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Self-propagating waves of crystallization in metallic glasses

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ABSTRACT

Self-propagating thermal waves of the amorphous-crystalline transformation in Fe-based metallic glasses, obtained by melt spinning, were observed using a high-speed infrared camera and reported here. Some experimental results are also reported concerning oscillating waves in the CuTi glassy foils. The thermal characteristics and wave propagating velocities, as well as the microstructure and atomic structure transformations were studied. A comparison of the results with exothermic reaction waves and explosive crystallization shows that the self-propagating waves in metallic glasses are slower and less violent than classical explosive crystallization in deposited films; thus we suggest naming this phenomenon “soft explosive crystallization”. The experimental data was confirmed by molecular dynamics simulation of the crystallization phenomenon.

Metastable amorphous materials can be produced by several routes, including electrolytic or vacuum deposition^{1,2}, melt quenching^{3,4} and mechanical amorphization^{5,6}. Amorphous phases, obtained by rapid cooling of metallic melts, called metallic glasses (MG), possess unique magnetic, mechanical, corrosion resistive, and other properties⁷⁻⁹. Fe-based amorphous alloys, such as Fe-B, Fe-Si-B doped with Cu, Nb, etc., are widely used as soft magnetic materials (e.g., Metglas®)^{10,11}. The Cu-Ti, Cu-Zr, or Cu-Nb based materials show outstanding mechanical properties^{12,13} and also are used as model systems for fundamental studies of atomic structure and ordering in MG¹⁴⁻¹⁶. A combination of several alloying elements resulted in a decrease of the critical quenching rate required to obtain MG from 10^6 K s^{-1} to 1 K s^{-1} , which permits production of bulk MGs for a variety of engineering applications^{9,17,18}. On the other hand, recently, ultrafast quenching at 10^{14} K s^{-1} has allowed for the production of pure glassy metal (tantalum)¹⁹. Thus, metallic glasses have formed a new, rapidly expanding niche in materials science and technologies, and already involve a large number of metals and alloys.

Some deposited amorphous metal or semiconductor films exhibit peculiar phenomena named “explosive crystallization”, i.e. rapid self-propagation of the amorphous-crystalline transformation in the form of a thermal wave^{1, 20-24}. However, this phenomenon has not been observed in metallic glasses produced by melt quenching. Though explosive crystallization was posited as a possible crystallization mechanism in Fe-Zr metallic glasses²⁵, no evidence of self-propagating waves was reported. Recently, we have found self-propagating waves of crystallization in the CuTi glassy tapes²⁶. In the present work, we report the existence of self-propagating crystallization waves in Fe-based amorphous alloys (Fe-B, Fe-Si-B), and present some results on the wave-like crystallization of Cu-Ti metallic glass.

We have studied four compositions of MG tapes produced by melt spinning and are listed in Table 1 (see Supplementary Materials for methods of production and research).

Table 1. Features of the soft explosive crystallization process.

Composition	Infrared Camera measurements			DSC measurements	
	Pre-heating temperature T_0 , K	Max temperature T_{max} , K	Propagating velocity U , cm/s	Heat release Q , J/g	Energy of Activation E , kJ/mol (by Kissinger method)
$Fe_{84}B_{16}$	580	700	7.00 ± 0.04	190 ± 22	250 ± 5
$Fe_{85}B_{15}$	570	720	4.95 ± 0.02	122 ± 13 (sum of two stages)	217 ± 7 (first stage) 252 ± 15 (second stage)
$Fe_{76}Si_{13}B_{11}$	640	740	1.26 ± 0.01	-	-
$Cu_{50}Ti_{50}$	500	640	1.26 ± 0.01	144 ± 7	251 ± 30

Unlike deposited amorphous films, the MGs do not exhibit explosive crystallization (EC) at room temperature. For this reason, the tapes were pre-heated up to 500-640 K by passing an electric current (DC) through the tape in Ar flow. The preheating process to T_0 was accomplished with heating rate > 200 K/min and thus its duration was less than 1 min. Immediately (~ 3 s) after the uniform preheating of the tape, an additional short heat pulse was applied locally at one end of the tape to initiate the amorphous-crystalline transition. Values of the pre-heating temperature T_0 were chosen well below the crystallization onset temperature T_x , which was found from the separate DSC measurements and appeared to be in the range 690 – 720 K for the Fe-B system and 680 K for Cu-Ti system. The lowest possible values of T_0 and short pre-heating time prevented partial crystallization (devitrification) of the samples during pre-heating stage. The EC wave, which is visible in infrared radiation, self-propagates along the tape (fig. 1 and multimedia view).

