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Microseconds Dynamics Simulations of the Outer–Membrane Protease T

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Abstract

Conformational fluctuations of enzymes may play an important role for Substrate recognition and/or cayalysis, as it has been suggested in the case of the protease enzymatic superfamily (Carnevale et al. J .Am. Chem. Soc. 128: 9766–9772). Unfortunately, theoretically addressing this issue is a problem of formidable complexity, as the number of the involved degrees of freedom is enormous: indeed, the biological function of a protein depends, in principle, on all its atoms and on the surrounding water molecules. Here we investigated a membrane protease enzyme, the OmpT from *E. coli*, by hybrid molecular mechanics/coarse–grained approach (MM/CG), in which the active site is treated with the GROMOS force field, whilst the protein scaffold is described with a Go–Model. The method has been previously tested against results obtained with all–atom simulations. Our results show that the large scale motions and fluctuations of the electric field in the μs time scale may impact on the biological function and suggest that OmpT employs the same catalytic strategy as aspartic proteases. Such a conclusion can not be drawn within the 10–100 ns time scale typical of current molecular dynamics simulations. In addition, our studies provide a structural explanation for the drop in the catalytic activity of two known mutants (S99A and H212A), suggesting that the CG approach is a fast and reliable tool for providing structure/function relationships for both wild type OmpT and mutants.

Key words: molecular dynamics; conformational fluctuations; Coarse–grained; microseconds; hybrid model; Omptin.

Introduction

Recently, molecular dynamics (MD) studies have lead to the suggestion that enzymatic function results from a subtle interplay between chemical kinetics and molecular motion. Examples include the dihydrofolate reductase enzyme (1), whose completion of the catalytic cycle has been suggested to require coupled motions, and the proteases superfamily (2–4), whose biological function has been proposed to be affected by conformational fluctuations. These conclusions are being corroborated by a variety of experimental biophysical studies (5–10).

A major limitation of MD in addressing this issue is obviously the accessible time scale (currently up to the subμs timescale). In principle, one could achieve much longer time scales by coarse–grained (CG) potential based MD approaches (11, 12). However, while CG potentials are useful to understand large–scale phenomena such as the folding process (13) they can not describe the exquisite molecular recognition events among enzymes and their substrates, which are key for the enzymatic function.

Recently, we have presented a hybrid approach, the MM/CG model (14), with the goal to preserve the advantages of both MM and CG approaches, *i.e.* the necessary details associated with the biological activity and the "large" accessible time scale. Specifically, the amino acid residues involved in the ligand binding are treated with atomic detail (MM region) by means of a molecular mechanics force field, whereas the rest of the protein is treated at the CG level. This approach allows a fast and efficient description of the mechanical coupling between the dynamics of the active site with the

enzymatic substrate (the MM region) and that of the protein environment (the CG region).

So far, MM/CG simulations have been applied to selected members of the aspartic protease family. These calculations were meant as a validation of the approach *versus* all–atom MD simulations (14). The MM/CG approach turned out to correctly reproduce both the local and the global features of two cytoplasmatic proteins (the aspartic proteases from HIV-1 PR (2) and β–secretase (BACE) (3)). Furthermore, the calculations reproduced the structural fluctuations of the substrate in the binding cavity, which play an important role for the enzymatic activity (2, 3). These test calculations covered the same timescale as reference MD simulations used for comparison (10 ns for HIV-1 PR and 8 ns for BACE). However, the computational cost of this scheme turned out to be about two orders of magnitude smaller than that of the corresponding all–atom MD (14).

In this paper, we use the MM/CG approach to study the relevant case of the Michaelis complex of a membrane protease, the outer–membrane protease T (OmpT) (15) on the μs time scale.

OmpT is a defense protein expressed by Gram–negative bacteria belonging to the omptin protein family (16), which has been shown to be important for the virulence of *Y. pestis* and clinical *E. coli* isolates (17). The protein cleaves peptides preferentially between two consecutive basic amino acid residues (18, 19).

The isoenzyme from *E. coli*, for which the X–ray structure is available (15), features a deep groove formed by loops L4 and L5 on one side and L1, L2, L3 on the other one (Figure 1a). This groove constitutes the

active site of the enzyme. Initially, OmpT was classified as a novel–type serine protease; however, the relatively large distance between the putative catalytic S99 and H212 (∼9Å) observed in the X–ray structure (15) has suggested that the protease may work by a novel mechanism involving the D210–H212 and D83–D85 pairs. H212 and D83 groups have been suggested to activate a water molecule for the nucleophilic attack, while D85 and D210 contribute to polarizing the substrate scissile peptide bond (18, 19). A peculiar H–bond network formed by these residues orients both the substrate and the nucleophilic water, promoting the cleavage of the peptide bond. MD simulations of the enzyme on a 10–ns time scale have further supported this scenario (20, 21).

A comparison of MM/CG and all–atom MD simulations of the enzyme in the free state shows that the MM/CG approach is equally well suited for membrane proteins (22) and suggests that fluctuations of the active site cleft may play a role for substrate recognition. Here, we extend our investigations to the Michaelis complex of the protein with a model substrate (Ala–Arg–Arg–Ala). We carried out simulations on four OmpT/ARRA Michaelis complexes (**A**–**D**), which differ for: (i) the protonation state of the putative catalytic residue D83; (ii) for the N– and C– terminal tails of the substrate, which are considered in the Zwitterionic form (charged state) or capped with acetyl and N–methyl groups (neutral state), respectively, following all–atom MD calculations on the same system.

The aim of our MM/CG simulations is twofold. First, we investigated the role of conformational fluctuations of the substrate in the active site on the μs time–scale. Results show that both large–scale motions of the

protein and the electrostatic field, evaluated on such a time–scale, impact on the function of the enzyme. Second, we used the MM/CG approach to investigate the effects of mutations for the enzymatic function. We focused on protein mutants S99A and H212A, which are experimentally known to be much less efficient than the wild–type (wt) (residual activity within 0% and 4% relative to the wt (18)). Our calculations, on the ~ 0.1 μs time scale, provide a structural basis for the dramatic decrease in the catalytic activity. Because of its relatively cheap computational cost (two order of magnitude faster than standard all–atom MD), the methodology may be helpful to investigate structure/function relationships of high–throughput site–directed mutagenesis data.

Methods

The MM/CG model

A detailed description of the MM/CG approach is reported elsewhere (14), therefore we only summarize its principles here.

A small part of the protein (*e.g.* the enzymatic active site) is investigated in atomic detail, while the rest is treated with a CG approach using a modified Go–model (23) by only considering C_α centroids. An interface region (I) is located between the two MM and CG regions, bridging the large discontinuity between full–atom and CG descriptions. The total potential energy of the system reads:

$$V = E_{MM} + E_{CG} + E_I + E_{MM/I} + E_{CG/I} + E_{SD} \ ,$$

where the first three terms represent the interactions within the MM, CG and I regions, respectively, whereas the fourth and fifth represent the cross-terms potentials. The last term, E_{SD}, mimics stochastic and frictional forces acting on the system due to the solvent (24, 25).

In regions MM and I, all atoms are explicitly considered and, consequently, E_{MM}, E_I and $E_{MM/I}$ energy terms have all the same formulation (*i.e.* the GROMOS96 43a1 force field (26)).

E_{CG} takes the following form:

$$E_{CG} = \frac{1}{4}\sum_i K_b(|\mathbf{R}_i - \mathbf{R}_{i+1}|^2 - b_{i\ i+1}^2)^2 + \sum_{i>j} V_0\left[1 - \exp\left(-B_{ij}(|\mathbf{R}_i - \mathbf{R}_j| - b_{ij})\right)\right]^2 \ . \quad (1)$$

The first term in Eq. (1) takes into account bonded interactions between consecutive CG (C_α) centroids, identified by the position vectors $\mathbf{R}_i$ and $\mathbf{R}_{i+1}$, and K_b is the relative bond force constant. b_{ij} is the equilibrium distance, corresponding to the native distance between CG atoms. The second term in Eq. (1) describes the non-bonded interactions between CG atoms. V_0 is the interaction well depth and B_{ij} is the modulating exponent of the Morse potential. The latter, along with V_0, have been obtained so as to reproduce features of all–atom MD calculations on a test system (14).

