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Comparative kinetic analysis of redox flow battery electrolytes: from micro-fibers to macro-felts

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Abstract

In the context of new aqueous organic electrolytes development for redox flow batteries, two different approaches of charge transfer kinetic constant determination on carbon-based felts are compared. The first one consists in direct electrochemical measurement on centimetric pieces of felts, referred to as macro-felts, and the second is based on electrochemical measurement on carbon fibers extracted from felts and conditioned as microelectrodes, referred to as micro-fibers. Quantitative cyclic voltammetry (CV) involving experimental data fitting or simulation and electrochemical impedance spectroscopy (EIS) are carried out on the two supports, with two typical aqueous organic electrolytes, and determined values for charge transfer kinetic constants are compared. k^0 values analyzed from micro-fibers are systematically lower than that analyzed from macro-felts where the differences may reach *ca.* two orders of magnitude in some cases. CV and EIS analysis on micro-fibers show very good consistency. With the simplicity of physical analysis, micro-fibers are considered as a reliable and powerful tool for kinetic determination, especially when confronted to fast kinetics.

Keywords: Carbon felt – Microelectrodes – EIS – CV – Batteries

1. Introduction

Redox flow batteries are a promising energy storage technology that may play an important role in the integration of intermittent renewable energies in grid networks. Typical electrodes used in redox flow batteries are porous materials that offer high specific surface area. Carbon-based felts are commonly used as electrode material in redox flow batteries due to their high specific surface area, high conductivity, resistance to corrosion and wide operation potential range in addition to their relatively low cost [1].

Although carbon felts allow easy permeation of electrolyte, yielding high current density in practical applications, their intrinsic electrochemical behavior is not easy to analyze. The apparent electrode performance depends not only on the properties of the electrode material but also on their geometry which is often not well defined for porous felts. Most electroanalytical methods are based on regular geometry (planar, cylindrical, spherical, micro-disk, *etc.*), and model equations are usually developed for planar electrodes. In the case of porous electrodes, the mass transfer within their structure shall be considered, which greatly complicates the mathematical analysis. The high surface area also makes the differentiation of apparent and intrinsic charge transfer kinetics difficult, especially when the latter could also be very fast in most energy storage systems. On the other hand, for guiding the material development, it is always of high interest to assess the intrinsic charge transfer kinetics of porous carbon felts.

Cyclic voltammetry (CV) and electrochemical impedance spectroscopy (EIS) are the most common electrochemical methods to determine charge transfer kinetics. The CV analysis for planar electrodes is highly matured [2] and one may derive kinetic information by routinely

fitting of experimental data with appropriate models using commercial software. For analyzing porous electrodes, the geometry needs to be properly approximated. One approach is considering arrays of electrode units, with the geometry chosen according to the type of porous material studied [3],[4]. Very recently, Tichter and co-authors improved this formalism by representing porous electrodes as randomly distributed arrays of microelectrodes for wider variety of geometries [5], [6]; they developed a software, Polarographica, for simulating CV of porous electrodes.

Considerable work has also been made to analyze kinetics from EIS of porous electrodes. DeLevie is often referred as pioneer of the field in the 1970s [7], [8] with the introduction of transmission line analogy. Following this work, Candy and co-authors studied gold powders' impedance in KOH [9], [10] and obtained good fit from numerical calculations considering their catalyst as an array of identical cylindrical pores. In the 1990s, Lasia established a similar numerical model including a potential gradient within a porous electrode and later added a concentration gradient [11]–[13]. Paasch and co-authors developed a “macrohomogeneous material” model [14], [15], and similar approaches were further complexified in cases of problems related to Lithium rechargeable batteries [16], [17].

As the challenge of directly analyzing porous felts from both CV and EIS mainly arises from the undefined geometry, an alternative way is to extract the fibers from the felt and study them as microelectrodes. As the diameter of fiber is in the range of a few microns while the length of fiber is usually in millimeter range, the contribution of the tip can be neglected and the electrode can be regarded as a 1D cylindrical system which eases the mathematical treatment. Kovach *et al.* used such electrodes with various electrochemical systems and

analyzed in particular their results by fitting experimental CV [18]. Bourke *et al.*[19] and Miller *et al.*[20] carried out CV and EIS on carbon fibers to measure the effect of anodic and cathodic treatments on charge transfer kinetics with typical vanadium electrolytes. Recently, Landon-Lane *et al.* carried out EIS on carbon fiber microelectrodes with ferrocyanide-ferricyanide and proposed a simplified approach to fit data for cylindrical geometry [21]. The author further developed the use of microelectrodes by fitting voltammograms obtained with both ferrocyanide-ferricyanide and vanadium electrolytes [22], minimizing the difference in peak potential separation between experimental and simulated data, conforming to methods suggested by other authors [23], [24].

In this work, we carried out comprehensive studies of intrinsic charge transfer kinetics of two typical electrolytes promoted in the development of aqueous organic electrolyte redox flow batteries on activated and non-activated carbon felts: ammonium ferrocyanide and a derivative of TEMPO. The analysis was achieved by fitting CV and EIS data on both macro-felts and extracted micro-fibers and the results were compared to identify the best experimental and modelling approach to quantify charge transfer kinetics of redox flow electrolyte on such porous material.

2. Materials and Methods

2.1 Chemical products

Two aqueous organic electrolytes are presented in this study: ammonium hexacyanoferrate, received as pure ammonium ferrocyanide $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$, and a derivative of TEMPO (one of the cycle's hydrogen is substituted with an organic group), further referred to as “TEMPO-based molecule”, both provided by CIC energiGUNE. Solutions were prepared by dissolving the right amount of salt in its corresponding supporting electrolyte: 1 M NaCl for TEMPO-based molecule, prepared from pure NaCl (Prolabo, France), and 1 M NH_4Cl for ammonium hexacyanoferrate, prepared from pure 99.5% NH_4Cl (Fisher Scientific, USA). The concentrations of the studied solutions were 10^{-4} M for “macro-felts” and 10^{-2} M for microelectrodes.

2.2. Experimental apparatus and protocols for macro-felts

All electrochemical experiments were performed with a VSP-300 potentiostat (Bio-Logic, France) interfaced by EC-Lab[®] software (Bio-Logic, France).

For “macro-felts”: a custom-made PVC cell to perform cyclic voltammetry and electrochemical impedance spectroscopy on pieces of graphite felts was designed and built (see **Figure S1**). Graphite felts were used as electrodes (SGL 4.6 GFD, thickness of 4.6 mm): two large cylindrical pieces (\varnothing 5 cm) served as counter electrodes (CE) while one smaller cylindrical piece (\varnothing 1 cm) was the working electrode (WE). The felts were cut with punches for better reproducibility. Both commercially available non-activated (n.a.) and activated versions of this felt were tested. The electrodes were positioned inside the cell using 3D-printed frames so that the WE was centered between the two CE and set so their respective

centers were aligned with one another on a horizontal line. A 3D-representation in **Figure S1** shows the positioning of the electrodes. Graphite felt electrodes were connected to the leads of the potentiostat by piercing through each of them with a composite rectangular carbon rod pultruded (DPP™, R&G Faserverbundwerkstoffe GmbH) attached to banana connectors with copper wire and silver-epoxy resin (silver conductive epoxy adhesive, MG Chemicals). The reference electrode was chosen to be an Ag/AgCl quasi-reference electrode (QRE) in order to mitigate interference with the measurements. QRE was made by coating a Ø 0.25 mm silver wire (99.9%, chemPUR Feinchemikalien und Forschungsbedarf GmbH) with AgCl by applying 1 V between the wire and a large piece of graphite felt as the counter electrode in 1 M HCl for one minute. This wire was then inserted in the cell via one of the frames to avoid contact with any of the other electrodes. The cell was closed with a 3D-printed lid. Before each experiment, felt cylinders were hand pressed in the studied electrolytic solution in the manner of a sponge in order to expel the air contained in the pores and to fill them with electroactive solution, as non-activated materials have poor water wettability.

EIS experiments for graphite felts were carried out applying the following protocol: first, the formal potential of studied redox couple was estimated from cyclic voltammetry by taking the average of oxidative and reductive peak potentials. Then, staircase potential electrochemical impedance spectroscopy (SPEIS) was performed with EC-Lab® software in a range of potentials centered on formal potential and including it: at each potential step, potential is applied for enough time to reach pseudo-equilibrium (typically ten minutes), then oscillates around its pseudo-equilibrium value while measuring the response current at different frequencies. For ferrocyanide, 50 mV potential steps and five minutes of buffer time were used. For TEMPO-based compound, 20 mV potential steps were applied using a buffer time of two minutes. In both cases, oscillations amplitude was set at 10 mV. Screening different

potentials allowed us to verify that a minimum impedance indeed occurred at formal potential for our two electrolytes, therefore validating the assumption of a Nernstian system. Performing the experiments around formal potential rather than equilibrium potential (or open-circuit potential), as often found in literature, was motivated by the fact that the TEMPO-based molecule received was only present in its reduced form and was extremely stable: it was thus impossible to use i-E equations for analysis as the oxidant surface concentration was unknown, and could not be approximated to 0 since the system's kinetics were very fast compared to mass transfer. Frequencies screened ranged from 7 MHz to 50 mHz, with six frequencies screened per decade, and for a total impedance measurement duration of 2.5 minutes.

The analysis of EIS spectra was done by performing Complex Non-linear Least Squares, that is to say by fitting the experimental Nyquist diagram with a model. The models used are detailed in a subsequent part. The optimization was done by programming a small custom-made routine with MATLAB® software based on the use of the optimization function “fmincon”. Constraint tolerance, function tolerance and optimality tolerance were all set to 10^{-8} . The maximum number of function evaluations was set to 200 and algorithm used was “sqp”. Before optimization, the parameters were weighted with their approximate order of magnitude to ease convergence. The relevant range of frequencies usable for analysis was typically below 5 kHz for graphite felts.

BET surface of graphite felts was measured using 3 g of graphite felt. The reproduction of the experiment showed a variation around 10% of estimated value.

