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# WHEN DO TWO RATIONAL FUNCTIONS HAVE LOCALLY BIHOLOMORPHIC JULIA SETS? 

ROMAIN DUJARDIN, CHARLES FAVRE, AND THOMAS GAUTHIER


#### Abstract

In this note we address the following question, whose interest was recently renewed by problems arising in arithmetic dynamics: under which conditions does there exist a local biholomorphism between the Julia sets of two given one-dimensional rational maps? In particular we find criteria ensuring that such a local isomorphism is induced by an algebraic correspondence. This extends and unifies classical results due to Baker, Beardon, Eremenko, Levin, Przytycki and others. The proof involves entire curves and positive currents.


## 1. Introduction

The problem of determining when two rational maps have the same Julia set has been considered by many authors in the holomorphic dynamics literature [1, 3, 9, 14, 21, 22, 37, in relation with some classical functional equations. In certain situations (e.g., if the Julia set is the whole sphere), it is preferable to ask when two rational maps have the same measure of maximal entropy. The conclusion is that these rational maps satisfy an algebraic relation whose analysis is quite delicate (this was explored in [30, 41). Such rigidity issues have recently played an important role in arithmetic dynamics (see e.g., [2, 17, 19], and also Remark 1.1 below).

In this article we consider the following problem: is any local biholomorphism preserving the Julia set (or the measure of maximal entropy) induced by an algebraic correspondence? The case of polynomials of the same degree with connected and locally connected Julia sets was recently addressed by Luo [24]. Here we deal with general rational maps, possibly of different degrees, and obtain a rather satisfactory answer when the maximal entropy measure is (quasi-)preserved (Theorem A ). We are also able to completely solve the problem for polynomials satisfying some mild expansion properties on their Julia sets (Theorem B, see also Remark 4.3 for a discussion of the relationship between our results and those of [24]). Related results were obtained in [6, 20, 23], under the stronger assumption of the existence of a partial analytic conjugacy.

To be more specific, for a complex rational map $f$ of degree $d \geqslant 2$, we denote by $J_{f} \subset \mathbb{P}^{1}$ its Julia set and by $\mu_{f}$ its measure of maximal entropy. It is the unique atomless probability measure such that $f^{*} \mu_{f}=d \mu_{f}$. Following the terminology of [17], we say that $f$ is integrable if it is either a Chebychev, monomial, or Lattès map. Likewise, we say that $J_{f}$ is smooth if it is equal to $\mathbb{P}^{1}$, a circle, or a segment. Any integrable rational map has a smooth Julia set, but there are many more examples (see e.g. [15]).

Given two positive measures $\mu_{1}$ and $\mu_{2}$, we write $\mu_{1}=\mu_{2}$ if $c^{-1} \mu_{2} \leqslant \mu_{1} \leqslant c \mu_{2}$ for some positive constant $c$. In other words, $\mu_{1}$ is absolutely continuous with respect to $\mu_{2}$
and $c^{-1} \leqslant \frac{d \mu_{1}}{d \mu_{2}} \leqslant c$. We use the more precise notation $\mu_{1} \asymp_{c} \mu_{2}$ if there is a need to specify the constant $c$. We also write $\mu_{1} \propto \mu_{2}$ if $\mu_{1}$ is proportional to $\mu_{2}$, that is, if there exists $\alpha>0$ such that $\mu_{1}=\alpha \mu_{2}$.

Here is our first main result:

Theorem A. Let $f_{1}$ and $f_{2}$ be non-integrable rational maps of degree larger than 1 , and $U$ be any open subset of $\mathbb{P}^{1}$ intersecting $J_{f_{1}}$. Let $\sigma: U \rightarrow \mathbb{P}^{1}$ be any non-constant holomorphic map satisfying $\sigma^{*} \mu_{f_{2}}=\mu_{f_{1}}$ on $U$. When $J_{f_{1}}$ is smooth we further require that $\sigma^{*} \mu_{f_{2}} \propto \mu_{f_{1}}$.

Suppose in addition that there exists a repelling periodic point $p_{1}$ for $f_{1}$ such that $\sigma\left(p_{1}\right)$ is preperiodic under $f_{2}$.

Then there exist positive integers $a, b \in \mathbb{N}^{*}$ and an irreducible algebraic curve $Z \subset$ $\mathbb{P}^{1} \times \mathbb{P}^{1}$ which is preperiodic under $\left(f_{1}^{a}, f_{2}^{b}\right)$ and contains the graph of $\sigma$.

In particular we have that $d_{1}^{a}=d_{2}^{b}$, and $\sigma^{*} \mu_{f_{2}} \propto \mu_{f_{1}}$ in all cases.

In plain words, the local measure class preserving morphism $\sigma$ between $J_{f_{1}}$ and $J_{f_{2}}$ is induced by an algebraic correspondence between $f_{1}$ and $f_{2}$. It is easy to see that such a result cannot be true in the integrable case (see Remark 2.13).

Preperiodic algebraic curves under $\left(f_{1}, f_{2}\right)$ in $\mathbb{P}^{1} \times \mathbb{P}^{1}$ were classified by Pakovich [29]: the upshot is that for any such curve there exists rational functions $X_{1}, X_{2}, Y_{1}, Y_{2}$ such that the curve is a component of $Y_{1}(x)-Y_{2}(y)=0, Y_{1} \circ X_{1}=Y_{2} \circ X_{2}, X_{1} \circ Y_{1}=f_{1}^{n}$ and $X_{2} \circ Y_{2}=f_{2}^{n}$. If $f=f_{1}=f_{2}$ is a generic map of degree $d \geqslant 3$, this curve must be a component of $\left\{f^{k}(x)=f^{\ell}(y)\right\}$ (see [31, 41]). We conclude that in this case any algebraic correspondence is induced by a branch of $f^{-\ell} \circ f^{k}$. In the polynomial case, more precise results were obtained in 26].

The proof of Theorem A is given in Section 2. A natural strategy to establish such a result is to use the expansion induced by the repelling point $p_{1}$ to propagate the local morphism $\sigma$ to the whole of $\mathbb{P}^{1}$. The difficulty is that the limiting objects will be transcendental and highly multivalued. Our approach is geometric and based on entire curves and positive currents: notable ingredients include the Ahlfors Five Islands Theorem (in its strong, quantitative form) and Siu's decomposition theorem for positive closed currents.

Another key tool in the proof is a deep rigidity theorem due to Levin [21], which asserts that, for a given rational map $f$, there are in a sense only finitely many nontrivial local symmetries of $J_{f}$ defined in a fixed open set $U$ intersecting $J_{f}$ (see $\$ 2.2$ for details). A posteriori, Theorem A may be viewed as a refinement of Levin's theorem: these symmetries are local branches of some global algebraic symmetry of $f$.

Note that the notion of local symmetry in Levin's theorem is less restrictive than ours: for a non-smooth Julia set it is just a local holomorphic map such that $\sigma^{-1}\left(J_{f}\right) \cap$ $U=J_{f} \cap U$. In view of this, it is natural to expect that, for non-smooth Julia sets, the assumption that $\sigma$ preserves the measure class of the maximal entropy measure is superfluous in Theorem A (see Remark 3.4 for a related discussion). More ambitiously, we may ask the following:

Problem. Let $f_{1}$ and $f_{2}$ be non-integrable rational maps, and $\sigma$ be a holomorphic map defined in some open set $U$ intersecting $J_{f_{1}}$, such that the equality $\sigma^{-1}\left(J_{f_{2}}\right) \cap U=J_{f_{1}} \cap U$ holds if $J_{f_{1}}$ is not smooth, and $\sigma^{*} \mu_{f_{2}}=\mu_{f_{1}}$ when $J_{f_{1}}$ is smooth.

Then is $\sigma$ induced by an algebraic correspondence between $f_{1}$ and $f_{2}$ ?
In Sections 3, 4, and 5 we take some further steps towards the resolution of this problem. We first prove that the assumption that $\sigma$ maps a repelling point to a preperiodic point can be dropped under suitable expansion properties for $f_{1}$ and $f_{2}$. Indeed, in this case the geometry of the Julia set (resp. of $\mu_{f}$ ) alone is enough to detect preperiodic points (see Section 3, and in particular Corollary 3.2, for details).

For polynomials, by using the fact that the maximal entropy measure coincides with the harmonic measure of the Julia set (viewed from infinity), we are able to show that, under reasonable assumptions, the measure class preservation $\sigma^{*} \mu_{f_{2}}=\mu_{f_{1}}$ is automatically implied by the geometric condition $\sigma^{-1}\left(J_{f_{2}}\right) \cap U=J_{f_{1}} \cap U$. Some non-trivial potentialtheoretic arguments are developed in Section 4 to deal with the delicate interplay between the local and global properties of the harmonic measure.

In Section 5, we deal with one specific issue concerning polynomial Julia sets that are Jordan curves: how to make locally the distinction between the inside and the outside of $J_{f}$ ? As a matter of fact, in Proposition 5.1 we prove that if $f_{1}$ and $f_{2}$ are polynomials whose Julia sets are non-smooth Jordan curves, there does not exist a local biholomorphism mapping $J_{1}$ to $J_{2}$ and mapping the bounded component of $\mathbb{C} \backslash J_{1}$ to the bounded component of $\mathbb{C} \backslash J_{2}$. This is particularly delicate for quasicircles, which have no preferred side from the conformal point of view. J1 to J2 and mapping the bounded component of C- J1 to the unbounded component of C- J2

Combining this to the results of Section 3 yields the following streamlined version of Theorem A, which solves the above problem for polynomials satisfying the so-called Topological Collet-Eckmann (TCE) condition (see Section 3 for this notion; Theorem B is proven in Section (4).

Theorem B. Let $f_{1}$ and $f_{2}$ be polynomials, such that either of $f_{1}$ or $f_{2}$ is non-integrable and satisfies the Topological Collet-Eckmann property. Let $\sigma$ be a holomorphic map defined in some open set $U$ intersecting $J_{f_{1}}$, such that the equality $\sigma^{-1}\left(J_{f_{2}}\right) \cap U=J_{f_{1}} \cap U$ holds. Then $\sigma$ is induced by an algebraic correspondence between $f_{1}$ and $f_{2}$.
Remark 1.1. Our investigations were partially motivated by a question that arose in the work of the first two authors on the dynamical Manin-Mumford problem for plane polynomial automorphisms [12]. The context is as follows:

Suppose $f$ is a complex plane polynomial automorphisms of positive entropy, and let $p$ be any hyperbolic fixed point. Suppose moreover that there exists a local biholomorphism $\sigma$ from the stable manifold of $p$ to its unstable manifold that maps the backward invariant current to the forward invariant current (these objects are higher dimensional analogs of the maximal entropy measure). Showing that $\sigma$ extends as a global algebraic involution of $\mathbb{C}^{2}$ would imply [12, Conjecture 1]; see Remark 4.4 of op.cit.

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## 2. Proof of Theorem A

2.1. Notation and conventions. Any positive measure $\mu$ on $\mathbb{P}^{1}$ can be locally defined by $\mu=d d^{c} u$ for some subharmonic function $u$. We say that $\mu$ has continuous potentials when $u$ is continuous. Let $f: U \rightarrow \mathbb{P}^{1}$ be any holomorphic map defined on a connected open subset of the Riemann sphere. For any positive measure $\mu=d d^{c} u$ as before, we locally define $f^{*} \mu$ by $f^{*} \mu=d d^{c}(u \circ f)$. Alternatively, if $\mu$ gives no mass to points, we may set

$$
\left.f^{*} \mu\right|_{A}=\left(\left.f\right|_{f(A)} ^{-1}\right)_{*} \mu
$$

on any Borel set $A$ on which $f$ is injective.
When two rational maps $f_{1}$ and $f_{2}$ are given, to ease notation we often write $\mu_{i}$ for $\mu_{f_{i}}$, and likewise $J_{i}$ for the Julia sets, etc.

We use the notation $D(z, r)$ for the open disk of center $z \in \mathbb{C}$ and radius $r$ in $\mathbb{C}$, write $\mathbb{D}=D(0,1)$ and identify $\partial \mathbb{D}$ with $\mathbb{S}^{1}=\mathbb{R} / \mathbb{Z}$. Likewise, $B(z, r)$ refers to the spherical disk of center $z$ and radius $r$ in $\mathbb{P}^{1}$. We often denote by $C$ a "constant" which may change from line to line, independently of some asymptotic quantity which should be clear from the context, and write $a \lesssim b$ for $a \leqslant C b$ and $a=b$ for $C^{-1} b \leqslant a \leqslant C b$.
2.2. Levin's theorem. By definition the Julia set of a rational map $f$ is said to be smooth if it contains an open set or if on some open subset it coincides with a smooth arc. In this case it was proved by Fatou that $J_{f}$ is respectively equal to $\mathbb{P}^{1}$, or it is a circle or a segment (see [16, §56 p. 250]).

Theorem 2.1 (Levin [21]). Let $f$ be a rational map of degree greater than 1. Suppose there exists a connected open set $U \subset \mathbb{P}^{1}$ intersecting $J_{f}$ and an infinite family of holomorphic maps $\sigma_{n}: U \rightarrow \mathbb{P}^{1}$ such that $\sigma_{n}^{-1}\left(J_{f}\right) \cap U=J_{f} \cap U$. If in addition $J_{f}$ is smooth we require that $\sigma_{n}^{*} \mu_{f} \propto \mu_{f}$.

If the family $\left(\sigma_{n}\right)$ is normal and all its limit functions are non-constant, then $f$ is integrable.

In the following, we will refer to a local map $\sigma$ satisfying these assumptions simply as a local symmetry of $J_{f}$. Let us also pinpoint an intermediate step in the proof of Theorem 2.1 which will also be useful.