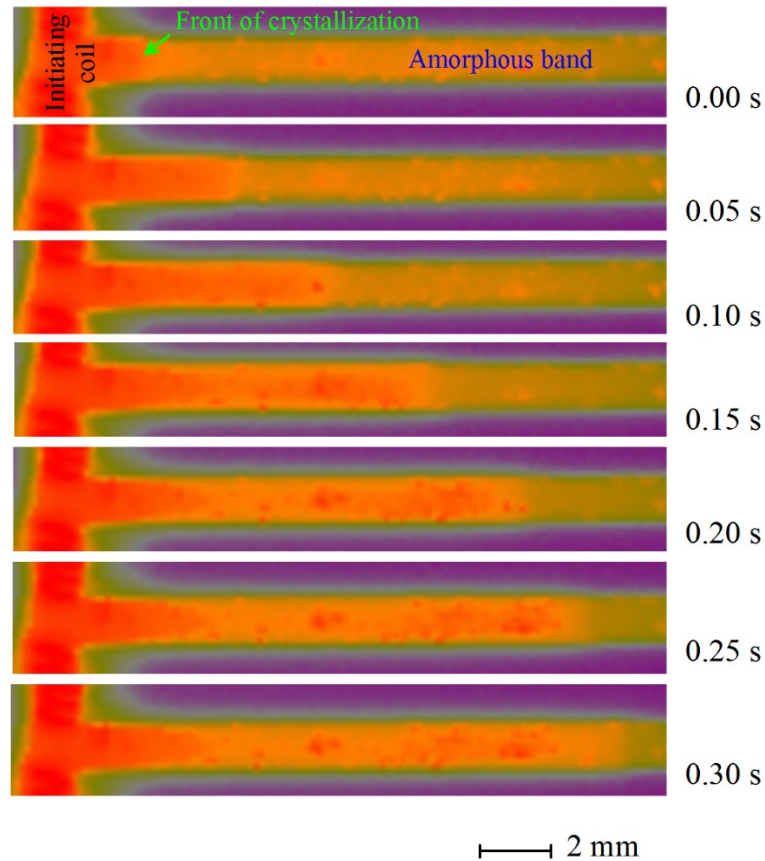


Figure 1. Infrared images of the propagating crystallization wave in $\text{Fe}_{84}\text{B}_{16}$ glassy metallic tape. (Multimedia view).

Instant temperature distributions along the direction of wave propagation reveal that, with the existence of the wave front, the temperature drops 100-150 K (fig. 2(a)). The front propagates from left to right (the arrows depict the propagation direction). It should be noted that the temperature remains approximately constant behind the front. Therefore, heat flux from the initiating wire does not affect the front propagation. While the samples do not glow, the front propagating was also observed in the visible light spectrum due to variations of the surface reflectivity in amorphous and crystalline phases.