At the interface between the I and CG regions, bonds between consecutive C_α belonging to the I and CG regions ensure backbone connectivity. In addition, a term describing the non–bonded interactions is added: the interface atom, that is either a C_α or a C_β, interacts with the C_α in the CG

region with the same potential as in Eq. (1).

A simple solvation model using explicit water molecules is also implemented: if the active site of the enzyme is solvent exposed a water drop is constructed around the MM and I regions. If a water molecule exits from the drop, its velocity is reflected toward the inside (22) to avoid evaporation.

Investigated Systems

We focused on two OmpT protomers in complex with their substrate ARRA. The simulations differ in the protonation state of D83 at the active site and in the N– and C– termini of the peptide (see Table 1). The protonated aspartate might orient the putative attacking water molecule. The neutral termini help to reduce artefacts relative to the short length of the peptide substrate. The catalytic histidine residue H212 was assumed to be δ-protonated as in (20) (Figure 1bc). Prior to the MM/CG simulations we equilibrated the systems **A**, **B**, **C** and **D** via 24 ns of standard all–atom MD starting from the complex model described in (20). S99A and H212A mutants were built by simply replacing these residues with alanine. We also performed MD simulations of a reference system (diglycine) to compare the electrostatic properties of key residues in the active site of OmpT. The systems underwent MM/CG simulations for the time scales summarized in Table 1.

Computational Details

Preliminary MM simulations on four protomers of the OmpT–ARRA complex described in (20) were performed for 24 ns using the Yasara software

Table 1: Systems and mutants of OmpT/ARRA complexes investigated in this work.

Complex id	**Simulated Time (μs)**
Production runs at the MM/CG level	
A: protonated D83, neutral N– and C–termini (AceARRANMe)	1.00
B: deprotonated D83, neutral N– and C–termini (AceARRANMe)	1.00
C: protonated D83, charged N– and C–termini (NH_3^+ARRACOO$^-$)	0.05
D: deprotonated D83, charged N– and C–termini (NH_3^+ARRACOO$^-$)	0.05
S99A–**A**	0.16
S99A–**B**	0.16
H212A–**A**	0.15
H212A–**B**	0.15
All-atom reference simulations	
Aeq: equilibration	0.024
Beq: equilibration	0.024
Ceq: equilibration	0.024
Deq: equilibration	0.024
Diglycine (AceGGNme)	0.015

and the associated Yamber2 force field (27). Details of these equilibration simulations are provided as supplementary material.

Hybrid MM/CG simulations were performed using a modified version of the Gromacs 3.2.1 program (14). The enzymatic active site was treated at atomic detail with the GROMOS96 43a1 force field (26), as aforementioned (MM and I regions in Figure 1a). $\sim$ 400 SPC water molecules (28) were added. This constitutes a water layer of $\sim$ 15 Å around the MM region (Figure 1a). This approach has been shown to accurately describe structural and dynamic features of the active site of two aspartic proteases in complex with their substrate (14) and of OmpT in the free state (22). The entire systems were composed of $\sim 4,400$ particles.

The analyses of the trajectories of **A** and **B** were performed over the last 0.95 μs, that is after the equilibration phase (Figure 1SIab). The analyses of the trajectories of **C** and **D** were not carried out, as we will discuss later in the results section.

The leap–frog stochastic dynamics algorithm was used to integrate the equations of motion with a time step $\Delta t = 2$ fs and a friction coefficient $\gamma_i = m_i/\tau$, where $\tau = 0.5$ ps is the time constant for the coupling and m_i is the mass of the ith particle. A cut–off distance of 14 Å was used for the electrostatics [1]. A cut–off of 14 Å was also used for the van der Waals interactions. The pair list was updated every 10 steps. The SHAKE algorithm (29) was used to keep bonds containing hydrogens at a fixed length.

Following the 24 ns of equilibration at an all–atom MM level, the systems were relaxed by a 1–ns MM/CG simulation with positional restraints on the OmpT/ARRA complexes to minimize the energy of the solvent. Then further 1 ns with positional restraints on OmpT were performed to allow the ligand to accommodate itself inside the binding pocket under the MM/CG potentials.

Finally, we performed an atomic–force field based MD simulation of a reference system. This is diglycine (AceGGNMe) in a periodic box of 779 water molecules. The GROMOS96 43a1 (26) and SPC (28) force fields were used for the dipeptide and water, respectively. Room conditions (T=300 K, P=1 bar) were achieved by coupling the system with a Berendsen thermo-

[1] This very crude assumption in the treatment of the electrostatics appears to be justified by the simplicity of the model used. Careful checks were made on energy conservation. In addition, test calculations with a longer cut–off for electrostatics (36 Å) provided very similar results to those with the shorter cut-off (data not shown).

stat (30) with τ=1.0 ps and a Berendsen barostat (30) with compressibility of 4.5 10^{-10}bar^{-1} in all three dimensions. The time step of the integration was 2 fs. Electrostatic and van der Waals interactions were calculated using a cut–off of 18 Å and 14 Å, respectively. 0.015 μs of trajectory were collected.

The following properties were calculated: (i) large scale motions as eigenvectors of the covariance matrix of the C_α's (31); (ii) the analysis of the cosine content (32, 33); (iii) the electric field in the MM region, using the electrostatic term in the GROMOS96 43a1 force field (26). This analysis has been used for qualitative comparisons between different systems.

Results

MM/CG of systems A and B

The structure of **A** and **B** (Figure 1bc) is well maintained within the time scale investigated (1 μs): the C_α root mean square deviations (RMSD) of **A** and **B** rise during the first $\sim$0.05μs and then fluctuate around an average value for the remainder of the simulations (Figures 1SIab), suggesting that this specific protein requires large sampling to equilibrate. The analysis of the cosine content (C_i) (32, 33) ascertains the convergence of the simulations: C_i computed for the first 8 eigenvectors of the covariance matrix is lower than 0.5 (see Figure 1SIcd), meaning that the largest fluctuations are related to the potential.

A and **B** appear to be productive Michaelis complexes: a water molecule bridges H212 and D83, thus pointing towards the substrate carbonyl carbon.

Table 2: MM/CG simulations of **A** and **B**, H212A–**A**, H212A–**B**, S99A–**A** and S99A–**B**: selected MD–averaged distances (Å) within the active site of the protein. Standard deviations (STD) are reported in parenthesis.

