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A 3D Lead Iodide Hybrid Based on a 2D Perovskite Subnetwork

Maroua Ben Haj Salah¹, Justine Tessier¹, Nicolas Mercier^{1,*}, Magali Allain¹, Antonin Leblanc¹, Xiaoyang Che², Claudine Katan² and Mikael Kepenekian^{2,*}

¹ MOLTECH-Anjou, UMR-CNRS 6200, Université d'Angers, 2 Bd Lavoisier, 49045 Angers, France; maroua.benhajsalah@univ-angers.fr (M.B.H.S.); justine72.tessier@gmail.com (J.T.); magali.allain@univ-angers.fr (M.A.); antonin.leblanc49@gmail.com (A.L.)

² Univ Rennes, ENSCR, INSA Rennes, CNRS, ISCR (Institut des Sciences Chimiques de Rennes), UMR 6226, F-35000 Rennes, France; xiaoyang.che@gmail.com (X.C.); claudine.katan@univ-rennes1.fr (C.K.)

* Correspondence: nicolas.mercier@univ-angers.fr (N.M.); mikael.kepenekian@univ-rennes1.fr (M.K.)

Abstract: Lead halide perovskites have emerged as promising materials for various optoelectronic applications. For photovoltaics, the reference compound is the 3D perovskite (MA)PbI₃ (MA⁺ = methylammonium). However, this material suffers from instabilities towards humidity or light. This makes the search of new stable 3D lead halide materials very relevant. A strategy is the use of intermediate size cations instead of MA, which are not suitable to form the 3D ABX₃ perovskites or 2D perovskites. Here, we report on a novel 3D metal halide hybrid material based on the intermediate size cation hydroxypropylammonium (HPA⁺), (HPA)₆(MA)Pb₅I₁₇. We will see that extending the carbon chain length from two CH₂ units (in the hydroxyethylammonium cation, HEA⁺) to three (HPA⁺) precludes the formation of a perovskite network as found in the lead and iodide deficient perovskite (HEA,MA)_{1+x}Pb_xI_{3-x}. In (HPA)₆(MA)Pb₅I₁₇ the 3D lead halide network results from a 2D perovskite subnetworks linked by a PbI₆ octahedra sharing its faces. DFT calculations confirm the direct band gap and reveal the peculiar band structure of this 3D network. On one hand the valence band has a 1D nature involving the p orbitals of the halide. On the other, the conduction band possesses a clear 2D character involving hybridization between the p orbitals of the metal and the halide.

Keywords: 3D network; halide perovskite; electronic structure; lead iodide hybrid



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1. Introduction

Halide perovskites became superstar semiconductors in recent years, especially for photovoltaic applications [1–4]. The reference compound is (MA)PbI₃ (MA⁺ = methylammonium) which adopts an ABX₃ perovskite structure. This structure can be described by corner-sharing BX₆ octahedra which extend along the three directions of space (3D perovskite), the A⁺ cation being located in the middle of space defined by eight adjacent corner-sharing octahedra. The electronic structure has also a 3D character, meaning that charge transport can occur effectively in the whole space which is key for an efficient photovoltaic effect [1]. However, (MA)PbI₃ suffers from instabilities towards humidity or light, mainly due to the volatile MA⁺ cation [5]. Different strategies have been considered to improve the stability of such materials. One is to partially or completely substitute MA⁺ by other cations. This has led to the multiple cation family (A, A', A'')PbI₃, the A site of the perovskite being occupied by three types of cations such as Cs⁺, Rb⁺, MA⁺, FA⁺ [6–8]. Another strategy is the use of layered perovskites instead of the 3D perovskite as active material [9–11]. The general formula of the most familiar family is (RNH₃)₂(A)_{n-1}Pb_nX_{3n+1}, also known as Ruddlesden–Popper (RP) phases [12,13]. In such structures, long-size organic cations RNH₃⁺ separate layers (n = 1) or multilayers (n > 1) of perovskite type (PbX₆ corner-sharing mode). These materials appear much more stable than their 3D congeners due to the hydrophobic effect of RNH₃⁺ cations [9]. However, while in photovoltaic devices, the charge transport takes place along the direction perpendicular to the substrate,

these layered compounds naturally grow with perovskite layers parallel to the substrate which is therefore very unfavorable for efficient solar cells. Noteworthy, a hot-casting technique has been shown to favor the growth of perovskite layers perpendicular to the substrate [14], but this hinders further upscaling and large area thin film fabrication.

