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# Multi-probe analysis of the galaxy cluster CL J1226.9+3332

## Hydrostatic mass and hydrostatic-to-lensing bias

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### ABSTRACT

The precise estimation of the mass of galaxy clusters is a major issue for cosmology. Large galaxy cluster surveys rely on scaling laws that relate cluster observables to their masses. From the high-resolution observations of ~45 galaxy clusters with the NIKA2 and *XMM-Newton* instruments, the NIKA2 Sunyaev-Zel'dovich Large Program should provide an accurate scaling relation between the thermal Sunyaev-Zel'dovich effect and the hydrostatic mass. In this paper we present an exhaustive analysis of the hydrostatic mass of the well-known galaxy cluster CL J1226.9+3332, the highest-redshift cluster in the NIKA2 Sunyaev-Zel'dovich Large Program at  $z = 0.89$ . We combined the NIKA2 observations with thermal Sunyaev-Zel'dovich data from the NIKA, Bolocam, and MUSTANG instruments and *XMM-Newton* X-ray observations, and tested the impact of the systematic effects on the mass reconstruction. We conclude that slight differences in the shape of the mass profile can be crucial when defining the integrated mass at  $R_{500}$ , which demonstrates the importance of the modelling in the mass determination. We prove the robustness of our hydrostatic mass estimates by showing the agreement with all the results found in the literature. Another key factor for cosmology is the bias of the masses estimated assuming the hydrostatic equilibrium hypothesis. Based on the lensing convergence maps from the Cluster Lensing And Supernova survey with *Hubble* (CLASH) data, we obtain the lensing mass estimate for CL J1226.9+3332. From this we are able to measure the hydrostatic-to-lensing mass bias for this cluster, which spans from  $1 - b_{\text{HSE}/\text{lens}} \sim 0.7$  to 1, presenting the impact of data sets and mass reconstruction models on the bias.

**Key words.** galaxies: clusters: intracluster medium – galaxies: clusters: individual: CL J1226.9+3332 – techniques: high angular resolution – cosmology: observations

## 1. Introduction

Galaxy clusters are formed by gravitational collapse at the last step of the hierarchical structure formation process (Kravtsov & Borgani 2012). Thus, they are tracers of the large-scale structure formation physics. Their abundance in mass and redshift is sensitive to the initial conditions in the primordial Universe and its expansion history and matter content (Huterer et al. 2015). Therefore, galaxy clusters are probes of the underlying cosmology (Allen et al. 2011).

Large catalogues of clusters (e.g., Planck Collaboration XVIII 2015; Rykoff et al. 2016; Adami et al. 2018; Bleem et al. 2020) have allowed us, in the past few years, to constrain cosmological parameters. Nevertheless, these results show some tension with respect to the cosmology obtained from the cosmic microwave background (CMB) power spectrum analysis (Planck Collaboration XXI 2013; Planck Collaboration XXIV 2015; Salvati et al. 2018), in line with a more general problem, in that early and late Universe probes are giving different results (Verde et al. 2019). Cluster-based cosmological analyses rely on their distribution in mass and one source of the discrepancy may be the inaccuracy of those mass estimates (e.g., Pratt et al. 2019; Salvati et al. 2020).

About 85% of the total mass of clusters of galaxies is composed of dark matter. The remaining 15% corresponds to the hot intracluster medium (ICM) and the galaxies in the cluster. For this reason, most of the mass content of the clusters is not directly observable and it needs to be estimated either from the gravitational potential reconstruction or via scaling relations that link cluster observables to their masses (Pratt et al. 2019). The gravitational potential of clusters can be inferred from their gravitational lensing effect on background sources (Bartelmann 2010), from the dynamics of member galaxies (Biviano & Girardi 2003; Aguado-Barahona et al. 2022), or from the combination of the thermodynamical properties of the gas in the ICM under the hydrostatic equilibrium (HSE) hypothesis. Scaling relations that allow us to recover the mass have been measured and calibrated for various cluster observables at different wavelengths: in the optical and infrared domains the galaxy member richness (Andreon & Hurn 2010), in X-ray the emitted luminosity of the hot gaseous ICM (Pratt et al. 2009), and at millimetre wavelengths the amplitude of the Sunyaev-Zel'dovich (SZ) effect (Sunyaev & Zeldovich 1972; Planck Collaboration XI 2011) proportional to the thermal energy in the ICM.

The thermal SZ (tSZ) effect (Sunyaev & Zeldovich 1972), corresponding to the spectral distortion of the CMB photons due to the inverse-Compton scattering on the hot thermal electrons of the ICM, is an excellent way of detecting clusters as its amplitude is not affected by cosmological dimming. Large surveys, in particular the *Planck* satellite observations (Planck Collaboration XVIII 2015) and the ground-based Atacama Cosmology Telescope (ACT, Hilton et al. 2018, 2021) and South Pole Telescopes (SPT, Bleem et al. 2020), have obtained large SZ-detected galaxy cluster catalogues. However, they need to rely on the SZ-mass scaling relation in order to carry out cosmological analyses. To date, most of the SZ-mass scaling relations have been mainly determined from low-redshift ( $z < 0.5$ ) cluster samples with masses obtained from X-ray observations (Arnaud et al. 2010; Planck Collaboration XI 2011). Other scaling relations, with optical data for example (Saro et al. 2017), are also used.

In any case, it is essential to study the redshift evolution of these scaling relations as they have an impact the cosmological results (Salvati et al. 2020). When building scaling relations, another key aspect is the impact of the assumptions on which the mass computations rely, such as the HSE hypothesis. High spatial resolution cluster observations help us to assess these assumptions as we study the impact of the dynamical states of clusters on the mass-observable scaling relations.

The NIKA2 SZ Large Program (Mayet et al. 2020; Perotto et al. 2022) described in Sect. 3.1.1 seeks to address the above-mentioned issues. The work presented in this paper constitutes the third analysis of a cluster in the NIKA2 SZ Large Program. The first analysis on PSZ2 G144.83+25.11 comprised a science verification study, as well as the proof of the impact of substructures in the reconstruction of the physical cluster properties (Ruppin et al. 2018). The second, the worst-case scenario for the NIKA2 SZ Large Program, analysed the ACT-CL J0215.4+0030 galaxy cluster, proving the quality of NIKA2 camera (Adam et al. 2018a; Bourrion et al. 2016; Calvo et al. 2016; Perotto et al. 2020) in the most challenging case of a high-redshift and low-mass cluster (Kéruzoré et al. 2020). In this work we present a study on the HSE mass of the CL J1226.9+3332 galaxy cluster and the impact of different systematic effects on the recovered mass, and make a thorough comparison to previous works. Moreover, we re-estimate its lensing mass and compare it to the mass obtained under the hydrostatic assumption, which allows us to compute the hydrostatic-to-lensing mass bias.

This paper is organised as follows. In Sect. 2 we present a summary of the results in the literature for CL J1226.9+3332. In Sect. 3 we describe the observations of the ICM with NIKA2 and *XMM-Newton* instruments. The reconstruction of the thermodynamical properties of the ICM is presented in Sect. 4. We detail the pressure profile reconstruction from NIKA2 maps, accounting for systematic effects due to the data reduction process and point source contamination. We compare these profiles to previous results and to the X-ray pressure profile. In Sect. 5 we describe the HSE mass estimates from the combination of SZ and X-ray data and test the robustness of the results against data processing and modelling effects. In Sect. 6 we present the lensing masses that we estimate for CL J1226.9+3332. All the results are discussed together in Sect. 7, where we compute the hydrostatic-to-lensing mass bias for this cluster. The conclusions are given in Sect. 8.

Throughout this paper we assume a flat  $\Lambda$ CDM cosmology with  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$  and  $\Omega_{m,0} = 0.3$ . With this assumption, 1 arcmin corresponds to a distance of 466 kpc at the cluster redshift  $z = 0.89$ .

## 2. The CL J1226.9+3332 galaxy cluster

This paper focuses on the CL J1226.9+3332 galaxy cluster, also known as PSZ2-G160.83+81.66. Discovered by the Wide Angle ROSAT Pointed Survey (WARPS, Ebeling et al. 2001), it has already been studied at different wavelengths: X-ray (Maughan et al. 2004; Bonamente et al. 2006), visible (Jee & Tyson 2009), and millimetre (Bonamente et al. 2006; Mroczkowski et al. 2009; Korngut et al. 2011; Adam et al. 2015) wavelengths. Located at redshift 0.89 (Planck Collaboration XVIII 2015; Aguado-Barahona et al. 2022), it is the highest-redshift cluster of the NIKA2 SZ Large Program sample, with the X-ray peak at  $(\text{RA}, \text{Dec})_{J2000} = (12\text{h}26\text{m}58.37\text{s}, +33\text{d}32\text{m}47.4\text{s})$  according to Cavagnolo et al. (2009). Less than 2 arcsecond away from this peak, its brightest cluster galaxy (BCG) is located at  $(\text{RA}, \text{Dec})_{J2000} = (12\text{h}26\text{m}58.25\text{s}, +33\text{d}32\text{m}48.57\text{s})$  according to Holden et al. (2009).

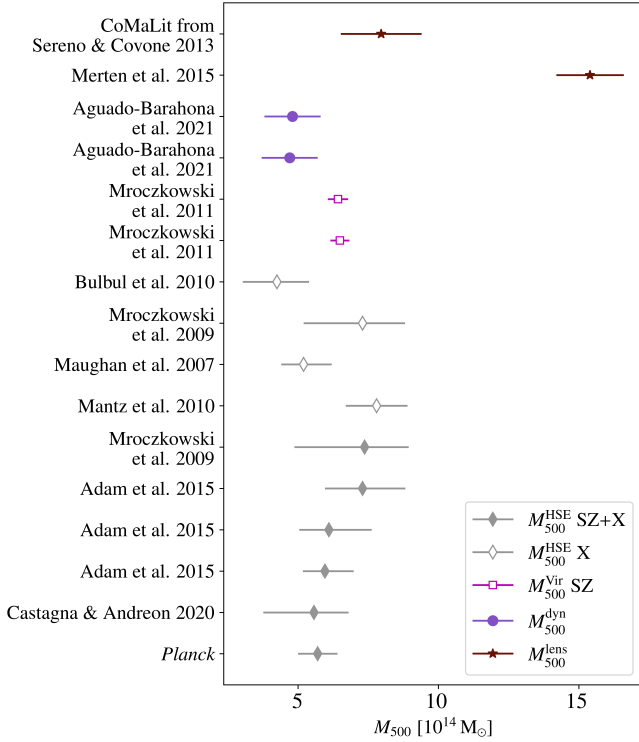
### 2.1. Previous observations

Since the first SZ observations with BIMA (Joy et al. 2001), the projected morphology of CL J1226.9+3332 appeared to have a quite circular symmetry. Nevertheless, the combination of *XMM-Newton* and *Chandra* X-ray data (Maughan et al. 2007) showed a region, at  $\sim 40''$  to the south-west of the X-ray peak with much higher temperature than the average in the ICM. This substructure was also confirmed by posterior SZ analyses with MUSTANG (Korngut et al. 2011) and NIKA (Monfardini et al. 2011; Adam et al. 2015, 2018b). Romero et al. (2018; hereafter R18) performed a study combining SZ data from the NIKA, MUSTANG, and Bolocam instruments. Their different capabilities allowed us to probe different angular scales in the reconstruction of ICM properties and agreed with a non-relaxed cluster core description for CL J1226.9+3332. In this work we make use of the pressure profiles obtained from the NIKA, MUSTANG, and Bolocam data summarised in Table 2 in Romero et al. (2018).

Lensing data from the Cluster Lensing And Supernova survey with *Hubble* (CLASH, Zitrin et al. 2015), as well as the galaxy distribution in the cluster (Jee & Tyson 2009), agree on the existence of a main clump centred on the BCG and a secondary clump to the south-west. However, this second region does not appear as a structure in X-ray surface brightness (Maughan et al. 2007). One hypothesis presented in Jee & Tyson (2009) suggests that the mass of the southwestern galaxy group is not big enough to be observed as an X-ray overdensity. Motivated by the slight elongation of the X-ray peak towards the south-west, Jee & Tyson (2009) also hypothesise that the two-halo system is being observed after the less massive cluster has passed through the central one. A previous study (Maughan et al. 2004) also showed a region of cooler emission on the west side of the BCG, that is, in the north of the mentioned hot region. This was seen using *Chandra* data, and it was explained as a possible infall of some cooler body. Additionally, from the diffuse radio emission analysis with LOFAR data, Di Gennaro et al. (2020) showed that CL J1226.9+3332 hosts the most distant radio halo discovered to date: a radio emission with a size of 0.7 Mpc that follows the thermal gas distribution. In brief, CL J1226.9+3332 shows evidence of disturbance in the core, but a relaxed morphology at large scales.

### 2.2. The mass of CL J1226.9+3332

Regarding the mass of CL J1226.9+3332, which constitutes the main topic in this study, we present here the results obtained in



**Fig. 1.**  $M_{500}$  estimates for CL J1226.9+3332 in the literature. Filled grey diamonds represent HSE masses from the combination of SZ and X-ray data and empty ones correspond to X-ray-only results. Pink squares show the SZ-only mass assuming virial relation, purple circles are dynamical mass estimates, and brown stars correspond to lensing  $M_{500}$ .

previous works (summarised in Table A.1). These masses have not been homogenised or scaled to the same cosmology, and are the values extracted directly from different analyses. We define  $M_{\Delta}$  as the mass of the cluster inside a sphere of radius  $R_{\Delta}$ ;  $R_{\Delta}$  is the radius at which the mean mass density of the cluster is  $\Delta$  times the critical density of the Universe at its redshift  $\rho_c = 3H(z)^2/(8\pi G)$ , where  $H(z)$  is the *Hubble* function.

The first SZ mass analysis of this cluster was done in Joy et al. (2001) and they estimated  $M(r < 340 h_{100}^{-1} \text{ kpc}) = (3.9 \pm 0.5) \times 10^{14} M_{\odot}$ . Using *Chandra* X-ray data from Cagnoni et al. (2001) and assuming an isothermal  $\beta$ -model, Jee & Tyson (2009) obtained the hydrostatic projected mass  $M(r < 1 \text{ Mpc}) = 1.4^{+0.6}_{-0.4} \times 10^{15} M_{\odot}$ . Also assuming an isothermal  $\beta$ -model and hydrostatic equilibrium, Maughan et al. (2004) obtained  $M_{1000} = 6.1^{+0.9}_{-0.8} \times 10^{14} M_{\odot}$  and  $M_{200} = (1.4 \pm 0.5) \times 10^{15} M_{\odot}$  with *XMM-Newton* data. The subsequent analysis of three-dimensional hydrodynamical properties with *Chandra* and *XMM-Newton* by Maughan et al. (2007), again under the assumptions of spherical symmetry and hydrostatic equilibrium, concluded that  $M_{500} = 5.2^{+1.0}_{-0.8} \times 10^{14} M_{\odot}$ . According to the X-ray analysis in Mantz et al. (2010),  $M_{500} = (7.8 \pm 1.1) \times 10^{14} M_{\odot}$ .

From the combination of Sunyaev-Zel'dovich Array (SZA, Muchovej et al. 2007) interferometric data and the *Chandra* X-ray observations, under the hydrostatic equilibrium hypothesis, Mroczkowski et al. (2009) obtained  $M_{500} = 7.37^{+2.50}_{-1.57} \times 10^{14} M_{\odot}$  and  $M_{2500} = 2.67^{+0.29}_{-0.27} \times 10^{14} M_{\odot}$ . This was compared to the results using only the X-ray data and assuming an isothermal  $\beta$ -model:  $M_{500} = 7.30^{+2.10}_{-1.51} \times 10^{14} M_{\odot}$ ,

$M_{2500} = 2.98^{+0.90}_{-0.63} \times 10^{14} M_{\odot}$ . Using a new approach that instead relies on the virial relation and the Navarro-Frenk-White (NFW, Navarro et al. 1996) density profile, Mroczkowski (2011, 2012) estimated the mass for CL J1226.9+3332 using only SZ data from SZA:  $M_{500} = 6.49^{+0.34}_{-0.34} \times 10^{14} M_{\odot}$  and  $M_{2500} = 2.35^{+0.15}_{-0.16} \times 10^{14} M_{\odot}$  assuming a pressure described by a generalised Navarro-Frenk-White (gNFW, Nagai et al. 2007) profile with  $(a, b, c) = (0.9, 5.0, 0.4)$  parameters and  $M_{500} = 6.42^{+0.36}_{-0.36} \times 10^{14} M_{\odot}$  and  $M_{2500} = 2.53^{+0.14}_{-0.15} \times 10^{14} M_{\odot}$  with  $(a, b, c) = (1.0510, 5.4905, 0.3081)$  as in Arnaud et al. (2010). Some years before, Muchovej et al. (2007) fitted the temperature decrement due to the cluster's SZ effect to the SZA data and assuming hydrostatic equilibrium and isothermality estimated  $M_{200} = 7.19^{+1.33}_{-0.92} \times 10^{14} M_{\odot}$  and  $M_{2500} = 1.68^{+0.37}_{-0.26} \times 10^{14} M_{\odot}$ .