	A	**B**	H212A–**A**	H212A–**B**	S99A–**A**	S99A–**B**
Hδ@H212–Oδ@D210	2.6(±0.8)	2.5(±0.7)	//	//	7.0(±3.5)	2.8(±0.7)
Hδ@H212–O@D210	5.2(±0.8)	4.3(±1.2)	//	//	7.0(±2.1)	5.0(±0.8)
Nϵ@H212–H@R2	4.7(±0.7)	5.0(±1.6)	//	//	6.4(±2.2)	5.0(±0.7)
Nϵ@H212–C@R2	4.4(±0.5)	6.2(±1.4)	//	//	6.5(±1.9)	4.4(±0.7)
Nϵ@H212–Oδ@D83	5.8(±1.0)	8.0(±1.5)	//	//	11.3(±1.8)	7.9(±1.4)
Cζ@R2–Cδ@E27	4.8(±0.5)	5.0(±0.8)	6.7(±2.4)	5.0(±0.7)	6.2(±1.7)	6.3(±1.4)
Cζ@R2–Cγ@D208	5.0(±0.7)	5.2(±0.9)	9.0(±2.3)	8.4(±0.3)	7.6(±2.1)	6.6(±0.7)
Cζ@R2–Cγ@D210	8.5(±0.7)	8.0(±0.9)	10.1(±3.0)	4.0(±0.4)	9.5(±1.6)	6.8(±1.2)
Cζ@R3–Cγ@D85	6.0(±1.0)	6.5(±1.3)	9.7(±1.6)	7.0(±1.7)	9.4(±1.4)	6.7(±1.2)
Cζ@R3–Cγ@D97	7.0(±1.5)	8.0(±1.8)	11.6(±2.1)	9.0(±1.6)	11.8(±3.4)	15.0(±3.1)
Hϵ@H101–Oγ@S99	3.1(±0.9)	4.2(±1.2)	4.8(±1.0)	3.0(±0.7)	//	//
Oδ@D97–Hγ@S99	5.5(±1.3)	6.6(±2.1)	3.5(±2.3)	8.2(±2.6)	//	//
Hδ@D83–O@R2	4.5(±1.1)	//	5.5(±1.8)	//	7.5(±1.4)	//
Hδ@D83–Oδ@D97	8.8(±1.2)	//	7.9(±1.3)	//	4.0(±1.2)	//
Hϵ@H101–Oδ@D83	3.7(±1.0)	3.5(±1.1)	4.6(±2.0)	3.1(±0.9)	5.0(±1.1)	3.5(±0.9)
H@D83–O@S(A)99	2.0(±0.2)	2.0(±0.3)	2.0(±0.2)	2.0(±0.2)	2.2(±0.4)	4.0(±0.6)
O@D83–H@S(A)99	2.0 (±0.3)	2.4(±0.5)	2.5(±0.6)	2.0(±0.2)	2.6(±0.7)	2.5(±0.9)
H@D85–O@D97	2.1(±0.4)	2.3(±0.3)	2.3(±0.7)	2.4(±0.7)	3.0(±1.2)	5.0(±1.8)
O@D85–H@D97	2.0(±0.4)	2.6(±1.0)	2.1(±0.4)	2.7(±0.7)	3.5(±1.6)	5.3(±1.7)

The average distance during the simulations between the catalytic water and the carbon of the scissile CO-NH bond in **A** and **B** is 3.91(±0.35)Å and 3.50(±0.30)Å, respectively. This water molecule is located in a position prone for a nucleophilic attack on the carbonyl group of the substrate (see Figure 2ab). At times a second water molecule bridges H212 and D83 and the H212–water(s)–D83 interaction is replaced by a direct H–bond between D83 and S99. Most of the key contacts at the active sites are maintained during the dynamics of both systems (Table 2, Figure 1bc): (i) Hδ@H212 H–bonding to Oδ@D210 or to O@D210; (ii) the salt bridge between substrate R2 and E27 and D208; (iii) the salt bridge between substrate R3 and D97 and D85; (iv) H S99 backbone H–bonding to D83 backbone; (v) D85 backbone H–bonding to D97 backbone. In addition (Table 2, Figure 1bc),

(i) at times, S99 H–bonds to D85 and D97 and Nϵ@H212 H–bonds with H@R3 of the substrate; (ii) Hϵ@H101H-bonds to either S99 or D83.

Structural fluctuations of the substrate in the binding cavity have been shown to play a functional role for cytoplasmatic proteases (4) such as HIV-1 protease (2) and BACE (3). For these enzymes it was found that the distance between the catalytic dyad and the substrate fluctuates around characteristic values corresponding to different mutual positions of the catalytic water relative to the substrate carbonyl carbon: only conformations in which the distance between the enzyme and the substrate is at a minimum turned out to be catalytically efficient (2, 3). Because those motions are correlated to the the large–scale motions of the proteins, the enzyme might play a role for the reaction by steering the substrate into its appropriate reactive conformation.

In order to ascertain whether this is also the case for OmpT, we have chosen to monitor the substrate motion within the β–barrel: (i) the distance δ of the center of mass of the ARRA peptide from the center of mass of the D83–D85 pair of residues; δ is affected by the width of the cleft and, therefore, its fluctuations also modulate the position of the water inside the catalytic cleft. (ii) the distance ξ of the center of mass of ARRA from the center of mass of the β–barrel (see Figure 3, left panel). ξ is affected by the distance between the substrate and the catalytic water. Thus, δ and ξ can be used as suitable descriptors of the enzyme "active" conformations (4) to evaluate the presence of a functional mechanical coupling between the substrate and large–scale conformational fluctuations of the entire enzyme.

In **A**, $\langle\delta\rangle \sim 7.8(0.5)$ Å and $\langle\xi\rangle \sim 24.1(0.4)$ Å, showing a sharp Gaussian–

like distribution (see Figure 5SIab). In **B**, both quantities feature bimodal distributions: δ fluctuates from $\delta_1 = 7.3$ Å to $\delta_2 = 9.8$ Å (Figure 3a) and ξ from $\xi_1 = 25.9$ Å to ξ_2 27.5 Å (Figure 3b). Interestingly, the transitions from ξ_1 to ξ_2 and from δ_1 to δ_2 occur at the same time, suggesting that the oscillation of the cleft is correlated with the oscillation of the substrate along the axis of the β–barrel (the linear correlation between the two data sets δ and ξ is 0.60 over $\sim 10^3$ data points). No such transitions were observed in the all–atom simulations used for equilibration, where the δ / ξ space is explored to a much lower extent (see supplementary material for details). The presence of relevant correlations in OmpT is investigated by computing the projection of the top 10 eigenvectors of the covariance matrix on the trajectory ($\mathbf{X} \cdot \mathbf{V}_i$) (31). The correlation can be assessed in a quantitative way by analyzing the scatter plot in which the projection at each time step is plotted *versus* the δ and ξ distances. We find a significant correlation only with the second largest eigenvector ($\mathbf{V}_2$, Figure 3a'b'). In fact, two mostly populated regions corresponding to δ_1 and δ_2 and to ξ_1 and ξ_2, are visible.

We conclude that $\mathbf{V}_2$ induces relevant variations in the relative distances of the active site, showing high correlations with both distances. In contrast, no correlations are found with the first eigenvector $\mathbf{V}_1$, which represents the first most dominant large–scale motion. Because the cosine content associated to $\mathbf{V}_1$ is relatively larger ($C_1 \sim 0.5$, while $C_2 \sim 0.1$; see Figure 1SId), $\mathbf{V}_1$ might not represent a real "coherent" motion, but, rather, a random diffusion of less structured parts of the protein.

The large–scale motion described by $\mathbf{V}_2$ mostly affects the solvent–exposed loops embracing the active site, as found in previous work (22).

In addition, the loops L2, L3, L4 and L5 as well as the substrate are more mobile in system **B** with respect to **A** (Figure 4SIa). This is consistent with the fact that the substrate is significantly anticorrelated to the motion of the loops in system **B** (Figure 4SIb).

Thus, large–scale fluctuations in **B** allow a well defined motion of the substrate within the catalytic cleft concerted with that of loops, similarly to what is found for several proteases (4), as HIV-1 PR (2) and BACE (3).

Our approach also allows to investigate the polarization of the reactants relative to a refence system (here the Gly-Gly dipeptide) in water. To this aim, we monitor the electric field along the substrate carbonyl group C@R2 of the substrate and the catalytic water C_2 axis (Figure 1de), which has been proposed to be the nucleophile agent for the protein.

As for the substrate, the field of **A** does not show any preferential direction with respect to the C=O bond: the angle α between the field and the vector identified by the C=O bond (Figure 1d) is spread within 0° and 180°: the standard deviation value is large around its average value ($\langle\alpha\rangle = 125°(44°)$), and not much smaller than that of the reference system here investigated, for which the field is completely isotropic ($\langle\alpha\rangle = 100°(59°)$). In contrast, in **B** the field is partially aligned with the C=O bond (Figure 1d) with far lower spread ($\langle\alpha\rangle = 140°\ (10°)$). We conclude that the carbonyl carbon is expected to be more electrophilic and thus reactive than in **A** and in aqueous solution.

