overline{\square})$. Moreover, (F.12) already shows that $g|_{\partial \square} = 0$ if $k|_{\partial \square} = 0$.

To check the Neumann condition on the diagonal and the fact that g satisfies the main equation (2.61), let us write derivatives of k in terms of g . We henceforth drop the symbol \circ . Using the inverse operator $C_k^{-1} = (\bar{\sigma}^{-1} - g)$, it holds that:

$$(\bar{\sigma} + k)(\bar{\sigma}^{-1} - g) = \text{id} \quad \Rightarrow \quad k\bar{\sigma}^{-1} = C_k g \quad \Rightarrow \quad k = C_k g \bar{\sigma}. \quad (\text{F.13})$$

In the same way $k = \bar{\sigma} g C_k$. Differentiating k with respect to the second variable (or alternatively integrating against the derivative ϕ' of a test function ϕ on $(-1, 1)$ and integrating by parts) yields:

$$\partial_2 k = \partial_2(C_k g \bar{\sigma}) = C_k \partial_2(g \bar{\sigma}). \quad (\text{F.14})$$

Note that since k is symmetric, $\partial_2 k(x, y) = \partial_1 k(y, x)$ for $(x, y) \in \square$ means that $\partial_1 k$ is the operator adjoint to $\partial_2 k$ in $\mathbb{L}^2((-1, 1))$, thus one has also:

$$\partial_1 k = \partial_1(\bar{\sigma} g) C_k. \quad (\text{F.15})$$

In particular, this gives the Neumann condition for g on the diagonal:

$$(\partial_1 - \partial_2)k(x_+, x) = \bar{\sigma}(x)^2(\partial_1 - \partial_2)g(x_+, x), \quad x \in (-1, 1). \quad (\text{F.16})$$

Let us now prove that g satisfies the main equation (2.61). Recall the convention: if $\phi : \square \rightarrow \mathbb{R}$ and $q : (-1, 1) \rightarrow \mathbb{R}$, then ϕq is the function $(x, y) \mapsto \phi(x, y)q(y)$, while $(q\phi)(x, y) := q(x)\phi(y)$. The Euler-Lagrange equation (F.5) then reads:

$$\Delta k - \left[\partial_1(\bar{\sigma} \partial_1 h) C_k + C_k \partial_2(\partial_2 h \bar{\sigma}) \right] - [\bar{\sigma} \delta_h k + k \bar{\sigma} \delta_h] = 0, \quad (\text{F.17})$$

where the above is an equality between functions on \square . In (F.17), the right-hand side in (F.5) has been included in the term in the bracket by using the operator $C_k = \bar{\sigma} + k$. The partial derivatives $\partial_2 k, \partial_1 k$ were obtained in (F.14)–(F.15). As k is not C^2 across the diagonal, taking the second derivative is more subtle, and we do so against a test function $\phi \in C^\infty(\square)$ vanishing on $\partial \square$:

$$-\partial_2 k \partial_2 \phi = \partial_2^2 k \phi + \delta_k \phi = -C_k \partial_2(g \bar{\sigma}) \partial_2 \phi = C_k \partial_2^2(g \bar{\sigma}) + C_k \bar{\sigma} \delta_g \phi. \quad (\text{F.18})$$

Above, recall that, for a symmetric function $\psi : \square \rightarrow \mathbb{R}$, δ_ψ is the function operator:

$$\begin{aligned} \delta_\psi(x) &= (\partial_2 - \partial_1)\psi(x, x_+) = (\partial_1 - \partial_2)\psi(x_+, x) \\ &= \partial_1\psi(x_+, x) - \partial_1\psi(x_-, x) = \partial_2\psi(x, x_+) - \partial_2\psi(x, x_-), \quad x \in (-1, 1). \end{aligned} \quad (\text{F.19})$$

Since $\delta_k = \bar{\sigma}^2 \delta_g$, we get:

$$\partial_2^2 k = C_k \partial_2^2(g \bar{\sigma}) + k \bar{\sigma} \delta_g. \quad (\text{F.20})$$

By symmetry, one has also

$$\partial_1 k = \partial_1(\bar{\sigma} g) C_k \quad \Rightarrow \quad \partial_1^2 k = \partial_1^2(\bar{\sigma} g) C_k + \delta_g \bar{\sigma} k. \quad (\text{F.21})$$

Thus:

$$\Delta k = (\bar{\sigma} \partial_1^2 g + 2\bar{\sigma}' \partial_1 g + \bar{\sigma}'' g) C_k + C_k (\partial_2^2 g \bar{\sigma} + 2\partial_2 g \bar{\sigma}' + g \bar{\sigma}'') + \delta_g \bar{\sigma} k + k \bar{\sigma} \delta_g. \quad (\text{F.22})$$

Compose (F.17) by C_k^{-1} on each side. Using $C_k^{-1} = \bar{\sigma}^{-1} - g$, this yields for Δk , still as an identity between functions on \mathbb{Z} :

$$C_k^{-1} \Delta k C_k^{-1} = \bar{\sigma}^{-1} \partial_1^2 (\bar{\sigma} g) - g \partial_1^2 (\sigma g) + \partial_2^2 (g \bar{\sigma}) \bar{\sigma}^{-1} - \partial_2^2 (g \bar{\sigma}) g + C_k^{-1} (\delta_g \bar{\sigma} k + k \bar{\sigma} \delta_g) C_k^{-1}. \quad (\text{F.23})$$

From $k C_k^{-1} = \bar{\sigma} g$ and the following identities obtained by integration by parts

$$-g \partial_1^2 (\sigma g) = \partial_2 g \bar{\sigma} \partial_1 g + g \bar{\sigma} \delta_g + \partial_2 g \bar{\sigma}' g, \quad (\text{F.24})$$

we obtain:

$$\begin{aligned} C_k^{-1} \Delta k C_k^{-1} &= \Delta g + 2 \frac{\bar{\sigma}'}{\bar{\sigma}} \partial_1 g + 2 \partial_2 g \frac{\bar{\sigma}'}{\bar{\sigma}} + 2 \partial_2 g \bar{\sigma} \partial_1 g + g \bar{\sigma} \delta_g + \bar{\sigma} \delta_g g + g \frac{\bar{\sigma}''}{\bar{\sigma}} + \frac{\bar{\sigma}''}{\bar{\sigma}} g \\ &\quad + \partial_2 g \bar{\sigma}' g + g \bar{\sigma}' \partial_1 g + \delta_g \bar{\sigma} g + g \delta_g \bar{\sigma} - 2g \delta_g \bar{\sigma}^2 g. \end{aligned} \quad (\text{F.25})$$

