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# Mathematical encoding within multiresonant planetary systems as SETI beacons 

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#### Abstract

How might an advanced alien civilization manipulate the orbits within a planetary system to create a durable signpost that communicates its existence? While it is still debated whether such a purposeful advertisement would be prudent and wise, we propose that mean-motion resonances between neighbouring planets - with orbital periods that form integer ratios - could in principle be used to encode simple sequences that one would not expect to form in nature. In this letter, we build four multiresonant planetary systems and test their long-term orbital stability. The four systems each contain six or seven planets and consist of (i) consecutive integers from 1 to 6; (ii) prime numbers from 2 to 11 ; (iii) the Fibonacci sequence from 1 to 13; and (iv) the Lazy Caterer sequence from 1 to 16 . We built each system using $N$-body simulations with artificial migration forces. We evaluated the stability of each system over the full 10 Gyr integration of the Sun's main-sequence phase. We then tested the stability of these systems for an additional 10 Gyr , during and after post-main-sequence evolution of the central stars (assumed to be Sun-like) to their final, white dwarf phase. The only system that was destabilized was the consecutive integer sequence (system i). The other three sequences therefore represent potential Search for Extraterrestrial Intelligence (SETI) beacons.


Key words: extraterrestrial intelligence - astrobiology - planets and satellites: dynamical evolution and stability - stars: evolution.

## 1 BACKGROUND

Over the last few decades, the search for intelligent life within the cosmos has evolved from a primarily radio-centric enterprise (e.g. Cocconi \& Morrison 1959; Drake 1961; Tarter 2001) to one that includes an increasingly broad range of possible technosignatures (NASA Technosignatures Workshop Participants 2018). The longevity of such signatures is directly proportional to their detectability (Balbi \& Ćirković 2021) and thus those that could persist for Myr, or even Gyr, are of particularly keen interest. However, this raises questions about what kind of technosignature could feasibly produce a detectable signature for so long, especially when one recognizes that most indications of human technology would likely be undetectable within a few million years (Schmidt \& Frank 2019).

Perhaps a sufficiently advanced civilization could engineer an active mechanical/electrical system able to reliably function over a Gyr or more, but an alternative approach is to consider a passive system inherently more robust to entropy's penchant for destruction. In the case of deliberate messaging, one proposed example is the construction of artificial transiting screens (Arnold 2005). Although the longevity of such structures is unclear, and surely

[^0]sensitive to the specific design and environment considered, the transmission energy is essentially provided by the star. By exploiting this astrophysical object, the power source is thus inherently longlived. However, the screens themselves may not be, potentially suffering degradation from dust abrasion (Hoang et al. 2017; Müller et al. 2019), radiation damage, or orbital instability (Batygin \& Laughlin 2008). This raises the question - in the same way that the power source can be astrophysical, and thus long-lived, could one use other astrophysical objects in place of the screens to foster longevity?

To serve as a technosignature though, the astrophysical object(s) in question would need to somehow clearly encode a signature of non-natural origin. Since the objects themselves are by definition astrophysical, inspection of their individual nature would not be sufficient. Instead, a degree of non-naturalness would have to be achieved through their arrangement. The orbital architecture of companion planets thus emerges as a possible solution. In particular, we propose that multiresonant planetary systems could serve this role, since such configurations can be persistent over Gyr and can encode a series of integers. Each pair of neighbouring planets forms a mean-motion resonance with an orbital period ratio that is the fraction of two integers. A chain of resonances thus encodes a sequence, one that could potentially signal intelligence (e.g. a mathematical sequence; Sagan 1975) and be arranged through local manipulation (see e.g. Korycansky, Laughlin \& Adams 2001).

