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1	Efficient Mass Conservative Numerical Model for
2	Solving Variably Saturated Groundwater Flow
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9	
10	Summary
11	A multi-dimensional mass conservative numerical method, particularly suitable for limited
12	computational resources, is developed for solving transient variably saturated groundwater
13	flow problems. The Richards equation is discretized spatially with a finite element method
14	and temporally with an implicit Euler scheme, in which mass-conservative and mass-lumping
15	techniques are used to keep the numerical simulation stable. In addition, the stiffness and
16	mass matrices involved are approximated in a way to guarantee less computational effort. To
17	confirm the accuracy and the efficiency of this code, we verified it using benchmark tests
18	using one, two and three-dimensional problems. The present model is also applied to a real
19	field case problem, where its superiority is clearly demonstrated. The code achieved reliable
20	results for each problem.
21	Keywords

22 Groundwater Flow, Finite Element, Modeling, Richards Equation, Variably Saturated Flow

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#### 23 1 Introduction

Flow in a variably saturated porous medium is governed by the Richards equation, a highly
nonlinear PDE. Few analytical solutions using simplifying conditions or restrictive
hypotheses were presented (Polubarinova, 1962; Teloglou et al, 2008; Moutsopoulos, 2013;
Tracy, 1995, 2006, 2007; Hayek, 2015). Nevertheless, the most attempts to solve saturatedunsaturated flow problems are based on numerical techniques.

Several models, using the numerical approach of Celia et al. (1990), were developed (Clement 29 et al, 1994; Van dam and Feddes, 2000; Kuznetsov et al, 2012; Herada, 2014). Belfort et al. 30 31 (2013) presented a review of several estimations of the equivalent conductivity and analyzed 32 the efficiency of the different averaging fashions when solving unsaturated flow. Fahs et al (2009) used the method of lines combined with the mixed finite element discretization to 33 solve variably saturated flow in porous media. Younes et al (2013) studied the monotonicity 34 of a finite volume-based method for solving flow in variably saturated and heterogeneous 35 porous media. Hassane Maina and Ackerer (2017) discussed Newton-Raphson method and 36 time-stepping strategies for solving the mixed form of Richards equation. Koohbor et al 37 (2020) solved Richards equation in the context of fractured porous media. Zha et al. (2017) 38 39 presented a modified Picard iteration scheme for overcoming drawbacks arising when simulating infiltration into dry soils. Farthing and Ogden (2018) and Zha et al (2019) 40 presented a state-of-the-art review of the numerical solution of the Richards equation. 41 42 Numerous aspects were discussed and challenges underlined related to equivalent hydraulic conductivity, computational efficiency or choice of primary variable among others. 43 Other models were also developed (Thoms et al, 2006; Šimůnek et al, 2006) in view to solve 44 3D flow in variably saturated porous media. These numerical models proceed by solving the 45 resulting equations system subsequent to various approximation methods. However, these 46 codes may be costing in CPU time when much finer temporal and spatial discretizations are 47

necessary, or in case of applications at the watershed scale, e. g., for climate change purposes 48 49 when huge equations systems are to be solved. Drawbacks to the numerical resolution attempt of the Richards equation may be exacerbated, when computational resources are 50 limited. Alternatives are proposed to avoid high CPU time costs of fully 3D numerical modes 51 52 (Yakirevich et al, 1998; Paulus et al, 2013; Kuznetsov et al, 2012; Xu et al, 2012). These techniques act by lowering dimensions of the fully 3D Richards equation in both unsaturated 53 54 and saturated domains or only in the unsaturated one when the saturated flow is modeled in its whole dimensionality. As a consequence the number of unknowns to be solved decrease 55 significantly and the CPU time as well. However, doing so may constitute a source of 56 57 inaccuracies that can affect the hydraulic heads as well as the fluxes in the separated domains. 58 Indeed, as it is pointed out by Twarakavi et al. (2008), only solving the fully 3D Richards equation for the whole continuous flow domain is able to provide accurate solutions for 59 hydraulic heads and subsequent fluxes in both unsaturated and saturated compartments. 60 However, the most obstacle that limits the applicability of such models for regional scales is 61 their computational demand and time cost. 62

In this paper, a multi-dimensional numerical model based on the mixed form of Richards 63 equation, particularly suitable for limited computational resources, is developed to solve flow 64 65 equation in both unsaturated and saturated zones under unsteady and steady regimes. The numerical model is Galerkin FE based and uses the Celia et al's (1990) first-order time 66 integration method to solve the derived algebraic equations system. The use the Richrads 67 equation may be inappropriate with coarse grids because of violation of the assumption of 68 representative elementary volume which can result in loss of accuracy (Zha et al, 2019). The 69 use of the fine grids, however, may complicate the resolution or simply make it infeasible due 70 to high computational coast. In order to circumvent these difficulties, the proposed numerical 71

numerical scheme approximates the matrices involved in a newly fashion to guarantee lesscomputational effort, when solving the involved PDE.

The paper is organized as follows: section 2 describes the flow model used to solve variably saturated porous media and reports the most used retention curves equations. Section 3 is dedicated to the numerical procedure adopted and the proposed approximation of the matrices obtained. In section 4 benchmark tests against analytical solutions show the present model capabilities to provide high level of accuracy. In addition, other tests available in the literature are also examined and the efficiency of the present model is demonstrated. Finally, section 5 summarizes the concluding remarks.

#### 81 2. Mathematical statement

The Richards equation can be expressed on three usual forms: the head-based form, the moisture-based form and under a mix form using the moisture and the head for the temporal and the spatial terms respectively. Discussion about these forms can be found in Fahs et al (2009), Farthing and Ogden (2018) and Zha et al (2019).

The mixed form of Richards' equation, modeling water flow in a variably saturated porous
medium, is obtained by combining Darcy's law and continuity equation and is expressed as:

88 
$$\nabla \cdot \left(\mathbf{K}_{s}K_{r}(h)\nabla h\right) - q = \frac{\partial\theta}{\partial t}$$
(1)

89 Where  $\mathbf{K}_s$  is the saturated hydraulic conductivity tensor,  $K_r(h)$  is the relative conductivity, q90 is the source/sink per volume of aquifer,  $\theta$  stands for the volumetric moisture content and 91 h is the hydraulic head.

Solving Richards' equation necessitates eliminating one of the two dependent variables, water content and the pressure head  $\Psi = h - z$  (where z is elevation). This is often done by introducing constitutive equations, namely the water retention curves that express the water content and the relative conductivity, as functions of the pressure head. Various formulations
can be found in the literature. In this work some of them will be used ranging from simplest to
very complex ones and are listed below.

