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WELL-POSEDNESS AND SINGULARITY FORMATION FOR THE KOLMOGOROV TWO-EQUATION MODEL OF TURBULENCE IN 1-D

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Abstract

We study the Kolomogorov two-equation model of turbulence in one space dimension. Two are the main results of the paper. First of all, we establish a local well-posedness theory in Sobolev spaces even in the case of vanishing mean turbulent kinetic energy. Then, we show that, in general, those solutions must blow up in finite time.

To the best of our knowledge, these results are the first establishing the well-posedness of the system for vanishing initial data and the occurrence of finite time singularities for the model under study.

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Keywords: Kolmogorov two-equation model of turbulence; vanishing turbulent kinetic energy; local well-posedness; finite time blow-up.

1 Introduction

In 1941, in a series of papers Kolmogorov and Obukhov laid the foundations of the celebrated mathematical theory of turbulence which now goes under the name of *Komogorov's 1941 theory*. We refer *e.g.* to Chapter 6 of [10] and to Chapters 4 and 14 of [4] for an account of that theory, as well as to references quoted therein.

As an outcome of that theory, in [7] (see the Appendix of [16] for an English translation) Kolmogorov formulated the following system of partial differential equations to describe a fully developed homogeneous isotropic turbulence:

$$(1) \quad \begin{cases} \partial_t u + (u \cdot \nabla)u + \nabla \pi - \nu \operatorname{div} \left(\frac{k}{\omega} \mathbb{D}u \right) = f \\ \partial_t \omega + u \cdot \nabla \omega - \alpha_1 \operatorname{div} \left(\frac{k}{\omega} \nabla \omega \right) = -\alpha_2 \omega^2 \\ \partial_t k + u \cdot \nabla k - \alpha_3 \operatorname{div} \left(\frac{k}{\omega} \nabla k \right) = -k \omega + \alpha_4 \frac{k}{\omega} |\mathbb{D}u|^2 \\ \operatorname{div} u = 0. \end{cases}$$

In the previous system, the three unknowns u , ω and k are functions of time and space variables $(t, x) \in \mathbb{R}_+ \times \Omega$, where Ω is a domain in \mathbb{R}^d , with $d = 2, 3$. The vector field $u \in \mathbb{R}^d$ represents the average velocity of the fluid: it is the time-averaged velocity in the original paper by Kolmogorov (see again [16]), but it is interpreted as an ensemble average in the modern approach (see [14], [10]). The velocity field u is assumed to be incompressible, whence the last equation in (1) and the presence of the lagrangian multiplier $\nabla\pi$ in the first relation of the system, where π is the mean pressure of the fluid. On the other hand, ω and k are positive scalar functions, representing respectively the mean frequency of the turbulent fluctuations and the average turbulent kinetic energy (called “pulsation velocity” in [16]). The vector field $f \in \mathbb{R}^d$ is a known (average) external force acting on the fluid, which we will set soon equal to 0. Finally, all the parameters $\nu, \alpha_1 \dots \alpha_4$ are physical adimensional parameters which are strictly positive numbers: we have

$$\nu > 0, \quad \alpha_1, \alpha_2, \alpha_3, \alpha_4 > 0.$$

In general, their values depend on the characteristics of the flow; however, for some of them Kolmogorov is able to specify the precise numerical value (we refer once more to [7] and [16]).

From a mathematical point of view, system (1) is a parabolic system, with viscosity (diffusion coefficient) proportional to k/ω . As a matter of fact, the symbol \mathbb{D} appearing in the first and third equations stands for the symmetric part of the gradient of u :

$$\mathbb{D}u := \frac{1}{2} (Du + \nabla u),$$

where we have denoted by Du the Jacobian matrix of u and by ∇u its transpose matrix. Notice that, despite the presence of the two damping terms $-\alpha_2\omega^2$ and $-\omega k$ in the second and third equations respectively, the mean motion is feeding energy from the larger scales to the small scales, namely to k , through the α_4 -term in the third equation.

To begin with, let us spend a few words on the derivation of the previous model.

Derivation of the model

As well explained in [10] (see especially Chapter 4) and [14] (see Chapter 3 therein), for instance, the common approach to the turbulence theory starts by decomposing the incompressible velocity $u = u(t, x)$ into its mean part $\langle u \rangle$ and fluctuating part \tilde{u} : namely, in the turbulent regime we write $u = \langle u \rangle + \tilde{u}$. Notice that the new velocity fields $\langle u \rangle$ and \tilde{u} are still incompressible. Injecting the decomposition $u = \langle u \rangle + \tilde{u}$ into the incompressible Navier-Stokes equations solved by u , we derive the so-called *Reynolds averaged Navier-Stokes equations*, denoted with the acronym RANS. These are a cascade of differential equations for $\langle u \rangle$, \tilde{u} and their higher order correlations. The main problem is that RANS models are not closed.

For instance, one can start by taking the average of the Navier-Stokes system to derive a differential relation for $\langle u \rangle$; however, owing to the non-linearities coming into play in the original system, when doing that an additional term appears, often called *Reynolds stress*, which depends on the oscillating component \tilde{u} . More precisely, one needs an information on the correlation $\langle \tilde{u} \otimes \tilde{u} \rangle$ in order to close the system and derive an equation for $\langle u \rangle$. This reflects the physical principle that small-scale processes affect the mean motion at larger scales. The problem is that such an information is not given *a priori*; in order to find it, one can use the equation for \tilde{u} and write an equation for the Reynolds stress, but it is easily seen that a third-order correlation appears. In this way, one can generate a cascade of equations for the correlation functions of any order, but again, the equation for the n -th order correlation will depend on the moments up to order $n + 1$.

Thus, a *closure hypothesis* is needed in order to obtain a closed system and work on it. From this point of view, Kolmogorov’s system (1) belongs to the class of two-equation models of turbulence, for which, as the name suggests, the closure is given by imposing two additional equations. After having been almost left aside for a long time in favour of the one-equation models

(whose main example is the k -model proposed by Prandtl in [15], and where the closure is given by imposing only one equation, typically for the turbulent kinetic energy k), two-equation models have recently gained a renewed interest, essentially because in those theories the local length scale ℓ of turbulence has not to be prescribed arbitrarily, but can be deduced from the solution of the system.

The most known two-equation models of turbulence are probably the so-called $k - \varepsilon$ models. Here, ε typically¹ denotes the turbulent energy dissipation rate. In fact, any pair of variables of the form $k^a \varepsilon^b$, with integer values of a and b , can be used to derive a different but equivalent $k - \varepsilon$ model. In Kolmogorov's system (1), the small-scale unknowns are instead k and ω : in this case, the energy dissipation rate ε and the length scale ℓ can be recovered via the formulas

$$(2) \quad \varepsilon = k\omega \quad \text{and} \quad \ell = c \frac{\sqrt{k}}{\omega},$$

where $c > 0$ is an adimensional physical constant.

Previous mathematical results

The systematic mathematical analysis of equations (1) is quite recent. Let us give a short account of previous works devoted to their study.

In [13] (see also [12] for an announcement of the results), Mielke and Naumann proved the existence of global in time weak solutions to system (1) set on the torus \mathbb{T}^3 . The solutions constructed therein are analogous in spirit to Leray's finite energy solutions for the incompressible Navier-Stokes equations, except for the fact that k cannot really belong to L^2 , because of the last term appearing on the right-hand side of its equation, which is merely in L^1 (yet, the authors are able to prove a gain of integrability for that quantity). Also, due to the lack of space compactness for that L^1 term, the solutions are in fact weak solutions of the system, but one has to admit the presence of a defect measure in the weak formulation of the equation for k . We refer to work [13] for more details.

An important point of the analysis of [13] is that the initial frequency function ω_0 and the initial mean turbulent kinetic energy k_0 are taken strictly positive: more precisely, the authors assume that

$$(3) \quad 0 < \omega_{0,*} \leq \omega_0 \leq \omega_0^* \quad \text{and} \quad k_0 \geq k_{0,*} > 0,$$

for suitable positive constants $\omega_{0,*}$, ω_0^* and $k_{0,*}$. As a matter of fact, those positivity conditions are preserved by the flow, in the sense that, for later times $t > 0$ and for any $x \in \mathbb{T}^3$, one gets

$$0 < \omega_*(t) \leq \omega(t, x) \leq \omega^*(t) \quad \text{and} \quad k(t, x) \geq k_*(t) > 0,$$

for suitable time-dependent functions $\omega_*(t)$, $\omega^*(t)$ and $k_*(t)$. See also our discussion in Subsection 3.1 below. Thus, system (1) becomes really parabolic under conditions (3), and this is a key ingredient in the analysis of [13].

The same positivity assumption (3) is formulated (and plays a fundamental role) in works [8] and [9] by Kosewski and Kubica. There, the authors prove respectively local in time well-posedness, global in time well-posedness for small data, for the Kolmogorov model (1) set in \mathbb{T}^3 , at H^2 level of regularity on the initial data.

At this point, let us comment more on assumption (3). As far as our understanding of turbulence theory goes, we observe that, while the first condition on ω_0 is somewhat expected from the physical point of view, the strict positivity of k_0 may look as an *ad hoc* assumption, formulated in order to guarantee the non-degeneracy of the parabolic smoothing for the solutions.

¹Be aware that, in the original Kolmogorov paper [7], the symbol ε was used instead to represent the term $|\mathbb{D}u|^2$ appearing in the third equation of (1).

To the best of our knowledge, the only work which is able to deal with a possible vanishing of the function k_0 is paper [3] by Bulíček and Málek. There, the authors study the existence of global in time finite energy (in the same sense as the one described above for [13]) weak solutions to system (1). Here, the word “weak” has to be intended in a classical sense, since no defect measures are needed. Remarkably, the problem is set on a bounded domain $\Omega \subset \mathbb{R}^3$, and suitable boundary conditions are formulated for the three unknowns u , k (called b in [3]) and ω . However, the authors reformulate the system in terms of the energy function E , which is defined (in our notation) as

$$E := \frac{1}{2} |u|^2 + \frac{\nu}{\alpha_4} k.$$

The energy function E satisfies a better relation than k itself, inasmuch as the higher order terms appearing in its equation are in a divergence form. Moreover, for reasons linked again to the lack of integrability for k , also the equation for ω is reformulated: essentially, one integrates by parts the non-linearity occurring in the parabolic term of the second equation of (1). Thus, the analysis is performed on a system collecting the (new) equations for u , ω and E (and not on the original system), and the existence of weak solutions is proved for that system. It is worth noticing that, as pointed out in [3], the two systems of equations are equivalent for smooth solutions, but they are not in a weak setting. However, the authors show that, recovering k from u and E , the triplet of functions (u, ω, k) are a kind of *suitable weak solutions* for the original system (1), where the equation for ω is yet the modified one.

Concerning the possible vanishing of the mean turbulent kinetic energy, in [3] k_0 is assumed to satisfy

$$k_0 > 0, \quad \text{with} \quad \log k_0 \in L^1(\Omega),$$

instead of (3). Thus, k_0 is not assumed to possess a strictly positive lower bound, but the logarithmic hypothesis gives a control on its possible vanishing. Correspondingly, the solutions constructed in that paper are proved to satisfy the same type of bounds for k , uniformly in time.

Contents of the paper and main results

In this paper, we study the 1-D version of the *Kolmogorov two-equation model of turbulence*. Dropping the divergence-free condition on the velocity field, and correspondingly erasing the pressure function from the equation for u , system (1) reduces to

$$(4) \quad \begin{cases} \partial_t u + u \partial_x u - \nu \partial_x \left(\frac{k}{\omega} \partial_x u \right) = 0 \\ \partial_t \omega + u \partial_x \omega - \alpha_1 \partial_x \left(\frac{k}{\omega} \partial_x \omega \right) = -\alpha_2 \omega^2 \\ \partial_t k + u \partial_x k - \alpha_3 \partial_x \left(\frac{k}{\omega} \partial_x k \right) = -k \omega + \alpha_4 \frac{k}{\omega} |\partial_x u|^2, \end{cases}$$

supplemented with suitable initial data

$$u|_{t=0} = u_0, \quad \omega|_{t=0} = \omega_0, \quad k|_{t=0} = k_0.$$

The precise assumptions on those initial data will be specified below.

We avoid here any consideration about boundary effects, and we set the previous system (4) in a very simple geometry, namely in the spacial domain

$$\Omega = \mathbb{T} := [-\pi, \pi] / \sim,$$

where \sim denotes the equivalence relation which identifies $-\pi$ and π . Of course, this is a quite important simplification, because it completely removes all boundary-induced turbulent phenomena in the dynamics. However, our focus here is rather on understanding the well-posedness vs ill-posedness problem of Kolmogorov’s model in the case of vanishing turbulent kinetic energy k .

The main contribution of this paper is twofold. On the one hand, we prove that, even for vanishing k_0 (which implies that $k(t)$ also vanishes at any later time), the system is locally in time well-posed in spaces H^m , for any $m \in \mathbb{N}$ such that $m \geq 2$. We refer to Theorem 1.1 below for the precise statement. On the other hand, we show also that those solutions, in general, blow up in finite time: this is the content of Theorem 1.3.

We will comment more below about our results and their meaning. For the time being, in order to explain the main difficulties linked to our study and to motivate the assumptions we will formulate in our main theorems, let us introduce a toy-model.

A toy-model. At a first sight, the local in time well-posedness result may appear a trivial consequence of classical quasi-linear hyperbolic theory applied to system (4). Actually, this is not completely true, owing to the possible vanishing of the function k , which in fact creates some complications and makes the argument for the local well-posedness more tricky than that.

In order to see this, let us introduce the following toy-model:

$$(5) \quad \begin{cases} \partial_t u + u \partial_x u - \partial_x (\gamma \partial_x u) = 0 \\ \partial_t \gamma + u \partial_x \gamma - \partial_x (\gamma \partial_x \gamma) = \gamma |\partial_x u|^2, \end{cases}$$

supplemented with initial conditions $(u, \gamma)|_{t=0} = (u_0, \gamma_0)$, for some $\gamma_0 \geq 0$. Notice that a similar system has been recently introduced in [11].

In system (5), the function γ essentially plays the role of the viscosity function k/ω appearing in the original system (4). Of course, we have retained only the main structural terms of the equation for k/ω , namely transport by u and diffusion with diffusion coefficient equal to γ itself, as well as the apparently most complicated term $\gamma |\partial_x u|^2$. In addition, we have set all the parameters ν and α_j equal to 1.

It is easy to see (for instance, by repeating the computation presented in Subsection 3.1 below) that system (5) preserves the minimal value of γ , exactly as equations (4) do for k . So, we can deduce that $\gamma(t) \geq 0$ for any time $t \geq 0$ for which the solution exists. Moreover, we can reproduce L^2 energy estimates for u and L^1 estimates for γ , obtaining similar controls as the ones encoded in the original system. Actually, the toy-model (5) preserves much more properties of the solutions of the original equations, in particular all the properties which allow us to deduce the singularity formation in finite time, but at this stage we do not need that in our discussion.