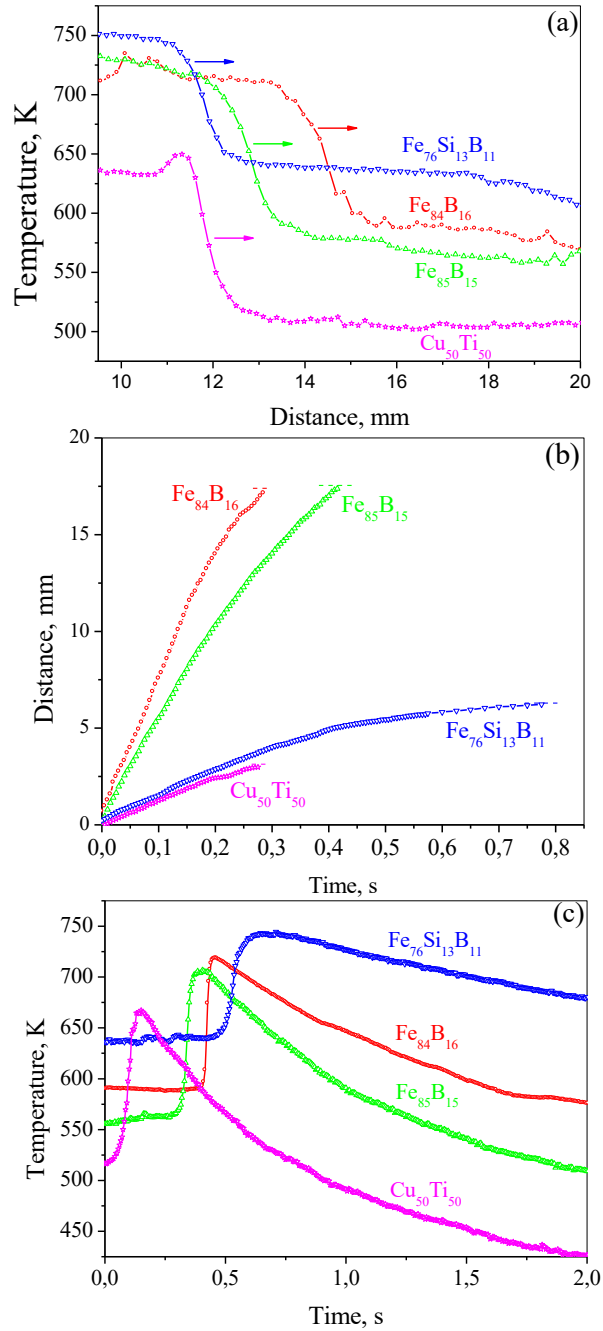


Figure 2. Characteristics of the crystallization wave: (a) temperature distribution along the tape (arrows depict propagation direction); (b) position of the wave front versus time; (c) temperature-time history measured in one spot ($\sim 0.16 - 0.18$ mm) in the middle of each sample.

The front position time dependence was well fit by a linear function in the middle of the tapes (fig. 2(b)), where the influence of the initiating hot wire from one end and cold steel clip (electrode) from another end are negligible. The macroscopic propagating velocity at this region

is essentially constant and in the range 1-7 cm s⁻¹ (Table 1). When the crystallization front approaches the cold electrode, it typically slows down and propagation becomes auto-oscillating. It is worth noting that the oscillations are more developed in the Cu₅₀Ti₅₀ and Fe₇₆Si₁₃B₁₁ tapes, while the Fe₈₄B₁₆ and Fe₈₅B₁₅ tapes exhibit a steady state propagation mode, almost until the end of the tape. Analysis of the temperature-time profiles, measured within the small surface areas in the middle of the samples (fig. 2(c)), shows that maximum self-heating rate in the amorphous-crystalline transition front is ~4300 K s⁻¹ and was observed for Fe₈₄B₁₆, while Fe₇₆Si₁₃B₁₁ has the minimum self-heating rate of ~1400 K s⁻¹. A comparison of the self-heating rates and propagating velocity with those measured for classical EC in electrolytically deposited Sb (~ 28000 K s⁻¹ heating rate, > 30 cm s⁻¹ propagating velocity)²⁷ shows that EC in metallic glasses has much “softer” characteristics. In addition, the MGs have much less sensitivity to external stimuli. In our experiments, mechanical impact, scratching, or local heating at room temperature did not initiate the self-propagating crystallization wave. Only the local thermal impulse in the pre-heated sample or uniform heating above 620K was found as proper methods for EC initiating²⁶. These differences allow us to distinguish between classical EC in deposited films and the soft EC in metallic glasses produced by fast melt quenching. Based on the comparison of the propagating parameters in different systems, we may suggest that EC-waves become weaker and slower with increasing number of elements in the system. Pure elements, such as explosive Sb, exhibit violent EC; binary systems, Fe-B or Cu-Ti, demonstrate ‘softer’ self-propagating transformation; ternary system Fe-Si-B possesses the weakest EC wave. This trend relates, probably, to increasing the glass forming ability and stability of metallic glasses with increasing complexity of their composition. For example, it was shown that adding of Si into the binary Fe-B alloy expands concentration region of amorphous phase²⁸. In addition, increasing number of elements may retard crystallization, since long-range diffusion becomes critical for formation of primary ordered clusters²⁹. Further research is required to determine the mechanism of EC in the multicomponent bulk metallic glasses.

