



HAL
open science

Observation of the decay $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$

R. Aaij, A.S.W. Abdelmotteleb, C. Abellán Beteta, F. Abudinén, T. Ackernley, B. Adeva, M. Adinolfi, H. Afsharnia, C. Agapopoulou, C.A. Aidala, et al.

► To cite this version:

R. Aaij, A.S.W. Abdelmotteleb, C. Abellán Beteta, F. Abudinén, T. Ackernley, et al.. Observation of the decay $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$. Physical Review Letters, 2022, 128 (19), pp.191803. 10.1103/PhysRevLett.128.191803 . hal-03537001

HAL Id: hal-03537001

<https://hal.science/hal-03537001v1>

Submitted on 7 Sep 2023

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution 4.0 International License



Observation of the decay

$$\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$$

LHCb Collaboration[†]

Abstract

The first observation of the semileptonic b -baryon decay $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$, with a significance of 6.1σ , is reported using a data sample corresponding to 3 fb^{-1} of integrated luminosity, collected by the LHCb experiment at centre-of-mass energies of 7 and 8 TeV at the LHC. The τ^- lepton is reconstructed in the hadronic decay to three charged pions. The ratio $\mathcal{K} = \mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau) / \mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^- \pi^+ \pi^-)$ is measured to be $2.46 \pm 0.27 \pm 0.40$, where the first uncertainty is statistical and the second systematic. The branching fraction $\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau) = (1.50 \pm 0.16 \pm 0.25 \pm 0.23)\%$ is obtained, where the third uncertainty is from the external branching fraction of the normalization channel $\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^- \pi^+ \pi^-$. The ratio of semileptonic branching fractions $\mathcal{R}(\Lambda_c^+) \equiv \mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau) / \mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu)$ is derived to be $0.242 \pm 0.026 \pm 0.040 \pm 0.059$, where the external branching fraction uncertainty from the channel $\Lambda_b^0 \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu$ contributes to the last term. This result is in agreement with the standard model prediction.

Published in Phys. Rev. Lett. 128, 191803.

© 2023 CERN for the benefit of the LHCb Collaboration. CC BY 4.0 licence.

[†]Authors are listed at the end of this Letter.

In the standard model of particle physics (SM) flavor-changing processes, such as semileptonic decays of b hadrons, are mediated by W^\pm bosons with universal coupling to leptons. Differences in the rates of decays involving the three lepton families are expected to arise only from the different masses of the charged leptons. Lepton flavor universality can be violated in many extensions of the SM with non-standard flavor structure. Since the uncertainty due to hadronic effects cancels to a large extent, the SM predictions for the ratios between branching fractions of semileptonic decays of b hadrons, such as $\mathcal{R}(D^{(*)}) \equiv \mathcal{B}(\bar{B} \rightarrow D^{(*)}\tau^-\bar{\nu}_\tau)/\mathcal{B}(\bar{B} \rightarrow D^{(*)}\mu^-\bar{\nu}_\mu)$ [1–3], where $D^{(*)}$ and B mesons can be either charged or neutral, and $\mathcal{R}(\Lambda_c^+) \equiv \mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+\tau^-\bar{\nu}_\tau)/\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+\mu^-\bar{\nu}_\mu)$, are known with uncertainties at the per cent level [4–6]. These ratios therefore provide a sensitive probe of SM extensions [6, 7].

Measurements of $\mathcal{R}(D^{0,+})$ and $\mathcal{R}(D^{*+,0})$ with τ^- decay final states involving electrons or muons have been reported by the BaBar [8, 9] and Belle [10–12] Collaborations. The LHCb Collaboration published a determination of $\mathcal{R}(D^{*+})$ [13], where the τ lepton is reconstructed using leptonic decays to a muon. The LHCb experiment has also reported a measurement of $\mathcal{R}(D^{*+})$ using the three-prong decay $\tau^- \rightarrow \pi^-\pi^+\pi^-(\pi^0)\nu_\tau$ [14]. These $\mathcal{R}(D^{(*)+,0})$ measurements yield values that are larger than the SM predictions with a combined significance of 3.4 standard deviations (σ) to date [15].

This Letter reports the observation of the decay $\Lambda_b^0 \rightarrow \Lambda_c^+\tau^-\bar{\nu}_\tau$ and the first determination of $\mathcal{R}(\Lambda_c^+)$ using $\tau^- \rightarrow \pi^-\pi^+\pi^-(\pi^0)\nu_\tau$ decays. The inclusion of charge-conjugate modes is implied throughout. The present work closely follows the strategy of Ref. [14]. Measurements in the baryonic sector provide complementary constraints on a potential lepton flavor universality violation because of the half-integer spin of the initial state [5, 6]. The $\Lambda_b^0 \rightarrow \Lambda_c^+$ transition is determined by a different set of form factors with respect to the mesonic decays probed so far. Likewise, new physics couplings can also be different, resulting in different scenarios regarding deviations from SM expectations of $\mathcal{R}(\Lambda_c^+)$ and $\mathcal{R}(D^{(*)})$ [7].

A data sample of proton-proton (pp) collisions at centre-of-mass energies $\sqrt{s} = 7$ and 8 TeV, corresponding to an integrated luminosity of 3 fb^{-1} , collected with the LHCb detector is used. The LHCb detector is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, described in detail in Refs. [16, 17]. The detector includes a high-precision tracking system consisting of a silicon-strip vertex detector surrounding the pp interaction region [18], and large-area silicon-strip detectors located upstream and downstream of the 4 Tm dipole magnet. The minimum distance of a track to a primary pp collision vertex (PV), the impact parameter (IP), is measured with a resolution of $(15 + 29/p_T)\text{ }\mu\text{m}$, where p_T is the component of the momentum transverse to the beam direction, in GeV/c . The online event selection is performed by a trigger system [19], which consists of a hardware stage based on information from the calorimeter and muon systems, followed by a software stage that performs a full event reconstruction. Events are selected at the hardware stage if the particles forming the signal candidate satisfy a requirement on the energy deposited in the calorimeters or if any other particles pass any trigger algorithm. The software trigger requires a two-, three-, or four-track secondary vertex with significant displacement from any PV and consistent with the decay of a b hadron, or

a three-track vertex with a significant displacement from any PV and consistent with the decay of a Λ_c^+ baryon. A multivariate algorithm is used for the identification of secondary vertices consistent with the decay of a b hadron, while secondary vertices consistent with the decay of a Λ_c^+ baryon are identified using topological criteria. In the simulation, pp collisions are generated using PYTHIA 8 [20] with a specific LHCb configuration [21]. Decays of hadronic particles are described by EVTGEN [22], in which final-state radiation is generated using PHOTOS [23]. The TAUOLA package [24] is used to simulate the decays of the τ^- lepton into $3\pi\nu_\tau$ and $3\pi\pi^0\nu_\tau$ final states, where $3\pi \equiv \pi^-\pi^+\pi^-$, according to the resonance chiral Lagrangian model [25] with a tuning based on the results from the BaBar Collaboration [26]. The interaction of the generated particles with the detector, and its response, are implemented using the GEANT4 toolkit [27] as described in Ref. [28]. The signal decays are simulated using form factors that are derived from heavy-quark effective theory [29].

The Λ_c^+ baryon candidates are reconstructed using the $\Lambda_c^+ \rightarrow pK^-\pi^+$ decay mode, by combining three charged tracks compatible with proton, kaon and pion hypotheses. The τ^- candidates are formed by $\pi^-\pi^+\pi^-$ combinations and include contributions from the $\tau^- \rightarrow 3\pi\nu_\tau$ and $\tau^- \rightarrow 3\pi\pi^0\nu_\tau$ decay modes, as neutral pions are not reconstructed. The Λ_c^+ and τ^- candidates are selected based on kinematic, geometric, and particle-identification criteria. The Λ_b^0 candidate is formed by combining a Λ_c^+ and a τ^- candidate. Background due to misreconstructed b hadrons, where at least one additional particle originates from either the 3π vertex or the b -hadron vertex, is suppressed by requiring a single Λ_b^0 candidate per event. Tracks other than those used for the signal candidate are exploited in a multivariate algorithm to assess the signal isolation, *i.e.* the absence of extra tracks compatible with the 3π vertex [30]. The algorithm is trained on simulated samples of $\Lambda_b^0 \rightarrow \Lambda_c^+\tau^-\bar{\nu}_\tau$ and $\Lambda_b^0 \rightarrow \Lambda_c^+\bar{D}^0K^-$ decays for signal and background, respectively; its efficiency is 20% higher than the cut-based algorithm used in Ref. [14] for the same rejection factor. Likewise, the neutral-particle energy contained in a cone centred around the direction of the τ^- candidates is used to further separate signal and background processes. The τ^- momentum can be determined, up to a two-fold ambiguity, from the momentum vector of the 3π system and the flight direction of the τ^- candidate. The average of the two solutions is used, as discussed in Ref. [14]. The same method is used to compute the Λ_b^0 momentum. This enables the computation of the invariant mass squared of the $\tau^-\bar{\nu}_\tau$ lepton pair (q^2), and the pseudo decay time of the τ^- candidate (t_τ). The variables q^2 and t_τ are reconstructed with a resolution of roughly 15%, providing good discrimination between the signal and background processes.

The finite τ^- lifetime causes the 3π vertex to be detached from the Λ_b^0 vertex. This key feature allows the suppression of the large background, called prompt background hereafter, from b hadrons decaying to a Λ_c^+ baryon accompanied by a 3π system being produced promptly at the b -hadron decay vertex, plus any other unreconstructed particles (X). The difference of the positions of the 3π and the Λ_c^+ vertices along the beam direction is required to be at least 5 times larger than its uncertainty. This requirement suppresses the prompt background 10 times more than the selection used in Ref. [14], reducing this initially dominant background to a negligible level, at a price of 50% reduction of the

signal efficiency.

Double-charm background processes due to Λ_b^0 baryon decays into a Λ_c^+ baryon plus another charmed hadron that subsequently decays into a final state containing three charged pions, are topologically similar to the signal and constitute the largest background source. The main contribution originates from $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-(X)$ decays, with D_s^- decays to $3\pi Y$ final states, where Y stands for any set of extra particles, such as one or two π^0 mesons. Such D_s^- decays have a large branching fraction ($\sim 30\%$) [31]. This background is reduced principally by taking into account the resonant structure of the 3π system. The $\tau^- \rightarrow 3\pi\nu_\tau$ decays proceed predominantly through the $a_1(1260)^- \rightarrow \rho^0\pi^-$ decay. By contrast, the $D_s^- \rightarrow 3\pi Y$ decays occur mainly through the η and η' resonances. This feature, captured by the shapes of the distributions of the smaller and larger mass of the two $\pi^+\pi^-$ combinations extracted from each 3π candidate, the energy carried by neutral particles within the cone around the 3π direction, and kinematic variables from partial reconstruction are exploited by means of a boosted decision tree (BDT) classifier [32, 33], as described in Ref. [14]. Fig. 4 of Supplemental Material [34] displays the markedly different distributions of the three main input variables to the BDT classifier obtained for signal and D_s^- background, respectively. The partial reconstruction of the Λ_b^0 decay kinematics is performed under the background hypothesis where the Λ_b^0 particle decays to $\Lambda_c^+ D_s^-(\rightarrow 3\pi Y)$. The BDT response in simulation is validated using three control samples: the $\Lambda_b^0 \rightarrow \Lambda_c^+ 3\pi$ normalization sample; a $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0(X)$ data sample with the subsequent $\bar{D}^0 \rightarrow K^+ 3\pi$ decay, which is obtained by removing the charged-particle isolation criterion and requiring an additional charged kaon originating from the 3π vertex; and a $\Lambda_b^0 \rightarrow \Lambda_c^+ D^-(X)$ data sample, using $D^- \rightarrow K^+ \pi^- \pi^-$ decays, which is obtained by assigning a kaon mass to the positively-charged pion of the τ^- candidate. For all these samples, good agreement between data and simulation is observed in the distributions of the variables used in the BDT classifier.

