Local to continental scale coupled fire-atmosphere simulation of large industrial fire plume

³ Roberta Baggio^a, Jean Baptiste Filippi^a, Benjamin Truchot^b and Flavio T. Couto^c

⁴ ^aSPE—Sciences Pour l'Environnement, Université de Corse, CNRS, Campus Grimaldi, 20250 Corte, France

5 ^bINERIS, Parc Technologique ALATA, 60550 Verneuil-en-Halatte, France

6 CInstituto de Ciências da Terra - ICT (Polo de Évora) & Earth Remote Sensing Laboratory (EaRS Lab), Universidade de Évora, Rua Romão

7 Ramalho, 59, 7000-671 Évora, Portugal.

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ABSTRACT

Keywords: This paper presents an approach to reproduce small to large scale dispersion of smoke induced by 12 13 Coupled fire modelling industrial fires based on the atmospheric code Meso-NH coupled with the ForeFire surface fire Fire plume code. Meso-NH can solve explicitly micro to meso-scale meteorology, while ForeFire allows 14 Multi-scale smoke dispersion for an efficient parametrization of the injected heat and vapour fluxes. The combination of 15 Industrial fire these two models allows for a straightforward implementation of the heat source term and 16 Air pollution simultaneously accounts for its effect on local meteorology (as convection and latent heat), thus 17 Atmosphere leading to a prognostic representation of the developing plume. The context of the study imposed 18 19 Simulation to set model parameters that allow for faster than real time forecast along with a qualitatively realistic representation of the plume. Model evaluation in the vicinity of the fire is performed 20 by comparing results with a combustion CFD solver on an idealized situation, followed by a 21 mesh sensitivity study. Then, the proposed coupled method is applied to a real-case multi-scale 22 simulation of smoke dispersion from the local (city) to the continental scale in a twice faster than 23 real-time computation time. 24

²⁶ 1. Introduction

Many industrial sites, when exposed to uncontrolled fires, are susceptible to provoke immediate and acute impact 27 to the local communities and to heavily contaminate the surrounding environment. An uncontrolled combustion of 28 the stored chemicals may result in the production of heat and pollution creating a rising plume that can affect nearby 20 population immediately through the dispersion of acute toxic products. On longer time scales, spanning from days 30 to years following the accident, the accumulation of pollutants in the environment may result in chronic poisoning. 31 Improving both short and medium term predictions of this type of hazards can help to redefine valuable tools that 32 supports both the short-term emergency response and the post-crisis intervention to assess and circumscribe the 33 environmental damage. Such tools are also possible thanks to recent developments in numerical simulations and 34 increasingly powerful computational resources. However, there are still obstacles to a more systematic use of such 35 numerical simulations when coordinating the short and long-time emergency responses. Firstly, there is an high 36 level of uncertainty characterizing all the different scales involved in the process, especially when considering the 37 short-time scale involved in emergency response. In particular, a relevant forecast of the accident requires a good 38 estimation of chemical emissions (the source term) and of the dynamics of the heat release rate (HRR), including fire 39 suppression activity. Secondly, the time required to compute the desired solutions must be smaller than forecasting 40 time. Satisfying those constraints in an emergency setting imposes that several models and codes are based on many 41 simplifying assumptions. However, this may limit the reliability of the provided predictions, especially when one aims 42 to have an overview on all the involved scales, spanning from the immediate surroundings of the site of the accident, 43 up to the regional scale. A typical example of fast-computing, simplified modelling tools are Gaussian dispersion 44 models, such as the British software ADMS [CHH⁺94] or the American AERMOD [CPV⁺05]. Such models assume 45 a Gaussian distribution of the concentration of pollutants in both downwind and cross-wind directions, given by an 46 47 analytical solution of the transport equation for an homogeneous, steady flow and a constant point-wise source term, along with hypothesis on the characteristics of turbulent diffusion. This type of modelling provides a fast solution 48

Seggio_r@univ-corse.fr (R. Baggio); filippi@univ-corse.fr (J.B. Filippi); benjamin.truchot@ineris.fr (B. Truchot); fcouto@uevora.pt (F.T. Couto)

