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Impact of initial states on the vapor hydration of iodine-bearing borosilicate glass

Haohan Zhang^{a,*}, Jean-Pierre Guin^b, Tomo Suzuki-Muresan^a, Michael Paris^c, Stéphane Gin^d,
 Abdesselam Abdelouas^a

^a SUBATECH, CNRS/IN2P3, IMT Atlantique, Université de Nantes, BP 20722, 44307 Nantes
cedex 3, France

- ^b Institut de Physique de Rennes, Univ Rennes, CNRS IPR-UMR 62051, F-35000 Rennes,
 France
- ^c Université de Nantes, CNRS, Institut des Matériaux Jean Rouxel, IMN, F-44000 Nantes,
 France
- ^d CEA, DES, ISEC, DE2D, University of Montpellier, Marcoule, F-30207 Bagnols sur Cèze,
 France
- 13

14 Abstract

15 Under the context of long-term geological disposal of nuclear waste glass, the present work 16 investigates the effect of initial states on the glass chemical durability in the vapor phase and 17 during subsequent leaching. Glass with different finishing steps of polishing and thermal 18 annealing were prepared, vapor-hydrated, and aqueous-altered. Multiple characterizations of 19 mechanical properties and structural modifications were performed to understand the 20 reduction of vapor hydration rates with annealed glasses. Discussions on the degradation of 21 mechanical properties provide insights to understanding the variations of glass chemical 22 durability, which shows a feasible approach to link mechanical and chemical studies together. 23 By looking into the release of boron in the vapor phase, similarities between glass vapor 24 hydration and glass aqueous alteration at residual rate are noticed, which allow the proposition 25 of mechanisms of boron release. Characterizations of the hydrated-then-altered glass reveal 26 the transport-limiting property of the vapor hydrated layer.

27 Keywords

28 Borosilicate glass; Vapor hydration; Chemical durability; Mechanical properties

29 1. Introduction

30 Hydration of vitrified nuclear glass in the unsaturated vapor phase is considered to be a realistic 31 scenario within the context of nuclear glass long-term disposal. The humid repository 32 environment in the deep geological disposal facilities could make the vapor hydration of 33 nuclear glass an important aging process which would influence the subsequent release of 34 radionuclides confined in the vitrified glass matrix. Various studies have been carried out to 35 investigate the nuclear glass hydration in the unsaturated vapor phase by using the same 36 autoclaves in which relative humidity is controlled by saline solution that is separated from the 37 glass sample by a Teflon holder [1–7]. Different affecting parameters have been evaluated, 38 including temperature, relative humidity (RH), pH (that is controlled by the vapor composition), 39 and glass composition [1-6]. Recent studies draw attention to the effect of initial surface state 40 (roughness, glass stress state) on the glass vapor hydration [7,8]. In the studies of hydration 41 of SON68 glass (French inactive nuclear glass) in the vapor phase, Jegou et al. [8] found a higher apparent activation energy than the one reported by Bouakkaz, revealing the impact of
 glass initial states on vapor hydration [9]. The difference in activation energies implies also the
 difference in the mechanisms of glass vapor hydration which have not been clearly established
 yet.

46 In terms of elemental releases during glass vapor hydration, previous studies showed that the 47 vapor hydration of iodine-bearing borosilicate glass (VNal glass) was accompanied by the 48 partial releases of B, Na, and I in the vapor phase, while the other network formers (Si and Al) 49 were totally retained [7]. The continuous release of boron in the vapor phase makes it possible 50 to trace the glass vapor hydration kinetics by following the boron release. In the studies of borosilicate glass alteration in the aqueous phase, boron is typically used as a tracer of glass 51 52 dissolution because of its high mobility and low retention in the altered layer [10]. However, a 53 study on the mechanism controlling the residual rate of glass aqueous alteration proposed to 54 reconsider the role of boron as a tracer because the residual rate was found to be controlled by the behavior of boron, which can, in turn, be retained in the alteration layer [11]. In the 55 borosilicate glass structure, boron could be found in either three-coordinated BO₃ or four-56 57 coordinated BO₄. As a cross-linking network former, the departure of boron is considered to 58 be responsible for both the creation of nanopores, and the formation of silanol groups which 59 may promote the re-polymerization of the remaining silicate network and thus affects 60 significantly the glass corrosion [12]. Considering the important role of boron during glass 61 aqueous alteration, its behavior may also influence the hydration of glass in the vapor phase. 62 During the alteration of borosilicate glass in the aqueous phase, the transformation from BO₄ 63 to BO₃ via hydrolysis process and the preferential release of BO₃ have been revealed by both 64 solid sample characterizations and simulation works [13–17]. Though the preferential release 65 of three-coordinated BO_3 was also suggested in the study of glass vapor hydration [7], the mechanism of boron release in the vapor phase is still unclear. On the other hand, previous 66 67 study showed a retention of about 90% iodine in the vapor-hydrated layer [7]. Considering the 68 fact that the vapor-hydrated glass will eventually end up being altered in an aqueous solution, 69 the stability of hydrated glass in contact with aqueous solution should be investigated to have 70 a comprehensive evaluation of the glass chemical durability.

71 In this study, the surfaces of monoliths of iodine-bearing borosilicate glass (VNaI) are prepared 72 with different finishing steps of polishing and thermal annealing in order to study the effect of 73 what we defined as the "initial sate" on mechanical properties, hydration and alteration kinetics. 74 The so prepared samples are firstly hydrated in the vapor phase, then altered in deionized (DI) 75 water to evaluate their chemical durability. By comparing the behavior of samples in the vapor 76 phase, the impacts of initial states on the glass hydration kinetics is discussed, which bring 77 lights to the rate-limiting mechanisms of glass vapor hydration (especially the mechanisms of 78 boron release). Solid state characterizations are performed to obtain elemental profiles and 79 structural information of the hydrated and altered layers, which suggests a transport-limiting 80 property of the pre-hydrated layer.

81 2. Materials and Methods

82 2.1. Sample preparation

83 The borosilicate glass sample VNal is provided by CEA (the samples used in the present study

- are from the same batch used in a previous study [7], with compositions given in **Table A1**, in
- 85 Supplementary data). Glass monoliths with a dimension of 10mm×5mm×1mm were cut. Glass

86 monoliths were firstly polished to 3µm by using SiC papers from P800 grade to P4000 grade 87 then with diamond suspension of 3 µm. The polished glass samples were cleaned with acetone, deionized (DI) water, and ethanol in the ultrasonic bath. The samples polished to 3µm are 88 89 denoted as P-3 type samples. To study the effect of the glass surface initial state on mechanical properties and chemical durability, some of the P-3 samples were further polished 90 91 by CeO₂ slurry to optical quality (Chemical mechanical polishing (CMP) with CeO₂). After 92 cleaning the glass by the same procedure described above, the sample was annealed at 93 540 °C for 2h then cooled down slowly to room temperature. The prepared glass samples are 94 denoted as P-3-C-A type samples. To investigate the effect of annealing on glass vapor 95 hydration, a third type of glass surface initial state was prepared by annealing the P-3 samples 96 with the above-mentioned annealing procedure (denoted P-3-A). Table 1 summarizes the samples considered in this study and the corresponding preparation procedures. P-3, P-3-A, 97 and P-C-3-A represent three initial surface states and define what is called later the pristine 98 state (or pristine glass). The term 'hydrated glass' denotes the glass samples hydrated in the 99 100 vapor phase, and 'altered glass' for samples altered in the aqueous phase.

Туре	3µm	CeO ₂	Anneal	Name	Hydrated	Characterizations/Tests
				PG-3-1	No	ToF-SIMS
				PG-3-2	No	AFM/Nano-indentation
DЭ	Voc	No	No	PG-3-3	No	Chemical durability (60 days)
F-3	165	INU	INU	PG-3-4	No	Chemical durability (2 days)
				VNal-22	Yes	ToF-SIMS/Chemical durability
				VNal-24	Yes	
P-3-C-A	Yes	Yes	Yes	PG-C-A-1	No	ToF-SIMS
				PG-C-A-2	No	AFM/Nano-indentation
				PG-C-A-3	No	Chemical durability (60 days)
				PG-C-A-4	No	Chemical durability (2 days)
				VNal-31	Yes	TEM//Nano-indentation
				VNal-32	Yes	
D_3_A	Voc	No	Voc	VNal-26	Yes	
F-3-A	165	INU	162	PG-3-A	No	AFM/Nano-indentation
Dowdor			No	PWD-PG	No	NMR
rowder			INO	PWD-H	Yes	NMR

101	Table 1	I Summary	∕ of	glass s	amples,	pre	paration	process,	and	charact	erizatior	าร
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102 **2.2. Chemical durability tests**

103 Chemical durability of pristine glass samples was tested in both vapor phase and aqueous 104 phase. In the vapor phase, hydration tests were performed at 90 °C in a stainless autoclave 105 (same experimental set-up and method as in our previous study is described in [7]). 8 mL of 106 NaCl saline solution (1.75 wt%) was used to control the relative humidity fixed to 99%. The 107 hydration kinetics was followed by applying the physical and chemical methods as detailed in 108 the previous study [7]. The physical method is based on the analysis of glass samples by FTIR 109 Spectroscopy (Bruker-Tensor 27), through which the thickness of hydrated layer could be 110 estimated (as detailed in Supplementary Note 1 in Supplementary data). The chemical method 111 is based on the analysis of saline solution by an inductively coupled plasma mass spectrometry 112 (Quadrupole ICP-MS Xseries 2, Thermoelectron).

In the aqueous phase, the chemical durability was tested by leaching the glass samples in deionized (DI) water at 50 °C in static mode. Alteration tests with two durations were carried out to determine the initial alteration rate and the elemental release behaviors during longer 116 term. The short-term alteration tests lasted for 2 days with pristine glass samples altered in 15 117 mL of DI water at 50 °C. Samplings of 0.6 mL of solution were performed after 2, 4, 8, 24 and 118 48 h of alteration. The longer-term alteration tests lasted for 60 days with pristine glass samples 119 altered in 27 mL of DI water at 50 °C. Samplings of 1 mL of solution were performed after 1, 7, 120 14, 28 and 60 days of alteration. The initial surface area to solution volume (SA/V) ratio is 3.2 121 m⁻¹ and 8.5 m⁻¹ for alteration tests of 60 days and 2 days, respectively. The ratios are chosen 122 to avoid important variations in SA/V ratios resulting from the samplings and to obtain sufficient 123 elemental concentrations in the leaching solution to facilitate ICP-MS analyses.