In this context, the quest for original 3D metal halide materials is of great interest. If a few small-size cations can form the 3D perovskite network, and large-size organic cation afford 2D layered structures, the case of intermediate size cations remains almost unexplored. Too large to form 3D ABX_3 perovskites and too small to form 2D perovskites, they can lead to unique structure types [15–20]. Thus, the hydroxyethylammonium (HEA^+) cation leads to maintain a corner-sharing 3D connectivity in which some $(PbI)^+$ units are replaced by HEA^+ in materials called *d*-HP (lead and iodide deficient-perovskites), $(HEA, MA)_{1+x}Pb_xI_{3-x}$ which appear more stable than the reference compounds $(MA)PbI_3$ or $(Cs,FA)PbI_3$ [15,16]. Dications such as ethylenediammonium or propylenediammonium also afford such compounds called *hollow* perovskites (disordered structures) [17,18]. In most cases, such intermediate size cations lead to 2D (or 1D) networks exhibiting two kinds of connectivity among edge-, face-, and corner-sharing [19,20], which consequently are not perovskite networks [21]. 3D metal halide networks are rarer. In the tin iodide system, a series of $ASnI_3$ (A^+ = guanidinium, ethylammonium) compounds has been found, the 3D iodostannate network being based on corner- and face-sharing octahedra [22]. More recently, a series of pyrazinium-based dications (A'^{2+}) has afforded the $(A'')Pb_2X_6$ compounds whose structures are based on 3D networks of corner- and edge-sharing lead halide octahedra [23].

Here we report on a new metal halide material that has a 3D network of interconnected PbI_6 octahedra, using the hydroxypropylammonium (HPA^+) cation, $(HPA)_6(MA)Pb_5I_{17}$. It is demonstrated that extending the carbon chain length from two CH_2 units (HEA^+) to three (HPA^+) precludes the formation of a perovskite network as found in HEA -based *d*-HP [15,16]. In $(HPA)_6(MA)Pb_5I_{17}$ the 3D lead halide framework results from 2D perovskite subnetworks linked by PbI_6 face-sharing octahedra. The properties of the newly synthesized material are further investigated based on DFT electronic structure calculations and optical absorption measurements revealing a direct band gap and reduced electronic dimensionality.

2. Materials and Methods

PbI_2 (99%), methylamine solution 40% in water, 3-amino-1-propanol, hydroiodic acid 57% in water, were purchased from commercial sources (Sigma-Aldrich) and were used as received without any purification. The organic salts 3-ammonium-1-propanol iodide (HPA , I) and methylammonium iodide (MA , I) were first prepared from HI solution containing organic molecules. The iodoplumbate salt $(HPA)_6(MA)Pb_5I_{17}$ was then prepared by a liquid–gas slow diffusion technique. Details of syntheses are given in the supplementary materials. The purity of the obtained compound was checked by both PXRD (D8 diffractometer from Bruker (Karlsruhe, Germany), Supplementary Materials, Figure S1, all experimental lines well fit with calculated ones from single-crystal X-ray data) and 1H NMR (Bruker ultrashieldTM 300 (Germany), Figure S2, the expected HPA/MA ratio of 6 is found).

X-ray single-crystal diffraction data of $(HPA)_6(MA)Pb_5I_{17}$ were collected at 149K on a Rigaku Oxford Diffraction (Rigaku Europe, Neu-Isenburg, Germany) SuperNova diffractometer. The structure was solved by direct methods, expanded and refined on F^2 by full matrix least-squares techniques using SHELX programs (SHELXS 2013/1 and SHELXL 2013/4), leading to R1 factor of 0.0387. A summary of crystallographic data and refinement results for this structure is given in the Supplementary Materials (Table S1). CCDC-2122380.

First-principles calculations are based on density functional theory (DFT) as implemented in the SIESTA package (v5.0.0-alpha) [24,25]. The nonlocal van der Waals density functional of Dion et al. corrected by Cooper (C09) [26,27] was used for full geometry opti-

mizations, i.e., of atomic positions and cell vectors, while the centrosymmetry is preserved. The electronic structure of the relaxed geometry is described including spin-orbit coupling through the on-site approximation as proposed by Fernández-Seivane et al. [28]. To prevent conflicts between the on-site treatment and the nonlocality of C09, single point calculations were conducted with the revPBE functional on which C09 is based [29]. Core electrons are described with Troullier–Martins pseudopotentials [30], while valence wavefunctions are developed over double- ζ polarized basis set of finite-range numerical pseudoatomic orbitals [31]. In all cases, an energy cutoff of 150 Ry for real-space mesh size was used. Electronic structures were converged using k-point samplings of the Brillouin zone of $5 \times 3 \times 3$.

