



HAL
open science

Bistable pulsating fronts in slowly oscillating environments *

Weiwei Ding, François Hamel, Xing Liang

► **To cite this version:**

Weiwei Ding, François Hamel, Xing Liang. Bistable pulsating fronts in slowly oscillating environments
*. 2022. hal-03642527

HAL Id: hal-03642527

<https://hal.science/hal-03642527>

Preprint submitted on 15 Apr 2022

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Bistable pulsating fronts in slowly oscillating environments*

Weiwei Ding[†] François Hamel[‡] Xing Liang[§]

Abstract

We consider reaction-diffusion fronts in spatially periodic bistable media with large periods. Whereas the homogenization regime associated with small periods had been well studied for bistable or Fisher-KPP reactions and, in the latter case, a formula for the limit minimal speeds of fronts in media with large periods had also been obtained thanks to the linear formulation of these minimal speeds and their monotonicity with respect to the period, the main remaining open question is concerned with fronts in bistable environments with large periods. In bistable media the unique front speeds are not linearly determined and are not monotone with respect to the spatial period in general, making the analysis of the limit of large periods more intricate. We show in this paper the existence of and an explicit formula for the limit of bistable front speeds as the spatial period goes to infinity. We also prove that the front profiles converge to a family of front profiles associated with spatially homogeneous equations. The main results are based on uniform estimates on the spatial width of the fronts, which themselves use zero number properties and intersection arguments.

1 Introduction and main results

In this paper, we study the following reaction-diffusion equation

$$u_t = (a_L(x)u_x)_x + f_L(x, u), \quad t \in \mathbb{R}, \quad x \in \mathbb{R}, \quad (1.1)$$

with $L > 0$, where $u_t = \partial_t u = \frac{\partial u}{\partial t}$ and $(a_L(x)u_x)_x = \partial_x(a_L(x)\partial_x u)$ and $\partial_x = \frac{\partial}{\partial x}$. The diffusion and reaction coefficients a_L and f_L are given by

$$a_L(x) = a\left(\frac{x}{L}\right) \quad \text{and} \quad f_L(x, u) = f\left(\frac{x}{L}, u\right),$$

where the function $a : \mathbb{R} \rightarrow \mathbb{R}$ is positive, of class $C^2(\mathbb{R})$, 1-periodic, that is, $a(x+1) = a(x)$ for all $x \in \mathbb{R}$, and the function $f : \mathbb{R} \times [0, 1] \rightarrow \mathbb{R}$, $(x, u) \mapsto f(x, u)$ is of class C^1 , 1-periodic in x , and $(x, u) \mapsto \partial_u f(x, u)$ is Lipschitz-continuous in $u \in [0, 1]$ uniformly with respect to $x \in \mathbb{R}$. We also assume the following conditions on f :

*This work has received funding from Excellence Initiative of Aix-Marseille Université - A*MIDEX, a French “Investissements d’Avenir” programme, and from the French ANR RESISTE (ANR-18-CE45-0019) project. The first author is partly supported by NSFC (12001206). The second author acknowledges support of the Institut Henri Poincaré (UAR 839 CNRS-Sorbonne Université), and LabEx CARMIN (ANR-10-LABX-59-01). The third author is partially supported by the NSFC (11971454).

[†]School of Mathematical Sciences, South China Normal University, Guangzhou 510631, China (dingweiwei@m.scnu.edu.cn).

[‡]Aix-Marseille Univ, CNRS, I2M, Marseille, France (francois.hamel@univ-amu.fr).

[§]School of Mathematical Sciences and Wu Wen-Tsun Key Laboratory of Mathematics, University of Science and Technology of China, Hefei, Anhui, 230026, China (xliang@ustc.edu.cn).

(A1) there is a continuous function $b : \mathbb{R} \rightarrow (0, 1)$, $x \mapsto b(x)$, such that for every $x \in \mathbb{R}$, the profile $f(x, \cdot)$ is of the bistable type in the following sense

$$f(x, 0) = f(x, 1) = f(x, b(x)) = 0, \quad f(x, \cdot) < 0 \text{ on } (0, b(x)), \quad f(x, \cdot) > 0 \text{ on } (b(x), 1);$$

(A2) 0 and 1 are uniformly (in x) stable zeroes of $f(x, \cdot)$, in the sense that there exist $\gamma_0 > 0$ and $\delta_0 \in (0, 1/2)$ such that

$$\begin{cases} f(x, u) \leq -\gamma_0 u & \text{for all } (x, u) \in \mathbb{R} \times [0, \delta_0], \\ f(x, u) \geq \gamma_0(1 - u) & \text{for all } (x, u) \in \mathbb{R} \times [1 - \delta_0, 1], \end{cases} \quad (1.2)$$

and $\partial_u f(\cdot, 0)$ and $\partial_u f(\cdot, 1)$ are assumed to be of class C^1 in \mathbb{R} .

Notice that the assumptions (A1)-(A2) imply in particular that $\max(\partial_u f(x, 0), \partial_u f(x, 1)) \leq -\gamma_0$ for all $x \in \mathbb{R}$. Hence, 0 and 1 are two linearly stable L -periodic steady states of (1.1). Here, an L -periodic steady state \bar{u} is said to be linearly stable (resp. linearly unstable) if $\lambda_1(\bar{u}) < 0$ (resp. $\lambda_1(\bar{u}) > 0$), where $\lambda_1(\bar{u})$ is the principal eigenvalue of the operator $\varphi \mapsto \mathcal{L}\varphi := \partial_x(a_L(x)\partial_x\varphi) + \partial_u f_L(x, \bar{u}(x))\varphi$ in the space of L -periodic functions $\varphi \in C^2(\mathbb{R})$.

In the paper, for mathematical convenience, we extend f in $\mathbb{R} \times (\mathbb{R} \setminus [0, 1])$ as follows:

$$\begin{cases} f(x, u) = \partial_u f(x, 0)u & \text{for } (x, u) \in \mathbb{R} \times (-\infty, 0), \\ f(x, u) = \partial_u f(x, 1)(u - 1) & \text{for } (x, u) \in \mathbb{R} \times (1, +\infty). \end{cases} \quad (1.3)$$

Thus, it is clear that f is of class C^1 in \mathbb{R}^2 , 1-periodic in x , $\min_{x \in \mathbb{R}} f(x, u) > 0$ for all $u < 0$ and $\max_{x \in \mathbb{R}} f(x, u) < 0$ for all $u > 1$, and $f(x, u)$, $\partial_u f(x, u)$ are globally Lipschitz-continuous in u uniformly in $x \in \mathbb{R}$. It is also easily seen that this extension does not affect the behavior of the pulsating front connecting 0 and 1 defined below.

By a *pulsating front of (1.1) connecting 0 and 1*, we mean a classical entire solution $U_L : \mathbb{R} \times \mathbb{R} \rightarrow (0, 1)$ of (1.1) for which there exist a real number c_L and a function $\phi_L : \mathbb{R} \times \mathbb{R} \rightarrow (0, 1)$ satisfying

$$\begin{cases} U_L(t, x) = \phi_L\left(x - c_L t, \frac{x}{L}\right) & \text{for all } (t, x) \in \mathbb{R} \times \mathbb{R}, \\ \phi_L(\xi, y) \text{ is 1-periodic in } y, \\ \phi_L(-\infty, y) = 1, \quad \phi_L(+\infty, y) = 0 & \text{uniformly in } y \in \mathbb{R}. \end{cases} \quad (1.4)$$

The constant c_L is called the *front speed* and ϕ_L is the *front profile*. Clearly, if $c_L \neq 0$, then

$$U_L(t, x) = U_L\left(t + \frac{L}{c_L}, x + L\right) \quad \text{for all } (t, x) \in \mathbb{R} \times \mathbb{R}, \quad (1.5)$$

that is, the spatial profiles $x \mapsto U_L(t, x + c_L t)$ of the solution in the frame moving with speed c_L repeat periodically in time with period L/c_L . In the case $c_L = 0$, (1.4) simply means that U_L does not depend on t (that is, it is a steady solution) and $U_L(-\infty) = 1$, $U_L(+\infty) = 0$. *Throughout this paper, when we refer to a pulsating front for (1.1), we mean a solution $U_L : \mathbb{R} \times \mathbb{R} \rightarrow (0, 1)$ satisfying (1.4) or, equivalently, a pair (ϕ_L, c_L) given as in (1.4), with $\phi_L : \mathbb{R} \times \mathbb{R} \rightarrow (0, 1)$ and U_L solving (1.1).*

Fronts are of great importance in reaction-diffusion equations, which are the most-used equations in population dynamics models in biology and ecology, as fronts are the main mathematical notion for the description of the invasion of a state (here, the trivial state 0) by another one

(here, the constant state 1). The notion of pulsating front in spatially periodic media was first introduced in [51, 54]. It is a natural generalization of the classical notion of traveling front in spatially homogeneous media where the coefficients a and f are independent of x . In that case, the function $x \mapsto b(x)$ is a constant which is the only unstable zero of f in $(0, 1)$. For the homogeneous bistable equation

$$u_t = au_{xx} + f(u),$$

it is well known [4, 28] that there exist a unique speed $c \in \mathbb{R}$ and a unique traveling front $u(t, x) = \phi(x - ct)$ with $0 < \phi < 1$ in \mathbb{R} and $\phi(-\infty) = 1$, $\phi(+\infty) = 0$. The speed c has the sign of the integral $\int_0^1 f$, and the front profile ϕ is decreasing, unique up to shifts, and globally asymptotically stable, as in Theorem 1.1 (iii) below. The global stability of the traveling fronts reinforces their fundamental role in reaction-diffusion equations.

Media are however rarely homogeneous, and reaction-diffusion equations with spatially periodic coefficients, such as (1.1), are some of the most important classes of non-homogeneous equations. Their study has attracted much attention in the mathematical literature. One of the reasons for that interest is that the question of the existence of pulsating fronts is rather subtle. Indeed, the possible presence of multiple ordered L -periodic steady states may prevent the existence of pulsating fronts connecting the extremal steady states 0 and 1, see [16, 23, 31, 32]. On the other hand, it is known from [23, 27] that such a possibility does not happen provided that equation (1.1) admits a bistable structure in the sense that any L -periodic steady state strictly between 0 and 1 is linearly unstable (we recall that 0 and 1 are linearly stable steady states). With such a bistable structure, pulsating fronts have been proved to exist by using different approaches [22, 23, 27]. Furthermore, in an earlier paper by the first two authors [18], the bistable structure was verified under various explicit conditions on the functions $a(x)$ and $f(x, u)$ when the spatial period L is large or small. In particular, if in addition to (A1)-(A2) the function f is assumed to satisfy

$$(A3) \quad \int_0^1 f(x, u) du > 0 \quad \text{and} \quad \partial_u f(x, b(x)) > 0 \quad \text{for all } x \in \mathbb{R},^1$$

then a pulsating front with a positive speed c_L exists when L is large (i.e., the medium oscillates slowly). On the other hand, if the function $x \mapsto \int_0^1 f(x, u) du$ is not of a constant sign, it was shown in [17] that, when L is large, non-stationary pulsating fronts do not exist, while multiple stationary fronts do exist (see also [17, 34, 38, 55, 57] for further results on the existence of stationary fronts for bistable equations of the type (1.1)). We point out that, apart from the aforementioned works, the existence of pulsating fronts was investigated in [43, 46, 54, 56, 58]. References [43, 46, 58] were also concerned with generalized transition fronts in the sense of [9, 10], for which the media can be non-periodic (see [3, 14, 20, 24, 25, 41, 45, 47, 49, 52] for further propagation or blocking results in non-periodic bistable media).

For clarity, we collect some known results from [17, 18] concerning the existence and qualitative properties of pulsating front which will be frequently used in the present work.

Theorem 1.1. [17, 18] *Let (A1)-(A3) hold. Then there exists $L_* > 0$ such that, for any $L \geq L_*$, problem (1.1) admits a pulsating front $U_L(t, x) = \phi_L(x - c_L t, x/L)$ with speed c_L . Furthermore:*

- (i) $c_L > 0$, c_L is the unique front speed and the map $L \mapsto c_L$ is bounded in $[L_*, +\infty)$;
- (ii) the front U_L is unique up to shifts in time t , and $\partial_t U_L(t, x) > 0$ for all $(t, x) \in \mathbb{R}^2$;

¹Together with (A1) and the C^1 smoothness of f , condition (A3) implies that $b \in C^1(\mathbb{R})$.

(iii) the front U_L is globally asymptotic stable in the sense that, for any $u_0 \in C(\mathbb{R}) \cap L^\infty(\mathbb{R})$ with

$$\liminf_{x \rightarrow -\infty} u_0(x) > 1 - \delta_0 \quad \text{and} \quad \limsup_{x \rightarrow +\infty} u_0(x) < \delta_0, \quad (1.6)$$

where $\delta_0 > 0$ is the constant given in (A2), we have

$$\lim_{t \rightarrow +\infty} \|u(t, \cdot; u_0) - U_L(t + \tau_0, \cdot)\|_{L^\infty(\mathbb{R})} = \lim_{t \rightarrow +\infty} \|u(t, \cdot; u_0) - \phi_L(\cdot - c_L t - c_L \tau_0, \cdot/L)\|_{L^\infty(\mathbb{R})} = 0,$$

where $(t, x) \mapsto u(t, x; u_0)$ is the solution of the Cauchy problem for (1.1) with initial condition u_0 , and $\tau_0 \in \mathbb{R}$ is a constant depending on u_0 ;

(iv) the front speed c_L and the front profile ϕ_L are continuous with respect to $L \geq L_*$ in the sense that $c_L \rightarrow c_{L_0}$ as $L \rightarrow L_0$ in $[L_*, +\infty)$ and, up to translations of ϕ_L in such a way that $\phi_L(0, 0) = \phi_{L_0}(0, 0)$, there holds $\phi_L - \phi_{L_0} \rightarrow 0$ in $H^1(\mathbb{R} \times (0, 1))$ and $U_L \rightarrow U_{L_0}$, $\phi_L \rightarrow \phi_{L_0}$ locally uniformly in \mathbb{R}^2 .

More precisely, we refer to [18, Theorems 1.1, 1.5, 1.9] for the existence, uniqueness, monotonicity and global stability of pulsating fronts as well as the sign property of c_L . The continuity of (c_L, ϕ_L, U_L) with respect to L follows directly from [18, Theorem 1.8] and its proof, together with the uniform boundedness of $\|\partial_t U_L\|_{L^\infty(\mathbb{R}^2)} + \|\partial_x U_L\|_{L^\infty(\mathbb{R}^2)}$ with respect to $L \in [L_*, +\infty)$ (which itself follows from standard parabolic estimates). The boundedness of c_L is an easy consequence of [17, Theorem 1.4] which established a uniform bound for the propagation rates of all generalized transition fronts.

In this paper, we are chiefly interested in the asymptotic behavior of the front speeds c_L and the front profiles ϕ_L as $L \rightarrow +\infty$. Before stating our main results, let us first recall some related results in the spatially periodic Fisher-KPP case,² i.e., when, instead of (A1)-(A2), $f(x, u)$ is periodic in x with $f(\cdot, 0) = f(\cdot, 1) = 0$ in \mathbb{R} and

$$0 < f(x, u) \leq \partial_u f(x, 0)u \quad \text{for all } (x, u) \in \mathbb{R} \times (0, 1). \quad (1.7)$$

In this case, for each $L > 0$, there is a minimal front speed c_L^* of pulsating fronts [8, 53] (namely, pulsating fronts exist for all speeds $c \geq c_L^*$, and do not exist if $c < c_L^*$), this minimal speed c_L^* is positive and non-decreasing with respect to $L > 0$ [44], the limit of c_L^* as $L \rightarrow +\infty$ was determined in [34, 35],³ and the homogenization limit of c_L^* as $L \rightarrow 0$ was characterized in [26, 39, 50]. The proofs in the aforementioned papers strongly rely on the fact that the minimal speed is linearly determined and has a variational characterization in terms of principal eigenvalues of parameterized elliptic operators [11, 12, 42, 53]. However, for equation (1.1) with the bistable assumptions (A1)-(A2) or (A1)-(A3), the unique front speed c_L is not linearly determined and it is not monotone in L in general. More precisely, first of all, the limit of c_L as $L \rightarrow 0$ was determined in [18, 37] under certain assumptions (see also Remark 1.5 below). Moreover, in the case where the diffusion rate a is a constant and b in (A1) is piecewise constant, numerical results [34] showed that c_L can in some cases be increasing in small L and then decreasing for larger L . On the other hand, if $a(x)$ truly depends on x while the reaction is homogeneous and

²This case is named after the seminal works [29, 40] on homogeneous equations.

³In multidimensional periodic media, even if the minimal speed $c_L^*(e)$ of pulsating fronts in a given direction e is not monotone in general with respect to L for equations with periodic drifts, the existence of a finite limit of $c_L^*(e)$ as $L \rightarrow +\infty$ is still known, see [35, 44]. We also refer to [1, 16, 32, 33] for further results on the dependence of the minimal or unique speeds of multidimensional pulsating monostable or bistable fronts with respect to the direction e .

has the special form $f(u) = u(1-u)(u-b)$ with $0 < b < 1/2$, it was proved in [19] that the speed c_L for $L \simeq 0$ is smaller than its limit as $L \rightarrow 0$.

In short, the question whether, in the bistable case (A1)-(A3), the speeds c_L of the pulsating fronts of (1.1) (which exist for all large L by Theorem 1.1) converge as $L \rightarrow +\infty$ is the main remaining open problem in the theory of bistable pulsating fronts. This question, which is definitely more challenging than in the Fisher-KPP case (1.7), is solved in this paper (see the main Theorem 1.2 below). Furthermore, a formula for the limit speed is established, as well as further properties on the front profiles as $L \rightarrow +\infty$.

Convergence of the front speeds c_L as $L \rightarrow +\infty$

In this section, we state our main result on the convergence of c_L as $L \rightarrow +\infty$. To do so, we need some notations. For each $y \in \mathbb{R}$, let $u(t, x) = \psi(x - c(y)t, y)$ be the traveling front connecting 0 and 1, with speed $c(y)$, of the following equation with coefficients frozen at y :

$$u_t = a(y)u_{xx} + f(y, u) \quad \text{for } t \in \mathbb{R} \text{ and } x \in \mathbb{R}. \quad (1.8)$$

Namely, $(\psi(\xi, y), c(y))$ satisfies

$$\begin{cases} a(y)\partial_{\xi\xi}\psi(\xi, y) + c(y)\partial_{\xi}\psi(\xi, y) + f(y, \psi(\xi, y)) = 0 & \text{for } \xi \in \mathbb{R}, \\ \psi(-\infty, y) = 1, \quad \psi(+\infty, y) = 0. \end{cases} \quad (1.9)$$

In (1.8)-(1.9), y is viewed as a parameter, since the derivatives only concern the variables (t, x) or ξ . Under conditions (A1)-(A3), it is well known [4, 28] that the function $c : y \mapsto c(y)$ is 1-periodic and positive, and that for each $y \in \mathbb{R}$, the wave profile $\psi(\xi, y)$ is unique up to shifts in ξ (see Section 4.1 below for more properties of the speed $c(y)$ and the wave profile $\psi(\xi, y)$, in particular the C^1 smoothness of $y \mapsto c(y)$). Throughout this paper, we normalize $\psi(\xi, y)$ uniquely in such a way that

$$\psi(0, y) = \frac{1}{2} \quad \text{for each } y \in \mathbb{R}, \quad (1.10)$$

and we denote by c_* the harmonic mean of the 1-periodic function $y \mapsto c(y)$, that is,

$$c_* = \left(\int_0^1 c^{-1}(y) dy \right)^{-1}. \quad (1.11)$$

Theorem 1.2. *Let (A1)-(A3) hold, and for any $L \geq L_*$ (with $L_* > 0$ as in Theorem 1.1), let c_L be the unique front speed of the pulsating front of (1.1). Then, $c_L \rightarrow c_*$ as $L \rightarrow +\infty$.*

Several comments are in order. Let us first point out that Theorem 1.2 coincides with the known result for the minimal front speeds c_L^* of equations with Fisher-KPP reactions (1.7) in the special case where the function $r := \partial_u f(x, 0)$ is a constant, see [35, Corollary 2.4]. Indeed, in the latter case, the limit of the minimal front speeds c_L^* as $L \rightarrow +\infty$ is given by

$$2\sqrt{r} \left(\int_0^1 (a(y))^{-1/2} dy \right)^{-1},$$

which is equal to the harmonic mean of the 1-periodic function $y \mapsto w^*(y) := 2\sqrt{r a(y)}$, where, for each $y \in \mathbb{R}$, $2\sqrt{r a(y)}$ is the minimal front speed of the associated homogeneous equation (1.8). In the case where $\partial_u f(x, 0)$ truly depends on x , the limit of the minimal front speeds c_L^* has a

variational formula in terms of the underlying parameters (see [35, Theorem 2.3]), but it may be not equal to the harmonic mean of the function w^* (see [34] for the case where $\partial_u f(x, 0)$ is a piecewise constant function). It is noteworthy that, in our bistable case (A1)-(A3), Theorem 1.2 shows that the unique wave speed c_L converges to the harmonic mean c_* of the function $y \mapsto c(y)$, no matter whether f depends on x or not.

Theorem 1.2 is proved by taking the large-time and large-space scaling, i.e., $t \rightarrow Lt$, $x \rightarrow Lx$, and determining the singular limit of the Cauchy problem of the rescaled equation

$$\partial_t u_L = L^{-1} \partial_x (a(x) \partial_x u_L) + Lf(x, u_L), \quad t \in \mathbb{R}, x \in \mathbb{R}. \quad (1.12)$$

For a given initial condition independent of L , the singular limit of the solutions u_L of (1.12) as $L \rightarrow +\infty$ has been widely investigated, even in the general case where the coefficients are allowed to be non-periodic. Due to the smallness of the diffusion and the largeness of the reaction in (1.12), an interface separating the regions where $u_L(t, x) \rightarrow 1$ and $u_L(t, x) \rightarrow 0$ is developed in a short time. This property has been proved by different approaches, see e.g. [30] for a probabilistic approach and [5, 6, 7, 15] for a viscosity solution approach. Particularly, in our spatially periodic case, Gärtner's result [30] suggested that, as $L \rightarrow +\infty$, the interface propagates at a mean speed equal to c_* . Thanks to this observation and further crucial uniform estimates on the boundedness of the wave length (more precisely, see Theorem 1.7 below), we show that c_* is the limit of c_L by constructing suitable sub- and super-solutions of (1.12).

Notice that, if the assumption $\int_0^1 f(x, u) du > 0$ in (A3) is replaced by $\int_0^1 f(x, u) du < 0$, equation (1.1) admits a pulsating front with negative speed $c_L < 0$ for all large L , see [18]. In this situation, the quantity c_* is negative, and similarly to Theorem 1.2, we can prove that $c_L \rightarrow c_*$ as $L \rightarrow +\infty$. Furthermore, as an immediate consequence of this observation and [17, Theorem 1.7], we have the following classification of the asymptotic behavior of c_L as $L \rightarrow +\infty$.

Corollary 1.3. *Let (A1)-(A2) hold. Then, we have*

- (i) *if (A3) holds, then $c_L > 0$ for all large L , and $c_L \rightarrow c_*$ as $L \rightarrow +\infty$;*
- (ii) *if $\int_0^1 f(x, u) du < 0$ and $\partial_u f(x, b(x)) > 0$ for all $x \in \mathbb{R}$, then $c_L < 0$ for all large L , and $c_L \rightarrow c_*$ as $L \rightarrow +\infty$;*
- (iii) *if the functions a , $\partial_u f(\cdot, 0)$ and $\partial_u f(\cdot, 1)$ are constants and if*

$$\min_{x \in \mathbb{R}} \int_0^1 f(x, u) du < 0 < \max_{x \in \mathbb{R}} \int_0^1 f(x, u) du,$$

then (1.1) admits pulsating fronts for all large L , with necessarily $c_L = 0$, that is, the fronts are stationary.

Remark 1.4. *By [18], the conditions (A1)-(A3) not only imply the existence of pulsating fronts of (1.1) satisfying (1.4), but they also provide the existence of pulsating fronts propagating in the opposite direction for the same range of periods L . Namely, for any $L \geq L_*$, (1.1) admits a pulsating front of the type $\tilde{U}_L(t, x) = \tilde{\phi}_L(x + \tilde{c}_L t, x/L)$, where $\tilde{\phi}_L$ is 1-periodic in the second variable, and satisfies the reversed limiting conditions:*

$$\tilde{\phi}_L(-\infty, y) = 0, \quad \tilde{\phi}_L(+\infty, y) = 1 \quad \text{uniformly in } y \in \mathbb{R}. \quad (1.13)$$

Moreover, both c_L and \tilde{c}_L are positive. It is worthy pointing out that if the functions $a(x)$ and $f(x, u)$ are not even in x , the speeds c_L and \tilde{c}_L may be dramatically different [16]. Nevertheless, since the function $y \mapsto c(y)$ given in (1.9) is 1-periodic and c_* in (1.11) can be expressed as

$$c_* = \left(\int_{-1/2}^{1/2} c^{-1}(y) dy \right)^{-1} = \left(\int_{-1/2}^{1/2} c^{-1}(-y) dy \right)^{-1},$$

it follows from Theorem 1.2 that c_L and \tilde{c}_L have the same limit c_* as $L \rightarrow +\infty$.

Remark 1.5. For Fisher-KPP functions f satisfying (1.7), for each $L > 0$, pulsating fronts of (1.1) exist, with speeds not smaller than a minimal speed $c_L^* > 0$, and the map $L \mapsto c_L^*$ is non-decreasing [44], hence $\lim_{L \rightarrow 0} c_L^* \leq \lim_{L \rightarrow +\infty} c_L^*$ (and the inequality is strict if the coefficients $a(x)$ and $\partial_u f(x, 0)$ are not constant in x).

In the bistable case with assumptions (A1)-(A3), the situation is in general different and more complicated. Consider for instance a function f defined by $f(x, u) = u(1 - u)(u - b(x))$ for $(x, u) \in \mathbb{R} \times [0, 1]$, where $b : \mathbb{R} \rightarrow (0, 1/2)$ is a given 1-periodic function of class C^1 . Such a function f satisfies (A1)-(A3).

On the one hand, if a is constant, then it follows from [18, Theorems 1.2 and 1.4] that equation (1.1) admits a pulsating front with speed c_L for every small $L > 0$, and that $c_L \rightarrow c_0$ as $L \rightarrow 0$, where $c_0 = \sqrt{2a}(1/2 - \bar{b})$ is the unique speed of a traveling front connecting 0 and 1 for the homogeneous equation $u_t = au_{xx} + u(1 - u)(u - \bar{b})$, with $\bar{b} = \int_0^1 b(x) dx \in (0, 1/2)$. With the notations (1.9), one has $c(y) = \sqrt{2a}(1/2 - b(y))$ for each $y \in \mathbb{R}$, and Theorem 1.2 implies that $c_L \rightarrow c_* = \sqrt{2a} \left(\int_0^1 (1/2 - b(y))^{-1} dy \right)^{-1}$ as $L \rightarrow +\infty$. Therefore, the Cauchy-Schwarz inequality yields $c_0 \geq c_*$ and even $c_0 > c_*$ as soon as b is not constant. This is in sharp contrast with the inequality $\lim_{L \rightarrow 0} c_L^* \leq \lim_{L \rightarrow +\infty} c_L^*$ under the Fisher-KPP assumption (1.7).

On the other hand, if $b \equiv \bar{b} \in (0, 1/2)$ is constant, it follows from [18, Theorems 1.2 and 1.4] that equation (1.1) admits a pulsating front with speed c_L for every small $L > 0$,⁴ and that $c_L \rightarrow c_0$ as $L \rightarrow 0$, where $c_0 = \sqrt{2a_H}(1/2 - \bar{b})$ is the unique speed of a traveling front connecting 0 and 1 for the homogeneous equation $u_t = a_H u_{xx} + u(1 - u)(u - \bar{b})$, with $a_H = \left(\int_0^1 a(y)^{-1} dy \right)^{-1}$ being the harmonic mean of the diffusion coefficient a . With the notations (1.9), one has $c(y) = \sqrt{2a(y)}(1/2 - \bar{b})$ for each $y \in \mathbb{R}$, and Theorem 1.2 implies that $c_L \rightarrow c_* = \sqrt{2} (1/2 - \bar{b}) \left(\int_0^1 \sqrt{a(y)}^{-1} dy \right)^{-1}$ as $L \rightarrow +\infty$. Therefore, the Cauchy-Schwarz inequality yields $c_0 \leq c_*$ and even $c_0 < c_*$ as soon as a is not constant.

To sum up, the influence of the oscillations of the heterogeneities in (1.1) under the bistable assumptions (A1)-(A3) turns out to be more complicated than under the Fisher-KPP assumption (1.7). Moreover, the oscillations in the diffusion and reaction terms can lead to opposite effects on the bistable front speed.

Convergence of the front profiles ϕ_L as $L \rightarrow +\infty$

The second main result of the present paper is concerned with the convergence as $L \rightarrow +\infty$ of the front profiles $\phi_L(\xi, y)$ given in (1.4). It is difficult to address this problem by considering the singular equation (1.12), since the solutions $\phi_L(L(x - c_L t), x)$ of (1.12) converge to a step function as $L \rightarrow +\infty$ for any $t \in \mathbb{R}$ (this will actually be a consequence of Theorem 1.8 below), hence the information of $\phi_L(\xi, y)$ is lost when passing to the limit as $L \rightarrow +\infty$. We instead show the

⁴Notice that, in this case, pulsating fronts are even known to exist for all $L > 0$ if the function a is uniformly close to a constant [27, 54].

convergence of $\phi_L(\xi, y)$ by constructing sub- and super-solutions for the original equation (1.1) when L is large (as a matter of fact, this approach also provides the convergence of c_L as $L \rightarrow +\infty$, as will follow from the proof of Theorem 1.6 below).

