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AN ODE CHARACTERISATION OF MULTI-MARGINAL OPTIMAL TRANSPORT WITH PAIRWISE COST FUNCTIONS

LUCA NENNA AND BRENDAN PASS

ABSTRACT. The purpose of this paper is to introduce a new numerical method to solve multi-marginal optimal transport problems with pairwise interaction costs. The complexity of multi-marginal optimal transport generally scales exponentially in the number of marginals m . We introduce a one parameter family of cost functions that interpolates between the original and a special cost function for which the problem's complexity scales linearly in m . We then show that the solution to the original problem can be recovered by solving an ordinary differential equation in the parameter ε , whose initial condition corresponds to the solution for the special cost function mentioned above; we then present some simulations, using both explicit Euler and explicit higher order Runge-Kutta schemes to compute solutions to the ODE, and, as a result, the multi-marginal optimal transport problem.

1. INTRODUCTION

The theory of optimal transport plays an important role in many applications (see [52, 54, 56, 57]). Its generalization to the multi-marginal case consists in minimizing the functional

$$\gamma \mapsto \int_{X^1 \times \dots \times X^m} c(x^1, \dots, x^m) d\gamma$$

among all probability measures $\gamma \in \mathcal{P}(X^1 \times \dots \times X^m)$ having fixed $\mu^i \in \mathcal{P}(X^i)$ with $i = 1, \dots, m$ as marginals, for a given cost function $c(x^1, \dots, x^m)$. This problem has been at the center of growing interest in recent years since it arises naturally in many different areas of applications, including Economics [22], Financial Mathematics [7, 29–31], Statistics [11, 21], Image Processing [53], Tomography [1], Machine Learning [36, 55], Fluid Dynamics [13] and Quantum Physics and Chemistry, in the framework of Density Functional Theory [17, 26].

The structure of solutions to the multi-marginal optimal transport problem is a notoriously delicate issue, and is still not well understood, despite substantial efforts on the part of many researchers [18, 23–25, 35, 37, 40, 41, 43, 46, 47, 49–51]; see also the surveys [48] and [28]. In many of the aforementioned applications, it is therefore pertinent to develop efficient numerical algorithms to compute solutions. This, however, represents a significant challenge, since the problem amounts to a linear

(or convex, in a popular regularized variant discussed below), yet high dimensional optimization problem: the complexity scales exponentially in the number m of marginals. For instance a crude discretization of each of 5 marginals (notice that in many applications the number of marginals could be dramatically large, e.g. in quantum mechanics where m is the number of electrons in a molecule) using 100 Dirac masses would mean that the coupling γ between the 5 marginals is supported over $100^5 = 10^{10}$ Dirac masses, rendering the problem practically intractable. There have been recently some attempts to tackle this problem by using different approaches: entropic regularization [8, 10, 44], relaxation via moment constraints approximation [3, 4], genetic column generation algorithm exploiting the existence of a sparse solution in the discrete case [32, 33], Wasserstein penalisation of the marginal constraints [42] and semidefinite relaxation [38, 39].

In many cases of interest, the cost function $c(x^1, \dots, x^m) = \sum_{i,j=1, i \neq j}^m w(x^i, x^j)$ is given by a sum of two marginal cost functions; when $w(x^i, x^j) = |x^i - x^j|^2$, for instance the multi-marginal problem is equivalent to the well known Wasserstein barycenter problem (see Proposition 4.2 in [2]), while the Coulomb cost $w(x^i, x^j) = \frac{1}{|x^i - x^j|}$ plays a central role in the quantum chemistry applications pioneered in [26] and [17]. Here, for such pairwise interaction costs, our aim is to develop a continuation method which, by introducing a suitable one parameter family of cost functions, establishes a link between the original multi-marginal problem and a simpler one whose complexity scales linearly in the number of marginals. For discrete marginals, we show that, after the addition of an entropic regularization term, the solution of the original multi-marginal problem can be recovered by solving an ordinary differential equation (ODE) whose initial condition is the solution to the simpler problem. This method is actually inspired by the one introduced in [19] to compute the Monge solution of the two marginal problem, starting from the Knothe-Rosenblatt rearrangement; note, however, note that since we apply this strategy to a regularized problem, our resulting ODE enjoys better regularity than the one in [19], which, in turn, makes it amenable to higher order numerical schemes (see the description of numerical results in Section 4 below). The above mentioned differential equation will be derived by differentiating the optimality conditions of the dual problem; in particular by penalizing the constraints with the soft-max function we will obtain a well defined ODE for which existence and uniqueness of a solution can be established.

When developing the ODE approach in Section 3 below, we restrict our attention to the case when the marginals μ^i are all identical. This has the significant advantage of reducing the Kantorovich dual problem to a maximization over a single potential function, while also capturing important applications arising in density functional theory. Though we do not pursue this direction here, our approach, with minor modifications,

will also work with distinct marginals. In this case, if each measure is discretized using N points, one would need to solve $(m - 1)N$ coupled, real-valued ODEs (rather than the N coupled ODEs dealt with here in Section 3), reflecting the $m - 1$ independent Kantorovich potentials needed to fully capture the solution. In this setting, again with minor modifications, the pairwise interactions may in fact differ as well; that is, we may consider $c(x^1, \dots, x^m) = \sum_{i,j=1, i \neq j}^m w_{ij}(x^i, x^j)$ where the w_{ij} differ for different pairs i, j . Although we do not present the detailed calculations for these non symmetric problems (since they are mathematically very close to the symmetric case, though more notationally cumbersome), we do present some numerical simulations for one such problem (relating to Brenier's relaxation of the variational formulation of the incompressible Euler equation) to illustrate how the numerics can be extended. In addition, a separate manuscript on a related but somewhat different approach, which deals with completely general (not necessarily pairwise) cost functions is currently under preparation.

The remainder of this manuscript is organized as follows. In Section 2, we recall some basic facts about multi-marginal optimal transport, as well as its entropic regularization and the duals of both problems, and prove that for a particular, simple cost function, the solutions to the regularized problem and its dual can be computed by solving $m - 1$ two marginal problems. This solution will serve as the initial condition for an ODE, which is introduced, and proven to be well-posed, in Section 3. In Section 4, algorithms, based on this ODE, to compute the solution to the multi-marginal optimal transport problem are described and some resulting numerical simulations are presented.

2. MULTI-MARGINAL OPTIMAL TRANSPORT AND ENTROPIC REGULARIZATION

Given m probability measures μ^i on bounded domains $X^i \subseteq \mathbb{R}^n$ for $i = 1, 2, \dots, m$ and a lower semi-continuous cost function $c : X^1 \times X^2 \times \dots \times X^m \rightarrow \mathbb{R} \cup \{+\infty\}$, the multi-marginal optimal transport problem consists in solving the following optimization problem

$$\inf_{\gamma \in \Pi(\mu^1, \dots, \mu^m)} \int_{X^1 \times X^2 \times \dots \times X^m} c(x^1, \dots, x^m) d\gamma \quad (1)$$

where $\Pi(\mu^1, \dots, \mu^m)$ denotes the set of probability measures on $X^1 \times X^2 \times \dots \times X^m$ whose marginals are the μ^i . One can easily show by means of the direct method of calculus of variations that this problem admits at least a solution, which will be referred as the *optimal transport plan*. It is well known that under some mild assumptions the above problem is dual to the following

$$\sup \left\{ \sum_{i=1}^m \int_{X^i} \varphi^i(x^i) d\mu^i \mid \varphi^i \in L^1(\mu^i), \sum_{i=1}^m \varphi^i(x^i) \leq c(x^1, \dots, x^m) \right\}. \quad (2)$$

We will also consider a common variant of (1), known as *entropic optimal transport* which consists in adding an entropy regularization term. For a parameter $\eta > 0$, this is to minimize

$$\inf_{\gamma \in \Pi(\mu^1, \dots, \mu^m)} \left\{ \int_{X^1 \times X^2 \dots \times X^m} c(x^1, \dots, x^m) d\gamma + \eta H_{\otimes_{i=1}^m \mu^i}(\gamma) \right\} \quad (3)$$

where the relative entropy $H_{\otimes_{i=1}^m \mu^i}(\gamma)$ with respect to product measure $\otimes_{i=1}^m \mu^i$ is defined by

$$H_{\otimes_{i=1}^m \mu^i}(\gamma) = \int_{X^1 \times \dots \times X^m} \frac{d\gamma}{d(\otimes_{i=1}^m \mu^i)} \log\left(\frac{d\gamma}{d(\otimes_{i=1}^m \mu^i)}\right) d(\otimes_{i=1}^m \mu^i),$$

if γ is absolutely continuous with respect to the product measure $\otimes_{i=1}^m \mu^i$ and $+\infty$ otherwise. The regularized transport is, in turn, dual to the following unconstrained optimization problem

$$\begin{aligned} & \sup \sum_{i=1}^m \int_{X^i} \varphi^i(x^i) d\mu^i \\ & - \eta \int_{X^1 \times \dots \times X^m} e^{\frac{\sum_{i=1}^m \varphi^i(x^i) - c(x^1, \dots, x^m)}{\eta}} d(\otimes_{i=1}^m \mu^i)(x^1, \dots, x^m). \end{aligned} \quad (4)$$

The regularized problem (3) and its dual (4) arise frequently in computational work. We note in particular that (3) is the minimization of a strictly convex functional and therefore admits a unique solution. It is well known that as $\eta \rightarrow 0$, solutions of (3) and (4) converge to solutions of (1) and (2), respectively. When each X^i is a finite set (a case of particular interest in this paper), we obtain discrete versions of (1) and its dual, which amount to linear programs, taking the forms

$$\inf \left\{ \sum_{\bar{x} \in \times_{i=1}^m X^i} c(\bar{x}) \gamma_{\bar{x}} \mid \gamma \in \Pi(\mu^1, \dots, \mu^m) \right\}, \quad (5)$$

where $\bar{x} = (x^1, \dots, x^m) \in X^1 \times \dots \times X^m$ and

$$\sup \left\{ \sum_{i=1}^m \sum_{x \in X^i} \varphi_x^i \mu_x^i \mid (\varphi^1, \dots, \varphi^m) \in \mathcal{T} \right\} \quad (6)$$

where, if we identify functions $\varphi^i : X^i \rightarrow \mathbb{R}$ with points in $\mathbb{R}^{|X^i|}$,

$$\mathcal{T} := \left\{ \varphi^i \in \mathbb{R}^{|X^i|} \mid \forall i = 1, \dots, m, \sum_{i=1}^m \varphi_{x^i}^i \leq c(x^1, \dots, x^m), \forall (x^1, \dots, x^m) \in \times_{i=1}^m X^i \right\}.$$

Notice that, if each $|X^i| = N$, in the case of the primal problem we have to deal with N^m unknowns and mN constraints, whereas in the dual problem there are mN unknowns and N^m constraints. In both cases we have to deal with the so called “curse of dimensionality,” namely the complexity of the problem increases exponentially with the number of marginals.