Similar conclusions can be drawn for the catalytic water: the field in **A** is orthogonal to the water C_2 axis (Figure 1e, the average angle β between the field and the C_2 axis of the catalytic water is $\sim 104°(14°)$); hence it

is not expected to significantly affect the nucleophilic power of the water molecule. Notice that in this case the spread is much smaller than that of the reference system ($\langle\beta\rangle \sim 93°(60°)$). In **B**, in contrast, the angle β shows a bimodal distribution, either orthogonal ($\langle\beta\rangle \sim 90°(12°)$), or aligned to the C_2 axis ($\langle\beta\rangle \sim 20°(8°)$). Thus, in the latter case the field in **B** might render the catalytic water more nucleophilic than in **A** and in aqueous solution.

The fluctuation of the β angle between β_1 and β_2 values in **B** is driven by $\mathbf{V}_2$ (see Figure 6SI). This means that the functional oscillations of the substrate allow the catalytic water molecule to take a well defined geometric configuration, able to stabilize a more negative charge distribution on the oxygen and resulting in an enhanced nucleophilicity of the oxygen atom itself.

We conclude that the polarization effect may be used by the enzyme in order to enhance the electrophilicity of the carbonyl carbon and the nucleophilicity of the water oxygen. The same calculation in the time scale typical of all–atom MD simulations ($\sim$ 50 ns) shows no preferential directions of the field in both systems (the α and β distributions are spread between 0° and 180°). This is somewhat confirmed in the all–atom equilibration simulations where no significant spread in the α and β distributions is reached (see supplementary material Figure 11SI) due to insufficient sampling.

MM/CG of systems C and D

The Michaelis complex in **C** and **D** disrupts already after 10 ns (Figure 2SIab): The NH_3^+ terminal groups, which formed salt bridges with D210 and E27 at the beginning of the dynamics, rotate and form a salt bridge with

D83 and an H–bond with Nϵ@H212 (Figure 2SIab). Consequently, the carbonyl carbon, which undergoes the nucleophilic attack, moves away from the putative catalytic dyad composed by H212 and D83 (see Figure 2SIab), whilst the oxygen of such a group flips, pointing toward the D210 and H212 dyad. Such a non–productive Michaelis complex, in which there is no putative nucleophilic agent in the close proximity of the carbonyl carbon, is maintained for further $\sim$0.04μs, after which we decided to stop the simulation. We conclude that the presence of the charged termini in addition to the two positive arginines affects the structure of the Michaelis complexes in both **C** and **D**.

MM/CG of H212A and S99A mutants

Here we compare MM/CG simulations of H212A and S99A OmpT, which show a residual activity ranging within 0% and 4% (18), with respect to wt. We ran the simulations on the mutants for protomers **A** and **B** for a shorter timescale than that of the wt (0.15 μs), as we are solely interested in constructing structural models.

In H212A, the mutation disrupts the H–bonds with Nϵ@H212 and the amide group of the substrate, which contributes to maintain the position of the substrate inside the enzymatic cleft. In the wt, R2 and R3 of the substrate form salt bridges with E27 and D208 and with D85 and D97, respectively. These two residues rotate in both **A** and **B** and their side chains face the solvent (Table 1 and Figure 7SIab). At the end of the simulation, the carbonyl carbon, which is cleaved in the wt enzyme, moves away from the putative catalytic dyad H212–D83 and the cleft is filled by

water.

In S99A, the replacement of S99 with alanine disrupts the S99–Hϵ@His101 H–bond (Figure 8SI). The H atom of the same histidine H–bonds to D83 in **A** and **B**.

In **A**, the breaking of the S99–H101 H–bond causes a rearrangement of H101 and consequently D83, which is protonated, rearranges and H–bonds to D97 (Table 2 and Figure 8SIa). Consequently, the water–mediated interaction between the proton of D83 and the substrate is lost, and the latter fluctuates allowing a rotation of the side chain of R2 of the substrate. As a result, R2 H–bonds with D210 and the catalytic residue H212 is permitted to move further apart from the active site (Table 2). This causes a drastic change of the substrate configuration and its partial detachment (Figure 8SIa).

In **B**, D83, which is ionized, forms a stable H–bond with Nϵ@H101. However, because of the lack of S99–D97 H-bonding, D97 moves away, in turn causing the loss of H–bond interactions between the backbone of D85 and D97 (Figure 8SIb). As a result, the salt bridge between D85, D97 and R3 is lost, allowing the side chain of the latter residue to rotate and become solvated. In this case, a partial detachment of the substrate occurs (Table 2, Figure 8SIb).

Discussion

MM/CG simulations have been used to investigate the fluctuations of OmpT in complex with its substrate ARRA on the μs time scale as well as the effect

of key point mutations at the active site. We have focused on four systems (**A**–**D**), which differ for the protonation state of D83 and the charge of the substrate (AceARRANMe or NH_3^+ARRA COO^-), which have first been simulated by all–atom MD simulations on a shorter time scale.

The complexes with substrates with charged tails (**C** and **D**) evolved to non–productive Michaelis complexes, because the charged tail groups interact with the residues among the catalytic cleft causing a distortion of the ARRA peptide and its detachment from the binding pocket. Interestingly the preliminary 24 ns all-atom simulations did not show a disruption of the Michaelis complex, which points to a less efficient exploration of conformational space compared to the MM/CG approach, where the systems are described by smoother potential energy surfaces. Obviously, our investigation does not rule out the presence of longer substrates with charged tail groups, as those substrates may have their termini located outside the catalytic cleft.

In contrast, the complexes with the substrates with neutral tails provided productive Michaelis complexes and both remained stable over 1–μs dynamics simulations. The two complexes (**A** and **B** in Figure 1bc) exhibit different protonation states of D83: in **A**, the aspartic acid is protonated while in **B** it is ionized. The systems are characterized by significant differences in the electrostatic polarization of the reactants. In **B**, the active site polarizes both the catalytic water and the carbonyl carbon, rendering the first more nucleophilic and the second more electrophilic relative to a reference system in water (diglycine). In contrast, in **A** no significant polarization is observed and the electric field acting on the carbonyl carbon

bond and on the catalytic water is very similar to that calculated for the reference system diglycine in water solution.

System **B** is also characterized by large conformational fluctuations of the substrate triggered by global large–scale motions (Figure 3a'b'), which populate significantly different conformations (see Figure 3ab and Figure 3SI). This contrasts to **A** in which the substrate fluctuates around a well defined conformation (see Figure 3SIa and Figure 5SIab). The large–scale fluctuations of complex **B** might have functional implications as found in the Aspartyl Proteases superfamily (2–4). However, quantum chemical calculations on the reaction mechanism are required to discriminate between the two protomers and further address this issue.

We have next used the MM/CG approach on H212A and S99A mutants (for both **A** and **B** protomers) to provide the structural basis for the much lower activity of these mutants relative to wt. This result, which is easy to rationalize for the H212A mutant, as it is part of the putative catalytic dyad, is rather intriguing for S99A, which is not involved directly in the catalytic cleft. Our simulations suggest that in H212A the ARRA peptide detached spontaneously in both systems due to the loss of the H–bond interaction between Nϵ@H212 and the amide group of the substrate, present in the wt. S99A, indeed, allowed a similar detachment of the peptide due to the disruption of the geometry of the active site in both systems, in spite of the fact that this residue is not located at the active site. In **A** the loss of the H–bond interaction between S99 and D83 allows the rotation of the latter residue causing the breaking of interactions between D83 and the substrate. In contrast, in **B** the detachment of the peptide is due to the loss

of the salt bridge between R3 and D85–D97 after the breaking of the S99–D97 interaction. We conclude that not only first–shell H–bond interactions (such as those formed by H212), but also second shell H–bonding (such as that of S99) play an important role for the stability of the geometry of the active site. Removing any of those stabilizing interactions may cause a high instability of the active site and therefore a reduced activity.

