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THE NONLINEAR HEAT EQUATION INVOLVING HIGHLY SINGULAR INITIAL VALUES AND NEW BLOWUP AND LIFE SPAN RESULTS

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ABSTRACT. In this paper we prove local existence of solutions to the nonlinear heat equation $u_t = \Delta u + a|u|^\alpha u$, $t \in (0, T)$, $x = (x_1, \dots, x_N) \in \mathbb{R}^N$, $a = \pm 1$, $\alpha > 0$; with initial value $u(0) \in L^1_{\text{loc}}(\mathbb{R}^N \setminus \{0\})$, anti-symmetric with respect to x_1, x_2, \dots, x_m and $|u(0)| \leq C(-1)^m \partial_1 \partial_2 \cdots \partial_m (|x|^{-\gamma})$ for $x_1 > 0, \dots, x_m > 0$, where $C > 0$ is a constant, $m \in \{1, 2, \dots, N\}$, $0 < \gamma < N$ and $0 < \alpha < 2/(\gamma + m)$. This gives a local existence result with highly singular initial values.

As an application, for $a = 1$, we establish new blowup criteria for $0 < \alpha \leq 2/(\gamma + m)$, including the case $m = 0$. Moreover, if $(N - 4)\alpha < 2$, we prove the existence of initial values $u_0 = \lambda f$, for which the resulting solution blows up in finite time $T_{\max}(\lambda f)$, if $\lambda > 0$ is sufficiently small. We also construct blowing up solutions with initial data $\lambda_n f$ such that $\lambda_n^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]}$ $T_{\max}(\lambda_n f)$ has different finite limits along different sequences $\lambda_n \rightarrow 0$. Our result extends the known “small lambda” blow up results for new values of α and a new class of initial data.

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1. INTRODUCTION

In this paper, which is a continuation of our previous article [27], we study the existence and uniqueness of local in time solutions to the semilinear heat equation

$$u_t = \Delta u + a|u|^\alpha u, \quad (1.1)$$

where $u = u(t, x) \in \mathbb{R}$, $t > 0$, $x \in \mathbb{R}^N$, $a = \pm 1$, $\alpha > 0$, with highly singular initial data. In addition, we obtain several new results on finite time blowup of solutions of (1.1) in the case $a = 1$.

In [27] we considered initial data of the form

$$u_0 = K(-1)^m \partial_1 \partial_2 \cdots \partial_m \delta, \quad (1.2)$$

in the case $a = 1$, where δ is the Dirac point mass at the origin. In this paper, our motivating example is the tempered distribution

$$u_0 = K(-1)^m \partial_1 \partial_2 \cdots \partial_m (|\cdot|^{-\gamma}) \in \mathcal{S}'(\mathbb{R}^N), \quad (1.3)$$

where m is an integer, $1 \leq m \leq N$, $0 < \gamma < N$, and $K \in \mathbb{R}$. In fact, we consider a more general class of initial data which are in some sense bounded by (1.3). By a local solution, we mean a function $u : (0, T] \rightarrow C_0(\mathbb{R}^N)$, a classical solution of (1.1) for $t \in (0, T]$, such that $u(t) \rightarrow u_0$ in $\mathcal{S}'(\mathbb{R}^N)$ as $t \rightarrow 0$. Here $\mathcal{S}'(\mathbb{R}^N)$ is the space of tempered distributions on \mathbb{R}^N , and

$$C_0(\mathbb{R}^N) = \left\{ f : \mathbb{R}^N \rightarrow \mathbb{R} \text{ continuous ; } \lim_{|x| \rightarrow \infty} f(x) = 0 \right\}.$$

As is standard practice, we study equation (1.1) via the associated integral formulation

$$u(t) = e^{t\Delta} u_0 + a \int_0^t e^{(t-\sigma)\Delta} [|u(\sigma)|^\alpha u(\sigma)] d\sigma, \quad (1.4)$$

where $e^{t\Delta}$ is the heat semigroup on \mathbb{R}^N defined by

$$e^{t\Delta} \varphi = G_t \star \varphi, \quad t > 0, \quad \varphi \in \mathcal{S}'(\mathbb{R}^N), \quad (1.5)$$

G_t being the Gauss kernel given by

$$G_t(x) = (4\pi t)^{-\frac{N}{2}} e^{-\frac{|x|^2}{4t}}, \quad t > 0, \quad x \in \mathbb{R}^N. \quad (1.6)$$

Throughout this paper, $\Omega \subset \mathbb{R}^N$ denotes the sector

$$\Omega = \left\{ x = (x_1, x_2, \dots, x_N) \in \mathbb{R}^N; x_1 > 0, x_2 > 0, \dots, x_m > 0 \right\}, \quad (1.7)$$

where $0 \leq m \leq N$ is an integer. The case $m = 0$ corresponds to $\Omega = \mathbb{R}^N$. We write Ω rather than Ω_m to simplify notation and since the value of m will always be clear in context. The important observation is that functions in $C_0(\mathbb{R}^N)$ which are anti-symmetric in x_1, x_2, \dots, x_m restrict to functions in $C_0(\Omega)$, and functions in $C_0(\Omega)$ extend uniquely to functions in $C_0(\mathbb{R}^N)$ which are anti-symmetric in x_1, x_2, \dots, x_m . Our basic approach, as in [27], is to study equation

(1.4) on Ω and extend the results by anti-symmetry to \mathbb{R}^N (in the case $m \geq 1$). One advantage of this approach is the ability to use positivity and comparison properties on Ω which would not hold directly on \mathbb{R}^N . Of course, in the case $m = 0$ the antisymmetry condition is vacuous.

Our first result directly complements the main result in [27] for the absorption case, i.e. (1.1) with $a = -1$.

Theorem 1.1. *Let $a = -1$ in (1.4) and let the positive integer m and the real number α be such that*

$$1 \leq m \leq N, \quad 0 < \alpha < \frac{2}{N+m}.$$

Let $K \in \mathbb{R}$, $K \neq 0$, and set

$$u_0 = K(-1)^m \partial_1 \partial_2 \cdots \partial_m \delta. \quad (1.8)$$

There exists a continuous solution $u : (0, \infty) \rightarrow C_0(\mathbb{R}^N)$ of the integral equation (1.4) with initial value u_0 and such that $u(t) \rightarrow u_0$ in $\mathcal{S}'(\mathbb{R}^N)$ as $t \rightarrow 0$. Furthermore, $u(t)$ is anti-symmetric in x_1, x_2, \dots, x_m , and if $K > 0$ then $u(t, x) > 0$ for all $t > 0$ and $x \in \Omega$.

For $m = 0$ the previous result is obtained in [3, Theorem 3 and Remark 3, p. 82]. For $m = 1$ and $\alpha < \frac{1}{N+1}$ it is obtained in [10, Corollary 11, p. 6110]. The result of [10] is stated for $a = 1$ but the proof is also valid for $a = -1$. See also [20, 24, 32] and references therein for the existence of solution to the nonlinear heat equation with other singular initial values.

The main purpose of this paper is to consider initial values of the form (1.3), as well as a more general class of singular initial values. We have obtained the following result for initial data given by (1.3).

Theorem 1.2. *Let the positive integer m and the real numbers α, γ be such that*

$$1 \leq m \leq N, \quad 0 < \gamma < N, \quad 0 < \alpha < \frac{2}{\gamma+m}.$$

Let $K \in \mathbb{R}$, $K \neq 0$, and set

$$u_0 = K(-1)^m \partial_1 \partial_2 \cdots \partial_m (|\cdot|^{-\gamma}) \in \mathcal{S}'(\mathbb{R}^N). \quad (1.9)$$

There exist $T > 0$ and a continuous solution $u : (0, T] \rightarrow C_0(\mathbb{R}^N)$ of the integral equation (1.4) with initial value u_0 and such that $u(t) \rightarrow u_0$ in $\mathcal{S}'(\mathbb{R}^N)$ as $t \rightarrow 0$. Moreover, $u(t)$ is anti-symmetric in x_1, x_2, \dots, x_m , and if $K > 0$ then $u(t, x) > 0$ for all $t \in (0, T]$ and $x \in \Omega$.

These solutions can be extended to maximal solutions $u : (0, T_{\max}(u_0)) \rightarrow C_0(\mathbb{R}^N)$. In the case $a = -1$, $T_{\max}(u_0) = \infty$. In the case $a = 1$, $T_{\max}(u_0) < \infty$.

Remark 1.3. The condition $\alpha < 2/(\gamma + m)$ is in some sense optimal. See Proposition 2.8 below for an explanation.

Remark 1.4. The upper bound $\alpha < 2/(\gamma + m)$ in Theorem 1.2 for which a solution exists with initial value given by (1.3) is precisely as predicted by a formal scaling argument. The initial value $\partial_1 \partial_2 \cdots \partial_m (|\cdot|^{-\gamma})$, $0 < \gamma < N$, is in the Sobolev space $H^{s,r}(\mathbb{R}^N)$, when

$$s + m + \gamma < \frac{N}{r} < \gamma.$$

On the other hand, the scaling critical exponent s_c is given by $s_c = \frac{N}{r} - \frac{2}{\alpha}$. To see this, we observe that the transformation $\lambda^{\frac{2}{\alpha}} u(\lambda^2 t, \lambda x)$, which leaves invariant the set of solutions to $u_t = \Delta u + a|u|^\alpha u$, acts on initial values as $\lambda^{\frac{2}{\alpha}} u_0(\lambda \cdot)$. Furthermore,

$$\|\lambda^{\frac{2}{\alpha}} u_0(\lambda \cdot)\|_{\dot{H}^{s,r}(\mathbb{R}^N)} = \lambda^{\frac{2}{\alpha} - \frac{N}{r} + s} \|u_0\|_{\dot{H}^{s,r}(\mathbb{R}^N)}, \quad \forall \lambda > 0.$$

In other words, the homogeneous Sobolev norm of $\dot{H}^{s,r}(\mathbb{R}^N)$ is invariant under the action of $\lambda^{\frac{2}{\alpha}} u_0(\lambda \cdot)$ precisely if $s = s_c$. We expect that solutions exist if the initial value is in some $H^{s,r}(\mathbb{R}^N)$ with $s > s_c$. With the initial value given by (1.3), this is possible when $\frac{N}{r} - \frac{2}{\alpha} < \frac{N}{r} - (\gamma + m)$ i.e. $\alpha < 2/(\gamma + m)$. A similar explanation for initial data given by (1.2) is given in [27].

We have also obtained a local existence result for more general class of initial values, in some sense “bounded by” the distribution (1.3).

Theorem 1.5. *Let the positive integer m and the real numbers α, γ be such that*

$$1 \leq m \leq N, \quad 0 < \gamma < N, \quad 0 < \alpha < \frac{2}{\gamma + m}.$$

Let $u_0 \in L^1_{\text{loc}}(\mathbb{R}^N \setminus \{0\})$ be anti-symmetric with respect to x_1, x_2, \dots, x_m and satisfy

$$|u_0(x)| \leq C(-1)^m \partial_1 \partial_2 \cdots \partial_m (|x|^{-\gamma}), \quad x \in \Omega, \quad (1.10)$$

where $C > 0$ is a constant. There exist $T > 0$ and a continuous solution $u : (0, T] \rightarrow C_0(\mathbb{R}^N)$ of the integral equation (1.4) with initial value u_0 and such that $u(t) \rightarrow u_0$ in $\mathcal{S}'(\mathbb{R}^N)$ as $t \rightarrow 0$. Moreover, $u(t)$ is anti-symmetric in x_1, x_2, \dots, x_m , and if $u_0(x) \geq 0$ ($u_0(x) \not\equiv 0$) for $x \in \Omega$, then $u(t, x) > 0$ for all $t \in (0, T]$ and $x \in \Omega$.

These solutions can be extended to maximal solutions $u : (0, T_{\max}(u_0)) \rightarrow C_0(\mathbb{R}^N)$. In the case $a = -1$, $T_{\max}(u_0) = \infty$. In the case $a = 1$, $\liminf_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda u_0) > 0$.

See Theorem 3.5 below (in Section 3) for a more general and precise version of the last result, including uniqueness and continuous dependence. Uniqueness of the solutions constructed in Theorems 1.2 and 1.5 holds in the space of functions anti-symmetric in x_1, x_2, \dots, x_m and bounded by a multiple of the solution of the linear heat equation with the initial value given by (1.9). See Sections 2 and 3 for more details.

Remark 1.6. The case $m = 0$ in Theorem 1.5 is known. See [7, Theorem 2.8, p. 313] for $a = 1$ and [4, Theorem 8.8, p. 536] for $a = -1$.

We now turn to the blowup results, and for this discussion we of course take $a = 1$ in (1.4). The first result is related to the decay of the initial data at infinity. Let $u \in C([0, T_{\max}(u_0)))$ be the solution of (1.1) with initial data $u_0 \in C_0(\mathbb{R}^N)$, $u_0 \geq 0$. It is shown in [15, Theorem 3.2 (i), p. 372] that there exists $C > 0$ such that if

$$\liminf_{|x| \rightarrow \infty} |x|^{2/\alpha} u_0(x) > C, \quad (1.11)$$

then $T_{\max}(u_0) < \infty$. This last condition is improved in [26, Theorem 2, p. 66] to be

$$\liminf_{|x| \rightarrow \infty, x \in D} |x|^{2/\alpha} u_0(x) > C, \quad (1.12)$$

where $C = C(D) > 0$ and D is a conic sector of \mathbb{R}^N . The result of [26] is valid for u_0 nonnegative belonging to the Sobolev space $W_0^{1,s}(O)$, s large and $O \subset \mathbb{R}^N$ a regular domain containing a sectorial domain D . In [19], the condition (1.12) is both improved and weakened. The improvement is that domains D which become thin compared to a sector as $|x| \rightarrow \infty$ are allowed, but the result only applies to $u_0 \in L^\infty(\mathbb{R}^N)$ such that

$$\liminf_{|x| \rightarrow \infty, x \in D} |x|^\gamma u_0(x) > C, \quad (1.13)$$

for some $0 < \gamma < 2/\alpha$, and $C > 0$ any positive constant. In this case, if $u_0 \geq 0$, then $T_{\max}(u_0) < \infty$. A similar result holds if $u_0 \in W^{1,\infty}(\mathbb{R}^N)$, u_0 not necessarily positive, but with an additional condition on ∇u_0 . See [19, Theorem 2.1, p. 1435], and the remark which follows. Finally, the result of [19] is improved in [25, Theorems 1 and 2, p. 1019] in the case of positive initial value by giving more precise conditions on the initial value, including the ‘‘critical case’’, and by extending the type of domains considered.

We are able to extend and improve some of these blowup criteria. As an example, it is a consequence of the next theorem that if $u_0 \in C_0(\mathbb{R}^N)$, $u_0 \geq 0$, is such that $u_0(x) \geq c|x|^{-2/\alpha} \sin^2(\log|x|)$ for $c > 0$ large enough and for all large $|x|$, then $T_{\max}(u_0) < \infty$. See Remark 1.8(iii) just below for details and Remark 4.1 for more exotic examples. As far as we are aware, this example is not covered by any previously known results. On the other hand, it seems that our methods do not recover the results in [19, 25] on domains which are asymptotically thinner than sectors.

Throughout this paper, we let D_λ denote the dilation operator

$$(D_\lambda f)(x) = f(\lambda x) \quad (1.14)$$

for all $\lambda > 0$. For future reference, we recall the commutation relation between the operator D_λ and the heat semigroup on \mathbb{R}^N ,

$$D_\lambda e^{\tau \lambda^2 \Delta} = e^{\tau \Delta} D_\lambda \quad (1.15)$$

for all $\lambda > 0$ and $\tau > 0$.