The problems come when we try to propagate higher regularity norms of the solution. In order to fix ideas, let us focus on H^2 estimates, which are the simplest ones, because they require to propagate the minimal integer regularity which is expected to be needed for well-posedness. Now, if we write an equation for $\partial_x^2 u$, we get

$$\partial_t \partial_x^2 u + u \partial_x \partial_x^2 u - \partial_x (\gamma \partial_x \partial_x^2 u) = [u, \partial_x^2] \partial_x u + \partial_x ([\gamma, \partial_x^2] \partial_x u),$$

where we have denoted by $[A, B] := AB - BA$ the commutator between two operators A and B . The bad term comes from the second commutator: indeed, when performing energy estimates for $\partial_x^2 u$, one has to deal with a term of the form

$$(6) \quad \int_{\Omega} \partial_x^2 \gamma \partial_x u \partial_x^3 u \, dx.$$

That integral is of course out of control, because any third order derivative can be estimated only if a prefactor $\sqrt{\gamma}$ appears in front of it, but we cannot simply do that because we have no strictly positive lower bound for $\sqrt{\gamma}$. It is also apparent that the situation does not improve when integrating by parts, for instance.

However, it is easy to see that the integral (6) can be estimated if one disposes of a L^∞ bound on $\partial_x \sqrt{\gamma}$. Indeed, it is enough to note that

$$\int_{\Omega} \partial_x^2 \gamma \partial_x u \partial_x^3 u \, dx = 2 \int_{\Omega} \sqrt{\gamma} \partial_x^2 \sqrt{\gamma} \partial_x u \partial_x^3 u \, dx + 2 \int_{\Omega} (\partial_x \sqrt{\gamma})^2 \partial_x u \partial_x^3 u \, dx :$$

the first term on the right can be immediately bounded by Cauchy-Schwarz inequality, whereas for the second we can integrate by parts once, and then apply Cauchy-Schwarz again. This consideration motivates us to work with the “good unknown” $\sqrt{\gamma}$ rather than with γ itself. Because of Sobolev inequalities, H^2 bounds for $\sqrt{\gamma}$ imply H^2 bounds for γ , whereas of course the contrary is not true in general. On the other hand, when reformulating system (5) in terms of $\sqrt{\gamma}$, the bad terms disappear, because whenever a third order derivative hits $\sqrt{\gamma}$ or u , it will be always possible to couple that term with another prefactor $\sqrt{\gamma}$.

Local well-posedness. The previous considerations constitute the main motivation to formulate our first main result, namely well-posedness of equations (4) in a hyperbolic functional setting, where however conditions will be formulated on (u, ω, \sqrt{k}) instead of (u, ω, k) . This is the matter of Theorem 1.1 below.

We consider solutions at H^m level of regularity, for any integer $m \geq 2$. Our proof will be based on elementary methods (see more details below); our goal is to give a simple to understand, but robust argument to construct solutions for vanishing k . Thus, we focus only on the case of regularity of integer order. As a matter of fact, we expect our result to hold true more in general, namely in H^s for any $s > 3/2$, or even in the class of Besov spaces $B_{2,r}^s$, for any $s > 3/2$ and $r \in [1, +\infty]$, up to the endpoint case $s = 3/2$ and $r = 1$. However, the case of fractional regularity looks a bit more involved from the technical point of view, because of the commutator estimates discussed above and the need of coupling the highest order terms with a $\sqrt{\gamma}$ in those computations. This is why we prefer to leave the fractional regularity case apart from our discussion.

On the other hand, as symmetry in L^2 of the parabolic part of the equation plays a major role in our argument (recall the arguments explained for the toy-model), the generalisation to integrability exponents different than 2 does not really look straightforward.

After those considerations, it is now time to state our first result. Notice that, differently from all previous results, no special assumptions on the possible vanishing of $\sqrt{k_0}$ are needed.

Theorem 1.1. *Fix $m \in \mathbb{N}$ such that $m \geq 2$. Let (u_0, ω_0, k_0) be a triplet of functions belonging to $H^m(\Omega)$. In addition, assume that the following conditions are verified:*

- (i) *there exist two constants such that $0 < \omega_* \leq \omega_0 \leq \omega^*$;*
- (ii) *$k_0 \geq 0$ is such that $\beta_0 := \sqrt{k_0} \in H^m(\Omega)$.*

Then, there exists a time $T > 0$ such that system (4), supplemented with (u_0, ω_0, k_0) as initial datum, admits a solution (u, ω, k) on $[0, T] \times \Omega$ with the following properties:

- (a) *for any time $t \in [0, T]$, one has $k(t) \geq 0$ and $\min_{x \in \Omega} \omega(t, x) > 0$;*
- (b) *the functions u , ω and \sqrt{k} belong to the space $L^\infty([0, T]; H^m(\Omega)) \cap \bigcap_{s < m} C([0, T]; H^s(\Omega))$;*
- (c) *the three functions $\sqrt{k/\omega} \partial_x^{m+1} u$, $\sqrt{k/\omega} \partial_x^{m+1} \omega$ and $\sqrt{k/\omega} \partial_x^{m+1} \sqrt{k}$ all belong to the space $L^2([0, T]; L^2(\Omega))$.*

In addition, such solution is unique in the class

$$\mathbb{X}_T(\Omega) := \left\{ (u, \omega, k) \mid \begin{array}{l} \omega, \omega^{-1}, k \in L^\infty([0, T] \times \Omega), \quad \omega > 0, \quad k \geq 0 \\ u, \omega, \sqrt{k} \in C([0, T]; L^2(\Omega)), \quad \partial_x u, \partial_x \omega, \partial_x \sqrt{k} \in L^\infty([0, T] \times \Omega) \end{array} \right\}.$$

Finally, denote by $T^ > 0$ the lifespan of the previous solution, and define*

$$\mathcal{E}_0 := \|u_0\|_{H^2}^2 + \|\omega_0\|_{H^2}^2 + \left\| \sqrt{k_0} \right\|_{H^2}^2.$$

Then there exists a constant $C > 0$, depending only on $(\nu, \alpha_1 \dots \alpha_4)$, such that one has

$$(7) \quad T^* \geq \min \left\{ 1, \frac{C}{(\max \{1, \mathcal{E}_0\})^2} \right\}.$$

Moreover, one has $T^* < +\infty$ if and only if

$$(8) \quad \int_0^{T^*} \left(1 + \|\sqrt{k}\|_{L^\infty}^2 \right) \left\| \left(\partial_x u, \partial_x \omega, \partial_x \sqrt{k} \right) \right\|_{L^\infty}^2 dt = +\infty.$$

A remark on the previous theorem is in order.

Remark 1.2. Notice that the continuation criterion (8) implies, as is classical in quasi-linear hyperbolic problems, that the lifespan of the solution does not depend on its regularity: if the datum is smooth, say in H^m with $m > 2$ large, the lifespan of the corresponding solution as a H^2 solution is the same as a H^m solution.

In particular, this justifies the lower bound (7) in terms only of the H^2 norm of the initial datum.

The proof of Theorem 1.1 is based on elementary arguments. We approximate problem (4) with vanishing initial k_0 with problems having initial $k_{0,\varepsilon}$'s which satisfy $k_{0,\varepsilon} \geq \varepsilon$. We thus construct approximate solutions by using the result from [8], and we pass to the limit in the approximation parameter $\varepsilon \rightarrow 0^+$ by using a compactness argument. This proves the existence, for which we refer to Subsection 4.1. The relevant uniform bounds in suitable norms will be shown before, in Section 3.

Uniqueness of solutions is based on a simple stability estimate in L^2 : see Subsection 4.2. Notice that our argument cannot be adapted to prove a weak-strong uniqueness principle for system (4), since higher regularity is required on both families of solutions.

To conclude this part, we point out that, in order to lighten the presentation, our proof will be limited to treat the lower regularity case only, namely the case $m = 2$. Generalisation to the case of $m > 2$ is rather straightforward, hence omitted.

Singularity formation. Next, we show that the solutions we previously constructed in Theorem 1.1 actually blow up in finite time, in general. This is the matter of the next result. That statement will be made more precise in Subsection 4.3, by formulating Theorem 4.5.

Theorem 1.3. *There are smooth initial data (u_0, ω_0, k_0) such that the corresponding solutions (u, ω, k) to system (4) blow up in finite time.*

Before moving on, let us spend a few words to comment the previous finite time singularity result. It is tempting to interpret it as an indication of the fact that the Kolmogorov model (1) is not a good model for describing turbulence. This remark would be pretty much in connection with some physical theories, which criticise $k - \varepsilon$ models and their pertinence for the theoretical understanding of turbulence in fluids.

However, we have to point out that system (4) is merely a 1-D reduction of the original Kolmogorov model (1), and, as such, it loses very much physics with respect to that system. In particular, one may object that real turbulence is expected to develop only in higher dimensions, namely in dimensions 2 and 3. Also, in system (4) we neglect the divergence-free condition on the velocity field, as well as the pressure term. This simplification kills completely non-local effects (often called long-range interactions in physical theories), which though are strongly supposed to play an important role in turbulence (see e.g. Chapter 7 of [10] and Chapter 16 of [4]).

More realistically, and as far as our understanding of turbulence goes, we can interpret Theorem 1.3 as a reflection of the fact that the model is formulated by Kolmogorov in order to describe

a fully developed turbulence, so it loses its meaning in the regime of vanishing turbulent kinetic energy. As a matter of fact, in Kolmogorov's theory of fully developed turbulence, the turbulent kinetic energy produced at small scales, whose typical length scale ℓ is defined in (2), is completely dissipated at large scales by viscosity. However, for vanishing $k \sim 0$ the definition of the typical length scale ℓ , at which turbulent phenomena act, degenerates. Therefore, we expect that, for vanishing k , one should rather change the model.

Organisation of the paper

Let us present now a short overview of the paper. First of all, in Section 2 we collect some useful tools from Fourier analysis which are needed in our study. In Section 3 we present *a priori* estimates for system (4). They will be the key step in the proof of our results, namely Theorems 1.1 and 1.3, which will be presented in Section 4.

2 Littlewood-Paley theory on the torus

We recall here a few elements of Littlewood-Paley theory and use them to derive some useful inequalities. We refer *e.g.* to Chapter 2 of [1] for details on the construction in the \mathbb{R}^d setting, to reference [5] for the adaptation to the case of a d -dimensional periodic box \mathbb{T}_a^d , where $a \in \mathbb{R}^d$ (this means that the domain is periodic in space with, for any $1 \leq j \leq d$, period equal to $2\pi a_j$ with respect to the j -th component).

For simplicity of presentation, we focus here on the one-dimensional case $d = 1$, the only relevant one for our study, and on the case where $a_1 = 1$. Hence, we simply write the spacial domain as $\Omega = \mathbb{T}$. We also denote by $|\mathbb{T}| = \mathcal{L}(\mathbb{T})$ the Lebesgue measure of the box \mathbb{T} .

First of all, let us recall that, for a tempered distribution $u \in \mathcal{S}'(\mathbb{T})$, we denote by $\mathcal{F}u = (\widehat{u}_k)_{k \in \mathbb{Z}}$ its Fourier series, so that we have

$$u(x) = \sum_{k \in \mathbb{Z}} \widehat{u}_k e^{ik \cdot x}, \quad \text{with} \quad \widehat{u}_k := \frac{1}{|\mathbb{T}|} \int_{\mathbb{T}} u(x) e^{-ik \cdot x} dx.$$

Next, we introduce the so called *Littlewood-Paley decomposition*, based on a non-homogeneous dyadic partition of unity with respect to the Fourier variable. We fix a smooth scalar function φ such that $0 \leq \varphi \leq 1$, φ is even and supported in the ring $\{r \in \mathbb{R} \mid 5/6 \leq |r| \leq 12/5\}$, and such that

$$\forall r \in \mathbb{R} \setminus \{0\}, \quad \sum_{j \in \mathbb{Z}} \varphi(2^{-j} r) = 1.$$

Let us define $|D| := (-\partial_x^2)^{1/2}$ as the Fourier multiplier² of symbol $|k|$, for $k \in \mathbb{Z}$. The dyadic blocks $(\Delta_j)_{j \in \mathbb{Z}}$ are then defined by

$$\forall j \in \mathbb{Z}, \quad \Delta_j u := \varphi(2^{-j} |D|) u = \sum_{k \in \mathbb{Z}} \varphi(2^{-j} |k|) \widehat{u}_k e^{ik \cdot x}.$$

Notice that, for $j < 0$ negative enough (in general, depending on the box \mathbb{T}_a^d), one has $\Delta_j \equiv 0$. In addition, one has the following Littlewood-Paley decomposition in $\mathcal{S}'(\mathbb{T})$:

$$\forall u \in \mathcal{S}'(\mathbb{T}), \quad u = \widehat{u}_0 + \sum_{j \in \mathbb{Z}} \Delta_j u \quad \text{in} \quad \mathcal{S}'(\mathbb{T}).$$

Recall that \widehat{u}_0 is simply the mean value of u on \mathbb{T} :

$$\widehat{u}_0 = \bar{u} = \frac{1}{|\mathbb{T}|} \int_{\mathbb{T}} u(x) dx.$$

²Throughout we agree that $f(D)$ stands for the pseudo-differential operator $u \mapsto \mathcal{F}^{-1}(f \mathcal{F}u)$.

We explicitly remark that, for any $j \in \mathbb{Z}$, the Fourier multipliers Δ_j are linear operators which are bounded on L^p for any $p \in [1, +\infty]$, with norm *independent* of j and p .

Now, we present a simplified version of the classical *Bernstein inequalities*, which turns out to be enough for our scopes. We refer to Chapter 2 of [1] for the statement in its full generality.