Another approach was considered in Bulbul et al. (2010) to compute the hydrostatic mass, with the polytropic equation of state and using only *Chandra* X-ray observations,  $M_{500} = 4.25^{+1.22}_{-1.14} \times 10^{14} M_{\odot}$  and  $M_{2500} = 2.16^{+0.69}_{-0.63} \times 10^{14} M_{\odot}$ . According to the Planck Collaboration XVIII (2015) results, the hydrostatic mass of the cluster is  $M_{500} = 5.70^{+0.63}_{-0.69} \times 10^{14} M_{\odot}$ . This mass was obtained using the SZ-mass scaling relation given in Eq. (7) of Planck Collaboration XX (2013).

In addition, combining SZ data from NIKA and *Planck* with the X-ray electron density from the *Chandra* ACCEPT data (Cavagnolo et al. 2009), Adam et al. (2015) obtained three hydrostatic mass estimates for different parameters in their gNFW pressure profile modelling:  $M_{500} = 5.96^{+1.02}_{-0.79} \times 10^{14} M_{\odot}$  using  $(a, b, c) = (1.33, 4.13, 0.014)$ ,  $M_{500} = 6.10^{+1.52}_{-1.06} \times 10^{14} M_{\odot}$  with  $(b, c) = (4.13, 0.014)$  and  $M_{500} = 7.30^{+1.52}_{-1.34} \times 10^{14} M_{\odot}$  with  $(a, b, c) = (0.9, 5.0, 0.4)$ . Also combining NIKA and *Chandra* data, Castagna & Andreon (2020) reconstructed  $M_{500} = 5.57^{+1.81}_{-1.23} \times 10^{14} M_{\odot}$ .

The weak-lensing analysis in Jee & Tyson (2009) realised by fitting a NFW density profile found that  $M_{200} = (1.38 \pm 0.20) \times 10^{15} M_{\odot}$ . Similarly, they computed the weak-lensing mass estimate at  $R_{500}$  from Maughan et al. (2007):  $M(r < (0.88 \pm 0.05) \text{ Mpc}) = (7.34 \pm 0.71) \times 10^{14} M_{\odot}$ , and found a 30% higher mass than the X-ray estimate in Maughan et al. (2007). This discrepancy was explained in Jee & Tyson (2009) as a sign of an ongoing merger in the cluster that would create an underestimation of the hydrostatic mass with X-rays without altering the lensing estimate. Jee & Tyson (2009) also estimated the projected mass in each of the two big substructures within  $r < 20''$ ; for the most massive and central clump they found  $M(r < 20'') = (1.3 \pm 0.1) \times 10^{14} M_{\odot}$ , and for the structure at  $\sim 40''$  to the south-west of the BCG  $M(r < 20'') = (8.5 \pm 0.6) \times 10^{13} M_{\odot}$ . Merten et al. (2015) performed a lensing analysis and obtained  $M_{200} = (2.23 \pm 0.14) \times 10^{15} M_{\odot}$ ,  $M_{500} = (1.54 \pm 0.12) \times 10^{15} M_{\odot}$ , and  $M_{2500} = (0.61 \pm 0.10) \times 10^{15} M_{\odot}$  by fitting a NFW density profile to the CLASH data. In addition, based on the weak and strong lensing analysis from Sereno & Covone (2013), Sereno (2015) followed the same procedure as for all clusters in the CoMaLit<sup>1</sup> sample and obtained  $M_{500} = (7.96 \pm 1.44) \times 10^{14} M_{\odot}$ .

Moreover, a recent study based on the velocity dispersion of galaxy members in Aguado-Barahona et al. (2022) obtained two dynamical mass estimates for CL J1226.9+3332:  $M_{500} = (4.7 \pm 1.0) \times 10^{14} M_{\odot}$  and  $M_{500} = (4.8 \pm 1.0) \times 10^{14} M_{\odot}$  from the velocities of 52 and 49 member galaxies, respectively.

We display in Fig. 1 the different  $M_{500}$  estimates found in the literature. Grey diamonds with error bars correspond to the HSE mass estimates. We distinguish the HSE masses obtained

<sup>1</sup> Comparing masses in the literature. Cluster lensing mass catalogue available at <http://pico.oabo.inaf.it/~sereno/CoMaLit/>

from the combination of SZ and X-ray data (filled diamonds) and the X-ray-only results (empty diamonds). The mass given by [Planck Collaboration XVIII \(2015\)](#) is considered here to be an SZ+X result, but it is important to keep in mind that this mass was obtained applying a scaling relation (derived from X-ray data) to the *Planck* SZ signal. The empty pink squares show the  $M_{500}$  assuming the virial relation and using only SZ data ([Mroczkowski 2011, 2012](#)). The purple circles are the dynamical masses from [Aguado-Barahona et al. \(2022\)](#) and the brown stars the lensing estimates from [Merten et al. \(2015\)](#) and [Sereno \(2015\)](#). We decided not to present in the same figure the projected masses, as it would be misleading to compare them to the masses integrated in a sphere. The figure shows that both HSE and lensing masses among them vary more than 40% from one analysis to another.

All these mass estimates for CL J1226.9+3332 are hindered by systematic effects, which are difficult to deal with. Properly comparing masses obtained from different observables, methods, or modelling approaches is crucial, but very challenging. Moreover, as the shape of the HSE mass profile varies depending on the data and the analysis procedure that is considered, the value of  $R_{500}$  is not the same for all estimates presented in Fig. 1. Comparisons are thus delicate due to the correlation between the mass and the radius at which it is estimated. As mentioned, an accurate knowledge of the mass of galaxy clusters is essential for cosmological purposes ([Pratt et al. 2019](#)). This motivates the study in this paper, which follows the previous work of [Ferragamo et al. \(2022\)](#).

### 3. ICM observations

#### 3.1. NIKA2

##### 3.1.1. NIKA2 and the SZ Large Program

The NIKA2 camera ([Adam et al. 2018a; Bourrion et al. 2016; Calvo et al. 2016](#)) is a millimetre camera operating at the IRAM 30-meter telescope on Pico Veleta (Sierra Nevada, Spain). It observes simultaneously in two bands centred at 150 and 260 GHz with three arrays of kinetic inductance detectors (KIDs, [Day et al. 2003; Shu et al. 2018](#)), two operating at 260 GHz and one at 150 GHz. These frequency bands are adapted to detect the characteristic distortion of the tSZ effect, which shows a decrement in the CMB brightness at frequencies lower than  $\sim 217$  GHz and an increment at higher frequencies. As described in the instrument performance analysis review ([Perotto et al. 2020](#)), NIKA2 maps the sky with a  $6.5'$  field of view (FoV) and high angular resolution,  $17.6'$  and  $11.1'$  full width at half maximum at 150 and 260 GHz, respectively. These capabilities, combined with its high sensitivity ( $9 \text{ mJy s}^{1/2}$  at 150 GHz and  $30 \text{ mJy s}^{1/2}$  at 260 GHz), enable us to study the ICM of galaxy clusters in detail, as has been proved in previous works ([Ruppin et al. 2018; Kérúzoré et al. 2020](#)). Substructures in clusters and contaminant sources can also be identified.

Part of the NIKA2 guaranteed time was allocated to the NIKA2 SZ Large Program. This programme is a high angular resolution follow-up of  $\sim 45$  galaxy clusters detected with *Planck* and ACT ([Planck Collaboration XVIII 2015; Hasselfield et al. 2013](#)). These clusters were chosen at intermediate to high redshift ranges ( $0.5 \leq z \leq 0.9$ ) and covering a wide range of masses ( $3 \leq M_{500}/10^{14} M_{\odot} \leq 10$ ) estimated according to their tSZ signal and mass-observable scaling relations. The improvement in resolution with respect to the previous instruments, going below the arcminute scale, allows us to resolve distant clusters for which the apparent angular sizes are small. The SZ Large Pro-

gram also benefits from the X-ray observations obtained with the *XMM-Newton* satellite. From the combination of tSZ and X-ray data, the NIKA2 SZ Large Program will be able to re-estimate precise HSE masses of the galaxy clusters in the sample and improve the current SZ-mass scaling relations. A more detailed explanation of the NIKA2 SZ Large Program is given in [Mayet et al. \(2020\)](#) and [Perotto et al. \(2022\)](#).

##### 3.1.2. NIKA2 observations and maps of CL J1226.9+3332

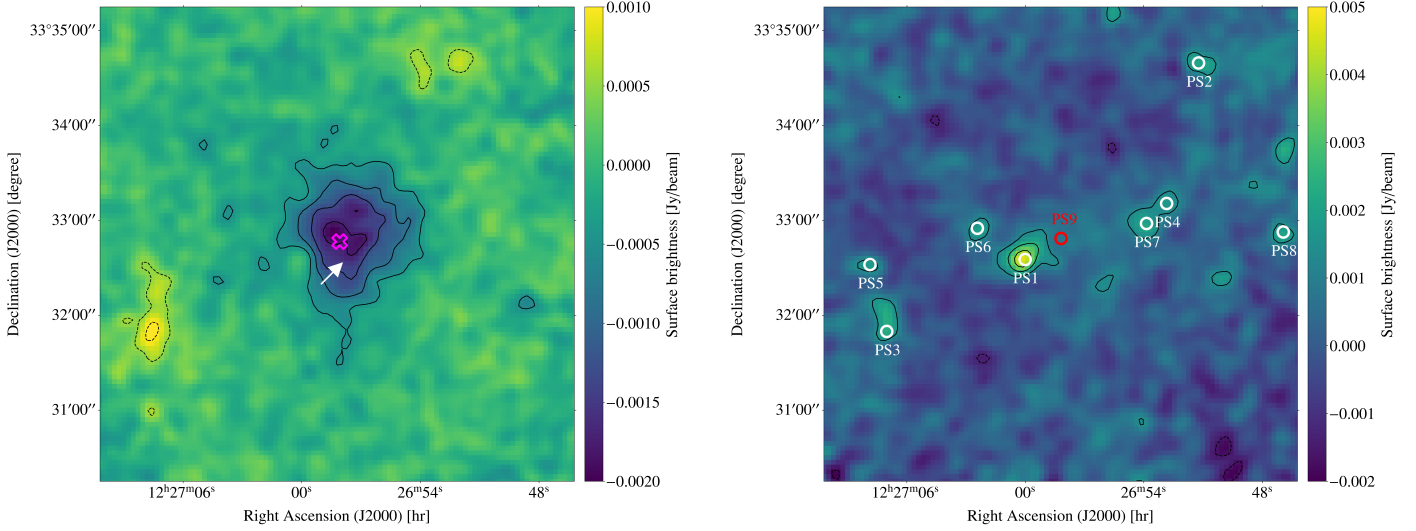
CL J1226.9+3332 was observed for 3.6 h during the 15th NIKA2 science-purpose observation campaign (13–20 February 2018) as part of the NIKA2 Guaranteed Time (under project number 199-16) at the IRAM 30m telescope. The data consists of 36 raster scans of  $8 \times 4$  arcmin in series of four scans with angles of 0, 45, 90, and 135 degrees with respect to the right ascension axis. The scans were centred at the *XMM-Newton* X-ray peak,  $(\text{RA, Dec})_{\text{J2000}} = (12\text{h}26\text{m}58.08\text{s}, +33\text{d}32\text{m}46.6\text{s})$ . The mean elevation of the scans is  $58.51^\circ$ . Data at 150 and 260 GHz were acquired simultaneously.

The raw data were calibrated and reduced following the *baseline* method developed for the performance assessment of NIKA2 ([Perotto et al. 2020](#)) and extended to diffuse emission, as in [Kérúzoré et al. \(2020\)](#) and [Rigby et al. \(2021\)](#). The data for the two frequency bands are processed independently. Common modes of the most-correlated detectors are estimated for data outside a circular mask covering the cluster. These modes are subtracted from the raw data, and a new mask is defined at the positions where the significance of the signal is above a threshold,  $S/N > 3$  in this case. At the next iteration the common modes are estimated outside the new mask and we repeat the procedure until convergence. We chose a 2 arcmin diameter disk centred at the centre of the scans as an initial mask.

We present in Fig. 2 the resulting NIKA2 surface brightness maps at 150 and 260 GHz for CL J1226.9+3332. Black contours indicate significance levels starting from  $3\sigma$  with a  $3\sigma$  spacing. The map at 150 GHz (left panel) shows the cluster as a negative decrement with respect to the background. This is the characteristic signature of the tSZ effect at frequencies lower than 217 GHz. In this map we also identify positive sources that can compensate the negative tSZ signal of the cluster, which is the case for the central south-eastern source. Moreover, in the 150 GHz map we observe an elongation of the tSZ peak towards the south-west. Similar structures were found by [Maughan et al. \(2007\)](#), [Korngut et al. \(2011\)](#), [Adam et al. \(2015\)](#), [Zitrin et al. \(2015\)](#), and [Jee & Tyson \(2009\)](#), as mentioned in Sect. 2. The 260 GHz map (right panel) is dominated by the signal of the point sources and can be used to identify them. Nonetheless, a tSZ signal of the cluster is not detected at this frequency. This is expected since the integrated tSZ signal is about three times weaker in the NIKA2 band centred at 260 GHz than at 150 GHz.

##### 3.1.3. Estimation of noise residuals in the NIKA2 maps

The residual noise in the final 150 GHz NIKA2 map of CL J1226.9+3332 needs to be quantified for the reconstruction of the pressure profile of the cluster. It is usually estimated on null maps, also known as jackknives (JKs), by computing half-differences of two statistically equivalent sets of scans in order to eliminate the astrophysical signal and recover the residuals. We developed two different noise estimates to evaluate possible systematic bias and uncertainties. The angle order (AO) noise map is computed from the half-differences of scans observed with the same angle with respect to the right ascension axis. This

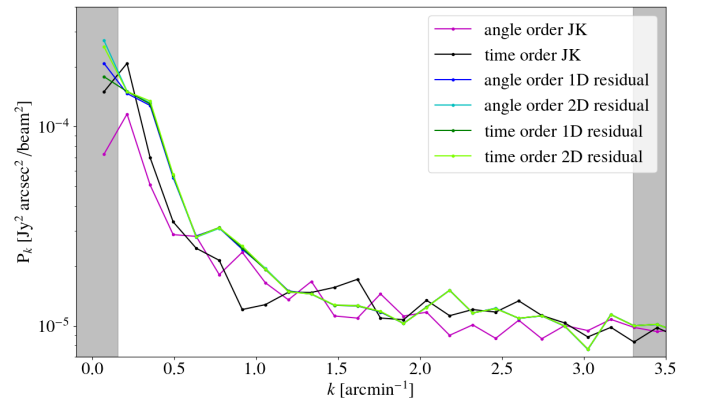


**Fig. 2.** NIKA2 maps of CL J1226.9+3332 at 150 GHz (left) and 260 GHz (right) in  $\text{Jy beam}^{-1}$  units. Contours show S/N levels in multiples of  $\pm 3\sigma$ . Both maps have been smoothed with a  $10'$  FWHM Gaussian kernel. The position of the X-ray centre is shown as a magenta cross in the 150 GHz map and the elongation of the tSZ signal towards the south-west is indicated by the white arrow. White and red circles in the 260 GHz map show the submillimetre and radio point sources, respectively.

ensures that signal residuals from differential filtering along the scan direction are minimised in the null maps. Alternatively, the time order (TO) noise map is calculated from the half-differences of consecutive scans. This minimises the time-dependent effects that may be induced by atmospheric residual fluctuations. We present in Fig. 3 the power spectra of the AO and TO null maps at 150 GHz for CL J1226.9+3332, respectively. The AO null map has a flatter power spectrum for small wave numbers, meaning that it contains less large-scale correlated noise than the TO null map. This suggests that the TO null map might be affected by signal or atmospheric residuals differently filtered for each scanning angle. The power spectra shown in Fig. 3 were used to compute 1000 Monte Carlo noise realisations to estimate the pixel-pixel noise covariance matrices used in Sect. 4.1.1 (following the method developed in Adam et al. 2016).