Theorem 1.7. *Let $a = 1$ in equation (1.4) and let the integer m and the real numbers α, γ be such that*

$$0 \leq m \leq N, \quad 0 < \gamma < N, \quad 0 < \alpha \leq \frac{2}{\gamma + m}.$$

Let $u_0 \in C_0(\mathbb{R}^N)$, anti-symmetric with respect to x_1, x_2, \dots, x_m , and let $u \in C([0, T_{\max}(u_0)), C_0(\mathbb{R}^N))$ be the maximal solution of (1.1) with initial data u_0 . Let $f \in C_0(\mathbb{R}^N)$ be such that

- (i) *f is anti-symmetric with respect to x_1, x_2, \dots, x_m ,*
- (ii) *$u_0(x) \geq f(x) \geq 0$ for $x \in \Omega$,*
- (iii) *there exists $K > 0$ such that*

$$f(x) \leq Kx_1x_2 \cdots x_m|x|^{-\gamma-2m}, \quad x \in \Omega. \quad (1.16)$$

Suppose there exist a sequence $\lambda_n \rightarrow \infty$ and $z \in L^1_{\text{loc}}(\mathbb{R}^N \setminus \{0\})$, $z \not\equiv 0$, anti-symmetric with respect to x_1, x_2, \dots, x_m , such that

$$\lambda_n^{\gamma+m} D_{\lambda_n} f|_{\Omega} \rightarrow z|_{\Omega} \quad (1.17)$$

in $\mathcal{D}'(\Omega)$, as $n \rightarrow \infty$. If $\alpha = \frac{2}{\gamma+m}$, assume in addition that

$$\|e^{\Delta} z\|_{L^{\infty}(\mathbb{R}^N)} > \left(\frac{1}{\alpha}\right)^{\frac{1}{\alpha}}. \quad (1.18)$$

Then $T_{\max}(u_0) < \infty$.

Remarks 1.8.

- (i) By comparison on Ω , it suffices to prove this result with $u_0 = f$. Furthermore, Theorem 1.7 applies not just to $u_0 \in C_0(\mathbb{R}^N)$, but to any anti-symmetric $u_0 \in L^1_{\text{loc}}(\mathbb{R}^N \setminus \{0\})$, with $u_0 \geq f$ on Ω , for which a local solution of equation (1.4), with $a = 1$, exists (and is nonnegative on Ω). This includes, for example, the initial data prescribed in Theorem 1.5.
- (ii) It follows from the considerations in [5, Section 2] and [21, Section 3] (see also [6, Section 5]) that under condition (1.16), then for any sequence $\lambda_n \rightarrow \infty$ there is a subsequence, which we still call λ_n , and $z \in L^1_{\text{loc}}(\mathbb{R}^N \setminus \{0\})$, anti-symmetric with respect to x_1, x_2, \dots, x_m , such that (1.17) holds. The issue in Theorem 1.7 is that we need $z \not\equiv 0$, or large in some sense if $\alpha = \frac{2}{\gamma+m}$. Also, in the case $m \geq 1$ we need to show that $e^{\Delta} z$ is well defined. This follows from the results in [21] and will be explained in Section 4.
- (iii) In the case $m = 0$, i.e. $\Omega = \mathbb{R}^N$ (and so the anti-symmetry condition is vacuous), taking $\alpha = 2/\gamma$, one can see that if $f \in C_0(\mathbb{R}^N)$ is such that $f(x) = c|x|^{-2/\alpha}$ for large x , for example $|x| \geq \rho$, then $\lambda^{2/\alpha} D_{\lambda} f(x) = c\lambda^{2/\alpha} |\lambda x|^{-2/\alpha} = c|x|^{-2/\alpha}$ for all $|x| \geq \rho/\lambda$. Thus $\lambda^{2/\alpha} D_{\lambda} f \rightarrow c|\cdot|^{-2/\alpha}$ in $\mathcal{D}'(\mathbb{R}^N)$, as $\lambda \rightarrow \infty$ and we recover the result of [15]. Furthermore, if $f(x) = c|x|^{-2/\alpha} \sin^2(\log|x|)$ for all $|x| \geq \rho$, then $\lambda^{2/\alpha} D_{\lambda} f(x) =$

$c|x|^{-2/\alpha} \sin^2(\log|x| + \log \lambda)$ for all $|x| \geq \rho/\lambda$, and so one may take $\lambda_n = e^{\pi n}$ in order to apply to above result.

- (iv) If $\alpha \leq 2/(N+m)$ then the anti-symmetry of u_0 and the condition $u_0 \geq 0$, $u_0 \not\equiv 0$ on Ω are enough to guarantee $T_{\max}(u_0) < \infty$. This is true thanks to Fujita-type results ([1, 17, 18, 14]) on the sector Ω . In other words, the above result is only of interest in the range $\frac{2}{N+m} < \alpha \leq \frac{2}{\gamma+m}$, and it is the case $\alpha = \frac{2}{\gamma+m}$ which is of most interest.

The other main blowup results in this article concern initial values $u_0 = \lambda f$, where $f \in C_0(\mathbb{R}^N)$ is anti-symmetric in x_1, x_2, \dots, x_m , but f is not necessarily positive on Ω . We give conditions on f which imply that $T_{\max}(u_0) = T_{\max}(\lambda f) < \infty$ for all sufficiently small $\lambda > 0$. Results of this type were first obtained by Dickstein [7], and our methods are inspired by those in [7]. In addition, we obtain information about the asymptotic behavior of $T_{\max}(\lambda f)$ as $\lambda \searrow 0$ and as $\lambda \nearrow \infty$.

Before stating the principal new result of this type, we complement our previous result of [27, Theorem 1.2, p. 508].

Theorem 1.9. *Let $a = 1$ in (1.4) and let the positive integer m and the real number α be such that*

$$1 \leq m \leq N, \quad 0 < \alpha < \frac{2}{N+m}.$$

Let $f \in C_0(\mathbb{R}^N)$ be anti-symmetric with respect to x_1, x_2, \dots, x_m , and suppose that there exist $t_0 > 0$ and $K > 0$ such that

$$|f(x)| \leq K |\partial_1 \partial_2 \cdots \partial_m G_{t_0}(x)|, \quad \forall x \in \mathbb{R}^N. \quad (1.19)$$

Moreover, assume that

$$K_0 = \int_{\mathbb{R}^N} x_1 x_2 \cdots x_m f(x) dx \neq 0. \quad (1.20)$$

Let $u_\lambda : [0, T_{\max}(\lambda f)) \rightarrow C_0(\mathbb{R}^N)$ be the maximal solution of the integral equation (1.4) with initial value λf . By [27, Theorem 1.2, p. 508], $T_{\max}(\lambda f) < \infty$ for $\lambda > 0$ sufficiently small. In addition, we have

$$\lim_{\lambda \searrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{N+m}{2})^{-1}]} T_{\max}(\lambda f) = T_{\max}(u_0),$$

where $u_0 = K_0(-1)^m \partial_1 \partial_2 \cdots \partial_m \delta$, and $T_{\max}(u_0) < \infty$ is the maximal existence time of the solution u of (1.4) with initial data u_0 , constructed by [27, Theorem 1.1, p. 506].

The previous result for $m = 0$ is obtained in [7, Theorem 1.4, p. 307].

We now consider initial data having polynomial decay at infinity. In [15] the asymptotic behavior of the life span $T_{\max}(\lambda f)$ is studied for small $\lambda > 0$ and with $f \in C_b(\mathbb{R}^N)$, $f \geq 0$ under the condition $\alpha \leq 2/N$. (In this case $T_{\max}(\lambda f) < \infty$ by the Fujita results.) It is shown

in [15, Theorem 3.15 (ii), p. 375] and [15, Theorem 3.21 (ii), p. 376], among other results, that if $\alpha < 2/N$, $0 < \gamma < N$, and

$$0 < \liminf_{|x| \rightarrow \infty} |x|^\gamma f(x) \leq \limsup_{|x| \rightarrow \infty} |x|^\gamma f(x) < \infty, \quad (1.21)$$

then

$$0 < \liminf_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma}{2})^{-1}]} T_{\max}(\lambda f) \leq \limsup_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma}{2})^{-1}]} T_{\max}(\lambda f) < \infty.$$

In [19, Theorems 2.3 and 2.4, p. 1436] it is shown, for a certain class of f (not necessarily positive) which satisfy the condition (1.13) with $0 < \gamma < 2/\alpha$ on the same type of “thin” domain used in [19, Theorem 2.1, p. 1435] cited earlier, that $T_{\max}(\lambda f) < \infty$ for all $\lambda > 0$ and

$$\lambda^\varepsilon \leq \lambda^{[(\frac{1}{\alpha} - \frac{\gamma}{2})^{-1}]} T_{\max}(\lambda f) \leq \lambda^{-\varepsilon} \text{ as } \lambda \searrow 0,$$

for all $\varepsilon > 0$. In [7, Theorem 1.3, p. 307] it shown that if $0 < \gamma < N$, $\alpha < 2/\gamma$, $(N-2)\alpha < 4$ and $\lim_{|x| \rightarrow \infty} |x|^\gamma f(x) = c \neq 0$ (without imposing that $f \geq 0$) then

$$\lim_{\lambda \searrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma}{2})^{-1}]} T_{\max}(\lambda f) = C > 0,$$

where $C > 0$ is a finite constant.

The following two theorems extend and improve some of the above mentioned results.

Theorem 1.10 (new Dickstein-type blow-up results). *Let $a = 1$ in (1.4) and let the integer m and the real numbers α , γ be such that*

$$0 \leq m \leq N, \quad 0 < \gamma < N, \quad 0 < \alpha < \frac{2}{\gamma + m}, \quad (N-2)\alpha < 4.$$

Let $f \in C_0(\mathbb{R}^N)$, not necessarily positive, be anti-symmetric with respect to x_1, x_2, \dots, x_m , and suppose that

$$0 < \liminf_{|x| \rightarrow \infty, x \in \Omega} \frac{|x|^{\gamma+2m}}{x_1 \cdots x_m} f(x) \leq \limsup_{|x| \rightarrow \infty, x \in \Omega} \frac{|x|^{\gamma+2m}}{x_1 \cdots x_m} f(x) < \infty. \quad (1.22)$$

Then there exists $\lambda_0 > 0$ such that for all $0 < \lambda < \lambda_0$, the maximal solution $u_\lambda : [0, T_{\max}(\lambda f)) \rightarrow C_0(\mathbb{R}^N)$ of the integral equation (1.4) with initial value λf blows up in finite time. Moreover,

$$0 < \liminf_{\lambda \searrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f) \leq \limsup_{\lambda \searrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f) < \infty. \quad (1.23)$$

Remarks 1.11.

- (i) Dickstein proved the first small- λ blow-up results in [7]. Theorem 1.1 in [7] assumes $\alpha < 2/N$ with $f \in L^1(\mathbb{R}^N) \cap C_0(\mathbb{R}^N)$ and $\int f \neq 0$. Subsequent results of this type were obtained in [10] for $\alpha < 1/(N+1)$ for $f \in L^1(\mathbb{R}^N) \cap C_0(\mathbb{R}^N)$ and $\int x_1 f \neq 0$ and later in [27] for $\alpha < 2/(N+m)$ and $f \in L^1(\mathbb{R}^N) \cap C_0(\mathbb{R}^N)$ with $\int x_1 \cdots x_m f \neq 0$.

- (ii) Theorem 1.10 broadens the range of initial data in the case $m = 0$ as treated in [7, Theorems 1.2 and 1.3, pp. 306-307] (but not for the single point blowup result in [7, Theorem 1.3, p. 307]). Furthermore, it extends the results to the cases $1 \leq m \leq N$. Theorem 1.10 also improves the range of allowable α compared to [27] since we may take $2/(N+m) \leq \alpha < 2/(\gamma+m)$. Note that both [7, Theorem 1.2, p. 306] and Theorem 1.10 above allow values of $\alpha > 2/N$.
- (iii) The hypothesis $(N-2)\alpha < 4$ implies type I blowup, a condition used in the proof in an essential way. This hypothesis also is needed in [7, Theorems 1.2 and 1.3, pp. 306-307].
- (iv) If $\alpha > \frac{2}{N}$ and $\varphi \in L^{q_c}(\mathbb{R}^N)$, $q_c = \frac{N\alpha}{2}$, then by [29, Theorem 3(b)], $T_{\max}(\lambda\varphi) = \infty$ for λ sufficiently small. (The positivity condition in [29] is clearly not needed.) If we consider Theorem 1.10 above in the case $\alpha > \frac{2}{N}$, it is clear that $f \notin L^{q_c}(\mathbb{R}^N)$. Indeed, for f to be in $L^{q_c}(\mathbb{R}^N)$ would require that $N(\gamma+m)\alpha/2 > N$, i.e. $\alpha > 2/(\gamma+m)$.
- (v) Let $\alpha > \frac{2}{\gamma+m}$. There exist $f \in C_0(\mathbb{R}^N)$, anti-symmetric with respect to x_1, x_2, \dots, x_m , satisfying (1.22) such that $T_{\max}(\lambda f) = \infty$ for λ sufficiently small. This follows by Proposition 2.8 below and a reflection argument. If in addition, $\alpha > 2/N$, it also follows since such f are in $L^{q_c}(\mathbb{R}^N)$ as described in the part (iv) above.

The following result gives a refinement of Theorem 1.10.

Theorem 1.12. *Let $a = 1$ in (1.4) and let the integer m and the real numbers α, γ be such that*

$$0 \leq m \leq N, 0 < \gamma < N, 0 < \alpha < \frac{2}{\gamma+m}, (N-2)\alpha < 4.$$

Let $f \in C_0(\mathbb{R}^N)$, not necessarily positive on Ω , be anti-symmetric with respect to x_1, x_2, \dots, x_m , and suppose that

$$|f(x)| \leq K x_1 x_2 \cdots x_m |x|^{-\gamma-2m}, \quad x \in \Omega, \quad (1.24)$$

for some $K > 0$. Suppose there exist a sequence $\mu_n \rightarrow \infty$ and $z \in L^1_{\text{loc}}(\mathbb{R}^N \setminus \{0\})$, anti-symmetric with respect to x_1, x_2, \dots, x_m , such that

$$\mu_n^{\gamma+m} D_{\mu_n} f|_{\Omega} \rightarrow z|_{\Omega} \quad (1.25)$$

in $\mathcal{D}'(\Omega)$, as $n \rightarrow \infty$, so that in particular, z satisfies the same estimate (1.24). Let $T_{\max}(z)$ be the maximal existence time of the solution to (1.4) with $u_0 = z$, as constructed in Theorem 1.5.

It follows that

$$\lambda_n^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda_n f) \rightarrow T_{\max}(z) \quad (1.26)$$

where $\lambda_n = \mu_n^{-[\frac{2}{\alpha} - (\gamma+m)]} \rightarrow 0$. In particular, if $T_{\max}(z) < \infty$, then $T_{\max}(\lambda_n f) < \infty$ for all sufficiently large n .

The following result shows that the second inequality in (1.23) may be strict.

Corollary 1.13. *Let $a = 1$. Let the integer m and the real numbers α, γ be such that*

$$0 \leq m \leq N, 0 < \gamma < N, 0 < \alpha < \frac{2}{\gamma + m}, (N - 2)\alpha < 4.$$

Then there exists $f \in C_0(\mathbb{R}^N)$ satisfying the hypotheses of Theorem 1.10 such that

$$\liminf_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f) < \limsup_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f).$$

Remarks 1.14.

(i) As noted in Remark 1.8(ii), it follows from the considerations in [5, Section 2] and [21, Sections 3 and 5] that under condition (1.24), for any sequence $\mu_n \rightarrow \infty$ there is a subsequence, which we still call μ_n , and $z \in L^1_{\text{loc}}(\mathbb{R}^N \setminus \{0\})$, anti-symmetric with respect to x_1, x_2, \dots, x_m , such that (1.25) holds. In fact, one can easily choose f which admit two such limits z_1 and z_2 with $0 < T_{\max}(z_1) < T_{\max}(z_2) < \infty$. See [5, Theorem 1.2 and Proposition 2.9] and [21, Theorem 1.4]. For such a function f , $\lim_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f)$ does not exist and takes different finite nonzero values for different sequences $\lambda_n \rightarrow \infty$. See the proof of Corollary 1.13 in Section 4 for more details.

