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Radio asymmetry in 3C 99

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Abstract. The *N*-galaxy 3C 99, a Compact Steep-spectrum Source showing a triple asymmetric structure, has been observed with several arrays of radio telescopes at sub-arcsecond resolution. New images from MERLIN, European VLBI Network and the Very Long Baseline Array show that the source components detected in the central region of 3C 99 have steep spectral indices, making it difficult to determine which of them is the core. The asymmetric radio structure of 3C 99 is explained using the Scheuer-Baldwin continuous streaming model with a rather large difference in the interstellar medium density on the two sides of the central region. The age of 3C 99 has been estimated; the source is found younger than 10^6 years.

Key words. galaxies: individual: 3C99 – galaxies: general – radio continuum: galaxies

1. Introduction

1.1. Radio properties of 3C 99

The radio source 3C 99 (0358+004) is a powerful, $\sim 10^{28}$ Watt/Hertz at 408 MHz *N*-galaxy at a redshift 0.426. On the arcsecond scale, imaged with the VLA and MERLIN, it shows a triple asymmetric structure (Mantovani et al. 1990, 1997). The outer components are located rather asymmetrically relative to the nuclear region and show quite different luminosities and surface brightness. Along the major axis, 3C 99 has an angular size of ≈ 6 arcsec and a linear size of ≈ 21 kpc ($H_0 = 100 \text{ km s}^{-1} \text{ Mpc}^{-1}$; $q_0 = 0$) almost fitting the selection criteria to be classified as Compact Steep-spectrum Source (CSS; see Fanti et al. 1995). The VLA images (Mantovani et al. 1997) show that 3C 99 has a rather asymmetric structure, with 4.9% and 9.5% of polarized emission at 5 GHz and 8.4 GHz respectively. The polarized emission comes mainly from the bright hot spot ~ 0.8 arcsec North-East to the central component. The jet region (i.e. the region between these two components) is also slightly polarized, with the electric vector perpendicular to the jet axis. In the hot spot region, a depolarization ratio of 0.52 between 8.4 GHz and 5 GHz is seen. The position angle of the electric vectors is the same to within the errors at the two frequencies, although there are $n\pi$ ambiguities.

The source 3C 99 was observed with the European VLBI Network plus Green Bank and Haystack in 1984 at $\lambda 18$ cm (Mantovani et al. 1990), with an angular resolution

of 11×3 mas along a PA of 169° . The central component was resolved into several components of emission and shows a distorted structure. A compact feature was also detected in the hot spot, at ~ 0.8 arcsec from the central component.

1.2. IR, Optical and X-ray properties of 3C 99

The *N*-galaxy associated with 3C 99 (Spinrad et al. 1985) lies close to the central radio component labelled *C* in Mantovani et al. (1990), which is expected to contain the nucleus of the source. Observations with the IDS spectrograph of the Isaac Newton Telescope with the slit of the spectrograph aligned along the radio source major axis, have shown strong [OIII] emission and weak H_β . The [OIII] emission extends 1 arcsec above and below the nuclear region, with a higher redshift in the South-West direction, suggesting that thermal plasma in the North-East region is moving towards the observer, with $\Delta z \approx 0.0025$ (Mantovani et al. 1990). 3C 99 has been observed again in 1993 with the ESO 2.2 m telescope. The spectrum shows only narrow H_α and H_β lines. While most of the *N*-galaxies are actually classified BLRG, 3C 99 is classified as a NLRG (Hes et al. 1996). Of particular interest is the superposition of the [OII] $\lambda 3727$ distribution, from which the continuum has been subtracted (Hes, Ph.D. Thesis, University of Groningen, 1996), on the VLA A-array image at 4.9 GHz (Mantovani et al. 1990). The [OII] nebular emission is about 7 arcsec in extent, roughly centered on the nuclear radio region. The image shows that the north-east radio components are completely embedded in the gas, while the south-west components are mostly outside the emitting gas (see Fig. 1).

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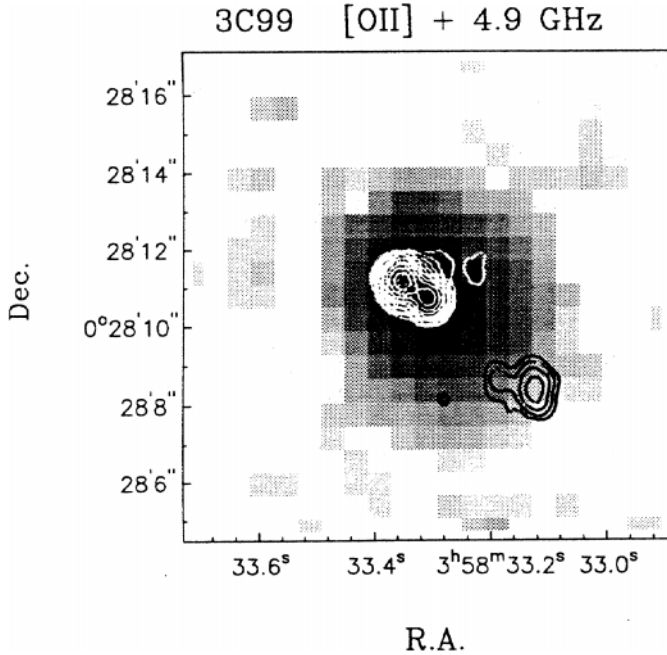


Fig. 1. The image of 3C 99 (contours) at 4.9 GHz (Mantovani et al. 1990) super imposed to the [OII] λ 3727 distribution (gray scale). Contours are at $0.3 \times (-1, 1, 2, 4, 8, 16, 32, 64, 128, 256, 512)$ mJy/beam. The peak flux density is 260 mJy. The beam is 0.48×0.42 arcsec² at PA 20°. The image is reproduced from Hes, Ph.D. Thesis, 1996).