The signal yield is measured using a three-dimensional binned maximum-likelihood fit to t_τ , the BDT output, and q^2 , which are shown in Fig. 1 and 2. The fit model includes a signal component; background components due to $B \rightarrow \Lambda_c^+ D_s^-(X)$, $\Lambda_b^0 \rightarrow \Lambda_c^+ D^-(X)$ and $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0(X)$ decays; background due to misreconstructed Λ_c^+ candidates; and combinatorial background. Template distributions for signal and background are obtained from simulation, with the exceptions of random $pK^-\pi^+$ combinations and the combinatorial background, which are constructed from data-based control samples. The signal template accounts for both $\tau^- \rightarrow 3\pi\nu_\tau$ and $\tau^- \rightarrow 3\pi\pi^0\nu_\tau$ decays, where the fraction of the former is fixed to 78% according to the branching fractions and selection efficiencies. A contribution from $\Lambda_b^0 \rightarrow \Lambda_c^{*+} \tau^- \bar{\nu}_\tau$ decays, where Λ_c^{*+} denotes any excited charmed baryon state decaying into final states involving Λ_c^+ baryon, constitutes a feeddown to the signal. Its yield fraction is constrained to be $(10 \pm 5)\%$ of the signal yield, derived from the Λ_c^{*+} relative abundance as measured in the $\Lambda_b^0 \rightarrow \Lambda_c^{*+} \pi^- \pi^+ \pi^-$ decays, their respective branching fractions in the $\Lambda_c^+ \pi^0 \pi^0$ and $\Lambda_c^+ \pi^+ \pi^-$ modes [31], and the corresponding selection efficiency obtained from simulation. The background originating from decays of $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-(X)$ is divided into contributions from $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-$, $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^{*-}$, $\Lambda_b^0 \rightarrow \Lambda_c^+ D_{s0}^*(2317)^-$, $\Lambda_b^0 \rightarrow \Lambda_c^+ D_{s1}(2460)^-$, and $\Lambda_b^0 \rightarrow \Lambda_c^{*+} D_s^-(X)$ decays. A control sample of $\Lambda_b^0 \rightarrow \Lambda_c^+ 3\pi$

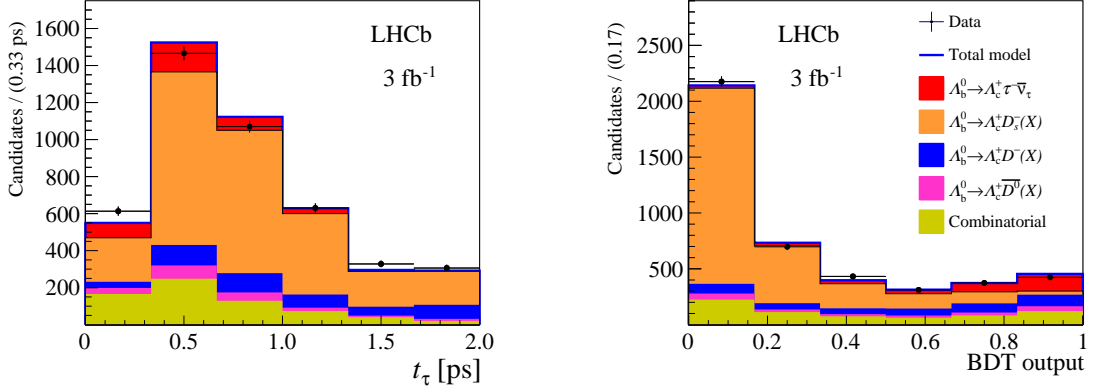


Figure 1: Distributions of (left) τ^- decay time and (right) BDT output for $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$ candidates. Projections of the three-dimensional fit results are overlaid. The various fit components are described in the legend.

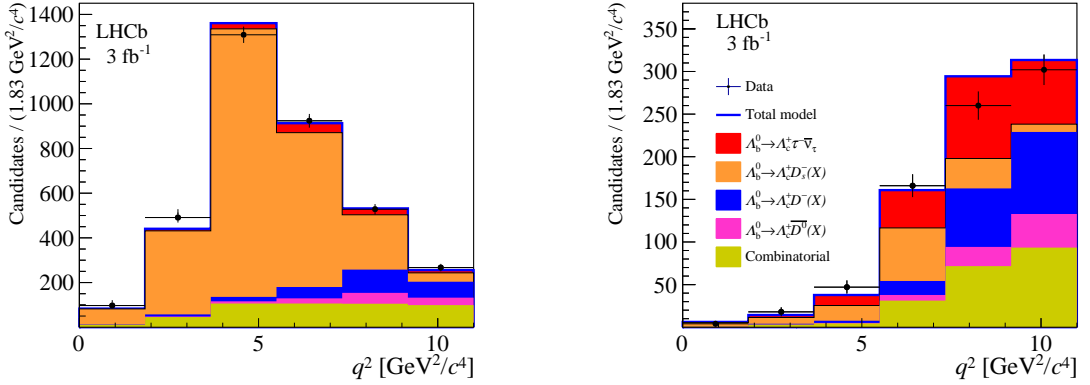


Figure 2: Distributions of q^2 for $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$ candidates having a BDT output value (left) below and (right) above 0.66. Projections of the three-dimensional fit are overlaid. The various fit components are described in the legend.

candidates, where the 3π invariant mass is selected within $45 \text{ MeV}/c^2$ of the known D_s^- mass [31] is shown in Fig. 3. The relative yield of each of above mentioned background processes is constrained using the results of a fit to the $\Lambda_c^+ \pi^- \pi^+ \pi^-$ mass distribution.

The D_s^- decay model used in the simulation does not accurately describe the data because of the limited knowledge of the D_s^- decay amplitudes to $3\pi Y$ final states. A correcting factor, taken from high precision $\bar{B}^0 \rightarrow D^{*+} D_s^-$ sample [14], is applied to each D_s^- branching fraction to match the $\pi^- \pi^+ \pi^-$ Dalitz distributions from simulation to those observed in data.

The background originating from $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0(X)$ decays is subdivided into two contributions, depending on whether the 3π system originates from the \bar{D}^0 vertex, or whether one pion originates from the \bar{D}^0 vertex and the other two from elsewhere. The

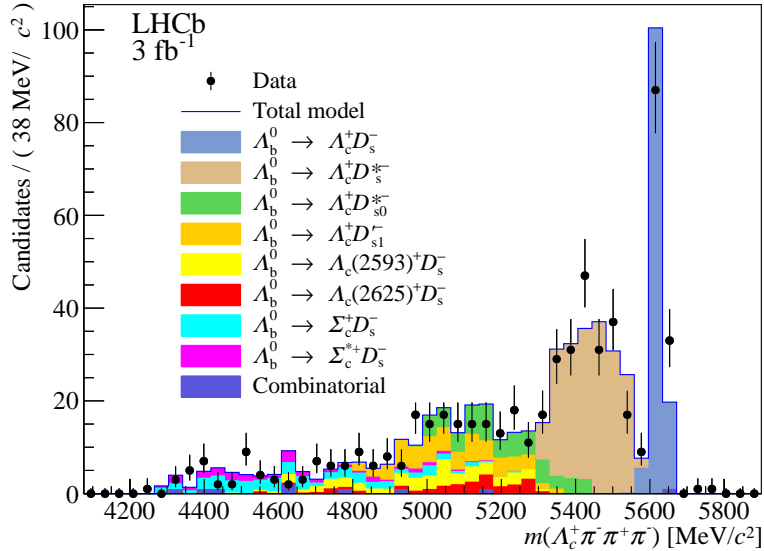


Figure 3: Distribution of the $\Lambda_c^+ \pi^- \pi^+ \pi^-$ invariant mass for the $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-(X)$ control sample, with $D_s^- \rightarrow \pi^- \pi^+ \pi^-$. The components contributing to the fit model are indicated in the legend.

former contribution is constrained by the yield obtained from the $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0(X)$ control sample. The template associated to $\Lambda_c^+ \bar{D}^0(X)$ background is also validated using the data-driven sample where the $\bar{D}^0 \rightarrow K^+ 3\pi$ decay is fully reconstructed. The yield of the other $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0(X)$ background component is a free parameter in the fit. The yield of the $\Lambda_b^0 \rightarrow \Lambda_c^+ D^-(X)$ background is also a free parameter and its template is validated using the data-driven sample with the D^- meson fully reconstructed in the $K^+ \pi^- \pi^-$ mode.

The combinatorial background is divided into two contributions, depending on whether the Λ_b^0 candidate contains a true Λ_c^+ baryon or a random $pK^- \pi^+$ combination. In the first case, the Λ_c^+ and the 3π system originate from different b -hadron decays. The data sample of wrong-sign Λ_b^0 candidates where the Λ_c^+ and the 3π system have the same electric charge is used to obtain a background template. Its yield is obtained by normalising to the right-sign data in the region where the reconstructed $\Lambda_c^+ 3\pi$ mass is significantly larger than the known Λ_b^0 mass [14]. The background not including a true Λ_c^+ baryon is parameterised using a specific data sample originating from Λ_b^0 candidates where the Λ_c^+ candidate has a mass outside a window of $15 \text{ MeV}/c^2$ around the known Λ_c^+ mass [31].

The projections of the fit on t_τ and the BDT output are shown in Fig. 1. The projections on q^2 in two different BDT output ranges are shown in Fig. 2. The signal yield is $N_{\text{sig}} = 349 \pm 40$. The fit is repeated with all nuisance parameters related to the template shapes varying freely, while the signal yield is fixed at zero. The χ^2 variation derived from the change of the fit maximum likelihood corresponds to an increase of 6.1σ with respect to the default fit with freely varying signal yield. This measurement signifies the

first observation of the decay $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$. A clear separation between signal and the main background originating from $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^+(X)$ decays is obtained, as demonstrated in the BDT distribution of Fig. 1. Figure 2 shows that the $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^+(X)$ background is dominant at low BDT values, while a good signal-to-background ratio is observed at high BDT output. Fig. 5 of Supplemental Material [35] shows similarly the τ decay time distribution for the same BDT intervals.

In order to reduce experimental systematic uncertainties, the $\Lambda_b^0 \rightarrow \Lambda_c^+ 3\pi$ decay is chosen as a normalization channel. This leads to a measurement of the ratio

$$\mathcal{K}(\Lambda_c^+) \equiv \frac{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ 3\pi)} = \frac{N_{\text{sig}} \varepsilon_{\text{norm}}}{N_{\text{norm}} \varepsilon_{\text{sig}}} \frac{1}{\mathcal{B}(\tau^- \rightarrow 3\pi(\pi^0)\nu_\tau)}, \quad (1)$$

where N_{sig} (N_{norm}) and ε_{sig} ($\varepsilon_{\text{norm}}$) are the yield and selection efficiency for the signal (normalization) channel, respectively. The normalization channel selection is identical to that of the signal channel, except the requirement that the 3π system has a larger flight distance than that of the Λ_c^+ candidate, which is not imposed. The yield of the normalization mode is determined by fitting the invariant-mass distribution of the $\Lambda_c^+ 3\pi$ candidates around the known Λ_b^0 mass [31], as shown in Fig. 6 of Supplemental Material [36]. A significant contribution from excited baryons which decay to $\Lambda_c^+ \pi^+ \pi^-$, $\Lambda_c^+ \pi^+$, or $\Lambda_c^+ \pi^-$ is explicitly vetoed from the normalization channel. As a result, the 3π dynamics resembles that of the signal, leading to a reduced systematic uncertainty.