According to X-ray diffraction of the tapes before and after passing the EC wave, the Fe-B amorphous alloys crystallized into α -Fe and tetragonal γ -Fe₃B phases and the Cu₅₀Ti₅₀ metallic glass transformed into a single intermetallic phase of polycrystalline tetragonal CuTi. The differential scanning calorimetry (DSC) results of the thermal effects and activation energies (Table 1) are consistent with those published earlier, e.g., 164-197 kJ mol⁻¹ for Fe-B³⁰ and 207±33 kJ mol⁻¹ for Cu₅₀Ti₅₀³¹. The composition of Fe₈₅B₁₅ has two exothermal peaks in the DSC curve that probably relate to the primary crystallization of α -Fe, followed by the crystallization of eutectic α -Fe + Fe₃B. Fe₈₄B₁₆, which is closer to eutectic point (83 at.% Fe, 17 at.% B), has one exothermal peak. Transmission electron microscopy (TEM) shows that the initial foils have a homogeneous microstructure (fig. 3, TEM BF, left column) that transforms into sub-micrometer polycrystalline material after passing the EC wave (fig. 3, TEM BF, right column). The atomic structure evolution, from practically stochastic dispositions of atoms to a long-distance ordered crystal lattice, have been observed by means of high-resolution TEM (fig. 3, HRTEM). These results were confirmed with selected area microdiffraction of electrons (fig.3, SAD).

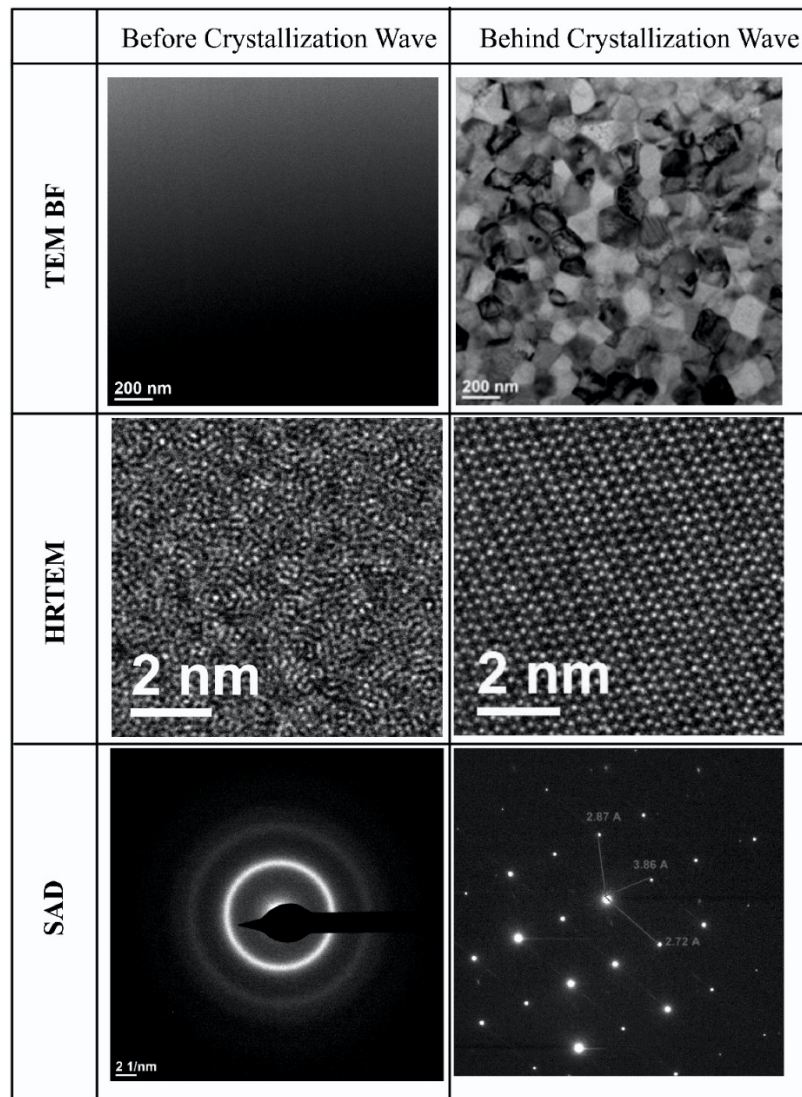


Figure 3. Structure transformations caused by the crystallization wave in the $\text{Fe}_{84}\text{B}_{16}$ metal glass.

Thus, the structural transformations induced by the self-propagating waves are similar to those that occur in DSC or annealing experiments³⁰⁻³².