Hes et al. (1996) also suggest that the emission lines are due to photoionization by the central source and that, if the B-model is adopted for photoionization, it is possible to infer from the H_{β} luminosity an electron density $n_e \leq 4 \times 10^2$ cm⁻³ with a filling factor $\approx 10^{-6}$. *N*-galaxies are occasionally much brighter at 25 μ m (van Bemmell & Barthel 2001). The high IR luminosity would classify 3C 99 as an unusually FIR-bright AGN. In addition certain CSS radio galaxies tend to be unusually FIR luminous (Hes et al. 1995), although this has not been found for the population as a whole (Fanti et al. 2000). On the other hand the X-ray luminosity is rather faint; it was not detected by ROSAT ($L_X < 10^{43}$ ergs⁻¹; Crawford & Fabian 1996). Together with the absence of broad lines, this seems to exclude a quasar classification. Table 1 summarizes the observed properties of 3C 99 at the various wavelengths.

1.3. The present paper

It has been suggested that the strong IR emission found in *N*-galaxies is mainly due to the interstellar medium (Hes et al. 1995), which in the case of CSSs sources might also be responsible for the radio plasma confinement. The asymmetry cannot be due to relativistic effects only, because the closer lobe is also the brighter one, and the enhanced emission in the northern lobe is also probably caused by a denser interstellar medium. We have carefully investigated the structure of 3C 99 with high resolution interferometric polarization observations. Polarization observations could be a sensitive probe of the effects of compression or shocks on the structure of the magnetic field.

Table 1. Summary of the observed properties at the various λ for 3C 99.

Optical Class	<i>N</i> -galaxy
Red-shift	0.426
Radio Class	CSS
Angular size	6 arcsec
Linear size	21 kpc
Polarization % $S_{5 \text{ GHz}}$	4.9
Polarization % $S_{8.4 \text{ GHz}}$	9.5
Radio Luminosity (408 MHz)	10^{28} Watts/Hz
IR Luminosity (60 μ m)	$10^{45.4}$ ergs/s
X Luminosity	$< 10^{43}$ ergs/s

In the following we will present observations done with the MERLIN and the European VLBI Network¹ at 5 GHz, and with the VLBA² at 1.6 GHz and at 5 GHz.

2. Observations and data analysis

The observations of 3C 99 are summarized in Table 2. The content of Table 2 is as follows: Col. 1, array name; Col. 2, observing stations; Col. 3, observing date; Col. 4, observing frequency; Cols. 5–7, beam major axis, minor axis and position angle respectively; Col. 8, total intensity rms noise; Col. 9, polarized intensity rms noise.

2.1. MERLIN observations

In 1995.38 the source was observed by the EVN+MERLIN array. In order to calibrate the MERLIN data set, initial values for the telescope and correlator gains were determined from a short observation of a bright, unresolved calibration source. The primary calibrator 3C 286 was included during the observations in order to fix the position angle of the electric vector. The absolute flux density was scaled assuming a flux density of 2325.6 mJy for OQ208. The source DA193 was used to calculate the leakage term. Images from the MERLIN LHC total intensity and polarization data, the latter without the Cambridge data due to the limits in bandwidth of the radio link, were produced using *AIPS* and shown in Figs. 2 and 3 respectively.

The images in Figs. 2 and 3 show that the core region (component *C*) and the North-East hot-spot (component *A*) are connected by a long, straight, collimated jet. The jet bends and slightly increases in transverse size before reaching the hot-spot region, which is resolved into a triple hot-spot. Component *C* is point-like at this resolution and is not polarized. The first knot along the jet is in fact polarized $\sim 8\%$, while in the remaining part of the jet, the polarization, if any, is below the detection

¹ The European VLBI Network is a joint facility of European and Chinese radio astronomy institutes funded by their national research Councils.

² The Very Long Baseline Array and the Very Large Array are facilities of the National Radio Astronomy Observatory, USA, operated by Associated Universities Inc., under cooperative agreement with the National Science Foundation.

Table 2. Observational parameters of 3C 99.

Array	Stations ^a	Date	Freq. MHz	Beam			σ_t mJy/beam	σ_p mJy/beam
				mas	mas	deg		
EVN	Ef, Mc, Nt, On-60, Jb-2, Wb	1990.87	4975	7.1	4.4	65	0.15	
EVN	Ef, Mc, Nt, On-60, Jb-2, Wb, Sm	1995.38	4975	8.4	4.5	37	0.16	
MERLIN	All	1995.38	4994	62	54	33	0.3	
MERLIN-Cm	All-Cm	1995.38	4994	100	100		0.5	0.5
VLBA	All-Sc+VLA1	1998.74	1655	17.5	4.9	-21	0.9	0.2
VLBA	All-Sc+VLA1	1998.74	4619	5.6	2.0	-18	0.2	0.3

^a The station labels listed are as follows: Ef = Effelsberg (D) 100-m, Mc = Medicina (I) 32-m, Nt = Noto (I) 32-m, On-60 = Onsala (S) 20-m, Jb-2 = Jodrell Bank (UK) 25-m, Cm = Cambridge (UK) 32-m, Wb = Westerbork array (NL) $n \times 25$ -m, Simeiz (Ukraine) 22-m; VLBA: Sc = St. Croix.

