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# Spatial distribution of far-infrared rotationally excited CH<sup>+</sup> and OH emission lines in the Orion Bar photodissociation region<sup>★</sup>

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## ABSTRACT

**Context.** The methylidyne cation (CH<sup>+</sup>) and hydroxyl (OH) are key molecules in the warm interstellar chemistry, but their formation and excitation mechanisms are not well understood. Their abundance and excitation are predicted to be enhanced by the presence of vibrationally excited H<sub>2</sub> or hot gas (~500–1000 K) in photodissociation regions (PDRs) with high incident far-ultraviolet (FUV) radiation field. The excitation may also originate in dense gas (>10<sup>5</sup> cm<sup>-3</sup>) followed by nonreactive collisions with H<sub>2</sub>, H, and electrons. Previous observations of the Orion Bar suggest that the rotationally excited CH<sup>+</sup> and OH correlate with the excited CO, which is a tracer of dense and warm gas, and that formation pumping contributes to CH<sup>+</sup> excitation.

**Aims.** Our goal is to examine the spatial distribution of the rotationally excited CH<sup>+</sup> and OH emission lines in the Orion Bar to establish their physical origin and main formation and excitation mechanisms.

**Methods.** We present spatially sampled maps of the CH<sup>+</sup>  $J = 3-2$  transition at 119.8  $\mu\text{m}$  and the OH  $\Lambda$  doublet at 84  $\mu\text{m}$  in the Orion Bar over an area of 110''  $\times$  110'' with *Herschel*/PACS. We compare the spatial distribution of these molecules with those of their chemical precursors, C<sup>+</sup>, O and H<sub>2</sub>, and tracers of warm and dense gas (high- $J$  CO). We assess the spatial variation of the CH<sup>+</sup>  $J = 2-1$  velocity-resolved line profile at 1669 GHz with *Herschel*/HIFI spectrometer observations.

**Results.** The OH and especially CH<sup>+</sup> lines correlate well with the high- $J$  CO emission and delineate the warm and dense molecular region at the edge of the Bar. While notably similar, the differences in the CH<sup>+</sup> and OH morphologies indicate that CH<sup>+</sup> formation and excitation are strongly related to the observed vibrationally excited H<sub>2</sub>. This, together with the observed broad CH<sup>+</sup> line widths, indicates that formation pumping contributes to the excitation of this reactive molecular ion. Interestingly, the peak of the rotationally excited OH 84  $\mu\text{m}$  emission coincides with a bright young object, proplyd 244–440, which shows that OH can be an excellent tracer of UV-irradiated dense gas.

**Conclusions.** The spatial distribution of CH<sup>+</sup> and OH revealed in our maps is consistent with previous modeling studies. Both formation pumping and nonreactive collisions in a UV-irradiated dense gas are important CH<sup>+</sup>  $J = 3-2$  excitation processes. The excitation of the OH  $\Lambda$  doublet at 84  $\mu\text{m}$  is mainly sensitive to the temperature and density.

**Key words.** ISM: individual objects: Orion Bar – ISM: lines and bands – photon-dominated region (PDR)

## 1. Introduction

The methylidyne cation (CH<sup>+</sup>) and hydroxyl (OH) have been observed in different environments from the local (e.g., Storey et al. 1981; Nagy et al. 2013; Dawson et al. 2014) to the extragalactic interstellar medium (ISM; e.g., Schmelz et al. 1986; Baan et al. 1992; Darling & Giovanelli 2002; Spinoglio et al. 2012; Rangwala et al. 2014). They are key molecules in the warm interstellar chemistry, and because they require ultraviolet (UV) radiation and high temperatures to form, they trace specific physical processes in the ISM (e.g., Gerin et al. 2016). Thus, it is expected that in addition to the physical conditions of the regions they are observed in, the spatial distribution of the rotationally

excited far-infrared (FIR) emission of these species might give clues to their formation and excitation processes. *Herschel* enables, for the first time, to study the spatial distribution of these lines. In this paper, we present the first spatially resolved maps of these lines in the Orion Bar.

Atoms and hydrogen molecules and radicals are the first chemical building blocks of the ISM. In warm gas CH<sup>+</sup> is believed to form mainly as a product of the C<sup>+</sup> + H<sub>2</sub> reaction. This formation route has a very high endothermicity of 0.374 eV (4300 K), and it has been suggested that this barrier could be overcome by reactive collisions with the vibrationally excited H<sub>2</sub> in strongly irradiated photodissociation regions (PDRs; e.g., White 1984; Lambert & Danks 1986; Jones et al. 1986; Agúndez et al. 2010; Naylor et al. 2010; Godard & Cernicharo 2013; Nagy et al. 2013; Zanchet et al. 2013). In diffuse interstellar clouds with low far-ultraviolet (FUV) radiation field and

<sup>★</sup> *Herschel* is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.

very low density, shocks and turbulence (e.g., Elitzur & Watson 1978; Pineau des Forets et al. 1986; Godard et al. 2009, 2012; Falgarone et al. 2010a,b) have been proposed to overcome the high endothermicity. In addition, Lim et al. (1999) found that the observed abundances of CH<sup>+</sup> by Cernicharo et al. (1997) are explained through thermal reaction between C<sup>+</sup> and H<sub>2</sub> ( $v = 0$ ) if the gas is hot enough ( $>1000$  K).

The reaction C<sup>+</sup> + H<sub>2</sub>( $v$ ) becomes exothermic when H<sub>2</sub> is in vibrationally excited states (Godard & Cernicharo 2013; Zanchet et al. 2013). On the other hand, CH<sup>+</sup> is highly reactive and easily destroyed by reactive collisions with H<sub>2</sub>, H, and electrons. Godard & Cernicharo (2013) and Zanchet et al. (2013) find that with the chemical destruction rates considered typical for PDR environments, all levels of CH<sup>+</sup> ( $J \geq 2$ ) are very sensitive to the formation pumping since the timescale for the chemical reaction becomes comparable or shorter than that of the nonreactive collision. Collisional excitation with H<sub>2</sub>, H, and electrons could, however, be important for the lowest rotational transitions. In contrast, radiative pumping is predicted to have only a marginal effect even in a strong FIR radiation field.

The OH radical is a key intermediary molecule in forming other important PDR tracers such as H<sub>2</sub>O, CO<sup>+</sup>, O<sub>2</sub>, NO, and SO<sup>+</sup> in the ISM. On the surface of high FUV-flux PDRs the OH formation is expected to be dominated by the endothermic reaction with H<sub>2</sub> and atomic oxygen O<sup>0</sup> (Goicoechea et al. 2011; Hollenbach et al. 2012). The route is endothermic by 0.08 eV ( $\sim 900$  K) and with an activation barrier of 0.4 eV ( $\sim 4800$  K). Similar to CH<sup>+</sup>, the possibility of the FUV-pumped vibrationally excited H<sub>2</sub> enhancing the abundance and excitation of OH has been suggested.

Goicoechea et al. (2011) made the first detection of rotationally excited CH<sup>+</sup> and OH emission lines with *Herschel*/PACS toward the CO<sup>+</sup> peak (Störzner et al. 1995) at the edge of the Orion Bar PDR, one of the nearest, nearly edge-on luminous PDRs (with a FUV radiation field of a few 10<sup>4</sup> in Draine units). While their pointed observations cover a small area (47''  $\times$  47''), they hint at a possible spatial correlation between these lines and suggest that the rotational OH emission originates in small irradiated dense structures ( $n_{\text{H}} \sim 10^{6-7}$  cm<sup>-3</sup> and  $T_{\text{k}} \sim 160-220$  K). The Orion Bar is thought to contain different physical structures with an interclump medium at medium density ( $n_{\text{H}} \sim 10^4-10^5$  cm<sup>-3</sup>) and high density clumps ( $n_{\text{H}} \sim 10^{6-7}$  cm<sup>-3</sup>) (e.g., Tielens et al. 1993; Tauber et al. 1994; Lis & Schilke 2003; Lee et al. 2013). A recent analysis of the high- $J$  CO, H<sub>2</sub>, OH, and CH<sup>+</sup> emission lines using the PDR Meudon code by Joblin et al. (in prep.) shows that these lines are very sensitive to the thermal pressure and suggests  $P \sim 2 \times 10^8$  K cm<sup>-3</sup> for the emitting structures. However, this study uses pointed observations and in order to understand the CH<sup>+</sup> and OH formation and excitation in detail, we need to understand their spatial distribution and how they compare with other tracers.

Using high spectral resolution *Herschel*/HIFI observations, Nagy et al. (2013) find that the CH<sup>+</sup>  $J = 1-0$  and  $J = 2-1$  lines have line widths of  $\Delta V \sim 5$  km s<sup>-1</sup> toward the Orion Bar that are larger than those of other molecular lines, which typically show line widths of 2–3 km s<sup>-1</sup>. This line broadening could result from formation pumping leading to a nonthermalized velocity distribution.

In this paper, we study the spatial distribution of the rotationally excited CH<sup>+</sup> and OH emission lines using fully sampled PACS maps of CH<sup>+</sup> transition  $J = 3-2$  at 119.8  $\mu\text{m}$  and OH 84  $\mu\text{m}$   $\Lambda$  doublet toward the Orion Bar. These are the first fully sampled maps of these emission lines in a PDR. Observational constraints on the excited CH<sup>+</sup> and OH spatial distribution and

comparison with tracers of warm and dense gas and vibrational excited H<sub>2</sub> are needed to establish their physical origin and their main excitation mechanisms (collisions versus pumping). Given the similar critical densities ( $\sim 10^{10}$  cm<sup>-3</sup>) and upper level energies ( $E/k \sim 250$  K) of the targeted lines, our observations allow us to study how the chemical reaction with H<sub>2</sub> ( $v > 0$ ) affects the formation and excitation of CH<sup>+</sup> and OH.

