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**Synthesis, crystal structure, and electrical and magnetic characterizations of the monoclinic compounds  $\text{Ln}_3\text{Mo}_{4+x}\text{Si}_{1-x}\text{O}_{14}$  ( $\text{Ln} = \text{La}, \text{Ce}, \text{Pr}, \text{and Nd}; X = 0$  and  $0.35$ ) containing  $\text{Mo}_3$  clusters and  $\text{Mo}_2$  dimers**

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**Synthesis, crystal structure, and electrical and magnetic characterizations of the monoclinic compounds  $\text{Ln}_3\text{Mo}_{4+x}\text{Si}_{1-x}\text{O}_{14}$  ( $\text{Ln} = \text{La, Ce, Pr, and Nd}$ ;  $x = 0$  and  $0.35$ ) containing  $\text{Mo}_3$  clusters and  $\text{Mo}_2$  dimers**

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## Abstract

Powder samples of the new monoclinic compounds  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  ( $\text{Ln} = \text{La}, \text{Ce}, \text{Pr}$  and  $\text{Nd}$ ) and single crystals of  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  were obtained by solid state reaction. The crystal structure of  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  was determined by single-crystal X-ray diffraction.  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  crystallizes in the monoclinic space group  $\text{P}2_1/\text{n}$  with unit-cell parameters  $a = 5.6361(2) \text{ \AA}$ ,  $b = 17.5814(8) \text{ \AA}$ ,  $c = 10.9883(4) \text{ \AA}$  and  $Z = 4$ . Full-matrix least-squares refinement on  $F^2$  using 7544 independent reflections for 203 refinable parameters results in  $R_1 = 0.0359$  and  $wR_2 = 0.0831$ . The structure contains chains of  $\text{Mo}_3\text{O}_{13}$  clusters and chains of edge-sharing  $\text{MoO}_6$  octahedra with alternately short ( $2.508 \text{ \AA}$ ) and long ( $3.162 \text{ \AA}$ ) Mo-Mo distances running parallel to the  $a$  axis that are separated by eight- or ten-coordinate Pr-O polyhedra. Magnetic susceptibility measurements on the  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  ( $\text{Ln} = \text{La}, \text{Ce}, \text{Pr}$  and  $\text{Nd}$ ) are in agreement with a trivalent state of the rare earths for the Ce, Pr and Nd compounds, while that on the lanthanum compound confirm the presence of one unpaired electron per  $\text{Mo}_3$  as expected. Resistivity measurements on a single crystal show that  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  is a small band gap semi-conductor.

## Keywords

Mo cluster oxides,  $\text{MoO}_2$  type chains,  $\text{Mo}_3$  triangle, rare earth, magnetic susceptibility properties, resistivity measurement.

## Introduction

The triangular  $\text{Mo}_3$  cluster has been observed for the first time in solid state chemistry in the series  $\text{M}_2\text{Mo}_3\text{O}_8$  ( $\text{M} = \text{Mg}, \text{Mn}, \text{Fe}, \text{Co}, \text{Ni}, \text{Zn}$  and  $\text{Cd}$ ) synthesized by McCarroll et al. in 1957 [1]. During the last decade, such compounds containing  $\text{Mo}_3$  triangle clusters have experienced a resurgence of interest both in terms of purely fundamental studies as well as for their potential applications. Owing to their non-centrosymmetric polar crystal structure (space group  $\text{P6}_3\text{mc}$ ), the  $\text{M}_2\text{Mo}_3\text{O}_8$  ( $\text{M} = \text{Mn}, \text{Fe}, \text{Co}$  and  $\text{Ni}$ ) compounds become multiferroics with very important magneto-electric coupling. Among the latter compounds,  $\text{Fe}_2\text{Mo}_3\text{O}_8$  has been particularly studied [2-7]. Thus, a large tunable magnetoelectric effect [3, 4, 6] as well as an optical diode effect [5] have been observed in Zn-doped  $\text{Fe}_2\text{Mo}_3\text{O}_8$ . A very large thermal Hall effect has been also measured on crystals of the solid solution  $(\text{Fe}_{1-x}\text{Zn}_x)_2\text{Mo}_3\text{O}_8$  [7], showing for the first time the coupling between magnetism and phonon in thermal transport properties.  $\text{Fe}_2\text{Mo}_3\text{O}_8$  either alone or with graphene oxide or reduced graphene oxides has also been studied for Li storage in batteries [8, 9]. The cluster magnet  $\text{LiZn}_2\text{Mo}_3\text{O}_8$  compound has also been the subject of very interesting studies by the group of T. M. McQueen [10-13]

In the 1980s, McCarroll reported the crystal structures of the isomorphous orthorhombic  $\text{La}_3\text{Mo}_4\text{XO}_{14}$  ( $\text{X} = \text{Si}$  or  $\text{Mo}_{0.33}\text{Al}_{0.67}$ ) compounds [14, 15]. Both compounds were prepared by electrolytic reduction at 1373 K of melts containing  $\text{Na}_2\text{MoO}_4$ ,  $\text{MoO}_3$ , and  $\text{La}_2\text{O}_3$  in silica or alumina crucibles. The composition of the crystals which grown at the platinum cathode depends upon the type of crucible used to contain the melt. Indeed, the silica crucible produces  $\text{La}_3\text{Mo}_4\text{SiO}_{14}$  while the alumina crucible yields  $\text{La}_3\text{Mo}_4(\text{Mo}_{0.33}\text{Al}_{0.67})\text{O}_{14}$ . The crystal structure of the latter two compounds contains triangular  $\text{Mo}_3\text{O}_{13}$  cluster units with Mo-Mo distances of about 2.56 Å and chains of edge-sharing  $\text{MoO}_6$  octahedral with alternately short (ca. 2.53 Å) and long (ca. 3.16 Å) Mo-Mo distances. The Si and the Al and Mo atoms are tetrahedrally

coordinated and the La atoms are surrounded by 7, 8 or 10 O atoms forming bi-, tri- or tetra-capped distorted trigonal prisms

We present here the crystal structure of  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  which constitutes a monoclinic variant of the  $\text{La}_3\text{Mo}_4\text{XO}_{14}$  structure-type. Similarly, to the previous compounds,  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  was obtained serendipitously by solid-state reaction. As for  $\text{La}_3\text{Mo}_4\text{SiO}_{14}$ , Si was abstracted from the silica tube used to contain the reactants since the starting products were  $\text{Pr}_6\text{O}_{11}$ ,  $\text{Gd}_2\text{O}_3$ ,  $\text{MoO}_3$ , and Mo. Subsequently, we have prepared, by solid state reaction, polycrystalline samples of the monoclinic series  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  ( $\text{X} = \text{La}, \text{Ce}, \text{Pr}$  and  $\text{Nd}$ ). Synthesis, and electrical resistivity and magnetic susceptibility measurements of the latter compounds are also presented in this paper.

## 1. Experimental

### 2.1 Syntheses

Some poor quality crystals of the monoclinic  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  compound were obtained serendipitously during the synthesis of  $\text{Pr}_3\text{MoO}_7$  by solid-state reaction in sealed silica tube at 1200 °C. As for the orthorhombic  $\text{La}_3\text{Mo}_4\text{SiO}_{14}$  compound synthesized by fused salt electrolysis, Si was abstracted from the silica tube used to contain the reactants. Crystals thus obtained have the shape of needles with quasi-rectangular cross-section. Qualitative microanalyses using a JEOL JSM 6400 scanning electron microscope equipped with an OXFORD link INCA energy-dispersive- X-ray analyser confirmed the presence of silicon in the crystals. Subsequently, powders of the monoclinic  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  ( $\text{Ln} = \text{La}, \text{Ce}, \text{Pr}$  and  $\text{Nd}$ )