124 The chemical durability of the hydrated VNal sample was also tested. To verify if there was 125 initial large release of elements at the beginning of the glass aqueous alteration, 4 cycles of 126 rinsing were performed with the hydrated sample. For the rinsing procedure, the hydrated 127 sample was firstly put in 5 mL of deionized water for 1 h at 50 °C; then the hydrated glass was 128 taken out of the first rinsing solution, put in 5 mL of fresh DI water, and altered at 50 °C for 129 another 3 h; the same process was repeated with increased alteration time steps (4 h then 16 130 h). After the rinsing process, the rinsed sample was altered in 27 mL of DI water at 50 °C. 131 Samplings of 1 mL of solution were performed after 1, 7, 14, 28 and 60 days of alteration.

132 The collected leaching solution was analyzed by ICP-MS for boron, silicon, and iodine 133 concentration. Boron and silicon were analyzed in HNO₃ with beryllium and scandium as an 134 internal standard. Iodine was analyzed in NH₄OH with indium as an internal standard. The 135 limits of quantification are 0.05 μ g·L⁻¹ for boron, 6.6 μ g·L⁻¹ for silicon, and 0.01 μ g·L⁻¹ for iodine 136 with relative uncertainties of 5 to10%.

137 The kinetics of glass alteration can be indicated by the normalized mass loss of element i (*NL(i*), 138 g m⁻²):

139
$$NL(i) = \frac{C_i}{X_i \times \frac{S_A}{V}}$$
(1)

140 with C_i the concentration of element *i* in solution (g m⁻³), X_i the weight fraction of element *i* in 141 pristine glass, and *SA/V* the ratio of the glass surface area to solution volume (m⁻¹).

142 From the normalized mass values, it is possible to calculate the equivalent thickness of the 143 altered layer:

$$144 \quad ET_i = \frac{NL_i}{\rho} \tag{2}$$

145 with ρ the density of pristine glass (provided in **Table A1**).

146 2.3. Solid characterization techniques

147 **2.3.1** Atomic force microscopy (AFM) and Instrumented Nano-indentation:

The 3 types of pristine glass samples were investigated for surface roughness, hardness and reduced elastic modulus characterization. The surface roughness of the surfaces was measured by atomic force microscopy (AFM ICON, from Bruker using TAP300AI probes) using tappingTM mode, which was calibrated with silicon atomic steps. For each type of surface preparation, 4 different zones of $5 \times 5 \,\mu\text{m}^2$ were imaged. Topographical data were post treated 153 with Gwyddion software [18] and reported values of roughness (Rq) are mean values (and 154 associated standard deviation values) computed over those 4 different zones. Nano-155 indentation tests were carried out with a nano-indenter testing device (TI-950 from Bruker, with 156 nano-DMA head) equipped with a Berkovich indenter at ambient conditions (21-22 °C, 42-46% 157 relative humidity). The partial unloading (PUL) mode with 20 partial unloading segments on 158 the loading curve and maximum forces (P_{max}) of 1 mN and 9 mN was used for probing the 159 evolution of mechanical properties as a function of depth at the nanometer scale (as shown in 160 Fig. A2 in Supplementary data). For each (pristine and hydrated) sample at 1 mN (9 mN) 161 indentation maximum load, 150 (50) indents having 20 measurement steps were performed at 162 a constant loading rate for statistic, and at two different randomly chosen areas on the sample. 163 Before that the frame compliance and the contact area (A_c) of the indenter were corrected and 164 calibrated respectively according to the well accepted method described in [19] and using a 165 silica sample that was glued along with the other tested samples of this study with phenyl 166 salicylate onto the stainless-steel plate of the nano-indenter. A_c was calibrated for a 5 to 80 nm 167 contact depth.

168 The Oliver and Pharr method was applied to determine, from the load-displacement curves, 169 the hardness (*H*) and the indentation modulus E_r [19]. In the Oliver and Pharr method, the 170 contact depth (h_c) at maximum load is provided by Eq. (3) where h_{max} is the maximal 171 penetration depth measured at maximum load, P_{max} , and the right term of the equation stands 172 for the elastic deformation of the indented surface where ε is a geometrical constant ($\varepsilon = 0.75$) 173 related to the indenter geometry and *S* is the contact stiffness at unloading defined by Eq. (4):

$$174 h_c = h_{max} - \varepsilon \frac{P_{max}}{S} aga{3}$$

175
$$S = \frac{dP}{dh} = \frac{2}{\sqrt{\pi}} E_r \sqrt{A_c}$$
(4)

176 with *h* the displacement of indenter.

From h_{max} , P_{max} and S, and knowing $A_c(h_c)$ by calibration, the hardness (*H*) equals to the mean contact pressure at P_{max} and the indentation reduced modulus (E_r) may be computed by Eqs. (5) & (6) (for diamond indenter, E = 1141 GPa and Poisson's ratio v = 0.07):

$$180 \quad H = P_{max} / A_c \tag{5}$$

181
$${}^{1}/E_{r} = {}^{1}/E_{r, indenter} + {}^{1}/E_{r, sample}$$
 (6)

Note that for partial unloading mode, a value of both H and E_r is computed from each of the 20 partial unloading segments lying between 0 and P_{max} . Due to the large amount of data, post processed indentation measurements were processed using Jupyter notebook environment.

185 2.3.2 MAS NMR spectroscopy

In order to obtain the glass structural information in the pristine and hydrated glass, ¹¹B, ²³Na, and ²⁷Al Solid State Nuclear Magnetic Resonance (NMR) analyses were performed. Glass powder was prepared by crushing and sieving particle sizes between 100 and 125 µm. The glass powder was hydrated during 150 days in a climatic chamber (WEISS WKL64) set at 90 °C and RH at 95%. NMR analyses were carried out with a Bruker Avance III 500 MHz

191 spectrometer using a 2.5 mm Magic Angle Spinning (MAS) probe. The glass powder was 192 loaded into ZrO₂ rotor with Teflon caps. MAS frequency was set to 30 kHz and recycle delay set to 10 s for ¹¹B and 1 s for ²³Na and ²⁷AI. MAS spectra were acquired with a radio frequency 193 field strength of 11 kHz and with a single pulse excitation of $\pi/10$ for ¹¹B spectra, $\pi/9$ for ²³Na 194 195 spectra, and $\pi/14$ for ²⁷Al spectra. The spectra were referenced against Al(NO₃)₃ and NaNO₃ aqueous solutions and against 1 mol L⁻¹ boric acid solution (19.6 ppm) for ²⁷Al and ²³Na and 196 197 11B, respectively. ¹¹B spectra were firstly processed to remove the probe background and the 198 contribution of satellite transitions to the central band of the spectrum [20,21]. The processed 199 spectra were decomposed by using DMFit2020 freeware [22]. The simplified decomposition 200 method proposed by Jolivet et al. was used to determine the N4 ratio (fraction of four-201 coordinated boron over total boron) [21].

202 2.3.3 Time-of-Flight Secondary Ion Mass Spectroscopy

203 Glass samples were analyzed by Time-of-Flight Secondary Ion Mass Spectroscopy (ToF-SIMS) 204 for elemental profiling. The profiles of negative secondary ions were obtained by alternating analysis and abrasion cycles. Primary ions of Bi_3^+ were applied for abrasion cycles at 25 keV. 205 206 The analyzed area was 50×50 μ m². At the end of the analyses, the depth of the crater was 207 measured by a mechanical profilometer. The sputtering rate was assumed to be constant over 208 time and calculated by dividing the depth of the crater by the total abrasion time. The elemental 209 profiles were then obtained by considering the constant sputtering rate. The analysis was 210 conducted with pristine P-3 and P-3-C-A samples, hydrated VNaI-22 sample, and hydrated-211 then-altered VNal-22 sample.

212 To investigate the passivating property of the altered layer on the hydrated-then-altered VNal-213 22 sample, after the first profiling analysis with ToF-SIMS, the sample was immersed during 24 h at room temperature in an ¹⁸O labeled water doped with methylene blue (molecule with a 214 1 nm hydrodynamic diameter). Initial isotopic ratio of ¹⁸O/¹⁶O was 0.86. After 24 h' alteration, 215 216 the sample traced with ¹⁸O and methylene blue was dried by blowing air to remove water 217 adsorbed on the external surfaces, then plunged in liquid N₂ (77 K) for 10 min to freeze the 218 sample (minimize the pore water mobility). The frozen sample was mounted on a cooled 219 sample holder for ToF-SIMS profiling.

220 **3. Results**

221 **3.1. Characterization of pristine state samples**

222 **3.1.1** Roughness and mechanical properties of pristine state samples

The 3 types of pristine surface states were characterized by AFM for the initial surface roughness (*Rq*) (**Fig. 1**), which gave 5.82 ± 0.12 nm, 5.79 ± 0.37 nm, and 1.46 ± 0.35 nm for P-3, P-3-A, and P-3-C-A samples, respectively. The polishing process with CeO₂ slurry reduced significantly the surface roughness, while the effect of annealing has no noticeable effect on the surface roughness.



228

Fig. 1. AFM topography of the VNal glass surfaces as a function of surface preparation: (a) polished with a 3µm sized diamond suspension (P-3); (b) same as (a) with a subsequent thermal annealing at T_g for 2 h (P-3-A); (c) chemical-mechanical polished with CeO₂, and annealed at T_g for 2 h (P-3-C-A). On figure (b), triangles illustrate the equivalent size of the indentation imprint at h_{max} for loads equal to 1000, 300 and 85 µN. Roughness values were averaged over 4 different areas of 5x5 µm² for each sample.



235

Fig. 2. Evolution of hardness as a function of penetration depth measured by instrumented Berkovich indentation test using a partial unloading set up (20 segments) with maximum loads of 1 mN and 9 mN made on P-3, P-3-A, and P-3-C-A samples. For the 1 mN maximum load data set two areas were tested with 150 indentations imprint each; for 9 mN 50 to 60 indentations were made. Data points represent the mean value of *H* from each unloading segment (at a constant load *P*) at the corresponding mean indentation depth both plotted along with the error bars (\pm 1 Standard deviation).