We first describe the *Herschel* observations and data reduction of the CH<sup>+</sup> and OH lines, as well as the lines observed for comparison, in Sect. 2. We discuss the line detection, spatial morphology and origin of the CH<sup>+</sup>  $J = 3-2$  and OH 84  $\mu\text{m}$  lines in Sect. 3. In Sect. 4 we investigate the CH<sup>+</sup> and OH formation and chemical pumping excitation mechanism via reaction with vibrationally excited H<sub>2</sub>. In this section we also compare OH with H<sub>2</sub>O, since OH can also be the product of H<sub>2</sub>O photodissociation in the gas unshielded against FUV radiation. In Sect. 5, we discuss the CH<sup>+</sup> velocity dispersion. The detection of a propylid in the OH map is described in Sect. 6. We conclude and summarize the findings of the paper in Sect. 7.

## 2. Observations and data reduction

### 2.1. CH<sup>+</sup> $J = 3-2$ and OH lines observed with PACS

We observed the Orion Bar over an area of 110''  $\times$  110'' with the Photoconductor Array Camera and Spectrometer (PACS, Poglitsch et al. 2010) on board *Herschel* Space Observatory. In particular, we mapped the following lines: CH<sup>+</sup>  $J = 3-2$  (120  $\mu\text{m}$ ), OH 84  $\mu\text{m}$   $\Lambda$  doublet (84.4 and 84.6  $\mu\text{m}$ , transitions  ${}^2\Pi_{3/2} J = 7/2^+ - 5/2^-$  and  ${}^2\Pi_{3/2} J = 7/2^- - 5/2^+$ , respectively), and OH 119  $\mu\text{m}$   $\Lambda$  doublet (119.2 and 119.4  $\mu\text{m}$ , transitions  ${}^2\Pi_{3/2} J = 5/2^- - 3/2^+$  and  ${}^2\Pi_{3/2} J = 5/2^+ - 3/2^-$ , respectively). Table 1 lists the observed transitions, their wavelengths, critical densities, upper level energy temperatures, and Einstein coefficients.

The lines were observed using the unchopped mode and due to the instantaneous wavelength window covered, the CH<sup>+</sup>  $J = 3-2$  observations at 119.8  $\mu\text{m}$  also contain the OH 119  $\mu\text{m}$  doublet. The observations were carried out on September 14 and 15, 2012 and the total observation time was 21 009 s and 13 795 s for OH and CH<sup>+</sup>, respectively. The observations consist of a fully Nyquist-sampled raster map of 4  $\times$  4 footprints of the CH<sup>+</sup>  $J = 3-2$  line, and a fully Nyquist-sampled 5  $\times$  5 raster map of the OH 84  $\mu\text{m}$  line. These footprints are composed of 5  $\times$  5 spatial pixels (spaxels). The line is observed in 16 different spectral scans for each spaxel, each with an up and down scan. The configuration is shown in Fig. 1, where the raster map is overlaid on top of the 8  $\mu\text{m}$  IRAC image of the Orion Bar. Because of the overlapping footprints, the middle of the map is better sampled and has (slightly) higher S/N than the edges of the map where we only have the data from one footprint. The typical line detection level outside the Bar is  $\sim 10-15\sigma$  and in the Bar  $\sim 15-30\sigma$  for CH<sup>+</sup>  $J = 3-2$  and OH 84  $\mu\text{m}$ . The typical detection level outside the Bar is  $\sim 50\sigma$  and in the Bar  $>100\sigma$  for OH 119  $\mu\text{m}$ .

The data were processed using the version 10.0.2843 of the reduction and analysis package HIPE. Since the PACS instrument has a spectral resolution of  $\sim 135$  km s<sup>-1</sup> at 85  $\mu\text{m}$  and  $\sim 300$  km s<sup>-1</sup> at 120  $\mu\text{m}$ , the line widths and the hyperfine structures of the OH  $\Lambda$  doublets are not resolved in our observations. The spatial resolution is 9'' for the CH<sup>+</sup> 120  $\mu\text{m}$  line and the OH 119  $\mu\text{m}$   $\Lambda$  doublets, and 6'' for the OH 84  $\mu\text{m}$   $\Lambda$  doublets.

The line fitting was performed using IDL-based software PACSman version 3.55 (Lebouteiller et al. 2012). Figure 2 gives an example of the observations and line fitting routine. Using a polynomial baseline, the lines were fitted with a

**Table 1.** Properties of the observed lines.

	$\lambda$ [ $\mu\text{m}$ ]	Transition	$n_{\text{crit}}$ [ $\text{cm}^{-3}$ ]	$E_u/k$ [K]	$A_{ij}$ [ $\text{s}^{-1}$ ]
OH	84.6 <sup>a</sup>	${}^2\Pi_{3/2} J = 7/2^- - 5/2^+$	1 (10) <sup>i</sup>	291	5.20 (−01)
OH	119.4 <sup>a</sup>	${}^2\Pi_{3/2} J = 5/2^+ - 3/2^-$	1 (09) <sup>i</sup>	120	1.38 (−01)
CH <sup>+</sup>	119.8 <sup>b</sup>	$J = 3-2$	4 (09) <sup>j</sup>	240	2.20 (−01)
CH <sup>+</sup>	179.6 <sup>b</sup>	$J = 2-1$	6 (08) <sup>j</sup>	120	6.10 (−02)
CH <sup>+</sup>	359.0 <sup>b</sup>	$J = 1-0$	7 (07) <sup>j</sup>	40	6.36 (−03)
CO	137.3 <sup>c</sup>	$J = 19-18$	5 (06) <sup>k</sup>	1050	6.65 (−04)
<sup>13</sup> CO	227.1 <sup>d</sup>	$J = 12-11$	2 (06) <sup>k</sup>	412	1.52 (−04)
C <sup>18</sup> O	341.5 <sup>e</sup>	$J = 8-7$	6 (05) <sup>k</sup>	190	4.47 (−05)
[CII]	157.7 <sup>f</sup>	${}^2P_{3/2} - {}^2P_{1/2}$	3 (03) <sup>l</sup>	91	2.29 (−06)
[OI]	145.5 <sup>g</sup>	${}^3P_0 - {}^3P_1$	2 (05) <sup>l</sup>	327	1.66 (−05)
H <sub>2</sub> O	269.3 <sup>h</sup>	$1_{11} - 0_{00}$	1 (09) <sup>m</sup>	53	1.84 (−02)
H <sub>2</sub> O	398.7 <sup>h</sup>	$2_{11} - 2_{00}$	3 (08) <sup>m</sup>	137	7.05 (−03)

**Notes.** Columns: wavelength ( $\lambda$ ), transition, critical density ( $n_{\text{crit}}$ ), upper level energy temperature ( $E_u/k$ ), and Einstein coefficient ( $A_{ij}$ ). Numbers in parenthesis are power of 10.

**References.** References for the rest frequencies are: *a*) Pickett et al. (1998); *b*) Müller (2010); *c*) Winnewisser et al. (1997); *d*) Cazzoli et al. (2004); *e*) Klapper et al. (2001); *f*) Cooksy et al. (1986); *g*) Zink et al. (1991); *h*) Yu et al. (2012). The critical densities were calculated using the collisional rates with *i*) ortho-H<sub>2</sub> (Offer et al. 1994); *j*) He (Hammami et al. 2008; Hammami et al. 2009); *k*) H<sub>2</sub> (Yang et al. 2010); *l*) H<sup>0</sup> (Launay & Roueff 1977); and *m*) para-H<sub>2</sub> (Dubernet & Grosjean 2002; Phillips et al. 1996).

**Table 2.** Results from the fitting of the *Herschel*/HIFI observations of the CH<sup>+</sup>  $J = 2-1$  line.

Offset	Area [K km s <sup>−1</sup> ]	$V_{\text{LSR}}$ [km s <sup>−1</sup> ]	$\Delta V$ [km s <sup>−1</sup> ]	$T_{\text{peak}}$ [K]
[−26'', 40'']	2.5 (0.2)	8.7 (0.1)	3.4 (0.3)	0.69
[−13'', 20'']	3.2 (0.2)	9.6 (0.1)	3.4 (0.2)	0.88
[0'', 0'']	6.8 (0.2)	10.6 (0.1)	4.0 (0.2)	1.57
[13'', −20'']	2.0 (0.2)	10.9 (0.1)	3.4 (0.3)	0.55
[26'', −40'']	0.9 (0.2)	11.1 (0.3)	3.1 (0.7)	0.28

**Notes.** The first column lists the offset from the central position of  $\alpha_{\text{J2000}}: 05^{\text{h}}35^{\text{m}}20.61^{\text{s}}$ ,  $\delta_{\text{J2000}}: -05^{\circ}25'14''$ , and subsequent columns give the area [K km s<sup>−1</sup>], velocity  $\Delta V$  [km s<sup>−1</sup>], and the peak temperature [K]. The uncertainties are given in parenthesis. For  $T_{\text{peak}}$  the uncertainty is approximately 10% (calibration). The rest frequency for CH<sup>+</sup>  $J = 2-1$  line is 1669.2813 GHz.