compounds were prepared from stoichiometric mixtures of the rare-earth oxides ( $\text{La}_2\text{O}_3$ ,  $\text{CeO}_2$ ,  $\text{Pr}_6\text{O}_{11}$  or  $\text{Nd}_2\text{O}_3$ ),  $\text{MoO}_3$  and Mo. Prior reactions, the molybdenum powder was heated at  $1000^\circ\text{C}$  under a  $\text{H}_2$  flow during 6 hours. The rare-earth oxides were heated at  $1000^\circ\text{C}$  during 12 hours and after, left at  $800^\circ\text{C}$  before syntheses. After homogenization, the powders were pelletized (approximately 5 g) and introduced into Mo crucibles, which were degassed beforehand at temperatures of approximately  $1500^\circ\text{C}$  for 15 min under a dynamic vacuum of approximately  $10^{-2}$  Pa. The Mo crucibles were then sealed under an Ar pressure of about 60 Pa by arc welding. The samples were subsequently heated to the temperature of  $1300^\circ\text{C}$  with a rate of rise of  $300^\circ\text{C}/\text{h}$ . The duration of the plateau at  $1300^\circ\text{C}$  was 48 hours and then the whole was subsequently cooled to  $1100^\circ\text{C}$  at the rate of  $100^\circ\text{C}/\text{h}$ , temperature at which the oven was stopped and cooled itself down to room temperature. The microcrystalline powder samples, thus obtained, were found to be monophasic based on their X-ray diffraction diagram collected on a Bruker D8 advance diffractometer equipped with a LynxEye detector ( $\text{CuK}\alpha_1$  radiation) and indexed in a monoclinic unit cell as exemplified by that of  $\text{Pr}_3\text{Mo}_4\text{SiO}_{14}$  in Figure 1. The lowering of the orthorhombic symmetry of the unit-cell to monoclinic results in splitting of some peaks as exemplified by the (101) and (111) reflections in Figure 2. Unit cell parameters of the  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  compounds determined by profile matching fits with the program JANA2006 [16] in the space group  $\text{P}2_1/\text{n}$  are given in table 1. Figure 3 shows a decrease of the unit-cell volume from the La to the Nd compound. This behaviour is in agreement with the well-known lanthanide contraction very often observed in the series of lanthanide-based compounds and confirms that all rare earths are in a trivalent state in our phases. Attempts to grow single crystals of the monoclinic  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  ( $\text{Ln} = \text{La}, \text{Ce}, \text{Pr}$  and  $\text{Nd}$ ) compounds by a solid-state procedure were unsuccessful up to now.

## 2.2 *Single-crystal study*

Several crystals of  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  were tested on a Nonius Kappa CCD diffractometer using graphite-monochromated Mo- $K\alpha$  radiation ( $\lambda = 0.71073 \text{ \AA}$ ). All of them were of poor quality, and present a substantial mosaic spread due to a twinning. Nevertheless, the totality of the reflections was identified with a primitive monoclinic unit-cell after introducing a twinning matrix corresponding to a twofold rotation about the c axis. The COLLECT program [17] was employed to determine the angular scan conditions ( $\varphi$  and  $\omega$  scans) for the intensity recording. The raw data were subsequently integrated with the package EvalCCD [18]. The analytical method as developed by de Meulenaar and Tompa was employed for the correction of the absorption after the indexation of the faces of the crystal investigated [19]. Analysis of the intensity data shows that the reflections (h0l) with  $h + l = 2n + 1$  as well as (0k0)  $k = 2n + 1$  were systematically absent, leading to the monoclinic space group  $P2_1/n$ . The crystal structure was solved by direct methods using the software SIR97 [20] and subsequently refined with SHELXL97 [21] in the monoclinic  $P2_1/n$  space group. Analysis of the data with the TwinRotMat procedure implemented in PLATON [22] confirmed that the crystal investigated was twinned with the two crystalline components related by a [001] twofold axis. Experimental crystal data and structure refinement results for  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  are recapitulated in Table 2, and the principal interatomic bond lengths are given in Table 3

### *2.3 Electrical Resistivity Study*

The electrical resistivity measurements were realized on a single crystal of  $\text{Pr}_3\text{Mo}_{4.34}\text{Si}_{0.66}\text{O}_{14}$  with an ac current of  $1 \mu\text{A}$  and a frequency of 80 Hz using the classical four-probe method. The four indium contacts were attached to the crystal by an ultrasonic soldering iron.

### *2.4 Magnetic susceptibility characterizations*

Magnetic susceptibility measurements were made on microcrystalline powder samples of about 100 mg on a *MPMS XL* Magnetometer from Quantum Design in the temperature range 4.2 - 300 K and under a magnetic field of 0.5 T.

## 2. Results and discussion

### 3.1 Crystal structure

The figure 4 represent a projected view of the crystal structure of  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  along the *a* axis. The Mo-O network is similar to that of the orthorhombic  $\text{La}_3\text{Mo}_4\text{XO}_{14}$  compounds. It is based on chains of the well-known triangular  $\text{Mo}_3\text{O}_{13}$  unit (Fig. 3) and chains of  $\text{MoO}_6$  octahedra of rutile-type as observed in the molybdenum dioxide  $\text{MoO}_2$  [23]. Each rutile type Mo chain is connected via the oxygen atoms O4 and O9 to two chains made up of  $\text{Mo}_3$  triangles to create a kind of ribbon, which develops along the *b* axis (fig.2). The most important difference between the orthorhombic and monoclinic variants arises from the loss of the mirror plane perpendicular to the smallest axis. Indeed, in the orthorhombic phase, it was necessary to move off the mirror plane the unique crystallographically Mo atom forming the edge sharing  $\text{MoO}_6$  chains with a site occupancy factor of 0.5 to describe the alternately long and short Mo-Mo distances encountered in this  $\text{MoO}_2$ -type chain. In the monoclinic compound, such an arrangement could be accommodated without disorder with one Mo atom, Mo4, in general position. The Mo-Mo distances along the chain are 2.5083(7) and 3.1609(7) Å in the monoclinic compound compared to 2.535(8) and 3.167(8) Å in the previous orthorhombic compounds. In the trimer  $\text{Mo}_3$ , the three Mo-Mo bond lengths are 2.5400(6), 2.5455(6) and 2.5584(4) Å leading to an average value of 2.548 Å in the range 2.47-2.58 Å of what is observed in the  $\text{Mo}_3$  triangles found in the reduced molybdenum oxides containing such units. The absence of the mirror plane also reduces the symmetry of the  $\text{Mo}_3\text{O}_{13}$  unit as well as that of the rare-earth sites from .m. to 1. The Mo-O bond lengths around the Mo1, Mo2 and Mo3 atoms of the  $\text{Mo}_3$  cluster



are in the range 1.928(4) - 2.132(4) Å compared to 1.918(3) - 2.107(4) Å around the Mo4 atoms of the Mo<sub>2</sub> dimer. From the lengths of the Mo-O bonds, we were able to obtain an estimate of the oxidation state of each independent Mo atom from the bond-length bond-strength equation,  $s(\text{Mo-O}) = [d(\text{Mo-O})/1.882]^{-6}$ , established by Brown and Wu [24]. In this formula,  $s(\text{Mo-O})$  is the bond strength in valence unit,  $d(\text{Mo-O})$  the observed Mo-O bond distance in Å, 1.882 Å corresponds to a Mo-O single bond distance and the exponential parameter -6 is characteristic of the Mo atom. Consequently, for the Mo1, Mo2 and Mo3 atoms forming the Mo<sub>3</sub> cluster, the values of +4.03(5), +3.49(5) and +3.71(5) were obtained, and for Mo4 of the MoO<sub>2</sub>-type chain we got +3.77(4). Consequently, the average valence for the Mo atoms of the Mo<sub>3</sub> clusters is +3.74, and +3.77 for the Mo atoms belonging to the MoO<sub>2</sub>-type chains. From these two values, we can deduce value of 6.8 e<sup>-</sup>/Mo<sub>3</sub> and 3.8 e<sup>-</sup>/Mo<sub>2</sub> in good agreement with the expected value of 7 e<sup>-</sup>/Mo<sub>3</sub> and 4 e<sup>-</sup>/Mo<sub>2</sub> deduced from the formula. For the orthorhombic La<sub>3</sub>Mo<sub>4</sub>SiO<sub>14</sub>, quasi-similar calculated values of +3.76 and +3.84 are found while in La<sub>3</sub>Mo<sub>4</sub>Al<sub>2/3</sub>Mo<sub>1/3</sub>O<sub>14</sub>, they slightly differ due the trivalent state of the aluminium with values of +3.90 for the Mo<sub>3</sub> cluster and +3.93 for those in the MoO<sub>2</sub>-type chains. In table 4, we have summarized the number of electrons per Mo<sub>3</sub> cluster deduced from stoichiometry and from the Brown and Wu relationship for compounds containing Mo<sub>3</sub> clusters the crystal structures of which were determined on single crystals. These calculations of bond valences have been used with success in numerous reduced molybdenum containing various Mo clusters and chains for many years. As we can, there is no correlation between the number of electrons bears by the triangular cluster and the distances between the atoms of molybdenum that form it. This is particularly highlighted with the compounds having clusters with 8 electrons and in which the Mo-Mo distances vary from one extreme to the other indicating that matrix effects are more important than electronic ones.

