



**HAL**  
open science

# Singularities of fractional Emden's equations via Caffarelli-Silvestre extension

Huyuan Chen, Laurent Véron

► **To cite this version:**

Huyuan Chen, Laurent Véron. Singularities of fractional Emden's equations via Caffarelli-Silvestre extension. 2022. hal-03689999v1

**HAL Id: hal-03689999**

**<https://hal.science/hal-03689999v1>**

Preprint submitted on 7 Jun 2022 (v1), last revised 3 Mar 2023 (v2)

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Singularities of fractional Emden’s equations via Caffarelli-Silvestre extension

**Huyuan Chen**

Department of Mathematics, Jiangxi Normal University  
Nanchang, Jiangxi 330022, PR China

**Laurent Véron**

Institut Denis Poisson, CNRS UMR 7013  
Université de Tours, 37200 Tours, France

## Abstract

We study the isolated singularities of functions satisfying

$$(E) \quad (-\Delta)^s v \pm |v|^{p-1} v = 0 \quad \text{in } \Omega \setminus \{0\}, \quad v = 0 \quad \text{in } \mathbb{R}^N \setminus \Omega,$$

where  $0 < s < 1$ ,  $p > 1$  and  $\Omega$  is a bounded domain containing the origin. We use the Caffarelli-Silvestre extension to  $\mathbb{R}_+ \times \mathbb{R}^N$ . We emphasize the obtention of *a priori* estimates, analyse the set of self-similar solutions via energy methods to characterize the singularities.

*AMS Subject Classification:* 35B40-35J60-35J70-45G05

*Keywords:* Fractional Laplacian; Isolated singularity; Lane-Emden-Fowler equation; Linear extension; Energy methods; Limit set.

## Contents

<b>1</b>	<b>Introduction</b>	<b>2</b>
<b>2</b>	<b>The lifting</b>	<b>7</b>
2.1	Caffarelli-Silvestre extension and its Poisson’s kernel . . . . .	7
2.2	The trace theorem . . . . .	8
2.3	Spherical coordinates . . . . .	11
<b>3</b>	<b>Self-Similar solutions</b>	<b>12</b>
3.1	Fractional Emden-Fowler equation . . . . .	13
3.2	Fractional Lane-Emden equation . . . . .	20
<b>4</b>	<b>A priori estimates</b>	<b>23</b>
4.1	The lifting estimates . . . . .	23
4.2	Integral estimates . . . . .	24
4.3	A priori estimate for Lane-Emden equation . . . . .	25
4.4	Some maximum principles . . . . .	27
4.5	Fractional Hardy operator . . . . .	29
4.6	Emden-Fowler equation with Dirac mass . . . . .	32
4.7	A priori estimate for Emden-Fowler equation . . . . .	37

<b>5 Singularity of solutions</b>	<b>40</b>
5.1 Energy method and limit set . . . . .	40
5.2 Singularities of Emden-Fowler equation . . . . .	41
5.3 Singularities for Lane-Emden equation . . . . .	42
5.4 Classification of fractional Emden-Fowler equation . . . . .	47

## 1 Introduction

In the past forty years a large series of articles have been devoted to study of the singular behaviour of solutions for the following two classes of semilinear equations

$$-\Delta v + \epsilon v^p = 0 \quad \text{in } B_1 \setminus \{0\}, \quad (1.1)$$

where  $p > 1$ ,  $\epsilon = \pm 1$  and  $v^p = |v|^{p-1}v$ . The first empirical studies of radial solutions to the Lane-Emden equation ( $\epsilon = -1$ ) are due to J. Lane and R. Emden when analyzing the structure of polytropic objects submitted to their gravitational fields. A comprehensive presentation can be found in Chandrasekhar's book [11, pp 84-182]. The so-called Emden-Fowler equation ( $\epsilon = 1$ ), was treated in details by R. Fowler in the radial case. For some specific values of  $p$  ( $\frac{3}{2}$  and  $\frac{5}{2}$ ) this equation appeared in the study of the density of the electronic cloud in the Thomas-Fermi-Dirac model of atoms (see e.g. [34]).

The study of non-radial solutions started in the heighties in connection with conformal deformation of metric to metric with constant scalar curvature ( $K = 1$  or  $K = -1$  and  $p = \frac{N+2}{N-2}$ ). The main results concerning the Lane-Emden equation are due to Lions [30] in the case  $1 < p < \frac{N}{N-2}$  who proved that any positive solution of

$$-\Delta v - v^p = 0 \quad \text{in } B_1 \setminus \{0\} \quad (1.2)$$

satisfies

$$v(x) = k|x|^{2-N} + O(1) \quad \text{as } x \rightarrow 0 \quad (1.3)$$

when  $N \geq 3$ , (with standard modification if  $N = 2$ ) and the following holds

$$-\Delta v - v^p = c_N k \delta_0 \in \mathcal{D}'(B_1).$$

In the case  $\frac{N}{N-2} < p < \frac{N+2}{N-2}$ , Gidas and Spruck [24] introduced new methods and proved that any positive solutions of (1.2) satisfies

$$v(x) \leq c_p |x|^{-\frac{2}{p-1}} \quad \text{for } x \in B_{\frac{1}{2}} \setminus \{0\}.$$

As a consequence, a positive solution is either regular at  $x = 0$  or has the following behaviour

$$v(x) = \alpha_{N,p} |x|^{-\frac{2}{p-1}} (1 + o(1)) \quad \text{as } x \rightarrow 0. \quad (1.4)$$

When  $p = \frac{N+2}{N-2}$  the energy method used in [24] failed because of the conformal invariance of the equation. In [9] a new method of asymptotic symmetry allowed to treat this case showing that the positive solutions behaved like the radial ones, which in turn were easy to describe because

of this invariance. Finally, when  $p = \frac{N}{N-2}$  the result of [24] was sharpened by Aviles [2] who proved that either the positive solutions are smooth or they satisfy

$$\lim_{x \rightarrow 0} |x|^{N-2} (-\ln |x|)^{\frac{N-2}{2}} v(x) = \left( \frac{N-2}{\sqrt{2}} \right)^{N-2}.$$

The non-radial study of the Emden-Fowler equations started with the works of Brezis and Lieb who obtained in [5] the universal *a priori* estimate (also called Keller-Osserman estimate), valid for signed solutions of

$$-\Delta v + v^p = 0 \quad \text{in } B_1 \setminus \{0\}, \quad (1.5)$$

and is expressed as follows:

$$|v(x)| \leq c'_p |x|^{-\frac{2}{p-1}} \quad \text{for any } x \in B_{\frac{1}{2}} \setminus \{0\}.$$

With this estimate, Brezis and Véron proved in [6] that if  $p \geq \frac{N}{N-2}$  any solution is at least  $C^2$  in  $B_1$ . When  $1 < p < \frac{N}{N-2}$  the precise description of singular solutions was obtained by Véron who obtained that any positive solution  $v$  is either smooth, either satisfies (1.3) or (1.4) (with a new constant positive  $\tilde{\alpha}_{N,p}$ ), all these behaviours being effective. Furthermore when  $\frac{N+1}{N-1} \leq p < \frac{N}{N-2}$  the assumption of positivity could be removed, the sign of  $k$  not rescribed in (1.3) and (1.4) replaced by there exists  $\ell \in \mathbb{R}$  such that

$$\lim_{x \rightarrow 0} |x|^{\frac{2}{p-1}} v(x) = \ell \in \{\tilde{\alpha}_{N,p}, -\tilde{\alpha}_{N,p}, 0\} \quad \text{as } x \rightarrow 0.$$

In the case  $N = 2$ , the requirement  $q \geq \frac{N+1}{N-1}$  (i.e.  $q \geq 3$ ) for signed solutions was removed by Chen, Matano and Véron [14] who proved that for any solution  $v$  of (1.5) there exists a  $2\pi$ -periodic solution of

$$-\omega_{\theta\theta} - \left( \frac{2}{p-1} \right)^2 \omega + \omega^p = 0$$

such that

$$\lim_{x \rightarrow 0} |x|^{\frac{2}{p-1}} v(x) = \lim_{r \rightarrow 0} r^{\frac{2}{p-1}} v(r, \theta) = \omega(\theta) \quad \text{as } r = |x| \rightarrow 0.$$

This panorama of the classical Emden equations intended to show what could be expected when the ordinary Laplacian is replaced by the fractional Laplacian in these equations. If  $s \in (0, 1)$  and  $(-\Delta)^s$  the fractional Laplacian in  $\mathbb{R}^N \setminus \{0\}$  is defined on functions  $u \in C^2(\mathbb{R}^N \setminus \{0\}) \cap L^1_{\mu_s}(\mathbb{R}^N)$  with  $\mu_s(x) = (1 + |x|)^{-N-2s}$  by the expression

$$(-\Delta)^s u(x) = c_{N,s} \lim_{\varepsilon \rightarrow 0} \int_{|x-y|>\varepsilon} \frac{u(x) - u(y)}{|x-y|^{N+2s}} dy \quad \text{for } x \in \mathbb{R}^N \setminus \{0\}, \quad (1.6)$$

where  $c_{N,s} = 2^{2s} \pi^{-\frac{N}{2}} \frac{\Gamma(\frac{N+2s}{2})}{\Gamma(1-s)}$ . The singularity problem for the fractional Emden equations in a punctured domain  $\Omega$  containing  $\bar{B}_1$  is

$$\begin{aligned} (-\Delta)^s v + \varepsilon v^p &= 0 \quad \text{in } \Omega \setminus \{0\} \\ v &= 0 \quad \text{in } \Omega^c, \end{aligned} \quad (1.7)$$

where  $\epsilon = \pm 1$  and  $\Omega^c = \mathbb{R}^N \setminus \Omega$ . Some results concerning this equation in the case  $\epsilon = -1$  (the fractional Lane-Emden equation) have already been obtained in particular in the case  $p = \frac{N+2s}{N-2s}$ , see [10]. Therein it is proved that either the solution is smooth at 0, or there exists some constant  $c \geq 1$  such that

$$\frac{1}{c}|x|^{-\frac{N-2s}{2}} \leq v(x) \leq c|x|^{-\frac{N-2s}{2}} \quad \text{for all } x \in B_{\frac{1}{2}} \setminus \{0\}.$$

Furthermore, there exists a radial solution  $\tilde{v}$  of (1.7) such that

$$v(x) = \tilde{v}(|x|)(1 + o(1)) \quad \text{as } x \rightarrow 0.$$

In the general case  $\frac{N}{N-2s} < p < \frac{N+2s}{N-2s}$ , Yang and Zou obtained in [36] the analogous of Gidas and Spruck estimates, namely any positive solution of

$$\begin{aligned} (-\Delta)^s v - v^p &= 0 & \text{in } \Omega \setminus \{0\} \\ v &= 0 & \text{in } \Omega^c, \end{aligned} \tag{1.8}$$

either is smooth or satisfies for some constant  $c > 0$

$$\frac{1}{c}|x|^{-\frac{2s}{p-1}} \leq v(x) \leq c|x|^{-\frac{2s}{p-1}} \quad \text{for all } x \in B_{\frac{1}{2}} \setminus \{0\}.$$

Concerning the fractional Emden-Fowler equation

$$\begin{aligned} (-\Delta)^s v + v^p &= 0 & \text{in } \Omega \setminus \{0\} \\ v &= 0 & \text{in } \Omega^c \end{aligned} \tag{1.9}$$

not so many results are known. The problem

$$\begin{aligned} (-\Delta)^s v + v^p &= \mu & \text{in } \mathcal{D}'(\Omega) \\ v &= 0 & \text{in } \Omega^c, \end{aligned}$$

where  $\mu$  is a positive bounded Radon measure in  $\Omega$ , is studied in [15] in a more general framework, that is by replacing  $u^p$  by a nondecreasing nonlinearity  $g(u)$ . The particular case where  $\mu = k\delta_0$  and  $g(u) = u^p$  is analyzed in [16]. In particular, the asymptotics of the solutions  $u_k$  (they are unique) when  $k \rightarrow \infty$  are thoroughly studied in [16].

The aim of this article is to present a unified analysis of the isolated singularities for problem (1.7). This analysis is based upon the Caffarelli-Silvestre lifting of the equation which associates to it the following degenerate elliptic equation in  $\mathbb{R}_+^{N+1} = \{\xi = (x, z) : x \in \mathbb{R}^N, z > 0\}$ ,

$$\begin{aligned} \operatorname{div}(z^{1-2s}\nabla u) &= 0 & \text{in } \mathbb{R}_+^{N+1} \\ u(\cdot, 0) &= v & \text{in } \mathbb{R}^N \sim \partial\mathbb{R}_+^{N+1}, \end{aligned} \tag{1.10}$$

with the property that

$$(-\Delta)^s v(x) = -\lim_{z \rightarrow 0} z^{1-2s} u_z(x, z) := \partial_{\nu^s} u(x, 0).$$

The study of (1.7) is replaced by

$$\begin{aligned} \operatorname{div}(z^{1-2s}\nabla u) &= 0 && \text{in } \mathbb{R}_+^{N+1} \\ \partial_{\nu^s} u(\cdot, 0) + \epsilon u(\cdot, 0)^p &= 0 && \text{in } \Omega \setminus \{0\} \\ u(\cdot, 0) &= 0 && \text{in } \mathbb{R}^N \setminus \Omega \end{aligned} \quad (1.11)$$

and

$$v = u(\cdot, 0) \quad \text{in } \mathbb{R}^N.$$

We prove that any positive solution of (1.10) admits a trace on  $\mathbb{R}^N$  which is a nonnegative Radon measure  $\mu$  satisfying that

$$\int_{\mathbb{R}^N} \frac{d\mu}{(1+|x|)^{N+2s}} < +\infty.$$

The first step for describing the behaviour of solutions of (1.11) near 0 is to obtain an *a priori* estimate of  $u(x, z)$ .

In the case  $\epsilon = 1$ , we use a blow-up technique combined with the trace theorem to prove that when  $p \geq 1 + \frac{2s}{N}$ , any positive solution  $u(x, z)$  of (1.11) satisfies

$$u(x, z) \leq c\rho^{-\frac{2s}{p-1}} \quad \text{for all } (x, z) \in \tilde{\mathbf{B}}_{\frac{1}{2}}^+ \setminus \{0\} \quad (1.12)$$

where  $\rho = \sqrt{|x|^2 + z^2}$  and  $\tilde{\mathbf{B}}_a^+ := \{(x, z) \in \overline{\mathbb{R}_+^{N+1}} : \rho < a\}$ . This estimate implies that any positive solution of (1.9) verifies

$$v(x) \leq c|x|^{-\frac{2s}{p-1}} \quad \text{for all } x \in \tilde{B}_{\frac{1}{2}}^+ \setminus \{0\}. \quad (1.13)$$

Note that this estimate has no interest if  $p \geq \frac{N}{N-2s}$  since the solution  $v$  is smooth.

In the case  $\epsilon = -1$ , we use the estimate (1.13) proved [36, Proposition 3.1] in the case  $1 < p < \frac{N+2s}{N-2s}$  combined with the expression of the Poisson kernel of the operator  $u \mapsto \operatorname{div}(z^{1-2s}\nabla u)$  in  $\mathbb{R}_+^{N+1}$  obtained by [8] to prove that positive solutions of (1.11) satisfies also (1.12). Furthermore, when  $v$  is a radially symmetric decreasing solution of (1.8), we prove that (1.13) holds for any  $p > 1$ .

The second step in our study is to consider the self-similar solutions of (1.11) in  $\overline{\mathbb{R}_+^{N+1}} \setminus \{0\}$ , with  $\Omega = \mathbb{R}^N$ . In spherical coordinates in  $\mathbb{R}_+^{N+1}$ , these solutions have the following expression

$$u(x, z) = u(\rho, \sigma) = \rho^{-\frac{2s}{p-1}} \omega(\sigma) \quad \text{for all } (\rho, \sigma) \in \mathbb{R}_+ \times \mathbb{S}_+^N$$

and, up to a rotation and a good choice of spherical variables on  $\mathbb{S}^N$ ,  $\omega$  satisfies

$$\begin{aligned} \mathcal{A}_s[\omega] + \Lambda_{s,p,N}\omega &= 0 && \text{in } \mathbb{S}_+^N \\ \frac{\partial \omega}{\partial \nu^s} + \epsilon|\omega|^{p-1}\omega &= 0 && \text{in } \mathbb{S}^{N-1}, \end{aligned} \quad (1.14)$$

Therein,  $\mathcal{A}_s$  is a degenerate elliptic operator on the  $N$ -sphere  $\mathbb{S}^N$ ,  $\frac{\partial}{\partial \nu^s}$  the corresponding conormal outward derivative on  $\partial\mathbb{S}_+^N = \mathbb{S}^{N-1}$  and

$$\Lambda_{s,p,N} = \frac{2s}{p-1} \left( \frac{2s}{p-1} + 2s - N \right).$$

The sign of  $\Lambda_{s,p,N}$  which is fundamental in the study of (1.14) depends on the value of  $p$  with respect to  $\frac{2s}{p-1}$ . The structure of the set  $\mathcal{E}_\epsilon$  (resp.  $\mathcal{E}_\epsilon^+$ ) of solutions (resp. positive solutions) plays a key role in our study.

**Theorem A** *Let  $s \in (0, 1)$ ,  $\epsilon = 1$  and  $p > 1$ .*

- 1- *If  $p \geq \frac{N}{N-2s}$ , then  $\mathcal{E}_1 = \{0\}$ .*
- 2- *If  $1 < p \leq 1 + \frac{2s}{N}$ , then  $\mathcal{E}_1^+ = \{0\}$ .*
- 3- *If  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$ , then  $\mathcal{E}_1^+ = \{0, \omega_1\}$ , where  $\omega_1$  is a positive solution of (1.14).*

Besides the case  $p \geq \frac{N}{N-2s}$  we can describe the set  $\mathcal{E}_1$  in another case.

**Theorem B** *There exists  $p^* = p^*(s, N) \in (1, \frac{N}{N-2s})$  such that if  $p^* \leq p < \frac{N}{N-2s}$ , then  $\mathcal{E}_1 = \{0, \omega_1, -\omega_1\}$ .*

The exponent  $p^*$  is explicitly given by (3.8) below and its algebraic expression is heavy. However if  $N \geq 3$  or  $N = 2$  and  $\frac{1}{2} \leq s < 1$  one has  $p^* \leq 1 + \frac{2s}{N}$ . The set of positive self-similar solutions of the Lane-Emden is characterized by the following statement.

**Theorem C** *Let  $s \in (0, 1)$ ,  $\epsilon = -1$  and  $p > 1$ .*

- 1- *If  $p \leq \frac{N}{N-2s}$ , then  $\mathcal{E}_{-1}^+ = \{0\}$ .*
- 2- *If  $p > \frac{N}{N-2s}$ , then  $\mathcal{E}_{-1}^+ = \{0, \omega_2\}$ , where  $\omega_2$  is a positive solution of (1.14) depending only on one variable.*

We use the classical  $\ln \rho$  variable in order to transform (1.14) into an autonomous equation, writing any  $u$  under the form

$$u(x, z) = \rho^{-\frac{2s}{p-1}} w(t, \sigma) \quad \text{with } t = \ln \rho. \quad (1.15)$$

We introduce standard tools of the dynamical systems theory such as energy functional and limit sets and prove the following general theorem.

**Theorem D** *Assume  $s \in (0, 1)$ ,  $\epsilon = \pm 1$ ,  $p \in (1, +\infty) \setminus \{\frac{N+2s}{N-2s}\}$ . Let  $u$  be a solution of (1.11) satisfying (1.12), even with  $u$  replaced by  $|u|$  and  $w$  be defined by (1.15). Then the limit set at  $-\infty$  of the negative trajectory  $\mathcal{T}_-[w] := \bigcup \{w(t, \cdot) : t < 0\}$  of  $w$  is a nonempty, compact and connected subset of  $\mathcal{E}_\epsilon^+$  in the  $C^{2s}(\overline{\mathbb{S}_+^N})$ -topology.*

Thanks to this result and the properties of  $\mathcal{E}_\epsilon^+$ , we can give the precise behaviour of positive solutions of (1.7) near the origin. This happens when  $\mathcal{E}_\epsilon$  is disconnected, the limit set is reduced to a single element usually  $\omega_j$  ( $j=1,2$ ) or 0. When this limit set is zero, the solutions have either a weak singularity or regular, according to the value of  $\epsilon$ .

**Theorem E** Let  $s \in (0, 1)$ ,  $\epsilon = 1$  and  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$ . If  $u$  is a positive solution of (1.11) satisfying

$$\lim_{(x,z) \rightarrow 0} \rho^{\frac{2s}{p-1}} u(x, z) = 0, \quad (1.16)$$

then there exists  $k > 0$  such that

$$u(x, z) = C_{N,s} k \frac{z^{2s}}{(|x|^2 + z^2)^{\frac{N}{2}+s}} + O(1) \quad \text{as } x \rightarrow 0,$$

and  $v_k = u(\cdot, 0)$  satisfies

$$(-\Delta)^s v + v^p = k \delta_0 \quad \text{in } \mathcal{D}'(\Omega), \quad (1.17)$$

where  $C_{N,s} = \frac{1}{4} \pi^{\frac{N+2-2s}{2}} \Gamma(\frac{N-2s}{2})$ .

Furthermore, if  $k = 0$ ,  $v_0$  is identically zero and  $u$  is smooth.

**Theorem F** Let  $s \in (0, 1)$ ,  $\epsilon = -1$  and  $p \in (\frac{N}{N-2s}, +\infty) \setminus \{\frac{N+2s}{N-2s}\}$ . If  $u$  is a positive solution of (1.11) satisfying (1.16), then  $u$  is smooth in  $\mathbb{R}_+^{N+1}$ .

Finally, we can give precise classification of nonnegative solutions to the fractional Emden-Fowler equation.

**Theorem G** Let  $s \in (0, 1)$ ,  $p > 1$  and  $v_k$  be the solution of (1.17) for  $k \geq 0$  vanishing in  $\Omega^c$ .

- 1- For  $p \in (1, 1 + \frac{2s}{N}]$ , the set of nonnegative solutions of (1.9) is the set of  $\{v_k\}_{k \geq 0}$ .
- 2- For  $p \in (1 + \frac{2s}{N}, \frac{N}{N-2s})$ , the set of nonnegative solutions of (1.9) is the set of  $\{v_k\}_{k \geq 0} \cup \{v_\infty\}$ , where  $v_\infty = \lim_{k \rightarrow +\infty} v_k$ .
- 3- For  $p \geq \frac{N}{N-2s}$ , the set of nonnegative solutions to (1.9) is reduced  $\{0\}$ .

The remainder of this paper is organized as follows. In Section 2, we provide preliminary tools: Caffarelli-Silvestre extension, the related Poisson kernel and basic setting. In Section 3, we classify singular self-similar solutions of (1.7). Section 4 is devoted to obtain some priori estimates the upper bounds of singular solutions (1.7) near the origin. Finally, we prove the classification of isolated singularities of (1.7).

## 2 The lifting

### 2.1 Caffarelli-Silvestre extension and its Poisson's kernel

Caffarelli and Silvestre [8] introduced a general method of lifting from  $\mathbb{R}^N$  to  $\mathbb{R}^N \times (0, +\infty)$  allowing to define the fractional Laplacian as the weighted trace of the normal derivative. Setting

$$\mathbb{R}_+^{N+1} = \{\xi = (x, z) : x \in \mathbb{R}^N, z > 0\},$$

they proved that if  $v \in C^2(\mathbb{R}^N) \cap L_{\mu_s}^1(\mathbb{R}^N)$  and  $u \in C^2(\mathbb{R}_+^{N+1})$  satisfy

$$\begin{aligned} \mathcal{D}_s u &:= \operatorname{div}(z^{1-2s} \nabla u) = 0 && \text{in } \mathbb{R}_+^{N+1} \\ u(\cdot, 0) &= v && \text{in } \mathbb{R}^N, \end{aligned} \quad (2.1)$$



then

$$(-\Delta)^s v(x) = -\lim_{z \rightarrow 0} z^{1-2s} u_z(x, z).$$

In the sequel we denote

$$\partial_{\nu^s} u(x, 0) := -\lim_{z \rightarrow 0} z^{1-2s} u_z(x, z).$$

Set

$$\mu_s(x) = (1 + |x|)^{-N-2s} \quad \text{for } x \in \mathbb{R}^N. \quad (2.2)$$

The Poisson's kernel for the operator  $u \mapsto \mathcal{D}_s u$  has been obtained in [8, (2.1)]

$$\mathcal{P}_s(x, z) = C_{N,s} \frac{z^{2s}}{(|x|^2 + z^2)^{\frac{N}{2}+s}}, \quad (2.3)$$

where  $C_{N,s} = \frac{1}{4} \pi^{\frac{N+2-2s}{2}} \Gamma(\frac{N-2s}{2})$ . If  $v \in L^1_{\mu_s}(\mathbb{R}^N)$ , then

$$u(x, z) := \mathbb{P}_s[v](x, z) = C_{N,s} z^{2s} \int_{\mathbb{R}^N} \frac{v(y) dy}{(|x-y|^2 + z^2)^{\frac{N}{2}+s}} \quad (2.4)$$

satisfies (2.1).

