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### Highlights

**Subject**: Submission to *Engineering Structures* of the manuscript entitled "Modal analysis of a bistable deployable module with a refined joint model" co-authored by M. V. B. Santana, M. Hjiaj and P.Z. Berke

- Assess of the dynamic behavior of a square bistable deployable scissor structure (BDS) module in the deployed state through modal analysis incorporating the influence of internal stresses (specially relevant when member length imperfections are considered)
- Identification of the main design and topological parameters that have a critical influence on the natural frequencies and vibration modes and quantify their effects
- Influence of the applied material pair for inner and outer scissor like elements, finite joint stiffness and member length imperfections (in a stochastic approach) and topological parameters of the module that a designer would potentially adjust for a modular structure to tame its dynamic behavior.

# Modal analysis of a bistable deployable module with a refined joint model

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#### Abstract

Bistable deployable scissor structures (BDS) are interesting solutions for temporary civil engineering structures having the requirement of rapid erection and use. The transformation and design of BDS have been studied in various works in the literature using numerical simulations based on a quasi-static loading assumption. In contrast, their dynamic behavior and factors influencing it have not yet been extensively investigated. This contribution presents a computational study in which the natural frequencies and vibration modes of a single square BDS module are evaluated, incorporating the effects of the variation of particular design parameters specific to BDS, i.e. the materials used for inner and outer scissors, finite joint stiffness and manufacturing imperfections on the member lengths (in a stochastic approach) and two geometrical parameters of the module. The most dominant parameters are identified, and their effect is quantified, yielding results of practical interest for design and manufacturing that can be used to reach a target dynamic behavior of BDS.

*Key words:* bistable deployable structures, structural joints, modal analysis, structural dynamics

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#### 1. Introduction

Employing transformable structures can be of great interest for temporary structures in civil engineering because of their ease of deployment and rapidity of entering service. A specific class of transformable structures are bistable deployable scissor structures (BDS) that are practically an assembly of beam and hub elements through revolute joints. They exhibit a click-open (snap-through) behavior due to intended geometric incompatibilities between beam members connected as scissors. These incompatibilities induce internal stresses during transformation that are released in the stable folded and deployed configurations. This bistable nature ideally allows for bearing the self-weight of BDS in the deployed state, requiring only limited further bracing and ensuring rapid serviceability [?].

The transformation behavior of BDS has been recently studied rather extensively [?], leading to recommendations on their design procedure [?] and the proposal of a methodology for their shape and topological optimization [??]. Among other findings of interest, BDS were observed to be rather sensitive to manufacturing imperfections, including the hinge axis misalignment and beam lengths variations through nonlinear finite element modeling [?]. The spurious effect of these imperfections was found to be attenuated using structural joints with a finite stiffness. This joint compliance is a result of the incorporation of deformable bushings in the hinge, as shown in [?] where a finite stiffness joint finite element, dedicated to BDS was developed.

A significant part of the efforts gathered in the dedicated literature investigate the transformation behavior and design of BDS based on a quasi-static numerical approach, however studies addressing their dynamic behavior in civil engineering applications are very scarce, to the best knowledge of the authors. Vibrations of deployable structures that are not bistable have been addressed

in the literature for aerospace [??????] and for civil engineering applications [?], none of these have the SLE (scissor like elements) based topology that generates bistability. Addressing structural failure is of prime importance for BDS and extension of the computational study to this aspect is planned as future work. However, the present contribution focuses rather on the vibrational behavior of bistable deployable structures. The dangers of the unwanted transformation in service state have been partially addressed in [??].

The specific goals and the main novel contributions here are: (i) to assess the dynamic behavior of a square BDS module in the deployed state through modal analysis incorporating the influence of internal stresses (specially relevant when member length imperfections are considered), (ii) to identify the main design and topological parameters that have a critical influence on the natural frequencies and vibration modes and quantify their effects, particularly focusing on (iii) the applied material pair for inner and outer scissor like elements, finite joint stiffness and member length imperfections (in a stochastic approach) and topological parameters of the module that a designer would potentially adjust for a modular structure to tame its dynamic behavior. This clearly defines a gap of knowledge, and giving a contribution to the understanding of the dynamic behavior of BDS, the identification and quantification of effects influencing it are the main original goals of this work.

This contribution is organized as follows: Section 2 presents the computational tool, the main ingredients of the numerical model employed for BDS, the considered design parameters and their range. Section 3 is dedicated to the computational results and their analysis. Finally, the main conclusions and future research paths are given in Section 4.

#### 2. Computational model for dynamic analysis

The computational model set up for the dynamic analysis of a single square bistable module is presented in this section. The necessary details for the understanding of the behavior of the hinge finite element developed in [?] are recalled first with an original procedure of identifying the joint FE parameters from realistic bushing geometries. The derivation of the mechanical response of the joints serves subsequently for feeding the stiffness relationship into the modal analysis. Next, the basic steps of the modal analysis incorporating internal stresses in the deployed state, as a consequence of manufacturing imperfections on the beam length are described. The modal analysis makes use of the stiffness and mass matrices in the deployed state to obtain, via the associated eigenvalue problem, the natural frequencies and vibration modes of the deployable module. Finally, the choice of the selected BDS module parameters and their range for the subsequent dynamic parametric study are given.