We next show that:

$$\begin{aligned} C_k^{-1} \left[\partial_1 (\bar{\sigma} \partial_1 h) C_k + C_k \partial_2 (\partial_2 h \bar{\sigma}) \right] C_k^{-1} &= \Delta h + \frac{\bar{\sigma}'}{\bar{\sigma}} \partial_1 h + \partial_2 h \frac{\bar{\sigma}'}{\bar{\sigma}} \\ &\quad + \partial_2 g \bar{\sigma} \partial_1 h + \partial_2 h \bar{\sigma} \partial_1 g + g \bar{\sigma} \delta_h + \delta_h \bar{\sigma} g. \end{aligned} \quad (\text{F.26})$$

This follows by noticing that

$$C_k^{-1} \partial_1 (\bar{\sigma} \partial_1 h) = \bar{\sigma}^{-1} (\bar{\sigma}' \partial_1 h + \bar{\sigma} \partial_1^2 h) - g \partial_1 (\bar{\sigma} \partial_1 h) = \frac{\bar{\sigma}'}{\bar{\sigma}} \partial_1 h + \partial_1^2 h + \partial_1 g \bar{\sigma} \partial_1 h + g \bar{\sigma} \delta_h, \quad (\text{F.27})$$

where in the last step, we used an integration by parts as in (F.24). Part of the boundary terms cancels thanks to the final identity:

$$\begin{aligned} C_k^{-1} [\bar{\sigma} \delta_h k + k \bar{\sigma} \delta_h] C_k^{-1} &= (\bar{\sigma}^{-1} - g) \bar{\sigma} \delta_h + \bar{\sigma} \delta_h (\bar{\sigma}^{-1} - g) - 2(\bar{\sigma}^{-1} - g) \bar{\sigma} \delta_h \bar{\sigma} (\bar{\sigma}^{-1} - g) \\ &= 2\delta_h - g \bar{\sigma} \delta_h - \bar{\sigma} \delta_h g - 2(1 - g \bar{\sigma}) \delta_h (1 - \bar{\sigma} g) = g \bar{\sigma} \delta_h + \bar{\sigma} \delta_h g - 2g \bar{\sigma}^2 \delta_h g. \end{aligned} \quad (\text{F.28})$$

Plugging the relations (F.25), (F.26), (F.28) in the Euler-Lagrange equation (F.17), we get:

$$\begin{aligned} 0 &= \Delta(g - h) + \frac{\bar{\sigma}'}{\bar{\sigma}} \partial_1 (2g - h) + \partial_2 (2g - h) \frac{\bar{\sigma}'}{\bar{\sigma}} + \partial_2 (g - h) \bar{\sigma} \partial_1 g + \partial_2 g \bar{\sigma} \partial_1 (g - h) \\ &\quad + 2g \bar{\sigma}^2 \delta_{h-g} g + g \bar{\sigma} \left(\frac{\bar{\sigma}''}{\bar{\sigma}^2} + 2\delta_g - 2\delta_h \right) + \bar{\sigma} \left(\frac{\bar{\sigma}''}{\bar{\sigma}^2} + 2\delta_g - 2\delta_h \right) g + \partial_2 g \bar{\sigma}' g + g \bar{\sigma}' \partial_1 g. \end{aligned} \quad (\text{F.29})$$

In addition, the boundary conditions of (F.5) imply that $\delta_{g-h} = -\frac{\bar{\sigma}''}{2\bar{\sigma}^2}$ (recall $\bar{\sigma}'' = -2(\bar{\rho}')^2$). Finally, an integration by parts gives:

$$\partial_2 g \bar{\sigma}' g = -g \bar{\sigma}'' g - g \bar{\sigma}' \partial_1 g. \quad (\text{F.30})$$

The second line in (F.29) thus vanishes, and the first line is precisely the main equation (2.61). \square

In the next sections, we focus on establishing existence, uniqueness and regularity for solutions of the Euler-Lagrange equation (F.5). In view of Proposition F.1, this will prove Proposition 2.5.

F.3 Existence and uniqueness

We now focus on solving the Euler-Lagrange equation (F.5) and the Poisson equation (5.6). To do so, we rewrite them in a common framework. Both equations are formulated as equations on \square involving symmetric functions. To solve them, it is therefore enough to look at the equation in a single triangle, say \triangleright . To do so, let us introduce some notations.

Recall that the function \mathcal{M} acts on $\phi, \psi \in \mathbb{L}^2(\square)$ according to:

$$\mathcal{M}(\phi, \psi)(x, y) = \int_{-1}^1 \phi(x, z) \bar{\sigma}(z) \psi(y, z) dz, \quad (x, y) \in \square. \quad (\text{F.31})$$

Define by extension \mathcal{M} on the triangle as follows: if $(\phi, \psi) \in \mathbb{L}^2(\triangleright)^2$,

$$\mathcal{M}(\phi, \psi) := \mathcal{M}(\phi_s, \psi_s), \quad \phi_s(x, y) = \begin{cases} \phi(x, y) & \text{if } x < y, \\ \phi(y, x) & \text{if } x > y. \end{cases} \quad (\text{F.32})$$

For $\phi \in \mathbb{L}^2(\triangleright)$ and $q : (-1, 1) \rightarrow \mathbb{R}$, we often write:

$$(qf)(x, y) := q(x)f(x, y), \quad (fq)(x, y) := f(x, y)q(y), \quad (x, y) \in \triangleright. \quad (\text{F.33})$$

Given symmetric $\phi, \psi, \xi : \square \rightarrow \mathbb{R}$ and $d : (-1, 1) \rightarrow \mathbb{R}$, we say that $f : \square \rightarrow \mathbb{R}$ solves the Poisson problem $(P) = (P_{\phi, \psi, \xi, d})$ if f is symmetric, and $f|_{\triangleright}$ is a classical solution of (P_{\triangleright}) , where:

$$(P_{\triangleright}) : \begin{cases} \Delta f(x, y) + [d(x) + d(y)]f(x, y) - \mathcal{M}(f, \xi) & \text{for } (x, y) \in \triangleright, \\ -\mathcal{M}(\xi, f) - \mathcal{M}(\partial_1 f, \partial_1 \psi) - \mathcal{M}(\partial_1 \psi, \partial_1 f) = \phi(x, y) \\ \partial_{\nu_{\triangleright}} f = 0 \quad \text{on } D, \\ f = 0 \quad \text{on } (\partial \triangleright) \setminus D. \end{cases} \quad (\text{F.34})$$

where $\partial_{\nu_{\triangleright}}$ stands for the normal derivative on the diagonal.