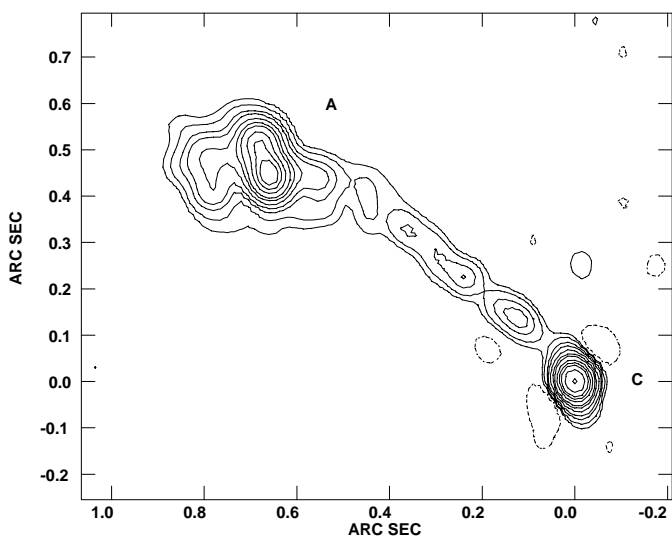


Fig. 2. The MERLIN (epoch 1995.38) image of 3C 99 at 5 GHz. Contours are at -1, 1, 2, 4, 6, 10, 15, 20, 30, 40, 50, 70, 100 mJy/beam. The peak flux density is 105.6 mJy. The beam is 62×54 mas at PA 33° .

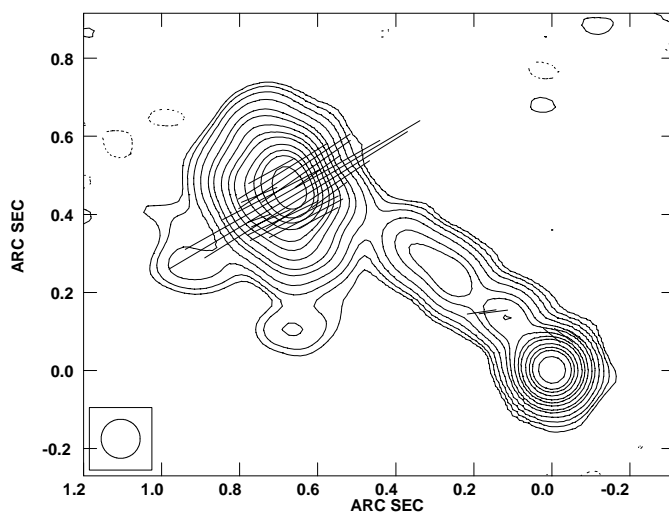


Fig. 3. The MERLIN minus Cambridge (epoch 1995.38) image of 3C 99 at 5 GHz. Contours are at -0.25, 0.25, 0.5, 1, 2, 4, 6, 10, 15, 20, 30, 40, 50, 70, 100 mJy/beam. The peak flux density is 96.4 mJy. The beam is 100×100 mas. An electric field vector length $1'' = 20$ mJy/beam.

limit. The region A shows $\sim 6.4\%$ polarized flux density in total. In the south-west region, the weak and extended emission seen with the VLA observations (Mantovani et al. 1990, 1997 and Fig. 1), was not detected by MERLIN.

Mantovani et al. (1997) using VLA observations measured a strong depolarization (0.52) between 8.4 GHz and 5 GHz and no indication of Faraday Rotation of the electric vector. Comparing the MERLIN image with the VLA images a similar percentage of polarized flux density is found at 5 GHz for the component A, confirming the depolarization between 8.4 GHz and 5 GHz. In the MERLIN image, the mean electric vector position angles around the peak of polarized emission are 120 ± 1 deg in agreement with previous results.

2.2. VLBA observations

The source was observed with 9 VLBA antennas (all but Saint Croix) and 1 VLA antenna for about 10 hours, switching between 1.6 GHz and 5 GHz. We made use of four 8 MHz bands at each frequencies centered at 1651.49, 1659.49, 1667.49 and 1675.49 and 4615.49, 4653.49, 4850.49, 5090.49 respectively. The target source was observed together with the

calibrator sources BL Lac, 3C 84, 3C 138 and OJ 287. The data have been analyzed using *AIPS*. The data have been calibrated in amplitude and phase, the LHC-RHC delay difference and the instrumental polarization have been corrected by fringe fitting a segment of the cross-hand data from a strong calibrator (3C 84).

The two data sets were then averaged in frequency over the four IF channels for a total observing bandwidth of 32 MHz. The images obtained at 1.6 GHz are shown in Figs. 4 and 5. The latter figure is a closeup of the central region of 3C 99. The hot-spot region has been resolved into three components which appear resolved and off-axis by $\sim 15^\circ$ with respect to the central region, as also apparent in the MERLIN map.

The central region shows two bright and two weak outer components. At 5 GHz we detected the central component only, which is dominated by the two bright components seen at 1.6 GHz. A similar image is seen at 4.6 GHz in Fig. 6 where they appear extended due to the better resolution. Again the structure is dominated by two bright blobs. It is worth noting that none of the components of Fig. 5 is point-like.

