

An optimal control-based numerical method for scalar transmission problems with sign-changing coefficients

Patrick Ciarlet, David Lassounon, Mahran Rihani

▶ To cite this version:

Patrick Ciarlet, David Lassounon, Mahran Rihani. An optimal control-based numerical method for scalar transmission problems with sign-changing coefficients. SIAM Journal on Numerical Analysis, 2023, 61 (3), pp.1316-1339. 10.1137/22M1495998 . hal-03666913v2

HAL Id: hal-03666913 https://hal.science/hal-03666913v2

Submitted on 13 Oct 2022

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution 4.0 International License

AN OPTIMAL CONTROL-BASED NUMERICAL METHOD FOR SCALAR TRANSMISSION PROBLEMS WITH SIGN-CHANGING COEFFICIENTS

2 3

19

1

PATRICK CIARLET JR.*, DAVID LASSOUNON[†], AND MAHRAN RIHANI[‡]

Abstract. In this work, we present a new numerical method for solving the scalar transmission problem with sign-changing coefficients. In electromagnetism, such a transmission problem can occur if the domain of interest is made of a classical dielectric material and a metal or a metamaterial, with for instance an electric permittivity that is strictly negative in the metal or metamaterial. The method is based on an optimal control reformulation of the problem. Contrary to other existing approaches, the convergence of this method is proved without any restrictive condition. In particular, no condition is imposed on the a priori regularity of the solution to the problem, and no condition is imposed on the meshes, other than that they fit with the interface between the two media. Our results are illustrated by some (2D) numerical experiments.

11 Key words. transmission problem, sign-changing coefficients, fictitious domain methods, optimal control.

12 **MSC codes.** 65N30, 78A48

13 **1** Introduction In the present paper, we study the numerical approximation of the scalar transmission 14 problem with sign-changing coefficients in \mathbb{R}^d , for $d \in \{2, 3\}$. To fix ideas, let Ω be an open, bounded, connected 15 subset of \mathbb{R}^d with a Lipschitz boundary, in other words a *domain* of \mathbb{R}^d . Further, consider that Ω is equal to 16 the union of two disjoint (sub)domains Ω_1 , Ω_2 . We denote the interface by $\Sigma = \partial \Omega_1 \cap \partial \Omega_2$ (see Figure 1 for 17 an example), and we assume that meas $\partial_{\Omega}(\partial \Omega_2 \setminus \Sigma) > 0$. The case of an inclusion corresponds to $\partial \Sigma \cap \partial \Omega = \emptyset$ 18 and, since meas $\partial_{\Omega}(\partial \Omega_2 \setminus \Sigma) > 0$, Ω_1 is always the inclusion in this configuration. When $\partial \Sigma \cap \partial \Omega$ is non-empty, 17 we assume that it is a Lipschitz submanifold of $\partial \Omega_1$ and of $\partial \Omega_2$.

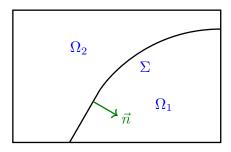


FIG. 1. Example of geometry.

We also introduce a coefficient $\varepsilon \in L^{\infty}(\Omega)$ such that $\varepsilon_1 = \varepsilon_{|\Omega_1|} \ge \varepsilon_+ > 0$ a.e. in Ω_1 and $\varepsilon_2 = \varepsilon_{|\Omega_2|} \le \varepsilon_- < 0$ in a.e. in Ω_2 . Here ε_+ and ε_- are two real constants. It will be useful to introduce the contrasts $\kappa_{\varepsilon}^1 := \varepsilon_1^-/\varepsilon_1^+$ and $\kappa_{\varepsilon}^2 := \varepsilon_2^-/\varepsilon_1^+$ where ε_1^{\pm} and ε_2^{\pm} are defined as follows:

$$\varepsilon_1^+ := \sup_{\Omega_1} \varepsilon_1, \quad \varepsilon_1^- := \inf_{\Omega_1} \varepsilon_1, \quad \varepsilon_2^+ := \sup_{\Omega_2} |\varepsilon_2| \text{ and } \varepsilon_2^- := \inf_{\Omega_2} |\varepsilon_2|.$$

Note that in the particular case where ε is piecewise constant, we have $\kappa_{\varepsilon}^1 = 1/\kappa_{\varepsilon}^2$.

25 Remark 1.1. In principle, ε could be a symmetric tensor-valued coefficient, i.e., $\varepsilon = (\varepsilon_{ij})_{1 \le i,j \le d}$ with $\varepsilon_{ij} \in L^{\infty}(\Omega)$ for all $1 \le i, j \le d$, and such that

$$\exists \varepsilon_{+} > 0, \ \forall \boldsymbol{z} \in \mathbb{R}^{d}, \quad \varepsilon_{+} |\boldsymbol{z}|^{2} \leq \varepsilon \boldsymbol{z} \cdot \boldsymbol{z} \text{ a.e. in } \Omega_{1}; \\ \exists \varepsilon_{-} > 0, \ \forall \boldsymbol{z} \in \mathbb{R}^{d}, \quad \varepsilon_{-} |\boldsymbol{z}|^{2} \leq -\varepsilon \boldsymbol{z} \cdot \boldsymbol{z} \text{ a.e. in } \Omega_{2}.$$

However, for the sake of conciseness, we consider a scalar-valued coefficient.

^{*}POEMS, CNRS, INRIA, ENSTA Paris, Institut Polytechnique de Paris, 91120 Palaiseau, France (patrick.ciarlet@ensta-paris.fr).

[†]Institut de Recherche Mathématique de Rennes (IRMAR), Université Européenne de Bretagne, 20 avenue des Buttes de Coësmes, CS 70839, 35708 Rennes Cédex 7, France (enagnon-david.lassounon@insa-rennes.fr).

[‡]CMAP, Ecole Polytechnique, Institut Polytechnique de Paris, 91128 Palaiseau, France (mahran.rihani@polytechnique.edu).

For a given source term $f \in L^2(\Omega)$, we consider the problem

30 (1.1) Find
$$u \in H_0^1(\Omega)$$
 such that $-\operatorname{div}(\varepsilon \nabla u) = f \in L^2(\Omega)$.

31 The equivalent variational formulation to (1.1) writes

32 (1.2) Find
$$u \in \mathrm{H}^{1}_{0}(\Omega)$$
 such that $\int_{\Omega} \varepsilon \nabla u \cdot \nabla v \, \mathrm{d}\boldsymbol{x} = \int_{\Omega} f v \, \mathrm{d}\boldsymbol{x}, \quad \forall v \in \mathrm{H}^{1}_{0}(\Omega).$

Because of the change of sign of ε , the well-posedness of this problem does not fit into the classical theory of elliptic PDEs and it can be ill-posed. On the other hand, one can show that when κ_{ε}^{1} or κ_{ε}^{2} is large enough, Problem 1.2 is T-coercive (for instance see [6]), i.e., there exists an operator $T : H_{0}^{1}(\Omega) \to H_{0}^{1}(\Omega)$ such that $(u, v) \mapsto \int_{\Omega} \varepsilon \nabla u \cdot \nabla(T(v))$ is coercive, and then it is well-posed. For the case of polygonal interfaces, the construction of such operator T is based on the use of local isometric geometrical transformations (such as reflections, rotations, ...) near the interface, see [3].

The implementation of a general conforming finite element method to discretize (1.2) leads us to consider the problem

41 (1.3) Find
$$u_h \in \mathcal{V}_h(\Omega)$$
 such that $\int_{\Omega} \varepsilon \nabla u_h \cdot \nabla v_h \, \mathrm{d}\boldsymbol{x} = \int_{\Omega} f v_h \, \mathrm{d}\boldsymbol{x}, \quad \forall v_h \in \mathcal{V}_h(\Omega),$

where $V_h(\Omega)$ is a well-chosen subspace of $H_0^1(\Omega)$, and the parameter h > 0 is the so-called meshsize. Even in the case where (1.2) is T-coercive, one can not guaranty that Problem (1.3) is also T-coercive. Indeed, it may happen that for some $v_h \in V_h(\Omega)$, there holds $T(v_h) \notin V_h(\Omega)$. To overcome this difficulty, an interesting idea is to try to construct meshes such that the approximation spaces $V_h(\Omega)$ are stable by operators T for which Problem (1.2) is T-coercive. This type of meshes are called T-conform meshes. Such an approach has been investigated in [26, 12, 10]. It works quite well but presents two main drawbacks:

- The construction of well-suited meshes for curved interfaces, interfaces with corners or 3D interfaces is not a straightforward task [10, 3].
- Sometimes the operator T for which the problem is T-coercive is constructed by abstract tools and therefore is not explicit. In these situations, one cannot find adapted meshes.

On general meshes, three alternatives have already been proposed. The first one was introduced in [6] and 52was based on the use of discrete trace liftings, with quasi-uniform meshes on the interface. In addition to this 53constraint on the mesh, one of the limitation of this approach is that, for interfaces with general shapes, the 54convergence can not be assured in all the configurations in which Problem (1.2) is well-posed, because it is 56based on a particular (non-optimal) T-coercivity operator. The second one is developed in [23] and is based on the use of interpolation techniques. Its essential limitation lies again in the fact that, for interfaces with general 57shapes, the convergence can not be assured for all configurations in which Problem (1.2) is well-posed. The 58 third one, presented in [14], consists in adding some dissipation to the problem (considering $\varepsilon + i\delta$ instead of ε 59in (1.2) where δ depends on the meshsize). Unfortunately, this methods has a sub-optimal order of convergence 60 even in the case where the solution and the interface are smooth (see [14]). 61

After that, in 2017, a new technique relying on the use of an optimal control reformulation has been introduced by Abdulle et al in [1], where the auxiliary control function is defined over Σ . Introducing

48

49

50

51

$$PH^{1+s}(\Omega) := \{ u \in H^1(\Omega) \, | \, u_{|\Omega_1} \in H^{1+s}(\Omega_1) \text{ and } u_{|\Omega_2} \in H^{1+s}(\Omega_2) \} \text{ for } s > 0,$$

their method is proved to be convergent for general meshes (that respect the interface) as soon as the exact solution to (1.1) belongs to the space $PH^{1+s}(\Omega)$ for some s > 1/2. Unfortunately, this regularity condition is not always satisfied, especially when Σ has corners in 2D or conical points in 3D. See the numerical illustration in Section 6.3 below.

69

⁷⁰ In this work, we present a new strategy which relies on the use of a different optimal control reformulation and ⁷¹ which converges without any restriction on the mesh (except the fact of being conforming to the interface).

and without any restriction on the regularity of the exact solution. In our approach, the auxiliary control

73 function is defined over one subdomain. This method is inspired by the smooth extension method that was

- ⁷⁴ used (without proof of convergence) in [19] to approximate the solution to some classical scalar transmission
- ⁷⁵ problems. The key idea is that, given a control, one can construct a pseudo-solution to the problem (1.1), and

to note that, as soon as one can relate the control to some extension of the solution, then one recovers exactly the solution.

78

The article is organized as follows. In Section 2, we start by giving a detailed description of the problem. Then, in Section 3, we explain how to derive an equivalent optimal control reformulation. Section 4 is dedicated to the study of some basic properties of the optimization problem and its regularization. The proposed numerical method and the proof of its convergence are given in Section 5. Our results are then illustrated by some numerical experiments in Section 6. Finally we give concluding remarks, including some possible extensions.

84 **2** Main assumption on ε and reformulation of the problem Introduce the bounded operator 85 $A_{\varepsilon}: \mathrm{H}^{1}_{0}(\Omega) \to (\mathrm{H}^{1}_{0}(\Omega))^{*}$ such that

86

$$(\mathrm{H}^{1}_{0}(\Omega))^{*}\langle A_{\varepsilon}u,v\rangle_{\mathrm{H}^{1}_{0}(\Omega)}=\int_{\Omega}\varepsilon\nabla u\cdot\nabla v\,\,\mathrm{d}\boldsymbol{x},\qquad\forall u,v\in\mathrm{H}^{1}_{0}(\Omega).$$

⁸⁷ Obviously A_{ε} is an isomorphism if, and only if, Problem (1.1) is well-posed in the Hadamard sense. In this ⁸⁸ article, we shall work under the following

Assumption 2.1. Assume that the coefficient ε is such that A_{ε} is an isomorphism.

90 If ε is piecewise constant by subdomain, the previous assumption is satisfied when the contrast $\kappa_{\varepsilon} := \varepsilon_2/\varepsilon_1$

91 does not belong to the so-called critical interval. The expression of this interval is in general not known

⁹² analytically, except for particular geometries like symmetric domains, simple 2D interface with corners, simple

⁹³ 3D interfaces with circular conical tips (see [24, Chapter 2]). Under assumption 2.1, one is able to prove the

accompanying shift theorem. We refer to [18, 7, 13, 12, 5].