In order to validate the possibility of EC in metallic glasses at the atomic scale, a molecular dynamics simulation (MDS) of this process in initially amorphous $\text{Cu}_{50}\text{Ti}_{50}$ alloy has been accomplished (see Methods). The $\text{Cu}_{50}\text{Ti}_{50}$ composition was selected for MDS because it undergoes polymorphous transformation, i.e., the amorphous alloy transforms into one crystalline phase of the same composition. This makes the MDS procedure easier, and the patterns obtained more obvious. The Fe-based compositions undergo eutectic or two-stage (primary & secondary) crystallization, when two or more crystalline phases form, including metastable intermediate

phases. For this reason, MDS of self-propagating waves is a matter of nearest future work. A rectangular, amorphous block ($7760 \times 123 \times 31$ angstroms) composed of 2000400 Cu and Ti atoms in 50:50 ratio was prepared by virtual rapid melt quenching (see supplementary files with Methods). After quenching, the temperature was 500 K and the amorphous structure remained steady at these conditions. When a local heat pulse (650 K) was applied to one end of the block, a self-propagating amorphous-to-crystalline transition wave appeared in the model sample (fig. 4, supplementary video file 2). The temperature rose up to ~ 630 K in that wave (fig. 4(a)), which correlates excellently with experimental results (fig. 2). The atomic ordering shows a crystal lattice close to the tetragonal CuTi structure, however, Cu and Ti atoms were randomly distributed in the lattice sites (fig. 4(b)). Thus, according to the molecular dynamic simulation, a crystalline intermetallic phase with disordered atom positions appears behind the EC wave front. The propagation velocity of the simulated virtual waves is about $20\text{-}40 \text{ m s}^{-1}$, which is considerably higher than experimentally observed. This discrepancy probably relates to very different spatial and temporal scales in the experimental and modelling research of EC. Experimentally, we observe propagation of the EC wave at distances of up to several cm, while the model considers the propagation at distance of $\sim 0.7 \text{ }\mu\text{m}$. Since the crystalline product consists of polycrystalline grains with size $\sim 1 \text{ }\mu\text{m}$ (e.g., see fig.3 TEM BF; polycrystalline structure was also observed for explosively crystallized CuTi), the MDS evidently shows crystallization inside one grain. We can assume here that the real EC wave propagates faster inside the crystalline grain and stays at the grain boundaries, which decreases the average macroscopic propagation velocity. A similar effect of rapid propagation of short distances followed by delays at the boundaries between structural elements was also observed for propagation of the crystallization wave along a chain of amorphous Ni particles³³, and periodic “leading wave crystallization” in amorphous Si films²⁴. Good correlations between MDS and experimental results not only give more evidence for the existence of EC waves in the metallic glasses, but also provide a tool for understanding the intimate

features of the process that cannot be currently investigated by experimental methods, for example movements of individual atoms in the crystallization wave (fig. 4(b) and multimedia view).