Remark F.4. (i) If $\psi = h \in \mathcal{S}(\infty)$ (recall (2.44)), $d = 0$, $\xi = 0$ and $\phi \in \mathcal{T} \cap C^2(\bar{\triangleright})$ has norm 2, then (P) corresponds to the Poisson problem encountered in the proof of large deviations, in (5.6).

(ii) Let $h \in \mathcal{S}(\infty)$, $\psi = 0$, $d = \bar{\sigma} \delta_h$ with δ_h given in (F.19). Take ξ as follows:

$$\xi(x, y) = \partial_1(\bar{\sigma} \partial_1 h)(x, y) = \partial_x(\bar{\sigma}(x) \partial_1 h(x, y)), \quad (x, y) \in \square. \quad (\text{F.35})$$

Define then $\phi(x, y)$ for $(x, y) \in \square$ as:

$$\begin{aligned} \phi(x, y) &= -\partial_1(\bar{\sigma} h)(x, y) - \partial_2(h \bar{\sigma}')(x, y) + \mathcal{M}(k_0 + \bar{\sigma} h \bar{\sigma}, \xi)(x, y) \\ &\quad + \mathcal{M}(k_0 + \bar{\sigma} h \bar{\sigma}, \xi)(x, y) - [k_0(x, y) + (\bar{\sigma} h \bar{\sigma})(x, y)] [\bar{\sigma}(x) \delta_h(x) + \bar{\sigma}(y) \delta_h(y)]. \end{aligned} \quad (\text{F.36})$$

Then (P) is the Euler-Lagrange equation (F.5) written with unknown $f = k - k_0 - \bar{\sigma} h \bar{\sigma}$. Note that f is chosen so that the boundary conditions in (F.34) imply the Neumann condition $\partial_{\nu_{\triangleright}} f = 0$. ■

In the remainder of the section, we study existence, uniqueness and regularity of solutions of (P) . For $\star \in \{\triangleright, \triangleleft\}$, we write $\langle \cdot, \cdot \rangle_\star$ for the usual scalar product on $\mathbb{L}^2(\star)$, and simply $\langle \cdot, \cdot \rangle$ for the scalar product on $\mathbb{L}^2(\square) = \mathbb{L}^2(\square)$. The norm on $\mathbb{L}^2(\square)$ is denoted by $\|\cdot\|_2$. Let also tr denote the trace operator on the boundary of \square . When interested only in a portion Γ of the boundary, we may write tr_Γ .

We will use the fact that the Laplacian with 0 Dirichlet boundary condition on $(\partial\triangleright) \setminus D$ and 0 Neumann boundary condition on the diagonal D has a gap $\alpha > \pi^2/4 > 0$, see e.g. Equation 5 in Section 3.3. of [Siu16]. This means that, for any symmetric $f \in \mathcal{T}$ satisfying the boundary conditions of (P_\triangleright) , one has:

$$\|f|_{\triangleright}\|_{\triangleright}^2 \leq \alpha^{-1} \|\nabla f|_{\triangleright}\|_{\triangleright}^2 \quad \Rightarrow \quad \|f\|_2^2 \leq \alpha^{-1} \|\nabla f\|_2^2. \quad (\text{F.37})$$

We first obtain existence and uniqueness of solutions of (P) in the set $\mathcal{T}_{(P)} \subset \mathcal{T} = \mathbb{H}^2(\square)$ of functions satisfying the boundary conditions of (P) by a fixed point argument. The set $\mathcal{T}_{(P)}$ and its counterpart $\mathcal{T}_{(P_\triangleright)}$ for functions on \triangleright are defined as follows:

$$\begin{aligned} \mathcal{T}_{(P)} &= \mathbb{H}^2(\square) \cap \{f : \text{tr}_D(\partial_{\nu_\triangleright} f) = 0, \text{tr}_{\partial\square}(f) = 0\}, \\ \mathcal{T}_{(P_\triangleright)} &= \mathbb{H}^2(\triangleright) \cap \{f : \text{tr}(f) = 0 \text{ on } \partial(\triangleright) \setminus \bar{D}, \text{tr}_D(\partial_{\nu_\triangleright} f) = 0\} = \{f|_{\triangleright} : f \in \mathcal{T}_{(P)}\}. \end{aligned} \quad (\text{F.38})$$

Proposition F.5 (Solving (P)). *Let $\phi, \xi \in \mathbb{L}^2(\square)$, $\psi \in \mathcal{T}$ be symmetric functions. Let also $d : (-1, 1) \rightarrow \infty$ be bounded. For $f \in \mathcal{T}_{(P)}$, define Sf as the symmetric function such that, for $(x, y) \in \triangleright$,*

$$\begin{aligned} Sf(x, y) &= (-\Delta_\triangleright)^{-1} \left[-\phi + df + fd - \mathcal{M}(f, \xi) - \mathcal{M}(\xi, f) \right. \\ &\quad \left. - \mathcal{M}(\partial_1 f, \partial_1 \psi) - \mathcal{M}(\partial_1 \psi, \partial_1 f) \right] (x, y). \end{aligned} \quad (\text{F.39})$$

Above, $\Delta_\triangleright^{-1}$ is the inverse of the Laplacian on \triangleright with 0 Dirichlet condition on $(\partial\triangleright) \setminus D$, and 0 Neumann conditions on the diagonal D .

Then $Sf \in \mathcal{T}_{(P)}$. Moreover, if $d = 0$, $\xi = 0$ and $\|\nabla\psi\|_2 \leq 1$, $\|\phi\|_2 \leq 2$, then S has a unique fixed point $f_{\phi, \psi} \in \mathcal{T}_{(P)}$ with $\|f_{\phi, \psi}\|_{\mathbb{H}^1(\square)} \leq C$ and $\|f_{\phi, \psi}\|_\infty \leq C'$, for constants $C, C' > 0$ independent of ϕ, ψ .

If instead $\psi = 0$, there is a fixed point $f_{d, \phi, \xi} \in \mathcal{T}_{(P)}$ provided $\|\xi\|_2, \|d\|_\infty$ are small enough, with $\|f_{d, \phi, \xi}\|_{\mathbb{H}^1(\square)} \leq \delta(\|d\|_\infty, \|\xi\|_2, \|\phi\|_2)$ and δ vanishes when $(\|d\|_\infty, \|\xi\|_2, \|\phi\|_2)$ vanishes.