Gaussian profile, and the corresponding line flux was integrated. PACSman measures the line fluxes for each spatial pixel independently. To produce the final map, it creates an oversampled pixelated grid of the observations with a 3'' pixel resolution and calculates the average fractional contribution of the given spatial pixels to the relevant position. PACSman also calculates the statistical uncertainties, including the dispersion in the reduction process and the rms of the fit. In our observations, these uncertainties are small and usually amount to at most 5–20% for the CH<sup>+</sup> and OH lines. The relative accuracy between spatial pixels given in the manual is 10%<sup>1</sup>. We assume a conservative total uncertainty of 22% for the integrated line intensities in the fainter regions, that is the upper limit of a combination of the calibration (10%) and line fitting uncertainties (5–20%). For the Bar, the uncertainty is less than 11%.

We also observed the CO  $J = 19-18$  line at 137.3  $\mu\text{m}$  with PACS (Parikka et al., in prep.). The observations consist of a fully Nyquist-sampled raster map of 4×4 footprints. For CO  $J = 19-18$ , the typical line detection level outside the Bar is  $\sim 10\sigma$  and in the Bar is  $\sim 100\sigma$ . The uncertainty for the integrated line intensities is less than 22% with less than 11% in the Bar. The spatial resolution is 10''.

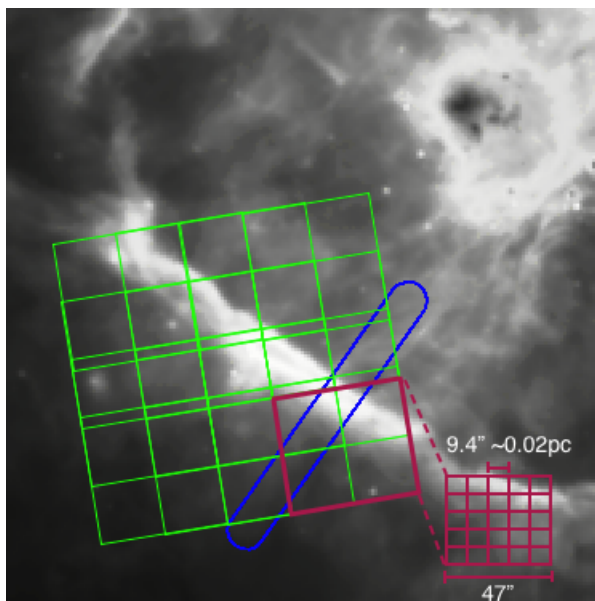
## 2.2. HIFI and SPIRE observations

In order to obtain information about the velocity structure of the CH<sup>+</sup> lines, we used the Heterodyne Instrument for the Far Infrared (HIFI; de Graauw et al. 2010) to observe the CH<sup>+</sup>  $J = 2-1$  line at 180  $\mu\text{m}$  along a cut perpendicular to the Orion Bar (see Fig. 1). The cut is centered at the position  $\alpha_{\text{J2000}} = 05^{\text{h}}35^{\text{m}}20.61^{\text{s}}$ ,  $\delta_{\text{J2000}} = -05^{\circ}25'14''$  and extends across 2' with an inclination of PA = 52°. The data were obtained with the high resolution spectrometer (HRS). We used the on-the-fly observing mode using half-beam sampling and the nominal HIFI beam size at this frequency (12.3''). After porting the FITS data into GILDAS format<sup>2</sup>, the data processing consisted of a scan-by-scan linear baseline subtraction and resampling to a common frequency grid with a resolution of 0.3 km s<sup>−1</sup>. The intensity scale is antenna temperature ( $T_{\text{A}}^*$ ), and the typical rms is 0.12 K calculated at a velocity resolution of 0.3 km s<sup>−1</sup>. The offset (from the central position of  $\alpha_{\text{J2000}} = 05^{\text{h}}35^{\text{m}}20.61^{\text{s}}$ ,  $\delta_{\text{J2000}} = -05^{\circ}25'14''$ ) of the spectra, the calculated area, the velocity ( $V_{\text{LSR}}$ ), FWHM line width ( $\Delta V$ ) and the peak temperature [K] are listed in Table 2.

We also observed CH<sup>+</sup>  $J = 1-0$  at 359  $\mu\text{m}$  and water lines  $1_{11}-0_{00}$  at 269  $\mu\text{m}$  and  $2_{11}-2_{02}$  at 399  $\mu\text{m}$  with the Spectral and Photometric Imaging Receiver (SPIRE; Griffin et al. 2010). The

<sup>1</sup> From the PACS spectroscopy performance and calibration manual. This can be found at <http://herschel.esac.esa.int/twiki/bin/view/Public/PacsCalibrationWeb>

<sup>2</sup> Institut de Radioastronomie Millimétrique (IRAM): <http://iram.fr/IRAMFR/GILDAS/>



**Fig. 1.** Overlay of the PACS observations and HIFI observations on a *Spitzer* IRAC 8  $\mu\text{m}$  image of the Orion Bar. A  $4 \times 4$  raster map of  $\text{CH}^+ J=3-2$  with 16 overlapping footprints, is shown in green and an example of a footprint is shown in red. The strip of  $\text{CH}^+ J=2-1$  is shown in blue.

observations were taken on September 20, 2010 in the high resolution SPIRE FTS full-sampling mode. We include the water lines in our observations as OH can also form via the photodissociation of  $\text{H}_2\text{O}$ . Within the SPIRE range of these observations, a wealth of rotational lines of CO and its isotopologues were detected (Habart et al. 2010, Parikka et al., in prep.). The  $^{13}\text{CO } J=12-11$  line at 227  $\mu\text{m}$  and  $\text{C}^{18}\text{O } J=8-7$  line at 342  $\mu\text{m}$  maps are presented.

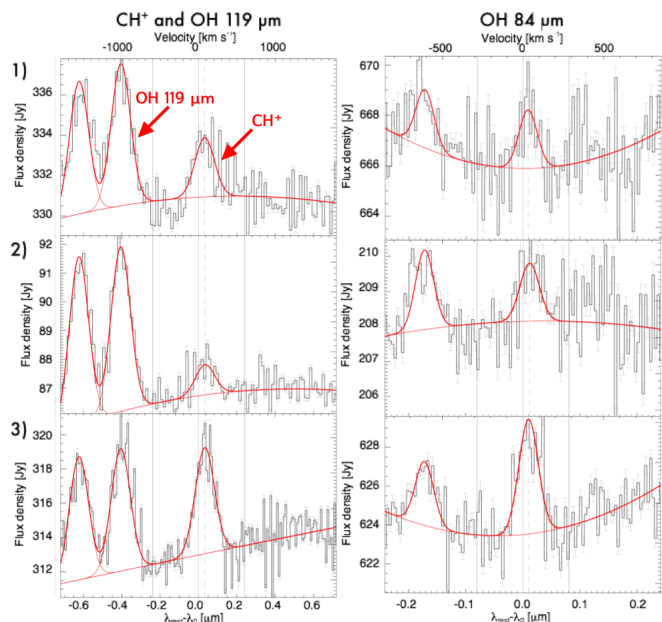
The SPIRE data were processed using HIPE 11.0.1. To achieve a better angular resolution, the super-resolution method SUPREME was applied to the FTS data as in Köhler et al. (2014). In Figs. 3 and 8 we give a FWHM using a Gaussian fit with the same bandwidth of the equivalent beam. It should, however, be noted that the point spread function (PSF) is non-Gaussian and a FWHM does not account for the shape of the beam. The gain in resolution with the SUPREME method depends on the wavelength. The resolution improves, for example, from 32.6'' to 19.0'' at 400  $\mu\text{m}$  or from 40.1'' to 25.8'' at 650  $\mu\text{m}$ . After SUPREME processing, the spatial resolution for the  $\text{CH}^+ J=1-0$ ,  $\text{H}_2\text{O } 399 \mu\text{m}$ , and  $\text{C}^{18}\text{O } J=8-7$  lines is between 17'' and 19'', and for the  $\text{H}_2\text{O } 269 \mu\text{m}$  and  $^{13}\text{CO } J=12-11$  line it is between 13'' and 14''. For a detailed description of the reduction procedure method, fitting routines, and PSF, see the SUPREME web site<sup>3</sup>. The lines were fitted with the HIPE Spectrometer Cube fitting procedure. We assume a total uncertainty of 36% for the integrated line intensities, which includes the calibration and line fitting uncertainties.

### 3. Line emission detection and spatial morphology

#### 3.1. Origin of the $\text{CH}^+$ emission

In this section we present the observations of  $\text{CH}^+ J=3-2$  transition at 120  $\mu\text{m}$  and  $J=1-0$  transition at 359  $\mu\text{m}$ . Maps of the

<sup>3</sup> <http://www.ias.u-psud.fr/supreme/hipeplugin.php>, see also a draft available at: <http://hal.archives-ouvertes.fr/hal-00765929>



**Fig. 2.** Spectra of  $\text{CH}^+ J=3-2$  (120  $\mu\text{m}$ ) and OH 119  $\mu\text{m}$  (left) and OH 84  $\mu\text{m}$  (right) at three positions indicated in panel a of Fig. 3.

integrated line intensities are shown in Fig. 3. We show line spectra for  $\text{CH}^+ J=3-2$  at three positions specified in the maps (2 in the Bar, 1 in front, Fig. 2).

