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Non trivial zeros of the Zeta function using the differential equations

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Abstract

Using the differential equations, we obtain a more flexible expression for the Riemann Zeta function on the right half-plan except the point s=1. Thanks to the Riemann functional equation, we suggest that the point $s\in\mathbb{C}$ such that $\Re(s)\in(0,2^{-1}]$ is a non-trivial zero of the function ζ implies $|s^2\Gamma(s)|\geq |(1-s)^2\Gamma(1-s)|$, where Γ is the Gamma function.

Keywords: Zeta function, Non trivial zeros, Riemann functional equation, Bernoulli numbers, Differential equations.

AMS subject classifications: 00A05

1 Main results

Consider the representation of the Riemann Zeta function ζ defined by the Abel summation formula [[1], page 14 Equation 2.1.5] as

$$\zeta(s) := -\frac{s}{1-s} - s \int_{1}^{+\infty} u^{-1-s} \{u\} du, \quad s \neq 1, \ \Re(s) > 0, \ \Im(s) \in \mathbb{R}^*, \quad (1)$$

where $\{u\}$ is the fractional part of the real u. In order to simplify the notation, denote $B\subset\mathbb{C}$ the right half-plan except the point s=1, defined as

$$B := \{ s \in \mathbb{C} : \quad s \neq 1, \quad \Re(s) > 0, \quad \Im(s) \in \mathbb{R}^* \},$$

In this manuscript, we denote $(B_k)_{k\in\mathbb{N}}$ the Bernoulli numbers. We prove the following theorem,

Theorem 1. Consider the Zeta function given by the Equation (1). For every $s \in B$ we have

$$\zeta(s) = -\frac{s}{1-s} - \frac{1}{4} \lim_{n \to +\infty} \left(\prod_{k=0}^{n+2} (k+s) \right) \eta_s(n),$$

where

$$\eta_s(n) := \int_n^{+\infty} \frac{(u+2)^{n+s}}{(u+1)^{n+2+s}} \int_{u+1}^{+\infty} \frac{q(n-1,v)}{(v+1)^{n+1+s}} dv du$$

$$- (n+1)^{-n-1-s} \int_n^{n+1} (u+1)^{n+s} \int_u^{+\infty} \frac{q(n-1,v)}{(v+1)^{n+1+s}} dv du$$

$$+ \int_n^{+\infty} (u+1)^{-n-2-s} \Big(q(n,u) - q(n,u+1) \Big) du, \quad \forall n \ge 1,$$

and where

$$q(n,u) := \sum_{k=1}^{n+1} \frac{B_k}{k!} \frac{u^{n+2-k}}{(n+2-k)!} \quad \forall n \ge 0, \quad \forall u \ge 0.$$

2 Main Proposition

For every $s \in B$, the Equation (1) is equivalent to,

$$\frac{\zeta(s)}{s} = -\frac{1}{1-s} - \int_{1}^{+\infty} u^{-1-s} \{u\} du.$$

The aim is to studies the differential equation of solutions the functions

$$t \mapsto \psi_s(z,t) := t^s \Big[z + \int_1^t u^{-1-s} \{u\} du \Big], \quad z \in \mathbb{C}, \ t \ge 1.$$

We focus only on the bounded solutions (there is a unique bounded solution. All other solutions are oscillating and diverge to infinity in norm). More precisely, the strategy to prove the Theorem 1, is to find the initial condition $z_s \in \mathbb{C}$ such that $\sup_{t\geq 1} |\psi_s(z_s,t)| < +\infty$. Since $\sup_{t\geq 1} |\psi_s(z,t)| < +\infty$ if and only if $z = -\int_1^{+\infty} u^{-1-s} \{u\} du$, we obtain $z_s = -\int_1^{+\infty} u^{-1-s} \{u\} du$. For every $s \in B$ we consider the following differential equation

$$\frac{d}{dt}x = st^{-1}x + t^{-1}\{t\},$$

$$t \in [1, +\infty)/\mathbb{N}, \quad x(1) = z, \ z \in \mathbb{C}, \quad x : [1, +\infty) \to \mathbb{C}.$$

$$(2)$$

Lemma 2. Let be $s \in B$. For every $z \in \mathbb{C}$ there exists a unique continuous solution $\psi_s(z,t): [1,+\infty) \to \mathbb{C}$ of the differential equation (2). Further,

$$\psi_s(z,t) = t^s \Big[z + \int_1^t u^{-1-s} \{u\} du \Big], \quad \forall t \ge 1.$$

