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# Artificially smart optimization of crash cushion device sustained by experimental and numerical study of re-entrant auxetic honeycomb

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## ABSTRACT

Honeycomb structures are widely used in energy absorption, and more recently auxetic honeycombs have been studied in order to improve absorption capabilities of such structures. The hexagonal re-entrant (HR) honeycomb is foreseen to be a promising structure under impact velocities. An experimental analysis of the re-entrant honeycomb under impact velocity has led to a finite elements model validation at scale one, i.e. scale of current car crash cushions. A new objective function based on the European Standard has been developed in order to improve crash cushions capabilities while avoiding peak deceleration by using a meta-heuristic optimization algorithm. The global optimization process has been performed using Inverse-PageRank-PSO algorithm. The algorithm has led to an optimal geometrical configuration of HR honeycomb improving the performance of current road safety devices. The optimal structure presents a quasi-linear absorption curve, as recommended by European standards.

## ARTICLE HISTORY

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## KEYWORDS

Optimization; energy absorption; hexagonal re-entrant honeycomb; road safety

## 1. Introduction

Honeycomb structures are widely used, solicited in-plane, in energy absorption [1–4] and constitute crash cushion systems used along roads. The aim of these devices is to efficiently absorb kinetic energy, while preserving the vehicle occupants integrity. European and French standards [5,6] give the impact conditions and the absorption capabilities of the structure to be validated and installed along roads. To do so, optimization methodologies can be used, to develop the most adapted structure, while considering the manufacturing capabilities of the road safety devices industry. Usually, simple folded metal sheets are assembled [7]. In this paper, a parametric optimization algorithm is used, to find the best geometric parameters of the structure, to improve the structural absorption capabilities, while meeting the European standards and guaranteeing manufacturing constraints.

Cellular materials, such as honeycombs, have been used as new engineered materials because of their very efficient energy absorption capabilities [3,8–10]. In the last few years, many works have been conducted on auxetic materials [11–13], which exhibit a negative Poisson's ratio, since they have been introduced by Lakes et al. [14]. Usually manufactured as foams [15,16], they can also be manufactured at bigger scales, and are then called cellular materials [17]. However, if honeycomb structures can not be considered as a Representative Volume Element, because the scale ratio conditions [18] are not met, the structure has to be entirely considered and no homogenization method is needed [19].

In this paper, the total structure will be studied and optimized. Numerous shapes of auxetic microstructures have been proposed in the literature, such as star honeycomb, arrow-head honeycomb, chiral honeycomb and hexagonal re-entrant honeycomb [13]. Optimization methodologies are sometimes used to design new microstructures exhibiting high negative Poisson's ratio effects [20,21].

The hexagonal re-entrant (HR) honeycomb has shown high efficient capabilities in energy absorption under impact compared to conventional honeycomb [8,10,19,22]. Under quasi-static conditions, many studies have been conducted to predict either the elastic constants [23], or the dependence of hexagonal re-entrant honeycombs' Poisson's ratio and Young's modulus on the cell geometric parameters, using Finite Elements Method (FEM) analysis [17,22]. A theoretical approach to predict negative Poisson's ratios of auxetic honeycombs has been developed, which is based on the large deflection model [2]. Experimental and numerical studies have been performed on the hexagonal re-entrant honeycomb under quasi-static loading to study its crushing behavior. They have investigated the negative Poisson ratio effects on the crushing stress and energy absorption efficiency [19]. If the elastic and quasi-static behaviors of hexagonal re-entrant honeycombs have been largely studied, the dynamic behavior under impact velocities is less studied. Hou et al. [8] and Hu et al. [9] have worked on the dynamic response of re-entrant honeycombs under high velocity crush [8,9]. In these works, the energy absorption efficiency of the structure is computed by using the propagation waves theory. The studied velocities are over

118 100 *m/s* and thus, do not correspond to car limitation  
 119 speeds which are limited to 30 *m/s* on highways in France  
 120 [5]. In fact, European standards impose a certain structural  
 121 volume and geometry adapted to the impact velocity, which  
 122 will then be studied in this paper.

123 In the literature, some studies have been conducted on  
 124 the optimization of HR honeycomb structures subjected to  
 125 crash. Wang et al. [24] have optimized a crash-box with 3 D  
 126 re-entrant auxetic core under 15 and 40 *km/h* crash condi-  
 127 tions. However, the studied volume is limited to the  
 128 imposed crash-box size, i.e. (130 × 133 × 73 *mm*) and does  
 129 not correspond to crash cushion device volumes  
 130 (1200 × 800 × 700 *mm* under 50 *km/h*) usually used by car  
 131 cushion manufacturers [7]. On another hand, Qi et al. [25]  
 132 have achieved a multi-scale optimization of the hexagonal  
 133 re-entrant honeycomb under blast impact. In order to  
 134 improve honeycomb capabilities, a multi-objective function  
 135 have been designed, to investigate the optimal thickness  
 136 value of a cell-wall material as an independent variable of  
 137 the problem. The obtained HR honeycomb structure exhib-  
 138 its an improvement about 5% of the energy absorption com-  
 139 pared to conventional honeycomb.

140 Most of the road networks are lined with many advanced  
 141 road safety systems. These devices absorb the kinetic energy  
 142 of crashing cars mostly by plastic deformation. The volume  
 143 of these safety systems varies as a function of the limitation  
 144 speed. Currently, these types of cushions are made of con-  
 145 ventional honeycomb structures, which absorb kinetic  
 146 energy essentially by in-plane plastic deformation [7,26]. In  
 147 this paper, the hexagonal re-entrant honeycomb deforma-  
 148 tion under car crash impact velocities is investigated to  
 149 improve the performance of current cushions. The major  
 150 criterion in car cushion design is about the deceleration felt  
 151 by the crashed-car-user and is called the Acceleration  
 152 Sensitivity Index (*ASI*) [5]. The European standard estimates  
 153 at  $9.8 \times g$  [*m/s*<sup>-2</sup>] the deceleration that can be sustained  
 154 during 50 *ms* by a human being before severe injuries [5],  
 155 where *g* is the gravity acceleration. Indeed, in crashworthi-  
 156 ness optimization, many works have investigated the  
 157 Specific Energy Absorption (*SEA*) based on the stress-strain  
 158 curve optimization [25,27–29]. *SEA* optimization consists of  
 159 maximizing the integral of the impact force, over the dis-  
 160 placement under quasi-static loading while minimizing the  
 161 mass of the structure. Here, according to the *ASI* criterion,  
 162 the striker limit force is defined, allowing to characterize the  
 163 capability of a cushion to absorb kinetic energy. In this  
 164 paper, this criterion is used to define a new appropriate  
 165 objective function based on the force-time curve, to  
 166 be optimized.

167 The remaining part of this paper is organized as follows.  
 168 Section 2 presents the experimental study conducted on the  
 169 HR honeycomb under impact conditions. The experimental  
 170 work allows to validate the FE model presented in Section  
 171 3. The new objective function is defined in Section 4.1,  
 172 which depends on the design variables defined in Section  
 173 4.2. Then, an optimization methodology is performed and  
 174 presented in Section 5. Finally, the discussion and conclu-  
 175 sion of this paper are outlined in Section 6.

## 2. Experimental analysis of the hexagonal re-entrant honeycomb absorbing cushion device

### 2.1. Experimental setup

177 In order to study the mechanical behavior of HR honey-  
 178 combs under impact conditions, an experimental campaign  
 179 has been conducted at the LBMC-crash-track [30]. The  
 180 device is composed of a track propelling a vehicle into a  
 181 car-cushion-specimen fixed on a concrete wall, as shown in  
 182 Figure 1. The tested vehicle is a 4-wheel-rigid vehicle with a  
 183 mass *m* of 690 *kg*. The vehicle is freed from the propeller  
 184 7 *m* before the concrete wall, and thus is guided neither dur-  
 185 ing nor after the impact. The 4-wheel-rigid vehicle is pro-  
 186 pelled at 14 *m/s* corresponding to the speed limitation of  
 187 50 *km/h* defined in the European standards [5,6].

188 The tested vehicle is instrumented with front and rear  
 189 accelerometers recording the acceleration of the vehicle in  
 190 *X*, *Y* and *Z* directions, defined in Figure 1. Acceleration sig-  
 191 nals are filtered with a Butterworth filter with 4 poles at  
 192 180 Hz (BW 4p 180 Hz) in conformity with European stand-  
 193 ard ISO6487 [31]. The front sensor is able to measure an  
 194 acceleration up to 500 *g* whereas the rear sensor has a cap-  
 195 acity of 100 *g*, where *g* is the gravity acceleration. The sam-  
 196 pling interval is  $10^{-5}$  *s*<sup>-1</sup> for all the accelerometers.

197 Three high speed video cameras are installed to film dif-  
 198 ferent shots. Their corresponding resolutions are described  
 199 in Table 1. The sampling interval is set to 3000 *images/sec*  
 200 for all cameras. Cameras are connected to the trigger identi-  
 201 fying the beginning of the contact between the vehicle and  
 202 the specimen. This allows to record 500 images on the cam-  
 203 eras before the contact.

