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# DISCRETE APPROXIMATION OF THE GRIFFITH FUNCTIONAL BY ADAPTATIVE FINITE ELEMENTS 

JEAN-FRANÇOIS BABADJIAN AND ÉLISE BONHOMME


#### Abstract

This paper is devoted to show a discrete adaptative finite element approximation result for the isotropic two-dimensional Griffith energy arising in fracture mechanics. The problem is addressed in the geometric measure theoretic framework of generalized special functions of bounded deformation which corresponds to the natural energy space for this functional. It is proved to be approximated in the sense of $\Gamma$-convergence by a sequence of discrete integral functionals defined on continuous piecewise affine functions. The main feature of this result is that the mesh is part of the unknown of the problem, and it gives enough flexibility to recover isotropic surface energies.


## 1. Introduction

1.1. The variational approach to fracture. The Griffith functional has been introduced in the context of brittle fracture. It finds its roots in the seminal work of Griffith [28] whose main ideas have been revisited in [27] (see also the monograph [12]) into a variational evolution formulation. The main point is that, in a quasi-static setting and in presence of irreversibility, a constrained global minimization principle together with an energy balance select equilibrium states of an elastic body experiencing brittle fracture. In a nutshell, the Griffith energy is defined by

$$
\begin{equation*}
\mathcal{G}(u, K):=\int_{\Omega \backslash K}|e(u)|^{2} d x+\mu \mathcal{H}^{N-1}(K) \tag{1.1}
\end{equation*}
$$

where $\Omega \subset \mathbb{R}^{N}$, a bounded open set, stands for the reference configuration of an elastic material, $K \subset \Omega$ is a codimension-one set representing the crack, $u: \Omega \backslash K \rightarrow \mathbb{R}^{N}$ is the displacement field which might be discontinuous across $K$, and its symmetric gradient $e(u):=\left(\nabla u+\nabla u^{T}\right) / 2$ is the linearized elastic strain. The constant $\mu>0$ is a material parameter called toughness. This energy puts in competition a bulk energy, representing the elastic energy stored in the body outside the crack, and a surface energy penalizing the presence of the crack $K$ through its ( $N-1$ )-dimensional Hausdorff measure, henceforth denoted by $\mathcal{H}^{N-1}$.

This problem falls within the framework of so-called free discontinuity problems (according to De Giorgi's terminology), and it presents many formal analogies with its scalar counterpart, the Mumford-Shah functional. Although, thanks to geometric measure theory, the existence theory for the latter is by now quite well understood (see e.g. [3] and references therein), the minimization of the Griffith functional had to face serious additional difficulties. In particular, a satisfactory existence theory has only recently been solved. As for the Mumford-Shah functional, it passes through the introduction of a "weak formulation" where the crack is replaced by the jump set $J_{u}$ of $u$. A convenient functional setting to investigate this problem is that of functions of bounded deformation, $B D(\Omega)$, which correspond to (integrable) vector fields $u: \Omega \rightarrow \mathbb{R}^{N}$ whose distributional symmetric gradient $E u$ is a bounded Radon measure. This space has been introduced in [40] (see also $[41,39]$ ) as a natural space to formulate problems of perfect plasticity. Brittle fracture however requires a finer understanding of this space and especially the introduction of the subspace $S B D(\Omega)$ of special functions of bounded deformation in $[2,10]$, for which the singular
part of $E u$ with respect to the Lebesgue measure is concentrated on the jump set. Unfortunately, this step forward was still not enough because of lack of control of the values of $u$ (due to the failure of Poincaré-Korn and/or Korn type inequalities in that space). It is only recently that the introduction of the space $G S B D(\Omega)$ of generalized special functions of bounded deformation in [26] (see Section 2 for the precise definition) has given a satisfactory mathematical framework to investigate a well founded existence theory for the weak formulation, as well as for the original one. Some further compactness properties of that space have been investigated in $[18,17]$ which has led to prove the existence of minimizers of the Griffith functional under Dirichlet boundary conditions (formulated in a relaxed sense).
1.2. Approximation of the Griffith energy. The $\Gamma$-convergence approximation of free discontinuity problems (e.g. by more tractable ones from a numerical point of view) is of fundamental importance in applications. It has been proven in [13] that it is not possible to approximate free discontinuity functionals by means of local integral functionals. To overcome this difficulty, a first possibility is to introduce an additional variable like, e.g., in phase field approximations where the sharp discontinuity is smoothened into a diffuse discontinuity. It represents one of the most popular methods which have already proven to be successful in other contexts such as the Modica-Mortola approximation of the perimeter functional [32], or the Ambrosio-Tortorelli approximation of the Mumford-Shah functional [4]. In the context of brittle fracture, such approximations, which have a founded mechanical interpretation as a gradient damage model, have only recently been established in full generality in [19] (see also [16, 30]). The main drawback is that, an additional numerical approximation would give rise to a multiscale problem with on the one hand the parameter of approximation, and on the other hand the mesh size (see e.g. [9, 7, 22]). Another possibility is to use nonlocal integral functionals as e.g. in [13, 35].

For what concerns the numerical treatment of free discontinuity problems, the main difficulty is related to the fact that the jump set is part of the unknowns and that standard discontinuous finite element methods do not in general apply in this context. Having this problematic in mind as well as the multiscale issues arising in phase field or nonlocal approximations, one is thus tempted to find single scale discrete approximations of free discontinuity problems. There is a huge literature on this subject and, without being exhaustive, we refer to discrete-to-continuous approximations results $[1,14,15,29,33,34,36]$, nonlocal finite elements approximations [31, 38] or discrete approximations based on stochastic meshes in $[8,37]$.

Let us focus on the discrete approximation result obtained in [21] for the Mumford-Shah functional in dimension $N=2$. In that work, the classical Mumford-Shah functional

$$
F(u):=\int_{\Omega}|\nabla u|^{2} d x+\mu \mathcal{H}^{1}\left(J_{u}\right)
$$

is approximated in the sense of $\Gamma$-convergence by a functional of the form

$$
F_{\varepsilon}(u):=\int_{\Omega} f_{\varepsilon}(\nabla u) d x
$$

putting a restriction on the functional space on which $F_{\varepsilon}$ is defined. The functional $F_{\varepsilon}$ is discrete in the sense that $u$ is (a scalar-valued) continuous function and piecewise affine on suitable $\varepsilon$ dependent meshes (see Definition 1.1). It consists in an adaptative finite element approximation because there is an implicit mesh optimization whose numerical implementation has been carried out in [11]. The function $f_{\varepsilon}(\nabla u)$ takes the form $\frac{1}{\varepsilon} f\left(\varepsilon|\nabla u|^{2}\right)$ where $f$ is a nondecreasing function satisfying the standard properties (1.4). Typical examples of functions $f$ are, on the one hand the arctan function (as e.g. in [29] following a conjecture of De Giorgi) and, on the other hand $f(t)=t \wedge \kappa$. The main feature of this result is that, allowing the mesh to move gives enough flexibility to approximate isotropic surface energies. The constant $\mu$ appearing in the functional $F$ is explicit and only depends on $\kappa$ and the geometry of the triangulation. An analogous analysis
has been carried out in [33], where the author constraints the mesh to be made either of equilateral triangles, or of right isosceles ones. In that case, the result is that the functional $F_{\varepsilon} \Gamma$-converges to an anisotropic version of the Mumford-Shah functional

$$
\int_{\Omega}|\nabla u|^{2} d x+\int_{J_{u}} \phi\left(\nu_{u}\right) d \mathcal{H}^{1}
$$

for some function $\phi: \mathbb{S}^{1} \rightarrow \mathbb{R}$, which can be explicitly computed, depending on the normal $\nu_{u}$ to the jump set $J_{u}$.

In [36], the same problem is addressed in the two-dimensional vectorial setting. If $f_{\varepsilon}$ is as before, the following approximating energy is considered

$$
\int_{\Omega} f_{\varepsilon}(e(u)) d x
$$

As in [33], the $\varepsilon$-dependent mesh is fixed and made either of equilateral triangles, or of right isosceles triangles, and the result is that this functional $\Gamma$-converges to an anisotropic version of the Griffith functional

$$
\int_{\Omega}|e(u)|^{2} d x+\int_{J_{u}} \phi\left(\nu_{u}\right) d \mathcal{H}^{1}
$$

where $\phi: \mathbb{S}^{1} \rightarrow \mathbb{R}$ is as in [33]. Note that if $f(t)=t \wedge \kappa$, then

$$
f_{\varepsilon}(e(u))= \begin{cases}\varepsilon|e(u)|^{2} & \text { if } \varepsilon|e(u)|^{2} \leq \kappa  \tag{1.2}\\ \kappa & \text { if } \varepsilon|e(u)|^{2}>\kappa\end{cases}
$$

In order to recover the isotropic Griffith energy (1.1), a similar approximation result is considered in [34] where, now, the meshes are allowed to move as in [21], but the function $f_{\varepsilon}$ now depends on the full gradient $\nabla u$ (instead of the symmetric gradient) and behaves like

$$
f_{\varepsilon}(\nabla u) \sim \begin{cases}\varepsilon|e(u)|^{2} & \text { if } \varepsilon|\nabla u|^{2} \leq \kappa  \tag{1.3}\\ \kappa & \text { if } \varepsilon|\nabla u|^{2}>\kappa\end{cases}
$$

(compare with (1.2)). In that case, the analysis of [21] can be adapted to show a $\Gamma$-convergence result towards the isotropic Griffith energy (1.1) with the same geometric multiplicative constant $\mu$ (as in [21]) in front of the surface energy.
1.3. Our result. The objective of the present work is to generalize the previous results in the two-dimensional vectorial case to show an analogous statement as in [21], namely an adaptative discrete finite element approximation of the isotropic Griffith functional. To state precisely our main result, Theorem 1.2, we need to introduce some notation (we refer to Section 2 regarding functional spaces).

Let $\Omega$ be a bounded open set of $\mathbb{R}^{2}$ with Lipschitz boundary. As in [21], we introduce the following class of admissible meshes.

Definition 1.1. A triangulation of $\Omega$ is a finite family of closed triangles intersecting $\Omega$, whose union contains $\Omega$, and such that, given any two triangles of this family, their intersection, if not empty, is exactly a vertex or an edge common to both triangles. Given some angle $\theta_{0}$ with $0<\theta_{0} \leq$ $45^{\circ}-\arctan (1 / 2)$, and a function $\varepsilon \mapsto \omega(\varepsilon)$ with $\omega(\varepsilon) \geq 6 \varepsilon$ for any $\varepsilon>0$ and $\lim _{\varepsilon \rightarrow 0^{+}} \omega(\varepsilon)=0$, we define, for any $\varepsilon>0$

$$
\mathcal{T}_{\varepsilon}(\Omega):=\mathcal{T}_{\varepsilon}\left(\Omega, \omega, \theta_{0}\right)
$$

as the set of all triangulations of $\Omega$ made of triangles whose edges have length between $\varepsilon$ and $\omega(\varepsilon)$, and whose angles are all greater than or equal to $\theta_{0}$. Then we consider the finite element space $V_{\varepsilon}(\Omega)$ of all continuous functions $u: \Omega \rightarrow \mathbb{R}^{2}$ for which there exists $\mathbf{T} \in \mathcal{T}_{\varepsilon}(\Omega)$ such that $u$ is affine on each triangle $T \in \mathbf{T}$.

Let us consider a nondecreasing continuous function $f:[0,+\infty) \rightarrow[0,+\infty)$ satisfying

$$
\begin{equation*}
f(0)=0, \quad \lim _{t \rightarrow 0^{+}} \frac{f(t)}{t}=1 \quad \text { and } \quad \lim _{t \rightarrow \infty} f(t)=\kappa \tag{1.4}
\end{equation*}
$$

for some constant $\kappa>0$, and a symmetric fourth order tensor $\mathbf{A} \in \mathscr{L}\left(\mathbb{M}_{\text {sym }}^{2 \times 2}, \mathbb{M}_{\text {sym }}^{2 \times 2}\right)$ such that

$$
\begin{equation*}
\alpha|\xi|^{2} \leq \mathbf{A} \xi: \xi \leq \beta|\xi|^{2} \quad \text { for all } \xi \in \mathbb{M}_{\mathrm{sym}}^{2 \times 2} \tag{1.5}
\end{equation*}
$$

for some constants $\alpha, \beta>0$.
Our main result is the following $\Gamma$-convergence approximation of the Griffith functional.
Theorem 1.2. The functional $\mathcal{F}_{\varepsilon}: L^{0}\left(\Omega ; \mathbb{R}^{2}\right) \rightarrow[0,+\infty]$ defined by

$$
\mathcal{F}_{\varepsilon}(u)= \begin{cases}\frac{1}{\varepsilon} \int_{\Omega} f(\varepsilon \mathbf{A} e(u): e(u)) d x & \text { if } u \in V_{\varepsilon}(\Omega)  \tag{1.6}\\ +\infty & \text { otherwise }\end{cases}
$$

$\Gamma$-converges, with respect to the $L^{0}\left(\Omega ; \mathbb{R}^{2}\right)$-topology of convergence in measure, to the Griffith functional $\mathcal{F}: L^{0}\left(\Omega ; \mathbb{R}^{2}\right) \rightarrow[0,+\infty]$ given by

$$
\mathcal{F}(u)= \begin{cases}\int_{\Omega} \mathbf{A} e(u): e(u) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{u}\right) & \text { if } u \in G S B D^{2}(\Omega) \\ +\infty & \text { otherwise }\end{cases}
$$

Remark 1.3. As explained above, a meaningful choice is the function $f(t)=t \wedge \kappa$, for which the energy reduces to

$$
\int_{\Omega} \frac{\kappa}{\varepsilon} \wedge \mathbf{A} e(u): e(u) d x
$$

It corresponds to the brittle damage energy of a linearly elastic material composed of two phases: an undamaged one whose elasticity coefficients are represented by the Hooke tensor A, and a damaged one whose elasticity coefficients are set to 0 . The constant $\kappa / \varepsilon$ stands for the toughness of the material whose diverging character as $\varepsilon \rightarrow 0$ forces the damaged zones to concentrate on vanishingly small sets (see [6]).
1.4. Strategy of proof. As usual in $\Gamma$-convergence, the proof is achieved by combining a compactness result, a lower bound and an upper bound inequality. In order to describe our argument, let us assume for simplicity that $f(t)=t \wedge \kappa$ and $\mathbf{A}=\mathrm{id}$.

Our compactness result, Proposition 4.3 rests on the general $G S B D$ compactness result of [20]. Given a sequence $\left\{u_{\varepsilon}\right\}_{\varepsilon>0}$ with uniformly bounded energy, one can apply [20, Theorem 1.1] to the modified function $v_{\varepsilon}:=u_{\varepsilon} \mathbb{1}_{\left\{\left|e\left(u_{\varepsilon}\right)\right|^{2} \leq \kappa / \varepsilon\right\}}$ which consists in putting the value zero inside each triangle $T$ where the (symmetric) gradient of $u_{\varepsilon}$ is "large". It might thus create a jump on the boundary of $T$ whose perimeter can be estimated by $\mathcal{L}^{2}(T) / \varepsilon$. It leads to compactness in measure for the sequence $\left\{u_{\varepsilon}\right\}_{\varepsilon>0}$ (up to subtracting a sequence of piecewise rigid motions, leaving the energy unchanged), which thus justifies why it is natural to consider $\Gamma$-convergence with respect to this topology.

The upper bound causes no particular difficulty. It consists in using known density results in $G S B D^{2}(\Omega)$ (see $[19,24]$ ) to reduce to the case where the jump set of $u$ is made of finitely many pairwise disjoint closed line segments, and $u$ is smooth outside. Then, considering a similar optimal triangulation of $\Omega$ as in [21] (whose vertices do not cross the jump set) and a piecewise affine Lagrange interpolation of $u$, it leads to the desired upper bound (see Proposition 3.11).

The proof of the lower bound inequality is much more delicate to address and it represents, to our opinion, the main achievement of this work. First of all, the blow-up method allows one to identify separately the bulk part and the singular part. The bulk part can be easily recovered by
modifying $u_{\varepsilon}$ into a new function which vanishes in all triangles where $e\left(u_{\varepsilon}\right)$ is "too large" as in the compactness argument (see Proposition 3.4). The main difficulty is to get a lower bound for the singular part of the energy.

Before describing our strategy of proof, let us briefly explain why the methods of [21] (and similarly [11]) fail in our situation. The idea of [21] consists in modifying every minimizing sequence $\left\{u_{\varepsilon}\right\}_{\varepsilon>0}$ inside each triangle $T$ of the associated triangulation $\mathbf{T}^{\varepsilon} \in \mathcal{T}_{\varepsilon}(\Omega)$ according to its variations along each edge of $T$. It rests on the introduction of a jump criterion which stipulates that if the variation of $u_{\varepsilon}$ is large enough, it is convenient to create a jump along the edge. More precisely, if $x_{1}, x_{2}$ and $x_{3}$ stand for the vertices of the triangle $T$, it will be energetically favorable to create a jump at the middle point of the segment $\left[x_{i}, x_{j}\right]$ if

$$
\frac{\left|u_{\varepsilon}\left(x_{i}\right)-u_{\varepsilon}\left(x_{j}\right)\right|}{\left|x_{i}-x_{j}\right|}>\frac{\sigma}{\sqrt{\varepsilon}}
$$

for some constant $\sigma>0$, while $u_{\varepsilon}$ remains unchanged on $\left[x_{i}, x_{j}\right]$ otherwise. This criterion has to be defined in such a way that:
(i) the new function, say $w_{\varepsilon}$, has a jump set in each triangle $T$ which satisfies $\mathcal{H}^{1}\left(J_{w_{\varepsilon}} \cap T\right) \leq$ $\mathcal{L}^{2}(T) /\left(\varepsilon \sin \theta_{0}\right)$, where $\theta_{0}$ is as in Definition 1.1, and $w_{\varepsilon}$ does not jump across $\partial T$;
(ii) the absolutely continuous part of the gradient, $\nabla v_{\varepsilon}$, is controlled in $L^{2}(T)$ by the energy restricted to $T$.
This construction ensures that the variation of the new discontinuous and piecewise affine function $w_{\varepsilon}$ is always controlled along at least two edges of each triangle $T$, and it yields a control of the full gradient $\nabla w_{\varepsilon}$ of $w_{\varepsilon}$ inside $T$. In [21], this is possible thanks of the scalar nature of the problem because the gradient $\nabla w_{\varepsilon \mid T}$ is a (constant) vector in $\mathbb{R}^{2}$ (see [21, Remark 3.5]).

In our case, $u_{\varepsilon}$ is not scalar-valued anymore, but vector-valued and the energy only depends on its symmetric gradient $e\left(u_{\varepsilon}\right)$. If one uses the same criterion than in [21], then condition (i) above will be satisfied for the new function $w_{\varepsilon}$ on $T$. However, one will only be able to estimate the $L^{2}$ norm of the (symmetric) gradient of $w_{\varepsilon}$ by that of the full gradient of $u_{\varepsilon}$ which, unfortunately, is not controlled by the energy $\mathcal{F}_{\varepsilon}\left(u_{\varepsilon}\right)$. Note that in [34], such a control is artificially made possible thanks to the particular form of the energy (see (1.3) above). This is however not natural in this linearized elasticity setting where the energy should be expressed in terms of the symmetric gradient of the displacement.

As a consequence, the jump criterion has to be modified. As the energy only depends on the symmetric part of the gradient of $u_{\varepsilon}$, it would be natural to consider a criterion involving the longitudinal variation of $u_{\varepsilon}$ along the edges of the triangle instead of the full variation. In other words, one could modify the criterion by asking that if

$$
\frac{\left|\left(u_{\varepsilon}\left(x_{i}\right)-u_{\varepsilon}\left(x_{j}\right)\right) \cdot\left(x_{i}-x_{j}\right)\right|}{\left|x_{i}-x_{j}\right|^{2}}>\frac{\sigma}{\sqrt{\varepsilon}},
$$

then we create a jump at the middle point of $\left[x_{i}, x_{j}\right]$, while $u_{\varepsilon}$ remains unchanged on $\left[x_{i}, x_{j}\right]$ otherwise. In that case, it is again not possible to control the symmetric gradient $e\left(w_{\varepsilon}\right)$ of the new function $w_{\varepsilon}$ by that of $u_{\varepsilon}$. Indeed, in a similar way as in [21], the previous criterion ensures that the longitudinal variation of $w_{\varepsilon}$ along at least two edges of each triangle is controlled by the energy restricted to $T$. If we call $\xi_{1}$ and $\xi_{2} \in \mathbb{S}^{1}$ both (linearly independent) directions associated to these "good" edges, it shows that $e\left(w_{\varepsilon}\right)_{\mid T}:\left(\xi_{1} \otimes \xi_{1}\right)$ and $e\left(w_{\varepsilon}\right)_{\mid T}:\left(\xi_{2} \otimes \xi_{2}\right)$ are controlled by $e\left(u_{\varepsilon}\right)_{\mid T}$ which is not enough to control the full $2 \times 2$ symmetric matrix $e\left(w_{\varepsilon}\right)_{\mid T}$ which has three degrees of freedom. In addition, some (uncontrolled) discontinuities can also be created at the interface $I:=\partial T \cap \partial T^{\prime}$ between two adjacent triangles $T$ and $T^{\prime}$ so that condition (i) fails as well.

Overcoming these difficulties seems to be a very serious issue so that we decided to attack this problem from a different angle. First of all, the use of the blow-up method allows one to reduce to
the case where $\Omega=B$ is the unit ball, $u$ is a step function of the form

$$
u(x)= \begin{cases}a & \text { if } x \cdot \nu<0 \\ b & \text { if } x \cdot \nu>0\end{cases}
$$

for some $a, b \in \mathbb{R}^{2}$ with $a \neq b$ and $\nu \in \mathbb{S}^{1}$ (with a jump set corresponding to the diameter of $B$ orthogonal to $\nu$ ), and, see Lemma 3.6, such that

$$
\begin{equation*}
\int_{\left\{\left|e\left(u_{\varepsilon}\right)\right|^{2} \leq \kappa / \varepsilon\right\}}\left|e\left(u_{\varepsilon}\right)\right|^{2} d x \rightarrow 0 \tag{1.7}
\end{equation*}
$$

To make our strategy of proof more transparent, we assume that $a \cdot \nu \neq b \cdot \nu$. A standard argument based on Fubini's Theorem shows that the one-dimensional section of $u_{\varepsilon}$ in the direction $\nu$ passing through the point $y$, namely $t \mapsto\left(u_{\varepsilon}\right)_{y}^{\nu}(t):=u_{\varepsilon}(y+t \nu) \cdot \nu$ converges (in measure) to the step function

$$
t \mapsto u_{y}^{\nu}(t)=a \cdot \nu \mathbb{1}_{\mathbb{R}^{-}}+b \cdot \nu \mathbb{1}_{\mathbb{R}^{+}}
$$

Let us denote by $\mathbf{T}^{\varepsilon}$ the triangulation on which $u_{\varepsilon}$ is (continuous and) piecewise affine. We further denote by $\mathbf{T}_{b}^{\varepsilon}$ the familly of all triangles $T \in \mathbf{T}^{\varepsilon}$ such that $\left|e\left(u_{\varepsilon}\right)_{\mid T}\right|^{2}>\kappa / \varepsilon$. Thanks to (1.7), we show that almost every line orthogonal to $J_{u} \cap B$ must cross at least one triangle $T \in \mathbf{T}_{b}^{\varepsilon}$ (see Lemma 3.7). The reason is that if, for some $y \in J_{u} \cap B$, the line $y+\mathbb{R} \nu$ intersects no such triangles, then $\left(u_{\varepsilon}\right)_{y}^{\nu}$ would be bounded in $H^{1}$ (because $\left.\left|\left(\left(u_{\varepsilon}\right)_{y}^{\nu}\right)^{\prime}\right| \leq\left|e\left(u_{\varepsilon}\right)(y+t \nu)\right|\right)$ and thus, it would converge weakly in that space to a constant function, contradicting that $a \cdot \nu \neq b \cdot \nu$. This information allows one to get a bad lower bound for the surface energy with $1 / 2$ multiplicative factor. It suggests to improve the previous argument by showing that "many" lines $y+\mathbb{R} \nu$ passing through $y \in J_{u} \cap B$ must actually cross at least two triangles in $\mathbf{T}_{b}^{\varepsilon}$, which is the object of Lemma 3.8. To do that, we show in Lemma 3.9 that there are very few points $y$ in $J_{u} \cap B$ such that the line $y+\mathbb{R} \nu$ crosses exactly one triangle $T \in \mathbf{T}_{b}^{\varepsilon}$. Indeed, in that case, up to a small error, the function $\left(u_{\varepsilon}\right)_{y}^{\nu}$ would have to pass from the value $a \cdot \nu$ to $b \cdot \nu$ inside $T$. Due to the particular shape of a triangle and of the fact that $u_{\varepsilon}$ is affine inside $T$, this could only happen for at most two values of $y$. Moreover, if $y$ is far away from these two values, the variation of $\left(u_{\varepsilon}\right)_{y}^{\nu}$ across the triangle $T$ is not sufficient, and it becomes necessary to cross an additional triangle $T^{\prime}$ in $\mathbf{T}_{b}^{\varepsilon}$. With this improvement, we can now construct two disjoint families of triangles with the property that both families project onto $J_{u} \cap B$ into two sets of almost full $\mathcal{H}^{1}$ measure (see Lemma 3.10). It enables one to compensate the bad multiplicative factor $1 / 2$ in the previous argument, and obtain the expected lower bound with the correct constant corresponding to $\kappa \sin \theta_{0}$ (see Proposition 3.5).

To conclude this introduction, let us mention that the originality of this work is twofold. First of all, we are able to provide a deterministic discrete finite element approximation result of the Griffith functional with isotropic surface energies. In particular, our approach does not require any unnatural dependence of the approximating energy with respect to the skew symmetric part of the gradient (in the context of linear elasticity) nor the use of stochastic meshes. Second, our method relies on an unusual application of the slicing method, which is rather employed in $\Gamma$-convergence analysis to reduce the dimension of the problem to a one-dimension study. Here, we instead use this method as a tool to enumerate in a non trivial way the number of triangles needed to derive the correct multiplicity in the surface energy.
1.5. Organisation of the paper. In Section 2, we collect useful notation and preliminary results that will be useful in the subsequent sections. Section 3 is devoted to show our main result, Theorem 1.2. It is divided into three parts: a first one consisting in a compactness result, Proposition 3.1, a second one corresponding to the lower bound inequality, Proposition 3.3, and a last one for the upper bound inequality, Proposition 3.11 through the construction of a recovery sequence. Eventually, in Section 4, we extend the previous $\Gamma$-convergence analysis allowing for Dirichlet
boundary conditions formulated in a suitable way at the discrete and continuum levels (see Theorem 4.1). We then deduce the fundamental property of $\Gamma$-convergence, Corollary 4.2 , in our specific setting, i.e., the convergence of minimizers as well as the minimum value.

## 2. Notation and Preliminaries

Vectors. The Euclidean scalar product between two vectors $x$ and $y \in \mathbb{R}^{n}$ is denoted by $x \cdot y$, and the associated Euclidean norm by $|x|:=\sqrt{x \cdot x}$. For $x \in \mathbb{R}^{n}$ and $\varrho>0$, we denote by $B_{\varrho}(x):=\left\{y \in \mathbb{R}^{n}:|x-y|<\varrho\right\}$ the open ball centered at $x$ with radius $\varrho$. If $x=0$, we simply write $B_{\varrho}$ instead of $B_{\varrho}(0)$. The notation $\mathbb{S}^{n-1}$ stands for the unit sphere $\partial B_{1}$.

Matrices. The set of all real $m \times n$ matrices is denoted by $\mathbb{M}^{m \times n}$, and the subset of symmetric real $n \times n$ matrices by $\mathbb{M}_{\mathrm{sym}}^{n \times n}$. It will be endowed with the Froebenius scalar product $A: B:=\operatorname{tr}\left(A^{T} B\right)$ and the associated Froebenius norm $|A|:=\sqrt{A: A}$.

Given two vectors $a$ and $b \in \mathbb{R}^{n}$, the tensor product between $a$ and $b$ is defined as $a \otimes b:=a b^{T} \in$ $\mathbb{M}^{n \times n}$ and the symmetric tensor product by $a \odot b:=(a \otimes b+b \otimes a) / 2 \in \mathbb{M}_{\text {sym }}^{n \times n}$.
Measures. The Lebesgue and the $k$-dimensional Hausdorff measures in $\mathbb{R}^{n}$ are respectively denoted by $\mathcal{L}^{n}$ and $\mathcal{H}^{k}$. If $U$ is a bounded open set of $\mathbb{R}^{n}$ and $Y$ is an Euclidean space, we denote by $\mathcal{M}(U ; Y)$ the space of $Y$-valued bounded Radon measures in $U$ which, according to the Riesz Representation Theorem, can be identified to the dual of $C_{0}(U ; Y)$ (the closure of $C_{c}(U ; Y)$ for the sup-norm in $U)$. For $\mu \in \mathcal{M}(U ; Y)$, its total variation is denoted by $|\mu|$.

Functional spaces. We use standard notation for Lebesgue and Sobolev spaces. If $U$ is a bounded open subset of $\mathbb{R}^{n}$, we denote by $L^{0}\left(U ; \mathbb{R}^{m}\right)$ the set of all $\mathcal{L}^{n}$-measurable functions from $U$ to $\mathbb{R}^{m}$. We recall that a sequence $\left\{g_{k}\right\}_{k \in \mathbb{N}}$ in $L^{0}\left(U ; \mathbb{R}^{m}\right)$ converges in measure to $g \in L^{0}\left(U ; \mathbb{R}^{m}\right)$ if for all $\varepsilon>0$,

$$
\mathcal{L}^{n}\left(\left\{x \in U:\left|g_{k}(x)-g(x)\right|>\varepsilon\right\}\right) \rightarrow 0
$$

Note that, for any fixed constant $M>0$, we can define the following mapping

$$
\begin{equation*}
d_{M}:(g, h) \in L^{0}\left(U ; \mathbb{R}^{m}\right) \times L^{0}\left(U ; \mathbb{R}^{m}\right) \mapsto \int_{U} M \wedge|g-h| d x \in \mathbb{R}^{+} \tag{2.1}
\end{equation*}
$$

which turns out to be a distance over $L^{0}\left(U ; \mathbb{R}^{m}\right)$, with the property that $g_{k}$ converges in measure to $g$ if and only if $d_{M}\left(g_{k}, g\right) \rightarrow 0$. It confers to $L^{0}\left(U ; \mathbb{R}^{m}\right)$ a metric space structure.

Functions of bounded variation and sets of finite perimeter. We refer to [3] for an exhaustive treatment on that subject and just recall few notation. Let $U \subset \mathbb{R}^{n}$ be a bounded open set. A function $u \in L^{1}\left(U ; \mathbb{R}^{m}\right)$ is a function of bounded variation in $U$, and we write $u \in B V\left(U ; \mathbb{R}^{m}\right)$, if its distributional derivative $D u$ belongs to $\mathcal{M}\left(U ; \mathbb{M}^{m \times n}\right)$. We use standard notation for that space, referring to [3] for details. We just recall that a function $u$ belongs to $S B V^{2}\left(U ; \mathbb{R}^{m}\right)$ if $u \in S B V\left(U ; \mathbb{R}^{m}\right)$ (the distributional derivative $D u$ has no Cantor part), its approximate gradient $\nabla u$ belongs to $L^{2}\left(U ; \mathbb{M}^{m \times n}\right)$ and its jump set $J_{u}$ satisfies $\mathcal{H}^{n-1}\left(J_{u}\right)<\infty$.

A Lebesgue measurable set $A \subset \mathbb{R}^{n}$ is a set of finite perimeter in $U$ if its characteristic function $\mathbb{1}_{A}$ belongs to $B V\left(U ; \mathbb{R}^{n}\right)$. The reduced boundary of $A$ is denoted by $\partial^{*} A$ and the essential (or measure theoretic) boundary is denoted by $\partial_{*} A$. For every $t \in[0,1]$, we denote by $A^{(t)}$ the set of points where $A$ has density $t$.

We also recall that a partition $\mathcal{P}=\left\{P_{i}\right\}_{i \in \mathbb{N}}$ of an open set $U$ is a Cacciopoli partition if each $P_{i}$ have finite perimeter in $U$, and $\sum_{i \in \mathbb{N}}\left|D \mathbb{1}_{P_{i}}\right|(U)<\infty$. In that case,

$$
\bigcup_{i \in \mathbb{N}}\left(P_{i}\right)^{(1)} \cup \bigcup_{i, j \in \mathbb{N}, i \neq j} \partial^{*} P_{i} \cap \partial^{*} P_{j}
$$

contains $\mathcal{H}^{n-1}$-almost all of $U$ (see [3, Section 4.4]). In the sequel (as in [20, Theorem 2.5]), we will sometimes use the following notation for Caccioppoli partitions:

$$
\mathcal{P}^{(1)}:=\bigcup_{i \in \mathbb{N}} P_{i}^{(1)}, \quad \partial^{*} \mathcal{P}:=\bigcup_{i \in \mathbb{N}} \partial^{*} P_{i}
$$

(Generalized) functions of bounded deformation. A function $u \in L^{1}\left(U ; \mathbb{R}^{n}\right)$ is a function of bounded deformation, and we write $u \in B D(U)$, if its distributional symmetric gradient $E u:=$ $\left(D u+D u^{T}\right) / 2$ belongs to $\mathcal{M}\left(U ; \mathbb{M}_{\text {sym }}^{n \times n}\right)$. We refer to [40, 41, 39, 2, 10] for the main properties and notation of that space. The space $S B D^{2}(U)$ is made of all functions $u \in S B D(U)$ ( $E u$ has no Cantor part) such that the approximate symmetric gradient $e(u)$ (the absolutely continuous part of $E u$ with respect to $\left.\mathcal{L}^{n}\right)$ belongs to $L^{2}\left(U ; \mathbb{M}_{\text {sym }}^{n \times n}\right)$ and its jump set $J_{u}$ satisfies $\mathcal{H}^{n-1}\left(J_{u}\right)<\infty$.