243 The hardness H and the reduced modulus E_r of the 3 types of pristine glass samples were 244 determined by instrumented nano-indentation as a function of maximum penetration depth 245 (h_{max}) into the surface. Regarding the elastic property ($E_{r,sample}$, shown in **Fig. A3** in 246 Supplementary data), the surface preparation has no detectable effect on its value, without 247 any differences being largely within the error bars (± 1 standard deviation) from 20 nm up to 248 more than 250 nm for h_{max} . For depth values below 20 nm, a slight difference of few GPa is 249 noticeable between the unannealed sample P-3 and the annealed ones (the shallower the 250 depth, the larger the difference). For hardness (Fig. 2), the trend is different: P-3-C-A and P-3251 A exhibit the same behavior on the whole range of tested penetration depth, while P-3 presents 252 a rather different one. From more than 250 nm down to 150 nm, H_{P-3} has the same and 253 constant value ($H = 8.2 \pm 0.6$ GPa) as the ones measured for P-3-A and P-3-C-A samples. For 254 smaller penetration depths and as penetration depth decreases, $H_{P,3}$ deviates from the 255 constant H values measured for the 2 other samples and decreases in a non-linear way down 256 to 4.9 \pm 0.9 GPa (40% decrease). It is also worth noting that the variability of both H and E_r 257 measurements on P-3-C-A sample is much lower than the one of P-3 and P-3-A samples, it is 258 certainly a direct effect of a lower surface roughness for P-3-C-A when compared to the two 259 others.

260 **3.1.2** Elemental profiles of pristine state samples

To characterize the surface of the pristine glass chemically disturbed by the preparation process, the pristine P-3 and P-3-C-A samples were analyzed by ToF-SIMS. **Fig. 3** gives the profiles of H, B, Na, AI, and I, normalized firstly to the intensity of Siin the disturbed zone then to the intensity ratio with Si in the pristine glass (PG) (Eq. (7)).

265 Normalized intensity of element i
$$\frac{Intensity of element i}{(Intensity of element i)}$$
 (7)



266

Fig. 3. ToF-SIMS profiles of pristine (a) PG-3-1, the glass sample polished to 3 µm (type P-3)
 and (b) PG-C-A-1, the glass sample polished to 3 µm and with CeO₂ slurry then annealed
 (type P-3-C-A).

270 Si is chosen as the reference for normalization because of its high concentration in VNal glass. 271 For both samples, it can be seen from Fig. 3 that the first 20-30 nm (estimation mainly based 272 on the depth of H penetration, from a mechanical viewpoint) of the surface has been altered, 273 which is characterized by a depletion in B, Na and I, and an enrichment in H. This pre-altered 274 layer is usually referred to the Beilby layer formed during the sample polishing [23]. By looking 275 at the profiles of different elements, the estimation of the Beilby layer (or the pre-altered layer) 276 thickness could give different values. In the studies of glass vapor hydration, the commonly 277 used altered layer thickness estimation method considers the thickness to be the depth at which the concentration of B is half of its concentration in the pristine glass [5,7]. The Beilby 278 279 layer thickness based on boron profiles is about 11.4 nm for polished-only PG-3-1 sample and 280 about 9.2 nm for polished-then-annealed PG-C-A-1 sample. Considering the 10% error for the estimated thickness by ToF-SIMS [24,25], the thicknesses of the Beilby layer are similar 281 282 between the two analyzed samples. The profiles of H show a larger depth of water penetration in the polished-only sample than in the annealed sample. Besides, within the thicknesses of 283 284 the Beilby layers (11.4 nm and 9.2 nm, respectively), it could be calculated that about 75% of B, 76% of Na, and 90% of I are lost from the outmost surface of P-3 sample (PG-3-1), while 285 84% of B, 41% of Na, and 90% of I from the annealed P-3-C-A sample (PG-C-A-1). On the 286 287 other hand, the recovery of Na can be observed after sample annealing by comparing the profiles of B and Na between P-3 and P-3-C-A samples. 288

289 **3.2.** Chemical durability of glass in the vapor phase and in the aqueous phase

290 Fig. 4a shows the variations of vapor hydrated layer thickness estimated by FTIR 291 measurements and by the analysis of boron concentration in saline solution. As a general 292 observation, a good repeatability could be found between samples of the same initial surface 293 state. Vapor hydration rate can be estimated through the increase of hydrated layer thickness. 294 An initial hydration rate is calculated by considering the increase of hydrated layer thickness 295 from the beginning of the tests (t = 0 day) to the first FTIR measurement. A reduced initial 296 hydration rate is calculated between the first FTIR measurement and the end of the experiment. 297 Fig. 4b gives the vapor hydration rates for the tested samples. For both initial and reduced 298 rates, the polished-only P-3 samples showed significantly larger hydration rate than the 299 annealed P-3-C-A and P-3-A samples.



300

Fig. 4. Vapor hydration kinetics of the 3 types of glass samples: (a) variation of hydrated
 layer thickness; (b) Initial rate and reduced initial rate. Initial rate is calculated by considering
 the increase of hydrated layer thickness from the beginning of the tests (t = 0 day) to the first

IR measurement. The reduced initial hydration rate is calculated between the first IR
 measurement and the end of the experiment.

306

315

307 Fig. 5a, b & c shows the evolution of the alteration layer thickness estimated based on B, I, 308 and Si release from the longer-term (60-day) alteration tests. Similar variations of thickness could be found between P-3 and P-3-C-A samples despite of their difference in initial states. 309 310 However, for the short-term (2-day) alteration tests (Fig. 5d), greater quantity of boron was released from the non-annealed P-3 sample. Table 2 summarizes the vapor hydration rates 311 312 (based on FTIR measurements) and the aqueous alteration rates (based on boron release) for 313 the tested samples. The aqueous alteration rates are calculated by dividing the NL_B at 60 days 314 and at 2 days by the total alteration times, respectively.



Fig. 5. Data from alteration of pristine glass in DI water at 50 °C. Variation of alteration layer thickness during 60 days of alteration of PG-3-3 (type P-3) and PG-C-A-3 (type P-3-C-A), with thickness estimated based on the release of (a) boron, (b) lodine, and (c) Silicon. (d)
Variation of alteration layer thickness during 2 days of alteration of PG-3-4 (type P-3), PG-3-4 (type P-3-A), and PG-C-A-4 (type P-3-C-A), with thickness estimated based on boron release. The initial SA/V ratio is 3.2 m⁻¹ and 8.5 m⁻¹ for the longer-term alteration during 60 days and for the 2-days shorter6-term alteration, respectively.

After 126 days of hydration in the vapor phase, the hydrated VNaI-22 sample was cut into two pieces. One piece of the hydrated sample was used for the characterization of vapor hydrated layer. The other piece was rinsed and altered in DI water for 140 days. **Fig. 6** shows the release of B, I, and Si from the hydrated glass. Comparing the release of elements from the pristine glass and from the hydrated VNaI-22, important reduction in iodine and silicon release is noticed from the hydrated sample, while the release of boron from the hydrated sample is not 329 affected. The restricted release of iodine and silicon from the hydrated sample indicates a 330 passivating property of the pre-hydrated layer.





332 Fig. 6. Data from alteration of hydrated VNaI-22 in DI water at 50 °C. Normalized mass loss 333 of (a) boron and iodine, and (b) silicon from the alteration of pristine glass (PG-3-3, type P-3) 334 and from hydrated VNal-22.

Table 2 Summary of vapor myulation rates and aqueous alteration ra	Ty of vapor hydration rates and aqueous alteration rates.
---	---

Sample type	Sample name	Initial vapor hydration rate (nm d ⁻¹)	Reduced initial vapor hydration rate (nm d ⁻¹)	Short-term aqueous alteration rate (nm d ⁻¹)	Longer-term aqueous alteration rate (nm d ⁻¹)	
P-3	VNal-22	14.8 ± 2.5	5.9 ± 1.4	446 ± 50	109 ± 13	
	VNal-24	13.7 ± 2.3	5.1 ± 2.2			
P-3-C-A	VNal-31	5.2 ± 0.9	1.1 ± 0.2	26 ± 2	100 ± 12	
	VNal-32	6.7 ± 1.1	0.7 ± 0.2			
P-3-A	VNal-26	1.8 ± 0.3	0.1 ± 0.2	39 ± 3	Non analyzed	

336 Vapor hydration rates are calculated by considering the variation of physical thickness (based on FTIR measurements). Aqueous alteration rates are calculated based on boron release.

3.3. Characterization of hydrated and hydrated-then altered samples 338

339 3.3.1 Degradation of mechanical properties of the hydrated layer

340 After 209 days of hydration in the vapor phase, FTIR spectrum taken on the hydrated VNaI-31 341 gave a hydrated layer thickness of 264 ± 44 nm (FTIR spectra of hydrated and pristine glass 342 samples are given in Fig. A1 in Supplementary Data). The hydrated VNaI-31 was cut then 343 characterized for hardness and reduced elastic modulus by performing nano-indentation tests 344 at different penetration depths. The hydrated glass samples of type P-3 and P-3-A were not 345 characterized by nano-indentation tests because of the important surface roughness, which 346 would introduce a tremendously large uncertainty on the measurements.

347 Partial unloading (PUL) instrumented nano-indentation tests allow for exploring the evolution 348 of mechanical properties (H and $E_{r, sample}$) as a function of indentation depth. Results are 349 reported in Fig. 7 and in Fig. A3 in Supplementary data. For the data sets (P-3-C-A & VNal-350 31), matrices of 10x15 indentations with a maximum load of 1 mN were made at two different 351 and randomly chosen areas. 10x10 indentations matrices were also carried on at a third spot 352 location this time with a maximum load of 9 mN. H and Er, sample values computed from each

³³⁷

353 PUL segment of each of the 400 curves (8000 data points in total) are reported in Fig. 7 and

Fig. A3, respectively. On **Fig. 7** was also added previous measurements obtained for P-3-C-A and P-3 pristine surface states.



356

357 **Fig. 7.** Evolution of the sample's hardness (*H*) as a function of penetration depth (h_{max}), at each partial unloading segment (20 segments, P_{max}=1 and 9 mN) performed on P-3-C-A and 358 359 VNaI-31 samples. While only the mean value and error bars (±1 standard deviation) are reported for P-3-C-A and P-3 pristine state samples, which sets up an upper limit to the 360 361 possible values of H for the hydrated glass. For the homogeneous area of the VNaI-31 sample (purple points), all the measured values from the data set are reported along with the 362 363 mean value and error bars computed for each segment (same maximum load). Scattered data points reported in green and in blue stand for different tested areas of VNaI-31 with a 364 365 maximum load Pmax equal to 1 mN and 9 mN respectively.

