



**HAL**  
open science

# On Some Model Problem for the Propagation of Interacting Species in a Special Environment

Michel Chipot, Mingmin Zhang

► **To cite this version:**

Michel Chipot, Mingmin Zhang. On Some Model Problem for the Propagation of Interacting Species in a Special Environment. DCDS-A, 2021. hal-03662715

**HAL Id: hal-03662715**

**<https://hal.science/hal-03662715>**

Submitted on 11 May 2022

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# On Some Model Problem for the Propagation of Interacting Species in a Special Environment

Michel Chipot<sup>(a)</sup> & Mingmin Zhang<sup>(b)</sup>

<sup>(a)</sup>Institut für Mathematik  
Universität Zürich  
Winterthurerstrasse 190  
8057 Zürich, Switzerland  
m.m.chipot@math.uzh.ch

<sup>(b)</sup>School of Mathematical Sciences  
University of Science and Technology of China  
Hefei, Anhui 230026, China  
and  
Aix Marseille Université  
CNRS, Centrale Marseille, I2M  
UMR 7373, 13453 Marseille, France  
lang925@mail.ustc.edu.cn

## Abstract

The purpose of this note is to study the existence of a nontrivial solution for an elliptic system which comes from a newly introduced mathematical problem so called Field-Road model. Specifically, it consists of coupled equations set in domains of different dimensions together with some interaction of non classical type. We consider a truncated problem by imposing Dirichlet boundary conditions and an unbounded setting as well.

**AMS 2020 Subject Classification:** 35J47, 35J57

**Key words:** Nonlinear elliptic systems; Coupled equations; Field-Road model; Truncated domain.

## 1 Introduction and notation

It has been observed in diffusion or propagation problems that the topography of the environment is playing a crucial rôle. We refer for instance to [2], [3] and the references there. In particular the roads are not only essential for human beings but also can be useful for other species. One can easily imagine that mosquitos could be interested in socialising with human beings along a road but many other situations seem to have occurred due to this element of civilisation. In this note we consider a model problem where the living space of our species consists in a field and one or several roads that we will assume to be unidimensional. We will also assume the roads to be straight but several extensions can be addressed very easily with our approach. On each domain -field or road- we will consider nonlinear diffusion equations which include for instance the Fisher-KPP types. We restrict ourselves to two species. The interaction with the two species is modelled by a particular flux condition (see [2], [3] and the set of equation below).

Let  $\Omega_\ell$  be the open set of  $\mathbb{R}^2$ , defined for  $\ell, L > 0$  as

$$\Omega_\ell = (-\ell, \ell) \times (0, L).$$

We denote by  $\Gamma_0$  the part of the boundary of  $\Omega_\ell$  located on the  $x_1$ -axis i.e.

$$\Gamma_0 = (-\ell, \ell) \times \{0\}$$

and by  $\Gamma_1$  the rest of the boundary that is to say

$$\Gamma_1 = \partial\Omega_\ell \setminus \Gamma_0.$$

When convenient we will identify  $\Gamma_0$  to  $(-\ell, \ell)$ . In this setting  $\Omega_\ell$  stands for a field and  $\Gamma_0$  for a portion of a road.

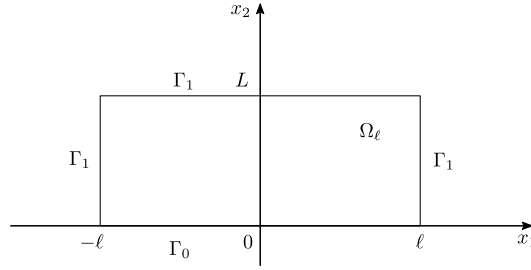


Figure 1: The domain  $\Omega_\ell$  for one-road problem

Set

$$V = \{v \in H^1(\Omega_\ell) \mid v = 0 \text{ on } \Gamma_1\}.$$

We would like to find a solution to the problem

$$\begin{cases} -D\Delta v = f(v) \text{ in } \Omega_\ell, \\ v = 0 \text{ on } \Gamma_1, \quad D \frac{\partial v}{\partial n} = \mu u - \nu v \text{ on } \Gamma_0, \\ -D'u'' + \mu u = g(u) + \nu v \text{ in } \Gamma_0, \\ u = 0 \text{ on } \partial\Gamma_0 = \{-\ell, \ell\}. \end{cases}$$

( $n$  denotes the outward unit normal to  $\Omega_\ell$ ).

In the weak form we would like to find a couple  $(u, v)$  such that

$$\begin{cases} (u, v) \in H_0^1(\Gamma_0) \times V, \\ \int_{\Omega_\ell} D\nabla v \cdot \nabla \varphi \, dx + \int_{\Gamma_0} \nu v(x_1, 0) \varphi \, dx_1 = \int_{\Omega_\ell} f(v) \varphi \, dx + \int_{\Gamma_0} \mu u \varphi \, dx_1 \quad \forall \varphi \in V, \\ \int_{\Gamma_0} D'u' \psi' + \mu u \psi \, dx_1 = \int_{\Gamma_0} \nu v(x_1, 0) \psi \, dx_1 + \int_{\Gamma_0} g(u) \psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0). \end{cases} \quad (1.1)$$

Here we assume that

$$D, D', m, \mu, \nu \text{ are positive constants,}$$

such that

$$m \geq \frac{\nu}{\mu}. \quad (1.2)$$

$f, g$  are Lipschitz continuous functions i.e. such that for some positive constants  $L_f, L_g$  it holds

$$|f(x) - f(y)| \leq L_f|x - y|, \quad |g(x) - g(y)| \leq L_g|x - y| \quad \forall x, y \in \mathbb{R}. \quad (1.3)$$

Note that this implies that for  $\lambda \geq L_f$  (respectively  $\eta \geq L_g$ ) the functions

$$x \rightarrow \lambda x - f(x), \quad \eta x - g(x) \quad (1.4)$$

are nondecreasing. In addition we will assume

$$f(0) = f(1) = 0, \quad f > 0 \text{ on } (0, 1), \quad f \leq 0 \text{ on } (1, +\infty). \quad (1.5)$$

$$g(0) = 0, \quad g(m) \leq 0. \quad (1.6)$$

One should remark that under the assumptions above,  $(0, 0)$  is a solution to (1.1). We are interested in finding a nontrivial solution to the problem (1.1). For the notation and the usual properties on Sobolev spaces we refer to [4], [5], [7], [8], [9]. We will also consider some extension in the case of a two-road problem which consists of three coupled equations with two interaction conditions on the upper- and lower- boundaries. Finally, we will address the case of an unbounded setting for the one-road problem.

## 2 Preliminary results

**Lemma 2.1.** *Suppose that  $w$  is a measurable function on  $\Gamma_0$  such that*

$$0 \leq w \leq m.$$

*Then under the assumptions above the problem*

$$\begin{cases} v \in V, \\ \int_{\Omega_\ell} D\nabla v \cdot \nabla \varphi \, dx + \int_{\Gamma_0} \nu v(x_1, 0) \varphi(x_1, 0) \, dx_1 \\ = \int_{\Omega_\ell} f(v) \varphi \, dx + \int_{\Gamma_0} \mu w \varphi \, dx_1 \quad \forall \varphi \in V, \end{cases} \quad (2.1)$$

*possesses a minimal and a maximal solution with values in  $(0, \frac{\mu}{\nu}m)$ .*

*Proof.* Let us remark first that any nonnegative solution to (2.1) takes its values in  $(0, \frac{\mu}{\nu}m)$ . Indeed if  $v$  is solution to (2.1) taking as test function  $\varphi = (v - k)^+$ ,  $k = \frac{\mu}{\nu}m \geq 1$  one gets

$$\begin{aligned} \int_{\Omega_\ell} D|\nabla(v - k)^+|^2 \, dx &= \int_{\Omega_\ell} D\nabla(v - k) \cdot \nabla(v - k)^+ \, dx = \int_{\Omega_\ell} D\nabla v \cdot \nabla(v - k)^+ \, dx \\ &= \int_{\Omega_\ell} f(v)(v - k)^+ \, dx + \int_{\Gamma_0} \{\mu w - \nu v(x_1, 0)\}(v - k)^+ \, dx_1 \\ &\leq \int_{\Gamma_0} \{\mu w - \nu v\}(v - k)^+ \, dx_1 \leq 0, \end{aligned}$$

since on the set where  $v \geq k$  one has  $v \geq \frac{\mu}{\nu}m$  and  $\{\mu w - \nu v\} \leq \{\mu w - \mu m\} \leq 0$ .

Next let us note that 0 is a subsolution to (2.1). Indeed this follows trivially from

$$\int_{\Omega_\ell} D\nabla 0 \cdot \nabla \varphi \, dx + \int_{\Gamma_0} \nu 0 \varphi \, dx_1 \leq \int_{\Omega_\ell} f(0) \varphi \, dx + \int_{\Gamma_0} \mu w \varphi \, dx_1 \quad \forall \varphi \in V, \varphi \geq 0.$$

On the other hand, for  $\varphi \in V$ ,  $\varphi \geq 0$ , one has also for  $k = \frac{\mu}{\nu}m$ , since  $k \geq 1$

$$\int_{\Omega_\ell} D\nabla k \cdot \nabla \varphi \, dx + \int_{\Gamma_0} \nu k \varphi \, dx_1 \geq \int_{\Omega_\ell} f(k) \varphi \, dx + \int_{\Gamma_0} \mu w \varphi \, dx_1$$

and the function constant equal to  $k = \frac{\mu}{\nu}m$  is a supersolution to (2.1).