We now recall the definition and the main properties of the space of generalized functions of bounded deformation introduced in [26]. We first need to introduce some notation. Let $\xi \in \mathbb{S}^{n-1}$, we denote by $\Pi_{\xi}:=\left\{y \in \mathbb{R}^{n}: y \cdot \xi=0\right\}$ the orthogonal space to $\xi$ and by $p_{\xi}$ the orthogonal projection onto $\Pi_{\xi}$. For every set $B \subset \mathbb{R}^{n}$, we define for $\xi \in \mathbb{S}^{n-1}$ and $y \in \mathbb{R}^{n}$,

$$
B_{y}^{\xi}:=\{t \in \mathbb{R}: y+t \xi \in B\}, \quad B^{\xi}:=p_{\xi}(B)
$$

and, for every (vector-valued) function $u: B \rightarrow \mathbb{R}^{n}$ and (scalar-valued) function $f: B \rightarrow \mathbb{R}$,

$$
u_{y}^{\xi}(t):=u(y+t \xi) \cdot \xi, \quad f_{y}^{\xi}(t)=f(y+t \xi) \quad \text { for all } y \in \mathbb{R}^{n} \text { and all } t \in B_{y}^{\xi}
$$

Definition 2.1. Let $U \subset \mathbb{R}^{n}$ be a bounded open set and $u \in L^{0}\left(U ; \mathbb{R}^{n}\right)$. Then, $u \in G B D(U)$ if there exists a nonnegative measure $\lambda \in \mathcal{M}(U)$ such that one of the following equivalent conditions holds true for every $\xi \in \mathbb{S}^{n-1}$ :
(1) for every $\tau \in C^{1}(\mathbb{R})$ with $-\frac{1}{2} \leq \tau \leq \frac{1}{2}$ and $0 \leq \tau^{\prime} \leq 1$, the partial derivative $D_{\xi}(\tau(u \cdot \xi))=$ $D(\tau(u \cdot \xi)) \cdot \xi$ belongs to $\mathcal{M}(U)$, and

$$
\left|D_{\xi}(\tau(u \cdot \xi))\right|(B) \leq \lambda(B) \quad \text { for every Borel set } B \subset U
$$

(2) $u_{y}^{\xi} \in B V_{\mathrm{loc}}\left(U_{y}^{\xi}\right)$ for $\mathcal{H}^{n-1}$-a.e. $y \in U^{\xi}$, and

$$
\int_{\Pi_{\xi}}\left(\left|D u_{y}^{\xi}\right|\left(B_{y}^{\xi} \backslash J_{u_{y}^{\xi}}^{1}\right)+\mathcal{H}^{0}\left(B_{y}^{\xi} \cap J_{u_{y}^{\xi}}^{1}\right)\right) d \mathcal{H}^{n-1}(y) \leq \lambda(B) \quad \text { for every Borel set } B \subset U
$$

where $J_{u_{y}^{\xi}}^{1}:=\left\{t \in J_{u_{y}^{\xi}}:\left|\left[u_{y}^{\xi}\right](t)\right| \geq 1\right\}$.
The function $u$ belongs to $G S B D(U)$ if $u \in G B D(U)$ and $u_{y}^{\xi} \in S B V_{\text {loc }}\left(U_{y}^{\xi}\right)$ for every $\xi \in \mathbb{S}^{n-1}$ and for $\mathcal{H}^{n-1}$-a.e. $y \in U^{\xi}$.

Every $u \in G B D(U)$ has an approximate symmetric gradient $e(u) \in L^{1}\left(U ; \mathbb{M}_{\text {sym }}^{n \times n}\right)$ such that for every $\xi \in \mathbb{S}^{n-1}$ and for $\mathcal{H}^{n-1}$-a.e. $y \in U^{\xi}$,

$$
e(u)(y+t \xi) \xi \cdot \xi=\left(u_{y}^{\xi}\right)^{\prime}(t) \quad \text { for } \mathcal{L}^{1} \text {-a.e. } t \in U_{y}^{\xi}
$$

Moreover, the jump set $J_{u}$ of $u \in G B D(U)$, defined as the set of all $x_{0} \in U$ for which there exist $\left(u^{+}\left(x_{0}\right), u^{-}\left(x_{0}\right), \nu_{u}\left(x_{0}\right)\right) \in \mathbb{R}^{n} \times \mathbb{R}^{n} \times \mathbb{S}^{n-1}$ with $u^{+}\left(x_{0}\right) \neq u^{-}\left(x_{0}\right)$ such that the function

$$
y \in B_{1} \mapsto u_{x_{0}, \varrho}:=u\left(x_{0}+\varrho y\right)
$$

converges in measure in $B_{1}$ as $\varrho \searrow 0$ to

$$
y \in B_{1} \mapsto \begin{cases}u^{+}\left(x_{0}\right) & \text { if } y \cdot \nu_{u}\left(x_{0}\right)>0 \\ u^{-}\left(x_{0}\right) & \text { if } y \cdot \nu_{u}\left(x_{0}\right) \leq 0\end{cases}
$$

is countably $\left(\mathcal{H}^{n-1}, n-1\right)$-rectifiable. Finally, the energy space $G S B D^{2}(U)$ is defined as

$$
G S B D^{2}(U):=\left\{u \in G S B D(U): e(u) \in L^{2}\left(U ; \mathbb{M}_{\mathrm{sym}}^{n \times n}\right), \mathcal{H}^{n-1}\left(J_{u}\right)<\infty\right\}
$$

## 3. Proof of the main result

Let us introduce the $\Gamma$-lower and upper limits (with respect to the topology of convergence in measure) $\mathcal{F}^{\prime}$ and $\mathcal{F}^{\prime \prime}: L^{0}\left(\Omega ; \mathbb{R}^{2}\right) \rightarrow[0,+\infty]$ defined by

$$
\mathcal{F}^{\prime}(u):=\inf \left\{\liminf _{\varepsilon \rightarrow 0} \mathcal{F}_{\varepsilon}\left(u_{\varepsilon}\right): u_{\varepsilon} \rightarrow u \text { in measure in } \Omega\right\}
$$

and

$$
\mathcal{F}^{\prime \prime}(u):=\inf \left\{\limsup _{\varepsilon \rightarrow 0} \mathcal{F}_{\varepsilon}\left(u_{\varepsilon}\right): u_{\varepsilon} \rightarrow u \text { in measure in } \Omega\right\}
$$

for all $u \in L^{0}\left(\Omega ; \mathbb{R}^{2}\right)$.
3.1. Domain of the $\Gamma$-limit. We begin our analysis by identifying the domain of finiteness of the $\Gamma$-limit.

Proposition 3.1. Let $\left\{\varepsilon_{k}\right\}_{k \in \mathbb{N}}$ satisfying $\varepsilon_{k} \rightarrow 0, u \in L^{0}\left(\Omega ; \mathbb{R}^{2}\right)$ and $\left\{u_{k}\right\}_{k \in \mathbb{N}} \subset L^{0}\left(\Omega ; \mathbb{R}^{2}\right)$ be such that $M:=\sup _{k} \mathcal{F}_{\varepsilon_{k}}\left(u_{k}\right)<\infty$ and $u_{k} \rightarrow u$ in measure in $\Omega$. Then, $u \in G S B D^{2}(\Omega)$.
Proof. According to the properties (1.4) satisfied by $f$, for all $\delta>0$, there exists a constant $0<K<\kappa$ such that

$$
f(t) \geq K \wedge[(1-\delta) t] \quad \text { for all } t \geq 0
$$

Indeed, since $f(t) / t \rightarrow 1$ as $t \rightarrow 0^{+}$, there exists $t^{*}>0$ such that $f(t) / t \geq 1-\delta$ for all $t \in\left[0, t^{*}\right]$ and $K:=(1-\delta) t^{*}<\kappa$. Hence, for all $t \in\left[0, t^{*}\right]$, we have $f(t) \geq(1-\delta) t$, while for all $t>t^{*}$, as $f$ is nondecreasing, $f(t) \geq f\left(t^{*}\right) \geq K$.

By definition of $\mathcal{F}_{\varepsilon_{k}}$, there exists a triangulation $\mathbf{T}^{k} \in \mathcal{T}_{\varepsilon_{k}}(\Omega)$ such that $u_{k} \in V_{\varepsilon_{k}}(\Omega)$ is affine on each triangle $T \in \mathbf{T}^{k}$. We introduce the characteristic functions

$$
\chi_{k}:=\mathbb{1}_{\left\{(1-\delta) \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) \geq \frac{K}{\varepsilon_{k}}\right\}} \in L^{\infty}(\Omega ;\{0,1\})
$$

which are constant on each triangle $T \in \mathbf{T}^{k}$, so that

$$
D_{k}:=\left\{\chi_{k}=1\right\} \cap \Omega=\bigcup_{i=1}^{N_{k}}\left(T_{i}^{k} \cap \Omega\right)
$$

for some triangles $T_{i}^{k} \in \mathbf{T}^{k}$. Remark that this choice of $\chi_{k}$ implies that

$$
M \geq \mathcal{F}_{\varepsilon_{k}}\left(u_{k}\right) \geq(1-\delta) \int_{\Omega}\left(1-\chi_{k}\right) \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) d x+\frac{K}{\varepsilon_{k}} \int_{\Omega} \chi_{k} d x
$$

forcing $\chi_{k}$ to converge to 0 in $L^{1}(\Omega)$ since $0 \leq \int_{\Omega} \chi_{k} d x \leq K^{-1} M \varepsilon_{k} \rightarrow 0$.
Let $v_{k}:=\left(1-\chi_{k}\right) u_{k}$ so that, by $\left[3\right.$, Theorem 3.84], $v_{k} \in S B V^{2}\left(\Omega ; \mathbb{R}^{2}\right)$ with $\nabla v_{k}=\left(1-\chi_{k}\right) \nabla u_{k}$ and

$$
J_{v_{k}} \subset \Omega \cap \partial D_{k} \subset \bigcup_{i=1}^{N_{k}} \partial T_{i}^{k}
$$

Note that

$$
\begin{equation*}
v_{k} \rightarrow u \text { in measure in } \Omega \text { and } A:=\left\{x \in \Omega:\left|u_{k}(x)\right| \rightarrow \infty\right\} \text { is } \mathcal{L}^{2} \text {-negligible. } \tag{3.1}
\end{equation*}
$$

Indeed, since $u_{k} \rightarrow u$ in measure in $\Omega$ and $\left\{u_{k} \neq v_{k}\right\} \subset D_{k}$ with $\mathcal{L}^{2}\left(D_{k}\right) \rightarrow 0$, for all $\eta>0$, we get that $\mathcal{L}^{2}\left(\left\{\left|v_{k}-u\right|>\eta\right\}\right) \leq \mathcal{L}^{2}\left(\left\{\left|u_{k}-u\right|>\eta\right\}\right)+\mathcal{L}^{2}\left(D_{k}\right) \rightarrow 0$. Additionally, up to a subsequence (not relabeled), $u_{k}(x) \rightarrow u(x) \in \mathbb{R}^{2}$ for $\mathcal{L}^{2}$-a.e. $x \in \Omega$.

On the one hand, using the energy bound $\mathcal{F}_{\varepsilon_{k}}\left(u_{k}\right) \leq M$ and the ellipticity property (1.5) of $\mathbf{A}$, we infer that

$$
\begin{equation*}
\int_{\Omega}\left|e\left(v_{k}\right)\right|^{2} d x \leq \frac{M}{(1-\delta) \alpha} \tag{3.2}
\end{equation*}
$$

On the other hand, by definition of an admissible triangulation, the edges of each triangle $T_{i}^{k}$ have length greater than or equal to $\varepsilon_{k}$ and their angles are all greater than or equal to $\theta_{0}$, so that the heights of such triangles must be greater than or equal to $\varepsilon_{k} \sin \theta_{0}$. Therefore, for all $1 \leq i \leq N_{k}$,

$$
\mathcal{L}^{2}\left(T_{i}^{k}\right) \geq \frac{1}{2}\left(\varepsilon_{k} \sin \theta_{0}\right) \frac{\mathcal{H}^{1}\left(\partial T_{i}^{k}\right)}{3}
$$

which implies that for all open subset $U \subset \subset \Omega$ :

$$
\mathcal{H}^{1}\left(J_{v_{k}} \cap U\right) \leq \frac{6}{\sin \theta_{0}} \sum_{i \in\left\{1, \ldots, N_{k}\right\}, T_{i}^{k} \cap U \neq \emptyset} \frac{\mathcal{L}^{2}\left(T_{i}^{k}\right)}{\varepsilon_{k}}
$$

Let $k_{U} \geq 1$ (depending on $U$ ) be such that for all $k \geq k_{U}$, any triangle $T \in \mathbf{T}^{k}$ intersecting $U$ is contained in $\Omega$, then it follows that for all $k \geq k_{U}$,

$$
\begin{equation*}
\mathcal{H}^{1}\left(J_{v_{k}} \cap U\right) \leq \frac{6}{\varepsilon_{k} \sin \theta_{0}} \int_{\Omega} \chi_{k} d x \leq \frac{6 M}{K \sin \theta_{0}} \tag{3.3}
\end{equation*}
$$

where we used once more the energy bound $\mathcal{F}_{\varepsilon_{k}}\left(u_{k}\right) \leq M$.
Gathering (3.2) and (3.3), we can apply the $G S B D^{2}$-compactness Theorem ([18, Theorem 1.1]). Together with (3.1), it ensures the existence of a subsequence (depending on the open subset $U$, which we do not relabel) such that $u_{\mid U} \in G S B D^{2}(U)$,

$$
e\left(v_{k}\right)_{\mid U} \rightharpoonup e\left(u_{\mid U}\right) \text { weakly in } L^{2}\left(U ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right) \text { and } \mathcal{H}^{1}\left(J_{u} \cap U\right) \leq \liminf _{k \rightarrow \infty} \mathcal{H}^{1}\left(J_{v_{k}} \cap U\right)
$$

We then consider an exhaustion of $\Omega$ by a sequence of open subsets $\left\{U_{m}\right\}_{m \in \mathbb{N}}$ satisfying $U_{m} \subset \subset$ $U_{m+1} \subset \subset \Omega$ for all $m \in \mathbb{N}$ and $\bigcup_{m} U_{m}=\Omega$. Using a diagonal extraction argument, we can find a subsequence (still denoted by $\left\{v_{k}\right\}_{k \in \mathbb{N}}$ ) such that for all $m \in \mathbb{N}, u_{\mid U_{m}} \in G S B D^{2}\left(U_{m}\right)$ and

$$
\begin{equation*}
e\left(v_{k}\right)_{\mid U_{m}} \rightharpoonup e\left(u_{\mid U_{m}}\right) \text { weakly in } L^{2}\left(U_{m} ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right) \text { and } \mathcal{H}^{1}\left(J_{u} \cap U_{m}\right) \leq \liminf _{k \rightarrow \infty} \mathcal{H}^{1}\left(J_{v_{k}} \cap U_{m}\right) \tag{3.4}
\end{equation*}
$$

Let us now check that $u$ belongs to $G S B D^{2}(\Omega)$. Indeed, let $\xi \in \mathbb{S}^{1}$ and $\tau \in C^{1}(\mathbb{R})$ be such that $|\tau| \leq \frac{1}{2}$ and $0 \leq \tau^{\prime} \leq 1$. For all test function $\phi \in C_{c}^{\infty}(\Omega)$, there exists $m \in \mathbb{N}$ such that $\operatorname{supp} \phi \subset U_{m}$ so that, owing to the dominated convergence Theorem,

$$
\left\langle D_{\xi}(\tau(u \cdot \xi)), \phi\right\rangle=-\int_{U_{m}} \tau(u \cdot \xi) D_{\xi} \phi d x=-\lim _{k \rightarrow \infty} \int_{U_{m}} \tau\left(v_{k} \cdot \xi\right) D_{\xi} \phi d x=\lim _{k \rightarrow \infty}\left\langle D_{\xi}\left(\tau\left(v_{k} \cdot \xi\right)\right), \phi\right\rangle
$$

Since $v_{k} \cdot \xi \in S B V^{2}(\Omega)$, using the chain rule formula in $B V$ ([3, Theorem 3.96]), we get that $\tau\left(v_{k} \cdot \xi\right) \in S B V^{2}(\Omega)$ with

$$
D_{\xi}\left(\tau\left(v_{k} \cdot \xi\right)\right)=\tau^{\prime}\left(v_{k} \cdot \xi\right) e\left(v_{k}\right):(\xi \otimes \xi) \mathcal{L}^{2}\left\llcorner\Omega+\left(\tau\left(v_{k}^{+} \cdot \xi\right)-\tau\left(v_{k}^{-} \cdot \xi\right)\right) \nu_{v_{k}} \cdot \xi \mathcal{H}^{1}\left\llcorner J_{v_{k}}\right.\right.
$$

Taking the variation, we infer that

$$
\left|D_{\xi}\left(\tau\left(v_{k} \cdot \xi\right)\right)\right| \leq\left|e\left(v_{k}\right)\right| \mathcal{L}^{2}\left\llcorner\Omega+\mathcal{H}^{1}\left\llcorner J_{v_{k}}=: \lambda_{k}\right.\right.
$$

As a consequence of (3.2) together with (3.3), the sequence $\left\{\lambda_{k}\right\}_{k \in \mathbb{N}}$ is bounded in $\mathcal{M}\left(U_{m}\right)$ for all $m \in \mathbb{N}$, with

$$
\sup _{k \geq k_{U_{m}}} \lambda_{k}\left(U_{m}\right) \leq \frac{M}{(1-\delta) \alpha}+\frac{6 M}{K \sin \theta_{0}}=: M_{\delta}<+\infty
$$

so that, up to a further diagonal extraction, $\lambda_{k}\left\llcorner U_{m} \rightharpoonup \lambda^{(m)}\right.$ weakly* in $\mathcal{M}(\Omega)$ for some nonnegative measure $\lambda^{(m)} \in \mathcal{M}(\Omega)$ satisfying, for all $m \in \mathbb{N}$,

$$
\lambda^{(m)}(\Omega) \leq \liminf _{k \rightarrow \infty} \lambda_{k}\left(U_{m}\right) \leq M_{\delta}
$$

Therefore, we can introduce the following nonnegative measure $\lambda \in \mathcal{M}(\Omega)$ defined by

$$
\lambda(B):=\sup _{m \in \mathbb{N}} \lambda^{(m)}(B)=\lim _{m \rightarrow \infty} \lambda^{(m)}(B) \text { for all Borel subset } B \subset \Omega
$$

We thus obtain that

$$
\left|\left\langle D_{\xi}(\tau(u \cdot \xi)), \phi\right\rangle\right| \leq \lim _{k \rightarrow \infty}\left\langle\lambda_{k}\left\llcorner U_{m},|\phi|\right\rangle=\left\langle\lambda^{(m)},\right| \phi \mid\right\rangle \leq\langle\lambda,| \phi| \rangle
$$

implying both that $D_{\xi}(\tau(u \cdot \xi)) \in \mathcal{M}(\Omega)$ according to Riesz Representation Theorem and that

$$
\left|D_{\xi}(\tau(u \cdot \xi))\right| \leq \lambda \quad \text { in } \mathcal{M}(\Omega)
$$

which shows that $u \in G B D(\Omega)$. Using next that $u \in G S B D\left(U_{m}\right)$ for all $m \in \mathbb{N}$ and [26, Definition 4.2], we deduce that $u \in G S B D(\Omega)$. Eventually, by locality of the definition of the approximate symmetric gradient $e(u)$ (see [26, Formula (9.1)]), as a consequence of (3.2) together with (3.4), we infer that $e\left(v_{k}\right) \rightharpoonup e(u)$ weakly in $L^{2}\left(\Omega ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right)$ with $e(u) \in L^{2}\left(\Omega ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right)$. Passing to the limit as $m \rightarrow \infty$ in the last property of (3.4) and using (3.3) shows that $\mathcal{H}^{1}\left(J_{u}\right)<\infty$. All of this establishes that $u \in G S B D^{2}(\Omega)$ and completes the proof of the Proposition.

Remark 3.2. We will later improve the previous result (see Proposition 4.3) by getting rid-off the a priori knowledge that $u_{k}$ converges in measure in $\Omega$. The price to pay will be to subtract a sequence of piecewise rigid body motions. Proposition 4.3 will a posteriori justify why the topology of convergence in measure is the natural one to address the $\Gamma$-convergence analysis.
3.2. The lower bound. The proof of the lower bound inequality relies on the blow up method which consists in identifying separately the Lebesgue and jump parts of the energy.
Proposition 3.3. For all $u \in L^{0}\left(\Omega ; \mathbb{R}^{2}\right)$,

$$
\mathcal{F}(u) \leq \mathcal{F}^{\prime}(u)
$$

Proof. Without loss of generality, we can assume that $\mathcal{F}^{\prime}(u)<\infty$. For any $\zeta>0$, there exists a sequence $\left\{u_{\varepsilon}\right\}_{\varepsilon>0}$ such that $u_{\varepsilon} \rightarrow u$ in measure in $\Omega$ and

$$
\liminf _{\varepsilon \rightarrow 0} \mathcal{F}_{\varepsilon}\left(u_{\varepsilon}\right) \leq \mathcal{F}^{\prime}(u)+\zeta<\infty
$$

Let us extract a subsequence $\left\{u_{k}\right\}_{k \in \mathbb{N}}:=\left\{u_{\varepsilon_{k}}\right\}_{k \in \mathbb{N}}$ from $\left\{u_{\varepsilon}\right\}_{\varepsilon>0}$ such that $u_{k} \rightarrow u \mathcal{L}^{2}$-a.e. in $\Omega$ and

$$
\lim _{k \rightarrow \infty} \mathcal{F}_{\varepsilon_{k}}\left(u_{k}\right)=\liminf _{\varepsilon \rightarrow 0} \mathcal{F}_{\varepsilon}\left(u_{\varepsilon}\right)<\infty
$$

This implies that, for $k$ large enough, $u_{k} \in V_{\varepsilon_{k}}(\Omega)$ and $\sup _{k} \mathcal{F}_{\varepsilon_{k}}\left(u_{k}\right)<\infty$. By definition of the finite element space $V_{\varepsilon_{k}}(\Omega)$, there exists a triangulation $\mathbf{T}^{k} \in \mathcal{T}_{\varepsilon_{k}}(\Omega)$ such that $u_{k}$ is affine on each $T \in \mathbf{T}^{k}$.

We first note that, according to Proposition 3.1, $u \in G S B D^{2}(\Omega)$. Let us show the lower bound inequality $\mathcal{F}^{\prime}(u) \geq \mathcal{F}(u)$. To this aim, we introduce the following sequence of Radon measures on $\Omega$

$$
\lambda_{k}:=\frac{1}{\varepsilon_{k}} f\left(\varepsilon_{k} \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right)\right) \mathcal{L}^{2}\llcorner\Omega
$$

Since the sequence $\left\{\lambda_{k}\right\}_{k \in \mathbb{N}}$ is uniformly bounded in $\mathcal{M}(\Omega)$, up to a subsequence (not relabeled), we have $\lambda_{k} \stackrel{*}{\rightharpoonup} \lambda$ weakly* in $\mathcal{M}(\Omega)$ for some nonnegative measure $\lambda \in \mathcal{M}(\Omega)$. Thanks to the lower semicontinuity of weak* convergence in $\mathcal{M}(\Omega)$ along open sets, we have that

$$
\begin{equation*}
\mathcal{F}^{\prime}(u)+\zeta \geq \lim _{k \rightarrow \infty} \lambda_{k}(\Omega) \geq \lambda(\Omega) \tag{3.5}
\end{equation*}
$$

Using that the measures $\mathcal{L}^{2}\left\llcorner\Omega\right.$ and $\mathcal{H}^{1}\left\llcorner J_{u}\right.$ are mutually singular, it is enough to show that

$$
\begin{equation*}
\frac{d \lambda}{d \mathcal{L}^{2}} \geq \mathbf{A} e(u): e(u) \quad \mathcal{L}^{2} \text {-a.e. in } \Omega \tag{3.6}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.} \geq \kappa \sin \theta_{0} \quad \mathcal{H}^{1} \text {-a.e. on } J_{u} \tag{3.7}
\end{equation*}
$$

Indeed, once (3.6) and (3.7) are satisfied, it follows from the Radon-Nikodým decomposition and the Besicovitch differentiation Theorems that

$$
\lambda=\frac{d \lambda}{d \mathcal{L}^{2}} \mathcal{L}^{2}\left\llcorner\Omega+\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.} \mathcal{H}^{1}\left\llcorner J_{u}+\lambda^{s}\right.\right.
$$

for some nonnegative measure $\lambda^{s}$ which is singular with respect to both $\mathcal{L}^{2}\left\llcorner\Omega\right.$ and $\mathcal{H}^{1}\left\llcorner J_{u}\right.$. Thus, after integration over $\Omega$ and recalling (3.5), we get that

$$
\mathcal{F}^{\prime}(u)+\zeta \geq \int_{\Omega} \mathbf{A} e(u): e(u) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{u}\right)=\mathcal{F}(u)
$$

Taking the limit as $\zeta \rightarrow 0$, we obtain the desired lower bound inequality.
The rest of this section is devoted to the establishment of (3.6) and (3.7). We start by identifying the lower bound for the bulk energy.

Proposition 3.4 (Lower bound for the Lebesgue part). For $\mathcal{L}^{2}$-a.e. $x_{0} \in \Omega$,

$$
\frac{d \lambda}{d \mathcal{L}^{2}}\left(x_{0}\right) \geq \mathbf{A} e(u)\left(x_{0}\right): e(u)\left(x_{0}\right)
$$

Proof. Let $x_{0} \in \Omega$ be such that

$$
\frac{d \lambda}{d \mathcal{L}^{2}}\left(x_{0}\right)=\lim _{\varrho \searrow 0} \frac{\lambda\left(\overline{B_{\varrho}\left(x_{0}\right)}\right)}{\pi \varrho^{2}}
$$

exists and is finite, and

$$
\lim _{\varrho \searrow 0} \frac{1}{\varrho^{2}} \int_{B_{\varrho}\left(x_{0}\right)}\left|e(u)(y)-e(u)\left(x_{0}\right)\right|^{2} d y=0 .
$$

According to Besicovitch and Lebesgue differentiation Theorems, $\mathcal{L}^{2}$-almost every point $x_{0}$ in $\Omega$ satisfies these properties. We next consider a sequence of radii $\left\{\varrho_{j}\right\}_{j \in \mathbb{N}}$ such that $\varrho_{j} \searrow 0$ and $\lambda\left(\partial B_{\varrho_{j}}\left(x_{0}\right)\right)=0$ for all $j \in \mathbb{N}$.

As in the proof of Proposition 3.1, according to the properties (1.4) satisfied by $f$, for all $\delta>0$, there exists a constant $0<K<\kappa$ such that $f(t) \geq K \wedge[(1-\delta) t]$ for all $t \geq 0$. Moreover, using the characteristic functions

$$
\chi_{k}:=\mathbb{1}_{\left\{(1-\delta) \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) \geq \frac{K}{\varepsilon_{k}}\right\}} \in L^{\infty}(\Omega ;\{0,1\})
$$

we have for every Borel set $B \subset \Omega$,

$$
\lambda_{k}(B) \geq(1-\delta) \int_{B}\left(1-\chi_{k}\right) \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) d x+\frac{K}{\varepsilon_{k}} \int_{B} \chi_{k} d x
$$

Note that because $u_{k}$ is affine on each triangle $T \in \mathbf{T}^{k}, \chi_{k}$ is constant on each triangle $T \in \mathbf{T}^{k}$. Following the proof of Proposition 3.1, the sequence $v_{k}:=\left(1-\chi_{k}\right) u_{k} \in S B V^{2}\left(\Omega ; \mathbb{R}^{2}\right)$ satisfies $v_{k} \rightarrow u$ in measure in $\Omega$ and $e\left(v_{k}\right) \rightharpoonup e(u)$ weakly in $L^{2}\left(\Omega ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right)$. Then, for all $j \in \mathbb{N}$,

$$
\begin{aligned}
\lambda\left(B_{\varrho_{j}}\left(x_{0}\right)\right)=\lim _{k \rightarrow \infty} \lambda_{k}\left(B_{\varrho_{j}}\left(x_{0}\right)\right) & \geq(1-\delta) \liminf _{k \rightarrow \infty} \int_{B_{\varrho_{j}}\left(x_{0}\right)} \mathbf{A} e\left(v_{k}\right): e\left(v_{k}\right) d x \\
& \geq(1-\delta) \int_{B_{\varrho_{j}}\left(x_{0}\right)} \mathbf{A} e(u): e(u) d x
\end{aligned}
$$

Dividing the previous inequality by $\pi \varrho_{j}^{2}$ and passing to the limit as $j \rightarrow \infty$ implies by the choice of the point $x_{0}$ that

$$
\begin{aligned}
\frac{d \lambda}{d \mathcal{L}^{2}}\left(x_{0}\right)=\lim _{j \rightarrow \infty} \frac{\lambda\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{\pi \varrho_{j}^{2}} & \geq(1-\delta) \lim _{j \rightarrow \infty} \frac{1}{\pi \varrho_{j}^{2}} \int_{B_{\varrho_{j}}\left(x_{0}\right)} \mathbf{A} e(u): e(u) d x \\
& =(1-\delta) \mathbf{A} e(u)\left(x_{0}\right): e(u)\left(x_{0}\right)
\end{aligned}
$$

Taking the limit as $\delta \rightarrow 0^{+}$completes the proof of the lower bound for the Lebesgue part.
We next pass to the lower bound inequality for the jump part of the energy which represents the most difficult and original part of our result.
Proposition 3.5 (Lower bound for the jump part). For $\mathcal{H}^{1}$-a.e. $x_{0} \in J_{u}$,

$$
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \geq \kappa \sin \theta_{0}
$$

The proof of Proposition 3.5 is quite long and involved. It necessitates the introduction of some tools in order to carry out the blow-up analysis coupled with the slicing method.

Let $x_{0} \in J_{u}$ be such that

$$
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right)=\lim _{\varrho \searrow 0} \frac{\lambda\left(\overline{B_{\varrho}\left(x_{0}\right)}\right)}{\mathcal{H}^{1}\left(J_{u} \cap \overline{B_{\varrho}\left(x_{0}\right)}\right)}
$$

exists and is finite, and

$$
\lim _{\varrho \searrow 0} \frac{\mathcal{H}^{1}\left(J_{u} \cap B_{\varrho}\left(x_{0}\right)\right)}{2 \varrho}=1
$$

According to the Besicovitch differentiation Theorem and the countably ( $\mathcal{H}^{1}, 1$ )-rectifiability of $J_{u}$ (see [3, Theorem 2.83]), it follows that $\mathcal{H}^{1}$-almost every point $x_{0}$ in $J_{u}$ fulfills these conditions.

By definition of the jump set $J_{u}$, there exist $\nu:=\nu_{u}\left(x_{0}\right) \in \mathbb{S}^{1}$ and $u^{ \pm}\left(x_{0}\right) \in \mathbb{R}^{2}$ with $u^{+}\left(x_{0}\right) \neq$ $u^{-}\left(x_{0}\right)$ such that the function $u_{x_{0}, \varrho}:=u\left(x_{0}+\varrho \cdot\right)$ converges in measure in $B:=B_{1}(0)$ to the jump function

$$
\bar{u}: y \in B \mapsto \begin{cases}u^{+}\left(x_{0}\right) & \text { if } y \cdot \nu>0 \\ u^{-}\left(x_{0}\right) & \text { if } y \cdot \nu<0\end{cases}
$$

as $\varrho \searrow 0$ (see [26, Definition 2.3]). Note that, the jump set $J_{\bar{u}}$ in $B$ coincides with the diameter $B^{\nu}=p_{\nu}(B)$ orthogonal to $\nu$. Moreover, since $\mathcal{H}^{1}\left(\left\{\xi \in \mathbb{S}^{1}:[u]\left(x_{0}\right) \cdot \xi=0\right\}\right)=0$, for any $\eta>0$, there exists $\xi \in \mathbb{S}^{1}$ such that

$$
\begin{equation*}
|\nu-\xi| \leq \eta, \quad \nu \cdot \xi \geq \frac{1}{2}, \quad\left|\nu \cdot \xi^{\perp}\right| \leq \eta \quad \text { and } \quad[u]\left(x_{0}\right) \cdot \xi \neq 0 \tag{3.8}
\end{equation*}
$$

where $[u]\left(x_{0}\right):=u^{+}\left(x_{0}\right)-u^{-}\left(x_{0}\right)$. If $[u]\left(x_{0}\right) \cdot \nu \neq 0$, we can simply take $\xi=\nu$. We then set

$$
\begin{equation*}
M_{\eta}:=\left|u^{+}\left(x_{0}\right) \cdot \xi\right|+\left|u^{-}\left(x_{0}\right) \cdot \xi\right|>0 \tag{3.9}
\end{equation*}
$$

From now on, when working with the convergence in measure, we will use the distance $d_{M_{\eta}}$ defined in (2.1) associated to this precise value of $M_{\eta}$. As before, we consider a sequence of radii $\left\{\varrho_{j}\right\}_{j \in \mathbb{N}}$ such that $\varrho_{j} \searrow 0$ and $\lambda\left(\partial B_{\varrho_{j}}\left(x_{0}\right)\right)=0$ for all $j \in \mathbb{N}$.

By our choice of $x_{0}$, we have

$$
\left\{\begin{array}{l}
\lim _{j \rightarrow \infty} \lim _{k \rightarrow \infty} u_{k}\left(x_{0}+\varrho_{j} \cdot\right)=\lim _{j \rightarrow \infty} u_{x_{0}, \varrho_{j}}=\bar{u} \quad \text { in measure in } B \\
\lim _{j \rightarrow \infty} \lim _{k \rightarrow \infty} \frac{\lambda_{k}\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{2 \varrho_{j}}=\lim _{j \rightarrow \infty} \frac{\lambda\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{2 \varrho_{j}}=\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \\
\lim _{j \rightarrow \infty} \lim _{k \rightarrow \infty} \frac{\varepsilon_{k}}{\varrho_{j}}=\lim _{j \rightarrow \infty} \lim _{k \rightarrow \infty} \frac{\omega\left(\varepsilon_{k}\right)}{\varrho_{j}}=0
\end{array}\right.
$$

The metrizability of the convergence in measure in $B$ shows the existence of an increasing sequence $\left\{k_{j}\right\}_{j \in \mathbb{N}}$ (depending on $\eta$ ) such that $k_{j} \nearrow \infty$ as $j \rightarrow \infty$ and

$$
\left\{\begin{array}{l}
v_{j}:=u_{k_{j}}\left(x_{0}+\varrho_{j} \cdot\right) \rightarrow \bar{u} \quad \text { in measure in } B  \tag{3.10a}\\
\frac{\lambda_{k_{j}}\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{2 \varrho_{j}} \rightarrow \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \\
\frac{\varepsilon_{k_{j}}}{\varrho_{j}} \rightarrow 0, \quad \frac{\omega\left(\varepsilon_{k_{j}}\right)}{\varrho_{j}} \rightarrow 0
\end{array}\right.
$$

In particular, using a change of variables, we get that

$$
\begin{aligned}
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) & =\lim _{j \rightarrow \infty} \frac{1}{\varrho_{j} \varepsilon_{k_{j}}} \int_{B_{\varrho_{j}}\left(x_{0}\right)} f\left(\varepsilon_{k_{j}} \mathbf{A} e\left(u_{k_{j}}\right): e\left(u_{k_{j}}\right)\right) d x \\
& =\lim _{j \rightarrow \infty} \frac{\varrho_{j}}{\varepsilon_{k_{j}}} \int_{B} f\left(\frac{\varepsilon_{k_{j}}}{\varrho_{j}^{2}} \mathbf{A} e\left(v_{j}\right): e\left(v_{j}\right)\right) d y \\
& \geq \limsup _{j \rightarrow \infty} \frac{\varrho_{j}}{\varepsilon_{k_{j}}} \int_{B} f\left(\frac{\varepsilon_{k_{j}}}{\varrho_{j}^{2}} \alpha\left|e\left(v_{j}\right) \xi \cdot \xi\right|^{2}\right) d y
\end{aligned}
$$

where, in the last inequality, we used the ellipticity property (1.5) of $\mathbf{A}$, the nondecreasing character of $f$ and that $\xi \in \mathbb{S}^{1}$.

