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Molecular Dynamics Simulations of Gas-Expanded Liquids

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Abstract

Gas-expanded Liquids (GXL's) are liquids whose volume is expanded by a pressurized gas. The incorporation of the gaseous compound and the expansion of the liquid phase lead to changes in solvation and transport properties of the new expanded phase.

This work aims at developing a Molecular Dynamics (MD) computational study for liquid methyl and ethyl acetates when expanded by CO_2 at pressures up to 120 bars and at temperatures between 298 and 393 K. Especially, density and viscosity of the liquid expanded phase were calculated by MD simulations. Good agreement with existing literature data has been obtained ascertained the validity of the method

Keywords: molecular dynamics, alkyl acetates, viscosity, density

1. Introduction

With the perspective of finding alternative benign media for different applications, solvent engineering represents an interesting alternative to propose innovative solutions to industrial problems. Gas-expanded liquids (GXLs), in particular CO₂-expanded liquids (CXLs), which are liquids whose volume is expanded by pressurized CO₂, represent a very promising strategy to perform solvent engineering.

A number of properties, including solvation properties, can be tuned by the addition of CO_2 to a liquid solvent (Jessop and Subramaniam, 2007) allowing the use of CXLs as "solubility switchers". In many cases, solutes that are soluble in CO_2 -expanded solvents will precipitate when CO_2 is removed and vice versa. For example, the addition of CO_2 to a biphasic system of a fluorinated solvent and an organic solvent causes the phases to become miscible (West et al., 2004). Opposite effect is seen in cases of water-soluble organics and an aqueous phase, where the addition of CO_2 forces the phases to split (Lu et al., 2004). Both systems may have applications in catalyst recycle and/or product recovery. Knowledge of phase equilibrium properties but also of transport properties is evidently important to design processes exploiting the peculiar behavior of these two phase systems.

Moreover, high-pressure experimental determinations of density and viscosity of systems involving hazardous compounds and/or at harsh conditions of pressure and temperature are usually problematic. As an alternative, numerical experiments with the help of molecular simulation, like Molecular Dynamics (MD) have gained interest in recent years as the performance of computational tools is continuously improving. In particular, MD

simulations have gained enormous progress in terms of 1) simulation methods, 2) understanding of complex phenomena, and 3) ability to model extremely complex systems.

MD simulations accuracy is highly dependent on the proper representation of the interaction forces acting on the molecules. This depends on the choice of a suitably parametrized force field and on a statistically representative sampling of the states of the system, with the help of Newton's equations integrator algorithm. Previous works on CO₂-expanded acetates (Smith et al., 1998; Byun et al., 2006; Kwon et al., 2007; Aida et al., 2010) describe some of these system properties, such as the Vapor-Liquid Equilibrium curves and the expanded phase density. However, no data were found about the viscosity of the alkyl acetates/CO₂ expanded phase, probably due to the difficulty of high-pressure expanded phase determinations.

Hence, in this work MD simulations were performed to numerically predict the value of the viscosity of CO₂-expanded acetates. The MD method and its accuracy were first evaluated for other CO₂-expanded solvents before predicting the densities and the viscosities of systems involving methyl and ethyl acetate. Indeed, Acetone (C_3H_6O) and Decane ($C_{10}H_{22}$) where chosen for this task for two main reasons: The available CO₂-expanded phase viscosity data and the presence of hydrocarbon chain in decane and the carbonyl group in acetone, both groups found in alkyl acetates.

2. Methods

2.1. Theory background

The fluctuation-dissipation theorem explains how transport coefficients associated to irreversible processes can be described using reversible microscopic fluctuations (Kubo, 1966). The Green-Kubo (Green, 1954; Kubo, 1957) relations give the exact mathematical expression for transport coefficients in terms of integrals of time-correlation functions. In the standard Green-Kubo method shear viscosity can be calculated by integrating the autocorrelation function of the stress tensor over time (1):

$$\eta = \frac{-\frac{1}{kBT}}{\frac{1}{kT}} \lim_{\tau \to \infty} \int_{0}^{\tau} \frac{\mathrm{iot}}{\mathrm{d}t \langle P_{xy}(t) P_{xy}(0) \rangle}$$
(1)

Where $\langle P_{xy}(t)P_{xy}(0)\rangle$ is the correlation function of the *xy* component of the stress tensor, *V* is the volume of the simulation box, k_B is the Boltzmann constant, *T* is the absolute temperature and *t* is the time.