Table 1. Summary of resonant chains tested in our investigation. The columns are as follows: (1) the simulation name, (2) the total number of planets in the chain, (3) the orbital period ratios of each consecutive pair of planets from inside out (i.e. $P_{1}: P_{2}: P_{3}$, etc.), and (4) the mass of each planet starting with the innermost body in Earth units.

| Name | $N_{\mathrm{pln}}$ | Resonant chain | $M_{\mathrm{pln}}\left(\mathrm{M}_{\oplus}\right)$ |
| :--- | :---: | :---: | :---: |
| Prime 1 | 6 | $1: 2: 3: 5: 7: 11$ | $20,20,20,20,2,2$ |
| Prime 2 | 6 | $1: 2: 3: 5: 7: 11$ | $2,2,2,2,2,2$ |
| Fibonacci 1 | 7 | $1: 1: 2: 3: 5: 8: 13$ | $20,20,20,20,20,2,2$ |
| Fibonacci 2 | 7 | $1: 1: 2: 3: 5: 8: 13$ | $2,2,2,2,2,2$ |
| Lazy Caterers ${ }^{1}$ | 6 | $1: 2: 4: 7: 11: 16$ | $20,20,20,20,2,2$ |
| Integer | 6 | $1: 2: 3: 4: 5: 6$ | $20,20,20,20,2,2$ |

${ }^{1}$ https://oeis.org/A000124
Of course, one challenge is that chains of orbital resonances frequently form in nature. Jupiter's Galilean moons are the firstdiscovered example, and the Trappist-1 exoplanet system is the longest currently known (seven-planet) resonant chain (Gillon et al. 2017; Luger et al. 2017). Resonant chains are naturally produced during planet formation as a consequence of orbital migration in gaseous protoplanetary discs (Lee \& Peale 2002; Kley \& Nelson 2012). However, simulations show that not all resonances are populated equally: first-order resonances (in which the integer sequence differs by 1 , such as the $4: 3,3: 2$, or $2: 1$ resonances) are by far the most common outcomes (Terquem \& Papaloizou 2007; Pierens \& Nelson 2008; Izidoro et al. 2017). Chains of 3:2 or $2: 1$ resonances are thus expected to form naturally and examples have indeed been discovered (e.g. the K2-138 chain of 3:2 resonances and the HR8799 chain of $2: 1$ resonances; Marois et al. 2010; Christiansen et al. 2018). Extended chains encoding abstract mathematical sequences have not been discovered and thus their detection would be quite curious. Two basic questions emerge concerning such sequences: (i) how often would they naturally form (the false-positive rate to a technosignature search) and (ii) how long-lived could we expect them to be?

In this work, we limit our scope to the second of these questions. After all, if such sequences are not long-lived, their spurious formation would be inconsequential anyway. Our approach is to first build four resonant chain planetary systems (Section 2) that each encode a specific integer sequence (see Table 1). While we assume that such a signpost could in principle have been produced by a highly advanced civilization, we expect its later evolution to be passive and governed by natural laws. However, it is worth acknowledging that an alien species advanced enough to harvest the substantial quantity of energy necessary to assemble such a resonant signpost might also be capable of devising a long-lasting system that would prevent orbital instabilities by applying slight deviations to the planets' orbits. Nevertheless, we proceed by testing the dynamical stability of each system during its central star's 10 Gyr main-sequence lifetime (Section 3) and the next 10 Gyr of post-mainsequence evolution (Section 4). We discuss the interpretation of such signatures in Section 5.

## 2 BUILDING RESONANT CHAINS

Resonant planetary systems are naturally produced during gas-driven migration (Lee \& Peale 2002; Batygin \& Morbidelli 2013; Pichierri, Morbidelli \& Crida 2018). We leverage the standard approach of modifying the equations of motion to mimic gas disc-induced orbital migration by including forced migration $(\dot{a})$ and eccentricity damping ( $\dot{e}$ ) in the MERCURY $N$-body integration package (Chambers 1999). We made use of MERCURY's Bulirsch-Stoer algorithm (with
a tolerance setting of $1.0 \times 10^{-15}$ ), as the force on any given particle is a function of both its position and momentum. The time-step was set to 1 per cent of the shortest orbital period in the simulation, and the central star's mass was set to $1.0 \mathrm{M}_{\odot}$.