According to the properties (1.4) satisfied by $f$, for all $\delta \in(0,1)$, there exists a constant $A>0$ such that

$$
f(t) \geq(A t) \wedge[(1-\delta) \kappa] \quad \text { for all } t \geq 0
$$

Indeed, since $f(t) \rightarrow \kappa$ as $t \rightarrow \infty$, there exists $t^{*} \geq 0$ such that for all $t \geq t^{*}, f(t) \geq(1-\delta) \kappa$. The function $f(t) / t$ being continuous over $\left[0, t^{*}\right]$ (extended by the value 1 at $t=0$ ), it reaches its minimum value $A>0$ over this segment so that $f(t) \geq A t$ for all $t \in\left[0, t^{*}\right]$.

Let us introduce the characteristic functions

$$
\chi_{j}:=\mathbb{1}_{\left\{\frac{A \alpha \varepsilon_{k_{j}}}{\varrho_{j}^{2}}\left|e\left(v_{j}\right) \xi \cdot \xi\right|^{2} \geq(1-\delta) \kappa\right\}} \in L^{\infty}(B ;\{0,1\}),
$$

so that

$$
\begin{equation*}
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \geq \limsup _{j \rightarrow \infty}\left\{\frac{\alpha A}{\varrho_{j}} \int_{B}\left(1-\chi_{j}\right)\left|e\left(v_{j}\right) \xi \cdot \xi\right|^{2} d y+\frac{(1-\delta) \kappa \varrho_{j}}{\varepsilon_{k_{j}}} \int_{B} \chi_{j} d y\right\} \tag{3.11}
\end{equation*}
$$

We then introduce the translated and rescaled triangulations

$$
\begin{equation*}
\mathbf{T}^{x_{0}, j}:=\frac{1}{\varrho_{j}}\left(\mathbf{T}^{k_{j}}-x_{0}\right), \quad \mathbf{T}_{b}^{x_{0}, j}:=\left\{T \in \mathbf{T}^{x_{0}, j}: \frac{\alpha A}{\varrho_{j}}\left|e\left(v_{j}\right)_{\mid T} \xi \cdot \xi\right|^{2} \geq \frac{(1-\delta) \kappa \varrho_{j}}{\varepsilon_{k_{j}}}\right\} \tag{3.12}
\end{equation*}
$$

Note that $v_{j}$ is affine on each $T \in \mathbf{T}^{x_{0}, j}$. Let us point out that

$$
\chi_{j \mid T}:= \begin{cases}1 & \text { if } T \in \mathbf{T}_{b}^{x_{0}, j}  \tag{3.13}\\ 0 & \text { otherwise }\end{cases}
$$

Since $\left(\left(v_{j}\right)_{z}^{\xi}\right)^{\prime}(t)=e\left(v_{j}\right)(z+t \xi) \xi \cdot \xi$, then for $\mathcal{H}^{1}$-a.e. $z \in B^{\xi}$,

$$
\left(\chi_{j}\right)_{z}^{\xi}(t)=\left\{\begin{array}{ll}
1 & \text { if } \frac{\alpha A}{\varrho_{j}}\left|\left(\left(v_{j}\right)_{z}^{\xi}\right)^{\prime}(t)\right|^{2} \geq \frac{(1-\delta) \kappa \varrho_{j}}{\varepsilon_{k_{j}}},  \tag{3.14}\\
0 & \text { otherwise }
\end{array} \quad \text { for } \mathcal{L}^{1} \text {-a.e. } t \in B_{z}^{\xi}\right.
$$

The triangles belonging to the collection $\mathbf{T}_{b}^{x_{0}, j}$ correspond to the sets where the longitudinal slope of $v_{j}$ in the direction $\xi$ is "very large". They, roughly speaking, represent the places where it will be energetically convenient to introduce a jump because of the sharp transition.

The following result, which will play a major role in the proof of Proposition 3.5, shows that for many points $y \in J_{\bar{u}} \cap B$, the one-dimensional energy on $B_{y}^{\xi}$ is arbitrarily small uniformly with respect to $y$.
Lemma 3.6. For all $\eta>0$, there exist a subset $Z \subset J_{\bar{u}} \cap B$ with $\mathcal{H}^{1}(Z) \leq \eta$ and a subsequence (not relabeled) such that the following property holds : for all $\gamma>0$, there exists $j_{0}=j_{0}(\gamma) \in \mathbb{N}$ such that for all $y \in J_{\bar{u}} \cap B \backslash Z$ and all $j \geq j_{0}$,

$$
\left\{\begin{array}{l}
\int_{B_{y}^{\xi}}\left(\chi_{j}\right)_{y}^{\xi} d t \leq \gamma  \tag{3.15a}\\
\int_{B_{y}^{\xi}}\left(1-\left(\chi_{j}\right)_{y}^{\xi}\right)\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \gamma^{2} \\
\int_{B_{y}^{\xi}} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{y}^{\xi}\right| d t \leq \gamma
\end{array}\right.
$$

Proof. According to Fubini's Theorem, the convergence in measure (3.10a) and (3.11), we infer that

$$
\begin{aligned}
\int_{B^{\xi}}\left(\int_{B_{z}^{\xi}} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{z}^{\xi}\right|\right. & \left.d t+\int_{B_{z}^{\xi}}\left(1-\left(\chi_{j}\right)_{z}^{\xi}\right)\left|\left(\left(v_{j}\right)_{z}^{\xi}\right)^{\prime}\right|^{2} d t+\int_{B_{z}^{\xi}}\left(\chi_{j}\right)_{z}^{\xi} d t\right) d \mathcal{H}^{1}(z) \\
\leq & \int_{B} M_{\eta} \wedge\left|v_{j}-\bar{u}\right| d x+\int_{B}\left(1-\chi_{j}\right)\left|e\left(v_{j}\right) \xi \cdot \xi\right|^{2} d x+\int_{B} \chi_{j} d x \rightarrow 0
\end{aligned}
$$

As a consequence, up to a subsequence (not relabeled), there exists an $\mathcal{H}^{1}$-negligible set $N \subset B^{\xi}$ such that

$$
\int_{B_{z}^{\xi}} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{z}^{\xi}\right| d t+\int_{B_{z}^{\xi}}\left(1-\left(\chi_{j}\right)_{z}^{\xi}\right)\left|\left(\left(v_{j}\right)_{z}^{\xi}\right)^{\prime}\right|^{2} d t+\int_{B_{z}^{\xi}}\left(\chi_{j}\right)_{z}^{\xi} d t \rightarrow 0 \quad \text { for all } z \in B^{\xi} \backslash N
$$

In order to pass from arbitrary points $z \in B^{\xi}$ to arbitrary points $y \in J_{\bar{u}} \cap B=B^{\nu}$, let us consider the following mapping (see Figure 1)

$$
\begin{equation*}
\Phi: z \in \mathbb{R}^{2} \longmapsto z-\frac{\nu \cdot z}{\nu \cdot \xi} \xi \in \Pi_{\nu} \tag{3.16}
\end{equation*}
$$

which corresponds to the linear projection onto $\Pi_{\nu}$ in the direction $\xi$. Thanks to (3.8), we can check that the Lipschitz constant of $\Phi$ is bounded by $\sqrt{1+4 \eta^{2}}$. Moreover, since for all $z \in B^{\xi}$ we have $B_{z}^{\xi}+\frac{\nu \cdot z}{\nu \cdot \xi}=B_{\Phi(z)}^{\xi}$, we deduce that

$$
\begin{aligned}
\int_{B_{z}^{\xi}} & M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{z}^{\xi}\right| d t+\int_{B_{z}^{\xi}}\left(1-\left(\chi_{j}\right)_{z}^{\xi}\right)\left|\left(\left(v_{j}\right)_{z}^{\xi}\right)^{\prime}\right|^{2} d t+\int_{B_{z}^{\xi}}\left(\chi_{j}\right)_{z}^{\xi} d t \\
& =\int_{B_{\Phi(z)}^{\xi}} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{\Phi(z)}^{\xi}\right| d s+\int_{B_{\Phi(z)}^{\xi}}\left(1-\left(\chi_{j}\right)_{\Phi(z)}^{\xi}\right)\left|\left(\left(v_{j}\right)_{\Phi(z)}^{\xi}\right)^{\prime}\right|^{2} d s+\int_{B_{\Phi(z)}^{\xi}}\left(\chi_{j}\right)_{\Phi(z)}^{\xi} d s
\end{aligned}
$$

thanks to the change of variables $s=t+\frac{\nu \cdot z}{\nu \cdot \xi}$. Since $B^{\nu} \subset \Phi\left(B^{\xi}\right)$, setting $N^{\prime}:=\Phi(N) \subset \Pi_{\nu}$, we get that $\mathcal{H}^{1}\left(N^{\prime}\right)=0$ and

$$
\int_{B_{y}^{\xi}} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{y}^{\xi}\right| d s+\int_{B_{y}^{\xi}}\left(1-\left(\chi_{j}\right)_{y}^{\xi}\right)\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d s+\int_{B_{y}^{\xi}}\left(\chi_{j}\right)_{y}^{\xi} d s \rightarrow 0 \quad \text { for all } y \in B^{\nu} \backslash N^{\prime}
$$

Applying Egoroff's theorem, for all $\eta>0$, there exists a subset $Z \subset B^{\nu}$ such that $\mathcal{H}^{1}(Z) \leq \eta$ and the above convergence is uniform with respect to $y \in B^{\nu} \backslash Z$.

For all $y \in\left(B_{1-\frac{\eta}{4}}\right)^{\nu}=J_{\bar{u}} \cap B_{1-\frac{\eta}{4}}$, we define the end points of the section passing through $y$ in the direction $\xi$ (see the Figure 1) :

$$
\left\{\begin{array}{l}
a(y):=\min \left\{t \in[-2,2]: y+t \xi \in \overline{B_{1-\frac{\eta}{4}}}\right\} \in[-2,0],  \tag{3.17}\\
b(y):=\max \left\{t \in[-2,2]: y+t \xi \in \overline{B_{1-\frac{\eta}{4}}}\right\} \in[0,2],
\end{array}\right.
$$

so that $\left(B_{1-\frac{\eta}{4}}\right)_{y}^{\xi}=(a(y), b(y))$. Note that, for all $y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \subset\left(B_{1-\frac{\eta}{4}}\right)^{\nu}$,

$$
\begin{equation*}
0<L_{\eta}:=\sqrt{\left(1-\frac{\eta}{2}\right)^{2}\left|\xi \cdot \nu^{\perp}\right|^{2}+\frac{\eta(8-3 \eta)}{16}}-\left(1-\frac{\eta}{2}\right)\left|\xi \cdot \nu^{\perp}\right| \leq|a(y)|,|b(y)| \leq 2 \tag{3.18}
\end{equation*}
$$



Figure 1
We introduce the family

$$
\mathbf{T}_{b, \text { int }}^{x_{0}, j}:=\left\{T \in \mathbf{T}_{b}^{x_{0}, j}: T \cap \overline{B_{1-\frac{\eta}{4}}} \neq \emptyset\right\}
$$

of triangles which intersect $\overline{B_{1-\frac{\eta}{4}}}$ and where $v_{j}$ varies enough in the direction $\xi$. Note that for $j \in \mathbb{N}$ large enough (depending on $\eta$ ), each $T \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ is contained in $B$, since the lengths of all triangles's edges are controlled by $\omega\left(\varepsilon_{k_{j}}\right) / \varrho_{j} \rightarrow 0$. The collection $\mathbf{T}_{b, i n t}^{x_{0}, j}$ is introduced for technical reasons to deal with triangles which intersect the boundary of the ball $B$.

In the following result, we show that, for some subset of $Z^{\prime} \subset J_{\bar{u}} \cap B_{1-\frac{\eta}{2}}$ of arbitrarily small $\mathcal{H}^{1}$ measure, and along a subsequence (only depending on $\eta$ ), all the sections in the direction $\xi$ passing through $J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime}$ must cross at least one triangle $T \in \mathbf{T}_{b, \text { int }}^{x_{0}, j}$ contained in $B$, and on which the longitudinal slope of $v_{j}$ in the direction $\xi$ is "large". The formal idea of the proof consists in observing that, if for some $y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}}$ the one-dimensional section $B_{y}^{\xi}$ intersects no triangle in the collection $\mathbf{T}_{b, i n t}^{x_{0}, j}$ for infinitely many $j^{\prime}$ s, then the function $\left(v_{j}\right)_{y}^{\xi}$ would be bounded in $H^{1}\left(B_{y}^{\xi}\right)$. Lemma 3.6 would then entail that $\left(v_{j}\right)_{y}^{\xi}$ converges (weakly in $H^{1}\left(B_{y}^{\xi}\right)$ and also $\mathcal{L}^{1}$-a.e. in $\left.B_{y}^{\xi}\right)$ to a constant function. This property contradicts the fact that $\left(v_{j}\right)_{y}^{\xi} \rightarrow \bar{u}_{y}^{\xi} \mathcal{L}^{1}$-a.e. in $B_{y}^{\xi}$, where $\bar{u}_{y}^{\xi}$ is a step function taking two different values $u^{ \pm}\left(x_{0}\right) \cdot \xi$.

Lemma 3.7. For all $\eta>0$, there exist a subset $Z^{\prime} \subset J_{\bar{u}} \cap B$ containing $Z$ with $\mathcal{H}^{1}\left(Z^{\prime}\right) \leq \eta$, and a subsequence (not relabeled) such that the following property holds : for all $y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime}$ and all $j \in \mathbb{N}$, there exists a triangle $T=T(y, j) \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ such that $(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset$.

Proof. Let $Z$ be the exceptional set given by Lemma 3.6. We first show the weaker result that there exists an increasing mapping $\phi: \mathbb{N} \rightarrow \mathbb{N}$ with the following property: for all $y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z$ and all $j \in \mathbb{N}$, there exists a triangle $T=T(y, \phi(j)) \in \mathbf{T}_{b, i n t}^{x_{0}, \phi(j)}$ such that $(T \cap B)_{y}^{\xi} \neq \emptyset$.

Suppose by contradiction that such is not the case, and define

$$
\gamma_{1}^{*}:=L_{\eta} M_{\eta}>0, \quad \gamma_{2}^{*}:=\frac{L_{\eta}\left|[u]\left(x_{0}\right) \cdot \xi\right|}{1+2 L_{\eta}}>0 \quad \text { and } \quad \gamma^{*}=\gamma^{*}(\eta):=\frac{\gamma_{1}^{*} \wedge \gamma_{2}^{*}}{4}>0
$$

where we recall that the constants $M_{\eta}$ and $L_{\eta}$ are defined in (3.9) and (3.18), respectively. Thanks to Lemma 3.6, there exists a rank $j^{*}=j^{*}\left(\gamma^{*}\right) \in \mathbb{N}$ such that for all $y \in J_{\bar{u}} \cap B \backslash Z$ and all $j \geq j^{*}$,

$$
\int_{B_{y}^{\xi}}\left(1-\left(\chi_{j}\right)_{y}^{\xi}\right)\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \gamma^{* 2} \text { and } \int_{B_{y}^{\xi}} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{y}^{\xi}\right| d t \leq \gamma^{*}
$$

We then consider the extraction $\phi: j \in \mathbb{N} \longmapsto j+j^{*} \in \mathbb{N}$ which only depends on $\eta$. By assumption, there exists $y=y(\phi) \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z$ and $j=j(\phi) \in \mathbb{N}$ such that $(T \cap B)_{y}^{\xi}=\emptyset$ for all $T \in \mathbf{T}_{b, i n t}^{x_{0}, j+j^{*}}$. Remembering (3.13), we deduce that $\left(\chi_{j+j^{*}}\right)_{y}^{\xi} \equiv 0$ on $(a(y), b(y))$. Moreover, since $\phi(j)=j+j^{*} \geq j^{*}$, we have

$$
\int_{a(y)}^{b(y)}\left|\left(\left(v_{j+j^{*}}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \gamma^{* 2} \text { and } \int_{a(y)}^{b(y)} M_{\eta} \wedge\left|\left(v_{j+j^{*}}-\bar{u}\right)_{y}^{\xi}\right| d t \leq \gamma^{*}
$$



Figure 2

By continuity of $\left(v_{j+j^{*}}\right)_{y}^{\xi}$ on the compact $[a(y), b(y)],\left(v_{j+j^{*}}\right)_{y}^{\xi}$ being in $H^{1}\left(B_{y}^{\xi}\right)$, there exist two points $t^{ \pm} \in[a(y), b(y)] \cap \mathbb{R}^{ \pm}$such that

$$
\min _{[a(y), b(y)] \cap \mathbb{R}^{ \pm}}\left(M_{\eta} \wedge\left|\left(v_{j+j^{*}}\right)_{y}^{\xi}-u^{ \pm}\left(x_{0}\right) \cdot \xi\right|\right)=M_{\eta} \wedge\left|\left(v_{j+j^{*}}\right)_{y}^{\xi}\left(t^{ \pm}\right)-u^{ \pm}\left(x_{0}\right) \cdot \xi\right| .
$$

Hence,

$$
\begin{aligned}
\frac{\gamma^{*}}{L_{\eta}} & \geq \frac{1}{L_{\eta}} \int_{a(y)}^{0} M_{\eta} \wedge\left|\left(v_{j+j^{*}}\right)_{y}^{\xi}-u^{-}\left(x_{0}\right) \cdot \xi\right| d t+\frac{1}{L_{\eta}} \int_{0}^{b(y)} M_{\eta} \wedge\left|\left(v_{j+j^{*}}\right)_{y}^{\xi}-u^{+}\left(x_{0}\right) \cdot \xi\right| d t \\
& \geq M_{\eta} \wedge\left|\left(v_{j+j^{*}}\right)_{y}^{\xi}\left(t^{-}\right)-u^{-}\left(x_{0}\right) \cdot \xi\right|+M_{\eta} \wedge\left|\left(v_{j+j^{*}}\right)_{y}^{\xi}\left(t^{+}\right)-u^{+}\left(x_{0}\right) \cdot \xi\right| \\
& \geq M_{\eta} \wedge\left(\left|\left(v_{j+j^{*}}\right)_{y}^{\xi}\left(t^{-}\right)-u^{-}\left(x_{0}\right) \cdot \xi\right|+\left|\left(v_{j+j^{*}}\right)_{y}^{\xi}\left(t^{+}\right)-u^{+}\left(x_{0}\right) \cdot \xi\right|\right) \\
& \geq M_{\eta} \wedge\left(\left|[u]\left(x_{0}\right) \cdot \xi\right|-\left|\int_{t^{-}}^{t^{+}}\left(\left(v_{j+j^{*}}\right)_{y}^{\xi}\right)^{\prime}(t) d t\right|\right) \\
& \geq M_{\eta} \wedge\left(\left|[u]\left(x_{0}\right) \cdot \xi\right|-2 \gamma^{*}\right)
\end{aligned}
$$

which is impossible thanks to of our choice of $\gamma^{*}$.
We are now in position to complete the proof of Lemma 3.7. For all $j \in \mathbb{N}$, let

$$
\begin{array}{r}
Z_{j}:=\left\{y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}}: \text { there exists } T \in \mathbf{T}_{b, i n t}^{x_{0}, j}\right. \text { such that } \\
\left.(T \cap B)_{y}^{\xi} \text { is contained in an edge or a vertex of } T\right\}
\end{array}
$$

and

$$
Z^{\prime}:=Z \cup \bigcup_{j \in \mathbb{N}} Z_{j}
$$

We notice that $\bigcup_{j} Z_{j}$ is $\mathcal{H}^{1}$-negligible (each $Z_{j}$ being finite), hence $\mathcal{H}^{1}\left(Z^{\prime}\right) \leq \eta$. Moreover, for all $y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime}$ and all $j \in \mathbb{N}$, there exists a triangle $T \in \mathbf{T}_{b, i n t}^{x_{0}, \phi(j)}$ such that $(T \cap B)_{y}^{\xi}$ is non-empty, and it is neither reduced to a vertex of $T$ nor contained in an edge of it. It thus implies that $(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset$.

As a consequence of Lemma 3.7, for all $j \in \mathbb{N}$, the family of triangles

$$
\begin{equation*}
\mathscr{F}_{j}:=\left\{T \in \mathbf{T}_{b, i n t}^{x_{0}, j}: \text { there exists } y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \text { such that }(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset\right\} \tag{3.19}
\end{equation*}
$$

is nonempty. Thanks to Lemma 3.7, it is possible to obtain a bad lower bound. Indeed, from that result, we infer that $J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime} \subset \bigcup_{T \in \mathscr{F}}^{j}$ $\Phi\left(p_{\xi}(T)\right)$ with $\Phi$ the projection onto $\Pi_{\nu}$ in the direction $\xi$ defined in (3.16). Using next that $\mathcal{L}^{2}(T) \geq \mathcal{H}^{1}\left(p_{\xi}(T)\right)\left(\varepsilon_{k_{j}} / \varrho_{j}\right) \sin \theta_{0} / 2$ and that the Lipschitz constant of $\Phi$ is bounded by $\sqrt{1+4 \eta^{2}}$, we deduce from (3.11) and our choice of $x_{0}$ that

$$
\begin{aligned}
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) & \geq \liminf _{j \rightarrow \infty} \frac{(1-\delta) \kappa \varrho_{j}}{\varepsilon_{k_{j}}} \int_{B} \chi_{j} d y \geq \liminf _{j \rightarrow \infty} \sum_{T \in \mathscr{F}_{j}} \frac{(1-\delta) \kappa \varrho_{j} \mathcal{L}^{2}(T)}{\varepsilon_{k_{j}}} \\
& \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{2 \sqrt{1+4 \eta^{2}}} \liminf _{j \rightarrow \infty} \mathcal{H}^{1}\left(\bigcup_{T \in \mathscr{F}_{j}} \Phi\left(p_{\xi}(T)\right)\right) \\
& \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{2 \sqrt{1+4 \eta^{2}}} \mathcal{H}^{1}\left(J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime}\right) \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{\sqrt{1+4 \eta^{2}}}(1-\eta)
\end{aligned}
$$

Letting $\eta \rightarrow 0$ and $\delta \rightarrow 0$ leads to

$$
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \geq \frac{\kappa \sin \theta_{0}}{2}
$$

which corresponds to a too low lower bound because of the factor $1 / 2$ in the right-hand side of the previous inequality. In order to improve the previous argument, we need to establish that many lines $B_{y}^{\xi}$ parallel to $\xi$ and passing through the jump set at some point $y \in J_{\bar{u}} \cap B$ must actually intersect at least two triangles of the collection $\mathbf{T}_{b, i n t}^{x_{0}, j}$, where the longitudinal variation of $v_{j}$ in the direction $\xi$ is "large". This idea is precisely formulated in the following result which is an improvement of Lemma 3.7.

Lemma 3.8. For all $\eta>0$, there exist $Z^{\prime \prime} \subset J_{\bar{u}} \cap B$ containing $Z^{\prime}$ with $\mathcal{H}^{1}\left(Z^{\prime \prime}\right) \leq 3 \eta$, and a (not relabeled) subsequence such that for all $j \in \mathbb{N}$ and for all $y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}$,

$$
\#\left\{T \in \mathbf{T}_{b, i n t}^{x_{0}, j}:(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset\right\} \geq 2
$$

The proof of Lemma 3.8 consists in constructing both $Z^{\prime \prime}$ and the subsequence inductively by means of the following technical result, Lemma 3.9. It stipulates that the set of all points $y \in J_{\bar{u}} \cap B$ such that $B_{y}^{\xi}$ intersects exactly one triangle $T$ in the collection $\mathbf{T}_{b, i n t}^{x_{0}, j}$, has arbitrarily small $\mathcal{H}^{1}$ measure. To establish this property, we first show that if such situation arises, then the function $\left(v_{j}\right)_{y}^{\xi}$ is uniformly close (with respect to $y$ ) to the step function $\bar{u}_{y}^{\xi}$ taking the values $u^{ \pm}\left(x_{0}\right) \cdot \xi$. Thus, up to a small error which is uniform in $y$, the function $\left(v_{j}\right)_{y}^{\xi}$ must pass from the value $u^{-}\left(x_{0}\right) \cdot \xi$ to $u^{+}\left(x_{0}\right) \cdot \xi$ in an affine way inside the only triangle $T \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ which is crossed by $B_{y}^{\xi}$. However, due to the shape of a triangle, this can happen for at most two different values of $y$, say $z_{1}$ and $z_{2}$. Then, if $y \in J_{\bar{u}} \cap B$ is far away from these two values $z_{1}$ and $z_{2}$, the variation of $\left(v_{j}\right)_{y}^{\xi}$ is not sufficient to connect the values $u^{ \pm}\left(x_{0}\right) \cdot \xi$ in an affine way. It thus becomes necessary for $B_{y}^{\xi}$ to intersect an additional triangle $T^{\prime} \in \mathbf{T}_{b, i n t}^{x_{0}, j}$, where the variation of $\left(v_{j}\right)_{y}^{\xi}$ is substantial, in order to recover the full jump.
Lemma 3.9. For all $\eta>0$, there exist constants $C_{*}=C_{*}(\eta)>0, \gamma_{*}=\gamma_{*}(\eta)>0$ and a subset $Z_{*}=Z_{*}(\eta) \subset J_{\bar{u}} \cap B$ containing $Z^{\prime}$ and satisfying $\mathcal{H}^{1}\left(Z_{*}\right) \leq 2 \eta$ such that the following property holds: for all $0<\gamma<\gamma_{*}$, there exists $j(\gamma) \in \mathbb{N}$ such that for all $j \geq j(\gamma)$, the set

$$
\begin{equation*}
Y_{j}:=\left\{y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime}: \text { there exists a unique } T \in \mathbf{T}_{b, \text { int }}^{x_{0}, j} \text { such that }(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset\right\} \tag{3.20}
\end{equation*}
$$

satisfies

$$
\mathcal{H}^{1}\left(Y_{j} \backslash Z_{*}\right) \leq C_{*} \gamma
$$

Proof of Lemma 3.9. The proof is divided into three steps.
Step 1. In this first step, we show that for $j$ large enough and for many points $y \in Y_{j}$, the set $(B \cap T)_{y}^{\xi}$ (where $T$ is the only triangle in $\mathbf{T}_{b, i n t}^{x_{0}, j}$ which crosses $B_{y}^{\xi}$ ) is close to $\left(J_{\bar{u}}\right)_{y}^{\xi}$, uniformly with respect to $y$.

For all $j \in \mathbb{N}$ and all $y \in Y_{j}$, let $T_{j}(y) \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ be the unique triangle such that $\left(\stackrel{\circ}{T}_{j}(y) \cap B\right)_{y}^{\xi} \neq \emptyset$. We define the end points of the section in the direction $\xi$ passing through $y$ inside $T_{j}(y)$ (see the Figure 3) by

$$
\left\{\begin{array}{l}
a_{j}(y):=\min \left\{t \in[-2,2]: y+t \xi \in T_{j}(y)\right\}  \tag{3.21}\\
b_{j}(y):=\max \left\{t \in[-2,2]: y+t \xi \in T_{j}(y)\right\}
\end{array}\right.
$$

so that $\left(T_{j}(y)\right)_{y}^{\xi}=\left[a_{j}(y), b_{j}(y)\right]$. Note that $T_{j}(y) \subset B\left(\right.$ since $\left.T_{j}(y) \cap \overline{B_{1-\frac{\eta}{4}}} \neq \emptyset\right)$, hence $-2 \leq$ $a_{j}(y) \leq b(y)$ and $2 \geq b_{j}(y) \geq a(y)$. Let us show that

$$
\begin{equation*}
f_{j}(y):=\left(\left|a_{j}(y)\right|+\left|b_{j}(y)\right|\right) \mathbb{1}_{Y_{j}}(y) \rightarrow 0 \quad \text { for all } y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime} \tag{3.22}
\end{equation*}
$$

Let $y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime}$ and set $\ell:=\limsup _{j} f_{j}(y) \in[0,4]$. Assume by contradiction that $\ell>0$ and extract a subsequence depending on $y$ (not relabeled) such that $f_{j}(y) \rightarrow \ell$. Then, there exists


Figure 3
a rank $j_{0} \in \mathbb{N}$ such that $y \in Y_{j}$ for all $j \geq j_{0}$. Moreover, according to Lemma 3.6 and setting $I_{j}(y):=(a(y), b(y)) \backslash\left(a_{j}(y), b_{j}(y)\right) \subset B_{y}^{\xi}$, we have

$$
\left\{\begin{array}{l}
\left|b_{j}(y)-a_{j}(y)\right| \leq \int_{B_{y}^{\xi}}\left(\chi_{j}\right)_{y}^{\xi} d t \rightarrow 0  \tag{3.23}\\
\int_{I_{j}(y)}\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \int_{B_{y}^{\xi}}\left(1-\left(\chi_{j}\right)_{y}^{\xi}\right)\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \rightarrow 0 \\
\int_{a(y)}^{b(y)} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{y}^{\xi}\right| d t \leq \int_{B_{y}^{\xi}} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{y}^{\xi}\right| d t \rightarrow 0
\end{array}\right.
$$

Up to another subsequence (still not relabeled), the first condition in (3.23) ensures that $a_{j}(y) \rightarrow m$ and $b_{j}(y) \rightarrow m$ for some $m \in[a(y), b(y)]$. Thus, for all $\tau>0$, there exists a rank $j_{1}=j_{1}(\tau) \geq j_{0}$ such that for all $j \geq j_{1}$,

$$
I_{\tau}:=(a(y), m-\tau) \cup(m+\tau, b(y)) \subset I_{j}(y)
$$

with the convention that $(x, y)=\emptyset$ if $y<x$. We set

$$
I_{\tau}^{-}:=(a(y), m-\tau), \quad I_{\tau}^{+}:=(m+\tau, b(y))
$$

so that $\left(v_{j}\right)_{y \mid I_{\tau}^{ \pm}}^{\xi} \in H^{1}\left(I_{\tau}^{ \pm}\right)$and the truncated function $w_{j}:=\left(M_{\eta} \wedge\left(v_{j}\right)_{y}^{\xi}\right) \vee\left(-M_{\eta}\right) \in H^{1}\left(I_{\tau}^{ \pm}\right)$ satisfies $w_{j}^{\prime}=\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime} \mathbb{1}_{\left\{\left|\left(v_{j}\right)_{y}^{\xi}\right| \leq M_{\eta}\right\}}$. According to the second condition in (3.23), the sequence $\left\{w_{j}\right\}_{j \in \mathbb{N}}$ is bounded in $H^{1}\left(I_{\tau}^{ \pm}\right)$and $w_{j}^{\prime} \rightarrow 0$ in $L^{2}\left(I_{\tau}^{ \pm}\right)$. As a consequence, up to a subsequence, there exist constants $c^{ \pm} \in \mathbb{R}$ such that $w_{j} \rightarrow c^{ \pm}$in $H^{1}\left(I_{\tau}^{ \pm}\right)$and $\mathcal{L}^{1}$-a.e. in $I_{\tau}^{ \pm}$. Yet, as $\left(v_{j}\right)_{y}^{\xi}$ converges in measure to $\bar{u}_{y}^{\xi}$ in $I_{\tau}^{ \pm}$, up to another subsequence (still not relabeled), we have that $\left(v_{j}\right)_{y}^{\xi}$ pointwise converges to $\bar{u}_{y}^{\xi} \mathcal{L}^{1}$-a.e. in $I_{\tau}^{ \pm}$. Hence $c^{ \pm}=\left(M_{\eta} \wedge u^{ \pm}\left(x_{0}\right) \cdot \xi\right) \vee\left(-M_{\eta}\right)=u^{ \pm}\left(x_{0}\right) \cdot \xi$ by our choice (3.9) of $M_{\eta}$. Thus, for all $\tau>0$,

$$
u^{-}\left(x_{0}\right) \cdot \xi \mathbb{1}_{(a(y), m-\tau)}+u^{+}\left(x_{0}\right) \cdot \xi \mathbb{1}_{(m+\tau, b(y))}=\bar{u}_{y \mid I_{\tau}}^{\xi} \quad \mathcal{L}^{1} \text {-a.e. in } I_{\tau}
$$

Taking the limit as $\tau \rightarrow 0^{+}$, we obtain that

$$
u^{-}\left(x_{0}\right) \cdot \xi \mathbb{1}_{(a(y), m)}+u^{+}\left(x_{0}\right) \cdot \xi \mathbb{1}_{(m, b(y))}=\bar{u}_{y}^{\xi} \quad \mathcal{L}^{1} \text {-a.e. in }(a(y), b(y))
$$

leading to $m=0$ since $[u]\left(x_{0}\right) \cdot \xi \neq 0$ by our choice (3.8) of $\xi$. As a consequence $f_{j}(y)=$ $\left(\left|a_{j}(y)\right|+\left|b_{j}(y)\right|\right) \mathbb{1}_{Y_{j}}(y) \rightarrow 0$ which is against $\ell>0$.

