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Quantization of Acoustic Modes in Dumbbell Nanoparticles

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The vibrational eigenmodes of dumbbell-shaped polystyrene nanoparticles are recorded by Brillouin light spectroscopy (BLS), and the full experimental spectra are calculated theoretically. Different from spheres with a degeneracy of $(2l + 1)$, with l being the angular momentum quantum number, the eigenmodes of dumbbells are either singly or doubly degenerate owing to their axial symmetry. The BLS spectrum reveals a new, low-frequency peak, which is attributed to the out-of-phase vibration of the two lobes of the dumbbell. The quantization of acoustic modes in these molecule-shaped dumbbell particles evolves from the primary colloidal spheres as the separation between the two lobes increases.

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Spatially confined objects have quantized acoustic modes with signature frequencies that depend on the shape, size, elastic properties, and boundary conditions of the object [1–7]. Understanding the modes of particles with complex physical, chemical, and geometric parameters not only facilitates the development of acoustic quantization theories, but also provides a unique access to the elastic properties of small objects inaccessible by conventional methods [8,9] and allows exploring resonant cavities in phononic materials [10]. For objects with dimensions in the submicrometer length scales, the frequencies of the acoustic resonance modes fall in the gigahertz (GHz) range [3,9,11–16] and can be effectively studied by Brillouin light spectroscopy (BLS).

So far, BLS studies have focused mainly on spheres [12,14,17–19] due to the availability of theory, the challenge of synthesizing uniform, nonspherical particles, and the degeneracy of the resonance modes at a given angular momentum quantum number l . In recent extension to anisotropic shapes, the BLS spectra of ellipsoidal colloids are observed to be more complex due to the lifting of the

degeneracy, $m = 2l + 1$ [6,16]. Yet, it is still barely understood how the reduced symmetry affects the eigenmodes of nanoparticles (NPs). In addition, interactions between spherical NPs (SPs) have been demonstrated to induce two prominent features: (1) the appearance of a contact-induced low-frequency mode ($n = 1, l = 1$), and (2) the blueshift and peak splitting of the fundamental quadrupolar ($n = 1, l = 2$) mode of SPs [9,18,20,21]. For plasmonic nanoparticles embedded in a polymer matrix, the optomechanically enhanced BLS spectra reveal the quadrupolar (1, 2) mode and an additional rattling mode, due to the coupling of the acoustic vibrations of the two spheres forming a dimer in close proximity [21–23]. However, the proximity in these studies is ill-defined, and the theoretical account of the vibration spectrum is only semiquantitative.

Dumbbell-shaped NPs (DBs) and their assemblies have been of great interest due to the anisotropic lattice structures, which can potentially provide direction-dependent optical and phononic properties [6,24]. The recent development of methods for synthesizing highly monodisperse DBs [25–27] has created opportunities to explore the shape dependence of vibrational modes. While the diameter is an exclusive parameter for a SP, a DB has three independent parameters: the diameters of the two lobes, D_1 and D_2 , and the center-to-center distance between the lobes d_{12} . DBs offer a versatile platform for investigating the effect of particle shape symmetry on the eigenfrequencies and BLS vibration spectra, bridging colloidal atoms (spheres) and molecules (dumbbells). As d_{12} increases, the shape of a DB

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changes from SP to DB and finally to two independent spheres with mechanical contact. The well-controlled proximity d_{12} is an additional advantage compared to the metallic nanocomposites, and the use of submicron particles enables a strong BLS signal [15].

In this work, we measure the BLS spectrum of polystyrene DBs using BLS and compare their vibrational eigenmodes with those of SPs. The incident light (a Nd:YAG laser with a wavelength of 532 nm) is directed to the SP and DB sample films to induce photon-phonon interactions, and the scattered light is collected by an aperture and focused by lenses to enter a tandem Fabry-Perot interferometer (JRS Scientific Instruments), with the light transmitted through the interferometer recorded by an avalanche photo diode. As the phonon wave vector \mathbf{q} is ill-defined for opaque samples due to the multiple light scattering, the BLS spectrum reveals the resonance modes of the particles that are uniquely defined by the particles' geometric and elastic characteristics. We also develop a numerical method for calculating BLS spectra by combining the finite element method (FEM) and the theory of BLS [see Supplemental Material (SM), Sec. S2 for details [28]]. The calculation results reproduce accurately the measured BLS spectra of SPs and DBs, and the vibrational patterns of the DB modes are interpreted in terms of the vibrational modes of the two lobes. A systematic tuning of d_{12} provides an understanding of how near-frequency modes of the two composing spheres combine into a DB mode. A new low-frequency, d_{12} -sensitive peak is revealed and attributed to the out-of-phase rigid translations and rotations of the two lobes.