A normalization yield of $N_{\text{norm}} = 8584 \pm 102$ is found, after subtraction of a small contribution of 168 ± 20 $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-(\rightarrow 3\pi)$ decays. This component is estimated by fitting the 3π mass distribution in the D_s^- mass region for candidates with a reconstructed $\Lambda_c^+ 3\pi$ mass in a window around the known Λ_b^0 mass [31]. The normalization sample is also used to correct for differences in the Λ_b^0 production kinematics between data and simulation. The reconstruction efficiencies for the $\tau^- \rightarrow 3\pi\nu_\tau$, $\tau^- \rightarrow 3\pi\pi^0\nu_\tau$ signal modes and normalization channel are determined using the simulation and found to be $(1.37 \pm 0.04) \times 10^{-5}$, $(0.82 \pm 0.05) \times 10^{-5}$, and $(11.21 \pm 0.11) \times 10^{-5}$, respectively. The ratio of branching fractions is derived from Eq. 1 as

$$\mathcal{K}(\Lambda_c^+) = 2.46 \pm 0.27 \pm 0.40,$$

where the first uncertainty is statistical and the second systematic.

Using $\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ 3\pi) = (6.14 \pm 0.94) \times 10^{-3}$ [31] corresponding to an average of measurements by the CDF [37], and LHCb [38] experiments, the signal branching fraction is determined as

$$\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau) = (1.50 \pm 0.16 \pm 0.25 \pm 0.23)\%,$$

where the first uncertainty is statistical, the second systematic and the third is due to the external branching fraction measurement. The branching fraction $\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu) = (6.2 \pm 1.4)\%$ from the DELPHI experiment [39] updated in Ref. [31] is used to obtain the ratio of semileptonic branching fractions $\mathcal{R}(\Lambda_c^+)$ as

$$\mathcal{R}(\Lambda_c^+) = 0.242 \pm 0.026 \pm 0.040 \pm 0.059,$$

Table 1: Relative systematic uncertainties in $\mathcal{K}(\Lambda_c^+)$.

Source	$\delta\mathcal{K}(\Lambda_c^+)/\mathcal{K}(\Lambda_c^+)[\%]$
Simulated sample size	3.8
Fit bias	3.9
Signal modelling	2.0
$\Lambda_b^0 \rightarrow \Lambda_c^{*+} \tau^- \bar{\nu}_\tau$ feeddown	2.5
$D_s^- \rightarrow 3\pi Y$ decay model	2.5
$\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^- X$, $\Lambda_b^0 \rightarrow \Lambda_c^+ D^- X$, $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0 X$ background	4.7
Combinatorial background	0.5
Particle identification and trigger corrections	1.5
Isolation BDT classifier and vertex selection requirements	4.5
D_s^- , D^- , \bar{D}^0 template shapes	13.0
Efficiency ratio	2.8
normalization channel efficiency (modelling of $\Lambda_b^0 \rightarrow \Lambda_c^+ 3\pi$)	3.0
Total uncertainty	16.5

where the first uncertainty is statistical, the second systematic and the third is due to the external branching fractions measurements. The measured value of $\mathcal{R}(\Lambda_c^+)$ is lower than but in agreement with the Standard Model prediction of 0.324 ± 0.004 [5].

The sources of systematic uncertainty of $\mathcal{K}(\Lambda_c^+)$ are reported in Table 1. For $\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau)$ and $\mathcal{R}(\Lambda_c^+)$, the systematic uncertainties related to the external branching fractions are added in quadrature. The uncertainty due to the limited size of the simulated samples is computed by repeatedly sampling each template with a bootstrap procedure, performing the fit, and taking the standard deviation of the resulting spread of N_{sig} values. The limited size of the simulated samples also contributes to the systematic uncertainty in the efficiencies for signal and normalization modes. The systematic uncertainty associated with the signal decay model originates from the limited knowledge of the form factors and the τ^- polarization. The form factor distributions are varied in their range allowed by measurements from $\Lambda_b^0 \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu$ decays. The contribution from the relative branching fractions and selection efficiencies of $\tau^- \rightarrow 3\pi\pi^0\nu_\tau$ and $\tau^- \rightarrow 3\pi\nu_\tau$ decays are computed by varying their ratio within their uncertainties. Potential contribution from other τ^- decay modes is investigated through a dedicated simulation including all known τ^- decay modes. Feed-down contribution where the τ^- is produced in association with an excited charmed baryon is computed varying in the fit the relative amount of such decays in their allowed range of $(10 \pm 5)\%$. The uncertainty due to the knowledge of the D_s^- decay model is estimated by repeatedly varying the correction factors of the templates within their uncertainties, as determined from the associated control sample, and performing the fit. The spread of the fit results is assigned as the corresponding systematic uncertainty.

The template shapes of the $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-(X)$, $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0 X$ and $\Lambda_b^0 \rightarrow \Lambda_c^+ D^-(X)$ background modes depend on the dynamics of the corresponding decays. A range of template deformations [14] is performed, and the spread of the fit results is taken as a

systematic uncertainty. The resulting uncertainty of 13% represents the largest single source. A similar procedure is applied to the template for the combinatorial background. The contribution from a potential bias in the fit is explored by fitting pseudoexperiments where the signal strength is varied from its SM value to a negligible amount. Other sources of systematic uncertainty arise from the inaccuracy on the yields of the various background contributions, and from the limited knowledge of the normalization channel modelling. The contribution from the removal of Λ_c^{*+} modes from the normalization channel is taken into account by varying the branching fractions of the various excited baryons decays within their measured range.

Systematic effects in the efficiencies for signal and normalization channels partially cancel in the ratio, with the remaining uncertainty being mostly due to the limited size of simulated sample. The trigger efficiency depends on the distributions of the decay time of the τ^- candidates and the invariant mass of the $\Lambda_c^+ 3\pi$ system. These distributions differ between the signal and normalization modes, and the corresponding difference of the trigger efficiencies is taken into account.

In conclusion, the first observation of the semileptonic decay $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$ is reported with a significance of 6.1σ , using a data sample of pp collisions, corresponding to 3fb^{-1} of integrated luminosity, collected by the LHCb experiment. The measurement exploits the three-prong hadronic τ^- decays with the technique pioneered by the LHCb experiment for the $\mathcal{R}(D^{*+})$ measurement [14]. The ratio $\mathcal{K} = \mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau) / \mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^- \pi^+ \pi^-)$ is measured to be $2.46 \pm 0.27 \pm 0.40$, where the first uncertainty is statistical and the second systematic. The branching fraction $\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau)$ is measured to be $(1.50 \pm 0.16 \pm 0.25 \pm 0.23)\%$, where the third uncertainty is due to external branching fraction measurements. A measurement of $\mathcal{R}(\Lambda_c^+) = 0.242 \pm 0.026 \pm 0.040 \pm 0.059$ is reported. The $\mathcal{R}(\Lambda_c^+)$ ratio is found to be in agreement with the SM prediction. This measurement provides constraints on new physics models, such as some of those described in Ref. [6], for which large values of $\mathcal{R}(\Lambda_c^+)$ are allowed by existing $\mathcal{R}(D^{(*)})$ measurements.

Acknowledgements

We express our gratitude to our colleagues in the CERN accelerator departments for the excellent performance of the LHC. We thank the technical and administrative staff at the LHCb institutes. We acknowledge support from CERN and from the national agencies: CAPES, CNPq, FAPERJ and FINEP (Brazil); MOST and NSFC (China); CNRS/IN2P3 (France); BMBF, DFG and MPG (Germany); INFN (Italy); NWO (Netherlands); MNiSW and NCN (Poland); MEN/IFA (Romania); MSHE (Russia); MICINN (Spain); SNSF and SER (Switzerland); NASU (Ukraine); STFC (United Kingdom); DOE NP and NSF (USA). We acknowledge the computing resources that are provided by CERN, IN2P3 (France), KIT and DESY (Germany), INFN (Italy), SURF (Netherlands), PIC (Spain), GridPP (United Kingdom), RRCKI and Yandex LLC (Russia), CSCS (Switzerland), IFIN-HH (Romania), CBPF (Brazil), PL-GRID (Poland) and NERSC (USA). We are indebted to the communities behind the multiple open-source software packages on which

we depend. Individual groups or members have received support from ARC and ARDC (Australia); AvH Foundation (Germany); EPLANET, Marie Skłodowska-Curie Actions and ERC (European Union); A*MIDEX, ANR, IPhU and Labex P2IO, and Région Auvergne-Rhône-Alpes (France); Key Research Program of Frontier Sciences of CAS, CAS PIFI, CAS CCEPP, Fundamental Research Funds for the Central Universities, and Sci. & Tech. Program of Guangzhou (China); RFBR, RSF and Yandex LLC (Russia); GVA, XuntaGal and GENCAT (Spain); the Leverhulme Trust, the Royal Society and UKRI (United Kingdom).

Supplemental material for the paper ”Observation of the decay $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$ ”

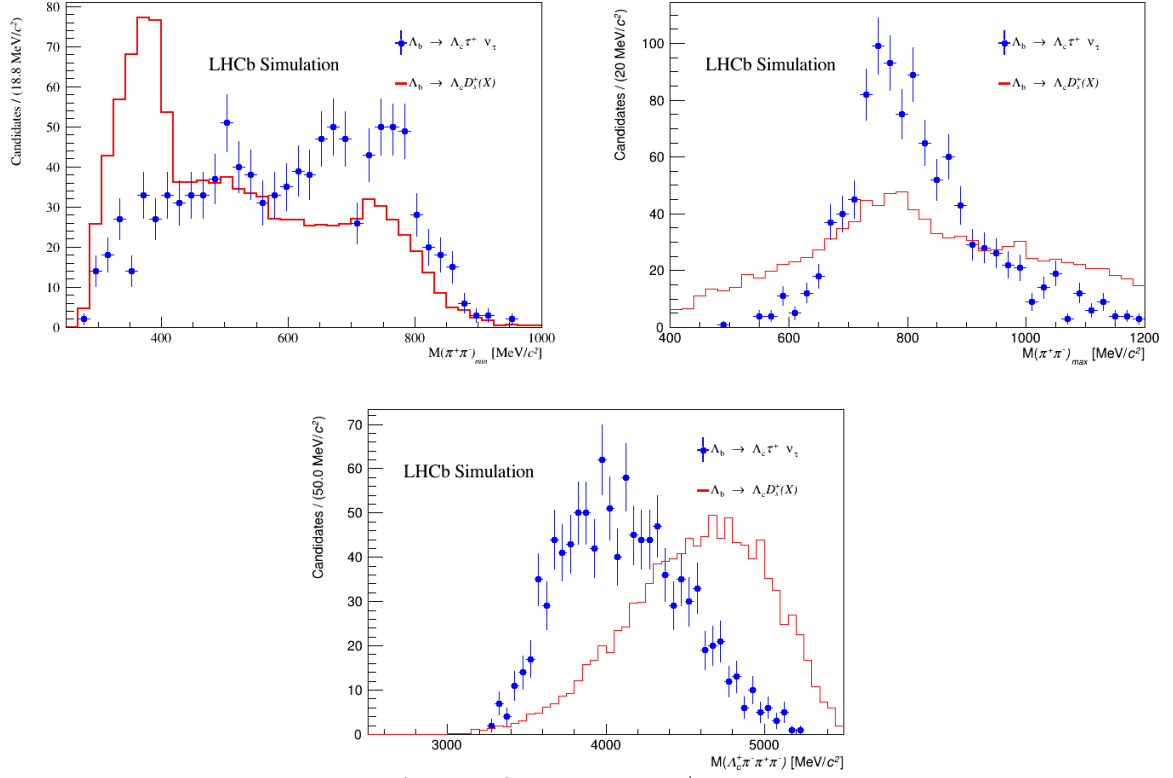


Figure 4: Distribution of the (top left) minimum $\pi^+\pi^-$ mass combination formed from the pion triplet, (top right) maximum $\pi^+\pi^-$ mass combination and (bottom) $\Lambda_c^+\pi^-\pi^+\pi^-$ mass, for simulated samples of $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$ (blue points) and $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^- (X)$ (red line) decays.

