



HAL
open science

Magnetisation Moment of a Bounded 3D Sample: Asymptotic Recovery from Planar Measurements on a Large Disk Using Fourier Analysis

Dmitry Ponomarev

► **To cite this version:**

Dmitry Ponomarev. Magnetisation Moment of a Bounded 3D Sample: Asymptotic Recovery from Planar Measurements on a Large Disk Using Fourier Analysis. 2022. hal-03813559v1

HAL Id: hal-03813559

<https://hal.science/hal-03813559v1>

Preprint submitted on 13 Oct 2022 (v1), last revised 9 Jan 2024 (v3)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

MAGNETISATION MOMENT OF A BOUNDED 3D SAMPLE: ASYMPTOTIC RECOVERY FROM PLANAR MEASUREMENTS ON A LARGE DISK USING FOURIER ANALYSIS

DMITRY PONOMAREV ^{1,2,3}

ABSTRACT. We consider the problem of reconstruction of the overall magnetisation vector (net moment) of a sample from partial data of the magnetic field. Namely, motivated by a concrete experimental set-up, we deal with a situation when the magnetic field is measured on a portion of the plane in vicinity of the sample and only one (normal to the plane) component of the field is available. We assume the measurement area to be a sufficiently large disk (lying in a horizontal plane above the sample) and we obtain a set of estimates for the components of the net moment vector with the accuracy asymptotically improving with the increase of the radius of the measurement disk. Compared to our previous preliminary results, the asymptotic estimates are now rigorously justified and higher-order estimates are given. The presented approach also elucidates the derivation of asymptotic estimates of an arbitrary order. The obtained results are illustrated numerically and their robustness with respect to noise is discussed.

1. INTRODUCTION

Constant advances in magnetometry allow measurements of magnetic fields of very low intensities with high spatial resolution. In particular, this opens new horizons in paleomagnetic contexts. Ancient rocks and meteorites possess remanent magnetisation and thus might preserve valuable records of a past magnetic activity on Earth and other planets, asteroids and satellites. Extraction of this relict magnetic information is a lucrative but challenging task. Deducing magnetisation of a geosample hinges on effective processing of the measurements of the magnetic field available in a nearest neighbourhood of the sample since the informative part of the field further away is very weak and significantly deteriorated by noise. In particular set-ups of scanning SQUID (Superconducting Quantum Interference Device) magnetometer or QDM (Quantum Diamond Microscope), measurements are available in a planar area above the sample, in a close vicinity of it, and such measurements typically feature only one component of the magnetic field. This is in contrast with more common settings that deal with magnetic fields of higher intensity and hence could, on a methodological level, rely on the classical dipole approximation of a sample valid in a far-away region.

In the present work, we are concerned with recovery of the overall magnetisation (the so-called net moment) of a sample rather than dealing with reconstruction of the entire magnetisation distribution. The latter is known to be an ill-posed problem, in particular, due to the presence of “silent” sources, i.e. magnetisations that do not produce magnetic field, see [3]. However, as it was shown in [1] for planar (thin-plate) magnetisation distributions, compactly supported silent sources do not contribute to the net moment. This statement fixes the non-uniqueness issue and makes the problem of net moment recovery a feasible task. In theory, this problem is even solvable in a closed form when measurements are available on the entire plane above the sample. In reality, however, the measurements are very limited and corrupted by the presence of noise which dominates the signal in distant regions. Therefore, we arrive at the problem of estimating the net moment of a sample from a magnetic field component available on some portion of the plane in proximity of the sample. Assuming circular geometry, we establish a set of estimates of the net moment components in terms of the size of the measurement area (disk) under condition that the radius of the disk is sufficiently large.

We do not intend here to provide neither physical nor mathematical description of the problem in any detailed fashion. Instead, we refer the reader to the set of previous publications [3, 8, 1, 5, 10, 7, 15, 9] and briefly introduce basic concepts that allow us to be more specific in formulating our main result and comparing it with relevant works.

We assume that the magnetic sample is described by a compactly supported vector distribution

$$\vec{\mathcal{M}}(\vec{x}) \equiv (\mathcal{M}_1(x_1, x_2, x_3), \mathcal{M}_2(x_1, x_2, x_3), \mathcal{M}_3(x_1, x_2, x_3))^T, \quad \text{supp } \vec{\mathcal{M}} =: Q,$$

¹FACTAS team, Centre Inria d’Université Côte d’Azur, France

²St. Petersburg Department of Steklov Mathematical Institute of Russian Academy of Sciences, Russia

³Contact: dmitry.ponomarev@asc.tuwien.ac.at

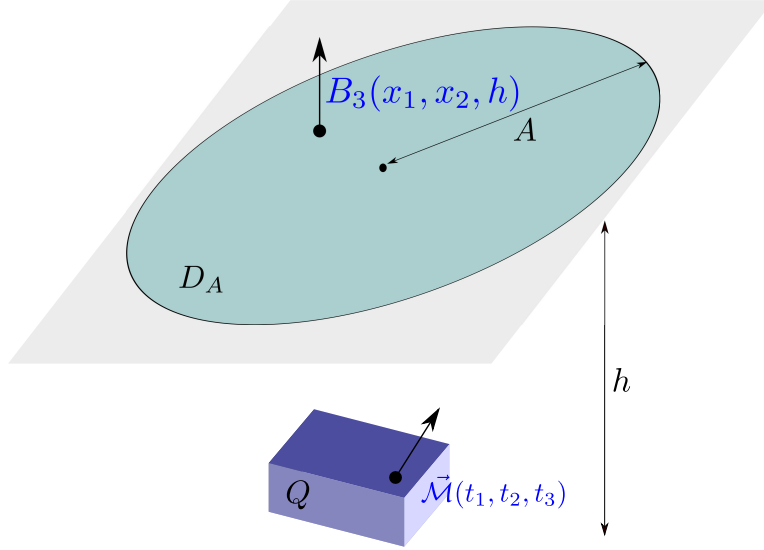


FIGURE 1.1. Schematic illustration of geometry of the problem

with some compact set $Q \subset \mathbb{R}^3$. The magnetic field produced by the sample can be expressed as $\vec{B}(\vec{x}) = -\nabla\Phi(\vec{x})$, where the scalar potential $\Phi(\vec{x})$ satisfies, in a distributional sense, the Poisson equation [3]

$$\Delta\Phi(\vec{x}) = \nabla \cdot \vec{\mathcal{M}}(\vec{x}), \quad \vec{x} \in \mathbb{R}^3.$$

Consequently, we can write [6, Sect. 2.4 Thm 1]

$$\Phi(\vec{x}) = -\frac{1}{4\pi} \iiint_Q \frac{1}{|\vec{x} - \vec{t}|} \nabla \cdot \vec{\mathcal{M}}(\vec{t}) d^3t, \quad \vec{x} \in \mathbb{R}^3 \setminus Q,$$

and (upon integration by parts) obtain

$$\Phi(\mathbf{x}, x_3) = \iiint_Q \frac{\mathcal{M}_1(\mathbf{t}, t_3)(x_1 - t_1) + \mathcal{M}_2(\mathbf{t}, t_3)(x_2 - t_2) + \mathcal{M}_3(\mathbf{t}, t_3)(x_3 - t_3)}{4\pi \left(|\mathbf{x} - \mathbf{t}|^2 + (x_3 - t_3)^2 \right)^{3/2}} d^3t, \quad \vec{x} \in \mathbb{R}^3 \setminus Q.$$

A physically measurable quantity is not the potential, but the magnetic field. In particular, the measurements are available for the vertical component (normal) of the magnetic field on the horizontal plane $x_3 = h$

$$(1.1) \quad B_3(\mathbf{x}, h) = -\partial_{x_3}\Phi(\mathbf{x}, h),$$

where $h > 0$ is some fixed constant. Explicitly, we have

$$(1.2) \quad B_3(\mathbf{x}, h) = \iiint_Q \frac{3(h - t_3)[\mathcal{M}_1(\mathbf{t}, t_3)(x_1 - t_1) + \mathcal{M}_2(\mathbf{t}, t_3)(x_2 - t_2)] + \mathcal{M}_3(\mathbf{t}, t_3)(2(h - t_3)^2 - |\mathbf{x} - \mathbf{t}|^2)}{4\pi \left(|\mathbf{x} - \mathbf{t}|^2 + (h - t_3)^2 \right)^{5/2}} d^3t.$$

Here, we employed bold symbols to denote \mathbb{R}^2 vectors, e.g. $\mathbf{x} \equiv (x_1, x_2)^T$. Also, the choice of the origin $\vec{x} = \vec{0}$ was assumed such that $\mathbf{x} = \mathbf{0}$ corresponds to the horizontal center of a minimal rectangular parallelepiped embedding Q , and $x_3 = 0$ corresponds to the (vertically) lowest point of Q .

The geometry of this setting is schematically shown in Figure 1.1.

A quantity of basic physical interest is the net magnetisation moment (which, as discussed above, is a constant vector equal to the overall magnetisation of the sample):

$$(1.3) \quad \vec{m} \equiv (m_1, m_2, m_3)^T := \iiint_Q \vec{\mathcal{M}}(\vec{x}) d^3x.$$

In what follows, we may use slightly different names for this quantity interchangeably: net moment, net magnetisation, zeroth-order algebraic moment of a magnetisation distribution.

Denote $D_A := \{\mathbf{x} \in \mathbb{R}^2 : |\mathbf{x}| < A\}$, the (horizontal) disk of radius A centered at the origin.

The main result of this work is summarised in the following theorem.

Theorem 1.1. Assume that $\vec{\mathcal{M}}$ is a distribution with a compact support $Q \subset \mathbb{R}^3$ producing a magnetic field whose vertical component B_3 is given by (1.2) for $x_3 = h$. Suppose that the measurement area disk D_A is sufficiently large so that the following condition holds

$$(1.4) \quad \sup_{\vec{t} \in Q, \mathbf{x} \in \mathbb{R}^2 \setminus D_A} \left| \frac{t_1^2 + t_2^2 + (h - t_3)^2}{x_1^2 + x_2^2} - 2 \frac{x_1 t_1 + x_2 t_2}{x_1^2 + x_2^2} \right| < 1.$$

Then, the components of the net moment (1.3) can be asymptotically estimated, with different orders of accuracy, using the formulas below.

First-order estimates:

$$(1.5) \quad m_j = 2 \iint_{D_A} x_j B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A}\right), \quad j \in \{1, 2\}.$$

Second-order estimates:

$$(1.6) \quad m_j = 2 \iint_{D_A} \left(1 + \frac{4x_j^2}{3A^2}\right) x_j B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A^2}\right), \quad j \in \{1, 2\},$$

$$(1.7) \quad m_3 = 2A \iint_{D_A} B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A^2}\right).$$

Third-order estimates:

$$(1.8) \quad m_j = \frac{2}{5} \iint_{D_A} \left[5 + 24 \left(\frac{x_j}{A}\right)^4\right] x_j B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A^3}\right), \quad j \in \{1, 2\},$$

$$(1.9) \quad m_3 = \frac{A}{4} \iint_{D_A} \left[5 + 40 \left(\frac{x_j}{A}\right)^4 - 128 \left(\frac{x_j}{A}\right)^6\right] B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A^3}\right), \quad j \in \{1, 2\}.$$

Fourth-order estimates:

$$(1.10) \quad m_j = \frac{2}{105} \iint_{D_A} \left[315 - 2016 \left(\frac{x_j}{A}\right)^4 + 19200 \left(\frac{x_j}{A}\right)^6 - 22400 \left(\frac{x_j}{A}\right)^8\right] x_j B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A^4}\right), \quad j \in \{1, 2\},$$

$$(1.11) \quad m_3 = \frac{A}{24} \iint_{D_A} \left[35 + 1792 \left(\frac{x_j}{A}\right)^6 - 3200 \left(\frac{x_j}{A}\right)^8\right] B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A^4}\right), \quad j \in \{1, 2\}.$$

Fifth-order estimates:

$$(1.12) \quad m_j = \frac{2}{693} \iint_{D_A} \left[693 - 158400 \left(\frac{x_j}{A}\right)^6 + 35200 \left(\frac{x_j}{A}\right)^8 - 677376 \left(\frac{x_j}{A}\right)^{10}\right] x_j B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A^5}\right), \quad j \in \{1, 2\}.$$

Remark 1.2. Here and onwards (except numerics in Section 3), for the sake of simplicity, we have assumed the system of physical units such that the constant of magnetic permeability of a vacuum μ_0 is 1. In general, the right-hand sides of expressions (1.1)–(1.2) should have the factor μ_0 , and, in Si units, $\mu_0 = 4\pi \cdot 10^{-7}$ N / A². Consequently, the right-hand sides of all the formulas (1.5)–(1.12) should, in principle, have the factor $1/\mu_0$.

Remark 1.3. It is clear that condition (1.4) can be written in an equivalent form as:

$$(1.13) \quad \sup_{\vec{t} \in Q, \mathbf{x} \in \mathbb{R}^2 \setminus D_A} \frac{(x_1 - t_1)^2 + (x_2 - t_2)^2 + (h - t_3)^2}{2(x_1^2 + x_2^2)} < 1.$$

The most relevant papers to the result of Theorem 1.1 are [1, 2] as well as preliminary works [12, Part 3], [4]. Namely, the current work stems from one of the methods introduced in [12, Part 3] and briefly summarised in [4]. There, the second-order asymptotic estimates given by (1.6)–(1.7) have already appeared, but have not been rigorously justified. The case of the leading-order asymptotic estimates (in particular, analogs of (1.5) and (1.7)) and their justification were considered for the case of a rectangular measurement area. The roots of the Fourier analysis of the magnetic field aimed to recover the net moment of a planar magnetisation (i.e. when support set Q is two-dimensional and parallel to the measurement plane) can already be found in [9, 3, 8]. In [1], the authors considered a more general problem of finding linear net moment estimators (analogous of polynomials appearing in the integrals of (1.5)–(1.12)), without resorting to the asymptotic analysis (hence with no assumption on the largeness

of the measurement area), but under restriction that the magnetisation distribution is planar and regular (square-integrable). Ill-posedness of the problem (due to sensitivity of the result to the magnetic field data) was recognised and its regularised version was proposed and solved numerically.

It should be mentioned that even though recovery of the net moment of a sample, our primary concern, is an important practical problem on its own right, it can also serve, under appropriate assumptions, as a preliminary step for finding a magnetisation distribution. Indeed, while it is unrealistic to retrieve 3 magnetisation components (functions or, more generally, distributions): $\mathcal{M}_1, \mathcal{M}_2, \mathcal{M}_3$ from the knowledge of only one function $B_3(\cdot, h)$ (see (1.2)), the problem simplifies significantly for a class of samples which are unidirectionally magnetised. Since the knowledge of the net moment of a sample implies a definite magnetisation direction, it is thus essential in this complete reconstruction procedure, see [8].

The structure of the paper is as follows. Section 2 has a twofold goal. First, it is meant to show how one idea based on elementary Fourier analysis can yield the simplest version of asymptotic net moment estimates for both tangential and normal components. Second, this section illustrates that, by means of a careful asymptotic analysis, the explicit estimates can be not only proved rigorously but also extended to higher orders. Hence, this material exactly constitutes the proof of Theorem 1.1. Then, in Section 3, we illustrate the results numerically on the case where the magnetisation has a singular support (i.e. the magnetisation distribution is a collection of dipoles) and deal with some practical aspects of the obtained estimates. Finally, we conclude with Section 4 summarising the work, discussing the results and outlining further research directions.

2. PROOF OF THEOREM 1.1

2.1. Some notation and preparatory transformations. Before we proceed with deriving rigorously formulas (1.5)–(1.12) and thus proving Theorem 1.1, let us introduce some useful notations. We shall use the following shortcut for the integral of the magnetisation distribution against monomials

$$(2.1) \quad \left\langle x_1^{j_1} x_2^{j_2} x_3^{j_3} M_n \right\rangle := \iiint_Q x_1^{j_1} x_2^{j_2} x_3^{j_3} \mathcal{M}_n(\vec{x}) d^3x, \quad n \in \{1, 2, 3\}, \quad j_1, j_2, j_3 \in \mathbb{N}_0,$$

where we denoted $\mathbb{N}_0 := \{0, 1, 2, \dots\}$, the set of natural numbers including zero. Following this convention, for the sake of brevity, we may also write, for example,

$$\langle (h - x_3) M_n \rangle := hm_n - \langle x_3 M_n \rangle = \iiint_Q (h - x_3) \mathcal{M}_n(\vec{x}) d^3x, \quad n \in \{1, 2, 3\}.$$

Note that, when magnetisation is not a function but merely a distribution, the integrals above should be understood as distributional pairing between a compactly supported distribution $\vec{\mathcal{M}}$ and smooth monomials $x_1^{j_1} x_2^{j_2} x_3^{j_3}$, $j_1, j_2, j_3 \in \mathbb{N}_0$.

Let us $\hat{\cdot}$ denote the two-dimensional Fourier transform which, by our convention, is defined as

$$\hat{f}(\mathbf{k}) \equiv \mathcal{F}[f](\mathbf{k}) := \iint_{\mathbb{R}^2} e^{2\pi i \mathbf{k} \cdot \mathbf{x}} f(\mathbf{x}) d^2x,$$

where $i = \sqrt{-1}$ stands for the imaginary unit, and $\mathbf{k} \cdot \mathbf{x} = k_1 x_1 + k_2 x_2$ is the Euclidean inner product. With this definition, the differentiation and convolution properties of Fourier transform have the form

$$(2.2) \quad \mathcal{F}[\partial_{x_j} f](\mathbf{k}) = -2\pi i k_j \hat{f}(\mathbf{k}), \quad \mathcal{F}[x_j f](\mathbf{k}) = \frac{1}{2\pi i} \partial_{k_j} \hat{f}(\mathbf{k}), \quad j \in \{1, 2\},$$

$$(2.3) \quad \mathcal{F}[f \star g](\mathbf{k}) := \iint_{\mathbb{R}^2} e^{2\pi i \mathbf{k} \cdot \mathbf{x}} \iint_{\mathbb{R}^2} f(\mathbf{x} - \mathbf{t}) g(\mathbf{t}) dt_1 dt_2 dx_1 dx_2 = \hat{f}(\mathbf{k}) \hat{g}(\mathbf{k}).$$

We also note that the Fourier transform of the two-dimensional Poisson kernel is well-known (see e.g. [14, Sect. 4.2]), that is, for any $H > 0$, we have

$$(2.4) \quad \mathcal{F} \left[\frac{H}{2\pi (|\mathbf{x}|^2 + H^2)^{3/2}} \right] (\mathbf{k}) = e^{-2\pi H |\mathbf{k}|}, \quad \mathbf{k} \in \mathbb{R}^2.$$

Let us now rewrite (1.2) as

$$(2.5) \quad B_3(\mathbf{x}, h) = -\frac{1}{4\pi} \iiint_Q \left[\left(\mathcal{M}_1(\mathbf{t}, t_3) \frac{\partial}{\partial x_1} + \mathcal{M}_2(\mathbf{t}, t_3) \frac{\partial}{\partial x_2} \right) \frac{h - t_3}{\left(|\mathbf{x} - \mathbf{t}|^2 + (h - t_3)^2 \right)^{3/2}} \right. \\ \left. + \mathcal{M}_3(\mathbf{t}, t_3) \frac{\partial}{\partial x_3} \Big|_{x_3=h} \frac{x_3 - t_3}{\left(|\mathbf{x} - \mathbf{t}|^2 + (x_3 - t_3)^2 \right)^{3/2}} \right] d^3t.$$

Taking Fourier transform of (2.5) in x_1, x_2 variables, we use (2.2), (2.3) and employ (2.4) twice: with $H := h - t_3$ in the first line of (2.5), and with $H := x_3 - t_3$ in the second one. We thus arrive at

$$(2.6) \quad \hat{B}_3(\mathbf{k}, h) = \pi \int_{Q_3} e^{-2\pi(h-t_3)|\mathbf{k}|} \left(ik_1 \hat{\mathcal{M}}_1(\mathbf{k}, t_3) + ik_2 \hat{\mathcal{M}}_2(\mathbf{k}, t_3) + |\mathbf{k}| \hat{\mathcal{M}}_3(\mathbf{k}, t_3) \right) dt_3,$$

where Q_3 denotes the vertical projection of the set Q .