THEOREM 2.2. Assume that Σ is smooth (of class \mathscr{C}^2), polygonal (in 2D) or polyhedral (in 3D) and that Problem (1.1) is well-posed in the Hadamard sense. Then, there exists $\sigma_D(\varepsilon) \in (0, 1]$ such that

97
$$\forall f \in L^2(\Omega), \text{ the solution } u \text{ to Problem (1.1) is such that} \begin{vmatrix} u \in \bigcap_{s \in [0, \sigma_D(\varepsilon))} PH^{1+s}(\Omega) & \text{if } \sigma_D(\varepsilon) < 1 \\ u \in PH^2(\Omega) & \text{if } \sigma_D(\varepsilon) = 1 \end{vmatrix}$$

98 with continuous dependence.

⁹⁹ The number $\sigma_D(\varepsilon)$ in the shift theorem is called the *(limit) regularity exponent*. For instance, when the ¹⁰⁰ interface is smooth and when it does not intersect with the boundary, then $\sigma_D(\varepsilon) = 1$ (cf. [18]).

101 Remark 2.3. In Problem (1.1), we consider homogeneous Dirichlet boundary conditions. Let us mention that

102 the results below extend quite straightforwardly to other situations, for example with Neumann or Robin-

103 Fourier boundary conditions which can be homogeneous or not, as long as the associated operator is an

104 isomorphism.

To introduce the method, we start by writing an equivalent version of (1.1) in which the unknown $u \in H_0^1(\Omega)$ is split into two unknowns defined in Ω_1 and Ω_2 : $(u_1, u_2) := (u_{|\Omega_1}, u_{|\Omega_2})$. To do so, we observe that since $f \in L^2(\Omega)$, the solution u to (1.1) is such that the vector field $\varepsilon \nabla u$ belongs to the space $H(\operatorname{div}, \Omega) = \{ \boldsymbol{v} \in (L^2(\Omega))^d \text{ such that } \operatorname{div} \boldsymbol{v} \in L^2(\Omega) \}$. Consequently, the pair of functions (u_1, u_2) satisfies the problem

(2.1) Find
$$(u_1, u_2) \in V_1(\Omega_1) \times V_2(\Omega_2)$$
 such that

$$\begin{aligned} -\operatorname{div}(\varepsilon_1 \nabla u_1) &= f_1 =: f_{|\Omega_1} \\ -\operatorname{div}(\varepsilon_2 \nabla u_2) &= f_2 =: f_{|\Omega_2} \\ \varepsilon_1 \partial_n u_1 &= \varepsilon_2 \partial_n u_2 \text{ in } (H_{00}^{1/2}(\Sigma))^* \\ u_1 &= u_2 \text{ in } H_{00}^{1/2}(\Sigma) \end{aligned}$$

in which n stands for the unit normal vector to Σ oriented to the exterior of Ω_2 (see Figure 1), the spaces $V_1(\Omega_1), V_2(\Omega_2)$ are given by

$$V_1(\Omega_1) := \{ v \in H^1(\Omega_1), v = 0 \text{ on } \partial\Omega_1 \setminus \Sigma \}, \qquad V_2(\Omega_2) := \{ v_2 \in H^1(\Omega_2), v_2 = 0 \text{ on } \partial\Omega_2 \setminus \Sigma \},$$

and the space $H_{00}^{1/2}(\Sigma)$ is defined as follows

$$H_{00}^{1/2}(\Sigma) = \begin{cases} H^{1/2}(\Sigma) & \text{if } \partial \Sigma \cap \partial \Omega = \emptyset \text{ (inclusion)}, \\ \{\lambda \in H^{1/2}(\Sigma), \ \tilde{\lambda} \in H^{1/2}(\partial \Omega_2)\} & \text{else.} \end{cases}$$

Above, in the definition of the space $H_{00}^{1/2}(\Sigma)$, $\tilde{\lambda}$ denotes the continuation of λ by 0 to $\partial\Omega_2$ (one can also 110 consider the continuation by 0 to $\partial \Omega_1$). 111

Since meas_{$\partial\Omega$} ($\partial\Omega_2 \setminus \Sigma$) > 0, all elements of V₂(Ω_2) fulfill a homogeneous boundary condition on a part of 112the boundary $\partial \Omega_2$. On the other hand, one can check that if (u_1, u_2) is a solution to (2.1), then the function 113u defined by $u_{|\Omega_i|} = u_i$ for j = 1, 2 solves (1.1). The equations satisfied by u_1 and u_2 are elliptic but they 114 are coupled by the transmission conditions on Σ . As a consequence, we cannot solve them independently. 115The purpose of the next paragraph is to explain how to proceed to write an alternative formulation (an 116 optimization-based one), which can be solved via an iterative procedure such that at each step one has to solve 117 118 a set of elliptic problems.

3 The smooth extension method and optimal control reformulation of the problem The 119smooth extension method was proposed in [21] and can be considered as a special case of fictitious domain 120 methods (see [2]). It has been adapted to study the classical scalar transmission problem, i.e., with constant 121sign coefficients, in [19]. In this section, we explain how to apply it to our problem. 122

123 3.1Presentation of the smooth extension method The idea behind the smooth extension method is the following: instead of looking for $(u_1, u_2) \in V_1(\Omega) \times V_2(\Omega_2)$ solution to (2.1), we search for a pair of 124functions $(\tilde{u}, u_2) \in \mathrm{H}^1_0(\Omega) \times \mathrm{V}_2(\Omega_2)$ such that $(\tilde{u}_{|\Omega_1}, u_2)$ is a solution to (2.1).¹ The function \tilde{u} is then a 125particular continuous extension of u_1 to the whole domain Ω . The difficulty is to find a "good" way to define 126 the function \tilde{u} so that it can be approximated by the classical FEM. The function u_2 can then be approximated 127by solving the elliptic problem satisfied by u_2 in Ω_2 completed by $\tilde{u}_{|\Sigma}$ (resp. $\varepsilon_1 \partial_n \tilde{u}_{|\Sigma}$) as a Dirichlet (resp. 128Neumann) boundary condition on Σ . Note that at first sight the construction of such \tilde{u} is not straightforward. 129This will be achieved thanks to an optimal control reformulation of (2.1). This is the main goal of the next 130paragraph in which we also reformulate the idea presented above in a more rigorous way. 131

3.2 An optimal control reformulation of the problem Before getting into details, let us first 132 introduce $\tilde{\varepsilon}_1 \in L^{\infty}(\Omega)$ such that $\tilde{\varepsilon}_1 \geq \tilde{\varepsilon}^+ > 0$ a. e. in Ω and $\tilde{\varepsilon}_1 = \varepsilon_1$ in Ω_1 . Then, let $E : V_1(\Omega_1) \to H_0^1(\Omega)$ be 133 an arbitrary continuous extension operator. By making use of (2.1), one can show easily that 134

135

$$\int_{\Omega} \tilde{\varepsilon}_1 \nabla E(u_1) \cdot \nabla v \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_1} f_1 v \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla E(u_1) \cdot \nabla v \, \mathrm{d}\boldsymbol{x} - \langle \varepsilon_1 \partial_n u_1, v \rangle_{\Sigma} \quad \forall v \in \mathrm{H}_0^1(\Omega),$$
$$\int_{\Omega_2} \varepsilon_2 \nabla u_2 \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} f_2 \, v_2 \, \mathrm{d}\boldsymbol{x} + \langle \varepsilon_1 \partial_n u_1, v_2 \rangle_{\Sigma} \quad \forall v_2 \in \mathrm{V}_2(\Omega_2).$$

136

Here and elsewhere, $\langle \cdot, \cdot \rangle_{\Sigma}$ denotes the duality product between $(H_{00}^{1/2}(\Sigma))^*$ and $H_{00}^{1/2}(\Sigma)$. Now, given that the linear form $v_2 \mapsto \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla E(u_1) \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} - \langle \varepsilon_1 \partial_n u_1, v_2 \rangle_{\Sigma}$ is continuous on $V_2(\Omega_2)$ one can define, thanks to the Riesz representation theorem, for each $E(u_1)$ a unique $w_{E(u_1)} \in V_2(\Omega_2)$ such that 137138

139 (3.1)
$$\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla E(u_1) \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} - \langle \varepsilon_1 \partial_n u_1, v_2 \rangle_{\Sigma} = \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w_{E(u_1)} \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} \quad \forall v_2 \in \mathcal{V}_2(\Omega_2).$$

Above we have used the fact that $(u, v) \mapsto (\tilde{\varepsilon}_1 \nabla u, \nabla v)_{L^2(\Omega)^d}$ is an inner product on $V_2(\Omega_2)$. As a consequence, 140141 we have

$$\int_{\Omega} \tilde{\varepsilon}_1 \nabla E(u_1) \cdot \nabla v \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_1} f_1 v \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w_{E(u_1)} \cdot \nabla v \, \mathrm{d}\boldsymbol{x} \qquad \forall v \in \mathrm{H}^1_0(\Omega),$$
$$\int_{\Omega_2} \varepsilon_2 \nabla u_2 \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} f_2 \, v_2 \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla (E(u_1) - w_{E(u_1)}) \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} \quad \forall v_2 \in \mathrm{V}_2(\Omega_2).$$

Since the coefficients $\tilde{\varepsilon}_1$ and ε_2 have fixed signs, the forms

$$(u,v)\mapsto \int_{\Omega} \tilde{\varepsilon}_1 \nabla u \cdot \nabla v \, \mathrm{d}\boldsymbol{x} \text{ and } (u_2,v_2)\mapsto -\int_{\Omega_2} \varepsilon_2 \nabla u_2 \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x},$$

¹In the text below, we choose an extension from Ω_1 to $\Omega^* = \Omega_2$. Obviously, one could choose an extension from Ω_2 to $\Omega^* = \Omega_1$ so that $(u_1, \tilde{u}_{|\Omega_2})$ is a solution to (2.1). In this case, the condition meas $\partial_{\Omega}(\partial \Omega_1 \setminus \Sigma) > 0$ must hold.

are coercive, respectively on $H_0^1(\Omega)$ and on $V_2(\Omega_2)$. With this in mind, we define for all $w \in V_2(\Omega_2)$, the couple of functions $(u^w, u_2^w) \in H_0^1(\Omega) \times V_2(\Omega_2)$ that are solution to the well-posed system of equations:

145 (3.2)
$$\int_{\Omega} \tilde{\varepsilon}_1 \nabla u^w \cdot \nabla v \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_1} f_1 v \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w \cdot \nabla v \, \mathrm{d}\boldsymbol{x} \qquad \forall v \in \mathrm{H}_0^1(\Omega),$$
$$\int_{\Omega_2} \varepsilon_2 \nabla u_2^w \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} f_2 \, v_2 \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla (u^w - w) \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} \quad \forall v_2 \in \mathrm{V}_2(\Omega_2)$$

146 Well-posedness is achieved by solving the elliptic problem in u^w first, and then the elliptic problem in u_2^w .

147 Remark 3.1. Observe that if we choose $w = w_{E(u_1)}$, then it follows that $(u^w|_{\Omega_1}, u_2^w)$ is the solution to (2.1). 148 Indeed, one finds first that $u^w = E(u_1)$, and then that $u_2^w = u_2$.

149 Given any auxiliary "control" function w, the solutions to (3.2) enjoy the properties listed below.

150 PROPOSITION 3.2. For all $w \in V_2(\Omega_2)$, the functions $u_1^w := u^w|_{\Omega_1}$ and u_2^w are such that

151
$$\begin{aligned} -\operatorname{div}(\varepsilon_{1}\nabla u_{1}^{w}) &= f_{1} \quad in \ \Omega_{1}, \\ -\operatorname{div}(\varepsilon_{2}\nabla u_{2}^{w}) &= f_{2} \quad in \ \Omega_{2} \\ \varepsilon_{1}\partial_{n}u_{1}^{w} &= \varepsilon_{2} \ \partial_{n}u_{2}^{w} \quad on \ \Sigma. \end{aligned}$$

Remark 3.3. In other words, the introduction of an auxiliary "control" function w allows us to construct pseudo-solutions to the equation (2.1) for which the condition on the normal derivatives is automatically satisfied. However we do not have in general continuity across the interface.