366 The first remark is the large distribution of the data points for each considered depth, which 367 may be understood in terms of heterogeneity of the surface in response to indentation. Nevertheless, this large distribution lies between an upper limit set by the properties of the 368 369 pristing glass (especially mean values of H_{P-3} at low penetration depth and a lower limit. The 370 latter is defined from the measurements made at low loads on a quite homogeneous area of 371 the hydrated surface and is illustrated in Fig.7 by the purple data points for which H does not 372 depend on the indentation depth. This zone is considered to provide the mechanical properties 373 of the hydrated zone of the glass sample (H_{hyd} = 3.2 GPa and $E_{r,hyd}$ = 40.0 GPa).

374 Following the increase in penetration depth, higher H and E_{r,sample} values are found. Thus, the 375 large scattering of the data points between the upper and lower limit values of H and $E_{r, sample}$ 376 reflects the heterogeneity of the hydrated surface in terms of thickness of the hydrated layer. Nonetheless, only few measurements asymptotically reach the hardness of the pristine glass 377 378 at larger penetration depth (200-300 nm). Regarding $E_{r,sample}$, the same trend is observed (Fig. A3), yet not as clear as for H. This is not unexpected as the hydrated sample may be crudely 379 described as a soft and compliant thin hydrated layer on top of a harder and stiffer substrate. 380 381 As plasticity is more confined around the indenter than elasticity is, E_r is expected to be more affected (at a shallower indentation depth) by a substrate effect than hardness values. In other words, for very large penetration depth, when compared to the hydrated layer thickness, values of the pristine glass are expected for *H* and $E_{r,sample}$. The plots of *H*- $E_{r,sample}$ provides another way of looking at the data (see **Fig. A5** in Supplementary data), which confirms this behaviour

The inhomogeneity of the surface hydrated layer reported through mechanical properties characterization is compatible with surface observations from TEM-EDX analysis performed on a thin foil extracted from the hydrated surface (**Fig. A6 and A7** in Supplementary data). The later revealed also the inhomogeneity of the hydrated surface of the VNaI-31 sample (**Fig. A6**). Besides, no-significant change in chemical composition was observed (**Fig. A7**).

391 3.3.2 ¹¹B-NMR of pristine and hydrated glass

MAS NMR spectroscopy was used to characterize the pristine and hydrated glass powder (with
 initial particle size between 100 and 125 μm). As seen in Fig. 8.



394

Fig. 8. MAS NMR spectra of pristine and vapor hydrated VNaI glass: (a) ¹¹B spectra of
 pristine and hydrated glass; (b) ²³Na, and (c) ²⁷Al spectra normalized to the surface area, with
 hydrated glass spectra subtracted by a fraction of the pristine glass spectra.

398 The ²⁷AI MAS NMR spectrum of the pristine VNaI shows a broad asymmetric resonance due 399 to both distributions of electric field gradient (EFG) and of isotropic chemical shifts. The ²⁷Al 400 MAS NMR spectrum can be fitted by a single line ('Czsimple' model of DMFit freeware) at 62 401 ppm (δ_{iso} , isotropic chemical shift), C_Q = 4.9 MHz (mean quadrupolar constant) and FWHM(CS) 402 = 11 ppm (full width at half maximum of the Gaussian distribution of isotropic shifts). The 403 spectrum of the hydrated glass shows a slightly narrower resonance, which is due to the 404 presence of neighboring water [26]. Similar reduction of linewidths of ²⁷AI MAS NMR spectra was also observed for glass samples altered in the aqueous phase [26,27]. Spectral 405 406 decomposition gives δ_{iso} = 62 ppm, C_Q = 4.1 MHz and FWHM(CS) = 7 ppm. Hence, for both 407 pristine and hydrated glass samples, AI only occurs as AIO₄ and no evidence of the presence 408 of AIO₆ is found (expected at ~0 ppm) after alteration.

409 The ²³Na MAS NMR spectrum of the pristine VNal also shows a broad featureless asymmetric 410 resonance. The ²³Na MAS NMR spectrum can be fitted by a single line ('Czsimple' model) at 411 $\delta_{iso} = -5$ ppm with C_Q = 3.0 MHz and FWHM(CS) = 21 ppm. Here again, the ²³Na MAS NMR 412 spectrum of the hydrated glass sample exhibits a more symmetrical line shape with reduced 413 linewidth due to the presence of neighboring water. A new narrow peak is observed at about
414 2.1 ppm, which may indicate the formation of sodium-containing secondary (possibly
415 crystalline) phases.

¹¹B MAS NMR spectra show narrow and broad resonances for BO₄ (^[4]B) and BO₃ (^[3]B), 416 417 respectively [20,28,29]. The fraction of BO₄ over total boron (N_4 ratio) can be determined rather 418 easily through the area under the MAS peaks. Therefore, the simplified peak decomposition 419 approach proposed by Jolivet et al. was applied [21]. The N_4 ratio is determined to be 0.71 for 420 pristine glass and 0.83 for hydrated glass powder. Though an increase of N₄ ratio from 0.71 to 421 0.83 is found, it should be noted that the pristine glass could still be present beneath the 422 hydrated layer that covers the external part of the glass particles. As a result, the N_4 ratio of 423 the hydrated layer could be higher than 0.83. The experimental spectra of the partially hydrated 424 glass sample could be considered as the sum of the pristine one and the totally hydrated one. 425 However, the NMR spectra of a totally hydrated sample is unknown in our study. In this case, 426 only a maximum contribution from the pristine glass could be estimated by subtracting the 427 spectrum of the partially hydrated sample by the spectrum of pristine sample while maintaining 428 the positive intensity for the subtracted spectrum. The NMR spectra were firstly normalized to 429 the surface area. With the ²³Na NMR spectra, a maximum fraction of 46% is found. Interestingly, the same fraction is obtained with the ²⁷Al spectra. Taking into account the 46% of the pristine 430 glass in the partially hydrated glass sample, the N₄ ratio for the totally hydrated layer is 431 432 calculated to be 0.93, which indicates the predominance of four-coordinated boron in the 433 hydrated layer.

434 **3.3.3** Elemental profiles of hydrated and hydrated-then-altered P-3 samples

435 The hydrated VNaI-22 and the hydrated-then-altered VNaI-22 were characterized by ToF-436 SIMS for elemental profiling in the degraded layers. The profile of hydrated VNaI-22 has been 437 presented in the previous study, showing a hydrated layer of 830 nm with 9% of boron and 92% 438 of iodine retained therein [7]. To further investigate the passivating properties of the hydrated layer, the hydrated-then-altered glass was immersed during 24 h at room temperature in an 439 ¹⁸O labeled water doped with methylene blue. The molecules of methylene blue have a 440 441 hydrodynamic diameter of 1 nm and can be traced by the profile of sulfur. The penetration of 442 methylene blue could imply the presence of pores or channels with diameters larger than 1nm. 443 The change of isotopic ratio of ¹⁸O/¹⁶O could indicate the fraction of pores accessible to water 444 molecules.

445 Fig. 9a shows the ToF-SIMS profiles of hydrated-then-altered VNaI-22 glass. The glass 446 sample was firstly hydrated in the vapor phase during 126 days then altered in DI water for 447 140 days. The elemental profiles in the hydrated-then-altered layer show depletion of boron 448 and sodium and retention of iodine, which are similar to the profiles obtained on the hydrated 449 VNaI-22. Moreover, the great retention of iodine in the hydrated-then-altered layer is not 450 surprising because iodine was merely released during the alteration tests (Fig. 6). At about 451 1300 nm, a local maximum concentration of boron and iodine is observed indicating a possible 452 accumulation of boron and iodine at this depth. Based on the profile of boron, the thickness of the hydrated-then-altered layer is estimated to be 8.3 μ m. The NL_B obtained during the 453 454 alteration of the hydrated VNaI-22 indicates that an altered layer of 8.7 µm has been developed 455 after 140 days of alteration in DI water. However, it should be noted that the surface of the

456 hydrated-then-altered glass shows a non-negligible roughness, which could induce great457 uncertainties in estimating the thickness of the degraded layer.





459 Fig. 9. ToF-SIMS profiles of hydrated-then-altered VNaI glass (a) before and (b) after tracing
 460 with ¹⁸O and methylene blue.

Fig. 9b gives the ToF-SIMS profiles obtained from the hydrated-then-altered glass sample 461 462 after tracing with ¹⁸O and methylene blue. The blocking barrier of boron and iodine release 463 (implied by the local maximum concentration of the two elements) is located at about 550 nm. The thickness of the hydrated-then-altered layer is about 3.3 µm. The difference in thickness 464 465 of the hydrated layer (or the depth of the blocking barrier) and of the hydrated-then-altered 466 layer could be due to the inhomogeneity of the degraded surface layer as the elemental profiles 467 are taken at different positions on the glass sample. The profile of methylene blue in the glass 468 sample is marked by the profile of sulfur. It can be seen that the methylene blue molecules did 469 not penetrate into the hydrated layer (or the hydrated-then-altered layer), which indicates that 470 pore sizes of the hydrated layer should be smaller than 1 nm.

471 By applying the method proposed by Gin et al., the fraction of pores accessible to water 472 molecules can be calculated based on the elemental and isotopic profiles obtained on the 473 hydrated-then-altered layer traced with ¹⁸O [25]. By supposing that the pores (assumed to be 474 spherical) in the hydrated-then-altered layer were created by the release of soluble elements 475 (B, Na, and I for the current VNal sample), the total volume of pores could be calculated based 476 on the elemental profiles. The ionic volumes considered in the calculation are given in Table 477 A2 in Supplementary data. The co-existence of three-coordinated BO₃ and four-coordinated 478 BO_4 has been considered when calculating the total pore volume, with the N_4 ratio (equals to 479 0.71) determined by NMR analysis of pristine VNal glass powder. However, the modification 480 of N₄ ratio in the altered layer was not considered. With the calculated total pore volume in the 481 hydrated-then-altered layer, a hypothetical isotopic ratio of ¹⁸O/¹⁶O could be determined by 482 assuming that all the pores are filled with the ¹⁸O labelled water. On the other hand, the ToF-SIMS analysis provided experimental profile of ¹⁸O/¹⁶O in the hydrated-then-altered layer. The 483 isotopic profile was firstly corrected by fixing the isotopic ratio in the pristine glass to the natural 484 abundance of ¹⁸O (0.002). The fraction between the corrected isotopic ratio and the 485 486 hypothetical isotopic ratio could then be considered as ratio of pores that are accessible to 487 water molecules during 24 h of alteration. Fig. 10 shows the low fraction of accessible pores, 488 about 3-4%, in the first 1800 nm of the hydrated-then-altered layer. It should be noted that the

489 calculation of the accessible pores fraction considered only the oxygen exchange among water 490 molecules. In the altered layer, oxygen of water molecules may exchange continuously with 491 oxygen atoms from silanols and structural (bridging) oxygen. Nevertheless, the low exchange 492 ratio of oxygen is attributed to the self-reorganization of the altered layer, explaining its 493 passivating property [25]. For the small value of fraction of accessible pores observed in this 494 study, especially within the first 500 nm (corresponding to the vapor-hydrated layer), it suggests 495 also the blocking properties of the hydrated layer in limiting the transport of boron and iodine 496 ions with water molecules.