In Pr<sub>3</sub>Mo<sub>4.35</sub>Si<sub>0.65</sub>O<sub>14</sub>, the Si atom is tetrahedrally coordinated by O atoms with Si-O distance ranging from 1.460(5) to 1.750(5) Å. The relationship of Brown and Wu was also applied to

the Si-O bonds ( $s = [d(\text{Si-O})/1.622]^{-4.429}$ ). The oxidation state of +4 for the Si atoms was confirmed with a calculated value of +3.9.

The Pr atoms are surrounded by 7, 8 or 10 oxygen atoms forming complex polyhedral based, in each case, on a trigonal prism of oxygen atoms. For the atoms Pr3 and Pr1, one or two pseudo-rectangular faces are capped, respectively, and, for Pr2, the two distorted triangular faces are also capped. In these three sites, the Pr-O distances range from 2.334(5) to 2.936(6) Å with average distances of 2.50 Å for the Pr1 and Pr3 sites (coordination number (CN) 8 and 7, respectively) and 2.63 Å for the Pr2 site (CN 10). The latter values compare well with those deduced from the sum of the ionic radii of  $\text{O}^{2-}$  and  $\text{Pr}^{3+}$  reported by Shannon, respectively 2.52 Å for the Pr1 and Pr3 sites and 2.56 Å for the Pr2 site [34]. From the bond length-bond strength formula, the oxidation states of +3.16, +2.91 and +2.70 were calculated for the Pr1, Pr2 and Pr3 atoms, respectively (average value +2.93).

### *3.2 Electrical and magnetic characterizations*

Electrical resistivity measurements were performed on a single-crystal of  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$ . The thermal variation of the electrical resistivity was measured along the a axis corresponding to the direction of propagation of the cluster chains.  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$  presents a semiconductor behaviour between 100 and 290 K as depicted by the curves shown in figure 8 that represent the thermal variation of the resistivity and  $\log(\rho)$  as a function of  $1000/T$ . The resistivity measured at 290K is 0.6  $\Omega\cdot\text{cm}$  and the activation energies deduced from the Arrhenius plot is 0.06 eV. Those values agree well with those published by Betteridge et al. in 1984 and McCarroll et al. in 1986 for the orthorhombic variants with room temperature resistivities in the range 0.2 - 0.9  $\Omega\cdot\text{cm}$  and activation energies around 0.05 eV.

The reciprocal magnetic susceptibility of  $\text{La}_3\text{Mo}_4\text{SiO}_{14}$  as a function of the temperature is shown in Fig. 9. Above 150 K, the curve of the reciprocal magnetic susceptibility as a function of the temperature follows a modified Curie-Weiss law that could be fitted with a constant of Curie  $C = 0.248$  leading to a magnetic moment of  $1.41 \mu_B$  per  $\text{Mo}_3$  cluster if we assume that no magnetic moment arises from the  $\text{MoO}_2$  rutile-type chains. The latter value, lower than the expected one of  $1.73 \mu_B$  for a single unpaired spin, confirms a configuration of seven electrons on each  $\text{Mo}_3$  clusters in  $\text{La}_3\text{Mo}_4\text{SiO}_{14}$ . The experimental magnetic moments observed for the reduced molybdenum oxides with non-magnetic cations based on triangular  $\text{Mo}_3$  clusters having 7 electrons are given in table 5. Chen et al. [35] introduce a phenomenological parameter  $\lambda$  to characterize the anisotropy of the Mo Kagome lattice such that  $\lambda = d(\text{Mo-Mo})_{\text{inter}}/d(\text{Mo-Mo})_{\text{intra}}$  where  $d(\text{Mo-Mo})_{\text{inter}}$  and  $d(\text{Mo-Mo})_{\text{intra}}$  are the shortest Mo-Mo distances between and within the  $\text{Mo}_3$  clusters, respectively. Later, Haraguchi et al. [36, 37] have demonstrated that the effective magnetic moment ( $\mu_{\text{eff}}$ ) is correlated to  $\lambda$  that they rename parameter of clusterization. Thus, when  $\lambda$  is sufficiently greater than one, the system can be categorized as a localized spin model while when  $\lambda$  is close to one, a weak Mott regime close to a metallic state is envisaged. Although, in our case, the  $\text{Mo}_3$  clusters did not form a 2D-Kagome lattice but a kind of “1D”, we can see on figure 10 that the value  $\mu_{\text{eff}}$  observed for  $\text{La}_3\text{Mo}_4\text{SiO}_{14}$  fits well with those of previous studies when we plot  $\mu_{\text{eff}}$  as a function of  $\lambda$  for the reduced Mo compounds containing magnet  $\text{Mo}_3$  aggregates, known to date. Among all of these compounds, it is interesting to note that  $\text{LiZn}_2\text{Mo}_3\text{O}_8$  has a magnetic moment of  $1.39 \mu_B$ , correspondent to constant of Curie  $C = 0.24$ , at temperatures greater than 96 K while at lower temperatures, the magnetic moment is reduced to  $0.8 \mu_B$  ( $C = 0.08$ ). According to Sheckelton et al. in 2012 [10], the diminution of the moment arises from the creation of an exotic condensed valence-bond ground state by 2/3 of the paramagnetic spins. At 14 K, the plot of the thermal variation of the magnetic susceptibility as a function of the temperature of  $\text{La}_3\text{Mo}_4\text{SiO}_{14}$  (Fig.

9) presents a maximum arising probably from the apparition of a short range magnetic ordering between the triangular  $\text{Mo}_3$  clusters as previously observed in  $\text{Li}_2\text{InMo}_3\text{O}_8$  that was the first reduced Mo oxide presenting a long range magnetic ordering between the triangular clusters. In solid-state metal cluster chemistry, long-range ordering between unpaired electron clusters was first observed in the  $\text{GaMo}_4\text{X}_8$  ( $\text{X} = \text{S}, \text{Se}$ ) compounds [38] made in our laboratory and based on tetrahedral  $\text{Mo}_4$  clusters with 11 electrons and that present ferromagnetic or antiferromagnetic interactions. A little more lately, antiferromagnetic interactions between octahedral  $\text{Nb}_6$  cluster clusters containing 15 electrons were clearly evidenced by  $^{19}\text{F}$  NMR and EPR measurements in the  $\text{LuNb}_6\text{Cl}_{18}$  and  $\text{Nb}_6\text{F}_{15}$  compounds [39, 40]. well as for the compound  $\text{NdMo}_8\text{O}_{14}$  [41] containing  $\text{Mo}_8$  clusters with 23 electrons.

The temperature dependence of the inverse of the molar magnetic susceptibility of  $\text{Ce}_3\text{Mo}_4\text{SiO}_{14}$ ,  $\text{Pr}_3\text{Mo}_4\text{SiO}_{14}$  and  $\text{Nd}_3\text{Mo}_4\text{SiO}_{14}$  are presented in Fig. 11. The susceptibility data for the Ce, Pr and Nd analogues show strong temperature dependencies and could be fit to modified Curie-Weiss like behaviours in the temperature range 100 – 300K. A least squares fitting of the observed after subtracting the contribution of the  $\text{Mo}_3$  clusters resulted in a  $C = 2.11$  emu.K/mole,  $\theta = -14.72$  K and  $\chi_0 = 2.46 \cdot 10^{-3}$ ,  $C = 4.74$  emu.K/mole,  $\theta = -31.9$  K and  $\chi_0 = 2.52 \cdot 10^{-3}$ , and  $C = 4.65$  emu.K/mole,  $\theta = -27.7$  K and  $\chi_0 = 3.01 \cdot 10^{-3}$  for the Ce, Pr and Nd compounds, respectively. The observed effective magnetic moments ( $\mu_{\text{eff.}} = 2.37 \mu_{\text{B}}/\text{Ce}$ ,  $3.55 \mu_{\text{B}}/\text{Pr}$  and  $3.52 \mu_{\text{B}}/\text{Nd}$ ) are in good agreement with the theoretically values of 2.54, 3.58 and  $3.62 \mu_{\text{B}}$  expected from Hund's rule [42]. In the case of the Pr compound, although the low symmetry of the  $\text{Pr}^{3+}$  sites, we did not observe an almost constant susceptibility below 10 K, indicating a non-magnetic single ground state as expected from the non-Kramer nature of the  $\text{Pr}^{3+}$  ions ( $^3\text{H}_4$  ground state).