Lemma 2.2 (see [21, Proposition 1]). Let $p$ be a repelling fixed point of the rational map $f$. Assume that $\sigma: U \rightarrow \mathbb{P}^{1}$ is a holomorphic map fixing $p$ such that $\sigma^{\prime}(p) \neq 0$ and $\sigma^{-1}\left(J_{f}\right) \cap U=J_{f} \cap U$; if in addition $J_{f}$ is smooth we further require that $\sigma^{*} \mu_{f} \propto \mu_{f}$. Then $\sigma$ and $f$ commute.
2.3. A normal families lemma. The normality assumption in Levin's theorem will be deduced from a uniform bound on the Radon-Nikodym derivative $\frac{d \sigma_{n}^{*} \mu_{f}}{d \mu_{f}}$, thanks to the following normal families criterion.

Lemma 2.3. Let $U \subset \mathbb{P}^{1}$ be a connected open set, $\nu_{1}$ a non-zero positive measure on $U$ and $\nu_{2}$ a positive measure on $\mathbb{P}^{1}$ with continuous local potentials. If $\sigma_{n}$ is a sequence of holomorphic mappings $U \rightarrow \mathbb{P}^{1}$ such that $\left(\sigma_{n}\right)^{*} \nu_{2} \asymp_{c} \nu_{1}$ for some uniform $c>0$, then $\left(\sigma_{n}\right)$ is a normal family and all its limiting maps are non-constant.

From Lemma 2.3 and Levin's theorem we get:
Corollary 2.4. Let $f$ be a non-integrable rational map of degree greater than 1 , and $\left(\sigma_{n}\right)$ be a sequence of local symmetries of $J_{f}$ as in Theorem 2.1. If in addition $\left(\sigma_{n}\right)^{*} \mu 七_{c} \mu$ for some uniform $c>0$, then the family $\left(\sigma_{n}\right)$ is finite.

Proof of Lemma 2.3. We may assume that $U$ is a disk. Note that our assumption implies that $\nu_{2}$ gives no mass to points, so neither does $\nu_{1}$ on $U$.

For the first assertion, assume by contradiction that $\left(\sigma_{n}\right)$ is not normal in $U$. Then by the Zalcman reparameterization lemma there exists a sequence $\left(a_{n}\right)$ converging to some $a^{*} \in U$, a sequence $r_{n} \rightarrow 0$ and an extraction $n_{j}$ such that the sequence or meromorphic functions $\left(\zeta \mapsto \sigma_{n_{j}}\left(a_{n_{j}}+r_{n_{j}} \zeta\right)\right)$ converges uniformly on compact subsets to a non-constant entire mapping $\sigma_{\infty}: \mathbb{C} \rightarrow \mathbb{P}^{1}$. Since $\nu_{2}$ gives no mass to points there is a regular value of $\sigma_{\infty}$ in $\operatorname{Supp}\left(\nu_{2}\right)$. In particular there is a disk $D^{\prime}$ on which $\sigma_{\infty}$ is univalent, together with a smaller disk $D \Subset D^{\prime}$ such that $\nu_{2}\left(\sigma_{\infty}(D)\right)>0$. Now $\sigma_{n_{j}}$ is univalent on $a_{n_{j}}+r_{n_{j}} D$ for large $j$, and $\sigma_{n_{j}}\left(a_{n_{j}}+r_{n_{j}} D\right)$ converges to $\sigma_{\infty}(D)$. Therefore if $D$ was further chosen so that $\nu_{2}\left(\partial\left(\sigma_{\infty}(D)\right)\right)=0$ we infer that

$$
\nu_{1}\left(\left\{a^{*}\right\}\right)=\lim _{n} \nu_{1}\left(a_{n_{j}}+r_{n_{j}} D\right) \geqslant c \nu_{2}\left(\sigma_{\infty}(D)\right)>0
$$

which is the desired contradiction.
For the second assertion we again argue by contradiction and assume that some subsequence $\left(\sigma_{n_{j}}\right)$ converges to a constant $a$ on $U$. Let $\varphi$ be a non-negative test function in $U$ such that $\int \varphi d \nu_{1}>0$. Let $g_{1}$ be a subharmonic potential for $\nu_{1}$ in $U$ and $g_{2}$ be a subharmonic potential for $\nu_{2}$ defined in a neighborhood of $a$. By assumption, we have $d d^{c}\left(g_{2} \circ \sigma_{n}\right) \geqslant c d d^{c} g_{1}$. Substracting a constant we may assume that $g_{2}(a)=0$. Then we have

$$
0<c \int_{U} \varphi d d^{c} g_{1}=\int_{U} \varphi d d^{c}\left(g_{2} \circ \sigma_{n}\right)=\int_{U}\left(g_{2} \circ \sigma_{n}\right) d d^{c} \varphi
$$

and

$$
\left|\int_{U}\left(g_{2} \circ \sigma_{n}\right) d d^{c} \varphi\right| \leqslant\|\varphi\|_{C^{2}(U)}\left\|g_{2} \circ \sigma_{n}\right\|_{L^{\infty}(\operatorname{Supp}(\varphi))}
$$

which tends to zero since $\sigma_{n}(\operatorname{Supp}(\varphi))$ converges to $\{a\}$ and $g_{2}$ is continuous. This contradiction finishes the proof.

Remark 2.5. The continuity of the potential of $g_{2}$ is essential in the second part of the proof, in particular assuming that $\nu_{2}$ gives no mass to points is not enough to conclude. Indeed the measure $\nu=d d^{c}(-\log |\log | z| |)$ gives no mass to points and satisfies $\left(\sigma_{n}\right)^{*} \nu=$ $\nu$ for $\sigma_{n}(z)=z^{n}$, while $z^{n}$ converges uniformly to 0 in a neighborhood of the origin.
2.4. Algebraization. The core of the proof of Theorem A is the following algebraization result, which will later be applied to (generalized) Poincaré-Koenigs linearization mappings. At this stage we do not claim any invariance for the implied algebraic curve.

Proposition 2.6. Let $f_{1}$ and $f_{2}$ be two non-integrable rational maps on $\mathbb{P}^{1}$. Assume that $\psi_{1}$ and $\psi_{2}$ are entire maps $\mathbb{C} \rightarrow \mathbb{P}^{1}$ such that $\left(\psi_{2}\right)^{*} \mu_{f_{2}}=\left(\psi_{1}\right)^{*} \mu_{f_{1}}$; if either $J_{f_{1}}$ or $J_{f_{2}}$ is smooth then we further require that $\left(\psi_{2}\right)^{*} \mu_{f_{2}} \propto\left(\psi_{1}\right)^{*} \mu_{f_{1}}$. Define the entire map $\Psi: \mathbb{C} \rightarrow \mathbb{P}^{1} \times \mathbb{P}^{1}$ by $\Psi=\left(\psi_{1}, \psi_{2}\right)$. Then $\overline{\Psi(\mathbb{C})}$ is an irreducible algebraic curve which is neither a vertical nor a horizontal line.

Notice that under our assumptions, $J_{f_{1}}$ is smooth if and only if $J_{f_{2}}$ is also smooth.
Proof. As a preliminary step, let us observe that if $\psi_{1}$ and $\psi_{2}$ are rational, then $\Psi(\mathbb{C})$ is an algebraic curve since the transcendence degree of $\mathbb{C}(T)$ over $\mathbb{C}$ is 1 . Another argument goes by using Remmert's Proper Mapping Theorem and the GAGA principle. So without loss of generality we may assume that $\Psi$ is transcendental.
Step 1: construction of inverse branches and geometry of the Ahlfors currents.
Let $\omega_{\mathbb{P}^{1}}$ be a Fubini-Study form on $\mathbb{P}^{1}$, normalized by $\int_{\mathbb{P}^{1}} \omega_{\mathbb{P}^{1}}=1$ and set $\omega=\pi_{1}^{*} \omega_{\mathbb{P}^{1}}+$ $\pi_{2}^{*} \omega_{\mathbb{P}^{1}}$. For any $R>0$, set

$$
\operatorname{Area}(\Psi(D(0, R))):=\int_{D(0, R)} \Psi^{*} \omega, \text { and Length }(\partial \Psi(D(0, R))):=\int_{\partial \mathbb{D}}\left|\Psi^{\prime}\left(R e^{i \theta}\right)\right|_{\omega} R d \theta
$$

Since $\Psi$ is transcendental, Area $(\Psi(D(0, R))) \rightarrow \infty$ when $R \rightarrow \infty$. By the Ahlfors isoperimetric inequality (see [40, §VI.5] or [5]) there exists a sequence $R_{j} \rightarrow \infty$ such that

$$
\text { Length }\left(\partial \Psi\left(D\left(0, R_{j}\right)\right)\right)=\mathrm{o}\left(\operatorname{Area}\left(\Psi\left(D\left(0, R_{j}\right)\right)\right)\right)
$$

Any cluster value of the sequence of positive currents

$$
T_{j}:=\frac{1}{\operatorname{Area}\left(\Psi\left(D\left(0, R_{j}\right)\right)\right)}\left[\Psi\left(D\left(0, R_{j}\right)\right)\right]
$$

is by definition an Ahlfors current associated to $\Psi$. Fix such an Ahlfors current $T$. Then $T$ is a positive closed $(1,1)$ current in $\mathbb{P}^{1} \times \mathbb{P}^{1}$ satisfying $\int T \wedge \omega=1$ so there exists $i \in\{1,2\}$ such that $\left(\pi_{i}\right)_{*} T \neq 0$, or equivalently $\left\langle[T],\left[\pi_{i}^{*} \omega_{\mathbb{P}^{1}}\right]\right\rangle \neq 0$, where $[\cdot]$ denotes the class in $H^{1,1}\left(\mathbb{P}^{1} \times \mathbb{P}^{1}\right)$ and $\langle\cdot, \cdot\rangle$ is the intersection pairing. Without loss of generality we may assume $i=1$.

We now apply Ahlfors' theory of covering surfaces, in the spirit of [4, §7], with an additional twist inspired from [11]. Fix any integer $q \geqslant 5$ and consider $q$ disks $D_{i}$ with disjoint closures, intersecting $J_{1}$. For every $1 \leqslant i \leqslant q$, let $N_{j}\left(D_{i}\right)$ be the number of univalent inverse branches ("good islands") of $\psi_{1}$ over $D_{i}$ contained in $D\left(0, R_{j}\right)$. We label the corresponding components of $\psi_{1}^{-1}\left(D_{i}\right)$ as $\left(\Omega_{i, n}\right)_{n \geqslant 1}$ in such a way that for $1 \leqslant n \leqslant N_{j}\left(D_{i}\right), \Omega_{i, n} \subset D\left(0, R_{j}\right)$. Note that at most one of the $\Omega_{i, n}$ contains the origin so we may assume that $0 \notin \Omega_{i, n}$. Then by Ahlfors' theorem,

$$
\sum_{i=1}^{q} N_{j}\left(D_{i}\right) \geqslant(q-4) \operatorname{Area}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)-h \operatorname{Length}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)
$$

where the area and length are computed with respect to $\omega_{\mathbb{P}^{1}}$, and $h$ is a geometric constant depending only on the disks $D_{i}$ (see [40, Theorem VI.4]).

Since

$$
0<\int\left(\pi_{1}\right)_{*} T \wedge \omega_{\mathbb{P}^{1}}=\lim _{j \rightarrow \infty} \frac{\operatorname{Area}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)}{\operatorname{Area}\left(\Psi\left(D\left(0, R_{j}\right)\right)\right)}
$$

there exists a constant $C_{1}$ such that for every $j$,

$$
\operatorname{Area}\left(\Psi\left(D\left(0, R_{j}\right)\right)\right) \leqslant C_{1} \operatorname{Area}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)
$$

in particular $\psi_{1}$ is transcendental. The number of good islands contained in $D\left(0, R_{j}\right)$ whose volume (relative to $\omega$ ) is at least $1 / 2$ is bounded from above by 2 Area $\left(\Psi\left(D\left(0, R_{j}\right)\right)\right.$, which is itself bounded by $2 C_{1} \operatorname{Area}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)$. Let us discard these components and denote by $N_{j}^{\prime}\left(D_{i}\right)$ the number of remaining ones. Since these components have volume bounded by $1 / 2$, by Bishop's compactness theorem (see, e.g., [11, Lemma 3.5]) they form a normal family. If $q$ was chosen so that $q>4+2 C_{1}$ we infer that

$$
\sum_{i=1}^{q} N_{j}^{\prime}\left(D_{i}\right) \geqslant\left(q-4-2 C_{1}\right) \operatorname{Area}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)-h \operatorname{Length}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)
$$

Extract a further subsequence of $\left(R_{j}\right)$ (still denoted by $\left(R_{j}\right)$ ) so that for a fixed $i \in$ $\{1, \ldots, q\}$

$$
\begin{equation*}
N_{j}^{\prime}\left(D_{i}\right) \geqslant \frac{q-4-2 C_{1}}{q} \operatorname{Area}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)-\frac{h}{q} \operatorname{Length}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right) \tag{1}
\end{equation*}
$$

and put $D=D_{i}$ and $\Omega_{n}=\Omega_{i, n}$. Let

$$
S_{j}=\frac{1}{\operatorname{Area}\left(\Psi\left(D\left(0, R_{j}\right)\right)\right)} \sum_{n=0}^{N_{j}^{\prime}(D)}\left[\Psi\left(\Omega_{n}\right)\right]
$$

which is a sum of integration currents of graphs over $D$. By (1), we have that $S_{j} \leqslant T_{j}$ and we may estimate the mass $\mathbf{M}\left(S_{j}\right):=\int S_{j} \wedge \omega$ as follows:
$\mathbf{M}\left(S_{j}\right) \geqslant \int S_{j} \wedge \pi_{1}^{*} \omega \geqslant \frac{q-4-2 C_{1}}{q} \cdot \frac{\operatorname{Area}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)}{\operatorname{Area}\left(\Psi\left(D\left(0, R_{j}\right)\right)\right)}-\frac{h}{q} \cdot \frac{\operatorname{Length}_{\mathbb{P}^{1}}\left(\psi_{1}\left(D\left(0, R_{j}\right)\right)\right)}{\operatorname{Area}\left(\Psi\left(D\left(0, R_{j}\right)\right)\right)}$
hence $\liminf _{j} \mathbf{M}\left(S_{j}\right) \geqslant \frac{q-4-2 C_{1}}{q} \mathbf{M}\left(\left(\pi_{1}\right)_{*} T\right)>0$ and any cluster value $S$ of the sequence $\left(S_{j}\right)$ satisfies $0<S \leqslant T$.