Based on the model of Irmay (1954), the moisture content, as in Tracy (2006), is expressed
as:

100 
$$\theta = \begin{cases} \theta_r + (\theta_s - \theta_r) e^{\alpha \psi} & \psi < 0 \\ \theta_s & \psi \ge 0 \end{cases}$$
(2)

101 Where  $\theta_s$  and  $\theta_r$  are saturated and residual moisture content respectively.

102 Gardner (1958) expressed the relative conductivity as an exponential function of the pressure103 head:

104 
$$k_r(\psi) = \begin{cases} e^{\alpha \psi} & \psi < 0\\ 1 & \psi \ge 0 \end{cases}$$
(3)

105 The moisture content and the relative conductivity formulated by Van Genuchten (1980) are106 very popular and are written as:

107 
$$\theta = \begin{cases} \frac{(\theta_s - \theta_r)}{\left(1 + (\alpha |\psi|)^n\right)^m} + \theta_r & \psi < 0 \\ \theta_s & \psi \ge 0 \end{cases}$$
(4)

108

 $k_{r}(\boldsymbol{\psi}) = \begin{cases} \Theta^{\frac{1}{2}} \left( 1 - \left( 1 - \Theta^{\frac{1}{m}} \right) \right)^{2} & \boldsymbol{\psi} < 0 \\ 1 & \boldsymbol{\psi} \ge 0 \end{cases}$ (5)

109 Where  $\Theta = \frac{\theta - \theta_r}{\theta_s - \theta_r}$  expresses the effective saturation and  $\alpha$ , *n* and  $m = \frac{n-1}{n}$  are constants

110 depending upon the soil properties.

111 Brooks and Corey (1964) provided also equations describing hydraulic properties of soils,

112 expressed originally as functions of the moisture content  $\theta$ :

113 
$$\psi = \psi_s \left(\frac{\theta - \theta_r}{\theta - \theta_s}\right)^{-\frac{1}{\lambda}}$$
(6)

114 
$$k_r(\theta) = \left(\frac{\theta - \theta_r}{\theta - \theta_s}\right)^{\frac{2+3\lambda}{\lambda}}$$
(7)

115 Where  $\Psi_s = \Psi(\theta_s)$ , is the pressure head at saturation,  $\theta_r$  is the residual water moisture 116 content and  $\lambda$  is a parameter describing the pore-size distribution index of the medium. 117 Other relationships, linking moisture content and relative conductivity to the pressure head, 118 are presented by Haverkamp et al. (1977):

119 
$$\theta = \begin{cases} \frac{A(\theta_s - \theta_r)}{1 + (100|\psi|)^B} + \theta_r & \psi < 0\\ \theta_s & \psi \ge 0 \end{cases}$$
(8)

120 
$$k_r(\psi) = \begin{cases} \frac{C}{C + (100|\psi|)^D} & \psi < 0\\ 1 & \psi \ge 0 \end{cases}$$
(9)

121 Where A, B, C and D are constants to be determined experimentally for the soil type.

#### 122 3. Numerical solution

#### 123 3.1. Finite element model

To solve the flow equation (1), the finite element Galerkin method with linear basis functions
is used to discretize the whole domain of flow (saturated and unsaturated), along with an
implicit finite difference scheme for time. Linear, rectangular and hexahedral elements are
used in the codes for 1D, 2D and 3D features. Discretization leads to the following system of
nonlinear equations:

129 
$$\mathbf{G}(\mathbf{h}^{k+1})\mathbf{h}^{k+1} + \mathbf{M}(\mathbf{h}^{k+1}) - \mathbf{q}(t^{k+1}) = 0$$
(10)

where  $\mathbf{h}^{k+1}$  is the vector of nodal head values at time  $t^{k+1}$ , *k* is a time step index. **G** is the conductance (stiffness) matrix, **M** is a vector containing the moisture content derivative with respect to time and **q** is a vector containing the boundary conditions and flow rates pumped or injected.

134 3.2. Linearization technique

Applied to (10), the Picard scheme, where the soil moisture at the current iteration level is approximated as a truncated expansion of Taylor series about the pressure head at the last iteration level (Celia et al, 1990), gives the following system of linear equations to be solved for  $\mathbf{h}^{k+1,m+1}$ , the vector of nodal hydraulic heads at the current iteration level; superscript *m* is the iteration index:

140 
$$\mathbf{G}(\mathbf{h}^{k+1,m})\mathbf{h}^{k+1,m+1} + \mathbf{C}(\mathbf{h}^{k+1,m})\frac{\mathbf{h}^{k+1,m+1} - \mathbf{h}^{k+1,m}}{\Delta t^{k+1}} + \mathbf{M}(\mathbf{h}^{k+1,m}) - \mathbf{q}(t^{k+1}) = 0$$
(11)

141 Stiffness matrix **G**, mass matrix **C** and vector **M** depend upon the hydraulic head.

142 The entries of these matrices and the elements of the vector are defined as follows:

143 
$$G_{IJ} = \sum_{e} \int_{V^{e}} k_{r}^{\ t} \nabla b_{i} \mathbf{K}_{s} \nabla b_{j} dx dy dz$$
(12)

144 
$$C_{IJ} = \sum_{e} \int_{V^e} S_g(\psi) b_i b_j dx dy dz$$
(13)

145 
$$C_{II} = \sum_{e} \int_{V^e} S_g(\psi) b_i dx dy dz$$
(14)

146 Where  $S_g(\psi) = \frac{\theta}{\theta_s} S_s + \frac{d\theta}{d\psi}$  indicates the general storage coefficient and  $S_s$  stands for the

147 specific storage coefficient.

148 Equation (14) gives the coefficients of the mass matrix **C** if mass lumping technique is used.

149 
$$M_{I} = \sum_{e} \int_{V^{e}} \frac{d\theta}{dt} b_{i} dx dy dz$$
(15)

150 Chord slope is used to perform the moisture content derivative with respect to time as151 follows:

152 
$$\frac{d\theta}{dt} = \frac{\theta^{k+1,m} - \theta^k}{\Delta t^{k+1}}$$
(16)

153 *I* and *J* represent the global nodes, while  $b_i$  and  $b_j$  are basis functions related, respectively, to 154 local nodes *i* and *j* belonging to the element *e*. Gaussian quadrature rule is applied to perform 155 the integrals.

#### 156 3.3. Solution procedure

To solve iteratively the equation (11), the matrices involved have to be updated at each 157 iteration level, making the solution process very costly on a CPU time basis, particularly for 158 159 large scale problems. To make the procedure fast, approximations are proposed for the matrix coefficients (Aharmouch and Amaziane, 2012). The idea is to split the coefficients into to 160 quantities, one constant, computed once and stored, and the other varying with iteration, 161 162 needing to be updated. The authors applied the method to solve 3D steady state groundwater flow problems only in the saturated zone and neglect the flow in the unsaturated region. 163 However, the flow modeling in the vadose zone constitutes an important issue in hydrology 164 and plays a key role for sustainable groundwater use, particularly in dry regions, by 165 evaluating the water transfer rates to the water table during recharge and also by assessing the 166 167 water losses during dry periods. For these purposes, retention curves equations have to be used to characterize the unsaturated medium. 168

In the present work, the stiffness entries are approximated as in the work of Aharmouch and Amaziane (2012), but with incorporating the retention curves equations, aforementioned, to solve correctly the flow in the vadose zone as well as for the saturated one.