### 3.1.4. Pipeline transfer function estimation

The filtering induced by the observations and the data reduction process on the cluster signal needs to be evaluated to be accounted for when estimating the pressure profile of the cluster. The transfer function (TF) is calculated by repeating the data processing steps discussed above on a simulation of the cluster signal. It is computed in Fourier space as the ratio of the power spectrum of the simulation filtered by the processing and that of the input simulation. To simulate the cluster tSZ signal we used the universal pressure profile described in Arnaud et al. (2010) with the integrated tSZ signal<sup>2</sup> and redshift of CL J1226.9+3332 given in the Planck Collaboration XVIII (2015) catalogue,  $Y_{500}^{\text{Planck}} \sim 3.82 \times 10^{-4} \text{ arcmin}^2$  and  $z = 0.89$ . We added to the simulated cluster a Gaussian signal with flat spectrum (i.e. random white noise) to explore angular scales at which the cluster signal is negligible. Up to now the NIKA2 SZ Large Program and NIKA analyses (Adam et al. 2015; Ruppin et al. 2018; Kéruzoré et al. 2020) considered one-

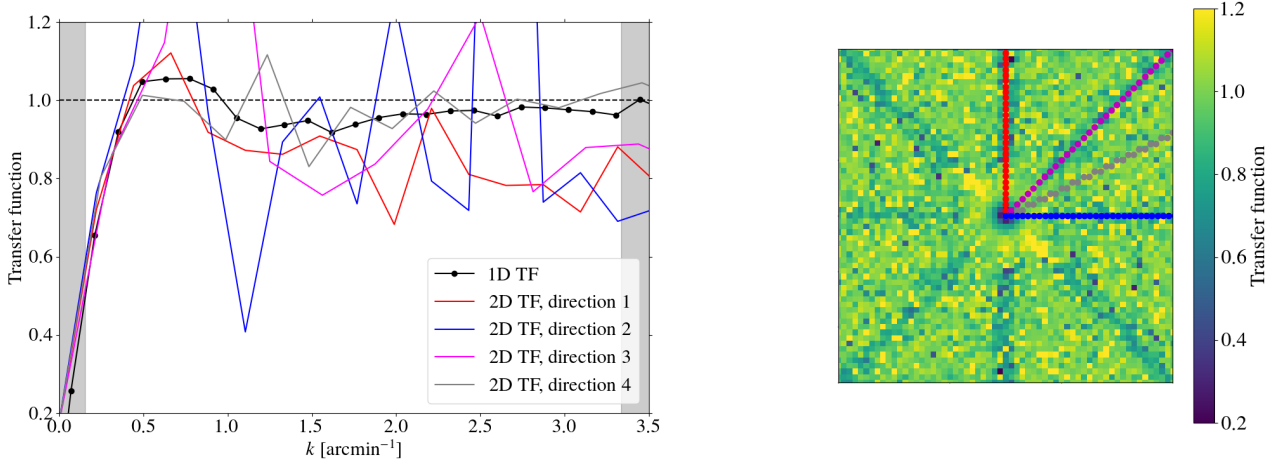


**Fig. 3.** Power spectra of noise map estimates for NIKA2 150 GHz data. The spectra for the JK maps estimated with angle ordered and time ordered scans are in pink and black, respectively. The power spectra of the different residual maps for the best-fit models shown in Fig. 5 are in blue and green. The grey regions represent the NIKA2 instrumental limits given by the field of view (for small angular frequencies) and the beam FWHM (for large angular frequencies).

dimensional transfer functions (1D TF). In these cases, circular symmetry is assumed and the 1D TF is obtained by averaging the power spectra ratio in Fourier-domain annuli at a fixed angular scale. Nevertheless, it is expected that the filtering is not isotropic in the map as it might depend on the scanning direction. This motivates the use of the two-dimensional transfer function (2D TF). In the right panel of Fig. 4 we present the 2D TF describing the filtering in the NIKA2 150 GHz map of Fig. 2. The black line in the left panel of the figure shows the 1D TF, whereas the coloured lines correspond to the one-dimensional cuts of the 2D TF for the different directions represented in the right plot.

Except for the scanning directions, the 2D TF is compatible with the 1D TF, and is greater than 0.8 at large angular scales, meaning that the signal is well preserved. On the contrary, filtering is significant for angular frequencies below  $\sim 0.5 \text{ arcmin}^{-1}$ . At  $0.4 \text{ arcmin}^{-1} \lesssim k \lesssim 0.8 \text{ arcmin}^{-1}$  the transfer

<sup>2</sup>  $Y_{500}$  is the integrated tSZ signal within a sphere of radius  $R_{500}$  with  $\mathcal{D}_A$  the angular diameter distance at the cluster redshift. It is defined as  $Y_{500} = 4\pi \frac{\sigma_T}{m_e c^2 \mathcal{D}_A} \int_0^{R_{500}} r^2 P_c(r) dr$ .



**Fig. 4.** Transfer functions, 1D (left) and 2D (right), describing the filtering induced by data processing for the 150 GHz map in Fig. 2. The coloured lines in the left panel represent the values of the 2D transfer function for the directions shown with the same colours in the right panel. Grey-shaded areas correspond to the NIKA2 field of view (for small angular frequencies) and beam FWHM (for large angular frequencies) instrumental limits.

function is larger than unity, meaning that the signal has been slightly enhanced by the data analysis process at these scales. In order to evaluate the impact of considering the anisotropy of the filtering, the analyses presented in the following sections will be carried out with both the 1D and 2D transfer functions (see Sect. 4.1.1).

### 3.1.5. Point source contamination

Point sources in the 150 GHz map contaminate the tSZ signal and also need to be considered. We start by identifying submillimetre sources by blindly searching for point sources in the NIKA2 260 GHz map of CL J1226.9+3332. By cross-checking the detections with a S/N greater than 3 with *Herschel* SPIRE and PACS catalogues, seven submillimetre sources were identified in the region covered by the NIKA2 maps. The position and fluxes from the above-mentioned catalogues for each submillimetre point source (PS1 to PS5, PS7, and PS8) are summarised in Table 1, and the corresponding *Herschel* names are given in Table 2.

For each of these sources, a modified black-body spectrum model is adjusted to the fluxes in Table 1 together with the measurement of the flux at 260 GHz from the NIKA2 map. The spectra are then extrapolated to 150 GHz to obtain an estimate of the flux of each source at 150 GHz. In Table 2 we summarise the fluxes at 260 GHz obtained from the NIKA2 map and the extrapolated values at 150 GHz. The contribution of the noise and the filtering in the 260 GHz map are considered when measuring the fluxes. The estimates obtained with the AO and TO noise approaches give compatible results for all point sources. A more detailed explanation of the method used to deal with submillimetre point sources is given in Kérusoré et al. (2020). The obtained probability distributions of the fluxes at 150 GHz are used as priors for the joint fit of the cluster pressure profile and the point sources fluxes described in Sect. 4.1.1.

Comparing results in Table 2, we find very large uncertainties for the extrapolated flux of PS1 at 150 GHz when estimated with the TO noise map, but this will not substantially affect the following results for the simultaneous fit of the PS1 flux (see Table B.1) when estimating the cluster pressure profile. Moreover, we can compare its flux at 260 GHz to the measurement in Adam et al. (2015), where the source is called PS260.

Using NIKA data, they obtain  $F_{260\text{ GHz}} = 6.8 \pm 0.7$  (stat.)  $\pm 1.0$  (cal.) mJy, which is consistent with our estimates.

Source PS6 does not have a counterpart in the *Herschel* SPIRE and PACS catalogues, but it appears as a weak signal in the *Herschel* maps and as a  $3\sigma$  detection in the 260 GHz NIKA2 map; moreover, it compensates the extended tSZ signal at 150 GHz (also clearly observed in Adam et al. 2015). For this source the modified black body is used to obtain prior knowledge of the flux at 150 GHz for the assumed prior distributions of the spectral index and temperature (Kérusoré et al. 2020). Another tricky point source is PS7. The extrapolated 150 GHz values shown in Table 2 clearly overestimate the flux of the source. This is understandable since we do not have enough constraints for the low-frequency slope of the spectral energy distribution. We choose to use the obtained values as upper limits of a flat prior for the flux of PS7 in the estimation of the cluster pressure profile.

In addition to submillimetre sources, according to the VLA FIRST Survey catalogue (White et al. 1997), a radio source of  $3.60 \pm 0.13$  mJy at 1.4 GHz is present in (RA, Dec)<sub>J2000</sub> = (12h26m58.19s, +33d32m48.61s), hereafter PS9. This galaxy corresponds to the BCG identified in Holden et al. (2009) and the compact radio source detected with LOFAR in Di Gennaro et al. (2020). We know beforehand that the contribution of this radio source at 150 GHz is small, but given its central position, it is important to consider it. Assuming a synchrotron spectrum  $F(\nu) = F_0(\nu/\nu_0)^\alpha$  with  $\alpha = -0.7 \pm 0.2$ , which describes the spectral energy distribution for an average radio source (Condon 1984), at 150 GHz we obtain  $0.1 \pm 0.2$  mJy. The obtained probability distribution of the flux at 150 GHz is also used as a prior for the fit in Sect. 4.1.1. The extrapolation of fluxes from radio to millimetre wavelengths can be dangerous and lead to biasing the electron pressure reconstruction. However, this is not the case for our analyses (see results in Sect. 4.1.1 and Table B.2).

### 3.2. XMM-Newton

Regarding the X-ray data, CL J1226.9+3332 was observed by *XMM-Newton* (Obs ID 0200340101) for a total observation time of 90/74 ks (MOS/pn), reducing to 63/47 ks after cleaning. The data were reduced following the standard procedures described

**Table 1.** Submillimetre point source coordinates and fluxes identified within a radius of 2' around the centre of CL J1226.9+3332.

Source	Coordinates J2000	600 GHz <sup>(a)</sup> [mJy]	860 GHz <sup>(a)</sup> [mJy]	1200 GHz <sup>(a)</sup> [mJy]	1870 GHz <sup>(b)</sup> [mJy]	3000 GHz <sup>(b)</sup> [mJy]
PS1	12h27m00.01s +33d32m35.29s	100.3 ± 10.0	121.2 ± 10.0	109.8 ± 7.6	55.7 ± 6.0	14.6 ± 2.1
PS2	12h26m51.22s +33d34m39.61s	37.8 ± 9.0	46.4 ± 9.9	29.1 ± 7.5	24.9 ± 7.4	8.0 ± 1.7
PS3	12h27m07.02s +33d31m49.79s	34.8 ± 7.9	32.4 ± 8.7	25.6 ± 7.7	31.1 ± 6.6	25.6 ± 2.5
PS4	12h26m52.84s +33d33m10.74s		33.0 ± 10.3	41.9 ± 9.7	31.5 ± 7.0	17.7 ± 1.8
PS5	12h27m07.87s +33d32m32.08s			30.0 ± 9.4		
PS6	12h27m02.43s +33d32m55.06s					
PS7	12h26m53.86s +33d32m58.10s				21.8 ± 1.5	14.4 ± 3.0
PS8	12h26m46.93s +33d32m52.66s			19.8 ± 5.5		

**Notes.** <sup>(a)</sup>Fluxes at 600, 860, and 1200 GHz are obtained from the SPIRE catalogue. <sup>(b)</sup>Fluxes at 1870 and 3000 GHz are given in the PACS catalogue. <sup>(a)</sup>European Space Agency, *Herschel* SPIRE Point Source Catalogue, Version 1.0, 2007, <https://doi.org/10.5270/esa-6gfkpzh> <sup>(b)</sup>European Space Agency, *Herschel* PACS Point Source Catalogue, Version 1.0, 2007, <https://doi.org/10.5270/esa-rw7rbo7>

**Table 2.** Submillimetre point sources fluxes and their corresponding names in *Herschel* SPIRE or PACS catalogue.

Source	<i>Herschel</i> name	150 GHz [extrap. AO] [mJy]	150 GHz [extrap. TO] [mJy]	260 GHz [AO] [mJy]	260 GHz [TO] [mJy]
PS1	HSPSC250A_J1227.00+3332.5	2.8 ± 0.3	2.7 ± 2.3	8.2 ± 0.5	8.1 ± 0.5
PS2	HSPSC250A_J1226.86+3334.7	1.1 ± 0.3	1.0 ± 0.3	3.7 ± 0.6	3.6 ± 0.6
PS3	HSPSC250A_J1227.12+3331.9	1.4 ± 0.4	1.4 ± 0.4	3.9 ± 0.6	3.8 ± 0.6
PS4	HSPSC250A_J1226.85+3333.2	0.5 ± 0.2	0.4 ± 0.2	2.6 ± 0.5	2.5 ± 0.5
PS5	HSPSC250A_J1227.13+3332.4	0.6 ± 0.2	0.5 ± 0.2	3.5 ± 0.6	3.4 ± 0.6
PS6		0.4 ± 2.0	0.4 ± 2.0	3.1 ± 0.5	2.9 ± 0.5
PS7	HPPSC160A_J122654.1+333253	1.2 ± 0.2 <sup>(a)</sup>	1.2 ± 0.2 <sup>(a)</sup>	3.2 ± 0.5	3.1 ± 0.5
PS8	HSPSC250A_J1226.78+3332.8	0.5 ± 0.2	0.5 ± 0.2	3.4 ± 0.6	3.3 ± 0.6

**Notes.** Fluxes at 260 GHz are estimated from the NIKA2 maps and fluxes at 150 GHz obtained from the extrapolation of the fitted spectral energy distributions. <sup>(a)</sup>For PS7 the extrapolated 150 GHz fluxes are too high, and we take these values as upper limits of flat priors for the fit in Sect. 4.1.1.

in Bartalucci et al. (2017). The raw data were processed using the *XMM-Newton* standard pipeline Science Analysis System (SAS version 16). Only standard events from the EMOS1, 2, and EPN detectors with PATTERN <4 and <13, respectively, were kept. The data were further filtered for badtime events and solar flares. Vignetting was accounted for following the weighting scheme described in Arnaud et al. (2001). Point sources were detected on the basis of wavelet-filtered images in low-energy ([0.3–2] keV) and high-energy ([2–5] keV) bands, then subsequently masked from the events list. This process was controlled by a visual check, allowing us to further extract obvious (sub-)structures present in the field. The instrumental background was modelled through the use of stacked filter-wheel closed observations, whilst the astrophysical contamination due to the Galaxy and the cosmic X-ray background were accounted for as a constant background in the 1D surface brightness analysis. This component was modelled in the spectral analysis following the method outlined in Pratt et al. (2010), using an annulus external to the target of radii  $300'' < \theta < 480''$ .

## 4. ICM thermodynamical profiles

### 4.1. Electron pressure reconstruction from tSZ

The spectral distortion of the CMB caused by the thermal energy in the cluster (i.e. the tSZ effect) is characterised by its amplitude or Compton parameter,  $y$  (Sunyaev & Zeldovich 1972). This is directly proportional to the thermal pressure of the electrons in

the ICM,  $P_e$ , integrated along the line of sight,

$$y = \frac{\sigma_T}{m_e c^2} \int P_e dl, \quad (1)$$

where  $\sigma_T$ ,  $m_e$ , and  $c$  are the Thomson cross-section, the electron rest mass, and the speed of light, respectively. Hence, the tSZ surface brightness is proportional to the Compton parameter integrated over the tSZ spectrum convolved by the NIKA2 bandpass, and therefore proportional to the integrated thermal pressure of the ICM in the cluster.

#### 4.1.1. Pressure profile reconstruction with NIKA2

##### Reconstruction procedure

To reconstruct the electron pressure in the ICM of CL J1226.9+3332 we fit a model map of the surface brightness of the cluster to the NIKA2 150 GHz map.

The model map is obtained from the pressure profile of the galaxy cluster integrated along the line of sight in Compton parameter ( $y$ ) units, following Eq. (1). We describe the pressure of the galaxy cluster with a radially binned spherical model (also called the non-parametric model in Ruppin et al. 2017; Romero et al. 2018),

$$P_e(r_i < r < r_{i+1}) = P_i \left( \frac{r}{r_i} \right)^{-\alpha_i}, \quad (2)$$

where  $P_i$  and  $\alpha_i$  are the values of the pressure and the slope at the radial bin  $r_i$ . The slope is directly calculated as



$$\alpha_i = -\frac{\log P_{i+1} - \log P_i}{\log r_{i+1} - \log r_i}. \quad (3)$$

We initialise the pressure bin values by taking random values from a normal distribution centred at the corresponding pressure from the universal profile of [Arnaud et al. \(2010\)](#) at each radial bin. The radial bins are chosen to cover mainly the range between the NIKA2 resolution and field of view capabilities. We centre the pressure profile at the coordinates of the X-ray peak, as determined from *XMM-Newton* data analysis (Sects. 3.1.2 and 3.2). The derived  $y$ -map is convolved with the NIKA2 beam, which is approximated by a two-dimensional Gaussian with  $FWHM = 17.6''$  ([Perotto et al. 2020](#)). In order to account for the attenuation or filtering effects due to data processing in the NIKA2 150 GHz map, the model map is also convolved with the transfer function (Sect. 3.1.4). We repeat this procedure for the 1D and 2D transfer functions. Finally, the  $y$ -map is converted into surface brightness units with a conversion coefficient, accounting for the tSZ spectra shape convolved by the NIKA2 bandpass, which is also left as a parameter of the fit (as done in [Kéruzoré et al. 2020](#)).