155 *Proof.* Take $\varphi_1 \in \mathscr{C}_0^{\infty}(\Omega_1)$ and extend it by 0 to the whole Ω to obtain the function $\varphi \in \mathscr{C}_0^{\infty}(\Omega)$. Take $v = \varphi$ 156 in the problem satisfied by u^w . One finds that $-\operatorname{div}(\varepsilon_1 \nabla u_1^w) = f_1$ in Ω_1 . Next, take some $\varphi_2 \in \mathscr{C}_0^{\infty}(\Omega_2)$, 157 extend it by 0 in Ω_1 and denote by φ the new function. By taking $v = \varphi$ in the problem satisfied by u^w and 158 $v_2 = \varphi_2$ in the problem satisfied by u_2^w one finds that

$$\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla u^w \cdot \nabla \varphi_2 \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w \cdot \nabla \varphi_2 \, \mathrm{d}\boldsymbol{x}, \int_{\Omega_2} \varepsilon_2 \nabla u_2^w \cdot \nabla \varphi_2 \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} f_2 \varphi_2 \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla (u^w - w) \cdot \nabla \varphi_2 \, \mathrm{d}\boldsymbol{x}.$$

By considering the sum of the two formulations, we conclude that $-\operatorname{div}(\varepsilon_2 \nabla u_2^w) = f_2$ in Ω_2 . To end the proof, it remains to show that $\varepsilon_1 \partial_n u^w = \varepsilon_2 \partial_n u_2^w$. For this, let $v \in \mathrm{H}^1_0(\Omega)$ and define $v_2 = v_{|\Omega_2|} \in \mathrm{V}_2(\Omega_2)$. By taking v and v_2 as test functions in (3.2), considering the sum of the two equations, integrating by parts in both formulations and then, using the equations satisfied by u_1^w and u_2^w , we infer that

164
$$-\langle \varepsilon_1 \partial_n u_1^w, v \rangle_{\Sigma} = -\langle \varepsilon_2 \partial_n u_2^w, v \rangle_{\Sigma}, \qquad v \in \mathrm{H}^1_0(\Omega)$$

According to the surjectivity of the trace mapping on Σ , this gives $\varepsilon_1 \partial_n u_1^w = \varepsilon_2 \partial_n u_2^w$ on Σ .

166 It follows that

159

167 LEMMA 3.4. If there exists $w^* \in V_2(\Omega_2)$ such that the solution to (3.2) satisfies $u^{w^*} = u_2^{w^*}$ on Σ , then 168 $(u^{w^*}|_{\Omega_1}, u_2^{w^*})$ solves (2.1).

169 Thanks to what we have explained in Remark 3.1, we know that to every continuous extension of u_1 to Ω , one 170 can define $w^* \in V_2(\Omega_2)$ for which $u^{w^*} = u_2^{w^*}$ on Σ . This leads us to the following result.

171 LEMMA 3.5. Let u_1 be the first part of the solution to (2.1). Then, the set of $w^* \in V_2(\Omega_2)$ such that the 172 solution to (3.2) satisfies $u^{w^*} = u_2^{w^*}$ on Σ is isomorphic to the set of all possible continuous extensions of u_1 173 to Ω . Furthermore, w^* and u^{w^*} are linked by relation

174 (3.3)
$$\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla u^{w^*} \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} - \langle \varepsilon_1 \partial_n u_1, v_2 \rangle_{\Sigma} = \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w^* \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} \quad \forall v_2 \in \mathrm{V}_2(\Omega_2).$$

Now, we have all the tools to introduce the optimal control reformulation of the problem (2.1). As a matter of fact, in order to find a function $w^* \in V_2(\Omega_2)$ for which $u^{w^*} = u_2^{w^*}$ on Σ , it is enough to solve the following

177 optimal control problem:

178 (3.4) Find
$$w^* = \underset{w \in V_2(\Omega_2)}{\operatorname{argmin}} J(w)$$
 with $J(w) = \frac{1}{2} \int_{\Sigma} |u^w - u_2^w|^2 \, \mathrm{d}\sigma$,

where $(u^w, u_2^w) \in H_0^1(\Omega) \times V_2(\Omega_2)$ is the solution to (3.2). Note that in (3.4), the functional J plays the role of the cost functional, while (3.2) plays the role of the state equation. Obviously, thanks to Lemma 3.5, the problem (3.4) has an infinite number of solutions. As a result, one may need to use a regularization technique in order to be able to construct a convergent discretization of the problem: this will be the subject of §4.3 where we will study the classical Tikhonov regularization method applied to Problem (3.4).

4 Basic properties of the optimization problem and its regularization In this section, we present in §4.1 some useful properties of the cost functional J and of the set of its minimizers in §4.2. After that in §4.3, we study the Tikhonov regularization of the problem. Furthermore, we explain, in §4.4, how to use the the adjoint approach in order to find an explicit expression of the gradient of J.

4.1 Properties of the cost functional Since we have used the $L^2(\Sigma)$ norm instead of the $H^{1/2}(\Sigma)$ norm in the definition of J, one has the following results.

190 PROPOSITION 4.1. The cost functional J satisfies the following properties:

191 1. Let $(w_n)_n$ be a sequence of elements of $V_2(\Omega_2)$ that converges weakly to $w_0 \in V_2(\Omega_2)$. Then, $(J(w_n))_n$ 192 converges to $J(w_0)$.

193 2. The functional J is continuous and convex on $V_2(\Omega_2)$.

194 Proof. 1. For all $n \in \mathbb{N}$, denote by $(u^n, u_2^n) \in \mathrm{H}_0^1(\Omega) \times \mathrm{V}_2(\Omega)$ the solution to (3.2) with $w = w_n$. From the 195 ellipticity of the problems involved in (3.2), it follows that $(u^n)_n$ (resp. $(u_2^n)_n$) converges weakly in $\mathrm{H}_0^1(\Omega)$ 196 (resp. $\mathrm{V}_2(\Omega_2)$) to some $u \in \mathrm{H}^1(\Omega)$ (resp. $u_2 \in \mathrm{V}_2(\Omega_2)$) such that (u, u_2) is the solution to (3.2) with $w = w_0$. 197 The continuity of the trace operator from $\mathrm{H}^1(\Omega)$ to $\mathrm{H}^{1/2}(\Sigma)$ ensures that $(u_{|\Sigma}^n - u_2^n|_{\Sigma})_n$ converges weakly to 198 $u_{|\Sigma} - u_{2|\Sigma}$ in $\mathrm{H}^{1/2}(\Sigma)$. Given that the embedding of $\mathrm{H}^{1/2}(\Sigma)$ into $\mathrm{L}^2(\Sigma)$ is compact, it actually converges 199 strongly to $u_{|\Sigma} - u_{2|\Sigma}$ in $\mathrm{L}^2(\Sigma)$. Thus $(J(w_n))$ converges to $J(w_0)$. The result is proved.

2. While the continuity is a direct consequence of the first item, the convexity follows from the fact that J: 20. $V_2(\Omega_2) \to \mathbb{R}$ is the composition of the affine map $j_1 : V_2(\Omega_2) \to L^2(\Sigma)$ and of the convex map $j_2 : L^2(\Sigma) \to \mathbb{R}$ 20. such that for all $w \in V_2(\Omega_2), g \in L^2(\Sigma)$ we have

(4.1)
$$\begin{aligned} j_1(w) &= (u^w - u_2^w)_{|\Sigma} \text{ where } (u^w, u_2^w) \in \mathrm{H}^1_0(\Omega) \times \mathrm{V}_2(\Omega_2) \text{ is the solution to } (3.2), \\ j_2(g) &= \frac{1}{2} \int_{\Sigma} |g|^2 \, \mathrm{d}\sigma. \end{aligned}$$

4.2 The set of minimizers of the functional J Thanks to Lemma 3.5, we know that J has an infinite number of minimizers. This (non-empty) set will be denoted by M_J . Without any difficulty, one can see that M_J coincides with the set of zeros of the functional J. As a result, since J is continuous, convex and positive, the set M_J is closed and convex in $V_2(\Omega_2)$. This allows us to say that the following minimization problem:

$$\min_{w\in M_J}\int_{\Omega_2} ilde{arepsilon}_1|
abla w|^2\;\mathrm{d}oldsymbol{x}$$

203

209

has a unique solution, as a consequence of the strict convexity of $v_2 \mapsto \int_{\Omega_2} \tilde{\varepsilon}_1 |\nabla v_2|^2 \, \mathrm{d} \boldsymbol{x}$ in $V_2(\Omega_2)$, and of the fact that M_J is a closed, convex subset of $V_2(\Omega_2)$. In the following, we shall denote by w_J^* the smallest minimizer of the functional J:

213 (4.2)
$$w_J^* = \operatorname*{argmin}_{w \in M_J} \int_{\Omega_2} \tilde{\varepsilon}_1 |\nabla w|^2 \, \mathrm{d}\boldsymbol{x}.$$

The goal of the rest of this paragraph is to find a characterization of $E_{w_J^*}(u_1)$, the continuous extension of u_1 that is associated with w_J^* . Note that the link between $E_{w_J^*}(u_1)$ and w_J^* is given by the following (see relation (3.3)):

217 (4.3)
$$\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla E_{w_J^*}(u_1) \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} - \langle \varepsilon_1 \partial_n u_1, v_2 \rangle_{\Sigma} = \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w_J^* \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} \quad \forall v_2 \in \mathcal{V}_2(\Omega_2).$$

To proceed, we define $E_H(u_1) \in H^1_0(\Omega)$ the continuous extension of u_1 that satisfies

219 (4.4)
$$\operatorname{div}(\tilde{\varepsilon}_1 \nabla E_H(u_1)) = 0 \text{ in } \Omega_2.$$

220 In particular, we have

$$\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla E_H(u_1) \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} = 0 \quad \forall v_2 \in \mathrm{H}^1_0(\Omega_2).$$

222 Denote by $w_H \in M_J$ the minimizer associated with $E_H(u_1)$. Thanks to (3.3), we know that

223 (4.5)
$$\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla E_H(u_1) \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} - \langle \varepsilon_1 \partial_n u_1, v_2 \rangle_{\Sigma} = \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w_H \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} \quad \forall v_2 \in \mathrm{V}_2(\Omega_2).$$

We infer that

225 (4.6)
$$\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w_H \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} = 0 \quad \forall v_2 \in \mathrm{H}^1_0(\Omega_2).$$

By taking the difference between (4.3) and (4.5), taking $v_2 = w_H$, using the fact that $E_H(u_1) - E_{w_J^*}(u_1) \in H_0^1(\Omega_2)$ and owing to (4.6), we infer that

221

$$\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla (w_H - w_J^*) \cdot \nabla w_H \, \mathrm{d}\boldsymbol{x} = 0,$$

229 so230

$$\int_{\Omega_2} \tilde{\varepsilon}_1 |\nabla w_J^*|^2 \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} \tilde{\varepsilon}_1 |\nabla w_J^* - \nabla w_H|^2 \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 |\nabla w_H|^2 \, \mathrm{d}\boldsymbol{x} \ge \int_{\Omega_2} \tilde{\varepsilon}_1 |\nabla w_H|^2 \, \mathrm{d}\boldsymbol{x}$$

231 Hence, from the definition of w_J^* , we then obtain the following

232 PROPOSITION 4.2. The functions w_H and w_J^* coincide.

Remark 4.3. It is worth noting that, thanks to (4.5) and using the definition of $E_H(u_1)$, the function w_H satisfies the problem:

235 (4.7)
$$\operatorname{div}(\tilde{\varepsilon} \nabla w_H) = 0 \text{ in } \Omega_2 \text{ and } \tilde{\varepsilon}_1 \partial_n w_{H|\Sigma} = \tilde{\varepsilon}_1 \partial_n E_H(u_1)_{|\Sigma} - \varepsilon_1 \partial_n u_{1|\Sigma}.$$

236 Recall that \boldsymbol{n} is the unit normal vector to $\boldsymbol{\Sigma}$ oriented to the exterior of Ω_2 .

4.3 Tikhonov regularization of the problem Tikhonov regularization, which was originally introduced in [25], is a classical method to regularize a convex optimization problem. Classically, this method is used in the context of regularization of ill-posed inverse problems (see [20] and the references therein). In this paragraph, we study the convergence of such regularization when it is applied to our problem. For $\delta > 0$, we introduce the functional $J^{\delta} : V_2(\Omega_2) \to \mathbb{R}$ defined by

242
$$J^{\delta}(w) = J(w) + \delta \int_{\Omega_2} \tilde{\varepsilon}_1 |\nabla w|^2 \, \mathrm{d}\boldsymbol{x} \qquad \forall w \in \mathrm{V}_2(\Omega_2).$$

243 To simplify notation, we will denote by $\|\cdot\|_{\tilde{\varepsilon}_1}: V_2(\Omega_2) \to \mathbb{R}_+$ the norm that is defined as follows:

244
$$\|w\|_{\tilde{\varepsilon}_1} := \left(\int_{\Omega_2} \tilde{\varepsilon}_1 |\nabla w|^2 \, \mathrm{d}\boldsymbol{x}\right)^{1/2}, \quad \forall w \in \mathcal{V}_2(\Omega_2)$$

Endowed with the associated scalar product $(\cdot, \cdot)_{\tilde{\varepsilon}_1}$, the space $V_2(\Omega_2)$ is a Hilbert space. Since J is convex and $\delta > 0$, the functional J^{δ} is strictly convex and coercive. Therefore the minimization problem

247 $\min_{w \in \nabla_2(\Omega_2)} J^{\delta}(w)$

has a unique solution that we denote by w_{δ}^* . Our goal is to study the behaviour of w_{δ}^* as δ tends to zero. One expects $(w_{\delta}^*)_{\delta}$ to converge to one of the solutions (3.4). If this is the case and because the problem (3.4) has an infinite number of solutions, it will be interesting to characterize the particular solution towards which $(w_{\delta}^*)_{\delta}$ converges. Our findings are summarized in the following

252 PROPOSITION 4.4. When $\delta \to 0$, the sequence $(w^*_{\delta})_{\delta}$ converges towards w^*_J , the smallest minimizer of J.