ORCID(s): 0000-0003-1531-0530 (R. Baggio); 0000-0002-6244-0648 (J.B. Filippi); 0000-0003-0175-2100 (F.T. Couto)

describing the diffusion of the pollutants on a medium range spanning from about 10 to 100 km, they can account 49 for terrain topography trough suitable parametrizations [Per92], but are known to give poor results when the external 50 wind velocities are low or complex [LMI+14]. Lagrangian models are another type of widely used dispersion models 51 accounting for more realistic wind fields which can be taken directly from ongoing meteorological forecasts (see for 52 instance the Micro Swift Spray software [TBO⁺07]). The advective term is determined by the imposed wind field while 53 diffusion is given by a stochastic model of turbulence. With respect to Gaussian models, dispersion models are more 54 expensive but allow to consider realistic meteorological fields and complex terrains. A major drawback characterizing 55 all the above methods is that the convective effect of the heat released by the fire on the local atmosphere is not explicitly 56 resolved. Indeed, industrial fires can release in the atmosphere considerable power (with peaks of several GWs) which 57 modifies locally the wind profile, causing an increase of the vertical lift and of the turbulent flow in correspondence of 58 the fire site. This fire-to-atmosphere interaction is therefore very important for describing properly the height reached 59 by the fire plume and the corresponding diffusion of pollutants. In many cases, the effect of the vertical lift on the rise 60 of the plume is modeled by simple empirical formulas which give an estimate for the height of injection [PWFM16]. 61

As a numerically intense alternative, a detailed description of the combustion process and its influence on the 62 local environment can be evaluated using Computational Fluid Dynamics (CFD) codes that account for combustion. 63 A typical example is the Fire Dynamics Simulator code FDS, developed by the joint efforts of the National Institute of 64 Standards and Technology (NIST) and the VTT Technical Research Centre of Finland [MHM⁺13]. Not surprisingly, 65 the forecast time and spatial dimension are limited to some hundreds of meters maximum, due to the implied 66 computational resources. At the opposite end of the spectrum of involved length scales, spanning from tens of 67 kilometers up to the synoptic scale, numerical weather prediction models (NWP), solving explicitly the equations 68 governing the physics of the atmosphere, can be used to state the diffusion and the interaction of the releasing pollutants 69 with other meteorological and chemical variables. This type of simulations is important in the case of big industrial 70 accidents. For instance, the French code CHIMERE [MMK⁺17] has been used for evaluating the impact of the NL 71 Logistique / Lubrizol industrial fire event in 2019 [RTM⁺21]. However, this type of models is too coarse to represent 72 the effects closer to the accident, in particular they do not solve the convection over a potential fire heat source. At 73 the intermediate (meso-scale) range, several studies used the Weather Research Forecasting model (WRF) [SKD⁺05] 74 coupled with a city sub-grid pollutant dispersion model [NTY21] or with a Lagrangian particle model [KSR+22] albeit 75 none explicitly resolving convection at the highest resolution. Wildfire is a specific application for which such micro 76 to meso-scale approach has been already applied successfully from small to large scale events and with models such as 77 WRF/SFire [KMF⁺19] or Meso-NH/ForeFire [FBP⁺11, FBML18]. An example is the simulation performed to predict 78 the evolution of the Aullene Mediterranean wildfire [FBML18], whose computational time did not exceeds the time 79 taken by the dynamics of the fire spreading. 80

In this paper, we propose to develop this latest wildfire approach to model small (near the flame) to large 81 (continental) scale smoke dispersion in case of major industrial fires. At first, a dynamical parametrization of the heat 82 release profile is modeled in the surface fire code ForeFire to account for the typical combustion in such accident. Then, 83 suitable settings, as well as an optimized mesh, are configured for the Meso-NH atmospheric code. For simplicity, we 84 do not inject source terms of chemical species and do not consider the physics of aereosols, despite the fact that suitable 85 models exists in the current versions of Meso-NH [TCS⁺03, TCC⁺05, TGGP06]. We keep track of the evolution of 86 the smoke cloud trough the injection of a passive scalar which is emitted at constant rate from the burning domain. 87 Despite its simplicity, such approach allows for a straightforward identification of the effects of the injected heat and 88 vapour on the local meteorological variables, and interactions with the different water phases are observed. The chosen 89 discretization is based on three nested domains, whose resolution varies from 80 m up to 2000 m. This choice allows to 90 model fine meteorological effects near the accident site while simultaneously solving the consequences of the fire on 91 larger scales. To choose the most suitable discretization at the smallest scale, we performed a comparison with a the 92 combustion computational fluid dynamics FDS code on an idealized scenario. Then, the modelling tool is applied to 93 the study of a real industrial fire and show that, despite the many length scales involved, it is able to model large scale 94 smoke dispersion. 95

