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Sea-ice detection from near-nadir Ku-band echoes from CFOSAT/SWIM scatterometer

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Key Points:

- For the first time, sea-ice is detected globally from Ku band off-nadir up to 11 degrees incidence thanks to the SWIM/CFOSAT data sets.
- Fully analytical Geophysical Model Functions are derived for the Normalized Radar Cross-Section over open water and sea-ice.
- A comparison of this new sea-ice flag with SSMI ice concentration shows a high accuracy of our algorithm.

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Abstract

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SWIM on board CFOSAT is the first spaceborne, low incidence, rotating scatterometer, aiming at measuring in near-real time ocean waves spectra. With 5 off-nadir beams at incidences between 2° and 10° plus one nadir beam, it covers the Earth in 13 days, including polar regions thanks to its polar orbit. This work aims at exploiting the SWIM data over ice regions with two objectives. An off-nadir data-based sea-ice flag is here proposed that allows, first, to eliminate sea-ice polluted echoes for improving the wave spectrum retrieval, and second, to open perspectives for application of sea-ice monitoring with near nadir Ku-band active sensors. To this end, the signature of both open water and sea-ice radar backscatter is parameterized into Geophysical Model Functions. Then, comparisons with observed profiles through a Bayesian scheme provide a probability of sea-ice presence. After comparison with both model (ECMWF-IFS) and radiometer (SSMI) derived reference data sets, the proposed flag is found to be ready for operational use. At latitudes greater than 40° in absolute value, the proposed flagging algorithm exhibits accuracies of approximately 98% for all beams compared to SSMI data. Beam to beam performances are characterized and show potential for the characterization of sea-ice at Ku band.

1 Introduction

The CFOSAT program (Hauser et al., 2016) is carried out through a cooperation between the French and Chinese Space Agencies (CNES and CNSA respectively). Its primary objective is the characterization of the ocean surface to better model and predict the ocean states and improve the knowledge in ocean/atmosphere exchanges. The CFOSAT satellite was launched on 2018 October 29^{th} on a sun synchronous repetitive orbit with a 13-day cycle. Its altitude and inclination are respectively 520km and 97.4. Two scientific Ku-Band radars with new technical concepts are embarked on-board the satellite: SWIM, a nadir and near-nadir wave scatterometer (Hauser et al., 2017) and SCAT, a wind fan-beam scatterometer (Xiaolong et al., 2010). Thanks to SWIM, directional ocean waves spectra are produced systematically for the very first time, with a real-aperture scanning radar system (Hauser et al., 2021). SWIM operates at small incidence angles from 0 to 10, contrary to other in-flight scatterometers (ASCAT, QuikSCAT,...) which are at higher incidence angles.

Since its launch, the CNES Wind and Wave Instrument Center (CWWIC) has generated in Near Real Time (NRT), *i.e.*, less than three hours from the measurement, SWIM products of levels 0, 1A, 1B, and 2. SWIM L1B and L2 products are freely available since 2019 April 25th on the AVISO+ website (https://www.aviso.altimetry.fr/en/home.html), after a commissioning and calibration period. All the SWIM products are clearly focused on the ocean. With the simultaneous observation of oceanic wind and wave fields, the primary objectives are definitely oriented towards a better understanding of ocean physics and climates, marine meteorology and all associated applications (Hauser et al., 2017). Nonetheless, CFOSAT also offers an opportunity to provide data for the estimation of land surface parameters and polar ice sheet characteristics as both payloads acquired data on all kinds of surfaces. Normalized Radar Cross-Section (NRCS) are provided over all surfaces in the L1 products, opening great opportunities for specific studies and potential products upgrade or complement.

The study presented here focuses on sea-ice. Sea-ice represents approximately from 7 to 15% of the ocean surface. Its study is important for observing and modeling the global warming, especially due to the interaction of sea-ice with atmosphere and ocean (Johannessen et al., 2007; Shokr & Sinha, 2015), and economic activities enhancement (fisheries, freight transport). Our goal is twofold: improve wave products in the marginal ice zone and exploit SWIM data over sea-ice to get additional information on its characteristics (extension, thickness, age, etc.). Furthermore, a good discrimination between open water and sea-ice covered areas would allow to further study the evolution of waves in the vicinity of the sea ice pack.

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Thanks to its original viewing angle configuration, SWIM complements other existing concepts such as altimetry, scatterometry or radiometry. As far as scatterometers are concerned, sea-ice was extensively studied, mostly at much larger incidence angles (typically above 30°) to detect and characterize sea ice and icebergs. Girard-Ardhuin and Ezraty (2012) merged passive (SSMI) and active (ASCAT) data into a sea ice drift product using a maximum cross-correlation method. Lots of sea-ice classification works have been performed on ASCAT and QuikSCAT data, especially allowing the distinction between multi-year and first-year ice (MYI and FYI) (Kwok, 2004; Lindell & Long, 2016). The synergy between radiometer and scatterometer allows for even better sea-ice characterization (Lindell & Long, 2016). For example, Zhang et al. (2019) merged ASCAT, QuikSCAT, AMSR and SSMI data on the 2012-2017 period, flagging open water from the radiometer brightness temperatures, and using machine learning algorithms to detect clustering resulting from different sea-ice types over the combination of brightness temperatures and NRCS. Belmonte Rivas et al. (2018) were able to estimate sea ice extents from the combination of ERS, QuikSCAT and ASCAT scatterometer data from 1992 to 2018.

The ice induces a clear change in the radar backscattering, not only in its level, but also in the shape of NRCS profiles with incidence. Near nadir, specular reflexions with facets oriented towards the instrument are considered as the main contributors to the radar backscatter (Barrick & Peake, 1968). From this perspective, open water and sea-ice covered areas NRCS profiles inherently differ: while open water has high coherence roughness, dominated by cm-wave length ocean waves, sea-ice roughness is more randomly distributed and its echoes vanish faster with greater incidence. At nadir, sea-ice reflects more energy than open water surfaces, but this tendency switches at slightly more off-nadir angles, where both the contribution from surface and volume scattering need to be taken into account (Remund & Long, 2003). Especially, the radar backscatter of sea-ice decreases very fast with incidence angle (Giles et al., 2007; Kurtz et al., 2013), while open water surfaces generally maintain a certain number of specular reflexion points due to its generally larger correlation length, especially caused by the wavy nature of its surface geometry. Although most mechanisms building up the sea-ice backscatter are known (Landy et al., 2019), the variability of sea ice physical characteristics make it complicated to summarize the NRCS behavior into a simple parametric model, contrary to the open water case. The analysis of the NRCS profiles is interesting not only for separating sea ice from open water, but also potentially for characterizing the ice itself (in particular ice age), thanks to its relation with sea-ice roughness.