Proof. Let be $s \in B$ fixed. The function

$$t \mapsto t^s \int_1^t u^{-1-s} \{u\} du, \ t \ge 1,$$

is C^{∞} on $[1, +\infty)/\mathbb{N}$ and continuous on $[1, +\infty)$. The Equation (2) is a non-homogeneous linear differential equation. For every $z \in \mathbb{C}$ the unique continuous solution $\psi_s(z,t): [1, +\infty) \to \mathbb{C}$ such that $\psi_s(z,1) = z$ is given by

$$\psi_s(z,t) = t^s \Big[z + \int_1^t u^{-1-s} \{u\} du \Big], \ \forall t \ge 1.$$

Proposition 3. Let be $s \in B$ and consider the continuous solution $\psi_s(z,t): [1,+\infty) \to \mathbb{C}$ of the differential equation (2). Then $\sup_{t\geq 1} |\psi_s(z,t)| < +\infty$ if and only if the initial condition $z \in \mathbb{C}$ satisfies the following equation

$$z_s = -\frac{1}{4} \lim_{n \to +\infty} \left(\prod_{k=0}^{n+1} (k+1+s) \right) \eta_s(n).$$

where $(\eta_s(n))_{n\in\mathbb{N}}$ is defined as in the Theorem 1.

Let us introduce the following notations,

Notation 4. Let $g: \mathbb{R}_+ \to \mathbb{C}$ be a continuous function, we denote the function $\Phi[g]: \mathbb{R}_+ \to \mathbb{C}$ as

$$\Phi[g](t) := \int_0^t g(u)du, \quad \forall t \ge 0,$$

Notation 5. For every $k \in \mathbb{N}$ define the function $p_k : \mathbb{R} \to \mathbb{R}$ as

$$p_{2k+1}(t) := (-1)^k \sum_{j \in \mathbb{Z}^*} \frac{1}{(j2\pi)^{2k+3}} \sin(j2\pi t),$$
$$p_{2k}(t) := (-1)^k \sum_{j \in \mathbb{Z}^*} \frac{1}{(j2\pi)^{2k+2}} \Big(\cos(j2\pi t) - 1\Big),$$

Lemma 6. We have

$$\int_0^t \left(\{u\} - \frac{1}{2} \right) du = p_0(t), \quad \forall t \ge 0,$$

and for every $n \in \mathbb{N}^*$ we have

$$\Phi^{n}[p_{0}](t) = -\frac{1}{(n+2)!} \sum_{k=2}^{n+1} B_{k} \binom{n+2}{k} t^{n-k+2} + p_{n}(t), \quad \forall t \ge 0, \ \forall n \ge 1.$$

Proof. Prove that

$$\int_0^t \left(\{u\} - \frac{1}{2} \right) du = p_0(t), \quad \forall t \ge 0,$$

The function $u \mapsto \{u\}$ is 1-periodic, then there exists a continuous 1-periodic function $p: \mathbb{R} \to \mathbb{R}$ such that

$$\int_{0}^{t} \{u\} du = t \int_{0}^{1} \{u\} du + p(t), \quad \forall t \ge 0.$$

Since

$$\int_0^1 \{u\} du = \int_0^1 u \ du = \frac{1}{2},$$

we get

$$\int_0^t \{u\} du = t\frac{1}{2} + p(t), \quad \forall t \ge 0.$$

The function p is a piecewise C^{∞} , continuous on \mathbb{R} and 1-periodic. By Dirichlet Theorem, the Fourier series

$$n \mapsto \sum_{k=-n}^{n} a_k \exp(ik2\pi t),$$

converge uniformly on \mathbb{R}_+ to the function $t \mapsto p(t)$, where $(a_k)_k \subset \mathbb{C}$ are the Fourier coefficients of the function p.

$$\int_0^t \left(\{v\} - \frac{1}{2} \right) dv = \sum_{j \in \mathbb{Z}} a_j \exp(ij2\pi t), \quad \forall t \ge 0.$$