Proof. We prove that S is a contraction on $\mathcal{T}_{(P)}$ for the norm $\|\nabla \cdot\|_2$. Let us start by showing that S is well defined. The inverse operator $\Delta_\triangleright^{-1}$ exists by Lemma 4.4.3.1 in [Gri11] and, by Theorem 4.4.3.7 in [Gri11], maps $\mathbb{L}^2(\triangleright)$ onto $\mathcal{T}_{(P_\triangleright)}$. It follows that $S(\mathcal{T}_{(P)}) \subset \mathcal{T}_{(P)}$.

We now prove that S is a contraction. For $f \in \mathcal{T}_{(P)}$, one has:

$$\begin{aligned} \|\nabla Sf\|_2^2 &= \langle (\nabla Sf)|_{\triangleright}, (\nabla Sf)|_{\triangleright} \rangle_{\triangleright} + \langle (\nabla Sf)|_{\triangleleft}, (\nabla Sf)|_{\triangleleft} \rangle_{\triangleleft} \\ &= -\langle (Sf)|_{\triangleright}, (\Delta Sf)|_{\triangleright} \rangle_{\triangleright} - \langle (Sf)|_{\triangleleft}, (\Delta Sf)|_{\triangleleft} \rangle_{\triangleleft} = -\langle Sf, \Delta Sf \rangle. \end{aligned} \quad (\text{F.40})$$

The integration by parts is legitimate by Theorem 1.5.3.1 in [Gri11]. Let us compute the right-hand

side. One has, using $\bar{\sigma} \leq 1/4$:

$$\begin{aligned} \forall(x, y) \in \square, \quad \mathcal{M}(f, \xi)(x, y) &= \int_{(-1,1)} f(z, x) \bar{\sigma}(z) \xi(z, y) dz \\ &\leq \frac{1}{4} \left(\int_{(-1,1)} f(z, x)^2 dz \right)^{1/2} \left(\int_{(-1,1)} \psi(z, y)^2 dz \right)^{1/2}, \end{aligned} \quad (\text{F.41})$$

and the same holds for $\mathcal{M}(\partial_1 f, \partial_1 \psi)$. As a result, by Cauchy-Schwarz inequality and using $\|\partial_1 f\|_2 = 2^{-1/2} \|\nabla f\|_2$ as implied by the symmetry of f :

$$\begin{aligned} \left| \langle Sf, \mathcal{M}(f, \xi) \rangle \right| &\leq \frac{1}{4} \|Sf\|_2 \|f\|_2 \|\xi\|_2, \\ \left| \langle Sf, \mathcal{M}(\partial_1 f, \partial_1 \psi) \rangle \right| &\leq \frac{1}{8} \|Sf\|_2 \|\nabla f\|_2 \|\nabla \psi\|_2. \end{aligned} \quad (\text{F.42})$$

Recalling the expression of Sf from (F.39), (F.40) is therefore bounded as follows:

$$\|\nabla Sf\|_2^2 \leq \|Sf\|_2 \left(\|\phi\|_2 + 2\|d\|_\infty \|f\|_2 + \frac{\|f\|_2 \|\xi\|_2}{4} + \frac{\|\nabla f\|_2 \|\nabla \psi\|_2}{8} \right). \quad (\text{F.43})$$

Since $f, Sf \in \mathcal{T}_{(P)}$, the Poincaré inequality F.37 can be applied and yields:

$$\|\nabla Sf\|_2 \leq \alpha^{-1/2} \left(\|\phi\|_2 + 2\alpha^{-1/2} \|d\|_\infty \|\nabla f\|_2 + \frac{\alpha^{-1/2} \|\nabla f\|_2 \|\xi\|_2}{4} + \frac{\|\nabla f\|_2 \|\nabla \psi\|_2}{8} \right). \quad (\text{F.44})$$

By similar computations, if $f_1, f_2 \in \mathcal{T}_{(P)}$, one obtains:

$$\|\nabla(Sf_1 - Sf_2)\|_2 \leq \alpha^{-1/2} \left(2\alpha^{-1/2} \|d\|_\infty + \frac{\alpha^{-1/2} \|\xi\|_2}{4} + \frac{\|\nabla \psi\|_2}{8} \right) \|\nabla(f_1 - f_2)\|_2. \quad (\text{F.45})$$

In particular, S is a contraction as soon as:

$$\alpha^{-1/2} \left(2\alpha^{-1/2} \|d\|_\infty + \frac{\alpha^{-1/2} \|\xi\|_2}{4} + \frac{\|\nabla \psi\|_2}{8} \right) < 1. \quad (\text{F.46})$$

Recall that $\alpha > \pi^2/4$. If $d = 0$, $\xi = 0$, $\|\nabla \psi\|_2 \leq 1$ and $\|\phi\|_2 \leq 2$ (which includes case (i) of Remark F.4), then:

$$\|\nabla S(f_1 - f_2)\|_2 \leq \frac{1}{4\pi} \|\nabla(f_1 - f_2)\|_2, \quad \|\nabla Sf\|_2 \leq \frac{1}{4\pi} \|\nabla f\|_2 + \frac{4}{\pi}. \quad (\text{F.47})$$

There is thus a unique fixed point $f_{\phi, \psi} \in \mathcal{T}_{(P)}$, and since $\|\phi\|_2 \leq 2$, it belongs to the ball $B(0, c) = \{u \in \mathbb{H}^1(\square) : \text{tr}(u) = 0 \text{ on } \partial \square, \|\nabla u\|_2 \leq c\}$ with $c = 8/(2\pi - 1) \leq 2$. Poincaré inequality (F.37) yields $\|f_{\phi, \psi}\|_{\mathbb{H}^1(\square)} \leq (4 + 16/\pi^2)^{1/2} \leq 2$, and Theorem 4.3.1.4 in [Gri11] yields $\|f_{\phi, \psi}\|_{\mathbb{H}^2(\square)} \leq C$ for some universal $C > 0$. The Sobolev embedding $\mathbb{H}^2(\square) \subset C^0(\bar{\square}) \cap C^0(\bar{\square})$ then implies $\|f_{\phi, \psi}\|_\infty \leq C'$ for a universal $C' > 0$ as claimed.