Before stating the result, we point out that, since equation (1.1) is invariant by translation in time t , we have to suitably shift the fronts in order to pass to the limit as $L \rightarrow +\infty$. Namely, with $L_* > 0$ as in Theorem 1.1, since for each $L \geq L_*$ the pulsating front U_L is unique up shifts in time, with $U_L(-\infty, x) = 0$ and $U_L(+\infty, x) = 1$ for all $x \in \mathbb{R}$, we normalize for definiteness U_L by requiring

$$U_L(0, 0) = \frac{1}{2}. \quad (1.14)$$

Clearly, under such a normalization, $U_L(t, x)$ is uniquely determined, and the front profile $\phi_L(\xi, y)$ is uniquely determined and satisfies $\phi_L(0, 0) = 1/2$. Now, for each $L \geq L_*$, since $\partial_\xi \phi_L < 0$ in \mathbb{R}^2 and ϕ_L satisfies (1.4) and is at least of class $C^1(\mathbb{R}^2)$, the implicit function theorem yields the existence of a uniquely determined $C^1(\mathbb{R})$ function $y \mapsto \zeta_L(y)$ such that

$$\phi_L(\zeta_L(y), y) = \frac{1}{2}. \quad (1.15)$$

Clearly, $\zeta_L(0) = 0$ and $\zeta_L(y)$ is 1-periodic in $y \in \mathbb{R}$. Moreover, it follows from Theorem 1.1 (ii) and (iv) that the function $(L, y) \mapsto \zeta_L(y)$ is continuous in $[L_*, +\infty) \times \mathbb{R}$.

For the convergence of the front profiles as $L \rightarrow +\infty$, in addition to (A1)-(A3), the following condition will also be assumed:

(A4) the function a is a constant, and there exists a constant $\delta'_0 \in (0, 1/2)$ such that $\partial_x f(x, u) = 0$ for all $x \in \mathbb{R}$ and $u \in [0, \delta'_0] \cup [1 - \delta'_0, 1]$, and then for all $x \in \mathbb{R}$ and $u \in [-\delta'_0, \delta'_0] \cup [1 - \delta'_0, 1 + \delta'_0]$ by (1.3).

Theorem 1.6. *Let (A1)-(A4) hold. For each $L \geq L_*$, let $U_L(t, x) = \phi_L(x - c_L t, x/L)$ be the pulsating front of (1.1) given by Theorem 1.1, and let ζ_L be the function given by (1.15). Let also $(\xi, y) \mapsto \psi(\xi, y)$ and $y \mapsto c(y)$ be given in (1.9)-(1.10). Then,*

$$\sup_{y \in \mathbb{R}, \xi \in \mathbb{R}} \left| \phi_L \left(\xi + \zeta_L(y), \frac{\xi}{L} + y \right) - \psi(\xi, y) \right| \xrightarrow{L \rightarrow +\infty} 0. \quad (1.16)$$

Moreover,

$$\sup_{y \in \mathbb{R}, x \in [-A, A]} \left| \zeta_L \left(y + \frac{x}{L} \right) - \zeta_L(y) - L \int_y^{y+x/L} \left(1 - \frac{c_L}{c(s)} \right) ds \right| \xrightarrow{L \rightarrow +\infty} 0 \quad (1.17)$$

for every $A > 0$, and

$$\sup_{L \geq L_*, y \in \mathbb{R}, x \in [-L, L]} \left| \zeta_L \left(y + \frac{x}{L} \right) - \zeta_L(y) - L \int_y^{y+x/L} \left(1 - \frac{c_L}{c(s)} \right) ds \right| < +\infty. \quad (1.18)$$

Several remarks are in order. First of all, by choosing $y = 0$ and $x = y_0 L$ for some $y_0 \in (0, 1]$ in (1.18), one has

$$\sup_{L \geq L_*} \left| \zeta_L(y_0) - L \int_0^{y_0} \left(1 - \frac{c_L}{c(s)} \right) ds \right| < +\infty.$$

In the special case $y_0 = 1$, since $\zeta_L(1) = \zeta_L(0) = 0$, we rediscover that $c_L \rightarrow c_* = \left(\int_0^1 c(s)^{-1} ds \right)^{-1}$ as $L \rightarrow +\infty$, here under the additional assumption (A4). Furthermore, since the function $y \mapsto$

$c(y)$ is in general not constant even under the assumption (A4), it then follows that $\zeta_L(y_0)$ is unbounded in $L \geq L_*$ for each $y_0 \in (0, 1)$ such that $y_0 \int_0^1 c(s)^{-1} ds \neq \int_0^{y_0} c(s)^{-1} ds$. On the other hand, (1.18) implies that $\zeta_L(y_0) = O(L)$ as $L \rightarrow +\infty$ for all $y_0 \in \mathbb{R}$ (that will coincide with the observation in Theorem 1.8 (iii) below).

Secondly, we point out that (1.16) can be equivalently rephrased as

$$\sup_{y \in \mathbb{R}, \xi \in \mathbb{R}} \left| U_L \left(\frac{Ly - \zeta_L(y)}{c_L}, \xi + Ly \right) - \psi(\xi, y) \right| \xrightarrow{L \rightarrow +\infty} 0.$$

As we will see, this last formulation implies that, as $L \rightarrow +\infty$, the spatial profiles $U_L(t, \cdot)$ look, at any time t and uniformly in space, like shifts of the profiles $\psi(\cdot, y)$ of traveling fronts of the homogeneous equations (1.9) for some values of y . Indeed, by (1.18) and $\lim_{L \rightarrow +\infty} c_L = c_*$, the continuous functions $z \mapsto z - \zeta_L(z)/L$ converge as $L \rightarrow +\infty$ locally uniformly in $z \in \mathbb{R}$ to the function $z \mapsto c_* \int_0^z c(s)^{-1} ds$ (by applying (1.18) with $y = 0$ and $x = Lz$). Therefore, for every $L > 0$ large enough and for every $t \in [0, L/c_L]$, there is $z_{L,t} \in \mathbb{R}$ such that $z_{L,t} - \zeta_L(z_{L,t})/L = c_L t/L$. Hence $U_L(t, \cdot) = U_L((Lz_{L,t} - \zeta_L(z_{L,t}))/c_L, \cdot)$ and $\|U_L(t, \cdot) - \psi(\cdot - Lz_{L,t}, z_{L,t})\|_{L^\infty(\mathbb{R})} \rightarrow 0$ as $L \rightarrow +\infty$ uniformly in $t \in [0, L/c_L]$. Since the profiles $x \mapsto U_L(t, x + c_L t)$ in the frames moving with speeds c_L are periodic in time with period L/c_L , one finally gets that

$$\sup_{t \in \mathbb{R}} d \left(U_L(t, \cdot), \{ \psi(\cdot + a, y) : a \in \mathbb{R}, y \in \mathbb{R} \} \right) \xrightarrow{L \rightarrow +\infty} 0, \quad (1.19)$$

where d is the distance associated with the norm in $L^\infty(\mathbb{R})$. Actually, the convergence of the profiles of the pulsating fronts to profiles of homogeneous traveling fronts as $L \rightarrow +\infty$ is quite natural since equation (1.1) looks like, as $L \rightarrow +\infty$, families of locally spatially homogeneous equations. But the difficulty in Theorem 1.6 and in the consequent formula (1.19) is to show rigorously this convergence and its uniformity with respect to time and space.

The strategy of the proof of Theorem 1.6 can be briefly summarized as follows. First, by virtue of Theorem 1.1 (iii) on the global stability of the pulsating front ϕ_L , we show that, up to some shift in ξ , $\phi_L(\cdot, y)$ is close to $\psi(\cdot, y)$ when L is sufficiently large. Next, we use Theorem 1.8 below on the uniform boundedness of the front width to determine the limit shifts appearing in (1.16). The assumption (A4) is used in the construction of suitable sub- and super-solutions (see Lemmas 4.2-4.3 below), which is the key point of the proof of Theorem 1.6. It is unclear whether this assumption can be relaxed.

Uniform boundedness of front width

As a matter of fact, Theorems 1.2 and 1.6 on the limit speeds and profiles of the pulsating fronts as $L \rightarrow +\infty$ are based on the uniform boundedness of the width of the front profiles, in a sense to be made more precise below. Under conditions (A1)-(A3), with $L_* > 0$ as in Theorem 1.1, it follows from the definition (1.4) of pulsating fronts that, for any $\delta \in (0, 1/2]$ and $L \geq L_*$, there exists a constant $C > 0$ such that

$$\sup_{x \in \mathbb{R}} \text{diam} \left(\{ t \in \mathbb{R} : \delta \leq U_L(t, x) \leq 1 - \delta \} \right) \leq C \quad (1.20)$$

and

$$\sup_{t \in \mathbb{R}} \text{diam} \left(\{ x \in \mathbb{R} : \delta \leq U_L(t, x) \leq 1 - \delta \} \right) \leq C, \quad (1.21)$$

where $\text{diam}(E)$ denotes the diameter of a set $E \subset \mathbb{R}$. Notice that, for each $\delta \in (0, 1/2]$, $L \geq L_*$ and $x \in \mathbb{R}$, the set $\{ t \in \mathbb{R} : \delta \leq U_L(t, x) \leq 1 - \delta \}$ is a compact interval, from (1.4) and the

continuity of U_L and its monotonicity with respect to t . Hence the diameter of the set in (1.20) is equal to its length. We also point out that property (1.21) could be rephrased equivalently by saying that, for any $\lambda \in (0, 1)$, the distance between $c_L t$ and the level set $\{x \in \mathbb{R} : U_L(t, x) = \lambda\}$ is uniformly bounded in t .

Our next result, which is a key stone in determining the convergences of front profiles ϕ_L and front speeds c_L as $L \rightarrow +\infty$, and which has its own interest, shows that the constants C in (1.20)-(1.21) can be taken independently of $L \geq L_*$.

Theorem 1.7. *Let (A1)-(A3) hold. For each $L \geq L_*$, with L_* as in Theorem 1.1, let U_L be a pulsating front of (1.1). Then, for any $\delta \in (0, 1/2]$, there exist constants $C > 0$ and $\beta > 0$ (both are independent of L) such that (1.20) and (1.21) hold true for all $L \geq L_*$, and*

$$\partial_t U_L(t, x) \geq \beta \quad \text{for all } L \geq L_* \text{ and } (t, x) \in \mathbb{R}^2 \text{ such that } \delta \leq U_L(t, x) \leq 1 - \delta. \quad (1.22)$$

The following theorem is a consequence of Theorem 1.7.

Theorem 1.8. *Let (A1)-(A3) hold. For each $L \geq L_*$, with L_* as in Theorem 1.1, let ϕ_L be a front profile of (1.1) and let ζ_L be defined by (1.15) and $\zeta_L(0) = 0$. Then,*

(i) *the following two convergences hold uniformly in $L \geq L_*$ and $y \in \mathbb{R}$:*

$$\lim_{\xi \rightarrow -\infty} \phi_L(\xi + \zeta_L(y), y) = 1 \quad \text{and} \quad \lim_{\xi \rightarrow +\infty} \phi_L(\xi + \zeta_L(y), y) = 0;$$

(ii) *for any $\delta \in (0, 1/2]$, there exists $\beta > 0$ (independent of L) such that*

$$\partial_\xi \phi_L(\xi, y) \leq -\beta \quad \text{for all } L \geq L_* \text{ and } (\xi, y) \in \mathbb{R}^2 \text{ such that } \delta \leq \phi_L(\xi, y) \leq 1 - \delta;$$

(iii) *the function $(L, y) \mapsto \zeta_L(y)/L$ is continuous and bounded in $[L_*, +\infty) \times \mathbb{R}$.*

Notice that the continuity of $(L, y) \mapsto \zeta_L(y)/L$ in $[L_*, +\infty) \times \mathbb{R}$ follows from the comments after (1.15), while the quantities $\zeta_L(y)$ are in general unbounded with respect to $L \geq L_*$, from the discussion after Theorem 1.6.

Theorems 1.7 and 1.8 mean that the fronts do not flatten as $L \rightarrow +\infty$. Their proofs are based on some intersection number arguments inspired by [2] and from the comparison of the pulsating fronts with some well chosen stationary solutions.

Outline of the paper

Section 2 is devoted to the proof of Theorems 1.7 and 1.8 on the uniform boundedness of the wave lengths of the pulsating fronts U_L with respect to $L \geq L_*$. The proof of Theorem 1.2 on the limit of the speeds c_L as $L \rightarrow +\infty$ is done in Section 3, and that of Theorem 1.6 on the limit of the profiles is done in Section 4. Lastly, Section 5 is an appendix on some properties of the profiles $\psi(\xi, y)$ and speeds $c(y)$ of the family of homogeneous equations (1.8)-(1.9).

2 Uniform boundedness of front width: proofs of Theorems 1.7 and 1.8

This section is devoted to the proof of Theorems 1.7 and 1.8. Our proof strongly relies on a zero number argument, and the classification of all solutions to the ordinary differential equations

$$(a_L w')' + f_L(x, w) = 0 \text{ in } \mathbb{R}, \quad w(0) \in (0, 1), \quad (2.23)$$

according to the number of intersections of their graphs with those of the constant solutions 0 and 1. First, after recalling in Section 2.1 some properties on intersection number arguments, we investigate in Section 2.2, for any given $L \geq L_*$ with $L_* > 0$ as in Theorem 1.1 and for any time $t \in \mathbb{R}$, the number and type of intersection points of the spatial profiles $U_L(t, \cdot)$ of the pulsating fronts U_L with all solutions to (2.23). Then, using a passage to the limit as $L \rightarrow +\infty$ and some properties of the stationary solutions to the y -frozen equation (1.8), we prove Theorem 1.7 by contradiction in Section 2.3. Lastly, we show in Section 2.4 that Theorem 1.7 implies Theorem 1.8.

2.1 Zero number properties

We begin with the collection of some zero number properties which will be used later. We refer the reader to [2, 21, 23, 48] for a more detailed overview of the general arguments and their applications.

Let us first introduce some fundamental notions. For any function $w : \mathbb{R} \rightarrow \mathbb{R}$, let $\mathcal{Z}[w]$ denote the number of sign changes of w , namely, the supremum of all $k \in \mathbb{N}$ such that there exist real numbers $x_1 < x_2 < \dots < x_{k+1}$ with

$$w(x_i) \cdot w(x_{i+1}) < 0 \quad \text{for all } i = 1, 2, \dots, k,$$

and we set $\mathcal{Z}[w] = 0$ if w has a constant sign and is not identically 0, and $\mathcal{Z}[w] = -1$ if $w \equiv 0$ in \mathbb{R} . Furthermore, if $0 < \mathcal{Z}[w] < +\infty$, we use $SGN[w]$ to denote the $(k+1)$ -uple of the type $[\varepsilon_1, \dots, \varepsilon_{k+1}]$ consisting of $+$ and $-$, which stand for the signs of $w(x_1), w(x_2), \dots, w(x_{k+1})$. When $\mathcal{Z}[w] = 0$, we set $SGN[w] = [sgn(w)]$ (where $sgn(w) \in \{-, +\}$ is the classical sign function), and we set $SGN[w] = []$ (the empty word) if $w \equiv 0$. By definition, the length of the word $SGN[w]$ is equal to $\mathcal{Z}[w] + 1$.

For any two words A, B consisting of $+$ and $-$, we write $A \triangleright B$ (or equivalently, $B \triangleleft A$) if B is a subword of A .

It follows from the above definitions that \mathcal{Z} is semi-continuous with respect to the pointwise convergence, as stated below.

Lemma 2.1. *Let $(w_n)_{n \in \mathbb{N}}$ be a sequence of real-valued functions converging to w pointwise on \mathbb{R} . Then,*

$$\mathcal{Z}[w] \leq \liminf_{n \rightarrow +\infty} \mathcal{Z}[w_n], \quad \text{and} \quad SGN[w] \triangleleft \liminf_{n \rightarrow +\infty} SGN[w_n].$$

The following lemma is an easy application of the zero number properties for solutions to a linear parabolic equation of the form

$$\partial_t w = a(t, x) \partial_{xx} w + b(t, x) \partial_x w + c(t, x) w \quad \text{in } (t_1, t_2) \times \mathbb{R}, \quad (2.24)$$

where $-\infty \leq t_1 < t_2 \leq +\infty$ and the coefficients $a > 0, a^{-1}, a_t, a_x, a_{xx}, b, b_t, b_x, c$ belong to $L^\infty((t_1, t_2) \times \mathbb{R})$.

Lemma 2.2. [2, 21, 23] *Let $L > 0$, let $u_1(t, x)$ be a bounded solution of the Cauchy problem associated with (1.1) in $(0, +\infty) \times \mathbb{R}$, with a piecewise continuous bounded initial condition $u_1(0, \cdot)$, and let $u_2(x)$ be a $C^2(\mathbb{R})$ stationary solution of (1.1). Assume also that $u_1(0, \cdot) - u_2$ changes sign at most finitely many times on \mathbb{R} , that is, $\mathcal{Z}[u_1(0, \cdot) - u_2] < +\infty$. Then,*

(i) *For any $0 \leq t < t' < +\infty$, we have*

$$\mathcal{Z}[u_1(t, \cdot) - u_2] \geq \mathcal{Z}[u_1(t', \cdot) - u_2]$$

and

$$SGN[u_1(t, \cdot) - u_2] \triangleright SGN[u_1(t', \cdot) - u_2].$$

- (ii) If $u_1(t_*, x_*) = u_2(x_*)$ and $\partial_x u_1(t_*, x_*) = u_2'(x_*)$ for some $t_* > 0$ and $x_* \in \mathbb{R}$, and if $u_1(t, x) \not\equiv u_2(x)$ in $(0, +\infty) \times \mathbb{R}$, then

$$\mathcal{Z}[u_1(t, \cdot) - u_2] - 2 \geq \mathcal{Z}[u_1(s, \cdot) - u_2] \geq 0 \quad \text{for all } t \in (0, t_*) \text{ and } s \in (t_*, +\infty).$$

Since the function $(x, u) \mapsto f(x, u)$ is globally Lipschitz-continuous with respect to $u \in \mathbb{R}$ uniformly in $x \in \mathbb{R}$, it is easily seen that $w(t, x) := u_1(t, x) - u_2(x)$ satisfies an equation of the form (2.24) in $(0, +\infty) \times \mathbb{R}$ with a bounded continuous coefficient c defined by $c(t, x) := (f_L(x, u_1(t, x)) - f_L(x, u_2(x)))/(u_1(t, x) - u_2(x))$ if $u_1(t, x) \neq u_2(x)$ and $c(t, x) := \partial_u f_L(x, u_1(t, x))$ if $u_1(t, x) = u_2(x)$. Then, the proof of Lemma 2.2 follows from that of [23, Lemma 2.4] with some obvious modifications; therefore, we omit the details.

2.2 Intersections of pulsating fronts with stationary solutions

Now, for any given $L \geq L_*$, with $L_* > 0$ as in Theorem 1.1, we investigate $\mathcal{Z}[U_L(t, \cdot) - w]$ and $SGN[U_L(t, \cdot) - w]$ at any time $t \in \mathbb{R}$, where $U_L : \mathbb{R}^2 \rightarrow (0, 1)$ is the pulsating front of (1.1) and w is an arbitrary $C^2(\mathbb{R})$ solution to (2.23). Let us first classify the solutions of (2.23) as follows.

Lemma 2.3. *Let $L \geq L_*$ be fixed, and let $w \in C^2(\mathbb{R})$ be a solution of (2.23). Then, w must be one of the following types:*

- (a) $\mathcal{Z}[w - 1] = 1$ and $\mathcal{Z}[w] = 1$;
- (b) $\mathcal{Z}[w - 1] = 0$, $\mathcal{Z}[w] = 0$, and $0 < w < 1$ in \mathbb{R} ;
- (c) $\mathcal{Z}[w - 1] = 0$, $\mathcal{Z}[w] = 1$, $w < 1$ in \mathbb{R} , and $SGN[w] = [-, +]$;
- (d) $\mathcal{Z}[w - 1] = 0$, $\mathcal{Z}[w] = 1$, $w < 1$ in \mathbb{R} , and $SGN[w] = [+ , -]$;
- (e) $\mathcal{Z}[w - 1] = 0$, $\mathcal{Z}[w] = 2$, $w < 1$ in \mathbb{R} , and $SGN[w] = [-, +, -]$.

Proof. First of all, since $w(0) \in (0, 1)$, we have $\mathcal{Z}[w - 1] \geq 0$ and $\mathcal{Z}[w] \geq 0$. Suppose now that there exists some $x_1 \in \mathbb{R}$ such that $w(x_1) = 1$. It is then clear that $w'(x_1) \neq 0$ (otherwise, since $f_L(\cdot, 1) \equiv 0$, the Cauchy-Lipschitz theorem would imply that $w \equiv 1$ in \mathbb{R} , a contradiction with $w(0) \in (0, 1)$).

Then, we observe that, if $w'(x_1) > 0$, then $w(x) > 1$ for all $x > x_1$. Indeed, otherwise, there would exist some $x_2 > x_1$ such that $w(x_2) = 1$, $w'(x_2) < 0$ and $w(x) > 1$ for all $x \in (x_1, x_2)$. Since $\max_{\mathbb{R}} f_L(\cdot, u) < 0$ for all $u > 1$, one would have $(a_L w')' > 0$ in (x_1, x_2) , and hence, $a_L(x_2)w'(x_2) > a_L(x_1)w'(x_1) > 0$. Since a is a positive function, there holds $w'(x_2) > 0$, which is a contradiction. Similarly, we can conclude that, if $w'(x_1) < 0$, then $w(x) > 1$ for all $x < x_1$. As a consequence, we obtain that $\mathcal{Z}[w - 1] \leq 2$, and if $\mathcal{Z}[w - 1] = 2$, then $SGN[w - 1] = [+ , - , +]$.

Since $\min_{\mathbb{R}} f_L(\cdot, u) > 0$ for all $u < 0$, with similar arguments as above, one gets that $\mathcal{Z}[w] \leq 2$, and if $\mathcal{Z}[w] = 2$, then $SGN[w] = [-, +, -]$.

Consequently, we always have

$$\mathcal{Z}[w - 1] + \mathcal{Z}[w] \leq 2.$$

Observe that the Cauchy-Lipschitz theorem also implies that, if $w \leq 1$ in \mathbb{R} (respectively $w \geq 0$ in \mathbb{R}), then $w < 1$ in \mathbb{R} (respectively $w > 0$ in \mathbb{R}). Therefore, since $0 < w(0) < 1$, one gets that $w < 1$ in \mathbb{R} if $\mathcal{Z}[w - 1] = 0$ (respectively $w > 0$ in \mathbb{R} if $\mathcal{Z}[w] = 0$).

Now, to complete the proof, it remains to exclude the following three cases:

- $\mathcal{Z}[w - 1] = 1$, $\mathcal{Z}[w] = 0$ and $SGN[w - 1] = [+ , -]$;
- $\mathcal{Z}[w - 1] = 2$, $\mathcal{Z}[w] = 0$ and $SGN[w - 1] = [+ , - , +]$;
- $\mathcal{Z}[w - 1] = 1$, $\mathcal{Z}[w] = 0$ and $SGN[w - 1] = [- , +]$.

In the first two cases, one can find a continuous function $g : \mathbb{R} \rightarrow [0, 1]$ satisfying (1.6) such that $g < w$ in \mathbb{R} . Let $u(t, x)$ be the solution of the Cauchy problem associated with (1.1) with initial condition $u(0, x) = g(x)$. Then, the comparison principle immediately implies that

$$u(t, x) < w(x) \quad \text{for all } t \geq 0 \text{ and } x \in \mathbb{R}. \quad (2.25)$$

On the other hand, it follows from Theorem 1.1 (iii) that $\|u(t, \cdot) - U_L(t + \tau_0, \cdot)\|_{L^\infty(\mathbb{R})} \rightarrow 0$ as $t \rightarrow +\infty$, for some $\tau_0 \in \mathbb{R}$. Remember that the wave speed c_L is positive. This implies in particular that $u(t, 0) \rightarrow 1$ as $t \rightarrow +\infty$. Since $w(0) \in (0, 1)$, we have $u(t, 0) > w(0)$ for all large t , which is a contradiction with (2.25).

Finally, if the last case occurs, then one reaches a similar contradiction by noticing that equation (1.1) admits a pulsating front of the type $\tilde{U}_L(t, x) = \tilde{\phi}_L(x + \tilde{c}_L t, x/L)$ satisfying the asymptotic conditions (1.13), and that such a front is also globally asymptotic stable, with positive speed \tilde{c}_L (see Remark 1.4). The proof of Lemma 2.3 is thus complete. \square

Lemma 2.4. *Let $L \geq L_*$ be fixed, let U_L be a pulsating front of (1.1), and let $w \in C^2(\mathbb{R})$ be a solution of $(a_L w)' + f_L(x, w) = 0$ in \mathbb{R} . Then $\mathcal{Z}[U_L(t, \cdot) - w] \leq 2$ for all $t \in \mathbb{R}$. Furthermore, if $\mathcal{Z}[U_L(t, \cdot) - w] = 2$ for some $t \in \mathbb{R}$, then there must hold*

$$SGN[U_L(t, \cdot) - w] = [+ , - , +]. \quad (2.26)$$

Proof. Let $t_* \in \mathbb{R}$ be an arbitrary time such that $U_L(t_*, x_*) = w(x_*)$ for some $x_* \in \mathbb{R}$. Denote $\alpha := w(x_*)$. Notice that $\alpha \in (0, 1)$ since $0 < U_L < 1$ in \mathbb{R}^2 , hence $w(\cdot + x_*)$ solves (2.23) and Lemma 2.3 applies to $w(\cdot + x_*)$ and then to w as well. To show Lemma 2.4, it suffices to prove that either $\mathcal{Z}[U_L(t_*, \cdot) - w] \leq 1$, or $\mathcal{Z}[U_L(t_*, \cdot) - w] = 2$ and (2.26) holds with $t = t_*$.

For any $z \in \mathbb{R}$, let $\hat{u}(t, x; z)$ be the solution of the Cauchy problem associated with (1.1) with initial condition

$$\hat{u}(0, x; z) = H(z - x),$$

where H denotes the Heaviside function defined by $H(y) = 0$ if $y < 0$ and $H(y) = 1$ if $y \geq 0$. When $z < x_*$, we define

$$\tau(z) := \min \{t > 0 : \hat{u}(t, x_*; z) = \alpha\}. \quad (2.27)$$

Let us first observe that $\tau(z)$ is well defined and $0 < \tau(z) < +\infty$. Indeed, since $c_L > 0$, it follows from Theorem 1.1 (iii) that $\lim_{t \rightarrow +\infty} \hat{u}(t, x_*; z) = 1 > \alpha$. This, together with the continuity of $\hat{u}(\cdot, x_*; z)$ in $[0, +\infty)$ and the fact that $\hat{u}(0, x_*; z) = 0 < \alpha$, immediately implies that the minimum in (2.27) is well defined and $0 < \tau(z) < +\infty$. Moreover, since $\hat{u}(t, x; z) \rightarrow 0$ as $z \rightarrow -\infty$ locally uniformly in $(t, x) \in [0, +\infty) \times \mathbb{R}$, we have

$$\lim_{z \rightarrow -\infty} \tau(z) = +\infty.$$

Furthermore, by the proof of [23, Lemma 3.1], there is an entire solution $\hat{u}_\infty : \mathbb{R}^2 \rightarrow (0, 1)$ of (1.1) that is steeper than any other entire solution between 0 and 1,⁵ and such that the following limit

⁵The steepness is understood in the following sense: for any two entire solutions $u_1 : \mathbb{R} \times \mathbb{R} \rightarrow (0, 1)$ and $u_2 : \mathbb{R} \times \mathbb{R} \rightarrow (0, 1)$ of (1.1), we say that u_1 is steeper than u_2 provided that $SGN[u_1(t + t', \cdot) - u_2(t, \cdot)] \triangleleft [+ , -]$ for any t and t' in \mathbb{R} .

exists in the topology of $C_{t;x;loc}^{1;2}(\mathbb{R}^2)$, up to extraction of a subsequence:

$$\lim_{k \rightarrow +\infty} \widehat{u}(t + \tau(-kL), x; -kL) = \widehat{u}_\infty(t, x). \quad (2.28)$$

Next, we claim that

$$\widehat{u}_\infty(t, x) \equiv U_L(t + t_*, x). \quad (2.29)$$

Indeed, since the pulsating front of (1.1) is unique up to shifts in time (see Theorem 1.1 (ii)) and since U_L is steeper than any other entire solution between 0 and 1 by [31, Theorem 1.8], the function U_L is then identically equal to \widehat{u}_∞ up to a shift in time. Since $\widehat{u}_\infty(0, x_*) = U_L(t_*, x_*) = \alpha$, we obtain (2.29).

For clarity, we divide the remaining arguments into two parts, according to the behavior of w . First, if w is one of the types (a)-(c) listed in Lemma 2.3, then there exists $k_* \in \mathbb{N}$ sufficiently large such that

$$\mathcal{Z}[\widehat{u}(0, \cdot; -kL) - w] = 1 \quad \text{for all } k \geq k_*.$$

It then follows from Lemma 2.2 that

$$\mathcal{Z}[\widehat{u}(t + \tau(-kL), \cdot; -kL) - w] \leq 1 \quad \text{for all } t \geq -\tau(-kL) \text{ and } k \geq k_*.$$

Hence, by Lemma 2.1 and (2.28)-(2.29), we obtain that $\mathcal{Z}[U_L(t, \cdot) - w] \leq 1$ for all $t \in \mathbb{R}$. Notice actually that, since $U_L(t_*, x_*) = w(x_*) = \alpha$ and U_L is not stationary, Lemma 2.2 implies that $\mathcal{Z}[U_L(t_*, \cdot) - w] = 1$.

Next, we consider the case where w is one of the types (d)-(e) in Lemma 2.3. In this situation, it is also easily seen that

$$\mathcal{Z}[\widehat{u}(0, \cdot; -kL) - w] = 2, \quad \text{and} \quad \text{SGN}[\widehat{u}(0, \cdot; -kL) - w] = [+ , - , +] \quad \text{for all large } k \in \mathbb{N}.$$

By using Lemmas 2.1 and 2.2 again, we have

$$\mathcal{Z}[U_L(t_*, \cdot) - w] \leq 2, \quad \text{and} \quad \text{SGN}[U_L(t_*, \cdot) - w] \triangleleft [+ , - , +].$$

In particular, this implies that if $\mathcal{Z}[U_L(t_*, \cdot) - w] = 2$, then $\text{SGN}[U_L(t_*, \cdot) - w] = [+ , - , +]$. The proof of Lemma 2.4 is thus complete. \square

In the following lemma, we give a special solution of (2.23) that decays to 0 as $x \rightarrow +\infty$, with a decay rate that is controlled uniformly with respect to $L \geq L_*$.