In this discrete setting, the entropy regularized problem (3) and its dual (4) then become finite dimensional convex optimization problems:

$$\inf \left\{ \sum_{\bar{x} \in \times_{i=1}^m X^i} c(\bar{x}) \gamma_{\bar{x}} + \eta [H(\gamma) - H(\otimes^m \mu^i)] \mid \gamma \in \Pi(\mu^1, \dots, \mu^m), \right\} \quad (7)$$

where $\eta > 0$ and H is the entropy with respect to uniform measure on the finite set X (we suppress the subscript on H indicating the reference measure in the finite case, as we will only deal with entropy relative to the uniform measure), $H(\gamma) = \sum_{\bar{x} \in \times_{i=1}^m X^i} h(\gamma_{\bar{x}})$, with

$$h(t) = \begin{cases} t(\log(t) - 1), & t > 0 \\ 0, & t = 0 \\ +\infty, & t < 0, \end{cases}$$

and

$$\sup \left\{ \sum_{i=1}^m \sum_{x \in X^i} \varphi_x^i \mu_x^i - \eta \sum_{\bar{x} \in \times_{i=1}^m X^i} \exp \left(\frac{\sum_i \varphi_{x^i}^i - c(\bar{x})}{\eta} \right) (\otimes^m \mu^i)_{\bar{x}} \right\}. \quad (8)$$

We note in particular that (8) is an *unconstrained* finite dimensional concave maximization problem. Solutions may be computed using a multi-marginal version of the Sinkhorn algorithm [9, 27, 34, 52], and one can then recover the optimal γ in (7) from the solutions $\varphi^1, \dots, \varphi^m$ to (8) via the well known formula:

$$\gamma_{\bar{x}} = \exp \left(\frac{\sum_{i=1}^m \varphi_{x^i}^i - c(\bar{x})}{\eta} \right) \mu_{x^1}^1 \mu_{x^2}^2 \dots \mu_{x^m}^m$$

where $\bar{x} = (x^1, \dots, x^m)$.

2.1. Pairwise costs. We are especially interested in this paper in cost functions $c(x^1, \dots, x^m)$ involving pair-wise interactions, that is

$$c(x^1, \dots, x^m) = \sum_{i < j}^m w(x^i, x^j).$$

Such costs are ubiquitous in applications: for example, for systems of interacting classical particles in [17, 26], c is a pair-wise cost, with $w(x - y) = \frac{1}{|x - y|}$, known as the Coulomb cost. The case where $w(x, y) = |x - y|^2$ is the quadratic distance is well known to be equivalent to the Wasserstein barycenter problem (see Proposition 4.2 in [2]), which has a wide variety of applications in statistics, machine learning and image processing, among other areas.¹

¹Pairwise costs with pair-dependent interactions, that is, costs of the form $c(x^1, \dots, x^m) = \sum_{i < j}^m w_{ij}(x^i, x^j)$ where the w_{ij} are not necessarily the same for each choice of i and j , also arise in the time discretization of Arnold's variational formulation of incompressible Euler equation [6] and Brenier's relaxation of it [13], and in

Let us now consider costs c_ε of the form

$$c_\varepsilon(x^1, \dots, x^m) := \varepsilon \sum_{i=2}^m \sum_{j=i+1}^m w(x^i, x^j) + \sum_{i=2}^m w(x^1, x^i). \quad (9)$$

It is clear that when $\varepsilon = 1$ we retrieve a pair-wise cost as defined above whereas in the limit $\varepsilon \rightarrow 0$ we obtain a cost involving only the interactions between x^1 and the other x^i individually. Later on, we will develop an ordinary differential equation that governs the evolution with ε of the solutions to the regularized dual problem (8); the results below assert that the initial condition for that equation (that is, the solutions when $\varepsilon = 0$) can be recovered by solving each of the individual two marginal problems between μ^1 and μ^i .

In what follows, we will assume that each marginal μ^i is absolutely continuous with respect to a fixed based measure ν^i with density given by $\frac{d\mu^i}{d\nu^i}$.

Proposition 1. *Assume that each marginal μ^i is absolutely continuous with respect to a fixed based measure ν^i with density given by $\frac{d\mu^i}{d\nu^i}(x^i)$: $d\mu^i(x^i) = \frac{d\mu^i}{d\nu^i}(x^i)d\nu^i(x^i)$. Consider the regularized problem (3) with limiting pairwise cost; that is, set $\varepsilon = 0$ in (9) to obtain:*

$$\min_{\gamma \in \Pi(\mu^1, \mu^2, \dots, \mu^m)} \int \sum_{i=2}^m w(x^1, x^i) d\gamma + \eta H_{\otimes_{i=1}^m \mu^i}(\gamma). \quad (10)$$

Let $\frac{d\bar{\pi}^i}{d(\nu^1 \otimes \nu^i)}$ be the density with respect to product measure $\nu^1(x^1) \otimes \nu^i(x^i)$ of the minimizer $\bar{\pi}^i = \frac{d\bar{\pi}^i}{d(\nu^1 \otimes \nu^i)} \nu^1 \otimes \nu^i$ in the regularized two marginal problem:

$$\min_{\pi^i \in \Pi(\mu^1, \mu^i)} \int w(x^1, x^i) d\pi^i(x^1, x^i) + \eta H_{\mu^1 \otimes \mu^i}(\pi^i). \quad (11)$$

Then the density $\frac{d\bar{\gamma}}{d(\otimes_{i=1}^m \nu^i)}$ of the optimal $\bar{\gamma} = \frac{d\bar{\gamma}}{d(\otimes_{i=1}^m \nu^i)}(\otimes_{i=1}^m \nu^i)$ in (10) is given by

$$\frac{d\bar{\gamma}}{d(\otimes_{i=1}^m \nu^i)}(x^1, \dots, x^m) = \frac{\frac{d\bar{\pi}^2}{d(\nu^1 \otimes \nu^2)}(x^1, x^2)}{\frac{d\mu^1}{d\nu^1}(x^1)} \frac{\frac{d\bar{\pi}^3}{d(\nu^1 \otimes \nu^3)}(x^1, x^3)}{\frac{d\mu^1}{d\nu^1}(x^1)} \dots \frac{\frac{d\bar{\pi}^m}{d(\nu^1 \otimes \nu^m)}(x^1, x^m)}{\frac{d\mu^1}{d\nu^1}(x^1)} \frac{d\mu^1}{d\nu^1}(x^1).$$

Proof. Choose any $\gamma = \frac{d\gamma}{d(\otimes_{i=1}^m \nu^i)}(\otimes_{i=1}^m \nu^i) \in \Pi(\mu^1, \mu^2, \dots, \mu^m)$ which is absolutely continuous with respect to $\otimes_{i=1}^m \nu^i$ and let $\pi^i(x^1, x^i) = \left((x^1, \dots, x^m) \mapsto \right.$

inference problems for probabilistic graphical models [36]. The results in this section adapt immediately to such costs.

$(x^1, x^i)_{\#} \gamma \in \Pi(\mu^1, \mu^i)$ be its twofold marginals. Then

$$\begin{aligned} H_{\otimes_{i=1}^m \mu^i}(\gamma) &= \int_{X^1 \times \dots \times X^m} \left[\log\left(\frac{d\gamma}{d(\otimes_{i=1}^m \nu^i)}(x^1, \dots, x^m)\right) - \sum_{i=1}^m \log\left(\frac{d\mu^i}{d\nu^i}(x^i)\right) \right] d\gamma(x^1, \dots, x^m) \\ &= \int_{X^1 \times \dots \times X^m} \left[\log\left(\frac{d\gamma}{d(\otimes_{i=1}^m \nu^i)}(x^1, \dots, x^m)\right) \right] d\gamma(x^1, \dots, x^m) - \sum_{i=1}^m H_{\nu^i}(\mu^i), \end{aligned} \quad (12)$$

where each $H_{\nu^i}(\mu^i) = \int_{X^i} \log\left(\frac{d\mu^i}{d\nu^i}(x^i)\right) d\mu^i(x^i)$ is constant throughout $\Pi(\mu^1, \dots, \mu^m)$.