2.2. Simulation details

The required liquid phase composition data at equilibrium was extracted from literature and from our own determinations.

Bond	Angle	Improper	Dihedral (hybrid)
$\frac{\mathbf{ble 1. Amb}}{\mathbf{bond}} = \frac{\mathbf{r}}{\mathbf{r}}$	$E = K(\theta - \theta^{0})$	$\frac{d \text{ pairwise}}{m \text{ proper}} OS \begin{pmatrix} d \text{ pairwise} \\ m \text{ proper} \end{pmatrix}$	$\overline{K} = K(X - X_o)^2$
VdW : L-J 12-6		Coulombic	$E = \sum k_i [1.0 + \cos(nib - di)]$
$\mathbf{W}: \mathbf{L} - \mathbf{J}_{12-6}$ $LI(r) = 4\epsilon \left[\left(\right. \right] \right]$	$\left(\frac{\sigma}{2}\right)^{\frac{12}{12}} = \left(\frac{\sigma}{2}\right)^{\frac{6}{12}} + s(r)$	$\frac{d}{dt} = \frac{cq_iq_j}{cq_iq_j}$	i=1,m

MD computations were carried out using the Materials Processes and Simulations (MAPS) Suite from Scienomics and in the Large-scale Atomic/Molecular Massively Parallel Simulator (LAMMPS) code from Sandia National Laboratories. The "Amber_Cornell_Extended" force field was selected (Table 1).

Atom	AM1	ESP	Atom	AM1	ESP
Acetone			Decane		
С	-0,2041	-0,522162	С	-0,21	-0,326006
C1	0,5631	0,811667	C1	-0,159	0,213133
C2	-0,2041	-0,512053	C2	-0,158	0,05664
0	-0,5311	-0,580332	C3	-0,158	-0,085466
Н	0,063033	0,141684	C4	-0,158	0,11713
H1	0,063033	0,131853	C5	-0,158	0,12006
H2	0,063033	0,131853	C6	-0,158	-0,087939
H3	0,063033	0,138492	C7	-0,158	0,057575
H4	0,063033	0,129499	C8	-0,159	0,213154
H5	0,063033	0,129499	C9	-0,21	-0,324668
			Н	0,072	0,068357
Methyl Acetate (MA)			H1	0,072	0,068324
С	-0,1501	-0,664509	H2	0,072	0,075387
C1	0,6291	0,915207	H3	0,078	-0,040085
0	-0,4449	-0,433813	H4	0,078	-0,040075
C2	0,1267	0,0011	H5	0,079	-0,018183
01	-0,541	-0,594817	H6	0,079	-0,018193
Н	0,0757	0,180868	H7	0,079	-0,001635
H1	0,0757	0,178806	H8	0,079	-0,001689
H2	0,0757	0,17922	H9	0,079	-0,034637
H3	0,0617	0,084701	H10	0,079	-0,034645
H4	0,0447	0,076499	H11	0,079	-0,035151
H5	0,0447	0,076736	H12	0,079	-0,035188
			H13	0,079	-0,001125
Ethyl Acetate (EA)			H14	0,079	-0,001192
С	-0,1501	-0,603183	H15	0,079	-0,018348
C1	0,6311	0,910216	H16	0,079	-0,018373
0	-0,4469	-0,51646	H17	0,078	-0,040213
C2	0,1384	0,356665	H18	0,078	-0,040178
C3	-0,1001	-0,34453	H19	0,072	0,075719
01	-0,544	-0,613058	H20	0,072	0,068759
Н	0,0517	0,011202	H21	0,072	0,06875
H1	0,0517	0,011164			
H2	0,0757	0,16283			
H3	0,0757	0,167183			
H4	0,0757	0,167011			
H5	0,0467	0,078193			
H6	0,0477	0,106401			
H7	0,0477	0,106364			