For  $z \in L^2(\Omega_\ell)$  we denote by  $y = S(z)$  the solution to

$$\begin{cases} y \in V, \\ \int_{\Omega_\ell} D\nabla y \cdot \nabla \varphi \, dx + \int_{\Omega_\ell} \lambda y \varphi \, dx + \int_{\Gamma_0} \nu y(x_1, 0) \varphi \, dx_1 \\ \qquad \qquad \qquad = \int_{\Omega_\ell} f(z) \varphi \, dx + \int_{\Omega_\ell} \lambda z \varphi \, dx + \int_{\Gamma_0} \mu w \varphi \, dx_1 \quad \forall \varphi \in V, \end{cases} \quad (2.2)$$

where we have chosen  $\lambda \geq L_f$ . Note that the existence of a unique solution  $y$  to the problem above is an immediate consequence of the Lax-Milgram theorem.

First we claim that the mapping  $S$  is continuous from  $L^2(\Omega_\ell)$  into itself. Indeed setting  $y' = S(z')$  one has by subtraction of the equations satisfied by  $y$  and  $y'$

$$\begin{aligned} \int_{\Omega_\ell} D\nabla(y - y') \cdot \nabla \varphi \, dx + \int_{\Omega_\ell} \lambda(y - y') \varphi \, dx + \int_{\Gamma_0} \nu(y - y') \varphi \, dx_1 \\ = \int_{\Omega_\ell} \{f(z) - f(z')\} \varphi \, dx + \int_{\Omega_\ell} \lambda(z - z') \varphi \, dx \quad \forall \varphi \in V. \end{aligned} \quad (2.3)$$

Taking  $\varphi = y - y'$  one derives easily

$$\begin{aligned} \lambda \int_{\Omega_\ell} |y - y'|^2 \, dx &\leq \int_{\Omega_\ell} |f(z) - f(z')| |y - y'| \, dx + \lambda \int_{\Omega_\ell} |z - z'| |y - y'| \, dx \\ &\leq (L_f + \lambda) \int_{\Omega_\ell} |z - z'| |y - y'| \, dx. \end{aligned}$$

By the Cauchy-Schwarz inequality we obtain then

$$|S(z) - S(z')|_{2, \Omega_\ell} \leq \frac{L_f + \lambda}{\lambda} |z - z'|_{2, \Omega_\ell}$$

which shows the continuity of the mapping  $S$  ( $|\cdot|_{2, \Omega_\ell}$  denotes the usual  $L^2(\Omega_\ell)$ -norm).

We show now that the mapping  $S$  is monotone. Indeed suppose that  $z \geq z'$  and as above denote by  $y'$  the function  $S(z')$ . Taking  $\varphi = -(y - y')^-$  in (2.3) we get

$$\begin{aligned} \int_{\Omega_\ell} D|\nabla(y - y')^-|^2 \, dx + \int_{\Omega_\ell} \lambda((y - y')^-)^2 \, dx + \int_{\Gamma_0} \nu((y - y')^-)^2 \, dx_1 \\ = - \int_{\Omega_\ell} [\lambda(z - z') + \{f(z) - f(z')\}] (y - y')^- \, dx \leq 0, \end{aligned}$$

since  $-\{f(z) - f(z')\} \leq L_f|z - z'| \leq \lambda(z - z')$  (see (1.3), (1.4)). This shows that  $(y - y')^- = 0$  and the monotonicity of the mapping  $S$ .

We consider now the following sequences (Cf. [1]) :

$$\begin{aligned} \underline{y}_0 &= 0, & \overline{y}_0 &= \frac{\mu}{\nu}m = k, \\ \underline{y}_n &= S(\underline{y}_{n-1}), & \overline{y}_n &= S(\overline{y}_{n-1}), \quad n \geq 1. \end{aligned} \quad (2.4)$$

One has

$$\underline{y}_0 = 0 \leq \underline{y}_1 \leq \dots \leq \underline{y}_n \leq \overline{y}_n \leq \dots \leq \overline{y}_1 \leq \overline{y}_0 = \frac{\mu}{\nu}m = k. \quad (2.5)$$

Indeed since  $\underline{y}_1 = S(\underline{y}_0) = S(0)$  one has for  $\forall \varphi \in V, \varphi \geq 0$ ,

$$\begin{aligned} & \int_{\Omega_\ell} D\nabla \underline{y}_1 \cdot \nabla \varphi \, dx + \int_{\Omega_\ell} \lambda \underline{y}_1 \varphi \, dx + \int_{\Gamma_0} \nu \underline{y}_1(x_1, 0) \varphi \, dx_1 \\ &= \int_{\Omega_\ell} f(\underline{y}_0) \varphi \, dx + \int_{\Omega_\ell} \lambda \underline{y}_0 \varphi \, dx + \int_{\Gamma_0} \mu w \varphi \, dx_1 \\ &\geq \int_{\Omega_\ell} D\nabla \underline{y}_0 \cdot \nabla \varphi \, dx + \int_{\Omega_\ell} \lambda \underline{y}_0 \varphi \, dx + \int_{\Gamma_0} \nu \underline{y}_0(x_1, 0) \varphi \, dx_1, \end{aligned}$$

since  $\underline{y}_0$  is a subsolution to (2.1). Using this inequality with  $\varphi = (\underline{y}_1 - \underline{y}_0)^-$  one derives easily

$$\int_{\Omega_\ell} D\nabla(\underline{y}_1 - \underline{y}_0) \cdot \nabla(\underline{y}_1 - \underline{y}_0)^- \, dx + \int_{\Omega_\ell} \lambda(\underline{y}_1 - \underline{y}_0)(\underline{y}_1 - \underline{y}_0)^- \, dx + \int_{\Gamma_0} \nu(\underline{y}_1 - \underline{y}_0)(\underline{y}_1 - \underline{y}_0)^- \, dx_1 \geq 0.$$

Thus it follows that  $\underline{y}_1 \geq \underline{y}_0$ . With a similar proof one gets that  $\overline{y}_1 \leq \overline{y}_0$ . Applying  $S^{n-1}$  to these inequalities leads to

$$S^{n-1}(\underline{y}_0) = \underline{y}_{n-1} \leq \underline{y}_n = S^{n-1}(\underline{y}_1) \quad , \quad S^{n-1}(\overline{y}_0) = \overline{y}_{n-1} \geq \overline{y}_n = S^{n-1}(\overline{y}_1).$$

Furthermore from  $\underline{y}_0 \leq \overline{y}_0$  one derives by applying  $S^n$  to both sides of the inequality

$$\underline{y}_n \leq \overline{y}_n.$$

This completes the proof of (2.5). Then for some functions  $\underline{v}, \overline{v}$  in  $L^2(\Omega_\ell)$  one has

$$\underline{y}_n \rightarrow \underline{v}, \quad \overline{y}_n \rightarrow \overline{v} \quad \text{in } L^2(\Omega_\ell).$$

Clearly  $\underline{v}$  and  $\overline{v}$  are fixed point for  $S$  and thus (see (2.2)) solutions to (2.1). This completes the proof of the lemma.  $\square$

We denote by  $\lambda_1 = \lambda_1(\Omega_\ell)$  the first eigenvalue of the Dirichlet problem in  $\Omega_\ell$  and by  $\varphi_1$  the corresponding first eigenfunction positive and normalised. More precisely  $(\lambda_1, \varphi_1)$  is such that

$$\begin{cases} -\Delta \varphi_1 = \lambda_1 \varphi_1 \text{ in } \Omega_\ell, \\ \varphi_1 = 0 \text{ on } \partial\Omega_\ell, \\ \varphi_1(0, \frac{L}{2}) = 1. \end{cases} \quad (2.6)$$

We suppose that for  $s > 0$  small enough one has

$$\lambda_1 \leq \frac{f(s)}{Ds}. \quad (2.7)$$

Then one has :

**Lemma 2.2.** *Under the assumptions of the preceding lemma and (2.7), for  $\epsilon > 0$  small enough, the maximal solution  $\bar{v}$  to (2.1) satisfies*

$$\epsilon\varphi_1 \leq \bar{v}.$$

In particular  $\bar{v}$  is bounded away from 0.

*Proof.* Due to (2.7) one has for  $\epsilon > 0$  small enough

$$D\lambda_1\epsilon\varphi_1 \leq f(\epsilon\varphi_1).$$

This allows us to show that  $\epsilon\varphi_1$  is a subsolution to (2.1). Indeed, for  $\varphi \in V$ ,  $\varphi \geq 0$  it holds after integration by parts

$$\begin{aligned} \int_{\Omega_\ell} D\nabla(\epsilon\varphi_1) \cdot \nabla\varphi \, dx + \int_{\Gamma_0} \nu\epsilon\varphi_1(x_1, 0)\varphi \, dx_1 \\ &= \int_{\Omega_\ell} D\nabla(\epsilon\varphi_1) \cdot \nabla\varphi \, dx = \int_{\Omega_\ell} \nabla \cdot (D\nabla(\epsilon\varphi_1)\varphi) - D\Delta(\epsilon\varphi_1)\varphi \, dx \\ &= \int_{\partial\Omega_\ell} D\partial_n(\epsilon\varphi_1)\varphi d\sigma + \int_{\Omega_\ell} D\lambda_1(\epsilon\varphi_1)\varphi \, dx \\ &\leq \int_{\Omega_\ell} f(\epsilon\varphi_1)\varphi \, dx + \int_{\Gamma_0} \mu w\varphi \, dx_1. \end{aligned}$$

( $n$  denotes the outward unit normal to  $\Omega_\ell$ , note that  $\partial_n(\epsilon\varphi_1) \leq 0$ ). Thus  $\epsilon\varphi_1$  is a positive subsolution to (2.1).

