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### SESAM mode-locked Yb:SrLaAlO<sub>4</sub> laser 1

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16 Abstract: We report on a detailed investigation of continuous-wave and SESAM mode-locked 17 (ML) laser operation of an Yb:SrLaAlO<sub>4</sub> laser. Pumped with a high-brightness fiber laser at 976 nm, the ML Yb:SrLaAlO<sub>4</sub> laser delivered soliton pulses as short as 38 fs at 1071 nm with 18 19 an average output power of 86 mW and a pulse repetition rate of ~59 MHz. The maximum 20 average output power reached 1.05 W at 1054 nm for longer pulses (94 fs), corresponding to a peak power of 146 kW and an optical efficiency of 38.6%. To the best of our knowledge, this 21 22 is the first demonstration of SESAM ML operation of the Yb:SrLaAlO<sub>4</sub> laser.

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#### 24 1. Introduction

25 Ytterbium ion (Yb<sup>3+</sup>) doped tetragonal aluminates, e.g., Yb:CaGdAlO<sub>4</sub> (abbreviated 26 Yb:CALGO) and Yb:CaYAlO<sub>4</sub> (abbreviated Yb:CALYO), represent a type of broadband, 27 structurally disordered laser gain media which are extremely suitable for the generation of high-28 power femtosecond pulses from passively mode-locked (ML) lasers in the spectral region of  $\sim 1 \mu m$  [1-4]. Yb:CALGO and Yb:CALYO crystals belong to the ABCO<sub>4</sub> crystal family, where 29 A = Ca or Sr, B = Y, Gd, La, etc., and C = Al or Ga. Due to the inhomogeneous spectral 30 broadening originating from their local structure disorder, Yb:CALGO and Yb:CALYO 31 32 crystals exhibit extremely broad, smooth and flat gain spectral profiles which, combined with 33 attractive thermo-optical properties [5], such as relatively high thermal conductivity and nearly 34 "athermal" behavior, determine their successful applications in the generation of high-average 35 power sub-100 fs soliton pulses at  $\sim 1 \ \mu m$  [6-12]. The shortest pulse duration (<30 fs) were 36 achieved via the soft-aperture Kerr-lens mode locking (KLM) technique from Yb:CALGO and 37 Yb:CALYO lasers [13-16].

38 The success of the Yb:CALGO and Yb:CALYO crystals paves the way to further explore 39 other Yb3+-doped disordered aluminate crystals as femtosecond laser gain media. Very 40 recently, we developed a novel Yb<sup>3+</sup>-doped strontium lanthanum aluminate crystals grown by 41 the Czochralski method, i.e., Yb:SrLaAlO<sub>4</sub>, abbreviated Yb:SALLO [17]. It is also an ABCD<sub>4</sub>-42 type compound belonging to the tetragonal system with a K<sub>2</sub>NiF<sub>4</sub> structure (space group 43 14/mmm). The Yb:SALLO crystal exhibits relatively large absorption bandwidths of 24 nm for 44  $\pi$ - and 14 nm for  $\sigma$ -polarization (full width half maximum, FWHM). The broadband absorption 45 releases the requirement of wavelength stabilization for the pump sources and facilitates power 46 scaling by using commercially available high-power InGaAs laser diodes. The lattice disorder 47 in the structure of the Yb:SALLO crystal provides a variation of the local crystal field around 48 the Yb<sup>3+</sup> dopants ions resulting in a considerable inhomogeneous spectral broadening. As a 49 result, Yb:SALLO exhibits also broad, flat and smooth spectral gain profiles supporting sub-50 50 fs pulse generation from ML lasers. Recently, we demonstrated the first passively ML 51 operation of an Yb:SALLO laser with soliton pulses as short as 44 fs via soft-aperture KLM 52 technique [18]. Compared to its Yb:CALGO and Yb:CALYO counterparts, Yb:SALLO features: (i) lower melting point (~1650°C) [19], (ii) higher crystallinity and better optical 53 54 quality and (iii) slightly broader spectra of Yb3+ dopant ions. The drawback of SALLO is the 55 lower segregation coefficient for Yb<sup>3+</sup> ions,  $K_{Yb} = 0.235$  [17].

The promising spectroscopic features, as well as the previous promising ML results 56 57 motivated us to further explore the potential of Yb:SALLO. Implementing a SEmiconductor 58 Saturable Absorber Mirror (SESAM) for starting and stabilizing the soliton pulse generation 59 and dispersive mirrors (DMs) for intracavity group delay dispersion (GDD) management, we 60 demonstrate sub-40 fs soliton pulse generation from a SESAM ML Yb:SALLO laser.