The geometries for fitting the partial charges for all the alkyl-acetates molecules were optimized at the HF/6-31G+** level by using DFT quantum calculations. The partial charges were then fitted using two separate methods. The first one used the Electrostatic Potential method (ESP) on the optimized geometries at the HF/6-31G+** level. The second one used the AM1-BBC method, which includes the optimization of the initial geometry at the AM1 level, and after the calculation of the Mulliken charges, the BCC bond correction was applied (Semi empirical Quantum Mechanics). The fitted charges are displayed in Table 2.

After repeated tests, a number of molecules per box and a finite time of equilibration, production and integration were defined. The simulated boxes were set with periodic boundaries in all directions containing a total of 1000 molecules in the Isobaric-Isothermal ensemble. The simulations were performed as follows: NPT equilibration and density production runs of 1,200,000 steps. Afterwards, NVT equilibration runs for 400,000 steps were followed by NVE viscosity production runs for 1,200,000 steps. All runs were performed with a 1fs timestep. The L-J and Coulombic cut-off was set to 8 Å with a switching distance of r =12 Å. The simulations were started with randomly assigned velocities, the total momentum of the system was set to zero and Newton's equations of motion were solved with the EOS super-calculator of CALMIP.

3. Results and Discussion

3.1. Density

Densities obtained by MD are shown in Figure 1 as function of CO_2 molar fraction at different temperatures and pressures.



Figure 1. CO₂-expanded Density (g/cm³) vs CO₂ molar fraction. Refs a) (Sih and Foster 2008) b) (Cullick and Mathis 1984).

In Figures 1 c) and d) two distinct zones can be observed. Firstly, the density increases monotonically as the CO₂ solubilizes. Then, after a certain point ($x_{CO2} = 0.7$), higher CO₂ molar fraction produces a drastic volume increment and the consequent decrease in density. Error bars correspond to Standard Deviation. This behavior has been explained by an increase in free volume observed in some CXLs (Li and Maroncelli, 2006), which would be responsible for their fluidity and solvation properties. Indeed, the initial solvent intermolecular spaces can first host the CO₂ to some extent, producing an imperceptible

increment in volume. Then, after a certain threshold in CO₂, its solubilisation induces a rapid volume expansion.

3.2. Viscosity

Calculated viscosities are shown in Figure 2 as function of CO₂ molar fraction at different temperatures and pressures. Literature data is plotted when available. Figure 2 shows a) CO₂+Acetone and b) CO₂+Decane systems for a single T with both charges sets; and c) CO₂+MA and d) CO₂+EA systems viscosity at different temperatures as a function of x_{CO2} . For these two systems, no literature data was found concerning viscosity. Standard deviation for all systems falls under 5 10⁻¹⁰ due to integration over time smoothing. The behavior is nevertheless similar to other of CXL's (Li and Maroncelli 2006). In these systems, the viscosity decreases as a quasi linear function of CO₂ mole fraction. This general behavior is explained by the decrease of electrostatic forces between the solvent molecules due to the interposing CO₂ molecules, clustering around carbonyl groups. Calculating the intermolecular distances reveal that CO₂ molecules act as a Lewis acid and tend to gather around the carbonyl group of the acetates, at the free electrons pair of the oxygen (Aida et al. 2010).



Figure 2. CO₂-expanded Viscosity (cP) vs CO₂ molar fraction. Refs a) (Sih and Foster 2008) b) (Cullick and Mathis 1984), c) (Lide 2004), d) (Trostin 1989).

4. Conclusion

MD simulations have been used to compute densities and viscosities of CO_2 -expanded methyl and ethyl acetates; good agreement with existing literature data has been obtained. In this work, prediction of the densities and viscosities of the expanded liquid phases were

possible by this technique after good parameterization and by using the adapted model for CO₂.

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