Then one argues as in the preceding lemma introducing the sequence defined for  $\epsilon$  small by :

$$\begin{aligned} \underline{y}_0 = \epsilon\varphi_1 \leq \frac{\mu}{\nu}m, \quad \overline{y}_0 = \frac{\mu}{\nu}m, \\ \underline{y}_n = S(\underline{y}_{n-1}), \quad \overline{y}_n = S(\overline{y}_{n-1}), \quad n \geq 1. \end{aligned}$$

One has with the same proof as above

$$\underline{y}_0 = \epsilon\varphi_1 \leq \underline{y}_1 \leq \cdots \leq \underline{y}_n \leq \overline{y}_n \leq \cdots \leq \overline{y}_1 \leq \overline{y}_0 = \frac{\mu}{\nu}m.$$

The result follows from the fact that  $\overline{y}_n \rightarrow \bar{v}$ . This completes the proof of the lemma.  $\square$

One has also :

**Lemma 2.3.** *Suppose that*

$$\frac{f(s)}{s} \text{ is decreasing on } (0, +\infty). \quad (2.8)$$

*If  $v_1, v_2$  are positive solutions to (2.1) corresponding to  $w_1, w_2$  respectively then*

$$w_1 \leq w_2 \text{ implies } v_1 \leq v_2.$$

*In particular (2.1) has a unique positive solution.*

*Proof.* Denote by  $\theta$  a smooth function such that

$$\theta(t) = 0 \quad \forall t \leq 0, \quad \theta(t) = 1 \quad \forall t \geq 1, \quad \theta'(t) \geq 0.$$

Set  $\theta_\epsilon(t) = \theta(\frac{t}{\epsilon})$ . Clearly

$$v_1\theta_\epsilon(v_1 - v_2), v_2\theta_\epsilon(v_1 - v_2) \in V.$$

From the equations satisfied by  $v_1, v_2$  one gets, setting  $\theta_\epsilon = \theta_\epsilon(v_1 - v_2)$ ,

$$\begin{aligned} \int_{\Omega_\ell} D\nabla v_1 \cdot \nabla(v_2\theta_\epsilon) \, dx + \int_{\Gamma_0} \nu v_1(x_1, 0)(v_2\theta_\epsilon) \, dx_1 &= \int_{\Omega_\ell} f(v_1)(v_2\theta_\epsilon) \, dx + \int_{\Gamma_0} \mu w_1(v_2\theta_\epsilon) \, dx_1, \\ \int_{\Omega_\ell} D\nabla v_2 \cdot \nabla(v_1\theta_\epsilon) \, dx + \int_{\Gamma_0} \nu v_2(x_1, 0)(v_1\theta_\epsilon) \, dx_1 &= \int_{\Omega_\ell} f(v_2)(v_1\theta_\epsilon) \, dx + \int_{\Gamma_0} \mu w_2(v_1\theta_\epsilon) \, dx_1. \end{aligned}$$

By subtraction we obtain

$$\begin{aligned} \int_{\Omega_\ell} D\{\nabla v_2 \cdot \nabla(v_1\theta_\epsilon) - \nabla v_1 \cdot \nabla(v_2\theta_\epsilon)\} \, dx &= \int_{\Omega_\ell} f(v_2)(v_1\theta_\epsilon) - f(v_1)(v_2\theta_\epsilon) \, dx \\ &\quad + \int_{\Gamma_0} \mu(w_2v_1 - w_1v_2)\theta_\epsilon(v_1 - v_2) \, dx_1. \end{aligned}$$

Clearly the last integral above is nonnegative so that one has

$$\int_{\Omega_\ell} D\{\nabla v_2 \cdot \nabla(v_1\theta_\epsilon) - \nabla v_1 \cdot \nabla(v_2\theta_\epsilon)\} \, dx \geq \int_{\Omega_\ell} f(v_2)(v_1\theta_\epsilon) - f(v_1)(v_2\theta_\epsilon) \, dx.$$

By a simple computation writing  $\theta'_\epsilon$  for  $\theta'_\epsilon(v_1 - v_2)$  one derives

$$\begin{aligned} \int_{\Omega_\ell} f(v_2)(v_1\theta_\epsilon) - f(v_1)(v_2\theta_\epsilon) \, dx &\leq \int_{\Omega_\ell} D\{\nabla v_2 \cdot \nabla(v_1\theta_\epsilon) - \nabla v_1 \cdot \nabla(v_2\theta_\epsilon)\} \, dx \\ &= \int_{\Omega_\ell} D\{\nabla v_2 \cdot \nabla(v_1 - v_2)\theta'_\epsilon v_1 - \nabla v_1 \cdot \nabla(v_1 - v_2)\theta'_\epsilon v_2\} \, dx \\ &= \int_{\Omega_\ell} D\{v_1 \nabla v_2 - v_2 \nabla v_1\} \cdot \nabla(v_1 - v_2)\theta'_\epsilon \, dx \\ &= \int_{\Omega_\ell} D\{v_1 \nabla v_2 - v_2 \nabla v_2 + v_2 \nabla v_2 - v_2 \nabla v_1\} \cdot \nabla(v_1 - v_2)\theta'_\epsilon \, dx \\ &= \int_{\Omega_\ell} D\nabla v_2 \cdot \nabla(v_1 - v_2)(v_1 - v_2)\theta'_\epsilon \, dx - \int_{\Omega_\ell} Dv_2 |\nabla(v_1 - v_2)|^2 \theta'_\epsilon \, dx \\ &\leq \int_{\Omega_\ell} D\nabla v_2 \cdot \nabla(v_1 - v_2)(v_1 - v_2)\theta'_\epsilon \, dx. \end{aligned}$$



Let us set  $\gamma_\epsilon(t) = \int_0^t s\theta'_\epsilon(s)ds$  in such a way that the inequality above reads

$$\int_{\Omega_\ell} f(v_2)(v_1\theta_\epsilon) - f(v_1)(v_2\theta_\epsilon) dx \leq \int_{\Omega_\ell} D\nabla v_2 \cdot \nabla \gamma_\epsilon(v_1 - v_2) dx.$$

From the equation satisfied by  $v_2$ , since  $\gamma_\epsilon(v_1 - v_2) \in V$  and  $v_2, \gamma_\epsilon$  are nonnegative one has

$$\int_{\Omega_\ell} D\nabla v_2 \cdot \nabla \gamma_\epsilon(v_1 - v_2) dx \leq \int_{\Omega_\ell} f(v_2)\gamma_\epsilon(v_1 - v_2) dx + \int_{\Gamma_0} \mu w_2 \gamma_\epsilon(v_1 - v_2) dx_1.$$

Since for some constant  $C$

$$\gamma_\epsilon(t) \leq \int_0^\epsilon s\theta'\left(\frac{s}{\epsilon}\right)\frac{1}{\epsilon}ds \leq C\epsilon$$

the right hand side of the two inequalities above goes to 0 when  $\epsilon \rightarrow 0$ . Since when  $\epsilon \rightarrow 0$  one has  $\theta_\epsilon(v_1 - v_2) \rightarrow \chi_{\{v_1 > v_2\}}$  the characteristic function of the set  $\{v_1 > v_2\} = \{x \in \Omega_\ell \mid v_1(x) > v_2(x)\}$  one gets

$$\int_{\{v_1 > v_2\}} f(v_2)v_1 - f(v_1)v_2 dx \leq 0.$$

But on the set of integration thanks to (2.8) one has  $f(v_2)v_1 - f(v_1)v_2 > 0$  hence the set of integration is necessarily of measure 0, i.e.  $v_1 \leq v_2$ . This completes the proof of the lemma.  $\square$

### 3 The main result

**Theorem 3.1.** *Suppose that (1.2)-(1.6),(2.7),(2.8) hold, then the problem (1.1) admits a non-trivial solution.*

*Proof.* As mentioned above it is of course clear that  $(0, 0)$  is solution to (1.1). Set

$$K = \{v \in L^2(\Gamma_0) \mid 0 \leq v \leq m\}.$$

For  $u \in K$ , let  $\bar{v}$  be the unique positive solution to (2.1) associated to  $w = u$ . For  $\eta \geq L_g$  let  $U = T(u)$  the solution to

$$\begin{cases} U \in H_0^1(\Gamma_0), \\ \int_{\Gamma_0} D'U'\psi' + \mu U\psi + \eta U\psi dx_1 = \int_{\Gamma_0} \nu \bar{v}(x_1, 0)\psi + g(u)\psi + \eta u\psi dx_1 \quad \forall \psi \in H_0^1(\Gamma_0). \end{cases} \quad (3.1)$$

The existence of  $U$  is a consequence of the Lax-Milgram theorem.

We claim that  $T$  is continuous on  $K \subset L^2(\Gamma_0)$ . Indeed suppose that  $u_n \rightarrow u$  in  $K$ . Denote by  $\bar{v}_n$  the solution to (2.1) associated to  $u_n$  i.e. satisfying

$$\int_{\Omega_\ell} D\nabla \bar{v}_n \cdot \nabla \varphi dx + \int_{\Gamma_0} \nu \bar{v}_n(x_1, 0)\varphi dx_1 = \int_{\Omega_\ell} f(\bar{v}_n)\varphi dx + \int_{\Gamma_0} \mu u_n \varphi dx_1 \quad \forall \varphi \in V. \quad (3.2)$$

Since  $\bar{v}_n$  and  $u_n$  are bounded, taking  $\varphi = \bar{v}_n$  in the equality above one gets easily

$$\int_{\Omega_\ell} D|\nabla \bar{v}_n|^2 dx + \int_{\Gamma_0} \nu \bar{v}_n(x_1, 0)^2 dx_1 \leq C,$$

where  $C$  is a constant independent of  $n$ . Thus, up to a subsequence, there exists  $v \in V$  such that

$$\bar{v}_n \rightharpoonup v \text{ in } H^1(\Omega_\ell), \quad \bar{v}_n \rightarrow v \text{ in } L^2(\Omega_\ell), \quad \bar{v}_n(\cdot, 0) \rightarrow v(\cdot, 0) \text{ in } L^2(\Gamma_0).$$

Passing to the limit in (3.2) it follows from Lemma 2.3 that  $v = \bar{v}$  the solution to (2.1) corresponding to  $w = u$ . By uniqueness of the limit one has convergence of the whole sequence and in particular

$$\bar{v}_n(\cdot, 0) \rightarrow \bar{v}(\cdot, 0) \text{ in } L^2(\Gamma_0).$$

Passing to the limit in (3.1) written for  $u = u_n$  one derives  $T(u_n) \rightarrow T(u)$  in  $L^2(\Gamma_0)$ .