Consider now the case $\psi = 0$ which includes item (ii) of Remark F.4. Then:

$$\begin{aligned} \|\nabla S(f_1 - f_2)\|_2 &\leq \frac{4}{\pi^2} \left(2\|d\|_\infty + \frac{\|\xi\|_2}{4} \right) \|\nabla(f_1 - f_2)\|_2, \\ \|\nabla Sf\|_2 &\leq \frac{4}{\pi^2} \left(2\|d\|_\infty + \frac{\|\xi\|_2}{4} \right) \|\nabla f\|_2 + \frac{2\|\phi\|_2}{\pi}. \end{aligned} \quad (\text{F.48})$$

If $\|d\|_\infty, \|\xi\|_2$ are small enough, S is a contraction and the norm of its fixed point vanishes when $\|d\|_\infty, \|\xi\|_2$ and $\|\phi\|_2$ vanish. This concludes the proof. \square

F.4 Regularity estimates

In Proposition F.5, the solution of (P) has been shown to be in $\mathbb{H}^2(\square)$. In this section, we use results of [Gri11] to argue that the solution of (P) is more regular if the data ϕ, ψ are regular. This concludes the proof of Proposition 2.5. The study of regularity is made very complicated by the presence of corners.

Proposition F.6 (Theorem 5.1.3.1. in [Gri11]). *Let $b \in \mathbb{N}$, $p > 2$ and let $\zeta \in \mathbb{W}^{b,p}(\triangleright)$. Let S_1, S_2, S_3 denote the corners of \triangleright numbered in a counter-clockwise fashion, with S_1 the upper left corner. Consider on \triangleright the equation $\Delta f = \zeta$, with the boundary conditions of (P_\triangleright) . If $b = 0$, then $f \in \mathbb{W}^{2,p}(\triangleright)$. If $b \leq 3$, $f \in \mathbb{W}^{b+2,p}(\triangleright)$ provided ζ vanishes at the corners, i.e. provided:*

$$\forall j \in \{1, 2, 3\}, \quad \zeta(S_j) = 0. \quad (\text{F.49})$$

Remark F.7. *Though the statement of Proposition F.6 makes no mention of them, we recall notations from [Gri11] so that the reader may check that Theorem 5.1.3.1 applies to our case. Label by $j \in \{1, \dots, 3\}$ the line segments composing $\partial\triangleright$ in a counter clockwise fashion, with the convention that $j = 1$ for the $y = -1$ segment. S_j is then the point joining segments $j, j + 1$ in $\partial\triangleright$. Let ω_j be the counter-clockwise measure of the inwards angle at S_j :*

$$\omega_j = \begin{cases} \pi/2 & \text{if } j = 1, \\ \pi/4 & \text{if } j \in \{2, 3\}. \end{cases} \quad (\text{F.50})$$

Let $\nu_j = \mu_j$ denote the unit outwards normal and τ_j be the (counter clockwise) unit tangent vector on the line segment j . Define also $\Phi_j = \pi/2$ if $j \in \{1, 2\}$, $\Phi_j = 0$ if $j = 3$ and $\Phi_{3+1} := \Phi_1$. Finally, for $m \in \mathbb{Z}$ and each j , define:

$$\lambda_{j,m} = \frac{\Phi_j - \Phi_{j+1} + m\pi}{\omega_j} = \begin{cases} 2m & \text{if } j = 1, \\ 2 + 4m & \text{if } j = 2, \\ -2 + 4m & \text{if } j = 3. \end{cases} \quad (\text{F.51})$$

■

In our context, Proposition F.6 translates to the following result.

Proposition F.8 (Regularity of solutions of (P_\triangleright)). *Let ϕ, ψ, d, ξ be such that the solution f given by Proposition F.5 exists.*

- (i) *Assume $d = 0$, $\xi = 0$, let $\phi \in C^2(\bar{\triangleright}) \cap C^2(\bar{\triangleleft})$ be symmetric and let $\psi \in \mathcal{S}(\infty)$, corresponding to item (i) of Remark F.4. If $\phi|_{\partial\square} = 0$, then $f \in \mathbb{W}^{4,p}(\square)$ for any $p > 2$.*
- (ii) *Recall the definition (2.44) of $\mathcal{S}(\varepsilon)$ for $\varepsilon > 0$. Take $\psi = 0$, and d, ϕ, ξ defined in terms of $h \in \mathcal{S}(\varepsilon)$ as in item (ii) of Remark F.4. Then $f \in \mathbb{W}^{4,p}(\square)$ for any $p > 2$, and $k = f + k_0 + \bar{\sigma}h\bar{\sigma} \in k_0 + \mathcal{S}(\varepsilon')$ for some $\varepsilon' > 0$ depending only on ρ_\pm and ε .*

Proof. Since f is symmetric, we work only on \triangleright . Let us first assume f has the alleged regularity and treat all claims of item 2 that do not have to do with the regularity of f .

Define:

$$\|k - k_0\|_{C^1} := \max \left\{ \|k - k_0\|_\infty, \|\partial_1(k - k_0)\|_\infty \right\}, \quad (\text{F.52})$$

Notice first that $k - k_0 = f + \bar{\sigma}h\bar{\sigma}$ with $h \in \mathcal{S}(\varepsilon)$ implies that $k - k_0 \in \mathbb{W}^{4,p}(\square)$ for all $p > 2$, thus k also as k_0 is regular.

Let us now prove that $\|k - k_0\|_{C^1}$ vanishes when d, ϕ, ξ vanish. By Sobolev embedding (see Proposition E.2), it is enough to prove the same for $\|k - k_0\|_{\mathbb{W}^{2,p}(\square)}$. Since $k - k_0 = f + \bar{\sigma}h\bar{\sigma}$, it is enough to bound $\|f\|_{\mathbb{W}^{2,p}(\square)}$. Theorem 4.3.2.4 in [Gri11] implies that, for a universal constant $C > 0$:

$$\|f\|_{\mathbb{W}^{2,p}} \leq C \left(\|\Delta f\|_{\mathbb{W}^{0,p}} + \|f\|_{\mathbb{W}^{1,p}} \right). \quad (\text{F.53})$$

Recalling the expression (F.39) of the mapping S , it holds that:

$$\begin{aligned} \|\Delta f\|_{\mathbb{W}^{0,p}(\square)} &\leq (\|\phi\|_p + 2\|d\|_\infty \|f\|_p + \|\xi\|_\infty) \|f\|_{\mathbb{W}^{1,p}(\square)} \\ &\leq C' (\|\phi\|_p + 2\|d\|_\infty \|f\|_p + \|\xi\|_\infty) \|f\|_{\mathbb{H}^2(\square)}, \end{aligned} \quad (\text{F.54})$$

where the second inequality is again a Sobolev embedding. The fact that $\|f\|_{\mathbb{H}^2(\square)}$ vanishes with d, ϕ, ξ now follows from Proposition F.5.