In this study, a Bayesian approach is selected. Sea-ice detection and characterization is particularly suited to machine-learning based algorithms (Alhumaidi et al., 1997; Zhang et al., 2019), but this approach, although quite efficient, has the disadvantage of remaining a black box. More classical approaches are also usual, especially Bayesian ones (Meier & Stroeve, 2008; Belmonte Rivas & Stoffelen, 2011; Breivik et al., 2012; Belmonte Rivas et al., 2012; Lindell & Long, 2016; Otosaka et al., 2018). Most of them consist in designing a statistical estimator building up on the difference between open water and sea-ice GMFs. Over open water, near-nadir Ka/Ku band GMF are now well established both theoretically (Chapron et al., 2000; Boisot et al., 2015) and empirically (Tran et al., 2007; Guerin et al., 2017), especially for missions TRMM (Freilich & Vanhoff, 2003) and GPM (Chu et al., 2012; Nouguier et al., 2016; Gressani et al., 2018; Yan et al., 2019; Hossan & Jones, 2021). In this study, we propose for the first time, a global characterization of the sea-ice backscatter from near-nadir Ku-band observations. Preliminary studies had been performed on GPM data (Mouche et al., 2015), followed by more quantitative results at regional scale (Panfilova et al., 2020; Karaev et al., 2021). The SASS scatterometer onboard Seasat mission is also worth mentioning, operating in Ku-band, from which sea-ice backscatter is characterized (Swift, 1999), but at slightly higher incidence angles than SWIM (above 23°).

The present study analyzes and demonstrates the possibility to distinguish sea-ice from open water based on SWIM data solely at the range gate resolution. After the description of data used for the study in section 2, the algorithm is described and illustrated over a test case in section 3. The resulting flag is then compared to other reference data and discussed in section 4.

2 Data

2.1 SWIM data

SWIM NRT level 1 (L1) and 2 (L2) products provided by the CWWIC are used for the present study (Hauser et al., 2021). A description of SWIM products used for this study is provided in Figure 1. Calibrated and georeferenced NRCS for each off-nadir beam (2, 4, 6, 8 and 10° central incidence) are delivered in the L1a products. The L1a algorithm inverts the radar equation to compensate for all instrument gain and losses and for the radar geometry. The mean thermal noise is also accounted in this inversion. The NRCS is then computed for each radar gate in all available azimuths. Note that along this study, all NRCS values are corrected from atmospheric effects (both L1a and L2). Atmospheric corrections are also embedded in SWIM auxiliary products. At a given beam central incidence θ_i , incidence angles range from $\theta_i - 2^\circ$ to $\theta_i + 2^\circ$, except for the 10° beam, for which the incidence only runs up to 11°. The NRCS values analyzed here are provided within the 3 dB footprint in elevation around θ_i . The theoretical radar range resolution is 0.47 m (see Hauser et al. (2017) or Hauser and Tison (2018)), which correspond, after the on-board processing, which includes range gates averaging, to the horizontal resolutions given in Table 1. Due to oversampling applied in the real-time on-board processing, this is 1.25 times greater than the range pixel spacing of the L1 products used here.

Table 1. Table of theoretical horizontal resolutions for each beam central incidence.

Incidence angle θ_i	Horizontal resolution
2°	$53.7\mathrm{m}$
4°	$26.9\mathrm{m}$
6°	$9.0\mathrm{m}$
8°	$10.1\mathrm{m}$
10°	$8.1\mathrm{m}$

Note that these NRCS observations although discretized at a fine scale in the elevation direction, are integrated over about 18 km in the direction perpendicular to the elevation plane. Alternatively, the L2 products proposes a simplified information which consists in mean profiles of NRCS with incidence, estimated as averaged values over incidence bins of 0.5 and azimuth bins of 15 over the geographical boxes presented in Figure 1. This variable is called *mini-profile*. L1a products being quite expensive in storage space, L2 *mini-profiles* are more convenient to handle when dealing with data sets spanning several months or years. Here, when dealing with L2 profiles, we work on NRCS profiles averaged azimuthally over 15° and in incidence over 0.5° .

NRCS profiles come together with Earth's surface state variables such as land-sea mask LSM, significant wave height H_S , 10 m wind speed U_{10} , sea-ice concentration SIC and seasurface temperature SST. They are obtained from colocation and interpolation from a 6-hourly ECMWF high resolution forecast, with approximately 9 km resolution. For the present study, version 5.1.2 processed SWIM data is used both for the GMFs computation and the flag validation. Only data from March 2020 to March 2021 are used here. Particular subsets are used either for GMF design or for flag validation. This is precised further down at each step.



Figure 1. Schematic of SWIM products used in this study. SWIM footprint (left panel) is an approximately 18 km diameter disc, with range incidences centered on the beam nominal incidence, and spanning $\pm 2^{\circ}$ at most around it. Such NRCS profile acquisition sequence is called *cycle*. Then (center panel), each beam is illuminated successively in a sequence called *macrocycle*. Along SWIM path, macrocycles are repeated, leading to such cycloid-type periodic pattern. SWIM swath is periodically divided into 70×88 km boxes, over which NRCS profiles averages (right panel), the so-called *mini-profiles* are performed at level 2.

2.2 Other data

Micro-wave radiometer derived sea-ice concentration data from SSMI are used as reference for the validation of the sea-ice flagging algorithm. SSMI constellation covers the Earth's surface in almost 1 day. Especially, high resolution (12.5 km) brightness temperature data are retrieved and processed by Ifremer CERSAT into daily maps of sea-ice concentration over both the Arctic and Antarctic over stereo polar grids (Kaleschke et al., 2001; Ezraty et al., 2007).

Sentinel-1 NRCS extra wide swath images are also used for comparison. Sentinel-1 is a C-band Synthetic Aperture Radar providing high resolution (GRDM products, 40 m) NRCS images in dual polarization at incidence angles between 19 and 47°.

However, due to their limited spatial and temporal resolutions, the previous references must be analyzed with a critical eye, especially in the marginal ice zone, where sea-ice spatio-temporal variability can be significant (drift, melt, break-up by ocean waves, etc.).

3 Algorithm description

The sea-ice flagging algorithm is summarized in Figure 2. By definition, the sea-ice flag is activated when the probability of sea-ice presence P(ice) exceeds 0.5:

$$flag = H[P(ice) - 0.5] \times (1 - LSM), \tag{1}$$

where H is the Heaviside function:

$$H(x) = \begin{cases} 1 & \text{if } x > 0\\ 0 & \text{otherwise} \end{cases},$$
(2)

and LSM is the land-sea mask (1 over land and 0 elsewhere), so that the flag equals 1 in the presence of sea-ice and 0 elsewhere, especially on land. Note that the land-sea border

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Figure 2. SWIM sea-ice flagging algorithm summary. A NRCS and incidence angle pair from L1a products are compared to GMFs both assuming either sea-ice or open water thanks to a likelihood estimate, given the measurement context (wind speed, beam number, etc.). A prior probability based on sea surface temperature is combined to the likelihoods into a Bayesian scheme, from which a probability is derived and translated into a binary flag.

may sometimes not coincide with coastlines. Especially around Antarctica, for example in the Ross and Weddel Sea, glaciers or areas of long-lasting continental ice LSM is set to 1 (continent). The introduction of this LSM correction prevents them to be detected as open water.