Step 2: using the local symmetries to conclude.
To simplify notation, write $\mu_{1}=\mu_{f_{1}}$ and $\mu_{2}=\mu_{f_{2}}$. For every $n$, define

$$
\psi_{1, n}^{-1}:=\left(\psi_{1} \mid \Omega_{n}\right)^{-1}: D \xrightarrow{\sim} \Omega_{n}
$$

and let $\sigma_{n}:=\psi_{2} \circ \psi_{1, n}^{-1}$. By construction $\sigma_{n}$ is defined in $D$ with values in $\mathbb{P}^{1}$. Writing $\psi_{2}^{*} \mu_{2}=h \psi_{1}^{*} \mu_{1}$, with $c^{-1} \leqslant h \leqslant c$ for some $c>0$, we infer that

$$
\begin{align*}
\sigma_{n}^{*} \mu_{2} & =\left(\psi_{1, n}\right)_{*}\left(\psi_{2}^{*} \mu_{2}\right)=\left(\psi_{1, n}\right)_{*}\left(h \psi_{1}^{*} \mu_{1}\right)  \tag{2}\\
& =\left(h \circ \psi_{1, n}^{-1}\right)\left(\psi_{1, n}\right)_{*} \psi_{1}^{*} \mu_{1}=\left(h \circ \psi_{1, n}^{-1}\right) \mu_{1} \asymp_{c} \mu_{1}
\end{align*}
$$

so by Lemma 2.3, $\left(\sigma_{n}\right)$ is a normal family and its limiting maps are non-constant.

Remark 2.7. Note that the normality of the family $\left(\sigma_{n}\right)$ was already obtained in Step 1, so the quasi-preservation of the measure is only used to guarantee that its normal limit are non-constant.

Now observe that the maps $\sigma_{n}$ give rise to local symmetries of $J_{2}$ : indeed we can pick a subdisk $D^{\prime}$ intersecting $J_{1}$ on which $\sigma_{1}$ is univalent, and define a sequence of local symmetries of $J_{2}$ by putting $\tau_{n}=\sigma_{n} \circ\left(\left.\sigma_{1}\right|_{D^{\prime}}\right)^{-1}$. These are holomorphic map from $\sigma_{1}\left(D^{\prime}\right)$ to $\mathbb{P}^{1}$ satisfying the relation $\tau_{n}^{*} \mu_{2} 二_{c^{2}} \mu_{2}$, where $c$ is as in 2). If in addition $J_{f_{1}}$ and $J_{f_{2}}$ are smooth, arguing as in (2) we further deduce that $\tau_{n}^{*} \mu_{2} \propto \mu_{2}$. Thus it follows from Corollary 2.4 that the family $\left(\tau_{n}\right)$ is finite, hence so does the family $\left(\sigma_{n}\right)$. From this we infer that the graphs $\Psi\left(\Omega_{n}\right)$ are contained in finitely many graphs over $D$, therefore $S_{j}$ is an integration current over a fixed finite union of graphs $\left(\Delta_{\ell}\right)$ over $D$, independent of $j$, namely $S_{j}=\sum_{\ell} s_{\ell, j}\left[\Delta_{\ell}\right]$. Extracting a converging subsequence, we get a current $S=\sum s_{\ell}\left[\Delta_{\ell}\right]$ supported on the same family of graphs and from Step 1 we know that $0<S \leqslant T$. Note that none of these graphs is horizontal because $\psi_{2}^{*} \mu_{2} \approx \psi_{1}^{*} \mu_{1}$.

With the above notation, fix $\ell$ such that $s_{\ell}>0$. Then $T \geqslant s_{\ell}\left[\Delta_{\ell}\right]$. By Siu's decomposition theorem (see [7, (2.18)]) there exists an analytic, hence algebraic, subvariety $\Gamma$ of $\mathbb{P}^{1} \times \mathbb{P}^{1}$, extending $\Delta_{\ell}$, such that $T \geqslant s_{\ell}[\Gamma]$. Since $\Delta_{\ell} \subset \Psi(\mathbb{C})$ by construction and $\Delta_{\ell} \subset \Gamma$, by analytic continuation $\Psi(\mathbb{C})$ is contained in $\Gamma$. Therefore we conclude that $\overline{\Psi(\mathbb{C})}$ is an algebraic curve, which is obviously irreducible, and which cannot be neither a vertical line because it contains a graph over the first coordinate, nor a horizontal line because this graph was shown to be non-horizontal. The proof is complete.

Remark 2.8. Note that if $\psi_{1}$ is transcendental, then so does $\psi_{2}$. Indeed by Proposition 2.6 there exists a polynomial $P \in \mathbb{C}[x, y]$ such that $P\left(\psi_{1}, \psi_{2}\right)=0$. If $\psi_{2}$ is algebraic, it follows that $\psi_{1}^{-1}\left(z_{1}\right)$ is finite for any $z_{1} \in \mathbb{C}$ which is contradictory.
2.5. Local isomorphisms and Poincaré-Koenigs functions. Let $f$ be a rational map of degree $d \geqslant 2$, and $p$ be a repelling fixed point. Denote by $\lambda=f^{\prime}(p)$ its multiplier. Then $f$ is linearizable in a neighborhood of $p$, consequently there exists a unique holomorphic map $\psi_{(f, p)}: \mathbb{C} \rightarrow \mathbb{P}^{1}$ such that $\psi_{(f, p)}(0)=p, \psi_{(f, p)}^{\prime}(0)=1$ and for every $\zeta \in \mathbb{C}$,

$$
\begin{equation*}
f \circ \psi_{(f, p)}(\zeta)=\psi_{(f, p)}(\lambda \zeta) \tag{3}
\end{equation*}
$$

This map is called the Poincaré-Koenigs linearizing map of $f$.
Lemma 2.9. Let $\chi:(\mathbb{C}, 0) \rightarrow\left(\mathbb{P}^{1}, p\right)$ be a germ of non-constant holomorphic map satisfying the functional equation

$$
\begin{equation*}
f \circ \chi(\zeta)=\chi(\kappa \zeta) \text { for all } \zeta \in \mathbb{C} \text { and some } \kappa \in \mathbb{C} . \tag{4}
\end{equation*}
$$

If we have $\chi(\zeta)=\beta \zeta^{l}+O\left(\zeta^{l+1}\right)$ with $l \geqslant 1$ and $\beta \neq 0$, then $\kappa^{\ell}=\lambda$ and $\chi(\zeta)=\psi_{(f, p)}\left(\beta \zeta^{\ell}\right)$.
Any function $\chi$ satisfying (4) will be referred to as a generalized Poincaré-Koenigs map.

Proof. The expansion of $f \circ \chi(\zeta)$ at the origin together with (4) force $\lambda=\kappa^{l}$. Locally at $0, \psi_{(f, p)}$ is invertible, so that we may consider the holomorphic germ $\tilde{\chi}:=\psi_{(f, p)}^{-1} \circ \chi$.

Observe that

$$
\tilde{\chi}(\kappa \zeta)=\psi_{(f, p)}^{-1} \circ f \circ \chi(\zeta)=\lambda \tilde{\chi}(\zeta)
$$

Expanding $\tilde{\chi}$ in power series at 0 yields $\tilde{\chi}(\zeta)=\beta \zeta^{\ell}$. The proof is complete.
Recall that for two rational maps $f_{i}, i=1,2$ we write $\mu_{i}=\mu_{f_{i}}, J_{i}=J_{f_{i}}$, etc.
Proposition 2.10. Let $f_{1}$ and $f_{2}$ be two non-integrable rational maps of respective degrees $d_{1}, d_{2} \geqslant 2$, and $U$ be any connected open set intersecting $J_{1}$. Suppose $\sigma: U \rightarrow \mathbb{P}^{1}$ is a non-constant holomorphic map sending a repelling fixed point $p_{1}$ for $f_{1}$ to a fixed point $p_{2}$ for $f_{2}$. Let $\lambda_{1}$ and $\lambda_{2}$ be the respective multipliers of $p_{1}$ and $p_{2}$, and set $\ell=\operatorname{deg}_{p_{1}}(\sigma) \geqslant 1$.

Suppose that:
(1) either $J_{2}$ is not smooth;
(2) or $J_{2}$ is smooth and $\sigma^{*}\left(\mu_{2}\right) \propto \mu_{1}$.

Then the point $p_{2}$ is repelling, and there exist $a, b \in \mathbb{N}^{*}$ such that $\lambda_{1}^{a \ell}=\lambda_{2}^{b}$ and $f_{2}^{b} \circ \sigma=$ $\sigma \circ f_{1}^{a}$. Moreover, for $\chi_{1}=\psi_{\left(f_{1}, p_{1}\right)}, \chi_{2}:=\sigma \circ \chi_{1}$ extends to a generalized Poincaré-Koenigs map for $f_{2}$ satisfying $\chi_{1}^{-1}\left(J_{1}\right)=\chi_{2}^{-1}\left(J_{2}\right)$.

If in Case (1) we further assume:
(1') $J_{2}$ is not smooth and $\sigma^{*} \mu_{2}=\mu_{1}$
then we have the identities $d_{1}^{a}=d_{2}^{b}$ and $\chi_{1}^{*} \mu_{1}=\chi_{2}^{*} \mu_{2}$ (resp. $\chi_{1}^{*} \mu_{1} \propto \chi_{2}^{*} \mu_{2}$ in Case (2)).
Remark 2.11. If $p_{1}$ and $p_{2}$ are periodic of respective periods $m_{1}$ and $m_{2}$, applying this result to $f^{m_{1}}$ and $f^{m_{2}}$ we get a similar conclusion, where the relations become $\lambda_{1}^{a \ell}=\lambda_{2}^{b}$, $f_{2}^{m_{2} b} \circ \sigma=\sigma \circ f_{1}^{m_{1} a}$, and $d_{1}^{m_{1} a}=d_{2}^{m_{2} b}$.

Proof. Choose local coordinates such that $p_{1}=p_{2}=0$ and $\sigma(z)=z^{\ell}$ for some $\ell \in \mathbb{N}^{*}$. Fix any $\ell$-th root $\kappa_{2}$ of $\lambda_{2}$. Then we can write

$$
f_{2} \circ \sigma(z)=\lambda_{2} z^{\ell}+\text { h.o.t. }=\left(\kappa_{2} z+\sum_{j \geqslant 2} a_{j} z^{j}\right)^{\ell}
$$

and we set $g_{1}(z):=\kappa_{2} z+\sum_{j \geqslant 2} a_{j} z^{j}$ so that $f_{2} \circ \sigma=\sigma \circ g_{1}$. Note that $g_{1}$ is a local isomorphism at $p_{1}$ which locally satisfies $g_{1}^{-1}\left(J_{1}\right)=J_{1}$ in Case (1) and $g_{1}^{*} \mu_{1} \propto \mu_{1}$ in Case (2).

Lemma 2.2 implies that $f_{1}$ and $g_{1}$ commute. In the linearizing coordinate of $f_{1}$, the map $\widetilde{g}_{1}$ corresponding to $g_{1}$ is a local biholomorphism satisfying $\widetilde{g}_{1}\left(\lambda_{1} \zeta\right)=\lambda_{1} \widetilde{g}_{1}(\zeta)$. Expanding $g_{1}$ in power series, and since $\lambda_{1}$ is not a root of unity, we obtain that $\widetilde{g}_{1}$ is linear: $\widetilde{g}_{1}(\zeta)=\kappa_{2} \zeta$. The subgroup generated by $\lambda_{1}$ and $\kappa_{2}$ in $\mathbb{C}^{*}$ must be discrete otherwise by taking sequences $\left(k_{j}\right)$ and $\left(\ell_{j}\right)$ such that $\lambda_{1}^{k_{j}} \kappa_{2}^{\ell_{j}} \rightarrow 1$ we would create an infinite normal family of local symmetries of $J_{1}$ contradicting the fact that $f_{1}$ is not integrable. It follows that there is a relation of the form $\lambda_{1}^{a}=\kappa_{2}^{b}$ for some $a \in \mathbb{N}$ and $b \in \mathbb{Z} \backslash\{0\}$. Since $p_{2} \in J_{2},\left|\lambda_{2}\right| \geqslant 1$ so $\left|\kappa_{2}\right| \geqslant 1$. Since $\widetilde{g}_{1}$ has infinite order, $\kappa_{2}$ is not a root of unity so $a$ is positive. This implies that $b$ is positive as well, hence $\left|\lambda_{2}\right|>1$, i.e. $p_{2}$ is repelling. Thus we have shown that there is a relation of the form $\lambda_{1}^{\ell a}=\lambda_{2}^{b}$, with
$a, b>0$, as asserted. Back to the initial coordinates, this means that $f_{1}^{a}=g_{1}^{b}$ so that $\sigma \circ f_{1}^{a}=\sigma \circ g_{1}^{b}=f_{2}^{b} \circ \sigma$.