#### 61 2. **Experimental setup**

The schematic of the Yb:SALLO laser is shown in Fig. 1. The high quality Yb:SALLO crystal 62 with an Yb<sup>3+</sup> doping of  $1.36 \times 10^{20}$  cm<sup>-3</sup>, corresponding to 1.175 at.% was employed. A cubic 63 64 sample with an aperture of 3 mm × 3 mm and a thickness of 3 mm was cut for light propagation 65 along the *a*-axis (*a*-cut). This crystal orientation was selected to ensure access to the desirable 66  $\pi$ -polarization corresponding to stronger pump absorption. The two sides of the crystal were polished to laser quality and remained uncoated. 67



Fig. 1. Schematic of the Yb:SALLO laser. FL: fiber laser; L: spherical focusing lens; M<sub>1</sub>, M<sub>2</sub> and M<sub>4</sub>: concave mirrors (RoC = -100 mm); M<sub>3</sub>: flat rear mirror used in the CW regime; DM<sub>1</sub> -DM<sub>4</sub>: flat dispersive mirrors; OC: output coupler; SESAM: SEmiconductor Saturable Absorber Mirror.

73 An X-folded astigmatically compensated linear cavity was constructed to evaluate the 74 performance of the Yb:SALLO crystal both in the continuous-wave (CW) and ML regimes. 75 The laser crystal was mounted in a water-cooled copper holder (coolant temperature: 20°C) and 76 placed at the Brewster's angle between the two concave mirrors  $M_1 - M_2$  (radius of curvature, 77 RoC = -100 mm) to minimize the insertion loss. The pump source was a CW, narrow-linewidth 78 (50 kHz) fiber laser at 976 nm. It emitted a nearly diffraction-limited beam with a propagation 79 factor ( $M^2$ ) of ~1.03. The pump beam was focused into the laser crystal through the dichroic 80 mirror  $M_1$  by a spherical focusing lens with a focal length of 75 mm, which resulted in a beam 81 waist of 16.2  $\mu$ m × 29.5  $\mu$ m in the sagittal and tangential planes, respectively.

82 In the CW regime, a four-mirror cavity was used: one cavity arm was terminated by a flat 83 rear mirror  $M_3$  and another arm – by a flat-wedged output coupler (OC) having a transmission at the laser wavelength  $T_{\rm OC}$  in the range 1% - 10%, see Fig. 1. The cavity mode size in the 84 85 crystal by the ABCD formalism yielded the radii of  $24 \,\mu\text{m} \times 47 \,\mu\text{m}$  in the sagittal and the 86 tangential planes, respectively.

87 For ML operation, the mirror  $M_3$  was replaced by a curved mirror  $M_4$  (RoC = -100 mm) to 88 create a second beam waist on the SESAM to ensure its deep bleaching. A commercial SESAM 89 (BATOP, GmbH) with a modulation depth of 1.2%, a relaxation time of  $\sim 1$  ps and a non-90 saturable loss of  $\sim 0.8\%$  was implemented for starting and stabilizing the ML operation. The 91 calculated radius of this second beam waist was ~77 µm. The intracavity GDD was optimized 92 by inserting two or four different flat dispersive mirrors (DMs) in the other cavity arm having 93 a negative GDD per bounce of:  $DM_1 = -250 \text{ fs}^2$ ,  $DM_2 = -250 \text{ fs}^2$ ,  $DM_3 = -100 \text{ fs}^2$  and  $DM_4 = -250 \text{ fs}^2$ . 94 55 fs<sup>2</sup>. The group velocity dispersion (GVD) of the Yb:SALLO crystal was estimated from the 95 dispersion curves [20] to be  $220 \pm 50$  fs<sup>2</sup>/mm at 1050 nm for  $\pi$ -polarization.

