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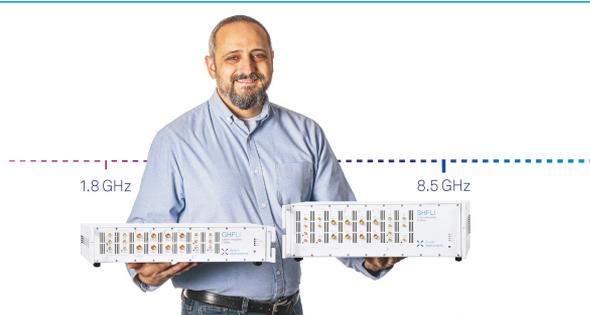
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ABSTRACT

X-ray photodesorption yields of $^{15}\text{N}_2$ and ^{13}CO are derived as a function of the incident photon energy near the N (~ 400 eV) and O K-edge (~ 500 eV) for pure $^{15}\text{N}_2$ ice and mixed $^{13}\text{CO}:\text{N}_2$ ices. The photodesorption spectra from the mixed ices reveal an indirect desorption mechanism for which the desorption of $^{15}\text{N}_2$ and ^{13}CO is triggered by the photoabsorption of ^{13}CO and $^{15}\text{N}_2$, respectively. This mechanism is confirmed by the x-ray photodesorption of ^{13}CO from a layered $^{13}\text{CO}/^{15}\text{N}_2$ ice irradiated at 401 eV on the N $1s \rightarrow \pi^*$ transition of $^{15}\text{N}_2$. This latter experiment enables us to quantify the relevant depth involved in the indirect desorption process, which is found to be 30–40 monolayers in that case. This value is further related to the energy transport of Auger electrons emitted from the photoabsorbing $^{15}\text{N}_2$ molecules that scatter toward the ice surface, inducing the desorption of ^{13}CO . The photodesorption yields corrected from the energy that can participate in the desorption process (expressed in molecules desorbed by eV deposited) do not depend on the photon energy; hence, they depend neither on the photoabsorbing molecule nor on its state after Auger decay. This demonstrates that x-ray induced electron stimulated desorption, mediated by Auger scattering, is the dominant process explaining the desorption of $^{15}\text{N}_2$ and ^{13}CO from the ices studied in this work.

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I. INTRODUCTION

Diatomic molecules such as CO and N_2 have been extensively studied in the laboratory at very low temperature ($T < 20$ K), in their solid phase. Indeed, these molecules are expected to play an important role in astrophysical environments. For instance, they are observed directly (especially for CO) or indirectly (especially for N_2) in the interstellar medium (ISM) and the solar system in a diversity of icy solids, such as interstellar ices (Boogert *et al.*, 2015), comets (Mumma and Charnley, 2011; Rubin *et al.*, 2015), and icy moons and planets (Owen *et al.*, 1993; Cruikshank *et al.*, 1993; Bennett *et al.*, 2013). Interaction of these icy solids with photons coming from nearby or distant stars can be at the origin of exchanges between solid and gas phases via the so-called photodesorption phenomenon. Photodesorption is a nonthermal process through which the energy of incident photons is absorbed by the ice and partly converted to the desorption of molecules

from its surface. In the last decades, it has been extensively studied in the vacuum ultraviolet (VUV) range (7–13.6 eV) for N_2 and CO ices (Öberg *et al.*, 2007, 2009; Muñoz Caro *et al.*, 2010; Fayolle *et al.*, 2011, 2013; Bertin *et al.*, 2013).

More recently, the desorption induced by the irradiation of x rays, referred to as x-ray photodesorption, was also highlighted as a significant process for ices composed of simple (Dupuy *et al.*, 2018, 2021b) and complex molecules (Jiménez-Escobar *et al.*, 2018; Ciaravella *et al.*, 2020; Basalgète *et al.*, 2021a). More importantly, it was suggested that x-ray photodesorption of neutral molecules was strongly dependent on the ice composition via indirect desorption mechanisms (Basalgète *et al.*, 2021b). A detailed comprehension of these indirect processes on a microscopic scale is of primary importance to understand their implication in astrophysical environments. For a proper modeling of such processes, it is also required to access quantitative data concerning desorption yields.

The x-ray photodesorption from condensed N_2 and CO has been previously studied for thin layers deposited on top of transition and noble metals (Frigo *et al.*, 1998; Feulner *et al.*, 2000; Romberg *et al.*, 2000; Feulner *et al.*, 2002). In that case, desorption of neutral and ionic molecules (N_2 , CO, CO^+) as well as fragments (N , N^+ , N^{2+} , C, O, O^+ , C^+) was found to strongly depend on singly or multiply valence-excited states, after decay of the core hole state. These states were associated with specific desorption mechanisms. In some cases, ultrafast processes occurring during the core hole state lifetime were also suggested to play a role for the desorption of ionic fragments (Frigo *et al.*, 1998).

For thick molecular ices, characterized by a higher number of layers deposited onto a substrate (generally more than 50 layers), the processes leading to desorption may be different. For instance, after x-ray absorption by molecules in the ice, the relaxation of the core hole state results in the emission of an Auger electron with a probability close to 1 for low Z elements (Walters and Bhalla, 1971; Krause, 1979). This Auger electron carries most of the initial photon energy and scatters inelastically in the ice, creating secondary events (e.g., ionization and excitation of the surrounding molecules) and, especially, a cascade of secondary low energy (<20 eV) electrons, which may lead to the desorption of molecules from the ice surface. This process is known as x-ray induced electron stimulated desorption (XESD). XESD was not suggested to be the dominant process for the x-ray photodesorption of ions from pure N_2 and CO ice (Feulner *et al.*, 1992; Dupuy *et al.*, 2021b). However, for the desorption of neutrals, it was assumed to be the case for pure CO ice (Dupuy *et al.*, 2021b), pure H_2O ice (Dupuy *et al.*, 2020), and pure and mixed methanol-containing ices (Basalgète *et al.*, 2021a,b).

Based on quantitative data, x-ray photodesorption of ions from molecular ices is expected to be negligible compared to the desorption of neutrals (Dupuy *et al.*, 2018, 2021b). In the present work, we experimentally study the x-ray photodesorption process for neutral desorption from pure $^{15}N_2$ ice and from mixed and layered $^{13}CO:^{15}N_2$ ices (isotopes are chosen to differentiate the molecules by mass spectrometry). Resonant 1s core excitation (or ionization) near the N (~ 400 eV) or O (~ 500 eV) K-edge can be achieved thanks to tunable and high spectral resolution x rays from synchrotron facilities, allowing thereby to selectively photoexcite $^{15}N_2$ or ^{13}CO to trigger desorption from ices containing these molecules. We first provide a quantitative study of the x-ray photodesorption of neutral molecules from pure $^{15}N_2$ ice. Then, we quantify this process for mixed $^{13}CO:^{15}N_2$ and layered $^{13}CO/^{15}N_2$ (^{13}CO on top of $^{15}N_2$) ices. This latter set of experiments brings important information to discuss the desorption mechanisms.

II. EXPERIMENT AND METHODOLOGY

A. Setup and x-ray beamline

To run the experiments, the Surface Processes and ICES (SPICES) setup was connected to the SEXTANTS beamline of the SOLEIL synchrotron facility at Saint-Aubin, France (Sachi *et al.*, 2013). SPICES consists of an ultrahigh vacuum (UHV) chamber (base pressure of $\sim 10^{-10}$ mbar) at the center of which a copper substrate (polycrystalline oxygen-free high-conductivity copper, surface area of 1.5×1.5 cm²) can be cooled down to ~ 15 K by a closed-cycle Helium cryostat. The substrate is electrically insulated

from its sample holder by a Kapton foil. This enables the measurement of the drain current generated by electrons escaping the ice after x-ray absorption, which we refer to as the total electron yield (TEY) in the following. The TEY can be measured while scanning the energy of the incident photons and it is representative of the x-ray absorption profile of the ices studied in this work. It is expressed in electrons per incident photon (displayed as e^- photon⁻¹ for more simplicity) by correcting the drain current from the photon flux.

The photon flux is measured by a calibrated silicon photodiode mounted on the beamline. The SEXTANTS beamline provides x-ray photons, the energy of which can be finely tuned by means of a monochromator. The energy resolution can also be controlled by adjusting the slit apertures at the exit of the monochromator. In the present study, we have used photons in the 390–460 eV range (nitrogen K-edge) and in the 530–560 eV range (oxygen K-edge). Spectral resolutions of ~ 100 and ~ 350 meV were used to run the experiments. The photon flux varies with the energy and the spectral resolution between 10^{11} and 10^{12} photon s⁻¹. The beam was sent at a 47° incidence relative to the normal of the substrate surface and the spot area at the surface was ~ 0.1 cm².

The calibration of the energy scale of the beamline was performed near the N K-edge and the O K-edge as follows:

- Near the N K-edge, the TEY was measured on a pure $^{15}N_2$ ice as a function of the energy of the incident photons and it was compared to gas phase x-ray absorption spectroscopy experiments (e.g., Chen *et al.*, 1989; Feifel *et al.*, 2004). The TEY peak corresponding to the $N\ 1s \rightarrow \pi^*$ ($v' = 0$) transition of $^{15}N_2$ was set to 400.868 eV according to Chen *et al.* (1989) (see Fig. 1 and Sec. III A). We, therefore, assumed that this energy position was not significantly shifted when going from gas phase to solid phase, as observed for a similar molecule: CO (Jugnet *et al.*, 1984). Note that the energy of the transition is slightly varying (± 20 meV) from one gas phase experiment to another (King *et al.*, 1977; Hitchcock and Brion, 1980; Sodhi and Brion, 1984; Kempgens *et al.*, 1996; Feifel *et al.*, 2004; Kato *et al.*, 2007).
- Near the O K-edge, the $O\ 1s \rightarrow \pi^*$ transition of ^{13}CO observed on the TEY of a mixed $^{13}CO:^{15}N_2$ ice was centered at 534.4 eV (see Fig. 2 and Sec. III A) according to pure CO ice spectroscopy experiments (Jugnet *et al.*, 1984), assuming that the intermolecular interactions between ^{13}CO and $^{15}N_2$ are not significantly shifting the energy position of the electronic transition, which is reasonable considering such weakly interacting solids.