Using (3.22), Lemma 3.6 and owing to Egoroff's Theorem, we can find a set $Z_{*} \subset J_{\bar{u}} \cap B$ containing $Z^{\prime}$ with $\mathcal{H}^{1}\left(Z_{*}\right) \leq 2 \eta$ such that for all $\gamma>0$, there exists $j_{0}(\gamma) \in \mathbb{N}$ satisfying

$$
\left\{\begin{array}{l}
\int_{B_{y}^{\xi}}\left(1-\left(\chi_{j}\right)_{y}^{\xi}\right)\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \gamma^{2},  \tag{3.24}\\
\int_{B_{y}^{\xi}}^{\xi} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{y}^{\xi}\right| d t \leq \gamma, \\
\left(\left|a_{j}(y)\right|+\left|b_{j}(y)\right|\right) \mathbb{1}_{Y_{j}}(y) \leq \gamma
\end{array} \quad \text { for all } y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z_{*} \text { and all } j \geq j_{0}(\gamma)\right.
$$

Step 2. In this step, we show that for many points $y \in Y_{j}$, the variation of $\left(v_{j}\right)_{y}^{\xi}$ inside the only triangle $T$ in $\mathbf{T}_{b, i n t}^{x_{0}, j}$ which is crossed by $B_{y}^{\xi}$, is uniformly close with respect to $y$ to the jump of $\bar{u}_{y}^{\xi}$. More precisely, let

$$
\begin{equation*}
C_{\eta}:=8\left(1+\frac{1}{L_{\eta}}\right)>0, \quad \gamma_{*}=\gamma_{*}(\eta):=\frac{1}{2} \min \left(1, \frac{M_{\eta}}{C_{\eta}}, L_{\eta}, \frac{\left|[u]\left(x_{0}\right) \cdot \xi\right|}{4 C_{\eta}}\right)>0 \tag{3.25}
\end{equation*}
$$

Let us show that for all $0<\gamma<\gamma_{*}$, there exists a rank $j_{1}(\gamma) \in \mathbb{N}$ such that

$$
\begin{equation*}
\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)-[u]\left(x_{0}\right) \cdot \xi\right| \leq C_{\eta} \gamma \quad \text { for all } j \geq j_{1}(\gamma) \text { and all } y \in Y_{j} \backslash Z_{*} \tag{3.26}
\end{equation*}
$$

Fix $0<\gamma<\gamma_{*}$ and, by (3.24), let $j_{0}(\gamma) \in \mathbb{N}$ be such that

$$
\left\{\begin{array}{l}
\int_{(a(y), b(y)) \backslash\left(a_{j}(y), b_{j}(y)\right)}\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \gamma^{2} \\
\int_{B_{y}^{\xi}} M_{\eta} \wedge\left|\left(v_{j}-\bar{u}\right)_{y}^{\xi}\right| d t \leq \gamma, \\
\left(\left|a_{j}(y)\right|+\left|b_{j}(y)\right|\right) \leq \gamma
\end{array} \quad \text { for all } j \geq j_{0}(\gamma) \text { and all } y \in Y_{j} \backslash Z_{*}\right.
$$

In particular, recalling $(3.17),(3.18)$ and by the choice (3.25) of $\gamma_{*}$, we get that $2 \geq|b(y)| \geq L_{\eta}>$ $\gamma_{*}>\gamma \geq\left|b_{j}(y)\right|$ and $2 \geq|a(y)| \geq L_{\eta}>\gamma_{*}>\gamma \geq\left|a_{j}(y)\right|$, hence

$$
a(y)<a_{j}(y) \leq b_{j}(y)<b(y)
$$

Writing

$$
\begin{aligned}
M_{\eta} \wedge\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)-[u]\left(x_{0}\right) \cdot \xi\right| \leq & M_{\eta} \wedge\left|\left(v_{j}\right)_{y}^{\xi}\left(0 \vee b_{j}(y)\right)-u^{+}\left(x_{0}\right) \cdot \xi\right| \\
& +M_{\eta} \wedge\left|\left(v_{j}\right)_{y}^{\xi}\left(0 \vee b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)\right| \\
& +M_{\eta} \wedge\left|\left(v_{j}\right)_{y}^{\xi}\left(0 \wedge a_{j}(y)\right)-u^{-}\left(x_{0}\right) \cdot \xi\right| \\
& +M_{\eta} \wedge\left|\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(0 \wedge a_{j}(y)\right)\right| \\
=: \quad & J_{1}+J_{2}+J_{3}+J_{4},
\end{aligned}
$$

it remains to control each of the last four terms.
Let us first estimate the terms $J_{2}$ and $J_{4}$. If $b_{j}(y) \geq 0, J_{2}=0$. Otherwise, by the CauchySchwarz inequality,

$$
J_{2}=M_{\eta} \wedge\left|\int_{b_{j}(y)}^{0}\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime} d t\right| \leq \sqrt{\left|b_{j}(y)\right|}\left(\int_{b_{j}(y)}^{0}\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t\right)^{\frac{1}{2}} \leq \gamma^{3 / 2} \leq \gamma
$$

Similarly, we have that $J_{4} \leq \gamma$.
Let us now estimate the term $J_{1}$. We consider the function

$$
z_{j}:=M_{\eta} \wedge\left|\left(v_{j}\right)_{y}^{\xi}-u^{+}\left(x_{0}\right) \cdot \xi\right| \in H^{1}\left(B_{y}^{\xi}\right)
$$

with

$$
z_{j}^{\prime}=\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime} \mathbb{1}_{\left\{0 \leq\left(v_{j}\right)_{y}^{\xi}-u^{+}\left(x_{0}\right) \cdot \xi \leq M_{\eta}\right\}}-\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime} \mathbb{1}_{\left\{0 \leq u^{+}\left(x_{0}\right) \cdot \xi-\left(v_{j}\right)_{y}^{\xi} \leq M_{\eta}\right\}},
$$

and the nonempty open interval $I^{+}:=\left(0 \vee b_{j}(y), b(y)\right)$. By the Sobolev embedding and (3.25), we have that for all $t \in I^{+}$,

$$
\begin{aligned}
\left|z_{j}(t)\right| & \leq \sqrt{b(y)-0 \vee b_{j}(y)}\left\|z_{j}^{\prime}\right\|_{L^{2}\left(I^{+}\right)}+\frac{1}{b(y)-0 \vee b_{j}(y)}\left\|z_{j}\right\|_{L^{1}\left(I^{+}\right)} \\
& \leq \sqrt{2}\left\|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right\|_{L^{2}\left(I^{+}\right)}+\frac{2}{L_{\eta}}\left\|M_{\eta} \wedge\left|\left(v_{j}\right)_{y}^{\xi}-u^{+}\left(x_{0}\right) \cdot \xi\right|\right\|_{L^{1}\left(I^{+}\right)} \\
& \leq\left(\sqrt{2}+\frac{2}{L_{\eta}}\right) \gamma
\end{aligned}
$$

By continuity of $\left(v_{j}\right)_{y}^{\xi}$ in $B_{y}^{\xi}$, the above inequality remains true up to the end point $0 \vee b_{j}(y)$ of $I^{+}$, so that $J_{1} \leq\left(\sqrt{2}+\frac{2}{L_{\eta}}\right) \gamma$. A similar argument shows that $J_{3} \leq\left(\sqrt{2}+\frac{2}{L_{\eta}}\right) \gamma$, and thus $J_{1}+J_{2}+J_{3}+J_{4} \leq 8\left(1+\frac{1}{L_{\eta}}\right) \gamma=C_{\eta} \gamma$, which shows that

$$
M_{\eta} \wedge\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)-[u]\left(x_{0}\right) \cdot \xi\right| \leq C_{\eta} \gamma
$$

Eventually, as $C_{\eta} \gamma<M_{\eta}$ for all $0<\gamma<\gamma^{*}$ by (3.25), we conclude the validity of (3.26).
Step 3. We now show that it is possible to include $Y_{j} \backslash Z_{*}$ inside a finite union of arbitrarily small segments contained in $J_{\bar{u}} \cap B$ (see Figure 6).

Let $0<\gamma<\gamma_{*}$ and $j_{1}(\gamma) \in \mathbb{N}$ be given by (3.26). For all $j \geq j_{1}(\gamma)$, we define

$$
\widehat{\mathbf{T}}_{j}:=\left\{T \in \mathbf{T}_{b, \text { int }}^{x_{0}, j}: \text { there exists } y \in Y_{j} \backslash Z_{*} \text { such that }(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset\right\}
$$

and, for all $T \in \widehat{\mathbf{T}}_{j}$, we introduce both following quantities :

$$
\left\{\begin{array}{l}
L^{\mathrm{ref}}(T):=\frac{\left|[u]\left(x_{0}\right) \cdot \xi\right|-4 C_{\eta} \gamma}{\left|e\left(v_{j}\right)_{\mid T}:(\xi \otimes \xi)\right|} \text { the reference length of } T  \tag{3.27}\\
L^{\max }(T):=\max _{z \in p_{\xi}(T)} \mathcal{L}^{1}\left(T_{z}^{\xi}\right) \text { the maximal section's length of } T \text { along the direction } \xi
\end{array}\right.
$$

Note that because $T \in \mathbf{T}_{b}^{x_{0}, j}$, see (3.12), then $\left|e\left(v_{j}\right)_{\mid T} \xi \cdot \xi\right|^{2} \geq(1-\delta) \kappa \varrho_{j}^{2} /\left(\alpha A \varepsilon_{k_{j}}\right)>0$, so that $L^{\text {ref }}(T)$ is well defined, and positive by (3.25) since $\gamma<\gamma_{*}$. The quantity $L^{\text {ref }}(T)$ stands for the required length of the section $T_{y}^{\xi}$ in order for the (affine) function $\left(v_{j}\right)_{y}^{\xi}$ to pass exactly from the values $u^{-}\left(x_{0}\right) \cdot \xi$ to $u^{+}\left(x_{0}\right) \cdot \xi$ across $T$, up to the error $4 C_{\eta} \gamma$. Note that $\mathcal{L}^{1}\left(T_{y}^{\xi}\right)=L^{\text {ref }}(T)$ for at most two values of $y$, say $z_{1}$ and $z_{2}$, only depending on $j$ and $T$. If $y \in Y_{j} \backslash Z_{*}$ is such that $(T \circ \cap B)_{y}^{\xi} \neq \emptyset$, we know from Step 2 that the variation of $\left(v_{j}\right)_{y}^{\xi}$ across $T$ is close to $[u]\left(x_{0}\right) \cdot \xi$, up to a small error of order $O(\gamma)$ which is uniform with respect to $y$. Therefore, we will show that if $y$ is far away from $z_{1}$ and $z_{2}$, then the variation of $\left(v_{j}\right)_{y}^{\xi}$ across $T$ is not sufficient to recover the full jump $[u]\left(x_{0}\right) \cdot \xi$.

Let $x_{1}, x_{2}$ and $x_{3} \in T$ be the three vertices of $T$ and $X_{i}:=p_{\xi}\left(x_{i}\right) \in B^{\xi}$. We easily see that there exists $i_{0} \in\{1,2,3\}$ such that $X_{i_{0}}=\underset{z \in p_{\xi}(T)}{\arg \max } \mathcal{L}^{1}\left(T_{z}^{\xi}\right)$. Up to a permutation of $\left\{x_{1}, x_{2}, x_{3}\right\}$, there is no loss of generality to assume that $i_{0}=3$ and $X_{3} \cdot \xi^{\perp} \leq X_{1} \cdot \xi^{\perp}$, with $\xi^{\perp} \in \mathbb{S}^{1}$ being one of the two orthogonal vectors to $\xi$.

Let $h_{T} \geq \frac{\varepsilon_{k_{j}} \sin \left(\theta_{0}\right)}{\varrho_{j}}>0$ be the smallest height of $T$. We claim that, for all $z, z^{\prime} \in p_{\xi}(T)$ be such that $\stackrel{\circ}{T}_{z}^{\xi} \neq \emptyset, \stackrel{\circ}{T}_{z^{\prime}}^{\xi} \neq \emptyset$ and either $z \cdot \xi^{\perp}, z^{\prime} \cdot \xi^{\perp} \geq X_{3} \cdot \xi^{\perp}$ or $z \cdot \xi^{\perp}, z^{\prime} \cdot \xi^{\perp} \leq X_{3} \cdot \xi^{\perp}$, then,

$$
\begin{equation*}
\left|z-z^{\prime}\right| \leq \frac{2 \mathcal{L}^{2}(T)}{h_{T}} \frac{\left|\mathcal{L}^{1}\left(T_{z}^{\xi}\right)-\mathcal{L}^{1}\left(T_{z^{\prime}}^{\xi}\right)\right|}{\max \left(\mathcal{L}^{1}\left(T_{z}^{\xi}\right), \mathcal{L}^{1}\left(T_{z^{\prime}}^{\xi}\right)\right)} \tag{3.28}
\end{equation*}
$$

Indeed, consider for instance the case where

$$
X_{1} \cdot \xi^{\perp} \geq z \cdot \xi^{\perp}>z^{\prime} \cdot \xi^{\perp} \geq X_{3} \cdot \xi^{\perp}
$$

(see Figure 4). Let

$$
L:=\mathcal{L}^{1}\left(T_{z}^{\xi}\right), \quad L^{\prime}:=\mathcal{L}^{1}\left(T_{z^{\prime}}^{\xi}\right), \quad d:=\left|X_{1}-z\right|, \quad d^{\prime}:=\left|X_{1}-z^{\prime}\right|
$$



Figure 4
Then, $L^{\prime}>L>0, d^{\prime}>d>0$ and using Thalès' Theorem, we have that

$$
\frac{d}{d^{\prime}}=\frac{d^{\prime}-\left|z-z^{\prime}\right|}{d^{\prime}}=\frac{L}{L^{\prime}}
$$

Since $d^{\prime} \leq\left|X_{1}-X_{3}\right|=\mathcal{H}^{1}\left(p_{\xi}\left(\left[x_{1}, x_{3}\right]\right)\right) \leq\left|x_{1}-x_{3}\right| \leq \frac{2 \mathcal{L}^{2}(T)}{h_{T}}$, we obtain that

$$
\left|z-z^{\prime}\right|=d^{\prime} \frac{L^{\prime}-L}{L^{\prime}} \leq \frac{2 \mathcal{L}^{2}(T)}{h_{T}} \frac{\left|L-L^{\prime}\right|}{L^{\prime}}
$$

so that (3.28) holds in that case. The proof of the other case $X_{2} \cdot \xi^{\perp} \leq z \cdot \xi^{\perp}, z^{\prime} \cdot \xi^{\perp} \leq X_{3} \cdot \xi^{\perp}$ is similar and we omit it.

For all $j \geq j_{1}(\gamma)$ and all $T \in \widehat{\mathbf{T}}_{j}$, we have $L^{\max }(T)>L^{\text {ref }}(T)$. Indeed, if such would not be the case, denoting by $y \in Y_{j} \backslash Z_{*}$ a point such that $(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset$, then $\mathcal{L}^{1}\left(T_{p_{\xi}(y)}^{\xi}\right)=\mathcal{L}^{1}\left(T_{y}^{\xi}\right)=$ $b_{j}(y)-a_{j}(y) \leq L^{\max }(T) \leq L^{\text {ref }}(T)$, entailing that

$$
\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)\right|=\left|e\left(v_{j}\right)_{\mid T}:(\xi \otimes \xi)\right|\left(b_{j}(y)-a_{j}(y)\right) \leq\left|[u]\left(x_{0}\right) \cdot \xi\right|-4 C_{\eta} \gamma
$$

by definition (3.27) of $L^{\text {ref }}(T)$. Therefore, we would obtain that

$$
4 C_{\eta} \gamma \leq\left|[u]\left(x_{0}\right) \cdot \xi\right|-\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)\right| \leq\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)-[u]\left(x_{0}\right) \cdot \xi\right|
$$

which is against (3.26). Applying the Intermediate Value Theorem to the continuous functions $y \in\left[X_{1}, X_{3}\right] \mapsto \mathcal{L}^{1}\left(T_{y}^{\xi}\right) \in\left[0, L^{\max }(T)\right]$ and $y \in\left[X_{2}, X_{3}\right] \mapsto \mathcal{L}^{1}\left(T_{y}^{\xi}\right) \in\left[0, L^{\max }(T)\right]$, there are at least one and at most two points $z_{\text {ref }}^{1}, z_{\text {ref }}^{2} \in p_{\xi}(T)$ (according to whether $T$ has an edge along the direction $\xi$ or not, see Figure 5), only depending on $j$ and $T$, such that

$$
\mathcal{L}^{1}\left(T_{z_{\mathrm{ref}}^{1}}^{\xi}\right)=\mathcal{L}^{1}\left(T_{z_{\mathrm{ref}}^{2}}^{\xi}\right)=L^{\mathrm{ref}}(T)
$$

(We set $z_{\mathrm{ref}}^{1}=z_{\mathrm{ref}}^{2}$ in the case where $T$ has an edge along the direction $\xi$ ). Without loss of generality, we can assume that $z_{\text {ref }}^{1} \cdot \xi^{\perp} \geq z_{\text {ref }}^{2} \cdot \xi^{\perp}$.


Figure 5. Two possible configurations of $T$.
Let us introduce the following segments (orthogonal to $\xi$ ) associated to $T$ (see Figure 6),

$$
\begin{equation*}
\mathfrak{T}_{i}(T):=\left\{z \in \Pi_{\xi}:\left|z-z_{\mathrm{ref}}^{i}\right| \leq C_{\eta}^{\prime} \frac{\varrho_{j} \mathcal{L}^{2}(T)}{\varepsilon_{k_{j}}} \gamma\right\} \quad \text { for } i \in\{1,2\} \tag{3.29}
\end{equation*}
$$

where

$$
C_{\eta}^{\prime}:=\frac{20 C_{\eta}}{\sin \theta_{0}\left|[u]\left(x_{0}\right) \cdot \xi\right|}
$$

is a constant only depending on $\eta$.
For every $j \geq j_{1}(\gamma)$ and every $y \in Y_{j} \backslash Z_{*}$, let $T \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ be such that $(T \cap B)_{y}^{\xi} \neq \emptyset$. Note that $T \in \widehat{\mathbf{T}}_{j}$. If $p_{\xi}(y) \cdot \xi^{\perp} \geq X_{3} \cdot \xi^{\perp}$ and $z_{\text {ref }}^{1} \cdot \xi^{\perp} \geq X_{3} \cdot \xi^{\perp}$ (the other cases being treated similarly), applying (3.28) above, with $z=p_{\xi}(y)$ and $z^{\prime}=z_{\text {ref }}^{1}$, we get that

$$
\begin{aligned}
\left|p_{\xi}(y)-z_{\mathrm{ref}}^{1}\right| & \leq \frac{2 \mathcal{L}^{2}(T)}{h_{T}} \frac{\left|\left(b_{j}(y)-a_{j}(y)\right)-L^{\mathrm{ref}}(T)\right|}{\max \left(b_{j}(y)-a_{j}(y), L^{\text {ref }}(T)\right)} \\
& \leq \frac{2 \mathcal{L}^{2}(T)}{h_{T}} \frac{| |\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)\left|-\left|[u]\left(x_{0}\right) \cdot \xi\right|+4 C_{\eta} \gamma\right|}{\left|e\left(v_{j}\right)_{\mid T}:(\xi \otimes \xi)\right| L^{\text {ref }}(T)} \\
& \leq \frac{2 \mathcal{L}^{2}(T)}{h_{T}} \frac{5 C_{\eta} \gamma}{\left|[u]\left(x_{0}\right) \cdot \xi\right|-4 C_{\eta} \gamma} \\
& \leq C_{\eta}^{\prime} \frac{\underline{\mathcal{J}}_{j} \mathcal{L}^{2}(T)}{\varepsilon_{k_{j}}} \gamma,
\end{aligned}
$$

where we also used (3.26) and (3.25).
We have just shown that for all $j \geq j_{1}(\gamma)$ and all $y \in Y_{j} \backslash Z_{*}$, there exists $T \in \widehat{\mathbf{T}}_{j}$ such that $p_{\xi}(y) \in \mathfrak{T}_{1}(T) \cup \mathfrak{T}_{2}(T)$. Since $y \in \Pi_{\nu}$, then $y=\Phi\left(p_{\xi}(y)\right) \in \Phi\left(\mathfrak{T}_{1}(T) \cup \mathfrak{T}_{2}(T)\right)$, with $\Phi$ introduced in (3.16). Recalling that the Lipschitz constant of $\Phi$ is less than $\sqrt{1+4 \eta^{2}} \leq 2$ for $\eta$ small enough, we deduce that

$$
\mathcal{H}^{1}\left(\Phi\left(\mathfrak{T}_{1}(T) \cup \mathfrak{T}_{2}(T)\right)\right) \leq 2 \mathcal{H}^{1}\left(\mathfrak{T}_{1}(T) \cup \mathfrak{T}_{2}(T)\right) \leq 8 C_{\eta}^{\prime} \frac{\varrho_{j} \mathcal{L}^{2}(T)}{\varepsilon_{k_{j}}} \gamma .
$$



Figure 6. The length of $\mathfrak{T}_{i}(T)$ is given by $L=\frac{2 C_{\eta}^{\prime} \varrho_{j} \mathcal{L}^{2}(T)}{\varepsilon_{k_{j}}} \gamma$.

Together with the fact that each triangle in $\widehat{\mathbf{T}}_{j} \subset \mathbf{T}_{b, \text { int }}^{x_{0}, j}$ is contained in $B$, we obtain that for all $j \geq j_{1}(\gamma)$,

$$
\begin{aligned}
\mathcal{H}^{1}\left(Y_{j} \backslash Z_{*}\right) & \leq \sum_{T \in \widehat{\mathbf{T}}_{j}} \mathcal{H}^{1}\left(\Phi\left(\mathfrak{T}_{1}(T) \cup \mathfrak{T}_{2}(T)\right)\right) \\
& \leq 8 C_{\eta}^{\prime} \gamma \frac{\varrho_{j}}{\varepsilon_{k_{j}}} \sum_{T \in \widehat{\mathbf{T}}_{j}} \mathcal{L}^{2}(T) \leq \frac{8 C_{\eta}^{\prime} \gamma}{\kappa(1-\delta)} \frac{(1-\delta) \kappa \varrho_{j}}{\varepsilon_{k_{j}}} \int_{B} \chi_{j} d x \leq \frac{8 C_{\eta}^{\prime} \gamma}{\kappa(1-\delta)} \frac{\lambda_{k_{j}}\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{\varrho_{j}} .
\end{aligned}
$$

Possibly taking a larger $j_{1}(\gamma) \in \mathbb{N}$, we finally get that for all $j \geq j_{1}(\gamma)$,

$$
\mathcal{H}^{1}\left(Y_{j} \backslash Z_{*}\right) \leq \frac{8 C_{\eta}^{\prime}}{\kappa(1-\delta)}\left(2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right)+1\right) \gamma=: C_{*} \gamma
$$

for some constant $C_{*}>0$ only depending on $\eta$.
We are now in position to prove Lemma 3.8.
Proof of Lemma 3.8. Let $j_{0}=j_{0}(\eta) \in \mathbb{N}$ such that $\frac{\eta}{2^{j 0} C_{*}}<\gamma_{*}$, where $C_{*}$ and $\gamma_{*}$ are given by Lemma 3.9. For all $j \geq j_{0}$, as $0<\gamma_{j}:=\frac{\eta}{2^{j} C_{*}}<\gamma_{*}$, Lemma 3.9 ensures the existence of a rank $i\left(\gamma_{j}\right) \geq j_{0}$ such that for all $i \geq i\left(\gamma_{j}\right)$,

$$
\mathcal{H}^{1}\left(Y_{i} \backslash Z_{*}\right) \leq C_{*} \gamma_{j}=\frac{\eta}{2^{j}} .
$$

Thereby, we define the following extraction

$$
\left\{\begin{array}{l}
\phi\left(j_{0}\right):=i\left(\gamma_{j_{0}}\right) \geq j_{0}  \tag{3.30}\\
\phi(j+1):=\max \left(\phi(j)+1, i\left(\gamma_{j+1}\right)\right) \quad \text { for all } j \geq j_{0}
\end{array}\right.
$$

Since $\phi(j) \geq i\left(\gamma_{j}\right)$, then $\mathcal{H}^{1}\left(Y_{\phi(j)} \backslash Z_{*}\right) \leq \frac{\eta}{2^{j}}$. Hence, we set

$$
\begin{equation*}
Z^{\prime \prime}:=Z_{*} \cup \bigcup_{j=j_{0}}^{+\infty} Y_{\phi(j)} \tag{3.31}
\end{equation*}
$$

which satisfies

$$
\mathcal{H}^{1}\left(Z^{\prime \prime}\right) \leq 2 \eta+\sum_{j \geq j_{0}} \frac{\eta}{2^{j}} \leq 3 \eta .
$$

Moreover, for all $j \geq j_{0}$ and all $y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}$, Lemma 3.7 ensures that

$$
\#\left\{T \in \mathbf{T}_{b, i n t}^{x_{0}, \phi(j)}:(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset\right\} \geq 1
$$

and since $y \notin Y_{\phi(j)}$ for all $j \geq j_{0}$, it actually follows that

$$
\#\left\{T \in \mathbf{T}_{b, i n t}^{x_{0}, \phi(j)}:(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset\right\} \geq 2
$$

concluding the proof of Lemma 3.8.
In order to derive a lower bound for the surface energy without the factor $1 / 2$, we now construct two disjoint subfamilies $\mathscr{F}_{j}^{1}$ and $\mathscr{F}_{j}^{2}$ from $\mathscr{F}_{j}$ (see (3.19)) with the property that both sets

$$
\bigcup_{T \in \mathscr{F}_{j}^{1}} T, \quad \bigcup_{T \in \mathscr{F}_{j}^{2}} T
$$

project onto $B^{\nu}=J_{\bar{u}} \cap B$, thanks to the mapping $\Phi \circ p_{\xi}$, into two sets of almost full $\mathcal{H}^{1}$ measure in $J_{\bar{u}} \cap B$. This is the object of the following technical result.

Lemma 3.10. Let $K \subset J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}$ be a compact set. For all $j \in \mathbb{N}$, there exist two disjoint subfamilies $\mathscr{F}_{j}^{1}$ and $\mathscr{F}_{j}^{2}$ of $\mathscr{F}_{j}$ such that

$$
K \subset \Phi\left(\bigcup_{T \in \mathscr{F}_{j}^{1}} p_{\xi}(\stackrel{\circ}{T})\right) \cap \Phi\left(\bigcup_{T \in \mathscr{F}_{j}^{2}} p_{\xi}(\stackrel{\circ}{T})\right)
$$

Proof. For the sake of clarity, we omit to write the explicit dependance on $j$ for the different objects considered herafter (triangles, intervals, and so forth).

For all $y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}$, we consider a pair of distinct triangles of $\mathscr{F}_{j}$ satisfying

$$
\begin{align*}
\left\{T^{1}(y), T^{2}(y)\right\} \in & \arg \min \left\{\mathcal{H}^{1}\left(\Phi\left(p_{\xi}\left(\stackrel{\circ}{T}^{1}\right) \cap p_{\xi}\left(\stackrel{\circ}{2}^{2}\right)\right)\right):\right. \\
& \left.T^{1}, T^{2} \in \mathbf{T}_{b, \text { int }}^{x_{0}, j}, \stackrel{\circ}{T^{1}} \cap \stackrel{\circ}{2}_{T^{2}}^{=} \emptyset,\left(\stackrel{\circ}{T}^{1} \cap B\right)_{y}^{\xi} \neq \emptyset,\left(\stackrel{\circ}{T}^{2} \cap B\right)_{y}^{\xi} \neq \emptyset\right\} \tag{3.32}
\end{align*}
$$

Note that Lemma 3.8 ensures that the set

$$
\left\{\left\{T^{1}, T^{2}\right\} \subset \mathbf{T}_{b, i n t}^{x_{0}, j}:\left(\stackrel{\circ}{T}^{i} \cap B\right)_{y}^{\xi} \neq \emptyset \text { for all } i \in\{1,2\}\right\}
$$

is nonempty and finite, hence the minimum in (3.32) is achieved and we have at our disposal such a pair of distinct triangles $\left\{T^{1}(y), T^{2}(y)\right\}$. Then, we introduce the following open segment in $B^{\nu}=J_{\bar{u}} \cap B$

$$
\begin{equation*}
I(y):=\Phi\left(p_{\xi}\left(\stackrel{\circ}{T}^{1}(y)\right) \cap p_{\xi}\left(\stackrel{\circ}{T}^{2}(y)\right)\right) \subset B^{\nu}=J_{\bar{u}} \cap B . \tag{3.33}
\end{equation*}
$$

Since $y \in I(y)$, it follows that

$$
K \subset J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime} \subset \bigcup_{y \in J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}} I(y)
$$

Furthermore, $I(y)$ is optimal in the sense that any other triangle $T \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ satisfying $(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset$ is such that

$$
\begin{equation*}
I(y) \subset \Phi\left(p_{\xi}(\stackrel{\circ}{T})\right) \tag{3.34}
\end{equation*}
$$



Figure 7
Indeed, setting $J:=\Phi\left(p_{\xi}\left({ }^{\circ}\right)\right)$, there exist points $a, b, \tilde{a}$ and $\tilde{b}$ in $B^{\nu}=J_{\bar{u}} \cap B$ such that $I(y)=(a, b)$, $J=(\tilde{a}, \tilde{b})$ with $a \cdot \nu^{\perp}<b \cdot \nu^{\perp}$ and $\tilde{a} \cdot \nu^{\perp}<\tilde{b} \cdot \nu^{\perp}$. By construction, there exist $i_{a}$ (resp. $i_{b}$ ) $\in\{1,2\}$ such that $a$ (resp. b) is the image by $\Phi \circ p_{\xi}$ of a vertex of $T^{i_{a}}(y)$ (resp. $T^{i_{b}}(y)$ ), see Figure 7 . Assume by contradiction that there exists a point $z \in I(y) \backslash J$. If $y \cdot \nu^{\perp}<z \cdot \nu^{\perp}$, then $J \subset(\tilde{a}, z)$ since $J$ is a segment containing $y$. In particular, $\Phi\left(p_{\xi}\left(T^{i_{a}}(y)\right)\right) \cap J \subset(a, z)$. Together with (3.32) and recalling that $J:=\Phi\left(p_{\xi}\left({ }^{\circ}\right)\right)$, it ensures that

$$
\mathcal{H}^{1}((a, z)) \geq \mathcal{H}^{1}\left(\Phi\left(p_{\xi}\left(\stackrel{\circ}{T}^{i_{a}}(y)\right)\right) \cap \Phi\left(p_{\xi}(\stackrel{\circ}{T})\right)\right) \geq \mathcal{H}^{1}(I(y))=\mathcal{H}^{1}((a, z))+\mathcal{H}^{1}((z, b))>\mathcal{H}^{1}((a, z))
$$

which is impossible. A similar argument shows that the other situation $y \cdot \nu^{\perp}>z \cdot \nu^{\perp}$ is also impossible. This shows the validity of (3.34).

By compactness of $K$, there exist an integer $N=N(j, K) \geq 1$ and points $y_{1}, \ldots, y_{N} \in J_{\bar{u}} \cap$ $B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}$ such that

$$
\begin{equation*}
K \subset \bigcup_{i=1}^{N} I\left(y_{i}\right) \tag{3.35}
\end{equation*}
$$

Up to relabeling the points $y_{i}$, we can assume that $y_{1} \cdot \nu^{\perp}<\cdots<y_{N} \cdot \nu^{\perp}$ (See Figure 8). Let us now construct two disjoint subfamilies $\mathscr{F}_{j}^{1}$ and $\mathscr{F}_{j}^{2}$ of $\mathscr{F}_{j}$ by induction in $N$ iterations.

Iteration 1. Set $\mathscr{F}^{1}(1):=\left\{T^{1}\left(y_{1}\right)\right\}$ and $\mathscr{F}^{2}(1):=\left\{T^{2}\left(y_{1}\right)\right\}$. Clearly $\mathscr{F}^{1}(1) \cap \mathscr{F}^{2}(1)=\emptyset$ and, for all $k \in\{1,2\}$, there is $T \in \mathscr{F}^{k}(1)$ such that $(B \cap \stackrel{\circ}{T})_{y_{1}}^{\xi} \neq \emptyset$ and $I\left(y_{1}\right) \subset \Phi\left(p_{\xi}(\stackrel{\circ}{T})\right)$.

Iteration 2. We distinguish two cases:
i) If $\left\{T^{1}\left(y_{2}\right), T^{2}\left(y_{2}\right)\right\} \cap\left(\mathscr{F}^{1}(1) \cup \mathscr{F}^{2}(1)\right)=\emptyset$, then we set $\mathscr{F}^{1}(2):=\mathscr{F}^{1}(1) \cup\left\{T^{1}\left(y_{2}\right)\right\}$ and $\mathscr{F}^{2}(2):=\mathscr{F}^{2}(1) \cup\left\{T^{2}\left(y_{2}\right)\right\}$. We have that $\mathscr{F}^{1}(2) \cap \mathscr{F}^{2}(2)=\emptyset$ and, for all $i, k \in\{1,2\}$, there exists $T \in \mathscr{F}^{k}(2)$ such that $(B \cap \stackrel{\circ}{T})_{y_{i}}^{\xi} \neq \emptyset$ and $I\left(y_{i}\right) \subset \Phi\left(p_{\xi}(\stackrel{\circ}{T})\right)$.
ii) Otherwise, there exist $i, k \in\{1,2\}$ such that $T^{i}\left(y_{2}\right) \in \mathscr{F}^{k}(1)$, i.e. $T^{i}\left(y_{2}\right)=T^{k}\left(y_{1}\right)$, and $T^{3-i}\left(y_{2}\right) \notin \mathscr{F}^{k}(1)$. In that case, we set $\mathscr{F}^{k}(2):=\mathscr{F}^{k}(1)$ and $\mathscr{F}^{3-k}(2):=\mathscr{F}^{3-k}(1) \cup$ $\left\{T^{3-i}\left(y_{2}\right)\right\}$. Note that, it might be the case that $T^{3-i}\left(y_{2}\right) \in \mathscr{F}^{3-k}(1)$. We have that $\mathscr{F}^{1}(2) \cap \mathscr{F}^{2}(2)=\emptyset$ and, for all $i, k \in\{1,2\}$, there exists $T \in \mathscr{F}^{k}(2)$ such that $(B \cap T \times)_{y_{i}}^{\xi} \neq \emptyset$ and $I\left(y_{i}\right) \subset \Phi\left(p_{\xi}(\overleftarrow{T})\right)$.