Lemma 2.5. *For any $L \geq L_*$ and any $\delta \in (0, \delta_0]$, where $\delta_0 \in (0, 1/2)$ is the constant provided by (A2), there exists a solution $w \in C^2(\mathbb{R})$ of (2.23) such that*

$$w(0) = \delta, \quad \mathcal{Z}[w - 1] = 0, \quad w < 1 \text{ in } \mathbb{R}, \quad \mathcal{Z}[w] \leq 1, \quad (2.30)$$

and

$$0 < w(x) \leq \delta e^{-\mu x} \quad \text{for all } x \geq 0, \quad (2.31)$$

for some constant $\mu > 0$ independent of L .

Proof. For any $L \geq L_*$ and any $\delta \in (0, \delta_0]$, we show the existence of the desired solution by an approximation argument. For each $n \in \mathbb{N}$, let us first consider the following problem in bounded interval :

$$\begin{cases} (a_L w)' + f_L(x, w) = 0 & \text{in } [0, n], \\ w(0) = \delta, \quad w(n) = 0. \end{cases} \quad (2.32)$$

Since $f_L(x, 0) \equiv 0$ and since $f_L(x, \delta) < 0$ for all $x \in \mathbb{R}$, it follows that $w = 0$ and $w = \delta$ are, respectively, a sub-solution and a super-solution of (2.32). Then, a classical iteration argument implies that there exists a solution $w_n \in C^2([0, n])$ such that $0 \leq w_n \leq \delta$ in $[0, n]$. Furthermore, $0 < w_n < \delta$ in $(0, n)$ from the strong maximum principle and, from the sliding method [13], each w_n is unique and the sequence $(w_n)_{n \in \mathbb{N}}$ is nondecreasing in $n \in \mathbb{N}$, in the sense that $w_n \leq w_m$ in $[0, n]$ if $n \leq m$. Define now

$$w(x) := \lim_{n \rightarrow +\infty} w_n(x) \text{ for } x \in [0, +\infty).$$

It then follows from the standard elliptic estimates that w is a $C^2([0, +\infty))$ solution of

$$(a_L w')' + f_L(x, w) = 0 \text{ in } [0, +\infty), \quad w(0) = \delta, \quad 0 \leq w \leq \delta \text{ in } [0, +\infty),$$

hence $0 < w < \delta$ in $(0, +\infty)$ from the strong maximum principle. Since $f_L(x, u)$ is continuous in $\mathbb{R} \times \mathbb{R}$ and globally Lipschitz-continuous in $u \in \mathbb{R}$ uniformly in $x \in \mathbb{R}$, the function w can be extended to $x \in \mathbb{R}$ so that it is a $C^2(\mathbb{R})$ solution of (2.23) with $w(0) = \delta$. It further follows from Lemma 2.3 that this solution must satisfy $\mathcal{Z}[w - 1] = 0$, $w < 1$ in \mathbb{R} , and $\mathcal{Z}[w] \leq 1$ (more precisely, only cases (b) or (c) are possible).

Now, to complete the proof, it remains to show the estimate (2.31). Since $0 \leq w_n \leq \delta$ in $[0, n]$, by the first line of (1.2), we have, for each $n \in \mathbb{N}$,

$$(a_L w_n')' - \gamma_0 w_n \geq 0 \text{ for } x \in [0, n].$$

Since $a_L(x) = a(x/L)$ is L -periodic and a is positive and at least of class $C^1(\mathbb{R})$, we can find a small constant $\mu > 0$ (independent of $L \geq L_*$) such that $a_L(x)\mu^2 - a_L'(x)\mu - \gamma_0 \leq 0$ for all $x \in \mathbb{R}$. Letting $\bar{w}(x) := \delta e^{-\mu x}$ for $x \geq 0$, we compute that

$$(a_L \bar{w}')'(x) - \gamma_0 \bar{w}(x) = (a_L(x)\mu^2 - a_L'(x)\mu - \gamma_0)\bar{w}(x) \leq 0 \text{ for all } x \geq 0.$$

It then follows from the elliptic weak maximum principle that $w_n(x) \leq \bar{w}(x) = \delta e^{-\mu x}$ for all $x \in [0, n]$ and $n \in \mathbb{N}$, hence $w(x) \leq \delta e^{-\mu x}$ for all $x \geq 0$. This ends the proof of Lemma 2.5. \square

In the last lemma of this subsection, we consider the intersection of the pulsating front with the stationary solution of (2.23) obtained in Lemma 2.5.

Lemma 2.6. *Let $L \geq L_*$ be fixed, let U_L be a pulsating front of (1.1), let $\delta \in (0, \delta_0]$ and let $w \in C^2(\mathbb{R})$ be the solution of (2.23) provided by Lemma 2.5. Then, $\mathcal{Z}[U_L(t, \cdot) - w] \leq 1$ for every $t \in \mathbb{R}$. Furthermore, if $\mathcal{Z}[U_L(t, \cdot) - w] = 1$ for some $t \in \mathbb{R}$, then $SGN[U_L(t, \cdot) - w] = [+ , -]$.*

Proof. The proof is similar to that of Lemma 2.4; therefore, we only give its outline. Let $t_* \in \mathbb{R}$ be arbitrary. Without loss of generality, we suppose that there exists some $x_* \in \mathbb{R}$ such that $U_L(t_*, x_*) = w(x_*)$, hence $w(x_*) \in (0, 1)$. To show the lemma, it suffices to prove that $\mathcal{Z}[U_L(t_*, \cdot) - w] \leq 0$, or $\mathcal{Z}[U_L(t_*, \cdot) - w] = 1$ and $SGN[U_L(t_*, \cdot) - w] = [+ , -]$.

For any $z < x_*$, let $\hat{u}(t, x; z)$ be the solution of (1.1) with Heaviside type initial condition $H(z - \cdot)$ and let $\tau(z)$ be defined as in (2.27) with $\alpha = w(x_*)$. By the proof of Lemma 2.4, we see that $0 < \tau(z) < +\infty$, $\lim_{z \rightarrow -\infty} \tau(z) = +\infty$, and that, up to extraction of a subsequence,

$$\hat{u}(t + \tau(-kL), x; -kL) \rightarrow U_L(t + t_*, x) \text{ as } k \rightarrow +\infty \text{ in } C_{t; x; loc}^{1; 2}(\mathbb{R}^2).$$

Since w satisfies (2.30) and $\lim_{x \rightarrow +\infty} w(x) = 0$, it is straightforward to check that

$$\mathcal{Z}[\hat{u}(0, \cdot; -kL) - w] = 1 \quad \text{and} \quad SGN[\hat{u}(0, \cdot; -kL) - w] = [+ , -]$$

for all large $k \in \mathbb{N}$. Passing to the limit as $k \rightarrow +\infty$, it follows from Lemmas 2.1 and 2.2 that $\mathcal{Z}[U_L(t_*, \cdot) - w] \leq 1$ and $SGN[U_L(t_*, \cdot) - w] \triangleleft [+ , -]$, yielding the desired results. \square

2.3 Proof of Theorem 1.7

For clarity, we divide the proof into three steps. In our arguments below, $\delta \in (0, 1/2]$ is given.

Step 1: (1.20) holds true with $C > 0$ independent of $L \geq L_*$. For any $L \geq L_*$ and $x \in \mathbb{R}$, since $U_L(t, x)$ is continuous and increasing in t with $U_L(-\infty, x) = 0$ and $U_L(+\infty, x) = 1$, the set

$$I_L(x) := \{t \in \mathbb{R} : \delta \leq U_L(t, x) \leq 1 - \delta\}$$

is a compact interval in \mathbb{R} . Denote by $m(I_L(x))$ the length of this interval. It follows from Theorem 1.1 (ii) and (iv) that $m(I_L(x))$ is continuous with respect to $(L, x) \in [L_*, +\infty) \times \mathbb{R}$.

Assume by contradiction that the desired result (1.20) is not true. Then, there exist sequences $(L_n)_{n \in \mathbb{N}} \subset [L_*, +\infty)$ and $(x_n)_{n \in \mathbb{N}} \subset \mathbb{R}$ such that

$$m(I_{L_n}(x_n)) = m(\{t \in \mathbb{R} : \delta \leq U_{L_n}(t, x_n) \leq 1 - \delta\}) \rightarrow +\infty \text{ as } n \rightarrow +\infty. \quad (2.33)$$

From (1.5), one easily sees that $m(I_n(x_n)) = m(I_n(x_n + L_n))$. Therefore, without loss of generality, we may assume that $(x_n)_{n \in \mathbb{N}} \subset [0, L_n]$. Moreover, it follows from the continuity of $m(I_L(x))$ in (L, x) that $L_n \rightarrow +\infty$ as $n \rightarrow +\infty$.

For each $n \in \mathbb{N}$, since $U_{L_n}(\cdot, x_n)$ is increasing and continuous, with $U_{L_n}(-\infty, x_n) = 0$ and $U_{L_n}(+\infty, x_n) = 1$, there exists a unique $t_n \in I_{L_n}(x_n)$ such that $U_{L_n}(t_n, x_n) = 1/2$. Define

$$\tilde{U}_n(t, x) := U_{L_n}(t + t_n, x + x_n) \text{ for } (t, x) \in \mathbb{R}^2.$$

Clearly, each function \tilde{U}_n is an entire solution of

$$\partial_t \tilde{U}_n = \partial_x \left(a \left(\frac{x + x_n}{L_n} \right) \partial_x \tilde{U}_n \right) + f \left(\frac{x + x_n}{L_n}, \tilde{U}_n \right) \text{ in } \mathbb{R}^2, \quad (2.34)$$

with $\tilde{U}_n(0, 0) = 1/2$. Up to extraction of some subsequence, one can assume that $x_n/L_n \rightarrow x_\infty \in [0, 1]$, and that, from standard parabolic estimates, there exists a function $\tilde{U}_\infty \in C_{t;x}^{1;2}(\mathbb{R}^2, [0, 1])$ such that

$$\tilde{U}_n \rightarrow \tilde{U}_\infty \text{ in } C_{t;x;loc}^{1;2}(\mathbb{R}^2) \text{ as } n \rightarrow +\infty. \quad (2.35)$$

Moreover, \tilde{U}_∞ satisfies $\tilde{U}_\infty(0, 0) = 1/2$ and it is a solution to

$$\partial_t \tilde{U}_\infty = a(x_\infty) \partial_{xx} \tilde{U}_\infty + f(x_\infty, \tilde{U}_\infty) \text{ in } \mathbb{R}^2. \quad (2.36)$$

By the strong maximum principle, we have $0 < \tilde{U}_\infty < 1$ in \mathbb{R}^2 . Furthermore, since $\partial_t \tilde{U}_n > 0$ in \mathbb{R}^2 , we have $\partial_t \tilde{U}_\infty \geq 0$ in \mathbb{R}^2 . Finally, one infers from (2.33) that

$$m(\{t \in \mathbb{R} : \delta \leq \tilde{U}_\infty(t, 0) \leq 1 - \delta\}) = +\infty. \quad (2.37)$$

By the monotonicity of \tilde{U}_∞ in t and standard parabolic estimates applied to equation (2.36), it follows that there exist two steady states $0 \leq p_- \leq p_+ \leq 1$ of (2.36) such that

$$\tilde{U}_\infty(t, \cdot) \rightarrow p_\pm \text{ as } t \rightarrow \pm\infty \text{ in } C_{loc}^2(\mathbb{R}). \quad (2.38)$$

It is clear that p_\pm satisfy

$$a(x_\infty) p_\pm'' + f(x_\infty, p_\pm) = 0 \text{ in } \mathbb{R}, \quad (2.39)$$

and $0 \leq p_-(0) \leq 1/2 \leq p_+(0) \leq 1$. Moreover, due to (2.37), one can conclude that either $\delta \leq p_-(0) \leq 1/2$ or $1/2 \leq p_+(0) \leq 1 - \delta$. Then, by the strong maximum principle, we have

$$\text{either } 0 < p_- < 1 \text{ in } \mathbb{R}, \text{ or } 0 < p_+ < 1 \text{ in } \mathbb{R}.$$

Without loss generality, we assume that the former case happens, since the latter case can be handled in a similar way. Then, thanks to the assumptions (A1) and (A3), according to the phase diagrams of equation (2.39), the solution p_- can only be one of the following three types: either a constant function, or a non-constant periodic function, or a ground state solution such that $p_-(x) \rightarrow 0$ as $x \rightarrow \pm\infty$. We will derive a contradiction in each of these three cases.

Case 1: p_- is a constant solution, that is, $p_- \equiv b(x_\infty)$ in \mathbb{R} . In this case, for any non-constant periodic solution $0 < q < 1$ of (2.39), one finds some points $y_1 < z_1 < y_2 < z_2$ and some constant $\varepsilon_0 > 0$ such that

$$q(y_i) - p_-(y_i) \geq \varepsilon_0, \quad \text{and} \quad q(z_i) - p_-(z_i) \leq -\varepsilon_0 \quad \text{for } i = 1, 2. \quad (2.40)$$

Remember that function $u \mapsto f(x, u)$ is globally Lipschitz-continuous in $u \in \mathbb{R}$ uniformly in $x \in \mathbb{R}$. Then, for each $n \in \mathbb{N}$, it follows from the classical ODE theory that the following problem

$$\left(a \left(\frac{x + x_n}{L_n} \right) w' \right)' + f \left(\frac{x + x_n}{L_n}, w \right) = 0 \quad \text{in } \mathbb{R}, \quad (2.41)$$

with $w(y_1) = q(y_1)$ and $w'(y_1) = q'(y_1)$, admits a unique solution $w_n \in C^2(\mathbb{R})$. Moreover, by standard elliptic estimates, we have $w_n \rightarrow q$ as $n \rightarrow +\infty$ in $C_{loc}^2(\mathbb{R})$. This implies in particular that there exists $N \in \mathbb{N}$ such that

$$|w_n(x) - q(x)| \leq \frac{\varepsilon_0}{4} \quad \text{for all } x \in [y_1, z_2] \text{ and } n \geq N.$$

On the other hand, it follows from (2.35) and (2.38) that there exists $s_* < 0$, with $|s_*|$ sufficiently large, such that, replacing N by a larger integer if necessary,

$$\left| \tilde{U}_n(s_*, x) - p_-(x) \right| \leq \frac{\varepsilon_0}{4} \quad \text{for all } x \in [y_1, z_2] \text{ and } n \geq N.$$

Combining the above, we obtain that, for all large n ,

$$w_n(y_i) - \tilde{U}_n(s_*, y_i) \geq \frac{\varepsilon_0}{2}, \quad \text{and} \quad w_n(z_i) - \tilde{U}_n(s_*, z_i) \leq -\frac{\varepsilon_0}{2} \quad \text{for } i = 1, 2.$$

By continuity, the function $x \mapsto \tilde{U}_n(s_*, x) - w_n(x)$ has at least three sign changes in $[y_1, z_2]$, which is a contradiction with the fact that $\mathcal{Z}[\tilde{U}_n(s_*, \cdot) - w_n] \leq 2$ (by Lemma 2.4 applied with $\tilde{U}_n(\cdot, \cdot - x_n)$ and $w_n(\cdot - x_n)$). Hence, Case 1 is ruled out.

Case 2: p_- is a non-constant periodic solution. In this case, letting $q(x) \equiv b(x_\infty)$, we see that q is a constant solution of (2.39). Then, we can find some points $y_1 < z_1 < y_2 < z_2$ and a constant $\varepsilon_0 > 0$ such that (2.40) holds true, and hence, the same reasoning as in Case 1 yields a contradiction. Therefore, Case 2 is ruled out too.

Case 3: p_- is a ground state solution with $p_-(x) \rightarrow 0$ as $x \rightarrow \pm\infty$. In the case, $p_-(x)$ is symmetrically decreasing and there exist some points $\tilde{y}_1 < \tilde{y}_0 < \tilde{y}_2$ and a constant $\tilde{\varepsilon}_0 > 0$ such that

$$b(x_\infty) - p_-(\tilde{y}_i) \geq \tilde{\varepsilon}_0 \quad \text{for } i = 1, 2, \quad \text{and} \quad b(x_\infty) - p_-(\tilde{y}_0) \leq -\tilde{\varepsilon}_0.$$

Similarly as in Case 1, for each $n \in \mathbb{N}$, problem (2.41) with $w(\tilde{y}_1) = b(x_\infty)$ and $w'(\tilde{y}_1) = 0$ has a unique solution $\tilde{w}_n \in C^2(\mathbb{R})$, and from standard elliptic estimates, $\tilde{w}_n \rightarrow b(x_\infty)$ as $n \rightarrow +\infty$ in $C_{loc}^2(\mathbb{R})$. As a consequence, there exists a large integer $\tilde{N} \in \mathbb{N}$ such that

$$|\tilde{w}_n(x) - b(x_\infty)| \leq \frac{\tilde{\varepsilon}_0}{4} \quad \text{for all } x \in [\tilde{y}_1, \tilde{y}_2] \text{ and } n \geq \tilde{N}.$$

On the other hand, by using (2.35) and (2.38) again and by making some adjustment to \tilde{N} if necessary, one finds some negative time $\tilde{s}_* < 0$, with $|\tilde{s}_*|$ large enough, such that

$$\left| \tilde{U}_n(\tilde{s}_*, x) - p_-(x) \right| \leq \frac{\tilde{\varepsilon}_0}{4} \quad \text{for all } x \in [\tilde{y}_1, \tilde{y}_2] \text{ and } n \geq \tilde{N}.$$

Therefore, we obtain that, for each $n \geq \tilde{N}$,

$$\tilde{w}_n(\tilde{y}_i) - \tilde{U}_n(\tilde{s}_*, \tilde{y}_i) \geq \frac{\tilde{\varepsilon}_0}{2} \quad \text{for } i = 1, 2, \quad \text{and} \quad \tilde{w}_n(\tilde{y}_0) - \tilde{U}_n(\tilde{s}_*, \tilde{y}_0) \leq -\frac{\tilde{\varepsilon}_0}{2}. \quad (2.42)$$

Remember that by Lemma 2.4 (applied with $\tilde{U}_n(\cdot, \cdot - x_n)$ and $\tilde{w}_n(\cdot - x_n)$), we have $\mathcal{Z}[\tilde{U}_n(\tilde{s}_*, \cdot) - \tilde{w}_n] \leq 2$. Then, (2.42) implies that

$$\mathcal{Z}[\tilde{U}_n(\tilde{s}_*, \cdot) - \tilde{w}_n] = 2, \quad \text{and} \quad \text{SGN}[\tilde{U}_n(\tilde{s}_*, \cdot) - \tilde{w}_n] = [-, +, -].$$

This last property contradicts property (2.26) of Lemma 2.4, hence Case 3 is ruled out too.

Consequently, the case where $0 < p_- < 1$ cannot happen. Furthermore, with similar arguments, one can exclude the case where $0 < p_+ < 1$. Therefore, our assumption (2.33) at the beginning of the proof is unreasonable. This ends the proof of Step 1.

Step 2: (1.22) holds true with $\beta > 0$ independent of $L \geq L_*$. Assume by contradiction that there exist sequences $(\bar{L}_n)_{n \in \mathbb{N}} \subset [L_*, +\infty)$ and $(\bar{t}_n, \bar{x}_n)_{n \in \mathbb{N}} \subset \mathbb{R}^2$ such that $U_{\bar{L}_n}(\bar{t}_n, \bar{x}_n) \in [\delta, 1 - \delta]$ for each $n \in \mathbb{N}$, and

$$\partial_t U_{\bar{L}_n}(\bar{t}_n, \bar{x}_n) \rightarrow 0 \quad \text{as } n \rightarrow +\infty.$$

Due to (1.5), we may assume without loss of generality that $0 \leq \bar{x}_n < \bar{L}_n$. By Theorem 1.1 (ii) (iv) and standard parabolic estimates, it follows that $\bar{L}_n \rightarrow +\infty$ as $n \rightarrow +\infty$.

Define now

$$\bar{U}_n(t, x) := U_{\bar{L}_n}(t + \bar{t}_n, x + \bar{x}_n) \quad \text{for } (t, x) \in \mathbb{R}^2.$$

Up to extraction of some subsequence, we may assume that the sequence $(\bar{x}_n/\bar{L}_n)_{n \in \mathbb{N}} \subset [0, 1]$ converges as $n \rightarrow +\infty$. By a slight abuse of notation, we still denote by x_∞ this limit. Then, standard parabolic estimates imply that, possibly up to extraction of a further subsequence, $\bar{U}_n \rightarrow \bar{U}_\infty$ in $C_{t;x;loc}^{1;2}(\mathbb{R}^2)$, where \bar{U}_∞ is an entire solution of (2.36). It is clear that $\partial_t \bar{U}_\infty \geq 0$ in \mathbb{R}^2 and $\partial_t \bar{U}_\infty(0, 0) = 0$. Furthermore, the strong maximum principle applied to the equation satisfied by $\partial_t \bar{U}_\infty$ immediately gives that $\partial_t \bar{U}_\infty \equiv 0$ in \mathbb{R}^2 . In other words, $\bar{U}_\infty(t, x) = \bar{U}_\infty(x)$ is independent of t and it obeys $a(x_\infty)\bar{U}_\infty'' + f(x_\infty, \bar{U}_\infty) = 0$ in \mathbb{R} . Notice that $0 \leq \bar{U}_\infty \leq 1$ in \mathbb{R} and $\delta \leq \bar{U}_\infty(0) \leq 1 - \delta$. It then follows from the strong maximum principle that $0 < \bar{U}_\infty < 1$ in \mathbb{R} . Furthermore, it follows from (A1) and (A3) that either \bar{U}_∞ is a constant solution (i.e., $\bar{U}_\infty \equiv b(x_\infty)$), or it is a non-constant periodic solution, or it is a ground state solution decaying to 0 as $x \rightarrow \pm\infty$. Proceeding similarly as in Step 1, we can find a contradiction in each of these cases: namely, for each large $n \in \mathbb{N}$, there exists then a solution \bar{w}_n of (2.41) (with \bar{L}_n and \bar{x}_n instead of L_n and x_n) such that either

$$\mathcal{Z}[\bar{U}_n(0, \cdot) - \bar{w}_n] \geq 3,$$

or

$$\mathcal{Z}[\bar{U}_n(0, \cdot) - \bar{w}_n] = 2 \quad \text{and} \quad \text{SGN}[\bar{U}_n(0, \cdot) - \bar{w}_n] = [-, +, -]$$

(the proof is even simpler than that in Step 1, as the limit function \bar{U}_∞ is stationary; therefore we do not repeat the details). Both situations are impossible, due to Lemma 2.4 (applied with $\bar{U}_n(\cdot, \cdot - \bar{x}_n)$ and $\bar{w}_n(\cdot - \bar{x}_n)$). As a consequence, the proof of Step 2 is complete.

Step 3: (1.21) holds true with $C > 0$ independent of $L \geq L_*$. Assume by contradiction that there does not exist such a bound independent of L . Then, we can find sequences $(L_n)_{n \in \mathbb{N}} \subset [L_*, +\infty)$, $(s_n)_{n \in \mathbb{N}} \subset \mathbb{R}$, $(\alpha_n)_{n \in \mathbb{N}}$, $(\gamma_n)_{n \in \mathbb{N}} \subset [\delta, 1 - \delta]$ and $(x_n)_{n \in \mathbb{N}}$, $(y_n)_{n \in \mathbb{N}} \subset \mathbb{R}$ such that

$$U_{L_n}(s_n, x_n) = \alpha_n \quad \text{and} \quad U_{L_n}(s_n, y_n) = \gamma_n \quad \text{for each } n \in \mathbb{N},$$

and

$$|x_n - y_n| \rightarrow +\infty \quad \text{as } n \rightarrow +\infty.$$

Without loss of generality, for each $n \in \mathbb{N}$, we may assume that $y_n > x_n$, and thanks to (1.5), we may also assume that $0 \leq x_n < L_n$. Moreover, by Theorem 1.1 (ii) (iv) again, we have $L_n \rightarrow +\infty$ as $n \rightarrow +\infty$.

Define now

$$\tilde{U}_n(t, x) := U_{L_n}(t + s_n, x + x_n) \quad \text{for } (t, x) \in \mathbb{R}^2.$$

Clearly, each function $\tilde{U}_n(t, x)$ is an entire solution of (2.34) with $\tilde{U}_n(0, 0) = \alpha_n$ and $\tilde{U}_n(0, z_n) = \gamma_n$, where $0 < z_n = y_n - x_n \rightarrow +\infty$ as $n \rightarrow +\infty$. Up to extraction of some subsequence, we have

$$\frac{x_n}{L_n} \rightarrow x_\infty \in [0, 1], \quad \alpha_n \rightarrow \alpha_\infty \in [\delta, 1 - \delta], \quad \gamma_n \rightarrow \gamma_\infty \in [\delta, 1 - \delta], \quad \text{as } n \rightarrow +\infty,$$

and by standard parabolic estimates, $\tilde{U}_n \rightarrow \tilde{U}_\infty$ in $C_{t;x;loc}^{1;2}(\mathbb{R}^2)$ as $n \rightarrow +\infty$, where \tilde{U}_∞ is a solution of (2.36) with $\tilde{U}_\infty(0, 0) = \alpha_\infty$.

Set $\delta_* := \min\{\delta_0, \delta\}$, where δ_0 is the constant provided by (A2). It follows from Lemma 2.5 that, for each $n \in \mathbb{N}$, (2.23) with $L = L_n$ admits a solution $w_n \in C^2(\mathbb{R})$ such that $w_n(0) = \delta_*$, $\mathcal{Z}[w_n - 1] = 0$, $\mathcal{Z}[w_n] \leq 1$ and $0 < w_n(x) \leq \delta_* e^{-\mu x}$ for all $x \geq 0$, where $\mu > 0$ is independent of $n \in \mathbb{N}$. Notice that for each $n \in \mathbb{N}$, $\tilde{U}_n(t, x)$ is decreasing and continuous in $t \in \mathbb{R}$, with $\tilde{U}_n(-\infty, x) = 0$ and $\tilde{U}_n(+\infty, x) = 1$ locally uniformly in $x \in \mathbb{R}$. This implies in particular that there exists a unique time $\tau_n \in \mathbb{R}$ such that $\tilde{U}_n(\tau_n, z_n) = w_n(z_n)$. Since $w_n(z_n) < \delta_* \leq \delta \leq \gamma_n = \tilde{U}_n(0, z_n)$ for each $n \in \mathbb{N}$, it is clear that $\tau_n < 0$.

Now, we claim that

$$\tau_n \rightarrow -\infty \quad \text{as } n \rightarrow +\infty. \tag{2.43}$$

Suppose to the contrary that $(\tau_n)_{n \in \mathbb{N}}$ converges, up to extraction of a subsequence, to $\tau_\infty \in (-\infty, 0]$. For each $n \in \mathbb{N}$, let us write

$$z_n := z'_n + z''_n, \quad \text{and} \quad V_n(t, x) := \tilde{U}_n(t, x + z_n) \quad \text{in } \mathbb{R}^2,$$

where $z'_n \in L_n \mathbb{N}$ and $z''_n \in [0, L_n)$. It is easily checked that each V_n is an entire solution of

$$\partial_t V_n = \partial_x \left(a \left(\frac{x + x_n + z''_n}{L_n} \right) \partial_x V_n \right) + f \left(\frac{x + x_n + z''_n}{L_n}, V_n \right) \quad \text{in } \mathbb{R}^2,$$

with $V_n(0, 0) = \gamma_n$ and $V_n(\tau_n, 0) = w_n(z_n)$. Since $z_n \rightarrow +\infty$ as $n \rightarrow +\infty$ and since $w_n(x)$ decays to 0 as $x \rightarrow +\infty$ uniformly in $n \in \mathbb{N}$, it follows that $\lim_{n \rightarrow +\infty} V_n(\tau_n, 0) = 0$. Next, up

to extraction of another subsequence, we may assume that $z_n''/L_n \rightarrow z_\infty \in [0, 1]$ as $n \rightarrow +\infty$, and that, by standard parabolic estimates, $V_n \rightarrow V_\infty$ as $n \rightarrow +\infty$ in $C_{t;x;loc}^{1;2}(\mathbb{R}^2)$, where V_∞ is an entire solution of

$$\partial_t V_\infty = a(x_\infty + z_\infty) \partial_{xx} V_\infty + f(x_\infty + z_\infty, V_\infty) \quad \text{in } \mathbb{R}^2.$$

Since $0 \leq V_\infty \leq 1$ and since $V_\infty(\tau_\infty, 0) = \lim_{n \rightarrow +\infty} V_n(\tau_n, 0) = 0$, it follows from the strong maximum principle that $V_\infty \equiv 0$ in \mathbb{R}^2 . This contradicts the fact that $V_\infty(0, 0) = \lim_{n \rightarrow +\infty} \gamma_n = \gamma_\infty \geq \delta$. Therefore, our claim (2.43) is proved.

We are now ready to derive a contradiction. By the proof of Step 1, it follows that the sequence $(m(\{t \in \mathbb{R} : \delta_* \leq \tilde{U}_n(t, 0) \leq 1 - \delta_*\}))_{n \in \mathbb{N}}$ is bounded uniformly in $n \in \mathbb{N}$. This together with $\tilde{U}_n(0, 0) = \alpha_n \in [\delta, 1 - \delta] \subset [\delta_*, 1 - \delta_*]$ and $\partial_t \tilde{U}_n > 0$ yields the existence of $A > 0$ such that $\{t \in \mathbb{R} : \delta_* \leq \tilde{U}_n(t, 0) \leq 1 - \delta_*\} \subset [-A, A]$ for all $n \in \mathbb{N}$. As a consequence, one finds $t_* < 0$ such that

$$\tilde{U}_n(t_*, 0) < \delta_* = w_n(0) \quad \text{for all } n \in \mathbb{N}.$$

Furthermore, by (2.43) and the positivity of $\partial_t \tilde{U}_n$, there exists $n_* \in \mathbb{N}$ sufficiently large such that

$$\tilde{U}_{n_*}(t_*, z_{n_*}) > w_{n_*}(z_{n_*}).$$

Combining the above, we immediately obtain that

$$\mathcal{Z}[\tilde{U}_{n_*}(t_*, \cdot) - w_{n_*}] \geq 1 \quad \text{and} \quad \text{SGN}[\tilde{U}_{n_*}(t_*, \cdot) - w_{n_*}] \triangleright [-, +],$$

which is a contradiction with the conclusions of Lemma 2.6 (applied with $\tilde{U}_{n_*}(\cdot, \cdot - x_{n_*})$ and $w_{n_*}(\cdot - x_{n_*})$).