The paper is structured as follows: The first section presents schematically the two coupled models. Description is limited to the main aspects of both and provide relevant bibliography for more detailed explanations. The second section explores a first simple, small scale application, where we compare mean features of a numerical fire plume simulated respectively with the proposed model and the FDS solver. In the same section, based on this simple setting, results of different Meso-NH grids are compared showing that it is possible to chose a relatively coarse discretization without degrading significantly the representation of the plume in the near filed. The third section demonstrates the feasibility of applying the Forefire/Meso-NH coupled setting to the simulation of a real case with the added advantage
of resolving multiple scales simultaneously. The NL Logistique / Lubrizol inductrial fire event in 2019 [RTM⁺21] is
selected as the study scenario for this purpose. General discussion and perspectives concludes the paper.

¹⁰⁶ 2. The coupled fire atmosphere model ForeFire/Meso-NH

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The French code Meso-NH (Mesoscale Non Hydrostatic) is a research mesoscale NWP model designed to simulate 107 the relevant physics on a wide range of scales spanning from tens of meters to several hundreds of kilometers 108 [LCM⁺18]. Developed by the joint effort of the Laboratoire de Aerologie and Meteo-France, it is a non-hydrostatic, 109 limited area-model based on the anelastic approximation. The solved equations are the conservation of momentum, 110 mass, water substances and user defined scalar variables, along with a suitably derived thermodynamic equation. The 111 anelastic approximation implies that sound waves are filtered out, thus avoiding the necessity of exceedingly fine time 112 steps in the time integration scheme. The turbulence model used for the parametrization of the unresolved eddies is the 113 scheme proposed by Redespelger and Sommeria [RS81, RS86] (see [CBR00] for details on the implementation). The 114 scheme uses a prognostic equation for the Turbulent Kinetic Energy (TKE), closed by adequate hypothesis on the form 115 of second order turbulent fluxes and the mixing length. The scheme is available both in a simplified 1D version, where 116 the horizontal components of turbulence fluxes are neglected) and in a fully 3D implementation, that is relevant in the 117 smaller scales of this work (lower than $\approx 1 \text{ km}$). Suitable modules allow to study a wide range of physical phenomena, 118 including, among others, the climatic impact of volcanic eruptions, and the electric activity during thunderstorms (see 119 120 Table 4 of [LCM⁺18] for a detailed overview). At the smaller length scales, the code runs in Large Eddy Simulation (LES) mode to allow an improved representation of turbulence. The meteorological fields imposed at the boundary 121 are taken directly from operational forecast NWP databases (examples of compatible data are the ones obtained by the 122 french model AROME (used in this study) or the European ECMWF). Moreover, the Meso-NH modelling tool allows 123 to account automatically for the local orography and is coupled with the surface model SURFEX. The latter include a 124 physical representation of soil-to-atmosphere heat, water and general exchanges for different surface types (land, town, 125 sea and inland water) [MLMM⁺13]. 126