To produce multiresonant planetary systems, we initialized each planet from inside out on a circular, co-planar orbit with a semimajor axis just outside of the desired resonant semimajor axis (with a period ratio $\sim 1$ per cent larger than the desired final period ratio; see Table 1). For consistency, each of our resonant chains places the innermost planet at 5.0 au . This selection ensures that none of the planets will be engulfed during the asymptotic giant branch (AGB) evolutionary phase. We then applied migration forces, integrating each system up to $t=10 \mathrm{Myr}$. We verified that pairs of planets are indeed in the desired resonance by inspecting resonant angles of the form
$\phi=p \lambda_{2}-q \lambda_{1}-r \varpi_{2}-s \varpi_{1}$,
where $\lambda_{i}$ and $\omega_{i}$ are the mean longitudes and longitudes of perihelia of each planet (with planet 2 being the outer one). The integers $p$ and $q$ define the resonance's period ratio, $p: q$, and the remaining integers satisfy
$p=q+r+s$.
In resonance, the values of $\phi$ librate about a specific value as the consequence of the fact that resonant entrapment shields objects from close approaches by coupling their orbital and precession frequencies (see Fig. 1).

We iterated through migration simulations until the desired resonance was established and remained intact throughout the initial 10 Myr simulation. Several migration simulation attempts are often required to determine input parameters that prevent the distant planets from being captured in powerful first-order resonances (e.g. 3:2 instead of 11:7); however, we do not see this as a major obstacle for an advanced civilization. For additional details on our numerical pipeline for generating unique resonant chains, we direct the reader to Clement et al. (2021).

## 3 MAIN-SEQUENCE STABILITY

We tested the dynamical stability of each of the resonant chains (see Table 1) over the full 10 Gyr duration of their (Sun-like) central star's main-sequence lifetime. We used the Hybrid integrator in the MERCURY integration package (Chambers 1999), again using a timestep that was 1 per cent of the innermost planet's orbital period.

Fig. 1 plots the 10 Gyr evolution of the relevant resonant angles (equation 1) of our four resonant chains (see Table 1). The only chain to destabilize was the more compact chain of consecutive integers. In this case, the outer four planets spontaneously dislodged from resonance around 7.2 Gyr into the simulation. Once the resonant phase protection mechanism terminated, all six planets began to interact strongly; thus exciting the eccentricity of each body to $\sim 0.20$ within the next 2 Myr . After this occurred, the first and fourth planets collided and the fifth and sixth planets were ejected after scattering off of the third planet.

Our five other chains each remained stable for 10 Gyr , with resonances maintained. Of course, while this demonstrates the general integrity of these configurations, it does not prove that they would evolve regularly in the presence of additional planets and small bodies. Indeed, perturbations from small bodies and additional non-resonant (or resonant) planets are known to drive instabilities in systems of otherwise stable planets both in the Solar system (Nesvorný \& Morbidelli 2012) and elsewhere (e.g. Raymond, Armitage \& Gorelick 2010; Izidoro et al. 2017).


Figure 1. Resonant angle evolution in our 10 Gyr simulations of our Prime 2 (upper left panel), Fibonacci 2 (upper right panel), Lazy Caterers (lower left panel), and Integer (lower right panel) resonant chains (Table 1). Each successive subplot depicts the dominant resonant angle (equation 1) for each respective pair of planets, starting with the innermost object. The top subplot in the upper right panel (Fibonacci chain) and the accompanying inset depict the position of the Trojan planet with respect to the inner planet (i.e. in the rotating reference frame: $\lambda_{1}-\lambda_{2}$ ), along with the positions of the L4 and L5 Lagrange points.

## 4 POST-MAIN-SEQUENCE EVOLUTION

To propagate our simulations beyond the main sequence, we applied the $N$-body code described in Mustill et al. (2018). This code incorporates stellar evolution output from the SSE code (Hurley, Pols \& Tout 2000) into the framework of the RADAU integrator from the MERCURY software package (Chambers 1999). Our code is the same one used to determine the post-main-sequence breaking
of the four-planet resonant chain in the HR 8799 system (Veras \& Hinkley 2021).

The physical and chemical properties of the parent star determine its evolutionary profile. Because here we are just seeking to demonstrate a proof of concept, we used a single set of fiducial parameters. In addition to adopting a main-sequence stellar mass of $1.0 \mathrm{M}_{\odot}$, we also assumed Solar metallicity, a Reimers mass-loss coefficient


Figure 2. Evolution during the GB and WD phases of the same systems from Fig. 1, but starting at the end of the main sequence. The semimajor axis evolution is shown in the Prime 2 (upper left panel) and Integer (lower right panel) systems, and the eccentricity evolution is shown in the Fibonacci 2 (upper right panel) and Lazy Caterer (lower left panel) systems. The semimajor axial ratio between each pair of planets remains nearly fixed during stellar mass-loss, and the corresponding eccentricity increases are of the order of $10^{-3}$. Only the Integer sequence becomes unstable, and just only about 100 Myr after the GB-WD transition.
of 0.5 (Reimers 1975), and included the superwind phase along the AGB (Vassiliadis \& Wood 1993).