3. Results

3.1. Crystal Structure, Thermal Stability, and Optical Properties

A slow diffusion of ethanol into an acidic solution (HI) containing hydroxypropylamine, methylamine, and lead iodide (see Section 2) affords needle like crystals of $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$. Figure 1 shows a general view of the structure. The compound crystallizes in the triclinic $P-1$ space group, the independent space containing 3 HPA^+ molecules, half MA^+ , 2 Pb^{2+} and 8 I^- in general positions, and 1 Pb^{2+} and 1 I^- located on symmetry centers. The metal iodide network consisting of connected PbI_6 octahedra has a 3D character. It can be described from a 2D perovskite subnetwork (purple octahedra, Figure 1) connected to each other by PbI_6 face-sharing octahedra (grey octahedra, Figure 1). The 2D perovskite subnetwork of corner-sharing octahedra is isolated on Figure 2a, the different views, along a , b , and c , highlighting the corner-sharing octahedra mode. The 2×2 octahedra building unit centered in the middle of the unit cell (when viewed along the a direction, Figures 1 and 2a) can be considered as a part of the $n = 2$ layered perovskite Pb_2I_7 . The repetition of these units along a leads to a 1D subnetwork Pb_4I_{18} . A similar situation is found in the structure of $(\text{PA})_5\text{Pb}_5\text{I}_{18}$ (PA^+ : propylammonium) compound [19]. However, while these 1D chains are connected through PbI_6 face-sharing octahedra in $(\text{PA})_5\text{Pb}_5\text{I}_{18}$ leading to an overall 2D lead iodide network in this structure, these 1D chains are connected to each other through corner-sharing mode leading to a 2D step-like Pb_4I_{17} ($\text{Pb}_4\text{I}_{16}\text{I}_{2/2}$, two shared iodides) perovskite network extending along the a and b directions of the unit cell. As already mentioned, these 2D perovskite subnetworks are linked by Pb^{2+} cations located on symmetry centers, with corresponding octahedra sharing faces with their two neighboring octahedra, leading finally to the overall 3D Pb_5I_{17} network.

While MA^+ cations are located near the shared iodides connecting adjacent 2×2 octahedra units within the 2D perovskite subnetwork, the HPA^+ are clearly located in channels defined by this open framework. Figure 2b shows a partial view of the structure focusing on one channel whose section is defined by 10 adjacent lead iodide octahedra, containing 6 HPA^+ cations. We observe 3 kinds of HPA^+ cations. They are numbered 1, 2, and 3 on Figure 2b. $\text{HPA}(1)$ and $\text{HPA}(3)$ at the surface of cavities clearly interact with iodides through both $\text{NH} \dots \text{I}$ and $\text{OH} \dots \text{I}$ interactions (see Figure 2b, such interactions with $\text{H} \dots \text{I}$ distances $< 3.05 \text{ \AA}$ are drawn as dashed lines). In contrast, $\text{HPA}(2)$ cations located at the center of channels do not interact with iodides through sizeable hydrogen bonding ($< 3.05 \text{ \AA}$). Meanwhile, two adjacent $\text{HPA}(2)$ molecules mainly interact together through strong $\text{NH} \dots \text{O}$ interactions ($d(\text{H} \dots \text{O}) = 2 \text{ \AA}$). The peculiar curled conformation of HPA^+ , related to the hydrogen bond ability of both the NH_3^+ and the OH ends, certainly hinders formation of 2D perovskites, while favoring here a structure having a 3D metal halide network, as already observed for other intermediate size cations [23].