B. Ice deposition

Within the SPICES setup, a tube connected to a gas introduction system, which can be translated a few millimeters in front of the substrate surface, enables the injection of a partial pressure of gas phase molecules that stick on the cold surface to form the molecular ices without significantly modifying the base pressure of the UHV chamber. The gaseous molecules used during this work were molecular nitrogen $^{15}N_2$ (98.6% ^{15}N purity, Eurisotop) and carbon monoxide ^{13}CO (99% ^{13}C purity, Eurisotop). The ice thickness is expressed in monolayer (ML), equivalent to a molecule surface

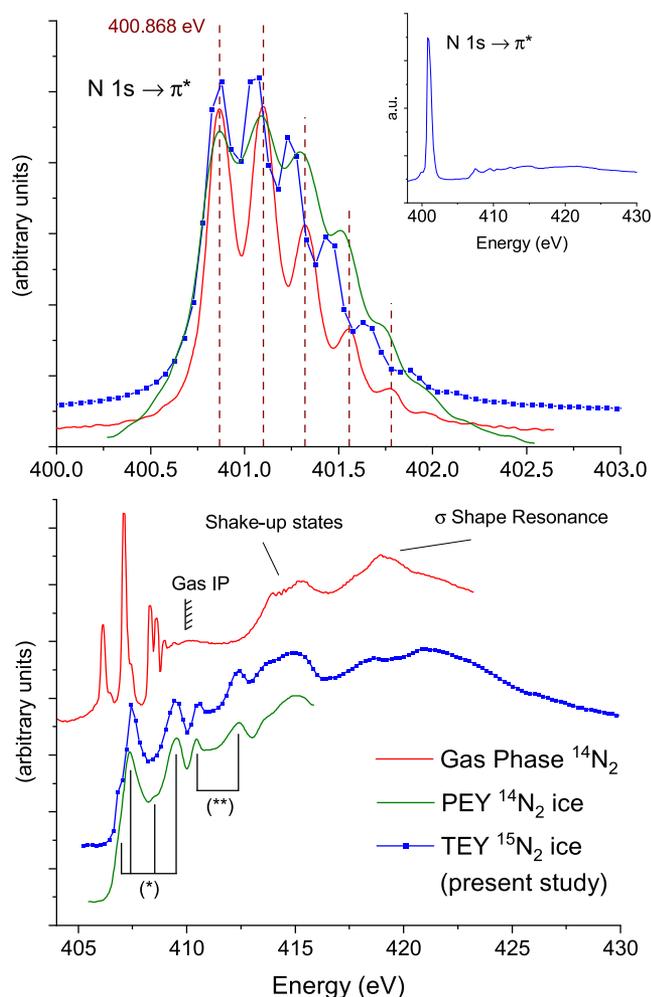


FIG. 1. Upper right panel: TEY of $^{15}\text{N}_2$ ice at 15 K (50 ML) from our experiments in the N K-edge region between 398 and 430 eV. The upper left panel shows the N $1s \rightarrow \pi^*$ excitation region and the lower panel shows the region near the ionization threshold between 405 and 430 eV. In solid red lines are the x-ray photoabsorption spectra of gas phase $^{14}\text{N}_2$ from [Chen et al. \(1989\)](#) with a spectral resolution of 40 meV. In solid green lines are the partial electron yields (PEYs) of 50 ML of $^{14}\text{N}_2$ ice at 15 K from [Feulner et al. \(1992\)](#) with a spectral resolution of 100 meV. In blue lines are the TEYs of 50 ML of $^{15}\text{N}_2$ ice at 15 K from our experiments with a spectral resolution of 100 meV for the upper left panel and 350 meV for the lower and upper right panels. The baselines of the different curves have been shifted for more clarity. The red dashed vertical lines in the upper left panel are centered on the gas phase vibrational bands.

density of $\sim 10^{15}$ molecule cm^{-2} . Prior to synchrotron experiments, the temperature-programmed desorption (TPD) technique was used to calibrate the rate and deposition time needed to grow a targeted ice thickness. The relative uncertainty on the number of ML deposited during a deposition procedure is about 10% (see, e.g., [Doronin et al., 2015](#)).

The ice samples studied in this work were deposited at 15 K. The mixed $^{13}\text{CO}:^{15}\text{N}_2$ ices were grown by controlling *in situ* the

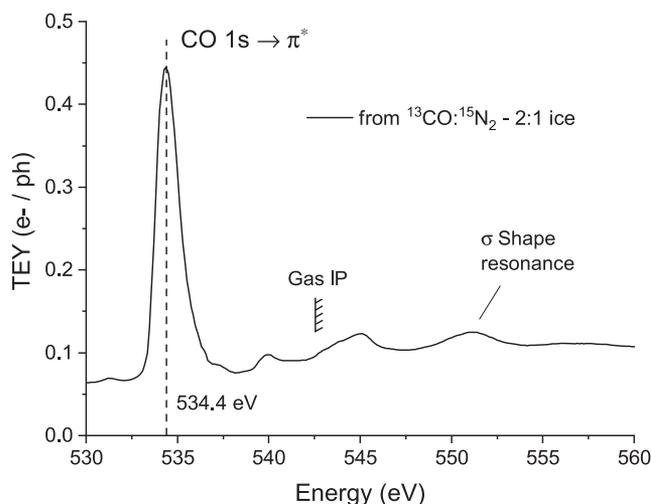


FIG. 2. TEY of a mixed $^{13}\text{CO}:^{15}\text{N}_2 - 2:1$ ice at 15 K (total of 100 ML) in the O K-edge region between 530 and 560 eV from our experiments, with a spectral resolution of 350 meV.

partial pressure of each species injected in the UHV chamber independently (two different valves connected to the deposition tube were used). In this way, the mixing ratio is set by adjusting the partial pressures corresponding to each species. The experiment regarding the layered $^{13}\text{CO}/^{15}\text{N}_2$ ice (see Sec. III B 3) was achieved by first depositing 50 ML of $^{15}\text{N}_2$ at 15 K. Then, on top of it, a given number of ML of ^{13}CO was deposited and the ice was irradiated to measure the photodesorption yields. This operation (deposition of ^{13}CO on top and irradiation) was sequentially repeated on the same ice until reaching ~ 90 ML of ^{13}CO on top of the $^{15}\text{N}_2$ ice. The number of ML of ^{13}CO deposited at each step is displayed in [Fig. 5](#) (see Sec. III B 3).

C. Derivation of the photodesorption yields

The photodesorbed molecules are probed in the gas phase during the x-ray irradiation of the ice by means of a quadrupole mass spectrometer (QMS) equipped with an electron impact ionization chamber at 70 eV. The photodesorption of ^{13}CO and $^{15}\text{N}_2$ was monitored by recording the QMS signal on the mass channel 29 and 30 u, respectively. The signal intensities are divided by the photon flux and multiplied by a factor k to obtain the x-ray photodesorption yields Γ_{raw} expressed in molecule desorbed per incident photon (displayed as molecules photon^{-1} in the following for more simplicity). The factor k was derived prior to the synchrotron runs to relate the QMS current to a calibrated number of $^{15}\text{N}_2$ molecules desorbed during a TPD experiment. For more details on the calibration procedure, see the work of [Basalgète et al. \(2021b\)](#). The apparatus function of our QMS can be considered constant between the mass channels 29 and 30 u and the non-dissociative electron impact ionization cross sections at 70 eV of ^{13}CO ([Tian and Vidal, 1998](#)) and $^{15}\text{N}_2$ ([Straub et al., 1996](#)) are nearly equal such that we considered the detection efficiency of our QMS to be similar for ^{13}CO and $^{15}\text{N}_2$. This results in using a similar factor k for quantifying the photodesorption yields of ^{13}CO and $^{15}\text{N}_2$. The systematic uncertainty

on the photodesorption yields provided by this calibration method is estimated to be $\sim 50\%$.

In Sec. IV D, we compute the photodesorption yields Γ_{eV} expressed in molecule desorbed per eV deposited by correcting the experimental photodesorption yields Γ_{raw} (expressed in molecules photon⁻¹) considering the absorption cross section. This results in the following formula:

$$\Gamma_{eV}(E) = \frac{\Gamma_{raw}(E)}{E a_{dil} (1 - e^{-\sigma(E)N})}, \quad (1)$$

where E is the energy of the incident photon, a_{dil} is a factor taking into account the dilution of the photodesorbing molecules in the case of mixed ices (i.e., $a_{dil} = 1/3$ and $2/3$ for the photodesorption yield of, respectively, ¹⁵N₂ and ¹³CO from a mixed ¹³CO:¹⁵N₂ – 2:1 ice), $\sigma(E)$ is the photoabsorption cross section at E , and N is the column density involved in the desorption process. N is equal to $N = \Lambda_{des} \times b_{dil} \times 10^{15}$ cm⁻², where Λ_{des} is the desorption-relevant depth expressed in ML (this notion is explained in more detail in Sec. IV), and b_{dil} is a factor taking into account the dilution of the photoabsorbing molecule in the case of mixed ices (i.e., $b_{dil} = 1/3$ and $2/3$ for the yields derived, respectively, near the N and O K-edge).