(ii) The situation described in Corollary 1.13 is quite different from previously known lifespan results. For example, if $f \in C_b(\mathbb{R}^N)$, $f \geq 0$, and $\lim_{|x| \rightarrow \infty} f(x) = f_\infty > 0$, then

$$\lim_{\lambda \rightarrow 0} \lambda^\alpha T_{\max}(\lambda f) = \frac{1}{\alpha} f_\infty^{-\alpha}.$$

Furthermore, when $f \in C_b(\mathbb{R}^N)$, $f \geq 0$, $f \not\equiv 0$, then

$$\lim_{\lambda \rightarrow \infty} \lambda^\alpha T_{\max}(\lambda f) = \frac{1}{\alpha} \|f\|_\infty^{-\alpha}.$$

For these two results, see [13, Theorem 1, p. 167]. See Proposition 4.6 below for a result analogous to Corollary 1.13, but with $\lambda \rightarrow \infty$, and for which $T_{\max}(\lambda f) \sim \lambda^{-(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}}$ as $\lambda \rightarrow \infty$, instead of $\lambda^{-\alpha}$.

(iii) If $\frac{2}{N} < \alpha < \frac{2}{\gamma+m}$, there exists f satisfying the hypotheses of Theorem 1.12, $f|_\Omega \geq 0$, for which two functions z_1 and $z_2 \not\equiv 0$ can be obtained as limits as in (1.25) and such that $0 < T_{\max}(z_1) < T_{\max}(z_2) = \infty$. In this case, there is a sequence $\lambda_n \rightarrow 0$ for which

$$\lambda_n^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda_n f) \rightarrow T_{\max}(z_1) < \infty,$$

and another sequence $\lambda'_n \rightarrow 0$ for which

$$(\lambda'_n)^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda'_n f) \rightarrow \infty,$$

Moreover, since $f \geq 0$ on Ω , it follows from Theorem 1.7 that $T_{\max}(\lambda f) < \infty$ for all $\lambda > 0$, and it must be a monotone function of $\lambda > 0$ (by comparison). This example shows that $\lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f)$ need not be monotone, and can exhibit qualitatively different behavior along different sequences of $\lambda \rightarrow 0$.

To justify the existence of such z_1 and z_2 , see [5, Lemma 2.7], as well as the proof of [5, Theorem 1.2, p. 1114] and the proof of [21, Theorem 1.4, p. 358]. In order to guarantee that $T_{\max}(z_2) = \infty$, it suffices to require that $\|z_2\|_{L^{q_c}}$ be small, where $q_c = N\alpha/2$.

(iv) It is shown in [7, Theorem 1.3, p. 307] that if $\lim_{|x| \rightarrow \infty} |x|^\gamma f(x) = c > 0$ then

$$\lim_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma}{2})^{-1}]} T_{\max}(\lambda f)$$

exists and is finite. This is a special case of Theorem 1.12 for $m = 0$. For $1 \leq m \leq N$ Theorem 1.12 implies that if $\lim_{|x| \rightarrow \infty} \frac{|x|^{\gamma+2m}}{x_1 x_2 \dots x_m} f(x) = c > 0$ then

$$\lim_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f)$$

exists and is finite.

(v) It is possible that, using the ideas described in [6, Section 7], one could construct initial values f for which $\lambda^\sigma T_{\max}(\lambda f)$ admits different finite, nonzero limits along different subsequences $\lambda_n \rightarrow 0$, and for different values of $\sigma > 0$.

See also Remark 4.4 below.

We close the introduction with a few words about the proofs. As mentioned earlier, the basic approach is to study equation (1.4) on the domain Ω defined by (1.7), and then to reformulate the results for solutions on \mathbb{R}^N which are anti-symmetric with respect to x_1, x_2, \dots, x_m .

Theorem 1.1 and Theorem 1.9 follow from straightforward adaptations of the methods in [27], using in addition in the case of Theorem 1.9, that the blowing up solutions are type I [11, 12] and the continuity of the maximal time of existence shown in [8].

To prove the existence results Theorems 1.2 and 1.5, the methods of [27] do not directly apply. Since the distribution (1.8) is in fact 0 as an element of $\mathcal{D}'(\Omega)$, in [27] the integral equation (1.4) was studied on Ω in the form

$$u(t) = \varphi(t) + \int_0^t e^{(t-\sigma)\Delta} [|u(\sigma)|^\alpha u(\sigma)] d\sigma, \tag{1.27}$$

for $t > 0$, without any specified initial value. In this paper, we need to make explicit the initial values described in Theorem 1.5, (in terms of their restriction to Ω). The key hypothesis is that

$$\int_0^T \|e^{t\Delta_\Omega} u_0\|_\infty^\alpha dt < \infty, \tag{1.28}$$

where u_0 is given by (1.9), but restricted to Ω , and $e^{t\Delta_\Omega}$ is the heat semigroup in $C_0(\Omega)$. This condition enables one to do a contraction mapping argument to prove existence of solutions to the integral equation (1.4) on Ω which are bounded by a multiple of $e^{t\Delta_\Omega} u_0$. We make extensive use of the results in [21], which is devoted to the study of $e^{t\Delta_\Omega}$ on the set of initial values of Theorem 1.5 (again, restricted to Ω).

As for the blowup results, the proof of Theorem 1.7 is a relatively straightforward application of [31, Theorem 1] along with [21, Theorem 1.1(ii)], or in the case $m = 0$, [5, Proposition 3.8(i)]. It is interesting to note that the current paper, along with [15, 26, 19], use four essentially different methods to obtain the same type of results. In [15], Kaplan's eigenvalue blowup method on a centered ball is used; in [26], self-similar subsolutions are used; and in [19] energy methods with self-similar variables are used.

The proofs of Theorem 1.10 and Theorem 1.12 follow the basic plan developed in [7], using some results from [11, 12, 8] on type I blowup and the continuity of the blowup time. In addition, the local theory of Theorem 1.5, including the continuous dependence result Theorem 3.5, below, is essential.

The rest of this paper is organized as follows. In Section 2, we prove well-posedness on sectors Ω for the nonlinear heat equation (1.4) with initial data bounded by the distribution (1.3). See the definition of the space \mathcal{X} given by (2.6) below. We also establish some nonexistence, blowup and global existence results. See Theorem 2.6 and Proposition 2.8 below. In Section 3 we apply the results of Section 2 to solutions on \mathbb{R}^N which are anti-symmetric in $x_1, x_2, x_3, \dots, x_m$. In particular, we establish the local existence with initial data of the form (1.3) and prove Theorems 1.1, 1.2 and 1.5. In Section 4, we prove Theorems 1.7, 1.9, 1.10, 1.12 and Corollary 1.13.

2. WELL-POSEDNESS ON SECTORS

The purpose of this section is to study well-posedness of the integral equation (1.4) on the sectors $\Omega \subset \mathbb{R}^N$ defined by (1.7), with the integer $m \in [1, N]$. To be more precise, we first recall that the heat semigroup on Ω , which we denote $e^{t\Delta_\Omega}$, is given by

$$e^{t\Delta_\Omega} f(x) = \int_{\Omega} K_t(x, y) f(y) dy, \quad t > 0, \quad (2.1)$$

where

$$K_t(x, y) = (4\pi t)^{-\frac{N}{2}} \prod_{j=m+1}^N e^{-\frac{|x_j - y_j|^2}{4t}} \prod_{i=1}^m \left[e^{-\frac{|x_i - y_i|^2}{4t}} - e^{-\frac{|x_i + y_i|^2}{4t}} \right]. \quad (2.2)$$

See, for example, [27]. It is well-known that $e^{t\Delta_\Omega}$ is a C_0 contraction, positivity preserving semigroup on $C_0(\Omega)$, where

$$C_0(\Omega) = \left\{ f \in C(\overline{\Omega}), f(x) = 0, \forall x \in \partial\Omega, \text{ and } \lim_{|x| \rightarrow \infty} f(x) = 0 \right\}.$$

Thus, instead of (1.4), we write

$$u(t) = e^{t\Delta_\Omega} v_0 + a \int_0^t e^{(t-\sigma)\Delta_\Omega} (|u(\sigma)|^\alpha u(\sigma)) d\sigma, \quad (2.3)$$

where $a = \pm 1$ and $\alpha > 0$. The initial value v_0 is in $\mathcal{D}'(\Omega)$, the space of distributions on Ω , and belongs to a specific class of distributions defined below. We study solutions $u \in C((0, T]; C_0(\Omega))$ such that $u(t) \rightarrow v_0$ in $\mathcal{D}'(\Omega)$ as $t \searrow 0$.

In order to define the class of initial values $v_0 = \psi$ which we consider in (2.3), we recall the following notation and definitions from [21] and [27]. For $0 < \gamma < N$ and integer $0 \leq m \leq N$, we let $\psi_0 : \Omega \rightarrow \mathbb{R}$ be given by

$$\psi_0(x) = c_{m,\gamma} x_1 \cdots x_m |x|^{-\gamma-2m}, \quad \forall x \in \Omega, \quad (2.4)$$

where

$$c_{m,\gamma} = \gamma(\gamma+2) \cdots (\gamma+2m-2). \quad (2.5)$$

In the case $m = 0$, then $c_{m,\gamma} = 1$ and $\psi_0(x) = |x|^{-\gamma}$, for $x \neq 0$. We are mainly interested in the case $m \geq 1$, but since some of our results are new for $m = 0$, i.e. $\Omega = \mathbb{R}^N$, we need to include $m = 0$ in the basic definitions. In what follows, some of the statements need to be modified in a trivial way if $m = 0$.

Clearly ψ_0 is homogenous of degree $-\gamma - m$, $\psi_0 \in C^\infty(\Omega)$ but $\psi_0 \notin C_0(\Omega)$, being singular at the origin. Also, $\psi_0(x) > 0$, $\forall x \in \Omega$. We define the Banach space

$$\mathcal{X} = \left\{ \psi \in L^1_{\text{loc}}(\Omega); \quad \frac{\psi}{\psi_0} \in L^\infty(\Omega) \right\}, \quad (2.6)$$

with the norm

$$\|\psi\|_{\mathcal{X}} = \left\| \frac{\psi}{\psi_0} \right\|_{L^\infty(\Omega)} \quad (2.7)$$

for all $\psi \in \mathcal{X}$, so that

$$|\psi| \leq \|\psi\|_{\mathcal{X}} \psi_0. \quad (2.8)$$

The set \mathcal{X} will be the space of initial data for the integral equation (2.3) on Ω . The norm used in [21] for \mathcal{X} differs from (2.7) by a factor. This has no effect on the results we will use from [21]. Note that with the choice of norm (2.7), we have $\|\psi_0\|_{\mathcal{X}} = 1$.

By [21, Theorem 1.1, p. 343] the heat semigroup $e^{t\Delta_\Omega}$ is well defined on \mathcal{X} . In particular, there exists a constant $\mathcal{C} \geq 1$ such that

$$\|e^{t\Delta_\Omega} \psi\|_{\mathcal{X}} \leq \mathcal{C} \|\psi\|_{\mathcal{X}}, \quad \forall t > 0, \quad \forall \psi \in \mathcal{X}. \quad (2.9)$$

Also $e^{t\Delta_\Omega}$ is a bounded map from \mathcal{X} to $C_0(\Omega)$ for all $t > 0$. Consider the Banach space

$$\mathcal{Y} := \left\{ \psi \in L^1_{\text{loc}}(\Omega); \quad \psi \psi_0 \in L^1(\Omega) \right\}. \quad (2.10)$$

Clearly, \mathcal{X} is the dual of \mathcal{Y} , that is $\mathcal{Y}' = \mathcal{X}$. Let

$$\mathcal{B}_K := \left\{ \psi \in \mathcal{X}; \quad \|\psi\|_{\mathcal{X}} \leq K \right\}, \quad (2.11)$$

where $K > 0$. Endowed with the weak* topology, \mathcal{B}_K is compact (see [2]) and metrizable since \mathcal{Y} is separable.

Definition 2.1. We denote by \mathcal{B}_K^* the ball \mathcal{B}_K endowed with the weak* topology on \mathcal{X} .

We know that \mathcal{B}_K^* is a compact metric space, hence complete and separable, for all $K > 0$. By [21, Proposition 3.1, p. 356], we also know that if $(\psi_n)_{n=1}^\infty$ is a sequence in \mathcal{B}_K and $\psi \in \mathcal{B}_K$, then $\psi_n \rightarrow \psi$ in \mathcal{B}_K^* if and only if $\psi_n \rightarrow \psi$ in $\mathcal{D}'(\Omega)$.

Remark 2.2. The function ψ_0 is simply the distribution u_0 given by (1.3), with $K = 1$, realized as a point function on Ω . The relationship between $e^{t\Delta_\Omega}\psi_0$ and $e^{t\Delta}u_0$ was studied in [21] and will be crucial in the next section to interpret the results of this section as results about (1.4) on \mathbb{R}^N .

We now introduce the space in which we will construct the solutions of (2.3). We set

$$\Psi(t, x) = e^{t\Delta_\Omega}\psi_0(x), \quad t > 0, \quad x \in \Omega, \quad (2.12)$$

and we sometimes write $\Psi(t) = e^{t\Delta_\Omega}\psi_0$. By [21, Theorem 1.1, p. 343], it follows that $\Psi : (0, \infty) \rightarrow C_0(\Omega)$ is a continuous map. Since $\psi_0(x) > 0$ for $x \in \Omega$, we also have

$$\Psi(t, x) > 0, \quad \forall t > 0, \quad \forall x \in \Omega.$$

Moreover, it follows from (2.8), (2.9) and (2.7) that

$$\Psi(t) = e^{t\Delta_\Omega}\psi_0 \leq \|e^{t\Delta_\Omega}\psi_0\|_{\mathcal{X}}\psi_0 \leq C\|\psi_0\|_{\mathcal{X}}\psi_0 = C\psi_0, \quad \forall t > 0. \quad (2.13)$$

For a given $T > 0$, let X be defined by

$$X = X_T = \left\{ u \in C((0, T]; C_0(\Omega)); \exists M > 0, |u(t, x)| \leq M\Psi(t, x), \forall t \in (0, T], \forall x \in \Omega \right\}.$$

The space $X = X_T$ equipped with the norm

$$\| \|u\| \| = \| \|u\| \|_{X_T} = \sup_{0 < t \leq T, x \in \Omega} \frac{|u(t, x)|}{\Psi(t, x)},$$

is a Banach space. Clearly,

$$|u(t)| \leq \| \|u\| \| e^{t\Delta_\Omega}\psi\Psi(t), \quad \text{for all } 0 < t \leq T. \quad (2.14)$$

If needed for clarity, we may denote X by X_T or $X_{T, \Psi}$ and the norm $\| \|u\| \|$ by $\| \|u\| \|_{X_T}$ or even $\| \|u\| \|_{X_{T, \Psi}}$. In addition, we observe that if $\psi \in \mathcal{X}$, then

$$\begin{aligned} \| \|e^{t\Delta_\Omega}\psi\| \|_{X_T} &= \sup_{0 < t \leq T, x \in \Omega} \frac{|e^{t\Delta_\Omega}\psi|}{\Psi(t, x)} = \sup_{0 < t \leq T, x \in \Omega} \frac{|e^{t\Delta_\Omega}(\frac{\psi}{\psi_0})\psi_0|}{\Psi(t, x)} \\ &\leq \sup_{0 < t \leq T, x \in \Omega} \frac{|e^{t\Delta_\Omega}\psi_0|}{\Psi(t, x)} \left\| \frac{\psi}{\psi_0} \right\|_{L^\infty(\Omega)} = \| \|\psi\| \|_{\mathcal{X}}. \end{aligned} \quad (2.15)$$

It also follows from [21, Theorem 1.1, p. 343] that $\|\Psi(t)\|_{L^\infty(\Omega)} \leq Ct^{-(\gamma+m)/2}$ for all $t > 0$. In fact, by [21, Relation (1.15), p. 344] and since ψ_0 is homogeneous of degree $-\gamma - m$, we have that

$$D_{\sqrt{t}}(e^{t\Delta_\Omega}\psi_0) = t^{-(\gamma+m)/2}e^{\Delta_\Omega}\psi_0, \quad \forall t > 0, \quad (2.16)$$

where the operator $D_{\sqrt{t}}$ is given by (1.14), so that

$$\|\Psi(t)\|_{L^\infty(\Omega)} = \|e^{t\Delta_\Omega}\psi_0\|_{L^\infty(\Omega)} = \|D_{\sqrt{t}}(e^{t\Delta_\Omega}\psi_0)\|_{L^\infty(\Omega)} = \|e^\Delta\psi_0\|_{L^\infty(\Omega)}t^{-(\gamma+m)/2}, \quad \forall t > 0. \quad (2.17)$$

It follows, and this is fundamental in our existence proofs, that for any $0 < \alpha < 2/(\gamma + m)$,

$$\int_0^A \|\Psi(t)\|_{L^\infty(\Omega)}^\alpha dt < \infty, \quad \forall A > 0. \quad (2.18)$$

We can now state and prove our existence result for (2.3) with more general singular initial data.