Iteration $n+1$ for some $n \in\{1, \ldots, N-1\}$. Assume that we have constructed two disjoint subfamilies $\mathscr{F}^{1}(n)$ and $\mathscr{F}^{2}(n)$ of $\mathscr{F}_{j}$ with the following properties: for all $k \in\{1,2\}$ and all $i \in\{1, \ldots, n\}$, there exists $T \in \mathscr{F}^{k}(n)$ such that $\left(B \cap{ }_{T}\right)_{y_{i}}^{\xi} \neq \emptyset$ and $I\left(y_{i}\right) \subset \Phi\left(p_{\xi}(\stackrel{\circ}{T})\right)$. Let us now construct $\mathscr{F}^{1}(n+1)$ and $\mathscr{F}^{2}(n+1)$ :
i) If $\left\{T^{1}\left(y_{n+1}\right), T^{2}\left(y_{n+1}\right)\right\} \cap\left(\mathscr{F}^{1}(n) \cup \mathscr{F}^{2}(n)\right)=\emptyset$, then we set $\mathscr{F}^{1}(n+1):=\mathscr{F}^{1}(n) \cup$ $\left\{T^{1}\left(y_{n+1}\right)\right\}$ and $\mathscr{F}^{2}(n+1):=\mathscr{F}^{2}(n) \cup\left\{T^{2}\left(y_{n+1}\right)\right\}$. In that case, we have that $\mathscr{F}^{1}(n+$ 1) $\cap \mathscr{F}^{2}(n+1)=\emptyset$ and that for all $k \in\{1,2\}$ and all $i \in\{1, \ldots, n+1\}$, there exists $T \in \mathscr{F}^{k}(n+1)$ such that $\left(B \cap \grave{T}^{\circ}\right)_{y_{i}}^{\xi} \neq \emptyset$ and $I\left(y_{i}\right) \subset \Phi\left(p_{\xi}\left(\frac{\circ}{T}\right)\right)$. For $i \in\{1, \ldots, n\}$, it follows from the previous iteration $n$ and because $\mathscr{F}^{k}(n) \subset \mathscr{F}^{k}(n+1)$, while for $i=n+1$, it is a consequence of the fact that $T^{k}\left(y_{n+1}\right) \in \mathscr{F}^{k}(n+1)$.
ii) Otherwise, there exist $p, q \in\{1,2\}$ such that $T^{p}\left(y_{n+1}\right) \in \mathscr{F}^{q}(n)$. Let us further distinguish two subcases:
(a) If $T^{3-p}\left(y_{n+1}\right) \notin \mathscr{F}^{q}(n)$, then we set $\mathscr{F}^{q}(n+1):=\mathscr{F}^{q}(n)$ and $\mathscr{F}^{3-q}(n+1):=$ $\mathscr{F}^{3-q}(n) \cup\left\{T^{3-p}\left(y_{n+1}\right)\right\}$. Then $\mathscr{F}^{1}(n+1) \cap \mathscr{F}^{2}(n+1)=\emptyset$ and, for all $k \in\{1,2\}$ and all $i \in\{1, \ldots, n+1\}$, there exists $T \in \mathscr{F}^{k}(n+1)$ such that $\left(B \cap{ }_{T}\right)_{y_{i}}^{\xi} \neq \emptyset$ and $I\left(y_{i}\right) \subset \Phi\left(p_{\xi}(T)\right)$. Indeed, for $i \in\{1, \ldots, n\}$ this is a consequence of the previous iteration $n$ and of the fact that $\mathscr{F}^{k}(n) \subset \mathscr{F}^{k}(n+1)$, while, for $i=n+1$, it results from $T^{p}\left(y_{n+1}\right) \in \mathscr{F}^{q}(n+1)$ and $T^{3-p}\left(y_{n+1}\right) \in \mathscr{F}^{3-q}(n+1)$.
(b) If both $T^{1}:=T^{1}\left(y_{n+1}\right) \in \mathscr{F}^{q}(n)$ and $T^{2}:=T^{2}\left(y_{n+1}\right) \in \mathscr{F}^{q}(n)$, we introduce the indexes

$$
i_{1}:=\arg \min \left\{i \in\{1, \ldots, n+1\}:\left(B \cap \stackrel{\circ}{1}^{1}\right)_{y_{i}}^{\xi} \neq \emptyset\right\}
$$

and

$$
i_{2}:=\arg \min \left\{i \in\{1, \ldots, n+1\}:\left(B \cap \stackrel{\circ}{T}^{2}\right)_{y_{i}}^{\xi} \neq \emptyset\right\}
$$

Up to interchanging $T_{1}$ and $T_{2}$, there is no loss of generality to assume that $i_{1} \geq i_{2}$ (see Figure 8). We set $\mathscr{F}^{q}(n+1):=\mathscr{F}^{q}(n) \backslash\left\{T^{1}\right\}$ and $\mathscr{F}^{3-q}(n+1):=\mathscr{F}^{3-q}(n) \cup\left\{T^{1}\right\}$. Once more, we have $\mathscr{F}^{1}(n+1) \cap \mathscr{F}^{2}(n+1)=\emptyset$ and, for all $k \in\{1,2\}$ and all $i \in\{1, \ldots, n+1\}$, there exists $T \in \mathscr{F}^{k}(n+1)$ such that $(B \cap \stackrel{\circ}{T})_{y_{i}}^{\xi} \neq \emptyset$ and $I\left(y_{i}\right) \subset$ $\Phi\left(p_{\xi}\left({ }_{T}\right)\right)$. This is immediate for $i=n+1$ because $T^{2}=T^{2}\left(y_{n+1}\right) \in \mathscr{F}^{q}(n+1)$ and $T^{1}=T^{1}\left(y_{n+1}\right) \in \mathscr{F}^{3-q}(n+1)$. For $i \in\{1, \ldots, n\}$, there are two possibilities:

- for $k=3-q$, it follows from $\mathscr{F}^{3-q}(n) \subset \mathscr{F}^{3-q}(n+1)$.
- for $k=q$ and $i \in\left\{1, \ldots, i_{1}-1\right\}$, by the previous iteration $n$ there exists $T \in \mathscr{F}^{q}(n)$ such that $(B \cap \stackrel{\circ}{T})_{y_{i}}^{\xi} \neq \emptyset$ and $I\left(y_{i}\right) \subset \Phi\left(p_{\xi}(\stackrel{\circ}{T})\right)$. As $\left(B \cap \stackrel{\circ}{1}^{1}\right)_{y_{i}}^{\xi}=\emptyset$, by definition of $i_{1}$, this implies that $T \neq T^{1}$, so that actually $T \in \mathscr{F}^{q}(n+1)$ satisfies the above requirements. Assuming next that $i \in\left\{i_{1}, \ldots, n\right\}$, we deduce
that $n+1>i \geq i_{2}$. By definition of $i_{2}$, we have $\left(B \cap \overparen{T}^{2}\right)_{y_{i_{2}}}^{\xi} \neq \emptyset$ while $\left(B \cap \stackrel{\circ}{T}^{2}\right)_{y_{n+1}}^{\xi} \neq \emptyset$ so that the convexity of $\stackrel{\circ}{T}^{2}$ together with the ordering of the points $y_{i}$ lead to $\left(B \cap \stackrel{\circ}{T}^{2}\right)_{y_{i}}^{\xi} \neq \emptyset$. Hence, owing to (3.34), we infer that $I\left(y_{i}\right) \subset \Phi\left(p_{\xi}\left(T^{2}\right)\right)$ and $T^{2} \in \mathscr{F}^{q}(n+1)$ satisfies the above requirements.
We proceed this construction up to the $N^{\text {th }}$ iteration, and finally define

$$
\begin{equation*}
\mathscr{F}_{j}^{k}:=\mathscr{F}^{k}(N) \text { for } k \in\{1,2\} \tag{3.36}
\end{equation*}
$$

which define two disjoint subfamilies of $\mathscr{F}_{j}$ satisfying in particular, thanks to (3.35),

$$
\begin{aligned}
K & \subset \bigcup_{i=1}^{N} I\left(y_{i}\right) \subset\left(\bigcup_{T \in \mathscr{F}_{j}^{1}} \Phi\left(p_{\xi}(\stackrel{\circ}{T})\right)\right) \cap\left(\bigcup_{T \in \mathscr{F}_{j}^{2}} \Phi\left(p_{\xi}(\stackrel{\circ}{T})\right)\right) \\
& =\Phi\left(\bigcup_{T \in \mathscr{F}_{j}^{1}} p_{\xi}(\stackrel{\circ}{T})\right) \cap \Phi\left(\bigcup_{T \in \mathscr{F}_{j}^{2}} p_{\xi}(\stackrel{\circ}{T})\right) .
\end{aligned}
$$

The proof of Lemma 3.10 is now complete.


Figure 8

We are now in position to complete the proof of Proposition 3.3.
Proof of Proposition 3.3. Since all triangles $T$ in $\mathscr{F}_{j}$ are contained in $B$, we get from (3.10b) and (3.11),

$$
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right)=\lim _{j \rightarrow \infty} \frac{\lambda_{k_{j}}\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{2 \varrho_{j}} \geq \liminf _{j \rightarrow \infty} \frac{(1-\delta) \kappa \varrho_{j}}{2 \varepsilon_{k_{j}}} \int_{B} \chi_{j} d x \geq \liminf _{j \rightarrow \infty} \frac{(1-\delta) \kappa \varrho_{j}}{2 \varepsilon_{k_{j}}} \sum_{T \in \mathscr{F}_{j}} \mathcal{L}^{2}(T) .
$$

Using next the inequality $\mathcal{L}^{2}(T) \geq \mathcal{H}^{1}\left(p_{\xi}(T)\right)\left(\varepsilon_{k_{j}} / \varrho_{j}\right) \sin \theta_{0} / 2$, we deduce that

$$
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{2} \liminf _{j \rightarrow \infty} \sum_{T \in \mathscr{F}_{j}} \mathcal{H}^{1}\left(p_{\xi}(T)\right)
$$

Let $K \subset J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}$ be a compact set and $\mathscr{F}_{j}^{1}$ and $\mathscr{F}_{j}^{2}$ be two disjoint subfamilies of $\mathscr{F}_{j}$ given by Lemma 3.10. Thus

$$
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{2} \liminf _{j \rightarrow \infty}\left\{\sum_{T \in \mathscr{F}_{j}^{1}} \mathcal{H}^{1}\left(p_{\xi}(T)\right)+\sum_{T \in \mathscr{F}_{j}^{2}} \mathcal{H}^{1}\left(p_{\xi}(T)\right)\right\}
$$

and remembering that $\Phi$ has a Lipschitz constant bounded by $\sqrt{1+4 \eta^{2}}$,

$$
\begin{aligned}
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) & \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{2 \sqrt{1+4 \eta^{2}}} \liminf _{j \rightarrow \infty}\left\{\sum_{T \in \mathscr{F}_{j}^{1}} \mathcal{H}^{1}\left(\Phi\left(p_{\xi}(T)\right)\right)+\sum_{T \in \mathscr{F}_{j}^{2}} \mathcal{H}^{1}\left(\Phi\left(p_{\xi}(T)\right)\right)\right\} \\
& \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{2 \sqrt{1+4 \eta^{2}}} \liminf _{j \rightarrow \infty}\left\{\mathcal{H}^{1}\left(\bigcup_{T \in \mathscr{F}_{j}^{1}} \Phi\left(p_{\xi}(T)\right)\right)+\mathcal{H}^{1}\left(\bigcup_{T \in \mathscr{F}_{j}^{2}} \Phi\left(p_{\xi}(T)\right)\right)\right\} \\
& \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{\sqrt{1+4 \eta^{2}}} \mathcal{H}^{1}(K)
\end{aligned}
$$

By inner regularity of the Radon measure $\mathcal{H}^{1}\left\llcorner\left(J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}\right)\right.$, passing to the supremum with respect to all compact sets $K \subset J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}$, we get that

$$
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{\sqrt{1+4 \eta^{2}}} \mathcal{H}^{1}\left(J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}\right)
$$

Remembering that $J_{\bar{u}} \cap B=B^{\nu}$, we have

$$
2=\mathcal{H}^{1}\left(J_{\bar{u}} \cap B\right)=\mathcal{H}^{1}\left(J_{\bar{u}} \cap B_{1-\frac{\eta}{2}}\right)+\eta \leq \mathcal{H}^{1}\left(J_{\bar{u}} \cap B_{1-\frac{\eta}{2}} \backslash Z^{\prime \prime}\right)+4 \eta
$$

because $\mathcal{H}^{1}\left(Z^{\prime \prime}\right) \leq 3 \eta$. Hence

$$
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \geq \frac{(1-\delta) \kappa \sin \theta_{0}}{\sqrt{1+4 \eta^{2}}}(2-4 \eta)
$$

Finally passing to the limit as $\eta \rightarrow 0$ and $\delta \rightarrow 0$, we deduce that

$$
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner J_{u}\right.}\left(x_{0}\right) \geq \kappa \sin \theta_{0}
$$

which corresponds to the desired lower bound with the correct multiplicative constant.
3.3. The upper bound. The proof of the $\Gamma$-limsup inequality relies on suitable approximation results in $G S B D$ (see $[16,30,19,24]$ ) which allow us to reduce to the case where the jump set of $u$ is a finite union of pairwise disjoint closed line segments, and $u$ is smooth outside the jump set. Then, an explicit mesh construction introduced in [21], adapted to this simple geometrical situation, provides the desired upper bound.

Proposition 3.11. For all $u \in L^{0}\left(\Omega ; \mathbb{R}^{2}\right)$,

$$
\mathcal{F}^{\prime \prime}(u) \leq \mathcal{F}(u)
$$

Proof. We can assume that $\mathcal{F}(u)<\infty$, and thus that $u \in G S B D^{2}(\Omega)$. Using the density result for $G S B D$ functions (see [19, Theorem 1.1]) as well as the lower semicontinuity of $\mathcal{F}^{\prime \prime}$ with respect to the convergence in measure (see [25, Proposition 6.8]), we can further assume without loss of generality that $u \in S B V^{2}\left(\Omega ; \mathbb{R}^{2}\right) \cap L^{\infty}\left(\Omega ; \mathbb{R}^{2}\right)$.

Writing $u=\left(u_{1}, u_{2}\right)$, we can apply [21, Lemma 4.2] to both components $u_{1}$ and $u_{2} \in S B V^{2}(\Omega) \cap$ $L^{\infty}(\Omega)$ of $u$. For $\Omega^{\prime}:=(a, b) \times(c, d) \subset \mathbb{R}^{2}$ with $\Omega \subset \subset \Omega^{\prime}$, we can find an extension $v \in$ $S B V^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right) \cap L^{\infty}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ such that

$$
\begin{equation*}
v_{\mid \Omega}=u, \quad\|v\|_{L^{\infty}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)} \leq \sqrt{2}\|u\|_{L^{\infty}\left(\Omega ; \mathbb{R}^{2}\right)} \quad \text { and } \quad \mathcal{H}^{1}\left(\partial \Omega \cap J_{v}\right)=0 \tag{3.37}
\end{equation*}
$$

Next owing to the density result in $S B V$ (see [24, Theorem 3.1]), there exists a sequence $\left\{v_{k}\right\}_{k \in \mathbb{N}}$ in $S B V^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right) \cap L^{\infty}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ as well as $N_{k}$ disjoint closed segments $L_{1}^{k}, \ldots, L_{N_{k}}^{k} \subset \overline{\Omega^{\prime}}$ with the following properties:

$$
\overline{J_{v_{k}}}=\bigcup_{i=1}^{N_{k}} L_{i}^{k}, \quad \mathcal{H}^{1}\left(\overline{J_{v_{k}}} \backslash J_{v_{k}}\right)=0, \quad v_{k} \in W^{2, \infty}\left(\Omega^{\prime} \backslash \overline{J_{v_{k}}} ; \mathbb{R}^{2}\right)
$$

and

$$
\left\{\begin{array}{l}
v_{k} \rightarrow v \quad \text { strongly in } L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)  \tag{3.38}\\
\nabla v_{k} \rightarrow \nabla v \quad \text { strongly in } L^{2}\left(\Omega^{\prime} ; \mathbb{M}^{2 \times 2}\right), \\
\limsup _{k} \mathcal{H}^{1}\left(\bar{A} \cap J_{v_{k}}\right) \leq \mathcal{H}^{1}\left(\bar{A} \cap J_{v}\right) \text { for all open subset } A \subset \subset \Omega^{\prime}
\end{array}\right.
$$

Using (3.38) and the lower semicontinuity of $\mathcal{F}^{\prime \prime}$ in $L^{0}\left(\Omega ; \mathbb{R}^{2}\right)$ with respect to the convergence in measure, we obtain that

$$
\mathcal{F}^{\prime \prime}(u) \leq \liminf _{k \rightarrow \infty} \mathcal{F}^{\prime \prime}\left(v_{k \mid \Omega}\right)
$$

The proof is complete once we know that $\liminf _{k} \mathcal{F}^{\prime \prime}\left(v_{k \mid \Omega}\right) \leq \mathcal{F}(u)$. This follows from Lemma 3.12 below, applied to each function $v_{k}$. Indeed, using that result, we get that

$$
\liminf _{k \rightarrow \infty} \mathcal{F}^{\prime \prime}\left(v_{k \mid \Omega}\right) \leq \liminf _{k \rightarrow \infty}\left\{\int_{\Omega} \mathbf{A} e\left(v_{k}\right): e\left(v_{k}\right) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v_{k}} \cap \bar{\Omega}\right)\right\}
$$

Recalling the convergences (3.38), we conclude that

$$
\mathcal{F}^{\prime \prime}(u) \leq \int_{\Omega} \mathbf{A} e(v): e(v) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v} \cap \bar{\Omega}\right)=\int_{\Omega} \mathbf{A} e(u): e(u) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{u}\right)=\mathcal{F}(u)
$$

where we used $\mathcal{H}^{1}\left(J_{v} \cap \partial \Omega\right)=0$ and that $v=u$ in $\Omega$.
We are back to establishing the following result.
Lemma 3.12. Let $v \in S B V^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right) \cap L^{\infty}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ be such that

$$
\overline{J_{v}}=\bigcup_{i=1}^{N} L_{i}, \quad \mathcal{H}^{1}\left(\overline{J_{v}} \backslash J_{v}\right)=0, \quad v \in W^{2, \infty}\left(\Omega^{\prime} \backslash \overline{J_{v}} ; \mathbb{R}^{2}\right)
$$

for some pairwise disjoint closed segments $L_{1}, \ldots, L_{N} \subset \overline{\Omega^{\prime}}$. Then,

$$
\mathcal{F}^{\prime \prime}\left(v_{\mid \Omega}\right) \leq \int_{\Omega} \mathbf{A} e(v): e(v) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v} \cap \bar{\Omega}\right)
$$

Proof. Since $\Omega \subset \subset \Omega^{\prime}$, then $d:=\operatorname{dist}\left(\Omega, \mathbb{R}^{2} \backslash \Omega^{\prime}\right)>0$. For all $\delta \in(0, d)$, let us consider the open sets

$$
\Omega_{\delta}:=\left\{x \in \Omega^{\prime}: \operatorname{dist}\left(x, \mathbb{R}^{2} \backslash \Omega^{\prime}\right)>\delta\right\}
$$

which satisfy $\Omega \subset \subset \Omega_{\delta} \subset \subset \Omega^{\prime}$. We introduce a cut-off function $\phi_{\delta} \in C_{c}^{\infty}\left(\mathbb{R}^{2} ;[0,1]\right)$ which is supported in $\Omega^{\prime}$ and such that $\phi_{\delta}=1$ in $\Omega_{\delta}, \phi_{\delta}=0$ in $\mathbb{R}^{2} \backslash \Omega_{\frac{\delta}{2}}$. We next introduce the function $\bar{v}:=\phi_{\delta} v \in S B V^{2}\left(\mathbb{R}^{2} ; \mathbb{R}^{2}\right) \cap L^{\infty}\left(\mathbb{R}^{2} ; \mathbb{R}^{2}\right)$. We remark that

$$
\begin{equation*}
\bar{v} \in W^{2, \infty}\left(\mathbb{R}^{2} \backslash \bigcup_{i=1}^{N} \overline{L_{i}} ; \mathbb{R}^{2}\right), \quad J_{\bar{v}} \subset J_{v}, \quad \text { and } \quad J_{v} \backslash J_{\bar{v}} \subset J_{v} \backslash \Omega_{\delta} \tag{3.39}
\end{equation*}
$$

Since $J_{v} \subset \overline{J_{v}}$ and $\mathcal{H}^{1}\left(\overline{J_{v}}\right)<\infty$, the disjoint closed segments $L_{i} \subset \overline{\Omega^{\prime}}$ satisfy

$$
\mathcal{H}^{1}\left(J_{v}\right)=\mathcal{H}^{1}\left(\overline{J_{v}}\right)=\sum_{i=1}^{N} \mathcal{H}^{1}\left(L_{i}\right)
$$

Then according to [21, Appendix A], since $\theta_{0}$ is smaller than or equal to $\Theta_{0}:=45^{\circ}-\arctan (1 / 2)$, for all $\varepsilon>0$ there exists an admissible triangulation $\mathbf{T}_{\varepsilon} \in \mathcal{T}_{\varepsilon}\left(\mathbb{R}^{2}, \omega, \theta_{0}\right)$ such that, setting $\mathbf{T}_{\varepsilon}^{\prime}:=$ $\left\{T \in \mathbf{T}_{\varepsilon}: T \cap \bigcup_{i} L_{i} \neq \emptyset\right\}$,

- The vertices of $\mathbf{T}_{\varepsilon}$ are never situated on any $L_{i}$ :

$$
\text { for all } i \in\{1, \ldots, N\}, \quad L_{i} \cap \operatorname{Vertices}\left(\mathbf{T}_{\varepsilon}\right)=\emptyset
$$

- Using [21, Formula (4.9)],

$$
\begin{equation*}
\sum_{T \in \mathbf{T}_{\varepsilon}^{\prime}} \frac{\mathcal{L}^{2}(T)}{\varepsilon} \rightarrow \sin \theta_{0} \mathcal{H}^{1}\left(J_{v}\right) \tag{3.40}
\end{equation*}
$$

We define the set $D_{\varepsilon}:=\bigcup_{T \in \mathbf{T}_{\varepsilon}^{\prime}} T$ and $\chi_{\varepsilon}:=\mathbb{1}_{D_{\varepsilon}} \in L^{\infty}\left(\mathbb{R}^{2} ;\{0,1\}\right)$, while $\bar{v}_{\varepsilon}$ is the Lagrange interpolation of the values of $\bar{v}$ at the vertices of the triangulation $\mathbf{T}_{\varepsilon}$. Note that, if $x_{1}, x_{2}$ and $x_{3}$ are the vertices of $T \in \mathbf{T}_{\varepsilon}$, the values $\bar{v}\left(x_{i}\right)$ are well defined since, by construction of the triangulation $\mathbf{T}_{\varepsilon}$, the points $x_{1}, x_{2}$ and $x_{3}$ do not belong to $\bigcup_{i} L_{i}$. In particular, $\bar{v}_{\varepsilon} \in V_{\varepsilon}\left(\Omega^{\prime}, \omega, \theta_{0}\right)$
and

$$
\begin{cases}\chi_{\varepsilon} \rightarrow 0 & \text { strongly in } L^{1}\left(\Omega^{\prime}\right)  \tag{3.41}\\ \bar{v}_{\varepsilon} \rightarrow \bar{v} & \text { strongly in } L^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right) \\ e\left(\bar{v}_{\varepsilon}\right) \mathbb{1}_{\Omega^{\prime} \backslash D_{\varepsilon}} \rightarrow e(\bar{v}) & \text { strongly in } L^{2}\left(\Omega^{\prime} ; \mathbb{M}_{\text {sym }}^{2 \times 2}\right)\end{cases}
$$

Indeed, the first convergence is a consequence of (3.40) since

$$
\left\|\chi_{\varepsilon}\right\|_{L^{1}\left(\Omega^{\prime}\right)} \leq \mathcal{L}^{2}\left(D_{\varepsilon}\right)=\sum_{T \in \mathbf{T}_{\varepsilon}^{\prime}} \mathcal{L}^{2}(T) \rightarrow 0
$$

Next, noticing that every $T \in \mathbf{T}_{\varepsilon} \backslash \mathbf{T}_{\varepsilon}^{\prime}$ is contained in $\mathbb{R}^{2} \backslash \bigcup_{i} L_{i}$ and $\bar{v} \in W^{2, \infty}\left(\mathbb{R}^{2} \backslash \bigcup_{i} L_{i} ; \mathbb{R}^{2}\right)$, we infer that for all $\varepsilon>0$ and $T \in \mathbf{T}_{\varepsilon} \backslash \mathbf{T}_{\varepsilon}^{\prime}$,

$$
\begin{equation*}
\left\|\bar{v}_{\varepsilon}-\bar{v}\right\|_{H^{1}\left(T ; \mathbb{R}^{2}\right)} \leq C \varepsilon\left\|D^{2} \bar{v}\right\|_{L^{2}(T)} \tag{3.42}
\end{equation*}
$$

for some constant $C=C\left(\theta_{0}\right)>0$ depending only on $\theta_{0}$ (see e.g. [23, Theorem 3.1.5]). On the one hand, since $\left\|\bar{v}_{\varepsilon}\right\|_{L^{\infty}\left(T ; \mathbb{R}^{2}\right)} \leq\|\bar{v}\|_{L^{\infty}\left(T ; \mathbb{R}^{2}\right)}$ for all $T \in \mathbf{T}_{\varepsilon}$, we get that

$$
\left\|\bar{v}_{\varepsilon}-\bar{v}\right\|_{L^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)}^{2} \leq 2\|\bar{v}\|_{L^{\infty}\left(\mathbb{R}^{2} ; \mathbb{R}^{2}\right)}^{2} \sum_{T \in \mathbf{T}_{\varepsilon}^{\prime}} \mathcal{L}^{2}(T)+\sum_{T \in \mathbf{T}_{\varepsilon} \backslash \mathbf{T}_{\varepsilon}^{\prime}, T \cap \Omega^{\prime} \neq \emptyset} \int_{T}\left|\bar{v}_{\varepsilon}-\bar{v}\right|^{2} d x
$$

Then, using (3.42) yields

$$
\begin{aligned}
\left\|\bar{v}_{\varepsilon}-\bar{v}\right\|_{L^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)}^{2} & \leq 2\|\bar{v}\|_{L^{\infty}\left(\mathbb{R}^{2} ; \mathbb{R}^{2}\right)}^{2} \mathcal{L}^{2}\left(D_{\varepsilon}\right)+C^{2} \varepsilon^{2} \sum_{T \in \mathbf{T}_{\varepsilon} \backslash \mathbf{T}_{\varepsilon}^{\prime}, T \cap \Omega^{\prime} \neq \emptyset}\left\|D^{2} \bar{v}\right\|_{L^{2}(T)}^{2} \\
& \leq 2\|\bar{v}\|_{L^{\infty}\left(\mathbb{R}^{2} ; \mathbb{R}^{2}\right)}^{2} \mathcal{L}^{2}\left(D_{\varepsilon}\right)+C^{2} \varepsilon^{2}\left\|D^{2} \bar{v}\right\|_{L^{2}\left(\mathbb{R}^{2} \backslash \cup_{i} L_{i}\right)}^{2} \rightarrow 0
\end{aligned}
$$

leading to the second convergence in (3.41). Note in particular that $\bar{v}_{\varepsilon}$ converges in measure to $\bar{v}$ in $\Omega^{\prime}$. On the other hand, writing

$$
\left\|e\left(\bar{v}_{\varepsilon}\right) \mathbb{1}_{\Omega^{\prime} \backslash D_{\varepsilon}}-e(\bar{v})\right\|_{L^{2}\left(\Omega^{\prime} ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right)}^{2}=\int_{\Omega^{\prime} \cap D_{\varepsilon}}|e(\bar{v})|^{2} d x+\int_{\Omega^{\prime} \backslash D_{\varepsilon}}\left|e\left(\bar{v}_{\varepsilon}\right)-e(\bar{v})\right|^{2} d x
$$

and using that $\mathcal{L}^{2}\left(D_{\varepsilon}\right) \rightarrow 0$, that $\nabla \bar{v} \in L^{2}\left(\mathbb{R}^{2} ; \mathbb{M}^{2 \times 2}\right)$ (because $\bar{v} \in S B V^{2}\left(\mathbb{R}^{2} ; \mathbb{R}^{2}\right)$ ) as well as (3.42), we get that

$$
\left\|e\left(\bar{v}_{\varepsilon}\right) \mathbb{1}_{\Omega^{\prime} \backslash D_{\varepsilon}}-e(\bar{v})\right\|_{L^{2}\left(\Omega^{\prime} ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right)}^{2} \leq \int_{\Omega^{\prime} \cap D_{\varepsilon}}|e(\bar{v})|^{2} d x+C^{2} \varepsilon^{2}\left\|D^{2} \bar{v}\right\|_{L^{2}\left(\mathbb{R}^{2} \backslash \bigcup_{i} L_{i}\right)}^{2} \rightarrow 0
$$

which implies the third convergence in (3.41).
We next show that

$$
\begin{equation*}
\mathcal{F}^{\prime \prime}\left(v_{\mid \Omega}\right) \leq \int_{\Omega} \mathbf{A} e(v): e(v) d x+\kappa \sin \theta_{0}\left(\mathcal{H}^{1}\left(J_{v} \cap \bar{\Omega}\right)+\mathcal{H}^{1}\left(J_{v} \backslash \Omega_{\delta}\right)\right) \tag{3.43}
\end{equation*}
$$

Indeed, as $f \leq \kappa$ thanks to the growth properties (1.4), we get
$\int_{\Omega^{\prime}} \frac{1}{\varepsilon} f\left(\varepsilon \mathbf{A} e\left(\bar{v}_{\varepsilon}\right): e\left(\bar{v}_{\varepsilon}\right)\right) d x \leq \sum_{T \in \mathbf{T}_{\varepsilon} \backslash \mathbf{T}_{\varepsilon}^{\prime}, T \cap \Omega^{\prime} \neq \emptyset} \mathcal{L}^{2}\left(T \cap \Omega^{\prime}\right) \frac{1}{\varepsilon} f\left(\varepsilon \mathbf{A} e\left(\bar{v}_{\varepsilon}\right)_{\mid T}: e\left(\bar{v}_{\varepsilon}\right)_{\mid T}\right)+\frac{\kappa}{\varepsilon} \sum_{T \in \mathbf{T}_{\varepsilon}^{\prime}} \mathcal{L}^{2}(T)$.
On the one hand, (3.40) implies that

$$
\begin{equation*}
\frac{\kappa}{\varepsilon} \sum_{T \in \mathbf{T}_{\varepsilon}^{\prime}} \mathcal{L}^{2}(T) \rightarrow \kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v}\right) \tag{3.44}
\end{equation*}
$$

On the other hand, since every triangle $T \in \mathbf{T}_{\varepsilon} \backslash \mathbf{T}_{\varepsilon}^{\prime}$ is contained in $\mathbb{R}^{2} \backslash \bigcup_{i=1}^{N} L_{i}$, then

$$
\left|\nabla \bar{v}_{\varepsilon \mid T} \frac{x_{i}-x_{j}}{\left|x_{i}-x_{j}\right|}\right|=\frac{\left|\bar{v}\left(x_{i}\right)-\bar{v}\left(x_{j}\right)\right|}{\left|x_{i}-x_{j}\right|} \leq\|\nabla \bar{v}\|_{L^{\infty}\left(\mathbb{R}^{2} \backslash \bigcup_{i} L_{i} ; \mathbb{M}^{2 \times 2}\right)}
$$

where $x_{1}, x_{2}$ and $x_{3}$ are the vertices of $T$. Hence, applying [21, Remark 3.5], it results that

$$
\begin{equation*}
\left\|e\left(\bar{v}_{\varepsilon}\right)\right\|_{L^{\infty}\left(\mathbb{R}^{2} \backslash D_{\varepsilon} ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right)} \leq \frac{\sqrt{5}}{\sin \theta_{0}}\|\nabla \bar{v}\|_{L^{\infty}\left(\mathbb{R}^{2} \backslash \bigcup_{i} L_{i} ; \mathbb{M}^{2 \times 2}\right)}=: K<\infty \tag{3.45}
\end{equation*}
$$

Therefore, setting

$$
\delta_{\varepsilon}:=\sup _{0<t<\varepsilon \beta K^{2}} \frac{f(t)}{t}
$$

we deduce, using $f(0)=0$ and the property (1.5) of $\mathbf{A}$, that

$$
\frac{1}{\varepsilon} f\left(\varepsilon \mathbf{A} e\left(\bar{v}_{\varepsilon}\right): e\left(\bar{v}_{\varepsilon}\right)\left(1-\chi_{\varepsilon}\right)\right) \leq \delta_{\varepsilon} \mathbf{A} e\left(\bar{v}_{\varepsilon}\right): e\left(\bar{v}_{\varepsilon}\right)\left(1-\chi_{\varepsilon}\right) \quad \text { in } \Omega^{\prime}
$$

From the properties (1.4) of $f$, we infer that $\delta_{\varepsilon} \rightarrow 1$ as $\varepsilon \rightarrow 0$. Hence, using the third convergence in (3.41), it ensures that

$$
\begin{align*}
\sum_{T \in \mathbf{T}_{\varepsilon} \backslash \mathbf{T}_{\varepsilon}^{\prime}, T \cap \Omega^{\prime} \neq \emptyset} \mathcal{L}^{2}\left(T \cap \Omega^{\prime}\right) \frac{1}{\varepsilon} f\left(\varepsilon \mathbf{A} e\left(\bar{v}_{\varepsilon}\right)_{\mid T}: e\left(\bar{v}_{\varepsilon}\right)_{\mid T}\right) & =\int_{\Omega^{\prime}} \frac{1}{\varepsilon} f\left(\varepsilon \mathbf{A} e\left(\bar{v}_{\varepsilon}\right): e\left(\bar{v}_{\varepsilon}\right)\left(1-\chi_{\varepsilon}\right)\right) d x \\
& \leq \delta_{\varepsilon} \int_{\Omega^{\prime}} \mathbf{A} e\left(\bar{v}_{\varepsilon}\right): e\left(\bar{v}_{\varepsilon}\right) \mathbb{1}_{\Omega^{\prime} \backslash D_{\varepsilon}} d x \\
& \rightarrow \int_{\Omega^{\prime}} \mathbf{A} e(\bar{v}): e(\bar{v}) d x \tag{3.46}
\end{align*}
$$

Gathering (3.44) and (3.46), we obtain that

$$
\begin{equation*}
\limsup _{\varepsilon \rightarrow 0^{+}} \int_{\Omega^{\prime}} \frac{1}{\varepsilon} f\left(\varepsilon \mathbf{A} e\left(\bar{v}_{\varepsilon}\right): e\left(\bar{v}_{\varepsilon}\right)\right) d x \leq \int_{\Omega^{\prime}} \mathbf{A} e(\bar{v}): e(\bar{v}) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v}\right) . \tag{3.47}
\end{equation*}
$$

Besides, after decomposing the above integral over $\Omega^{\prime} \backslash \bar{\Omega}$ and $\bar{\Omega}$, we can apply the lower bound estimate of Propostion 3.3 to the open bounded set with Lipschitz boundary $\Omega^{\prime} \backslash \bar{\Omega}$ (for which $\mathbf{T}_{\varepsilon}$ is also an admissible triangulation, $\bar{v}_{\varepsilon \mid \Omega^{\prime} \backslash \bar{\Omega}} \in V_{\varepsilon}\left(\Omega^{\prime} \backslash \bar{\Omega}\right)$ and $\bar{v}_{\varepsilon}$ converges in measure to $\bar{v}$ in $\left.\Omega^{\prime} \backslash \bar{\Omega}\right)$, which leads to

$$
\begin{aligned}
\limsup _{\varepsilon \rightarrow 0^{+}} \int_{\Omega^{\prime}} & \frac{1}{\varepsilon} f\left(\varepsilon \mathbf{A} e\left(\bar{v}_{\varepsilon}\right): e\left(\bar{v}_{\varepsilon}\right)\right) d x \\
& \geq \limsup _{\varepsilon \rightarrow 0} \int_{\Omega} \frac{1}{\varepsilon} f\left(\varepsilon \mathbf{A} e\left(\bar{v}_{\varepsilon}\right): e\left(\bar{v}_{\varepsilon}\right)\right) d x+\liminf _{\varepsilon \rightarrow 0} \int_{\Omega^{\prime} \backslash \bar{\Omega}} \frac{1}{\varepsilon} f\left(\varepsilon \mathbf{A} e\left(\bar{v}_{\varepsilon}\right): e\left(\bar{v}_{\varepsilon}\right)\right) d x \\
& \geq \mathcal{F}^{\prime \prime}\left(\bar{v}_{\mid \Omega}\right)+\int_{\Omega^{\prime} \backslash \bar{\Omega}} \mathbf{A} e(\bar{v}): e(\bar{v}) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{\bar{v}} \cap \Omega^{\prime} \backslash \bar{\Omega}\right)
\end{aligned}
$$

Gathering (3.47) and (3.39), as by construction $\bar{v}_{\mid \Omega}=v_{\mid \Omega}$, we deduce that

$$
\begin{aligned}
\mathcal{F}^{\prime \prime}\left(v_{\mid \Omega}\right) & \leq \int_{\Omega} \mathbf{A} e(v): e(v) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v} \backslash J_{\bar{v}}\right)+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v} \cap \bar{\Omega}\right) \\
& \leq \int_{\Omega} \mathbf{A} e(v): e(v) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v} \backslash \Omega_{\delta}\right)+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v} \cap \bar{\Omega}\right)
\end{aligned}
$$

which settles (3.43). Passing to the limit as $\delta \searrow 0^{+}$thanks to the monotone convergence Theorem, we obtain that $\mathcal{H}^{1}\left(J_{v} \backslash \Omega_{\delta}\right) \rightarrow \mathcal{H}^{1}\left(J_{v} \backslash \Omega^{\prime}\right)=0$, hence

$$
\mathcal{F}^{\prime \prime}\left(v_{\mid \Omega}\right) \leq \int_{\Omega} \mathbf{A} e(v): e(v) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v} \cap \bar{\Omega}\right)
$$

which completes the proof of Lemma 3.12.