We can show also that  $T$  is monotone. Indeed, suppose that  $u_1 \geq u_2$  and set  $U_i = T(u_i)$ ,  $i = 1, 2$ . One has

$$\begin{aligned} \int_{\Gamma_0} D'U_1' \psi' + \mu U_1 \psi + \eta U_1 \psi \, dx_1 &= \int_{\Gamma_0} \nu \bar{v}_1(x_1, 0) \psi + g(u_1) \psi + \eta u_1 \psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0), \\ \int_{\Gamma_0} D'U_2' \psi' + \mu U_2 \psi + \eta U_2 \psi \, dx_1 &= \int_{\Gamma_0} \nu \bar{v}_2(x_1, 0) \psi + g(u_2) \psi + \eta u_2 \psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0). \end{aligned}$$

By subtraction it comes

$$\begin{aligned} \int_{\Gamma_0} D'(U_1 - U_2)' \psi' + \mu(U_1 - U_2) \psi + \eta(U_1 - U_2) \psi \, dx_1 \\ = \int_{\Gamma_0} \nu(\bar{v}_1(x_1, 0) - \bar{v}_2(x_1, 0)) \psi + (g(u_1) - g(u_2)) \psi + \eta(u_1 - u_2) \psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0). \end{aligned}$$

Choosing  $\psi = -(U_1 - U_2)^-$  and taking into account that, by Lemma 2.3,  $\bar{v}_1(x_1, 0) - \bar{v}_2(x_1, 0) \geq 0$  and that for  $\eta \geq L_g$ ,  $(g(u_1) - g(u_2)) + \eta(u_1 - u_2) \geq 0$  (Cf. (1.3), (1.4)), one gets

$$\int_{\Gamma_0} D'|\{(U_1 - U_2)^-\}'|^2 + \mu\{(U_1 - U_2)^-\}^2 + \eta\{(U_1 - U_2)^-\}^2 \, dx_1 \leq 0.$$

Thus  $(U_1 - U_2)^- = 0$  and  $T(u_1) \geq T(u_2)$ .

Next we assert that  $T$  maps  $K$  into itself. Indeed, if  $U_0 = T(0)$  one has, with an obvious notation for  $\bar{v}_0$

$$\int_{\Gamma_0} D'U_0' \psi' + \mu U_0 \psi + \eta U_0 \psi \, dx_1 = \int_{\Gamma_0} \nu \bar{v}_0(x_1, 0) \psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0).$$

Taking  $\psi = -U_0^-$  one deduces easily since  $\bar{v}_0 > 0$  that  $U_0 = T(0) \geq 0$ . Similarly if  $U_m = T(m)$  one has, with an obvious notation for  $\bar{v}_m$

$$\int_{\Gamma_0} D'U_m' \psi' + \mu U_m \psi + \eta U_m \psi \, dx_1 = \int_{\Gamma_0} \nu \bar{v}_m(x_1, 0) \psi + g(m) \psi + \eta m \psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0).$$

Thus choosing  $\psi = (U_m - m)^+$  it comes since  $g(m) \leq 0$ ,  $\bar{v}_m \leq \frac{\mu}{\nu} m$

$$\begin{aligned} \int_{\Gamma_0} D'|\{(U_m - m)^+\}'|^2 + (\mu + \eta)\{(U_m - m)^+\}^2 \, dx_1 \\ = \int_{\Gamma_0} (\nu \bar{v}_m - \mu m)(U_m - m)^+ \, dx_1 \leq 0. \end{aligned}$$

From which it follows that  $U_m \leq m$ . By the monotonicity of  $T$  it results that  $T$  maps the convex  $K$  into itself. But clearly  $T(K) \subset C^{\frac{1}{2}}(\Gamma_0)$  is relatively compact in  $L^2(\Gamma_0)$ . Thus, by the Schauder fixed point theorem (see [8], [9], [5]),  $T$  has a fixed point in  $K$  which leads to a nontrivial solution to (1.1). This completes the proof of the theorem.  $\square$

If it is clear at this point that the solution we constructed is non degenerate in  $v$  it is not clear that the same holds for  $u$ . In fact we have :

**Proposition 3.1.** *Let  $(u, v)$  be the solution constructed in Theorem 3.1. One has*

$$u \neq 0.$$

*Proof.* Suppose that  $u \equiv 0$ . Due to the second equation of (1.1) one has  $v(x_1, 0) = 0$  and from the first equation of (1.1) we get

$$\int_{\Omega_\ell} D\nabla v \cdot \nabla \varphi \, dx = \int_{\Omega_\ell} f(v)\varphi \, dx \quad \forall \varphi \in V. \quad (3.3)$$

Consider then a small ball  $B = B_{x_0}$  centered at  $x_0 \in \Gamma_0$ . Set

$$\tilde{v} = \begin{cases} v & \text{in } \Omega_\ell \cap B, \\ 0 & \text{in the rest of the ball.} \end{cases}$$

Let  $\varphi \in \mathcal{D}(B)$ . One has by (3.3),

$$\int_B D\nabla \tilde{v} \cdot \nabla \varphi \, dx = \int_{\Omega_\ell \cap B} D\nabla v \cdot \nabla \varphi \, dx = \int_{\Omega_\ell \cap B} f(v)\varphi \, dx = \int_B f(\tilde{v})\varphi \, dx \quad \forall \varphi \in \mathcal{D}(B).$$

Thus

$$-D\Delta \tilde{v} = f(\tilde{v}) \text{ in } B.$$

It is clear that  $\tilde{v}$  and thus  $f(\tilde{v})$  are bounded and one has  $f(\tilde{v}) \in L^\infty(B) \subset L^p(B) \forall p$ . From the usual regularity theory it follows that  $\tilde{v} \in W^{2,p}(B) \subset C^{1,\alpha}(B)$ . Since  $f(\tilde{v}) \geq 0, f(\tilde{v}) \neq 0$  it follows that  $\tilde{v} > 0$  in  $B$  (see [9]). Hence a contradiction. This shows the impossibility for  $u$  to be identical to 0 and this completes the proof of the proposition.  $\square$

**Remark 1.** *One can easily show (see [7], [6]) that*

$$\lambda_1 = \lambda_1(\Omega_\ell) = \left(\frac{\pi}{2\ell}\right)^2 + \left(\frac{\pi}{L}\right)^2.$$

*Thus for a smooth function  $f$  it is clear that (2.7) is satisfied if*

$$\lambda_1 = \lambda_1(\Omega_\ell) < \frac{f'(0)}{D},$$

*i.e. for  $\ell$  and  $L$  large enough.*

*Note that (2.8) (see also (1.5)) is satisfied in the case of the Fisher equation i.e. for*

$$f(v) = v(1 - v)$$

*the Lipschitz character of  $f$  being used only on a finite interval.*

## 4 Some extension

We would like to show that our results extend in the case of a so called two road problem. More precisely set

$$\begin{aligned}\Gamma'_0 &= (-\ell, \ell) \times \{L\}, \quad \Gamma_1 = \partial\Omega_\ell \setminus \{\Gamma_0 \cup \Gamma'_0\}, \\ V &= \{v \in H^1(\Omega_\ell) \mid v = 0 \text{ on } \Gamma_1\},\end{aligned}$$

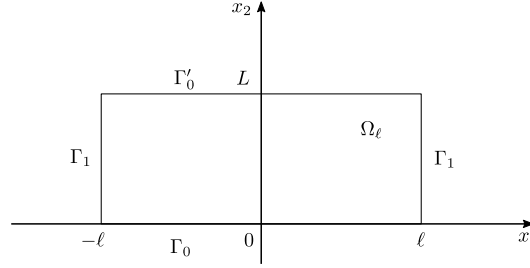


Figure 2: The domain  $\Omega_\ell$  for two-road problem

We consider the problem of finding a  $(u, v, w)$  solution to

$$\left\{ \begin{aligned} &(u, v, w) \in H_0^1(\Gamma_0) \times V \times H_0^1(\Gamma'_0), \\ &\int_{\Omega_\ell} D\nabla v \cdot \nabla \varphi \, dx + \int_{\Gamma_0} \nu v(x_1, 0) \varphi \, dx_1 + \int_{\Gamma'_0} \nu v(x_1, L) \varphi \, dx_1 \\ &\quad = \int_{\Omega_\ell} f(v) \varphi \, dx + \int_{\Gamma_0} \mu u \varphi \, dx_1 + \int_{\Gamma'_0} \mu' w \varphi \, dx_1 \quad \forall \varphi \in V, \\ &\int_{\Gamma_0} D' u' \psi' + \mu u \psi \, dx_1 = \int_{\Gamma_0} \nu v(x_1, 0) \psi \, dx_1 + \int_{\Gamma_0} g(u) \psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0), \\ &\int_{\Gamma'_0} D'' w' \phi' + \mu' w \phi \, dx_1 = \int_{\Gamma'_0} \nu' v(x_1, L) \phi \, dx_1 + \int_{\Gamma'_0} h(w) \phi \, dx_1 \quad \forall \phi \in H_0^1(\Gamma'_0). \end{aligned} \right. \quad (4.1)$$