## 2.2 The trace theorem

The counter part of Caffarelli-Silvestre theorem is the following.

**Theorem 2.1** *Let  $s \in (0, 1)$  and  $u \in L^1_{loc}(\mathbb{R}_+^{N+1})$  satisfies  $u \geq 0$  and*

$$\mathcal{D}_s u = 0 \quad \text{in } \mathbb{R}_+^{N+1}. \quad (2.5)$$

*Then  $u \in C^\infty(\mathbb{R}_+^{N+1})$  and  $u$  admits a boundary trace on  $\partial\mathbb{R}_+^{N+1}$ , which is a nonnegative Radon measure  $\mu$  in the sense that for any  $\zeta \in C_0^\infty(\mathbb{R}^N)$ ,*

$$\lim_{z \rightarrow 0} \int_{\mathbb{R}^N} u(x, z) \zeta(x) dx = \int_{\mathbb{R}^N} \zeta d\mu(x),$$

*and  $\mu$  satisfies*

$$\int_{\mathbb{R}^N} \frac{d\mu(x)}{(|x|+1)^{N+2s}} < +\infty. \quad (2.6)$$

*Furthermore, if  $\|u(z, \cdot)\|_{L^1(\mathbb{R}^N)} = o(z^{2s})$  when  $z \rightarrow \infty$ , then*

$$u(x, z) = C_{N,s} z^{2s} \int_{\mathbb{R}^N} \frac{d\mu(y)}{(|x-y|^2 + z^2)^{\frac{N}{2}+s}} \quad \text{for all } (x, z) \in \mathbb{R}_+^{N+1}. \quad (2.7)$$

**Proof.** The fact that  $u \in C^\infty(\mathbb{R}_+^{N+1})$  follows from the standard theory of elliptic operators. Furthermore, the equation (2.5) not only holds in the sense of distributions but also in the strong sense. The core of the proof is an adaptation of the Brezis-Lions wellknown note on isolated singularities of linear elliptic equations [7].

Let  $R > 0$ ,  $\lambda_{R,1} > 0$  be first eigenvalue of  $-\Delta$  in  $B_R \subset \mathbb{R}^N$  under the zero Dirichlet boundary condition and  $\phi_{R,1}$  be the related first eigenfunction. We set

$$X(z) = \int_{B_R} u(x, z) \phi_{R,1}(x) dx.$$

Then  $X$  satisfies the ODE

$$(z^{1-2s} X')'(z) - \lambda_{R,1} z^{1-2s} X(z) \leq 0 \quad \text{for } z \in (0, \infty).$$

Integrating the equation on  $(z, T)$ , we obtain

$$z^{1-2s} X'(z) \geq T^{1-2s} X'(T) - \lambda_{R,1} \int_z^T t^{1-2s} X(t) dt = T^{1-2s} X'(T) - \lambda_{R,1} F(z),$$

where we have denoted

$$F(z) = \int_z^T t^{1-2s} X(t) dt.$$

The function  $F$  is decreasing and

$$\begin{aligned} X(z) &\leq X(T) - \frac{1}{2s} (T - T^{1-2s} z^{2s}) X'(T) + \lambda_{R,1} \int_z^T t^{2s-1} F(t) dt \\ &\leq X(T) - \frac{1}{2s} (T - T^{1-2s} z^{2s}) X'(T) + \frac{\lambda_{R,1}}{2s} (T^{2s} - z^{2s}) F(z). \end{aligned} \quad (2.8)$$

Hence  $F'(z) = -z^{1-2s} X(z)$  and (2.8) becomes

$$\begin{aligned} -F'(z) &\leq z^{1-2s} X(T) - \frac{1}{2s} (T z^{1-2s} - T^{1-2s} z) X'(T) + \frac{\lambda_{R,1}}{2s} (T^{2s} z^{1-2s} - z) F(z) \\ &=: A(z) + B(z) F(z) \end{aligned} \quad (2.9)$$

for  $z \in (0, \infty)$ . Since  $A, B \in L^1(0, T)$  and  $F(T) = 0$ , we deduce by integration,

$$F(z) \leq e^{-\int_z^T B(t) dt} \int_z^T e^{\int_t^T B(\tau) d\tau} A(t) dt. \quad (2.10)$$

Therefore  $F$  is uniformly bounded on  $(0, T)$ , and from (2.8) the function  $X$  shares the same property. This implies that  $u \in L^1_{loc}(\mathbb{R}^N \times [0, +\infty))$ .

Next, let  $\zeta \in C_0^\infty(\mathbb{R}^N)$  with support in  $B_R$  for some  $R$ . We set

$$Y(z) = \int_{\mathbb{R}^N} u(x, z) \zeta(x) dx \quad \text{and} \quad \Phi(z) = \int_{\mathbb{R}^N} u(x, z) \Delta_x \zeta(x) dx.$$

Then

$$(z^{1-2s} Y'(z))' + z^{1-2s} \Phi(z) = 0.$$

Consequently, we have that

$$Y(z) = Y(1) - \frac{1 - z^{2s}}{2s} Y'(1) - \frac{1}{2s} \int_z^1 (t^{2s} - z^{2s}) \Phi(t) dt.$$

Since

$$\int_z^1 |(t^{2s} - z^{2s})\Phi(t)|dt \leq \|\Delta\zeta\|_{L^\infty} \int_0^1 \int_{B_R} u(x,t)dxdt < +\infty,$$

it follows that  $Y(z)$  admits a limit when  $z \rightarrow 0$ , which defines a positive linear functional on  $C_0^\infty(\mathbb{R}^N)$ , hence a Radon measure denoted by  $\mu$ .

Let  $R > 0$  and  $\varepsilon > 0$ . For  $k > 0$  we set

$$Q_{R,R+k,\varepsilon}^T = B_{R+k} \times (\varepsilon, T),$$

and we consider the problem

$$\begin{aligned} \mathcal{D}_s w &= 0 && \text{in } Q_{R,R+k,\varepsilon}^T \\ w &= 0 && \text{in } \partial B_{R+k} \times [\varepsilon, T] \cup \overline{B}_{R+k} \times \{T\} \\ w(\cdot, \varepsilon) &= u(\cdot, \varepsilon)\chi_{B_R} && \text{in } B_{R+k} \times \{\varepsilon\}. \end{aligned}$$

The solution  $w = w_{\varepsilon,R,k,T}$  is unique and satisfies

$$w_{\varepsilon,R,k,T} \leq u \quad \text{in } Q_{R,R+k,\varepsilon}^T.$$

The correspondence  $(R, k, T) \mapsto w_{\varepsilon,R,k,T}$  is increasing and we have that

$$w_{\varepsilon,R} := \lim_{k,T \rightarrow \infty} w_{\varepsilon,R,k,T} \leq u \quad \text{in } \mathbb{R}^N \times (\varepsilon, +\infty).$$

When  $\varepsilon \rightarrow 0$ ,  $w_{\varepsilon,R}(\cdot, \varepsilon)$  converges to  $\chi_{B_R}\mu$ . Hence its limit  $w_R$  is expressed by the Poisson kernel (2.3), then

$$w_R(x, z) = C_{N,s} z^{2s} \int_{\mathbb{R}^N} \frac{\chi_{B_R}(y) d\mu(y)}{(|x-y|^2 + z^2)^{\frac{N}{2}+s}} \quad \text{in } \mathbb{R}^N \times (0, +\infty).$$

As a consequence, letting  $R \rightarrow +\infty$ , we get

$$u(0, 1) \geq C_{N,s} \int_{B_R} \frac{d\mu(y)}{(|y|^2 + 1)^{\frac{N}{2}+s}},$$

which implies (2.6).

Finally, since  $R \mapsto w_R$  is increasing and  $w_R$  is upper bounded by  $u$ , we obtain that  $w_R$  admits a limit  $w$  when  $R \rightarrow +\infty$ , which is a positive solution of (2.5) in  $\mathbb{R}_+^{N+1}$  and satisfies

$$w(x, z) = C_{N,s} z^{2s} \int_{\mathbb{R}^N} \frac{d\mu(y)}{(|x-y|^2 + z^2)^{\frac{N}{2}+s}} \leq u(x, z) \quad \text{for } (x, z) \in \mathbb{R}^N \times (0, +\infty). \quad (2.11)$$

Thus  $w$  is the minimal positive solution of (2.5) in  $\mathbb{R}_+^{N+1}$  with trace  $\mu$  on  $\partial\mathbb{R}_+^{N+1}$ . Set  $\psi = u - w$ , then  $\psi$  is a nonnegative solution of (2.5) in  $\mathbb{R}_+^{N+1}$  with trace 0 on  $\partial\mathbb{R}_+^{N+1}$ . The even extension of  $\psi$  to whole  $\mathbb{R}^{N+1}$  defined by [8, Lemma 4.1] is a nonnegative solution of (2.5) in  $\mathbb{R}^{N+1}$ . It is therefore continuous and it satisfies Harnack inequality [26]. For proving uniqueness, we set  $t = (\frac{z}{2s})^{2s}$  and  $\tilde{\psi}(x, t) = \psi(x, z)$ . Then  $\tilde{\psi}$  satisfies (see[8, 1.8])

$$t^{\frac{2s-1}{s}} \tilde{\psi}_{tt} + \Delta_x \tilde{\psi} = 0.$$

Since the operator  $-\Delta_x$  is accretive in  $L^1(\mathbb{R}^N)$ , the function  $t \mapsto \left\| \tilde{\psi}(\cdot, t) \right\|_{L^1(\mathbb{R}^N)}$  is convex on  $(0, \infty)$ . If  $\left\| \tilde{\psi}(\cdot, 0) \right\|_{L^1(\mathbb{R}^N)} = 0$  and  $\left\| \tilde{\psi}(\cdot, t) \right\|_{L^1(\mathbb{R}^N)} = o(t)$  when  $t \rightarrow \infty$ ,  $\left\| \tilde{\psi}(\cdot, t) \right\|_{L^1(\mathbb{R}^N)} = 0$  for all  $t > 0$ . The condition  $\left\| \tilde{\psi}(\cdot, t) \right\|_{L^1(\mathbb{R}^N)} = o(t)$  is equivalent to  $\|\psi(\cdot, z)\|_{L^1(\mathbb{R}^N)} = o(z^{2s})$ .

Therefore, the uniqueness and (2.11) implies (2.7), which ends the proof.  $\square$

### 2.3 Spherical coordinates

We recall spherical coordinates  $(r, \sigma) \in (0, +\infty) \times \mathbb{S}^N$ . The parametrization of  $\mathbb{R}_+^{N+1}$  is the following

$$\mathbb{R}_+^{N+1} = \{\xi = (r, \sigma) : r > 0, \sigma \in \mathbb{S}_+^N\},$$

where

$$\mathbb{S}_+^N = \left\{ \sigma = (x, z) = (\sigma' \cos \phi, \sin \phi) : \sigma' \in \mathbb{S}^{N-1}, \phi \in \left[0, \frac{\pi}{2}\right] \right\}.$$

Without confusion, we also use the next notation

$$\mathbb{S}^{N-1} = \partial \mathbb{S}_+^N = \{(0, \sigma') \in \mathbb{R}^{N+1} : |\sigma'| = 1\}.$$

With this parametrization

$$\Delta_{\mathbb{S}^N} u = u_{\phi\phi} - (N-1)(\tan \phi) u_\phi + \frac{1}{\cos^2 \phi} \Delta_{\mathbb{S}^{N-1}} u.$$

We define the operator  $\mathcal{A}_s$  in  $C^2(\mathbb{S}_+^N)$  by

$$\mathcal{A}_s[w] = \frac{1}{\lambda_s(\phi)(\cos \phi)^{N-1}} (\lambda_s(\phi)(\cos \phi)^{N-1} w_\phi)_\phi + \frac{1}{\cos^2 \phi} \Delta_{\mathbb{S}^{N-1}} w,$$

where

$$\lambda_s(\phi) = (\sin \phi)^{1-2s}.$$

We denote by  $dS$  the invariant measure obtained by the isometric imbedding of  $\mathbb{S}^N$  into  $\mathbb{R}^{N+1}$ , then

$$dS(\sigma) = (\cos \phi)^{N-1} dS'(\sigma') d\phi,$$

and by  $\nabla'$  the covariant gradient in the canonical metric on  $\mathbb{S}^N$  identified with the tangential gradient on the unit sphere of  $\mathbb{R}^{N+1}$ . The bilinear form associated to  $\mathcal{A}_s$  is

$$\mathcal{B}[w, \tilde{w}] = \int_{\mathbb{S}_+^N} \left( w_\phi \tilde{w}_\phi + \frac{1}{\cos^2 \phi} \nabla' w \cdot \nabla' \tilde{w} \right) \lambda_s(\phi) dS = - \int_{\mathbb{S}_+^N} \tilde{w} \mathcal{A}_s(w) \lambda_s(\phi) dS,$$

and we have the corresponding Green's formula

$$\int_{\mathbb{S}_+^N} \tilde{w} \mathcal{A}_s(w) \lambda_s dS = \int_{\mathbb{S}^{N-1}} \tilde{w} \partial_{\phi^s} w dS' - \mathcal{B}[w, \tilde{w}].$$

If  $u = \mathbb{P}_s[v]$ , then

$$(-\Delta)^s v(r, \sigma') = - \lim_{\phi \rightarrow 0} (\sin \phi)^{1-2s} u_\phi(r, \sigma', \phi) := \partial_{\phi^s} u(r, \sigma', 0).$$

### 3 Self-Similar solutions

Let  $\Omega \subset \mathbb{R}^N$  be a bounded domain containing the origin. A function  $v \in C^2(\Omega \setminus \{0\}) \cap L^1_{\mu_s}(\mathbb{R}^N)$  satisfies the fractional Emden equations in  $\Omega \setminus \{0\}$  with zero exterior value if

$$\begin{aligned} (-\Delta)^s v + \epsilon v^p &= 0 \quad \text{in } \Omega \setminus \{0\} \\ v &= 0 \quad \text{in } \Omega^c, \end{aligned} \tag{3.1}$$

where  $p > 1$  and  $\epsilon = \pm 1$ . If  $\epsilon = 1$  the equation considered is called *the fractional Emden-Fowler equation* while if  $\epsilon = -1$  it is called *the fractional Lane-Emden equation*. We denote  $u = \mathbb{P}_s[v]$ , which is admissible since  $v \in L^1(\mathbb{R}^N)$ , then

$$\begin{aligned} \mathcal{D}_s u &= 0 \quad \text{in } \mathbb{R}_+^{N+1} \\ \partial_{\nu^s} u(\cdot, 0) + \epsilon u(\cdot, 0)^p &= 0 \quad \text{in } \Omega \setminus \{0\} \\ u(\cdot, 0) &= 0 \quad \text{in } \Omega^c. \end{aligned} \tag{3.2}$$

Hence (3.2) implies that

$$\begin{aligned} u_{rr} + \frac{N+1-2s}{r} u_r + \frac{1}{r^2} \mathcal{A}_s[u] &= 0 \quad \text{in } (0, \infty) \times \mathbb{S}_+^N, \\ - \lim_{\phi \rightarrow 0} (\sin \phi)^{1-2s} u_\phi(r, \theta, \phi) + \epsilon u(r, \theta, 0)^p &= 0 \quad \text{in } (0, \infty) \times \mathbb{S}^{N-1}. \end{aligned} \tag{3.3}$$

Equation (3.2) is equivariant under the family of transformations  $\mathcal{S}_\ell$ ,  $\ell > 0$ , defined by

$$\mathcal{S}_\ell[u](x, z) = \ell^{\frac{2s}{p-1}} u(\ell x, \ell z) = \ell^{\frac{2s}{p-1}} u(\ell \xi).$$

Therefore self-similar solutions of (3.2) have the form

$$u(r, \sigma) = r^{-\frac{2s}{p-1}} \omega(\sigma) \quad \text{with } (r, \sigma) \in (0, \infty) \times \mathbb{S}_+^N,$$

and  $\omega$  satisfies

$$\begin{aligned} \mathcal{A}_s[\omega] + \Lambda_{s,p,N} \omega &= 0 \quad \text{in } \mathbb{S}_+^N \\ \frac{\partial \omega}{\partial \nu^s} + \epsilon \omega^p &= 0 \quad \text{in } \mathbb{S}^{N-1}, \end{aligned} \tag{3.4}$$

where

$$\Lambda_{s,p,N} = \frac{2s}{p-1} \left( \frac{2s}{p-1} + 2s - N \right).$$

We denote by  $\mathcal{E}_\epsilon$  (resp.  $\mathcal{E}_\epsilon^+$ ) the set of functions (resp. positive functions)  $\omega \in C(\overline{\mathbb{S}_+^N}) \cap C^2(\mathbb{S}_+^N)$  satisfying (3.4). There are several critical exponents will play some role later on:

$$\begin{aligned} (i) \quad 1 < p \leq \frac{N}{N-2s} &\iff \Lambda_{s,p,N} \geq 0; \\ (ii) \quad p \geq 1 + \frac{2s}{N} &\iff \Lambda_{s,p,N} \leq 2sN; \\ (iii) \quad p = \frac{N+2s}{N-2s} &\iff \Theta_{s,p,N} := N - 2s \frac{p+1}{p-1} = 0. \end{aligned} \tag{3.5}$$

### 3.1 Fractional Emden-Fowler equation

The structure of  $\mathcal{E}_1$  is described as follows.

**Theorem 3.1** *Assume  $s \in (0, 1)$ ,  $\epsilon = 1$  and  $p > 1$ .*

1- *If  $p \geq \frac{N}{N-2s}$ , then  $\mathcal{E}_1 = \{0\}$ .*

2- *If  $1 < p \leq 1 + \frac{2s}{N}$ , then  $\mathcal{E}_1^+ = \{0\}$ .*

3- *If  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$ , then  $\mathcal{E}_1^+ = \{0, \omega_1\}$ , where  $\omega_1$  is a positive solution of (3.4) depending only on the variable  $\phi$ .*

**Proof.** 1- Let  $\omega$  be a solution, then

$$\int_{\mathbb{S}_+^N} \left( \omega_\phi^2 + \frac{1}{\cos^2 \phi} |\nabla' w|^2 - \Lambda_{s,p,N} \omega^2 \right) \lambda_s(\phi) dS = \int_{\mathbb{S}^{N-1}} \omega \partial_{\phi^s} \omega dS' = - \int_{\mathbb{S}^{N-1}} |\omega|^{p+1} dS'.$$

If  $p \geq \frac{N}{N-2s}$ , then  $\Lambda_{N,s,p} \leq 0$ . This implies that  $\omega = 0$ .

2- It is easy to check that the first eigenfunction  $\psi_1$  of  $\mathcal{A}_s$  in  $W_0^{1,2}(\mathbb{S}_+^N)$  is  $\phi \mapsto \psi_1(\theta, \phi) = (\sin \phi)^{2s}$  with corresponding eigenvalue  $\ell_1 = 2sN$ . Assume now that  $\omega$  is a positive solution of (3.4), then, multiplying the equation by  $\psi_1$  and integrating yields

$$(\Lambda_{s,p,N} - 2sN) \int_{\mathbb{S}_+^N} \omega \psi_1 \lambda_s(\phi) dS = \int_{\mathbb{S}^{N-1}} \frac{\partial \psi_1}{\partial \nu^s} \omega dS' = -2s \int_{\mathbb{S}^{N-1}} \omega dS'.$$

The claim follows by (3.5)-(ii).

3- *Existence.* We denote by  $W(\mathbb{S}_+^N)$  the space of functions  $w$  such that  $\mathcal{B}[w, w] < +\infty$ , where

$$\mathcal{B}[w, v] = \int_{\mathbb{S}_+^N} \left( w_\phi v_\phi + \frac{1}{\cos^2 \phi} \nabla' w \cdot \nabla' v \right) \lambda_s(\phi) dS.$$

Then  $w \mapsto \sqrt{\mathcal{B}[w, w]}$  is a semi-norm on  $W(\mathbb{S}_+^N)$  and it is a norm on the subspace  $C_c^\infty(\mathbb{S}_+^N)$ . Indeed

$$\mathcal{B}[w, w] \geq 2sN \int_{\mathbb{S}_+^N} w^2 \lambda_s(\phi) dS \quad \text{for all } w \in C_c^\infty(\mathbb{S}_+^N).$$

The closure of  $C_c^\infty(\mathbb{S}_+^N)$  in  $W(\mathbb{S}_+^N)$  is denoted by  $W_0(\mathbb{S}_+^N)$ . We define the functional

$$J(w) = \frac{1}{2} \int_{\mathbb{S}_+^N} \left( w_\phi^2 + \frac{1}{\cos^2 \phi} |\nabla' w|^2 - \Lambda_{s,p,N} w^2 \right) \lambda_s(\phi) dS + \frac{1}{p+1} \int_{\mathbb{S}^{N-1}} |\gamma_0(w)|^{p+1} dS',$$

where  $\gamma_0$  denotes the trace operator from  $W(\mathbb{S}_+^N)$  onto  $L^2(\mathbb{S}^{N-1})$  identified with  $w(0, \theta)$ . Then  $J$  is a quadratic perturbation of a convex lower semicontinuous functional defined in  $W(\mathbb{S}_+^N)$ , hence it is weakly lower continuous with domain of definition

$$D(J) = \{w \in W(\mathbb{S}_+^N) : \gamma_0(w) \in L^{p+1}(\mathbb{S}^{N-1})\}.$$

Let  $w \in W(\mathbb{S}_+^N)$  and

$$w = w_1 + w_2,$$

where  $w_1 \in W_0(\mathbb{S}_+^N)$  is defined by

$$\mathcal{B}[w_1, \zeta] = \mathcal{B}[w, \zeta] \quad \text{for all } \zeta \in W_0(\mathbb{S}_+^N).$$

Since  $\mathcal{B}$  is a continuous quadratic form in  $W(\mathbb{S}_+^N)$ ,  $\zeta \mapsto \mathcal{B}[w, \zeta]$  is a continuous linear form in  $W(\mathbb{S}_+^N)$  and there holds

$$\mathcal{B}[\zeta, \zeta] \geq (2sN - \Lambda_{s,p,N}) \int_{\mathbb{S}_+^N} \zeta^2 \lambda_s(\phi) dS \quad \text{for all } \zeta \in W_0(\mathbb{S}_+^N),$$

then the function  $w_1$  is uniquely defined. Thus, one has

$$J(w) = J(w_1 + w_2) = \frac{1}{2} \mathcal{B}[w_1, w_1] + \frac{1}{2} \mathcal{B}[w_2, w_2] + 2\mathcal{B}[w_1, w_2] + \frac{1}{p+1} \int_{\mathbb{S}^{N-1}} |\gamma_0(w)|^{p+1} dS'.$$

We have  $\mathcal{B}[w_1, w_1] = \mathcal{B}[w, w_1] = \mathcal{B}[w_1, w_1] + \mathcal{B}[w_2, w_1]$ , therefore

$$\mathcal{B}[w_2, w_1] = \mathcal{B}[w_1, w_2] = 0.$$

Furthermore  $w_2(\phi, \theta) = w_2(\phi, \theta) - w_2(0, \theta) + w_2(0, \theta)$ . Since for any  $\varepsilon > 0$ ,

$$w_2^2(\phi, \theta) \leq (1 + \varepsilon) (w_2(\phi, \theta) - w_2(0, \theta))^2 + \left(1 + \frac{1}{\varepsilon}\right) w_2^2(0, \theta),$$

and  $(\phi, \theta) \mapsto w_2(\phi, \theta) - w_2(0, \theta)$  belongs to  $W_0^{1,2}(\mathbb{S}_+^N)$ , we have the inequality

$$\begin{aligned} \mathcal{B}[w_2, w_2] &\geq \int_{\mathbb{S}_+^N} \frac{1}{\cos^2 \phi} |\nabla w_2|^2 \lambda_s(\phi) dS + \int_{\mathbb{S}_+^N} \left( (w_2(\phi, \theta) - w_2(0, \theta))_\phi \right)^2 \lambda_s(\phi) dS \\ &\quad - (1 + \varepsilon) \Lambda_{s,p,N} \int_{\mathbb{S}_+^N} (w_2(\phi, \theta) - w_2(0, \theta))^2 \lambda_s(\phi) dS \\ &\quad - \left(1 + \frac{1}{\varepsilon}\right) \Lambda_{s,p,N} \int_{\mathbb{S}_+^N} w_2^2(0, \theta) \lambda_s(\phi) dS. \end{aligned}$$

We choose  $\varepsilon$  such that

$$(1 + \varepsilon) \Lambda_{s,p,N} < 2sN$$

and we get

$$\begin{aligned} \mathcal{B}[w_2, w_2] &\geq \left(1 - \frac{(1 + \varepsilon) \Lambda_{s,p,N}}{2sN}\right) \int_{\mathbb{S}_+^N} \left( w_2^2_\phi + \frac{1}{\cos^2 \phi} |\nabla' w_2|^2 \right) \lambda_s(\phi) dS \\ &\quad - \left(1 + \frac{1}{\varepsilon}\right) \Lambda_{s,p,N} \int_{\mathbb{S}_+^N} w_2^2(0, \theta) \lambda_s(\phi) dS. \end{aligned}$$

Notice that

$$\begin{aligned} \int_{\mathbb{S}_+^N} w_2^2(0, \theta) \lambda_s dS &= \int_0^{\frac{\pi}{2}} (\sin \phi)^{1-2s} (\cos \phi)^{N-1} d\phi \int_{\mathbb{S}^{N-1}} w_2^2(0, \theta) dS' \\ &\leq c_1 \left( \int_{\mathbb{S}^{N-1}} |w_2(0, \theta)|^{p+1} dS' \right)^{\frac{2}{p+1}}, \end{aligned}$$

where  $c_1 = c_1(s, p, N) > 0$ . Because  $w_1 \in W_0^{1,2}(\mathbb{S}_+^N)$ , there holds also

$$\mathcal{B}[w_1, w_1] \geq \left( 1 - \frac{\Lambda_{s,p,N}}{2sN} \right) \int_{\mathbb{S}_+^N} \left( w_1^2_\phi + \frac{1}{\cos^2 \phi} |\nabla' w_1|^2 \right) \lambda_s(\phi) dS.$$

Combining the previous inequalities, we finally obtain

$$\begin{aligned} J(w) &\geq \left( 1 - \frac{(1+\varepsilon)\Lambda_{s,p,N}}{2sN} \right) \int_{\mathbb{S}_+^N} \left( w_\phi^2 + \frac{1}{\cos^2 \phi} |\nabla' w|^2 \right) \lambda_s(\phi) dS \\ &\quad - \left( 1 + \frac{1}{\varepsilon} \right) M \Lambda_{s,p,N} \left( \int_{\mathbb{S}^{N-1}} |\gamma_0 w|^{p+1} dS' \right)^{\frac{2}{p+1}} + \frac{2}{p+1} \int_{\mathbb{S}^{N-1}} |\gamma_0 w|^{p+1} dS'. \end{aligned}$$

This implies that  $J(w)$  tends to  $\infty$  when  $\min \left\{ \sqrt{\mathcal{B}[w, w]}, \|\gamma_0(w)\|_{L^{p+1}(\mathbb{S}^{N-1})} \right\}$  tends to infinity. Consequently the functional  $J$  admits a minimum  $\omega$  in  $W(\mathbb{S}_+^N)$ . Moreover, since  $J(|\omega|) = J(\omega)$  the minimum is achieved by a nonnegative function and it is therefore classical that  $\omega$  is a solution of (3.4), and it is positive by the strong maximum principle.