ForeFire was originally developed to simulate the propagation of wildfires [FBG14]. Its coupling with Meso-NH gives 127 the possibility to inject in the atmosphere arbitrary, time dependent, heat, vapour or any chemical species fluxes, while 128 solving explicitly the convective motions resulting from the interaction with the local atmosphere. The effects of the 129 injected fluxes on the calculated wind fields and on the other meteorological variables are considered explicitly. The 130 flexibility of the ForeFire code allows for a straightforward implementation of ad-hoc routines, suitably designed for 131 the selected release models. Different functions and parameters can be assigned to separate regions of the burning 132 domain, and the dynamics of fire spread can be represented by specifying a suitable *Time of Arrival* Matrix (TAM), 133 which associates a burning starting time at every point in the domain with user defined resolution. Alternatively, one 134 may use very simple maps but specify arbitrary time varying fluxes. Given the flexibility of this time-space tuning, the 135 resulting forcing term can be quite complex and is easily adjusted to assigned dynamic profiles. Once the source terms 136 and the TAM have been specified, the resulting fluxes enter the coupled simulations under the form of time dependent 137 ground forcing terms. We remark that in the case of wildfire spread, a feedback coupling may be necessary, that is, 138 the burning region is adjusted on the basis of atmospheric motions, which, themselves, are driven by the developing 139 fire. This two-way coupling is not relevant for the condition considered here, where the perimeter of the fire hazard is 140 limited by the one of the burning facilities. The most important effects are given by the elevation and the following 141 evolution of the developing plume. 142

Thanks to the multiscale setting allowed by Meso-NH, we study the diffusion of the plume on multiple length scales simultaneously by using simulations on nested grids. Each grid is characterized by its own with resolution and simulation parameters, in order to model coherently the underlying physics. Before discussing in detail the set-up we used for the multiscale NL Logistique / Lubrizol simulation, we present a simple study concerning the smallest scales of interest. These preliminary tests drove the choice of the parameters used in the finer grids.

3. Coupled simulations results

¹⁴⁹ **3.1.** Plume evolution in a simple setting

¹⁵⁰ We begin our study on the effectiveness of the ForeFire Meso-NH coupling by studying the development of a fire ¹⁵¹ plume under idealized atmospheric conditions at scales between $\approx 10 m$ and $\approx 100 m$. In the following subsection, we

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Figure 1: FDS simulation results. On the left: the $s \le 10^{-3}$ three dimensional subset is presented, colors indicates three different values of temperature. On the right we show the vertical velocity along a *X*-parallel section along the plane containing w_{max} (on top) the non-dimensional *s* on an analogous section (bottom), in this latter graph contours are in log scale to improve visualization.

draw a qualitative comparison of the plume evolutions calculated by FDS and the one modelled by Meso-NH/ForeFire. 152 We consider a Meso-NH grid that is much coarser than the one used in the FDS model, despite the fact that Meso-NH 153 is capable to handle finer grids (see for instance [CRLM21]). This choice is motivated by the fact that we are interested 154 in providing some confidence in the capability of Meso-NH in solving properly the evolution of the plume with grids 155 that remain compatible with relatively fast computational times. Indeed, the main purpose of this section is to illustrate 156 as a relatively coarse Meso-NH grid can model qualitatively well the plume dynamics, once a good estimate of the 157 HRR has been provided from previous modelling phases. We remark that detailed comparisons between an anelastic 158 atmospheric model analogous to Meso-NH and FDS has been already discussed in the literature [SJK⁺06]. We limit 159 our analysis on the main features of the plume, that is, its vertical velocity and the diffusion of the smoke tracer. In 160 the subsection that follows the representation of the same plume is considered on different Meso-NH grids, which are 161 made coarser and coarser. The goal is to select a grid that will not filter excessively the main structures inside the 162 plume. The following examples are based on the simple case of a quadratically increasing heat flux, distributed over 163 an area of 30 m^2 and reaching a maximum of 3 GW after t = 150 s. The flux then continues steadily. We are interested 164 in the plume development under a time steady, X - oriented wind profile, representative of a neutral atmosphere and 165 fixed to a 5D Pasquill class, the velocity then grows from $0 ms^{-1}$ at ground level up to $5.0 ms^{-1}$ at the reference height 166 Z = 10 m, it still grow after based on a power law. 167

3.2. small scale Meso-NH and FDS plume representations

To be coherent with the type of modelling adopted by both methods, the evolution of the plume needs to be analyzed at the scales of the order of some hundreds of meters surrounding the fire. In the coupled model, smoke is described as a passive scalar which is emitted at constant rate of $1s^{-1}m^{-2}$ from the heat source where the fire is active, at the resolution of the given TAM, and is then advected by the convective flow. To make a more relevant comparison with the "Soot" variable in FDS, we compare these two smoke tracers using their non-dimensional form, obtained by dividing them by the maximum reached value within the domain.