The vibrational modes of a free SP were first investigated by Lamb in 1881 [1] and later revisited and extended in 1982 [34]. They consist of two types: (1) spheroidal, $S(n, l, m)$, with $n \geq 1$, $l \geq 0$, and $-l \leq m \leq l$, and (2) torsional, $T(n, l, m)$, with $n \geq 1$, $l \geq 1$, and $-l \leq m \leq l$, where n , l , and m are the quantization numbers analogous to the atomic orbitals. For SPs, the quantization number m is often dropped, because the modes with the same n and l but different m are degenerate, although their displacement patterns differ. The $S(1, 1)$ and $T(1, 1)$ modes of individual SPs, both threefold degenerate, represent the particle translation and rotation at zero frequency, respectively. For mechanically interacting SPs, however, both $S(1, 1)$ and $T(1, 1)$ have nonzero frequencies, and provide information on the interaction strength and contact area [9,35]. The eigenfrequencies and displacement fields of the S and T modes of SPs can be calculated [34,36] as presented in SM, Sec. S1 [28]. The sequence of the eigenfrequencies is determined by c_L/c_T , while their absolute values depend on the longitudinal c_L (or transverse c_T) sound velocities in the NP and the NP diameter D . In SM Table S1 [28], we show the lowest ten eigenfrequencies of a SP with $D = 190$ nm, $c_L = 2300$ m/s, and $c_T = 1210$ m/s. For DBs, the eigenfrequencies and displacement fields are solved numerically. Besides the eigenfrequencies and eigenvectors of the

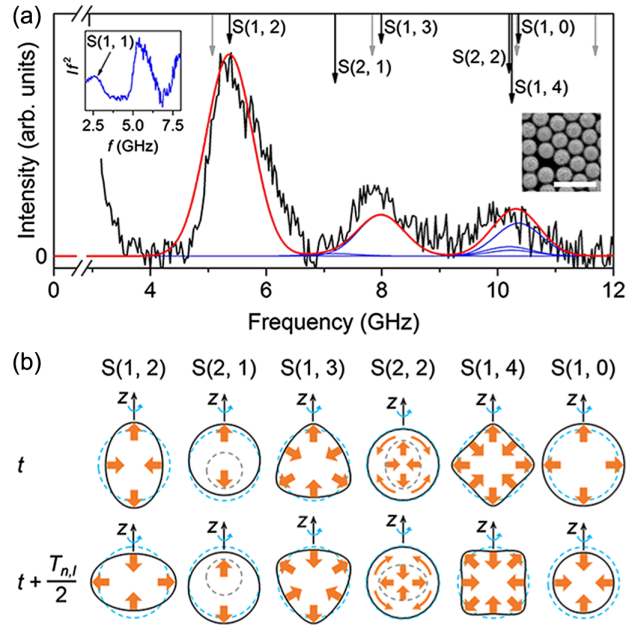


FIG. 1. Vibrational eigenmodes of polystyrene spherical nanoparticles. (a) Experimental (black line) and calculated (red line) BLS spectra of SP190. The blue lines represent the calculated intensities of the individual eigenmodes broadened to Gaussian profiles by the instrumental width (0.35 GHz). Six spheroidal eigenmodes are indicated by black arrows. Four torsional eigenmodes are indicated by gray arrows. Left inset: a reduced spectrum, $I f^2$. Right inset: a scanning electron microscopy (SEM) image of the nanoparticles (scale bar: 500 nm). The relatively low signal-to-noise ratio is due to the short accumulation time. (b) Illustration of the mode patterns (cross-sectional view, for $m = 0$) of the six spheroidal eigenmodes in (a). The cyan, dashed circles represent nanoparticles at rest. The black, solid lines represent the nanoparticles surface in the specific vibrational eigenmodes. The gray, dashed lines indicate displacements in internal volumes. The orange arrows indicate local deformations. t and $(t + T_{n,l}/2)$ represent two chosen time instants separated by half of the vibrational period $T_{n,l}$ of the (n, l) mode.

vibrational modes, we calculate the BLS scattering intensities [37] of these modes to fully account for the BLS experiment (SM, Sec. S2 [28]).