### 2.1. Computational tool and ingredients

The studied square bistable module, shown in Fig. 1 is meant to be part of modular BDS. This module was first proposed by [?] and is composed of beams forming the inner and outer scissors. Its geometry is selected in such a way that the inner SLE are subjected to compressive axial forces and bending moments in both directions of the transformation (opening and closing), which generates the bistable behavior of the system. The force vs displacements transformation curves of this system were studied by many researchers, including more recently in [?], where the effect of the finite stiffness of the joints was considered.

In the numerical simulations, each beam half span is discretized into five FE, which is confirmed to be a converged mesh. Each of these 2-noded elements takes into account shear deformations following the Timoshenko beam's theory, i.e. the contribution of the shear stresses in the internal energy and stiffness are considered. The formulation also takes into account large displacements and rotations via the corotational approach [??]. Since the joints are modeled as usual finite elements with finite stiffness and inertia, their contributions to the global stiffness and inertia matrices of the system are assembled traditionally, with no need to rely on an augmented system of nonlinear equations or penalization. The model in Sec. 3 has a total of 160 elements and 186 nodes, corresponding to 1116 degrees of freedom.



Figure 1: The studied bistable deployable square module in the deployed (a) and in the folded (b) states.

The finite size hubs that serve as connection of more than two beams, situated at the corners and at the intersection of the inner diagonal scissor like elements (SLEs) are also modeled using these beam FE, however a very high stiffness is associated to them to match the usual design assumption that hubs are rigid [? ?].

Revolute hinges incorporating deformable bushings connect beams to beams and beams to hubs (Fig. 2). These are made of rubber, undergoing small elastic deformations, and are modeled employing the dedicated finite element allowing for incorporating a finite stiffness for the radial direction and for rotations around axes perpendicular to the hinge axis, referred to as tilting modes (explained in details later). This finite stiffness models the revolute hinge system in which two deformable bushings are inserted, its stiffness along these free directions is a user set parameter that can vary between very large values compared to the scissor beam stiffness (rigid pinned connection) and low values allowing for the relative movements of the connected parts that can be useful to compensate for manufacturing imperfections [?]. The axial movement of the hinge is set strictly to zero, independently of the other stiffness terms to simplify the model, considering that the axial stiffness of the hinges is expected to not play a significant role in their structural behavior.



Figure 2: Three dimensional view of the embedded deformable hinge in a single scissor showing the cutting plane (a) in which the details of the proposed technological solution are sketched with the bushings shown in light blue (b).

An important question that was not addressed in [?] is the identification of the stiffness parameters from actual bushing geometries, which is an original contribution presented in the following.

#### 2.1.1. Flexible hinge model and parameter identification

The derivation of the radial and tilting stiffness (Fig. 3) of a single cylindrical bushing when linear elastic behavior and small strains are supposed has been established in the literature. The expressions from [?] and [?] serve in this work to obtain the single bushing radial and tilting stiffness, using  $E_{\text{bushing}} = 0.8$ GPa [?]. The inner and outer diameters are  $D_{\text{in}} = 0.01$  m and  $D_{\text{out}} = 0.03$  m and  $D_{\text{in}} = 0.02$  m and  $D_{\text{out}} = 0.03$  m, with a bushing width of  $w_{\text{s}} = 0.01$  m for the soft and hard bushings, respectively. The resulting soft and hard single bushing radial and tilting stiffness values are obtained as  $K_{\text{r}} = 3.96 \times 10^7$  N/m and  $K_{\text{t}} = 9.40 \times 10^2$  N m/rad and  $K_{\text{r}} = 1.87 \times 10^8$  N/m and  $K_{\text{t}} = 4.20 \times 10^3$  N m/rad, respectively. The compliant bushings are spaced by a distance  $e_{\text{s}} = 0.004$  m (Fig. 2), corresponding to the beams interdistance.



Figure 3: Allowed deformation modes of a single deformable bushing.

Note that in a single beam-beam and beam-hub joint two of these bushings are mounted, as shown in Fig. 2 and the derivation of the stiffness relationships of the associated joint finite element (relating the shearing force to the relative radial displacement  $u_{\rm h}$  in Fig. 4a and the tilting moment to the relative tilting

rotation  $\theta_{\rm h}$  in Fig. 4b) is a non-trivial task. The related main equations are given in the following, starting with the relationships describing the kinematics of the two-bushing joint and the force and moment equilibrium allowing to deduce the radial and rotational stiffness for the joint assembly.



Figure 4: Assumed deformations of the hinge composed of two compliant bushings for shearing load (a) and for a tilting/torsion moment (b).