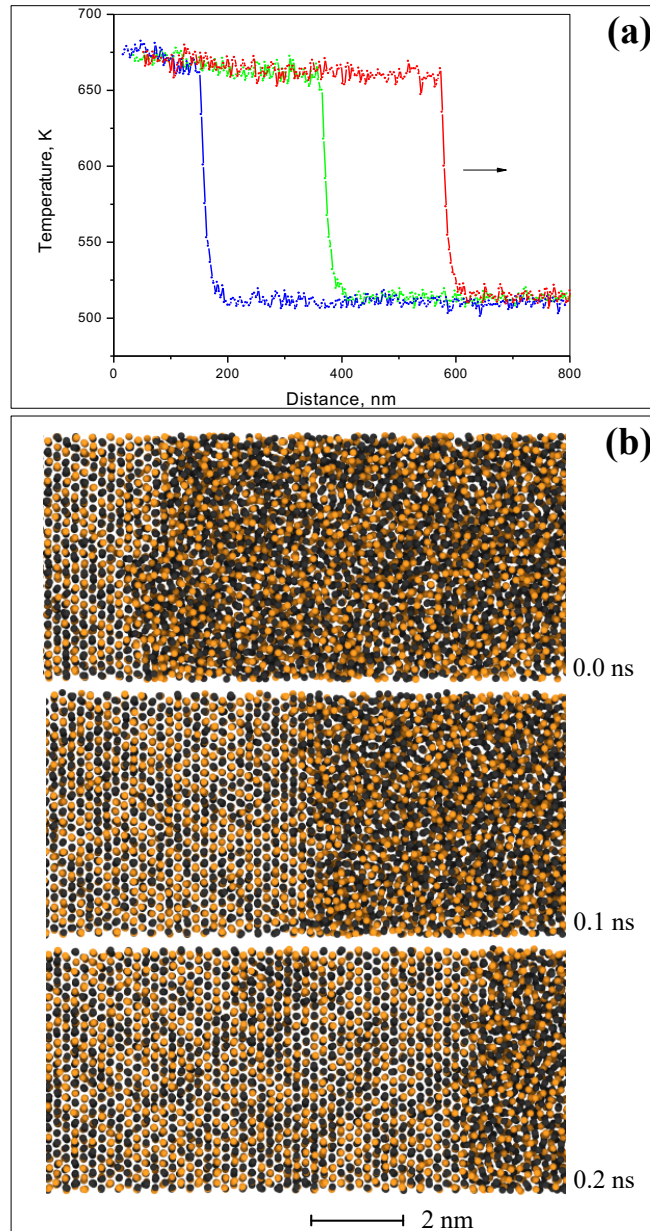


Figure 4. Molecular dynamic simulation of the self-propagating crystallization wave in the $\text{Cu}_{50}\text{Ti}_{50}$ metallic glass. Temperature profiles of the wave (a) and selected cross-sections of the virtual sample (b). The golden balls are Cu atoms and the black balls are Ti. (Multimedia view).

The practical meaning of the soft explosive crystallization phenomena is to be elucidated in the future. We should note that EC in the deposited semiconductor films is used for converting amorphous Si, Ge, etc. into polycrystalline materials required for solar cells²¹⁻²⁴. The SEC can play a similar role for metallic glasses. This conversion may be the most significant for bulk metal glasses, because self-propagating EC waves induce a soft, pre-determined self-heating in all points of the sample independently of the spatial location or depth within the sample. In such a way, SEC can create uniform (or lamellar in case of oscillating propagation²⁶) polycrystalline microstructures favorable for mechanical, magnetic, and other properties of materials.

SUPPLEMENTARY MATERIALS

See supplementary materials for detailed methods description, IR video of the waves (multimedia files “Figure_1.mov” and “supplemental_video_CuTi.wmv”) and Molecular Dynamic Simulation of the wave (multimedia file “Figure_4.mov”).

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References

1. Gore, G., On a peculiar phenomenon in the electro-deposition of antimony. *The London, Edinburg, and Dublin philosophical magazine and journal of science* **9**, 73-74 (1855).
2. Prins, J.A. Diffraction of electrons in amorphous and in crystalline antimony. *Nature* **131**, 760-761 (1933).

3. Duwez, P., Willens, R.H. & Klement Jr., W. Continuous series of metastable solid solutions in silver-copper alloys. *J. Appl. Phys.* **31**, 1136-1137 (1960).
4. Klement Jr., W., Willens, R.H. & Duwez, P. Non-crystalline structure in solidified gold-silicon alloys. *Nature* **187**, 869-870 (1960).
5. Weber, A.W. & Bakker, H. Amorphization by ball milling. A review. *Physica B* **153**, 93-135 (1988).
6. Politis, C. Bulk Amorphous and nanostructured materials by high energy ball milling. *Int. J. of Modern Physics B* **22**, 2905-2913 (2008).
7. Reer, A.L. Metallic Glasses. *Science* **267**, 1947-1953 (1995).
8. Wang, W.H. Metallic Glasses. Family traits. *Nature Materials* **11**, 275-276 (2012).
9. Johnson, W.L. & Plummer, J. Is metallic glass poised to come of age? *Nature Materials* **14**, 553-555 (2015).
10. Hasegawa, R. Applications of amorphous magnetic alloys. *Materials Science and Engineering: A* **375-377**, 90-97 (2004).
11. Si, J., Mei, J., Wang, R., Chen, X. & Hui, X. Fe-B-Si-Zr bulk metallic glasses with ultrahigh compressive strength and excellent soft magnetic properties. *Materials Letters* **181**, 282-284 (2016).
12. Ashby, M.F. & Greer, A.L. Metallic glasses as structural materials. *Scripta Materialia* **54**, 321-325 (2006).
13. Schuh, C.A., Hufnagel, T.C., & Ramamurty, U. Mechanical behavior of amorphous alloys. *Acta Materialia* **55**, 4067-4109 (2007).
14. Ma, E. Tuning order in disorder. *Nature Materials* **14**, 547-552 (2015).
15. Vauth, S., & Mayr, S.G. Atomic dynamics in molecular dynamics simulation of glassy CuTi thin films. *Appl. Phys. Lett.* **86**, 061913(1-3) (2005).
16. Yavari, A.R. A new order for metallic glasses. *Nature* **439**, 405-406 (2006).