¹⁷⁵ We use a Meso-NH grid characterized by an uniform horizontal discretization $\Delta X = \Delta Y = 40 m$, and a vertically ¹⁷⁶ varying grid with ground value $\Delta Z_0 = 10 m$. This grid is progressively enlarged in correspondence of each subsequent ¹⁷⁷ vertical element with constant ratio (here 4%), as typically done in atmospheric models. The simulation is started in



Figure 2: On the left we show w contours for the averaged FDS simulation (top) and Meso-NH (bottom). On the right we show $w_{max}(X)$ and its Z(X) altitude in the two models (FDS results are not averaged).



Figure 3: The altitude of s_{max} (continuous line), along with the $s \ge 10^{-3}$ threshold (dotted) for the two models.

¹⁷⁸ idealized, stable atmospheric conditions, the assigned wind profile is forced at the boundary with resolution given by the ¹⁷⁹ vertical discretzation. The FDS simulation is performed on a much finer grid, with uniform $\Delta X = \Delta Y = \Delta Z = 4 m$. ¹⁸⁰ The simulated domain is limited to $600 m \times 600 m$ horizontally, and to 300 m vertically, to keep low the computational ¹⁸¹ times. The simulated Meso-NH domain is of about one order of magnitude larger ($8 km \times 8 km \times 3 km$), to lower the ¹⁸² influence of lateral boundary effects. Moreover, such dimensions are comparable to the scales implied in the more ¹⁸³ complex models showed later on.

The FDS results are schematized on Figure 1, showing the vertical wind velocity w and the distribution of the scaled soot density s. The vertical sections are taken along planes parallel to X and such that the domain is cut in correspondence of the maximum w and s values respectively. These profiles are taken at time 300 s after ignition, at which the flame has stabilized. A global view of the plume is illustrated on the left of Figure 1, where only points such that $s \ge 10^{-3}$ are displayed. One observes clearly the slope imposed by the applied wind field, with the plume leaving the domain at about 225 m from the emission point. In order to compare the w velocity field with the one obtained on the coarser Meso-NH plume, the vertical velocity is averaged on a grid analogous to the one used in the Meso-NH model. Results are shown on the plot in the left of Figure 2. White contours corresponds to the two thresholds $w = 0.5, 2 m s^{-1}$.

The coarser grid used on Meso-NH is here clearly incapable to solve the highly turbulent motions, which makes the velocity field smoother. However, the two cases show a similar maximum value for w, and the two modelled plumes display comparable slopes. The global agreement between the two simulations is confirmed by the right plot in Figure 2, where the evolution of the maximum value of w with x of the (non-averaged) FDS and Meso-NH are displayed. More specifically, on the top right graph we illustrate the maximum w value encountered along the X direction, together with the altitude Z at which such value is encountered (top right graph).

In Figure 3, we consider the evolution of non-dimensional smoke tracer s. The continuous line shows the height 199 at which, in correspondence of each X, the maximum value s is encountered, while the dotted lines are the $s \ge 10^{-3}$ 200 boundaries. On this smoke diagnostic, it is possible to notice a qualitative overall agreement between the two plume 201 shapes. We conclude that, despite the coarser grid, the Meso-NH model can capture the global properties of the 202 developing plume, well represented by the trend of the vertical component of the velocity and the associated diffusion 203 of the smoke term. Both plumes show comparable peak values of w and bend similarly under the action of the imposed 204 horizontal wind. This aspect is fundamental to explicitly predict smoke dispersion at a larger scale since it governs the 205 initial height of the fire plume without the need of a parameterized plume-rise model. 206