Now disintegrating $\gamma = \gamma_{x^1}(x^2, \dots, x^m) \otimes \mu^1(x^1)$ with respect to its first marginal μ^1 , we note that $\gamma_{x^1}(x^2, \dots, x^m) = \frac{d\gamma}{d(\otimes_{i=1}^m \nu^i)}(x^1, x^2, \dots, x^m) \frac{1}{\frac{d\mu^1}{d\nu^1}(x^1)} (\otimes_{i=2}^m \nu^i)$ and so

$$\begin{aligned} & \int_{X^1 \times \dots \times X^m} \log\left(\frac{d\gamma}{d(\otimes_{i=1}^m \nu^i)}\right) d\gamma \\ &= \int_{X^1 \times \dots \times X^m} \log\left(\frac{d\gamma_{x^1}}{d(\otimes_{i=2}^m \nu^i)}\right) d\gamma + H_{\nu^1}(\mu^1) \\ &= \int_{X^1} \int_{X^2 \times \dots \times X^m} \log\left(\frac{d\gamma_{x^1}}{d(\otimes_{i=2}^m \nu^i)}\right) d\gamma_{x^1} d\mu^1 + H_{\nu^1}(\mu^1) \\ &= \int_{X^1} H_{\otimes_{i=2}^m \nu^i}(\gamma_{x^1}) d\mu^1 + H_{\nu^1}(\mu^1) \end{aligned} \quad (13)$$

where $H_{\otimes_{i=2}^m \nu^i}(\gamma_{x^1})$ is the entropy of γ_{x^1} with respect to $\otimes_{i=2}^m \nu^i$ and $H_{\nu^1}(\mu^1)$ is the entropy of μ^1 with respect to ν^1 . Now note that if we disintegrate each $\pi^i = \pi_{x^1}^i(x^i) \otimes \mu^1(x^1)$ with respect to μ^1 , then for each fixed x^1 , the conditional probability $\pi_{x^1}^i$ is the i th marginal of γ_{x^1} and so

$$\begin{aligned} \int_{X^1} H_{\otimes_{i=2}^m \nu^i}(\gamma_{x^1}) d\mu^1(x^1) &\geq \int_{X^1} \sum_{i=2}^m H_{\nu^i}(\pi_{x^1}^i) d\mu^1(x^1) \\ &= \sum_{i=2}^m \int_{X^1} H_{\nu^i}(\pi_{x^1}^i) d\mu^1(x^1) \\ &= \sum_{i=2}^m [H_{\nu^1 \otimes \nu^i}(\pi^i) - H_{\nu^1}(\mu^1)] \\ &= \sum_{i=2}^m [H_{\mu^1 \otimes \mu^i}(\pi^i) + H_{\nu^i}(\mu^i)], \end{aligned} \quad (14)$$

where the equality $H_{\nu^1 \otimes \nu^i}(\pi^i) = \int_{X^1} H_{\nu^i}(\pi_{x^1}^i) d\mu^1(x^1) + H_{\nu^1}(\mu^1)$ in the second to last line follows very similarly to the derivation of (13) above and the equality $H_{\nu^1 \otimes \nu^i}(\pi^i) - H_{\nu^1}(\mu^1) = H_{\mu^1 \otimes \mu^i}(\pi^i) + H_{\nu^i}(\mu^i)$ in the last line follows very similarly to the derivation of (12). Therefore, combining

(12), (13) and (14), we get

$$\begin{aligned} \int_{X^1 \times \dots \times X^m} \sum_{i=1}^m w(x^1, x^i) d\gamma + \eta H_{\otimes_{i=1}^m \mu_i}(\gamma) &= \int_{X^1 \times X^i} \sum_{i=1}^m w(x^1, x^i) d\pi^i + \eta H_{\otimes_{i=1}^m \mu_i}(\gamma) \\ &\geq \int_{X^1 \times X^i} \sum_{i=2}^m w(x^1, x^i) d\pi^i + \sum_{i=2}^m [H_{\mu^1 \otimes \mu^i}(\pi^i) + H_{\nu^i}(\mu^i)] + H_{\nu^1}(\mu^1) - \sum_{i=1}^m H_{\nu^i}(\mu^i) \\ &\geq \int_{X^1 \times X^i} \sum_{i=2}^m w(x^1, x^i) d\bar{\pi}^i + \sum_{i=2}^m H_{\mu^1 \otimes \mu^i}(\bar{\pi}^i) \end{aligned}$$

by optimality of $\bar{\pi}$. We have equality in the last line if and only if $\pi^i = \bar{\pi}^i$ for each i , and equality in the line above if and only if μ^1 almost every γ_{x^1} couples the $\pi_{x^i}^i$ independently; this yields the desired result. $\square$

Note in particular that this result allows us to recover the solution to problem (7) with cost (9), when $\varepsilon = 0$, by solving $m - 1$ individual regularized two marginal optimal transport problems. In the following section, we will develop a dynamical approach to solve the dual problem (8) to (7) for cost (9) with $\varepsilon > 0$. Our initial condition will be the dual potentials when $\varepsilon = 0$, which we can obtain from the corresponding two marginal dual potentials, as the following corollary confirms.

Corollary 2. *Assume each μ^i is absolutely continuous with respect to a given reference measure ν^i . For each $i = 2, \dots, m$, let $\psi^i(x^1), \varphi^i(x^i)$, solve the regularized two marginal dual problem (8) between marginals μ^1 and μ^i with cost function $w(x^1, x^i)$. Then $\varphi^1(x^1), \varphi^2(x^2), \dots, \varphi^m(x^m)$, with $\varphi^1(x^1) = \sum_{i=2}^m \psi^i(x^1)$ solve the regularized dual (8) with marginals $\mu^1, \mu^2, \dots, \mu^m$ and cost $c(x^1, \dots, x^m) = \sum_{i=2}^m w(x^1, x^i)$.*

Proof. We have that for each i , the optimizer $\bar{\pi}^i$ in the regularized two marginal primal problem satisfies

$$\frac{d\bar{\pi}^i}{d(\nu^1 \otimes \nu^i)}(x^1, x^i) = e^{\frac{\varphi^i(x^i) + \psi^i(x^1) - w(x^1, x^i)}{\eta}} \frac{d\mu^1}{d\nu^1} \frac{d\mu^i}{d\nu^i}.$$

By the preceding proposition, the optimizer in the regularized multi-marginal problem (10) satisfies

$$\begin{aligned} \frac{d\bar{\gamma}}{d(\otimes_{i=1}^m \nu^i)}(x^1, x^2, \dots, x^m) &= \frac{\frac{d\bar{\pi}^2}{d(\nu^1 \otimes \nu^2)}(x^1, x^2)}{\frac{d\mu^1}{d\nu^1}(x^1)} \frac{\frac{d\bar{\pi}^3}{d(\nu^1 \otimes \nu^3)}(x^1, x^3)}{\frac{d\mu^1}{d\nu^1}(x^1)} \dots \frac{\frac{d\bar{\pi}^m}{d(\nu^1 \otimes \nu^m)}(x^1, x^m)}{\frac{d\mu^1}{d\nu^1}(x^1)} \frac{d\mu^1}{d\nu^1}(x^1) \\ &= e^{\frac{\sum_{i=2}^m [\varphi^i(x^i) + \psi^i(x^1) - w(x^1, x^i)]}{\eta}} \frac{d\mu^2}{d\nu^2} \frac{d\mu_3}{d\nu^3} \dots \frac{d\mu^m}{d\nu^m} \frac{d\mu^1}{d\nu^1} \\ &= e^{\frac{\sum_{i=1}^m \varphi^i(x^i) - \sum_{i=2}^m w(x^1, x^i)}{\eta}} \frac{d\mu^1}{d\nu^1} \frac{d\mu^2}{d\nu^2} \frac{d\mu_3}{d\nu^3} \dots \frac{d\mu^m}{d\nu^m} \end{aligned}$$

This is exactly the first order condition identifying the regularized potentials for the multi-marginal regularized problem with cost $\sum_{i=2}^m w(x^1, x^i)$. $\square$

Remark 3. *We note that the cost (9) at $\varepsilon = 0$ falls under the class of tree-structured costs investigated in [36]. These problems are known to correspond to probabilistic graphical models; in our setting, the corresponding graph is in fact a star graph centered at x_1 . From an algorithmic perspective, such costs are desirable since one can solve the*

multi-marginal problem by solving a two marginal problem for each edge of the graph (which is $m-1$ in our case); as was shown in [36], considering the regularized multi-marginal optimal transport problem in fact has some computational advantages over the corresponding series of pairwise regularization. This suggests a connection between that line of research and Proposition 1 and Corollary 2, since these results essentially assert an equivalence between the regularized multi-marginal problem (10) and the m pairwise problems (11); it is not clear to us whether the techniques in [36] can be used to provide an alternative proof of these results.

3. AN ODE CHARACTERISATION OF DISCRETE MULTI-MARGINAL OPTIMAL TRANSPORT

We now turn our attention to developing an ODE for the Kantorovich potentials after discretizing the marginals. Working with the regularized discrete problem (7) and its dual (8) with pairwise cost (9), we make the following, standing assumptions throughout this section:

- (1) (Equal marginals) All the marginals are equal $\mu^i = \rho = \sum_{x \in X} \rho_x \delta_x$, where X is a finite subset of $\mathbb{R}^d$,
- (2) (Symmetric cost) The two body cost w is symmetric $w(x, y) = w(y, x)$.
- (3) (Finite cost) The two body cost function $w : X \times X \rightarrow \mathbb{R}$ is everywhere real-valued.

A motivating example of a pairwise, symmetric two body cost arises in Density Functional Theory where the cost is given by $w(x, y) = \frac{1}{|x - y|}$; in problems with this cost, the marginals are typically also identical. The cost does not satisfy the finiteness hypothesis, but one can consider a truncation $w(x, y) = \min\left(\frac{1}{|x - y|}, C\right)$ cost; it is known that the solution stays away from the diagonal, and for sufficiently large C , the solution with the truncated cost coincides with the solution for the original Coulomb cost (for instance see [16]).

Remark 4. *One could dispose of the equal marginal and symmetric cost assumptions. Analogues of the results proved below would still hold; one could characterize the solution to the regularized dual problem (8) by an ODE, and prove that this ODE is well-posed. In this setting, one could also work with pair-dependent costs of the form $c(x^1, \dots, x^m) = \sum_{i < j}^m w_{ij}(x^i, x^j)$, as discussed briefly in Section 2.1. As we will see below, however, solving the problem numerically becomes more feasible under the hypotheses above, as the solution can be characterized by a single Kantorovich potential, and so the resulting ODE is an equation on $\mathbb{R}^N$, where N is the number of points in the support of the marginal. With unequal marginals and a non-symmetric cost, one would require $m-1$ independent Kantorovich potentials to fully characterize the solution; if each marginal is supported on N points, this would lead to an $(m-1)N$*

dimensional system of ODEs. For the sake of simplicity we do not re-write the proof of the well-posedness since it works exactly as in the symmetric framework; however, we refer the reader to section 4.1 where we present numerical simulations for the (non-symmetric) case of the multi-marginal problem associated to the relaxed formulation of incompressible Euler equation introduced by Brenier.