172 The stiffness entries (equation 12) are approximated for 1D, 2D and 3D features as:

173 
$$G_{IJ} \cong k^* \sum_{e} \int_{-1}^{1} {}^t \nabla b_i \mathbf{K}_s b_j d\xi$$
(17)

174 
$$G_{IJ} \cong k^* \sum_{e} \int_{-1}^{1} \int_{-1}^{1} \nabla b_i \mathbf{K}_s b_j d\xi d\eta$$
(18)

175 
$$G_{IJ} \cong k^* \sum_{e} \int_{-1}^{1} \int_{-1}^{1} \int_{-1}^{1} \nabla b_i \mathbf{K}_s \nabla b_j d\xi d\eta d\zeta$$
(19)

176 
$$k^* = \frac{1}{N_G . n_{eS}} \sum_{e=1}^{n_{eS}} \sum_{n_g=1}^{N_G} k_r(e, n_g) . \det[\mathbf{J}](e, n_g)$$
(20)

The mass matrix entries for the distributed form (equation 13) can be readily approximated as
for the stiffness coefficients with making the necessary change. However, the lumped form
(equation 14) is used only, because of its property to guarantee non oscillatory solution as it
was demonstrated by Celia et al. (1990).

181 The mass matrix entries using mass lumping technique (equation 14) are approximated for182 1D, 2D and 3D features as:

183 
$$C_{II} \cong S_g^* \sum_e \int_{-1}^{1} b_i d\xi$$
 (21)

184 
$$C_{II} \cong S_g^* \sum_e \int_{-1}^{1} \int_{-1}^{1} b_i d\xi d\eta$$
 (22)

185 
$$C_{II} \cong S_g^* \sum_{e} \int_{-1}^{1} \int_{-1}^{1} \int_{-1}^{1} b_i d\xi d\eta d\zeta$$
(23)

186 
$$S_{g}^{*} = \frac{1}{N_{G}.n_{eS}} \sum_{e=1}^{n_{eS}} \sum_{n_{g}=1}^{N_{G}} S_{g}\left(e, n_{g}\right) \det[\mathbf{J}]\left(e, n_{g}\right)$$
(24)

187  $\xi, \eta$  and  $\zeta$  are local coordinates,  $n_{es}$  is the number of elements sharing the nodes *I* and *J*, 188  $n_s$  indicates Gauss integration points,  $N_G (= 2, 4, 8)$  stands for their numbers, depending on 189 the dimensionality of the problem and det[J] $(e, n_s)$  is the Jacobian matrix determinant related 190 to Gauss integration point  $n_s$  and element *e*.

As it was pointed out in the work of Aharmouch and Amaziane (2012),  $k^*$  expresses the arithmetic mean of the term  $k_r$ .det[J] over the integration points involved in the elements that share the nodes *I* and *J*. Similarly, it can be seen that  $S_g^*$  is the arithmetic mean of the term  $S_g$ .det[J] over the integration points involved in the elements that share the nodes *I* and *J*. Only the hydraulic properties (relative conductivity and general storage coefficient) have to be updated during the iteration process, which guarantees gain in CPU time.

Time stepping is performed as in the work of Paniconi and Putti (1994). Time step sizes 197 during transient simulations are dynamically recalculated depending on the convergence 198 199 behavior of the nonlinear procedure. The simulations start with a certain time step  $\Delta t_0$  which can be augmented, by a magnificence factor  $\Delta_{mag}$ , to a certain  $\Delta t_{max}$ , or reduced, using a 200 reduction coefficient  $\Delta_{red}$  to a minimum value,  $\Delta t_{min}$ , depending on the number of iterations 201 202 allowed in each nonlinear iteration and according to some performance criteria. If the convergence is achieved in less than maxit<sub>1</sub> iterations, the time step in increased and remains 203 unchanged if the number of iterations is within maxit<sub>1</sub> and maxit<sub>2</sub> iterations and reduced if 204 205 convergence iterations needed are more than maxit<sub>2</sub>. The solution can also be recomputed at the current time, if convergence is not accomplished among the specified criteria, when maxit 206 207 limit is exceeded, by reducing the time step and going back to the old status ("back stepping"). 208

209 The infinity norm  $(l_{\infty})$  of the convergence error is adopted as stopping criterion of the

210 nonlinear iterative process, thus the convergence is achieved when  $\|\mathbf{h}^{k+1,m+1} - \mathbf{h}^{k+1,m}\|_{\infty} < tol_{is}$ 

satisfied. In the field case example, the convergence is checked using a relative norm based

212 stopping criterion. The convergence is, therefore, achieved when  $\frac{\left\|\mathbf{h}^{k+1,m+1} - \mathbf{h}^{k+1,m}\right\|_2}{\left\|\mathbf{h}^{k+1,m}\right\|_2} < tol$  is

satisfied, where  $\|\cdot\|_2$  expresses the  $(l_2)$  norm (the square root of the sum of squares of a vector components).

Owing to that the systems of linear equations obtained with the Picard finite element approximation are sparse and symmetric, conjugate gradient solvers are particularly indicated in such situation. In this work preconditioned conjugate gradient method with incomplete Cholesky decomposition preconditioner (ICCG) is used to solve the resulting linear systems for 2D and 3D problems, whereas tridiagonal direct algorithm is used for solving those obtained for 1D configuration.

#### 4. Numerical tests and model assessment

For validation purposes, numerical examples are treated to verify the proposed model 222 described in section 3 either (equations17-24), designated as "M<sub>1</sub>" or equations (14 and 223 17—20), designated as "M<sub>2</sub>". The obtained results are then checked against some available 224 225 solutions measured or analytical to account for the model accuracy. Numerous examples are considered including 1D, 2D and 3D flow features and dealing with the most encountered 226 problems of groundwater in variably saturated porous media, ranging from infiltration-227 evaporation in a soil column, recharge-drainage of a free surface 2D aquifer and free surface 228 3D aquifer response to recharge and pumping. Special attention is paid to flow in unsaturated 229 flow in porous media. 230

- 231 The efficiency of the proposed numerical model is demonstrated against the standard model,
- labeled "Stand", where stiffness and mass matrices entries are expressed normally (equations
- 233 12 and 14).
- 234 Comparisons between the two models were focused on the following criteria:
- 235 The number of time steps, Nts,
- 236 The number of back steps, Nbs
- 237 The total computation time, CPU,
- 238 The total number of nonlinear iterations, Nnlit,
- 239 The mass balance error, MBE, computed as in Yeh (1981).
- All the simulations were performed on a 3.00 GHz Pentium Dual-Core E5700

241 microcomputer.