2.3. Experimental apparatus and protocols for microelectrodes

For microelectrodes: the cell used to perform EIS and CV experiments was made out of a simple flat-bottomed polystyrene sampling tube ($\text{\O} 3.5 \text{ cm}$) closed with a 3D-printed lid. The lid allowed to mount and fix the position of an Ag/AgCl QRE (prepared similarly to the one used with graphite felts), a counter electrode made from a DPPTM composite rectangular carbon rod pultruded and a custom-made 3D-printed support holding the carbon fiber microelectrode studied. The 3D-printed support for fiber microelectrodes consists in two parts: part A, longer, serves to allocate a carbon fiber and a copper wire ($\text{\O} 0.2 \text{ mm}$), and part B, shorter, seals the electrical contact fiber-copper and completes the support. The groove in part A is filled with silver-epoxy resin (Silver conductive epoxy adhesive, MG Chemicals). From one side, copper wire is inserted into the guiding hole until it reaches the Ag-epoxy. On the other side, carbon fiber is positioned on the groove ensuring the good contact with the copper wire. The piece is cured at 65°C for 2 h. Then, the exposed electrical contact is sealed by gluing part B and part A using UV-curing resin. The microelectrodes prepared as described show excellent chemical and mechanical stability over time in the tested electrolytes. Photograph as well as an exploded view are available in **Figure S2** for better visualization. Before each experiment, microelectrodes were cleaned by immersing the tip of the support in a 70% ethanol solution for ten minutes in order to remove possible presence of dust or other impurities on the surface of the fibers.

Carbon fibers' mean radius was estimated after observation of several carbon fibers with scanning electron microscopy (JSM-IT500HR, JEOL). They were found to range typically between 8 and 10 μm , thus an average diameter of 9 μm was chosen, therefore giving a mean radius of 4.5 μm . Carbon fibers' lengths were estimated by optical inspection under the microscope, with an estimated error of $\pm 0.25 \text{ mm}$.

The procedure applied for EIS experiments with carbon fiber microelectrodes was the same as the one with graphite felts, except for ammonium hexacyanoferrate on non-activated fiber: in this particular case, the kinetics were found very slow and an impedance minimum could no longer be found at formal potential, instead appearing at more oxidative potential. In this case only, a wider potential range was scanned, and the analysis was focused around the potential giving minimum impedance values. This choice was motivated, first by the impedance spectra shape at these potentials being more consistent with the one theoretically obtainable with the model used, and second by failure to obtain good fit with said model at equilibrium potential. In addition, due to the very small surface of microelectrodes, frequency could be analyzed up to very high frequencies, typically up to 200 kHz.

All CV and EIS experiments were triplicated for reproducibility.

3. Mathematical modelling

3.1 “Macro-felts” porous electrodes modelling

3.1.1. Cyclic voltammetry model for porous electrodes

In order to estimate kinetic constants based on cyclic voltammograms measured on porous macro-felts, the model developed by Tichter and co-authors was used. This model, available under the form of a free software called Polarographica, is based on the representation of porous materials as randomly spaced arrays of microelectrodes. The spacing of microelectrodes is handled with a statistical distribution that can be divided into several intervals of inter-electrode distances. Cyclic voltammograms are then calculated for each of these intervals, each time based on corresponding average distance, and their sum is weighted according to the statistical distribution to give the resulting voltammogram. Among the different geometries available, the array of cylindrical electrodes is the one that should best represent carbon and graphite felts. Although currently unable to automatically fit curves, it is possible to import experimental data and to manually change input parameters so for a simulated curve to approach the experimental one.

Using version 2.1.0 of this software with “statistically weighted cylindrical, cylindrical finite diffusion model”, CV curves were simulated with typical parameters determined by the authors and obtained with similar felts [6]. Hence, the distribution of inter-electrode distances was not changed from “default parameters”, and intervals of distribution as well as the number of cylinders per mm^2 were set to the same values as the ones used by the authors with the same type of felt (respectively 722 cylinders per mm^2 and 10 intervals). Then, the total surface was adjusted so for the peak magnitude to match the experimental one and the kinetic

constant was adjusted in order to find the minimum value allowing the peak-to-peak distance to match the experimental one.

3.1.2 Macrohomogeneous model for graphite felts EIS

For analysis of the impedance spectra with graphite felts, the macrohomogeneous model developed by Paasch, Micka and Gersdorf [15] was used, in addition to some minor modifications. This model assumes the porous electrode as a biphasic material, composed of a liquid phase (the electrolyte inside the pores) and a solid phase (the actual electrode). The porous electrode is in contact on one side with the bulk solution and on the other side by a metal connector. This model is one-dimensional: physical variables are assumed constant over one “slice” of the porous element. The concentration of active species is assumed constant through the whole electrode, but the model allows the addition of a “hindrance function” in order to represent for instance the influence of diffusion phenomena. Finally, double-layer capacitance per length of porous material is also assumed constant. This model considers the separation of the total electric current circulating through the electrode to the part circulating in solid phase and to the part circulating in the liquid phase. Their respective value changes with the axial position in the porous electrode, but their sum stays constant by virtue of charge continuity. A gradient of potential with position in solid and liquid phase therefore arises.

This model was originally derived relative to equilibrium potential, but it can be shown that the equation they established is basically identical when formal potential is used instead of equilibrium potential (development in supplementary material). The main difference is situated in the expression of parameter k . The complex impedance equation is given in

equation (1), the parameters are further explained in the following equations and symbols are given in **appendix**:

$$Z(\omega) = R_{ohm} + \frac{1}{A} \left[\frac{\rho_1^2 + \rho_2^2}{\rho_1 + \rho_2} \cdot \frac{\coth(d\beta)}{\beta} + \frac{2\rho_1\rho_2}{\rho_1 + \rho_2} \cdot \frac{1}{\beta \sinh(d\beta)} + \frac{d\rho_1\rho_2}{\rho_1 + \rho_2} \right] \quad (1)$$

$$\beta = \left(\frac{ky(\omega) + (j\omega)^\gamma}{K} \right) \quad (2)$$

$$K = \frac{1}{(\rho_1 + \rho_2)CS_c} \quad (3)$$

$$k = \frac{n^2 F^2 k^0 (C_O^* + C_R^*) e^{\frac{(1-\alpha)nF}{R_g T}(\bar{E} - E^{0'})}}{R_g T C \left(1 + e^{\frac{nF}{R_g T}(\bar{E} - E^{0'})} \right)} \quad (4)$$

$$y(\omega) = \frac{1}{1 + \left(\sqrt{\frac{\omega_2}{j\omega}} \right)^\phi \coth \left(\left(\sqrt{\frac{j\omega}{\omega_3}} \right)^\phi \right)} \quad (5)$$

$$\omega_2 = \frac{k^0}{\sqrt{D}} \left[\exp \left(-\frac{\alpha nF}{R_g T} (\bar{E} - E^{0'}) \right) + \exp \left(\frac{(1-\alpha)nF}{R_g T} (\bar{E} - E^{0'}) \right) \right] \quad (6)$$

$$\omega_3 = \frac{D}{l_p^2} \quad (7)$$

R_{ohm} is the ohmic resistance between the working electrode and the reference electrode. C_O^s and C_R^s are the surface concentrations of respectively the oxidant and the reductant of considered redox couple; studied salts were assumed to have a Nernstian behavior on graphite felts, and therefore these values were obtained through Nernst equation as follows:

$$C_O^s = (C_O^* + C_R^*) \cdot \frac{\exp\left(\frac{nF}{R_g T} (\bar{E} - E^{o'})\right)}{1 + \exp\left(\frac{nF}{R_g T} (\bar{E} - E^{o'})\right)} \quad (8)$$

$$C_R^s = (C_O^* + C_R^*) \cdot \frac{1}{1 + \exp\left(\frac{nF}{R_g T} (\bar{E} - E^{o'})\right)} \quad (9)$$

n is the number of electrons involved in the electrochemical process (equal to 1 for both salts), F is the Faraday's constant, R_g is the ideal gas constant, T the room temperature (298.15 K). d is the electrode thickness, equal to 4.6 mm for all cases studied her. ρ_1 and ρ_2 are the resistivities of respectively the solid phase (the actual electrode material) and the liquid phase (the electrolyte) within the porous material. ρ_2 is modified from the bulk electrolyte resistivity ρ_2^* due to the porosity ϵ and tortuosity f_p of the material according to equation (10):

$$\rho_2 = \rho_2^* \cdot \frac{f_p}{\epsilon} \quad (10)$$

k^0 is the intrinsic charge transfer constant while S_c is the specific surface area of the porous electrode, that is to say its total reactive surface relative to its geometric volume. A is the geometric area of the porous electrode cross section.

Parameter β was modified to include a constant phase element (CPE) instead of a capacitor to represent the double-layer of the porous material. This can be done simply by adding an exponent as shown in equation (2). The capacitance C in parameter k and K then becomes a pseudo-capacitance.

The mass transfer “hindrance function” was chosen to represent finite diffusion with reflective boundary. The finite diffusion reflective boundary, as reported by MacDonald [25],

corresponds to solving the diffusion equation with a specific boundary condition. The use of this boundary rather than a transmissive boundary was mainly motivated by the shape of our impedance spectra that tended towards infinite values at low frequencies for the imaginary part of the impedance. In addition, a modification to the finite diffusion hindrance function was performed in order to correctly represent our data at low frequencies: a dimensionless exponent ϕ was added to represent “non-uniform diffusion” [25]. This modification allows the slope in the Nyquist diagram to tend towards a finite value instead of infinity, which was observed in particular for ammonium hexacyanoferrate. A similar modification was performed by Sun *et al.* [26], although the parameter introduced was a factor rather than an exponent. The values of fixed physical parameters taken as input in the model are given in **Table 1**, parameters fitted were: R_{ohm} , C , γ , k^0 , l_p , S_c , f_p , ϕ .

The CPE exponent was not always needed for the convergence to occur, and it was fixed to 1 unless convergence was impossible without it. Failed convergence would typically manifest in the routine not being able to find a value for the kinetic constant, and constantly searching for a higher value, stopping at the upper boundary. Increasing the value of this upper boundary would result in the model reaching a higher value for k^0 while leaving the other parameters virtually unchanged. In these cases, the value found for ϕ during the failed convergence was kept and fixed, and γ was then optimized along with the other parameters. This procedure always allowed reaching convergence. It was preferred over an optimization with eight free parameters which would have been more difficult to realize: indeed, multiple optima were typically available with only seven free parameters as will be explained in the discussion part.