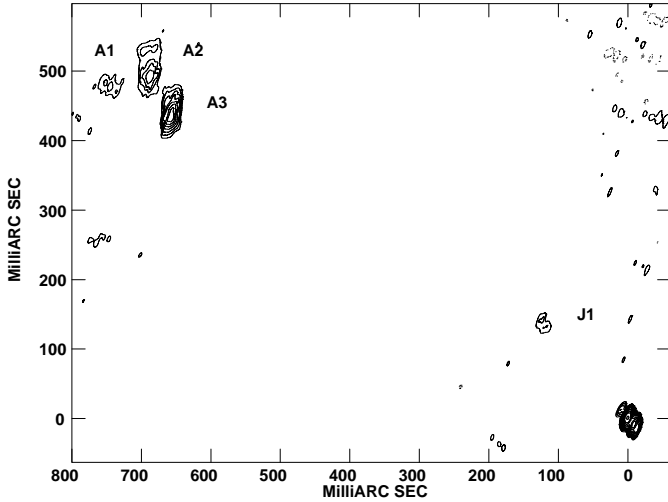


Fig. 4. The VLBA (epoch 1998.74) image of 3C 99 at 1.6 GHz. Contours are at $-2, 2, 3, 4, 5, 6, 7, 8, 10, 15, 20, 30, 40, 60$ mJy/beam. The peak flux density is 63.6 mJy. The beam is 17.5×4.9 mas at PA -21° .

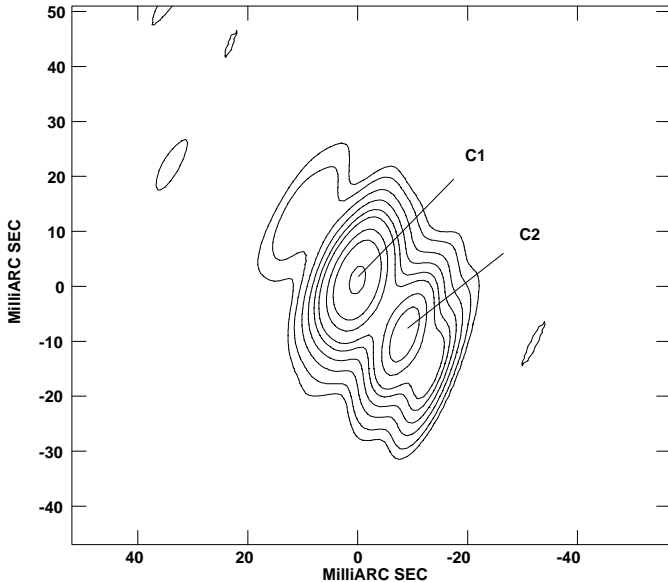


Fig. 5. The VLBA (epoch 1998.74) image of the central region of 3C 99 at 1.6 GHz. Contours are at $-1.5, 1.5, 3, 6, 10, 15, 20, 30, 40, 60$ mJy/beam. The peak flux density is 63.6 mJy. The beam is 17.5×4.9 mas at PA -21° .

No polarized emission was detected, at the two frequencies, 1.6 GHz and 5 GHz above the rms noise of 0.2 mJy/beam and 0.3 mJy/beam respectively.

2.3. EVN observations

The two EVN observations at 5 GHz were done recording in MarkIII Mode A (56 MHz bandwidth; date 1990.87) and recording in MarkIII Mode B (28 MHz bandwidth; date 1995.38). The data were processed at the Max-Planck-Institut für Radioastronomie Correlator in Bonn. The correlator output was calibrated in amplitude and phase using *AIPS*³ and

³ *AIPS* is the NRAO's *Astronomical Image Processing System*.

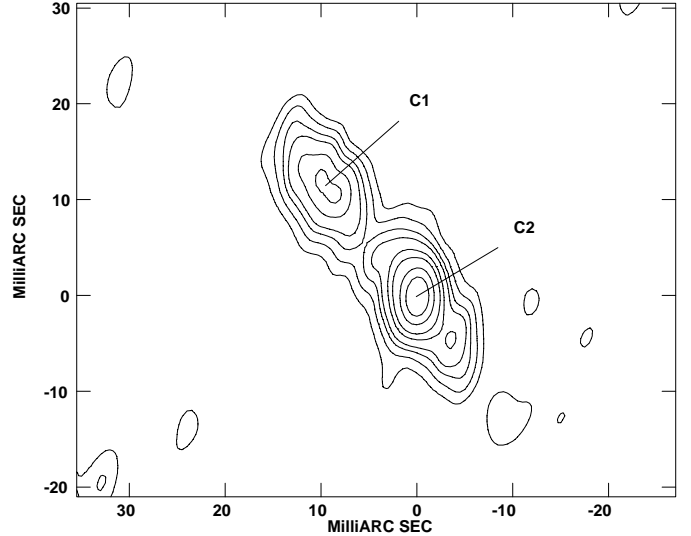


Fig. 6. The VLBA (epoch 1998.74) image of 3C 99 at 5 GHz. Contours are at $-0.6, 0.6, 1, 1.5, 2, 3, 4, 6, 8$ mJy/beam. The peak flux density is 10.5 mJy. The beam is 5.6×2.0 mas at PA -18° .

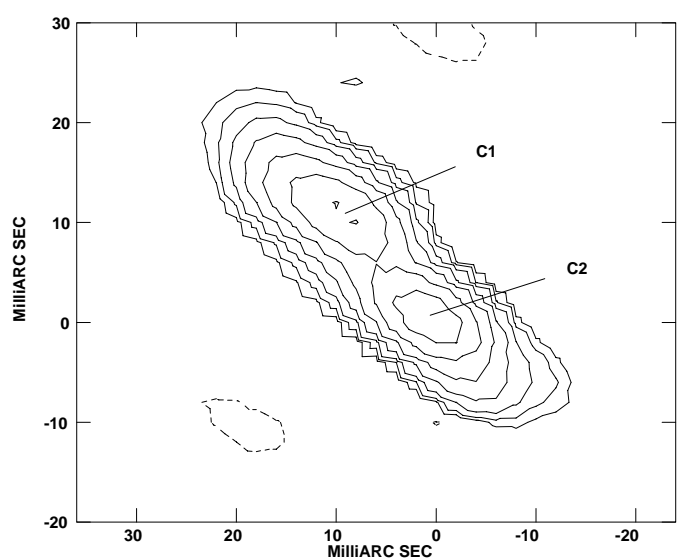


Fig. 7. The EVN (epoch 1990.87) image of 3C 99 at 5 GHz. Contours are at $-0.4, 0.4, 0.8, 1.6, 3.2, 6.4, 12.8, 25.6, 51.2$ mJy/beam. The peak flux density is 43.3 mJy. The beam is 7.1×4.4 mas at PA 65° .