Lemma 2.1. *There exists a universal constant $C > 0$, on the size of the torus \mathbb{T} and on the support of the function φ fixed above, such that the following properties hold true: for any $j \in \mathbb{Z}$, for any $m \in \mathbb{N}$, for any couple $(p, q) \in [1, +\infty]^2$ such that $p \leq q$, and for any smooth enough $u \in \mathcal{S}'(\mathbb{T})$, we have*

$$\|\Delta_j u\|_{L^q} \leq C 2^{j\left(\frac{1}{p} - \frac{1}{q}\right)} \|\Delta_j u\|_{L^p}$$

and $C^{-m-1} 2^{-jm} \|\Delta_j u\|_{L^p} \leq \|\partial_x^m \Delta_j u\|_{L^p} \leq C^{m+1} 2^{jm} \|\Delta_j u\|_{L^p}.$

It is well known that, for all $s \in \mathbb{R}$, the space H^s can be characterised in terms of Littlewood-Paley decomposition (see Section 2.7 of [1]). In particular, one has the equivalence

$$(9) \quad \|u\|_{H^s}^2 \sim |\widehat{u}_0|^2 + \sum_{j \in \mathbb{Z}} 2^{2sj} \|\Delta_j u\|_{L^2}^2.$$

Next, we use Littlewood-Paley decomposition to derive some useful inequalities. To begin with, we recall the following classical interpolation estimate:

$$\forall f \in H^2(\mathbb{T}), \quad \|\partial_x f\|_{L^2} \lesssim \|f\|_{L^2}^{1/2} \|\partial_x^2 f\|_{L^2}^{1/2}.$$

Moreover, we have the following inequality:

$$(10) \quad \forall f \in H^2(\mathbb{T}), \quad \|\partial_x f\|_{L^4} \lesssim \|f\|_{L^\infty}^{1/2} \|\partial_x^2 f\|_{L^2}^{1/2}.$$

Indeed, if f is constant, the inequality is easily verified. On the other hand, if $\partial_x f \neq 0$, the claimed bound is easily seen by writing, after an integration by parts,

$$\|\partial_x f\|_{L^4}^4 = \int_{\mathbb{T}} (\partial_x f)^4 dx = -3 \int_{\mathbb{T}} f \partial_x^2 f (\partial_x f)^2 dx.$$

In fact, observing that $\partial_x \bar{f} = 0$, we can obtain a refined version of (10), namely

$$\forall f \in H^2(\mathbb{T}), \quad \|\partial_x f\|_{L^4} \lesssim \|f - \bar{f}\|_{L^\infty}^{1/2} \|\partial_x^2 f\|_{L^2}^{1/2}.$$

Similarly, we can prove an interpolation estimate in the same spirit of Poincaré-Wirtinger inequality, but which allows us to precise the exponents of the interpolation:

$$(11) \quad \forall f \in H^1(\mathbb{T}) \quad \text{such that } \bar{f} = 0, \quad \|f\|_{L^\infty} \lesssim \|f\|_{L^p}^{p/(p+2)} \|\partial_x f\|_{L^2}^{2/(p+2)},$$

for any $p \in [1, +\infty]$. Indeed, with the help of the Bernstein inequalities, for any $N \in \mathbb{N}$ to be fixed later, we can estimate

$$\begin{aligned} \|f\|_{L^\infty} &\leq \sum_{j < 0} \|\Delta_j f\|_{L^\infty} + \sum_{j=0}^N \|\Delta_j f\|_{L^\infty} + \sum_{j \geq N+1} \|\Delta_j f\|_{L^\infty} \\ &\lesssim \sum_{j < 0} 2^{j/p} \|\Delta_j f\|_{L^p} + \sum_{j=0}^N 2^{j/p} \|\Delta_j f\|_{L^p} + \sum_{j \geq N+1} 2^{-j/2} \|\partial_x \Delta_j f\|_{L^2} \\ &\lesssim \left(1 + 2^{N/p}\right) \|f\|_{L^p} + 2^{-N/2} \|\partial_x f\|_{L^2} \lesssim 2^{N/p} \|f\|_{L^p} + 2^{-N/2} \|\partial_x f\|_{L^2}. \end{aligned}$$

At this point, choosing N such that

$$2^{N/p} \|f\|_{L^p} \approx 2^{-N/2} \|\partial_x f\|_{L^2} \quad \implies \quad 2^N \approx \left(\frac{\|\partial_x f\|_{L^2}}{\|f\|_{L^p}} \right)^{2p/(p+2)}$$

completes the proof of (11).

3 A priori estimates

In this section, we collect useful *a priori* estimates satisfied by smooth solutions (u, ω, k) to system (4), defined in $\mathbb{R}_+ \times \Omega$, where we recall that $\Omega = \mathbb{T}$. We will see in Section 4, and more precisely in Subsection 4.1 how to use those bounds in order to prove local well-posedness at the level of regularity claimed in Theorem 1.3.

In all the estimates below, we have to put special attention on the dependence of all the estimates on the lower bound of k , which may vanish, under our assumptions.

3.1 ODE analysis

We start by deriving pointwise upper and lower bounds for the functions ω and k . This is based on the following ODE analysis. Recall that, throughout this section, the solution (u, ω, k) is assumed to be smooth, defined on the whole $\mathbb{R}_+ \times \Omega$.

First of all, let us define λ and μ as the solutions to the following system of ODEs:

$$(12) \quad \begin{cases} \frac{d}{dt}\lambda = -\alpha_2 \lambda^2, & \lambda|_{t=0} = \lambda_0 \\ \frac{d}{dt}\mu = -\lambda \mu, & \mu|_{t=0} = \mu_0, \end{cases}$$

for some λ_0 and μ_0 to be fixed later. System (12) can be solved explicitly: we have

$$\lambda(t) = \frac{\lambda_0}{\lambda_0 \alpha_2 t + 1} \quad \text{and} \quad \mu(t) = \frac{\mu_0}{(\lambda_0 \alpha_2 t + 1)^{1/\alpha_2}}.$$

Next, let us introduce the quantities

$$\omega^* := \max_{x \in \Omega} \omega_0(x), \quad \omega_* := \min_{x \in \Omega} \omega_0(x), \quad k_* := \min_{x \in \Omega} k_0(x)$$

as the initial upper and lower bounds of the functions ω and k . We also define their counterparts at any time $t \geq 0$:

$$\omega_{\min}(t) := \min_{x \in \Omega} \omega(t, x) \quad \text{and} \quad \omega_{\max}(t) := \max_{x \in \Omega} \omega(t, x).$$

Due to the H^2 regularity of the solutions, these functions are Lipschitz and, as a consequence of Rademacher theorem, they are differentiable almost everywhere. Now, we observe that, from (4), we immediately infer that, at the point $x = x_{\max}(t)$ which realises the maximum value of $\omega(t)$, one has

$$\frac{d}{dt}\omega_{\max} = \alpha_1 \left(\frac{k}{\omega} \partial_x^2 \omega \right)_{|x=x_{\max}} - \alpha_2 \omega_{\max}^2 \leq -\alpha_2 \omega_{\max}^2.$$

Hence, solving the first equation in (12) with initial datum $\lambda_0 = \omega^*$, we deduce

$$\omega_{\max}(t) \leq \frac{\omega^*}{\omega^* \alpha_2 t + 1} \quad \forall t \geq 0,$$

where $T \in]0, +\infty]$ denotes the maximal existence time. In a pretty analogous way, we also obtain that

$$\omega_{\min}(t) \geq \frac{\omega_*}{\omega_* \alpha_2 t + 1} \quad \forall t \geq 0.$$

In the end, we get the following pointwise estimate for ω :

$$(13) \quad \forall (t, x) \in \mathbb{R}_+ \times \Omega, \quad 0 < \frac{\omega_*}{\omega_* \alpha_2 t + 1} \leq \omega_{\min}(t) \leq \omega(t, x) \leq \omega_{\max}(t) \leq \frac{\omega^*}{\omega^* \alpha_2 t + 1}.$$

In addition, we can define

$$k_{\min}(t) := \min_{x \in \Omega} k(t, x).$$

Denoting by $x = x_{\min}(t)$ the point where the minimum value of $k(t)$ is attained, we obtain that $k_{\min}(t)$ satisfies the ODE

$$\frac{d}{dt} k_{\min} = \alpha_3 \left(\frac{k}{\omega} \partial_x^2 k \right)_{|x_{\min}} - k_{\min} \omega_{|x_{\min}} + \alpha_4 \left(\frac{k}{\omega} |\partial_x u|^2 \right)_{|x_{\min}} \geq -k_{\min} \omega_{\max}.$$

Solving the second equation (12), where we take $\lambda(t) = \omega_{\max}(t)$ and $\mu_0 = k_*$, we infer

$$(14) \quad k_{\min}(t) \geq \frac{k_*}{(\omega^* \alpha_2 t + 1)^{1/\alpha_2}} \quad \forall t \geq 0.$$

Notice that, under the assumptions of Theorem 1.3, we have $k_* = 0$. Therefore, we deduce that

$$(15) \quad \forall (t, x) \in \mathbb{R}_+ \times \Omega, \quad k(t, x) \geq 0.$$

3.2 Basic energy estimates

We now perform basic L^2 estimates for the smooth solution (u, ω, k) , defined in $\mathbb{R}_+ \times \Omega$. Those estimates will provide us with a control on the low frequencies of the solution. The bounds for the high frequencies will be given in the next subsection.

To begin with, we test the equation for u against u itself, and then integrate over Ω : after performing suitable integrations by parts, we find

$$\frac{1}{2} \frac{d}{dt} \int_{\Omega} |u|^2 dx + \nu \int_{\Omega} \frac{k}{\omega} |\partial_x u|^2 dx = 0.$$

After an integration in time, we immediately deduce the following estimate: for all $t \geq 0$, one has

$$(16) \quad \|u(t)\|_{L^2}^2 + 2\nu \int_0^t \int_{\Omega} \frac{k}{\omega} |\partial_x u|^2 dx d\tau \leq \|u_0\|_{L^2}^2.$$

Repeating the same computations for the second equation in (4), tested this time against the function ω , we get

$$(17) \quad \frac{1}{2} \frac{d}{dt} \int_{\Omega} |\omega|^2 dx + \alpha_1 \int_{\Omega} \frac{k}{\omega} |\partial_x \omega|^2 dx + \alpha_2 \int_{\Omega} \omega^3 dx = \frac{1}{2} \int_{\Omega} \partial_x u \omega^2.$$

We have to control the last term in the right-hand side of the previous equality. There are several possibilities for that; for instance, the cheapest choice in terms of regularity consists in taking advantage of the last term on the left-hand side of (17), thus making³ a $L_T^3(L^3)$ norm of $\partial_x u$ appear. In turn, this would be fine, in view of the higher order estimates of Subsection 3.3 below. However, here we opt for a simpler control and employ a L^∞ bound for $\partial_x u$, which will be justified from the higher order estimates of Subsection 3.3: we get

$$\left| \int_{\Omega} \partial_x u \omega^2 \right| \leq \|\partial_x u\|_{L^\infty} \|\omega\|_{L^2}^2.$$

Plugging this bound into (17) and integrating in time, by Grönwall's lemma we get the following estimate: for any time $t \geq 0$, one has

$$(18) \quad \|\omega(t)\|_{L^2}^2 + \int_0^t \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x \omega|^2 dx d\tau + \|\omega\|_{L_t^3(L^3)}^3 \leq \|\omega_0\|_{L^2}^2 \exp \left(C \int_0^t \|\partial_x u(\tau)\|_{L^\infty} d\tau \right),$$

where the multiplicative constant $C > 0$ depends only on the values of α_1 and α_2 . Recall that we dispose also of L^∞ controls on ω , which are provided by (13).

³From now on, we will use the following notation: if X is a Fréchet space of functions over Ω , for any finite $T > 0$ we note $L_T^p(X) = L^p([0, T]; X)$ and $C_T(X) = C([0, T]; X)$.

Let us turn our attention to the estimates for k . Looking at the third equation of (4), it is easy to realise that a L^2 estimate for k has no chance to be closed, owing to the presence of a right-hand side which is merely L^1 . Therefore, we rather perform an L^1 estimate on k , obtaining, after an integration by parts, the relation

$$(19) \quad \frac{d}{dt} \int_{\Omega} k \, dx + \int_{\Omega} k \omega \, dx = \alpha_4 \int_{\Omega} \frac{k}{\omega} |\partial_x u|^2 \, dx + \int_{\Omega} k \partial_x u \, dx.$$

Notice that, following the approach of [13], it is possible to propagate better integrability for k and its derivative. However, the bounds coming from the previous relation (19) are enough for our scopes. Indeed, we can write

$$\left| \int_{\Omega} k \partial_x u \, dx \right| = \left| \int_{\Omega} \sqrt{k\omega} \sqrt{\frac{k}{\omega}} \partial_x u \, dx \right| \leq \frac{1}{2} \int_{\Omega} k \omega \, dx + \frac{1}{2} \int_{\Omega} \frac{k}{\omega} |\partial_x u|^2 \, dx.$$

Therefore, by use of (16) and noticing that the constant may depend on the parameters α_i and ν , from (19) we deduce, after an integration in time, that for any time $t \geq 0$, one has

$$(20) \quad \|k(t)\|_{L^1} + \int_0^t \int_{\Omega} k \omega \, dx \, d\tau \lesssim \|k_0\|_{L^1} + \|u_0\|_{L^2}^2,$$

for an implicit multiplicative constant, which depends only on the parameters ν and α_4 . Observe that, owing to our assumption on k_0 , we have

$$\|k_0\|_{L^1} = \left\| \sqrt{k_0} \right\|_{L^2}^2 < +\infty.$$

As it will appear clear from the computation in the next subsection, we also need some control on how the function k touches zero. This piece of information is obtained by propagating regularity for \sqrt{k} , rather than for k itself. For notational convenience, let us introduce the quantities

$$\beta := \sqrt{k} \quad \text{and} \quad \beta_0 := \sqrt{k_0}.$$

To begin with, we use again the fact that, for any $t \geq 0$, one has $\|\beta(t)\|_{L^2}^2 = \|k(t)\|_{L^1}$. Therefore, owing to estimate (20), we deduce, for all $t \geq 0$, the bound

$$(21) \quad \|\beta(t)\|_{L^2} \lesssim \|\beta_0\|_{L^2} + \|u_0\|_{L^2}.$$

Next, in view of the computations of Subsection 3.3 below, we derive an equation for β . For this, we start by multiplying the third equation in (4) by the factor $1/(2\sqrt{k})$; then, straightforward manipulations yield

$$(22) \quad \partial_t \beta + u \partial_x \beta - \alpha_3 \partial_x \left(\frac{\beta^2}{\omega} \partial_x \beta \right) + \frac{1}{2} \beta \omega = \frac{\alpha_4}{2} \frac{\beta}{\omega} |\partial_x u|^2 + \alpha_3 \frac{\beta}{\omega} |\partial_x \beta|^2.$$

To conclude this part, we also recast the first and second equations of (4) in terms of the new set of unknowns (u, ω, β) . We get the new system

$$(23) \quad \begin{cases} \partial_t u + u \partial_x u - \nu \partial_x \left(\frac{\beta^2}{\omega} \partial_x u \right) = 0 \\ \partial_t \omega + u \partial_x \omega - \alpha_1 \partial_x \left(\frac{\beta^2}{\omega} \partial_x \omega \right) = -\alpha_2 \omega^2 \\ \partial_t \beta + u \partial_x \beta - \alpha_3 \partial_x \left(\frac{\beta^2}{\omega} \partial_x \beta \right) + \frac{1}{2} \beta \omega = \frac{\alpha_4}{2} \frac{\beta}{\omega} |\partial_x u|^2 + \alpha_3 \frac{\beta}{\omega} |\partial_x \beta|^2. \end{cases}$$

The higher order estimates of the next subsection will be performed on the previous system (23). As a matter of fact, it is much more natural to work with the unknown $\beta = \sqrt{k}$ rather than with k itself.