3.1. Formulation of the ODE problem. Notice now that although the cost (9) at $\varepsilon = 1$ is symmetric in the variables $x^1, x^2, \dots, x^m$, the one at $\varepsilon < 1$ is not. It is, however, symmetric in the variables $x^2, \dots, x^m$; this means that the optimal φ^i in (8) satisfy $\varphi^i = \varphi^j = \varphi$ for $i, j \geq 2$ and so, setting $\varphi^1 = \psi$, we can rewrite (8) as

$$\inf_{\varphi, \psi: X \rightarrow \mathbb{R}} \{ \Phi(\varphi, \psi, \varepsilon) \}, \quad (15)$$

where

$$\Phi(\varphi, \psi, \varepsilon) := -(m-1) \sum_{x \in X} \varphi_x \rho_x - \sum_{x \in X} \psi_x \rho_x + \eta \sum_{\bar{x} \in X^m} e^{\left(\frac{\sum_{i=2}^m \varphi_{x^i} + \psi_{x^1} - c_\varepsilon(\bar{x})}{\eta} \right)} \otimes^m \rho.$$

Remark 5 (Notation). Recall that we use the notation $\bar{x}$ to represent a point in a product space, such as $\bar{x} = (x^1, \dots, x^m) \in X^m$, as above, or, as will often be the case in what follows, $\bar{x} = (x^1, \dots, x^{m-1}) \in X^{m-1}$. We introduce the following notation to represent corresponding products of the densities:

$$\tilde{\rho}_{\bar{x}} = (\otimes_{i=1}^{m-1} \rho)_{\bar{x}} = \otimes_{i=1}^{m-1} \rho_{x^i}$$

Since the functional $\Phi(\varphi, \psi, \varepsilon)$ is convex on the set $\{\varphi, \psi : X \rightarrow \mathbb{R}\} \approx \mathbb{R}^{2|X|}$, as the sum of a linear and an exponential function, optimal solutions (φ^*, ψ^*) can be characterized by the first order optimality conditions $\nabla_\varphi \Phi = \nabla_\psi \Phi = 0$, or (component-wise):

$$\varphi_z^* = -\eta \log \left(\sum_{\bar{x} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^{m-1} \varphi_{x^i}^* + \psi_{x^1}^* - c_\varepsilon(\bar{x}, z)}{\eta} \right) \tilde{\rho}_{\bar{x}} \right)$$

and

$$\psi_z^* = -\eta \log \left(\sum_{\bar{x} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^m \varphi_{x^i}^* - c_\varepsilon(z, \bar{x})}{\eta} \right) \tilde{\rho}_{\bar{x}} \right). \quad (16)$$

In particular, note that (16) allows us to express the optimal ψ^* in (15) in terms of the optimal φ^* , after which (15) reduces to the following optimization problem

$$\inf_{\varphi: X \rightarrow \mathbb{R}} \{ \tilde{\Phi}(\varphi, \varepsilon) \}, \quad (17)$$

where

$$\tilde{\Phi}(\varphi, \varepsilon) := -(m-1) \sum_{x \in X} \varphi_x \rho_x + \eta \sum_z \log \left(\sum_{\bar{x} \in X^{m-1}} e^{\left(\frac{\sum_{i=2}^m \varphi_{x^i} - c_\varepsilon(z, \bar{x})}{\eta} \right)} \tilde{\rho}_{\bar{x}} \right) \rho_z.$$

Remark 6 (LogSumExp and convexity). *The function*

$$\varphi \mapsto \log \left(\sum_{\bar{x} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^m \varphi_{x^i} - c_\varepsilon(z, \bar{x})}{\eta} \right) \tilde{\rho}_{\bar{x}} \right) := LSE_{c_\varepsilon}(\varphi)_z$$

is also known as Log-Sum-Exp function (LSE). By using the Hölder inequality one can easily show that the Log-Sum-Exp is convex.

It is well known that the solution to (8) is unique up to the addition of constants $\varphi^i \mapsto \varphi^i + C^i$ adding to 0, $\sum_{i=1}^m C^i = 0$; thus, solutions to (17) are unique up to the addition of a single constant, $\varphi \mapsto \varphi + C$. We therefore impose the normalization

$$\varphi_{x_0} = 0 \quad (18)$$

for all $\varepsilon \in [0, 1]$ and a fixed $x_0 \in X$.

The problem (17), restricted to φ 's satisfying (18) then has a unique solution; the function $\tilde{\Phi}(\cdot, \varepsilon)$ is strictly convex when restricted to this set, and the solution $\varphi^* = \varphi(\varepsilon)$ can be characterized by the optimality condition $\nabla_\varphi \tilde{\Phi}(\varphi^*, \varepsilon) = 0$, where each component of the gradient is given by

$$\frac{\partial}{\partial \varphi_z} \tilde{\Phi} = -(m-1)\rho_z + (m-1)e^{\varphi_z/\eta} \rho_z \sum_y \sum_{\bar{x} \in X^{m-2}} e^{\left(\frac{\sum_{i=3}^m \varphi_{x^i} - c_\varepsilon(y, z, \bar{x})}{\eta} \right)} (\otimes^{m-2} \rho_{\bar{x}}) \bar{\rho}_y. \quad (19)$$

where

$$\bar{\rho}_y = \frac{\rho_y}{\sum_{\bar{x} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^m \varphi_{x^i} - c_\varepsilon(y, \bar{x})}{\eta} \right) \tilde{\rho}_{\bar{x}}}.$$

Our numerical method consists then in solving an ODE for the evolution of $\varphi(\varepsilon)$ obtained by differentiating

$$\nabla_\varphi \tilde{\Phi}(\varphi(\varepsilon), \varepsilon) = 0 \quad (20)$$

with respect to ε :

$$\frac{\partial}{\partial \varepsilon} [\nabla_\varphi \tilde{\Phi}(\varphi, \varepsilon)]|_{\varphi=\varphi(\varepsilon)} + D_{\varphi, \varphi}^2 \tilde{\Phi}(\varphi(\varepsilon), \varepsilon) \frac{d\varphi}{d\varepsilon}(\varepsilon) = 0. \quad (21)$$

If the pure second derivatives with respect to φ as well as the mixed second derivatives with respect to φ and ε exist and are Lipschitz, and the Hessian with respect to φ is invertible, we will obtain a characterization of φ as the solution to the following well-posed Cauchy problem:

$$\begin{cases} \frac{d\varphi}{d\varepsilon}(\varepsilon) = -[D_{\varphi, \varphi}^2 \tilde{\Phi}(\varphi(\varepsilon), \varepsilon)]^{-1} \frac{\partial}{\partial \varepsilon} \nabla_\varphi \tilde{\Phi}(\varphi(\varepsilon), \varepsilon), \\ \varphi(0) = \varphi_w, \end{cases} \quad (22)$$

where, by Corollary 2, the initial value $\varphi(0)$ of φ when $\varepsilon = 0$ coincides with φ_w , the optimal potential for the two marginal optimal transport problem with cost w .

The next section is devoted to proving these properties.

3.2. Well posedness of the ODE. We refer the reader to appendix A for the computation of the second pure and mixed derivatives with respect to φ and the second mixed derivative with respect to φ and ε . In order to prove invertibility of $D_{\varphi,\varphi}^2 \tilde{\Phi}$ and well posedness of the ODE we need some lemmas giving uniform bounds on the potential φ and the eigenvalues of $D_{\varphi,\varphi}^2 \tilde{\Phi}$. We highlight that the following arguments are similar to (and largely inspired by) those in [20] (the main differences lie in the fact that we re-write the dual problem by using the Log-Sum-Exp function).

Lemma 7. *Let c_ε satisfy the boundedness assumption $\|c_\varepsilon\|_\infty \leq M$, $\forall \varepsilon \in [0, 1]$ ². Then the maximizer $\varphi(\varepsilon)$ of (17) subject to the normalization constraint (18) satisfies*

$$\|\varphi(\varepsilon)\|_\infty \leq 4M.$$

Proof. By the first order optimality condition $\nabla_\varphi \tilde{\Phi} = 0$ for (17) we deduce that each component of $\varphi(\varepsilon)$ is given by

$$\varphi_z = -\eta \log \left(\sum_y \sum_{\bar{x} \in X^{m-2}} \exp \left(\frac{\sum_{i=3}^m \varphi_{x^i} - c_\varepsilon(y, z, \bar{x})}{\eta} \right) (\otimes^{m-2} \rho)_{\bar{x}} \bar{\rho}_y \right)$$

It is easy to see that $\bar{\rho}_y$ can be bounded as follows

$$\frac{e^{-M/\eta} \rho_y}{\sum_{\bar{x} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^m \varphi_{x^i}}{\eta} \right) \tilde{\rho}_{\bar{x}}} \leq \bar{\rho}_y \leq \frac{e^{M/\eta} \rho_y}{\sum_{\bar{x} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^m \varphi_{x^i}}{\eta} \right) \tilde{\rho}_{\bar{x}}}.$$

Since we have imposed the normalization $\varphi_{x_0} = 0$ we get

$$\begin{aligned} \varphi_z = \varphi_z - \varphi_{x_0} &\leq -\eta \log \left(e^{-2M/\eta} \frac{\sum_{\bar{x} \in X^{m-2}} \exp \left(\frac{\sum_{i=3}^m \varphi_{x^i}}{\eta} \right) (\otimes^{m-2} \rho)_{\bar{x}}}{\sum_{\bar{x} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^m \varphi_{x^i}}{\eta} \right) \tilde{\rho}_{\bar{x}}} \right) \\ &\quad + \eta \log \left(e^{2M/\eta} \frac{\sum_{\bar{x} \in X^{m-2}} \exp \left(\frac{\sum_{i=3}^m \varphi_{x^i}}{\eta} \right) (\otimes^{m-2} \rho)_{\bar{x}}}{\sum_{\bar{x} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^m \varphi_{x^i}}{\eta} \right) \tilde{\rho}_{\bar{x}}} \right), \\ &\leq 4M, \end{aligned}$$

and the desired result immediately follows. $\square$

²Note that the boundedness $\|c_\varepsilon\|_\infty \leq M$ for some $M > 0$ follows immediately from our finite cost assumption on the finite set X^m .