## 4. Convergence of minimizers

In order to investigate the approximation of minimizers for the Griffith energy, it is natural to impose boundary conditions to avoid trivial minimizers such as rigid displacements. This section is devoted to an approximation of the Griffith functional under a Dirichlet boundary condition by means of brittle damage energies, where the density $f$ reduces to $f(t)=\kappa \wedge t$ for $t \in \mathbb{R}$, and the energy $\mathcal{F}_{\varepsilon}$ corresponds to

$$
\mathcal{F}_{\varepsilon}(u)=\int_{\Omega} \frac{\kappa}{\varepsilon} \wedge \mathbf{A} e(u): e(u) d x \quad \text { for } u \in V_{\varepsilon}(\Omega)
$$

4.1. Griffith energy with Dirichlet boundary condition. In order to formulate a Dirichlet boundary condition, we need to consider a larger bounded Lipschitz open set $\Omega^{\prime}$ such that $\bar{\Omega} \subset \Omega^{\prime}$. Let $w \in W^{2, \infty}\left(\mathbb{R}^{2} ; \mathbb{R}^{2}\right)$ be a prescribed boundary displacement. Given an admissible triangulation $\mathbf{T}_{\varepsilon} \in \mathcal{T}_{\varepsilon}\left(\Omega^{\prime}\right)$ of $\Omega^{\prime}$, we define $w_{\mathbf{T}_{\varepsilon}}$ as the piecewise affine Lagrange interpolation of $w$ on $\mathbf{T}_{\varepsilon}$. Note that by standard finite element estimates (see [23, Theorem 3.1.5]),

$$
\begin{equation*}
w_{\mathbf{T}_{\varepsilon}} \in V_{\varepsilon}\left(\Omega^{\prime}\right), \quad w_{\mathbf{T}_{\varepsilon}} \rightarrow w \text { strongly in } H^{1}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right) \quad \text { and } \quad \sup _{\varepsilon>0} \mathcal{F}_{\varepsilon}\left(w_{\mathbf{T}_{\varepsilon}}\right)<+\infty \tag{4.1}
\end{equation*}
$$

We define $V_{\varepsilon}^{\text {Dir }}\left(\Omega^{\prime}\right)$ to be the set of all continuous functions $u: \Omega^{\prime} \rightarrow \mathbb{R}^{2}$ for which there exists a triangulation $\mathbf{T}_{\varepsilon} \in \mathcal{T}_{\varepsilon}\left(\Omega^{\prime}\right)$ so that $u$ is affine on each triangle $T \in \mathbf{T}_{\varepsilon}$ and $u=w_{\mathbf{T}_{\varepsilon}}$ on each triangle $T \in \mathbf{T}_{\varepsilon}$ such that $T \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right) \neq \emptyset$. We consider the following discrete functionals

$$
\mathcal{G}_{\varepsilon}: u \in L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right) \mapsto \begin{cases}\int_{\Omega} \frac{\kappa}{\varepsilon} \wedge \mathbf{A} e(u): e(u) d x & \text { if } v \in V_{\varepsilon}^{\mathrm{Dir}}\left(\Omega^{\prime}\right) \\ +\infty & \text { otherwise }\end{cases}
$$

The Griffith energy with Dirichlet boundary condition $w$ is defined, for $u \in L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$, by

$$
\mathcal{G}(u):=\left\{\begin{array}{l}
\int_{\Omega} \mathbf{A} e(u): e(u) d x \\
\quad+\kappa \sin \theta_{0}\left[\mathcal{H}^{1}\left(J_{u} \cap \Omega\right)+\mathcal{H}^{1}(\partial \Omega \cap\{u \neq w\})\right] \\
+\infty
\end{array}\right.
$$

if $\left\{\begin{array}{l}u \in G S B D^{2}\left(\Omega^{\prime}\right), \\ u=w \mathcal{L}^{2} \text {-a.e. in } \Omega^{\prime} \backslash \bar{\Omega},\end{array}\right.$
otherwise.

Note that the additional boundary term accounts for possible jumps at the boundary, where the boundary condition fails to be satisfied. In the previous expression and in the sequel, we still denote by $u$ the trace of $u_{\mid \Omega} \in G S B D^{2}(\Omega)$ on $\partial \Omega$ (see [26, Theorem 5.5]).

We will show a compactness result for sequences of dispacements $u_{\varepsilon}$ with uniformly bounded energy, with respect to the $L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$-topology of convergence in measure, from which we will deduce the following result, generalizing Theorem 1.2 to the case of Dirichlet boundary conditions.

Theorem 4.1 ( $\Gamma$-convergence under Dirichlet boundary conditions). The family $\left\{\mathcal{G}_{\varepsilon}\right\}_{\varepsilon>0}$ $\Gamma$-converges, with respect to the $L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$-topology, to the Griffith functional $\mathcal{G}$.

Considering eventually a sequence of minimizers of $\mathcal{G}_{\varepsilon}$ (see 4.11 ), we will show that, up to a subsequence and up to subtracting a sequence of piecewise rigid body motions, it converges in measure in $\Omega$ to a minimizer of $\mathcal{G}$ and the minimal value of $\mathcal{G}_{\varepsilon}$ converges to the minimal value of $\mathcal{G}$. In other words, we obtain the fundamental theorem of $\Gamma$-convergence in our specific context.

Corollary 4.2 (Convergence of minimizers). Assume further that $\Omega$ and $\Omega^{\prime}$ are connected. For each $\varepsilon>0$ small, let $u_{\varepsilon} \in V_{\varepsilon}^{\operatorname{Dir}}\left(\Omega^{\prime}\right)$ be a minimizer of $\mathcal{G}_{\varepsilon}$. Then, there exist a subsequence (not relabeled), a sequence of piecewise rigid body motions $\left\{r_{\varepsilon}\right\}_{\varepsilon>0}$ and a function $u \in G S B D^{2}\left(\Omega^{\prime}\right)$ with $u=w \mathcal{L}^{2}$-a.e. in $\Omega^{\prime} \backslash \bar{\Omega}$, such that $u_{\varepsilon}-r_{\varepsilon} \rightarrow u$ in measure in $\Omega, \mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right) \rightarrow \mathcal{G}(u)$ and $u$ is a minimizer of $\mathcal{G}$.
4.2. Compactness for sequences with uniformly bounded energy. The following result shows a compactness and lower bound estimate for any sequence with uniformly bounded energy.
Proposition 4.3. Let $\left\{\varepsilon_{k}\right\}_{k \in \mathbb{N}}$ satisfying $\varepsilon_{k} \rightarrow 0$ and let $\left\{u_{k}\right\}_{k \in \mathbb{N}} \subset L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ be such that $M:=\sup _{k} \mathcal{G}_{\varepsilon_{k}}\left(u_{k}\right)<\infty$. Then there exist a subsequence (not relabeled), a Caccioppoli partition $\mathcal{P}=\left\{P_{j}\right\}_{j \in \mathbb{N}}$ of $\Omega^{\prime}$, a sequence of piecewise rigid motions $\left\{r_{k}\right\}_{k \in \mathbb{N}}$ with

$$
r_{k}:=\sum_{j \in \mathbb{N}} r_{k}^{j} \mathbb{1}_{P_{j}}
$$

and a function $u \in G S B D^{2}\left(\Omega^{\prime}\right)$ such that $u=w \mathcal{L}^{2}$-a.e. in $\Omega^{\prime} \backslash \bar{\Omega}$,

$$
\begin{gather*}
\left|r_{k}^{i}(x)-r_{k}^{j}(x)\right| \rightarrow+\infty \quad \text { for } \mathcal{L}^{2} \text {-a.e. } x \in \Omega^{\prime}, \quad \text { for all } i \neq j  \tag{4.2}\\
u_{k}-r_{k} \rightarrow u \quad \text { in measure in } \Omega^{\prime} \tag{4.3}
\end{gather*}
$$

and

$$
\begin{equation*}
\liminf _{k \rightarrow \infty} \int_{\Omega} \frac{\kappa}{\varepsilon_{k}} \wedge \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) d x \geq \int_{\Omega} \mathbf{A} e(u): e(u) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{u} \cup \partial^{*} \mathcal{P}\right) \tag{4.4}
\end{equation*}
$$

Proof. By definition of $V_{\varepsilon_{k}}^{\text {Dir }}\left(\Omega^{\prime}\right)$, there exists an admissible triangulation $\mathbf{T}_{k} \in \mathcal{T}_{\varepsilon_{k}}\left(\Omega^{\prime}\right)$ such that $u_{k}$ is affine on each triangle $T \in \mathbf{T}_{k}$ and $u_{k}=w_{\mathbf{T}_{k}}$ on each triangle $T \in \mathbf{T}_{k}$ intersecting $\Omega^{\prime} \backslash \bar{\Omega}$. We introduce the characteristic functions

$$
\chi_{k}:=\mathbb{1}_{\left\{\mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) \geq \frac{\kappa}{\varepsilon_{k}}\right\}} \in L^{\infty}\left(\Omega^{\prime} ;\{0,1\}\right)
$$

which are constant on each triangle $T \in \mathbf{T}_{k}$. Since $u_{k}=w_{\mathbf{T}_{k}}$ on each triangle $T \in \mathbf{T}_{k}$ intersecting $\Omega^{\prime} \backslash \bar{\Omega}$ and $w \in W^{2, \infty}\left(\mathbb{R}^{2} ; \mathbb{R}^{2}\right)$, we deduce that, for $k$ large enough, $\chi_{k}=0$ in $\Omega^{\prime} \backslash \bar{\Omega}$. Thus

$$
D_{k}:=\left\{\chi_{k}=1\right\}=\bigcup_{i=1}^{N_{k}} T_{i}^{k} \subset \bar{\Omega}
$$

for some triangles $T_{i}^{k} \in \mathbf{T}_{k}$, and $\mathcal{L}^{2}\left(D_{k}\right)=\int_{\Omega} \chi_{k} d x \rightarrow 0$.
Let $v_{k}:=\left(1-\chi_{k}\right) u_{k} \in S B V^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ with $\nabla v_{k}=\left(1-\chi_{k}\right) \nabla u_{k}$ and $J_{v_{k}} \subset \bigcup_{i=1}^{N_{k}} \partial T_{i}^{k} \subset \bar{\Omega}$. Arguing as in the proof of Proposition 3.1, we infer that

$$
\sup _{k \in \mathbb{N}}\left\{\int_{\Omega^{\prime}}\left|e\left(v_{k}\right)\right|^{2} d x+\mathcal{H}^{1}\left(J_{v_{k}}\right)\right\}<\infty
$$

In view of the $G S B D^{2}$-compactness Theorem ([20, Theorem 1.1]), there exist a subsequence (not relabeled), a Caccioppoli partition $\mathcal{P}=\left\{P_{j}\right\}_{j \in \mathbb{N}}$ of $\Omega^{\prime}$, a sequence of piecewise rigid motions $\left\{\tilde{r}_{k}\right\}_{k \in \mathbb{N}}$ with

$$
\tilde{r}_{k}:=\sum_{j \in \mathbb{N}} \tilde{r}_{k}^{j} \mathbb{1}_{P_{j}},
$$

and a function $\tilde{u} \in G S B D^{2}\left(\Omega^{\prime}\right)$ such that

$$
\left\{\begin{array}{l}
\left|\tilde{r}_{k}^{i}(x)-\tilde{r}_{k}^{j}(x)\right| \rightarrow+\infty \quad \text { for } \mathcal{L}^{2} \text {-a.e. } x \in \Omega^{\prime}, \text { for all } i \neq j \\
v_{k}-\tilde{r}_{k} \rightarrow \tilde{u} \quad \text { in measure in } \Omega^{\prime} \\
e\left(v_{k}\right) \rightharpoonup e(\tilde{u}) \quad \text { weakly in } L^{2}\left(\Omega^{\prime} ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right)
\end{array}\right.
$$

Since $\mathcal{L}^{2}\left(D_{k}\right) \rightarrow 0$, we deduce that $u_{k}-\tilde{r}_{k} \rightarrow \tilde{u}$ in measure in $\Omega^{\prime}$.
For all $j \in \mathbb{N}$ such that $\mathcal{L}^{2}\left(P_{j} \cap \Omega^{\prime} \backslash \bar{\Omega}\right)>0$, the convergence in measure of $u_{k}-\tilde{r}_{k}^{j}$ to $\tilde{u}$ together with the convergence in measure of $u_{k}$ to $w$ in $P_{j} \cap \Omega^{\prime} \backslash \bar{\Omega}=: V_{j}$ ensure that $\tilde{r}_{k}^{j} \rightarrow w-\tilde{u}$ in measure in $V_{j}$. Since the space of rigid body motions is a closed finite dimensional subspace of $L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$, we can find a rigid body motion $r^{j}$ such that $r_{\mid V_{j}}^{j}=w-\tilde{u} \mathcal{L}^{2}$-a.e. in $V_{j}$. Therefore, with

$$
r:=\sum_{j \in \mathbb{N}, \mathcal{L}^{2}\left(P_{j} \cap \Omega^{\prime} \backslash \bar{\Omega}\right)>0} r^{j} \mathbb{1}_{P_{j}}
$$

the piecewise rigid body motion $r_{k}:=\tilde{r}_{k}-r$ and the function $u=\tilde{u}+r \in G S B D^{2}\left(\Omega^{\prime}\right)$ are such that

$$
\left\{\begin{array}{l}
u_{k}-r_{k} \rightarrow u \text { in measure in } \Omega^{\prime}  \tag{4.5}\\
u=w \quad \mathcal{L}^{2} \text {-a.e. in } \Omega^{\prime} \backslash \bar{\Omega} \\
e\left(v_{k}\right) \rightharpoonup e(u) \quad \text { weakly in } L^{2}\left(\Omega^{\prime} ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right) .
\end{array}\right.
$$

We are now back to prove (4.4). As in the proof of Proposition 3.3, we define the following Radon measures on $\Omega^{\prime}$

$$
\lambda_{k}:=\frac{\kappa}{\varepsilon_{k}} \wedge \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) \mathcal{L}^{2}\left\llcorner\Omega^{\prime} \in \mathcal{M}\left(\Omega^{\prime}\right)\right.
$$

Using (4.1) and the energy bound assumption on $u_{k}$, we obtain that the sequence $\left\{\lambda_{k}\right\}_{k \in \mathbb{N}}$ is uniformly bounded in $\mathcal{M}\left(\Omega^{\prime}\right)$. Thus, up to a subsequence (not relabeled), we have $\lambda_{k} \stackrel{*}{\rightharpoonup} \lambda$ weakly* in $\mathcal{M}\left(\Omega^{\prime}\right)$ for some nonnegative measure $\lambda \in \mathcal{M}\left(\Omega^{\prime}\right)$. Thanks to the lower semicontinuity of weak* convergence in $\mathcal{M}\left(\Omega^{\prime}\right)$ along open sets, we have that

$$
\begin{equation*}
\liminf _{k \rightarrow \infty} \int_{\Omega^{\prime}} \frac{\kappa}{\varepsilon_{k}} \wedge \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) d x=\liminf _{k \rightarrow \infty} \lambda_{k}\left(\Omega^{\prime}\right) \geq \lambda\left(\Omega^{\prime}\right) \tag{4.6}
\end{equation*}
$$

Recalling that $\mathcal{P}^{(1)} \cup \partial^{*} \mathcal{P}$ contains $\mathcal{H}^{1}$-almost all of $\Omega^{\prime}$, and using that the measures $\mathcal{L}^{2}\left\llcorner\Omega^{\prime}\right.$, $\mathcal{H}^{1}\left\llcorner\left(\mathcal{P}^{(1)} \cap J_{u}\right)\right.$ and $\mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.$ are mutually singular, it is enough to show that

$$
\begin{gather*}
\frac{d \lambda}{d \mathcal{L}^{2}\left\llcorner\Omega^{\prime}\right.} \geq \mathbf{A} e(u): e(u) \quad \mathcal{L}^{2} \text {-a.e. in } \Omega^{\prime}  \tag{4.7}\\
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\left(\mathcal{P}^{(1)} \cap J_{u}\right)\right.} \geq \kappa \sin \theta_{0} \quad \mathcal{H}^{1} \text {-a.e. in } \mathcal{P}^{(1)} \cap J_{u} \tag{4.8}
\end{gather*}
$$

and

$$
\begin{equation*}
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.} \geq \kappa \sin \theta_{0} \quad \mathcal{H}^{1} \text {-a.e. in } \partial^{*} \mathcal{P} \tag{4.9}
\end{equation*}
$$

Indeed, once (4.7), (4.8) and (4.9) are satisfied, it follows from the Radon-Nikodým decomposition and the Besicovitch differentiation Theorems that

$$
\lambda=\frac{d \lambda}{d \mathcal{L}^{2}\left\llcorner\Omega^{\prime}\right.} \mathcal{L}^{2}\left\llcorner\Omega^{\prime}+\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\left(\mathcal{P}^{(1)} \cap J_{u}\right)\right.} \mathcal{H}^{1}\left\llcorner\left(\mathcal{P}^{(1)} \cap J_{u}\right)+\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.} \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}+\lambda^{s},\right.\right.\right.
$$

for some nonnegative measure $\lambda^{s}$ which is singular with respect to $\mathcal{L}^{2}\left\llcorner\Omega^{\prime}, \mathcal{H}^{1}\left\llcorner\left(\mathcal{P}^{(1)} \cap J_{u}\right)\right.\right.$ and $\mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.$. Thus, after integration over $\Omega^{\prime}$ and recalling (4.6), we would get that

$$
\liminf _{k \rightarrow \infty} \int_{\Omega^{\prime}} \frac{\kappa}{\varepsilon_{k}} \wedge \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) d x \geq \int_{\Omega^{\prime}} \mathbf{A} e(u): e(u) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(\left(J_{u} \cap \mathcal{P}^{(1)}\right) \cup \partial^{*} \mathcal{P}\right)
$$

On the one hand, the convergence in $H^{1}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ of $w_{\mathbf{T}_{k}}$ to $w$ (see (4.1)) ensures that

$$
\begin{align*}
\liminf _{k \rightarrow \infty} \int_{\Omega^{\prime}} \frac{\kappa}{\varepsilon_{k}} \wedge \mathbf{A} e\left(u_{k}\right): e\left(u_{k}\right) d x & \leq \limsup _{k \rightarrow \infty} \int_{\Omega^{\prime} \backslash \Omega} \mathbf{A} e\left(w_{\mathbf{T}_{k}}\right): e\left(w_{\mathbf{T}_{k}}\right) d x+\liminf _{k \rightarrow \infty} \mathcal{G}_{\varepsilon_{k}}\left(u_{k}\right) \\
& \leq \int_{\Omega^{\prime} \backslash \Omega} \mathbf{A} e(w): e(w) d x+\liminf _{k \rightarrow \infty} \mathcal{G}_{\varepsilon_{k}}\left(u_{k}\right) \tag{4.10}
\end{align*}
$$

On the other hand, using that $u=w$ in $\Omega^{\prime} \backslash \bar{\Omega}$ and that $\mathcal{P}^{(1)} \cup \partial^{*} \mathcal{P}$ covers $\mathcal{H}^{1}$ almost every $\Omega^{\prime}$, we obtain that

$$
\begin{align*}
\int_{\Omega^{\prime}} \mathbf{A} e(u): e(u) d x & +\kappa \sin \theta_{0} \mathcal{H}^{1}\left(\left(J_{u} \cap \mathcal{P}^{(1)}\right) \cup \partial^{*} \mathcal{P}\right) \\
= & \int_{\Omega^{\prime} \backslash \Omega} \mathbf{A} e(w): e(w) d x+\int_{\Omega} \mathbf{A} e(u): e(u) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{u} \cup \partial^{*} \mathcal{P}\right) . \tag{4.11}
\end{align*}
$$

Gathering (4.10) and (4.11) leads to (4.4), which completes the proof of Proposition 4.3.
Using the last convergence in (4.5), we easily get inequality (4.7) arguing in an identical manner than in the proof of Proposition 3.4. We do not reproduce the argument. The rest of this section is devoted to the establishment of (4.8) and (4.9). We start with the lower bound inequality for the jump part of the energy in the measure theoretic interior of $\mathcal{P}$.

Proposition 4.4 (Lower bound for the jump part in $\mathcal{P}^{(1)}$ ). For $\mathcal{H}^{1}$-a.e. $x_{0} \in \mathcal{P}^{(1)} \cap J_{u}$,

$$
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\left(\mathcal{P}^{(1)} \cap J_{u}\right)\right.}\left(x_{0}\right) \geq \kappa \sin \theta_{0} .
$$

Proof. The proof is very similar to that of Proposition 3.5. We just sketch it, underlying the main differences.

Let $x_{0} \in \mathcal{P}^{(1)} \cap J_{u}$ be such that

$$
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\left(\mathcal{P}^{(1)} \cap J_{u}\right)\right.}\left(x_{0}\right)=\lim _{\varrho \searrow 0} \frac{\lambda\left(\overline{B_{\varrho}\left(x_{0}\right)}\right)}{\mathcal{H}^{1}\left(\mathcal{P}^{(1)} \cap J_{u} \cap \overline{B_{\varrho}\left(x_{0}\right)}\right)}
$$

exists and is finite, and

$$
\lim _{\varrho \searrow 0} \frac{\mathcal{H}^{1}\left(\mathcal{P}^{(1)} \cap J_{u} \cap B_{\varrho}\left(x_{0}\right)\right)}{2 \varrho}=1 .
$$

According to the Besicovitch differentiation Theorem and the countably ( $\mathcal{H}^{1}, 1$ )-rectifiability of $\mathcal{P}^{(1)} \cap J_{u}$, it follows that $\mathcal{H}^{1}$-almost every point $x_{0}$ in $\mathcal{P}^{(1)} \cap J_{u}$ fulfills these conditions.

By definition of the jump set $J_{u}$, there exist $\nu:=\nu_{u}\left(x_{0}\right) \in \mathbb{S}^{1}$ and $u^{ \pm}\left(x_{0}\right) \in \mathbb{R}^{2}$ with $u^{+}\left(x_{0}\right) \neq$ $u^{-}\left(x_{0}\right)$ such that the function $u_{x_{0}, \varrho}:=u\left(x_{0}+\varrho \cdot\right)$ converges in measure in $B:=B_{1}(0)$ to the jump function

$$
\bar{u}: y \in B \mapsto \begin{cases}u^{+}\left(x_{0}\right) & \text { if } y \cdot \nu>0 \\ u^{-}\left(x_{0}\right) & \text { if } y \cdot \nu<0\end{cases}
$$

as $\varrho \searrow 0$. As before, we consider a sequence of radii $\left\{\varrho_{j}\right\}_{j \in \mathbb{N}}$ such that $\varrho_{j} \searrow 0$ and $\lambda\left(\partial B_{\varrho_{j}}\left(x_{0}\right)\right)=0$ for all $j \in \mathbb{N}$. Arguing as in Proposition 3.5, there exists an increasing sequence $\left\{k_{j}\right\}_{j \in \mathbb{N}}$ such that $k_{j} \nearrow \infty$ as $j \rightarrow \infty$ and

$$
\left\{\begin{array}{l}
\left(u_{k_{j}}-r_{k_{j}}\right)\left(x_{0}+\varrho_{j} \cdot\right) \rightarrow \bar{u} \quad \text { in measure in } B \\
\frac{\lambda_{k_{j}}\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{2 \varrho_{j}} \rightarrow \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\left(\mathcal{P}^{(1)} \cap J_{u}\right)\right.}\left(x_{0}\right) \\
\varepsilon_{k_{j}} / \varrho_{j} \rightarrow 0, \quad \omega\left(\varepsilon_{k_{j}}\right) / \varrho_{j} \rightarrow 0
\end{array}\right.
$$

By definition of $\mathcal{P}^{(1)}$, there exists $i_{0} \in \mathbb{N}$ such that $x_{0} \in P_{i_{0}}{ }^{(1)}$. We thus infer that the function $v_{j}:=\left(u_{k_{j}}-r_{k_{j}}^{i_{0}}\right)\left(x_{0}+\varrho_{j} \cdot\right) \in H^{1}\left(B ; \mathbb{R}^{2}\right)$ converges in measure to $\bar{u}$ in $B$. Indeed, for all $\eta>0$,

$$
\begin{aligned}
\mathcal{L}^{2}\left(B \cap \left\{\left|v_{j}-\bar{u}\right|\right.\right. & >\eta\}) \\
& \leq \mathcal{L}^{2}\left(B \cap\left\{\left|\left(u_{k_{j}}-r_{k_{j}}\right)\left(x_{0}+\varrho_{j} \cdot\right)-\bar{u}\right|>\eta\right\}\right)+\mathcal{L}^{2}\left(B \backslash\left(\frac{P_{i_{0}}-x_{0}}{\varrho_{j}}\right)\right) \rightarrow 0
\end{aligned}
$$

where we used that $x_{0}$ is a point of density 1 for $P_{i_{0}}$. We are now back to an analogous situation than (3.10), since $v_{j}$ is continuous on $B$ and piecewise affine on each triangle $T \in\left(\mathbf{T}_{k_{j}}-x_{0}\right) / \varrho_{j}$. Therefore, from here the conclusion of Proposition 4.4 results from the proof of Proposition 3.5.

We next pass to the lower bound inequality for the energy on the reduced boundary of $\mathcal{P}$, which presents some non trivial adaptations of the proof of Proposition 3.5.
Proposition 4.5 (Lower bound on the reduced boundary $\partial^{*} \mathcal{P}$ ). For $\mathcal{H}^{1}$-a.e. $x_{0} \in \partial^{*} \mathcal{P}$,

$$
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.}\left(x_{0}\right) \geq \kappa \sin \theta_{0}
$$

The rest of this subsection is devoted to prove Proposition 4.5, with essentially the same structure than the proof of Proposition 3.5.

Blow-up. Let $x_{0} \in \partial^{*} \mathcal{P}$ be such that

$$
\begin{gathered}
x_{0} \in \partial^{*} P_{i_{0}} \cap \partial^{*} P_{j_{0}} \quad \text { for some } i_{0} \neq j_{0}, \\
\nu:=\nu_{P_{i_{0}}}\left(x_{0}\right)=-\nu_{P_{j_{0}}}\left(x_{0}\right) \quad \text { where } \nu_{P_{k}}\left(x_{0}\right):=\lim _{\varrho \searrow 0} \frac{D \mathbb{1}_{P_{k}}\left(B_{\varrho}\left(x_{0}\right)\right)}{\left|D \mathbb{1}_{P_{k}}\right|\left(B_{\varrho}\left(x_{0}\right)\right)} \quad \text { for } k \in\left\{i_{0}, j_{0}\right\}, \\
\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.}\left(x_{0}\right)= \\
\lim _{\varrho \searrow 0} \frac{\lambda\left(\overline{B_{\varrho}\left(x_{0}\right)}\right)}{\mathcal{H}^{1}\left(\partial^{*} \mathcal{P} \cap \overline{B_{\varrho}\left(x_{0}\right)}\right)} \quad \text { exists and is finite, } \\
\\
\lim _{\varrho \searrow 0} \frac{\mathcal{H}^{1}\left(\partial^{*} \mathcal{P} \cap B_{\varrho}\left(x_{0}\right)\right)}{2 \varrho}=1,
\end{gathered}
$$

and there exist traces $u^{ \pm}\left(x_{0}\right) \in \mathbb{R}^{2}$ such that the function $u_{x_{0}, \varrho}:=u\left(x_{0}+\varrho \cdot\right)$ converges in measure in $B:=B_{1}(0)$ to

$$
y \in B \mapsto \bar{u}(y):=\left\{\begin{array}{ll}
u^{+}\left(x_{0}\right) & \text { if } y \cdot \nu>0, \\
u^{-}\left(x_{0}\right) & \text { if } y \cdot \nu<0,
\end{array} \quad \text { as } \varrho \searrow 0\right.
$$

The previous properties turn out to be satisfied for $\mathcal{H}^{1}$-a.e. $x_{0} \in \partial^{*} \mathcal{P}$. This is a consequence of the countably $\left(\mathcal{H}^{1}, 1\right)$-rectifiability of that set, the Besicovitch differentiation Theorem, the fact that $\mathcal{P}^{(1)} \cup \bigcup_{i \neq j}\left(\partial^{*} P_{i} \cap \partial^{*} P_{j}\right)$ covers $\mathcal{H}^{1}$ almost all of $\Omega^{\prime}([3$, Theorem 4.17]), and the existence of traces on ( $\mathcal{H}^{1}, 1$ )-rectifiable sets (see [26, Theorem 5.2] in the case of 1-dimensional $C^{1}$ submanifolds which may be extended to countably ( $\left.\mathcal{H}^{1}, 1\right)$-rectifiable sets arguing as in [5, Proposition 4.1]).

To simplify notation, let us denote by $P^{+}:=P_{i_{0}}$ and $P^{-}:=P_{j_{0}}$. According to De Giorgi's Theorem (see [3, Theorem 3.59]) we infer that

$$
\begin{equation*}
\mathbb{1}_{\frac{P^{ \pm}-x_{0}}{\varrho}} \rightarrow \mathbb{1}_{H^{ \pm}} \text {strongly in } L^{1}(B) \text { as } \varrho \searrow 0 \tag{4.12}
\end{equation*}
$$

where $H^{ \pm} \subset \mathbb{R}^{2}$ denote the halfspaces orthogonal to $\nu$ and containing $\pm \nu$. With these notation, we have that

$$
\bar{u}=u^{+}\left(x_{0}\right) \mathbb{1}_{H^{+} \cap B}+u^{-}\left(x_{0}\right) \mathbb{1}_{H^{-} \cap B} .
$$

Note also that contrary to Proposition 3.5 where jump points were considered, it might be the case that $u^{+}\left(x_{0}\right)=u^{-}\left(x_{0}\right)$, i.e. that $\bar{u}$ is constant.

Extraction of diagonal subsequences. As before, we consider a sequence of radii $\left\{\varrho_{j}\right\}_{j \in \mathbb{N}}$ such that $\varrho_{j} \searrow 0$ and $\lambda\left(\partial B_{\varrho_{j}}\left(x_{0}\right)\right)=0$ for all $j \in \mathbb{N}$. By our choice of $x_{0}$, (4.2) and (4.3), with $r_{k}^{ \pm}:=r_{k \mid P^{ \pm}}$, we have :

$$
\left\{\begin{array}{l}
\lim _{j \rightarrow \infty} \lim _{k \rightarrow \infty}\left(u_{k}-r_{k}\right)\left(x_{0}+\varrho_{j} \cdot\right)=\lim _{j \rightarrow \infty} u_{x_{0}, \varrho_{j}}=\bar{u} \quad \text { in measure in } B, \\
\lim _{j \rightarrow \infty} \lim _{k \rightarrow \infty} \arctan \left|r_{k}^{+}-r_{k}^{-}\right|\left(x_{0}+\varrho_{j} \cdot\right)=\frac{\pi}{2} \quad \text { in measure in } B \\
\lim _{j \rightarrow \infty} \lim _{k \rightarrow \infty} \frac{\lambda_{k}\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{2 \varrho_{j}}=\lim _{j \rightarrow \infty} \frac{\lambda\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{2 \varrho_{j}}=\frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.}\left(x_{0}\right) \\
\lim _{j \rightarrow \infty} \lim _{k \rightarrow \infty} \frac{\varepsilon_{k}}{\varrho_{j}}=\lim _{j \rightarrow \infty} \lim _{k \rightarrow \infty} \frac{\omega\left(\varepsilon_{k}\right)}{\varrho_{j}}=0
\end{array}\right.
$$

We can thus find an increasing sequence $\left\{k_{j}\right\}_{j \in \mathbb{N}}$ such that $k_{j} \nearrow \infty$ as $j \rightarrow \infty$ and

$$
\left\{\begin{array}{l}
\left(u_{k_{j}}-r_{k_{j}}\right)\left(x_{0}+\varrho_{j} \cdot\right) \rightarrow \bar{u} \quad \text { in measure in } B  \tag{4.13a}\\
\arctan \left|r_{k_{j}}^{+}-r_{k_{j}}^{-}\right|\left(x_{0}+\varrho_{j} \cdot\right) \rightarrow \frac{\pi}{2} \quad \text { in measure in } B \\
\frac{\lambda_{k_{j}}\left(B_{\varrho_{j}}\left(x_{0}\right)\right)}{2 \varrho_{j}} \rightarrow \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.}\left(x_{0}\right) \\
\frac{\varepsilon_{k_{j}}}{\varrho_{j}} \rightarrow 0, \quad \frac{\omega\left(\varepsilon_{k_{j}}\right)}{\varrho_{j}} \rightarrow 0
\end{array}\right.
$$

Let $v_{j}:=u_{k_{j}}\left(x_{0}+\varrho_{j} \cdot\right), r_{j}^{ \pm}:=r_{k_{j}}^{ \pm}\left(x_{0}+\varrho_{j} \cdot\right)$ and $r_{j}:=r_{j}^{+} \mathbb{1}_{H^{+} \cap B}+r_{j}^{-} \mathbb{1}_{H^{-} \cap B}$. By (4.12) and (4.13), we have for all $\eta>0$,

$$
\begin{aligned}
\mathcal{L}^{2}\left(\left\{\mid v_{j}-r_{j}^{ \pm}-\right.\right. & \left.\left.u^{ \pm}\left(x_{0}\right) \mid>\eta\right\} \cap B \cap H^{ \pm}\right) \\
& \leq \mathcal{L}^{2}\left(B \cap H^{ \pm} \backslash\left(P^{ \pm}-x_{0}\right) / \varrho_{j}\right)+\mathcal{L}^{2}\left(\left\{\left|v_{j}-r_{k_{j}}\left(x_{0}+\varrho_{j} \cdot\right)-\bar{u}\right|>\eta\right\} \cap B\right) \rightarrow 0
\end{aligned}
$$

Thus, up to a subsequence

$$
\left\{\begin{array}{l}
v_{j}-r_{j}^{ \pm} \rightarrow u^{ \pm}\left(x_{0}\right) \quad \text { in measure in } B \cap H^{ \pm}  \tag{4.14}\\
\left|r_{j}^{+}-r_{j}^{-}\right| \rightarrow+\infty \quad \mathcal{L}^{2} \text {-a.e. in } B
\end{array}\right.
$$

Selection of a slicing direction. According to [20, Lemma 2.8], there exist an $\mathcal{H}^{1}$-negligible set $N \subset B^{\nu}$ and a countable dense subset $D$ of $\mathbb{S}^{1}$ such that for all $\xi \in D$ and all $y \in B^{\nu} \backslash N$,

$$
\left|\left(r_{j}^{+}-r_{j}^{-}\right)(y) \cdot \xi\right| \rightarrow+\infty
$$

as $j \rightarrow+\infty$. Note that $\nabla r_{j}^{ \pm} \xi \cdot \xi=0$, so that the quantity $t \mapsto\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}(t)=\left(r_{j}^{+}-r_{j}^{-}\right)(y) \cdot \xi$ is independent of $t \in B_{y}^{\xi}$. Thus, for all $y \in B^{\nu} \backslash N$, we have

$$
\begin{equation*}
\arctan \left|\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}\right| \rightarrow \frac{\pi}{2} \tag{4.15}
\end{equation*}
$$

For any $\eta>0$, let $\xi \in \mathbb{S}^{1} \cap D$ be such that

$$
\begin{equation*}
|\nu-\xi| \leq \eta, \quad \nu \cdot \xi \geq \frac{1}{2}, \quad\left|\nu \cdot \xi^{\perp}\right| \leq \eta \tag{4.16}
\end{equation*}
$$

As in the proof of Proposition 3.5, using a change of variables and the ellipticity property (1.5) of $\mathbf{A}$ and (4.13d), we get

$$
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.}\left(x_{0}\right) \geq \limsup _{j \rightarrow \infty} \varrho_{j} \int_{B} \frac{\kappa}{\varepsilon_{k_{j}}} \wedge \frac{\alpha}{\varrho_{j}^{2}}\left|e\left(v_{j}\right) \xi \cdot \xi\right|^{2} d y
$$

so that, introducing the following characteristic functions,

$$
\chi_{j}:=\mathbb{1}_{\left\{\frac{\alpha}{e_{j}^{2}}\left|e\left(v_{j}\right) \xi \cdot \xi\right|^{2} \geq \frac{\kappa}{\varepsilon_{k_{j}}}\right\}} \in L^{\infty}(B ;\{0,1\}),
$$

we obtain that

$$
\begin{equation*}
2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.}\left(x_{0}\right) \geq \limsup _{j \rightarrow \infty}\left\{\frac{\alpha}{\varrho_{j}} \int_{B}\left(1-\chi_{j}\right)\left|e\left(v_{j}\right) \xi \cdot \xi\right|^{2} d y+\frac{\kappa \varrho_{j}}{\varepsilon_{k_{j}}} \int_{B} \chi_{j} d y\right\} \tag{4.17}
\end{equation*}
$$

We define the translated and rescaled triangulations:

$$
\mathbf{T}^{x_{0}, j}:=\frac{1}{\varrho_{j}}\left(\mathbf{T}^{k_{j}}-x_{0}\right), \quad \mathbf{T}_{b}^{x_{0}, j}:=\left\{T \in \mathbf{T}^{x_{0}, j}: \frac{\alpha}{\varrho_{j}}\left|e\left(v_{j}\right)_{\mid T} \xi \cdot \xi\right|^{2} \geq \frac{\kappa \varrho_{j}}{\varepsilon_{k_{j}}}\right\}
$$

and the family of triangles which intersect $\overline{B_{1-\frac{\eta}{4}}}$ :

$$
\mathbf{T}_{b, i n t}^{x_{0}, j}:=\left\{T \in \mathbf{T}_{b}^{x_{0}, j}: T \cap \overline{B_{1-\frac{\eta}{4}}} \neq \emptyset\right\}
$$

Note that $v_{j}-r_{j}^{ \pm}$is affine and $\chi_{j}$ is constant on each $T \in \mathbf{T}^{x_{0}, j}$, and that (4.13d) ensures that for $j \in \mathbb{N}$ large enough (depending on $\eta$ ), each $T \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ is contained in $B$. As in (3.17), for all $y \in\left(B_{1-\frac{\eta}{4}}\right)^{\nu}$, we denote by $a(y)$ and $b(y)$ the end points of the section passing through $y$ in the direction $\xi$ (see the Figure 1) in such a way that $\left(B_{1-\frac{\eta}{4}}\right)_{y}^{\xi}=(a(y), b(y))$. We also recall that $L_{\eta}$ is defined as in (3.18) and satisfies $0<L_{\eta} \leq|a(y)|,|b(y)| \leq 2$.