Sea-ice flagging can be formulated as a Bayesian hypothesis testing problem. The probability of presence of sea-ice is decomposed into conditional probabilities. Given a value σ of the NRCS, the probability of sea-ice presence is $p(\sigma|\text{ice})$. Assuming that the only alternative to *sea-ice* is *open water*, the maximum a posteriori estimator of sea-ice presence is:

$$P(\text{ice}) = \frac{p(\text{ice}|\sigma)}{p(\text{ice}|\sigma) + p(\text{water}|\sigma)},$$
(3)

which compares the two alternative posterior probabilities corresponding other to *sea-ice* or *open water*. The log-likelihood is defined by

$$\mathcal{L} = \ln \left[\frac{p(\text{ice}|\sigma)}{p(\text{water}|\sigma)} \right]$$
(4)

so that

$$P(\text{ice}) = \frac{1}{1 + e^{-\mathcal{L}}},\tag{5}$$

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and, using Bayes rule

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$$p(\mathbf{x}|\sigma) \propto p(\sigma|\mathbf{x})p(\mathbf{x}),$$
 (6)

where x is the hypothesis of either sea-ice or open water presence, we have

$$\mathcal{L} = \ln \left[\frac{p(\sigma | \text{ice})}{p(\sigma | \text{water})} \right] + \mathcal{L}_{\text{prior}},\tag{7}$$

where $p(\sigma|ice)$ and $p(\sigma|water)$ are the likelihoods and the prior log-likelihood is

$$\mathcal{L}_{\text{prior}} = \ln\left[\frac{p(\text{ice})}{p(\text{water})}\right] \tag{8}$$

where p(ice) and p(water) are the a priori probabilities. Comparing P(ice) to 0.5 in order to flag sea-ice contaminated data is equivalent to comparing \mathcal{L} to 0.

The Bayesian approach allows for the proper accounting of an a priori knowledge. Indeed, in order to avoid spurious sea-ice detection at low latitudes, a latitude dependent prior can be defined, that will be discussed in section 3.1. Then, the likelihood can be estimated from an empirical knowledge of the NRCS distribution (GMF). Here, we propose to first parametrize the mean behavior of NRCS (GMF), then quantify the uncertainty that needs to be allowed for an efficient and consistent flag definition.

3.1 Prior probability



Figure 3. Prior log-likelihood, with T = 276 K and $\Delta T = 1$ K (see equations (8) and (10)).

The chosen prior formulation is based on the SST. As the ECMWF forecast SST field is available in SWIM products, a simple prior can be derived based on a pseudo melting criterion, at temperature T set to 276 K. Allowing for an arbitrary standard deviation of $\Delta T = 1$ K for the ECMWF SST, and assuming Gaussian statistics, the a priori probability of sea-ice presence writes:

$$p(\text{ice}) = \frac{1}{2} \int_{-\infty}^{T} \frac{dT}{\sqrt{2\pi}\Delta T} \exp\left[-\left(\frac{\text{SST} - T}{\sqrt{2}\Delta T}\right)^2\right].$$
(9)

The 1/2 factor constrains the prior between 0 (no ice) and 1/2 (ice and open water equiprobability). (9) can be expressed with the error function

$$p(\text{ice}) = \frac{1}{4} \left[1 + \operatorname{erf}\left(\frac{276 \,\mathrm{K} - \mathrm{SST}}{\sqrt{2}\Delta T}\right) \right]. \tag{10}$$

For consistency, the prior probability of open water is defined as the complementary of the sea-ice probability:

$$p(\text{water}) = 1 - p(\text{ice}). \tag{11}$$

The corresponding prior log-likelihood is plotted in Figure 3 from equation (8). This a priori does not favor the detection of sea-ice but only penalizes it at low latitudes, where the SST exceeds 276 K, thus preventing spurious sea-ice detections.

3.2 Practical log-likelihood estimate

3.2.1 Statistical model

A log-normal distribution is assumed for the NRCS likelihood:

$$p(\sigma|\mathbf{x}) = \frac{\lambda}{\sigma\sqrt{2\pi}\Delta\sigma_{\mathbf{x}}} \exp\left\{-\left[\frac{\mathrm{dB}(\sigma) - \mu_{\mathbf{x}}}{\sqrt{2}\Delta\sigma_{\mathbf{x}}}\right]^{2}\right\},\tag{12}$$

where

$$dB(\sigma) = 10\log_{10}(\sigma) = \lambda \ln \sigma, \tag{13}$$

and $\lambda = 10/\ln 10 \simeq 4.3$. Although the log-normal distribution is a crude simplification of the empirical distributions, it seems to be a convenient choice in our case, with a good compromise between suitability and complexity. First, a normal distribution cannot be chosen, as NRCS values are defined only over positive reals. This assumption breaks down when looking at NRCS distributions in dB, which is the reason why a log-normal assumption seems reasonable. Moreover, the log-normal assumption is quite convenient in our case as it is based on only two parameters, μ_x and $\Delta \sigma_x$ which relate directly to NRCS mean and standard deviation.

Hence, the log-likelihood (7) writes:

$$\mathcal{L} = \left[\frac{\mathrm{dB}(\sigma) - \mu_{\mathrm{water}}}{\sqrt{2}\Delta\sigma_{\mathrm{water}}}\right]^2 - \left[\frac{\mathrm{dB}(\sigma) - \mu_{\mathrm{ice}}}{\sqrt{2}\Delta\sigma_{\mathrm{ice}}}\right]^2 + \ln\left(\frac{\Delta\sigma_{\mathrm{water}}}{\Delta\sigma_{\mathrm{ice}}}\right) + \mathcal{L}_{\mathrm{prior}}.$$
 (14)

3.2.2 Parameters estimation from data

Parameters appearing in the log-likelihood (14) are evaluated from ancillary meteorological data available in SWIM products. Their behavior must be constrained from past observations, under the so-called GMFs. They are estimated from NRCS average mini profiles (L2 products) in linear units:

$$\sigma_{\mathbf{x}} = \mathbf{E}(\sigma | \mathbf{x}),\tag{15}$$

where E is the expectation, estimated itself by ensemble average. Their empirical estimation and parameterization are detailed in section 3.4.