17. Johnson, W.L. Bulk metallic glasses - a new engineering material. *Current Opinion in Solid State & Materials Science* **1**, 383-386 (1996).
18. Ding, S., Liu, Y., Li, Y., Liu, Z., Sohn, S., Walker, F.J. & Schroers J. Combinatorial development of bulk metallic glasses. *Nature materials* **13**, 494-500 (2014).
19. Zhong, L., Wang, J., Sheng, H., Zhang, Z. & Mao, S.X. Formation of monoatomic metallic glasses through ultrafast liquid quenching. *Nature* **512**, 177-180 (2014).
20. Götzberger, A., Über die kristallisation aufgedampfter antimonschichten. *Zeitschrift für Physik* **142**, 182-200 (1955) (in German).
21. Leamy, H.J., Brown, W.L., Celler, G.K., Foti, G. & Gilmer, G.H. Explosive crystallization of amorphous germanium. *Appl. Phys. Lett.* **38(3)**, 137-139 (1981).
22. Aleksandrov, L.N. & Edelman, F.L. Shock crystallization of films. *Phys. Stat. Sol. (a)* **76**, 409-427 (1983).
23. Ohdaira, K., Fujiwara, T., Endo, Y., Nishizaki, S. & Matsumura, H. Explosive crystallization of amorphous silicon films by flash lamp annealing. *J. Appl. Phys.* **106(4)**, 044907(1-8) (2009).
24. Hayashi, S., Fujita, Y., Kamikura, T., Sakaike, K., Akazawa, M., Ikeda, M. & Higashi, S. Leading Wave Crystallization Induced by Micro-Thermal-Plasma-Jet Irradiation of Amorphous Silicon Films. *Japanese Journal of Applied Physics* **52**, 05EE02(1-6) (2013).
25. Z. Altounian, C. A. Volkert, and J. O. Strom-Olsen, Crystallization characteristics of Fe-Zr metallic glasses from Fe₄₃Zr₅₇ to Fe₂₀Zr₈₀. *Journal of Applied Physics* **57**, 1777-1782 (1985); doi: 10.1063/1.334455.
26. Rogachev, A.S., Vadchenko, S.G., Shchukin, A.S., Kovaleva, I.D. & Aronin, A.S. Self-propagating crystallization waves in the TiCu amorphous alloy. *JETP Letters* **104**, 726-729 (2016).

27. Rogachev, A.S., Vadchenko, S.G. & Shchukin, A.S. SHS reaction and explosive crystallization in thin films: resemblance and distinction. *Int. J. Self-Propag. High-Temp. Synth.* **26**, 44-48 (2017).
28. Shao, G., Lu, B., Liu, Y.Q. & Tsakirooulos P. Glass forming ability of multi-component metallic systems. *Intermetallics* **13**, 409-414 (2005).
29. Chen, N., Martin, L., Luzguine-Luzgin, D.V. & Inoue A. Role of alloying additions in glass formation and properties of bulk metallic glasses. *Materials* **3**, 5320-5339 (2010).
30. Abrosimova, G.E. & Aronin, A.S. Reversible structural changes on heat treatment of amorphous Fe-B alloys. *Int. J. of Rapid Solidification* **6**, 29-40 (1991).
31. Shanker Rao, T.L., Lad, K.N. & Pratap, A. Study of non-isothermal crystallization of amorphous Cu₅₀Ti₅₀ alloy. *Journal of Thermal Analysis and Calorimetry* **78**, 769-774 (2004).
32. Kovalev, D.Yu., Vadchenko, S.G., Rogachev, A.S., Aronin, A.S. & Alymov, M.I. Time-resolved X-ray diffraction study of the transition of an amorphous TiCu alloy to the crystalline state. *Doklady Physics* **62**, 111-114 (2017).
33. Manukyan, K.V., Shuck, C.E., Cherukara, M.J., Rouvimov, S., Kovalev, D.Y., Strachan, A. & Mukasyan, A.S. Exothermic Self-Sustained Waves with Amorphous Nickel. *J. Phys. Chem. C* **120**, 5827–5838 (2016).