We now prove that $k - k_0$ satisfies the boundary conditions of elements of $\mathcal{S}(\infty)$. Since f, k_0, h vanish on $\partial\square$, so does k . Moreover, $\partial_\triangleright h = 0$ at the extremities S_2, S_3 of the diagonal D and $\partial_\triangleright f = 0$ on D means $\partial_\triangleright(k - k_0)(S_2) = 0 = \partial_\triangleright(k - k_0)(S_3)$. We have thus shown that if f has the alleged regularity, then $k - k_0 \in \mathcal{S}(\varepsilon')$ for some $\varepsilon' > 0$ depending on d, ϕ, ξ , i.e. on ρ_\pm, ε recalling (F.36)–(F.35) and $h \in \mathcal{S}(\varepsilon)$.

Let us now prove the regularity of f . By definition, $\Delta f = \Delta S f$, with S defined in (F.39). Let $p > 2$. The idea is classical: if $f \in \mathbb{W}^{2+n,p}(\square)$, $n \in \mathbb{N}$, we want to prove that $\Delta S f \in \mathbb{W}^{1+n,p}(\square)$, from which $f \in \mathbb{W}^{3+n,p}(\square)$ by Proposition F.6 provided $\Delta S f$ satisfies suitable boundary conditions. To implement this recursion scheme, we first prove that $\Delta S f \in \mathbb{L}^p(\square)$. By assumption on ϕ, ψ in case of item (i) (using the Sobolev embedding $C^2(\bar{\triangleright}) \subset \mathbb{W}^{2,s}(\triangleright)$ for any $s > 2$), and from (F.35)–(F.36) for item (ii), we see that it is the regularity of f only that limits the regularity of $\Delta S f$. The fact that $\Delta S f \in \mathbb{L}^p(\square)$ then follows from the Sobolev embedding $\mathbb{H}^1(\square) \subset \mathbb{L}^s(\square)$, valid for any $s > 2$, see Proposition E.2. It follows that $f \in \mathbb{W}^{2,p}(\square)$ by Proposition F.6.

To obtain further regularity on f , let us check that $\Delta S f$ is in $\mathbb{W}^{b-1,p}(\square)$ whenever $f \in \mathbb{W}^{b,p}(\square)$ for $b \in \mathbb{N}^*$, and also that $\Delta S f$ satisfies the condition (F.49) in Proposition F.6. The regularity of $\Delta S f$ boils down to proving that $\mathcal{W}(u, v)$, defined in (F.31), is in $\mathbb{W}^{b,p}(\triangleright)$ ($b \in \mathbb{N}^*$) whenever $u, v \in \mathbb{W}^{b,p}(\triangleright)$. This is the claim of the following lemma.

Lemma F.9. *Let $p \geq 2$ and $1 \leq b \leq 4$. Let $u, v \in \mathbb{W}^{b,p}(\triangleright)$, and recall from (F.31) the definition of \mathcal{M} . Then $\mathcal{M}(u, v) \in \mathbb{W}^{b,p}(\triangleright)$.*

Lemma F.9 is easily proven by approximating u, v in $\mathbb{W}^{b,p}(\triangleright)$ by sequences in $C^\infty(\bar{\triangleright})$, and integrating by parts.

It remains to prove that $\Delta S f$ satisfies the condition of Proposition F.6, i.e. that $\Delta S f(S_j) = 0$ for $j \in \{1, 2, 3\}$. By assumption in the case of item (i), and from the expression (F.36) and the definition (2.44) of $\mathcal{S}(\varepsilon)$ for item (ii), we know that $\phi|_{\partial\square} = 0$. The fact that $f|_{\partial\square} = 0$ gives $\mathcal{M}(f, \xi)(S_j) = 0 = \mathcal{M}(\xi, f)(S_j)$ for each j . Integrating by parts and since $\psi \in \mathcal{S}(\infty)$, one has also $\mathcal{M}(\partial_1 f, \partial_1 \psi)(S_j) = 0 = \mathcal{M}(\partial_1 \psi, \partial_1 f)(S_j)$. It follows that $\Delta S f(S_j) = 0$ for $j \in \{1, 2, 3\}$, thus $f \in \mathbb{W}^{4,p}(\square)$. Since $p > 2$ was arbitrary, this concludes the proof. \square

F.5 Bounds on the solution and definition of ε_B

Let $h \in \mathcal{S}(\infty)$ and let $g_h = \bar{\sigma}^{-1} - (\bar{\sigma} + k_h)^{-1}$ be the solution of the main equation (2.61) obtained from the solution k_h of the Euler-Lagrange equation (F.5). In this section, we define the ε_B arising in Theorems 2.3–2.6, and show that the C^1 norm of $g_h - g_0$ can be controlled by $\|h\|_{C^1}$ and $\bar{\rho}'$ as claimed in Theorem 2.6. Writing $C_{k_0} = \bar{\sigma} + k_0 = (\bar{\sigma}^{-1} - g_0)^{-1}$, notice first that, as soon as $\|k_h - k_0\|_2$ is sufficiently small:

$$g_h - g_0 = C_{k_0}^{-1} - (\bar{\sigma} + k_h)^{-1} = C_{k_0}^{-1} - (C_{k_0} + (k_h - k_0))^{-1} = C_{k_0}^{-1} \sum_{n \geq 1} (-1)^n ((k_h - k_0) C_{k_0}^{-1})^{on}. \quad (\text{F.55})$$

For $k_h - k_0 \in \mathcal{S}(\varepsilon')$ for some suitably small $\varepsilon' > 0$, one can take derivatives inside the sum, which directly yields:

$$\|g_h - g_0\|_{C^1} \leq \tilde{\delta}(\rho_{\pm}, \|k_h - k_0\|_{C^1}), \quad \lim_{(x,y) \rightarrow 0} \tilde{\delta}(x, y) = 0. \quad (\text{F.56})$$

Since $f = k_h - k_0 - \bar{\sigma}h\bar{\sigma}$ is the solution of (P), $\|k_h - k\|_{C^1}$ in turn only depends on $\rho_{\pm}, \|h\|_{C^1}$, thus $\|g_h - g_0\|_{C^1}$ only depends on $\bar{\rho}', \|h\|_{C^1}$, and vanishes when they both vanish as claimed in Theorem 2.6.

We now define ε_B .