Furthermore, for the comparison with the 150 GHz NIKA2 map, we added the contribution of point sources to the model map. Point sources are modelled as two-dimensional Gaussian functions, and we repeat the procedure detailed in [Kéruzoré et al. \(2020\)](#) to fit the flux of each source at 150 GHz. Priors on the flux of the sources at 150 GHz are obtained from the results of the spectral fits presented in Sect. 3.1.5. The last component in the model map is a constant zero-level that we also adjust as a nuisance parameter.

The parameters ( $\vartheta$ ) of our fit are the pressure bins describing the ICM of the cluster, the fluxes of the contaminant point sources, the conversion factor from Compton to surface brightness units, and the zero-level. The likelihood that we use to compare our model  $\mathcal{M}$  pixel by pixel to the data  $\mathcal{D}$  is given by

$$\log \mathcal{L}(\vartheta) = -\frac{1}{2} \sum_{i=1}^{n_{\text{pixels}}} [(\mathcal{M}(\vartheta) - \mathcal{D})^T C_{\text{pix-pix}}^{-1} (\mathcal{M}(\vartheta) - \mathcal{D})]_i - \frac{1}{2} \left( \frac{Y_{500}(\vartheta) - Y_{500}^{\text{Planck}}}{\Delta Y_{500}^{\text{Planck}}} \right)^2. \quad (4)$$

Here  $C_{\text{pix-pix}}$  is the pixel-pixel noise covariance matrix accounting for the residual noise in the NIKA2 150 GHz map (Sect. 3.1.3). We repeat the fit with the covariance matrices from both AO and TO noise estimates.

We also compute the  $Y_{500}$  integrated Compton parameter and compare it in the likelihood to the integrated Compton parameter measured by [Planck Collaboration XVIII \(2015\)](#),  $Y_{500}^{\text{Planck}} = (3.82 \pm 0.79) \times 10^{-4}$  arcmin<sup>2</sup> within an aperture of  $\theta_{500} = 1.907$  arcmin. We do not compare the integrated Compton parameter at  $5\theta_{500}$  as measured by *Planck* because it would require extrapolating the pressure profile far beyond the NIKA2 data.

For the map fit we use the PANCO2 pipeline ([Kéruzoré et al. 2022](#)) and follow the procedure described in [Adam et al. \(2015\)](#), [Ruppin et al. \(2018\)](#), and [Kéruzoré et al. \(2020\)](#). This pipeline performs a Markov chain Monte Carlo (MCMC) fit using the emcee Python package ([Foreman-Mackey et al. 2019](#); [Goodman & Weare 2010](#)). The sampling is performed using 40 walkers and  $10^5$  steps, with a burn-in of  $10^3$  samples, and convergence is monitored following the  $\hat{R}$  test of [Gelman & Rubin \(1992\)](#) and chains autocorrelation. The PANCO2 code has been successfully tested on simulations.

## NIKA2 pressure radial profile

In order to estimate the robustness of the results of the above procedure, we performed the fit to the NIKA2 data in four different cases with respect to the choice of noise residuals and transfer function estimates. Thus, we consider AO1D (TO1D) and AO2D (TO2D) using the AO (TO) noise residual map and the 1D and 2D transfer functions, respectively. In Fig. 5 we compare the NIKA2 150 GHz map of CL J1226.9+3332 to the obtained best-fit models and their residuals for these four analyses. Comparing the power spectra of the residual maps to the power spectra of the noise estimate maps, we see in Fig. 3 that for the TO case the fit residuals and the noise estimates power spectra are consistent. However, for the AO cases there is an excess of power in the fit residuals, which could be interpreted as coming from the signal due to the differential filtering effects that are not captured in the AO noise, as discussed in Sect. 3.1.3. Regarding point sources, the reconstructed fluxes are consistent for the four analyses (see Tables B.1 and B.2).

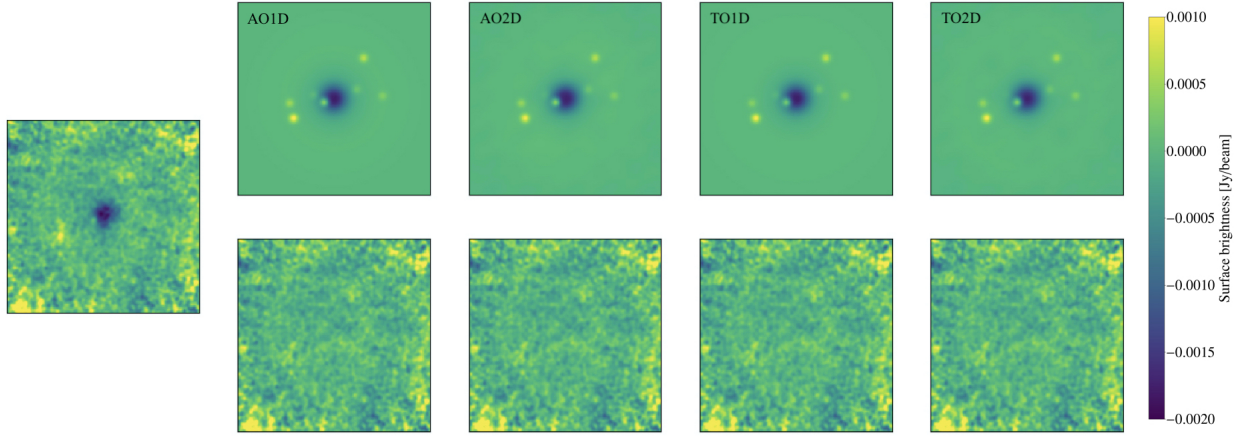
We present in Fig. 6 the radially binned best-fit pressure profiles obtained for the four tested cases. The blue and cyan (dark and light green) dots correspond to the AO (TO) 1D and 2D transfer function estimates, respectively. The plotted uncertainties correspond to  $1\sigma$  of the posterior distributions derived from the MCMC chains. Overall, the four NIKA2 analyses give consistent results, especially in the radial ranges where we expect the NIKA2 results to be reliable (i.e. between the beam and the FoV scales, both represented with dashed vertical lines in the figure). We give the HWHM of the NIKA2 beam ( $17.6''/2$ ) and half the diameter of the FoV ( $6.5''/2$ ) in the physical distances corresponding to the redshift of the cluster.

In terms of noise estimates we observe that the uncertainties on the pressure bin estimates are slightly larger for the time-ordered case, as expected. However, we note no significant bias between the time and angle ordered results. The effect of the transfer function is hard to evaluate: even if the 2D TF is a more precise description of the filtering in the map, when fitting a spherical cluster model the use of the 1D TF gives consistent results. In the following we use the results for the four analyses to evaluate possible systematic uncertainties induced by the NIKA2 processing.

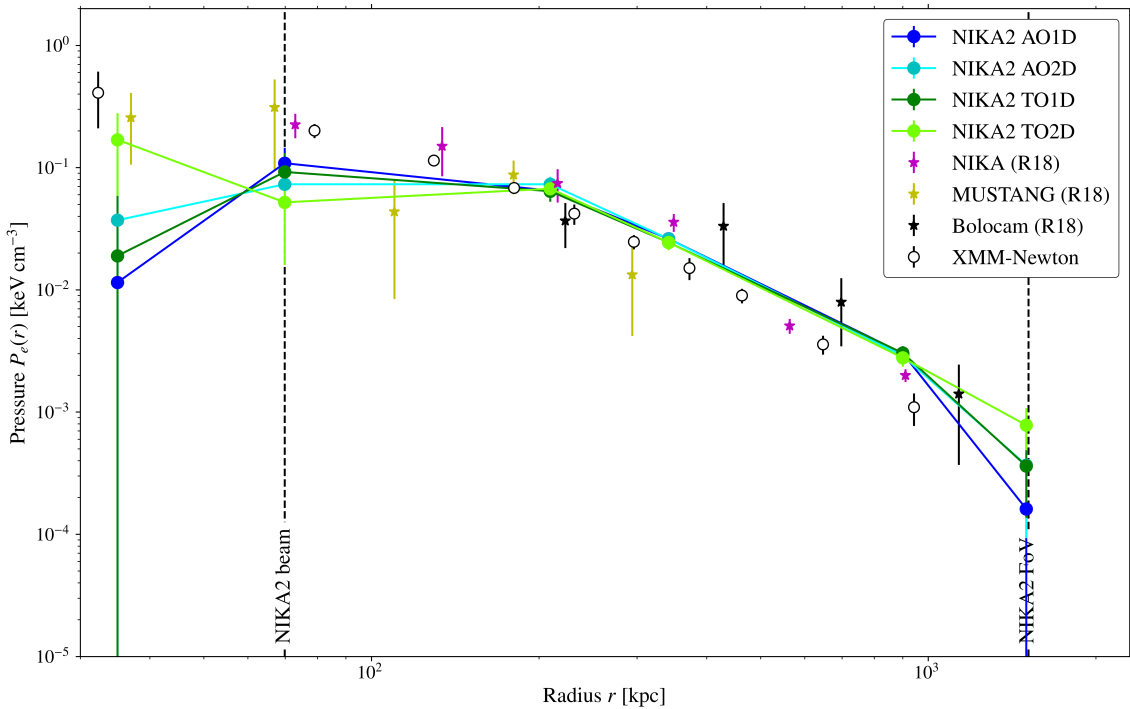
### 4.1.2. Comparison to previous results

In Fig. 6 we compare our results to the profiles obtained in R18 with tSZ data from NIKA (pink), Bolocam (black), and MUSTANG (yellow). MUSTANG's high angular resolution ( $9''$  FWHM at 90 GHz) enables us to map the core of the cluster, whereas Bolocam's large field of view ( $8'$  at 140 GHz) allows us to recover the large angular scales. NIKA and the improved NIKA2 camera are able to cover all the intermediate radii. The consistency of the different pressure bins in the radial range from the NIKA2 beam to the FoV proves the reliability of the reconstruction with NIKA2 data.

Before going further, we have to consider again the effect of the filtering on the NIKA2 data. The filtering due to the data processing affects mainly small angular frequencies, i.e. small  $k$  numbers (Sect. 3.1.4), which is translated into large angular scales in real space. In this case it means that the region at  $\sim 1000$  kpc from the centre of the cluster is strongly filtered. For this reason, we cast doubt on the results of our fits for the last NIKA2 bin in pressure.



**Fig. 5.** 150 GHz maps of CL J1226.9+3332. Left: NIKA2 150 GHz surface brightness map of CL J1226.9+3332. Top: best-fit models of the tSZ signal and point sources. Bottom: residual maps, difference between the data map and each best-fit model. Results obtained with different transfer function and noise estimates (from left to right): AO1D, AO2D, TO1D, and TO2D. All maps have been smoothed with a  $10'$  Gaussian kernel for display purposes and are shown in units of  $\text{Jy beam}^{-1}$ .



**Fig. 6.** Pressure profile of the ICM of CL J1226.9+3332. Blue and green symbols correspond to the results obtained in this work from the NIKA2 150 GHz map. The error bar edges represent the  $1\sigma$  uncertainties. Pink, yellow, and black stars show the reconstructed profiles in R18 for NIKA, MUSTANG, and Bolocam data, respectively. Empty symbols correspond to the pressure profile obtained from the combination of *XMM-Newton* electron density and temperature profiles. Vertical dashed lines indicate the instrumental limits of NIKA2 as radius of the beam and FoV.

#### 4.2. Thermodynamical profiles from X-rays

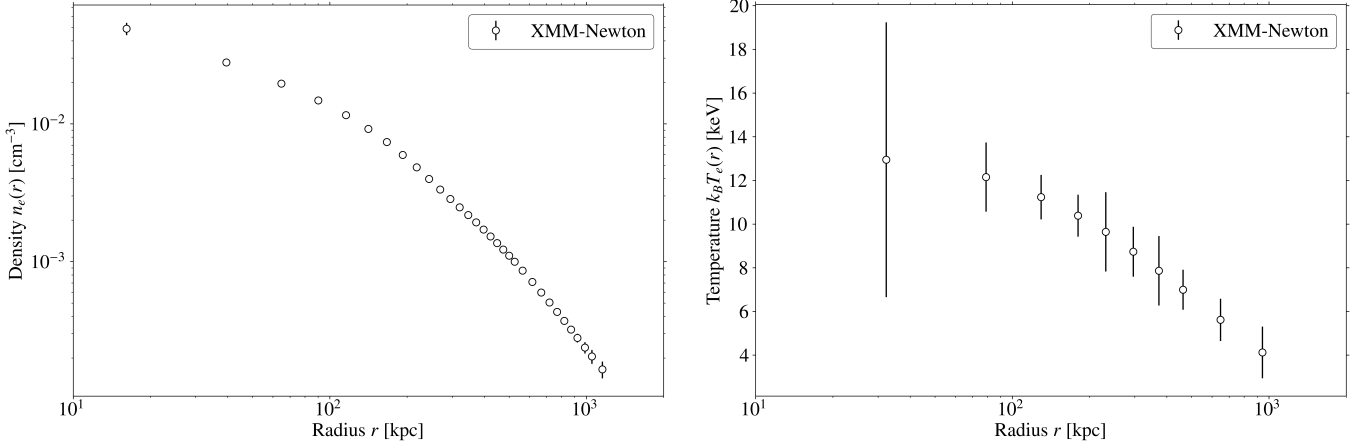
The electron density and temperature profiles were extracted following the methodology described by Pratt et al. (2010) and Bartalucci et al. (2017). In short, the vignettted-corrected and background-subtracted surface brightness profile obtained in concentric annuli from the X-ray peak is deconvolved from the point spread function (PSF) and geometrically deprojected assuming spherical symmetry using the regularisation technique described in Croston et al. (2006).

The temperature profile is derived in bins defined from the previously derived binned surface brightness profile through a spectral analysis modelling the ICM emission via an absorbed

MEKAL model under XSPEC<sup>3</sup> and accounting for both the instrumental and astrophysical backgrounds. The derived 2D temperature profile is then PSF-corrected and deprojected following the ‘non-parametric-like’ method presented in Bartalucci et al. (2018).

The gas pressure,  $P$ , and entropy,  $K$ , profiles were then derived from the deprojected density,  $n_e$ , and temperature,  $T$ , profiles assuming  $P = n_e \times T$  and  $K = T/n_e^{2/3}$ , respectively. In this paper we focus on the gas pressure profile, which is shown

<sup>3</sup> <https://heasarc.gsfc.nasa.gov/xanadu/xspec/>



**Fig. 7.** Electron density (left) and temperature (right) profiles reconstructed from *XMM-Newton* observations, with  $1\sigma$  error bars. The profiles are centred at the X-ray peak (RA, Dec.)<sub>J2000</sub> = (12h26m58.08s, +33d32m46.6s).

in Fig. 6 with open black circles. The electron density and temperature profiles are shown in Fig. 7.

## 5. Hydrostatic mass

### 5.1. Hydrostatic equilibrium

Under the hydrostatic equilibrium (HSE) hypothesis, for a spherical cluster we can compute the total cluster mass enclosed within the radius  $r$  as

$$M_{\text{HSE}}(< r) = -\frac{1}{\mu m_p G} \frac{r^2}{n_e(r)} \frac{dP_e(r)}{dr}, \quad (5)$$

where  $\mu$ ,  $m_p$ , and  $G$  are the mean molecular weight of the ICM gas, the proton mass and the gravitational constant, respectively. We assume  $\mu \approx 0.6$  (Pessah & Chakraborty 2013; Ettori et al. 2019) for the gas. Combining the pressure profiles obtained from the thermal SZ or X-ray data with the electron density from the X-ray, we can reconstruct the mass of the galaxy cluster as in Adam et al. (2015), Ruppin et al. (2018), and Kéruszoré et al. (2020).

### 5.2. Pressure profile modelling for mass estimation

Deriving the mass directly from the radially binned profiles presented above (see Fig. 6) leads to non-physical results (i.e. negative mass contributions) as no condition was imposed regarding the slope of the profile in the pressure reconstruction. This was done to prevent extra constraints on the pressure profile induced by assumptions on the model. To overcome this issue, we fit here pressure models ensuring physical mass profiles to the radially binned tSZ results from Sect. 4.1.1. We consider two different approaches: 1) a direct fit of a generalised Navarro-Frenk-White (gNFW) profile to the radially binned pressure, and 2) an indirect fit of a Navarro-Frenk-White (NFW) mass density model under the HSE assumption. Aiming for a precise reconstruction of the HSE mass, which requires having accurately constrained slopes for the pressure profile, in both cases we combine the NIKA2 pressure bins with the results obtained in R18<sup>4</sup>.

<sup>4</sup> The binned profiles in R18 and those in this work are centred on positions at 3 arcsec of distance, which is the typical rms pointing error for NIKA2 (Perotto et al. 2020), so we consider that combining them is a valid approach.