### 3. Conclusion

In summary, we have showed that the  $\text{Ln}_3\text{Mo}_{4+x}\text{Si}_{1-x}\text{O}_{14}$  ( $\text{Ln} = \text{La}, \text{Ce}, \text{Pr}, \text{and Nd}; x = 0 \text{ and } 0.33$ ) compounds can be synthesized by solid state chemistry and, crystallize in the monoclinic space group  $\text{P}2_1/\text{n}$  contrary to the  $\text{La}_3\text{Mo}_4\text{XO}_{14}$  ( $x = \text{Si or Mo}_{0.33}\text{Al}_{0.67}$ ) compounds prepared by electrolytic reduction at 1373 K of melts containing  $\text{Na}_2\text{MoO}_4$ ,  $\text{MoO}_3$ , and  $\text{La}_2\text{O}_3$  in silica or alumina crucibles which crystallize in the orthorhombic space group  $\text{Pbam}$ . The lowering of the space-group symmetry suppresses the disordering of the Mo atoms forming the rutile-type  $\text{Mo}_2$ -chains. An electrical semi-conducting behaviour as previously found for the orthorhombic phases was also observed for  $\text{Pr}_3\text{Mo}_{4.35}\text{Si}_{0.65}\text{O}_{14}$ . Magnetic susceptibility measurements as a function of the temperature show a paramagnetic spin behaviour that confirm the presence of one unpaired electron per  $\text{Mo}_3$  as expected from value of  $7 e^-/\text{Mo}_3$  and  $4 e^-/\text{Mo}_2$  deduced from the empirical bond length-bond strength relationship. Finally, the temperature dependence of the magnetic susceptibility shows a maximum at 14 K arising probably from an antiferromagnetic ordering between the  $\text{Mo}_3$  clusters.

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**Table Captions:**

Table 1. Unit cell parameters of the  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  compounds determined by profile matching fits.

Table 2. Crystal data and structure refinement for  $\text{Pr}_3\text{Mo}_4(\text{Mo}_{0.35}\text{Si}_{0.65})\text{O}_{14}$ .

Table 3. Selected Interatomic Distances for  $\text{Pr}_3\text{Mo}_4(\text{Mo}_{0.35}\text{Si}_{0.65})\text{O}_{14}$ .

Table 4. Number of electrons per  $\text{Mo}_3$  cluster deduced from stoichiometry and from the Brown and Wu relationship for the compounds containing  $\text{Mo}_3$  clusters.

Table 5. Effective magnetic moment for the compounds containing  $\text{Mo}_3$  clusters with 7 electrons.

## Figure Captions

Figure 1. Observed (dotted line), calculated (red line) and difference profiles for the refinement of  $\text{Pr}_3\text{Mo}_4\text{SiO}_{14}$  in profile-matching mode ( $\lambda = 1.5406 \text{ \AA}$ ).

Figure 2. Detail of the RX diagram of  $\text{Pr}_3\text{Mo}_4\text{SiO}_{14}$  showing the lowering the unit-cell symmetry.

Figure 3. Variation of the unit cell volume for the monoclinic  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  ( $\text{Ln} = \text{La}$  to  $\text{Nd}$ ) compounds as a function of the  $\text{Ln}^{3+}$  ionic radius in eight coordination.

Figure 4. The crystal structure of  $\text{Pr}_3\text{Mo}_4\text{SiO}_{14}$  as viewed down the  $a$  axis, parallel to the direction of the  $\text{Mo}_3\text{O}_{13}$  and  $\text{MoO}_6$  chain growth. Ellipsoids are drawn at the 97 % probability level.

Figure 5. View of the  $\text{Mo}_3\text{O}_{13}$  cluster (ellipsoids at the 97 % probability level).

Figure 6. Fragment of the  $\text{MoO}_2$ -type chain.

Figure 7. Oxygen environments for the Pr atoms (ellipsoids at the 97 % probability level).

Figure 8.

Figure 9. Temperature dependence of the inverse magnetic susceptibility for  $\text{La}_3\text{Mo}_4\text{SiO}_{14}$ .

Figure 10. Plot of  $\mu_{\text{eff}}$  as a function of the parameter  $\lambda = d(\text{Mo-Mo})_{\text{inter}}/d(\text{Mo-Mo})_{\text{intra}}$

Figure 11. Reciprocal magnetic susceptibility of the  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  ( $\text{Ln} = \text{Ce}$ ,  $\text{Pr}$ , and  $\text{Nd}$ ) compounds as a function of temperature. Data were taken under an applied field of 0.1 T.

	a (Å)	b (Å)	c (Å)	$\beta$ (°)	Vol(Å <sup>3</sup> )
La <sub>3</sub> Mo <sub>4</sub> SiO <sub>14</sub>	5.6383(1)	17.6818(3)	11.0456(2)	90.2810(7)	1101.19(6)
Ce <sub>3</sub> Mo <sub>4</sub> SiO <sub>14</sub>	5.6269(3)	17.6095(9)	10.9728(5)	90.2759(15)	1087.24(16)
Pr <sub>3</sub> Mo <sub>4</sub> SiO <sub>14</sub>	5.6185(1)	17.5598(2)	10.9201(2)	90.3118(6)	1077.36(4)
Nd <sub>3</sub> Mo <sub>4</sub> SiO <sub>14</sub>	5.6107(1)	17.5281(1)	10.8635(1)	90.3015(4)	1068.37(2)

Table 1. Unit cell parameters of the Ln<sub>3</sub>Mo<sub>4</sub>SiO<sub>14</sub> compounds determined by profile matching fits.

Table 2. Crystal data and structure refinements of Pr<sub>3</sub>Mo<sub>4</sub>(Mo<sub>0.35</sub>Si<sub>0.65</sub>)O<sub>14</sub>

Empirical formula	Pr <sub>3</sub> Mo <sub>4</sub> (Mo <sub>0.348(4)</sub> Si <sub>0.652(4)</sub> )O <sub>14</sub>
Formula weight (g mol <sup>-1</sup> )	1081.65
Crystal system, space group	monoclinic, P2 <sub>1</sub> /n
Unit cell dimensions (Å, deg)	$a = 5.6361 (2)$ $b = 17.5814 (8)$ $c = 10.9883 (4)$ $\beta = 90.1654 (15)$
Volume (Å <sup>3</sup> )	1088.83 (7)
Z, Calculated density (g/cm <sup>3</sup> )	4, 6.659
Absorption coefficient (mm <sup>-1</sup> )	18.140
Crystal color and habit	Black, needle like
Crystal size (mm <sup>3</sup> )	0.181 × 0.034 × 0.025
Theta range for data collection (deg)	3.61–40.00
Limiting indices	$-10 \leq h \leq 10$ , $-25 \leq k \leq 31$ , $-19 \leq l \leq 19$
Reflections collected/unique	36566/6749
R (int)	0.0454
Absorption correction	analytical
Max./min. transmission	0.6709/0.2026
Data/restraints/parameters	6749/0/202
Goodness-of-fit on $F^2$	1.121
R indices [ $I > 2\sigma(I)$ ]	R1=0.0302, wR2=0.0672
Extinction coefficient	0.00036(5)
Largest diff. peak and hole (eÅ <sup>-3</sup> )	3.716 and -3.663

Table 3. Selected Interatomic Distances for Pr<sub>3</sub>Mo<sub>4</sub>(Mo<sub>0.35</sub>Si<sub>0.65</sub>)O<sub>14</sub>.