In conclusion, our MM/CG approach emerges as a useful tool to investigate μs simulations of enzymes, which is presently difficult with all–atom MD. Our conclusions about the electric field acting on the reactants and about the motion of the substrate inside the catalytic cleft cannot be drawn if one focuses on the typical time scale of all-atom MD (0.01–0.1 μs). This suggests that this approach, within its limitations deriving from the use of a coarse–grained model for modeling the most of the protein in solution, may provide useful information –complementary to all atom MD– on phenomena occuring on relatively long time scales. In addition, it may be useful for computational molecular biology, allowing to test the effect of point mutations via a computationnally affordable method. Thus, MM/CG, by allowing to run more numerous and longer simulations than all–atom MD, is expected, on the one hand to improve our confidence in the results, and on the other one it may strengthen the interaction between molecular biology experiments and simulations.

Acknowledgments

The authors thank Alessandra Magistrato and Michele Cascella for discussions. We would like to acknowledge financial support by MURST–FIRB and IIT.

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Figure Legends

Figure 1.

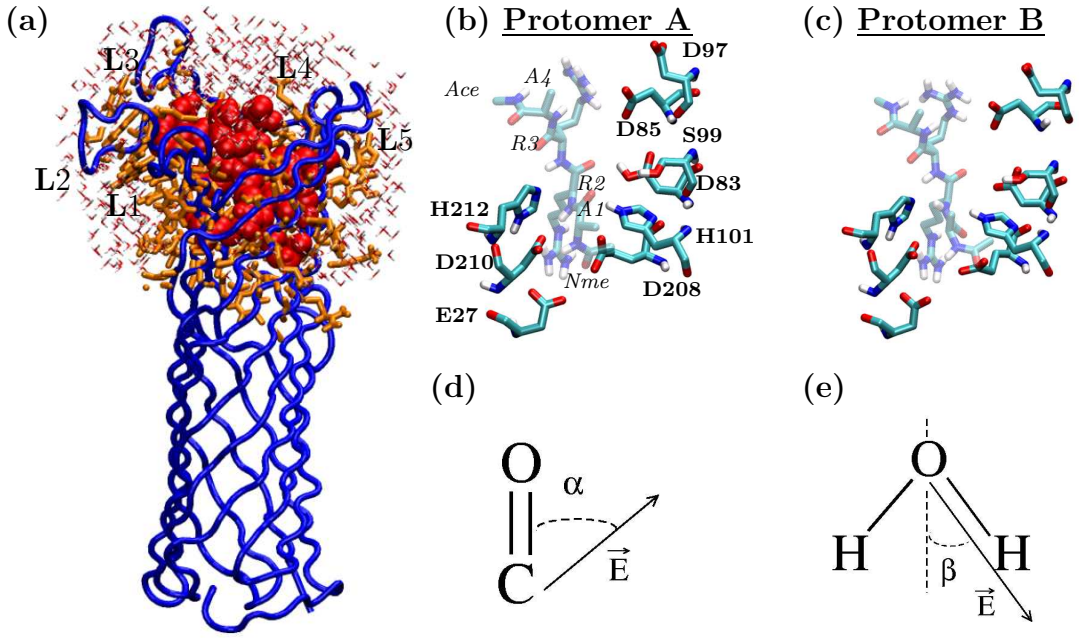
OmpT in complex with the model substrate ARRA. (a): atoms treated with MD potential are depicted as red van der Waals spheres. Atoms belonging to the interface region (also treated with MD potential) are represented as orange licorice and the CG region is depicted in blue tube representation. A shell of water centered around the MM region is also shown. (b) and (c): geometry of the active site of **A** and **B**, respectively, after 24 ns of MM equilibration simulations. The residues in licorice representation constitute the MM region. The substrate is depicted with a transparent effect. (d) and (e) represent the electrostatic field acting on the carbonyl carbon and on the dipole axes C_2 of the catalytic water, respectively.

Figure 2.

Significant snapshots of the active site of **A**: (a), and **B**: (b). The catalytic water is depicted in van der Waals representation with a transparent effect.

Figure 3.

Left of the panel: cartoon representation of the OmpT/ARRA complex. Dashed lines depict δ and ξ distances defined in the text. Right of the panel: time evolution in **B** of δ, (a), and ξ, (b), parameters (defined in the text), respectively. Right: plot of $\mathbf{X} \cdot \mathbf{V_2}$ *versus* δ, (a'), and ξ, (b'), distances for **B**.


(a)
L3
L4
L5
L2
L1
(b) Protomer A
D97
Ace
A4
D85
S99
R3
D83
R2
H212
A1
H101
D210
Nme
D208
E27
(c) Protomer B
(d)
O
α
C
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(e)
O
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H
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Figure 1:

(a) **(b)**

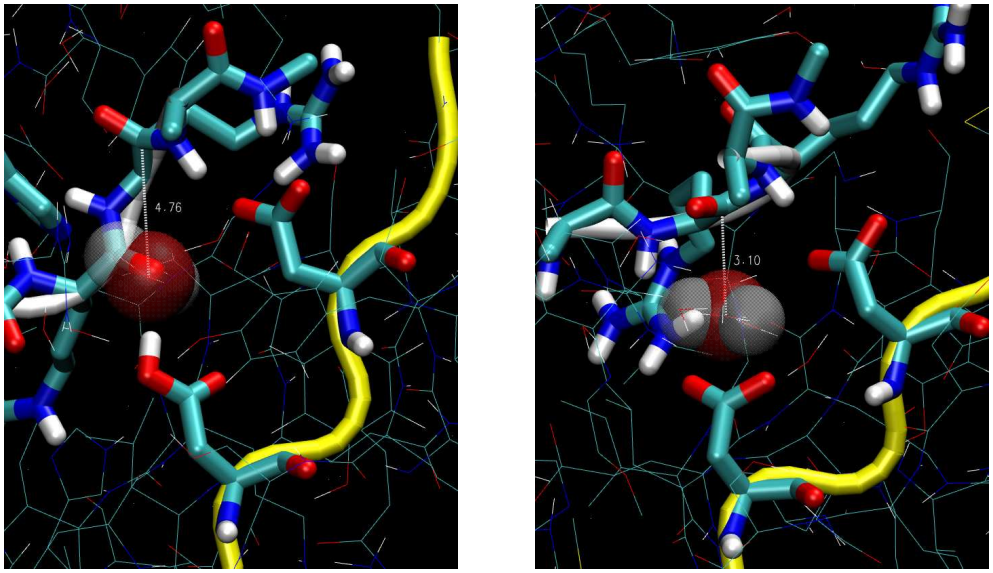

Figure 2:

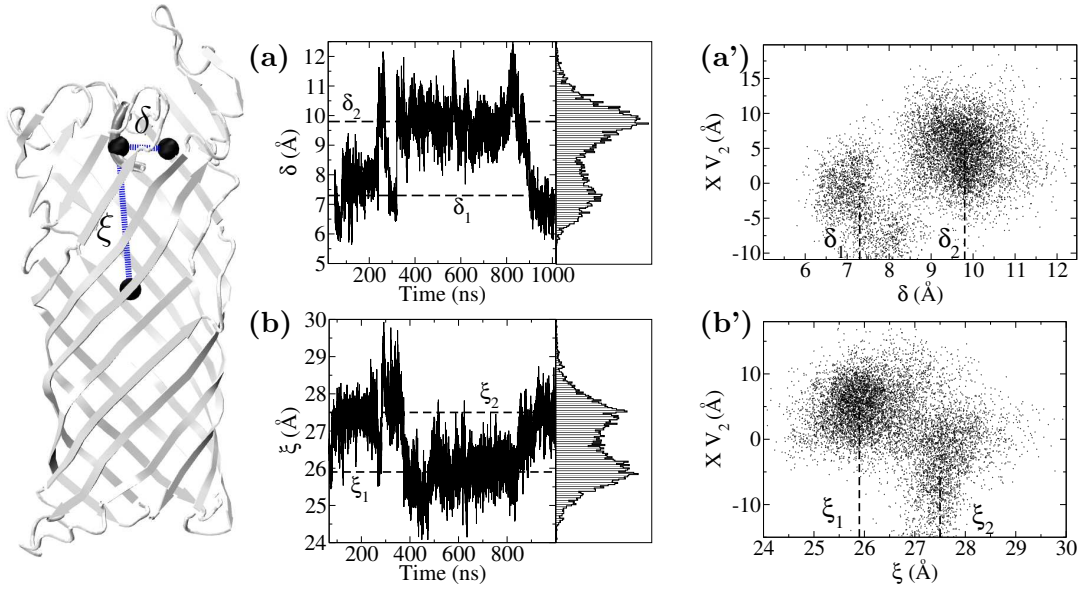
δ
ξ
(a)
δ (Å)
Time (ns)
δ2
δ1
(a')
X V2 (Å)
δ (Å)
δ1
δ2
(b)
ξ (Å)
Time (ns)
ξ2
ξ1
(b')
X V2 (Å)
ξ (Å)
ξ1
ξ2

Figure 3:

Supplementary Information (SI)

Microseconds Dynamics Simulation of the Outer–Membrane Proteases T

Marilisa Neri, Marc Baaden, Vincenzo Carnevale, Claudio Anselmi, Amos Maritan and Paolo Carloni

21 June 2007

Additional information on MM/CG simulations of systems A–B

Wild–type OmpT

The root mean square deviations (RMSD) of the C_α of **A** and **B** is stable after $0.05\mu s$ and fluctuates around an average value of $\sim$3Å, as shown in figure 1SIab. We evaluate the convergence of the simulations after 50ns by the analysis of the cosine content C_i of the principal component p_i, defined as:

$$C_i = \frac{2}{T}\left(\int_{T_0}^{T}\cos(4\pi t)p_i(t)dt\right)^2\left(\int_{T_0}^{T}p_i^2(t)dt\right)^{-1}, \qquad (1SI)$$

where T is the time of simulation and p_i is the principal component associated to the ith eigenvector. C_i can take values between 0, no cosine, and 1, a perfect cosine. It was demonstrated that insufficient sampling can lead to behaviors that resemble a functional motion, but describing a random motion [Hess(2000), Hess(2002)]. In Figure 1SIcd are reported the C_i values associated to the the first 8 eigenvectors of **A** and **B** computed after 50ns. The picture shows that the cosine content is lower than 0.5 for both the systems.

The RMSD values for the C_α atoms of the substrate relative to the MM/CG snapshot at 0.05 μs of the wt as a reference structure (Figure 3SIab), are either characterized by small fluctuations ($\sim$ 0.4 Å) around an average value of $\sim$ 2.3Å (system **A** , see Figure 3SIa) or they show many peaks featuring fluctuations of larger amplitude of $\sim$ 1.7Å (system **B** , see Figure 3SIb). In addition, based on a calculation of the root mean square fluctuations (RMSF) of each residue, we conclude that loops L2, L3 and L4 as well as the substrate in **B** are more mobile than in **A** (Figure 4SIa). This is consistent with the analysis of the large scale motions of the systems, which suggest that the motions of the substrate in the catalytic cleft are concerted with that of the loops, the correlation being larger for **B** (Figure 4SIb).

The δ and ξ distances, defined in the Main Text, are used in this work to characterize the motion of the substrate inside the catalytic cleft. In **A**, as shown in Figure 5SIab, δ and ξ oscillate around an average value ($\langle\delta\rangle \sim 7.8(0.5)$ Å and $\langle\xi\rangle \sim 24.1(0.4)$ Å) with a sharp Gaussian–like distribution. In contrast, in **B** the distances feature bimodal distributions (see Main Text).

H212A and S99A OmpT

For both systems, in H212A, the side chains of R2 and R3 of the substrate, which form a salt bridge with E27–D208 and D85–D97, respectively, rotate and face the solvent at the end of the simulation (0.05 μs, Figure 7SIab).

In S99A-system **A**, R2 side chains interacts at 0.05 μs with D210 breaking the D210–H212 H–bond. This causes a motion of H212 away from the catalytic site (see Figure 8SIa). In S99A-system **B**, the salt bridge between R3 and D85–D97 is disrupted after 0.02 μs. This causes, as for H212A and S99A of **A** mutants, a motion of the substrate carbonyl carbon away from the catalytic dyad D83–H212 (see Figure 8SIb).

Additional information on MM equilibration simulations of systems A, B, C and D

The OmpT–ARRA complex described in [Baaden and Sansom(2004)] was optimized using a molecular mechanics based approach with the Yasara software and the Yamber2 forcefield [Krieger et al.(2004)]. A limited region (see Figure 9SI) around the active site was free to move in order to speed up the calculations and be able to compare the four systems **A**, **B**, **C** and **D**. First a protocol was established where the complex is gradually solvated. The box is filled with water at a density of 0.997 g/l, then a steepest descent minimization is run followed by simulated annealing and short molecular dynamics at 298 K. After solvating the box, the density is adapted using an NPT algorithm within Yasara. All equilibrations are a combination of steepest descents followed by 500 steps of simulated annealing and 0.6 to 6 ns of molecular dynamics. During these steps the protein, peptide and catalytic water were first kept fixed, then parts were released, then instead of keeping them fixed some springs were used and finally all restraints were removed. Each of the systems **Aeq**, **Beq**, **Ceq** and **Deq** underwent 24 ns of unrestrained equilibration.

Characteristic structural parameters were analyzed for all four simulations (Figure 10SI). The simulations with neutral termini are clearly separated from the ones with ionized residues with respect to the δ and ξ parameters as shown in Figure 10SIabd. The evolution of the catalytic water–scissile carbon distance (Figure 10SIc) indicates that simulation **Beq** shows the shortest distances and might thus visit reactive transition states more often than the other protomers. Only a small fraction of the conformational δ–ξ space is explored compared to the MM/CG simulations.

The polarization of the active site is characterised by the α and β angles defined in the main text. Figure 11SIab shows their fluctuations and time evolution. Average values for α range from 134 to 151° with fluctuations of $\pm 10°$. Average values for β are spread between 158 and 165° with equal fluctuations of $\pm 10°$. Occasional large peaks are observed like for β in simulation **Aeq**, which hints at insufficient sampling on the short 24 ns simulation timescale.