- 242 4.1. Transient flow in a soil column under infiltration
- 243 4.1.1. Description
- 244 This example was chosen to verify the presented model when dealing with infiltration. The
- experimental results have been reported by Haverkamp et al. (1977). Later, Clement et al.
- 246 (1994) presented a 2-dimensional numerical study of the same case and checked the obtained
- results to those observed by Haverkamp et al. (1977). In the present study, the flow domain is
- 248 a one-dimensional sand column of 0.7 m length, where a constant pressure head of  $\psi = -0.615$
- m is specified at the outlet and a constant flux of 3.29 m/day is imposed at the inlet.
- 250 The soil properties are taken as in Clement et al. (1994). The saturated hydraulic conductivity
- is  $K_s = 8.16$  m/day, the porosity (saturated water content) is  $\theta_s = 0.287$  and the residual water
- 252 content is  $\theta_r = 0.075$ . The values of the soil properties used to parameterize the retention
- curve (equation 8) are  $A = 1.61 \times 10^6$  and B = 3.96 and those for the relative conductivity
- expression (equation 9) are  $C = 1.18 \times 10^6$  and D = 4.74.

4.1.2. Results and discussion

The domain of study is discretized in 71 nodes with a constant spacing of 0.01 m. Transient 256 simulations are performed during a period of 0.8 hours using a varying time step technique, 257  $\Delta t_0 = 0.1 \times 10^{-2}$  hr,  $\Delta t_{min} = 0.1 \times 10^{-5}$  hr,  $\Delta t_{max} = 0.1 \times 10^{-1}$  hr,  $\Delta t_{mag} = 1.2$ ,  $\Delta t_{red} = 0.5$ , along with 258  $maxit_1 = 10$ ,  $maxit_2 = 20$  and maxit = 30. A tolerance of  $10^{-3} m$  is fixed for non linear 259 iterations. The linear system of equations is solved directly using tridiagonal algorithm. 260 Transient solutions are then computed at each 0.1 hours during a total time simulation of 0.8 261 hours, both with the present model and the standard one. Simulation details are summarized in 262 263 Table 1.

Simulated and measured water content are shown in Fig.1. The numerical water content profiles are identical for the standard formulation and the new model. Both are in good agreement with the experimental results. Numerically, the two models perform quiet similarly as it can be seen in Table 1. For every time simulation, the time steps number, the nonlinear iterations number, the CPU times are almost the same, as well as the mass balance error. The new model (version  $M_1$ ) performs similarly, but takes 2 or 3 iterations more to converge to the same solution.

4.2. Two-dimensional, transient variably saturated water-table recharge

272 4.2.1. Description

273 This test case is taken from Vauclin et al. (1979) and involves a transient flow through a

homogeneous and isotropic rectangular soil embankment of 6.00 length and 2.00 m height.

The soil surface is recharged with a constant flux of q = 3.55 m/day, applied over a strip of

1.00 m width, located at the center of the domain. The remaining surface is covered to avoid

evaporation. Because of symmetry, only one part of the whole domain is to be modeled (for

instance, the right one). The subsequent domain is of 3.00 m by 2.00 m and the recharge flux

is then applied over a 0.5 m wide strip. The remaining soil surface is subject to no-flow
condition as well as the bottom and the left side of the resulting domain. Fixed head boundary
is used at the wetting zone of the right side of the domain by maintaining the water level at

282 0.65 m, whereas a no-flow condition is set to the remaining zone (Clement et al, 1994).

283 The soil properties are those of Vauclin et al. (1979). The saturated hydraulic conductivity is

284  $K_s = 8.40$  m/day, the porosity (saturated water content) is  $\theta_s = 0.30$  and the residual water

content is  $\theta_r = 0.30$ . The values of the soil properties used to fit Van Genuchten's (1980)

retention curve model (equations 4 and 5) are  $\alpha = 3.3 \text{ m}^{-1}$  and n = 4.1 (Clement et al, 1994;

Kuznetsov et al, 2012). Specific storage is set to zero (Clement et al, 1994).

288 4.2.2. Results and discussion

Two grids are used, one coarse, with nodal distance of  $\Delta x = 0.1$  m along the horizontal axis, 289 and  $\Delta z = 0.05$  m with respect to the vertical axis. For the fine grid, nodal distances are taken 290 to be the half of those of the coarse grid. For the initial conditions, the water level is fixed to 291 292 0.65 m throughout the whole domain. Transient simulations are performed during a period of 8 hours using a varying time step technique,  $\Delta t_0 = 0.1 \times 10^{-2}$  hr,  $\Delta t_{min} = 0.1 \times 10^{-6}$  hr,  $\Delta t_{max} = 0.1$ 293 hr,  $\Delta t_{mag} = 1.2$ ,  $\Delta t_{red} = 0.5$ , with maxit<sub>1</sub> = 10, maxit<sub>2</sub> = 10 and maxit = 30. A tolerance of 294  $10^{-3}$  m is fixed for non linear iterations and  $10^{-12}$  for inner ICCG iteration. Transient solutions 295 are then computed at 2, 3, 4 and 8 hr elapsed time, using the present model and the standard 296 one. Simulation details, related to the coarse grid, are exhibited in the Table 2. 297

298 Simulated water-table positions for different periods independently of the model are in good

agreement with the measured data of Vauclin et al. (1979), as it can be seen in Fig.2.

300 However, the solutions obtained with the fine grid are more close to the experimental data.

301 The proposed model  $(M_1 \text{ and } M_2)$  and the standard one provide identical solution. The two

versions of the present model, in all cases, converge in less CPU time when compared to the

standard model. Overall, the standard model takes more CPU time to achieve the convergence 303 304 to the fixed level of accuracy. The best performance is recorded by the version  $M_2$  of the proposed model. As an example, during the 2 hours duration of the first period, the 305 convergence of the version M<sub>2</sub> is accomplished in 6401 iterations for 436 s whereas the 306 version M<sub>1</sub> converges in 9552 iterations for 617 s and the standard one converges in 7839 307 iterations for 993 s. However, on a per-non linear iteration basis, the versions  $M_1$  and  $M_2$  of 308 309 the present model and the standard one needed respectively an average CPU time of 0.0645 s, 0.0681 s and 0.1266 to converge. In addition, the version M<sub>1</sub>, the standard model and, in a 310 lesser extent, the version M<sub>2</sub>, experience difficulties during the iterative process which is 311 312 reflected by the number of the needing back steps, 6 for the version M<sub>1</sub>, 5 for the standard model and only 2 for the version M<sub>2</sub>. The mass balance errors are of  $0.215 \times 10^{-1}$  % and 313  $0.208 \times 10^{-1}$  %, for the standard and the version M<sub>2</sub>, respectively, and of 0.398 % for the 314 version M<sub>1</sub>. The same performances of the competing models can be observed at the end of 315 the other periods of simulation. The overall CPU time required for the version M<sub>2</sub> to attain the 316 solution is about 2.3 times lesser than that needed by the standard model. 317 For the fine grid, the same outcome is observed regarding the rate of convergence and the 318

efficiency of the 3 models, with, however, the CPU time needed for model  $M_2$  about 2.5 times

320 lesser than that necessitated by the standard model.