499 **4. Discussions**

500 **4.1 Impacts of polishing and annealing on glass initial states**

501 Characterizations of the pristine glass samples of different initial states (P-3, P-3-A, and P-3-502 C-A) showed the existence of an outmost layer that is physically and chemically modified by 503 the polishing process and by the subsequent thermal treatment. This modified layer is often 504 referred as the Beilby layer [30,31]. For oxide glasses, this layer results from both the 505 mechanical action of the slurries on the glass surface (densification, microcracks) and the 506 mechanochemical interaction of the slurry with the glass surface which lead to a gentler 507 material's removal. The glass composition, the nature of the slurry as well as the supporting 508 medium (often a liquid) will affect the relative importance of the two above-mentioned 509 mechanisms and the properties of the resulting Beilby layer. Beilby layer is reported to be 510 composed of an outmost hydrated layer followed by a densified layer at deeper depth, 511 thicknesses of which will depend on the polishing parameters such as stress applied, liquid 512 used for polishing, and nature of the particles (diamond, alumina, CeO2...). Névot and Croce 513 reported a hydrated layer of 1.1 nm and a densified sublayer of several tens of nanometres' 514 thick having a maximum density at a depth of 5.1 nm for chemical mechanical polishing (CMP) 515 conditions with CeO₂[32]. Using TEM, Toyolo et al. reported that the thickness of the modified 516 layer, depending on polishing conditions, could be up to 150 nm for polished PECVD SiO₂ film 517 deposited onto Si [33]. Further TEM investigations revealed the presence of 2 sub zones in 518 this hydrated layer: the first one, which was chemically altered, extends 3-4 nm below the 519 polished surface, while the second one whose thickness varies between 15 and 20 nm was 520 attributed to a plastically densified region. This structure of layers agrees well with values 521 previously reported using small angle X rays by Nevot and Bensch [32,34]. Therefore, as a 522 rule of thumbs, the main difference between CMP polished glass surface with CeO₂ (P-3-C-A

523 sample) and diamond polished surfaces (P-3 samples) will be for the later: a larger thickness 524 of the pre-hydrated layer, the presence of a larger and deeper densified layer along with the 525 possible existence of sub surface damages (micro and sub microcracks).

526 The next point to be discussed is the effect of a high temperature (at transition temperature, 527 T_g) thermal treatment on the Beilby layer. First, it is well established that indentation-induced 528 densification is known to be totally cured by a sub-Tg (0.9*Tg, 2h) thermal treatment leading 529 to no difference in dissolution rate at a previously densified site once thermally cured [35–37]. 530 Moreover, thermal treatment at T_g allows for both the dehydration and the dihydroxylation of 531 oxides glasses [38–40].

532 Considering the thermal cure of annealed samples, the decrease of hardness values with 533 decreasing penetration depth observed for P-3 sample only is related to the presence of a 534 modified and hydrated layer at the outmost surface. Furthermore, for soda lime and borosilicate 535 glasses, it was shown that the effect of the hydrated layer on the apparent elastic modulus 536 ceased at a penetration depth equal to the thickness of the hydrated layer [41]. From Fig. A3, 537 P-3 behavior departs slightly from the one of P-3-A samples within a penetration depth of about 538 20 nm, which would set an approximate value of this layer thickness. The later value agrees 539 well with the P-3 ToF-SIMS measurements reported in Fig. 3a, especially when remembering 540 that water content is known for altering the mechanical properties of glass [42,43] and 541 considering that hydrogen concentration is at a maximum value down to a depth of 15 nm 542 followed by a smooth decrease reaching the bulk glass value at a depth of approximately 25 543 nm.

544 **4.2** Influence of initial states on glass vapor hydration

545 Glass samples (P-3-C-A) showed a net decrease in surface roughness after polishing with 546 CeO_2 slurry, while no effect of annealing was observed in modifying the surface roughness (P-547 3 and P-3-A samples). When comparing the chemical durability in the unsaturated vapor phase 548 and in the aqueous phase (with short-term alteration tests), a major difference was observed 549 between the annealed and the unannealed glass samples. The two types of the annealed 550 samples (P-3-C-A, P-3-A) showed no noticeable difference in hydration kinetics, although a 551 smaller surface roughness was found on the P-3-C-A glass. These observations suggest thus 552 minor effect of surface roughness on the kinetics of glass vapor hydration.

553 As discussed above (Section 4.1), the polished P-3 glass may present a larger and deeper 554 densified layer along with the possible existence of subsurface damages (micro and sub 555 microcracks). The plastically densified layer is considered to be formed due to the compressive 556 stress created by contacting with the abrasive particles [44,45]. Molecular dynamics simulation 557 works concerning the chemical mechanical polishing of silicon showed the generation of huge 558 hydrostatic pressure in the local area [46]. Such densification behavior induced by hydrostatic 559 stresses was found to be comparable with the behavior of glass under indentation [47]. 560 Consequently, we may consider that the densified layer formed during polishing process is similar to the densified zone created by indentation. Besides, locally enhanced dissolution 561 562 rates were found at indentation sites on the glass surfaces (silica and floated window glasses) 563 [37]. Therefore, the densified layer formed during glass polishing may contribute to the 564 enhanced short-term alteration rate (Fig. 5d). As mentioned above, the indentation-induced 565 densification can be totally cured by the thermal annealing, which is also compatible with the 566 reduction of alteration rates as shown in the short-terme alteration data obtained with annealed 567 glass (Fig. 5d).

568 Underlying the macroscopic densification are the modification in glass structural network. For 569 silicate glasses densified under high pressure or indentation, it has been found the decrease 570 of Si-O-Si angles between silicon tetrahedrons and the increase in Q₂ species (Si-Ostretching in a structural unit with n non-bridging O atoms) at the expense of Q₃ species [48] 571 572 (Taniguchi 2004). Hence, the decrease in inter-tetrahedral angle could increase effectively the 573 chemical activity of Si-O-Si bonds, resulting in a lower resistance to hydrolysis [49]. These 574 structural changes can be identified by Raman spectroscopy, where the changes in mean angle between silicon tetrahedrons could be characterized by a shift of Raman spectra towards 575 576 higher frequency at round 500 cm⁻¹ [48,50–52]. However, the structural modifications on the 577 annealed samples in this study were not successfully characterized with Raman spectrometry 578 (Fig. A8 in Supplementary data), which might be due to the relatively thin densified layer in front of the spatial resolution of micro-Raman probe. Besides, the presence of boron in the 579 580 glass network leads to the mixture of B-O-B, B-O-Si, and Si-O-Si bonds, which makes it more 581 complex for interpreting Raman spectra.

582 Moreover, the role of boron should not be neglected because the hydration of glass in the vapor phase is always accompanied by boron release. Studies showed that the thermal 583 annealing of the glass may lead to the increase of BO4 and decrease of non-ring BO3 in the 584 585 glass matrix [13,53]. The current work confirms the preferential release of three-coordinated BO₃ in the vapor phase by performing ¹¹B NMR analysis. Recent simulation work on 586 587 borosilicate glass-water interface showed that BO₃ units are more accessible to water attacks than BO₄ units, which is surrounded by a higher amount of sodium [17]. As a result, the lower 588 589 release of boron from the annealed glass may be induced by the increase in the fraction of 590 tetrahedral BO₄ whose release may require the prior hydrolysis of the B-O-Si bond.

591 **4.3 Insights into the mechanism of boron release**

592 Considering the possible hydrolysis process during glass vapor hydration, especially the 593 hydrolysis of B-O-Si bonds, linear regressions between chemical thickness (thickness of 594 hydrated layer estimated through boron release) and time (or square root of time) were 595 performed. As shown in Fig. 11, the increases of hydrated layer thickness for P-3 samples (VNaI-22 and VNaI-24) are proportional to the square root of time, which suggests a diffusion-596 597 controlled process of boron release [10,54]. By applying Fick's second law, the apparent 598 diffusion coefficients for boron were calculated to be $(2.5 \pm 1.2) \times 10^{-20} \text{ m}^2 \text{ s}^{-1}$ for P-3 glass 599 samples. For the P-3-C-A samples (VNaI-31 and VNaI-32), linear increases of hydrated layers 600 can be observed, which implies a hydrolysis-dominating process. As a matter of fact, the 601 constant rate of elemental release has been attributed to the contribution of hydrolysis reaction 602 in the studies of glass aqueous alteration [55]. It is therefore possible that the polishing process 603 has lowered the resistance for the hydrolysis of B-O-Si liaisons in the P-3 glass network, while cured subsequently by thermal annealing in the P-3-A and P-3-C-A glasses. Consequently, 604 the hydrolysis process became a rate-limiting step for the hydration of annealed glass in the 605 vapor phase. The fact that a higher initial aqueous alteration rate is observed with the non-606 607 annealed glass samples is compatible with the hypothesis of decrease of B-O-Si hydrolysis 608 barrier, because the hydrolysis process is considered to be a pre-dominant mechanism in 609 controlling the glass initial alteration rate [56,57]. To summarize, the linear regressions 610 performed on the P-3 and P-3-C-A samples suggest the co-presence of diffusion and hydrolysis processes in controlling the release of boron during glass vapor hydration. 611 612 Noteworthy, between these two processes, the pre-dominant one should depend strongly on 613 the surrounding environment of boron, which consequently controls the kinetics of its release.



614



Fig. 11. Linear fitting of chemical thickness: (a) chemical thickness of VNaI-22 and VNaI-24 616 versus square root of time; (b) chemical thickness of VNaI-31 and VNaI-32 versus time.

