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Simulation of accelerated deflagrations with the CALIF³S-P²REMICS software using RANS and LES approaches

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Background & Motivations

□ Safety issues related to explosion hazards

- in case of severe accident in the reactor building,
- inside or outside the reactor building in case of loss of sealing of pipes containing hydrogen.





Explosion at Fukushima reactor 3

Hydrogen jet deflagration

 \rightarrow The IRSN develops an in-house software for the simulation of explosions: CALIF³S-P²REMICS.



Background & Motivations

- Models currently used in CALIF³S-P²REMICS for deflagrations:
 - Turbulence: RANS approach,
 - Combustion: turbulent flame-speed closure combustion model.
- Drawbacks of this approach:
 - Rough description of the turbulence,
 - Results strongly dependent on the turbulent flame speed correlation.
- Difficulties:

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- Lack of experimental data for the validation of turbulence models,
- Cost of the experiments, variability of the configurations encountered, need for change of scale...

\rightarrow Development of a LES approach for deflagration

- □ Objective: use LES approach in order to
 - Obtain predictive and detailed data on turbulence,
 - Improve the RANS approach.



Volvo test rig: fuel mass fraction

Outline

Combustion models

LES approach RANS approach

An accelerated deflagration

Experimental setup

LES numerical results

Influence of obstacles location Constant vs dynamic wrinkling factor exponent

RANS numerical results

Summary & Perspectives



- A virtually thickened flame formalism (TFLES) is used [Butler & O'Rourke, 1977]
 - The flame thickness is increased by a factor *F* while keeping the laminar flame speed s_l constant.
- Increasing the flame thickness induces a reduced sensitivity of the flame to turbulent motions:
 - Eddies of size smaller than *F* δ⁰_l cannot wrinkle the thickened flame front, with δ⁰_l the laminar flame thickness,
 - A wrinkling factor \(\exists \Delta\) for the sub grid scales is introduced.

Chemical species balances:

$$\partial_t \overline{\rho y_k} + \operatorname{div}(\overline{\rho y_k} \widetilde{\boldsymbol{u}}) + \operatorname{div}\left(\boldsymbol{\mathcal{F}} \overline{\boldsymbol{j}}_k
ight) = rac{\dot{\omega}_k}{\boldsymbol{\mathcal{F}}}$$

Sensible energy balance:

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$$\partial_t(\overline{
ho}\,\widetilde{e}) + \operatorname{div}(\overline{
ho}\,\widetilde{u}\,\widetilde{e}) + \overline{p}\operatorname{div}(\widetilde{u}) + \operatorname{div}(\mathcal{F}\,\overline{q}) = rac{\dot{\omega}_T}{\mathcal{F}} + \overline{ au}(\widetilde{u}): \mathbf{
abla}\widetilde{u} + \epsilon_v$$



DNS of flame turbulence interactions



Flame artificially thickened by a factor ${\cal F}$

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- ➡ A wrinkling factor Ξ_Δ for the sub grid scales is introduced.

Chemical species balances:

$$\partial_t \widetilde{\rho y_k} + \operatorname{div}(\widetilde{\rho y_k} \widetilde{\boldsymbol{u}}) + \operatorname{div}\left(\Xi_\Delta \mathcal{F} \widetilde{\boldsymbol{j}}_k\right) = \Xi_\Delta \frac{\omega_k}{\mathcal{F}}$$

Sensible energy balance:

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$$\partial_t(\overline{\rho}\,\widetilde{e}) + \operatorname{div}(\overline{\rho}\,\widetilde{u}\,\widetilde{e}) + \overline{p}\operatorname{div}(\widetilde{u}) + \operatorname{div}(\Xi_\Delta\,\mathcal{F}\,\overline{q}) = \Xi_\Delta\,\frac{\dot{\omega}_T}{\mathcal{F}} + \overline{\tau}(\widetilde{u}): \boldsymbol{\nabla}\widetilde{u} + \epsilon_v$$



DNS of flame turbulence interactions



Flame artificially thickened by a factor ${\cal F}$

Interpretation in the subgrid scale wrinkling factor Ξ_Δ:

$$\Xi_\Delta = \left(rac{\Delta}{\delta_c}
ight)^{\ell}$$

 Δ being the combustion filter size, δ_c the inner cutoff scale (i.e. the lowest wrinkling scale).

- \square The model parameter eta can be:
 - fixed,

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- dynamically determined equating flame surfaces computed at filtered and test-filtered scales (Germano-like identity),
- corrected in order to take into account flame front interactions and flame/wall interactions [Mouriaux, 2011]
- \rightarrow avoid unphysical values of the subgrid scale wrinkling factor.