Having established the above bounds, we aim to prove the well posedness of the Cauchy problem (22) on the set

$$U := \{\varphi : X \rightarrow \mathbb{R} \mid \varphi_{x_0} = 0, \|\varphi\|_\infty \leq 4M\}. \quad (23)$$

Lemma 8. $D_{\varphi,\varphi}^2 \tilde{\Phi}(\varphi, \varepsilon)$ and $\frac{\partial}{\partial \varepsilon} \nabla_\varphi \tilde{\Phi}(\varphi, \varepsilon)$ are Lipschitz with respect to φ on U .

Proof. This immediately follows from the fact that the second pure and mixed derivatives computed in Appendix A are easily seen to be C^1 , and their derivatives are all bounded on U . $\square$

In order to prove the invertibility of $D_{\varphi,\varphi}^2 \tilde{\Phi}$ we need the following lemma assuring the strong convexity of the Log-Sum-Exp function on the set U .

Lemma 9. Let $\Psi : \tilde{U} \rightarrow \mathbb{R}$ be defined on

$$\tilde{U}_C = \{\theta : X^{m-1} \rightarrow \mathbb{R} \mid \theta_{\bar{x}_0} = 0, \|\theta\|_\infty < C\}.$$

where $\bar{x}_0 = (x_0, \dots, x_0) \in X^{m-1}$, by $\Psi(\theta) = \sum_{y \in X} \log \left(\sum_{\bar{x} \in X^{m-1}} e^{\theta_{\bar{x}} - c_\varepsilon(y, \bar{x})} \tilde{\rho}_{\bar{x}} \right) \rho_y$.

Then Ψ is β -strongly convex for some $\beta > 0$.

Proof. It is enough to show strong convexity on this set of the function

$$f_y : \theta \in \tilde{U}_C \mapsto \log \left(\sum_{\bar{x} \in X^{m-1}} e^{\theta_{\bar{x}} - c_\varepsilon(y, \bar{x})} \tilde{\rho}_{\bar{x}} \right) = \log \left(e^{-c_\varepsilon(y, \bar{x}_0)} \tilde{\rho}_{\bar{x}_0} + \sum_{\bar{x} \in X^{m-1} \setminus \{\bar{x}_0\}} e^{\theta_{\bar{x}} - c_\varepsilon(y, \bar{x})} \tilde{\rho}_{\bar{x}} \right)$$

for a fixed y .

Enumerating the set $X^{m-1} \setminus \{\bar{x}_0\}$ of independent variables as $\bar{x}_j$ for $j \in (1, \dots, K)$ with $K = |X|^{m-1} - 1$, and denoting $z^j = e^{\theta_{\bar{x}_j} - c_\varepsilon(y, \bar{x}_j)} \tilde{\rho}_{\bar{x}_j}$ the Hessian of f_y is

$$\frac{1}{\left(e^{-c_\varepsilon(y, \bar{x}_0)} \tilde{\rho}_{\bar{x}_0} + \sum_j z^j \right)^2} \left(-z \otimes z + \text{diag}(z) \left(\sum_j z^j + e^{-c_\varepsilon(y, \bar{x}_0)} \tilde{\rho}_{\bar{x}_0} \right) \right)$$

The first two terms together constitute a positive semi-definite matrix (see [12], p.74), while the third is positive definite, with lower bound

$$\beta = \frac{e^{-C-2M} \tilde{\rho}_{\bar{x}_0} \min_{\bar{x}} \tilde{\rho}_{\bar{x}}}{(e^{C+M} \tilde{\rho}_{\bar{x}_0} + \sum_{\bar{x} \in X^{m-1} \setminus \{\bar{x}_0\}} e^{C+M} \tilde{\rho}_{\bar{x}})^2} = e^{-4M-3C} \tilde{\rho}_{\bar{x}_0} \min_{\bar{x}} \tilde{\rho}_{\bar{x}}.$$

It follows that f_y , and therefore Ψ , is β -convex on $\tilde{U}$. $\square$

Lemma 10. Let $\Lambda : U \rightarrow \tilde{U}_{4(m-1)M}$ be the linear mapping defined by $\Lambda(\varphi)_{\bar{x}} = \varphi_{x^1} + \dots + \varphi_{x^{m-1}} \quad \forall \bar{x} \in X^{m-1}$. Then $\tilde{\Psi}(\varphi) := \Psi(\Lambda(\varphi))$ is α -strongly convex on U .

Proof. By the linearity of Λ , one gets that, for $\varphi \in U$, $D^2 \tilde{\Psi}(\varphi)(v, v) = D^2 \Psi(\Lambda(\varphi))(\Lambda(v), \Lambda(v))$ for all $v \in U$. Thus,

$$D^2 \tilde{\Psi}(\varphi)(v, v) = D^2 \Psi(\Lambda(\varphi))(\Lambda(v), \Lambda(v)) \geq \beta \|\Lambda(v)\|^2.$$

Since $\|\Lambda(v)\|^2 \geq \sum_{x \in X} \|(m-1)v_x\|^2$ we finally get

$$D^2\tilde{\Psi}(\varphi)(v, v) \geq \alpha \|v\|^2,$$

with $\alpha = \beta(m-1)^2 > 0$, proving the α -strong convexity of $\tilde{\Psi}$. $\square$

Remark 11. *The α obtained in the Lemma above is not optimal; indeed we would have obtained a better lower bound on the eigenvalues of $D_{\varphi, \varphi}^2 \tilde{\Psi}$ by computing the smallest eigenvalue of $\Lambda^* \Lambda$. Moreover, in the previous lemma we take, for simplicity, $\eta = 1$, otherwise the parameter α would have taken the form $\alpha = e^{(-4C-3M)/\eta}(m-1)^2$. Notice now that α approaches to 0 as $\eta \rightarrow 0$, meaning the the condition number of the Hessian of $\tilde{\Phi}$ explodes. Namely, this will produce numerical instabilities.*

It easily follows from the previous lemma that $D_{\varphi, \varphi}^2 \tilde{\Phi} = D_{\varphi, \varphi}^2 \tilde{\Psi}$ is invertible on the set U ; we can then state the following result on the well posedness of (22).

Theorem 12. *Let $\varphi(\varepsilon)$ be the solution to (17) for all $\varepsilon \in [0, 1]$. Then $\varepsilon \mapsto \varphi(\varepsilon)$ is C^1 and is the unique solution to the Cauchy problem*

$$\begin{cases} \frac{d\varphi}{d\varepsilon}(\varepsilon) = -[D_{\varphi, \varphi}^2 \tilde{\Phi}(\varphi(\varepsilon), \varepsilon)]^{-1} \frac{\partial}{\partial \varepsilon} \nabla_{\varphi} \tilde{\Phi}(\varphi(\varepsilon), \varepsilon), \\ \varphi(0) = \varphi_w, \end{cases} \quad (24)$$

where φ_w is the optimal solution to (8) with cost w and two marginals equal to ρ .

Proof. As $\varphi(\varepsilon)$ minimizes $\tilde{\Phi}(\cdot, \varepsilon)$ for each fixed ε , we clearly have (20). Since $\tilde{\Phi}$ is clearly twice differentiable with respect to φ and ε and $D_{\varphi, \varphi}^2 \tilde{\Phi}$ is invertible by Lemma 10, the Implicit Function Theorem then implies that $\varepsilon \mapsto \varphi(\varepsilon)$ is C^1 and satisfies (21), or equivalently, (24).

Since $D_{\varphi, \varphi}^2 \tilde{\Phi}$ and $\frac{\partial}{\partial \varepsilon} \nabla_{\varphi} \tilde{\Phi}$ are Lipschitz continuous with respect to φ on U by Lemma 8 and clearly continuous with respect to ε , and since $D_{\varphi, \varphi}^2 \tilde{\Phi}$ is uniformly positive definite by Lemma 10, we have that

$$(\varphi, \varepsilon) \mapsto -[D_{\varphi, \varphi}^2 \tilde{\Phi}(\varphi(\varepsilon), \varepsilon)]^{-1} \frac{\partial}{\partial \varepsilon} \nabla_{\varphi} \tilde{\Phi}(\varphi(\varepsilon), \varepsilon)$$

is Lipschitz with respect to φ and continuous with respect to ε on U . Since by Lemma 7 $\varphi(\varepsilon) \in U$ for all ε , the Cauchy-Lipschitz Theorem then implies uniqueness of the solution to (24) on $U \times [0, 1]$, as desired. $\square$

4. ALGORITHM AND SIMULATION

In this subsection we present some numerical simulations³ obtained by discretizing the above ODE.

The algorithm consists in discretizing (24) by an explicit Euler scheme (notice that one could also use some high order method for the ODEs).

³The simulations have been performed in Python on 2 GHz Quad-Core Intel Core i5 MacBook Pro.

Let h be the step size and set $\varphi(0) = \varphi_w$ the solution of a 2 marginal problem with cost w , then the φ can be defined inductively as detailed in 1.