#### 5.2.1. gNFW pressure model

The first approach, which was used in previous NIKA2 studies (Ruppin et al. 2018; Kéruszoré et al. 2020; Ferragamo et al. 2022), consists in fitting the widely used gNFW pressure profile model (Nagai et al. 2007) to the tSZ data,

$$P_e(r) = \frac{P_0}{\left(\frac{r}{r_p}\right)^c \left(1 + \left(\frac{r}{r_p}\right)^a\right)^{(b-c)/a}}, \quad (6)$$

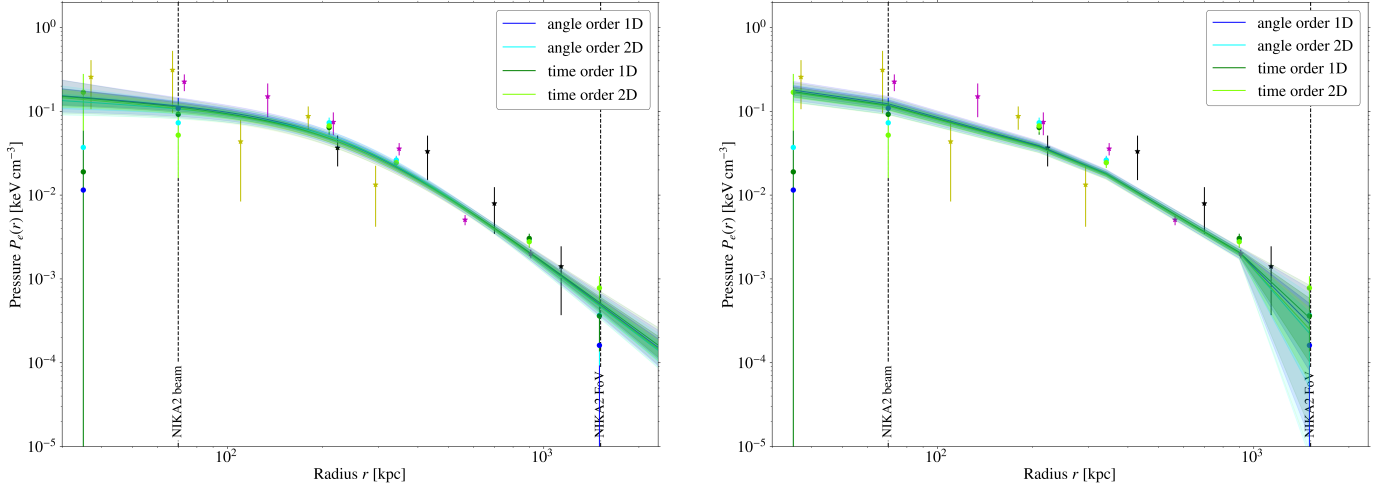
where  $P_0$  is the normalisation constant,  $b$  and  $c$  respectively the external and internal slopes,  $r_p$  the characteristic radius of slope change, and  $a$  the parameter describing the steepness of the slope's transition.

We perform a MCMC fit using *emcee* package. The likelihood function is given by

$$\begin{aligned} \log \mathcal{L}(\vartheta) = & -\frac{1}{2} \left( P^{\text{gNFW}}(\vartheta) - P^{N2} \right)^T C^{-1} \left( P^{\text{gNFW}}(\vartheta) - P^{N2} \right) \\ & - \frac{1}{2} \sum_{k=1}^{n_{\text{R18bins}}} \left( \frac{P_k^{\text{gNFW}}(\vartheta) - P_k^{\text{R18}}}{\Delta P_k^{\text{R18}}} \right)^2 \\ & - \frac{1}{2} \left( \frac{Y_{500}^{\text{gNFW}}(\vartheta) - Y_{500}^{\text{Planck}}}{\Delta Y_{500}^{\text{Planck}}} \right)^2, \end{aligned} \quad (7)$$

where  $P^{N2}$  and  $C$  represent the NIKA2 radially binned pressure profile bins and associated covariance matrix,  $P_k^{\text{R18}}$  and  $\Delta P_k^{\text{R18}}$  are the R18 pressure profile data and uncertainties, and  $P^{\text{gNFW}}(\vartheta)$  are the gNFW pressure profile values for a set of parameters  $\vartheta = [P_0, r_p, a, b, c]$ . As we do not rely on the value of the last NIKA2 pressure bin, we chose to modify the NIKA2 inverse covariance matrix  $C^{-1}$  by setting the last diagonal term to  $[C^{-1}]_{6,6} = 0$ , so that the correlation of the last bin with the others is taken into account, but not its value. We also set a constraint on the integrated Compton parameter of the model  $Y_{500}^{\text{gNFW}}(\vartheta)$ , again using the *Planck* satellite PSZ2 catalogue results,  $Y_{500}^{\text{Planck}}$  (Planck Collaboration XVIII 2015). However, we find that the impact of this constraint is completely negligible for this cluster. Furthermore, we added an extra condition to the fit to ensure increasing HSE mass profiles with radius:  $\frac{r^2}{n_e(r)} \frac{dP_e(r)}{dr} < 0$ .

The best-fit gNFW pressure profiles (solid lines) and uncertainties (shaded area) are presented in the left panel of Fig. 8 for



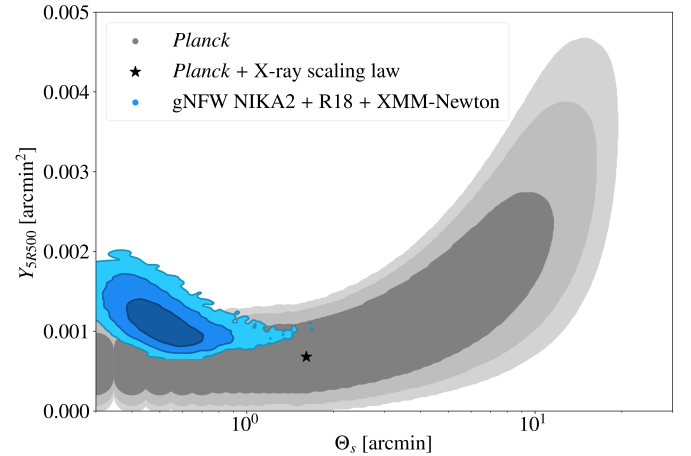
**Fig. 8.** Pressure profile and best fit for the gNFW (left) and NFW (right) models. The data points correspond to the NIKA2 radially binned results for the four data sets discussed above, and to the NIKA, MUSTANG, and Bolocam bins from R18. Blue and green solid lines represent the best-fit values for the four NIKA2 pressure estimates considered. The shaded regions show the 2.5th, 16th, 84th, and 97.5th percentiles.

the four sets of NIKA2 data discussed in Sect. 4.1.1. We observe that the best-fit models are a good representation of the data over the full range in radius, as demonstrated by the corresponding reduced  $\chi^2$ , which are close to 1 for all the cases (see solid lines in Fig. C.2). The posterior distributions of the  $\vartheta_{\text{gNFW}}$  parameters can be found in Appendix C.

In addition, it is interesting to compare our results to those from *Planck* for which a similar modelling was used. In Fig. 9 we present the 2D posterior distributions of the integrated Compton parameter at  $5R_{500}$  (with  $R_{500}$  calculated independently in each case) with respect to the  $\Theta_s$  parameter of the gNFW model, at a confidence level (C.L.) of 68%, 95%, and 99%. The parameters  $\Theta_s$  and  $r_p$  are related via the angular diameter distance at the cluster redshift:  $\tan(\Theta_s) = r_p/D_A$ . We compare the results obtained in *Planck* Collaboration XVIII (2015; with the MMF3 matched multi-filter, available in the *Planck* Legacy Archive<sup>5</sup>) to the constraints from the gNFW profiles obtained in this work with the NIKA2, R18, and *XMM-Newton* data. Our contours were obtained varying all the parameters in the gNFW model fit, while for *Planck*  $a$ ,  $b$ , and  $c$  the parameters were fixed. For simplicity, we only show the contours for the NIKA2 AO1D case. This figure illustrates the important gain in precision due to high-resolution observations: resolving the galaxy cluster allows us to determine, even at such high redshift, the  $\Theta_s$  characteristic radius. The contours are marginally in agreement.

### 5.2.2. NFW density model

We present here a different approach for the modelling of the pressure profile in the scope of estimating the HSE mass. The estimation of the pressure derivative (Eq. (5)) can be very problematic, first because it is very sensitive to local variations in the slope of the pressure profile and because it requires, as discussed above, additional constraints to ensure recovering physical masses. To overcome these issues we model the pressure profile starting from a mass density model and assuming HSE. An equivalent idea is the ‘backward process’ of fitting X-ray temperatures described in Ettori et al. (2013) and references therein. This method was used for the mass reconstruction in Ettori et al. (2019) and in Eckert et al. (2022). From the HSE



**Fig. 9.** Distribution of  $Y_{5R500}$  with respect to  $\Theta_s$  for the gNFW pressure model fits to *Planck* data (grey) in *Planck* Collaboration XVIII (2015) and to the NIKA2 + R18 + *XMM-Newton* data (blue) in this work. Different contours show the 68%, 95%, and 99% confidence intervals. The black star corresponds to the intersection between the *Planck* Collaboration XVIII (2015) distribution and the X-ray scaling law shown in Fig. 16 in *Planck* Collaboration XVIII (2015).

defined in Eq. (5), we can write

$$P(r_b) - P(r_a) = \int_{r_a}^{r_b} -\mu m_p G n_c(r) \frac{M_{\text{HSE}}(< r)}{r^2} dr. \quad (8)$$

Moreover, we can relate a radial mass density profile  $\rho(R)$  to the mass by

$$M(< r) = \int_0^r 4\pi R^2 \rho(R) dR, \quad (9)$$

which allows us to relate the pressure directly to a mass density profile. We use here the NFW model, which is a good description of dark matter halos (Navarro et al. 1996) and has been widely used in the literature (e.g., Ettori et al. 2019),

$$\rho_{\text{NFW}}(R) = \frac{\rho_c \delta_{c200} (c_{200})}{R/r_s (1 + R/r_s)^2}, \quad (10)$$

<sup>5</sup> <https://pla.esac.esa.int/#catalogues>

where  $\rho_c$  is the critical density of the Universe at the cluster redshift and  $\delta_{c_{200}}^6$  is a function that depends only on  $c_{200}$ , the concentration parameter. Here we switch from an overdensity of 500 to 200 in order to conform to most of previous works. Finally,  $r_s$  represents a characteristic radius and it is also a free parameter of the model. Using this definition, we obtain

$$P_{\text{zero}} - P(r_a) = -\mu m_p G 4\pi \rho_c \delta_{c_{200}}(c_{200}) r_s^3 \int_{r_a}^{r_{\text{zero}}} \frac{n_e(r)}{r^2} \left[ \frac{1}{1+r/r_s} + \ln(1+r/r_s) - 1 \right] dr, \quad (11)$$

where  $r_{\text{zero}}$  is the radius at which we are dominated by a zero-level component.

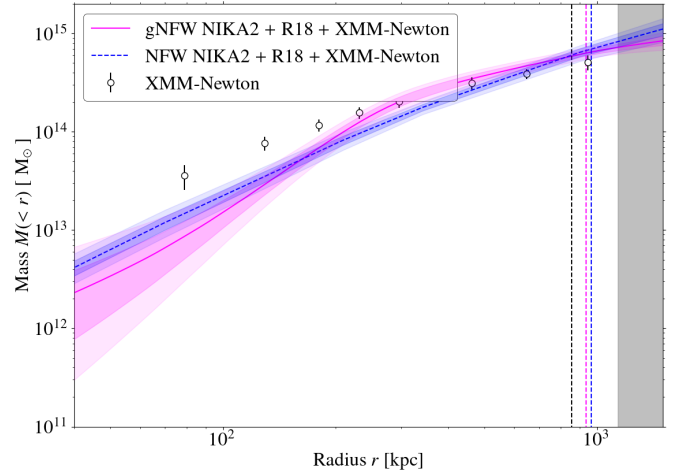
We perform a MCMC analysis similar to the one described above for the gNFW pressure profile model. In this case the free parameters of the model are  $\vartheta = [c_{200}, r_s, P_{\text{zero}}]$ . At each step of the MCMC we compute the integral in Eq. (11) to evaluate  $P(\vartheta)$  as needed for the likelihood function in Eq. (7). Calculating the integral can be computationally very expensive. As the result of this integral depends only on  $r_s$  and  $r_a$ , we create a grid of the integrals for a range of  $r_s$  values (from 100 to 2000 kpc) and  $r_a$  the radial bins of interest. We use this grid to interpolate the values of the integrals at each step. Flat priors are given for the concentration and the characteristic radius,  $0 < c_{200} < 8$  and  $100 \text{ kpc} < r_s < 2000 \text{ kpc}$ . The values of the electron density at different radii are obtained by logarithmic interpolation. We also make use of the Python NFW package<sup>7</sup>.

The best-fit pressure profiles and uncertainties are presented in the right panel of Fig. 8 for the four NIKA2 radially binned data sets discussed above. The posterior probability distributions of the free parameters of the models are shown in Appendix C. The posterior distributions of the  $c_{200}$  and  $r_s$  parameters can be compared to the results for the analysis of clusters in X-rays in Pointecouteau et al. (2005), Ettori et al. (2019), and Eckert et al. (2022). In these studies,  $c_{200}$  spans from 1 to 6 and  $r_s$  from 200 kpc to 1200 kpc, which is compatible with our results. We find that the NFW model is overall a good fit to the data as shown by the reduced  $\chi^2 \sim 1$  (see Fig. C.2 for the distributions). However, we observe that the uncertainties increase significantly in the outskirts of the cluster with respect to the gNFW pressure profile model discussed in the previous section. This can probably be explained by the flexibility of the NFW-based approach, which is high enough to show that the last point in the profile is not well constrained by the data.

### 5.3. HSE mass estimates

We present in Fig. 10 the HSE mass profiles inferred from the gNFW best-fit pressure profile, in combination with the *XMM-Newton* electron density, and from the NFW density best-fit model. Uncertainties (shaded areas) are obtained directly from the MCMC chains by computing the HSE mass profile for each sample from the model parameters. For the sake of clarity, we only present the masses obtained with NIKA2 AO1D estimates, but we changed the colour-coding for the gNFW profile so that we can differentiate between the two results.

The capability of the pressure model to describe the shape of the profile slopes is the key element for a good mass reconstruction. From the above results it seems that the NFW approach



**Fig. 10.** HSE mass profile estimates for CL J1226.9+3332 obtained with NIKA2 (angle order 1D) and R18 tSZ data combined with the *XMM-Newton* electron density profile. The solid magenta and dashed blue lines correspond to the gNFW and NFW methods, respectively. The shaded areas show the 2.5th, 16th, 84th, and 97.5th percentiles. Empty dots correspond to the HSE mass profile obtained from *XMM-Newton*-only data. Vertical dashed lines show the  $R_{500}^{\text{HSE}}$  obtained from each mass profile. The grey region represents the radial ranges at which the profiles are extrapolated.

does not have enough degrees of freedom to fully describe slope variations in the reconstructed pressure profile. Nevertheless, the resulting HSE mass profiles for the two models are compatible within  $2\sigma$ . The vertical dashed lines in the figure represent the  $R_{500}^{\text{HSE}}$  for each mass profile. Slight differences in the shape of the pressure profile at these radial ranges are critical for defining  $R_{500}^{\text{HSE}}$ .

These mass profiles, obtained from the combination of tSZ and X-ray data, are also compared to the X-ray-only HSE mass estimate in Fig. 10. Assuming spherical symmetry, the X-ray mass profile was derived, following the Monte Carlo procedure detailed in Démoclès et al. (2010) and Bartalucci et al. (2017) with the *XMM-Newton* electron density and temperature profiles presented in Sect. 4.2.

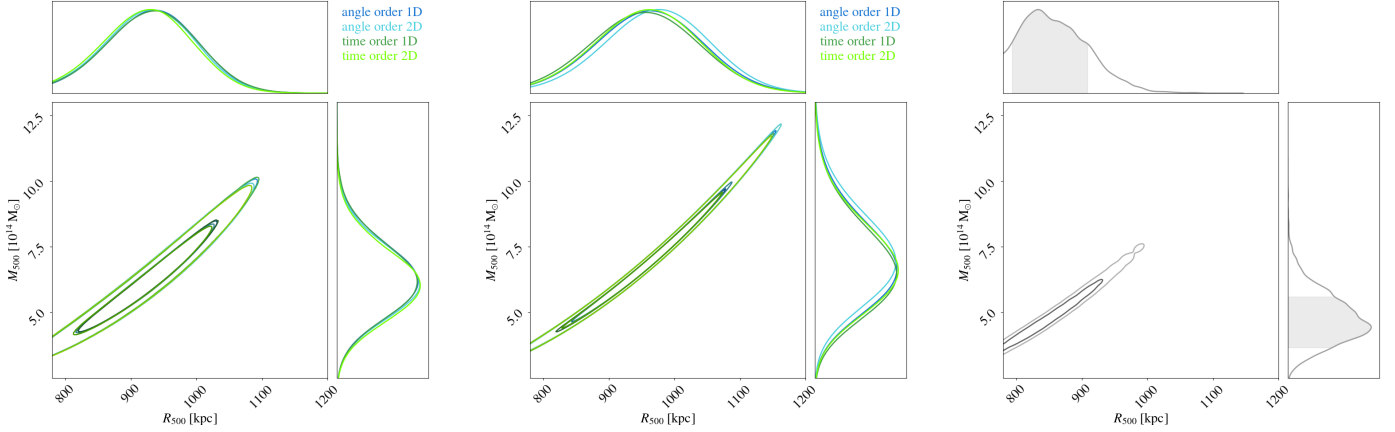
Despite the different behaviour of the X-ray-only profile in the cluster core, it is consistent with the tSZ+X estimates at around  $R_{500}^{\text{HSE}}$ . We note that the difference must come from the estimate of the derivative of the pressure, which for X-analyses is computed using the density and temperature profiles. The temperature profile from the X-rays is not used to compute the HSE mass in the tSZ+X analyses.