III. RESULTS

A. X-ray absorption spectra

1. ¹⁵N₂ ice photoabsorption near N K-edge

The x-ray absorption spectrum of pure ¹⁵N₂ ice at 15 K near the N K-edge, assimilated to our TEY measurements, is displayed in Fig. 1 (blue lines) for different energy ranges and spectral resolutions. The absorption is dominated by the N 1s $\rightarrow \pi^*$ transition of ¹⁵N₂ near 401 eV, whose vibrational structure has been resolved with a spectral resolution of 100 meV. Our TEY in the N 1s $\rightarrow \pi^*$ region (see upper panel of Fig. 1) compares well with other absorption spectra of gas phase ¹⁴N₂ (Chen *et al.*, 1989) and ¹⁴N₂ ice at 15 K (Feulner *et al.*, 1992). The N 1s⁻¹ π^* core electronic state of the free molecule is not significantly modified in the solid phase, as expected for weak van der Waals interactions between N₂ molecules in the solid. There are, however, slight discrepancies in the position of the vibrational levels between the three curves in the upper panel of Fig. 1. When going from gas phase (red line) to solid phase (green line), the small differences in the position of the vibrational levels of the N 1s⁻¹ π^* state of ¹⁴N₂ could be attributed either to experimental uncertainties on the relative position of the peaks or to intermolecular interactions in the solid phase. The comparison of our TEY of ¹⁵N₂ ice (blue line) with the PEY of ¹⁴N₂ ice (green line) shows a smaller vibrational level spacing for ¹⁵N₂. This is expected to come from an isotopic effect. A similar behavior is observed in the vibrational level spacing of the A¹Π state of ¹³CO vs ¹²CO ice and of the b¹Π_u state of ¹⁵N₂ vs ¹⁴N₂ ice probed by vacuum ultraviolet (VUV) photodesorption experiments (Bertin *et al.*, 2012; Fayolle *et al.*, 2013). In the lower panel of Fig. 1, the Rydberg states of gas phase ¹⁴N₂ seen below the ionization potential (IP) are significantly modified in the solid phase, as expected from the overlapping of the Rydberg orbitals with that of the neighboring molecules in the solid. Instead, series of broad peaks [labeled (*)

and (**) in Fig. 1], whose positions go above the gas phase IP, populate this region. There is no significant difference in this region between ¹⁴N₂ and ¹⁵N₂ in the solid phase. The first series of the four peaks labeled (*) has been attributed to the equivalent of the gas phase Rydberg states 3sσ, 3pπ, 4sσ, and 5sσ in solid phase by Feulner *et al.* (1992), based on polarization dependent x-ray photodesorption of N⁺ and based on the data from Chen *et al.* (1989). The attribution of the peaks labeled (**) above the gas phase IP in Fig. 1 has not been studied in the literature to our knowledge. The gas phase broad feature observed near 415 eV has been attributed by Chen *et al.* (1989) to shake-up states, tentatively associated with the state configurations (1s)⁻¹(5σ)⁻¹(1π_g^{*})⁻¹(3sσ)⁻¹ and (1s)⁻¹(1π)⁻¹(1π_g^{*})⁻¹(3pπ)⁻¹. As one can see in Fig. 1, a similar feature is observed in this region for ¹⁵N₂ and ¹⁴N₂ ice. The feature observed near 419 eV in the gas phase has been attributed to a σ shape resonance (Chen *et al.*, 1989). A similar feature is observed for ¹⁵N₂ ice.

2. ¹³CO:¹⁵N₂ ice photoabsorption near O K-edge

The TEY of a mixed ¹³CO:¹⁵N₂ ice (2:1) at 15 K near the O K-edge is displayed in Fig. 2. The absorption is dominated by the O 1s $\rightarrow \pi^*$ transition of ¹³CO centered at 534.4 eV. The spectrum is similar to that of a pure CO ice studied in similar experiments and for which more details are available regarding the attribution of the spectral features (Dupuy *et al.*, 2021b). The intermolecular interactions between ¹³CO and ¹⁵N₂ in the mixed ice do not modify the transitions observed in this energy range compared to a pure CO ice. The N 1s core electron photoionization of ¹⁵N₂ is expected to contribute to the absorption spectrum in this energy range, but it is reasonably assumed to be negligible compared to ¹³CO photoabsorption, and it should in any case account for a monotonically decreasing signal over the whole energy range.

B. X-ray photodesorption yields

1. ¹⁵N₂ pure ice

The photodesorption spectra of ¹⁵N₂ from a pure ¹⁵N₂ ice at 15 K are shown in Fig. 3 (green lines), with the upper panel corresponding to the N 1s $\rightarrow \pi^*$ transition and the lower panel corresponding to the near-IP region. The TEYs measured simultaneously during the irradiation procedure are also displayed in blue lines for comparison. The photodesorption yields are higher by more than one order of magnitude in the N 1s $\rightarrow \pi^*$ region than in the near-IP region, with a maximum of ~ 5 molecules photon⁻¹ at the N 1s $\rightarrow \pi^*$ ($v' = 0, 1$) transitions. The photodesorption spectra follow the same variations as the TEYs except for small discrepancies on the relative intensities of the features, which are likely due to background correction issues and for which we do not attribute any physical meaning. The photodesorption yields of ¹⁵N₂ and the TEYs are obtained from an ice having received a fluence $< 3.10^{16}$ ph cm⁻². They both decrease with increasing fluence due to the ice photoaging, that is, the structural and chemical modifications of the ice due to photoabsorption, but the shape of the spectra remains the same and no significant new absorption features are observed in the energy range probed (400–450 eV). The possible photodesorption of ¹⁵N₃ and ¹⁵N₄ was monitored during irradiation, but no desorption signals were found (with a detection limit of

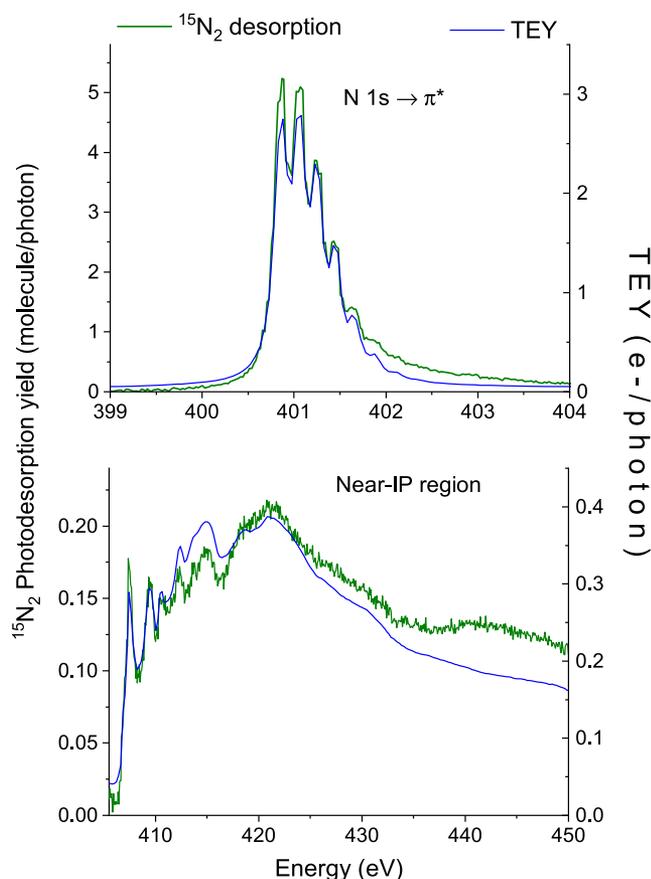


FIG. 3. X-ray photodesorption yields of $^{15}\text{N}_2$ (green lines) from pure $^{15}\text{N}_2$ ice at 15 K (50 ML) in the $\text{N } 1s \rightarrow \pi^*$ excitation region (upper panel, spectral resolution of 100 meV) and in the near-IP region (lower panel, spectral resolution of 350 meV). The TEYs measured during the irradiation are also displayed for comparison (in blue lines).

10^{-2} molecules photon $^{-1}$ considering the mass channels and photon flux). Photodesorption of atomic ^{15}N was only detectable (with a detection limit of 10^{-3} molecules photon $^{-1}$) at the $\text{N } 1s \rightarrow \pi^*$ transition, with a yield of ~ 0.01 molecules photon $^{-1}$, which is two orders of magnitude less than the desorption yield of $^{15}\text{N}_2$.

2. Mixed $^{13}\text{CO}:^{15}\text{N}_2$ ice

The photodesorption yields of ^{13}CO and $^{15}\text{N}_2$ from mixed $^{13}\text{CO}:^{15}\text{N}_2$ (2:1 and 20:1) ices are displayed in Fig. 4 in the N K-edge and O K-edge energy range. The TEYs measured simultaneously during the irradiation procedures are also displayed in dashed lines and vertically shifted for more clarity. The shape of the photodesorption spectra is well-correlated to the TEYs both at the N and O K-edge, similar to what is observed for a pure $^{15}\text{N}_2$ ice. There is, however, a small discrepancy: In the lower panels of Fig. 4, the photodesorption yield of $^{15}\text{N}_2$ from a mixed $^{13}\text{CO}:^{15}\text{N}_2$ (2:1) ice is found to decrease after 403 eV for the N K-edge and after 555 eV for the O K-edge whereas the corresponding TEY is constant. This is due to background correction issues.