Algorithm 1 Algorithm to compute the φ via explicit Euler method

Require: $\varphi(0) = \varphi_w$
1: **while** $\|\varphi^{(k+1)} - \varphi^{(k)}\| < \text{tol}$ **do**
2: $D^{(k)} := D_{\varphi, \varphi}^2 \tilde{\Phi}(\varphi^{(k)}, kh)$
3: $b^{(k)} := -\frac{\partial}{\partial \varepsilon} \nabla_{\varphi} \tilde{\Phi}(\varphi^{(k)}, kh)$
4: Solve $D^{(k)} z = b^{(k)}$
5: $\varphi^{(k+1)} = \varphi^{(k)} + hz$
6: **end while**

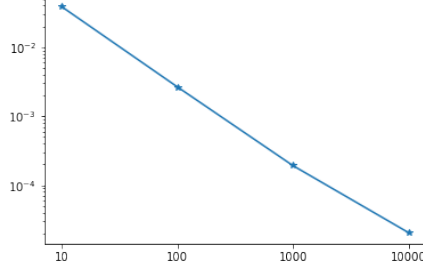
Notice that by the regularity we have proved above, we can conclude that the Euler scheme converges linearly. Moreover, the uniform error between the discretized solution obtained via the scheme and the solution to the ODE is $O(h)$.

Remark 13 (Complexity in the marginals). *We highlight that this ODE approach does not overpass the NP hardness (see [5]) of some multi-marginal problems, as the one with the Coulomb cost, and it still suffers the exponential complexity in the number of marginals. However, since the ODE is smooth, one can try apply an high order method (we will do a careful analysis of it later in the section) which can converge quickly in the number of iterations and have a computational time competitive with respect to Sinkhorn.*

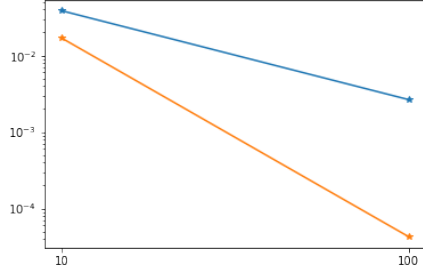
We also note that, if the symmetry assumptions are dropped, as discussed in Remark 4, one needs to solve a system of $(N-1)m$ ODEs. The matrix $D^{(k)}$ in Algorithm 1 is then be $(N-1)m \times (N-1)m$ instead of $m \times m$, increasing the leading term in the complexity by the relatively manageable multiplicative factor $(N-1)^2$. Numerical results for one such case are presented in Section 4.1 below.

In Figure 1a we plot the convergence order for the Euler scheme described above. The error is computed with respect to the solution to (17) computed via a gradient descent method with backtracking. Notice that the regularity of the objective function and the boundedness of the Hessian guarantee the convergence of the method. For these simulations we have taken $m = 3$, the uniform measure on $[0, 1]$ uniformly discretized with 100 gridpoints and the pairwise interaction $w(x, y) = -\log(0.1 + |x - y|)$. Moreover, since the RHS in (24) is regular one can try to apply an high order scheme to solve the Cauchy problem. In Figure 1b we compare the convergence of the Euler method and a Runge-Kutta of order 3; notice that with 100 time steps the RK method converges to a solution with an error of order 10^{-5} and by an estimation of the slope

of the two lines we obtain 3 and 1.16, respectively for RK and Euler (as expected).



(A) Linear convergence for an explicit Euler scheme.



(B) Comparison between explicit Euler (blue line) and an explicit Runge-Kutta (red line) of order 3

FIGURE 1. Order of convergence.

In Figure 2 we show the numerical result obtained with $\eta = 0.006$, $h = 1/1000$, $m = 3$, the uniform measure on $[0, 1]$ uniformly discretized with 100 gridpoints and the pairwise interaction $w(x, y) = -\log(0.1 + |x - y|)$. Notice that since we have developed our continuation method by the entropic regularization of optimal transport, we can easily reconstruct the optimal (regularized) plan at each time k by using the potential $\varphi^{(k)}$. Moreover, it is interesting to notice that the optimal plan at each time step of the ODE stays deterministic (taking into account the blurring effect of the entropic regularization); that is, it is concentrated on a low dimensional structure.

Remark 14 (Intermediate step of the ODE approach). *It is interesting to highlight that each i th-step of the ODE approach actually returns the optimal solution to the multi-marginal problem with the cost c_{kh} . This, in particular, implies that one can also use this approach in order to retrieve the solutions of many multi-marginal problems with different costs, by*

choosing a suitable c_ε , instead of solving these problems individually using Sinkhorn.

We are now interested in comparing the ODE approach and the Sinkhorn algorithm in terms of performance. In order to do this we consider the optimal solution of the regularized dual problem obtained with a gradient descent with backtracking and take it as the reference solution to compute the relative error $\frac{\|\varphi - \varphi_{ref}\|_\infty}{\|\varphi_{ref}\|_\infty}$. Concerning the ODE, since we have already remarked above that it is smooth, we consider different high order methods such as 3rd, 5th and 8th order Runge-Kutta methods. By looking at Table 1, it is clear that all the methods achieve almost the same relative error, but the number of iterations to reach it as well as the CPU time in seconds slightly differs. In particular we remark that a 3rd order RK is faster than a Sinkhorn in terms of computational time and iterations but less precise. The other RK methods achieve comparable precision to Sinkhorn with less iterations but the computational cost at the step of the ODE becomes now quite onerous implying a significant increase (especially 8th RK) in terms of CPU time.

	3rd RK	5th RK	8th RK	Sinkhorn
relative error	1.47×10^{-5}	7.8×10^{-6}	7.62×10^{-6}	5.46×10^{-6}
iterations	87	87	87	820
CPU time (sec)	72.39	158.9	385.1	102.8

TABLE 1. Comparison between the ODE approach and Sinkhorn for the uniform density and 400 gridpoints

In Figures 3-4 we have kept the same data as before, but we have chose the negative harmonic cost, that is $w(x, y) = -|x - y|^2$. In particular in Figures 3a and 3b we compare the potential obtained with Sinkhorn and the one at time 1 of the ODE for the log cost and the negative harmonic, respectively. We highlight that the solution at $\varepsilon = 0$ is $-Id$ and then the final coupling is supported, as expected, on the hyperplane $x + y + z = 1.5$. Looking at Figure 4, notice that in this case the optimal transport plan is deterministic at the initial time, and also shortly after, but then it is possible to notice an interesting transition to a spread transport plan.

4.1. The non-symmetric ODE: Euler case. In a series of papers [13–15] Brenier proposed a relaxation of the incompressible Euler equation with constrained initial and final data interpreted as a geodesic on the group of measure preserving diffeomorphisms. Brenier’s relaxed formulation consists in finding a probability measure over absolutely continuous paths which minimizes the average kinetic energy. In this framework the incompressibility is encoded by an additional constraint that at each

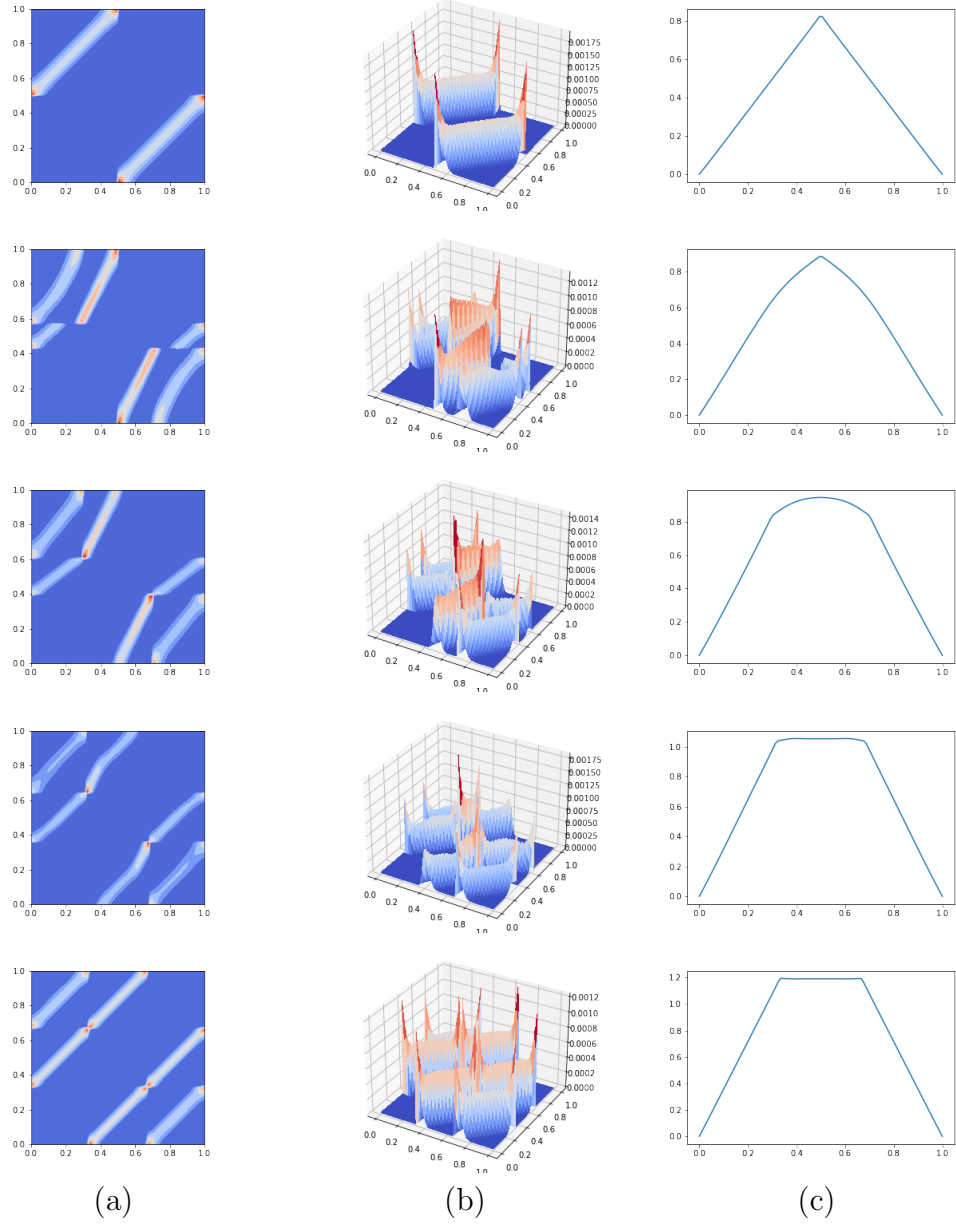


FIGURE 2. (Log cost) Column (a): support of the coupling $\gamma_{1,2}^\varepsilon$. Column (b): surface of the coupling $\gamma_{1,2}^\varepsilon$. Column (c): potential $\varphi(\varepsilon)$. For $\varepsilon = 0, 0.25, 0.5, 0.75, 1$