imaged using DIFMAP⁴ (Shepherd et al. 1995). Total power measurements and gain-curve of the telescopes, were applied in the amplitude calibration process. The self-calibration procedure, which uses closure amplitudes to determine telescope amplitude corrections, gave calibration factors that were within 10% of unity for all the telescopes.

The two images from the two observing sessions are shown in Figs. 7 and 8 respectively. They look very similar in structure, size and flux density despite almost five years between the two observations. We see two resolved, elongated components, aligned with the source major axis in PA $\sim 40^\circ$, of almost equal flux density. The properties of 3C 99, from the EVN and the VLBA observations are given in Table 3.

⁴ DIFMAP is part of the *Caltech VLBI software Package*.

Table 3. Observed properties of the compact components of 3C 99.

Array	Date	Freq. MHz	Beam			Comp.	S_{peak} mJy/b	S_{tot} mJy	Major mas	Minor mas	PA deg
			mas	mas	deg						
EVN	1990.87	4975	7.1	4.4	65	C1	27.2	44.4	7.5	1.4	52
						C2	43.6	53.1	4.9	0.7	46
EVN	1995.38	4975	8.4	4.5	37	C1	26.2	46.6	10.1	2.4	44
						C2	43.1	55.9	–	–	–
VLBA	1998.74	1655	17.5	4.9	–21	A1	3.8	20.4	49.1	24.0	54
						A2	4.6	84.3	95.5	29.7	178
						A3	8.5	98.8	58.5	23.0	169
						J1	3.3	18.2	49.1	17.1	34
						C1	62.0	138.5	8.7	5.3	51
VLBA	1998.74	4619	5.6	2.0	–18	C1	4.3	23.4	9.8	3.8	40
						C2	10.0	26.7	5.3	2.4	32
						C1	13.1	20.0	–	–	–
									–	–	–
									–	–	–
C2	22.4	34.9	–	–	–						

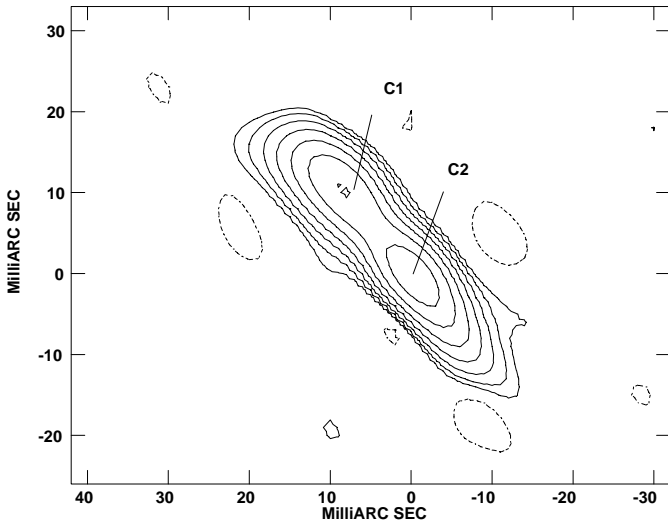


Fig. 8. The EVN (epoch 1995.38) image of 3C 99 at 5 GHz. Contours are at –0.4, 0.4, 0.8, 1.6, 3.2, 6.4, 12.8, 25.6, 51.2 mJy/beam. The peak flux density is 43.7 mJy. The beam is 8.0×4.5 mas at PA 62° .

2.4. Comparison of the VLBI images of the central component

The 5 GHz image has been convolved with the 1.6 GHz beam and the flux densities of the two main components compared. A spectral index of $\alpha_1 = 0.8$ is derived for component C1 and $\alpha_1 = 1.8$ for component C2 ($S \propto \nu^{-\alpha}$).

To check for any possible change of the relative position of the two components C1 and C2 with epoch, we made use of the images at 5 GHz (their maximum time separation is ~ 8 years), all convolved to the same beam size. The angular separation between the two components of ~ 14 mas did not change within the estimated errors of ≈ 0.5 mas. The errors are affected by the discrete size of the components (of the order of the beam size) which are extended in the same direction as the source major axis.

The derived parameters, resolved angular sizes, steep spectral indices, constant separation, suggest that nor C1 neither C2 can be considered as the core of 3C 99.

3. Discussion

The arcsecond resolution images of 3C 99 show that the source has a triple asymmetric structure. Even if we could not achieve a firm identification of the source core, the core location is doubtlessly in the component C region. Also, the position of the optical N-galaxy 3C 99 (Spinrad et al. 1985) is close to that component. Thus, we know approximately where the central engine is located and we can confirm the source asymmetry. The ratio between the distances of the hot-spots from the center of component C is 4.5, while that of the hot-spots brightnesses is ≈ 40 . From the optical observations we also know that the North-East side is approaching the observer being less redshifted than the South-West part (see Mantovani et al. 1990).