Mo(1)-Mo(3)	2.5400(6)	Mo(3)-O(10)	1.981(4)
Mo(1)-Mo(2)	2.5455(6)	Mo(3)-O(4)	2.011(4)
Mo(2)-Mo(3)	2.5584(4)	Mo(3)-O(2)	2.020(4)
Mo(1)-O(5)	1.928(4)	Mo(3)-O(9)	2.064(4)
Mo(1)-O(7)	1.946(4)	Mo(3)-O(1)	2.085(4)
Mo(1)-O(2)	2.021(4)	Mo(3)-O(11)	2.108(5)
Mo(1)-O(3)	2.027(4)	Mo4-Mo4	2.5083(7)
Mo(1)-O(1)	2.048(3)	Mo4-Mo4	3.1609(7)
Mo(1)-O(6)	2.132(4)	Mo(4)-O(9)	1.918(3)
Mo(2)-O(10)	2.001(4)	Mo(4)-O(13)	2.033(5)
Mo(2)-O(3)	2.034(5)	Mo(4)-O(13)	2.035(4)
Mo(2)-O(4)	2.044(3)	Mo(4)-O(4)	2.055(3)
Mo(2)-O(1)	2.083(4)	Mo(4)-O(12)	2.081(4)
Mo(2)-O(8)	2.085(6)	Mo(4)-O(12)	2.107(4)
Mo(2)-O(9)	2.094(4)		
Pr(1)-O(7)	2.334(5)	Pr(3)-O(7)	2.406(4)
Pr(1)-O(5)	2.358(5)	Pr(3)-O(13)	2.411(4)
Pr(1)-O(10)	2.454(3)	Pr(3)-O(12)	2.469(5)
Pr(1)-O(13)	2.473(4)	Pr(3)-O(8)	2.482(4)
Pr(1)-O(3)	2.529(4)	Pr(3)-O(11)	2.510(5)
Pr(1)-O(2)	2.546(4)	Pr(3)-O(5)	2.543(4)
Pr(1)-O(12)	2.569(4)	Pr(3)-O(6)	2.684(5)
Pr(1)-O(4)	2.742(4)		
Pr(2)-O(14)	2.405(5)	Si/Mo5-O14	1.460(5)
Pr(2)-O(2)	2.496(4)	Si/Mo5-O6	1.724(4)
Pr(2)-O(3)	2.504(4)	Si/Mo5-O11	1.741(5)
Pr(2)-O(5)	2.521(5)	Si/Mo5-O8	1.750(5)
Pr(2)-O(7)	2.571(5)		
Pr(2)-O(1)	2.578(3)		
Pr(2)-O(11)	2.761(6)		
Pr(2)-O(8)	2.775(6)		
Pr(2)-O(6)	2.802(6)		
Pr(2)-O(6)	2.936(6)		

Table 4. Number of electrons per Mo<sub>3</sub> cluster deduced from stoichiometry and from the Brown and Wu relationship as well as average or Mo-Mo distances (Å) in the Mo<sub>3</sub> clusters for compounds containing Mo<sub>3</sub> clusters the crystal structure of which was determined on single crystals.

		e <sup>-</sup> /Mo <sub>3</sub> (stoichiometry)	e <sup>-</sup> /Mo <sub>3</sub> (Brown and Wu)	Average or Mo-Mo distance (in Å) in the Mo <sub>3</sub> clusters	Space group	Reference
Zn <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	6	5.8	2.5326(2)	P6 <sub>3</sub> mc	25
Zn <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	6	6.0	2.5325(6)	P6 <sub>3</sub> mc	26
Mg <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	6	5.8	2.5271(2)	P6 <sub>3</sub> mc	26
Co <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	6	5.8	2.5274(2)	P6 <sub>3</sub> mc	26
Mn <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	6	6.0	2.5391(6)	P6 <sub>3</sub> mc	26
La <sub>5</sub> Mo <sub>6</sub> O <sub>21</sub>	Mo <sub>3</sub> + Mo <sup>5+</sup>	6	6.3	2.5598(6)	P2 <sub>1</sub> 2 <sub>1</sub> 2 <sub>1</sub>	27
Na <sub>2</sub> In <sub>2</sub> Mo <sub>5</sub> O <sub>16</sub>	Mo <sub>3</sub> +Mo <sup>6+</sup>	6	6.3	2.6164(5)	P3m1	28
La <sub>3</sub> Mo <sub>4</sub> Al <sub>2/3</sub> Mo <sub>1/3</sub> O <sub>14</sub>	Mo <sub>3</sub> +Mo <sub>2</sub>	6.33	6	2.568	Pnma	15
La <sub>3</sub> Mo <sub>4</sub> SiO <sub>14</sub>	Mo <sub>3</sub> +Mo <sub>2</sub>	7	6.5	2.555	Pnma	14
Na <sub>3</sub> Fe <sub>2</sub> Mo <sub>5</sub> O <sub>16</sub>	Mo <sub>3</sub> +Mo <sup>6+</sup>	7	6.5	2.567(1)	P-3m1	29
Pr <sub>3</sub> Mo <sub>4</sub> (Mo <sub>0.35</sub> Si <sub>0.65</sub> )O <sub>14</sub>	Mo <sub>3</sub> +Mo <sub>2</sub>	7	6.8	2.548	P2 <sub>1</sub> /n	This study
LiZn <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	7	7.4	2.578(1)	R-3m	30
Li <sub>2</sub> InMo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	7	7.4	2.5455(4)	P6 <sub>3</sub> mc	31
Li <sub>2</sub> SnMo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	8	8	2.5036(7)	P6 <sub>3</sub> mc	32
Li <sub>2</sub> GeMo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	8	7.9	2.4728(8)	P6 <sub>3</sub> mc	33
Zn <sub>3</sub> Mo <sub>3</sub> O <sub>8</sub>	Mo <sub>3</sub>	8	8.3	2.580(2)	R-3m	30





Table 5. Effective magnetic moment for the compounds containing Mo<sub>3</sub> clusters with 7 electrons and non-magnetic cations.

	$\mu_{\text{eff}}$ ( $\mu\text{B}$ )	Reference
La <sub>3</sub> Mo <sub>4</sub> SiO <sub>14</sub>	1.41	This work
LiZn <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	1.39 (T>96 K)	20
LiZn <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	1.2	15
ScZnMo <sub>3</sub> O <sub>8</sub>	1.50	15
Li <sub>2</sub> InMo <sub>3</sub> O <sub>8</sub>	1.50	16
Li <sub>2</sub> ScMo <sub>3</sub> O <sub>8</sub>	1.65	22
Na <sub>3</sub> In <sub>2</sub> (MoO <sub>4</sub> ) <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	1.43	21
Na <sub>3</sub> Sc <sub>2</sub> (MoO <sub>4</sub> ) <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	1.46	21
Na <sub>2</sub> In <sub>2</sub> (MoO <sub>4</sub> ) <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	0.366	21
Na <sub>2</sub> Sc <sub>2</sub> (MoO <sub>4</sub> ) <sub>2</sub> Mo <sub>3</sub> O <sub>8</sub>	0.340	21