This ends the proof of Step 3, and the proof of Theorem 1.7 is thus complete. \square

2.4 Proof of Theorem 1.8

We first show statements (i). From Theorem 1.1 (ii), it suffices to prove that for any $\delta \in (0, 1/2]$, the interval

$$E_L(y) := \left\{ \xi \in \mathbb{R} : \delta \leq \phi_L(\xi + \zeta_L(y), y) \leq 1 - \delta \right\}$$

is bounded uniformly with respect to $y \in \mathbb{R}$ and $L \geq L_*$. To do so, observe that the formula $U_L(t, x) = \phi_L(x - c_L t, x/L)$ implies that

$$U_L\left(-\frac{\xi}{c_L} + \frac{Ly - \zeta_L(y)}{c_L}, Ly\right) = \phi_L(\xi + \zeta_L(y), y) \quad \text{for all } \xi \in \mathbb{R} \text{ and } y \in \mathbb{R}. \quad (2.44)$$

Clearly, $U_L((Ly - \zeta_L(y))/c_L, Ly) = 1/2$ for all $y \in \mathbb{R}$, by (1.15). Then, by Theorem 1.7, the following interval

$$\tilde{I}_L(Ly) := \left\{ t \in \mathbb{R} : \delta \leq U_L\left(t + \frac{Ly - \zeta_L(y)}{c_L}, Ly\right) \leq 1 - \delta \right\}$$

is bounded uniformly in $L \geq L_*$ and $y \in \mathbb{R}$. Remember that c_L is bounded in $L \geq L_*$, by Theorem 1.1 (i). This together with (2.44) immediately implies that the interval $E_L(y)$ is bounded uniformly in $y \in \mathbb{R}$ and $L \geq L_*$. The proof of statements (i) of Theorem 1.8 is thus complete.

Next, we show statement (ii). It follows from (2.44) that for any $(\xi, y) \in \mathbb{R} \times \mathbb{R}$ with $\xi \in E_L(y)$, we have $-\xi/c_L \in \tilde{I}_L(Ly)$ and

$$\partial_\xi \phi_L(\xi + \zeta_L(y), y) = -\frac{1}{c_L} \partial_t U_L \left(-\frac{\xi}{c_L} + \frac{Ly - \zeta_L(y)}{c_L}, Ly \right).$$

Then by Theorem 1.7 and the fact that $c_L > 0$ is bounded in $L \geq L_*$, we immediately obtain statement (ii) (with a constant $\beta > 0$ which is in general different from the one appearing in (1.22)).

Finally, we prove statement (iii). Since $\zeta_L(y)$ is continuous in $(L, y) \in [L_*, +\infty) \times \mathbb{R}$, it is clear that $(L, y) \mapsto \zeta_L(y)/L$ is also continuous. It remains to show that this function is bounded. Since $\zeta_L(y)$ is 1-periodic in $y \in \mathbb{R}$, we only need to show that $\zeta_L(y)/L$ is bounded in $(L, y) \in [L_*, +\infty) \times [0, 1]$. It follows from (1.14) and Theorem 1.7 that the set $\{x \in \mathbb{R} : U_L(0, x) = 1/2\}$ is bounded uniformly in $L \geq L_*$, that is, there exists $\bar{L} \geq L_*$ such that

$$\{x \in \mathbb{R} : U_L(0, x) = 1/2\} \subset [-\bar{L}, \bar{L}] \quad \text{for all } L \geq L_*.$$

Remember that $\lim_{x \rightarrow +\infty} U_L(0, x) = 0$ for each $L \geq L_*$. This implies that $U_L(0, L(y+1)) < 1/2$ for all $y \in [0, 1]$ and $L > \bar{L}$. Thanks to (2.44) applied with $\xi := L(y+1) - \zeta_L(y+1)$ and $y+1$ instead of y , it follows that $\phi_L(L(y+1) - \zeta_L(y+1) + \zeta_L(y), y) = \phi_L(L(y+1) - \zeta_L(y+1) + \zeta_L(y), y+1) = U_L(0, L(y+1)) < 1/2$. Hence, the monotonicity of ϕ_L in its first variable implies that

$$L(y+1) - \zeta_L(y+1) > 0 \quad \text{for all } y \in [0, 1] \text{ and } L > \bar{L}.$$

Consequently, by the periodicity of $\zeta_L(y)$, we obtain $\zeta_L(y)/L < y+1$ for all $y \in [0, 1]$ and $L > \bar{L}$.

In a similar way, by using the fact that $\lim_{x \rightarrow -\infty} U_L(0, x) = 1$ and the monotonicity of U_L in its first variable, we can prove that

$$L(y-2) - \zeta_L(y-2) < 0 \quad \text{for all } y \in [0, 1] \text{ and } L > \bar{L},$$

hence $\zeta_L(y)/L > y-2$ for all $y \in [0, 1]$ and $L > \bar{L}$. Combining the above, we immediately obtain that $-2 < \zeta_L(y)/L < 2$ for all $y \in [0, 1]$ and $L > \bar{L}$. This together with the continuity of $\zeta_L(y)/L$ in $[L_*, +\infty) \times [0, 1]$ implies that $\zeta_L(y)/L$ is bounded in $[L_*, +\infty) \times [0, 1]$. The proof of Theorem 1.8 is thus complete. \square

3 Convergence of c_L as $L \rightarrow +\infty$: proof of Theorem 1.2

This section is devoted to the proof of Theorem 1.2. When L is large, we consider the Cauchy problem associated with equation (1.1) under the large-time and large-space scaling, $t \rightarrow Lt$, $x \rightarrow Lx$. More precisely, letting $v(t, x) = u(Lt, Lx)$ with $u(t, x)$ being a solution of (1.1), the rescaled Cauchy problem reads

$$\begin{cases} v_t = \frac{1}{L} (a(x)v_x)_x + Lf(x, v), & t > 0, x \in \mathbb{R}, \\ v(0, x) = g(x), & x \in \mathbb{R}. \end{cases} \quad (3.45)$$

We consider initial conditions $g \in C(\mathbb{R}, [0, 1])$ which are independent of L and front-like, in the sense that

$$\liminf_{x \rightarrow -\infty} (g(x) - b(x)) > 0, \quad \limsup_{x \rightarrow +\infty} (g(x) - b(x)) < 0, \quad (3.46)$$

where $x \mapsto b(x) \in (0, 1)$ is the zero of $f(x, \cdot)$ provided by (A1). For each $L > 0$, denote by $v_L(t, x)$ the solution of the Cauchy problem (3.45). By the strong maximum principle, we have $0 < v_L(t, x) < 1$ for all $t > 0$ and $x \in \mathbb{R}$. Moreover, it is easily seen from (1.4) that for each $L \geq L_*$, with $L_* > 0$ as in Theorem 1.1, $\phi_L(L(x - c_L t), x)$ is a solution of the first equation of (3.45), and it is defined for all $(t, x) \in \mathbb{R}^2$.

As we mentioned in Section 1, we use Gärtner's result [30] to show the convergence of c_L as $L \rightarrow +\infty$. Recall that the function $c : \mathbb{R} \rightarrow \mathbb{R}$, $y \mapsto c(y)$ denotes the unique front speeds of the family of homogeneous equations (1.8). The following lemma is a direct application of [30, Corollary, p. 140] to our one-dimensional spatially periodic equation (3.45) (see also Xin's review paper [56, Theorem 3.7] in the case where g is compactly supported):

Lemma 3.1. [30] *Let $g \in C(\mathbb{R}, [0, 1])$ satisfy (3.46) and let $\Gamma_0 = \{x \in \mathbb{R} : g(x) = b(x)\}$. For any $L > 0$, let $v_L(t, x)$ be the solution of (3.45) with initial condition g (g is independent of L). Then, as $L \rightarrow +\infty$,*

$$\begin{cases} v_L(t, x) \rightarrow 1 & \text{locally uniformly in } \{(t, x) \in (0, +\infty) \times \mathbb{R} : \rho(x, \Gamma_0) < t\}, \\ v_L(t, x) \rightarrow 0 & \text{locally uniformly in } \{(t, x) \in (0, +\infty) \times \mathbb{R} : \rho(x, \Gamma_0) > t\}, \end{cases}$$

where $\rho(x, \Gamma_0) := \inf\{\rho(x, x_2) : x_2 \in \Gamma_0\}$ and $\rho(\cdot, \cdot)$ denotes the signed distance function defined by

$$\rho(x_1, x_2) := \int_{x_2}^{x_1} c^{-1}(y) dy \quad \text{for } x_1 \in \mathbb{R}, x_2 \in \mathbb{R}.$$

Recall that c_* is the constant defined in (1.11). Since the function $y \mapsto c(y)$ is 1-periodic, the above result suggests that as $L \rightarrow +\infty$, the solution $v_L(t, x)$ propagates at a mean speed equal to c_* . This however does not imply directly the convergence of the speeds c_L to c_* , since the profiles of the pulsating fronts, say at time 0, depend on L , unlike the initial condition g in Lemma 3.1. To circumvent this difficulty, we establish in Lemma 3.2 some further comparisons between the solutions $v_L(t, x)$ and the pulsating fronts $\phi_L(L(x - c_L t), x)$ for $L \geq L_*$, which easily lead to the proof of Theorem 1.2, on the basis of the uniform estimates proved in Section 2.⁶ To do so, let δ_0 and γ_0 be the positive constants provided by the assumption (A2). Since $f(\cdot, 0) = f(\cdot, 1) = 0$ in \mathbb{R} , there is $\varepsilon_0 \in (0, \delta_0/2]$ such that

$$\partial_u f(x, u) \leq -\frac{\gamma_0}{2} \quad \text{for all } u \in [-2\varepsilon_0, 2\varepsilon_0] \cup [1 - 2\varepsilon_0, 1 + 2\varepsilon_0] \text{ and } x \in \mathbb{R}. \quad (3.47)$$

We now choose a special initial condition $g \in C(\mathbb{R}, [0, 1])$, independent of $L \geq L_*$, such that

$$\liminf_{x \rightarrow -\infty} g(x) > 1 - \varepsilon_0, \quad \limsup_{x \rightarrow +\infty} g(x) < \varepsilon_0, \quad \text{and} \quad \Gamma_0 = \{0\}. \quad (3.48)$$

Notice that it is for instance possible to choose g in such a way that $g = 1$ in $(-\infty, -A]$ and $g = 0$ in $[A, +\infty)$, for some large A .

Lemma 3.2. *For any $L \geq L_*$, with $L_* > 0$ as in Theorem 1.1, let (ϕ_L, c_L) be a pulsating front of (1.1) and let $v_L(t, x)$ be the solution of (3.45) with the initial condition g satisfying (3.48). Then there exist constants $C_{\pm} \geq 0$ and $K_0 > 0$ (all independent of L) such that*

$$\begin{cases} v_L(t, x) \geq \phi_L(L(x - c_L t + C_-) + K_0 \varepsilon_0, x) - \varepsilon_0 e^{-\gamma_0 L t / 2} \\ v_L(t, x) \leq \phi_L(L(x - c_L t - C_+) - K_0 \varepsilon_0, x) + \varepsilon_0 e^{-\gamma_0 L t / 2} \end{cases} \quad \text{for all } t \geq 0 \text{ and } x \in \mathbb{R}. \quad (3.49)$$

⁶An alternate approach could be to use [30, Theorem 4.1] with L -dependent initial conditions $\phi_L(L \cdot, \cdot)$ combined with Theorem 1.8. We here use Lemma 3.1 and the comparisons established in Lemma 3.2, since the method of proof based on sub- and super-solutions will serve as an archetype to get more involved comparisons in Section 4.2.

Proof. We use a Fife-McLeod [28] type sub- and super-solutions method to show this lemma. We only give the construction of a super-solution, as the analysis for a sub-solution is analogous.

First of all, since the initial function g satisfies (3.48), it follows from Theorem 1.8 (i) and (iii) that there exists a constant $C_+ > 0$ independent of L such that, for all $L \geq L_*$ and $x \in \mathbb{R}$,

$$\phi_L(L(x - C_+), x) + \varepsilon_0 = \phi_L\left(L\left(x - C_+ - \frac{\zeta_L(x)}{L}\right) + \zeta_L(x), x\right) + \varepsilon_0 \geq g(x). \quad (3.50)$$

For any $L \geq L_*$, we then set

$$\bar{v}_L(t, x) := \phi_L(L(x - c_L t - C_+) + \eta_L(t), x) + q_L(t) \quad \text{for } t \geq 0 \text{ and } x \in \mathbb{R},$$

where $t \mapsto q_L(t)$ and $t \mapsto \eta_L(t)$ are some $C^1([0, +\infty))$ functions satisfying

$$\begin{cases} q_L(0) = \varepsilon_0, & q'_L(t) < 0 < q_L(t) & \text{for all } t \geq 0, \\ \eta_L(0) = 0, & \eta'_L(t) < 0 & \text{for all } t \geq 0. \end{cases} \quad (3.51)$$

By choosing suitable functions η_L and q_L later, we will show that \bar{v}_L is a super-solution of (3.45).

Notice first that (3.50)-(3.51) imply that $\bar{v}_L(0, \cdot) \geq g = v_L(0, \cdot)$ in \mathbb{R} for each $L \geq L_*$. Now, for $(t, x) \in (0, +\infty) \times \mathbb{R}$, we define

$$N_L(t, x) := \partial_t \bar{v}_L(t, x) - \frac{1}{L} \partial_x (a(x) \partial_x \bar{v}_L(t, x)) - Lf(x, \bar{v}_L(t, x)).$$

Since $(t, x) \mapsto \phi_L(L(x - c_L t - C_+), x)$ is an entire solution of the first equation of (3.45), a straightforward calculation gives, for all $(t, x) \in (0, +\infty) \times \mathbb{R}$,

$$\begin{aligned} N_L(t, x) &= \eta'_L(t) \partial_\xi \phi_L(L(x - c_L t - C_+) + \eta_L(t), x) + q'_L(t) + Lf(x, \phi_L(L(x - c_L t - C_+) + \eta_L(t), x)) \\ &\quad - Lf(x, \bar{v}_L(t, x)). \end{aligned}$$

Since $\partial_\xi \phi_L < 0$ in \mathbb{R}^2 , we get $\eta'_L(t) \partial_\xi \phi_L(L(x - c_L t - C_+) + \eta_L(t), x) > 0$ and $0 < q_L(t) \leq \varepsilon_0$ for all $t \geq 0$ and $x \in \mathbb{R}$, provided q_L and η_L fulfill (3.51). Then, on the one hand, for any pair $(t, x) \in (0, +\infty) \times \mathbb{R}$ such that $\phi_L(L(x - c_L t - C_+) + \eta_L(t), x) \in (0, \varepsilon_0] \times [1 - \varepsilon_0, 1)$, it follows from (3.47) that

$$N_L(t, x) \geq q'_L(t) + Lf(x, \phi_L(L(x - c_L t - C_+) + \eta_L(t), x)) - Lf(x, \bar{v}_L(t, x)) \geq q'_L(t) + \frac{\gamma_0 L}{2} q_L(t).$$

On the other hand, by Theorem 1.8 (ii), there is $\beta > 0$ (independent of $L \geq L_*$) such that $\partial_\xi \phi_L(\xi, y) \leq -\beta$ for all $L \geq L_*$ and for all $(\xi, y) \in \mathbb{R}^2$ with $\varepsilon_0 \leq \phi_L(\xi, y) \leq 1 - \varepsilon_0$. It follows that, if $\phi_L(L(x - c_L t - C_+) + \eta_L(t), x) \in (\varepsilon_0, 1 - \varepsilon_0)$, then

$$\begin{aligned} N_L(t, x) &\geq -\eta'_L(t) \beta + q'_L(t) + Lf(x, \phi_L(L(x - c_L t - C_+) + \eta_L(t), x)) - Lf(x, \bar{v}_L(t, x)) \\ &\geq -\eta'_L(t) \beta + q'_L(t) - LC_1 q_L(t), \end{aligned}$$

where $C_1 := \|\partial_u f\|_{L^\infty(\mathbb{R} \times \mathbb{R})}$. Let us now choose $q_L(t)$ and $\eta_L(t)$ such that

$$\begin{cases} q_L(0) = \varepsilon_0, & q'_L(t) + \frac{\gamma_0 L}{2} q_L(t) = 0 & \text{for } t \geq 0, \\ \eta_L(0) = 0, & -\beta \eta'_L(t) + q'_L(t) - LC_1 q_L(t) = 0 & \text{for } t \geq 0. \end{cases}$$

Namely, we set

$$q_L(t) = \varepsilon_0 e^{-\gamma_0 L t / 2} \quad \text{and} \quad \eta_L(t) = -\frac{\varepsilon_0(\gamma_0 + 2C_1)}{\beta \gamma_0} (1 - e^{-\gamma_0 L t / 2}) \quad \text{for } t \geq 0.$$

These functions q_L and η_L satisfy (3.51). Consequently, $N_L(t, x) \geq 0$ for all $(t, x) \in (0, +\infty) \times \mathbb{R}$.

Finally, the comparison principle implies that, for any $L \geq L_*$, $\bar{v}_L(t, x) \geq v_L(t, x)$ for all $t \geq 0$ and $x \in \mathbb{R}$. Taking $K_0 = (\gamma_0 + 2C_1)/(\beta\gamma_0)$ and using the fact that $\phi_L(\xi, y)$ is decreasing in ξ , we obtain the second inequality of (3.49). As we have mentioned above, similar arguments imply the first one. The proof of Lemma 3.2 is thus complete. \square

We are now ready to complete the

Proof of Theorem 1.2. Assume by contradiction that c_L dose not converge to c_* as $L \rightarrow +\infty$. Since the family $(c_L)_{L \geq L_*} \subset (0, +\infty)$ is bounded by Theorem 1.1 (i), one finds a sequence $(L_n)_{n \in \mathbb{N}} \subset [L_*, +\infty)$ with $L_n \rightarrow +\infty$ as $n \rightarrow +\infty$ and a real number $c_\infty \geq 0$ such that

$$c_{L_n} \rightarrow c_\infty \text{ as } n \rightarrow +\infty, \quad \text{and} \quad c_\infty \neq c_*.$$

Without loss of generality, we assume that $c_\infty < c_*$ (as sketched below, the case where $c_\infty > c_*$ can be treated similarly). Then we have $c_{L_n} < (c_\infty + c_*)/2 < c_*$ for all large $n \in \mathbb{N}$. Now, we can choose a large time $t_* > 0$ and a positive integer x_* (both independently of n) such that

$$c_{L_n} t_* + C_+ + 1 < x_* < c_* t_* \quad \text{for all large } n \in \mathbb{N}, \quad (3.52)$$

where $C_+ \geq 0$ is the constant given by Lemma 3.2. Since each function $\phi_{L_n}(\xi, y)$ is decreasing in ξ , it follows from the notations of Lemma 3.2 and the second inequality of (3.49) that

$$\begin{aligned} 0 < v_{L_n}(t_*, x_*) &\leq \phi_{L_n}(L_n(x_* - c_{L_n} t_* - C_+) - K_0 \varepsilon_0, x_*) + \varepsilon_0 e^{-\gamma_0 L_n t_*/2} \\ &\leq \phi_{L_n}(L_n - K_0 \varepsilon_0, x_*) + \varepsilon_0 e^{-\gamma_0 L_n t_*/2} \end{aligned} \quad (3.53)$$

for all large $n \in \mathbb{N}$. Moreover, since $x_* \in \mathbb{N}$ and each function ζ_{L_n} given by (1.15) is 1-periodic, we have $\zeta_{L_n}(x_*) = \zeta_{L_n}(0) = 0$. Therefore, passing to the limit as $n \rightarrow +\infty$ in (3.53), we see from Theorem 1.8 that $\phi_{L_n}(L_n - K_0 \varepsilon_0, x_*) \rightarrow 0$ as $n \rightarrow +\infty$, hence $v_{L_n}(t_*, x_*) \rightarrow 0$ as $n \rightarrow +\infty$.

On the other hand, we had chosen g such that $\Gamma_0 = \{x \in \mathbb{R} : g(x) = b(x)\} = \{0\}$. Since x_* is a positive integer and since the function $y \mapsto c(y)$ is positive and 1-periodic, we have $\rho(x_*, \Gamma_0) = x_*/c_*$, whence $\rho(x_*, \Gamma_0) < t_*$ due to (3.52). It then follows from Lemma 3.1 that $v_{L_n}(t_*, x_*) \rightarrow 1$ as $n \rightarrow +\infty$, yielding a contradiction. Therefore, the case where $c_\infty < c_*$ is ruled out.

In the case where $c_\infty > c_*$, one can derive a similar contradiction by choosing a large time $t^* > 0$ and a positive integer x^* (both independently of n) such that

$$c_{L_n} t^* - C_- - 1 > x^* > c_* t^* \quad \text{for all large } n \in \mathbb{N}.$$

As a conclusion, $c_L \rightarrow c_*$ as $L \rightarrow +\infty$, and the proof of Theorem 1.2 is thus complete. \square

4 Convergence of the front profiles: proof of Theorem 1.6

This section is devoted to the proof of Theorem 1.6, that is, the convergence of the front profiles ϕ_L as $L \rightarrow +\infty$, under the assumptions (A1)-(A4). The key step is to determine the asymptotic behavior as $L \rightarrow +\infty$ of the solutions of Cauchy problem of the following equation:

$$\partial_t z_L = a \partial_{xx} z_L + f\left(y + \frac{x}{L}, z_L\right) \quad \text{for } t > 0, \quad x \in \mathbb{R}, \quad (4.54)$$

where $y \in \mathbb{R}$ is arbitrary. We point out that a is here a positive constant, thanks to (A4). We recall that, for each $y \in \mathbb{R}$, the couple $(\psi(\cdot, y), c(y))$ denotes the front profile and front speed

of the homogeneous equation (1.9)-(1.10). We first present some properties on $(\psi(\cdot, y), c(y))$ in Section 4.1. In Section 4.2, we choose the initial condition of (4.54) sufficiently close to the homogeneous front $\psi(\cdot, y)$, and then construct a pair of sub- and super-solutions of (4.54). This will ensure that, at a certain time t , $z_L(t, \cdot + L)$ is still close to $\psi(\cdot, y)$ provided that L is sufficiently large (see Proposition 4.7 in Section 4.3). We complete the proof of Theorem 1.6 in Section 4.4.

4.1 Preliminaries on homogeneous fronts

In this subsection, we present some properties regarding the homogeneous traveling front $(\psi(\cdot, y), c(y))$ of (1.9)-(1.10) with respect to the parameter y . These properties will be used in our construction of sub- and super-solutions later and are also of interest in themselves.

Proposition 4.1. *Let (A1)-(A3) hold. Then, the following statements hold true:*

- (i) *the function $y \mapsto c(y)$ is positive, 1-periodic, and the function $(\xi, y) \mapsto \psi(\xi, y) \in (0, 1)$ is 1-periodic in $y \in \mathbb{R}$ and decreasing in $\xi \in \mathbb{R}$;*
- (ii) *$\lim_{\xi \rightarrow +\infty} \psi(\xi, y) = 0$ and $\lim_{\xi \rightarrow -\infty} \psi(\xi, y) = 1$ uniformly in $y \in \mathbb{R}$; furthermore, there exist positive constants μ_1, μ_2, M, C_1 and C_2 (all are independent of y) such that*

$$\begin{cases} 0 < \psi(\xi, y) \leq C_1 e^{-\mu_1 \xi} & \text{for all } \xi \geq M \text{ and } y \in \mathbb{R}, \\ 0 < 1 - \psi(\xi, y) \leq C_2 e^{\mu_2 \xi} & \text{for all } \xi \leq -M \text{ and } y \in \mathbb{R}; \end{cases} \quad (4.55)$$

- (iii) *for any $\delta \in (0, 1/2]$, there exists $\gamma = \gamma(\delta) > 0$ (independent of y) such that*

$$\partial_\xi \psi(\xi, y) \leq -\gamma \text{ for all } (\xi, y) \in \mathbb{R}^2 \text{ such that } \delta \leq \psi(\xi, y) \leq 1 - \delta;$$

- (iv) *the function $y \mapsto c(y)$ is of class $C^1(\mathbb{R})$, and the function $(\xi, y) \mapsto \psi(\xi, y)$ is of class $C_{\xi; y}^{2;1}(\mathbb{R}^2)$, and satisfies*

$$\sup_{\xi \in \mathbb{R}, y \in \mathbb{R}} |\partial_y \psi(\xi, y)| < +\infty. \quad (4.56)$$

The proof of Proposition 4.1 is quite lengthy and is therefore postponed to the appendix in Section 5. Let us emphasize that this proposition holds without the extra assumption (A4).

4.2 Sub- and super-solutions

In the remaining part of Section 4, we always assume that (A1)-(A4) hold. To present our sub- and super-solutions, we need some notations. Since the function $z \mapsto c(z)$ is 1-periodic and of class $C^1(\mathbb{R})$ by Proposition 4.1 (iv), one infers that, for each $y \in \mathbb{R}$ and $L > 0$, the following ODE problem

$$X'_{y,L}(t) = c\left(y + \frac{X_{y,L}(t)}{L}\right) \text{ for } t \geq 0, \quad X_{y,L}(0) = 0, \quad (4.57)$$

admits a unique solution $X_{y,L} : [0, +\infty) \rightarrow \mathbb{R}$, with $X_{y,L} = X_{y+1,L}$ in $[0, +\infty)$ for each $y \in \mathbb{R}$ and $L > 0$. Furthermore, each function $X_{y,L}$ is increasing in $[0, +\infty)$ and $X_{y,L}(t) \rightarrow +\infty$ as $t \rightarrow +\infty$, since $\min_{\mathbb{R}} c > 0$. Let then $T_L > 0$ be the unique time such that

$$X_{y,L}(T_L) = L.$$

Since the function c is 1-periodic, it is then checked by integrating the function $t \mapsto X'_{y,L}(t)/c(y + X_{y,L}(t)/L) = 1$ over $[0, T_L]$ that T_L does not depend on y and is equal to

$$T_L = \frac{L}{c_*} \quad (4.58)$$

where $c_* > 0$ is the constant defined by (1.11). In particular, $T_L \rightarrow +\infty$ as $L \rightarrow +\infty$ (uniformly in $y \in \mathbb{R}$). Notice also that, again by 1-periodicity of the function c , one has

$$X_{y,L}(kT_L) = kL \quad \text{for all } k \in \mathbb{N}. \quad (4.59)$$

Our super-solution is stated in the following lemma.

Lemma 4.2. *There exists $\varepsilon_0 \in (0, \delta_0)$, with $\delta_0 \in (0, 1/2)$ as in assumption (A2), such that, for every $\varepsilon \in (0, \varepsilon_0]$, there exists $L_{1,\varepsilon} > 0$ such that, for every $y \in \mathbb{R}$ and $L \geq L_{1,\varepsilon}$, the function $v_{\varepsilon,y,L}^+ : [0, +\infty) \times \mathbb{R} \rightarrow \mathbb{R}$ defined by*

$$v_{\varepsilon,y,L}^+(t, x) := \psi \left(x - X_{y,L}(t) + \eta_{\varepsilon,L}(t), y + \frac{X_{y,L}(t)}{L} \right) + q_{\varepsilon,L}(t)$$

is a super-solution of (4.54) for $t \geq 0$ and $x \in \mathbb{R}$, where $q_{\varepsilon,L}$ and $\eta_{\varepsilon,L}$ are $C^1([0, +\infty))$ functions satisfying

$$\begin{cases} q_{\varepsilon,L}(0) = \varepsilon, & q'_{\varepsilon,L}(t) < 0 < q_{\varepsilon,L}(t) \quad \text{for all } t \geq 0, \\ \eta_{\varepsilon,L}(0) = 0, & \eta'_{\varepsilon,L}(t) < 0 \quad \text{for all } t \geq 0. \end{cases} \quad (4.60)$$

Proof. First of all, by the assumptions (A1)-(A2) and the $C^1(\mathbb{R} \times [0, 1])$ smoothness of f and its periodicity in x , together with (1.3), there exist $\delta_1 \in (0, 1/2)$ and $\gamma_1 > 0$ such that

$$\partial_u f(x, u) \leq -\gamma_1 \quad \text{for all } x \in \mathbb{R} \text{ and } u \in [-\delta_1, \delta_1] \cap [1 - \delta_1, 1 + \delta_1]. \quad (4.61)$$

Without loss of generality, one can assume that $\delta_1 \leq \delta'_0$, with $\delta'_0 > 0$ as in assumption (A4). Moreover, by Proposition 4.1 (iv) and the boundedness of the function $z \mapsto c(z)$, there is a constant $C_1 > 0$ such that, for all $y \in \mathbb{R}$ and $L > 0$,

$$\left| \frac{X'_{y,L}(t)}{L} \partial_y \psi(\xi, z) \right| \leq \frac{C_1}{L} \quad \text{for all } t \geq 0 \text{ and } (\xi, z) \in \mathbb{R} \times \mathbb{R}. \quad (4.62)$$

It further follows from Proposition 4.1 (i)-(ii) that there exist $M_1 > 0$ and

$$\varepsilon_0 \in (0, \min\{\delta_1/2, \delta_0\}) \subset (0, 1/4),$$

with $\delta_0 \in (0, 1/2)$ as in assumption (A2), such that, for all $(\xi, z) \in \mathbb{R} \times \mathbb{R}$,

$$\begin{cases} 0 < \psi(\xi, z) \leq \frac{\delta_1}{2} & \text{if } \xi \geq M_1, \\ 1 - \frac{\delta_1}{2} \leq \psi(\xi, z) < 1 & \text{if } \xi \leq -M_1, \\ 2\varepsilon_0 \leq \psi(\xi, z) \leq 1 - 2\varepsilon_0 & \text{if } -M_1 < \xi < M_1, \end{cases}$$

where the last inequality is actually a consequence of the continuity of $\psi : \mathbb{R}^2 \rightarrow (0, 1)$ and its periodicity in the second variable.