Fig. 4 shows the distribution of the three most significant input variables to the BDT classifier : minimum and maximum mass of the two $\pi^+\pi^-$ combinations formed from the pion triplet and $\Lambda_c^+\pi^-\pi^+\pi^-$ mass. Fig. 5 shows the τ decay time distribution for BDT output values below and above 0.66. Fig. 6 shows the invariant-mass distribution of selected $\Lambda_c^+3\pi$ candidates with a fit overlaid to extract the yield of the normalization mode.

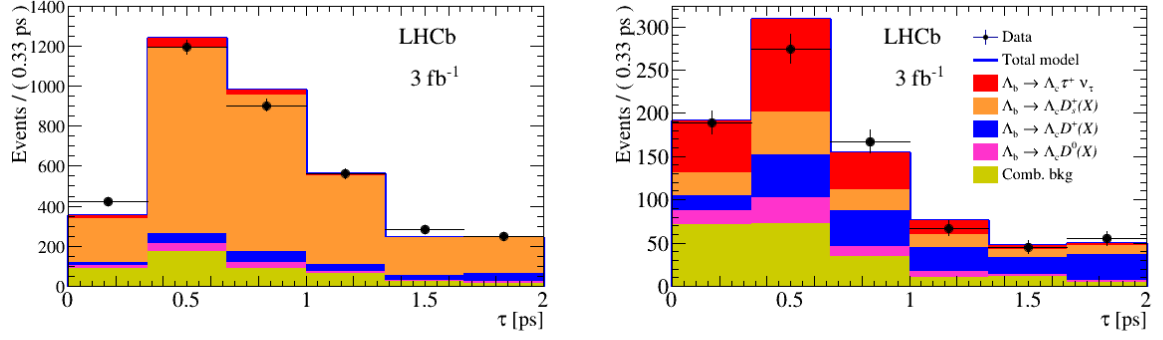


Figure 5: Distribution of the τ decay time for $\Lambda_b^0 \rightarrow \Lambda_c^+ \tau^- \bar{\nu}_\tau$ candidates with (top) BDT output value below 0.66 (bottom) BDT output value above 0.66. The various fit components are described in the legend.

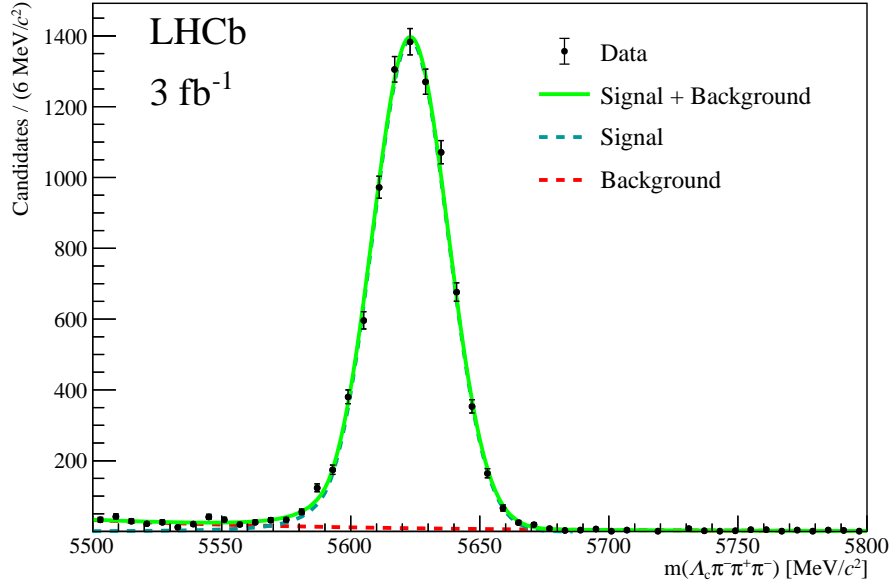


Figure 6: Distribution of the $\Lambda_c^+ \pi^- \pi^+ \pi^-$ invariant mass for all candidates in the normalization channel, after removal of the Λ_c^{*+} contributions. The fit components are indicated in the legend. The signal is described by a Crystal Ball (CB) function, and the background by an exponential term.

References

- [1] D. Bigi and P. Gambino, *Revisiting $B \rightarrow D\ell\nu$* , Phys. Rev. **D94** (2016) 094008, [arXiv:1606.08030](#).
- [2] S. Jaiswal, S. Nandi, and S. K. Patra, *Updates on extraction of $|V_{cb}|$ and SM prediction of $R(D^*)$ in $B \rightarrow D^*\ell\nu_\ell$ decays*, JHEP **06** (2020) 165, [arXiv:2002.05726](#).
- [3] F. U. Bernlochner, Z. Ligeti, M. Papucci, and D. J. Robinson, *Combined analysis of semileptonic B decays to D and D^* : $R(D^{(*)})$, $|V_{cb}|$, and new physics*, Phys. Rev. **D95** (2017) 115008, [arXiv:1703.05330](#), [Erratum: Phys.Rev. **D97**, 059902 (2018)].
- [4] F. U. Bernlochner, Z. Ligeti, D. J. Robinson, and W. L. Sutcliffe, *$\Lambda_b \rightarrow \Lambda_c$ semileptonic decays and tests of heavy quark symmetry*, Physical Review Letters **121** (2018) , [arXiv:1808.09464](#).
- [5] F. U. Bernlochner, Z. Ligeti, D. J. Robinson, and W. L. Sutcliffe, *Precise predictions for $\Lambda_b \rightarrow \Lambda_c$ semileptonic decays*, Phys. Rev. **D99** (2019) 055008, [arXiv:1812.07593](#).
- [6] A. Datta, S. Kamali, S. Meinel, and A. Rashed, *Phenomenology of $\Lambda_b \rightarrow \Lambda_c\tau\bar{\nu}_\tau$ using lattice QCD calculations*, JHEP **08** (2017) 131, [arXiv:1702.02243](#).
- [7] F. U. Bernlochner, M. F. Sevilla, D. J. Robinson, and G. Wormser, *Semitauponic b -hadron decays: A lepton flavor universality laboratory*, [arXiv:2101.08326](#).
- [8] BaBar collaboration, J. P. Lees *et al.*, *Evidence for an excess of $\bar{B} \rightarrow D^{(*)}\tau^-\bar{\nu}_\tau$ decays*, Phys. Rev. Lett. **109** (2012) 101802, [arXiv:1205.5442](#).
- [9] BaBar collaboration, J. P. Lees *et al.*, *Measurement of an excess of $\bar{B} \rightarrow D^{(*)}\tau^-\bar{\nu}_\tau$ decays and implications for charged Higgs bosons*, Phys. Rev. **D88** (2013) 072012, [arXiv:1303.0571](#).
- [10] Belle collaboration, M. Huschle *et al.*, *Measurement of the branching ratio of $\bar{B} \rightarrow D^{(*)}\tau^-\bar{\nu}_\tau$ relative to $\bar{B} \rightarrow D^{(*)}\ell^-\bar{\nu}_\ell$ decays with hadronic tagging at Belle*, Phys. Rev. **D92** (2015) 072014, [arXiv:1507.03233](#).
- [11] Belle collaboration, G. Caria *et al.*, *Measurement of $\mathcal{R}(D)$ and $\mathcal{R}(D^*)$ with a semileptonic tagging method*, Phys. Rev. Lett. **124** (2020) 161803, [arXiv:1910.05864](#).
- [12] Belle collaboration, S. Hirose *et al.*, *Measurement of the τ lepton polarization and $R(D^*)$ in the decay $\bar{B} \rightarrow D^*\tau^-\bar{\nu}_\tau$ with one-prong hadronic τ decays at Belle*, Phys. Rev. **D97** (2018) 012004, [arXiv:1709.00129](#).
- [13] LHCb collaboration, R. Aaij *et al.*, *Measurement of the ratio of branching fractions $\mathcal{B}(\bar{B}^0 \rightarrow D^{*+}\tau^-\bar{\nu}_\tau)/\mathcal{B}(\bar{B}^0 \rightarrow D^{*+}\mu^-\bar{\nu}_\mu)$* , Phys. Rev. Lett. **115** (2015) 111803, Publisher’s Note *ibid.* **115** (2015) 159901, [arXiv:1506.08614](#).

- [14] LHCb collaboration, R. Aaij *et al.*, *Test of lepton flavor universality by the measurement of the $B^0 \rightarrow D^{*-} \tau^+ \nu_\tau$ branching fraction using three-prong τ decays*, Phys. Rev. **D97** (2018) 072013, [arXiv:1711.02505](#).
- [15] Heavy Flavor Averaging Group, Y. Amhis *et al.*, *Averages of b-hadron, c-hadron, and τ -lepton properties as of 2018*, Eur. Phys. J. **C81** (2021) 226, [arXiv:1909.12524](#), updated results and plots available at <https://hflav.web.cern.ch>.
- [16] LHCb collaboration, A. A. Alves Jr. *et al.*, *The LHCb detector at the LHC*, JINST **3** (2008) S08005.
- [17] LHCb collaboration, R. Aaij *et al.*, *LHCb detector performance*, Int. J. Mod. Phys. **A30** (2015) 1530022, [arXiv:1412.6352](#).
- [18] R. Aaij *et al.*, *Performance of the LHCb Vertex Locator*, JINST **9** (2014) P09007, [arXiv:1405.7808](#).
- [19] R. Aaij *et al.*, *The LHCb trigger and its performance in 2011*, JINST **8** (2013) P04022, [arXiv:1211.3055](#).
- [20] T. Sjöstrand, S. Mrenna, and P. Skands, *PYTHIA 6.4 physics and manual*, JHEP **05** (2006) 026, [arXiv:hep-ph/0603175](#); T. Sjöstrand, S. Mrenna, and P. Skands, *A brief introduction to PYTHIA 8.1*, Comput. Phys. Commun. **178** (2008) 852, [arXiv:0710.3820](#).
- [21] I. Belyaev *et al.*, *Handling of the generation of primary events in Gauss, the LHCb simulation framework*, J. Phys. Conf. Ser. **331** (2011) 032047.
- [22] D. J. Lange, *The EvtGen particle decay simulation package*, Nucl. Instrum. Meth. **A462** (2001) 152.
- [23] N. Davidson, T. Przedzinski, and Z. Was, *PHOTOS interface in C++: Technical and physics documentation*, Comp. Phys. Comm. **199** (2016) 86, [arXiv:1011.0937](#).
- [24] N. Davidson *et al.*, *Universal interface of TAUOLA technical and physics documentation*, Comput. Phys. Commun. **183** (2012) 821, [arXiv:1002.0543](#).
- [25] I. M. Nugent *et al.*, *Resonance chiral Lagrangian currents and experimental data for $\tau^- \rightarrow \pi^- \pi^- \pi^+ \nu_\tau$* , Phys. Rev. **D88** (2013) 093012, [arXiv:1310.1053](#).
- [26] I. M. Nugent, *Invariant mass spectra of $\tau^- \rightarrow h^- h^- h^+ \nu_\tau$ decays*, Nucl. Phys. Proc. Suppl. **253-255** (2014) 38, [arXiv:1301.7105](#).
- [27] Geant4 collaboration, J. Allison *et al.*, *Geant4 developments and applications*, IEEE Trans. Nucl. Sci. **53** (2006) 270; Geant4 collaboration, S. Agostinelli *et al.*, *Geant4: A simulation toolkit*, Nucl. Instrum. Meth. **A506** (2003) 250.