The photodesorption of $^{15}\text{N}_2$ from mixed $^{13}\text{CO}:^{15}\text{N}_2$ ices is found to significantly decrease between the 2:1 and 20:1 ratio at both the N and O K-edges (lower panels of Fig. 4) due to the high dilution of $^{15}\text{N}_2$, whereas ^{13}CO desorption yield is in any case found in the same order of magnitude (upper panels). In the case of a mixed $^{13}\text{CO}:^{15}\text{N}_2$ (2:1) ice, the photodesorption yields of ^{13}CO and $^{15}\text{N}_2$ are found in the same order of magnitude at the $\text{N } 1s \rightarrow \pi^*$ ($\sim 1.4 - 1.5$ molecules photon $^{-1}$) and the $\text{O } 1s \rightarrow \pi^*$ ($\sim 0.5 - 0.7$ molecules photon $^{-1}$) transitions of $^{15}\text{N}_2$ and ^{13}CO , respectively. This is consistent with the fact that CO and N_2 have similar binding energies in mixed CO: N_2 ices (Öberg *et al.*, 2005; Bisschop *et al.*, 2006). The fact that the photodesorption yields are lower at the $\text{O } 1s \rightarrow \pi^*$ transition compared to the $\text{N } 1s \rightarrow \pi^*$ transition in the case of the 2:1 ratio is probably due to a higher absorption cross section (based on gas phase experiments) of $^{15}\text{N}_2$ at 401 eV [~ 26 Mbarn from the work of Kato *et al.* (2007) and when considering the spectral resolution of our photon source] compared to that of ^{13}CO at 534.4 eV [$\sim 2.8 - 3.2$ Mbarn from the work of Barrus *et al.* (1979)].

Despite a higher photoabsorption of ^{13}CO expected for a dilution ratio of 20:1 compared to 2:1, the TEY value at the $\text{O } 1s \rightarrow \pi^*$ transition of ^{13}CO is lower for the 20:1 ratio (see right panels of Fig. 4). We would have expected the photodesorption yield of ^{13}CO to also follow this trend. However, the photodesorption yield of ^{13}CO at this energy does not change significantly from a 2:1 to a 20:1 ratio (the yield is at ~ 0.7 molecules photon $^{-1}$). A higher fluence received by the $^{13}\text{CO}:^{15}\text{N}_2$ (20:1) ice before this specific measurement seems to be responsible for this behavior. The fluences corresponding to the other measurements displayed in Fig. 4 are similar ($< 2 \times 10^{16}$ ph cm $^{-2}$).

Photodesorption signals were observed on the mass channels 28 u and 44 u. They could correspond to the desorption of $^{13}\text{C}^{15}\text{N}$ and $^{13}\text{C}^{15}\text{NO}$, respectively. However, the signals were too weak ($\text{S/N} \gtrsim 1$, with a photodesorption yield of ~ 0.02 molecules photon $^{-1}$ at the $\text{N } 1s \rightarrow \pi^*$ and the $\text{O } 1s \rightarrow \pi^*$ for both mass channels) and more easily detected in the case of a 20:1 ratio such that we cannot totally exclude the contribution of ^{12}CO , $^{14}\text{N}_2$ (mass 28 u), and $^{12}\text{CO}_2$ (mass 44 u) photodesorption to these mass channels due to natural isotopes in our gas samples ($\sim 1\%$). Finally, a desorption signal was found on the mass channel 31 u, which is associated with the x-ray photodesorption of ^{15}NO . The photodesorption spectra of ^{15}NO are also following the same variations as the TEYs and the corresponding yields are 5.3×10^{-2} molecules photon $^{-1}$ at the $\text{N } 1s \rightarrow \pi^*$ transition of $^{15}\text{N}_2$ and 2.7×10^{-2} molecules photon $^{-1}$ at the $\text{O } 1s \rightarrow \pi^*$ transition of ^{13}CO for a $^{13}\text{CO}:^{15}\text{N}_2$ (2:1) ice. The formation of ^{15}NO is expected to originate from chemistry induced by the dissociation of $^{15}\text{N}_2$ and ^{13}CO . Its presence in the ice has also been evidenced, in very small quantities, by TPD of the photo-irradiated ice at the end of the experiment. The photodesorption yields of the previous photo-products (possibly $^{13}\text{C}^{15}\text{N}$ and $^{13}\text{C}^{15}\text{NO}$ for the mass 28 and 44 u, respectively, and unambiguously ^{15}NO for the mass 31 u) are at least an order of magnitude lower than the yields of ^{13}CO and $^{15}\text{N}_2$. This shows that, if some photochemical change occurs, its impact on the photodesorption process from mixed ices is limited. Photodesorption of atomic ^{15}N , ^{13}C , and O was not detected (with a detection limit of 10^{-3} molecules photon $^{-1}$).

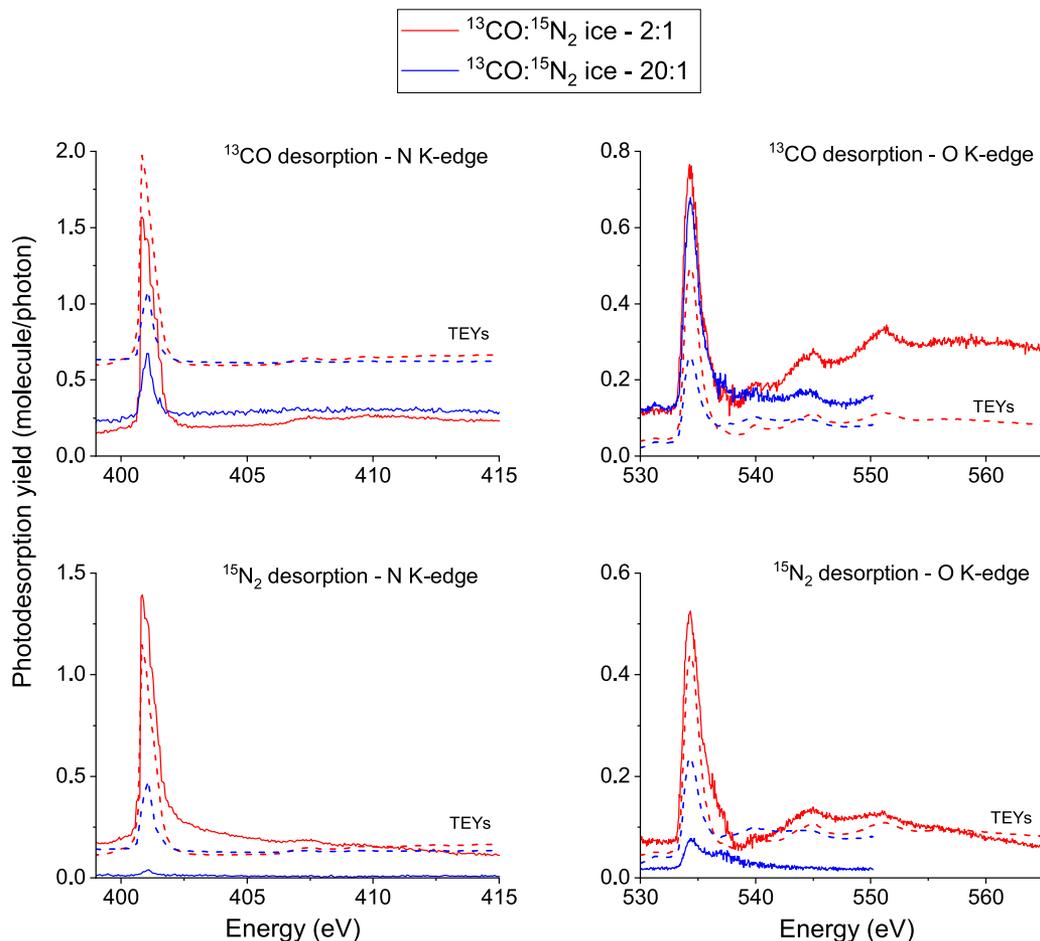


FIG. 4. X-ray photodesorption yields of ^{13}CO and $^{15}\text{N}_2$ (solid lines) in the N K-edge (left panels) and O K-edge (right panels) regions from mixed $^{13}\text{CO}:^{15}\text{N}_2$ ices at 15 K (total of 100 ML) with mixing ratio of 2:1 (red lines) and 20:1 (blue lines). The TEYs measured simultaneously during the corresponding irradiation procedure are also displayed in dashed lines in arbitrary units.

3. Layered $^{13}\text{CO}/^{15}\text{N}_2$ ice

The photodesorption yields at 401 eV of ^{13}CO and $^{15}\text{N}_2$ from a layered $^{13}\text{CO}/^{15}\text{N}_2$ ice are displayed in Fig. 5 as a function of the number of ML of ^{13}CO on top of a $^{15}\text{N}_2$ ice. For this experiment, the photon flux was low ($\sim 10^{11}$ photon s^{-1}) and the irradiation time was short (a few seconds) in order to ensure that the number of ^{13}CO molecules photodesorbed during one irradiation step is negligible compared to the number of ML of ^{13}CO deposited on top of the $^{15}\text{N}_2$ ice. Considering an x-ray photodesorption yield of ^{13}CO between 1.0 and 2.4 molecules photon $^{-1}$, an equivalent of 0.01–0.03 ML of ^{13}CO is desorbed during one irradiation step. In comparison, the first point measured in Fig. 5 is at 1.4 ML of ^{13}CO on top of the $^{15}\text{N}_2$ ice. We do not expect significant structural changes of the ice during each irradiation step at such low fluence. Chemical changes are also not expected based on the fact that no signs of significant chemistry are observed on both the TEYs and in desorption for the previous pure and mixed irradiated ices.