617 Considering the fact that both diffusion and hydrolysis are mechanisms that have been well established to explain the release of elements in the aqueous phase [11,12,25,54,58-62], the 618 619 proposition of the mechanisms of boron release during glass vapor hydration could be inspired 620 from these existing studies. Recent study on glass aqueous alteration in the silica saturated 621 condition also revealed the important role of boron in controlling the residual rate of glass 622 dissolution [11]. By looking into the experimental observations obtained from the experiments 623 of glass aqueous alteration in the silica-saturated condition and of glass hydration in the 624 unsaturated vapor condition, similarities in boron behaviour could be found. At residual 625 alteration rate, boron (about 12%) and alkali elements (about 25-30%) were found to be 626 retained in the altered gel layer [11,54]. Similar boron and alkali retention ratios were found 627 during the glass vapor hydration [7]. In the silica saturated condition, the release of Si from the 628 glass is largely limited while the release of boron continues [63]. In the unsaturated vapor 629 condition, only the release of boron in the vapor phase was detected. In a study of water 630 adsorption in the unsaturated vapor condition, Asay et al. suggested the presence of three 631 types of water layer structures at relative humidity (RH) higher than 60%: (1) the ice-like layer 632 covering the glass surface, preferably formed by interacting with the silanol groups; (2) a 633 transitional layer that links the ice-like layer and the outmost liquid water layer; (3) the outmost 634 liquid water layer with totally liquid-like configured water layer [64,65]. At 99%RH (the tested experimental condition), tens of monolayers of water molecules could be adsorbed at the 635 636 surface, including monolayers of liquid-like water [64]. Due to the limited quantity of liquid-like 637 water on the glass surface, the fast increase or even the saturation of boron and silica in the 638 liquid-like water should be expected. Previous study showed that boron could be easily 639 evaporated in the vapor condition [7]. Following the fast increase of boron concentration in the 640 liquid-like water layer, the evaporation of boron could be a rate-limiting factor controlling the 641 rate of glass vapor hydration. Considering the similarity of glass hydration in the vapor phase 642 and glass alteration at residual rate and the previous discussions on the effect annealing, we 643 propose the following mechanisms of glass vapor hydration involving the release of boron:

644 Adsorption of water on the glass surface; a)

645 Diffusion of water in the glass network and release of Na by inter-diffusion; b)

- 646 c) Proton (occurred after the inter-diffusion and loss of Na⁺) catalyses the hydrolysis of B-O647 Si bond (which is less resistant than Si-O-Si bond [58]);
- d) Diffusion of BO_{3, aq} towards the glass surface and accumulation in the liquid-like adsorbed
 water layer;
- e) Evaporation of boron from the liquid-like water layer.

651 **4.4 Transport-limiting property of vapor hydrated layer**

It could be implied from the ToF-SIMS profiles of the hydrated glass that the pores in the 652 653 hydrated layer are mainly formed by the departure of boron and sodium. The release of boron 654 through hydrolysis and diffusion processes results in the formation of silanol groups. The 655 depolymerization of the silicate structure was found in the previous study through Raman 656 analysis [7]. Besides, the degradation of the physical properties in hardness and elastic moduli 657 is coherent with the depolymerization of the glass network. Though the preferential release of 658 boron leads to the increase in the fraction of four-coordinated BO₄ in the hydrated glass, it is 659 not contradictory to the apparent network depolymerization because of the small concentration 660 of boron left in the hydrated layer.

661 Alteration of the pre-hydrated VNal glass in DI water shows reduced release of iodine and 662 silicon, which suggests the role of hydrated layer as a blocking barrier. ToF-SIMS 663 characterizations of the hydrated-then-altered layer confirm the accumulation and blockage of 664 boron and iodine at the interface between the pre-hydrated layer and the newly developed 665 altered layer during glass aqueous leaching. During the leaching of the pre-hydrated VNal 666 glass, the release of boron was not influenced by the presence of the pre-hydrated layer. 667 Besides, the release of boron from the pristine glass and from the hydrated glass are both 668 proportional to the square-root of time, suggesting a diffusion-controlled process. 669 Consequently, the pre-hydrated layer could be considered as a transport-limiting barrier for the 670 release of elements during the subsequent leaching: with the ingress of water molecules 671 through the hydrated layer into the pristine glass, the B-O-Si bonds are hydrolyzed to allow the 672 detachment of boron from the glass network; then the release of the detached boron into the 673 leaching solution is limited by the process of diffusion through the pre-hydrated layer. In terms 674 of iodine retention in the hydrated layer during glass leaching test, aside from the transport-675 limiting effect, it would be necessary to investigate the structural and chemical environment of 676 iodine, as well as sodium and aluminum which are usually found in the surrounding 677 environment of iodine [66,67].

678 **4.5 Degradation of glass mechanical properties: thickness of the hydrated layer**

679 The decrease of hardness with penetration depth of P-3 sample is well described by the 680 Korsunsky [68] model (Fig. A9 in Supplementary data) which has proven itself to be well suited 681 for describing the substrate effect on thin film mechanical properties measurements (H, E_{L}) 682 sample) by instrumented nano indentation. This confirms, if needed, the fact that a thicker softer 683 and certainly more compliant modified layer exists at the outmost surface of P-3 sample as 684 reported in a previous study [69]. Unfortunately, this model with the data in our possession does not allow for the computation of a layer thickness without a proper calibration, which can 685 686 be done from the thickness (≈ 20 nm) of the hydrated layer provided by the ToF-SIMS 687 technique.

688 Regarding the VNaI-31 sample for indentation depths ranging from 10 nm to 400 nm, the whole 689 range of possible H and E_r, sample values (from fully hydrated layer to bulk glass) are covered on the VNaI-31 sample surface (Fig. 7). This means that the thickness of the altered layer that 690 691 has developed at the surface of the sample is highly heterogeneous in thickness. Although the 692 Korsunsky [68] model stands for a homogeneous thin film on a homogeneous substrate 693 geometry, a rough estimation of a possible range of hydrated layer thicknesses is possible. 694 The procedure is explained, and results are presented in Supplementary data (Note 5). As 695 mentioned above, the Korsunsky model is calibrated with P-3 sample, then the model is 696 applied to VNaI-31 data, the film thickness (t) becoming in this case the adjusting parameter. 697 From the $H(h_{max})$ values coming from 16 indentation curves randomly selected in the data set, 698 values of t ranging from 30 to 800 nm are obtained (Fig. A10 in Supplementary data). The 699 majority of the values lies between 100 and 200 nm (~50%), the rest is distributed as follow: 700 20-50 nm (25%), 200-500 (19%), 500-800(8%). The same method was applied to the VNaI-31 701 homogeneous area, possible values of t were found to range between 600 and 800 nm. It is 702 interesting to note that P-3 H- h_{max} curve (Fig. 7) fits almost perfectly the outmost H_{sample} values 703 of VNal sample, the two having similar hydrated layer thicknesses (20-30 nm range), although 704 their hardness value is different 4.8 GPa for P-3 and 3.2 GPa for VNaI-31. Those values of t 705 compare well with other techniques used in this study in terms of: heterogeneity of the hydrated 706 layer that has developed on VNaI-31 sample; thicknesses of the hydrated layer for P-3 sample 707 but also for VNaI-31 sample for which the 264 nm obtained by FTIR compares relatively well 708 with the values reported above.

709 5. Conclusion

710 By studying the mechanical properties and boron release behavior between glass samples 711 possessing different initial states, we propose a mechanism of boron release involving the 712 hydrolysis of B-O-Si bonds followed by the diffusion of three-coordinated boron to the 713 glass/vapor interface. The preferential release of three-coordinated boron is confirmed by ¹¹B 714 NMR analysis. The reduction of glass vapor hydration rate observed for the annealed glass is 715 supposed to be caused by the increase in both the fraction of four-coordinated BO4 and the 716 energy barrier for B-O-Si hydrolysis. As a result, the surrounding environment of boron in the 717 glass matrix could strongly influence the kinetics of glass vapor hydration. Glass structural 718 modifications indicated by the variations in hardness and elastic modulus bring lights into the 719 understanding of the glass chemical durability, which provides a feasible approach to link the 720 mechanical studies and chemical studies together.

721 Characterizations of the hydrated glass indicate degradation of the mechanical properties by 722 glass hydration. In the subsequent leaching tests with the hydrated glass sample, the hydrated 723 layer is found to be passivating. The transport-limiting property of the hydrated-layer is verified 724 by the characterizations of the hydrated-then-altered layer. With the tested samples, the initial 725 surface state (including the surface roughness and the surface chemical composition changes 726 induced by annealing) shows negligible effect on the glass vapor hydration. Similarities 727 between glass vapor hydration and glass alteration at residual rate are discussed, which 728 provide insights into further studies to understand the mechanism of glass vapor hydration.

729 CRediT authorship contribution statement

Haohan Zhang: Conceptualization, Formal analysis, Investigation, Methodology, Writing –
 original draft, Writing – review & editing. Jean-Pierre Guin: Formal analysis, Investigation,
 Methodology, Funding acquisition, Writing – original draft, Writing – review & editing. Tomo

Suzuki-Muresan: Conceptualization, Supervision, Funding acquisition, Writing – review & editing.
Michael Paris: Methodology, Formal analysis, Validation, Writing – review & editing.
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738 **Declaration of Competing Interest**

739 The authors declare that they have no known competing financial interests or personal 740 relationships that could have appeared to influence the work reported in this paper.