Finally, if we consider the restriction  $J_{rad}$  of  $J$  to the space  $W_{rad}(\mathbb{S}_+^N)$  of function of  $W(\mathbb{S}_+^N)$  depending only on the variable  $\phi$  (they are called radial), then  $\gamma_0(w^{p+1})$  is a real number and  $J_{rad}(w)$  has the following form provided we write  $w = w(\phi)$ ,

$$J_{rad}(w) = \frac{|\mathbb{S}^{N-1}|}{2} \int_0^{\frac{\pi}{2}} (w_\phi^2 - \Lambda_{s,p,N} w^2) \lambda_s(\phi) (\cos \phi)^{N-1} d\phi + \frac{|\mathbb{S}^{N-1}|}{p+1} |w(0)|^{p+1}.$$

The functional  $J_{rad}$  is also a quadratic perturbation of a convex lower semicontinuous functional defined in  $W_{rad}(\mathbb{S}_+^N)$  and it tends to  $\infty$  at  $\infty$ . Thus it admits a minimum which is achieved by a nonnegative function due to  $J_{rad}(w) = J_{rad}(|w|)$ . Hence there exists a minimizing nonnegative solution of (3.4) which depends only on the variable  $\phi$ .

*3- Uniqueness.* Let  $\omega$  and  $\tilde{\omega}$  be two positive solutions of (3.4), then

$$\begin{aligned} 0 &= \int_{\mathbb{S}_+^N} \left( \frac{\mathcal{A}_s[\omega]}{\omega} - \frac{\mathcal{A}_s[\tilde{\omega}]}{\tilde{\omega}} \right) (\omega^2 - \tilde{\omega}^2) \lambda_s dS \\ &= \int_{\mathbb{S}^{N-1}} \left( \partial_{\phi^s} \omega \left( \omega - \frac{\tilde{\omega}^2}{\omega} \right) - \partial_{\phi^s} \tilde{\omega} \left( \tilde{\omega} - \frac{\omega^2}{\tilde{\omega}} \right) \right) dS' - \left( \mathcal{B} \left[ \omega, \omega - \frac{\tilde{\omega}^2}{\omega} \right] - \mathcal{B} \left[ \tilde{\omega}, \tilde{\omega} - \frac{\omega^2}{\tilde{\omega}} \right] \right) \\ &= A - B. \end{aligned}$$



Furthermore

$$\begin{aligned} A &= \int_{\mathbb{S}^{N-1}} \left( -\omega^p \left( \omega - \frac{\tilde{\omega}^2}{\omega} \right) + \tilde{\omega}^p \left( \tilde{\omega} - \frac{\omega^2}{\tilde{\omega}} \right) \right) dS' \\ &= - \int_{\mathbb{S}^{N-1}} (\omega^2 - \tilde{\omega}^2) (\omega^{p-1} - \tilde{\omega}^{p-1}) dS' \leq 0 \end{aligned}$$

and

$$B = \int_{\mathbb{S}^{N-1}} \left( \frac{1}{\omega^2} + \frac{1}{\tilde{\omega}^2} \right) \left( (\omega \tilde{\omega}_\phi - \tilde{\omega} \omega_\phi)^2 + \frac{1}{\cos^2 \phi} |\omega \nabla' \tilde{\omega} - \tilde{\omega} \nabla' \omega|^2 \right) \lambda_s(\phi) dS' \geq 0.$$

Therefore,  $A = B = 0$  and  $\omega = \tilde{\omega}$ .  $\square$

**Remark.** To the unique positive element  $\omega_1$  of  $\mathcal{E}_1^+$  corresponds a unique positive singular self-similar solution  $U_p$  of

$$(-\Delta)^s v + v^p = 0 \quad \text{in } \mathbb{R}^N \setminus \{0\},$$

where  $U_p(x) = c_p |x|^{-\frac{2s}{p-1}}$  with

$$c_p = \omega_1(1) = \left( -2^{2s} \frac{\Gamma(\frac{N}{2} - \frac{s}{p-1}) \Gamma(s + \frac{s}{p-1})}{\Gamma(\frac{s}{p-1}) \Gamma(\frac{N}{2} - s - \frac{s}{p-1})} \right)^{\frac{1}{p-1}}.$$

In order to prove Theorem 3.4, we need the following intermediate result.

**Proposition 3.2** *Assume  $\Lambda \neq 0$ , then for any  $a \neq 0$ , there exists a unique function  $\omega_a$  satisfying*

$$\omega_a(\phi) = a - \Lambda \int_\phi^{\frac{\pi}{2}} (\sin \sigma)^{2s-1} (\cos \sigma)^{1-N} \int_\sigma^{\frac{\pi}{2}} \omega_a(\theta) (\sin \theta)^{1-2s} (\cos \theta)^{N-1} d\theta d\sigma \quad \text{in } (0, \frac{\pi}{2}). \quad (3.6)$$

Furthermore  $a = \omega_a(\frac{\pi}{2})$  and  $\omega_a = a\omega_1$ . If  $a > 0$ , then  $\omega_a$  is positive and increasing (resp. decreasing) on  $(0, \frac{\pi}{2})$  if  $\Lambda < 0$  (resp.  $\Lambda > 0$ ).

**Proof.** Let  $a > 0$  and  $\tilde{\omega}_a(\phi) = \omega_a(\frac{\pi}{2} - \phi)$ . To find a function  $\omega_a$  satisfying (3.6) is equivalent to find  $\tilde{\omega}_a$  satisfying

$$\tilde{\omega}_a(\phi) = a - \Lambda \int_0^\phi (\cos \sigma)^{2s-1} (\sin \sigma)^{1-N} \int_0^\sigma \tilde{\omega}_a(\theta) (\cos \theta)^{1-2s} (\sin \theta)^{N-1} d\theta d\sigma \quad \text{in } (0, \frac{\pi}{2}).$$

We define the operator  $\mathcal{T}$  on  $C([0, \frac{\pi}{2}])$  by

$$\mathcal{T}[w](\phi) = a - \Lambda \int_0^\phi (\cos \sigma)^{2s-1} (\sin \sigma)^{1-N} \int_0^\sigma w(\theta) (\cos \theta)^{1-2s} (\sin \theta)^{N-1} d\theta d\sigma \quad \text{in } (0, \frac{\pi}{2}).$$

We have that

$$\begin{aligned} &|\mathcal{T}[w - w'](\phi)| \\ &\leq |\Lambda| \left( \int_0^\phi (\cos \sigma)^{2s-1} (\sin \sigma)^{1-N} \int_0^\sigma (\cos \theta)^{1-2s} (\sin \theta)^{N-1} d\theta d\sigma \right) \sup_{[0, \phi]} |(w - w')(\theta)|. \end{aligned}$$

1- We first assume that  $1 - 2s \geq 0$ , then

$$\begin{aligned} \int_0^\phi (\cos \sigma)^{2s-1} (\sin \sigma)^{1-N} \int_0^\sigma (\cos \theta)^{1-2s} (\sin \theta)^{N-1} d\theta d\sigma &\leq \int_0^\phi (\cos \sigma)^{2s-1} \sigma d\sigma \\ &\leq \frac{\pi}{2} \int_0^\phi (\cos \sigma)^{2s-1} \sin \sigma d\sigma \leq c_2, \end{aligned}$$

where  $c_2 = \frac{\pi}{4s}$ , we used here the monotonicity  $\theta \in (0, \frac{\pi}{2}) \mapsto (\sin \theta)^{N-1}$  and  $(\cos \theta)^{1-2s} \leq 1$ . Hence

$$|\mathcal{T}[w - w'](\phi)| \leq c_2 \sup_{[0, \phi]} |(w - w')(\theta)|.$$

It is therefore straightforward to check by induction that for  $k \in \mathbb{N}^*$ ,

$$|\mathcal{T}^k[w - w'](\phi)| \leq \frac{c_2^k}{k!} \sup_{[0, \phi]} |(w - w')(\theta)|, \quad (3.7)$$

then  $\mathcal{T}^k$  admits a unique fixed point  $\tilde{\omega}_a$ , and uniqueness implies that  $\tilde{\omega}_a$  is also a fixed point of  $\mathcal{T}$  and it is unique.

2- Assume that  $1 - 2s < 0$ . Then

$$\begin{aligned} \int_0^\phi (\cos \sigma)^{2s-1} (\sin \sigma)^{1-N} \int_0^\sigma (\cos \theta)^{1-2s} (\sin \theta)^{N-1} d\theta d\sigma \\ \leq \int_0^\phi \int_0^\sigma d\theta d\sigma = \frac{\phi^2}{2} \leq \frac{\pi^2}{8} \end{aligned}$$

by there the monotonicity  $\theta \in (0, \frac{\pi}{2}) \mapsto (\cos \theta)^{1-2s} (\sin \theta)^{N-1}$ . Then (3.7) holds and the conclusion follows.

If  $\Lambda < 0$  (resp.  $\Lambda > 0$ ), the function  $\omega_a$  is increasing (resp. decreasing). The fact that  $\omega_a = a\omega_1$  is a consequence of the linearity of  $\mathcal{T}$  and uniqueness.  $\square$

The following critical exponent corresponds to the sequel of isotropy of solutions of (3.4) in the case  $\epsilon = 1$ ,

$$p^* := \frac{N + 2s + \sqrt{N^2 + 4(N - s) + 4s^2 - 4}}{N - 2s + \sqrt{N^2 + 4(N - s) + 4s^2 - 4}}. \quad (3.8)$$

This exponent corresponds to the fact that  $\frac{2s}{p^*-1}$  is the positive root of the equation

$$Q(X) := X^2 + (2s - N)X + 1 - N = 0. \quad (3.9)$$

**Theorem 3.3** *Assume  $s \in (0, 1)$ ,  $\epsilon = 1$  and  $p^* \leq p < \frac{N}{N-2s}$ . Then  $\mathcal{E}_1 = \{0, \omega_1, -\omega_1\}$ , where  $\omega_1$  is a positive solution of (3.4) depending only on the variable  $\phi$ .*

**Proof.** *Step 1.* We first prove that under the condition  $p \geq p^*$  any element of  $\mathcal{E}_1$  does not depend on the variable  $\sigma' \in \mathbb{S}^{N-1}$ . If  $\psi$  is any function defined on  $\mathbb{S}^N$ , we set

$$\bar{\psi}(\phi) = \frac{1}{|\mathbb{S}^{N-1}|} \int_{\mathbb{S}^{N-1}} \psi(\phi, \sigma') dS'.$$

By averaging (3.4) we obtained that  $\bar{\omega}$  satisfies

$$\begin{aligned}\mathcal{A}_s[\bar{\omega}] + \Lambda_{s,p,N}\bar{\omega} &= 0 && \text{in } \mathbb{S}_+^N \\ \frac{\partial \bar{\omega}}{\partial \nu^s} + \bar{\omega}^p &= 0 && \text{in } \mathbb{S}^{N-1},\end{aligned}$$

where

$$\mathcal{A}_s[\bar{\omega}] = \frac{1}{\lambda_s(\phi)(\cos \phi)^{N-1}} (\lambda_s(\phi)(\cos \phi)^{N-1}\bar{\omega})_\phi.$$

Multiplying the equation by  $(\omega - \bar{\omega})\lambda_s$  and integrating over  $\mathbb{S}_+^N$ , we obtain

$$\begin{aligned}\int_{\mathbb{S}^{N-1}} (\omega - \bar{\omega})(|\omega|^{p-1}\omega - |\bar{\omega}|^{p-1}\bar{\omega})dS' + \int_0^{\frac{\pi}{2}} \int_{\mathbb{S}^{N-1}} \left( (\omega_\phi - \bar{\omega}_\phi)^2 + \frac{1}{\cos^2 \phi} |\nabla'(\omega - \bar{\omega})|^2 \right) \lambda_s(\phi) dS \\ = \Lambda_{s,p,N} \int_{\mathbb{S}_+^N} (\omega - \bar{\omega})^2 \lambda_s dS.\end{aligned}\tag{3.10}$$

There holds

$$\begin{aligned}\int_{\mathbb{S}^{N-1}} (\omega - \bar{\omega})(\omega^p - \bar{\omega}^p) dS' &= \int_{\mathbb{S}^{N-1}} (\omega - \bar{\omega})(\omega^p - \bar{\omega}^p) dS' + (\bar{\omega}^p - \bar{\omega}^p) \int_{\mathbb{S}^{N-1}} (\omega - \bar{\omega}) dS' \\ &= \int_{\mathbb{S}^{N-1}} (\omega - \bar{\omega})(\omega^p - \bar{\omega}^p) dS' \\ &\geq 2^{-p} \int_{\mathbb{S}^{N-1}} |\omega - \bar{\omega}|^{p+1} dS'.\end{aligned}$$

Furthermore, we see that

$$\begin{aligned}\int_0^{\frac{\pi}{2}} \int_{\mathbb{S}^{N-1}} \left( (\omega_\phi - \bar{\omega}_\phi)^2 + \frac{1}{\cos^2 \phi} |\nabla'(\omega - \bar{\omega})|^2 \right) \lambda_s(\phi) dS \\ \geq \int_0^{\frac{\pi}{2}} \left( \int_{\mathbb{S}^{N-1}} |\nabla'(\omega - \bar{\omega})|^2 dS' \right) (\cos \phi)^{N-3} \lambda_s(\phi) d\phi\end{aligned}\tag{3.11}$$

and the inequality is strict unless  $\omega_\phi = \bar{\omega}_\phi$ . By Wirtinger's inequality, there holds for any  $\phi \in (0, \frac{\pi}{2})$ ,

$$\int_{\mathbb{S}^{N-1}} |\nabla'(\omega - \bar{\omega})|^2 dS' \geq (N-1) \int_{\mathbb{S}^{N-1}} (\omega - \bar{\omega})^2 dS'.$$

Since  $(\cos \phi)^{N-3} \geq (\cos \phi)^{N-1}$ , we obtain

$$\int_0^{\frac{\pi}{2}} \int_{\mathbb{S}^{N-1}} \left( (\omega_\phi - \bar{\omega}_\phi)^2 + \frac{1}{\cos^2 \phi} |\nabla'(\omega - \bar{\omega})|^2 \right) \lambda_s(\phi) dS \geq (N-1) \int_{\mathbb{S}_+^N} (\omega - \bar{\omega})^2 \lambda_s(\phi) dS.$$

Finally we have that

$$(\Lambda_{s,p,N} + 1 - N) \int_{\mathbb{S}_+^N} (\omega - \bar{\omega})^2 \lambda_s(\phi) dS \geq 2^{-p} \int_{\mathbb{S}^{N-1}} |\omega - \bar{\omega}|^{p+1} dS'.$$

Introducing the polynomial  $Q$  defined in (3.9) and setting  $X = \frac{2s}{p-1}$ , then (3.10) becomes

$$Q(X) \int_{\mathbb{S}_+^N} (\omega - \bar{\omega})^2 \lambda_s(\phi) dS \geq 2^{-p} \int_{\mathbb{S}^{N-1}} |\omega - \bar{\omega}|^{p+1} dS'.$$

Hence if  $X < X^*$ , where  $X^*$  is the positive root of  $Q$ , then  $\omega - \bar{\omega} = 0$  in  $\mathbb{S}_+^N$ . If  $X = X^*$  we deduce that  $\omega(0, \theta) = \bar{\omega}(0) = 0$ . Then we use that fact that the inequality (3.11) is strict unless  $\omega_\phi = \bar{\omega}_\phi$ . If this holds, combining with  $\omega(\theta) - \bar{\omega}(0) = 0$  we deduce that  $\omega = \bar{\omega}$  in  $\mathbb{S}_+^N$ .

*Step 2.* We prove that  $\mathcal{E}_1 = \{0, \omega_1, -\omega_1\}$ , where  $\omega_1$  is a positive solution of (3.4) with  $\epsilon = 1$ , depending only on  $\phi$ . For this sake, it suffices to prove uniqueness among the radial solutions. If  $\phi \mapsto \omega(\phi)$  is a solution of (3.4) with  $\epsilon = 1$ , we have

$$\begin{aligned} -\frac{d}{d\phi} \left( (\sin \phi)^{1-2s} (\cos \phi)^{N-1} \frac{d\omega}{d\phi} \right) &= \Lambda_{s,p,N} (\sin \phi)^{1-2s} (\cos \phi)^{N-1} \omega \quad \text{in } \left(0, \frac{\pi}{2}\right) \\ \frac{d\omega(0)}{d\phi^s} + \omega(0)^p &= 0, \end{aligned} \tag{3.12}$$

where

$$\frac{d\omega(0)}{d\phi^s} = -\lim_{\phi \rightarrow 0} (\sin \phi)^{1-2s} \frac{d\omega(\phi)}{d\phi}.$$

Since  $\omega'(\frac{\pi}{2}) = 0$ , by integrating (3.12) we obtain that

$$\frac{d\omega(\phi)}{d\phi} = \Lambda_{s,p,N} (\sin \phi)^{2s-1} (\cos \phi)^{1-N} \int_{\phi}^{\frac{\pi}{2}} \omega(\theta) (\sin \theta)^{1-2s} (\cos \theta)^{N-1} d\theta \tag{3.13}$$

on  $(0, \frac{\pi}{2})$ . Therefore for  $\phi \in (0, \frac{\pi}{2})$ ,

$$\omega(\phi) = \omega\left(\frac{\pi}{2}\right) - \Lambda_{s,p,N} \int_{\phi}^{\frac{\pi}{2}} (\sin \sigma)^{2s-1} (\cos \sigma)^{1-N} \int_{\sigma}^{\frac{\pi}{2}} \omega(\theta) (\sin \theta)^{1-2s} (\cos \theta)^{N-1} d\theta d\sigma. \tag{3.14}$$

From (3.13),

$$\frac{d\omega}{d\phi^s}(0) = -\Lambda_{s,p,N} \int_0^{\frac{\pi}{2}} \omega(\theta) (\sin \theta)^{1-2s} (\cos \theta)^{N-1} d\theta. \tag{3.15}$$

Note that for  $p < \frac{N}{N-2s}$ ,  $\Lambda_{s,p,N} > 0$ . As a parameter we take  $a = \omega(\frac{\pi}{2})$ , and by linearity  $\omega_a = a\omega_1$ . We define the mapping  $F^* : \mathbb{R}_+ \mapsto \mathbb{R}$  by

$$F^*(a) = \frac{d\omega_a}{d\phi^s}(0) + (\omega_a(0))^p = a \left( \frac{d\omega_1}{d\phi^s}(0) + a^{p-1} (\omega_1(0))^p \right).$$

If  $\frac{d\omega_1}{d\phi^s}(0) \geq 0$ , then for any  $a > 0$  we have  $\frac{1}{a} F^*(a) > 0$ . Since  $F^*$  is increasing, there is no solution of (3.12), then no radial solution of (3.4) with  $\epsilon = 1$ , which contradicts Theorem 3.1, which states that there exists a radial solution, even a positive one. Therefore  $\frac{d\omega_1}{d\phi^s}(0) < 0$ . The mapping  $a \mapsto \frac{1}{a} F^*(a)$  is increasing and negative at  $a = 0$ . Since  $F^*$  tends to infinity as  $a \rightarrow +\infty$ , it follows that there exists a unique  $a^* > 0$  such that  $F^*(a^*) = 0$ , equivalently a unique solution  $\omega := \omega_{a^*}$  of (3.12) verifying  $a^* = \omega_{a^*}(\frac{\pi}{2})$ .

We end the proof as follows. From *Step 1* the solutions  $\omega$  depends only on  $\phi$ . Such a solution satisfying  $\omega(\frac{\pi}{2}) > 0$  is unique, thus it coincides with the radial positive solution obtained by minimization among radial functions.  $\square$

### 3.2 Fractional Lane-Emden equation

**Theorem 3.4** *Assume  $s \in (0, 1)$ ,  $\epsilon = -1$  and  $p > 1$ .*

1- *If  $p \leq \frac{N}{N-2s}$ , then  $\mathcal{E}_{-1}^+ = \{0\}$ .*

2- *If  $p > \frac{N}{N-2s}$ , then  $\mathcal{E}_{-1}^+ = \{0, \omega_2\}$ , where  $\omega_2$  is a positive solution of (3.4) depending only on  $\phi$ .*

**Proof.** *Part 1.* Direct computation implies that

$$\Lambda_{s,p,N} \int_{\mathbb{S}_+^N} \omega \lambda_s dS + \int_{\mathbb{S}_+^N} \partial_{\phi^s} \omega dS' = \Lambda_{s,p,N} \int_{\mathbb{S}_+^N} \omega \lambda_s dS + \int_{\mathbb{S}_+^N} \omega^p dS' = 0.$$

Since  $1 < p \leq \frac{N}{N-2s}$ ,  $\Lambda_{s,p,N} \geq 0$  and the assertion 1 follows.

2- *Existence of a unique positive solution depending only on  $\phi$ .* If  $\omega$  is a solution of (3.4) with  $\epsilon = -1$  depending only on  $\phi$ , then

$$-\frac{d}{d\phi} \left( (\sin \phi)^{1-2s} (\cos \phi)^{N-1} \omega_\phi \right) = \Lambda_{s,p,N} (\sin \phi)^{1-2s} (\cos \phi)^{N-1} \omega \quad \text{in } \left(0, \frac{\pi}{2}\right) \quad (3.16)$$

$$\frac{d\omega(0)}{d\phi^s} = \omega(0)^p.$$

Since  $\omega'(\frac{\pi}{2}) = 0$ , the identities (3.13), (3.14), and (3.15) obtained by integrating (3.16) are still valid. Because  $p > \frac{N}{N-2s}$ , we have that  $\Lambda_{s,p,N} < 0$ , hence  $\frac{d\omega}{d\phi^s}(0) > 0$  from identity (3.15). By linearity,  $\omega_a = a\omega_1$ . We define the mapping  $F$  from  $[0, \infty)$  to  $(-\infty, \infty)$  by

$$F(a) = \frac{d\omega_a(0)}{d\phi^s} - \omega_a(0)^p = a \left( \frac{d\omega_1(0)}{d\phi^s} - a^{p-1} \omega_1(0)^p \right).$$

This implies that  $a \mapsto \frac{1}{a} F(a)$  is decreasing, with a positive limit at  $a = 0$ , and tends to  $-\infty$  when  $a \rightarrow +\infty$ . Therefore,  $F$  admits a unique positive zero point for some  $a_0$ , and the function  $\omega_{a_0}$  is the unique positive solution of (3.4) which depends only on  $\phi$ .