We note that even though (1.5)–(1.6), (1.8), (1.10), (1.12) give estimates for tangential net moment components m_1, m_2 , we shall restrict ourselves to dealing only with m_1 . The situation with m_2 is completely analogous.

First, we are going to illustrate our strategy of the derivation of asymptotic estimates of the net magnetisation moment. Here, we shall be only concerned with the low-order formulas of Theorem 1.1 and we shall omit a rigorous justification step. Then, we shall proceed with formal justification and extension of the result to higher orders.

2.2. Illustration of the basic idea of the derivation of the net moment estimates. We are going to focus on deriving (1.5).

We take $k_2 = 0$ in expression (2.6) to obtain

$$(2.7) \quad \hat{B}_3(k_1, 0, h) = \pi \int_{Q_3} e^{-2\pi(h-t_3)|k_1|} \left(ik_1 \hat{\mathcal{M}}_1(k_1, 0, t_3) + |k_1| \hat{\mathcal{M}}_3(k_1, 0, t_3) \right) dt_3.$$

Since the magnetisation distribution $\vec{\mathcal{M}}$ is compactly supported, the Fourier transforms $\hat{\mathcal{M}}_j(\mathbf{k}, t_3)$, $j \in \{1, 2, 3\}$, are entire functions in $k_1, k_2 \in \mathbb{C}$ for each $t_3 \in Q_3$, according to the Paley-Wiener theory (see e.g. [13, Thm 4.1]). In particular, power-series expansion of $\hat{\mathcal{M}}_j(k_1, 0, t_3)$, $j \in \{1, 2, 3\}$, about the origin $k_1 = 0$ of the complex plane ($\text{Re } k_1, \text{Im } k_1$) gives

$$\hat{\mathcal{M}}_j(k_1, 0, t_3) = \hat{\mathcal{M}}_j(\mathbf{0}, t_3) + \partial_{k_1} \hat{\mathcal{M}}_j(\mathbf{0}, t_3) k_1 + \frac{1}{2} \partial_{k_1}^2 \hat{\mathcal{M}}_j(\mathbf{0}, t_3) k_1^2 + \mathcal{O}\left(|k_1|^3\right), \quad j \in \{1, 2, 3\}.$$

Combining this expansion with the straightforward identities

$$m_j = \int_{Q_3} \hat{\mathcal{M}}_j(\mathbf{0}, t_3) dt_3, \quad j \in \{1, 2, 3\}, \\ 2\pi i \langle M_j x_1 \rangle = \int_{Q_3} \partial_{k_1} \hat{\mathcal{M}}_j(\mathbf{0}, t_3) dt_3, \quad j \in \{1, 2, 3\},$$

and the Taylor expansion in $|k_1|$ of the exponential factor in (2.7)

$$e^{-2\pi(h-t_3)|k_1|} = 1 - 2\pi(h-t_3)|k_1| + 2\pi^2(h-t_3)^2|k_1|^2 + \mathcal{O}\left(|k_1|^3\right),$$

we obtain

$$\hat{B}_3(k_1, 0, h) = \text{Re } \hat{B}_3(k_1, 0, h) + i \text{Im } \hat{B}_3(k_1, 0, h), \quad k_1 \in \mathbb{R}, \\ \text{Re } \hat{B}_3(k_1, 0, h) = \pi m_3 |k_1| - 2\pi^2 (\langle x_1 M_1 \rangle + \langle (h-x_3) M_3 \rangle) |k_1|^2 \\ + 2\pi^3 \left(2 \langle (h-x_3) x_1 M_1 \rangle + \langle (h-x_3)^2 M_3 \rangle - \langle x_1^2 M_3 \rangle \right) |k_1|^3 + \mathcal{O}\left(|k_1|^4\right),$$

$$\text{Im } \hat{B}_3(k_1, 0, h) = \pi m_1 k_1 - 2\pi^2 (\langle (h-x_3) M_1 \rangle - \langle x_1 M_3 \rangle) k_1 |k_1| \\ - 2\pi^3 \left(\langle x_1^2 M_1 \rangle - \langle (h-x_3)^2 M_1 \rangle + 2 \langle (h-x_3) x_1 M_3 \rangle \right) k_1^3 + \mathcal{O}\left(|k_1|^4\right).$$

Here, the remainder term is uniformly small for all $t_3 \in Q_3$ due to the boundedness of the set Q_3 .

On the other hand, we can write

$$(2.10) \quad \hat{B}_3(k_1, 0, h) = \iint_{D_A} e^{2\pi i k_1 x_1} B_3(\mathbf{x}, h) d^2x + \iint_{\mathbb{R}^2 \setminus D_A} e^{2\pi i k_1 x_1} B_3(\mathbf{x}, h) d^2x.$$

We note that in the first term on the right-hand side of (2.10), the integration range is finite and hence an expansion in powers of k_1 simply follows from that of the exponential factor:

$$(2.11) \quad \iint_{D_A} e^{2\pi i k_1 x_1} B_3(\mathbf{x}, h) d^2x = 1 + 2\pi i k_1 \iint_{D_A} x_1 B_3(\mathbf{x}, h) d^2x - 2\pi^2 k_1^2 \iint_{D_A} x_1^2 B_3(\mathbf{x}, h) d^2x \\ - \frac{4\pi^3}{3} i k_1^3 \iint_{D_A} x_1^3 B_3(\mathbf{x}, h) d^2x + \mathcal{O}(|k_1|^4),$$

and hence

$$(2.12) \quad \iint_{D_A} \cos(2\pi k_1 x_1) B_3(\mathbf{x}, h) d^2x = \iint_{D_A} B_3(\mathbf{x}, h) d^2x - 2\pi^2 k_1^2 \iint_{D_A} x_1^2 B_3(\mathbf{x}, h) d^2x + \mathcal{O}(|k_1|^4),$$

$$(2.13) \quad \iint_{D_A} \sin(2\pi k_1 x_1) B_3(\mathbf{x}, h) d^2x = 2\pi k_1 \iint_{D_A} x_1 B_3(\mathbf{x}, h) d^2x - \frac{4\pi^3}{3} k_1^3 \iint_{D_A} x_1^3 B_3(\mathbf{x}, h) d^2x + \mathcal{O}(|k_1|^5).$$

Producing an expansion in powers of k_1 of the second term in (2.10) is much less straightforward and requires a preliminary simplification. More precisely, we expand the field $B_3(\mathbf{x}, h)$ for large $|\mathbf{x}|$ and, assuming largeness of the region D_A , retain only first few terms of this expansion. Namely, from (1.2), we have

$$(2.14) \quad B_3(\mathbf{x}, h) = -\frac{m_3}{4\pi |\mathbf{x}|^3} + \frac{3}{4\pi} \frac{(\langle (h-x_3) M_1 \rangle - \langle x_1 M_3 \rangle) x_1 + (\langle (h-x_3) M_2 \rangle - \langle x_2 M_3 \rangle) x_2}{|\mathbf{x}|^5} + \mathcal{O}\left(\frac{1}{|\mathbf{x}|^5}\right).$$

Consequently, passing to the polar coordinates using $x_1 = r \cos \theta$, $x_2 = r \sin \theta$, $d^2x = r dr d\theta$, we can write

$$\iint_{\mathbb{R}^2 \setminus D_A} e^{2\pi i k_1 x_1} B_3(\mathbf{x}, h) d^2x = -\frac{m_3}{4\pi} \int_A^\infty \int_0^{2\pi} e^{2\pi i k_1 r \cos \theta} d\theta \frac{dr}{r^2} \\ + \frac{3}{4\pi} (\langle (h-x_3) M_1 \rangle - \langle x_1 M_3 \rangle) \int_A^\infty \int_0^{2\pi} e^{2\pi i k_1 r \cos \theta} \cos \theta d\theta \frac{dr}{r^3} \\ + \frac{3}{4\pi} (\langle (h-x_3) M_2 \rangle - \langle x_2 M_3 \rangle) \int_A^\infty \int_0^{2\pi} e^{2\pi i k_1 r \cos \theta} \sin \theta d\theta \frac{dr}{r^3} \\ + \mathcal{R}_{A, k_1},$$

where the residue term \mathcal{R}_{A, k_1} is expected to be $\mathcal{O}(1/A^3)$ for sufficiently small values of $|k_1|$.

Furthermore, taking real and imaginary parts of both sides, we obtain, respectively,

$$(2.15) \quad \iint_{\mathbb{R}^2 \setminus D_A} \cos(2\pi k_1 x_1) B_3(\mathbf{x}, h) d^2x = -\frac{m_3}{4\pi} \int_A^\infty \int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) d\theta \frac{dr}{r^2} + \text{Re } \mathcal{R}_{A, k_1},$$

$$(2.16) \quad \iint_{\mathbb{R}^2 \setminus D_A} \sin(2\pi k_1 x_1) B_3(\mathbf{x}, h) d^2x = \frac{3}{4\pi} (\langle (h-x_3) M_1 \rangle - \langle x_1 M_3 \rangle) \int_A^\infty \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) \cos \theta d\theta \frac{dr}{r^3} + \text{Im } \mathcal{R}_{A, k_1},$$

where we took into account that

$$\int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) \sin \theta d\theta = \int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) \cos \theta d\theta = 0, \\ \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) d\theta = \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) \sin \theta d\theta = 0,$$

according to the results of Lemma A.4.

Next, as it turns out (see Subsection 2.3 for more details), the integrals on the right-hand sides of (2.15)–(2.16) can be evaluated explicitly in terms of some cylindrical functions. Known representations of these special functions lead to the desired asymptotic expansions in powers of k_1 . In particular, for small $|k_1|$, we can deduce that

$$(2.17) \quad \int_A^\infty \int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) d\theta \frac{dr}{r^2} = 2\pi |k_1| \int_{2\pi|k_1|A}^\infty \int_0^{2\pi} \cos(r \cos \theta) d\theta \frac{dr}{r^2} \\ = \frac{2\pi}{A} - 4\pi^2 |k_1| + 2\pi^3 A |k_1|^2 + \mathcal{O}(A^2 |k_1|^3),$$

$$(2.18) \quad \int_A^\infty \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) \cos \theta d\theta \frac{dr}{r^3} = 4\pi^2 k_1 |k_1| \int_{2\pi|k_1|A}^\infty \int_0^{2\pi} \sin(r \cos \theta) \cos \theta d\theta \frac{dr}{r^3} \\ = \frac{2\pi^2}{A} k_1 - \frac{8\pi^3}{3} k_1 |k_1| + \pi^4 A k_1^3 + \mathcal{O}\left(A^3 |k_1|^5\right),$$

where the notation $\mathcal{O}\left(A^2 |k_1|^3\right)$ also hides the terms of powers of $|k_1|$ higher than 3 regardless of the presence of the A factors such as $\mathcal{O}\left(A^3 |k_1|^4\right)$.

Therefore, by taking the real part of (2.10) and using (2.12), (2.15) and (2.17), we obtain

$$\operatorname{Re} \hat{B}_3(k_1, 0, h) = \iint_{D_A} B_3(\mathbf{x}, h) d^2x - 2\pi^2 k_1^2 \iint_{D_A} x_1^2 B_3(\mathbf{x}, h) d^2x \\ - \frac{m_3}{4\pi} \left(\frac{2\pi}{A} - 4\pi^2 |k_1| + 2\pi^3 A |k_1|^2 \right) + \operatorname{Re} \mathcal{R}_{A, k_1} + \mathcal{O}\left(A^2 |k_1|^3\right).$$

Comparing this with (2.8) and, in particular, evaluating both expressions at $k_1 = 0$, we arrive at the following identity:

$$(2.19) \quad \iint_{D_A} B_3(\mathbf{x}, h) d^2x - \frac{m_3}{2A} + \mathcal{O}\left(\frac{1}{A^3}\right) = 0,$$

where we took into account that $\operatorname{Re} \mathcal{R}_{A, k_1}|_{k_1=0} = \mathcal{O}\left(1/A^3\right)$.

Similarly, taking the imaginary part of (2.10), we combine (2.13), (2.16) and (2.18) to deduce that

$$\operatorname{Im} \hat{B}_3(k_1, 0, h) = 2\pi k_1 \iint_{D_A} x_1 B_3(\mathbf{x}, h) d^2x - \frac{4\pi^3}{3} k_1^3 \iint_{D_A} x_1^3 B_3(\mathbf{x}, h) d^2x \\ + \frac{3}{4\pi} (\langle (h - x_3) M_1 \rangle - \langle x_1 M_3 \rangle) \left(\frac{2\pi^2}{A} k_1 - \frac{8\pi^3}{3} k_1 |k_1| + \pi^4 A k_1^3 \right) + \operatorname{Im} \mathcal{R}_{A, k_1} + \mathcal{O}\left(A^3 |k_1|^5\right).$$

Comparison of this expression with (2.9) and matching the coefficients of the k_1 terms yields

$$(2.20) \quad 2\pi \iint_{D_A} x_1 B_3(\mathbf{x}, h) d^2x + \frac{3\pi}{2A} (\langle (h - x_3) M_1 \rangle - \langle x_1 M_3 \rangle) + \mathcal{O}\left(\frac{1}{A^3}\right) = \pi m_1,$$

where we assumed that, for sufficiently small $|k_1|$, we have $\operatorname{Im} \mathcal{R}_{A, k_1} = \mathcal{O}\left(1/A^3\right)$.

While (2.19), (2.20) imply estimates (1.7), (1.5), respectively, the derivation given above was not rigorous and required additional assumptions on the residue term \mathcal{R}_{A, k_1} which was reasonably deemed to be sufficiently small for large A but was not estimated uniformly in k_1 . We shall now proceed with rigorous analysis which will also make it possible to derive higher-order analogs of estimates (1.5), (1.7).

2.3. Rigorous analysis and higher-order asymptotic estimates. Let us start by improving estimate (2.14). To this effect, we use the following elementary Taylor expansions, convergent for $|z| < 1$,

$$\frac{1}{(1+z)^{3/2}} = 1 - \frac{3}{2}z + \frac{15}{8}z^2 - \frac{35}{16}z^3 + \mathcal{O}(z^4),$$

$$\frac{1}{(1+z)^{5/2}} = 1 - \frac{5}{2}z + \frac{35}{8}z^2 + \mathcal{O}(z^3),$$

to obtain, for fixed $t_1, t_2, t_3, h \in \mathbb{R}$,

(2.21)

$$\frac{1}{\left[(x_1 - t_1)^2 + (x_2 - t_2)^2 + (h - t_3)^2\right]^{3/2}} = \frac{1}{(x_1^2 + x_2^2)^{3/2}} \left[1 - 2 \frac{x_1 t_1 + x_2 t_2}{x_1^2 + x_2^2} + \frac{t_1^2 + t_2^2 + (h - t_3)^2}{x_1^2 + x_2^2} \right]^{-3/2} \\ = \frac{1}{(x_1^2 + x_2^2)^{3/2}} \left[1 + 3 \frac{x_1 t_1 + x_2 t_2}{x_1^2 + x_2^2} - \frac{3}{2} \frac{t_1^2 + t_2^2 + (h - t_3)^2}{x_1^2 + x_2^2} + \frac{15}{2} \frac{(x_1 t_1 + x_2 t_2)^2}{(x_1^2 + x_2^2)^2} \right. \\ \left. - \frac{15}{2} \frac{(x_1 t_1 + x_2 t_2) (t_1^2 + t_2^2 + (h - t_3)^2)}{(x_1^2 + x_2^2)^2} + \frac{35}{2} \frac{(x_1 t_1 + x_2 t_2)^3}{(x_1^2 + x_2^2)^3} \right] + \mathcal{O}\left(\frac{1}{|\mathbf{x}|^7}\right),$$

$$(2.22) \quad \frac{1}{\left[(x_1 - t_1)^2 + (x_2 - t_2)^2 + (h - t_3)^2\right]^{5/2}} = \frac{1}{(x_1^2 + x_2^2)^{5/2}} \left[1 + 5 \frac{x_1 t_1 + x_2 t_2}{x_1^2 + x_2^2} - \frac{5}{2} \frac{t_1^2 + t_2^2 + (h - t_3)^2}{x_1^2 + x_2^2} \right. \\ \left. + \frac{35}{2} \frac{(x_1 t_1 + x_2 t_2)^2}{(x_1^2 + x_2^2)^2} \right] + \mathcal{O}\left(\frac{1}{|\mathbf{x}|^8}\right),$$

where x_1, x_2 are sufficiently large so that

$$(2.23) \quad \left| \frac{t_1^2 + t_2^2 + (h - t_3)^2}{x_1^2 + x_2^2} - 2 \frac{x_1 t_1 + x_2 t_2}{x_1^2 + x_2^2} \right| < 1.$$

Expansions (2.21)–(2.22) imply that, for $|\mathbf{x}| \gg 1$, (1.2) can be written as

$$(2.24) \quad B_3(\mathbf{x}, h) = B_3^{\text{asympt}}(\mathbf{x}, h) + \mathcal{O}\left(\frac{1}{|\mathbf{x}|^7}\right),$$

with

$$(2.25) \quad B_3^{\text{asympt}}(\mathbf{x}, h) := \frac{a_0}{|\mathbf{x}|^3} + \frac{a_1^{(1)} x_1 + a_1^{(2)} x_2}{|\mathbf{x}|^5} + \frac{a_2}{|\mathbf{x}|^5} + \frac{a_3^{(1)} x_1^2 + a_3^{(2)} x_2^2 + a_3^{(3)} x_1 x_2}{|\mathbf{x}|^7} \\ + \frac{a_4^{(1)} x_1 + a_4^{(2)} x_2}{|\mathbf{x}|^7} + \frac{a_5^{(1)} x_1^3 + a_5^{(2)} x_2^3 + a_5^{(3)} x_1^2 x_2 + a_5^{(4)} x_1 x_2^2}{|\mathbf{x}|^9},$$

$$(2.26) \quad a_0 := -\frac{m_3}{4\pi},$$

$$(2.27) \quad a_1^{(1)} := \frac{3}{4\pi} [\langle (h - x_3) M_1 \rangle - \langle x_1 M_3 \rangle], \quad a_1^{(2)} := \frac{3}{4\pi} [\langle (h - x_3) M_2 \rangle - \langle x_2 M_3 \rangle],$$

$$(2.28) \quad a_2 := -\frac{3}{8\pi} \left[2 \langle (h - x_3) x_1 M_1 \rangle + 2 \langle (h - x_3) x_2 M_2 \rangle - 3 \langle (h - x_3)^2 M_3 \rangle - \langle x_1^2 M_3 \rangle - \langle x_2^2 M_3 \rangle \right],$$

$$(2.29) \quad a_3^{(1)} := \frac{15}{8\pi} [2 \langle (h - x_3) x_1 M_1 \rangle - \langle x_1^2 M_3 \rangle], \quad a_3^{(2)} := \frac{15}{8\pi} [2 \langle (h - x_3) x_2 M_2 \rangle - \langle x_2^2 M_3 \rangle],$$

$$(2.30) \quad a_3^{(3)} := \frac{15}{4\pi} [\langle (h - x_3) x_2 M_1 \rangle + \langle (h - x_3) x_1 M_2 \rangle - \langle x_1 x_2 M_3 \rangle],$$

$$(2.31) \quad a_4^{(1)} := -\frac{15}{8\pi} \left[3 \langle (h - x_3) x_1^2 M_1 \rangle + \langle (h - x_3) x_2^2 M_1 \rangle + \langle (h - x_3)^3 M_1 \rangle + 2 \langle (h - x_3) x_1 x_2 M_2 \rangle \right. \\ \left. - \langle x_1^3 M_3 \rangle - \langle x_1 x_2^2 M_3 \rangle - 3 \langle (h - x_3)^2 x_1 M_3 \rangle \right],$$

$$(2.32) \quad a_4^{(2)} := -\frac{15}{8\pi} \left[3 \langle (h - x_3) x_2^2 M_2 \rangle + \langle (h - x_3) x_1^2 M_2 \rangle + \langle (h - x_3)^3 M_2 \rangle + 2 \langle (h - x_3) x_1 x_2 M_1 \rangle \right. \\ \left. - \langle x_2^3 M_3 \rangle - \langle x_1^2 x_2 M_3 \rangle - 3 \langle (h - x_3)^2 x_2 M_3 \rangle \right],$$

$$(2.33) \quad a_5^{(1)} := \frac{35}{8\pi} [3 \langle (h - x_3) x_1^2 M_1 \rangle - \langle x_1^3 M_3 \rangle], \quad a_5^{(2)} := \frac{35}{8\pi} [3 \langle (h - x_3) x_2^2 M_2 \rangle - \langle x_2^3 M_3 \rangle],$$

$$(2.34) \quad a_5^{(3)} := \frac{105}{8\pi} [\langle (h - x_3) x_1^2 M_2 \rangle + 2 \langle (h - x_3) x_1 x_2 M_1 \rangle - \langle x_1^2 x_2 M_3 \rangle],$$

$$(2.35) \quad a_5^{(4)} := \frac{105}{8\pi} [\langle (h - x_3) x_2^2 M_1 \rangle + 2 \langle (h - x_3) x_1 x_2 M_2 \rangle - \langle x_1 x_2^2 M_3 \rangle],$$

and condition (2.23), for our particular context, rewrites as (1.4).