**321 4.3.1.** Description

This test deals with three-dimensional flow in an unsaturated porous media. The flow domain is taken as a very dry cubic block of soil of length a = 15 m and is discretized with hexahedral elements of size  $0.25 \times 0.25 \times 0.25$  m<sup>3</sup>. Because of the symmetry of the problem, only the half of the domain is discretized, resulting in 58621 nodes and 54000 elements. On the top and the bottom of the cubic domain, the pressure head is given by:

327 
$$\psi(x, y, z, t) = \frac{1}{\alpha} \ln \left( e^{\alpha \psi_r} + (1 - e^{\alpha \psi_r}) \sin \frac{\pi x}{a} \sin \frac{\pi y}{a} \right)$$
 and  $\psi = \psi_r$  respectively, where as a non-

flux condition is assigned to the vertical sides. To complete the problem the following initial conditions are taken,  $h(x, y, z, 0) = \psi_r + z$ , where z indicates elevation. The relative conductivity is modeled as  $k_r = e^{\alpha \psi}$ . and he moisture content as  $\theta = \theta_r + (\theta_s - \theta_r) e^{\alpha \psi}$ , where  $\theta_s$  is the saturated moisture content and  $\theta_r$  is the residual content. Parameters of the problem are :  $\alpha = 0.25$ ,  $\theta_s = 0.45$ ,  $\theta_r = 0.15$ ,  $\psi_r = -15$  m and the saturated hydraulic conductivity is taken  $K_s = 0.15$  m/day.

**334** 4.3.2. Results and discussion

Transient simulations are conducted during a period of 15 days using a varying time step technique,  $\Delta t_0 = 0.01$  day,  $\Delta t_{min} = 0.001$  day,  $\Delta t_{max} = 0.025$  day, with maxit<sub>1</sub> = 10, maxit<sub>2</sub> = 10 and maxit = 30. Transient solutions are then computed at 2.5, 5, 10, and 15 days elapsed time, using the present model and the standard one. Simulation details are listed in Table 3.

Fig. 3 shows the pressure head contours at different times obtained with the present numerical

 $M_1$  model (versions  $M_1$  and  $M_2$ ) and the standard one, together with those obtained with the

analytical solution (Tracy, 2006). Very good agreement can be noted between the solutions.

342 As for the precedent cases, the present model performs better than the standard one for all

343 time simulations. In addition, the version  $M_1$  ensures its superiority when compared to the

 $M_2$ . On a per-non linear iteration basis, for all the time simulations, the versions  $M_1$ 

and  $M_2$  of the present model needed an average of 8.23 s and 8.96 s of CPU time,

respectively, to converge, whereas an average time of 25.83 s is required for the standard

347 model, which represent more than 3 times the time needed by the present model, in the case

of the version  $M_1$ , and slightly less than 3 for the other version. Overall, with respect to the

total CPU time, the standard model takes about 3.11 times and 2.88 times the CPU time

needed for the versions M<sub>1</sub> and M<sub>2</sub>, respectively, to converge. The standard model is very
costly in terms of CPU time. As an example, during the 15 days period of simulation, the
CPU time needed for the standard model to attain the solution is 9 hr 51 min 40 s, whereas 3
hr 10 min and 3 hr 25 min were recorded for the versions M<sub>1</sub> and M<sub>2</sub>, respectively. Again the
version M<sub>1</sub> of the presented model shows its highest efficiency.

Regarding the rate of convergence, the three models show practically the same number of iterations for every time simulation, despite a little advantage of the standard model and the version  $M_2$  against the version  $M_1$ , which takes 7 more iterations to accomplish the convergence to the same solution. None of the 3 models had experienced difficulties to converge during all the periods of simulation, thus any back step was noticed. Finally, the Present model and the standard one show very satisfactory mass balance. The absolute values of the mass balance errors for almost all simulations are less than  $10^{-3}$  %.

362 4.4. Three-dimensional transient variably saturated infiltration and pumping

**363** 4.4.1. Description

Two tests are involved and are taken from Kuznetzov et al. (2012), both deals with three-364 dimensional flow in a variably saturated porous media. Recharge from the top is considered 365 first and pumping is studied afterwards, separately. The flow domain is a rectangular 366 parallelepiped of 1000 m in both lateral directions and of 10 m height. The domain is 367 discretized with hexahedral elements of size  $10 \times 10 \times 1$  m<sup>3</sup> resulting in 112211 nodes and 368 100000 elements. For the infiltration case, the boundaries of the domain are subject to no flow 369 condition except the top, which receives a recharging flux of 0.01 m /day from a square area 370 of 100 m×100 m, located at the centre of the surface domain, whereas the remainder part is 371 taken to be impervious. For the pumping case, the boundaries of the domain are subject to no 372 flow condition and a fully penetrating vertical well is situated in the centre of the domain and 373

screened over 2 m from the bottom, with constant pumping rate 50 m<sup>3</sup>/day, during 150 days.

For the both cases, the initial water table height is of 5 m from the bottom of the domain.

376 Parameters of the soil properties used along with the Van Genuchten's (1980) retention curve

model are  $\alpha = 0.145 \text{ m}^{-1}$  and n = 2.68 (Kuznetsov et al, 2012) and the saturated hydraulic

378 conductivity is taken  $K_s = 5$  m/day.

379 4.4.2. Results and discussion

380 Transient simulations are conducted during a period of 150 days using a varying time step

technique,  $\Delta t_0 = 0.1$  day,  $\Delta t_{min} = 0.001$  day,  $\Delta t_{max} = 1$  day, with maxit<sub>1</sub> = 10, maxit<sub>2</sub> = 10 and maxit = 30. A tolerance of  $10^{-3} m$  is fixed for non linear iterations and  $10^{-12}$  for inner PCG iteration. Transient solutions are then computed at 50, 100, and 150 days elapsed time, using the present model and the standard one. Tables 4 and 5 summarize simulation results and

simulated water-table positions obtained are shown in Fig.4.

386 As can be expected, the present model performs better than the standard one for all time simulations. Also, the version M<sub>1</sub> shows again its relative superiority when compared to the 387 version M<sub>2</sub>. On a per-non linear iteration basis, for all the time simulations, the versions M<sub>1</sub> of 388 the present model necessitated an average of 24.20 s to converge and the version M<sub>2</sub> needed 389 an average of 25.44 s of CPU time to achieve the convergence, while the standard model 390 391 required an average time of 93.24 s to accomplish the convergence to the same level of accuracy. For the pumping case the performances are of 24.10 s, 25.39 s and 90.68 s of CPU 392 393 time, respectively. Consequently, the versions M<sub>1</sub> and M<sub>2</sub> of the present model are more than 3 times and half faster than the standard model. Globally, with respect to the total CPU time, 394 the standard model takes about 3.85 times and 3.66 times the CPU time needed, respectively, 395 for the versions M<sub>1</sub> and M<sub>2</sub> to converge when dealing with infiltration and about 3.95 times 396 and 3.74 times the CPU time needed in case of pumping. This demonstrates the high 397

398 computational cost of the standard model and the highest efficiency of the version  $M_1$  of the 399 present model.