From now on, parameters Δ_x and μ_x can be estimated using the method of moments. First, the standard deviation in dB is:

$$\Delta \sigma_{\mathbf{x}} = \lambda \sqrt{\ln \left[1 + \frac{\operatorname{Var}\left(\sigma|\mathbf{x}\right)}{\operatorname{E}\left(\sigma|\mathbf{x}\right)^{2}}\right]},\tag{16}$$

with the mean $E(\sigma|\mathbf{x})$ and variance $Var(\sigma|\mathbf{x})$ being estimated from the L2 mini-profiles averages and standard deviation. This shortcut is acceptable as long as both mean and variance variability within each 0.5° mini profile incidence bin can be neglected, which is assumed here. Second, the $\mu_{\mathbf{x}}$ parameter, which does not only depend on the GMF, is expressed in dB as:

$$\mu_{\mathbf{x}} = \mathrm{dB}(\sigma_{\mathbf{x}}) - \Delta \sigma_{\mathbf{x}}^2 / (2\lambda). \tag{17}$$

3.3 Maximum Likelihood combination of neighboring flag estimates

For both visualization and application purposes, it is useful to have an optimal mean of combining neighboring estimates of the L1a flag. Indeed, a profile as provided by the L1 product is composed of no less than a few hundreds or thousands of NRCS measurements. For the present case, they are hard to plot on a single figure, and the comparison with lower resolution external references may not be adequate. Assuming homogeneity of the surface for N neighboring SWIM measurements (for example the N points of a profile), a maximum likelihood estimator of the combined flag is defined similarly to (5) :

$$P_{\{1\dots,N\}}(\text{ice}) = \frac{p(\text{ice}|\sigma_1,\dots,\sigma_N)}{p(\text{ice}|\sigma_1,\dots,\sigma_N) + p(\text{water}|\sigma_1,\dots,\sigma_N)}.$$
(18)

Assuming the σ_i to be i.i.d. random variables,

$$p(\mathbf{x}|\sigma_1, ..., \sigma_N) = \prod_{j=1}^N p(\mathbf{x}|\sigma_j).$$
(19)

Then, (18) can be rewritten as

$$P_{\{1\dots,N\}}(\text{ice}) = \frac{1}{1 + e^{-\mathcal{L}_{\{1\dots,N\}}}},\tag{20}$$

where

$$\mathcal{L}_{\{1\dots,N\}} = \sum_{j=1}^{N} \mathcal{L}_j,\tag{21}$$

and

$$\mathcal{L}_j = \ln\left[\frac{p(\text{ice}|\sigma_j)}{p(\text{water}|\sigma_j)}\right]$$
(22)

is the log-likelihood of measurement j. The sea-ice flag for the set $\{1, ..., N\}$ is then

$$flag_{\{1,\dots,N\}} = H(\mathcal{L}_{\{1,\dots,N\}}).$$
(23)

In practise, the mean log-likelihood is preferred, as it should only weakly depend on N:

$$flag_{\{1...,N\}} = H(\bar{\mathcal{L}}_{\{1...,N\}}),$$
 (24)

with

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$$\bar{\mathcal{L}}_{\{1...,N\}} = \frac{1}{N} \mathcal{L}_{\{1...,N\}}.$$
(25)

As a consequence, assuming homogeneity of the scene for a set of NRCS measurements, an optimal flag estimator (in the sense of maximum likelihood estimators) is obtained by averaging the individual log-likelihoods. Since the averaging operation is linear, one can even consider higher-level combinations, for example combinations between beams with different incidences, without having to know the detail of all the individual measurements log-likelihoods.

3.4 Geophysical Model Functions (GMF)

In the present algorithm, GMFs constrain flag estimates based on prior empirical distributions that are assumed to be stable along time. For this reason, the data base used for their estimation must be representative of the flag usage : it will be used all along the year and over both the northern and southern hemispheres. This data base is composed of all mini-profiles at latitudes above $50^{\circ}N$ and below $50^{\circ}S$ from March 1st 2020 to March 1^{st} 2021. The hypothesis open water or sea-ice is derived from the auxiliary meteorological files sea-ice concentration (SIC). The open water case corresponds to SIC = 0 while sea-ice is assumed when SIC > 0.9. In both cases, the land-sea mask is checked to be 0 (ocean).

3.4.1 Open water GMF

The L2 mini-profile GMF is plotted as a function of wind speed and incidence angle in Figure 4, from equations (15) and (16).

Part of the observed NRCS distribution deviation can be explained by an additional sea-state dependency. Sea state is usually measured by the significant wave height H_S , conventionally defined as 4 times the standard deviation of the elevations. A mean H_S related NRCS deviation can be defined by

$$\delta = \mathrm{dB}\left[\mathsf{E}(\sigma|U_{10}, H_S)\right] - \mathrm{dB}\left[\mathsf{E}(\sigma|U_{10})\right].$$
(26)

and plotted in Figure 5. The sea-state dependency increases with decreasing wind speed. Its relative contribution reaches up to 3 dB at low wind speeds and becomes negligible towards larger winds. Although not shown here, attempts to derive a physics-based NRCS parameterization as a function of H_S proved usefulness at low wind speeds, but introduced an artificial bias at larger ones due to fitting errors. Consequently, the GMF explicit dependency on H_S is dropped in this study. This effect however appears in the residual GMF standard deviation (16).





Figure 4. Distribution of L2 NRCS over open water and sea-ice as a function of incidence angle for each SWIM off-nadir beam. The nadir beam averages are plotted at 0.5° incidence. For the open water case, the GMF is further detailed by wind speed. The dots are GMF estimates from data using (15). Their least-squares fit for open water and for sea-ice is plotted in solid lines. The colored sector around the GMF corresponds to the $\pm 0.5\Delta\sigma_x$ interval. Its least-squares fit is plotted in thin dashed lines.

Following the Geometric Optics formulation (Boisot et al., 2015; Nouguier et al., 2016), the near-nadir NRCS dependence with incidence angle over open water is known to follow:

$$\sigma_W = \frac{|R|^2}{\cos^4 \theta \text{mss}_{\text{eff}}} \exp\left(-\frac{\tan^2 \theta}{\text{mss}_{\text{eff}}}\right)$$
(27)

More sophisticated models (Chapron et al., 2000; Boisot et al., 2015; Nouguier et al., 2016) can be considered, but which only bring improvements at higher incidences and/or lower frequency. Nouguier et al. (2016) used the same kind of parameterization for their analysis of GPM data in Ku band. Here, $|R|^2$ is to be interpreted as a Fresnel reflection coefficient, which by definition should be positive and lower than 1. The mss_{eff} parameter



Figure 5. Distribution of L2 NRCS deviation from equation (26) by H_S category over openwater as a function of incidence angle for all off-nadir beams for 4 different wind speeds.

can be interpreted as a filtered version of the Mean Square Slope (mss) derived from ocean waves spectra (Boisot et al., 2015).

Over open water, the GMF is assumed to depend on U_{10} , the incidence angle θ and the beam number. $|R|^2$, mss_{eff} and $\Delta \sigma_{\text{water}}$ (equation (16)) are empirically parameterized as a function of wind speed and beam number i = 1, ..., 5, with central incidences $\theta_i = 2i$ degrees, inspired by preliminary works on GPM (Yan et al., 2019):

$$|R|^{2} = |R|_{i}^{2} / \left(1 + \rho_{i} e^{-\lambda_{i} U_{10}}\right) - s_{i} U_{10}, \qquad (28)$$

$$mss_{eff} = M_i / \left(1 + \nu_i e^{-\xi_i U_{10}} \right) + t_i U_{10}, \tag{29}$$

$$\Delta \sigma_{\text{water}} = \Delta \Sigma_i + \alpha_i e^{-\beta_i U_{10}^2} + v_i U_{10}. \tag{30}$$

Least-squares fitting coefficients appearing in equations (28), (29) and (30) are summarized in tables 2, 3 and 4, and plotted in Figures 6, 7 and 8. The GMF dependency over azimuth angle relative to the wind direction is not investigated here.