The proof of the previous result is quite classical. However, for the convenience of the reader, we will detail it in Appendix A.

In conclusion, we can say that the Tikhonov regularization allows us to obtain a stabilized version of the

- optimization problem (3.4). This will be used in order to introduce a stabilization of the discretization of the
- problem (3.4), but in that case the stabilization parameter δ will be chosen as a function of the discretization

parameter. This will be detailed in 5.3. Note that the same idea was employed in 1.

4.4 Gradient of the functional J As indicated in the introduction, the main objective of this work is to propose a new numerical method for approximating the solution to (1.1). This method will be based on the numerical approximation of the solution to the optimization problem (3.4). In this section, we will explain how to obtain an explicit expression of J'(w) the gradient of J at some $w \in V_2(\Omega)$. We recall that the functional J is differentiable, because it can be written as a composition of the two differentiable maps j_1 and j_2 , cf. (4.1). Since the functional J is scalar valued, its differential at $w \in V_2(\Omega_2)$ can be represented by its gradient $J'(w) \in V_2(\Omega_2)$:

266 For all
$$h \in V_2(\Omega_2)$$
, $\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla J'(w) \cdot \nabla h \, \mathrm{d}\boldsymbol{x} = \lim_{t \to 0} \frac{J(w+th) - J(w)}{t}$.

To find an explicit expression of J'(w), we use the adjoint approach [11]. Details about the application of this approach to our problem are given in Appendix B (see also [19]). Here, we present final result. To do so, we start by introducing the so-called adjoint equations. For all $w \in V_2(\Omega_2)$, recalling that $(u^w, u_2^w) \in H_0^1(\Omega) \times V_2(\Omega_2)$ is the solution to (3.2), we introduce $(g^w, g_2^w) \in H_0^1(\Omega) \times V_2(\Omega_2)$ such that

271 (4.8)
$$\int_{\Omega} \tilde{\varepsilon}_1 \nabla g^w \cdot \nabla v \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla g_2^w \cdot \nabla v \, \mathrm{d}\boldsymbol{x} - \int_{\Sigma} (u^w - u_2^w) v \, \mathrm{d}\boldsymbol{\sigma} \qquad \forall v \in \mathrm{H}_0^1(\Omega)$$
$$\int_{\Omega_2} \varepsilon_2 \nabla g_2^w \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} = \int_{\Sigma} (u^w - u_2^w) v_2 \, \mathrm{d}\boldsymbol{\sigma} \qquad \forall v_2 \in \mathrm{V}_2(\Omega_2)$$

272 As observed before, the functions g^w , g_2^w are well-defined. In Appendix B, we prove the

273 LEMMA 4.5. For all $w \in V_2(\Omega_2)$, there holds $J'(w) = g_2^w - g_{|\Omega_2}^w$, where (g^w, g_2^w) solve (4.8).

274 We have the following optimality result

275 COROLLARY 4.6. We have the equivalence

276
$$\left[w^* \in \mathcal{V}_2(\Omega_2) \text{ is such that } J'(w^*) = 0\right] \iff w^* \in M_J$$

277 Proof. Let us start with the proof of the direct implication. Suppose that there exists some $w^* \in V_2(\Omega_2)$ such 278 that $g^{w^*}|_{\Omega_2} = g_2^{w^*}$. By taking the sum of the variational formulations of (4.8), we deduce that

279
$$\int_{\Omega} \varepsilon \nabla g^{w^*} \cdot \nabla v \, \mathrm{d}\boldsymbol{x} = 0 \qquad \forall v \in \mathrm{H}^1_0(\Omega).$$

This means $A_{\varepsilon}(g^{w^*}) = 0$ and then, thanks to Assumption 2.1, $g^{w^*} = 0$. This implies that $g_2^{w^*} = 0$ and then by using the second equation of (4.8), that $u^{w^*} = u_2^{w^*}$ on Σ . This shows that w^* is a minimizer of J. The reverse implication is a consequence of the fact that if $w^* \in M_J$ we have $J(w^*) = 0$ and then $u^{w^*} = u_2^{w^*}$ on Σ . This implies that $g_2^{w^*} = 0$ and that $g^{w^*} = 0$.

We end this paragraph with the following result that can be useful to prove the convergence of the classical gradient descent algorithm.

286 COROLLARY 4.7. The functional $J': V_2(\Omega_2) \to V_2(\Omega_2)$ is Lipschitz continuous.

Proof. Starting from (3.2), we deduce that $w \mapsto u^w$, $w \mapsto u_2^w$ are Lipschitz continuous. Inserting this into (4.8), we obtain the result.

5 Numerical discretization of the problem In this part, we are concerned with the numerical approximation of (3.4) by means of the Finite Element Method. To do so, we start by presenting some details and notations about the family of meshes that will be used. To simplify the presentation, the domain Ω and the subdomains $(\Omega_i)_{i=1,2}$ are supposed to have polygonal (when d = 2) or polyhedral (resp. d = 3) boundaries.

5.1 Meshes and discrete spaces Let $(\mathcal{T}_h)_h$ be a regular family of meshes of $\overline{\Omega}$ (see [15]), composed of (closed) simplices. The subscript $_h$ stands for the meshsize.

Assumption 5.1. We suppose that for all h, every simplex of \mathcal{T}_h belongs either to $\overline{\Omega_1}$ or to $\overline{\Omega_2}$.

According to Assumption 5.1, for i = 1, 2, one can consider the family of meshes $(\mathcal{T}_h^i)_h$ made of those simplices that belong to $\overline{\Omega_i}$.

298 For all $k \in \mathbb{N}^*$, we set

$$\mathbf{V}_h^k(\Omega) := \{ v_h \in \mathbf{H}_0^1(\Omega) \mid v_{h|T} \in \mathbf{P}^k(T), \ \forall T \in \mathfrak{T}_h \}.$$

Here $P^k(T)$ stands for the space of polynomials (of *d* variables) defined on *T* of degree at most equal to *k*. In the same way, we define for i = 1, 2,

302
$$\mathbf{V}_h^k(\Omega_i) := \{ v_{i,h} \in \mathbf{V}_i(\Omega_i) \mid v_{i,h}|_T \in \mathbf{P}^k(T), \forall T \in \mathbf{T}_h^i \}.$$

303

304 Remark 5.2. Note that for all h > 0, for i = 1, 2 the space $V_h^k(\Omega_i)$ coincides with $\{u_{|\Omega_i|} | u \in V_h^k(\Omega)\}$.

³⁰⁵ Finally, we recall the basic approximability properties

306 (5.1)
$$\forall v \in H_0^1(\Omega), \quad \lim_{h \to 0} \left(\inf_{\substack{v_h \in V_h^k(\Omega)}} \|v - v_h\|_{H_0^1(\Omega)} \right) = 0, \\ \forall v_2 \in V_2(\Omega_2), \quad \lim_{h \to 0} \left(\inf_{\substack{v_{2,h} \in V_h^k(\Omega_2)}} \|v_2 - v_{2,h}\|_{\tilde{\varepsilon}_1} \right) = 0.$$

307

325

308 **5.2** Discretization strategy For h > 0 and $w \in V_2(\Omega)$, define the functions $u_h^w \in V_h^k(\Omega)$ and $u_{2,h}^w \in V_h^k(\Omega_2)$ as the solutions to the following well-posed discrete problems:

(5.2)
$$\int_{\Omega} \tilde{\varepsilon}_1 \nabla u_h^w \cdot \nabla v_h \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_1} f v_h \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w \cdot \nabla v_h \, \mathrm{d}\boldsymbol{x}, \; \forall v_h \in \mathcal{V}_h^k(\Omega)$$
$$\int_{\Omega_2} \varepsilon_2 \nabla u_{2,h}^w \cdot \nabla v_{2,h} \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} f_2 v_{2,h} \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla (u_h^w - w) \cdot \nabla v_{2,h} \, \mathrm{d}\boldsymbol{x}, \; \forall v_{2,h} \in \mathcal{V}_h^k(\Omega_2).$$

Then introduce the projection operator $\pi_h^k : V_2(\Omega_2) \to V_h^k(\Omega_2)$ such that for all $w \in V_2(\Omega_2)$, $\pi_h^k w$ is defined as the unique element of $V_h^k(\Omega_2)$ that satisfies the problem

313
$$\int_{\Omega_2} \tilde{\varepsilon}_1 \nabla(\pi_h^k w) \cdot \nabla v_{2,h} \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w \cdot \nabla v_{2,h} \, \mathrm{d}\boldsymbol{x} \quad \forall v_{2,h} \in \mathrm{V}_h^k(\Omega_2).$$

314 Obviously, one has the estimate

315 (5.3)
$$\|\pi_h^k w\|_{\tilde{\varepsilon}_1} \le \|w\|_{\tilde{\varepsilon}_1}$$

From the definition of $\pi_h^k w$, one can easily see that for all $w \in V_2(\Omega_2)$ we have the identities

317 (5.4)
$$u_h^{\pi_h^k w} = u_h^w$$
 and $u_{2,h}^{\pi_h^k w} = u_{2,h}^w$.

Now, let us turn our attention to the discretization of the optimization problem (3.4). The natural way to do that is to replace it by the problem

320 (5.5)
$$\inf_{w_h \in \mathcal{V}_h^k(\Omega_2)} J_{0,h}(w_h) := \frac{1}{2} \int_{\Sigma} |u_h^{w_h} - u_{2,h}^{w_h}|^2 \, \mathrm{d}\sigma$$

One can proceed as in the proof of proposition 4.1 to show that the cost functional $J_{0,h} : V_h^k \to \mathbb{R}$ (defined in (5.5)) is convex and continuous. Unfortunately this result is not sufficient to justify that the problem (5.5) is well-posed. The difficulty comes from the fact that, even under Assumption 2.1, one can not guarantee that the problem

Find
$$u_h \in \mathcal{V}_h^k(\Omega)$$
 such that $\int_{\Omega} \varepsilon \nabla u_h \cdot \nabla v_h \, \mathrm{d}\boldsymbol{x} = \int_{\Omega} f v_h \, \mathrm{d}\boldsymbol{x} \quad \forall v_h \in \mathcal{V}_h^k(\Omega)$

is well-posed even for h small enough. To cope with this difficulty, an idea is to use again the Tikhonov regularization (see §4.3), with a regularization parameter that depends now on h. This idea was originally proposed in [22] for the case of elliptic equations and then, was used by Assyr Abdulle et al. in [1] for the case of problems with sign-changing coefficients. Here, we explain how to adapt it to our approach. The idea is to replace the cost functional $J_{0,h}$ in (5.5) by the functional $J_h: V_h^k(\Omega_2) \to \mathbb{R}_+$ such that for all $w_h \in V_h^k(\Omega_2)$, we have

332
$$J_h(w_h) := \frac{1}{2} \int_{\Sigma} |u_h^{w_h} - u_{2,h}^{w_h}|^2 \, \mathrm{d}\sigma + \lambda_h ||w_h||_{\tilde{\varepsilon}_1}^2,$$

where $\lambda_h > 0$ tends to zero as h goes to 0. Since $\lambda_h > 0$ for all h > 0, the functional J_h is strictly convex and 333 coercive. This guarantees that the optimization problem 334

335 (5.6)
$$\min_{w_h \in \mathbf{V}_h^k(\Omega_2)} J_h(w_h)$$

has a unique solution that we denote by $w_{k,h}^*$. All the difficulty now is to choose the parameter λ_h in order to 336 be able to ensure the convergence of $(w_{k,h}^*)_h$ towards a solution to (3.4) as h tends to zero. This is the main 337 goal of the next paragraph. 338

Convergence of the method The starting point of our discussion is the following 339 5.3

340 LEMMA 5.3. We have the estimate

341 (5.7)
$$J_h(w_{k,h}^*) \le \frac{1}{2} \int_{\Sigma} |u_h^{w_J^*} - u_{2,h}^{w_J^*}|^2 \, \mathrm{d}\sigma + \lambda_h ||w_J^*||_{\tilde{\varepsilon}_1}^2$$

where w_J^* is defined by (4.2). 342

349

371

Proof. Starting from the fact that $\pi_h^k w_J^* \in V_h^k(\Omega_2)$ and using that $w_{k,h}^*$ is the unique solution to the optimiza-343 tion problem (5.6), we conclude that $J_h(w_{k,h}^*) \leq J_h(\pi_h^k w_J^*)$. On the other hand, the identity (5.4) allows us to 344

write 345 1

346
$$J_h(\pi_h^k w_J^*) = \frac{1}{2} \int_{\Sigma} |u_h^{w_J^*} - u_{2,h}^{w_J^*}|^2 \, \mathrm{d}\sigma + \lambda_h \|\pi_h^k w_J^*\|_{\tilde{\varepsilon}_1}^2.$$

The Lemma is then proved by recalling the estimate (5.3).