Now observe that with $\chi_{1}=\psi_{\left(f_{1}, p_{1}\right)}$ we have

$$
\sigma \circ \chi_{1}\left(\lambda_{1}^{a} \zeta\right)=\sigma \circ f_{1}^{a}\left(\chi_{1}(\zeta)\right)=\sigma \circ g_{1}^{b}\left(\chi_{1}(\zeta)\right)=f_{2}^{b} \circ \sigma \circ \chi_{1}(\zeta)
$$

hence by Lemma 2.9 locally we have $\sigma \circ \chi_{1}(\zeta)=\psi_{\left(f_{2}, p_{2}\right)}\left(\zeta^{\ell}\right)$. Set $\chi_{2}(\zeta):=\psi_{\left(f_{2}, p_{2}\right)}\left(\zeta^{\ell}\right)$, which by definition is a generalized Poincaré-Koenigs map. Locally near the origin we have $\sigma \circ \chi_{1}=\chi_{2}$, hence

$$
\begin{equation*}
\chi_{2}^{-1}\left(J_{2}\right)=\left(\sigma \circ \chi_{1}\right)^{-1}\left(J_{2}\right)=\chi_{1}^{-1}\left(\sigma^{-1}\left(J_{2}\right)\right)=\chi_{1}^{-1}\left(J_{1}\right) \tag{5}
\end{equation*}
$$

Since $\chi_{2}\left(\lambda_{1}^{a} \zeta\right)=f_{2}^{b} \circ \chi_{2}(\zeta), \chi_{2}^{-1}\left(J_{2}\right)$ is invariant under multiplication by $\lambda_{1}^{a}$. The same holds evidently for $\chi_{1}^{-1}\left(J_{1}\right)$, so 5 propagates from a neighborhood of 0 to the whole complex plane.

Now assume that we are in Case ( $1^{\prime}$ ) so that $\sigma^{*} \mu_{2} \asymp_{c} \mu_{1}$. From the relation $f_{2}^{n} \circ \sigma=$ $\sigma \circ g_{1}^{n}$, for all $n \in \mathbb{N}$ we obtain

$$
d_{2}^{b n} \mu_{1} \asymp_{c^{2}}\left(g_{1}^{b n}\right)^{*} \mu_{1},
$$

so that

$$
d_{1}^{a n} \mu_{1}=\left(f_{1}^{a n}\right)^{*} \mu_{1}=\left(g_{1}^{b n}\right)^{*} \mu_{1} \asymp_{c^{2}} d_{2}^{b n} \mu_{1}
$$

which implies that $d_{1}^{a}=d_{2}^{b}$.
Locally near the origin we have that

$$
\begin{equation*}
\chi_{2}^{*} \mu_{2}=\chi_{1}^{*}\left(\sigma^{*} \mu_{2}\right) \asymp_{c} \chi_{1}^{*} \mu_{1} \tag{6}
\end{equation*}
$$

It remains to explain why the relation (6) propagates to $\mathbb{C}$. Write $M_{\kappa}(\zeta):=\kappa \zeta$, and define positive measures on $\mathbb{C}$ by $\widetilde{\mu}_{1}:=\chi_{1}^{*} \mu_{1}$ and $\widetilde{\mu}_{2}:=\chi_{2}^{*} \mu_{2}$. We have

$$
M_{\lambda_{1}^{a}}^{*} \widetilde{\mu}_{1}=\left(\chi_{1} \circ M_{\lambda_{1}^{a}}\right)^{*} \mu_{1}=\chi_{1}^{*}\left(f_{1}^{*} \mu_{1}\right)=d_{1}^{a} \widetilde{\mu}_{1}
$$

and likewise, since $\chi_{2} \circ M_{\lambda_{1}^{a}}=f_{2}^{b} \circ \chi_{2}$ we get

$$
M_{\lambda_{1}^{a}}^{*} \tilde{\mu}_{2}=M_{\lambda_{1}^{a}}^{*} \chi_{2}^{*} \mu_{2}=\chi_{2}^{*}\left(f_{2}^{b}\right)^{*} \mu_{2}=d_{2}^{b} \tilde{\mu}_{2}=d_{1}^{a} \tilde{\mu}_{2}
$$

Therefore $\widetilde{\mu}_{1}$ and $\widetilde{\mu}_{2}$ are positive measures on $\mathbb{C}$ satisfying the same relation $M_{\Lambda}^{*} \widetilde{\mu}_{i}=D \widetilde{\mu}_{i}$ (for $\Lambda=\lambda_{1}^{a}$ and $D=d_{1}^{a}$ ) and such that $\widetilde{\mu}_{1} \asymp_{c} \widetilde{\mu}_{2}$ in some small disk $D(0, r)$. Since $M_{\Lambda}$ is invertible on $\mathbb{C}$, it follows that $\widetilde{\mu}_{1} \asymp_{c} \widetilde{\mu}_{2}$ globally. Indeed, let $A$ be any Borel set and let $n$ be so large that $M_{\Lambda}^{-n}(A) \subset D(0, r)$. Then

$$
\widetilde{\mu}_{2}(A)=D^{n} \widetilde{\mu}_{2}\left(M_{\Lambda}^{-n}(A)\right) \leqslant c D^{n} \widetilde{\mu}_{1}\left(M_{\Lambda}^{-n}(A)\right)=c \widetilde{\mu}_{1}(A)
$$

and similarly for the reverse inequality, so we are done.
In Case (2), repeating the same argument with $\sigma^{*} \mu_{2} \propto \mu_{1}$ we arrive at $\widetilde{\mu}_{1} \propto \widetilde{\mu}_{2}$, and the proof is complete.

Remark 2.12. This argument is reminiscent from the work of Ghioca, Nguyen and Ye [19].
2.6. Conclusion of the proof of Theorem A. Recall that $\sigma: U \rightarrow \mathbb{P}^{1}$ is a nonconstant holomorphic map such that $\sigma^{*} \mu_{f_{2}}=\mu_{f_{1}}$ (resp. $\sigma^{*} \mu_{f_{2}} \propto \mu_{f_{1}}$ when $J_{1}$ is smooth), and $\sigma\left(p_{1}\right)=p_{2}$ where $p_{1}$ is a repelling periodic point for $f_{1}$ and $p_{2}$ a preperiodic point for $f_{2}$. Replacing $f_{1}$ by a suitable iterate, and $\sigma$ by $f_{2}^{k} \circ \sigma$ for a suitable $k$ we may suppose that $f_{1}\left(p_{1}\right)=p_{1}$ and $f_{2}\left(p_{2}\right)=p_{2}$. Also we write $\ell=\operatorname{deg}_{p_{1}}(\sigma)$.

By Proposition 2.10, the point $p_{2}$ is repelling and there exist a generalized PoincaréKoenigs maps $\chi_{1}, \chi_{2}$ associated to $\left(f_{1}, p_{1}\right)$ and $\left(f_{2}, p_{2}\right)$ such that $\chi_{1}^{*} \mu_{1}=\chi_{2}^{*} \mu_{2}$. We also have the relations $d_{1}^{a}=d_{2}^{b}$ for some $a, b \in \mathbb{N}^{*}$, and

$$
\begin{equation*}
\chi_{1}\left(\lambda_{1}^{a} \zeta\right)=f_{1}^{a}\left(\chi_{1}(\zeta)\right) \text { and } \chi_{2}\left(\lambda_{1}^{a} \zeta\right)=\sigma \circ \chi_{1}\left(\lambda_{1}^{a} \zeta\right)=\sigma \circ f_{1}^{a} \circ \chi_{1}(\zeta)=f_{2}^{b}\left(\chi_{1}(\zeta)\right) . \tag{7}
\end{equation*}
$$

Let $\Psi=\left(\chi_{1}, \chi_{2}\right): \mathbb{C} \rightarrow \mathbb{P}^{1} \times \mathbb{P}^{1}$ and $F=\left(f_{1}^{a}, f_{2}^{b}\right)$. By Proposition 2.6, $\overline{\Psi(\mathbb{C})}$ is an irreducible algebraic curve $Z$ which is neither a horizontal nor a vertical line, and from (7) we deduce that

$$
F(\Psi(\zeta))=\Psi\left(\lambda_{1}^{a} \zeta\right)
$$

so $Z$ is $F$-invariant.
Now note that the first (resp. second) projection $\pi_{1}: \mathbb{P}^{1} \times \mathbb{P}^{1} \rightarrow \mathbb{P}^{1}$ (resp. $\pi_{2}$ ) semiconjugates $\left.F\right|_{Z}$ to $f_{1}$ (resp. to $f_{2}$ ). This implies that the measure of maximal entropy $\left.\mu\right|_{F}$ is equal to

$$
\mu_{F \mid Z}=\frac{1}{\operatorname{deg}\left(\pi_{1}\right)} \pi_{1}^{*} \mu_{f_{1}}=\frac{1}{\operatorname{deg}\left(\pi_{2}\right)} \pi_{2}^{*} \mu_{f_{1}}
$$

which implies $\sigma^{*} \mu_{f_{2}} \propto \mu_{f_{1}}$ (see e.g., [19] for details). This concludes the proof.
Remark 2.13. Theorem $A$ fails in the integrable case.
Indeed, let $f_{1}$ and $f_{2}$ be arbitrary Lattès maps, associated to finite branched covers $\pi_{i}: \mathbb{C} / \Lambda_{i} \rightarrow \mathbb{P}^{1}$, for some lattices $\Lambda_{i} i=1,2$. Let $p_{i}$ be any point in $\mathbb{P}^{1}$ located outside the critical value locus of $\pi_{i}$. Then there is a local measure preserving isomorphism $\sigma$ mapping $p_{1}$ to $p_{2}$. Indeed let $q_{i}$ be a lift of $p_{i}$ in $\mathbb{C} / \Lambda_{i}, \tilde{q}_{i}$ be a lift of $q_{i}$ in $\mathbb{C}$, and $\tilde{\pi}_{i}:\left(\mathbb{C}, \tilde{q}_{i}\right) \rightarrow\left(\mathbb{P}^{1}, p_{i}\right)$ the natural germ of biholomorphism. Then $\left(\tilde{\pi}_{i}^{-1}\right)_{*} \mu_{i}$ is proportional to the Haar measure of the torus. Therefore, putting $\sigma=\tilde{\pi}_{2} \circ \tau \circ\left(\tilde{\pi}_{1}\right)^{-1}$, where $\tau$ is the translation mapping $\tilde{q}_{1}$ to $\tilde{q}_{2}$, we get $\sigma^{*} \mu_{2} \propto \mu_{1}$.

On the other hand, there is an algebraic correspondence between $f_{1}$ and $f_{2}$ when and only when there is an isogeny between the corresponding elliptic curves $\mathbb{C} / \Lambda_{1}$ and $\mathbb{C} / \Lambda_{2}$, and $\lambda_{1}^{a}=\lambda_{2}^{b}$ for some $a, b \in \mathbb{N}^{*}$ (where $\lambda_{i}$ is the derivative of any lift of $f_{i}$ to $\mathbb{C}$ ).

An analogous discussion can be made in the monomial case.

## 3. LOCAL CONTRACTIONS AND PREPERIODIC POINTS

It is natural to expect that pre-repelling points in the Julia set are geometrically characterized by the existence of a contracting local symmetry. We confirm this intuition when $f$ satisfies suitable expansion properties on its Julia set, namely when $f$ satisfies the topological Collet-Eckmann (TCE) condition. This condition can be defined in a number of equivalent ways, for instance by the following exponential shrinking property: there exists $\lambda>1$ and $r>0$ such that for every $x \in J_{f}$ and $n \in \mathbb{N}$, every connected component $W$ of $f^{-n}(B(x, r))$ satisfies $\operatorname{Diam}(W) \leqslant \lambda^{-n}$. We refer to [33] for a thorough discussion of this notion.

Theorem 3.1. Let $f$ be a non-integrable rational map satisfying the Topological ColletEckmann condition. Let $\sigma: U \rightarrow \mathbb{P}^{1}$ be a non-constant holomorphic map satisfying $\sigma^{-1}\left(J_{f}\right) \cap U=J_{f} \cap U$, and furthermore $\sigma^{*} \mu_{f} \propto \mu_{f}$ if $J_{f}$ is smooth.

Suppose that there exists $p \in J_{f} \cap U$ such that $\sigma(p)=p$ and $\left|\sigma^{\prime}(p)\right|<1$. Then $p$ is preperiodic to a repelling point.

Proof. Note that it is enough to show that $p$ is preperiodic: indeed for a TCE map all periodic points on the Julia set are hyperbolic.

To make the main idea more transparent, we sketch a proof under the stronger assumption that $f$ is hyperbolic. Then there exists $r>0$ such that for every $n \geqslant 0$, there is a univalent inverse branch $f_{-n}$ of $f^{n}$ on $B\left(f^{n}(p), 2 r\right)$ such that $f_{-n}\left(f^{n}(p)\right)=p$. Reducing $r$ if necessary we may assume that $B(p, r) \Subset U$, where $U$ is the domain of definition of $\sigma$. By the Koebe distortion theorem, $f_{-n}$ has uniformly bounded distortion on $B\left(f^{n}(p), r\right)$. Therefore $f_{-n}\left(B\left(f^{n}(p), r\right)\right)$ is up to uniformly bounded distortion a disk centered at $p$ and of radius $r_{n}$, where $r_{n}$ decreases exponentially with $n$. Let now $k=k(n)$ be the least integer such that $C^{2}\left(\sigma^{\prime}(0)\right)^{k} r<r_{n}$, where $C$ bounds the distortion of $f_{-n}$ on $B\left(f^{n}(p), r\right)$ and the distortion of $\sigma^{k}$ on $B(p, r)$. Then $f^{n} \circ \sigma^{k(n)}$ is a sequence of univalent symmetries of $J$ defined on $B(p, r)$, with derivative at $p$ bounded away from 0 and infinity. Levin's theorem entails that this sequence is finite, and we conclude that there exists $n_{1}<n_{2}$ such that $f^{n_{1}}(p)=f^{n_{2}}(p)$, as desired.