We are going to pursue the idea outlined in the previous subsection. Namely, comparing a series expansion of (2.7) about $k_1 = 0$ with that of (2.10), we shall deduce a set of identities which relate magnetisation moments to the integrals of the measured data $B_3(\mathbf{x}, h)$ on D_A . More precisely, using compactness of the support of the magnetisation \mathcal{M} , it follows from (2.7) that $\text{Re } \hat{B}_3(k_1, 0, h)$ is a convergent series in powers of $|k_1|$ whereas $\text{Im } \hat{B}_3(k_1, 0, h)$ is a power series in $|k_1|$ multiplied by k_1 . To facilitate the situation of matching the coefficients of

different representations, we shall focus on the region $k_1 > 0$. Consequently, in what follows, evaluation at $k_1 = 0^+$ will mean the limiting value at $k_1 = 0$ taken from the positive semiaxis ($k_1 > 0$).

2.3.1. *Tangential components of the net moment.* Expanding the integrand in (2.7) in power series in a positive neighborhood of $k_1 = 0$ and taking the imaginary part, we deduce:

$$(2.36) \quad \partial_{k_1} \left[\text{Im } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = \pi m_1 =: d_1,$$

$$(2.37) \quad \frac{1}{6} \partial_{k_1}^3 \left[\text{Im } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = 2\pi^3 \left[\langle (h-x_3)^2 M_1 \rangle - \langle x_1^2 M_1 \rangle - 2 \langle (h-x_3) x_1 M_3 \rangle \right] =: d_3,$$

$$(2.38) \quad \frac{1}{5!} \partial_{k_1}^5 \left[\text{Im } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = \frac{2\pi^5}{3} \left[\langle (h-x_3)^4 M_1 \rangle - 6 \langle (h-x_3)^2 x_1^2 M_1 \rangle + \langle x_1^4 M_1 \rangle - 4 \langle (h-x_3)^3 x_1 M_3 \rangle + 4 \langle (h-x_3) x_1^3 M_3 \rangle \right] =: d_5,$$

$$(2.39) \quad \frac{1}{7!} \partial_{k_1}^7 \left[\text{Im } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = \frac{4\pi^7}{45} \left[\langle (h-x_3)^6 M_1 \rangle - 15 \langle (h-x_3)^4 x_1^2 M_1 \rangle + 15 \langle (h-x_3)^2 x_1^4 M_1 \rangle - \langle x_1^6 M_1 \rangle - 6 \langle (h-x_3)^5 x_1 M_3 \rangle + 20 \langle (h-x_3)^3 x_1^3 M_3 \rangle - 6 \langle (h-x_3) x_1^5 M_3 \rangle \right] =: d_7,$$

$$(2.40) \quad \frac{1}{9!} \partial_{k_1}^9 \left[\text{Im } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = \frac{2\pi^9}{315} \left[\langle (h-x_3)^8 M_1 \rangle - 28 \langle (h-x_3)^6 x_1^2 M_1 \rangle + 70 \langle (h-x_3)^4 x_1^4 M_1 \rangle - 28 \langle (h-x_3)^2 x_1^6 M_1 \rangle + \langle x_1^8 M_1 \rangle - 8 \langle (h-x_3)^7 x_1 M_3 \rangle + 56 \langle (h-x_3)^5 x_1^3 M_3 \rangle - 56 \langle (h-x_3)^3 x_1^5 M_3 \rangle + 8 \langle (h-x_3) x_1^7 M_3 \rangle \right] =: d_9,$$

$$(2.41) \quad \frac{1}{11!} \partial_{k_1}^{11} \left[\text{Im } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = \frac{4\pi^{11}}{14175} \left[\langle (h-x_3)^{10} M_1 \rangle - 45 \langle (h-x_3)^8 x_1^2 M_1 \rangle + 210 \langle (h-x_3)^6 x_1^4 M_1 \rangle - 210 \langle (h-x_3)^4 x_1^6 M_1 \rangle + 45 \langle (h-x_3)^2 x_1^8 M_1 \rangle - \langle x_1^{10} M_1 \rangle - 10 \langle (h-x_3)^9 x_1 M_3 \rangle + 120 \langle (h-x_3)^7 x_1^3 M_3 \rangle - 252 \langle (h-x_3)^5 x_1^5 M_3 \rangle + 120 \langle (h-x_3)^3 x_1^7 M_3 \rangle - 10 \langle (h-x_3) x_1^9 M_3 \rangle \right] =: d_{11}.$$

On the other hand, from (2.10), we have, for $n \in \mathbb{N}_0$,

$$(2.42) \quad \partial_{k_1}^{2n+1} \left[\text{Im } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = (-1)^n (2\pi)^{2n+1} \iint_{D_A} x_1^{2n+1} B_3(\mathbf{x}, h) d^2x + \partial_{k_1}^{2n+1} \left(\iint_{\mathbb{R}^2 \setminus D_A} \sin(2\pi k_1 x_1) B_3^{\text{asympt}}(\mathbf{x}, h) d^2x \right) \Big|_{k_1=0^+} + \partial_{k_1}^{2n+1} \left(\iint_{\mathbb{R}^2 \setminus D_A} \sin(2\pi k_1 x_1) [B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h)] d^2x \right) \Big|_{k_1=0^+}.$$

We shall proceed in 3 steps. First, we evaluate the integral

$$(2.43) \quad \mathcal{I}^{\sin} \equiv \mathcal{I}^{\sin}(k_1, A) := \iint_{\mathbb{R}^2 \setminus D_A} \sin(2\pi k_1 x_1) B_3^{\text{asympt}}(\mathbf{x}, h) d^2x,$$

and compute its derivatives appearing on the second line of (2.42), hence obtaining a set of valuable identities. Second, we estimate the derivatives of the remainder

$$(2.44) \quad \mathcal{R}_{2n+1}^{\sin} \equiv \mathcal{R}_{2n+1}^{\sin}(A) := \partial_{k_1}^{2n+1} \left(\iint_{\mathbb{R}^2 \setminus D_A} \sin(2\pi k_1 x_1) [B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h)] d^2 x \right) \Big|_{k_1=0^+}, \quad n \in \mathbb{N}_0,$$

in order to show that the contribution of the term in the third line of (2.42) is not significant for large A (for the chosen order of the asymptotic expansion). Finally, at the last step, we combine the obtained identities and derive asymptotic formulas (1.5)–(1.6), (1.8), (1.10), (1.12) in a rigorously justified fashion.

Step 1: Derivation of the set of identities. Using (2.25) and passing to the polar coordinates using $x = r \cos \theta$, $y = r \sin \theta$, $d^2 x = r dr d\theta$, we obtain from (2.43)

$$(2.45) \quad \mathcal{I}^{\sin} = a_1^{(1)} \mathcal{I}_1^{\sin} + a_4^{(1)} \mathcal{I}_2^{\sin} + a_5^{(1)} \mathcal{I}_3^{\sin} + a_5^{(4)} \mathcal{I}_4^{\sin},$$

where

$$(2.46) \quad \mathcal{I}_1^{\sin} := \int_A^\infty \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) \cos \theta d\theta \frac{dr}{r^3},$$

$$(2.47) \quad \mathcal{I}_2^{\sin} := \int_A^\infty \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) \cos \theta d\theta \frac{dr}{r^5},$$

$$(2.48) \quad \mathcal{I}_3^{\sin} := \int_A^\infty \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) \cos^3 \theta d\theta \frac{dr}{r^5},$$

$$(2.49) \quad \mathcal{I}_4^{\sin} := \int_A^\infty \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) \cos \theta \sin^2 \theta d\theta \frac{dr}{r^5}.$$

Here, we used results (A.42)–(A.43) of Lemma A.4 multiple times to deduce vanishing of the integrals associated with the terms which involve $a_0, a_1^{(2)}, a_2, a_3^{(j)}, j \in \{1, 2, 3\}, a_4^{(2)}, a_5^{(2)}, a_5^{(3)}$.

We now employ the integral representation of Bessel functions, given in (A.5), to rewrite, for $k_1 > 0$,

$$(2.50) \quad \mathcal{I}_1^{\sin} = 2\pi \int_A^\infty J_1(2\pi k_1 r) \frac{dr}{r^3} = (2\pi)^3 k_1^2 \int_{2\pi k_1 A}^\infty \frac{J_1(x)}{x^3} dx,$$

$$(2.51) \quad \mathcal{I}_2^{\sin} = 2\pi \int_A^\infty J_1(2\pi k_1 r) \frac{dr}{r^5} = (2\pi)^5 k_1^4 \int_{2\pi k_1 A}^\infty \frac{J_1(x)}{x^5} dx,$$

$$(2.52) \quad \begin{aligned} \mathcal{I}_3^{\sin} &= -2\pi \int_A^\infty J_1''(2\pi k_1 r) \frac{dr}{r^5} = -(2\pi)^5 k_1^4 \int_{2\pi k_1 A}^\infty \frac{J_1''(x)}{x^5} dx \\ &= (2\pi)^5 k_1^4 \left(\frac{5J_1(\rho)}{\rho^6} + \frac{J_1'(\rho)}{\rho^5} - 30 \int_\rho^\infty \frac{J_1(x)}{x^7} dx \right) \Big|_{\rho=2\pi k_1 A}, \end{aligned}$$

$$(2.53) \quad \begin{aligned} \mathcal{I}_4^{\sin} &= 2\pi \int_A^\infty [J_1(2\pi k_1 r) + J_1''(2\pi k_1 r)] \frac{dr}{r^5} = (2\pi)^5 k_1^4 \int_{2\pi k_1 A}^\infty \frac{J_1(x) + J_1''(x)}{x^5} dx, \\ &= -(2\pi)^5 k_1^4 \left(\frac{5J_1(\rho)}{\rho^6} + \frac{J_1'(\rho)}{\rho^5} - \int_\rho^\infty \frac{J_1(x)}{x^5} dx - 30 \int_\rho^\infty \frac{J_1(x)}{x^7} dx \right) \Big|_{\rho=2\pi k_1 A}. \end{aligned}$$

Note that, in (2.52)–(2.53), we employed integration by parts twice using the asymptotic behavior of J_1 given in (A.6).

Using the results of Lemmas A.2–A.3, we have

$$(2.54) \quad \begin{aligned} \mathcal{I}_1^{\sin} &= \frac{(2\pi)^2 k_1}{A} \left(\rho \int_\rho^\infty \frac{J_1(x)}{x^3} dx \right) \Big|_{\rho=2\pi k_1 A} \\ &= \frac{(2\pi)^2 k_1}{3A} \left[J_0(\rho) + \frac{J_1(\rho)}{\rho} - \rho - \rho J_1(\rho) + \rho^2 J_0(\rho) - \frac{\pi}{2} \rho^2 J_0(\rho) H_1(\rho) + \frac{\pi}{2} \rho^2 J_1(\rho) H_0(\rho) \right] \Big|_{\rho=2\pi k_1 A}, \end{aligned}$$

$$\begin{aligned}
(2.55) \quad \mathcal{I}_2^{\sin} &= \frac{(2\pi)^2 k_1}{A^3} \left(\rho^3 \int_{\rho}^{\infty} \frac{J_1(x)}{x^5} dx \right) \Big|_{\rho=2\pi k_1 A} \\
&= \frac{7(2\pi)^2 k_1}{105A^3} \left[\frac{4J_1(\rho)}{\rho} + J_1'(\rho) - \frac{\rho^2 J_0(\rho)}{3} - \frac{\rho J_1(\rho)}{3} + \frac{\rho^3}{3} + \frac{\rho^3 J_1(\rho)}{3} \right. \\
&\quad \left. - \frac{\rho^4 J_0(\rho)}{3} + \frac{\pi}{6} \rho^4 J_0(\rho) H_1(\rho) - \frac{\pi}{6} \rho^4 J_1(\rho) H_0(\rho) \right] \Big|_{\rho=2\pi k_1 A},
\end{aligned}$$

$$\begin{aligned}
(2.56) \quad \mathcal{I}_3^{\sin} &= \frac{(2\pi)^2 k_1}{A^3} \left(\frac{5J_1(\rho)}{\rho^3} + \frac{J_1'(\rho)}{\rho^2} - 30\rho^3 \int_{\rho}^{\infty} \frac{J_1(x)}{x^7} dx \right) \Big|_{\rho=2\pi k_1 A} \\
&= \frac{(2\pi)^2 k_1}{105A^3} \left[-\frac{15J_1(\rho)}{\rho^3} + \frac{15J_1'(\rho)}{\rho^2} + \frac{24J_1(\rho)}{\rho} + 6J_1'(\rho) - 2\rho^2 J_0(\rho) - 2\rho J_1(\rho) \right. \\
&\quad \left. + 2\rho^3 + 2\rho^3 J_1(\rho) - 2\rho^4 J_0(\rho) + \pi\rho^4 J_0(\rho) H_1(\rho) - \pi\rho^4 J_1(\rho) H_0(\rho) \right] \Big|_{\rho=2\pi k_1 A},
\end{aligned}$$

$$\begin{aligned}
(2.57) \quad \mathcal{I}_4^{\sin} &= -\frac{(2\pi)^2 k_1}{A^3} \left(\frac{5J_1(\rho)}{\rho^3} + \frac{J_1'(\rho)}{\rho^2} - \rho^3 \int_{\rho}^{\infty} \frac{J_1(x)}{x^5} dx - 30\rho^3 \int_{\rho}^{\infty} \frac{J_1(x)}{x^7} dx \right) \Big|_{\rho=2\pi k_1 A} \\
&= \frac{(2\pi)^2 k_1}{105A^3} \left[\frac{15J_1(\rho)}{\rho^3} - \frac{15J_1'(\rho)}{\rho^2} + \frac{4J_1(\rho)}{\rho} + J_1'(\rho) - \frac{\rho^2 J_0(\rho)}{3} - \frac{\rho J_1(\rho)}{3} \right. \\
&\quad \left. + \frac{\rho^3}{3} + \frac{\rho^3 J_1(\rho)}{3} - \frac{\rho^4 J_0(\rho)}{3} + \frac{\pi\rho^4}{6} J_0(\rho) H_1(\rho) - \frac{\pi\rho^4}{6} J_1(\rho) H_0(\rho) \right] \Big|_{\rho=2\pi k_1 A}.
\end{aligned}$$

Therefore, (2.45) together with (2.54)–(2.57) furnishes an explicit form of (2.43). In particular, using (A.2), (A.11), we can compute

$$(2.58) \quad \partial_{k_1} \mathcal{I}^{\sin} \Big|_{k_1=0^+} = \frac{1}{2} (2\pi)^2 \left(\frac{a_1^{(1)}}{A} + \frac{4a_4^{(1)} + 3a_5^{(1)} + a_5^{(4)}}{12A^3} \right),$$

$$(2.59) \quad \partial_{k_1}^3 \mathcal{I}^{\sin} \Big|_{k_1=0^+} = \frac{3}{8} (2\pi)^4 A^2 \left(\frac{a_1^{(1)}}{A} - \frac{6a_4^{(1)} + 5a_5^{(1)} + a_5^{(4)}}{6A^3} \right),$$

$$(2.60) \quad \partial_{k_1}^5 \mathcal{I}^{\sin} \Big|_{k_1=0^+} = -\frac{5}{16} (2\pi)^6 A^4 \left(\frac{a_1^{(1)}}{3A} + \frac{8a_4^{(1)} + 7a_5^{(1)} + a_5^{(4)}}{8A^3} \right),$$

$$(2.61) \quad \partial_{k_1}^7 \mathcal{I}^{\sin} \Big|_{k_1=0^+} = \frac{7}{128} (2\pi)^8 A^6 \left(\frac{a_1^{(1)}}{A} + \frac{10a_4^{(1)} + 9a_5^{(1)} + a_5^{(4)}}{6A^3} \right),$$

$$(2.62) \quad \partial_{k_1}^9 \mathcal{I}^{\sin} \Big|_{k_1=0^+} = -\frac{21}{256} (2\pi)^{10} A^8 \left(\frac{3a_1^{(1)}}{7A} + \frac{12a_4^{(1)} + 11a_5^{(1)} + a_5^{(4)}}{20A^3} \right),$$

$$(2.63) \quad \partial_{k_1}^{11} \mathcal{I}^{\sin} \Big|_{k_1=0^+} = \frac{33}{14336} (2\pi)^{12} A^{10} \left(\frac{98a_1^{(1)}}{9A} + \frac{14a_4^{(1)} + 13a_5^{(1)} + a_5^{(4)}}{A^3} \right).$$

Taking into account (2.43), (2.44), we use (2.36)–(2.41) and (2.58)–(2.63) in (2.42) with $n = 0, \dots, 5$, respectively, and thus arrive at the following set of identities:

$$(2.64) \quad 2\pi \iint_{D_A} x_1 B_3(\mathbf{x}, h) d^2x + 2\pi^2 \left(\frac{a_1^{(1)}}{A} + \frac{4a_4^{(1)} + 3a_5^{(1)} + a_5^{(4)}}{12A^3} \right) + \mathcal{R}_1^{\sin} = d_1,$$

$$(2.65) \quad -\frac{4\pi^3}{3} \iint_{D_A} x_1^3 B_3(\mathbf{x}, h) d^2x + \pi^4 A^2 \left(\frac{a_1^{(1)}}{A} - \frac{6a_4^{(1)} + 5a_5^{(1)} + a_5^{(4)}}{6A^3} \right) + \frac{1}{6} \mathcal{R}_3^{\sin} = d_3,$$

$$(2.66) \quad \frac{4\pi^5}{15} \iint_{D_A} x_1^5 B_3(\mathbf{x}, h) d^2x - \frac{\pi^6 A^4}{6} \left(\frac{a_1^{(1)}}{3A} + \frac{8a_4^{(1)} + 7a_5^{(1)} + a_5^{(4)}}{8A^3} \right) + \frac{1}{5!} \mathcal{R}_5^{\sin} = d_5,$$

$$(2.67) \quad -\frac{(2\pi)^7}{5040} \iint_{D_A} x_1^7 B_3(\mathbf{x}, h) d^2x + \frac{(2\pi)^8 A^6}{92160} \left(\frac{a_1^{(1)}}{A} + \frac{10a_4^{(1)} + 9a_5^{(1)} + a_5^{(4)}}{6A^3} \right) + \frac{1}{7!} \mathcal{R}_7^{\sin} = d_7,$$

$$(2.68) \quad \frac{(2\pi)^9}{362880} \iint_{D_A} x_1^9 B_3(\mathbf{x}, h) d^2x - \frac{7(2\pi)^{10} A^8}{10321920} \left(\frac{a_1^{(1)}}{7A} + \frac{12a_4^{(1)} + 11a_5^{(1)} + a_5^{(4)}}{60A^3} \right) + \frac{1}{9!} \mathcal{R}_9^{\sin} = d_9,$$

$$(2.69) \quad -\frac{(2\pi)^{11}}{39916800} \iint_{D_A} x_1^{11} B_3(\mathbf{x}, h) d^2x + \frac{(2\pi)^{12} A^{10}}{17340825600} \left(\frac{98a_1^{(1)}}{9A} + \frac{14a_4^{(1)} + 13a_5^{(1)} + a_5^{(4)}}{A^3} \right) + \frac{1}{11!} \mathcal{R}_{11}^{\sin} = d_{11}.$$