We integrated the resulting stellar evolutionary mass profile with each of the systems in Table 1 in three separate simulations: one with a RADAU tolerance of $10^{-10}$, another with a tolerance of $10^{-11}$, and a third with $10^{-12}$. These tolerances determine the variable timesteps by which the planets are propagated, and hence our collection of three simulations per architectural set-up provides some small statistical measure of the fate of these systems.

We integrated the systems from the end of the main-sequence phase for 10 Gyr . This time-span includes the approximately 1.4 Gyr duration traversing the GB phases, about 0.2 Myr of which covers the AGB phase. In all cases, the planets begin sufficiently far away from the star such that there is no danger of engulfment into the expanding giant star unless one of the planets is scattered inward due to an instability.

The majority of architectures from Table 1 remained stable for the duration of our simulations. In fact, gravitational instability appeared in only two of the chains: (i) the Fibonacci 1 chain and (ii) the Integer chain. For (i), all three simulations remained stable throughout the giant branch (GB) phases, and then became unstable approximately 2.0, 2.4, and 5.2 Gyr into the white dwarf (WD) phase. For (ii), the instability time-scale and trigger were qualitatively different: here the end of the AGB phase triggered instabilities quickly, just 87, 148, and 464 Myr after the star became a WD.

Unlike the instability depicted in Fig. 1, each post-main-sequence instability resulted in a destruction of the entire resonant chain. We
also checked for libration of the relevant resonant angles in our stable evolution and found that the first-order resonances (i.e. 2:1, $3: 2,4: 3$, etc.) remained intact for the entire 10 Gyr calculation. While the libration of the higher order resonant angles was lost during the GB phases in the majority of our simulations, in all stable cases the planets' period ratios remained within $\sim 0.05$ per cent of their initial values for the duration of the integration. We note that, in any case, it would not be possible to detect the libration of a distant pair of planets in a higher order resonance with current technology.

Schematically, some of the evolutions are shown in Fig. 2. The figure shows how the semimajor axial ratio of each planet pair remains nearly fixed during stellar mass-loss, and the orbital eccentricity increases are of the order of $10^{-3}$. These variations are natural and reflect the true nature of non-adiabatic stellar mass-loss on orbiting objects (Veras et al. 2011), although the primary trigger for instability on these spatial scales probably represents the change in stability boundaries due to the altered gravitational potential (Debes \& Sigurdsson 2002).

## 5 DISCUSSION

It is worth considering how confident we can be that our chosen resonant chains unequivocally signify the presence of an intelligence. As discussed earlier, chains of orbital resonances - especially firstorder resonances like $2: 1$ and $3: 2$ - are a natural outcome of
both the planet (e.g. Masset \& Snellgrove 2001) and satellite (e.g. Madeira, Izidoro \& Giuliatti Winter 2021) formation processes. Notable examples include the three inner Galilean moons, the K2138 system of close-in exoplanets (Christiansen et al. 2018), and the HR 8799 system of wide-orbit gas giants (Marois et al. 2010). It is also likely that the Solar system's giant planets were in a primordial resonant chain when they emerged from the gaseous disc (Morbidelli et al. 2007; Nesvorný \& Morbidelli 2012; Clement et al. 2021). The only examples of higher order resonances in the Solar system are in small body and satellite reservoirs. For instance, Neptune's moons Naiad and Thalassa are in a 73:69 resonance (Brozović et al. 2020), and trans-Neptunian objects have been detected in resonances as distant as 9:1 (Volk et al. 2018) and 11:1 (Clement \& Sheppard 2021). High-order resonances are difficult to confirm in exoplanet systems; they are expected as an outcome of dynamical instabilities but only at a very low rate and in systems with large orbital eccentricities (Raymond et al. 2008).

