

# Individual heme a and heme a3 contributions to the Soret absorption spectrum of the reduced bovine cytochrome c oxidase

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1	Individual heme <i>a</i> and heme <i>a</i> <sub>3</sub> contributions to the Soret
2	absorption spectrum of the reduced bovine cytochrome <i>c</i> oxidase
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# **Graphical abstract:**





# 45 Abstract

Bovine cytochrome c oxidase (CcO) contains two hemes, a and a<sub>3</sub>, chemically identical 46 but differing in coordination and spin state. The Soret absorption band of reduced *aa*<sub>3</sub>-type 47 cytochrome *c* oxidase consists of overlapping bands of the hemes  $a^{2+}$  and  $a_{3}^{2+}$ . It shows a 48 peak at ~444 nm and a distinct shoulder at ~425 nm. However, attribution of individual 49 spectral lineshapes to hemes  $a^{2+}$  and  $a_{3}^{2+}$  in the Soret is controversial. In the present work, 50 we characterized spectral contributions of hemes  $a^{2+}$  and  $a_{3}^{2+}$  using two approaches. First, we 51 reconstructed bovine CcO heme  $a^{2+}$  spectrum using a selective Ca<sup>2+</sup>-induced spectral shift of 52 the heme  $a^{2+}$ . Second, we investigated photobleaching of the reduced *Thermus thermophilus* 53 ba3- and bovine aa3-oxidases in the Soret induced by femtosecond laser pulses in the Q-54 band. The resolved spectra show splitting of the electronic B<sub>0x</sub> -, B<sub>0y</sub>-transitions of both 55 reduced hemes. The heme  $a^{2+}$  spectrum is shifted to the red relative to heme  $a_{3}^{2+}$  spectrum. 56 The ~425 nm shoulder is mostly attributed to heme  $a_3^{2+}$ . 57 58 **Keywords:** Cytochrome *c* oxidase, absorption spectra, femtosecond pump-probe 59 60 spectroscopy. 61 Abbreviations: 62 CcO, cytochrome c oxidase 63 FWHH, full width at half-height 64 MCD, magnetic circular dichroism 65 SMR, succinate: menaquinone reductase 66 67 **Highlights:** 68 • Bovine CcO  $a^{2+}$  Soret band is a split electronic 447.5-nm peak with a minimal 69 short-wavelength satellite 70 •  $a^{2+} \alpha$ -band is also split 71 •  $a_{3^{2+}}$  Soret band is a split electronic 441-nm peak with a short-wavelength satellite 72 Ratio of integrals under  $a^{2+}$  and  $a_{3}^{2+}$  Soret bands is approximately 1 to 2 73 74

## 75 **1. Introduction**

In mitochondrial and some bacterial respiratory chains,  $aa_3$ -type cytochrome *c* oxidase (CcO) catalyzes oxidation of ferrous cytochrome *c* by oxygen, converting redox potential energy to the transmembrane difference of proton electrochemical potentials,  $\Delta \mu H^+$ . The exergonic electron transfer coupled to proton translocation across the membrane [1] via proton channels [2–6] is mediated by two A-type hemes, low-spin 6-coordinate *a* and highspin 5-coordinate  $a_3$ , and two copper centers: binuclear Cu<sub>A</sub> and mononuclear Cu<sub>B</sub> [5,7–9].

Absorption spectra of hemes a and  $a_3$ , sensitive to their redox and spin state, are 82 routinely used to monitor the reaction. Fig. 1 shows the absorption spectrum of dithionite-83 reduced bovine CcO. The heme absorption spectra (350-700 nm, Fig. 1a) display empirically 84 85 defined  $\alpha$ -,  $\beta$ - and y- (Soret) bands corresponding to the electronic and vibronic transitions from the same ground state to increasingly higher-energy excited states (Fig. 1b). Platt's 86 convention [10] denotes the Soret band as B-band and groups the α- and β-bands into the Q-87 ("visible") band. The  $\alpha$ - and  $\gamma$ -bands comprise two perpendicular electronic transitions each 88  $(Q_{0x}, Q_{0y} \text{ and } B_{0x}, B_{0y}, \text{ respectively})$  oriented in the porphyrin plane (**Fig. 1c**). Each of the four 89 purely electronic transitions may have a vibronic satellite due to transition to a non-zero 90 vibrational state (for example,  $B_{1y}$ , etc). The  $\beta$ -band comprises vibronic satellites of the  $\alpha$ -91 band. 92

While heme *a* absorption dominates the Q-band, overlapping heme *a* and  $a_3$  bands with 93 high and similar intensities challenge spectral interpretation in the Soret. The Soret band of 94 reduced bovine CcO shows a main peak at 443-444 nm accompanied by a shoulder at ~425 95 nm (Fig. 1a). Our current knowledge of individual absolute spectra of hemes  $a^{2+}$  and  $a_{3}^{2+}$ 96 relies on Vanneste's indirect derivation [11] based on variable-valence vs unliganded state 97 spectral comparison referenced to the action spectrum of CO photodissociation from yeast 98 CcO as the absolute spectrum of CO-bound heme  $a_{3^{2+}}$ . A 444-nm peak/~428-nm shoulder 99 pair and a symmetric 442.5-nm peak (with a weak shoulder at ~420 nm) are therein 100 respectively attributed to heme  $a^{2+}$  and heme  $a_{3}^{2+}$  (Fig. 5 in ref. [11] and Fig. A.1); these 101 spectra are included in the textbooks and monographs (e.g. [12]). However, bovine CcO 102 reduction kinetics analysis led Orii [13] to essentially opposite assignment of a featureless 103 ~445-nm peak and the asymmetric band with a short-wavelength shoulder at ~428 nm to 104 hemes  $a^{2+}$  and  $a_{3}^{2+}$ , respectively (see Fig. A.1). 105

<sup>106</sup> In this work, we resolve individual absolute absorption spectra of the reduced bovine <sup>107</sup> CcO hemes *a* and *a*<sub>3</sub> using two yet different approaches. First, we reconstructed the heme  $a^{2+}$  spectrum from Ca<sup>2+</sup>-induced spectral shift. Second, we selectively excited hemes  $a^{2+}$  and  $a_{3}^{2+}$ with femtosecond laser pulses and measured photobleaching spectra of the two hemes. Both approaches showed B<sub>0x</sub>/B<sub>0y</sub>-splitting in both hemes, heme  $a^{2+}$  transitions red-shifted relative to heme  $a_{3}^{2+}$  transitions, and the ~425 nm shoulder attributed mostly to heme  $a_{3}^{2+}$  spectrum.