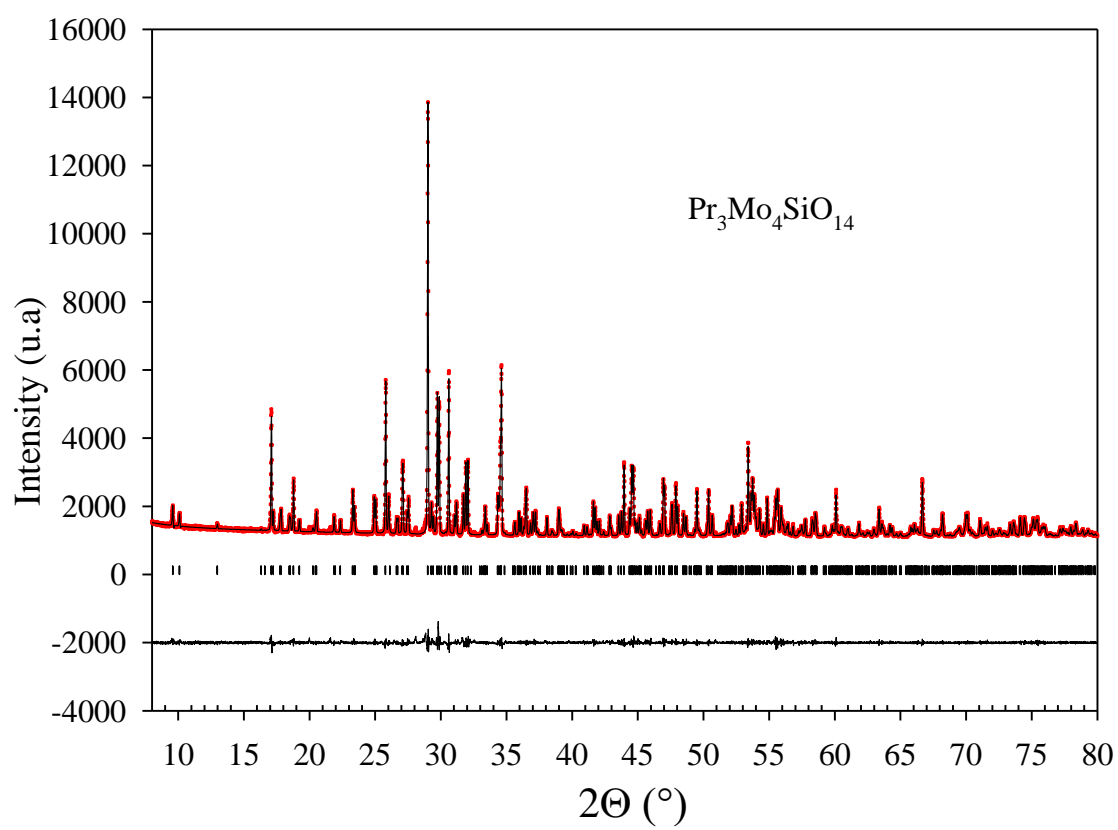
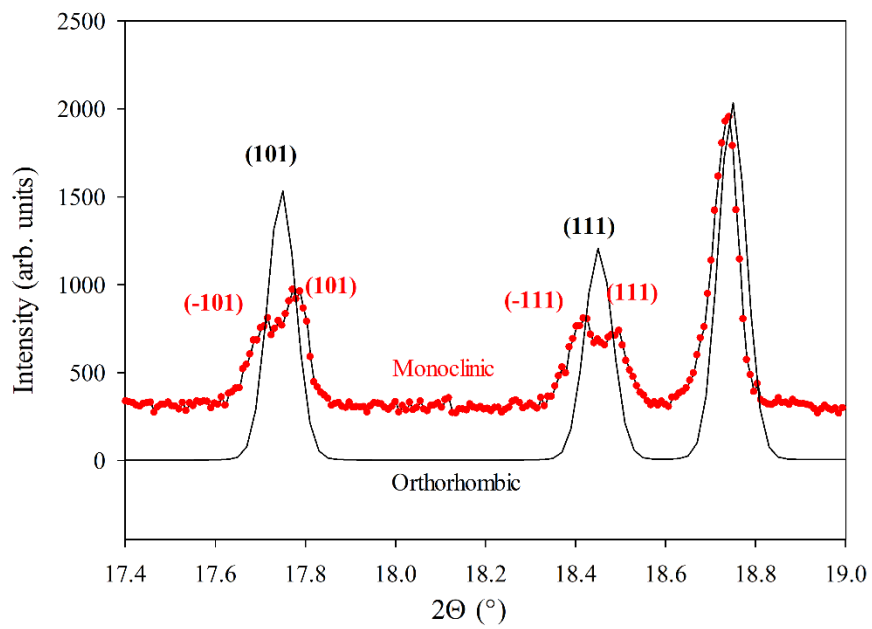


Figure 1



**Figure 2**

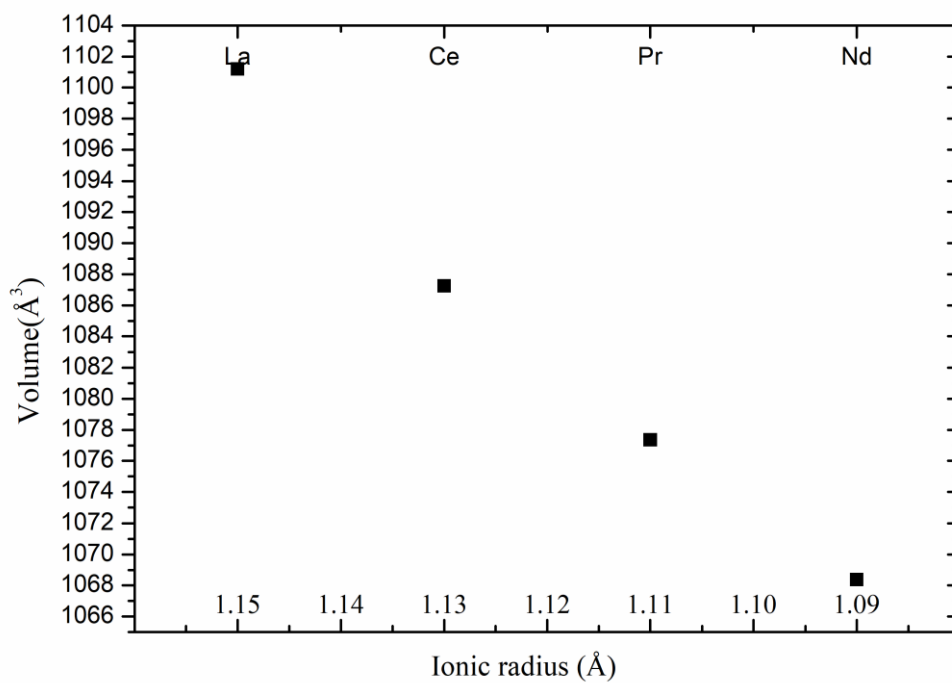


Figure 3. Variation of the unit cell volume for the monoclinic  $\text{Ln}_3\text{Mo}_4\text{SiO}_{14}$  ( $\text{Ln} = \text{La}$  to  $\text{Nd}$ ) compounds as a function of the  $\text{Ln}^{3+}$  ionic radius in eight coordination.