Before moving on, we notice that the basic L^2 parabolic estimates leading to a smoothing effect for the solution are not closed yet, because one still needs to control the last term appearing in the right-hand side of (18). This is a purely 1-D effect, since now it is no more meaningful to assume the velocity field u to be divergence-free.

Of course, those estimates will be closed by using higher order bounds on (u, ω, β) , which are the matter of the next subsection.

3.3 Higher order estimates

Here, we perform higher order Sobolev estimates for the smooth solution (u, ω, β) , defined on $\mathbb{R}_+ \times \Omega$. In the end, we will need to close the estimates in some (possibly small) time interval $[0, T]$. Let us recall that we limit ourselves to treat the case of solutions at H^2 level of regularity; it goes without saying that the same method applies also to give propagation of H^m norms, for any $m \geq 2$. We expect an extension to almost critical H^s regularities, for any $s > 3/2$ (or even to $B_{2,1}^{3/2}$ regularity), to be also possible, by use of Littlewood-Paley analysis and paradifferential calculus techniques. However, our goal here is to give a simple proof by elementary methods, and to put in evidence in a clear way the special structure of the system, for which the new quantity $\beta = \sqrt{k}$ plays a key role.

Finally, we stress that, once again, particular care has to be used while dealing with the function k , owing to possible regions where this quantity vanishes.

Before starting the estimates, let us observe that, for any $p \in [1, +\infty]$ and any $f \in L^p(\Omega)$, we have

$$|\widehat{f_0}| \lesssim \|f\|_{L^1} \lesssim \|f\|_{L^p}.$$

Thus, owing to the Bernstein inequalities of Lemma 2.1 and the Littlewood-Paley characterisation (9) of Sobolev spaces, we have

$$\|u\|_{H^2}^2 \sim \|u\|_{L^2}^2 + \|\partial_x^2 u\|_{L^2}^2,$$

and analogous equivalences for ω and β . Therefore, we only need to bound the second order derivatives of the solution in L^2 . We will do this in an elementary way, by performing simple energy estimates on the equations satisfied by $\partial_x^2 u$, $\partial_x^2 \omega$ and $\partial_x^2 \beta$.

We start by writing an equation for $\partial_x^2 u$. Differentiating the first equation in (23) twice with respect to the space variable, we find

$$(24) \quad \partial_t \partial_x^2 u + u \partial_x \partial_x^2 u - \nu \partial_x \left(\frac{\beta^2}{\omega} \partial_x^3 u \right) = f,$$

where we have defined

$$f = \sum_{j=1}^4 f_j := -3 \partial_x u \partial_x^2 u + 3 \nu \partial_x^2 \left(\frac{\beta^2}{\omega} \right) \partial_x^2 u + 2 \nu \partial_x \left(\frac{\beta^2}{\omega} \right) \partial_x^3 u + \nu \partial_x^3 \left(\frac{\beta^2}{\omega} \right) \partial_x u.$$

It is not hard to see that the last term appearing in the definition of f is the most dangerous one, when performing energy estimates on $\partial_x^2 u$.

Likewise, we get an equation for $\partial_x^2 \omega$:

$$(25) \quad \partial_t \partial_x^2 \omega + u \partial_x \partial_x^2 \omega - \alpha_1 \partial_x \left(\frac{\beta^2}{\omega} \partial_x^3 \omega \right) + 2 \alpha_2 \omega \partial_x^2 \omega = g,$$

where this time the external force g is defined by

$$g = \sum_{j=1}^6 g_j := -2 \partial_x u \partial_x^2 \omega - \partial_x^2 u \partial_x \omega$$

$$+ 3\alpha_1 \partial_x^2 \left(\frac{\beta^2}{\omega} \right) \partial_x^2 \omega + 2\alpha_1 \partial_x \left(\frac{\beta^2}{\omega} \right) \partial_x^3 \omega + \alpha_1 \partial_x^3 \left(\frac{\beta^2}{\omega} \right) \partial_x \omega - 2\alpha_2 (\partial_x \omega)^2 .$$

Finally, we repeat the same computations for β , finding

$$(26) \quad \partial_t \partial_x^2 \beta + u \partial_x \partial_x^2 \beta - \alpha_3 \partial_x \left(\frac{\beta^2}{\omega} \partial_x^3 \beta \right) + \frac{1}{2} \omega \partial_x^2 \beta = h ,$$

where we have denoted by h the quantity

$$\begin{aligned} h = \sum_{j=1}^{15} h_j := & -2 \partial_x u \partial_x^2 \beta - \partial_x^2 u \partial_x \beta - \partial_x \omega \partial_x \beta - \frac{1}{2} \beta \partial_x^2 \omega \\ & + 3\alpha_3 \partial_x^2 \left(\frac{\beta^2}{\omega} \right) \partial_x^2 \beta + 2\alpha_3 \partial_x \left(\frac{\beta^2}{\omega} \right) \partial_x^3 \beta + \alpha_3 \partial_x^3 \left(\frac{\beta^2}{\omega} \right) \partial_x \beta \\ & + \frac{\alpha_4}{2} \partial_x^2 \left(\frac{\beta}{\omega} \right) |\partial_x u|^2 + 2\alpha_4 \partial_x \left(\frac{\beta}{\omega} \right) \partial_x u \partial_x^2 u + \alpha_4 \frac{\beta}{\omega} (\partial_x^2 u)^2 + \alpha_4 \frac{\beta}{\omega} \partial_x u \partial_x^3 u \\ & + \alpha_3 \partial_x^2 \left(\frac{\beta}{\omega} \right) |\partial_x \beta|^2 + 2\alpha_4 \partial_x \left(\frac{\beta}{\omega} \right) \partial_x \beta \partial_x^2 \beta + 2\alpha_4 \frac{\beta}{\omega} (\partial_x^2 \beta)^2 + 2\alpha_4 \frac{\beta}{\omega} \partial_x \beta \partial_x^3 \beta . \end{aligned}$$

In the next paragraphs, we will perform L^2 estimates in (24), (25) and (26). Of course, the key point will be the control of the forcing terms f , g and h in the energy space.

3.3.1 The H^2 estimate for u

We start by testing equation (24): simple computations give us

$$(27) \quad \frac{1}{2} \frac{d}{dt} \|\partial_x^2 u\|_{L^2}^2 + \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 dx = \frac{1}{2} \int_{\Omega} \partial_x u (\partial_x^2 u)^2 dx + \int_{\Omega} f \partial_x^2 u dx .$$

The core of the argument consists in finding suitable bounds for the terms on the right-hand side. Notice that, owing to the fact that we do not dispose of a strictly positive lower bound for k , the contribution of the viscosity term may appear of no importance. Yet, the piece of information coming from that term turn out to be essential for closing the estimates.

Thus, let us bound the terms appearing on the right-hand side of the previous relation. We start by observing that

$$(28) \quad \left| \int_{\Omega} \partial_x u (\partial_x^2 u)^2 dx \right| \lesssim \|\partial_x u\|_{L^\infty} \|\partial_x^2 u\|_{L^2}^2 ,$$

which of course will be absorbed by a Grönwall-type argument.

It goes without saying that the control of the integral involving f_1 works absolutely in the same way. Thus, let us turn our attention to the integral involving f_2 : after an integration by parts, this term can be coupled with the one coming from f_3 , thus getting

$$\int_{\Omega} (f_2 + f_3) \partial_x^2 u dx = -4\nu \int_{\Omega} \partial_x \left(\frac{\beta^2}{\omega} \right) \partial_x^2 u \partial_x^3 u dx .$$

By computing

$$\partial_x \left(\frac{\beta^2}{\omega} \right) = 2 \frac{\beta}{\omega} \partial_x \beta - \frac{\beta^2}{\omega^2} \partial_x \omega ,$$

in the end we have to bound two integrals. We start by writing

$$\int_{\Omega} \frac{\beta}{\omega} \partial_x \beta \partial_x^2 u \partial_x^3 u dx = \int_{\Omega} \frac{1}{\sqrt{\omega}} \partial_x \beta \partial_x^2 u \frac{\beta}{\sqrt{\omega}} \partial_x^3 u dx ,$$

and an application of Cauchy-Schwarz and Young inequalities gives us, for any $\delta \in]0, 1[$ to be chosen later, the estimate

$$\nu \left| \int_{\Omega} \frac{\beta}{\omega} \partial_x \beta \partial_x^2 u \partial_x^3 u \, dx \right| \leq \delta \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx + C(\delta, \nu) \int_{\Omega} \frac{1}{\omega} |\partial_x \beta|^2 |\partial_x^2 u|^2 \, dx.$$

Similarly, we can control

$$\nu \left| \int_{\Omega} \frac{\beta^2}{\omega^2} \partial_x \omega \partial_x^2 u \partial_x^3 u \, dx \right| \leq \delta \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx + C(\delta, \nu) \int_{\Omega} \frac{\beta^2}{\omega^3} |\partial_x \omega|^2 |\partial_x^2 u|^2 \, dx.$$

Therefore, using also (13), in the end we get

$$(29) \quad \left| \int_{\Omega} (f_2 + f_3) \partial_x^2 u \, dx \right| \leq 2\delta \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx + C(1+t)^3 \left(\|\partial_x \beta\|_{L^\infty}^2 + \|\beta\|_{L^\infty}^2 \|\partial_x \omega\|_{L^\infty}^2 \right) \|\partial_x^2 u\|_{L^2}^2,$$

where the value of $\delta \in]0, 1[$ will be fixed later on, and the constant $C > 0$ depends only on the values of δ and of the various coefficients $(\nu, \alpha_1 \dots \alpha_4)$. Of course, for δ small enough, the first term on the right can be absorbed into the left-hand side of (27), whereas the second term can be controlled by the Grönwall lemma.

It remains us to control the term depending on f_4 . For this, we integrate by parts once, to get

$$\int_{\Omega} f_4 \partial_x^2 u \, dx = -\nu \int_{\Omega} \partial_x^2 \left(\frac{\beta^2}{\omega} \right) |\partial_x^2 u|^2 \, dx - \nu \int_{\Omega} \partial_x^2 \left(\frac{\beta^2}{\omega} \right) \partial_x u \partial_x^3 u \, dx.$$

Observe that the first term on the right is analogous to the one coming from f_2 , hence it can be controlled as we have done before (by computing a further integration by parts). On the other hand, for bounding the second term in the right-hand side of the previous equation, we compute explicitly the second order derivative of the function β^2/ω : we get

$$(30) \quad \partial_x^2 \left(\frac{\beta^2}{\omega} \right) = \frac{2}{\omega} (\partial_x \beta)^2 + 2 \frac{\beta}{\omega} \partial_x^2 \beta - 4 \frac{\beta}{\omega^2} \partial_x \beta \partial_x \omega - \frac{\beta^2}{\omega^2} \partial_x^2 \omega + 2 \frac{\beta^2}{\omega^3} (\partial_x \omega)^2.$$

So, we are reconducted to estimating five different integrals. Since the third-order term $\partial_x^3 u$ appears, the idea is to make a coefficient $\beta/\sqrt{\omega}$ appear in front of it, as we have done for (29) above. To begin with, we estimate

$$\begin{aligned} \nu \left| \int_{\Omega} \frac{\beta}{\omega} \partial_x^2 \beta \partial_x u \partial_x^3 u \, dx \right| &\leq \delta \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx + C(\delta, \nu) \int_{\Omega} \frac{1}{\omega} |\partial_x u|^2 |\partial_x^2 \beta|^2 \, dx \\ \nu \left| \int_{\Omega} \frac{\beta^2}{\omega^2} \partial_x^2 \omega \partial_x u \partial_x^3 u \, dx \right| &\leq \delta \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx + C(\delta, \nu) \int_{\Omega} \frac{\beta^2}{\omega^3} |\partial_x u|^2 |\partial_x^2 \omega|^2 \, dx, \end{aligned}$$

where $\delta > 0$ is as above. The same method would apply also to the third item in (30), but this would give a bound which is not optimal. For this reason, we prefer to integrate by parts once and write

$$(31) \quad \begin{aligned} \int_{\Omega} \frac{\beta}{\omega^2} \partial_x \beta \partial_x \omega \partial_x u \partial_x^3 u \, dx &= - \int_{\Omega} \frac{1}{\omega^2} (\partial_x \beta)^2 \partial_x^2 u \partial_x u \partial_x \omega \, dx + 2 \int_{\Omega} \frac{\beta}{\omega^3} \partial_x \beta (\partial_x \omega)^2 \partial_x^2 u \partial_x u \, dx \\ &\quad - \int_{\Omega} \frac{\beta}{\omega^2} (\partial_x^2 \beta \partial_x \omega \partial_x u + \partial_x \beta \partial_x \omega \partial_x^2 u + \partial_x \beta \partial_x^2 \omega \partial_x u) \partial_x^2 u \, dx. \end{aligned}$$

All these terms can be controlled in a simple way. For instance, using also (13) and the interpolation inequality (10), for the first term we have

$$\begin{aligned} \left| \int_{\Omega} \frac{1}{\omega^2} (\partial_x \beta)^2 \partial_x^2 u \partial_x u \partial_x \omega \, dx \right| &\lesssim \frac{1}{(\omega_{\min}(t))^2} \|\partial_x^2 u\|_{L^2} \|\partial_x u\|_{L^\infty} \|\partial_x \omega\|_{L^\infty} \|\partial_x \beta\|_{L^4}^2 \\ &\lesssim (1+t)^2 \|(\partial_x^2 u, \partial_x^2 \beta)\|_{L^2}^2 \|\partial_x u\|_{L^\infty} \|\beta\|_{L^\infty} \|\partial_x \omega\|_{L^\infty} . \end{aligned}$$

Similarly, we can bound the second term as follows:

$$\begin{aligned} \left| \int_{\Omega} \frac{\beta}{\omega^3} \partial_x \beta (\partial_x \omega)^2 \partial_x^2 u \partial_x u \, dx \right| &\lesssim \frac{1}{(\omega_{\min}(t))^3} \|\beta\|_{L^\infty} \|\partial_x^2 u\|_{L^2} \|\partial_x u\|_{L^\infty} \|\partial_x \beta\|_{L^\infty} \|\partial_x \omega\|_{L^4}^2 \\ &\lesssim \frac{\omega_{\max}(t)}{(\omega_{\min}(t))^3} \|\beta\|_{L^\infty} \|(\partial_x^2 u, \partial_x^2 \omega)\|_{L^2}^2 \|\partial_x u\|_{L^\infty} \|\partial_x \beta\|_{L^\infty} \\ &\lesssim (1+t)^2 \|\beta\|_{L^\infty} \|(\partial_x^2 u, \partial_x^2 \omega)\|_{L^2}^2 \|\partial_x u\|_{L^\infty} \|\partial_x \beta\|_{L^\infty} . \end{aligned}$$

The control of the terms in the third line of (31) is even more straightforward. All in all, we finally get

$$\left| \int_{\Omega} \frac{\beta}{\omega^2} \partial_x \beta \partial_x \omega \partial_x u \partial_x^3 u \, dx \right| \lesssim (1+t)^2 \|\beta\|_{L^\infty} \|(\partial_x u, \partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \|(\partial_x^2 u, \partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2 .$$

Finally, the terms coming from the first and fifth items of (30) can be treated as the previous one, after an integration by parts: we obtain

$$\begin{aligned} \left| \int_{\Omega} \frac{1}{\omega} (\partial_x \beta)^2 \partial_x u \partial_x^3 u \, dx \right| &\lesssim (1+t)^2 (1 + \|\beta\|_{L^\infty}) \|(\partial_x u, \partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \|(\partial_x^2 u, \partial_x^2 \beta)\|_{L^2}^2 \\ \left| \int_{\Omega} \frac{\beta^2}{\omega^3} (\partial_x \omega)^2 \partial_x u \partial_x^3 u \, dx \right| &\lesssim (1+t)^3 (1 + \|\beta\|_{L^\infty}^2) \|(\partial_x u, \partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \|(\partial_x^2 u, \partial_x^2 \omega)\|_{L^2}^2 . \end{aligned}$$

Putting all those inequalities together, we finally deduce the sought control for the f_4 term:

$$(32) \quad \left| \int_{\Omega} f_4 \partial_x^2 u \, dx \right| \leq C_1 \delta \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx + C_2 (1+t)^3 (1 + \|\beta\|_{L^\infty}^2) \|(\partial_x u, \partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \|(\partial_x^2 u, \partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2 ,$$

where the value of $\delta \in]0, 1[$ will be fixed later on, the constant $C_1 > 0$ is “universal” (independent of all the data and parameters appearing in the equations) and the constant $C_2 > 0$ depends only on the values of δ and of the various coefficients $(\nu, \alpha_1 \dots \alpha_4)$.