Step 2: Analysis of the remainder terms $\mathcal{R}_{2n+1}^{\sin}$ for $0 \leq n \leq 5$. We shall now show that the remainder terms $\mathcal{R}_{2n+1}^{\sin}$ with $n = 0, \dots, 5$ given by (2.44) (with $k_1 > 0$, as assumed before), can be estimated, for $A \gg 1$, as follows

$$(2.70) \quad \mathcal{R}_{2n+1}^{\sin} = \mathcal{O} \left(\frac{1}{A^{5-2n}} \right), \quad 0 \leq n \leq 5.$$

Proceeding with higher-order terms in the expansions in (2.21)–(2.22), and hence also in (2.24), we notice that we can write, for $N \geq 4$,

$$(2.71) \quad B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h) = \sum_{q=4}^N \sum_{\substack{l_1, l_2 \geq 0, \\ l_1 + l_2 = q}} c_{l_1, l_2} \frac{x_1^{l_1} x_2^{l_2}}{|\mathbf{x}|^{2q+3}} + \mathcal{O} \left(\frac{1}{|\mathbf{x}|^{N+4}} \right) =: L_N(\mathbf{x}) + \mathcal{O} \left(\frac{1}{|\mathbf{x}|^{N+4}} \right),$$

with some constants $c_{l_1, l_2} \in \mathbb{R}$ for $l_1, l_2 \in \mathbb{N}_0$. Consequently, we consider, for $n = 0, \dots, 5$,

$$(2.72) \quad \mathcal{R}_{2n+1}^{\sin} = \partial_{k_1}^{2n+1} \left(\iint_{\mathbb{R}^2 \setminus D_A} \sin(2\pi k_1 x_1) L_N(\mathbf{x}) d^2x \right) \Big|_{k_1=0^+} \\ + \partial_{k_1}^{2n+1} \left(\iint_{\mathbb{R}^2 \setminus D_A} \sin(2\pi k_1 x_1) [B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h) - L_N(\mathbf{x})] d^2x \right) \Big|_{k_1=0^+}.$$

First of all, we deal with the term on the second line. The integrand is regular and, for $N > 9$, it decays at infinity sufficiently fast so that the differential operator $\partial_{k_1}^{2n+1}$ with $0 \leq n \leq 5$ can be passed under the integral sign. We can thus estimate

$$\left| \partial_{k_1}^{2n+1} \left(\iint_{\mathbb{R}^2 \setminus D_A} \sin(2\pi k_1 x_1) [B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h) - L_N(\mathbf{x})] d^2x \right) \Big|_{k_1=0^+} \right| \\ = \left| \left(\iint_{\mathbb{R}^2 \setminus D_A} \frac{(2\pi x_1)^{2n+1}}{|\mathbf{x}|^{N+4}} |\mathbf{x}|^{N+4} [B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h) - L_N(\mathbf{x})] d^2x \right) \right| \\ \leq (2\pi)^{2(n+1)} C_N \int_A^\infty \frac{dr}{r^{N-2n+1}} = \frac{(2\pi)^{2(n+1)} C_N}{(N-2n)} \frac{1}{A^{N-2n}}$$

for some constant $C_N > 0$ such that

$$|\mathbf{x}|^{N+4} |B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h) - L_N(\mathbf{x})| \leq C_N, \quad |\mathbf{x}| \geq A,$$

and such a bound is possible due to the remainder estimate $\mathcal{O}(1/|\mathbf{x}|^{N+4})$ in (2.71). Here, in the third line of the estimates, we used the fact that the integral in r converges for $N > 2n - 1$, $0 \leq n \leq 5$, which is true for $N > 9$. For such N , the obtained estimate of order $\mathcal{O}(1/A^{N-2n})$ is clearly even better than was aimed for (recall (2.70)).

We now fix $N = 10$ and proceed with estimating the term in the first line of (2.72). Upon substitution of (2.71) in (2.44) and use of polar coordinates (with $x_1 = r \cos \theta$, $x_2 = r \sin \theta$, as before), let us observe that, due to Lemma A.4 (namely, identities (A.42)–(A.43)), the only non-vanishing terms stemming from the L_N part are those proportional to

$$(2.73) \quad \int_A^\infty \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) \cos^{2(p-l)+1} \theta \sin^{2l} \theta d\theta \frac{dr}{r^{2p+3}}, \quad 0 \leq l \leq p, \quad p \geq 2.$$

Since we can write

$$\sin^{2l} \theta = (1 - \cos^2 \theta)^l = \sum_{j=0}^l \binom{l}{j} (-1)^{l-j} \cos^{2(l-j)} \theta,$$

with $\binom{l}{j}$ denoting a binomial coefficient, we deduce that, to estimate $\mathcal{R}_{2n+1}^{\sin}$, it suffices only to consider the quantities

$$(2.74) \quad \mathcal{S}_{p,j} := \int_A^\infty \int_0^{2\pi} \sin(2\pi k_1 r \cos \theta) \cos^{2j+1} \theta \, d\theta \frac{dr}{r^{2p+3}}, \quad 0 \leq j \leq p, \quad p \geq 2,$$

and, in particular, their derivatives evaluated at $k_1 = 0$ from the right: $\partial_{k_1}^{2n+1} \mathcal{S}_{p,j} \big|_{k_1=0^+}$, $0 \leq n \leq 5$.

Note that the relation between p (2.73)–(2.74) and q in (2.71) is $q = 2p + 1$, $p \geq 2$. In other words, in the asymptotic expansion of the field B_3 at infinity, not every term contributes to $\mathcal{R}_{2n+1}^{\sin}$, but only the terms of every second order in $1/A$, i.e. $\mathcal{O}(1/A^8)$, $\mathcal{O}(1/A^{10})$ and so on.

Therefore, we can consider $\mathcal{S}_{p,j}$ only for $p \leq \lfloor \frac{N-1}{2} \rfloor = 8$, where $\lfloor x \rfloor$ designates the integer part of x .

To sum up, we need to show that, for all $0 \leq j \leq p$, $2 \leq p \leq 8$ and $0 \leq n \leq 5$, we are able to produce an estimate

$$(2.75) \quad \partial_{k_1}^{2n+1} \mathcal{S}_{p,j} \big|_{k_1=0^+} = \mathcal{O}\left(\frac{1}{A^{5-2n}}\right).$$

For $0 \leq n \leq p$, we have

$$\partial_{k_1}^{2n+1} \mathcal{S}_{p,j} \big|_{k_1=0^+} = (-1)^n (2\pi)^{2n+1} \int_A^\infty \int_0^{2\pi} \cos^{2j+1} \theta \, d\theta \frac{dr}{r^{2(p-n)+2}},$$

and hence

$$\left| \partial_{k_1}^{2n+1} \mathcal{S}_{p,j} \big|_{k_1=0^+} \right| \leq (2\pi)^{2(n+1)} \int_A^\infty \frac{dr}{r^{2(p-n)+2}} = \mathcal{O}\left(\frac{1}{A^{2(p-n)+1}}\right),$$

where the convergence of the last integral is due to $p \geq n$. The obtained estimate is in agreement with (2.75) since $p \geq 2$.

To treat the case $n \geq p + 1$, a more careful estimate is needed. To this end, it is convenient to make use of the integral representation of the Bessel function J_1 given by (A.5) and rewrite (2.74) as

$$\mathcal{S}_{p,j} = (-1)^j 2\pi \int_A^\infty J_1^{(2j)}(2\pi k_1 r) \frac{dr}{r^{2p+3}}, \quad 0 \leq j \leq p, \quad p \geq 2.$$

We then evaluate

$$(2.76) \quad \begin{aligned} \partial_{k_1}^{2n+1} \mathcal{S}_{p,j} &= (-1)^j (2\pi)^{2p+3} \partial_{k_1}^{2(n-p)-1} \int_A^\infty J_1^{(2j+2p+2)}(2\pi k_1 r) \frac{dr}{r} \\ &= (-1)^{j+1} (2\pi)^{2p+4} A \partial_{k_1}^{2(n-p)-2} \left. \frac{J_1^{(2j+2p+2)}(\rho)}{\rho} \right|_{\rho=2\pi k_1 A} \\ &= (-1)^{j+1} (2\pi)^{2(n+1)} A^{2(n-p)-1} \left. \frac{d^{2(n-p)-2}}{d\rho^{2(n-p)-2}} \left(\frac{J_1^{(2j+2p+2)}(\rho)}{\rho} \right) \right|_{\rho=2\pi k_1 A}, \end{aligned}$$

where, in passing the differential operator $\partial_{k_1}^{2p+2}$ under the integral sign, we took into account the asymptotic behavior at infinity of J_1 given by (A.6) and, on the second line, employed the following identity valid for $k_1 > 0$:

$$\begin{aligned} \partial_{k_1} \int_A^\infty J_1^{(2j+2p+2)}(2\pi k_1 r) \frac{dr}{r} &= \partial_{k_1} \int_{2\pi k_1 A}^\infty J_1^{(2j+2p+2)}(\rho) \frac{d\rho}{\rho} \\ &= -2\pi A \left. \frac{J_1^{(2j+2p+2)}(\rho)}{\rho} \right|_{\rho=2\pi k_1 A}. \end{aligned}$$

Now, recalling the analytic character of the function J_1 (see beginning of Appendix) and, more precisely, its series representation given by (A.2), it is clear that every derivative of J_1 of even order is also analytic and vanishes at zero. This implies analyticity of the function $J_1^{(2j+2p+2)}(\rho)/\rho$ and, consequently, a bound on its every derivative at the origin. Therefore, from (2.76), we deduce that

$$\left| \partial_{k_1}^{2n+1} \mathcal{S}_{p,j} \big|_{k_1=0^+} \right| \leq C A^{2(n-p)-1},$$

for some constant $C > 0$, and hence (2.75) follows due to the fact that $p \geq 2$.

Step 3: Asymptotic estimates for the net moment components. Recalling that $d_1 = \pi m_1$ (according to (2.36)) and using (2.70), we obtain from (2.64)

$$(2.77) \quad m_1 = 2 \iint_{D_A} x_1 B_3(\mathbf{x}, h) d^2x + 2\pi \left(\frac{a_1^{(1)}}{A} + \frac{4a_4^{(1)} + 3a_5^{(1)} + a_5^{(4)}}{12A^3} \right) + \mathcal{O}\left(\frac{1}{A^5}\right).$$

Similarly, using (2.70), we rewrite (2.65)–(2.69), respectively, as

$$(2.78) \quad -\frac{1}{6A^2} \iint_{D_A} x_1^3 B_3(\mathbf{x}, h) d^2x + \frac{\pi}{8} \left(\frac{a_1^{(1)}}{A} - \frac{6a_4^{(1)} + 5a_5^{(1)} + a_5^{(4)}}{6A^3} \right) = \frac{d_3}{(2\pi)^3 A^2} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^2}\right),$$

$$(2.79) \quad \frac{1}{120A^4} \iint_{D_A} x_1^5 B_3(\mathbf{x}, h) d^2x - \frac{\pi}{192} \left(\frac{a_1^{(1)}}{3A} + \frac{8a_4^{(1)} + 7a_5^{(1)} + a_5^{(4)}}{8A^3} \right) = \frac{d_5}{(2\pi)^5 A^4} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^4}\right),$$

$$(2.80) \quad -\frac{1}{5040A^6} \iint_{D_A} x_1^7 B_3(\mathbf{x}, h) d^2x + \frac{\pi}{46080} \left(\frac{a_1^{(1)}}{A} + \frac{10a_4^{(1)} + 9a_5^{(1)} + a_5^{(4)}}{6A^3} \right) = \frac{d_7}{(2\pi)^7 A^6} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^5}\right),$$

$$(2.81) \quad \frac{1}{362880A^8} \iint_{D_A} x_1^9 B_3(\mathbf{x}, h) d^2x - \frac{\pi}{737280} \left(\frac{a_1^{(1)}}{7A} + \frac{12a_4^{(1)} + 11a_5^{(1)} + a_5^{(4)}}{60A^3} \right) = \frac{d_9}{(2\pi)^9 A^8} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^5}\right),$$

$$(2.82) \quad -\frac{1}{39916800A^{10}} \iint_{D_A} x_1^{11} B_3(\mathbf{x}, h) d^2x + \frac{\pi}{8670412800} \left(\frac{98a_1^{(1)}}{9A} + \frac{14a_4^{(1)} + 13a_5^{(1)} + a_5^{(4)}}{A^3} \right) = \frac{d_{11}}{(2\pi)^{11} A^{10}} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^5}\right).$$

While the first-order estimate for m_1 given in (1.5) follows immediately from rigorously justified (2.77), the higher-order estimates require more work. Namely, we wish to combine (2.78)–(2.82) in order to eliminate in (2.77) the terms with

$$(2.83) \quad \tilde{a}_1^{(1)} := \frac{a_1^{(1)}}{A}, \quad \tilde{a}_4^{(1)} := \frac{a_4^{(1)}}{A^3}, \quad \tilde{a}_5^{(1)} := \frac{a_5^{(1)}}{A^3}, \quad \tilde{a}_5^{(4)} := \frac{a_5^{(4)}}{A^3},$$

and, at the same time, would not commit a larger error (in order of A) than that of the eliminated term.

Expressing $a_1^{(1)}/A$ in terms of $\mathcal{O}(1/A^2)$ quantities from (2.78) and inserting it into (2.77), we deduce the second-order estimate for m_1 given by (1.6). We note, however, that, for third or higher order estimates, identity (2.78) is not useful due to the fact that its right-hand side has an unknown quantity d_3 appearing of order $\mathcal{O}(1/A^2)$ which will block any further effort to increase the accuracy of estimates.

Derivation of the third-order estimate given by (1.8) is analogous to the previous one with the only difference that $a_1^{(1)}/A$ is expressed (now in terms of $\mathcal{O}(1/A^3)$ quantities) from (2.79) rather than from (2.78).

To proceed with derivation of estimates (1.10) and (1.12), it is convenient first to rewrite (2.79)–(2.82), respectively, as

$$(2.84) \quad 8\tilde{a}_4^{(1)} + 7\tilde{a}_5^{(1)} + \tilde{a}_5^{(4)} = \frac{64}{5\pi A^4} \iint_{D_A} x_1^5 B_3(\mathbf{x}, h) d^2x - \frac{8}{3}\tilde{a}_1^{(1)} + \mathcal{O}\left(\frac{1}{A^4}\right) =: \mathcal{T}_5,$$

$$(2.85) \quad 10\tilde{a}_4^{(1)} + 9\tilde{a}_5^{(1)} + \tilde{a}_5^{(4)} = \frac{384}{7\pi A^6} \iint_{D_A} x_1^7 B_3(\mathbf{x}, h) d^2x - 6\tilde{a}_1^{(1)} + \mathcal{O}\left(\frac{1}{A^5}\right) =: \mathcal{T}_7,$$

$$(2.86) \quad 12\tilde{a}_4^{(1)} + 11\tilde{a}_5^{(1)} + \tilde{a}_5^{(4)} = \frac{2560}{21\pi A^8} \iint_{D_A} x_1^9 B_3(\mathbf{x}, h) d^2x - \frac{60}{7}\tilde{a}_1^{(1)} + \mathcal{O}\left(\frac{1}{A^5}\right) =: \mathcal{T}_9,$$

$$(2.87) \quad 14\tilde{a}_4^{(1)} + 13\tilde{a}_5^{(1)} + \tilde{a}_5^{(4)} = \frac{64512}{297\pi A^{10}} \iint_{D_A} x_1^{11} B_3(\mathbf{x}, h) d^2x - \frac{98}{9}\tilde{a}_1^{(1)} + \mathcal{O}\left(\frac{1}{A^5}\right) =: \mathcal{T}_{11}.$$

It is easy to recognise a linear dependence of the left-hand sides of (2.84)–(2.87). Namely, we can write

$$(2.88) \quad \frac{1}{2}(\mathcal{T}_5 + \mathcal{T}_9) = \mathcal{T}_7, \quad \frac{1}{2}(\mathcal{T}_7 + \mathcal{T}_{11}) = \mathcal{T}_9.$$

To obtain the fourth-order estimate for m_1 given in (1.10), we shall use (2.84)–(2.86).

We start by using the first equation of (2.88) (together with definitions (2.84)–(2.86)) to express $\tilde{a}_1^{(1)} = a_1^{(1)}/A$ up to order $\mathcal{O}(1/A^4)$, namely,

$$(2.89) \quad \frac{a_1^{(1)}}{A} = \tilde{a}_1^{(1)} = -\frac{4}{\pi} \iint_{D_A} \left[\frac{21}{5} \left(\frac{x_1}{A}\right)^4 - 36 \left(\frac{x_1}{A}\right)^6 + 40 \left(\frac{x_1}{A}\right)^8 \right] x_1 B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A^4}\right).$$

Second, we observe that the quantity $4a_4^{(1)} + 3a_5^{(1)} + a_5^{(4)}$ appearing in (2.77) is related to (2.84)–(2.86) as follows:

$$(2.90) \quad \begin{aligned} \frac{4a_4^{(1)} + 3a_5^{(1)} + a_5^{(4)}}{A^3} &= 4\tilde{a}_4^{(1)} + 3\tilde{a}_5^{(1)} + \tilde{a}_5^{(4)} = 4(\mathcal{T}_5 - \mathcal{T}_7) + \mathcal{T}_9 \\ &= \frac{256}{\pi} \iint_{D_A} \left[\frac{1}{5} \left(\frac{x_1}{A}\right)^4 - \frac{6}{7} \left(\frac{x_1}{A}\right)^6 + \frac{10}{21} \left(\frac{x_1}{A}\right)^8 \right] x_1 B_3(\mathbf{x}, h) d^2x \\ &\quad + \frac{100}{21}\tilde{a}_1^{(1)} + \mathcal{O}\left(\frac{1}{A^4}\right). \end{aligned}$$

Finally, substitution of (2.90) in (2.77) followed by the use of (2.89) furnishes (1.10).

To arrive at the fifth-order estimate given by (1.12), we shall use identities (2.85)–(2.87). The second equation of (2.88) gives

$$(2.91) \quad \frac{a_1^{(1)}}{A} = \tilde{a}_1^{(1)} = \frac{24}{\pi} \iint_{D_A} \left[-9 \left(\frac{x_1}{A}\right)^6 + 40 \left(\frac{x_1}{A}\right)^8 - \frac{392}{11} \left(\frac{x_1}{A}\right)^{10} \right] x_1 B_3(\mathbf{x}, h) d^2x + \mathcal{O}\left(\frac{1}{A^5}\right).$$

We also have an analog of (2.83), namely,

$$(2.92) \quad \begin{aligned} \frac{4a_4^{(1)} + 3a_5^{(1)} + a_5^{(4)}}{A^3} &= 4\tilde{a}_4^{(1)} + 3\tilde{a}_5^{(1)} + \tilde{a}_5^{(4)} = 5(\mathcal{T}_7 - \mathcal{T}_9) + \mathcal{T}_{11} \\ &= \frac{128}{\pi} \iint_{D_A} \left[\frac{15}{7} \left(\frac{x_1}{A}\right)^6 - \frac{100}{21} \left(\frac{x_1}{A}\right)^8 + \frac{56}{33} \left(\frac{x_1}{A}\right)^{10} \right] x_1 B_3(\mathbf{x}, h) d^2x \\ &\quad + \frac{124}{63}\tilde{a}_1^{(1)} + \mathcal{O}\left(\frac{1}{A^5}\right). \end{aligned}$$

Therefore, inserting of (2.92) in (2.77) followed by the use of (2.91) gives (1.12).