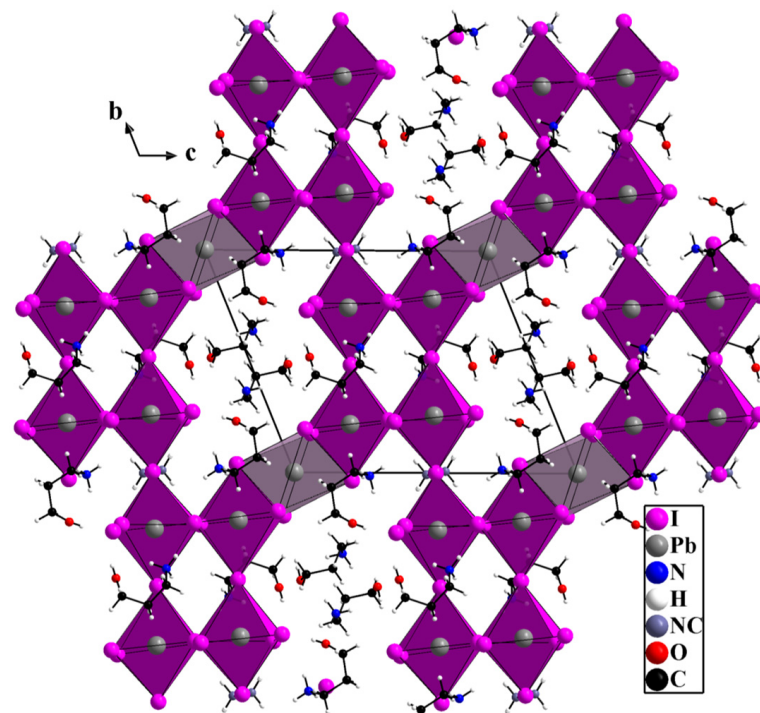


Figure 1. General view of the structure of $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$ showing the 3D lead iodide network resulting from 2D perovskite subnetwork (purple octahedra) linked together by PbI_6 face-sharing octahedra (grey octahedra).

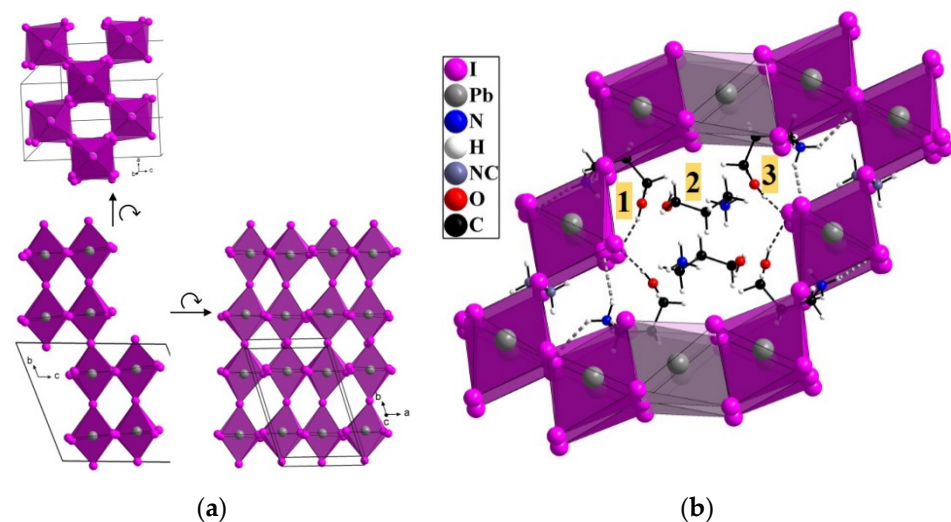


Figure 2. Structure of $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$. (a) Partial view showing the 2D perovskite subnetwork along a , b , and c highlighting the corner-sharing connectivity; (b) partial view showing organic cations HPA^+ and MA^+ in channels defined by corner- and face-sharing PbI_6 octahedra, highlighting the three types of HPA^+ cations (numbered 1, 2, and 3) and hydrogen bonding ($\text{OH} \dots \text{I}$ and $\text{NH} \dots \text{I}$ contacts $< 3.05 \text{ \AA}$, dashed lines).

We further explored the stability of $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$ upon heating. TGA experiment shows that the compound is stable up to $200 \text{ }^\circ\text{C}$ and decomposes into two steps (Figure S3). Its stability in ambient air appears improved as compared to the well-known 3D MAPbI_3 , since PXRD of a sample left in the air during two weeks is similar to the pristine one (Figure S4). Optical properties were investigated on $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$ obtained as a crystallized powder. The absorption spectrum shown in Figure 3a reveals a broad absorption band in the first part of the visible region (400–550 nm range). The corresponding Tauc plot

(Figure 3b) leads to a direct optical band gap of 2.33 eV. The $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$ crystallized powder also exhibits a weak broad band photoluminescence emission centered at 530 nm (Figure S5).