- [28] M. Clemencic *et al.*, *The LHCb simulation application, Gauss: Design, evolution and experience*, J. Phys. Conf. Ser. **331** (2011) 032023.
- [29] I. Caprini, L. Lellouch, and M. Neubert, *Dispersive bounds on the shape of $\bar{B} \rightarrow D^{(*)}$ lepton anti-neutrino form-factors*, Nucl. Phys. **B530** (1998) 153, [arXiv:hep-ph/9712417](#).
- [30] LHCb collaboration, R. Aaij *et al.*, *A precise measurement of the B^0 meson oscillation frequency*, Eur. Phys. J. **C76** (2016) 412, [arXiv:1604.03475](#).
- [31] Particle Data Group, P. A. Zyla *et al.*, *Review of particle physics*, Prog. Theor. Exp. Phys. **2020** (2020) 083C01.
- [32] L. Breiman, J. H. Friedman, R. A. Olshen, and C. J. Stone, *Classification and regression trees*, Wadsworth international group, Belmont, California, USA, 1984.
- [33] R. E. Schapire and Y. Freund, *A decision-theoretic generalization of on-line learning and an application to boosting*, Jour. Comp. and Syst. Sc. **55** (1997) 119.
- [34] See Supplemental Material for the figure showing the distributions of the three main input variables to the BDT classifier, for signal and background, respectively.
- [35] See Supplemental Material for the figure showing the τ decay time distribution for low and high BDT output values.
- [36] See Supplemental Material for the figure showing the $\Lambda_c^+ \pi^- \pi^+ \pi^-$ mass distribution in the normalisation sample.
- [37] CDF collaboration, T. Aaltonen *et al.*, *Measurement of the branching fraction $\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^- \pi^+ \pi^-)$ at CDF*, Phys. Rev. **D85** (2012) 032003, [arXiv:1112.3334](#).
- [38] LHCb collaboration, R. Aaij *et al.*, *Measurements of the branching fractions for $B_{(s)}^0 \rightarrow D_{(s)} \pi \pi \pi$ and $\Lambda_b^0 \rightarrow \Lambda_c^+ \pi \pi \pi$* , Phys. Rev. **D84** (2011) 092001, Erratum *ibid.* **D85** (2012) 039904, [arXiv:1109.6831](#).
- [39] DELPHI collaboration, J. Abdallah *et al.*, *Measurement of the Λ_b^0 decay form-factor*, Phys. Lett. **B585** (2004) 63, [arXiv:hep-ex/0403040](#).

LHCb collaboration

R. Aaij³², A.S.W. Abdelmotteleb⁵⁶, C. Abellán Beteta⁵⁰, F. Abudinén⁵⁶, T. Ackernley⁶⁰, B. Adeva⁴⁶, M. Adinolfi⁵⁴, H. Afsharnia⁹, C. Agapopoulou¹³, C.A. Aidala⁸⁷, S. Aiola²⁵, Z. Ajaltouni⁹, S. Akar⁶⁵, J. Albrecht¹⁵, F. Alessio⁴⁸, M. Alexander⁵⁹, A. Alfonso Albero⁴⁵, Z. Aliouche⁶², G. Alkhazov³⁸, P. Alvarez Cartelle⁵⁵, S. Amato², J.L. Amey⁵⁴, Y. Amhis¹¹, L. An⁴⁸, L. Anderlini²², M. Andersson⁵⁰, A. Andreianov³⁸, M. Andreotti²¹, F. Archilli¹⁷, A. Artamonov⁴⁴, M. Artuso⁶⁸, K. Arzymatov⁴², E. Aslanides¹⁰, M. Atzeni⁵⁰, B. Audurier¹², S. Bachmann¹⁷, M. Bachmayer⁴⁹, J.J. Back⁵⁶, P. Baladron Rodriguez⁴⁶, V. Balagura¹², W. Baldini²¹, J. Baptista Leite¹, M. Barbetti^{22,h}, R.J. Barlow⁶², S. Barsuk¹¹, W. Barter⁶¹, M. Bartolini⁵⁵, F. Baryshnikov⁸³, J.M. Basels¹⁴, S. Bashir³⁴, G. Bassi²⁹, B. Batsukh⁶⁸, A. Battig¹⁵, A. Bay⁴⁹, A. Beck⁵⁶, M. Becker¹⁵, F. Bedeschi²⁹, I. Bediaga¹, A. Beiter⁶⁸, V. Belavin⁴², S. Belin²⁷, V. Bellee⁵⁰, K. Belous⁴⁴, I. Belov⁴⁰, I. Belyaev⁴¹, G. Bencivenni²³, E. Ben-Haim¹³, A. Berezhnoy⁴⁰, R. Bernet⁵⁰, D. Berninghoff¹⁷, H.C. Bernstein⁶⁸, C. Bertella⁶², A. Bertolin²⁸, C. Betancourt⁵⁰, F. Betti⁴⁸, Ia. Bezshyiko⁵⁰, S. Bhasin⁵⁴, J. Bhom³⁵, L. Bian⁷³, M.S. Bieker¹⁵, N.V. Biesuz²¹, S. Bifani⁵³, P. Billoir¹³, A. Biolchini³², M. Birch⁶¹, F.C.R. Bishop⁵⁵, A. Bitadze⁶², A. Bizzeti^{22,l}, M. Bjørn⁶³, M.P. Blago⁵⁵, T. Blake⁵⁶, F. Blanc⁴⁹, S. Blusk⁶⁸, D. Bobulska⁵⁹, J.A. Boelhauve¹⁵, O. Boente Garcia⁴⁶, T. Boettcher⁶⁵, A. Boldyrev⁸², A. Bondar⁴³, N. Bondar^{38,48}, S. Borghi⁶², M. Borisyak⁴², M. Borsato¹⁷, J.T. Borsuk³⁵, S.A. Bouchiba⁴⁹, T.J.V. Bowcock^{60,48}, A. Boyer⁴⁸, C. Bozzi²¹, M.J. Bradley⁶¹, S. Braun⁶⁶, A. Brea Rodriguez⁴⁶, J. Brodzicka³⁵, A. Brossa Gonzalo⁵⁶, D. Brundu²⁷, A. Buonauro⁵⁰, L. Buonincontri²⁸, A.T. Burke⁶², C. Burr⁴⁸, A. Bursche⁷², A. Butkevich³⁹, J.S. Butter³², J. Buytaert⁴⁸, W. Byczynski⁴⁸, S. Cadeddu²⁷, H. Cai⁷³, R. Calabrese^{21,g}, L. Calefice^{15,13}, S. Cali²³, R. Calladine⁵³, M. Calvi^{26,k}, M. Calvo Gomez⁸⁵, P. Camargo Magalhaes⁵⁴, P. Campana²³, A.F. Campoverde Quezada⁶, S. Capelli^{26,k}, L. Capriotti^{20,e}, A. Carbone^{20,e}, G. Carboni^{31,q}, R. Cardinale^{24,i}, A. Cardini²⁷, I. Carli⁴, P. Carniti^{26,k}, L. Carus¹⁴, K. Carvalho Akiba³², A. Casais Vidal⁴⁶, R. Caspary¹⁷, G. Casse⁶⁰, M. Cattaneo⁴⁸, G. Cavallero⁴⁸, S. Celani⁴⁹, J. Cerasoli¹⁰, D. Cervenkov⁶³, A.J. Chadwick⁶⁰, M.G. Chapman⁵⁴, M. Charles¹³, Ph. Charpentier⁴⁸, C.A. Chavez Barajas⁶⁰, M. Chefdeville⁸, C. Chen³, S. Chen⁴, A. Chernov³⁵, V. Chobanova⁴⁶, S. Cholak⁴⁹, M. Chruszcz³⁵, A. Chubykin³⁸, V. Chulikov³⁸, P. Ciambone²³, M.F. Cicala⁵⁶, X. Cid Vidal⁴⁶, G. Ciezarek⁴⁸, P.E.L. Clarke⁵⁸, M. Clemencic⁴⁸, H.V. Cliff⁵⁵, J. Closier⁴⁸, J.L. Cobbledick⁶², V. Coco⁴⁸, J.A.B. Coelho¹¹, J. Cogan¹⁰, E. Cogneras⁹, L. Cojocariu³⁷, P. Collins⁴⁸, T. Colombo⁴⁸, L. Congedo^{19,d}, A. Contu²⁷, N. Cooke⁵³, G. Coombs⁵⁹, I. Corredoira⁴⁶, G. Corti⁴⁸, C.M. Costa Sobral⁵⁶, B. Couturier⁴⁸, D.C. Craik⁶⁴, J. Crkovská⁶⁷, M. Cruz Torres¹, R. Currie⁵⁸, C.L. Da Silva⁶⁷, S. Dadabaev⁸³, L. Dai⁷¹, E. Dall’Occo¹⁵, J. Dalseno⁴⁶, C. D’Ambrosio⁴⁸, A. Danilina⁴¹, P. d’Argent⁴⁸, A. Dashkina⁸³, J.E. Davies⁶², A. Davis⁶², O. De Aguiar Francisco⁶², K. De Bruyn⁷⁹, S. De Capua⁶², M. De Cian⁴⁹, U. De Freitas Carneiro Da Graca¹, E. De Lucia²³, J.M. De Miranda¹, L. De Paula², M. De Serio^{19,d}, D. De Simone⁵⁰, P. De Simone²³, F. De Vellis¹⁵, J.A. de Vries⁸⁰, C.T. Dean⁶⁷, F. Debernardis^{19,d}, D. Decamp⁸, V. Dedu¹⁰, L. Del Buono¹³, B. Delaney⁵⁵, H.-P. Dembinski¹⁵, V. Denysenko⁵⁰, D. Derkach⁸², O. Deschamps⁹, F. Dettori^{27,f}, B. Dey⁷⁷, A. Di Cicco²³, P. Di Nezza²³, S. Didenko⁸³, L. Dieste Maronas⁴⁶, H. Dijkstra⁴⁸, V. Dobishuk⁵², C. Dong³, A.M. Donohoe¹⁸, F. Dordei²⁷, A.C. dos Reis¹, L. Douglas⁵⁹, A. Dovbnaya⁵¹, A.G. Downes⁸, M.W. Dudek³⁵, L. Dufour⁴⁸, V. Duk⁷⁸, P. Durante⁴⁸, J.M. Durham⁶⁷, D. Dutta⁶², A. Dziurda³⁵, A. Dzyuba³⁸, S. Easo⁵⁷, U. Egede⁶⁹, V. Egorychev⁴¹, S. Eidelman^{43,u,†}, S. Eisenhardt⁵⁸, S. Ek-In⁴⁹, L. Eklund⁸⁶,