3.3. Meso-NH plume representation on varying grids

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In this section, we focus on different Meso-NH representations of the same simple setting discussed before. While 208 a complete grid sensitivity analysis is beyond the scope of this work, it is still important to estimate how the grid 209 resolutions affects the plume features, considering the strong impact on the computational time. In particular, the 210 $\Delta X = \Delta Y = 40 m$ discretization previously shown is still relatively costly when one aims to model the plume spreading 211 beyond the regional scale, even within a suitably designed embedded model. Moreover, the $\Delta Z_0 = 10 m$ grid becomes 212 problematic nearby the heat releasing source, since the ensuing violent convective motions imposes a strong reduction 213 of the implied time step. We thus consider three different grids with increasing discretization parameters $\Delta X = \Delta Y$, 214 ΔZ_0 (see labels in Figures 4 and 5), where the finer grid is the same as the one discussed in the subsection above. The 215 ratio controlling the growing of ΔZ with altitude is 4% in all models. We keep ΔZ_0 relatively low with respect of the 216 step used in the horizontal direction, to avoid an excessive deterioration in the representation of the plume. All three 217 simulations consider an horizontal domain of $8 km \times 8 km$ with elevation $\approx 3 km$. In Figure 4 we show the $s \ge 10^{-3}m$ 218 representation of the three plumes after the heat flux has become stable (t > 300s), while on Figure 5 we illustrate 219 a vertical cross-section at about 3 km from the fire source, showing the development of a vortex structure. Contours 220 are proportional to w, (a cutoff between -1 and $1 ms^{-1}$ has been set to improve the vortex visualization), while the 221 arrows show the direction of the in-plane component of the wind. We remark that all the three models present the 222 same qualitative features, with a rising plume which splits with the development of a symmetric structure of rotating 223 vortexes. While the coarser grid leads to a degraded plume representation, the choice $\Delta X = \Delta Y = 80 m$, $\Delta Z_0 = 30 m$ 224 is here a reasonable compromise. This is also confirmed by the graphs showed in Figure 6 (analogous to the ones 225 discussed before when comparing Meso-NH with FDS). On the right, we illustrate the altitude of $w_{max}(X)$ along with 226 the observed value. On the left we report a schematic representation of the s plumes. The continuous line indicates the 227 value for which s is maximum, while the dotted lines enclose the $s \le 10^{-3}$ region. 228

The previous comparison with the FDS simulation shows the ability of Meso-NH to model the initial plume behaviour in the near field, while the present analysis indicates that decreasing resolution does not strongly modify its main features. As one could expect, flow and smoke finest structures are filtered out due to the lower resolution but the upflow mechanism remains well represented and the initial height of the fire plume is in good accordance between the different cases.



Figure 4: Overview of the three $s \ge 10^{-3}$ plumes (from left to right, coarser to finer grid).

4. Towards real case simulations

In this section, we show the application of the coupled ForeFire/Meso-NH tool to the simulation of the industrial fire which affected the French city of Rouen the 26^{th} September 2019. Due to the chemicals involved in the combustion, firefight actions took almost a day to extinguish it completely, and the fire released in the atmosphere an estimated power of more than 83 *GW h*. The resulting fire plume traveled thousands of kilometers and raised air quality concerns. This choice is motivated by the fact that the Rouen accident is a typical example for which a full-scale atmospheric simulation may give an important support to the organization of the emergency response, and offers a complementary, multiscale tool to the assessment of the resulting environmental damage.

243 **4.1. Main Simulation Settings**

In order to resolve properly all the length scales involved in the disaster, we exploit the nesting capability of the 244 Meso-NH software and we subdivide the simulated area in three nested domains with different resolutions (see Figure 245 7). The horizontal grid size is 2000 m for the outer domain covering a total area of $600 \, km \times 600 \, km$. The inner 246 computational grids have horizontal discretizations of respectively 400 and 80 m, over areas of 120 km \times 120 km, and 247 $24 \text{ km} \times 24 \text{ km}$ for the innermost model. The vertical resolution is shared by all the nested domains, with 50 levels up 248 to 20 km and a first level above the ground at 30-m height. The time step is 10 s for the outermost model and decreases 249 to 2 s and 0.5 s for the finer models. The grid resolution and model parameterisation of the finest domain corresponds 250 to the one used in the idealized configuration analyzed in the previous subsection, and represents here a reasonable 251 compromise between the computational efficiency and the accuracy of the plume representation. 252

Initial and lateral boundary conditions are provided by the French AROME model (1.3-km horizontal resolution)
at 0000 UTC, with updates of the coupling files every 3 h. The simulation with the coarsest resolution began on 26
September 2019 at 0000 UTC, with a progressive downscaling up to the finest resolution beginning at 0020 UTC.