If $f$ only satisfies the TCE property the argument is similar, except that we can only map a small neighborhood of $p$ to the large scale with bounded degree and along a subsequence of integers. More specifically, the TCE condition of [33, p. 31] reads as follows. There exists a radius $r>0$ and an integer $\delta$ such that if $W_{n}$ denotes the connected component of $f^{-n}\left(B\left(f^{n}(p), r\right)\right)$ containing $p$, then there exists a sequence of integers $\left(n_{j}\right)$ of positive lower density such that

$$
\operatorname{deg}\left(f^{n_{j}}: W_{n_{j}} \rightarrow B\left(f^{n_{j}}(p), r\right)\right) \leqslant \delta
$$

We claim that if $r$ is small enough, $W_{n}$ is simply connected for all $n$. Then by 34, Lemma 2.1], $\left.f^{n_{j}}\right|_{W_{n_{j}}}$ satisfies some bounded distortion properties.

To prove our claim we make the following observations: first, the local structure of holomorphic maps shows that there exists $r_{0}=r_{0}(f)$ such that for $r \leqslant r_{0}$, for every $p$, every component $W$ of $f^{-1}(B(p, r))$ is simply connected. Then $W$ is biholomorphic to a disk in $\mathbb{C}$ so by the maximum principle, if $U \subset B(p, r)$ is a simply connected open set, $f^{-1}(U) \cap W$ is simply connected. Next, by the TCE property there exists $r_{1}$ such that for every $r \leqslant r_{1}$, every $p \in \mathbb{P}^{1}$ and every $n \geqslant 0$ and every component $W_{n}$ of $f^{-n}(B(p, r))$ has diameter smaller than $r_{0}$ (see the Backward Lyapunov Stability condition in [33, §5]). Then the simple connectivity of $W_{n}$ easily follows by induction.

Reduce $r$ if necessary so that $\sigma$ is well defined and univalent on $B(p, 2 r)$, and write $\lambda:=\left|\sigma^{\prime}(0)\right|<1$. By the Koebe distortion theorem, there exist constants $C_{1}$ and $C_{2}$ such that

$$
B\left(p, C_{1} \lambda^{k} r\right) \subset \sigma^{k}(B(p, r)) \subset B\left(p, C_{2} \lambda^{k} r\right)
$$

For any $0<\tau<1$, denote by $W_{n}(\tau)$ the connected component of $f^{-n}\left(B\left(f^{n}(p), \tau r\right)\right)$ containing $p$. To simplify notation we write $W_{n}^{\prime}=W_{n}(1 / 2)$.

Pick $\alpha<1$, and let $k=k\left(n_{j}\right)$ be the least integer such that $\lambda^{k} r \leqslant \alpha \operatorname{Diam}\left(W_{n_{j}}^{\prime}\right)$. Then $\alpha \operatorname{Diam}\left(W_{n_{j}}^{\prime}\right) \leqslant \lambda^{k-1} r$, and we get

$$
B\left(p, \alpha \lambda C_{1} \operatorname{Diam}\left(W_{n_{j}}^{\prime}\right)\right) \subset \sigma^{k}(B(p, r)) \subset B\left(p, \alpha C_{2} \operatorname{Diam}\left(W_{n_{j}}^{\prime}\right)\right)
$$

Now, by [34, Lemma 2.1 (2.3)], we have $B\left(p, \alpha C_{2} \operatorname{Diam}\left(W_{n}^{\prime}\right)\right) \subset W_{n}^{\prime}$ when $\alpha$ small enough, independently on $n$. Furthermore by [34, Lemma 2.1 (2.2)], when $\tau$ is small enough, then for every $n, W_{n}^{\prime}(\tau) \subset B\left(p, \alpha \lambda C_{1} \operatorname{Diam}\left(W_{n}^{\prime}\right)\right)$.

It follows that $f^{n_{j}} \circ \sigma^{k\left(n_{j}\right)}$ is a sequence of symmetries of $J$ defined on $B(p, r)$ which satisfies:

$$
\begin{gathered}
f^{n_{j}} \circ \sigma^{k\left(n_{j}\right)}(B(p, r)) \subset f^{n_{j}}\left(W_{n}^{\prime}\right) \subset B\left(f^{n}(p), r / 2\right), \text { and } \\
\operatorname{Diam}\left(f^{n_{j}} \circ \sigma^{k\left(n_{j}\right)}(B(p, r))\right) \geqslant \operatorname{Diam}\left(f^{n_{j}}\left(W_{n_{j}}(\tau)\right)\right) \geqslant \tau r
\end{gathered}
$$

The first estimate implies that $f^{n_{j}} \circ \sigma^{k\left(n_{j}\right)}$ forms a normal family on $B(p, r)$ and the second that no cluster value of this sequence is constant. At this stage we conclude as in the hyperbolic case: the sequence $f^{n_{j}} \circ \sigma^{k\left(n_{j}\right)}$ must be finite, and we find integers $n_{j_{1}}$ and $n_{j_{2}}$ such that $f^{n_{j_{1}}}(p)=f^{n_{j_{2}}}(p)$.

As a consequence we infer that the assumption that $\sigma$ maps a repelling point to a preperiodic point is superfluous in Theorem $A$ when $f_{2}$ satisfies the TCE property.

Corollary 3.2. Let $f_{1}$ and $f_{2}$ be two rational maps and assume $f_{2}$ is non-integrable and satisfies the TCE property. Let $\sigma: U \rightarrow \mathbb{P}^{1}$ be any non-constant holomorphic map satisfying $\sigma^{-1}\left(J_{2}\right) \cap U=J_{1} \cap U$ if $J_{f_{1}}$ is not smooth, and $\sigma^{*} \mu_{2} \propto \mu_{1}$ otherwise. Then $\sigma$ maps any repelling periodic point of $f_{1}$ to a pre-repelling point of $f_{2}$.

Proof. Note that $J_{2}$ is smooth if and only if $J_{1}$ is smooth. Fix a repelling periodic $p_{1}$ of $f_{1}$ of period $k$ outside the critical set of $\sigma$. Then $\sigma \circ f_{1}^{-k} \circ \sigma^{-1}$ defines a local holomorphic contraction of $J_{2}$ at $\sigma\left(p_{1}\right)$, which furthermore preserves $\mu_{2}$ up to a constant if $J_{2}$ is smooth, thus the previous proposition gives the result.

Remark 3.3. The TCE property is detected by the maximal entropy measure: indeed it is equivalent to the property that the measure of small balls satisfies an estimate of the form $\mu(B(x, r)) \gtrsim r^{\theta}$ for some $\theta>0$ and for every $x \in J$ (see [36]). It is not difficult to see that if such an estimate holds for every $x \in U$, where $U$ is any open set interesting $J_{f}$, then it holds everywhere (possibly with a different $\theta$ ). It follows that under the assumptions of Theorem A, $f_{1}$ is TCE if and only if $f_{2}$ is TCE.

Remark 3.4. Pick any local symmetry $\sigma$ of $J_{f}$. By precomposing with some inverse branch of $f$, we may always assume that it satisfies $\sigma(U) \Subset U$, so it has an attracting fixed point. The proof of Theorem 3.1 then gives the existence of integers $n_{1}<n_{2}$ and $k_{1}<k_{2}$ such that $f^{n_{1}} \circ \sigma^{k_{1}}=f^{n_{2}} \circ \sigma^{k_{2}}$, so $\sigma^{k_{2}-k_{1}}=f^{-n_{2}} \circ f^{n_{1}}$, and we infer that an iterate of $\sigma$ is the restriction of an algebraic correspondence.

Unfortunately, the algebraicity of $\sigma$ itself does not seem to follow from this relation, and the only route we know of to algebraicity goes through measure class preservation and Theorem A. Still, this gives additional credit to the problem stated in the Introduction.

## 4. The polynomial case

For polynomials the maximal entropy measure is determined by the Julia set: indeed it coincides with the harmonic measure of $K_{f}$ viewed from infinity. Thus in this case it is natural to expect that the measure class preservation hypothesis in Theorem A should follow from weaker property $\sigma^{-1}\left(J_{2}\right) \cap U=J_{1} \cap U$. However, since we are working locally, some technicalities arise and we are only able to confirm this expectation under a mild additional assumption.

Theorem 4.1. Let $f_{1}$ and $f_{2}$ be non-integrable polynomials such that:
(1) either $J_{1}$ and $J_{2}$ are disconnected;
(2) or $J_{1}$ and $J_{2}$ are connected and locally connected.

Let $U \subset \mathbb{C}$ be an open subset intersecting $J_{1}$ and $\sigma: U \rightarrow \mathbb{C}$ be a non-constant holomorphic map such that $\sigma^{-1}\left(J_{2}\right) \cap U=J_{1} \cap U$. Then there exists $\Omega \subset U$ intersecting $J_{1}$ on which $\sigma^{*} \mu_{2}=\mu_{1}$.

Remark 4.2. Under the assumptions of Theorem 4.1, $J_{1}$ is connected (resp. locally connected) iff $J_{2}$ is connected (resp. locally connected), so we could state the assumption only for one of $J_{1}$ or $J_{2}$.

Indeed suppose $J_{1}$ is disconnected. Then point components of $J_{1}$ accumulate the whole Julia set (see $\$ 4.1$ below), hence $J_{2} \cap \sigma(U)$ admits points components, and $J_{2}$ is disconnected.

When $J_{1}$ is locally connected, then it is clear that $J_{2} \cap \sigma(U)$ is also locally connected. Since $f_{2}$ is open and $f_{2}^{n}(\sigma(U)) \supset J_{2}$ for large $n$, we conclude that $J_{2}$ is also connected.

We now have all the necessary ingredients for Theorem B.
Proof of Theorem B. Restricting $U$ if necessary we may assume that $\sigma$ is a biholomorphism, and without loss of generality we may assume that $f_{2}$ is non-integrable and TCE.

When $J_{2}$ is disconnected, then $J_{1}$ is disconnected too by the preceding remark, and Theorem 4.1 implies $\sigma^{*} \mu_{2} \asymp \mu_{1}$ on some $\Omega \subset U$. By Remark 3.3, $f_{1}$ is TCE, and nonintegrable. Corollary 3.2 implies that $\sigma$ maps any repelling periodic point of $f_{1}$ to a preperiodic point of $f_{2}$. Thus the result follows from Theorem A.

When $J_{2}$ is connected, then it is not locally smooth since $f_{2}$ is not integrable, and it is locally connected by the TCE property; see [27]. Applying the local biholomorphism $\sigma$, we infer that $J_{1}$ is not smooth, hence $f_{1}$ is not integrable, and by Remark 4.2, $J_{1}$ is connected and locally connected. We conclude as in the previous case.

The remainder of this section will be devoted to the proof of Theorem 4.1. It relies on a localization principle for harmonic measure, which requires different arguments in the disconnected and connected cases; the latter is the most delicate, a more precise outline is given at the beginning of $\$ 4.2$.

Let us first introduce a few standard potential theoretic tools (see [28] for a gentle introduction to the probabilistic viewpoint on potential theory, [10] for a systematic account, and [18] for the planar case). If $\Omega$ is a domain on the Riemann sphere with non-polar complement, for $z \in \Omega$ the harmonic measure $\omega(z, \cdot, \Omega)$ is the measure on $\partial \Omega$
defined by declaring that $\int \varphi(w) \omega(z, d w, \Omega)$ is the value at $z$ of the solution of the Dirichlet problem with boundary values given by $\varphi \in \mathcal{C}(\partial \Omega)$. It is also the exit distribution of the Brownian motion issued from $z$, that is, if we denote by $B_{z}$ the Brownian motion issued from $z$ on $\mathbb{P}^{1}$, and $\tau_{z}=\inf \left\{t>0, B_{z}(t) \in \partial \Omega\right\}$ the hitting time of the boundary (which is a.s. finite), then for a (say closed) subset $E \subset \partial \Omega, \omega(z, E, \Omega)=\mathbb{P}\left(B_{z}\left(\tau_{z}\right) \in E\right)$. If $\Omega$ is simply connected, let $\phi: \mathbb{D} \rightarrow \Omega$ be a uniformizing map such that $\phi(0)=z$ (which is unique up to pre-composition with a rotation). Then $\phi$ extends radially outside a set of rays of zero capacity, and still denoting by $\phi$ this extension we have

$$
\omega(z, E, \Omega)=\omega\left(0, \phi^{-1}(E), \mathbb{D}\right)=\frac{1}{2 \pi}\left|\phi^{-1}(E)\right|
$$

where $|\cdot|$ denotes usual arclength. We will only use this fact when $\partial \Omega$ is locally connected, in which case by the Caratheodory theorem $\phi$ extends to a continuous surjection $\overline{\mathbb{D}} \rightarrow \bar{\Omega}$.

If $f$ is a polynomial, the properties of the Green function $G_{f}$ imply that the maximal entropy measure $\mu_{f}$ coincides with the harmonic measure of the basin of infinity: $\mu_{f}=$ $\omega\left(\infty, \cdot, K_{f}^{\complement}\right)$ (where the complement here is understood in $\mathbb{P}^{1}(\mathbb{C})$ ).