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Fig. 1. Absorption bands of reduced cytochrome c oxidase. (a) experimental

absorption spectrum of "activated" (see 2.1) CcO (1.7 µM) bovine CcO reduced by dithionite.

(*b*) electronic ( $Q_{0x}$ ,  $Q_{0y}$  and  $B_{0x}$ ,  $B_{0y}$ ) and first vibronic ( $Q_{1x}$ ,  $Q_{1y}$  and  $B_{1x}$ ,  $B_{1y}$ ) transitions from

the ground state S<sub>0</sub> corresponding to the  $\alpha$ -,  $\beta$ - and  $\gamma$ - absorption bands for single hemes. (*c*)

orientation of the x,y-transitions in the plane of porphyrin ring of heme A.

#### 119 **2. Materials and methods**

2.1. Reagents and preparations. aa<sub>3</sub>-type CcO was isolated from bovine heart 120 mitochondria according to [14] modified as described in [15] and dissolved in 25-50 mM 121 HEPES containing 0.05% DM. The concentration of the enzyme was determined from 122 "dithionite-reduced *minus* oxidized" difference absorption spectra using an extinction 123 124 coefficient  $\Delta \epsilon_{605-630}$  of 27 mM<sup>-1</sup>cm<sup>-1</sup>. Before KCN binding, CcO was "activated" as follows. Argon-deoxygenized ~30 µM CcO in 100mM HEPES/tris pH 8.0 was reduced by a small 125 amount of dithionite in the presence of 20  $\mu$ M of hexaammineruthenium (Ru(NH<sub>3</sub>)<sub>6</sub><sup>3+</sup>; RuAm) 126 for 30 minutes until full reduction. Dithionite and RuAm were subsequently removed using 127 centrifugal filter devices (Amicon Ultra-4, Millipore). ba<sub>3</sub>-type CcO purified from Thermus 128 thermophilus was prepared as described [16]. The concentration of the ba<sub>3</sub> CcO was 129 determined using extinction coefficient  $\Delta \varepsilon_{560-590} = 26 \text{ mM}^{-1} \text{ cm}^{-1}$  for the absolute spectrum of 130 the fully reduced oxidase [17]. Succinate-menaquinone reductase preparation was a 131 generous gift from Prof. Lars Hederstedt (Lund University, Sweden). The concentration of the 132 enzyme was determined from "dithionite-reduced minus oxidized" difference absorption 133 spectra using an extinction coefficient  $\Delta \varepsilon_{558-575}$  of 45 mM<sup>-1</sup>cm<sup>-1</sup> [18]. 134

2.2. Static spectroscopy. Absorption spectra were recorded in a Cary-300 Bio
 spectrophotometer in 1 ml quartz semi-micro cuvettes with a 1 cm optical path. A magnetic
 circular dichroism spectrum was recorded at room temperature in a quartz cuvette with a
 custom-built dichrograph (effective magnetic field ~ 0.7T for a 1 cm cell). The difference
 between direct and reverse magnetic field spectra was calculated to eliminate the contribution
 of the natural circular dirchroism.

2.3. Femtosecond selective photobleaching of the hemes. Multicolor pump-probe 141 spectroscopy was performed with a setup based on a Quantronix Integra-C Ti-Saph oscillator/ 142 amplifier system, delivering ~100 fs, 1.0 mJ pulses centered at 780 nm and operated at a 143 repetition rate of 500 Hz. Pump pulses with different spectral characteristics (maxima at 623, 144 145 608, and 570 nm for the experiments with aa<sub>3</sub>; 623 and 570 nm for the experiments with ba<sub>3</sub>) with an energy of ~1  $\mu$ J were generated using a home-built, BBO-based (1 mm thick,  $\theta$ =29.2°) 146 type I non-collinear optical amplifier (NOPA [19]). The NOPA was pumped with pulses 147 centered at 390 nm, generated by second harmonic generation of part of the 780-nm 148 fundamental beam in a 1-mm  $\theta$ =32° BBO crystal, and seeded with white light continuum 149 generated by focusing part of the 780-nm beam on a 1 mm sapphire plate; the continuum was 150 subsequently filtered appropriately using selected sharp edge interference filters (Semrock). 151

White light probe pulses were generated by focusing part of the fundamental on a 1 mm-thick CaF2 window that was continuously translated perpendicular to the beam. The probe beam was dispersion-minimized around 435 nm by passing through a prism compressor, and then separated into the test and reference beams. After passing through the sample, the beams were dispersed and detected, shot-by-shot, on a CCD camera (Roper Scientific, SPEC-10) configured as a dual-array detector.

To mimic isotropic conditions, the polarization of the pump beam was set at 54.7° 158 («magic angle») relative to polarization of the probe beam using a polarizer and a half-wave 159 plate. To vary delay between the pump and probe pulses, the pump beam was passed 160 through a delay line comprised of a retroreflector mounted on a step-motor driven translation 161 stage. The 1-mm optical pathlength sample cell was rastered in a Lissajous scanner to 162 ensure optimal exchange of the sample in the interaction volume. For every type of pump 163 pulse, 10 spectrokinetics sets were recorded in the range 400-500 nm (spectral resolution ~1 164 nm) and with temporal steps of 83 fs from -1 to 2 ps and of 2 ps from 2 to 50 ps. Due to 165 dispersion the effective zero time depended on the wavelength, varying the onset of the 166 absorption changes for up to 700 fs. We accounted for this dispersion during data processing. 167

2.4. Preparation of the samples for the femtosecond spectroscopy. Enzymes were 168 diluted in 200 mM HEPES/Tris buffer, pH ~ 7, with 0.05% DM dodecylmaltoside and 0.5 mM 169 Na-EDTA. Quartz cuvettes with a 1 mm optical pathway ("Hellma-Analytics"), type 110-QS 170 were filled with the sample and hermetically closed with plug-type rubber sleeve stoppers 171 ("Kimble", USA). CcO and succinate: menaquinone reductase (SMR) were reduced with a 172 fresh 0.2M sodium dithionite solution (final concentration ca. 1-5 mM dithionite). To prepare 173 the dithionite solution, degassed MilliQ was added to degassed dithionite powder into an 174 argon-filled vial. The sample in the cuvette was degassed with a vacuum pump and the 175 volume above the sample was saturated with argon. Dithionite was added through the 176 stopper using a Hamilton syringe, and the volume above the sample was again saturated with 177 argon. Reduction was considered complete if the absorption spectrum corresponded to the 178 reduced form (in the case of CcO showing maxima at 444 nm and 605 nm) and stayed 179 unchanged for at least 10 minutes. The absorption in the Soret band of the reduced form was 180 0.8-1.0 OD. 181

2.5. Data analysis. Global analysis of the pump-induced absorption changes was
 performed using Glotaran software [20] to extract the fastest evolution-associated spectra
 (EASs) with correction for the instrumental response function, zero-time dispersion and

coherent artifacts. Prior to global analysis, noise in the data was reduced using singular value
decomposition. Relaxation included typically three phases with characteristic times of ca. 0.2,
1 and 5 ps, the fastest component representing photobleaching. Additional components did
not improve the fit. Decompositions into gaussians were performed in MATLAB using the
'Interior-point' algorithm [21–23] to minimize the mean square difference between the
experimental and simulated curves.