S. Ely⁶⁸, A. Enc³⁷, E. Epple⁶⁷, S. Escher¹⁴, J. Eschle⁵⁰, S. Esen⁵⁰, T. Evans⁶², L.N. Falcao¹,
 Y. Fan⁶, B. Fang⁷³, S. Farry⁶⁰, D. Fazzini^{26,k}, M. Féo⁴⁸, A. Fernandez Prieto⁴⁶, A.D. Fernandez⁶⁶,
 F. Ferrari²⁰, L. Ferreira Lopes⁴⁹, F. Ferreira Rodrigues², S. Ferreres Sole³², M. Ferrillo⁵⁰,
 M. Ferro-Luzzi⁴⁸, S. Filippov³⁹, R.A. Fini¹⁹, M. Fiorini^{21,g}, M. Firlej³⁴, K.M. Fischer⁶³,
 D.S. Fitzgerald⁸⁷, C. Fitzpatrick⁶², T. Fiutowski³⁴, A. Fkiaras⁴⁸, F. Fleuret¹², M. Fontana¹³,
 F. Fontanelli^{24,i}, R. Forty⁴⁸, D. Foulds-Holt⁵⁵, V. Franco Lima⁶⁰, M. Franco Sevilla⁶⁶,
 M. Frank⁴⁸, E. Franzoso²¹, G. Frau¹⁷, C. Frei⁴⁸, D.A. Friday⁵⁹, J. Fu⁶, Q. Fuehring¹⁵,
 E. Gabriel³², G. Galati^{19,d}, A. Gallas Torreira⁴⁶, D. Galli^{20,e}, S. Gambetta^{58,48}, Y. Gan³,
 M. Gandelman², P. Gandini²⁵, Y. Gao⁵, M. Garau²⁷, L.M. Garcia Martin⁵⁶, P. Garcia Moreno⁴⁵,
 J. García Pardiñas^{26,k}, B. Garcia Plana⁴⁶, F.A. Garcia Rosales¹², L. Garrido⁴⁵, C. Gaspar⁴⁸,
 R.E. Geertsema³², D. Gerick¹⁷, L.L. Gerken¹⁵, E. Gersabeck⁶², M. Gersabeck⁶², T. Gershon⁵⁶,
 D. Gerstel¹⁰, L. Giambastiani²⁸, V. Gibson⁵⁵, H.K. Gienza³⁶, A.L. Gilman⁶³,
 M. Giovannetti^{23,q}, A. Gioventù⁴⁶, P. Gironella Gironell⁴⁵, C. Giugliano²¹, K. Gizdov⁵⁸,
 E.L. Gkougkousis⁴⁸, V.V. Gligorov^{13,48}, C. Göbel⁷⁰, E. Golobardes⁸⁵, D. Golubkov⁴¹,
 A. Golutvin^{61,83}, A. Gomes^{1,a}, S. Gomez Fernandez⁴⁵, F. Goncalves Abrantes⁶³, M. Goncerz³⁵,
 G. Gong³, P. Gorbounov⁴¹, I.V. Gorelov⁴⁰, C. Gotti²⁶, J.P. Grabowski¹⁷, T. Grammatico¹³,
 L.A. Granado Cardoso⁴⁸, E. Graugés⁴⁵, E. Graverini⁴⁹, G. Graziani²², A. Grecu³⁷,
 L.M. Greeven³², N.A. Grieser⁴, L. Grillo⁶², S. Gromov⁸³, B.R. Gruberg Cazon⁶³, C. Gu³,
 M. Guarise²¹, M. Guittiere¹¹, P. A. Günther¹⁷, E. Gushchin³⁹, A. Guth¹⁴, Y. Guz⁴⁴, T. Gys⁴⁸,
 T. Hadavizadeh⁶⁹, G. Haefeli⁴⁹, C. Haen⁴⁸, J. Haimberger⁴⁸, S.C. Haines⁵⁵,
 T. Halewood-leagas⁶⁰, P.M. Hamilton⁶⁶, J.P. Hammerich⁶⁰, Q. Han⁷, X. Han¹⁷, E.B. Hansen⁶²,
 S. Hansmann-Menzemer¹⁷, N. Harnew⁶³, T. Harrison⁶⁰, C. Hasse⁴⁸, M. Hatch⁴⁸, J. He^{6,b},
 M. Hecker⁶¹, K. Heijhoff³², K. Heinicke¹⁵, R.D.L. Henderson^{69,56}, A.M. Hennequin⁴⁸,
 K. Hennessy⁶⁰, L. Henry⁴⁸, J. Heuel¹⁴, A. Hicheur², D. Hill⁴⁹, M. Hilton⁶², S.E. Hollitt¹⁵,
 R. Hou⁷, Y. Hou⁸, J. Hu¹⁷, J. Hu⁷², W. Hu⁷, X. Hu³, W. Huang⁶, X. Huang⁷³, W. Hulsbergen³²,
 R.J. Hunter⁵⁶, M. Hushchyn⁸², D. Hutchcroft⁶⁰, D. Hynds³², P. Ibis¹⁵, M. Idzik³⁴, D. Ilin³⁸,
 P. Ilten⁶⁵, A. Inglessi³⁸, A. Ishteev⁸³, K. Ivshin³⁸, R. Jacobsson⁴⁸, H. Jage¹⁴, S. Jakobsen⁴⁸,
 E. Jans³², B.K. Jashal⁴⁷, A. Jawahery⁶⁶, V. Jevtic¹⁵, X. Jiang⁴, M. John⁶³, D. Johnson⁶⁴,
 C.R. Jones⁵⁵, T.P. Jones⁵⁶, B. Jost⁴⁸, N. Jurik⁴⁸, S.H. Kalavan Kadavath³⁴, S. Kandybei⁵¹,
 Y. Kang³, M. Karacson⁴⁸, D. Karpenkov⁸³, M. Karpov⁸², J.W. Kautz⁶⁵, F. Keizer⁴⁸,
 D.M. Keller⁶⁸, M. Kenzie⁵⁶, T. Ketel³³, B. Khanji¹⁵, A. Kharisova⁸⁴, S. Kholodenko⁴⁴,
 T. Kirn¹⁴, V.S. Kirsebom⁴⁹, O. Kitouni⁶⁴, S. Klaver³³, N. Kleijne²⁹, K. Klimaszewski³⁶,
 M.R. Kmiec³⁶, S. Koliiev⁵², A. Kondybayeva⁸³, A. Konoplyannikov⁴¹, P. Kopciwicz³⁴,
 R. Kopečna¹⁷, P. Koppenburg³², M. Korolev⁴⁰, I. Kostiuk^{32,52}, O. Kot⁵², S. Kotriakhova^{21,38},
 A. Kozachuk⁴⁰, P. Kravchenko³⁸, L. Kravchuk³⁹, R.D. Krawczyk⁴⁸, M. Kreps⁵⁶,
 S. Kretschmar¹⁴, P. Krokovny^{43,u}, W. Krupa³⁴, W. Krzemien³⁶, J. Kubat¹⁷, M. Kucharczyk³⁵,
 V. Kudryavtsev^{43,u}, H.S. Kuindersma^{32,33}, G.J. Kunde⁶⁷, T. Kvaratskheliya⁴¹, D. Lacarrere⁴⁸,
 G. Lafferty⁶², A. Lai²⁷, A. Lampis²⁷, D. Lancierini⁵⁰, J.J. Lane⁶², R. Lane⁵⁴, G. Lanfranchi²³,
 C. Langenbruch¹⁴, J. Langer¹⁵, O. Lantwin⁸³, T. Latham⁵⁶, F. Lazzari²⁹, R. Le Gac¹⁰,
 S.H. Lee⁸⁷, R. Lefèvre⁹, A. Leflat⁴⁰, S. Legotin⁸³, O. Leroy¹⁰, T. Lesiak³⁵, B. Leverington¹⁷,
 H. Li⁷², P. Li¹⁷, S. Li⁷, Y. Li⁴, Z. Li⁶⁸, X. Liang⁶⁸, T. Lin⁶¹, R. Lindner⁴⁸, V. Lisovskyi¹⁵,
 R. Litvinov²⁷, G. Liu⁷², H. Liu⁶, Q. Liu⁶, S. Liu⁴, A. Lobo Salvia⁴⁵, A. Loi²⁷, R. Lollini⁷⁸,
 J. Lomba Castro⁴⁶, I. Longstaff⁵⁹, J.H. Lopes², S. López Soliño⁴⁶, G.H. Lovell⁵⁵, Y. Lu⁴,
 C. Lucarelli^{22,h}, D. Lucchesi^{28,m}, S. Luchuk³⁹, M. Lucio Martinez³², V. Lukashenko^{32,52},
 Y. Luo³, A. Lupato⁶², E. Luppi^{21,g}, O. Lupton⁵⁶, A. Lusiani^{29,n}, X. Lyu⁶, L. Ma⁴, R. Ma⁶,
 S. Maccolini²⁰, F. Machefert¹¹, F. Maciuc³⁷, V. Macko⁴⁹, P. Mackowiak¹⁵,