By definition of the Fourier coefficients we have

$$a_{j} = \int_{0}^{1} \exp(-ij2\pi u) \left(\int_{0}^{u} \left(\{v\} - \frac{1}{2} \right) dv \right) du$$
$$= \frac{1}{2} \int_{0}^{1} \exp(-ij2\pi u) u(u - 1) du = \frac{1}{(2j\pi)^{2}}, \quad \forall j \in \mathbb{Z}^{*},$$

and

$$a_0 = \frac{1}{2} \int_0^1 u(u-1)du = -\frac{1}{12} = -\sum_{j \in \mathbb{Z}^*} \frac{1}{(j2\pi)^2}.$$

The function p satisfies

$$p(t) = \sum_{j \in \mathbb{Z}^*} \frac{1}{(j2\pi)^2} \Big(\exp(ij2\pi t) - 1 \Big) = p_0(t), \quad \forall t \ge 0.$$

Integrate successively to obtain

$$\Phi^{n}[p_{0}](t) = \sum_{k=0}^{n} \frac{c_{k}}{(n-k)!} t^{n-k} + p_{n}(t), \quad \forall t \ge 0, \ \forall n \ge 1,$$

where $(c_k)_{k\in\mathbb{N}}$ defined as

$$c_{2k+1} = 0$$
 and $c_{2k} = (-1)^k \sum_{j \in \mathbb{Z}^*} \frac{1}{(j2\pi)^{2(k+1)}}, \quad \forall k \ge 0.$

By definition of the Bernoulli numbers $(B_k)_{k\in\mathbb{N}}$, we get

$$\Phi^{n}[p_{0}](t) = -\frac{1}{(n+2)!} \sum_{k=2}^{n+1} B_{k} \binom{n+2}{k} t^{n-k+2} + p_{n}(t), \quad \forall t \ge 0, \ \forall n \ge 1,$$

Lemma 7. For every $s \in B$, there exists a unique $z_s \in \mathbb{C}$ such that the solution $\psi_s(z_s,t): [1,+\infty) \to \mathbb{C}$ of the differential equation (2) satisfies $\sup_{t\geq 1} |\psi_s(z_s,t)| < +\infty$. Further,

$$z_s = -\int_1^{+\infty} u^{-1-s} \{u\} du.$$

Proof. Let be $s \in B$ fixed. There is a unique bounded solution. All other solutions are oscillating and diverge to infinity in norm. Let be $z \in \mathbb{C}$ and suppose that

$$\sup_{t\geq 1}|\psi_s(z,t)|<+\infty.$$

Since $\Re(s) > 0$, by the Lemma 2, we have

$$\sup_{t>1} |\psi_s(z,t)| < +\infty \implies \left| z + \int_1^{+\infty} u^{-1-s} \{u\} du \right| = 0,$$

therefore

$$z_s = -\int_1^{+\infty} u^{-1-s} \{u\} du.$$

Prove that $\sup_{t\geq 1} |\psi_s(z_s,t)| < +\infty$. Since

$$\left| z_s + \int_1^{+\infty} u^{-1-s} \{u\} du \right| = 0,$$

by consequence,

$$z_s + \int_1^t u^{-1-s} \{u\} du = -\int_t^{+\infty} u^{-1-s} \{u\} du, \quad \forall t > 1.$$

By the Lemma 2, we get

$$\psi_s(z_s, t) = -t^s \int_t^{+\infty} u^{-1-s} \{u\} du, \quad \forall t \ge 1.$$

Then $\sup_{t\geq 1} |\psi_s(z_s,t)| \leq \frac{1}{\Re(s)}$.

Proof of the Proposition 3. Let be $s \in B$ and consider the bounded continuous solution $\psi_s(z_s,t):[1,+\infty)\to\mathbb{C}$ defined in the Lemma 7. In order to simplify the notation, denote

$$\psi_s(t) := \psi_s(z_s, t+1), \quad \forall t \ge 0.$$

Using the fact $t \mapsto \psi_s(z_s, t)$ is solution of the differential equation (2). By definition of Φ in the Notation 4, we get

$$(t+1)\frac{d}{dt}\Phi^{n+1}[\psi_s](t) = (n+1+s)\Phi^{n+1}[\psi_s](t)$$

$$+ z_s \frac{t^n}{n!} + h(n,t), \quad \forall t \ge 0, \quad \forall n \ge 0,$$

$$h(n,t) := \frac{1}{2} \frac{t^{n+1}}{(n+1)!} + \Phi^n[p_0](t), \quad \forall t \ge 0, \quad \forall n \ge 0,$$

$$\Phi^{n+1}[\psi_s](0) = 0, \quad \forall n \ge 0,$$

$$(3)$$

where the function p_0 is defined in the Notation 5. Integrate, for every $n \ge 0$ and $t \ge 0$ we have

$$\Phi^{n+1}[\psi_s](t) = (t+1)^{n+1+s} \int_0^t (u+1)^{-n-2-s} \left[z_s \frac{u^n}{n!} + h(n,u) \right] du. \tag{4}$$