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### **3. Results**

194 3.1. Reconstruction of heme  $a^{2+}$  absorption bands using Ca<sup>2+</sup>-induced spectral shift. In 195 CcO from vertebrate mitochondria, a specific cation-binding site in subunit I near heme *a* and 196 Cu<sub>A</sub> reversibly binds Ca<sup>2+</sup> and Na<sup>+</sup> ions [24–26]. Ca<sup>2+</sup> red-shifts the entire heme *a* absorption 197 spectrum by ~1 nm [24,27]. The shift is small compared to the linewidths of the heme  $a^{2+}$ 198 absorption bands, so the difference spectrum of the shift represents the differential of the 199 absolute *a*<sub>3</sub>-free heme  $a^{2+}$  spectrum, and can be used to derive the heme  $a^{2+}$  absolute 200 spectrum lineshape [28].

We first compared the Ca<sup>2+</sup>-induced spectral shift in the fully reduced unliganded 201 complex  $a^{2+}a_{3}^{2+}$  and the mixed-valence cyanide complex  $a^{2+}a_{3}^{3+}$ -CN (Fig. 2a,b). In contrast to 202 absolute spectra (Fig. 2a), the Ca<sup>2+</sup>-induced difference spectra (Fig. 2b) for these two states 203 are virtually identical both in the Soret and the visible. This finding confirms that heme  $a^{2+}$  is a 204 unique contributor to the Ca<sup>2+</sup>-induced spectral changes. Thus, the non-zero  $a_3^{2+}$  contribution 205 put forward earlier [24] is not supported by our data obtained using a more stable reduction 206 system. It also shows that the heme a spectrum is not affected by the redox and spin state of 207 heme *a*<sub>3</sub>. 208

We chose the  $a^{2+}a_3^{3+}$ -CN complex as the starting point to better separate the overlapping γ-peaks and minimize any possible interference from heme  $a_3$ . The absolute  $a^{2+}a_3^{3+}$ -CN spectrum (**Fig. 2***a*, black) showed a split Soret band with maxima at 442 nm and 430 nm (mainly heme  $a^{2+}$  and heme  $a_3^{3+}$ -CN, respectively) and rather symmetric α-band at 605 nm.

The corresponding difference spectrum (**Fig. 2b**, black) is dominated by narrow  $1^{st}$  – derivative-like features in the Soret and in the visible. The Soret band signal shows A<sub>min</sub> at 438 nm, A<sub>max</sub> at 460 nm and a zero-crossing point at 449 nm, that coincides with neither of the two maxima in the parent absolute spectrum. An inflection around the zero-crossing point suggests juxtaposition of two overlapping 1<sup>st</sup> derivative-shaped components [24,28]. In the visible, the difference spectrum shows the zero-crossing point at 605 nm, close to the maximum of heme  $a^{2+}$ . A small derivative-shaped feature centered around 520 nm may correspond to a red shift of the  $\beta$ -band of heme  $a^{2+}$ .

The integral of this differential spectrum in **Fig. 2***b* gives the reconstructed lineshape of the heme  $a^{2+}$  absolute spectrum (**Fig. 2***c*, black). A weak baseline (**Fig. 2***b*, green dotted line) was subtracted from the difference spectrum before integration to avoid negative values in the reconstructed spectrum. We decomposed the spectral lineshape into a sum of gaussians with the following assumptions:

(1) the Soret and α-band are comprised of two individual electronic transitions (cf. Fig.
 1*b,c*);

(2) the x and y transitions can be degenerate or non-degenerate in energy;

(3) The x and y transitions of the same band must have equal linewidths, but theirintensity ratio may vary;

(4) Bands in 500-580 nm region are approximated by two gaussians of variablelinewidths.

The reconstructed heme  $a^{2+}$  Soret band with a maximum at 447.5 nm is symmetric and much narrower than reported by Vanneste (full width at half-height (FWHH) 18 nm vs 24 nm, correspondingly [11]). It is described by the sum (**Fig. 2***c*, red) of two gaussians of equal FWHH of 15 nm at 442 nm and 450 nm, and an additional minor component at 420 nm with a FWHH of 18 nm (red dotted). We assign the gaussians at 442 and 450 nm to two nondegenerate B<sub>y</sub> and B<sub>x</sub> electronic transitions in heme  $a^{2+}$ , split by 8 nm, respectively.

Short-wavelength components of the heme  $a^{2+}$  Soret band could reside, in principle, 240 below 420 nm. However, most conventional reductants such as dithionite or ascorbate 241 strongly absorb in the near-UV possibly masking these spectral components. We therefore 242 attempted to obtain a "clean" spectrum of CcO in the Soret band by bubbling gaseous 243 hydrogen in statu nascendi released by sodium borohydride through the anaerobic solution of 244 the oxidized enzyme. Hydrogen bubbling resulted in partial (~85%) CcO reduction. The H<sub>2</sub>-245 reduced CcO spectrum corrected for the fraction of the oxidized form (Fig. 2d) reveals no 246 additional spectral components below 420 nm. 247

The splitting of the Soret band of heme  $a^{2+}$  implies that the  $\alpha$ -band of the heme should similarly contain split x- and y- components [29]. The  $\alpha$ -band differential Ca<sup>2+</sup>-induced spectrum shows certain asymmetry (**Fig. 2***b*), and, consistently, the reconstructed  $\alpha$ -band (maximum at 605 nm, FWHH 21 nm) has somewhat extended short wavelength side. The  $\alpha$ band fitting requires two gaussians of 17-18 nm FWHH separated by 13 nm (at 593 nm and 606 nm). However, in variance with the Soret band, the x and y components of the Q<sub>0</sub>-band differ much in intensity, so that the overall maximum at 605 nm almost coincides with the maximum of the dominating Q<sub>0x</sub>-transition at 606 nm.