The pin inserted in the bushings is assumed to move as a rigid body (Fig. 5a) with a displacement  $u_p$  and a rotation  $\theta_p$ , as a result of the deformation of the two compliant bushings. Points  $B_1$  and  $B_2$  are the intersection of the pin axis and the two beam center lines. Their displacements ( $u_1$  and  $u_2$ ) can be described by the following kinematics:

$$u_1 = u_p - \left(\frac{w_s + e_s}{2}\right)\theta_p \tag{1}$$

$$u_2 = u_{\rm p} + \left(\frac{w_{\rm s} + e_{\rm s}}{2}\right)\theta_{\rm p} \tag{2}$$

with the beam thickness  $w_{\rm s}$  and spacer size  $e_{\rm s}$  as shown in Fig. 2.

For the sake of simplicity it is assumed that a shearing force induces only relative radial displacement and a tilting moment induces only relative tilting

rotation in the finite element formulation, i.e. the coupling of shearing force and relative tilting rotation,  $\theta_{\rm h}$  is disregarded in [?].

The assumed deformation of the system for a shearing load is shown in Fig. 4a. The forces acting on the rigid pin are equilibrated by the internal forces due to the radial deformation of the two bushings (Fig. 6a). Considering the relative radial displacement  $u_{\rm h}$  between the two compliant bushings and the single bushing radial stiffness  $K_{\rm r}$ , the following force equilibrium relation can then be derived:

$$K_{\rm r}u_1 + K_{\rm r}(u_2 - u_{\rm h}) = 0 \tag{3}$$



Figure 5: Rigid pin in the deformable hinge of a single beam-beam or beam-hub connection and the associated joint FE and its axis.



Figure 6: Force and moment free body diagram of the deformable hinge for shear and torsion loads.

Combining Eqs. (1 - 3), the radial displacement of the pin  $u_{\rm p}$  can be obtained as:

$$u_{\rm p} = \frac{u_{\rm h}}{2} \tag{4}$$

The moments generated in the pin by the eccentric radial forces and the tilting moments associated to the single bushing tilting stiffness  $K_{\rm t}$  can be related in the following equilibrium equation:

$$2K_{\rm t}\theta_{\rm p} + \left(\frac{e_{\rm s} + w_{\rm s}}{2}\right)K_{\rm r}(u_2 - u_{\rm h}) - \left(\frac{e_{\rm s} + w_{\rm s}}{2}\right)K_{\rm r}u_1 = 0$$
(5)

Combining Eqs. (1 - 2) and Eq. (5), the pin rotation  $\theta_{\rm p}$  can be obtained as:

$$\theta_{\rm p} = \left[\frac{\left(e_{\rm s} + w_{\rm s}\right)K_{\rm r}}{4K_{\rm t} + \left(e_{\rm s} + w_{\rm s}\right)^2K_{\rm r}}\right]u_{\rm h} \tag{6}$$

Finally, the joint radial stiffness is obtained by taking the equilibrium of forces at the loaded bushing (Fig. 4a):

$$F = K_{\rm r} \left( u_{\rm h} - u_{\rm p} \right) = \frac{K_{\rm r}}{2} u_{\rm h}$$
 (7)

The assumed deformation of the system for a torsion moment is shown in Fig. 4b. The moments acting on the rigid pin are equilibrated by the internal moments due to the tilt deformation of the two bushings (Fig. 6b). The force equilibrium relation of the rigid pin can be written as:

$$K_{\rm r}u_1 + K_{\rm r}u_2 = 0 \tag{8}$$

Combining Eqs. (1 - 2) with Eq. (8), the pin radial displacement  $u_p$  can be obtained as:

$$u_{\rm p} = 0 \tag{9}$$

Considering the relative tilt rotation  $\theta_{\rm h}$  between the two compliant bushings (Fig. 4b) and the single bushing tilting stiffness  $K_{\rm t}$ , the following moment equilibrium relation can then be derived:

$$K_{\rm t}\theta_{\rm p} + K_{\rm t}\left(\theta_{\rm p} - \theta_{\rm h}\right) + \left(e_{\rm s} + w_{\rm s}\right)K_{\rm r}\left(\frac{e_{\rm s} + w_{\rm s}}{2}\right)\theta_{\rm p} = 0 \tag{10}$$

Combining Eqs. (1 - 2) with Eq. (10), the rigid pin tilt rotation  $\theta_p$  can be

obtained as:

$$\theta_{\rm p} = \left[\frac{2K_{\rm t}}{4K_{\rm t} + \left(e_{\rm s} + w_{\rm s}\right)^2 K_{\rm r}}\right]\theta_{\rm h} \tag{11}$$

Finally, the joint tilt stiffness is obtained by taking the equilibrium of moments at the loaded bushing (Fig. 4b):

$$M = K_{\rm t} \left(\theta_{\rm h} - \theta_{\rm p}\right) = K_{\rm t} \left[\frac{2K_{\rm t} + \left(e_{\rm s} + w_{\rm s}\right)^2 K_{\rm r}}{4K_{\rm t} + \left(e_{\rm s} + w_{\rm s}\right)^2 K_{\rm r}}\right] \theta_{\rm h}$$
(12)

From the results above, the joint combined radial stiffness in (Eq. 7) is half of the single bushing radial stiffness. The joint combined tilting stiffness (Eq. 12) depends on the single bushing radial  $K_r$  and tilting  $K_t$  stiffness. In the choice of dimensions and material parameters of the present work, the term on the single bushing radial stiffness  $K_r$  in Eq. (12) is very large when compared with the term on the single bushing tilting stiffness  $K_t$ , and so the joint combined tilting stiffness is approximately equal to  $K_t$ .