At 401 eV, photoabsorption of $^{15}\text{N}_2$ is due to the $\text{N } 1s \rightarrow \pi^*$ transition and photoabsorption of ^{13}CO is due to the ionization of its C 1s electron (IP ~ 296 eV). Quantitatively speaking, and based on gas phase data, the absorption cross section of $^{15}\text{N}_2$ and ^{13}CO at 401 eV is ~ 20 – 40 and ~ 0.3 Mbarn, respectively (Barrus *et al.*, 1979; Kato *et al.*, 2007). Hence, photoabsorption at 401 eV in our layered $^{13}\text{CO}/^{15}\text{N}_2$ ice is dominated by $^{15}\text{N}_2$ photoabsorption. The idea of this experiment is, therefore, to trigger the desorption of the above lying ^{13}CO molecules by photoexcitation of the supporting $^{15}\text{N}_2$ molecules and to estimate the depth of this indirect desorption mechanism (see Sec. IV) by sequentially depositing monolayers of ^{13}CO on top of the ice and irradiate it. The contribution of ^{13}CO photoabsorption to the desorption of ^{13}CO at 401 eV can be estimated by considering that it is quantitatively similar to the photodesorption of CO from a pure CO ice at 530 eV, which exhibits a photodesorption yield of ~ 0.1 molecules photon $^{-1}$ (Dupuy *et al.*, 2021b). This stems from the fact that the photoabsorption of ^{13}CO at 530 eV is also due to the ionization of its C 1s electron

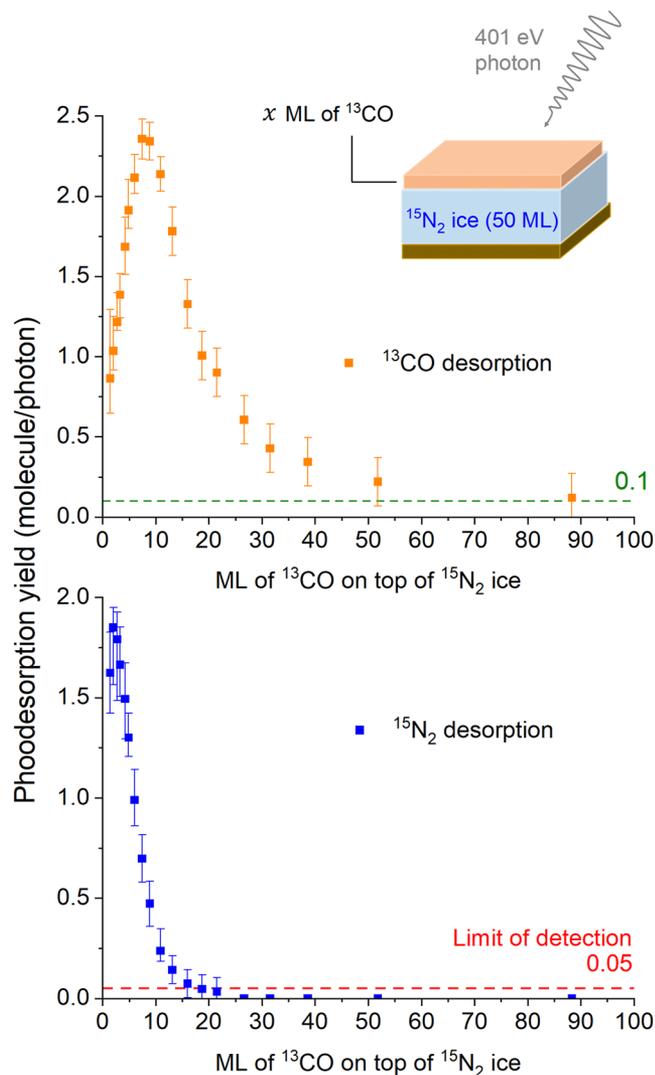


FIG. 5. X-ray photodesorption yields at 401 eV of ^{13}CO (upper panel) and $^{15}\text{N}_2$ (lower panel) from a layered $^{13}\text{CO}/^{15}\text{N}_2$ ice at 15 K, as a function of the number of ML of ^{13}CO deposited on top of a 50 ML $^{15}\text{N}_2$ ice.

and that we do not expect the ionization cross section to significantly differ between 401 and 530 eV. Consequently, the photodesorption yield of 0.1 molecules photon $^{-1}$ is used as a reference to estimate the contribution of ^{13}CO photoabsorption to the desorption of ^{13}CO at 401 eV in our layered ice experiment (it is displayed as a horizontal dashed green line in the upper panel of Fig. 5). From the upper panel of Fig. 5, we observe two different regimes for the photodesorption yield of ^{13}CO :

- A low to medium coverage regime, for a number of ML of ^{13}CO from 1 to roughly 30–40 ML: In this regime, the photodesorption yield of ^{13}CO is >0.1 – 0.2 molecules photon $^{-1}$ such that the desorption of ^{13}CO is dominantly triggered

by photoexcitation of $^{15}\text{N}_2$. In that case, the photodesorption yield of ^{13}CO increases until ~ 10 ML and then starts to decrease.

- A high coverage regime, for a number of ML of ^{13}CO > 40 ML: In this regime, the photodesorption yield of ^{13}CO is close to 0.1 molecules photon $^{-1}$ such that the desorption of ^{13}CO is dominantly triggered by photoionization of ^{13}CO . In that case, the possible photodesorption of ^{13}CO triggered by photoexcitation of $^{15}\text{N}_2$ can be neglected.

The photodesorption yield of $^{15}\text{N}_2$ from our layered ice is displayed in the lower panel of Fig. 5. It is found to decrease as the number of ML of ^{13}CO increases until it goes below our detection limit of ~ 0.05 molecules photon $^{-1}$ (considering the low photon flux for this experiment) for a number of ML of ^{13}CO > 15 ML. The fact that we still observe $^{15}\text{N}_2$ desorption even when the $^{15}\text{N}_2$ ice is covered by a few ML of ^{13}CO may be due to diffusion of $^{15}\text{N}_2$ from below the ML of ^{13}CO or due to a few $^{15}\text{N}_2$ molecules still present at the surface due to inhomogeneous deposition of ^{13}CO . This is further discussed in Sec. IV C.

IV. DISCUSSION

A. Indirect desorption and estimation of Λ_{des}

Our experiments on mixed $^{13}\text{CO}/^{15}\text{N}_2$ ices highlight an indirect desorption mechanism, in the sense that the photoabsorption of one molecule is inducing the desorption of another one. This is clearly seen in the mixed $^{13}\text{CO}/^{15}\text{N}_2$ (2:1) ice for which the photodesorption of ^{13}CO at the $^{15}\text{N}_2$ N 1s $\rightarrow \pi^*$ transition (upper left panel of Fig. 4) and that of $^{15}\text{N}_2$ at the ^{13}CO O 1s $\rightarrow \pi^*$ transition (lower right panel of Fig. 4) dominate the photodesorption spectrum of, respectively, ^{13}CO and $^{15}\text{N}_2$. More globally, the fact that the photodesorption spectrum of $^{15}\text{N}_2$ and ^{13}CO follows the x-ray absorption spectrum of, respectively, ^{13}CO near the O K-edge and $^{15}\text{N}_2$ near the N K-edge confirms this indirect mechanism. Our experiment on the layered $^{13}\text{CO}/^{15}\text{N}_2$ ice also sheds light on this process: The ice is irradiated at 401 eV (N 1s $\rightarrow \pi^*$ transition of $^{15}\text{N}_2$) and ^{13}CO desorption is observed.

This indirect desorption mechanism necessarily involves energy transport from the molecule that absorbs the incident photon to the molecule that desorbs from the ice surface. The length scale of such energy transport can be defined as the maximum distance from the ice surface (or the depth) at which an absorbed photon can induce desorption. This depth will be defined as Λ_{des} in the following and it will be expressed in ML. Λ_{des} depends on the indirect desorption mechanism and may vary with the photon energy and the ice composition.

Our experiment on the layered $^{13}\text{CO}/^{15}\text{N}_2$ ice nicely enables the estimation of Λ_{des} for the indirect desorption process of ^{13}CO induced by the photoabsorption of $^{15}\text{N}_2$ at 401 eV. In Fig. 5, we can see that $^{15}\text{N}_2$ photoabsorption can induce ^{13}CO desorption when the $^{15}\text{N}_2$ ice is covered by a number of ML of ^{13}CO from 1 to roughly 30–40 ML, with a maximum efficiency around 10 ML. Above 40 ML, the photodesorption yield of ^{13}CO from the layered ice is similar to the one induced by the photoionization of the core C 1s electron of ^{13}CO at 530 eV in pure CO ice (Dupuy *et al.*, 2021b) and the indirect desorption mechanism of ^{13}CO induced by photoabsorption of $^{15}\text{N}_2$ is negligible. After fitting, the exponential decay of the

^{13}CO photodesorption yield starting at ~ 10 ML is associated with a characteristic length of $\Lambda_c \sim 11$ ML. The exponential reaches 95% extinction at $3 \times \Lambda_c = 33$ ML, which corresponds to a good estimate of Λ_{des} and for which the indirect desorption of ^{13}CO becomes negligible. In order to take into account the uncertainty associated with the number of ML of ^{13}CO , which is $\sim 10\%$, we finally estimate that Λ_{des} is between 30 and 40 ML in the case of our layered $^{13}\text{CO}/^{15}\text{N}_2$ ice and at 401 eV. This shows that desorption can be triggered from the bulk of the ice, beyond the first few ML.