In the arguments below, $\varepsilon \in (0, \varepsilon_0]$ is arbitrary. To show Lemma 4.2, it suffices to find suitable $C^1([0, +\infty))$ functions $q_{\varepsilon, L}$ and $\eta_{\varepsilon, L}$ satisfying (4.60) such that, if L is sufficiently large and independently of y , there holds

$$N(t, x) := \partial_t v_{\varepsilon, y, L}^+(t, x) - a \partial_{xx} v_{\varepsilon, y, L}^+(t, x) - f\left(y + \frac{x}{L}, v_{\varepsilon, y, L}^+(t, x)\right) \geq 0$$

for all $(t, x) \in [0, +\infty) \times \mathbb{R}$. Since ψ is a solution of (1.9) and since here a is constant by assumption (A4), it is straightforward to check that, for any $y \in \mathbb{R}$, $L > 0$ and any $C^1([0, +\infty))$ functions $q_{\varepsilon, L}$ and $\eta_{\varepsilon, L}$, one has

$$N(t, x) = \frac{X'_{y, L}(t)}{L} \partial_y \psi + \eta'_{\varepsilon, L}(t) \partial_\xi \psi + q'_{\varepsilon, L}(t) + f\left(y + \frac{X_{y, L}(t)}{L}, \psi\right) - f\left(y + \frac{x}{L}, v_{\varepsilon, y, L}^+(t, x)\right)$$

for all $(t, x) \in [0, +\infty) \times \mathbb{R}$, where $\psi = \psi(x - X_{y, L}(t) + \eta_{\varepsilon, L}(t), y + X_{y, L}(t)/L)$, and $\partial_\xi \psi$, $\partial_y \psi$ stand for the partial derivatives of ψ with respect to the first variable and the second variable, respectively, evaluated at the same point $(x - X_{y, L}(t) + \eta_{\varepsilon, L}(t), y + X_{y, L}(t)/L)$. We complete the proof by considering three cases: (a) $x - X_{y, L}(t) + \eta_{\varepsilon, L}(t) \geq M_1$, (b) $x - X_{y, L}(t) + \eta_{\varepsilon, L}(t) \leq -M_1$, (c) $-M_1 < x - X_{y, L}(t) + \eta_{\varepsilon, L}(t) < M_1$.

In case (a), with $q_{\varepsilon, L}$ required to satisfy (4.60), we have

$$0 < \psi\left(x - X_{y, L}(t) + \eta_{\varepsilon, L}(t), y + \frac{X_{y, L}(t)}{L}\right) < v_{\varepsilon, y, L}^+(t, x) \leq \frac{\delta_1}{2} + \varepsilon \leq \frac{\delta_1}{2} + \varepsilon_0 \leq \delta_1.$$

Since ψ is decreasing in its first variable and since $\eta_{\varepsilon, L}$ is required to satisfy (4.60), we have

$$\eta'_{\varepsilon, L}(t) \partial_\xi \psi\left(x - X_{y, L}(t) + \eta_{\varepsilon, L}(t), y + \frac{X_{y, L}(t)}{L}\right) \geq 0.$$

Remember that $f(x, u)$ is independent of x for $u \in [0, \delta_1] \subset [0, \delta'_0]$ by assumption (A4). It then follows from (4.61)-(4.62) that

$$\begin{aligned} N(t, x) &\geq -\frac{C_1}{L} + q'_{\varepsilon, L}(t) + f\left(0, \psi\left(x - X_{y, L}(t) + \eta_{\varepsilon, L}(t), y + \frac{X_{y, L}(t)}{L}\right)\right) - f(0, v_{\varepsilon, y, L}^+(t, x)) \\ &\geq -\frac{C_1}{L} + q'_{\varepsilon, L}(t) + \gamma_1 q_{\varepsilon, L}(t). \end{aligned}$$

Let us choose

$$L_{1, \varepsilon} = \frac{2C_1}{\gamma_1 \varepsilon} > 0 \tag{4.63}$$

(notice that $L_{1, \varepsilon}$ is independent of $y \in \mathbb{R}$) and the function $q_{\varepsilon, L}$ such that

$$q_{\varepsilon, L}(0) = \varepsilon \quad \text{and} \quad -\frac{C_1}{L} + q'_{\varepsilon, L}(t) + \gamma_1 q_{\varepsilon, L}(t) = 0 \quad \text{for } t \geq 0, \tag{4.64}$$

namely,

$$q_{\varepsilon, L}(t) = \frac{C_1}{L\gamma_1} + \left(\varepsilon - \frac{C_1}{L\gamma_1}\right) e^{-\gamma_1 t} \quad \text{for } t \geq 0. \tag{4.65}$$

It is clear that, for any $L \geq L_{1, \varepsilon}$ and for any $y \in \mathbb{R}$, the function $q_{\varepsilon, L}$ satisfies (4.60) and $N(t, x) \geq 0$ for all $(t, x) \in [0, +\infty) \times \mathbb{R}$ such that $x - X_{y, L}(t) + \eta_{\varepsilon, L}(t) \geq M_1$, provided $\eta_{\varepsilon, L}$ satisfies (4.60) too.

Proceeding similarly as above, we can conclude that, for any $L \geq L_{1,\varepsilon}$ and $y \in \mathbb{R}$, $N(t, x) \geq 0$ for all $(t, x) \in [0, +\infty) \times \mathbb{R}$ such that $x - X_{y,L}(t) + \eta_{\varepsilon,L}(t) \leq -M_1$, as soon as $\eta_{\varepsilon,L}$ satisfies (4.60).

It remains to find a suitable function $\eta_{\varepsilon,L}$ such that $N(t, x) \geq 0$ in case (c), i.e., $-M_1 < x - X_{y,L}(t) + \eta_{\varepsilon,L}(t) < M_1$. In this case, we have

$$2\varepsilon_0 \leq \psi \left(x - X_{y,L}(t) + \eta_{\varepsilon,L}(t), y + \frac{X_{y,L}(t)}{L} \right) \leq 1 - 2\varepsilon_0$$

and $2\varepsilon_0 \leq v_{\varepsilon,y,L}^+(t, x) \leq 1 - \varepsilon_0$ (since $0 < q_{\varepsilon,L}(t) \leq \varepsilon \leq \varepsilon_0$). It then follows from Proposition 4.1 (iii) that there exists $\beta_1 > 0$ (independent of ε, y, L, t and x) such that $\partial_\xi \psi(x - X_{y,L}(t) + \eta_{\varepsilon,L}(t), y + X_{y,L}(t)/L) \leq -\beta_1$. Noticing that the function f is of class $C^1(\mathbb{R} \times [0, 1])$ and periodic in x , one finds some constants $K_1 > 0$ and $K_2 > 0$ (independent of ε, y, L, t and x) such that

$$\begin{cases} |f(x_1, u) - f(x_2, u)| \leq K_1|x_1 - x_2| & \text{for all } x_1, x_2 \in \mathbb{R} \text{ and } u \in [0, 1], \\ |f(z, u_1) - f(z, u_2)| \leq K_2|u_1 - u_2| & \text{for all } z \in \mathbb{R} \text{ and } u_1, u_2 \in [0, 1]. \end{cases}$$

Therefore, with $\eta_{\varepsilon,L}$ required to satisfy (4.60) (then being nonpositive in $[0, +\infty)$), we have

$$\begin{aligned} N(t, x) &\geq -\frac{C_1}{L} - \beta_1 \eta'_{\varepsilon,L}(t) + q'_{\varepsilon,L}(t) + f \left(y + \frac{X_{y,L}(t)}{L}, \psi \left(x - X_{y,L}(t) + \eta_{\varepsilon,L}(t), y + \frac{X_{y,L}(t)}{L} \right) \right) \\ &\quad - f \left(y + \frac{x}{L}, v_{\varepsilon,y,L}^+(t, x) \right) \\ &\geq -\frac{C_1}{L} - \beta_1 \eta'_{\varepsilon,L}(t) + q'_{\varepsilon,L}(t) - K_1 \frac{|X_{y,L}(t) - x|}{L} - K_2 q_{\varepsilon,L}(t), \\ &\geq -\frac{C_1}{L} - \beta_1 \eta'_{\varepsilon,L}(t) + q'_{\varepsilon,L}(t) - K_1 \frac{M_1 - \eta_{\varepsilon,L}(t)}{L} - K_2 q_{\varepsilon,L}(t). \end{aligned}$$

Then, with $q_{\varepsilon,L}(t)$ given by (4.64)-(4.65), by choosing $\eta_{\varepsilon,L}$ such that

$$\eta_{\varepsilon,L}(0) = 0 \quad \text{and} \quad -\beta_1 \eta'_{\varepsilon,L}(t) + \frac{K_1}{L} \eta_{\varepsilon,L}(t) - (\gamma_1 + K_2) q_{\varepsilon,L}(t) = \frac{K_1 M_1}{L} \quad \text{for } t \geq 0,$$

we have $N(t, x) \geq 0$ for $(t, x) \in [0, +\infty) \times \mathbb{R}$ such that $-M_1 < x - X_{y,L}(t) + \eta_{\varepsilon,L}(t) < M_1$. It is straightforward to compute that

$$\eta_{\varepsilon,L}(t) = -\frac{\gamma_1 + K_2}{\beta_1} \left[(C_2 + C_{3,\varepsilon,L}) e^{K_1 t / (L\beta_1)} - C_{3,\varepsilon,L} e^{-\gamma_1 t} - C_2 \right] \quad \text{for } t \geq 0, \quad (4.66)$$

where

$$C_2 = \frac{C_1 \beta_1}{\gamma_1 K_1} + \frac{M_1 \beta_1}{K_2 + \gamma_1} \quad \text{and} \quad C_{3,\varepsilon,L} = \frac{\varepsilon - C_1 / (L\gamma_1)}{\gamma_1 + K_1 / (L\beta_1)}.$$

It is clear that $C_2 > 0$ and $C_{3,\varepsilon,L} > 0$ if $L \geq L_{1,\varepsilon}$, whence the function $\eta_{\varepsilon,L}$ given in (4.66) satisfies (4.60).

Combining the above, we can conclude that for any $\varepsilon \in (0, \varepsilon_0]$, $y \in \mathbb{R}$ and $L \geq L_{1,\varepsilon}$, one has $N(t, x) \geq 0$ for all $(t, x) \in [0, +\infty) \times \mathbb{R}$. This ends the proof of Lemma 4.2. \square

The following lemma gives the sub-solution of problem (4.54).

Lemma 4.3. *Let $\varepsilon_0 \in (0, \delta_0)$ be given by Lemma 4.2 and, for any $\varepsilon \in (0, \varepsilon_0]$, let $L_{1,\varepsilon} > 0$ be given by (4.63) as in Lemma 4.2. Then, for every $\varepsilon \in (0, \varepsilon_0]$, $y \in \mathbb{R}$ and $L \geq L_{1,\varepsilon}$, the function $v_{\varepsilon,y,L}^- : [0, +\infty) \times \mathbb{R} \rightarrow \mathbb{R}$ defined by*

$$v_{\varepsilon,y,L}^-(t, x) := \psi \left(x - X_{y,L}(t) - \eta_{\varepsilon,L}(t), y + \frac{X_{y,L}(t)}{L} \right) - q_{\varepsilon,L}(t)$$

is a sub-solution of (4.54) for $t \geq 0$ and $x \in \mathbb{R}$, where $q_{\varepsilon,L}$ and $\eta_{\varepsilon,L}$ are $C^1([0, +\infty))$ functions satisfying (4.60), given by (4.65) and (4.66), respectively.

Proof. The proof is analogous to that of Lemma 4.2; therefore, we omit the details. \square

Before going further on, we give two remarks about the functions $q_{\varepsilon,L}$ and $\eta_{\varepsilon,L}$, which will be useful in the proof of our main result later.

Remark 4.4. *For any $\varepsilon \in (0, \varepsilon_0]$ and $L \geq L_{1,\varepsilon} = 2C_1/(\gamma_1\varepsilon)$, let $T_L = L/c_* > 0$ be the time provided by (4.58). For any given $k \in \mathbb{N}$ and $\tau > 0$, the function $\eta_{\varepsilon,L}$ given by (4.66) is bounded in $[0, kT_L + \tau]$ uniformly with respect to $\varepsilon \in (0, \varepsilon_0]$ and $L \geq L_{1,\varepsilon}$. More precisely, we can find some constants $0 < A_1 < A_2$ (independent of ε , L , k and τ) such that*

$$A_1 - A_2 e^{K_1 k / (c_* \beta_1) + K_1 \tau \gamma_1 \varepsilon_0 / (2C_1 \beta_1)} \leq A_1 - A_2 e^{K_1 k / (c_* \beta_1) + K_1 \tau / (L \beta_1)} \leq \eta_{\varepsilon,L}(t) \leq 0$$

for all $0 \leq t \leq kT_L + \tau$.

Remark 4.5. *We also point out that, for any $\varepsilon \in (0, \varepsilon_0]$, the functions $t \mapsto q_{\varepsilon,L}(t)$ and $t \mapsto \eta_{\varepsilon,L}(t)$ converge locally uniformly in $t \geq 0$ as $L \rightarrow +\infty$. More precisely, there holds*

$$\lim_{L \rightarrow +\infty} q_{\varepsilon,L}(t) = \varepsilon e^{-\gamma_1 t} \quad \text{and} \quad \lim_{L \rightarrow +\infty} \eta_{\varepsilon,L}(t) = -\varepsilon \frac{\gamma_1 + K_1}{\gamma_1 \beta_1} (1 - e^{-\gamma_1 t}),$$

locally uniformly in $t \in [0, +\infty)$.

Now, for any $\varepsilon \in (0, \varepsilon_0]$, with $\varepsilon_0 \in (0, \delta_0)$ provided by Lemmas 4.2-4.3, we consider any family of continuous functions $v_{\varepsilon,y}^0 : \mathbb{R} \rightarrow [0, 1]$ such that

$$v_{\varepsilon,y+1}^0 \equiv v_{\varepsilon,y}^0 \text{ in } \mathbb{R}, \quad v_{\varepsilon,y}^0(0) = \frac{1}{2} \quad \text{and} \quad \|v_{\varepsilon,y}^0 - \psi(\cdot, y)\|_{L^\infty(\mathbb{R})} \leq \varepsilon \text{ for all } y \in \mathbb{R}.^7 \quad (4.67)$$

Then, for any $y \in \mathbb{R}$ and $L > 0$, we denote by $v_{\varepsilon,y,L} : [0, +\infty) \times \mathbb{R} \rightarrow [0, 1]$ the solution of (4.54) with the initial condition $v_{\varepsilon,y}^0$. Since the function f is 1-periodic in its first variable, it follows that the functions $v_{\varepsilon,y,L}$ are 1-periodic with respect to the parameter $y \in \mathbb{R}$. It is also easily seen that

$$\max \{v_{\varepsilon,y,L}^-(0, x), 0\} \leq v_{\varepsilon,y}^0(x) \leq \min \{v_{\varepsilon,y,L}^+(0, x), 1\} \quad \text{for all } x \in \mathbb{R}, \quad (4.68)$$

with $v_{\varepsilon,y,L}^\pm$ are in Lemmas 4.2-4.3. Therefore, the next lemma is an immediate consequence of the comparison principle.

Lemma 4.6. *For any $\varepsilon \in (0, \varepsilon_0]$, with $\varepsilon_0 \in (0, \delta_0)$ as in Lemmas 4.2-4.3, any $y \in \mathbb{R}$, and any $L \geq L_{1,\varepsilon}$ with $L_{1,\varepsilon} > 0$ as in (4.63), there holds*

$$\max \{v_{\varepsilon,y,L}^-(t, x), 0\} \leq v_{\varepsilon,y,L}(t, x) \leq \min \{v_{\varepsilon,y,L}^+(t, x), 1\} \quad \text{for all } t \geq 0 \text{ and } x \in \mathbb{R}.$$

⁷A typical example is given by: $v_{\varepsilon,y}^0 = \psi(\cdot, y)$ for $y \in \mathbb{R}$, but other functions $v_{\varepsilon,y}^0$ will be used in the proof of Theorem 1.6 in Section 4.4.

4.3 Asymptotic behavior of the Cauchy problem as $L \rightarrow +\infty$

In this subsection and the following one, we always let $\varepsilon_0 \in (0, \delta_0)$ be provided by Lemmas 4.2-4.3. Now, for any $\varepsilon \in (0, \varepsilon_0]$, we consider the asymptotic behavior as $L \rightarrow +\infty$ of the solutions $v_{\varepsilon,y,L}$ of (4.54) with initial conditions $v_{\varepsilon,y}^0$ satisfying (4.67), at a time $\tilde{T}_{\varepsilon,y,L}$ defined by

$$\tilde{T}_{\varepsilon,y,L} := \inf \left\{ t > 0 : v_{\varepsilon,y,L}(t, L) = \frac{1}{2} \right\}. \quad (4.69)$$

Notice immediately that, since the functions $v_{\varepsilon,y,L}$ are 1-periodic with respect to y , so are the quantities $\tilde{T}_{\varepsilon,y,L}$'s. Furthermore, for each $y \in \mathbb{R}$ and $L \geq L_{1,\varepsilon}$, one has $v_{\varepsilon,y,L}(0, L) \leq v_{\varepsilon,y,L}^+(0, L) = \psi(L, y) + \varepsilon$ by (4.60) and (4.68), hence

$$\limsup_{L \rightarrow +\infty} \left(\sup_{y \in \mathbb{R}} v_{\varepsilon,y,L}(0, L) \right) \leq \varepsilon \leq \varepsilon_0 < \delta_0 < \frac{1}{2}$$

since $\psi(+\infty, y) = 0$ uniformly in $y \in \mathbb{R}$ by Proposition 4.1 (ii). Therefore, $\tilde{T}_{\varepsilon,y,L} > 0$ for all L large enough, uniformly in $y \in \mathbb{R}$. On the other hand, for each $y \in \mathbb{R}$ and $L \geq L_{1,\varepsilon}$, with T_L as in (4.58), one has $v_{\varepsilon,y,L}(2T_L, L) \geq v_{\varepsilon,y,L}^-(2T_L, L) \geq \psi(-L - \eta_{\varepsilon,L}(2T_L), y + 2) - \varepsilon$ by Lemma 4.6 together with (4.59)-(4.60). Hence,

$$\liminf_{L \rightarrow +\infty} \left(\inf_{y \in \mathbb{R}} v_{\varepsilon,y,L}(2T_L, L) \right) \geq 1 - \varepsilon \geq 1 - \varepsilon_0 > 1 - \delta_0 > \frac{1}{2}$$

by Remark 4.4 and since $\psi(-\infty, y) = 1$ uniformly in $y \in \mathbb{R}$ by Proposition 4.1 (ii). As a consequence, there is

$$L_{2,\varepsilon} \geq L_{1,\varepsilon} > 0$$

such that

$$0 < \tilde{T}_{\varepsilon,y,L} \leq 2T_L < +\infty \quad \text{for all } y \in \mathbb{R} \text{ and } L \geq L_{2,\varepsilon}, \quad (4.70)$$

and then

$$v_{\varepsilon,y,L}(\tilde{T}_{\varepsilon,y,L}, L) = \frac{1}{2}. \quad (4.71)$$

The following proposition shows that when L is sufficiently large, the profile $v_{\varepsilon,y,L}(t, \cdot + L)$ at time $t = \tilde{T}_{\varepsilon,y,L}$ is close to that of the initial condition $v_{\varepsilon,y,L}(0, \cdot) = v_{\varepsilon,y}^0$.

Proposition 4.7. *For any $\varepsilon \in (0, \varepsilon_0]$, with $\varepsilon_0 \in (0, \delta_0)$ as in Lemmas 4.2-4.3, there exists $L_{3,\varepsilon} \geq L_{2,\varepsilon}$, with $L_{2,\varepsilon} > 0$ as in (4.70), such that, for any $y \in \mathbb{R}$ and $L \geq L_{3,\varepsilon}$, there holds*

$$\left\| v_{\varepsilon,y,L}(\tilde{T}_{\varepsilon,y,L}, \cdot + L) - \psi(\cdot, y) \right\|_{L^\infty(\mathbb{R})} \leq \frac{\varepsilon}{2},$$

where $\tilde{T}_{\varepsilon,y,L}$ is given by (4.69), and $v_{\varepsilon,y,L}$ solves (4.54) with initial conditions $v_{\varepsilon,y}^0$ satisfying (4.67).

For the proof of Proposition 4.7, let us first show two lemmas which are concerned with the comparison of T_L and $\tilde{T}_{\varepsilon,y,L}$.

Lemma 4.8. *For any fixed $\varepsilon \in (0, \varepsilon_0]$, with $\varepsilon_0 \in (0, \delta_0)$ as in Lemmas 4.2-4.3, let T_L and $\tilde{T}_{\varepsilon,y,L}$ be given by (4.58) and (4.69), respectively, for $y \in \mathbb{R}$ and $L \geq L_{2,\varepsilon}$. Then, there is a constant $M_\varepsilon \geq 0$ such that*

$$\sup_{y \in \mathbb{R}, L \geq L_{2,\varepsilon}} |T_L - \tilde{T}_{\varepsilon,y,L}| \leq M_\varepsilon,$$

for all families of initial conditions $v_{\varepsilon,y}^0 \in C(\mathbb{R}, [0, 1])$ satisfying (4.67).

Proof. We fix $\varepsilon \in (0, \varepsilon_0]$ throughout the proof. We first prove the existence of some real numbers $L_{2,\varepsilon}^+ \in [L_{2,\varepsilon}, +\infty)$ and $M_\varepsilon^+ \geq 0$ such that

$$\sup_{y \in \mathbb{R}, L \geq L_{2,\varepsilon}^+} (\tilde{T}_{\varepsilon,y,L} - T_L) \leq M_\varepsilon^+ \quad (4.72)$$

for all families of initial conditions $v_{\varepsilon,y}^0 \in C(\mathbb{R}, [0, 1])$ satisfying (4.67). Assume by contradiction, there exist sequences $(y_n)_{n \in \mathbb{N}} \subset \mathbb{R}$, $(L_n)_{n \in \mathbb{N}} \subset [L_{2,\varepsilon}, +\infty)$, and initial conditions $(v_{\varepsilon,y_n}^0)_{n \in \mathbb{N}} \subset C(\mathbb{R}, [0, 1])$ satisfying (4.67), such that $L_n \rightarrow +\infty$ and $\tilde{T}_{\varepsilon,y_n,L_n} - T_{L_n} \rightarrow +\infty$ as $n \rightarrow +\infty$. Then for each $n \in \mathbb{N}$, by Lemma 4.6, we have

$$\begin{cases} v_{\varepsilon,y_n,L_n}(t, x) \geq \psi\left(x - X_{y_n,L_n}(t) - \eta_{\varepsilon,L_n}(t), y_n + \frac{X_{y_n,L_n}(t)}{L_n}\right) - q_{\varepsilon,L_n}(t) \\ v_{\varepsilon,y_n,L_n}(t, x) \leq \psi\left(x - X_{y_n,L_n}(t) + \eta_{\varepsilon,L_n}(t), y_n + \frac{X_{y_n,L_n}(t)}{L_n}\right) + q_{\varepsilon,L_n}(t) \end{cases} \quad (4.73)$$

for all $t \geq 0$ and $x \in \mathbb{R}$.

We already know from (4.70) that $\tilde{T}_{\varepsilon,y_n,L_n} \leq 2T_{L_n}$ for all $n \in \mathbb{N}$. Next, for any $n \in \mathbb{N}$, choosing $t = \tilde{T}_{\varepsilon,y_n,L_n}$ and $x = L_n$ in the first inequality of (4.73) yields

$$\begin{aligned} \frac{1}{2} = v_{\varepsilon,y_n,L_n}(\tilde{T}_{\varepsilon,y_n,L_n}, L_n) &\geq \psi\left(L_n - X_{y_n,L_n}(\tilde{T}_{\varepsilon,y_n,L_n}) - \eta_{\varepsilon,L_n}(\tilde{T}_{\varepsilon,y_n,L_n}), y_n + \frac{X_{y_n,L_n}(\tilde{T}_{\varepsilon,y_n,L_n})}{L_n}\right) \\ &\quad - q_{\varepsilon,L_n}(\tilde{T}_{\varepsilon,y_n,L_n}). \end{aligned}$$

Since we have assumed $\tilde{T}_{\varepsilon,y_n,L_n} - T_{L_n} \rightarrow +\infty$ as $n \rightarrow +\infty$, we may assume without loss of generality that $\tilde{T}_{\varepsilon,y_n,L_n} > T_{L_n} = L_n/c_*$ for all $n \in \mathbb{N}$ (hence, $\tilde{T}_{\varepsilon,y_n,L_n} \rightarrow +\infty$ as $n \rightarrow +\infty$). Then, we have

$$X_{y_n,L_n}(\tilde{T}_{\varepsilon,y_n,L_n}) = \underbrace{X_{y_n,L_n}(T_{L_n})}_{=L_n} + \int_{T_{L_n}}^{\tilde{T}_{\varepsilon,y_n,L_n}} c\left(y_n + \frac{X_{y_n,L_n}(t)}{L_n}\right) dt \geq L_n + c^-(\tilde{T}_{\varepsilon,y_n,L_n} - T_{L_n}),$$

where $c^- = \min_{z \in \mathbb{R}} c(z) > 0$. Remember that the function $(\xi, z) \mapsto \psi(\xi, z)$ is decreasing in $\xi \in \mathbb{R}$. It follows that

$$\frac{1}{2} \geq \psi\left(-c^-(\tilde{T}_{\varepsilon,y_n,L_n} - T_{L_n}) - \eta_{\varepsilon,L_n}(\tilde{T}_{\varepsilon,y_n,L_n}), y_n + \frac{X_{y_n,L_n}(\tilde{T}_{\varepsilon,y_n,L_n})}{L_n}\right) - q_{\varepsilon,L_n}(\tilde{T}_{\varepsilon,y_n,L_n}). \quad (4.74)$$

Since $\tilde{T}_{\varepsilon,y_n,L_n} \rightarrow +\infty$ as $n \rightarrow +\infty$, we see from (4.65) that $q_{\varepsilon,L_n}(\tilde{T}_{\varepsilon,y_n,L_n}) \rightarrow 0$ as $n \rightarrow +\infty$. Furthermore, since $0 < \tilde{T}_{\varepsilon,y_n,L_n} \leq 2T_{L_n}$ by (4.70), one infers from Remark 4.4 that the sequence $(\eta_{\varepsilon,L_n}(\tilde{T}_{\varepsilon,y_n,L_n}))_{n \in \mathbb{N}}$ is bounded. Passing to the limit as $n \rightarrow +\infty$ in (4.74) and using Proposition 4.1 (ii), we obtain $1/2 \geq 1$, which is impossible. Thus, (4.72) is proved.

Similarly as above, one can prove the existence of some real numbers $L_{2,\varepsilon}^- \in [L_{2,\varepsilon}, +\infty)$ and $M_\varepsilon^- \geq 0$ such that $\sup_{y \in \mathbb{R}, L \geq L_{2,\varepsilon}^-} (T_L - \tilde{T}_{\varepsilon,y,L}) \leq M_\varepsilon^-$ for all families of initial conditions $v_{\varepsilon,y}^0 \in C(\mathbb{R}, [0, 1])$ satisfying (4.67), by using this time in the previous paragraph the second inequality of (4.73) instead of the first one, together with $q_{\varepsilon,L_n} \leq \varepsilon < 1/2$ in $[0, +\infty)$ by (4.60).

Finally, since $0 < \tilde{T}_{\varepsilon,y,L} \leq 2T_L = 2L/c_*$ for all $y \in \mathbb{R}$ and $L \geq L_{2,\varepsilon}$ by (4.58) and (4.70), the desired conclusion of Lemma 4.8 follows, with $M_\varepsilon := \max\{M_\varepsilon^+, M_\varepsilon^-, 2L_{2,\varepsilon}^+/c_*, 2L_{2,\varepsilon}^-/c_*\}$. \square

Lemma 4.9. For any fixed $\varepsilon \in (0, \varepsilon_0]$, with $\varepsilon_0 \in (0, \delta_0)$ as in Lemmas 4.2-4.3, let T_L and $\tilde{T}_{\varepsilon, y, L}$ be given by (4.58) and (4.69), respectively, for $y \in \mathbb{R}$ and $L \geq L_{2, \varepsilon}$. Then there holds

$$\lim_{L \rightarrow +\infty} \frac{X_{y, L}(t + \tilde{T}_{\varepsilon, y, L})}{X_{y, L}(T_L)} = 1 \quad (4.75)$$

and

$$\lim_{L \rightarrow +\infty} |X_{y, L}(t + \tilde{T}_{\varepsilon, y, L}) - X_{y, L}(T_L) - c(y)(t + \tilde{T}_{\varepsilon, y, L} - T_L)| = 0 \quad (4.76)$$

locally uniformly with respect to $t \in \mathbb{R}$, and uniformly with respect to $y \in \mathbb{R}$ and with respect to the families of initial conditions $v_{\varepsilon, y}^0 \in C(\mathbb{R}, [0, 1])$ satisfying (4.67).

Proof. For any $y \in \mathbb{R}$, $L \geq L_{2, \varepsilon}$, and $t \in \mathbb{R}$, one has $X_{y, L}(T_L) = L$ and

$$\left| \frac{X_{y, L}(t + \tilde{T}_{\varepsilon, y, L})}{X_{y, L}(T_L)} - 1 \right| = \frac{\left| \int_{T_L}^{t + \tilde{T}_{\varepsilon, y, L}} c\left(y + \frac{X_{y, L}(s)}{L}\right) ds \right|}{X_{y, L}(T_L)} \leq \frac{c^+ |\tilde{T}_{\varepsilon, y, L} - T_L + t|}{L},$$

where $c^+ = \max_{z \in \mathbb{R}} c(z)$. Lemma 4.8 then immediately gives (4.75).