To further explore the proposed Q<sub>0</sub>-band splitting, we recorded room temperature 256 magnetic circular dichroism (MCD) spectrum of our CcO preparation in the visible range (Fig. 257 2e, black). The spectrum with positive and the negative lobes at 609 nm and 590 nm and 258 zero-crossing at 600 nm is very close to that in [29], except for a slight shift of our spectrum to 259 the red. The MCD spectrum decomposes into a sum of two gaussians of equal magnitude 260 and width but of opposite sign corresponding to the individual heme  $a^{2+}$  Q<sub>0x</sub> and Q<sub>0y</sub> terms, 261 and a third negative gaussian at 574 nm (Fig. 2e, red dashed: gaussians, red solid: their 262 sum). The extrema of the two components are observed at 594 (min) and 605.5 nm (max), 263 *i.e.*, very close to the maxima of the two components fitted to the heme  $a^{2+} \alpha$ -band (Fig. 2c), 264 supporting their attribution to the Q<sub>0x</sub>- and Q<sub>0y</sub>- electronic transitions. 265 266



26′.

Fig. 2. Reconstruction of the heme  $a^{2+}$  absorption spectrum using selective Ca<sup>2+</sup>-268 induced spectral shift. (a) Absorption spectra of the bovine CcO: black, mixed-valence CN-269 complex  $a^{2+}a_{3}^{3+}$ -CN; red, fully reduced unliganded enzyme  $a^{2+}a_{3}^{2+}$ . To obtain the mixed-270 valence CN-complex, "activated" (see 2.1) CcO (1.7 µM) in the basic medium was incubated 271 for 10 minutes with 4 mM KCN and then reduced by 5 mM Tris-ascorbate in the presence of 272 100 µM TMPD. To produce fully reduced form, CcO was incubated with 5 mM ascorbate and 273 50 µM RuAm, in a cell covered by a gastight cap. Anaerobic condition followed by full enzyme 274 reduction were achieved after 5-min incubation. (b) Difference spectra of the Ca<sup>2+</sup>-spectral 275 shift of heme a. The curves show the difference between the absolute spectra of the CcO 276 mixed-valence CN-complex (black) or fully reduced form (red) after the addition of 0.1 mM 277

CaCl<sub>2</sub> and after the subsequent addition of 1 mM EGTA. With the adventitious calcium in the 278 buffer (~ 30  $\mu$ M), the Ca<sup>2+</sup> concentration after CaCl<sub>2</sub> addition was ~130-fold the K<sub>d</sub> value [26], 279 i.e., sufficient to saturate the calcium binding site. (c) Reconstructed lineshape of the heme 280  $a^{2+}$  absorption spectrum. Black, integral over the black curve in *b*. Red, its approximation by 281 the sum of seven gaussians (dashed red). (d) H<sub>2</sub>-reduced bovine CcO absoption spectrum 282 corrected for the contribution (~15%) of resting oxidized form. The spectrum shows no 283 additional shoulder below 420 nm. Insert: zoom of the 360-420 nm region. (e) Black, MCD 284 spectrum of reduced bovine CcO. Red, its approximation with 3 gaussian components 285 (dashed red). To obtain reduced CcO, 50 µM Ru(NH<sub>3</sub>)<sub>6</sub><sup>3+</sup> and 5 mM ascorbate was added to 286 12 µM CcO in 0.1M HEPES pH 8.1 with 0.1 mM of added Ca<sup>2+</sup>, 5 minutes prior to the 287 recording of the spectrum. 288

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3.2. Individual spectra of hemes in the T. thermophilus ba<sub>3</sub> and mitochondrial aa<sub>3</sub> CcO 291 resolved with photoexcitation spectroscopy. Independently, we employed a selective 292 photobleaching approach [30-33] by exciting the hemes by a femtosecond light pump pulse 293 in the Q band to transiently depopulate the electronic ground state and thus to bleach the 294 corresponding ground state-based absorption. Spectral changes were followed using a white 295 continuum probe pulse. Since both Q and B bands originate from the same ground state, 296 excitation in the Q-band bleaches in the Soret as well. Wide separation of the two bands 297 prevents interference of the relatively narrow pump pulse in the Q-band with the probe pulses 298 in the Soret. 299

However, the heme  $a^{2+}$  and heme  $a_{3}^{2+}$  Q-band overlap and finite intrinsic width of the 300 pump pulses complicates selective excitation of one heme. Therefore, we turned initially to T. 301 thermophilus ba<sub>3</sub>-type CcO belonging to the same heme-copper superfamily as the aa<sub>3</sub>-type 302 oxidases but with heme b instead of heme a. In the ba<sub>3</sub>-type CcO, the Q-bands of high-spin 303 heme  $a_{3^{2+}}$  (~613 nm) and low-spin heme  $b^{2+}$  (559 nm) are separated better than the Q-bands 304 of bovine CcO hemes a and a<sub>3</sub> (Fig. 3a,b). The Soret bleaching only occurs if the pump pulse 305 overlaps with the visible absorption bands. Pumping at 570 nm is expected to excite both 306 hemes and results in photobleaching spectrum with maximum at 426 nm (heme  $b^{2+}$  band) 307 with a weak feature at 443 nm (putative heme  $a_{3^{2+}}$  band) (Fig. 3c). In contrast, the pump 308 pulse at 623 nm ("623exc" condition) (Fig. 3b, red) is expected to selectively excite heme  $a_{3^{2^{+}}}$ 309 since it does not overlap with the heme  $b^{2+}$  narrow  $\alpha$ -band at 559 nm. Compared to the 310 spectrum of heme  $a^{2+}$  in the bovine CcO (Fig. 2c), 623exc photobleaching spectrum (Fig. 3d) 311 is blue-shifted and strongly asymmetric with extended decay on the high-energy side. Besides 312 two equal-width (12 nm) gaussians at 435 and 444 nm assigned to the purely electronic B<sub>0y</sub>-313 and  $B_{0x}$ - transitions, the fitting requires a wider band at ~427 nm that may represent the sum 314 of vibronic components. The absence of 623exc photobleaching of *Bacillus subtilis* reduced 315 succinate: menaquinone reductase (SMR) containing heme b as the only bleaching target 316 317 [34] rules out potential heme b contribution to this 427-nm band (Fig. 3d).