Remark 4.3. Theorem B has some overlap with [24, Theorem 1.4] (see also [25, Prop. $4.5]$ ), which relies on completely different ideas (even if Levin's finiteness theorem also plays a key role there). In his paper, Luo assumes that $f_{1}$ and $f_{2}$ are polynomials of the same degree with connected and locally connected Julia sets, and all these assumptions are essential. In this setting, his result is stronger than ours since no no additional hyperbolicity assumption is required to guarantee that a periodic point is mapped to a preperiodic point. Note that by applying his methods, we can obtain the following generalization of [24, Theorem 1.1]: if $M_{d}$ denotes the degree $d$ Multibrot set, then for $d \neq d^{\prime}$, an open subset of $\partial M_{d}$ cannot be biholomorphic to an open subset of $\partial M_{d^{\prime}}$.
4.1. Proof of Theorem 4.1 in the disconnected case. The localization principle that we use in the disconnected case is the following:

Lemma 4.4. Let $f$ be a polynomial of degree $\geqslant 2$. Let $\Omega$ be a connected and simply connected open set with smooth boundary such that $\Omega \cap J_{f} \neq \varnothing$ and $\partial \Omega \cap K_{f}=\varnothing$. Then for every $z \in \Omega \backslash K_{f}$, there exists a constant $c>0$ such that $c^{-1} \mu_{f} \leqslant \omega\left(z, \cdot, \Omega \backslash K_{f}\right) \leqslant c \mu_{f}$ on $J_{f} \cap \Omega$.

This is most likely well-known, however we give a (probabilistic) proof because we have not been able to locate it in the literature.

Proof. Since $K_{f}$ is full, $\Omega \backslash K_{f}$ is connected. By the Harnack inequality, if $L \subset \Omega \backslash K_{f}$ then for any $\left(z, z^{\prime}\right) \in L^{2}$, there exists $c=c(L)$ such that

$$
c^{-1} \omega\left(z, \cdot, \Omega \backslash K_{f}\right) \leqslant \omega\left(z^{\prime}, \cdot, \Omega \backslash K_{f}\right) \leqslant c \omega\left(z, \cdot, \Omega \backslash K_{f}\right) .
$$

In particular if the conclusion of the lemma is true for some $z \in \Omega \backslash K_{f}$, then it is true for every $z \in \Omega \backslash K_{f}$. Reduce $\Omega$ a little bit to get a smoothly bounded $\Omega^{\prime} \Subset \Omega$ with the same properties as $\Omega$ and such that $\Omega \cap K_{f}=\Omega^{\prime} \cap K_{f}$. Pick $z \in \Omega^{\prime}$. Let $\tilde{\mu}=\omega\left(z, \cdot, \mathbb{P}^{1} \backslash K_{f}\right)$. By the Harnack inequality again, $\tilde{\mu}$ is equivalent to $\mu_{f}$ with uniform bounds, so it is enough to compare $\tilde{\mu}$ and $\omega\left(z, \cdot, \Omega \backslash K_{f}\right)$.

Denote by $\nu_{z}$ the restriction of $\omega\left(z, \cdot, \Omega \backslash K_{f}\right)$ to $J_{f} \cap \Omega$. The difference between $\left.\tilde{\mu}\right|_{\Omega \cap J_{f}}$ and $\nu_{z}$ is accounted for by the contributions of Brownian paths leaving $\Omega$ before reaching $\Omega \cap J_{f}$. If a Brownian path from $z$ eventually hits $\Omega \cap J_{f}$ without staying in $\Omega$, then it must $\operatorname{cross} \partial \Omega^{\prime}$. Let $B_{z}$ be the Brownian motion from $z$ in $K_{f}^{\complement}$ killed when hitting $J_{f}$, and $\tau$ be the hitting time of $\Omega \cap J_{f}$. Then $\tau<\infty$ if $B_{z}(t)$ exits $K_{f}^{\complement}$ in $\Omega \cap J_{f}$, and the distribution of $B_{z}(\tau)$ conditioned to $\tau<\infty$ is the harmonic measure $\left.\tilde{\mu}\right|_{J_{f} \cap \Omega}$. Introduce the following sequence of stopping times: $T_{0}^{\prime}=0$, and by induction $T_{i+1}=\inf \left\{t>T_{i}^{\prime}, B_{z}(t) \in \partial \Omega\right\}$ and $T_{i+1}^{\prime}=\inf \left\{t>T_{i+1}, B_{z}(t) \in \partial \Omega^{\prime}\right\}$, so that $T_{1}$ is the first hitting time of $\partial \Omega, T_{1}^{\prime}$ the first hitting time of $\partial \Omega^{\prime}$ after $T_{1}$, etc. The event $\tau<T_{1}$ holds if $B_{z}$ does not reach $\partial \Omega$ before hitting $J_{f}$ so the distribution of $B_{z}(\tau)$ conditioned to $\left\{\tau<T_{1}\right\}$ is $\nu_{z}$. Now conditioned to the event that $\left\{T_{i}^{\prime}<\tau<T_{i+1}\right\}$, the distribution of $B_{z}\left(T_{i}^{\prime}\right)$ is a certain probability measure $p_{i}$ on $\partial \Omega^{\prime}$. Hence by the strong Markov property of Brownian motion [28, §2.2] the conditional distribution of $B_{z}(\tau)$ is $\nu^{(i)}:=\int \nu_{w} d p_{i}(w)$, which satisfies $c^{-1} \nu_{z} \leqslant \nu^{(i)} \leqslant c \nu_{z}$ for a constant $c$ depending only on $\partial \Omega^{\prime}$ and $z$. Finally, decomposing the event $\{\tau<\infty\}$ as a disjoint union $\{\tau<\infty\}=\bigcup_{i \geqslant 0}\left\{T_{i}^{\prime}<\tau<T_{i+1}\right\}$, we express $\left.\tilde{\mu}\right|_{J_{f} \cap \Omega}$ (normalized by $\mathbb{P}(\tau<\infty)$ ) as an infinite convex combination of $\nu_{z}$ and of the $\nu^{(i)}$, and the lemma follows.

Conclusion of the proof of the theorem. Assume that the assumptions of Theorem 4.1 hold, with $J_{1}$ disconnected.

Since point components are dense in $J_{1}$ (see DeMarco-McMullen [8] or Emerson [13]), $J_{1}$ admits arbitrary small relatively compact components in $U$.

Therefore we can fix a smoothly bounded simply connected open set $\Omega$ intersecting $J_{1}$, relatively compact in $U$ and such that $\partial \Omega \cap J_{1}=\varnothing$. By choosing $\Omega$ small enough we can further assume that $\sigma$ is a biholomorphism in a neighborhood of $\bar{\Omega}$. Its image $\sigma(\Omega)$ under $\sigma$ satisfies the same properties relatively to $J_{2}$ and $\sigma(U)$. By the holomorphic invariance of harmonic measure, for $z \in \Omega \backslash K_{1}$ we have that $\sigma_{*}(\omega(z, \cdot, \Omega))=\omega(\sigma(z), \cdot, \sigma(\Omega))$, so the property $\sigma^{*} \mu_{2}=\mu_{1}$ follows from Lemma 4.4 applied to $\Omega \cap J_{1}$ and $\sigma\left(\Omega \cap J_{2}\right)$.
4.2. Proof of Theorem 4.1 in the connected case. To establish Theorem 4.1 for connected $J_{f}$ we face several difficulties. The first one is that we need to take care of possible boundary effects in Lemma 4.4. indeed the argument of Lemma 4.4 breaks down since we cannot assume that $\Omega \cap J_{f}=\varnothing$. For this, we uniformize $K_{f}^{\complement}$ and use some facts from Caratheodory theory (see $\S 4.2 .1$ as well as the notion of endpoint in $\$ 4.2 .2$ ). The other new difficulty is that given a small open set $\Omega$ intersecting $J_{f}$ and $x \in \Omega \backslash J_{f}$, we need to detect whether $x$ belongs to $K_{f}$ or not, by using only the data given by $J_{f} \cap \Omega$. If $J_{f}$ is not a Jordan curve this can be done by looking at the local topological properties of $J_{f}$ (the endpoints are also used here). In the Jordan curve case we cannot distinguish the inside from the outside of $J_{f}$ from topology, nor even from complex analysis when the Julia set is a quasicircle, so a completely different argument needs to be found, which is postponed to the next section (see Proposition 5.1). Note that this study is necessary because the harmonic measures viewed from the two sides of a non-smooth Jordan curve are typically mutually singular (see [18, Theorem VI.6.3]).
4.2.1. From local to global harmonic measure. We denote by $\arg (\cdot)$ the argument function $\mathbb{C} \backslash \mathbb{R}_{-}$, with values in $(-\pi, \pi)$.

Lemma 4.5. For $0<\theta<\pi$ and $0<\delta<1$, let $S_{\theta, \delta} \subset \mathbb{D}$ be the sector defined by

$$
S_{\theta, \delta}=\{z,|\arg (z)|<\theta, 1-\delta<|z|<1\}
$$

There exists a constant $c=c(\theta, \delta)$ such that for $\zeta_{0}=1-\delta / 2$, if $E$ is any measurable subset of $\partial \mathbb{D} \cap\{z,|\arg (z)| \leqslant \theta / 2\}$,

$$
\begin{equation*}
\omega\left(\zeta_{0}, E, S_{\theta, \delta}\right) \leqslant \omega\left(\zeta_{0}, E, \mathbb{D}\right) \leqslant c \omega\left(\zeta_{0}, E, S_{\theta, \delta}\right) \tag{8}
\end{equation*}
$$

Proof. The first inequality in (8) follows automatically from the fact that $S_{\theta, \delta} \subset \mathbb{D}$. To prove the second one, let $\phi: \mathbb{D} \rightarrow S_{\theta, \delta}$ be the uniformisation such that $\phi(0)=z_{0}$ and $\phi^{\prime}(0) \in \mathbb{R}_{+}$, which depends only on $(\theta, \delta)$ and extends as a homeomorphism from $\partial \mathbb{D}$ to $\partial S_{\theta, \delta}$. Then $\omega\left(z_{0}, E, S_{\theta, \delta}\right)=\frac{1}{2 \pi}\left|\phi^{-1}(E)\right|$. On the other hand $\phi^{-1}\left(\overline{S_{\theta, \delta}} \cap \partial \mathbb{D}\right)$ is a closed circular arc $I_{0}$, and for every sub-arc $I \Subset I_{0}, \phi$ extends by Schwarz reflection to a biholomorphism in a neighborhood of $I$. Fix $I=\phi^{-1}(\{z \in \partial \mathbb{D},|\arg (z)| \leqslant \theta / 2\})$. Then $\phi: I \rightarrow \phi(I)$ distorts lengths by a uniformly bounded amount, so $\left|\phi^{-1}(E)\right|=|E|$, where the implied constant depends only on $\theta$ and $\delta$. By the Harnack inequalities, we get that $|E|=2 \pi \omega(0, E, \mathbb{D}) \asymp \omega\left(\zeta_{0}, E, \mathbb{D}\right)$, and this concludes the proof.

Let $K$ be a full connected and locally connected compact set of $\mathbb{C}$, containing at least two points. Denote by $\phi_{K}: \mathbb{D} \rightarrow K^{\complement}$ the uniformisation map fixing $\infty$ and tangent to the identity at $\infty$. Recall that by the Caratheodory theorem, $\phi_{K}$ extends continuously to a $\operatorname{map} \phi_{K}: \overline{\mathbb{D}} \rightarrow K^{\complement} \cup \partial K$. A crosscut of $K$ is an open Jordan $\operatorname{arc} C$ in $\mathbb{C} \backslash K$ such that $\bar{C}=C \cup\{a, b\}$ with $a, b \in \partial K$. Note that we allow $a=b$. It follows that the open set $K^{\complement} \backslash C$ admits two connected components, see [32, Proposition 2.12].

Lemma 4.6. Let $K$ be any connected and locally connected full compact subset of the complex plane. Let $C$ be any crosscut of $K$, and denote by $W$ the bounded connected component of $K^{\complement} \backslash C$. Suppose that $\bar{W} \cap \partial K$ is not reduced to a singleton (that is, it is not reduced to $\bar{C} \cap \partial K)$.

Then for any point $z \in W$, and for any $x \in \partial K \cap \bar{W} \backslash \bar{C}$, there exist a neighborhood $\Omega$ of $x$ and a constant $c>0$ such that

$$
\left.c^{-1} \omega\left(\infty, \cdot, K^{\complement}\right)\right|_{\Omega} \leqslant\left.\omega(z, \cdot, W)\right|_{\Omega} \leqslant\left. c \omega\left(\infty, \cdot, K^{\complement}\right)\right|_{\Omega}
$$

Proof. Lift $C$ to $\hat{C}:=\phi_{K}^{-1}(C) \subset \mathbb{D}$. This is an arc in $\mathbb{D}$ whose closure in $\overline{\mathbb{D}}$ intersects $\partial \mathbb{D}$ in two points $e^{i \theta_{0}}$ and $e^{i \theta_{1}}$. Let $\hat{W}$ be the connected component of $\hat{C} \backslash \mathbb{D}$ which is mapped to $W$, and denote by $I:=\left[\theta_{0}, \theta_{1}\right]$ the $\operatorname{arc} \mathbb{S}^{1} \cap \partial \hat{W}$. Observe that $\phi_{K}(I)=\bar{W} \cap \partial K$. By assumption, $I$ contains $\phi_{K}^{-1}(x)$, so $I$ is non trivial and in particular, $\theta_{0} \neq \theta_{1}$.

Rotate the situation so that $1 \in I$ and $\phi_{K}(1)=x$. Fix a sector $S=S_{\theta, \delta}$ such that $S \subset \hat{W}$, and $\phi_{K}^{-1}(x) \subset(-\theta, \theta)$. Pick any $\zeta \in S$ and any open arc $A$ satisfying $\phi_{K}^{-1}(x) \subset A \subset \bar{A} \subset(-\theta, \theta)$. Choose $\Omega=D(x, r)$ with $r>0$ small enough so that the closure of $\phi_{K}^{-1}(\Omega)$ is included in $A$.