The burning region determining the site of ignition has been placed in correspondence of the estimated ignition site (see the inset in the right of Figure Figure 7). The total burning area has been set up to $57600 m^2$ (corresponding to the bottom of 9 ground elements), which is way more extended than the estimated one ($\approx 20000 m^2$). This modelling choice was motivated by the fact that it allows to release the total heat flux while avoiding excessive gradients in proximity of the heat source. This did not affect significantly the physical description, and at the same time allowed to use larger time steps, thus reducing the implied computational time. The assigned burned region is then subdivided in

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Figure 5: Y-parallel section (illustrated ion top) showing the vortex structure accompanying the plume separation (from left to right, coarser to finer).



Figure 6: On the left we show the height of $w_{max}(X)$ (top) with the corresponding value (bottom) for the three considered grids. On the right the three $s \ge 10^3$ plume thresholds (dotted) along with the height of s_{max} are showed.

two subdomains of different areas, characterized by the same flux model, but with different stating times t_0 (see left of Figure 8). This simple combination allowed us to obtain the HRR dynamics illustrated on the right of Figure 8, which is coherent with the estimated dynamics of the fire (including a roof collapse at T+3). The vapour follows an analogous profile and the total quantity released is coherent with the calculated $\approx 10^4$ t of mineral oil which burned in the fire (the parameters used in the flux model are: $\phi_{peak} = 250 \, kW m^{-2}$, $v_{peak} = 0.002 \, kg s^{-1} m^{-2}$, $t_1 = 1800 \, s$,



Figure 7: On the left we show the three used nested domains along with the respective discretizations. On the right we show in greater detail the smaller of the nested domain. The inset show the set burning region



Figure 8: On the left, the implemented burning map along with the time assigned for the arrival of the fire (spread is considered instantaneous within the area). The type of flux function associated with heat and vapor flux is also illustrated. On the right: the corresponding HRR.

 $t_2 = 18000 \text{ s and } t_3 = 57600 \text{ s}$). We simulated only the fist 10 *h* of the fire, time in which the developing plume reaches the boundaries of the outermost domain. The total energy released in this laps of time is $\approx 80 \text{ GW } h$.

4.2. Dynamical plume structure

One of the aims of the current study is to use numerical simulations to better understand fire and atmosphere interactions and directly solve the updraft caused by the presence of the fire.

In Figure 9a) we present two views of the $s \ge 10^{-5}$ and $s \ge 10^{-3}$ plume configurations when the heat release has reached its peak. As a consequence of the heat injection, the smoke plume is lifted significantly from the ground, raising up to above 1 km over the site of injection and evolving at similar altitudes. The strength of such lift is important as the altitude influences the direction of the plume in the zone surrounding the accident. The large scale evolution of plume is summarized on Figure 9b), which illustrates the $s \ge 10^{-5}$ plume at the end of the simulation, after 10 h from the ignition. On large distances, the plume rises above 3 km, height that may influence large scale transports of pollutants. White contours illustrate different s thresholds at the ground level. While the overall plume simulation is

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Figure 9: On the left: Plume structure on the smallest of the nested domains. the $s \ge 10^{-5}$ plume is showed using transparent blue while $s \ge 10^{-3}$ is illustrated with solid black. On the right: Evolution of the plume on large scales (colorscale is proportional to the plume altitude). White contours correspond to different *s* thresholds at the ground.

in agreement with previous simulations and observations [RTM⁺21], a complete evaluation of simulation results on 279 this specific case is beyond the scope of this paper. The main objective of this work is to demonstrate the ability of 280 the coupled approach to resolve multi-scale plume simulation in a reasonable computational time. An advantage of 281 Meso-NH with respect of other modelling tools is that it calculates explicitly the interaction between the injected heat 282 and vapour fluxes and the different water phases present in the atmosphere. We do not consider here the physics and 283 the chemistry of aereosols, which will be necessary to evaluate more properly the effect of the smoke cloud on the local 284 precipitation (with a strong impact on computation). However, even in this simpler setting, it is still possible to have an 285 idea of the regions impacted by fire-related precipitations. In Figure 10 we illustrate the difference between the coupled 286 simulation and a control simulation (where no heat and vapour fluxes were injected) in terms of total accumulated rain 287 precipitations at the ground during the 10 hours of the simulated fire, on the intermediate $120 \, km \times 120 \, km$ nested 288 domain. It is possible to observe the presence of fire induced precipitations on a region below the one covered by the 289 plume. It is noteworthy that such a result gives an indication about a possible environmental impact coming from the 290 deposition of pollutants. 291