2- *Any positive solution depends only on  $\phi$ .*

*Symmetry in the case  $s = \frac{1}{2}$ .* In that case  $s = \frac{1}{2}$ , the problem (3.4) reduces to

$$\begin{aligned} \Delta_{\mathbb{S}^N} \omega + \Lambda_{\frac{1}{2},p,N} \omega &= 0 & \text{in } \mathbb{S}_+^N \\ \frac{\partial \omega}{\partial \nu} - \omega^p &= 0 & \text{in } \mathbb{S}^{N-1}, \end{aligned} \quad (3.17)$$

where  $\Delta_{\mathbb{S}^N}$  is the Laplace-Beltrami operator on  $\mathbb{S}^N$  and

$$\Lambda_{\frac{1}{2},p,N} = \frac{1}{p-1} \left( \frac{1}{p-1} + 1 - N \right).$$

We denote the coordinates in  $\mathbb{R}^N$  by  $(x_1, \dots, x_N)$  and by  $(x_1, \dots, x_N, z)$ ,  $z > 0$ , the coordinates in  $\mathbb{R}_+^{N+1}$ . If  $\phi_0 \in (0, \frac{\pi}{2})$ , let  $\mathbb{H}_{\phi_0}$  be the hyperplane of  $\mathbb{R}^{N+1}$  passing through the  $(N-1)$ -plane  $D_{x_1} := \{(0, x_2, \dots, x_N, 0) \in \mathbb{R}^{N+1} : x_j \in \mathbb{R}, j = 2, \dots, N\}$  with angle  $\phi_0$  with the hyperplane  $z = 0$ . We denote by  $\Gamma_{\phi_0}$  the domain of  $\mathbb{S}^N$  with boundary  $\mathbb{H}_{\phi_0} \cap \partial \mathbb{R}_+^{N+1} \cap \{(x_1, \dots, x_N, z) :$

$x_1 \geq 0$  and by  $\tilde{\Gamma}_{\phi_0}$  the symmetric of this domain with respect to the plane  $\mathbb{H}_{\phi_0}$ . If  $\omega$  is a positive solution of (3.4), we define the function

$$(\phi, \sigma') \mapsto \varpi(\phi, \sigma') := \omega(\phi, \sigma') - \omega(2\phi_0 - \phi, \sigma') \quad \text{for } (\phi, \sigma') \in (0, \phi_0) \times \mathbb{S}^{N-1}. \quad (3.18)$$

Since  $\phi$  can be larger than  $\frac{\pi}{2}$ , we make the convention that  $\omega(\phi, \sigma') = \omega(\pi - \phi, -\sigma')$  if  $\frac{\pi}{2} < \phi < \pi$ . The symmetry with respect to  $\mathbb{H}_{\phi_0}$  is an isometry of  $\mathbb{R}^{N+1}$ , hence  $\varpi$  satisfies

$$\begin{aligned} \Delta_{\mathbb{S}^N} \varpi + \Lambda_{\frac{1}{2}, p, N} \varpi &= 0 && \text{in } \Gamma_{\phi_0} \\ \varpi &= 0 && \text{in } \mathbb{H}_{\phi_0} \cap \mathbb{S}_+^N \\ \varpi(0, \cdot) &= \omega(0, \cdot) - \omega(2\phi_0, \cdot) && \text{in } \mathbb{S}^{N-1} \cap \{(x_1, x_2, \dots, x_N, 0) : x_1 > 0\}. \end{aligned}$$

Since  $\frac{\partial \omega}{\partial \nu^s} = \omega^p$  on  $\partial \mathbb{S}_+^N$ , the function  $\phi \mapsto \omega(\phi, \sigma')$  is strictly decreasing if  $(\phi, \sigma') \in (0, \phi_0) \times \mathbb{S}^{N-1}$  for  $\phi_0$  small enough. Hence  $\varpi(0, \sigma') > 0$  in  $\mathbb{S}^{N-1} \cap \{x = (x_1, x_2, \dots, x_N, 0) : x_1 > 0\}$ , and by the maximum principle we have that  $\varpi > 0$  in  $(0, \phi_0) \times \mathbb{S}^{N-1}$ . We proceed as in the Gidas-Nirenberg paper [23]: we denote by  $\phi^*$  the maximum of the  $\phi_0 \in (0, \frac{\pi}{2}]$  such that

$$\omega(\phi, \sigma') - \omega(2\phi_0 - \phi, \sigma') > 0 \quad \text{for } (\phi, \sigma') \in (0, \phi_0) \times \mathbb{S}^{N-1}. \quad (3.19)$$

We assume by contradiction that  $\phi^* < \frac{\pi}{2}$ . Then there exists  $\sigma^* \in \mathbb{S}^{N-1}$  such that  $v(\phi^*, \sigma^*) = 0$ . By the Hopf's Lemma,  $(\phi^*, \sigma^*)$  cannot belong to  $\mathbb{S}_+^{N-1}$ , hence it belongs to  $\overline{\mathbb{S}_+^N} \cap \{z = 0\}$ . Since Hopf's lemma is also valid at this point, we get a contradiction. Therefore  $\phi^* = \frac{\pi}{2}$ . As a consequence,

$$\omega(0, \sigma') - \omega(2\phi^* - \phi, \sigma') \geq 0 \quad \text{for all } \sigma' \in \mathbb{S}^{N-1} \cap \{(0, x_2, \dots, x_N, z) : z > 0\}.$$

Similarly, starting the reflexion from  $x_1 < 0$ , we get

$$\omega(0, \sigma') - \omega(2\phi^* - \phi, \sigma') \leq 0 \quad \text{for all } \sigma' \in \mathbb{S}^{N-1} \cap \{(0, x_2, \dots, x_N, z) : z > 0\}.$$

This implies that the gradient of  $\omega$  alongside the ‘‘great circle’’  $\mathbb{S}_+^N \cap \{(0, x_2, \dots, x_N, z) : z > 0\}$  is zero, which can be written as

$$\nabla' \omega \cdot \mathbf{e}_1 = 0 \quad \text{on } \mathbb{S}_+^N \cap \{(0, x_2, \dots, x_N, z) : z > 0\}.$$

Performing a rotation we conclude that  $\nabla'(\omega(\phi, \sigma')) = 0$  for all  $\sigma' \in \mathbb{S}^{N-1}$ . This implies that  $\omega$  depends only on  $\phi$ .

*Sketch of the proof for  $s \in (0, 1)$ .* We perform the same reflection method as in the case  $s = \frac{1}{2}$ , defining  $\varpi$  as in (3.18), it satisfies

$$\begin{aligned} \mathcal{A}_s[\varpi] + \Lambda_{s, p, N} \varpi &= 0 && \text{in } \Gamma_{\phi_0} \\ \varpi &= 0 && \text{in } \mathbb{H}_{\phi_0} \cap \mathbb{S}_+^N \\ \varpi(0, \cdot) &= \omega(0, \cdot) - \omega(2\phi_0, \cdot) && \text{in } \mathbb{S}^{N-1} \cap \{(x_1, x_2, \dots, x_N, 0) : x_1 > 0\}, \end{aligned}$$

where

$$\mathcal{A}_s[\varpi] = \Delta_{\mathbb{S}^N} \varpi + (1 - 2s) \cot(\phi) \varpi_\phi.$$

Since  $\frac{\partial \omega}{\partial \nu^s} = \omega^p$  on  $\partial \mathbb{S}_+^N$ , the function  $\phi \mapsto \omega(\phi, \sigma')$  should be decreasing if  $(\phi, \sigma') \in (0, \phi_0) \times \mathbb{S}^{N-1}$  for  $\phi_0$  small enough. Let  $\phi^*$  be the maximum of  $\phi_0 \in (0, \frac{\pi}{2}]$  such that (3.19) holds.

If  $\phi^* = \frac{\pi}{2}$ , we are done. If not, we can assume that  $\phi^* \in (0, \frac{\pi}{2})$  and let  $(\phi^*, \sigma^*)$  be the point verifying  $\varpi(\phi^*, \sigma^*) = 0$ . Note that  $(1 - 2s) \cot(\cdot)$  is always bounded in a neighborhood of  $(\phi^*, \sigma^*)$ . Thus, the Hopf's lemma could be applied and the left proof is the same as the case  $s = \frac{1}{2}$ .  $\square$

**Proofs of Theorem A–C.** Theorem A–C are exact Theorem 3.1, Theorem 3.3 and Theorem 3.4 respectively.  $\square$

**Remark.** From the unique positive solution  $\omega_2$  obtained in 2 in above theorem, the fractional Lane-Emden equation

$$(-\Delta)^s v - v^p = 0 \quad \text{in } \mathbb{R}^N \setminus \{0\}$$

has an explicit self-similar solution  $U_p$ . From (4.11) and (4.10) below (also see [19, Lemma 3.1]), we have that

$$U_p(x) = c_p |x|^{-\frac{2s}{p-1}}, \quad \text{where } c_p = \omega_2(1) = \left( 2^{2s} \frac{\Gamma(\frac{N}{2} - \frac{s}{p-1}) \Gamma(s + \frac{s}{p-1})}{\Gamma(\frac{s}{p-1}) \Gamma(\frac{N}{2} - s - \frac{s}{p-1})} \right)^{\frac{1}{p-1}}.$$

**Remark.** Another approach could be to transform the problem on the half-sphere  $\mathbb{S}_+^N$  into a problem on the unit ball  $B_1$  of  $\mathbb{R}^N$  using a stereographic projection from the antipodal point  $\mathbf{S}$  of the North pole  $\mathbf{N}$  on  $\mathbb{S}^N$ . This stereographic projection  $\mathbb{P}_N$  is an isomorphism from  $\mathbb{S}^N \setminus \{\mathbf{S}\}$  onto  $\mathbb{R}^N$ . Then the image of  $\mathbb{S}_+^N$  is the unit ball  $B_1$ . It is well known that  $\mathbb{P}_N$  is a conformal diffeomorphism with conformal factor  $D^{\frac{4}{N-2}}$ , where

$$D(x) = \left( \frac{2}{1 + |x|^2} \right)^{\frac{N-2}{2}}.$$

If  $\omega$  is a solution of (3.17), and if

$$\tilde{w}(x) = \omega(\sigma) D(x) \quad \text{with } \sigma = \mathbb{P}_N^{-1}(x).$$

Then the above boundary value follows from the fact that

$$\tan \phi = \frac{1 - |x|^2}{2|x|}$$

and if  $\omega$  is defined on  $\overline{\mathbb{S}_+^N}$ , we have that

$$\frac{\partial \omega}{\partial \nu} \Big|_{\phi=0} = \frac{\partial \omega}{\partial \nu} \Big|_{|x|=1} + \frac{N-2}{2} w \Big|_{|x|=1}.$$

As a consequence,  $\tilde{w}$  satisfies

$$\begin{aligned} \Delta \tilde{w} + \frac{4}{(1 + |x|^2)^2} \left( \Lambda_{\frac{1}{2}, p, N} + \frac{N(N-2)}{4} \right) \tilde{w} &= 0 \quad \text{in } B_1 \\ \frac{\partial \tilde{w}}{\partial \nu} - \tilde{w}^p + \frac{N-2}{2} \tilde{w} &= 0 \quad \text{on } \partial B_1. \end{aligned}$$

Strong Maximum Principle and Hopf's Lemma could imply that  $\tilde{w} > \left(\frac{N-2}{2}\right)^{\frac{1}{p-1}}$  and the standard method of moving planes could be applied to obtain the radial symmetry.

## 4 A priori estimates

### 4.1 The lifting estimates

Our aim in this sections is to obtain the upper bounds for the positive solutions of

$$\begin{aligned} (-\Delta)^s v + \epsilon v^p &= 0 & \text{in } \Omega \setminus \{0\} \\ v &= 0 & \text{in } \Omega^c, \end{aligned} \quad (4.1)$$

under the form

$$|v(x)| \leq c_3 |x|^{-\frac{2s}{p-1}} \quad \text{for all } x \in \Omega \setminus \{0\}, \quad (4.2)$$

where  $\Omega \subset \mathbb{R}^N$  is a bounded regular domain containing the origin.

In order to study the asymptotics of singular solutions near the origin, we use the Caffarelli-Silvestre lifting of  $v$  and the Poisson kernel associated to the operator  $\mathcal{D}_s$  by formula (2.3), (2.4). Set  $\mu_s(x) = (|x| + 1)^{-N-2s}$ . If  $v \in L^1_{\mu_s}(\mathbb{R}^N) := L^1(\mathbb{R}^N, \mu_s dx)$ , then, with  $\xi = (x, z)$ ,

$$u(\xi) = u(x, z) = \int_{\mathbb{R}^N} \mathcal{P}_s(x - y, z) v(y) dy \quad (4.3)$$

satisfies

$$\mathcal{D}_s u = 0 \quad \text{in } \mathbb{R}_+^{N+1}, \quad u(\cdot, 0) = v \quad \text{on } \mathbb{R}^N.$$

**Proposition 4.1** *Assume that  $p > 1 + \frac{2s}{N}$  and  $u$  is a positive solution of  $\mathcal{D}_s u = 0$  in  $\mathbb{R}_+^{N+1}$  such that  $v(x) = u(x, 0)$  satisfies estimate (4.2). Then there exists  $c_4 > 0$  such that*

$$u(x, z) \leq c_4 (|x|^2 + z^2)^{-\frac{s}{p-1}} \quad \text{for all } (x, z) \in \mathbb{R}^N \times \mathbb{R}_+. \quad (4.4)$$

**Proof.** Assume that  $u(\cdot, 0)$  is nonnegative and satisfies (4.2). If  $p > 1 + \frac{2s}{N}$ , then  $v \in L^1_{\mu_s}(\mathbb{R}^N)$  and

$$\begin{aligned} u(x, z) &= \int_{\mathbb{R}^N} \mathcal{P}_s(x - y, z) v(y) dy \\ &\leq c_5 \int_{\mathbb{R}^N} \frac{z^{2s}}{(|x - y|^2 + z^2)^{\frac{N+2s}{2}}} |y|^{-\frac{2s}{p-1}} dy \\ &= c_5 z^{-\frac{2s}{p-1}} \int_{\mathbb{R}^N} \frac{|y|^{-\frac{2s}{p-1}}}{(|x_z - y|^2 + 1)^{\frac{N+2s}{2}}} dy, \end{aligned} \quad (4.5)$$

where we take  $x_z = \frac{x}{z}$ .

When  $0 \leq |x_z| \leq 8$ , we obtain that

$$\int_{\mathbb{R}^N} \frac{|y|^{-\frac{2s}{p-1}}}{(|x_z - y|^2 + 1)^{\frac{N+2s}{2}}} dy \leq \int_{\mathbb{R}^N} (|y|^2 + 1)^{-\frac{N+2s}{2}} |y|^{-\frac{2s}{p-1}} dy := c_6 < +\infty.$$

When  $|x_z| > 8$ , letting  $r = \frac{1}{2}|x_z|$ , we have that

$$\begin{aligned} \int_{B_r(0)} \frac{|y|^{-\frac{2s}{p-1}}}{(|x_z - y|^2 + 1)^{\frac{N+2s}{2}}} dy &\leq (1 + |x_z|^2)^{-\frac{N+2s}{2}} \int_{B_r(0)} |y|^{-\frac{2s}{p-1}} dy \\ &= c_7 |x_z|^{-N-2s} r^{N-\frac{2s}{p-1}}, \end{aligned}$$



$$\begin{aligned} \int_{B_r(x_z)} \frac{|y|^{-\frac{2s}{p-1}}}{(|x_z - y|^2 + 1)^{\frac{N+2s}{2}}} dy &\leq \left(\frac{r}{2}\right)^{-\frac{2s}{p-1}} \int_{B_r(x_z)} (|x_z - y|^2 + 1)^{-\frac{N+2s}{2}} dy \\ &\leq c_8 r^{-\frac{2s}{p-1}} (1 + r^{-2s}), \end{aligned}$$

and

$$\begin{aligned} \int_{\mathbb{R}^N \setminus (B_r(x_z) \cup B_r(0))} \frac{|y|^{-\frac{2s}{p-1}}}{(|x_z - y|^2 + 1)^{\frac{N+2s}{2}}} dy &\leq c_9 \int_{\mathbb{R}^N \setminus B_r(0)} (|y|^2 + 1)^{-\frac{N+2s}{2}} |y|^{-\frac{2s}{p-1}} dy \\ &\leq c_{10} r^{-\frac{2s}{p-1} - 2s}. \end{aligned}$$

Thus, we conclude that for  $|x_z| \geq 8$ ,

$$\int_{\mathbb{R}^N} \frac{|y|^{-\frac{2s}{p-1}}}{(|x_z - y|^2 + 1)^{\frac{N+2s}{2}}} dy \leq c_{11} |x_z|^{-\frac{2s}{p-1}}.$$

As a consequence, we obtain that for  $|x_z| > 0$ ,

$$u(x, z) \leq c_{12} z^{2s} (1 + |x_z|^2)^{-\frac{s}{p-1}} \leq c_{12} (z^2 + |x|^2)^{-\frac{s}{p-1}}.$$

Which ends the proof.  $\square$

**Remark.** Since (4.4) can imply (4.2) immediately by taking  $z = 0$ , the previous result shows that these two estimates are equivalent when  $p > 1 + \frac{2s}{N}$ .

## 4.2 Integral estimates

**Lemma 4.2** *Let  $p > 0$  and  $v \in C^{1,s}(\overline{\Omega} \setminus \{0\})$  be a nonnegative solution of (4.1) in  $\Omega \setminus \{0\}$ , which vanishes in  $\Omega^c$ . Then for any  $\xi \in C^{1,s}(\overline{\Omega} \setminus \{0\})$  with compact support in  $\overline{\Omega} \setminus \{0\}$ , there holds*

$$\int_{\Omega} (v(-\Delta)^s \xi + \epsilon v^p \xi) dx = 0. \quad (4.6)$$

Furthermore, if  $v^p \in L^1(\Omega)$ , which is always satisfied for  $\epsilon = -1$ , there exists  $k \geq 0$  such that

$$(-\Delta)^s v + \epsilon v^p = k \delta_0 \quad \text{in } \mathcal{D}'(\Omega). \quad (4.7)$$

**Proof.** Without loss of generality, we can assume that  $B_2 \subset \Omega$ . Let  $\eta_0 : \mathbb{R}^N \rightarrow [0, 1]$  be a nonnegative cut-off function in  $C_c^\infty(\mathbb{R}^N)$ , equal to 1 in  $B_1$ , vanishing in  $B_2^c$ . For  $0 < \epsilon < 1$  we set  $\tilde{\eta}_\epsilon(x) = \eta_0(\epsilon x)(1 - \eta_0(x/\epsilon))$  and denote  $v_\epsilon = \tilde{\eta}_\epsilon v$ . Then

$$\begin{aligned} (-\Delta)^s v_\epsilon(x) &= \tilde{\eta}_\epsilon(x) (-\Delta)^s v(x) + v(x) (-\Delta)^s \tilde{\eta}_\epsilon(x) \\ &\quad - c_{N,s} \int_{\mathbb{R}^N} \frac{(v(x) - v(y))(\tilde{\eta}_\epsilon(x) - \tilde{\eta}_\epsilon(y))}{|x - y|^{N+2s}} dy, \quad \forall x \in B_1 \setminus \{0\}. \end{aligned}$$

Furthermore,

$$(-\Delta)^s v_\epsilon(0) = \lim_{x \rightarrow 0} (-\Delta)^s v_\epsilon(x) = -c_{N,s} \int_{B_2 \setminus B_\epsilon} \frac{v(y) \eta_\epsilon(y)}{|y|^{N+2s}} dy.$$

We can assume that  $\xi$  vanishes in  $\overline{B_r}$  and we choose  $0 < \varepsilon < r/8$ . By integration, there holds

$$\int_{B_1} \xi(-\Delta)^s v_\varepsilon dx = \int_{B_1} v_\varepsilon(-\Delta)^s \xi dx.$$

Since  $v = v_\varepsilon + (1 - \tilde{\eta}_\varepsilon)v$  and  $\xi(1 - \tilde{\eta}_\varepsilon) = 0$ , we have that

$$\begin{aligned} \int_{B_1} [v(-\Delta)^s \xi + \varepsilon v^p \xi] dx &= \int_{B_1} [\xi v(-\Delta)^s \tilde{\eta}_\varepsilon + (1 - \tilde{\eta}_\varepsilon)v(-\Delta)^s \xi] dx \\ &\quad - c_{N,s} \int_{B_1 \setminus B_r} \left( \int_{\mathbb{R}^N} \frac{(v(x) - v(y))(\tilde{\eta}_\varepsilon(x) - \tilde{\eta}_\varepsilon(y))}{|x - y|^{N+2\alpha}} dy \right) \xi(x) dx \\ &= A_{1,\varepsilon} + A_{2,\varepsilon} + A_{3,\varepsilon}. \end{aligned}$$

Since  $v \in L^1(B_1)$  and  $|(-\Delta)^s \xi(x)| \leq c_{13}(1 + |x|)^{-N-2s}$  for any  $x \in \mathbb{R}^N$  and some  $c_{13} > 0$ ,  $A_{2,\varepsilon} \rightarrow 0$  when  $\varepsilon \rightarrow 0$ . We have also for  $r < |x| < 1$  (where  $\xi(x)$  may not vanish),

$$\begin{aligned} |(-\Delta)^s \tilde{\eta}_\varepsilon(x)| &\leq c_{N,s} \int_{B_{2\varepsilon}} \frac{1}{|x - y|^{N+2s}} dy + c_{N,s} \int_{B_{2/\varepsilon}} \frac{1}{|x - y|^{N+2s}} dy \\ &\leq c_{14} (\varepsilon^N (r - 2\varepsilon)^{-N-2s} + \varepsilon^{2s}), \end{aligned}$$

for some  $c_{14} > 0$  independent of  $\varepsilon$ . This implies that  $A_{1,\varepsilon} \rightarrow 0$  when  $\varepsilon \rightarrow 0$ . Finally, for  $x \in B_1 \setminus B_r$  and  $\varepsilon < \frac{r}{4}$ , there holds

$$\begin{aligned} |A_{3,\varepsilon}| &\leq \|\xi\|_{L^\infty} \left( \int_{B_1 \setminus B_r} \int_{B_{2\varepsilon}} \frac{|v(x) - v(y)|}{|x - y|^{N+2s}} dy dx + \int_{B_1 \setminus B_r} \int_{B_{\frac{1}{\varepsilon}}} \frac{|v(x) - v(y)|}{|x - y|^{N+2s}} dy dx \right) \\ &\leq c_{15} (r - 2\varepsilon)^{-N-2s} \|\xi\|_{L^\infty} \left( (\varepsilon^N + \varepsilon^{2s}) \int_{B_1} v(x) dx + \int_{B_{2\varepsilon}} v(y) dy + \int_{B_{\frac{1}{\varepsilon}}} v(y) d\mu_s(y) \right), \end{aligned}$$

where  $c_{15} > 0$  is independent of  $\varepsilon$ . Hence  $A_{3,\varepsilon} \rightarrow 0$  as  $\varepsilon \rightarrow 0$ . This implies (4.6).

If  $\varepsilon = -1$ , it follows from [29, Theorem 3] that  $v^p \in L^1(B_1)$  and there exists  $k \geq 0$  such that (4.7) holds.

If  $\varepsilon = 1$  and we assume now that  $F := v^p \in L^1(B_1)$ , the function  $w = v + \mathbb{G}_{s,1}[F]$  is  $s$ -harmonic and positive in  $B_1 \setminus \{0\}$ . By the extension to  $s$ -harmonic functions of Bôcher's theorem [29, Theorem 4], there exists  $k \geq 0$  such that

$$(-\Delta)^s w = k\delta_0 \quad \text{in } \mathcal{D}'(B_1(0)),$$

and this is (4.7). □

### 4.3 A priori estimate for Lane-Emden equation

The following result is proved in [36].

**Proposition 4.3** *Assume that  $0 < s < 1$ ,  $1 < p < \frac{N+2s}{N-2s}$  and  $u$  is a nonnegative solution of*

$$\begin{aligned} \mathcal{D}_s u &= 0 && \text{in } \mathbf{B}_1 \cap \mathbb{R}_+^{N+1} \\ \partial_{\nu^s} u(\cdot, 0) - u(\cdot, 0)^p &= 0 && \text{in } B_1 \cap \mathbb{R}^N \setminus \{0\}. \end{aligned}$$

*Then there exists a constant  $c_{16} = c_{16}(N, p, s) > 0$  such that*

$$u(x, 0) := v(x) \leq c_{16} |x|^{-\frac{2s}{p-1}} \quad \text{for all } x \in B_{\frac{1}{2}} \cap \mathbb{R}^N \setminus \{0\}. \quad (4.8)$$

The upper bound of  $v$  comes from the technique used in the proof which is based upon a blow-up argument and a Liouville type theorem only valid in this range of exponents. When  $v$  is a radial function, we have a much shorter proof.