3.2 Carbon fiber microelectrodes modelling

3.2.1. Fitting of microelectrodes' cyclic voltammograms

The representation of cyclic voltammograms was done according to the method presented by Compton, Laborda and Ward [2]: the diffusion equation in cylindrical geometry (**equation 11**) was numerically solved using finite element methods with backward implicit method following Thomas' algorithm, as described in the authors' book.

$$\frac{\partial C_R}{\partial t} = D \left(\frac{\partial^2 C_R}{\partial r^2} + \frac{1}{r} \cdot \frac{\partial C_R}{\partial r} \right) \quad (11)$$

Diffusion coefficient of reductant and oxidant were assumed to be of same value. An expanding grid was used for radial dimension. The first spatial boundary condition was equality of concentration with bulk concentration at maximum distance from the electrode:

$$(C_R)_{r=r_{max}} = C_R^* \quad (12)$$

The second spatial boundary condition was equality of concentration flux at the electrode's surface with consumption flux given by extended Butler-Volmer equation, as presented in equation (13):

$$D \left(\frac{\partial C_R}{\partial r} \right)_{r=r_0} = k^0 \left[(C_R)_{r=r_0} \exp \left(\frac{(1-\alpha)nF}{R_g T} (\bar{E} - E^{0'}) \right) - (C_R^* + C_O^* - (C_R)_{r=r_0}) \exp \left(-\frac{\alpha nF}{R_g T} (\bar{E} - E^{0'}) \right) \right] \quad (13)$$

The time boundary condition is given by equality of concentrations with bulk concentration at initial time.

$$(C_R)_{t=0} = C_R^* \quad (14)$$

Once the concentration profile is plotted along the spatial grid at different wanted times, the current profile can be deduced by the surface gradient according to equation (15):

$$I = nFAD \left(\frac{\partial C_R}{\partial r} \right)_{r=r_0} \quad (15)$$

The fitting of cyclic voltammetry experimental data was done as follows: experimental time array and measured potential array were taken as input variables, and corresponding response current was determined as a function of parameters to optimize. Then, typically the third cycle was used for the fitting: the calculated current points at corresponding input potentials were compared to experimental current points with least square minimization method. Optimized parameters were charge transfer kinetic constant k^0 , as well as geometric surface area A in order to include the measurement error of fibers' lengths as well as the variability of their radii. The whole optimization procedure was coded with MATLAB®.

3.2.2. Impedance model in cylindrical geometry for carbon fibers microelectrodes EIS

Carbon fibers microelectrodes can be reasonably represented as cylindrical electrodes and impedance at their surface can then simply be modelled with a Randles circuit. However, semi-infinite diffusion in cylindrical geometry leads to a Warburg impedance equation involving complicated Bessel functions [27]. Landon-Lane and co-authors mentioned problems in convergence when using this expression and suggested the use of a simplified equation with errors in real and imaginary parts not exceeding 4% of their original values [21]. The total impedance, as written by aforementioned authors, is presented in equation (16) and (17):

$$Z = R_{ohm} + \frac{1}{A \left(Y_0(j\omega)^\gamma + \frac{1}{R_{ct} \left(1 + \frac{2k^0 r_0}{D} \ln \left(\frac{1 + r_0 \sqrt{\frac{j\omega}{D}}}{r_0 \sqrt{\frac{j\omega}{D}}} \right) \right)} \right)} \quad (16)$$

$$R_{ct} = \frac{R_g T}{n^2 F^2 A k^0 \left[\alpha C_O^s e^{-\alpha \frac{nF}{RT}(\bar{E} - E^{o'})} + (1 - \alpha) C_R^s e^{(1-\alpha) \frac{nF}{RT}(\bar{E} - E^{o'})} \right]} \quad (17)$$

Most symbols refer to the same physical parameters than the ones used for the macrohomogeneous model, and will therefore not be detailed here. The only exception is the CPE pseudocapacitance Y_0 .

In most cases, the assumption of a Nernstian system was verified (impedance minimum at estimated formal potential, and reversible-shaped cyclic voltammogram) and surface concentrations were obtained using Nernst equation with **equations (8) and (9)**. However, in the case of non-activated fibres in $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$, that is for $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$ with non-activated fibers, the system did not appear Nernstian but quasi-reversible, based on the shape of the voltammograms and the impedance responses. Thus, kinetics became too slow for surface concentrations to be approximated with values given by Nernst equation, and mass transfer had to be considered to deduce their values.

In order to do so without introducing excessive complexity, steady-state and finite diffusion were assumed, semi-infinite diffusion assumption leading to no simple solution for diffusion equation in cylindrical geometry. After a development, presented in supplementary material,

surface concentrations can be expressed as a function of mass transfer coefficient m_O (assumed equal for oxidant and reductant):

$$C_O^s = \frac{C_O^* + (C_R^* + C_O^*) \frac{k_0}{m_O} e^{\frac{(1-\alpha)nF}{RgT}(\bar{E}-E^{0'})}}{1 + \frac{k_0}{m_O} \left[e^{-\frac{\alpha nF}{RgT}(\bar{E}-E^{0'})} + e^{\frac{(1-\alpha)nF}{RgT}(\bar{E}-E^{0'})} \right]} \quad (18)$$

$$C_R^s = \frac{C_R^* + (C_R^* + C_O^*) \frac{k_0}{m_O} e^{-\frac{\alpha nF}{RgT}(\bar{E}-E^{0'})}}{1 + \frac{k_0}{m_O} \left[e^{-\frac{\alpha nF}{RgT}(\bar{E}-E^{0'})} + e^{\frac{(1-\alpha)nF}{RgT}(\bar{E}-E^{0'})} \right]} \quad (19)$$

Parameters fitted were: R_{ohm} , Y_0 , γ , k^0 and D . m_O was also fitted in the case where surface concentrations could not be approximated to Nernst concentrations.

4. Results and discussion

4.1. Experimental results with porous “macro-felts”

4.1.1. Cyclic voltammetry on porous “macro-felts”

Typical cyclic voltammograms on graphite felts are presented for both studied salts in **Figure 1**, where the potential distance between oxidative and reductive peaks appear narrower with TEMPO-based molecule than with ferrocyanide. However, much dispersion was observed, especially with ferrocyanide, and distance between peaks could range from 20 to 100 mV with the same molecule (**Figure S3**). This dispersion is attributed to felts inhomogeneities. The height of the peaks, *i.e.* the peak currents, although similarly subject to some dispersion, was reproducibly different for the two molecules. Current peaks in TEMPO-based molecules are about 1.3 times higher in intensity. On the other hand, a much higher non-faradaic current is observable on the CV plotted with activated felt, which corresponds to

a higher amount of species adsorbed. This is well correlated by the increase in specific surface area measured by BET method: indeed, specific surfaces measured were about $0.5 \text{ m}^2 \text{ g}^{-1}$ for non-activated felt and $1.4 \text{ m}^2 \text{ g}^{-1}$ for activated one (about $500 \text{ cm}^2 \text{ cm}^{-3}$ and $1400 \text{ cm}^2 \text{ cm}^{-3}$). Care should be given however to the meaning of these values, as these correspond to values referring to N_2 adsorption, and might represent a surface not readily available to bigger molecules used in this study.

Using Polarographica software, manual fits of the cyclic voltammograms of the two systems on non-activated felts are presented in **Figure 2**, estimated values for k^0 and total surface area A_{total} are displayed below the graphs. The fit reproduces quite well the experimental data, especially around formal potential. Far from this value, important deviations are observed, which may arise from the non-faradaic current that is not considered in the model. Therefore, no fit was attempted for activated felts. Furthermore, the kinetic constants estimated here are actually a “minimum possible value” as the curves do not change shape anymore when inputting higher k^0 values. Total surface areas are different for the two systems and in agreement with observed differences in current peak values. The results of 45 cm^2 for ammonium hexacyanoferrate and 60 cm^2 for TEMPO derivative can be converted in specific surface areas using working electrode dimensions: a 1 cm diameter cylinder with length of 4.6 mm yields specific surface areas of $122 \text{ cm}^2 \text{ cm}^{-3}$ and $162 \text{ cm}^2 \text{ cm}^{-3}$ respectively. These values are of the same order of magnitude than the ones given by BET method, although sensibly lower, which could arise for example because of gases used with said method accessing extremely small pores on fibers surfaces. The total surfaces found here are also very close to the one found by other authors, although with felts from different providers [28], where 86.2 cm^2 was estimated for a $1 \text{ cm} \times 1 \text{ cm}$ square piece of felt while in this study a surface between 45 cm^2 and 60 cm^2 corresponds to a 1 cm diameter cylinder.

4.1.2. Electrochemical impedance spectroscopy on porous “macro-felts”

Nyquist and Bode phase diagrams at formal potential are displayed in **Figure 3** for TEMPO-based and ferrocyanide molecules. Each experimental diagram corresponds to the replication giving intermediate values, out of the three performed, along with its fit using macrohomogeneous model. The three replications' diagrams can be seen in **Figure S4**. A strong dispersion on dc resistance value was observed, which resulted in Nyquist diagrams having very close shapes despite being sometimes far from one another on the real numbers' axis, which is proven in **Figure S5** by translating the curves on the horizontal direction. Nyquist diagrams show fast redox systems, conformingly to what could be observed with cyclic voltammograms: it is very hard to discern an impedance part corresponding to the “kinetic arc” that should typically be observed [13], [15], especially for activated felts. Comparing the curves at other potentials around estimated formal potential allows to see more clearly that the area between the highest and intermediate frequencies ought to contain kinetic information (**Figure S6**).