It remains to show (4.76). Let $y \in \mathbb{R}$, $L \geq L_{2, \varepsilon}$ and $t \in \mathbb{R}$ be arbitrary. Without loss of generality, we may assume that $t + \tilde{T}_{\varepsilon, y, L} \geq T_L$ (the case $t + \tilde{T}_{\varepsilon, y, L} \leq T_L$ can be treated identically). Since

$$X_{y, L}(t + \tilde{T}_{\varepsilon, y, L}) - X_{y, L}(T_L) - c(y)(t + \tilde{T}_{\varepsilon, y, L} - T_L) = \int_{T_L}^{t + \tilde{T}_{\varepsilon, y, L}} \left(c\left(y + \frac{X_{y, L}(s)}{L}\right) - c(y) \right) ds,$$

and since the function $z \mapsto c(z)$ is 1-periodic and of class C^1 by Proposition 4.1 (iv), we have $c(y) = c(y + 1)$ and

$$\begin{aligned} & \left| X_{y, L}(t + \tilde{T}_{\varepsilon, y, L}) - X_{y, L}(T_L) - c(y)(t + \tilde{T}_{\varepsilon, y, L} - T_L) \right| \\ & \leq \|c'\|_{L^\infty(\mathbb{R})} \times |t + \tilde{T}_{\varepsilon, y, L} - T_L| \times \max_{s \in [T_L, t + \tilde{T}_{\varepsilon, y, L}]} \left| \frac{X_{y, L}(s)}{X_{y, L}(T_L)} - 1 \right|. \end{aligned}$$

Lemma 4.8 and (4.75) then yield (4.76). \square

Now, we are ready to give the

Proof of Proposition 4.7. Fix any $\varepsilon \in (0, \varepsilon_0]$, and assume by contradiction that the conclusion is not true. Then, there exist sequences $(y_n)_{n \in \mathbb{N}} \subset \mathbb{R}$, $(L_n)_{n \in \mathbb{N}} \subset [L_{2, \varepsilon}, +\infty)$, $(x_n)_{n \in \mathbb{N}} \subset \mathbb{R}$, and initial conditions $(v_{\varepsilon, y_n}^0)_{n \in \mathbb{N}} \subset C(\mathbb{R}, [0, 1])$ satisfying (4.67) such that $L_n \rightarrow +\infty$ as $n \rightarrow +\infty$, and

$$\left| v_{\varepsilon, y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}, x_n) - \psi(x_n - L_n, y_n) \right| > \frac{\varepsilon}{2} \quad \text{for all } n \in \mathbb{N}. \quad (4.77)$$

Remember that, for each $n \in \mathbb{N}$, the inequalities stated in (4.73) hold true. By 1-periodicity of the functions $v_{\varepsilon, y, L}$ and $\psi(\cdot, y)$ with respect to y , we may assume without loss of generality that $(y_n)_{n \in \mathbb{N}} \subset [0, 1]$.

We first claim that the sequence $(x_n - L_n)_{n \in \mathbb{N}}$ is bounded. Suppose by way of contradiction that $x_n - L_n \rightarrow +\infty$ as $n \rightarrow +\infty$ (as we will sketch below, the case where $x_n - L_n \rightarrow -\infty$

as $n \rightarrow +\infty$ can be treated analogously). Choosing $t = \tilde{T}_{\varepsilon, y_n, L_n}$ (notice that $\tilde{T}_{\varepsilon, y_n, L_n} \rightarrow +\infty$ as $n \rightarrow +\infty$ by (4.58) and Lemma 4.8) and $x = x_n$ in the second inequality of (4.73), we have

$$\begin{aligned} 0 &\leq v_{\varepsilon, y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}, x_n) \\ &\leq \psi\left(x_n - L_n + X_{y_n, L_n}(T_{L_n}) - X_{y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}) + \eta_{\varepsilon, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}), y_n + \frac{X_{y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n})}{L_n}\right) \\ &\quad + q_{\varepsilon, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}). \end{aligned}$$

Due to Lemma 4.8 and the definition of $X_{y_n, L_n}(t)$, the sequence $(X_{y_n, L_n}(T_{L_n}) - X_{y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}))_{n \in \mathbb{N}}$ is bounded. Moreover, by Remark 4.4 and (4.70), the sequence $(\eta_{\varepsilon, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}))_{n \in \mathbb{N}}$ is also bounded. Thus, since $\psi(\xi, y) \rightarrow 0$ as $\xi \rightarrow +\infty$ uniformly in $y \in \mathbb{R}$ by Proposition 4.1 (ii), we have

$$\psi\left(x_n - L_n + X_{y_n, L_n}(T_{L_n}) - X_{y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}) + \eta_{\varepsilon, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}), y_n + \frac{X_{y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n})}{L_n}\right) \rightarrow 0$$

as $n \rightarrow +\infty$. This together with $q_{\varepsilon, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}) \rightarrow 0$ as $n \rightarrow +\infty$ (from (4.65) and Lemma 4.8) yields $v_{\varepsilon, y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}, x_n) \rightarrow 0$ as $n \rightarrow +\infty$, hence $v_{\varepsilon, y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}, x_n) - \psi(x_n - L_n, y_n) \rightarrow 0$ as $n \rightarrow +\infty$, which is a contradiction with (4.77).

Similarly, if $x_n - L_n \rightarrow -\infty$ as $n \rightarrow +\infty$, by using the first inequality of (4.73), one can conclude that $v_{\varepsilon, y_n, L_n}(\tilde{T}_{\varepsilon, y_n, L_n}, x_n) \rightarrow 1$ and $\psi(x_n - L_n, y_n) \rightarrow 1$ as $n \rightarrow +\infty$, which is also a contradiction with (4.77). Therefore, the sequence $(x_n - L_n)_{n \in \mathbb{N}}$ is bounded.

Next, for each $n \in \mathbb{N}$, we set $w_n(t, x) = v_{\varepsilon, y_n, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}, x + L_n)$ for $(t, x) \in [-\tilde{T}_{\varepsilon, y_n, L_n}, +\infty) \times \mathbb{R}$. Clearly, $w_n(0, 0) = 1/2$ and w_n solves

$$\partial_t w_n = a \partial_{xx} w_n + f\left(y_n + \frac{x}{L_n}, w_n\right) \quad \text{for } t > -\tilde{T}_{\varepsilon, y_n, L_n} \text{ and } x \in \mathbb{R}.$$

It is also easily seen from (4.73) that, for any $t \in \mathbb{R}$, $x \in \mathbb{R}$ and any $n \in \mathbb{N}$ such that $t + \tilde{T}_{\varepsilon, y_n, L_n} \geq 0$, $w_n(t, x)$ satisfies

$$\begin{aligned} w_n(t, x) &\geq \psi\left(x + X_{y_n, L_n}(T_{L_n}) - X_{y_n, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}) - \eta_{\varepsilon, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}), y_n + \frac{X_{y_n, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n})}{L_n}\right) \\ &\quad - q_{\varepsilon, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}) \end{aligned}$$

and

$$\begin{aligned} w_n(t, x) &\leq \psi\left(x + X_{y_n, L_n}(T_{L_n}) - X_{y_n, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}) + \eta_{\varepsilon, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}), y_n + \frac{X_{y_n, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n})}{L_n}\right) \\ &\quad + q_{\varepsilon, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}). \end{aligned}$$

Thanks to Lemma 4.8 and the boundedness of the sequence $(x_n - L_n)_{n \in \mathbb{N}}$, one can find some $\tau_\infty \in \mathbb{R}$, $\xi_\infty \in \mathbb{R}$ and $y_\infty \in [0, 1]$, such that, up to extraction of some subsequence,

$$\tilde{T}_{\varepsilon, y_n, L_n} - T_{L_n} \rightarrow \tau_\infty, \quad x_n - L_n \rightarrow \xi_\infty \quad \text{and} \quad y_n \rightarrow y_\infty \quad \text{as } n \rightarrow +\infty.$$

Furthermore, by standard parabolic estimates and possibly up to extraction of a further subsequence, there is a function $w_\infty \in C_{t;x}^{1;2}(\mathbb{R}^2)$ such that $w_n \rightarrow w_\infty$ in $C_{t;x;loc}^{1;2}(\mathbb{R}^2)$ as $n \rightarrow +\infty$. Clearly, w_∞ is an entire solution of

$$\partial_t w_\infty = a \partial_{xx} w_\infty + f(y_\infty, w_\infty) \quad \text{for } t \in \mathbb{R} \text{ and } x \in \mathbb{R}. \quad (4.78)$$

Notice that for any $t \in \mathbb{R}$, $\lim_{n \rightarrow +\infty} q_{\varepsilon, L_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}) = 0$ and there exists a constant $C > 0$ independent of t such that

$$-C \leq \liminf_{n \rightarrow +\infty} \eta_{\varepsilon, y_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}) \leq \limsup_{n \rightarrow +\infty} \eta_{\varepsilon, y_n}(t + \tilde{T}_{\varepsilon, y_n, L_n}) \leq 0 \quad (4.79)$$

(indeed, since the sequence $(T_{L_n} - \tilde{T}_{\varepsilon, y_n, L_n})_{n \in \mathbb{N}}$ is bounded by Lemma 4.8 and since $T_{L_n} \rightarrow +\infty$ as $n \rightarrow +\infty$, it follows that, for any $t \in \mathbb{R}$, one has $0 \leq t + \tilde{T}_{\varepsilon, y_n, L_n} \leq 2T_{L_n}$ for all large n ; then (4.79) follows immediately from Remark 4.4). Passing to the limit as $n \rightarrow +\infty$ in the above inequalities on $w_n(t, x)$, it follows from Lemmas 4.8 and 4.9, together with the continuity of the map $z \mapsto c(z)$ and the 1-periodicity of $\psi(\xi, y)$ with respect to y , that

$$\psi(x - c(y_\infty)(t + \tau_\infty) + C, y_\infty) \leq w_\infty(t, x) \leq \psi(x - c(y_\infty)(t + \tau_\infty) - C, y_\infty)$$

for all $(t, x) \in \mathbb{R}^2$. Namely, w_∞ is an entire solution of (4.78) which is trapped between two shifts of the corresponding traveling front $\psi(x - c(y_\infty)t, y_\infty)$. Remember that the reaction $f(y_\infty, \cdot)$ is of the bistable type from (A1). It then follows from [9, Theorem 3.1] that there exists $x_0 \in \mathbb{R}$ such that $w_\infty(t, x) \equiv \psi(x - c(y_\infty)t + x_0, y_\infty)$ for all $(t, x) \in \mathbb{R}^2$. Furthermore, notice that $w_n(0, 0) = 1/2$ for each $n \in \mathbb{N}$, hence $w_\infty(0, 0) = 1/2$. This together with the normalization condition (1.10) implies that x_0 has to be 0. Therefore, we have

$$w_\infty(t, x) \equiv \psi(x - c(y_\infty)t, y_\infty)$$

for all $(t, x) \in \mathbb{R}^2$.

On the other hand, it follows from (4.77) that $|w_n(0, x_n - L_n) - \psi(x_n - L_n, y_n)| > \varepsilon/2$ for all $n \in \mathbb{N}$. Taking the limit as $n \rightarrow +\infty$ yields $|w_\infty(0, \xi_\infty) - \psi(\xi_\infty, y_\infty)| \geq \varepsilon/2$, which is a contradiction. The proof of Proposition 4.7 is thus complete. \square

In addition to the previous observations, the last step before doing the proof of Theorem 1.6 in Section 4.4 is the global stability of the pulsating fronts ϕ_L , as stated below.

Proposition 4.10. *For any $\varepsilon \in (0, \varepsilon_0]$ with $\varepsilon_0 \in (0, \delta_0)$ as in Lemmas 4.2-4.3, and for any $L \geq L_*$ with $L_* > 0$ as in Theorem 1.1, let $(v_{\varepsilon, y, L})_{y \in \mathbb{R}}$ be the family the solutions of (4.54) with initial conditions $(v_{\varepsilon, y}^0)_{y \in \mathbb{R}}$ satisfying (4.67). Then, there exists a 1-periodic real-valued function $y \mapsto \xi_{\varepsilon, y, L}$ such that*

$$\lim_{t \rightarrow +\infty} \left\| v_{\varepsilon, y, L}(t, \cdot) - \phi_L \left(\cdot - c_L t + \xi_{\varepsilon, y, L}, \frac{\cdot}{L} + y \right) \right\|_{L^\infty(\mathbb{R})} = 0 \quad \text{for every } y \in \mathbb{R}. \quad (4.80)$$

Proof. Let $\varepsilon \in (0, \varepsilon_0] \subset (0, \delta_0)$ and $L \geq L_*$ be fixed. For each $y \in \mathbb{R}$, the function

$$u_{\varepsilon, y, L} : (t, x) \mapsto u_{\varepsilon, y, L}(t, x) := v_{\varepsilon, y, L}(t, x - Ly)$$

is a solution of

$$\partial_t u_{\varepsilon, y, L} = a \partial_{xx} u_{\varepsilon, y, L} + f(x/L, u_{\varepsilon, y, L}) \quad \text{for } t > 0 \text{ and } x \in \mathbb{R},$$

with initial condition $u_{\varepsilon, y, L}(0, \cdot) = v_{\varepsilon, y}^0(\cdot - Ly) \in C(\mathbb{R}, [0, 1])$ such that $\|u_{\varepsilon, y, L}(0, \cdot) - \psi(\cdot - Ly, y)\| \leq \varepsilon \leq \varepsilon_0 < \delta_0$, by (4.67). It then follows from Theorem 1.1 (iii) that there exists a unique $\tilde{\xi}_{\varepsilon, y, L} \in \mathbb{R}$ such that

$$\lim_{t \rightarrow +\infty} \left\| u_{\varepsilon, y, L}(t, \cdot) - \phi_L \left(\cdot - c_L t + \tilde{\xi}_{\varepsilon, y, L}, \frac{\cdot}{L} \right) \right\|_{L^\infty(\mathbb{R})} = 0.$$

Choosing $\xi_{\varepsilon, y, L} = \tilde{\xi}_{\varepsilon, y, L} + Ly$, we see from the above that (4.80) holds.

Lastly, since $v_{\varepsilon, y, L}$ and $\phi_L(\cdot, y)$ are 1-periodic in y , and since $\phi_L(\xi, y)$ is decreasing with respect to ξ , it follows that $\xi_{\varepsilon, y, L}$ in (4.80) is unique and necessarily 1-periodic in y . The proof of Proposition 4.10 is thus complete. \square

4.4 Proof of Theorem 1.6

For clarity, we proceed with several steps.

Step 1: approximation of the front profiles ϕ_L by $\psi(\cdot, y)$. In this step, we fix any $\varepsilon \in (0, \varepsilon_0]$, with $\varepsilon_0 \in (0, \delta_0)$ as in Lemmas 4.2-4.3 and $\delta_0 \in (0, 1/2)$ as in assumption (A2). Let

$$L_{4,\varepsilon} := \max\{L_*, L_{3,\varepsilon}\},$$

with $L_* > 0$ and $L_{3,\varepsilon} > 0$ provided by Theorem 1.1 and Proposition 4.7, respectively.

We shall prove in this step that there exists a real number $L_{5,\varepsilon} \in [L_{4,\varepsilon}, +\infty)$ such that, for every $L \geq L_{5,\varepsilon}$, there is a 1-periodic real-valued function $y \mapsto \xi_{\varepsilon,y,L}^*$ such that

$$\left\| \phi_L \left(\cdot + \xi_{\varepsilon,y,L}^*, \frac{\cdot}{L} + y \right) - \psi(\cdot, y) \right\|_{L^\infty(\mathbb{R})} \leq \varepsilon \quad \text{for every } y \in \mathbb{R} \text{ and } L \geq L_{5,\varepsilon}. \quad (4.81)$$

To do so, for any $y \in \mathbb{R}$ and $L \geq L_{4,\varepsilon}$ ($\geq L_{3,\varepsilon} \geq L_{2,\varepsilon}$), let $v_{\varepsilon,y,L}$ be the solution of (4.54) with initial condition $v_{\varepsilon,y}^0$ satisfying (4.67) (for instance, say, take here $v_{\varepsilon,y}^0 = \psi(\cdot, y)$). By Proposition 4.7 and (4.71), we have

$$v_{\varepsilon,y,L}(\tilde{T}_{\varepsilon,y,L}, L) = \frac{1}{2} \quad \text{and} \quad \left\| v_{\varepsilon,y,L}(\tilde{T}_{\varepsilon,y,L}, \cdot + L) - \psi(\cdot, y) \right\|_{L^\infty(\mathbb{R})} \leq \frac{\varepsilon}{2} \leq \varepsilon,$$

where $\tilde{T}_{\varepsilon,y,L}$ is defined by (4.69)-(4.70). Since f is 1-periodic with respect to its first variable, the function $(t, x) \mapsto v_{\varepsilon,y,L}(t + \tilde{T}_{\varepsilon,y,L}, x + L)$ is still a solution of (4.54), now with the initial condition $v_{\varepsilon,y,L}(\tilde{T}_{\varepsilon,y,L}, \cdot + L) \in C(\mathbb{R}, [0, 1])$. Furthermore, since both $v_{\varepsilon,y,L}$ and $\tilde{T}_{\varepsilon,y,L}$ are 1-periodic with respect to y , one has $v_{\varepsilon,y+1,L}(\tilde{T}_{\varepsilon,y+1,L}, \cdot + L) \equiv v_{\varepsilon,y,L}(\tilde{T}_{\varepsilon,y,L}, \cdot + L)$ in \mathbb{R} for each $y \in \mathbb{R}$. By using Proposition 4.7 again, we infer that

$$\left\| v_{\varepsilon,y,L}(\tilde{T}_{\varepsilon,y,L}^2, \cdot + 2L) - \psi(\cdot, y) \right\|_{L^\infty(\mathbb{R})} \leq \frac{\varepsilon}{2} \leq \varepsilon,$$

where $\tilde{T}_{\varepsilon,y,L}^2 := \min\{t > \tilde{T}_{\varepsilon,y,L} : v_{\varepsilon,y,L}(t, 2L) = 1/2\}$ (the quantities $\tilde{T}_{\varepsilon,y,L}^2$ are well defined real numbers, as are the $\tilde{T}_{\varepsilon,y,L}$'s, from the same arguments as in (4.69)-(4.71)). Then, a simple induction argument implies that, for any $k \in \mathbb{N}$,

$$\left\| v_{\varepsilon,y,L}(\tilde{T}_{\varepsilon,y,L}^k, \cdot + kL, y) - \psi(\cdot, y) \right\|_{L^\infty(\mathbb{R})} \leq \frac{\varepsilon}{2}, \quad (4.82)$$

where $\tilde{T}_{\varepsilon,y,L}^k = \min\{t > \tilde{T}_{\varepsilon,y,L}^{k-1} : v_{\varepsilon,y,L}(t, kL) = 1/2\}$. Moreover, due to Lemma 4.8, there holds

$$\sup_{y \in \mathbb{R}, L \geq L_{4,\varepsilon}, k \in \mathbb{N}} \left| \tilde{T}_{\varepsilon,y,L}^{k+1} - \tilde{T}_{\varepsilon,y,L}^k - T_L \right| < +\infty,$$

where T_L is defined by (4.58). Since $T_L \rightarrow +\infty$ as $L \rightarrow +\infty$ (independently of $y \in \mathbb{R}$), there exists $L_{5,\varepsilon} \in [L_{4,\varepsilon}, +\infty)$ such that $\tilde{T}_{\varepsilon,y,L}^k \rightarrow +\infty$ as $k \rightarrow +\infty$ uniformly in $y \in \mathbb{R}$ and $L \geq L_{5,\varepsilon}$.

On the other hand, for each $L \geq L_{5,\varepsilon}$ ($\geq L_{4,\varepsilon} \geq L_*$), it follows from Proposition 4.10 that there exists a 1-periodic real-valued function $y \mapsto \xi_{\varepsilon,y,L}$ such that

$$\lim_{t \rightarrow +\infty} \left\| v_{\varepsilon,y,L}(t, \cdot) - \phi_L \left(\cdot - c_L t + \xi_{\varepsilon,y,L}, \frac{\cdot}{L} + y \right) \right\|_{L^\infty(\mathbb{R})} = 0$$

for every $y \in \mathbb{R}$. Therefore, for each $L \geq L_{5,\varepsilon}$ and $y \in [0, 1)$, there is $k_{\varepsilon,y,L} \in \mathbb{N}$ such that

$$\left\| v_{\varepsilon,y,L}(\tilde{T}_{\varepsilon,y,L}^{k_{\varepsilon,y,L}}, \cdot + k_{\varepsilon,y,L}L) - \phi_L\left(\cdot + \xi_{\varepsilon,y,L}^*, \frac{\cdot}{L} + y\right) \right\|_{L^\infty(\mathbb{R})} \leq \frac{\varepsilon}{2},$$

where $\xi_{\varepsilon,y,L}^* := k_{\varepsilon,y,L}L - c_L \tilde{T}_{\varepsilon,y,L}^{k_{\varepsilon,y,L}} + \xi_{\varepsilon,y,L}$ (we here use the 1-periodicity of ϕ in its second variable). Combining this with (4.82) at $k = k_{\varepsilon,y,L}$, we obtain that

$$\left\| \phi_L\left(\cdot + \xi_{\varepsilon,y,L}^*, \frac{\cdot}{L} + y\right) - \psi(\cdot, y) \right\|_{L^\infty(\mathbb{R})} \leq \varepsilon. \quad (4.83)$$

We can then extend the function $y \mapsto \xi_{\varepsilon,y,L}^*$ as a 1-periodic function, so that (4.83) holds for all $y \in \mathbb{R}$, since $\phi_L(\xi, y)$ and $\psi(\xi, y)$ are 1-periodic with respect to $y \in \mathbb{R}$. The proof of Step 1 is thus complete.

Step 2: proof of

$$\lim_{\varepsilon \rightarrow 0} \left(\sup_{y \in \mathbb{R}, L \geq L_{5,\varepsilon}} |\xi_{\varepsilon,y,L}^* - \zeta_L(y)| \right) = 0, \quad (4.84)$$

with ζ_L given by (1.15). Indeed, for each $\varepsilon \in (0, \varepsilon_0]$, since $\psi(0, y) = \phi_L(\zeta_L(y), y) = 1/2$ for all $y \in \mathbb{R}$ and $L \geq L_{5,\varepsilon}$, (4.81) implies in particular that

$$\sup_{y \in \mathbb{R}, L \geq L_{5,\varepsilon}} \left| \phi_L(\xi_{\varepsilon,y,L}^*, y) - \underbrace{\phi_L(\zeta_L(y), y)}_{=1/2} \right| \leq \varepsilon \leq \varepsilon_0 < \frac{1}{4}. \quad (4.85)$$

Then, by Theorem 1.8 (i)-(ii), it follows that there exists a constant $B_1 > 0$ such that

$$|\xi_{\varepsilon,y,L}^* - \zeta_L(y)| \leq B_1 \quad \text{for all } \varepsilon \in (0, \varepsilon_0], y \in \mathbb{R} \text{ and } L \geq L_{5,\varepsilon}, \quad (4.86)$$

and also that (4.84) holds.

Step 3: $\limsup_{L \rightarrow +\infty} |L - c_L T_L| < +\infty$, with T_L given by (4.58).⁸ Notice that for any $y \in \mathbb{R}$ and $L \geq L_{5,\varepsilon_0} \geq L_{4,\varepsilon_0} = \max\{L_*, L_{3,\varepsilon_0}\}$, the function

$$(t, x) \mapsto \phi_L\left(x - c_L t + \xi_{\varepsilon_0,y,L}^*, \frac{x}{L} + y\right) = U_L\left(t - \frac{\xi_{\varepsilon_0,y,L}^*}{c_L} + \frac{Ly}{c_L}, x + Ly\right)$$

is the solution of (4.54) with initial condition $\phi_L(\cdot + \xi_{\varepsilon_0,y,L}^*, \cdot/L + y)$. Since $L \geq L_{5,\varepsilon_0} \geq L_{1,\varepsilon_0}$, it follows from (4.81) and Lemmas 4.2-4.3 (with their notations) that

$$v_{\varepsilon_0,y,L}^-(t, x) \leq \phi_L\left(x - c_L t + \xi_{\varepsilon_0,y,L}^*, \frac{x}{L} + y\right) \leq v_{\varepsilon_0,y,L}^+(t, x) \quad \text{for all } t \geq 0 \text{ and } x \in \mathbb{R}. \quad (4.87)$$

In particular, choosing $t = T_L$ and $x = L$ yields

$$\psi(-\eta_{\varepsilon_0,L}(T_L), y) - q_{\varepsilon_0,L}(T_L) \leq \phi_L(L - c_L T_L + \xi_{\varepsilon_0,y,L}^*, y) \leq \psi(\eta_{\varepsilon_0,L}(T_L), y) + q_{\varepsilon_0,L}(T_L).$$

Since $\lim_{L \rightarrow +\infty} q_{\varepsilon_0,L}(T_L) = 0$ by (4.65) and since $\sup_{L \geq L_{5,\varepsilon_0}} |\eta_{\varepsilon_0,L}(T_L)| < +\infty$ by Remark 4.4, it then follows from Proposition 4.1 (namely, the continuity of $\psi : \mathbb{R}^2 \rightarrow (0, 1)$ and its 1-periodicity in its second variable) that there exist $\kappa_0 \in (0, 1/2)$ and $\bar{L} \geq L_{5,\varepsilon_0}$ such that

$$\kappa_0 \leq \phi_L(L - c_L T_L + \xi_{\varepsilon_0,y,L}^*, y) \leq 1 - \kappa_0 \quad \text{for all } y \in \mathbb{R} \text{ and } L \geq \bar{L}.$$

⁸With this property, we here recover that $c_L \rightarrow c_*$ as $L \rightarrow +\infty$.

Hence, by Theorem 1.8 (i), there exists $B_2 > 0$ such that

$$|L - c_L T_L + \xi_{\varepsilon_0, y, L}^* - \zeta_L(y)| \leq B_2 \quad \text{for all } y \in \mathbb{R} \text{ and } L \geq \bar{L}.$$

This together with (4.86) implies that $|L - c_L T_L| \leq B_1 + B_2$ for all $L \geq \bar{L}$, which yields the desired result.

Step 4: proof the convergence (1.16). We fix any $\sigma > 0$. We have to show the existence of $L_\sigma^* > 0$ such that for any $L \geq L_\sigma^*$,

$$\sup_{y \in \mathbb{R}} \left\| \phi_L \left(\cdot + \zeta_L(y), \frac{\cdot}{L} + y \right) - \psi(\cdot, y) \right\|_{L^\infty(\mathbb{R})} \leq \sigma. \quad (4.88)$$

To do so, we first recall that the map $L \mapsto c_L$ is continuous and positive from Theorem 1.1 (i) and (iv), and that $c_L \rightarrow c_* > 0$ by Theorem 1.2 (or by Step 3 of the present proof), hence $\sup_{L \geq L_*} c_L^{-1} < +\infty$. It then follows from standard parabolic estimates that $M := \sup_{L \geq L_*} \|\partial_\xi \phi_L\|_{L^\infty(\mathbb{R}^2)} = \sup_{L \geq L_*} c_L^{-1} \|\partial_t U_L\|_{L^\infty(\mathbb{R}^2)} \in [0, +\infty)$. Next, by Step 2, there is $\varepsilon_\sigma \in (0, \min\{\varepsilon_0, \sigma/2\}]$ such that

$$M \times |\xi_{\varepsilon_\sigma, y, L}^* - \zeta_L(y)| \leq \frac{\sigma}{2} \quad \text{for all } y \in \mathbb{R} \text{ and } L \geq L_{5, \varepsilon_\sigma}.$$

Hence,

$$\left\| \phi_L \left(\cdot + \zeta_L(y), \frac{\cdot}{L} + y \right) - \phi_L \left(\cdot + \xi_{\varepsilon_\sigma, y, L}^*, \frac{\cdot}{L} + y \right) \right\|_{L^\infty(\mathbb{R})} \leq \frac{\sigma}{2} \quad \text{for all } y \in \mathbb{R} \text{ and } L \geq L_{5, \varepsilon_\sigma}.$$

Then, (4.81) (applied with ε_σ) and the inequality $\varepsilon_\sigma \leq \sigma/2$ give (4.88) for all $L \geq L_\sigma^* := L_{5, \varepsilon_\sigma} > 0$. This ends the proof of Step 4.

Step 5: the convergence (1.17) holds locally uniformly in $x \in \mathbb{R}$. Let us first show that, for any given $A > 0$,

$$\sup_{y \in \mathbb{R}, L \geq L_{5, \varepsilon}, x \in [0, A]} \left| \xi_{\varepsilon, y, L}^* + L \int_y^{y+x/L} \left(1 - \frac{c_L}{c(s)} \right) ds - \zeta_L \left(y + \frac{x}{L} \right) \right| \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0. \quad (4.89)$$

For any $\varepsilon \in (0, \varepsilon_0]$, $y \in \mathbb{R}$, $L \geq L_{5, \varepsilon}$ and $x \in [0, A]$, define

$$t_{y, L, x} := L \int_y^{y+x/L} c^{-1}(s) ds. \quad (4.90)$$

Clearly, $t_{y, L, x} \geq 0$. It is easily seen from the definition of $X_{y, L}(t)$ in (4.57) that $X_{y, L}(t_{y, L, x}) = x$. Choosing x and $t = t_{y, L, x}$ in (4.87) (these inequalities actually hold as in Step 3 with ε_0 replaced by any $\varepsilon \in (0, \varepsilon_0]$, and for any $L \geq L_{5, \varepsilon}$) gives

$$\begin{aligned} & \psi \left(-\eta_{\varepsilon, L}(t_{y, L, x}), y + \frac{X_{y, L}(t_{y, L, x})}{L} \right) - q_{\varepsilon, L}(t_{y, L, x}) \\ & \leq \phi_L \left(L \int_y^{y+x/L} \left(1 - \frac{c_L}{c(s)} \right) ds + \xi_{\varepsilon, y, L}^*, \frac{x}{L} + y \right) \\ & \leq \psi \left(\eta_{\varepsilon, L}(t_{y, L, x}), y + \frac{X_{y, L}(t_{y, L, x})}{L} \right) + q_{\varepsilon, L}(t_{y, L, x}). \end{aligned} \quad (4.91)$$

Moreover, by the normalization conditions (1.10) and (1.15), we have

$$\psi\left(0, y + \frac{X_{y,L}(t_{y,L,x})}{L}\right) = \phi_L\left(\zeta_L\left(\frac{x}{L} + y\right), \frac{x}{L} + y\right) = \frac{1}{2}.$$

It then follows that

$$\begin{aligned} & \left| \phi_L\left(L \int_y^{y+x/L} \left(1 - \frac{c_L}{c(s)}\right) ds + \xi_{\varepsilon,y,L}^*, \frac{x}{L} + y\right) - \phi_L\left(\zeta_L\left(\frac{x}{L} + y\right), \frac{x}{L} + y\right) \right| \\ & \leq \max \left\{ \left| \psi\left(\eta_{\varepsilon,L}(t_{y,L,x}), y + \frac{X_{y,L}(t_{y,L,x})}{L}\right) - \psi\left(0, y + \frac{X_{y,L}(t_{y,L,x})}{L}\right) \right|, \right. \\ & \quad \left. \left| \psi\left(-\eta_{\varepsilon,L}(t_{y,L,x}), y + \frac{X_{y,L}(t_{y,L,x})}{L}\right) - \psi\left(0, y + \frac{X_{y,L}(t_{y,L,x})}{L}\right) \right| \right\} \\ & \quad + q_{\varepsilon,L}(t_{y,L,x}). \end{aligned} \quad (4.92)$$

By definition of $t_{y,L,x}$, we have $t_{y,L,x}/L \rightarrow 0$ as $L \rightarrow +\infty$ uniformly in $y \in \mathbb{R}$ and $x \in [0, A]$. Furthermore, thanks to (4.60) and (4.66), and since $L_{5,\varepsilon} \geq L_{1,\varepsilon} \rightarrow +\infty$ as $\varepsilon \rightarrow 0$ by (4.63), one infers that

$$\sup_{y \in \mathbb{R}, L \geq L_{5,\varepsilon}, x \in [0, A]} \underbrace{\left(q_{\varepsilon,L}(t_{y,L,x}) + |\eta_{\varepsilon,L}(t_{y,L,x})| \right)}_{\geq 0} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0.$$

Passing to the limit as $\varepsilon \rightarrow 0$ in (4.92) and using the boundedness of $\partial_\xi \psi$ in \mathbb{R}^2 (which itself follows from standard elliptic estimates), one infers that

$$\sup_{y \in \mathbb{R}, L \geq L_{5,\varepsilon}, x \in [0, A]} \left| \phi_L\left(L \int_y^{y+x/L} \left(1 - \frac{c_L}{c(s)}\right) ds + \xi_{\varepsilon,y,L}^*, \frac{x}{L} + y\right) - \underbrace{\phi_L\left(\zeta_L\left(\frac{x}{L} + y\right), \frac{x}{L} + y\right)}_{=1/2} \right| \rightarrow 0$$

as $\varepsilon \rightarrow 0$. This together with Theorem 1.8 (i)-(ii) implies (4.89) (the proof is actually similar to that of Step 2; therefore, we omit the details). Combining (4.84) and (4.89), we immediately obtain that (1.17) holds in the case where $x \in [0, A]$.