Table 2. Table of coefficients for $|\mathbf{R}|^2$, equation (28).

i	$1 (2^{\circ})$	$2 (4^{\circ})$	$3~(6^{\circ})$	$4 (8^{\circ})$	$5 (10^{\circ})$
$ R _{i}^{2}$	0.59	0.53	0.67	0.68	0.68
$ ho_i$	0.45	0.24	0.33	0.35	0.43
$\lambda_i [(\mathrm{m/s})^{-1}]$	0.56	0.53	0.51	0.52	0.51
$s_i [(m/s)^{-1}]$	0.012	0.010	0.0073	0.0081	0.0070

Table 3. Table of coefficients for mss_{eff} , equation (29).

i	$1 (2^{\circ})$	$2 (4^{\circ})$	$3~(6^{\circ})$	$4 (8^{\circ})$	$5 (10^{\circ})$
M_i	0.014	0.013	0.012	0.033	0.033
$ u_i$	$\frac{2.6}{2.6}$	$\frac{1.2}{1.2}$	0.022	0.86	0.60
$\xi_i [({ m m/s})^{-1}]$	0.59	0.55	0.23	0.37	0.37
$t_i [(m/s)^{-1}]$	0.00070	0.00077	0.0025	0.0015	0.0015

i	$1 (2^{\circ})$	$2 (4^{\circ})$	$3~(6^{\circ})$	$4 (8^{\circ})$	$5 (10^{\circ})$
$\Delta \Sigma_i [dB]$	0.44	0.48	0.45	0.44	0.44
$\alpha_i [\mathrm{dB}]$	$\frac{1.7}{1.7}$	1.0	0.66	0.87	0.96
$\beta_i \left[(\mathrm{m/s})^{-2} \right]$	0.057	0.061	0.098	0.11	0.089
$v_i \left[dB/(m/s) \right]$	0.0050	0.0067	0.0046	0.0085	0.0095

Table 4. Table of coefficients for $\Delta \sigma$ in dB over open water, equation (30).



Figure 6. $|R|^2$ parameter estimated from least-squares fitting of average NRCS profiles (15) (dots) and parameterization (28) (solid line) for SWIM off-nadir beams.



Figure 7. mss_{eff} parameter estimated from least-squares fitting of average NRCS profiles (equation (15)) (dots) and parameterization (equation (29)) (solid line) for SWIM off-nadir beams.



Figure 8. Standard deviation over open water in dB estimated from least-squares fitting of average NRCS profiles (15) (dots) and parameterization (equation 30)) (solid line) for SWIM offnadir beams.

To our knowledge, it is the first time that near-nadir NRCS profiles are estimated at various incidences independently. In the past, studies had been lead on Ku-band near-nadir scatterometers such as TRMM (Freilich & Vanhoff, 2003) or GPM (Nouguier et al., 2016). In both studies, incidences span the whole range from 0 to 18° . Effective mss estimates from these studies are in agreement with ours (Figure 7), especially with the 6° beam estimates. This evolution of the effective mss with the nominal incidence is also expected due to diffraction/curvature effects (see Appendix A, equation (A9)). Also, the fact that the mss does not fall to 0 at 0 m/s wind speed can be understood through the effect of the remaining surface waves, which is more important at low wind speeds (see Figure 5). As far as the reflection coefficient is concerned, only Freilich and Vanhoff (2003) and Chen et al. (2018) provided such estimate from TRMM data. The agreement with our figure is satisfying for the sea-ice flagging, especially again with the 6° beam. Diffraction/curvature effects predict a slight increase of the fitted reflection coefficient (see Appendix A, equation (A10)), on top of which inter-beam biases may come into play to explain the results of Figure 6. Especially, a positive inter-beam bias is visible in Figure 4 on the 2° that can explain the non-uniform evolution of the reflection coefficient with incidence in Figure 6. A NRCS overestimation by the 2° is also observed by Hauser et al. (2021) (their figure 10a) and Ren et al. (2021) (their table 2). Finally, the standard deviation $\Delta\sigma$ is here presented for the first time over open water. The increase of spread at low wind speeds corroborates the higher sea-state dependency observed in Figure 5. Then, the spread is essentially independent of incidence over open water, except for the 2° beam at low wind speeds.

3.4.2 Sea-ice GMF

The sea-ice backscatter is commonly modelled as the sum of surface scattering and a volume scattering terms (Remund & Long, 2003). Over sea-ice, exponential correlation is often assumed, which leads to a surface scattering of the type (Hagfors, 1970; Kurtz et al., 2014):

$$\sigma_{\rm surf} \propto \left(1 + \gamma \sin^2 \theta\right)^{-3/2},\tag{31}$$

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where γ (constant) can be interpreted as an angular scattering efficiency. The volume scattering term is (Remund & Long, 2003):

$$\sigma_{\rm vol} \propto \cos \theta,$$
 (32)

which is essentially constant at near-nadir incidences. The combination of a surface scattering (31) with a volume scattering term (32) is fitted over off-nadir average miniprofiles of Figure 4. In addition, a nadir correction is inspired from Giles et al. (2007):

$$\sigma_{\rm nadir} \propto e^{-(\theta/\theta_{\rm pr})^2},\tag{33}$$

which is fitted over the whole range of incidences, including nadir. Finally, the full sea-ice GMF reads:

$$\sigma_I = A \left(1 + \gamma \sin^2 \theta \right)^{-3/2} + B \cos \theta + C e^{-(\theta/\theta_{\rm pr})^2},\tag{34}$$

with variables numerical values obtained from the fitting procedure gathered in table 5. It is plotted in Figure 4. Although these coefficients are derived from physical parameterizations, it does not seem possible to interpret them in terms of the physical properties of the sea ice: the data on which the GMF was estimated includes multiple types of sea ice (e.g. FYI, MYI), and this effective GMF implicitly incorporates their statistical distribution in the training database. Similarly to the open water case, an empirical parameterization is found for the sea-ice standard deviation in dB:

$$\Delta\sigma_{\rm ice} = a_i \tan^2 \theta + b_i. \tag{35}$$

Least-squares fitting coefficients a_i and b_i are gathered in table 6 and the resulting parameterization plotted in Figure 9.

Table 5. Table of coefficients for SWIM GMF over sea ice, equation (34).

A	γ	B	C	$\theta_{\rm pr}[^\circ]$
17.2	401	1.4	202	0.7

Table 6. Table of coefficients for $\Delta \sigma$ in dB over sea-ice, equation (35).

i	$1 (2^{\circ})$	$2 (4^{\circ})$	$3~(6^{\circ})$	$4 (8^{\circ})$	$5(10^{\circ})$
$a_i \left[\mathrm{dB} \right]$	-70	-9.8	-31	-16	-4.7
$b_i [\mathrm{dB}]$	2.2	2.0	1.8	1.5	1.2



Figure 9. Standard deviation over sea-ice in dB estimated from least-squares fitting of average NRCS profiles (equation (15)) (dots) and parameterization (35) (solid line) for SWIM off-nadir beams.