2.3.2. Normal component of the net moment. Similarly to the case of tangential net moment components, we expand the integrand in (2.7) in power series in a positive neighborhood of $k_1 = 0$, but, in contrast to that previous situation, the attention will now be on the real part. This yields

$$(2.93) \quad \left[\operatorname{Re} \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = 0,$$

$$(2.94) \quad \begin{aligned} \frac{1}{2} \partial_{k_1}^2 \left[\operatorname{Re} \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} &= -2\pi^2 [\langle x_1 M_1 \rangle + \langle (h - x_3) M_3 \rangle] \\ &=: d_2, \end{aligned}$$

$$(2.95) \quad \begin{aligned} \frac{1}{24} \partial_{k_1}^4 \left[\operatorname{Re} \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} &= -\frac{4\pi^4}{3} \left[3 \langle (h - x_3)^2 x_1 M_1 \rangle - \langle x_1^3 M_1 \rangle + \langle (h - x_3)^3 M_3 \rangle - 3 \langle (h - x_3) x_1^2 M_3 \rangle \right] \\ &=: d_4, \end{aligned}$$

$$(2.96) \quad \frac{1}{6!} \partial_{k_1}^6 \left[\text{Re } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = -\frac{4\pi^6}{15} \left[5 \langle (h-x_3)^4 x_1 M_1 \rangle - 10 \langle (h-x_3)^2 x_1^3 M_1 \rangle + \langle x_1^5 M_1 \rangle \right. \\ \left. + \langle (h-x_3)^5 M_3 \rangle - 10 \langle (h-x_3)^3 x_1^2 M_3 \rangle + 5 \langle (h-x_3) x_1^4 M_3 \rangle \right] \\ =: d_6,$$

$$(2.97) \quad \frac{1}{8!} \partial_{k_1}^8 \left[\text{Re } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = -\frac{8\pi^8}{315} \left[7 \langle (h-x_3)^6 x_1 M_1 \rangle - 35 \langle (h-x_3)^4 x_1^3 M_1 \rangle + 21 \langle (h-x_3)^2 x_1^5 M_1 \rangle \right. \\ \left. - \langle x_1^7 M_1 \rangle + \langle (h-x_3)^7 M_3 \rangle - 21 \langle (h-x_3)^5 x_1^2 M_3 \rangle + 35 \langle (h-x_3)^3 x_1^4 M_3 \rangle \right. \\ \left. - 7 \langle (h-x_3) x_1^6 M_3 \rangle \right] \\ =: d_8,$$

$$(2.98) \quad \frac{1}{10!} \partial_{k_1}^{10} \left[\text{Re } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = -\frac{4\pi^{10}}{2385} \left[9 \langle (h-x_3)^8 x_1 M_1 \rangle - 84 \langle (h-x_3)^6 x_1^3 M_1 \rangle + 126 \langle (h-x_3)^4 x_1^5 M_1 \rangle \right. \\ \left. - 36 \langle (h-x_3)^2 x_1^7 M_1 \rangle + \langle x_1^9 M_1 \rangle + \langle (h-x_3)^9 M_3 \rangle - 36 \langle (h-x_3)^7 x_1^2 M_3 \rangle \right. \\ \left. + 126 \langle (h-x_3)^5 x_1^4 M_3 \rangle - 84 \langle (h-x_3)^3 x_1^6 M_3 \rangle + 9 \langle (h-x_3) x_1^8 M_3 \rangle \right] \\ =: d_{10},$$

On the other hand, (2.10) implies that, for $n \in \mathbb{N}_0$,

$$(2.99) \quad \partial_{k_1}^{2n} \left[\text{Re } \hat{B}_3(k_1, 0, h) \right] \Big|_{k_1=0^+} = (-1)^n (2\pi)^{2n} \iint_{D_A} x_1^{2n} B_3(\mathbf{x}, h) d^2x \\ + \partial_{k_1}^{2n} \left(\iint_{\mathbb{R}^2 \setminus D_A} \cos(2\pi k_1 x_1) B_3^{\text{asympt}}(\mathbf{x}, h) d^2x \right) \Big|_{k_1=0^+} \\ + \partial_{k_1}^{2n} \left(\iint_{\mathbb{R}^2 \setminus D_A} \cos(2\pi k_1 x_1) [B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h)] d^2x \right) \Big|_{k_1=0^+}.$$

As before, we continue in 3 steps. First, we evaluate explicitly

$$(2.100) \quad \mathcal{I}^{\cos} \equiv \mathcal{I}^{\cos}(k_1, A) := \iint_{\mathbb{R}^2 \setminus D_A} \cos(2\pi k_1 x_1) B_3^{\text{asympt}}(\mathbf{x}, h) d^2x,$$

that allow us to obtain from (2.99), a set of useful identities involving remainder terms. Second, we estimate the remainder terms, namely,

$$(2.101) \quad \mathcal{R}_{2n}^{\cos} \equiv \mathcal{R}_{2n}^{\cos}(A) := \partial_{k_1}^{2n} \left(\iint_{\mathbb{R}^2 \setminus D_A} \cos(2\pi k_1 x_1) [B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h)] d^2x \right) \Big|_{k_1=0^+}, \quad n \in \mathbb{N}_0,$$

for large A . At last, we rigorously derive asymptotic formulas (1.7), (1.9), (1.11) from the obtained set of identities.

Step 1: Derivation of the set of identities. Using (2.25) and passing to the polar coordinates using $x = r \cos \theta$, $y = r \sin \theta$, $d^2x = r dr d\theta$, we obtain from (2.100)

$$(2.102) \quad \mathcal{I}^{\cos} = a_0 \mathcal{I}_1^{\cos} + a_2 \mathcal{I}_2^{\cos} + a_3^{(1)} \mathcal{I}_3^{\cos} + a_3^{(2)} \mathcal{I}_4^{\cos},$$

where

$$(2.103) \quad \mathcal{I}_1^{\cos} := \int_A^\infty \int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) d\theta \frac{dr}{r^2},$$

$$(2.104) \quad \mathcal{I}_2^{\cos} := \int_A^\infty \int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) \theta d\theta \frac{dr}{r^4},$$

$$(2.105) \quad \mathcal{I}_3^{\cos} := \int_A^\infty \int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) \cos^2 \theta d\theta \frac{dr}{r^4},$$

$$(2.106) \quad \mathcal{I}_4^{\cos} := \int_A^\infty \int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) \sin^2 \theta d\theta \frac{dr}{r^4}.$$

Here, we used results (A.40)–(A.41) of Lemma A.4 multiple times to deduce vanishing of the integrals associated with the terms which involve $a_1^{(1)}$, $a_1^{(2)}$, $a_3^{(3)}$, $a_4^{(1)}$, $a_4^{(2)}$, $a_5^{(j)}$, $j \in \{1, 2, 3, 4\}$.

Using the integral representation of Bessel functions (due to (A.4)), integration by parts, the asymptotics of J_1 given by (A.6), and the relation $J_0'(x) = -J_1(x)$ (see (A.10)), we can write, for $k_1 > 0$,

$$(2.107) \quad \begin{aligned} \mathcal{I}_1^{\cos} &= 2\pi \int_A^\infty J_0(2\pi k_1 r) \frac{dr}{r^2} = (2\pi)^2 k_1 \int_{2\pi k_1 A}^\infty \frac{J_0(x)}{x^2} dx \\ &= (2\pi)^2 k_1 \left(\frac{J_0(\rho)}{\rho} - \int_\rho^\infty \frac{J_1(x)}{x} dx \right) \Big|_{\rho=2\pi k_1 A}, \end{aligned}$$

$$(2.108) \quad \begin{aligned} \mathcal{I}_2^{\cos} &= 2\pi \int_A^\infty J_0(2\pi k_1 r) \frac{dr}{r^4} = (2\pi)^4 k_1^3 \int_{2\pi k_1 A}^\infty \frac{J_0(x)}{x^4} dx \\ &= \frac{(2\pi)^4}{3} k_1^3 \left(\frac{J_0(\rho)}{\rho^3} - \int_\rho^\infty \frac{J_1(x)}{x^3} dx \right) \Big|_{\rho=2\pi k_1 A}, \end{aligned}$$

$$(2.109) \quad \begin{aligned} \mathcal{I}_3^{\cos} &= -2\pi \int_A^\infty J_0''(2\pi k_1 r) \frac{dr}{r^4} = (2\pi)^4 k_1^3 \int_{2\pi k_1 A}^\infty \frac{J_1'(x)}{x^4} dx \\ &= (2\pi)^4 k_1^3 \left(\frac{J_1(\rho)}{\rho^4} + 4 \int_\rho^\infty \frac{J_1(x)}{x^5} dx \right) \Big|_{\rho=2\pi k_1 A}, \end{aligned}$$

$$(2.110) \quad \begin{aligned} \mathcal{I}_4^{\cos} &= 2\pi \int_A^\infty [J_0(2\pi k_1 r) + J_0''(2\pi k_1 r)] \frac{dr}{r^4} = (2\pi)^4 k_1^3 \int_{2\pi k_1 A}^\infty \frac{J_0(x) + J_0''(x)}{x^4} dx, \\ &= (2\pi)^4 k_1^3 \left(\frac{J_0(\rho)}{3\rho^3} - \frac{J_1(\rho)}{\rho^4} - \frac{1}{3} \int_\rho^\infty \frac{J_1(x)}{x^3} dx - 4 \int_\rho^\infty \frac{J_1(x)}{x^5} dx \right) \Big|_{\rho=2\pi k_1 A}. \end{aligned}$$

Inserting here the results of Lemmas (A.2)–(A.3), namely, (A.22), (A.34)–(A.35), we arrive at

$$(2.111) \quad \begin{aligned} \mathcal{I}_1^{\cos} &= \frac{2\pi}{A} \left[2J_0(\rho) + \frac{J_1(\rho)}{\rho} - J_1'(\rho) - 2J_1(\rho) - \rho - \rho J_1(\rho) \right. \\ &\quad \left. + \rho^2 J_0(\rho) - \frac{\pi\rho^2}{2} J_0(\rho) H_1(\rho) + \frac{\pi\rho^2}{2} J_1(\rho) H_0(\rho) \right] \Big|_{\rho=2\pi k_1 A}, \end{aligned}$$

$$(2.112) \quad \begin{aligned} \mathcal{I}_2^{\cos} &= \frac{2\pi}{3A^3} \left[J_0(\rho) - \frac{\rho^2 J_0(\rho)}{3} - \frac{\rho J_1(\rho)}{3} + \frac{\rho^3}{3} + \frac{\rho^3}{3} J_1(\rho) \right. \\ &\quad \left. - \frac{\rho^4 J_0(\rho)}{3} + \frac{\pi\rho^4}{6} J_0(\rho) H_1(\rho) - \frac{\pi\rho^4}{6} J_1(\rho) H_0(\rho) \right] \Big|_{\rho=2\pi k_1 A}, \end{aligned}$$

$$(2.113) \quad \begin{aligned} \mathcal{I}_3^{\cos} &= \frac{2\pi}{15A^3} \left[-\frac{J_0(\rho)}{3} + 4J_1'(\rho) - \frac{4\rho^2 J_0(\rho)}{3} - \frac{4\rho J_1(\rho)}{3} + \frac{4\rho^3}{3} + \frac{4\rho^3 J_1(\rho)}{3} \right. \\ &\quad \left. - \frac{4\rho^4 J_0(\rho)}{3} + \frac{2\pi\rho^4}{3} J_0(\rho) H_1(\rho) - \frac{2\pi\rho^4}{3} J_1(\rho) H_0(\rho) \right] \Big|_{\rho=2\pi k_1 A}, \end{aligned}$$

$$(2.114) \quad \begin{aligned} \mathcal{I}_4^{\cos} &= \frac{2\pi}{3A^3} \left[J_0(\rho) + \frac{J_1(\rho)}{15\rho} - \frac{4J_1'(\rho)}{5} - \frac{\rho^2 J_0(\rho)}{15} - \frac{\rho J_1(\rho)}{15} + \frac{\rho^3}{15} + \frac{\rho^3 J_1(\rho)}{15} \right. \\ &\quad \left. - \frac{\rho^4 J_0(\rho)}{15} + \frac{\pi\rho^4}{30} J_0(\rho) H_1(\rho) - \frac{\pi\rho^4}{30} J_1(\rho) H_0(\rho) \right] \Big|_{\rho=2\pi k_1 A}. \end{aligned}$$

Substitution of (2.111)–(2.114) into (2.102), we employ (A.2) and (A.11) to compute

$$(2.115) \quad \mathcal{I}^{\cos}|_{k_1=0+} = 2\pi \left(\frac{a_0}{A} + \frac{30a_2 + 11a_3^{(1)} + 19a_3^{(2)}}{90A^3} \right),$$

$$(2.116) \quad \partial_{k_1}^2 \mathcal{I}^{\cos} \Big|_{k_1=0^+} = \frac{1}{4} (2\pi)^3 A^2 \left(\frac{a_0}{A} - \frac{180a_2 + 131a_3^{(1)} + 49a_3^{(2)}}{180A^3} \right),$$

$$(2.117) \quad \partial_{k_1}^4 \mathcal{I}^{\cos} \Big|_{k_1=0^+} = -\frac{1}{8} (2\pi)^5 A^4 \left(\frac{a_0}{A} + \frac{270a_2 + 229a_3^{(1)} + 41a_3^{(2)}}{90A^3} \right),$$

$$(2.118) \quad \partial_{k_1}^6 \mathcal{I}^{\cos} \Big|_{k_1=0^+} = \frac{1}{16} (2\pi)^7 A^6 \left(\frac{a_0}{A} + \frac{120a_2 + 109a_3^{(1)} + 11a_3^{(2)}}{72A^3} \right),$$

$$(2.119) \quad \partial_{k_1}^8 \mathcal{I}^{\cos} \Big|_{k_1=0^+} = -3136 (2\pi)^9 A^8 \left(\frac{a_0}{7A} + \frac{18a_2 + 17a_3^{(1)} + a_3^{(2)}}{90A^3} \right),$$

$$(2.120) \quad \partial_{k_1}^{10} \mathcal{I}^{\cos} \Big|_{k_1=0^+} = \frac{7}{256} (2\pi)^{11} A^{10} \left(\frac{a_0}{A} + \frac{540a_2 + 523a_3^{(1)} + 17a_3^{(2)}}{420A^3} \right).$$

Taking into account (2.100), (2.101), we use (2.93)–(2.98) and (2.115)–(2.120) in (2.99) with $n = 0, \dots, 5$, respectively, and thus arrive at the following set of identities:

$$(2.121) \quad \iint_{D_A} B_3(\mathbf{x}, h) d^2x + 2\pi \left(\frac{a_0}{A} + \frac{30a_2 + 11a_3^{(1)} + 19a_3^{(2)}}{90A^3} \right) + \mathcal{R}_0^{\cos} = 0,$$

$$(2.122) \quad -\frac{(2\pi)^2}{2} \iint_{D_A} x_1^2 B_3(\mathbf{x}, h) d^2x + \frac{(2\pi)^3}{4} A^2 \left(\frac{a_0}{A} - \frac{180a_2 + 131a_3^{(1)} + 49a_3^{(2)}}{180A^3} \right) + \frac{1}{2} \mathcal{R}_2^{\cos} = d_2,$$

$$(2.123) \quad \frac{(2\pi)^4}{24} \iint_{D_A} x_1^4 B_3(\mathbf{x}, h) d^2x - \frac{1}{192} (2\pi)^5 A^4 \left(\frac{a_0}{A} + \frac{270a_2 + 229a_3^{(1)} + 41a_3^{(2)}}{90A^3} \right) + \frac{1}{4!} \mathcal{R}_4^{\cos} = d_4,$$

$$(2.124) \quad -\frac{(2\pi)^6}{720} \iint_{D_A} x_1^6 B_3(\mathbf{x}, h) d^2x + \frac{1}{11520} (2\pi)^7 A^6 \left(\frac{a_0}{A} + \frac{120a_2 + 109a_3^{(1)} + 11a_3^{(2)}}{72A^3} \right) + \frac{1}{6!} \mathcal{R}_6^{\cos} = d_6,$$

$$(2.125) \quad \frac{(2\pi)^8}{40320} \iint_{D_A} x_1^8 B_3(\mathbf{x}, h) d^2x - \frac{7}{90} (2\pi)^9 A^8 \left(\frac{a_0}{7A} + \frac{18a_2 + 17a_3^{(1)} + a_3^{(2)}}{90A^3} \right) + \frac{1}{8!} \mathcal{R}_8^{\cos} = d_8,$$

$$(2.126) \quad -\frac{(2\pi)^{10}}{3628800} \iint_{D_A} x_1^{10} B_3(\mathbf{x}, h) d^2x + \frac{1}{132710400} (2\pi)^{11} A^{10} \left(\frac{a_0}{A} + \frac{540a_2 + 523a_3^{(1)} + 17a_3^{(2)}}{420A^3} \right) + \frac{1}{10!} \mathcal{R}_{10}^{\cos} = d_{10}.$$

Step 2: Analysis of the remainder terms \mathcal{R}_{2n}^{\cos} for $0 \leq n \leq 5$. We shall show that

$$(2.127) \quad \mathcal{R}_{2n}^{\cos} = \mathcal{O} \left(\frac{1}{A^{5-2n}} \right), \quad 0 \leq n \leq 5.$$

The reasoning will be identical to that of Step 2 of Subsection 2.3.1, therefore, we omit repetition of some details.

Using previously introduced notation L_N (see (2.71)), we can write

$$(2.128) \quad \begin{aligned} \mathcal{R}_{2n}^{\cos} &= \partial_{k_1}^{2n} \left(\iint_{\mathbb{R}^2 \setminus D_A} \cos(2\pi k_1 x_1) L_N(\mathbf{x}) d^2x \right) \Big|_{k_1=0^+} \\ &\quad + \partial_{k_1}^{2n} \left(\iint_{\mathbb{R}^2 \setminus D_A} \cos(2\pi k_1 x_1) [B_3(\mathbf{x}, h) - B_3^{\text{asymp}}(\mathbf{x}, h) - L_N(\mathbf{x})] d^2x \right) \Big|_{k_1=0^+}. \end{aligned}$$

Assuming $N > 9$, we estimate

$$\left| \partial_{k_1}^{2n} \left(\iint_{\mathbb{R}^2 \setminus D_A} \cos(2\pi k_1 x_1) [B_3(\mathbf{x}, h) - B_3^{\text{asymp}}(\mathbf{x}, h) - L_N(\mathbf{x})] d^2x \right) \Big|_{k_1=0^+} \right|$$

$$\begin{aligned}
&= \left| \left(\iint_{\mathbb{R}^2 \setminus D_A} \frac{(2\pi x_1)^{2n}}{|\mathbf{x}|^{N+4}} |\mathbf{x}|^{N+4} [B_3(\mathbf{x}, h) - B_3^{\text{asympt}}(\mathbf{x}, h) - L_N(\mathbf{x})] d^2x \right) \right| \\
&\leq (2\pi)^{2n+1} \tilde{C}_N \int_A^\infty \frac{dr}{r^{N-2n+2}} = \frac{(2\pi)^{2(n+1)} \tilde{C}_N}{(N-2n+1) A^{N-2n+1}},
\end{aligned}$$

for some constant $\tilde{C}_N > 0$, and note that this complies with (2.127).

We now fix $N = 10$ and proceed with estimating the term on the first line of (2.127). Identities (A.40)–(A.41) of Lemma A.4 entail that the only non-vanishing terms are those of the form

$$(2.129) \quad \int_A^\infty \int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) \cos^{2(p-l)} \theta \cos^{2l} \theta d\theta \frac{dr}{r^{2(p+1)}}, \quad 0 \leq l \leq p, \quad p \geq 2,$$

and hence it is sufficient to only deal with the quantities

$$(2.130) \quad \mathcal{C}_{p,j} := \int_A^\infty \int_0^{2\pi} \cos(2\pi k_1 r \cos \theta) \cos^{2j} \theta d\theta \frac{dr}{r^{2(p+1)}}, \quad 0 \leq j \leq p, \quad p \geq 2.$$

As before, p in (2.129)–(2.130) is related to q from (2.71) as $q = 2p + 1$, $p \geq 2$: only the terms $\mathcal{O}(1/A^7)$, $\mathcal{O}(1/A^9)$, ..., in (2.71) contribute to $\mathcal{R}_{2n}^{\text{cos}}$. Hence, we consider $\mathcal{C}_{p,j}$ only for $p \leq \lfloor \frac{N-1}{2} \rfloor = 8$.

We are intending to show that the following estimate holds for all $0 \leq n \leq 5$, $0 \leq j \leq p$, $2 \leq p \leq 8$:

$$(2.131) \quad \partial_{k_1}^{2n} \mathcal{C}_{p,j} \Big|_{k_1=0^+} = \mathcal{O}\left(\frac{1}{A^{5-2n}}\right).$$

For $0 \leq n \leq p$, we have

$$\partial_{k_1}^{2n} \mathcal{C}_{p,j} \Big|_{k_1=0^+} = (-1)^n (2\pi)^{2n} \int_A^\infty \int_0^{2\pi} \cos^{2j} \theta d\theta \frac{dr}{r^{2(p-n)+2}},$$

and hence

$$\left| \partial_{k_1}^{2n+1} \mathcal{C}_{p,j} \Big|_{k_1=0^+} \right| \leq (2\pi)^{2n+1} \int_A^\infty \frac{dr}{r^{2(p-n)+2}} = \mathcal{O}\left(\frac{1}{A^{2(p-n)+1}}\right),$$

where the convergence of the last integral is due to $p \geq n$. The obtained estimate satisfies (2.131) due to $p \geq 2$.