4.3. Computational time

Computations were performed on the Orsu computer at the University of Corsica (Intel Xeon) composed of 20 293 nodes of 52 cores each. Faster than real-time (and operational compatible) computation was reached here by running 294 on 625 processors dispatched in 15 nodes (780 cores mobilized). In this case, it was possible to obtain a 0.5 calculation 295 ratio with 10 hours of forecast available in 5 hours of computation. This ratio do not take into the generation and 296 interpolation of boundary conditions that may be already available at the time of alert but would otherwise requires an 297 additional hour before the simulation can be started (computational time being highly dependent of the resolution of 298 the data used for these boundary conditions). This 0.5 ratio may limit operational uses, with the first really predictive 299 output (T+4 forecast) that will technically be available only 3 hours after alert (alert T+3). Nevertheless, as the first 300 output (forecast T+0) may be available in the first two hours after alert, it may still be useful early, to get an overall 301 representation of the phenomenon. Moreover, in case of a contained industrial fire, these 3 hours can also be in line 302 with the time before the breakout of the fire containment (such as a roof collapsing) that is likely to drastically change 303 plume structure heat release. 304



Figure 10: Surplus in total accumulated precipitation (mm) after 10 h of simulation. Landscape features are shown in grey in the background.

5. Conclusions and perspectives

In this paper we applied the atmospheric code Meso-NH coupled with the ForeFire surface fire code to the simulation of large industrial fires. The implied Meso-NH model solves micro to meso-scale meteorology, including convection, turbulence and clouds, in nested domains configurations with resolutions spanning from 80 meters to 2km. Combustion is resolved here as a sub-mesh fluxes, injected in the atmospheric fields using the ForeFire model, that computes heat and vapour release at 20 meters resolution using user-defined behavioural functions. Ideal

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simulation cases were used to compare to the more detailed, and computationally intensive, FDS combustion code. 311 The comparison shows that the coupled Meso-NH/Forefire simulates coherently the plume convective behaviour near 312 the heat source, at least for heat fluxes of the order of some GWs, which typically characterize industrial fires. 313 The performed real-case simulation illustrates the capability of the proposed MesoNH/ForeFire coupled approach of 31/ simulating major industrial accidents, although with some arguable simplifications in the implied physical processes 315 (as for instance the absence of particle injection), and the relatively raw parametrizations of the source terms compared 316 to the complexity of the accident. This multi scale simulation allowed to explore the dynamical structure of the 317 surrounding atmosphere, the strong convection induced by the fire and the related plume lifting, and the interaction 318 with air moisture, creating a rain anomaly under the plume some kilometers away from the source. Although only 319 tested on one real case example, models have been configured, along with grid geometry to be generic for any large 320 industrial fire incident of the same scale. In terms of computational time, the simulation reached a 0.5 ratio (5 hours of 321 computation for a 10 hour forecast) that is therefore predictive, but not instantaneous as simplified tools can provide. 322

The major point here was to demonstrate the feasibility of such computation with operational forecast constraints, 323 as well as build the platform and configuration to simulate large industrial fires, from local fire release near the source 324 to continental scale smoke transport. Future work is directed to evaluate these specific simulation results by comparing 325 it with in situ observations and existing numerical simulations of the same accident [RTM⁺21]. The proposed approach 326 can be improved by accounting for aerosols and their chemistry, which are known to play an important in precipitation 327 and pollution. On the one hand, the simplified approach used here (a passive scalar) may not be appropriate for precise 328 quantitative estimates of impacts on rain and air quality. On the other hand, the implementation of aereosols physics 329 requires as input an accurate source term with information enough about the characteristics of the combustion products 330 (initial size of the aerosols, gas composition including condensable ones), which may require additional analysis. 331 Finally, the current method could be improved by introducing an automatic procedure to generate the heat release 332 functions or profiles based on the characteristics of the industrial facilities and combustible products involved in the 333 accident. 334

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