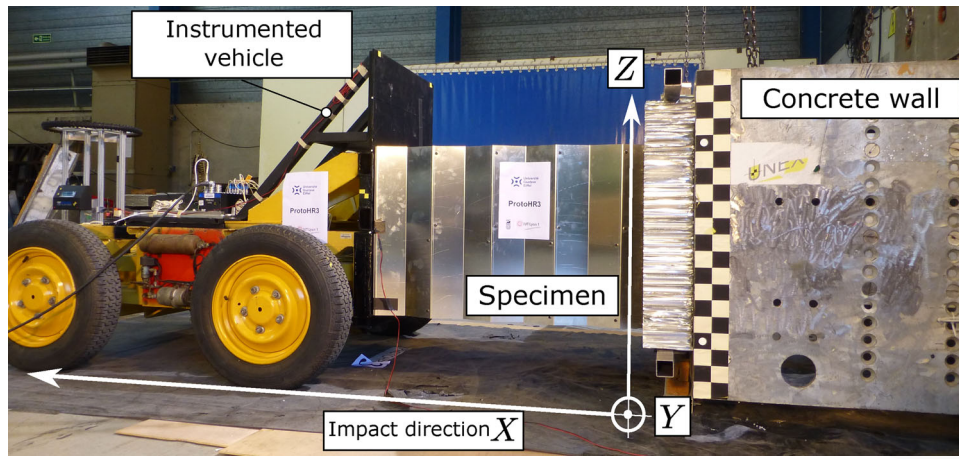
204 Three test patterns are installed on the vehicle in order  
 205 to automatically track the displacement. The Tracker 5.1.3  
 206 software is used with the automated object tracking position  
 207 functionality to fit the displacements of the test patterns.

### 2.2. Geometric configuration and manufacturing process of the absorbing cushion device

208 The unit cell configuration adopted in this study is a re-  
 209 entrant honeycomb which can be described with three geo-  
 210 metric parameters *l*, *h* and  $\theta$  as shown in Figure 2. The cell  
 211 wall thickness is set to *t*. By duplicating the unit cell in both  
 212 *X* and *Y* directions, the complete structure is constituted of  
 213 *n* and *m* cells in the *X* and *Y* directions respectively. The  
 214 total lengths of the structure are called *L<sub>X</sub>*, *L<sub>Y</sub>* and *L<sub>Z</sub>* in the  
 215 *X*, *Y* and *Z* directions respectively.

216 The specimen dimensions are selected among a set of  
 217 parameters leading to an absorbing crash-cushion design  
 218 preserving the experimental devices during impact tests, and  
 219 are described in Table 2.

220 The constitutive material used to manufacture the speci-  
 221 men is a certificated (EN 485-2) aluminum (SHEET EN  
 222 AW 5754 H111) exhibiting a linear elastic isotrope behavior.  
 223 The Young modulus *E*, the offset yield point *R<sub>p0.2</sub>*, the fail-  
 224 ure stress *R<sub>m</sub>* and the maximal elongation *A%* are given in  
 225 Table 3. The specimen is composed of two major parts. The  
 226 first one is called wavy sheet and is *L<sub>X</sub>* long and *L<sub>Z</sub>* high.



Q4 Figure 1. LBMC-crash-track experimental device and the car-crash-cushion specimen.

Table 1. Cameras setup.

Camera	Resolution (pixels)	Shot
1	1024 × 576	(X, Y) plan
2	768 × 768	(X, Y) plan
3	1024 × 576	(X, Z) plan

Table 3. Mechanical properties of aluminum AW5754H111.

	E (MPa)	$R_{p0.2}$ (MPa)	$R_m$ (MPa)	A%
AW5754H111	69 000	118	211	26.5

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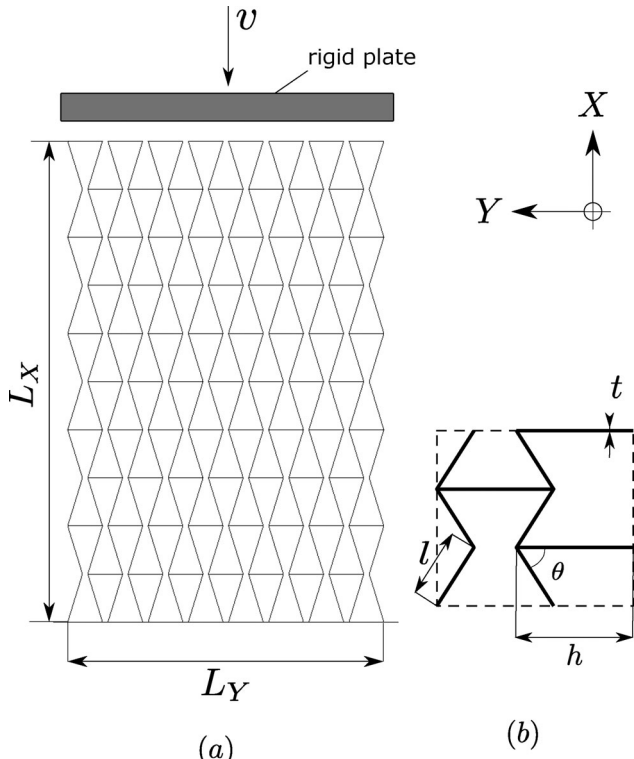


Figure 2. Hexagonal re-entrant honeycomb specimen: (a) whole structure and (b) unit cell detailed with geometric parameters.

Table 2. Geometrical parameters of HR honeycomb specimen.

$L_x$ (mm)	$L_y$ (mm)	$L_z$ (mm)	$t$ (mm)	$l$ (mm)	$h$ (mm)	$\theta$ (°)
1240	810	690	2	130	90	73

Wavy sheets are bent with a metal press using a V16mm vee. The second part is the joint (as shown in Figure 3 and illustrated by the blue line) allowing the assembling of the wavy sheets together. The whole specimen is composed of 16 wavy sheets and 67 joints. Joints and wavy sheets are

bolted together with 8 bolts per joint. Each bolt is composed of 1 hexagonal head screw M6x16 ISO4017 class 10.9, 1 hexagonal nut M6 (DIN934) class 10 and 2 plain washers normal type (DIN125-1A) 140HV made of stainless steel. Bolt are tighten at a torque of 11N.m. On the front side of the HR honeycomb specimen, 3 metal sheets are riveted to the specimen in order to avoid expansion in the Y direction of the front cells during impact (see Figure 3).

The experimental deformation of the HR honeycomb structure under  $v = 13.89m/s$  impact velocity is shown in Figure 5. Before  $t = 0.02s$  the Negative Poisson Ratio (NPR) effect is highlighted in the deformation process where the upper layers are crushed. Then, the structure wears down in the Y direction from  $t = 0.05s$  to the end of the crash at  $t = 0.15s$ . The manufacturing process allows to keep the HR honeycomb in one piece all over the impact crashing. The specimen can be considered as a one piece honeycomb without failure or wrenching. These experimental results will then be used to validate the numerical model presented in next Section.

### 3. Hexagonal re-entrant honeycomb numerical study

#### 3.1. Finite elements analysis

A numerical model of the impacted hexagonal re-entrant honeycomb crash-cushion has been computed, using LS-DYNA explicit solver. In a step of validation, a crashed car of mass  $m = 690kg$  is modelled by a rigid plate with an equivalent mass, impacting the HR structure with an initial velocity set to  $v = 13.89m/s$ , as presented in Figure 2 to match with the measured velocity in the experimental study previously presented in Section 2.

The acceleration  $\gamma_x$  in the X direction on the rigid wall is computed as the ratio of (i) the force ( $F_x$ ) in the X direction computed in the rigid wall during the crash and (ii) the mass ( $m$ ) of the wall, as follow:

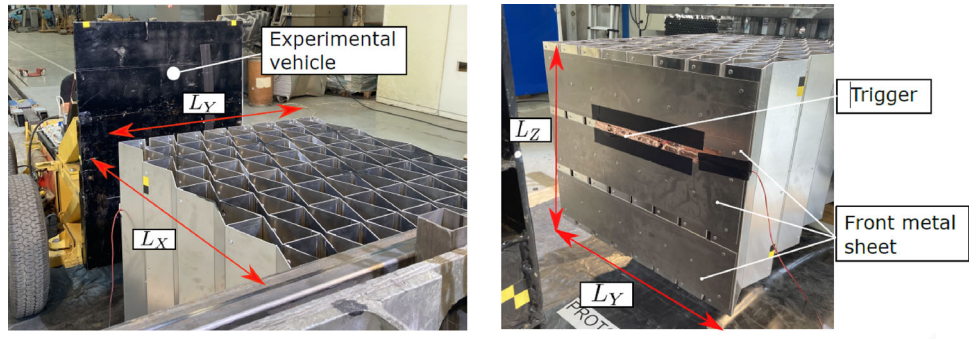


Figure 3. Zoom in on unit HR cell and a joint in the HR honeycomb specimen.