Additional Figures

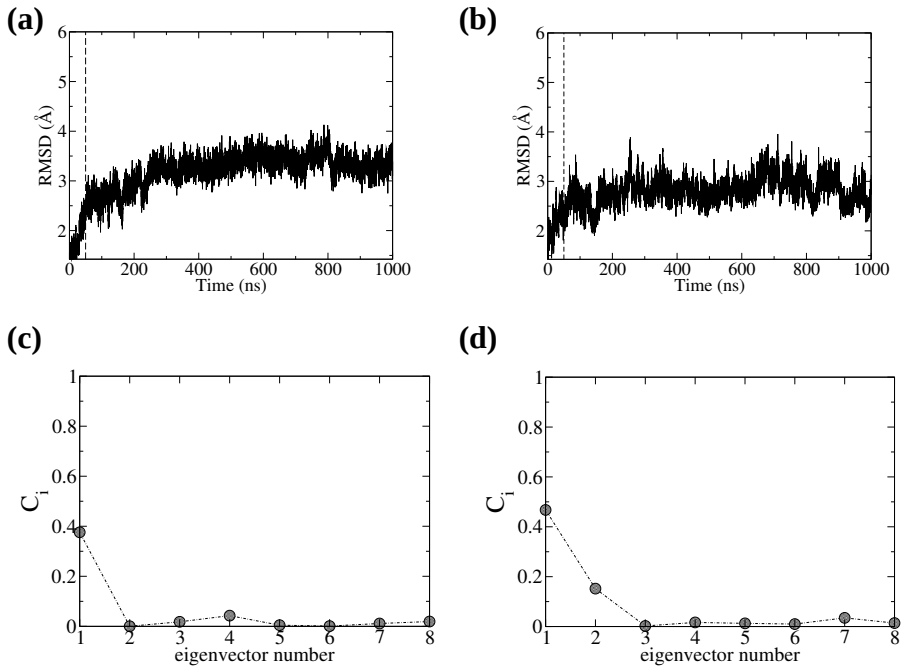

Figure 1SI: RMSD of C_α atoms plotted as a function of simulated time for **A**: (a) and **B**: (b). Pictures (c) and (d) depict the cosine content, defined in Eq. 1SI, of the first 8 eigenvectors of the covariance matrix after 50ns, for **A** and **B**, respectively.

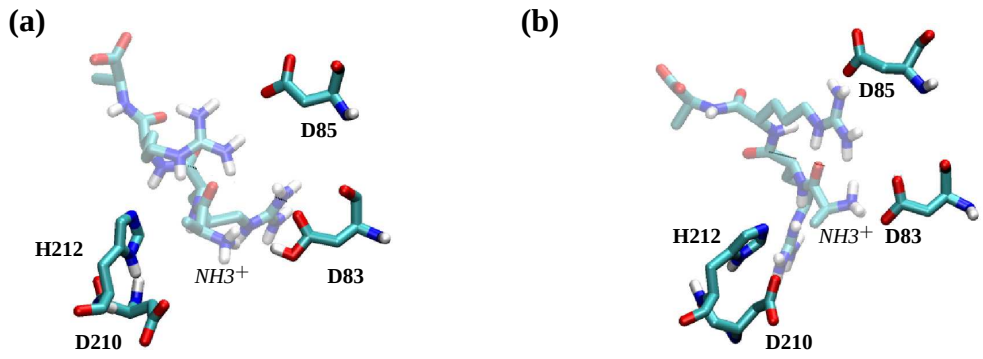

Figure 2SI: MM/CG snapshots of the active site after 10 ns for **C**: (a) and **D**: (b). The substrate is depicted with a transparent effect.

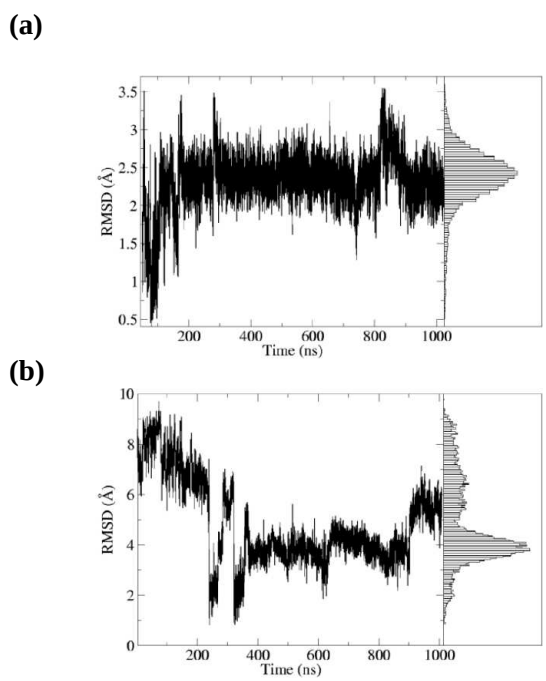

Figure 3SI: Left: RMSD of the C_α atoms of the ARRA substrate in **A**: (a), and **B**: (b) plotted as a function of time. Right: correspondent distribution of RMSD values.

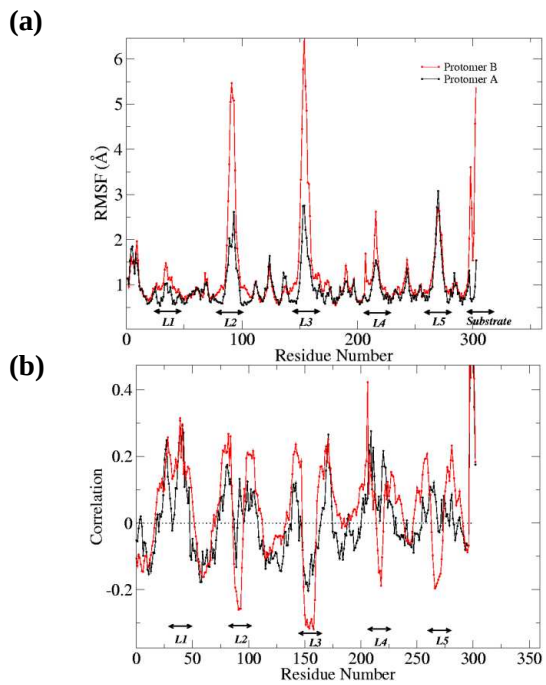

Figure 4SI: (a) Root mean square fluctuations of **A** and **B** (black and red colors, respectively). (b) Degree of correlation between the instantaneous displacement of the C_α atom of residue R3 and that of the other C_α's. This degree of correlation is calculated from the normalized reduced covariance matrix C. The covariance matrix is calculated from the instantaneous displacement vector $\delta\vec{x}_i$ (after the optimal roto–translation) of the ith C_α atom from the reference (time-averaged) position. The matrix element is defined as $C_{ij,\alpha\beta} = \langle \delta x_i^\alpha \, \delta x_j^\beta \rangle$ where α, β denote the Cartesian components and the brackets indicate the time average. Finally the reduced normalized covariance matrix is calculated as: $C = \frac{C_{ij,\nu\nu}}{\sqrt{C_{ii,\nu\nu} C_{jj,\mu\mu}}}$.

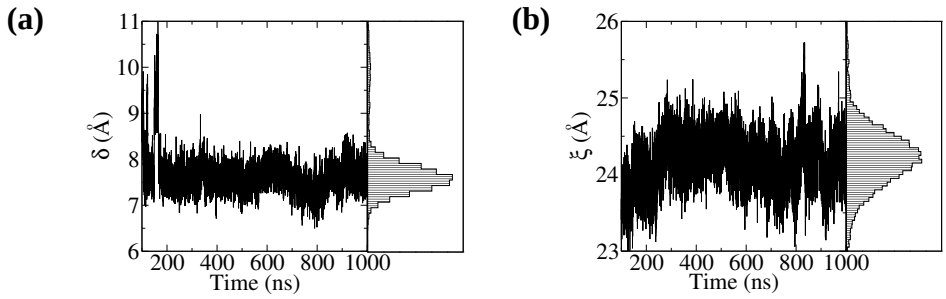

Figure 5SI: Structural parameters plotted as function of simulated time of **A**. Time evolution of the distance between the center of mass of the C_α atoms of the substrate and the center of mass of the dyad D83–D85: (a) and the center of mass of the β-Barrel: (b).

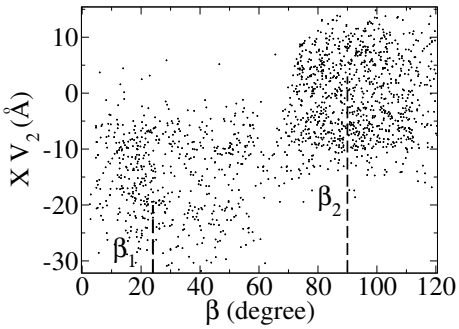

Figure 6SI: Projection of the trajectory **X** of **B** on the eigenvector $\mathbf{V}_2$ of the covariance matrix plotted *versus* the time evolution of the angle β defined in the main text. The calculation was carried out every 100 ps.