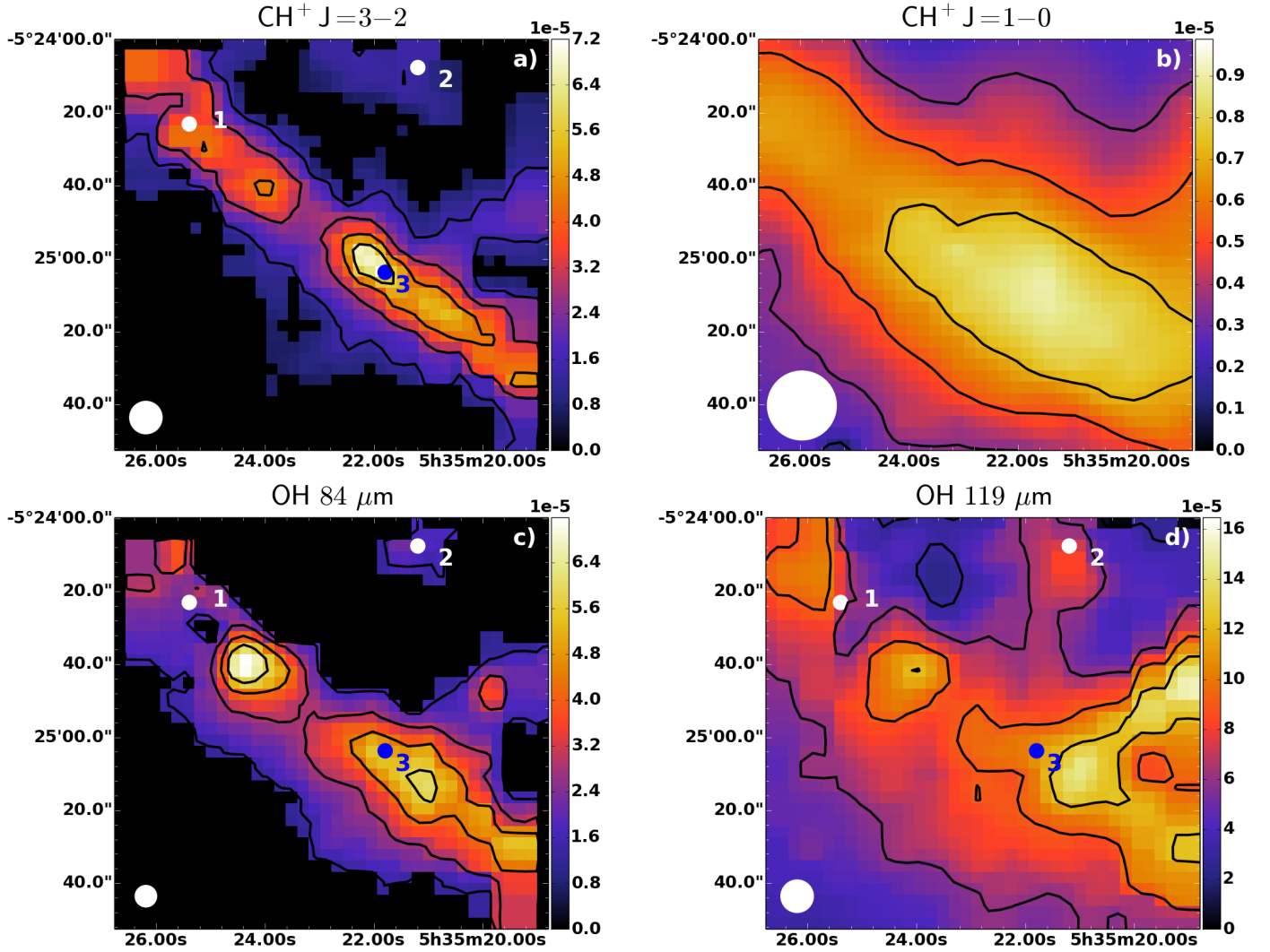
Figure 3 shows that the  $\text{CH}^+ J=3-2$  line emission is fully detected in the Bar, but we do not detect it behind the Bar. The emission delineates the edge of the Bar, which indicates that it is a good tracer of the warm molecular zone and is very sensitive to the physical conditions prevailing in the Bar. The  $\text{CH}^+ J=1-0$  emission is much broader and originates both from the Bar and from the face-on Orion Molecular Cloud surface in the background. The  $J=1-0$  transition is also more easily excited which could affect the breadth of the emission. Emission from the UV-illuminated surface of Orion Molecular Cloud Complex behind the Bar has been seen in several tracers, such as  $\text{C}^+$  (Bernard-Salas et al. 2012; Goicoechea et al. 2015). The  $\text{CH}^+ J=3-2$  is less extended than the  $J=1-0$  line, and may originate from higher density structures in the Bar. In our maps, the  $\text{CH}^+ J=3-2$  line emission that is not associated with the Bar comes from the molecular cloud complex in the northern and western regions of the map (Orion Ridge). Emission in these regions have also been observed in [CII] and [OI] (Bernard-Salas et al. 2012).

Figure 4 shows that the warmest PDR molecular lines between the [NII] 122  $\mu\text{m}$  emission tracing the hot ionization front<sup>4</sup> (Bernard-Salas et al. 2012), and the  $\text{C}^{18}\text{O } J=8-7$  line that traces colder gas. Most of the molecules are photodissociated in the ionized layer in front of the Bar. Behind the Bar and further away, the conditions (low FUV flux, low temperature) are not favorable to form and/or excite  $\text{CH}^+$ .

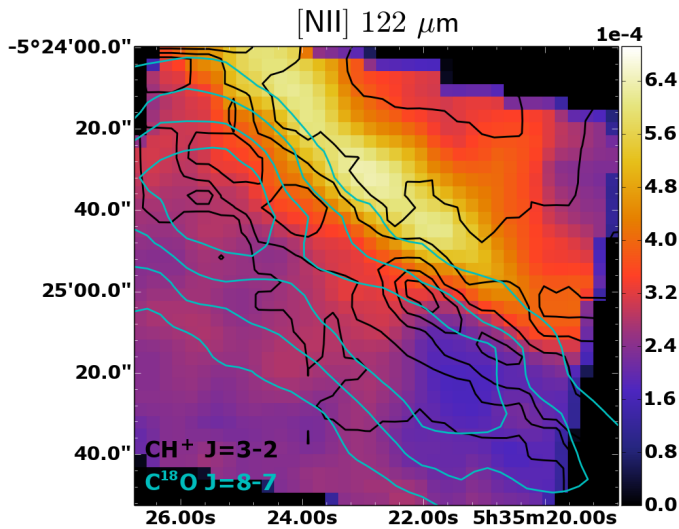
Figure 5 shows the comparison of CO  $J=19-18$  map with  $^{13}\text{CO } J=12-11$ , which is a tracer of the warm and dense phase. Both of these lines show clumps that correspond to those detected from the ground in  $\text{H}^{13}\text{CN}$  (Lis & Schilke 2003) and in  $\text{CS } J=2-1$  (Lee et al. 2013). We can clearly see the clumps<sup>5</sup> 1 and 3, respectively at the north and south part of the Bar. The clumps have diameters of about  $\approx 9''$  or 0.02 pc

<sup>4</sup> [NII] 122  $\mu\text{m}$  peaks in the same place where other ionization front tracers peak, e.g., OI 1.32  $\mu\text{m}$  (Lis & Schilke 2003).

<sup>5</sup> We have adopted the numbering of Lis & Schilke (2003).



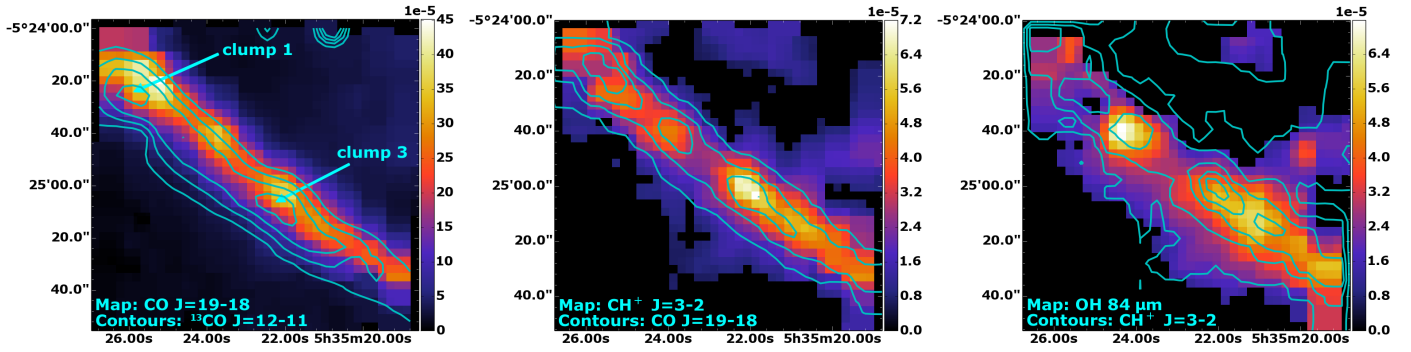
**Fig. 3.** Morphology of different lines: a) CH<sup>+</sup>  $J=3-2$ ; b) CH<sup>+</sup>  $J=1-0$ ; c) OH  $84.6\ \mu\text{m}$ ; and d) OH  $119.4\ \mu\text{m}$ . Contours with contour steps of 20% of the peak emission in the Bar. The intensities are in units of  $\text{erg s}^{-1}\ \text{cm}^{-2}\ \text{sr}^{-1}$ . The black pixels are pixels where lines are not detected. The beam sizes ( $6''$  for OH  $84\ \mu\text{m}$ ,  $9''$  for CH<sup>+</sup>  $J=3-2$  and OH  $119\ \mu\text{m}$ , and  $19''$  for CH<sup>+</sup>  $J=1-0$ ) are shown in the bottom left corner. The positions from which the spectra in Fig. 2 are also indicated in the pictures with corresponding numbers.