We first investigate the vibrational eigenmodes of polystyrene SPs with $D = 190 \pm 4$ nm (SP190) [right inset to Fig. 1(a); SM, Sec. S3 and Fig. S1(a)]. The experimental BLS spectrum of SP190, obtained from a measurement in the transmission geometry (SM, Sec. S4), features three peaks in the range of 4–12 GHz [Fig. 1(a)]. An additional, broad low-frequency Brillouin peak, centered at about 2.5 GHz, appears in the plot of the reduced intensity $I f^2$ [left inset to Fig. 1(a)]. This peak, which is absent in the spectrum of a free sphere, is designated as $S(1, 1)$, as it stems from the zero-frequency rigid translations of the sphere and can be attributed to the propagation of longitudinal sound waves due to physical contact between adjacent spheres [9]. We further calculate the

BLS spectrum [red line in Fig. 1(a)] to identify the contributions of the individual eigenmodes (SM, Sec. S2). The first ten eigenfrequencies (up to 12 GHz), including six spheroidal and four torsional modes, agree well with those in Table S1 computed using Lamb's theory [1]. The vibrational patterns of the spheroidal modes with $m = 0$ at the six frequencies are shown in Fig. 1(b), whereas those of the four torsional modes are shown in Fig. S2. A complete list of the eigenfrequencies (up to 12 GHz) and their displacement fields are shown in Fig. S3.

The BLS intensities of the torsional modes are zero, as these modes cause no change in the NP volume or shape (i.e., no change in the polarizability of the particle). The nonzero BLS intensity of the spheroidal modes confirms that both even- and odd- l eigenmodes contribute to BLS [15]. While the first peak is contributed solely by the $S(1, 2)$ quadrupolar mode, the second peak is contributed by two modes, $S(2, 1)$ and $S(1, 3)$. The third peak includes contributions from three modes, $S(2, 2)$, $S(1, 4)$, and $S(1, 0)$, although the intensities of the first two modes are significantly lower than the third one. Figure 1(a) implies that a quantitative analysis of a BLS spectrum requires considering both the eigenfrequencies and intensities of the modes, because the peaks in the BLS spectrum may each include contributions from multiple eigenmodes [9,12]. Furthermore, the relative intensities of the higher-frequency peaks with respect to the first one increase, as D/λ increases (Fig. S4) [15]. We note that the spherical particles possess high sphericity [Figs. S1(a) and S1(c)]. A slight asphericity and small irregularities on the particle surface have a minor effect on the vibrational eigenfrequencies and eigenmodes, as revealed by the FEM calculations (Figs. S5 and S6).

The eigenmodes of DBs are more complex due to their lower symmetry. The SEM image [right inset to Fig. 2(a); SM, Sec. S3 and Fig. S1(b)] is used to estimate the three geometric parameters, including $D_1 = 232 \pm 9$ nm, $D_2 = 214 \pm 10$ nm, and $d_{12} = 151 \pm 9$ nm. For these DBs, denoted as DB232-151-214, the recorded BLS spectrum [black line in Fig. 2(a)] reveals five peaks labeled as *A*, *B*, *C*, *D*, and *E*. We calculate the BLS spectrum of a single DB232-151-214 (SM, Sec. S2) with the same sound velocities of SP190. The reduced symmetry and the consequent m splitting renders the number of observed peaks in the BLS spectrum [Fig. 2(a)] much larger than for the corresponding SP190 [Fig. 1(a)]. The characteristic eigenmodes under peaks *A*–*E* are labeled, with the focus on peaks *A* and *B*. Peak *A* includes the contributions from three modes, denoted as $(A1)_2$, $(A2)$, and $(A3)_2$, which are doubly (indicated by the subscript “2”), singly, and doubly degenerate, respectively. Peak *B* is contributed by eight modes, denoted as $(B1)_2$ to $(B8)_2$, while peaks *C*–*E* include contributions from even more modes. The vibrational patterns of the five singly degenerate modes indicated as $(A2)$, $(B6)$, $(C1)$, $(D1)$, and $(E1)$ in Fig. 2(a) are