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752 **References**

- J. Neeway, A. Abdelouas, B. Grambow, S. Schumacher, C. Martin, M. Kogawa, S.
 Utsunomiya, S. Gin, P. Frugier, Vapor hydration of SON68 glass from 90 °C to 200 °C:
 A kinetic study and corrosion products investigation, J. Non. Cryst. Solids. 358 (2012)
 2894–2905. https://doi.org/10.1016/j.jnoncrysol.2012.07.020.
- A. Ait Chaou, A. Abdelouas, Y. El, C. Martin, A.A. Chaou, A. Abdelouas, Y. El Mendili,
 R. Bouakkaz, C. Martin, The French SON68 glass vapor hydration under different
 atmospheres, Procedia Mater. Sci. 7 (2014) 179–185.
 https://doi.org/http://dx.doi.org/10.1016/j.mspro.2014.10.024.
- A. Ait Chaou, A. Abdelouas, Y. El Mendili, R. Bouakkaz, S. Utsunomiya, C. Martin, X.
 Bourbon, Vapor hydration of a simulated borosilicate nuclear waste glass in unsaturated conditions at 50 °C and 90 °C, RSC Adv. 5 (2015) 64538–64549.
 https://doi.org/10.1039/c5ra12384d.
- A. Ait Chaou, A. Abdelouas, Y. El Mendili, C. Martin, The role of pH in the vapor
 hydration at 175 °C of the French SON68 glass, Appl. Geochemistry. 76 (2017) 22–
 35. https://doi.org/10.1016/j.apgeochem.2016.11.006.
- S. Narayanasamy, P. Jollivet, N. Godon, F. Angeli, S. Gin, M. Cabié, J.
 Cambedouzou, C. Le Guillou, A. Abdelouas, Influence of composition of nuclear waste
 glasses on vapor phase hydration, J. Nucl. Mater. 525 (2019) 53–71.
 https://doi.org/10.1016/j.jnucmat.2019.07.015.
- [6] S. Narayanasamy, P. Jollivet, L. Sessegolo, F. Angeli, A. Abdelouas, Influence of temperature and relative humidity on vapor hydration of an AVM nuclear waste glass, J. Nucl. Mater. 543 (2021) 152571. https://doi.org/10.1016/j.jnucmat.2020.152571.

- H. Zhang, T. Suzuki-Muresan, Y. Morizet, S. Gin, A. Abdelouas, Investigation on
 boron and iodine behavior during nuclear glass vapor hydration, Npj Mater. Degrad. 5
 (2021) 10. https://doi.org/10.1038/s41529-021-00157-6.
- [8] C. Jegou, S. Narayanasamy, F. Angeli, Short communication on the Influence of the temperature between 30 and 70°C on the hydration of SON68 nuclear waste glass in a vapour phase, J. Nucl. Mater. 545 (2021) 152738.
 https://doi.org/10.1016/j.jnucmat.2020.152738.
- R. Bouakkaz, Altération aqueuse et hydratation en phase vapeur duverre SON68 à basse température (35-90 ° C), Ecole des Mines de Nantes, 2014.
- P. Frugier, S. Gin, Y. Minet, T. Chave, B. Bonin, N. Godon, J.E. Lartigue, P. Jollivet, A.
 Ayral, L. De Windt, G. Santarini, SON68 nuclear glass dissolution kinetics: Current
 state of knowledge and basis of the new GRAAL model, J. Nucl. Mater. 380 (2008) 8–
 21. https://doi.org/10.1016/j.jnucmat.2008.06.044.
- [11] S. Gin, X. Guo, J.M. Delaye, F. Angeli, K. Damodaran, V. Testud, J. Du, S. Kerisit,
 S.H. Kim, Insights into the mechanisms controlling the residual corrosion rate of
 borosilicate glasses, Npj Mater. Degrad. 4 (2020) 41. https://doi.org/10.1038/s41529020-00145-2.
- M. Collin, M. Fournier, P. Frugier, T. Charpentier, M. Moskura, L. Deng, M. Ren, J. Du,
 S. Gin, Structure of International Simple Glass and properties of passivating layer
 formed in circumneutral pH conditions, Npj Mater. Degrad. 2 (2018) 4.
 https://doi.org/10.1038/s41529-017-0025-y.
- F. Angeli, T. Charpentier, P. Jollivet, D. de Ligny, M. Bergler, A. Veber, S. Gin, H. Li,
 Effect of thermally induced structural disorder on the chemical durability of
 International Simple Glass, Npj Mater. Degrad. 2 (2018).
 https://doi.org/10.1038/s41529-018-0052-3.
- 800 [14] G. Geneste, F. Bouyer, S. Gin, Hydrogen-sodium interdiffusion in borosilicate glasses
 801 investigated from first principles, J. Non. Cryst. Solids. 352 (2006) 3147–3152.
 802 https://doi.org/10.1016/j.jnoncrysol.2006.04.023.
- 803 [15] H. Aréna, R. Podor, H.P. Brau, J. Nelayah, N. Godon, M. Cabié, E. Garcès, C.
 804 Mansas, D. Rébiscoul, Characterization of the boron profile and coordination in altered
 805 glass layers by EEL spectroscopy, Micron. 141 (2021).
 806 https://doi.org/10.1016/j.micron.2020.102983.
- P. Zapol, H. He, K.D. Kwon, L.J. Criscenti, First-principles study of hydrolysis reaction barriers in a sodium borosilicate glass, Int. J. Appl. Glas. Sci. 4 (2013) 395–407.
 https://doi.org/10.1111/ijag.12052.
- 810 [17] H. Jabraoui, T. Charpentier, S. Gin, J.-M. Delaye, R. Pollet, Atomic Insights into the
 811 Events Governing the Borosilicate Glass–Water Interface, J. Phys. Chem. C. 125
 812 (2021) 7919–7931. https://doi.org/10.1021/acs.jpcc.1c00388.
- 813 [18] D. Nečas, P. Klapetek, Gwyddion: an open-source software for SPM data analysis,
 814 Open Phys. 10 (2012) 181–188. https://doi.org/doi:10.2478/s11534-011-0096-2.
- 815 [19] W.C. Oliver, G.M. Pharr, An improved technique for determining hardness and elastic modulus using load and displacement sensing indentation experiments, J. Mater. Res.
 817 7 (1992) 1564–1583. https://doi.org/10.1557/jmr.1992.1564.

- F. Angeli, T. Charpentier, D. De Ligny, C. Cailleteau, Boron speciation in soda-lime
 borosilicate glasses containing zirconium, J. Am. Ceram. Soc. 93 (2010) 2693–2704.
 https://doi.org/10.1111/j.1551-2916.2010.03771.x.
- [21] V. Jolivet, L. Jossé, M. Rivoal, M. Paris, Y. Morizet, L. Carole, T. Suzuki-Muresan,
 Quantification of boron in aluminoborosilicate glasses using Raman and 11 B NMR, J.
 Non. Cryst. Solids. 511 (2019) 50–61.
 https://doi.org/10.1016/j.jnoncrysol.2018.12.038.
- Bujoli, Z. Gan, G. Hoatson, Modelling one- and two-dimensional solid-state NMR
 spectra, Magn. Reson. Chem. 40 (2002) 70–76. https://doi.org/10.1002/mrc.984.
- T. Suratwala, W. Steele, L. Wong, M.D. Feit, P.E. Miller, R. Dylla-Spears, N. Shen, R.
 Desjardin, Chemistry and Formation of the Beilby Layer during Polishing of Fused
 Silica Glass, J. Am. Ceram. Soc. 98 (2015) 2395–2402.
 https://doi.org/10.1111/jace.13659.
- 832 [24] M. Py, Study of interfaces and nanometric structures by ToF-SIMS : upon a spatially 833 resolved quantitative analysis, Université de Grenoble, 2011.
- S. Gin, M. Collin, P. Jollivet, M. Fournier, Y. Minet, L. Dupuy, T. Mahadevan, S. Kerisit,
 J. Du, Dynamics of self-reorganization explains passivation of silicate glasses, Nat.
 Commun. 9 (2018) 1–9. https://doi.org/10.1038/s41467-018-04511-2.
- F. Angeli, M. Gaillard, P. Jollivet, T. Charpentier, Influence of glass composition and alteration solution on leached silicate glass structure: A solid-state NMR investigation, Geochim. Cosmochim. Acta. 70 (2006) 2577–2590.
 https://doi.org/10.1016/j.gca.2006.02.023.
- 841 [27] T. Ohkubo, Y. Iwadate, K. Deguchi, S. Ohki, T. Shimizu, Changes in surface structure 842 of sodium aluminoborosilicate glasses during aqueous corrosion analyzed by using 843 NMR, J. Phys. Chem. Solids. 77 (2015) 164–171. 844 https://doi.org/10.1016/j.jpcs.2014.10.018.
- 845 [28] W.J. Dell, P.J. Bray, S.Z. Xiao, 11B NMR studies and structural modeling of
 846 Na2OB2O3SiO2 glasses of high soda content, J. Non. Cryst. Solids. 58 (1983) 1–16.
 847 https://doi.org/10.1016/0022-3093(83)90097-2.
- L.S. Du, J.F. Stebbins, Nature of silicon-boron mixing in sodium borosilicate glasses: A
 high-resolution 11B and 17O NMR study, J. Phys. Chem. B. 107 (2003) 10063–10076.
 https://doi.org/10.1021/jp034048I.
- [30] F.W. Preston, The structure of abraded glass surfaces, Trans. Opt. Soc. 23 (1922)
 141. https://doi.org/10.1088/1475-4878/23/3/301.
- 853 [31] Beilby G. T., Surface flow in crystalline solids under mechanical disturbance, Proc. R.
 854 Soc. London. 72 (1904) 218–225. https://doi.org/10.1098/rspl.1903.0045.
- L. Névot, P. Croce, Caractérisation des surfaces par réflexion rasante de rayons X.
 Application à l'étude du polissage de quelques verres silicates, Rev. Phys. Appliquée.
 15 (1980) 761–779. https://doi.org/10.1051/rphysap:01980001503076100.
- J.A. Trogolo, K. Rajan, Near surface modification of silica structure induced by
 chemical/mechanical polishing, J. Mater. Sci. 1994 2917. 29 (2004) 4554–4558.
 https://doi.org/10.1007/BF00376278.