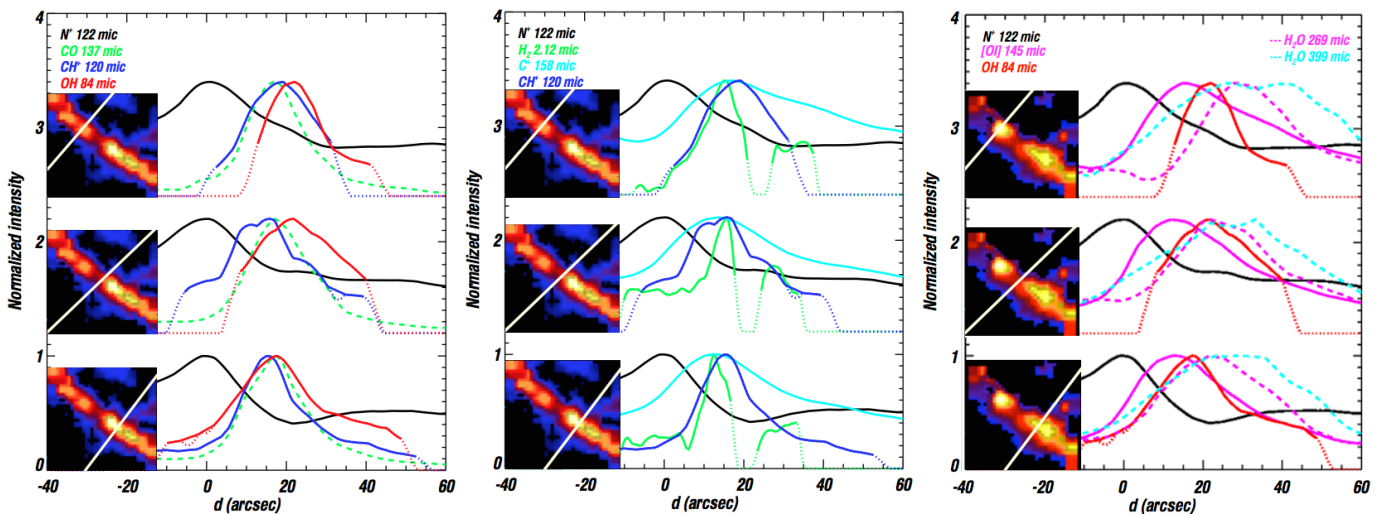
**Fig. 4.** [NII]  $122\ \mu\text{m}$  map with contours of CH<sup>+</sup>  $J=3-2$  (black) and C<sup>18</sup>O  $J=8-7$  (cyan, peak emission  $5.6 \times 10^{-6}$ ). Contours are with contour steps of 20% of the peak emission in the Bar. The intensities are in units of  $\text{erg s}^{-1}\ \text{cm}^{-2}\ \text{sr}^{-1}$ .

and are marginally resolved by *Herschel*. The clumps seen in  $^{13}\text{CO}\ J=12-11$  correspond exactly to those seen in  $\text{H}^{13}\text{CN}$  and  $\text{CS}\ J=2-1$ . The  $\text{CO}\ J=19-18$  map shows the same clumps, but shifted by  $\sim 5''$  toward the PDR front. The CO peaks could, thus, be associated with the surface of the high density zones in the Bar. Figure 5 also compares CH<sup>+</sup>  $J=3-2$  with  $\text{CO}\ J=19-18$  and OH  $84\ \mu\text{m}$  with CH<sup>+</sup>  $J=3-2$ . The clumps can be seen in CH<sup>+</sup>, while in OH we see only clump 3. The intensity peak of CH<sup>+</sup>  $J=3-2$  (and  $J=1-0$ ) coincides with the clump 3, while  $\text{CO}\ J=19-18$  (and C<sup>18</sup>O  $J=8-7$ ) peaks at clump 1. Clump 1 has a larger column density, suggesting that the optically thin high- $J$  CO lines are more sensitive to the total column density of warm gas than CH<sup>+</sup>. Like CH<sup>+</sup>, most of the surface tracers, such as  $\text{H}_2\ v=1-0\ \text{S}(1)$  or optically thick lines (e.g.,  $\text{CO}\ J=6-5$ ), peak at the southern part of the Bar.

Figure 6 presents the profile of the emission of various lines along a cut intersecting the Bar at three different positions. The profile of the CH<sup>+</sup>  $J=3-2$  line emission is resolved with an average FWHM of  $20''$ , which is wider than the PSF ( $9''$ ). This FWHM is similar to that of the  $\text{CO}\ J=19-18$  line. The observed width is most likely the result of the Bar being tilted to the observer, an effect already inferred in other studies



**Fig. 5.** Map of CO  $J = 19-18$  ( $137 \mu\text{m}$ ) integrated intensity with  $^{13}\text{CO } J = 12-11$  contours (*left*). Map of  $\text{CH}^+ J = 3-2$  integrated intensity with CO  $J = 19-18$  contours (*middle*) and OH  $84.6 \mu\text{m}$  with  $\text{CH}^+ J = 3-2$  contours (*right*). All contours are with contour steps of 20% of the peak emission in the Bar starting from 10%, except  $^{13}\text{CO}$  starting from 30%. The intensities are in units of  $\text{erg s}^{-1} \text{cm}^{-2} \text{sr}^{-1}$ . The maps are not convolved as the resolutions are similar. Clumps 1 and 3 (Lis & Schilke 2003) are indicated in the figure on the left.



**Fig. 6.** Profile cuts across the Bar at positions indicated in the small  $\text{CH}^+ J = 3-2$  and OH  $84 \mu\text{m}$  maps. The intensity of the lines have been normalized to the maximum emission in the Bar and the lines are not convolved to the same beam, so that the structures seen in the unconvolved maps are not lost. The distance is from the ionization front indicated by the peak of [NII] emission. The dashed part of the  $\text{H}_2$  profile in the cuts indicates the slit seen in the map (see Fig. 7) and the positions of the cuts are indicated in the small maps next to each cut.

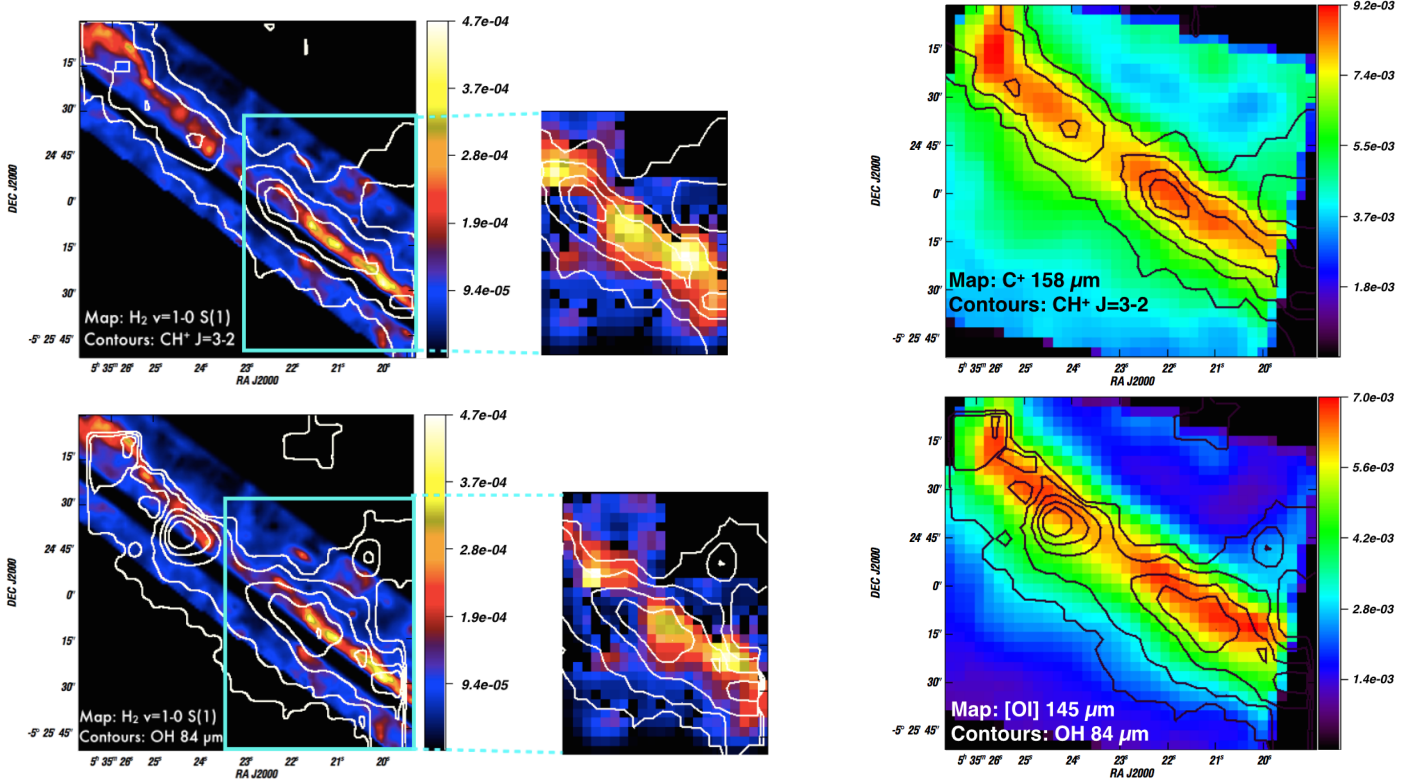
(e.g., Hogerheijde et al. 1995; Allers et al. 2005; Pellegrini et al. 2009). Using an inclination angle of  $\theta = 7^\circ$  (Hogerheijde et al. 1995; Pellegrini et al. 2009), and a length of the PDR along the line of sight of 0.35 pc (Bernard-Salas et al. 2012) needed to reproduce the [CII] and [OI] emission, we find that a width of  $\sin(\theta) \times l_{\text{PDR}} \sim 0.04 \text{ pc}$  (or  $\sim 20''$ ) is in agreement with the observed width. The  $\text{CH}^+$  and high- $J$  CO emission come from the inclined thin surface of the Bar. We stress that higher angular resolution ( $1''$ ) observations of rotationally and vibrationally excited  $\text{H}_2$  originating from the thin PDR surface show a smaller width of  $\sim 10''$  (see Fig. 6) than those observed for  $\text{CH}^+$  and high- $J$  CO. Thus, while our data shows a connection with the clumps seen in  $\text{H}^{13}\text{CO}$  and CS, the  $\text{CH}^+$  and high- $J$  CO emission may also originate from unresolved structures close to the cloud surface. Indeed, small-scale, high density  $\text{HCO}^+ J = 4-3$  emission structures have been recently resolved by ALMA (Goicoechea et al. 2016).