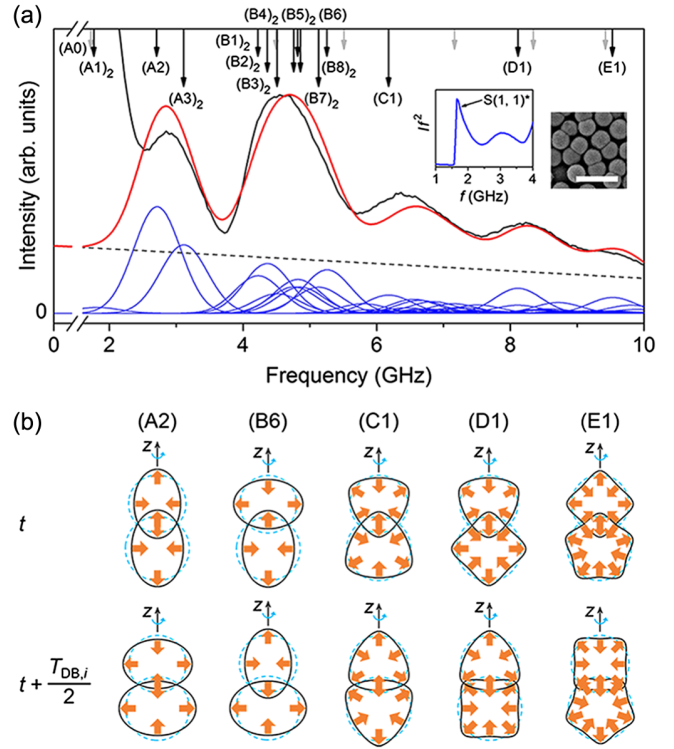


FIG. 2. Vibrational eigenmodes of polystyrene dumbbell-shaped nanoparticles. (a) Experimental (black line) and theoretical (red line) BLS spectra of DB232-151-214. The blue lines represent the theoretical intensities of the individual eigenmodes broadened by the instrumental width (0.35 GHz) and the dashed line denotes a linear background added to the calculated total spectrum to account for the Rayleigh wing in the experimental spectrum. Characteristic eigenmodes in the five peaks (*A* to *E*) are indicated by black arrows. The subscript “2” means the mode is doubly degenerate. Six pure torsional eigenmodes are indicated by gray arrows. Left inset: a reduced spectrum, If^2 (the abrupt intensity drop at around 1.6 GHz is due to the shutter). Right inset: an SEM image of the nanoparticles (scale bar: 500 nm). (b) Illustration of the mode patterns (cross-sectional view) of the six singly degenerate spheroidal eigenmodes in (a). See the Fig. 1(b) caption for more details. $T_{DB,i}$ represents the vibrational period of the corresponding eigenmode.

shown in Fig. 2(b), whereas those of the six singly degenerate modes indicated by gray arrows are shown in Fig. S7.

It is reasonable to envisage the vibrational patterns of the DBs as those of the two lobes, seen as spheres. For example, $(A2)$ and $(B6)$ have patterns that mimic the patterns of two $S(1, 2)$ modes of the individual lobes vibrating in phase and out of phase [Fig. 2(b)], respectively. Modes of the individual lobes with different l 's interact too, as seen in the vibrational patterns of $(D1)$ and $(E1)$. $(D1)$ can be considered as a hybridization of the $S(1, 3, 0)$ mode of the top (smaller) lobe with the $S(1, 4, 0)$ of the bottom (larger) lobe [Fig. 2(b)]. $(E1)$ can be considered as a hybridization of the $S(1, 4, 0)$ mode of the top (smaller) lobe with the $S(1, 5, 0)$ of bottom (larger) lobe [Fig. 2(b)].

This hybridization is favored by a small mismatch between the frequency of a mode with a smaller l in the smaller sphere and that of a mode with a larger l in the larger sphere. The similar patterns of the two sphere modes in the overlap region also facilitate the hybridization. We note that this schematic representation in terms of the sphere modes is not general. It works well only for DBs with high d_{12} values, where the two spheres are well evident and interact via a disk with a small radius r_0 (i.e., $r_0 \ll \min\{D_1, D_2\}$). Based on the displacement fields of the DB232-151-214 eigenmodes (Fig. S8), only a few modes resemble those of the two individual lobes. The resemblance occurs, when only one sphere vibrates, while the second sphere has no mode with a comparable frequency.