Due to spectral overlap, the photobleaching of individual  $aa_3$  CcO hemes cannot be resolved with single excitation. We therefore used three excitation wavelengths to vary the heme  $a^{2+}$  vs heme  $a_3^{2+}$  excitation degree. **Fig. 4a** shows dithionite-reduced  $aa_3$ photobleaching spectra obtained by excitation with pump pulses centered at 570, 608 and 623 nm ("570exc", "608exc" and "623exc"). The 608exc spectrum qualitatively resembles the CcO photobleaching spectrum reported recently for a similar pump-probe protocol with the

excitation at 600 nm [35,36]. The "570exc", "608exc" and "623exc" spectra differ in amplitude 324 since the photoexcitation yield depends on the absorption strength in the region of excitation 325 and on the pump energy. However, their lineshapes also differ, suggesting varying degree of 326  $a^{2+}$  vs.  $a_{3}^{2+}$  excitation selectivity. The 623exc lineshape is almost symmetrical with a minor 327 shoulder below 430 nm. The shoulder contributes more to 608exc and especially 570exc 328 spectra. The maximum of the 570exc spectrum is slightly blue-shifted (see Fig. 4a, color 329 traits). As the heme  $a_{3^{2+}}$  ground state absorption spectrum is blue-shifted relative to heme  $a^{2+}$ 330 spectrum (cf. Fig. 2c and Fig. 3d), the blue shift of the maximum on the 570exc spectrum 331 provides a first indication that heme  $a_3^{2+}$  contributes more to the 570exc spectrum than to the 332 608exc and 623exc spectra. 333

To estimate the excitation selectivity, we convoluted the pump pulse positions with the 334 Q-bands of the hemes  $a^{2+}$  and  $a_{3}^{2+}$  (Fig. 4b). The spectrum reconstructed from the Ca<sup>2+</sup>-335 induced shift assay was used to model the heme  $a^{2+}$  Q-band, assuming that heme  $a^{2+}$ 336 contributes about 80% of absorption in the  $\alpha$ -band maximum. The heme  $a_{3^{2+}}$  Q-band was 337 modeled by subtraction of the heme  $a^{2+}$  spectrum from the smoothened  $aa_3$  Q-band. The 338 resulting  $a_3^{2+}$  spectrum has a maximum at ~600 nm and a shoulder below 580 nm, consistent 339 with the reduced minus oxidized heme  $a_3$  spectrum [37] (see also Fig. A.5). We then 340 estimated the heme  $a^{2+}$  vs. heme  $a_{3}^{2+}$  selectivity as the convolution of the pulse profile with 341 the heme  $a^{2+}$  absorption band divided by a convolution of the same pulse profile with the 342 heme  $a_{3^{2+}}$  absorption band: 343

344 
$$a: a_3 \ selectivity = \frac{\int I_a(\omega) I_{pulse}(\omega)}{\int I_{a_3}(\omega) I_{pulse}(\omega)},$$

where  $I_a$  and  $I_{a3}$  are heme  $a^{2+}$  and heme  $a_{3}^{2+}$  Q-bands, respectively,  $I_{pulse}$  is the corresponding pump pulse profile (570exc, 608exc or 623exc), and the integrals are taken over the 15391-19992 cm<sup>-1</sup> spectral range (500-650 nm).

(1)

The estimated  $a^{2+}$  vs.  $a_3^{2+}$  excitation yield ratios (**Fig. 4***b*) suggest a higher contribution of heme  $a^{2+}$  to the photobleaching spectra at 623exc and 608exc and approximately equal  $a_3^{2+}$  and  $a^{2+}$  contributions at 570exc. However, under any excitation conditions, it was not possible to assign the photobleaching spectra entirely to either heme  $a^{2+}$  or  $a_3^{2+}$ . To extract "pure" hemes  $a^{2+}$  and  $a_3^{2+}$  spectral lineshapes we performed simultaneous decomposition of the photobleaching spectra obtained at all three excitations (**Fig. 4***c***-e**) with the following constraints:

356 .(1) The Soret (*B*-) band of each heme is modeled by 3 gaussian components: two electronic 357 transitions ( $B_{0x}$ - and  $B_{0y}$ -) with equal widths plus a putative vibronic satellite of variable width. 358 The three gaussians describing heme  $a^{2+}$  and heme  $a_3^{2+}$  contributions are shown in **Fig. 4***c***-e** 359 by orange and green, respectively.

360 (2) Position and bandwidth in cm<sup>-1</sup> (but not amplitude) of a specific transition within a 361 heme is identical for all three photobleaching spectra. For example, the heme  $a^{2+}$  B<sub>0x</sub>-362 associated gaussians are centered at the same wavelength and have the same width in each 363 of the three sets of gaussians but may differ in amplitude.

(3) The resulting shape of the spectrum of an individual heme ( $a^{2+}$  or  $a_{3}^{2+}$ ) is preserved across the excitation conditions. For example, the sums of the three gaussians attributed to heme  $a^{2+}$  have the same lineshape for each of the three sets of gaussians but may differ in amplitude.

(4) The photoinduced absorption is simulated by two negative gaussians, one at either
 side of the spectra, whose positions and bandwidths in cm<sup>-1</sup> were identical across the
 excitation conditions.

The "pure" normalized heme  $a^{2+}$  photobleaching spectrum (**Fig. 4***f*, solid orange) is in agreement with the spectrum obtained using Ca<sup>2+</sup>-induced shift approach (**Fig. 4***f*, gray, same as **Fig. 2***c*, black). It reveals almost symmetrical lineshape with minor shoulder at 421 nm and contains two individual gaussian components with maxima at 442.5 nm and 450 nm, with a bandwidth of 14 nm each. In contrast, the "pure" normalized heme  $a_3^{2+}$  photobleaching spectrum (**Fig. 4***g*) has a pronounced shoulder described by 423 nm band with a 25 nm width. In both hemes *a* and  $a_3$ , the  $B_{x-}$  and  $B_{y-}$  electronic transitions are separated by ~8 nm.