This conclusion can also be drawn for the case of mixed $^{13}\text{CO}:^{15}\text{N}_2$ ices. In that case, the photodesorption yield of ^{13}CO at the $\text{N } 1s \rightarrow \pi^*$ transition of $^{15}\text{N}_2$ exactly follows the variations of the TEY from a 2:1 to a 20:1 ratio (see upper left panel of Fig. 4). This is not the case for the desorption of $^{15}\text{N}_2$, which is barely visible in the case of a 20:1 ratio (see lower left panel of Fig. 4). This could be explained as follows. In the case of the 2:1 ratio, $^{15}\text{N}_2$ molecules are sufficiently present near the ice surface and within its bulk to observe desorption of both $^{15}\text{N}_2$ and ^{13}CO at the $\text{N } 1s \rightarrow \pi^*$ transition of $^{15}\text{N}_2$. However, in the case of the 20:1 ratio, $^{15}\text{N}_2$ molecules are not sufficiently present near the ice surface to significantly desorb, but they are sufficiently photoabsorbing from the surface to the bulk of the ice to produce a non-negligible TEY value and to induce the desorption of ^{13}CO . In the case of mixed $^{13}\text{CO}:^{15}\text{N}_2$ ices, x-ray photodesorption can, therefore, also be triggered from the bulk of the ice.

B. Indirect desorption mechanisms

In this section, we discuss the possible mechanisms responsible for the indirect desorption of ^{13}CO by excitation of $^{15}\text{N}_2$ observed in our layered ice experiment. Extrapolation to the case of the pure and mixed ices will be discussed afterward. One possible mechanism that has been suggested for the x-ray photodesorption of neutrals from ices of simple (CO , H_2O) and more complex (CH_3OH) molecules (Dupuy *et al.*, 2018, 2020, 2021b; Basalgète *et al.*, 2021a,b) is x-ray induced electron stimulated desorption (XESD). XESD is a consequence of the relaxation of the core hole photoexcited or photoionized state that decays by emission of an Auger electron with a probability close to 1 for low Z elements (Walters and Bhalla, 1971; Krause, 1979). This Auger electron scatters inelastically in the ice, creating secondary events (e.g., ionizations and excitations) and a cascade of low energy (< 20 eV) secondary electrons that may lead to desorption.

On the other hand, mechanisms decorrelated from Auger scattering may also play a role. This includes all the mechanisms related to the energy left to the photoabsorbing molecule after Auger decay. Quantitatively, most of the photon energy is given to the Auger electron in the form of kinetic energy (~ 380 eV for the $\text{N } 1s \rightarrow \pi^*$ transition of $^{15}\text{N}_2$; Moddeman *et al.*, 1971) and a few tens of eV is left to the photoabsorbing molecule. This latter found itself in a singly or multiply valence-excited state, singly or doubly ionized depending on the energy of the incident photon. For instance, in the specific case of our layered ice experiment at 401 eV, $^{15}\text{N}_2$ is expected to be left, after Auger decay, in an excited ionized state ($^{15}\text{N}_2^+)^*$. The energy left to the molecule in the form of electronic excitation may be transported toward the ice surface via different mechanism to cause indirect desorption (i.e., desorption of other molecules). These mechanisms possibly include the following:

- (1) (a) Indirect desorption due to migration of the electronic excitation through the ice (exciton migration): This mechanism concerns the transport of exciton in ices, a process that stems from the work of Kimmel and co-workers on electron-stimulated desorption (ESD) from D_2O ice (Orlando and Kimmel, 1997; Petrik and Kimmel, 2003). It has also been suggested to play a role in the ESD of C_6H_6 on top of water ice by migration of long-lived excitons through the ice (Marchione *et al.*, 2016; Marchione and McCoustra, 2017).
(b) Relaxation of the excited state to modes coupled with desorption similar to what is observed in VUV photodesorption experiments (Bertin *et al.*, 2013) for which it has been shown that excitation of ^{13}CO (to the $\text{A}^1\Pi$ state) and $^{15}\text{N}_2$ (to the $\text{b}^1\Pi_u$ state) can induce the desorption of, respectively, $^{15}\text{N}_2$ and ^{13}CO from layered ices: In that case, an energy transfer between the excited and the desorbing molecule via excitation of collective vibrational modes (phonons) was suggested. To simplify the following discussion, we will consider that this process occurs without exciton migration, but one should note that it could take place after exciton migration [in that case, it is included in process (1) (a)].
- (2) Indirect desorption due to diffusion of energetic fragments if the ionized excited state is dissociative, similar to the kick-out mechanism suggested in the case of indirect desorption of simple molecules on top of water ice, induced by energetic H atoms that are produced from photodissociation of H_2O in the VUV range (Dupuy *et al.*, 2021a): In some extreme cases, the dissociation may occur before Auger decay.

The aim of our discussion is then to identify the main mechanism responsible for the indirect desorption observed in our layered $^{13}\text{CO}/^{15}\text{N}_2$ ice experiment. We will especially discriminate between the XESD process and the other processes cited previously that are decorrelated from the Auger scattering and discuss the processes with respect to the value derived for Λ_{des} (30–40 ML).

For a XESD process, the energy absorbed in the bulk of the ice by $^{15}\text{N}_2$ can be transported toward the ice surface by the scattering of the Auger electron. Considering the energy deposition profile of keV electrons in light materials ($1 \leq Z \leq 18$) from the work of Valkealahti *et al.* (1989), we can reasonably expect that the Auger electron and the secondary electrons can deposit their energy through a few tens of ML. In that sense, the value of Λ_{des} previously estimated (30–40 ML) is consistent with a XESD process. The migration distance of excitons in water ice is potentially tens of ML (Petrik and Kimmel, 2003). Migration of excitons in the case of weakly interacting van der Waals solids such as in $^{15}\text{N}_2$ ice and migration between $^{15}\text{N}_2$ and ^{13}CO layers has not been studied yet, so the process (1)(a) cannot be discussed regarding the value derived for Λ_{des} . However, as discussed in Sec. IV D, we expect this process to depend on the exciton state and, hence, on the incident photon energy.

Process (1) (b) has been characterized by Bertin *et al.* (2013) for the indirect desorption of ^{13}CO and $^{15}\text{N}_2$ by excitation of, respectively, $^{15}\text{N}_2$ and ^{13}CO in layered ice experiments. In that case, it has been shown that the energy can be efficiently transferred over ~ 5 ML to the desorbing molecules at the surface (this value was similar for excitation of ^{13}CO to the $\text{A}^1\Pi$ state and for excitation of $^{15}\text{N}_2$

to the $b^1\Pi_u$ state). Therefore, we expect process (1)(b) to be a near-surface process. Process (2) involves the diffusion of ionic or neutral fragments (in that case N or N^+) toward the ice surface and collision with ^{13}CO molecules followed by desorption. After possible dissociation of $(^{15}\text{N}_2^+)^*$, we expect the fragments to have less than 10 eV of kinetic energy. Although no studies are available on the typical diffusion length of such fragments in the bulk of molecular ices, we do not expect it to be comparable with the value of Λ_{des} derived. Instead, we expect process (2) to be a near-surface process involving just the first few ML of the ice. This is also supported by the very low level of desorbing ^{15}N fragments from pure $^{15}\text{N}_2$ ice (see Sec. III B 1).

Processes (1)(b) and (2) are then expected to possibly occur in the low coverage regime of our layered ice experiment, for a few ML of ^{13}CO on top of the $^{15}\text{N}_2$ ice. Regarding the photodesorption yields of ^{13}CO in Fig. 5, these near-surface processes seem to be negligible compared to the processes involving deeper layers of the $^{15}\text{N}_2$ ice, such as XESD. In fact, the yields of ^{13}CO at low coverage (between 1 and 10 ML) and higher coverage are similar. In the low coverage regime, both near-surface processes and bulk processes can participate in indirect desorption, whereas for higher coverage, only

bulk processes participate. The fact that the yields are similar in both regimes excludes a significant contribution of the near-surface processes compared to the bulk ones. Therefore, processes (1)(b) and (2) can be neglected.

In Sec. IV C, we focus on characterizing the XESD process for our layered ice experiment and we conclude on the dominant mechanism expected for indirect desorption of neutral $^{15}\text{N}_2$ and ^{13}CO from the ices studied in this work.