Theorem 2.3 (well-posedness in \mathcal{X}). *Let the integer m and the real numbers α, γ be such that*

$$0 \leq m \leq N, \quad 0 < \gamma < N, \quad 0 < \alpha < \frac{2}{\gamma + m}.$$

Let $K > 0$, $M > 0$ and $T > 0$ be such that

$$K + 2(\alpha + 1)M^{\alpha+1} \int_0^T \|\Psi(\sigma)\|_{L^\infty(\Omega)}^\alpha d\sigma \leq M \quad (2.19)$$

and

$$2(\alpha + 1)M^\alpha \int_0^T \|\Psi(\sigma)\|_{L^\infty(\Omega)}^\alpha d\sigma < 1, \quad (2.20)$$

where Ψ is given by (2.12). Let $a = \pm 1$. For every $\psi \in \mathcal{X}$ with $\|\psi\|_{\mathcal{X}} \leq K$ there exists a unique solution $u \in X_T$ of

$$u(t) = e^{t\Delta_\Omega}\psi + a \int_0^t e^{(t-\sigma)\Delta_\Omega} (|u(\sigma)|^\alpha u(\sigma)) d\sigma \quad (2.21)$$

such that

$$\|u\|_{X_T} \leq M.$$

The integral in (2.21) is absolutely convergent in $C_0(\Omega)$.

Furthermore, we have the following additional properties.

- (i) $u(t) \rightarrow \psi$ as $t \rightarrow 0$ in $L^1_{\text{loc}}(\Omega)$, hence in $\mathcal{D}'(\Omega)$ and in $\mathcal{B}_{M'}^*$, where $M' = CM$, and C is the constant appearing in (2.9).
- (ii) Let ψ_1 and ψ_2 be in \mathcal{B}_K . Let $u_1 \in X_T$, with $\|u_1\|_{X_T} \leq M$, and $u_2 \in X_T$, with $\|u_2\|_{X_T} \leq M$, be the solutions of (2.21) with respectively initial values ψ_1 and ψ_2 . It follows that

$$\|u_1 - u_2\|_{X_T} \leq C \|e^{t\Delta_\Omega}(\psi_1 - \psi_2)\|_{X_T} \leq C \|\psi_1 - \psi_2\|_{\mathcal{X}}, \quad (2.22)$$

where $C > 0$ is the constant given by (2.25) below.

- (iii) Let $(\psi_k)_{k=1}^\infty \subset \mathcal{B}_K^*$ and $\psi \in \mathcal{B}_K^*$. Let $u_k \in X_T$, with $\|u_k\|_{X_T} \leq M$, be the corresponding solution of (2.21) with initial data ψ_k and let $u \in X_T$, with $\|u\|_{X_T} \leq M$ be the corresponding solution of (2.21) with initial data ψ . If $\psi_k \rightarrow \psi$ as $k \rightarrow \infty$ in $\mathcal{D}'(\Omega)$ (that is also in \mathcal{B}_K^*) then $u_k(t) \rightarrow u(t)$ as $k \rightarrow \infty$ in $C_0(\Omega)$ for all $t \in (0, T]$, and uniformly in $[t_0, T]$, $\forall t_0 \in (0, T)$.
- (iv) Let $\psi \in \mathcal{B}_K^*$, let u be the solution of (2.21) with initial value ψ and v be the solution of (2.21) with initial value $|\psi|$. Then

$$|u(t)| \leq v(t), \quad \forall t \in (0, T].$$

(v) If $\psi \geq 0$ then $u(t) \geq 0$ on Ω for all $t \in (0, T]$.

- (vi) Let ψ_1 and ψ_2 be in \mathcal{B}_K . Let u_1 and u_2 be the solutions of (2.21) with respectively initial values ψ_1 and ψ_2 . If $\psi_1 \leq \psi_2$ then $u_1(t) \leq u_2(t)$, $\forall t \in (0, T]$.

Proof. The proof uses and adapts some arguments of [27] in conjunction with some results established in [21, 5]. For simplicity, we consider only the case $m \geq 1$. For $m = 0$, the results cited below in [21] need to be replaced by corresponding results in [5]. Let $K > 0$, $M > 0$ and $T > 0$ satisfying (2.19) and (2.20). Let

$$B(0, M) := \{u \in X_T : \|u\|_{X_T} \leq M\},$$

which is a complete metric space, with the metric

$$\|u - v\|_{X_T}.$$

Fix $\psi \in \mathcal{X}$ such that $\|\psi\|_{\mathcal{X}} \leq K$. Define, for $u \in B(0, M)$,

$$\mathcal{F}_\psi u(t) = e^{t\Delta_\Omega} \psi + a \int_0^t e^{(t-\sigma)\Delta_\Omega} (|u(\sigma)|^\alpha u(\sigma)) d\sigma.$$

We will show that \mathcal{F}_ψ is a contraction on $B(0, M)$. (For simplicity, we denote $\|\cdot\|_{L^\infty(\Omega)}$ by $\|\cdot\|_\infty$.) Let $\psi_1, \psi_2 \in \mathcal{B}_K$ and let $u_1, u_2 \in B(0, M)$. Then

$$\mathcal{F}_{\psi_1}(u_1)(t) - \mathcal{F}_{\psi_2}(u_2)(t) = e^{t\Delta_\Omega}(\psi_1 - \psi_2) + a \int_0^t e^{(t-\sigma)\Delta_\Omega} (|u_1|^\alpha u_1(\sigma) - |u_2|^\alpha u_2(\sigma)) d\sigma.$$

Using the fact that $e^{t\Delta_\Omega}$ is positivity preserving, and (2.14), we get that

$$\begin{aligned}
|\mathcal{F}_{\psi_1}(u_1)(t) - \mathcal{F}_{\psi_2}(u_2)(t)| &\leq |e^{t\Delta_\Omega}(\psi_1 - \psi_2)| + \\
&\quad (\alpha + 1) \int_0^t e^{(t-\sigma)\Delta_\Omega} \left(|u_1(\sigma) - u_2(\sigma)| (|u_1(\sigma)|^\alpha + |u_2(\sigma)|^\alpha) \right) d\sigma \\
&\leq e^{t\Delta_\Omega} |\psi_1 - \psi_2| + (\alpha + 1) \int_0^t \left(e^{(t-\sigma)\Delta_\Omega} \Psi(\sigma) \frac{|u_1(\sigma) - u_2(\sigma)|}{\Psi(\sigma)} \right) \times \\
&\quad \left(\|u_1(\sigma)\|_\infty^\alpha + \|u_2(\sigma)\|_\infty^\alpha \right) d\sigma \\
&\leq e^{t\Delta_\Omega} |\psi_1 - \psi_2| + (\alpha + 1) \int_0^t \left(e^{(t-\sigma)\Delta_\Omega} \Psi(\sigma) \right) \| |u_1 - u_2| \|_{X_T} \times \\
&\quad \left(\|u_1(\sigma)\|_\infty^\alpha + \|u_2(\sigma)\|_\infty^\alpha \right) d\sigma \\
&\leq e^{t\Delta_\Omega} |\psi_1 - \psi_2| + \Psi(t) \left(2M^\alpha (\alpha + 1) \int_0^t \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right) \| |u_1 - u_2| \|_{X_T}. \quad (2.23)
\end{aligned}$$

Let now $\psi_1 = \psi \in \mathcal{B}_K$, $u_1 = u \in B(0, M)$ and $\psi_2 = 0$, $u_2 = 0$. Then, by the previous inequality and (2.8), we get

$$\begin{aligned}
|\mathcal{F}_\psi u(t)| &\leq e^{t\Delta_\Omega} |\psi| + \Psi(t) \left(2M^\alpha (\alpha + 1) \int_0^t \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right) \| |u| \|_{X_T} \\
&\leq e^{t\Delta_\Omega} \|\psi\|_{\mathcal{X}} \psi_0 + \Psi(t) \left(2M^{\alpha+1} (\alpha + 1) \int_0^t \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right) \\
&\leq K \Psi(t) + \Psi(t) \left(2M^{\alpha+1} (\alpha + 1) \int_0^t \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right) \\
&\leq \Psi(t) \left[K + 2M^{\alpha+1} (\alpha + 1) \int_0^t \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right].
\end{aligned}$$

That is,

$$\frac{|\mathcal{F}_\psi u(t)|}{\Psi(t)} \leq K + 2M^{\alpha+1} (\alpha + 1) \int_0^t \|\Psi(\sigma)\|_\infty^\alpha d\sigma, \quad \forall t \in (0, T].$$

Thus by (2.19), $\| |\mathcal{F}_\psi u| \|_{X_T} \leq M$, and hence \mathcal{F}_ψ maps $B(0, M)$ into itself.

Let now $\psi_1 = \psi_2 = \psi \in \mathcal{B}_K$ in (2.23). This gives

$$\begin{aligned}
\| |\mathcal{F}_\psi(u_1) - \mathcal{F}_\psi(u_2)| \|_{X_T} &= \sup_{0 < t \leq T} \frac{|\mathcal{F}_\psi(u_1)(t) - \mathcal{F}_\psi(u_2)(t)|}{\Psi(t)} \\
&\leq \left(2(\alpha + 1) M^\alpha \int_0^T \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right) \| |u_1 - u_2| \|_{X_T}.
\end{aligned}$$

Then by (2.20), \mathcal{F} is a strict contraction from $B(0, M)$ into itself. It follows, by the Banach fixed point theorem that there exists a unique solution $u \in X_T$ of the integral equation (2.21) such that $\| |u| \|_{X_T} \leq M$.

Let u be the solution of the integral equation constructed by the above fixed point argument. Then

$$\begin{aligned} \int_0^t \|e^{(t-\sigma)\Delta_\Omega} |u(\sigma)|^\alpha u(\sigma)\|_\infty d\sigma &\leq \int_0^t \left\| (e^{(t-\sigma)\Delta_\Omega} |u(\sigma)|) \|u(\sigma)\|_\infty^\alpha \right\|_\infty d\sigma \\ &\leq M^{\alpha+1} \|\Psi(t)\|_\infty \int_0^t \|\Psi(\sigma)\|_\infty^\alpha d\sigma, \end{aligned} \quad (2.24)$$

which shows that the integral in (2.21) is absolutely convergent in $C_0(\Omega)$.

We now prove the additional properties.

(i) By (2.14) and (2.13),

$$|u(t)| \leq \|u\|_{X_T} \Psi(t) \leq M\mathcal{C}\psi_0, \quad \text{for all } 0 < t \leq T.$$

Hence $u(t) \in \mathcal{X}$ for all $t \in (0, T]$ and $\|u(t)\|_{\mathcal{X}} \leq M\mathcal{C}$, where \mathcal{C} is given by (2.9). On the one hand we have that

$$u(t) - e^{t\Delta_\Omega}\psi = a \int_0^t e^{(t-\sigma)\Delta_\Omega} (|u(\sigma)|^\alpha u(\sigma)) d\sigma,$$

and, again using (2.14) and (2.13),

$$\begin{aligned} \left| \int_0^t e^{(t-\sigma)\Delta_\Omega} (|u(\sigma)|^\alpha u(\sigma)) d\sigma \right| &\leq M^{\alpha+1} \int_0^t e^{(t-\sigma)\Delta_\Omega} (\Psi(\sigma) |\Psi(\sigma)|^\alpha) d\sigma \\ &\leq M^{\alpha+1} \int_0^t (e^{(t-\sigma)\Delta_\Omega} \Psi(\sigma)) \|\Psi(\sigma)\|_\infty^\alpha d\sigma \\ &= M^{\alpha+1} \Psi(t) \int_0^t \|\Psi(\sigma)\|_\infty^\alpha d\sigma \\ &\leq \mathcal{C} M^{\alpha+1} \psi_0 \int_0^t \|\Psi(\sigma)\|_\infty^\alpha d\sigma \rightarrow 0, \quad \text{as } t \rightarrow 0, \end{aligned}$$

in $L_{\text{loc}}^1(\Omega)$, hence in $\mathcal{D}'(\Omega)$. On the other hand, by [21, Proposition 2.5, p. 353] and Remark 2.4 below, $e^{t\Delta_\Omega}\psi \rightarrow \psi$ as $t \rightarrow 0$ in $L_{\text{loc}}^1(\Omega)$, hence in $\mathcal{D}'(\Omega)$. We thus conclude that $u(t) \rightarrow \psi$ as $t \rightarrow 0$ in $L_{\text{loc}}^1(\Omega)$, hence in $\mathcal{D}'(\Omega)$. The convergence in $\mathcal{B}_{M'}^*$ follows by [21, Proposition 3.1 (i), p. 356].

(ii) By (2.23), (2.14) and the fact that $\mathcal{F}_{\psi_1}(u_1) = u_1$ and $\mathcal{F}_{\psi_2}(u_2) = u_2$, we have that

$$\begin{aligned} |u_1(t) - u_2(t)| &\leq |e^{t\Delta_\Omega}(\psi_1 - \psi_2)| + \Psi(t) \left(2(\alpha + 1) M^\alpha \int_0^T \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right) \|u_1 - u_2\|_{X_T} \\ &\leq \Psi(t) \|e^{t\Delta_\Omega}(\psi_1 - \psi_2)\|_{X_T} + \Psi(t) \left(2(\alpha + 1) M^\alpha \int_0^T \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right) \|u_1 - u_2\|_{X_T} \\ &= \Psi(t) \left[\|e^{t\Delta_\Omega}(\psi_1 - \psi_2)\|_{X_T} + \left(2(\alpha + 1) M^\alpha \int_0^T \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right) \|u_1 - u_2\|_{X_T} \right]. \end{aligned}$$

Consequently,

$$\| \|u_1 - u_2\| \|_{X_T} \leq \| \|e^{t\Delta_\Omega}(\psi_1 - \psi_2)\| \|_{X_T} + \left(2(\alpha + 1)M^\alpha \int_0^T \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right) \| \|u_1 - u_2\| \|_{X_T}.$$

Thus, by (2.20), the first inequality in (2.22) follows with

$$C = \frac{1}{1 - \left(2(\alpha + 1)M^\alpha \int_0^T \|\Psi(\sigma)\|_\infty^\alpha d\sigma \right)}. \quad (2.25)$$

The second inequality in (2.22) then follows from (2.15).