It remains to consider the case where $x \in [-A, 0]$. Clearly, the convergence (4.89) holds with x replaced by A and by $x + A \in [0, A]$ as well. Since $\xi_{\varepsilon,y,L}^*$ is independent of the variable x , it then follows that

$$\sup_{y \in \mathbb{R}, L \geq L_{5,\varepsilon}, x \in [-A, 0]} \left| \zeta_L\left(y + \frac{x+A}{L}\right) - \zeta_L\left(y + \frac{A}{L}\right) - L \int_{y+A/L}^{y+(x+A)/L} \left(1 - \frac{c_L}{c(s)}\right) ds \right| \rightarrow 0$$

as $\varepsilon \rightarrow 0$. Replacing y by $y + A/L$, we obtain (1.17) in the case where $x \in [-A, 0]$. The proof of Step 5 is thus complete.

Step 6: proof of the estimate (1.18). By (4.84) and the continuity of the map $(L, y) \mapsto \zeta_L(y)$ in $[L_*, +\infty) \times \mathbb{R}$ and its 1-periodicity in y , in order to show (1.18), it suffices to prove the existence of $\varepsilon_1 \in (0, \varepsilon_0]$ such that

$$\sup_{\varepsilon \in (0, \varepsilon_1], y \in \mathbb{R}, L \geq L_{5,\varepsilon}} \left\| \xi_{\varepsilon,y,L}^* + L \int_y^{y+1/L} \left(1 - \frac{c_L}{c(s)}\right) ds - \zeta_L\left(y + \frac{\cdot}{L}\right) \right\|_{L^\infty([-L, L])} < +\infty. \quad (4.93)$$

Consider any $\varepsilon \in (0, \varepsilon_0]$, $y \in \mathbb{R}$, $L \geq L_{5,\varepsilon}$, and $x \in [0, L]$. Let $t_{y,L,x}$ be as in (4.90). There holds $0 \leq t_{y,L,x} \leq T_L$ and $X_{y,L}(t_{y,L,x}) = x$, since $X_{y,L}(T_L) = L \geq x$ and $X_{y,L}$ is increasing. It is also easily seen that the inequalities (4.91) remain valid. Moreover, by (4.60) and Remark 4.4, we have $0 < q_{\varepsilon,L}(t_{y,L,x}) \leq \varepsilon$ and $-B_3 \leq \eta_{\varepsilon,L}(T_L) \leq \eta_{\varepsilon,L}(t_{y,L,x}) \leq 0$, where B_3 is a positive constant independent of ε , y , L and x . Then, by the continuity of the map $\psi : \mathbb{R}^2 \rightarrow (0, 1)$ and its 1-periodicity with respect to its second variable, there exist $\varepsilon_1 \in (0, \varepsilon_0]$ and $\kappa_1 \in (0, 1/2)$ such that

$$\kappa_1 \leq \phi_L \left(L \int_y^{y+x/L} \left(1 - \frac{c_L}{c(s)} \right) ds + \xi_{\varepsilon,y,L}^*, \frac{x}{L} + y \right) \leq 1 - \kappa_1$$

for all $\varepsilon \in (0, \varepsilon_1]$, $y \in \mathbb{R}$, $L \geq L_{5,\varepsilon}$, and $x \in [0, L]$. It further follows from Theorem 1.8 that there exists $B_4 \in [0, +\infty)$ such that

$$\sup_{\varepsilon \in (0, \varepsilon_1], y \in \mathbb{R}, L \geq L_{5,\varepsilon}, x \in [0, L]} \left| \xi_{\varepsilon,y,L}^* + L \int_y^{y+x/L} \left(1 - \frac{c_L}{c(s)} \right) ds - \zeta_L \left(y + \frac{x}{L} \right) \right| \leq B_4. \quad (4.94)$$

This means that (4.93) holds true with $[-L, L]$ replaced by $[0, L]$.

Next, for any $x \in [-L, 0]$, we have $x+L \in [0, L]$. Since $z \mapsto \zeta_L(z)$ and $z \mapsto c(z)$ are 1-periodic, it then follows from (4.94) and the formula $T_L = L/c_*$ that

$$\begin{aligned} & \left| \xi_{\varepsilon,y,L}^* - \zeta_L \left(y + \frac{x}{L} \right) + L \int_y^{y+x/L} \left(1 - \frac{c_L}{c(s)} \right) ds \right| \\ &= \left| \xi_{\varepsilon,y,L}^* - \zeta_L \left(y + \frac{x+L}{L} \right) + L \int_y^{y+(x+L)/L} \left(1 - \frac{c_L}{c(s)} \right) ds - (L - c_L T_L) \right| \\ &\leq B_4 + |L - c_L T_L|, \end{aligned}$$

for all $\varepsilon \in (0, \varepsilon_1]$, $y \in \mathbb{R}$, $L \geq L_{5,\varepsilon}$, and $x \in [-L, 0]$. By the conclusion of Step 3 and the continuity of c_L and T_L with respect to L , the last quantity is then bounded uniformly with respect to $\varepsilon \in (0, \varepsilon_1]$, $y \in \mathbb{R}$, $L \geq L_{5,\varepsilon}$, and $x \in [-L, 0]$. Combining the above, we obtain (4.93), hence (1.18) is proved. This ends the proof of Theorem 1.6. \square

5 Appendix

The appendix is devoted to the proof of Proposition 4.1. The periodicity of the functions $y \mapsto c(y)$ and $(\xi, y) \mapsto \psi(\xi, y)$ with respect to y follows easily from the assumption that $a(y)$ and $f(y, u)$ are 1-periodic in y and the fact that $(\psi(\cdot, y), c(y))$ is the unique solution to (1.9)-(1.10). Since $c(y)$ and $\int_0^1 f(y, u) du$ have the same sign, the assumption (A3) immediately implies that $c(y) > 0$ for all $y \in \mathbb{R}$. Moreover, it is well known (see e.g., [4]) that $\psi(\xi, y)$ is decreasing in $\xi \in \mathbb{R}$. Thus, statement (i) of Proposition 4.1 is proved.

Now, we show statements (ii)-(iv) of Proposition 4.1. Let us first establish the continuity of $c(y)$ and $\psi(\cdot, y)$ with respect to y , as stated in the following lemma.

Lemma 5.1. *For any sequence $(y_n)_{n \in \mathbb{N}} \subset \mathbb{R}$ such that $y_n \rightarrow y_0$ as $n \rightarrow +\infty$ for some $y_0 \in \mathbb{R}$, we have $c(y_n) \rightarrow c(y_0)$ as $n \rightarrow +\infty$, and $\psi(\cdot, y_n) \rightarrow \psi(\cdot, y_0)$ in $C_{loc}^2(\mathbb{R})$ as $n \rightarrow +\infty$.*

Proof. The continuity of the speeds $c(y)$ with respect to y could alternatively be derived from [30, Proposition 2.3]. We here provide a short proof based on explicit estimates and we show with the

same token the continuity of the functions $\psi(\cdot, y)$ in $C_{loc}^2(\mathbb{R})$ with respect to y . More precisely, let us first show that the family $(c(y))_{y \in \mathbb{R}}$ is bounded. Since the function $y \mapsto \int_0^1 f(y, u) du$ is positive, continuous and 1-periodic, and since f is bounded in $\mathbb{R} \times [0, 1]$, there exists $s_0 \in (0, 1)$ such that $\int_0^{s_0} f(y, u) du > 0$ for all $y \in \mathbb{R}$. For each $y \in \mathbb{R}$, let $\xi_0(y)$ be the unique real number such that $\psi(\xi_0(y), y) = s_0$. Then, integrating the first equation of (1.9) against $\partial_\xi \psi(\cdot, y)$ over $[\xi_0(y), +\infty)$, we obtain

$$(\partial_\xi \psi(\xi_0(y), y))^2 \geq \frac{2}{a(y)} \int_0^{s_0} f(y, u) du > 0.$$

Letting $\xi_1(y) \in \mathbb{R}$ be such that $\partial_\xi \psi(\xi_1(y), y) = \min_{\xi \in \mathbb{R}} \partial_\xi \psi(\xi, y)$ (such a $\xi_1(y)$ exists since $\partial_\xi \psi(\pm\infty, y) = 0$ and the function $\partial_\xi \psi(\cdot, y)$ is negative and continuous), we get $\partial_{\xi\xi} \psi(\xi_1(y), y) = 0$. Since $\psi(\xi, y)$ is decreasing in $\xi \in \mathbb{R}$ and since $c(y) > 0$, it follows that

$$0 < c(y) = \frac{f(y, \psi(\xi_1(y), y))}{\partial_\xi \psi(\xi_1(y), y)} \leq \frac{\max_{x \in \mathbb{R}, u \in [0, 1]} |f(x, u)|}{-\partial_\xi \psi(\xi_0(y), y)} \leq \frac{\max_{x \in \mathbb{R}, u \in [0, 1]} |f(x, u)|}{\sqrt{\min_{x \in \mathbb{R}} 2a^{-1}(x) \int_0^{s_0} f(x, u) du}}.$$

Since the last quantity is a real number independent of $y \in \mathbb{R}$, this completes the proof of the boundedness of $(c(y))_{y \in \mathbb{R}}$.

Next, we show that $c(y_n) \rightarrow c(y_0)$ as $n \rightarrow +\infty$. Due to the periodicity, we may assume, without loss of generality, that $(y_n)_{n \in \mathbb{N}} \subset [0, 1]$ and $y_0 \in [0, 1]$. Assume by contradiction that $c(y_n)$ does not converge to $c(y_0)$ as $n \rightarrow +\infty$. Then, we can find some $c_\infty \in [0, c^+]$ such that $c_\infty \neq c(y_0)$ and, up to extraction of some sequence, $c(y_n) \rightarrow c_\infty$ as $n \rightarrow +\infty$.

Remember that $x \mapsto b(x) \in (0, 1)$ is the function given in (A1) such that $f(x, b(x)) = 0$ for all $x \in \mathbb{R}$. Pick any $\varepsilon_0 \in (0, b(y_0))$. For each $n \in \mathbb{N}$, let z_n be the unique real number such that $\psi(z_n, y_n) = \varepsilon_0$. Writing $\psi_n(\xi) = \psi(\xi + z_n, y_n)$, we obtain that $\psi_n(\xi)$ is decreasing in $\xi \in \mathbb{R}$ and

$$\begin{cases} a(y_n) \psi_n'' + c(y_n) \psi_n' + f(y_n, \psi_n) = 0 & \text{in } \mathbb{R}, \\ 0 < \psi_n < 1 & \text{in } \mathbb{R} \quad \text{and} \quad \psi_n(0) = \varepsilon_0. \end{cases}$$

By the standard elliptic estimates, we find a $C^2(\mathbb{R})$ nonincreasing function $0 \leq \psi_\infty \leq 1$ such that, up to extraction of some subsequence, $\psi_n \rightarrow \psi_\infty$ in $C_{loc}^2(\mathbb{R})$ as $n \rightarrow +\infty$, and the function ψ_∞ solves

$$a(y_0) \psi_\infty'' + c_\infty \psi_\infty' + f(y_0, \psi_\infty) = 0 \text{ in } \mathbb{R}, \quad \text{and} \quad \psi_\infty(0) = \varepsilon_0.$$

By the strong maximum principle, we have $0 < \psi_\infty < 1$ in \mathbb{R} . Since ψ_∞ is nonincreasing, the limits $\psi_\infty(\pm\infty) := \lim_{\xi \rightarrow \pm\infty} \psi_\infty(\xi)$ exist and they are zeros of $f(y_0, \cdot)$. Since $0 < \psi_\infty(0) = \varepsilon_0 < b(y_0)$, and since $0, 1, b(y_0)$ are the only three zeros of $f(y_0, \cdot)$, it is easily checked that $\psi_\infty(+\infty) = 0$ and that either $\psi_\infty(-\infty) = 1$ or $\psi_\infty(-\infty) = b(y_0)$. In other words, either $\psi_\infty(x - c_\infty t)$ is a traveling front of (1.8) with $y = y_0$ connecting 0 and 1, or it is a traveling front connecting 0 and $b(y_0)$. The former case is impossible, since the uniqueness of front speed of bistable traveling front [4] would give $c_\infty = c(y_0)$, a contradiction with our assumption $c_\infty \neq c(y_0)$. The latter case leads to a contradiction as well, since equation (1.8) (with $y = y_0$) restricted to $u \in [0, b(y_0)]$ is of the monostable type and it is well known [4] that there exists a constant $c^* > 0$ such that a (nonincreasing) traveling front with speed c connecting $b(y_0)$ and 0 exists if and only if $c \leq -c^*$. This contradicts the fact that $c_\infty \geq 0$. Therefore, $c(y_n) \rightarrow c(y_0)$ as $n \rightarrow +\infty$.

Finally, we also notice that, from the strong maximum principle applied to the nonpositive function ψ_∞' , one has $\psi_\infty' < 0$ in \mathbb{R} . Let then $z_\infty \in \mathbb{R}$ be the unique point such that $\psi(z_\infty, y_0) =$

ε_0 . Then, by the uniqueness of traveling wave of (1.8) connecting 0 and 1, it is easily seen from the above proof that the whole sequence $(\psi_n)_{n \in \mathbb{N}}$ converges to $\psi(\cdot + z_\infty, y_0)$ in $C_{loc}^2(\mathbb{R})$ as $n \rightarrow +\infty$. This implies in particular that $z_n \rightarrow z_\infty$ as $n \rightarrow +\infty$. Therefore, by the normalization condition (1.10), we obtain $\psi(\cdot, y_n) \rightarrow \psi(\cdot, y_0)$ in $C_{loc}^2(\mathbb{R})$ as $n \rightarrow +\infty$. This ends the proof of Lemma 5.1. \square

We are now ready to give the

Proof of Proposition 4.1 (ii)-(iii). We first prove that $\psi(\xi, y)$ tends to its limits as $\xi \rightarrow \pm\infty$ uniformly in $y \in \mathbb{R}$. We only consider the convergence as $\xi \rightarrow +\infty$, since the proof of the other one is identical. Assume by contradiction that there exist $\varepsilon_1 \in (0, 1)$ and a sequence $(\xi_n, y_n)_{n \in \mathbb{N}} \subset \mathbb{R}^2$ such that $\xi_n \rightarrow +\infty$ as $n \rightarrow +\infty$ and $\psi(\xi_n, y_n) \geq \varepsilon_1$ for each $n \in \mathbb{N}$. Due to the periodicity, without loss of generality, we may assume that $(y_n)_{n \in \mathbb{N}} \subset [0, 1]$ and that $y_n \rightarrow y_\infty$ as $n \rightarrow +\infty$ for some $y_\infty \in [0, 1]$. Let $\xi_\infty \in \mathbb{R}$ be the unique point such that $\psi(\xi_\infty, y_\infty) = \varepsilon_1/2$. Since $\xi_n \rightarrow +\infty$ as $n \rightarrow +\infty$, we have $\xi_n > \xi_\infty$ for all large n , whence $\psi(\xi_\infty, y_n) > \psi(\xi_n, y_n) \geq \varepsilon_1$, since $\psi(\cdot, y_n)$ is decreasing in \mathbb{R} . By Lemma 5.1, passing to the limit as $n \rightarrow +\infty$, we obtain $\psi(\xi_\infty, y_\infty) \geq \varepsilon_1$, which is impossible. Therefore, $\psi(\xi, y) \rightarrow 0$ as $\xi \rightarrow +\infty$ uniformly in $y \in \mathbb{R}$.

Next, we show that $\psi(\xi, y)$ decays to its limiting states as $\xi \rightarrow \pm\infty$ exponentially fast uniformly in $y \in \mathbb{R}$. We only give the proof of the first inequality of (4.55), since that of the second one is analogous. Since $\psi(\xi, y) \rightarrow 0$ as $\xi \rightarrow +\infty$ uniformly in y , there exists some $M > 0$ (independent of y) such that $0 < \psi(\xi, y) \leq \delta_0$ for all $\xi \geq M$ and $y \in \mathbb{R}$, where $\delta_0 \in (0, 1/2)$ is the positive constant provided by (A2). Then, by the first line of (1.2), for each $y \in \mathbb{R}$, we have

$$a(y)\partial_{\xi\xi}\psi(\xi, y) + c(y)\partial_\xi\psi(\xi, y) - \gamma_0\psi(\xi, y) \geq 0 \quad \text{for all } \xi \geq M.$$

Since the functions $y \mapsto a(y)$ and $y \mapsto c(y)$ are continuous (see Lemma 5.1), positive and 1-periodic, we can choose a constant μ_1 such that

$$0 < \mu_1 \leq \min_{y \in \mathbb{R}} \frac{c(y) + \sqrt{c^2(y) + 4\gamma_0 a(y)}}{2a(y)}.$$

It is then easily checked that $a(y)(-\mu_1)^2 + c(y) \times (-\mu_1) - \gamma_0 \leq 0$ for all $y \in \mathbb{R}$. Define $v(\xi) = C_1 e^{-\mu_1 \xi}$, where $C_1 = \delta_0 e^{\mu_1 M}$. Clearly, for each $y \in \mathbb{R}$, one has $a(y)v''(\xi) + c(y)v'(\xi) - \gamma_0 v(\xi) \leq 0$ for all $\xi \geq M$. Then, by the elliptic weak maximum principle, we obtain the first inequality of (4.55). The second one can be proved in a similar way by using the second line of (1.2).

Finally, statement (iii) is an easy consequence of statements (i)-(ii) and the fact that $\psi : \mathbb{R}^2 \rightarrow (0, 1)$ is continuous in \mathbb{R}^2 (due to Lemma 5.1). This ends the proof. \square

Remark 5.2. *Before proceeding with the proof of statement (iv), we collect some easy corollaries of the estimate (4.55), which will be used frequently later. First of all, by (4.55) and standard elliptic estimates applied to (1.9) and its derivative, we can find a constant $\tilde{C}_1 > 0$ (independent of $y \in \mathbb{R}$) such that*

$$\|1 - \psi(\cdot, y)\|_{L^2((-\infty, 0))} + \|\psi(\cdot, y)\|_{L^2((0, +\infty))} + \|\partial_\xi \psi(\cdot, y)\|_{H^2(\mathbb{R})} \leq \tilde{C}_1 \quad (5.95)$$

for all $y \in \mathbb{R}$. Furthermore, for $\xi \in \mathbb{R}$ and $y \in \mathbb{R}$, letting

$$w^*(\xi, y) := \exp\left(\frac{c(y)}{a(y)}\xi\right) \partial_\xi \psi(\xi, y), \quad (5.96)$$

one observes from the proof of (4.55) and standard elliptic estimates that there is a constant $\tilde{C}_2 > 0$ such that

$$|\partial_\xi \psi(\xi, y)| \leq \tilde{C}_2 \min\{e^{-((c(y) + \sqrt{c^2(y) + 4\gamma_0 a(y)})/(2a(y)))\xi}, 1\} \text{ for all } (\xi, y) \in \mathbb{R}^2. \quad (5.97)$$

Therefore, $w^*(\xi, y)$ decays no slower than $e^{-\tilde{\mu}_1 \xi}$ as $\xi \rightarrow +\infty$, where $\tilde{\mu}_1$ is given by

$$\tilde{\mu}_1 = \min_{y \in \mathbb{R}} \frac{-c(y) + \sqrt{c^2(y) + 4\gamma_0 a(y)}}{2a(y)} > 0,$$

and, as above, there is a constant $\tilde{C}_3 > 0$ such that $\|w^*(\cdot, y)\|_{H^2(\mathbb{R})} \leq \tilde{C}_3$ for all $y \in \mathbb{R}$.

Now, we turn to the proof of statement (iv) of Proposition 4.1. The proof shares some similarities with the proof of [18, Theorem 1.2] which established the existence of pulsating fronts of (1.1) when L is small and the convergence of those fronts as $L \rightarrow 0$ by using the implicit function theorem (see also [36, Theorem 1.1] for the study of pulsating fronts in perforated domains). Here, we will apply the implicit function theorem under a similar setting to show the C^1 -smoothness of the homogeneous fronts $(\psi(\cdot, y), c(y))$ with respect to $y \in \mathbb{R}$.

Let us first introduce a family of auxiliary operators. In the sequel, we fix a real number $\beta > 0$. For any $c \in \mathbb{R}$ and $y \in \mathbb{R}$, we define

$$M_{c,y}(v) = a(y)v'' + cv' - \beta v \text{ for } v \in H^2(\mathbb{R}).$$

Clearly, each $M_{c,y}$ maps $H^2(\mathbb{R})$ into $L^2(\mathbb{R})$. In the following lemma, we present some basic properties of this operator.

Lemma 5.3. *Fix $\beta > 0$. We have*

- (i) *For any $c \in \mathbb{R}$ and $y \in \mathbb{R}$, the operator $M_{c,y} : H^2(\mathbb{R}) \rightarrow L^2(\mathbb{R})$ is invertible. Furthermore, for every $A > 0$, there exists a constant $C > 0$ (depending on β and A , but independent of y) such that for any $c \in [-A, A]$, $y \in \mathbb{R}$ and $g \in L^2(\mathbb{R})$,*

$$\|M_{c,y}^{-1}(g)\|_{H^2(\mathbb{R})} \leq C \|g\|_{L^2(\mathbb{R})}. \quad (5.98)$$

- (ii) *For any $g \in L^2(\mathbb{R})$,*

$$M_{c_n, y_n}^{-1}(g_n) \rightarrow M_{c,y}^{-1}(g) \text{ in } H^2(\mathbb{R}) \quad (5.99)$$

as $n \rightarrow +\infty$ for all sequences $(g_n)_{n \in \mathbb{N}} \subset L^2(\mathbb{R})$, $(c_n)_{n \in \mathbb{N}} \subset \mathbb{R}$, $(y_n)_{n \in \mathbb{N}} \subset \mathbb{R}$ such that $\|g_n - g\|_{L^2(\mathbb{R})} \rightarrow 0$, $c_n \rightarrow c$ and $y_n \rightarrow y$ as $n \rightarrow +\infty$. Furthermore, the above convergence is uniform in $(g, c, y) \in B_A \times \mathbb{R}$ for any $A > 0$, where B_A is the ball given by $B_A = \{(g, c) \in L^2(\mathbb{R}) \times \mathbb{R} : \|g\|_{L^2(\mathbb{R})} + |c| \leq A\}$.

Proof. The proof of the invertibility of $M_{c,y}$ follows from similar arguments used in the proof of [18, Lemma 3.1]; therefore, we omit the details. To obtain the estimate (5.98), let $v = M_{c,y}^{-1}(g)$. Integrating $M_{c,y}(v) = g$ against v over \mathbb{R} gives $\int_{\mathbb{R}} a(y)(v')^2 + \beta v^2 = - \int_{\mathbb{R}} gv$, whence

$$\int_{\mathbb{R}} a(y)(v')^2 + \frac{\beta}{2} v^2 \leq \frac{1}{2\beta} \int_{\mathbb{R}} g^2.$$

This together with $M_{c,y}(v) = g$ implies that

$$\|v\|_{L^2(\mathbb{R})} \leq \frac{1}{\beta} \|g\|_{L^2(\mathbb{R})}, \quad \|v'\|_{L^2(\mathbb{R})} \leq \sqrt{\frac{1}{2\beta \min_{x \in \mathbb{R}} a(x)}} \|g\|_{L^2(\mathbb{R})}$$

and

$$\|v''\|_{L^2(\mathbb{R})} \leq \frac{1}{\min_{x \in \mathbb{R}} a(x)} \left(\sqrt{\frac{c^2}{2\beta \min_{x \in \mathbb{R}} a(x)}} + 1 \right) \|g\|_{L^2(\mathbb{R})}.$$

Thus, part (i) of Lemma 5.3 follows.

Next, we prove the convergence (5.99). For each $n \in \mathbb{N}$, let $u_n = M_{c_n, y_n}^{-1}(g_n)$ and $u = M_{c, y}^{-1}(g)$. Let $A > 0$ be fixed and $(g, c, y) \in B_A \times \mathbb{R}$. Without loss of generality, we may assume that $(g_n, c_n, y_n) \in B_{2A} \times \mathbb{R}$ for all $n \in \mathbb{N}$. By the proof of (5.98), we have $\|u\|_{H^2(\mathbb{R})} \leq C_1 \|g\|_{L^2(\mathbb{R})}$, where C_1 is a positive constant depending only on β and A . Notice that

$$M_{c_n, y_n}(u_n - u) = (a(y_n) - a(y))u'' + (c - c_n)u' + (g_n - g)$$

for all $n \in \mathbb{N}$. By the proof of (5.98) again, we find a positive constant C_2 depending only on β and A such that $\|u_n - u\|_{H^2(\mathbb{R})} \leq C_2 \|(a(y_n) - a(y))u'' + (c_n - c)u' + (g_n - g)\|_{L^2(\mathbb{R})}$, whence

$$\|u_n - u\|_{H^2(\mathbb{R})} \leq C_1 C_2 \left(\max_{x \in \mathbb{R}} |a'(x)| \right) |y_n - y| \|g\|_{L^2(\mathbb{R})} + C_1 C_2 |c_n - c| \|g\|_{L^2(\mathbb{R})} + C_2 \|g_n - g\|_{L^2(\mathbb{R})},$$

for all $n \in \mathbb{N}$. This implies that (5.99) holds uniformly in $(g, c, y) \in B_A \times \mathbb{R}$. The proof of Lemma 5.3 is thus complete. \square

In addition to $\beta > 0$, we also consider in the sequel an arbitrary real number y_0 . Now, for any $(v, c, y) \in L^2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R}$, we define

$$K(v, c, y) : \xi \mapsto K(v, c, y)(\xi) := a(y) \partial_{\xi\xi} \psi(\xi, y_0) + c \partial_{\xi} \psi(\xi, y_0) + \beta v(\xi) + f(y, v(\xi) + \psi(\xi, y_0)).$$

Clearly, $K(0, c(y_0), y_0)(\xi) = 0$ for all $\xi \in \mathbb{R}$, by (1.9). By Remark 5.2, we know that $\psi(\cdot, y_0) \in L^2((0, +\infty))$, $1 - \psi(\cdot, y_0) \in L^2((-\infty, 0))$, and $\partial_{\xi} \psi(\cdot, y_0) \in H^2(\mathbb{R})$. Moreover, since the function $f(x, u)$ satisfies (A1) and is globally Lipschitz continuous with respect to u uniformly in $x \in \mathbb{R}$, it follows that the function $\xi \mapsto f(y, v(\xi) + \psi(\xi, y_0))$ belongs to $L^2(\mathbb{R})$. Therefore, for any $(v, c, y) \in L^2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R}$, we have $K(v, c, y) \in L^2(\mathbb{R})$.

Now, for any $(v, c, y) \in H^2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R}$, we set $G(v, c, y) = (G_1, G_2)(v, c, y)$ with

$$G_1(v, c, y) = v + M_{c, y}^{-1}(K(v, c, y)) \quad \text{and} \quad G_2(v, c, y) = v(0).$$

It is easily seen from Lemma 5.3 (i) that the function G maps $H^2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R}$ into $H^2(\mathbb{R}) \times \mathbb{R}$. Remember that $(\psi(\cdot, y), c(y))$ is the unique solution of (1.9)-(1.10). It is straightforward to check that for any $y \in \mathbb{R}$, we have $\psi(\cdot, y) - \psi(\cdot, y_0) \in H^2(\mathbb{R})$ by (5.95), and

$$G(\psi(\cdot, y) - \psi(\cdot, y_0), c(y), y) = (0, 0).$$

On the other hand, thanks to Lemma 5.3 (ii) and the assumption that the function $f(x, u)$ is of class C^1 in $(x, u) \in \mathbb{R}^2$, and $f(x, u)$, $\partial_u f(x, u)$ are globally Lipschitz continuous in u uniformly in $x \in \mathbb{R}$, by similar arguments to those used in the proof of [18, Lemma 3.4], one can check that the function $G : H^2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R} \rightarrow H^2(\mathbb{R}) \times \mathbb{R}$ is continuously Fréchet differentiable (the verification is simpler in our case, since no singularity occurs) and that the derivatives are given by

$$\partial_{(v, c, y)} G(v, c, y)(\tilde{v}, \tilde{c}, \tilde{y}) = \begin{pmatrix} \partial_v G_1(v, c, y)(\tilde{v}) + \partial_c G_1(v, c, y)(\tilde{c}) + \partial_y G_1(v, c, y)(\tilde{y}) \\ \partial_v G_2(v, c, y)(\tilde{v}) + \partial_c G_2(v, c, y)(\tilde{c}) + \partial_y G_2(v, c, y)(\tilde{y}) \end{pmatrix}$$

for all $(\tilde{v}, \tilde{c}, \tilde{y}) \in H^2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R}$, where

$$\left\{ \begin{array}{l} \partial_v G_1(v, c, y)(\tilde{v}) = \tilde{v} + M_{c,y}^{-1} [\partial_u f(y, v + \psi(\cdot, y_0)) \tilde{v} + \beta \tilde{v}], \\ \partial_c G_1(v, c, y)(\tilde{c}) = -\tilde{c} M_{c,y}^{-1} \left\{ \partial_\xi [M_{c,y}^{-1}(K(v, c, y)) - \psi(\cdot, y_0)] \right\}, \\ \partial_y G_1(v, c, y)(\tilde{y}) = -\tilde{y} M_{c,y}^{-1} \left\{ a'(y) \partial_{\xi\xi} [M_{c,y}^{-1}(K(v, c, y)) - \psi(\cdot, y_0)] - \partial_x f(y, v + \psi(\cdot, y_0)) \right\}, \end{array} \right.$$

and

$$\partial_v G_2(v, c, y)(\tilde{v}) = \tilde{v}(0), \quad \partial_c G_2(v, c, y)(\tilde{c}) = 0, \quad \partial_y G_2(v, c, y)(\tilde{y}) = 0.$$

We will apply the implicit function theorem for the function $G : H^2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R} \rightarrow H^2(\mathbb{R}) \times \mathbb{R}$ and show the C^1 -smoothness of the functions $y \mapsto c(y)$ and $y \mapsto \psi(\cdot, y)$ at $y = y_0$. To do so, we need the following lemma.