C. Modeling of the energy deposited in the ice by Auger electrons

XESD is mediated by the scattering of Auger electrons. These electrons deposit their energy through the ice via ionization and excitation of the molecules in their path. For a XESD process, the photodesorption yields are expected to be correlated to the amount of energy deposited near the ice surface by the scattering of Auger electrons. This amount of energy can be estimated based on the work of Valkealahti *et al.* (1989). They showed, by Monte Carlo simulations, that the energy deposition profile of a keV electron in a light material ($1 \leq Z \leq 18$) as a function of depth can be approximated

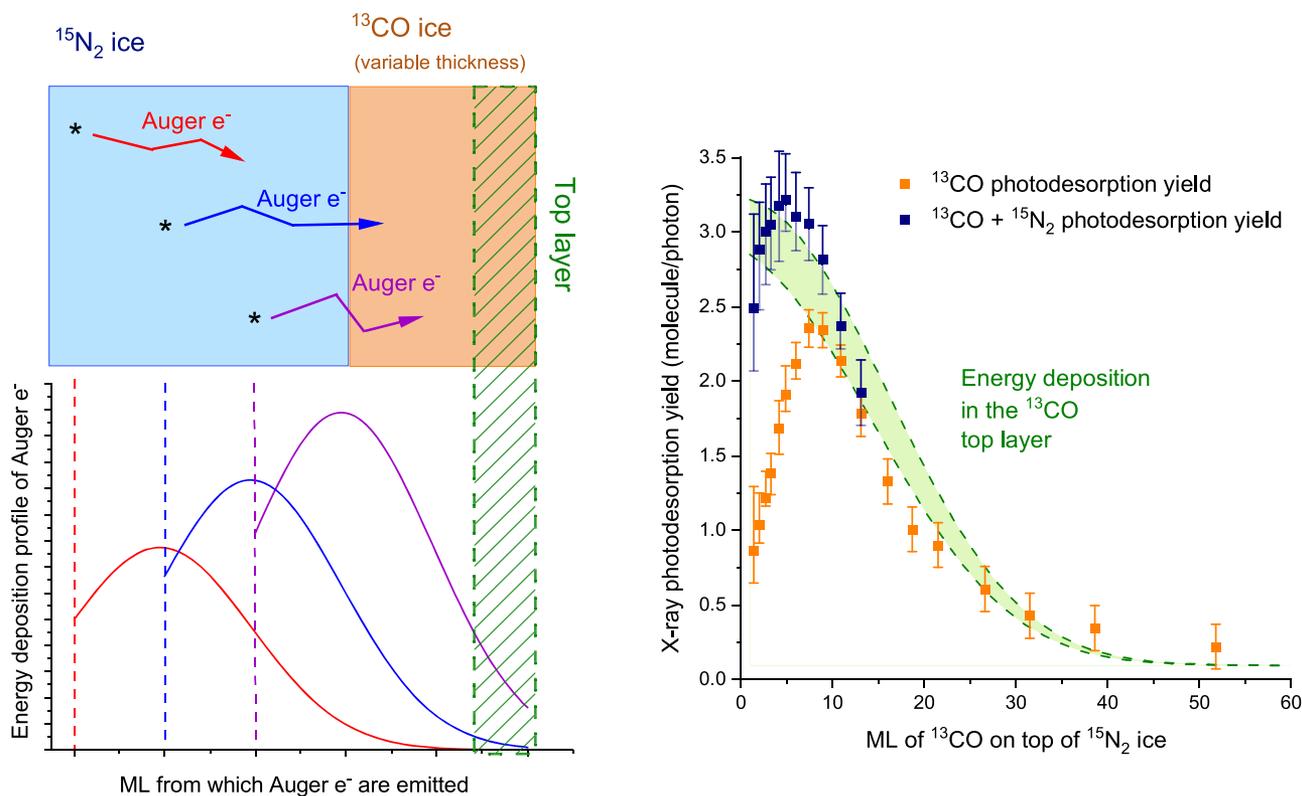


FIG. 6. Left side: scheme of our simple model to estimate the energy deposited to the top ^{13}CO layer of a layered $^{13}\text{CO}/^{15}\text{N}_2$ ice by scattering of Auger electrons emitted after photoabsorption by $^{15}\text{N}_2$ molecules. Right side: in green is displayed the estimation of this energy (in arbitrary units and where the zero of the right side distribution has been scaled with a photodesorption yield of 0.1 molecules photon $^{-1}$), compared with the x-ray photodesorption yields of ^{13}CO (orange squares) from our layered $^{13}\text{CO}/^{15}\text{N}_2$ ice (from Fig. 5). The blue squares show the sum of the photodesorption yields of ^{13}CO plus $^{15}\text{N}_2$ from our layered $^{13}\text{CO}/^{15}\text{N}_2$ ice. The upper and lower green curves were computed by taking a photoabsorption cross section at 401 eV of 40 and 20 Mbarn, respectively.

by a Gaussian distribution (in 1D). The parameters of the distribution depend on the energy of the electron and on the stopping power of the material at this energy. These parameters have been fitted for a N_2 ice. We can, therefore, estimate the amount of energy deposited near the ice surface by the scattering of Auger electrons in the case of our layered $^{13}\text{CO}/^{15}\text{N}_2$ ice experiment using the following methodology and assumptions (this is also sketched in the left side of Fig. 6):

- We consider a layered $^{13}\text{CO}/^{15}\text{N}_2$ ice: 50 ML of $^{15}\text{N}_2$ on top of which an increasing number of ML of ^{13}CO is deposited. We neglect the photoabsorption of the ^{13}CO layers at 401 eV.
- We assume that the initial kinetic energy of the Auger electron emitted by relaxation of the $\text{N}(1s)^{-1}\pi^*$ state of $^{15}\text{N}_2$ is ~ 380 eV, based on Auger spectra of gas phase N_2 (Moddeman *et al.*, 1971).
- We only consider the emitted Auger electrons that statistically contribute the most to the energy deposited near the ice surface, i.e., the Auger electrons that are emitted from the bulk toward the surface of the ice.
- The energy deposition profile of the Auger electrons is approximated by a Gaussian distribution with parameters from Valkealahti *et al.* (1989), in 1D. This also takes into account the energy loss of secondary electrons and it is computed as the sum of the average kinetic energy of all electrons in a depth element. According to the Born–Bethe approximation and Bragg’s rule, the stopping power of ^{13}CO and $^{15}\text{N}_2$ at 380 eV is similar, so we reasonably assumed that the energy deposition profile does not significantly differ in $^{15}\text{N}_2$ or ^{13}CO layers. The shape of the energy deposition profiles is plotted in arbitrary units in the left panel of Fig. 6.
- We consider that each ML of $^{15}\text{N}_2$ is emitting a number of Auger electrons that are equal to the number of photons absorbed in that ML. Thus, the relative intensity of the energy deposition distributions from two different ML is driven by a Beer–Lambert law, with an absorption cross section at 401 eV taken between 20 and 40 Mbarn (Kato *et al.*, 2007).

The energy deposited in the top layer of the ^{13}CO ice is then computed by summing the contributions from each ML of $^{15}\text{N}_2$. This operation is repeated for an increasing number of ML of ^{13}CO deposited on top of the $^{15}\text{N}_2$ ice. The results are presented in the right panel of Fig. 6: The energy deposited on the top layer of the ^{13}CO ice is compared to the experimental photodesorption yields of ^{13}CO from our layered $^{13}\text{CO}/^{15}\text{N}_2$ ice (taken from Fig. 5). A comparison with the sum of the photodesorption yields of ^{13}CO plus $^{15}\text{N}_2$, which have similar photodesorption yields at 401 eV when mixed in similar stoichiometries (according to our results from mixed ices), is also done for the following discussion. At low coverage (number of ML of $^{13}\text{CO} \lesssim 10$ ML), the ice surface is not homogeneously covered by ^{13}CO molecules and the photodesorption yield of ^{13}CO increases linearly with the number of ^{13}CO molecules available for desorption. Therefore, in this regime, the photodesorption yield of ^{13}CO is not correlated with the energy deposition of Auger electrons since it actually depends on the ice surface homogeneity. This lack of homogeneity is supported by the fact that $^{15}\text{N}_2$ still desorbs in

this regime. Furthermore, in this low coverage regime, we observe a good agreement between the energy deposition profile and the sum of the photodesorption yields of ^{13}CO plus $^{15}\text{N}_2$ (see blue squares in Fig. 6). This is consistent with the ice surface being composed of a mix of ^{13}CO and $^{15}\text{N}_2$ molecules in the low coverage regime and for which photodesorption of neutral molecules is correlated with the energy deposited at the top layer and dominated by ^{13}CO plus $^{15}\text{N}_2$ desorption (no significant desorption is found for other possible neutral molecules). For medium coverage (number of ML of ^{13}CO between 10 and 40 ML), the exponential decay of the ^{13}CO desorption yield is in good agreement with the modeled deposited energy. For high coverage (number of ML of $^{13}\text{CO} > 40$ ML), the photodesorption is triggered by the photoionization of ^{13}CO (see Sec. III B 3).

These results strongly suggest that XESD is the dominant mechanism responsible for the indirect desorption of ^{13}CO at 401 eV from our layered $^{13}\text{CO}/^{15}\text{N}_2$ ice. It also confirms our estimated value for Λ_{des} . In that case, excitation of ^{13}CO to non-dissociative states by low energy secondary electrons (produced by the Auger scattering) may be the exact mechanism explaining ^{13}CO desorption. For comparison, ESD of CO from pure CO ice has an energy threshold around 6 eV (Rakhovskaia *et al.*, 1995), which is accessible to the secondary electrons, as suggested in Dupuy *et al.* (2021b). This energy threshold is around 7 eV for the ESD of N_2 from pure N_2 ice.

Auger emission occurs with a probability close to one for each core hole state of $^{15}\text{N}_2$ or ^{13}CO above and below their respective IP. The energy deposition profile due to the Auger scattering is also reasonably expected to be similar in the different ices tested in this work (pure $^{15}\text{N}_2$ ice, mixed $^{13}\text{CO}:^{15}\text{N}_2$ ices, and layered $^{13}\text{CO}/^{15}\text{N}_2$ ice), based on the fact that ^{13}CO and $^{15}\text{N}_2$ have similar stopping power in the energy range considered. It is, therefore, straightforward to consider that XESD should also play a role in the desorption of $^{15}\text{N}_2$ from pure $^{15}\text{N}_2$ ice and in the desorption of ^{13}CO and $^{15}\text{N}_2$ from mixed $^{13}\text{CO}:^{15}\text{N}_2$ ices, at each photon energy. The simple model previously implemented nicely enables the estimation of Λ_{des} for a XESD process at different photon energies as long as the photoabsorption cross section and the initial kinetic energy of the Auger electron emitted after core hole decay are known. This is used in Sec. IV D.