(iii) Since $u_k \in X_T$, with $\| \|u_k\| \|_{X_T} \leq M$, it follows from [21, Lemma 2.6, p. 355] that

$$|u_k(t, x)| \leq M\Psi(t, x) \leq C(t + |x|^2)^{-\frac{\gamma+m}{2}}, \quad t \in (0, T], x \in \Omega, \quad (2.26)$$

for some $C > 0$ independent of k . In particular, for any fixed $t \in (0, T]$, the functions $u_k(t)$ are uniformly bounded in $C_0(\Omega)$. It follows by standard parabolic regularity arguments, along with (2.26), that the sequence $(u_k)_{k=1}^\infty$ is relatively compact in $C((t_0, T]; C_0(\Omega))$, for any $0 < t_0 < T$. By a diagonal argument, it follows that there exist $v \in C((0, T]; C_0(\Omega))$, and a subsequence u_{k_n} such that $u_{k_n} \rightarrow v$ in $C((t_0, T]; C_0(\Omega))$, for all $0 < t_0 < T$. It follows that $v \in X_T$, with $\| \|v\| \|_{X_T} \leq M$.

We next observe that v satisfies the integral equation (2.21). This follows easily since

$$u_{k_n}(t) = e^{t\Delta_\Omega}\psi_{k_n} + a \int_0^t e^{(t-\sigma)\Delta_\Omega} (|u_{k_n}(\sigma)|^\alpha u_{k_n}(\sigma)) d\sigma,$$

for all $0 < t \leq T$, and letting $n \rightarrow \infty$. Indeed, $e^{t\Delta_\Omega}\psi_{k_n} \rightarrow e^{t\Delta_\Omega}\psi$ as $n \rightarrow \infty$, by [21, Proposition 4.1(ii), p. 359]. Moreover, by the dominated convergence theorem, since

$$\| \|e^{(t-\sigma)\Delta_\Omega}|u_{k_n}(\sigma)|^\alpha u_{k_n}(\sigma)\| \|_\infty \leq M^{\alpha+1} \|\Psi(t)\|_\infty \|\Psi(\sigma)\|_\infty^\alpha,$$

as in (2.24), the integral terms converge to $a \int_0^t e^{(t-\sigma)\Delta_\Omega} (|v(\sigma)|^\alpha v(\sigma)) d\sigma$.

Since $v \in X_T$, $\| \|v\| \|_{X_T} \leq M$, and v satisfies the integral equation (2.21), it follows from uniqueness that $v = u$, so $u_{k_n} \rightarrow u$ on $C((t_0, T]; C_0(\Omega))$, for any $0 < t_0 < T$. A standard subsequence argument shows that in fact $u_k \rightarrow u$ on $C((t_0, T]; C_0(\Omega))$, for any $0 < t_0 < T$.

(iv) By [21, Proposition 3.2, p. 357], we know that $\mathcal{B}_K \cap \mathcal{D}(\Omega)$ is dense in \mathcal{B}_K^* . Let $(\psi_n)_{n=1}^\infty \subset \mathcal{B}_K \cap \mathcal{D}(\Omega_m)$ such that $\psi_n \rightarrow \psi$ as $n \rightarrow \infty$ in $\mathcal{B}_{M'}^*$, hence in $\mathcal{D}'(\Omega)$. In fact, the proof of [21, Proposition 3.2, p. 357] shows that we may choose $(\psi_n)_{n=1}^\infty$ so that in addition $\psi_n \rightarrow \psi$ as $n \rightarrow \infty$ in $L_{\text{loc}}^1(\Omega)$. Thus, $(|\psi_n|) \subset \mathcal{B}_K \cap C_0(\Omega)$ and $|\psi_n| \rightarrow |\psi|$ as $n \rightarrow \infty$ in $L_{\text{loc}}^1(\Omega)$, hence in $\mathcal{D}'(\Omega)$ and \mathcal{B}_K^* . Let u_n the solution of (2.3) with initial data $\psi_n \in C_0(\Omega)$ and let v_n be the solution of (2.3) with initial data $|\psi_n| \in C_0(\Omega)$. It is well known, by the $C_0(\Omega)$ local well-posedness and comparison argument, that

$$|u_n(t)| \leq v_n(t), \quad \forall t \in (0, T].$$

Letting $n \rightarrow \infty$ we get by Part (iii) that $u_n(t) \rightarrow u(t)$, $v_n(t) \rightarrow v(t)$ in $C_0(\Omega)$ for $t \in (0, T]$ and hence

$$|u(t)| \leq v(t), \quad \forall t \in (0, T].$$

(v) This is an immediate consequence of Part (iv).

(vi) The proof is analogous to the proof of Part (iv). \square

Remark 2.4. There is a small error in the proof of [21, Proposition 2.5, p. 353]. It claims to show that if $\psi \in \mathcal{X}$ then $e^{t\Delta_\Omega}\psi \rightarrow \psi$ a.e. as $t \rightarrow 0$. The proof in fact shows that $e^{t\Delta_\Omega}\psi \rightarrow \psi$ in $L^1_{\text{loc}}(\Omega)$ as $t \rightarrow 0$, which is indeed sufficient to guarantee convergence in $\mathcal{D}'(\Omega)$. The authors of [21] regret this error.

Remark 2.5. By well-posedness of the integral equation in $C_0(\Omega)$, the solution u of (2.3) with initial data $\psi \in \mathcal{B}_K$, constructed in Theorem 2.3 above, can be continued to a maximal solution $u \in C((0, T_{\max}(\psi)); C_0(\Omega))$. If $T_{\max}(\psi) < \infty$ then $\lim_{t \rightarrow T_{\max}(\psi)} \|u(t)\|_\infty = \infty$.

We have also obtained the following result.

Theorem 2.6. *Assume the hypotheses of Theorem 2.3. Let $\psi \in \mathcal{B}_K$ and let $u \in C((0, T_{\max}(\psi)); C_0(\Omega))$ be the solution of (2.3) with initial data ψ constructed by Theorem 2.3, extended to its maximal existence time. Then we have the following.*

- (i) If $a = -1$ then $T_{\max}(\psi) = \infty$.
- (ii) If $a = 1$, then there exists $c > 0$ such that

$$\lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda\psi) > c \quad (2.27)$$

for all $\lambda > 0$.

- (iii) If $a = 1$, then $T_{\max}(\psi_0) < \infty$. Moreover,

$$\lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda\psi_0) = T_{\max}(\psi_0) \quad (2.28)$$

for all $\lambda > 0$.

Proof of Theorem 2.6. (i) If $a = -1$, we claim that

$$|u(t)| \leq e^{t\Delta_\Omega} |\psi| \quad \text{for all } 0 < t < T_{\max}(\psi). \quad (2.29)$$

To prove this, let $0 < t_0 < T$, where T is as in the fixed point argument. Since $u(t_0) \in C_0(\Omega)$, we know (by Kato's parabolic inequality) that $|u(t+t_0)| \leq e^{t\Delta_\Omega} |u(t_0)|$ for all $0 < t < T_{\max}(\psi) - t_0$. We next let $t_0 \rightarrow 0$, so that $u(t_0) \rightarrow \psi$, and likewise $|u(t_0)| \rightarrow |\psi|$, in $L^1_{\text{loc}}(\Omega)$ by Theorem 2.3(i), and therefore in \mathcal{B}_K^* . Finally, by [21, Proposition 4.1, p. 359], it follows that $e^{t\Delta_\Omega} |u(t_0)| \rightarrow e^{t\Delta_\Omega} |\psi|$ in $C_0(\Omega)$ as $t \rightarrow 0$, which establishes (2.29). Global existence now follows since $e^{t\Delta_\Omega} |\psi| \in C_0(\Omega)$ for all $t > 0$.

(ii) It is clear that $T_{\max}(\psi)$ is bigger than the T obtained from the fixed point argument. Hence, by condition (2.19), which implies (2.20), we must have that

$$K + 2(\alpha + 1)M^{\alpha+1} \left(1 - \frac{\gamma + m}{2}\alpha\right)^{-1} \|e^{\Delta_{\Omega}}\psi_0\|_{\infty}^{\alpha} T_{\max}(\psi)^{1 - \frac{\gamma+m}{2}\alpha} > M.$$

for all $M > K$, where we have used (2.17) to simplify $\int_0^T \|\Psi(\sigma)\|_{L^{\infty}(\Omega)}^{\alpha} d\sigma$. Likewise, it must be that

$$\lambda K + 2(\alpha + 1)M^{\alpha+1} \left(1 - \frac{\gamma + m}{2}\alpha\right)^{-1} \|e^{\Delta_{\Omega}}\psi_0\|_{\infty}^{\alpha} T_{\max}(\lambda\psi)^{1 - \frac{\gamma+m}{2}\alpha} > M.$$

for all $M > \lambda K$, for all $\lambda > 0$. If we set $M = 2\lambda K$, we obtain the desired relationship.

(iii) We first show that $T_{\max}(\psi_0) < \infty$. Let u be the solution of (1.4) with initial data ψ_0 , extended to its maximal solution in $C_0(\Omega)$. Since $\psi_0 \geq 0$, u is positive on Ω . By [31, Theorem 1, p. 546], we have that

$$\|e^{t\Delta_{\Omega}}\psi_0\|_{\infty} \leq (\alpha t)^{-1/\alpha}, \quad t \in (0, T_{\max}(\psi_0)). \quad (2.30)$$

Since $\alpha < 2/(\gamma + m)$, formula (2.17) implies that $T_{\max}(\psi_0) < \infty$.

The formula for $T_{\max}(\lambda\psi_0)$ follows by a rescaling argument, using the fact that ψ_0 is homogeneous of degree $-(\gamma + m)$. More precisely, we know that

$$\lambda\psi_0(x) = \mu^{2/\alpha}\psi_0(\mu x),$$

where $\lambda = \mu^{\frac{2}{\alpha} - (\gamma + m)}$. Furthermore, if $u(t, x)$ is the maximal solution of (2.3) with initial value ψ_0 , then $\mu^{2/\alpha}u(\mu^2 t, \mu x)$ is the maximal solution with initial value $\mu^{2/\alpha}\psi_0(\mu \cdot)$, so that

$$T_{\max}(\lambda\psi_0) = T_{\max}(\mu^{2/\alpha}\psi_0(\mu \cdot)) = \mu^{-2}T_{\max}(\psi_0) = \lambda^{-[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]}T_{\max}(\psi_0). \quad (2.31)$$

□

Remark 2.7. It is well-known that a fixed point argument used to prove local existence to a nonlinear integral equation can yield lower estimates for the blowup rate. The idea is to observe that the fixed point argument can not work on the interval $[0, T_{\max} - t_0]$ with initial value $u(t_0)$. To our knowledge, this idea was first used in [29, Section 4 and Remark 6(2)] and then formalized more clearly in [22, Proposition 5.3, p. 901]. We are unable to apply this method in the present context since it is not clear that the solution with initial data in \mathcal{X} will remain in \mathcal{X} throughout its entire $C_0(\Omega)$ trajectory. On the other hand, as the proof of Theorem 2.6(ii) shows, a similar idea can be used to obtain lower estimates for $T_{\max}(\lambda u_0)$, for all $\lambda > 0$. The idea is to observe that $T_{\max}(\lambda u_0)$ is larger than the T obtained by the fixed point argument, and then apply the fixed point argument to initial values which are multiples λu_0 of a given function u_0 . This idea is developed more fully in [28].

We have the following proposition showing in particular that the local existence result of Theorem 1.2 is optimal. We also give initial data which give rise to global solutions for $\alpha > 2/(\gamma + m)$.

Proposition 2.8. *Let the positive integer m and the real numbers α, γ be such that*

$$1 \leq m \leq N, \quad 0 < \gamma < N, \quad \alpha > 2/(\gamma + m).$$

Then the following holds.

- (i) *Assume $a = 1$. Let $\psi \in \mathcal{X}$ such that $\psi \geq 0$ and there exist $c > 0$ and $L > 0$ such that*

$$\psi(x) \geq c\psi_0(x), \quad \forall x \in \Omega, \quad |x| \leq L.$$

Then there is no local nonnegative solution of (2.3) on Ω with initial data ψ .

- (ii) *Assume $a = \pm 1$. Let $\psi \in \mathcal{X}$. Assume that there exists $t_0 > 0$ such that*

$$|\psi| \leq e^{t_0 \Delta_\Omega} \psi_0.$$

Then for $\lambda > 0$ sufficiently small there exists a global solutions u of (2.3) on Ω with initial data $\lambda\psi$. Moreover, there exists $M > 0$, such that

$$|u(t, x)| \leq M e^{(t+t_0)\Delta_\Omega} \psi_0(x), \quad \forall t > 0, \quad \forall x \in \Omega.$$

Remarks 2.9.

- (i) The case $m = 0$ of (i) in the previous proposition is known, see [23, Theorem 15.2, p. 76 and pp. 85-86].
- (ii) In [16, 17, 18, 14] global solutions of (2.3) are proved to exist in the case $\alpha > 2/(N+m)$. The initial data giving global existence, considered in these papers, are dominated by Gaussian functions. In the part (ii) of the previous proposition, the initial data may have polynomial decay.
- (iii) It is known that if $\alpha > 2/N$ and for $\psi \in C_0(\mathbb{R}^N)$, and

$$|\psi(x)| \leq C(1 + |x|^2)^{-\sigma/2}, \quad \sigma > 2/\alpha \tag{2.32}$$

with C is sufficiently small, then the solution of (1.4) with initial value $u_0 = \psi$ is global. See [29, Theorem 3, p. 32]. Indeed, $\psi \in L^{q_c}(\mathbb{R}^N)$ with $q_c = N\alpha/2 > 1$. By the previous proposition part (ii), using [21, Lemma 2.6, p. 355], the initial value ψ satisfies (2.32) with $\sigma = \gamma + m > 2/\alpha$. But here we do not impose the condition $\alpha > 2/N$. In other words, this proposition allows the possibility that $\psi \in L^{q_c}(\mathbb{R}^N)$ with $q_c = N\alpha/2 \leq 1$.

Proof of Proposition 2.8. (i) Assume that a local nonnegative solution u of (1.4) exists on $(0, T]$ with initial data ψ . Then we have

$$u(t) = e^{t\Delta_\Omega} \psi + \int_0^t e^{(t-\sigma)\Delta_\Omega} (u(\sigma)^{\alpha+1}) d\sigma$$

on $(0, T] \times \Omega$. As in the proof of Theorem 2.6 (iii), by [31, Theorem 1, p. 546], we have that

$$e^{t\Delta_\Omega} \psi \leq (\alpha t)^{-1/\alpha}, \quad \text{on } (0, T] \times \Omega, \quad (2.33)$$

for all $T < T_{\max}(\psi)$. Write

$$\psi = \psi 1_{\{|x| \leq L\}} + \psi 1_{\{|x| > L\}}.$$

Then

$$t^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega} \psi\|_\infty \geq t^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega}(\psi 1_{\{|x| \leq L\}})\|_\infty - t^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega}(\psi 1_{\{|x| > L\}})\|_\infty. \quad (2.34)$$

Using the relation (2.16), we get first that

$$\begin{aligned} t^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega}(\psi 1_{\{|x| \leq L\}})\|_\infty &\geq ct^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega}(\psi_0 1_{\{|x| \leq L\}})\|_\infty \\ &= ct^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega}(\psi_0 1_{\{|x| \leq L\}})(\sqrt{t} \cdot)\|_\infty \\ &= c \|e^{\Delta_\Omega}(\psi_0 1_{\{|x| \leq \frac{L}{\sqrt{t}}\}})\|_\infty \rightarrow c \|e^{\Delta_\Omega} \psi_0\|_\infty \quad \text{as } t \rightarrow 0, \end{aligned} \quad (2.35)$$

and second, we have

$$\begin{aligned} t^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega}(\psi 1_{\{|x| > L\}})\|_\infty &\leq \|\psi\|_{\mathcal{X}} t^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega}(\psi_0 1_{\{|x| > L\}})\|_\infty \\ &= \|\psi\|_{\mathcal{X}} t^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega}(\psi_0 1_{\{|x| > L\}})(\sqrt{t} \cdot)\|_\infty \\ &= \|\psi\|_{\mathcal{X}} \|e^{\Delta_\Omega}(\psi_0 1_{\{|x| > \frac{L}{\sqrt{t}}\}})\|_\infty \rightarrow 0 \quad \text{as } t \rightarrow 0. \end{aligned} \quad (2.36)$$

The convergence results follow by [21, Proposition 4.1 (ii), p. 359]. Then for $t > 0$ sufficiently small, by (2.34), (2.35) and (2.36), there exists $C > 0$ a constant such that

$$\|e^{t\Delta_\Omega} \psi\|_\infty \geq Ct^{-\frac{\gamma+m}{2}}. \quad (2.37)$$

By (2.30) and (2.37), we have

$$0 < C \leq t^{(\gamma+m)/2} \|e^{t\Delta_\Omega} \psi\|_\infty \leq (\alpha t)^{-1/\alpha} t^{(\gamma+m)/2}, \quad t \in (0, T].$$

This leads to a contradiction for t sufficiently small and $\alpha > 2/(\gamma + m)$. Then there is no local nonnegative solution.