time t , the distribution of position need be uniform. If we consider a uniform discretization of $[0, T]$ (where T is the final time) with m steps in time, we recover a multi-marginal formulation of the Brenier principle with the specific cost function

$$c(x^1, \dots, x^m) = \frac{m^2}{2T^2} \sum_{i=1}^{m-1} |x^{i+1} - x^i|^2$$

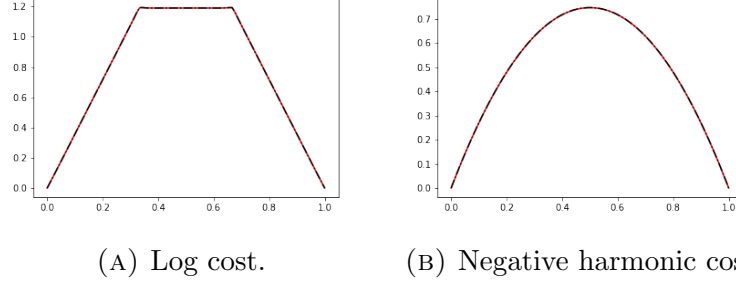


FIGURE 3. Optimal potential computed via Sinkhorn (red line). Potential computed via the ODE (black dot-dashed line).

representing the discretized, in time, kinetic energy ($|\cdot|$ denotes the standard euclidean norm). The coupling γ is the probability to find a *generalized particle* on the discrete path $x^1, \dots, x^m$. Because the fluid is incompressible, the i th marginal μ^i of γ is the uniform on the d -dimensional cube $[0, 1]^d$. For each $i \in \{1, \dots, m\}$ the transition probability from time 1 to time i is given by the coupling

$$\gamma_{1,i}\pi^i(x^1, x^i) = \left((x^1, \dots, x^m) \mapsto (x^1, x^i) \right)_{\#} \gamma,$$

which represents the probability of finding a *generalized particle* initially at x^1 to be at x^i at time i . In order to impose the initial and final constraint we include by penalization, that is by adding a term to the cost function which now reads as

$$c(x^1, \dots, x^m) = \frac{m^2}{2T^2} \sum_{i=1}^{m-1} |x^{i+1} - x^i|^2 + \beta |F(x^1) - x^m|^2,$$

where $\beta > 0$ is a penalization parameter in order to enforce the initial-final constraint. $F(x_1)$ represents the prescribed final position of the particle initially at position x_1 , and the coupling γ can be interpreted as a (generalization of) a discrete time geodesic between the identity mapping and F on the space of measure preserving maps.

For an efficient implementation of Sinkhorn in this case we refer the reader to [9, 10, 44]. If we consider now the ODE setting, we have now to deal with a non symmetric case (the cost is not symmetric in the marginals anymore) and so to solve a system, still well posed, of ODEs. In particular we consider the following c_ε cost

$$c_\varepsilon(x^1, \dots, x^m) = \frac{m^2}{2T^2} |x^2 - x^1|^2 + \varepsilon \left(\frac{m^2}{2T^2} \sum_{i=2}^{m-1} |x^{i+1} - x^i|^2 \right) + \beta |F(x^1) - x^m|^2.$$

For the numerical simulations we took 100 gridpoints discretization of $[0, 1]$, $m = 9$ time marginals, $\beta = 20$ and $\eta = 0.002$. We solved the ODE system by using a 5th order Runge-Kutta method and $h = 1/100$. In

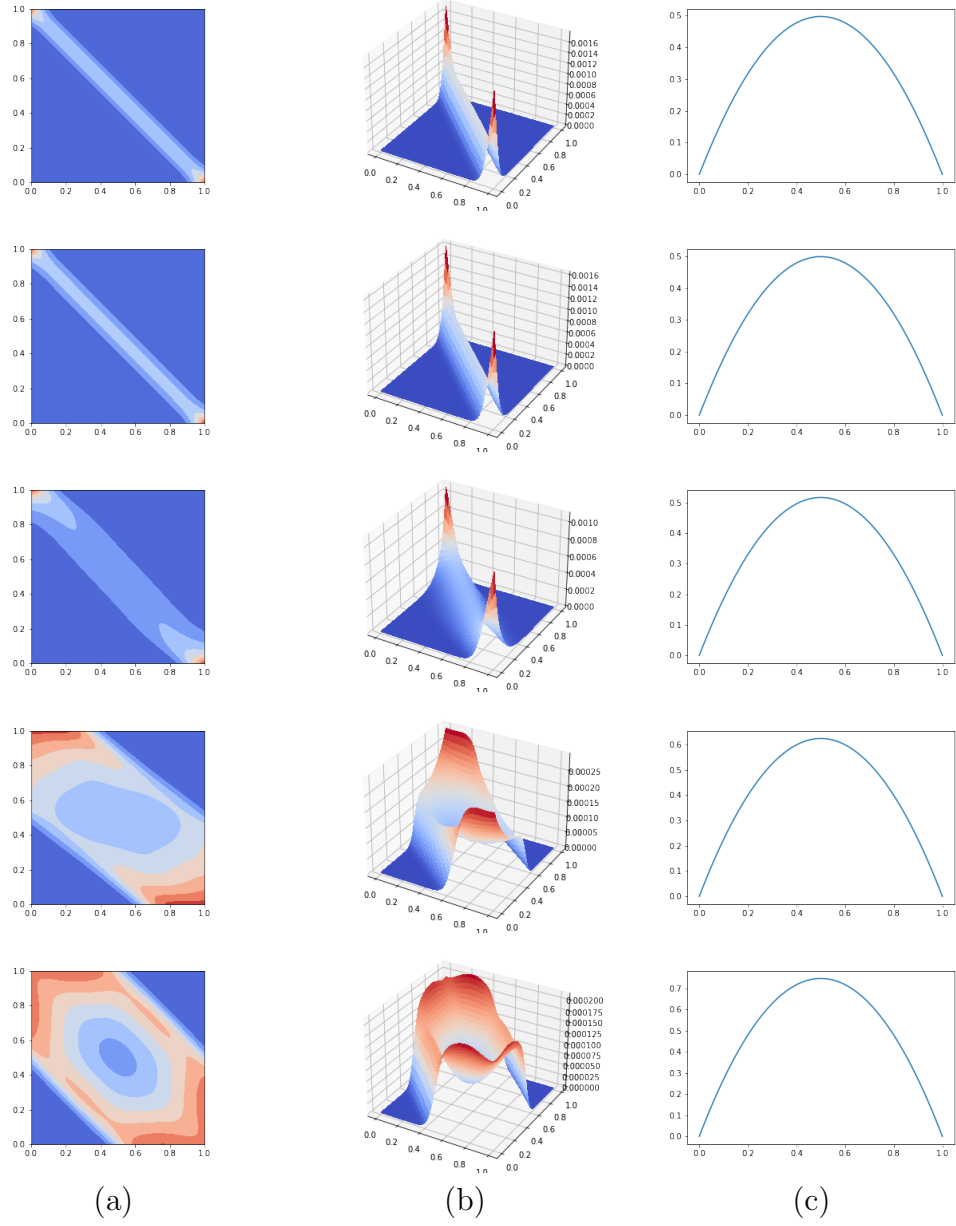


FIGURE 4. (Negative Harmonic cost) Column (a): support of the coupling $\gamma_{1,2}^\varepsilon$. Column (b): surface of the coupling $\gamma_{1,2}^\varepsilon$. Column (c): potential $\varphi(\varepsilon)$. For $\varepsilon = 0, 0.25, 0.5, 0.75, 1$

figure 5 we compare the potentials obtained via the ODE approach at $\varepsilon = 1$ and Sinkhorn for two different initial-final configurations : $F(x) = 1 - x$ figure 5a and $F(x) = (x + 1/2) \bmod 1$ in figure 5b. Notice in figure 5a that for the particular choice of initial-final configuration the potentials φ_1 and φ_m coincides as well as φ_i for $i = 1, \dots, m - 1$. In figure 6 we plot the transition couplings for the initial-final configurations considered

above, at $\varepsilon = 1$. As for the symmetric case we compare the number of iteration and the CPU time of the ODE approach and Sinkhorn in table 2. Notice that Sinkhorn performs better than the 3rd order RK in terms of relative error (the reference solution for the error has been computed again via a gradient descent with backtracking).

	3rd RK	5th RK	8th RK	Sinkhorn
relative error	8.33×10^{-4}	3.5×10^{-5}	3.51×10^{-6}	3.5×10^{-6}
iterations	92	92	92	6649
CPU time (sec)	321.58	790.17	1225.45	327.82

TABLE 2. Comparison between the ODE approach and Sinkhorn for the Euler case with $F(x) = (x + 1/2) \bmod 1$, 100 gridpoints and 9 time marginals.