The details of the joint finite element formulation, including the kinematic transformations via the corotational method and the explicit expressions of the internal force vector and tangent stiffness matrix, are given in [?]. A flowchart of the procedure is shown in Fig. 7. Provided the displacements and rotations of the nodes connected to the bushings, the corotational method is used to filter the rigid body motion of the joint and capture the portion of the motion that causes radial  $u_{\rm h}$  and tilting  $\theta_{\rm h}$  movements. A nonlinear constitutive model, that uses as parameters the combined bushing stiffness of Eqs. (7) and (12) to feed the local stiffness( $\overline{k}_0^{\rm n}$  and  $\overline{k}_0^{\rm m}$  in [?]), is then applied to obtain the corresponding forces and moments in the local system. Finally, these quantities are rotated back to the global system, where the equilibrium of the structure is evaluated.



Figure 7: Joint finite element procedure.

#### 2.1.2. Modal analysis accounting for internal stresses

Manufacturing imperfections naturally exist in all engineering structures and BDS were shown to be particularly sensitive to them in [??]. Imperfections on the hinge alignment and on the beam length generate spurious deformations and internal stresses in the deployed configuration of BDS, that can be mitigated by the use of deformable joints [?]. These internal stresses are expected to influence the dynamic behavior of BDS as well and should ideally be incorporated in a rigorous investigation when imperfections are considered.

In this work, this is attempted by a two-step approach in which fictitious thermal strains representing unequal beam lengths in the structure are first applied in a static equilibrium problem. The resulting internal stresses are incorporated as geometric stiffness in the subsequent modal analysis that is therefore performed in a deformed equilibrium state in the deployed configuration when member length imperfections are considered. Due to the complexity of the bistable deployable module, specially when considering the flexibility of the joints, analytical models aren't a feasible option and one must then rely on FE models for the modal analysis. The expressions for the beams mass and stiffness matrices and joints stiffness matrices are directly taken from the original works describing their corresponding FE formulations [? ?]. Considering that these references contain exhaustive information, no details are presented here for the sake of focus and brevity.

Once the stiffness  $\mathbf{K}_e$  and inertia  $\mathbf{M}_e$  matrices of each element are assembled into the global stiffness  $\mathbf{K}$  and inertia  $\mathbf{M}$  matrices of the structure, respectively, the modal analysis of the deployable unit is performed evaluating the eigenvalues,  $\omega$  (natural frequencies) and eigenvectors,  $\mathbf{v}$  (vibration modes) of the system, via the following symmetric generalized eigenvalue problem:

$$\left(\mathbf{K}(\mathbf{d}) - \omega^2 \mathbf{M}\right) \mathbf{v} = \mathbf{0} \tag{13}$$

Here, the stiffness matrix  $\mathbf{K}(\mathbf{d}) = \mathbf{K}_m(\mathbf{d}) + \mathbf{K}_g(\mathbf{d})$  depends on the system degrees of freedom  $\mathbf{d}$  and so incorporates the effects of internal stresses generated by the gravity load and by initial imperfections. In the present work it's assumed that the damping characteristics of the system are of traditional Aluminum structures (around 1.5% of the critical damping [?]) and therefore having a negligible influence on the natural frequencies of the deployable module. The routine employed to solve the generalized eigenvalue problem in Eq. (13) is *dsygv* provided in the classical numerical library *Lapack* [?]. The term  $\mathbf{K}_m(\mathbf{d})$  corresponds to the material stiffness (depending on the displacements because of the geometrically nonlinear corotational formulation employed here), while  $\mathbf{K}_g(\mathbf{d})$  incorporates internal stress induced stiffness, often referred to as geometric stiffness in the literature [? ?]:

$$\mathbf{K}(\mathbf{d}) = \underbrace{\mathbf{T}^T \overline{\mathbf{K}} \mathbf{T}}_{\mathbf{K}_{\mathrm{m}}} + \underbrace{\frac{\partial \mathbf{T}^T}{\partial \mathbf{d}} : \overline{\mathbf{f}}}_{\mathbf{K}_{\mathrm{g}}}$$
(14)

with  $\mathbf{\overline{f}}$  and  $\mathbf{\overline{K}} = \partial \mathbf{\overline{f}} / \partial \mathbf{\overline{d}}$  being the internal force vector and (elastic) material stiffness matrix in the corotational system, respectively, and  $\mathbf{T}$  is the transformation operator between structural and corotational axes.