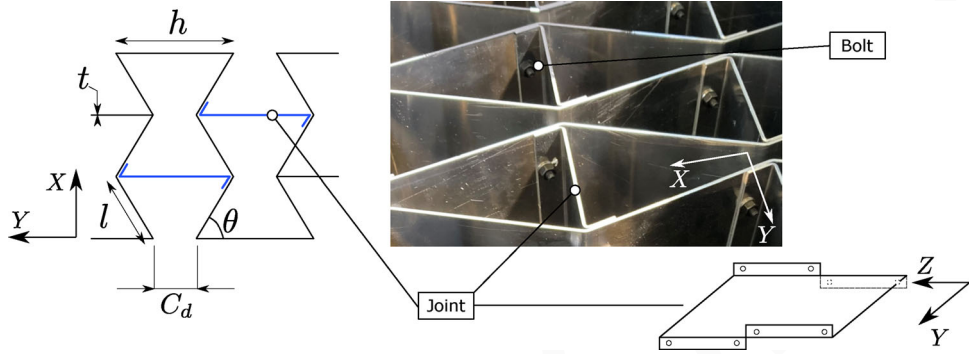


Figure 4. Global view of the HR honeycomb specimen.

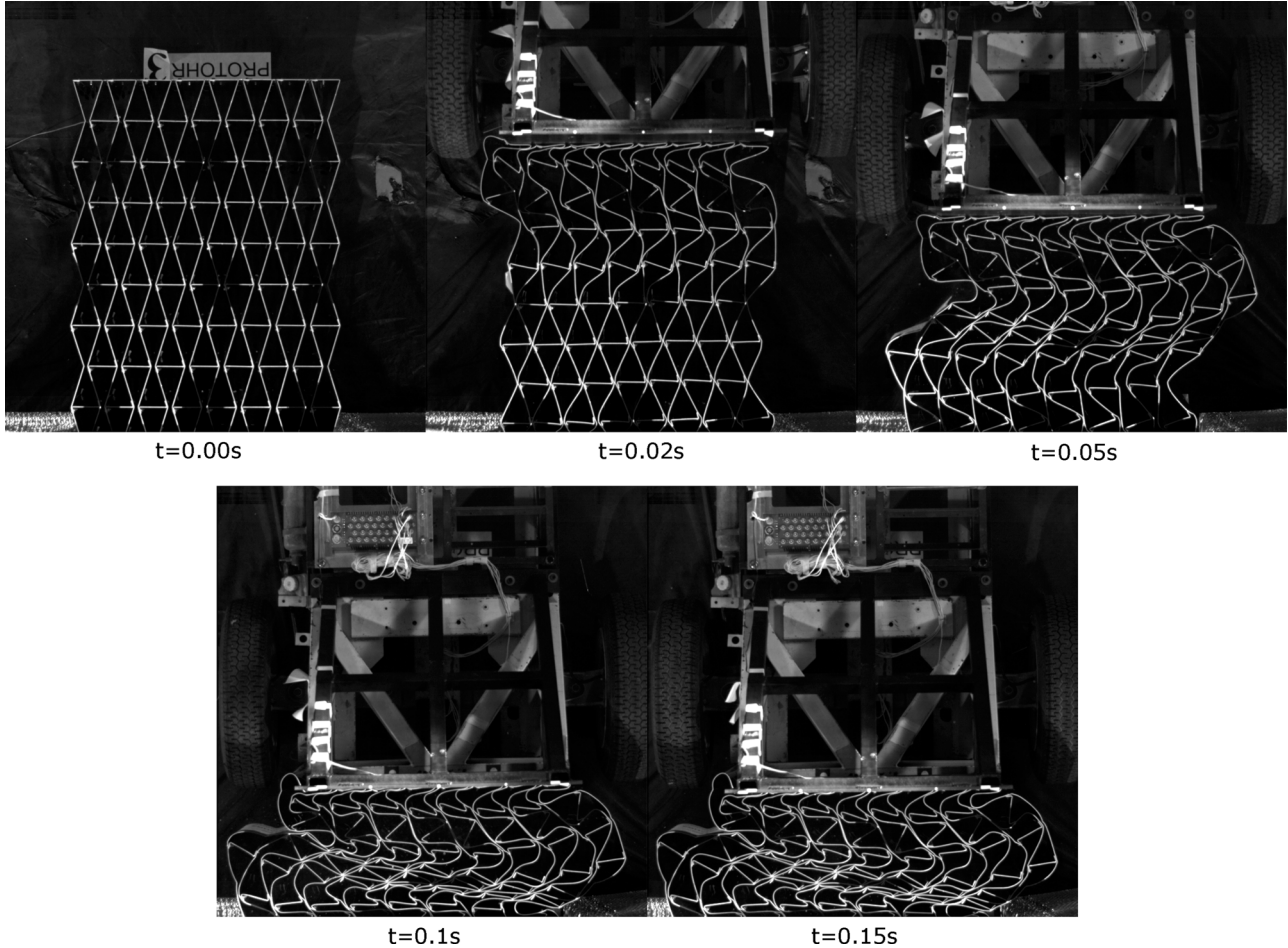


Figure 5. Global deformation of the experimental HR honeycomb specimen over time under crashed condition  $v = 13.89m/s$ .

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$$\gamma_x = \frac{F_x}{m} \quad (1)$$

The acceleration signal is filtered according to European standard (ISO6487 [31]) with a butterworth filter 4 poles at 180 Hz. The honeycomb cell walls are meshed using 4-nodes fully-integrated shell elements (Shell 163 in LS-DYNA). Whatever the value of  $l$  and  $h$  (see Figure 3), the element size is set to  $h/4$  to insure a constant number of elements in each cell. An automatic single-self contact (AUTOMATIC\_SINGLE\_SURFACE) is applied on the whole model by considering a friction coefficient of 0.35. The contact between the rigid plate and the honeycomb structure is considered frictionless. The constitutive material (aluminum alloy) of the cell-wall honeycomb is assumed to be elastic-linear-plastic (MAT\_PIECEWISE\_LINEAR\_PLASTICITY) with Young's modulus  $E$ , yield stress  $\sigma_{ys}$ , modulus tangent  $E_{tan}$  and constitutive material density  $\rho$  given in Table 4, equivalent to the AW 5754 H111 aluminum used for the experimental study presented in Section 2. The mass of the numerical honeycomb structure is equal to the whole structure mass of the experimental specimen, including bolts. As presented before, the failure does not occur during experimental tests, and thus is not considered in the numerical model.

The material of the rigid plate is assumed to be steel and not deformable and thus, neither yield stress nor tangent modulus are needed, as summed up in Table 4.

### 3.2. Model validation

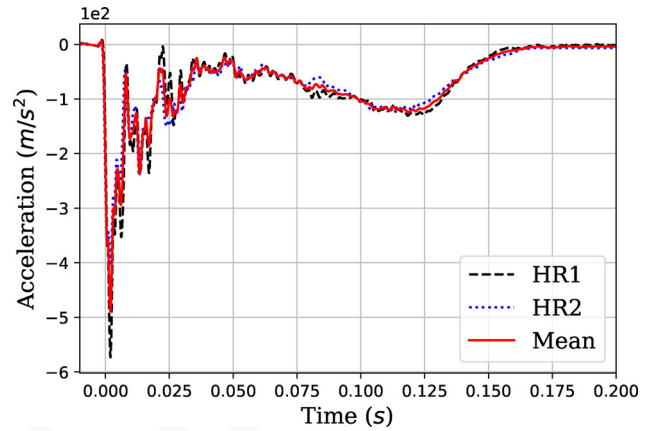
The experimental and computed results are compared in order to validate the finite elements model presented in Section 3.1. As explained in previous Section 3.1, the numerical rigid wall and experimental vehicle accelerations are compared. Both experimental and numerical signals are filtered using a Butterworth filter with 4 poles at 180 Hz as recommended in the European standard ISO6487 [31].

Firstly, in order to verify the reproducibility of the experimental setup, two experimental crash tests have been conducted on two theoretically identical HR honeycomb specimens called HR1 and HR2. Both experimentations are compared regarding the filtered acceleration signals and the striker displacement over time. The good match between HR1 and HR2 specimen crash tests of the measured accelerations and displacements of the striker are shown in Figures 6 and 7 respectively. These figures express the good reproducibility of the experimental setup. Thus, the mean values of the acceleration and the displacement of the HR1 and HR2 specimens are computed and considered as the experimental data used to validate the numerical model in following sections.

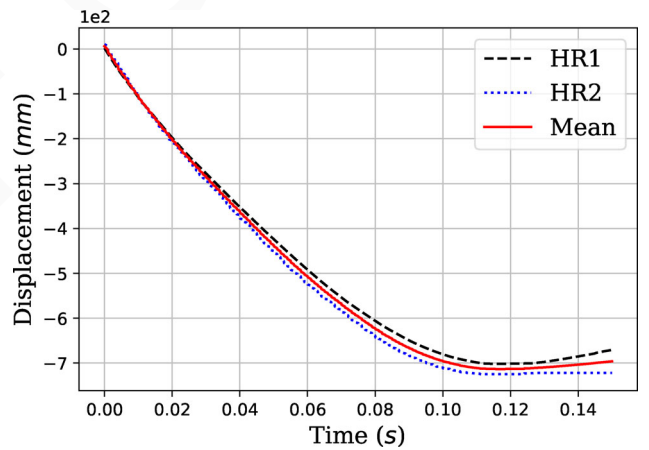
**Table 4.** Material properties of modelled aluminum alloy and rigid wall steel.