Here we assume that

$$D, D', D'', \mu, \nu, \mu', \nu' \text{ are positive constants,}$$

$f, g, h$  are Lipschitz continuous functions with Lipschitz constants  $L_f, L_g, L_h$  respectively (Cf. (1.3)), which implies that for  $\lambda \geq L_f$ ,  $\eta \geq L_g$  and  $\xi \geq L_h$  the functions

$$x \rightarrow \lambda x - f(x), \quad \eta x - g(x), \quad \xi x - h(x)$$

are nondecreasing. We will suppose that  $f$  satisfies (1.5) and that

$$g(0) = 0, \quad h(0) = 0.$$

Since  $\Gamma_0$  and  $\Gamma'_0$  are playing exactly identical roles there is no loss of generality in assuming for instance

$$\frac{\mu}{\nu} \geq \frac{\mu'}{\nu'}.$$

Then for

$$m \geq \frac{\nu}{\mu}, \quad m' = \frac{\nu'}{\mu'} \frac{\mu}{\nu} m,$$

we will assume

$$g(m) \leq 0, \quad h(m') \leq 0. \quad (4.2)$$

One should notice the following properties

$$\begin{aligned} \frac{\mu}{\nu}m &\geq 1, \quad \frac{\mu'}{\nu'}m' = \frac{\mu}{\nu}m. \\ m' &= \frac{\nu'}{\mu'}\frac{\mu}{\nu}m \geq \frac{\nu'}{\mu'}. \end{aligned}$$

Then with small variants we can reproduce the results we had in the preceding sections. First we have

**Lemma 4.1.** *Suppose that  $\tilde{u}, \tilde{w}$  are measurable functions on  $\Gamma_0$  and  $\Gamma'_0$  respectively such that*

$$0 \leq \tilde{u} \leq m, \quad 0 \leq \tilde{w} \leq m'.$$

*Then under the assumptions above the problem*

$$\begin{cases} v \in V, \\ \int_{\Omega_\ell} D\nabla v \cdot \nabla \varphi \, dx + \int_{\Gamma_0} \nu v(x_1, 0) \varphi(x_1, 0) \, dx_1 + \int_{\Gamma'_0} \nu' v(x_1, L) \varphi(x_1, L) \, dx_1 \\ \quad = \int_{\Omega_\ell} f(v) \varphi \, dx + \int_{\Gamma_0} \mu \tilde{u} \varphi(x_1, 0) \, dx_1 + \int_{\Gamma'_0} \mu' \tilde{w} \varphi(x_1, L) \, dx_1 \quad \forall \varphi \in V, \end{cases} \quad (4.3)$$

*possesses a minimal and a maximal solution with values in  $(0, \frac{\mu}{\nu}m)$ .*

*Proof.* Let us remark first that any nonnegative solution to (4.3) takes its values in  $(0, \frac{\mu}{\nu}m)$ . Indeed if  $v$  is solution to (4.3) taking as test function  $\varphi = (v - k)^+$ ,  $k = \frac{\mu}{\nu}m \geq 1$  one gets

$$\begin{aligned} \int_{\Omega_\ell} D|\nabla(v - k)^+|^2 \, dx &= \int_{\Omega_\ell} D\nabla(v - k) \cdot \nabla(v - k)^+ \, dx = \int_{\Omega_\ell} D\nabla v \cdot \nabla(v - k)^+ \, dx = \\ &= \int_{\Omega_\ell} f(v)(v - k)^+ \, dx + \int_{\Gamma_0} \{\mu \tilde{u} - \nu v(x_1, 0)\}(v - k)^+ \, dx_1 + \int_{\Gamma'_0} \{\mu' \tilde{w} - \nu' v(x_1, L)\}(v - k)^+ \, dx_1 \\ &\leq \int_{\Gamma_0} \{\mu \tilde{u} - \nu v(x_1, 0)\}(v - k)^+ \, dx_1 + \int_{\Gamma'_0} \{\mu' \tilde{w} - \nu' v(x_1, L)\}(v - k)^+ \, dx_1 \leq 0, \end{aligned}$$

since on the set where  $v \geq k = \frac{\mu}{\nu}m = \frac{\mu'}{\nu'}m'$  one has  $\{\mu \tilde{u} - \nu v(x_1, 0)\} \leq \{\mu \tilde{u} - \mu m\} \leq 0$  and  $\{\mu' \tilde{w} - \nu' v(x_1, L)\} \leq \{\mu' \tilde{w} - \mu' m'\} \leq 0$ .

Next let us note that 0 is a subsolution to (4.3). Indeed  $\forall \varphi \in V, \varphi \geq 0$ , one has

$$\begin{aligned} \int_{\Omega_\ell} D\nabla 0 \cdot \nabla \varphi \, dx + \int_{\Gamma_0} \nu 0 \varphi \, dx_1 + \int_{\Gamma'_0} \nu' 0 \varphi \, dx_1 \\ \leq \int_{\Omega_\ell} f(0) \varphi \, dx + \int_{\Gamma_0} \mu \tilde{u} \varphi \, dx_1 + \int_{\Gamma'_0} \mu' \tilde{w} \varphi \, dx_1. \end{aligned}$$

On the other hand,  $k = \frac{\mu}{\nu}m$  is a supersolution since for  $\varphi \in V, \varphi \geq 0$ ,

$$\int_{\Omega_\ell} D\nabla k \cdot \nabla \varphi \, dx + \int_{\Gamma_0} \nu k \varphi \, dx_1 + \int_{\Gamma'_0} \nu' k \varphi \, dx_1 \geq \int_{\Omega_\ell} f(k) \varphi \, dx + \int_{\Gamma_0} \mu \tilde{u} \varphi \, dx_1 + \int_{\Gamma'_0} \mu' \tilde{w} \varphi \, dx_1.$$

For  $z \in L^2(\Omega_\ell)$  we denote by  $y = S(z)$  the solution to

$$\begin{cases} y \in V, \\ \int_{\Omega_\ell} D\nabla y \cdot \nabla \varphi \, dx + \int_{\Omega_\ell} \lambda y \varphi \, dx + \int_{\Gamma_0} \nu y(x_1, 0) \varphi \, dx_1 + \int_{\Gamma'_0} \nu' y(x_1, L) \varphi \, dx_1 \\ \quad = \int_{\Omega_\ell} f(z) \varphi \, dx + \int_{\Omega_\ell} \lambda z \varphi \, dx + \int_{\Gamma_0} \mu \tilde{w} \varphi \, dx_1 + \int_{\Gamma'_0} \mu' \tilde{w} \varphi \, dx_1 \quad \forall \varphi \in V, \end{cases}$$

where  $\lambda \geq L_f$ . The existence of a unique solution  $y$  to the problem above follows from the Lax-Milgram theorem. Then reproducing the arguments of Lemma 2.1 it is easy to show that  $S$  is continuous and monotone. Introducing the sequence defined in (2.4) one concludes as in the Lemma 2.1 to the existence of a minimal and maximal solution  $\underline{v}$  and  $\bar{v}$ .  $\square$

Then one has :

**Lemma 4.2.** *Under the assumptions of the preceding lemma and (2.7), for  $\epsilon > 0$  small enough, the maximal solution  $\bar{v}$  to (4.3) satisfies*

$$\epsilon \varphi_1 \leq \bar{v}.$$

In particular  $\bar{v}$  is bounded away from 0.

*Proof.* Due to (2.7) one has for  $\epsilon > 0$  small enough

$$D\lambda_1 \epsilon \varphi_1 \leq f(\epsilon \varphi_1).$$

Then for  $\varphi \in V$ ,  $\varphi \geq 0$  it holds after integration by parts

$$\begin{aligned} & \int_{\Omega_\ell} D\nabla(\epsilon \varphi_1) \cdot \nabla \varphi \, dx + \int_{\Gamma_0} \nu \epsilon \varphi_1(x_1, 0) \varphi \, dx_1 + \int_{\Gamma'_0} \nu' \epsilon \varphi_1(x_1, L) \varphi \, dx_1 \\ &= \int_{\Omega_\ell} D\nabla(\epsilon \varphi_1) \cdot \nabla \varphi \, dx = \int_{\Omega_\ell} \nabla \cdot (D\nabla(\epsilon \varphi_1) \varphi) - D\Delta(\epsilon \varphi_1) \varphi \, dx \\ &= \int_{\partial\Omega_\ell} D\partial_n(\epsilon \varphi_1) \varphi \, d\sigma + \int_{\Omega_\ell} D\lambda_1(\epsilon \varphi_1) \varphi \, dx \\ &\leq \int_{\Omega_\ell} f(\epsilon \varphi_1) \varphi \, dx + \int_{\Gamma_0} \mu \tilde{w} \varphi \, dx_1 + \int_{\Gamma'_0} \mu' \tilde{w} \varphi \, dx_1. \end{aligned}$$

( $n$  denotes the outward unit normal to  $\Omega_\ell$ , note that  $\partial_n(\epsilon \varphi_1) \leq 0$ ). Thus  $\epsilon \varphi_1$  is a positive subsolution to (4.3) and one concludes as in the proof of Lemma 2.2.  $\square$

Analogous to Lemma 2.3 one has :

**Lemma 4.3.** *Suppose that  $f$  satisfies (2.8). If  $v_1, v_2$  are positive solutions to (4.3) corresponding to  $(u_1, w_1)$  and  $(u_2, w_2)$  respectively then*

$$u_1 \leq u_2 \text{ and } w_1 \leq w_2 \text{ implies } v_1 \leq v_2.$$

In particular (4.3) has a unique positive solution.