When gravity loads and imperfections are considered, the elements internal forces  $\overline{\mathbf{f}}$  are first computed via a static nonlinear analysis. The modal analysis of the system is then performed with the tangent stiffness in the deformed configuration and the natural frequencies and vibration modes are computed.

### 2.2. Study parameters

The topology, geometrical parameters and dimensions of the studied square bistable module are shown in Fig. 8, where a diagonal cut plane (CT - T - B - CB) is shown. In the subsequent parametric study the thickness of the module (H =0.40 m) and the hubs size (r = 0.05 m) are kept constant, while the edge length L and top center point height h are varied. The considered boundary condition is setting all degrees of freedom of the bottom central point (CB) to zero, as often applied in transformation simulations [? ?]. Note that establishing more realistic support conditions to mimic the effect of other neighboring modules in a modular structure (e.g. by adding an equivalent stiffness along the square edges) is part of future work.

When the transformation of the module is studied, four horizontal transformation forces are applied on the top nodes of the module in a direction along the square diagonals. In the figures reporting transformation load vs displacement data the magnitude of a single diagonal transformation force F is plotted as a function of the top center point's (CT) vertical displacement w, systematically.

The reference case is defined as a single square module excluding geometric imperfections, made entirely of Aluminum (specific mass 2700 kg /  $m^3$ , elastic modulus 70 GPa) with rigid joints, i.e. sufficiently high stiffness to concentrate the deformation in the scissors only, while ensuring that no numerical ill-conditioning appears. The relative variation in the modal analysis results (e.g. natural frequencies) as a consequence of the variation of different model parameters is measured against this case (which corresponds to the design stud-



Figure 8: Half-diagonal plane cut of the studied square BDS and corresponding geometric variables.

ied in [?]). All other parameters of the module are given in Tab. 1. The hub size (r = 0.05 m) and beams spacing (s = 0.014 m) are as proposed in [?] and kept the same in all simulations. The inner and outer scissor beam cross-sections are chosen to be  $0.04 \text{ m} \times 0.01 \text{ m}$  with 0.001 m wall thickness. The dimensions chosen for the studied model are the same as the ones presented in references [?] and [?], where it was shown that the module is suitable as a deployable system.

Geometry		
Hub's size $(r)$	0.05 m	
Edge length $(L)$	1.00 m	
Module height $(H)$	0.40 m	
Apex hub height $(h)$	0.50 m	
Joints		
Spacer thickness	0.004 m	
Combined tilt stiffness max	$3.54 \times 10^3$ N m/rad	
Combined tilt stiffness min	$7.80 \times 10^2$ N m/rad	
Combined radial stiffness max	$1.98 \times 10^7 \text{ N/m}$	
Combined radial stiffness min	$9.34 \times 10^7 \text{ N/m}$	

Table 1: Reference model parameters.

The choice of the materials used for inner and outer scissor beams is a basic degree of freedom in BDS design [?]. In these structures the buckling that

induces the snap-through behavior during transformation is limited to the inner scissors. Increasing the elastic stiffness of these beams by the appropriate choice of the material leads to a higher peak force required for module transformation and a higher structural stiffness in the deployed service state. Choosing a material with low stiffness for the inner scissor beams can also be a viable strategy when aiming for a reduction of this peak force to allow human force induced transformation. Considering the above, a case in which the inner scissor beams are made of HDPE [?] (specific mass 940 kg /  $m^3$ , elastic modulus 0.8 GPa), while the outer ones are Aluminum will be studied to assess the effect of the material pair choice on the dynamic behavior of the module.

A given modular structure, with fixed volume, can be constructed with different number of modules, by controlling the size of the modules. Reducing the size of a single square module (i.e. the length of the square edge, L) naturally implies shorter scissor beams resulting in the variation of the module stiffness and consequently of its natural frequencies. This effect is considered in this work by varying the square module edge length with respect to the reference case.

Another geometrical design parameter incorporated here is the top center point (CT) height at the intersection of the inner diagonal scissors, h, shown in Fig. 8. Increasing this height was shown to result in a higher transformation peak force in [?] with a high sensitivity on this parameter. Changing h is a viable manipulation to tune the peak force of transformation in the design and since it modifies the orientation of the inner scissors it may also have an impact on the module's dynamic behavior, which motivates incorporating it in this study.

The cross-section geometry influences the mass and stiffness matrices via its area (A) and moments of inertia  $(J, I_{22} \text{ and } I_{33})$ . In the order to keep the focus

on external geometrical parameters, the influence of the section geometry on the natural frequencies and vibration modes is not considered in the present work. The procedure to incorporate them would be analogous to the one used for the other parameters, computing the natural frequencies and vibration modes via the obtained mass and stiffness matrices.