From the reconstructed HSE mass profiles we can obtain  $R_{500}^{\text{HSE}} - M_{500}^{\text{HSE}}$  probability distributions for each of the considered cases. We present in the left and central panel of Fig. 11 the  $R_{500}^{\text{HSE}} - M_{500}^{\text{HSE}}$  distributions for the gNFW and NFW models. They were obtained from the MCMC chains in the same way the uncertainties in Fig. 10 were computed. The results also account for the additional  $\sim 20\%$  uncertainty introduced by the electron density profile to  $M_{500}$ . These uncertainties were computed by combining the best-fit NFW and gNFW profiles with random realisations of the electron density profile following a Gaussian distribution centred in the central values of  $n_e$  and with the error bars as standard deviation. The width of the ellipses is an artefact from the display procedure, and each value of  $M_{500}^{\text{HSE}}$  is associated with a single value of  $R_{500}^{\text{HSE}}$ .

We present here the results for the four NIKA2 analyses (AO1D/2D and TO1D/2D). The results are consistent, with little dependence on the chosen TF estimate. From the comparison of

<sup>6</sup> Defined as  $\delta_{c_{200}} = \frac{200}{3} \frac{c_{200}^3}{\ln(1+c_{200}) - c_{200}/(1+c_{200})}$ .

<sup>7</sup> Jörg Dietrich, 2013–2017, <https://doi.org/10.5281/zenodo.50664>



**Fig. 11.** Probability distributions (1D and 2D) for  $M_{500}^{\text{HSE}}$  and  $R_{500}^{\text{HSE}}$  from the gNFW (left) and NFW (centre) models in the combined *XMM-Newton* and NIKA2 and R18 data, and from the *XMM-Newton* X-ray-only data (right). The different blue and green lines correspond to results for the four NIKA2 test cases considered.

**Table 3.** HSE masses for different estimates at  $R_{500}^{\text{HSE}}$ .

HSE mass estimates	$M_{500}^{\text{HSE}} [10^{14} M_{\odot}]$
(tSZ+X-ray) <sub>gNFW</sub>	$6.26^{+1.38}_{-1.38}$
(tSZ+X-ray) <sub>NFW</sub>	$7.00^{+1.76}_{-1.76}$
X-ray	$4.83^{+0.98}_{-0.96}$

the left and central panels in Fig. 11 we verify that the largest uncertainty in the HSE mass estimates comes from the modelling of the pressure profile. In spite of this effect, the reconstructed HSE mass profiles are compatible within  $1\sigma$ . The right panel of Fig. 11 shows the  $R_{500}^{\text{HSE}} - M_{500}^{\text{HSE}}$  probability distribution obtained with *XMM-Newton*-only data. Even if it is compatible with the gNFW and NFW results, the X-ray-only results favour lower HSE masses. A similar effect was observed for the ACT-CL J0215.4+0030 cluster (Kéruzoré et al. 2020), but not for PSZ2 G144.83+25.11 (Ruppin et al. 2018). We summarise in Table 3 the marginalised  $M_{500}^{\text{HSE}}$  masses obtained in this work. We give the mean value and the 84th and 16th percentiles. For gNFW and NFW we combine the probability distributions obtained for the four NIKA2 results so that the results account for the systematic effects from the NIKA2 data processing. The error bars also account for the uncertainties related to the electron density.

## 6. Lensing mass

Lensing masses, in contrast to HSE ones, probe the total mass without assumptions on the dynamical state of the cluster. For this reason it is of great interest to compare the HSE masses to lensing estimates. In this section we present the lensing mass estimate for CL J1226.9+3332.

### 6.1. Lensing data

We use the CLASH convergence maps (hereafter  $\kappa$ -maps) obtained from the weak and strong lensing analysis by Zitrin et al. (2015). In the analysis the authors reconstructed the  $\kappa$ -maps for the 25 massive CLASH clusters (Postman et al. 2012) using two different lensing models: 1) Light Traces Mass (LTM) and 2) Pseudo Isothermal Elliptical Mass Distribution

plus an elliptical NFW dark matter halo (PIEMD+eNFW). These models have been detailed in previous studies (Zitrin et al. 2009, 2013a,b, 2015). In this work we compute the mass estimate for both  $\kappa$ -maps in order to account for differences in the convergence map modelling.

### 6.2. Lensing mass density profile

For the lensing mass profile reconstruction we followed a similar approach to the one described in Ferragamo et al. (2022). The convergence maps describe the projected mass density of the cluster,  $\Sigma$ , in critical density units,  $\kappa = \Sigma/\Sigma_{\text{crit}}$ , with

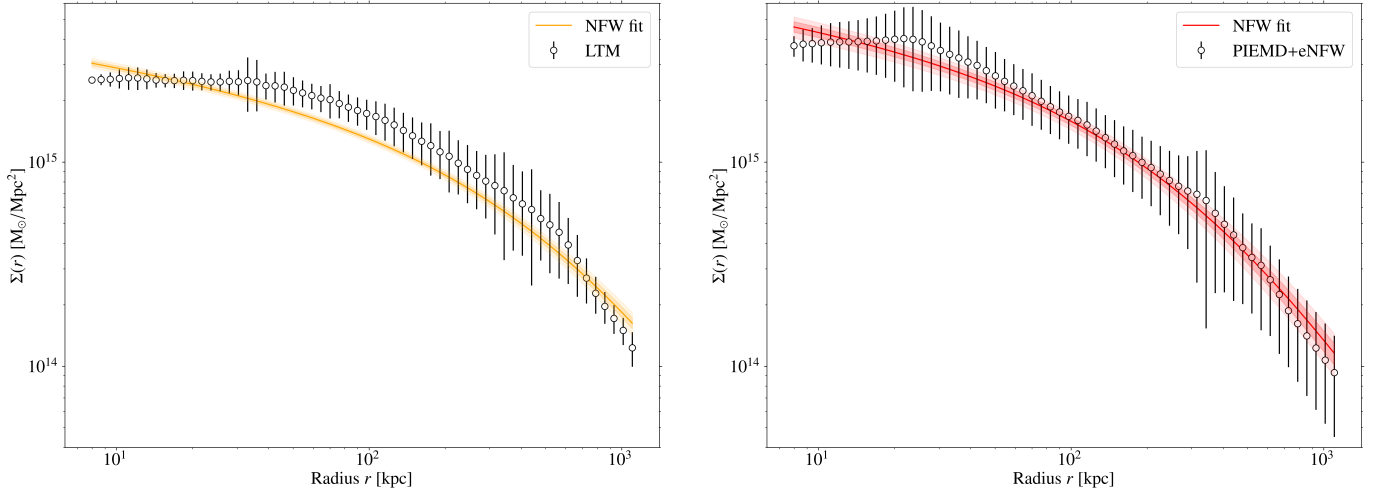
$$\Sigma_{\text{crit}} = \frac{c^2}{4\pi G} \frac{D_s}{D_l D_{ls}}. \quad (12)$$

Here  $D_s$ ,  $D_l$ , and  $D_{ls}$  correspond to the angular diameter distance between the observer and the source, between the observer and the lens (the cluster), and between the source and the lens, respectively. The publicly available CLASH  $\kappa$ -maps (Zitrin et al. 2015) have been normalised to  $D_s/D_{ls} = 1$ .

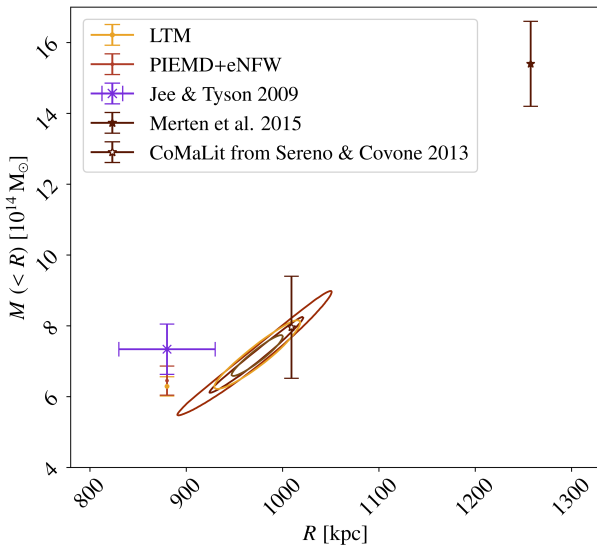
To estimate the lensing mass profile of CL J1226.9+3332, we fit a mass density model to the  $\Sigma$ -map. We assume spherical symmetry and a NFW density profile. We choose to directly fit the analytical projected NFW density profile (Eq. (5) in Ferragamo et al. 2022, derived from Navarro et al. 1996; Bartelmann 1996) to the radially averaged projected profiles of the  $\Sigma$ -maps. We consider as free parameters  $r_s$  and  $c_{200}$  (see Eq. (10)). The fit is performed via a MCMC analysis using the emcee software and the NFW Python package. We centre the projected mass profiles at the same position as for the pressure and electron density profiles in Sects. 4.1 and 4.2, the X-ray centre. The covariance matrix of the radial profile bins is built accounting for the 100 realisations provided for each model<sup>8</sup> and the uncertainties from the dispersion in each radial bin.

We show in Fig. 12 the radial profiles of the projected mass density for CL 1226.9+3332 obtained from the CLASH LTM (left) and PIEMD+eNFW (right) convergence maps. We present for both profiles the projected best-fit NFW density model and percentiles (shaded area). We observe that for the

<sup>8</sup> The 100 maps are available at <https://archive.stsci.edu/missions/hlsp/clash/clj1226/models/zitrin/lrm-gauss/v2/range/> and <https://archive.stsci.edu/missions/hlsp/clash/clj1226/models/zitrin/nfw/v2/range/>



**Fig. 12.** Projected mass density profiles obtained from CLASH convergence maps for the LTM (left) and PIEMD+eNFW (right) models. We also show the best-fit NFW model (orange and red lines) and the 2.5th, 16th, 84th and 97.5th percentiles (shaded area).



**Fig. 13.** Probability distribution for  $M_{500}^{\text{lens}}$  and  $R_{500}^{\text{lens}}$  obtained from the fit of the NFW density profile model on the CLASH PIEMD+eNFW (red) and LTM (orange) convergence maps. Shown are the 1 and  $2\sigma$  contours. The purple cross corresponds to the lensing mass estimate in Jee & Tyson (2009). For comparison, also shown are lensing masses estimated at 880 kpc for the PIEMD+eNFW (red) and LTM (orange) as vertical error bars at  $1\sigma$ . The brown stars correspond to  $M_{500}^{\text{lens}}$  from literature: Merten et al. (2015, filled) and Sereno (2015), Sereno & Covone (2013, empty).

LTM convergence map the best-fit NFW model underestimates the data except for cluster core and that the uncertainties in the model do not fully account for this. By contrast, the fit for the PIEMD+eNFW succeeds in representing the data.

From these results we conclude that the main uncertainties in the reconstruction of the lensing density profile come from the reconstruction of the convergence map. Thus, in the following we account for those.

### 6.3. Lensing mass estimates

From the obtained NFW density profiles we can reconstruct the lensing mass profiles (Eq. (9)) and subsequently the  $M_{500}^{\text{lens}}$

and  $R_{500}^{\text{lens}}$  probability distributions. We present in Fig. 13 the main results of this analysis for both convergence map models: LTM and PIEMD+eNFW. We obtain consistent results, and accounting for their combined posterior distributions we measure  $M_{500}^{\text{lens}} = 7.15^{+0.52}_{-0.51} \times 10^{14} M_{\odot}$  mean value and the 84th and 16th percentiles.

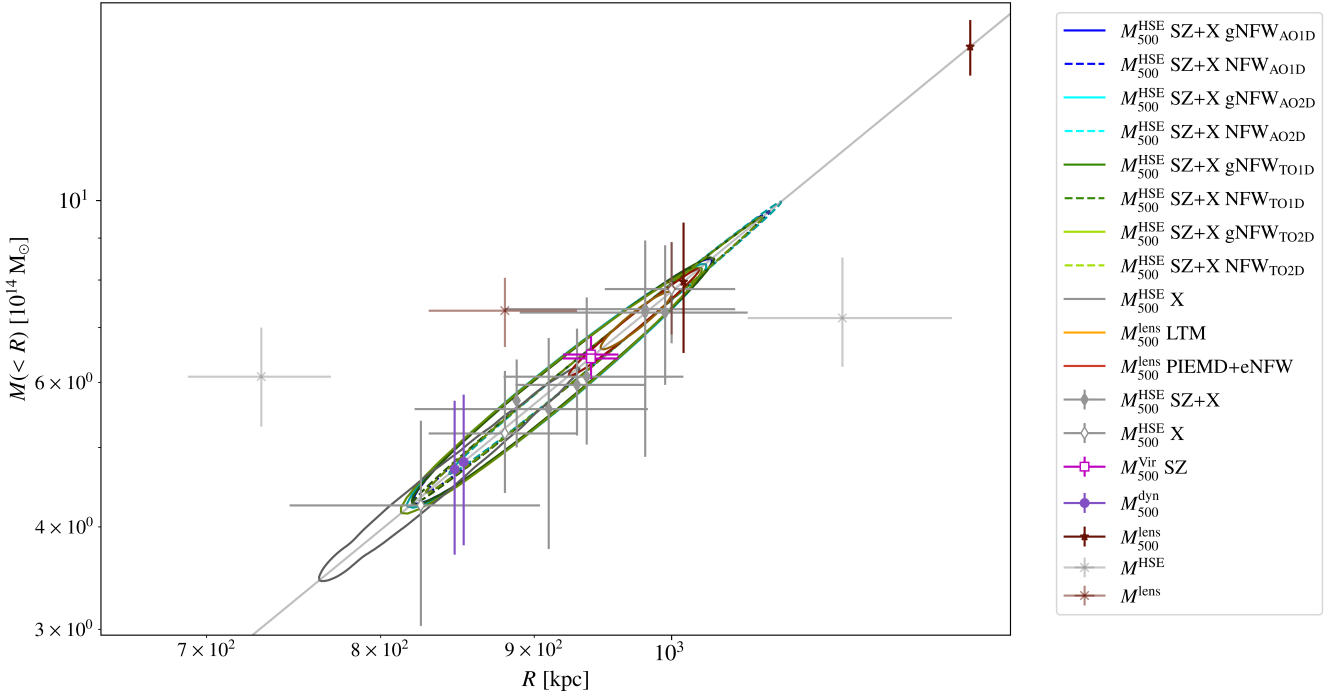
The Jee & Tyson (2009) weak-lensing analysis did not provide the direct  $M_{500}^{\text{lens}}$ , but evaluated the lensing mass at the  $R_{500}$  from Maughan et al. (2007). This result, shown in purple in Fig. 13, is consistent within  $1\sigma$  with our PIEMD+eNFW analysis when evaluating the lensing mass at the same radius (see red error bars in the figure) and less than  $2\sigma$  away from LTM (orange error bars). Merten et al. (2015) performed an independent analysis of the CLASH data, reconstructing their own convergence map. The projected mass density profile presented in Fig. 16 in Merten et al. (2015) shows a denser cluster than the profiles from the convergence maps used in this work (Fig. 12). For this reason, Merten et al. (2015) obtained, also with a NFW density fit, 35% higher masses than Jee & Tyson (2009). The corresponding  $M_{500}^{\text{lens}}$  is shown as a filled brown star in Fig. 13. The disturbed state of CL J1226.9+3332 could be the reason, according to Merten et al. (2015), for the different lensing mass estimates. Moreover, the high redshift of the cluster makes the precise reconstruction of the convergence map more difficult. The empty brown star in Fig. 13 corresponds to the lensing mass from CoMaLit (Sereno 2015) obtained from the analysis in Sereno & Covone (2013). In this case, the result is compatible with our mass estimates.

## 7. Comparison of mass estimates

### 7.1. CL J1226.9+3332 mass at $R_{500}$

The comparison of different mass estimates is difficult and can lead to incorrect physical conclusions. In particular, when comparing integrated masses the radius at which the mass is computed has a significant impact:  $R_{500}$  and  $M_{500}$  being constrained at the same time, our data are affected by a degeneracy. For this reason, in Fig. 14 we show, in the  $R_{500} - M_{500}$  plane, the results from the literature compared with those obtained in this work.