**Proposition 4.4** *Let  $0 < s < 1$ ,  $p > 1$ ,  $\Omega = B_1$  and  $v$  be a positive and radially decreasing solution of (4.1) with  $\epsilon = -1$ . Then there exists a constant  $c_{17} = c_{17}(N, p, s) > 0$  such that*

$$v(x) \leq c_{17} |x|^{-\frac{2s}{p-1}} \quad \text{for all } x \in B_{\frac{1}{2}} \setminus \{0\}.$$

**Proof.** Let  $\zeta \in C_c^\infty(\overline{B_1})$  be a radially symmetric function with value in  $[0, 1]$  and such that

$$\zeta(x) = \begin{cases} 1 & \text{for } x \in \overline{B_{\frac{1}{2}}} \setminus B_{\frac{1}{4}} \\ 0 & \text{for } x \in B_{\frac{1}{8}} \cup B_1^c \end{cases}$$

and for  $n > 1$ ,  $\zeta_n(x) = \zeta(nx)$ . By Lemma 4.2, there holds

$$\int_{B_1} v^p \zeta_n^{p'}(x) dx = \int_{B_1} v(-\Delta)^s \zeta_n^{p'}(x) dx.$$

From [15, Lemma 2.3], there holds

$$(-\Delta)^s \zeta_n^{p'} = p' \zeta_n^{p'-1} (-\Delta)^s \zeta_n - K^2,$$

where

$$K^2 := K^2(x) = \frac{p'(p'-1)}{2} \zeta_n^{p'-2}(z_x) \int_{B_1} \frac{(\zeta_n(x) - \zeta_n(y))^2}{|x-y|^{N+2s}} dy$$

for some  $z_x \in \overline{B_1}$ . Hence

$$\int_{B_1} v^p \zeta_n^{p'}(x) dx = \int_{B_1} v(-\Delta)^s \zeta_n^{p'}(x) dx \leq p' \int_{B_1} v \zeta_n^{p'-1} (-\Delta)^s \zeta_n(x) dx.$$

Now

$$\begin{aligned} \int_{B_1} v \zeta_n^{p'-1} (-\Delta)^s \zeta_n(x) dx &= \int_{\frac{1}{8n} < |x| < \frac{1}{n}} v \zeta_n^{p'-1} (-\Delta)^s \zeta_n(x) dx \\ &\leq \left( \int_{\frac{1}{8n} < |x| < \frac{1}{n}} v^p \zeta_n^{p'}(x) dx \right)^{\frac{1}{p}} \left( \int_{\frac{1}{8n} < |x| < \frac{1}{n}} |(-\Delta)^s \zeta_n(x)|^{p'} dx \right)^{\frac{1}{p'}}. \end{aligned}$$

Therefore

$$\int_{B_1} v^p \zeta_n^{p'}(x) dx = \int_{\frac{1}{8n} < |x| < \frac{1}{n}} v^p \zeta_n^{p'}(x) dx \leq p^{p'} \int_{\frac{1}{8n} < |x| < \frac{1}{n}} |(-\Delta)^s \zeta_n(x)|^{p'} dx.$$

Now

$$\int_{\frac{1}{8n} < |x| < \frac{1}{n}} v^p \zeta_n^{p'}(x) dx \geq \int_{\frac{1}{4n} < |x| < \frac{1}{2n}} v^p \zeta_n^{p'}(x) dx \geq \omega_N \frac{4^N - 1}{4^N} n^{-N} v^p \left( \frac{1}{2n} \right),$$

and

$$\int_{\frac{1}{8n} < |x| < \frac{1}{n}} |(-\Delta)^s \zeta_n(x)|^{p'} dx = n^{2sp' - N} \int_{\frac{1}{8} < |y| < 1} |(-\Delta)^s \zeta(y)|^{p'} dy.$$

Since  $\zeta$  vanishes in  $B_1^c$ ,

$$(-\Delta)^s \zeta(x) = \begin{cases} c_{N,s} \text{p.v.} \int_{B_1} \frac{\zeta(x) - \zeta(y)}{|x - y|^{N+2s}} dy + c_{N,s} \zeta(x) \int_{B_1^c} \frac{dy}{|x - y|^{N+2s}} & \text{for } x \in B_1 \\ -c_{N,s} \int_{B_1} \frac{\zeta(y)}{|x - y|^{N+2s}} dy & \text{for } x \in B_1^c. \end{cases}$$

Therefore,  $\|(-\Delta)^s \zeta\|_{L^k(B_n)}$  is bounded independently of  $k$  and  $n$ . Combining the previous inequalities, we obtain

$$\omega_N \frac{4^N - 1}{4^N} n^{-N} v^p \left( \frac{1}{2n} \right) \leq n^{2sp' - N} \int_{\frac{1}{8} < |y| < 1} |(-\Delta)^s \zeta(y)|^{p'} dy,$$

which implies that there holds, for some  $c_{18} > 0$ ,

$$v \left( \frac{1}{2n} \right) \leq c_{18} n^{\frac{2s}{p-1}}.$$

Taking  $r = \frac{1}{2n}$ , we infer (2.1). □

**Remark.** When  $\Omega = B_1$ , it is easy to prove that any positive solution  $v$  of (4.1) satisfying

$$\lim_{x \rightarrow 0} v(x) = +\infty$$

is radially symmetric and decreasing.

#### 4.4 Some maximum principles

In this section we recall some fundamental maximum principles which will be used in the sequel. Their proof can be found in [12].

**Lemma 4.5** *Let  $\mathcal{O} \subset \mathbb{R}^N$  be any domain and  $v \in C(\overline{\mathcal{O}}) \cap L^1_{\mu_s}(\mathcal{O})$ . If  $v$  achieves the maximum at some point  $x_0 \in \mathcal{O}$ , then*

$$(-\Delta)^s v(x_0) \geq 0.$$

Furthermore equality holds if and only if

$$v(x) = v(x_0) \quad \text{for all } x \in \mathcal{O}.$$

The following corollaries follow from the strong maximum principle.

**Corollary 4.6** *Let  $\mathcal{O} \subset \mathbb{R}^N$  be a bounded regular domain and  $v \in C(\overline{\mathcal{O}})$  satisfy*

$$(-\Delta)^s v \leq 0 \quad \text{in } \mathcal{O}, \quad v = h \quad \text{in } \mathcal{O}^c,$$

where  $h$  is continuous and bounded in  $\mathcal{O}^c$ . Then

$$\sup_{x \in \mathcal{O}} v(x) \leq \sup_{x \in \mathcal{O}^c} h(x).$$

**Corollary 4.7** *Let  $\mathcal{O} \subset \mathbb{R}^N$  be a bounded regular domain,  $g : \mathbb{R} \rightarrow \mathbb{R}$  be a continuous nondecreasing function and  $v_j \in C(\overline{\mathcal{O}})$  ( $j=1,2$ ) satisfy the two inequalities*

$$(-\Delta)^s v_1 + g(v_1) \geq 0 \quad \text{and} \quad (-\Delta)^s v_2 + g(v_2) \leq 0 \quad \text{in } \mathcal{O}.$$

Assume furthermore that  $v_1 \geq v_2$  a.e. in  $\mathcal{O}^c$ . Then either  $v_1 > v_2$  in  $\mathcal{O}$  or  $v_1 \equiv v_2$  in  $\mathbb{R}^N$ .

The next statement proved in [15, Lemma 3.1] is the analogue of the classical local regularity result for the Laplace operator, see also [28, Theorem 2.2] for local Schauder estimates.

**Lemma 4.8** *Assume that  $w \in C^{2s+\varepsilon}(\overline{B_1}) \cap L^1_{\mu_s}(\mathbb{R}^N)$  for some  $\varepsilon \in (0,1)$  satisfies*

$$(-\Delta)^s w = f \quad \text{in } B_1,$$

where  $f \in C^1(\overline{B_1})$ . Then for  $\beta \in (0, 2s)$ , there exist  $c_{19}, c_{20} > 0$  such that

$$\|w\|_{C^\beta(\overline{B_{1/2}})} \leq c_{19} \left( \|w\|_{L^\infty(B_1)} + \|f\|_{L^\infty(B_1)} + \|w\|_{L^1_{\mu_s}(\mathbb{R}^N)} \right)$$

and

$$\|w\|_{C^{2s+\beta}(\overline{B_{1/4}})} \leq c_{20} \left( \|w\|_{C^\beta(B_{\frac{1}{2}})} + \|f\|_{C^\beta(B_{\frac{1}{2}})} + \|w\|_{L^1_{\mu_s}(\mathbb{R}^N)} \right).$$

The next general result proved in [20] is an important tool for obtaining the existence of solutions of semilinear equations.

**Proposition 4.9** *Assume that  $\Omega \subset \mathbb{R}^N$  is a bounded regular domain,  $g : \mathbb{R} \rightarrow \mathbb{R}$  is a continuous nondecreasing function,  $f \in C^\beta(\Omega)$  with  $\beta \in (0,1)$ . If there exist a super-solution  $\bar{v}$  and a sub-solution  $\underline{v}$  of*

$$\begin{aligned} (-\Delta)^s v + g(v) &= f && \text{in } \Omega \setminus \{0\} \\ v &\geq 0 && \text{in } \Omega \setminus \{0\} \\ v &= 0 && \text{in } \Omega^c \end{aligned} \tag{4.9}$$

such that  $\bar{v}, \underline{v} \in C^2(\Omega \setminus \{0\}) \cap L^1(\Omega)$ ,

$$\bar{v} \geq \underline{v} \geq 0 \quad \text{in } \Omega \setminus \{0\} \quad \text{and} \quad \bar{v} \geq 0 \geq \underline{v} \quad \text{in } \Omega^c.$$

Then there exists at least one solution  $v$  of (4.9) which satisfies  $\underline{v} \leq v \leq \bar{v}$  in  $\Omega \setminus \{0\}$ .

We end this subsection with some comparison statements.

**Proposition 4.10** *Let  $v$  be a nonnegative  $s$ -subharmonic function in  $\mathbb{R}^N \setminus \{0\}$  satisfying*

$$\lim_{x \rightarrow 0} |x|^{N-2s} v(x) = 0 \quad \text{and} \quad \lim_{|x| \rightarrow +\infty} v(x) = 0,$$

then

$$v \equiv 0 \quad \text{in} \quad \mathbb{R}^N.$$

**Proof.** The function  $x \mapsto w(x) := |x|^{2s-N} + 1$  is  $s$ -harmonic in  $\mathbb{R}^N \setminus \{0\}$  and for any  $\varepsilon > 0$  the function  $v - \varepsilon w$  is  $s$ -subharmonic in  $\mathbb{R}^N \setminus \{0\}$  and negative in  $B_R^c$  for  $R > 0$  large enough and in  $B_\delta$  for  $\delta > 0$  small enough. Taking  $\mathcal{O} = B_R \setminus \overline{B_\delta}$  and  $h = 0$  in Corollary 4.6, it implies that for any  $\varepsilon > 0$ ,

$$0 \leq v \leq \varepsilon w \quad \text{in} \quad \mathbb{R}^N \setminus \{0\}.$$

Thus passing to the limit as  $\varepsilon \rightarrow 0$  implies the claim.  $\square$

As a variant of this result we have

**Corollary 4.11** *Let  $\Omega \subset \mathbb{R}^N$  be a bounded regular domain containing the origin,  $g : \mathbb{R}_+ \rightarrow \mathbb{R}_+$  be a continuous function such that  $g(0) = 0$  and  $h$  be a nonnegative function bounded in  $\mathbb{R}^N \setminus \Omega$ . If  $v$  is a nonnegative solution of*

$$(-\Delta)^s v + g(v) = 0 \quad \text{in} \quad \Omega \setminus \{0\}, \quad v = h \quad \text{in} \quad \Omega^c,$$

satisfying  $\lim_{x \rightarrow 0} v(x)|x|^{N-2s} = 0$ . Then  $v \leq \|h\|_{L^\infty}$  in  $\mathbb{R}^N$ .

## 4.5 Fractional Hardy operator

For  $\tau \in (-N, 2s)$ , we define  $\mathcal{C}_s(\tau)$  by the expression

$$\mathcal{C}_s(\tau) = 2^{2s} \frac{\Gamma(\frac{N+\tau}{2})\Gamma(\frac{2s-\tau}{2})}{\Gamma(-\frac{\tau}{2})\Gamma(\frac{N-2s+\tau}{2})} = -\frac{c_{N,s}}{2} \int_{\mathbb{R}^N} \frac{|\mathbf{e}_1 + z|^\tau + |\mathbf{e}_1 - z|^\tau - 2}{|z|^{N+2s}} dz,$$

where  $\mathbf{e}_1 = (1, 0, \dots, 0) \in \mathbb{R}^N$ . Since the function  $\Gamma$  is infinite at 0,  $\mathcal{C}_s$  vanishes at 0 and  $2s - N$ , i.e.

$$\mathcal{C}_s(0) = \mathcal{C}_s(2s - N) = 0.$$

The role of  $\mathcal{C}_s(\tau)$  is enlightened by the following identity [17]

$$(-\Delta)^s |x|^\tau = \mathcal{C}_s(\tau) |x|^{\tau-2s} \quad \text{for all } x \in \mathbb{R}^N \setminus \{0\}.$$

The function  $\mathcal{C}_s$  is concave and achieves its maximum in  $(-N, 2s)$  for  $\tau = \frac{2s-N}{2}$  with corresponding maximal value  $2^{2s} \frac{\Gamma^2(\frac{N+2s}{4})}{\Gamma^2(\frac{N-2s}{4})}$  (see [17, Lemma 2.3]). Furthermore,

$$\mathcal{C}_s(\tau) = \mathcal{C}_s(2s - N - \tau) \quad \text{for } \tau \in (-N, 2s),$$

and

$$\lim_{\tau \rightarrow -N} \mathcal{C}_s(\tau) = \lim_{\tau \rightarrow 2s} \mathcal{C}_s(\tau) = -\infty.$$

Let  $\tau_p = -\frac{2s}{p-1}$ , then we have that for  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$ ,

$$\tau_p \in (-N, 2s - N) \quad \text{and} \quad \mathcal{C}_s(\tau_p) > 0$$

and for  $p > \frac{N}{N-2s}$

$$\tau_p \in (2s - N, 0) \quad \text{and} \quad \mathcal{C}_s(\tau_p) < 0.$$

Let

$$c_p = \begin{cases} \mathcal{C}_s(\tau_p)^{\frac{1}{p-1}} & \text{if } 1 + \frac{2s}{N} < p < \frac{N}{N-2s} \\ (-\mathcal{C}_s(\tau_p))^{\frac{1}{p-1}} & \text{if } p > \frac{N}{N-2s}, \end{cases} \quad (4.10)$$

then the function  $U_p$  expressed by

$$U_p(x) := c_p |x|^{-\frac{2s}{p-1}} \quad \text{in } \mathbb{R}^N \setminus \{0\}, \quad (4.11)$$

satisfies, for  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$ ,

$$(-\Delta)^s U_p + U_p^p = 0 \quad \text{in } \mathbb{R}^N \setminus \{0\},$$

and for  $p > \frac{N}{N-2s}$

$$(-\Delta)^s U_p = U_p^p \quad \text{in } \mathbb{R}^N \setminus \{0\}.$$

The fractional Hardy operator is defined by

$$\mathcal{L}_\mu^s v(x) = (-\Delta)^s v(x) + \frac{\mu}{|x|^{2s}} v(x),$$

under the condition

$$\mu \geq \mu_0 := -2^{2s} \frac{\Gamma^2(\frac{N+2s}{4})}{\Gamma^2(\frac{N-2s}{4})}.$$

It is shown in [17] that the linear equation

$$\mathcal{L}_\mu^s v = 0 \quad \text{in } \mathbb{R}^N \setminus \{0\}$$

has two distinct radial solutions  $\Phi_\mu$  and  $\Gamma_\mu$  defined by

$$\Phi_\mu(x) = \begin{cases} |x|^{\tau_-(\mu)} & \text{if } \mu > \mu_0 \\ |x|^{-\frac{N-2s}{2}} \ln\left(\frac{1}{|x|}\right) & \text{if } \mu = \mu_0 \end{cases} \quad \text{and} \quad \Gamma_\mu(x) = |x|^{\tau_+(\mu)},$$

where the exponents  $\tau_-(\mu) \leq \tau_+(\mu)$  verify the relations

$$\begin{aligned} \tau_-(\mu) + \tau_+(\mu) &= 2s - N \quad \text{for all } \mu \geq \mu_0, \\ \tau_-(\mu_0) = \tau_+(\mu_0) &= \frac{2s-N}{2}, \quad \tau_-(0) = 2s - N, \quad \tau_+(0) = 0, \\ \lim_{\mu \rightarrow +\infty} \tau_-(\mu) &= -N \quad \text{and} \quad \lim_{\mu \rightarrow +\infty} \tau_+(\mu) = 2s. \end{aligned}$$

In the remaining of the paper and when there is no ambiguity, we use the notations  $\tau_+ = \tau_+(\mu)$ ,  $\tau_- = \tau_-(\mu)$ . The following inequalities are proved in [17],

$$\mathcal{C}_s(\tau) + \mu > 0 \quad \text{for } \tau \in (\tau_-, \tau_+)$$

and

$$\mathcal{C}_s(\tau) + \mu < 0 \quad \text{for } \tau \in (-N, \tau_-) \cup (\tau_+, 2s).$$

**Definition 3.1** We denote by  $W^s(\Omega)$  the space of functions  $v \in L^1(\mathbb{R}^N, d\mu_s)$  such that  $v|_{\Omega} \in L^2(\Omega)$  which satisfy

$$\mathcal{E}_{\Omega, \Omega'}^s(v, v) < +\infty \quad \text{for some domain } \Omega' \subset \mathbb{R}^N \text{ with } \Omega \subset\subset \Omega',$$

where we have set

$$\mathcal{E}_{A,B}^s(v, w) = \frac{c_{N,s}}{2} \int_{A \times B} \frac{(v(x) - v(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy.$$

We define  $W_*^s(\Omega)$  as the space of functions  $v \in L^1(\mathbb{R}^N, d\mu_s)$  such that  $v|_{\Omega \setminus B_\epsilon} \in W^s(\Omega \setminus B_\epsilon)$  for every  $\epsilon > 0$ .

The comparison principle for the fractional Hardy operator is the following, see [17, Lemma 4.11],

**Lemma 4.12** *Let  $v \in W_*^s(\Omega)$  satisfy  $\mathcal{L}_\mu^s v \geq 0$  in  $\Omega \setminus \{0\}$  and*

$$v \geq 0 \quad \text{in } \mathbb{R}^N \setminus \Omega, \quad \liminf_{x \rightarrow 0} \frac{v(x)}{\Phi_\mu(x)} \geq 0.$$

*Then  $v \geq 0$  in  $\mathbb{R}^N$ .*

As an application we have the following Liouville type result.

**Corollary 4.13** *Let  $0 < s < 1$  and  $p > 0$ . If  $v$  is a solution of*

$$(-\Delta)^s v + |v|^{p-1} v = 0 \quad \text{in } \mathbb{R}^N \setminus \{0\}, \tag{4.12}$$

*satisfying*

$$\lim_{x \rightarrow 0} |x|^{N-2s} v_+(x) = 0 \quad \text{and} \quad \lim_{|x| \rightarrow +\infty} |x|^{\tau_0} v_+(x) = 0, \tag{4.13}$$

*for some  $\tau_0 \in (0, 2s) \cap [2s(p+1) - Np, +\infty)$ , then  $v_+ \equiv 0$ .*

**Proof.** Let  $\varepsilon > 0$  and

$$w_\varepsilon(x) = \varepsilon(|x|^{2s-N} + 1) + \varepsilon^{p+1}|x|^{\tau_0} \quad \text{for } x \in \mathbb{R}^N \setminus \{0\}.$$

Then

$$w_\varepsilon^p(x) \geq \varepsilon^p(|x|^{(2s-N)p} + 1),$$

and

$$\begin{aligned} (-\Delta)^s w_\varepsilon(x) + w_\varepsilon^p(x) &= \varepsilon^{p+1} \mathcal{C}_s(\tau_0) |x|^{\tau_0 - 2s} + w_\varepsilon^p(x) \\ &\geq \varepsilon^p (\varepsilon \mathcal{C}_s(\tau_0) |x|^{\tau_0 - 2s} + |x|^{(2s-N)p} + 1), \end{aligned} \tag{4.14}$$

where  $\mathcal{C}_s(\tau_0) < 0$ . Since  $\tau_0 \in (0, 2s) \cap [2s(p+1) - Np, +\infty)$ , we have that  $\tau_0 - 2s \geq (2s - N)p$ . Thus, taking  $\varepsilon > 0$  small enough we infer that the right-hand side of (4.14) is nonnegative. Hence  $w_\varepsilon$  is a supersolution of (4.12), larger than  $v$  at 0 and at  $\infty$ . Thus  $v \leq w_\varepsilon$ . Letting  $\varepsilon \rightarrow 0$  yields the claim.  $\square$



## 4.6 Emden-Fowler equation with Dirac mass

The Emden-Fowler equation with right-hand side measure is studied in [15] and [16]. A particular case which is important in our present study is the following problem

$$\begin{aligned} (-\Delta)^s v + v^p &= k\delta_0 & \text{in } \mathcal{D}'(\Omega) \\ v &= 0 & \text{in } \Omega^c, \end{aligned} \quad (4.15)$$

where  $\Omega \subset \mathbb{R}^N$  is a bounded domain containing the origin. We set

$$\mathbb{X}_{s,\Omega} = \left\{ \zeta \in C(\mathbb{R}^N) : \zeta = 0 \text{ in } \Omega^c, |(-\Delta)^s \zeta| \in L^\infty(\mathbb{R}^N), |(-\Delta)_\varepsilon^s \zeta| \leq \phi \text{ if } 0 < \varepsilon \leq \varepsilon_0 \right\},$$

where

$$(-\Delta)_\varepsilon^s \zeta(x) = \int_{|x-y|>\varepsilon} \frac{\zeta(x) - \zeta(y)}{|x-y|^{N+2s}} dy$$

and  $\phi \in L^1(\Omega, d_{\partial\Omega} dx)$  with  $d_{\partial\Omega}(x) = \text{dist}(x, \Omega^c)$ . The following result is proved in [15].

**Proposition 4.14** *Let  $0 < p < \frac{N}{N-2s}$ , then for any  $k \geq 0$  there exists a unique function  $v = v_k$  belonging to  $L^1(\Omega)$  such that  $v_k^p \in L^1(\Omega, d_{\partial\Omega} dx)$ , satisfying*

$$\begin{aligned} \int_{\Omega} (v_k (-\Delta)^s \zeta + v_k^p \zeta) dx &= k\zeta(0) & \text{for all } \zeta \in \mathbb{X}_{s,\Omega} \\ v_k &= 0 & \text{in } \Omega^c \end{aligned}$$

and the following estimate holds

$$kG_{s,\Omega}(x, 0) - c_{21}k^p \leq v_k(x) \leq kG_{s,\Omega}(x, 0) \leq kG_s(x, 0) \quad \text{for all } x \in \Omega \setminus \{0\}, \quad (4.16)$$

where  $c_{21} > 0$ ,  $G_{s,\Omega}$  and  $G_s$  denote the Green functions of  $(-\Delta)^s$  in  $\Omega$  and  $\mathbb{R}^N$  respectively. Furthermore, the mapping  $k \mapsto v_k$  is increasing and it is stable in the sense that if  $\{\rho_\varepsilon\}$  is a sequence of nonnegative functions with compact support in  $\Omega$  converging to  $k\delta_0$  in the sense of distributions in  $\Omega$ , then the sequence of solutions  $v := v_{\rho_\varepsilon}$  of

$$(-\Delta)^s v + v^p = \rho_\varepsilon \text{ in } \mathcal{D}'(\Omega), \quad v = 0 \text{ in } \Omega^c,$$

converges to  $v_{\Omega,k}$  uniformly on any compact subset of  $\bar{\Omega} \setminus \{0\}$ , and  $v_{\rho_\varepsilon}$  converges to  $v_{\Omega,k}$  in  $L^p(\Omega, d_{\partial\Omega} dx)$ .

Since  $k \mapsto v_k$  is increasing, we set

$$v_\infty(x) = \lim_{k \rightarrow +\infty} v_k(x) \quad \text{for all } x \in \Omega \setminus \{0\}.$$

**Proposition 4.15** *Let  $0 < p < \frac{N}{N-2s}$ , then*

1- *If  $0 < p \leq 1 + \frac{2s}{N}$ , then  $v_\infty(x) = \infty$  for all  $x \in \Omega$ .*

2- *If  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$ , then  $v_\infty$  is a positive solution of*

$$(-\Delta)^s v + v^p = 0 \text{ in } \Omega \setminus \{0\}, \quad v = 0 \text{ in } \Omega^c, \quad (4.17)$$

satisfying

$$\lim_{x \rightarrow 0} |x|^{\frac{2s}{p-1}} v_\infty(x) = c_p, \quad (4.18)$$

and

$$U_p(x) - c^* \leq v_\infty(x) \leq U_p(x) \quad \text{for all } x \in \Omega \setminus \{0\}, \quad (4.19)$$

where  $c^* > 0$ ,  $c_p$  and  $U_p$  are defined in (4.10) and (4.11) respectively.