Previous work [?] shows that the behavior of deformable hinges and manufacturing imperfections [?] have a large influence on the transformation response of BDS. It is naturally expected that the hinge behavior and such imperfections impact the dynamic structural response as well, which is why they are assessed computationally in a stochastic approach in Sec. 3.2. Note that adjusting the hinge stiffness impacts directly the global structural stiffness, but a very low hinge stiffness may also allow for new vibration modes to appear, resulting in more freedom of deformation of the module. In all the analyses presented in the following a single model parameter is varied while keeping the others at their nominal values given for the reference model. The transformation load vs displacement curves of the module will therefore be reported in this work as well, when relevant.

#### 3. Modal analysis of bistable deployable structures

First the results for the reference pure Aluminum model, i.e. excluding manufacturing imperfections and supposed to be built with rigid hinges are presented first. In the present work, all the transformation load vs displacement curves are obtained with a displacements driven algorithm, where the vertical degree of freedom of the top node is controlled. The transformation load vs top central point vertical displacement curve is shown in Fig. 9, exhibiting clearly a bistable nature: the force required to fold the initially deployed structure increases first with increasing displacement, followed by a force decrease leading to a range

where this force is negative and ending at the compact folded equilibrium configuration with a zero external load required.



Figure 9: Transformation load vs displacement curves for different material pairs.

The presentation of the modal analysis results is limited to the first three modes in this work, since after careful numerical investigation these have been shown to be the most critical, having the lowest natural frequencies that can be close to the ones of possible excitations encountered in the service state (e.g. wind and/or machinery induced vibrations). For reference, the forth natural frequency of the pure Aluminum reference case is 3.8 times larger the third natural frequency, being therefore associated to higher energy deformation modes and less likely to be activated by usual loads occurring in civil engineering applications. The mode shapes are shown in different views on Fig. 10. The two first modes with the same natural frequencies are bending modes (bending around the x and y axes), followed by a third torsional mode (around the z axis). This torsional mode is of special interest, since a similar buckling deformation of the inner scissors is observed during the transformation [?]. Hence, exciting this torsional mode by external dynamic loading may lead to unwanted interference with the transformation behavior of the structure. These same modes, i.e. two bending modes with bending axes perpendicular to each other and a torsion mode characterize the modules studied in this work and the focus will

be therefore set on the variation of the natural frequencies of the structure.



(a) 1st vibration mode: bending around first diagonal.



(b) 2nd vibration mode: bending around second diagonal.



(c) 3rd vibration mode: torsion around vertical axis.

Figure 10: Top view of the first three vibration modes of the module in the deployed state.

The natural frequencies of the pure Aluminum reference case are reported in Tab. 2. It is noteworthy that the two first modes have the same natural frequency. This comes as a consequence of the symmetry of the model, as they are related to the bending of the module with respect to the two symmetry axes in the horizontal plane. The natural frequencies of the first three modes of the reference case are expected to be the highest among all studied cases of the same geometry without imperfections, since this setup corresponds to the stiffest material used for the beams that are linked with a rigid hinge (this is confirmed in Sec. 3.2). The obtained natural frequencies for the studied class of deployable modules falls in the range of usual civil engineering applications ([? ? ]).

f [Hz]			
Al	3.18	3.18	5.02
Al-HDPE	0.48	0.48	0.71

Table 2: First three natural frequencies of the structure built from different material pairs.

#### 3.1. Material and geometrical parameters

This section investigates the effect of changing the material to HDPE for the inner scissors that buckle during transformation and the variation of the square module's edge length and the height of the top center point. Fig. 9 shows the dominant impact of the HDPE inner scissors on the transformation response of the module. The peak force magnitude is roughly fifty-fold decreased as a result of the reduction of the elastic modulus of the inner scissors by a factor of 59.7, which clearly illustrates the nonlinear nature of the structural behavior during transformation. The eigenvalues of the module are reported in Tab. 2 and are 6.6 to 7 times lower than the pure Aluminum reference case. A variation of lesser degree in the natural frequencies than in the transformation peak force (both being related to stiffness quantities) can be explained partially by the decrease in the structural stiffness and in part by the difference between the linear modal analysis and the geometrically nonlinear transformation analysis

assumptions.

The first two natural frequencies are the same, corresponding to similar bending vibration modes (around perpendicular axes x and y) and the third vibration mode is torsional, as for the pure Aluminum module. The choice of the material pair thus impacts significantly the transformation and vibrational behavior of the module, as expected, and could be used to target a desired range of natural frequencies. In view of the very low transformation forces that make the structure sensitive to self deployment from the folded configuration, the pure Aluminum structure is studied exclusively in the following.

The transformation load vs top center point displacement curves as a function of the edge length L and of the top center point height h are plotted in Fig. 11. The influence of the geometric parameters on the transformation of the module has been studied in [?]. The present work focus on the influence of these parameters on the modal analysis of the structure and its link with the transformation peak load. Reducing the edge length leads to shorter inner scissor length and therefore a higher peak force required for their buckling, as it can be observed on Fig. 11a. Halving the edge length results in an increase of 200% in the peak load, while a 50% increase in L roughly halves it. When h is decreased, the magnitude of the geometrical incompatibility also reduces, which was shown to have a dominant effect on the transformation behavior in [?]. This trend is confirmed in Fig. 11b, i.e. reducing h to half results in a roughly tenfold decrease in the peak load and in a transformation curve of a different nature exhibiting a close to flat plateau in the positive range.