(ii) Let $t_0 > 0$ and $\psi \in \mathcal{X}$ be as in the statement of the proposition. Let

$$\Psi_{t_0}(t) = \Psi(t_0 + t), \quad t > 0.$$

Then by (2.17), we have

$$\|\Psi_{t_0}(t)\|_\infty = \|\Psi(t_0 + t)\|_\infty = \|e^{\Delta_\Omega} \psi_0\|_\infty (t_0 + t)^{-(\gamma+m)/2}, \quad \forall t > 0.$$

Hence, for any $\alpha > 2/(\gamma + m)$,

$$\int_0^\infty \|\Psi_{t_0}(t)\|_\infty^\alpha dt < \infty. \quad (2.38)$$

That is, (2.18) is satisfied by Ψ_{t_0} with $A = \infty$, in place of Ψ . We will apply Theorem 2.3 with Ψ replaced by Ψ_{t_0} and $T = \infty$. Let $K > 0$, $M > 0$ be such that

$$K + 2(\alpha + 1)M^{\alpha+1} \int_0^\infty \|\Psi_{t_0}(\sigma)\|_\infty^\alpha d\sigma \leq M$$

and

$$2(\alpha + 1)M^\alpha \int_0^\infty \|\Psi_{t_0}(\sigma)\|_\infty^\alpha d\sigma < 1.$$

Let $\lambda > 0$ be such that $\|\lambda\psi\|_{\mathcal{X}} \leq K$. Then by Theorem 2.3, there exists a unique solution $u \in X_{\infty, \Psi_{t_0}}$ of (2.3) on $(0, \infty)$ such that

$$\|u\|_{X_{\infty, \Psi_{t_0}}} \leq M.$$

This proves the existence of global solution of (2.3) satisfying the estimate

$$|u(t)| \leq M\Psi_{t_0}(t), \quad \forall t > 0,$$

on Ω . This finishes the proof of the proposition. \square

3. WELL-POSEDNESS ON THE WHOLE SPACE

In this section we use the results of the previous section and an anti-symmetric reflection argument to construct a solution of the nonlinear heat equation (1.4) on \mathbb{R}^N with singular initial values. In particular, we prove Theorems 1.1, 1.2 and 1.5. As in [27], for $1 \leq i \leq N$, let T_i be the operator defined on the space of functions f by

$$[T_i f](x_1, \dots, x_{i-1}, x_i, x_{i+1}, \dots, x_N) = f(x_1, \dots, x_{i-1}, -x_i, x_{i+1}, \dots, x_N).$$

A function f is anti-symmetric in x_1, \dots, x_m if it satisfies

$$T_1 f = T_2 f = \dots = T_m f = -f. \quad (3.1)$$

We denote the space of functions which are anti-symmetric in x_1, \dots, x_m by

$$\mathcal{A} := \{f : \mathbb{R}^N \rightarrow \mathbb{R}; f \text{ satisfies (3.1)}\}. \quad (3.2)$$

Definition 3.1. Let $\Omega \subset \mathbb{R}^N$ be the domain given by (1.7), and suppose $g : \bar{\Omega} \rightarrow \mathbb{R}$ is such that $g|_{\partial\Omega} \equiv 0$. We denote by $\tilde{g} : \mathbb{R}^N \rightarrow \mathbb{R}$ the anti-symmetric extension of g , i.e. $\tilde{g} \in \mathcal{A}$ and $\tilde{g}|_\Omega = g$.

If $g \in C_0(\Omega)$, then $\tilde{g} \in C_0(\mathbb{R}^N)$. It is proved in [27, formula (3.6), p. 514] that

$$e^{t\Delta_\Omega} g = (e^{t\Delta} \tilde{g})|_{\bar{\Omega}}, \quad \forall t > 0, \quad \forall g \in C_0(\mathbb{R}^N). \quad (3.3)$$

Moreover, the commutation relation (1.15) is valid with $e^{t\Delta_\Omega}$ instead of $e^{t\Delta}$, and in particular

$$D_{\sqrt{t}} e^{t\Delta_\Omega} = e^{\Delta_\Omega} D_{\sqrt{t}} \quad (3.4)$$

for all $t > 0$, where the dilation operators D_λ are given by (1.14).

The following proposition is immediate.

Proposition 3.2. *Let $g \in C_0(\Omega)$ and let $v : [0, T_{\max}(g)) \rightarrow C_0(\Omega)$ be the maximal solution of (2.3) with initial value $v_0 = g$. It follows that $\tilde{v} : [0, T_{\max}(g)) \rightarrow C_0(\mathbb{R}^N) \cap \mathcal{A}$ is the maximal solution of (1.4) with initial value $u_0 = \tilde{g} \in C_0(\mathbb{R}^N)$. In particular, $T_{\max}(g) = T_{\max}(\tilde{g})$.*

Similarly, let $f \in C_0(\mathbb{R}^N) \cap \mathcal{A}$ and let $u : [0, T_{\max}(f)) \rightarrow C_0(\mathbb{R}^N)$ be the maximal solution of (1.4) with initial value $u_0 = f$. It follows that $u(t) \in \mathcal{A}$ for all $0 \leq t < T_{\max}(f)$ and that $u|_{\overline{\Omega}} : [0, T_{\max}(f)) \rightarrow C_0(\Omega)$ is the maximal solution of (2.3) with initial value $v_0 = f|_{\overline{\Omega}} \in C_0(\Omega)$. In particular, $T_{\max}(f) = T_{\max}(f|_{\overline{\Omega}})$.

The analogous result is clearly true in other spaces where the two integral equations are well-defined. The point of this paper is to use the solutions on Ω to construct solutions on \mathbb{R}^N , which is not quite the same as the above proposition.

If $\psi \in \mathcal{X}$, given by (2.6), it is proved in [21, Definition 1.6, p. 346 and Proposition 5.1, p. 361] that the pointwise anti-symmetric extension $\tilde{\psi}$ of ψ to \mathbb{R}^N has a natural interpretation as a tempered distribution. It is natural to define

$$\tilde{\mathcal{X}} = \left\{ \varphi \in L^1_{\text{loc}}(\mathbb{R}^N \setminus \{0\}) \cap \mathcal{A}, \varphi|_{\Omega} \in \mathcal{X} \right\} = \{ \tilde{\psi} \mid \psi \in \mathcal{X} \} \subset \mathcal{S}'(\mathbb{R}^N). \quad (3.5)$$

and

$$\widetilde{\mathcal{B}}_K := \{ \tilde{\psi} \mid \psi \in \mathcal{B}_K \}, \quad (3.6)$$

where $\mathcal{B}_K, K > 0$, is given by (2.11). Also $\widetilde{\mathcal{B}}_K$ inherits the metric space structure from \mathcal{B}_K^* (recall Definition 2.1). If $m = 0$, then $\tilde{\mathcal{X}} = \mathcal{X}$ and $\widetilde{\mathcal{B}}_K = \mathcal{B}_K$.

We recall the following result from [21] and [5] showing the equivalence of various kinds of convergence.

Proposition 3.3 (Propositions 3.1(i) and 5.1 in [21] and Proposition 2.1(i) in [5]). *Let $0 < \gamma < N$ and m an integer with $0 \leq m \leq N$, and let $(\psi_k)_{k=1}^{\infty} \subset \mathcal{B}_K$ and $\psi \in \mathcal{B}_K$. The following are equivalent:*

- (i) $\psi_k \rightarrow \psi$ in \mathcal{B}_K^* as $k \rightarrow \infty$;
- (ii) $\psi_k \rightarrow \psi$ in $\mathcal{D}'(\Omega)$ as $k \rightarrow \infty$;
- (iii) $\tilde{\psi}_k \rightarrow \tilde{\psi}$ in $\mathcal{S}'(\mathbb{R}^N)$ as $k \rightarrow \infty$.

Remark 3.4. If $(\psi_k)_{k=1}^{\infty} \subset \mathcal{B}_K$, $\psi \in L^1_{\text{loc}}(\Omega)$, and $\psi_k \rightarrow \psi$ in $\mathcal{D}'(\Omega)$ as $k \rightarrow \infty$, then in fact $\psi \in \mathcal{B}_K$ and the above conditions hold. This follows either from the compactness of \mathcal{B}_K^* (so that the $(\psi_k)_{k=1}^{\infty}$ must have a convergent subsequence in \mathcal{B}_K^*), or from the fact that convergence in $\mathcal{D}'(\Omega)$ preserves a uniform bound.

Moreover, ([21, Proposition 5.1]),

$$\widetilde{e^{t\Delta}\psi} = e^{t\Delta}\tilde{\psi}, \quad \forall t > 0, \forall \psi \in \mathcal{X}.$$

For $0 < \gamma < N$, we put

$$\varphi_0 = (-1)^m \partial_1 \partial_2 \cdots \partial_m (|\cdot|^{-\gamma}) \in \mathcal{S}'(\mathbb{R}^N). \quad (3.7)$$

By [21, Proposition 2.1 p. 347], it follows that

$$\tilde{\psi}_0 = \varphi_0, \quad (3.8)$$

in $\mathcal{S}'(\mathbb{R}^N)$.

By (3.7), it follows that

$$\Phi(t) = e^{t\Delta}\varphi_0 = G_t \star \varphi_0 = \Phi_0(t) \star |\cdot|^{-\gamma} \quad (3.9)$$

$$= (-1)^m \partial_1 \partial_2 \cdots \partial_m (e^{t\Delta} |\cdot|^{-\gamma}), \quad (3.10)$$

in $\mathcal{S}'(\mathbb{R}^N)$, where Φ_0 is given by

$$\Phi_0(t) = (-1)^m \partial_1 \partial_2 \cdots \partial_m G_t, \quad t > 0. \quad (3.11)$$

We have that $\Phi : (0, \infty) \rightarrow C_0(\mathbb{R}^N)$ is a solution of the linear heat equation such that $T_i \Phi(t) = -\Phi(t)$ for all $1 \leq i \leq m$. The key observation, to use the results of the previous section, is that, by [21, Proposition 2.1, p. 347], we have

$$\Phi(t)|_{\bar{\Omega}} = (e^{t\Delta}\varphi_0)|_{\bar{\Omega}} = e^{t\Delta\Omega}\psi_0 = \Psi(t), \quad \forall t > 0.$$

Also

$$\tilde{\Psi}(t) = \Phi(t), \quad \forall t > 0,$$

where $\tilde{\Psi}(t)$ is the unique anti-symmetric extension of $\Psi(t)$ to \mathbb{R}^N . In particular, we have

$$\|\Phi(t)\|_{L^\infty(\mathbb{R}^N)} = \|\Psi(t)\|_\infty, \quad \forall t > 0.$$

We now wish to re-formulate Theorem 2.3 in terms of a result on the whole space \mathbb{R}^N . This reformulation is possible thanks to [27, Proposition 3.3, p. 515], which is analogous to Proposition 3.2. The following theorem is thus an immediate consequence of Theorem 2.3 and Proposition 3.3. Of course in the case $m = 0$, the following result is the same as Theorem 2.3.

Theorem 3.5 (well-posedness in $\tilde{\mathcal{X}}$). *Let the positive integer m and the real numbers α, γ be such that*

$$0 \leq m \leq N, \quad 0 < \gamma < N, \quad 0 < \alpha < \frac{2}{\gamma + m}.$$

Let $K, M, T > 0$ be satisfying (2.19) and (2.20). Then for every $u_0 \in \tilde{\mathcal{X}}$ with $\|u_0\|_{\tilde{\mathcal{X}}} \leq K$, there exists a solution $u \in C((0, T]; C_0(\mathbb{R}^N) \cap \mathcal{A})$ of (1.4) such that $u(t) \rightarrow u_0$ in $\mathcal{S}'(\mathbb{R}^N)$ as $t \rightarrow 0$. Furthermore,

- (i) $|u(t)| \leq M e^{t\Delta}\psi_0, \quad \forall t \in [0, T], \quad \forall x \in \Omega.$
- (ii) *If $v \in C((0, T]; C_0(\mathbb{R}^N) \cap \mathcal{A})$ is a solution of (1.4) such that $|v(t)| \leq M e^{t\Delta}\psi_0, \quad \forall t \in [0, T], \quad \forall x \in \Omega,$ with $v(0) = u_0$ then $u(t) = v(t), \quad \forall t \in [0, T].$*
- (iii) *Let $(u_{0,n}) \subset \widetilde{\mathcal{B}}_K, \quad u_0 \in \widetilde{\mathcal{B}}_K$ with $u_{0,n} \rightarrow u_0$ in $\mathcal{S}'(\mathbb{R}^N)$. Let u_n be the solution of (1.4) with initial data $u_{0,n}$ and u be the solution of (1.4) with initial data u_0 . Then $u_n(t) \rightarrow u(t)$ in $C_0(\mathbb{R}^N), \quad \forall t \in (0, T],$ and uniformly in $[t_0, T], \quad \forall t_0 \in (0, T).$*

Moreover u can be continued to a maximal solution $u \in C((0, T_{\max}); C_0(\mathbb{R}^N))$ and $u(t) \in \mathcal{A}$ for all $0 < t < T_{\max}$.

We now give the proofs of Theorems 1.1, 1.2 and 1.5.

Proof of Theorem 1.1. The proof is the same as in [27, Theorem 1.1, p. 506] which is valid also for minus sign in front of the nonlinearity. The last statement in the theorem follows by the fact that in the absorption case the solutions are global. This completes the proof of the theorem. \square

Proof of Theorem 1.2. Theorem 1.2 is a particular case of Theorem 3.5 with $u_0 = K\varphi_0$. Here we fix K and we take M, T such that (2.19) and (2.20) are satisfied. Hence the proof of the existence follows. The others statements follows by Theorem 2.6. \square

Proof of Theorem 1.5. The proof follows by Theorem 3.5 with $u_0 \in \tilde{\mathcal{X}}$ and by Theorem 2.6. \square

4. BLOW-UP RESULTS

This section is devoted to the proofs of Theorems 1.7, 1.9, 1.10, 1.12 and Corollary 1.13. Let $K > 0$. For $\psi \in \mathcal{B}_K$, given by (2.11), where $0 < \gamma < N$ and m an integer with $0 \leq m \leq N$, we recall from [21, Formula (1.18), p. 344] for $1 \leq m \leq N$ and [5, Formula (1.17), p. 1107] for $m = 0$, the following definition

$$\mathcal{Z}(\psi) := \left\{ z \in \mathcal{B}_K, \exists \lambda_n \rightarrow \infty, n \rightarrow \infty, \text{ such that } \lim_{n \rightarrow \infty} \lambda_n^{\gamma+m} D_{\lambda_n} \psi = z \text{ in } \mathcal{B}_K^* \right\}, \quad (4.1)$$

where \mathcal{B}_K^* denotes \mathcal{B}_K endowed with respect to the weak* topology on \mathcal{X} and the dilation operator D_λ is given by (1.14). By [21, Proposition 3.1 (iii), p. 356] and [5, Remark 2.4(iii), p. 1112], $\mathcal{Z}(\psi)$ is a nonempty, compact, and connected subset of the compact metric space \mathcal{B}_K^* .