An effective SETI beacon would encode a mathematically meaningful sequence that includes multiple resonances more exotic than the common $2: 1$ and $3: 2$. Ideally, such a beacon would include at least one extremely high-order resonance (such as $11: 7$ or 13:8 in our Prime and Fibonacci chains) and avoid sequences that grow exponentially (thus making the distant bodies extremely difficult to detect). In this manner, a sequence like our Fibonacci chain that includes a pair of co-orbital (Trojan) planets would be extremely effective as a signpost of an alien civilization, as such configurations are not expected to be likely end states of planet formation. ${ }^{2}$ Given that the exact period ratios of planets can deviate from perfect integer ratios, it is important that an SETI beacon system be characterized by librating resonant angles (as in Fig. 1).

Might any currently known resonant chains qualify as possible SETI beacons? The best current candidate is the TRAPPIST-1 system, whose seven planets form an $8: 5,5: 3,3: 2,3: 2,4: 3,3: 2$ resonant chain (Gillon et al. 2017; Luger et al. 2017). If we require the innermost planet's orbital period to be an integer, this constitutes the following sequence from the innermost planet outward: 15 , $24,40,60,90,120,180$. While this sequence is of no obvious mathematical significance, a simpler sequence can be constructed from the outermost planet inward: $2,3,4,6,9,15,24$. This represents consecutive values in the sequence representing the number of compositions of odd numbers into primes of the form $x^{2}+y^{2}$. While this is not an integer sequence ${ }^{3}$ of note, we cannot rule out the possibility of an SETI beacon. However, it has been proposed that each pair of planets was initially captured into first-order (3:2 or 4:3) resonances (Papaloizou, Szuszkiewicz \& Terquem 2018). It is also worth noting that this sequence could be extended in either direction. While the existence of an additional, undetected interior planet in the system is unlikely, an outer sub-giant planet in a $2: 1$ resonance with TRAPPIST- 1 h cannot be ruled out by currently available radial velocity data (e.g. Boss et al. 2017). However, such a planet would have to be less than a few Earth masses so as to not overly excite the inner planets' eccentricities (Dencs \& Regály 2019) beyond the observed values (Quarles et al. 2017; Grimm et al. 2018).

While the energy required to modify the orbits of a collection of planetary-mass objects by several astronomical units is substantial, a civilization capable of exploiting the complete energy output of its host star (e.g. a Kardashev type II civilization; Kardashev 1964)

[^1]could acquire the necessary energy in a reasonable amount of time. Indeed, the orbital energy needed to modulate the semimajor axes of each of the six planets in our Prime 2 sequence by $\sim 10$ per cent is around $2.8 \times 10^{32} \mathrm{~J}$. Even if an advanced civilization around a 1.0 $\mathrm{M}_{\odot}$ star were only 10 per cent efficient at capturing solar energy (for instance, by utilizing a Dyson sphere), it would only require 2.3 yr to build up the necessary energy. While this timeline would be lengthened by around two orders of magnitude if the civilization happened to inhabit an early M-dwarf system, it is also possible that such a species would employ more passive means of orbital modification. Indeed, it is feasible to devise orbital trajectories for smaller, approximately asteroid-mass objects that slowly transfer energy between two planets in a system over longer periods of time, thus modifying the orbits of both planets in opposite directions with relatively minimal energy input to the three-body system (Korycansky et al. 2001).

We conclude by emphasizing the remarkable scientific value of resonant chain planetary systems. Multiresonant systems induce larger-amplitude transit-timing variations than their non-resonant counterparts (Agol et al. 2005), enabling precise mass determination (Agol et al. 2021) and constraining the planets' bulk compositions (Dorn et al. 2015) and bombardment histories (Raymond et al. 2022). As SETI beacons, resonant chains encode simple integer sequences that can be deciphered by stars able to see those systems edge-on (see Kaltenegger \& Faherty 2021). Finally, as we have shown, a carefully constructed resonant beacon can in principle remain dynamically stable for longer than the current age of the Universe.

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## DATA AVAILABILITY

The simulation inputs and results discussed in this paper are available upon reasonable request to the corresponding author.

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[^1]:    ${ }^{2}$ The Fibonacci sequence was indeed the one transmitted by intelligent aliens in Contact (Sagan 1985).
    ${ }^{3}$ See https://oeis.org/A287148.