Asymmetry in the radio lobe arm length ratio is expected when the hot spots move at a relativistic speed as a consequence of the difference in travel time of the radiation emitted by the approaching and receding hot spot. The approaching lobe should be the one seen at a later stage of evolution and therefore more distant from the core as compared to the receding one. Furthermore some relativistic boosting (de-boosting) should also be expected for the approaching (receding) hot spot, giving rise as well to an asymmetry in luminosity. In 3C 99 the lobe closer to the core is the brighter and approaching one. It is clear that relativistic effects alone cannot fully explain the observed asymmetry.

Alternatively, we attribute the asymmetry to an inhomogeneity in the ambient medium on the two sides of the source. The brighter and closer lobe would be the one moving through a more dense medium having a small density gradient, being slowed down and remaining bright as a consequence of a more effective confinement. An indication that this could be the case comes from the image of Fig. 1, a superposition of the radio and [OII] images. If we assume that they actually overlap in

space, it follows that the north–east bright component of the source, i.e. the part closest to the core, is completely immersed in the cloud of gas, while the south–west hot–spot lies in a region at the border where the gas density is clearly much lower.

3.1. A model for asymmetry

To explain the asymmetry in the radio structure of 3C 99 we made use of the Scheuer-Baldwin model (Scheuer 1974; Baldwin 1982), one of the so called continuous streaming model. In the model, a relativistic twin-jet propagates in an external medium with a non uniform density.

We assume that the jets carry an energy flux F_e and a thrusts F_e/c . The advance speed of the jets head (hot spots), v_h , is controlled by the balance between the jet thrust and the ram-pressure of the ambient medium, viz. $v_h \propto \frac{F_e}{c} D^{-2}$. Modeling the external density as

$$\rho \propto D^{-n}$$

where D is the distance from the nucleus. By integration of these relations, we get the instantaneous and average advance speed

$$v_h \propto D^{-(2-n)/2} \quad \langle v_h \rangle = v_h \times (4 - n)/2.$$

The energy carried by the jet is assumed to accumulate into the lobe under equipartition conditions. This allows the computation of the lobe luminosity (Begelman 1996) as

$$L_l \propto [F_e \times \tau \times f(n)]^{7/4} \times V_l^{-3/4}$$

where τ is the lobe age, $f(n)$ a factor accounting for the energy lost in the source expansion (of the order of 0.5) and V_l the lobe volume.

Finally, assuming again equipartition, we evaluate the hot spot luminosity as

$$L_h \propto F_e^{7/4} \times r_h^{-0.5}$$

where r_h is the hot spot radius, that we assume, for sake of simplicity, proportional to D .

We remark that, because of the large advance speeds, the travel time effects would cause the two lobes to be seen at different ages, thus

$$\tau_1/\tau_2 = (1 + \langle \beta_2 \rangle \cos \theta) / (1 - \langle \beta_1 \rangle \cos \theta).$$

Therefore, an important item for the following discussion is the orientation of the source with respect to the line of sight.

3.2. Physical parameters from jets, lobes and hot spots

In the above framework, we will interpret the asymmetry in 3C 99 making use of the physical quantities, listed in Table 4, derived from the present observations and from those presented in Mantovani et al. (1990).

The ratio of the distances of the lobes at the two opposite sides of the nucleus is:

$$\frac{D_1}{D_2} = \frac{\langle \beta_1 \rangle}{\langle \beta_2 \rangle} \times \frac{1 + \langle \beta_{2h} \rangle \cos \theta}{1 - \langle \beta_{1h} \rangle \cos \theta} \quad (1)$$

where D_1 and D_2 are the distances of the approaching and receding lobe respectively, $\langle \beta_1 \rangle$ and $\langle \beta_2 \rangle$ their mean velocities relative to the speed of light and θ the angle between the source major axis and the observer's line of sight. The second factor accounts for differences in the travel time from the lobes.

The lobe luminosity ratio is given by:

$$\frac{L_{l1}}{L_{l2}} = \left[\frac{F_{e1}}{F_{e2}} \times \frac{1 + \langle \beta_2 \rangle \cos \theta f(n_1)}{1 - \langle \beta_1 \rangle \cos \theta f(n_2)} \right]^{7/4} \times \left(\frac{V_{l2}}{V_{l1}} \right)^{3/4}. \quad (2)$$

The factor $(1 + \langle \beta_2 \rangle \cos \theta)/(1 - \langle \beta_1 \rangle \cos \theta)$ accounts for the fact that the two lobes are seen at different times of their evolution. A similar equation holds for the hot spot luminosity ratio:

$$\frac{L_{hs1}}{L_{hs2}} = \left(\frac{D_2}{D_1} \right)^{1/2} \left(\frac{F_{e1}}{F_{e2}} \right)^{7/4} \left[\frac{\gamma_2(1 + \beta_2 \cos \theta)}{\gamma_1(1 - \beta_1 \cos \theta)} \right]^{3 + \alpha_h} \quad (3)$$

where $\gamma = (\sqrt{1 - \beta^2})^{-1}$ is the Lorentz factor, α_h the hot spot spectral index. The parameters needed to satisfy the above equations can be compared to those derived from the existing observations of 3C 99, summarized in Tables 4 and 5.

3.3. Search for the unknown parameters

We make use of the above equations to evaluate the unknown parameters β_1 , β_2 , n_1 , n_2 and θ .

We first constrain the range of values for θ . Assuming a symmetric two-sided jet, the lower limit to the jet to counter-jet brightness ratio implies $\beta_j \cos \theta \simeq 0.63$, which requires $\theta \leq 50^\circ$. Moreover, in order to fit the unified models scheme (see Saikia et al. 1995) a CSS source optically classified as a N -galaxy is expected to have the radio axis at $30^\circ < \theta < 50^\circ$ to the observers line of sight.