400 Concerning the rate of convergence, the three models show exactly the same number of iterations for every time simulation. No resize was necessary for the time step already fixed, 401 depending on the convergence history, during all the time simulations. Finally, the two 402 versions of the present model show very satisfactory mass balance, contrarily to the standard 403 model. Indeed, for all the simulations, the mass balance errors are under  $10^{-2}$  % for the present 404 405 model and are more than 5% for the standard model, in the infiltration case and are under 0.1 % for the present model and more than 10.0% for the standard model, when dealing with 406 407 pumping.

As a comparison with an already published work, Kuznetzov et al. (2012) studied the same
test (medium scale) and reported that their Quasi-3D model performs almost as a standard
full-3D model. The CPU time needed for the standard full-3D is almost about 1.16 times
superior to the CPU time needed for the Quasi-3D model.

412 4.5.1. Description

413 This example deals with a hypothetical large scale groundwater flow problem in a semi-arid context that is described by Prudic et al. (2004). The flow domain is a closed alluvial basin in 414 which recharge is tributary to interactions between the aquifer and the rivers that cross it, 415 416 broadly, from northwest to southeast and from east to west. The aquifer is constituted of unconsolidated deposits of mostly sand and gravel overlying a substratum that is many times 417 less permeable than the filling material of the valley, which thickness increase toward the 418 419 centre. A three dimensional structured mesh of this hypothetical aquifer is generated and adjusted to fit aquifer geometry and wells location. Maps of the aquifer top and bottom are 420 presented in Fig.5 below. The studied domain is discretized into non uniform finite 421

hexahedral elements by a 51×54×11 mesh, respectively in x-, y- and z-direction, resulting in
30294 nodes and 26500 elements.

The boundaries of the domain are subject to no flow condition except the top where 424 425 evaporation occurs, mostly in the vicinity of the streams network, as described in Prudic et al. (2004), when the depth of the free surface is less than a certain value (6.5 m for the present 426 427 simulations) and with variable rates, to a maximum of 0.0021 m/d. The streambed thickness is of 0.915 m and the corresponding leakance is globally constant of 0.864 d<sup>-1</sup>, but can reach 428 locally 1.728 d<sup>-1</sup> in the blue river (Prudic et al, 2004). The elevation of the water in different 429 430 streams (sum of the streambed thickness and the stream depth) is taken constant and equal to 431 1.25 m.

432 For the retention curve model, Brooks and Corey's model (1964) is chosen (equations 6 and 7) with the sets of values of Herrada et al. (2014) and which are used as follows: in the 433 vicinity of the channels ( $\theta_s = 0.396$ ,  $\theta_r = 0$ .,  $\psi_s = -0.131$  m and  $\lambda = 0.127$ ) and ( $\theta_s = 0.477$ ,  $\theta_r =$ 434 435 0.1,  $\psi_s = -0.45$  m and  $\lambda = 1.20$ ) for the sediments in the remainder part of the valley. The hydraulic conductivity is taken as in Prudic et al. (2004): 52.8 m/d in the vicinity of the 436 streams and 10.7 m/d elsewhere. Ten pumping wells are considered; from each one, an 437 amount of 40000  $m^3/d$  is extracted during a period of 50 years. The initial conditions are 438 439 provided as a result of a steady-state simulation of the model without pumping.

440 4.8.2. Results and discussion

Two simulations were conducted to provide the results presented in the framework of this study. First, a steady-state simulation without pumping, but taking into account leakage and evaporation, was conducted to provide initial conditions inputs necessary to the transient simulation. For this purpose, a constant guess water head of 320.25 m was used for all the nodes. A tolerance of 10<sup>-4</sup>, using the relative norm, is fixed for non linear iterations together

with a tolerance of  $10^{-12}$  for inner PCG iteration. The steady-state simulation, performed with 446 447 the present model, required 18 iterations to converge to the specified criteria with an error mass balance of 0.211%. The CPU time needed for the model to converge in steady-state 448 regime is 44.3 s. The standard model, however, fails to converge to the same level of solution 449 mentioned above. The convergence of the standard method is also examined with alleviated 450 criteria and did not converge even with non linear tolerance reduced to  $10^{-2}$ . Indeed, the inner 451 iterations (i, e), the PCG iterations, fail to converge when the inner tolerance is still  $10^{-12}$ . 452 With less drastic conditions (non linear tolerance of  $10^{-2}$  and inner tolerance of  $10^{-5}$ ) the 453 standard method reaches convergence in 17 iterations within 125 s of CPU time and with a 454 455 mass balance error of 34.6%, whereas, the present new model converges in 2 iterations within only 11.8 s of CPU time and with a mass balance error of 0.866%. Fig.6 (left) shows steady-456 state water-table contours. 457

458 Transient simulations are conducted during a period of 50 years using a varying time step

technique,  $\Delta t_0 = 15 \text{ day}$ ,  $\Delta t_{min} = 5 \text{ day}$ ,  $\Delta t_{max} = 365 \text{ day}$ , with maxit<sub>1</sub> = 10, maxit<sub>2</sub> = 10 and maxit = 30. A tolerance of 10<sup>-4</sup> is fixed for non linear iterations and 10<sup>-12</sup> for inner PCG iteration. Transient solutions are then computed at 50 years elapsed time, using the present model and the standard one. Table 6 displays obtained simulation results.

463 Fig.6 (right) exhibits water level contours obtained with the version M<sub>1</sub> of the present model
464 after 50 years of pumping.

For this mimicking real world groundwater flow problem, the present model shows its absolute superiority compared to the standard model. The versions  $M_1$  and  $M_2$  of the present model achieve the convergence in 167 and 301 iterations with CPU times, respectively, of 584 s and 1040 s, whereas, for the standard model, the convergence is achieved in 450 iterations within 5680 s. On a per-iteration basis, standard model, version  $M_1$  and version  $M_2$ 

of the present model need, respectively, 12.62 s, 3.45 sand 3.5 s CPU time. Thus, the present 470 471 model (versions M<sub>1</sub> and M<sub>2</sub>) is about 3.6 times faster than the standard one. In addition, regarding the total CPU time, the version M<sub>1</sub> of the present model is almost 10 times faster 472 than the standard model and almost 2 times faster than the version M<sub>2</sub> of the present model, 473 which demonstrates the high convergence rate of the version M<sub>1</sub> compared to the standard 474 model and the version  $M_2$  of the present model. This is confirmed by the iterations numbers 475 476 needed for convergence which are of 167, 301 and 450, respectively for the versions M<sub>1</sub> and M<sub>2</sub> of the present model and the standard one. As it can be seen, particularly for this real 477 world test, the version M<sub>1</sub> of the present model shows the highest numerical efficiency to 478 479 solve large and complex groundwater flow problems.