Also, it is important to remark that the conditions on $z \in L_{\text{loc}}^1(\mathbb{R}^N \setminus \{0\})$ in the statements of Theorems 1.7 and 1.12 can be interpreted to say that $z|_{\overline{\Omega}} \in \mathcal{Z}(f|_{\overline{\Omega}})$.

Proof of Theorem 1.7. This result concerns solutions in $C_0(\mathbb{R}^N) \cap \mathcal{A}$, where \mathcal{A} is defined by (3.2), i.e. solutions which are anti-symmetric in x_1, x_2, \dots, x_m . Moreover, the solutions in question are positive on Ω . It is clear from Proposition 3.2 that it suffices to consider the restrictions of these solutions to $\overline{\Omega}$, as solutions of (2.3), rather than as solutions on \mathbb{R}^N of (1.4). This will enable us to use the positivity.

To simplify the notation, we will use, *by abuse of notation*, the same letters to denote functions on \mathbb{R}^N and their restrictions to $\overline{\Omega}$. For example, if $f \in C_0(\mathbb{R}^N)$ is as in the statement of the theorem, and therefore in \mathcal{A} , in this proof f will denote its restriction to $\overline{\Omega}$, as an element of $C_0(\Omega)$. Also, we consider the case $m \geq 1$, the case $m = 0$ being entirely analogous.

Let u_0 and f be as in the statement of the theorem (restricted to $\overline{\Omega}$). By comparison, it suffice to show that $T_{\max}(f) < \infty$. We argue by contradiction and assume that $T_{\max}(f) = \infty$.

By [31, Theorem 1, p. 546], since $f \geq 0$, it follows that

$$\|e^{t\Delta_\Omega} f\|_{L^\infty(\Omega)} \leq (\alpha t)^{-1/\alpha}, \quad \forall t > 0. \quad (4.2)$$

Moreover, by (3.4) we have

$$\|e^{t\Delta_\Omega} f\|_{L^\infty(\Omega)} = \|D_{\sqrt{t}}(e^{t\Delta_\Omega} f)\|_{L^\infty(\Omega)} = \|e^{\Delta_\Omega} D_{\sqrt{t}} f\|_{L^\infty(\Omega)}.$$

Hence,

$$t^{\frac{\gamma+m}{2}} \|e^{t\Delta_\Omega} f\|_{L^\infty(\Omega)} = \left\| e^{\Delta_\Omega} \left(t^{\frac{\gamma+m}{2}} D_{\sqrt{t}} f \right) \right\|_{L^\infty(\Omega)}. \quad (4.3)$$

By the hypotheses of the theorem, $f \in C_0(\Omega) \cap \mathcal{B}_K$, where \mathcal{B}_K is defined by (2.11). Let $z \in L^1_{\text{loc}}(\Omega)$, $z \not\equiv 0$, and $\lambda_n \rightarrow \infty$ be such that

$$\lambda_n^{\gamma+m} D_{\lambda_n} f \rightarrow z$$

in $\mathcal{D}'(\Omega)$, as $n \rightarrow \infty$. We claim that $z \in \mathcal{Z}(f)$ and $\lambda_n^{\gamma+m} D_{\lambda_n} f \rightarrow z$ in \mathcal{B}_K^* , where \mathcal{Z} is defined in (4.1). In fact, since \mathcal{B}_K^* is compact, there exists a subsequence $(\lambda_{n_k})_{k=1}^\infty$ and ζ in \mathcal{B}_K^* such that $\lambda_{n_k}^{\gamma+m} D_{\lambda_{n_k}} f \rightarrow \zeta$ in \mathcal{B}_K^* . But convergence in \mathcal{B}_K^* implies convergence in $\mathcal{D}'(\Omega)$, so $z = \zeta \in \mathcal{B}_K^*$. It then follows from Proposition 3.3 that $\lambda_n^{\gamma+m} D_{\lambda_n} f \rightarrow z$ in \mathcal{B}_K^* as $n \rightarrow \infty$, and so $z \in \mathcal{Z}(f)$.

Let now $t_n = \lambda_n^2 \rightarrow \infty$, as $n \rightarrow \infty$. Then $t_n^{\frac{\gamma+m}{2}} D_{\sqrt{t_n}} f \rightarrow z$ in \mathcal{B}_M^* as $n \rightarrow \infty$. By [21, Proposition 4.1(ii), p. 359], it follows that

$$e^{t\Delta_\Omega} \left(t_n^{\frac{\gamma+m}{2}} D_{\sqrt{t_n}} f \right) \rightarrow e^{\Delta_\Omega} z, \quad \text{in } C_0(\Omega), \quad (4.4)$$

as $n \rightarrow \infty$, and so

$$\lim_{n \rightarrow \infty} t_n^{\frac{\gamma+m}{2}} \|e^{t_n \Delta_\Omega} f\|_{L^\infty(\Omega)} = \|e^{\Delta_\Omega} z\|_{L^\infty(\Omega)} \neq 0. \quad (4.5)$$

Using (4.2), we get

$$t_n^{\frac{\gamma+m}{2}} \|e^{t_n \Delta_\Omega} f\|_{L^\infty(\Omega)} \leq \alpha^{-1/\alpha} t_n^{\frac{\gamma+m}{2} - \frac{1}{\alpha}}, \quad \forall n > 0. \quad (4.6)$$

By (4.5) and (4.6), we see that for any $\varepsilon > 0$,

$$\|e^{\Delta_\Omega} z\|_{L^\infty(\Omega)} - \varepsilon < \alpha^{-1/\alpha} t_n^{\frac{\gamma+m}{2} - \frac{1}{\alpha}} \quad \text{for large } n.$$

Letting $n \rightarrow \infty$ in the last inequality, we get a contradiction if $\alpha < 2/(\gamma + m)$ with the fact that $\|e^{\Delta_\Omega} z\|_{L^\infty(\Omega)} > 0$. We also get a contradiction if $\alpha = 2/(\gamma + m)$ by letting $n \rightarrow \infty$, since in this case $\|e^{\Delta_\Omega} z\|_{L^\infty(\Omega)} > \alpha^{-1/\alpha}$. It follows that $T_{\max}(f) < \infty$. This completes the proof of Theorem 1.7. \square

The following remark, inspired by [5, Proposition 2.12, p. 1117], gives some examples of functions f satisfying the conditions of Theorem 1.7 such that $\mathcal{Z}(f|_{\overline{\Omega}}) \neq \{0\}$, i.e. for which an appropriate $z \in L^1_{\text{loc}}(\mathbb{R}^N \setminus \{0\}) \cap \mathcal{A}$ exists. In particular, it generalizes the example discussed in Remark 1.8(iii).

Note first that if $\gamma + m < N$ and $f \in L^p(\Omega)$ for some $1 \leq p < N/(\gamma + m)$ then $\mathcal{Z}(f) = \{0\}$. This follows by scaling argument. See [5, Proposition 2.5 (iii), p. 1112] for the case $m = 0$. Also, if $f \in C_0(\mathbb{R}^N) \cap \mathcal{A}$ has compact support, then $\mathcal{Z}(f) = \{0\}$.

Remark 4.1. Let $0 \leq m \leq N$ be an integer and $0 < \gamma < N$. Let $g \in C(\mathbb{R}, \mathbb{R})$ be a bounded function such that $g(t) \geq 0$ for $t \in \mathbb{R}$. Let $\zeta \in C(S^{N-1}) \cap \mathcal{A}$, and suppose that $f \in C_0(\mathbb{R}^N) \cap \mathcal{A}$ is such that $f_{|\bar{\Omega}} \geq 0$ and satisfies

$$f(x) = \psi_0(x)g(\log(|x|))\zeta\left(\frac{x}{|x|}\right), \quad x \in \Omega, |x| \geq \rho, \quad (4.7)$$

where ψ_0 is given by (2.4) and \mathcal{A} is given by (3.2). Then there exists $K > 0$ such that $f_{|\bar{\Omega}} \in C_0(\Omega) \cap \mathcal{B}_K$. Moreover, we have

$$\lambda^{(\gamma+m)}D_\lambda f(x) = \psi_0(x)g(\log(|x|) + \log \lambda)\zeta\left(\frac{x}{|x|}\right), \quad x \in \Omega, |x| \geq \rho/\lambda.$$

It is now a simple matter to choose a wide variety of functions g such that $\mathcal{Z}(f_{|\bar{\Omega}}) \neq \{0\}$. For example, if g is periodic, then the function z , given by expression in (4.7) is an element of $\mathcal{Z}(f)$.

Remark 4.2. We wish to point out that the proof of Theorem 1.7, which gives a genuine improvement of known blowup results for equation (1.1), including in the case $m = 0$, depends only on calculations with the (linear) heat semigroup. The first ingredient of the proof is [31, Theorem 1, p. 546], which gives a necessary condition on the heat semigroup for existence of nonnegative solutions to (1.1) on a certain time interval. The second ingredient is [21, Proposition 4.1(ii), p. 359], and [5, Proposition 3.8, p. 1123] in the case $m = 0$. In particular, the example discussed in Remark 1.8(iii) could have been observed after the appearance of [31] and [5].

Both Theorem 1.10 and Corollary 1.13 are consequences of Theorem 1.12. The key arguments of the proof of Theorem 1.12 i.e. the rescaling and the continuity of blowup times, are inspired by the paper [7]. Furthermore, in order to apply these methods, we need the well-posedness of the integral equation (1.4) on the space $\tilde{\mathcal{X}}$ given by (3.5), which was shown in Theorem 3.5. The following theorem is a slight reformulation and generalization of Theorem 1.12. (Recall Proposition 3.3 and Remark 3.4 above.)

Theorem 4.3. *Let the integer m and the real numbers α, γ be such that*

$$0 \leq m \leq N, \quad 0 < \gamma < N, \quad 0 < \alpha < \frac{2}{\gamma + m}, \quad (N - 2)\alpha < 4.$$

Let $u_0 \in \tilde{\mathcal{X}}$ and let $u \in C((0, T_{\max}(u_0)); C_0(\mathbb{R}^N) \cap \mathcal{A})$ be the maximal solution of (1.4) given by Theorem 3.5 such that $u(t) \rightarrow u_0$ in $\mathcal{S}'(\mathbb{R}^N)$ as $t \rightarrow 0$. Let $f \in \tilde{\mathcal{B}}_K$, for some $K > 0$. For

$\lambda > 0$, let $u^\lambda \in C((0, T_{\max}(\lambda f)); C_0(\mathbb{R}^N))$ be the maximal solution of (1.4) given by Theorem 3.5 such that $u(t) \rightarrow \lambda f$ in $\mathcal{S}'(\mathbb{R}^N)$ as $t \rightarrow 0$.

Suppose that there exists a sequence $(\tau_n)_{n=1}^\infty$, with $\tau_n > 0$, such that

$$\tau_n^{-\frac{\gamma+m}{2}} D_{\frac{1}{\sqrt{\tau_n}}} f \rightarrow u_0 \text{ in } \mathcal{S}'(\mathbb{R}^N) \text{ as } n \rightarrow \infty. \quad (4.8)$$

It follows that

$$\lim_{n \rightarrow \infty} \lambda_n^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda_n f) = T_{\max}(u_0), \quad (4.9)$$

where $\lambda_n = \tau_n^{\frac{1}{\alpha} - \frac{\gamma+m}{2}}$. In particular, if $T_{\max}(u_0) < \infty$ then $T_{\max}(\lambda_n f) < \infty$ for all n sufficiently large.

Furthermore, if $\tau_n \rightarrow 0$ as $n \rightarrow \infty$, then $\lambda_n \rightarrow 0$ as $n \rightarrow \infty$, and if $\tau_n \rightarrow \infty$ as $n \rightarrow \infty$, then $\lambda_n \rightarrow \infty$ as $n \rightarrow \infty$.

Proof. By the standard invariance properties of solutions to (1.1), for any $\tau > 0$,

$$u_\tau^\lambda(t, x) = \tau^{-1/\alpha} u^\lambda\left(\frac{t}{\tau}, \frac{x}{\sqrt{\tau}}\right) \quad (4.10)$$

is the maximal solution of (1.1) with initial value

$$u_{0,\tau}^\lambda = \lambda \tau^{-1/\alpha} D_{\frac{1}{\sqrt{\tau}}} f \in \widetilde{\mathcal{B}}_{\rho K}$$

in the sense of Theorem 3.5, where $\rho = \lambda \tau^{-(\frac{1}{\alpha} - \frac{\gamma+m}{2})}$. For $\tau > 0$ we let $\lambda = \lambda(\tau)$ be given by

$$\lambda = \lambda(\tau) = \tau^{\frac{1}{\alpha} - \frac{\gamma+m}{2}}.$$

With this choice of $\lambda = \lambda(\tau)$, we refer to the solutions $u_\tau^{\lambda(\tau)}$ as u_τ . It follows that the initial data of u_τ , is $u_{0,\tau} = f_\tau$ where

$$f_\tau := \tau^{-\frac{\gamma+m}{2}} D_{\frac{1}{\sqrt{\tau}}} f \in \widetilde{\mathcal{B}}_K. \quad (4.11)$$

By (4.10),

$$T_{\max}(u_{0,\tau}) = \tau T_{\max}(\lambda f) = \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f). \quad (4.12)$$

Moreover, $f_\tau, u_0 \in \widetilde{\mathcal{B}}_K$, given by (3.6).

Since $u_{0,\tau_n} = f_{\tau_n} \rightarrow u_0$ in $\mathcal{S}'(\mathbb{R}^N)$, by (4.8), it follows from Theorem 3.5 that there exists $T > 0$ such that $u_{\tau_n}(t) \rightarrow u(t)$ in $C_0(\mathbb{R}^N)$ for all $t \in (0, T]$. Fix $0 < \delta < T$ and let $u_{\tau_n,\delta}(t) = u_{\tau_n}(t + \delta)$ and $u_\delta(t) = u(t + \delta)$ with maximal existence times respectively $T_{\max}(u_{0,\tau_n}) - \delta$ and $T_{\max}(u_0) - \delta$. Since $(N-4)\alpha < 2$, we know by [11, 12] that blowing up solutions are always type I. Hence, by continuity of the maximal time of existence in $C_0(\mathbb{R}^N)$ for type I blow-up solutions (see [8]), in the case where $T_{\max}(u_0) < \infty$, and by the lower semi-continuity of the maximal time of existence in the case $T_{\max}(u_0) = \infty$, it follows that

$$T_{\max}(u_{\tau_n}(\delta)) \rightarrow T_{\max}(u(\delta)),$$

which implies that

$$T_{\max}(u_{0,\tau_n}) \rightarrow T_{\max}(u_0).$$

The relation (4.9) now follows from (4.12).

Finally, since $\alpha < 2/(\gamma + m)$, we have that as $\tau_n \rightarrow 0$ (respectively ∞) if and only if $\lambda_n \rightarrow 0$ (respectively ∞). \square

As mentioned above, Theorem 1.12 is included in Theorem 4.3, and has thus been proved. We next turn to Theorem 1.10.