- 861 [34] W. Bensch, W. Bergholz, An FT-IR study of silicon dioxides for VLSI microelectronics,
 862 Semicon Sci. Tech. 5 (1990) 421. https://doi.org/10.1088/0268-1242/5/5/008.
- 863 [35] S. Yoshida, J.-C. Sanglebœuf, T. Rouxel, Quantitative evaluation of indentationinduced densification in glass, J. Mater. Res. 20 (2005) 3404–3412.
 865 https://doi.org/10.1557/jmr.2005.0418.
- [36] J.D. MACKENZIE, High-Pressure Effects on Oxide Glasses: II, Subsequent Heat
 Treatment, J. Am. Ceram. Soc. 46 (1963) 470–476. https://doi.org/10.1111/J.11512916.1963.TB13777.X.
- 869 [37] Y.-F. Niu, K. Han, J.-P. Guin, Locally Enhanced Dissolution Rate as a Probe for
 870 Nanocontact-Induced Densification in Oxide Glasses, Langmuir. 28 (2012) 10733–
 871 10740. https://doi.org/10.1021/la300972j.
- [38] L.T. Zhuravlev, The surface chemistry of amorphous silica. Zhuravlev model, Colloids
 Surfaces A Physicochem. Eng. Asp. 173 (2000) 1–38. https://doi.org/10.1016/S09277757(00)00556-2.
- [39] H. Liu, D. Ngo, M. Ren, J. Du, S.H. Kim, Effects of surface initial condition on aqueous corrosion of glass—A study with a model nuclear waste glass, J. Am. Ceram. Soc. 102 (2019) 1652–1664. https://doi.org/10.1111/jace.16016.
- [40] J. Banerjee, V. Bojan, C.G. Pantano, S.H. Kim, Effect of heat treatment on the surface
 chemical structure of glass: Oxygen speciation from in situ XPS analysis, J. Am.
 Ceram. Soc. 101 (2018) 644–656. https://doi.org/10.1111/JACE.15245.
- [41] D.R. Tadjiev, R.J. Hand, P. Zeng, Comparison of glass hydration layer thickness
 measured by transmission electron microscopy and nanoindentation, Mater. Lett. 64
 (2010) 1041–1044. https://doi.org/10.1016/J.MATLET.2010.02.004.
- W. -T Han, M. Tomozawa, Indentation Creep of Na2O · 3SiO2 Glasses with Various
 Water Contents, J. Am. Ceram. Soc. 73 (1990) 3626–3632.
 https://doi.org/10.1111/J.1151-2916.1990.TB04268.X.
- [43] N.M. Keulen, Indentation Creep of Hydrated Soda–Lime Silicate Glass Determined by Nanoindentation, J. Am. Ceram. Soc. 76 (1993) 904–912.
 https://doi.org/10.1111/J.1151-2916.1993.TB05314.X.
- In Section 2018
 In Structure induced by chemical/mechanical polishing, J. Mater. Sci. 29 (1994) 4554–4558.
 In Structure induced by chemical/mechanical polishing, J. Mater. Sci. 29 (1994) 4554–4558.
 In Structure induced by chemical/mechanical polishing, J. Mater. Sci. 29 (1994) 4554–4558.
- [45] J.C. Lambropoulos, S. Xu, T. Fang, Constitutive law for the densification of fused
 silica, with applications in polishing and microgrinding, J. Am. Ceram. Soc. 79 (1996)
 1441–1452. https://doi.org/10.1111/j.1151-2916.1996.tb08748.x.
- [46] X. Han, Y. Hu, S. Yu, Investigation of material removal mechanism of silicon wafer in the chemical mechanical polishing process using molecular dynamics simulation method, Appl. Phys. A Mater. Sci. Process. 95 (2009) 899–905.
 https://doi.org/10.1007/s00339-009-5097-2.
- 900 [47] H. Ji, V. Keryvin, T. Rouxel, T. Hammouda, Densification of window glass under very
 901 high pressure and its relevance to Vickers indentation, Scr. Mater. 55 (2006) 1159–
 902 1162. https://doi.org/10.1016/j.scriptamat.2006.08.038.

- 903 [48] T. Deschamps, C. Martinet, J.L. Bruneel, B. Champagnon, Soda-lime silicate glass
 904 under hydrostatic pressure and indentation: a micro-Raman study, J. Phys. Condens.
 905 Matter. 23 (2011) 035402. https://doi.org/10.1088/0953-8984/23/3/035402.
- 906 [49] T.A. Michalske, B.C. Bunker, Slow fracture model based on strained silicate
 907 structures, J. Appl. Phys. 56 (1984) 2686–2693. https://doi.org/10.1063/1.333789.
- 908 [50] A. Winterstein-Beckmann, D. Möncke, D. Palles, E.I. Kamitsos, L. Wondraczek, A
 909 Raman-spectroscopic study of indentation-induced structural changes in technical
 910 alkali-borosilicate glasses with varying silicate network connectivity, J. Non. Cryst.
 911 Solids. 405 (2014) 196–206. https://doi.org/10.1016/j.jnoncrysol.2014.09.020.
- 912 [51] S. Peuget, J.M. Delaye, C. Jégou, Specific outcomes of the research on the radiation
 913 stability of the French nuclear glass towards alpha decay accumulation, J. Nucl. Mater.
 914 44 (2014) 76–91. https://doi.org/10.1016/j.jnucmat.2013.09.039.
- 915 [52] D. Vandembroucq, T. Deschamps, C. Coussa, A. Perriot, E. Barthel, B. Champagnon,
 916 C. Martinet, Density hardening plasticity and mechanical ageing of silica glass under
 917 pressure: a Raman spectroscopic study, J. Phys. Condens. Matter. 20 (2008) 485221.
 918 https://doi.org/10.1088/0953-8984/20/48/485221.
- 919 [53] N. Stone-Weiss, E.M. Pierce, R.E. Youngman, O. Gulbiten, N.J. Smith, J. Du, A. Goel,
 920 Understanding the structural drivers governing glass–water interactions in borosilicate
 921 based model bioactive glasses, Acta Biomater. 65 (2018) 436–449.
 922 https://doi.org/10.1016/j.actbio.2017.11.006.
- P. Frugier, S. Gin, J.E. Lartigue, E. Deloule, SON68 glass dissolution kinetics at high
 reaction progress: Mechanisms accounting for the residual alteration rate, Mater. Res.
 Soc. Symp. Proc. 932 (2006) 305–312.
- 926 [55] M.I. Ojovan, R.J. Hand, N. V. Ojovan, W.E. Lee, Corrosion of alkali-borosilicate waste
 927 glass K-26 in non-saturated conditions, J. Nucl. Mater. 340 (2005) 12–24.
 928 https://doi.org/10.1016/j.jnucmat.2004.10.095.
- J.D. Vienna, J. V. Ryan, S. Gin, Y. Inagaki, Current understanding and remaining
 challenges in modeling long-term degradation of borosilicate nuclear waste glasses,
 Int. J. Appl. Glas. Sci. 4 (2013) 283–294. https://doi.org/10.1111/ijag.12050.
- 932 [57] S. Gin, J.-M. Delaye, F. Angeli, S. Schuller, Aqueous alteration of silicate glass: state
 933 of knowledge and perspectives, Npj Mater. Degrad. 5 (2021) 42.
 934 https://doi.org/10.1038/s41529-021-00190-5.
- 935 [58] B.C. Bunker, G.W. Arnold, D.E. Day, P.J. Bray, The effect of molecular structure on borosilicate glass leaching, J. Non. Cryst. Solids. 87 (1986) 226–253.
 937 https://doi.org/10.1016/S0022-3093(86)80080-1.
- B.C. Bunker, Molecular mechanisms for corrosion of silica and silicate glasses, J. Non.
 Cryst. Solids. 179 (1994) 300–308. https://doi.org/10.1016/0022-3093(94)90708-0.
- [60] S. Gin, P. Jollivet, G. Barba Rossa, M. Tribet, S. Mougnaud, M. Collin, M. Fournier, E. Cadel, M. Cabie, L. Dupuy, Atom-Probe Tomography, TEM and ToF-SIMS study of borosilicate glass alteration rim: A multiscale approach to investigating rate-limiting mechanisms, Geochim. Cosmochim. Acta. 202 (2017) 57–76.
 https://doi.org/10.1016/j.gca.2016.12.029.
- 945 [61] P. Frugier, T. Chave, S. Gin, J.E. Lartigue, Application of the GRAAL model to
 946 leaching experiments with SON68 nuclear glass in initially pure water, J. Nucl. Mater.

- 947 392 (2009) 552–567.
- 948 http://www.sciencedirect.com/science/article/pii/S0022311509005996.
- 949 [62] S. Gin, A.H. Mir, A. Jan, J.M. Delaye, E. Chauvet, Y. De Puydt, A. Gourgiotis, S.
 950 Kerisit, A General Mechanism for Gel Layer Formation on Borosilicate Glass under
 951 Aqueous Corrosion, J. Phys. Chem. C. 124 (2020) 5132–5144.
 952 https://doi.org/10.1021/acs.jpcc.9b10491.
- J. Neeway, A. Abdelouas, B. Grambow, S. Schumacher, Dissolution mechanism of the
 SON68 reference nuclear waste glass: New data in dynamic system in silica
 saturation conditions, J. Nucl. Mater. 415 (2011) 31–37.
 https://doi.org/10.1016/j.jnucmat.2011.05.027.
- 957 [64] D.B. Asay, S.H. Kim, Evolution of the adsorbed water layer structure on silicon oxide
 958 at room temperature, J. Phys. Chem. B. 109 (2005) 16760–16763.
 959 https://doi.org/10.1021/jp0530420.
- 960 [65] D.B. Asay, A.L. Barnette, S.H. Kim, Effects of surface chemistry on structure and 961 thermodynamics of water layers at solid-vapor interfaces, J. Phys. Chem. C. 113 962 (2009) 2128–2133. https://doi.org/10.1021/jp806815p.
- 963 [66] B.J. Riley, M.J. Schweiger, D.S. Kim, W.W. Lukens, B.D. Williams, C. Iovin, C.P.
 964 Rodriguez, N.R. Overman, M.E. Bowden, D.R. Dixon, J. V. Crum, J.S. McCloy, A.A.
 965 Kruger, Iodine solubility in a low-activity waste borosilicate glass at 1000 °c, J. Nucl.
 966 Mater. 452 (2014) 178–188. https://doi.org/10.1016/j.jnucmat.2014.04.027.
- 967 [67] D.A. McKeown, I.S. Muller, I.L. Pegg, Iodine valence and local environments in
 968 borosilicate waste glasses using X-ray absorption spectroscopy, J. Nucl. Mater. 456
 969 (2015) 182–191. https://doi.org/10.1016/j.jnucmat.2014.09.033.
- 970 [68] A.M. Korsunsky, M.R. McGurk, S.J. Bull, T.F. Page, On the hardness of coated
 971 systems, Surf. Coatings Technol. 99 (1998) 171–183. https://doi.org/10.1016/S0257 972 8972(97)00522-7.
- 973 [69] B.L. Hackett, A.A. Wereszczak, G.M. Pharr, Effect of aqueous-based mechanical
 974 polishing on the nanoindentation response of borosilicate glasses, Int. J. Appl. Glas.
 975 Sci. 10 (2019) 302–306. https://doi.org/10.1111/IJAG.13108.

976