In the end, inserting the bounds we have obtained in (28), (29) and (32) into (27) and choosing $\delta > 0$ small enough, we find a bound for $\partial_x^2 u$:

$$(33) \quad \frac{1}{2} \frac{d}{dt} \|\partial_x^2 u\|_{L^2}^2 + \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx \lesssim (1+t)^3 (1 + \|\beta\|_{L^\infty}^2) \left(1 + \|(\partial_x u, \partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \right) \|(\partial_x^2 u, \partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2 .$$

3.3.2 The H^2 estimate for ω

We now perform H^2 estimates on ω . For this, it is enough to perform L^2 estimates on equation (25) for $\partial_x^2 \omega$. In this way, we easily get

$$\frac{1}{2} \frac{d}{dt} \|\partial_x^2 \omega\|_{L^2}^2 + \alpha_1 \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 \omega|^2 \, dx + 2\alpha_2 \int_{\Omega} \omega |\partial_x^2 \omega|^2 \, dx = \frac{1}{2} \int_{\Omega} \partial_x u (\partial_x^2 \omega)^2 \, dx + \int_{\Omega} g \partial_x^2 \omega \, dx .$$

In the estimate of the right-hand side of the equation above, only few terms change with respect to the ones we have already treated in the previous paragraph. First of all, we put together the transport term, *i.e.* the first term appearing on the right, with g_1 : we find

$$\left| \int_{\Omega} \left(\frac{1}{2} \partial_x u - 2 \partial_x u \right) |\partial_x^2 \omega|^2 dx \right| \lesssim \|\partial_x u\|_{L^\infty} \|\partial_x^2 \omega\|_{L^2}^2 .$$

Next, we easily see that the control of the terms g_3 , g_4 and g_5 is analogous to the one of f_2 , f_3 and f_4 respectively, up to change u by ω and ν by α_1 everywhere in those computations. We infer the bound

$$\begin{aligned} \left| \int_{\Omega} (g_3 + g_4 + g_5) \partial_x^2 \omega dx \right| &\leq C_1 \delta \alpha_1 \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 \omega|^2 dx \\ &+ C_2 (1+t)^3 (1 + \|\beta\|_{L^\infty}^2) \|(\partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \|(\partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2 , \end{aligned}$$

where $\delta \in]0, 1[$ can be taken arbitrarily small and the constants $C_1 > 0$ and $C_2 > 0$ are as in (32).

Thus, it remains us to bound the terms coming from g_2 and g_6 ; however, both controls are fairly straightforward. As a matter of fact, we have

$$\begin{aligned} \left| \int_{\Omega} g_2 \partial_x^2 \omega dx \right| &\lesssim \|\partial_x \omega\|_{L^\infty} \|(\partial_x^2 u, \partial_x^2 \omega)\|_{L^2}^2 \\ \left| \int_{\Omega} g_6 \partial_x^2 \omega dx \right| &\lesssim \|\partial_x \omega\|_{L^4}^2 \|\partial_x^2 \omega\|_{L^2} \lesssim \omega_{\max}(t) \|\partial_x^2 \omega\|_{L^2}^2 , \end{aligned}$$

where, in the last line, we have also used the interpolation inequality (10).

In the end, we have obtained the following bound, analogous to (33):

$$(34) \quad \begin{aligned} \frac{1}{2} \frac{d}{dt} \|\partial_x^2 \omega\|_{L^2}^2 + \alpha_1 \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 \omega|^2 dx + 2 \alpha_2 \int_{\Omega} \omega |\partial_x^2 \omega|^2 dx \\ \lesssim (1+t)^3 (1 + \|\beta\|_{L^\infty}^2) \left(1 + \|(\partial_x u, \partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \right) \|(\partial_x^2 u, \partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2 . \end{aligned}$$

3.3.3 The H^2 estimate for β

Let us switch to the control of $\partial_x^2 k$. Once again, we start with an energy estimate for equation (26), thus finding

$$\frac{1}{2} \frac{d}{dt} \|\partial_x^2 \beta\|_{L^2}^2 + \alpha_3 \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 \beta|^2 dx + \frac{1}{2} \int_{\Omega} \omega |\partial_x^2 \beta|^2 dx = \frac{1}{2} \int_{\Omega} \partial_x u (\partial_x^2 \beta)^2 dx + \int_{\Omega} h \partial_x^2 \beta dx .$$

As before, the transport term has to be joined with h_1 : this gives us

$$\left| \int_{\Omega} \left(\frac{1}{2} \partial_x u - 2 \partial_x u \right) |\partial_x^2 \beta|^2 dx \right| \lesssim \|\partial_x u\|_{L^\infty} \|\partial_x^2 \beta\|_{L^2}^2 .$$

On the other hand, the group of terms h_5 , h_6 and h_7 have to be handled in the same way as f_2 , f_3 and f_4 above. *Mutatis mutandis*, we gather the control

$$\begin{aligned} \left| \int_{\Omega} (h_5 + h_6 + h_7) \partial_x^2 \beta dx \right| &\leq C_1 \delta \alpha_3 \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 \beta|^2 dx \\ &+ C_2 (1+t)^3 (1 + \|\beta\|_{L^\infty}^2) \|(\partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \|(\partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2 . \end{aligned}$$

The previous bound holds true for any $\delta \in]0, 1[$; the precise value of δ will be fixed later on. In addition, the positive constants C_1 and C_2 satisfy the same properties as in (32) above.

Next, arguing as for g_6 and using the interpolation inequality (10), we easily get

$$\begin{aligned} \left| \int_{\Omega} (h_2 + h_4) \partial_x^2 \beta \, dx \right| &\lesssim (\|\partial_x \beta\|_{L^\infty} + \|\beta\|_{L^\infty}) \|(\partial_x^2 u, \partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2 \\ \left| \int_{\Omega} h_3 \partial_x^2 \beta \, dx \right| &\lesssim \|\partial_x \omega\|_{L^4} \|\partial_x \beta\|_{L^4} \|\partial_x^2 \beta\|_{L^2} \lesssim (\omega_{\max}(t) + \|\beta\|_{L^\infty}) \|(\partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2. \end{aligned}$$

Hence, it remains us to bound the terms appearing in the last two lines of the definition of h . We start from h_{11} : using the usual trick based on Cauchy-Schwarz and Young inequalities, for $\delta \in]0, 1[$ as above, we bound

$$\begin{aligned} \left| \int_{\Omega} h_{11} \partial_x^2 \beta \, dx \right| &\leq \delta \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx + C(\delta, \nu, \alpha_4) \int_{\Omega} \frac{1}{\omega} |\partial_x u|^2 |\partial_x^2 \beta|^2 \, dx \\ &\leq \delta \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx + C(1+t) \|\partial_x u\|_{L^\infty}^2 \|\partial_x^2 \beta\|_{L^2}^2. \end{aligned}$$

Next, we consider h_{10} : for it, we can use an integration by parts to write

$$\int_{\Omega} h_{10} \partial_x^2 \beta \, dx = -2\alpha_4 \int_{\Omega} \frac{\beta}{\omega} \partial_x^2 u \partial_x^3 u \partial_x \beta \, dx - \alpha_4 \int_{\Omega} \partial_x \left(\frac{\beta}{\omega} \right) (\partial_x^2 u)^2 \partial_x \beta \, dx.$$

It is easily seen that the first term on the right-hand side can be handled in the same way as for h_{11} , thus producing

$$\alpha_4 \left| \int_{\Omega} \frac{\beta}{\omega} \partial_x^2 u \partial_x^3 u \partial_x \beta \, dx \right| \leq \delta \nu \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 u|^2 \, dx + C(\delta, \nu, \alpha_4) (1+t) \|\partial_x \beta\|_{L^\infty}^2 \|\partial_x^2 u\|_{L^2}^2,$$

whereas for the second term one can explicitly compute the space derivative of β/ω and obtain

$$\left| \int_{\Omega} \partial_x \left(\frac{\beta}{\omega} \right) (\partial_x^2 u)^2 \partial_x \beta \, dx \right| \lesssim (1+t)^2 (1 + \|\beta\|_{L^\infty}) \|(\partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \|\partial_x^2 u\|_{L^2}^2.$$

The control of h_9 works in a pretty similar way: after computing $\partial_x(\beta/\omega)$, one easily gets

$$\left| \int_{\Omega} h_9 \partial_x^2 \beta \, dx \right| \lesssim (1+t)^2 (1 + \|\beta\|_{L^\infty}) \|(\partial_x u, \partial_x \omega, \partial_x \beta)\|_{L^\infty}^2 \|(\partial_x^2 u, \partial_x^2 \beta)\|_{L^2}^2.$$

Next, in order to estimate the term presenting h_8 , we compute

$$\partial_x^2 \left(\frac{\beta}{\omega} \right) = \frac{1}{\omega} \partial_x^2 \beta - \frac{2}{\omega^2} \partial_x \beta \partial_x \omega - \frac{\beta}{\omega^2} \partial_x^2 \omega + 2 \frac{\beta}{\omega^3} (\partial_x \omega)^2$$

and reduce our problem to bound four different integrals. For this, we are going to argue in a very similar way as done for estimating the terms appearing in (31). On the one hand, we have

$$\begin{aligned} \left| \int_{\Omega} \frac{1}{\omega} |\partial_x u|^2 (\partial_x^2 \beta)^2 \, dx \right| &\lesssim (1+t) \|\partial_x u\|_{L^\infty}^2 \|\partial_x^2 \beta\|_{L^2}^2 \\ \left| \int_{\Omega} \frac{\beta}{\omega^2} \partial_x^2 \omega |\partial_x u|^2 \partial_x^2 \beta \, dx \right| &\lesssim (1+t)^2 \|\beta\|_{L^\infty} \|\partial_x u\|_{L^\infty}^2 \|(\partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2. \end{aligned}$$

On the other hand, for the terms presenting a mixing of first order derivatives, we use the interpolation inequality (10) to deduce

$$\left| \int_{\Omega} \left(\frac{1}{\omega^2} \partial_x \beta \partial_x \omega + \frac{\beta}{\omega^3} (\partial_x \omega)^2 \right) |\partial_x u|^2 \partial_x^2 \beta \, dx \right| \lesssim (1+t)^2 (1 + \|\beta\|_{L^\infty}) \|\partial_x u\|_{L^\infty}^2 \|(\partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2.$$

Finally, we see that the terms arising from the last line of the definition of h can be handled as $h_8 \dots h_{11}$, up to replace u by β and α_4 by α_3 everywhere in the computations.

In the end, putting all those inequalities together and choosing $\delta > 0$ small enough, we arrive at the following estimate for β :

$$(35) \quad \frac{1}{2} \frac{d}{dt} \|\partial_x^2 \beta\|_{L^2}^2 + \alpha_3 \int_{\Omega} \frac{\beta^2}{\omega} |\partial_x^3 \beta|^2 dx + \int_{\Omega} \omega |\partial_x^2 \beta|^2 dx \\ \lesssim (1+t)^3 (1 + \|\beta\|_{L^\infty}^2) \left(1 + \|(\partial_x u, \partial_x \omega, \partial_x \beta)\|_{L^\infty}^2\right) \|(\partial_x^2 u, \partial_x^2 \omega, \partial_x^2 \beta)\|_{L^2}^2,$$

for an implicit multiplicative constant $C > 0$, depending only on the chosen value of $\delta > 0$ and on the various parameters $(\nu, \alpha_1 \dots \alpha_4)$.

3.4 Closing the estimates

In order to close the estimates in some small time interval $[0, T]$, we introduce some convenient notations. For all $t \geq 0$, we set

$$E(t) := \|u(t)\|_{H^2}^2 + \|\omega(t)\|_{H^2}^2 + \|\beta(t)\|_{H^2}^2$$

to be the energy of the solution. Our goal is to bound it in terms of $E(0)$, which is defined as the same quantity, computed on the initial datum (u_0, ω_0, β_0) . We also define

$$E_0(t) := \|u(t)\|_{L^2}^2 + \|\omega(t)\|_{L^2}^2 + \|\beta(t)\|_{L^2}^2 \\ E_2(t) := \|\partial_x^2 u(t)\|_{L^2}^2 + \|\partial_x^2 \omega(t)\|_{L^2}^2 + \|\partial_x^2 \beta(t)\|_{L^2}^2,$$

and we denote by $E_0(0)$ and $E_2(0)$ the same quantities, when computed on the initial data. Therefore, in view of what we have said at the beginning of Subsection 3.3, we have

$$(36) \quad \forall t \geq 0, \quad E(t) \sim E_0(t) + E_2(t).$$

First of all, we observe that, from (16), (18) and (21), we deduce that

$$(37) \quad \forall t \geq 0, \quad E_0(t) \leq C_1 E_0(0) \exp\left(C_1 \int_0^t \|\partial_x u(\tau)\|_{L^\infty} d\tau\right),$$

for a ‘‘universal’’ numerical constant $C_1 > 0$ independent of any parameters and data of the problem.