Values of fitted parameters are available in **Table 2** and all three replication's data is given in supplementary material. These values are presented under the form of the average of each fitted parameter calculated from each replication's fit, as well as the corresponding absolute standard deviation. Charge transfer kinetic constants k^0 obtained by the fit have values of the order of 10^{-3} cm s⁻¹ for both molecules on non-activated felt, with significantly higher values for ferrocyanide, and values of the order of $3.0-5.0 \times 10^{-2}$ cm s⁻¹ for both molecules with activated felt. The relative deviation is high, between 33% and 50% of its value, except for TEMPO-based molecule on non-activated felt where it is very small. Values for non-activated

felts are in good agreement with the order of magnitude found in the literature for similar materials [21], [28], [29]. Some authors found extremely higher kinetic constant values in the order of $10^{-1} \text{ cm.s}^{-1}$ on the same felt without any mention of activation [30], but this measurement was based on Nicholson analysis which assumes flat electrodes. Specific surface areas range between 150 and $350 \text{ cm}^2 \text{ cm}^{-3}$, which is in good agreement with BET measurements on non-activated felts and in even better agreement with estimation from manual fit of cyclic voltammograms. However, no increase in fitted surface area value is observed for activated felts, which does not agree with experimental BET results as well as cyclic voltammetry.

CPE exponent, when needed, would only slightly deviate from 1. Yet, this parameter was necessary for the convergence of kinetic constant to occur for activated felt. With such a high exponent value, typically above 0.97, capacitance value can be approximated by fitted pseudo-capacitance without the need of Brug relation. Its value is very high, ranging from $70 \mu\text{F cm}^{-2}$ to $100 \mu\text{F cm}^{-2}$. Values of the order of $10 \mu\text{F cm}^{-2}$ or lower are expected for non-activated graphite felt. Tortuosity coefficient f_p was fitted, as its value was not known, and it also permitted to correct values of the felt porosity and the electrolyte resistivity. Its value was expected to be close to 1, but significantly higher values were needed to obtain a good fit at higher frequencies for non-activated felts. The parameters l_p and ϕ were mainly considered adjustment parameters, much like f_p , and no regard is given to their possible physical meaning. It is nevertheless possible to say that their values are very consistent: l_p for the non-activated felt present a value of about $10 \mu\text{m}$ whereas for activated one it shows a value of about $20 \mu\text{m}$ and the non-uniform diffusion coefficient ϕ is around 0.91 except for ferrocyanide on activated felt where it is about 0.82 with a higher deviation. We do not have any good explanation on why the slope of complex impedance at low frequencies is different

for ferrocyanide/ferricyanide system on non-activated felt. This result was nevertheless systematically reproduced.

Therefore, Paasch and co-authors' model allows a very good fit of impedance data at the cost of many parameters to optimize. Consistent values for heterogeneous charge transfer kinetic constant and specific surface estimates are obtained but resulting double-layer capacitance is most probably overestimated. Other optima could have been selected, starting from different initial points in the optimization routine. Some optima were found to give a double-layer capacitance value for non-activated felt about ten times lower than the ones presented here, thus much closer to what should be expected. Optimized kinetic constant was then typically higher by one order of magnitude, and tortuosity coefficient closer to one, but the fit at lower frequencies was significantly worse. It came to mind that the ferrocyanide peculiarity at low frequency could actually mean a much slower kinetic constant, and it was indeed possible to find optima with kinetic constant of the order of 10^{-4} cm s⁻¹ for this system, but the specific surface area would then typically converge to a very high value that seemed absurd based on BET measurements and thus this result was rejected.

4.2. Experimental results with carbon fiber microelectrodes

4.2.1. Cyclic voltammetry on carbon fiber microelectrodes

Cyclic voltammograms on carbon fiber microelectrodes are presented for both studied salts in **Figure 4**. While the curves for TEMPO-based molecule are close in shape for non-activated and activated fibers, a tremendous difference appears for ammonium hexacyanoferrate: activated fiber exhibits a typically reversible curve, much like the curves obtained from TEMPO-based molecule, but non-activated fiber displays a slower system behavior. This

result, corroborated with EIS measurements at different potentials, motivated the use of an adapted quasi-reversible model for the latter case.

In **Figure 5** an example of experimental voltammogram with its fit is shown for each system on activated and non-activated fibers. The agreement between the fit and experimental data is overall very good, especially for activated fibers with fast kinetics, but some error can be seen typically for the reverse scan. Optimized parameters, k^0 and A , are presented in **Table 3**, as the average of fitted values for the three replications along with standard deviation. Kinetic constant values are indeed very low, of the order of 10^{-5} cm s⁻¹, for ferrocyanide on non-activated fibers. However, the activation permits a dramatic increase of charge transfer kinetic constant with this system, where kinetic constants reach several 10^{-3} cm s⁻¹. Such a low kinetic constant with ferrocyanide/ferricyanide on carbon fiber microelectrodes was also observed by Landon-Lane, along with the dramatic increase after thermal activation [22]. The kinetic constants obtained for TEMPO-based molecule on non-activated fibers range from 3×10^{-4} cm s⁻¹ to almost 10^{-3} cm s⁻¹, and after activation these values increase by an order of magnitude. The surface area was fitted to almost the same value as the geometric one in the case of TEMPO-based molecule, and for ammonium hexacyanoferrate in the case of activated fiber. However, its fitted value was slightly lower in the case of non-activated fibers for the latter system. This could not be the conjunction of an overestimated measurement of the fiber length and a fiber radius deviating from its average value, as these electrodes were the same as the ones for measurements with TEMPO-based system. As it was not the purpose of this study, this discrepancy was not further investigated.

4.2.2. Electrochemical impedance spectroscopy on carbon fiber microelectrodes

Nyquist and Bode phase diagrams for the two systems are presented in **Figure 6**. Here impedance was first normalized to the corresponding microelectrode geometric area for standard deviations to be meaningful. Otherwise, formalism used for the presentation of data in Nyquist diagrams is identical to the one previously described for graphite felts. The three replications' diagrams can be seen in **Figure S7**. The shape of all Nyquist diagrams is the one expected and reported in other works [21], [27]: A half-circle containing charge transfer kinetics information is present and clearly identifiable at higher frequencies, while at lower frequencies complex impedance represents mass transfer and the imaginary part appears to tend towards a constant.

Optimized parameters' values, averaged over the three replications, are displayed in **Table 4** and data for all replications is given in supplementary material. Optimized parameters with activated fibers all present significantly higher standard deviation values than non-activated felts. All k^0 obtained from EIS are in very good agreement with the ones obtained by fitting cyclic voltammograms: charge transfer kinetic constants k^0 is only slightly lower than $10^{-3} \text{ cm.s}^{-1}$ for TEMPO-based system with non-activated fibers and of the order of several $10^{-3} \text{ cm.s}^{-1}$ with activated ones, while the diffusion coefficient D is in both cases about $4 \times 10^{-6} \text{ cm}^2.\text{s}^{-1}$ which is basically the value that was determined with other methods (results not shown in this work). For ferrocyanide/ferricyanide system, k^0 is of the order of $10^{-5} \text{ cm.s}^{-1}$ with non-activated fibers whereas with activated fibers it has about the same value and standard deviation as the TEMPO-based molecule, which makes for a dramatic increase of a factor 500 due to activation, as was observed when fitting cyclic voltammograms. This validates the specific formalism that was used to fit the spectra for ammonium ferrocyanide.

The diffusion coefficient is optimized at a value close to its measured one for activated fibers, but is far from the correct value of $4.5 \times 10^{-6} \text{ cm}^2 \cdot \text{s}^{-1}$ in the case of non-activated fibers. It was attempted to fix this value, but resulting optima would yield poor fitting of experimental data. On the other hand, the mass transfer coefficient fitted with quasi-reversible model in the case of $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$ with non-activated fiber is $1.0 \times 10^{-3} \text{ cm} \cdot \text{s}^{-1}$, and this value is not absurd when compared to typical values of this parameter [31],[32]. Furthermore, using Szabo's approximation for semi-infinite diffusion in cylindrical geometry [33] in the limiting case of quasi-steady-state as reported by Bard and Faulkner [31], and as presented in **equation (20)**, it is possible to estimate a mass transfer coefficient of about $2.9 \times 10^{-3} \text{ cm} \cdot \text{s}^{-1}$ taking a value of five minutes for t (buffer time applied before each potential step in EIS). Taking a time of one hour only changes the value to $2.3 \times 10^{-3} \text{ cm} \cdot \text{s}^{-1}$ and these values agree very well with the one optimized here.

$$m_{O,qss} = \frac{2D}{r_0 \ln\left(\frac{4Dt}{r_0^2}\right)} \quad (20)$$

The fitted values of CPE element's parameters are similar for the two redox systems studied here: the exponent ν is very close to 1 for non-activated fiber but decreases to a value of about 0.8 for activated fibers. The deviation of this parameter from unity is usually attributed to a more inhomogeneous electrode surface [34],[35], which is expected for a felt having undergone for instance a thermal treatment. Non-activated fibers have a pseudo-capacitance of several $\mu\text{F} \cdot \text{cm}^{-2}$, which is very small but not unexpected for this material [18], [20], [21]. Activated ones show dramatically higher values for this parameter, with about $25 \mu\text{F} \cdot \text{cm}^{-2}$ for TEMPO-based system and going above $100 \mu\text{F} \cdot \text{cm}^{-2}$ for ammonium hexacyanoferrate. One may recombine Brug relation [34] to access actual capacitance values from those estimated:

$$C = \left[Y_0 \left(\frac{1}{R_{ohm}} + \frac{1}{R_{ct}} \right)^{\nu-1} \right]^{\frac{1}{\nu}} \quad (21)$$

Capacitances on activated felts obtained with this method are between 6 and 25 $\mu\text{F}\cdot\text{cm}^{-2}$ for the TEMPO-based compound and between 40 and 130 $\mu\text{F}\cdot\text{cm}^{-2}$ for ferrocyanide. This increase in double-layer capacitance represents a higher number of charges adsorbed at the electrode's surface, which could arise from an increase of this surface. This is also observed and described in Landon-Lane's work [22].

4.3. Comparison between macro-felts and micro-fibers

The charge transfer kinetics of two redox systems on activated and non-activated carbon felts are studied at two different scales: macro-felt and micro-fiber. For each scale, CV and EIS are carried out and analyzed by corresponding models. At macro-felts, CV analysis is based on Polarographica software and EIS data are fitted with a modified macrohomogeneous model. At micro-fibers, CV is fitted by considering non-Nernstian charge transfer kinetics and 1D semi-infinite diffusion in cylindrical coordinates, and EIS is fitted with a modified Randles circuit adapted to cylindrical geometry.