As a consistency check, we attempted to reconstruct the overall  $a^{2+}a_{3}^{2+}$  Soret band as a 378 weighted sum of the individual  $a^{2+}$  and  $a_{3}^{2+}$  contributions. First, we estimated the proportion of 379 the heme  $a^{2+}$  vs. heme  $a_{3}^{2+}$  contribution by assuming that the  $a:a_{3}$  excitation selectivity is 380 most reliably estimated for 608exc condition, where the pumping is close to the heme  $a^{2+}$ 381 band maximum (Fig. 4b). Selectivity modeling suggests that pumping at 608 nm excites 382 heme  $a^{2+}$  3.15 times stronger than heme  $a_{3^{2+}}$ . At the same time, the ratio of the areas under 383 the sum of the gaussians attributed to heme  $a^{2+}$  (orange) to that under the sum of the 384 gaussians attributed to heme  $a_{3^{2+}}$  (green) on the 608exc is not 3.15:1 but 1.70:1 (**Fig. 4***d*). 385 This could mean that the contributions of the two hemes to the total extinction are not 1:1 but 386 rather 0.54:1 (heme  $a^{2+}$ : 35%, heme  $a_{3^{2+}}$ : 65%). The ratio of the areas under the sums of the 387 heme  $a^{2+}$  and heme  $a_{3}^{2+}$ -assigned gaussians for 570exc and 623exc fits (predicted 388

selectivities 0.77 and 4.95, respectively) are 0.58:1 and 2.55:1. This gives the following estimates of the hemes' contributions to the extinction: 43% (heme *a*) and 57% (heme  $a_3$ ) for 570exc fit; 34% (heme *a*) and 66% (heme  $a_3$ ) for 623exc fit. The both estimations are fairly close to that obtained above for the 608exc data which confirms the consistency of our modeling.

Finally, we obtained the reconstructed *a*a<sub>3</sub> Soret lineshape (Fig. 4*h*, red) by adding the 394 normalized heme  $a^{2+}$  spectrum multiplied by 0.54 to the normalized heme  $a_{3}^{2+}$  spectrum 395 multiplied by 1, and normalizing the sum to its maximum. For comparison, the dithionite-396 reduced aa<sub>3</sub> spectrum is shown, also normalized to its maximum (Fig. 4h, gray). The 397 reconstructed Soret lineshape reproduces the position of the maximum and the shoulder at 398 399 the short-wavelength side. We then investigated whether scattering may explain higher experimental absorption below 430 nm compared to the reconstructed spectrum. For this, we 400 added  $\sim \lambda^{-4}$  absorption (**Fig. 4***h*, dotted dark red curve) to the sum of  $a^{2+}$  and  $a_{3}^{2+}$  bands. The 401 maximum-normalized resulting spectrum (Fig. 4h, dotted blue curve) is fairly close to the 402 experimental Soret lineshape. The remaining differences between the reconstructed band 403 and the experimental  $a^{2+}a_3^{2+}$  spectrum can be due to imprecise description of the vibronic 404 satellite(s) by unique gaussian band, unaccounted higher-energy heme transitions, and the  $\beta$ -405 band tails. Attempting more precise Soret band reconstruction, as well as determination of 406 absolute heme absorbance values, is out of the scope of the present work. 407

408



Fig. 3. Photobleaching of reduced ba<sub>3</sub> CcO from T. thermophilus. (a) Dithionite-reduced 411 *T. thermophilus ba*<sub>3</sub> (7.9 µM; solid) and *Bacillus subtilis* succinate: menaquinone reductase 412 (SMR, 4.1 µM; dashed) absorption spectra. (b) Pump laser pulse profiles (right vertical axis) 413 centered at 570 nm (blue) and 623 nm (red) relative to the  $\alpha$ -band of the reduced bas 414 spectrum (black, left vertical axis). The B. subtilis SMR spectrum (dashed) does not overlap 415 with the 623exc pulse profile suggesting that heme  $b^{2+}$  is not excited by this pulse. (*c*,*d*) 416 Inverted fastest EASs after pumping at 570exc and 623exc, respectively. In d, Curve 1 (solid 417 black), experimental spectrum that has been fitted by 5 gaussians. The 3 positive gaussians 418 (dashed red) are assigned to heme  $a_{3^{2+}}$  photobleaching. Two negative gaussians (dashed 419 gray) approximate the increase of absorption following the excitation that likely represent the 420 contribution of the electronically excited states (typically weak and diffuse) [38-40]. Curve 2 421 (solid red) shows the sum of the 5 gaussians mostly coinciding with the experimental curve 1. 422 Curve 3 (black with open circles), control experiment showing that heme  $b^{2+}$  in *B. subtilis* 423 SMR is not photobleached by excitation at 623 nm. 424



428 Fig. 4. Photobleaching of reduced bovine CcO. (a) Inverted fastest EASs after pumping at 429 570, 608 and 623 nm. The traits at the bottom show corresponding maxima positions. (b) 430 Heme  $a^{2+}$  vs. heme  $a_{3}^{2+}$  excitation selectivity estimation. Heme  $a^{2+}$  spectrum (red area) 431 reconstructed from the Ca2+-induced shift (Fig. 2c) was subtracted from the smoothened 432 reduced  $aa_3$  spectrum (gray) to model heme  $a_3^{2+}$  spectrum in the visible range (green area). 433 Heme  $a^{2+}$  was assumed to provide ~80% absorption at the  $\alpha$ -band maximum. The pump 434 pulse profiles for 570exc, 608exc and 623exc are shown in blue, green and red, respectively. 435 Excitation selectivity was estimated using Eqn (1). (c-e) Approximation of the EASs (thick 436 gray lines) by sets of 6 positive (dashed) and 2 negative (dotted) gaussians. In each fit, three 437 positive gaussians were attributed to heme  $a_{3^{2+}}$  (green) and the other three to heme  $a^{2+}$ 438 (orange). The two negative gaussians describe photoinduced absorption. The sum of the 439 gaussians is shown by a solid red line. (*f*,*g*) Heme  $a^{2+}$  and heme  $a_{3}^{2+}$  Soret lineshapes (solid 440 lines) determined from the EASs fitting as sums of the three corresponding gaussians 441

- (dashed lines). The lineshapes are normalized to their maxima. In *f*, heme  $a^{2+}$  Soret lineshape
- 443 determined with  $Ca^{2+}$ -induced shift technique (gray) is shown for comparison. (*h*) Bovine
- 444  $a^{2+}a_{3}^{2+}$  Soret lineshape reconstruction. The ratio of the areas under scaled Soret bands of
- hemes  $a^{2+}$  and  $a_{3}^{2+}$  (orange and green) was estimated to be  $a^{2+}:a_{3}^{2+}=0.54:1$  (see Section
- 446 3.2). Red, the sum of the individual  $a^{2+}$  and  $a_{3}^{2+}$  contributions normalized to its maximum.
- 447 Dotted dark red, simulated  $\sim \lambda^{-4}$  scattering. Dotted blue, the sum of reconstructed  $a^{2+}$  and  $a_{3}^{2+}$
- 448 bands and the scattering normalized to its maximum. Gray, Soret band of the dithionite-
- reduced bovine CcO normalized to its maximum.