### 3.2. Origin of the OH emission

In this section we present the observations of OH  $84 \mu\text{m}$  and  $119 \mu\text{m}$   $\Lambda$  doublets observed with PACS. The two lines of the  $\Lambda$  doublets have the same distribution and almost the

same intensities, and, therefore, one line map is presented for each wavelength, namely  $84.6 \mu\text{m}$  ( $^2\Pi_{3/2} J = 7/2^- - 5/2^+$ ) and  $119.4 \mu\text{m}$  ( $^2\Pi_{3/2} J = 5/2^+ - 3/2^-$ ). The intensities vary only slightly (by  $\sim 30\%$ , see Fig. 2) between the two  $\Lambda$  doublets for each wavelength. Asymmetries are predicted to be small when collisions with ortho- $\text{H}_2$  dominate (i.e., in the warm gas) and when FIR radiative pumping plays a role (Offer & van Dishoeck 1992; Offer et al. 1994; Goicoechea et al. 2011).

The first rotationally excited OH lines at  $84 \mu\text{m}$  are fully detected in the Bar and also in the Orion Ridge and northern region in front of the Bar (see Fig. 3). Like  $\text{CH}^+ J = 3-2$  in Fig. 5, the OH  $84 \mu\text{m}$   $\Lambda$  doublet traces the warm molecular layer, but the OH emitting zone is more extended behind the Bar (see also Fig. 6). This could result from the fact that the formation of OH does not depend on  $\text{H}_2 v > 0$  like  $\text{CH}^+$  (see Sect. 4) and that the excitation of OH is mainly sensitive to the high density as suggested by previous studies. Goicoechea et al. (2011) suggest that the OH  $84 \mu\text{m}$  emission is coming from unresolved structures with  $n(\text{H}) \sim 10^{6-7} \text{ cm}^{-3}$  and  $T_{\text{kin}} \sim 160-220 \text{ K}$  whereas  $\text{CH}^+$  emission can be explained by a density of  $n(\text{H}) \sim 10^5 \text{ cm}^{-3}$  and  $T_{\text{kin}} \sim 500-1000 \text{ K}$  (Nagy et al. 2013, Joblin et al., in prep.). In the OH  $84 \mu\text{m}$  map, there are two main density enhancements in the Bar, one of which may be associated with a dense structure



**Fig. 7.** Top row: CH<sup>+</sup> 120  $\mu\text{m}$  compared with H<sub>2</sub>  $v = 1-0$  S(1) and [CII] 158  $\mu\text{m}$ . CH<sup>+</sup> 120  $\mu\text{m}$  contours overlaid on the H<sub>2</sub> map (left) and [CII] map (right). Zoom: map of vibrationally excited H<sub>2</sub> convolved to the spatial resolution of CH<sup>+</sup>  $J = 3-2$  map including contours of CH<sup>+</sup>  $J = 3-2$ . Bottom row: OH 84  $\mu\text{m}$  compared with H<sub>2</sub>  $v = 1-0$  S(1) and [OI] 145  $\mu\text{m}$ . OH 84  $\mu\text{m}$  contours overlaid on the H<sub>2</sub> map (left) and [OI] map (right). All contours are with contour steps of 20% of the peak emission, except for the zoomed picture where the contours are with contour steps of 10%. The maps are in units of  $\text{erg s}^{-1} \text{cm}^{-2} \text{sr}^{-1}$ . The H<sub>2</sub>  $v = 1-0$  S(1) data taken from Walmsley et al. (2000, big map) and van der Werf et al. (1996, zoomed map). The [CII] 158  $\mu\text{m}$  and [OI] 145  $\mu\text{m}$  data were taken from Bernard-Salas et al. (2012).

(referred to as clump 3), and another that coincides with an externally illuminated protoplanetary disk, which is discussed in Sect. 6.

The OH ground-state rotational lines at 119  $\mu\text{m}$  are clearly detected everywhere in the mapped area (panel d in Fig. 3). Since the OH 119  $\mu\text{m}$   $\Lambda$  doublets are more easily excited than the 84  $\mu\text{m}$   $\Lambda$  doublets, they show a more extended spatial distribution (including the Orion Ridge and the Molecular Cloud Complex in the north). Goicoechea et al. (2011) used nonlocal, non-LTE radiative transfer model, including FIR pumping, to fit several FIR OH rotational lines (at  $\sim 119 \mu\text{m}$ ,  $\sim 84 \mu\text{m}$ ,  $\sim 163 \mu\text{m}$ ,  $\sim 79 \mu\text{m}$ , and  $\sim 65 \mu\text{m}$ ) detected with *Herschel*/PACS toward a small region<sup>6</sup> in the Orion Bar. The best fit was found for unresolved structures with  $T_{\text{kin}} \sim 200$  K,  $n(\text{H}) \sim 10^6 \text{cm}^{-3}$  and a source-averaged OH column density of  $N(\text{OH}) = 10^{15} \text{cm}^{-2}$ . For these conditions, the OH 119  $\mu\text{m}$   $\Lambda$  doublets are optically thick ( $\tau > 10$ ), but the OH 84  $\mu\text{m}$   $\Lambda$  doublets are optically thin (or effectively optically thin,  $\tau < 5$  and  $T_{\text{ex}} \ll T_{\text{kin}}$ ). Therefore, OH 84  $\mu\text{m}$  line intensities are expected to be proportional to  $N(\text{OH})$  in the Orion Bar.

In conclusion, the OH 84  $\mu\text{m}$  and CH<sup>+</sup>  $J = 3-2$  line emissions correlate well, albeit with some dispersion reflecting the difference in emission discussed above. The observed general correlation was expected given the similar critical densities and upper level energies and a main formation route with H<sub>2</sub> for both lines. However, owing to our higher angular resolution and much larger map, we also observe differences in the morphology.

#### 4. Formation and excitation of CH<sup>+</sup> and OH via vibrationally excited H<sub>2</sub>

In this section, we compare the spatial distribution of the CH<sup>+</sup>  $J = 3-2$  line with those of H<sub>2</sub>  $v = 1-0$  S(1) and [CII] 158  $\mu\text{m}$ , tracing the reactants of  $\text{C}^+ + \text{H}_2 \rightarrow \text{CH}^+ + \text{H}$ . We also compare the spatial distribution of the OH 84  $\mu\text{m}$  emission to those of H<sub>2</sub>  $v = 1-0$  S(1) and [OI] 145  $\mu\text{m}$ , which are the components of reaction  $\text{O} + \text{H}_2 \rightarrow \text{H} + \text{OH}$ .

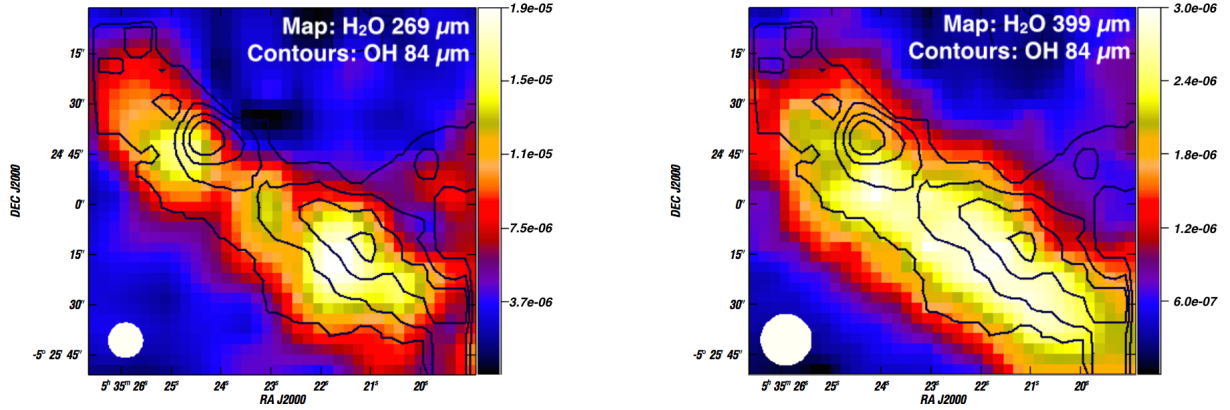
##### 4.1. Spatial comparison of CH<sup>+</sup> with H<sub>2</sub> ( $v = 1$ ) and C<sup>+</sup>

Figure 7 shows a map of H<sub>2</sub>  $v = 1-0$  S(1) at 2.12  $\mu\text{m}$  and the CH<sup>+</sup>  $J = 3-2$  line. The H<sub>2</sub>  $v = 1-0$  S(1) was observed from the ground with a resolution of  $\sim 1''$  (Walmsley et al. 2000). The maps are not convolved, so that we do not blur the detailed structures. At the edge of the Bar, the contours of the CH<sup>+</sup> emission trace the structures seen in H<sub>2</sub> in detail. In the same figure, we also include a smaller H<sub>2</sub> map (van der Werf et al. 1996) which is convolved to the same spatial scale. Again, we find a good correlation between the spatial distribution of the CH<sup>+</sup>  $J = 3-2$  emission and the vibrationally excited H<sub>2</sub> emission, except for clump 3 and in the lower bright structure seen in H<sub>2</sub>. The correlation supports that CH<sup>+</sup> is formed by  $\text{C}^+ + \text{H}_2(v > 0)$  and/or that the excitation of CH<sup>+</sup> is affected by formation pumping (Agúndez et al. 2010; Godard & Cernicharo 2013; Nagy et al. 2013; Zanchet et al. 2013).