In order to simplify notations, for h > 0 and $w \in V_2(\Omega_2)$, we denote by $A_h(w)$ the real number 348

$$\mathcal{A}_h(w) = \frac{1}{2} \int_{\Sigma} |u_h^w - u_{2,h}^w|^2 \, \mathrm{d}\sigma.$$

From (5.4), we know that for all $w \in V_2(\Omega_2)$, we have $A_h(w) = J_0^h(\pi_h^k w)$. The main result of this paragraph 350 is the following theorem. 351

THEOREM 5.4. Assume that the parameter λ_h can be chosen such that the sequences $(\lambda_h)_h$ and $(A_h(w_J^*)/\lambda_h)_h$ 352 tend to zero as h tends to zero. Then, as h goes to 0: 353

- the sequence $(w_{k,h}^*)_h$ converges to w_J^* in $V_2(\Omega_2)$; 354
- the sequence $(u_h^{w_{k,h}^*})_h$ converges to $E_H(u_1)$ in $\mathrm{H}_0^1(\Omega)$, resp. the sequence $(u_{2,h}^{w_{k,h}^*})_h$ converges to u_2 in $\mathrm{V}_2(\Omega_2)$, where (u_1, u_2) is the solution to (2.1) and $E_H(u_1)$ is the extension of u_1 defined in (4.4). 355 356

Proof. The strategy of proof is similar to the one of proposition 4.4. To simplify notations, we denote by $u^{k,h} \in V_h^k(\Omega)$ and $u_2^{k,h} \in V_h^k(\Omega_2)$ the functions 357 358

359
$$u^{k,h} = u_h^{w_{k,h}^*}$$
 and $u_2^{k,h} = u_{2,h}^{w_{k,h}^*}$

360

In order to make the proof as clear as possible, we divide it into four steps. Step 1: weak convergence of $(w_{k,h}^*)_h$, $(u^{k,h})_h$ and $(u_2^{k,h})_h$. Starting from the estimate 361

362 (5.8)
$$\|w_{k,h}^*\|_{\tilde{\varepsilon}_1}^2 \le J_h(w_{k,h}^*)/\lambda_h \le A_h(w_J^*)/\lambda_h + \|w_J^*\|_{\tilde{\varepsilon}_1}^2$$

and using the fact that $(A_h(w_J^*)/\lambda_h)_h$ tends to 0 as h goes to 0, we infer that $(w_{k,h}^*)_h$ is bounded in $V_2(\Omega_2)$. 363 This implies that, up to a sub-sequence, $(w_{k,h}^*)_h$ converges weakly to some $w_0 \in V_2(\Omega)$. For the reader's 364 convenience, this sub-sequence is still denoted by $(w_{k,h}^*)_h$. 365

Since the problems in (5.2) are uniformly elliptic with respect to h, we know that the sequence $(u^{k,h})_h$ (resp. 366 $(u_2^{k,h})_h)$ converges weakly in $\mathrm{H}_0^1(\Omega)$ (resp. in $\mathrm{V}_2(\Omega_2)$) to some $u \in \mathrm{H}_0^1(\Omega)$ (resp. $u_2 \in \mathrm{V}_2(\Omega_2)$). Using the basic approximability property (5.1), we infer that $u = u^{w_0}$ and $u_2 = u_2^{w_0}$. 367 368

Step 2: w_0 is a minimizer of J. The continuity of the trace operator and the compactness of the embedding 369 $\mathrm{H}^{1/2}(\Sigma) \subset \mathrm{L}^2(\Sigma)$ ensure that 370

$${u^{k,h}}_{|\Sigma} - {u^{k,h}_2}_{|\Sigma} \to {u^{w_0}}_{|\Sigma} - {u^{w_0}_2}_{|\Sigma}$$

372 in $L^2(\Sigma)$ as $h \to 0$. By noticing that

373
$$\frac{1}{2} \int_{\Sigma} |u^{k,h} - u_2^{k,h}|^2 \, \mathrm{d}\sigma = J_0^h(w_{k,h}^*) \le J_h(w_{k,h}^*) \le \lambda_h(\mathrm{A}_h(w_J^*)/\lambda_h + \|w_J^*\|_{\tilde{\varepsilon}_1}^2)$$

and using that $\lambda_h, A_h(w_J^*)/\lambda_h \to 0$ as h goes to zero, we deduce that $u^{w_0}|_{\Sigma} - u_2^{w_0}|_{\Sigma} = 0$. This shows that w_0 is a minimizer of J.

Step 3: strong convergence of $(w_{k,h}^*)_h$ to w_J^* . Thanks to the fact that $A_h(w_J^*)/\lambda_h \to 0$ as $h \to 0$ and by means of the estimate (5.8), we can write

$$\limsup_{h \to 0} \|w_{k,h}^*\|_{\tilde{\varepsilon}_1} \le \|w_j^*\|_{\tilde{\varepsilon}_1}.$$

379 On the other hand, since $(w_{k,h}^*)_h$ converges weakly to w_0 as $h \to 0$, we infer that

380
$$\|w_0\|_{\tilde{\varepsilon}_1} \le \liminf_{k \to 0} \|w_{k,h}^*\|_{\tilde{\varepsilon}_1},$$

which is a consequence of the fact that the norm of a Banach space is weakly lower semicontinuous, see [9, Proposition III.5 (iii)]. This implies that $||w_0||_{\tilde{\varepsilon}_1} \leq ||w_J^*||_{\tilde{\varepsilon}_1}$. Since w_0 is a minimizer of J, we conclude that $w_0 = w_J^*$. Furthermore, we also deduce that

$$\lim_{k \to \infty} \|w_{k,h}^*\|_{\tilde{\varepsilon}_1} = \|w_0\|_{\tilde{\varepsilon}_1}$$

As a result, by applying [9, Proposition III.32], we infer that $(w_{k,h}^*)_h$ converges, strongly, in $V_2(\Omega_2)$ to $w_0 = w_J^*$. **Step 4: strong convergence of** $(u^{k,h})_h$ and $(u_2^{k,h})_h$. The ellipticity of the problems in (3.2), combined with the strong convergence of $(w_{k,h}^*)_h$ to w_J^* , imply the convergence of $(u^{k,h})_h$ in $H_0^1(\Omega)$ to $u^{w_J^*}$ and of $(u_2^{k,h})_h$ in $V_2(\Omega_2)$ to $u_2^{w_J^*}$.

The result is then proved by using that $u^{w_J^*} = E_H(u_1)$ and by observing that these limits are independent of the chosen sub-sequences.

The rest of this paragraph is devoted to explain why it is possible to choose the parameter λ_h in such a way that $(\lambda_h)_h$ and $(A_h(w_J^*)/\lambda_h)_h$ both converge to 0 as h tends to 0. To do so, one needs to study the behaviour of $A_h(w_J^*)$ as h tends to 0. We recall that, according to theorem 2.2, we know that $u \in \bigcap_{s \in [0,\sigma_D(\varepsilon))} PH^{1+s}(\Omega)$, where $\sigma_D(\varepsilon) \in (0,1]$ is the so-called regularity exponent. Let us start with the following

PROPOSITION 5.5. Suppose that the coefficients $\tilde{\varepsilon}_1$ and ε_2 are smooth, or piecewise smooth. Assume that the solution u to (1.1) belongs to $\text{PH}^{1+s}(\Omega)$ for some s > 0. Then there exists $s' \in (0, s]$ that depends only on the geometry of Ω_2 and on the coefficient ε_2 , and there exists $\sigma \in (0, 1]$ that depends only on the geometry of Ω and of Ω_2 such that

378

$$\begin{aligned} \|u^{w_{J}^{*}} - u_{h}^{w_{J}^{*}}\|_{\mathrm{H}_{0}^{1}(\Omega)} &\leq Ch^{p'} \|u\|_{\mathrm{PH}^{1+p'}(\Omega)} \quad and \quad \|u_{2}^{w_{J}^{*}} - u_{2,h}^{w_{J}^{*}}\|_{\tilde{\varepsilon}_{1}} \leq Ch^{p'} \|u_{2}\|_{\mathrm{H}^{1+p'}(\Omega_{2})}, \\ \|u^{w_{J}^{*}} - u_{h}^{w_{J}^{*}}\|_{\mathrm{L}^{2}(\Omega)} &\leq Ch^{p'+\sigma} \|u\|_{\mathrm{PH}^{1+p'}(\Omega)} \quad and \quad \|u_{2}^{w_{J}^{*}} - u_{2,h}^{w_{J}^{*}}\|_{\mathrm{L}^{2}(\Omega_{2})} \leq Ch^{p'+\sigma} \|u_{2}\|_{\mathrm{H}^{1+p'}(\Omega_{2})}. \end{aligned}$$

- 400 with C independent of h and $p' = \min(s', k)$.
- 401 *Proof.* Along this proof, C denotes a positive constant whose value can change from one line to the next but 402 does not depend on h.

Given that $u^{w_j^*} = E_H(u_1)$ solves (4.4), and since $u_1 \in \mathrm{H}^{1+s}(\Omega_1)$, it follows that $E_H(u_1)_{|\Omega_2}$ exhibits some extra-regularity because ε_2 is (piecewise) smooth (via a classical shift theorem). In other words, there exists $s' \in (0, s]$ such that $u^{w_j^*} \in \mathrm{PH}^{1+s'}(\Omega)$.

Given that $u_2^{w_j^*} = u_2 \in PH^{1+s}(\Omega_2) \subset PH^{1+s'}(\Omega_2)$ and since the problems in (3.2) are elliptic with (piecewise) smooth coefficients $\tilde{\varepsilon}_1$ and ε_2 , we obtain the estimates (see [15])

411

$$\|u^{w_{J}^{*}} - u_{h}^{w_{J}^{*}}\|_{\mathrm{H}_{0}^{1}(\Omega)} \leq Ch^{p'} \|u\|_{\mathrm{PH}^{1+p'}(\Omega)} \text{ and } \|u_{2}^{w_{J}^{*}} - u_{2,h}^{w_{J}^{*}}\|_{\tilde{\varepsilon}_{1}} \leq Ch^{p'} \|u_{2}\|_{\mathrm{H}^{1+p'}(\Omega_{2})},$$

where $p' = \min(s', k)$. By applying the classical Aubin-Nitsche's lemma (see [15, Theorem 3.2.4]), we infer that there exists $\sigma \in (0, 1]$ such that

$$\|u^{w_{J}^{*}} - u_{h}^{w_{J}^{*}}\|_{L^{2}(\Omega)} \leq Ch^{p'+\sigma} \|u\|_{\mathrm{PH}^{1+p'}(\Omega)} \text{ and } \|u_{2}^{w_{J}^{*}} - u_{2,h}^{w_{J}^{*}}\|_{L^{2}(\Omega_{2})} \leq Ch^{p'+\sigma} \|u_{2}\|_{\mathrm{H}^{1+p'}(\Omega_{2})}.$$

- Remark 5.6. It is worth to note that the value of s' is prescribed by the regularity of $E_H(u_1)$, the harmonic-like 412
- extension u_1 that satisfies (4.4). More precisely, s' depends both on the regularity exponent $\sigma_D(\tilde{\varepsilon}_1) \in (0,1]$, 413 and on the regularity of the boundary data on $\partial \Omega_2$, because $u_{2H} := E_H(u_1)_{|\Omega_2} \in V_2(\Omega_2)$ solves the Dirichlet
- 414
- problem (cf. (4.4)): $-\operatorname{div}(\tilde{\varepsilon}_1 \nabla u_{2H}) = 0$ in Ω_2 , with $u_{2H} = u_1$ in $H_{00}^{1/2}(\Sigma)$. Assume for instance that Ω and Ω_2 are convex domains, and that the coefficients ε_1 and ε_2 are constant. In this case, one can choose $\tilde{\varepsilon}_1$ 415 416
- to be constant over Ω . We recall that the solution u to (1.1) belongs to $\mathrm{PH}^{1+s}(\Omega)$ for all $s \in (0, \sigma_D(\varepsilon))$. 417