To understand the vibrational eigenmodes of DB232-151-214, we calculate the eigenfrequencies of a series of DBs with varied d_{12} [Fig. 3(a)]. At $d_{12} \leq 9$ nm, the two lobes of the DB overlap, and the eigenmodes of the “spherical” DB are similar to those of SP232. The corresponding eigenfrequencies of the torsional and spheroidal modes are indicated on the left y axis in Fig. 3(a) by empty and filled triangles, respectively. As d_{12} increases, the eigenfrequencies separate into $(2l + 1)$ branches due to the lifting of the degeneracy. Specifically, a $T(n, l)$ or $S(n, l)$ mode separates into one singly and l doubly degenerate modes, which correspond to $m = 0$ and $m = \pm 1, \pm 2, \dots, \pm l$, respectively. In Fig. 3(a), the singly degenerate modes originating from $T(n, l, m = 0)$ [$S(n, l, m = 0)$] are represented by dashed (solid), blue lines, whereas the remaining doubly degenerate modes are shown by solid, red lines. Among the DB eigenmodes, only those originating from $T(n, l, m = 0)$ have zero Brillouin intensity, while the others are all Brillouin active. As d_{12} approaches $(D_1 + D_2)/2 = 223$ nm, the eigenmodes become a combination of those of the two lobes (i.e., SP232 and SP214). Concomitantly, the frequencies of the lowest six eigenmodes, shown by the lowest four lines in Fig. 3(a), vanish. Among the six vanishing modes, the lower three and higher three modes can be assigned to the rotational and translational modes of the two spheres, respectively. We note that a single DB has three rotational and three translational modes, whereas two free SPs have six rotational and six translational modes. The six disappearing modes are the zero-frequency translations and rotations of the whole DB. Moreover, the DB eigenfrequencies collapse into those of SP232 and SP214, as indicated by the black and green triangles on the right y axis in Fig. 3(a), respectively.

The modes contributing to peaks A and B in Fig. 2(a) are highlighted in Fig. 3(a). At large d_{12} , the mode (A2) is well represented by the nearly rigid translation of the two spheres along the z axis in opposite directions [i.e., by the $S(1, 1, 0)$ modes of the individual lobes vibrating out-of-phase]. As d_{12} decreases, (A2) gains nature of the $S(1, 2)$ modes of the two lobes. For even smaller d_{12} , (A2) converts

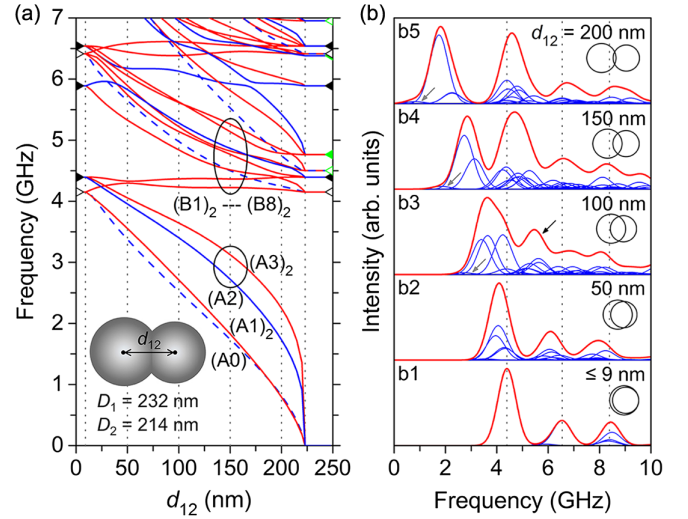


FIG. 3. Vibrational eigenmodes and BLS spectra of asymmetric dumbbell-shaped nanoparticles. (a) Eigenmodes of DB232- d_{12} -214 at varying d_{12} 's. The red lines denote doubly degenerate modes. The blue, solid (dashed) lines represent singly degenerate modes with non-zero (zero) Brillouin intensities. The eigenmodes contributing to peaks A and B in Fig. 2(a) are circled. The black, filled (open) triangles denote spheroidal (torsional) eigenmodes of SP232. From bottom to top: $T(1, 2)$, $S(1, 2)$, $S(2, 1)$, $T(1, 3)$, and $S(1, 3)$. The green, filled (open) triangles denote spheroidal (torsional) eigenmodes of SP214. From bottom to top: $T(1, 2)$, $S(1, 2)$, $S(2, 1)$, and $T(1, 3)$. Inset: a schematic of the DB. (b) Calculated BLS spectra of DB232- d_{12} -214 at five d_{12} 's. The gray arrows in (b3)–(b5) indicate the lowest-frequency mode. The black arrow in (b3) indicates the new, intermediate-frequency peak at $d_{12} = 100$ nm. See the caption of Fig. 2(a) for more details. Insets: schematics of DBs.