	Material	$E$ (MPa)	$\sigma_{ys}$ (MPa)	$E_{tan}$ (MPa)	$\rho$ (kg/m <sup>3</sup> )
Specimen	Aluminum alloy	69000	118	633	3080
Rigid plate	Not deformable steel	210000	-	-	7800

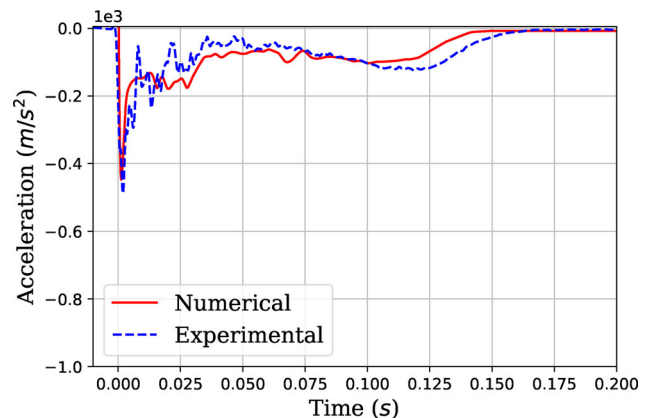
Figure 8 compares the filtered deceleration computed by the FE model, and measured on the rigid wall (red line) and the measured filtered deceleration of the crashed vehicle (blue line). The mean plateau deceleration  $\bar{\gamma}_x$  is computed as the mean deceleration between  $t = 0.04s$  and  $t = 0.12s$  and values are summed up in Table 5 for both experimental and numerical signals. Furthermore, the maximal deceleration (i.e. the minimal acceleration) is identified as  $\min(\gamma_x)$



**Figure 6.** Comparison of measured accelerations on the striker under 13.89 m/s crash conditions for both specimens and the corresponding mean acceleration.



**Figure 7.** Comparison of measured displacement on the striker under 14 m/s crash conditions for both specimens and the corresponding mean displacement.



**Figure 8.** Comparison of the acceleration on the striker under 13.89 m/s crashed conditions computed by FE analysis and experimentally measured.

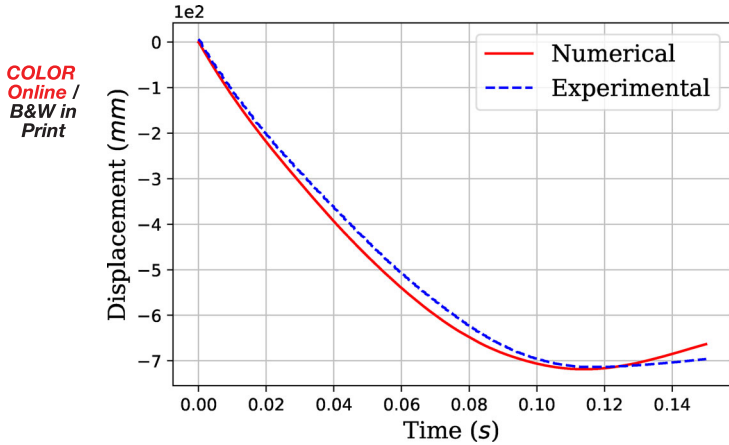
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**Table 5.** Comparison of numerical and experimental results.

	$\min(\gamma_x)$ ( $m/s^2$ )	$\bar{\gamma}_x$ ( $m/s^2$ )	$\max(D_x)$ (mm)
Experimental	-487.89	-78.59	702.2
Numerical	-447.74	-85.13	718.6
$r_e$ (%)	8.41	8.32	2.3

**Figure 9.** Comparison of the computed and measured striker displacements under 14 m/s crashed conditions.

and occurs at  $t = 7.89E-5$  seconds on both numerical and experimental acceleration curves.

The displacements of the rigid wall and the crashed vehicle in the  $X$  direction are compared in Figure 9. The maximum displacement  $\max(D_x)$  is defined as the maximal crushing distance of the striker in the  $X$  direction during the impact.

The relative errors  $r_{e,\bullet}$  between the experimental and the computed numerical values given in Table 5 are defined as:

$$r_{e,\bullet} = \frac{\bullet_{expe} - \bullet_{simu}}{\bullet_{expe}} \quad (2)$$

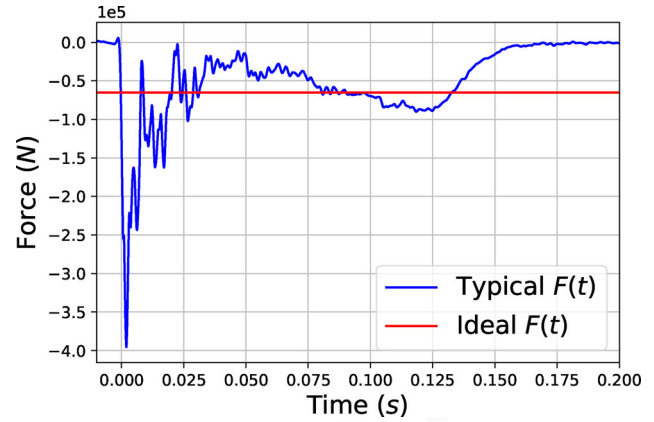
where  $\bullet_{expe}$  and  $\bullet_{simu}$  refer to the considered experimental and numerical values respectively.

The behavior of the whole structure in the FE model is similar to the experimental data given in Section 2. The relative errors around 8% on the acceleration values and less than 3% on the crushing displacement (see Table 5) are acceptable and allow to validate the proposed FE model. The errors can be explained because (i) bolts are not represented in the FE model and (ii) the tested vehicle is not guided during the impact and thus touches the specimen with a tilt angle. The developed and validated FE model will then be used within a parametric optimization process, presented in the next Section.

## 4. Optimization setup

### 4.1. Objective function

The force-displacement curve resulting from the numerical model is studied to evaluate the capability of the structure to absorb kinetic energy. One of the crashworthiness criteria is the energy absorption (EA) defined as  $EA(x) = \int_0^l F(x)dx$  where  $l$  is the total crushing distance,  $x$  represents the crushing displacement, and  $F$  is the corresponding impact

**Figure 10.** Force-time curve for a typical car crash (blue solid line) and the ideal curve (red solid line) for a mass  $m = 690$ kg.

force as described in Section 3.1. The ideal force-displacement curve to be observed over car crash should be a rectangular function (see Figure 10) [32]. However, focusing only on the force-displacement curve is not safe enough to prevent human being irreversible injuries. Indeed, the European standard is a time-based criterion, due to human tolerances [5]. The European standard uses the Acceleration Index Severity (ASI) to classify the severity of a car crash defined as follows:

$$ASI(t) = \sqrt{\left(\frac{\bar{\gamma}_x}{\hat{\gamma}_x}\right)^2 + \left(\frac{\bar{\gamma}_y}{\hat{\gamma}_y}\right)^2 + \left(\frac{\bar{\gamma}_z}{\hat{\gamma}_z}\right)^2} \quad (3)$$

where  $\hat{\gamma}_x$ ,  $\hat{\gamma}_y$  and  $\hat{\gamma}_z$  are the thresholds defined in the European standard [5] in the corresponding directions and are equal to 7g, 5g and 6g respectively.  $\bar{\gamma}_x$ ,  $\bar{\gamma}_y$  and  $\bar{\gamma}_z$  are the mean acceleration values of the crashed car over 50 ms in the corresponding directions. For a frontal crash test, the value of ASI has to be lower than 1.4 over a time period of 0.05 s to avoid human irreversible damage. The impact direction being  $x$  in this study, the ASI criterion can be written here as:

$$\bar{\gamma}_x < 1.4\hat{\gamma}_x \quad (4)$$

where  $\hat{\gamma}_x$  is set to 7g [5] with  $g$  the gravity acceleration. From Eq. (4), the striker limit force  $\bar{F}_L$  can be expressed as:

$$\bar{F}_L = m\bar{\gamma}_x \quad (5)$$

By replacing  $\hat{\gamma}_x$  by 7g in Eq. (4) it follows:

$$\begin{aligned} \bar{F}_L &< 1.4m\hat{\gamma}_x \\ \bar{F}_L &< 9.8mg \end{aligned} \quad (6)$$

where  $m$  is the mass of the striker.

In this way, it is interesting to investigate the force-time curve. The impact force  $F(t)$  should attempt to fit a rectangular function where the threshold corresponds to the striker limit force  $F_L$  inducing human severe injuries over a period of 0.05 s [5]. Hence, the occupant of the damaged vehicle avoids to suffer a force peak during the crash. Figure 10 illustrates the ideal and typical force-time curves obtained during a car crash.



Commonly in optimization, problems are formulated as follows:

$$\begin{cases} \min f(\alpha_1, \alpha_2, \dots, \alpha_p) \\ g_i(\alpha_1, \alpha_2, \dots, \alpha_p) \leq 0 \quad i = 1, 2, \dots, D \end{cases} \quad (7)$$

where  $(\alpha_1, \alpha_2, \dots, \alpha_p)$  are the  $p$  design variables to be optimized,  $f$  the objective function defined in a domain called the research domain  $\mathcal{D}_f$ ,  $g_i$  the  $D$  inequality constraints to be verified.

The newly proposed objective function  $f$  is based on the force-time curve measured during the impact. This function measures the difference between the ideal force-time curve and the force-time curve computed in the finite elements model. The objective function is calculated as follows:

$$f(\alpha_1, \alpha_2, \dots, \alpha_p) = (\bar{F}_L - \text{mean}(F_{simu}))^2 \quad (8)$$

where  $F_{simu}(\alpha_1, \alpha_2, \dots, \alpha_p)$  is the value of the computed force over time.