### 96 3. Continuous-wave laser operation

97 The Yb:SALLO laser generated a maximum CW output power of 1.42 W at 1051.6 nm with a 98 high slope efficiency of 67.3% and a low laser threshold of 121 mW (for  $T_{\rm OC} = 2.5\%$  and an 99 absorbed pump power of 2.22 W), see Fig. 2(a). The measured single-pass pump absorption 100 under lasing conditions depended on the transmission of the OCs ranging from 25% to 49%, 101 indicating a certain ground-state bleaching being suppressed by the recycling effect. The laser 102 threshold gradually increased with the transmission of the OC, from 66 mW ( $T_{OC} = 1\%$ ) to 103 285 mW ( $T_{\rm OC} = 10\%$ ). The laser spectra varied with the output coupling: for low  $T_{\rm OC}$  of 1%, 104 the laser operated at two lines, 1050 and 1068 nm, for intermediate  $T_{\rm OC} = 1.6\%$  - 2.5% - only 105 at ~1051 nm and for high  $T_{\rm OC} = 7.5\%$  - 10% - at yet shorter wavelength of ~1016 nm, see 106 Fig. 2(b). Such a blue-shift of the laser wavelength with increasing the output coupling is 107 typical for quasi-three-level Yb<sup>3+</sup> lasers with inherent reabsorption at the laser wavelength and 108 it agrees with the gain spectra of Yb:SALLO for  $\pi$ -polarization [17].



**Fig. 2.** (a) CW performance of the Yb:SALLO laser with different OCs,  $\eta$  – slope efficiency; (b) laser spectra; (c) Caird analysis: slope efficiency plotted as a function of the OC reflectivity ( $R_{\rm OC} = 1 - T_{\rm OC}$ ); (d) spectral tuning curve obtained with an intracavity SF10 prism and an OC with  $T_{\rm OC} = 0.6\%$ .

114 The total round-trip resonator losses  $\delta$  (reabsorption losses excluded) as well as the intrinsic 115 slope efficiency  $\eta_0$  (accounting for the mode-matching efficiency and the quantum efficiency) 116 were estimated using the Caird analysis by fitting the laser slope efficiency as a function of the 117 OC reflectivity ( $R_{\rm OC} = 1 - T_{\rm OC}$ ) [21]. The best fitting curve gave  $\eta_0 = 69.9\%$  and  $\delta = 0.13\%$ , as 118 shown in Fig. 2(c). Such low value of  $\delta$  is an evidence for the excellent optical quality of the 119 Yb:SALLO crystal. The spectral tuning curve in the CW regime was studied at an absorbed 120 pump power of 0.77 W using a SF10 prism and a 0.6% OC. The wavelength was continuously tunable between 1012 - 1092 nm corresponding to a range as broad as 80 nm, Fig. 2(d). 121

## 122 4. Mode-locked laser operation

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123 Stable and self-starting SESAM ML operation was readily achieved by implementing two flat

124 DMs  $(DM_1 - DM_2)$  into the cavity arm terminated by the OC, Fig. 1, which provided a total

# round-trip negative GDD of -2000 fs<sup>2</sup> to balance the material dispersion, as well as the induced self-phase modulation (SPM).





For 2.5% OC, the measured optical spectrum and the second harmonic generation (SHG) based intensity autocorrelation trace are shown in Fig. 3(a) and Fig. 3(b), respectively. The observed excellent sech<sup>2</sup>-shaped spectral and temporal profiles (although in theory in both cases this is an approximation) are an evidence for soliton-like pulse generation. The pulse duration of 94.3 fs at a central wavelength of 1054 nm and the emission bandwidth (FWHM) of 12.7 nm correspond to a time-bandwidth product (TBP) of 0.323 indicating nearly bandwidth-limited pulses. The inset of Fig. 3(b) shows an intensity autocorrelation trace on a longer time scale (50 ps) indicating single-pulse mode-locking without multiple pulse instabilities. The average output power of the ML laser amounted to 1.05 W for an absorbed pump power of 2.72 W. This corresponds to an optical efficiency of 38.6%. The recorded radio frequency (RF) spectrum of the laser output is shown in Fig. 3(c) and (d). A sharp peak at the fundamental beat note of 67.08 MHz exhibited a high extinction ratio of 76 dBc above the noise level. The uniform harmonics on a 1-GHz frequency span reveal high stability of the ML operation.



A rapid spectral broadening was observed by translating the folding mirror  $M_2$  by few hundred of micrometers away from the laser crystal. After carefully aligning the cavity, the bandwidth (FWHM) of the laser spectrum increased from 12.7 to 22.6 nm, the average output power dropped to 404 mW and the central laser wavelength experienced a red-shift to 1062.8 nm, see Fig. 4(a). The recorded autocorrelation trace gave a pulse duration of 53 fs (FWHM), as shown in Fig. 4(b). The corresponding TBP of 0.318 was even closer to the Fourier-transform limit. The long-time autocorrelation trace (Fig. 4(b), inset) and the RF spectra, Fig. 4(c-d), confirmed similarly stable and single pulse mode-locking performance.