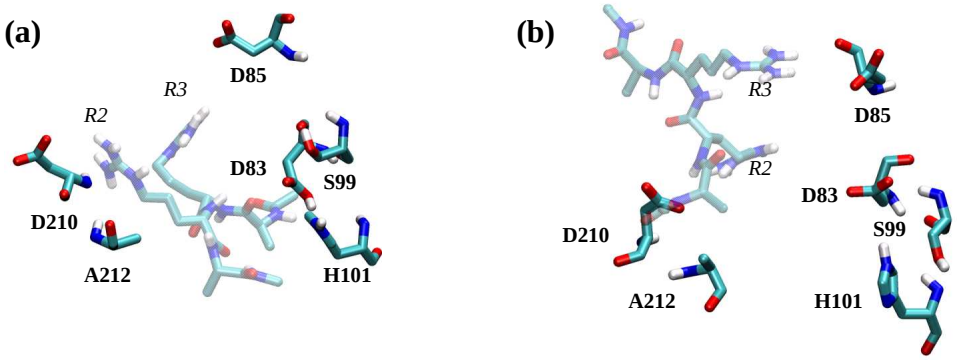

Figure 7SI: H212A OmpT MM/CG simulation: Active site structure after 150 ns for **A**: (a) and **B**: (b). The substrate is depicted with a transparent effect.

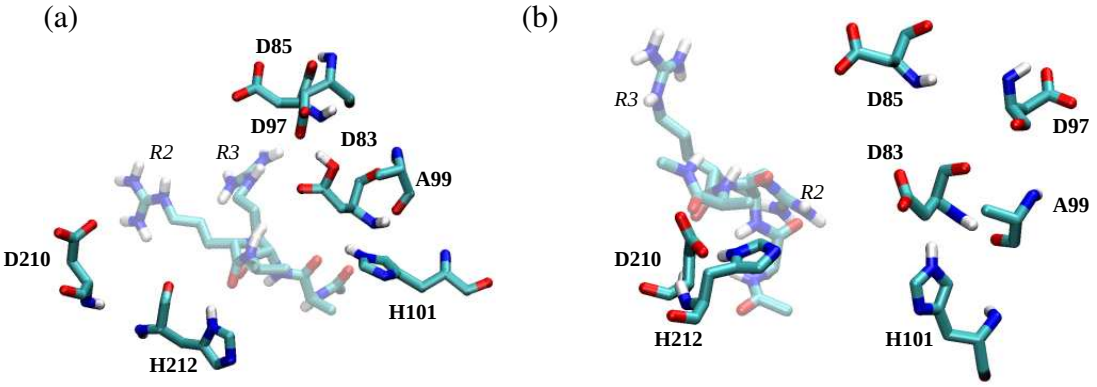

Figure 8SI: S99A OmpT MM/CG simulation: Active site structure after 150 ns of the S99A mutant for **A**: (a) and **B**: (b).

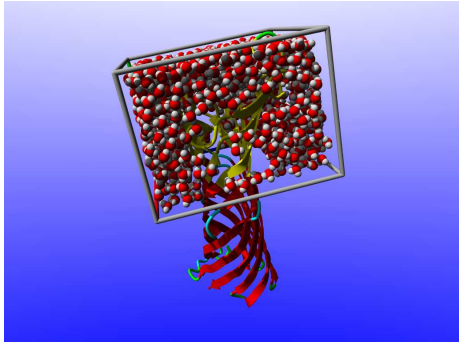

Figure 9SI: OmpT MM simulation: simulation system highlighting the box around the active site that is free to move. It is hydrated using a cell with the size 40 x 30 x 40 Å.

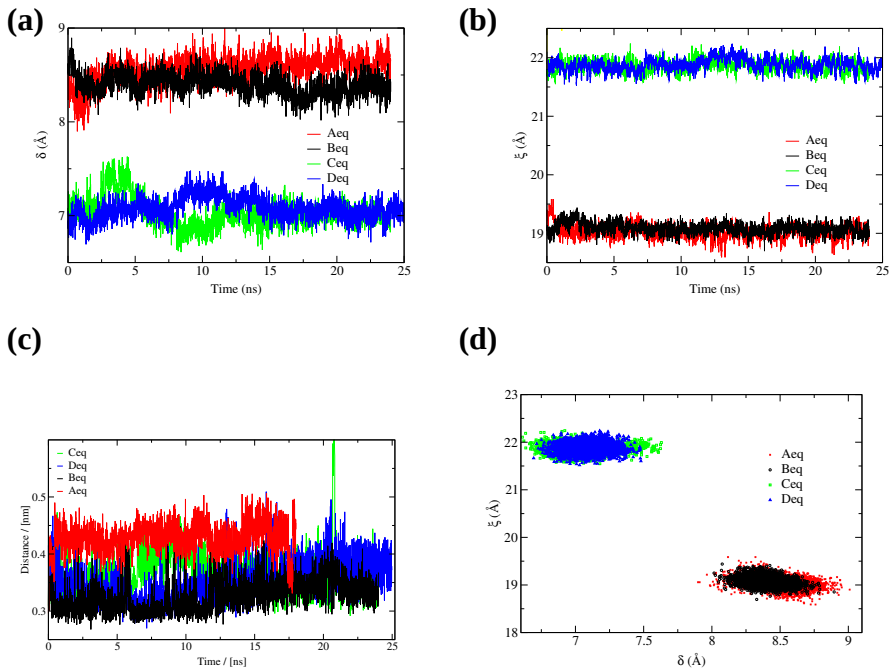

Figure 10SI: OmpT MM simulations **Aeq**, **Beq**, **Ceq** and **Deq**: (a-c) Structural parameters plotted as function of simulated time. Time evolution of the distance between the center of mass of the C_α atoms of the substrate and the center of mass of the dyad D83–D85: (a) and the center of mass of the β-Barrel: (b). Time evolution of the distance between the catalytic water molecule and the scissile C: (c). ξ versus δ plot illustrating the exploration of conformational space: (d).

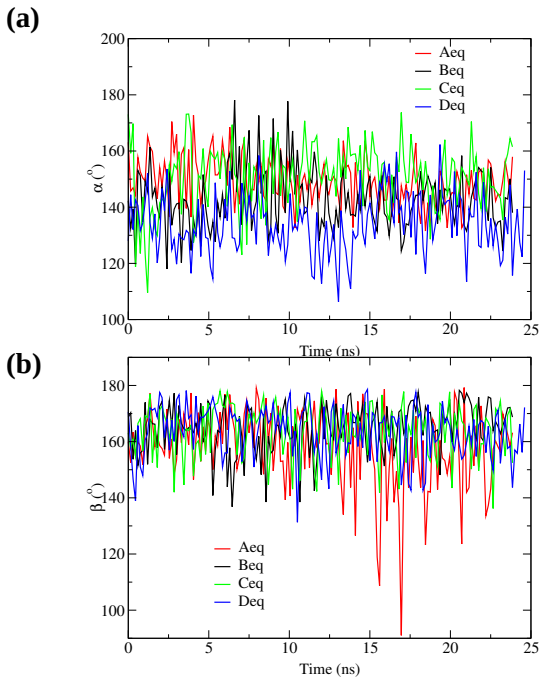

Figure 11SI: Time evolution of the electric field angles with catalytically important groups for OmpT MM simulations **Aeq**, **Beq**, **Ceq** and **Deq**. (a) α angle between the electric field and the carbonyl bond. (b) β angle between the electric field and the water C2 axis.

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