For macro-felts, the oxidation and reduction peak separation is an indication of charge transfer kinetics on the electrode. From the results in **Figure 1**, it is seen that the differences between activated and non-activated felts in both TEMPO-based compound and ferrocyanide electrolytes are insignificant. This is confirmed by Polarographica simulations where voltammograms almost do not change when increasing k^0 to values higher than $10^{-2} \text{ cm}\cdot\text{s}^{-1}$. EIS is more sensitive to differences in charge transfer kinetics. Qualitatively, differences can be observed from the Nyquist plot in high-frequency range (**Figure 3**). By fitting with the

macrohomogenous model, k^0 value is one order of magnitude higher for activated felts than for non-activated ones in both TEMPO-based and ferrocyanide electrolytes and this is consistent with what can be found for other systems in literature [36].

For micro-fibers, the difference in kinetics yields a clear change in the shape of CV, especially for ferrocyanide electrolyte. The k^0 values derived from fitting voltammograms show an increase of 300 and 10 times by activation in ferrocyanide and TEMPO electrolytes, respectively. Furthermore, they are in good agreement with the values analyzed from EIS measurements (**Table 4**). This suggests that both CV and EIS may provide reliable kinetic information of the micro-fibers. We consider that the micro-fibers have two main advantages over the macro-felts for the quantitative analysis: well-defined geometry that gets rid of the approximations in felt structure and simplifies the mathematical treatment in 1D, and much lower capacitive contribution in CV due to much lower active surface area. However, it must be pointed out that k^0 values analyzed from micro-fibers are systematically lower than that analyzed from macro-felts where the differences may reach *ca.* two orders of magnitude in some cases. The reason for these discrepancies is unclear but one possible explanation could be the analytical model for macro-felts, where the geometrical approximations might not always correspond well to the experimental systems. From the practical point of view, direct measurements on macro-felts are more convenient and straightforward. On the other hand, measurements on extracted micro-fibers may reveal more sensitively the kinetic difference for electrode materials undergoing different treatments. Moreover, they may also provide information on the heterogeneity of the felt by extracting fibers randomly from different areas.

5. Conclusion

Charge transfer kinetic constant was measured using CV fitting or simulation and EIS on both macro-felts and micro-fibers, activated or not, for two aqueous organic electrolytes: ferrocyanide and TEMPO-based molecule. For macro-felts, EIS is more sensitive to k^0 determination than CV analysis in the range of k^0 considered, and an increase by one order of magnitude is measured for activated material. For micro-fibers, differences can already be observed on CV curves, and analysis using CV fitting and EIS give k^0 values with very good consistency. However, k^0 measured on micro-fibers are typically lower than the ones obtained on macro-felts. The factor by which these values are decreased is usually about one order of magnitude, but reaches about two orders of magnitude in the case of ferrocyanide on non-activated felt. Since results given by the two electrochemical methods are consistent with each other for the felt and for the micro-fiber, the origin of this discrepancy is not well understood. Since physical analysis on microelectrodes is much straightforward and subject to less approximations, it is reasonable to attribute said discrepancy to possible misrepresentation of the complex porous structure by the models used. Because of the high sensitivity of methods used on micro-fibers to detect k^0 variation with material treatment, this support makes for an excellent tool to detect the effect of various treatments on charge transfer kinetics for carbon-based felt materials.

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References

- [1] T. Liu, X. Li, H. Zhang, and J. Chen, "Progress on the electrode materials towards vanadium flow batteries (VFBs) with improved power density," *J. Energy Chem.*, vol. 27, no. 5, pp. 1292–1303, 2018, doi: 10.1016/j.jechem.2018.07.003.
- [2] R. G. Compton, E. Laborda, and K. R. Ward, *Understanding Voltammetry*. 2014.
- [3] E. O. Barnes, X. Chen, P. Li, and R. G. Compton, "Voltammetry at porous electrodes: A theoretical study," *J. Electroanal. Chem.*, 2014, doi: 10.1016/j.jelechem.2014.03.028.
- [4] R. E. G. Smith, T. J. Davies, N. D. B. Baynes, and R. J. Nichols, "The electrochemical characterisation of graphite felts," *J. Electroanal. Chem.*, 2015, doi: 10.1016/j.jelechem.2015.03.029.
- [5] T. Tichter, D. Andrae, J. Mayer, J. Schneider, M. Gebhard, and C. Roth, "Theory of cyclic voltammetry in random arrays of cylindrical microelectrodes applied to carbon felt electrodes for vanadium redox flow batteries," *Phys. Chem. Chem. Phys.*, vol. 21, no. 18, pp. 9061–9068, 2019, doi: 10.1039/c9cp00548j.
- [6] T. Tichter, J. Schneider, D. Andrae, M. Gebhard, and C. Roth, "Universal Algorithm for Simulating and Evaluating Cyclic Voltammetry at Macroporous Electrodes by Considering Random Arrays of Microelectrodes," *ChemPhysChem*, vol. 21, no. 5, pp. 428–441, 2020, doi: 10.1002/cphc.201901113.
- [7] R. de Levie, "On porous electrodes in electrolyte solutions. I. Capacitance effects,"

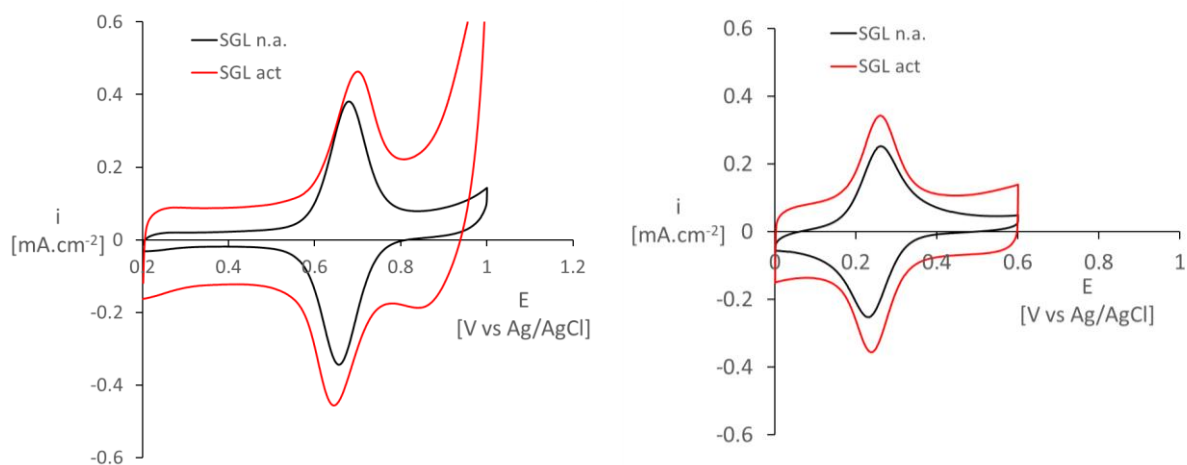
- Electrochim. Acta*, 1963, doi: 10.1016/0013-4686(63)80042-0.
- [8] R. de Levie, "On porous electrodes in electrolyte solutions-IV," *Electrochim. Acta*, 1964, doi: 10.1016/0013-4686(64)85015-5.
- [9] J. P. Candy, P. Fouilloux, M. Keddam, and H. Takenouti, "The characterization of porous electrodes by impedance measurements," *Electrochim. Acta*, vol. 26, no. 8, pp. 1029–1034, 1981, doi: 10.1016/0013-4686(81)85072-4.
- [10] J. P. Candy, P. Fouilloux, M. Keddam, and H. Takenouti, "The pore texture of raney-nickel determined by impedance measurements," *Electrochim. Acta*, vol. 27, no. 11, pp. 1585–1593, 1982, doi: 10.1016/0013-4686(82)80084-4.
- [11] A. Lasia, "Impedance of porous electrodes," *J. Electroanal. Chem.*, vol. 397, no. 1–2, pp. 27–33, 1995, doi: 10.1016/0022-0728(95)04177-5.
- [12] A. Lasia, "Porous electrodes in the presence of a concentration gradient," *J. Electroanal. Chem.*, vol. 428, no. 1–2, pp. 155–164, 1997, doi: 10.1016/S0022-0728(96)05071-1.
- [13] A. Lasia, "Modeling of Impedance of Porous Electrodes," 2008.
- [14] G. Paasch, K. Micka, M. Schwarzenberg, K. Jobst, and L. Swatschenko, "Studies on thermally structurized polyacrylonitrile by the electrochemical impedance method," *Electrochim. Acta*, vol. 37, no. 13, pp. 2453–2458, 1992, doi: 10.1016/0013-4686(92)87083-C.
- [15] G. Paasch, K. Micka, and P. Gersdorf, "Theory of the electrochemical impedance of macrohomogeneous porous electrodes," *Electrochim. Acta*, vol. 38, no. 18, pp. 2653–2662, 1993, doi: 10.1016/0013-4686(93)85083-B.
- [16] J. P. Meyers, M. Doyle, R. M. Darling, and J. Newman, "The Impedance Response of a Porous Electrode Composed of Intercalation Particles," *J. Electrochem. Soc.*, vol. 147, no. 8, p. 2930, 2000, doi: 10.1149/1.1393627.
- [17] M. Doyle, J. P. Meyers, and J. Newman, "Computer Simulations of the Impedance Response of Lithium Rechargeable Batteries," *J. Electrochem. Soc.*, vol. 147, no. 1, p. 99, 2000, doi: 10.1149/1.1393162.
- [18] P. M. Kovach, M. R. Deakin, and R. M. Wightman, "Electrochemistry at partially blocked carbon-fiber microcylinder electrodes," *J. Phys. Chem.*, vol. 90, no. 19, pp. 4612–4617, 1986, doi: 10.1021/j100410a028.
- [19] A. Bourke, M. A. Miller, R. P. Lynch, J. S. Wainright, R. F. Savinell, and D. N. Buckley, "Effect of Cathodic and Anodic Treatments of Carbon on the Electrode Kinetics of V^{IV}/V^V Oxidation-Reduction," *J. Electrochem. Soc.*, vol. 162, no. 8, pp. A1547–A1555, 2015, doi: 10.1149/2.0671508jes.
- [20] M. A. Miller *et al.*, "Kinetic Study of Electrochemical Treatment of Carbon Fiber Microelectrodes Leading to In Situ Enhancement of Vanadium Flow Battery Efficiency," *J. Electrochem. Soc.*, vol. 163, no. 9, pp. A2095–A2102, 2016, doi: 10.1149/2.1091609jes.
- [21] L. Landon-Lane, A. T. Marshall, and D. A. Harrington, "EIS at carbon fiber cylindrical microelectrodes," *Electrochem. commun.*, vol. 109, no. September, p. 106566, 2019, doi: 10.1016/j.elecom.2019.106566.