S. Maddrell-Mander⁵⁴, O. Madejczyk³⁴, L.R. Madhan Mohan⁵⁴, O. Maev³⁸, A. Maevskiy⁸²,
 D. Maisuzenko³⁸, M.W. Majewski³⁴, J.J. Malczewski³⁵, S. Malde⁶³, B. Malecki³⁵, A. Malinin⁸¹,
 T. Maltsev^{43,u}, H. Malygina¹⁷, G. Manca^{27,f}, G. Mancinelli¹⁰, D. Manuzzi²⁰, D. Marangotto^{25,j},
 J. Maratas^{9,s}, J.F. Marchand⁸, U. Marconi²⁰, S. Mariani^{22,h}, C. Marin Benito⁴⁸,
 M. Marinangeli⁴⁹, J. Marks¹⁷, A.M. Marshall⁵⁴, P.J. Marshall⁶⁰, G. Martelli⁷⁸, G. Martellotti³⁰,
 L. Martinazzoli^{48,k}, M. Martinelli^{26,k}, D. Martinez Santos⁴⁶, F. Martinez Vidal⁴⁷,
 A. Massafferri¹, M. Materok¹⁴, R. Matev⁴⁸, A. Mathad⁵⁰, V. Matiunin⁴¹, C. Matteuzzi²⁶,
 K.R. Mattioli⁸⁷, A. Mauri³², E. Maurice¹², J. Mauricio⁴⁵, M. Mazurek⁴⁸, M. McCann⁶¹,
 L. Mcconnell¹⁸, T.H. Mcgrath⁶², N.T. Mchugh⁵⁹, A. McNab⁶², R. McNulty¹⁸, J.V. Mead⁶⁰,
 B. Meadows⁶⁵, G. Meier¹⁵, D. Melnychuk³⁶, S. Meloni^{26,k}, M. Merk^{32,80}, A. Merli^{25,j},
 L. Meyer Garcia², M. Mikhasenko^{75,c}, D.A. Milanes⁷⁴, E. Millard⁵⁶, M. Milovanovic⁴⁸,
 M.-N. Minard⁸, A. Minotti^{26,k}, S.E. Mitchell⁵⁸, B. Mitreska⁶², D.S. Mitzel¹⁵, A. Mödden¹⁵,
 R.A. Mohammed⁶³, R.D. Moise⁶¹, S. Mokhnenko⁸², T. Mombächer⁴⁶, I.A. Monroy⁷⁴,
 S. Monteil⁹, M. Morandin²⁸, G. Morello²³, M.J. Morello^{29,n}, J. Moron³⁴, A.B. Morris⁷⁵,
 A.G. Morris⁵⁶, R. Mountain⁶⁸, H. Mu³, F. Muheim⁵⁸, M. Mulder⁷⁹, K. Müller⁵⁰, C.H. Murphy⁶³,
 D. Murray⁶², R. Murta⁶¹, P. Muzzetto²⁷, P. Naik⁵⁴, T. Nakada⁴⁹, R. Nandakumar⁵⁷,
 T. Nanut⁴⁸, I. Nasteva², M. Needham⁵⁸, N. Neri^{25,j}, S. Neubert⁷⁵, N. Neufeld⁴⁸,
 R. Newcombe⁶¹, E.M. Niel⁴⁹, S. Nieswand¹⁴, N. Nikitin⁴⁰, N.S. Nolte⁶⁴, C. Normand⁸,
 C. Nunez⁸⁷, A. Oblakowska-Mucha³⁴, V. Obraztsov⁴⁴, T. Oeser¹⁴, D.P. O’Hanlon⁵⁴,
 S. Okamura²¹, R. Oldeman^{27,f}, F. Oliva⁵⁸, M.E. Olivares⁶⁸, C.J.G. Onderwater⁷⁹, R.H. O’Neil⁵⁸,
 J.M. Otalora Goicochea², T. Ovsianikova⁴¹, P. Owen⁵⁰, A. Oyanguren⁴⁷, O. Ozcelik⁵⁸,
 K.O. Padeken⁷⁵, B. Pagare⁵⁶, P.R. Pais⁴⁸, T. Pajero⁶³, A. Palano¹⁹, M. Palutan²³, Y. Pan⁶²,
 G. Panshin⁸⁴, A. Papanestis⁵⁷, M. Pappagallo^{19,d}, L.L. Pappalardo^{21,g}, C. Pappenheimer⁶⁵,
 W. Parker⁶⁶, C. Parkes⁶², B. Passalacqua²¹, G. Passaleva²², A. Pastore¹⁹, M. Patel⁶¹,
 C. Patrignani^{20,e}, C.J. Pawley⁸⁰, A. Pearce^{48,57}, A. Pellegrino³², M. Pepe Altarelli⁴⁸,
 S. Perazzini²⁰, D. Pereima⁴¹, A. Pereiro Castro⁴⁶, P. Perret⁹, M. Petric^{59,48}, K. Petridis⁵⁴,
 A. Petrolini^{24,i}, A. Petrov⁸¹, S. Petrucci⁵⁸, M. Petruzzo²⁵, T.T.H. Phan⁶⁸, A. Philippov⁴²,
 R. Piandani⁶, L. Pica^{29,n}, M. Piccini⁷⁸, B. Pietrzyk⁸, G. Pietrzyk¹¹, M. Pili⁶³, D. Pinci³⁰,
 F. Pisani⁴⁸, M. Pizzichemi^{26,48,k}, P.K. Resmi¹⁰, V. Placinta³⁷, J. Plews⁵³, M. Plo Casasus⁴⁶,
 F. Polci^{13,48}, M. Poli Lener²³, M. Poliakov⁶⁸, A. Poluektov¹⁰, N. Polukhina^{83,t}, I. Polyakov⁶⁸,
 E. Polcarpo², S. Ponce⁴⁸, D. Popov^{6,48}, S. Popov⁴², S. Poslavskii⁴⁴, K. Prasanth³⁵,
 L. Promberger⁴⁸, C. Prouve⁴⁶, V. Pugatch⁵², V. Puill¹¹, G. Punzi^{29,o}, H. Qi³, W. Qian⁶,
 N. Qin³, R. Quagliani⁴⁹, N.V. Raab¹⁸, R.I. Rabadan Trejo⁶, B. Rachwal³⁴, J.H. Rademacker⁵⁴,
 R. Rajagopalan⁶⁸, M. Rama²⁹, M. Ramos Pernas⁵⁶, M.S. Rangel², F. Ratnikov^{42,82},
 G. Raven^{33,48}, M. Reboud⁸, F. Redi⁴⁸, F. Reiss⁶², C. Remon Alepuz⁴⁷, Z. Ren³, V. Renaudin⁶³,
 R. Ribatti²⁹, A.M. Ricci²⁷, S. Ricciardi⁵⁷, K. Rinnert⁶⁰, P. Robbe¹¹, G. Robertson⁵⁸,
 A.B. Rodrigues⁴⁹, E. Rodrigues⁶⁰, J.A. Rodriguez Lopez⁷⁴, E.R.R. Rodriguez Rodriguez⁴⁶,
 A. Rollings⁶³, P. Roloff⁴⁸, V. Romanovskiy⁴⁴, M. Romero Lamas⁴⁶, A. Romero Vidal⁴⁶,
 J.D. Roth⁸⁷, M. Rotondo²³, M.S. Rudolph⁶⁸, T. Ruf⁴⁸, R.A. Ruiz Fernandez⁴⁶, J. Ruiz Vidal⁴⁷,
 A. Ryzhikov⁸², J. Ryzka³⁴, J.J. Saborido Silva⁴⁶, N. Sagidova³⁸, N. Sahoo⁵³, B. Saitta^{27,f},
 M. Salomoni⁴⁸, C. Sanchez Gras³², R. Santacesaria³⁰, C. Santamarina Rios⁴⁶, M. Santimaria²³,
 E. Santovetti^{31,q}, D. Saranin⁸³, G. Sarpis¹⁴, M. Sarpis⁷⁵, A. Sarti³⁰, C. Satriano^{30,p}, A. Satta³¹,
 M. Saur¹⁵, D. Savrina^{41,40}, H. Sazak⁹, L.G. Scantlebury Smead⁶³, A. Scarabotto¹³, S. Schael¹⁴,
 S. Scherl⁶⁰, M. Schiller⁵⁹, H. Schindler⁴⁸, M. Schmelling¹⁶, B. Schmidt⁴⁸, S. Schmitt¹⁴,
 O. Schneider⁴⁹, A. Schopper⁴⁸, M. Schubiger³², S. Schulte⁴⁹, M.H. Schune¹¹, R. Schwemmer⁴⁸,
 B. Sciascia^{23,48}, S. Sellam⁴⁶, A. Semennikov⁴¹, M. Senghi Soares³³, A. Sergi^{24,i}, N. Serra⁵⁰,

L. Sestini²⁸, A. Seuthe¹⁵, Y. Shang⁵, D.M. Shangase⁸⁷, M. Shapkin⁴⁴, I. Shchemerov⁸³, L. Shchutska⁴⁹, T. Shears⁶⁰, L. Shekhtman^{43,u}, Z. Shen⁵, S. Sheng⁴, V. Shevchenko⁸¹, E.B. Shields^{26,k}, Y. Shimizu¹¹, E. Shmanin⁸³, J.D. Shupperd⁶⁸, B.G. Siddi²¹, R. Silva Coutinho⁵⁰, G. Simi²⁸, S. Simone^{19,d}, N. Skidmore⁶², R. Skuza¹⁷, T. Skwarnicki⁶⁸, M.W. Slater⁵³, I. Slazyk^{21,g}, J.C. Smallwood⁶³, J.G. Smeaton⁵⁵, E. Smith⁵⁰, M. Smith⁶¹, A. Snoch³², L. Soares Lavra⁹, M.D. Sokoloff⁶⁵, F.J.P. Soler⁵⁹, A. Solovev³⁸, I. Solovyev³⁸, F.L. Souza De Almeida², B. Souza De Paula², B. Spaan¹⁵, E. Spadaro Norella^{25,j}, P. Spradlin⁵⁹, F. Stagni⁴⁸, M. Stahl⁶⁵, S. Stahl⁴⁸, S. Stanislaus⁶³, O. Steinkamp^{50,83}, O. Stenyakin⁴⁴, H. Stevens¹⁵, S. Stone^{68,48,†}, D. Strelakina⁸³, F. Suljik⁶³, J. Sun²⁷, L. Sun⁷³, Y. Sun⁶⁶, P. Svihra⁶², P.N. Swallow⁵³, K. Swientek³⁴, A. Szabelski³⁶, T. Szumlak³⁴, M. Szymanski⁴⁸, S. Taneja⁶², A.R. Tanner⁵⁴, M.D. Tat⁶³, A. Terentev⁸³, F. Teubert⁴⁸, E. Thomas⁴⁸, D.J.D. Thompson⁵³, K.A. Thomson⁶⁰, H. Tilquin⁶¹, V. Tisserand⁹, S. T'Jampens⁸, M. Tobin⁴, L. Tomassetti^{21,g}, X. Tong⁵, D. Torres Machado¹, D.Y. Tou³, E. Trifonova⁸³, S.M. Trilov⁵⁴, C. Trippel⁴⁹, G. Tuci⁶, A. Tully⁴⁹, N. Tuning^{32,48}, A. Ukleja^{36,48}, D.J. Unverzagt¹⁷, E. Ursov⁸³, A. Usachov³², A. Ustyuzhanin^{42,82}, U. Uwer¹⁷, A. Vagner⁸⁴, V. Vagnoni²⁰, A. Valassi⁴⁸, G. Valenti²⁰, N. Valls Canudas⁸⁵, M. van Beuzekom³², M. Van Dijk⁴⁹, H. Van Hecke⁶⁷, E. van Herwijnen⁸³, M. van Veghel⁷⁹, R. Vazquez Gomez⁴⁵, P. Vazquez Regueiro⁴⁶, C. Vázquez Sierra⁴⁸, S. Vecchi²¹, J.J. Velthuis⁵⁴, M. Veltri^{22,r}, A. Venkateswaran⁶⁸, M. Veronesi³², M. Vesterinen⁵⁶, D. Vieira⁶⁵, M. Vieites Diaz⁴⁹, H. Viemann⁷⁶, X. Vilasis-Cardona⁸⁵, E. Vilella Figueras⁶⁰, A. Villa²⁰, P. Vincent¹³, F.C. Volle¹¹, D. Vom Bruch¹⁰, A. Vorobyev³⁸, V. Vorobyev^{43,u}, N. Voropaev³⁸, K. Vos⁸⁰, R. Waldi¹⁷, J. Walsh²⁹, C. Wang¹⁷, J. Wang⁵, J. Wang⁴, J. Wang³, J. Wang⁷³, M. Wang³, R. Wang⁵⁴, Y. Wang⁷, Z. Wang⁵⁰, Z. Wang³, Z. Wang⁶, J.A. Ward^{56,69}, N.K. Watson⁵³, D. Websdale⁶¹, C. Weisser⁶⁴, B.D.C. Westhenry⁵⁴, D.J. White⁶², M. Whitehead⁵⁴, A.R. Wiederhold⁵⁶, D. Wiedner¹⁵, G. Wilkinson⁶³, M. Wilkinson⁶⁸, I. Williams⁵⁵, M. Williams⁶⁴, M.R.J. Williams⁵⁸, F.F. Wilson⁵⁷, W. Wislicki³⁶, M. Witek³⁵, L. Witola¹⁷, G. Wormser¹¹, S.A. Wotton⁵⁵, H. Wu⁶⁸, K. Wyllie⁴⁸, Z. Xiang⁶, D. Xiao⁷, Y. Xie⁷, A. Xu⁵, J. Xu⁶, L. Xu³, M. Xu⁵⁶, Q. Xu⁶, Z. Xu⁹, Z. Xu⁶, D. Yang³, S. Yang⁶, Y. Yang⁶, Z. Yang⁵, Z. Yang⁶⁶, Y. Yao⁶⁸, L.E. Yeomans⁶⁰, H. Yin⁷, J. Yu⁷¹, X. Yuan⁶⁸, O. Yushchenko⁴⁴, E. Zaffaroni⁴⁹, M. Zavertyaev^{16,t}, M. Zdybal³⁵, O. Zenaiev⁴⁸, M. Zeng³, D. Zhang⁷, L. Zhang³, S. Zhang⁷¹, S. Zhang⁵, Y. Zhang⁵, Y. Zhang⁶³, A. Zharkova⁸³, A. Zhelezov¹⁷, Y. Zheng⁶, T. Zhou⁵, X. Zhou⁶, Y. Zhou⁶, V. Zhovkovska¹¹, X. Zhu³, X. Zhu⁷, Z. Zhu⁶, V. Zhukov^{14,40}, Q. Zou⁴, S. Zucchelli^{20,e}, D. Zuliani²⁸, G. Zunica⁶².