Using (4.14), (4.15), (4.17), Fubini's and Egoroff's Theorem, we can adapt the proof of Lemma 3.6 to show the following result.

Lemma 4.6. For all $\eta>0$, there exist a subset $Z \subset B^{\nu}$ containing $N$ with $\mathcal{H}^{1}(Z) \leq \eta$, and a subsequence (not relabeled) such that the following property holds : for all $\gamma>0$, there exists $j_{0}=j_{0}(\gamma) \in \mathbb{N}$ such that for all $y \in B^{\nu} \backslash Z$ and all $j \geq j_{0}$,

$$
\int_{B_{y}^{\xi}}\left(\chi_{j}\right)_{y}^{\xi} d t \leq \gamma, \quad \int_{B_{y}^{\xi}}\left(1-\left(\chi_{j}\right)_{y}^{\xi}\right)\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \gamma^{2}
$$

and

$$
\int_{B_{y}^{\xi} \cap \mathbb{R}^{ \pm}} 1 \wedge\left|\left(v_{j}-r_{j}^{ \pm}\right)_{y}^{\xi}-u^{ \pm}\left(x_{0}\right) \cdot \xi\right| d t \leq \frac{\gamma}{2}, \quad\left|\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}\right| \geq\left|[u]\left(x_{0}\right) \cdot \xi\right|+1 .
$$

As in the proof of Proposition 3.5, we next show that, for some subset $Z^{\prime} \subset\left(B_{1-\frac{\eta}{2}}\right)^{\nu}$ of arbitrarily small $\mathcal{H}^{1}$ measure, and along a subsequence (only depending on $\eta$ ), all the sections in the direction $\xi$ passing through $\left(B_{1-\frac{\eta}{2}}\right)^{\nu} \backslash Z^{\prime}$ must cross at least one triangle $T \in \mathbf{T}_{b, \text { int }}^{x_{0}, j}$ contained in $B$.
Lemma 4.7. For all $\eta>0$, there exist a subset $Z^{\prime} \subset B^{\nu}$ containing $Z$ with $\mathcal{H}^{1}\left(Z^{\prime}\right) \leq \eta$, and a subsequence (not relabeled) such that the following property holds : for all $y \in\left(B_{1-\frac{\eta}{2}}\right)^{\nu} \backslash Z^{\prime}$ and all $j \in \mathbb{N}$, there exists a triangle $T=T(y, j) \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ such that $(T \cap B)_{y}^{\xi} \neq \emptyset$.
Proof. Let $Z$ be the exceptional set given by Lemma 4.6. We first show that there exists an increasing mapping $\phi: \mathbb{N} \rightarrow \mathbb{N}$ such that : for all $y \in\left(B_{1-\frac{\eta}{2}}\right)^{\nu} \backslash Z$ and all $j \in \mathbb{N}$, there exists a triangle $T=T(y, \phi(j)) \in \mathbf{T}_{b, \text { int }}^{x_{0}, \phi(j)}$ such that $(T \cap B)_{y}^{\xi} \neq \emptyset$. Suppose by contradiction that such is not the case, and define

$$
\gamma_{1}^{*}:=L_{\eta}>0, \quad \gamma_{2}^{*}:=\frac{L_{\eta}}{1+2 L_{\eta}}>0 \quad \text { and } \quad \gamma^{*}=\gamma^{*}(\eta):=\frac{\gamma_{1}^{*} \wedge \gamma_{2}^{*}}{4}>0 .
$$

Thanks to Lemma 4.6, there exists a rank $j^{*}=j^{*}\left(\gamma^{*}\right) \in \mathbb{N}$ such that for all $y \in B^{\nu} \backslash Z$ and all $j \geq j^{*},\left|\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}\right| \geq\left|[u]\left(x_{0}\right) \cdot \xi\right|+1$ and

$$
\int_{B_{y}^{\xi}}\left(1-\left(\chi_{j}\right)_{y}^{\xi}\right)\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \gamma^{* 2}, \quad \int_{B_{乡 勹}^{\xi} \cap \mathbb{R}^{ \pm}} 1 \wedge\left|\left(v_{j}-r_{j}^{ \pm}\right)_{y}^{\xi}-u^{ \pm}\left(x_{0}\right) \cdot \xi\right| d t \leq \frac{\gamma^{*}}{2} .
$$

As in Lemma 3.7, we consider the extraction $\phi: j \in \mathbb{N} \longmapsto j+j^{*} \in \mathbb{N}$. By assumption, there exist $y \in\left(B_{1-\frac{\eta}{2}}\right)^{\nu} \backslash Z$ and $j \in \mathbb{N}$ such that $(T \cap B)_{y}^{\xi}=\emptyset$ for all $T \in \mathbf{T}_{b, \text { int }}^{x_{0}, j+j^{*}}$, implying that $\left(\chi_{j+j^{*}}\right)_{y}^{\xi} \equiv 0$ on $(a(y), b(y)),\left|\left(r_{j+j^{*}}^{+}-r_{j+j^{*}}^{-}\right)_{y}^{\xi}\right| \geq\left|[u]\left(x_{0}\right) \cdot \xi\right|+1$ and

$$
\int_{a(y)}^{b(y)}\left|\left(\left(v_{j+j^{*}}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \gamma^{* 2}, \quad \int_{[a(y), b(y)] \cap \mathbb{R}^{ \pm}} 1 \wedge\left|\left(v_{j+j^{*}}-r_{j+j^{*}}^{ \pm}\right)_{y}^{\xi}-u^{ \pm}\left(x_{0}\right) \cdot \xi\right| d t \leq \frac{\gamma^{*}}{2}
$$

since $\phi(j)=j+j^{*} \geq j^{*}$. By continuity of $\left(v_{j+j^{*}}-r_{j+j^{*}}^{ \pm}\right)_{y}^{\xi}$ on the compact sets $[a(y), b(y)] \cap \mathbb{R}^{ \pm}$, there exist two points

$$
t^{ \pm} \in \underset{[a(y), b(y)] \cap \mathbb{R}^{ \pm}}{\arg \min }\left(1 \wedge\left|\left(v_{j+j^{*}}-r_{j+j^{*}}^{ \pm}\right)_{y}^{\xi}-u^{ \pm}\left(x_{0}\right) \cdot \xi\right|\right) .
$$

Hence, recalling (3.18)

$$
\begin{aligned}
\frac{\gamma^{*}}{L_{\eta}} & \geq 1 \wedge\left|\left(v_{j+j^{*}}-r_{j+j^{*}}^{-}\right)_{y}^{\xi}\left(t^{-}\right)-u^{-}\left(x_{0}\right) \cdot \xi\right|+1 \wedge\left|\left(v_{j+j^{*}}-r_{j+j^{*}}^{+}\right)_{y}^{\xi}\left(t^{+}\right)-u^{+}\left(x_{0}\right) \cdot \xi\right| \\
& \geq 1 \wedge\left(\left|[u]\left(x_{0}\right) \cdot \xi+\left(r_{j+j^{*}}^{+}-r_{j+j^{*}}^{-}\right)_{y}^{\xi}\right|-\left|\int_{t^{-}}^{t^{+}}\left(\left(v_{j+j^{*}}\right)_{y}^{\xi}\right)^{\prime}(t) d t\right|\right) \geq 1-2 \gamma^{*},
\end{aligned}
$$

which is impossible thanks to of our choice of $\gamma^{*}$. We conclude the proof of Lemma 4.7 in the same way as for Lemma 3.7.

As a consequence of Lemma 4.7, introducing the family of triangles

$$
\mathscr{F}_{j}:=\left\{T \in \mathbf{T}_{b, \text { int }}^{x_{0}, j}: \text { there exists } y \in\left(B_{1-\frac{n}{2}}\right)^{\nu} \text { such that }(T \circ \cap B)_{y}^{\xi} \neq \emptyset\right\}
$$

for all $j \in \mathbb{N}$, it is possible to obtain a too low lower bound, roughly speaking because Lemma 4.7 does not exhibit enough triangles in $\mathbf{T}_{b, i n t}^{x_{0}, j}$, as explained after (3.19). Therefore, we need to establish that many lines $B_{y}^{\xi}$ parallel to $\xi$ and passing through $B^{\nu}$ must actually intersect at least
two triangles of the collection $\mathbf{T}_{b, i n t}^{x_{0}, j}$. To this aim, we show that the set of points $y \in B^{\nu}$ such that $B_{y}^{\xi}$ intersects exactly one triangle $T$ in the collection $\mathbf{T}_{b, i n t}^{x_{0}, j}$, has arbitrarily small $\mathcal{H}^{1}$ measure.

Lemma 4.8. For all $\eta>0$, there exist constants $C_{*}=C_{*}(\eta)>0, \gamma_{*}=\gamma_{*}(\eta)>0$ and a subset $Z_{*}=Z_{*}(\eta) \subset B^{\nu}$ containing $Z^{\prime}$ and satisfying $\mathcal{H}^{1}\left(Z_{*}\right) \leq 3 \eta$ such that the following property holds: for all $0<\gamma<\gamma_{*}$, there exists $j(\gamma) \in \mathbb{N}$ such that for all $j \geq j(\gamma)$, the set

$$
Y_{j}:=\left\{y \in\left(B_{1-\frac{\eta}{2}}\right)^{\nu} \backslash Z^{\prime}: \text { there exists a unique } T \in \mathbf{T}_{b, \text { int }}^{x_{0}, j} \text { such that }(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset\right\}
$$

satisfies

$$
\mathcal{H}^{1}\left(Y_{j} \backslash Z_{*}\right) \leq C_{*} \gamma
$$

Proof of Lemma 4.8. We follow the same three steps structuring the proof of Lemma 3.9.
Step 1. We start by showing that for $j$ large enough and for many points $y \in Y_{j}$, the only triangle $T$ in $\mathbf{T}_{b, i n t}^{x_{0}, j}$ crossing $B_{y}^{\xi}$ is drawing closer to the diameter $B^{\nu}$.

For all $j \in \mathbb{N}$ and all $y \in Y_{j}$, let $T_{j}(y) \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ be the unique triangle such that $\left(\stackrel{\circ}{T}_{j}(y) \cap B\right)_{y}^{\xi} \neq \emptyset$. We keep the notation (3.21) for the end points $a_{j}(y)$ and $b_{j}(y)$ of the section in the direction $\xi$ passing through $y$ inside $T_{j}(y)$ (see the Figure 3 ). Let us show that

$$
f_{j}(y):=\left(\left|a_{j}(y)\right|+\left|b_{j}(y)\right|\right) \mathbb{1}_{Y_{j}}(y) \rightarrow 0 \quad \text { for all } y \in\left(B_{1-\frac{\eta}{2}}\right)^{\nu} \backslash Z^{\prime}
$$

Let $y \in\left(B_{1-\frac{\eta}{2}}\right)^{\nu} \backslash Z^{\prime}$. Assume by contradiction that $\ell:=\lim \sup _{j} f_{j}(y)>0$ and extract a subsequence (depending on $y$, not relabeled) such that $f_{j}(y) \rightarrow \ell$. Then, there exists a rank $j_{0} \in \mathbb{N}$ such that $y \in Y_{j}$ for all $j \geq j_{0}$. Moreover, according to Lemma 4.6 and setting $I_{j}(y):=$ $(a(y), b(y)) \backslash\left(a_{j}(y), b_{j}(y)\right) \subset B_{y}^{\xi}$, we infer that

$$
\begin{equation*}
\left|b_{j}(y)-a_{j}(y)\right|+\int_{I_{j}(y)}\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t+\int_{[a(y), b(y)] \cap \mathbb{R}^{ \pm}} 1 \wedge\left|\left(v_{j}-r_{j}^{ \pm}\right)_{y}^{\xi}-u^{ \pm}\left(x_{0}\right) \cdot \xi\right| d t \rightarrow 0 \tag{4.18}
\end{equation*}
$$

Up to another subsequence (still not relabeled), it ensures that $a_{j}(y), b_{j}(y) \rightarrow m$ for some $m \in$ $[a(y), b(y)]$. Thus, for all $\tau>0, I_{\tau}:=(a(y), m-\tau) \cup(m+\tau, b(y)) \subset I_{j}(y)$ for $j \in \mathbb{N}$ sufficiently large. In particular, we deduce that $m-\tau \leq 0$. Indeed, assuming that $m-\tau>0$, by continuity of $\left(v_{j}-r_{j}^{+}\right)_{y}^{\xi}$ on $(0, m-\tau)$ and of $\left(v_{j}-r_{j}^{-}\right)_{y}^{\xi}$ on $(a(y), 0)$, there exist

$$
t_{j}^{+} \in \underset{(0, m-\tau)}{\arg \min } 1 \wedge\left|\left(v_{j}-r_{j}^{+}\right)_{y}^{\xi}-u^{+}\left(x_{0}\right) \cdot \xi\right| \quad \text { and } \quad t_{j}^{-} \in \underset{(a(y), 0)}{\arg \min } 1 \wedge\left|\left(v_{j}-r_{j}^{-}\right)_{y}^{\xi}-u^{-}\left(x_{0}\right) \cdot \xi\right|
$$

which satisfy

$$
\begin{aligned}
& 1 \wedge\left|\left(v_{j}-r_{j}^{+}\right)_{y}^{\xi}\left(t_{j}^{+}\right)-u^{+}\left(x_{0}\right) \cdot \xi\right|+1 \\
& \wedge\left|\left(v_{j}-r_{j}^{-}\right)_{y}^{\xi}\left(t_{j}^{-}\right)-u^{-}\left(x_{0}\right) \cdot \xi\right| \\
& \geq \wedge\left(\left|[u]\left(x_{0}\right) \cdot \xi+\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}\right|-2 \sqrt{\int_{I_{j}(y)}\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t}\right) \rightarrow 1
\end{aligned}
$$

according to (4.18) and (4.15). However,

$$
\begin{aligned}
& 1 \wedge\left|\left(v_{j}-r_{j}^{+}\right)_{y}^{\xi}\left(t_{j}^{+}\right)-u^{+}\left(x_{0}\right) \cdot \xi\right|+1 \wedge\left|\left(v_{j}-r_{j}^{-}\right)_{y}^{\xi}\left(t_{j}^{-}\right)-u^{-}\left(x_{0}\right) \cdot \xi\right| \\
& \leq \frac{1}{L_{\eta}} \int_{a(y)}^{0} 1 \wedge\left|\left(v_{j}-r_{j}^{-}\right)_{y}^{\xi}-u^{-}\left(x_{0}\right) \cdot \xi\right| d t+\frac{1}{m-\tau} \int_{0}^{m-\tau} 1 \wedge\left|\left(v_{j}-r_{j}^{+}\right)_{y}^{\xi}-u^{+}\left(x_{0}\right) \cdot \xi\right| d t \rightarrow 0
\end{aligned}
$$

according again to (4.18), which leads to a contradiction. We similarly show that $m+\tau \geq 0$, leading to $|m| \leq \tau$. Taking the limit as $\tau \rightarrow 0^{+}$, we obtain that $m=0$ which is against $\ell>0$.

Therefore, owing to Lemma 4.6 and Egoroff's Theorem, we can find a set $Z_{*}^{1} \subset B^{\nu}$ containing $Z^{\prime}$ with $\mathcal{H}^{1}\left(Z_{*}^{1}\right) \leq 2 \eta$ such that for all $\gamma>0$, there exists $j_{0}(\gamma) \in \mathbb{N}$ satisfying

$$
\left\{\begin{array}{l}
\int_{B_{y}^{\xi}}\left(1-\left(\chi_{j}\right)_{y}^{\xi}\right)\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|^{2} d t \leq \gamma^{2}  \tag{4.19}\\
\int_{B_{y}^{\xi} \cap \mathbb{R}^{ \pm}} 1 \wedge\left|\left(v_{j}-r_{j}^{ \pm}\right)_{y}^{\xi}-u^{ \pm}\left(x_{0}\right) \cdot \xi\right| d t \leq \frac{\gamma}{2} \\
\left(\left|a_{j}(y)\right|+\left|b_{j}(y)\right|\right) \mathbb{1}_{Y_{j}}(y) \leq \gamma
\end{array}\right.
$$

for all $y \in\left(B_{1-\frac{\eta}{2}}\right)^{\nu} \backslash Z_{*}^{1}$ and all $j \geq j_{0}(\gamma)$.
Step 2. In this step, we show that for many points $y \in Y_{j}$, the variation of $\left(v_{j}-r_{j}\right)_{y}^{\xi}$ inside the only triangle $T$ in $\mathbf{T}_{b, i n t}^{x_{0}, j}$ which is crossed by $B_{y}^{\xi}$, is close to that of $\bar{u}_{y}^{\xi}$. We consider the constants $C_{\eta}$ and $\gamma_{*}$ defined by

$$
C_{\eta}:=8\left(1+\frac{1}{L_{\eta}}\right)>0, \quad \gamma_{*}=\gamma_{*}(\eta):=\frac{1}{2} \min \left(1, \frac{1}{C_{\eta}}, L_{\eta}\right)>0
$$

Let us show that for all $0<\gamma<\gamma_{*}$, there exists a rank $j_{1}(\gamma) \in \mathbb{N}$ such that for all $j \geq j_{1}(\gamma)$ and all $y \in Y_{j} \backslash Z_{*}^{1}$,

$$
\begin{equation*}
\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)-\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}-[u]\left(x_{0}\right) \cdot \xi\right| \leq C_{\eta} \gamma \tag{4.20}
\end{equation*}
$$

Fix $0<\gamma<\gamma_{*}$ and let $j_{0}(\gamma) \in \mathbb{N}$ be such that (4.19) holds for all $j \geq j_{0}(\gamma)$ and all $y \in Y_{j} \backslash Z_{*}^{1}$. As before, the choice of $\gamma_{*}$ leads to $a(y)<a_{j}(y) \leq b_{j}(y)<b(y)$, and we have

$$
\begin{aligned}
\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)-\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}-[u]\left(x_{0}\right) \cdot \xi\right| \leq & \left|\left(v_{j}\right)_{y}^{\xi}\left(0 \vee b_{j}(y)\right)-\left(r_{j}^{+}\right)_{y}^{\xi}-u^{+}\left(x_{0}\right) \cdot \xi\right| \\
& +\left|\left(v_{j}\right)_{y}^{\xi}\left(0 \vee b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)\right| \\
& +\left|\left(v_{j}\right)_{y}^{\xi}\left(0 \wedge a_{j}(y)\right)-\left(r_{j}^{-}\right)_{y}^{\xi}-u^{-}\left(x_{0}\right) \cdot \xi\right| \\
& +\left|\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(0 \wedge a_{j}(y)\right)\right| \\
= & : J_{1}+J_{2}+J_{3}+J_{4} .
\end{aligned}
$$

The terms $J_{2}$ and $J_{4}$ are controlled by $\gamma$ thanks to the Cauchy-Schwarz inequality and (4.19), as shown in the proof of (3.26). As for the terms $J_{1}$ and $J_{3}$, they are being treated almost exactly as in the proof of (3.26), except that the function $z_{j}$ is replaced by $z_{j}^{+}:=1 \wedge \mid\left(v_{j}-\right.$ $\left.r_{j}^{+}\right)_{y}^{\xi}-u^{+}\left(x_{0}\right) \cdot \xi \mid \in H^{1}\left(B_{y}^{\xi}\right)$ for $J_{1}$, and $z_{j}^{-}:=1 \wedge\left|\left(v_{j}-r_{j}^{-}\right)_{y}^{\xi}-u^{-}\left(x_{0}\right) \cdot \xi\right| \in H^{1}\left(B_{y}^{\xi}\right)$ for $J_{3}$, both satisfying $\left|\left(z_{j}^{ \pm}\right)^{\prime}\right| \leq\left|\left(\left(v_{j}\right)_{y}^{\xi}\right)^{\prime}\right|$. Using the Sobolev embedding, (4.19) and our definition of $C_{\eta}$ and $\gamma^{*}$, it leads to $z_{j}^{+}\left(0 \vee b_{j}(y)\right) \leq\left(\sqrt{2}+\frac{2}{L_{\eta}}\right) \gamma \leq C_{\eta} \gamma_{*}<1$, which eventually shows that $J_{1}=z_{j}^{+}\left(0 \vee b_{j}(y)\right) \leq\left(\sqrt{2}+\frac{2}{L_{\eta}}\right) \gamma$. Similarly, $J_{3} \leq\left(\sqrt{2}+\frac{2}{L_{\eta}}\right) \gamma$, which concludes the validity of (4.20).

Step 3. We now show that, after enlarging slightly the set $Z_{*}^{1}$ into a set $Z_{*} \subset B^{\nu}$ with $\mathcal{H}^{1}\left(Z_{*}\right) \leq 3 \eta$, it is possible to include $Y_{j} \backslash Z_{*}$ inside a finite union of arbitrarily small segments contained in $B^{\nu}$ (see Figure 6).

By definition of rigid body motions, there exist skew symmetric matrices $M_{j} \in \mathbb{M}_{\text {skew }}^{2 \times 2}$ and vectors $m_{j} \in \mathbb{R}^{2}$ such that

$$
r_{j}^{+}-r_{j}^{-}=M_{j} \mathrm{id}_{\mathbb{R}^{2}}+m_{j}
$$

for all $j \in \mathbb{N}$. Recalling (4.15), we get that for all $y \in B^{\nu} \backslash Z_{*}^{1} \subset B^{\nu} \backslash N$,

$$
\left|\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}\right|=\left|\left(M_{j}^{T} \xi\right) \cdot y+m_{j} \cdot \xi\right| \rightarrow+\infty
$$

as $j \rightarrow+\infty$. In particular, setting $\alpha_{j}:=M_{j}^{T} \xi \in \mathbb{R}^{2}$ and $\beta_{j}:=m_{j} \cdot \xi \in \mathbb{R}$, we get that $\mu_{j}:=$ $\left|\alpha_{j}\right|+\left|\beta_{j}\right| \rightarrow+\infty$ as $j \rightarrow+\infty$. Hence, up to a subsequence (depending only on $\xi$, not relabeled), there exist $\alpha \in \mathbb{R}^{2}$ and $\beta \in \mathbb{R}$ such that

$$
\frac{\alpha_{j}}{\mu_{j}} \rightarrow \alpha, \quad \frac{\beta_{j}}{\mu_{j}} \rightarrow \beta, \quad \text { and } \quad|\alpha|+|\beta|=1
$$

In particular,

$$
\begin{equation*}
\frac{1}{\mu_{j}}\left(\alpha_{j} \cdot \operatorname{id}_{\mathbb{R}^{2}}+\beta_{\mathrm{j}}\right) \rightarrow \alpha \cdot \mathrm{id}_{\mathbb{R}^{2}}+\beta \text { uniformly on } B \tag{4.21}
\end{equation*}
$$

Notice that the affine line $\Delta:=\left\{z \in \mathbb{R}^{2}: \alpha \cdot z+\beta=0\right\}$ cannot coincide with $\Pi_{\nu}$. Indeed, if such would be the case, it would entail that $\beta=0$ and $\alpha= \pm \nu \in \mathbb{S}^{1}$. Yet, $M_{j} \in \mathbb{M}_{\text {skew }}^{2 \times 2}$ being skew symmetric, we would obtain that

$$
0=\frac{M_{j} \xi \cdot \xi}{\mu_{j}}=\frac{\alpha_{j}}{\mu_{j}} \cdot \xi \rightarrow \alpha \cdot \xi= \pm \nu \cdot \xi
$$

which is against our choice (4.16) of $\xi \in \mathbb{S}^{1} \cap D$. As a consequence $\Delta$ intersects $B^{\nu}$ in at most one point $z_{*}$.

If $\Delta \cap \Pi_{\nu}=\left\{z_{*}\right\}$, we define $Z_{*}^{2}:=B^{\nu} \cap B_{\frac{\eta}{2}}\left(z_{*}\right)$ while if $\Delta \cap \Pi_{\nu}=\emptyset$, we define $Z_{*}^{2}=\emptyset$. The continuity of $\alpha \cdot \operatorname{id}_{\mathbb{R}^{2}}+\beta$ on the compact set $B^{\nu} \backslash Z_{*}^{2}$ entails that

$$
0<m_{\alpha, \beta}(\eta):=\min \left\{|\alpha \cdot y+\beta|: y \in B^{\nu} \backslash Z_{*}^{2}\right\}
$$

Set $Z_{*}:=Z_{*}^{1} \cup Z_{*}^{2} \subset B^{\nu}$, which satisfies $\mathcal{H}^{1}\left(Z_{*}\right) \leq 3 \eta$, and for all $j \in \mathbb{N}$, we define

$$
\widehat{\mathbf{T}}_{j}:=\left\{T \in \mathbf{T}_{b, i n t}^{x_{0}, j}: \text { there exists } y \in Y_{j} \backslash Z_{*} \text { such that }(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset\right\}
$$

Thus, for all $j \in \mathbb{N}$ and for each triangle $T \in \widehat{\mathbf{T}}_{j}$, there exists a point $y_{T} \in \Phi \circ p_{\xi}(\stackrel{\circ}{T}) \backslash Z_{*}^{2} \subset B^{\nu} \backslash Z_{*}^{2}$ which satisfies

$$
\left|\alpha \cdot y_{T}+\beta\right| \geq m_{\alpha, \beta}(\eta)>0
$$

with $\Phi$ introduced in (3.16).
Remembering that $\omega\left(\varepsilon_{k_{j}}\right) / \varrho_{j} \rightarrow 0$ and that the Lipschitz constant of $\Phi$ is less than $\sqrt{1+4 \eta^{2}} \leq 2$ for $\eta$ small enough, together with the uniform convergence (4.21) and (4.20), it follows that for all $\gamma>0$, there exists a rank $j_{2}(\gamma) \geq j_{1}(\gamma)$ such that for all $j \geq j_{2}(\gamma)$,

$$
\begin{cases}\left|\frac{1}{\mu_{j}}\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}-(\alpha \cdot y+\beta)\right| \leq \frac{m_{\alpha, \beta}}{8} \gamma \quad \text { for all } y \in B^{\nu} \backslash Z_{*}  \tag{4.22a}\\ \mathcal{H}^{1}\left(\Phi \circ p_{\xi}(T)\right) \leq 2 \omega\left(\varepsilon_{k_{j}}\right) / \varrho_{j} \leq \frac{m_{\alpha, \beta}}{8} \gamma \quad \text { for all } T \in \mathbf{T}^{x_{0}, j}\end{cases}
$$

Therefore, for all $j \geq j_{2}(\gamma)$ and all $T \in \widehat{\mathbf{T}}_{j}$, we introduce the following quantities :

$$
\left\{\begin{array}{l}
L^{\mathrm{ref}}(T):=\frac{\left|[u]\left(x_{0}\right) \cdot \xi+\mu_{j}\left(\alpha \cdot y_{T}+\beta\right)\right|-\left(C_{\eta}+\frac{m_{\alpha, \beta}}{2} \mu_{j}\right) \gamma}{\left|e\left(v_{j}\right)_{\mid T}:(\xi \otimes \xi)\right|} \text { the reference length of } T,  \tag{4.23}\\
L^{\max }(T):=\max _{z \in p_{\xi}(T)} \mathcal{L}^{1}\left(T_{z}^{\xi}\right) \text { the maximal section's length of } T \text { along the direction } \xi
\end{array}\right.
$$

Note that $L^{\mathrm{ref}}(T)$ is well defined (since $\left|e\left(v_{j}\right)_{\mid T} \xi \cdot \xi\right|^{2} \geq \kappa \varrho_{j}^{2} /\left(\alpha \varepsilon_{k_{j}}\right)>0$ as $T \in \mathbf{T}_{b}^{x_{0}, j}$ ) and positive for $j$ large enough since $\mu_{j} \rightarrow+\infty$ and $\left|\alpha \cdot y_{T}+\beta\right|>m_{\alpha, \beta} \gamma / 2>0$. Moreover, we have
$L^{\max }(T)>L^{\text {ref }}(T)$. Indeed, if such would not be the case, denoting by $y \in Y_{j} \backslash Z_{*}$ a point such that $(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset$, then $\mathcal{L}^{1}\left(T_{p_{\xi}(y)}^{\xi}\right)=\mathcal{L}^{1}\left(T_{y}^{\xi}\right)=b_{j}(y)-a_{j}(y) \leq L^{\max }(T) \leq L^{\text {ref }}(T)$, entailing that

$$
\begin{aligned}
&\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)\right|=\left|e\left(v_{j}\right)_{\mid T}:(\xi \otimes \xi)\right|\left(b_{j}(y)-a_{j}(y)\right) \\
& \leq\left|[u]\left(x_{0}\right) \cdot \xi+\mu_{j}\left(\alpha \cdot y_{T}+\beta\right)\right|-\left(C_{\eta}+\frac{m_{\alpha, \beta}}{2} \mu_{j}\right) \gamma
\end{aligned}
$$

by definition (4.23) of $L^{\text {ref }}(T)$. Therefore, we would obtain that

$$
\begin{aligned}
C_{\eta} \gamma+\frac{m_{\alpha, \beta}}{2} \mu_{j} \gamma \leq & \left|[u]\left(x_{0}\right) \cdot \xi+\mu_{j}\left(\alpha \cdot y_{T}+\beta\right)\right|-\left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)\right| \\
\leq & \left|\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)-[u]\left(x_{0}\right) \cdot \xi-\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}\right| \\
& \quad+\left|\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}-\mu_{j}(\alpha \cdot y+\beta)\right|+\left|\mu_{j} \alpha \cdot\left(y-y_{T}\right)\right| \\
\leq & C_{\eta} \gamma+\frac{m_{\alpha, \beta}}{8} \mu_{j} \gamma+\frac{m_{\alpha, \beta}}{8} \mu_{j} \gamma
\end{aligned}
$$

where we used $(4.20),(4.22 \mathrm{a})$ and $(4.22 \mathrm{~b})$ (since $y, y_{T} \in \Phi \circ p_{\xi}(T)$ ), leading to a contradiction.
Therefore, arguing as in the proof of Lemma 3.9, there are exactly one or two points $z_{\text {ref }}^{1}$, $z_{\text {ref }}^{2} \in p_{\xi}(T)$, only depending on $j$ and $T$, such that $\mathcal{L}^{1}\left(T_{z_{\text {ref }}^{1}}^{\xi}\right)=\mathcal{L}^{1}\left(T_{z_{\text {ref }}^{2}}^{\xi}\right)=L^{\text {ref }}(T)$. Then, as in (3.29), we introduce the following segments (orthogonal to $\xi$ ) associated to $T$ (see Figure 6),

$$
\mathfrak{T}_{i}(T):=\left\{z \in \Pi_{\xi}:\left|z-z_{\mathrm{ref}}^{i}\right| \leq C_{\eta}^{\prime} \frac{\varrho_{j} \mathcal{L}^{2}(T)}{\varepsilon_{k_{j}}} \gamma\right\} \quad \text { for } i \in\{1,2\}
$$

where the constant $C_{\eta}^{\prime}$, only depending on $\eta$, now changes into

$$
C_{\eta}^{\prime}:=\frac{8}{\sin \theta_{0}}\left(\frac{2 C_{\eta}}{m_{\alpha, \beta}}+1\right)
$$