- [22] L. Landon-lane, “Investigating Electron Transfer Kinetics in Redox Flow Batteries,” University of Canterbury, Christchurch, New-Zealand, 2020.
- [23] K. Aoki and H. Kaneko, “Theory of irreversible cyclic voltammograms at microcylinder electrodes,” *J. Electroanal. Chem.*, vol. 247, no. 1–2, pp. 17–27, 1988, doi: 10.1016/0022-0728(88)80127-X.
- [24] A. Neudeck and J. Dittrich, “The determination of diffusion coefficients and rate constants from the dependence of the peak separation and peak current on the scan rate of cyclic voltammograms at micro-cylindrical electrodes,” *J. Electroanal. Chem.*, vol. 313, no. 1–2, pp. 37–59, 1991, doi: 10.1016/0022-0728(91)85170-T.
- [25] J. R. Macdonald, *Impedance Spectroscopy – Emphasizing solid materials and systems*. New York: John Wiley & Sons, 1987.
- [26] C.-N. Sun, F. M. Delnick, D. S. Aaron, A. B. Papandrew, M. M. Mench, and T. A. Zawodzinski, “Resolving Losses at the Negative Electrode in All-Vanadium Redox Flow Batteries Using Electrochemical Impedance Spectroscopy,” *J. Electrochem. Soc.*, vol. 161, no. 6, pp. A981–A988, 2014, doi: 10.1149/2.045406jes.
- [27] T. Jacobsen and K. West, “Diffusion impedance in planar, cylindrical and spherical symmetry,” *Electrochim. Acta*, vol. 40, no. 2, pp. 255–262, 1995, doi: 10.1016/0013-4686(94)E0192-3.
- [28] L. Landon-Lane, A. J. Downard, and A. T. Marshall, “Single fibre electrode measurements – A versatile strategy for assessing the non-uniform kinetics at carbon felt electrodes,” *Electrochim. Acta*, vol. 354, p. 136709, 2020, doi: 10.1016/j.electacta.2020.136709.
- [29] R. L. McCreery, “Advanced carbon electrode materials for molecular electrochemistry,” *Chemical Reviews*. 2008, doi: 10.1021/cr068076m.
- [30] J. Luo *et al.*, “Unraveling pH dependent cycling stability of ferricyanide/ferrocyanide in redox flow batteries,” *Nano Energy*, vol. 42, no. October, pp. 215–221, 2017, doi: 10.1016/j.nanoen.2017.10.057.
- [31] A. J. Bard and L. R. Faulkner, *Electrochemical Methods*, 2nd ed. John Wiley & Sons, Inc., 2001.
- [32] X. Gao and H. S. White, “Rotating Microdisk Voltammetry,” *Anal. Chem.*, vol. 67, no. 22, pp. 4057–4064, 1995, doi: 10.1021/ac00118a005.
- [33] A. Szabo, D. K. Cope, D. E. Tallman, P. M. Kovach, and R. M. Wightman, “Chronoamperometric current at hemicylinder and band microelectrodes: Theory and experiment,” *J. Electroanal. Chem.*, vol. 217, no. 2, pp. 417–423, 1987, doi: 10.1016/0022-0728(87)80233-4.
- [34] G. J. Brug, A. L. G. van den Eeden, M. Sluyters-Rehbach, and J. H. Sluyters, “The analysis of electrode impedances complicated by the presence of a constant phase element,” *J. Electroanal. Chem.*, 1984, doi: 10.1016/S0022-0728(84)80324-1.
- [35] M. E. Orazem and B. Tribollet, *Electrochemical Impedance Spectroscopy*. 2008.
- [36] P. Mazúr *et al.*, “Performance evaluation of thermally treated graphite felt electrodes for vanadium redox flow battery and their four-point single cell characterization,” *J. Power Sources*, vol. 380, no. November 2017, pp. 105–114, 2018, doi:

Appendix: List of symbols

Parameter	Symbol	Unit
Number of electrons exchanged	n	[-]
Faraday constant	F	[C mol ⁻¹]
Ideal gas constant	R_g	[J mol ⁻¹ K ⁻¹]
Temperature	T	[K]
Charge transfer coefficient	α	[-]
Oxidant concentration near the electrode surface	C_O^s	[mol cm ⁻³]
Reductant concentration near the electrode surface	C_R^s	[mol cm ⁻³]
Oxidant concentration in bulk electrolyte	C_O^*	[mol cm ⁻³]
Reductant concentration in bulk electrolyte	C_R^*	[mol cm ⁻³]
Formal potential of the redox system	E^0	[V]
dc potential applied before EIS measurement	\bar{E}	[V]
Carbon felt porosity	ε	[-]
Carbon felt resistivity	ρ_l	[Ω cm]
Bulk electrolyte resistivity	ρ_2^*	[Ω cm]
Macrohomogeneous model parameter 1	k	[s ⁻¹]
Macrohomogeneous model parameter 2	K	[cm ² s ⁻¹]
Macrohomogeneous model parameter 3	β	[-]
Hindrance function	y	[-]
Double-layer capacitance	C	[F cm ⁻²]
Specific surface area	S_c	[cm ² cm ⁻³]
Heterogeneous charge transfer kinetic constant	k^0	[cm s ⁻¹]
Diffusion coefficient	D	[cm ² s ⁻¹]
Diffusion layer thickness	l_p	[cm]
Porous electrode thickness	d	[cm]
Electrode geometric area (cross section for carbon felt, cylinder area for microelectrode)	A	[cm ²]
Time	t	[s]
Angular frequency	ω	[s ⁻¹]
Imaginary number	j	[-]
CPE exponent	γ	[-]
Non-uniform diffusion exponent	ϕ	[-]
Ohmic resistance	R_{ohm}	[Ω]
Charge transfer resistance	R_{ct}	[Ω]
Carbon fiber microelectrode radius	r_0	[cm]
Pseudo-capacitance (for microelectrode model)	Y_0	[F cm ⁻² s ^($\nu-1$)]
Mass transfer coefficient	m_0	[cm s ⁻¹]

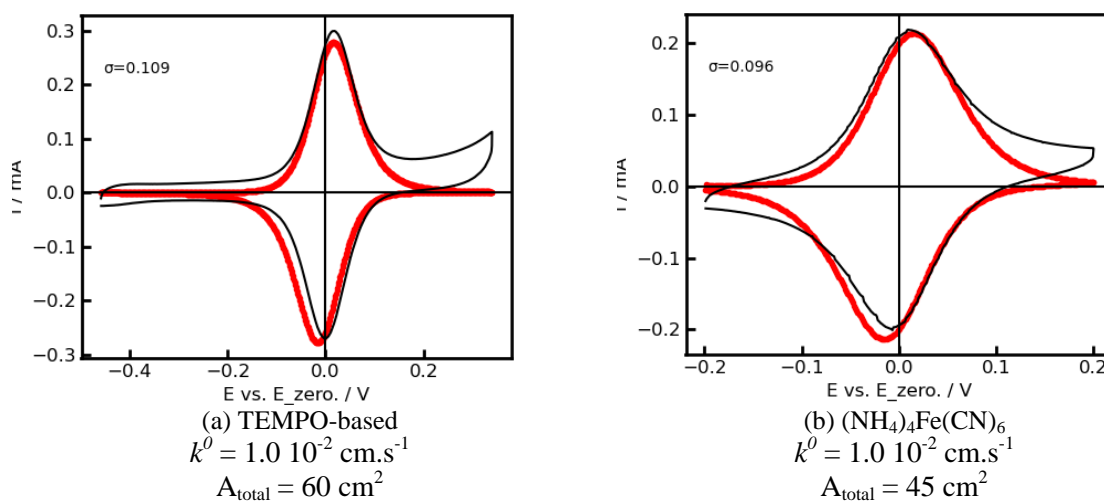
Figures:



(a) TEMPO-based molecule

(b) Ammonium hexacyanoferrate

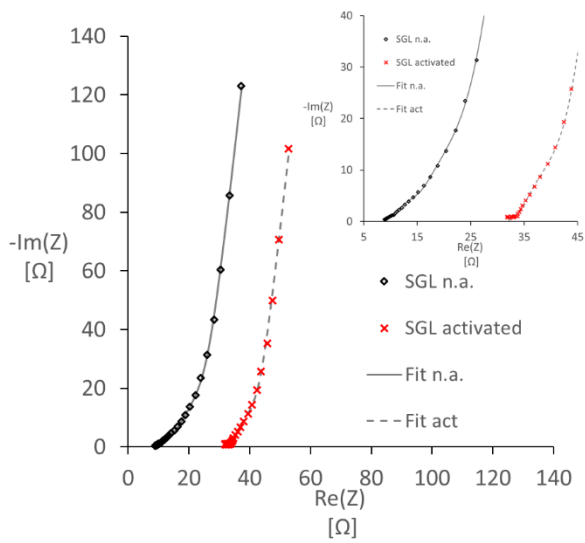
Figure 1: Typical cyclic voltammograms on SGL graphite felts with (a) TEMPO-based molecule and (b) ammonium hexacyanoferrate. Scan rate: 10 mV.s^{-1}