to $S(1, 2)$ of SP232. Similar arguments hold for the doubly degenerate $(A3)_2$ modes, but in this case the $S(1, 1, 1)$ and $S(1, 1, -1)$ modes (i.e., translations of the two spheres along the x and y directions) are present for high d_{12} values. As d_{12} decreases, $(A3)_2$ gain nature of $S(1, 2)$ of the two spheres, while for even smaller d_{12} , they convert to $S(1, 2)$ of SP232. On the other hand, for high d_{12} values, the $(A0)$ and $(A1)_2$ modes correspond to pure, out-of-phase rotations of the two spheres with respect to the z axis and x (or y) axis, respectively. They become three $T(1, 1, m)$ modes of SP232 at very small d_{12} . For intermediate d_{12} , the $(A1)_2$ modes are not purely torsional. They are hybridized with the $S(1, 1)^*$ modes [left inset to Fig. 2(a)] and thereby gain some Brillouin activity that produces the weak, lowest-frequency mode [blue line indicated by a gray arrow in Figs. 3(b3)–3(b5)]. These examples demonstrate that a perturbative approach based on the knowledge of the modes of the two spheres helps to understand the nature of the DB vibrations. The eight modes contributing to peak B originate from the $T(1, 2)$, $S(1, 2)$, $S(2, 1)$, $T(1, 3)$, and $S(1, 3)$ modes of SP232. Alternatively, they could be considered to result from the $T(1, 2)$ and $S(1, 2)$ modes

of SP232 and SP214. We note that the (*B6*) mode is closely related to the $S(2, 1)$ mode of SP232 at small d_{12} and to the $S(1, 2)$ of SP232 at large d_{12} . Therefore, as d_{12} increases from $(D_1 - D_2)/2$ to $(D_1 + D_2)/2$, the nature of the DB eigenmodes changes.

To investigate how d_{12} affects the vibration spectrum of DBs, we calculate the BLS spectra at five d_{12} values [Fig. 3(b)]. For $d_{12} \leq 9$ nm, the dumbbell shape becomes spherical, and the BLS spectrum is similar to that of SP190 [Fig. 1(a)]; due to the different diameters, the eigenfrequencies are shifted ($f \sim D^{-1}$) [13]. At $d_{12} = 50$ nm, the NP has a prolate ellipsoidal shape, and the anticipated degeneracy lifting is observed. The total spectrum resembles that of SP232, except for a slight redshift. At $d_{12} = 100$ nm, the spectrum becomes more complicated with a broadened first peak and the appearance of a new peak [indicated by the black arrow in Fig. 3(b3)] with a frequency between those of the first two peaks of SP232. As d_{12} further increases, a strong, low-frequency peak appears, and it gradually becomes separated from the higher-frequency peaks. With increasing d_{12} , its frequency decreases, and its relative intensity increases. The three higher-frequency peaks in the DB spectrum resemble those of SP232 with slight modifications caused by the different diameters of the two DB lobes. We also conduct calculations on symmetric DBs with $D_1 = D_2 = 200$ nm and d_{12} varying from 0 to 200 nm. The peak splitting of DB200- d_{12} -200 (Fig. S9) is similar to that of DB232- d_{12} -214 (Fig. 3). In the limit of $d_{12} = (D_1 + D_2)/2$, the new, low-frequency peak resembles the low-frequency, rattling mode in the BLS spectrum of symmetric metallic dimers with mechanical contact [21,23].

In conclusion, we report the measured and calculated BLS spectra of DBs with well-defined proximity d_{12} . As d_{12} increases, the spectrum evolves from that of a SP to that of a DB, revealing a new, low-frequency peak that originates from the out-of-phase vibration of the two lobes. The BLS spectrum of a DB includes contributions from all spheroidal and torsional eigenmodes, except for those singly degenerate torsional modes that cause no change in the particle volume or shape. The calculation method, which combines the FEM and theory of BLS, and predicts BLS spectra in close agreement with experiments, could be extended to other particle shapes and thus advance the particle vibration spectroscopy. The elucidation of the vibration eigenmodes of DBs and the well-controlled tunability of their geometric parameters (i.e., D_1 , D_2 , d_{12}) are envisaged to promote the development of a wide range of DB-based phononic crystals with direction-dependent, tunable band gaps, which could have potential applications in areas such as nanomaterials and on-chip devices by providing more degrees of freedom [38] in wave-based filtering and computing.

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