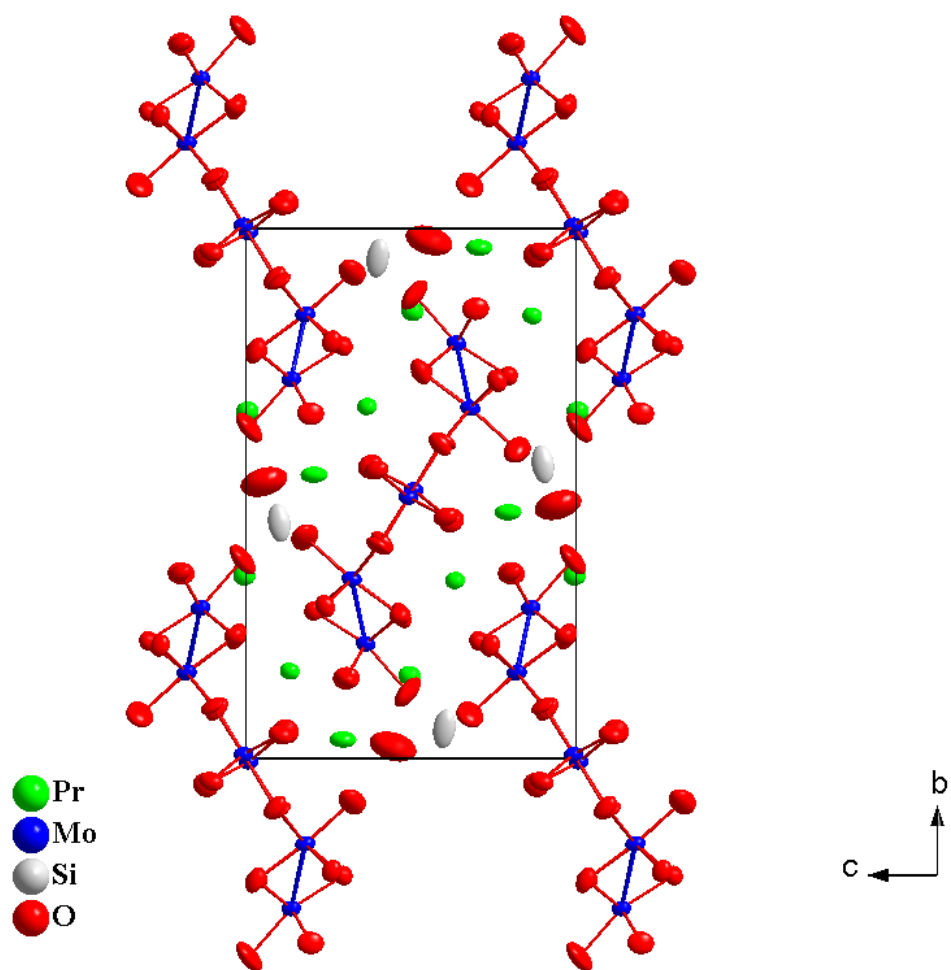


Figure 4

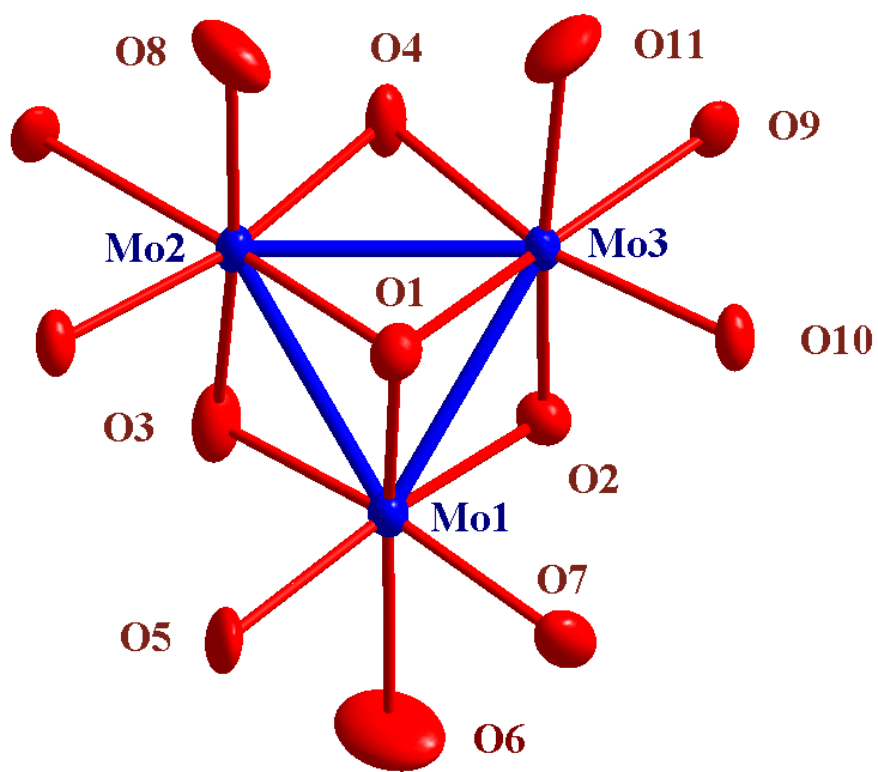


Figure 5

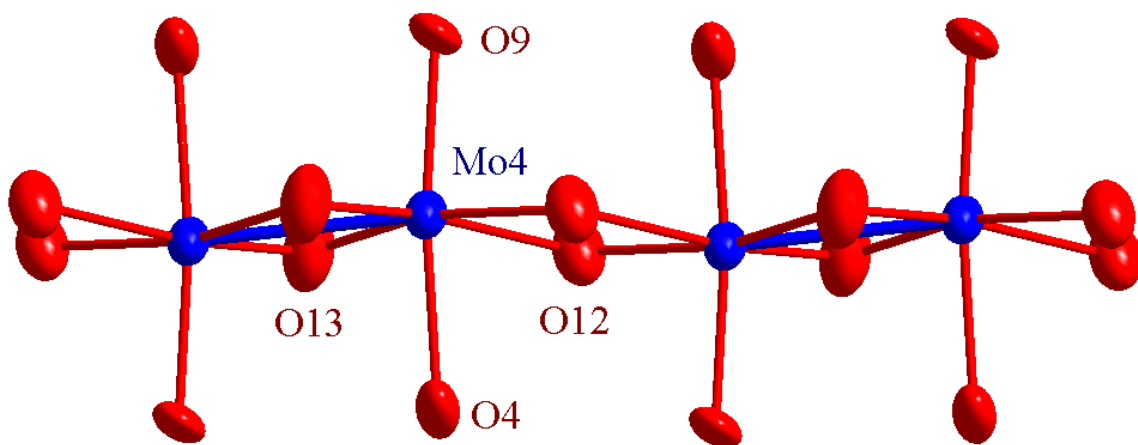


Figure 6



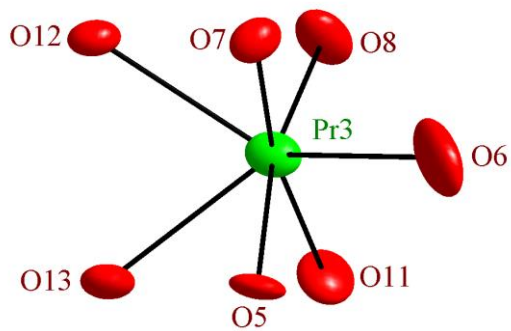
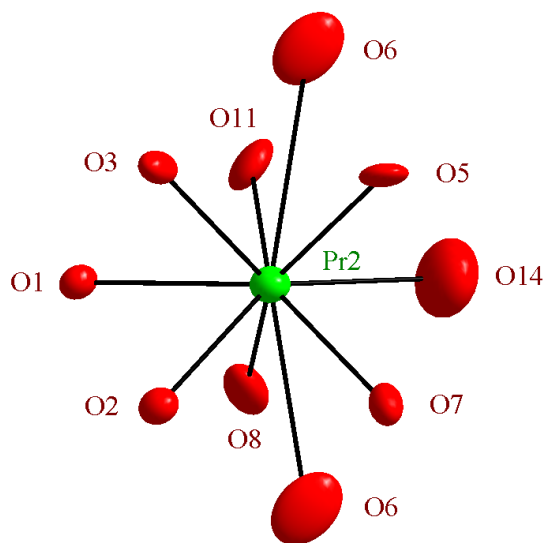
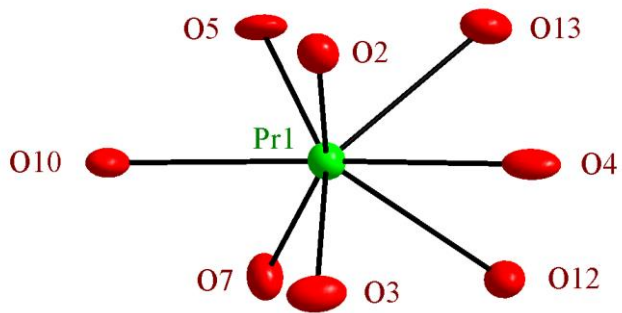