Definition F.10 (Definition of ε_B). *Let $h \in \mathcal{S}(\infty)$. Let $\rho_- \in (0, 1)$, and choose $\varepsilon_B = \varepsilon_B(\rho_-) > 0$ and $\rho_+ \in [\rho_-, 1)$ such that, if $\bar{\rho}' := \frac{\rho_+ - \rho_-}{2} \leq \varepsilon_B$ and $h \in \mathcal{S}(\varepsilon_B)$, then:*

- *The contraction S has a fixed point f .*
- *g_h is a negative kernel (as defined in (A.2)).*
- *To ensure good concentration properties in Section 3, $\|g_h - h\|_{C^1} \|g_h\|_{C^1} \leq (2^{10} C_{LS})^{-1}$, $\bar{\rho}' \|g_h - h\|_{C^1} \leq (2^{10} C_{LS})^{-1}$ and $\bar{\rho}' \varepsilon_B \leq (2^{10} C_{LS})^{-1}$, where C_{LS} is the log-Sobolev constant appearing in Lemma 3.4.*

References

- [AF03] R. A. Adams and J. J. Fournier. *Sobolev Spaces*. Elsevier, 2003 (cit. on p. 87).
- [Ass07] S. Assing. “A Limit Theorem for Quadratic Fluctuations in Symmetric Simple Exclusion”. In: *Stochastic Processes and their Applications* 117.6 (6 2007), pp. 766–790. ISSN: 0304-4149. DOI: 10.1016/j.spa.2006.10.005 (cit. on p. 4).
- [ACR21] M. Ayala, G. Carinci, and F. Redig. “Higher Order Fluctuation Fields and Orthogonal Duality Polynomials”. In: *Electronic Journal of Probability* 26.none (none 2021), pp. 1–35. ISSN: 1083-6489, 1083-6489. DOI: 10.1214/21-EJP586 (cit. on p. 4).
- [BB19] R. Bauerschmidt and T. Bodineau. “A Very Simple Proof of the LSI for High Temperature Spin Systems”. In: *Journal of Functional Analysis* 276.8 (8 2019), pp. 2582–2588. ISSN: 0022-1236. DOI: 10.1016/j.jfa.2019.01.007 (cit. on p. 72).
- [BGL22a] L. Bertini, D. Gabrielli, and C. Landim. *Metastable $\mathbb{S} \setminus \Gamma$ -Expansion of Finite State Markov Chains Level Two Large Deviations Rate Functions*. 2022. DOI: 10.48550/arXiv.2207.02588. arXiv: 2207.02588 [cond-mat] (cit. on p. 4).

- [Ber+03] L. Bertini, A. De Sole, D. Gabrielli, G. Jona-Lasinio, and C. Landim. “Large Deviations for the Boundary Driven Symmetric Simple Exclusion Process”. In: *Mathematical Physics, Analysis and Geometry* 6.3 (3 2003), pp. 231–267 (cit. on pp. 8–9, 62).
- [Ber+15] L. Bertini, A. De Sole, D. Gabrielli, G. Jona-Lasinio, and C. Landim. “Macroscopic Fluctuation Theory”. In: *Reviews of Modern Physics* 87.2 (2 2015), pp. 593–636. DOI: 10.1103/RevModPhys.87.593 (cit. on pp. 1–2, 5).
- [BGL21] L. Bertini, D. Gabrielli, and C. Landim. *Concurrent Donsker-Varadhan and Hydrodynamical Large Deviations*. 2021. DOI: 10.48550/arXiv.2111.05892. arXiv: 2111.05892 [cond-mat] (cit. on pp. 4, 9).
- [BGL22b] L. Bertini, D. Gabrielli, and C. Landim. *Large Deviations for Diffusions: Donsker-Varadhan Meet Freidlin-Wentzell*. 2022. DOI: 10.48550/arXiv.2211.02593. arXiv: 2211.02593 [cond-mat] (cit. on p. 4).
- [Bod+08] T. Bodineau, B. Derrida, V. Lecomte, and F. van Wijland. “Long Range Correlations and Phase Transitions in Non-equilibrium Diffusive Systems”. In: *Journal of Statistical Physics* 133.6 (6 2008), pp. 1013–1031. ISSN: 1572-9613. DOI: 10.1007/s10955-008-9647-3 (cit. on p. 18).
- [BKM96] A. V. Bukhvalov, V. B. Korotkov, and B. M. Makarovt. “Integral Operators”. In: *Vector Lattices and Integral Operators*. Ed. by S. S. Kutateladze. Mathematics and Its Applications. Dordrecht: Springer Netherlands, 1996, pp. 279–346. ISBN: 978-94-009-0195-7. DOI: 10.1007/978-94-009-0195-7_4 (cit. on p. 16).
- [CS21] J. P. Chen and F. Sau. “Higher-Order Hydrodynamics and Equilibrium Fluctuations of Interacting Particle Systems”. In: *Markov Processes and Related Fields* 27.3 (2021), pp. 339–380. ISSN: 1024-2953 (cit. on p. 4).
- [Dag22] B. Dagallier. “Fluctuations and Correlations in the SSEP on a Ring Subject to a Macroscopic Current”. In: *ongoing work* (2022+) (cit. on pp. 9, 16, 18).
- [DLS02] B. Derrida, J. L. Lebowitz, and E. R. Speer. “Large Deviation of the Density Profile in the Steady State of the Open Symmetric Simple Exclusion Process”. In: *Journal of statistical physics* 107.3-4 (3-4 2002), pp. 599–634 (cit. on pp. 6, 15).
- [Der+05] B. Derrida, C. Enaud, C. Landim, and S. Olla. “Fluctuations in the Weakly Asymmetric Exclusion Process with Open Boundary Conditions”. In: *Journal of Statistical Physics* 118.5 (5 2005), pp. 795–811. ISSN: 1572-9613. DOI: 10.1007/s10955-004-1989-x (cit. on pp. 2–3, 6).
- [DLS07] B. Derrida, J. L. Lebowitz, and E. R. Speer. “Entropy of Open Lattice Systems”. In: *Journal of Statistical Physics* 126.4-5 (4-5 2007), pp. 1083–1108 (cit. on p. 2).
- [DR13] B. Derrida and M. Retaux. “Finite Size Corrections to the Large Deviation Function of the Density in the One Dimensional Symmetric Simple Exclusion Process”. In: *Journal of Statistical Physics* 152.5 (5 2013), pp. 824–845 (cit. on p. 2).
- [Der07] B. Derrida. “Non-Equilibrium Steady States: Fluctuations and Large Deviations of the Density and of the Current”. In: *Journal of Statistical Mechanics: Theory and Experiment* 2007.07 (07 2007), P07023 (cit. on pp. 2, 5).