Next, we put estimates (33), (34) and (35) together: in this way, after setting

$$(38) \quad A(t) := (1+t)^3 (1 + \|\beta\|_{L^\infty}^2) \left(1 + \|(\partial_x u, \partial_x \omega, \partial_x \beta)\|_{L^\infty}^2\right),$$

we get, for all $t > 0$, the inequality

$$\frac{d}{dt} E_2(t) + \int_{\Omega} \frac{\beta^2}{\omega} \left(|\partial_x^3 u|^2 + |\partial_x^3 \omega|^2 + |\partial_x^3 \beta|^2\right) dx + \int_{\Omega} \omega \left(|\partial_x^2 \omega|^2 + |\partial_x^2 \beta|^2\right) dx \leq C_2 A(t) E_2(t),$$

for a new constant $C_2 = C_2(\nu, \alpha_1 \dots \alpha_4) > 0$ only depending on the quantity in the brackets. An application of the Grönwall lemma thus yields

$$(39) \quad \forall t \geq 0, \quad E_2(t) \leq E_2(0) \exp\left(C_2 \int_0^t A(\tau) d\tau\right).$$

In view of the equivalence stated in (36), summing up inequalities (37) and (39) thus yields

$$(40) \quad \forall t \geq 0, \quad E(t) \leq E(0) \exp\left((C_1 + C_2) \int_0^t A(\tau) d\tau\right).$$

At this point, if we define the time $T > 0$ such that

$$T := \sup \left\{ t > 0 \mid \int_0^t A(\tau) d\tau \leq 2 \log 2 \right\},$$

we deduce from (40) that, for a new suitable constant $C_3 > 0$, only depending on the parameters $(\nu, \alpha_1 \dots \alpha_4)$, we have

$$(41) \quad \sup_{t \in [0, T]} E(t) \leq C_3 E(0).$$

As a direct consequence of the previous computations, we also get the blow-up criterion. Indeed, from (40) we see that, if $T^* < +\infty$ and if we have

$$\int_0^{T^*} A(\tau) d\tau < +\infty,$$

then $E(T^*) < +\infty$. Thus the solution cannot blow up at time $T^* > 0$, and hence it can be continued, by a standard argument, into a solution of system (4) with the same regularity (notice that, for this step, we also need the uniqueness of the solutions in the considered functional framework). At this point, we can remove the factor $(1+t)^3$ in the continuation criterion, by an easy argument. In addition, we observe that, thanks to (11), where we take $p = 2$, we can control

$$\int_0^{T^*} \|\beta\|_{L^\infty}^2 dt \lesssim \int_0^{T^*} \left(\bar{\beta}^2 + \|\beta\|_{L^2} \|\partial_x \beta\|_{L^2} \right) dt \lesssim \|\beta_0\|_{L^2}^2 T^* + \|\beta_0\|_{L^2} \int_0^{T^*} \|\partial_x \beta\|_{L^\infty} dt,$$

where we have used also (21). Therefore, we see that the previous integral is finite, if the integral appearing in (8) is finite. The continuation criterion is then proved.

Finally, let us derive the claimed lower bound on the lifespan of the solutions. For this, we recall the definition of the function $A(t)$, given in (38). Arguing as before and making use of inequality (11) with $p = 2$ for $\partial_x u$, $\partial_x \omega$ and $\partial_x \beta$, we find that

$$\begin{aligned} \|\beta\|_{L^\infty}^2 &\lesssim \bar{\beta}^2 + \|\beta\|_{L^2} \|\partial_x \beta\|_{L^2} \lesssim E(t) \\ \|\partial_x u\|_{L^\infty}^2 &\lesssim \|\partial_x u\|_{L^2} \|\partial_x^2 u\|_{L^2} \lesssim E(t), \end{aligned}$$

and analogous inequalities for $\partial_x \omega$ and $\partial_x \beta$. From those estimates, we deduce that, for any $t \in [0, T]$, one has

$$\begin{aligned} \int_0^t A(\tau) d\tau &\leq (1+t)^3 \int_0^t \left(1 + \|\beta\|_{L^\infty}^2 \right) \left(1 + \|(\partial_x u, \partial_x \omega, \partial_x k)\|_{L^\infty}^2 \right) d\tau \\ &\lesssim (1+t)^3 \int_0^t (1 + E(\tau))^2 d\tau. \end{aligned}$$

Owing to (41), in the end we get the following bound:

$$\forall t \in [0, T], \quad \int_0^t A(\tau) d\tau \lesssim (1+t)^3 t \left(\max \{1, E(0)\} \right)^2.$$

Therefore, using the definition of T , by continuity we discover that

$$C(1+T)^3 \left(\max \{1, E(0)\} \right)^2 T \geq \int_0^T A(\tau) d\tau = 2 \log 2,$$

for a suitable positive constant C . At this point, if $T \geq 1$, we are done. Assume, instead, that $T < 1$: then, the previous inequality implies that

$$\tilde{C} \left(\max \{1, E(0)\} \right)^2 T \geq \int_0^T A(\tau) d\tau = 2 \log 2,$$

for a new constant $\tilde{C} = 8C$, which implies the claimed lower bound on the time T .

4 Proof of the main results

In this section, we prove our main results, namely Theorems 1.1 and 1.3. First we deal with the local in time well-posedness issue for system (4). Then, we show that, under suitable symmetry assumptions, the solutions we have previously constructed in fact blow up in finite time.

4.1 Proof of the local existence

Here we prove the local in time well-posedness of system (4), contained in Theorem 1.1. Once again, we recall that we only show existence and uniqueness of solutions at H^2 level of regularity; the case of higher H^m regularity follows by similar arguments.

For notational convenience, throughout all this subsection we adopt the following convention: if $(f_\varepsilon)_{\varepsilon>0}$ is a sequence of functions in a normed space X , such that $(f_\varepsilon)_{\varepsilon>0}$ is *bounded* in X , we simply write $(f_\varepsilon)_{\varepsilon>0} \subset X$.

Construction of approximated solutions. To construct solutions to system (4) under the assumptions of Theorem 1.1, the first step is to remove the singularity created by the possible vanishing of the mean turbulent kinetic energy k .

Therefore, for any $\varepsilon \in]0, 1]$, we define

$$k_{0,\varepsilon} := \left(\sqrt{k_0} + \varepsilon \right)^2 = k_0 + 2\varepsilon \sqrt{k_0} + \varepsilon^2.$$

Observe that, since $\sqrt{k_0} \in H^2(\Omega)$, then also $k_0 \in H^2(\Omega)$. Therefore, we have

$$(42) \quad (k_{0,\varepsilon})_\varepsilon \subset H^2(\Omega), \quad \text{with} \quad k_{0,\varepsilon} \geq \varepsilon^2 > 0.$$

Next, for any fixed $\varepsilon \in]0, 1]$, we solve system (4) with respect to the initial datum

$$(u_0, \omega_0, k_{0,\varepsilon}) \in H^2(\Omega) \times H^2(\Omega) \times H^2(\Omega).$$

By adapting, for instance, the method of [8] to the one-dimensional setting, we construct a local in time solution $(u_\varepsilon, \omega_\varepsilon, k_\varepsilon)$ at the H^2 level of regularity, defined in some time interval $[0, T_\varepsilon]$ (see Theorem 1 in [8] for a precise statement).

The construction of solutions in [8] passes through Galerkin approximation. Thus, the approximated solutions are smooth, and our computations of Section 3 are fully justified for them. By lower semi-continuity of the norms and strong convergence of the approximate initial data towards the “true” initial datum $(u_0, \omega_0, k_{0,\varepsilon})$, we can establish some *uniform* properties for the sequence of solutions $(u_\varepsilon, \omega_\varepsilon, k_\varepsilon)_{\varepsilon \in]0, 1]}$.

First of all, owing to the lower bound (7), we observe that all the solutions are defined in a common time interval $[0, T]$, where

$$T := \inf_{\varepsilon \in]0, 1]} T_\varepsilon > 0.$$

Next, from (13), we gather that

$$(43) \quad (\omega_\varepsilon)_\varepsilon \subset L_T^\infty(L^\infty), \quad \text{with} \quad 0 < \frac{\omega_*}{\omega_* \alpha_2 t + 1} \leq \omega_\varepsilon(t, x) \leq \frac{\omega^*}{\omega^* \alpha_2 t + 1},$$

whereas (15) implies that $k_\varepsilon(t, x) \geq 0$ for all $(t, x) \in [0, T] \times \Omega$.

Moreover, from (41), we infer that

$$(44) \quad (u_\varepsilon, \omega_\varepsilon, \sqrt{k_\varepsilon})_\varepsilon \subset L_T^\infty(H^2) \times L_T^\infty(H^2) \times L_T^\infty(H^2)$$

and also that

$$(45) \quad \left(\sqrt{\frac{k_\varepsilon}{\omega_\varepsilon}} \partial_x^3 u_\varepsilon \right)_\varepsilon, \quad \left(\sqrt{\frac{k_\varepsilon}{\omega_\varepsilon}} \partial_x^3 \omega_\varepsilon \right)_\varepsilon, \quad \left(\sqrt{\frac{k_\varepsilon}{\omega_\varepsilon}} \partial_x^3 \sqrt{k_\varepsilon} \right)_\varepsilon \subset L_T^2(L^2).$$

As a consequence of (44), we also get that

$$(k_\varepsilon)_\varepsilon \subset L_T^\infty(H^2).$$

Stability analysis and convergence. Thanks to the previous uniform embedding properties, we can extract a weakly converging subsequence. More precisely, owing to (44), we deduce the existence of functions u , ω and k , all belonging to the space $L_T^\infty(H^2)$, such that, up to the extraction of a subsequence (not relabelled here), one has the convergences

$$u_\varepsilon \overset{*}{\rightharpoonup} u, \quad \omega_\varepsilon \overset{*}{\rightharpoonup} \omega \quad \text{and} \quad k_\varepsilon \overset{*}{\rightharpoonup} k \quad \text{in} \quad L_T^\infty(H^2).$$

The basic idea is that the triplet (u, ω, k) identified above is the sought solution to the Kolmogorov system (4). Nonetheless, of course the previous weak convergence properties are not enough to assert this, owing to the non-linear terms appearing in the equations. Therefore, we need some compactness properties for the sequence of solutions $(u_\varepsilon, \omega_\varepsilon, k_\varepsilon)_\varepsilon$: this is the goal of the next paragraph.

We start by analysing the equation for u_ε , which reads

$$\partial_t u_\varepsilon = -u_\varepsilon \partial_x u_\varepsilon + \nu \partial_x \left(\frac{k_\varepsilon}{\omega_\varepsilon} \partial_x u_\varepsilon \right).$$

From the previous uniform bounds (43) and (44), it is easy to deduce that $(\partial_t u_\varepsilon)_\varepsilon$ is uniformly bounded in the space $L_T^\infty(H^{-1})$, hence

$$(u_\varepsilon)_\varepsilon \subset L_T^\infty(H^2) \cap W_T^{1,\infty}(H^{-1}).$$

An application of the Ascoli-Arzelà theorem (see e.g. Lemma 3.6 in [17]) thus implies that $(u_\varepsilon)_\varepsilon$ is compact in the space $C_T(H^{-1})$. Then, up to the extraction of a further subsequence, one gathers the strong convergence

$$(46) \quad u_\varepsilon \longrightarrow u \quad \text{in} \quad C_T(H^s) \quad \forall 0 \leq s < 2,$$

where we have also used an interpolation argument between the strong convergence in $C_T(H^{-1})$ and the uniform bounds in $L_T^\infty(H^2)$.

It goes without saying that a similar argument applies also to the equations satisfied for ω_ε and k_ε , yielding strong convergence of $(\omega_\varepsilon)_\varepsilon$ and $(k_\varepsilon)_\varepsilon$ to, respectively, ω and k in $C_T(H^s)$, for any $s < 2$.

In particular, from (46) and Sobolev embeddings, we deduce the pointwise convergences

$$(47) \quad u_\varepsilon \longrightarrow u \quad \text{and} \quad \partial_x u_\varepsilon \longrightarrow \partial_x u \quad \text{almost everywhere in } [0, T] \times \Omega,$$

and analogous pointwise convergence properties also for the sequences $(\omega_\varepsilon)_\varepsilon$, $(k_\varepsilon)_\varepsilon$ and their first order space derivatives.

At this point, by making use of (47) and the strong convergence of the initial datum $k_{0,\varepsilon} \longrightarrow k_0$, it is easy to pass to the limit in the weak formulation of the equations for $(u_\varepsilon, \omega_\varepsilon, k_\varepsilon)_\varepsilon$. We thus discover that the triplet (u, ω, k) indeed solves system (4).

To conclude, we remark that, owing to (44) again, there exists a function $\beta \in L_T^\infty(H^2)$ such that $\sqrt{k_\varepsilon} \overset{*}{\rightharpoonup} \beta$ in that space. However, owing to the pointwise convergence of the k_ε 's, we deduce that $\beta = \sqrt{k}$. Studying the equation for the quantities $\sqrt{k_\varepsilon}$ and using an argument similar to the one employed above, we can get that \sqrt{k} belongs to $C_T(H^{-1})$, hence by interpolation it belongs to $C_T(H^s)$ for any $s < 2$. Thanks to this property, (44) and (45), we obtain also that the triplet (u, ω, \sqrt{k}) solves system (23).