*Proof.* Denote by  $\theta$  a smooth function such that

$$\theta(t) = 0 \quad \forall t \leq 0, \quad \theta(t) = 1 \quad \forall t \geq 1, \quad \theta'(t) \geq 0.$$

Set  $\theta_\epsilon(t) = \theta(\frac{t}{\epsilon})$ . Clearly

$$v_1\theta_\epsilon(v_1 - v_2), v_2\theta_\epsilon(v_1 - v_2) \in V.$$

From the equations satisfied by  $v_1, v_2$  one gets, setting  $\theta_\epsilon = \theta_\epsilon(v_1 - v_2)$ ,

$$\begin{aligned} & \int_{\Omega_\ell} D\nabla v_1 \cdot \nabla(v_2\theta_\epsilon) \, dx + \int_{\Gamma_0} \nu v_1(x_1, 0)(v_2\theta_\epsilon) \, dx_1 + \int_{\Gamma'_0} \nu' v_1(x_1, L)(v_2\theta_\epsilon) \, dx_1 \\ &= \int_{\Omega_\ell} f(v_1)(v_2\theta_\epsilon) \, dx + \int_{\Gamma_0} \mu u_1(v_2\theta_\epsilon) \, dx_1 + \int_{\Gamma'_0} \mu' w_1(v_2\theta_\epsilon) \, dx_1, \\ & \int_{\Omega_\ell} D\nabla v_2 \cdot \nabla(v_1\theta_\epsilon) \, dx + \int_{\Gamma_0} \nu v_2(x_1, 0)(v_1\theta_\epsilon) \, dx_1 + \int_{\Gamma'_0} \nu' v_2(x_1, L)(v_1\theta_\epsilon) \, dx_1 \\ &= \int_{\Omega_\ell} f(v_2)(v_1\theta_\epsilon) \, dx + \int_{\Gamma_0} \mu u_2(v_1\theta_\epsilon) \, dx_1 + \int_{\Gamma'_0} \mu' w_2(v_1\theta_\epsilon) \, dx_1. \end{aligned}$$

By subtraction we obtain

$$\begin{aligned} \int_{\Omega_\ell} D\{\nabla v_2 \cdot \nabla(v_1\theta_\epsilon) - \nabla v_1 \cdot \nabla(v_2\theta_\epsilon)\} \, dx &= \int_{\Omega_\ell} f(v_2)(v_1\theta_\epsilon) - f(v_1)(v_2\theta_\epsilon) \, dx \\ &+ \int_{\Gamma_0} \mu(u_2v_1 - u_1v_2)\theta_\epsilon(v_1 - v_2) \, dx_1 + \int_{\Gamma'_0} \mu'(w_2v_1 - w_1v_2)\theta_\epsilon(v_1 - v_2) \, dx_1. \end{aligned}$$

Clearly the last two integrals above are nonnegative so that one has

$$\int_{\Omega_\ell} D\{\nabla v_2 \cdot \nabla(v_1\theta_\epsilon) - \nabla v_1 \cdot \nabla(v_2\theta_\epsilon)\} \, dx \geq \int_{\Omega_\ell} f(v_2)(v_1\theta_\epsilon) - f(v_1)(v_2\theta_\epsilon) \, dx.$$

Then the rest of the proof is like in Lemma 2.3.  $\square$

Then we can now show :

**Theorem 4.1.** *Under the assumptions above the problem (4.1) admits a nontrivial solution.*

*Proof.* It is of course clear that  $(0, 0, 0)$  is solution to (4.1). Set

$$K = \{u \in L^2(\Gamma_0) \mid 0 \leq u \leq m\}, \quad K' = \{w \in L^2(\Gamma'_0) \mid 0 \leq w \leq m'\}.$$

For  $(u, w) \in K \times K'$ , let  $\bar{v}$  be the unique positive solution to (4.3) associated to  $(\tilde{u}, \tilde{w}) = (u, w)$ .

For  $\eta \geq L_g, \xi \geq L_h$ , let  $(U, W) = T(u, w)$  be the solution to

$$\begin{cases} (U, W) \in H_0^1(\Gamma_0) \times H_0^1(\Gamma'_0), \\ \int_{\Gamma_0} D'U'\psi' + \mu U\psi + \eta U\psi \, dx_1 = \int_{\Gamma_0} \nu \bar{v}(x_1, 0)\psi + g(u)\psi + \eta u\psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0), \\ \int_{\Gamma'_0} D''W'\phi' + \mu' W\phi + \xi W\phi \, dx_1 \\ \quad = \int_{\Gamma'_0} \nu' \bar{v}(x_1, L)\phi + h(w)\phi + \xi w\phi \, dx_1 \quad \forall \phi \in H_0^1(\Gamma'_0). \end{cases}$$

The existence of  $(U, W)$  is a consequence of the Lax-Milgram theorem.

We show as in Theorem 3.1 that  $T$  is continuous on  $K \times K' \subset L^2(\Gamma_0) \times L^2(\Gamma'_0)$ . Indeed suppose that  $u_n \rightarrow u$  in  $K$  and  $w_n \rightarrow w$  in  $K'$ . Denote by  $\bar{v}_n$  the solution to (4.3) associated to  $(u_n, w_n)$  i.e. satisfying

$$\begin{aligned} \int_{\Omega_\ell} D\nabla\bar{v}_n \cdot \nabla\varphi \, dx + \int_{\Gamma_0} \nu\bar{v}_n(x_1, 0)\varphi \, dx_1 + \int_{\Gamma'_0} \nu'\bar{v}_n(x_1, L)\varphi \, dx_1 \\ = \int_{\Omega_\ell} f(\bar{v}_n)\varphi \, dx + \int_{\Gamma_0} \mu u_n \varphi \, dx_1 + \int_{\Gamma'_0} \mu' w_n \varphi \, dx_1 \quad \forall \varphi \in V. \end{aligned} \quad (4.4)$$

Since  $\bar{v}_n$ ,  $u_n$  and  $w_n$  are bounded, taking  $\varphi = \bar{v}_n$  in the equality above one gets easily

$$\int_{\Omega_\ell} D|\nabla\bar{v}_n|^2 \, dx + \int_{\Gamma_0} \nu\bar{v}_n(x_1, 0)^2 \, dx_1 + \int_{\Gamma'_0} \nu'\bar{v}_n(x_1, L)^2 \, dx_1 \leq C,$$

where  $C$  is a constant independent of  $n$ . Thus, up to a subsequence, there exists  $v \in V$  such that

$$\begin{aligned} \bar{v}_n \rightharpoonup v \text{ in } H^1(\Omega_\ell), \quad \bar{v}_n \rightarrow v \text{ in } L^2(\Omega_\ell), \\ \bar{v}_n(\cdot, 0) \rightarrow v(\cdot, 0) \text{ in } L^2(\Gamma_0), \quad \bar{v}_n(\cdot, L) \rightarrow v(\cdot, L) \text{ in } L^2(\Gamma'_0). \end{aligned}$$

Passing to the limit in (4.4) one derives as in Theorem 3.1 that  $T(u_n, w_n) \rightarrow T(u, w)$  in  $L^2(\Gamma_0) \times L^2(\Gamma'_0)$ .

We can show also that  $T$  is monotone. Indeed, suppose that  $(u_1, w_1) \geq (u_2, w_2)$  in the sense that  $u_1 \geq u_2$  and  $w_1 \geq w_2$  and set  $(U_i, W_i) = T(u_i, w_i)$ ,  $i = 1, 2$ . First, for  $U_i$  one has

$$\begin{aligned} \int_{\Gamma_0} D'U_1'\psi' + \mu U_1\psi + \eta U_1\psi \, dx_1 &= \int_{\Gamma_0} \nu\bar{v}_1(x_1, 0)\psi + g(u_1)\psi + \eta u_1\psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0), \\ \int_{\Gamma_0} D'U_2'\psi' + \mu U_2\psi + \eta U_2\psi \, dx_1 &= \int_{\Gamma_0} \nu\bar{v}_2(x_1, 0)\psi + g(u_2)\psi + \eta u_2\psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0). \end{aligned}$$

By subtraction it comes

$$\begin{aligned} \int_{\Gamma_0} D'(U_1 - U_2)'\psi' + \mu(U_1 - U_2)\psi + \eta(U_1 - U_2)\psi \, dx_1 \\ = \int_{\Gamma_0} \nu(\bar{v}_1(x_1, 0) - \bar{v}_2(x_1, 0))\psi + (g(u_1) - g(u_2))\psi + \eta(u_1 - u_2)\psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0). \end{aligned}$$

Choosing  $\psi = -(U_1 - U_2)^-$  and taking into account that, by Lemma 4.3,  $\bar{v}_1(x_1, 0) - \bar{v}_2(x_1, 0) \geq 0$  and that for  $\eta \geq L_g$ ,  $(g(u_1) - g(u_2)) + \eta(u_1 - u_2) \geq 0$  (Cf. (1.3), (1.4)), one gets

$$\int_{\Gamma_0} D'|\{(U_1 - U_2)^-\}'|^2 + \mu\{(U_1 - U_2)^-\}^2 + \eta\{(U_1 - U_2)^-\}^2 \, dx_1 \leq 0.$$

Thus  $(U_1 - U_2)^- = 0$  and  $U_1 \geq U_2$ . Similarly one shows that  $W_1 \geq W_2$ .



Next we assert that  $T$  maps  $K \times K'$  into itself. Indeed, if  $(U_0, W_0) = T(0, 0)$  one has, with an obvious notation for  $\bar{v}_0$

$$\int_{\Gamma_0} D'U_0'\psi' + \mu U_0\psi + \eta U_0\psi \, dx_1 = \int_{\Gamma_0} \nu \bar{v}_0(x_1, 0)\psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0).$$

$$\int_{\Gamma'_0} D''W_0'\phi' + \mu W_0\phi + \xi W_0\phi \, dx_1 = \int_{\Gamma'_0} \nu' \bar{v}_0(x_1, L)\phi \, dx_1 \quad \forall \phi \in H_0^1(\Gamma'_0).$$

Taking  $\psi = -U_0^-$ ,  $\phi = -W_0^-$ , one deduces easily since  $\bar{v}_0(x_1, 0) \geq 0$  and  $\bar{v}_0(x_1, L) \geq 0$  that  $U_0 \geq 0$  and  $W_0 \geq 0$ . Similarly, if  $(U_1, W_1) = T(m, m')$  one has, with an obvious notation for  $\bar{v}_1$

$$\int_{\Gamma_0} D'U_1'\psi' + \mu U_1\psi + \eta U_1\psi \, dx_1 = \int_{\Gamma_0} \nu \bar{v}_1(x_1, 0)\psi + g(m)\psi + \eta m\psi \, dx_1 \quad \forall \psi \in H_0^1(\Gamma_0),$$

$$\int_{\Gamma'_0} D''W_1'\phi' + \mu' W_1\phi + \xi W_1\phi \, dx_1 = \int_{\Gamma'_0} \nu' \bar{v}_1(x_1, L)\phi + h(m')\phi + \xi m'\phi \, dx_1 \quad \forall \phi \in H_0^1(\Gamma'_0).$$