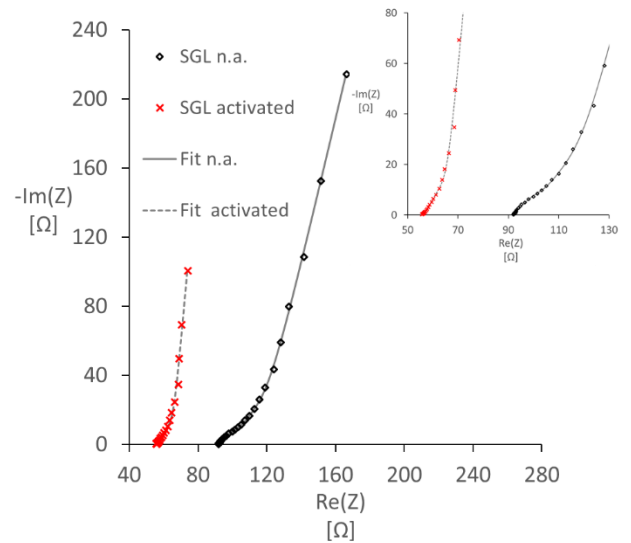
(a) TEMPO-based
 $k^0 = 1.0 \cdot 10^{-2} \text{ cm.s}^{-1}$
 $A_{\text{total}} = 60 \text{ cm}^2$

(b) $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$
 $k^0 = 1.0 \cdot 10^{-2} \text{ cm.s}^{-1}$
 $A_{\text{total}} = 45 \text{ cm}^2$

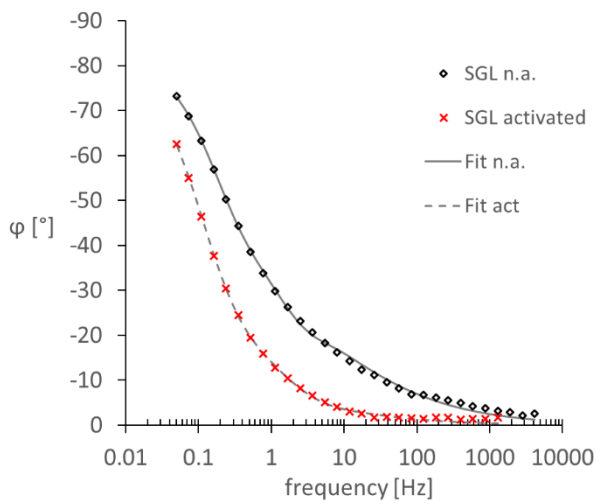
Figure 2: Simulation of cyclic voltammetry with “Statistically Weighted External Cylindrical Finite Diffusion” model using Polarographica software (in red) compared with compensated experimental data (in black) on non-activated SGL graphite felt with TEMPO-based system (left) and ammonium hexacyanoferrate (right). Scan rate: 10 mV.s^{-1} . Concentration: 10^{-4} M in supporting electrolyte.



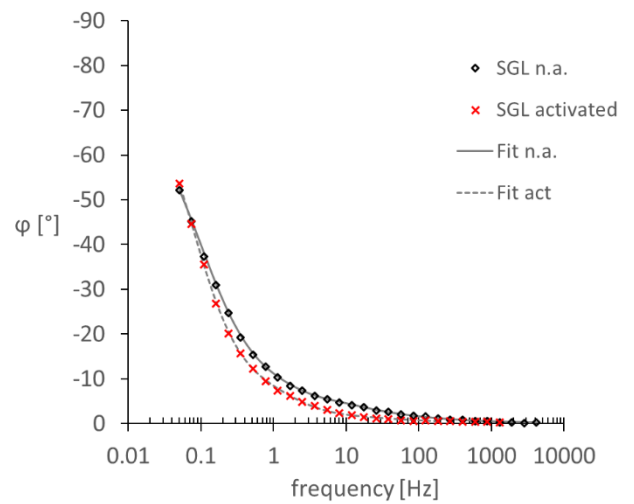
(a) TEMPO-based Nyquist diagram



(b) $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$ Nyquist diagram

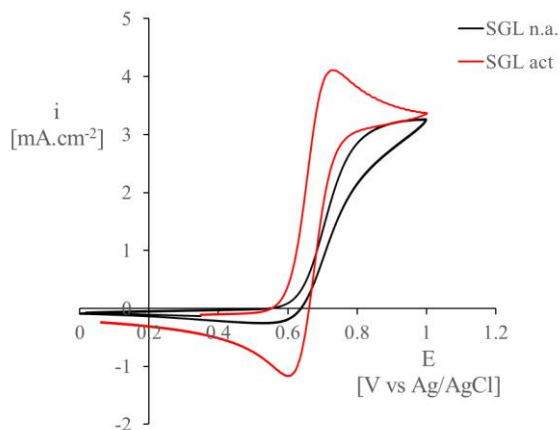


(c) TEMPO-based Bode phase diagram

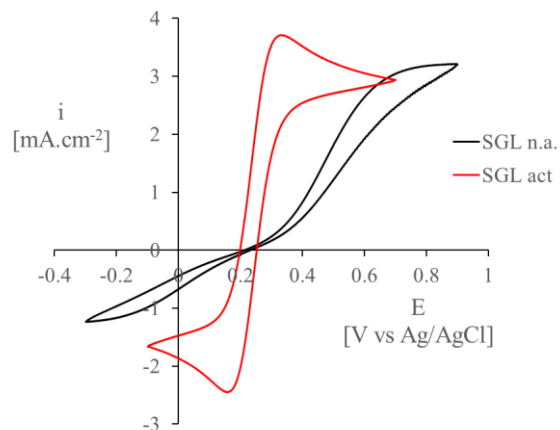


(d) $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$ Bode phase diagram

Figure 3: Nyquist diagrams (up) and Bode phase diagrams (down) of EIS experiments on SGL graphite felt with TEMPO-based molecule (left) at estimated formal potential ($E^{0'} = 0.66 \text{ V}$ vs Ag/AgCl quasi-reference) and with $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$ molecule (right) at estimated formal potential ($E^{0'} = 0.25 \text{ V}$ vs Ag/AgCl quasireference)

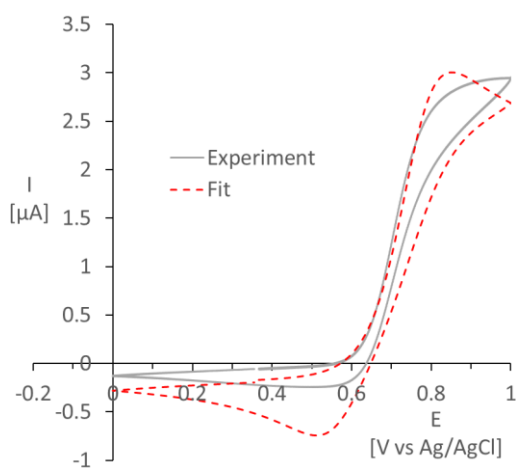


(a) TEMPO-based molecule

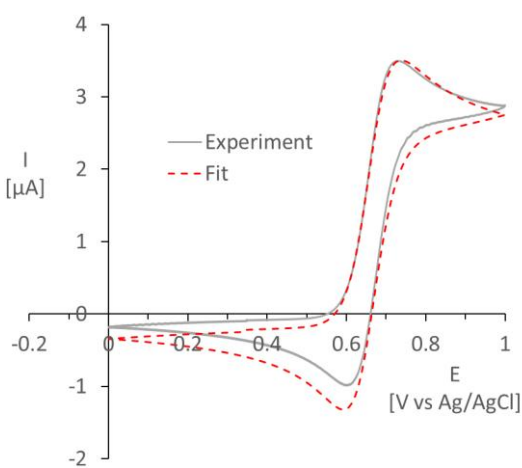


(b) Ammonium hexacyanoferrate

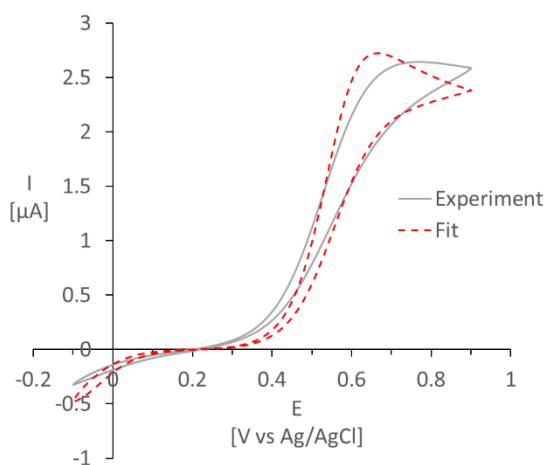
Figure 4: Typical cyclic voltammograms on carbon fiber microelectrode plotted on micro-fibers extracted from SGL graphite felt with (a) TEMPO-based molecule and (b) ammonium hexacyanoferrate. Scan rate: 100 mV.s^{-1}



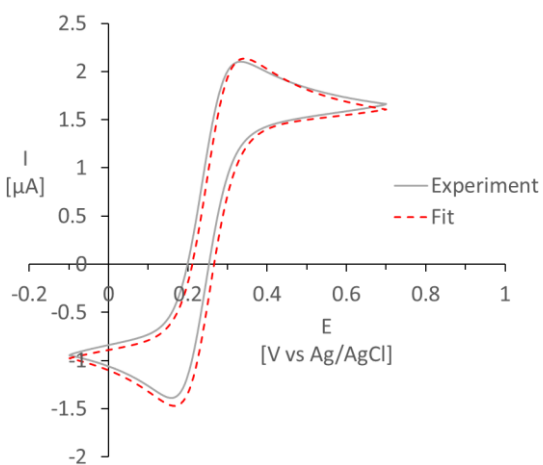
(a) TEMPO-based on non-activated fiber



(b) TEMPO-based on activated fiber

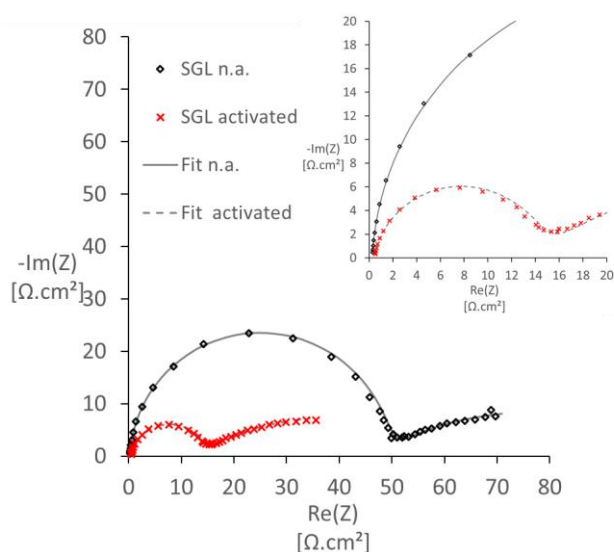


(c) $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$ on non-activated fiber