#### 4.2. Input data and design variables

The optimization is performed under an impact velocity  $v$  of 14 m/s. The rigid plate striking the top of the structure in the FE model has an initial kinetic energy  $K_{simu}$  defined as  $K_{car} = \frac{1}{2}mv^2$ . From European standard [6] the car mass for light vehicles is set to  $m = 900\text{kg}$ . The studied volume is equal to the volume of actual car cushions located along roads [7,26]. Thus, the volume of the re-entrant honeycomb structure in the FE model is  $V_{simu} = L_X \times L_Y \times L_Z$ .  $L_X$  and  $L_Y$  can slightly fluctuate around the defined value to maintain a set of complete cells in the structure. The input data are given in Table 6.

As previously presented in Section 4.1, the computation of the objective function using Eq. (8) depends on the variable set  $(\alpha_1, \alpha_2, \dots, \alpha_p)$ .  $\alpha_p$  are normalized design variables defined in  $[0; 1]$ . The normalized mathematical variables used in Eq. (8) correspond to the HR honeycomb geometrical parameters  $t$ ,  $l$ ,  $h$  and  $\theta$  respectively (see Table 7).

These variables, presented in Figure 2, are defined in their definition domain  $\mathcal{D}_{\beta_p}$  and take into account manufacturing constraints. However, a constraint on variables conducts to a research domain reduction. In the studied case a geometrical constraint must be considered to preserve the hexagonal geometry and is expressed in Eq. (9) according to the geometrical variables previously defined in Figure 2.

$$h - 2l\cos(\theta) > 0 \quad (9)$$

Whatever the variable set, the FE model is automatically generated and performed using LS-DYNA (see Section 3.1) considering all the problem data input.

Table 6. Input data for optimization.

$v(m/s)$	$K_{car}(kJ)$	$L_X(mm)$	$L_Y(mm)$	$L_Z(mm)$	$\bar{F}_L(N)$
14	66.64	800	1200	690	65 373

Table 7. Geometrical correspondence and design domain of all design variables  $\alpha_p$ .

Normalized variables $\alpha_p$	Geometrical variables $\beta_p$	$\mathcal{D}_{\beta_p}$
$\alpha_1$	$t$	$[0.5; 3.5] mm$
$\alpha_2$	$l$	$[30; 150] mm$
$\alpha_3$	$h$	$[30; 150] mm$
$\alpha_4$	$\theta$	$[15; 90]^\circ$

## 5. Optimization process: Inverse-PageRank-PSO

### 5.1. Methodology

The optimization method has to be carefully chosen because of the potential non-linearities of the objective function  $f$  and the potential interactions between the design variables. The studied problem is intrinsically complex and thus, leads to use a meta-heuristic optimization method. Inverse-PageRank-PSO (I-PR-PSO) is an optimization method particularly adapted to solve non-linear problems [33,34]. This algorithm has shown great global research abilities, that are interesting for solving non-linear problems. The algorithm combines two well known methods by coupling (i) Particle Swarm Optimization (PSO) and (ii) the PageRank (PR) algorithm used by the search engine Google ©.

PSO was first introduced by Kennedy and Eberhart [35], and is inspired by the analysis of bird flock movements. In this method, a population constituted of  $N_{part}$  particles, investigates the research domain in order to converge together to the global minimum of the considered objective function. The particles move over the research domain and communicate with each other about their position in the research domain. The communication strategy inside the swarm is called the population topology. Many works have investigated population topologies and suggested improvements such as PSO/ACO, PSO-SQP, PSO-GWO or I-PR-PSO [33,36–38].

In I-PR-PSO, the particles connectivity can be seen as an oriented graph. In this way, the influence of the particles on the others is weighed regarding their respective fitness values. A particle which exhibits a low value of fitness influences more the next position of other particles than a particle exhibiting a high value of fitness (for a minimization problem). The position  $\mathbf{X}_i^{k+1}$  of the particle  $i$  at the iteration  $k+1$  is calculated as follows:

$$\begin{cases} \mathbf{V}_i^{k+1} = \omega \times \mathbf{V}_i^k + c_1 \times \text{rand}_1 \times (\mathbf{P}_{i,best}^{k+1} - \mathbf{X}_i^k) + c_2 \times \text{rand}_2 \times \sum_{j=1}^n C_{ij} \times (\mathbf{P}_{j,best}^{k+1} - \mathbf{X}_i^k) \\ \mathbf{X}_i^{k+1} = \mathbf{X}_i^k + \mathbf{V}_i^{k+1} \end{cases} \quad (10)$$

where  $\mathbf{P}_{i,best}$  is the best personal position of the particle  $i$ ,  $\mathbf{V}_i^k$  is the velocity<sup>1</sup> of particle  $i$  at iteration  $k$ ,  $\omega$  is the inertia weight of the particle,  $rand_1$  and  $rand_2$  are random numbers in  $[0; 1]$ ,  $c_1$  and  $c_2$  are coefficients representing the social and individual behavior of the particle and  $\mathbf{C}$  is the stochastic connectivity matrix defined by the Inverse-PageRank algorithm. Particle  $i$  is influenced by all the particles of the swarm, and their respective influences are given by the components of the  $j$ th line of  $\mathbf{C}$ . The component  $C_{ij}$  corresponds to the influence of all particles on particle  $i$ . Indeed, the particle swarm can be seen as a Markov chain, in which each link represents the influence of a particle on another one. In the same way as in the search engine Google, a PageRank score is attributed to each particle according to its own success, that allows to build the stochastic connectivity matrix  $\mathbf{C}$ . The reader can refer to [34] for more details about I-PR-PSO.

I-PR-PSO is applied to the FE model of the HR honeycomb crashing (see Section 3.1) in order to find the best combination of the geometrical parameters to efficiently absorb the kinetic energy in car-crashed conditions. In other words, the algorithm is employed to find a parameter set for which the objective function  $f$  defined in Section 4.1 reaches a minimum. In the considered problem, each particle  $i$  is represented by a vector  $\mathbf{X}_i^k$  containing the position of the particle in the research domain  $\mathcal{D}_f$ , at iteration  $k$ . The particle coordinates correspond to normalized values of the design variables in their respective domains as follows:

$$\mathbf{X}_i^k = \begin{pmatrix} \alpha_{1,i}^k \\ \alpha_{2,i}^k \\ \alpha_{3,i}^k \\ \alpha_{4,i}^k \end{pmatrix} \quad (11)$$

For each particle  $i$ , at each iteration  $k$ , a FE calculation is launched, and the value of the objective function  $f(\mathbf{X}_i^k)$  is calculated by using Eq. (8).

As explained in Section 4.2, the problem is highly constrained because of the optimization constraint imposing  $h - 2l \cos(\theta) > 0$ . A simple penalty method is applied to a particle that does not respect the condition [39]. In the studied case, if Eq. 9 is not verified, the fitness value is not computed in order to save computational time and because the structure might be unmanufacturable. A penalty value, set to  $10E30$ , is attributed to the considered particle to isolate it from the swarm. Thus, the particle which does not verify Eq. (9) has an insignificant influence on its pairs.

### 5.2. I-Pr-PSO parameters

As explained in the latter Section, I-PR-PSO parameters have to be carefully chosen to ensure the algorithm convergence. As for classical PSO, a large diversity of parameter

**Table 8.** I-PR-PSO parameters used in Eq. (10).

$\omega$	$c_1$	$c_2$	$\Delta_c$	$N_{part}$	$it_{PSO,MAX}$	$it_{conv}$
0.6	1.5	1.5	0.033	8	45	30

sets are discussed [40–45]. Indeed, the parameter set is strongly problem-dependent as demonstrated in [46].

The values of I-PR-PSO parameters ( $\omega$ ,  $c_1$  and  $c_2$ ) are summed up in Table 8. The algorithm can be considered as converged either when the number of iterations  $k$  reaches the maximum number of iterations (set by the user)  $it_{PSO,MAX}$  or when the value of the best position does not improve after  $it_{conv}$  iterations.

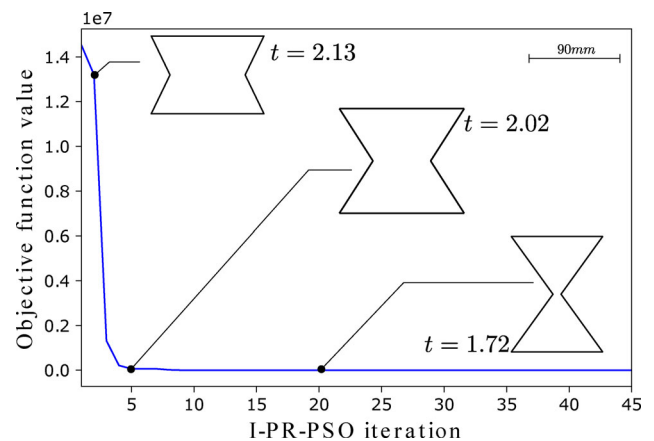
Parameters can be considered as a good selection when the algorithm presents an efficient convergence curve (i.e. the evolution of the best fitness value found over iterations). So, in the following optimization process, a convergence study is presented, where I-PR-PSO parameters have been empirically found.