**Proof.** When  $1 + \frac{2s}{N} \geq \frac{2s}{N-2s}$ , the above result is proved in [16, Theorem 1.1], but when  $1 + \frac{2s}{N} < \frac{2s}{N-2s}$  the situation is not yet completely clarified [16, Theorem 1.2]. In this proof, we give a proof to overcome this gap.

To distinguish the effect of domain, we replace  $v_k$  by  $v_{\Omega,k}$  and for simplicity in the proof of this Proposition, we set

$$v_{0,k} = v_{\mathbb{R}^N,k} \quad \text{and} \quad v_{0,\infty} = v_{\mathbb{R}^N,\infty},$$

and

$$v_{1,k} = v_{B_1,k} \quad \text{and} \quad v_{1,\infty} = v_{B_1,\infty}.$$

From [18, Theorem 1.4] for  $\Omega = \mathbb{R}^N$ , there holds

(i) If  $p \in (0, 1 + \frac{2s}{N}]$ , then  $v_{0,\infty}(x) = \infty$  for all  $x \in \mathbb{R}^N \setminus \{0\}$ ;

(ii) If  $p \in (1 + \frac{2s}{N}, \frac{N}{N-2s})$ , then  $v_{0,\infty}(x) = U_p(x) = c_p |x|^{-\frac{2s}{p-1}}$ ,  $x \in \mathbb{R}^N \setminus \{0\}$ .

Assume that  $B_1 \subset \Omega$ , then it follows by the comparison principle that for any  $k > 0$ ,

$$v_{1,k} \leq v_{\Omega,k} \leq v_{0,k} \quad \text{in } \mathbb{R}^N \setminus \{0\}. \quad (4.20)$$

Moreover, we see that  $v_{1,k}$  is radially symmetric and decreasing with respect to  $|x|$ , and

$$\lim_{x \rightarrow 0} v_{1,k}(x) |x|^{N-2s} = c_N k.$$

*Step 1: we prove that  $v_{\Omega,\infty}$  is infinite everywhere in  $\Omega$  when  $p \in (0, 1 + \frac{2s}{N}]$ .*

Since the mapping  $k \mapsto v_{1,k}$  is monotone, if there is one point  $x_0 \notin B_1 \setminus \{0\}$  such that  $\lim_{k \rightarrow +\infty} v_{1,k}(x_0)$  is finite, then up to changing  $x_0$  we can assume that  $\lim_{k \rightarrow +\infty} v_{1,k}(z) < \infty$  for all  $z$  such that  $|z| \geq |x_0| - \epsilon$  for some  $\epsilon > 0$ . Hence

$$\begin{aligned} v_{1,k}^p(x_0) &= c_{N,s} \int_{|y| < |x_0| - \epsilon} \frac{v_{1,k}(y) - v_{1,k}(x_0)}{|y - x_0|^{N+2s}} dy + c_{N,s} P.V. \int_{|y| > |x_0| - \epsilon} \frac{v_{1,k}(y) - v_{1,k}(x_0)}{|y - x_0|^{N+2s}} dy \\ &= A_k + B_k. \end{aligned}$$

Since  $v_{1,k}(y)$  is uniformly bounded in  $B_1 \setminus B_{|x_0| - \epsilon}$ , the term  $B_k$  remains bounded too. If there exists  $y_0 \in B_{|x_0| - \epsilon} \setminus \{0\}$  such that  $\lim_{k \rightarrow +\infty} v_{1,k}(y_0) = \infty$  then  $\lim_{k \rightarrow +\infty} v_{1,k}(y) = \infty$  for all  $y \in B_{y_0} \setminus \{0\}$ ,

hence  $A_k \rightarrow \infty$  which contradicts the fact that  $v_{1,k}^p(x_0)$  is uniformly bounded. Consequently,

(i) either  $\lim_{k \rightarrow +\infty} v_{1,k}(x) = \infty$  for all  $x \in B_1 \setminus \{0\}$ ,

(ii) or  $\lim_{k \rightarrow +\infty} v_{1,k}(x) := v_{1,\infty}(x) < \infty$  for all  $x \in B_1 \setminus \{0\}$ .

Let us assume that (ii) occurs. Then  $v_{1,\infty}$  is a classical solution of

$$\begin{aligned} (-\Delta)^s v + v^p &= 0 \quad \text{in } B_1 \setminus \{0\} \\ v &= 0 \quad \text{in } B_1^c. \end{aligned} \tag{4.21}$$

Let  $\eta_0 : \mathbb{R}^N \rightarrow [0, 1]$  be a smooth, radially symmetric function such that

$$\eta_0 = 1 \quad \text{in } B_{\frac{1}{2}} \quad \text{and} \quad \eta_0 = 0 \quad \text{in } B_{\frac{3}{4}}^c.$$

Set

$$\tilde{v} = v_{1,k} \eta_0 \quad \text{for } k \leq \infty.$$

A straightforward calculation shows that there exists  $c_{22} > 0$ , independent of  $k$  such that

$$|(-\Delta)^s \tilde{v}(x) + \tilde{v}^p(x)| \leq c_{22}(1 + |x|)^{-N-2s} \quad \text{for all } x \in \mathbb{R}^N \setminus \{0\}.$$

Therefore there exists  $l_0 = l_0(N, s, p) > 0$  such that

$$(-\Delta)^s(2\tilde{v} + l_0) + (2\tilde{v} + l_0)^p > 0 \quad \text{in } \mathbb{R}^N \setminus \{0\}.$$

Hence  $2\tilde{v} + l_0$  is a positive super solution of

$$(-\Delta)^s v + v^p = 0 \quad \text{in } \mathbb{R}^N \setminus \{0\}.$$

From ([16]) we have that for any finite  $k > 0$

$$\lim_{x \rightarrow 0} \frac{v_{0,k}(x)}{v_{1,k}(x)} = 1 \quad \text{and} \quad \lim_{|x| \rightarrow +\infty} v_{0,k}(x) = 0.$$

The comparison principle applies and we obtain that for any  $k > 0$ , there holds

$$v_{0,k} \leq 2\tilde{v} + l_0 \quad \text{in } \mathbb{R}^N \setminus \{0\}.$$

Since  $p \in (0, 1 + \frac{2s}{N}]$  it is proved in ([16]) that  $\lim_{k \rightarrow +\infty} v_{0,k}(x) = \infty$  for all  $x \in \mathbb{R}^N \setminus \{0\}$ , which is a contradiction. Therefore  $v_{1,\infty} \equiv \infty$ .

Since  $v_{1,k} \leq v_{\Omega,k} \leq v_{0,k}$  we infer that  $v_{\Omega,\infty} \equiv \infty$ .

*Step 2: Asymptotics of  $v_{\Omega,\infty}$  for  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$ .*

If  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$ , it is proved in ([16]) that

$$\lim_{k \rightarrow +\infty} v_{0,k} = U_p.$$

Let  $H_s$  be the  $s$ -harmonic extension of  $U_p$  in  $B_1$ , i.e.

$$(-\Delta)^s H_s = 0 \quad \text{in } B_1, \quad H_s = U_p \quad \text{in } B_1^c.$$

Then  $H_s$  is bounded and positive in  $B_1$ . Note that  $H_s \geq v_{0,k}$  in  $B_1^c$  for any  $k > 0$ . We define  $W_k$  in  $\mathbb{R}^N \setminus \{0\}$  by

$$W_k = v_{0,k} - H_s.$$

Then there exist a sequence  $\{r_k\} \subset (0, 1)$  such that  $r_k \rightarrow 1$  and  $W_k(x) > 0$  for  $0 < |x| < r_k$ . Actually, this is due to the fact that for any  $k > 0$

$$v_{0,k} = c_N k |x|^{2s-N} (1 + o(1)) \quad \text{as } x \rightarrow 0 \quad \text{and} \quad \lim_{k \rightarrow +\infty} v_k = 0.$$

Moreover,  $W_k$  is a subsolution of (4.21). Therefore, comparison principle implies that for any  $k > 0$

$$v_{\Omega,k} \geq v_{1,k} \geq W_k \quad \text{in } B_1 \setminus \{0\}.$$

Passing to the limit as  $k \rightarrow +\infty$ , we have that

$$v_{\Omega,\infty} \geq v_{1,\infty} \geq U_p - H_s \quad \text{in } B_1 \setminus \{0\}.$$

Passing to the limit of (4.20) as  $k \rightarrow +\infty$ , we obtain that  $v_{\Omega,\infty} \leq v_{0,\infty} = U_p$  in  $B_1 \setminus \{0\}$ .

As a consequence, we have that

$$U_p - \sup_{\mathbb{R}^N} H_s \leq v_{\Omega,\infty} \leq U_p \quad \text{in } B_1 \setminus \{0\},$$

which also implies (4.18). □

The next result leads to an integral criteria for characterizing positive solutions of Emden-Fowler equations with strong singularity without *a priori* estimate.

**Proposition 4.16** *Assume  $p \in (0, \frac{N}{N-2s})$  and  $B_1 \subset \Omega$ . If  $v$  is a positive solution of (4.17) such that*

$$\int_{B_1} v^p dx = +\infty, \tag{4.22}$$

*then  $v \geq v_{1,k}$  in  $B_1 \setminus \{0\}$  for any  $k > 0$ , where  $v_{1,k}$  is the solution of (4.17). As a consequence,  $v \geq v_{1,\infty}$ .*

**Proof.** Let  $\alpha > 1$ , then  $w_k = \alpha^{-\frac{1}{p-1}} v_{1,k}$  satisfies

$$\begin{aligned} (-\Delta)^s w_k + \alpha w_k^p &= \alpha^{-\frac{1}{p-1}} k \delta_0 & \text{in } B_1 \\ w_k &= 0 & \text{in } B_1^c. \end{aligned} \tag{4.23}$$

The equation satisfied by  $v$  can be written under the form

$$\begin{aligned} (-\Delta)^s v + \alpha v^p &= (\alpha - 1)v^p & \text{in } \Omega \setminus \{0\} \\ v &= 0 & \text{in } \Omega^c. \end{aligned}$$

From (4.22) and regularity of  $v$  in  $\Omega \setminus \{0\}$ , for any  $\varepsilon \in (0, 1)$  there holds

$$\int_{B_\varepsilon} v^p dx = +\infty.$$

Hence for any  $k > 0$ , there exists  $\ell = \ell(\varepsilon) > 0$  such that

$$(\alpha - 1) \int_{B_\varepsilon} (\min\{v, \ell\})^p dx = k.$$

Let  $\tilde{v}_\varepsilon$  be the unique solution of

$$\begin{aligned} (-\Delta)^s v + \alpha v^p &= (\alpha - 1) (\min\{v, \ell\})^p \chi_{B_\varepsilon} && \text{in } B_1 \\ v &= 0 && \text{in } B_1^c. \end{aligned}$$

Note that  $\tilde{v}_\varepsilon$  is bounded in  $B_1$ .

The comparison principle (see Corollary 4.7) shows that for any  $l > 0$ ,

$$v \geq \tilde{v}_\varepsilon \quad \text{in } \mathbb{R}^N \setminus \{0\}.$$

Since  $\ell_\varepsilon \rightarrow \infty$  when  $\varepsilon \rightarrow 0$  and

$$(\alpha - 1) (\min\{v, \ell\})^p \chi_{B_\varepsilon} \rightarrow \alpha^{-\frac{1}{p-1}} k \delta_0,$$

then it follows by Proposition 4.14 that  $\tilde{v}_\varepsilon$  converges to the solution  $w_k$  of (4.23). Thus, we deduce that  $v \geq \alpha^{-\frac{1}{p-1}} v_{1,k}$ . Since  $k > 0$  is arbitrary, we infer the claim.  $\square$

In the following we establish a local integral *a priori* estimate

**Proposition 4.17** *Let  $p > 1 + \frac{2s}{N}$  and  $v$  be a nonnegative solution of (4.17). Then for any  $1 \leq q < \frac{p}{1 + \frac{2s}{N}}$ , there holds*

$$\left( \int_{B_{\frac{1}{2}}} v^q dx \right)^{\frac{1}{q}} \leq c_{23} + c_{24} \int_{\Omega \setminus B_{\frac{1}{2}}} v dx,$$

where  $c_{23}$  and  $c_{24}$  are positive constants depending on  $p, q, s$  and  $N$ .

**Proof.** We use the same test function  $\zeta$  as in the proof of Proposition 4.4 and we set  $\tilde{\zeta}_n(x) = \zeta(2^n x)$  and  $\Gamma_n = B_{2^{-n-1}} \setminus B_{2^{-n-2}}$ ,

$$\int_{\Gamma_n} v^p dx \leq \int_{\Gamma_n} v^p \tilde{\zeta}_n dx = \left| \int_{\Omega} (-\Delta)^s \tilde{\zeta}_n v dx \right| \leq 2^{2ns} \|(-\Delta)^s \zeta\|_{L^\infty(\mathbb{R}^N)} \|v\|_{L^1(\Omega)}.$$

Then

$$\int_{\Gamma_n} v^q dx \leq \left( \int_{\Gamma_n} v^p dx \right)^{\frac{q}{p}} \left( \int_{\Gamma_n} dx \right)^{1 - \frac{q}{p}} \leq c_{25} 2^{-\frac{N}{p}(p - q(1 + \frac{2s}{N}))n} \|(-\Delta)^s \zeta\|_{L^\infty(\mathbb{R}^N)}^{\frac{q}{p}} \|v\|_{L^1(\Omega)}^{\frac{q}{p}}.$$

Since  $1 \leq q < \frac{p}{1 + \frac{2s}{N}}$ , we set  $\alpha_0 = \frac{N}{p}(p - q(1 + \frac{2s}{N})) > 0$  by our assumption that  $q < \frac{p}{1 + \frac{2s}{N}}$ . Summing the above inequalities from  $n = 0$  to  $+\infty$ , we derive

$$\int_{B_{\frac{1}{2}}} v^q dx \leq \frac{2^{\alpha_0} c_{25}}{2^{\alpha_0} - 1} \|(-\Delta)^s \zeta\|_{L^\infty(\mathbb{R}^N)}^{\frac{q}{p}} \|v\|_{L^1(\Omega)}^{\frac{q}{p}},$$

therefore

$$\begin{aligned} \left( \int_{B_{\frac{1}{2}}} v^q dx \right)^{\frac{p}{q}} &\leq \left( \frac{2^{\alpha_0} c_{25}}{2^{\alpha_0} - 1} \right)^{\frac{p}{q}} \|(-\Delta)^s \zeta\|_{L^\infty(\mathbb{R}^N)} \|v\|_{L^1(\Omega)} \\ &\leq c_{26} \|v\|_{L^1(\Omega \setminus B_{\frac{1}{2}})} + c_{27} \left( \int_{B_{\frac{1}{2}}} v^q dx \right)^{\frac{1}{q}}, \end{aligned} \quad (4.24)$$

where  $c_{26}$  and  $c_{27}$  are positive constants depending on  $N$ ,  $p$ ,  $q$  and  $s$ . If we set  $X = \|v\|_{L^q(B_{\frac{1}{2}})}^{\frac{q}{p}}$ , (4.24) becomes

$$X^p \leq c_{26} \|v\|_{L^1(\Omega \setminus B_{\frac{1}{2}})} + c_{27} X.$$

Standard algebraic computations combined with a convexity argument lead to

$$\|v\|_{L^q(B_{\frac{1}{2}})} \leq c_{27}^{\frac{1}{p-1}} + \frac{c_{26}}{(p-1)c_{27}} \|v\|_{L^1(\Omega \setminus B_{\frac{1}{2}})}.$$

We complete the proof.  $\square$

#### 4.7 A priori estimate for Emden-Fowler equation

Without loss of generality, we can assume that  $\overline{B_1} \subset \Omega \subset \overline{\Omega} \subset B_\ell$  for some  $\ell > 1$ . The main estimate is the following

**Proposition 4.18** *Assume that  $s \in (0, 1)$ ,  $p > 1 + \frac{N}{2s}$  and  $v$  is a nonnegative function satisfying (4.17). Then there exists a constant  $c_{28} = c_{28}(N, p, s) > 0$  such that*

$$v(x) \leq c_{28} |x|^{-\frac{2s}{p-1}} \quad \text{for all } x \in B_{\frac{1}{2}}. \quad (4.25)$$

**Proof.** The proof follows the technique introduced by Gidas and Spruck in [24] dealing with the classical Lane-Emden equation and extended to the fractional Lane-Emden equation in [28] and [36]. Only the final argument is different and specific to the Emden-Fowler equation. We assume by contradiction that there exists a sequence  $\{x_k\} \subset B_1 \setminus \{0\}$  converging to 0 such that

$$|x_k|^{\frac{2s}{p-1}} v(x_k) = \max_{|x_k| \leq |x| \leq 1} |x|^{\frac{2s}{p-1}} v(x) \rightarrow +\infty \quad \text{as } k \rightarrow +\infty.$$

We set

$$\phi_k(x) = \left( \frac{|x_k|}{2} - |x - x_k| \right)^{\frac{2s}{p-1}} v(x), \quad \phi_k(\bar{x}_k) = \max_{|x - x_k| \leq \frac{|x_k|}{2}} \phi_k(x).$$

Since

$$(2\nu_k)^{\frac{2s}{p-1}} v(\bar{x}_k) = \phi_k(\bar{x}_k) \geq \phi_k(x) \geq \nu_k^{\frac{2s}{p-1}} v(x) \quad \text{for all } x \in B_{\nu_k}(\bar{x}_k),$$

which implies

$$2^{\frac{2s}{p-1}} v(\bar{x}_k) \geq v(x) \quad \text{if } |x - \bar{x}_k| \leq \nu_k.$$

Since  $\frac{1}{2}|\bar{x}_k| \geq \frac{1}{2}|x_k| - |\bar{x}_k - x_k|$ , we obtain

$$|\bar{x}_k|^{\frac{2s}{p-1}} v(\bar{x}_k) \geq |x_k|^{\frac{2s}{p-1}} v(x_k) \geq |x|^{\frac{2s}{p-1}} v(x) \quad \text{for all } x \in B_1 \setminus B_{|x_k|}$$

and

$$(2\nu_k)^{\frac{2s}{p-1}} v(\bar{x}_k) = \phi(\bar{x}_k) \geq \phi(x_k) \geq \left(\frac{|x_k|}{2}\right)^{\frac{2s}{p-1}} v(x_k) \rightarrow +\infty \quad \text{as } k \rightarrow +\infty.$$

Therefore,

$$w_k(y) \leq 2^{\frac{2s}{p-1}} \quad \text{for all } |y| \leq r_k := (v(\bar{x}_k))^{\frac{p-1}{2s}} \nu_k.$$

As in [36] we define the function  $W_k$  by

$$W_k(y, z) = \frac{1}{v(\bar{x}_k)} v\left(\frac{y}{(v(\bar{x}_k))^{\frac{p-1}{2s}}} + \bar{x}_k, \frac{z}{(v(\bar{x}_k))^{\frac{p-1}{2s}}}\right) \quad \text{for all } (y, z) \in \Omega_k \setminus \{(\xi_k, 0)\},$$

where

$$\Omega_k = \left\{ (y, z) \in \mathbb{R}_+^{N+1} : \left(\frac{y}{(v(\bar{x}_k))^{\frac{p-1}{2s}}} + \bar{x}_k, \frac{z}{(v(\bar{x}_k))^{\frac{p-1}{2s}}}\right) \in \Omega \times (0, 1) \right\}$$

and

$$\xi_k = -(v(\bar{x}_k))^{\frac{p-1}{2s}} \bar{x}_k.$$

We set  $w_k = W_k(\cdot, 0)$  and  $\mathcal{O}_k = \bar{\Omega}_k \cap \partial\mathbb{R}_+^{N+1}$ . The function  $W_k$  satisfies

$$\begin{aligned} \mathcal{D}_s W_k &= 0 && \text{in } \Omega_k \setminus \{(\xi_k, 0)\} \\ \partial_{\nu^s} W_k(\cdot, 0) + W_k(\cdot, 0)^p &= 0 && \text{in } \mathcal{O}_k \setminus \{(\xi_k, 0)\} \\ W_k(\cdot, 0) &= 0 && \text{in } \mathcal{O}_k^c, \end{aligned}$$

where  $\mathcal{O}_k^c$  denotes the complement of  $\mathcal{O}_k$  in  $\partial\mathbb{R}_+^{N+1}$ . Since  $w_k^{p-1} \leq 2^{2s}$  in  $B_{r_k}$ , there exists a constant  $c_{29} > 0$  depending on  $N, p$  and  $s$  such that

$$0 \leq W_k(y, z) \leq c_{29} \quad \text{for all } (y, z) \in B_{\frac{r_k}{2}} \times [0, \frac{r_k}{2}]. \quad (4.26)$$

It is a consequence of the Harnack inequality [36, Proposition 2.2] and the invariance of the problem by scaling. From this inequality follows a local regularity estimate of  $W_k$ , independent of  $k$  already obtained in [36] that we recall

$$\|W_k\|_{W^{1,2}(B_R \times (0, R), z^{1-2s})} + \|W_k\|_{C^\alpha(\bar{B}_R \times [0, R])} + \|w_k\|_{C^{2,\alpha}(\bar{B}_R)} \leq C(R)$$

for some  $\alpha > 0$  and any  $R > 0$ , where  $W^{1,2}(B_R \times (0, R), z^{1-2s})$  denotes the weighted Sobolev space with weight  $z^{1-2s} dx dz$ . Therefore there exists a subsequence of  $\{W_k\}$  still denoted by  $\{W_k\}$  and a nonnegative bounded function  $W$  such that  $W_k \rightarrow W$  weakly in  $W^{1,2}(\mathbb{R}_+^{N+1}, z^{1-2s})$  and in  $C_{loc}^s(\mathbb{R}_+^{N+1})$  and such that  $w_k \rightarrow w := W(\cdot, 0) \in C_{loc}^2(\mathbb{R}^N)$ . Therefore the functions  $W, w$  satisfy  $w(0) = 1$  and, because of (4.26), it is a nonnegative and bounded solution of

$$\begin{aligned} \mathcal{D}_s W &= 0 && \text{in } \mathbb{R}_+^{N+1} \\ \partial_{\nu^s} W(\cdot, 0) + W(\cdot, 0)^p &= 0 && \text{in } \mathbb{R}^N. \end{aligned}$$

By Theorem 2.1  $w$  belongs to  $L_{\mu_s}^1(\mathbb{R}^N)$ , thus  $\partial_{\nu^s} W(\cdot, 0) = (-\Delta)^s w$  and  $w$  satisfies (4.12). By Corollary 4.13 such a function cannot exist. This implies that (4.25) holds.  $\square$

The next statement is a surprising consequence of Proposition 4.18.