Then, because $u_{2H} \in \mathrm{H}^1(\Omega_2)$ is governed by: $-\Delta u_{2H} = 0$ in the convex domain Ω_2 with Dirichlet data in $\mathrm{H}^{1/2+s}(\partial\Omega_2)$, one has $u_{2H} \in \mathrm{H}^{1+s}(\Omega_2)$. In other words, s' = s, and $p' = \min(s', k) = s' = s$. Finally, because 418

419 Ω and Ω_2 are convex, one finds that $\sigma = 1$. 420

- Now we have all the tools to study the behavior $A_h(w_I^*)$ as h goes to 0. 421
- COROLLARY 5.7. Under the same assumptions as in proposition 5.5, one has 422

423
$$A_h(w_J^*) \le Ch^{2p'+\sigma}$$

- with C independent of h and $p' = \min(s', k)$. 424
- *Proof.* Applying the multiplicative trace inequality (recalled in proposition A.1) and using the estimates of 425proposition 5.5 yield the estimates 426

427
$$\|u^{w_{J}^{*}} - u_{h}^{w_{J}^{*}}\|_{L^{2}(\Sigma)}^{2} \leq Ch^{2p'+\sigma} \|u\|_{PH^{1+p'}(\Omega)} \text{ and } \|u_{2}^{w_{J}^{*}} - u_{2,h}^{w_{J}^{*}}\|_{L^{2}(\Sigma)}^{2} \leq Ch^{2p'+\sigma} \|u_{2}\|_{H^{1+p'}(\Omega_{2})}.$$

By design, one has $u^{w_{J}^{*}|\Sigma} = u_{2}^{w_{J}^{*}|\Sigma}$. So, observing that 428

$$\|u_h^{w_J^*} - u_{2,h}^{w_J^*}\|_{\mathrm{L}^2(\Sigma)}^2 \le 2(\|u^{w_J^*} - u_h^{w_J^*}\|_{\mathrm{L}^2(\Sigma)}^2 + \|u_2^{w_J^*} - u_{2,h}^{w_J^*}\|_{\mathrm{L}^2(\Sigma)}^2),$$

we conclude that $A_h(w_I^*) \leq Ch^{2p'+\sigma}$. 430

The previous result gives us a simple way to choose the parameter λ_h in order to ensure that both $(\lambda_h)_h$ and 431 $(A_h(w_J^*)/\lambda_h)_h$ tend to 0 as h tends to 0. 432

COROLLARY 5.8. Under the same assumptions as in proposition 5.5, any parameter λ_h of the form $\lambda_h = Ch^q$ 433 with C > 0 independent of h and $q \in (0, 2p' + \sigma)$ satisfies the conditions of theorem 5.4. 434

Remark 5.9. Within the framework of remark 5.6, one may choose $q \in (0, 2\sigma_D(\varepsilon) + 1)$ in the statement of 435corollary 5.8. 436

Thanks to theorem 5.4, using the conditions of corollary 5.8, one obtains the convergence of the discrete 437 solutions to the exact solution. 438

On the one hand, convergence is guaranteed even on meshes that are not T-conforming. Compared to [1], 439

- convergence holds in very general situations, namely as soon as there is a shift theorem for problem (1.1), cf. 440 theorem 2.2, even with a regularity exponent $\sigma_D(\varepsilon) < 1/2$. 441
- On the other hand, there is no associated convergence rate. Assuming a Céa lemma-like result, and using the 442 same notations as above, the *expected* convergence rate is $h^{p'}$ in H_0^1 -norm, and $h^{p'+\sigma}$ in L^2 -norm. Whereas, classically, the *optimal* convergence rate is h^k in H_0^1 -norm, and h^{k+1} in L^2 -norm. 443 444

Numerical experiments In this section we turn our attention to the validation of the numerical 445 6 method that we have proposed. We limit ourselves to the case of 2D domains and use P^1 Lagrange finite 446 elements. The numerical results that we present below have been obtained with the help of the library 447 **FreeFem**++². Since the well-posedness of (1.1) depends on the shape of the interface Σ , we test the performance 448 of our method in three different configurations. In the first one, Σ is flat, in the second one, Σ is a circular 449interface and in the last one, Σ has a "corner", in the sense that the angle at the intersection with the boundary 450is not a right angle. In all these experiments, we suppose that the coefficients ε_1 and ε_2 are constant with 451 $\varepsilon_1 = 1$. We denote by κ_{ε} the contrast $\kappa_{\varepsilon} = \varepsilon_2/\varepsilon_1$. 452

The shape, smoothness and (respective) volumes of Ω_1 and Ω_2 are taken into account to choose the domain 453 $\Omega^{\star} \in \{\Omega_1, \Omega_2\}$ to which the extension is performed (we recall that one must have meas $\partial_{\Omega}(\partial \Omega^{\star} \setminus \Sigma) > 0$, see 454 footnote¹ on page 4). Indeed, to have a better convergence rate, one should choose Ω^* convex, or with as 455

²See https://freefem.org/.

smooth a boundary as possible. Also, in order to speed up the convergence of the optimization algorithm, 456we must choose Ω^* as small as possible. Once Ω^* is fixed, one has to extend the function ε_1 or ε_2 to all 457the domain Ω . Because the coefficients are constant, we extend ε_1 (resp. ε_2) by ε_1 in Ω_2 (resp. in Ω_1). 458In the case where Σ is flat or circular, we take $\Omega^* = \Omega_2$. In the third configuration, we take $\Omega^* = \Omega_1$. To 459solve the optimization problem, we will use two different algorithms that are available in FreeFem++. The 460 first one is the algorithm BFGS (Broyden-Fletcher-Goldfarb-Shanno) and the second one is the algorithm NLCG 461 (Nonlinear Conjugate Gradient). Compared to the NLCG algorithm which uses only the gradient of the cost 462 function to solve the optimization problem, the BFGS algorithm, which belongs to the class of quasi-Newton 463 methods, uses a particular approximation of the hessian of the objective function. As already mentioned in 464the documentation of FreeFem++ (see page 65 of the third version), when the unknown of the optimization 465problem is a finite element function with a large size, it is preferable to work with the NLCG algorithm because 466the BFGS algorithm can be very memory consuming and its convergence rate can be low. In our numerical 467 experiments, we observed that in general both algorithms work similarly for the case where the interface is flat 468 or circular, however for the case of the interface with corner, the algorithm NLCG performs better (the observed 469convergence rate is better). Below we present the numerical results we obtained using the BFGS algorithm 470for the case where the interface is flat or circular and the results we obtained using the NLCG algorithm for 471the case where the interface has a corner. In our numerical experiments, the BFGS function was used with 472the following parameters: eps=1.e-11, nbiter=10, nbiterline=1 and the NLCG function with the following: 473nbiter=10,eps=1.e-11. 474

475 476

6.1 Flat interface In this paragraph, we take

$$\Omega_1 = \{(x,y) \in (0;1/2) \times (0;1)\}$$
 and $\Omega_2 = \{(x,y) \in (1/2;1) \times (0;1)\}$

(a flat interface and a domain which is symmetric with respect to Σ). We consider a family of meshes of $\overline{\Omega}$ satisfying Assumption 5.1 (see Figure 2). In the rest of this paragraph we suppose that $\kappa_{\varepsilon} \neq -1$. To test the performance of our method, we work with the same example considered in [1, 14]. Define the function $u_{\kappa_{\varepsilon}}$ such that

481
$$u_{\kappa_{\varepsilon}}(x,y) = \begin{cases} (x^2 + bx)\sin(\pi y) & \text{if } x < 1/2\\ a(x-1)\sin(\pi y) & \text{if } 1/2 < x \end{cases}, \text{ where } a = \frac{1}{2(\kappa_{\varepsilon} + 1)} \text{ and } b = -\frac{\kappa_{\varepsilon} + 2}{2(\kappa_{\varepsilon} + 1)}.$$

and consider it as an exact solution to (1.1). This is possible because $\operatorname{div}(\varepsilon \nabla u_{\kappa_{\varepsilon}}) \in \operatorname{L}^{2}(\Omega)$. The source term *f* is computed accordingly. As observed in remark 5.6, by choosing $\lambda_{h} = Ch^{q}$ with $q \in (0,3)$, the method is convergent. In our experiment, we take $\lambda_{h} = 0.002h^{2}$. We work with $\kappa_{\varepsilon} = -1.001$. The behaviors of the relative L²-norm error ($||e_{h}^{r}||_{0}$) and the relative H¹₀-norm error ($||e_{h}^{r}||_{1}$) between the exact solution and the numerical one are reported in Figure 2. We observe that both rates of convergence are equal to 2.

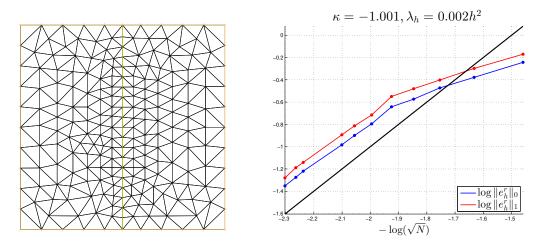


FIG. 2. A given mesh (left). Behavior of the relative L^2 and H_0^1 errors with respect to the meshsize $h \sim \sqrt{N}$, where N is the total number of nodes of the mesh (right).

Remark 6.1. The constant C in $\lambda_h = Ch^q$ must be adjusted by the user according to the contrast κ_{ε} in order to obtain a fast convergence of the method. Clearly this depends on $||w_J^*||_{\tilde{\varepsilon}_1}$. Using the fact that $w_J^* = w_H$ and owing to (4.7) we see that this depends on the jump of the normal derivative (across Σ) between u_1 and its harmonic extension. It is also important to note that, once q is fixed and when h is small enough, the choice of C has little influence on the convergence of the method.

6.2 The case of a circular interface In this paragraph, we consider the case of a circular inclusion, precisely the domains Ω_1 and Ω_2 are such that $\Omega_1 = \{ \boldsymbol{x} \in \mathbb{R}^2 \mid |\boldsymbol{x}| < 1 \}$ and $\Omega_2 = \{ \boldsymbol{x} \in \mathbb{R}^2 \mid 1 < |\boldsymbol{x}| < 2 \}$. In proposition A.2, we prove that A_{ε} is an isomorphism $\kappa_{\varepsilon} \notin \{-1\} \cup \mathscr{S}$ with $\mathscr{S} := \{-(1 - (1/2)^{2n})/(1 + (1/2)^{2n}) \mid n \in \mathbb{N}^*\}$. For this reason, we consider the case where $\kappa_{\varepsilon} = -2 \notin \mathscr{S}$. Given that both Ω_2 and Ω have smooth boundaries, we infer that $\sigma = 1$ and s' = s. By taking f as the source term associated to the function

497
$$u_{\kappa_{\varepsilon}}(x,y) = \begin{cases} r^2 + b & \text{if } r < 1\\ a(r-2)^2 & \text{if } 1 < r < 2. \end{cases}, \text{ with } r = \sqrt{x^2 + y^2}, a = -1/\kappa_{\varepsilon} \text{ and } b = a - 1 \end{cases}$$

and by taking $\lambda_h = 0.002h^2$, we obtain the results displayed in Figure 3. We observe that the method converges with optimal rate (i.e., the relative L²-norm error ($||e_h^r||_0$) is of order 2, while the relative H¹₀-norm error is of order 1), even though the exterior boundary and the interface are curved.