Lemma 5.4. *The operator $Q_{y_0} = \partial_{(v,c)} G(0, c(y_0), y_0) : H^2(\mathbb{R}) \times \mathbb{R} \rightarrow H^2(\mathbb{R}) \times \mathbb{R}$ is invertible. Furthermore, there exists a constant $C > 0$ independent of y_0 such that for any $(\tilde{g}, \tilde{d}) \in H^2(\mathbb{R}) \times \mathbb{R}$,*

$$\|Q_{y_0}^{-1}(\tilde{g}, \tilde{d})\|_{H^2(\mathbb{R}) \times \mathbb{R}} \leq C \|(\tilde{g}, \tilde{d})\|_{H^2(\mathbb{R}) \times \mathbb{R}}, \quad (5.100)$$

where the space $H^2(\mathbb{R}) \times \mathbb{R}$ is endowed with norm $\|(\tilde{g}, \tilde{d})\|_{H^2(\mathbb{R}) \times \mathbb{R}} = \|\tilde{g}\|_{H^2(\mathbb{R})} + |\tilde{d}|$.

Proof. Let H be the linearization of (1.9) with respect to ψ at $(\psi(\cdot, y), c(y))$ with $y = y_0$. Namely, we define

$$H(v) = a(y_0)v'' + c(y_0)v' + \partial_u f(y_0, \psi(\cdot, y_0))v \quad \text{for } v \in H^2(\mathbb{R}).$$

By [36], the operator H and the adjoint operator H^* , given by $H^*(v) = a(y_0)v'' - c(y_0)v' + \partial_u f(y_0, \psi(\cdot, y_0))v$ for $v \in H^2(\mathbb{R})$, have algebraically simple eigenvalue 0. Moreover, the kernel $\ker(H)$ is equal to $\mathbb{R}\partial_\xi \psi(\cdot, y_0)$ and $\ker(H^*)$ is equal to $\mathbb{R}w^*$, with $w^* := w^*(\cdot, y_0)$ given by (5.96). Moreover, it is also known from [36] that the range of H is closed in $L^2(\mathbb{R})$. Based on these properties, one can conclude that the operator $Q_{y_0} : H^2(\mathbb{R}) \times \mathbb{R} \rightarrow H^2(\mathbb{R}) \times \mathbb{R}$ is invertible. The proof follows the same lines as those used in [18, Lemma 3.4] and [36, Lemma 2.4]; therefore, we do not repeat the details here.

Next, we show the estimate (5.100) (with a constant C independent of y_0). To do so, by linearity, it is sufficient to show that

$$\sup_{y_0 \in \mathbb{R}, (\tilde{g}, \tilde{d}) \in \mathcal{S}} \|Q_{y_0}^{-1}(\tilde{g}, \tilde{d})\|_{H^2(\mathbb{R}) \times \mathbb{R}} < +\infty,$$

where $\mathcal{S} := \{(\tilde{g}, \tilde{d}) \in H^2(\mathbb{R}) \times \mathbb{R} : \|(\tilde{g}, \tilde{d})\|_{H^2(\mathbb{R}) \times \mathbb{R}} = 1\}$. For any $y_0 \in \mathbb{R}$ and $(\tilde{g}, \tilde{d}) \in \mathcal{S}$, set $(\tilde{v}, \tilde{c}) = Q_{y_0}^{-1}(\tilde{g}, \tilde{d})$ and $\tilde{w} = \tilde{v} - \tilde{g}$. It then follows from the definition of Q_{y_0} that

$$\tilde{w} = -M_{c(y_0), y_0}^{-1} [\partial_u f(y_0, \psi(\cdot, y_0)) \tilde{v} + \beta \tilde{v}] - \tilde{c} M_{c(y_0), y_0}^{-1} [\partial_\xi \psi(\cdot, y_0)],$$

whence, owing to the definition of H ,

$$H(\tilde{w}) = -\partial_u f(y_0, \psi(\cdot, y_0)) \tilde{g} - \beta \tilde{g} - \tilde{c} \partial_\xi \psi(\cdot, y_0).$$

Testing this equation with $w^*(\cdot, y_0) \in \ker(H^*)$ (which is given by (5.96)), we obtain

$$\tilde{c} \int_{\mathbb{R}} \partial_\xi \psi(\xi, y_0) w^*(\xi, y_0) d\xi = - \int_{\mathbb{R}} [\partial_u f(y_0, \psi(\xi, y_0)) \tilde{g}(\xi) + \beta \tilde{g}(\xi)] w^*(\xi, y_0) d\xi.$$

From Remark 5.2, the map $y \mapsto \int_{\mathbb{R}} \partial_{\xi} \psi(\xi, y) w^*(\xi, y) d\xi$ is positive, periodic, continuous in \mathbb{R} , and $\|w^*(\cdot, y)\|_{L^2(\mathbb{R})}$ is bounded uniformly in $y \in \mathbb{R}$. Therefore, there is a constant $C_1 > 0$ (independent of $y_0 \in \mathbb{R}$ and $(\tilde{g}, \tilde{d}) \in \mathcal{S}$) such that

$$|\tilde{c}| \leq C_1 \|\tilde{g}\|_{L^2(\mathbb{R})} \leq C_1. \quad (5.101)$$

It remains to estimate $\|\tilde{v}\|_{H^2(\mathbb{R})}$. Let $\tilde{w} = \tilde{w}_1 + \tilde{w}_2$, where \tilde{w}_1 is orthogonal to $\partial_{\xi} \psi(\cdot, y_0)$ in $L^2(\mathbb{R})$ and $\tilde{w}_2 \in \mathbb{R} \partial_{\xi} \psi(\cdot, y_0)$. Since $\tilde{w}_2 \in \text{Ker}(H)$, we have

$$H(\tilde{w}_1) = H(\tilde{w}) = -\partial_u f(y_0, \psi(\cdot, y_0)) \tilde{g} - \beta \tilde{g} - \tilde{c} \partial_{\xi} \psi(\cdot, y_0).$$

We first claim that $\|\tilde{w}_1\|_{L^2(\mathbb{R})}$ is bounded uniformly with respect to $y_0 \in \mathbb{R}$ and $(\tilde{g}, \tilde{d}) \in \mathcal{S}$. Assume by contradiction that this is not true. Then, thanks to (5.101) and the 1-periodicity with respect to y_0 , there exist a sequence $(y_n)_{n \in \mathbb{N}} \subset [0, 1]$, a sequence $(\tilde{g}_n, \tilde{d}_n)_{n \in \mathbb{N}} \subset \mathcal{S}$, a bounded sequence $(\tilde{c}_n)_{n \in \mathbb{N}} \subset \mathbb{R}$ and a sequence of $(p_n)_{n \in \mathbb{N}} \subset H^2(\mathbb{R})$ with $\|p_n\|_{L^2(\mathbb{R})} \rightarrow +\infty$ as $n \rightarrow +\infty$ such that p_n is orthogonal to $\partial_{\xi} \psi(\cdot, y_n)$ for each $n \in \mathbb{N}$ and

$$H_n(p_n) := a(y_n) p_n'' + c(y_n) p_n' + \partial_u f(y_n, \psi(\cdot, y_n)) p_n = -\partial_u f(y_n, \psi(\cdot, y_n)) \tilde{g}_n - \beta \tilde{g}_n - \tilde{c}_n \partial_{\xi} \psi(\cdot, y_n).$$

For each $n \in \mathbb{N}$, writing

$$q_n = \frac{p_n}{\|p_n\|_{L^2(\mathbb{R})}} \quad \text{and} \quad g_n = \frac{-\partial_u f(y_n, \psi(\cdot, y_n)) \tilde{g}_n - \beta \tilde{g}_n - \tilde{c}_n \partial_{\xi} \psi(\cdot, y_n)}{\|p_n\|_{L^2(\mathbb{R})}},$$

we have $\|g_n\|_{L^2(\mathbb{R})} \rightarrow 0$ as $n \rightarrow +\infty$ by (5.95), and $M_{c(y_n), y_n}(q_n) = -\partial_u f(y_n, \psi(\cdot, y_n)) q_n - \beta q_n + g_n$ is uniformly bounded in $L^2(\mathbb{R})$. Up to extraction of a sequence, we can assume that $y_n \rightarrow y_* \in [0, 1]$ as $n \rightarrow +\infty$. It follows from Lemma 5.3 and the boundedness of the sequence $(c(y_n))_{n \in \mathbb{N}}$ that the sequence $(q_n)_{n \in \mathbb{N}}$ is bounded in $H^2(\mathbb{R})$. Hence, up to extraction of another subsequence, the sequence $(q_n)_{n \in \mathbb{N}}$ converges in $H^2(\mathbb{R})$ weakly and in $C_{loc}^1(\mathbb{R})$ to some $q_* \in H^2(\mathbb{R})$. This implies that $a(y_*) q_*'' + c(y_*) q_*' + \partial_u f(y_*, \psi(\cdot, y_*)) q_* = 0$. Furthermore, q_* is orthogonal to $\partial_{\xi} \psi(\cdot, y_*)$ in $L^2(\mathbb{R})$, by (5.97) and Lemma 5.1. As a consequence, $q_* = 0$.

Therefore, $q_n \rightarrow 0$ in $C_{loc}^1(\mathbb{R})$, hence in $L_{loc}^2(\mathbb{R})$, as $n \rightarrow +\infty$. In order to get a contradiction, we further show that this convergence holds in $L^2(\mathbb{R})$. Indeed, by Proposition 4.1 (ii) and the assumption (A2), there exists a constant $M' > 0$ (independent of $n \in \mathbb{N}$) such that $\partial_u f(y_n, \psi(\xi, y_n)) \leq -\gamma_0/2$ for all $\xi \in (-\infty, -M'] \cup [M', +\infty)$ and $n \in \mathbb{N}$. Integrating the equation $H_n(q_n) = g_n$ against q_n over $(M', +\infty)$, we get

$$-a(y_n) \int_{M'}^{+\infty} (q_n')^2(\xi) d\xi + \int_M^{+\infty} \partial_u f(y_n, \psi(\xi, y_n)) q_n^2(\xi) d\xi \rightarrow 0 \quad \text{as } n \rightarrow +\infty.$$

This implies in particular that $\|q_n\|_{L^2((M', +\infty))} \rightarrow 0$ as $n \rightarrow +\infty$. The same analysis over $(-\infty, -M')$ gives $\|q_n\|_{L^2((-\infty, -M'))} \rightarrow 0$ as $n \rightarrow +\infty$. Finally, the sequence $(q_n)_{n \in \mathbb{N}}$ tends to 0 in $L^2(\mathbb{R})$, which is a contradiction with the fact that $\|q_n\|_{L^2(\mathbb{R})} = 1$ for each $n \in \mathbb{N}$. We can thus conclude that $\|\tilde{w}_1\|_{L^2(\mathbb{R})}$ is bounded uniformly with respect to $y_0 \in \mathbb{R}$ and $(\tilde{g}, \tilde{d}) \in \mathcal{S}$.

Furthermore, since $M_{c(y_0), y_0}(\tilde{w}_1) = -\partial_u f(y_0, \psi(\cdot, y_0)) (\tilde{w}_1 + \tilde{g}) - \beta (\tilde{w}_1 + \tilde{g}) - \tilde{c} \partial_{\xi} \psi(\cdot, y_0)$, it follows from (5.95), (5.101) and Lemma 5.3, together with the boundedness of $\tilde{c}(y_0)$ with respect to $y_0 \in \mathbb{R}$, that there exists a positive constant C_2 (independent of y_0 and $(\tilde{g}, \tilde{d}) \in \mathcal{S}$) such that

$$\|\tilde{w}_1\|_{H^2(\mathbb{R})} \leq C_2. \quad (5.102)$$

By the Sobolev inequality, $\max_{\xi \in \mathbb{R}} |\tilde{w}_1(\xi)| \leq C_3 \|\tilde{w}_1\|_{H^2(\mathbb{R})} \leq C_3 C_2$ for some constant $C_3 > 0$ independent of y_0 and $(\tilde{g}, \tilde{d}) \in \mathcal{S}$. In particular, we have $|\tilde{w}_1(0)| \leq C_3 C_2$. Similarly, we have $|\tilde{g}(0)| \leq C_3 \|\tilde{g}\|_{H^2(\mathbb{R})} \leq C_3$. Finally, remembering that $\tilde{v}(0) = \tilde{d} \in [-1, 1]$ and $\tilde{v} - \tilde{g} = \tilde{w} = \tilde{w}_1 + \tilde{w}_2$, we have $\tilde{w}_2(0) = \tilde{d} - \tilde{g}(0) - \tilde{w}_1(0)$. This together with Proposition 4.1 (iii), (1.10), (5.95) and the fact that $\tilde{w}_2 \in \mathbb{R} \partial_{\xi} \psi(\cdot, y_0)$ implies that

$$\|\tilde{w}_2\|_{H^2(\mathbb{R})} = \left| \frac{\tilde{w}_2(0)}{\partial_{\xi} \psi(0, y_0)} \right| \|\partial_{\xi} \psi(\cdot, y_0)\|_{H^2(\mathbb{R})} \leq C_4$$

for some constant $C_4 > 0$ independent of y_0 and $(\tilde{g}, \tilde{d}) \in \mathcal{S}$. Combining this with (5.101)-(5.102) and the fact that $\tilde{v} = \tilde{w} + \tilde{g}$, we obtain that $\|Q_{y_0}^{-1}(\tilde{g}, \tilde{d})\|_{H^2(\mathbb{R}) \times \mathbb{R}} \leq C_5$ for some constant $C_5 > 0$ independent of y_0 and $(\tilde{g}, \tilde{d}) \in \mathcal{S}$. This completes the proof of Lemma 5.4. \square

Based on the above preparations, we are now ready to complete the

Proof of Proposition 4.1 (iv). For any $y_0 \in \mathbb{R}$, we define

$$T(y) = (\psi(\cdot, y) - \psi(\cdot, y_0), c(y)) \quad \text{for } y \in \mathbb{R}.$$

Clearly, T maps \mathbb{R} into $H^2(\mathbb{R}) \times \mathbb{R}$, and $G(T(y), y) = (0, 0)$ for all $y \in \mathbb{R}$, by (1.9) and the definitions of T and G . Applying the implicit function theorem to the function $G : H^2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R} \rightarrow H^2(\mathbb{R}) \times \mathbb{R}$, one infers that, for each $y_0 \in \mathbb{R}$, there is $\delta > 0$ such that the function $T : (y_0 - \delta, y_0 + \delta) \rightarrow H^2(\mathbb{R}) \times \mathbb{R}$ is continuously Fréchet differentiable, namely of class C^1 . Denote the derivative operator at $y \in (y_0 - \delta, y_0 + \delta)$ by $A_y : \mathbb{R} \rightarrow H^2(\mathbb{R}) \times \mathbb{R}$. Thus, the function $y \mapsto A_y$ is continuous from $(y_0 - \delta, y_0 + \delta)$ to $\mathcal{L}(\mathbb{R}, H^2(\mathbb{R}) \times \mathbb{R})$. At $y = y_0$, we have $A_{y_0}(\tilde{y}) = -Q_{y_0}^{-1}(\partial_y G(0, c(y_0), y_0)(\tilde{y}))$ for every $\tilde{y} \in \mathbb{R}$, that is,

$$A_{y_0}(\tilde{y}) = -\tilde{y} Q_{y_0}^{-1} \left(M_{c(y_0), y_0}^{-1} (a'(y_0) \partial_{\xi\xi} \psi(\cdot, y_0) + \partial_x f(y_0, \psi(\cdot, y_0))), 0 \right).$$

Since the function f is of class C^1 in \mathbb{R}^2 and satisfies (A1), it follows from (5.95) that there exists a constant $C_1 > 0$ independent of y_0 such that $\|a'(y_0) \partial_{\xi\xi} \psi(\cdot, y_0) + \partial_x f(y_0, \psi(\cdot, y_0))\|_{L^2(\mathbb{R})} \leq C_1$. Remember that the periodic function $y_0 \mapsto c(y_0)$ is bounded by Lemma 5.1. Then, by Lemmas 5.3 and 5.4, we find a constant $C_2 > 0$, independent of $y_0 \in \mathbb{R}$, such that $\|A_{y_0}(\tilde{y})\|_{H^2(\mathbb{R}) \times \mathbb{R}} \leq C_2 |\tilde{y}|$ for all $\tilde{y} \in \mathbb{R}$, that is,

$$\|A_{y_0}\|_{H^2(\mathbb{R}) \times \mathbb{R}} \leq C_2, \tag{5.103}$$

where we identify $A_{y_0} \in \mathcal{L}(\mathbb{R}, H^2(\mathbb{R}) \times \mathbb{R})$ to an element of $H^2(\mathbb{R}) \times \mathbb{R}$, with a slight abuse of notation.

On the other hand, notice that, for each $y_0 \in \mathbb{R}$,

$$T(y_0 + \tilde{y}) - T(y_0) = (\psi(\cdot, y_0 + \tilde{y}) - \psi(\cdot, y_0), c(y_0 + \tilde{y}) - c(y_0)) = A_{y_0} \tilde{y} + \omega(y_0, \tilde{y}) \quad \text{for all } \tilde{y} \in \mathbb{R},$$

where $\omega(y_0, \tilde{y}) \in H^2(\mathbb{R}) \times \mathbb{R}$ satisfies $\|\omega(y_0, \tilde{y})\|_{H^2(\mathbb{R}) \times \mathbb{R}} = o(|\tilde{y}|)$ as $\tilde{y} \rightarrow 0$. Hence, for each $y_0 \in \mathbb{R}$,

$$\left\| \left(\frac{\psi(\cdot, y_0 + \tilde{y}) - \psi(\cdot, y_0)}{\tilde{y}}, \frac{c(y_0 + \tilde{y}) - c(y_0)}{\tilde{y}} \right) - A_{y_0} \right\|_{H^2(\mathbb{R}) \times \mathbb{R}} \rightarrow 0 \quad \text{as } \tilde{y} \xrightarrow{\neq} 0. \tag{5.104}$$

Writing $A_{y_0} = (A_{y_0}^1, A_{y_0}^2) \in H^2(\mathbb{R}) \times \mathbb{R}$ for each $y_0 \in \mathbb{R}$, we see from (5.104) and the previous paragraph that the functions $y_0 \mapsto c(y_0)$ and $y_0 \mapsto \psi(\cdot, y_0)$ are of class C^1 in \mathbb{R} with derivatives $\partial_y \psi(\cdot, y_0) = A_{y_0}^1$ and $c'(y_0) = A_{y_0}^2$ at each $y_0 \in \mathbb{R}$. It further follows from (5.103) that $\|\partial_y \psi(\cdot, y_0)\|_{H^2(\mathbb{R})} \leq C_2$ for all $y_0 \in \mathbb{R}$, with C_2 independent of y_0 . Using the Sobolev inequality, we immediately obtain (4.56). Finally, by standard elliptic estimates and Lemma 5.1, the function $(\xi, y) \mapsto \psi(\xi, y)$ is of class $C_{\xi, y}^{2;1}(\mathbb{R}^2)$. This completes the proof of Proposition 4.1. \square

References

- [1] M. Alfaro, T. Giletti, *Varying the direction of propagation in reaction-diffusion equations in periodic media*, Netw. Heterog. Media 11 (2016), 369-393.
- [2] S. B. Angenent, *The zero set of a solution of a parabolic equation*, J. Reine Angew. Math. 390 (1988), 79-96.
- [3] D. G. Aronson, N. V. Mantzaris, H. G. Othmer, *Wave propagation and blocking in inhomogeneous media*, Disc. Cont. Dyn. Syst. 13 (2005), 843-876.
- [4] D. G. Aronson, H. F. Weinberger, *Multidimensional nonlinear diffusions arising in population genetics*, Adv. Math. 30 (1978), 33-76.
- [5] G. Barles, L. Bronsard, P. E. Souganidis, *Front propagation for reaction-diffusion equations of a bistable type*, Ann. Inst. Henri Poincaré, Anal. Non Linéaire 9 (1992), 479-496.
- [6] G. Barles, G. Soner, P. E. Souganidis, *Front propagation and phase field theory*, SIAM J. Control. Optim. 31 (1993), 489-469.
- [7] G. Barles, P. E. Souganidis, *A new approach to front propagation problems: theory and applications*, Arch. Ration. Mech. Anal. 141(1998), 237-296.
- [8] H. Berestycki, F. Hamel, *Front propagation in periodic excitable media*, Comm. Pure Appl. Math. 55 (2002), 949-1032.
- [9] H. Berestycki, F. Hamel, *Generalized travelling waves for reaction-diffusion equations*, Perspectives in Nonlinear Partial Differential Equations, Contemp. Math. 446, Amer. Math. Soc., Providence, RI, 2007, 101-123.
- [10] H. Berestycki, F. Hamel, *Generalized transition waves and their properties*, Comm. Pure Appl. Math. 65 (2012), 592-648.
- [11] H. Berestycki, F. Hamel, N. Nadirashvili, *The speed of propagation for KPP type problems. I - Periodic framework*, J. Europ. Math. Soc. 7 (2005), 173-213.
- [12] H. Berestycki, F. Hamel, L. Roques, *Analysis of the periodically fragmented environment model : II - Biological invasions and pulsating travelling fronts*, J. Math. Pures Appl. 84 (2005), 1101-1146.
- [13] H. Berestycki, L. Nirenberg, *On the method of moving planes and the sliding method*, Bol. Soc. Bras. Mat. 22 (1991), 1-37.
- [14] J.-G. Caputo, B. Sarels, *Reaction-diffusion front crossing a local defect*, Phys. Rev. E 84 (2011), 041108.
- [15] C. De Zan, P. Soravia, *Singular limits of reaction diffusion equations and geometric flows with discontinuous velocity*, Nonlinear Anal. 200 (2020), 111989.
- [16] W. Ding, T. Giletti, *Admissible speeds in spatially periodic bistable reaction-diffusion equations*, Adv. Math. 398 (2021), 107889.
- [17] W. Ding, F. Hamel, X.-Q. Zhao, *Propagation phenomena for periodic bistable reaction-diffusion equations*, Calc. Var. Part. Diff. Equations 54 (2015), 2517-2551.
- [18] W. Ding, F. Hamel, X.-Q. Zhao, *Bistable pulsating fronts for reaction-diffusion equations in a periodic habitat*, Indiana Univ. Math. J. 66 (2017), 1189-1265.
- [19] F. Dkhil, A. Stevens, *Traveling wave speeds in rapidly oscillating media*, Disc. Cont. Dyn. Syst. A, 25 (2009), 89-108.

- [20] J. Dowdall, V. LeBlanc, F. Lutscher, *Invasion pinning in a periodically fragmented habitat*, J. Math. Biol. 77 (2018), 55-78.
- [21] Y. Du, H. Matano, *Convergence and sharp thresholds for propagation in nonlinear diffusion problems*, J. Europ. Math. Soc. 12 (2010), 279-312.
- [22] A. Ducrot, *A multi-dimensional bistable nonlinear diffusion equation in a periodic medium*, Math. Ann. 366 (2016), 783-818.
- [23] A. Ducrot, T. Giletti, H. Matano, *Existence and convergence to a propagating terrace in one-dimensional reaction-diffusion equations*, Trans. Amer. Math. Soc. 366 (2014), 5541-5566.
- [24] S. Eberle, *Front blocking in the presence of gradient drift*, J. Diff. Equations 267 (2019), 7154-7166.
- [25] S. Eberle, *Front blocking versus propagation in the presence of drift term varying in the direction of propagation*, Nonlinear Anal. 197 (2020), 111836.
- [26] M. El Smailly, F. Hamel, L. Roques, *Homogenization and influence of fragmentation in a biological invasion model*, Disc. Cont. Dyn. Syst. A 25 (2009), 321-342.
- [27] J. Fang, X.-Q. Zhao, *Bistable traveling waves for monotone semiflows with applications*, J. Europ. Math. Soc. 17 (2015), 2243-2288.
- [28] P. C. Fife, J. B. McLeod, *The approach of solutions of non-linear diffusion equations to traveling front solutions*, Arch. Ration. Mech. Anal. 65 (1977), 335-361.
- [29] R. A. Fisher, *The advance of advantageous genes*, Ann. Eugenics 7 (1937), 335-369.
- [30] J. Gärtner, *Bistable reaction-diffusion equations and excitable media*, Math. Nachr. 112 (1983), 125-152.
- [31] T. Giletti, H. Matano, *Existence and uniqueness of propagating terraces*, Commun. Contemp. Math. 22 (2020), 1950055.
- [32] T. Giletti, L. Rossi, *Pulsating solutions for multidimensional bistable and multistable equations*, Math. Ann. 378 (2020), 1555-1611.
- [33] H. Guo, *Propagating speeds of bistable transition fronts in spatially periodic media*, Calc. Var., Part. Diuff. Equations 57 (2018), 47.
- [34] F. Hamel, J. Fayard, L. Roques, *Spreading speeds in slowly oscillating environments*, Bull. Math. Biol. 72 (2010), 1166-1191.
- [35] F. Hamel, G. Nadin, L. Roques, *A viscosity solution method for the spreading speed formula in slowly varying media*, Indiana Univ. Math. J. 60 (2011), 1229-1247.
- [36] S. Heinze, *Wave solutions to reaction-diffusion systems in perforated domains*, Z. Anal. Anwendungen 20 (2001), 661-670.
- [37] S. Heinze, G. Papanicolaou, A. Stevens, *Variational principles for propagation speeds in inhomogeneous media*, SIAM J. Appl. Math. 62 (2001), 129-148.
- [38] J.P. Keener, *Homogenization and propagation in the bistable equation*, Physica D 136 (2000), 1-17.
- [39] N. Kinezaki, K. Kawasaki, N. Shigesada, *Spatial dynamics of invasion in sinusoidally varying environments*, Popul. Ecol. 48 (2006), 263-270.

- [40] A. N. Kolmogorov, I. G. Petrovsky, N. S. Piskunov, *Étude de l'équation de la diffusion avec croissance de la quantité de matière et son application à un problème biologique*, Bjul. Moskowskogo Gos. Univ. A 1 (1937), 1-26.
- [41] T. J. Lewis, J. P. Keener, *Wave-block in excitable media due to regions of depressed excitability*, SIAM J. Appl. Math. 61 (2000), 293-316.
- [42] X. Liang, X.-Q. Zhao, *Asymptotic speeds of spread and traveling waves for monotone semiflows with applications*, Comm. Pure Appl. Math. 60 (2007), 1-40.
- [43] H. Matano, *Traveling waves in spatially random media*, RIMS Kokyuroku 1337 (2003), 1-9.
- [44] G. Nadin, *The effect of the Schwarz rearrangement on the periodic principal eigenvalue of a non-symmetric operator*, SIAM J. Math. Anal. 4 (2010) 2388-2406.
- [45] G. Nadin, *Critical travelling waves for general heterogeneous one-dimensional reaction-diffusion equations*, Ann. Inst. H. Poincaré, Analyse Non Linéaire 32 (2015), 841-873.
- [46] J. Nolen, L. Ryzhik, *Traveling waves in a one-dimensional heterogeneous medium*, Ann. Inst. H. Poincaré, Analyse Non Linéaire 26 (2009), 1021-1047.
- [47] J. P. Pauwelussen, *Nerve impulse propagation in a branching nerve system: a simple model*, Physica D 4 (1981), 67-88.
- [48] P. Poláčik, *Propagating terraces and the dynamics of front-like solutions of reaction-diffusion equations on \mathbb{R}* , Mem. Amer. Math. Soc. 264 (2020), no. 1278.
- [49] A. Scheel, S. Tikhomirov, *Depinning asymptotics in ergodic media*, Patterns of Dynamics, P. Gurevich, J. Hell, B. Sandstede, and A. Scheel, editors, Springer International Publishing, Cham, 2017, 88-108.
- [50] N. Shigesada, K. Kawasaki, *Biological Invasions: Theory and Practice*, Oxford Series in Ecology and Evolution, Oxford: Oxford UP, 1997.
- [51] N. Shigesada, K. Kawasaki, E. Teramoto, *Traveling periodic waves in heterogeneous environment*, Theor. Popul. Biol. 30 (1986) 143-160.
- [52] J. Sneyd, J. Sherratt, *On the propagation of calcium waves in an inhomogeneous medium*, SIAM J. Appl. Math. 57 (1997), 73-94.
- [53] H. F. Weinberger, *On spreading speeds and traveling waves for growth and migration in periodic habitat*, J. Math. Biol. 45 (2002), 511-548.
- [54] J. X. Xin, *Existence and stability of travelling waves in periodic media governed by a bistable nonlinearity*, J. Dyn. Diff. Equations 3 (1991), 541-573.
- [55] J. X. Xin, *Existence and nonexistence of traveling waves and reaction-diffusion front propagation in periodic media*, J. Stat. Phys. 73 (1993), 893-926.
- [56] J. X. Xin, *Front propagation in heterogeneous media*, SIAM Review 42 (2000), 161-230.
- [57] J. X. Xin, J. Zhu, *Quenching and propagation of bistable reaction-diffusion fronts in multidimensional periodic media*, Physica D 81 (1995), 94-110.
- [58] A. Zlatoš, *Existence and non-existence of transition fronts for bistable and ignition reactions*, Ann. Inst. H. Poincaré, Analyse Non Linéaire 34 (2017), 1687-1705.