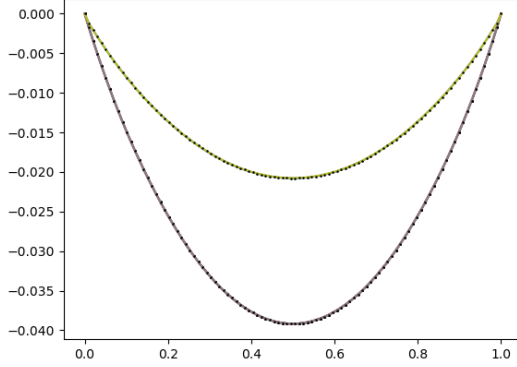
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APPENDIX A. SECOND DERIVATIVES OF $\tilde{\Phi}$

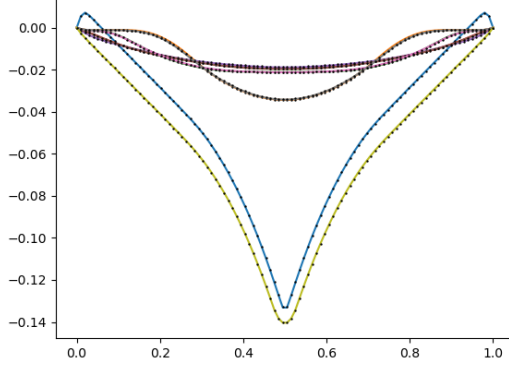
In this appendix we detail the second derivatives of $\tilde{\Phi}$ with respect to φ and ε . Let us consider firstly the term (21)

$$\frac{\partial}{\partial \varepsilon} \nabla_{\varphi} \Phi(\varphi, \varepsilon),$$

notice that $\nabla_{\varphi} \Phi(\varphi, \varepsilon)$ it is a composition of an exponential with a linear function in ε , meaning that it is differentiable with respect to ε . We



(A) Comparison of potentials for $F(x) = 1 - x$.



(B) Comparison of potentials for $F(x) = x + 1/2 \bmod 1$.

FIGURE 5. Optimal potentials computed via Sinkhorn (black dots). Potential computed via the ODE (colored solid lines).

obtain then the following

$$\frac{\partial}{\partial \varepsilon} (\nabla_{\varphi} \Phi(\varphi, \varepsilon))_{\varphi_z} = -\frac{\exp(\varphi_z/\eta)}{\eta} \sum_{\bar{x} \in X^{m-1}} \partial_{\varepsilon} (c_{\varepsilon}(z, \bar{x})) \exp \left(\frac{\sum_{i=1}^{m-1} \varphi_{x^i} - c_{\varepsilon}(z, \bar{x})}{\eta} \right),$$

$$\begin{aligned} \frac{\partial}{\partial \varepsilon} \nabla_{\varphi} \Phi(\varphi(\varepsilon), \varepsilon) &:= -\frac{1}{\eta} \exp(\varphi/\eta) \rho \left(\sum_{(x, \bar{y}) \in X^{m-1}} \partial_{\varepsilon} c_{\varepsilon}(x, z, \bar{y}) \exp \left(\frac{\sum_{i=3}^m \varphi_{y^i} - c_{\varepsilon}(x, z, \bar{y})}{\eta} \right) \tilde{\rho}_{\bar{y}} \bar{\rho}_x \right. \\ &\quad \left. + \sum_{(x, \bar{y}) \in X^{m-1}} \exp \left(\frac{\sum_{i=3}^m \varphi_{y^i} - c_{\varepsilon}(x, z, \bar{y})}{\eta} \right) \tilde{\rho}_{\bar{y}} \frac{\sum_{\bar{w} \in X^{m-1}} \partial_{\varepsilon} c_{\varepsilon}(x, \bar{w}) \exp \left(\frac{\sum_{i=2}^m \varphi_{w^i} - c_{\varepsilon}(x, \bar{w})}{\eta} \right) \tilde{\rho}_{\bar{w}}}{\sum_{\bar{w} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^m \varphi_{w^i} - c_{\varepsilon}(x, \bar{w})}{\eta} \right) \tilde{\rho}_{\bar{w}}} \bar{\rho}_x \right) \end{aligned}$$

where

$$\partial_{\varepsilon} c_{\varepsilon}(z, \bar{x}) = \sum_{i,j=2, i \neq j}^{m-1} w(x^i, x^j).$$

Concerning the second derivative with respect to φ . it is again quite easy to see that Φ is twice differentiable, then we have

$$D_{\varphi, \varphi}^2 \tilde{\Phi} = \frac{1}{\eta} \text{diag}(I_1) + \frac{m-2}{\eta} (e^{\varphi/\eta} \rho) \otimes (e^{\varphi/\eta} \rho) I_2 - \frac{m-1}{\eta} \sum_{y \in X} I_3^y \otimes I_3^y \rho_y,$$

where

$$\begin{aligned} (I_1)_z &= e^{\varphi_z/\eta} \rho_z \sum_y \sum_{\bar{x} \in X^{m-2}} \exp \left(\frac{\sum_{i=3}^m \varphi_{x^i} - c_{\varepsilon}(y, z, \bar{x})}{\eta} \right) (\otimes^{m-2} \rho)_{\bar{x}} \bar{\rho}_y, \\ (I_2)_{z,w} &= \sum_y \sum_{\bar{x} \in X^{m-3}} \exp \left(\frac{\sum_{i=4}^m \varphi_{x^i} - c_{\varepsilon}(y, z, w, \bar{x})}{\eta} \right) (\otimes^{m-3} \rho)_{\bar{x}} \bar{\rho}_y, \\ (I_3^y)_z &= \frac{e^{\varphi_z/\eta} \rho_z \sum_{\bar{x} \in X^{m-2}} \exp \left(\frac{\sum_{i=3}^m \varphi_{x^i} - c_{\varepsilon}(y, z, \bar{x})}{\eta} \right) (\otimes^{m-2} \rho)_{\bar{x}}}{\sum_{\bar{x} \in X^{m-1}} \exp \left(\frac{\sum_{i=2}^m \varphi_{x^i} - c_{\varepsilon}(y, \bar{x})}{\eta} \right) \tilde{\rho}_{\bar{x}}} \end{aligned}$$

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UNIVERSITÉ PARIS-SACLAY, CNRS, LABORATOIRE DE MATHÉMATIQUES D’ORSAY, 91405, ORSAY, FRANCE.

Email address: `luca.nenna@universite-paris-saclay.fr`

DEPARTMENT OF MATHEMATICAL AND STATISTICAL SCIENCES, 632 CAB, UNIVERSITY OF ALBERTA, EDMONTON, ALBERTA, CANADA, T6G 2G1

Email address: `pass@ualberta.ca`

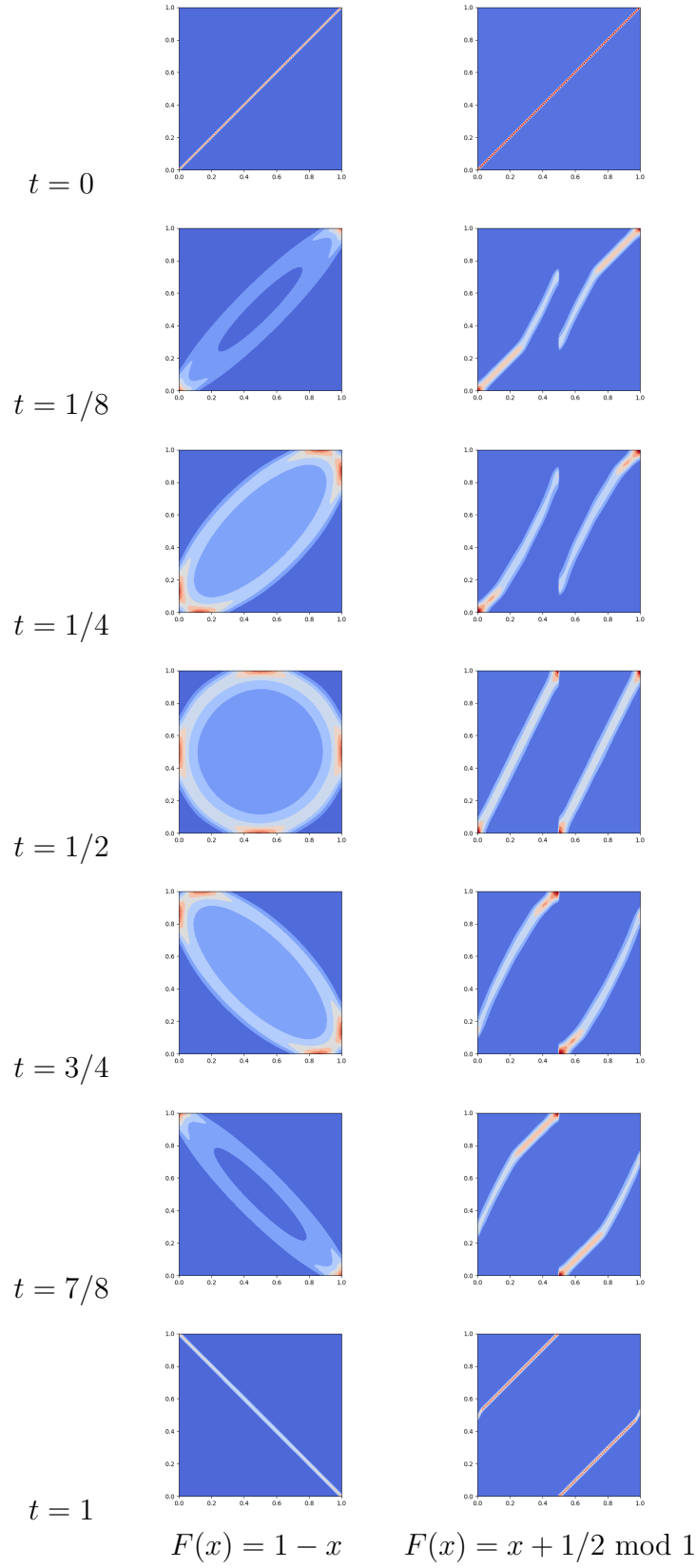


FIGURE 6. Transition plan $\gamma_{1,i}$ for different final configuration of Brenier relaxed formulation of incompressible Euler equations.