Lemma 4.5 now asserts that $\left.\omega(\zeta, \cdot, S)\right|_{A}=\left.\omega(\zeta, \cdot, \mathbb{D})\right|_{A}$. Pushing forward by $\phi_{K}$ and using the conformal invariance of harmonic measure (see, e.g., [28, Theorem 7.22]), we
get that $\left.\omega\left(z, \cdot, \phi_{K}(S)\right)\right|_{\Omega}=\left.\omega\left(z, \cdot, K^{\complement}\right)\right|_{\Omega}$, where $z=\phi_{K}(\zeta)$. Thus we infer that

$$
\left.\omega(z, \cdot, W)\right|_{\Omega} \leqslant\left.\omega\left(z, \cdot, K^{\complement}\right)\right|_{\Omega} \leqslant\left. c \omega\left(z, \cdot, \phi_{K}(S)\right)\right|_{\Omega} \leqslant\left. c \omega(z, \cdot, W)\right|_{\Omega}
$$

where the first and last inequality follow directly from the inclusions $W \subset K^{\complement}$ and $\phi_{K}(S) \subset W$ respectively. Finally by the Harnack inequality, $\omega\left(z, \cdot, K^{\complement}\right)=\omega\left(\infty, \cdot, K^{\complement}\right)$, and the proof is complete.
4.2.2. Endpoints. Assume that $J_{f}$ is connected and locally connected, and write $\phi_{f}=$ $\phi_{K_{f}}$. We say that $x=\phi_{f}\left(e^{i \theta}\right)$ is an endpoint of $J_{f}$ if there exists a sequence of intervals $\left(\theta_{n}^{1}, \theta_{n}^{2}\right)$ in the circle $\mathbb{S}^{1}$, decreasing to $\{\theta\}$ and such that for every $n, \phi\left(\theta_{n}^{1}\right)=\phi\left(\theta_{n}^{2}\right)$. The following lemma is essentially contained in 42].

Lemma 4.7. Let $f$ is a polynomial with connected and locally connected $J_{f}$ and such that $J_{f}$ is neither a Jordan curve nor an interval. Then endpoints are dense in $J_{f}$.

Remark 4.8. If $x \in J_{f}$ is an endpoint, then for any $\delta>0$, for large enough $n$ the image $C$ under $\phi_{f}$ of the chord joining $\theta_{n}$ to $\theta_{n}^{\prime}$ is a crosscut of $K_{f}^{\complement}$ contained in $D(x, \delta)$.

Proof. Recall that the identification of external angles in $K_{f}$ can be encoded in the socalled Thurston lamination $\mathcal{L}_{f}$, which is the lamination by hyperbolic geodesics in $\partial \mathbb{D}$ such that a leaf joins $\theta$ and $\theta^{\prime}$ in $\mathbb{S}^{1}$ whenever $\phi_{f}\left(e^{i \theta}\right)=\phi_{f}\left(e^{i \theta^{\prime}}\right)$ (see [39, Appendix] for details). When $J_{f}$ is not a Jordan curve, this lamination is non trivial, so by backward invariance the set of endpoints of leaves of $\mathcal{L}_{f}$ is dense in $\mathbb{S}^{1}$. A gap of this lamination is the closure of a connected component of $\mathbb{D} \backslash \operatorname{Supp}\left(\mathcal{L}_{f}\right)$. A gap $P$ is the closure of the (hyperbolic) convex hull of its intersection with $\partial \mathbb{D}$, and if we write $\partial P=P \cap \partial \mathbb{D}$, then $\partial \mathbb{D} \backslash \partial P$ has at least three connected components. Since $f$ is not integrable, it follows from [39, Proposition II.6.1] (see also the discussion in [24, §4]) that the union of gaps is dense in $\overline{\mathbb{D}}$.

Pick any gap $P$. Since the leaves of the lamination do not cross, for each connected component $I$ of $\partial \mathbb{D} \backslash \partial P$ we can find a gap $P_{I}$ contained in the convex hull of $I$, and $I \backslash \partial P_{I}$ admits at least 4 components, so two them, say $I_{1}$ and $I_{2}$ satisfy $\left|I_{j}\right| \leqslant|I| / 2$. Proceeding inductively, we construct a sequence of disjoint intervals $I_{\varepsilon_{1} \cdots \varepsilon_{n}}, \varepsilon_{i} \in\{1,2\}$ such that $\left|I_{\varepsilon_{1} \cdots \varepsilon_{n}}\right| \leqslant 2^{-n}$. Note that if $\theta$ is the decreasing intersection of such a sequence of intervals, then $\phi_{f}\left(e^{i \theta}\right)$ is an endpoint. Furthermore by construction these endpoints are separated by crosscuts so they are disjoint.

This argument thus produces a Cantor set, hence an uncountable set, of endpoints. In particular there exists an endpoint $x$ which does not lie in the post-critical set of $f$, so every preimage of $x$ is an endpoint, and we conclude that endpoints are dense in $J_{f}$.
4.2.3. Conclusion of the proof of Theorem 4.1 when $J_{f}$ is not a Jordan curve. Let $\sigma$ be as in the statement of the theorem and reduce $U$ if necessary so that $\sigma$ is a biholomorphism in a neighborhood of $\bar{U}$. Since by assumption $f_{1}$ and $f_{2}$ are not integrable, their Julia set cannot be an interval.

Suppose $J_{1}$ is not a Jordan curve. Then by Lemma 4.7 (see also Remark 4.8) there exists $a \in J_{1} \cap U$ and a crosscut $C$ of $K_{1}^{\complement}$ with $\bar{C} \cap K_{1}=\{a\}$ (so that $C$ is a Jordan curve), and $C \subset U$. Let $W$ be the bounded component of $K_{1}^{C} \backslash C$ and pick $z \in W$. Then $\sigma(C)$ is a crosscut of $J_{2}$ with $\sigma(\bar{C}) \cap J_{2}=\{\sigma(a)\}$. Now by the maximum principle $\sigma(C)$ must be
contained in $K_{2}^{\complement}$ otherwise its interior would be disjoint from $J_{2}$, which is contradictory. By applying Lemma 4.6 to $z$ in $W$ and $\sigma(z)$ in $\sigma(W)$, and using the conformal invariance of harmonic measure, we conclude that $\mu_{1}=\sigma^{*} \mu_{2}$ on some open set $\Omega$ intersecting $J_{1}$, as was to be shown.
4.2.4. The case of Jordan curves. Assume now that $J_{1}$ is a Jordan curve. Then $J_{1}$ has no endpoints in $U$, so that $J_{2}$ has no endpoints either in $\sigma(U)$. It follows from Lemma 4.7 that $J_{2}$ is a Jordan curve as well.

Take a crosscut $C$ of $J_{1}$ in $U$ such that $\bar{C} \cap J_{1}$ consists now of two distinct points, and denote by $W$ the bounded connected component of $K_{1}^{\complement} \backslash C$. When $\sigma(W) \subset K_{2}^{\complement}$, then the arguments of $\$ 4.2 .3$ applies ad litteram. Proposition 5.1 below proves that the possibility that $\sigma(W) \subset K_{2}$ does not occur. This finishes the proof.

## 5. Preserving the sides of a Jordan curve Julia set

Proposition 5.1. Let $f_{1}$ and $f_{2}$ be non-integrable polynomials such that $J_{1}$ is a Jordan curve. Let $U$ be a connected open subset intersecting $J_{1}$ and $\sigma$ be a biholomorphism defined on $U$ such that $\sigma\left(U \cap J_{1}\right)=\sigma(U) \cap J_{2}$. Then $J_{2}$ is a Jordan curve and $\sigma$ maps $K_{1} \cap U$ to $K_{2} \cap \sigma(U)$.

Proof. We already saw in $\S 4.2 .4$ that under the assumptions of the proposition, $J_{2}$ is a Jordan curve, too. We argue by contradiction and assume that $\sigma$ flips the interior of $J_{1}$ to the exterior of $J_{2}$.
Step 1: reduction to the hyperbolic case.
There is no critical point of $f_{1}$ on $J_{1}$, otherwise by pulling back a neighborhood of the corresponding critical value, $J_{1}$ would have several branches at the critical point, and would not be a Jordan curve. The same applies to $J_{2}$. Thus the unique bounded Fatou component of $f_{1}$ (resp. $f_{2}$ ) must be the basin of an attracting or parabolic fixed point.

If $f_{1}$ has an attracting point, then it attracts all critical points, and we conclude that $f_{1}$ is hyperbolic. In particular $J_{1}$ is a quasi-circle, and so is $J_{2}$. In particular $J_{2}$ has no cusps, so $f_{2}$ has no parabolic points, and we infer that $f_{2}$ is hyperbolic as well.

If $f_{1}$ has a parabolic fixed point, then $J_{1}$ admits a dense set of cusps, hence $J_{2}$ cannot be a quasi-circle. It follows that $f_{2}$ is not hyperbolic, hence $J_{2}$ admits a dense set of cusps as well. Since all critical points of $f_{1}$ belong to $\operatorname{Int}\left(K_{1}\right)$, 11, Theorem 3] implies that all cusps are preimages of the cusp at the parabolic fixed point $p$. From this and the fact that all critical points are attracted to $p$, all cusps of $J_{1}$ point inwards, that is, towards the interior of $K_{1}$. The same must be true of $J_{2}$. However, our assumption on $\sigma$ implies that the cusps of $J_{2}$ should point outwards, which is a contradiction. Therefore $f_{1}$ and $f_{2}$ are hyperbolic.

Step 2: using the uniform Levin theorem.
A straightforward compactness argument yields the following uniformity statement in Levin's Theorem:

Corollary 5.2 (of Theorem 2.1). Let $f$ be a rational map with a non-smooth Julia set. Fix $r>0$ and $r_{1}<r_{2}$. Then there exists $M=M\left(f, r, r_{1}, r_{2}\right)$ such that for every $x \in J_{f}$
there are at most $M$ local symmetries $\sigma$ of $J_{f}$ defined in $B(x, 2 r)$ and such that

$$
\begin{equation*}
r_{1} \leqslant \operatorname{Diam}(\sigma(B(x, r))) \leqslant r_{2} \tag{9}
\end{equation*}
$$

For local biholomorphisms between Julia sets, this yields:
Lemma 5.3. Let $f_{1}$ and $f_{2}$ be rational maps with non-smooth Julia sets. Then, given $r>0$ and $0<c_{1}<c_{2}<1$, there exists $M^{\prime}=M^{\prime}\left(f_{1}, f_{2}, r, c_{1}, c_{2}\right)$ such that for every $x \in J_{1}$ there are at most $M^{\prime}$ local biholomorphisms $\sigma$ defined in $B(x, 2 r)$, mapping $J_{1}$ to $J_{2}$, and such that

$$
\begin{equation*}
c_{1} \leqslant \operatorname{Diam}(\sigma(B(x, r))) \leqslant c_{2} \tag{10}
\end{equation*}
$$

Proof. By the Koebe Distortion Theorem, there exist $r_{1}$ and $r_{2}$ depending only on $r, c_{1}$ and $c_{2}$ such that for any $\sigma$ as in the statement of the lemma

$$
B\left(\sigma(x), 2 r_{1}\right) \subset \sigma(B(x, r)) \subset B\left(\sigma(x), r_{2}\right)
$$

Likewise, there exists $r_{3}=r_{3}\left(r, c_{1}, c_{2}\right)$ such that $\sigma^{-1}\left(B\left(\sigma(x), r_{1}\right)\right) \supset B\left(x, r_{3}\right)$ and finally there exists $r_{4}=r_{4}\left(r, c_{1}, c_{2}\right)$ such that $\sigma\left(B\left(x, r_{3}\right)\right) \supset B\left(\sigma(x), r_{4}\right)$. Now fix such a local biholomorphism $\sigma_{0}$ and let $y=\sigma(x) \in J_{2}$. For any other such $\sigma$, we infer that $\tau=\sigma \circ \sigma_{0}^{-1}$ is a symmetry of $J_{2}$, defined in $B\left(y, 2 r_{1}\right)$, and satisfying

$$
B\left(\tau(y), r_{4}\right) \subset \tau\left(B\left(y, r_{1}\right)\right) \subset B\left(\tau(y), r_{2}\right)
$$

and from Corollary 5.2 we conclude that there are only $M^{\prime}=M^{\prime}\left(f_{2}, r_{1}, r_{4}, 2 r_{2}\right)$ such maps $\tau$, and we are done.

The next result plays a key role in our argument.
Lemma 5.4. Let $f_{1}$ and $f_{2}$ be non-integrable hyperbolic polynomials such that $J_{1}$ and $J_{2}$ are Jordan curves. Let $U$ be a connected open subset intersecting $J_{1}$ and $\sigma$ be a biholomorphism defined on $U$ such that $\sigma\left(U \cap J_{1}\right)=\sigma(U) \cap J_{2}$.

There exists a constant $B$ depending only on $f_{1}, f_{2}$ and $\sigma$ such that the following holds. For any periodic point $p_{1} \in J_{1}$ of period $k_{1}$, there exists a local biholomorphism $\tilde{\sigma}$ defined in a neighborhood of $p_{1}$ such that $\tilde{\sigma}\left(p_{1}\right)=p_{2}$ is periodic under $f_{2}$ and $\tilde{\sigma} \circ f_{1}^{k_{1} b}=f_{2}^{b^{\prime}} \circ \tilde{\sigma}$ for some $b \leqslant B$ and some integer $b^{\prime}$.

Proof. Since $f_{1}$ and $f_{2}$ are hyperbolic, we may suppose that $\left|f_{1}^{\prime}\right|>1$ on $J_{1}$, and $\left|f_{2}^{\prime}\right|>1$ on $J_{2}$. We may also find $r=r\left(f_{1}\right)>0$ such that for any point $p \in J_{1}$ and any $q \in f_{1}^{-n}(p)$ there exists a univalent branch $f_{1,-n}$ of $f_{1}^{n}$ defined in $B(p, 2 r)$ and mapping $p$ to $q$ (we fix $\rho>0$ such that the analogous property holds for $f_{2}$ ).