Since $\sup_{v\geq 0} |\psi_s(v)| = \sup_{v\geq 1} |\psi_s(z_s,v)| < +\infty$, then for every fixed $n\geq 0$ we have

$$\left| \Phi^{n+1}[\psi_s](t) \right| \le \Phi^{n+1}[1](t) \sup_{v \ge 0} |\psi_s(v)| = \frac{t^{n+1}}{(n+1)!} \sup_{v \ge 1} |\psi_s(z_s, v)|, \quad \forall t \ge 0,$$

The Equation (12) implies

$$\int_0^{+\infty} (u+1)^{-n-2-s} \left[z_s \frac{u^n}{n!} + h(n,u) \right] du = 0, \quad \forall n \ge 0,$$

in other words

$$\frac{z_s}{\prod_{k=1}^{n+1}(k+s)} = -\int_0^{+\infty} (u+1)^{-n-2-s} h(n,u) du, \quad \forall n \ge 0,$$
 (5)

To know, the previous equality allows us to obtain the representation of the Zeta function obtained by the Euler-Maclaurin formula. Let be $\psi_s(0,t):[1,+\infty)\to\mathbb{C}$ the unique continuous solution of the differential equation (2) of initial condition $|\psi_s(0,1)|=0$. We recall that $t\mapsto \psi_s(z,t)$ is defined in the Lemma 2 for every initial condition $z\in\mathbb{C}$. In order to simplify the notation, denote

$$\psi_{0,s}(t) := \psi_s(0, t+1), \quad \forall t \ge 0.$$

Using the fact $t \mapsto \psi_s(0,t)$ is solution of the differential equation (2). As in the Equation (3), we get

$$(t+1)\frac{d}{dt}\Phi^{n+1}[\psi_{0,s}](t) = (n+1+s)\Phi^{n+1}[\psi_{0,s}](t) + h(n,t),$$

$$\Phi^{n+1}[\psi_{0,s}](0) = 0, \quad \forall n \ge 0, \forall t_q eq 0,$$
(6)

where we recall that

$$h(n,t) = \frac{1}{2} \frac{t^{n+1}}{(n+1)!} + \Phi^n[p_0](t), \quad \forall t \ge 0, \quad \forall n \ge 0.$$

Integrate the last differential equation,

$$\Phi^{n+1}[\psi_{0,s}](t) = (t+1)^{n+1+s} \int_0^t (u+1)^{-n-2-s} h(n,u) du, \quad \forall t \ge 0, \quad \forall n \ge 0.$$

which can be written as

$$\Phi^{n+1}[\psi_{0,s}](t) = (t+1)^{n+1+s} \int_0^{+\infty} (u+1)^{-n-2-s} h(n,u) du$$
$$-(t+1)^{n+1+s} \int_t^{+\infty} (u+1)^{-n-2-s} h(n,u) du, \quad \forall t \ge 0, \quad \forall n \ge 0.$$

The Equation (5) implies

$$\Phi^{n+1}[\psi_{0,s}](t) = -\frac{z_s}{\prod_{k=1}^{n+1} (k+s)} (t+1)^{n+1+s}$$

$$-(t+1)^{n+1+s} \int_t^{+\infty} (u+1)^{-n-2-s} h(n,u) du.$$
(7)

In order to simplify the notation, denote

$$\delta_s(n,t) := \Phi^{n+1}[\psi_{0,s}](t+1) - \Phi^{n+1}[\psi_{0,s}](t), \quad \forall n \ge 0, \ \forall t \ge 0.$$
 (8)

By the Lemma 6, we have

$$h(n,u) = p_n(u) - q(n,u). \tag{9}$$

where p_n are given by the Notation 5 and where in order to simplify the notation, we denoted q(n, u) the following polynomial

$$q(n,u) := -\frac{1}{2} \frac{u^{n+1}}{(n+1)!} + \frac{1}{(n+2)!} \sum_{k=2}^{n+1} B_k \binom{n+2}{k} u^{n-k+2}.$$

or since $B_1 = -\frac{1}{2}$, the function q_s can be written as

$$q(n,u) := \frac{1}{(n+2)!} \sum_{k=1}^{n+1} B_k \binom{n+2}{k} u^{n-k+2} = \sum_{k=1}^{n+1} \frac{B_k}{k!} \frac{u^{n+2-k}}{(n+2-k)!}.$$
 (10)