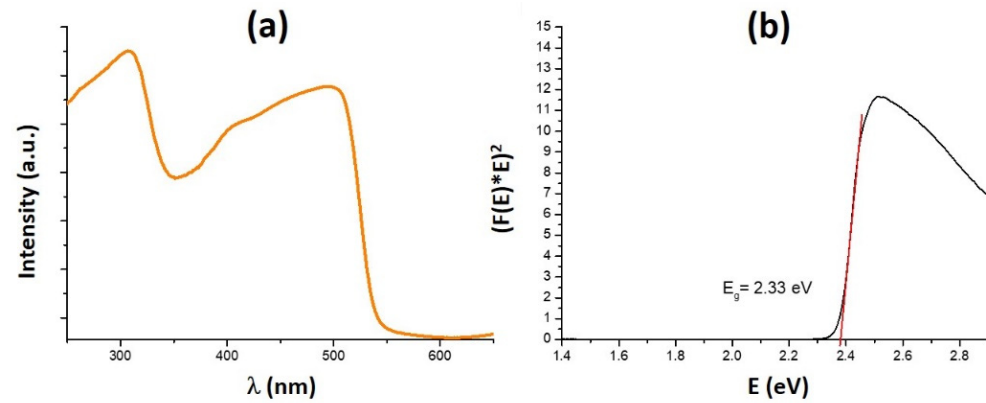


Figure 3. Absorption spectrum of a crystallized powder of $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$ (a) and corresponding Tauc plot (black curve) assuming a direct band gap ((b), red line: fitted line giving E_g).

3.2. Electronic Structure

Insights into the electronic structure of $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$ is provided by first-principles calculations based on density functional theory (DFT). Figure 4a shows the band structure computed at the GGA level (see Section 2 for computational details) that provides reliable description of the electronic features despite the well-known band gap underestimation. The band structure presents a direct band gap of 1.60 eV at T. We verify that the valence band maximum (VBM) is governed by I(5*p*) orbitals, while the conduction band minimum (CBM) is due to Pb(6*p*) and I(5*p*) orbitals (Figure S7), i.e., no cation, HPA or MA, state is found near the band edges. Interestingly, the band structure does not adopt the typical shape found for 3D, 2D, nor 1D halide perovskites. Indeed, the lower energy conduction band is flat in one direction, along *c*, and strong dispersion is observed along the two others. Such features are typically found in 2D perovskites, where flat bands correspond to the stacking direction. On the other hand, the top valence band presents sizeable dispersion only along *a*, a characteristic expected for 1D systems.

The negligible dispersion along *c*, observed both in the conduction and valence band can be assigned to the presence of the 2D perovskite subnetwork (vide supra) separated by the large cation (HPA) and the face-sharing octahedra that efficiently block the electronic coupling between layers as it has been described in previous cases [19,32]. The partial charge density plotted at the CBM (Figure 4b) adopts the form typically observed in layered perovskites with a large contribution from Pb(6*p*) and I(5*p*) orbitals along *a* and *b*, i.e., the two directions with corner-sharing octahedra. It leads to the important bandwidth following the line Z-T-R (Figure 4a). The case is different with partial charge density taken at the VBM (Figure 4c,d) that illustrates the strong interaction between iodide atoms along *a* responsible for the dispersion observed along the T-R line. As opposed to the typical shape of valence band in halide perovskites, almost no contribution arises from Pb(6*s*) orbitals. Instead, I(5*p_x*) orbitals directly overlap thanks to the zigzag connectivity of the octahedra in the (*a*, *c*) plane. The direct overlap between I(5*p*) orbitals along *a* leads to greater interaction and therefore to a higher energy of this anti-bonding combination. Since, the coupling in the *b* direction is of the classical I(5*p*)-Pb(6*s*)-I(5*p*) form, it lies lower in energy. As a result, the VBM is formed by couplings only along the *a* direction, giving a 1D-like character to the valence states.