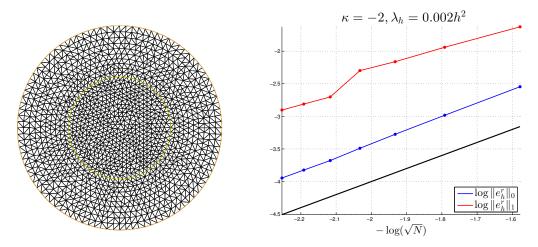


FIG. 3. A given mesh (left). Behavior of the relative L^2 and H_0^1 errors with respect to the meshsize $h \sim \sqrt{N}$, where N is the total number of nodes of the mesh (right)

6.3 The case of an interface with corner Now, we consider the configuration where the interface 502 Σ has a corner. More precisely, we assume that $\Omega := \{ \boldsymbol{x} \in \mathbb{R}^2 \mid |\boldsymbol{x}| < 1 \text{ and } \arg(\boldsymbol{x}) \in (0; \pi) \}$ and $\Omega_1 := \{ \boldsymbol{x} \in \Omega \mid \arg(\boldsymbol{x}) \in (0; \pi/4) \}$ (see Figure 4). In such configuration, it can be proved (see [4]) that A_{ε} is an isomorphism 504 if and only if $\kappa_{\varepsilon} \in \mathbb{R}^* \setminus [-3, -1]$. Furthermore, contrarily to the two previous cases, in this configuration the 505 solution to (1.1) can be very singular near the origin. Indeed, it was proved in [12, Chapter 2] that the 506 regularity of the solution to (1.1) depends in κ_{ε} and can be very low as κ_{ε} approaches [-3, -1]: more precisely,

507
$$\lim_{\kappa_{\varepsilon} \to -3^{-}} \sigma_{D}(\varepsilon) = \lim_{\kappa_{\varepsilon} \to -1^{+}} \sigma_{D}(\varepsilon) = 0$$

As a matter of fact, the value of the regularity exponent $\sigma_D(\varepsilon)$ is $\Re e(\lambda_0)$, where λ_0 is the solution to

509 (6.1)
$$\kappa_{\varepsilon} = -\tan(3\lambda\pi/4)/\tan(\lambda\pi/4)$$

that has the smallest positive real part. Note that one can show (see [12, Chapter 3]) that all the solutions to (6.1) are real-valued. In the particular cases where $\kappa_{\varepsilon} = -5$ and $\kappa_{\varepsilon} = -3.1$, one finds, respectively that $\lambda_0 \approx 0.458$ and $\lambda_0 \approx 0.139$. As mentioned previously this regularity result is optimal. Indeed, one can check that the function

514
$$u_{\lambda_0}(r,\theta) := (1-r)r^{\lambda_0} \begin{cases} \sin(\lambda_0\theta) / \sin(\lambda_0\pi/4) & \theta \in (0;\pi/4), \\ \sin(\lambda_0(\pi-\theta)) / \sin(3\lambda_0\pi/4) & \theta \in (\pi/4;\pi) \end{cases}$$

satisfies div $(\varepsilon \nabla u_{\lambda_0}) \in L^2(\Omega)$. Observe that $u_{\lambda_0} \notin PH^{\lambda_0}(\Omega)$. This means that $u_{\lambda_0} \notin PH^{3/2}(\Omega)$. Now, given that Ω and Ω_2 are both convex, owing to proposition 5.4, we can say that by choosing $\lambda_h = Ch^q$ with $q \in (0, 1 + 2\lambda_0)$, the convergence of the method can be guaranteed. In the case $\kappa_{\varepsilon} = -5$ (resp. $\kappa_{\varepsilon} = -3.1$), we work with $\lambda_h = 6h^{1.8}$ (resp. $\lambda_h = 1.5h^{1.2}$).

The behaviors of the relative L²-norm error and the relative H_0^1 -norm error (for the cases $\kappa_{\varepsilon} = -5$ and $\kappa_{\varepsilon} = -3.1$) are given in Figure 4. In either case, the expected rate of convergence is equal to λ_0 (≈ 0.458 when $\kappa_{\varepsilon} = -5$ and ≈ 0.139 when $\kappa_{\varepsilon} = -3.1$) for the case of the H_0^1 -norm error, while it is equal to $2\lambda_0$ for the case of the L²-norm error. We observe that, in both cases, the method converges with optimal rate of convergence

523 for the H^1 -norm and the L^2 one.

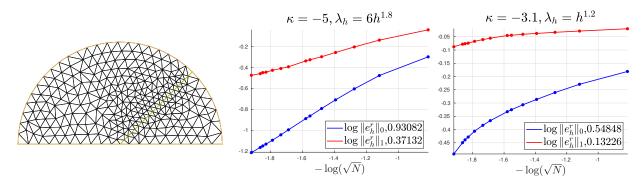


FIG. 4. A given mesh (left). Behavior of the relative L^2 and H_0^1 errors with respect to the meshsize $h \sim \sqrt{N}$, where N is the total number of nodes of the mesh, with the observed convergence rates, when $\kappa_{\varepsilon} = -5$ (center) and $\kappa_{\varepsilon} = -3.1$ (right).

7 Concluding remarks In this work, we have presented a new numerical method to approximate the solution to the scalar transmission problem with sign-changing coefficients. We proved that the method converges without any restriction on the mesh sequence, nor on the regularity of the solution. This result has been illustrated by several 2D numerical experiments. The convergence rate of our method seems to be optimal. In order to improve the performance of the method, several questions can be studied:

- 1. Choose the parameter λ_h in order to accelerate convergence. An interesting idea would be to find an adaptive approach to fit its value. Also, one could use adaptive mesh refinement, together with a posteriori estimates. We refer to [17] for estimators that deliver guaranteed error bounds, and that are robust with respect to the sign-changing coefficient ε . Work with other regularization approaches i.e. other choices for the coefficient $\tilde{\varepsilon}_1$ and/or an alter-
 - 2. Work with other regularization approaches, i.e., other choices for the coefficient $\tilde{\varepsilon}_1$, and/or an alternative to the Tikhonov regularization method.
- In the case where the interface is not regular, it would be interesting to combine our approach with
 other existing methods for solving PDE with singular solution such as the Singular Complement
 Method [16].

Besides that, it will be also interesting to extend this approach to other models involving sign-changing coefficients.

Acknowledgments The authors would like to thank Anne-Sophie Bonnet-Ben Dhia and Lucas Chesnel for their helpful comments on the first draft of this article. We also thank Farah Chaaban, for her constructive remarks during her internship, regarding the implementation of the method.

543 Appendix A. Missing results.

545

544 Proof of Proposition 4.4. From the definition of w_{δ}^* , we can write that

$$\delta \|w_{\delta}^*\|_{\tilde{\varepsilon}_1}^2 \leq J^{\delta}(w_{\delta}^*) \leq J^{\delta}(w_J^*) = J(w_J^*) + \delta \|w_J^*\|_{\tilde{\varepsilon}_1}^2 = \delta \|w_J^*\|_{\tilde{\varepsilon}_1}^2.$$

This means that for all $0 < \delta$, there holds $||w_{\delta}^*||_{\tilde{\varepsilon}_1} \leq ||w_J^*||_{\tilde{\varepsilon}_1}$. As a result (w_{δ}^*) is bounded in $V_2(\Omega_2)$. This implies that, up to a sub-sequence, $(w_{\delta}^*)_{\delta}$ converges, as δ tends to 0, weakly in $V_2(\Omega_2)$ to some $w_0 \in V_2(\Omega_2)$. For the reader's convenience, this sequence is also denoted by $(w_{\delta}^*)_{\delta}$. Now, let us prove that w_0 is a minimizer of J. To do that, we start by observing that for all $\delta > 0$, we have

550
$$0 \le J(w_{\delta}^*) \le J^{\delta}(w_{\delta}^*) \le J^{\delta}(w_J^*) = \delta \|w_J^*\|_{\tilde{e}_1}^2.$$

- This shows that $(J(w_{\delta}^*))_{\delta}$ converges to zero as δ tends to zero. On the other hand, by using the result of proposition 4.1, we know that $(J(w_{\delta}^*))_{\delta}$ converges to $J(w_0)$. Consequently, $J(w_0) = 0$ and then w_0 is a minimizer of J.
- The next step is to show that the convergence of $(w_{\delta}^*)_{\delta}$ to w_0 occurs in the strong sense and that $w_0 = w_J^*$. To do so, we observe that

556
$$\|w_{\delta}^*\|_{\tilde{\varepsilon}_1} \le \|w_J^*\|_{\tilde{\varepsilon}_1} \ \forall \delta \Longrightarrow \limsup_{\delta \to 0} \|w_{\delta}^*\|_{\tilde{\varepsilon}_1} \le \|w_J^*\|_{\tilde{\varepsilon}_1}, \text{ and } w_{\delta}^* \rightharpoonup w_0 \text{ in } \mathcal{V}_2(\Omega_2) \Longrightarrow \|w_0\|_{\tilde{\varepsilon}_1} \le \liminf_{\delta \to 0} \|w_{\delta}^*\|_{\tilde{\varepsilon}_1},$$

where the latter is a consequence of the fact that the norm of a Banach space is weakly lower semicontinuous, see [9, Proposition III.5 (iii)]. This implies that $||w_0||_{\tilde{\varepsilon}_1} \leq ||w_J^*||_{\tilde{\varepsilon}_1}$. Thanks to the definition of w_J^* , we deduce that $w_0 = w_J^*$. With this in mind and with the help of the previous inequality, we conclude that

560
$$\lim_{\delta \to 0} \|w_{\delta}^*\|_{\tilde{\varepsilon}_1} = \|w_J^*\|_{\tilde{\varepsilon}_1}.$$

Since $V_2(\Omega_2)$ is a Hilbert space, it follows (see [9, Proposition III.32]) that $w_{\delta} \to w_J^*$ in $V_2(\Omega_2)$. By noticing that w_J^* is independent of the considered sub-sequence, the result is then proved.

From Proposition A.1. [8, Theorem 1.6.6] Let Ω be an open, bounded, connected subset of \mathbb{R}^d (d = 2, 3) with a Lipschitz boundary. Then the estimate

565
$$||u||_{L^{2}(\partial\Omega)} \leq C||u||_{L^{2}(\Omega)}^{1/2} ||u||_{H^{1}(\Omega)}^{1/2} \quad \forall u \in H^{1}(\Omega)$$

566 holds with 0 < C independent of u.

567 PROPOSITION A.2. Let $\Omega_1 = \{ \boldsymbol{x} \in \mathbb{R}^2 \mid |\boldsymbol{x}| < 1 \}$ and $\Omega_2 = \{ \boldsymbol{x} \in \mathbb{R}^2 \mid 1 < |\boldsymbol{x}| < 2 \}$. Assume that $\kappa_{\varepsilon} := \varepsilon_2 / \varepsilon_1 \notin \{ -1 \} \cup \mathscr{S}$ with

$$\mathscr{S} := \left\{ -\frac{1 - (1/2)^{2n}}{1 + (1/2)^{2n}} \mid n \in \mathbb{N}^* \right\}.$$

570 Then the operator $A_{\varepsilon} : \mathrm{H}^{1}_{0}(\Omega) \to \mathrm{H}^{1}_{0}(\Omega)$ is an isomorphism.

Remark A.3. Note that in accordance with the results concerning the Neumann-Poincaré operator [24, Chapter 1], we observe that -1 is an accumulation point of \mathscr{S} .

Proof. [12, Theorem 1.3.3] guarantees that A_{ε} is Fredholm of index 0 when $\kappa_{\varepsilon} \neq -1$. Therefore it suffices to study its kernel. Let $u \in \mathrm{H}_{0}^{1}(\Omega)$ be such that $A_{\varepsilon}u = 0$. Then $u_{1} := u_{|\Omega_{1}|}$ and $u_{2} = u_{|\Omega_{2}|}$ satisfy

575

569

$$\begin{cases} \Delta u_1 = 0 \quad \text{in } \Omega_1 \\ \Delta u_2 = 0 \quad \text{in } \Omega_2 \\ u_1(1,\theta) = u_2(1,\theta) \quad \text{and} \quad \partial_r u_1(1,\theta) = \kappa_{\varepsilon} \partial u_2(1,\theta) \quad \forall \theta \in [0;2\pi]. \end{cases}$$

Since the problem is invariant with respect to θ , by Fourier decomposition for u_1 , u_2 we have the representations:

578
$$u_1(r,\theta) = \sum_{n \in \mathbb{N}} a_n r^n e^{in\theta} \quad \text{and} \quad u_2(r,\theta) = b_0 \ln(r/2) + \sum_{n \in \mathbb{Z}^*} b_n ((r/2)^n - (r/2)^{-n}) e^{in\theta},$$

579 where $a_n, b_n \in \mathbb{C}$. Using the transmission conditions, we get

$$\begin{vmatrix} a_0 = b_0 \ln(1/2), & 0 = b_0 \kappa_{\varepsilon} \\ a_n = b_n ((1/2)^n - (1/2)^{-n}), & a_n = b_n ((1/2)^n + (1/2)^{-n}) \kappa_{\varepsilon}, & n \in \mathbb{N}^* \\ 0 = b_n ((1/2)^n - (1/2)^{-n}), & 0 = b_n ((1/2)^n + (1/2)^{-n}) \kappa_{\varepsilon}, & -n \in \mathbb{N}^*. \end{vmatrix}$$

581 Therefore we deduce that A_{ε} is injective when $\kappa_{\varepsilon} \notin \mathscr{S}$.

Appendix B. On the use of the adjoint approach to compute the gradient of the cost functional J. The adjoint approach was introduced in [11] as a method for computing the gradient of cost functions that depend in non-explicit way on the main variable of the problem, namely via the solution of PDEs (the state equations) in which the main variable plays the role of a parameter. Here, we are going to explain how

to apply this method to our case. The idea is to introduce a Lagrangian functional $\mathscr{L}: V_2(\Omega_2) \times H^1_0(\Omega) \times H^1_0(\Omega)$ 586 $V_2(\Omega_2) \times H_0^1(\Omega) \times V_2(\Omega_2) \to \mathbb{R}$ such that 587

$$\mathscr{L}(w, u, u_2, g, g_2) = \frac{1}{2} \int_{\Sigma} |u - u_2|^2 \, \mathrm{d}\sigma + a_1(w, u, g) + a_2(w, u, u_2, g_2)$$

in which $a_1(w, u_1, g)$ and $a_2(w, u_2, g_2)$ are respectively given by 589ſ

I.