From the Equation (6), the periodicity of the function p_n implies

$$(t+1)\frac{d}{dt}\delta_s(n,t) = (n+1+s)\delta_s(n,t) - q(n,t+1) + q(n,t) - \frac{d}{dt}\Phi^{n+1}[\psi_{0,s}](t+1), \quad \forall t \ge 0, \quad \forall n \ge 0,$$
$$\delta_s(n,0) = \Phi^{n+1}[\psi_{0,s}](1), \quad \forall n \ge 0.$$

Since

$$\frac{d}{dt}\Phi^{n+1}[\psi_{0,s}](t) = \Phi^n[\psi_{0,s}](t), \quad \forall t > 0, \quad \forall n \ge 0,$$

the previous differential Equation can be written as

$$\frac{d}{dt}\delta_{s}(n,t) = (t+1)^{-1} \Big[(n+1+s)\delta_{s}(n,t) + \tilde{q}(n,t) - \Phi^{n}[\psi_{0,s}](t+1) \Big],
\tilde{q}(n,t) := q(n,t) - q(n,t+1), \ \forall t \ge 0, \ \forall n \ge 0,
\delta_{s}(n,0) = \Phi^{n+1}[\psi_{0,s}](1), \quad \forall n \ge 0.$$
(11)

Integrate, for every $n \geq 0$ and $t \geq 0$ we have

$$\frac{\delta_s(n,t)}{(t+1)^{n+1+s}} = \Phi^{n+1}[\psi_{0,s}](1) + \int_0^t \frac{\tilde{q}(n,u) - \Phi^n[\psi_{0,s}](u+1)}{(u+1)^{n+2+s}} du.$$
 (12)

Since $\sup_{v\geq 0} |(v+1)^{-\Re(s)}\psi_{0,s}(v)| = \sup_{v\geq 1} |v^{-\Re(s)}\psi_s(0,v)| < +\infty$, then for every fixed $n\geq 0$ we have

$$\left| \Phi^{n+1}[\psi_{0,s}](t) \right| \le (t+1)^{\Re(s)} \Phi^{n+1}[1](t) \sup_{v \ge 0} |(v+1)^{-\Re(s)} \psi_{0,s}(v)| \qquad (13)$$

$$= (t+1)^{\Re(s)} \frac{t^{n+1}}{(n+1)!} \sup_{v > 1} |v^{-\Re(s)} \psi_{s}(0,v)|, \quad \forall t \ge 0,$$

By the notation of δ_s in the Equation (8) and by the definition of Φ in the Notation 4, we have

$$\delta_s(n,t) := \int_t^{t+1} \Phi^n[\psi_{0,s}](u) du, \quad \forall n \ge 0, \ \forall t \ge 0.$$
 (14)

The Equation (13) implies

$$|\delta_s(n,t)| = \left| \int_t^{t+1} \Phi^n[\psi_{0,s}](u) du \right|$$

$$\leq (t+2)^{\Re(s)} \frac{(t+1)^n}{n!} \sup_{v>1} |v^{-\Re(s)} \psi_s(0,v)|, \ \forall t \geq 0, \quad \forall n \geq 0.$$

By consequence

$$\lim_{t \to +\infty} \left| (t+1)^{-n-1-s} \delta_s(n,t) \right| = 0, \quad \forall n \ge 0.$$

From the Equations (12) we obtain

$$\Phi^{n+1}[\psi_{0,s}](1) + \int_0^{+\infty} \frac{\tilde{q}(n,u) - \Phi^n[\psi_{0,s}](u+1)}{(u+1)^{n+2+s}} du = 0, \quad \forall n \ge 0.$$
 (15)

and the function $t \mapsto \delta_s(n,t)$ can be written as

$$\frac{\delta_s(n,t)}{(t+1)^{n+1+s}} = -\int_t^{+\infty} \frac{\tilde{q}(n,u) - \Phi^n[\psi_{0,s}](u+1)}{(u+1)^{n+2+s}} du.$$