Figure 7 also shows that CH<sup>+</sup>  $J = 3-2$  correlates well with [CII] 158  $\mu\text{m}$  (Bernard-Salas et al. 2012). Nevertheless,

<sup>6</sup> One PACS footprint,  $47'' \times 47''$ .





**Fig. 8.** Maps of  $\text{H}_2\text{O } 1_{11}-0_{00}$  at  $269 \mu\text{m}$  (left) and  $\text{H}_2\text{O } 2_{11}-2_{02}$   $399 \mu\text{m}$  (right) with contours of OH  $84.6 \mu\text{m}$  with contour steps of 20% of the peak emission. The maps are in units of  $\text{erg s}^{-1} \text{cm}^{-2} \text{sr}^{-1}$ .

[CII]  $158 \mu\text{m}$  is more extended in front and behind the Bar than  $\text{CH}^+ J=3-2$  (see Figs. 6 and 7). As expected,  $\text{C}^+$  exists in the outermost layers of the PDR, where most of the hydrogen is in atomic form and, thus, little  $\text{CH}^+$  has formed. In the envelope surrounding the Bar, the physical conditions are still favorable for [CII] emission, but not for excited  $\text{CH}^+$ . The [CII]  $158 \mu\text{m}$  line is easily excited because of its low excitation energy of 91 K. The background PDR contributes to 30% of the [CII] emission (Bernard-Salas et al. 2012). The  $\text{CH}^+ J=3-2$  emission is too faint to be detected from the background PDR.

Since the  $\text{C}^+$  is more extended, it appears that the  $\text{CH}^+$  emission is not limited by the  $\text{C}^+$  abundance but by the abundance of vibrationally excited  $\text{H}_2$ . For PDRs with lower FUV radiation field, such as moderately excited PDRs (e.g., Horsehead or NGC 7023 East with radiation fields  $\sim 100$  times lower than in the Orion Bar) the  $\text{CH}^+$  emission, and particularly the excited rotational transitions ( $J > 2$ ), is expected to decrease significantly. Toward the sample of PDRs observed by PACS and SPIRE as a part of the Evolution of Interstellar Dust key program (Abergel et al. 2010, data available on the *Herschel* Idoc Database HESI0D webpage<sup>7</sup>), [CII]  $158 \mu\text{m}$  line is detected in all the PDRs while  $\text{CH}^+ J=1-0$  line is only detected in PDRs associated with a large amount of vibrationally excited  $\text{H}_2$ . The line intensity upper limit for the nondetection of  $\text{CH}^+ (J=1-0)$  in the moderate excited PDRs, such as the Horsehead and NGC 7023 East, is  $2 \times 10^{-7} \text{ erg s}^{-1} \text{cm}^{-2} \text{sr}^{-1}$  with SPIRE (e.g., Köhler et al. 2014). In the Horsehead PDRs, the  $\text{CH}^+ (J=1-0)$  has not even been detected with HIFI within the rms limit of 4 mK. Low FUV radiation field and thermal pressure are clearly not favorable for the  $\text{CH}^+$  line emission.

In conclusion, from our comparison of  $\text{CH}^+$  with vibrationally excited  $\text{H}_2$  and high- $J$  CO,  $\text{CH}^+ J=3-2$  emission is due to both formation pumping and collisions. This agrees with Godard & Cernicharo (2013) and Zanchet et al. (2013), who found that formation pumping dominates the excitation of the  $J \geq 2$  levels in the physical conditions of the Orion Bar.

#### 4.2. Spatial comparison of OH with $\text{H}_2 (v > 0)$ and $\text{O}^0$

In general, the OH  $84 \mu\text{m}$  line map follows the  $\text{H}_2$  emission (see Fig. 7), but the OH emission is more extended and less correlated with  $\text{H}_2 (v=1)$  than  $\text{CH}^+$ . An energy activation barrier exists in the  $\text{O} + \text{H}_2$  reaction even for  $\text{H}_2 (v=3)$

(Sultanov & Balakrishnan 2005; Agúndez et al. 2010), thus, the OH formation is expected to be less sensitive to the presence of  $\text{H}_2 (v=1)$ .

In Fig. 7, we also compare OH  $84 \mu\text{m}$  with [OI]  $145 \mu\text{m}$  line, the other component of the proposed formation route. Similar to [CII], the [OI] emission is more extended and shifted by  $5''$  to the edge of the PDR compared to the OH emission, although the general features are similar. The differences can be explained because [OI] is present in the atomic PDR layers with low molecular gas fractions where  $\text{H}_2$  is photodissociated and cannot participate in forming OH. The background PDR contributes to 15% of the [OI] emission (Bernard-Salas et al. 2012). Like  $\text{CH}^+ J=3-2$ , we do not detect OH  $84 \mu\text{m}$  line from the background PDR.

In conclusion, the spatial comparison of OH with [OI] and  $\text{H}_2 1-0 \text{ S}(1)$  indicates that the reaction  $\text{O}^0 + \text{H}_2 (v=0)$  is the likely OH formation route at the physical conditions of the Orion Bar as predicted by Agúndez et al. (2010). Thus, OH excitation is more sensitive to the temperature and density variations than to the presence of  $\text{H}_2 (v > 0)$ .

#### 4.3. OH and $\text{H}_2\text{O}$ distribution comparison

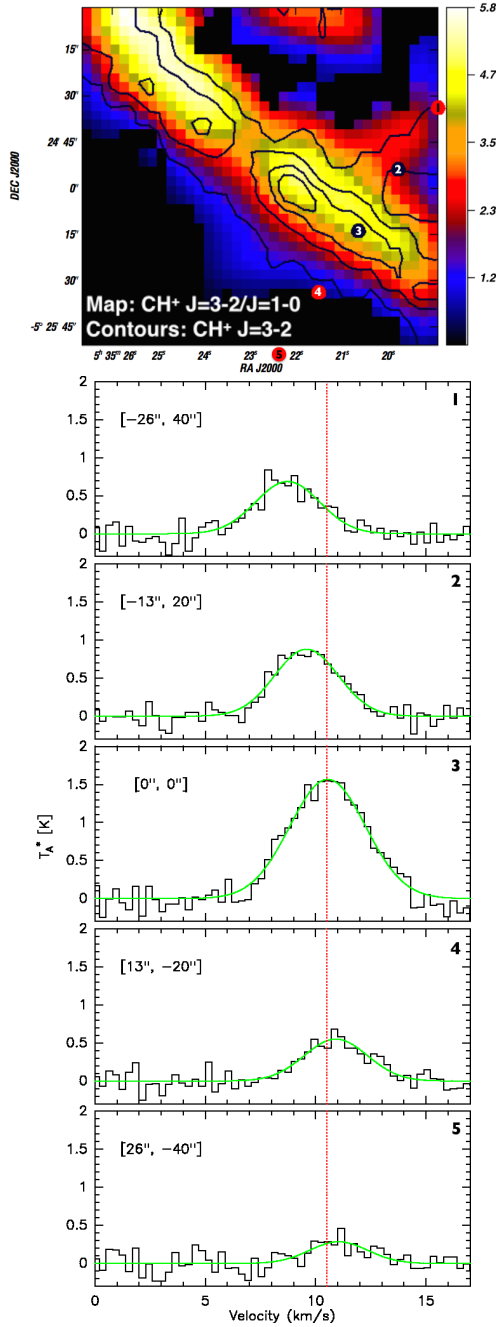
In the warm gas unshielded against FUV radiation, there is an alternative formation route for OH via the photodissociation of  $\text{H}_2\text{O}$  (Hollenbach et al. 2009, 2012). Photodissociation is the main destruction route of  $\text{H}_2\text{O}$  and OH. The photodissociation is predicted to limit the  $\text{H}_2\text{O}$  abundance in the warm zone and the peak of  $\text{H}_2\text{O}$  emission is expected to be deeper in the PDR (Hollenbach et al. 2009). Here, we compare the OH  $84 \mu\text{m}$  line with two water lines at  $269 \mu\text{m}$  ( $1_{11}-0_{00}$ ,  $E_u/k = 53 \text{ K}$ ) and  $399 \mu\text{m}$  ( $2_{11}-2_{02}$ ,  $E_u/k = 137 \text{ K}$ ) in Fig. 8.

Both  $\text{H}_2\text{O}$  lines, more easily excited than OH  $84 \mu\text{m}$ , peak behind the OH by  $\sim 15''$  and are more extended (see the cut in Fig. 6). This could be explained by the lower gas temperature in the regions in which the bulk of the  $\text{H}_2\text{O}$  arises compared to OH. The photodissociation is likely to limit the water abundance in the warm zone. Thus, the bulk of the OH and water emission arises from different depths. It also implies that this formation route is less important than the  $\text{O}^0 + \text{H}_2 \rightarrow \text{H} + \text{OH}$  reaction.