The green and blue contours show the  $R_{500}^{\text{HSE}} - M_{500}^{\text{HSE}}$  results obtained in this work for the gNFW (solid lines) and NFW



**Fig. 14.**  $R - M(<R)$  plane summarising the  $R_{500} - M_{500}$  results for CL J1226.9+3332. In the case of the HSE mass estimates the blue and green  $1\sigma$  contours show the results obtained in this work combining tSZ and X-ray data for the four NIKA2 analyses. The solid and dashed lines are for the gNFW pressure and the NFW density models, respectively. The grey contour corresponds to the HSE mass estimate for the *XMM-Newton*-only data. In the case of literature data the filled (empty) grey diamonds represent HSE masses from the combination of tSZ and X-ray data (X-ray-only results). The red and orange contours correspond to the lensing mass estimates obtained from the CLASH LTM and PIEMD+eNFW convergence maps in this work, respectively. Pink squares show the tSZ-only mass assuming virial relation, purple circles are dynamical mass estimates, and brown stars the lensing estimates. Grey and brown crosses show HSE and lensing masses from the literature close to  $M_{500}$ . The diagonal bright grey line defines the  $R_{500} - M_{500}$  relation. Slight deviations from this line are due to differences in the cosmological model used in each work.

(dashed lines) tSZ and X-ray data combined analyses. For comparison the filled grey diamonds correspond to the results from the literature presented in Fig. 1, also for combined tSZ and X-ray data. We observe that the results in this paper are compatible with previous analyses within  $1\sigma$ , centred around  $\sim 7 \times 10^{14} M_{\odot}$ .

Regarding X-ray-only results, the HSE mass estimates obtained in this work with *XMM-Newton* data (grey contours) suggest mass values centred at  $\sim 5 \times 10^{14} M_{\odot}$ . This is in agreement with the lowest estimates from the literature (empty grey diamonds) presented in Bulbul et al. (2010) and Maughan et al. (2007). On the contrary, the results from Mantz et al. (2010) and Mroczkowski et al. (2009) show higher masses. However, the  $M_{500}^{\text{HSE}}$  in Mantz et al. (2010) is not a direct measurement, but an extrapolation from a gas mass measured at  $R_{2500}$  converted into total mass, making this result less reliable. Overall, for CL J1226.9+3332 the HSE masses obtained only from X-ray data tend to lower values than those from the combination of tSZ and X-rays.

The result from Planck Collaboration XVIII (2015) is also a special case, as it is not a direct mass measurement, but a mass obtained from the X-ray-derived scaling relation (Eq. (7) in Planck Collaboration XX 2013) applied to the SZ measurement. This may explain why it lies at the border between the X-ray-only data and the tSZ+X combined results. The differences observed between X-ray-only and the combined tSZ+X results could have a physical and observational origin. For such a high-redshift cluster X-ray observations become challenging. If the south-western sub-clump in the cluster is really a hot but not dense structure (as suggested by Jee & Tyson 2009), the electron

density measurements from X-ray observations might be difficult to perform.

We also show in Fig. 14 the lensing, dynamical and virial mass estimates. The lensing mass estimates from this work for the CLASH LTM and PIEMD+eNFW convergence maps are presented as dark orange and red contours, respectively. We observe that they are consistent with the lensing mass from Sereno (2015) and Sereno & Covone (2013) and with HSE mass estimates, but very different from the Merten et al. (2015) lensing estimate for the reasons explained in Sect. 6. The virial masses estimated in Mroczkowski (2011, 2012) without X-ray data are shown as pink squares. They rely on the virial relation and on given pressure and density profile models to relate directly the integrated tSZ flux to the mass (Eq. (15) in Mroczkowski 2011). This kind of analysis seems a good alternative to the HSE mass for clusters without X-ray data. The dynamical mass estimates (purple circles), which we would expect to be larger than the HSE estimate, appear particularly low for CL J1226.9+3332 (Aguado-Barahona et al. 2022). According to the  $M_{500}^{\text{SZ}} - M_{500}^{\text{dyn}}$  scaling relation obtained from the analysis of 297 *Planck* galaxy clusters in Aguado-Barahona et al. (2022, Eq. (8) and Table 2) and considering  $M_{500}^{\text{SZ}}$  the value in Planck Collaboration XVIII (2015), the dynamical mass corresponding to CL J1226.9+3332 should be in the range  $6 - 7.5 \times 10^{14} M_{\odot}$ , thus more in agreement with our lensing mass estimates. Similar problems are reported in Ettori et al. (2019) and Logan et al. (2022), the latter showing that a large number of galaxies with spectroscopic redshift measurements ( $>200$ ) are needed. The orientation of the merger could also be an explanation for the underestimation: if the merger is happening in



the plane of the sky, the dispersion, and thus the mass, is lower. Finally, the brown and grey crosses show that the lensing and HSE mass estimates from the literature summarised in Table A.1 are close to  $M_{500}$ , but have not been evaluated at  $R_{500} - M_{500}$ : the  $M_{1000}$  from Maughan et al. (2004), the  $M_{200}$  from Muchovej et al. (2007), the mass at 880 kpc from Jee & Tyson (2009), and the mass at 1000 kpc from Sereno & Covone (2013) and Sereno (2015).

## 7.2. Hydrostatic-to-lensing mass bias

### 7.2.1. Hydrostatic mass bias problem

We do not expect the HSE to be fulfilled by all galaxy clusters in the Universe. We define the hydrostatic mass bias as

$$b = (M^{\text{true}} - M^{\text{HSE}})/M^{\text{true}}, \quad (13)$$

where  $M^{\text{true}}$  is the total true mass of the cluster.

From the observational point of view there are hints of a non-null hydrostatic bias. As mentioned in Sect. 1, one example is the tension observed between the cosmological parameters derived from *Planck* cluster number counts and those from the CMB analyses (Planck Collaboration XXI 2013). A possible explanation for this tension is that cluster masses are underestimated.

According to Planck Collaboration XXIV (2015) the bias needed to reconcile the cosmological constraints obtained from the CMB power spectra to the cluster counts is  $1 - b = M_{500}^{\text{HSE}}/M_{500}^{\text{true}} = 0.58 \pm 0.04$ . A compatible value was obtained from the updated analysis in Salvati et al. (2019):  $1 - b = M_{500}^{\text{HSE}}/M_{500}^{\text{true}} = 0.62 \pm 0.05$ .

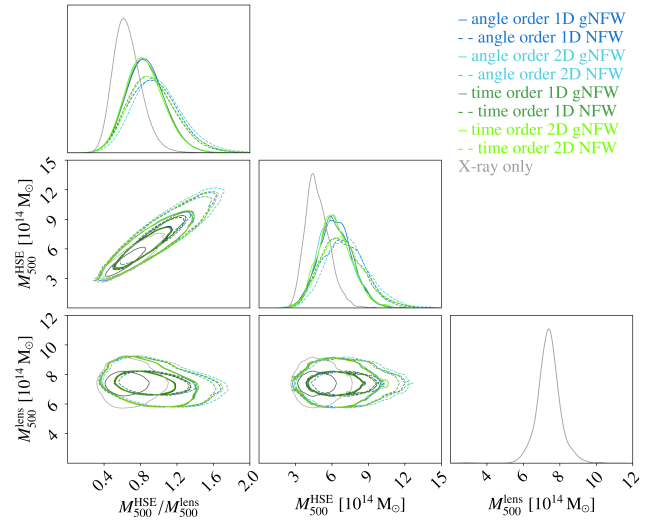
The hydrostatic bias has also been the topic of a large number of studies on numerical simulations (see Ansarifard et al. 2020; Gianfagna et al. 2021, and references therein). However, simulation-based analyses agree on values in the range of 0.75–0.9 for  $1 - b$ , and thus not able to reconcile the mentioned tension.

In a hierarchical formation scenario one would expect clusters at higher redshift to be more disturbed, and therefore the hydrostatic equilibrium hypothesis to be less valid (Neto et al. 2007; Angelinelli et al. 2020). Thus, this possible redshift evolution has been studied in the literature; for example, Salvati et al. (2019) find a modest hint of redshift dependence for the bias. However, other works based on cluster observations (McDonald et al. 2017) do not find traces of evolution of the morphological state with redshift. To date, the possible bias dependence with redshift has not been confirmed in simulations (see Gianfagna et al. 2021, and references therein).

### 7.2.2. Hydrostatic-to-lensing mass bias estimates

From the observational point of view, the real HSE bias is unachievable as one cannot determine the true mass of a cluster. However, it can be approximated using mass estimates that do not rely on the HSE hypothesis and trace the total mass of the cluster, for instance the lensing mass. In this work we compute the hydrostatic-to-lensing mass bias using the results obtained in Sects. 5.3 and 6 (see Sereno & Etti 2015, for an analysis of the CoMaLit samples). For the lensing mass  $M_{500}^{\text{lens}}$  we combine the probability distributions of both lensing models in Fig. 13. On the contrary, we consider the different HSE mass estimates independently. Assuming that HSE and lensing masses are uncorrelated estimates, we combined their probability distributions and computed the ratio  $M_{500}^{\text{HSE}}/M_{500}^{\text{lens}} = 1 - b_{\text{HSE}/\text{lens}}$ .

We present in Fig. 15 the hydrostatic-to-lensing mass ratio at  $R_{500}$ . The same colour-coding as in previous figures is used



**Fig. 15.** HSE mass estimates with respect to lensing estimates at  $R_{500}$ . Blue and green solid lines correspond to the HSE masses obtained from the gNFW pressure model fit and dashed lines to the NFW method. The grey line corresponds to the HSE mass obtained with *XMM-Newton*-only data. We present the  $1\sigma$  and  $2\sigma$  contours. The lensing mass distribution is the combination of the results for the PIEMD+eNFW and LTM analyses.

**Table 4.** Hydrostatic-to-lensing mass ratio for different HSE mass estimates.

HSE mass estimates	$1 - b_{\text{HSE}/\text{lens}}$
(tSZ+X-ray) <sub>gNFW</sub>	$0.86^{+0.20}_{-0.20}$
(tSZ+X-ray) <sub>NFW</sub>	$0.96^{+0.25}_{-0.25}$
X-ray	$0.66^{+0.14}_{-0.14}$

**Notes.** Shown are the mean value and the 84th and 16th percentiles.

to distinguish the HSE estimates that were obtained with each of the NIKA2 noise and filtering estimators. We draw with solid lines the results obtained modelling the pressure profile with the gNFW model and with dashed lines the results from the NFW fit. All of them are inferred from the NIKA2, R18, and *XMM-Newton* data. The grey contours show the bias using the X-ray-only HSE mass results. We present in Table 4 the marginalised hydrostatic-to-lensing mass ratio and uncertainties for the different cases considered. The results for the gNFW and NFW cases correspond to a combination of the four NIKA2 analyses. As concluded in Ferragamo et al. (2022), we observe that the value of the bias is sensitive to the considered data sets and modelling choices.

## 8. Summary and conclusions

The precise estimation of the mass of single clusters appears extremely complicated and affected by multiple systematic effects (Pratt et al. 2019), but it is the key to building accurate scaling relations for cosmology. The NIKA2 SZ Large Program provides an SZ-mass scaling relation from the combination of NIKA2 and *XMM-Newton* data. Within this programme, we presented a thorough study on the mass of the CL J1226.9+3332 galaxy cluster.

We obtained NIKA2 150 and 260 GHz maps, which allowed us to reconstruct the radially binned pressure profile for the ICM

from the tSZ data. To characterise the impact of the data processing, we repeated the whole analysis for two pipeline-filtering transfer functions and noise estimates for the 150 GHz map. We accounted for the presence of point sources that contaminate the negative tSZ signal at 150 GHz. The reconstructed NIKA2 pressure bins are compatible, within the angular scales accessible to NIKA2, with the profiles obtained from three independent instruments in R18. This validates the pressure reconstruction procedure that will be used for the whole sample analysis in the NIKA2 SZ Large Program.

We compared two approaches to estimate the HSE mass from the combination of tSZ-obtained pressure and X-ray electron density profiles. We considered either a gNFW pressure model (traditionally used for this kind of analysis) or an integrated NFW density model. The second seems a promising approach to ensure radially increasing HSE mass estimates. However, other density models should be tested in order to describe more satisfactorily the shape of the pressure profile. Both methods give completely compatible HSE mass profiles and integrated  $M_{500}^{\text{HSE}}$ . From the comparison of the different mass estimates, we also conclude that for the moment, when estimating the HSE mass of the CL J1226.9+3332 galaxy cluster in the NIKA2 SZ Large Program, the error budget is dominated by model dependence rather than by the instrumental and data processing systematic effects that we investigated. In addition, these results are in agreement with the X-ray-only HSE mass estimate obtained in this paper from the *XMM-Newton* electron density and temperature profiles. Nevertheless, the X-ray-only estimate favours lower mass values than the combined tSZ+X-ray results.

We also showed that our results are compatible with all the HSE mass estimates found in the literature within uncertainties, which are large. We think that the only way to reduce the current uncertainties is to precisely constrain the slope of the mass profile at  $\sim R_{500}$  (as already done in Etori et al. 2019) since we have proved that very similar mass profiles overall can result in significant differences at  $M_{500}$ .

From two differently modelled CLASH convergence maps, we reconstructed the lensing mass profile of CL J1226.9+3332 and measured the hydrostatic-to-lensing mass bias. We find that for CL J1226.9+3332 the bias is  $b_{\text{HSE}/\text{lens}} \sim 0.1$  for tSZ+X-ray combined HSE masses and  $\sim 0.4$  for X-ray-only estimates, with  $\sim 0.1$  uncertainties depending on the models. Despite the large uncertainties, we have the sensitivity to measure the bias of a single cluster, even for the highest-redshift cluster of the NIKA2 SZ Large Program. This is the second cluster in the NIKA2 SZ Large Program with such a measurement (the first was studied in Ferragamo et al. 2022) and the analysis of the HSE-to-lensing bias for a larger sample will be the topic of a forthcoming work (as done in Bartalucci et al. 2018 with X-ray data). Measuring the hydrostatic-to-lensing mass bias associated with LPSZ clusters will bring a new perspective to the bias from the tSZ side for spatially resolved clusters at intermediate and high redshifts.

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## Appendix A: Masses from the literature

We present in Table A.1 all the mass estimates found in the literature for CL J1226.9+3332 described in Sect. 2.2. We differentiate the masses reconstructed from ICM observables from the

lensing effect on background sources and from the study of the dynamics of member galaxies.

Most of the masses were computed from spherical models, and we give the radius at which each mass is evaluated when available. When the mass has been evaluated at a given  $R = R_\Delta$  we also present the value of the density contrast  $\Delta$ .