### **4**50 **4. Discussion**

**Table 1** resumes heme  $a^{2+}$  and  $a_3^{2+}$  spectral properties reported here and compares 451 them with previous studies. Only two papers treat explicitly the lineshape of the absolute 452 spectra of bovine CcO hemes  $a^{2+}$  and  $a_{3}^{2+}$ : Vanneste's [11] and Orii's [13]. Our spectra 453 coincide with neither of the two reports but would appear closer to Orii's conclusions. 454 However, the latter depend critically on the choice of the model parameters and can be 455 converted into Vanneste's attributions by changing the baseline slope (Fig. A.2). It has been 456 shown by using improved reduction technique that Vanneste's heme  $a_3$  spectra need to be 457 corrected at least in the  $\alpha$ -region [37]. In the following we will discuss the spectral properties 458 of the two hemes in detail. 459

4.1. Soret splitting. Both hemes  $a^{2+}$  and  $a_{3}^{2+}$  showed split Soret bands with  $B_{0x-}$  and  $B_{0y-}$ 460 components separated by 8-9 nm and close electronic transition widths. The heme  $a^{2+}$  Soret 461 band splitting has been inferred earlier [28,29,41–43], but sometimes misinterpreted as 462 evidence for two heme *a* conformations [42]. The hemes *a* and *a*<sub>3</sub> in C*c*O have the same 463 chemical structure - heme A, - but differ in the protein surroundings and in iron coordination. 464 Therefore, the simplest explanation is that the similar splitting observed for the two hemes is 465 intrinsic to heme A structure. Indeed, the electronic transition along the pyrrole II  $\rightarrow$  pyrrole IV 466 axis with the two electron-withdrawing substituents (vinyl and formyl) should differ in energy 467 from the perpendicular transition along the pyrrole I  $\rightarrow$  pyrrole III axis with electron donating 468 substituents (methyl, propionate, hydroxyethylfarnesyl group). Conversion of the formyl 469 substituent to the Schiff-base results in drastic narrowing of the heme A Soret band [44] (see 470 also Fig. A.3) suggesting restoration of degeneracy or at least reduction of splitting. Inherent 471 heme A splitting is supported by the MCD spectrum of the bis-1-methyl imidazole-6c complex 472 of heme  $A^{2+}$  showing the same lineshape (except for the position) as the CcO heme  $a^{2+}$  MCD 473 spectrum [29]. However, it seemingly contradicts the absorption spectra second derivative 474 analysis (cf. [41] and refs. therein) suggesting that the Soret splitting is specific for heme  $a^{2+}$ 475 476 in CcO and is not observed in the hexacoordinate (6c-) low-spin (ls-) complexes of isolated heme A<sup>2+</sup> nor in non-aa<sub>3</sub>-CcO proteins reconstituted with heme a. 477

To clarify this issue, we first tested whether deconvolution into gaussians can reveal the Soret splitting of the isolated heme  $A^{2+}$  spectrum by analyzing the bis-imidazole-heme  $A^{2+}$ spectra reported in [44]. The Is-6c-heme  $A^{2+}$  spectra in solution are broader than in CcO (35 nm in aqueous buffer or 27 nm in CH<sub>2</sub>Cl<sub>2</sub> *vs* 18 nm in the CcO protein) and, at this width, single- and two-gaussian approximations (both including one additional gaussian describing a

vibronic shoulder) are equally valid (Fig. A.3). Further, simple modeling (Fig. A.4) shows that 483 the ~10-nm splitting of two gaussians with equal heights and widths typical for CcO requires 484 their FWHH being below ~8 nm to be detected on the second derivative curve. The exact 485 lineshape of the second derivative depends also on the relative heights of the two gaussians. 486 Thus, the second-derivative method may resolve the  $B_{0x}$  and  $B_{0y}$  bands of heme *a* in the CO-487 complex of reduced CcO, with heme  $a_3$  bands shifted far to the blue (cf. Fig. 1B in [41]) and to 488 a lesser extent in the fully reduced CcO (Fig. 1A in ref. [41]), but is not expected to resolve the 489  $B_{0x}$ - and  $B_{0y}$ - bands in the broad spectra of the isolated 6c-heme  $A^{2+}$ . 490

The same applies to the recent second-derivative analysis of the Soret CcO spectra confirming ferrous heme *a* Soret splitting but arguing against splitting of the ferric hemes *a* and  $a_3$  Soret bands [45]. On this basis, the authors proposed that the ferric hemes are symmetric but their reduction changes the protein conformation distorting heme planes and splitting the B<sub>0x</sub>- and B<sub>0y</sub>-transitions. However, the Soret bands of the ferric hemes are known to be broader than those of the reduced hemes and their splitting may be unresolved by the second derivative assay.

We conclude that although not detectable by second-derivative analysis, the inherent splitting of the heme A<sup>2+</sup> Soret band is in the origin of the Soret band splitting of both hemes  $a^{2+}$  and  $a_3^{2+}$ .

4.2. Heme  $a^{2+} \alpha$ -band splitting. Spectral forms of heme  $a^{2+}$  differing in the  $\alpha$ -peak were 501 described [46,47] but the possibility of a split  $\alpha$ -band of the heme was not considered in these 502 studies. Our analysis revealed the presence of a second minor component at 593 nm in 503 addition to the dominating band at 606 nm (Fig. 2c). Their attribution to the split  $Q_{0y}$  and  $Q_{0x}$ 504 bands of heme  $a^{2+}$ , respectively, is supported by MCD spectroscopy that allows for better 505 resolution of the split bands. The  $\alpha$ -band heme  $a^{2+}$  MCD spectrum can be interpreted as a 506 sum of two opposite B-terms belonging to the  $Q_{0x}$  and  $Q_{0y}$  transitions in heme  $a^{2+}$  [29]. The 507 MCD spectrum of our preparation is indeed described as the sum of two identical B-terms of 508 509 opposite signs (Fig. 2e).