Proof of Theorem 1.10. Let $f \in C_0(\mathbb{R}^N) \cap \mathcal{A}$ satisfying (1.22). It follows that there exist two constants, $0 < c < C < \infty$ and $L > 0$, such that

$$c\psi_0 \leq f(x) \leq C\psi_0, \quad \forall |x| \geq L, x \in \Omega,$$

where ψ_0 is given by (2.4), and so

$$c\psi_0 \leq f_\tau(x) \leq C\psi_0, \quad \forall |x| \geq L\sqrt{\tau}, x \in \Omega,$$

where f_τ is given by (4.11). Hence, if u_0 and τ_n satisfy the hypotheses of Theorem 4.3, and if $\tau_n \rightarrow 0$, we must have that

$$c\psi_0 \leq u_0(x) \leq C\psi_0, \quad \forall x \in \Omega.$$

By comparison, and Theorems 2.6 and 4.3, that

$$T_{\max}(C\psi_0) \leq T_{\max}(u_0) = \lim_{n \rightarrow \infty} \lambda_n^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda_n f) \leq T_{\max}(c\psi_0) < \infty, \quad (4.13)$$

where $\lambda_n = \tau_n^{(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}} \rightarrow 0$ as $n \rightarrow \infty$.

It follows from (4.13) that

$$T_{\max}(C\psi_0) \leq \liminf_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f) \leq \limsup_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f) \leq T_{\max}(c\psi_0) \quad (4.14)$$

To see this, suppose to the contrary there is a sequence $\lambda_n \rightarrow 0$ such that $\lambda_n^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda_n f)$ does not ultimately fall in the range given by (4.14). Passing to a subsequence, we may assume, by the compactness of $\mathcal{Z}(f)$, defined by (4.1), and by [21, Proposition 5.1 p. 361], that there exists u_0 such that the hypotheses of the Theorem 4.3 are satisfied with $\tau_n = \lambda_n^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]}$, and hence we obtain (4.13). This completes the proof of the theorem. \square

We now give the proof of Corollary 1.13.

Proof of Corollary 1.13. As in the proof of Theorem 1.7, it suffices to consider the restrictions of functions to $\bar{\Omega}$. By [5, Proposition 2.9, p. 1115] and [21, Section 3], given $0 < c_1 < c_2 \leq K < \infty$, there exists $f \in \mathcal{B}_K^* \cap C_0(\Omega)$, such that $\mathcal{Z}(f) = \{c\psi_0; c_1 \leq c \leq c_2\}$. (A result analogous to [5, Proposition 2.9, p. 1115] was not explicitly given in [21, Section 3] but is easily proved.) It follows that

$$\liminf_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f) = T_{\max}(c_2\psi_0) < T_{\max}(c_1\psi_0) = \limsup_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f),$$

where we have used formulas (1.26) and (2.28). \square

Remarks 4.4.

- (i) Let $f \in C_0(\mathbb{R}^N) \cap \mathcal{A}$ be a universal solution in as in [21, Theorem 1.4, p. 345], or [5, Theorem 1.2, p. 1108] in the case $m = 0$. More precisely, for some $K > 0$, $f|_\Omega := g \in \mathcal{B}_K \cap C^\infty(\Omega) \cap C_0(\Omega)$ and

$$\mathcal{Z}(g) = \mathcal{B}_K,$$

where $\mathcal{Z}(g)$ is defined by (4.1). Hence $\mathcal{Z}(f) = \widetilde{\mathcal{Z}(g)} = \widetilde{\mathcal{B}_K}$. It follows that for every $z \in \widetilde{\mathcal{B}_K}$ there exists a sequence $\mu_n \rightarrow \infty$ such that (1.25) holds. Thus, assuming the conditions on α in Theorem 1.12, it follows that

$$\lambda_n^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda_n f) \rightarrow T_{\max}(z)$$

where $\lambda_n = \mu_n^{-[\frac{2}{\alpha} - (\gamma+m)]} \rightarrow 0$.

This adds to the examples described in the previous remarks.

- (ii) If the initial data f in Theorem 1.10 satisfies also

$$\lim_{|x| \rightarrow \infty, x \in \Omega} \frac{f(x)}{\psi_0(x)} = c > 0,$$

Then

$$\mathcal{Z}(f|_\Omega) = \{c\psi_0\}.$$

Hence

$$\lim_{\lambda \searrow 0} \lambda^{(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}} T_{\max}(\lambda f) = T_{\max}(c\psi_0) < \infty.$$

In particular, the last limit exists and is finite. The case $m = 0$ is known, see [7, Theorem 1.3, p. 307]. See also [13, Theorem 2, p. 172] for the case $m = 0$, but only for a more restrictive class of initial data.

The following two results are analogues of Theorem 1.10 and Corollary 1.13 where the asymptotic behavior of $T_{\max}(\lambda f)$ as $\lambda \rightarrow \infty$ is studied, as opposed to $\lambda \rightarrow 0$.

Proposition 4.5. *Let $a = 1$ in (1.4) and let the integer m and the real numbers α, γ be such that*

$$0 \leq m \leq N, \quad 0 < \gamma < N, \quad 0 < \alpha < \frac{2}{\gamma + m}, \quad (N - 2)\alpha < 4.$$

Let $f \in \tilde{\mathcal{X}}$, not necessarily positive, be anti-symmetric with respect to x_1, x_2, \dots, x_m , and suppose that

$$0 < \liminf_{|x| \rightarrow 0, x \in \Omega} \frac{|x|^{\gamma+2m}}{x_1 \cdots x_m} f(x) \leq \limsup_{|x| \rightarrow 0, x \in \Omega} \frac{|x|^{\gamma+2m}}{x_1 \cdots x_m} f(x) < \infty. \quad (4.15)$$

Then there exists $\lambda_0 > 0$ such that for all $\lambda > \lambda_0$, the maximal solution $u_\lambda : (0, T_{\max}(\lambda f)) \rightarrow C_0(\mathbb{R}^N)$ of the integral equation (1.4) with initial value λf blows up in finite time. Moreover,

$$0 < \liminf_{\lambda \nearrow \infty} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f) \leq \limsup_{\lambda \nearrow \infty} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]} T_{\max}(\lambda f) < \infty. \quad (4.16)$$

Proof. The proof follows exactly the same steps as the proof of Theorem 1.10 with the obvious modifications. For example, the first inequality in the proof is true for $|x| \leq L$, instead of $|x| \geq L$, etc. □

Proposition 4.6. *Let $a = 1$. Let the integer m and the real numbers α, γ be such that*

$$0 \leq m \leq N, 0 < \gamma < N, 0 < \alpha < \frac{2}{\gamma + m}, (N - 2)\alpha < 4.$$

Then there exists $f \in \widetilde{\mathcal{X}}$ satisfying the hypotheses of Proposition 4.5 such that

$$\liminf_{\lambda \rightarrow \infty} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]T_{\max}(\lambda f)} < \limsup_{\lambda \rightarrow \infty} \lambda^{[(\frac{1}{\alpha} - \frac{\gamma+m}{2})^{-1}]T_{\max}(\lambda f)}.$$

Proof. The proof is same as that of Corollary 1.13, however taking into account Remark 2.4 (vi) of [5]. More precisely, if we set

$$\mathcal{J}f(x) = |x|^{-2(\gamma+m)} f(x/|x|^2)$$

for all $f \in \widetilde{\mathcal{B}}_K$, then \mathcal{J} is a homeomorphism of $\widetilde{\mathcal{B}}_K^*$, \mathcal{J} leaves invariant all functions which are homogeneous of degree $-(\gamma + m)$,

$$\lambda^{\gamma+m} D_\lambda \mathcal{J} = (1/\lambda)^{\gamma+m} \mathcal{J} D_{1/\lambda}$$

and

$$\mathcal{H}(f) = \mathcal{J}^{-1} \mathcal{Z}(\mathcal{J}(f)),$$

where

$$\mathcal{H}(f) := \left\{ z \in \mathcal{B}_K, \exists \lambda_n \rightarrow 0, n \rightarrow \infty, \text{ such that } \lim_{n \rightarrow \infty} \lambda_n^{\gamma+m} D_{\lambda_n} f = z \in \mathcal{B}_K^* \right\}. \quad (4.17)$$

Thus, there exists $f \in \widetilde{\mathcal{B}}_K^*$ such that $\mathcal{H}(f) = \{c\psi_0; c_1 \leq c \leq c_2\}$. □

Remarks 4.7.

- (i) The functions $f \in \widetilde{\mathcal{X}}$ satisfying the hypotheses of Propostion 4.5, and in particular $f \in \widetilde{\mathcal{X}}$ such that $\mathcal{H}(f) = \{c\psi_0; c_1 \leq c \leq c_2\}$, constructed in the previous proof, clearly can *not* be bounded in any neighborhood of 0.
- (ii) On the other hand, if $\gamma + m < N$, there exist $f \in \widetilde{\mathcal{X}}$ satisfying the hypotheses of Propostion 4.5 with $f \in L^q(\Omega)$, $q > 1$ and $q_c < q < \frac{N}{\gamma+m}$, where $q_c = \frac{N\alpha}{2}$. In other words, the ‘‘classical’’ theory of local wellposedness for (1.1) in L^q includes such initial values.

We conclude with the proof of Theorem 1.9.

Proof of Theorem 1.9. This result is just a small addition to [27, Theorem 1.2, p. 508] and the proof is almost completely given there. To complete the proof, it suffice to observe, *using the notation established in* [27, pp. 521-523], that $T_{\max}(v_\tau(t_0)) \rightarrow T_{\max}(u(t_0))$ as $\tau \rightarrow 0$. This

is true by continuity of the maximal time of existence in $C_0(\mathbb{R}^N)$ for type I blow-up solutions (see [8]). This implies, *continuing with the notation established in [27, pp. 521-523]*, that

$$\lim_{\lambda \rightarrow 0} \lambda^{[(\frac{1}{\alpha} - \frac{N+m}{2})^{-1}] T_{\max}(\lambda f)} = \lim_{\tau \rightarrow 0} T_{\max}(f_\tau) = T_{\max}(v_\tau(t_0)) + t_0 = T_{\max}(u(t_0)) + t_0 = T(u_0).$$

This finishes the proof of the theorem. \square

REFERENCES

- [1] C. Bandle and H. A. Levine, *On the existence and nonexistence of global solutions of reaction-diffusion equations in sectorial domains*, Trans. AMS, 316 (1989), 595–622.
- [2] H. Brézis, *Functional analysis, Sobolev spaces and partial differential equations*, Universitext, Springer, New York, 2011.
- [3] H. Brézis and A. Friedman, *Nonlinear parabolic equations involving measures as initial conditions*, J. Math. Pures Appl., 62 (1983), 73–97.
- [4] T. Cazenave, F. Dickstein, M. Escobedo and F. B. Weissler, *Self-similar solutions of a nonlinear heat equation*, J. Math. Sci. Univ. Tokyo, 8 (2001), 501–540.
- [5] T. Cazenave, F. Dickstein and F. B. Weissler, *Universal solutions of the heat equation in \mathbb{R}^N* , Discrete Contin. Dynam. Systems, 9 (2003), 1105–1132.
- [6] T. Cazenave, F. Dickstein and F. B. Weissler, *Multi-scale multi-profile global solutions of parabolic equations in \mathbb{R}^N* , Discrete Contin. Dynam. Systems - Series S, 5 (2012), 449–472.
- [7] F. Dickstein, *Blowup stability of solutions of the nonlinear heat equation with a large life span*, J. Differential Equations, 223 (2006), 303–328.
- [8] C. Fermanian Kammerer, F. Merle and H. Zaag, *Stability of the blow-up profile of nonlinear heat equations from the dynamical system point of view*, Math. Ann., 317 (2000), 347–387.
- [9] H. Fujita, *On the blowing up of solutions of the Cauchy problem for $u_t = \Delta u + u^{1+\alpha}$* , J. Fac. Sci. Univ. Tokyo Sect. I, 13 (1966), 109–124.
- [10] T. Ghoul, *An extension of Dickstein’s “small lambda” theorem for finite time blowup*, Nonlinear Analysis, T.M.A., 74 (2011), 6105–6115.
- [11] Y. Giga, S. Matsui and S. Sasayama, *Blow up rate for semilinear heat equations with subcritical nonlinearity*, Indiana Univ. Math. J., 53 (2004), 483–514.
- [12] Y. Giga, S. Matsui and S. Sasayama, *On blow up rate for sign-changing solutions in a convex domain*, Math. Methods Appl. Sci., 27 (2004), 1771–1782.
- [13] C. Gui and X. Wang, *Life span of solutions of the Cauchy problem for a semilinear heat equation*, J. Differential Equations, 115 (1995), 166–172.
- [14] O. Kavian, *Remarks on the large time behaviour of a nonlinear diffusion equation*, Ann. I. H. Poincaré, Anal. non linéaire, 4 (1987), 423–452.
- [15] T. Y. Lee and W. M. Ni, *Global existence, large time behavior and life span of solutions of a semilinear parabolic Cauchy problem*, Trans. Amer. Math. Soc., 333 (1992), 365–378.
- [16] H. A. Levine and P. Meier, *The value of the critical exponent for reaction-diffusion equations in cones*, Arch. Ration. Mech. Anal., 109 (1990), 73–80.
- [17] P. Meier, *Existence et non-existence de solutions globales d’une équations de la chaleur semi-linéaire: extention d’un théoreme de Fujita*, C. R. Acad. Sci., Paris Ser. I, 303 (1986), 635–637.
- [18] P. Meier, *Blow up of solutions of semilinear parabolic differential equations*, Z. Angew. Math. Phys., 39 (1988), 135–149.

- [19] N. Mizoguchi and E. Yanagida, *Blowup and life span of solutions for a semilinear parabolic equation*, SIAM J. Math. Anal., 29 (1998), 1434–1446.
- [20] L. Molinet and S. Tayachi, *Remarks on the Cauchy problem for the one-dimensional quadratic (fractional) heat equation*, J. Funct. Anal., 269 (2015), 2305–2327.
- [21] H. Mouajria, S. Tayachi and F. B. Weissler, *The heat equation on sectorial domains, highly singular initial values and applications*, J. Evol. Equ., 16 (2016), 341–364.
- [22] C. E. Mueller and F. B. Weissler, *Single point blow-up for a general semilinear heat equation*, Indiana Univ. Math. J., 34 (1985), 881–913.
- [23] P. Quittner and Ph. Souplet, “*Superlinear parabolic problems. Blow-up, global existence and steady states*”, Birkhäuser, Basel, 2007.
- [24] F. Ribaud, *Semilinear parabolic equations with distributions as initial values*, Discrete Contin. Dynam. Systems, 3 (1997), 305–316.
- [25] P. Rouchon, *Blow-up of solutions of nonlinear heat equations in unbounded domains for slowly decaying initial data*, ZAMP, 52 (2001), 1017–1032.
- [26] Ph. Souplet and F. B. Weissler, *Self-similar sub-solutions and blow-up for nonlinear parabolic equations*, J. Math. Anal. Appl., 212 (1997), 60–74.
- [27] S. Tayachi and F. B. Weissler, *The nonlinear heat equation with high order mixed derivatives of the Dirac delta as initial values*, Trans. AMS, 366 (2014), 505–530.
- [28] S. Tayachi and F. B. Weissler, *Some remarks on life span results*, in preparation.
- [29] F. B. Weissler, *Existence and nonexistence of global solutions for a semilinear heat equation*, Israel J. Math., 38 (1981), 29–40.
- [30] F. B. Weissler, *Rapidly decaying solutions of an ordinary differential equation with applications to semilinear elliptic and parabolic partial differential equations*, Arch. Rach. Mech. Anal., 91 (1985), 247–266.
- [31] F. B. Weissler, *L^p -Energy and blow-up for a semilinear heat equation*, Proceedings of Symposia in Pure Mathematics, 45 (1986), Part 2, 545–551.
- [32] J. Wu, *Well-posedness of a semilinear heat equation with weak initial data*, Jnl. Fourier Anal. Appl., 4 (1998), 629–642.