Thus choosing  $\psi = (U_1 - m)^+$  and  $\phi = (W_1 - m')^+$ , due to (4.2) and  $\bar{v}_1 \leq \frac{\mu}{\nu}m = \frac{\mu'}{\nu'}m'$  it comes

$$\begin{aligned} \int_{\Gamma_0} D'\{|(U_1 - m)^+\}'|^2 + (\mu + \eta)\{(U_1 - m)^+\}^2 \\ \leq \int_{\Gamma_0} (\nu \bar{v}_1(x_1, 0) - \mu m)(U_1 - m)^+ dx_1 \leq 0. \end{aligned}$$

$$\begin{aligned} \int_{\Gamma'_0} D''\{|(W_1 - m')^+\}'|^2 + (\mu' + \xi)\{(W_1 - m')^+\}^2 \\ \leq \int_{\Gamma'_0} (\nu' \bar{v}_1(x_1, L) - \mu' m')(W_1 - m')^+ dx_1 \leq 0. \end{aligned}$$

From which it follows that  $U_1 \leq m$ ,  $W_1 \leq m'$ . By the monotonicity of  $T$  it results that  $T$  maps the convex  $K \times K'$  into itself. But clearly  $T(K \times K') \subset C^{\frac{1}{2}}(\Gamma_0) \times C^{\frac{1}{2}}(\Gamma'_0)$  is relatively compact in  $L^2(\Gamma_0) \times L^2(\Gamma'_0)$ . Thus, by the Schauder fixed point theorem (see [8], [9], [5]),  $T$  has a fixed point in  $K \times K'$  which leads to a nontrivial solution  $(u, v, w)$  to (4.1). This completes the proof of the theorem.  $\square$

**Remark 2.** *One can show as in Proposition 3.1 that  $u, w$  are also non degenerate in the sense that*

$$u \not\equiv 0, \quad w \not\equiv 0.$$

## 5 The case of an unbounded domain

The goal of this section is to show that when  $\ell = +\infty$  it remains possible to define and find a nontrivial solution to problem (1.1). Let us introduce some notation. For convenience we will denote by  $V_\ell$  the space  $V$  defined in section 1. Similarly we will indicate the dependence in  $\ell$  for  $\Gamma_0$  i.e

$$\Gamma_0 = \Gamma_0^\ell = (-\ell, \ell) \times \{0\}.$$

When convenient we will set  $I_\ell = (-\ell, \ell)$ . In addition we set

$$\Omega_\infty = \mathbb{R} \times (0, L), \quad \Gamma_0^\infty = \mathbb{R} \times \{0\}, \quad \Gamma_1^\infty = \mathbb{R} \times \{L\},$$

$$V_\infty = \{v \in H_{loc}^1(\overline{\Omega_\infty}) \mid v = 0 \text{ on } \Gamma_1^\infty\},$$

where

$$H_{loc}^1(\overline{\Omega_\infty}) = \{v \mid v \in H^1(\Omega_{\ell_0}) \forall \ell_0 > 0\}.$$

Then we have

**Theorem 5.1.** *Under the assumption above, in particular (1.2), (1.3), (1.5), (1.6), (2.7), (2.8) there exists  $(u, v)$  nontrivial solution to*

$$\begin{cases} (u, v) \in H_0^1(\Gamma_0^\infty) \times V_\infty, \\ \int_{\Omega_{\ell_0}} D\nabla v \cdot \nabla \varphi \, dx + \int_{I_{\ell_0}} \nu v(x_1, 0) \varphi \, dx_1 \\ \qquad \qquad \qquad = \int_{\Omega_{\ell_0}} f(v) \varphi \, dx + \int_{I_{\ell_0}} \mu u \varphi \, dx_1 \quad \forall \varphi \in V_{\ell_0}, \forall \ell_0, \\ \int_{I_{\ell_0}} D'u' \psi' + \mu u \psi \, dx_1 = \int_{I_{\ell_0}} \nu v(x_1, 0) \psi \, dx_1 + \int_{I_{\ell_0}} g(u) \psi \, dx_1 \quad \forall \psi \in H_0^1(I_{\ell_0}), \forall \ell_0. \end{cases} \quad (5.1)$$

(We identify  $\Gamma_0^\infty$  with  $\mathbb{R}$ . Recall that  $I_\ell = (-\ell, \ell)$ ).

*Proof.* Let  $(u_\ell, v_\ell)$  be a solution to (1.1). We can find such a solution for every  $\ell > 0$  (Cf. Theorem 3.1). One notices that for  $\ell' \geq \ell$  one has

$$\Omega_\ell \subset \Omega_{\ell'}, \quad H_0^1(\Omega_\ell) \subset H_0^1(\Omega_{\ell'}),$$

(we suppose the functions of  $H_0^1(\Omega_\ell)$  extended by 0 outside  $\Omega_\ell$ ). By definition of  $\lambda_1 = \lambda_1(\Omega_\ell)$  one has

$$\lambda_1(\Omega_\ell) = \inf_{H_0^1(\Omega_\ell) \setminus \{0\}} \frac{\int_{\Omega_\ell} |\nabla v|^2 \, dx}{\int_{\Omega_\ell} v^2 \, dx},$$

and thus clearly

$$\lambda_1(\Omega_{\ell'}) \leq \lambda_1(\Omega_\ell) \quad \forall \ell' \geq \ell.$$

Let us assume for some  $\ell_1 > 0$  (Cf. (2.7))

$$\lambda_1(\Omega_{\ell_1}) \leq \frac{f(s)}{Ds} \quad \text{for } s > 0 \text{ small enough.} \quad (5.2)$$

Then for any  $\ell \geq \ell_1$  one has for for  $s > 0$  small enough

$$\lambda_1(\Omega_\ell) \leq \frac{f(s)}{Ds}.$$

Moreover, since it is easy to see that  $\varphi_1$  defined in (2.6) is given by

$$\varphi_1 = \sin \frac{\pi}{2\ell}(x_1 + \ell) \sin \frac{\pi}{L} x_2,$$

one has  $0 \leq \varphi_1 \leq 1$  and if (5.2) holds one has

$$D\lambda_1(\Omega_\ell)\epsilon\varphi_1 \leq f(\epsilon\varphi_1)$$

for  $\epsilon > 0$  small enough independently of  $\ell \geq \ell_1$ . We suppose from now on that this  $\epsilon$  is fixed such that if  $(u_\ell, v_\ell)$  is a solution to (1.1) constructed as in Theorem 3.1 one has

$$\epsilon\varphi_1 \leq v_\ell$$

and in particular for every  $\ell \geq \ell_1$

$$\epsilon(\sin \frac{\pi}{4})^2 \leq \epsilon\varphi_1 \leq v_\ell \quad \text{a.e. } x \in (-\frac{\ell}{2}, \frac{\ell}{2}) \times (\frac{L}{4}, \frac{3L}{4}). \quad (5.3)$$

One should also notice that independently of  $\ell$  one has

$$0 \leq u_\ell \leq m, \quad \epsilon\varphi_1 \leq v_\ell \leq \frac{\mu}{\nu}m. \quad (5.4)$$

We assume from now on  $\ell \geq \ell_1$  and for  $\ell_0 \leq \ell - 1$  we define  $\rho$  by

$$\rho = \rho(x_1) = \begin{cases} 1 & \text{on } I_{\ell_0}, \\ x_1 + \ell_0 + 1 & \text{on } (-\ell_0 - 1, -\ell_0), \\ -x_1 + \ell_0 + 1 & \text{on } (\ell_0, \ell_0 + 1), \\ 0 & \text{outside } I_{\ell_0+1}, \end{cases}$$

whose graph is depicted below.

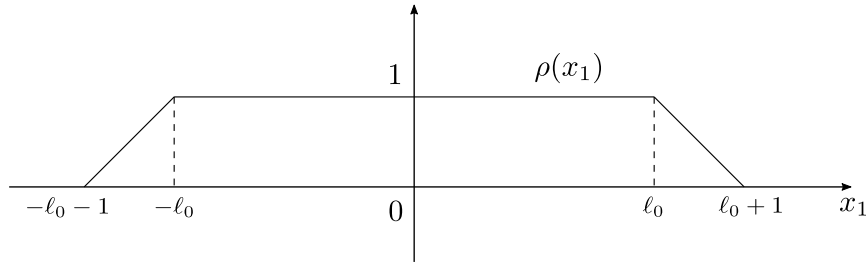


Figure 3: The graph of the function  $\rho(x_1)$

Clearly  $\rho^2 v_\ell = \rho^2(x_1)v_\ell \in V_\ell$  and from the first equation of (1.1) one gets

$$\int_{\Omega_\ell} D\nabla v_\ell \cdot \nabla(\rho^2 v_\ell) dx + \int_{\Gamma_0} \nu \rho^2 v_\ell^2(x_1, 0) dx_1 = \int_{\Omega_\ell} f(v_\ell) \rho^2 v_\ell dx + \int_{\Gamma_0} \mu u_\ell \rho^2 v_\ell(x_1, 0) dx_1. \quad (5.5)$$

One should notice that in the integrals in  $\Omega_\ell$  one integrates only on  $\Omega_{\ell_0+1}$  and for the ones in  $\Gamma_0$  on  $I_{\ell_0+1}$ . Then remark that

$$\int_{\Omega_\ell} \nabla v_\ell \cdot \nabla(\rho^2 v_\ell) dx = \int_{\Omega_\ell} |\nabla v_\ell|^2 \rho^2 + 2\rho v_\ell \nabla v_\ell \cdot \nabla \rho dx,$$

and

$$\int_{\Omega_\ell} |\nabla(\rho v_\ell)|^2 dx = \int_{\Omega_\ell} |\rho \nabla v_\ell + v_\ell \nabla \rho|^2 dx = \int_{\Omega_\ell} |\nabla v_\ell|^2 \rho^2 + 2\rho v_\ell \nabla v_\ell \cdot \nabla \rho + v_\ell^2 |\nabla \rho|^2 dx.$$