The Soret splitting is more extensive than that of the  $\alpha$ -band (381-416 vs 373 cm<sup>-1</sup>) probably due to the larger Soret vs  $\alpha$ -band transition moment implying stronger interaction with the permanent local electric field [48]. In contrast to the Soret where the B<sub>0x</sub>/ B<sub>0y</sub> intensity ratio is about 1.4 (see **Table 1**), the Q<sub>0y</sub> band at 593 nm is much weaker than the Q<sub>0x</sub> component at 606 nm (absorption intensity ratio Q<sub>0x</sub>/Q<sub>0y</sub> is ~4). The weaker high-energy (y) Q<sub>0</sub>-component is typical of metal-porphyrin compounds with split absorption bands [49–53] and agrees with calculations [50,54]. Therefore, the commonly observed  $\alpha$ -peak of heme  $a^{2+}$ at 605 nm largely corresponds to the Q<sub>0x</sub> transition of the iron-porphyrin ring of heme  $a^{2+}$ .

4.3. Position of the Soret absorption peaks of the hemes  $a^{2+}$  and  $a_3^{2+}$ . It is commonly assumed that the Soret absorption maxima of hemes  $a^{2+}$  and  $a_3^{2+}$  are close to each other (Vanneste: 444 nm  $a^{2+}$  vs 443 nm  $a_3^{2+}$  [11]; Orii: 445 nm  $a^{2+}$  vs 443 nm  $a_3^{2+}$  [13]). This fact has been long considered surprising in view of the difference in iron coordination and hemeprotein interaction.

Our data, in contrast, suggest that the Soret maxima (447-447.5 nm  $a^{2+}$  vs 441 nm  $a_{3}^{2+}$ ) 523 and the B<sub>0x</sub>-, B<sub>0y</sub>-transition energies of hemes  $a^{2+}$  and  $a_{3}^{2+}$  do actually differ. The heme  $a_{3}^{2+}$ 524 B<sub>0x</sub>- and B<sub>0y</sub>-transitions are shifted by 6-7 nm to the blue relative to the corresponding 525 transitions of heme  $a^{2+}$ , in agreement with the blue shift of the reduced minus oxidized 526 difference spectrum of CcO heme a<sub>3</sub> relative to that of heme a [37]. In apparent contradiction, 527 Soret maxima of high spin mono-His-coordinated (5c-hs) cytochromes b are usually situated 528 to the red from those of low spin, bis-His-coordinated (6c-ls) cytochromes b (Table A.1). 529 Congruently, the wide  $\alpha$ -band of 5c-hs 1,2-diMelM-heme A<sup>2+</sup> (heme  $a_{3^{2+}}$  model) locates to the 530 red of the  $\alpha$ -band of 6c-ls bis-1-Melm-heme A<sup>2+</sup> (heme  $a^{2+}$  model) [29,41,44]. The inverted 531 maxima positions in CcO may be explained by the opposite shifts of the CcO hemes relative 532 to the corresponding isolated complexes. The spectrum of CcO heme  $a^{2+}$  is red-shifted by 9-533 11 nm in the Soret and by 12-16 nm in the  $\alpha$ -band compared to the 6c-ls bis-1-Melm-heme 534 A<sup>2+</sup> (**Table 1**). This shift is probably induced by a strong hydrogen bond between the in-plane 535 oriented heme a formyl group and the conserved arginine R38<sub>bt</sub> as proposed originally in [55] 536 and confirmed by mutagenesis of the homologous arginine in bacterial CcOs [56,57]. Indeed, 537 the replacement of the homologous arginine to methionine causes a 605 nm to 589 nm shift 538 of the  $\alpha$ -band that would correspond to ~10 nm shift in the Soret [56]. On the other hand, the 539 reduced *minus* oxidized difference spectrum of the CcO heme *a*<sub>3</sub> appears to be blue-shifted 540 by about 7 nm compared to that of 5c-hs heme  $A^{2+}$  while preserving the spectral shape ([37]; 541 **Fig. A.5**) suggesting ~3-4 nm blue shift in the Soret. 542

4.4. *B-band short-wavelength shoulder*. Heme  $a_3^{2+}$  shows a pronounced high-energy shoulder of the Soret peak (**Fig. 4***g*), presumably a vibronic satellite of the B<sub>0</sub> transitions. Isolated 6c-ls heme A<sup>2+</sup> also exhibits a shoulder ([44]; **Fig. A.3***a*,*b*). However, no vibronic satellite can be seen in the spectrum of heme  $a^{2+}$  in C*c*O, except for a tiny band at 420-421 nm (**Fig. 2***c*, *4f*). The vibronic β-band shows two peaks at ~517 nm and ~560 nm. The one at ~517 nm [11] or both peaks [58] were attributed to heme  $a^{2+}$  features. Both peaks are red-

shifted by Ca<sup>2+</sup> and appear in the integral of the Ca<sup>2+</sup>-induced difference spectrum (Fig. 2c). 549 The first vibrational progression peak, if any, of the  $a^{2+}$  B band would therefore be expected at 550 1500 cm<sup>-1</sup> higher energy than the electronic peak, i.e., at ~419 nm. Therefore, the small 420-551 nm band in the reconstructed heme  $a^{2+}$  spectrum could indeed indicate the presence of a 552 strongly dampened vibronic mode. The reason of the dampening is not yet clear. However, 553 the nearly symmetric lineshape of the heme  $a^{2+}$  spectrum is consistent with Soret spectra of 554 other reduced 6c-ls cytochromes, for example, cytochrome c<sub>5</sub> [59–61], neuroglobin [62–66], 555 cytochrome b<sub>559</sub> [62], superoxide dismutase [62], cytochrome b<sub>560</sub> [67], cytoglobin [68], and 556 SMR (Fig. 3a). The absence of the shoulder in the spectrum of the 6c-ls heme integrated into 557 the protein could be explained by silencing of the corresponding heme plane vibrations by 558 axial histidines attached to the polypeptide backbone so that out-of-plane movements may be 559 restricted. Heme  $a_{3^{2+}}$  and the free bis-imidazole 6c-ls heme  $A^{2+}$  would have this vibrational 560 mode permitted because the heme is either attached to only one histidine or to axial 561 imidazoles that are not connected to any larger scaffold. 562

4.5. Conclusion and outlook. Upon reinvestigating the decomposition of the reduced aa<sub>3</sub> optical absorption spectrum we propose different contributions of the *a* and *a*<sub>3</sub> hemes that may resolve long-lasting confusion in the literature. We foresee that this work will be valuable for (re-evaluation of) spectroscopic studies of the CcO mechanism. Extension of our approach may come from polarization-sensitive studies in particular on macroscopically ordered systems and true 2-dimensional pump-probe approaches.

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# 6. Data availability

578 Numeric data and all scripts used to analyze data and produce figures are downloadable 579 from Mendeley Data: <u>https://data.mendeley.com/datasets/dmrkn9pfrr/3</u> [69].

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