**Table A.1.** Mass estimates found in the literature for CL J1226.9+3332.

Observable	$R$ [kpc]	$\Delta$	$M(< R)$ [ $10^{14} M_\odot$ ]	Reference	Notes
<b>ICM</b>					
	$340 h_{100}^{-1}$	-	$3.9 \pm 0.5$	Joy et al. (2001)	
	1000	-	$14^{+6}_{-4}$	Jee & Tyson (2009)	Projected
	$730 \pm 40$	1000	$6.1^{+0.9}_{-0.8}$	Maughan et al. (2004)	
	$1660 \pm 340$	200	$14 \pm 4$	Maughan et al. (2004)	
	$880 \pm 50$	500	$5.2^{+1.0}_{-0.8}$	Maughan et al. (2007)	
	$1000 \pm 50$	500	$7.8 \pm 1.1$	Mantz et al. (2010)	
	$980^{+100}_{-70}$	500	$7.37^{+2.50}_{-1.57}$	Mroczkowski et al. (2009)	
	$410^{+10}_{-10}$	2500	$2.67^{+0.29}_{-0.27}$	Mroczkowski et al. (2009)	
	$980^{+90}_{-70}$	500	$7.30^{+2.10}_{-1.51}$	Mroczkowski et al. (2009)	
	$420^{+40}_{-30}$	2500	$2.98^{+0.90}_{-0.63}$	Mroczkowski et al. (2009)	
	$940^{+20}_{-20}$	500	$6.49^{+0.34}_{-0.34}$	Mroczkowski (2011, 2012)	
	$390^{+10}_{-10}$	2500	$2.35^{+0.15}_{-0.16}$	Mroczkowski (2011, 2012)	
	$940^{+20}_{-20}$	500	$6.42^{+0.36}_{-0.36}$	Mroczkowski (2011, 2012)	
	$400^{+10}_{-10}$	2500	$2.53^{+0.14}_{-0.15}$	Mroczkowski (2011, 2012)	
	$1140^{+100}_{-80}$	200	$7.19^{+1.33}_{-0.92}$	Muchovej et al. (2007)	
	$310^{+30}_{-20}$	2500	$1.68^{+0.37}_{-0.26}$	Muchovej et al. (2007)	
	$812^{+71}_{-81}$	500	$4.25^{+1.22}_{-1.14}$	Bulbul et al. (2010)	
	$379^{+37}_{-41}$	2500	$2.16^{+0.69}_{-0.63}$	Bulbul et al. (2010)	
	-	500	$5.7^{+0.63}_{-0.69}$	Planck Collaboration XVIII (2015)	Scaling relation
	$930^{+50}_{-43}$	500	$5.96^{+1.02}_{-0.79}$	Adam et al. (2015)	
	$937^{+72}_{-58}$	500	$6.10^{+1.52}_{-1.06}$	Adam et al. (2015)	
	$995^{+65}_{-65}$	500	$7.30^{+1.52}_{-1.34}$	Adam et al. (2015)	
	$910^{+89}_{-72}$	500	$5.57^{+1.81}_{-1.23}$	Castagna & Andreon (2020)	
<b>Lensing</b>					
	$1640 \pm 100$	200	$13.8 \pm 2.0$	Jee & Tyson (2009)	
	$880 \pm 50$	-	$7.34 \pm 0.71$	Jee & Tyson (2009)	
	155	-	$1.3 \pm 0.1$	Jee & Tyson (2009)	Projected
	155	-	$0.85 \pm 0.06$	Jee & Tyson (2009)	Projected (SW clump)
	$1680^{+100}_{-90}$	200	$13.7^{+2.4}_{-2.0}$	Jee et al. (2011)	
	-	200	$22.3 \pm 1.4$	Merten et al. (2015)	
	-	500	$15.4 \pm 1.2$	Merten et al. (2015)	
	-	2500	$6.1 \pm 1.0$	Merten et al. (2015)	
	-	200	$10.0 \pm 2.4$	Sereno & Covone (2013)	
	-	200	$11.114 \pm 2.442$	Sereno & Covone (2013), Sereno (2015) <sup>1</sup>	
	-	500	$7.96 \pm 1.44$	Sereno & Covone (2013), Sereno (2015) <sup>1</sup>	
	-	2500	$3.45 \pm 0.37$	Sereno & Covone (2013), Sereno (2015) <sup>1</sup>	
	500	-	$3.947 \pm 0.285$	Sereno & Covone (2013), Sereno (2015) <sup>1</sup>	
	1000	-	$7.882 \pm 1.013$	Sereno & Covone (2013), Sereno (2015) <sup>1</sup>	
	1500	-	$10.938 \pm 1.784$	Sereno & Covone (2013), Sereno (2015) <sup>1</sup>	
<b>Galaxy dynamics</b>					
	-	500	$4.7 \pm 1.0$	Aguado-Barahona et al. (2022)	
	-	500	$4.8 \pm 1.0$	Aguado-Barahona et al. (2022)	

## Appendix B: Pressure bins and point source fluxes at 150 GHz from joint fit

We present in Tables B.1 and B.2 the obtained point source fluxes at 150 GHz from the joint fit of the point sources with the pressure profile of the cluster in Sect. 4.1.1. We show the results

for the nine sources (PS1 to PS9), and each column corresponds to the value obtained for each of the four analyses. The reconstructed fluxes are consistent within all the considered cases.

The binned pressure reconstructions are also listed in Table B.3 and the bin correlation matrices are shown in Fig. B.1.

**Table B.1.** Submillimetre point sources fluxes. Fluxes at 150 GHz obtained from the joint point sources and pressure profile fits in Sect. 4.1.1.

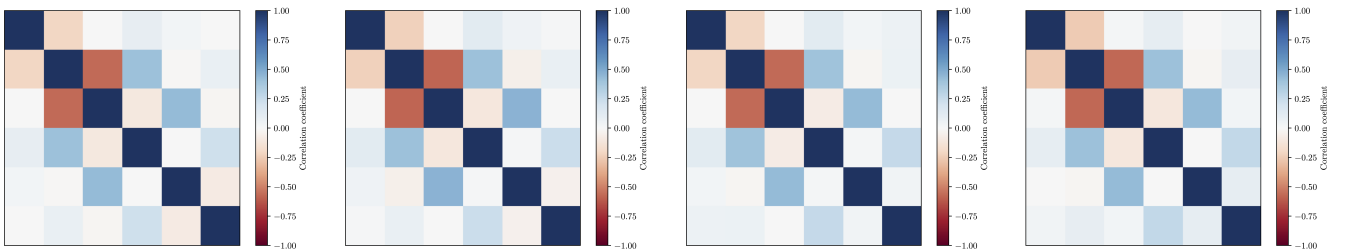
Source	150 GHz [AO 1D] [mJy]	150 GHz [TO 1D] [mJy]	150 GHz [AO 2D] [mJy]	150 GHz [TO 2D] [mJy]
PS1	$2.0^{+0.1}_{-0.2}$	$1.9^{+0.2}_{-0.2}$	$2.0^{+0.2}_{-0.2}$	$1.9^{+0.2}_{-0.2}$
PS2	$0.9^{+0.5}_{-0.1}$	$0.9^{+0.2}_{-0.1}$	$0.9^{+0.1}_{-0.2}$	$0.9^{+0.2}_{-0.1}$
PS3	$1.3^{+0.2}_{-0.1}$	$1.3^{+0.2}_{-0.1}$	$1.4^{+0.2}_{-0.2}$	$1.3^{+0.2}_{-0.2}$
PS4	$0.4^{+0.1}_{-0.1}$	$0.38^{+0.10}_{-0.07}$	$0.38^{+0.11}_{-0.07}$	$0.39^{+0.09}_{-0.08}$
PS5	$0.6^{+0.1}_{-0.1}$	$0.5^{+0.2}_{-0.1}$	$0.6^{+0.2}_{-0.1}$	$0.5^{+0.2}_{-0.1}$
PS6	$0.3^{+0.2}_{-0.1}$	$0.2^{+0.2}_{-0.1}$	$0.3^{+0.1}_{-0.2}$	$0.2^{+0.2}_{-0.1}$
PS7	$0.03^{+0.08}_{-0.03}$	$0.03^{+0.08}_{-0.03}$	$0.03^{+0.08}_{-0.03}$	$0.04^{+0.08}_{-0.04}$
PS8	$0.5^{+0.1}_{-0.1}$	$0.45^{+0.12}_{-0.09}$	$0.45^{+0.13}_{-0.09}$	$0.44^{+0.14}_{-0.07}$

**Table B.2.** Radio point source identified in the centre of the cluster. Fluxes obtained from the joint fits with the pressure profile of the cluster.

Source	150 GHz [AO 1D] [mJy]	150 GHz [TO 1D] [mJy]	150 GHz [AO 2D] [mJy]	150 GHz [TO 2D] [mJy]
PS9	$0.06^{+0.11}_{-0.05}$	$0.07^{+0.11}_{-0.06}$	$0.06^{+0.09}_{-0.05}$	$0.06^{+0.10}_{-0.04}$

**Table B.3.** Radially binned pressure profile fits. Mean values and the  $1\sigma$  uncertainties. The last bin results are trimmed.

$r$ [kpc]	$P_e$ [AO 1D] [ $10^4$ keV cm $^{-3}$ ]	$P_e$ [TO 1D] [ $10^4$ keV cm $^{-3}$ ]	$P_e$ [AO 2D] [ $10^4$ keV cm $^{-3}$ ]	$P_e$ [TO 2D] [ $10^4$ keV cm $^{-3}$ ]
35.0	$115 \pm 1088$	$190 \pm 1161$	$371 \pm 1054$	$1689 \pm 1103$
69.8	$1086 \pm 367$	$921 \pm 395$	$730 \pm 329$	$519 \pm 360$
209.5	$648 \pm 113$	$641 \pm 115$	$730 \pm 111$	$671 \pm 110$
341.7	$262 \pm 32$	$244 \pm 33$	$261 \pm 32$	$244 \pm 33$
900.0	$30 \pm 4$	$30 \pm 4$	$29 \pm 4$	$28 \pm 4$
1500.0	$2 \pm 2$	$4 \pm 2$	$4 \pm 3$	$8 \pm 3$



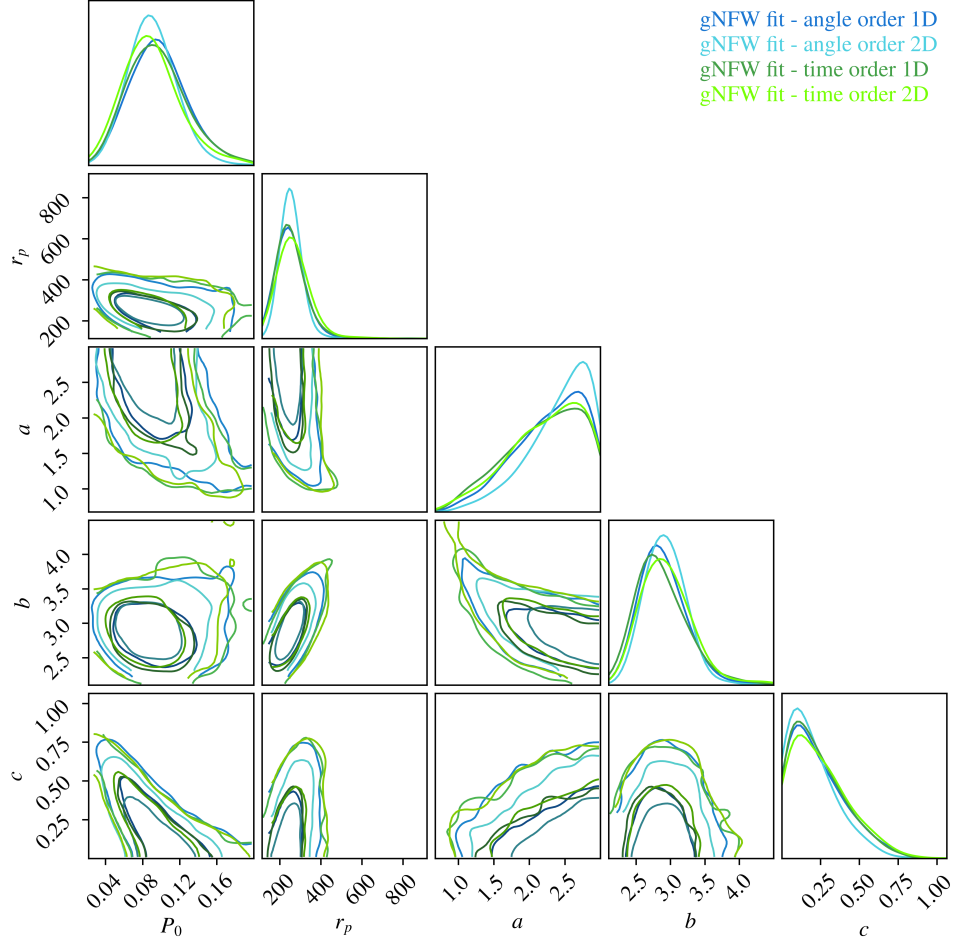
**Fig. B.1.** Correlation matrices of radially binned electron pressure bins. Results obtained with different transfer function and noise estimates (from left to right): AO1D, TO1D, AO2D, and TO2D.

**Appendix C: Pressure profile fits**

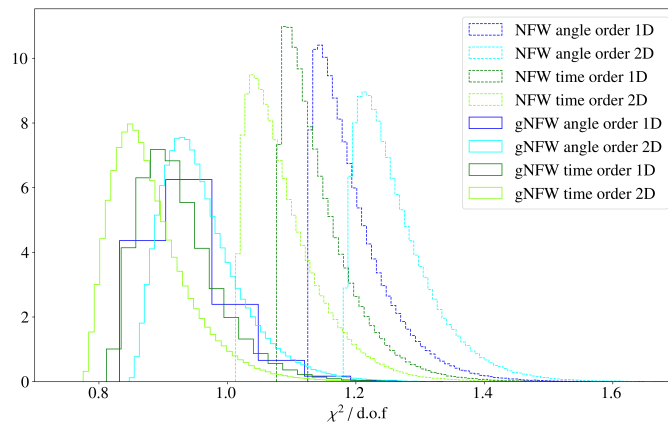
We show in figures C.1 and C.3 the posterior distributions of the parameters obtained from the fit of the pressure bins in

Sect. 5.2.1 and 5.2.2. The reduced  $\chi^2$  distributions corresponding to these fits are shown in Fig. C.2.

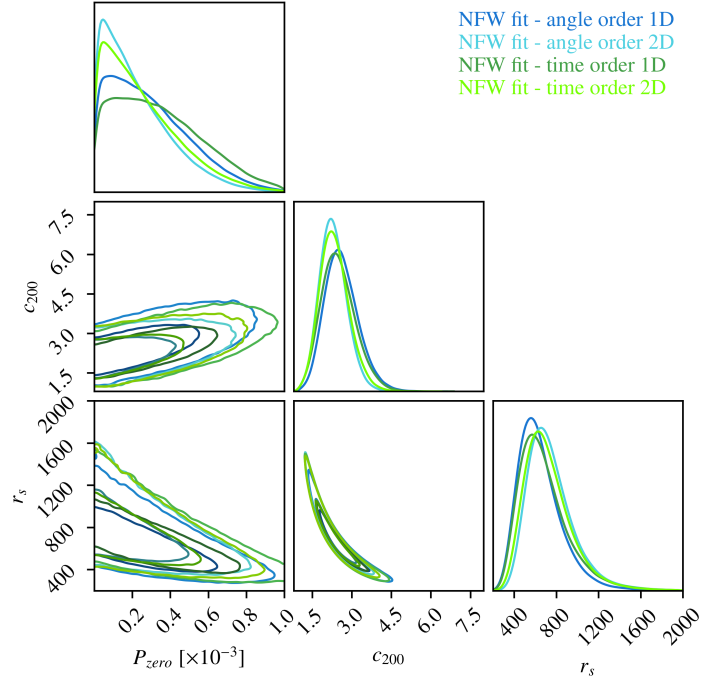
In Fig. C.4 we present the HSE mass profiles resulting from each of the fits.



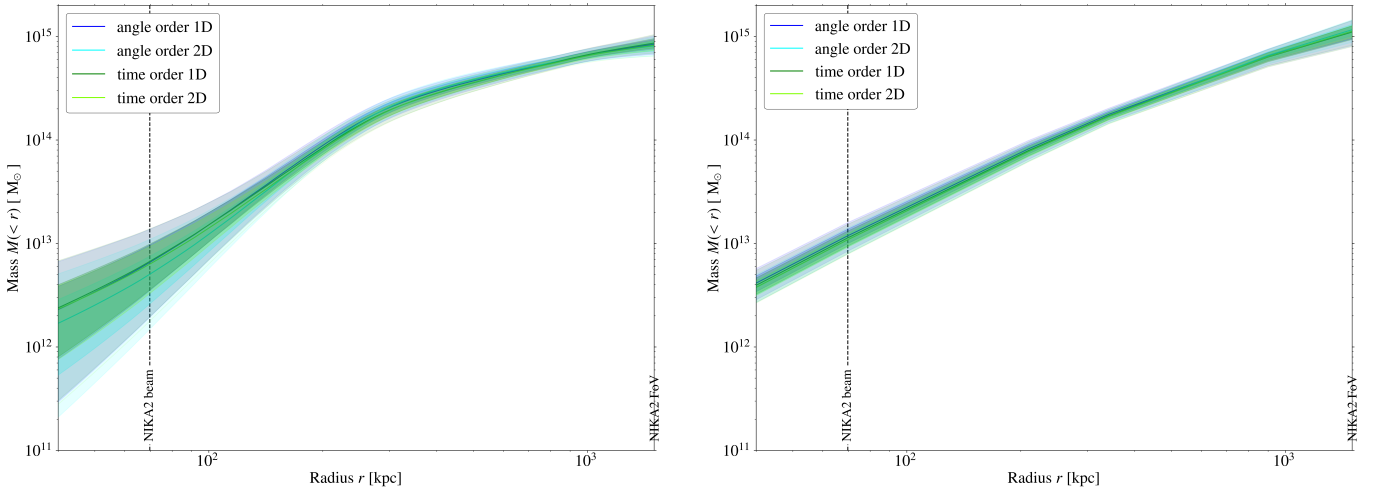
**Fig. C.1.** Posterior distributions of the parameters obtained in the fit of the NIKA2 and R18 pressure bins for the gNFW pressure model.



**Fig. C.2.** Reduced  $\chi^2$  of the gNFW (solid) and NFW (dashed) model fits to the NIKA2 and R18 pressure bins. The different colours show the NIKA2 bins used.



**Fig. C.3.** Posterior distributions of the parameters obtained in the fit of the NIKA2 and R18 pressure bins, combined with the XMM-Newton electron density, for the NFW density model fitted to the pressure.



**Fig. C.4.** HSE mass profiles obtained with the gNFW (left) and NFW (right) fit methods. Each colour represents the profile obtained using the corresponding NIKA2 pressure bins.