For $n \geq p + 1$, we use (A.4) to rewrite (2.130) as

$$\mathcal{C}_{p,j} = (-1)^j 2\pi \int_A^\infty J_0^{(2j)}(2\pi k_1 r) \frac{dr}{r^{2(p+1)}}, \quad 0 \leq j \leq p, \quad p \geq 2.$$

We then evaluate

$$\begin{aligned}
(2.132) \quad \partial_{k_1}^{2n} \mathcal{C}_{p,j} &= (-1)^j (2\pi)^{2(p+1)} \partial_{k_1}^{2(n-p)-1} \int_A^\infty J_0^{(2j+2p+1)}(2\pi k_1 r) \frac{dr}{r} \\
&= (-1)^{j+1} (2\pi)^{2p+3} A \partial_{k_1}^{2(n-p)-2} \frac{J_0^{(2j+2p+1)}(\rho)}{\rho} \Big|_{\rho=2\pi k_1 A} \\
&= (-1)^{j+1} (2\pi)^{2n+1} A^{2(n-p)-1} \frac{d^{2(n-p)-2}}{d\rho^{2(n-p)-2}} \left(\frac{J_0^{(2j+2p+1)}(\rho)}{\rho} \right) \Big|_{\rho=2\pi k_1 A},
\end{aligned}$$

where the following identity, for $k_1 > 0$, was used:

$$\begin{aligned}
\partial_{k_1} \int_A^\infty J_0^{(2j+2p+1)}(2\pi k_1 r) \frac{dr}{r} &= \partial_{k_1} \int_{2\pi k_1 A}^\infty J_0^{(2j+2p+1)}(\rho) \frac{d\rho}{\rho} \\
&= -2\pi A \frac{J_0^{(2j+2p+1)}(\rho)}{\rho} \Big|_{\rho=2\pi k_1 A}.
\end{aligned}$$

Due to analyticity J_0 (see beginning of Appendix), we see from (A.2) that its every derivative of odd order is also analytic and vanishes at zero. Hence, $J_0^{(2j+2p+1)}(\rho)/\rho$ is analytic as well, and, in particular, has bounded derivatives at the origin. Therefore, (2.132) entails that

$$\left| \partial_{k_1}^{2n} \mathcal{C}_{p,j} \Big|_{k_1=0^+} \right| \leq \tilde{C} A^{2(n-p)-1},$$

for some constant $\tilde{C} > 0$, and hence (2.131) follows due to the fact that $p \geq 2$.

Step 3: Asymptotic estimates for the net moment component. Since, according to (2.26), we have $a_0 = -\frac{m_3}{4\pi}$, and hence using (2.127), it follows from (2.121) that

$$(2.133) \quad m_3 = 2 \iint_{D_A} B_3(\mathbf{x}, h) d^2x + \frac{\pi}{45A^2} \left(30a_2 + 11a_3^{(1)} + 19a_3^{(2)} \right) + \mathcal{O}\left(\frac{1}{A^4}\right).$$

Employing (2.70), we rewrite (2.65)–(2.69), respectively, as

$$(2.134) \quad -\frac{1}{2A^2} \iint_{D_A} x_1^2 B_3(\mathbf{x}, h) d^2x - \frac{\pi}{8} \left(\frac{m_3}{\pi A} + \frac{180a_2 + 131a_3^{(1)} + 49a_3^{(2)}}{45A^3} \right) = \frac{d_2}{(2\pi A)^2} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^2}\right),$$

$$(2.135) \quad \frac{1}{24A^4} \iint_{D_A} x_1^4 B_3(\mathbf{x}, h) d^2x + \frac{\pi}{192} \left(\frac{m_3}{2\pi A} - \frac{270a_2 + 229a_3^{(1)} + 41a_3^{(2)}}{45A^3} \right) = \frac{d_4}{(2\pi A)^4} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^4}\right),$$

$$(2.136) \quad -\frac{1}{720A^6} \iint_{D_A} x_1^6 B_3(\mathbf{x}, h) d^2x - \frac{\pi}{2} \left(\frac{m_3}{\pi A} - \frac{120a_2 + 109a_3^{(1)} + 11a_3^{(2)}}{18A^3} \right) = \frac{d_6}{(2\pi A)^6} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^5}\right),$$

$$(2.137) \quad \frac{1}{40320A^8} \iint_{D_A} x_1^8 B_3(\mathbf{x}, h) d^2x + \frac{\pi}{147456} \left(\frac{m_3}{14\pi A} - \frac{18a_2 + 17a_3^{(1)} + a_3^{(2)}}{45A^3} \right) = \frac{d_8}{(2\pi A)^8} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^5}\right),$$

$$(2.138) \quad -\frac{1}{3628800A^{10}} \iint_{D_A} x_1^{10} B_3(\mathbf{x}, h) d^2x - \frac{\pi}{66355200} \left(\frac{m_3}{\pi A} - \frac{540a_2 + 523a_3^{(1)} + 17a_3^{(2)}}{105A^3} \right) = \frac{d_{10}}{(2\pi A)^{10}} + \mathcal{O}\left(\frac{1}{A^5}\right) \\ = \mathcal{O}\left(\frac{1}{A^5}\right).$$

We see that (2.133) already provides the second-order estimate of m_3 given in (1.7). The higher-order estimates given by (1.9) and (1.11) can be obtained by combining (2.135)–(2.137). Note that elimination of m_3 using either (2.134) would incur a rather large error $\mathcal{O}(1/A)$, reducing the final estimate to the first order. Similarly, any use of (2.138) would result in an estimate of order $\mathcal{O}(1/A^4)$, but such an estimate could already be deduced using the other listed above relations.

For the sake of simplification, let us set

$$(2.139) \quad \tilde{a}_2 := \frac{a_2}{A^3}, \quad \tilde{a}_3^{(1)} := \frac{a_3^{(1)}}{A^3}, \quad \tilde{a}_3^{(2)} := \frac{a_3^{(2)}}{A^3},$$

and rewrite (2.133) and (2.135)–(2.137), respectively, as

$$(2.140) \quad 30\tilde{a}_2 + 11\tilde{a}_3^{(1)} + 19\tilde{a}_3^{(2)} = -\frac{45}{\pi} \iint_{D_A} B_3(\mathbf{x}, h) d^2x + \frac{45m_3}{2\pi A} + \mathcal{O}\left(\frac{1}{A^5}\right) =: \mathcal{T}_0,$$

$$(2.141) \quad 180\tilde{a}_2 + 131\tilde{a}_3^{(1)} + 49\tilde{a}_3^{(2)} = -\frac{180}{\pi A^2} \iint_{D_A} x_1^2 B_3(\mathbf{x}, h) d^2x - \frac{45m_3}{\pi A} + \mathcal{O}\left(\frac{1}{A^2}\right) =: \mathcal{T}_2,$$

$$(2.142) \quad 270\tilde{a}_2 + 229\tilde{a}_3^{(1)} + 41\tilde{a}_3^{(2)} = \frac{360}{\pi A^4} \iint_{D_A} x_1^4 B_3(\mathbf{x}, h) d^2x + \frac{45m_3}{2\pi A} + \mathcal{O}\left(\frac{1}{A^4}\right) =: \mathcal{T}_4,$$

$$(2.143) \quad 120\tilde{a}_2 + 109\tilde{a}_3^{(1)} + 11\tilde{a}_3^{(2)} = \frac{576}{\pi A^6} \iint_{D_A} x_1^6 B_3(\mathbf{x}, h) d^2x + \frac{18m_3}{\pi A} + \mathcal{O}\left(\frac{1}{A^5}\right) =: \mathcal{T}_6,$$

$$(2.144) \quad 18\tilde{a}_2 + 17\tilde{a}_3^{(1)} + \tilde{a}_3^{(2)} = \frac{1152}{7\pi A^8} \iint_{D_A} x_1^8 B_3(\mathbf{x}, h) d^2x + \frac{45m_3}{14\pi A} + \mathcal{O}\left(\frac{1}{A^5}\right) =: \mathcal{T}_8.$$

Inspection of the left-hand sides of (2.140)–(2.144) reveals a number of linear dependencies (the rank of matrix of coefficients in front of \tilde{a}_2 , $\tilde{a}_3^{(1)}$, $\tilde{a}_3^{(2)}$ is merely 2). In particular, we should use only the following two of these relations which are straightforward to verify:

$$(2.145) \quad \mathcal{T}_0 = \mathcal{T}_4 - 2\mathcal{T}_6, \quad \mathcal{T}_0 = 4\mathcal{T}_6 - 25\mathcal{T}_8.$$

Substitution of (2.140), (2.142) and (2.143) in the first relation of (2.145) allows us to solve for m_3/A . This leads to the third-order estimate of m_3 given by (1.9).

Similarly, plugging (2.140), (2.143) and (2.144) into the second relation of (2.145) gives (1.11), a fourth-order estimate of m_3 .

3. NUMERICAL VALIDATION AND PRACTICAL CONSIDERATIONS

We demonstrate results by performing a numerical simulation on a synthetic example. In this example, we choose magnetisation distribution to consist of few magnetic dipoles (i.e. point-supported magnetisation sources) with positions $\vec{x}^{(1)} = (3.5, 3.0, 1.0)^T 10^{-5}$ m, $\vec{x}^{(2)} = (0.0, 0.0, 7.0)^T 10^{-5}$ m, $\vec{x}^{(3)} = (4.0, -5.5, 11.5)^T 10^{-5}$ m, $\vec{x}^{(4)} = (-4.0, 5.5, 2.5)^T 10^{-5}$ m, and magnetic moments $\vec{m}^{(1)} = (4.5, 3.5, 1.0)^T 10^{-12}$ A·m², $\vec{m}^{(2)} = (2.5, 4.5, 0.5)^T 10^{-12}$ A·m², $\vec{m}^{(3)} = (-3.0, 2.0, 2.5)^T 10^{-12}$ A·m², $\vec{m}^{(4)} = (-1.0, 2.0, 1.5)^T 10^{-12}$ A·m². The net moment of this magnetisation distribution is thus simply equal to

$$(3.1) \quad \vec{m}^{\text{true}} = \sum_{j=1}^4 \vec{m}^{(j)} = (3.0, 12.0, 5.5)^T 10^{-12} \text{ A} \cdot \text{m}^2.$$

The produced magnetic field B_3 is given by

$$(3.2) \quad B_3(\mathbf{x}, h) = \frac{\mu_0}{4\pi} \sum_{j=1}^4 \frac{3 \left(h - x_3^{(j)} \right) \left[\left(x_1 - x_1^{(j)} \right) m_1^{(j)} + \left(x_2 - x_2^{(j)} \right) m_2^{(j)} \right] + \left(2 \left(h - x_3^{(j)} \right)^2 - |\mathbf{x} - \mathbf{x}^{(j)}|^2 \right) m_3^{(j)}}{\left(|\mathbf{x} - \mathbf{x}^{(j)}|^2 + \left(h - x_3^{(j)} \right)^2 \right)^{5/2}},$$

and is measured on the disk $D_A = \{\mathbf{x} \in \mathbb{R}^2 : |\mathbf{x}| < A\}$ at the height $x_3 = h = 2.5 \cdot 10^{-4}$ m. Since we now work in Si units, we should recall Remark 1.2 and take into account the previously omitted factor $\mu_0 = 4\pi \cdot 10^{-7}$ N / A².

In order to check robustness of the moment estimates obtained in Theorem 1.1, we also perform simulations on data with a synthetic noise. Namely, we modify B_3 using additive Gaussian white noise with the amplitude $\sqrt{10^{-\text{SNR}/10} \cdot \text{Var}(B_3)}$, where SNR is the signal-to-noise ratio (in decibels) and $\text{Var}(B_3)$ is the variance of B_3 on D_A . For our simulations, we choose SNR = 20 dB which corresponds to the 10% noise level.

In Figure 3.1, we illustrate the field (3.2) and its noise component on the disk D_A of radius $A = 7.5 \cdot 10^{-4}$ m.

We shall now compute the integrals on the right-hand sides of (1.5)–(1.12) for different values of A . According to our asymptotic result for large A , we expect to see that, as A grows, the values of each of these integrals converge, with a different rate, to the value of a component of the net moment given by (3.1). Figures 3.2–3.3 show exactly that for the tangential and normal net moment components, respectively. We note that in Figure 3.3 and further figures involving the normal net moment component m_3 , a pair of the estimates are used: one with $x_j = x_1$ in (1.9), (1.11) and the other with $x_j = x_2$.

To illustrate the different convergence rates better, in Figures 3.4–3.5, we plot the differences $\log |m_j^{\text{true}} - m_j|$, $j \in \{1, 2, 3\}$, against $\log A$. This, in general, shows agreement with the estimates of the remainder term given in (1.5)–(1.12).

Finally, in Figure 3.6, we directly test the estimates of the net moment components when the magnetic field is contaminated by noise (with the noise model discussed above). We observe the persistence of the estimates for the tangential net moment components m_1 , m_2 , even though higher-order estimates perform worse. The situation is different for the estimates for the normal component m_3 in that the estimates of any order clearly have a growing trend with the increase of A , see Figure (3.7). This is deemed to be due to the presence of the factor A in all estimates (1.7), (1.9), (1.11) resulting in the noise amplification for large A . The issue, however, can be dealt with. One of possible remedies is demonstrated in Figure (3.8). There, a simple 11-point linear backward regression was used: namely, for each fixed value A , 11 (nearest) points with smaller values of A (hence lying within the same original measurement disk D_A) were used to establish and subtract a linear trend in A .

Let us now briefly comment on the fact that we chose to illustrate the results on a magnetisation distribution with a singular support. Besides its simplicity, this choice is physically motivated as any magnetisation can be

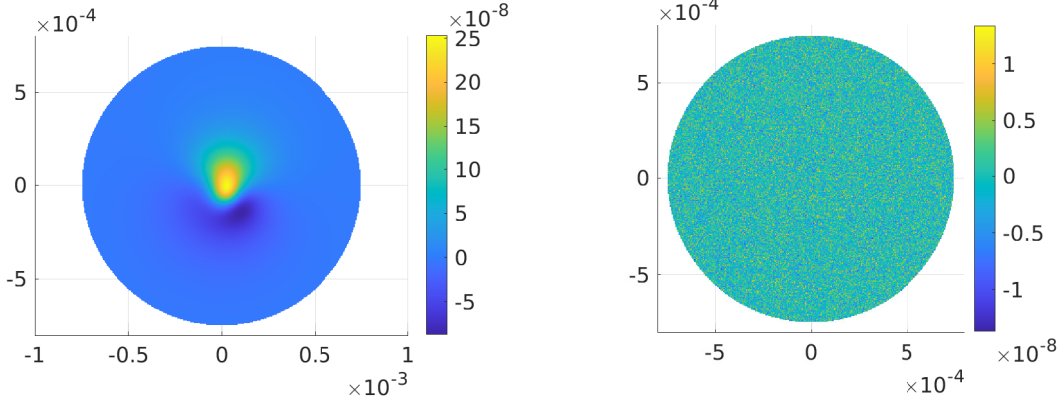


FIGURE 3.1. Magnetic field $B_3(\mathbf{x}, h)$ (left) and added noise (right) on D_A for $A = 7.5 \cdot 10^{-4}$ m

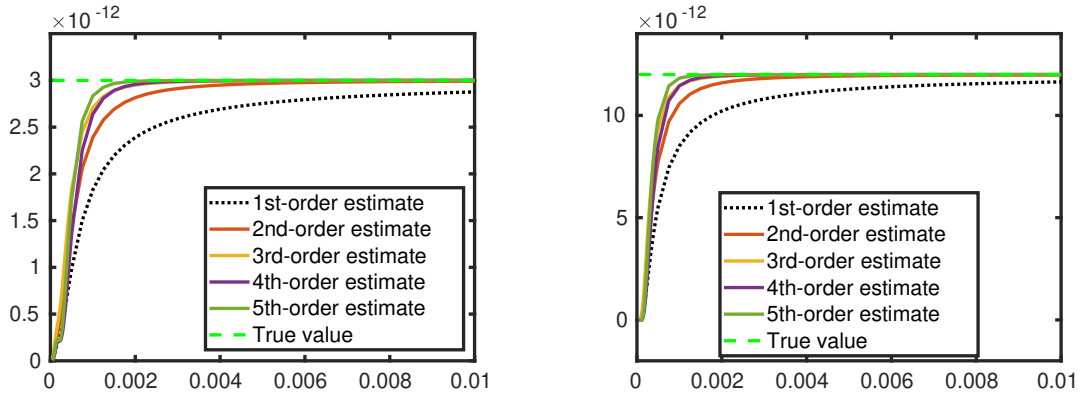


FIGURE 3.2. Estimates of the tangential net moment components m_1 (left) and m_2 (right) versus A

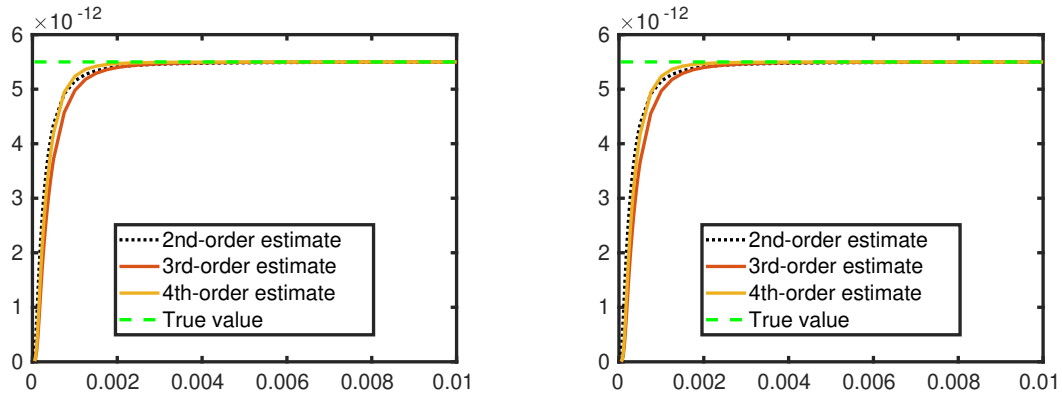
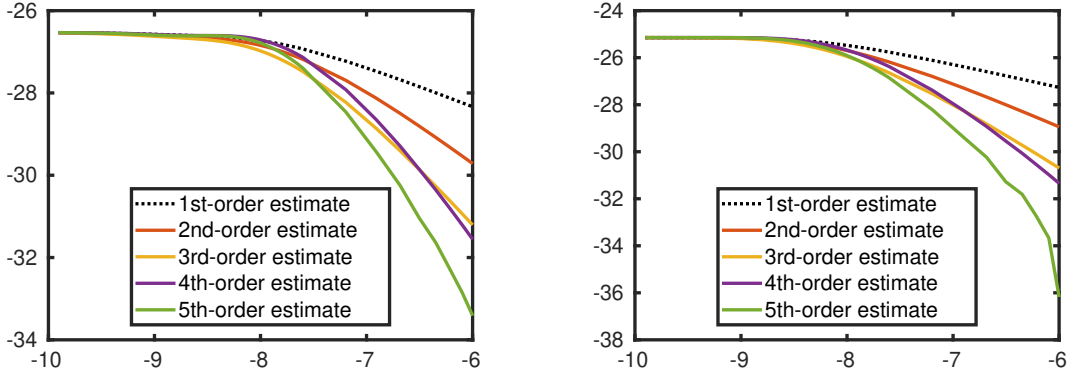
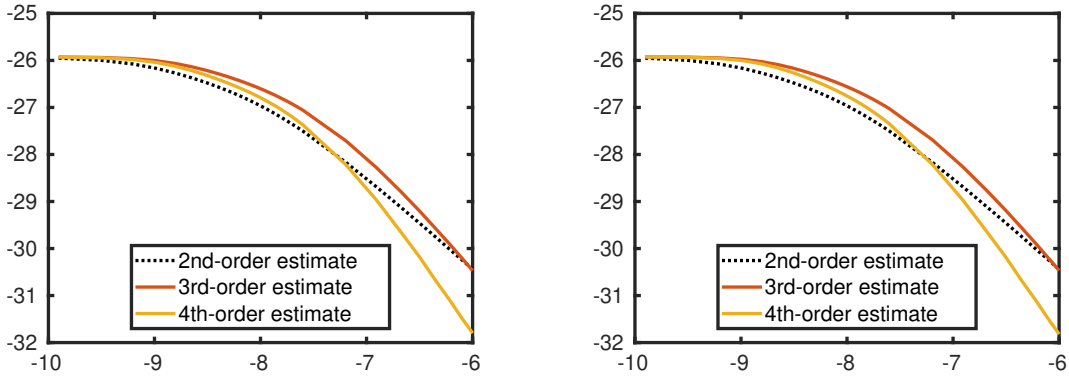
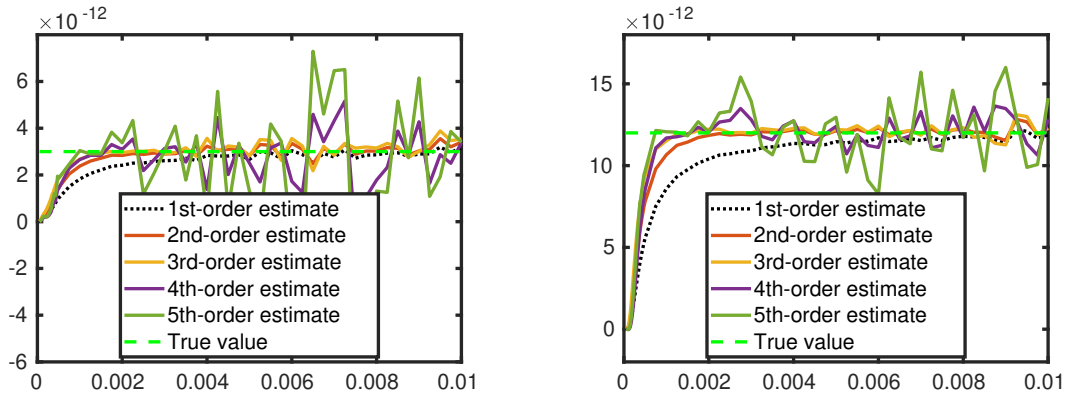


FIGURE 3.3. Estimates of the normal net moment component m_3 versus A : using x_1 (left) and x_2 (right) formulas

thought of a combination of dipole sources. From mathematical (numerical) viewpoint, a magnetic field produced by continuous magnetisation distribution is given by the integral whose numerical approximation (quadrature rule) is nothing but a weighted sum of dipoles. Consequently, we do not expect results for continuous magnetisation distributions to be of any drastical difference. On the other hand, this highlights the applicability of our methodology to the magnetisations that could be much more singular than square-integrable functions.