**Corollary 4.19** *Assume that  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$ . If  $v$  is a positive function satisfying (4.17) and*

$$\int_{\Omega} v^p dx = +\infty.$$

*Then there holds*

$$\liminf_{x \rightarrow 0} |x|^{\frac{2s}{p-1}} v(x) = c_p.$$

**Proof.** By Proposition 4.15 and Proposition 4.16 we have

$$\liminf_{x \rightarrow 0} |x|^{\frac{2s}{p-1}} v(x) \geq c_p.$$

Therefore we proceed by contradiction assuming that there exists a nonnegative solution  $v$  and  $\bar{\mu} > 0$  such that

$$c_p < \bar{\mu}^{1/(p-1)} < \liminf_{x \rightarrow 0} |x|^{\frac{2s}{p-1}} v(x) < +\infty. \quad (4.27)$$

We can write the equation satisfied by  $v$  under the form

$$\mathcal{L}_{\bar{\mu}}^s v = g(x) := \frac{\bar{\mu}}{|x|^{2s}} v - v^p.$$

The mapping  $\mu \mapsto \tau_-(\mu)$  is continuous and monotone decreasing from  $[\mu_0, +\infty)$  onto  $(-\infty, \frac{2s-N}{2}]$  and  $\tau_-(c_p^{p-1}) = -\frac{2s}{p-1}$ . Therefore, there exists  $\bar{\tau} > \frac{2s}{p-1}$  such that  $\tau_-(\bar{\mu}) = -\bar{\tau}$ . By (4.25) we have

$$\liminf_{x \rightarrow 0} |x|^{\bar{\tau}} v(x) = 0$$

for any  $\tau > \frac{2s}{p-1}$ . Choosing  $\tau = \bar{\tau}$  we see that there exist  $r, c_{30} > 0$  such that

$$0 < -g(x) \leq c_{30} |x|^{-p\bar{\tau}} \quad \text{for any } 0 < |x| < r.$$

Set  $m = r^{-\tau+(\bar{\mu})} v(r)$ . For any  $\varepsilon > 0$  the function  $w_{\varepsilon, m} := \varepsilon \Phi_{\bar{\mu}} + m \Gamma_{\bar{\mu}}$  satisfies  $\mathcal{L}_{\bar{\mu}}^s w_{\varepsilon, m} = 0$ . Therefore the function  $w_{\varepsilon, m} - v$  satisfies in turn

$$\mathcal{L}_{\bar{\mu}}^s (w_{\varepsilon, m} - v) = -g \geq 0 \quad \text{in } B_r \setminus \{0\}.$$

This function is nonnegative in  $\mathbb{R}^N \setminus B_r$  and there holds

$$\liminf_{x \rightarrow 0} \frac{(w_{\varepsilon, m} - v)(x)}{\Phi_{\bar{\mu}}(x)} = \varepsilon > 0.$$

By Lemma 4.12,  $w_{\varepsilon, m} - v \geq 0$ . Letting  $\varepsilon \rightarrow 0$  yields  $v(x) \leq m \Gamma_{\bar{\mu}}(x)$ . Thus (4.27) does not hold and no such  $\bar{\mu}$  exists, which implies the claim.  $\square$



## 5 Singularity of solutions

### 5.1 Energy method and limit set

The estimates (4.8) and (4.25) combined with the lifting estimate (4.4) allow to use an energy method to describe the behaviour of signed solutions of (3.2). Furthermore, this energy method does not depend of the signs of solutions, even if the bounds do. If we set

$$u(\xi) = u(r, \sigma) = r^{-\frac{2s}{p-1}} w(t, \sigma) \quad \text{with } t = \ln r \quad (5.1)$$

and  $w(t, \sigma) = w(t, \sigma', \phi)$ . Then  $w$  satisfies

$$\begin{aligned} w_{tt} + \Theta_{s,p,N} w_t + \Lambda_{s,p,N} w + \mathcal{A}_s[w] &= 0 \quad \text{in } \mathbb{R} \times \mathbb{S}_+^N \\ \frac{\partial w}{\partial \nu^s} + \epsilon |w|^{p-1} w &= 0 \quad \text{in } \mathbb{R} \times \mathbb{S}^{N-1}, \end{aligned}$$

where  $\Theta_{s,p,N}$  and  $\Lambda_{s,p,N}$  are defined in (3.5) and

$$\frac{\partial w}{\partial \nu^s}(t, \sigma') = - \lim_{\phi \rightarrow 0} (\sin \phi)^{1-2s} w_\phi(t, \sigma', \phi).$$

We define the negative (resp. positive) trajectory of  $w$  by

$$\mathcal{T}_-[w] = \bigcup \{w(t, \cdot), t \leq 0\},$$

respectively

$$\mathcal{T}_+[w] = \bigcup \{w(t, \cdot), t \geq 0\}.$$

The trajectory of  $w$  denoted by  $\mathcal{T}[w] = \mathcal{T}_-[w] \cup \mathcal{T}_+[w]$ . We also define the  $\alpha$ -limit set  $\Gamma_-[w]$  and the  $\omega$ -limit set  $\Gamma_+[w]$  of a trajectory by

$$\Gamma_-[w] = \bigcap_{t \leq 0} \left( \overline{\bigcup \{w(\tau, \cdot), \tau \leq t\}}^{\mathbf{X}} \right) \quad \text{and} \quad \Gamma_+[w] = \bigcap_{t \geq 0} \left( \overline{\bigcup \{w(\tau, \cdot), \tau \geq t\}}^{\mathbf{X}} \right),$$

where  $\mathbf{X} = \left\{ \zeta \in C^s(\mathbb{S}_+^N) : \zeta(\cdot, 0) \in C^2(\mathbb{S}^{N-1}) \right\}$ . If we assume that the set  $\{w(t, \cdot)\}$  is bounded in  $\mathbb{R}_- \times \mathbb{S}_+^N$  (resp. in  $\mathbb{R}_+ \times \mathbb{S}_+^N$ ), its trajectory  $\mathcal{T}_-[w]$  (resp.  $\mathcal{T}_+[w]$ ) is uniformly bounded in  $\left\{ \zeta \in C^\alpha(\mathbb{S}_+^N) : \zeta(\cdot, 0) \in C^{2+\alpha}(\mathbb{S}^{N-1}) \right\}$ . By Ascoli's theorem they are relatively compact in  $\mathbf{X}$ . By the general theory of limit sets,  $\Gamma_\pm[w]$  are compact and connected subsets of  $\mathbf{X}$ . We first give a general result about these limit sets under the assumption that the solutions are signed and satisfy *a priori* estimate which is indeed only proved for nonnegative solutions.

**Theorem 5.1** *Assume that  $s \in (0, 1)$ ,  $p \in (1, +\infty) \setminus \left\{ \frac{N+2s}{N-2s} \right\}$ ,  $\epsilon = \pm 1$  and  $u \in C(\overline{\mathbb{R}_+^{N+1}} \setminus \{(0, 0)\}) \cap C^2(\mathbb{R}_+^{N+1})$  is a solution of (3.2). If  $u$  satisfies*

$$|u(x, z)| \leq c_{31} (|x|^2 + z^2)^{-\frac{s}{p-1}} \quad \text{for } 0 < |x| < 1, \quad (\text{resp. for } |x| > 1) \quad (5.2)$$

*for some  $c_{31} > 0$ , then  $\Gamma_-[w]$  (resp.  $\Gamma_+[w]$ ) is a nonempty compact connected subset of the set  $\mathcal{E}_\epsilon$ , which is the set of solutions of (3.4).*

**Proof.** We give the proof for  $\Gamma_-[w]$ , the case  $\Gamma_+[w]$  being the same up to changing  $t$  in  $-t$ . We define the energy function

$$\mathcal{I}_\epsilon[w](t) = \frac{1}{2} \int_{\mathbb{S}_+^N} (|\nabla' w|^2 - \Lambda_{s,p,N} w^2 - w_t^2) d\mu_s - \frac{\epsilon}{p+1} \int_{\mathbb{S}^{N-1}} |w|^p dS'.$$

The function  $w$  is bounded in  $C^2(\overline{\mathbb{S}_+^N} \times (-\infty, -1])$  and there holds

$$\frac{d}{dt} \mathcal{I}_\epsilon[w](t) = \left( p - \frac{2s(p+1)}{p-1} \right) \int_{\mathbb{S}_+^N} w_t^2 d\mu_s = \Theta_{s,p,N} \int_{\mathbb{S}_+^N} w_t^2 d\mu_s. \quad (5.3)$$

Since  $p \neq \frac{N+2s}{N-2s}$ ,  $\Theta_{s,p,N} \neq 0$ , and we have the damping estimate

$$\int_{\mathbb{S}_+^N} w_t^2 d\mu_s < +\infty.$$

We conclude by using the uniform continuity of  $w_t$  that  $w_t(t, \cdot) \rightarrow 0$  uniformly on  $\overline{\mathbb{S}_+^N}$  when  $t \rightarrow -\infty$ . Let  $\omega \in \Gamma_-$ . There exists a sequence  $\{t_n\}$  converging to  $-\infty$  such that  $w(t_n, \cdot) \rightarrow w$  uniformly on  $\overline{\mathbb{S}_+^N}$ , and since  $w_t(t, \cdot) \rightarrow 0$  uniformly on  $\overline{\mathbb{S}_+^N}$ , for any  $T > 0$ ,  $w(t, \cdot) \rightarrow w$  uniformly on  $[t_n - T, t_n + T] \times \overline{\mathbb{S}_+^N}$ . Let  $\zeta \in C^2(\overline{\mathbb{S}_+^N})$  such that  $\partial_{\nu^s} \zeta = 0$ , then

$$\begin{aligned} & \int_{\mathbb{S}_+^N} (w_t(T + t_n, \cdot) - w_t(T - t_n, \cdot)) \zeta d\mu_s + \Theta_{s,p,N} \int_{\mathbb{S}_+^N} (w(T + t_n, \cdot) - w(T - t_n, \cdot)) \zeta d\mu_s \\ & - \Lambda_{s,p,N} \int_{-T}^T \int_{\mathbb{S}_+^N} w(t_n + t, \cdot) \zeta d\mu_s dt - \int_{-T}^T \int_{\mathbb{S}_+^N} w(t_n + t, \cdot) \mathcal{A}_s[\zeta] d\mu_s dt \\ & + \epsilon \int_{-T}^T \int_{\mathbb{S}^{N-1}} (|w|^{p-1} w)(t_n + t, \cdot) \zeta dS' dt = 0. \end{aligned}$$

Letting  $t_n \rightarrow -\infty$  yields

$$\int_{-T}^T \int_{\mathbb{S}_+^N} (\Lambda_{s,p,N} \omega \zeta + \omega \mathcal{A}_s[\zeta]) d\mu_s dt - \epsilon \int_{-T}^T \int_{\mathbb{S}^{N-1}} |\omega|^{p-1} \omega \zeta dS' dt = 0.$$

This implies that  $\omega$  is a weak solution of (3.4) and therefore a strong one.  $\square$

**Remark.** The case  $p = \frac{N+2s}{N-2s}$  is not covered since the damping coefficient  $\Theta_{s,p,N}$  in (5.3) vanishes.

## 5.2 Singularities of Emden-Fowler equation

The next result concerns the singularities of solutions to the Emden-Fowler equations.

**Theorem 5.2** *Assume that  $s \in (0, 1)$ ,  $\epsilon = 1$  and  $u \in C^1(\overline{\mathbb{R}_+^{N+1}} \setminus \{(0, 0)\}) \cap C^2(\mathbb{R}_+^{N+1})$  is a solution of (3.2) satisfying (5.2),  $v = u(\cdot, 0)$  and  $w$  is defined by (5.1).*

*1-a If  $p \geq \frac{N}{N-2s}$ , then  $w(t, \cdot)$  converges uniformly in  $\mathbb{S}_+^N$  to 0,  $v \equiv 0$ , hence  $u$  is a smooth solution;*

1-b If  $1 + \frac{2s}{N} < p < \frac{N}{N-2s}$  and  $u$  is nonnegative, then either  $w(t, \cdot)$  converges uniformly in  $\mathbb{S}_+^N$  to  $\omega_1$  or it converges to 0;

1-c If  $p^* \leq p < \frac{N}{N-2s}$ , then  $w(t, \cdot)$  converges uniformly in  $\mathbb{S}_+^N$  to some  $\ell \in \{0, \omega_1, -\omega_1\}$ ;

1-d If  $1 < p \leq 1 + \frac{2s}{N}$ , and  $u$  is nonnegative, then  $w(t, \cdot)$  converges uniformly in  $\mathbb{S}_+^N$  to 0.

2. Let  $1 < p < \frac{N}{N-2s}$ ,  $u$  is nonnegative, if  $w(t, \cdot)$  converges uniformly in  $\mathbb{S}_+^N$  to 0, then there exists  $k \in \mathbb{R}_+$  such that

$$e^{(N-\frac{2s}{p-1})t} w(t, \sigma', \phi) \rightarrow k \sin^{2s}(\phi) \quad \text{uniformly in } \mathbb{S}_+^N.$$

If  $k = 0$ , then  $v \equiv 0$  and  $u$  is smooth.

**Proof.** 1-a. If  $p \geq \frac{N}{N-2s}$ , the set  $\mathcal{E}_1$  is reduced to the zero function. Therefore  $w(t, \cdot)$  converges to 0 when  $t \rightarrow -\infty$ . By (3.3) the function  $\tilde{u}(t, \sigma) = u(r, \sigma)$  with  $t = \ln r$  satisfies

$$\begin{aligned} \tilde{u}_{tt} + (N-2s)\tilde{u}_t + \mathcal{A}_s[\tilde{u}] &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}_+^N \\ \frac{\partial \tilde{u}}{\partial \nu^s} + |\tilde{u}|^{p-1}\tilde{u} &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}^{N-1}. \end{aligned} \quad (5.4)$$

Since  $\frac{2s}{p-1} \leq N-2s$ , we have that  $\tilde{u}(t, \cdot) = o(e^{(2s-N)t})$  when  $t \rightarrow -\infty$ . Let  $m = \max |\tilde{u}(0, \cdot)|$ . For any  $\varepsilon > 0$ ,  $t \mapsto \varepsilon e^{(2s-N)t} + m$  is a supersolution of (5.4), larger than  $\tilde{u}$  at  $t = 0$  and  $t \rightarrow -\infty$ . Hence  $\tilde{u}(t, \sigma) \leq \varepsilon e^{(2s-N)t} + m$ . Letting  $\varepsilon \rightarrow 0^+$  yields  $\tilde{u}(t, \sigma) \leq m$ . Similarly  $\tilde{u}(t, \sigma) \geq -m$ , hence  $\tilde{u}$  is bounded and so is  $u$ . Therefore  $v (= u(\cdot, 0)) \in L^1_{\mu_s}(\mathbb{R}^N)$ ,

$$\frac{\partial \tilde{u}}{\partial \nu^s} = (-\Delta)^s v,$$

and  $v$  is a bounded solution of the fractional Emden-Fowler equation in  $\Omega \setminus \{0\}$  vanishing on  $\Omega^c$ . It is therefore identically 0. By the even extension of  $u$  to  $\mathbb{R}^{N+1}$ , it follows that  $u$  is smooth.

1-b Since  $\mathcal{E}_1^+ = \{\omega_1, 0\}$  and is disconnected, the result follows.

1-c Since  $\mathcal{E}_1 = \{\omega_1, -\omega_1, 0\}$  and is disconnected, the result follows.

1-d Since  $\mathcal{E}_1^+ = \{0\}$  the result follows.

2. If  $u$  is nonnegative, then  $v \in L^1_{\mu_s}(\mathbb{R}^N)$ , thus (4.17) holds. If (4.22) holds, then

$$v \geq v_\infty = \lim_{k \rightarrow +\infty} v_k$$

by Proposition 4.16, which is impossible if  $w(t, \cdot)$  converges to 0. Therefore  $v \in L^p(B_1)$  and by Lemma 4.2 there exists  $k \geq 0$  such that (4.15) holds and therefore  $v = v_k$ . Furthermore, (4.16) is satisfied. If  $k = 0$ , then  $v$  is identically 0 and  $u$  is smooth as in the previous case.  $\square$

### 5.3 Singularities for Lane-Emden equation

In the case of Lane-Emden equation we have the following.

**Theorem 5.3** Assume that  $s \in (0, 1)$ ,  $\epsilon = -1$ ,  $p > 1$  and  $p \neq \frac{N+2s}{N-2s}$  and  $u \in C(\overline{\mathbb{R}_+^{N+1}} \setminus \{(0, 0)\}) \cap C^2(\mathbb{R}_+^{N+1})$  is a nonnegative solution of (2.2) satisfying (5.2).

1-a If  $p > \frac{N}{N-2s}$ , then  $w(t, \cdot)$  converges uniformly in  $\mathbb{S}_+^N$  either to  $\omega_1$  or to 0.

1-b If  $1 < p \leq \frac{N}{N-2s}$ , then  $w(t, \cdot)$  converges uniformly to 0 in  $\mathbb{S}_+^N$ .

2-a If  $p > \frac{N}{N-2s}$  and  $w(t, \cdot)$  converges to 0, then  $u$  is a smooth function in  $\overline{\mathbb{R}_+^{N+1}}$  and so is  $v$ .

2-b If  $1 < p < \frac{N}{N-2s}$ , there exists  $k \geq 0$  such that

$$e^{(N-\frac{2s}{p-1})t} v(t, \sigma', \phi) \rightarrow k(\sin \phi)^{2s} \quad \text{uniformly in } \mathbb{S}_+^N.$$

**Proof.** Assertions 1-a and 1-b follow from Theorem 5.1 and the description of  $\mathcal{E}^+$  given in Theorem 3.4. Furthermore, by Lemma 4.2 there exists  $k \geq 0$  such that (4.7) holds with  $\epsilon = -1$ . Observe that  $v(x) \geq kG_{s,\Omega}(x, 0)$ , where  $G_{s,\Omega}$  is the Green function of  $(-\Delta)^s$  in  $\Omega$  subject to zero Dirichlet condition in  $\Omega^c$ . If  $p \geq \frac{N}{N-2s}$ , then  $G_{s,\Omega} \notin L^p(B_1)$ , therefore  $k = 0$  in that case.

2-a. *Step 1:* If  $w(t, \cdot)$  converges to 0, we first prove that there exist  $\varepsilon > 0$  and  $c_{32} > 0$  such that

$$0 \leq w(t, \sigma) \leq c_{32}e^{\varepsilon t} \quad \text{for all } (t, \sigma) \in (-\infty] \times \overline{\mathbb{S}_+^N}. \quad (5.5)$$

Set  $\rho(t) = \|w(t, \cdot)\|_{C^0(\overline{\mathbb{S}_+^N})}$ . If (5.5) does not hold we have that

$$\limsup_{t \rightarrow -\infty} e^{-\varepsilon t} \rho(t) = +\infty \quad \text{for all } \varepsilon > 0.$$

By [14, Lemma 2.1] there exists a function  $\eta \in C^\infty((-\infty, 0])$  such that

$$\begin{aligned} (i) \quad & \eta > 0, \eta' > 0, \lim_{t \rightarrow -\infty} \eta(t) = 0, \\ (ii) \quad & 0 < \limsup_{t \rightarrow -\infty} \frac{\rho(t)}{\eta(t)} < +\infty; \\ (iii) \quad & \lim_{t \rightarrow -\infty} e^{-\varepsilon t} \eta(t) = +\infty \quad \text{for all } \varepsilon > 0; \\ (iv) \quad & \left(\frac{\eta'}{\eta}\right)', \left(\frac{\eta''}{\eta}\right)' \in L^1((-\infty, 0)); \\ (v) \quad & \lim_{t \rightarrow -\infty} \frac{\eta'(t)}{\eta(t)} = \lim_{t \rightarrow -\infty} \frac{\eta''(t)}{\eta(t)} = 0. \end{aligned} \quad (5.6)$$

We define  $\psi$  by  $w(t, \cdot) = \eta(t)\psi(t, \cdot)$ , then

$$\begin{aligned} \psi_{tt} + \varrho_1 \psi_t + \varrho_2 \psi + \mathcal{A}_s[\psi] &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}_+^N \\ \frac{\partial \psi}{\partial \nu^s} - \eta^{p-1} \psi^p &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}^{N-1}, \end{aligned} \quad (5.7)$$

where

$$\varrho_1 = \Theta_{s,p,N} + 2\frac{\eta'}{\eta} \quad \text{and} \quad \varrho_2 = \Lambda_{s,p,N} + \frac{\eta'}{\eta} \Theta_{s,p,N} + \frac{\eta''}{\eta}.$$

The function  $\psi$  is uniformly bounded. Then by standard regularity theory for degenerate elliptic operator (see e.g. [26]),  $\psi$  is bounded in  $C^2(-\infty, -1] \times \overline{\mathbb{S}_+^N}$ ). Therefore, the trajectory of

$\mathcal{T}_-[\psi] := \{\psi(t, \cdot) : t \leq 0\}$  is relatively compact in  $\mathbf{X}$  ( $= \{\phi \in C^s(\mathbb{S}_+^N) : \phi(\cdot, 0) \in C^2(\mathbb{S}^{N-1})\}$ ). By (5.6)-(ii) there exists a nonzero nonnegative function  $\theta \in C^2(\overline{\mathbb{S}_+^N})$  and a sequence  $\{t_n\}$  tending to  $-\infty$  such that  $\psi(t_n, \cdot) \rightarrow \theta$  in  $C^2(\overline{\mathbb{S}_+^N})$ . If we multiply the equation (5.7) that  $\psi$  satisfies by  $\psi_t$  and integrate over  $\mathbb{S}_+^N$ , we obtain the following energy formula

$$\frac{d}{dt}\mathcal{F}[\psi](t) = \varrho_1 \int_{\mathbb{S}_+^N} \psi_t^2 d\mu_s - \frac{1}{2}\varrho_2' \int_{\mathbb{S}_+^N} \psi^2 d\mu_s - \frac{p-1}{p+1}\eta'\eta^{p-2} \int_{\mathbb{S}^{N-1}} \psi^{p+1} dS', \quad (5.8)$$

where

$$\varrho_2' = \left(\frac{\eta'}{\eta}\right)' \Theta_{s,p,N} + \left(\frac{\eta''}{\eta}\right)'$$

and

$$\mathcal{F}[\psi](t) := \frac{1}{2} \int_{\mathbb{S}_+^N} (|\nabla' w|^2 - \varrho_2) \psi^2 d\mu_s - \frac{\eta^{p-1}}{p+1} \int_{\mathbb{S}^{N-1}} \psi^{p+1} dS'.$$

By properties (i), (iv) and (v) in (5.6), we have that

$$\eta'\eta^{p-2} \left( = \frac{\eta'}{\eta}\eta^{p-1} \right), \varrho_2', \frac{\eta'(t)}{\eta(t)} \in L^1((-\infty, 0)) \quad \text{as } t \rightarrow \infty.$$

Since  $\mathcal{F}[\psi](t)$  is bounded independently of  $t$ , it follows that

$$\int_{-\infty}^0 \int_{\mathbb{S}_+^N} \psi_t^2 d\mu_s dt < +\infty$$

using the regularity estimates it implies that

$$\lim_{t \rightarrow -\infty} \int_{\mathbb{S}_+^N} \psi_t^2 d\mu_s = 0.$$

This implies that the convergence of  $\psi(t_n, \cdot)$  to  $\theta$  holds in the stronger following way: for any  $T > 0$ ,

$$\psi(t, \cdot) \rightarrow \theta \quad \text{uniformly in } [t_n - T, t_n + T] \times \overline{\mathbb{S}_+^N}.$$

Let  $\zeta \in C^2(\overline{\mathbb{S}_+^N})$  such that  $\partial_{\nu^s} \zeta = 0$ , then from (5.7),

$$\begin{aligned} & \int_{\mathbb{S}_+^N} (w_t(T + t_n, \cdot) - w_t(T - t_n, \cdot)) \zeta d\mu_s + \Theta_{s,p,N} \int_{\mathbb{S}_+^N} (w(T + t_n, \cdot) - w(T - t_n, \cdot)) \zeta d\mu_s \\ & - \Lambda_{s,p,N} \int_{-T}^T \int_{\mathbb{S}_+^N} \psi(t_n + t, \cdot) \zeta d\mu_s dt - \int_{-T}^T \int_{\mathbb{S}_+^N} \psi(t_n + t, \cdot) \mathcal{A}_s[\zeta] d\mu_s dt \\ & = 2 \int_{-T}^T \int_{\mathbb{S}_+^N} \left( \frac{\eta'}{\eta} \psi_t(t_n + t, \cdot) + \varrho_2' \psi(t_n + t, \cdot) \right) \zeta d\mu_s dt + \int_{-T}^T \int_{\mathbb{S}^{N-1}} \eta^{p-1} \psi^p(t_n + t, \cdot) \zeta dS' dt. \end{aligned} \quad (5.9)$$

When  $t_n \rightarrow -\infty$ , the right-hand side of (5.9) tends to 0. Since  $\partial_{\nu^s} \theta = 0$  on  $\mathbb{S}^{N-1}$  from the equation satisfied by  $\psi$  and the convergence of  $\psi$ , it follows, as in the proof of Theorem 5.1, that  $\theta$  satisfies

$$\begin{aligned} \mathcal{A}_s[\theta] + \Lambda_{s,p,N} \theta &= 0 \quad \text{in } \mathbb{S}_+^N \\ \frac{\partial \theta}{\partial \nu^s} &= 0 \quad \text{in } \mathbb{S}^{N-1}. \end{aligned}$$

As  $\Lambda_{s,p,N} < 0$  we derive  $\theta = 0$  by integration, contradiction. Hence (5.5) holds for some  $\varepsilon > 0$ .