In order to maximize the effects due to the asymmetric distribution of the ambient medium, we assume $n_1 = 0$ and $n_2 = 2$. Following the models of Scheuer (1974) modified by Baldwin (1982) and of Begelman (1996), we compute $f(n_1)/f(n_2) \approx 1.25$. Furthermore $\beta_2 = \langle \beta_2 \rangle$ and $\beta_1 = 0.5 \langle \beta_1 \rangle$.

The observed parameters which the model uses are:

$$D_1/D_2 = 0.21 \pm 0.02$$

$$L_{l1}/L_{l2} = 23 \pm 4$$

$$V_{l1}/V_{l2} = 25 \pm 8$$

$$L_{hs1}/L_{hs2} = 35 \pm 10$$

$$r_{h1}/r_{h2} = 8 \pm 2.$$

One of the most uncertain input parameter is the ratio of the lobe volumes, deduced from the 1.7 GHz map of Mantovani et al. (1990), which is the best in order to evaluate the lobe's luminosity, but suffers from poor angular resolution for a proper measurement of the lobe's size. On the other hand, the above mentioned source models (Scheuer-Baldwin and Begelman) allow the computation of the volume ratio, which is consistent with the observed one.

Table 4. Physical quantities from the images of 3C 99.

Region	Comp.	$\log P_{1.4 \text{ GHz}}$	$AS_1 \times AS_2$	$LS_1 \times LS_2$	Volume	U_{\min}	u_{\min}	H_{eq}
			mas	kpc	kpc^3	erg	erg cm^{-3}	mGauss
Lobe _{North}		26.4	750×470	4.1×2.8	17	1.8×10^{57}	3.5×10^{-9}	0.2
Lobe _{South}		25.1	1500×1500	5.9×5.9	111	0.8×10^{57}	2.4×10^{-10}	0.05
Hot-spot _{North}	I	25.5	52×42	0.32×0.25	0.01	2.2×10^{55}	7.0×10^{-8}	0.9
	II	25.7	48×42	0.29×0.23	0.01	2.6×10^{55}	9.6×10^{-8}	1.0
	III	25.2	148×87	0.90×0.50	0.12	4.2×10^{55}	1.1×10^{-8}	0.35
Hot-spot _{South}	I	24.2	531×290	3.15×1.72	4.9	5.0×10^{55}	3.6×10^{-10}	0.06
	II	23.9	300×231	1.0×0.76	1.75	2.2×10^{55}	4.4×10^{-10}	0.07
Jet (mas)		25.3	700×70	4.14×0.21	0.09	4.5×10^{55}	1.6×10^{-8}	0.4
5 GHz	C1	25.0	10.4×4.2	0.05×0.02	1.8×10^{-5}	8.0×10^{53}	1.5×10^{-6}	4.1
1.6 GHz	C1	25.4	8.6×3.1	0.03×0.01	1.8×10^{-5}	1.4×10^{54}	2.8×10^{-6}	5.5
5 GHz	C2	25.3	10.5×3.6	0.05×0.02	0.9×10^{-5}	8.6×10^{53}	3.3×10^{-6}	6.0
1.6 GHz	C2	25.3	29.4×4.3	0.18×0.01	5.0×10^{-5}	1.8×10^{54}	1.2×10^{-6}	3.6

The quantities have been derived assuming equipartition, a filling factor of 1, electron to proton ratio of 1 and the behavior of the radio spectrum between 10 MHz and 100 GHz. To determine the flux densities of the north-east lobe, we subtracted the core, jet and hot spot flux densities from images at higher resolutions (for example the MERLIN image in Fig. 2). Since those images were also at different frequencies, the flux densities were scaled accordingly, adopting a mean spectral index 0.5 ($S \propto \nu^{-\alpha}$). The linear dimensions were derived from the angular dimensions obtained using the AIPS task JMFIT, deconvolved from the beam. Table 4 is organized as follows: Col. 1, region; Col. 2, component (the hot-spot north components I, II and III correspond to the components A1, A2 and A3 of Fig. 4); Col. 3, radio power at 1.4 GHz; Col. 4, linear dimensions in kpc; Col. 5, volume in kpc^3 ; Col. 6, minimum energy in erg; Col. 7, energy density erg/kpc^3 ; Col. 8, magnetic field in mGauss.

Table 5. Physical parameters of 3C 99.

Parameter	Image Freq.	Approaching jet	Receding jet	Ratio
Distance nucleus-lobe	1.4 GHz	$D_1 = 1260 \text{ mas}$	$D_2 = 5940 \text{ mas}$	0.21
Lobe Flux Density	0.4 GHz	$S_{11} = 2900 \text{ mJy}$	$S_{12} = 170 \text{ mJy}$	17.1
Hot-spots Flux Density	5 GHz	$S_{\text{hs1}} = 159 \text{ mJy}$	$S_{\text{hs2}} = 4 \text{ mJy}$	37.9
Jet Brightness		$b_j \approx 4 \text{ mJy/beam}$	$b_{\text{cj}} \leq 0.32 \text{ mJy/beam}$	> 10

The jet counter-jet brightness ratio in the last row of Table 5 has been derived from the MERLIN image at 5 GHz considering the jet as having a conical structure with longitudinal dimension of $\sim 700 \text{ mas}$, a transverse dimension of $\sim 70 \text{ mas}$, an angle $\leq 1^\circ$ and a total flux density of 43 mJy. The jet mean brightness $b_j \approx 4 \text{ mJy/beam}$ is obtained dividing the total flux density by the beam area to jet area ratio. Since the counter-jet is below the detection limit of the MERLIN image, we have $b_{\text{cj}} \leq 3\sigma = 0.32 \text{ mJy/beam}$, and a ratio between the brightnesses ≥ 12 .