FIGURE 3.4. Logarithmic scale of the convergence for the estimates of m_1 (left) and m_2 (right)FIGURE 3.5. Logarithmic scale of the convergence for the m_3 estimates: using x_1 (left) and x_2 (right) formulasFIGURE 3.6. Estimates of the tangential net moment components m_1 (left) and m_2 (right) versus A . Noisy data.

4. DISCUSSION AND CONCLUSION

Motivated by a concrete experimental setup, we considered a problem of estimating of net magnetisation of a sample from one component of the magnetic field available in the limited measurement area in the plane above the sample. We approached this problem asymptotically, assuming the size of the measurement area to be large. We derived a set of explicit formulas for the asymptotic estimates of all three components of the net moment. For simplicity, we considered only the case of the circular geometry, i.e. where the measurement area is a disk. There is no doubt that results analogous to those in Theorem 1.1 could be deduced for the rectangular geometry

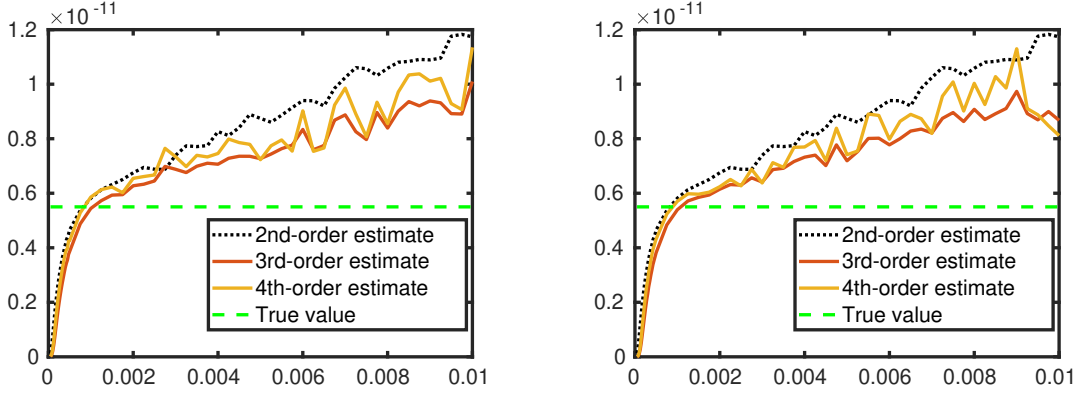


FIGURE 3.7. Estimates of the normal net moment component m_3 versus A : using x_1 (left) and x_2 (right) formulas. Noisy data.

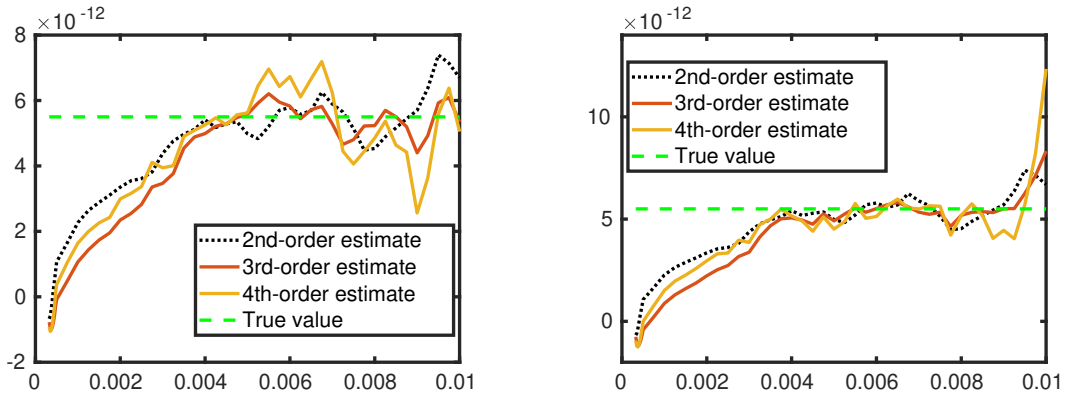


FIGURE 3.8. Estimates of the normal net moment component m_3 versus A : using x_1 (left) and x_2 (right) formulas. Noisy data with the correction.

which technically is even closer to what is used in Paleomagnetism lab at EAPS (Earth, Atmospheric & Planetary Sciences) department of MIT (Massachusetts Institute of Technology), USA. The main difference in obtaining such results with the present approach would be only in a technique of the asymptotic estimation of the Fourier integrals.

In this paper, we have obtained and proved asymptotic estimates up to the order 5 for the tangential net moment components m_1 , m_2 and up to the order 4 for the normal component m_3 . The main purpose was, however, to introduce a machinery that can generate asymptotic estimates of an arbitrary order. It is clear from the proof of Theorem 1.1 and auxiliary computations in Appendix that the asymptotic order of the estimates can be upgraded by proceeding in an established manner. This would require integration of the field against polynomials of a higher order. Practical advantages of it, however, are not yet obvious. First, such estimates are expected to be sensitive to any noise in B_3 : it was demonstrated in Section 3 that while lower- and middle-order estimates lead to expected results, estimates of higher order are much more prone to the noise. Second, when pursuing an estimate of a higher order, one should not forget that the obtained result is of only asymptotic nature: while a high-order estimate would be advantageous for very large values of A , it may not be so for a smaller A due to a potentially large value of a multiplicative constant (in A) in the remainder term. In particular, while Figure 3.2 shows a significant improvement due to the use of the second-order estimates compared to those of the first-order, the same figure together with Figure 3.3 demonstrate that, even without any noise, a higher-order estimate may be suboptimal for mid-range values of A . A relevant issue to bear in mind is that, apart from the basic asymptotic condition given by (1.4), proceeding to higher-order estimates assumes implicitly (but it is evident due to the form of the remainder terms) that the magnetisation is sufficiently localised so that its higher algebraic moments $L_{j_1, j_2, j_3}^{(n)} := \langle x_1^{j_1} x_2^{j_2} x_3^{j_3} M_n \rangle$, $n \in \{1, 2, 3\}$, (see (2.1)) do not grow too fast with respect to their order $j := j_1 + j_2 + j_3$, namely, that the quantities $L_{j_1, j_2, j_3}^{(n)} / A^j$ are not large for the value of A in question. A conclusion to draw from these observations is that the middle-order estimates (e.g. those of orders 2-3) are perhaps the best: both from practical perspective of proximity

to the true value of the net moment and from the viewpoint of the robustness to imperfect measurements. One can also consider a possibility of choosing an estimate which is best possible for a given magnitude of A . This is pertinent to the problem of finding an optimal truncation of asymptotic series.

In future works, we should study relation of our asymptotic results to the stable estimates of the net moment using constrained optimisation approach, in particular those obtained in [1] for the planar and regular (square-integrable) magnetisation. The strength of that approach is in its generality (potentially giving more than just a net moment estimate) and, in particular, that the result does not require any largeness of the measurement area size. However, it does not furnish explicit formulas for the net moment estimates: the obtained estimators (i.e. functions that the field needs to be integrated against in order to estimate a net moment component; in our case these are polynomials) could only be computed numerically by solving an integral-partial-differential equation. Moreover, these estimators, despite being a solution of a regularised version of the problem, could be seen to have rather irregular behavior with a rapid growth near the end of the measurement area and, as the authors noted, this feature is not the most practical factor of their result. Studying of other ways of stabilisation of the estimates and regularisation is also important. In this respect, it is natural to ask a question: what is the best way of using redundancy of a set of the asymptotic formulas for the net moment components to arrive at an estimate with the minimal effect of noise or to have the fastest (non-asymptotic) convergence to the true value of a net moment? For example, in the present approach, we have a natural redundancy for all estimates of m_3 of order 3 and higher due to the freedom of choice $x_j = x_1$ and $x_j = x_2$ in formulas (1.9), (1.11) and so on. In Figure 3.3, these estimates give almost indistinguishable results (it can be checked that the difference between the two is not zero but very small), and hence may not look directly useful. However, since having more relations than quantities to determine is always better when dealing with noise, this is still an advantage. Also, it is clear from Step 3 of the proof of Theorem 1.1 that, upon involving higher order polynomials, a richer set of lower-order estimates could be obtained.

Finally, the results of this work naturally connect to the issue of the asymptotic field extension. Indeed, the asymptotic field expansion at infinity (2.24) is seen to feature m_3 at the leading order and quantities $a_1^{(1)}$, $a_1^{(2)}$ at the next order. While the estimates of m_3 are given, to a different order, in (1.7), (1.9) and (1.11), it is evident from the proof of Theorem 1.1, that the quantities $a_1^{(1)}$, $a_1^{(2)}$ can also be asymptotically estimated, see e.g. (2.89) and (2.91) (and their versions with x_1 replaced by x_2 , as well as their lower-order analogs). Thus, this furnishes an explicit 3-term expansion of B_3 at infinity which, in general, when solving inverse magnetisation problems, should serve as a better alternative to a simple prolongation of the field by zero outside of the measurement area. Such a strategy can also potentially be used in setting up an iterative scheme. Of course, all of this is meaningful only when the actual measurement area is already sufficiently large so that the asymptotic estimates for m_3 , $a_1^{(1)}$, $a_1^{(2)}$ are sufficiently accurate.

ACKNOWLEDGEMENTS

The author acknowledges inspiring discussions on the topic of inverse magnetisation problems with L. Baratchart, S. Chevillard, J. Leblond, C. Villalobos-Guillen (Inria SAM, France), D. Hardin (Vanderbilt University, USA) and E. A. Lima (MIT, USA).

APPENDIX A

We collect here several useful results about cylindrical functions and some relevant integrals. In what follows we will use the notation \mathbb{N}_+ to denote natural numbers, $\mathbb{N}_0 := \mathbb{N}_+ \cup \{0\}$, and the notation \mathbb{Z} for integer numbers.

Basic facts about Bessel, Neumann and Struve functions. The Bessel function J_n of order $n \in \mathbb{Z}$ is an entire function satisfying the differential equation [11, (10.2.1)]

$$(A.1) \quad z^2 J_n''(z) + z J_n'(z) + (z^2 - n^2) J_n(z) = 0, \quad z \in \mathbb{C}.$$

We have the following series expansion [11, (10.2.2)]

$$(A.2) \quad J_n(z) = \left(\frac{z}{2}\right)^n \sum_{k=0}^{\infty} \frac{(-1)^k}{k! \Gamma(n+k+1)} \left(\frac{z}{2}\right)^{2k},$$

which, due to the entire character of J_n , is absolutely convergent for every $z \in \mathbb{C}$. Here, Γ denotes the Euler gamma function, for which we have, in particular, $\Gamma(k) = (k-1)!$ for $k \in \mathbb{N}_+$. Moreover, it is worth noting that $1/\Gamma(k) = 0$ for $k \in \mathbb{Z} \setminus \mathbb{N}_+$, which implies vanishing of negative powers of z in expansion (A.2) even for negative orders n .

The following integral representation holds [11, (10.9.1)]

$$(A.3) \quad J_n(x) = \frac{1}{\pi} \int_0^\pi \cos(nt - x \sin t) dt, \quad x \in \mathbb{R}.$$

In particular,

$$(A.4) \quad J_0(x) = \frac{1}{\pi} \int_0^\pi \cos(x \sin t) dt = \frac{1}{\pi} \int_0^\pi \cos(x \cos t) dt = \frac{1}{2\pi} \int_0^{2\pi} \cos(x \cos t) dt, \quad x \in \mathbb{R},$$

$$(A.5) \quad J_1(x) = \frac{1}{\pi} \int_0^\pi \cos(t - x \sin t) dt = \frac{1}{2\pi} \int_0^{2\pi} \sin(x \cos t) \cos t dt, \quad x \in \mathbb{R}.$$

For $x \gg 1$, $n \in \mathbb{N}_0$, the leading order asymptotics reads [11, (10.17.2)]

$$(A.6) \quad J_n(x) = \left(\frac{2}{\pi x}\right)^{1/2} \cos\left(x - \frac{n\pi}{2} - \frac{\pi}{4}\right) + \mathcal{O}\left(\frac{1}{x^{3/2}}\right),$$

where the estimate of the remainder term is due to the discussion in [11, Sect. 10.17(iii)].

The Bessel functions J_n , $n \in \mathbb{Z}$, satisfy the connection formula [11, (10.4.1)]

$$(A.7) \quad J_{-n}(x) = (-1)^n J_n(x),$$

as well as simple recurrence relations [11, (10.6.1)]

$$(A.8) \quad \frac{1}{x} J_n(x) = \frac{1}{2n} (J_{n-1}(x) + J_{n+1}(x)), \quad n \neq 0,$$

$$(A.9) \quad J'_n(x) = \frac{1}{2} (J_{n-1}(x) - J_{n+1}(x)).$$

In particular, (A.7) and (A.9) entail that

$$(A.10) \quad J'_0(x) = -J_1(x).$$

The Struve function H_n of order $n \in \mathbb{N}_0 \cup \{-1\}$ is an entire function defined by the absolutely convergent power series [11, (11.2.1)]

$$(A.11) \quad H_n(z) = \left(\frac{z}{2}\right)^{n+1} \sum_{k=0}^{\infty} \frac{(-1)^k}{\Gamma\left(k + \frac{3}{2}\right) \Gamma\left(k + n + \frac{3}{2}\right)} \left(\frac{z}{2}\right)^{2k}.$$

The companion Struve function K_n of order $n \in \mathbb{Z}$ is defined [11, (11.2.5)] as

$$(A.12) \quad K_n(z) = H_n(z) - Y_n(z),$$

where Y_n is the Neumann function.

For $x \gg 1$, $n \in \mathbb{N}_0$, the following asymptotics hold true [11, (11.6.1)]

$$(A.13) \quad K_n(x) = \frac{1}{\pi} \frac{2^{n+1} n!}{(2n)!} x^{n-1} + \mathcal{O}(x^{n-3}),$$

$$(A.14) \quad Y_n(x) = \left(\frac{2}{\pi x}\right)^{1/2} \sin\left(x - \frac{n\pi}{2} - \frac{\pi}{4}\right) + \mathcal{O}\left(\frac{1}{x^{3/2}}\right),$$

and the remainder terms are discussed in [11, Sect. 11.6(i)] and [11, Sect. 10.17(iii)], respectively. The asymptotic behavior of $H_n(x)$ for $x \gg 1$, $n \in \mathbb{N}_0$, hence follows from (A.12)–(A.14). In particular, for $x \gg 1$,

$$(A.15) \quad H_0(x) = \left(\frac{2}{\pi x}\right)^{1/2} \sin\left(x - \frac{\pi}{4}\right) + \mathcal{O}\left(\frac{1}{x}\right),$$

$$(A.16) \quad H_1(x) = \frac{2}{\pi} + \left(\frac{2}{\pi x}\right)^{1/2} \sin\left(x - \frac{3\pi}{4}\right) + \mathcal{O}\left(\frac{1}{x^{3/2}}\right).$$

Moreover, we have the following connection formula

$$(A.17) \quad H_{-1}(x) = \frac{2}{\pi} - H_1(x).$$

Some useful integrals pertinent to the cylindrical functions. The following lemmas establish several integral relations which are also crucial for the proof in Section 2.

Lemma A.1. For $n \in \mathbb{N}_+$, $\rho > 0$, the following identity holds

$$(A.18) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n+1}} dx = \frac{1}{4n^2 - 1} \left[2n \frac{J_1(\rho)}{\rho^{2n}} + \frac{J_1'(\rho)}{\rho^{2n-1}} - \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n-1}} dx \right].$$

Proof. First, upon integration by parts (using the asymptotic behavior at infinity of J_1 given by (A.6)), we have

$$(A.19) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n+1}} dx = \frac{1}{2n} \frac{J_1(\rho)}{\rho^{2n}} + \frac{1}{2n} \int_{\rho}^{\infty} \frac{J_1'(x)}{x^{2n}} dx.$$

Then, if we employ (A.1) to express J_1' in terms of J_1 and J_1'' , we obtain

$$\int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n+1}} dx = \frac{1}{2n} \frac{J_1(\rho)}{\rho^{2n}} + \frac{1}{2n} \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n+1}} dx - \frac{1}{2n} \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n-1}} dx - \frac{1}{2n} \int_{\rho}^{\infty} \frac{J_1''(x)}{x^{2n-1}} dx,$$

and hence

$$(A.20) \quad \int_{\rho}^{\infty} \frac{J_1''(x)}{x^{2n-1}} dx = \frac{J_1(\rho)}{\rho^{2n}} - (2n-1) \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n+1}} dx - \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n-1}} dx.$$

On the other hand, returning to (A.19) and integrating it by parts again, we arrive at

$$(A.21) \quad \begin{aligned} \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n+1}} dx &= \frac{1}{2n} \frac{J_1(\rho)}{\rho^{2n}} + \frac{1}{2n(2n-1)} \frac{J_1'(\rho)}{\rho^{2n-1}} + \frac{1}{2n(2n-1)} \int_{\rho}^{\infty} \frac{J_1''(x)}{x^{2n-1}} dx \\ &= \frac{1}{2n-1} \frac{J_1(\rho)}{\rho^{2n}} + \frac{1}{2n(2n-1)} \frac{J_1'(\rho)}{\rho^{2n-1}} - \frac{1}{2n} \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n+1}} dx - \frac{1}{2n(2n-1)} \int_{\rho}^{\infty} \frac{J_1(x)}{x^{2n-1}} dx. \end{aligned}$$

Here, in the second line, we eliminated the intergral term involving J_1'' using (A.20).