4.2 Uniqueness of solutions

In order to complete the proof of Theorem 1.1, it remains us to prove uniqueness of solutions at the previous level of regularity. First of all, we need the following simple lemma, stating the equivalence of the equations for k and \sqrt{k} in the considered functional framework. We recall that, for $T > 0$, we have defined the functional space

$$\mathbb{X}_T(\Omega) := \left\{ (u, \omega, k) \mid \begin{array}{l} \omega, \omega^{-1}, k \in L^\infty([0, T] \times \Omega), \quad \omega > 0, \quad k \geq 0 \\ u, \omega, \sqrt{k} \in C([0, T]; L^2(\Omega)), \quad \partial_x u, \partial_x \omega, \partial_x \sqrt{k} \in L^\infty([0, T] \times \Omega) \end{array} \right\}.$$

Lemma 4.1. *Let $T > 0$ be a given time.*

- (a) *Let (u, ω, k) be a triplet in $\mathbb{X}_T(\Omega)$ and assume that k solves the third equation in (4) in the weak sense in $[0, T] \times \Omega$. Then $\beta := \sqrt{k}$ is a weak solution to equation (22), related to the initial datum $\beta_0 := \sqrt{k(0)}$.*
- (b) *Let the triplet of functions (u, ω, β) be such that $(u, \omega, \beta^2) \in \mathbb{X}_T(\Omega)$ and β is a weak solution to (22) in $[0, T] \times \Omega$, related to the initial datum β_0 . Set $k := \beta^2$. Then k is a weak solution to the third equation in (4) on $[0, T] \times \Omega$, related to the initial datum $k_0 := \beta_0^2$.*

Proof. The proof is based on simple verifications. To see that statement (b) holds true, it is enough to multiply equation (22) by the quantity 2β . Next, we check that all the integrations by parts are justified in the sense of $\mathcal{D}'([0, T] \times \Omega)$, owing to the regularity of the different quantities. For instance, from equation (22), one gets $\partial_t \beta \in L_T^2(H^{-1})$, thus the duality pairing

$$\langle \partial_t \beta, 2\beta \varphi \rangle_{L_T^2(H^{-1}) \times L_T^2(H^1)}$$

for $\varphi \in \mathcal{D}([0, T] \times \Omega)$, is well-justified; from this computation, one obtains that

$$(48) \quad 2\beta \partial_t \beta = \partial_t \beta^2 \quad \text{in the sense of } \mathcal{D}'([0, T] \times \Omega).$$

Notice that the previous equality makes sense, as (owing to the assumption $(u, \omega, \beta^2) \in \mathbb{X}_T(\Omega)$) one has $\sqrt{k} = \beta \in C_T(L^2)$. In order to see (48), let us take some $\varphi \in \mathcal{D}([0, T] \times \Omega)$, as above. First of all, by a density argument we check that the equality

$$\beta \partial_t \varphi = \partial_t(\beta \varphi) - \varphi \partial_t \beta$$

holds true in $L_T^2(H^{-1})$. Next, using that identity, we can compute

$$\begin{aligned} \langle \partial_t(\beta^2), \varphi \rangle_{\mathcal{D}' \times \mathcal{D}} &= -\langle \beta^2, \partial_t \varphi \rangle_{\mathcal{D}' \times \mathcal{D}} = -\langle \beta, \beta \partial_t \varphi \rangle_{L_T^2(H^1) \times L_T^2(H^{-1})} \\ &= -\langle \beta, \partial_t(\beta \varphi) \rangle_{L_T^2(H^1) \times L_T^2(H^{-1})} + \langle \beta, \varphi \partial_t \beta \rangle_{L_T^2(H^1) \times L_T^2(H^{-1})}. \end{aligned}$$

But $\beta \in W_T^{1,2}(H^{-1})$ and $\beta \varphi \in L_T^2(H^1)$, thus a duality argument (proceed e.g. by approximating β by smooth functions again) shows that

$$-\langle \beta, \partial_t(\beta \varphi) \rangle_{L_T^2(H^1) \times L_T^2(H^{-1})} = \langle \partial_t \beta, \beta \varphi \rangle_{L_T^2(H^{-1}) \times L_T^2(H^1)},$$

which in the end proves (48), as claimed. One can argue similarly for the terms presenting the space derivatives (using, for instance, an approximation argument as the one used in Proposition 9.4 of [2]). Finally, for the initial datum, we make use of the fact that $\beta \in C_T(L^2)$ to infer that $\beta^2 = k \in C_T(L^1)$. In the end, statement (b) of the lemma is proved.

Now, let us prove assertion (a). First of all, for $\varepsilon > 0$, we define the functions $k_\varepsilon := k + \varepsilon$. Then $k_\varepsilon \geq \varepsilon$ in $[0, T] \times \Omega$ and $(k_\varepsilon)_\varepsilon \subset C_T(L^1) \cap L_T^\infty(H^1)$. Moreover $k_\varepsilon \rightarrow k$ for $\varepsilon \rightarrow 0^+$, uniformly on $[0, T] \times \Omega$. Notice also that k_ε verifies the equation

$$\partial_t k_\varepsilon + u \partial_x k_\varepsilon - \alpha_3 \partial_x \left(\frac{k_\varepsilon}{\omega} \partial_x k_\varepsilon \right) + k_\varepsilon \omega = \alpha_4 \frac{k_\varepsilon}{\omega} |\partial_x u|^2 - \varepsilon \alpha_3 \partial_x \left(\frac{1}{\omega} \partial_x k_\varepsilon \right) - \alpha_4 \frac{\varepsilon}{\omega} |\partial_x u|^2.$$

Owing to the strictly positive lower bound on k_ε , we can multiply the previous equation by $k_\varepsilon^{-1/2}/2$, thus finding (by the same manipulations as in the previous step) an equation for $\sqrt{k_\varepsilon}$. At this point, passing to the limit for $\varepsilon \rightarrow 0^+$ is easy, thanks to the convergence properties of $(k_\varepsilon)_\varepsilon$. This finally implies that $\beta = \sqrt{k}$ solves (22) in the weak sense. The claimed equality for the initial datum is also satisfied, owing to the time continuity of \sqrt{k} with values in L^2 .

The lemma is now completely proven. \square

With this lemma at hand, we can prove a L^2 stability estimate for strong solutions of system (23) possessing vanishing mean turbulent kinetic energy. This is provided by the next statement. The uniqueness of solutions claimed in Theorem 1.1 will be a direct consequence of it.

Theorem 4.2. *Let the triplets (u_1, ω_1, k_1) and (u_2, ω_2, k_2) be two solutions to the Kolmogorov two-equation model (4). Assume that, for some time $T > 0$, the following hypotheses are satisfied for any $j \in \{1, 2\}$:*

- (i) *one has $\omega_j > 0$ and $k_j \geq 0$ on $[0, T] \times \Omega$;*
- (ii) *the quantities ω_j, ω_j^{-1} and $\sqrt{k_j}$ belong to $L^\infty([0, T]; L^\infty(\Omega))$;*
- (iii) *the functions $\partial_x u_j, \partial_x \omega_j$ and $\partial_x \sqrt{k_j}$ all belong to $L^\infty([0, T]; L^\infty(\Omega))$.*

After defining the quantities

$$U := u_1 - u_2, \quad \Sigma := \omega_1 - \omega_2 \quad \text{and} \quad B := \sqrt{k_1} - \sqrt{k_2},$$

assume also that U, Σ and B all belong to $C([0, T]; L^2(\Omega))$.

Then, there exist a positive constant $C = C(\nu, \alpha_1 \dots \alpha_4)$ and a function $\Theta \in L^1([0, T])$ such that, for all times $t \in [0, T]$, one has

$$\|U(t)\|_{L^2}^2 + \|\Sigma(t)\|_{L^2}^2 + \|B(t)\|_{L^2}^2 \leq \left(\|U(0)\|_{L^2}^2 + \|\Sigma(0)\|_{L^2}^2 + \|B(0)\|_{L^2}^2 \right) \exp \left(C \int_0^t \Theta(\tau) d\tau \right).$$

Proof. Before starting the computations, let us introduce some more notation and define

$$(49) \quad K := k_1 - k_2 = B(\beta_1 + \beta_2),$$

where, as in Section 3, we have set $\beta_i := \sqrt{k_i}$ for $i \in \{1, 2\}$. However, for lightening the notation, we will keep using the symbols $k_{1,2}$ and K when convenient.

Thanks to Lemma 4.1, we know that, for any $j \in \{1, 2\}$, the triplet (u_j, ω_j, β_j) solves system (23). Thus, we can perform the stability estimates on system (23). We point out that, in view of assumption (iii), the three functions U, Σ and B belong in particular to $L^2([0, T]; H^1(\Omega))$, with $\partial_t U, \partial_t \Sigma$ and $\partial_t B$ in $L^2([0, T]; H^{-1}(\Omega))$. Then, our energy estimates below are rigorously justified, in light of the regularity assumptions formulated in the statement of the theorem: this is a consequence of Theorem 3 in Section 5.9 of [6], a standard regularisation-in-space process (which we omit here) and related commutator estimates.

To begin with, we see that U solves the equation

$$(50) \quad \partial_t U + u_1 \partial_x U - \nu \partial_x \left(\frac{k_1}{\omega_1} \partial_x U \right) = -U \partial_x u_2 + \nu \partial_x (\Delta \partial_x u_2),$$

where we have defined

$$(51) \quad \Delta := \frac{k_1}{\omega_1} - \frac{k_2}{\omega_2} = -\frac{k_1}{\omega_1 \omega_2} \Sigma + \frac{1}{\omega_2} K.$$

A simple energy estimate for (50) yields

$$(52) \quad \frac{1}{2} \frac{d}{dt} \|U\|_{L^2}^2 + \nu \int_{\Omega} \frac{k_1}{\omega_1} |\partial_x U|^2 dx = \int_{\Omega} \left(\frac{1}{2} \partial_x u_1 - \partial_x u_2 \right) |U|^2 dx \\ + \nu \int_{\Omega} \frac{k_1}{\omega_1 \omega_2} \Sigma \partial_x u_2 \partial_x U dx + \nu \int_{\Omega} \frac{1}{\omega_2} K \partial_x u_2 \partial_x U dx.$$

It is apparent that the problem relies in the control of the last term on the right-hand side of the previous equality. However, we remark that U solves not only (50), but also the equivalent equation

$$\partial_t U + u_2 \partial_x U - \nu \partial_x \left(\frac{k_2}{\omega_2} \partial_x U \right) = -U \partial_x u_1 + \nu \partial_x (\Delta \partial_x u_1).$$

If we now perform an energy estimate on this relation, and we sum up the resulting expression with (52), we find

$$(53) \quad \frac{d}{dt} \|U\|_{L^2}^2 + \nu \int_{\Omega} \left(\frac{\beta_1^2}{\omega_1} + \frac{\beta_2^2}{\omega_2} \right) |\partial_x U|^2 dx \\ \leq 2 \left(\|\partial_x u_1\|_{L^\infty} + \|\partial_x u_2\|_{L^\infty} \right) \|U\|_{L^2}^2 \\ + \nu \left| \int_{\Omega} \frac{\beta_1^2}{\omega_1 \omega_2} \Sigma (\partial_x u_2 + \partial_x u_1) \partial_x U dx \right| + \nu \left| \int_{\Omega} \frac{1}{\omega_2} B (\beta_1 + \beta_2) (\partial_x u_2 + \partial_x u_1) \partial_x U dx \right|.$$

where, now, we have used (49) for expressing K and we have replaced k_j by β_j^2 .

At this point, we easily estimate

$$\nu \left| \int_{\Omega} \frac{\beta_1^2}{\omega_1 \omega_2} \Sigma (\partial_x u_2 + \partial_x u_1) \partial_x U dx \right| \\ \leq \nu \left(\|\partial_x u_1\|_{L^\infty} + \|\partial_x u_2\|_{L^\infty} \right) \left\| \frac{\beta_1}{\sqrt{\omega_1} \omega_2} \right\|_{L^\infty} \|\Sigma\|_{L^2} \left\| \frac{\beta_1}{\sqrt{\omega_1}} \partial_x U \right\|_{L^2} \\ \leq \delta \nu \int_{\Omega} \frac{\beta_1^2}{\omega_1} |\partial_x U|^2 dx + C(\delta, \nu) \left(\|\partial_x u_1\|_{L^\infty}^2 + \|\partial_x u_2\|_{L^\infty}^2 \right) \left\| \frac{\beta_1}{\sqrt{\omega_1} \omega_2} \right\|_{L^\infty}^2 \|\Sigma\|_{L^2}^2,$$

where $\delta > 0$ can be taken arbitrarily small and, like in Section 3, the constant $C(\delta, \nu) > 0$ depends only on the quantities in the brackets.

Similarly, for $j \in \{1, 2\}$, we can bound

$$\nu \left| \int_{\Omega} \frac{1}{\omega_2} B \beta_j (\partial_x u_2 + \partial_x u_1) \partial_x U dx \right| \\ \leq \nu \left(\|\partial_x u_1\|_{L^\infty} + \|\partial_x u_2\|_{L^\infty} \right) \left\| \frac{\sqrt{\omega_j}}{\omega_2} \right\|_{L^\infty} \|B\|_{L^2} \left\| \frac{\beta_j}{\sqrt{\omega_j}} \partial_x U \right\|_{L^2} \\ \leq \delta \nu \int_{\Omega} \frac{\beta_j^2}{\omega_j} |\partial_x U|^2 dx + C(\delta, \nu) \left(\|\partial_x u_1\|_{L^\infty}^2 + \|\partial_x u_2\|_{L^\infty}^2 \right) \left\| \frac{\sqrt{\omega_j}}{\omega_2} \right\|_{L^\infty}^2 \|B\|_{L^2}^2,$$

where $\delta > 0$ is again as small as one wants. Therefore, taking δ small enough and inserting those bounds into (53), we finally find

$$(54) \quad \frac{d}{dt} \|U\|_{L^2}^2 + \nu \int_{\Omega} \left(\frac{\beta_1^2}{\omega_1} + \frac{\beta_2^2}{\omega_2} \right) |\partial_x U|^2 dx \lesssim \Theta_1(t) \mathbb{E}(t),$$

where the implicit multiplicative constant depends only on $\nu > 0$, and where we have defined the energy function

$$\mathbb{E}(t) := \|U\|_{L^2}^2 + \|\Sigma\|_{L^2}^2 + \|B\|_{L^2}^2$$

and the function

$$\Theta_1(t) := \|(\partial_x u_1, \partial_x u_2)\|_{L^\infty} + \|(\partial_x u_1, \partial_x u_2)\|_{L^\infty}^2 \left(\left\| \frac{\beta_1}{\sqrt{\omega_1} \omega_2} \right\|_{L^\infty}^2 + \left\| \frac{\sqrt{\omega_1}}{\omega_2} \right\|_{L^\infty}^2 + \left\| \frac{1}{\omega_2} \right\|_{L^\infty} \right).$$

Next, we consider the quantity $\Sigma := \omega_1 - \omega_2$. Easy computations yield that Σ satisfies equivalently the two equations

$$\begin{aligned} \partial_t \Sigma + u_1 \partial_x \Sigma - \alpha_1 \partial_x \left(\frac{k_1}{\omega_1} \partial_x \Sigma \right) + \alpha_2 (\omega_1 + \omega_2) \Sigma &= -U \partial_x \omega_2 + \alpha_1 \partial_x (\Delta \partial_x \omega_2) \\ \partial_t \Sigma + u_2 \partial_x \Sigma - \alpha_1 \partial_x \left(\frac{k_2}{\omega_2} \partial_x \Sigma \right) + \alpha_2 (\omega_1 + \omega_2) \Sigma &= -U \partial_x \omega_1 + \alpha_1 \partial_x (\Delta \partial_x \omega_1). \end{aligned}$$

An energy estimate on both equations, as before, thus gives us

$$\begin{aligned} \frac{d}{dt} \|\Sigma\|_{L^2}^2 + \alpha_1 \int_{\Omega} \left(\frac{\beta_1^2}{\omega_1} + \frac{\beta_2^2}{\omega_2} \right) |\partial_x \Sigma|^2 dx \\ + 2\alpha_2 \int_{\Omega} (\omega_1 + \omega_2) |\Sigma|^2 dx \leq \frac{1}{2} \|(\partial_x u_1, \partial_x u_2)\|_{L^\infty} \|\Sigma\|_{L^2}^2 \\ + \|(\partial_x \omega_1, \partial_x \omega_2)\|_{L^\infty} \|\Sigma\|_{L^2} \|U\|_{L^2} + \alpha_1 \left| \int_{\Omega} (\partial_x \omega_1 + \partial_x \omega_2) \Delta \partial_x \Sigma dx \right|. \end{aligned}$$