¹Centro Brasileiro de Pesquisas Físicas (CBPF), Rio de Janeiro, Brazil

²Universidade Federal do Rio de Janeiro (UFRJ), Rio de Janeiro, Brazil

³Center for High Energy Physics, Tsinghua University, Beijing, China

⁴Institute Of High Energy Physics (IHEP), Beijing, China

⁵School of Physics State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China

⁶University of Chinese Academy of Sciences, Beijing, China

⁷Institute of Particle Physics, Central China Normal University, Wuhan, Hubei, China

⁸Université Savoie Mont Blanc, CNRS, IN2P3-LAPP, Annecy, France

⁹Université Clermont Auvergne, CNRS/IN2P3, LPC, Clermont-Ferrand, France

¹⁰Aix Marseille Université, CNRS/IN2P3, CPPM, Marseille, France

¹¹Université Paris-Saclay, CNRS/IN2P3, IJCLab, Orsay, France

¹²Laboratoire Leprince-Ringuet, CNRS/IN2P3, Ecole Polytechnique, Institut Polytechnique de Paris, Palaiseau, France

¹³LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris, France

- ¹⁴*I. Physikalisches Institut, RWTH Aachen University, Aachen, Germany*
- ¹⁵*Fakultät Physik, Technische Universität Dortmund, Dortmund, Germany*
- ¹⁶*Max-Planck-Institut für Kernphysik (MPIK), Heidelberg, Germany*
- ¹⁷*Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
- ¹⁸*School of Physics, University College Dublin, Dublin, Ireland*
- ¹⁹*INFN Sezione di Bari, Bari, Italy*
- ²⁰*INFN Sezione di Bologna, Bologna, Italy*
- ²¹*INFN Sezione di Ferrara, Ferrara, Italy*
- ²²*INFN Sezione di Firenze, Firenze, Italy*
- ²³*INFN Laboratori Nazionali di Frascati, Frascati, Italy*
- ²⁴*INFN Sezione di Genova, Genova, Italy*
- ²⁵*INFN Sezione di Milano, Milano, Italy*
- ²⁶*INFN Sezione di Milano-Bicocca, Milano, Italy*
- ²⁷*INFN Sezione di Cagliari, Monserrato, Italy*
- ²⁸*Università degli Studi di Padova, Università e INFN, Padova, Padova, Italy*
- ²⁹*INFN Sezione di Pisa, Pisa, Italy*
- ³⁰*INFN Sezione di Roma La Sapienza, Roma, Italy*
- ³¹*INFN Sezione di Roma Tor Vergata, Roma, Italy*
- ³²*Nikhef National Institute for Subatomic Physics, Amsterdam, Netherlands*
- ³³*Nikhef National Institute for Subatomic Physics and VU University Amsterdam, Amsterdam, Netherlands*
- ³⁴*AGH - University of Science and Technology, Faculty of Physics and Applied Computer Science, Kraków, Poland*
- ³⁵*Henryk Niewodniczanski Institute of Nuclear Physics Polish Academy of Sciences, Kraków, Poland*
- ³⁶*National Center for Nuclear Research (NCBJ), Warsaw, Poland*
- ³⁷*Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest-Magurele, Romania*
- ³⁸*Petersburg Nuclear Physics Institute NRC Kurchatov Institute (PNPI NRC KI), Gatchina, Russia*
- ³⁹*Institute for Nuclear Research of the Russian Academy of Sciences (INR RAS), Moscow, Russia*
- ⁴⁰*Institute of Nuclear Physics, Moscow State University (SINP MSU), Moscow, Russia*
- ⁴¹*Institute of Theoretical and Experimental Physics NRC Kurchatov Institute (ITEP NRC KI), Moscow, Russia*
- ⁴²*Yandex School of Data Analysis, Moscow, Russia*
- ⁴³*Budker Institute of Nuclear Physics (SB RAS), Novosibirsk, Russia*
- ⁴⁴*Institute for High Energy Physics NRC Kurchatov Institute (IHEP NRC KI), Protvino, Russia, Protvino, Russia*
- ⁴⁵*ICCUB, Universitat de Barcelona, Barcelona, Spain*
- ⁴⁶*Instituto Galego de Física de Altas Enerxías (IGFAE), Universidade de Santiago de Compostela, Santiago de Compostela, Spain*
- ⁴⁷*Instituto de Física Corpuscular, Centro Mixto Universidad de Valencia - CSIC, Valencia, Spain*
- ⁴⁸*European Organization for Nuclear Research (CERN), Geneva, Switzerland*
- ⁴⁹*Institute of Physics, Ecole Polytechnique Fédérale de Lausanne (EPFL), Lausanne, Switzerland*
- ⁵⁰*Physik-Institut, Universität Zürich, Zürich, Switzerland*
- ⁵¹*NSC Kharkiv Institute of Physics and Technology (NSC KIPT), Kharkiv, Ukraine*
- ⁵²*Institute for Nuclear Research of the National Academy of Sciences (KINR), Kyiv, Ukraine*
- ⁵³*University of Birmingham, Birmingham, United Kingdom*
- ⁵⁴*H.H. Wills Physics Laboratory, University of Bristol, Bristol, United Kingdom*
- ⁵⁵*Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom*
- ⁵⁶*Department of Physics, University of Warwick, Coventry, United Kingdom*
- ⁵⁷*STFC Rutherford Appleton Laboratory, Didcot, United Kingdom*
- ⁵⁸*School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom*
- ⁵⁹*School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom*
- ⁶⁰*Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom*

- ⁶¹ *Imperial College London, London, United Kingdom*
- ⁶² *Department of Physics and Astronomy, University of Manchester, Manchester, United Kingdom*
- ⁶³ *Department of Physics, University of Oxford, Oxford, United Kingdom*
- ⁶⁴ *Massachusetts Institute of Technology, Cambridge, MA, United States*
- ⁶⁵ *University of Cincinnati, Cincinnati, OH, United States*
- ⁶⁶ *University of Maryland, College Park, MD, United States*
- ⁶⁷ *Los Alamos National Laboratory (LANL), Los Alamos, New Mexico, United States*
- ⁶⁸ *Syracuse University, Syracuse, NY, United States*
- ⁶⁹ *School of Physics and Astronomy, Monash University, Melbourne, Australia, associated to ⁵⁶*
- ⁷⁰ *Pontifícia Universidade Católica do Rio de Janeiro (PUC-Rio), Rio de Janeiro, Brazil, associated to ²*
- ⁷¹ *Physics and Micro Electronic College, Hunan University, Changsha City, China, associated to ⁷*
- ⁷² *Guangdong Provincial Key Laboratory of Nuclear Science, Guangdong-Hong Kong Joint Laboratory of Quantum Matter, Institute of Quantum Matter, South China Normal University, Guangzhou, China, associated to ³*
- ⁷³ *School of Physics and Technology, Wuhan University, Wuhan, China, associated to ³*
- ⁷⁴ *Departamento de Física, Universidad Nacional de Colombia, Bogota, Colombia, associated to ¹³*
- ⁷⁵ *Universität Bonn - Helmholtz-Institut für Strahlen und Kernphysik, Bonn, Germany, associated to ¹⁷*
- ⁷⁶ *Institut für Physik, Universität Rostock, Rostock, Germany, associated to ¹⁷*
- ⁷⁷ *Eotvos Lorand University, Budapest, Hungary, associated to ⁴⁸*
- ⁷⁸ *INFN Sezione di Perugia, Perugia, Italy, associated to ²¹*
- ⁷⁹ *Van Swinderen Institute, University of Groningen, Groningen, Netherlands, associated to ³²*
- ⁸⁰ *Universiteit Maastricht, Maastricht, Netherlands, associated to ³²*
- ⁸¹ *National Research Centre Kurchatov Institute, Moscow, Russia, associated to ⁴¹*
- ⁸² *National Research University Higher School of Economics, Moscow, Russia, associated to ⁴²*
- ⁸³ *National University of Science and Technology “MISIS”, Moscow, Russia, associated to ⁴¹*
- ⁸⁴ *National Research Tomsk Polytechnic University, Tomsk, Russia, associated to ⁴¹*
- ⁸⁵ *DS4DS, La Salle, Universitat Ramon Llull, Barcelona, Spain, associated to ⁴⁵*
- ⁸⁶ *Department of Physics and Astronomy, Uppsala University, Uppsala, Sweden, associated to ⁵⁹*
- ⁸⁷ *University of Michigan, Ann Arbor, Michigan, United States, associated to ⁶⁸*

^a *Universidade Federal do Triângulo Mineiro (UFMT), Uberaba-MG, Brazil*

^b *Hangzhou Institute for Advanced Study, UCAS, Hangzhou, China*

^c *Excellence Cluster ORIGINS, Munich, Germany*

^d *Università di Bari, Bari, Italy*

^e *Università di Bologna, Bologna, Italy*

^f *Università di Cagliari, Cagliari, Italy*

^g *Università di Ferrara, Ferrara, Italy*

^h *Università di Firenze, Firenze, Italy*

ⁱ *Università di Genova, Genova, Italy*

^j *Università degli Studi di Milano, Milano, Italy*

^k *Università di Milano Bicocca, Milano, Italy*

^l *Università di Modena e Reggio Emilia, Modena, Italy*

^m *Università di Padova, Padova, Italy*

ⁿ *Scuola Normale Superiore, Pisa, Italy*

^o *Università di Pisa, Pisa, Italy*

^p *Università della Basilicata, Potenza, Italy*

^q *Università di Roma Tor Vergata, Roma, Italy*

^r *Università di Urbino, Urbino, Italy*

^s *MSU - Iligan Institute of Technology (MSU-IIT), Iligan, Philippines*

^t *P.N. Lebedev Physical Institute, Russian Academy of Science (LPI RAS), Moscow, Russia*

^u *Novosibirsk State University, Novosibirsk, Russia*

[†] *Deceased*