- [DS19] B. Derrida and T. Sadhu. “Large Deviations Conditioned on Large Deviations II: Fluctuating Hydrodynamics”. In: *Journal of Statistical Physics* 177.1 (2019), pp. 151–182. ISSN: 1572-9613. DOI: 10.1007/s10955-019-02363-8 (cit. on p. 18).
- [DS96] P. Diaconis and L. Saloff-Coste. “Logarithmic Sobolev Inequalities for Finite Markov Chains”. In: *The Annals of Applied Probability* 6.3 (3 1996), pp. 695–750. ISSN: 1050-5164, 2168-8737. DOI: 10.1214/aoap/1034968224 (cit. on p. 68).
- [DV75] M. D. Donsker and S. R. S. Varadhan. “Asymptotic Evaluation of Certain Markov Process Expectations for Large Time, I”. In: *Communications on Pure and Applied Mathematics* 28.1 (1 1975), pp. 1–47. ISSN: 1097-0312. DOI: 10.1002/cpa.3160280102 (cit. on pp. 2, 7).
- [ELS90] G. Eyink, J. L. Lebowitz, and H. Spohn. “Hydrodynamics of Stationary Non-Equilibrium States for Some Stochastic Lattice Gas Models”. In: *Communications in mathematical physics* 132.1 (1 1990), pp. 253–283 (cit. on p. 2).
- [FLM11] J. Farfan, C. Landim, and M. Mourragui. “Hydrostatics and Dynamical Large Deviations of Boundary Driven Gradient Symmetric Exclusion Processes”. In: *Stochastic Processes and their Applications* 121.4 (4 2011), pp. 725–758. ISSN: 0304-4149. DOI: 10.1016/j.spa.2010.11.014 (cit. on p. 2).
- [Gar+90] P. L. Garrido, J. L. Lebowitz, C. Maes, and H. Spohn. “Long-Range Correlations for Conservative Dynamics”. In: *Physical Review A* 42.4 (4 1990), pp. 1954–1968. DOI: 10.1103/PhysRevA.42.1954 (cit. on p. 1).
- [Gon+22] P. Gonçalves, M. Jara, R. Marinho, and O. Menezes. *CLT for NESS of a Reaction-Diffusion Model*. 2022. DOI: 10.48550/arXiv.2211.04835. arXiv: 2211.04835 [math-ph] (cit. on p. 3).
- [Gon+20] P. Gonçalves, M. Jara, O. Menezes, and A. Neumann. “Non-Equilibrium and Stationary Fluctuations for the SSEP with Slow Boundary”. In: *Stochastic Processes and their Applications* 130.7 (7 2020), pp. 4326–4357. ISSN: 0304-4149. DOI: 10.1016/j.spa.2019.12.006 (cit. on pp. 3, 15).
- [GJ19] P. Gonçalves and M. Jara. “Quadratic Fluctuations of the Symmetric Simple Exclusion”. In: *ALEA* 16 (2019), pp. 605–632 (cit. on pp. 4, 6, 15).
- [Gon+21] P. Gonçalves, M. Jara, R. Marinho, and O. Menezes. *Sharp Convergence to Equilibrium for the SSEP with Reservoirs*. 2021. DOI: 10.48550/arXiv.2110.06353. arXiv: 2110.06353 [cond-mat, physics:math-ph] (cit. on p. 72).
- [GSS+19] F. Götze, H. Sambale, A. Simulis, et al. “Higher Order Concentration for Functions of Weakly Dependent Random Variables”. In: *Electronic Journal of Probability* 24 (2019) (cit. on pp. 67–68).
- [Gri11] P. Grisvard. *Elliptic Problems in Nonsmooth Domains*. Classics in Applied Mathematics. Society for Industrial and Applied Mathematics, 2011. ISBN: 978-1-61197-202-3. DOI: 10.1137/1.9781611972030 (cit. on pp. 13, 87, 89, 95–98).
- [JM18] M. Jara and O. Menezes. “Non-Equilibrium Fluctuations of Interacting Particle Systems”. 2018. arXiv: 1810.09526 (cit. on pp. 1, 3–4, 6, 15, 18–21, 40, 74–75).

- [JM20] M. Jara and O. Menezes. “Non-Equilibrium Fluctuations for a Reaction-Diffusion Model via Relative Entropy”. In: *Markov Processes and Related Fields* 26.1 (2020), pp. 95–124. ISSN: 1024-2953 (cit. on pp. 3, 18–19).
- [KOV89] C. Kipnis, S. Olla, and S. Varadhan. “Hydrodynamics and Large Deviation for Simple Exclusion Processes”. In: *Communications on Pure and Applied Mathematics* 42.2 (2 1989), pp. 115–137 (cit. on p. 3).
- [KL99] C. Kipnis and C. Landim. *Scaling Limits of Interacting Particle Systems*. Springer Berlin Heidelberg, 1999. DOI: 10.1007/978-3-662-03752-2 (cit. on pp. 3, 6, 12, 14, 42, 44, 47, 49, 56, 90).
- [Lan22] C. Landim. *Metastability from the Large Deviations Point of View: A Γ -Expansion of the Level Two Large Deviations Rate Functional of Non-Reversible Finite-State Markov Chains*. 2022. DOI: 10.48550/arXiv.2209.11532. arXiv: 2209.11532 [math] (cit. on p. 4).
- [LMO06] C. Landim, A. Milanés, and S. Olla. “Stationary and Nonequilibrium Fluctuations in Boundary Driven Exclusion Processes”. In: *Markov processes and related fields* 14.2 (2 2006), pp. 165–184 (cit. on pp. 2–3, 6).
- [SS20] H. Sambale and A. Sinulis. “Logarithmic Sobolev Inequalities for Finite Spin Systems and Applications”. In: *Bernoulli* 26.3 (3 2020), pp. 1863–1890. DOI: 10.3150/19-BEJ1172 (cit. on p. 67).
- [Siu16] B. Siudeja. “On Mixed Dirichlet-Neumann Eigenvalues of Triangles”. In: *Proceedings of the American Mathematical Society* 144.6 (6 2016), pp. 2479–2493. ISSN: 0002-9939, 1088-6826. DOI: 10.1090/proc/12888 (cit. on p. 95).
- [Spo83] H. Spohn. “Long Range Correlations for Stochastic Lattice Gases in a Non-Equilibrium Steady State”. In: *Journal of Physics A: Mathematical and General* 16.18 (18 1983), p. 4275 (cit. on pp. 1–3, 5–6, 18).
- [Yau91] H.-T. Yau. “Relative Entropy and Hydrodynamics of Ginzburg-Landau Models”. In: *Letters in Mathematical Physics* 22.1 (1 1991), pp. 63–80. ISSN: 1573-0530. DOI: 10.1007/BF00400379 (cit. on pp. 3, 15).