3.5 Test case: Weddell Sea

A test case is selected over an open water sea-ice transition in the marginal ice zone offshore of the eastern Weddell Sea on October 7th 2020, when the sea-ice extent is usually maximum around Antarctica (see Figure 10a). CFOSAT sampled this region around 07 : 30 UTC while Sentinel-1 flew over it approximately 13 hours later. The HV channel NRCS of Sentinel-1 (Figure 10b) is used as a qualitative ground-truth illustration for the situation at that time, which enables direct visualization of the scene measured by SWIM. Especially, from this SAR image, sea-ice appears very variable at that location, with the presence of heterogeneous sea-ice meso-scale features and melting sea-ice

First, the presence of open water and sea-ice on both sides of the margin are correctly identified, according to the sign of the beam level log-likelihood (Figure 10g). Its color (blue or red) corroborates the observed similarity between each of the GMFs (Figure 10e and f) and the NRCS (averaged over profiles). The SWIM L1 NRCS profile corresponding to beam 10° (see red square in Figure 10d-g) is analyzed across an open-water/sea-ice transition (detected by the SAR data) in Figure 10c. Between 9 and 9.8° incidence, the profile is quite stable and close to the open water GMF. Consistently, the estimated probability of sea-ice presence is constantly equal to 0, indicating an open-water area. The transition to sea-ice occurs in two steps: after a sharp decrease of the NRCS (black curve) around 10° incidence, the NRCS undergoes a second decrease after 10.7° incidence with a steepest slope. The algorithm should flag these two areas as sea-ice, with slightly lower probability over the $10-10.7^{\circ}$ interval.

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Figure 10. Test case at an open water/sea-ice transition. (a) Location of the center of the scene (red star). (b) Sentinel-1 NRCS in HV polarization. (c) L1 NRCS profile example. The raw NRCS profile is plotted in black (with y-axis labels on the right), together with the open water (red) and sea-ice GMF (blue), with their standard deviations underlayed (light blue and light red). The sea-ice probability is plotted in green (y-axis labels on the left). The dashed line, represents the limit P(ice) = 0.5. At that geographical location $U_{10} = 4.7 \text{ m/s}$ and $H_S = 2.5 \text{ m}$. The red arrow link the subfigure (c) with the subfigure (d) to show which position is plotted. (d) L1 NRCS averaged over each profile (1 value per azimuth position of a beam). The colors show clearly two classes (ice / no ice) matching the boundaries of the S1 images. (e) Open water GMF averaged over each profile(1 value per azimuth position of a beam). (f) same as (e) for sea-ice (1 value per azimuth position of a beam). (g) Profile level log-likelihood, equation (25).

4 Flag qualification

In the absence of fully reliable ground truth simultaneously colocated with SWIM, the qualification of the proposed flag against other sea-ice data must be carried out with a critical eye. Here, sea-ice flag references are derived from sea-ice concentration maps. As there is no consensual sea-ice concentration threshold for this flagging, this variability is also explored in what follows. A first reference is the one embedded in SWIM NRT products, *i.e.*, ECMWF ice concentrations. A second reference is the SSMI sea-ice concentration products processed at Ifremer (see section 2.2).

Sea-ice is flagged in SWIM L1a profiles for the period between Jan. 15^{th} 2021 and Jan. 28^{th} 2021 included, this period corresponding to SWIM cycle 63 at latitudes above $40^{\circ}N$ and below $40^{\circ}S$. It is interesting noticing that the data used for the GMF derivation in section 3.4 are taken from the L2 products, while they are here taken from L1a products. The latitudes are limited in order to focus on sea-ice covered regions. The latitude limit is lowered to $40^{\circ}N$ in order to include the Caspian and Black Sea which undergo a sea water freezing episode at that time of the year. Nevertheless, the sea-ice flag design, through the inclusion of a latitude dependent a priori prevents any spurious sea-ice detection at low latitudes. L1 flags are combined according to the method in section 3.3 before being compared to references. References are colocated with SWIM profiles using a nearest neighbor algorithm. In January, situations are quite contrasted between the poles: Arctic is characterized by intense sea-ice formation while the melting is already quite advances around Antarctica. Thus, sea-ice types differ, with quite young and dry ice in the Arctic marginal ice, and older wet sea ice around Antarctic.

Table 7. Confusion matrix for the 8° beam profile level flag 24 compared to SSMI forecast with 0.5 concentration threshold for latitudes above $40^{\circ}N$ and below $40^{\circ}S$ for SWIM cycle 63. The number of data points entering in each category is represented in the table.

SWIM flag SSMI flag	0	1
$\begin{array}{c c} 0 & (\text{no sea-ice}) \\ 1 & (\text{sea-ice}) \end{array}$	TN = 2309896 FN = 33329	FP = 29802 $TP = 411407$

A convenient way of qualifying a classification against a ground truth data sets is the confusion matrices formalism. For a given sea-ice threshold, a confusion matrix provides an inventory of the profiles according to their labels with respect to SWIM on the one hand and the reference data set on the other hand. An example of such a matrix is provided in table 7 for the 8° beam. Sea-ice detections are qualified of positive: these are either True Positives (TP) when sea-ice is detected by both methods, and False Positives (FP) when sea-ice is only detected by SWIM. Cases when no sea-ice is present are referred as negatives: True Negatives (TN) when both methods are in agreement, and False Negatives when sea-ice is detected by the reference but not by SWIM. If the reference and SWIM agree perfectly, only the matrix diagonal elements are non-zero. Looking at confusion matrix in table 7, the largest populations are found along diagonal elements. According to SSMI with a threshold at 0.5 sea-ice concentration, 15.97% of the echoes above 40° latitude are classified as sea-ice, against 15.84% for SWIM. The false negative rate is the percentage of spurious open water detections by SWIM when considering the ground truth as reference :

$$FNR = \frac{FN}{TP + FN}.$$
(36)

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In the present case, this situation occurs for 7.49% of the macrocycles over sea-ice. The false negative rate is an analog to the false positive rate when considering open water instead of sea-ice :

$$FPR = \frac{FP}{FP + TN}.$$
(37)

Here, 1.27% of the open water macrocycles are mistakenly labeled as sea-ice considering SSMI as reference. Although these rates are quite meaningful, neither of them combines all the confusion matrix elements into a single number. Among others, the accuracy metrics seems adapted:

$$ACC = \frac{TP + TN}{TP + TN + FP + FN}$$
(38)

which quantifies the rate of agreement between classifications. The accuracy can be thought as an optimal indicator of the agreement between two classifications: a 100% accuracy corresponds to a perfect agreement between both data sets, and hence to zero off-diagonal elements. Here, it reaches 97.73%. At that level, the remaining 2.27% of data points where SWIM and SSMI disagree is hard to distinguish with separate figures. This is the reason why alse Positives and False Negatives were highlighted on figures 14 and 15. It can be concluded from a first look at classification results that both SWIM and SSMI are in excellent agreement overall.