By the Equation (14), for every $n \ge 1$ and $t \ge 0$, the last equality implies

$$\int_{t}^{t+1} \Phi^{n}[\psi_{0,s}](u)du = -(t+1)^{n+1+s} \int_{t}^{+\infty} \frac{\tilde{q}(n,u) - \Phi^{n}[\psi_{0,s}](u+1)}{(u+1)^{n+2+s}} du.$$

The Equation (7) implies

$$\frac{z_s}{\prod_{k=1}^n (k+s)} \left[\int_t^{+\infty} \frac{(u+2)^{n+s}}{(u+1)^{n+2+s}} du - (t+1)^{-n-1-s} \int_t^{t+1} (u+1)^{n+s} du \right]
= (t+1)^{-n-1-s} \int_t^{t+1} (u+1)^{n+s} \int_u^{+\infty} (v+1)^{-n-1-s} h(n-1,v) dv du
- \int_t^{+\infty} (u+1)^{-n-2-s} (u+2)^{n+s} \int_{u+1}^{+\infty} (v+1)^{-n-1-s} h(n-1,v) dv du
- \int_t^{+\infty} (u+1)^{-n-2-s} \tilde{q}(n,u) du, \quad \forall n \ge 1, \ \forall t \ge 0.$$

In particular, for every $n \ge 1$ we have

$$\frac{z_s}{\prod_{k=1}^n (k+s)} \left[\int_n^{+\infty} \frac{(u+2)^{n+s}}{(u+1)^{n+2+s}} du - \frac{1}{n+1+s} \left(\left(\frac{n+2}{n+1} \right)^{n+1+s} - 1 \right) \right] \\
= (n+1)^{-n-1-s} \int_n^{n+1} (u+1)^{n+s} \int_u^{+\infty} (v+1)^{-n-1-s} h(n-1,v) dv du \\
- \int_n^{+\infty} (u+1)^{-n-2-s} (u+2)^{n+s} \int_{u+1}^{+\infty} (v+1)^{-n-1-s} h(n-1,v) dv du \\
- \int_n^{+\infty} (u+1)^{-n-2-s} \tilde{q}(n,u) du.$$

Using the notation of h given by the Equation (9), we have obtained

$$z_s = \frac{1}{\omega_s(n)} \Big(\prod_{k=1}^{n+1} (k+s) \Big) \Big[\eta_s(n) - \epsilon_s(n) \Big], \quad \forall n \ge 1.$$

where

$$\omega_{s}(n) := 1 - \left(\frac{n+2}{n+1}\right)^{n+1+s} + (n+1+s) \int_{n}^{+\infty} \frac{(u+2)^{n+s}}{(u+1)^{n+2+s}} du,$$

$$\eta_{s}(n) := \int_{n}^{+\infty} \frac{(u+2)^{n+s}}{(u+1)^{n+2+s}} \int_{u+1}^{+\infty} \frac{q(n-1,v)}{(v+1)^{n+1+s}} dv du$$

$$- (n+1)^{-n-1-s} \int_{n}^{n+1} (u+1)^{n+s} \int_{u}^{+\infty} \frac{q(n-1,v)}{(v+1)^{n+1+s}} dv du$$

$$+ \int_{n}^{+\infty} (u+1)^{-n-2-s} \tilde{q}(n,u) du,$$

$$\epsilon_{s}(n) := \int_{n}^{+\infty} \frac{(u+2)^{n+s}}{(u+1)^{n+2+s}} \int_{u+1}^{+\infty} \frac{p_{n-1}(v)}{(v+1)^{n+1+s}} dv du$$

$$- (n+1)^{-n-1-s} \int_{n}^{n+1} (u+1)^{n+s} \int_{u}^{+\infty} \frac{p_{n-1}(v)}{(v+1)^{n+1+s}} dv du.$$

By the notation of p_n in the Notation 5 we have $\sup_{n>0} \max_{v\in[0,1]} |p_n(v)| < +\infty$. By Stirling formula, we get

$$\lim_{n \to +\infty} \left(\prod_{k=1}^{n+2} (k+s) \right) |\epsilon_s(n)| = 0.$$