The values of  $c_1$  and  $c_2$  gradually decrease while I-PR-PSO process progresses.  $\Delta_c$  is subtracted to  $c_1$  and  $c_2$ , at each iteration  $k$ .

## 5.3. Obtained results

### 5.3.1. Optimal structure

This section presents the results obtained by the optimization process. The numerical model used within the optimization has been previously presented in Section 3.1, and validated based on experimental tests presented in Section 3.2. Four design variables have been optimized, corresponding to the geometrical parameters of the HR honeycomb structure. The optimization algorithm, Inverse-PageRank-PSO, proposes optimized values of these parameters, effectively minimizing the objective function defined in Section 4.1. The convergence curve, presenting the best values of the objective function found so far during the optimization process and the corresponding cell dimensions, is plotted in Figure 11. Moreover, the tested values of the design variables, which are the positions of the particles during the process, are scattered in Figure 12. As one can see in these figures: (i) the research domain is explored until the 20<sup>th</sup> iteration, where the swarm switches from exploration to exploitation, the particles converging together, by following each other, to promising zones of the research domain. (ii) The objective function is optimized, switching from an



**Figure 11.** Convergence curve of the objective function for HR honeycomb optimization and corresponding cell dimensions.

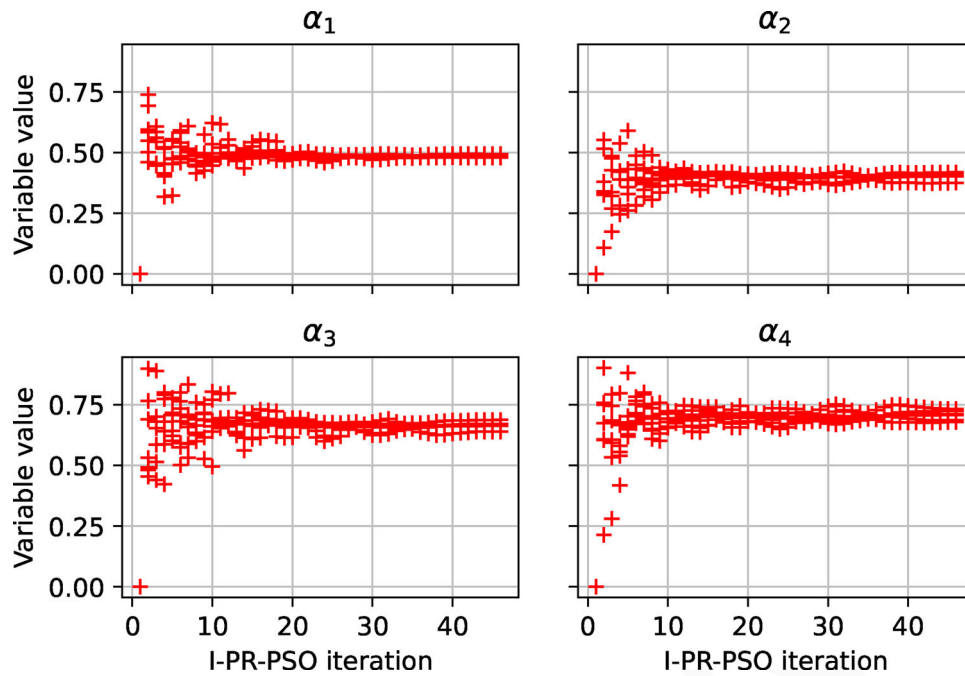


Figure 12. Evolution of particles position in the research domain for HR honeycomb optimization.

Table 9. HR honeycomb optimal geometrical parameters for energy absorption.

$t(\text{mm})$	$l(\text{mm})$	$h(\text{mm})$	$\theta(^{\circ})$
1.95	77.90	109.27	60.705

initial value of  $G_{best} = 1.45E + 07$ , to a minimal value of  $G_{best} = 68.94$  at the end of the optimization process.

The optimal values of the design variables are given in Table 9, corresponding to the best values of the geometrical parameters of the HR honeycomb.

A FE analysis is carried out on the optimal structure. The evolution of the force computed on the striker over time  $F_{opt}(t)$  is presented in Figure 13. As one can see, due to the fitting process induced by the using of the objective function presented in Section 4.1, the force actually fits the limit force  $\bar{F}_L$ , and remains at this plateau value during the whole crash. Therefore, the optimization process, by minimizing the objective function, permits to smooth the force suffered by the vehicle occupant, limiting the human damage. The plateau force obtained in Figure 13 induces a quasi-linear decreasing of the striker kinetic energy, with a slope of  $6.46E + 08$ , as presented in Figure 14. As one can see, at the crushing beginning, the kinetic energy slope is steeper ( $1.17E + 11$ ), due to the force  $F_{opt}(t)$  exhibiting a peak at the crushing instant. However, the acceleration criterion, previously defined in Eq. (3) and being proportional to  $F_{opt}(t)$ , is equal to  $9.79g$ , which is inferior to the limit of  $9.8g$  imposed by European standards [5]. The optimal HR honeycomb behavior is then smoothly absorbing kinetic energy as a cushion device. The optimal HR honeycomb structure absorbs 79% of the initial energy. The optimization of the mean force over time, as presented in Eq. (8), allows the kinetic energy absorption to be spread over a sufficiently long time to minimize the damage to the vehicle occupants, as well as verifying the European standard

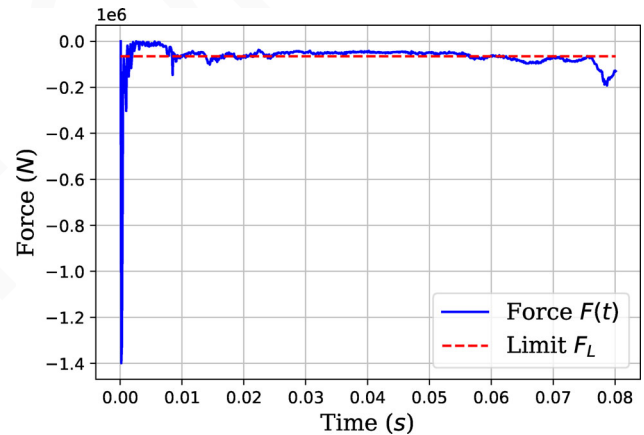


Figure 13. Force curve  $F(t)$  and limit force  $\bar{F}_L$  of the optimal HR honeycomb.

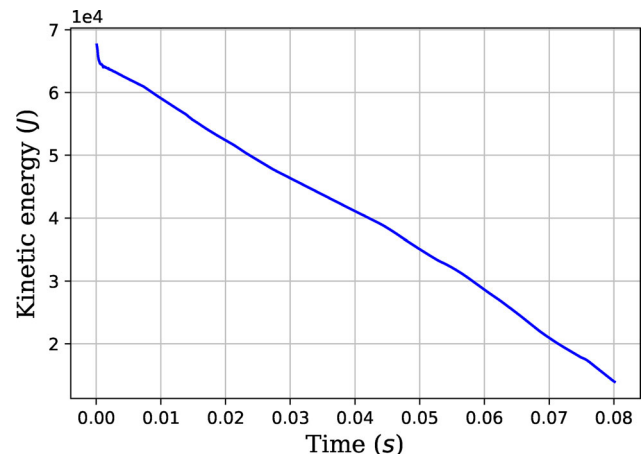


Figure 14. Kinetic energy absorption over the time for optimal HR honeycomb.

criterion. As one can see in Figure 15 presenting the structure deformation over time, the negative Poisson ratio effect occurs. At  $t = 0.018s$ , a V-mode deformation is observed on