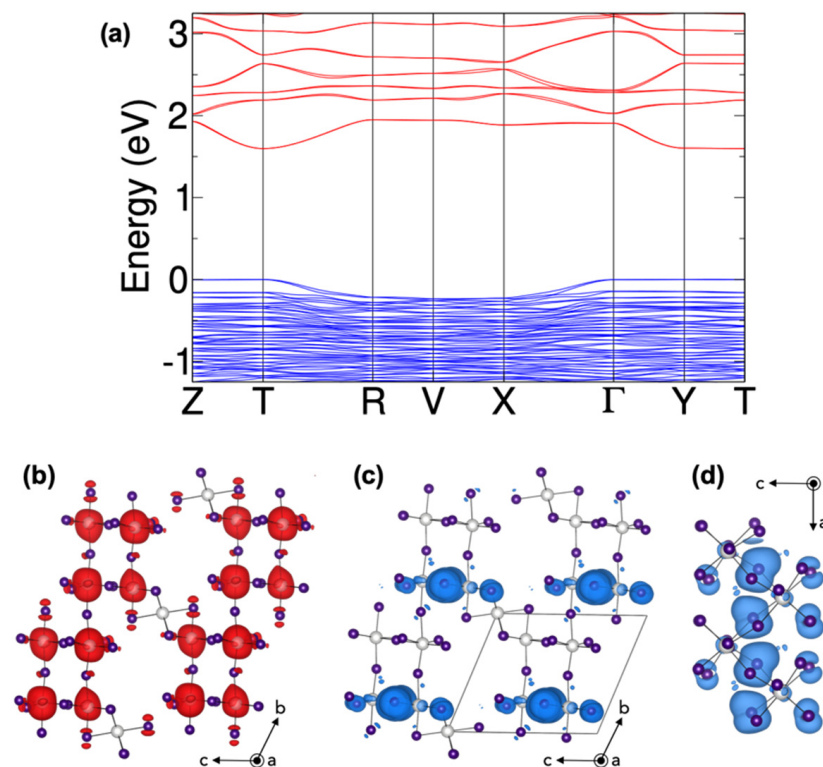


Figure 4. (a) Computed band structure (blues lines: valence band, red lines: conduction band) of $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$. The system presents a direct band gap at T with $E_G = 1.60$ eV. (b,c) Partial charge densities computed at the conduction band minimum and the valence band maximum, respectively. (d) Same as (c) from a different viewpoint. Organic cations are not depicted for clarity (purple atoms: I, white atoms: Pb).

Even though not quantitative, it is worth noticing that the here-computed bandgap (1.60 eV) is greater than the ones computed at the same level of theory not only for 3D halide perovskites, but also for typical 2D ones (1.1–1.3 eV). Still, it remains smaller than bandgaps of related 1D-like compounds (2.0 eV) [33]. This is indicative that the electronic dimensionality of the 3D $(\text{HPA})_6(\text{MA})\text{Pb}_5\text{I}_{17}$ structure is intermediate between 1D and 2D.

4. Conclusions

Using the hydroxypropylammonium (HPA^+) intermediate size cation, a new metal halide organic–inorganic material was synthesized. The X-ray structure determination reveals a rare and unprecedented 3D inorganic network. The structure consists of a 2D perovskite subnetwork of corner-sharing PbI_6 octahedra that are in turn connected to each other by PbI_6 octahedra sharing their faces. This compound absorbs in the 400–550 nm visible range leading to an optical band gap of 2.33 eV. DFT calculations further reveal that this 3D structural network has a peculiar band structure that shows a direct band gap but reduced electronic dimensionality. On one hand, the valence band has a 1D nature and is mainly composed of p orbitals of the halide. On the other hand, the conduction band possesses a clear 2D character with contributions of p orbitals of both the metal and the halide.

Supplementary Materials: The following are available online at <https://www.mdpi.com/article/10.3390/cryst11121570/s1>, Figure S1: Calculated and experimental PXRD; Figure S2: ^1H NMR in deuterated DMSO; Figure S3: TGA; Figure S4: PXRD of aged sample; Figure S5: Photoluminescence; Figure S6: Brillouin zone for $P-1$ space group; Figure S7: Density of states; Table S1: Summary of crystal data and structure refinement.

Author Contributions: Conceptualization, N.M. and C.K.; methodology, M.K. and A.L.; investigation, M.B.H.S., J.T., X.C. and M.A.; writing, N.M., C.K. and M.K. All authors have read and agreed to the published version of the manuscript.

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