Although the previous example is quite representative of SWIM sea-ice flag performances, it is interesting to look in details at the sensitivity of the previous metrics, defined in equations (36), (37) and (38), to the sea-ice concentration threshold, beam by beam, and for both reference data sets, as it is presented in Figures 11 and 12. Overall, performances are similar from beam to beam. Looking at the accuracy, it is worth noting that the accuracy of SWIM sea-ice flag is greater with SSMI than with the ECMWF derived flag (currently in SWIM product), thus confirming the added value of such a flag compared to the current one. Now looking at disagreements between data set in Figure 11, it is interesting looking at their sensitivity to sea-ice concentration. Looking at the FNR slope, it can be seen that the sensitivity to the presence of water among sea-ice is decreased with increasing beam nominal incidence : the 10° beam is more inclined to detect low sea-ice concentration areas than is the 2° beam. On the other hand, the false positive rate slightly increases with incidence: the 10° beam could also be more inclined to spurious sea-ice detections.

Back to Figure 12, it is interesting to look at the SIC threshold for which the accuracy is maximum (vertical solid and dashed lines). This information is summarized in table 8. First, as was already observed previously, the overall agreement between SWIM and SSMI is better than with SWIM and ECMWF forecast: accuracies are all greater in the first case. Second, the SIC threshold of best agreement decreases with the beam incidence angle. This corroborates the already observed increased sensitivity of beams with incidence. Second, the optimum SIC threshold goes down to 0.10 for the 10° beam. It must be noticed that the sea-ice category was defined during the GMF conception with a threshold of 0.9 on the colocated sea-ice concentration (see section 3.4), thus not guarantying intermediate concentrations to be properly accounted for.

At that stage, validity limits of the reference data sets are probably reached. The disagreement between SWIM and the references could not only be caused by spurious detections, but also probably reaches the limit of reference data sets. SWIM profile level flag results are compared to a flag built from SSMI data with a concentration threshold set at 1%, both at the North (Figure 14) and South poles (Figure 15). A first visual inspection informs that most of the sea-ice pack is indeed detected as such on all beams (blue dots), both in Arctic and Antarctic. Another look at the data shows the relatively low number of



Figure 11. Probability of disagreement between SWIM and ECMWF classification (dashed lines) or SSMI sea-ice (solid lines) reference data sets for latitudes above $40^{\circ}N$ and below $40^{\circ}S$. These disagreements can be decomposed into False Negatives (crosses, decreasing with SIC), equation (36), and into False Positives (dots, increasing with SIC), equation (37)



Figure 12. Same as Figure 11 for accuracy, eq. (38). The vertical lines locate the SIC value corresponding to the maximum accuracy for each reference data set (dashed: ECMWF, solid: SSMI).

spurious detections at low latitudes (confirmed by a global map of flagged data), thanks to the definition of the prior and the application of the land-sea mask.

False negatives (sea-ice areas missed by SWIM) are colored according to their seaice concentration in the SSMI product. Most of them correspond to low ice concentrations (purple patches), but some locations with higher concentration are missed. These are located either along the coastlines or occasionally inside the ice pack for low incidence beams. False negatives close to the coastlines can be attributed to the presence of ice shelves near the continents, which radar signature may be further from the one of sea-ice than from the one of open water. For the second kind, this can be explained by the presence of melting sea-ice in the South pole. Indeed, at that time of the year, the sea-ice extent in the Arctic is still growing, while it is almost at its minimum in the Antarctic. Thus, sea-ice types differ in Figures 14 and 15, with a tendency to freezing in the North and to melting in the South. Melting sea-ice, as it mixes with water, may have a less distinct signature from the one of open water.

Looking now at spurious sea-ice detections (golden triangles), two main observations can be drawn. First, most of these detections occur in the marginal ice zone (close to Svalbard in the North), where low sea-ice concentrations are likely to occur. Second, the number of these detections increases with the beam nominal incidence. The case of the Amundsen Sea and Ross Sea (90 to $180^{\circ}W$) is particularly instructive. As the number of false positives increases, the number of false negatives decreases as well, especially for the 8° and 10° beams, meaning that these false positives, may not be all spurious and may contain some signal, which is complementary to SSMI. Finally, coherent structures (present on 6 and 10° beams) appear offshore of Antarctica, which, at that time of the year, could correspond to icebergs.

It is worth noting that there is an increased sensitivity to sea-ice with beam incidence. Two main reasons could probably be at its roots. First, the sensitivity to sea-ice detection decreases with incidence (above 2 degrees), as can be seen in Figure 13. The average NRCS ratio in dB

$$\mu = \mu_{\rm ice} \mu_{\rm water} \tag{39}$$

is plotted in Figure 13 at various wind speeds. On the other hand, the sea-ice response gets closer and closer to the noise level with increased central incidence. Two effects must be considered. The figure shows the first one, namely the theoretical separability of the open water/sea ice echoes. We can take the comparison of the 4 and 10° beams as an example. At 10°, the mu ratio is lower than at 4°. For this reason, the open water/sea ice transition will be more pronounced at 4° than at 10°, thus limiting the number of false detections at 4°. The analysis in Figure 13 does not take into account the noise level of each beam. Again, it can be argued that the detection quality is better at 4 degrees than at 10° because of the higher signal to noise ratio at low incidence. However, this analysis can be weighted by observing that the wind dependence of the mu ratio is more important at low incidence, which suggests different performances with different wind speeds. Further studies need to be conducted in order to better characterize it. Further studies need to be conducted in order to better characterize it.



Figure 13. Average sea-ice/open water NRCS ratio in dB from equation (39) at various wind speeds (35) (solid line) for SWIM off-nadir beams.

5 Conclusion

A Bayesian approach is proposed to discriminate sea-ice from open water in near-nadir microwave NRCS observations. It is specifically designed for Ku-Band SWIM/CFOSAT data sets, but intrinsically, the method is more general. The main principle is to compare in a statistical sense the measured values of NRCS to GMFs which relate the NRCS to wind speed and incidence angle. First the GMF models have been established. Then,



Figure 14. Maps of flagged SWIM profiles location for each beam and comparison with SSMI data with concentrations above 1% around the North pole for cycle 63 (Jan. 2021).



Figure 15. Same as Figure 14 for Antarctica.

the inversion method has been applied to SWIM data. The comparison of its results with model and remote sensing based reference data sets exhibits performances that reach up

Table 8. Values of the maximum accuracy and corresponding SSMI SIC threshold from Figure12. The same results with ECMWF SIC are printed in parentheses.