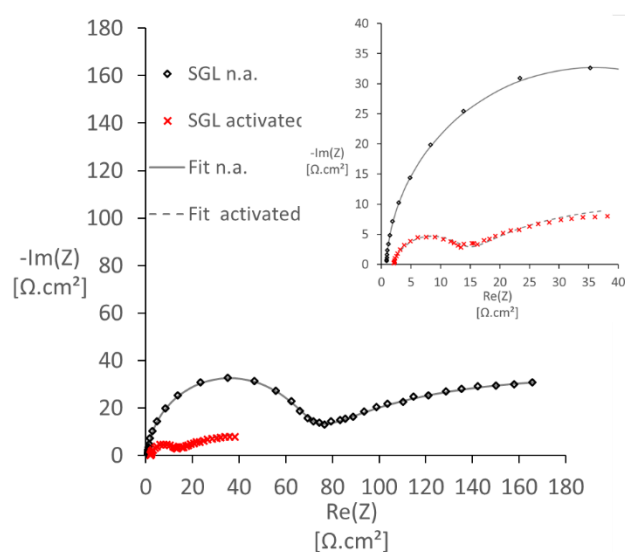


(d) $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$ on activated fiber

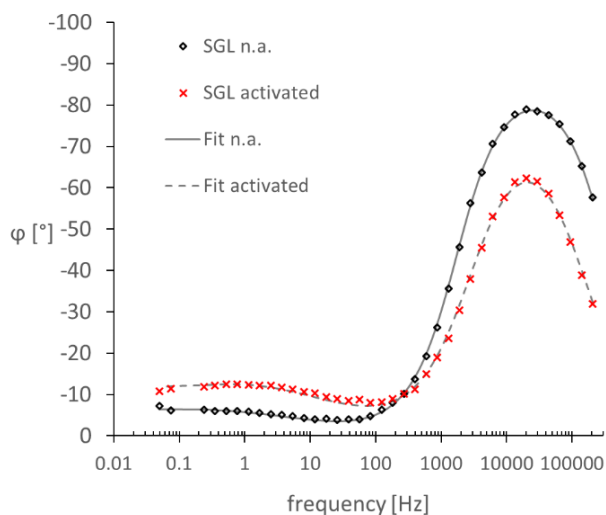
Figure 5: Examples of microelectrodes' cyclic voltammograms and their fit with k^0 and A optimized. Scan rate: 100 mV.s^{-1} .



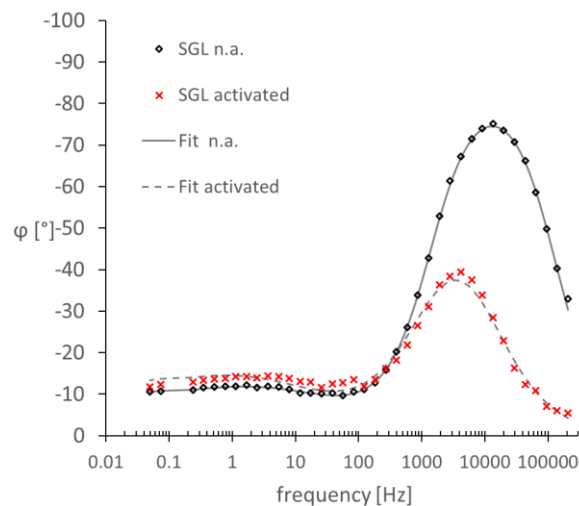
(a) TEMPO-based Nyquist diagram



(b) $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$ Nyquist diagram



(c) TEMPO-based Bode phase diagram



(d) $(\text{NH}_4)_4\text{Fe}(\text{CN})_6$ Bode phase diagram

Figure 6: Nyquist diagrams (up) and Bode phase diagrams (down) of EIS experiments on carbon fiber microelectrode plotted on micro-fibers extracted from SGL graphite felt with ammonium hexacyanoferrate. Potential of study was the estimated formal potential for fiber from activated felt ($E^0 = 0.25 \text{ V}$ vs Ag/AgCl quasi-reference), whereas for fiber from non-activated felt, it was chosen close to impedance minimum (situated around $E = 0.6 \text{ V}$ vs Ag/AgCl quasi-reference).

Tables:

Table 1: Physical parameters fixed for the macrohomogeneous model

Parameter	Symbol	Unit	TEMPO-based	(NH ₄) ₄ Fe(CN) ₆
Charge transfer coefficient	α	[-]	0.5	0.5
Carbon felt Porosity	ε	[-]	0.95 ^a	0.95 ^a
Carbon felt resistivity	ρ_1	[Ω cm]	0.5 ^b	0.5 ^b
Bulk electrolyte resistivity	ρ_2^*	[Ω cm]	12.5 ^c	10 ^d
Diffusion coefficient	D	[cm ² s ⁻¹]	3.8 10 ^{-6f}	4.5 10 ^{-6f}

^{a,b} taken from providers' estimations

^{c,d} Gross estimation based on literature conductivity values of support electrolyte [37]

^f measured from Linear Scan Voltammetry on a glassy carbon rotating electrode (not shown in this work)

Table 2: Averaged optimized parameters' values (calculated from the three replications) and their standard deviation values for EIS experiments carried out on SGL graphite felt with TEMPO-based molecule (left) at estimated formal potential of 0.66 V vs Ag/AgCl and with ammonium hexacyanoferrate (right) at estimated formal potential of 0.25 V vs Ag/AgCl

	TEMPO-based				(NH ₄) ₄ Fe(CN) ₆			
	Non-activated		Activated		Non-activated		Activated	
	Average	Std Dev	Average	Std Dev	Average	Std Dev	Average	Std Dev
R_{ohm} [Ω]	19.1	17.7	31.1	2.3	64.0	48.2	51.2	41.1
C [F cm ⁻² s ^{γ-1}]	1.02 $\times 10^{-4}$	3.51 $\times 10^{-6}$	9.56 $\times 10^{-5}$	1.34 $\times 10^{-5}$	5.27 $\times 10^{-5}$	7.34 $\times 10^{-6}$	9.04 $\times 10^{-5}$	2.21 $\times 10^{-6}$
γ [-]	0.97	0.002	0.99	0.01	1.00 [*]	0.00	1.00	0.00
k^0 [cm s ⁻¹]	1.57 $\times 10^{-3}$	7.26 $\times 10^{-5}$	5.77 $\times 10^{-2}$	3.97 $\times 10^{-2}$	2.20 $\times 10^{-3}$	3.70 $\times 10^{-4}$	3.27 $\times 10^{-2}$	2.26 $\times 10^{-2}$
l_p [cm]	0.0011	0.00003	0.0020	0.0001	0.0010	0.00009	0.0019	0.00004
S_c [cm ² cm ⁻³]	331.8	4.1	286.4	23.4	294.6	9.6	278.1	35.7
f_p [-]	3.01	0.24	1.21	0.11	6.44	0.87	1.31	0.21
ϕ [-]	0.91 [*]	0.00	0.91 [*]	0.01	0.81	0.02	0.91 [*]	0.01

^{*} value fixed during optimization

Table 3: Average of optimized parameters' values (calculated from the three replications) and their standard deviation values for microelectrodes' cyclic voltammograms with TEMPO-based compound (left) and ammonium hexacyanoferrate (right). The resulting ratio between optimized surface area and measured geometric area is also given.

	TEMPO-based				(NH) ₄ Fe(CN) ₆			
	Non-activated		Activated		Non-activated		Activated	
	Average	Std Dev	Average	Std Dev	Average	Std Dev	Average	Std Dev
k^0 [cm s ⁻¹]	6.87×10⁻⁴	3.22×10 ⁻⁴	6.58×10⁻³	4.16×10 ⁻³	1.89×10⁻⁵	9.32×10 ⁻⁶	4.76×10⁻³	2.46×10 ⁻³
A [cm ²]	7.10×10⁻⁴	1.41×10 ⁻⁴	8.08×10⁻⁴	2.03×10 ⁻⁴	6.42×10⁻⁴	8.82×10 ⁻⁵	7.79×10⁻⁴	1.32×10 ⁻⁴
A_{fit}/A_{geo} [-]	0.96	0.14	1.07	0.05	0.85	0.19	1.05	0.12

Table 4: Averaged optimized parameters' values (calculated from the three replications) and their standard deviation values for EIS experiments carried out on carbon fiber microelectrode drawn from SGL graphite felt with TEMPO-based molecule (left) at estimated formal potential of 0.66 V vs Ag/AgCl, and for ammonium hexacyanoferrate (right) at 0.60 V vs Ag/AgCl for non-activated fibers and at estimated formal potential of 0.25 V vs Ag/AgCl for activated fibers.

	TEMPO-based				(NH) ₄ Fe(CN) ₆			
	Non-activated		Activated		Non-activated		Activated	
	Average	Std Dev	Average	Std Dev	Average	Std Dev	Average	Std Dev
R_{ohm} [Ω]	534.4	172.2	2605.0	3736.3	569.0	187.4	2090.4	1555.5
Y_0 [F cm ⁻² s ⁻¹]	2.68×10⁻⁶	5.10×10 ⁻⁷	2.56×10⁻⁵	1.58×10 ⁻⁵	3.64×10⁻⁶	2.05×10 ⁻⁶	1.08×10⁻⁴	5.63×10 ⁻⁵
γ [-]	0.96	0.01	0.84	0.04	0.96	0.02	0.79	0.02
k^0 [cm s ⁻¹]	9.00×10⁻⁴	1.75×10 ⁻⁴	5.81×10⁻³	3.63×10 ⁻³	1.76×10⁻⁵	6.44×10 ⁻⁶	5.95×10⁻³	3.84×10 ⁻³
D [cm ² s ⁻¹]	4.47×10⁻⁶	1.44×10 ⁻⁶	3.88×10⁻⁶	5.59×10 ⁻⁷	1.26×10⁻⁵	5.58×10 ⁻⁶	4.38×10⁻⁶	9.17×10 ⁻⁷
m_0 [cm s ⁻¹]	-	-	-	-	1.00×10⁻³	1.63×10 ⁻⁴	-	-