Regarding the mass balance analysis, the versions  $M_1$  and  $M_2$  provide satisfactory mass balance, 0.144% and -0.633%, contrarily to the standard model which shows an excessively high mass balance error of 131%. Thus, the standard model (equations 12 and 14) may be inadequate to solve complex groundwater flow problems, in addition to its high numerical cost. Fig.7 shows the solutions obtained with the version  $M_1$  and  $M_2$  of the present model together with the solution obtained with the standard one.

486 5. Conclusions

487 An efficient transient numerical mass-conservative model is constructed to solve variably saturated groundwater flow problems. A Galerkin linear finite element mesh is used to 488 discretize the entire domain, saturated and unsaturated compartments. Picard iteration 489 490 technique, together with the fully implicit Euler method, is used to solve the obtained discretized equations system. The present model is particularly suitable for limited 491 492 computational resources when 3D groundwater flow problems are tackled. The stiffness and mass matrices involved are approximated in a way that only hydraulic properties are updated 493 during the iterative process to guarantee less computational effort, when solving the PDE. The 494

geometric part of the matrices remains unchanged, so computed once and stored, which
ensures important savings in CPU time. The new model is presented with or without
approximate mass matrix (versions M<sub>1</sub> and M<sub>2</sub>). Both are tested and compared to a standard
numerical model obtained with non altered matrices.

The proposed model capabilities are tested and compared to the standard model across benchmark tests for which analytical solutions or experimental results are available. A "real world" groundwater flow problem is also solved to examine the performance of the competing numerical tools. Different soil water curve models, conferring variable non linearity degrees to the Richards equation, are also implemented. Emphasis is made when reproducing these examples on two major aspects, that are accuracy and efficiency assessed, respectively, by the mass balance error and the CPU time to achieve convergence.

506 Along with the studied examples, the developed new model shows its superiority to solve variably saturated flow problems in 1D, 2D and 3D configurations. The developed new model 507 508 exhibits high computational performance regarding time cost and solution accuracy. These 509 outcomes are clearly revealed with high dimensionality problems. In effect, the present model and the standard one perform almost similarly when dealing with 1D flow problems, may be, 510 because of the relatively small number of nodes and elements. The numerical water content 511 curves are identical and the CPU times are very close as well as the mass balance errors. 512 Since the dimensionality increases, the present model  $(M_1 \text{ and } M_2)$  performs increasingly 513 514 better than the standard model. The latter is computationally intensive due to the fact that the involved matrices should be, totally rebuilt, at every nonlinear iteration, contrarily to the 515 proposed new model, where only the hydraulic parts are to be updated; the geometric parts are 516 computed once and stored. As a consequence, the CPU time needed to attain the convergence 517 518 becomes dramatically highest for the standard model. It is clear that the required CPU time increases with high dimensionality and the total number of elements and nodes. 519

For the 2D recharge test, the CPU time needed by the standard model is about 2.3–2.5 times 520 521 higher than the CPU time needed for the present model, whereas the mass balance errors are satisfactory for the competing models. When dealing with the 3D flow examples the CPU 522 523 time needed by the standard model to achieve convergence is about 3.66–3.95 times greater 524 than that needed by the present model to converge to the same level of accuracy. On the other hand, the mass balance error is very satisfactory for the present model, less than 0.01% for the 525 526 recharge test and less than 0.1% for test of pumping, whereas for the standard model, these mass balance errors are more than 5% and more than 10% for the two tests mentioned, 527 respectively. It is to notice that the mass balance error provided by the two competing model 528 529 are very satisfactory (less than 0.001%), when dealing with the totally unsaturated flow 530 example. Indeed, the numerical solutions obtained with the two models (present and standard) match very well. However, the CPU time needed by the standard model to converge under the 531 532 same constraints is globally about 3 times greater than that needed for the present model. 533 For the hypothetical real word test, the present model shows its absolute superiority when compared to the standard model. Also, the version M<sub>1</sub> confirms its superiority when 534 535 compared to the version M<sub>2</sub>. Globally, with respect to CPU time, the standard model takes almost 10 times and more than 5 times the CPU time needed respectively by the versions M<sub>1</sub> 536 537 and M<sub>2</sub> of the present model to converge. This, clearly, shows the highest efficiency of the 538 version M<sub>1</sub> of the present model. The version M<sub>2</sub> take about 2 times the CPU time needed to the version M<sub>1</sub> to converge, indeed. Regarding the mass balance analysis, the versions M<sub>1</sub> and 539 M<sub>2</sub> provide satisfactory mass balance, 0.144% and -0.633%, contrarily to the standard model 540 which shows an excessively high mass balance error of 131%. 541

542 As it can be seen, particularly for this real world test, the version  $M_1$  of the present model

543 shows the highest numerical efficiency to solve large and complex groundwater flow

544 problems. Consequently, the use of the standard model to solve groundwater flow problems

may be difficult and almost prohibitive to perform more complicated and repetitive operations
as the calibration, for instance, or to solve coupled groundwater flow and solute transport
problems, particularly with poor or insufficient computational resources. The present model,
however, particularly the version M<sub>1</sub>, is shown to be suitable for solving large and complex
groundwater flow problems and, probably, to be used usefully for the anticipated related
issues.

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657	Table 1. Simula	ation resu	Its for the t	transien	t flow in a s	soil column ເ	under i	nfiltra	ation		
	Model	$\Delta z$	Time	Nts	Nbs C	PU Nnli	it N	/IBE	(%)		

Stand	0.01	0.8	89	0	10	366	0.204E-2
New, M <sub>2</sub>	0.01	0.8	89	0	9.94	366	0.525E-3
Stand	0.01	0.7	79	0	8.96	326	0.228E-2
New, M <sub>2</sub>	0.01	0.7	79	0	8.84	326	0.894E-3
Stand	0.01	0.6	69	0	7.85	286	0.224E-2
New, M <sub>2</sub>	0.01	0.6	69	0	7.80	286	0.924E-3
Stand	0.01	0.5	59	0	6.76	246	0.212E-2
New, M <sub>2</sub>	0.01	0.5	59	0	6.70	246	0.912E-3
Stand	0.01	0.4	49	0	5.72	206	0.192E-2
New, M <sub>2</sub>	0.01	0.4	49	0	5.66	206	0.940E-3
Stand	0.01	0.3	39	0	4.62	166	0.162E-2
New, M <sub>2</sub>	0.01	0.3	39	0	4.56	166	0.103E-2
Stand	0.01	0.2	29	0	3.51	126	0.927E-3
New, M <sub>2</sub>	0.01	0.2	29	0	3.52	126	0.583E-3
Stand	0.01	0.1	19	0	2.36	82	-0.165E-3
New, M <sub>2</sub>	0.01	0.1	19	0	2.31	82	-0.261E-3