For $p_{1}$ as in the statement, choose any integer $N \in \mathbb{N}^{*}$ and a univalent inverse branch $f_{1,-N}$ of $f_{1}^{N}$ defined on $B\left(p_{1}, 2 r\right)$ with values in $U$. For notational ease we denote by $f_{1}^{-k_{1} n}$ the branch of $\left(f_{1}^{k_{1} n}\right)^{-1}$ fixing $p_{1}$.

Write $\kappa=\left|\left(f^{k_{1}}\right)^{\prime}\left(p_{1}\right)\right|>1$ and for each $n$ consider the map $F_{n}:=\sigma \circ f_{1,-N} \circ f_{1}^{-k_{1} n}$. This map is defined on $B\left(p_{1}, 2 r\right)$, and by the Koebe Distortion Theorem, it satisfies

$$
B\left(F_{n}\left(p_{1}\right), c_{1}^{\prime} \kappa^{n}\right) \subset F_{n}\left(B\left(p_{1}, r\right)\right) \subset B\left(F_{n}\left(p_{1}\right), c_{2}^{\prime} \kappa^{n}\right)
$$

for some uniform constants $c_{1}^{\prime}<c_{2}^{\prime}$ depending only on $f_{1}$ and $\sigma$.

For any $n$, let $m=m_{n}$ be the largest integer such that the diameter of $f_{2}^{m}\left(F_{n}\left(B\left(p_{1}, r\right)\right)\right.$ is bounded by $\rho$. Note that $n \mapsto m_{n}$ is non-decreasing. Write $q_{m}=f_{2}^{m}\left(F_{n}\left(p_{1}\right)\right)$. Let $f_{2,-m}$ be the univalent branch defined on $B\left(q_{m}, 2 \rho\right)$ and mapping $q_{m}$ to $F_{n}\left(p_{1}\right)$. Set $L=\sup _{J_{2}}\left|f_{2}^{\prime}\right|$. By Koebe distortion, we get

$$
f_{2,-m}\left(B\left(q_{m}, \frac{\rho}{L}\right)\right) \supset B\left(F_{n}\left(p_{1}\right), \frac{\rho}{4 L\left|\left(f_{2}^{m}\right)^{\prime}\left(F_{n}\left(p_{1}\right)\right)\right|}\right)
$$

Now observe that by maximality of $m, f_{2}^{m}\left(F_{n}\left(B\left(p_{1}, r\right)\right)\right)$ is not included in $B\left(q_{m}, \rho / L\right)$ so that $\frac{\rho}{4 L\left|\left(f_{2}^{m}\right)^{\prime}\left(F_{n}\left(p_{1}\right)\right)\right|} \leqslant c_{2}^{\prime} \kappa^{n}$. Thus by applying Koebe to $f_{2}^{m}$ on the disk $B\left(F_{n}\left(p_{1}\right), c_{1}^{\prime} \kappa^{n}\right)$, we conclude that $\tau_{n}:=f_{2}^{m} \circ F_{n}$ is a sequence of univalent maps defined on $B\left(p_{1}, 2 r\right)$ such that

$$
\frac{c_{1}^{\prime}}{16 L c_{2}^{\prime}} \rho \leqslant \operatorname{Diam}\left(\tau_{n}\left(B\left(p_{1}, r\right)\right)\right) \leqslant \rho
$$

and $\sigma_{n}\left(J_{1} \cap B\left(p_{1}, 2 r\right)\right) \subset J_{2}$.
Lemma 5.3 yields an integer $B$ depending only of $f_{1}, f_{2}$ and $\sigma$, and a pair of integers $0<n<n^{\prime} \leqslant B$ such that $\tau_{n}=\tau_{n^{\prime}}$. Expanding this equality gives

$$
f_{2}^{m_{n}} \circ \sigma \circ f_{1,-N} \circ f_{1}^{-k_{1} n}=f_{2}^{m_{n^{\prime}}} \circ \sigma \circ f_{1,-N} \circ f_{1}^{-k_{1} n^{\prime}}
$$

that is,

$$
\tilde{\sigma}=f_{2}^{m_{n^{\prime}}-m_{n}} \circ \tilde{\sigma} \circ f_{1}^{-k_{1}\left(n^{\prime}-n\right)}, \text { where } \tilde{\sigma}=f_{2}^{m_{n}} \circ \sigma \circ f_{1,-N} \circ f_{1}^{-k_{1} n}
$$

and the result follows.
Remark 5.5. By pushing the argument further it is possible to prove that $b^{\prime} \leqslant k_{1} B^{\prime}$ for some uniform $B^{\prime}$.

Step 3: multipliers and smooth rigidity for expanding maps on the circle.
In this paragraph we prove some rigidity results for Blaschke products based on periodic points multipliers. If $p$ is a periodic point of period $n$, its Lyapunov exponent is by definition $\frac{1}{n} \log \left|\left(f^{n}\right)^{\prime}(p)\right|$. The following lemma is presumably well-known. We provide a proof for the convenience of the reader.

Lemma 5.6. Let $g$ be a uniformly hyperbolic Blaschke product of degree d such that $J_{g}=\partial \mathbb{D}$. Assume that the Lyapunov exponents of periodic points of $g$ take only one value $\log e$. Then $e=d$ and $g$ is conjugate to $z \mapsto z^{d}$ by a Möbius transformation.

Proof. Since $g$ is uniformly expanding, any ergodic invariant measure has a positive Lyapunov exponent. By approximating it by periodic orbits we infer that this Lyapunov exponent is equal to $\log e$. Applying the dimension formula (see e.g. [35, Theorem 11.4.1]) to the unique measure of maximal entropy $\mu_{g}$, we get that

$$
1 \geqslant \operatorname{HD}\left(\mu_{g}\right)=\frac{h_{\mu_{g}}(g)}{\chi_{\mu_{g}}(g)}=\frac{\log d}{\log e}, \text { hence } \log e \geqslant \log d
$$

Likewise, applying it to the unique smooth invariant measure $\nu$, we get

$$
1=\mathrm{HD}(\nu)=\frac{h_{\nu}(g)}{\chi_{\nu}(g)} \leqslant \frac{\log d}{\log e}, \text { hence } \log e \leqslant \log d
$$

From this we conclude that $e=d$ and that $\nu$ is the measure of maximal entropy. Then, if we let $h$ be the conjugacy between $g$ and $M_{d}, h \circ g=M_{d} \circ h$, by uniqueness of the maximal entropy measure, we infer that $h_{*} \nu=\operatorname{Leb}_{\partial \mathbb{D}}$, so $h$ is smooth. Finally, by [38, Theorem 4], $\left.g\right|_{\partial \mathbb{D}}$ is conjugate to $z^{d}$ by a Möbius transformation, and the proof is complete.

Lemma 5.7. For any uniformly expanding $C^{1}$-map of the circle, the closure of the set of Lyapunov exponents of periodic orbits is an interval.

From the two previous lemmas we immediately get:
Corollary 5.8. Suppose $g$ is a uniformly hyperbolic Blaschke product of degree $d$ whose set of Lyapunov exponents of periodic orbits is discrete. Then $g$ is Möbius conjugate to $z \mapsto z^{d}$.

Proof of Lemma 5.7. It is enough to show that if there are two periodic orbits $x_{1}$ of period $n_{1}$, and $x_{2}$ of period $n_{2}$, of respective Lyapunov exponents $\chi_{1}$ and $\chi_{2}$, then there is a periodic orbit whose Lyapunov exponent is approximately $\frac{1}{2}\left(\chi_{1}+\chi_{2}\right)$. This follows from the periodic specification property, which holds for any expanding map of the circle.

Choose $\eta>0$, and pick $\varepsilon>0$ so small that

$$
\chi_{1}-\eta \leqslant \frac{1}{n_{1}} \log \left|\left(f^{n_{1}}\right)^{\prime}(x)\right| \leqslant \chi_{1}+\eta
$$

(resp. $\left.\chi_{2}-\eta \leqslant \frac{1}{n_{2}} \log \left|\left(f^{n_{2}}\right)^{\prime}(x)\right| \leqslant \chi_{2}+\eta\right)$ for any $x \in \partial \mathbb{D}$ such that $d\left(x, x_{1}\right) \leqslant \varepsilon$ (resp. $\left.d\left(x, x_{2}\right) \leqslant \varepsilon\right)$.

By the periodic specification property, there exists an integer $M \geqslant 1$ such that for any $q, q^{\prime} \geqslant 1$ there is a periodic orbit $x, f(x), \cdots, f^{q n_{1}+q^{\prime} n_{2}+2 M}(x)=x$ such that $d\left(f^{k}(x), f^{k}\left(x_{1}\right)\right) \leqslant \varepsilon$ for all $0 \leqslant k \leqslant q n_{1}-1$, and $d\left(f^{k+q n_{1}+M}(x), f^{k}\left(x_{2}\right)\right) \leqslant \varepsilon$ for all $0 \leqslant k \leqslant q^{\prime} n_{2}-1$.

The Lyapunov exponent of $x$ then satisfies:

$$
\begin{aligned}
& \chi \leqslant \frac{1}{q n_{1}+q^{\prime} n_{2}+2 M}\left(n_{1} q\left(\chi_{1}+\eta\right)+n_{2} q^{\prime}\left(\chi_{2}+\eta\right)+(2 M) \log \left(\sup \left|f^{\prime}\right|\right)\right) \\
& \chi \geqslant \frac{1}{q n_{1}+q^{\prime} n_{2}+2 M}\left(n_{1} q\left(\chi_{1}-\eta\right)+n_{2} q^{\prime}\left(\chi_{2}-\eta\right)+(2 M) \log \left(\inf \left|f^{\prime}\right|\right)\right)
\end{aligned}
$$

So if we choose $q$ and $q^{\prime}$ very large compared to $M$ and satisfying $q n_{1}=q^{\prime} n_{2}$, we conclude that $\chi$ is very close to $\frac{1}{2}\left(\chi_{1}+\chi_{2}\right)$, as announced.

Step 4: uniformization and conclusion.
We return to the original situation. Let $f_{1}$ and $f_{2}$ be non-integrable hyperbolic polynomials such that $J_{1}$ and $J_{2}$ are Jordan curves. Let $U$ be a connected open subset intersecting $J_{1}$ and $\sigma$ be a biholomorphism defined on $U$ such that $\sigma\left(U \cap J_{1}\right)=\sigma(U) \cap J_{2}$ and $\sigma\left(K_{1} \cap U\right) \subset K_{2}^{\complement}$.

Let $\phi_{1}: \mathbb{D} \rightarrow \operatorname{Int}\left(K_{1}\right)$ be a uniformization with $\phi_{1}(0)$ the attracting fixed point of $f_{1}$, and $\phi_{2}: \mathbb{D} \rightarrow \mathbb{P}^{1} \backslash K_{2}$ be a uniformization with $\phi_{2}(0)=\infty$. Since $J_{1}$ and $J_{2}$ are Jordan curves, both $\phi_{1}$ and $\phi_{2}$ extend to respective homeomorphisms $\partial \mathbb{D} \rightarrow J_{1}$ and $\partial \mathbb{D} \rightarrow J_{2}$. Actually, since $J_{1}$ and $J_{2}$ are quasi-circles, these homeomorphisms are biHölder (indeed they are quasi-symmetric, see [32, Chapter 5]). Let $g_{1}=\phi_{1}^{-1} \circ f_{1} \circ \phi_{1}$
(resp. $g_{2}=\phi_{2}^{-1} \circ f_{2} \circ \phi_{2}$ ). Up to conjugating by a rotation, $g_{2}(z)=M_{d_{2}}(z):=z^{d_{2}}$. Observe that $g_{1}(z)$ extends to a Blaschke product of degree $d_{1}$. Indeed $g_{1}$ extends by Schwarz reflexion to a rational map which satisfies $g_{1}(\mathbb{D})=\mathbb{D}$ and $g_{1}\left(\mathbb{D}^{\complement}\right) \subset \mathbb{D}^{\complement}$, hence $\mathbb{D}$ is totally invariant so $g_{1}$ is a Blaschke product of the same degree as $f_{1}$, which is uniformly hyperbolic on $\partial \mathbb{D}$ because no critical orbit of $g_{1}$ approaches $\partial \mathbb{D}$.

We claim that the set of Lyapunov exponents of periodic orbits of $g_{1}$ is discrete. Taking this claim for granted, by Corollary 5.8, we obtain that $g_{1}$ is Möbius conjugate to $M_{d_{1}}$ thus $f_{1}$ admits a totally invariant fixed point in $K_{1}$, so it is integrable, which is a contradiction.

To justify our claim, we proceed as follows. Pick any periodic point $q_{1} \in \partial \mathbb{D}$ for $g_{1}$ of period $k_{1}$. Then $p_{1}:=\phi_{1}^{-1}\left(q_{1}\right)$ is $f_{1}$-periodic of the same period, and Lemma 5.4 implies the existence of a local biholomorphism $\tilde{\sigma}$ sending $p_{1}$ to a periodic point $p_{2}$ for $f_{2}$ such that $\tilde{\sigma} \circ f_{1}^{k_{1} b}=f_{2}^{b^{\prime}} \circ \tilde{\sigma}$ for some $b \leqslant B=B\left(f_{1}, f_{2}, \sigma\right)$ and some $b^{\prime} \geqslant 1$. By the Schwarz reflection principle, the map $\tau:=\phi_{2}^{-1} \circ \tilde{\sigma} \circ \phi_{1}$ is defined in a neighborhood of $q_{1}$ and satisfies $\tau \circ g_{1}^{k_{1} b}=g_{1}^{b^{\prime}} \circ \tau=M_{d_{2}}^{b^{\prime}} \circ \tau$. Computing derivatives, this implies that the Lyapunov exponent of $g_{1}$ at $q_{1}$ belongs to $\bigcup_{1 \leqslant b \leqslant B} \frac{\mathbb{N}}{b} \log d_{2}$ which is a discrete set.

The proof is complete.

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