D. Conversion of the absorbed energy to desorption

In this section, we correct our x-ray photodesorption yields by taking into account only the part of the absorbed energy that can participate in the desorption process. This results in expressing the yields in molecules $(\text{eV deposited})^{-1}$ using Eq. (1). We consider that the number of the top monolayers involved in the desorption process (Λ_{des}) should be at least equal to the one corresponding to a XESD process. The energy absorbed in the layers beyond Λ_{des} is “lost” in the sense that it cannot be converted to desorption. This results in taking $\Lambda_{des} = 30\text{--}40$ ML near the N K-edge, where the photoabsorption is dominated by $^{15}\text{N}_2$ absorption. Near the O K-edge, the photoabsorption is dominated by ^{13}CO absorption and the Auger electron is expected to have an initial kinetic energy between 500 and 520 eV (Moddeman *et al.*, 1971). Based on our simple modeling in Sec. IV C, this results in a slightly higher Λ_{des} (40–50 ML). The results are presented in Table I (first four rows) for

TABLE I. Photodesorption yields expressed in molecule desorbed by eV deposited, of ^{13}CO and $^{15}\text{N}_2$ from different ices and at different energies, computed via Eq. (1) and by using the experimental yields derived in this work. The lower and higher values of the yields are derived by taking the higher and lower value of Λ_{des} , respectively.

Photon energy	σ (Mbarn) ^a	Λ_{des} (ML)	Photodesorption yield [$\times 10^{-2}$ molecules (eV deposited) ⁻¹]			
			$^{15}\text{N}_2$ from		^{13}CO from	
			Pure $^{15}\text{N}_2$	$^{13}\text{CO}:^{15}\text{N}_2$ (2:1)	Pure ^{13}CO	$^{13}\text{CO}:^{15}\text{N}_2$ (2:1)
401 eV ($^{15}\text{N}_2$; N 1s $\rightarrow \pi^b$)	26	30–40	1.9–2.3	3.2–4.1		1.9–2.5
425 eV ($^{15}\text{N}_2$; N 1s ionization)	1.8	30–40	0.6–0.8	^b		2.1–2.8
534.2 eV (^{13}CO ; O 1s $\rightarrow \pi^b$)	3.0	40–50		2.8–3.5		2.1–2.6
570 eV (^{13}CO ; O 1s ionization)	0.8	40–50		^b	1.5 ^c	2.4–3.0
8.3 eV [^{13}CO ; $A^1\Pi - X^1\Sigma^+ (2, 0)$] ^d	15	5		18	8.3	21

^aThe photoabsorption cross section is taken from gas phase measurements (Barrus et al., 1979; Kato et al., 2007) and we took into account the spectral resolution of our photon source.

^bBackground issues on the mass channel 30 around 425 eV and 570 eV during the x-ray irradiation of the $^{13}\text{CO}:^{15}\text{N}_2$ -2:1 ice prevent us from using the experimental yields of $^{15}\text{N}_2$ at these energies.

^cTaken from Dupuy et al. (2021b) where it was assumed that $\Lambda_{des} = 30$ ML.

^dFrom Bertin et al. (2013).

the desorption of ^{13}CO and $^{15}\text{N}_2$ from their respective pure ice and from mixed $^{13}\text{CO}:^{15}\text{N}_2$ (2:1) ices at different photon energies.

The yields displayed in Table I represent a quantitative estimate of how efficiently the absorbed energy (only the part that can participate in the desorption process) is converted to the desorption channel. This takes into account all the desorption processes that may occur due to photoabsorption in the first Λ_{des} ML of the ice. The estimated yields are very similar for the desorption of ^{13}CO and $^{15}\text{N}_2$ from pure or mixed ices and at each energy considered. They range between ~ 0.02 and ~ 0.04 molecules (eV deposited)⁻¹. There is, however, a small discrepancy with the photodesorption yield of $^{15}\text{N}_2$ from pure $^{15}\text{N}_2$ ice at 425 eV, which is lower than the previous values. This is assumed to be due to a higher fluence received by the ice before this specific measurement compared to the other ones.

These results indicate that there is presumably a unique dominant process explaining the desorption of ^{13}CO and $^{15}\text{N}_2$ from the ices tested and whose efficiency does not strongly depend on the photon energy; hence, it depends neither on the photoabsorbing molecule nor on its state after Auger decay (singly or doubly ionized excited state). This is fully consistent with a XESD process for which desorption depends only on the Auger scattering, the latter being expected to not significantly differ between pure and mixed ices of ^{13}CO and $^{15}\text{N}_2$. The similarity of the yields between ^{13}CO and $^{15}\text{N}_2$ is also consistent with their similar binding energies in pure or mixed ice (Öberg et al., 2005; Bisschop et al., 2006). A possible significant desorption due to exciton migration from the photoabsorbing molecule (process (1)(a) defined in Sec. IV B) can also be excluded by the fact that the efficiency of this process is expected to depend on the exciton state after Auger decay and, therefore, on the photon energy, which is not what we observe in Table I.

Finally, in Table I, we also computed the VUV photodesorption yields in molecules (eV deposited)⁻¹ at 8.3 eV, which corresponds to the $A^1\Pi - X^1\Sigma^+ (2, 0)$ electronic transition of ^{13}CO , from the data published in Bertin et al. (2013). In this study, Λ_{des} was estimated to be ~ 5 ML for the indirect desorption mechanism from layered

$^{15}\text{N}_2/^{13}\text{CO}$ experiments. The yields in molecules (eV deposited)⁻¹ are found higher in the VUV range (at 8.3 eV) than in the x-ray range. This means that the energy absorbed by a valence excitation of ^{13}CO (to the $A^1\Pi$ state) is more efficiently converted to desorption, by a process similar to the one described in Sec. IV B [process (1)(b)], than the energy absorbed by a core hole excitation or ionization (of ^{13}CO or $^{15}\text{N}_2$) and for which a XESD process is expected to be dominant according to our results. This is true for both the VUV photodesorption of ^{13}CO (direct and/or indirect process) and $^{15}\text{N}_2$ (indirect process) from pure ^{13}CO ice and mixed $^{13}\text{CO}:^{15}\text{N}_2$ ice.

V. CONCLUSION

X-ray photodesorption of $^{15}\text{N}_2$ from pure $^{15}\text{N}_2$ ice and of $^{15}\text{N}_2$ and ^{13}CO from mixed $^{13}\text{CO}:^{15}\text{N}_2$ ices is found to follow the x-ray absorption profile of the ices in the N and O K-edge energy range (near 400 and 500 eV, respectively). Photodesorption experiments on mixed $^{13}\text{CO}:^{15}\text{N}_2$ ices revealed an indirect desorption mechanism triggered by the absorption of $^{15}\text{N}_2$ near the N K-edge and ^{13}CO near the O K-edge, leading to both ^{13}CO and $^{15}\text{N}_2$ desorption. This indirect mechanism is also highlighted by the desorption of ^{13}CO from an x-ray irradiated layered $^{13}\text{CO}/^{15}\text{N}_2$ ice at 401 eV. This latter experiment enables the derivation of the relevant depth of desorption Λ_{des} , which is a quantitative measure of depth involved in the indirect desorption process. It was estimated to be 30–40 ML near the N K-edge. The photodesorption yields of ^{13}CO from the layered $^{13}\text{CO}/^{15}\text{N}_2$ ice as a function of ^{13}CO coverage are found to be well-correlated to the amount of energy deposited at the top layer by the scattering of Auger electrons after x-ray absorption, the latter being estimated by a simple model based on the work of Valkealahti, Schou, and Nieminen (1989). The desorption yields of ^{13}CO and $^{15}\text{N}_2$ corrected from the absorbed energy that can participate in desorption do not significantly depend on the photon energy (hence on the photoabsorbing molecule, ^{13}CO or $^{15}\text{N}_2$) and on the nature of the ice (pure or mixed). These results indicate that

XESD, mediated by the scattering of Auger electrons, is the dominant mechanism explaining the x-ray photodesorption of $^{15}\text{N}_2$ and ^{13}CO from the ices tested in this work.

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AUTHOR DECLARATIONS

Conflict of Interest

The authors have no conflicts to disclose.

Author Contributions

R. Basalgète: Data curation (equal); Formal analysis (equal); Methodology (equal); Writing – original draft (lead); Writing – review & editing (equal). **M. Bertin:** Formal analysis (equal); Writing – review & editing (equal). **D. Torres-Díaz:** Formal analysis (equal); Writing – review & editing (equal). **A. Lafosse:** Formal analysis (equal); Writing – review & editing (equal). **L. Amiaud:** Formal analysis (equal); Writing – review & editing (equal). **G. Féraud:** Formal analysis (equal); Writing – review & editing (equal). **P. Jeseck:** Formal analysis (equal); Writing – review & editing (equal). **L. Philippe:** Formal analysis (equal); Writing – review & editing (equal). **X. Michaut:** Formal analysis (equal); Writing – review & editing (equal). **J.-H. Fillion:** Formal analysis (equal); Writing – review & editing (equal).

DATA AVAILABILITY

The data that support the findings of this study are available within the article.

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