Using the expression of Δ found in (51), it is easy to see (repeating the computations above) that, for any $\delta > 0$ arbitrarily fixed, we can bound

$$\begin{aligned} \alpha_1 \left| \int_{\Omega} (\partial_x \omega_1 + \partial_x \omega_2) \Delta \partial_x \Sigma dx \right| \\ \leq \delta \alpha_1 \int_{\Omega} \frac{\beta_1^2}{\omega_1} |\partial_x \Sigma|^2 dx + C \|(\partial_x \omega_1, \partial_x \omega_2)\|_{L^\infty}^2 \left\| \frac{\beta_1}{\sqrt{\omega_1} \omega_2} \right\|_{L^\infty}^2 \|\Sigma\|_{L^2}^2 \\ + \delta \alpha_1 \int_{\Omega} \left(\frac{\beta_1^2}{\omega_1} + \frac{\beta_2^2}{\omega_2} \right) |\partial_x \Sigma|^2 dx + C \|(\partial_x \omega_1, \partial_x \omega_2)\|_{L^\infty}^2 \left(\left\| \frac{\sqrt{\omega_1}}{\omega_2} \right\|_{L^\infty}^2 + \left\| \frac{1}{\omega_2} \right\|_{L^\infty} \right) \|B\|_{L^2}^2, \end{aligned}$$

for a suitable constant $C = C(\delta, \alpha_1) > 0$. This estimate immediately implies that, for a (implicit) positive multiplicative constant depending only on α_1 , one has

$$(55) \quad \frac{d}{dt} \|\Sigma\|_{L^2}^2 + \alpha_1 \int_{\Omega} \left(\frac{\beta_1^2}{\omega_1} + \frac{\beta_2^2}{\omega_2} \right) |\partial_x \Sigma|^2 dx + \alpha_2 \int_{\Omega} (\omega_1 + \omega_2) |\Sigma|^2 dx \lesssim \Theta_2(t) \mathbb{E}(t),$$

where the function Θ_2 is analogous to Θ_1 and is defined by

$$\Theta_2(t) := \|(\partial_x u_1, \partial_x u_2)\|_{L^\infty} + \|(\partial_x \omega_1, \partial_x \omega_2)\|_{L^\infty}^2 \left(\left\| \frac{\beta_1}{\sqrt{\omega_1} \omega_2} \right\|_{L^\infty}^2 + \left\| \frac{\sqrt{\omega_1}}{\omega_2} \right\|_{L^\infty}^2 + \left\| \frac{1}{\omega_2} \right\|_{L^\infty} \right).$$

The last step is to write an equation for $B := \beta_1 - \beta_2$ and perform an energy estimate for it. First of all, after setting

$$\tilde{\Delta} := \frac{\beta_1}{\omega_1} - \frac{\beta_2}{\omega_2} = -\frac{\beta_1}{\omega_1 \omega_2} \Sigma + \frac{1}{\omega_2} B,$$

in analogy with Δ in (51), from the last equation of (23) we see that B solves

$$\partial_t B + u_1 \partial_x B - \alpha_3 \partial_x \left(\frac{k_1}{\omega_1} \partial_x B \right) + \frac{1}{2} \omega_1 B = -U \partial_x \beta_2 + \alpha_3 \partial_x (\Delta \partial_x \beta_2) - \frac{1}{2} \Sigma \beta_2$$

$$\begin{aligned}
& + \frac{\alpha_4}{2} \tilde{\Delta} |\partial_x u_2|^2 + \frac{\alpha_4}{2} \frac{\beta_1}{\omega_1} \partial_x U (\partial_x u_1 + \partial_x u_2) \\
& + 2 \alpha_3 \tilde{\Delta} |\partial_x \beta_2|^2 + 2 \alpha_3 \frac{\beta_1}{\omega_1} \partial_x B (\partial_x \beta_1 + \partial_x \beta_2).
\end{aligned}$$

We now test the previous relation on B and integrate by parts whenever useful. First of all, we notice that, in order to handle the Δ term, we need to resort again to a symmetrization of the previous equation. Arguing as for estimating the corresponding term in the equation for U , we see that

$$\alpha_3 \left| \int_{\Omega} (\partial_x \beta_1 + \partial_x \beta_2) \Delta \partial_x B \, dx \right| = \alpha_3 \left| \int_{\Omega} (\partial_x \beta_1 + \partial_x \beta_2) \left(-\frac{k_1}{\omega_1 \omega_2} \Sigma + \frac{1}{\omega_2} B (\beta_1 + \beta_2) \right) \partial_x B \, dx \right|$$

can be controlled, for $\delta > 0$ arbitrarily small, by the quantity

$$\begin{aligned}
& \delta \alpha_3 \int_{\Omega} \frac{\beta_1^2}{\omega_1} |\partial_x B|^2 \, dx + C \|(\partial_x \beta_1, \partial_x \beta_2)\|_{L^\infty}^2 \left\| \frac{\beta_1}{\sqrt{\omega_1 \omega_2}} \right\|_{L^\infty}^2 \|\Sigma\|_{L^2}^2 \\
& + \delta \alpha_3 \int_{\Omega} \left(\frac{\beta_1^2}{\omega_1} + \frac{\beta_2^2}{\omega_2} \right) |\partial_x B|^2 \, dx + C \|(\partial_x \beta_1, \partial_x \beta_2)\|_{L^\infty}^2 \left(\left\| \frac{\sqrt{\omega_1}}{\omega_2} \right\|_{L^\infty}^2 + \left\| \frac{1}{\omega_2} \right\|_{L^\infty} \right) \|B\|_{L^2}^2,
\end{aligned}$$

for a suitable constant $C = C(\delta, \alpha_3) > 0$. Next, we have to control all the other terms appearing in the equation for B . Recall that we have symmetrized the equations: then, for $j \in \{1, 2\}$, we start to bound

$$\left| \int_{\Omega} \left(-u_j \partial_x B - U \partial_x \beta_j - \frac{1}{2} \Sigma \beta_j \right) B \, dx \right| \lesssim \mathbb{E} (\|\partial_x u_j\|_{L^\infty} + \|\partial_x \beta_j\|_{L^\infty} + \|\beta_j\|_{L^\infty}).$$

Next, using the definition of $\tilde{\Delta}$, we easily obtain

$$\left| \int_{\Omega} \left(\frac{\alpha_4}{2} \tilde{\Delta} |\partial_x u_j|^2 + 2 \alpha_3 \tilde{\Delta} |\partial_x \beta_j|^2 \right) B \, dx \right| \lesssim \mathbb{E} \|(\partial_x u_j, \partial_x \beta_j)\|_{L^\infty}^2 \left(\left\| \frac{\beta_1}{\omega_1 \omega_2} \right\|_{L^\infty} + \left\| \frac{1}{\omega_2} \right\|_{L^\infty} \right),$$

where the (implicit) multiplicative constant depends on the parameters α_3 and α_4 . Finally, for $\delta > 0$ as above, we can estimate

$$\begin{aligned}
& \alpha_3 \left| \int_{\Omega} \frac{\beta_j}{\omega_j} \partial_x B (\partial_x \beta_1 + \partial_x \beta_2) B \, dx \right| \\
& \leq \delta \alpha_3 \int_{\Omega} \frac{\beta_j^2}{\omega_j} |\partial_x B|^2 \, dx + C(\delta, \alpha_3) \left\| \frac{1}{\omega_j} \right\|_{L^\infty} \|(\partial_x \beta_1, \partial_x \beta_2)\|_{L^\infty}^2 \|B\|_{L^2}^2,
\end{aligned}$$

whereas, for any $\eta > 0$ arbitrarily small, we have

$$\begin{aligned}
& \alpha_4 \left| \int_{\Omega} \frac{\beta_j}{\omega_j} \partial_x U (\partial_x u_1 + \partial_x u_2) B \, dx \right| \\
& \leq \eta \nu \int_{\Omega} \frac{\beta_j^2}{\omega_j} |\partial_x U|^2 \, dx + C(\eta, \nu, \alpha_4) \left\| \frac{1}{\omega_j} \right\|_{L^\infty} \|(\partial_x u_1, \partial_x u_2)\|_{L^\infty}^2 \|B\|_{L^2}^2.
\end{aligned}$$

In the end, if we choose $\delta > 0$ small enough and we fix (say) $\eta = 1/4$, putting all the previous estimates together yields, for a suitable multiplicative constant $C = C(\nu, \alpha_1 \dots \alpha_4) > 0$, the inequality

$$\begin{aligned}
(56) \quad \frac{d}{dt} \|B\|_{L^2}^2 + \alpha_3 \int_{\Omega} \left(\frac{\beta_1^2}{\omega_1} + \frac{\beta_2^2}{\omega_2} \right) |\partial_x B|^2 \, dx + \int_{\Omega} (\omega_1 + \omega_2) |B|^2 \, dx \\
\leq C \Theta_3(t) \mathbb{E}(t) + \frac{\nu}{4} \int_{\Omega} \left(\frac{\beta_1^2}{\omega_1} + \frac{\beta_2^2}{\omega_2} \right) |\partial_x U|^2 \, dx,
\end{aligned}$$

where we have defined the function

$$\begin{aligned} \Theta_3(t) := & \|(\partial_x u_1, \partial_x u_2, \partial_x \beta_1, \partial_x \beta_2)\|_{L^\infty} + \|(\beta_1, \beta_2)\|_{L^\infty} \\ & + \|(\partial_x u_1, \partial_x u_2, \partial_x \beta_1, \partial_x \beta_2)\|_{L^\infty}^2 \\ & \times \left(\left\| \frac{\beta_1}{\sqrt{\omega_1 \omega_2}} \right\|_{L^\infty}^2 + \left\| \frac{\sqrt{\omega_1}}{\omega_2} \right\|_{L^\infty}^2 + \left\| \left(\frac{1}{\omega_1}, \frac{1}{\omega_2} \right) \right\|_{L^\infty} + \left\| \frac{\beta_1}{\omega_1 \omega_2} \right\|_{L^\infty} \right). \end{aligned}$$

We are now in the position of closing our stability estimates. Indeed, summing up (54), (55) and (56), after absorbing the last term on the right-hand side of (56) into the left-hand side, we finally get

$$\frac{d}{dt} \mathbb{E}(t) \lesssim \left(\Theta_1(t) + \Theta_2(t) + \Theta_3(t) \right) \mathbb{E}(t),$$

and an application of the Grönwall lemma concludes the proof. \square

We end this section by formulating a couple of remarks.

Remark 4.3. In light of estimates (54), (55) and (56), and of the definition of the functions $\Theta_{1,2,3}$, one could relax the L_T^∞ assumptions formulated in items (ii) and (iii) of Theorem 4.2. This would lead to uniqueness of solutions in a class which is actually larger than the one claimed in Theorem 1.1.

However, the conditions to require would be much more in number, and more complex to formulate. For this reason, we have decided to keep simple assumptions, and to limit ourselves to the present form of Theorem 4.2.

Remark 4.4. Observe that Theorem 4.2 is *not* a weak-strong uniqueness type result, because assumption (iii) has to be made on both triplets of solutions.

4.3 The blow-up mechanism

In this section, we prove Theorem 1.3 and show that, in general, the solutions constructed in the previous part blow up in finite time. More precisely, we are able to prove that, for initial data which possess a suitable symmetry, the corresponding solutions to the Kolmogorov system (4) exhibit a growth in the space derivative of the velocity field. The precise statement is the following one.

Theorem 4.5. *Let (u_0, ω_0, k_0) be a triplet of functions belonging to $H^3(\Omega)$, with $\omega_* \leq \omega_0 \leq \omega^*$, for some strictly positive constants $0 < \omega_* \leq \omega^*$, and with $k_0 \geq 0$ such that $\sqrt{k_0} \in H^3(\Omega)$. Assume the following conditions:*

- (i) $k_0(0) = 0$;
- (ii) u_0 is odd with respect to 0, while ω_0 and k_0 are even with respect to 0;
- (iii) $\partial_x u_0(0) < 0$.

Then, the corresponding H^3 solution (u, ω, k) , whose existence is guaranteed by Theorem 1.1, blows up in finite time. More precisely, if the solution does not blow up first at a different place, there exists a time $t_0 > 0$ such that

$$\lim_{t \rightarrow t_0^-} \partial_x u(t, 0) = +\infty.$$

Proof. We are going to show that the solution (u, ω, k) cannot be global arguing by contradiction: more precisely, we will prove that, assuming that the solution exists and remains smooth, then there will be a blow-up of the solution in a finite time $t = t_0$ exactly at $x = 0$.

Recall that we have supposed u_0 to be odd, and k_0 and ω_0 to be even with respect to the origin. As a consequence, besides the assumptions of Theorem 4.5, we also know that

$$\omega_0(0) > 0 \quad \text{and} \quad \partial_x k_0(0) = \partial_x \omega_0(0) = 0.$$

In addition, we notice that the equations preserve the previous symmetries. Therefore, by uniqueness of solutions, we deduce that

$$\forall t \geq 0, \quad u(t) \quad \text{is odd} \quad , \quad \omega(t), \quad k(t) \quad \text{are even}.$$

Since the functions (u, k, ω) are smooth in space, repeating the ODE analysis performed in Subsection 3.1 we easily find that

$$\forall t > 0, \quad \frac{d}{dt} k(t, 0) = 0 \quad \implies \quad \forall t \geq 0, \quad k(t, 0) = 0.$$

Therefore, at any time $t \geq 0$, the point $x = 0$ remains a point of minimum for the function $k(t)$. As a consequence, by symmetry we also get that $\partial_x k(t, 0) \equiv 0$ for all times.

Now, we differentiate the equation for u with respect to x and compute the resulting expression at the point $x = 0$. By virtue of the previous properties, we find that $\xi(t) := \partial_x u(t, 0)$ solves

$$(57) \quad \frac{d}{dt} \xi(t) = -\xi^2(t) + a(t) \xi(t), \quad \text{with} \quad a(t) := \frac{\partial_x^2 k(t, 0)}{\omega(t, 0)},$$

related to the initial datum $\xi|_{t=0} = \xi_0 := \partial_x u_0(0)$. Notice that $a(t) \geq 0$ for all times, since $x = 0$ is a point of minimum for $k(t)$, for all $t \geq 0$.

Observe that, by assumption, one has $\xi_0 < 0$, so $d\xi/dt$ starts negative; by (57), ξ is decreasing, thus it stays negative for all times. Then one deduces that

$$\frac{d}{dt} \xi \leq -\xi^2 \quad \implies \quad \xi(t) \leq \frac{\partial_x u_0(0)}{1 + \partial_x u_0(0) t} = -\frac{|\partial_x u_0(0)|}{1 - |\partial_x u_0(0)| t},$$

where the last equality holds, owing to the fact that $\partial_x u_0(0) < 0$. The previous inequality obviously implies blow-up in finite time for the quantity $\xi(t) = \partial_x u(t, 0)$.

Theorem 4.5 is thus completely proved, and so does Theorem 1.3. □

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