		(m)	(m)	(hr)			(s)		
	New, M <sub>1</sub>	0.1	0.05	2	779	6	617	9552	0.398
	New, M <sub>2</sub>	0.1	0.05	2	541	2	436	6401	0.208E-1
	Stand	0.1	0.05	2	685	5	993	7839	0.215E-1
	New, M <sub>1</sub>	0.1	0.05	3	810	6	636	9873	0.287
	New, M <sub>2</sub>	0.1	0.05	3	555	2	445	6540	0.142E-1
	Stand	0.1	0.05	3	707	5	1020	8061	0.201E-1
	New, M <sub>1</sub>	0.1	0.05	4	834	6	651	10106	0.240
	New, M <sub>2</sub>	0.1	0.05	4	565	2	453	6636	0.111E-1
	Stand	0.1	0.05	4	718	5	1030	8170	0.201E-1
	New, M <sub>1</sub>	0.1	0.05	8	875	6	676	10478	0.186
	New, M <sub>2</sub>	0.1	0.05	8	605	2	473	6937	0.622E-2
	Stand	0.1	0.05	8	758	5	1080	8517	0.350E-1
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74	Table 3. Sim	ulation	results fo	<u>r the tra</u> r	nsient, th	ree-dime	nsional unsa	aturated flo	w

Model	$\Delta x$	$\Delta y$	$\Delta z$	Time	Nts	Nbs	CPU	Nnlit	MBE (%)
	(m)	(m)	(m)	(days)			(s)		
New, M <sub>1</sub>	0.25	0.25	0.25	2.5	103	0	2660	352	-0.232E-3
New, M <sub>2</sub>	0.25	0.25	0.25	2.5	103	0	2850	345	-0.210E-3
Stand	0.25	0.25	0.25	2.5	103	0	8150	345	-0.449E-3
New, M <sub>1</sub>	0.25	0.25	0.25	5	203	0	4420	552	-0.867E-3
New, M <sub>2</sub>	0.25	0.25	0.25	5	203	0	4790	545	-0.106E-2
Stand	0.25	0.25	0.25	5	203	0	13600	545	-0.159E-2
New, M <sub>1</sub>	0.25	0.25	0.25	10	403	0	7930	952	-0.444E-3
New, M <sub>2</sub>	0.25	0.25	0.25	10	403	0	8570	945	-0.660E-3
Stand	0.25	0.25	0.25	10	403	0	24600	945	-0174E-2
New, M <sub>1</sub>	0.25	0.25	0.25	15	603	0	11400	1352	-0.231E-3
New, M <sub>2</sub>	0.25	0.25	0.25	15	603	0	12300	1345	-0.401E-3
Stand	0.25	0.25	0.25	15	603	0	35800	1345	-0.200E-2

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Model	$\Delta x$	$\Delta y$	$\Delta z$	Time	Nts	Nbs	CPU	Nnlit	MBE (%)
	(m)	(m)	(m)	(days)			(s)		
New, M <sub>1</sub>	10.0	10.0	1.0	50	59	0	2860	117	0.880E-2
New, M <sub>2</sub>	100	10.0	1.0	50	59	0	3010	117	0.470E-2
Stand	10.0	10.0	1.0	50	59	0	10600	117	5.40
New, M <sub>1</sub>	10.0	10.0	1.0	100	109	0	5310	217	0.508E-2
New, M <sub>2</sub>	10.0	10.0	1.0	100	109	0	5530	217	0.403E-2
Stand	10.0	10.0	1.0	100	109	0	19500	217	5.46
New, M <sub>1</sub>	10.0	10.0	1.0	150	159	0	7580	317	0.385E-2
New, M <sub>2</sub>	10.0	10.0	1.0	150	159	0	8020	317	0.374E-2
Stand	10.0	10.0	1.0	150	159	0	30600	317	5.53

686 <u>Table 4</u>. Simulation results for the transient, three-dimensional water-table recharge.

Model	$\Delta x$	$\Delta y$	$\Delta z$	Time	Nts	Nbs	CPU	Nnlit	MBE (%)
	(m)	(m)	(m)	(days)			(s)		
New, M <sub>1</sub>	10.0	10.0	1.0	50	59	0	3060	130	-0.255E-1
New, M <sub>2</sub>	10.0	10.0	1.0	50	59	0	3230	131	-0.237E-1
Stand	10.0	10.0	1.0	50	59	0	12500	132	-12.8
New, M <sub>1</sub>	10.0	10.0	1.0	100	109	0	5440	230	-0.291E-1
New, M <sub>2</sub>	10.0	10.0	1.0	100	109	0	5750	231	-0.323E-1
Stand	10.0	10.0	1.0	100	109	0	20700	232	-12.8
New, M <sub>1</sub>	10.0	10.0	1.0	150	159	0	7360	298	-0.265E-1
New, M <sub>2</sub>	10.0	10.0	1.0	150	159	0	7780	298	-0.343E-1
Stand	10.0	10.0	1.0	150	159	0	29100	323	-12.7

700 <u>Table 5. Simulation res</u>ults for the transient, three-dimensional water-table with pumping.

$\Delta x$	$\Delta y$	$\Delta z$	Time	Nts	Nbs	CPU	Nnlit	MBE (%)
(m)	(m)	(m)	(years)			(s)		
-	-	-	50	64	0	584	167	0.144
-	-	-	50	70	0	1040	301	-0.633
-	-	-	50	84	7	5680	450	131
	Δx (m) - -	Δx    Δy     (m)     (m)     -     -     -     -     -     -		$\Delta x$ $\Delta y$ $\Delta z$ Time           (m)         (m)         (m)         (years)           -         -         -         50           -         -         -         50           -         -         -         50	$\Delta x$ $\Delta y$ $\Delta z$ Time         Nts           (m)         (m)         (m)         (years)         -           -         -         -         50         64           -         -         -         50         70           -         -         -         50         84	$\Delta x$ $\Delta y$ $\Delta z$ Time         Nts         Nbs           (m)         (m)         (m)         (years)         -         -         -         -         -         -         0           -         -         -         50         64         0           -         -         -         50         70         0           -         -         -         50         84         7	$\Delta x$ $\Delta y$ $\Delta z$ Time         Nts         Nbs         CPU           (m)         (m)         (m)         (years)         (s)         (s)           -         -         -         50         64         0         584           -         -         -         50         70         0         1040           -         -         -         50         84         7         5680	$\Delta x$ $\Delta y$ $\Delta z$ Time         Nts         Nbs         CPU         Nnlit           (m)         (m)         (m)         (years)         (s)         (s)

## 715 <u>Table 6</u>. Simulation results for the transient, hypothetical river-connected aquifer.

