Apart from the transformation behavior, the natural frequencies of the module are also very sensitive to the variation of these geometrical parameters (Tab. 3). The vibration modes remain similar to the reference case and in the same order, independently of the value of L in the considered range. Reducing



Figure 11: Transformation load vs displacement curves for different geometrical parameters.

the edge length to half results in a 3.4 to 3.8 times increase in the first three natural frequencies, while increasing its value by 1.5 gives natural frequencies that are 41% to 46% of the ones of the reference case; these trends are in line with the one observed for the peak transformation load. Considering the magnitude of these variations, a special attention is thus required when the edge length of such a square bistable module is modified in the design, potentially requiring recomputing not only its transformation but also its dynamic behavior.

An interesting phenomena occurs when h equals 0.325 m, as shown in Fig. 12a:

h[m]	L[m]		f [Hz]	
0.250	1.00	2.89	3.81	3.81
0.325	1.00	3.66	3.66	3.73
0.500	1.00	3.16	3.16	5.02
0.500	0.50	12.17	12.17	17.35
0.500	1.50	1.29	1.29	2.32

Table 3: First three natural frequencies of the module for different values of the edge length L and of the top center point height h.

at this geometry the three first natural frequencies become the same, which from a practical point of view means that the three first vibration modes are excited simultaneously by a harmonic load at 3.7 Hz. This design should be avoided, because it could induce internal resonance and thereby large magnitude vibrations [? ? ? ]. It is noteworthy that for h < 0.325 the order of the vibration modes changes, i.e. the torsional vibration mode is first followed by the two bending modes of equal natural frequencies (Tab. 3). Fig. 12a shows the highly nonlinear evolution of the natural frequencies as a function of h exhibiting an abrupt change below 0.37 m and a smooth evolution above. It is emphasized that even though the variation of h, induces a small relative variation in the natural frequencies, it should not be disregarded in a rigorous design approach, since it also impacts the eigenmode order and hence the dynamic structural behavior. The observed small sensitivity of the natural frequencies accompanied by a substantial variation of the transformation peak force can be explained by the different phenomena governing these quantities: structural stiffness in the deployed state and buckling during the transformation, respectively.

The findings above illustrate the complex nonlinear nature of the dynamic behavior of bistable deployable structures and how natural frequencies of such BDS can be adjusted by varying two geometrical parameters of the design, e.g. decreasing the edge length increases the natural frequencies.



Figure 12: First three natural frequencies of the module as a function of different geometrical parameters.

#### 3.2. Finite stiffness joints and imperfect member length

The positive influence of having compliant joints in bistable deployable structures in order to decrease the impact of imperfections on the transformation behavior has been established in [?]. In this work finite joint stiffness effects on the *dynamic/vibrational* response of a single square BDS module with the reference geometrical parameter set as nominal values are considered, first in an idealized design geometry lacking manufacturing imperfections, second when incorporating beam length imperfections in the study.



Figure 13: Transformation load vs displacement curves for different hinge stiffness.

For the sake of completeness, the transformation force vs top center point displacement curves for soft and hard joints (as defined in Sec. 2.1.1) are compared to the reference case with rigid joints in Fig. 13. It is clearly apparent that incorporating compliance in the hinges is beneficial, since it results in the reduction of the peak transformation load, while the shape of the curve and the part in the negative force range remain practically the same. This implies that additionally to compensating for manufacturing imperfections in the transformation behavior, employing compliant hinges is a practically relevant way in a design for decreasing the force required for transforming BDS without changing the global nature of the curve (as opposed to the top center point height in Fig. 11).

Considering the natural frequencies, a logical decreasing trend is observed in Tab. 4 when the hinge compliance is increased. The natural frequencies of the soft joint case are 74% to 77% of the rigid case (for a peak force decrease of around 68%), which is a moderate variation when compared to the dominant effect of L (Sec. 3.1). It is noteworthy that no vibration mode order change was observed here.

In the following the coupled effect of manufacturing imperfections and joint

Hinge	f [Hz]		
Rigid	3.18	3.18	5.02
Hard	2.52	2.52	4.15
Soft	2.35	2.35	3.90

Table 4: First three natural frequencies of the module for different hinge stiffness.

stiffness on the first three natural frequencies of the BDS module in the deployed state are focused upon. The study is performed in a stochastic framework, considering the length of every beam in the structure as random variable with a uniform distribution. The range of these geometric imperfections in the uniform PDF is chosen to be -0.1 % and +0.1 %, similar to [?]. They are incorporated as fictitious thermal strains resulting in internal stresses in a preliminary step to the modal decomposition, as explained in Sec. 2. For achieving statistical representativity a total of 50.000 simulations were run for each studied case, i.e. rigid, hard and soft joints.