For the limit of $(n+2+s)\omega_s(n)$, we have

$$\omega_{s}(n) = -\int_{1}^{+\infty} v^{-2} \exp(v^{-1}) dv + \exp(1) - \left(\frac{n+2}{n+1}\right)^{n+1+s}$$

$$+ \int_{1}^{+\infty} \left(\frac{nv+2}{nv+1}\right)^{s} \frac{n(n+1+s)}{(nv+1)^{2}} \exp\left(n\ln\left(\frac{nv+2}{nv+1}\right)\right) dv$$

$$= \exp(1) \left[1 - \exp\left((n+1)\ln\left(\frac{n+2}{n+1}\right) - 1\right)\right]$$

$$- \left(\left(\frac{n+2}{n+1}\right)^{s} - 1\right) \left(\frac{n+2}{n+1}\right)^{n+1}$$

$$+ \int_{1}^{+\infty} \left(\frac{nv+2}{nv+1}\right)^{s} \left(\frac{n(n+1+s)}{(nv+1)^{2}} - v^{-2}\right) \exp\left(n\ln\left(\frac{nv+2}{nv+1}\right)\right) dv$$

$$+ \int_{1}^{+\infty} \frac{\exp(v^{-1})}{v^{2}} \left(\frac{nv+2}{nv+1}\right)^{s} \left[\exp\left(n\ln\left(\frac{nv+2}{nv+1}\right) - v^{-1}\right) - 1\right] dv$$

$$+ \int_{1}^{+\infty} v^{-2} \exp(v^{-1}) \left(\left(\frac{nv+2}{nv+1}\right)^{s} - 1\right) dv.$$

Use the Taylor formula, we get

$$\lim_{n \to +\infty} (n+2+s)\omega_s(n) = (\frac{1}{2} - s) \exp(1) + (1+s) \int_1^{+\infty} v^{-2} \exp(v^{-1}) dv$$
$$+ s \int_1^{+\infty} v^{-3} \exp(v^{-1}) dv - \frac{3}{2} \int_1^{+\infty} v^{-4} \exp(v^{-1}) dv$$
$$= -4.$$

By consequence

$$z_s = -\frac{1}{4} \lim_{n \to +\infty} \left(\prod_{k=0}^{n+1} (k+1+s) \right) \eta_s(n).$$

3 Proof of the Theorem 1

Proof of the Theorem 1. Let be $s \in B$. By the Equation (1) we have,

$$\frac{\zeta(s)}{s} + \frac{1}{1-s} = -\int_{1}^{+\infty} u^{-1-s} \{u\} du.$$

Thanks to the Main Proposition 3 and the Lemma 7 we get

$$\frac{\zeta(s)}{s} + \frac{1}{1-s} = -\frac{1}{4} \lim_{n \to +\infty} \left(\prod_{k=0}^{n+1} (k+1+s) \right) \eta_s(n).$$

4 Conclusion

Let be $s \in B$ such that $\Re(s) \in (0,1)$ and suppose that $|\zeta(s)| = 0$. By Theorem 1, we have

$$\lim_{n \to +\infty} \left(\prod_{k=0}^{n+2} (k+s) \right) \eta_s(n) = -4 \frac{s}{1-s},\tag{16}$$

Thanks to the Riemann functional equation, the point 1-s is also a zero of the function ζ , then

$$\lim_{n \to +\infty} \left(\prod_{k=0}^{n+2} (k+1-s) \right) \eta_{1-s}(n) = -4 \frac{1-s}{s}.$$
 (17)

The terms $\eta_x(n)$ are integrable for every $n \ge 1$ and $x \in B$. For $2\Re(s) \le 1$, we suggest that

$$\lim_{n \to +\infty} \frac{n^{\Re(s)}}{n^{1-\Re(s)}} \left| \frac{\eta_s(n)}{\eta_{1-s}(n)} \right| \ge 1,$$

The Equations (16) and (17) implies

$$\left| \frac{s^2}{(1-s)^2} \right| = \lim_{n \to +\infty} \left| \left(\prod_{k=0}^{n+2} \frac{k+s}{k+1-s} \right) \frac{\eta_s(n)}{\eta_{1-s}(n)} \right|$$

$$\geq \lim_{n \to +\infty} \left| \frac{n^{1-s}}{n^s} \prod_{k=0}^{n+2} \frac{k+s}{k+1-s} \right| = \left| \frac{\Gamma(1-s)}{\Gamma(s)} \right|,$$

where Γ is the Gamma function.

References

[1] E.C. Titchmarsh, The Theory of the Riemann Zeta-Function (revised by D.R. Heath-Brown), Clarendon Press, Oxford. (1986).