E_s(K) resolved part modeled part





The [Veynante & Moureau 2015] formulation is used for the subgrid scale wrinkling factor Ξ_{Δ} :

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ight)^{eta}$$

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IRSI

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E.(K) resolved part



Germano-like procedure: resolved scales are used to model subgrid scales



RANS approach

Geometrical description of the flame front (G-equation, [Peters, 2000]):

$$\partial_t(\overline{\rho}\,\widetilde{G}) + \operatorname{div}(\overline{\rho}\,\widetilde{G}\,\widetilde{u}) + \rho_u u_f |\mathbf{\nabla}\widetilde{G}| = 0$$

□ but, for partially premixed flow, G cannot be identified to a progress variable...
→ keep the usual mass balance equations for the chemical species, and compute the reaction rate as a function of G:

 \blacksquare The turbulent flame speed u_{f} is given by correlations (Zimont, Peters, \dots)

 $u_{\mathbf{f}} = f(s_l, u', l_t, \dots),$

and the laminar flame speed is defined from unburnt mixture local conditions

$$s_l(\phi, p, \theta) = s_l^{ref}(\phi) \left(\frac{T}{T_{ref}}\right)^{\alpha} \left(\frac{p}{p_{ref}}\right)^{\beta}.$$



RANS approach

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ight)^{lpha} \, \left(rac{p}{p_{ref}}
ight)^{eta}.$$



An accelerated deflagration - Experimental setup

Context:

 Accelerating flame in an obstructed square channel.

Geometry [Wen et al., 2013]:

- Obstructed semi-confined chamber (150 mm × 150 mm × 500 mm),
- Top of the chamber sealed with a thin polyvinyl chloride membrane,
- Three configurations of obstacles (75 mm × 150 mm × 10 mm).

Configuration:

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- 😁 Stoichiometric methane-air mixture,
- Ignition at the bottom of the chamber.



LES numerical results

- 2D numerical results are performed in order to setup 3D simulations.
- 3D simulations have been performed on the supercomputer Topaze at the Research and Technology Computing Center (CCRT).
- 3D numerical results confirm the results obtained in 2D.
- Only 2D numerical results are presented here.



Configuration 3: flame front structure



Experimental images



Influence of obstacles location - Flame front position



- The transition from laminar to turbulent well recovered (before the first obstacle)
- Good reproduction of the flame acceleration, but flame faster than experiments between the first and the second obstacles
- Staggered obstacles configuration: faster flame.

Influence of obstacles location - Overpressure evolution

- First experimental peak due to sealing film disintegration: not simulated, thus not recovered.
- Configurations 1 and 2:
 - Overpressure peaks reached when the flame front is located between the third obstacle and the exit,
 - Overpressure peaks well recovered.
- Configuration 3:
 - Overpressure peak reached when the flame front is located outside the chamber,
 - Overpressure peaks overestimated.

Staggered obstacles configuration: most penalizing.



Configuration 3 - Constant vs dynamic wrinkling factor exponent



Flame front position

Overpressure evolution

igsquirin Numerical results obtained with a fixed or dynamically computed wrinkling factor are compared.

The dynamical formulation of the wrinkling factor is essential to catch the transition from laminar to turbulent regimes and recover the good flame behaviour.

RANS numerical results - Modelling

 \Box Turbulence: k- ω SST model,

Combustion: turbulent flame-speed closure combustion model

- [Bray, 1990] correlation:
$$u_f = 0.875 \,\mathrm{Ka}^{-0.392} \,u'$$
, with $\mathrm{Ka} = 0.157 \left(rac{u'}{s_l}
ight)^2 \mathrm{Re}_T^{-0.5}$,

- [Zimont, 2000] correlation:
$$u_f = A u' \operatorname{Da}^{1/4}$$
,

· [Goulier, 2015] correlation:
$$u_f = s_l \, 1.613 \, \left(rac{r}{l_t}
ight)^{0.333} \, \left(rac{u'}{s_l}
ight)^{0.526} \, {
m Le}^{-0.14}$$

 \rightarrow Numerical simulations are performed on a 2D computational domain with the three turbulent flame speed correlations.

 \rightarrow Only configuration 3 (staggered obstacles) is simulated.

Flame front structure (RANS vs LES) - Goulier correlation



Snapshots of the temperature field and velocity vectors - RANS approach



Snapshots of the temperature field and velocity vectors - LES approach



Flame front structure (RANS vs Experiment) - Goulier correlation



Snapshots of the temperature field and velocity vectors



Experimental images



Turbulent flame speed correlation impact

Flame front position:

 Goulier correlation: good approximation of the flame front position,

 \rightarrow the transition from laminar to turbulent flame is well reproduce probably due to the term $(r/l_t)^{0.333}$ [Gostintsev *et al.*, 1988].

 Zimont and Bray correlations: underestimation of the flame speed

 \rightarrow these correlations are designed for situations where the turbulence is fully developped.

Overpressure evolution:

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- Underestimation of the pressure peak with the three correlations,
- Pressure fluctuations due to pressure waves reflections on walls and on obstacles.

Flame front position



Overpressure evolution



Summary & Perspectives

LES approach for deflagration developed and implemented in CALIF³S-P²REMICS software.

Simulation of an accelerated deflagration with LES and RANS approaches:

- Good agreement between LES numerical results and experimental data;
- Dynamical formulation of the wrinkling factor necessary for a good reproduction of the flame front propagation;
- 🐡 RANS: dispersion of numerical results depending on the turbulent flame speed correlation used.

Next step

- Use LES approach (turbulent characteristics of the flow) in order to perform in indepth the interpretation of experiments and to support RANS model validation.
- 🗯 LES approach more predictive than RANS approach but too expensive for industrial purpose:
 - 2D simulations: 13 hours on 180 processors (6.5×10^5 cells, 2.6×10^5 time steps),
 - 3D simulations: 20 days on 2500 processors (160 \times 10^{6} cells, 2.6×10^{5} time steps),
 - \rightarrow Development of an "hybrid" RANS/LES approach?