From this it follows that

$$\int_{\Omega_\ell} \nabla v_\ell \cdot \nabla(\rho^2 v_\ell) dx = \int_{\Omega_\ell} |\nabla(\rho v_\ell)|^2 dx - \int_{\Omega_\ell} v_\ell^2 |\nabla \rho|^2 dx.$$

Thus, since the second integral of (5.5) is nonnegative, it comes

$$D \int_{\Omega_{\ell_0+1}} |\nabla(\rho v_\ell)|^2 dx \leq D \int_{\Omega_{\ell_0+1}} v_\ell^2 |\nabla \rho|^2 dx + \int_{\Omega_{\ell_0+1}} f(v_\ell) \rho^2 v_\ell dx + \int_{I_{\ell_0+1}} \mu u_\ell \rho^2 v_\ell(x_1, 0) dx_1.$$

Using the definition of  $\rho$  and in particular the fact that  $\rho = 1$  on  $\Omega_{\ell_0}$  we get easily by (5.4)

$$\int_{\Omega_{\ell_0}} |\nabla v_\ell|^2 dx \leq C \tag{5.6}$$

where  $C$  is independent of  $\ell$ . Taking now  $\psi = \rho^2 u_\ell$  in the second equation of (1.1) we get

$$\int_{I_{\ell_0+1}} D' u'_\ell (\rho^2 u_\ell)' + \mu \rho^2 u_\ell^2 dx_1 = \int_{I_{\ell_0+1}} \nu v_\ell(x_1, 0) \rho^2 u_\ell + g(u_\ell) \rho^2 u_\ell dx_1.$$

Arguing as above we derive easily

$$\int_{I_{\ell_0+1}} u'_\ell (\rho^2 u_\ell)' dx_1 = \int_{I_{\ell_0+1}} |(\rho u_\ell)'|^2 dx_1 - \int_{I_{\ell_0+1}} u_\ell^2 \rho'^2 dx_1.$$

This leads to

$$\int_{I_{\ell_0+1}} D' |(\rho u_\ell)'|^2 + \mu \rho^2 u_\ell^2 dx_1 \leq \int_{I_{\ell_0+1}} D' u_\ell^2 \rho'^2 dx_1 + \nu v_\ell(x_1, 0) \rho^2 u_\ell + g(u_\ell) \rho^2 u_\ell dx_1.$$

Integrating only on  $I_{\ell_0}$  in the first integral i.e. where  $\rho = 1$  we obtain

$$\int_{I_{\ell_0}} (u'_\ell)^2 + u_\ell^2 dx_1 \leq C \tag{5.7}$$

where  $C$  is some other constant independent of  $\ell$ . It results from (5.6), (5.7) that  $(u_\ell, v_\ell)$  is bounded in  $H^1(I_{\ell_0}) \times V_{\ell_0}$  independently of  $\ell$ . Thus there exists a subsequence of  $(u_\ell, v_\ell)$  that we will denote by  $(u_{n,0}, v_{n,0})$  such when  $n \rightarrow \infty$

$$u_{n,0} \rightharpoonup u^0 \text{ in } H^1(I_{\ell_0}), \quad v_{n,0} \rightharpoonup v^0 \text{ in } V_{\ell_0}, \quad u_{n,0} \rightarrow u^0 \text{ in } L^2(I_{\ell_0}), \quad v_{n,0} \rightarrow v^0 \text{ in } L^2(\Omega_{\ell_0}),$$

$$v_{n,0}(\cdot, 0) \rightarrow v^0(\cdot, 0) \text{ in } L^2(I_{\ell_0}).$$

Considering the equations

$$\int_{\Omega_{\ell_0}} D \nabla v_\ell \cdot \nabla \varphi dx + \int_{I_{\ell_0}} \nu v_\ell(x_1, 0) \varphi dx_1 = \int_{\Omega_{\ell_0}} f(v_\ell) \varphi dx + \int_{I_{\ell_0}} \mu u_\ell \varphi dx_1 \quad \forall \varphi \in V_{\ell_0},$$

$$\int_{I_{\ell_0}} D' u_{\ell}' \psi' + \mu u_{\ell} \psi \, dx_1 = \int_{I_{\ell_0}} \nu v_{\ell}(x_1, 0) \psi \, dx_1 + \int_{I_{\ell_0}} g(u_{\ell}) \psi \, dx_1 \quad \forall \psi \in H_0^1(I_{\ell_0}),$$

with  $(u_{\ell}, v_{\ell})$  replaced by  $(u_{n,0}, v_{n,0})$ , one can pass to the limit in  $n$  and see that  $(u^0, v^0) \in H^1(I_{\ell_0}) \times V_{\ell_0}$  satisfies

$$\int_{\Omega_{\ell_0}} D \nabla v^0 \cdot \nabla \varphi \, dx + \int_{I_{\ell_0}} \nu v^0(x_1, 0) \varphi \, dx_1 = \int_{\Omega_{\ell_0}} f(v^0) \varphi \, dx + \int_{I_{\ell_0}} \mu u^0 \varphi \, dx_1 \quad \forall \varphi \in V_{\ell_0},$$

$$\int_{I_{\ell_0}} D' u^{0'} \psi' + \mu u^0 \psi \, dx_1 = \int_{I_{\ell_0}} \nu v^0(x_1, 0) \psi \, dx_1 + \int_{I_{\ell_0}} g(u^0) \psi \, dx_1 \quad \forall \psi \in H_0^1(I_{\ell_0}).$$

(Note that a function of  $V_{\ell_0}$  extended by 0 belongs to  $V_{\ell}$ ). Clearly -as a subsequence of  $(u_{\ell}, v_{\ell})$ - the sequence  $(u_{n,0}, v_{n,0})$  is bounded in  $H^1(I_{\ell_0+1}) \times V_{\ell_0+1}$  independently of  $n$  and one can extract a subsequence that we still label by  $n$  and denote by  $(u_{n,1}, v_{n,1})$  such that

$$u_{n,1} \rightharpoonup u^1 \text{ in } H^1(I_{\ell_0+1}), \quad v_{n,1} \rightharpoonup v^1 \text{ in } V_{\ell_0+1}, \quad u_{n,1} \rightarrow u^1 \text{ in } L^2(I_{\ell_0+1}), \quad v_{n,1} \rightarrow v^1 \text{ in } L^2(\Omega_{\ell_0+1})$$

$$v_{n,1}(\cdot, 0) \rightarrow v^1(\cdot, 0) \text{ in } L^2(I_{\ell_0+1}).$$

Note that  $(u^1, v^1) = (u^0, v^0)$  on  $I_{\ell_0} \times \Omega_{\ell_0}$ . Clearly  $(u^1, v^1)$  satisfies

$$\int_{\Omega_{\ell_0+1}} D \nabla v^1 \cdot \nabla \varphi \, dx + \int_{I_{\ell_0+1}} \nu v^1(x_1, 0) \varphi \, dx_1 = \int_{\Omega_{\ell_0+1}} f(v^1) \varphi \, dx + \int_{I_{\ell_0+1}} \mu u^1 \varphi \, dx_1 \quad \forall \varphi \in V_{\ell_0+1},$$

$$\int_{I_{\ell_0+1}} D' u^{1'} \psi' + \mu u^1 \psi \, dx_1 = \int_{I_{\ell_0+1}} \nu v^1(x_1, 0) \psi \, dx_1 + \int_{I_{\ell_0+1}} g(u^1) \psi \, dx_1 \quad \forall \psi \in H_0^1(I_{\ell_0+1}).$$

By induction one constructs a sequence  $(u_{n,k}, v_{n,k})$  extracted from the preceding converging toward  $(u^k, v^k)$  satisfying the equations above where we have replaced  $\ell_0 + 1$  by  $\ell_0 + k$ . Then using the usual diagonal process it is clear that the sequence  $(u_{n,n}, v_{n,n})$  will converge toward a solution to (5.1) nontrivial thanks to (5.3). This completes the proof of the theorem.  $\square$

**Acknowledgement** : Part of this work was performed when the first author was visiting USTC in Hefei. We thank the university and the school of mathematical sciences for their hospitality and their support.

## References

- [1] H. AMANN : Fixed point equations and nonlinear eigenvalue problems in ordered Banach spaces. SIAM review, 18, 620-709, 1976.
- [2] H. BERESTYCKI, J-M. ROQUEJOFFRE, L. ROSSI : The influence of a line with fast diffusion on Fisher-KPP propagation. J. Math. Biol., 66 (2013), 743-766.
- [3] H. BERESTYCKI, J-M. ROQUEJOFFRE, L. ROSSI : Travelling waves, spreading and extinction for Fisher-KPP propagation driven by a line with fast diffusion. Nonlinear Analysis, 137, 2016, 171-189

- [4] P. G. CIARLET: *Linear and Nonlinear Functional Analysis with Applications*. SIAM, Philadelphia, 2013.
- [5] M. CHIPOT: *Elliptic Equations: An Introductory Course*. Birkhäuser, Basel, Birkhäuser Advanced Texts, 2009.
- [6] M. CHIPOT: *Asymptotic Issues for Some Partial Differential Equations*. Imperial College Press, London, 2016.
- [7] R. DAUTRAY, J.-L. LIONS: *Analyse mathématique et calcul numérique pour les sciences et les techniques*. Tome 1, Masson, Paris, 1984.
- [8] L. C. EVANS: *Partial Differential Equations*, Volume 19 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, 1998.
- [9] D. GILBARG, N. S. TRUDINGER: *Elliptic Partial Differential Equations of Second Order*. Springer, Berlin, Heidelberg, New-York, Classics in Mathematics, Reprint of the 1998 edition, 2001.