2-a. *Step 2.* Next we prove that there exist  $c_{33} > 0$  such that

$$w(t, \sigma) \leq c_{33} e^{\frac{2s}{p-1}t} \quad \text{for all } (t, \sigma) \in (-\infty, 0] \times \overline{\mathbb{S}_+^N}.$$

The set  $\mathcal{Q}$  of  $\varepsilon > 0$  such that (5.5) holds is a non-empty interval and it admits an upper bound  $\varepsilon^*$  which is finite if  $v$  is not identically 0. If  $\varepsilon \in \mathcal{Q}$ , we set  $\psi_\varepsilon(t, \cdot) = e^{-\varepsilon t} w(t, \cdot)$

If  $\varepsilon^* \in \mathcal{Q}$ , we take  $\varepsilon = \varepsilon^*$  and denote  $\psi_{\varepsilon^*} = \psi^*$ . Then  $\psi^*$  is bounded and satisfies

$$\begin{aligned} \psi_{tt}^* + \varrho_3^* \psi_t^* + \varrho_4^* \psi^* + \mathcal{A}_s[\psi^*] &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}_+^N \\ \frac{\partial \psi^*}{\partial \nu^s} - e^{(p-1)\varepsilon^* t} \psi^{*p} &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}^{N-1}, \end{aligned}$$

where

$$\varrho_3^* = \Theta_{s,p,N} + 2\varepsilon^* \quad \text{and} \quad \varrho_4^* = \Lambda_{s,p,N} + \varepsilon^* \Theta_{s,p,N} + \varepsilon^{*2}.$$

Because  $\psi^*$  is bounded, the same analysis as in the previous cases shows that any element  $\theta^*$  in the limit set of the trajectory of  $\{\psi^*\}$  is a nonnegative function satisfying

$$\varrho_4^* \theta^* + \mathcal{A}_s[\theta^*] = 0 \quad \text{in } \mathbb{S}_+^N, \quad \frac{\partial \theta^*}{\partial \nu^s} = 0 \quad \text{in } \mathbb{S}^{N-1}. \quad (5.10)$$

The polynomial  $P(z) = z^2 + \Theta_{s,p,N}z + \Lambda_{s,p,N}$  admits two roots  $z_2 < 0 < z_1$ , where

$$z_1 = \frac{2s}{p-1} \quad \text{and} \quad z_2 = 2s - N + \frac{2s}{p-1}.$$

Notice that  $0 < \varepsilon^* \leq z_1$  since  $v(0) > 0$ . If  $\varepsilon^* < z_1$ , the limit set is reduced to  $\{0\}$ , hence we proceed as in Step 1, by introducing a function  $\eta^*$  which has the same properties as the function  $\eta$  shown in (5.6), except that (5.6)-(ii) is replaced by

$$0 < \limsup_{t \rightarrow -\infty} \frac{\rho^*(t)}{\eta^*(t)} < +\infty \quad \text{with} \quad \rho^*(t) = \|\psi^*(t, \cdot)\|_{L^\infty(\mathbb{S}_+^N)}.$$

Defining the function  $v^* = \frac{w^*}{\eta^*}$  which satisfies

$$\begin{aligned} v_{tt}^* + \varrho_5^* v_t^* + \varrho_6^* v^* + \mathcal{A}[v^*] &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}_+^N \\ \frac{\partial v^*}{\partial \nu^s} - e^{(p-1)\varepsilon^* t} \eta^{*p-1} v^{*p} &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}^{N-1}, \end{aligned}$$

where

$$\varrho_5^* = \Theta_{s,p,N} + 2\varepsilon^* + 2\frac{\eta^{*'}}{\eta^*} \quad \text{and} \quad \varrho_6^* = \Lambda_{s,p,N} + \left(\varepsilon^* + \frac{\eta^{*'}}{\eta^*}\right) \Theta_{s,p,N} + \varepsilon^{*2} + \frac{\eta^{*''}}{\eta^*}.$$

The same analysis carried on in Step 1 shows that there exists a nonnegative nonzero element in the limit set of the trajectory of  $v^*$  which satisfies the same equation as  $\theta^*$  in (5.10), contradiction. Therefore  $\varepsilon^* = z_1$  which implies that  $v$  is bounded and therefore regular.

If  $\varepsilon^* \notin \mathcal{Q}$ , then for any  $\varepsilon < \varepsilon^*$  the function  $\psi_\varepsilon(t, \cdot) = e^{-\varepsilon t} w(t, \cdot)$  is bounded, tends to 0 when  $t \rightarrow -\infty$  and satisfies

$$\begin{aligned} \psi_\varepsilon tt + \varrho_3 \psi_\varepsilon t + \varrho_4 \psi_\varepsilon + \mathcal{A}_s[\psi_\varepsilon] &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}_+^N \\ \frac{\partial \psi_\varepsilon}{\partial \nu^s} - e^{(p-1)\varepsilon t} \psi_\varepsilon^p &= 0 \quad \text{in } (-\infty, 0) \times \mathbb{S}^{N-1}, \end{aligned}$$

where

$$\varrho_3 = \Theta_{s,p,N} + 2\varepsilon \quad \text{and} \quad \varrho_4 = \Lambda_{s,p,N} + \varepsilon \Theta_{s,p,N} + \varepsilon^2.$$

Put  $X_\varepsilon(t) = \int_{\mathbb{S}_+^N} \psi_\varepsilon(t, \cdot) d\mu_s$  and  $F_\varepsilon(t) = \int_{\mathbb{S}^{N-1}} \psi_\varepsilon^p dS'$ , then

$$X_\varepsilon'' + \varrho_3 X_\varepsilon' + \varrho_4 X_\varepsilon + e^{(p-1)\varepsilon t} F_\varepsilon = 0.$$

The equation

$$Y'' + \varrho_3 Y' + \varrho_4 Y = 0$$

admits the two linearly independent solutions

$$Y_1(t) = e^{\left(\frac{2s}{p-1} - \varepsilon\right)t} \quad \text{and} \quad Y_2(t) = e^{\left(2s-N + \frac{2s}{p-1} - \varepsilon\right)t}.$$

The function  $X_\varepsilon$  satisfies

$$-X_\varepsilon'' - \varrho_3 X_\varepsilon' - \varrho_4 X_\varepsilon \leq c_{34} e^{(p-1)\varepsilon t} \quad \text{on } (-\infty, 0].$$

Since it tends to 0 when  $t \rightarrow -\infty$ ,  $X_\varepsilon$  is bounded from above by the solution  $\tilde{X}_\varepsilon$  of

$$-\tilde{X}_\varepsilon'' - \varrho_3 \tilde{X}_\varepsilon' - \varrho_4 \tilde{X}_\varepsilon = c_{34} e^{(p-1)\varepsilon t} \quad \text{on } (-\infty, 0]$$

$$\tilde{X}_\varepsilon(0) = X_\varepsilon(0),$$

$$\lim_{t \rightarrow -\infty} \tilde{X}_\varepsilon(t) = 0.$$

The solution of this last differential equation is easily computable as we can always assume that  $(p-1)\varepsilon \neq \frac{2s}{p-1} - \varepsilon$  (the equality is called the resonance phenomenon) and we have

$$\tilde{X}_\varepsilon(t) = X_\varepsilon(0) e^{\left(\frac{2s}{p-1} - \varepsilon\right)t} + c_{35} e^{(p-1)\varepsilon t}$$

for some explicit constant  $c_{35}$  depending on the coefficients. This implies

$$\int_{\mathbb{S}_+^N} w(t, \cdot) d\mu_s \leq X_\varepsilon(0) e^{\frac{2s}{p-1}t} + c_{33} e^{p\varepsilon t}.$$

Applying Harnack inequality [36, Proposition 3.1] we obtain that

$$w(t, \cdot) \leq c_{36} e^{\frac{2s}{p-1}t} + c_{37} e^{p\varepsilon t} \quad \text{for all } t \leq T_0.$$

Because  $\varepsilon$  can be taken arbitrarily close to  $\varepsilon^*$ , we have that  $\min\{p\varepsilon, \frac{2s}{p-1}\} > \varepsilon^*$ , which contradicts the definition of  $\varepsilon^*$ . Hence  $\varepsilon^* \in \mathcal{Q}$ , which ends the proof of 2-a.

2-b. If  $1 < p < \frac{N}{N-2s}$ , there exists  $k \geq 0$  such that (4.7) holds, and

$$kG_{s,\Omega}(x, 0) \leq v(x) \leq c_{38}|x|^{2s-N},$$

where  $c_{38} > 0$ . □

**Proofs of Theorem E–F.** Theorem E is Theorem 5.1, Theorem F follows Theorem 5.2 part 2 and Theorem E does Theorem 5.3 part 1-a, 2-a. □

## 5.4 Classification of fractional Emden-Fowler equation

**Proof of Theorem G.** Let  $v$  be a nontrivial nonnegative solution of (1.9). It suffices to prove that

1. for  $p \in (0, 1 + \frac{2s}{N}]$ ,

$$v = v_k \quad \text{in } B_1 \setminus \{0\} \quad \text{for some } k \geq 0;$$

2. for  $p \in (1 + \frac{2s}{N}, \frac{N}{N-2s})$ , either

$$v = v_k \quad \text{in } B_1 \setminus \{0\} \quad \text{for some } k \geq 0$$

or

$$v = v_\infty \quad \text{in } B_1 \setminus \{0\}.$$

3. for  $p \geq \frac{N}{N-2s}$ ,  $v \equiv 0$ .

*Part 1.* For  $p \in (1, \frac{N}{N-2s})$ , if  $\int_\Omega v^p dx < +\infty$ , it follows by Lemma 4.2 that  $v = v_k$  for some  $k \geq 0$ . Moreover, if  $k = 0$ ,  $v = v_0 \equiv 0$  in  $\Omega \setminus \{0\}$ .

If  $\int_\Omega v^p dx = +\infty$ , from Proposition 4.16, we have that

$$v \geq v_\infty \quad \text{in } \mathbb{R}^N \setminus \{0\}.$$

If  $p \in (0, 1 + \frac{2s}{N}]$ , then  $v_\infty = +\infty$  in  $\Omega$ , which implies that  $v$  can be taken in this case.

*Part 2.* If  $p \in (1 + \frac{2s}{N}, \frac{N}{N-2s})$ ,  $v_\infty$  is a solution of (1.7) satisfying

$$\lim_{x \rightarrow 0} v_\infty(x)|x|^{\frac{2s}{p-1}} = c_p$$

by Corollary 4.19. Moreover, by the strong maximum principle, we have that

$$v \equiv v_\infty \quad \text{or} \quad v > v_\infty \quad \text{in } \Omega \setminus \{0\}.$$

If  $v > v_\infty$ , we let

$$V = \frac{v^p - v_\infty^p}{v - v_\infty}$$

and then

$$v_\infty^{p-1} \leq V \leq v^{p-1} \quad \text{in } \Omega \setminus \{0\}.$$



Let  $\bar{v} = v - v_\infty$  in  $\mathbb{R}^N \setminus \{0\}$ , then  $\bar{v}$  satisfies that

$$\lim_{x \rightarrow 0} \bar{v}(x) |x|^{\frac{2s}{p-1}} = 0$$

and

$$(-\Delta)^s \bar{v} + V \bar{v} = 0 \quad \text{in } \Omega \setminus \{0\}, \quad \bar{v} = 0 \quad \text{in } \Omega^c. \quad (5.11)$$

Note that  $\varepsilon v_\infty$  is a super solution of (5.11) for any  $\varepsilon > 0$  and by the comparison principle, we have that

$$\bar{v} \leq \varepsilon v_\infty \quad \text{in } \Omega.$$

By the arbitrary of  $\varepsilon > 0$ , we have that  $\bar{v} \leq 0$ . This is impossible from our setting that  $v > v_\infty$ . Therefore,  $v \equiv v_\infty$ .

*Part 3.* From Proposition 4.18 we have that

$$v(x) \leq c_{28} |x|^{-\frac{2s}{p-1}} \quad \text{in } \Omega \setminus \{0\}.$$

When  $p > \frac{N}{N-2s}$ , we have that  $-\frac{2s}{p-1} > 2s - N$  and  $\varepsilon |\cdot|^{2s-N}$  is a super solution of (5.11) for any  $\varepsilon > 0$ , then comparison principle implies that  $v \leq \varepsilon |\cdot|^{2s-N}$ , which by the arbitrary of  $\varepsilon > 0$ , implies that  $v = 0$ .

When  $p = \frac{N}{N-2s}$ , there holds  $-\frac{2s}{p-1} = 2s - N$ . For  $m \in \mathbb{R}$  let  $\nu_m$  be a smooth, radially symmetric function such that

$$\nu_m(x) = (-\ln |x|)^m \quad \text{for } 0 < |x| < \frac{1}{e^2} \quad \text{and} \quad \nu_m(x) = 0 \quad \text{for } |x| > 1. \quad (5.12)$$

Moreover,  $\nu_m$  is non-increasing in  $|x|$  if  $m > 0$ . Denote

$$w_m(x) = |x|^{2s-N} \nu_m(x) \quad \text{for } x \in \mathbb{R}^N \setminus \{0\}. \quad (5.13)$$

From [13, Proposition 2.1] for  $m \neq 0$ , we denote

$$\mathcal{I}_m = \mathcal{C}'_s(0)m, \quad \mathcal{J}_m = \mathcal{C}''_s(0) \frac{m(m-1)}{2}, \quad (5.14)$$

where  $\mathcal{C}'_s(0) < 0$  and  $\mathcal{C}''_s(0) < 0$ . Then  $\mathcal{I}_m > 0$  for  $m < 0$ , and there exist  $r_0 \in (0, \frac{1}{e^2}]$  and  $c_5 > 0$  such that for  $0 < |x| < r_0$

$$\left| (-\Delta)^s w_m(x) - \mathcal{I}_m |x|^{-N} (-\ln |x|)^{m-1} - \mathcal{J}_m |x|^{-N} (-\ln |x|)^{m-2} \right| \leq c_5 |x|^{-N} (-\ln |x|)^{m-3}.$$

Now we take  $m_0 = -\frac{N-2s}{2s} < 0$  and for some  $r_1 > 0$  there holds

$$(-\Delta)^s w_{m_0}(x) \geq 0 \quad \text{in } B_{r_1}(0) \setminus \{0\}.$$

Now we assume that  $\Omega \subset B_{r_1}$ , then  $\varepsilon w_{m_0}$  is a super solution of (5.11) for any  $\varepsilon > 0$  in  $B_{r_1}$ , then comparison principle implies that  $v \leq \varepsilon w_{m_0}$ , which by the arbitrary of  $\varepsilon > 0$ , implies that  $v = 0$ .

When  $\Omega$  is general domain, if  $v$  is a nontrivial singular solution of (5.11) in  $\Omega$ , then by the super and sub solution method, (5.11) in  $B_{r_1}$  has a solution  $\tilde{v}$  such that

$$\lim_{x \rightarrow 0} \frac{\tilde{v}}{v} = 1.$$

Then a contradiction could be obtained. □

## Open problems

**Problem 1.** Non-radial elements of  $\mathcal{E}_1$  can be formally constructed in the following way: set

$$\mathbb{S}_{+,1}^N = \mathbb{S}^N \cap \{(x_1, \dots, x_N, z) : x_1 > 0, z > 0\}.$$

Construct by minimization a positive function  $\omega$  satisfying

$$\begin{aligned} \mathcal{A}_s[\omega] + \Lambda_{s,p,N}\omega &= 0 && \text{in } \mathbb{S}_{+,1}^N \\ \frac{\partial \omega}{\partial \nu^s} + \omega^p &= 0 && \text{in } \mathbb{S}^{N-1} \cap \{x_1 > 0\} \\ \omega &= 0 && \text{in } \mathbb{S}_+^N \cap \{x_1 = 0\}. \end{aligned} \tag{5.15}$$

Then the function  $\tilde{\omega}$  defined in  $\mathbb{S}_+^N$  by

$$\tilde{\omega} = \begin{cases} \omega(\sigma) & \text{in } \mathbb{S}_{+,1}^N \\ -\omega(\tilde{\sigma}) & \text{in } \mathbb{S}_+^N \cap \{x_1 < 0\}, \end{cases}$$

where  $\tilde{\sigma}$  is the symmetric of  $\sigma$  with respect to the plane  $x_1 = 0$ , is a signed solution of (2.2) with  $\epsilon = 1$ . It vanishes on the plane  $\{(x, z) : x_1 = 0\}$ . It could be interesting to see if the condition for existence of such a solution coincides with the condition  $p < p^*$ .

**Problem 2.** It could be interesting to investigate the type of operators in  $\mathbb{R}_+^N$  which could lead by a trace process somehow similar to the one of [8] to perturbations of  $(-\Delta)^s$  such as  $(-\Delta + mI)^s$  or to Hardy type operators  $(-\Delta)^s + \mu|x|^{-2s}$ . This could be usefull for studying nonlinear semilinear equations involving fractional type-Laplacians.

**Problem 3.** The description of the singular behaviour in the case  $p = \frac{N}{N-2s}$  is more complicated due to the fact that  $\Lambda_{s,p,N} = 0$  and the linear blow-up rate  $2s - N$  coincides with the nonlinear blow-up rate  $\frac{2s}{p-1}$ . This phenomenon has been described by Aviles [2] in the case  $s = 1$ . We conjecture that the following result holds: either the function  $v$  can be extended as a smooth solution at 0, or

$$\lim_{x \rightarrow 0} |x|^{N-2s} (-\ln |x|)^{\frac{N-2s}{2s}} v(x) = C(N, s)$$

for some positive constant.

**Problem 4.** In an imporatnt article, Simon [33] studied the unique limit problem for bounded solutions of analytic functional. His results were used in [3] to obtained the precise asymptotics of solutions of

$$-\Delta u = \lambda e^u \quad \text{in } B_1 \setminus \{0\} \subset \mathbb{R}^3$$

with  $\lambda > 0$  satisfying  $|x|^2 e^{u(x)} \in L^\infty(B_{\frac{1}{2}})$ , or positive solutions of

$$-\Delta u = u^p + \frac{\ell}{|x|^2} u \quad \text{in } B_1 \setminus \{0\} \subset \mathbb{R}^N$$

with  $p > 1$  satisfying  $|x|^{\frac{2}{p-1}} \in L^\infty(B_{\frac{1}{2}})$ . The technique was to reduced the problem to a semilinear elliptic equations in an infinite cylinder and to prove that the limit set of the renormalized

corresponding solutions was reduced to a single element whatever is the structure of the set of possible limits (or self-similar solutions). It is natural to investigate similar questions concerning

$$\begin{aligned} u_{zz} + \frac{1-2s}{z}u_z + \Delta_x u + \frac{\ell}{|x|^2 + z^2}u &= 0 \quad \text{in } \mathbb{R}_+^{N+1}, \\ -\lim_{z \rightarrow 0} z^{1-2s}u_z(x, z) - u^p(x, 0) &= 0 \quad \text{in } B_1 \setminus \{0\} \subset \mathbb{R}^N. \end{aligned}$$

Assuming that a solution  $u$  satisfies  $(|x|^2 + z^2)^{\frac{s}{p-1}} \in L^\infty(B_{\frac{1}{2}})$  and using the variable (5.1) the new function  $w(t, \sigma) = r^{\frac{2s}{p-1}}u(r, \sigma)$ ,  $t = \ln r$ , satisfies

$$\begin{aligned} w_{tt} + \Theta_{s,p,N}w_t + (\Lambda_{s,p,N} + \ell)w + \mathcal{A}_s[w] &= 0 \quad \text{in } \mathbb{R} \times \mathbb{S}_+^N \\ \frac{\partial w}{\partial \nu^s} + w^p &= 0 \quad \text{in } \mathbb{R} \times \mathbb{S}^{N-1}, \end{aligned} \tag{5.16}$$

where  $\Theta_{s,p,N}$  and  $\Lambda_{s,p,N}$  are unchanged.

**Acknowledgements:** H. Chen is supported by the Natural Science Foundation of China, No. 12071189, by Jiangxi Province Science Fund, No. 20212ACB211005 and 20202ACBL201001.

## References

- [1] D. R. Adams, L. I. Hedberg, Function spaces and potential theory. *Grundlehren Math. Wissen. 314 Springer* (1996).
- [2] P. Aviles, Local behavior of solutions of some elliptic equations. *Comm. Math. Phys. 108*, 177–192 (1987).
- [3] M.F. Bidaut-Véron, L. Véron, Nonlinear elliptic equations on compact Riemannian manifolds and asymptotics of Emden equations. *Invent. Math. 106*, 489–539 (1991).
- [4] O. Boukarabila, L. Véron, Nonlinear boundary value problems relative to harmonic functions. *Nonlinear Analysis 201*, 112090 (2020).
- [5] H. Brezis, E. H. Lieb, Long range atomic potentials in Thomas-Fermi theory. *Comm. Math. Phys. 65*, 231–246 (1979).
- [6] H. Brezis, L. Véron, Removable Singularities for Some Nonlinear Elliptic Equations. *Arch. Rat. Mech. Anal. 75*, 1–6 (1981).
- [7] H. Brezis, P. L. Lions, A note on isolated singularities for linear elliptic equations. *Math. Anal. Appl. Part A. Adv. in Math. Supplementary studies Vol 7A*, 363–366 (1981).
- [8] L. Caffarelli, L. Silvestre, An extension problem related to the fractional Laplacian. *Comm. Part. Diff. Eq. 32*, 1245–1260 (2007).
- [9] L. Caffarelli, B. Gidas, J. Spruck, Asymptotic symmetry and local behavior of semilinear elliptic equations with critical Sobolev growth. *Comm. Pure Appl. Math. 42*, 271–297 (1989).
- [10] L. Caffarelli, T. Jin, Y. Sire, J. Xiong, Local analysis of solutions of fractional semi-linear elliptic equations with isolated singularities. *Arch. Rat. Mech. Anal. 213*, 245–268 (2014).

- [11] S. Chandrasekhar, Introduction to Stellar Structure. *University of Chicago Press* (1939).
- [12] H. Chen, P. Felmer, A. Quaas, Large solutions to elliptic equations involving fractional Laplacian. *Ann. Inst. H. Poincaré–AN* 32, 1199–1228 (2015).
- [13] H. Chen, H. Hajaiej, On isolated singularities for fractional Lane-Emden equation in the Serrin’s critical case, *arXiv:2109.07085* (2022).
- [14] X. Chen, H. Matano, L. Véron, Anisotropic Singularities of Solutions of Nonlinear Elliptic Equations in  $\mathbb{R}^2$ . *J. Funct. Anal.* 83, 50–97 (1989).
- [15] H. Chen, L. Véron, Semilinear fractional elliptic equations involving measures. *J. Diff. Eq.* 257, 1457–1486 (2014).
- [16] H. Chen, L. Véron, Weakly and strongly singular solutions of semilinear fractional elliptic equations. *Asymp. Anal.* 88, 165–184 (2014).
- [17] H. Chen, T. Weth, The Poisson problem for the fractional Hardy operator: distributional identities and singular solutions. *Trans. Amer. Math. Soc.* 374, 6881–6925 (2021).
- [18] H. Chen, J. Yang, Semilinear fractional elliptic equations with measures in unbounded domain. *Nonlinear Anal.* 145, 118–142 (2016).
- [19] M. Fall, Semilinear elliptic equations for the fractional Laplacian with Hardy potential. *Nonlinear Analysis* 193, 111311 (2020).
- [20] Y. Fang, D. Tang, Method of sub-super solutions for fractional elliptic equations. *Discr. Cont. Dyn. Sys.* 23, 3153–3165 (2018).
- [21] P. Felmer, A. Quass, J. Tan, Positive solutions of the nonlinear Schrödinger equation with the fractional Laplacian. *Proceed. Royal Society of Edinburgh*, 142A, 1237–1262 (2012).
- [22] N. Garofalo, Fractional thoughts in New Developments in the Analysis of Nonlocal Operators. *Contemp. Math.* 723, 1–136 (2019).
- [23] B. Gidas, W. Ni, L. Nirenberg, Symmetry and related properties via the maximum principle. *Comm. Math. Phys.* 68, 209–243 (1979).
- [24] B. Gidas, J. Spruck, A priori bounds for positive solutions of nonlinear elliptic equations. *Comm. Part. Diff. Eq.* 6, 883–901 (1981).
- [25] B. Gidas, J. Spruck, Global and local behaviour of positive solutions of nonlinear elliptic equations. *Comm. Pure Appl. Math.* 34, 525–598 (1981).
- [26] E. Fabes, C. Kenig, R. Serapioni, The local regularity of solutions of degenerate elliptic equations. *Comm. Part. Diff. Eq.* 7, 77–116 (1982).
- [27] R. Fowler, Further studies on Emden’s and similar differential equations. *Quart. J. Math.* 2, 259–288 (1931).
- [28] T. Jin, Y.Y. Li, J. Xiong, On a fractional Nirenberg problem, part I: blow up analysis and compactness of solutions. *J. Eur. Math. Soc.* 16, 1111–1171 (2014).
- [29] C. Li, Z. Wu, H. Xu, Maximum principles and Bôcher type theorems. *Proc Natl Acad Sci U S A.* 115, 6976–6979 (2018).
- [30] P. L. Lions, Isolated singularities in semilinear problems. *J. Diff. Eq.* 38, 441–450 (1980).

- [31] M. Marcus, L. Véron, Uniqueness and asymptotic behaviour of solutions with boundary blow-up for a class of nonlinear elliptic equations. *Ann. Inst. H. Poincaré-AN* 14, 237–274 (1997).
- [32] R. Menon, Integro-differential operators: connections to degenerate elliptic equations and some free boundary problems. *PhD Thesis, Michigan State University* (2020).
- [33] L. Simon, Isolated singularities of extrema of geometric variational problems, In: Giusti E. (ed.) Harmonic Mappings and Minimal Immersions. *Lect. Notes Math.* 1161, 206–277, Berlin, Heidelberg, New-York (1985).
- [34] A. Sommerfeld, Asymptotische integration der differential-gleichung des Thomas-Fermischen atoms. *Zeitschrift für Physik* 78, 283–308 (1932).
- [35] L. Véron, Singular solutions of some nonlinear elliptic equations. *Nonlinear Analysis, Th. Meth. Appl.* 5, 225–242 (1981),
- [36] H. Yang, W. Zou, On isolated singularities of fractional semi-linear elliptic equations. *Ann. Inst. H. Poincaré-AN.* 38, 403–420 (2021).