Figure 7

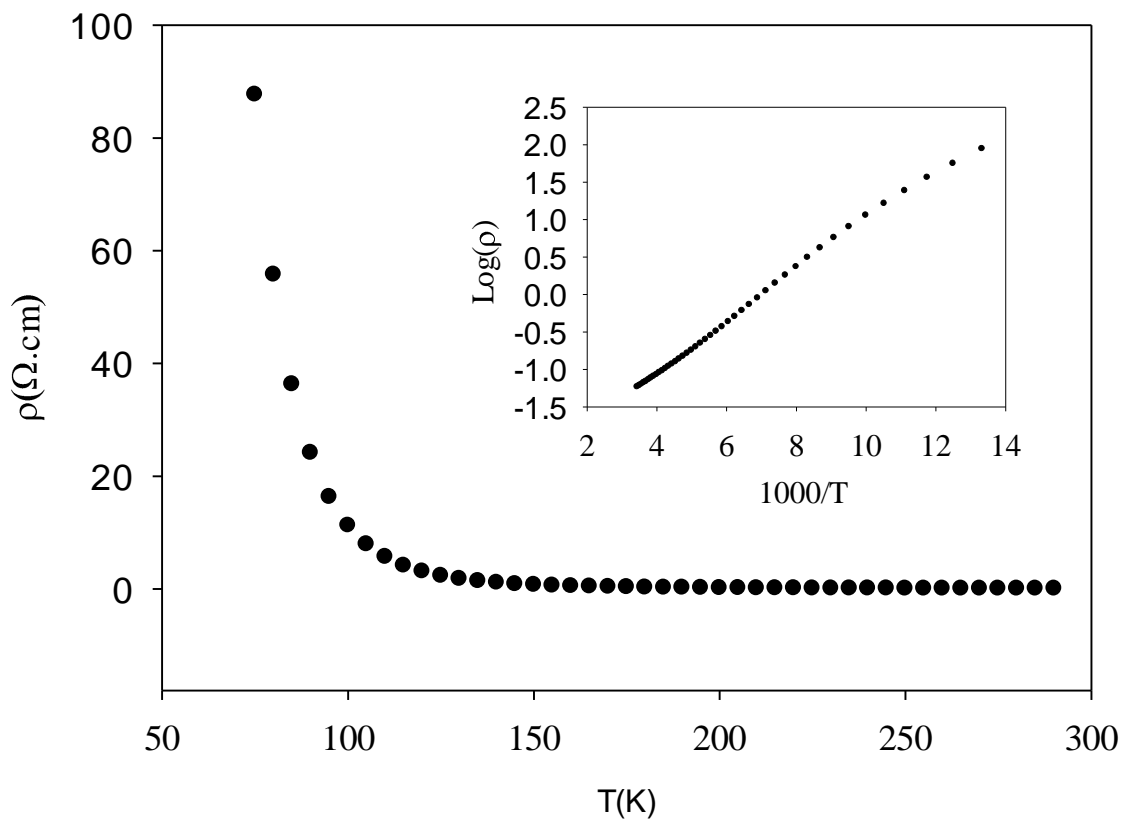


Figure 8

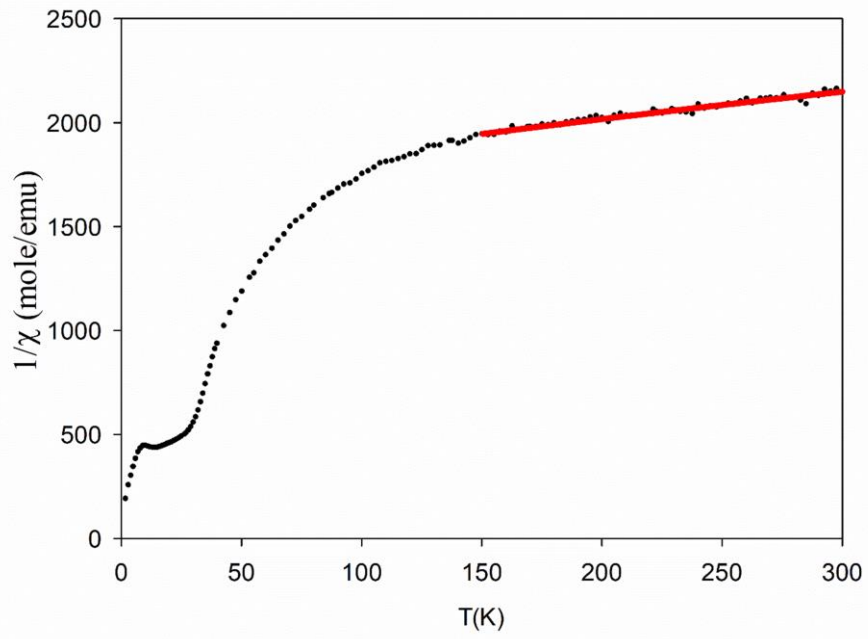
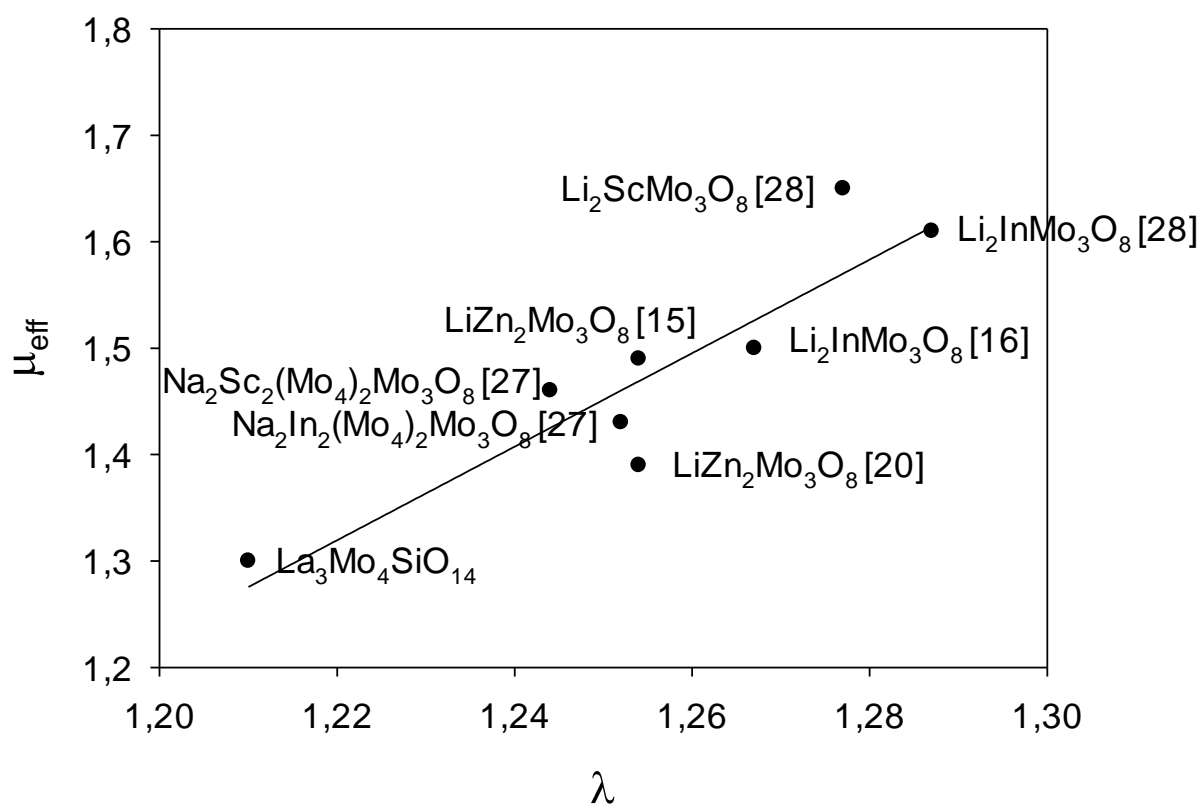


Figure 9

Figure 10



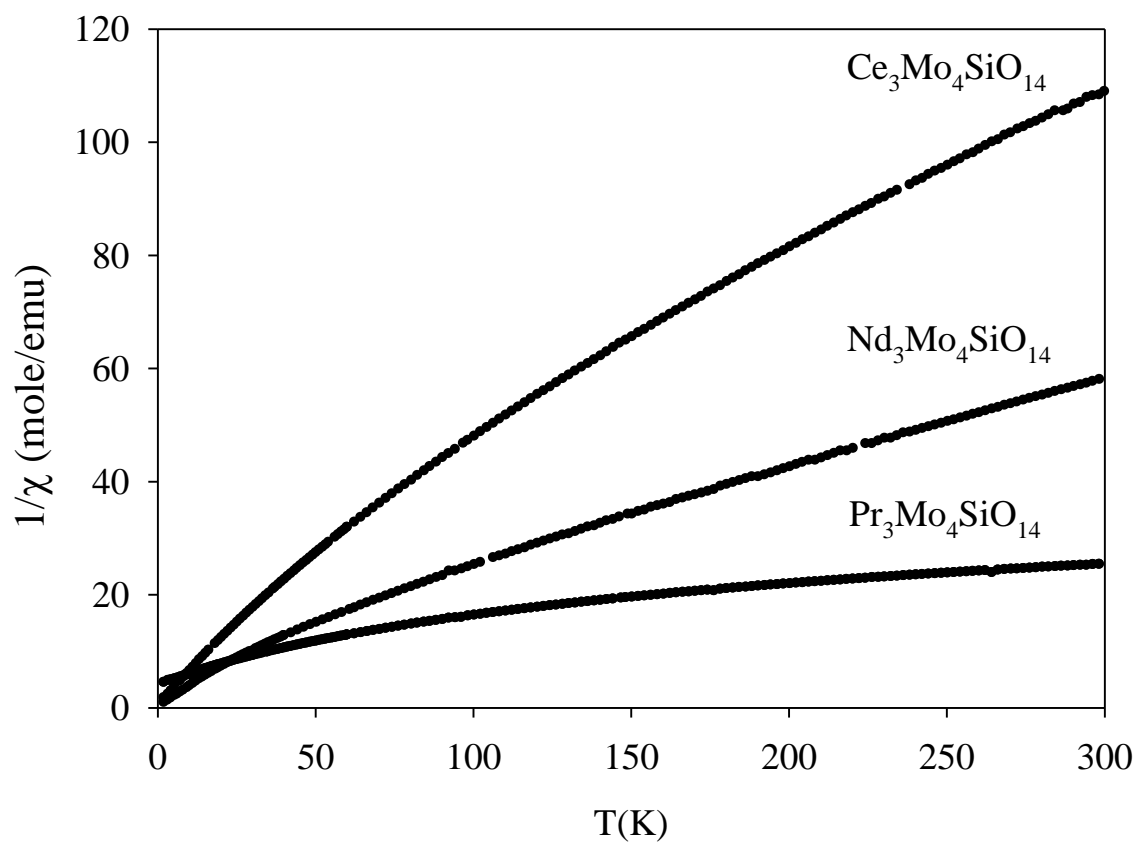


Figure 11