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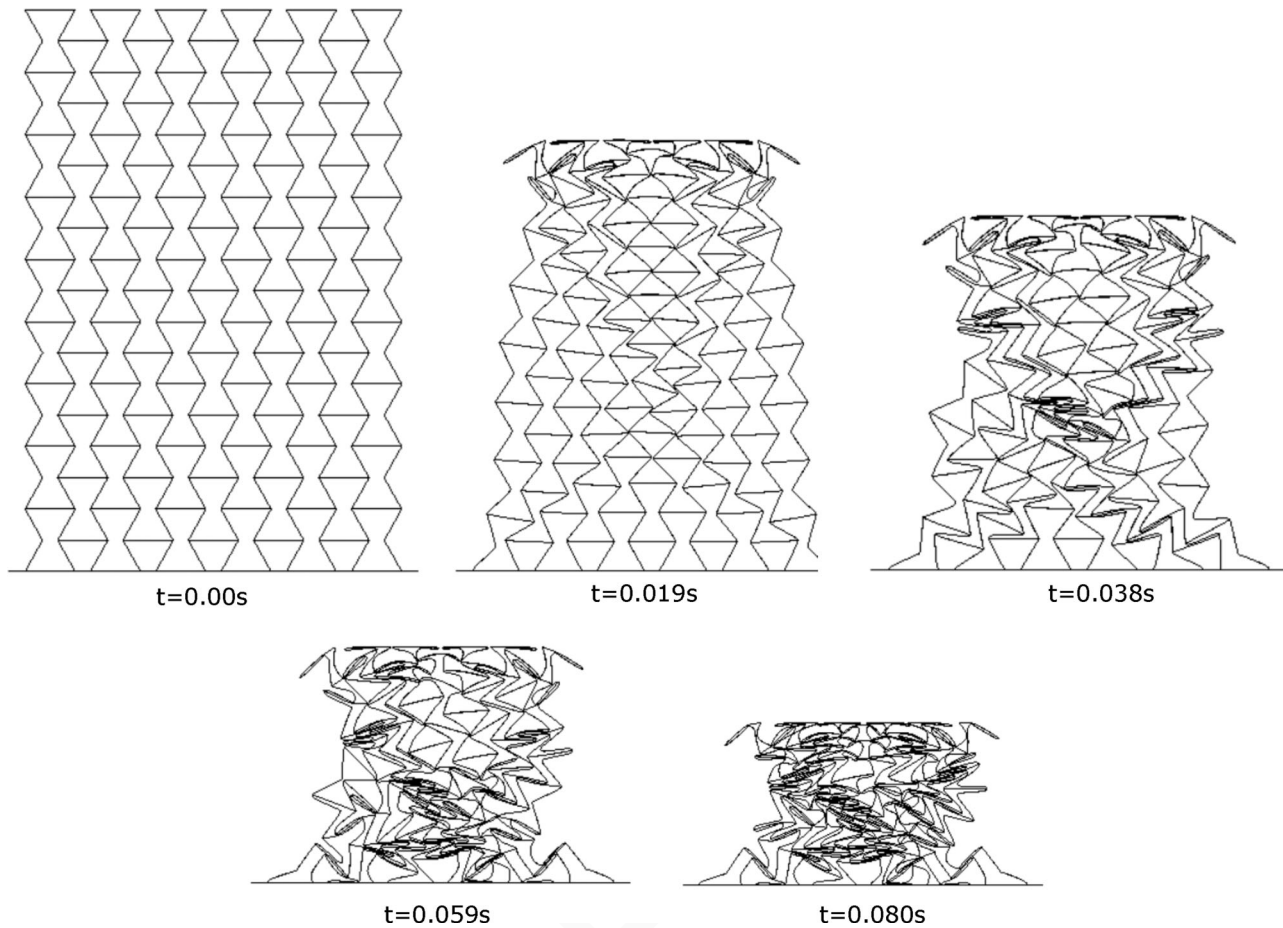


Figure 15. Finite elements simulation performed on the optimal HR honeycomb structure.

the top of the structure, while a double horizontal V-mode deformation occurs at  $t=0.06s$ . This behavior induces the efficient kinetic energy absorption, as it was demonstrated in [47]

### 5.3.2. Off-centered and titled crash test on the optimal structure

The optimal structure presents a good efficiency under frontal impact conditions as presented in Section 5.3.1. In road safety area, the probability to crash a car-cushion device in a perfectly frontal impact condition is very low and thus, the behavior of the HR honeycomb structure under off-centered and titled impact conditions must be verified. In this Section, two impact configurations described in the European Standard EN1317-3 [6] are considered on the optimal HR honeycomb structure obtained in the previous Section 5.3.1. The impact configurations can be described with two parameters, the titled angle and the off-set distance applied to the rigid wall in the FE model defined in Section 3.1. The off-set distance is defined as  $\frac{1}{4} \times L_Y$  in EN1317-3 leading to the impact parameters given in Table 10. The trajectories of the center of gravity of the rigid wall in the three configurations are described in Figure 16.

Two new FE computations are added to the FE analysis carried out on the optimal HR honeycomb structure (Section 5.3.1) in the  $N^\circ 1$  and  $N^\circ 2$  impact configurations

Table 10. Impact configuration.

Configuration	Titled angle ( $^\circ$ )	Off-set (mm)
$N^\circ 0$	0	0
$N^\circ 1$	0	205
$N^\circ 2$	15	0

given in Table 10. Energy absorption curves and force curves are presented and compared in Figures 17 and 18 respectively. These figures demonstrate that the global behavior of the optimal HR honeycomb is not sensitive to the impact conditions. Indeed, whatever the impact configuration considered, the kinetic energy absorbed after 0.08 s is in a range of 70 to 79% of the initial kinetic energy. Moreover, the energy absorption curves remain quasi-linear for all impact configurations. For both frontal impacts, the curves are similar in the first 0.03 s of the crash, with a steeper slope at the beginning. The tilted impact presents no steeper slope at the beginning, and thus is much smoother than the two other impact conditions. The soft decreasing of energy absorption is also noticeable on the not-filtered force-time curve given in Figure 18. The part of kinetic energy absorbed by the HR honeycomb structure after  $t=0.08s$  are given in Table 11 for the three configurations. The tilted-force is a quasi-linear plateau. The centered and off-centered force curves present a plateau that reflects the quasi-linear behavior of the energy absorption curve while there is a peak force at the beginning of the contact. The plateau force under tilted configuration is lower than the

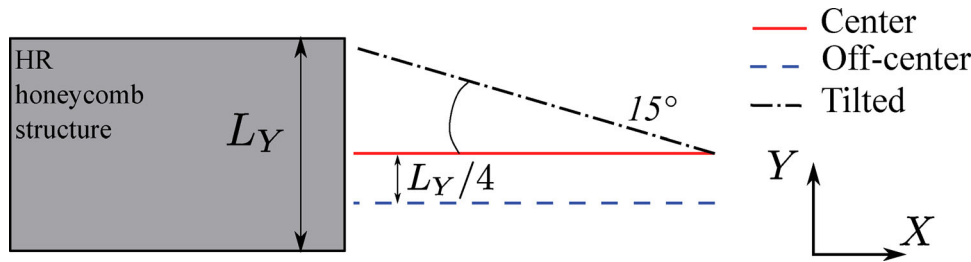


Figure 16. Center of gravity trajectories of the three impact configurations.

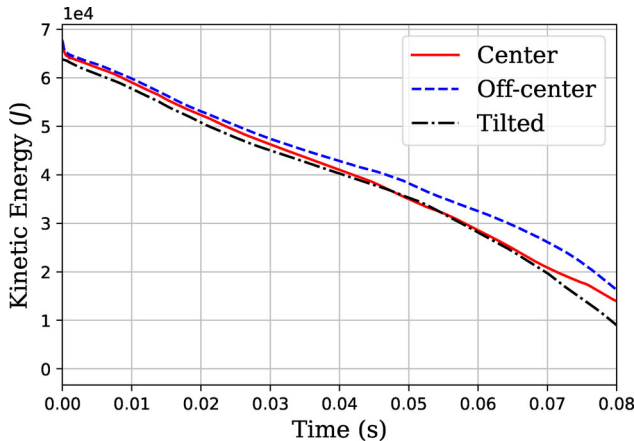


Figure 17. Energy absorption curves computed on the optimal HR honeycomb structure under i) frontal ii) off-centered frontal and iii) tilted impact configurations.

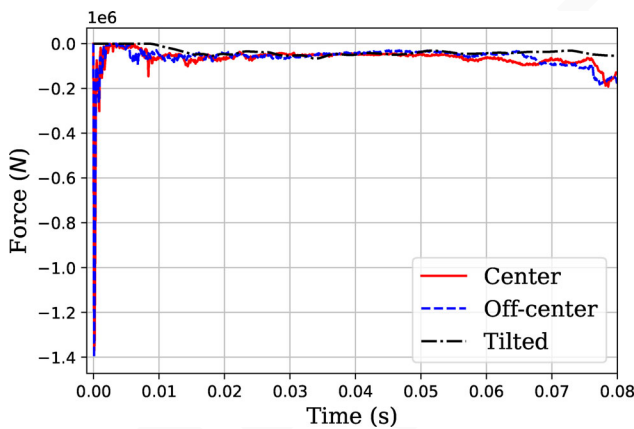


Figure 18. Force-time curves computed on the optimal HR honeycomb structure under i) frontal ii) off-centered frontal and iii) tilted impact configurations.

Table 11. Results on the optimal HR structure under different impact configurations.

Configuration	Kinetic energy absorbed (%)	ASI
N ° 0	79.0	1.311
N ° 1	75.9	1.232
N ° 2	70.5	0.804

frontal off-centered and the centered configurations. Indeed, the frontal centered configuration is the most critical human being configuration and thus presents a higher value of the plateau force. To illustrate these curves, the ASI (defined in Section 4.1) are computed for the three configurations and given in Table 11. For all the configurations, the ASI value is lower than the limit set to 1.4 by the European standard [5].

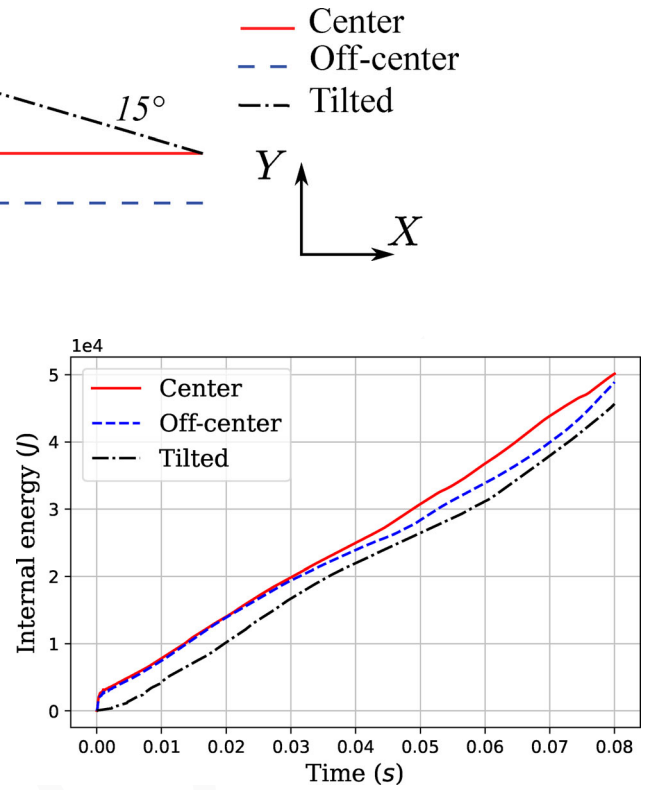


Figure 19. Finite elements simulation performed on the optimal HR honeycomb structure in the tilted impact configuration.