The resulting probability distribution functions of the first three natural frequencies are plotted in Fig. 14 and indicators of the stochastic data are given in Tab. 5, obtained by fitting a normal PDF to the raw data. As expected, the mean natural frequencies for the three first modes decrease with increasing joint compliance: they are 77%, 80%, 82% and 71%, 75% 77% of the reference case for the hard and soft joint, respectively. Note that the mean values are lower than the one obtained in a deterministic computation without geometrical imperfections (Tab. 4) and the mean natural frequencies of the first and second mode are not the same in the probabilistic approach. It is interesting to point out that the variance in the results systematically increases with the decrease of the joint stiffness, which is an expected result considering that a compliant joint accommodates easier manufacturing imperfections leading to a larger dispersion in the equilibrated configuration used as the basis of the modal analysis. Additionally, the variance in the third, torsional vibration mode appeared to have the largest

relative increase when incorporating finite joint stiffness in the model. It must be emphasized that the initial order of the first two bending modes followed by a third torsional mode was observed systematically in the stochastic study.

Mode	Mean [Hz]	Variance	
Rigid			
1	2.99	$1.95 \times 10^{-2}$	
2	3.30	$1.10 \times 10^{-2}$	
3	5.00	$4.96 \times 10^{-3}$	
Hard			
1	2.30	$3.03 \times 10^{-2}$	
2	2.66	$1.42 \times 10^{-2}$	
3	4.12	$8.37 \times 10^{-3}$	
Soft			
1	2.13	$3.24 \times 10^{-2}$	
2	2.49	$1.47 \times 10^{-2}$	
3	3.86	$9.59 \times 10^{-3}$	

Table 5: Statistical data on the first three natural frequencies of the module for different hinge stiffness.

The findings above show that employing compliant hinges in BDS can be recommended as a practical means to (i) decrease the peak transformation force, (ii) manage the spurious effects of geometrical imperfections on the transformation behavior [?], while (iii) maintaining a similar vibrational behavior than when using rigid hinges, only resulting in the decrease of the mean natural frequencies and in an expected increase in the variance in the results when the hinge stiffness is decreased.

#### 4. Conclusions

A computational study was presented in which the natural frequencies and vibration modes of a single square BDS module were evaluated, incorporating the effects of the variation of particular design parameters (material choice, geometric parameters, finite joint stiffness) and manufacturing imperfections. The first three vibration modes were two bending and a third torsional mode in the



Figure 14: Probability distribution of the natural frequencies of the first three vibration modes when incorporating random member length imperfections for different hinge stiffness.

vast majority of the cases. The torsional mode was observed to be similar to the deformation mode of the structure when transformation takes place and the inner scissors buckle. The first three natural frequencies of the pure Aluminum structure were found between 2 Hz and 5 Hz. The study allowed to identify the dominant influence of the material pair choice for inner and outer scissors and of the edge length of the square BDS module. The top center point height was observed to affect the natural frequencies to a lesser extent, however decreasing this parameter led to a reversal in the order of the vibration modes, i.e. the torsional mode became the one with the lowest natural frequency. Incorporating finite joint stiffness led to the decrease of the natural frequencies, as expected. When it was coupled to taking geometric imperfections into account in a probabilistic approach, it resulted in a larger variance in the results, while maintaining a dynamic response similar to the reference case with rigid hinges.

The presented results illustrate the complex nonlinear nature of the dynamic behavior of bistable deployable structures and show how their natural frequencies can be adjusted by making practical design choices. Incorporating compliant hinges in the design can be clearly recommended, since (i) it was shown to compensate for manufacturing imperfections in the transformation behavior, (ii) it is a practically relevant way in a design for decreasing the force required for transforming BDS, while (iii) resulting in a moderate decrease in the natural frequencies and keeping a dynamic behavior similar to the one for rigid hinges. Examples of realistic hinges incorporating two bushings for BDS were presented and the procedure to derive the corresponding hinge FE stiffness was explained. Another practical recommendation is to pay a special attention to the edge length variation in iterative BDS design, potentially requiring recomputing not only the transformation but also the dynamic behavior.

The present work leaves of course space for future improvement. Note that

establishing more realistic support conditions in the modal analysis to mimic the effect of other neighboring modules in a modular structure (e.g. by adding an equivalent stiffness) is part of future work, as well as considering larger, multimodular structures. The present work focuses on a pure numerical approach for obtaining the modal properties of the system (natural frequencies and vibration modes) yielding results of practical interest, that could indeed be verified and coupled to an experimental study in a future work. Addressing the transformation (deployment) behavior of BDS, incorporating friction as energy dissipation mechanism is also part of future work of immediate interest. For this, a full nonlinear dynamic analysis is required and the effects of snap-through and instabilities, present in bistable deployable structures, must be investigated.

### **Data Availability Statement**

The data that support the findings of this study are available from the corresponding author upon reasonable request.

#### References

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