Rearranging the terms in (A.21) (solving for the quantity on the left-hand side), we deduce (A.18). \square

Lemma A.2. For $\rho > 0$, we have

$$(A.22) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x^3} dx = \frac{J_0(\rho)}{3\rho} + \frac{J_1(\rho)}{3\rho^2} - \frac{1}{3} - \frac{J_1(\rho)}{3} + \frac{\rho J_0(\rho)}{3} - \frac{\pi\rho}{6} [J_0(\rho) H_1(\rho) - J_1(\rho) H_0(\rho)].$$

Proof. Note that using (A.8), we have

$$(A.23) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x^3} dx = \frac{1}{2} \int_{\rho}^{\infty} \frac{J_0(x)}{x^2} dx + \frac{1}{2} \int_{\rho}^{\infty} \frac{J_2(x)}{x^2} dx.$$

Let us start by transforming the first term in (A.23), namely,

$$(A.24) \quad \begin{aligned} \int_{\rho}^{\infty} \frac{J_0(x)}{x^2} dx &= \frac{J_0(\rho)}{\rho} + \int_{\rho}^{\infty} \frac{J_0'(x)}{x} dx \\ &= \frac{J_0(\rho)}{\rho} + J_0'(\rho) - \int_{\rho}^{\infty} J_0(x) dx. \end{aligned}$$

Here, in the first line, we employed integration by parts (together with the asymptotic behavior at infinity of J_0 given by (A.6)). To arrive at the second line, we used the identity

$$\frac{J_0'(x)}{x} = -J_0(x) - J_0''(x), \quad x \neq 0,$$

implied by (A.1).

Performing the same procedure with the second term in (A.23), we have

$$\begin{aligned} \int_{\rho}^{\infty} \frac{J_2(x)}{x^2} dx &= \frac{J_2(\rho)}{\rho} + \int_{\rho}^{\infty} \frac{J_2'(x)}{x} dx \\ &= \frac{J_2(\rho)}{\rho} + J_2'(\rho) - \int_{\rho}^{\infty} J_2(x) dx + 4 \int_{\rho}^{\infty} \frac{J_2(x)}{x^2} dx, \end{aligned}$$

and, hence,

$$(A.25) \quad \int_{\rho}^{\infty} \frac{J_2(x)}{x^2} dx = -\frac{1}{3} \left[\frac{J_2(\rho)}{\rho} + J_2'(\rho) - \int_{\rho}^{\infty} J_2(x) dx \right].$$

Expressions (A.24) and (A.25) imply that the integral on the left-hand side of (A.23) is expressible in terms of two integral quantities: $\int_{\rho}^{\infty} J_0(x) dx$ and $\int_{\rho}^{\infty} J_2(x) dx$. Let us now show that these quantities are simply related:

$$(A.26) \quad \int_{\rho}^{\infty} J_2(x) dx = \int_{\rho}^{\infty} J_0(x) dx + 2J_1(\rho),$$

and, moreover,

$$(A.27) \quad \int_{\rho}^{\infty} J_0(x) dx = 1 - \rho J_0(\rho) - \frac{\pi}{2} \rho [J_1(\rho) H_0(\rho) - J_0(\rho) H_1(\rho)].$$

Recalling asymptotics (A.6), we note that the integrals on the left-hand sides of (A.26), (A.27) are not absolutely convergent, and hence an additional care with technical manipulations is needed. Namely, we shall first replace the integration range (ρ, ∞) with (ρ, R) for arbitrary finite $R > \rho$, and then pass to the limit as $R \rightarrow +\infty$. We shall proceed in several steps.

Step 1: Establishing of (A.26)

We start with (A.26) and use integral representation (A.3):

$$J_2(x) = \frac{1}{\pi} \int_0^{\pi} \cos(2t - x \sin t) dt = \frac{1}{\pi} \int_0^{\pi} \cos(x \sin t - 2t) dt.$$

Consequently, exchanging the order of integration (permissible due to the regularity of the integrand and finiteness of the integration limits), we obtain

$$(A.28) \quad \begin{aligned} \int_{\rho}^R J_2(x) dx &= \frac{1}{\pi} \int_0^{\pi} \frac{\sin(R \sin t - 2t) - \sin(\rho \sin t - 2t)}{\sin t} dt \\ &= \frac{1}{\pi} \int_0^{\pi} \frac{\sin(R \sin t) - \sin(\rho \sin t)}{\sin t} dt - \frac{2}{\pi} \int_0^{\pi} [\sin(R \sin t) \sin t + \cos(R \sin t) \cos t] dt \\ &\quad + \frac{2}{\pi} \int_0^{\pi} [\sin(\rho \sin t) \sin t + \cos(\rho \sin t) \cos t] dt \\ &= \int_{\rho}^R J_0(x) dx - 2(J_1(R) - J_1(\rho)), \end{aligned}$$

where we used the identities

$$\sin(R \sin t - 2t) = (1 - 2 \sin^2 t) \sin(R \sin t) - 2 \sin t \cos t \cos(R \sin t),$$

$$\sin(R \sin t) \sin t + \cos(R \sin t) \cos t = \cos(R \sin t - t),$$

$$\frac{1}{\pi} \int_0^{\pi} \cos(R \sin t - t) dt = J_1(R),$$

$$\frac{1}{\pi} \int_0^{\pi} \frac{\sin(R \sin t) - \sin(\rho \sin t)}{\sin t} dt = \int_{\rho}^R J_0(x) dx.$$

and, except for the last one, also their analogs with ρ instead of R . The first two of these identities are purely trigonometrical whereas the last two are due to (A.4)–(A.5).

Passing to the limit $R \rightarrow +\infty$ in (A.28) using asymptotics (A.6), we thus conclude with (A.26).

Step 2: Establishing of (A.27)

To deduce relation (A.27), we use [11, (10.22.2)] to write

$$(A.29) \quad \begin{aligned} \int_{\rho}^R J_0(x) dx &= \frac{\pi}{2} R [J_0(R) H_{-1}(R) - J_{-1}(R) H_0(R)] - \frac{\pi}{2} \rho [J_0(\rho) H_{-1}(\rho) - J_{-1}(\rho) H_0(\rho)] \\ &= R J_0(R) - \rho J_0(\rho) + \frac{\pi}{2} R [J_1(R) H_0(R) - J_0(R) H_1(R)] - \frac{\pi}{2} \rho [J_1(\rho) H_0(\rho) - J_0(\rho) H_1(\rho)], \end{aligned}$$

where we used (A.7), (A.17) in passing to the second line.

Note that, employing asymptotics (A.6), (A.15)–(A.16), we have, for $R \gg 1$,

$$\begin{aligned} RJ_0(R) + \frac{\pi}{2}R[J_1(R)H_0(R) - J_0(R)H_1(R)] &= \cos\left(R - \frac{3\pi}{4}\right)\sin\left(R - \frac{\pi}{4}\right) - \cos\left(R - \frac{\pi}{4}\right)\sin\left(R - \frac{3\pi}{4}\right) + \mathcal{O}\left(\frac{1}{R^{1/2}}\right) \\ &= 1 + \mathcal{O}\left(\frac{1}{R^{1/2}}\right). \end{aligned}$$

Therefore, by passing to the limit as $R \rightarrow +\infty$ in (A.29), we obtain (A.27).

Step 3: Conclusion of the proof

Plugging (A.27) into (A.26) and (A.24), we obtain

$$(A.30) \quad \int_{\rho}^{\infty} J_2(x) dx = 1 + 2J_1(\rho) - \rho J_0(\rho) - \frac{\pi}{2}\rho[J_1(\rho)H_0(\rho) - J_0(\rho)H_1(\rho)],$$

$$(A.31) \quad \int_{\rho}^{\infty} \frac{J_0(x)}{x^2} dx = \frac{J_0(\rho)}{\rho} + J_0'(\rho) - 1 + \rho J_0(\rho) + \frac{\pi}{2}\rho[J_1(\rho)H_0(\rho) - J_0(\rho)H_1(\rho)],$$

respectively.

Substitution of (A.30) into (A.25) gives

$$(A.32) \quad \int_{\rho}^{\infty} \frac{J_2(x)}{x^2} dx = -\frac{J_2(\rho)}{3\rho} - \frac{1}{3}J_2'(\rho) + \frac{1}{3} + \frac{2}{3}J_1(\rho) - \frac{1}{3}\rho J_0(\rho) - \frac{\pi}{6}\rho[J_1(\rho)H_0(\rho) - J_0(\rho)H_1(\rho)].$$

Finally, using (A.31)–(A.32) in (A.23), we arrive at

$$(A.33) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x^3} dx = \frac{J_0(\rho)}{2\rho} - \frac{J_2(\rho)}{6\rho} + \frac{1}{2}J_0'(\rho) - \frac{1}{6}J_2'(x) - \frac{1}{3} + \frac{\rho}{3}J_0(\rho) + \frac{1}{3}J_1(\rho) + \frac{\pi}{6}\rho[J_1(\rho)H_0(\rho) - J_0(\rho)H_1(\rho)].$$

Using (A.10) and the recursive identities (due to (A.8)–(A.9))

$$\begin{aligned} J_2(\rho) &= \frac{2}{\rho}J_1(\rho) - J_0(\rho), & J_3(\rho) &= \frac{4}{\rho}J_2(\rho) - J_1(\rho) = \frac{8}{\rho^2}J_1(\rho) - \frac{4}{\rho}J_0(\rho) - J_1(\rho), \\ J_2'(\rho) &= \frac{1}{2}[J_1(\rho) - J_3(\rho)] = -\frac{4}{\rho^2}J_1(\rho) + \frac{2}{\rho}J_0(\rho) + J_1(\rho), \end{aligned}$$

we transform (A.33) into the desired relation (A.22). □

Lemma A.3. For $\rho > 0$, in addition to (A.22), we have the following identities

$$(A.34) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x} dx = \frac{J_1(\rho)}{\rho^2} + \frac{J_1'(\rho)}{\rho} - \frac{1}{\rho}J_0(\rho) + 1 + J_1(\rho) - \rho J_0(\rho) + \frac{\pi\rho}{2}[J_0(\rho)H_1(\rho) - J_1(\rho)H_0(\rho)],$$

$$(A.35) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x^5} dx = \frac{4J_1(\rho)}{15\rho^4} + \frac{J_1'(\rho)}{15\rho^3} - \frac{J_0(\rho)}{45\rho} - \frac{J_1(\rho)}{45\rho^2} + \frac{1}{45} + \frac{J_1(\rho)}{45} - \frac{\rho J_0(\rho)}{45} + \frac{\pi\rho}{90}[J_0(\rho)H_1(\rho) - J_1(\rho)H_0(\rho)],$$

$$(A.36) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x^7} dx = \frac{6}{35} \frac{J_1(\rho)}{\rho^6} + \frac{1}{35} \frac{J_1'(\rho)}{\rho^5} - \frac{4J_1(\rho)}{525\rho^4} - \frac{J_1'(\rho)}{525\rho^3} + \frac{J_0(\rho)}{1575\rho} + \frac{J_1(\rho)}{1575\rho^2} - \frac{1}{1575} - \frac{J_1(\rho)}{1575} + \frac{\rho J_0(\rho)}{1575} - \frac{\pi\rho}{3150}[J_0(\rho)H_1(\rho) - J_1(\rho)H_0(\rho)].$$

Proof. We shall prove identities (A.34)–(A.36) sequentially.

Applying Lemma A.1 with $n = 1$, we have

$$\int_{\rho}^{\infty} \frac{J_1(x)}{x^3} dx = \frac{2}{3} \frac{J_1(\rho)}{\rho^2} + \frac{1}{3} \frac{J_1'(\rho)}{\rho} - \frac{1}{3} \int_{\rho}^{\infty} \frac{J_1(x)}{x} dx,$$

and, hence,

$$(A.37) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x} dx = 2 \frac{J_1(\rho)}{\rho^2} + \frac{J_1'(\rho)}{\rho} - 3 \int_{\rho}^{\infty} \frac{J_1(x)}{x^3} dx.$$

Using Lemma A.2, (A.37) becomes (A.34).

Application of Lemma A.1 with $n = 2$ yields

$$(A.38) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x^5} dx = \frac{4}{15} \frac{J_1(\rho)}{\rho^4} + \frac{1}{15} \frac{J_1'(\rho)}{\rho^3} - \frac{1}{15} \int_{\rho}^{\infty} \frac{J_1(x)}{x^3} dx.$$

With the help of Lemma A.2, (A.38) transforms into (A.35).

Finally, we use Lemma A.1 with $n = 3$ to obtain

$$(A.39) \quad \int_{\rho}^{\infty} \frac{J_1(x)}{x^7} dx = \frac{6}{35} \frac{J_1(\rho)}{\rho^6} + \frac{1}{35} \frac{J_1'(\rho)}{\rho^5} - \frac{1}{35} \int_{\rho}^{\infty} \frac{J_1(x)}{x^5} dx.$$

Substitution of the already proven identity (A.35) into (A.39) furnishes (A.36). \square

Lemma A.4. For $\alpha \in \mathbb{R}$, $m, n \in \mathbb{N}_0$, we have the following identities

$$(A.40) \quad \int_0^{2\pi} \cos(\alpha \cos \theta) \cos^{2m+1} \theta \sin^n \theta d\theta = 0,$$

$$(A.41) \quad \int_0^{2\pi} \cos(\alpha \cos \theta) \cos^m \theta \sin^{2n+1} \theta d\theta = 0,$$

$$(A.42) \quad \int_0^{2\pi} \sin(\alpha \cos \theta) \cos^m \theta \sin^{2n+1} \theta d\theta = 0,$$

$$(A.43) \quad \int_0^{2\pi} \sin(\alpha \cos \theta) \cos^{2m} \theta \sin^n \theta d\theta = 0.$$

Proof. We shall prove identities (A.40)–(A.43) one after another.

Using periodicity of the integrand, we can shift the interval from $(0, 2\pi)$ to $(-\pi/2, 3\pi/2)$ and further split it in two:

$$(A.44) \quad \begin{aligned} \int_0^{2\pi} \cos(\alpha \cos \theta) \cos^{2m+1} \theta \sin^n \theta d\theta &= \int_{-\pi/2}^{3\pi/2} \cos(\alpha \cos \theta) \cos^{2m+1} \theta \sin^n \theta d\theta \\ &= \int_{-\pi/2}^{\pi/2} \dots + \int_{\pi/2}^{3\pi/2} \dots \end{aligned}$$

Performing the change of variable $y = \sin \theta$, $dy = \cos \theta d\theta$ in each of the integrals on the second line of (A.44), we have

$$\begin{aligned} \int_{-\pi/2}^{\pi/2} \cos(\alpha \cos \theta) \cos^{2m+1} \theta \sin^n \theta d\theta &= \int_{-1}^1 \cos(\alpha \sqrt{1-y^2}) (1-y^2)^m dy, \\ \int_{\pi/2}^{3\pi/2} \cos(\alpha \cos \theta) \cos^{2m+1} \theta \sin^n \theta d\theta &= \int_1^{-1} \cos(-\alpha \sqrt{1-y^2}) (-\sqrt{1-y^2})^{2m} dy \\ &= - \int_{-1}^1 \cos(\alpha \sqrt{1-y^2}) (1-y^2)^m dy. \end{aligned}$$

Here, we used the fact that $\cos \theta = \sqrt{1-\sin^2 \theta} \geq 0$ for $\theta \in [-\pi/2, \pi/2]$ and $\cos \theta = -\sqrt{1-\sin^2 \theta} \leq 0$ for $\theta \in [\pi/2, 3\pi/2]$ which makes the integrands in each of both terms in the decomposition single-valued. Therefore, both integral quantities on the second line of (A.44) are opposite to each other in sign, and this entails (A.40).

To show (A.41), a useful change of variable is $y = \cos \theta$, $dy = -\sin \theta d\theta$. The integrand is single-valued on the whole interval, and hence

$$\int_0^{2\pi} \cos(\alpha \cos \theta) \cos^m \theta \sin^{2n+1} \theta d\theta = - \int_1^{-1} \cos(\alpha y) y^m (1-y^2)^n dy.$$

Recognising that the lower and upper limits of the integration are the same, we deduce the desired vanishing of the integral.

Similarly, we have

$$\int_0^{2\pi} \sin(\alpha \cos \theta) \cos^m \theta \sin^{2n+1} \theta d\theta = - \int_1^{-1} \sin(\alpha y) y^m (1-y^2)^n dy = 0,$$

which proves (A.42).

Finally, to show (A.43), we use the fact that identity (A.40) holds, in particular, for any $\alpha \in [0, \beta]$ with arbitrary $\beta \geq 0$. Integrating it in α over this interval and interchanging the order of integration (permissible by the regularity of the integrand and the finite integration range), we obtain

$$0 = \int_0^\beta \int_0^{2\pi} \cos(\alpha \cos \theta) \cos^{2m+1} \theta \sin^n \theta d\theta d\alpha = \int_0^{2\pi} \sin(\beta \cos \theta) \cos^{2m} \theta \sin^n \theta d\theta.$$

Since the same reasoning also works $\beta \leq 0$ by working with an interval $[\beta, 0]$, identity (A.43) is thus proved up to a change of the notation β to α . \square

REFERENCES

- [1] Baratchart, L., Chevillard, S., Hardin, D., Leblond, J., Lima, E., Marmorat, J. P., "Magnetic moment estimation and bounded extremal problems", *Inverse Problems and Imaging*, 13 (1), 2019.
- [2] Baratchart, L., Chevillard, S., Leblond, J., Lima, E., Ponomarev, D., "Asymptotic method for estimating magnetic moments from field measurements on a planar grid" (preprint), HAL Id: hal-01421157, 2018.
- [3] Baratchart, L., Hardin, D. P., Lima, E. A., Saff, E. B., Weiss, B. P., "Characterizing kernels of operators related to thin-plate magnetizations via generalizations of Hodge decompositions", *Inverse Problems*, 29 (1), 2013.
- [4] Baratchart, L., Leblond, J., Lima, E., Ponomarev, D., "Magnetization moment recovery using Kelvin transformation and Fourier analysis", *J. Phys. Conf. Ser.*, 904, 2017.
- [5] Baratchart, L., Villalobos-Guillen, C., Hardin, D. P., Northington, M. C., Saff, E. B., "Inverse potential problems for divergence of measures with total variation regularization", *Foundations of Computational Mathematics*, 20 (5), 1273–1307, 2020.
- [6] Evans, L. C., "Partial differential equations", American Mathematical Society, Second edition, 2010.
- [7] Fu, R. R., Lima, E. A., Volk, M. W., Trubko, R., "High-sensitivity moment magnetometry with the quantum diamond microscope", *Geochemistry, Geophysics, Geosystems*, 21 (8), 2020.
- [8] Lima, E. A., Weiss, B. P., Baratchart, L., Hardin, D. P., Saff, E. B., "Fast inversion of magnetic field maps of unidirectional planar geological magnetization", *J. Geophys. Res.: Solid Earth*, 118 (6), 2723–2752, 2013.
- [9] Lima, E. A., Weiss, B. P., "Obtaining vector magnetic field maps from single-component measurements of geological samples", *Journal of Geophysical Research: Solid Earth*, 114 (B6), 2009.
- [10] Lima, E. A., Weiss, B. P., "Ultra-high sensitivity moment magnetometry of geological samples using magnetic microscopy", *Geochemistry, Geophysics, Geosystems*, 17 (9), 3754–3774, 2016.
- [11] Olver, F. W. J., Lozier, D. W., Boisvert, R. F., Clark, C. W., "NIST Handbook of Mathematical Functions", Cambridge University Press, 2010.
- [12] Ponomarev, D., "Some inverse problems with partial data", Doctoral thesis, Université de Nice - Sophia Antipolis (University of Côte d'Azur), 2016.
- [13] Stein, E. M., Weiss, G., "Introduction to Fourier Analysis on Euclidean Spaces", Princeton University Press, 2016.
- [14] Strichartz, R., "Guide to Distribution Theory and Fourier Transforms", CRC Press, 1994.
- [15] Weiss, B. P., Lima, E. A., Fong, L. E., Baudenbacher, F. J., "Paleomagnetic analysis using SQUID microscopy", *Journal of Geophysical Research: Solid Earth*, 112 (B9), 2007.