Beam	2°	4°	6°	8°	10°
SIC threshold Accuracy [%]	$\begin{array}{c} 0.65 \ (0.55) \\ 98.0 \ (97.9) \end{array}$	$\begin{array}{c} 0.50 \ (0.50) \\ 97.9 \ (97.9) \end{array}$	$\begin{array}{c} 0.35 \ (0.40) \\ 97.8 \ (97.9) \end{array}$	$\begin{array}{c} 0.35 \ (0.40) \\ 97.8 \ (97.7) \end{array}$	$\begin{array}{c} 0.10 \ (0.25) \\ 97.7 \ (97.5) \end{array}$

to 98% accuracy. A prior probability based on SST is defined that prevents any spurious sea-ice detection at low latitudes, *i.e.*, where the SST exceeds 276 K. Moreover, this new lag performs better than the existing flag, as it compares more satisfyingly with external remote sensing based data. More quantitatively, at polar latitudes, the present flag is in disagreement with reference data in less than 4% of the time, and without any combination of beam incidences. Sensitivity increases with incidence, and the intermediate incidence beams (4 and 6°) exhibit the highest agreement with reference data sets.

Yet, performances assessment provides interesting insights into the scientific content of SWIM data for sea-ice studies. In addition, there are evidences that SWIM may be able to identify more particular features such as icebergs or particular sea-ice types, that may not be detected by radiometers such as SSMI. Especially, the larger incidence beams (8 and 10°) detect features away from the marginal ice zone that are consistent from a beam to another. In addition, GMFs are derived that can be used for the improvement of SWIM data processing. Also, one could think of properly diagnosing the flag performances seasonality by looking at a whole year of such processing.

Perspectives opened by this study are vast and would benefit of additional validation efforts. At the moment, comparisons have been carried out at the profile level with typically 10 km resolution reference data. The present flag being defined at very high resolution (see Table 1), high resolution ground truths are needed. One can think of optical or SAR images for example. The exact nature of false positive detections at large incidences could be qualified using iceberg data bases (Tournadre et al., 2008).

This algorithm is a base-line algorithm that could be improved. For example, the dependency of the standard deviations on sea-state could be taken into account. Sea-ice GMF could be refined by a proper accounting of sea-ice types (age, thickness, wetness, etc.), and thus its uncertainty diminished. More specifically, the distinction between FYI and MYI would be a first way to account for all these effects. Spatio-temporal dependencies of the GMF can be considered, for example on the hemisphere or seasonal variation. Similarly the azimuthal dependency of the open water GMF is known to exist and could be taken into account. It could also be adapted to GPM Dual Polarization Radar data in the same radar band and at similar incidence angles.(Hauser & Tison, 2018)

Based on this study, the definition of sea-ice products from the SWIM observations is in progress. Combining incidences seems an interesting way to further characterize seaice, that builds up on SWIM acquisition geometry. In a second step, the scatterometer on board CFOSAT could be use in fusion with SWIM, and next generation satellite missions operating near nadir such as SWOT could benefit from this work.

6 Open research

All SWIM data used for this study, e.g. L1A and auxiliary meteorological files, are freely available upon registration on AVISO+ website: https://aviso-data-center.cnes.fr/. Spatio-temporal ranges as well as processing version and variable names are detailed all along the text when needed. Sentinel 1 data are freely available on Copernicus Scihub upon registra-

tion: https://scihub.copernicus.eu/. Colocated NRCS map in Figure 10 is obtained from Level 2 file S1B_EW_OCN__2SDH_20201007T202743_20201007T202837_023711_02D0EF_E580.SAFE without denoising from processing version IPF 3.31. SSMI data for this research are freely available and described in Ezraty et al. (2007). They were accessed upon request to Ifremer CERSAT facility. Sea-ice concentration data sets can be accessed from CERSAT web site: http://cersat.ifremer.fr/oceanography-from-space/our-domains-of-research/sea-ice/radar-backscatterof-sea-ice.

Appendix A Curvature effect with incidence on effective parameters mss_{eff} and $|R|^2$ over open water

For the present study, Geometrical Optics model (27) is assumed. However, this does not account for curvature and diffraction related effects (Guimbard, 2010; Nouguier et al., 2016), for which a more general model is more appropriate:

$$\sigma_W = \frac{n+2}{n+3} \frac{|R|_T^2}{\cos^4 \theta \operatorname{mss}_T(n+1)} \left[1 + \frac{\tan^2 \theta}{\operatorname{mss}_T(n+1)} \right]^{-(n+3)}$$
(A1)

Equation (27) is obtained in the limit $n \to \infty$. *n* is a curvature parameter lying between 1 and 3 in Ku-band (Guimbard, 2010), and mss_T is the total Mean Square Slope, *i.e.*, the one which was measured by Cox and Munk (1954), which inherently differ from mss_{eff}. Around nominal incidence $\theta = \theta_i$, the function $\psi = \ln (\sigma_W \cos^4 \theta)$ can be expanded around $x_i = \tan^2 \theta_i$:

$$\psi(x) \simeq \psi(x_i) + (x - x_i) \left. \partial_x \psi \right|_{x = x_i} \tag{A2}$$

where $x = \tan^2 \theta$. According to geometric optics (27):

$$\psi \simeq \ln\left(\frac{|R|^2}{\mathrm{mss}_{\mathrm{eff}}}\right) - \frac{x_i}{\mathrm{mss}_{\mathrm{eff}}} - \frac{x - x_i}{\mathrm{mss}_{\mathrm{eff}}}.$$
 (A3)

Taking curvature effects (A1) into account leads to:

$$\psi \simeq \ln\left(|R|_T^2 \frac{n+2}{\mathrm{mss}_T(n+1)}\right) - (n+3)\ln\left(1+y_i\right) - \frac{(x-x_i)(n+3)}{\mathrm{mss}_T(n+1)+x_i} \tag{A4}$$

where $y_i = x_i / [(n+1) \text{mss}_T]$. Equations (A3) and (A4) lead to

$$\mathrm{mss}_{\mathrm{eff}} \simeq \frac{x_i + (n+1)\mathrm{mss}_T}{n+3} \tag{A5}$$

and

$$R|^{2} \simeq |R|_{T}^{2} \frac{n+2}{n+3} \left(1+y_{i}\right)^{-(n+2)} \exp\left(\frac{n+3}{1+y_{i}^{-1}}\right)$$
(A6)

More specifically, the evolution of the effective parameters with nominal incidence θ_i can be deduced:

$$\partial_{x_i} \mathrm{mss}_{\mathrm{eff}} = \frac{1}{n+3} \tag{A7}$$

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$$\partial_{x_i} |R|^2 = \frac{|R|^2}{(n+1) \text{mss}_T (1+y_i)^2} \left[1 - (n+2)y_i \right]$$
(A8)

As in near-nadir Ku-band $n \ 1 > 0$ and $y_i \ll 1$, then

$$\partial_{\theta_i} \mathrm{mss}_{\mathrm{eff}} > 0,$$
 (A9)

and

$$\partial_{\theta_i} |R|^2 > 0. \tag{A10}$$

This trend is indeed the one found in Figure 7 for mss_{eff} . For $|R|^2$, it is also found in Figure 6 with the exception of beam 2° whose results may be affected by a relative bias in NRCS between the beams.

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