For every $j \geq j_{2}(\gamma)$ and every $y \in Y_{j} \backslash Z_{*}$, let $T \in \mathbf{T}_{b, i n t}^{x_{0}, j}$ be such that $(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset$. In particular, note that $T \in \widehat{\mathbf{T}}_{j}$. Arguing in the same way as in the proof of Lemma 3.9, we get that there exists $i \in\{1,2\}$ such that

$$
\begin{aligned}
& \left|p_{\xi}(y)-z_{\mathrm{ref}}^{i}\right| \\
& \quad \leq \frac{2 \mathcal{L}^{2}(T)}{h_{T}} \frac{\left|\left(b_{j}(y)-a_{j}(y)\right)-L^{\mathrm{ref}}(T)\right|}{L^{\mathrm{ref}}(T)} \\
& \quad=\frac{2 \mathcal{L}^{2}(T)}{h_{T}} \frac{| |\left(v_{j}\right)_{y}^{\xi}\left(b_{j}(y)\right)-\left(v_{j}\right)_{y}^{\xi}\left(a_{j}(y)\right)\left|-\left|[u]\left(x_{0}\right) \cdot \xi+\mu_{j}\left(\alpha \cdot y_{T}+\beta\right)\right|+\left(C_{\eta}+\frac{m_{\alpha, \beta}}{2} \mu_{j}\right) \gamma\right|}{\left|[u]\left(x_{0}\right) \cdot \xi+\mu_{j}\left(\alpha \cdot y_{T}+\beta\right)\right|-\left(C_{\eta}+\frac{m_{\alpha, \beta}}{2} \mu_{j}\right) \gamma} \\
& \quad \leq \frac{2 \mathcal{L}^{2}(T)}{h_{T}} \frac{C_{\eta} \gamma+\left|\left(r_{j}^{+}-r_{j}^{-}\right)_{y}^{\xi}-\mu_{j}\left(\alpha \cdot y_{T}+\beta\right)\right|+\left(C_{\eta}+\frac{m_{\alpha, \beta}}{2} \mu_{j}\right) \gamma}{\mu_{j} m_{\alpha, \beta} / 4} \\
& \quad \leq \frac{2 \mathcal{L}^{2}(T)}{h_{T}} \frac{\left(2 C_{\eta}+m_{\alpha, \beta} \mu_{j}\right) \gamma}{\mu_{j} m_{\alpha, \beta} / 4} \leq C_{\eta}^{\prime} \frac{\varrho_{j} \mathcal{L}^{2}(T)}{\varepsilon_{k_{j}}} \gamma,
\end{aligned}
$$

where we used (4.20), (4.22a), (4.22b) and the fact that

$$
\left|[u]\left(x_{0}\right) \cdot \xi+\mu_{j}\left(\alpha \cdot y_{T}+\beta\right)\right|-\left(C_{\eta}+\frac{\mu_{j} m_{\alpha, \beta}}{2}\right) \gamma \geq \frac{\mu_{j} m_{\alpha, \beta}}{4}
$$

up to enlarging $j_{2}(\gamma) \in \mathbb{N}$. As in the proof of Lemma 3.9, we deduce that for all $j \geq j_{2}(\gamma)$,

$$
\mathcal{H}^{1}\left(Y_{j} \backslash Z_{*}\right) \leq \sum_{T \in \widehat{\mathbf{T}}_{j}} \mathcal{H}^{1}\left(\Phi\left(\mathfrak{T}_{1}(T) \cup \mathfrak{T}_{2}(T)\right)\right) \leq 8 C_{\eta}^{\prime} \gamma \frac{\varrho_{j}}{\varepsilon_{k_{j}}} \sum_{T \in \widehat{\mathbf{T}}_{j}} \mathcal{L}^{2}(T) \leq \frac{8 C_{\eta}^{\prime} \gamma}{\kappa} \frac{\kappa \varrho_{j}}{\varepsilon_{k_{j}}} \int_{B} \chi_{j} d x
$$

Recalling (4.17) and possibly taking a larger $j_{2}(\gamma) \in \mathbb{N}$, we finally get that for all $j \geq j_{2}(\gamma)$,

$$
\mathcal{H}^{1}\left(Y_{j} \backslash Z_{*}\right) \leq \frac{8 C_{\eta}^{\prime}}{\kappa}\left(2 \frac{d \lambda}{d \mathcal{H}^{1}\left\llcorner\partial^{*} \mathcal{P}\right.}\left(x_{0}\right)+1\right) \gamma=: C_{*} \gamma
$$

for some constant $C_{*}>0$ only depending on $\eta$, which settles Lemma 4.8.
Arguing exactly as in the proof of Lemma 3.8, having Lemma 4.8 at hand, we deduce the following result.
Lemma 4.9. For all $\eta>0$, there exist $Z^{\prime \prime} \subset B^{\nu}$ containing $Z^{\prime}$ with $\mathcal{H}^{1}\left(Z^{\prime \prime}\right) \leq 4 \eta$, and a (not relabeled) subsequence such that for all $j \in \mathbb{N}$ and for all $y \in\left(B_{1-\frac{\eta}{2}}\right)^{\nu} \backslash Z^{\prime \prime}$,

$$
\#\left\{T \in \mathbf{T}_{b, i n t}^{x_{0}, j}:(\stackrel{\circ}{T} \cap B)_{y}^{\xi} \neq \emptyset\right\} \geq 2
$$

Finally, owing to Lemma 4.9, the proof of Proposition 4.5 is identical to that of Proposition 3.5.
4.3. $\Gamma$-limit and convergence of minimizers. Let us introduce the $\Gamma$-lower and upper limits $\mathcal{G}^{\prime}$ and $\mathcal{G}^{\prime \prime}$ defined, for all $u \in L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$, by

$$
\mathcal{G}^{\prime}(u):=\inf \left\{\liminf _{\varepsilon \rightarrow 0} \mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right): u_{\varepsilon} \rightarrow u \text { in measure in } \Omega^{\prime}\right\}
$$

and

$$
\mathcal{G}^{\prime \prime}(u):=\inf \left\{\limsup _{\varepsilon \rightarrow 0} \mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right): u_{\varepsilon} \rightarrow u \text { in measure in } \Omega^{\prime}\right\}
$$

Proof of Theorem 4.1. Lower bound. Let $u \in L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$. Without loss of generality, we can assume that $\mathcal{G}^{\prime}(u)<+\infty$. For any $\zeta>0$, there exists a sequence $\left\{u_{\varepsilon}\right\}_{\varepsilon>0}$ such that $u_{\varepsilon} \rightarrow u$ in measure in $\Omega^{\prime}$ and

$$
\liminf _{\varepsilon \rightarrow 0} \mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right) \leq \mathcal{G}^{\prime}(u)+\zeta<+\infty
$$

Let us extract a subsequence $\left\{u_{k}\right\}_{k \in \mathbb{N}}:=\left\{u_{\varepsilon_{k}}\right\}_{k \in \mathbb{N}}$ such that $u_{k} \rightarrow u \mathcal{L}^{2}$-a.e. in $\Omega^{\prime}$ and

$$
\lim _{k \rightarrow \infty} \mathcal{G}_{\varepsilon_{k}}\left(u_{k}\right)=\liminf _{\varepsilon \rightarrow 0} \mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right)<+\infty
$$

This implies that, for $k$ large enough, $u_{k} \in V_{\varepsilon_{k}}^{\operatorname{Dir}}\left(\Omega^{\prime}\right)$ and $\sup _{k} \mathcal{G}_{\varepsilon_{k}}\left(u_{k}\right)<+\infty$. Hence, according to Proposition 4.3, there exist a subsequence (not relabeled), a Caccioppoli partition $\mathcal{P}=\left\{P_{j}\right\}_{j \in \mathbb{N}}$ of $\Omega^{\prime}$, a sequence of piecewise rigid body motions $\left\{r_{k}\right\}_{k \in \mathbb{N}}$ with

$$
r_{k}:=\sum_{j \in \mathbb{N}} r_{k}^{j} \mathbb{1}_{P_{j}},
$$

and a function $v \in G S B D^{2}\left(\Omega^{\prime}\right)$ such that $u_{k}-r_{k} \rightarrow v$ in measure in $\Omega^{\prime}, v=w \mathcal{L}^{2}$-a.e. in $\Omega^{\prime} \backslash \bar{\Omega}$, and

$$
\liminf _{k \rightarrow \infty} \mathcal{G}_{\varepsilon_{k}}\left(u_{k}\right) \geq \int_{\Omega} \mathbf{A} e(v): e(v) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v} \cup \partial^{*} \mathcal{P}\right)
$$

In particular, $r_{k} \rightarrow u-v$ in measure in $\Omega^{\prime}$, implying that for all $j \in \mathbb{N}, r_{k}^{j} \rightarrow u-v$ in measure in $P_{j}$. In particular, for all $j \in \mathbb{N}$, there exists a rigid body motion $r^{j}$ such that $r_{\mid P_{j}}^{j}=u-v \mathcal{L}^{2}$-a.e. in $P_{j}$. Therefore, the piecewise rigid body motion

$$
r:=\sum_{j \in \mathbb{N}} r^{j} \mathbb{1}_{P_{j}}
$$

is such that $u=r+v \quad \mathcal{L}^{2}$-a.e. in $\Omega^{\prime}$, hence $u=r+v \in G S B D^{2}\left(\Omega^{\prime}\right)$. It follows that $e(u)=e(v)$ $\mathcal{L}^{2}$-a.e. in $\Omega$ and that $J_{u} \subset J_{v} \cup \partial^{*} \mathcal{P}$ up to a $\mathcal{H}^{1}$-negligible set. Thus,

$$
\liminf _{k \rightarrow \infty} \mathcal{G}_{\varepsilon_{k}}\left(u_{k}\right) \geq \int_{\Omega} \mathbf{A} e(u): e(u) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{u}\right)
$$

Recalling that $J_{u} \cap \partial \Omega=\{u \neq w\} \cap \partial \Omega$ and $J_{u} \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right)=J_{w} \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right)=\emptyset$, it entails that $\zeta+\mathcal{G}^{\prime}(u) \geq \mathcal{G}(u)$, and the conclusion follows letting $\zeta \searrow 0$.

Upper bound. Let $u \in L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$. We can assume that $\mathcal{G}(u)<+\infty$ so that $u \in \operatorname{GSBD}^{2}\left(\Omega^{\prime}\right)$ and $u=w \mathcal{L}^{2}$-a.e. in $\Omega^{\prime} \backslash \bar{\Omega}$. According to the density results for $G S B D$ functions (see [19, Theorem 1.1] and [19, Formula (5.11)]), there exists a sequence of functions $u_{n} \in S B V^{2}\left(\Omega ; \mathbb{R}^{2}\right) \cap L^{\infty}\left(\Omega ; \mathbb{R}^{2}\right)$ such that

$$
\left\{\begin{array}{l}
u_{n} \rightarrow u \text { in measure in } \Omega  \tag{4.24}\\
u_{n}=w \text { in an open bounded neighborhood of } \partial \Omega \\
{\lim \sup _{n} \mathcal{G}\left(u_{n}\right) \leq \mathcal{G}(u)}^{\text {l }}
\end{array}\right.
$$

Extending (continuously) $u_{n}$ by $w$ on $\Omega^{\prime} \backslash \bar{\Omega}$, we get that $u_{n} \rightarrow u$ in measure in $\Omega^{\prime}$. The proof is thus complete once we know that $\mathcal{G}^{\prime \prime}\left(u_{n}\right) \leq \mathcal{G}\left(u_{n}\right)$, for all $n \in \mathbb{N}$, as it would imply $\mathcal{G}^{\prime \prime}(u) \leq \mathcal{G}(u)$, using the lower semicontinuity of $\mathcal{G}^{\prime \prime}$ in $L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ with respect to the convergence in measure together with the last point of (4.24).

Therefore, we can assume without loss of generality that $u \in S B V^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right) \cap L^{\infty}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ and $u=w$ on $V \cup\left(\Omega^{\prime} \backslash \bar{\Omega}\right)$ with $V$ an open bounded neighborhood of $\partial \Omega$. Next owing to the density result in $S B V$ (see [24, Theorem 3.1]), there exist a sequence $\left\{u_{k}\right\}_{k \in \mathbb{N}}$ in $S B V^{2}\left(\Omega ; \mathbb{R}^{2}\right) \cap L^{\infty}\left(\Omega ; \mathbb{R}^{2}\right)$ as well as $N_{k}$ disjoint closed segments $L_{1}^{k}, \ldots, L_{N_{k}}^{k} \subset \bar{\Omega}$ satisfying :

$$
\begin{aligned}
& \overline{J_{u_{k}}}=\bigcup_{i=1}^{N_{k}} L_{i}^{k}, \quad \mathcal{H}^{1}\left(\overline{J_{u_{k}}} \backslash J_{u_{k}}\right)=0, \quad u_{k} \in W^{2, \infty}\left(\Omega \backslash \overline{J_{u_{k}}} ; \mathbb{R}^{2}\right), \\
& \left\{\begin{array}{l}
u_{k} \rightarrow u \quad \text { strongly in } L^{1}\left(\Omega ; \mathbb{R}^{2}\right), \\
\nabla u_{k} \rightarrow \nabla u \text { strongly in } L^{2}\left(\Omega ; \mathbb{M}^{2 \times 2}\right), \\
\lim _{\sup }^{k}
\end{array} \mathcal{H}^{1}\left(\bar{A} \cap J_{u_{k}}\right) \leq \mathcal{H}^{1}\left(\bar{A} \cap J_{u}\right) \text { for all open subset } A \subset \subset \Omega,\right.
\end{aligned}
$$

and with the additional property that, because $u=w$ is smooth on $V$, there exists a slightly smaller neighborhood of $\partial \Omega$, say $W \subset V$, on which the approximating functions $u_{k}$ simply are the convolution between $w$ and a given standard smooth mollifier $\phi_{k}$, i.e. $u_{k}=w * \phi_{k}$ on $W \cap \Omega$. We can thus consider the corrected approximating functions $\hat{u}_{k}:=u_{k}+w-w * \phi_{k}$, which still satisfy

$$
\begin{gather*}
\overline{J_{\hat{u}_{k}}}=\bigcup_{i=1}^{N_{k}} L_{i}^{k}, \quad \mathcal{H}^{1}\left(\overline{J_{\hat{u}_{k}}} \backslash J_{\hat{u}_{k}}\right)=0, \quad \hat{u}_{k} \in W^{2, \infty}\left(\Omega \backslash \overline{J_{\hat{u}_{k}}} ; \mathbb{R}^{2}\right), \\
 \tag{4.25}\\
\begin{cases}\hat{u}_{k} \rightarrow u & \text { strongly in } L^{1}\left(\Omega ; \mathbb{R}^{2}\right), \\
\nabla \hat{u}_{k} \rightarrow \nabla u & \text { strongly in } L^{2}\left(\Omega ; \mathbb{M}^{2 \times 2}\right),\end{cases}
\end{gather*}
$$

and now

$$
\hat{u}_{k}=w \text { on } W \cap \Omega \text {. }
$$

Moreover, with $A \subset \subset \Omega$ an open subset such that $\Omega \backslash W \subset A$, since $J_{u}$ and $J_{\hat{u}_{k}}=J_{u_{k}}$ are contained in $\Omega \backslash W$, we infer that

$$
\begin{equation*}
\limsup _{k \rightarrow \infty} \mathcal{H}^{1}\left(J_{\hat{u}_{k}}\right)=\limsup _{k \rightarrow \infty} \mathcal{H}^{1}\left(\bar{A} \cap J_{\hat{u}_{k}}\right) \leq \mathcal{H}^{1}\left(\bar{A} \cap J_{u}\right)=\mathcal{H}^{1}\left(J_{u}\right) . \tag{4.26}
\end{equation*}
$$

We extend continuously $\hat{u}_{k}$ to $\Omega^{\prime}$ by setting $\hat{u}_{k}=w$ on $\Omega^{\prime} \backslash \Omega$, so that $\hat{u}_{k} \in S B V^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right) \cap$ $L^{\infty}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$, its jump set $J_{\hat{u}_{k}} \subset \Omega \backslash W$ is unchanged, $\hat{u}_{k} \in W^{2, \infty}\left(\Omega^{\prime} \backslash \overline{J_{\hat{u}}} ; \mathbb{R}^{2}\right)$ and $\hat{u}_{k} \rightarrow u$ in $L^{1}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$. Using again the lower semicontinuity of $\mathcal{G}^{\prime \prime}$ in $L^{0}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ with respect to the convergence in measure together with the convergences in (4.25) and (4.26), we are back to establishing that $\mathcal{G}^{\prime \prime}\left(\hat{u}_{k}\right) \leq \mathcal{G}\left(\hat{u}_{k}\right)$, as the conclusion follows letting $k \rightarrow+\infty$. Arguing almost word for word as in the proof of Lemma 3.12, we show the following Lemma 4.10, which leads to the desired inequality.

Lemma 4.10. Let $v \in S B V^{2}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right) \cap L^{\infty}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ be such that $v=w$ on $W \cup\left(\Omega^{\prime} \backslash \bar{\Omega}\right)$ with $W$ an open bounded neighbourhood of $\partial \Omega$, and

$$
\overline{\Omega \backslash W} \supset \overline{J_{v}}=\bigcup_{i=1}^{N} L_{i}, \quad \mathcal{H}^{1}\left(\overline{J_{v}} \backslash J_{v}\right)=0, \quad v \in W^{2, \infty}\left(\Omega^{\prime} \backslash \overline{J_{v}} ; \mathbb{R}^{2}\right)
$$

for some pairwise disjoint closed segments $L_{1}, \ldots, L_{N} \subset \overline{\Omega \backslash W}$. Then,

$$
\mathcal{G}^{\prime \prime}(v) \leq \int_{\Omega} \mathbf{A} e(v): e(v) d x+\kappa \sin \theta_{0} \mathcal{H}^{1}\left(J_{v}\right)=\mathcal{G}(v)
$$

We do not detail the proof of Lemma 4.10. We only stress that, following Lemma 3.12, for $\varepsilon>0$ small enough, if $\mathbf{T}_{\varepsilon}$ is the admissible triangulation given by [21, Appendix A$]$ and $v_{\varepsilon}$ is the Lagrange interpolation of the values of $v$ at the vertices of $\mathbf{T}_{\varepsilon}$, each triangle $T \in \mathbf{T}_{\varepsilon}$ such that $T \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right) \neq \emptyset$ is contained in $W \cup \mathbb{R}^{2} \backslash \Omega$, so that $v_{\varepsilon}=w_{\mathbf{T}_{\varepsilon}}$ on $T$. In particular, it ensures that $v_{\varepsilon} \in V_{\varepsilon}^{\operatorname{Dir}}\left(\Omega^{\prime}\right)$.

In the following result, we prove the existence of minimizers of the discrete brittle damage energy $\mathcal{G}_{\varepsilon}$ on $V_{\varepsilon}^{\operatorname{Dir}}\left(\Omega^{\prime}\right)$.

Lemma 4.11. Assume that $\Omega$ and $\Omega^{\prime}$ are connected. For $\varepsilon>0$ sufficiently small, there exists $a$ minimizer $u_{\varepsilon} \in V_{\varepsilon}^{\operatorname{Dir}}\left(\Omega^{\prime}\right)$ of $\mathcal{G}_{\varepsilon}$.
Proof. Let $\varepsilon_{0}:=\kappa /\left(\beta\|\nabla w\|_{L^{\infty}\left(\mathbb{R}^{2} ; \mathbb{M}^{2 \times 2}\right)}^{2}\right)$ and fix $\varepsilon<\varepsilon_{0}$. Since $\mathcal{G}_{\varepsilon}\left(w_{\mathbf{T}_{\varepsilon}}\right)<+\infty$, we can consider a minimizing sequence $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subset V_{\varepsilon}^{\text {Dir }}\left(\Omega^{\prime}\right)$ satisfying

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \mathcal{G}_{\varepsilon}\left(u_{n}\right)=\inf _{L^{0}\left(\Omega ; \mathbb{R}^{2}\right)} \mathcal{G}_{\varepsilon} \in[0,+\infty) . \tag{4.27}
\end{equation*}
$$

By definition of the finite element space $V_{\varepsilon}^{\operatorname{Dir}}\left(\Omega^{\prime}\right)$, there exists a triangulation $\mathbf{T}^{n} \in \mathcal{T}_{\varepsilon}\left(\Omega^{\prime}\right)$ such that $u_{n}$ is affine on each $T \in \mathbf{T}^{n}$ and $u_{n}=w_{\mathbf{T}^{n}}$ on every triangle $T \in \mathbf{T}^{n}$ such that $T \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right) \neq \emptyset$.

Let $\Omega^{\prime \prime}$ be a bounded open set such that $\Omega^{\prime} \subset \subset \Omega^{\prime \prime}$ and $\bigcup_{T \in \mathbf{T}_{n}} T \subset \Omega^{\prime \prime}$ for all $n \in \mathbb{N}$. Since, for all $T \in \mathbf{T}^{n}, \mathcal{L}^{2}(T) \geq \varepsilon^{2} \sin \theta_{0} / 2$, it is easily seen that

$$
\# \mathbf{T}^{n} \leq \frac{2 \mathcal{L}^{2}\left(\Omega^{\prime \prime}\right)}{\varepsilon^{2} \sin \theta_{0}}
$$

As a consequence, the sequence of integers $\left\{\# \mathbf{T}^{n}\right\}_{n \in \mathbb{N}}$ admits a subsequence converging as $n \rightarrow+\infty$ to an integer $N \in \mathbb{N}$. We can thus assume, without loss of generality, that

$$
\# \mathbf{T}^{n}=N \quad \text { for all } n \in \mathbb{N}
$$

We write $\mathbf{T}^{n}=\left\{T_{1}^{n}, \ldots, T_{N}^{n}\right\}$ for all $n \in \mathbb{N}$. Up to a subsequence, we can check that for all $i \in\{1, \ldots, N\}$, the closed triangle $T_{i}^{n}$ converges to a closed limit triangle $T_{i}$ in the sense of Hausdorff, with the property that the limit triangulation $\mathbf{T}:=\left\{T_{1}, \ldots, T_{N}\right\} \in \mathcal{T}_{\varepsilon}\left(\Omega^{\prime}\right)$ remains an admissible triangulation of $\Omega^{\prime}$.

Introducing the characteristic functions $\chi_{n}:=\mathbb{1}_{\left\{\varepsilon \mathbf{A} e\left(u_{n}\right): e\left(u_{n}\right) \geq \kappa\right\}} \in L^{\infty}\left(\Omega^{\prime} ;\{0,1\}\right)$, we can write the energy as

$$
\begin{equation*}
\mathcal{G}_{\varepsilon}\left(u_{n}\right)=\int_{\Omega}\left(1-\chi_{n}\right) \mathbf{A} e\left(u_{n}\right): e\left(u_{n}\right) d x+\frac{\kappa}{\varepsilon} \int_{\Omega} \chi_{n} d x \tag{4.28}
\end{equation*}
$$

First, by definition of $\varepsilon_{0}$ and since $\varepsilon<\varepsilon_{0}$, we have that $\chi_{n}=0$ in $\Omega^{\prime} \backslash \bar{\Omega}$ for all $n \in \mathbb{N}$, since

$$
\mathbf{A} e\left(u_{n}\right): e\left(u_{n}\right)=\mathbf{A} e\left(w_{\mathbf{T}_{n}}\right): e\left(w_{\mathbf{T}_{n}}\right) \leq \beta\left|e\left(w_{\mathbf{T}_{n}}\right)\right|^{2} \leq \beta\left|\nabla w_{\mathbf{T}_{n}}\right|^{2} \leq \beta\|\nabla w\|_{L^{\infty}\left(\mathbb{R}^{2} ; \mathbb{M}^{2 \times 2}\right)}^{2}<\frac{\kappa}{\varepsilon}
$$

on that set. Being constant equal to 1 or 0 on each triangle of $\mathbf{T}^{n}, \chi_{n}$ can be identified with a vector $V_{n} \in\{0,1\}^{N}$. Hence, up to a subsequence, there exists $V \in\{0,1\}^{N}$ such that $V_{n} \rightarrow V$ in
$\mathbb{R}^{N}$. In particular, there exists $n_{0} \in \mathbb{N}$ such that $V_{n}=V$ for all $n \geq n_{0}$. Up to reordering the triangles, we can thus find a integer $0 \leq M<N$ such that

$$
\left\{\chi_{n}=1\right\}=\bigcup_{i=1}^{M} T_{i}^{n}, \quad\left\{\chi_{n}=0\right\}=\bigcup_{i=M+1}^{N} T_{i}^{n} \quad \text { for all } n \geq n_{0}
$$

By the Hausdorff convergence property, we infer that

$$
\begin{equation*}
\chi_{n} \rightarrow \chi:=\mathbb{1}_{\bigcup_{i=1}^{M} T_{i}} \quad \text { strongly in } L^{1}\left(\Omega^{\prime}\right) \tag{4.29}
\end{equation*}
$$

We next show some compactness on the sequence of displacements $\left\{u_{n}\right\}_{n \in \mathbb{N}}$. Remembering that $\left(1-\chi_{n}\right)\left|e\left(u_{n}\right)\right|^{2} \leq \kappa /(\alpha \varepsilon)$ for all $n \geq n_{0}$ and that the sequence $\left\{\left(1-\chi_{n}\right) e\left(u_{n}\right)\right\}_{n \in \mathbb{N}}$ lives in the finite dimensional space $\left(\mathbb{M}^{2 \times 2}\right)^{N}$, up to a new subsequence (not relabeled), there exists a function $\xi \in L^{\infty}\left(\Omega^{\prime} ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right)$ which is constant on each triangle $T \in \mathbf{T}$ such that

$$
\begin{equation*}
\left(1-\chi_{n}\right) e\left(u_{n}\right) \rightarrow \xi \text { strongly in } L^{2}\left(\Omega^{\prime} ; \mathbb{M}_{\mathrm{sym}}^{2 \times 2}\right) \tag{4.30}
\end{equation*}
$$

and $\xi=0$ on $\bigcup_{i=1}^{M} T_{i}$. Let us define the set

$$
\omega_{0}:=\bigcup_{i=M+1}^{N} T_{i}
$$

Note that $\Omega^{\prime} \backslash \bar{\Omega} \subset \omega_{0}$. Indeed, if $x \in \Omega^{\prime} \backslash \bar{\Omega}$, then for all $n \geq n_{0}$, there exists $M+1 \leq i_{n} \leq N$ such that $x \in T_{i_{n}}^{n}$. At the expense of extracting a further subsequence, there is no loss of generality to assume that $i_{n}=i$ is independent of $n$. By the Hausdorff convergence of $T_{i}^{n}$ to $T_{i}$, we infer that $x \in T_{i} \subset \omega_{0}$. By connectedness of $\Omega^{\prime} \backslash \bar{\Omega} \subset \omega_{0}$, we can consider $\omega$ the connected component of $\omega_{0}$ containing $\Omega^{\prime} \backslash \bar{\Omega}$. Let $M \leq K<N$ be such that $\omega=\bigcup_{i=K+1}^{N} T_{i}$, up to reordering the triangles again.

Observe that for all $T \in \mathbf{T}$ such that $T \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right) \neq \emptyset$, then $T \in\left\{T_{K+1}, \ldots, T_{N}\right\}$. Thus, for all $T \in\left\{T_{1}, \ldots, T_{K}\right\}, T \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right)=\emptyset$ so that $T \subset \bar{\Omega}$. Therefore, for all open set $W \subset \subset \Omega^{\prime}$ with $\bigcup_{i=1}^{K} T_{i} \subset W$, having that

$$
\bigcup_{i=1}^{K} T_{i}^{n} \rightarrow \bigcup_{i=1}^{K} T_{i} \quad \text { in the sense of Hausdorff }
$$

there exists $n_{1} \geq n_{0}$ such that $\bigcup_{i=1}^{K} T_{i}^{n} \subset W$ for all $n \geq n_{1}$. Since

$$
\Omega^{\prime} \backslash \bar{W} \subset \bigcup_{i=K+1}^{N} T_{i}^{n} \subset \bigcup_{i=M+1}^{N} T_{i}^{n}=\left\{\chi_{n}=0\right\}
$$

owing to (4.27), (4.28) and that $u_{n}=w_{\mathbf{T}^{n}}$ in $\Omega^{\prime} \backslash \bar{\Omega}$, we infer that

$$
\int_{\Omega^{\prime} \backslash \bar{W}}\left|e\left(u_{n}\right)\right|^{2} d x \leq C_{*},
$$

for some constant $C_{*}>0$ independent of $n$ and $W$. Using that $u_{n}-w_{\mathbf{T}^{n}} \in H^{1}\left(\Omega^{\prime} \backslash \bar{W} ; \mathbb{R}^{2}\right)$ is equal to 0 on the open set $\left(\Omega^{\prime} \backslash \bar{W}\right) \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right) \neq \emptyset$, the Poincaré-Korn inequality ensures that (up to a subsequence) there exists $u \in H^{1}\left(\Omega^{\prime} \backslash \bar{W} ; \mathbb{R}^{2}\right)$ such that

$$
u_{n} \rightharpoonup u \text { weakly in } H^{1}\left(\Omega^{\prime} \backslash \bar{W} ; \mathbb{R}^{2}\right)
$$

$u=w_{\mathbf{T}}$ on $\left(\Omega^{\prime} \backslash \bar{W}\right) \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right)$ since $w_{\mathbf{T}^{n}} \rightarrow w_{\mathbf{T}}$ strongly in $H^{1}\left(\Omega^{\prime} ; \mathbb{R}^{2}\right)$ and, thanks to (4.30), $e(u)=\xi$ in $\Omega^{\prime} \backslash \bar{W}$. In addition, by weak lower semicontinuity of the norm, we get that

$$
\int_{\Omega^{\prime} \backslash \bar{W}}|e(u)|^{2} d x \leq \liminf _{n \rightarrow \infty} \int_{\Omega^{\prime} \backslash \bar{W}}\left|e\left(u_{n}\right)\right|^{2} d x \leq C_{*}
$$

Considering a decreasing sequence of open sets $\left\{W_{j}\right\}_{j \in \mathbb{N}}$ such that $\bigcup_{i=1}^{K} T_{i} \subset W_{j} \subset \subset \Omega^{\prime}$ for each $j \in \mathbb{N}$, and $\bigcap_{j} W_{j}=\bigcup_{i=1}^{K} T_{i}$, we deduce through a diagonalisation argument that there exists $u \in H^{1}\left(\stackrel{\circ}{\omega} \cap \Omega^{\prime} ; \mathbb{R}^{2}\right)$ such that $u=w_{\mathbf{T}}$ on $\stackrel{\circ}{\omega} \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right)=\Omega^{\prime} \backslash \bar{\Omega}$ and $e(u)=\xi$ in $\stackrel{\circ}{\omega} \cap \Omega^{\prime}$. In particular, since $\xi$ is constant in each triangle of $\omega$, we infer that $u$ is affine in the interior of each triangle of $\omega$. Being in $H^{1}\left(\stackrel{\circ}{\omega} \cap \Omega^{\prime} ; \mathbb{R}^{2}\right)$, we get that $u$ is continuous at the interfaces of each triangle in $\omega$. Moreover, since $u=w_{\mathbf{T}}$ on $\Omega^{\prime} \backslash \bar{\Omega}$, we deduce that $u_{\mid T}=w_{\mathbf{T}}$ on each triangle $T \in \mathbf{T}$ such that $T \cap\left(\Omega^{\prime} \backslash \bar{\Omega}\right) \neq \emptyset$. Note that $u$ is defined on such triangles $T$, as they are included in $\omega$.

In order to extend $u$ outside $\omega$, we introduce the family of triangles which are at a distance of at least one triangle from $\omega$, i.e.

$$
\mathbf{T}^{\mathrm{far}}:=\{T \in \mathbf{T}: T \cap \omega=\emptyset\} \subset\left\{T_{1}, \ldots, T_{K}\right\}
$$

so that every remaining triangle $T \notin \mathbf{T}^{\mathrm{far}}$ and such that $T \not \subset \omega$, has its three vertices in $\operatorname{Vertices}(\omega) \cup \operatorname{Vertices}\left(\mathbf{T}^{\text {far }}\right)$. Note that $\left\{T_{M+1}, \ldots T_{K}\right\} \subset \mathbf{T}^{\text {far }}$ since, by construction of the connected component $\omega$ of $\omega_{0}$, each triangle $T \in \mathbf{T}$ included in $\omega_{0} \backslash \dot{\omega}$ is at a distance of at least one triangle from $\omega$. We extend the function $u$ to all triangles by setting $u \equiv 0$ on every triangle $T \in \mathbf{T}^{\text {far }}$, and by interpolating on each remaining triangle which happens to have its three vertices' values imposed. It defines a function $u \in V_{\varepsilon}^{\operatorname{Dir}}\left(\Omega^{\prime}\right)$ which satisfies $e(u)=\xi$ on $\omega$ and $e(u) \equiv 0$ on each triangle $T \in\left\{T_{M+1}, \ldots T_{K}\right\} \subset \mathbf{T}^{\text {far }}$.

On the one hand, $\xi=0$ in $\{\chi=1\}$, hence $\xi=(1-\chi) \xi$. On the other hand, $e(u)=\xi$ in $\omega$ and $e(u)=0$ in $\omega_{0} \backslash \omega$, so that $(1-\chi) \mathbf{A} \xi: \xi \geq(1-\chi) \mathbf{A} e(u): e(u)$ by positivity of $\mathbf{A}$. Thus, by (4.29) together with (4.30),

$$
\begin{aligned}
\inf _{L^{0}\left(\Omega ; \mathbb{R}^{2}\right)} \mathcal{G}_{\varepsilon} & =\lim _{n \rightarrow \infty}\left\{\int_{\Omega}\left(1-\chi_{n}\right) \mathbf{A} e\left(u_{n}\right): e\left(u_{n}\right) d x+\frac{\kappa}{\varepsilon} \int_{\Omega} \chi_{n} d x\right\} \\
& =\int_{\Omega}(1-\chi) \mathbf{A} \xi: \xi d x+\frac{\kappa}{\varepsilon} \int_{\Omega} \chi d x \geq \int_{\Omega}(1-\chi) \mathbf{A} e(u): e(u) d x+\frac{\kappa}{\varepsilon} \int_{\Omega} \chi d x=\mathcal{G}_{\varepsilon}(u)
\end{aligned}
$$

which settles that $u$ is a minimizer of $\mathcal{G}_{\varepsilon}$.
We are now in position to prove the fundamental property of $\Gamma$-convergence.
Proof of Corollary 4.2. On the one hand, for all $\varepsilon>0$, we remark that $\mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right) \leq \mathcal{G}_{\varepsilon}\left(w_{\mathbf{T}_{\varepsilon}}\right)$ is uniformly bounded due to (4.1). Therefore, Proposition 4.3 implies that, up to a subsequence, there exist a sequence of piecewise rigid motions $\left\{r_{\varepsilon}\right\}_{\varepsilon>0}$ and a function $u \in G S B D^{2}\left(\Omega^{\prime}\right)$ with $u=w \mathcal{L}^{2}$-a.e. in $\Omega^{\prime} \backslash \bar{\Omega}$, such that $u_{\varepsilon}-r_{\varepsilon} \rightarrow u$ in measure in $\Omega^{\prime}$ and $\liminf _{\varepsilon} \mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right) \geq \mathcal{G}(u)$.

On the other hand, the $\Gamma$-convergence of $\mathcal{G}_{\varepsilon}$ to $\mathcal{G}$ ensures that, for all $v \in G S B D^{2}\left(\Omega^{\prime}\right)$ with $v=w \mathcal{L}^{2}$-a.e. in $\Omega^{\prime} \backslash \bar{\Omega}$, there exists a recovery sequence $v_{\varepsilon} \in L^{0}\left(\Omega ; \mathbb{R}^{2}\right)$ such that $v_{\varepsilon} \rightarrow v$ in measure in $\Omega^{\prime}$ and $\mathcal{G}_{\varepsilon}\left(v_{\varepsilon}\right) \rightarrow \mathcal{G}(v)$. Hence

$$
\mathcal{G}(v)=\lim _{\varepsilon \rightarrow 0} \mathcal{G}_{\varepsilon}\left(v_{\varepsilon}\right) \geq \limsup _{\varepsilon \rightarrow 0} \mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right) \geq \liminf _{\varepsilon \rightarrow 0} \mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right) \geq \mathcal{G}(u)
$$

implying both that $u \in \arg \min \mathcal{G}$ and $\mathcal{G}_{\varepsilon}\left(u_{\varepsilon}\right) \rightarrow \mathcal{G}(u)$.

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(J.-F. Babadjian, E. Bonhomme) Université Paris-Saclay, CNRS, Laboratoire de mathématiques d'Orsay, 91405, Orsay, France.

Email address: jean-francois.babadjian@universite-paris-saclay.fr
Email address: elise.bonhomme@universite-paris-saclay.fr