## 5. $\text{CH}^+$ kinematics

To investigate gas kinematics and  $\text{CH}^+$  line width variations, we examine the  $\text{CH}^+ J=2-1$  ( $180 \mu\text{m}$ ) line. We have velocity-resolved line profiles of a strip across the southwestern part of

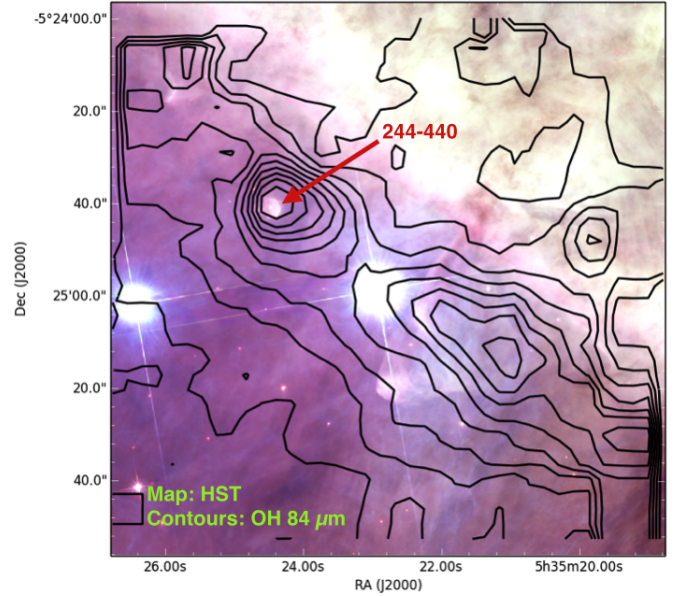
<sup>7</sup> <http://idoc-herschel.ias.u-psud.fr/sitools/client-user/>



**Fig. 9.** Upper panel: ratio of CH<sup>+</sup>  $J=3-2$  and CH<sup>+</sup>  $J=1-0$  integrated intensity with contours of CH<sup>+</sup>  $J=3-2$  in steps of 20% of the peak emission to indicate the position of the Bar. Lower panels: HIFI spectra of CH<sup>+</sup>  $J=2-1$  at the positions across the Bar indicated in the map with corresponding numbering.

the Bar observed with HIFI (see Fig. 1 for configuration). A selection of the spectra obtained along the cut are shown in Fig. 9. The selected positions are  $\sim 23''$  apart from each other. The spectra were obtained by averaging all the scans that fall within  $4''$  of the selected position.

The line centroid velocity varies from  $9 \text{ km s}^{-1}$  in front of the Bar to  $11 \text{ km s}^{-1}$  in the Bar and behind the Bar (see Table 2). These centroid velocities are similar to what Peng et al. (2012) found for CO transitions from  $J=6-5$  to  $J=8-7$ . We find a broad line width ranging from  $\Delta V = 3.4 \pm 0.3$ – $3.1 \pm 0.7$  (in front and behind the Bar) to  $4 \pm 0.2 \text{ km s}^{-1}$  (in the Bar). This



**Fig. 10.** Contours of OH  $84.6 \mu\text{m}$  at every 10 % of the emission peak overlaid on an optical image from the HST (Robberto et al. 2013).

is consistent with the observations of Nagy et al. (2013) on the CO<sup>+</sup> peak who found line widths of  $\Delta V = 5.46 \pm 0.04 \text{ km s}^{-1}$  and  $4.57 \pm 0.11 \text{ km s}^{-1}$  for CH<sup>+</sup>  $J=1-0$  and  $J=2-1$ , respectively. The CH<sup>+</sup>  $J=2-1$  line appears narrower than the  $J=1-0$  line, which could be due to opacity broadening and to the larger beam size at the lower frequency. CH<sup>+</sup> has a broader line width than any other molecular tracer of dense gas in the Orion Bar (typically showing line widths of  $2-3 \text{ km s}^{-1}$ ; Joblin et al., in prep.). The CH<sup>+</sup>  $J=2-1$  line appears to be broad all over the region, even though the ratio of the CH<sup>+</sup>  $J=3-2$  to  $J=1-0$  shown in Fig. 9 points to different physical conditions.

The broad width of CH<sup>+</sup> line could originate in a low density interclump component ( $n \sim 10^4 \text{ cm}^{-3}$ ), where low density tracers, such as [CII], have been observed to have similar line widths of  $\sim 4 \text{ km s}^{-1}$  (Nagy et al. 2013). However, our maps show that the excited CH<sup>+</sup> and high- $J$  CO both originate both in the dense and warm region of the Bar, while their line widths are different ( $3 \text{ km s}^{-1}$  for the high- $J$  CO; Joblin et al., in prep). Alternatively, based on reactivity of the molecule, Nagy et al. (2013) suggest that the broadening could be explained by formation pumping. If this is the case, there is no time for the CH<sup>+</sup> to thermalize, and the line width remains unchanged. In conclusion, our results agree with Nagy et al. (2013) that the broad line width of CH<sup>+</sup> is most likely a signature of formation pumping.

## 6. Proplyd in the Orion Bar

The OH  $84 \mu\text{m}$  map reveals a clear emission peak that is not detected in the other lines. In Fig. 10, we compare the OH  $84 \mu\text{m}$  emission to an optical image from the *Hubble Space Telescope* (HST). The HST image reveals an optical source at the OH peak position. This source has been identified as the proplyd 244-440 (e.g., Bally et al. 2000). Proplyds, or externally illuminated dense protoplanetary disks, are a special class of low-mass young stellar objects modulated by a strong, external UV field. They are found embedded within or near a HII region and are identified by their typical cometary photoionized envelopes.

This is, to our knowledge, the first time that OH emission has been tentatively associated with a proplyd, although OH

has been observed in a protoplanetary disk in the DIGIT survey (Sturm et al. 2010). In Fig. 10, the highest OH contour level at the proplyd position is  $8''$ , as our resolution is  $6''$ , which corresponds to 3000 AU at a distance of 415 pc. While disks in proplyds are small (100–500 AU), their photoevaporating envelopes can extend up to a couple of 1000 AU. The OH  $84 \mu\text{m}$  line emission at this position could come from UV-radiation heating or shocks. The HST optical image of the proplyd 244–440 reveals the presence of microjets. Alternatively, in proplyds, the surface of the disk is illuminated and is, thus, similar to a very dense and warm PDR, and is also expected to produce FIR OH emission. If the OH origin is UV heating, the proplyd detection would confirm that the OH does originate in hot irradiated structures.

Like OH,  $\text{CH}^+ J=3-2$  has been detected toward protoplanetary flared disks such as HD 100546 and HD 97048 (Thi et al. 2011; Fedele et al. 2013). In the upper layers of the disks, which are directly exposed to the stellar radiation field, the reaction  $\text{C}^+ + \text{H}_2 (v > 0) \rightarrow \text{CH}^+ + \text{H}$  is expected to be dominant. We do not see any excess of  $\text{CH}^+ J=3-2$  in the Orion Bar, nor any excess of high- $J$  CO emission. As opposed to OH,  $\text{CH}^+$  and high- $J$  CO lines are fainter in disks (e.g., Fedele et al. 2013; Meeus et al. 2013) and their emission is likely overshadowed by the strong nebular emission in the region.

## 7. Summary and conclusions

We have presented *Herschel* fully sampled maps of the FIR rotationally excited emission lines of  $\text{CH}^+$  and OH toward the Orion Bar over a large area ( $\sim 110'' \times 110''$ ). The  $\text{CH}^+$  and OH lines delineate the edge of the Bar indicating that they are good tracers of the warm molecular zone and very sensitive to the physical conditions. Our main results are summarized as follows:

1. We confirm the correlation of OH and  $\text{CH}^+$  with high- $J$  CO emission (a tracer of dense gas). We find that the spatial thickness of the observed line emission layers results from the Bar being tilted toward the observer, an effect already inferred in other studies. OH,  $\text{CH}^+$ , and high- $J$  CO originate in the inclined thin irradiated surface of the Bar with embedded dense structures. The emission peaks of these lines are associated with the surface of the largest clumps (density enhancements) seen with high density tracers, such as HCN or CS.
2. Although excited  $\text{CH}^+$  and OH display similar overall spatial distribution, there are also relevant differences in their morphology. The good correlation of  $\text{CH}^+$  with vibrationally excited  $\text{H}_2$  supports that  $\text{CH}^+$  is formed via the reaction  $\text{C}^+ + \text{H}_2 (v > 0)$ . We provide observational evidence to the fact that the excitation of the  $\text{CH}^+$  rotational levels is affected by formation pumping, which dominates the excitation for  $J \geq 2$  levels, as suggested by previous models. The  $\text{CH}^+$  emission is limited by the abundance of vibrationally excited  $\text{H}_2$  and not by the  $\text{C}^+$  abundance. We find the  $\text{C}^+$  emission to be more extended than the  $\text{CH}^+$  emission.
3. OH is less well correlated with the vibrationally excited  $\text{H}_2$ , which is a fact pointed out by theoretical studies predicting that the reaction  $\text{O}^0 + \text{H}_2 (v = 0) \rightarrow \text{H} + \text{OH}$  dominates the OH formation. The other formation route for OH via the photodissociation of  $\text{H}_2\text{O}$  appears not to be as important since the bulk of the OH and water emission seem to arise from different cloud depths.
4. We find a broad line width of  $\sim 3-4 \text{ km s}^{-1}$  for  $\text{CH}^+ J=2-1$ , which is broader than other molecular line tracing dense gas in the Orion Bar. The line width most likely reflects the excess energy transferred to translational energy during the formation pumping of  $\text{CH}^+$ .
5. Interestingly, the peak of the OH  $84 \mu\text{m}$  emission corresponds to the position of the proplyd 244–440. Far-infrared OH emission could be a good diagnostic of hot gas in externally illuminated protoplanetary disks.

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