590

588

$$\begin{vmatrix} a_1(w, u, g) = \int_{\Omega} \tilde{\varepsilon}_1 \nabla u \cdot \nabla g \, \mathrm{d}\boldsymbol{x} - \int_{\Omega_1} fg \, \mathrm{d}\boldsymbol{x} - \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla w \cdot \nabla g \, \mathrm{d}\boldsymbol{x} \\ a_2(w, u, u_2, g_2) = \int_{\Omega_2} \varepsilon_2 \nabla u_2 \cdot \nabla g_2 \, \mathrm{d}\boldsymbol{x} - \int_{\Omega_2} f_2 g_2 \, \mathrm{d}\boldsymbol{x} + \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla (w - u) \cdot \nabla g_2 \, \mathrm{d}\boldsymbol{x} \end{vmatrix}$$

ſ

ſ

The functions $g \in H_0^1(\Omega), g_2 \in V_2(\Omega_2)$ are the adjoint variables associated to u, u_2 respectively. Let (u^w, u_2^w) 591be the solution to (3.2). By design, $a_1(w, u^w, g) = 0$ for all $g \in H_0^1(\Omega)$, and $a_2(w, u^w, u_2^w, g_2) = 0$ for all 592 $g_2 \in V_2(\Omega_2)$, so one has 593

594 (B.1)
$$\mathscr{L}(w, u^w, u_2^w, g, g_2) = J(w) \qquad \forall g \in \mathrm{H}^1_0(\Omega), \ \forall g_2 \in \mathrm{V}_2(\Omega_2).$$

Clearly, the functional \mathscr{L} is differentiable with respect to all its variables. For all $\mathbf{v} = (w, u, u_2, g, g_2) \in$ 595 $V_2(\Omega_2) \times H_0^1(\Omega) \times V_2(\Omega_2) \times H_0^1(\Omega) \times V_2(\Omega_2)$, the partial derivatives of \mathscr{L} at v belong respectively to 596

 $\partial_{w} \mathscr{L}(\mathbf{v}) \in (\mathcal{V}_{2}(\Omega_{2}))^{*}, \ \partial_{u} \mathscr{L}(\mathbf{v}) \in (\mathcal{H}_{0}^{1}(\Omega))^{*}, \ \partial_{u_{2}} \mathscr{L}(\mathbf{v}) \in (\mathcal{V}_{2}(\Omega_{2}))^{*}, \\ \partial_{g} \mathscr{L}(\mathbf{v}) \in (\mathcal{H}_{0}^{1}(\Omega))^{*}, \ \partial_{g_{2}} \mathscr{L}(\mathbf{v}) \in (\mathcal{V}_{2}(\Omega_{2}))^{*}. \\ \text{Let } g \in \mathcal{H}_{0}^{1}(\Omega) \text{ and } g_{2} \in \mathcal{V}_{2}(\Omega_{2}) \text{ be given, and } \mathbf{v}^{w} = (w, u^{w}, u^{w}_{2}, g, g_{2}). \\ \text{By taking the derivative of the relation}$ 598(B.1) with respect to w, we find that, by applying the chain rule formula, 599

600
$$(J'(w),h)_{\tilde{\varepsilon}_1} = \langle \partial_w \mathscr{L}(\mathbf{v}^w),h\rangle + \langle \partial_u \mathscr{L}(\mathbf{v}^w),\frac{du^w}{dw}(h)\rangle + \langle \partial_{u_2} \mathscr{L}(\mathbf{v}^w),\frac{du_2^w}{dw}(h)\rangle, \quad \forall h \in \mathcal{V}_2(\Omega_2).$$

Now, if there exists $(g^w, g_2^w) \in \mathrm{H}^1_0(\Omega) \times \mathrm{V}_2(\Omega_2)$ for which the equations 601

602
$$\partial_u \mathscr{L}(w, u^w, u^w_2, g^w, g^w_2) = 0$$
 and $\partial_{u_2} \mathscr{L}(w, u^w, u^w_2, g^w, g^w_2) = 0$

are satisfied for all $w \in V_2(\Omega_2)$, this yields 603

$$(J'(w),h)_{\tilde{\varepsilon}_1} = \langle \partial_w \mathscr{L}(w,u^w,u_2^w,g_1^w,g_2^w),h \rangle \qquad \forall w \in \mathcal{V}_2(\Omega_2), \ \forall h \in \mathcal{V}_2(\Omega_2).$$

To investigate the existence of (g^w, g_2^w) , we need to write down the expression of $\partial_u \mathscr{L}(\mathbf{v}^w)$ and $\partial_{u_2} \mathscr{L}(\mathbf{v}^w)$: By 605 a direct calculus, one checks that 606

$$\langle \partial_u \mathscr{L}(\mathbf{v}^w), v \rangle = \int_{\Omega} \tilde{\varepsilon}_1 \nabla g \cdot \nabla v \, \mathrm{d}\boldsymbol{x} - \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla g_2 \cdot \nabla v \, \mathrm{d}\boldsymbol{x} + \int_{\Sigma} (u^w - u_2^w) v \, \mathrm{d}\sigma \qquad \forall v \in \mathrm{H}^1_0(\Omega)$$

$$\langle \partial_{u_2} \mathscr{L}(\mathbf{v}^w), v_2 \rangle = \int_{\Omega_2} \varepsilon_2 \nabla g_2 \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} - \int_{\Sigma} (u^w - u_2^w) v_2 \, \mathrm{d}\sigma \qquad \forall v_2 \in \mathrm{V}_2(\Omega_2).$$

607

608

614

604

$$J_{\Omega_2}$$
 J_{Σ}
Hence, the functions $(g^w, g_2^w) \in \mathrm{H}^1_0(\Omega) \times \mathrm{V}_2(\Omega_2)$ are governed by the following system of equations:

609 (B.2)
$$\begin{cases} \int_{\Omega} \tilde{\varepsilon}_1 \nabla g^w \cdot \nabla v \, \mathrm{d}\boldsymbol{x} = \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla g_2^w \cdot \nabla v \, \mathrm{d}\boldsymbol{x} - \int_{\Sigma} (u^w - u_2^w) v \, \mathrm{d}\sigma & \forall v \in \mathrm{H}^1_0(\Omega) \\ \int_{\Omega_2} \varepsilon_2 \nabla g_2^w \cdot \nabla v_2 \, \mathrm{d}\boldsymbol{x} = \int_{\Sigma} (u^w - u_2^w) v_2 \, \mathrm{d}\sigma & \forall v_2 \in \mathrm{V}_2(\Omega_2). \end{cases}$$

Clearly the previous system of equations is well-posed. Therefore the functions g^w , g_2^w are well-defined. We 610 then have all the tools to prove the result stated in Lemma 4.5. 611

Proof of Lemma 4.5. Take $w \in V_2(\Omega_2)$. From the characterization (B.2) of g^w and g_2^w , we deduce that for all 612 $h \in V_2(\Omega_2)$, we have 613

 $(J'(w),h)_{\tilde{\varepsilon}_1} = \langle \partial_w \mathscr{L}(w,u^w,u^w_2,g^w_1,g^w_2),h \rangle.$

On the other hand, one can compute explicitly the value of $\langle \partial_w \mathscr{L}(w, u, u_2, g, g_2), h \rangle$: 615

616
$$\langle \partial_w \mathscr{L}(w, u, u_2, g, g_2), h \rangle = \int_{\Omega_2} \tilde{\varepsilon}_1 \nabla h \cdot \nabla (g_2 - g_{|\Omega_2}) \, \mathrm{d}\boldsymbol{x}.$$

This shows that $J'(w) = g_2^w - g_{|\Omega_2}^w$ and then the result is proved. 617

17

REFERENCES

- [1] A. Abdulle, M.E. Huber, and S. Lemaire. An optimization-based numerical method for diffusion problems with sign-changing
 coefficients. C. R. Acad. Sci. Paris, Série I, 355:472–478, 2017.
- [2] C. Atamian and P. Joly. Une analyse de la méthode des domaines fictifs pour le problème de Helmholtz extérieur. ESAIM:
 Math. Modell. Numer. Anal., 27:251–288, 1993.
- [3] A.-S. Bonnet-Ben Dhia, C. Carvalho, and P. Ciarlet, Jr. Mesh requirements for the finite element approximation of problems
 with sign-changing coefficients. *Numer. Math.*, 138:801–838, 2018.
- [4] A.-S. Bonnet-Ben Dhia, L. Chesnel, and P. Ciarlet Jr. T-coercivity for scalar interface problems between dielectrics and metamaterials. *ESAIM: Math. Modell. Numer. Anal.*, 46:1363–1387, 2012.
- [5] A.-S. Bonnet-Ben Dhia, L. Chesnel, and X. Claeys. Radiation condition for a non-smooth interface between a dielectric and
 a metamaterial. Math. Model. Meth. App. Sci., 23:1629–1662, 2013.
- [6] A.-S. Bonnet-Ben Dhia, P. Ciarlet Jr., and C.-M. Zwölf. Time harmonic wave diffraction problems in materials with
 sign-shifting coefficients. J. Comput. Appl. Math., 234:1912–1919, Corrigendum 2616, 2010.
- [7] A.-S. Bonnet-Ben Dhia, M. Dauge, and K. Ramdani. Spectral analysis and singularities of a non-coercive transmission
 problem (in French). C. R. Acad. Sci. Paris, Ser. I, 328:717-720, 1999.
- [8] S. Brenner and L.R. Scott. The mathematical theory of finite element methods, 3rd Edition. Springer Verlag, 2008.
- [9] H. Brezis. Functional analysis, Sobolev spaces and partial differential equations. Springer, 2011.
- [10] C. Carvalho. Mathematical and numerical study of plasmonic structures with corners. PhD thesis, Ecole Polytechnique,
 2015.
- [11] G. Chavent. Nonlinear least squares for inverse problems: theoretical foundations and step-by-step guide for applications.
 Springer, 2010.
- [12] L. Chesnel. Investigation of some transmission problems with sign-changing coefficients. Application to metamaterials (in French). PhD thesis, Ecole Polytechnique, 2012.
- [13] L. Chesnel and P. Ciarlet, Jr. Compact imbeddings in electromagnetism with interfaces between classical materials and meta-materials. SIAM J. Math. Anal., 43:2150–2169, 2011.
- [14] L. Chesnel and P. Ciarlet, Jr. T-coercivity and continuous Galerkin methods: application to transmission problems with
 sign changing coefficients. Numer. Math., 124:1–29, 2013.
- [15] P.G. Ciarlet. The finite element method for elliptic problems. Classics in Applied Mathematics. SIAM, 2002.
- [16] P. Ciarlet, Jr. and J. He. The singular complement method for 2D scalar problems. C. R. Acad. Sci. Paris, Série I,
 336:353–358, 2003.
- [17] P. Ciarlet Jr. and M. Vohralík. Localization of global norms and robust a posteriori error control for transmission problems
 with sign-changing coefficients. ESAIM: Math. Modell. Numer. Anal., 52:2037–2064, 2018.
- [18] M. Dauge and B. Texier. non-coercive transmission problems in polygons (in French). Technical Report 97-27, IRMAR,
 Rennes 1 University, Rennes, France, 1997.
- [19] A. Decoene, S. Martin, and F. Vergnet. A smooth extension method for transmission problems. Technical Report hal 02146271v2, HAL, 2019.
- [20] H.W. Engl. Regularization methods for the stable solution of inverse problems. Surv. Math. Ind., 3:71–143, 1993.
- [21] B. Fabrèges, L. Gouarin, and B. Maury. A smooth extension method. C. R. Acad. Sci. Paris, Série I, 351:361–366, 2013.
- [22] M.D. Gunzburger, J.S. Peterson, and H. Kwon. An optimization based Domain Decomposition method for Partial Differential
 Equations. Computers Math. Applic., 37:77–93, 1999.
- [23] S. Nicaise and J. Venel. A posteriori error estimates for a finite element approximation of transmission problems with sign changing coefficients. J. Comput. Appl. Math., 235:4272–4282, 2011.
- 660 [24] M. Rihani. Maxwell's equations in presence of negative index materials. PhD thesis, Institut Polytechnique de Paris, 2022.
- [25] A.N. Tikhonov, A.V. Goncharsky, V.V. Stepanov, and A.G. Yagola. Numerical methods for the solution of ill-posed problems.
 Springer Verlag, 1995.
- [26] C.-M. Zwölf. Méthodes variationnelles pour la modélisation des problèmes de transmission d'onde électromagnétique entre diélectrique et méta-matériau. PhD thesis, Université de Versailles Saint Quentin en Yvelines, 2007.