An additional uncertainty comes from the identification of the *hotspot* of the north lobe. If we take the brighter one (named I in Table 4, labelled A1 in Fig. 4) we get a luminosity ratio of 30, while using the second brightest (named II in Table 4, labelled A2 in Fig. 4) we get a ratio of 20. The previous equations allow a search over a range of parameters to account for the observed asymmetry of the radio source. For $30^\circ \leq \theta \leq 40^\circ$, the first two equations require $\langle \beta_2 \rangle \approx 0.3 \pm 0.1$ and $\langle \beta_2 \rangle / \langle \beta_1 \rangle \approx 5.5$.

The above numbers allow the computation of the ratio of hot spot luminosities, with Eq. (3). We get: $10 < L_{\text{h1}}/L_{\text{h2}} < 15$. This range of values is somewhat lower than the one we have from the observations (≈ 35). In order to increase the $L_{\text{h1}}/L_{\text{h2}}$ computed ratio, one should increase the β s, which in turns implies an increase of the ratio L_{11}/L_{12} .

3.4. The age of 3C 99

The radio source age can be derived from the age of the emitting relativistic electrons providing we know the break

frequency in the spectral index of the source. From the existing images, we derive straight steep spectral indices for the two lobes ($\alpha \sim 1.2$), thus we cannot infer any value for break frequencies. The only suggestion we can make is that $\nu_{\text{break}} \leq 408 \text{ MHz}$. For the equipartition magnetic fields of Table 4, following Carilli et al. (1991) and Murgia et al. (1999) we have $\tau_{\text{N}} \leq 7.5 \times 10^5 \text{ years}$ for the North lobe, which accounts for most of the total radio luminosity.

A second way to determine the age of 3C 99 is a dynamical approach via the following relation:

$$\langle \tau \rangle = \frac{D}{c \langle \beta \rangle \sin \theta} \quad (4)$$

and for example for $\langle \beta_1 \rangle = 0.05$, $\langle \beta_2 \rangle = 0.3$ and $\theta = 30^\circ$ dynamical ages of 7×10^5 and 4×10^5 for each lobe respectively, are derived. The two ages differ due to the time delay in the propagation of the radiation.

4. Conclusions

The present observations confirm the asymmetry of 3C 99. A one-sided long, thin jet was detected with MERLIN. The 5-GHz MERLIN observations show that the source is polarized in the hot spot region and confirm the depolarization previously measured, and a PA of the electric vector in agreement to previous results. The hot spot region contains three compact features detected by the 1.6 GHz VLBA observations, the jet was resolved out and the central area shows a double structure in both 1.6 GHz and 5 GHz VLBA images. The observations did not allow us to detect polarized emission from any components in 3C 99 above the detection limit of the VLBA observations (0.3 mJy/beam at 5 GHz). The comparison between the 1.6 GHz and the 5 GHz shows that the detected source components do show a steep spectral index, making it difficult to suggest which of them is the core. We suggest that the two components detected in the central region of 3C 99, are co-moving jet components along the north–east side of the source.

The asymmetry in the radio structure of 3C 99 can be explained with the Scheuer-Baldwin model requiring a rather large difference in the interstellar medium density on the two sides of the central component. Such an assumption makes the model predictions consistent with the observational parameters. We have assumed that in the south–west direction, where the hot-spot is less bright and more distant, the ambient gas density goes as r^{-2} , where r is the radial distance from the source nucleus. On the opposite side, i.e. in north–east direction, we have assumed that the ambient gas density is constant with r , with a stronger interaction with the advancing jet. This should also explain the strong far infrared emission detected. The interaction of the radio plasma with the interstellar medium may trigger the star formation and the heating of dust will generate strong infrared emission. However, relativistic effects were also required to fit the parameters derived from the observations. An angle of 30° – 40° between the source major axis and the observers line of sight is needed.

The age of 3C 99 has been derived using both the age of the emitting relativistic electrons with a dynamical approach. The source is found younger than 10^6 years.

A time scale of $\sim 10^5$ – 10^6 has been suggested recently by Baker et al. (2002) for the ionization cone to emerge from their dusty cocoons in quasars which is comparable to those of many CSS and, of course, to that of 3C 99. They show that the stronger C IV absorbers are found in small steep-spectrum radio sources and argue that the absorption column density changes with time. Since it is unlikely that CSS sources are confined by dense clouds of gas and since the environmental density cannot account for the full range of radio source size in quasars, they suggest that the age must be important. Models of dust grains in the vicinity of powerful radio sources

requires special conditions for grain survival up to $\sim 10^7$ years. According to such interpretation, 3C 99 can be understood to be in a phase in which the jets are emerging from the clouds of dense gas.

We note that, though 3C 99 has been classified as a *N*-galaxy, an overview of the spectroscopy available at several wavelength rises new puzzling questions about its nature in the light of the unified models. Its high infrared luminosity, which is a factor of 20 greater than that of a typical radio galaxy, and its optical spectrum which shows high luminosity narrow lines, would classify 3C 99 as FIR-bright AGN. The lack of broad band optical emission and the rather weak X-ray emission, on the other hand, are against such interpretation. New observations are needed to confirm the alternative suggestion for the star formation process as origin of the infrared emission.

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