The numerical deformation of the optimal HR honeycomb under off-centered and tilted impact configurations are illustrated in Figures 19 and 20 respectively. Around  $t = 0.02$ s, the negative Poisson ratio effect allows a densification of the structure at the centre of the honeycomb while the upper part of the structure is also crushed. The quasi-linear energy decreasing described in Section 5.3.1 is allowed thanks to the deformation occurring in the whole structure at the same time. Indeed, the optimization process carried out on the FE model has led to the design of a structure able to absorb energy mostly by plastic deformation spread on the whole structure at the same time. Around 74% of the initial kinetic energy is absorbed by plastic deformation as illustrated in Figure 21 in which the internal energy in the HR honeycomb over time is plotted for the three different impact configurations.

Whatever the impact conditions, the optimal HR honeycomb structure obtained in this optimization study is efficient in car crash energy absorption, and meets the European standard criterion defined in Section 4.1.

## 6. Discussion and conclusion

The hexagonal re-entrant honeycomb is a promising structure in energy absorption partly because of its auxetic behavior. This study investigates an optimal geometry of HR honeycomb structure under car-crashed impact velocity which corresponds to speed limitations in the European standard [5].

An experimental study is conducted on a HR honeycomb specimen which is impacted by a rigid vehicle launched at  $v = 14$ m/s. The specimen, made of aluminum and

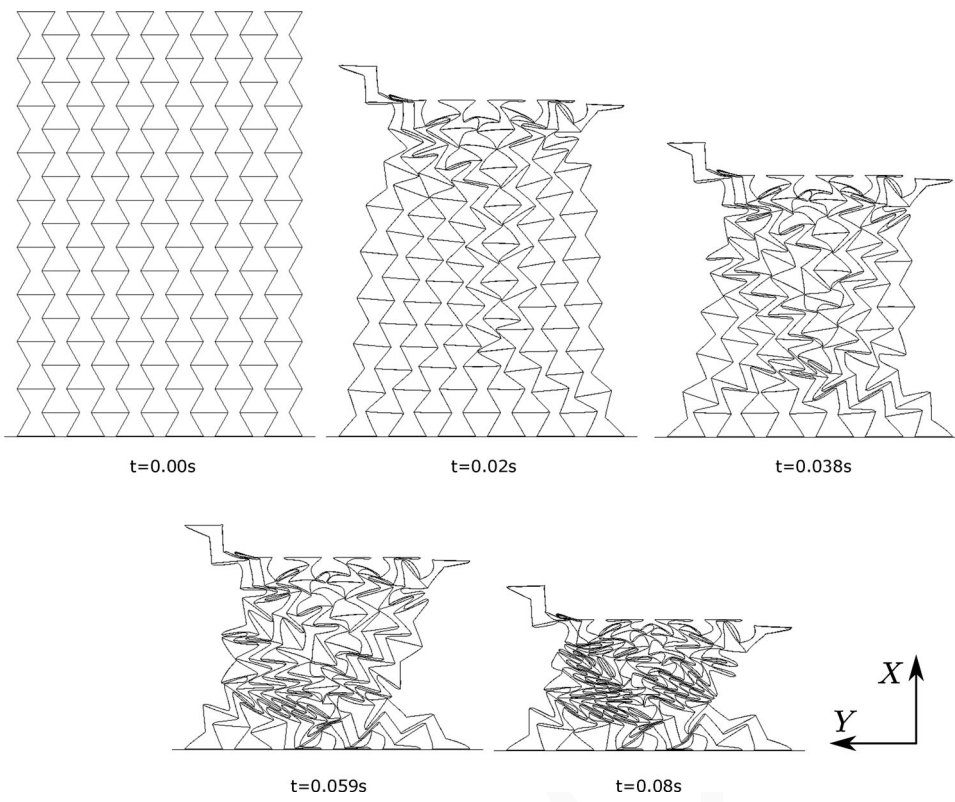


Figure 20. Internal energy curves computed on the optimal HR honeycomb structure under i) frontal ii) off-centered frontal and iii) tilted impact configurations.

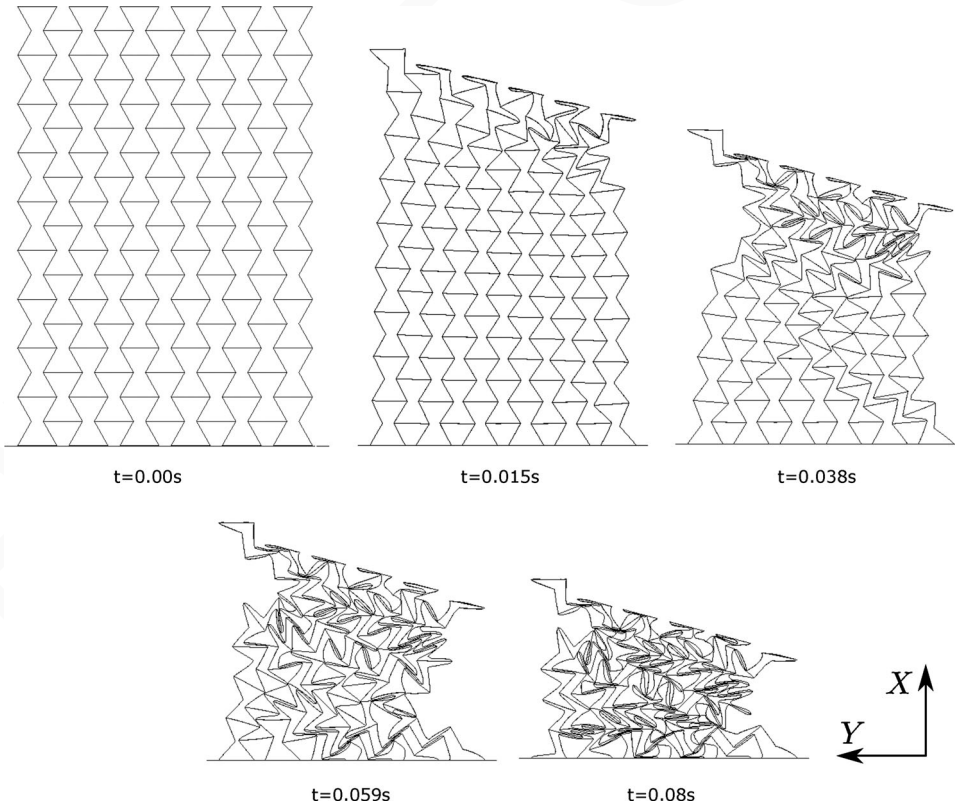


Figure 21. Finite elements simulation performed on the optimal HR honeycomb structure in the off-centered impact configuration.

conventionally assembled, is on the scale of car-cushion devices. The experimental analysis led to the FE model validation by comparing measured and computed signals of displacement and acceleration. The relative differences

around 8% are considered as acceptable and can be due to assembly simplifications in the FEM.

The objective function proposed shows a good efficiency to measure the absorbing capabilities of the considered

structure to be optimized. This new objective function is based on the European standard criteria [6], and conducts to optimize the force-time curve, by trying to fit an ideal rectangular force-time curve. This objective function is computed with the validated FE model, and depends on the HR honeycomb geometry. Four parameters,  $t$ ,  $l$ ,  $h$  and  $\theta$  are defined to describe the honeycomb structure. Moreover, a geometrical constraint is modeled and corresponds to manufacturing considerations.

A global optimization process is applied to the HR honeycomb structure and performed with the Inverse-PageRank-PSO (I-PR-PSO) algorithm. This meta-heuristic method has shown high efficiency in complex real-problem optimization. A penalty method is applied to I-PR-PSO in order to apprehend the geometrical constraint of the HR honeycomb structure. The optimal structure found presents a quasi-linear energy absorption curve, that can be considered as ideal, with 80% of the initial kinetic energy absorbed. This performance will be increased in real-world application because the crashed car also absorbs energy due to its own deformation not considered in this study. Finally, the optimal HR honeycomb structure obtained in the optimization process is tested under three different impact configurations as defined in the European standard EN1317-3. The optimal HR structure presents high performance in energy absorption whatever the impact configuration considered. This study gives an optimization process with a validated objective function to design new car crash cushion devices, and could be used in other crash contexts such as different materials and impact velocities, or to design truck-cushion devices by adjusting the mass of the striker.

## Note

1. Called "velocity" in the literature,  $V$  is actually a displacement imposed to particles.

## Disclosure statement

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