Fig. 5. Measured far-field beam profiles of the ML Yb:SALLO laser: transition of (a) SESAM ML to (b) Kerr-lens assisted SESAM ML regimes.

Far-field beam profiles were measured to reveal the dominating ML mechanism for such a transition. An IR camera was placed ~1.1 m away from the OC. In the SESAM ML regime (cf. Fig. 3), the laser beam diameters were 2.10 mm × 1.96 mm, see Fig. 5(a). After spectral broadening (cf. Fig. 4), a significant change appeared in the measured far-field beam profile with a beam diameter of 1.89 mm × 1.89 mm, see Fig. 5(b). The shrinking of the far-field beam profile validated the underlying spectral broadening mechanism dominated by Kerr-lens mode-locking (KLM) stabilized by the SESAM, which is supported by the high-brightness of the pump laser. Such soft-aperture Kerr-lens effect occurred due to the gradually increasing laser mode size inside the Yb:SALLO crystal when varying the separation between the pump mirror  $M_1$  and the folding mirror  $M_2$ . Nevertheless, pulse shaping by the SESAM cannot be ruled out in view of the much wider stability zone and the self-starting mode-locking operation compared to our previous work [18]. Such hybrid pulse shaping mechanism introduced an enhanced self-amplitude modulation (SAM) by the Kerr-lens effect, which assisted the SESAM ML operation for shorter pulse duration, i.e., Kerr-lens assisted SESAM ML operation.



**Fig. 6.** Characterization of the shortest pulses from the ML Yb:SALLO laser. (a) Optical spectrum and (b) SHG-based intensity autocorrelation trace with a sech<sup>2</sup>-fit. *Inset*: simultaneously measured long-scale background free intensity autocorrelation trace for the time span of 50 ps. Radio-frequency (RF) spectra of the ML Yb:SALLO laser: (a) fundamental beat note at ~59.4 MHz recorded with a resolution bandwidth (RBW) of 100 Hz, and (b) harmonics on a 1-GHz frequency span, measured with a RBW of 100 kHz.

The pulse duration could be further shortened by reducing the total intracavity negativeGDD in the Kerr-lens assisted SESAM ML regime. By implementing two extra bounces on flat

DMs (DM<sub>3</sub> and DM<sub>4</sub>, see Fig. 1) with a total negative intracavity GDD of -1620 fs<sup>2</sup>, the 184 185 Yb:SALLO laser delivered soliton pulses as short as 38 fs at 1071 nm for the same 2.5% OC, 186 see Fig. 6(a) and (b). The optical spectrum of the ML laser had a bandwidth (FWHM) of 187 32.2 nm assuming a sech<sup>2</sup>-shape spectral profile (due to the very strong SPM, the measured 188 spectrum exhibited a slight deviation from such an ideal sech<sup>2</sup>-shaped spectral profile). The 189 average output power amounted to 86 mW at an absorbed pump power of 0.78 W. The resulting 190 TBP was 0.32, very close to the Fourier transform limit value (0.315). The steady-state ML 191 pulse train corresponding to the shortest pulse duration was characterized by long scale 192 autocorrelation, inset Fig. 6(b), and RF spectra, as shown in Fig. 6(c) and (d). The sharp first 193 beat note at ~59.4 MHz with a signal-noise-ratio above 71 dBc, and the uniform harmonics 194 again proved stable CW mode-locking without any Q-switched instability.

#### 195 5. Conclusion

196 In conclusion, we demonstrate for the first time, to the best of our knowledge, SESAM ML 197 operation of an Yb:SALLO laser. Assisted by enhanced self-amplitude modulation (SAM) 198 originating from the Kerr-lens effect, the laser produced soliton pulses as short as 38 fs at 199 1071 nm with an average output power of 86 mW. Watt-level sub-100 fs pulse generation was 200 obtained in the SESAM ML regime at the expense of longer pulse duration of 94.3 fs at 201 1054 nm: the average output power reached 1.05 W at an absorbed power of 2.72 W, which 202 corresponded to a peak power of 146 kW and an optical efficiency of 38.6%. The excellent 203 spectroscopic and thermo-mechanical properties of the Yb:SALLO crystal reveals a great 204 potential for power scalable operation in the sub-50 fs time domain via Kerr-lens assisted 205 SESAM ML technique in combination with commercially available high-power InGaAs pump diode lasers. 206

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213 214 Data availability. Data underlying the results presented in this paper are not publicly available at this time but may be obtained from the authors upon reasonable request.

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